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THE NON-STATIONARY FLOW OF AN ANOMALOUS VISCOUS LIQUID

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FOREIGN TECHNOLOGY DIVISION



THE NON-STATIONARY FLOW OF AN ANOMALOUS VISCOUS LIQUID

bу

O. K. Fedorov



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properti: at displacement speed gamma-less than pamma, than - . . rammag with a viscosity of Mul but at gamma greater than rammag it has a viscosity of Mul. It is assumed that the liquid was initially at rest. At the initial moment of time the surface, on top of which the limuid is found, begins to move with constant speed. Thus a boundary arises separating liquids with viscosities. Mu_l, and Mu_p. On this boundary forward speed and drift speed are to be continuous. The solution is expressed through a Laplace function. In the statement of the problem considered by the author surface forces should act on the separation

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^{*} ye initially, after vowels, and after \$, \$; e elsewhere.
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The use of diacritical marks is preferred, but such marks may be omitted when expediency dictates.

THE NON-STATIONARY FLOW OF AN ANOMALOUS VISCOUS LIQUID

O. K. Fedorov

During the flow of a dispersed system the effect of a sharp drop in viscosity is observed at a certain rate of deformation of displacement [1]. Such an anomaly in viscosity can lead to a unique situation when two areas of liquid flow are formed: an area of low velocities of displacement, characterized by a high value of viscosity, and a low-viscosity area, corresponding to high velocities of displacement. Below an anlaysis is made of the problem of the plane-parallel flow of such a liquid caused by the movement of flat wall.

During the flow of an anomalous viscous liquid the stress of the displacement τ is connected with the velocity of displacement $\dot{\gamma}$ by the dependence

$$\tau = \mu(\dot{\tau})\dot{\tau}. \tag{1}$$

where $\mu(\dot{\gamma})$ - viscosity, depending on the velocity of displacement.

At $\dot{\gamma}=\dot{\gamma}_0$ there is a drop in viscosity, since the flow of a liquid at $\dot{\gamma}\leq\dot{\gamma}_0$ can be characterized by the viscosity μ_1 (the greatest Newtonian viscosity), and at $\dot{\gamma}>\dot{\gamma}_0$ — by the viscosity μ_2 (the least Newtonian viscosity).

The liquid occupies the half space y > 0 and is in contact with the wall y = 0. The wall moves uniformly at a constant rate u directed along the axis X. The velocity v of plane-parallel flow of a liquid will be a function of time t and coordinate y. The field of velocities of displacement deformation of the liquid corresponds to two values of viscosity μ_1 and μ_2 , characterizing the two areas of flow. The surface of separation between the areas is determined by coordinate Y(t), whereas on this surface the conditions of continuity of the field of velocities $v_{1,2}$ and the field of velocities of displacement $\dot{\gamma}_{1,2}$ should be satisfied (subscripts 1, 2 are taken for areas with the corresponding values of viscosity μ_1 and μ_2). On the moving wall $\dot{\gamma}$ = 0, therefore the area with viscosity μ_1 adjoins it.

It follows from the Navier-Stokes equations that without a calculation for the pressure gradient

$$\frac{\partial v_i}{\partial I} = a_i \frac{\partial^2 v_i}{\partial y^2} (i = 1, 2), \tag{2}$$

where $v_1 = \frac{\mu_1}{\rho}$, ρ - density of the liquid.

Let us present the boundary conditions in the form

$$v_1 = u$$
 when $y = 0$, $v_2 = 0$ when $y \to \infty$, (3)

$$v_1 = v_2$$
, $\frac{\partial v_1}{\partial y} = \frac{\partial v_2}{\partial y}$, when $y = Y(i)$. (4)

Serving for the determination of the surface of separation is the condition that the velocity of displacement on it is equal to $\dot{\gamma}_0$:

$$\frac{\partial v_1}{\partial y} = -i_0 \text{ at } y = Y(t). \tag{5}$$

In such a formulation the examined problem is close to the known Attefan problem concerning the freezing of Lee [2].

The solution of equation (2), satisfying the boundary conditions (3), we present in the form

$$v_1 = u - A\Phi\left(\frac{y}{2\sqrt{a_1 l}}\right)$$

and

$$v_2 = B \left[1 - \Phi \left(\frac{y}{2 \sqrt{a_2 t}} \right) \right].$$

Here

$$\Phi(x) = \frac{2}{V_{\pi}} \int_{0}^{x} e^{-it} d\xi$$

is the Laplace function; A and B - arbitrary constants. Using conditions (4) and (5) we obtain

$$v_i = \mu - \frac{Y_{10}^*}{\lambda_0 z} \frac{\Phi(z\lambda)}{\Phi'(z\lambda_0)} \tag{6}$$

and

$$v_2 = \frac{Y_{70}^{\prime}}{\lambda_0} \frac{1 - \Phi(\lambda)}{\Phi'(\gamma_0)}, \qquad (7)$$

where

$$\lambda = \frac{y}{2\sqrt[3]{a_2t}}, \quad \lambda_0 = \frac{y}{2\sqrt[3]{a_2t}}, \quad z = \sqrt[3]{\frac{a_2}{a_1}},$$

$$\Phi'(x) = \frac{\lambda \Phi(x)}{dx},$$

whereas Y(t) is determined from the equation

$$\frac{1 - \Psi(\lambda_0)}{\Psi'(\lambda_0)} - \left| -\frac{1}{2} \frac{\Phi(z\lambda_0)}{\Psi'(z\lambda_0)} \right| = \frac{u\lambda_0}{Y_{10}^*}.$$
 (8)

Expressions (6) and (7) are the solutions of the problem which satisfy the following conditions: f = 0, Y = 0, $v_2 = 0$, i.e., at the initial moment all the liquid is stationary.

We will make use of the fact that usually $\mu_1 >> \mu_2$; then $\alpha << 1$. Using for the Lapalce function the concept in the form of a series and retaining in it only the first member, we obtain the following expression for the velocity of flow of a liquid in a highly viscous area

$$v_i = u - \dot{\gamma}_{oy}. \tag{9}$$

We arrive at such an expression if we examine the flow of a liquid at $f \to \infty$. From here it follows that in a highly viscous wall zone the flow is established very rapidly and it depends on the rate of movement of the wall u and the velocity of displacement $\dot{\gamma}_0$.

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