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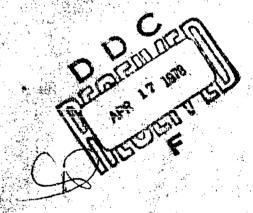


FEBRUARY 1978

PREDICTION OF ACCELERATION
HYDRODYNAMIC COEFFICIENTS FOR
UNDERWATER VEHICLES FROM
GEOMETRIC PARAMETERS

ODC FILE COPY

D. E. HUMPHREYS K. W. WATKINSON



MAVAL CONSTAL SYSTEMS LABORATORY

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ADMINISTRATIVE INFORMATION

This work was performed under the joint sponsorship of the projects Vehicle Hydrodynamics, Task Area Number ZF 61-112-601, Work Unit Number 00112-50-1A; Analytical Investigations of Countermeasure Launchers, Task Area Number SF 34-371-208, Work Unit Number 16436; and the Advance! Submarine Control Program, Subproject Number S0207, Work Unit 20259.

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Released by M. J. Wynn, Head Coastal Technology Department February 1978

Under Authority of G. G. Gould Technical Director

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SECURITY CLASSIFICATION OF THIS PAGE (When Date Entered)

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:		S S	6. PERPORMING ORG. REPORT NUMBER
(10)	D. E. Humphreys K. W. Watkinson		6. CONTRACT OR GRANT NUMBER(*)
1	9. PERFORMING ORGANIZATION NAME AND ADDRE: Naval Coastal Systems Laboratory Panama City, Florida 32407		10. PROGRAM ELEMENT, PROJECT, TASK AREA & WORK UNIT NUMBERS
62 7 N	11. COLTROLLING OFFICE NAME AND ADDRESS	(1)	February 1978
OF WA	14. MONITORING AGENCY NAME & ADDRESS(II dilla) (16) F6111d, (17) ZF6	1112000	108 15. SECURITY CLASS. (of this report) UNCLASSIFIED
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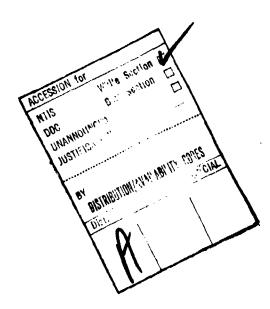
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Task Area No. ZF 61-112-001, Work Unit No. 00112-50-1A Task Area No. SF 34-371-208, Work Unit Number 16436 Subproject No. SO207, Work Unit 20259

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the value of the coefficient. Relationships for cross-coupled and non-linear hydrodynamic coefficients arising from potential flow theory are derived. All of the techniques have been implemented as part of the computer program GEORGE, a program for computing underwater vehicle dynamic stability characteristics in use at NCSL since 1972. Comparison of experimental and computed results is presented, along with a physical description of the generalized underwater vehicle used as the basis for the calculations and graphic representations contained in the report.



GLOSSALY OF TERMS

All symbols, subscripts, and nondimensionalized hydrodynamic coefficient expressions appearing in the body of this report are defined below. Any dimensions involved will be consistent with the foot-pound-second system of units. All angles are in degrees.

SYMBOLS

Symbol	<u>Definition</u>
AR	Aspect ratio of surface = b ² /s.
В	Buoyancy force which is positive upwards, $B = \rho \Psi$ (except where used in fluid kinetic energy equations).
Ъ	Exposed root-to-tip span of surface.
С	Chord.
c	Mean chord of surface = s/b = b/AR.
c _r	Root chord of surface.
ct	Tip chord of surface.
I _a	Added moment of inertia.
I _x	Moment of inertia of an underwater vehicle about the x-axis.
$^{\mathrm{I}}{_{\mathrm{a}}}$	Added moment of inertia about the x-axis.
^I y	Moment of inertia of an underwater vehicle about the y-axis.
$^{\mathrm{I}}\mathrm{y}_{\mathrm{df}}$	Moment of inertia about the y-axis of the fluid displaced by the body.
$\mathbf{I}_{\mathbf{z}}$	Moment of inertia of an underwater vehicle about the z-axis.
К	Moment about the x-axis exerted on the body by the dynamic pressure of the surrounding fluid.
^k 1	Lamb's inertial coefficient for a prolate ellipsoid in axial flow.

Symbol	(Continued) <u>Definition</u>
k ₂	Lamb's inertial coefficient for a prolate ellipsoid in cross flow.
k'p	Coefficient for added moment of inertia for a flat plate
k't b	Lamb's coefficient for added moment of inertia for a prolate ellipsoid body.
L	Overall length of the vehicle.
М	Moment about the y-axis exerted on the body by the dynamic pressure of the surrounding fluid.
m	Mass of the vehicle including the water in the free-flooding spaces.
m a	Additional mass.
^m df	Mass of fluid displaced by hull envelope.
N	Noment about the z-axis exerted on the body by the dynamic pressure of the surrounding fluid.
p	Component of angular velocity about the body fixed x-axis.
p	Component of angular velocity about the body fixed y-axis.
r	Component of angular velocity about the body fixed z-axis.
s	Exposed planform area of surface.
s _b	Maximum cross-sectional area of the body.
T	Total kinetic energy of the body-fluid system.
Tf	Kinetic energy of the fluid representing the system of impulsive pressures exerted by the surface of the body on the fluid during acceleration.
^T veh	Kinetic energy of the vehicle representing linear and angular momentum of the body.
t	Time.
U	Linear velocity of origin of body axes relative to a fluid-fixed axis system.

Symbol	(Continued) <u>Definition</u>
u	Component of U along the body x-axis.
₽	Displaced volume of the hull envelope.
v	Component of U along the body y-axis.
w	Component of U along the body z-axis.
x	Force along the x-axis exerted on the body by the dynamic pressure of the surrounding fluid.
X _{ncb}	x-distance from nose to cb.
$\mathbf{x}_{\mathtt{ncg}}$	x-distance from nose to cg.
X nht	x-distance from nose to centroid of the horizontal tail fins.
x	Longitudinal axis of the body fixed coordinate axis system.
Y	Force along the y-axis exerted on the body by the dynamic pressure of the surrounding fluid.
у	Transverse axis of the body fixed coordinate axis system.
Z	Force along the z-axis exerted on the body by the dynamic pressure of the surrounding fluid.
Z	Vertical axis of the body fixed coordinate axis system.
^z cb	The z-distance to the center of buoyancy (cb) from the center of gravity (cg), positive down.
α.	Angle of attack.
β	Angle of side-slip.
$^{\delta}$ bp	Deflection angle of bowplane (or sailplane).
δr	Deflection angle of rudder.
δs	Deflection angle of sternplane.

Symbol	(Continued) <u>Definition</u>
λ	Planform taper ratio of surface = c _t /c _r .
ω	Angular velocity.
ψ	Yaw angle.
ф	Roll angle.
ρ	Mass density of seawater.
θ	Pitch angle.
ı	The component in the plane of a surface of the perpendicular distance between the axis of rotation and the centroid of the area of the surface.
•	A prime over any symbol signifies nondimensionalization.
•	A dot over any symbol signifies differentiation with respect to time.

SUBSCRIPTS

Subscript	Definition
a	Added or additional inertia.
ъ	Body.
bp	Bow plane.
сb	Distance from cg to cb.
cg	Distance to cg.
df	Displaced fluid.
f	Fluid.
ht	Horizontal tail.
lvt	Lower vertical tail section.

Subscript (Continued) Definition n() Referenced from nose to some point, (), e.g., ncb implies nose-to-cb. Plate. Rudder. Sternplane. Shroud. sh Tip. Upper vertical tail section. uvt veh Vehicle. Pertaining to the x-axis. x Pertaining to the y-axis. У Pertaining to the z-axis.

NONDIMENSIONALIZED HYDRODYNAMIC COEFFICIENT EXPRESSIONS

Note that $X'\dot{u} = \partial X'/\partial \dot{u}'$, $X'\dot{p} = \partial X'/\partial \dot{p}'$, etc.

TABLE OF CONTENTS

	P.ge No	2
INTRODUCTION	. 1	
Background	. 1	
Acceleration Coefficients	. 3	
ACCELERATION COEFFICIENTS IN POTENTIAL FLOW THEORY	. 4	
METHODS FOR COMPUTING ACCELERATION HYDRODYNAMIC COEFFICIENTS.	. 9	
Introduction	. 9	
Longi idinal Coefficients	. 9	
X'	. 9	
Z'	. 14	
M'	. 16	
X'	. 16	
z'	. 17	
м',		
x'		
Z'		
м',		
Lateral Coefficients		
Y'	. 38	
Y',		
p		

TABLE OF CONTENTS (Cont'd)

																											P	age No	<u>,</u>
	Y'	ť	•		•		•	•		•	•	•	•	•	•	•	•	•	•	•	•	•	•	•	•	•	•	44	
	K'	ÿ	•	•	•	•	•	•	•	•	•	•	•		•	•	•	•	•	•	•	•	•	•	•	•	•	44	
	K'	Þ	•	•	•	•	•	•	•	•	•	•	•		•	•	•	•	•	•	•	•	•	•	•	•	•	50	
	K'	t	•	•		•	•	•	•	•	•	•	•	•	•	•		•	•	•	•	•	•	•	•	•	•	54	
	N'	'n	•	•	•	•	•	•		•	•		•	•	•	•	•	•	•	•	•	•	•	•	•	•	•	56	
	N'	, Þ	•	•	•	•	•	•	•	•	•	•	•	•		•	•	•	•	•	•	•	•	•	•	•	•	56	
	N'	' t	•	•	•	•	•	•	•	•	•	•	•	•	•	•	•	•		•	•	•	•	•	•	•	•	59	
SUMMARY .	•	•	•		•		•		•	•	•	•	•	•	•	•	•	•	•		•	•	•	•	•	•	•	59	
APPENDIX	Α ·	- :	DEI REI	RI'	VA:	TIC SNC	ON SH	OI PS	F 1	HYI FR(OR(OM	DD! El	(NA	AM] RGY	[C	CC EQU	EI JA	FF: TI(CI	EI	T.		•	•	•		•	A-1	
APPENDIX	В -	-	CO! EXI	MP/	AR:	IS(NC VTA	OI AL	D.	AN/ AT/	AL'	YT:	C.		P1	REI		CT:	101	NS •	w.	[T]	I .	•	•	•	•	B-1	
APPENDIX	C ·	-	DE: USI	SC! ED	RI F	PT:	[O] S]	n (Ens)F	T!	YP VI	IC. IY	AL Al	H LAN	ľD! LY	ROI SIS	OYI	IAN	AI(Æ!	HI(CL:	E.	•	•	•	C-1	
APPENDIX	D ·	-	DE:	SCI AJI	RII EC'	PT: TOI	IOI RY	N ()F [M	M UL	AN AT	EU'	VE)	RS •	u	SEI) :	IN •	•	•			•		•	•		D-1	

LIST OF ILLUSTRATIONS

Figure No.		Page No.
1	Positive Directions of Axes, Angles, Velocities, Forces, and Moments	10
2	Trajectory Simulation Plot of X $^{\circ}$ Contribution to X-Force During a $\delta_{\rm S}$ = ±30 $^{\circ}$ Dive/Climb Maneuver	12
3	Percent Contribution of X_{u}^{t} to X-Force During a $\delta_{g} = \pm 30^{\circ}$ Dive/Climb Maneuver	12
4	Lamb's Inertia Coefficients k_1 , k_2 and k^{\dagger} Plotted as a Function of Vehicle ℓ/d	15
5	Trajectory Simulation Plot of Z** Contribution to Z-Force During a $\delta_r = 30^\circ$ Steady Turn	
	(Automatic Depth-Keeping)	18
6	Percent Contribution of Z_{W}^{\dagger} to Z-Force During a $\delta_{r} = 30^{\circ}$ Steady Turn (Automatic Depth-Keeping)	18
7	Trajectory Simulation Plot of Z'. Contribution to Z-Force During a δ_s = $\pm 30^{\circ}$ Dive/Climb Maneuver	19
8	Percent Contribution of Z' to Z-Force During a $\delta_{\rm g} = \pm 30^{\circ}$ Dive/Climb Maneuver	19
9	Trajectory Simulation Plot of Z_w^* Contribution to Z-Force During a $\delta_r = \pm 30^\circ$ Turn Reversal Maneuver (Automatic Depth-Keeping)	20
10	Percent Contribution of Z_{W}^{\dagger} to Z-Force During a $\delta_{r} = \pm 30^{\circ}$ Turn Reversal Maneuver (Automatic Depth-Keeping)	
11	Root Locus Plot of u/o for a Generalized Underwater Vehicle, Varying Z':	21

LIST OF ILLUSTRATIONS (Cont'd)

Figure No.		Page No.
12	Root Locus Plot of w/δ_s for a Generalized Underwater Vehicle, Varying Z^{\bullet} .	21
13	Root Locus Plot of θ/δ_s for a Generalized Underwater Vehicle, Varying Z'.	21
14	Coefficients of Additional Mass for Rectangular Plates	25
15	Top View of Body and Horizontal Tail Fin Configuration with Descriptive Dimensions	26
16	Trajectory Simulation Plot of M' Contribution to Pitch Moment During a $\delta_r = 30^\circ$ Steady Turn (Automatic Depth-Keeping	31
17	Percent Contribution of M' to Pitch Moment During a $\delta_r = 30^{\circ}$ Steady Turn (Automatic Depth-Keeping)	31
18	Trajectory Simulation Plot of M $^{\circ}$ Contribution to Pitch Moment During a $\delta_s = \pm 30^{\circ}$ Dive/Climb Maneuver	32
19	Percent Contribution of M' to Pitch Moment During a $\delta_s = \pm 30^{\circ}$ Dive/Climb Maneuver	32
20	Trajectory Simulation Plot of M' Contribution to Pitch Moment During a $\delta_r = \pm 30^\circ$ Turn Reversal Maneuver (Automatic Depth-Keeping)	33
21	Percent Contribution of M_q^* to Pitch Moment During a $\delta_r = \pm 30^\circ$ Turn Reversal Maneuver (Automatic Depth-Keeping)	33
22	Root Locus Plot of u/δ_g for a Generalized Underwater Vehicle, Varying M_{g}^{\bullet}	34

LIST OF ILLUSTRATIONS (Cont'd)

Figure No.		Page No.
23	Root Locus Plot of w/δ_s for a Generalized	•
	Underwater Vehicle, Varying M'.	34
24	Root Locus Plot of θ/δ_s for a Generalized	
	Underwater Vehicle, Varying M'.	34
25	Coefficient of Additional Moment of Inertia for Rectangular Plates	37
26	Trajectory Simulation Plot of Y' Contribution to	
	Y-Force During a $\delta_r = 30^{\circ}$ Steady Turn	39
27	Percent Contribution of Y' to Y-Force During a $\delta_r = 30^{\circ}$ Steady Turn	39
28	Trajectory Simulation Plot of Y'. Contribution to	
	Y-Force During a $\delta_r = \pm 30^{\circ}$ Turn Reversal Maneuver	40
29	Percent Contribution of Y' to Y-Force During a	
	$\delta_{\mathbf{r}} = \pm 30^{\circ}$ Turn Reversal Maneuver	40
30	Root Locus Plot of v/δ_r for a Generalized	
	Underwater Vehicle, Varying Y'.	41
31	Root Locus Plot of ψ/δ_r for a Generalized	
	Underwater Vehicle, Varying Y'.	41
32	Root Locus Plot of ϕ/δ_r for a Generalized	
	Underwater Vehicle, Varying Y'.	41
33	Trajectory Simulation Plot of K' Contribution	
	to Roll Moment During a $\delta_r = \pm 30^{\circ}$ Turn Reversal Maneuver	47
34	Percent Contribution of K' to Roll Moment During	
	a δ = ±30° Turn Reversal Maneuver	47

LIST OF ILLUSTRATIONS (Cont'd)

Figure No.		Page No.
35	Root Locus Plot of ϕ/δ_r for a Generalized	
	Underwater Vehicle, Varying K'.	48
36 .	Dependence of the Additional Moment of Inertia on Taper Ratio	53
37	Variation of the Additional Moments of Inertia of a Single Plate With Dihedral Angle	54
38	Trajectory Simulation Plot of N $^{\circ}$ Contribution to Yaw Moment During a $\delta_r = \pm 30^{\circ}$ Turn Reversal Maneuver	. 60
39	Percent Contribution of N_r^{\dagger} to Yaw Moment During a $\delta_r = \pm 30^{\circ}$ Turn Reversal Maneuver	60
40	Root Locus Plot of v/δ_r for a Generalized Underwater Vehicle, Varying N':	61
41	Root Locus Plot of ψ/δ_r for a Generalized Underwater Vehicle, Varying N';	61
42	Root Locus Plot of ϕ/δ_r for a Generalized Vehicle, Varying N':	61

LIST OF TABLES

Table No.		Page No
1A	Longitudinal Sensitivity Analysis of Characteristic Equation Non-Dimensional Roots for a Generalized Underwater Vehicle, Varying X'. ±100%	13
1B	Longitudinal Sensitivity Analysis of U Numerator Non-Dimensional Roots for a Generalized Underwater Vehicle, Varying X': ±100%	13
1C	Longitudinal Sensitivity Analysis of W Numerator Non-Dimensional Roots for a Generalized Underwater Vehicle, Varying X'. ±100%	13
1D	Longitudinal Sensitivity Analysis of θ Numerator Non-Dimensional Roots for a Generalized Underwater Vehicle, Varying X'. \pm 100%	13
2A	Longitudinal Sensitivity Analysis of Characteristic Equation Non-Dimensional Roots for a Generalized Underwater Vehicle, Varying Z'. ±100%	22
2В	Longitudinal Sensitivity Analysis of U Numerator Non-Dimensional Roots for a Generalized Underwater Vehicle, Varying Z' ±100%	22
2C	Longitudinal Sensitivity Analysis of W Numerator Non-Dimensional Roots for a Generalized Underwater Vehicle, Varying Z' ±100%	22
2D	Longitudinal Sensitivity Analysis of θ Numerator Non-Dimensional Roots for a Generalized Underwater Vehicle, Varying Z $^{\rm t}_{\dot{w}}$ $\pm 100\%$	22
3A	Longitudinal Sensitivity Analysis of Characteristic Equation Non-Dimensional Roots for a Generalized Underwater Vehicle, Varying M'. ±100%	28
3B	Longitudinal Sensitivity Analysis of U Numerator Non-Dimensional Roots for a Generalized Underwater Vehicle, Varying M'. ±100%	28
3C	Longitudinal Sensitivity Analysis of W Numerator Non-Dimensional Roots for a Generalized Underwater Vehicle. Varving M'. ±100%	28

Table	No.		Page No
3D		Longitudinal Sensitivity Analysis of θ Numerator Non-Dimensional Roots for a Generalized Underwater Vehicle, Varying M¹. $\pm 100\%$	28
4A		Longitudinal Sensitivity Analysis of Characteristic Equation Non-Dimensional Roots for a Generalized Underwater Vehicle, Varying Z', $\pm 100\%$	30
4B		Longitudinal Sensitivity Analysis of U Numerator Non-Dimensional Roots for a Generalized Underwater Vehicle, Varying Z' ±100%	30
4C		Longitudinal Sensitivity Analysis of W Numerator Non-Dimensional Roots for a Generalized Underwater Vehicle, Varying Z $^{\prime}_{\dot{q}}$ $\pm 100\%$	30
4D		Longitudinal Sensitivity Analysis of θ Numerator Non-Dimensional Roots for a Generalized Underwater Vehicle, Varying Z $^{\prime}_{\dot{q}}$ $\pm 100\%$	30
5A		Longitudinal Sensitivity Analysis of Characteristic Equation Non-Dimensional Roots for a Generalized Underwater Vehicle, Varying M' $_{\hat{q}}$ $\pm 100\%$	35
5В		Longitudinal Sensitivity Analysis of U Numerator Non-Dimensional Roots for a Generalized Underwater Vehicle, Varying M', ±100%	35
5C		Longitudinal Sensitivity Analysis of W Numerator Non-Dimensional Roots for a Generalized Underwater Vehicle, Varying M', ±100%	35
5D		Longitudinal Sensitivity Analysis of θ Numerator Non-Dimensional Roots for a Generalized Underwater Vehicle, Varying M' $_{\hat{\bf q}}$ $\pm 100\%$	35
6 A		Lateral Sensitivity Analysis of Characteristic Equation Non-Dimensional Roots for a Generalized Underwater Vehicle, Varying Y' ±100%	42
6B		Lateral Sensitivity Analysis of V Numerator Non-Dimensional Roots for a Generalized Underwater Vehicle, Varying Y'. ±100%	42

Table No.		Page No.
6C	Lateral Sensitivity Analysis of ψ Numerator Non-Dimensional Roots for a Generalized Underwater Vehicle, Varying Y' $_{\mathring{\mathbf{V}}}$ ±100%	42
6D	Lateral Sensitivity Analysis of ϕ Numerator Non-Dimensional Roots for a Generalized Underwater Vehicle, Varying Y $^{\dagger}_{\dot{V}}$ ±100%	· 42
7 A	Lateral Sensitivity Analysis of Characteristic Equation Non-Dimensional Roots for a Generalized Underwater Vehicle, Varying Y', ±100%	45
7B	Lateral Sensitivity Analysis of V Numerator Non-Dimensional Roots for a Generalized Underwater Vehicle, Varying Y'. ±100%	45
7C	Lateral Sensitivity Analysis of ψ Numerator Non-Dimensional Roots for a Generalized Underwater Vehicle, Varying Y'. $\pm 100\%$	45
70	Lateral Sensitivity Analysis of ϕ Numerator Non-Dimensional Roots for a Generalized Underwater Vehicle, Varying Y', ±100%	45
8 A	Lateral Sensitivity Analysis of Characteristic Equation Non-Dimensional Roots for a Generalized Underwater Vehicle, Varying Y'; ±100%	46
8B	Lateral Sensitivity Analysis of V Numerator Non-Dimensional Roots for a Generalized Underwater Vehicle, Varying Y': ±100%	46
8C	Lateral Sensitivity Analysis of ψ Numerator Non-Dimensional Roots for a Generalized Underwater Vehicle, Varying Y': ±100%	46
8D	Lateral Sensitivity Analysis of ϕ Numerator Non-Dimensional Roots for a Generalized Underwater Vehicle, Varying Y': ±100%	46
9A	Lateral Sensitivity Analysis of Characteristic Equation Non-Dimensional Roots for a Generalized Underwater Vehicle, Varying K', ±100%	49

Table No.		Page No
9В	Lateral Sensitivity Analysis of V Numerator Non-Dimensional Roots for a Generalized Underwater Vehicle, Varying K'. ±100%	49
9C	Lateral Sensitivity Analysis of ψ Numerator Non-Dimensional Roots for a Generalized Under Vehicle, Varying K' $_{\hat{\mathbf{V}}}$ ±100%	49
9D	Lateral Sensitivity Analysis of ϕ Numerator Non-Dimensional Roots for a Generalized Underwater Vehicle, Varying K' $_{\dot{v}}$ ±100%	49
10A	Lateral Sensitivity Analysis of Characteristic Equation Non-Dimensional Roots for a Generalized Underwater Vehicle, Varying K'. ±100%	52
10B	Lateral Sensitivity Analysis of V Numerator Non-Dimensional Roots for a Generalized Underwater Vehicle, Varying K'. ±100%	52
10C	Lateral Sensitivity Analysis of ψ Numerator Non-Dimensional Roots for a Generalized Underwater Vehicle, Varying K'. $\pm 100\%$	52
100	Lateral Sensitivity Analysis of \$\phi\$ Numerator Non-Dimensional Roots for a Generalized Underwater Vehicle, Varying K': \$\pm\$100%	52
11A	Lateral Sensitivity Analysis of Characteristic Equation Non-Dimensional Roots for a Generalized Underwater Vehicle, Varying K'; ±100%	55
118	Lateral Sensitivity Analysis of V Numerator Non-Dimensional Roots for a Generalized Underwater Vehicle, Varying K': ±100%	55
110	Lateral Sensitivity Analysis of ψ Numerator Non-Dimensional Roots for a Generalized Underwater Vehicle, Varying K'; $\pm 100\%$	55
110	Lateral Sensitivity Analysis of \$\phi\$ Numerator Non-Dimensional Roots for a Generalized Underwater Vehicle, Varying K'. \pm 100%	55

Table No.		Page No.
12A	Lateral Sensitivity Analysis of Characteristic Equation Non-Dimensional Roots for a Generalized Underwater Vehicle, Varying N', ±100%	57
12B	Lateral Sensitivity Analysis of V Numerator Non-Dimensional Roots for a Generalized Underwater Vehicle, Varying N'. ±100%	. 57
12C	Lateral Sensitivity Analysis of ψ Numerator Non-Dimensional Roots for a Generalized Underwater Vehicle, Varying N' $_{\tilde{\mathbf{V}}}^{\star}$ $\pm 100\%$	57
1.20	Lateral Sensitivity Analysis of ϕ Numerator Non-Dimensional Roots for a Generalized Underwater Vehicle, Varying N' $_{\mathring{V}}$ ±100%	57
13A	Lateral Sensitivity Analysis of Characteristic Equation Non-Dimensional Roots for a Generalized Underwater Vehicle, Varying N'. ±100%	58
13B	Lateral Sensitivity Analysis of V Numerator Non-Dimensional Roots for a Generalized Underwater Vehicle, Varying N'. ±100%	58
13C	Lateral Sensitivity Analysis of ψ Numerator Non-Dimensional Roots for a Generalized Underwater Vehicle, Varying N'. $\pm 100\%$	58
13D	Lateral Sensitivity Analysis of ϕ Numerator Non-Dimensional Roots for a Generalized Underwater Vehicle, Varying N' , $\pm 100\%$	58
14A	Lateral Sensitivity Analysis of Characteristic Equation Non-Dimensional Roots for a Generalized Underwater Vehicle, Varying N'; ±100%	62
14B	Lateral Sensitivity Analysis of V Numerator Non-Dimensional Roots for a Generalized Underwater Vehicle, Varying N'. ±100%	62

Table No.		Page No.
14C	Late al Sensitivity Analysis of ψ Numerator Non-Dimensional Roots for a Generalized Underwater Vehicle, Varying N' \dot{r} ±100%	62
14D	Lateral Sensitivity Analysis of ϕ Numerator Non-Dimensional Roots for a Generalized Underwater Vehicle, Varying N' ±100%	62
15	Root Sensitivity Analyses for Derivative Variation of ±100%	64

INTRODUCTION

BACKGROUND

This report is primarily a compilation of methods for analytically predicting the acceleration hydrodynamic coefficients from the geometric and mass distribution characteristics of an underwater vehicle. It represents the current results of an effort to fin' the most generally applicable and accurate techniques for calculating these coefficients.

The need for such analytical techniques grew from the many unique underwater vehicle requirements inherent to the Naval Coastal Systems Laboratory mission areas. Typically these requirements include design, modification, or simulation of towed, tethered, or free vehicles for which analytical calculation of the hy rodynamic coefficients is necessary. This necessity arises from the need for some combination of the following:

- 1. Flexibility. The user's project is in a design phase where vehicle requirements imposed by hardware subsystems are subject to change and it is desirable to assess the effect of many more vehicle configurations than would be practical for experimental determination.
- Responsiveness. The project is of a research and development nature such that there is a requirement for rapid and frequent updating of the vehicle's predicted performance in response to changes in performance goals and/or subsystem hardware.
- 3. Economy. Cost constraints on the project will not permit use (or extensive use) of model basin testing, but the project still requires some degree of assurance that there are no potential catastrophic failures due to problems such as dynamic instability.

To more effectively meet each of these requirements, all of the methods presented herein, as well as methods for computing all of the remaining hydrodynamic coefficients, have been computer-implemented in the complete stability and control analysis program GEORGE. The GEORGE computer program uses vehicle geometric and mass distribution information to compute vehicle dynamic characteristics.

It should be pointed out that the procedure yielding the highest degree of confidence for an underwater vehicle design involves the use of both model basin testing and analytical techniques. An example of a desirable sequence of design steps would be the implementation of a computerized hydrodynamic math model of the proposed vehicle using conceptual information available in the very early design phases. Using this model and incorporating updated information as it becomes available, the effect of various sizing, distributing, and shaping ideas on the vehicle dynamics can be assessed.

As soon as a reasonably firm design configuration is laid out, model basin tests can be performed to validate the computer modeling. These tests can be much less extensive than those necessary to produce the parametric information required as a basis for new design techniques—techniques sufficiently general to maintain predictive accuracy for the redesigns or modifications frequently necessary in vehicle design. This reduction in model testing requirements is primarily the result of having a computer model to use in defining what needs to be verified experimentally, and in providing a sophisticated tool for describing complicated trends, thus reducing data requirements.

The validated computer simulation model can then be used in many ways to support the project. Trajectory simulations can be made for given operational conditions so that motion information can be generated for performance evaluation or definition of hardware specifications. Examples of work performed using simulations in this manner are trajectory analyses of submarine-launched acoustic countermeasure devices and motion prediction of a swimmer delivery vehicle to assess the motion influence on the design of a high-resolution, side-looking sonar.*

Another use has been to provide the hydrodynamic coefficients for analog or hybrid simulation models to study man-machine interface problems or to provide training. An example of recent work in this area is the implementation of a hybrid model of the Mine Neutralization Vehicle (MNV) as a performance evaluation and training device. The MNV simulation solves the nonlinear equations of motion for a tethered vehicle (including four propulsors, bowplanes, and cable tension) and produces a motion display for either body or inertial reference frame, all in real time. Of course, the computer simulation model is most often used to assess the influence of design variations of the vehicle hydrodynamic and control system parameters on dynamic stability and control characteristics.

^{*}Trajectory simulations have usually been performed using suitable modified versions of the standard six-degrees-of-freedom submarine simulation computer program, ZZMN, developed by DTNSRDC.

Having briefly stated the background for the prediction and simulation capability regarding the stability and control of underwater vehicles at NCSL, attention will now be focused on the analytical techniques for assigning a value to acceleration, nonlinear, and crosscoupled hydrodynamic coefficients.* These techniques have been applied to a wide variety of body types, including torpedoes, submarines, swimmer delivery vehicles (SDV's), and many other special purpose vehicles such as the Deep Submergence Rescue Vehicle (DSRV), the Shipboard Minehunting System (SMS) towed sonar body, and submarine-launched torpedo countermeasures devices.

Generality in applicability has been sought; that is, use of a method that yields accurate results for only a specific class of vehicles, such as torpedoes, has been avoided. A more generally applicable and, in the experience of the authors, more accurate, formulation based on theoretical derivation and experimental verification is presented. The relative importance of each hydrodynamic coefficient as it contributes to the responsive motion of the vehicle is noted. In addition, important relationships between acceleration, cross-coupled, and nonlinear hydrodynamic coefficients are derived from potential flow theory.

ACCELERATION COEFFICIENTS

The forces represented by acceleration hydrodynamic coefficients are attributed to the acceleration of fluid particles surrounding a submerged body. Some of the forces appear in the equations of motion as additions to terms representing the inertia of the body. They are known as acided mass or mass accession terms and, when combined with the inertial parameters of a body, have been called apparent, virtual, and hydrodynamic mass.

Analytical methods for estimating the acceleration coefficients of the bare hull differ from the methods used for the contributions of appendages in that they are not estimated using semi-empirical formulae. Rather, the body is assumed to be moving through an infinite, inviscid, circulation-free fixed fluid. For these conditions, Lamb⁽¹⁾ derives the complete expression for the kinetic energy of the body-fluid system which, coupled with inertia coefficients for prolate ellipsoids⁽¹⁾, yields estimates of body bare-hull acceleration derivatives.

^{*}Although techniques are also available for the prediction of static and dynamic velocity terms with viscous effects included, the scope of this report is limited to the acceleration coefficients.

⁽¹⁾Lamb, Sir Horace, Hydrodynamics, Dover Publications, New York, Sixth Edition, 1932. pps 172 and 154.

The basis for the added masses and moments due to body appendages, such as fins and bowplanes, is largely semi-empirical. Theoretical expressions for thin, flat plates moving in a perfect fluid have been derived from hydrodynamic theory with experimentally determined correction factors for non-ideal conditions and shapes. A great deal of the primary work done in this area was accomplished by $\operatorname{Munk}^{(2)(3)}$, with later significant experimental work by $\operatorname{Gracey}^{(4)}$, and $\operatorname{Malvestuto}$ and $\operatorname{Gale}^{(5)}$.

ACCELERATION COEFFICIENTS IN POTENTIAL FLOW THEORY

The forces and moments represented by the acceleration hydrodynamic coefficients can, to a very great extent, be modeled as potential flow phenomena. Neglecting the details of the boundary layer in modeling acceleration-dependent forces and moments acting on a submerged body yields quite satisfactory results for most stability and control simulation. Coefficient contributions due to appendages are most noticeably influenced by boundary layer effects. Acceleration coefficient values predicted using potential flow theory are probably most suspect in their application to fins that have a significant portion of their planform area in separated flow. There are a few considerations concerning these sources of error that should be mentioned before launching into further development.

Acceleration-dependent forces can arise from changes in the boundary layer, resulting, in turn, from changes in angle-of-attack or drift angle. There are several interesting facts in this regard. One is that the speed/power relationship for underwater vehicles generally encourages a designer to fair his vehicle so that the boundary layer has a significant thickness over as little of the vehicle as possible. The result

⁽²⁾ National Advisory Committee for Aeronautics T.N. No. 197, Some Tables of the Factor of Apparent Additional Mass, by Max M. Munk, 1924.

⁽³⁾ Munk, Max M., Fluid Mechanics, Part II, Aerodynamic Theory, Vol. 1, edited by W. F. Durand, Dover Publications, Inc., New York, 1934.

⁽⁴⁾ National Advisory Committee for Aeronautics T.N. No. 707, The Additional-Mass Effect of Plates as Determined by Experiments, by William Gracey. 1941.

⁽⁵⁾ National Advisory Committee for Aeronautics T.N. No. 1187, Formulas for Additional-Mass Corrections to the Moments of Inertia of Airplanes, by Frank S. Malvestuto, Jr. and Lawrence J. Gale, 1946.

is that the pressure distribution over 80 percent of the surface area of a typical underwater vehicle can be described very accurately by potential flow theory. Consequently, the acceleration coefficients most important to the description of underwater vehicle dynamics can usually be accurately computed as potential flow phenomena. Some less important acceleration coefficients, especially those resulting from the imbalance of added mass effects distributed fore and aft of the cg, can be significantly erroneous if the boundary layer is neglected. Unfortunately, those coefficients influenced by boundary layer effects suffer from similar problems in experimental model testing; e.g., improper Reynolds number scaling leads to nondynamic similarity in the modeled boundary layer.

Describing acceleration-dependent effects on fins is simplified somewhat by the vehicle designer's avoidance of the sizing and placement of fins so that they lose effectiveness in regions of separated flow. Consequently, more often than not, treatment of the added mass effects of fins as if they are in unspoiled flow provides good results. Tube-launched devices are often exceptions to this rule. They are generally constrained by tail size and after-body taper so that large portions of the fins are in separated flow.

Decisions on "effective" fin aspect ratios based on past experience with similar vehicles may be necessary to achieve a curacy in those coefficients with significant fin contributions. Since the fins on a torpedo are generally small, their contributions to the major coefficients of importance are less significant than is the case with other types of underwater vehicles. Roll coefficients are important exception since they are due almost completely to fin coefficients.

The acceleration, cross-coupled, and nonlinear coefficient equalities derived in Appendix A must be used with great caution when simulating vehicle trajectories. The NCSL approach has been to assess the importance of coefficient accuracy using two computer-implemented tools. One computer program performs root locus sensitivity studies for the linear coefficients on the vehicle of interest. The required accuracy is determined on a coefficient-by-coefficient basis, taking advantage of the insight that can be gained by a control system analysis approach to the vehicle system.

A more thorough evaluation of the math model simulating the vehicle fluid system, including the nonlinear and cross-coupled terms, must be done on a maneuver-by-maneuver basis. To assess the relative importance of each contributing force and moment, the authors have modified the standard six-degrees-of-freedom submarine simulation computer program (6)

⁽⁶⁾ Naval Ship Research and Development Center Test and Evaluation Report P-433-H-01, User's Guide NSRDC Digital Program for Simulating Submarine Motion ZZMN-Revision 1.0, by Ronald W. Richards, June 1971.

to plot the following on any maneuver: (1) time histories of the force (or moment) of each contributing coefficient, propulsor, buoyancy/gravity, and control surface term together with the total force (or moment); and (2) time histories of the motions of the vehicle with the contribution of a single term removed, and repeated for each term. Knowing the degree of accuracy expected in the computed value for each coefficient, one can analyze the plots on the maneuvers of interest and assign a level of confidence. This procedure will give an indication of which coefficients are most critical in achieving accurate simulation on a given maneuver.

Using these techniques and a multitude of comparisons with planar motion mechanism and rotating arm results, the methods presented in this report are, on the basis of the authors' experience, considered to be quite accurate on any hydrodynamically-faired vehicle for all but extreme maneuvers. This is especially true for linear models, as will be shown by examples in Appendix B.

The modeling accuracy achieved for acceleration terms under these assumptions is quite contrary to the extremely poor results obtained when the same assumptions are applied to velocity-dependent terms. The assumptions are: (1) the body is submerged so that there are no near-surface or near-bottom effects; (2) the fluid is inviscid; and (3) there is zero circulation.

Under these conditions the momentum equations for the body-fluid system may be written in the six Lagrange equations:

$$\frac{d}{dt}\frac{\partial T}{\partial u} = r\frac{\partial T}{\partial v} - q\frac{\partial T}{\partial w} - X,$$
(1)

$$\frac{\mathrm{d}}{\mathrm{d}\mathbf{r}}\frac{\partial \mathbf{T}}{\partial \mathbf{v}} = \mathbf{p}\,\frac{\partial \mathbf{T}}{\partial \mathbf{w}} - \mathbf{r}\,\frac{\partial \mathbf{T}}{\partial \mathbf{u}} - \mathbf{Y},\tag{1}$$

$$\frac{d}{dt}\frac{\partial T}{\partial w} = q \frac{\partial T}{\partial u} - p \frac{\partial T}{\partial v} - Z; \qquad (1)$$

$$\frac{d}{dt}\frac{\partial T}{\partial p} = w\frac{\partial T}{\partial v} - v\frac{\partial T}{\partial w} + r\frac{\partial T}{\partial q} - q\frac{\partial T}{\partial r} - K,$$
(2)

$$\frac{d}{dt}\frac{\partial T}{\partial q} = u\frac{\partial T}{\partial w} - w\frac{\partial T}{\partial u} + p\frac{\partial T}{\partial r} - r\frac{\partial T}{\partial p} - M,$$
 (2)

$$\frac{d}{dt}\frac{\partial T}{\partial r} = v \frac{\partial T}{\partial u} - u \frac{\partial T}{\partial v} + q \frac{\partial T}{\partial p} - p \frac{\partial T}{\partial q} - N, \qquad (2)$$

where X, Y, Z, K, M, and N represent the forces and moments exerted on the body by the pressure of the surrounding fluid, and the total kinetic energy T may be written

$$T = T_{veh} + T_{f}$$

with $T_{\rm veh}$ representing the linear and angular momentum of the vehicle itself, and $T_{\rm f}$ representing the system of impulsive pressures exerted by the surface of the solid on the fluid in the supposed instantaneous generation of motion from rest.

If we rewrite Equations (1) and (2), isolating the terms due to T_f , we obtain expressions for the forces exerted on the body by the dynamic pressure of the surrounding fluid, e.g.:

$$X = -\frac{d}{dt} \frac{\partial T_f}{\partial u} + r \frac{\partial T_f}{\partial v} - q \frac{\partial T_f}{\partial w}, \qquad (3)$$

and

$$K = -\frac{d}{dt} \frac{\partial T_f}{\partial p} + w \frac{\partial T_f}{\partial v} - v \frac{\partial T_f}{\partial w} + r \frac{\partial T_f}{\partial q} - q \frac{\partial T_f}{\partial r}, \qquad (3)$$

etc.

The fluid kinetic energy, T_f , can be expressed as a quadratic form of the body axis velocities u, v, w, p, q, and r with no explicit time dependence.* The quadratic expression for T_f may be written⁽¹⁾:

$$2T_{f} = Au^{2} + Bv^{2} + Cw^{2} + 2A'vw + 2B'wu + 2C'vu + Pp^{2} + Qq^{2}$$

$$+ Rr^{2} + 2P'qr + 2Q'rp + 2R'pq + 2Lup + 2Mvq + 2Nwr$$

$$+ 2F(vr + wq) + 2G(wp + ur) + 2H(uq + vp) + 2F'(vr - wq)$$

$$+ 2G'(wp - ur) + 2H'(uq - vp). \tag{4}$$

The body axes are generally chosen so that the xz-plane is a plane of symmetry (port-starboard symmetry). This eliminates nine of the coefficients, A', C', P', R', L, M, N, G, and G', since by symmetry arguments, sign changes in v, p, or r should not change T_f . Letting

$$F + F' = F_1$$
, $F - F' = F_2$, $H + H' = H_1$, and $H - H' = H_2$, (5)

⁽¹⁾ibid.

^{*}Recent investigations indicate that there can be significant time variant terms arising from vortices shed from the forward appendages and hull. The mathematical modeling of this phenomenon is accomplished using force and moment terms in the equations of motion that are convoluted in time. These terms are outside the set of acceleration, nonlinear, and cross-coupled hydrodynamic coefficients under consideration here.

the remaining 12 coefficients are included in the equation

$$2T_{f} = Au^{2} + Bv^{2} + Cw^{2} + 2B'wu + Pp^{2} + Qq^{2} + Rr^{2}$$

$$+ 2Q'rp + 2F_{1}vr + 2F_{2}wq + 2H_{1}uq + 2H_{2}vp .$$
(6)

The complete expressions, assuming no planes of symmetry, are presented in Appendix A.

Proceeding as indicated in Equation (3), the full six-degrees-of-freedom momentum equations can be written in terms of the 12 remaining coefficients (7). The coefficients are written in hydrodynamic coefficient notation as follows.

Added inertia terms:

$$A = -X_{\dot{u}}, \qquad P = -K_{\dot{p}}, \qquad B = -Y_{\dot{v}}, \qquad (7)$$

$$Q = -M_{\dot{q}}, \qquad C = -Z_{\dot{w}}, \text{ and } R = -N_{\dot{r}}.$$
 (7)

Inertial intermodal coupling terms:

$$F_1 = -Y_{\dot{r}} = -N_{\dot{v}}, \quad F_2 = -Z_{\dot{q}} = -M_{\dot{w}}, \text{ and } H_2 = -Y_{\dot{p}} = -K_{\dot{v}}.$$
 (8)

The remaining three are

$$B' = -X_{\dot{u}} = -Z_{\dot{u}}, Q' = -N_{\dot{p}} = -K_{\dot{r}}, \text{ and } H_1 = -X_{\dot{q}} = -M_{\dot{u}}.$$
 (9)

If, as shown in the complete expressions in Appendix A, there is also fore-aft symmetry (symmetry about the yz-plane), then, in addition to the aforementioned zeroed terms

$$B' = Q' = F_1 = F_2 = 0$$
.

Vehicles having symmetry about both the xz-plane and the xy-plane have the additional zeroed terms

$$B' = Q' = H_1 = H_2 = 0$$
.

⁽⁷⁾ Massachusetts Institute of Technology, Instrumentation Laboratory Report R-570-A, Deep Submergence Rescue Vehicle Simulation and Ship Control Analysis, by Charles Broxmeyer, Pierre P. Dogan, et al, 1967, p. 31.

These relations will be used extensively in the following coefficient formulations. It should be noted that the coefficient identities, such as Equations (8) and (9), result from derivation under no assumptions of symmetry. The existence of symmetry results in some of the terms being equal to zero. The complete set of identities of acceleration, nonlinear and cross-coupled coefficients is presented in Appendix A.

METHODS FOR COMPUTING ACCELERATION HYDRODYNAMIC COEFFICIEN ACCELERATION ACCELERATION HYDRODYNAMIC COEFFICIEN ACCELERATION HYDRODYNAMIC COEFFICIEN ACCELERATION HYDRODYNAMIC ACCELERATION HYDRODYNAMIC COEFFICIEN ACCELERATION ACCELERATION HYDRODYNAMIC COEFFICIEN ACCELERATION ACC

INTRODUCTION

Methods for estimating the acceleration-dependent hydrodynamic coefficients are given in this section. Each coefficient is described using the following format:

- 1. Introductory description.
- 2. Relative importance major or minor, and under what circumstances.
 - 3. Causal relationship important assumptions.
- 4. Analytical expression including normal range of coefficients in the expression and necessary references.

Appendix C provides a physical description of the underwater vehicle used as the basis for calculations illustrated graphically in this report, and Appendix D contains time history plots of the four maneuvers from which the trajectory simulation and percent contribution plots were obtained. Reference frame conventions are as shown in Figure 1.

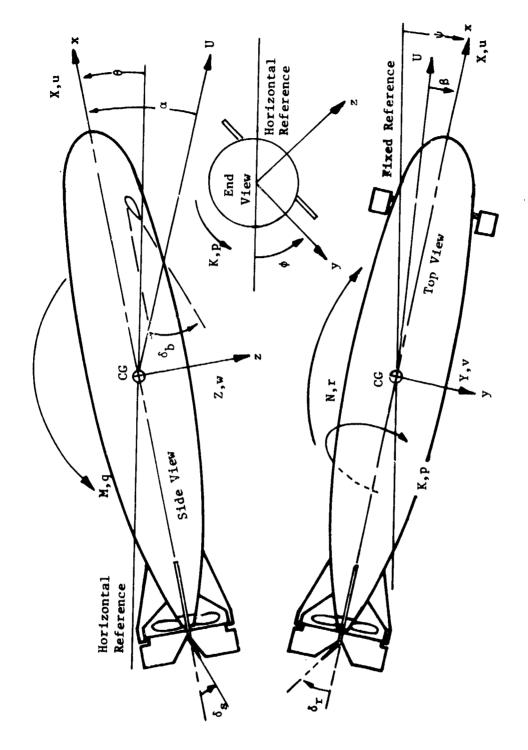
A word of caution is appropriate here. When comparing computed acceleration coefficient values with experimentally determined values, attention must be given to the location of the reference frame origin. If the origin during the experimental tests is not the vehicle cg, then erroneous comparisons will result for those coefficients where distances from the cg are involved.

LONGITUDINAL COEFFICIENTS

X'ڼ

1. X^{\dagger} is an acceleration coefficient resulting from the resisting force of the fluid on a vehicle accelerating in the x-direction. This

(Text Continued on Page 11)



POSITIVE DIRECTIONS OF AXES, ANGLES, VELOCITIES, FORCES, AND MOMENTS FIGURE 1.

force is required to change the flow field about the vehicle, whereas the static coefficients represent the forces necessary to maintain a flow field configuration at a constant vehicle velocity. More simply, X' may be thought of as a mass of fluid that must be accelerated along with the vehicle, thus giving the vehicle an "added" or "apparent" mass which is manifest during vehicle acceleration.

2. Insight to the relative significance of X', can be gained by noting that X', appears in the linear equations of motion in the mass term

$$(m' - X'_{\dot{i}})$$
.

Thus, in cases where the value of X' represents a significant contribution to this term, accuracy in predicting X' is important. The X' contribution to vehicular motion during a control surface step input is generally minor, and dissipates rapidly as the vehicle attains a steady state. In control surface reversal maneuvers, as illustrated in trajectory simulation and x-force percent contribution plots in Figures 2 and 3, X' becomes slightly more significant and longer-acting, but still remains a relatively minor influence. For nearly neutrally buoyant faired vehicles having an \$\ellsymbol{l}/d\text{ ratio higher than 8.0, X' typically contributes less than 3 percent to the mass term. X' has increasing relative significance for lower \$\ellsymbol{l}/d\text{ ratios.}\text{ Root locus studies for typical submarine and torpedo-like bodies has corraborated the foregoing, indicating very minor sensitivity to large percentage errors in predicting X' ...

When performing a two-degree-of-freedom linear analysis in the longitudinal plane (M and Z equations), three roots of the characteristic equation are obtained. The investigator may find a fourth root arising from the x-degree of freedom equation by assuming that the only significant terms are X' and X' (this is often acceptable for an axisymmetric vehicle). The fourth root is found by the relation

$$\sigma_4' = \frac{X'_u}{m' - X'_u} . \qquad (10)$$

Tables 1A through 1D present the results of a sensitivity analysis of nondimensional roots of the characteristic equation (transfer function denominator) and u, w, and θ (numerators) for a generalized underwater vehicle, varying X' $_{\dot{\bf u}}$ from -100 percent to +100 percent.

3. Mathematically, given the linearizing truncation of the Taylor series expansion to the first term, $X'_{\dot{u}}\dot{u}$ is the force additional to the inertial force occurring with a vehicular x-acceleration. In the computation of $X'_{\dot{u}}$, the body is approximated by a prolate ellipsoid having

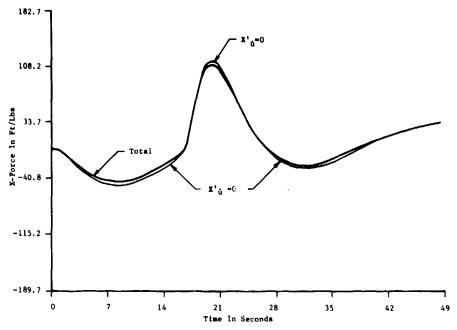


FIGURE 2. TRAJECTORY SIMULATION PLOT OF X' CONTRIBUTION TO X-FORCE DURING A $\delta_{_S}$ = $\pm 30\,^{\circ}$ DIVE/CLIMB MANEUVER

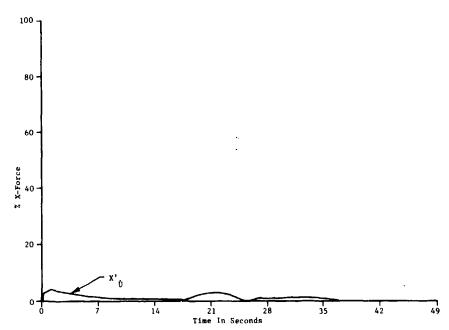


FIGURE 3. PERCENT CONTRIBUTION OF X' $_{\dot{u}}$ TO X-FORCE DURING A $_{\delta_{_{\bm{S}}}}$ = ±30° DIVE/CLIMB MANEUVER

TABLE 1D

LONGITUDINAL SENSITIVITY ANALYSIS OF CHARACTERISTIC EQUATION KIN-DIMENSIONAL ROOTS FOR A GENERALIZED UNDERMATER YEHICLE, VARYING \mathbf{x}_{1} \mathbf{x}_{1} \mathbf{x}_{2} 1005

LONGITUDINAL SENSITIVITY ANALYSIS OF B numerator non-dimensional right for a generalized imdernater vehicle, varying x_0^{-} ±100%

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like length and maximum diameter, and neglecting appendages to the body due to their comparatively small frontal area. Using the results of Lamb⁽¹⁾, derived assuming potential flow, we have:

4.
$$X^{\dagger}_{\hat{\mathbf{u}}} = X^{\dagger}_{\hat{\mathbf{u}}_{b}} = \frac{k_{1}^{m} df}{2\rho \ell^{3}}$$
, (11)

where

$$k_1 = \frac{\alpha_0}{2 - \alpha_0} , \qquad (12)$$

with

$$\alpha_0 = \frac{2(1-e^2)}{e^3} (l_2 \ln \left[\frac{1+e}{1-e}\right] - e) ,$$
 (13)

e being the eccentricity of the rotated ellipse expressed in common terms as

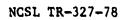
$$e = \frac{2}{\ell} \left(\frac{\ell^2}{4} - \frac{s_b}{\pi} \right)^{\frac{1}{2}}$$
 (14)

The inertia coefficient for axial flow, k_1 , ranges from 0.5 to 0 for fineness ratios of 1.0 to ∞ . k_1 is graphically presented in Figure 4.

۲'_ن

- 1. The acceleration coefficient Z' represents a z-force resulting from an acceleration in the x-direction.
- 2. Root locus sensitivity studies on typical submarine and torpedolike vehicles show extreme insensitivity to large magnitude excursions and sign changes of Z'. . Thus the coefficient is generally assumed negligible and set to zero.
- 3. Vehicles exhibiting symmetry about the xy-plane produce balanced z-forces during acceleration along x, indicating $Z'_{\dot{u}} = 0$. In potential theory, symmetry about the xy-plane or yz-plane causes $Z'_{\dot{u}} = 0$ as shown in Appendix A.
- 4. In the vast majority of cases, the combination of near symmetry about the xy-plane and minor influence on roots of the characteristic equation leads to setting $Z^{\dagger}_{\dot{u}} = 0$.

⁽¹⁾ibid.



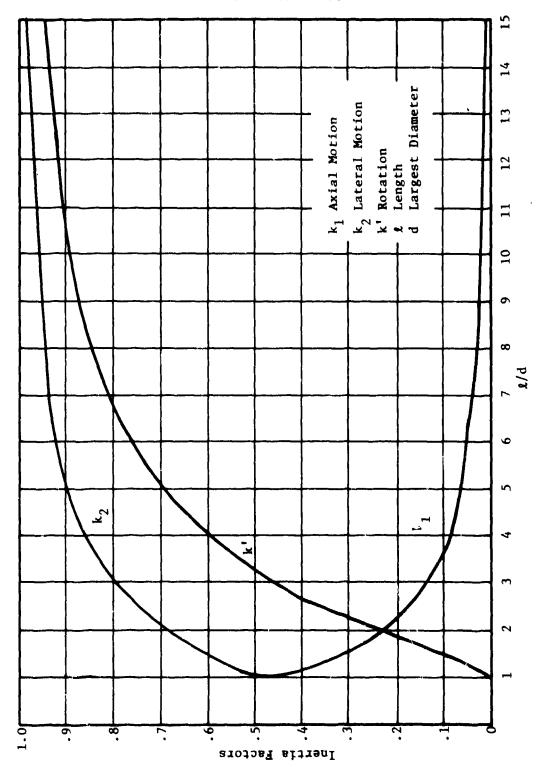


FIGURE 4. LAMB'S INERTIA COEFFICIENTS k_1 , k_2 AND k' PLOTTED AS A FUNCTION OF VEHICLE ℓ/d

M' ů

- 1. The term M'. represents the pitching moment resulting from an acceleration parallel to the x-axis.
- 2. Root locus sensitivity studies on typical submarine and torpedolike bodies show extreme insensitivity to large magnitude changes in M'_{ii} .
- 3. As was the case with Z', vehicles having symmetry or near symmetry about the xy-plane will create a negligible M-moment due to x-acceleration. Potential theory predicts M = 0 for yz-plane symmetry (see Appendix A). Extreme vertical separation of the cb and cg, coupled with a large X' force (arising from low l/d ratio), could conceivably cause a significantly large M'. For free-flooded vehicles, launcher devices, and torpedo-like bodies, this is a highly unlikely occurrence.
- 4. The combination of near symmetry about the xy-plane and minor influence on the roots of the characteristic equation usually leads to setting M' = 0. For a vehicle with an extremely large vertical cg cb separation and/or an X' value that is significant when compared to M' is calculated as

$$M'_{\dot{u}} = X'_{\dot{u}} \frac{z_{cb}}{\ell} ,$$

where z_{ch} = z-distance from the cg to the cb, positive down.

X'<u>ŵ</u>

- 1. The acceleration derivative X' represents the x-force resulting from an acceleration in the z-direction."
- 2. As with Z', and M', the roots of the characteristic equation show negligible change with X', set equal to zero, then varied positively and negatively over a range much greater than would be likely in reality.
- 3. Near balanced z-force during z-acceleration due to near foreaft symmetry yields resultants of a large z-force component and some pitching moment, but very little x-force. It is significant to note that Z' = X' according to potential theory (Equation (9)). If the cg and cb are significantly separated along both the x- and z-axes then, since those forces arising from potential flow theory are in a reference frame with a cb origin, the pitching moment term M' will yield an X' term when transferred to the cg. As with Z, though, even variation through extreme ranges showed that setting X' = 0 was justified for application to the vast majority of underwater vehicles.
 - 4. $X'_{\dot{w}} = 0$.

Z'w

- 1. The resisting hydrodynamic force resulting from a z-acceleration is represented in the equation of motion by Z'. The term (m'-Z') is the nondimensional "apparent mass" of the vehicle when one is considering heave acceleration.
- 2. Z' is one of the most important acceleration derivatives. Trajectory Simulation and z-force percent contribution plots in Figures 5 through 10 illustrate its significance to vehicular motion during various maneuvers. The root locus plots in Figures 11, 12, and 13 illustrate the wide variation in the locations of the poles due to variation in Z'. Tables 2A through 2D present the results of a sensitivity analysis of nondimensional roots of the characteristic equation (transfer function denominator) and u, w, and θ (numerators) for a generalized underwater vehicle, varying Z'. from -100 percent to +100 percent. As in X' and M', insight into the influence of Z' can be gained by noting the significance of Z' in the hydrodynamic mass term (m' Z'). Vehicles with l/d ratios greater than 8.5 will usually have nearly a 50 percent Z' contribution to the apparent mass. That is, the added mass will be roughly equal to the mass. Horizontal fins will contribute to Z' and requently cause it to be larger than m'.
- 3. As with the body is approximated by a prolate ellipsoid having like length and maximum diameter. Lamb's (1) inertial coefficient for such a body in crasswise flow is used as follows:

$$Z_{\dot{w}_{b}}^{\dagger} = \frac{k_{2}^{m} df}{k_{2} \rho \ell^{3}} = \frac{k_{2} \Psi}{k_{2} \ell^{3}}$$
 (15)

where

$$k_2 = \frac{\beta_0}{2 - \beta_0} , \qquad (16)$$

with

$$\beta_0 = \frac{1}{e^2} - \frac{1 - e^2}{2e^3} \ln \frac{1 + e}{1 - e} , \qquad (17)$$

e being the eccentricity of the rotated ellipse as noted in Equation (14). k_2 is graphically presented in Figure 14. Fortunately, this very important body coefficient has proved to be accurate to ± 5 percent in comparison with experimental data on the majority of underwater vehicles

⁽¹⁾ibid.

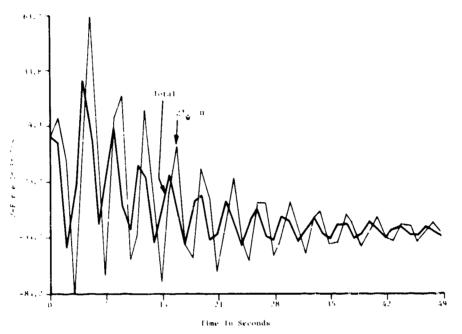


FIGURE 5. TRAJECTORY SIMULATION PLOT OF 21 CONTRIBUTION TO 2-FORCE DURING A $\delta_{\mu} \approx 30^\circ$ STEADY TURN (AUTOMATIC DEPTH-KEEPING)

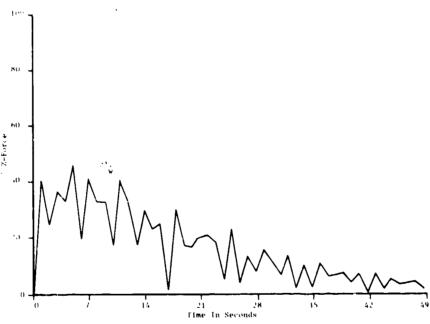


Figure 6. Percent contribution of 2 $_{\rm W}^{\star}$ to 2-force during a $_{\rm F}$ = 30 $^{\circ}$ Steady turn (automatic Depth-Keeping)

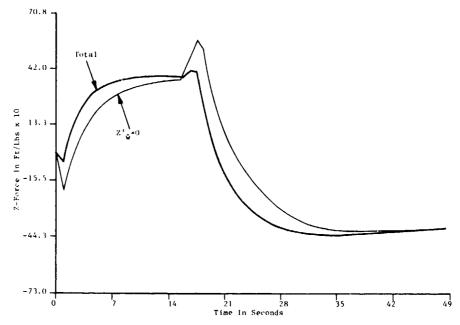


FIGURE 7. TRAJECTORY SIMULATION PLOT OF Z $_{\dot w}$ CONTRIBUTION TO Z-FORCE DURING A δ_g = ±30° DIVE/CLIMB MANEUVER

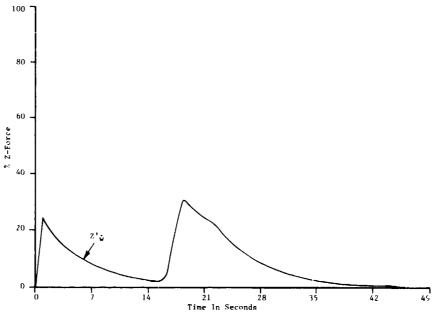


FIGURE 8. PERCENT CONTRIBUTION OF Z $_{\rm M}^{\star}$ TO Z-FORCE DURING A $_{\rm 6}$ = $\pm 30^{\circ}$ DIVE/CLIMB MANEUVER

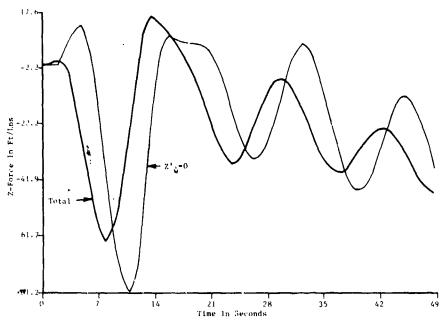


FIGURE 9. TRAJECTORY SIMULATION PLOT OF Z' CONTRIBUTION TO Z-FORCE DURING A $\delta_{\mu}=\pm30^{\circ}$ Turn reversal maneuver (automatic depth-keeping)

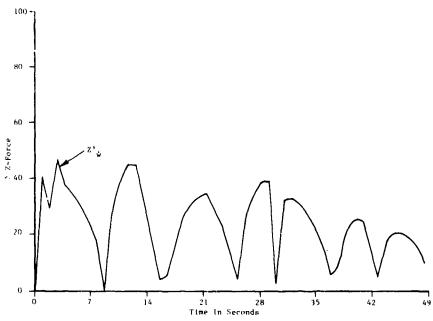


FIGURE 10. PERCENT CONTRIBUTION OF Z' $_{\dot w}$ TO Z-FORCE DURING A δ_{μ} = ±30° TURN REVERSAL MANEUVER (AUTOMATIC DEPTH-KEEPING)

NCSL TR-327-78

 Z^{*}_{W} variation using the formula $\frac{(100+Constant)}{100} \times (Z^{*}_{W}),$

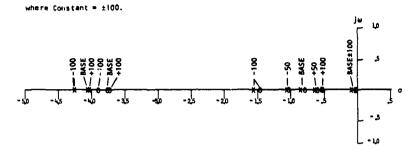


FIGURE 11. ROOT LOCUS PLOT OF u/δ_s for a generalized underwater vehicle, varying $Z^i_{\dot{w}}$

 $Z^{\dagger}_{-\omega}$ variation using the formula

$$\frac{(100+Constant)}{100} \times (Z'_{w}),$$

where Constant = ± 100 .

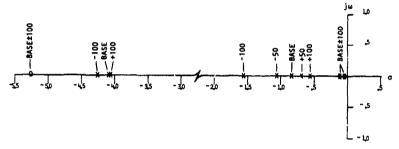


FIGURE 12. ROOT LOCUS PLOT OF w/ $\delta_{_{\rm S}}$ FOR A GENERALIZED UNDERWATER VEHICLE, VARYING Z' $_{\dot{W}}$

 $Z^{\bullet}_{\ \omega}$ variation using the formula

$$\frac{(100+Constant)}{100} \times (Z^{\dagger}_{\dot{w}}),$$

where Constant = ± 100 .

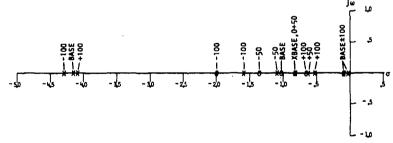


FIGURE 13. ROOT LOCUS PLOT OF $\theta/\delta_{_{\mathbf{S}}}$ FOR A GENERALIZED UNDERWATER VEHICLE, VARYING Z'

+100 - 019957 - 23 - 53578 - 32.3 - 27194 - 32.

-3,7421 -.34

23.- 5555.5- -3.61- 05553.3

005650 - 105+ +251 -. 049373

LONGITUDINAL SENSITIVITY ANALYSIS OF CHARALTERISTIC EQUATION NON-DIMENSIONAL REGIS FOR A GENERALIZED UNDERWITER VEHICLE, VARYING Z'_W ±100%

LONGITUDINAL SENSITIVITY ANALYSIS OF U NUMERATOR NON-DIMENSIONAL RODIS FOR A GENERALIZED UNDERWATER VEHICLE, VARYING 2', ±100?

TABLE 28

8 VAR ROOT 1 8 CMANGE ROOT 2 8 CHANGE ROOT 3 8 CHA -100 -0.45221 -.21 -1.4502 88.3 -3.2026 4.0

-251 -.049319 -.05 -.33602 :2.4 -3.2725 -27

-101 -. 049335

+101 -: 049356 MSE -- 049345

-50 -049283 -- 11 -1.0368 31.5 -5.2262

ŀ	-4.2485 2.4	0		-4.1731 .67	0	}	-4.1616 .28	0		-4.1543 .10	0		-4.1521		-4 7465 09	t		-4.1417 1 - 20	0		-4.135235	0
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	-1.5711	0		-1.0805	0		93360	0		86304	0		-, 82160	2	20202	4	,	-, 73345			65223	0
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LONGITUDINAL SENSITIVITY ANALYSIS OF 8 NUKERATOR NON-DYNENSIONAL ROOTS FOR A GENERALIZED UNDERNATER VEHICLE, VARYING 2', ±100% TABLE 20 LONGITUDINAL SENSITIVITY ANALYSIS OF 11 NUMERATOR NON-DIMENSIONAL ROOTS FOR A GENERALIZED UNDERMATER VEHICLE, VARYING $Z_{ij}^{\prime} \pm 100\%$

TABLE 2C

S CHANGE	82.0	32.7	14.0	2.2		-4.7	-11.0	-19.8	-33.0
NOOT 2	-2,3088	-1.3529	-1.1630	-1.0199	-1.0133 0	97201	01806 -	0	6. 1\$£89*-
S CHANGE									
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4	+	$\frac{1}{1}$	-	+	$\frac{1}{1}$	-	+	4		-	4		1		-	+	$\frac{1}{1}$	1	+	4		4	\dashv	-	+
S CHANGE				l																				l	
R00T 2		1]	1	1				7		10885	<u></u>	-	1]	_			\prod	_	Н		-
_		4		4	4		4	4				١			ļ	4	4			_			H		+
S CHANGE																									
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S VAR	1			5	3		7	2		[In-					Α	?		100	5		L	¥.		÷ 100;-

compared. In application to a wide variety of underwater vehicles, it has been our experience that if a body is nearly axisymmetric then Lamb's inertial coefficient is a reliable formulation for Z'. This is especially true for large ℓ/d ratios.

This is the first hydrodynamic coefficient with significant contribution from appendages to the body. The discussion and formulations that follow can be applied to any surface projecting from the body that can be approximated as a flat plate in accelerated fluid flow normal to its surface. The equations given will be for horizontal tail fins, but they are applicable to any horizontal surfaces such as bowplanes, sailplanes, fairwater planes, depressors, wings, faired struts, sails, etc. Tail contributions to other longitudinal hydrodynamic coefficients are formulated using the expressions given below. Theoretical values of the additional mass of a number of types of bodies (bodies of infinite length, bodies of revolution, ellipsoids or elliptic plates of finite dimensions) have been derived by Lamb⁽¹⁾ and Munk⁽²⁾⁽³⁾. Values for bodies of finite dimensions not covered by the theory; for example, rectangular plates and wings, have been the object of many experimental research programs on the phenomenon.

Extensive test programs on additional-mass and additional-moment effects were conducted by the National Advisory Committee for Aeronautics (NACA) from 1933 to 1940. The coefficients given here are based on the results of these tests as reported by $\operatorname{Gracey}^{(4)}$, and Malvestuto and $\operatorname{Gale}^{(5)}$. A very similar formulation is given by Pastor and Abkowitz⁽⁸⁾.

For a thin flat plate of infinite span moving in a perfect fluid at constant velocity, V, along the normal to its surface, the momentum imparted to the fluid per unit span is given by hydrodynamic theory as

$$\frac{\pi\rho c^2V}{4}$$
.

For plates of finite span, this expression must be corrected by the introduction of coefficients whose values depend on the dimensions of

⁽¹⁾ibid.

⁽a)ibid.

⁽a)ibid.

⁽⁴⁾ibid.

⁽⁵⁾ibid.

⁽⁸⁾ Naval Underwater Ordnance Station T.N. No. 120, Hydrodynamic Stability and Control Derivatives, by D. L. Pastor and M. A. Abkowitz, 1957.

the plate. The additional mass, m, for translation of a plate of span b, V is thus determined from the equation of linear momentum

$$m_{a}V = \frac{k\pi\rho c^{2}bV}{4} ,$$

so that

$$m_a = \frac{k\pi\rho c^2 b}{4} , \qquad (18)$$

where k is the coefficient of additional mass shown in Figure 14. This formulation is for one fin. So, with the proper nondimensionalizing terms added, the expression for two identical horizontal fin surfaces in a tail configuration is written

$$Z'_{\dot{\mathbf{w}}_{ht}} = \frac{1}{\frac{1}{2\rho} \ell^3} \frac{k\pi\rho c^2 b}{2} = \frac{k\pi c^2 b}{\ell^3}$$
 (19)

Note: Pastor and Abkowitz substituted a curve fitted expression for k:

$$k = \frac{1}{(1 + 1/AR)^{\frac{1}{2}}}$$
.

An effective chord, c, is usually used since most surfaces are not simple rectangular shapes. If a sternplane or any movable horizontal tail surface is present and is an extension of the fixed horizontal tail surface, then it should be included in the area for calculating c and s. If it is not a part of the fin, then its contribution to Z'. can be calculated separately, as can the bowplane contribution, by using Equation 19 to yield Z'. and Z'. For completeness and clarification, Figure 15 and bp the following relationships are given.

 $AR = aspect ratio of surface = b^2/s$

b = root-to-tip span of surface

s = planform area of surface

c_r = root chord of surface

 $c_t = tip chord of surface$

c = mean chord of surface = s/b = b/AR

 λ = planform taper ratio of surface = c_t/c_r .

(Text Continued on Page 26)

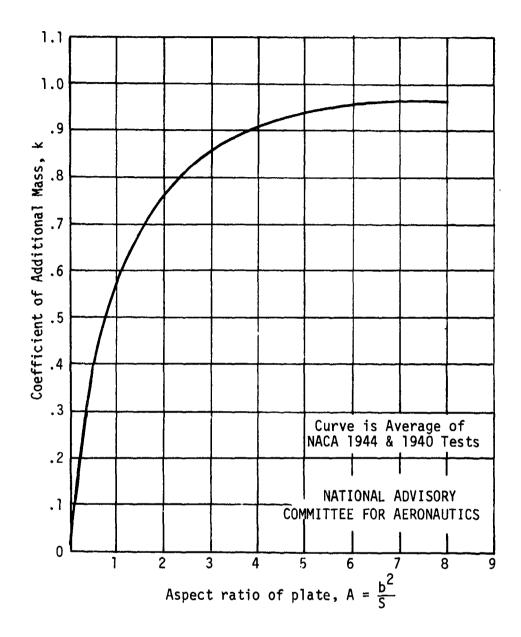


FIGURE 14. COEFFICIENTS OF ADDITIONAL MASS FOR RECTANGULAR PLATES

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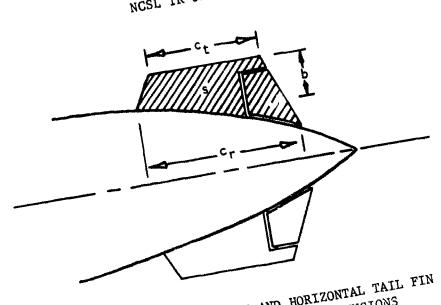


FIGURE 15. TOP VIEW OF BODY AND HORIZONTAL TAIL FIN CONFIGURATION WITH DESCRIPTIVE DIMENSIONS

A common way of computing the shroud contribution to Z' is to z' is to Z' is to Z' is z treat it as a reat prace in projection. That is, project the into the xy-plane and treat its projection as a flat plate.

Z'
$$\frac{k\pi c^2 r}{\sqrt{3}}$$
 (20)

where c = chord of the shroud and r = radius of the shroud. The complete where c = chord of the shroud and r = radius of the shroud. The shroud and r = radius of the shroud and r = radius of the shroud and r = radius of the shroud. The shroud and r = radius of the shroud. The shroud and r = radius of the shroud and r =

of the sum of the components
$$z'$$
 is z' is

M' w

- tiences a pitching moment proportional to the acceleration this experiences a pitching moment proportional to the acceleration, this experiences a piccning moment proportional to the acceptance moment is represented in the equations of motion as M' ...
- 2. M' usually has a minor influence on the characteristic equation of the vehicles that have large stern to ignored. It most important on vehicles that have vehicles and/or cg locations well forward on the vehicle. De ignored. It most important on ventures that have large of configurations and/or cg locations well forward on the vehicle.

Tables 3A through 3D present the results of a sensitivity analysis of nondimensional roots of the characteristic equation (transfer function denominator) and u, w, and θ (numerators) for a generalized underwater vehicle, varying M' from -100 percent to +100 percent.

3. The cg and cb of a submerged vehicle are generally in vertical alignment to schieve level trim at zero speed. Since the body additional mass effects are lumped as acting through the body cb (caution must be exercised when the cb is not coincident with the centroid of the volume displaced by the outer hull surface—as with "wet" vehicles) and the vehicle equations of motion are typically written about a cg origin, the resisting force to a z-acceleration is a pure Z' w term.

That is, there is no x-moment arm to cause the generation of a moment M' \dot{w} . The added mass and moment effects of the horizontal surfaces generally are the primary contributors to M' . Just as any force on a moment arm is calculated to produce a moment, so with M' \dot{w} we have

4.
$$M'_{\dot{\mathbf{w}}_{\mathbf{b}}} = Z'_{\dot{\mathbf{w}}_{\mathbf{b}}} \frac{(\mathbf{X}_{\mathbf{ncb}} - \mathbf{X}_{\mathbf{ncg}})}{\ell}$$
 (21)

$$M'_{\dot{w}_{t}} = Z'_{\dot{w}_{ht}} \frac{(X_{nht} - X_{neg})}{\ell}$$
 (22)

where

 $X_{ncb} = x-distance from nose to cb$

 $X_{ncg} = x$ -distance from nose to cg

X_{nht} = x-distance from nose to centroid
 of the horizontal tail fins .

Note: The distances are all positive numbers.

Substituting the x-distance to bowplanes or shroud into Equation (22) yields M' and M' . The complete expression is the sum of the component contributions; e.g.,

$$M'_{\dot{w}} = M'_{\dot{w}_b} + M'_{\dot{w}_{ht}} + M'_{\dot{w}_{bp}} + H'_{\dot{w}_s}.$$

(Text Continued on Page 29)

-3.7641 .22

+100 - 649356 ... 22 - ... 42 -... 3.7723

TABLE 30

+50 -.049351 .01 -.78955 -.21

+25 -. 049348 .006 -. 72072 -. 11

-3.7630 .1

-3.7576

.002 -.79122

INSE -. 049345 +10 -. 049346

TABLE 3A

LOMGITUDINAL SEMSI.IVITY ANALYSIS OF CHARACTERISTIC EQUATION NOM-DIMENSIONAL ROOTS FOR A GENERALIZED UNDERWATER VEHICLE, VARYING M. 100%

ROOT 3 \$ (HANGE

ROOT 1 X CHANGE ROOT 2 X CHANGE -- 059335 -- 02 -- 735493 -55

5 VAR

LONGITUDINAL SENSITIVITY ANALYSIS OF U NUMERATOR NOW-DIMENSIONAL POOTS FOR A GENERALIZED UNDERWATER VEHICLE, VARVING $\mathbf{M}_{\mathbf{w}'}^{*}$ ±103:

-3.7543 -- 35

-.002 -.79189

-101 -. 049344

-25<mark>1 -.048343 -.004 -.78240 .11 -2.7812 -.11</mark>

-50 -.049350 -.01 -.79325

200 20	1 0	1-4.13.82	67
10883		0	
10865	82277 .13 0	-4.1442 9	-14
-: 10885	82215 .07 0	-4.147:	07
- 10885	5218203	-4.7482	20
	82160	-4.1507	
	8213803	0	. 03
	82105 07	-4.1537	20.
	8205013	-4.2567	.14
18'-	-, 8134027	-4.152:	62.

TABLE 3C LONGITUDINAL SENSITIVITY ANALYSIS OF W NUMERATOR NON-DIMENSIONAL ROOTS FOR A GENERALIZED UNDERWATER VEHICLE, VARVING M_W ±100%

S CHANGE]	
R00T 3	1			-					-5,3461	-						-
S CHANGE																
R007 2			$\left \cdot \right $	+		+			10885		-		$\frac{1}{1}$			_
I CHANGE																
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s var	ā]	103		7		1		3548		+10+	1367		7	1.100	

LONGITUDIMAL SENSITIVITY AMALYSIS OF 8 NUMERATOR NON-DIMENSIONAL ROOTS FOR A GENERALIZED UNDERMATER VEHICLE, VARYING M'_w±100° ROUT 2 % CHANGE -1.0158 -1,0179 -1.0191 -1.0207 ROOT 1 S CHANGE MSE - 10885 -25¹ 1001+ 105-101+ 105+ **1**01 -100k +25^R

X'ġ

- 1. Forces along the x-direction caused by angular acceleration in pitch are represented by X' $_{\mathring{\textbf{q}}}$.
- 2. & 3. X' is usually an insignificant term as is its near-equal, M' (Equation (9)).

4.
$$X^{\dagger}_{\dot{q}} = 0$$
.

Z'ġ

- 1. z-forces arising from pitch angle accelerations are termed Z'.
- 2. Z', as its counterpart M', is a minor derivative, yet it can have some influence over the characteristic equation of a vehicle. It should therefore be calculated. Tables 4A through 4D present the results of a sensitivity analysis of nondimensional roots of the characteristic equation (transfer function denominator) and u, w, and θ (numerators) for a generalized underwater vehicle, varying Z', from -100 percent to ± 100 percent.

3. & 4.
$$Z^{\dagger}_{\dot{q}} = M^{\dagger}_{\dot{w}}$$
.

M'ġ

- 1. When a submerged body is accelerated in pitch, it must change the flow pattern in the surrounding fluid. The added moment needed to effect this change in flow pattern is represented in the equations of motion as an apparent moment of inertia, $(I'_v M'_o)$.
- 2. M' is a major derivative in influence on the characteristic equation. As with Z' and X', M' enters the equations of motion as an addition to the vehicle inertia term, in this case the moment of inertia about the y-axis. Trajectory simulation and pitch moment percent contribution plots in Figures 16 through 21 illustrate its significance to vehicular motion during various maneuvers. The root locus plots in Figures 22, 23, and 24 illustrate the wide variation in the locations of the poles due to variations in M'. Tables 5A through 5D present the results of a sensitivity analysis of nondimensional roots of the characteristic equation (transfer function denominator) and u, w, and θ (numerators) for a generalized underwater vehicle, varying M'. from -100 percent to +100 percent.

(Text Continued on Page 36)

25. 0101. -121. -123. -121. -121. 02. -23.

11. 1097. 2- 01.- 1006. - 100

TABLE 4A

LOWGITUDINAL SENSITIVITY ANALYSIS OF CHARACTERISTIC EQUATION NON-DIMENSIONAL ROOTS FOR A GENERALIZED IMDERMATER VEHICLE, VARYING Σ_{ij}^{*} ±100%

LONGITUDINAL SENSITIVITY ANALYSIS OF U RUNERATOR NON-SIMENSIONAL ROOTS FOR A GENERALIZED UNDERMATER VEHICLE, VARVING Z^{\prime}_{ij} = 100%

TABLE 48

S VAR ROOT 1 S CHANGE ROOT 2 S CHANGE 300T 3 S CHANGE -1007 -24252 -123

ROOT 4 % CHANGE	-4.141521	0	-4.1458 10	0	ł	-4.148005	0	Ì	-4.145302	0		-4.1501		ŀ	-4.1510 .02	0	ŀ	4.15%	0		4.1544	0		4.1587 .21
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ROOT 3 S CHANGE	62311	0	82235	0		82137	0		82175	0		82160	0		82145	c		82122	0		82085	0		00000
NOT 2 S CHANGE												_	_											
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1 1000			- 043024	c		- 043025	0		240200	0.0000	1	- 043026	0		043026	0		043027	0		- 043027			
K VAR. R		100	L	ş		_	-252		-	2		_	L		_	<u></u>		Ь	Ļ		_	\$	1	

+10f --648346 --002 --78139 --62 -3.7369 -02

255 - 043357 - 004 - 72115 - 05 - 317591 - 06

-101 - 1257. 2 - 20. 27127. 0 28128. -1.22

-3.7560

-- 75156 0

DASE# -- 049345

-56 -043343 -.004 -.72338 .10 -3.7517 -.11

TABLE 40 LOMGITUDINAL SENSITIVITY ANALYSIS OF 0 NUMERATOR NON-DIMENSIONAL ROOTS FOR A GENERALIZED UNDERRATER VEHICLE, VARYING $\mathbf{Z}_{ij}^{*} \pm 100^{2}$

TABLE 4C LONGITUDIAM, SENSITIVITY ANALYSIS OF W MUMERATOR MOM-DIMENSIONAL ROOTS FOR A GENERALIZED UNDERWATER VEHICLE, VARYIMG $Z_{\rm q}^{\prime}$ ±100%

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X CHANGE	0.5			-2.0			-1.0			-, 41						25.			1.0			2. I			4.3	
NOOT 3 % CHANGE	-5.1334	9		-5.2376	0		-5.2913	0		-5.3240	0		-5.3461	0		-5,3683	0		-5.4020	0		-5.4591	0		-5.5770	0
S CHANGE													_	_												
ROOT 2	_	Ц				L		F	-	F		L	10885	0				L	F		-			-		-
-	-	L		ŀ	L		-	L		L			Ľ	L		L	L		L	L		F	L		ŀ	H
A CHANGE	20.			70,	L		900			0			_			003	L		900-			20:-			03	
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s va.		8			ş.			Ş		-	-101			3			101			+2			ş			÷ 1

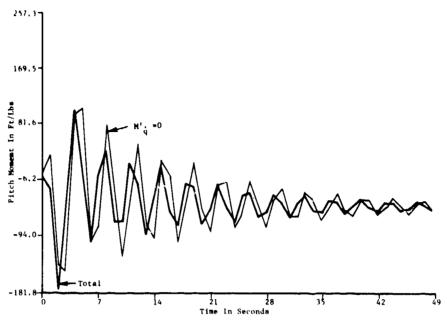


FIGURE 16. TRAJECTORY SIMULATION PLOT OF M' CONTRIBUTION TO PITCH MOMENT DURING A $\delta_p = 30^\circ$ STEADY TURN (AUTOMATIC DEPTH-KEEPING)

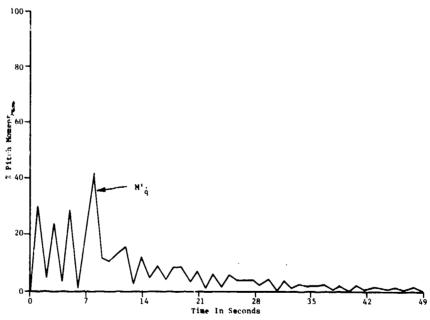


FIGURE 17. PERCENT CONTRIBUTION OF M' TO PITCH MCMENT DURING A $\rm d_p$ = 30° STEADY TURN (AUTOMATIC DEPTH-KEEPING)

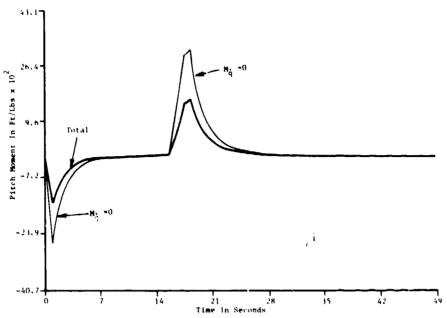


FIGURE 18. TRAJECTORY SIMULATION PLOT OF H' $_{q}$ CONTRIBUTION TG PITCH MOMENT DURING A $_{0}$ = $\pm 30^{\circ}$ DIVE/CLIMB MANEUVER

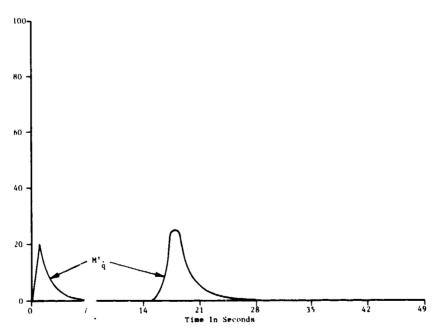


FIGURE 19. PERCENT CONTRIBUTION OF N' TO PITCH MOMENT DURING A δ_g = $\pm 30^\circ$ DIVE/CLIMB MANEUVER

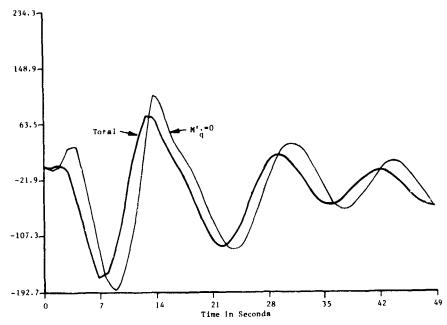


FIGURE 20. TRAJECTORY SIMULATION PLOT OF M' CONTRIBUTION TO PITCH MOMENT DURING A σ_{μ} = ±30° TURN REVERSAL MANEUVER (AUTOMATIC DEPTH-KEEPING)

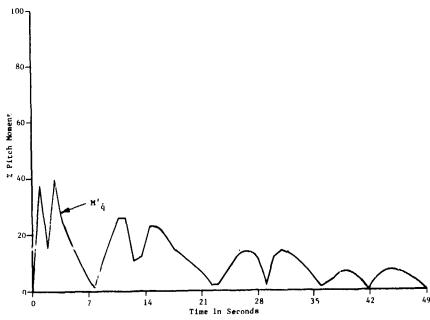


FIGURE 21. PERCENT CONTRIBUTION OF M $_{\dot{q}}$ TO PITCH MOMENT DURING A $_{r}$ = ±30° TURN REVERSAL MANSUVER (AUTOMATIC DEPTH-KEEPING)

NCSL TR-327-78

 $M'_{\dot{q}}$ variation using the formula $\frac{(100+Constant)}{100} \times (M'_{\dot{q}}),$

where Constant = ±100.

FIGURE 22. ROOT LOCUS PLOT OF u/δ_s FOR A GENERALIZED UNDERWATER VEHICLE, VARYING M' $_{\dot{\mathbf{q}}}$

 $\mathbf{H}^{\bullet}_{-q}$ variation using the formula

 $\frac{\text{(100+Constant)}}{100} \times (\text{M}^{1};),$

where Constant = ±100.

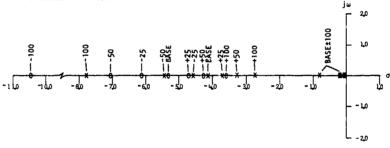


FIGURE 23. ROOT LOCUS PLOT OF W/6 $_{\rm S}$ FOR A GENERALIZED UNDERWATER VEHICLE, VARYING M $_{\rm Q}$

M' variation using the formula

 $\frac{(100+Constant)}{100}$ x (M':),

where Constant - ±100.

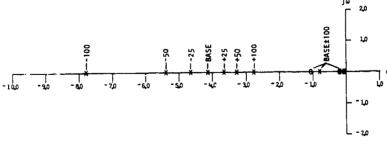


FIGURE 24. ROOT LOCUS PLOT OF θ/δ_s FOR A GENERALIZED UNDERWATER VEHICLE, VARYING M $^{\prime}_{,0}$

+50 0 -: 049519 35

-.049432

+25.

LONGITUDINAL SENSITIVITY ANALYSIS OF ${\bf e}$ numerator non-dimensional roots for a generalized underwater vehicle, varying ${\bf w}_{({\bf d})} \pm 100{\bf s}$

TABLE 5A LOWGITUDINAL SENSITIVITY ANALYSIS OF CARRACTERISTIC EQUATION NON-DIMENSIONAL ROOTS FOR A GENERALIZED UNDERWATER VEHICLE, VARYING $M_{\gamma}^{*}\pm 100\%$

LONGITUDIMAL SENSITIVITY ANALYSIS OF U NUMERATOR NGN-DIMENSIONAL ROOTS FOR A GENERALIZED UNDERWATER VEHICLE, VARYING $\mathbf{H}_{\mathbf{q}}^{+}$ ±100°2

* CHANGE ROOT 2 * CHANGE -- 68 -- 80122 1.2

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-. 79419

-25<mark>1 - 049260</mark>

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<u>ا</u>	0		-		0		0	_

TABLE 5C
LONGITUDINAL SENSITIVITY ANALYSIS OF H HUMERATOR HOM-DIMENSIONAL ROOTS
FOR A GENERALIZED UNDERBATER VEHICLE, VARYING M° ±100%

S CHANGE ,	97.1		32.6		14.1	7	5.2					[- 4.7	7	[-71.0		8.6		-33.1	7
F 1001 3	-10,539	0	- 7.0370	0	- 6.0990	0	- 5.6239	0		-5.3461	0		-5.0942	0		-4.7578	0	-4.2855	0	-3.5743	0
ROOT 2 % CHANGE	-									- 10885	7										
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	-100			17.		110-		101-		MSE	+10]	M364	163	No.	100		1001÷	

3. & 4. The body contribution to $M'_{\tilde{q}}$ is calculated assuming the body is a prolate ellipsoid. Thus, we use Lamb's coefficient for added moment of inertia and the parallel axis theorem

$$M'_{\dot{q}_{b}} = \frac{-k'_{b}^{I}y_{df}}{\frac{1}{2\rho}\ell^{5}} + Z'_{\dot{w}_{b}} \frac{(X_{ncb} - X_{ncg})^{2}}{\ell^{2}}$$

where

$$k'_{b} = \frac{e^{4}(\beta_{0} - \alpha_{0})}{(2 - e^{2}) \{2e^{2} - (2 - e^{2})(\beta_{0} - \alpha_{0})\}}$$

 α_0 and β_0 are defined in Equations (13) and (17). These equations are repeated here for convenience:

$$\alpha_0 = \frac{2(1 - e^2)}{e^3} \left(\frac{1}{2} \ln \left[\frac{1 + e}{1 - e} \right] - e \right) , \qquad (13)$$

$$\beta_0 = \frac{1}{e^2} - \frac{1 - e^2}{2e^3} \ln \frac{1 + e}{1 - e} , \qquad (17)$$

I = moment of inertia about the y-axis of the
ydf mass of the fluid displaced by the body.

The tail contribution to M^{\dagger} is in the same form as the body terms. That is, there is an added moment of inertia term and a term representing an added mass acting on a moment arm from the cg. The added moment of inertia of a flat plate resulting from rotation about an axis in the plane of the plate and parallel to the span is given by

$$I_a = \frac{k_p^1 \pi \rho b^2 c^3}{48}$$

where k' is the coefficient of added moment of inertia for a flat plate plotted on Figure 25. When c/b is sufficiently small (AR sufficiently large), this term becomes negligible compared with the second term in M', given as q_t

$$z'_{\text{wht}} \frac{(x_{\text{nht}} - x_{\text{neg}})^2}{\ell^2}$$
.

(Text Continued on Page 38)

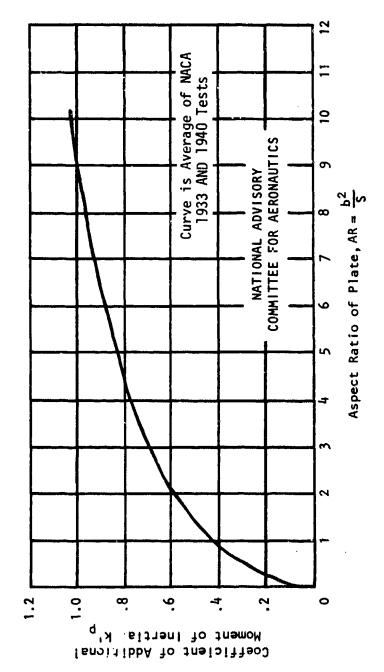


FIGURE 25. COEFFICIENT OF ADDITIONAL MOMENT OF INERTIA FOR RECTANGULAR PLATES

The horizontal tail surface contribution is then

$$M'_{\dot{q}_{ht}} = -\frac{1}{3^{0.25}} \frac{(k'_p \pi \rho b^2 c^{-3})}{48} + Z'_{\dot{w}_{ht}} \frac{(x_{nht} - x_{ncg})^2}{\ell^2}.$$

Similar expressions may be written for bowplanes and a shroud so that complete M' formula would be

$$\mathbf{M'}\,\dot{\mathbf{q}} = \mathbf{M'}\,\dot{\mathbf{q}}_{\mathbf{b}} + \mathbf{M'}\,\dot{\mathbf{q}}_{\mathbf{ht}} + \mathbf{M'}\,\dot{\mathbf{q}}_{\mathbf{bp}} + \mathbf{M'}\,\dot{\mathbf{q}}_{\mathbf{sh}} \ .$$

LATERAL COEFFICIENTS

Υ'_v

- 1. Forces along the y-axis proportional to accelerations along the same axis are characterized by the added mass term Y'. . This term is the lateral counterpart of $Z'_{\dot{\dot{v}}}$ in the longitudinal equations.
- 2. Just as Z' was seen to be a major stability derivative in the longitudinal plane, ${}^{\dot{W}}{}^{\dot{Y}}$ ' is a very important coefficient in the lateral characteristic equation. Trajectory simulation and y-force percent contribution plots in Figures 26 through 29 illustrate its significance to vehicular motion during various maneuvers. The root locus plots in Figures 30, 31, and 32 illustrate the wide variation in the locations of the poles due to variation in Y' . Tables 6A through 6D present the results of a sensitivity analysis of nondimensional roots of the characteristic equation (transfer function denominator) and v, ψ , and ϕ (numerators) for a generalized underwater vehicle, varying Y' from -100 percent to +100 percent.
 - 3. & 4. For axisymmetric bodies

$$Y'_b = Z'_{\dot{w}_b}$$
.

Most underwater vehicles have nearly axisymmetric hull forms. Of those whose shapes are not a body of revolution, quite often good approximations can be made by treating the body as everywhere circular in cross-section with a cross-sectional area equal to the actual body. This approximation is often used in the analysis of underwater vehicles.

Appendages that contribute to Y', are those with significant projections onto the xz-plane. Usually these consist of vertical tail fins

(Text Continued on Page 43)

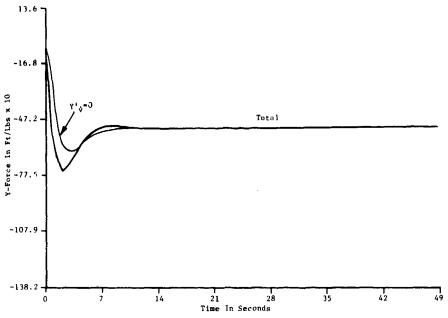


FIGURE 26. TRAJECTORY SIMULATION PLOT OF Y' CONTRIBUTION TO Y-FORCE DURING A ϵ_{μ} = 30° STEADY TURN

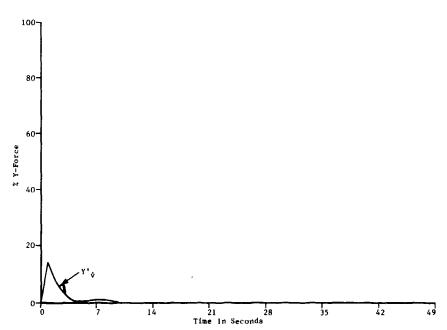


FIGURE 27. PERCENT CONTRIBUTION OF Y' $_{\hat{\mathbf{v}}}$ TO Y-FORCE DURING A $_{6_{_{\mathbf{P}}}}$ = 30° STEADY TURN

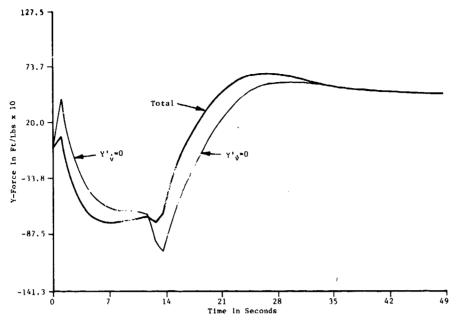


FIGURE 28. TRAJECTORY SIMULATION PLOT OF Y' CONTRIBUTION TO Y-FORCE DURING A $\varepsilon_{\mu}=\pm30^{\circ}$ Turn reversal maneuver

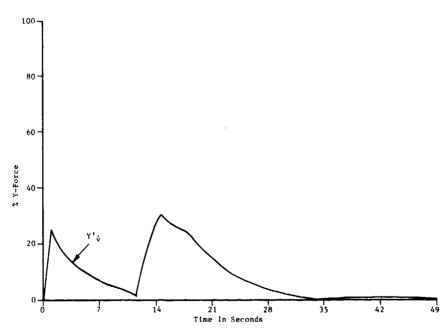


FIGURE 29. PERCENT CONTRIBUTION OF Y' $_{\tilde{V}}$ TO Y-FORCE DURING A $\delta_{_{\rm P}}$ = ±30° TURN REVERSAL MANEUVER

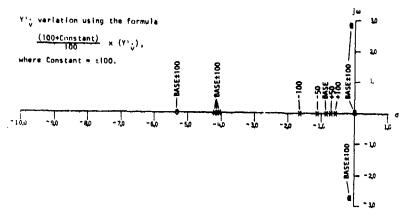


FIGURE 30. ROOT LOCUS PLOT OF $v/\delta_{_{\bf P}}$ FUR A GENERALIZED UNDERWATER VEHICLE, VARYING Y' $_{\dot{\bf V}}$

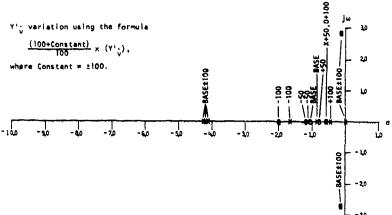


FIGURE 31. ROOT LOCUS PLOT OF ψ/ϕ_{μ} FOR A GENERALIZED UNDERWATER VEHICLE, VARYING Y'

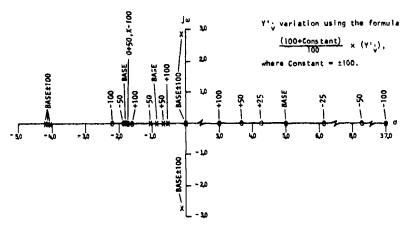


FIGURE 32. ROOT LOCUS PLOT OF ϕ/δ_{μ} FOR A GENERALIZED UNDERNATER VEHICLE, VARYING Y',

LATERAL SENSITIVITY AMALYSIS OF V MUMERATOR MOM-DIMENSIONAL ROOTS FOR A GENERALIZED UNDERMATER YEHICLE, VARYING $V_{\gamma}^{+}\pm 100x$ LASE -5.3392 - P +505+ 1352+ 35 152-+1001 *50 - 18478 .66 - 18478 .66 - 4.1839 - 51 - 5828 - 19.4 | 100 | -1853 | 1.1 - 18550 | 1.1 - 1755 | -51 - 5770 - 32.5 | -51 | -5770 | -32.5 | -50 - 18130 -1.2 - 18130 -1.2 - 4.2218 .59 -1.11177 31.6 -.08 -.80828 - 4.6 -251 -18263 - .51 - .18265 - .53 -4.2072 .24 - .96285 12.7 -2.8382 - .03 2.8882 - .03 -101 - 18323 - .19 - .18323 - .19 -4.2007 .09 - .88956 5.1 +28 - 18426 .39 - 18426 .38 -4.1695 - 18 - 75618 -10.7 . . 84717 LATERAL SENSITIVITY ANALYSIS OF CHARACTERISTIC EQUATION NON-DIMENSIONAL ROOTS FOR A GENERALIZED UNDERMATER VEHICLE, VARYING Y'; ±100\$ -4.1970 +101 - 2.8596 .01 - 2.8596 .01 0 2.8592 MSE - 18357

* A fifth set of real and imaginary roots exists, which has a base \$100% variation value of sero

LATERAL SENSITIVITY ANALYSIS OF ϕ numerator nom-dimensional roots for a generalized undermater vehicle, varying $Y^*_{\psi} \pm 100\%$

TABLE 60 LATERAL SENSITIVITY AWALYSIS OF ⊕ NÜMERATOR MON-DIMENSIONAL ROOTS FOR A GENERALIZED UNDERBAIER VEHICLE, VANYING Y'∵±100% -92 R 22.926 -88 I 19.363 -84 I 16.805

> 4.1 6.1751 23.6 0 1.6 5.4013 8.1

-50 -1.9349 0 -25 -1.8462

-101 -1.8010

MSC^I -1.7735

9.1 8.2590

-2.1989 24.0 36.992

100 ×

S CHANGE	8.38	32.7	14.0	5.2		- 4.7	-11.0	-19.8	-33.0
R00T 3	-2.0114	-1.3551	-1.1648	-1.0743	-1.0214	97344	0.80939	1 13618	68425
S CHANGE	-1.5	33	12	10		.03	:05 :03	.08	60.
R00T 2	2.8603	2,1637	2.8650	18516 2.8656	2.8660	2.8663	2.8688	2.8675	2,8685
S CHAIGE	-1.5	53	12	01		. 01	.05	.05	90:
X VAR, MOOT 1	-1001 - 18251 -2.8603	-501 -: 18431 -2.8637	-251 18501 -3.8610	-101 - 18516 -101 -2.8656	MSE 18523	+10 18528	+25 - 18533 -2.8668	+50 -18537	+100 - 18536 -2.8685

 $\frac{100 \frac{1}{4} - \frac{27.875}{0}}{0} - \frac{-11.8}{0} \frac{2.8777}{0} - \frac{40.4}{0}$ 4. A third set of real and imaginary roots exists, which has a base 100% variation value of serc.

-60 R 9.6334

- 3.5 4.2320 - 15.3

+25 -1.7123 +50 -1.8598

4.6548

+101 -1.7478

3.6930 - 26.

-80 I 11.676 -70 I 17.654

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(or stabilizers), rudders, canards, and shrouds. The contributions of these and any other surfaces that can be treated as flut plates in accelerated fluid flow normal to its surface can be calculated using Equation (18). It is repeated here for convenience.

$$m_{a} = \frac{k\pi\rho c^{2}b}{4} \tag{18}$$

Since it is common to have unequal upper and lower vertical fins in the tail, the two will not be combined as were the two horizontal tail fins. The expressions for the upper and lower vertical tail fins are written as

$$Y'_{\text{uvt(lvt)}} = -\frac{k\pi^{-2}b}{2k^3} .$$

If the rudder is an extension of the vertical tail fins, then it should be included in the area for calculation of c. If the rudder is separate, then it can be treated separately using the same expression.

Since the bowplanes have no significant vertical surface area, they do not contribute to Y'. The shroud has the same projection in the vertical plane as in the horizontal plane, so

$$Y'_{\dot{v}_{s}} = Z'_{\dot{w}_{sh}} .$$

The complete expression of Y' $_{\mathring{\mathbf{V}}}$ for a typical underwater vehicle is written

$$Y' = Y' + Y' + Y' + Y' + Y'$$

$$\dot{v}_b \quad \dot{v}_{uvt} \quad \dot{v}_{lvt} \quad \dot{v}_{sh} \quad .$$

Υ',

- 1. y-forces on a submerged vehicle that are due to angular acceleration in roll are represented in the equations of motion as $Y'_{\dot{D}}$.
- 2. This term can be significant on vehicles with large Z_{cb}/ℓ , Z_{uct}/ℓ , and Z_{lct}/ℓ ratios, where
 - Z_{cb} = z-distance from body cg to cb
 - Z_{uct} = z-distance from the x-axis to the centroid
 of the upper vertical tail, and
 - Z_{lct} = z-distance from the x-axis to the centroid
 of the lower vertical tail.

Tables 7A through 7D present the results of a sensitivity analysis of nondimensional roots of the characteristic equation (transfer function denominator) and v, ψ , and ϕ (numerators) for a generalized underwater vehicle, varying Y', from -100 percent to +100 percent.

- 3. Due to intermodal coupling, Y'. = K'. as shown in Equation (8). Since the conceptual description of the physical phenomenon causing K'. moments is more easily understood than that for Y'., only the equality relationship is given here; that is
 - 4. $Y^{\dagger}_{\dot{D}} = K^{\dagger}_{\dot{V}}$ (See below).

Y';

- 1. When an angular acceleration in yaw results in a force in the y-direction, the y-force is represented in the equations of motion by the term Y'. The longitudinal counterpart of Y'. is Z'.
- 2. Tables 8A through 8D present the results of a sensitivity analysis of nondimensional roots of the characteristic equation (transfer function denominator) and v, ψ , and ϕ (numerators) for a generalized underwater vehicle, varying Y', from -100 percent to +100 percent.
- 3. Intermodal coupling yields the relationships in Equation (8). That is,
 - 4. $Y'_{\dot{r}} = N'_{\dot{v}}$.

Κ'_v

- 1. Rolling moment due to acceleration along the y-axis is termed $\textbf{k'}_{\dot{\textbf{v}}}$.
- 2. K' couples the roll and yaw motions of a vehicle. As a vehicle enters a turn, the changes in the drift angle, β , will appear as a change in the sway velocity, v; i.e., an acceleration v. The rolling moment produced by v in this situation can be significant on a vehicle with unsymmetrical vertical fin arrangements, like a submarine's sail. Trajectory simulation and roll moment percent contribution plots in Figures 33 and 34 illustrate its significance to vehicular motion during a turn reversal maneuver. The root locus plot in Figure 35 illustrates the variation in the locations of the poles due to variation in K'. Tables 9A through 9D present the results of a sensitivity analysis of nondimensional roots of the characteristic equation (transfer function denominator) and v, ψ , and ϕ (numerators) for a generalized underwater vehicle, varying K' from -100 percent to +100 percent.

(Text Continued on Page 50)

LATERAL SENSITIVITY MAALYSIS OF 4 MUMERATOR WOM-DIMENSIONAL ROOTS FOR A GENERALIZED UNDERMATER YEHICLE, VARYING $V_{\hat{p}}^{*}$ ±100% +56 - .17459 -1.7 - .17459 -1.7 -5.3205 -.35 +1007 - .17155 -3.3 - .17165 -3.3 -5.3020 -.70 +251 - 12607 - . 84 - . 17607 - . 84 - 5.3299 - . 17 1.7 -.18034 1.7 -£.3592 .36 .06 2.8720 .06 0 .84 -.17905 .84 -5.34±7 .18 .03 .03 2.8711 .33 - 17815 .33 -5.3430 .07 .01 .01 - .34 - .17696 - .34 -5.3355 - .07 - .01 2.8698 - .01 -5.3392 \$ CHANCE ROOT 2 \$ CHUNCE 3.4 3.4 3.4 3.4 3.4 3.4 3.4 2.8702 -50 -- 18054 +10] -.17638 -251 -: 17905 -2: 3711 -101 - 17815 MSE^R - 17756 A fifth set of real and imaginary roots exists, which has a base 1100% variation value of zero *1001 - 18540 1.0 - 18540 1.0 - 4.1957 - 0.05 - 84769 .06 +50 - 18449 .50 - 18443 .50 - 4.1964 - 01 - 184743 .03 .002 - .84711 - .007 . . 34730 . . 02 .02 - .84690 - .03 .007 - .04704 - .02 - 34717 -.002 - :4722 LATERAL SENSITIVITY ANALYSIS OF CHARACTERISTIC EQUATION NOW-DIMENSIONAL ROOTS FOR A GENERALIZED UNDERMATER YEHICLE, VARYING Y $_p^{\prime}$ ±100% -.007 -4.1970 25 -4.1967 -50 - 18265 - .50 - .19265 - .50 -4.1977 - .2.9627 .12 0 -.25 -.18311 -.25 -4.1973 .06 2.8610 .06 0 .02 -4.1971 -4.1969 S VAR ROOT 1 S CHANGE ROOT 2 S CHANGE -100 - 18172 -1.0 - 18172 -1.0 - 18172 -1.0 . 25 - . 18403 - . 06 2. 6575 2.8592 ..10 -.18375 -..02 2.8585 - 10 - 18338 .02 2.8599 -251 -: 18311 -101 -2.8538 +101 - 18375 +25 -18403 MSC -: 18357

LATERAL SENSITYITY ANALYSIS OF # NUMERATOR NON-DIMENSIONAL ROOTS

S VAR ROOT 1 S CHANGE ROOT 2 S CHANGE ROOT 3 S CHANGE

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 ROOT 3
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 -.27

 **207
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 -.28
 -.1633
 -.1631
 -.03
 -.07

 **207
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 -.28
 -.1621
 -.03
 -.03
 -.03

 **107
 -.1852
 -.08
 -.1853
 -.09
 -.1021
 -.01

 **107
 -.1852
 -.09
 -.1852
 -.09
 -1.021
 -.01

 **107
 -.1852
 -.03
 -.1852
 -.03
 -1.021
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 **107
 -.1852
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 -.1021
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 -.1865</

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LATERAL SENSITIVITY ANALYSIS OF ϕ numerator non-dimensional roots for a generalized undermater vehicle, warting Y'_{D} =100%

LATERAL SENSITIVITY ANALYSIS OF V NUMERATOR NON-DIMENSIONAL ROOTS FOR A GENERALIZED UNDERNATER VEHICLE, VARYING Y'; ±100% LATERAL SENSITIVITY AMALYSIS OF & NUMERATOR NON-DIMENSICYAL F FOR A GENERALIZED UNDERMATER YEHICLE, VARYING Y'; ±100% -5c1 -17762 .03 -17762 .03 -5.2327 -2.0 .02 -.17759 .02 -5,2854 -1.9 0 2,3702 0 C .006 -.17757 .006 -5.3176 -.40 0 2.8702 0 0 +100 -: 12744 -: 07 -: 17744 -: 97 -: 5.5659 4.2 -: 100 -: 1003 0.003 +50 -.17750 -.03 -.17750 -.03 -5.4502 2.1 -5.3392 0 -50, -1.7711 -.14 5.0877 1.8 100 2 CHWE ROOT 1 S CHWE ROOT 2 S CHWEE 255 -1.7723 -- 07 8.0418 -10t -1.7730 -.03 5.0143 2.8702 -255 -1.7746 +101 -1-7739 MSE -1.7735 -251 -2.8702 -101 -2.8702 +101 -. 17755 A fifth set of real and imaginary roots exists, which has a base :100% variation value of nero *100T = .18365 .05 = .18385 .05 = -4.2050 .19 = .84586 = .15 -.01 -.18355 -.01 -4.1950 -.05 -.84749 ..04 0 2.8592 0 0 .05 -.84684 -.04 -.03 -.18352 -.03 -4,1830 -.10 -.84782 .03 .003 2.8583 .003 .02 -.84704 -.02 -. 84730 . . 02 LATERAL SENSITIVITY ANALYSIS OF CHARACTERISTIC EQUATION NON-DIMENSIONAL RODIS FOR A GEMERALIZED UNDERMATER VEHICLE, WARYING Y'; ±100% LATERAL SENSITIVITY ANALYSIS OF 4 NUMERATOR NON-DIMENSIONAL ROOTS FOR A GEMERALIZED UNDERMATER VEHICLE, VARYING Y'; ±100% -.005 -.18356 -.005 -4.1962 0 2.8592 0 0 +5c - 38561 . 02 - 18361 . 02 - 4.2010 -4.1970 -4.1990 9261.9--. 18357 2.8592 MSE -19523 101-1,52+ - 1 5 5 7 -50 -. 18352 +25<mark>8 - 18359</mark> +101 -2.8582 -251 -. 18355 -251 -2.8592 -101 -. 18356 12 MSE -: 18357

1001 -1.2781 .26 4.8224 -3.5

100

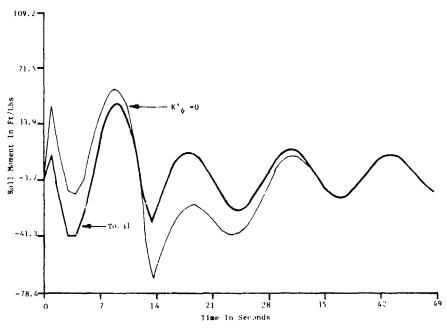


FIGURE 33. TRAJECTORY SIMULATION PLOT OF K1. CONTRIBUTION TO ROLL MOMENT DURING A δ_{μ} = ±30° TURN REVERSAL MANEUVER

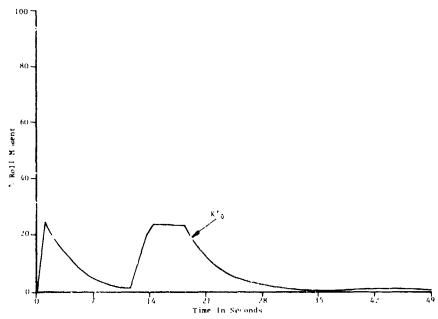
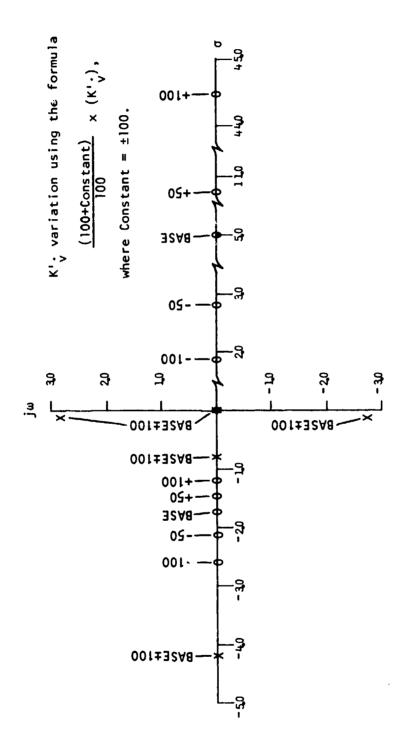


FIGURE 34. PERCENT CONTRIBUTION OF K' TO ROLL MOMENT DURING A δ_{μ} = :30° TURN REVERSAL MANEUVER



ROOT LOCUS PLOT OF $\phi/\delta_{\mathbf{r}}$ FOR A GENERALIZED UNDERWATER VEHICLE, VARYING K'. FIGURE 35.

LATFIAL SENSITIVITY ANALYSIS OF V NUMERATOR NON-DIMENSIONAL ROOTS FOR A GENERALIZED UNDERNATER VEHICLE, VARYING F'; =100% MSE - 17756 s. -251 喜 +25.F TOS. -. 84703 -. 02 -. 84692 -- 04 -- 95648 -- 08 7.1 = 18557 1.1 -4.1865 -01 - 84751 .04 2.8603 .04 0 -- 64717 - 34785 LATERAL SENSITIVITY ANALYSIS OF CHAKACTERISTIC EQUATION MON-DIMENSIONAL ROOTS FOR A GENERALIZED UNDERWATER VEHICLE, VARYING K'_y ±100\$ -50 -18757 2.2 -18757 2.2 -1.1860 -.02 *1001 -12562 -4.5 -17562 -4.3 -6.1990 -15 - 1.35 - 6.1990 - 1.5 - 1.35 --4,1972 +25k --18158 -1.1 --18158 -1.1 +501 -17959 -2.2 -17959 -2.2 3 YAR RGJT 1 S CHANGE 1-100 4.4 -25^t = 1952 -107 -2.8536 +10 -2,8588 MSE - 18357

LATERAL SENSITIVITY ANALYSIS OF & HUMERATOR NON-DIMENSIONAL ROOTS
FOR A GENERALIZED UNDERNATER YEHICLE, VARTING K.; =100X
ROOT 1 % CHANGE ROOT 2 % CHANGE
22.5587 44.3 1.8986 --22.0
0
0
22.2734 22.5 2.8880 -42.2

+70 R 16.211

11.2 3.7265

+84 K 22.028 +88 K 27.028

4.9963

7.0373

4135 -20.3 10.654

1 TOOM	S CHAPCE	2 100	T CHANGE	R007 3	X CHANGE
13836	1.7	18836	1,1	-1.0220	90.
-2.8715	ē.	2.8715	ŝ.	Q	
- 18673	16.	. 18679	189	-1.0217	.03
	60.	2.8687	69:	0	Ц
10801	57	18807	27	-7 0275	10
-2.8674	.05		.05	0	
-10	-11/2	- 18554 - 0665	- 17	-1.0215	
		2000			
18523	_	. 18523	_	-1.0214	_
	_	2.5660	_	0	_
18491	1	18431	11	1.0213	10:-
+101 -2.8654	32	2.8654	20	0	
18445	1.25	18445	20:-	-1.0212	20:-
-251 -2.8646	05	2.3646	05	0	
79281 - 18367	84	- 18367	84	-1.0211	- 93
	11.	2.9632	10	0	
21281 -	-1.7	18212	-1.7	-1.3208	JE
2032 6 1001	0	2030 6	104 -	¢	

• A chird set of real and imaginary roots exists, units has a base 1905 pario

+92 K 37.245

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A Afrik set of real and imaginary roots exists, which has a base 1100% variation value of ser

LATERAL SENSITIVITY ANALYSIS OF ψ numerator non-dimensional roots for a generalized undermater vehicle, warting $\kappa^{}_{i,j} \pm 100x$

NCSL TR-327-78

3. & 4. This coefficient is evaluated in a similar fashion to that of M'. . It is the representation of added mass effects acting on moment arms from the cg.

$$K'\dot{v}_b = -Y'\dot{v}_b \frac{Z_{cb}}{\ell} .$$

As noted before, it is not uncommon for the upper and lower vertical tail fins to be of different size and shape. Also, since the cg is generally below the body axis of symmetry, there are unequal moment arms to the centroids of the upper and lower fins even if the fins are symmetric about the center line. The vertical tail surfaces' contribution to K', is written

$$K'\dot{v}_{t} = -\frac{1}{\ell} \{Y'\dot{v}_{uvt} Z_{uct} + Y'\dot{v}_{\ell vt} Z_{ct}\}$$

$$= K'\dot{v}_{uvt} + K'\dot{v}_{\ell vt}.$$

 $Z_{uct} = distance from the x-axis to the centroid of the upper vertical tail$

 Z_{lct} = distance from the x-axis to the centroid of the lower vertical tail.

Note: The vertical tail fins are assumed to lie in the xz-plane.

The complete expression, then, is written

$$K'_{\dot{\mathbf{v}}} = K'_{\dot{\mathbf{v}}_{\mathbf{b}}} + K'_{\dot{\mathbf{v}}_{\mathbf{t}}}.$$

Any other vertical surfaces that can be approximated as flat plates lying in the xz-plane can be treated similarly to the upper and lower vertical tail surfaces.

Κ',

1. Moments additional to the I_X moment of inertia resisting angular acceleration in roll are included in the term K'. This added moment of inertia is one that has no Lamb coefficient (1) for the body contribution. This is because a homogeneous body of revolution can rotate about its axis of symmetry and not experience any forces from an inviscid fluid.

⁽¹⁾ibid.

- 2. Tables 10A through 10D present the results of a sensitivity analysis of nondimensional roots of the characteristic equation (transfer function denominator) and v, ψ , and ϕ (numerators) for a generalized underwater vehicle, varying K' \dot{p} from -100 percent to +100 percent.
- 3. Since \dot{p} , angular acceleration in roll, is defined about the x-axis passing through the cg, there will be some added moment effect of the body due to an added mass effect acting at the cb. It is written as

$$K'_{\dot{p}_b} = Y'_{\dot{v}_b} \frac{z^2_{cb}}{\ell^2}.$$

Any surface projecting radially from the vehicle (such as tail fins, bowplanes, etc.) will contribute to the added moment effect represented by K'. Each surface contributes in two ways. One is as an added mass acting on a moment arm from the x-axis. The general form for this is the same as the body contribution above. Considering any "flat plate" surface, such as the upper vertical tail fin, this term would be written

$$K'_{p_{uvt}} = Y'_{uvt} \frac{z^2_{uct}}{\ell^2}.$$

The other form of contribution to K', is by the individual added moment of inertia of the flat plate rotating about an axis parallel to the x-axis but passing through the centroid of the flat plate. Again, we turn to Malvestuto and $\operatorname{Gale}^{(5)}$ for experimental data on flat plates with varying dihedral angle and taper ratio. For rotation about an axis in the plane of the plate, parallel to the chord, and passing through the centroid, the equation is

$$I_{x_{a}} = -\frac{\pi \rho}{48} D_{\lambda} D_{\Gamma} k_{p}^{\dagger} c^{2} b^{3}$$

where k' is the coefficient of added moment of inertia for a flat plate plotted in Figure 25, D is the correction factor for taper ratio plotted in Figure 36, and D is the correction factor for dihedral angle plotted in Figure 37. Thus, for any "flat plate" surface such as the upper vertical tail fin

$$K'_{\dot{p}_{uvt}} = \frac{1}{\frac{1}{2\rho} \ell^5} I_{x_a} + Y'_{\dot{v}} \frac{z^2_{uct}}{\ell^2}.$$

⁽b) ibid.

LATERAL SENSITIVITY ANALYSIS OF ϕ NUMERATOR NON-DIMENSIONAL ROUTS FOR A GENERALIZED UNDERMATER VEHICLE, VARYING $K^{'}_{\rm p}$ $\pm 100\%$ LATERAL SENSITIVITY AMALYSIS OF V NUMERATOR NOM-DIMENSIONAL ROOTS FOR A GENERALIZED UNDERMATER VEHICLE, VARYING $K^{*}_{
m p}$ $\pm 100\%$ +1001 - 17579 -1.0 - 17579 -1.0 -5.5396 .007 +50k - 17667 - .50 - .17667 - .50 -5.3394 .004 ROOT 3 % CHANGE -5.3389 -.006 -5.3393 .51 - .17846 .51 .26 2.3776 .26 +10R - 17738 - 10 - 17738 - 10 +25₁ - 17711 - .25 - .17711 - .25 - .15 - .13 1.0 - 17936 1.0 52 2.3851 .52 R00T 2 -101 - 17774 . 10 - 17774 - 101 - 17774 -25k - 17801 .25 - 17801 - 25 - 28739 2.8702 ROOT 1 S CHANGE MSE^K -1.7735 -122-+25 1001 -100 -So. -50 - 17846 ASCI - 17756 VAR 1 ROOT 1 . A fifth set of real and imaginary roots exists, which has a base 1100% variation value of sero. 100 -- 912 LATERAL SENSITIVITY ANALYSIS OF CHARACTERISTIC EQUATION KON-DIMENSIONAL ROOTS FOR A GENERALIZED UNDERWATER VEHICLE, VARING K' $_{\rm F}$ $_{\rm 5}$ 1005 LATERAL SENSITIVITY AMALYSIS OF ψ NUMERATOR NON-DIMENSIONAL ROOTS FOR A GENERALIZED UNDERMATER VEHICLE, VARVING $K'_{\hat{p}} \pm 100\Sigma$ ROOT 3 S CHANGE -4.1971 .002 -4.1970 ROOT 2 % CHANGE -.18716 1.0 2.8608 .52 +10d -:18167 -1.0 -:18167 -1.0 -4.1969 -50 - 18619 .52 - 18619 .52 - 50 - 26 - 26 - 26 - 26 - 26 - 26 -25k - .18571 .26 - .18571 .26 - .13 2.5697 .13 +50| -.18427 -.52 -.18427 -.52 +50| .10 - 18542 .10 +100 - 18333 -1.0 - 18333 -1.0 .10 - .18376 .10 .05 2.8607 .05 - . 18523 2.8860 +10[- 16503 - 11 - 18503 +10[-2.8645 - .05 2.8645 TABLE 10A .26 - .18405 .26 -.26 -.18309 -.26 -.13 2.8556 -.13 +50 - 18261 - .52 - .18261 - .52 -2.8519 - .26 2.8519 - .26 ROOT 2 % CHANGE 2.8592 -.10 -.18338 -.05 2.3578 1 VAR ROOT 1 1 CHANGE -100 1 - 18716 1.0 -101 - : 18542 -2: 8675 +251 - . 18475 -2.8623 DASE - 18523 -108 - 18376 -2.8607 +10k - 18338 +25k - . 18309 -50 -2.8666 -251 - 18405 -251 -2.8629 MSET - 18357

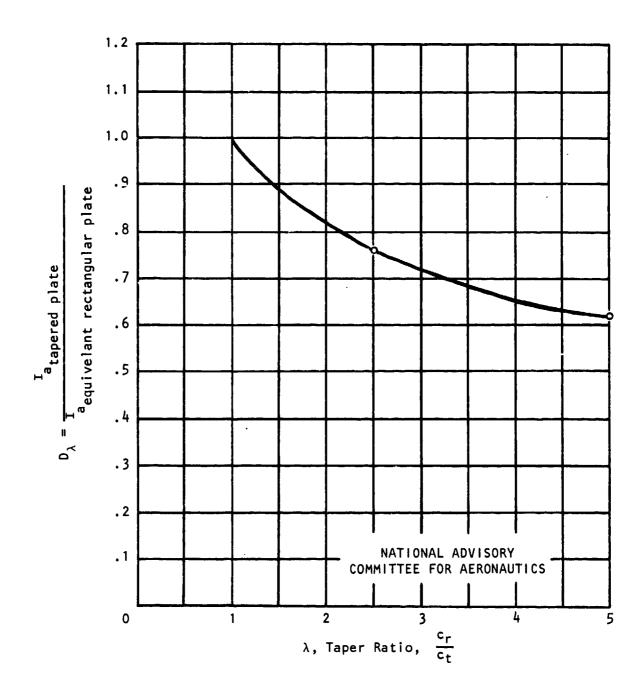


FIGURE 36. DEPENDENCE OF THE ADDITIONAL MOMENT OF INERTIA ON TAPER RATIO

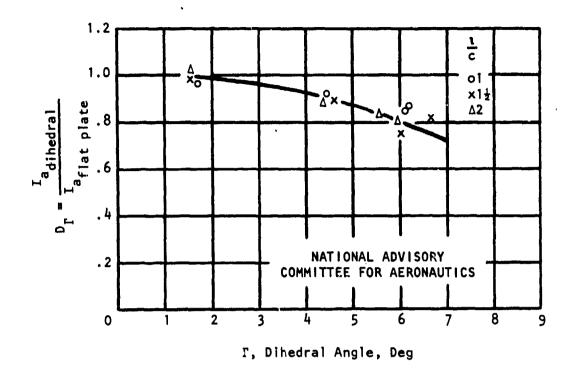


FIGURE 37. VARIATION OF THE ADDITIONAL MOMENTS OF INERTIA OF A SINGLE PLATE WITH DIHEDRAL ANGLE

The first term is often negligible compared to the second. The complete expression for a common vehicle would be

$$K'_{\dot{p}} = K'_{\dot{p}_{\dot{b}}} + K'_{\dot{p}_{uvt}} + K'_{\dot{p}_{lvt}}$$
.

 $\frac{\texttt{K'}_{\dot{\textbf{r}}}}{}$

- 1. Rolling moments caused by angular accelerations in yaw are represented by K' $_{\dot{\tau}}$.
- 2. Tables 11A through 11D present the results of a sensitivity analysis of nondimensional roots of the characteristic equation (transfer function denominator) and v, ψ , and ϕ (numerators) for a generalized underwater vehicle, varying K' \dot{r} from -100 percent to +100 percent.
- 3. K' arises in the same manner as the K' v term except that the linear acceleration normal to the vertical surfaces is the result of an angular acceleration in yaw acting about the z-axis.

(Text Continued on Page 56)

LATERAL SENSITIVITY ANALYSIS OF CHARACTERISTIC EQUATION NON-DIMENSIONAL ROOTS FOR A GENERALIZED UNDERMATER VEHICLE, VARYING K'; ±100°

LATERAL SENSITIVITY ANALYSIS OF V NUMERATOR NON-DIMENSIONAL ROOTS FOR A GENERALIZED UNDERWATER VEHICLE, VARYING K'; =100%

2 CHANGE										
*CO01 *	4					g G	-			+
ROOF 3 T CHANGE	.03	1	20.	800	700	-	200:	002	- 01	03
2001	-5.3409		-5.3401	-5.3397	-5.3394	-5.3392 0	-5.3321 2	-5.3388	-5.3384	-5.3376
KOOT 2 THANKS	.19	.003	27	03	27		20	04	0	15
1001 z	-, 17781	2.8703	2, 8702	2.8702	- 12758 2.8702	2,8702	2,8702	2.8702	2.8702	2.8701
A CHARGE	2.24	.003	.37 0	202	19.		0 02	Š 0	20	
F001	17783	-2.8703	17769	177 62 -2. 8702	-2.8702	-2.8702 -2.8702	-2.8702	- 17749	-2.0702	-2,4701
Z VAR	1		35-	12:		- Swa			\$	1001+
	-بر		001							
R007	847.6	0	2.84716		$\frac{1}{1}$	84717		$\frac{1}{1}$	 	
Z CHANGE	01		005	002	0	ب	. 032	200	. 005	0.
R007 3	Γ	0	-4.1968	-4.1969	-4.1970	-4.1970	-4.1971	-4.1971	-4.1972 0	.00 -4.1975 .01
T CHANGE	30	.00.	-:03	20	005		.005	20.	003	
. ROOT 2		2.8594	2,8333	-, 18354	2.8592	2.8592	18358 2.8592	2.6592	2.6591	1001 -18368 -006 -18368
K CHANGE		.007	-:03	02	905		. 005	20.	.003	007
100	101	-2.8594	18351	-, 10354 -2, 8593	18356	- 18357	-,18358	-,18360	18362	-2.8590
VAD	Ī	-100;	-50	1.52-	101-	1354	ŽĮ.	152+	ş	+100.

LATERAL SENSITIVITY ANALYSIS OF ψ MUMERATOR NON-DIMENSIONAL ROOTS FOR A GENERALIZED UNDERMATER YEHICLE, WARYING K^+_{τ} ±100%

LATERAL SENSITIVITY ANALYSIS OF 4 NUMERATOR NON-DIMENSICHAL ROOTS FOR A GENERALIZED UNDERWATER VEHICLE, VARYING K'₇ ±100% +255 -1.7201 --19 4.3211 -1.5 -101 -1.7746 -07 5.0271 +101 -2.7721 -.09 4.9659 -501 -1.7803 .39 5.1545 \$ VAR RD0T 1 \$ CHANGE -100T -2.7874 .78 -251 -1.7769 $\operatorname{MSE}_{\mathrm{I}}^{\mathrm{K}} = 1.7735$

1001 -1.7602 -.75 4.7036 -5.7

ASE -18523 152 101 105 1,52 喜 1001

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$$K'_{\dot{r}} = K'_{\dot{v}_{uvt}} \frac{(X_{nuvt} - X_{neg})}{\ell} + K'_{\dot{v}_{lvt}} \frac{(X_{n\ell vt} - X_{neg})}{\ell}$$

where

X
nuvt = distance from nose to upper vertical
tail surface centroid, positive aft

X_{nlvt} = distance from nose to lower vertical
tail surface centroid, positive aft.

N'_v

- l. Linear accelerations in the y-direction causing yawing moments give rise to the stability derivative N'. to describe these resulting moments.
- 2. Tables 12a through 12D present the results of a sensitivity analysis of nondimensional roots of the characteristic equation (transfer function denominator) and v, ψ , and ϕ (numerators) for a generalized underwater vehicle, varying N' $_{\mathring{\mathbf{v}}}$ from -100 percent to +100 percent.
- 3. Primarily due to the vertical tail surfaces, a vehicle changing drift angle will experience a restoring force proportional to $\dot{\mathbf{v}}$ (resulting from $\dot{\boldsymbol{\beta}}$). N', is the term representing this force.

4.
$$N'_{\dot{v}_{b}} = -Y'_{\dot{v}_{b}} \frac{(X_{\text{ncb}} - X_{\text{ncg}})}{\ell}$$

$$N'_{\dot{v}_{b}} = -Y'_{\dot{v}_{\text{nuvt}}} \frac{(X_{\text{nuvt}} - X_{\text{ncg}})}{\ell} - Y'_{\dot{v}_{\text{nuvt}}} \frac{(X_{\text{nkvt}} - X_{\text{ncg}})}{\ell}.$$

N'ė

- 1. Yawing moments caused by angular accelerations in roll are represented in the vehicular equation of motion by N':
- 2. Tables 13A through 13D present the results of a sensitivity analysis of nondimensional roots of the characteristic equation (transfer function denominator) and v, ψ , and ϕ (numerators) for a generalized underwater vehicle, varying N' $_{\mathring{\mathbf{D}}}$ from -100 percent to +100 percent.
 - 3. & 4. Equation (9) shows

+1001+

A fifth set of real and imaginary roots exists, which has a base ±100% variation value of zero

+1001 -18351 .02 -18351 .02 -4.2085

LATERAL SENSITIVITY ANALYSES OF V NUMERATOR NON-DIMENSIONAL ROOTS FOR A GENERALIZED UNDERMATER VEHICLE, VARYING N°, ±100% 900T 3 ROOT 2 S CHANGE MSEI -. 17755 -101-201+ +25 Ş 1₅₂-35 100 -,84523 .13 -. 84695 -. 03 - 81770 .06 -. 84664 -. 06 -.84739 .02 -. 84717 R007 4 LATERAL SENSITIVITY AVALYSIS OF CHARACIERISTIC EQUATION NON-DIMENSIONAL ROOIS FOR A GENERALIZED UNDERWATER VEHICLE, VARYING N'₂ ±100° ROOT 3 S CHANGE -2.8591 -.003 2.8591 -.03 -4.855 -.27 -4.1982 -4.1970 +251 -2,6593 .003 2,8593 .003 0 -50 -18355 -.01 -18355 -.01 -4.1813 - 0 0 0 0 -.005 -,18356 -.005 -4,1941 0 2.8592 0 0 -10 - 18356 - 005 - 18356 - 005 - 41959 - 0 +50 -18359 .01 -18359 .01 -4.2028 .05 -2.8593 .003 0 ROOT 2 S CHANGE 2.8592 -. 18357 2.8592 +101 = 1835 0

-251 -- 18356 -2.8592

-190°

MSE -. 18357 -2.8592

LATERAL SENSITIVITY ANALYSIS OF & NUMERATOR NON-DIMENSIONAL POOTS FOR A GENERALIZED UNDERMATER VEHICL: VARYING N', ±100% +50 -1.7819 .47 4.9753 -.42 100 -1,7305 . 35 4,5543 -.84 NOOT 1 S CHANGE NOOT 2 S CHANGE -21.7567 -.85 5.0284 .81 -50 -1.7651 -- 47 5.0173 *251 -1.7777 . 24 4.9858 -251 -1.7693 -.24 5.0069 -101 -1.7718 -.10 5.0005 4.9921 +101 -1.7752 1,7735 * 100 -*501 = 18525 .01 = 18525 .01 = 1.0255 .42 *1001 - 28527 .02 - 18527 .02 -1.0295 .87 -25<mark>1 - 18522 - .005 - .18522 - .005 - 1.0194 - .20</mark> -10t -2.8660 0 2.6660 0 0 0 LATERAL SENSITIVITY ANALYSIS OF \$\psi numerator non-dimensional roots for a generalized underwater vehicle, varying n*; \didsi -501 -1825; -01 -18521 -01 -1.0173 +251 -2,8660 0 2,8660 0 0 2,8660 1 VAR ROTT 1 CHANGE ROTT 2 5 CHANGE 100T 2 5 CHANGE 100T 2 2.8659 - 02 2.8659 0 2.8659 0 +101 -18523 0 -18523 0 2,8660

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MSE -. 18523

S VAR 1 NOOT 1

LATERAL SENSITIVITY ANALYSIS OF V MUMERATOR NON-DIMENSIONAL RODIS FOR A GENERALIZED UNDERMATER VEHICLE, VARYING M $^{\circ}_{0}$ $^{\circ}$ =100% LATERAL SENSITIVITY ANALYSIS OF 4 NUMERATOR NON-DIMENSIONAL ROOTS FOR A GEMERALIZED UMCERANTER VEHICLE, VARYING N° = 100% 0 2,8702 0 -5.7384 .004 -501 -12759 02 -12759 02 -5.3594 007 -5, 2393 , 002 +501 -2,8703 -.02 -.18752 -.02 -5,8389 -.006 +100 - 12729 - 64 -5.3486 - 01 - 01 - 01 - 01 - 01 - 01 +25, -1,27254 --01 -,17754 --01 +10 -17755 -005 -17755 -0 ROOT 1 S CHANGE MSE, -1.77.75 -101 -25, į -25K - 17,58 -101 -- 17757 -2.8702 \$ 55 -25. 105+ ASE - 17756 -2.8702 \$ P A fifth set of real and imaginary roote exiets, which has a base 1100% variation value of zero 1001 - 15847 - 20 - 1884 - 20 - 10 - 201 -LATERAL SENSITIVITY ANALYSIS OF CHARACTERISTIC EQUATION MON-DIMENSIONAL ROOTS FOR ℓ Generalized underhater vehicle, varying $\mathbf{n}^*_{\mathbf{p}}$ ±1005 LATERAL SENSITIVITY ANALYSIS OF ψ WUMERATOR NON-DIMENSIONAL ROOTS FOR A GENERALIZED UNDERWATER VEHICLE, VARVING N' $_{\rm p}$ =1005 +501 -18383 -.03 -18383 -.03 -4.1868 -.025 25. - 125. - 2. 2. - 125. - 125. - 1.02 - 1.05. - 1.05 -50\ -18382 .03 -18362 .03 -4.1372 -101 -18358 .005 -18358 .005 -4.1271 -4.1970 +10[-2.8582 -005 -18256 -005 -4.1870 0 -251 -18359 01 -18359 07 -4.1821 0 -50 -18520 -.02 -.18520 -.02 -100 -18517 - 03 -18517 - 03 -251 -18521 -01 -18521 -01 0 2,8660 0 2,8660 -101 -18522 -.005 -.18522 -2.8660 0 2.8660 +254 -, 18524 , 005 -, 18524 -2,8660 0 2,8660 -501 - 18525 .01 - 18525 0 2.8660 +10 -18523 0 -18523 -2,8660 0 2,8660 2.8592 ASE -. 18523 -2,8660 LASE -. 18357

N'r

- 1. Yawing moments additional to I are denoted N'. This is the added moment of inertia term whose longitudinal counterpart is M'.
- 2. Trajectory simulation and yaw moment percent contribution plots in Figures 38 and 39 illustrate its significance to vehicular motion during various maneuvers. The root locus plots in Figures 40, 41 and 42 illustrate the wide variation in the locations of the poles due to variation in N'. Tables 14A through 14D present the results of a sensitivity analysis of nondimensional roots of the characteristic equation (transfer function denominator) and v, ψ , and ϕ (numerators) for a generalized underwater vehicle varying N'. from -100 percent to +100 percent.
- 3. & 4. Like M' , there are two parts to the N' body contribution

$$N'_{\dot{r}_{b}} = \frac{-K_{b}' I_{z_{df}}}{\frac{1}{2} \rho \ell^{5}} + Y'_{\dot{v}_{b}} \frac{(X_{ncb} - X_{ncg})^{2}}{\ell^{2}}$$

for an axisymmetric vehicle body I $_z$ = I $_d$. Again, similar to M' $_d$, the upper vertical tail fin contribution is

$$N'_{iuvt} = -\frac{1}{\frac{1}{2\rho \ell^5}} \frac{(K_p^{\dagger} \pi \rho b^2 c^3)}{48} + Y'_{iuvt} \frac{(X_{nuvt} - X_{ncg})^2}{\ell^2}$$

The lower vertical tail fin is identical in form. The complete expression, then, is written $N'\dot{r}=N'\dot{r}+N'\dot{r}+N'\dot{r}$.

SUMMARY

A compilation of methods for analytically predicting the acceleration hydrodynamic coefficients from the geometric and mass distribution characteristics of underwater vehicles has been presented. A theoretical development from potential flow theory is given to show the equivalences among acceleration and coupled terms. The detailed discussion of each coefficient presented includes descriptions intended to provide the reader with an understanding of the physical relationships involved.

(Text Continued on page 63)

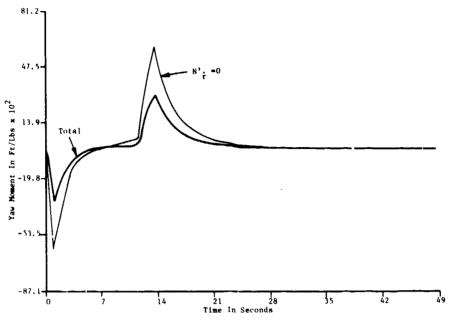


FIGURE 38. TRAJECTORY SIMULATION PLOT OF N' CONTRIBUTION TO YAM MOMENT DURING A 6, = $\pm 30^{\circ}$ TURN REVERSAL MANEUVER

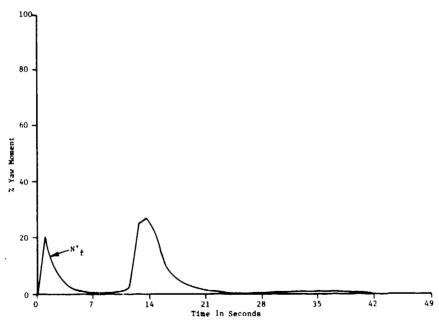


FIGURE 39. PERCENT CONTRIBUTION OF N', TO YAW MOMENT DURING A δ_{μ} = ±30° TURN REVERSAL MANEUVER

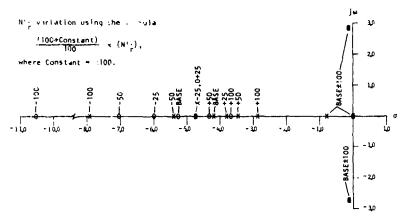


FIGURE 40. ROOT LOCUS PLOT OF v/\circ_p FOR A GENERALIZED UNDERMATER VEHICLE, VARYING N' $_{\hat{r}}$

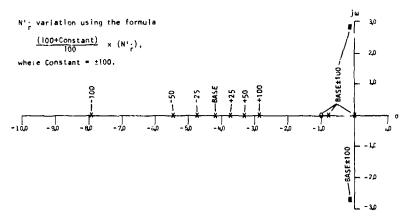


FIGURE 41. ROOT LOCUS PLOT OF $\phi/\phi_{\bf r}$ FOR A GENERALIZED UNDERWATER VEH! VARYING N'.

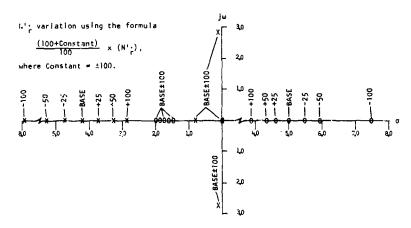


FIGURE 42. ROOT LOCUS PLOT OF ${\it p/6}_{\it p}$ FGR A GENERALIZED UNDERWATER VEHICLE. ARYING N'

100k -1.5917 -10.5 3.9106 -21.7

1001

-25

1,05+

+50 -1.6746 -5.6 4.3665 -12.6

LATERAL SENSITIVITY AMALYSIS OF & MUMERATOR NCX-DIMENSIONAL ROOTS FOR A GENERALIZED UNDERMATER VEHICLE, VARVING N°; ±100% LATERAL SENSITIVITY ANALYSIS OF V NUMERATOR NON-DIMENSIONAL ROOTS FOR A GENERALIZED UNDERWATER VEHICLE, VARYING N°; ±100% +100 -1251 -2.855 -12561 -65 -32.9 -32.9 -32.9 -32.9 +50 -12886 - 39 -12886 - 39 -4.2802 -19.6 -25| -17802 .26 -.17822 .26 -6.0840 13.3 25. - 127.7 - 22 - 10.9 2. 20. - 20 -7.0719 32.4 +10 = 12744 = .29 = .17746 = .39 = .5.0902 = 4.7 -10 -12775 -10 -17775 -10 -5.6141 -5.6141 -5.3392 105- 10845 L.1 SEESS 10.4 S VAR ROOT 1 S CHANGE ROOT 2 S CHANGE -1007 -2.5445 15.5 2.5533 59.2 -101 -1.7956 1.2 5.1526 3.1 -50 -17865 .56 -.17855 .56 -.56 -.56 -.56 4.9983 *25, -1,7217 -2.9 4.6537 +101 -1.7522 -1.2 4.8515 S VAR ROOT 1 S CHANGE ROOT 2 S CHANGE 12.00 -1.17981 1.3 -1.17981 1.3 -1.17981 1.3 -1.17981 1.3 - 17756 5.8702 MSF -1,7735 A fifth set of real and imaginary roots exists, which has a base 1908 variation value of sero. 95. 64.40. --4.0104 - 4.4 - 34656 -.00 +251 - 18405 - 26 - 18405 - 26 - 3.7601 - 10.4 - 15187 - 15 3.5. -56 - 1835 - 183 - 1835 £3. C'158. -11.00 ROOT 3 S CHANGE LATERAL SENSITIVITY ANALYSIS OF ϕ NUMERATOR NON-CINENSIONAL MOOTS FIRE A GENERALIZED UNDERNATER VEHICLE, VARING N $_{7}^{+}$ +130% LATERAL SENSITIVITY ANALYSIS OF CHARACTERISTIC EQUATION NON-DIPENSIONAL ROOTS FOR A GENERALIZED UNDENWATER VEHICLE. VARYING N ; : :1001 <u>6</u> MODT 1 1 CHANGE NOT 2 1 CHANGE ROT 3 1 Ch. "E -1.0] - 18337 - 11 - 18357 - 17 -4.402. 4.5 S CHWEE -4.1970 -25 - 18307 - 187 - 18307 - 177 - 17525 - 1752 TABLE 14C REDUT 1 S CHANGE 1 ROOT 2 . 18169 -1.1 - 18169 -1.1 . 8 05 . 05 2.8605 .05 2.3392 .10 -.18376 MSC - 1.662 100 1 +101 - 18376 DASE - 18357

¥ 8

Results of censitivity analyses are presented to show the relative contributions of each coefficient to the equations of motion. A discussion of the important assumptions involved in mathematically computing each coefficient is given, along with the analytical expression used to compute each coefficient. An example case (Appendix B) is given to show a comparison between computed coefficients and experimental data.

Table 15 presents a summary of the sensitivity analyses conducted for each acceleration coefficient. These results are for ± 100 percent coefficient variation. It can be seen in Table 15 that five coefficients make major contributions to the vehicle dynamics. These are Z', M', Y', K', and l. . Variations of ± 100 percent change both the numerator and denominator roots by better than 50 percent, except for K'. Changes in K' significantly effect only the roll numerator, N $^{\circ}$. It should be noted that the \dot{v} terms, such as K' and Y', do not affect the v-numerator because no v-term appears in the v-numerator. This is easily seen when one applies Cramer's rule to form the transfer functions (9). This same observation can be made about each of the other numerators; i.e. no w-term in the w-numerator.

⁽⁹⁾ Naval Coastal Systems Laboratory Technical keport 287-76, Development of the Equations of Motion and Transfer Functions for Underwater Vehicles, by D. E. Humphreys, July 1976.

TABLE 15

ROOT SENSITIVITY ANALYSES FOR DERIVATIVE VARIATION OF +100% (Sheet 1 of 2)

LONGITUDINAL DERIVATIVE ROOT VARIATION PERCENTAGES

7	CHAR. EC	UATION	U NUME	RATOR	W NUME	RATOR	0 NUME	RATOR
	Derivative -100%	Variation +1002	Derivative -100%	Variation +100%	Derivative -100%	Variation +100%	Derivative -100%	Variation +100%
x, º	0 3.30%	-3.10% 0	0	0 0	0 3.30%	-3,10% 0	3.30 % 0	-3.10% 0
min.	0	0						
	-0.12% 0 91.20% 2.40%	0.12% 0 -32.50% ~0.58%	-0.21% 88.30% 4.00%	0.23% -32.30% -0.98%	0 0 0	0 0 0	0 97.0 0%	0 -33.00\$
maj.	<u> </u>				 			
H'ù	-0.01% 0 0.27% -0.29%	0.01% 0 -0.27% 0.29%	-0.02% 0.43% -0.44%	0.02% -0.42% 0.43%	0 0	0 0 0	0 -0.79%	0 0.79 %
neg.	L		l		<u> </u>		<u></u>	

LONGITUDINAL DERIVATIVE ROOT VARIATION PERCENTAGES

7	CHAR. EC	DATION	U NUME	RATOR	W NUME	RATOR	0 NUME	RATOR
	Derivative -100%	Variation +100%	Derivative -100%	Variation +100%	Derivative -100%	Variation +100%	Derivative -100%	Variation +100%
Z'ą	-0.007% 0 0.18% -0.21%	0.007% 0 -0.18% 0.21%	-0.008% 0.21% -0.23%	0.01% -0.21% 0.23%	0.02% 0 -4.00%	-0.03% 0 4.30%	0	0
Μ'α maj.	-0.52% 0 0.68% 89.10%	0.53% 0 -0.85% -31.80%	-0.68% 1.20% 88.40%	3,71% -1,50% -31,50%	-0.30% 0 97.10%	0.30% 0 -33.10%	0	0

LATERAL DERIVATIVE ROOT VARIATION PERCENTAGES

ſ	CHAR. E	UATION	V NUME	RATOR		Y NUME	RATOR	♦ NUME	RATOR
	Derivative -100%	Variation +100%	Derivative -100%	Variation +100%	Dei	rivative -100%	Variation +100%	Derivative -100%	Variation +100%
۷';	1 -3.90%	1.10%	0	0	*	-1.50% -1.50%	0.07%	24.00% 640.40%	-11.30% -40.40%
ĺ	2.10% 91.50%	-0.51% -32.50%	0	0		96.90%	-33.00%	0	O
maj.		0							
K' ن	++ 4.40%	-4.30% -4.30%	0	0	**	1.70%	-1.70% -1.70%	44.30 2 -62.00 2	-36.40% 789.00%
	-0.05% 0.16%	0.05%	C O	0		0.06%	-0.06%	0	0
maj.	0	0			_				
N';	+++ -0.02% +++ -0.02%	0.02% 0.02%	0	0		-0.02% -0.02%	0.02% 0.02%	-0.95% 0.84%	0.96%
	-0.27%	0.27% -0.25%	0	0		-0.79%	0.80%	0.544	0
neg.	0	0							

TABLE 15

(Sheet 2 of ?)

LATERAL DERIVATIVE ROOT VARIATION PERCENTAGES

1	CHAR. EC	UATION	V NUME	RATOR	W NUME	RATOR	♦ NUME	MATOR
	Derivative -100%	Variation +100%	Derivative -100%	Variation +100%	Derivative -100%	Variation +100%	Derivative -100%	Variation +100%
Y';	† -1.00% † -1.00% 0.03% -0.06%	1.00% 1.00% -0.03% 0.06%	* 3.40% * 3.40% 0.71%	-3.30% -3.30% -0.70%	# -0.91% # -0.91% -0.07%	0.90% 0.90% 0.07%	0 0 0	0 0 0
neg.	0	0						
, p	++ 1.10% ++ 1.10% 0.002% 0	-1.00% -1.00% -0.002% -0.001%	†† 1.00% †† 1.00% -0.006%	-1.00% -1.00% 0.007% 0	†† 1.00% †† 1.00% 0	-1.00% -1.00% 0	0 0	0
neg.	 							
N'è	+++ 0.05% +++ 0.05% 	-0.05% -0.05% -0.01% -0.001%	*** 0.04% *** 0.04% 0.01%	-0.04% -0.04% -0.01% 0	### -0.03% ### -0.03% 0	0.03% 0.03% 0	0 0 0	0 0 0
neg.	0	0	1		1		İ	

IMAGINARY ROOT VARIATIONS: + 0.24 to -0.24; * 0.13 to -0.13; # 0.26 to -0.26; ++ 0.52 to -0.51; +++-0.003 to 0.007; *** -0.007 to 0.003; ### -0.003 to 0.003

LATERAL DERIVATIVE ROOT VARIATION PERCENTAGES

5	CHAR. E	UATION	V NUME	RATOR	Y NUMES	RATOR	♦ NUMER	RA TOR
	Derivative -100%	Variation +100%	Derivative -100%	Variation +100%	Derivative -100%	Variation +100%	Derivative -100%	Variation +100%
Y' ÷	+ -0.05% + -0.05% -0.19%	0.05% 0.05% 0.19%	* 0.06% * 0.06% -3.90%	-0.07% -0.07% 4.20%	0 0	0	-0.27% 3.70% 0	0.26% -3.50% 0
min.	0.15% 0	-0.15% 0	0	· · · · · · · · · · · · · · · · · · ·				
K' ÷	++ -0.06% ++ -0.06% -0.01% -0.001	0.06% 0.06% 0.01%	** 0.14% ** 0.14% 0.03%	-0.15% -0.15% -0.03%	0 0	0 0 0	0.78% 6.60% 0	-0.75% -5.70% 0
min.	0	0						
N';	88.20% 0.53%	0.95% 0.95% -31.60% -0.67%	*** 1.30% *** 1.30% 96.20%	-0.65% -0.65% -32.80%	0 0	0 0 0	15.50% 50.20% 0	-10.30% -21.70% 0
maj.	0	0						

IMAGINARY ROOT VARIATIONS: + 0.003 to 0; * -0.003 to 0.003; ++ 0.007 to -0.007; *** 0.003 to -0.003; ++ 0.007; *** 0.05 to -0.007

APPENDIX A

DERIVATION OF HYDRODYNAMIC COEFFICIENTS RELATIONSHIPS FROM ENERGY EQUATIONS

Beginning with the momentum equations given in Equations (1) and (2) in the body of this report, a derivation is given here for the complete expressions, showing the equalities and relationships between all of the acceleration and nonlinear hydrodynamic coefficients possible for a vehicle moving in an infinite, inviscid, fixed fluid with zero circulation.

A summary of the identities and relationships derived is presented on Page A-6 with simplifications arising from vehicle symmetry on Page A-8.

Note: The hydrodynamic coefficients in parenthesis with a subscript are added together to get the value of the parenthesized hydrodynamic coefficient, e.g.,

$$X_{rq} = (X_{rq})_{1} + (-X_{rq})_{2}$$
 $X_{rq} = (M) + (-N)$
 $X_{rq} = M - N$.

X-FORCE, SURGE

$$-X = \frac{d}{dt} \frac{\delta T_{f}}{\delta u} - r \frac{\delta T_{f}}{\delta v} + q \frac{\delta T_{f}}{\delta w}$$

$$\frac{d}{dt} \frac{\delta T_{f}}{\delta u} = A\dot{u} + B'\dot{w} + C'\dot{v} + L\dot{p} + G\dot{r} - G'\dot{r} + H\dot{q} + H'\dot{q}$$

$$-r \frac{\delta T_{f}}{\delta v} = -Brv - A'rw - C'ru - Mrq - Frr - F'rr - Hrp + H'rp$$

$$q \frac{\delta T_{f}}{\delta w} = Cqw + A'qv + B'qu + Nqr + Fqq - F'qq + Gqp + G'qp$$

$$1etting F_{1} = F + F', F_{2} = F - F', G_{1} = G + G', G_{2} = G - G', H_{1} = H + H',$$

$$H_{2} = H - H'$$

Y-FORCE, SWAY

$$-Y = \frac{d}{dt} \frac{\delta^{T}f}{\delta v} - p \frac{\delta^{T}f}{\delta w} + r \frac{\delta^{T}f}{\delta u}$$

$$\frac{d}{dt} \frac{\delta^{T}f}{\delta v} = B\dot{v} + A'\dot{w} + C'\dot{u} + M\dot{q} + F\dot{r} + F'\dot{r} + H\dot{p} - H'\dot{p}$$

$$-p \frac{\delta^{T}f}{\delta w} = -Cpw - A'pv - B'pu - Npr - Fpq + F'pq - Gpp - G'pp$$

$$r \frac{\delta^{T}f}{\delta u} = Aru + B'rw + C'rv + Lrp - Grr - G'rr + Hrq + H'rq$$

$$Y = -B\dot{v} - A'\dot{w} - C'\dot{u} - M\dot{q} - F_{1}\dot{r} - H_{2}\dot{p} + Cpw + A'pv + B'pw + Npr + F_{2}pq + G_{1}pp - Aru - B'rw - C'rv - Lrp - G_{2}rr - H_{1}rg$$

$$A = -Y_{ru} \quad A' = -Y_{\dot{w}} = Y_{pv} \quad M = -Y_{\dot{q}} \quad F_{1} = -Y_{\dot{r}} \quad F_{2} = Y_{pq}$$

$$B = -Y_{\dot{v}} \quad B' = -Y_{rw} = Y_{pu} \quad N = (Y_{pr})_{1} \quad G_{1} = Y_{pp} \quad G_{2} = -Y_{rr}$$

$$C = Y_{pw} \quad C' = -Y_{\dot{u}} = -Y_{rv} \quad L = (-Y_{rp})_{2} \quad H_{1} = -Y_{rq} \quad H_{2} = -Y_{\dot{p}}$$

Other resulting relations: $Y_{rp} = N - L$.

Z-FORCE, HEAVE

$$-Z = \frac{d}{dt} \frac{\delta T_{f}}{\delta w} - q \frac{\delta T_{f}}{\delta u} + p \frac{\delta T_{f}}{\delta v}$$

$$\frac{d}{dt} \frac{\delta T_{f}}{\delta w} = C\dot{w} + A'\dot{v} + B'\dot{u} + N\dot{r} + F'\dot{q} + G\dot{p} + G'\dot{p}$$

$$- q \frac{\delta T_{f}}{\delta u} = -Aqu - B'qw - C'qv - Lqp - Gqr + G'qr - Hqq - H'qq$$

$$p \frac{\delta T_{f}}{\delta v} = Bpv + A'pw + C'pu + Mpq + Fpr + F'pr + Hpp - H'pp$$

$$Z = -C\dot{w} - A'\dot{v} - B'\dot{u} - N\dot{r} - F_{2}\dot{q} - G_{1}\dot{p} + Aqu + B'qw + C'qv$$

$$+ Lqp + G_{2}qr + H_{1}qq - Bpv - A'pw - C'pu - Mpq - F_{1}pr - H_{2}pp$$

$$A = Z_{qu} \qquad A' = -Z_{\dot{v}} = -Z_{pw} \qquad L = (Z_{qp})_{1} \qquad F_{1} = -Z_{pr} \qquad F_{2} = -Z_{\dot{q}}$$

$$B = -Z_{pv} \qquad B' = Z_{qw} = -Z_{\dot{u}} \qquad M = (-Z_{pq})_{2} \qquad G_{1} = -Z_{\dot{p}} \qquad G_{2} = -Z_{qr}$$

$$C = -Z_{\dot{w}} \qquad C' = Z_{qv} = -Z_{pu} \qquad N = -Z_{\dot{r}} \qquad H_{1} = Z_{qq} \qquad H_{2} = -Z_{pp}$$

$$Other resulting relations: Z_{pq} = L - M.$$

K-MOMENT, ROLL

$$K = -\frac{d}{dt} \frac{\delta T_f}{\delta p} + w \frac{\delta T_f}{\delta v} - v \frac{\delta T_f}{\delta w} + r \frac{\delta T_f}{\delta q} - q \frac{\delta T_f}{\delta r}$$

$$-\frac{d}{dt} \frac{\delta T_f}{\delta p} = -P\dot{p} - Q'\dot{r} - R'\dot{q} - L\dot{u} - G_1\dot{w} - H_2\dot{v}$$

$$w \frac{\delta T_f}{\delta v} = Bwv + A'ww + C'wu + Mwq + F_1wr + H_2wp$$

$$-v \frac{\delta T_f}{\delta w} = Cvw - A'vv - B'vu - Nvr - F_2vq - G_1vp$$

$$r \frac{\delta T_f}{\delta q} = Qrq + P'rr + R'rp + Mrv + F_2rw + H_1ru$$

$$-q \frac{\delta T_f}{\delta r} = Rqr - P'qq - Q'qp - Nqw - F_1qv - G_2qu$$

$$A-3$$

$$A' = K_{ww} = -K_{vv} \qquad L = -K_{\dot{u}} \qquad F_1 = (K_{wr})_1 = (-K_{qv})_1$$

$$B = (K'_{wv})_1 \qquad B' = -K_{vu} \qquad M = (K_{wq})_1 = (K_{rv})_1 \qquad G_1 = K_{vp} = -K_{\dot{w}}$$

$$C = (-K_{vw})_2 \qquad C' = K_{wu} \qquad N = (-K_{vr})_2 = (-K_{qw})_2 \qquad H_1 = K_{ru}$$

$$F_2 = (-K_{vq})_2 = (K_{rw})_2 \qquad P = -K_{\dot{p}} \qquad P' = K_{rr} = -K_{qq}$$

$$G_2 = -K_{qu} \qquad Q = (K_{rq})_1 \qquad Q' = -K_{\dot{r}} = -K_{qp}$$

$$H_2 = -K_{\dot{v}} = K_{wp} \qquad R = (-K_{qr})_2 \qquad R' \qquad -K_{\dot{q}} = K_{rp}$$

Other resulting relations: $K_{wv} = B - C$ $K_{vr} = M - N$ $K_{qw} = M - N$ $K_{rw} = F_1 + F_2 = 2F$ $K_{vq} = -F_1 - F_2 = -2F$ $K_{qr} = Q - R$

M-MOMENT, PITCH

$$M = \frac{d}{dt} \frac{\delta T_f}{\delta q} + u \frac{\delta T_f}{\delta w} - w \frac{\delta T_f}{\delta u} + p \frac{\delta T_f}{\delta r} - r \frac{\delta T_f}{\delta p}$$

$$- \frac{d}{dt} \frac{\delta T_f}{\delta q} = -Q\dot{q} - P'\dot{r} - R'\dot{p} - M\dot{v} - F_2\dot{w} - H_1\dot{u}$$

$$u \frac{\delta T_f}{\delta w} = Cuw + A'uv + B'uu + Nur + F_2uq + G_1up$$

$$- w \frac{\delta T_f}{\delta u} = -Awu - B'ww - C'wv - Lwp - G_2wr - H_1wq$$

$$\begin{array}{c} \frac{\delta T_{f}}{\delta r} = Rpr + P'pq + Q'pp + Npw + F_{1}pv + G_{2}pu \\ \\ -r \frac{\delta T_{f}}{\delta p} = - Prp - Q'rr - R'rq - Lru - G_{1}rw - H_{2}rv \\ \\ A = (-M_{wu})_{1} \quad A' = M_{uv} \quad L = (-M_{wp})_{1} = (-M_{ru})_{1} \quad F_{1} = M_{pv} \\ \\ B' = -M_{ww} = M_{uu} \quad M = -M_{\tilde{v}} \quad G_{1} = (M_{up})_{1} \\ \\ = (-M_{rw})_{1} \\ \\ C = (M_{uw})_{2} \quad C' = -M_{wv} \quad N = (M_{pw})_{2} = (M_{ur})_{2} \quad H_{1} = -M_{\tilde{u}} \\ \\ = -M_{wq} \\ \\ F_{2} = -M_{\tilde{w}} = M_{uq} \quad P = (-M_{rp})_{1} \quad P' = -M_{\tilde{r}} = M_{pq} \\ \\ G_{2} = (-M_{wr})_{2} = (M_{pu})_{2} \quad Q = -M_{\tilde{q}} \quad Q' = -M_{rr} = M_{pp} \\ \\ H_{2} = -M_{rv} \quad R = (M_{pr})_{2} \quad R' = -M_{\tilde{p}} = -M_{rq} \\ \\ Other resulting relations: \quad M_{wv} = C - A \\ \\ M_{pw} = N - L \\ \\ M_{ru} = N - L \\ \\ M_{pu} = G_{1} + G_{2} = 2G \\ \\ M_{rp} = G_{1} - G_{2} = -2G \\ \\ M_{rp} = R - P \end{array}$$

N-MOMENT, YAW

$$N = -\frac{d}{dt} \frac{\delta T_f}{\delta r} + v \frac{\delta T_f}{\delta u} - u \frac{\delta T_f}{\delta v} + q \frac{\delta T_f}{\delta p} - p \frac{\delta T_f}{\delta q}$$

$$-\frac{d}{dt} \frac{\delta^{T}f}{\delta r} = -R\hat{r} - P'\hat{q} - Q'\hat{p} - N\hat{w} - F_{1}\hat{v} - G_{2}\hat{u}$$

$$v \frac{\delta^{T}f}{\delta u} = Avu + B'vw + C'vv + Lvp + G_{2}vr + H_{1}vq$$

$$-u \frac{\delta^{T}f}{\delta v} = -Buv - A'uw - C'uu - Muq - F_{1}ur - H_{2}up$$

$$q \frac{\delta^{T}f}{\delta p} = Pqp + Q'qr + P'qq + Lqu + G_{1}qw + H_{2}qv$$

$$-p \frac{\delta^{T}f}{\delta q} = -Qpq - P'pr - R'pp - Mpv - F_{2}pw - H_{1}pu$$

$$A = (N_{vu})_{1} \qquad A' = N_{uw} \qquad L = (N_{vp})_{1} = (N_{qu})_{1} \qquad F_{1} = -N_{v}$$

$$= -N_{ur}$$

$$B = (-N_{uv})_{2} \qquad B' = N_{vw} \qquad M = (-N_{uq})_{2} = (-N_{pv})_{2} \qquad G_{1} = N_{qw}$$

$$C' = N_{vv} = -N_{uu} \qquad N = -N_{w} \qquad H_{1} = (N_{vq})_{1}$$

$$F_{2} = -N_{pw} \qquad P = (N_{qp})_{1} \qquad P' = -N_{q} = -N_{p}$$

$$H_{2} = (-N_{up})_{2} = (N_{qv})_{2} \qquad R = -N_{r} \qquad R' = N_{qq} = -N_{pp}$$

$$Other Resulting relations: N_{vu} = A - B$$

$$N_{vp} = L - M$$

$$N_{qu} = L - M$$

$$N_{vq} = H_1 + H_2 = 2H$$
 $N_{pu} = -H_1 - H_2 = -2H$
 $N_{pq} = P - Q$

$$A = -X_{\dot{u}} = -Y_{ru} = Z_{qu} = (-M_{wu})_1 = (N_{vu})_1$$

$$B = -Y_{\dot{v}} = X_{rv} = -Z_{pv} = (K_{wv})_1 = (-N_{uv})_2$$

$$C = -Z_{\dot{w}} = -X_{q\dot{w}} = Y_{p\dot{w}} = (-K_{v\dot{w}})_2 = (M_{u\dot{w}})_2$$

$$A' = -Z_{\dot{v}} = Y_{\dot{w}} = X_{rw} = -X_{qv} = Y_{pv} = -Z_{pw} = X_{ww} = -X_{vv} = M_{uv} = N_{uw}$$

$$B' = -X_{\dot{\mathbf{w}}} = -Z_{\dot{\mathbf{u}}} = -X_{\mathbf{q}\mathbf{u}} = -Y_{\mathbf{r}\mathbf{w}} = Y_{\mathbf{p}\mathbf{u}} = Z_{\mathbf{q}\mathbf{w}} = -K_{\mathbf{v}\mathbf{u}} = -M_{\mathbf{w}\mathbf{w}} = M_{\mathbf{u}\mathbf{u}} = N_{\mathbf{v}\mathbf{w}}$$

$$C^{\dagger} = -X_{\hat{\mathbf{v}}} = -Y_{\hat{\mathbf{u}}} = X_{\mathbf{r}\mathbf{u}} = -Y_{\mathbf{r}\mathbf{v}} = Z_{\mathbf{q}\mathbf{v}} = -Z_{\mathbf{p}\mathbf{u}} = X_{\mathbf{w}\mathbf{u}} = -M_{\mathbf{w}\mathbf{v}} = N_{\mathbf{v}\mathbf{v}} = -N_{\mathbf{u}\mathbf{u}}$$

$$L = -X_{\dot{p}} = -K_{\dot{u}} = (Z_{qp})_{1} = (-Y_{rp})_{1} = (-M_{wp})_{1} = (-M_{ru})_{1} = (N_{vp})_{1} = (N_{qu})_{1}$$

$$M = -Y_{\dot{q}} = -M_{\dot{v}} = (X_{rq})_{\dot{1}} = (-Z_{pq})_{\dot{2}} = (K_{wq})_{\dot{1}} = (K_{rv})_{\dot{1}} = (-N_{uq})_{\dot{2}} = (-N_{pv})_{\dot{1}}$$

$$N = -Z_{r} = -N_{w} = (-X_{qr})_{2} = (Y_{rp})_{2} = (-K_{vr})_{2} = (-K_{qw})_{2} = (M_{pw})_{2} = (M_{ur})_{2}$$

$$F_1 = Y_{\dot{r}} = -N_{\dot{v}} = X_{rr} = -Z_{pr} = (K_{wr})_1 = (K_{qv})_1 = M_{pv} = -N_{ur}$$

$$G_1 = -Z_{\dot{p}} = -K_{\dot{w}} = -X_{qp} = Y_{pp} = K_{vp} = (M_{up})_1 = (-M_{rw})_1 = N_{qw}$$

$$H_1 = -X_{\dot{q}} = -M_{\dot{u}} = -Y_{rq} = Z_{qq} = K_{ru} = -M_{wq} = (N_{vq})_1 = (-N_{pu})_1$$

$$F_2 = -Z_{\dot{q}} = -M_{\dot{w}} = -X_{qq} = Y_{pq} = (K_{vq})_2 = (K_{rw})_2 = M_{uq} = -N_{pw}$$

$$G_2 = -X_{\dot{r}} = -N_{\dot{u}} = -Y_{rr} = -Z_{qr} = -K_{qu} = (-M_{wr})_2 = (M_{pu})_2 = N_{vr}$$

$$H_2 = -Y_{\dot{p}} = -K_{\dot{v}} = X_{rp} = -Z_{pp} = K_{wp} = -M_{rv} = (-N_{up})_2 = (N_{qv})_2$$

$$P = -K_{\dot{p}} = (-M_{rp})_{1} = (N_{qp})_{1}$$

$$Q = -M_{\dot{q}} = (K_{rq})_{1} = (-N_{pq})_{2}$$

$$R = -N_{\dot{r}} = (-K_{qr})_{2} = (M_{pr})_{2}$$

$$P' = -M_{\dot{r}} = -N_{\dot{q}} = K_{rr} = -K_{qq} = M_{pq} = -N_{pr}$$

$$Q' = -K_{\dot{r}} = -N_{\dot{p}} = -K_{qp} = -M_{rr} = M_{pp} = N_{qr}$$

$$R' = -K_{\dot{q}} = -M_{\dot{p}} = K_{rp} = -M_{rq} = N_{qq} = -N_{pp}$$

When xz-plane symmetry exists, then the following coefficients are zero; A', C', P', R', L, M, N, G, G'.

Table Al gives a summary of force and moment coefficients for a body with xz-plane of symmetry. When xy-plane symmetry exists, then the following coefficients are zero; A', B', P', Q', L, M, N, H, H'.

When the yz-plane symmetry exists, then the following coefficients are zero; $B^{:}$, $C^{:}$, $Q^{:}$, $R^{:}$, L, M, N, F, $F^{:}$.

When two planes of symmetry exist, all of the terms that are zero in the single plane cases are zero.

TABLE Al (Sheet 1 of 2)

POTENTIAL FLOW FORCE AND MOMENT HYDRODYNAMIC COEFFICIENTS FOR BODY WITH XZ PLANE SYMMETRY

	X	Y	Z	K	<u> </u>	<u> </u>
ů	x _å	Ø	X.	Ø	x •	Ø
ů	ø×	Y v	Ø	ģ	Ø	Y
ŵ	X.	Ø	Z.	Ø	z	Ø
p	Ø	Y. P	Ø	K,	Ø	ĸ
ģ	х _ф	Ø	Z ġ	Ø	M. q	Ø
ř	Ø	Y _r	Ø	K _{ř}	Ø	N _ř
u^2	**			_	-x ₩	ø
uv		~		X.	Ø	Y, - X _ů
uw			******	Ø	X _ů - Z _ŵ	ø
up		-X ₩	Ø		Ø	X + Y p
uq	X,		-X _ů	ø	−Z ġ	Ø
ur	Ø	X _ů		-x _å	Ø	Ϋ́r̀
v^2		_		Ø		Ø
VW				$z_{\dot{\mathbf{w}}} - y_{\dot{\mathbf{v}}}$	Ø	−X _ŵ
v p		Ø	${\tt Y}_{\bf \mathring{v}}$	ø	-Y _r	Ø
νq	Ø	-	Ø	Y + Z		-(x + y)
vr	`-Y *	Ø		ø :	Ϋ́ṕ	Ø
w ²				Ø	X.	ø
wp		−Z _ŵ	Ø	-Y	Ø	z ġ
wq	Z.		-X	Ø	x.	ø
Wr	Ø	X,	_	$-(Y_{\dot{r}} + Z_{\dot{q}})$	Ø	
				A_Q		

TABLE Al (Sheet 2 of 2)

	<u>x</u>	<u> </u>	2	K	<u> </u>	N
p^2		Ø	Y. P		-K _r	Ø
pq	Ø	−z _•	Ø	K.	Ø	M K.
pr	-Y p	Ø	Y	Ø	$K_{\dot{p}} - N_{\dot{r}}$	Ø
q^2	Z †		-X	Ø	_	Ø
qr	Ø	x,	Ø	N _r - M _q	Ø	-K
r^2	-Y	Ø		Ø	ĸţ	

^{*}Coefficient zero because of xz symmetry

^{**}Coefficient not present for arbitrary body with no planes of symmetry

APPENDIX B

COMPARISON OF ANALYTICAL PREDICTIONS WITH EXPERIMENTAL DATA

In this appendix hydrodynamic coefficients of an operational U.S. Navy submersible, obtained at the David Taylor Naval Ship Research and Development Center are compared to corresponding data calculated by the methods set forth in the body of this report. Table Bl lists those 15 acceleration coefficients generally considered to be important, their experimental and calculated values, and the percent difference between the two. Percent difference in this table was obtained as follows:

Experimental - Calculated Experimental

As can be seen in Table Bl, percent difference is significant in only four of the 15 cases, that is, those above 12 percent.

Three of those six coefficients $(M_{\mathring{\mathbf{w}}},\ Z_{\mathring{\mathbf{q}}},\ \text{and}\ K_{\mathring{\mathbf{p}}})$ have been shown in sensitivity analyses illustrated in the body of this report to be of relatively minor significance.

Some questions exist concerning the validity of using the PMM experimental technique to obtain roll derivatives, and specifically, in this case, the fourth problem coefficient listed in table, K. .

It should be noted that not all of the coefficients obtained by DTNSRDC were measured. Some by necessity were calculated. Those marked with an asterisk in Table Bl were calculated by some DTNSRDC method and not by NCSL methods.

TABLE B1

COMPARISON OF CALCULATED VERSUS EXPERIMENTALLY-DETERMINED COEFFICIENTS OF AN OPERATIONAL SUBMERSIBLE

Coef.	Experiment	Calculation	% Error
Mù	16800E-03	60300E-04	64.0
Ζŵ	10110E-01	10100E-01	0.1
Mė	4990≎≣-03	54880E-03	10.0
Z,	17100E-03	60300E-04	65.0
K.*	33000E-05	21500E-06	93.0
Yċ	10880E-01	100G1E-01	8.0
x.*	20900E-03	18500E-03	11.5
K;	51000E-04	28500E-04	44.0
N÷	48400E-03	53570E-03	10.7
х *	10110E-01	10100E-01	0.1
X _{vr}	.10880E-01	.10001E-01	8.0
Ywp	.10110E-01	.10100E-01	0.1
Npq	49600E-03	54900E-03	10.7
Z *	10880E-01	10001E-01	8.1
M _{rp} *	.48070E-03	.53550E-03	11.4

APPENDIX C

DESCRIPTION OF TYPICAL HYDRODYNAMIC VEHICLE USED FOR SENSITIVITY ANALYSIS

This appendix contains a physical description of the underwater vehicle used as the subject of this report. Table C1 is a listing of pertinent geometric facts describing the vehicle, and Tables C2 through C6 list calculated derivatives, moments and symmetry terms. Figure C1 is a set of orthographic projections illustrating the vehicle.

TABLE C1

SUBJECT UNDERWATER VEHICLE GEOMETRY

BARE HULL Vehicle Total Length, Ft.23534E 02 TAIL STRUCTURE Overall Exposed Span, Root to Tip, of Fins (each of 4), Ft. .13200E 01 Horizontal Tail Vertical Tail

Total Exposed Planform Area (each of 2 fins), Ft² . .14510E 01

TABLE C2

S N A M E NON-DIMENSIONAL LONGITUDINAL STABILITY DERIVATIVES

XU	=21408E-02	ZU	0	MU	0
XW	0	ZW	=32464E-01	MW	25379E-02
ΧQ	15275E-03	ZQ	=15275E-01	MQ	=73062E-02
хтн	0	ZTH	0	нтн	=29769E-03
DUX	=62385E-03	ZUD	0	MUD	- .0
XWD	0	ZWD	18182E-01	MWD	=14340E-03
арх	0	ZQD	=14340E-03	MQD	=83506E-03
xx	0	ZX	0	MX	0
XZ	0	ZZ	™ .0	MZ	- .0
XDELT	=57867E-03	ZDELT	- 57867E-02	MDELT	28238E-02
	M = .19043	E-01	IY = .	93512E-	03

TABLE C3
NONLINEAR X-Z SYMMETRY TERMS

XUU	=10704E-02	ZUU	0	MUU =0
XWW	=11127E 00	ZWW	0	O. = WWM
XWQ	=18182E-01			
XQQ	=14340E-03			
xvv	=11127E 00	ZVV	0	MVV = .0
		ZVP	=18182E-01	MVP =14350E-03
XVR	18182E-01	ZVR	0	MVR = .77258E-04
XRP	=77258E-04	ZRP	= .14350E-03	MRP = .83473E-03
XRR	=14350E-03	ZRR	0	MRR =60977E-06

TABLE C4

S N A M E NON-DIMENSIONAL LATERAL STABILITY DERIVATIVES

YV	= -	.33046F-01	ΚV		15537E-03	NV	=22682E-02
ΥP	= -	·.17294E-03	KP	=	13896E-04	NP	= .80153E-04
YR	=	.15280E-01	KR	•	.53435E-04	NR	=73083E-02
YPHI	*	.0	KPHI	3	29769E-03	NPHI	0
YPSI	-	.0	KPSI	*	• .0	NPSI	= .0
YVD	-	.18182E-01	KVD		77258E-04	NVD	= .14350E-03
YPD	=	.77258E-04	KPD	=	37211E-06	NPD	= .60977E-06
YRD	=	.14350E-03	KRD		60977E-06	NRD	=83511E-03
YY	м	.0	KY		.0	NY	0
YDELT	=	.13285E-01	KDELT	u	56451E-04	NDELT	=64829E-02

TABLE C5
MASS MOMENTS OF INERTIA

М	==	. 19043E-01	IXZ	-	.0
ΙX	-	.35537E-04	12	•	.93512E-03

TABLE C6
NONLINEAR X-Z SYMMETRY TERMS

YUU	-	.0	KUU == .0	NUU = .0
YWP	=	.18182E-01	KWP =77258E-04	NWP =14340E-03
			KQR =45882E-07	
			KVW = .21335E-06	
YPQ	-	.14340E-03		NPQ =83469E-03

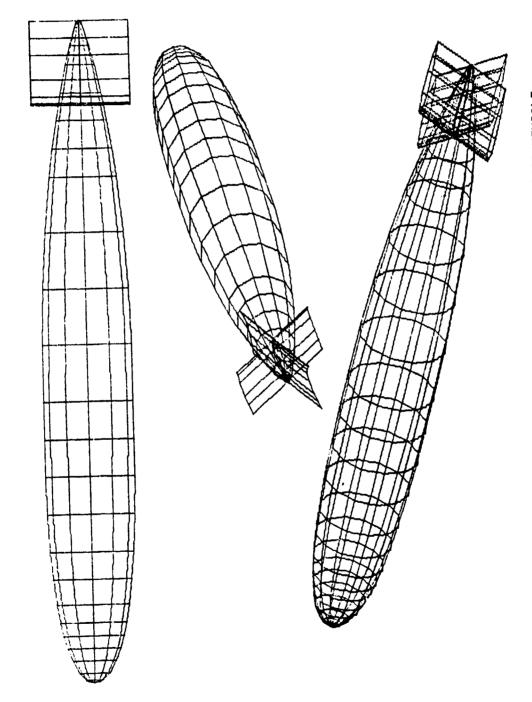


FIGURE C1. ORTHOGRAPHIC PROJECTIONS OF SUBJECT UNDERWATER YEHICLE

APPENDIX D .

DESCRIPTION OF MANEUVERS USED IN TRAJECTORY SIMULATION

Figures D1 through D4 are time history plots of the four basic maneuvers used with the subject underwater vehicle to generate trajectory simulation and percent contribution plots of the various hydrodynamic coefficients discussed in the body of this report. These maneuvers were performed at a simulated vehicle speed of 14.35 feet per second.

The following definitions are provided to assist the reader in the interpretation of Figures D1 through D4.

DELS = Stern plane deflection, in degrees.

PITCH = Pitch angle, in degrees.

Z = Depth, in feet.

ALPHA = Angle of attack, in degrees.

DELR = Rudder deflection, in degrees.

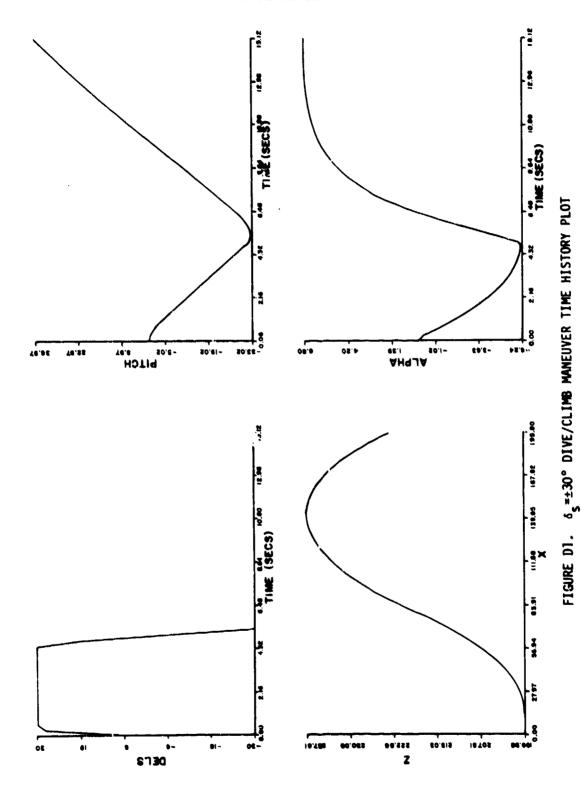
BETA = Sideslip angle, in degrees.

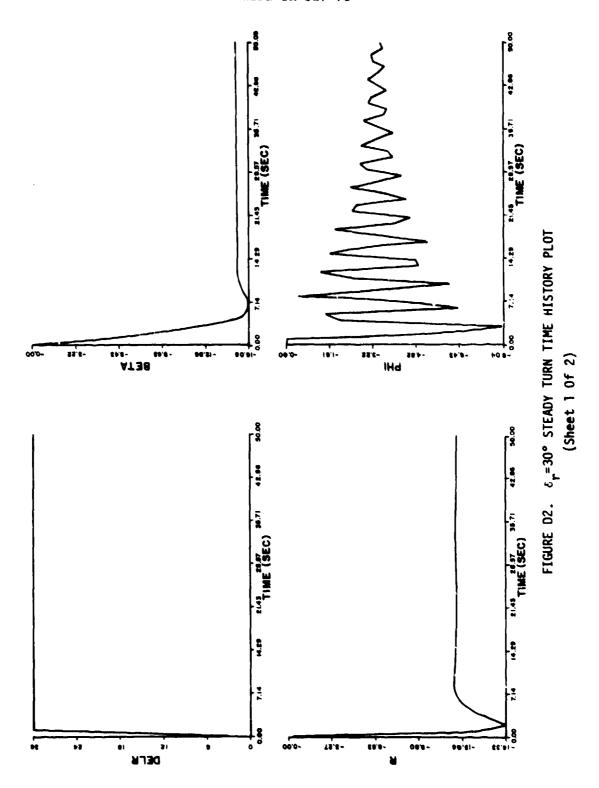
R = Yaw rate, in degrees per second.

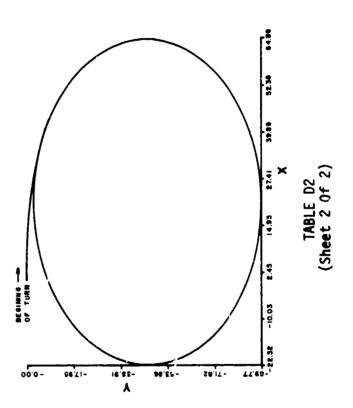
PHI = Roll angle, in degrees.

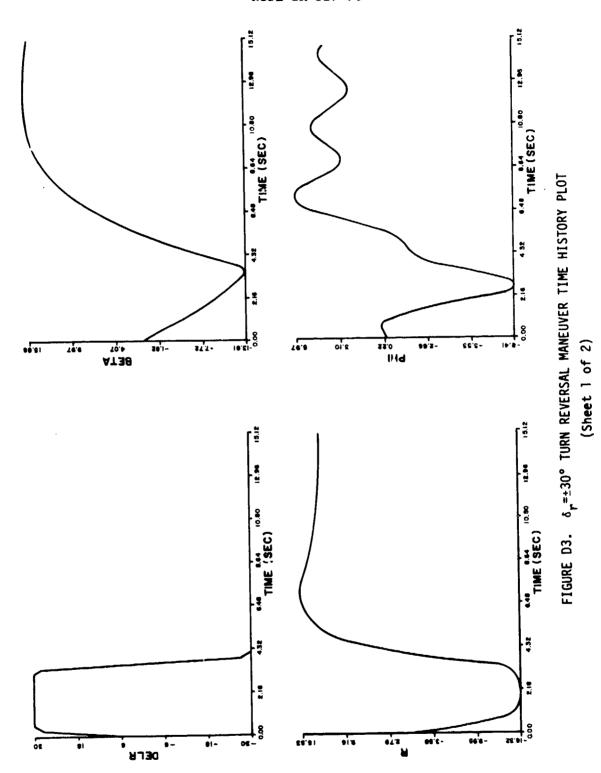
Y = Off-track distance, in feet.

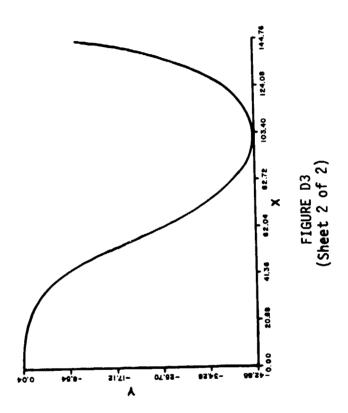
X = Along-track distance, in feet.

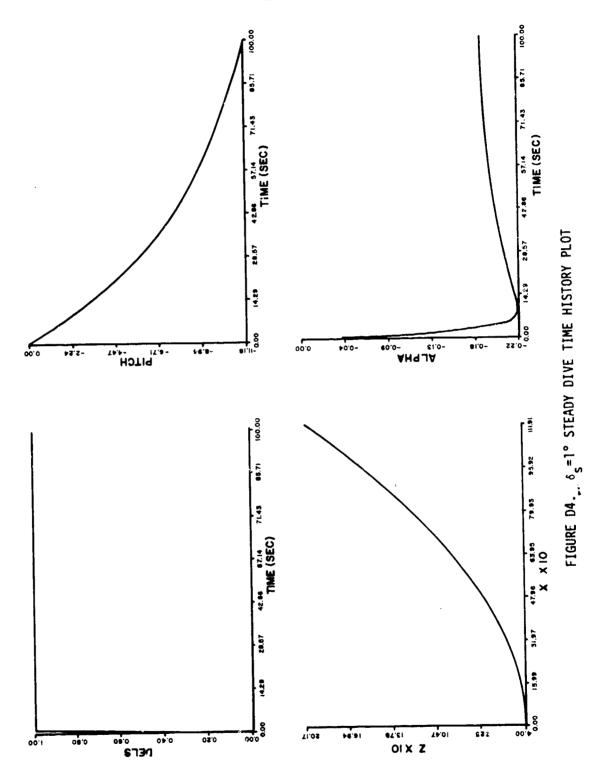












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