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## Reverse Breakdown in Large Area, Field-Plated, Vertical $\beta$ -Ga<sub>2</sub>O<sub>3</sub> Rectifiers

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There is interest in developing large area Ga<sub>2</sub>O<sub>3</sub> rectifiers for applications in hybrid power converters. Vertical geometry, Schottky rectifiers with area  $1.2 \times 1.2 \text{ mm}^2$  fabricated on thick ( $8 \mu\text{m}$ ), undoped ( $n = 4.4 \times 10^{15} \text{ cm}^{-3}$ )  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> epitaxial layers on N conducting bulk substrates exhibit both high forward current (1 A in pulsed mode) and reverse breakdown voltage ( $V_B = 760 \text{ V}$ ). This breakdown voltage was  $\sim 200 \text{ V}$  higher than rectifiers without the presence of a bilayer SiO<sub>2</sub>/SiN<sub>x</sub> field plate. This edge termination is critical for obtaining high breakdown voltage by reducing electric field crowding around the metal contact periphery. Optimization of the field plate design is still needed, since devices are observed experimentally to breakdown at the contact periphery. When purposely driven to failure at high reverse bias, pits are observed in the high field regions at the edge of the contact. The specific on-resistance ( $R_{\text{on}}$ ) for these large area rectifiers was  $22 \text{ m}\Omega\cdot\text{cm}^{-2}$ , with a figure-of-merit  $V_B^2/R_{\text{on}}$  of  $26 \text{ MW}\cdot\text{cm}^{-2}$ . The potential of Ga<sub>2</sub>O<sub>3</sub> for power electronics is clear when it is realized that these values are still an order of magnitude lower than theoretical values. The diode on-off ratio was in the range  $2.7 \times 10^7$ – $2.2 \times 10^9$  when switching from  $+1.5 \text{ V}$  forward bias to  $1$ – $100 \text{ V}$  reverse bias.

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Wide-bandgap (WBG) devices are promising candidates for next-generation power electronics converters with higher efficiency and higher power conversion densities.<sup>1–6</sup> These inverters have applications in a range of power conditioning and control systems, including pulsed power for avionics and electric ships, solid-state drivers for heavy electric motors and in advanced power management and control electronics.<sup>1–10</sup> In addition, high-power ( $\sim 50 \text{ kW}$ ) is required for fast wireless charging systems in transportation applications.<sup>11</sup> Highly efficient energy conversion for these systems is critical and this depends mainly on the ability of power switching transistors to provide low resistance in the on-state and highly resistive off-state conditions. To date, the focus has been on SiC and GaN to provide performance beyond Si.<sup>1–5</sup> High cost, challenging fabrication of practical devices, demonstrated reliability, and system integration remain important barriers to the widespread adoption of WBG devices.

A possible solution is hybrid switches, a combination of Si MOSFETs and wide bandgap devices. For example, hybrid modules containing Si MOSFETs and SiC rectifiers are commercially available.<sup>12</sup> One component of a switching module is a diode, most commonly Schottky rectifiers. These have fast switching speed, low forward voltage drop and high temperature operability.<sup>8–10</sup> Their advantage over p-n diodes is shorter switching times due to absence of minority carriers, but a disadvantage is higher on-state resistance ( $R_{\text{ON}}$ ). Employing material with a wider bandgap than Si improves rectifier performance, with lower on-state resistance at a given reverse voltage.<sup>1–5</sup> Recently, Ga<sub>2</sub>O<sub>3</sub> vertical geometry rectifiers have shown promising performance in terms of high reverse breakdown voltage ( $V_B > 1 \text{ kV}$ ) and low  $R_{\text{ON}}$ , leading to good power figure-of-merits ( $V_B^2/R_{\text{ON}}$ ).<sup>11,13–46</sup> Notable are reverse breakdown voltages ( $V_B$ ) of  $2300 \text{ V}$  for a  $150 \mu\text{m}$  diameter device (area =  $1.77 \times 10^{-4} \text{ cm}^2$ )<sup>47</sup> and  $2440 \text{ V}$  breakdown in trench structures.<sup>48</sup> Another potential advantage of Ga<sub>2</sub>O<sub>3</sub> are the lower substrate costs compared to SiC.<sup>32–42</sup> Reese et al.<sup>40</sup> used techno-economic modeling of Ga<sub>2</sub>O<sub>3</sub> wafer cost based on future large-size and high-volume manufacturing scenarios and estimated there could be a  $> 3\times$

cost advantage compared to SiC wafers. This could lead to a significant cost reduction for wide-bandgap power electronics.<sup>49–52</sup>

In current Ga<sub>2</sub>O<sub>3</sub> rectifiers, breakdown resulting from impact ionization preferentially occurs at the contact periphery unless the electric field is reduced by edge termination.<sup>5,27</sup> Baliga's figure of merit (FOM) ( $\epsilon_r \mu E_{\text{br}}^3$ ) where  $\epsilon_r$  is relative dielectric constant and  $\mu$  is mobility<sup>7,8,42–46</sup> is  $4\times$  larger for  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> than SiC and GaN.<sup>5–10</sup> The theoretical breakdown field for Ga<sub>2</sub>O<sub>3</sub> is between  $5$ – $9 \text{ MV/cm}$ , with peak experimental values of  $5.3 \text{ MV}\cdot\text{cm}^{-1}$ .<sup>31</sup> Lateral structures have shown breakdown voltages of over  $3 \text{ kV}$ .<sup>29</sup> The switching characteristics of Ga<sub>2</sub>O<sub>3</sub> vertical Schottky rectifier show recovery times of  $20$ – $30 \text{ ns}$ ,<sup>45,46,48,49,53</sup> faster than Si or SiC diodes.

Vertical geometry, large area planar devices can produce large total forward currents, while maintaining adequate reverse breakdown.<sup>45–58</sup> This is a stringent test of material quality, since large diodes increase the probability of incorporating defects into the active region, degrading reverse breakdown voltage.<sup>37,39,43,45</sup> Lateral devices can achieve high breakdown voltage using large contact separations, but cannot simultaneously achieve high forward current due to the low total conduction thickness. They also have high on-state resistance.<sup>45,48</sup>

In this paper, we show vertical geometry Schottky rectifiers with a minimum of process steps can achieve forward currents of  $1 \text{ A}$  at  $2.3 \text{ V}$  and reverse breakdown voltages of  $760 \text{ V}$  for large ( $0.14 \text{ cm}^2$ ) devices, emphasizing the potential of Ga<sub>2</sub>O<sub>3</sub> for majority carrier switching devices.

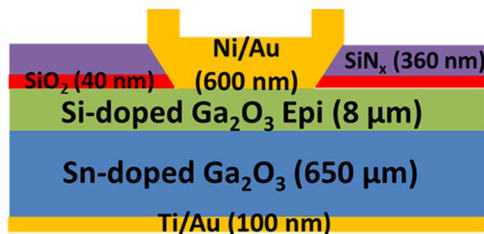
### Experimental

Epitaxial layers ( $\sim 8 \mu\text{m}$  final thickness) of lightly Si-doped n-type ( $4.36 \times 10^{15} \text{ cm}^{-3}$ ) Ga<sub>2</sub>O<sub>3</sub> were grown by Halide Vapor Phase Epitaxy (HVPE) on  $n^+$  ( $3.6 \times 10^{18} \text{ cm}^{-3}$ ),  $\beta$ -phase Sn-doped Ga<sub>2</sub>O<sub>3</sub> wafers ( $\sim 650 \mu\text{m}$  thick) with (001) surface orientation grown by the edge-defined film-fed method.<sup>17,18</sup> The design of the edge termination was guided by device simulations using the MEDICI code of breakdown voltage with various thicknesses, overlap and type of dielectric used in the field plate. The main findings were that the use of an optimized field plate edge termination can increase the reverse breakdown voltage of vertical Ga<sub>2</sub>O<sub>3</sub> rectifiers by up to a factor of two compared

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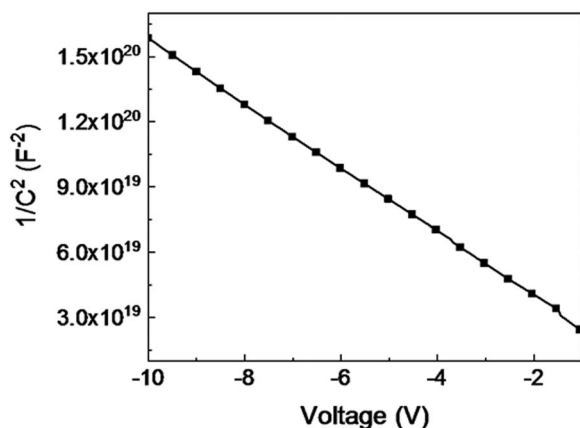
**Figure 1.** Schematic of vertical  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> Schottky rectifier utilizing a bilayer SiO<sub>2</sub>/SiN<sub>x</sub> field plate.

to unterminated devices. Moreover, the dielectric material, thickness (and ramp angle if using a bevel edge termination) all influence the resulting  $V_B$  of the rectifier by determining where the maximum field strength occurs in the device structure. The key aspect in designing the field plate edge termination is to shift the region of the high field region away from the periphery of the rectifying contact. Ar implantation,<sup>58</sup> beveling<sup>42</sup> or use of trenches<sup>48</sup> have also been shown to reduce field crowding in Ga<sub>2</sub>O<sub>3</sub> rectifiers. However, these add complexity to the fabrication relative to simple field plates and in the case of trenches, reduce the current capability.

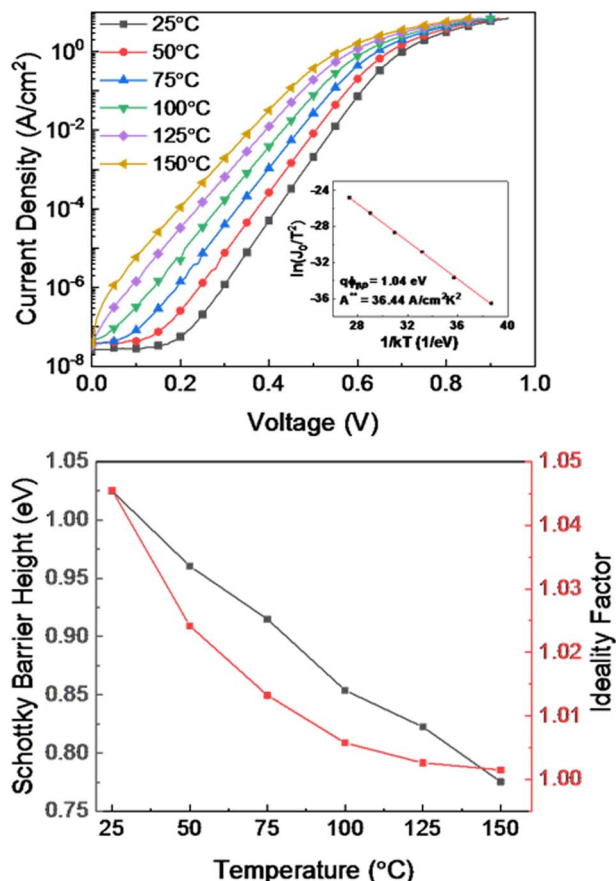
The rectifiers employed back ohmic contacts (20 nm Ti/80 nm Au) deposited by E-beam evaporation. These contacts were annealed at 550°C for 30 s in N<sub>2</sub>. The epi surface was treated with ozone for 10 min to remove adventitious carbon contamination, followed by deposition of 40 nm of SiO<sub>2</sub> and 360 nm of SiN<sub>x</sub> by plasma enhanced chemical vapor deposition at 300°C using silane and ammonia precursors. This field plate reduces the maximum electric field around the rectifying contact periphery.<sup>1-4,27,55</sup> The SiO<sub>2</sub>/SiN<sub>x</sub> contact windows were lithographically patterned and opened with 1:10 buffered oxide etch (BOE) at 25°C. The front Schottky contacts were overlapped 10 μm on the SiO<sub>2</sub>/SiN<sub>x</sub> window openings by lift-off of E-beam deposited Ni/Au (120 nm/480 nm). This geometry was guided by the simulations discussed above. The size of these contacts was fixed at 0.12 cm × 0.12 cm. In a few cases, we made smaller diodes and these were found to have higher reverse breakdown (in the range 800–1100V), but their total forward currents were also smaller and in this paper we focused on achieving high numbers for both. Figure 1 shows a schematic of the rectifier. The current-voltage (I-V) and capacitance-voltage (C-V) characteristics were measured in air at 25–150°C on an Agilent 4145B parameter analyzer and 4284A Precision LCR Meter. For reverse voltages > 100 V and forward currents > 100 mA, a Tektronix 370A curve tracer was used due to the limitation of the Agilent analyzer.

### Results and Discussion

Figure 2 shows the C<sup>-2</sup>-V characteristics used to obtain the drift layer n-type donor concentrations ( $N_D$ ). The value of  $4.36 \times$



**Figure 2.** Plot of C<sup>-2</sup>-V to determine carrier density in drift region.



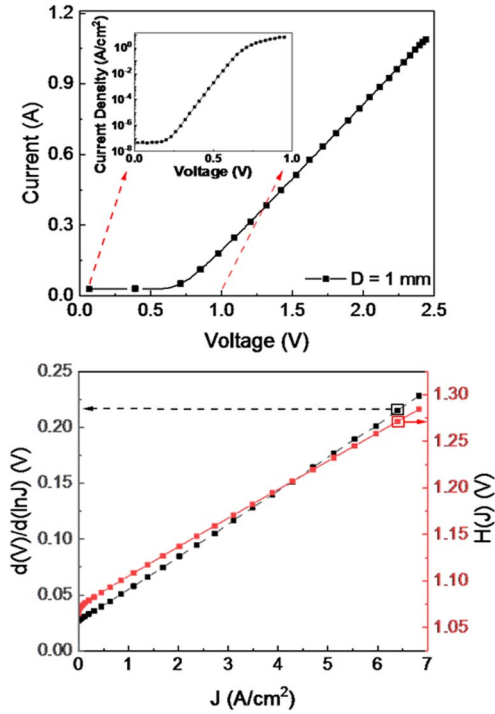
**Figure 3.** (top) Temperature dependence of forward current. The inset shows the Richardson plot derived from the forward J-V-T data. (bottom) temperature dependence of  $\Phi_b$  and ideality factor.

$10^{15} \text{ cm}^{-3}$  shows epitaxial layers of  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> can be controllably doped at low enough levels to sustain a large reverse bias while still having good forward characteristics.

Figure 3 (top) shows the temperature dependence of forward current (I-V) characteristics. The zero voltage barrier height  $e\Phi_{b0}$  (e is electronic charge and  $\Phi_b$  is barrier height) was determined from the forward current density (J)-voltage-temperature (J-V-T) characteristics by linear fitting of the Richardson's plot (inset of Figure 3). This was 1.04 eV, with a Richardson's constant of  $36.44 \text{ A} \cdot \text{cm}^{-2} \cdot \text{K}^{-2}$ , comparable to the values reported<sup>23,24</sup> for Pt of 1.15 eV and  $55 \text{ A} \cdot \text{cm}^{-2} \cdot \text{K}^{-2}$ . Figure 3 (bottom) shows the temperature dependence of  $\Phi_b$  and ideality factor  $n$  for 25–125°C. The barrier height decreases with temperature, as expected in pure thermionic emission.<sup>55–57</sup> For these lightly doped layers,  $n$  should be close to unity, with a small increase due to the image force effect.<sup>27</sup> The ideality factor improves with temperature, reaching 1.00 at 150°C, useful for elevated temperature operation.

The single sweep forward current density (J-V) characteristic is shown in Figure 4 (top). The forward current reached 1 A at 2.3 V. In this mode, the collector supply sweeps from 0 V to its preselected value. During the sweep, 1% duty cycle of a 280 μsec pulse width was employed. The on-state resistance was  $22.3 \text{ m}\Omega \cdot \text{cm}^2$ .

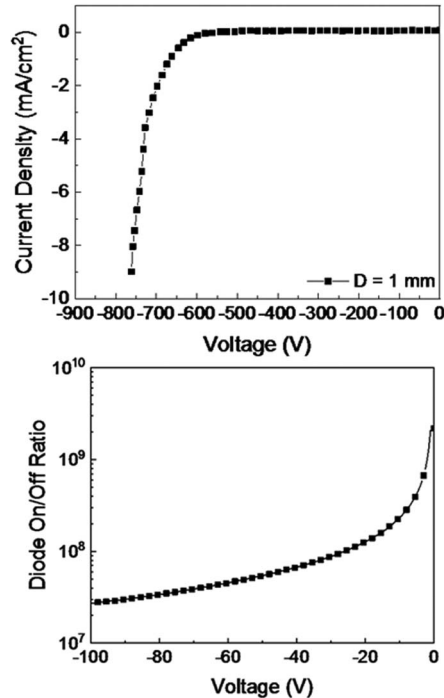
As an independent check on this data, the model of Cheung and Cheung<sup>56</sup> was applied to extract  $\Phi_b$ ,  $n$  and series resistance  $R$ . Since  $V = RAJ + n\Phi_B + (nkT/e)\ln(J/A^{**}T^2)$ , then  $d(V)/d(\ln J) = RAJ + nkT/e$ , where  $A$  is the rectifier area,  $A^{**}$  is Richardson's constant,  $k$  is Boltzmann's constant,  $T$  is measurement temperature. Thus, a plot of  $d(V)/d(\ln J)$  vs  $J$  will give  $RA$  as the slope and  $nkT/e$  as the y-axis intercept.<sup>56</sup> This data is shown in Figure 4 (bottom). The barrier height is extracted by defining the function  $H(J) = V - (nkT/e)\ln(J/A^{**}T^2)$  which is also equal to  $RAJ + n\Phi_b$ . Using the  $n$  value determined from



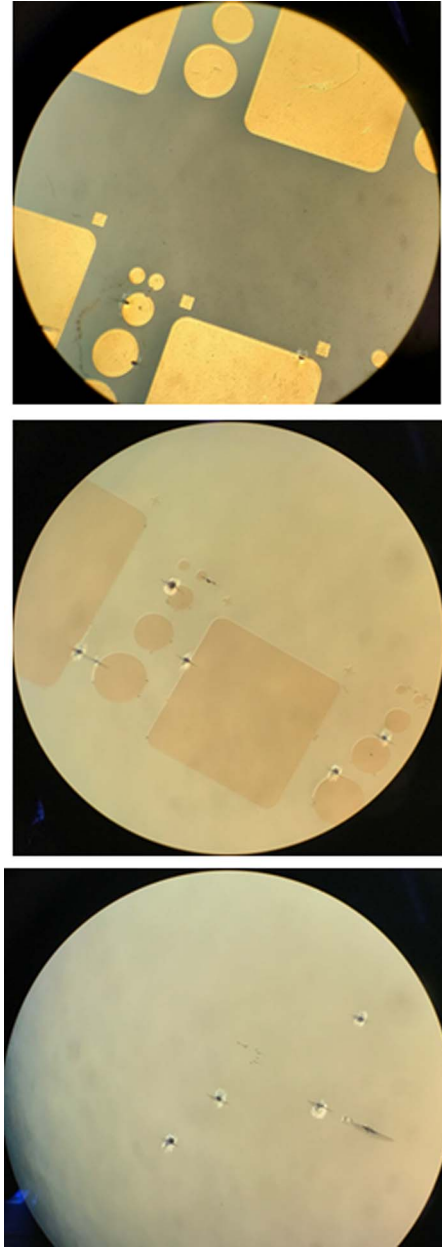
**Figure 4.** (top) Forward J-V characteristic of rectifier. The inset shows the data on a log scale. (bottom) plot of  $d(V)/d(\ln J)$  vs  $J$  to extract barrier height.

the plot of  $d(V)/d(\ln J)$  vs  $J$ , a plot of  $H(J)$  vs  $J$  will also give a straight line with y-axis intercept equal to  $n \Phi_b$ . We obtained  $R_{ON}$  values of  $29.2 \text{ m}\Omega\cdot\text{cm}^2$  from the  $d(V)/d(\ln J)$  data and  $32.4 \text{ m}\Omega\cdot\text{cm}^2$  from the  $H(J)$  analysis, in agreement with the thermionic emission analysis.

Figure 5 (top) shows the reverse breakdown voltage is 760 V, so the figure of merit ( $V_B^2/R_{ON}$ ) is  $26 \text{ MW}\cdot\text{cm}^{-2}$ . Rectifiers with much smaller areas ( $\sim 10^{-5} \text{ cm}^2$ ) exhibit higher values of

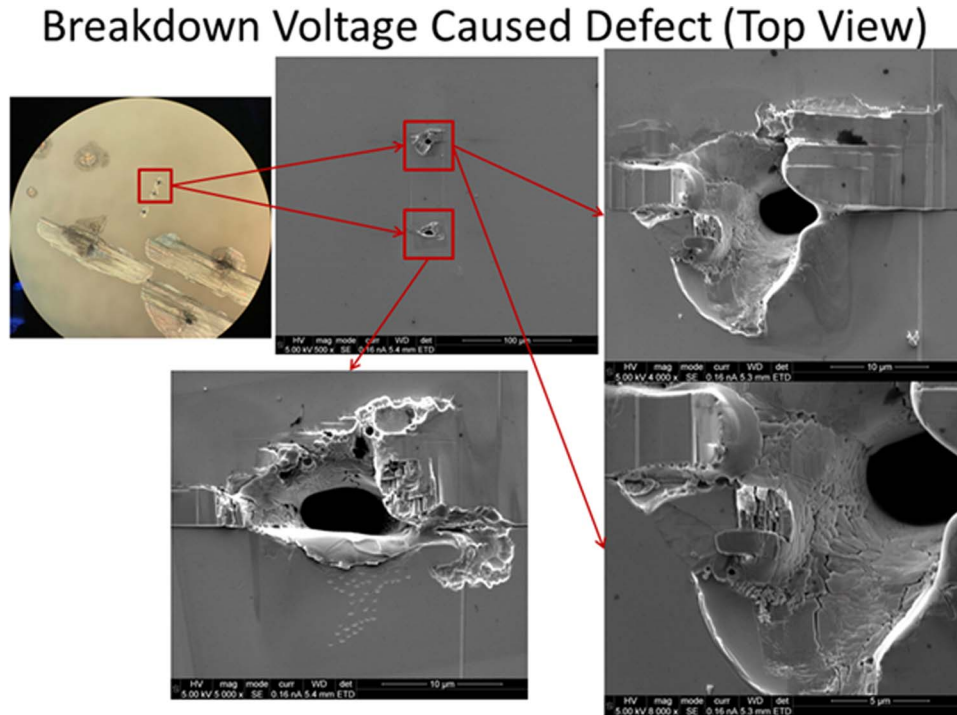


**Figure 5.** (top) Reverse current density as a function of voltage. (bottom) Rectifier on-off ratio as a function of reverse bias. The on-current was 1 A at 2.3 V and the on-off ratio range measured was  $3.3 \times 10^9$ – $5.7 \times 10^6$ .



**Figure 6.** Optical microscope images of rectifiers after reverse breakdown failure. The top image shows the sample with contacts in place, the middle image after Au/Ni etching to remove the contact and the bottom image shows the same area after BOE soaking.

$102$ – $154 \text{ MW}\cdot\text{cm}^{-2}$ ,<sup>15,16,24–30</sup> but those devices had lower total forward currents. The 760 V breakdown voltage is applicable to efficient power switching in systems for photovoltaic, wind energy and motor drives.<sup>11,12,40</sup> Figure 5 (bottom) shows the on/off current ratio was in the range  $2.7 \times 10^7$ – $2.2 \times 10^9$  when switching from +1.5 V forward bias to 1–100 V reverse bias. This is better than the previous report,<sup>53</sup> due to continued improvement in epitaxial doping control. The reverse recovery time when switching from +2 V to –10 V, was 33.5 ns, comparable to devices with much smaller rectifier dimensions.<sup>14,26,49</sup> More detailed measurements of the diode recovery time to the current level of 25% of the reverse recovery current using a clamped inductive load test circuit showed that for switching from 1 A forward current to –300 V reverse bias, the recovery time was 64 ns with  $I_r$  of 0.82 A, and the  $dI/dt$  was  $24.7 \text{ A}/\mu\text{s}$ .<sup>58</sup> These values are comparable to previously reported  $\text{Ga}_2\text{O}_3$  trench MOS Schottky diodes, Si-fast recovery



**Figure 7.** Optical microscope image and close-up SEM images of the pits formed as a result of failure under reverse bias breakdown.

diodes (FRDs, Rohm RF1005TF6S) and SiC Schottky barrier diodes from Cree (part number C3D0260A) for turn-off from  $I_F = 1$  A.<sup>49</sup>

It is also worth noting the rectifiers when pushed to failure under reverse bias conditions, showed formation of pits along the periphery of the Schottky contact. An example is shown in Figure 7, which shows optical microscope images of the pits formed at the contact edge. The pits are clearer after etch removal of the Au/Ni contact. Figure 8 shows scanning electron microscopy images of these pits, which result from avalanche failure of the  $\text{Ga}_2\text{O}_3$  under the high field generated at the edge of the rectifying contact. These results indicate that further optimization of the edge termination material and geometry is needed to reduce field crowding.

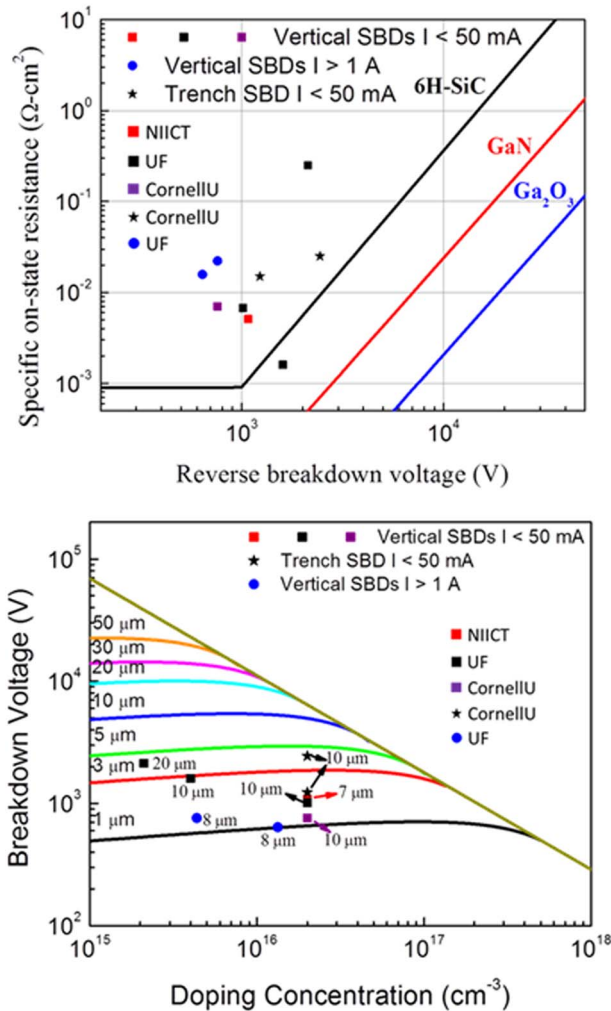
To put our results in context with the previous literature on  $\text{Ga}_2\text{O}_3$  and also the performance of SiC and GaN rectifiers, our  $\beta\text{-Ga}_2\text{O}_3$  rectifiers are benchmarked in the plot of specific  $R_{\text{ON}}$  versus  $V_b$  in Figure 8 (top). This also shows the relation between breakdown voltage, electric

field and doping in vertical geometry rectifier consisting of a lightly doped drift region on a more heavily doped layer on a conducting substrate of these respective materials. Experimental points for  $\text{Ga}_2\text{O}_3$  from different groups<sup>19,21,22,25–28,34,35,46–48,57</sup> are also shown-these are not yet at the values achieved by the smaller bandgap SiC and GaN, where the theoretical limits are now being approached. Continued development of low defect substrates, optimized epi growth and surface treatments and improved device design and processing methods for  $\text{Ga}_2\text{O}_3$  are still required to push the experimental results closer to their theoretical values. Table I also shows a detailed compilation of vertical rectifier results from the literature.

Figure 8 (bottom) is a plot of theoretical breakdown voltage of  $\text{Ga}_2\text{O}_3$  vertical punchthrough diodes as a function of doping concentration and drift region thickness. This assumes the breakdown voltage is given by  $V_B = E_C W / [e N_B W^2 / 2 \epsilon \epsilon_r]$ , where  $E_C$  is the critical field for breakdown,  $W$  is the depletion depth,  $N_B$  the doping concentration

**Table I. Summary of vertical geometry  $\text{Ga}_2\text{O}_3$  rectifiers reported in literature.**

Reference	Epi Thickness( $\mu\text{m}$ )	Drift Layer Doping ( $\text{cm}^{-3}$ )	Edge Termination	$V_B$ (V)	$R_{\text{ON}}$ ( $\Omega\cdot\text{cm}^2$ )
Konishi et al. <sup>27</sup>	10	$1.8 \times 10^{16}$	Yes-field plate	1076	$5.1 \times 10^{-3}$
Yang et al. <sup>25</sup>	10	$4.02 \times 10^{15}$	No	1600	$25 \times 10^{-3}$
Yang et al. <sup>26</sup>	10	$2 \times 10^{16}$	No	1016	$6.7 \times 10^{-3}$
Sasaki et al. <sup>21</sup>	Unintentionally doped substrate	$3 \times 10^{16}$	No	150	$4.3 \times 10^{-3}$
Li et al. <sup>28</sup>	10	$2 \times 10^{16}$	Trench	1350	$15 \times 10^{-3}$
Li et al. <sup>48</sup>	10	$2 \times 10^{16}$	Trench	2440	$25 \times 10^{-3}$
Oh et al. <sup>38</sup>	2	Undoped, $< 3 \times 10^{16}$	No	210	2582
He et al. <sup>22</sup>	Unintentionally doped substrate	$2 \times 10^{14}$	No	>40	$12.5 \times 10^{-3}$
Tadger et al. <sup>34</sup>	$\sim 10$	$8 \times 10^{12}$	No	2380	n/a
Fu et al. <sup>35</sup>	Sn-doped EFG substrates	$4 \times 10^{18}$	No	low	$0.77 \times 10^{-3}$
Joishi et al. <sup>42</sup>	2	$2.5 \times 10^{17}$	Bevel	190	$3.9 \times 10^{-3}$
Yang et al. <sup>46</sup>	10	$1.33 \times 10^{16}$	Yes-field plate	650	$1.58 \times 10^{-2}$
Yang et al. <sup>47</sup>	20	$2.1 \times 10^{15}$	Yes-field plate	2300	0.25
Yang et al. <sup>48</sup>	7	$2 \times 10^{16}$	No	466	$0.26\text{--}5.9 \times 10^{-4}$
Li <sup>57</sup>	15, exfoliated	Sn doped	No	97	$2.1 \times 10^{-3}$
Gao et al. <sup>19</sup>	10, exfoliated	Sn doped	Ar implantation	550	$1.7 \times 10^{-3}$
This work	8	$4.4 \times 10^{15}$	Yes-field plate	760	$22.2 \times 10^{-3}$



**Figure 8.** (top) The reverse breakdown voltage of punch-through junctions for  $\text{Ga}_2\text{O}_3$  as a function of doping concentration and drift region thickness. (bottom) Specific  $R_{\text{ON}}$  versus  $V_B$  of state-of-the-art vertical  $\beta\text{-Ga}_2\text{O}_3$  rectifiers.

and  $\epsilon_r$  and  $\epsilon$  the relative and absolute permittivity. Even 3  $\mu\text{m}$  epi layers with doping concentration of  $10^{16} \text{ cm}^{-3}$  should have a theoretical breakdown voltage of  $\sim 1800\text{V}$ . The actual experimental value of  $V_B$  is well below the theoretical predictions. Our case of 8  $\mu\text{m}$  with doping  $4.4 \times 10^{15} \text{ cm}^{-3}$  has a theoretical breakdown more than an order of magnitude larger than the experimental value.<sup>43,44</sup>

### Conclusions

The initial thrust on  $\text{Ga}_2\text{O}_3$  electronics is targeted toward high power converters for both DC/DC and DC/AC applications. Schottky barrier diodes on  $\beta\text{-Ga}_2\text{O}_3$  have achieved a breakdown strength of  $\sim 4 \text{ MV/cm}$ . The question remains as to whether  $\text{Ga}_2\text{O}_3$  will have commercial advantages over the more mature SiC and GaN technology for power switching and power amplifier applications. While the initial device performance looks promising, challenges remain, including growth maturity, thermal limits, cost, and device reliability. The results summarized here show the potential of  $\beta\text{-Ga}_2\text{O}_3$  for fast-switching power devices, capable of simultaneously achieving both high on-state currents and breakdown voltages. The use of thick, lightly doped epitaxial drift regions enable realization of large dimension ( $0.014 \text{ cm}^2$ )  $\beta\text{-Ga}_2\text{O}_3$  Schottky rectifiers with large forward current (1A),  $V_{\text{BR}}$  values of 760 V and power density figures of merit of  $26 \text{ MW} \cdot \text{cm}^{-2}$ . Since power converters require the power device to switch at high frequen-

cies for improved dynamic response capability and reduced passive component size and weight, the performance of these  $\text{Ga}_2\text{O}_3$  rectifiers is consistent with these goals. It is likely that  $\text{Ga}_2\text{O}_3$  will not displace materials such as SiC and GaN, but possibly supplement them at high voltages in hybrid systems.

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