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A FLUX CIRCUIT ANALYSIS FOR THE
MAGNETIC TRANSDUCER OF A FLUIDIC
REED GENERATOR

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Harry Diamond Laboratories

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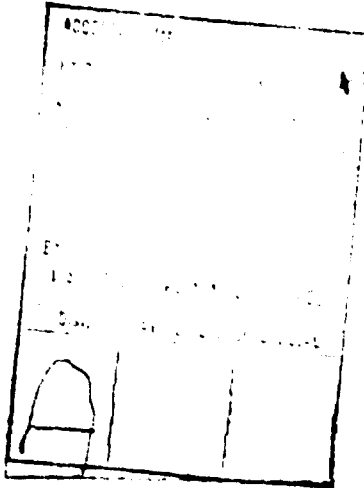
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20. Abstract (cont'd)

Experimental data verify this conclusion, as agreement with the calculated value of power is better than 1%. Losses due to eddy currents and hysteresis are found to be much less serious amounting to no more than a few percent. Recommendations of possible remedies for the saturation problem are made. A computer study is made to determine the effect of variation of the relevant design parameters on the peak voltage output of the generator. The resulting design matrices provide a useful guide to design optimization as well as a clear delineation between the parameter combinations which result in saturation of the reed material and those which do not. Recommendations are made for design changes to improve the generator performance.

10

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CONTENTS

	<u>Page</u>
INTRODUCTION	1
THE MAGNETIC CIRCUIT EQUATIONS	2
THE PERMEANCE CALCULATIONS	4
CALCULATIONS OF GAP FLUX DENSITIES IN THE ABSENCE OF A REED	17
SAMPLE CALCULATION FOR POWER OUTPUT IN THE LOAD RESISTANCE R_L	19
EXPERIMENTAL CHECK OF CALCULATIONS	26
LOSSES IN THE REED	27
LOSSES IN THE KEEPER	28
SUMMARY AND RECOMMENDATIONS	29
APPENDIX I	32
APPENDIX II	35

TABLES

I. Fixed Dimensions of Generator Configuration and Values of Critical Constants	7
II. Voltage Outputs for Various Gap Lengths, Reed Thicknesses and Reed Displacement Amplitudes	41-49

FIGURES

1. Schematic View of Fluidic Generator with Reed Type Magnetic Transducer	1
2. Schematic Representation of the Magnetic Circuit and its Analogue	2
3. Fluidic Generator	5
4. Sectional View of Fluidic Generator in Projectile Ogive	6
5. Calculation of the Relevant Permeances	8
6. Calculation of the Permeance P_{M4}	9
7. Determination of the Width y_F in the Formula for Permeance P_{L3}	14

	<u>Page</u>
8. Determination of the Inner Flux Paths Comprising Permeance P_{M4}	18
9. Determination of B for a Fluidic Generator with a .070' Gap Length	20
10. Comparison of Experimental Points with Theoretical EMF-Displacement Curve	27
11. Inductance of Sensing Coil as a Function of Reed Thickness	31
12. Derivation of the Eddy Current Loss Formula	33

A FLUX CIRCUIT ANALYSIS FOR THE MAGNETIC TRANSDUCER OF A FLUIDIC REED GENERATOR

INTRODUCTION

There is considerable interest in the utilization of fluidic generators in various fuzing applications.¹⁻³ These devices are mounted in the ogives of various projectiles and they essentially convert some of the energy of the relative motion of projectile and air first into acoustical energy and thence to electrical energy. A typical device is pictured schematically in Fig. 1.

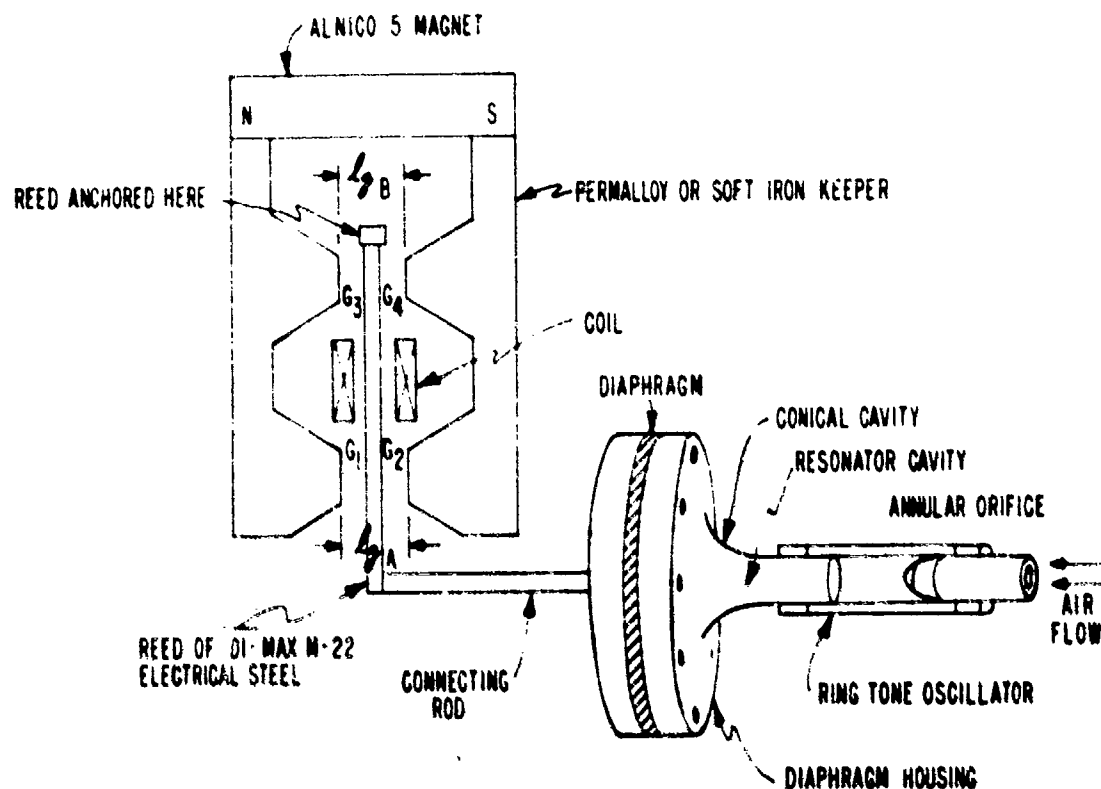


Fig. 1. Schematic View of Fluidic Generator with Reed Type Magnetic Transducer

1. Carl J. Campagnoulo, "The Fluidic Generator," NDL Technical Report-1328, September 1966.
2. Carl J. Campagnoulo, "Fluidic Power Generators for Ordnance Application," NDL Technical Report-1423, December 1968.
3. Carl J. Campagnoulo and Richard N. Gottron, "The Fluidic Generator: A New Electrical Power Source," 24th Annual Proceedings Power Sources Symposium, May 19-21, 1970.

As the projectile moves forward, air is forced through the annular orifice to excite the resonant cavity beyond. The oscillations of the cavity in turn cause the diaphragm to vibrate. This vibration is then transmitted to a permeable reed of Armco electrical steel. As can be seen from Fig. 1, when the reed is in its equilibrium position it is exactly midway between the pole pieces of both magnetic gaps and, therefore, carries no magnetic flux. The vibration of the reed, however, changes the relative lengths of gaps 1 and 2 and flux passes through the reed, alternating direction with the frequency of vibration. Gaps 3 and 4 are taken to remain constant and equal. The alternating flux through the reed induces an ac voltage in the coil around it and causes a current to flow in a circuit with a load resistance R_L . To optimize the design parameters we need an analysis of the magnetic circuit of the transducer. This will enable us to calculate the energy dissipated in R_L as a function of the gaps between pole pieces (l_g), the reed thickness (t), the reed displacement amplitude (a), and the frequency of vibration of the reed (f).

THE MAGNETIC CIRCUIT EQUATIONS

The magnetic circuit can be replaced by the analogous electric circuit as shown in Fig. 2, where the various electrical admittances have the values

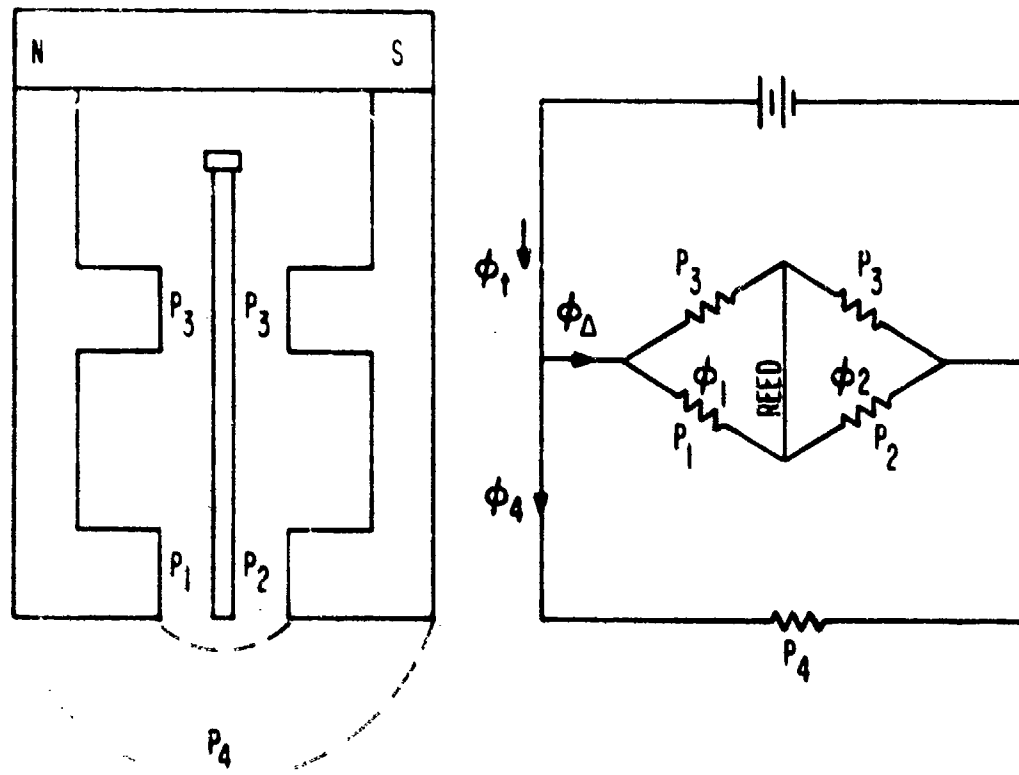


Fig. 2. Schematic Representation of the Magnetic Circuit and its Electrical Analogue

of the relevant magnetic permeances (P) of the reed generator, and the magnitudes of the currents deduced will then be the same as those of the fluxes (Φ) carried in the corresponding branches of the magnetic circuit. The circuit analysis is straightforward:

$$\Phi_{\Delta} = \frac{P_{\Delta}}{P_4 + P_{\Delta}} \Phi_t \quad (1)$$

$$\Phi_1 = \frac{P_1}{P_1 + P_3} \Phi_{\Delta} \quad (2)$$

$$\Phi_2 = \frac{P_2}{P_2 + P_3} \Phi_{\Delta} \quad (3)$$

$$\Phi_R = \Phi_1 - \Phi_2 \quad (4)$$

Substituting (1), (2), and (3) in (4) yields

$$\Phi_R = \left(\frac{P_1}{P_1 + P_3} - \frac{P_2}{P_2 + P_3} \right) \frac{\Phi_t}{1 + \frac{P_4}{P_{\Delta}}} \quad (5)$$

$$\frac{1}{P_{\Delta}} = \frac{1}{P_1 + P_3} + \frac{1}{P_2 + P_3} \quad (6)$$

$$P_{\Delta} = \frac{(P_1 + P_3)(P_2 + P_3)}{P_1 + P_2 + 2P_3} \quad (7)$$

Substitution of (7) in (5) finally gives us:

$$\Phi_R = \Phi_t \left(\frac{P_1}{P_1 + P_3} - \frac{P_2}{P_2 + P_3} \right) \frac{1}{1 + \frac{P_4(P_1 + P_2 + 2P_3)}{(P_1 + P_3)(P_2 + P_3)}} \quad (8)$$

$$\Phi_R = \Phi_t \left(\frac{P_1}{P_1 + P_3} - \frac{P_2}{P_2 + P_3} \right) \left(\frac{(P_1 + P_3)(P_2 + P_3)}{(P_1 + P_3)(P_2 + P_3) + P_4(P_1 + P_2 + 2P_3)} \right) \quad (9)$$

$$\Phi_R = \Phi_t \frac{P_1(P_2 + P_3) - P_2(P_1 + P_3)}{(P_1 + P_3)(P_2 + P_3) + P_4(P_1 + P_2 + 2P_3)} \quad (10)$$

$$\Phi_R = \Phi_t \frac{P_3(P_1 - P_2)}{(P_1 + P_3)(P_2 + P_3) + P_4(P_1 + 2P_3 + P_2)} \quad (11)$$

THE PERMEANCE CALCULATIONS

At this point it is necessary to obtain expressions for the various permeances P in terms of the reed displacement amplitude (a) and the configurational parameters which can be conveniently varied, namely the two gap lengths i_{gA} and i_{gB} , and the reed thickness t . Since i_{gA} is always kept equal to i_{gB} we denote both gap lengths by i_g .

These permeances will be calculated by essentially the same techniques used in Ref. 4. These consist principally of approximating the actual flux paths by circular arcs and then applying the permeance formulae for these paths as listed in the standard text by Roters.⁵ Suitable modifications of these general formulae must sometimes be made to deal with the peculiarities of a particular configuration. The actual arrangements for the prototype generator in both the isolated and in-shell conditions are shown in Figs. 3 and 4 respectively. The numerical values of the pertinent configurational dimensions are summarized in Table 1.

The permeances associated with all of the flux paths which do not pass through the reed are collectively denoted by P_4 . There are five such paths which are schematically illustrated in Fig. 5, and are calculated immediately below along with the other permeances illustrated in Fig. 5. All calculated permeances are in inches.

(1) P_{p4} (Fig. 5-A)

The flux path here is in the form of a parallel plate magnetic capacitor and the permeance is simply given by

$$P_{p4} = \frac{\text{Plate Area}}{\text{Plate Separation}} = \frac{A_p}{S} = \frac{(.7)(.365)}{.395} = .647'' \quad (12)$$

P_{p4} will be slightly modified by the presence of the reed, but since the reed is very thin compared to the separation of the keeper plates this modification is very small and will be neglected.

(2) P_{y4} (Fig. 5-B)

This is a standard permeance path for which the formula is

$$.318 y \ln \left(1 + \frac{2w}{l} \right) \quad (13)$$

where y is the common perimeter length of the keeper edges linked by the flux path. In the fluidic generator $y = 2.2$ inches. w is the edge thickness $.07''$ and l is the diameter of the smaller of the two semicircular boundaries of the flux path $(.395'')$. Insertion of these values into Eq. (13) yields

4. H.A. Leupold, F. Rothwarf, C.G. Campagnoulo, J. Fine, and H. Lee, "Magnetic Circuit Design Studies for an Inductive Sensor," ECOM Technical Report-4158, October 1973.

5. H. Roters, Electromagnetic Devices, (John Wiley & Sons, Inc., 1941).

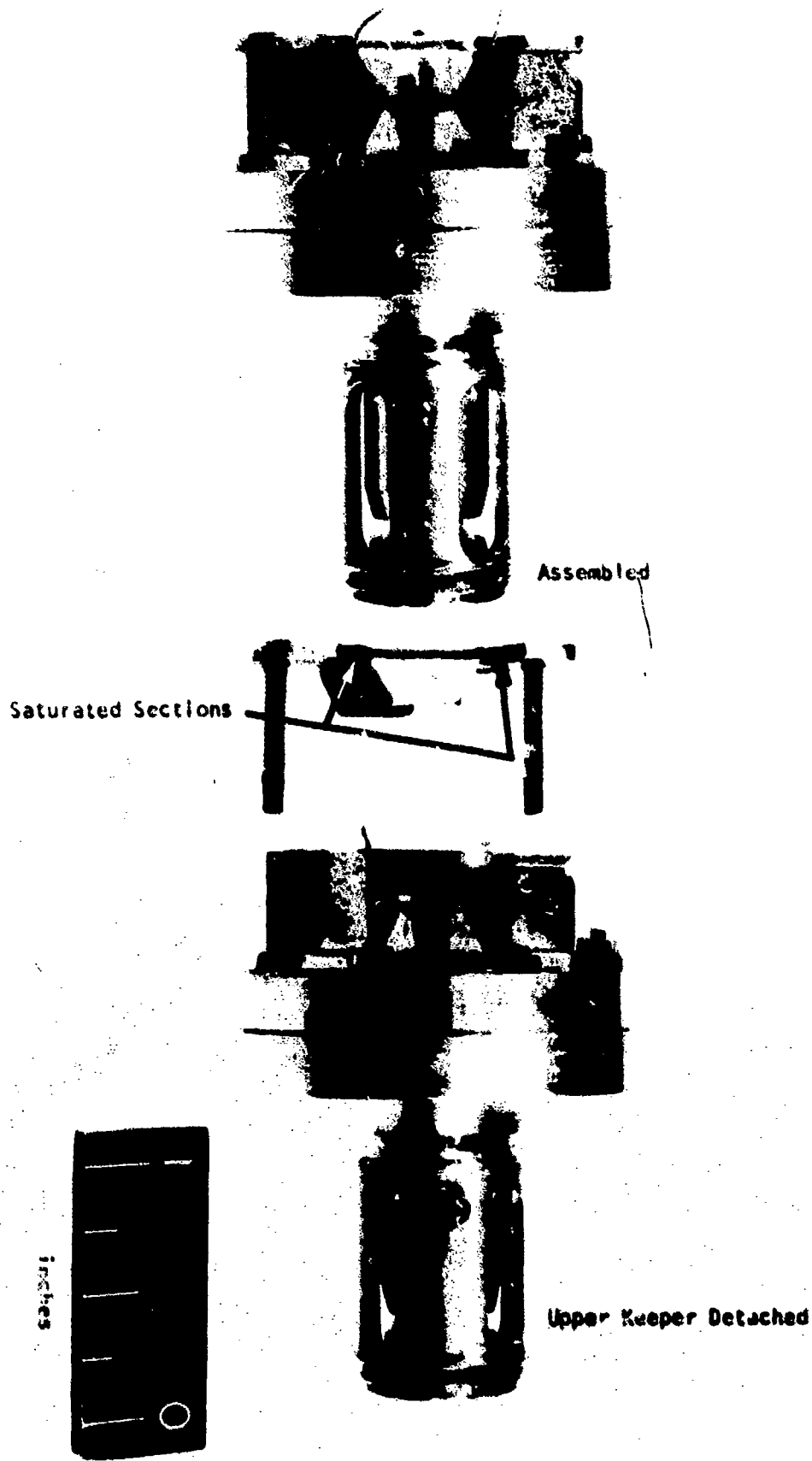
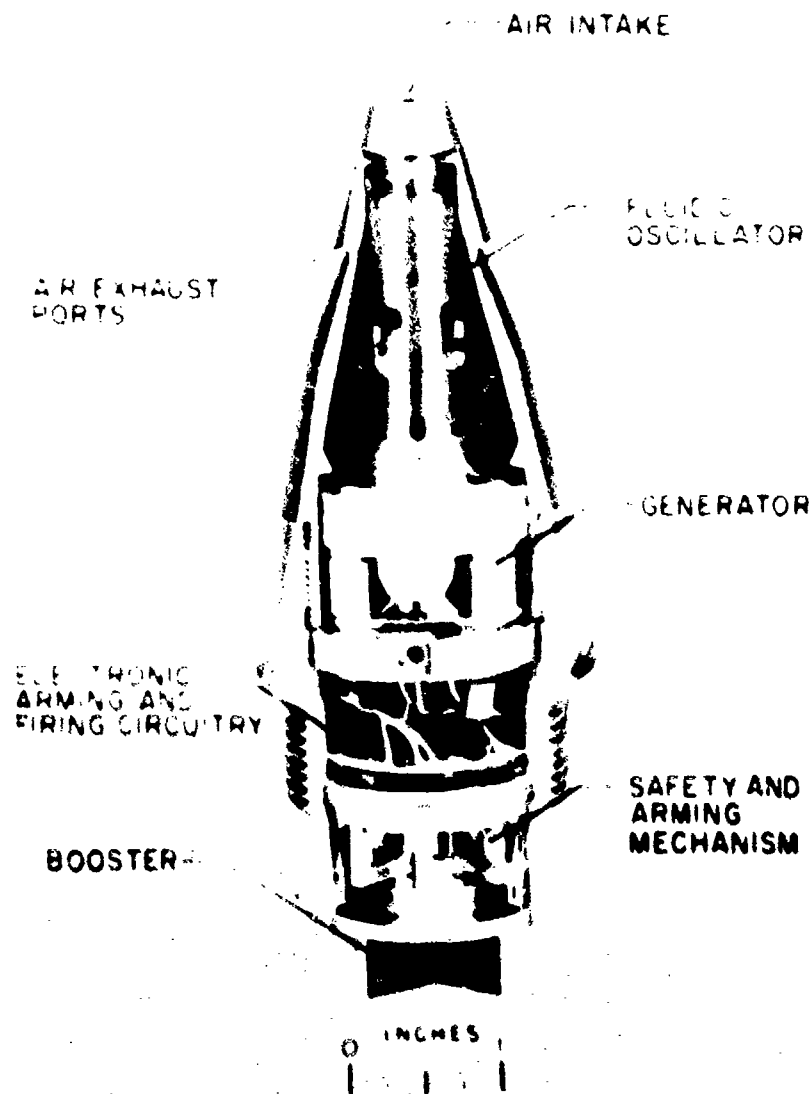


Fig. 3. Fluidic Generator.



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Fig. 4. Sectional View of Fluidic Generator
in Projectile Ogive

TABLE I

Fixed Dimensions of Generator Configuration and Values of Critical Constants

<u>Keeper</u>	<u>Inches</u>	<u>cm</u>	<u>Magnet</u>	<u>Inches</u>	<u>cm</u>
Plate area for P_{p_2} calculation	0.256 in ²	1.65 cm ²	Length	0.395	1.00
Plate separation	0.395	1.00	Cross-section area	0.337 in ²	2.17 cm ²
Thickness	0.070	0.178	Effective outer leakage perimeter	2.3	5.8
Width (narrowest)	0.360	0.914			
Perimeter (effective)	2.2	5.6			
Saturation flux density	16,000 gauss				
Resistivity ρ	45 $\mu\Omega$ cm				
Hysteresis constant η	0.0001				
<u>Foles</u>	<u>Inches</u>	<u>cm</u>	<u>Reed</u>	<u>Inches</u>	<u>cm</u>
Face thickness	0.070	0.178	Width	0.270	0.686
Face width	0.313	0.795	Effective length	0.500	1.27
Face perimeter	0.766	1.95	Thickness *	0.0320	0.0812
Face area	0.0219 in ²	0.141 cm ²	Projection beyond gap	0.150	0.381
Average width	0.475	1.21	Vibrational amplitude *	0.017	0.043
Height (inner)	0.150	0.381	Volume *	0.00432 in ³	0.0706 cm ³
Height (outer)	0.235	0.597	Saturation flux density B_m	19,900 gauss	
Gap length *	0.070	0.178	Resistivity ρ	50 $\mu\Omega$ cm	
			Hysteresis constant η	.002	
			Projection beyond pole pieces	0.15	0.38

* Values are for sample calculations only. They are variable for computer calculated values in Table II (see Appendix II).

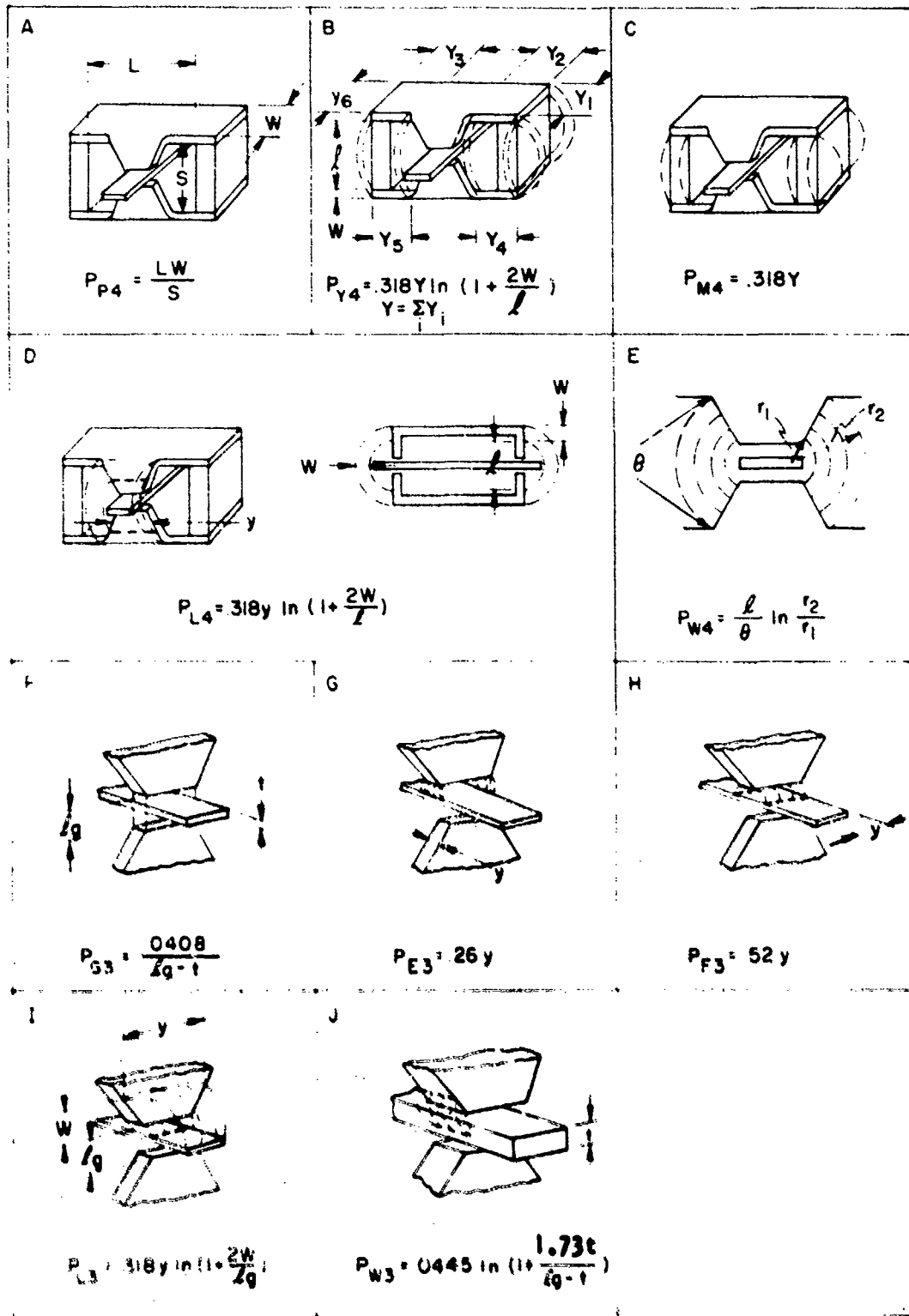


Fig. 5. Calculation of the Relevant Permeances

$$P_{y4} = .318 (2.2) \ln \left(1 + \frac{2(.07)}{.395} \right) = (.318)(2.2)(.303) = .212''$$

(3) P_{M4} (Fig. 5-C)

This is not a standard permeance but an expression for it can be derived as follows: Consider for the time being only that portion of the permeance path formed by the annular ring of width W as shown in Fig. 6. As we shall

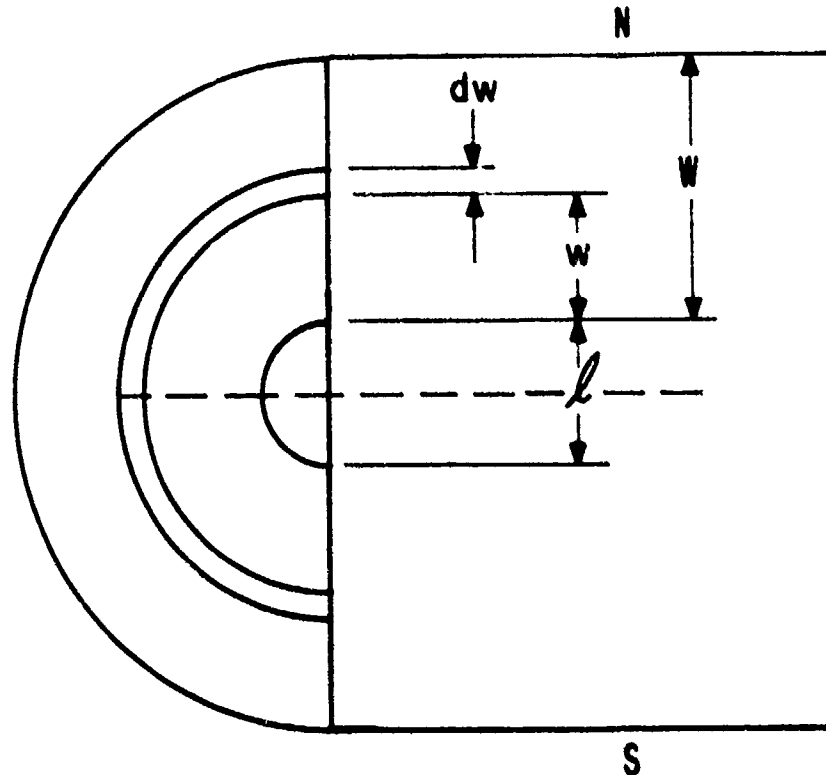


Fig. 6. Calculation of the Permeance P_{M4}

see a simple extension of the resulting formula can then be used to determine the remaining permeance consisting of the semicircular region of diameter l . The permeance of a path such as the annular ring is of the same standard form as Eq. (13), i.e.

$$P = .318 y \ln \left(1 + \frac{2W}{l} \right) \quad (14)$$

Similarly, the portion of the permeance within the annular ring of width w is

$$P = .318 y \ln \left(1 + \frac{2w}{l} \right) \quad (15)$$

and the permeance of element dw is

$$dP = \frac{.318y \frac{2}{l} dw}{1 + \frac{2w}{l}} \quad (16)$$

However, since the flux lines are emanating from the sides of a magnet, the magnetomotive force driving an element such as dw is only a fraction of the full MMF of the magnet. Since this MMF is in direct proportion to the distance of the path terminals from the magnet center, this fraction is equal to

$$f = \frac{\frac{l}{2} + w}{\frac{l}{2} + W} \quad (17)$$

and the effective permeance of dP is reduced by the factor f , viz.,

$$dP_{\text{eff}} = fdP = .318y \frac{(\frac{l}{2} + w) \frac{2}{l}}{(\frac{l}{2} + W) (1 + \frac{2w}{l})} dw$$

$$dP_{\text{eff}} = \frac{.318y}{(\frac{l}{2} + W)} \quad (18)$$

and the total permeance of the annular ring W is

$$P_{\text{eff}} = \int_0^W \frac{.318y}{W + \frac{l}{2}} dw = \frac{.318y W}{W + \frac{l}{2}} \quad (19)$$

To obtain the total effective permeance P_{M4} , one needs only to set $l = 0$ in the expression for P_{eff} and we are left with

$$P_{M4} = .318y \quad (20)$$

y is just the perimeter of the magnet minus the length of the inner wall. The latter length is subtracted because it is rendered largely ineffective by competition from the flux paths of permeance P_{p4} . The effective y is somewhat awkward to estimate because of irregularities such as the screw slots cut into the magnet walls, but 2.3" seems a reasonable value for the two magnets together. Accordingly, our value for P_{M4} is

$$P_{M4} = (.318)(2.3) = .731''$$

(4) P_{L4} (Fig. 5-0)

This again is a standard permeance of the form used in P_{y4} . There is one such path for the outside surface of each pole piece so we must multiply Eq. (13) by two. The flux lines emanating from the inner surfaces

go to the reed and hence contribute to P_1 , P_2 , and P_3 rather than to P_{L4} . For y we take the average width, (.475") of the effective portion of the trapezoidal outer surface of each pole piece. l , (.30") is twice the length of the reed projection beyond the gap as is clear from Fig. 5-D. As .235" is the height of the trapezoidal pole piece plus the gap from the pole face to reed, it also follows from Fig. 5-D that $w = (.235 - .150) = .085$ ". Inserting these values in Eq. (13) yields

$$P_{L4} = 2(.318)(.475) \ln \left(1 + \frac{2(.085)}{.300} \right)$$

$$P_{L4} = .302 \ln 1.57 = .136''$$

(5) P_{w4} (Fig. 5-E)

The formula for a path of this kind is

$$P = \frac{l}{\theta} \ln \frac{r_2}{r_1} \quad (21)$$

and since we have four such paths (2 for each pair of pole pieces)

$$P_{w4} = 4 \frac{l}{\theta} \ln \frac{r_2}{r_1} \quad (22)$$

l is just the thickness of our pole pieces which is .070" and θ is $\frac{2}{3} \pi$ radians. r_2 is the distance from the vertex of angle θ to the larger base of the trapezoidal pole piece. Thus, r_2 is equal to the length of the pole piece edge plus one-half the gap width divided by the sine of $\frac{\pi}{3}$ radians. So since the pole piece edge is .2" long

$$r_2 = .2 + \frac{l_g}{2} \csc \frac{\pi}{3} = .2 + .577 l_g \quad (23)$$

r_1 is equal to $\frac{l_g}{2} \csc \frac{\pi}{3}$ plus the distance along the pole piece edge from which the flux lines go to the edge of the reed. The latter distance is just one-half the reed thickness (t). So finally we have

$$r_1 = \frac{t}{2} + .577 l_g \quad (24)$$

and

$$P_{w4} = \frac{(4)(.070)(3)}{2\pi} \ln \frac{.2 + .577 l_g}{\frac{t}{2} + .577 l_g} \quad (25)$$

$$P_{w4} = .134 \ln \frac{.4 + 1.16 l_g}{t + 1.16 l_g} \quad (26)$$

(6) P_4

We now have expressions for all of the components of the "leakage permeance" P_4 , so P_4 is given by:

$$\begin{aligned} P_4 &= P_{p4} + P_{y4} + P_{M4} + P_{L4} + P_{w4} \\ P_4 &= .647 + .212 + .731 + .136 + .134 \ln \frac{.4 + 1.16 \ell_g}{t + 1.16 \ell_g} \\ P_4 &= 1.72 + .134 \ln \frac{.4 + 1.16 \ell_g}{t + 1.16 \ell_g} \end{aligned} \quad (27)$$

If the gap is varied in the usual manner, i.e. by placing permalloy shims between the magnets and the permalloy keeper, P_{w4} , P_{p4} , and P_{y4} will be slightly affected, but since the shim thicknesses are only of the order of 10 or 20 thousandths of an inch the change in these quantities will be negligible.

(7) P_{G3} (Fig. 5-f)

This is the permeance of the gap between the reed and either pole piece at the pinned end of the reed. The length of this gap L is given by

$$L = \frac{\ell_g - t}{2} \quad (28)$$

Since the widths of the pole piece (.313") and the reed .270" are not the same, we take an average (.291") in our determination of the gap's cross-sectional area. The pole face thickness is .070" so

$$P_{G3} = \frac{A_{c.s.}}{L} = \frac{2(.070)(.291)}{\ell_g - t} \quad (29)$$

$$P_{G3} = \frac{.0408}{\ell_g - t} \quad (30)$$

If the gap sizes are varied by machining the pole faces the constant in this formula will change due to the alteration in cross-sectional area of the pole faces.

(8) P_{E3} (Fig. 5-G)

P_{E3} is the permeance of the path extending from the edges of the reed to the pole face edges parallel to the former. The standard expression for this type of permeance path is

$$P = .26y \quad (31)$$

where y is the edge length which here is the thickness of the pole piece (.070") so

$$P = (.26)(.070) = .0182'' \quad (32)$$

Since we have two such edges the total permeance is

$$P_{E3} = 2P = (2)(.0182) = .0364'' \quad (33)$$

(9) P_{F3} (Fig. 5-H)

These are the paths from the edges of the pole pieces to the flat portion of the reed. These are standard forms whose permeances are given by

$$P = .52y = .52(.291) \quad (34)$$

Since there are two paths we have

$$P_{F3} = 2P = 2(.52)(.291) = .303'' \quad (35)$$

(10) P_{L3} (Fig. 5-I)

These permeances involve the paths between the walls of the pole pieces and the wide faces of the reed. The general expressions are the same as those for P_{L4} except for a factor of 2 to take into account that the flux lines here are quarter circles rather than semicircles as in P_{L4} . The gap length l used in the formula must be taken as twice the separation between the reed and the pole piece, i.e.,

$$l = 2 \left(\frac{l_g - t}{2} \right) = l_g - t.$$

The quantity w_i for the inner flux path is just the altitude of the trapezoidal pole piece which is .15". For the effective y one chooses an average of the reed width (.270") and the width of the trapezoid midway between its bases (.406"). This results in a value $y = .338''$.

The outer path must be treated slightly differently because the reed projects only .15" beyond the gap edge. Accordingly, w_0 is just taken as

$$w_0 = .15 - \frac{(l_g - t)}{2} \quad (36)$$

and the average value y is also slightly altered.

As is seen in Fig. 7 the flux lines from the reed extend only up to line AB which is .15" above the top surface of the reed. The average width of the flux-emitting region of the pole face is then

$$y_F = \frac{AB + .313}{2} \quad (37)$$

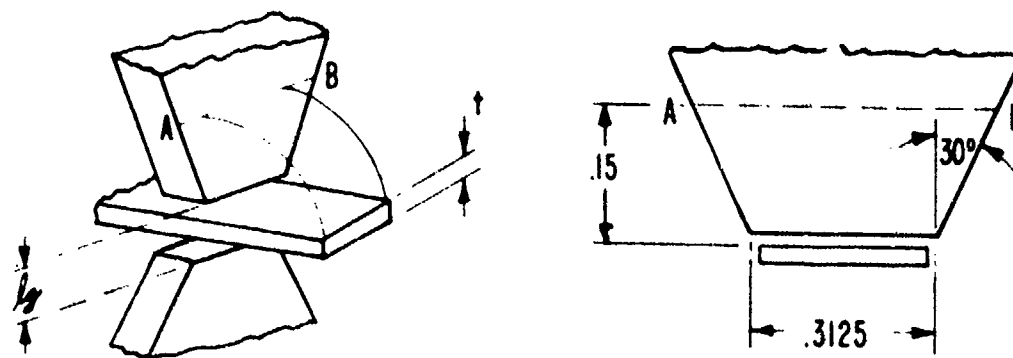


Fig. 7. Determination of the width y_F in the Formula for Permeance P_{L3}

and

$$y_0 = \frac{y_F + y_R}{2} \quad (38)$$

where $y_R = .270$ is the reed width.

$$AB = .313 + 2 \left(.15 - \left(\frac{l_g - t}{2} \right) \right) \tan 30^\circ \quad (39)$$

$$AB = .313 + (.30 - (l_g - t)) .577$$

$$AB = .486 - .577 (l_g - t) \quad (40)$$

$$y_F = \frac{.486 - .577 (l_g - t) + .313}{2} = .399 - .289 (l_g - t) \quad (41)$$

$$y_0 = \frac{.399 - .289 (l_g - t) + .270}{2} \quad (42)$$

$$y_0 = .335 - .144 (l_g - t) \quad (43)$$

Substituting these values in the general expression yields

$$P_{L3} = 2(.318)(.338) \ln \left(1 + \frac{2(.15)}{l_g - t} \right) + 2(.318)[.335 - .144 (l_g - t)] \ln \left(1 + \frac{2 \left(.15 - \frac{l_g - t}{2} \right)}{l_g - t} \right) \quad (44)$$

$$P_{L3} = .215 \ln \left(1 + \frac{.30}{l_g - t} \right) + (.213 - .0916 (l_g - t)) \ln \left(1 + \frac{.30 - (l_g - t)}{l_g - t} \right) \quad (45)$$

$$P_{L3} = .215 \ln \left(1 + \frac{.30}{l_g - t} \right) + (.213 - .0916 (l_g - t)) \ln \left(\frac{.30}{l_g - t} \right) \quad (46)$$

(11) P_{W3} (Fig. 5-J)

If the narrow side face of the reed were coplanar with the narrow face of the pie piece as indicated by the dashed line in Fig. 5-J, the expression for permeance P_{W3} would be the often used standard form

$$P_{W3} = .318y \ln \left(1 + \frac{2w}{l} \right) \quad (47)$$

Although it is possible to correct for the fact that the actual flux lines are shorter than in this ideal case, the correction would be very small, and since the total permeance P_{W3} is itself so small as to barely affect the calculations with one significant figure, this correction is completely

negligible and will not be made here. So w is taken as $\frac{t}{2}$, $l = \frac{l_g - t}{2 \cos 30^\circ}$.

$y = .070$, and we precede the expression with the factor 2 because there are two paths of this type

$$P_{W3} = 2(.318)(.070) \ln \left(1 + \frac{2t \cos 30^\circ}{l_g - t} \right) \quad (48)$$

$$P_{W3} = .0445 \ln \left(1 + \frac{1.73t}{l_g - t} \right) \quad (49)$$

(12) P_{G1} and P_{G2}

These permeances are of the same form as P_{G3} with allowance made for the spacing changes due to the displacement of the reed. When the reed is at equilibrium, that is, at zero displacement, $P_{G1} = P_{G2} = P_{G3}$. Since we are interested in the maximum and minimum values of P_{G1} and P_{G2} , we write them in terms of the displacement amplitude (a) of the reed.

$$P_{G1}^{\max} = P_{G2}^{\max} = \frac{.0408}{l_g - t - 2a} \quad (50)$$

$$P_{G1}^{\min} = P_{G2}^{\min} = \frac{.0408}{l_g - t + 2a} \quad (51)$$

where P_{G1} is a maximum when P_{G2} is a minimum and vice versa.

(13) $\underline{P_{E1}}$

$$P_{E1} = P_{E3} = .0364'' \quad (52)$$

(14) $\underline{P_{F1}}$

$$P_{F1} = P_{F3} = .303'' \quad (53)$$

(15) $\underline{P_{w1} \text{ and } P_{w2}}$

P_{w1} is of the same form as P_{w3} when the reed is in equilibrium. When P_{w1} is a maximum, $l = (l_g - t)$ must be replaced by $(l_g - t - 2a)$ and by $(l_g - t + 2a)$ when P_{w1} is a minimum. The same holds for P_{w2} . Thus we have

$$P_{w1}^{\max} = P_{w2}^{\max} = .0445 \ln \left(1 + \frac{1.73t}{l_g - t - 2a} \right) \quad (54)$$

$$P_{w1}^{\min} = P_{w2}^{\min} = .0445 \ln \left(1 + \frac{1.73t}{l_g - t + 2a} \right) \quad (55)$$

(16) $\underline{P_{L1} \text{ and } P_{L2}}$

The term in P_{L1} arising from the inner flux path is the same as the corresponding term in P_{L3} except that again the quantity $(l_g - t)$ is replaced by $(l_g - t - 2a)$ or $(l_g - t + 2a)$ to obtain P_{L1}^{\max} and P_{L1}^{\min} , respectively.

The term arising from the outer flux path is similarly derived from the corresponding term in P_{L3} by the same substitution but an additional alteration is necessary. The portion of the reed projecting beyond the limits of the gap is triangular rather than rectangular as is the projection at the anchored end of the reed. For this reason the flux lines emanating from the walls of the pole piece have only one half the reed area in which to terminate. To take this into account it seems reasonable to substitute in the expression for y_0 , one half of the reed width for the full reed width. Thus, in Eq. (42) the value of .270 is replaced by .135. After all these adjustments are made the expression for P_{L1} and P_{L2} becomes

$$P_{L1}^{\max} = P_{L2}^{\max} = .215 \ln \left(1 + \frac{.30}{l_g - t - 2a} \right) + .170 - .0916 (l_g - t - 2a) \ln \frac{.30}{l_g - t - 2a} \quad (56)$$

There are other permeances associated with flux paths going from corner to corner, edge to corner and face to corner. In this problem, these are not of standard form and are very difficult to estimate. They tend however, to be small compared to most of the calculated permeances, and the labor of any attempt to estimate them would not be justified by the precision attainable by our computational techniques.

CALCULATIONS OF GAP FLUX DENSITIES IN THE ABSENCE OF A REED

We now have expressions for all of the permeances relevant to the desired power calculations. Before a sample calculation for particular values of l_g , t and a are made, it would be instructive to calculate the field in the gaps when the reed is not present. Since a measured field value exists for a gap width of .060, a comparison with the result of our calculation would give an indication as to the general efficacy of the computational methods used. To do this we need the permeance P_{c4} which is given by:

$$P_{c4} = \frac{(2)(\text{Area Pole Faces})}{\text{Gap Length}} = \frac{(2)(.0219)}{.06} = .730'' \quad (57)$$

We now also have the permeances P_{E4} between the corresponding edges of the pole faces and P_{c4} between the corresponding corners on either side of the gap. P_{E4} is given by the usual formula $P_{E4} = 2(.26)y$, where y is the perimeter of a pole face and the factor of 2 enters because there are two gaps. P_{c4} is given by $(8)(.077)l_g$, where the factor 8 occurs because there are four corner-to-corner permeances for each of the two gaps. Therefore, P_{c4} is given by

$$P_{c4} = (8)(.077)(.060) = .0369'' \quad (58)$$

$P_{E4} = (.52)(.766) = .398$. P_{p4} , P_{y4} and P_{m4} are the same as when the reed is present. We, therefore, need only to find the new values of P_{w4} and P_{L4} . P_{w4} is given by Eq. (26) with t set equal to zero. Thus

$$P_{w4} = .134 \ln \frac{.4 + (1.16)(.06)}{1.16(.06)} = .134 \ln 6.75$$

$$P_{w4} = (.134)(1.91) = .256'' \quad (59)$$

The contributions of the two outer surfaces to P_{L4} are given by the same formula as when the reed is present, but now $l = l_g = .06$ and

$$\bar{y} = \frac{.5 + .313}{2} = .406. \text{ So these contributions}$$

$$P_{L4}^0 = 2(.318)(.406) \ln \left(1 + \frac{2(.2)}{.06} \right) = .258 \ln 7.67$$

$$P_{L4}^0 = (.258)(2.04) = .526'' \quad (60)$$

In the absence of the reed the inner surfaces also contribute. The longest semicircular arc of this path, however, must be no longer than the distance d as shown in Fig. 8, so

$$\pi \left(\frac{l}{2} + w \right) = d \quad (61)$$

$$w = \frac{d}{\pi} - \frac{l}{2} \quad (62)$$

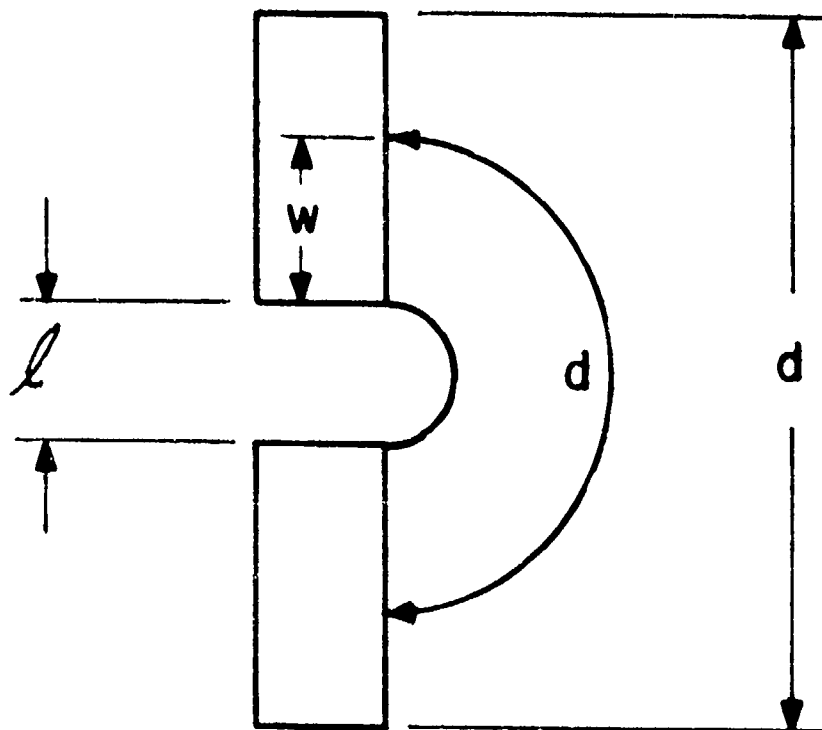


Fig. 8. Determination of the Inner Flux Paths Comprising Permeance P_{M4}

$$w = \frac{.395}{3.14} - \frac{.060}{2}$$

$$w = .126 - .030 = .0958''$$

and

$$\bar{y} = \frac{2(.313) + (2)(.0958) \tan 30^\circ}{2} = .313 + .055 = .368'' \quad (63)$$

Thus

$$P_{L4}^1 = 2(.318)(.368) \ln \left(1 + \frac{2(.0958)}{.06} \right) = .234 \ln (4.03) = .325 \quad (64)$$

$$P_{L4} = P_{L4}^1 + P_{L4}^0 = .325 + .526 = .851''$$

Our total permeance P_4 in the absence of a reed is then

$$P_4 = P_{G4} + P_{M4} + P_{Y4} + P_{H4} + P_{W4} + P_{L4} + P_{E4} + P_{C4} \quad (65)$$

$$P_4 = .730 + .647 + .212 + .731 + .256 + .851 + .398 + .037$$

$$P_4 = 3.86 \quad (66)$$

In this case P_4 is our total permeance P_t and from it we can obtain our load line slope

$$\frac{B}{H} = \frac{L}{A} P_t = \frac{(.395)(3.86)}{(.337)} = 4.52 \quad (67)$$

where $L = .395$ in and $A = .337$ in². And from the intersection of the Alnico 5 demagnetization curve with the load line $B/H = 4.52$ in Fig. 9 we obtain

$$\bar{B}_{mag} = 2670 \text{ gauss} \quad (68)$$

Since the magnetic configuration is of irregular shape, B is not perfectly uniform over the entire magnet, therefore, \bar{B}_{mag} represents an effective average flux density. So the total flux output of the magnet is

$$\Phi_t = \bar{B}_{mag} A_{mag} = (2670)(2.17) = 5790 \text{ Maxwells} \quad (69)$$

The flux in one gap Φ_G is given by

$$\Phi_G = \frac{P_{G4}}{2P_t} \Phi_t = \frac{.730}{2(3.86)} 5790 = 547 \text{ Maxwells} \quad (70)$$

and

$$B_G = \frac{\Phi_G}{A_G} = \frac{547}{.141} = 3880 \text{ gauss} \quad (71)$$

This value is about 29% higher than the measured value of 3000 gauss which is satisfactory agreement considering the various approximations made in the course of the calculations. One must also keep in mind that such comparisons should not be taken too seriously. The calculated B_G represents an average flux density obtained by dividing the total flux emanating from a pole face by the area of that pole face. The measured B_G can represent either a much more localized value of B_G or an average that includes fluxes not even associated with the gap, depending on the type, positioning and size of the probe used. The wider the gap compared to the pole face dimensions, the more difficult it is to obtain a meaningful measured value of B_G . Considering these limitations, we conclude that the general computational procedure yields reasonable agreement with the measured results, and that its judicious application to the determination of expected power outputs should not lead us too far astray.

SAMPLE CALCULATION FOR POWER OUTPUT IN THE LOAD RESISTANCE R_L

We will now perform a sample calculation for the most efficient configuration measured thus far. The relevant parameters for this arrangement are: $l_g = .070$ ", $t = .032$ ", $a = .017$ ".

We will first calculate P_4 . We need only the components, P_{w4} and P_{L4} . Using Eq. (26) we obtain

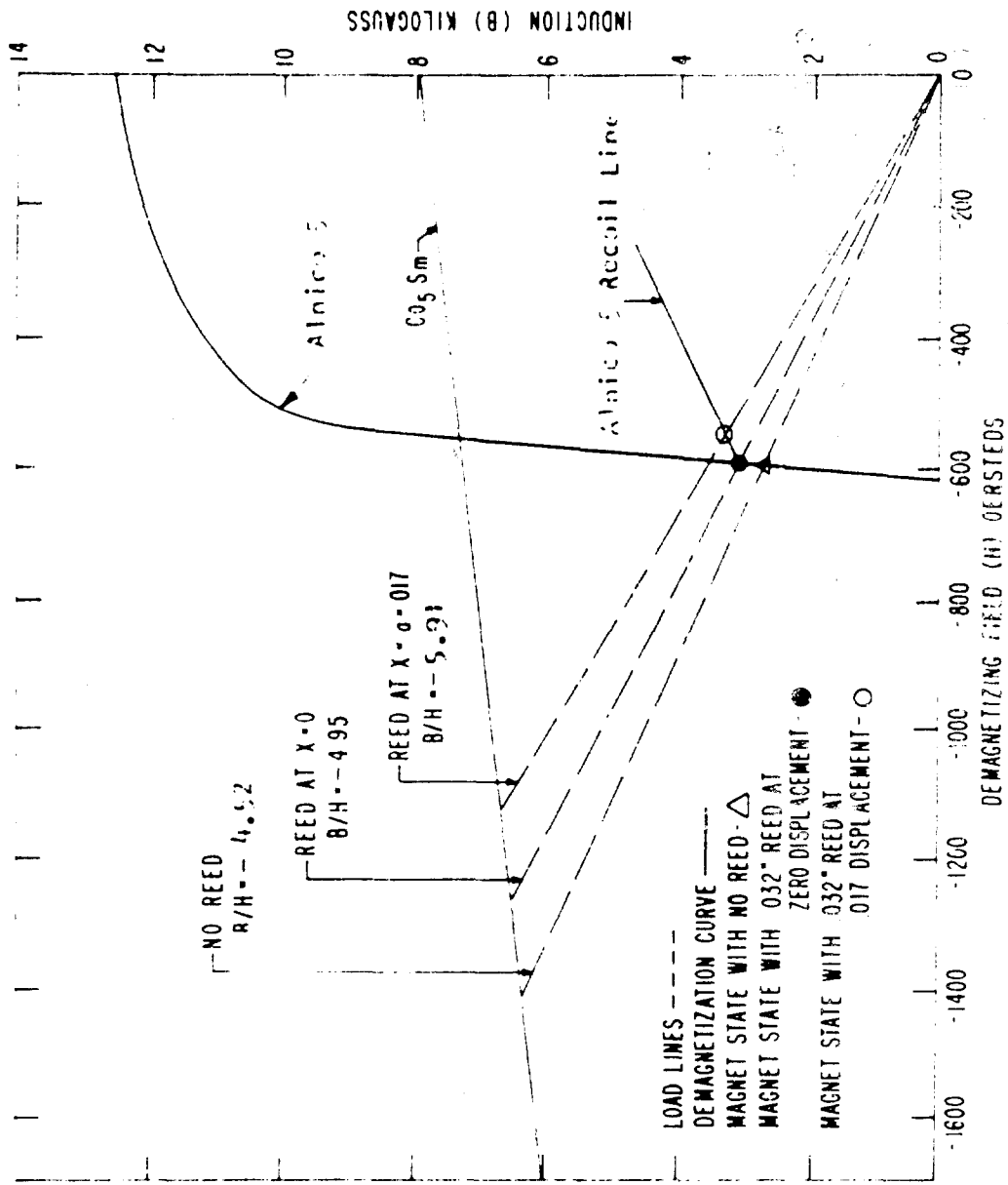


FIG. 9 DETERMINATION OF B FOR A FLUIDIC GENERATOR WITH A .070 " GAP LENGTH.

$$P_{w4} = .134 \ln \frac{.4 + 1.16 l_g}{t + 1.16 l_g} = .134 \ln \frac{.4 + (1.16)(.07)}{.032 + (1.16)(.07)}$$

$$P_{w4} = .134 \ln 4.26 = .134(1.45) = .194'' \quad (72)$$

The appropriate P_{L4} has already been calculated on page 11 and found to be .136. The other components of P_4 are the same as were used to calculate the gap fields except that now there is no P_{G4} , P_{E4} or P_{C4} . Thus

$$P_4 = P_{p4} + P_{y4} + P_{M4} + P_{L4} + P_{w4} \quad (73)$$

$$P_4 = .647 + .212 + .731 + .136 + .194$$

$$P_4 = 1.92'' \quad (74)$$

P_3 has five components

$$P_{G3} = \frac{.0408}{.070 - .032} = \frac{.0408}{.038} = 1.07'' \quad (75)$$

$$P_{E3} = .0364'' \quad (76)$$

$$P_{F3} = .303'' \quad (77)$$

$$P_{L3} = .215 \ln \left(1 + \frac{.30}{l_g - t} \right) + (.213 - .0916 (l_g - t)) \ln \left(\frac{.30}{(l_g - t)} \right) \quad (78)$$

$$P_{L3} = .215 \ln \left(1 + \frac{.30}{.038} \right) + (.213 - .0916(.038)) \ln \left(\frac{.30}{.038} \right)$$

$$P_{L3} = .215 \ln 8.90 + (.213 - .0035) \ln 7.90 = (.215)(2.19) + (2.09)(2.07)$$

$$P_{L3} = .470 + .432''$$

$$P_{L3} = .902'' \quad (79)$$

$$P_{w3} = .0445 \ln \left(1 + \frac{1.73t}{l_g - t} \right) = .0445 \ln \left(1 + \frac{(1.73)(.032)}{.038} \right) \quad (80)$$

$$P_{w3} = .0445 \ln (2.46) = (.0445)(.900) = .0400'' \quad (81)$$

$$P_3 = 1.07 + .0364 + .303 + .902 + .0400$$

$$\underline{P_3 = 2.35''} \quad (82)$$

P_1 and P_2 , when the reed is at its equilibrium position, are obtained by substituting $a = 0$ in the expression for the components of P_1 and P_2 and then adding these components. This procedure yields $P_1 = P_2 = 2.27''$. We need to find P_t^E , that is, the total permeance at equilibrium to locate the point on the Alnico 5 demagnetization curve at which the recoil line begins. (Point E in Fig. 9)

$$P_t^E = P_4 + \frac{P_3}{2} + \frac{P_1}{2} = 4.23 \quad (83)$$

$$\left(\frac{B}{H}\right)_E = -\frac{l}{A} P_t^E = - (1.17)(4.23) = - 4.95 \quad (84)$$

$$B_H = 2950 \text{ gauss (from Fig. 9)} \quad (85)$$

$$H = -\frac{2950}{4.95} = - 596 \text{ Oe} \quad (86)$$

The slope of recoil lines for Alnico 5 is 4.3 so the equation of the recoil line is

$$B = 4.3H + B_0 \quad (87)$$

$$2950 = (4.3)(- 596) + B_0$$

$$2950 + (4.3)(596) = B_0$$

$$2950 + 2560 = B_0$$

$$5510 = B_0 \quad (88)$$

So we have for our recoil line

$$B = 4.3H + 5510 \text{ gauss} \quad (89)$$

We now need only P_1 and P_2 to complete our analysis. Using Eqs. (50) and (51) we have

$$P_{G1}^{\max} = P_{G2}^{\max} = \frac{.0408}{l_g - t - 2a} = \frac{.0408}{.070 - .032 - 2(.017)} = \frac{.0408}{.004} \quad (90)$$

$$P_{G1}^{\max} = P_{G2}^{\max} = 10.2'' \quad (91)$$

$$P_{G2}^{\min} = P_{G1}^{\min} = \frac{.0408}{l_g - t + 2a} = \frac{.0408}{.070 - .032 + .034} = \frac{.0408}{.072} = .566'' \quad (92)$$

From (33) and (35) respectively, we have

$$P_{E1} = P_{E2} = .0364'' \quad (93)$$

$$P_{F1} = P_{F2} = .303'' \quad (94)$$

From (54)

$$\left. \begin{array}{l} p_{w1}^{\max} \\ p_{w2}^{\min} \end{array} \right| = .0445 \ln \left(1 + \frac{1.73t}{L_g - t \mp 2a} \right) = .0445 \ln \left(1 + \frac{(1.73)(.032)}{.070 - .032 \mp 2(.017)} \right)$$

$$p_{w1}^{\max} = .0445 \ln \left(1 + \frac{0.0554}{.004} \right) = (.0445) \ln 14.85$$

$$p_{w1}^{\max} = (.0445)(2.70) = .120 \quad (95)$$

$$p_{w2}^{\min} = .0445 \ln \left(1 + \frac{0.0554}{.072} \right) = .0445 \ln 1.77$$

$$p_{w2}^{\min} = (.0445)(.571) = .0254 \quad (96)$$

From (56)

$$p_{L1}^{\max} = .215 \ln \left(1 + \frac{.30}{.004} \right) + [.170 - (.0916)(.004)] \ln \left(\frac{.30}{.004} \right)$$

$$p_{L1}^{\max} = .215 \ln 76 + .170 \ln 75 = (.215)(4.33) + (.170)(4.32)$$

$$p_{L1}^{\max} = .931 + .734 = 1.66 \quad (97)$$

Similarly

$$p_{L2}^{\min} = .215 \ln \left(1 + \frac{.30}{.072} \right) + [.170 - (.0916)(.072)] \ln \left(\frac{.30}{.072} \right)$$

$$p_{L2}^{\min} = .215 \ln 5.17 + .163 \ln 4.17 = .215(1.64) + (.163)(1.43)$$

$$p_{L2}^{\min} = .353 + .233 = .586'' \quad (98)$$

Finally, we have for P_1^{\max} and P_2^{\min}

$$P_1^{\max} = 10.2 + .036 + .303 + .120 + 1.663$$

$$P_1^{\max} = 12.32 \quad (99)$$

$$p_2^{\min} = .566 + .036 + .303 + .025 + .586$$

$$p_2^{\min} = 1.52'' \quad (100)$$

$$p_t^{\max} = P_4 + \frac{(p_1^{\max} + P_3)(p_2^{\min} + P_3)}{p_1^{\max} + p_2^{\min} + 2P_3} \quad (101)$$

$$p_t^{\max} = 1.92 + \frac{(12.32 + 2.35)(1.52 + 2.35)}{12.32 + 2.35 + 1.52 + 2.35}$$

$$p_t^{\max} = 1.92 + \frac{(14.67)(3.87)}{18.5} = 4.98'' \quad (102)$$

$$\left(\frac{B}{H}\right)^{\max} = -\frac{i}{A} p_t^{\max} = - (1.17)(4.98) = - 5.83 \quad (103)$$

From (89)

$$B^{\max} = 4.3H + 5510$$

$$B^{\max} = \frac{5510}{1 + \frac{4.3}{5.83}} = \frac{5510}{1.73} = 3170 \text{ gauss} \quad (104)$$

$$\Phi_t^{\max} = B^{\max} A_{\text{magnet}} = (3170)(2.18) = 6880 \text{ Maxwells} \quad (105)$$

Substituting all of our calculated permeance values and $\Phi_t^{\max} = 6880$ in Eq. (11) yields for the amount of flux carried by the reed

$$\Phi_R^{\max} = 6880 \frac{2.35(12.32 - 1.52)}{(14.7)(3.87) + (1.92)(18.5)}$$

$$\Phi_R^{\max} = \frac{(6880)(2.35)(10.8)}{92.4} = 1890 \text{ Maxwells} \quad (106)$$

The induced voltage amplitude is then given by

$$V^{\max} = \frac{N L \Phi_R^{\max}}{10^8} \quad (107)$$

where ω is the angular frequency of vibration of the reed, which in this case is $2\pi(1500)$ and $N = (1500)$ is the number of turns in the coil.

$$V^{\max} = \frac{(1500)(2\pi)(1500)(1890)}{10^8} = 267 \text{ volts} \quad (108)$$

Attached to the coil is a 600- Ω load resistor and the resistance of the coil is 600 Ω as well. The inductance of the coil with a .017" reed is about 23 mhenrys. To find the power dissipated in the load, we need the total impedance Z of the circuit.

$$Z^2 = R^2 + X_L^2 = (1200)^2 + X_L^2 \quad (109)$$

The inductive reactance X_L is just

$$X_L = 2\pi fL = \frac{(6.28)(1500)(23)}{1000} = 217 \Omega \quad (110)$$

$$Z^2 = (1200)^2 + (217)^2 = 1.49 \times 10^6 \quad (111)$$

$$Z = 1220 \Omega \quad (112)$$

The root mean square current \bar{i} is given by $\bar{i} = \frac{i_{\max}}{\sqrt{2}} = \frac{V_{\max}}{Z/\sqrt{2}}$ and the power dissipated in the load p_L is then

$$p_L = \bar{i}^2 R_L = \frac{V_{\max}^2 R_L}{2Z^2} = \frac{(267)(267)(600)}{(2)(1220)(1220)} = 14.4 \text{ watts} \quad (113)$$

The actual power measured, however, was only 5 watts. The foregoing calculations were made on the assumption that the reed had infinite permeance. The electrical steel of which the reed is composed saturates at 19,900 gauss and its cross-section area is

$$A = (.27 \times .032 \times 2.54 \times 2.54) \text{ cm}^2 \quad (114)$$

Therefore, the maximum amount of flux it can carry is

$$\Phi_s = (19,900)(.27)(.032)(2.54)(2.54) = 1110 \text{ Maxwells} \quad (115)$$

and if, as was shown, 1890 Maxwells in the reed would yield 14.4 watts then the maximum allowable flux of 1110 Maxwells would give us

$$p_s = \left(\frac{1110}{1890} \right)^2 14.4 = 4.97 \text{ watts which is almost exactly the power observed.}$$

It, therefore, appears that to utilize the full flux output of the magnets used, the reed thickness must be increased by about 70%. Whether this can be done without drastically affecting the vibrational response to the oscillator is a mechanical problem beyond the scope of this study.³ It is clear, however, that increased output with the present configuration depends largely upon increasing the flux carrying capability of the reed. If, on the other

3. Carl J. Campagnuolo and Richard N. Gottron, "The Fluidic Generator: A New Electrical Power Source," 24th Annual Proceedings Power Sources Symposium, May 19-21, 1970.

hand, the present output of 5 watts is satisfactory, it can be attained with considerably less magnetic material since with the present arrangement 70% more flux is available than can be profitably used. Of course, if the reed thickness is increased, the gaps must be lengthened by an appropriate amount to keep the load line of the magnets the same and to maintain reasonable space in which the reed can oscillate. Also the expected power output would be altered slightly since the inductance of the coil is a function of reed thickness. This effect, however, would not be significant.

Parts of the keeper may also be saturated under the given operating conditions. The cross sections which would saturate first are the ones indicated in Fig. 3. Under the given conditions, the flux that must be carried by the narrowest sections of the keeper is given by

$$\Phi_s = \frac{\Phi_t^{\max}}{2} \frac{p_t^{\max} - p_m^4 - p_y^4}{p_t^{\max}} \quad (116)$$

$$\Phi_s = \frac{6880}{2} \frac{4.98 - .731 - .212}{4.98}$$

$$\Phi_s = \frac{6880}{2} \frac{4.04}{4.98} = 2790 \text{ Maxwells} \quad (117)$$

and

$$B = \frac{\Phi}{A} = \frac{2790}{(.36)(.07)(2.54)(2.54)} = 17,200 \text{ gauss} \quad (118)$$

Since permalloy saturates at 16,000 gauss, the keeper cross section need be increased by only 7% to take full advantage of the available flux at the given reed displacement. This change alone, of course, would do no good unless the reed thickness were doubled as well, as we have already shown.

EXPERIMENTAL CHECK OF CALCULATIONS

To test the validity of the foregoing calculations, a series of measurements was made to determine the maximum voltage output as a function of reed displacement under conditions for which there is no magnetic saturation. The reed displacement was measured by 'stopping' the vibrating reed at its maximum displacement (a) with a strobe light and then measuring (a) with a telemicroscope equipped with a calibrated micrometer screw.

The results of this experiment are shown in Fig. 10 where they are compared with the theoretical curve plotted from results of the computer calculations listed in Appendix II. As is apparent from Fig. 10, agreement is excellent considering the many approximations used in the calculations and that (a) can be measured with a precision of no better than ten to fifteen percent.

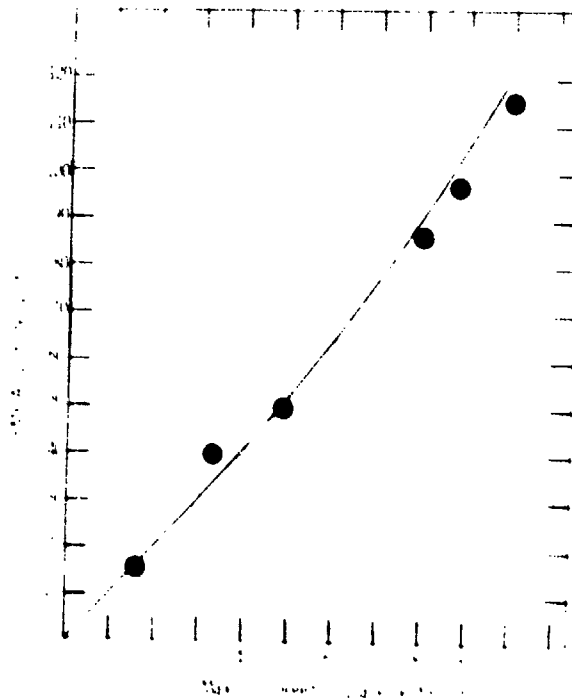


Fig. 10. Comparison of Experimental Points (dots) with Theoretical EMF-Displacement Curve (solid curve)

LOSSES IN THE REED

It is of interest to know how much electrical energy is dissipated in the reed through eddy current losses and hysteresis. The expression for the eddy current power dissipation in watts (see Appendix I) is

$$\frac{P_E}{n} = \frac{\pi^2 f^2 B_m^2 \tau^3 / w}{60 \cdot 10^{16}} = \frac{.2 f^2 B_m^2 \tau^2 V}{60 \cdot 10^{16}} \quad (119)$$

where n is number of laminations, f is the frequency, B_m is the flux density amplitude in gauss, τ is the lamination thickness, l and w are the reed length and width respectively, V is the lamination volume, and ρ is the resistivity in ohm-cm. The reed dimensions are all measured in centimeters and the power is in watts. Since the reed is composed of two laminations each of the thickness $\tau = .0406$ cm, we multiply (119) by two and insert .0406 for τ as well as the appropriate values of the other parameters under saturation conditions.

$$P_E = .34 \text{ watts} \quad (120)$$

The expression for hysteresis power loss in the reed is given by

$$P_H = \pi f^2 \left(\frac{B_m}{10^8} \right)^{1.6} \Delta wt = \pi f^2 \left(\frac{B_m}{10^8} \right)^{1.6} V \quad (121)$$

where $V = .0706 \text{ cc.}$ is a constant depending on the material and is .002 for our reed. Making the proper insertions gives us

$$P_H = (.002)(1500)(1500) \left(\frac{19.9}{10^5} \right)^{1.6} (.0706)$$

$$P_H = \frac{(.002)(1500)(1500)(1200)(.0706)}{10^{10}}$$

$$P_H = .00004 \text{ watts which is negligible.} \quad (122)$$

The eddy current losses, however, could become quite significant if the reed were not properly laminated. For example, an unlaminated reed of the same thickness as the laminated one would have four times the eddy current loss. It is, therefore, important that if the reed thickness is increased as recommended, it be done by adding laminations. A solid unlaminated reed of twice the thickness of the one presently used would have a power loss of about 1 watt, whereas a fourfold lamination of such a reed would reduce the loss to about .68 watts.

LOSSES IN THE KEEPER

Eddy current losses in the keeper come about because of fluctuations in the flux density as the magnet load line changes with the vibration of the reed. If the keeper thickness is increased by the required 7%, or better still, widened at the narrow points which saturate by 7%, the maximum flux density would be no more than 16,000 gauss. The operating conditions would then be optimized and it would be useful to have some estimate as to the eddy current losses. Such an estimate would be very rough since it is very difficult to determine the flux pattern in the keeper. We already know that when the reed is at its greatest displacement the flux carried by the narrowest portions of the keeper would be 16,000 gauss. We also need the flux density when the reed is in its equilibrium position. This is obtained by substituting the appropriate equilibrium values of the parameters in Eqs. (116 - 118). In this fashion we obtain

$$\Phi_K^E = \frac{\Phi_t^E}{2} \frac{P_t^E - P_m^E - P_y^E}{P_t^E}$$

$$\Phi_K^E = \frac{(2950)(2.17)}{2} \frac{4.23 - .731 - .212}{4.23} = 2490 \text{ Maxwells} \quad (123)$$

The cross-section dimensions of the keeper are .914 cm and .178 cm, so

$$B_K^E = \frac{2490}{(.914)(.178)} = 15,300 \text{ gauss}$$

Therefore, the field at the narrowest portion of the keeper fluctuates from 15,300 gauss to 16,000 gauss with an amplitude of $\frac{700}{2} = 350$ gauss. The fluctuation in other parts of the keeper would be somewhat smaller, but 350 gauss can be used to get an upper limit for the expected losses. To this end we use Eq. (119) with the appropriate constants for the keeper whose volume is 2.44 cc and whose thickness would be .178 cm. Further, the fluctuation frequency of B_K^E is double that of the reed since B_K^E is a maximum at both ends of the reed's vibrational cycle. Using these values in Eq. (119) yields:

$$P_E = \frac{(3.14)(3.14)(3000)(3000)(350)(350)(.178)(.178)(2.44) 10^6}{(6)(45) 10^{16}}$$

$$P_E = .31 \text{ watts} \tag{124}$$

which is about 2% of the expected optimum output. This loss of electrical energy can be avoided by a double lamination of the keeper to cut the loss to 0.08 watts or even a quadruple lamination which would result in the low eddy current loss of 0.02 watts. Since these numbers represent only upper limits, the actual values would probably be much smaller. The use of high-resistivity, high-permeability composite castings for the keepers would probably eliminate the eddy current losses entirely.

The hysteresis loss in the keeper cannot be obtained directly from Eq. (121) because the field in the keeper varies only over positive values of B, while (121) applies only to complete cycles about the principal hysteresis loops. We can, however, get some idea as to the order of magnitude of the loss by substituting our amplitude 350 gauss into (121). Using the value $\eta = .0001$ for permalloy we have

$$P_H = (.0001)(3000)(3000) \left(\frac{350}{10^8} \right)^{1.6} (2.44) \tag{125}$$

$$P_H = 4.1 \times 10^{-6} \text{ watts} \tag{126}$$

which as in the case of the reed is negligible.

SUMMARY AND RECOMMENDATIONS

(1) Under the present conditions of operation the fluidic generator is producing only about 35% of its potential power output. Doubling the cross-section of the reed by adding lamination or increasing its width would augment its flux carrying capability to where the potential power output could be fully realized with only negligible eddy current and hysteresis losses in the reed. Appropriate adjustments in gap length would also be necessary to prevent saturation and to allow sufficient latitude for vibration.

(2) Keeper hysteresis losses are also negligible but eddy current losses can be appreciable unless the keeper is laminated or made of a suitable composite material. Further, the narrowest part of the keeper cross section must be increased by about 7%, to avoid saturation with the given reed displacement amplitude of .017".

(3) To facilitate design optimization of future generators a computer study was done to obtain a series of design matrices which list voltage amplitude as a function of gap length (l_g), reed thickness (τ), and reed displacement amplitude (a). The computer program uses the mathematical procedure of the sample calculation made in this report and it is written in Fortran IV for the Burroughs 5500 Computer. The program is reproduced in Appendix II together with a summary of the equations used in the calculations.

The computed voltage amplitudes are listed in Table II in Appendix II. All of the listed voltages are calculated for a 1500 turn coil, a reed width of .27", a reed vibrational frequency of 1500 hertz, and the standard magnet and keeper dimensions used in the sample calculation and summarized in Table I. The assumption is also made that the keeper does not saturate for any of the parameter choices listed. This assumption, as we have shown in the sample, is not valid for the larger values of (a) unless the keeper thickness is increased by an appropriate amount. This thickening of the keeper would, of course, affect the value of P_{y4} and hence the permeance calculations. The effect on the calculated voltage output, however, is not likely to be significant because in the sample calculation P_{y4} is only of the order of 10% of P_4 , and P_4 in turn constitutes only about 40% of P_t . Furthermore, P_{y4} itself is relatively insensitive to keeper thickness (W) since W appears only in a small correction to the logarithmic factor in Eq. (13). Hence, thickening of the keeper to required values will probably not affect the calculated voltages by more than the order of a percent or two even in the worst cases.

The calculations also assume infinite reed permeability and the values not actually attainable with the reed material used are marked with an asterisk. Thus the design can easily be optimized by choosing the parameters yielding the largest non-asterisked voltage in the Tables. This, of course, is always subject to the condition that the reed of requisite thickness can be driven with the stated amplitude and frequency. Perhaps this can be assured in many cases by making appropriate adjustments of the reed stiffness through variation of the reed material and dimensions at the point where it is anchored or by widening the reed rather than thickening it.

The Tables can be used to obtain the power output in any load resistance R_L from Eq. (126), viz.

$$P_L = \frac{V_{\max}^2 R_L}{2Z^2} \quad (127)$$

where Z is the total circuit impedance obtained from Eqs. (109) and (110). The coil inductance is a function of reed thickness and the empirical curve of the relationship is shown in Fig. 11. As is evident from (111), however, the reactive portion of the impedance is only about a percent and a half of the total, and the variation of reactance with reed thickness will not have a significant effect upon power output at the frequency and resistance used.

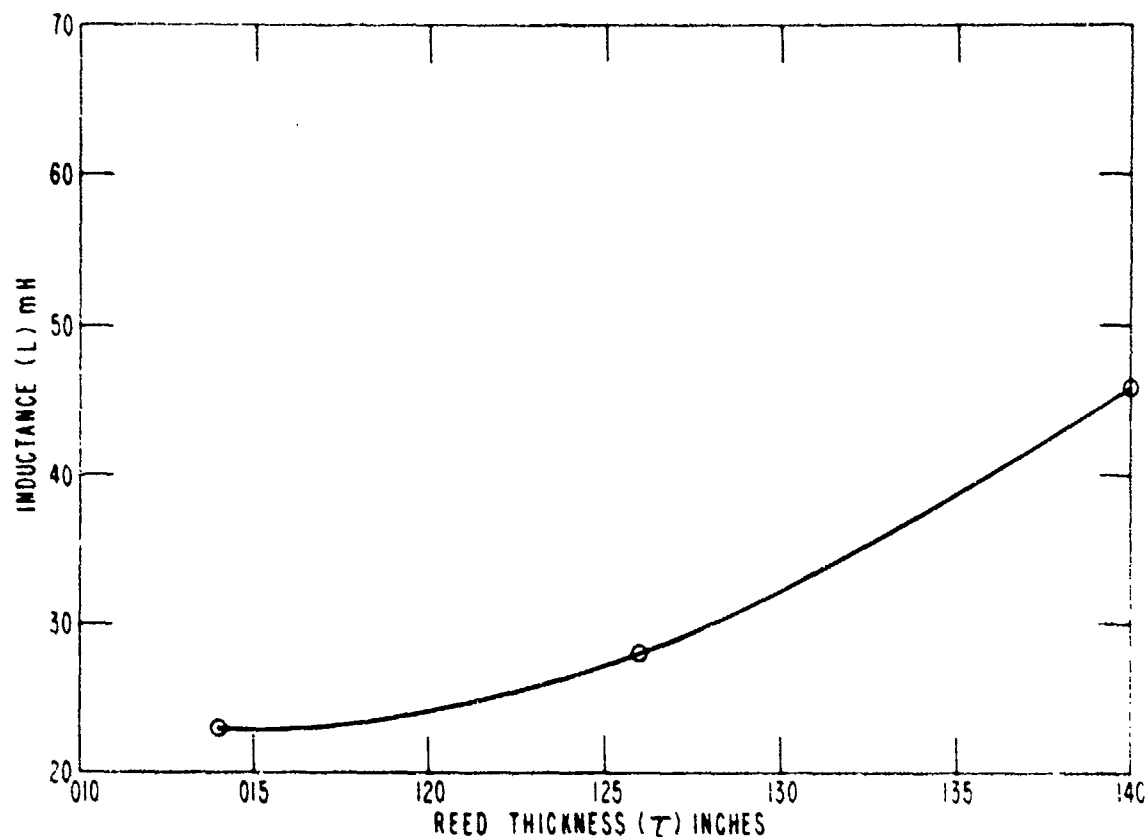


Fig. 11. Inductance of Sensing Coil as a Function of Reed Thickness.

(4) The present configuration of the magnetic circuit results in relatively low-sloped load lines which intersect the Alnico 5 demagnetization curve well below its 'knee' (Fig. 9). Note that if Co_5Sm were used in place of Alnico 5 much higher flux densities would result for any given load line with a concomitant increase in power output. This, of course, is always assuming that neither reed nor keeper are saturated. It probably would be most practical to keep the power output at about 10-12 watts with the use of much less magnetic material through substitution of Co_5Sm . A more compact generator operating at lower load lines with the same power output could probably be designed using the latter material.

(5) A composite material consisting of permalloy particles densely packed in a nonconducting binder might be desirable for the keeper. Providing a sufficiently high packing density is obtainable, eddy current losses could be virtually eliminated without too great a loss in flux-carrying capability. Such a composite could also be press molded into the desired shape thus improving the efficiency of fabrication.

APPENDIX I

DERIVATION OF THE FORMULA FOR ENERGY LOSS THROUGH EDDY CURRENTS

Consider a sheet of metal such as the one pictured in Fig. 12. The length l is much larger than the thickness τ and the width w is arbitrary. A sinusoidal field B is acting in the y direction as shown. The varying field B will induce EMFs which will give rise to circulating or eddy currents. We wish to find the power expended by these currents. We proceed by considering a current loop indicated by the arrows in Fig. 12. The EMF around this loop is then given by

$$V = \frac{d\Phi}{10^8 dt} = \frac{A dB}{10^8 dt} = \frac{2lx}{10^8} \frac{dB}{dt} \quad (128)$$

If we neglect the small voltage drops across the width of the loop at the ends we can write for the voltage drop per unit length in the z direction

$$v = \frac{V}{2l} = \frac{x dB}{10^8 dt} \quad (129)$$

The current density i is then obtained from the relation

$$i = \frac{v}{\rho} \quad (130)$$

where ρ is the resistivity of the material. Substituting (129) in (130) yields

$$i = \frac{x}{10^8 \rho} \frac{dB}{dt} = \frac{x B_{\max}}{\rho 10^8} \cos \omega t \quad (131)$$

The instantaneous power, dP , dissipated in the unit lamina element dx is then

$$dP = w \rho i^2 dx = \frac{w x^2 \omega^2 B_{\max}^2 dx}{\rho 10^{16}} \cos^2 \omega t$$

$$dP = \frac{w \omega^2 B_{\max}^2}{10^{16} \rho} x^2 dx \cos^2 \omega t \quad (132)$$

and the total power loss, P , of a unit length of sheet is

$$P = \int_{-\frac{\tau}{2}}^{\frac{\tau}{2}} dP = \frac{w \omega^2 B_{\max}^2}{3 \rho 10^{16}} x^3 \Big|_{-\frac{\tau}{2}}^{\frac{\tau}{2}} \cos^2 \omega t \quad (133)$$

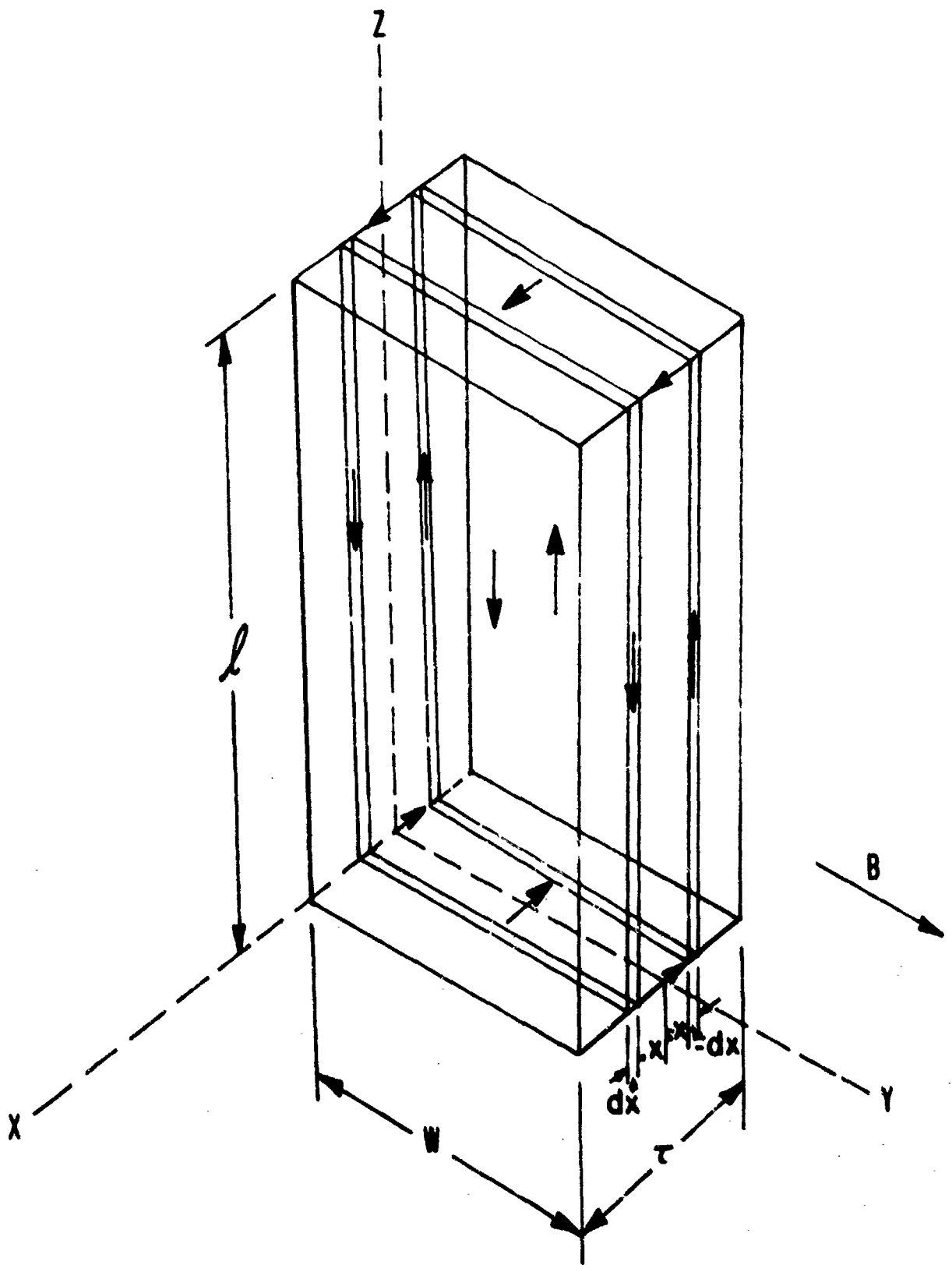


Fig. 12. Derivation of the Eddy Current Loss Formula.

$$P = \frac{w w_B^2 \tau^3}{3 \rho 10^{16} 4} \cos^2 \omega t = \frac{w w_B^2 \tau^3}{12 \rho 10^{16}} \cos^2 \omega t \quad (134)$$

The energy loss per cycle is then given by

$$w_c = \frac{w w_B^2 \tau^3}{12 \rho 10^{16}} \int_0^{2\pi} \cos^2(\omega t) (\omega dt) \quad (135)$$

$$w_c = \frac{\pi w w_B^2 \tau^3}{12 \rho 10^{16}} \quad (136)$$

The average power per unit length of sheet is

$$\bar{P} = f w_c = \frac{f \pi w w_B^2 \tau^3}{12 \rho 10^{16}} = \frac{\pi^2 w f^2 \tau^3}{6 \rho 10^{16}} \quad (137)$$

The average power for the entire length of the sheet is then

$$\bar{P}_c = l \bar{P} = \frac{\pi^2 f^2 \tau^3 B^2 l w}{6 \rho 10^{16}} = \frac{\pi^2 f^2 \tau^3 B^2 V}{6 \rho 10^{16}} \quad (138)$$

where V is the volume of the sheet.

APPENDIX II

SUMMARY OF COMPUTATIONAL PROCEDURE IN CALCULATION OF PEAK VOLTAGE AND ADAPTATION TO COMPUTER PROGRAM

All relevant permeances are calculated for the generator with the reed in the equilibrium position. There are five such permeances of which four are associated with the four air gaps and the fifth with the magnetic leakage paths. The two gaps at the fixed end of the reed are always equal and the associated permeances are equal and denoted by P_3 . P_1 and P_2 refer to the gaps at the vibrating end of the reed. P_4 is the permeance of the composite of all leakage paths. P_3 has five components and we can write for the equilibrium condition:

$$P_3 = P_{G3} + P_{E3} + P_{F3} + P_{L3} + P_{W3}$$

all of which are illustrated in Fig. 5 and are computed by the following formulae as functions of l_g , t and a . All other dimensions are fixed and are the same as those of the sample calculations.

$$P_{G3} = \frac{.0408}{l_g - t}$$

$$P_{E3} = .0364$$

$$P_{F3} = .303$$

$$P_{L3} = .215 \ln \left(1 + \frac{.30}{l_g - t} \right) + [.213 - .0916(l_g - t)] \ln \left(\frac{.30}{l_g - t} \right)$$

$$P_{W3} = .0445 \ln \left(1 + \frac{1.73 t}{l_g} \right)$$

P_4 has five components and is given by

$$P_4 = P_{p4} + P_{y4} + P_{M4} + P_{L4} + P_{W4}$$

and

$$P_{p4} = .647''$$

$$P_{y4} = .212''$$

$$P_{M4} = .731''$$

$$P_{L4} = .136''$$

$$P_{W4} = .134 \ln \frac{.4 + 1.16 t}{t + 1.16 t \frac{g}{l_g}}$$

P_1 and P_2 have the same form with either always reaching its maximum when the other is at its minimum. Each has 5 components and

$$P_1 = P_{G1} + P_{E1} + P_{F1} + P_{L1} + P_{W1}$$

$$P_{G1 \min}^{\max} = P_{G2 \min}^{\max} = \frac{.0408}{l_g - t \mp 2a}$$

$$P_{E1} = P_{E2} = .0364''$$

$$P_{F1} = P_{F2} = .303''$$

$$P_{L1 \min}^{\max} = P_{L2 \min}^{\max} = .215 \ln \left(1 + \frac{.30}{l_g - t \mp 2a} \right)$$

$$+ [.170 - .0918(l_g - t \mp 2a)] \frac{.30}{l_g - t \mp 2a}$$

$$P_{W1 \min}^{\max} = P_{W2 \min}^{\max} = .0445 \ln \left(1 + \frac{1.73 t}{l_g - t \mp 2a} \right)$$

The total permeance of the magnetic circuit P_t is then given by the formula:

$$P_t = P_4 + \frac{(P_1 + P_3)(P_2 + P_3)}{P_1 + P_2 + 2P_3}$$

and both the equilibrium and full reed displacement values of P_t are calculated. We use the equilibrium value P_t^E to obtain our load line slope B/H via the formula

$$\left(\frac{B}{H} \right)^E = - \frac{l}{A} P_t^E$$

where l and A are respectively the length and cross-sectional area of the magnet.

The intersection of this load line with the demagnetization curve of Alnico 5 is the base point of the recoil line (slope 4.3 for Alnico 5) along which the system then operates. The recoil line is thus of the form:

$$B = 4.3 H + B_0$$

where B_0 is determined from the base point. The maximum value of P_t corresponding to full displacement is then used in the relation

$$\left(\frac{B}{H} \right)^{\max} = - \frac{l}{A} P_t^{\max}$$

to obtain the new load line. The intersection of this load line with the recoil line already determined yields the maximum value of B in the magnet.

The corresponding total available flux is then given by

$$\Phi_t^{\max} = B^{\max} A_{\text{magnet}}$$

and the flux amplitude in the reed can then be calculated from the relation

$$\Phi_R^{\max} = \Phi_t^{\max} \frac{P_3(P_1 - P_2)}{(P_1 + P_3)(P_2 + P_3) + P_4(P_1 + 2P_3 + P_2)}$$

The induced voltage amplitude is then obtained from

$$V^{\max} = \frac{N\omega \Phi_R^{\max}}{10^8}$$

The computer program which follows this procedure is reproduced below in Fortran IV. In addition to calculating the expected voltage for a given set of parameters, the program also instructs the computer to calculate

$$B_R^{\max} = \frac{\Phi_R^{\max}}{10^8}$$

If the result is larger than the saturation flux density of the reed material (19,900 gauss), the calculated voltage is not attainable and is marked with an asterisk in Table I.

COMPUTER PROGRAM USED TO OBTAIN TABLE II

```

FILE 1=FILE1,UNIT=PRINTER
REAL MAX,PM4,PM5,PP4,PP5,PYA,PL4,PA,PG3,PF3,PF3,PI 3,PM3,P3,PTF,
=
  K,M,P,PL,LG
DIMENSION K(6),R(6),P(3),F(3),VOLT(35),RFFEN(35),ASTK(35)
DATA (R(1),Y(1),A)/1000,0000,0000,10000,11000,12000/
- (K(1),Y(1),A)/7617.,7543.,7546.,7420.,7817.,7200/
Z=1210.
RI=600.
TIMNS=1500.
NMEGAT=300003.1415926
PI 110 200.000,001,005
WRITE(1,900) IG
500 FORMAT(17,15(1X,/),1Y,"GENERATOR OUTPUT IN VOLTS AC FOR GAP "
- "LENGTH = ",F5.3,2Y," INCHES",//)
WRITE(1,907)
97 FORMAT(1Y,"RFFEN",3Y,/,1Y,"TIME PER FISS",15Y,
- "AVERAGE TIME IN INCHES",/,1Y,"INCHES",1X,".001,002,003,004"
- ".005,006,007,008,009,010,011,012,013,014,015,016,017,018,019"
- ".020,021,022,023,024,025,026,027,028,029",/3
PI 110 10.01,100,005,002
N=0
YN=(L0+1)/2,=.001
IF(YA.GT,0.0291) Y1=0.020
PI 120 10.001,YN,001
N=4
ASTK(1)=0
PM4=.1330*ALOG((.0+1.155*IG)/(Y+1.155*IG))
PM4=.7310
PP4=.4450
PY4=.2170
PL4=.1340
PA=PP4+PM4+PP4+PY4+PL4
Z=LG*Y
PG3=.00070/7
PF3=.0340
FF3=.7070
PR=.215*ALOG(1.+30/Z)
C=(.2120+.0014*Z)*ALOG(.30/7)
PI 3=PR+C
PM3=.0440*ALOG(1.+(1.732*Y)/7)
P3=PG3+PF3+FF3+PL3+PM3
F(1)=IG*Y-2.*A
F(2)=IG*Y+2.*A
F(3)=IG*Y
PI 35 1F,3
G=.3303+.04450*ALOG(1.+(1.732*Y)/7)/F(1)+.04070/F(1)
R=.215*ALOG(1.+3/F(1))
P(1)=R+R*(.1400+.0016*F(Y))+ALOG(.3/F(Y))
35 CONTINUE
PTE=PR+P4/2.+P(3)/2.
RPHF=-1.172*PTE
K=ARS(BDWE)
IF(K.LT.04.18) GO TO 5
IF(K.LE.04.18 .AND. K.LT.16.19) GO TO 10
IF(K.LE.04.19 .AND. K.LT.19.23) GO TO 15
IF(K.LE.04.23 .AND. K.LT.26.30) GO TO 20
IF(K.LE.04.30 .AND. K.LT.40.) GO TO 25
I=7
GO TO 50
5 I=2
GO TO 50
10 I=3

```


TABLE II

Voltage outputs for various gap lengths, reed thicknesses and reed displacement amplitudes. The values marked with an asterisk are not attainable with the reed material presently in use. They are the values to be expected with the use of reed material that does not saturate under the given operating conditions.

GENERATED OUTPUT IN VOLTS AC FOR GAP LENGTH = 0.040 INCHES

REED THICKNESS (INCHES)	AMPLITUDE IN INCHES	OUTPUT IN VOLTS AC
.010	.25	39
.010	.30	51*
.010	.35	67*
.010	.40	87*
.010	.45	107*
.010	.50	128*
.010	.55	143*
.010	.60	158*
.010	.65	170*
.010	.70	178*
.010	.75	183*
.010	.80	187*
.010	.85	190*
.010	.90	192*
.010	.95	193*
.010	1.00	193*
.012	.25	44
.012	.30	58*
.012	.35	76*
.012	.40	98*
.012	.45	124*
.012	.50	154*
.012	.55	188*
.012	.60	226*
.012	.65	268*
.012	.70	315*
.012	.75	367*
.012	.80	424*
.012	.85	486*
.012	.90	552*
.012	.95	622*
.012	1.00	696*
.014	.25	51
.014	.30	67*
.014	.35	87*
.014	.40	111*
.014	.45	140*
.014	.50	174*
.014	.55	213*
.014	.60	257*
.014	.65	306*
.014	.70	360*
.014	.75	419*
.014	.80	483*
.014	.85	552*
.014	.90	626*
.014	.95	705*
.014	1.00	789*
.016	.25	59
.016	.30	77*
.016	.35	99*
.016	.40	126*
.016	.45	159*
.016	.50	198*
.016	.55	243*
.016	.60	294*
.016	.65	351*
.016	.70	414*
.016	.75	483*
.016	.80	558*
.016	.85	639*
.016	.90	726*
.016	.95	819*
.016	1.00	918*
.018	.25	67
.018	.30	87*
.018	.35	111*
.018	.40	140*
.018	.45	174*
.018	.50	213*
.018	.55	257*
.018	.60	306*
.018	.65	360*
.018	.70	419*
.018	.75	483*
.018	.80	552*
.018	.85	626*
.018	.90	705*
.018	.95	789*
.018	1.00	878*
.020	.25	77
.020	.30	99*
.020	.35	126*
.020	.40	159*
.020	.45	198*
.020	.50	243*
.020	.55	294*
.020	.60	351*
.020	.65	414*
.020	.70	483*
.020	.75	558*
.020	.80	639*
.020	.85	726*
.020	.90	819*
.020	.95	918*
.020	1.00	1023*
.022	.25	87
.022	.30	111*
.022	.35	140*
.022	.40	174*
.022	.45	213*
.022	.50	257*
.022	.55	306*
.022	.60	360*
.022	.65	419*
.022	.70	483*
.022	.75	552*
.022	.80	626*
.022	.85	705*
.022	.90	789*
.022	.95	878*
.022	1.00	973*
.024	.25	97
.024	.30	126*
.024	.35	159*
.024	.40	198*
.024	.45	243*
.024	.50	294*
.024	.55	351*
.024	.60	414*
.024	.65	483*
.024	.70	552*
.024	.75	626*
.024	.80	705*
.024	.85	789*
.024	.90	878*
.024	.95	973*
.024	1.00	1074*
.026	.25	107
.026	.30	138*
.026	.35	174*
.026	.40	213*
.026	.45	257*
.026	.50	306*
.026	.55	360*
.026	.60	419*
.026	.65	483*
.026	.70	552*
.026	.75	626*
.026	.80	705*
.026	.85	789*
.026	.90	878*
.026	.95	973*
.026	1.00	1074*
.028	.25	117
.028	.30	150*
.028	.35	187*
.028	.40	226*
.028	.45	270*
.028	.50	319*
.028	.55	374*
.028	.60	435*
.028	.65	501*
.028	.70	570*
.028	.75	642*
.028	.80	718*
.028	.85	798*
.028	.90	882*
.028	.95	970*
.028	1.00	1062*
.030	.25	127
.030	.30	163*
.030	.35	201*
.030	.40	243*
.030	.45	290*
.030	.50	342*
.030	.55	399*
.030	.60	462*
.030	.65	530*
.030	.70	603*
.030	.75	681*
.030	.80	764*
.030	.85	852*
.030	.90	945*
.030	.95	1042*
.030	1.00	1144*
.032	.25	137
.032	.30	176*
.032	.35	217*
.032	.40	261*
.032	.45	310*
.032	.50	364*
.032	.55	423*
.032	.60	488*
.032	.65	558*
.032	.70	633*
.032	.75	714*
.032	.80	800*
.032	.85	892*
.032	.90	989*
.032	.95	1092*
.032	1.00	1200*
.034	.25	147
.034	.30	188*
.034	.35	231*
.034	.40	277*
.034	.45	328*
.034	.50	384*
.034	.55	446*
.034	.60	513*
.034	.65	586*
.034	.70	664*
.034	.75	748*
.034	.80	837*
.034	.85	932*
.034	.90	1033*
.034	.95	1140*
.034	1.00	1254*

TABLE II (cont'd)

GENERATOR JUMPING IN VOLTAGE FOR END CAP (FACTORY @ 0.050 INCHES)

OFFICE NUMBER	001	002	003	004	005	006	007	008	009	010	011	012	013	014	015	016	017	018	019	020	021	022	023	024	025	026	027	028	029	
017	0	15	23	31	39	48	57	67	77	88	100	114	129	145	161	177	193	209	225	241	257	273	289	305	321	337	353	369	385	
018	0	10	20	30	40	50	60	70	80	90	100	110	120	130	140	150	160	170	180	190	200	210	220	230	240	250	260	270	280	290
019	0	10	20	30	40	50	60	70	80	90	100	110	120	130	140	150	160	170	180	190	200	210	220	230	240	250	260	270	280	290
020	0	10	20	30	40	50	60	70	80	90	100	110	120	130	140	150	160	170	180	190	200	210	220	230	240	250	260	270	280	290
021	0	10	20	30	40	50	60	70	80	90	100	110	120	130	140	150	160	170	180	190	200	210	220	230	240	250	260	270	280	290
022	0	10	20	30	40	50	60	70	80	90	100	110	120	130	140	150	160	170	180	190	200	210	220	230	240	250	260	270	280	290
023	0	10	20	30	40	50	60	70	80	90	100	110	120	130	140	150	160	170	180	190	200	210	220	230	240	250	260	270	280	290
024	0	10	20	30	40	50	60	70	80	90	100	110	120	130	140	150	160	170	180	190	200	210	220	230	240	250	260	270	280	290
025	0	10	20	30	40	50	60	70	80	90	100	110	120	130	140	150	160	170	180	190	200	210	220	230	240	250	260	270	280	290
026	0	10	20	30	40	50	60	70	80	90	100	110	120	130	140	150	160	170	180	190	200	210	220	230	240	250	260	270	280	290
027	0	10	20	30	40	50	60	70	80	90	100	110	120	130	140	150	160	170	180	190	200	210	220	230	240	250	260	270	280	290
028	0	10	20	30	40	50	60	70	80	90	100	110	120	130	140	150	160	170	180	190	200	210	220	230	240	250	260	270	280	290
029	0	10	20	30	40	50	60	70	80	90	100	110	120	130	140	150	160	170	180	190	200	210	220	230	240	250	260	270	280	290
030	0	10	20	30	40	50	60	70	80	90	100	110	120	130	140	150	160	170	180	190	200	210	220	230	240	250	260	270	280	290
031	0	10	20	30	40	50	60	70	80	90	100	110	120	130	140	150	160	170	180	190	200	210	220	230	240	250	260	270	280	290
032	0	10	20	30	40	50	60	70	80	90	100	110	120	130	140	150	160	170	180	190	200	210	220	230	240	250	260	270	280	290
033	0	10	20	30	40	50	60	70	80	90	100	110	120	130	140	150	160	170	180	190	200	210	220	230	240	250	260	270	280	290
034	0	10	20	30	40	50	60	70	80	90	100	110	120	130	140	150	160	170	180	190	200	210	220	230	240	250	260	270	280	290
035	0	10	20	30	40	50	60	70	80	90	100	110	120	130	140	150	160	170	180	190	200	210	220	230	240	250	260	270	280	290
036	0	10	20	30	40	50	60	70	80	90	100	110	120	130	140	150	160	170	180	190	200	210	220	230	240	250	260	270	280	290
037	0	10	20	30	40	50	60	70	80	90	100	110	120	130	140	150	160	170	180	190	200	210	220	230	240	250	260	270	280	290
038	0	10	20	30	40	50	60	70	80	90	100	110	120	130	140	150	160	170	180	190	200	210	220	230	240	250	260	270	280	290

TABLE II (cont'd)

GENERATOR OUTPUT IN VOLTS AC FOR GAP FACTOR OF 0.040 INCHES

REFLECTOR POSITION	APPLYING 10 INCHES	APPLYING 15 INCHES	APPLYING 20 INCHES	APPLYING 25 INCHES	APPLYING 30 INCHES	APPLYING 35 INCHES	APPLYING 40 INCHES	APPLYING 45 INCHES	APPLYING 50 INCHES	APPLYING 55 INCHES	APPLYING 60 INCHES	APPLYING 65 INCHES	APPLYING 70 INCHES	APPLYING 75 INCHES	APPLYING 80 INCHES	APPLYING 85 INCHES	APPLYING 90 INCHES	APPLYING 95 INCHES	APPLYING 100 INCHES	
0.010	5	10	15	20	25	30	35	40	45	50	55	60	65	70	75	80	85	90	95	100
0.012	4	10	17	24	31	38	45	52	59	66	73	80	87	94	101	108	115	122	129	136
0.014	3	10	18	26	34	42	50	58	66	74	82	90	98	106	114	122	130	138	146	154
0.016	4	12	21	30	39	48	57	66	75	84	93	102	111	120	129	138	147	156	165	174
0.018	5	14	24	34	44	54	64	74	84	94	104	114	124	134	144	154	164	174	184	194
0.020	6	16	27	38	49	60	71	82	93	104	115	126	137	148	159	170	181	192	203	214
0.022	7	18	30	42	54	66	78	90	102	114	126	138	150	162	174	186	198	210	222	234
0.024	8	20	33	46	59	72	85	98	111	124	137	150	163	176	189	202	215	228	241	254
0.026	9	22	36	50	64	78	92	106	120	134	148	162	176	190	204	218	232	246	260	274
0.028	10	24	39	54	69	84	99	114	129	144	159	174	189	204	219	234	249	264	279	294
0.030	11	26	42	58	74	90	106	122	138	154	170	186	202	218	234	250	266	282	298	314
0.032	12	28	46	63	80	97	114	131	148	165	182	199	216	233	250	267	284	301	318	335
0.034	13	30	50	68	86	104	122	140	158	176	194	212	230	248	266	284	302	320	338	356
0.036	14	32	54	73	92	111	130	149	168	187	206	225	244	263	282	301	320	339	358	377
0.038	15	34	58	78	98	118	138	158	178	198	218	238	258	278	298	318	338	358	378	398
0.040	16	36	62	83	104	125	146	167	188	209	230	251	272	293	314	335	356	377	398	419
0.042	17	38	66	88	110	132	154	176	198	220	242	264	286	308	330	352	374	396	418	440
0.044	18	40	70	93	116	139	162	185	208	231	254	277	300	323	346	369	392	415	438	461
0.046	19	42	74	98	122	146	170	194	218	242	266	290	314	338	362	386	410	434	458	482
0.048	20	44	78	103	128	153	178	203	228	253	278	303	328	353	378	403	428	453	478	503
0.050	21	46	82	108	134	160	186	212	238	264	290	316	342	368	394	420	446	472	498	524
0.052	22	48	86	113	140	167	194	221	248	275	302	329	356	383	410	437	464	491	518	545
0.054	23	50	90	118	146	174	202	230	258	286	314	342	370	398	426	454	482	510	538	566
0.056	24	52	94	123	152	181	210	239	268	297	326	355	384	413	442	471	500	529	558	587

TABLE II (cont'd)

GENERATOR OUTPUT IN VOLTS AC FOR GAP FACTOR = 0.075 INCHES

REFC THICKNESS (INCHES)	AMPLITUDE IN INCHES																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																									
.013	3	7	10	13	17	20	24	27	31	37	39	43	47	51	56	60	64	68	72	77	81	85	89	93	97	101	105	109	113	117	121	125	129	133	137	141	145	149	153	157	161	165	169	173	177	181	185	189	193	197	201	205	209	213	217	221	225	229	233	237	241	245	249	253	257	261	265	269	273	277	281	285	289	293	297	301	305	309	313	317	321	325	329	333	337	341	345	349	353	357	361	365	369	373	377	381	385	389	393	397	401	405	409	413	417	421	425	429	433	437	441	445	449	453	457	461	465	469	473	477	481	485	489	493	497	501	505	509	513	517	521	525	529	533	537	541	545	549	553	557	561	565	569	573	577	581	585	589	593	597	601	605	609	613	617	621	625	629	633	637	641	645	649	653	657	661	665	669	673	677	681	685	689	693	697	701	705	709	713	717	721	725	729	733	737	741	745	749	753	757	761	765	769	773	777	781	785	789	793	797	801	805	809	813	817	821	825	829	833	837	841	845	849	853	857	861	865	869	873	877	881	885	889	893	897	901	905	909	913	917	921	925	929	933	937	941	945	949	953	957	961	965	969	973	977	981	985	989	993	997	1001	1005	1009	1013	1017	1021	1025	1029	1033	1037	1041	1045	1049	1053	1057	1061	1065	1069	1073	1077	1081	1085	1089	1093	1097	1101	1105	1109	1113	1117	1121	1125	1129	1133	1137	1141	1145	1149	1153	1157	1161	1165	1169	1173	1177	1181	1185	1189	1193	1197	1201	1205	1209	1213	1217	1221	1225	1229	1233	1237	1241	1245	1249	1253	1257	1261	1265	1269	1273	1277	1281	1285	1289	1293	1297	1301	1305	1309	1313	1317	1321	1325	1329	1333	1337	1341	1345	1349	1353	1357	1361	1365	1369	1373	1377	1381	1385	1389	1393	1397	1401	1405	1409	1413	1417	1421	1425	1429	1433	1437	1441	1445	1449	1453	1457	1461	1465	1469	1473	1477	1481	1485	1489	1493	1497	1501	1505	1509	1513	1517	1521	1525	1529	1533	1537	1541	1545	1549	1553	1557	1561	1565	1569	1573	1577	1581	1585	1589	1593	1597	1601	1605	1609	1613	1617	1621	1625	1629	1633	1637	1641	1645	1649	1653	1657	1661	1665	1669	1673	1677	1681	1685	1689	1693	1697	1701	1705	1709	1713	1717	1721	1725	1729	1733	1737	1741	1745	1749	1753	1757	1761	1765	1769	1773	1777	1781	1785	1789	1793	1797	1801	1805	1809	1813	1817	1821	1825	1829	1833	1837	1841	1845	1849	1853	1857	1861	1865	1869	1873	1877	1881	1885	1889	1893	1897	1901	1905	1909	1913	1917	1921	1925	1929	1933	1937	1941	1945	1949	1953	1957	1961	1965	1969	1973	1977	1981	1985	1989	1993	1997	2001	2005	2009	2013	2017	2021	2025	2029	2033	2037	2041	2045	2049	2053	2057	2061	2065	2069	2073	2077	2081	2085	2089	2093	2097	2101	2105	2109	2113	2117	2121	2125	2129	2133	2137	2141	2145	2149	2153	2157	2161	2165	2169	2173	2177	2181	2185	2189	2193	2197	2201	2205	2209	2213	2217	2221	2225	2229	2233	2237	2241	2245	2249	2253	2257	2261	2265	2269	2273	2277	2281	2285	2289	2293	2297	2301	2305	2309	2313	2317	2321	2325	2329	2333	2337	2341	2345	2349	2353	2357	2361	2365	2369	2373	2377	2381	2385	2389	2393	2397	2401	2405	2409	2413	2417	2421	2425	2429	2433	2437	2441	2445	2449	2453	2457	2461	2465	2469	2473	2477	2481	2485	2489	2493	2497	2501	2505	2509	2513	2517	2521	2525	2529	2533	2537	2541	2545	2549	2553	2557	2561	2565	2569	2573	2577	2581	2585	2589	2593	2597	2601	2605	2609	2613	2617	2621	2625	2629	2633	2637	2641	2645	2649	2653	2657	2661	2665	2669	2673	2677	2681	2685	2689	2693	2697	2701	2705	2709	2713	2717	2721	2725	2729	2733	2737	2741	2745	2749	2753	2757	2761	2765	2769	2773	2777	2781	2785	2789	2793	2797	2801	2805	2809	2813	2817	2821	2825	2829	2833	2837	2841	2845	2849	2853	2857	2861	2865	2869	2873	2877	2881	2885	2889	2893	2897	2901	2905	2909	2913	2917	2921	2925	2929	2933	2937	2941	2945	2949	2953	2957	2961	2965	2969	2973	2977	2981	2985	2989	2993	2997	3001	3005	3009	3013	3017	3021	3025	3029	3033	3037	3041	3045	3049	3053	3057	3061	3065	3069	3073	3077	3081	3085	3089	3093	3097	3101	3105	3109	3113	3117	3121	3125	3129	3133	3137	3141	3145	3149	3153	3157	3161	3165	3169	3173	3177	3181	3185	3189	3193	3197	3201	3205	3209	3213	3217	3221	3225	3229	3233	3237	3241	3245	3249	3253	3257	3261	3265	3269	3273	3277	3281	3285	3289	3293	3297	3301	3305	3309	3313	3317	3321	3325	3329	3333	3337	3341	3345	3349	3353	3357	3361	3365	3369	3373	3377	3381	3385	3389	3393	3397	3401	3405	3409	3413	3417	3421	3425	3429	3433	3437	3441	3445	3449	3453	3457	3461	3465	3469	3473	3477	3481	3485	3489	3493	3497	3501	3505	3509	3513	3517	3521	3525	3529	3533	3537	3541	3545	3549	3553	3557	3561	3565	3569	3573	3577	3581	3585	3589	3593	3597	3601	3605	3609	3613	3617	3621	3625	3629	3633	3637	3641	3645	3649	3653	3657	3661	3665	3669	3673	3677	3681	3685	3689	3693	3697	3701	3705	3709	3713	3717	3721	3725	3729	3733	3737	3741	3745	3749	3753	3757	3761	3765	3769	3773	3777	3781	3785	3789	3793	3797	3801	3805	3809	3813	3817	3821	3825	3829	3833	3837	3841	3845	3849	3853	3857	3861	3865	3869	3873	3877	3881	3885	3889	3893	3897	3901	3905	3909	3913	3917	3921	3925	3929	3933	3937	3941	3945	3949	3953	3957	3961	3965	3969	3973	3977	3981	3985	3989	3993	3997	4001	4005	4009	4013	4017	4021	4025	4029	4033	4037	4041	4045	4049	4053	4057	4061	4065	4069	4073	4077	4081	4085	4089	4093	4097	4101	4105	4109	4113	4117	4121	4125	4129	4133	4137	4141	4145	4149	4153	4157	4161	4165	4169	4173	4177	4181	4185	4189	4193	4197	4201	4205	4209	4213	4217	4221	4225	4229	4233	4237	4241	4245	4249	4253	4257	4261	4265	4269	4273	4277	4281	4285	4289	4293	4297	4301	4305	4309	4313	4317	4321	4325	4329	4333	4337	4341	4345	4349	4353	4357	4361	4365	4369	4373	4377	4381	4385	4389	4393	4397	4401	4405	4409	4413	4417	4421	4425	4429	4433	4437	4441	4445	4449	4453	4457	4461	4465	4469	4473	4477	4481	4485	4489	4493	4497	4501	4505	4509	4513	4517	4521	4525	4529	4533	4537	4541	4545	4549	4553	4557	4561	4565	4569	4573	4577	4581	4585	4589	4593	4597	4601	4605	4609	4613	4617	4621	4625	4629	4633	4637	4641	4645	4649	4653	4657	4661	4665	4669	4673	4677	4681	4685	4689	4693	4697	4701	4705	4709	4713	4717	4721	4725	4729	4733	4737	4741	4745	4749	4753	4757	4761	4765	4769	4773	4777	4781	4785	4789	4793	4797	4801	4805	4809	4813	4817	4821	4825	4829	4833	4837	4841	4845	4849	4853	4857	4861	4865	4869	4873	4877	4881	4885	4889	4893	4897	4901	4905	4909	4913	4917	4921	4925	4929	4933	4937	4941	4945	4949	4953	4957	4961	4965	4969	4973	4977	4981	4985	4989	4993	4997	5001	5005	5009	5013	5017	5021	5025	5029	5033	5037	5041	5045	5049	5053	5057	5061	5065	5069	5073	5077	5081	5085	5089	5093	5097	5101	5105	5109	5113	5117	5121	5125	5129	5133	5137	5141	5145	5149	5153	5157	5161	5165	5169	5173	5177	5181	5185	5189	5193	5197	5201	5205	5209	5213	5217	5221	5225	5229	5233	5237	5241	5245	5249	5253	5257	5261	5265	5269	5273	5277	5281	5285	5289	5293	5297	5301	5305	5309	5313	5317	5321	5325	5329	5333	5337	5341	5345	5349	5353	5357	5361	5365	5369	5373	5377	5381	5385	5389	5393	5397	5401	5405	5409	5413	5417	5421	5425	5429	5433	5437	5441	5445	5449	5453	5457	5461	5465	5469	5473	5477	5481	5485	5489	5493	5497	5501	5505	5509	5513	5517	5521	5525	5529	5533	5537	

