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THE ORBIT OF ARIEL 3

(1967-42A)

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by

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### SUMMARY

The definitive orbit for Ariel 3 has been computed, from Minitrack observations, for a period of  $27\frac{1}{2}$  months from the launch of the satellite. The orbit was represented by a model with seven independent orbital parameters and the values of these parameters were determined, and are listed, at three-day intervals. Typical accuracies are  $10^{-5}$  in eccentricity and  $4^{6}$  in angular parameters, that is, about  $\frac{1}{2}$  km in position.

A curious feature of the secular variation of orbital inclination, viz. that the expected decrease of about 0.02° appeared to occur over a three-month period instead of the full 27½-month period, is discussed but has not been explained.

Departmental Reference: Space 327

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#### 1 INTRODUCTION

Ariel 3, the third of the series of satellites being launched in the scientific programme of Anglo-American co-operation, was the first spacecraft to be built entirely in Britain. The five experiments in the payload were concerned with electron density and temperature (University of Birmingham), VLF radio waves (University of Sheffield), cosmic radio noise (Nuffield Radio Astrohomy Laboratory, Jodrell Bank), molecular oxygen (Meteorological Office) and terrestrial noise sources (Radio and Space Research Station).

Known as S53 or UK3 before launch, and as Ariel 3 or 1967-42A afterwards, the satellite was placed in a near-polar, near-circular orbit at 18<sup>h</sup> UT on 5 May 1967 by a Scout rocket launched from the Western Test Range, California.

The definitive orbit of Ariel 3, as for Ariel 2, has been derived at R.A.E. from Minitrack (interferometer) data provided by NASA. Orbital parameters have been obtained at three-day intervals by the use of the new computer program PROP<sup>3</sup>. They are tabulated, for the first 27½ months of the satellite's lifetime, in this Report. For epochs up to 1968 JAN 16 the work was done on London University's Atlas computer and for subsequent epochs it was done on the ICL 1907 computer at R.A.E.

#### 2 OBSERVATIONS

The STADAN Minitrack network now consists of ten stations - five of the original twelve stations have disappeared and there are three new ones. These are listed in Table 1, with their assumed positions in standard geocentric co-ordinates (x-axis towards the Greenwich meridian).

Observations consist of pairs of direction cosines. Their a priori accuracy (s.d.) has been taken, as usual<sup>3,4</sup>, at 0.00029, equivalent to 1' in angular measure, though their true accuracy is believed to be worse than this.

Times are given in the UTC system, i.e. the system defined by WWV time transmissions from America, and have not been corrected during the orbital determination, except that the times of observations made on 1 February 1968 and used at epoch 1968 JAN 31 had to be reduced by 0.1<sup>8</sup> to allow for the step advance of UTC at Feb 1.0. Times should be accurate to about 1 ms; no allowance for timing error was made by PROP.

In total, about 10000 Minitrack observations were used, covering the period from launch until, about a week before the satellite's transmitter was switched off (1 September 1969); i.e. these were about 12 per day. They were received from NASA on punched cards, suitable for direct input to PROP. The epochs for orbit determination were taken at three-day intervals, and always at midnights (unlike the Ariel 2 epochs which were at ascending nodes). The orbit determination at each epoch used observations over a four-day period, allowing one-day overlaps in the periods of validity of the resulting orbital parameters, but observations were not (in general) used twice; observations on an 'overlap day' were divided into two sets, by alternate allocation, for use with the epochs before and after the overlap day. About 650 of the observations were rejected during analysis, but this includes (a) a rather high rejection rate during the first 3½ months - all the observations from Orroral were being rejected at one stage - and (b) nearly 100 observations, over a period of a month (Jan-Feb, 1968), which all had a one wavelength error in the north-south direction cosine due to a temporary error in the NASA program for processing the raw data; the normal rejection rate was about 4½%.

The number of Minitrack observations per day varied, of course, but there was at least one on every day of the period covered, apart from the week

28 November-4 December 1967, for which there was a complete absence of data.

A few observations of Ariel 3 were made by the Hewitt camera at Malvern. Among these, 8 observations came from a pass on the evening of 10 April 1969 and 12 observations came from two passes close to midnight on 19 June 1969 and 21 June 1969. It was decided to incorporate these 20 observations into re-runs of the orbit determinations at the appropriate two epochs, to see how fit and accuracy were affected. The remaining Hewitt camera observations, and the many visual observations of Ariel 3, have not been used.

### 3 ORBITAL MODEL

The orbital model of the program PROP is not the same as that used in the analysis of the orbit of Ariel 2. Eccentricity, inclinations etc. are defined slightly differently in the two programs, and the connecting relations are given in Appendix C of Ref. 3. The model allows some choice as to the set of orbital parameters which represent the orbit and which are determined from fitting to observations. The set chosen for the Ariel 3 orbit contained seven parameters, viz, e (eccentricity), i (inclination),  $\Omega_{\rm O}$  (right ascension of the node),  $\omega_{\rm O}$  (argument of perigee), M (mean anomaly), M (mean motion) and M (half the mean acceleration). The first four parameters are epoch values of mean elements, as defined in Refs.7 and 8 and the last three are the coefficients in the polynomial representation of (mean) mean anomaly:

$$M = M_0 + M_1 t + M_2 t^2$$
,

where t is measured from epoch.

Secular rates of change of e, i,  $\Omega$  and  $\omega$  (i.e. the polynomial coefficients  $e_1$ ,  $i_1$ ,  $\Omega_1$  and  $\omega_1$ ) were computed, inside PROP at the beginning of each iteration of the differential-correction process, as functions of the seven independent parameters  $^{3,8}$ . These quantities, together with the long-periodic and short-periodic terms computed at each observation time, represented orbit perturbations due to drag and to the earth's zonal harmonics up to  $J_9$ . The along-track effect of the tesseral harmonic  $J_{2,2}$  was represented as usual  $^{3,8}$  using the value  $1.8 \times 10^{-6}$ , but, apart from this, tesseral harmonics were neglected. Luni-solar perturbations were ignored; their effect on Ariel 3 over a period of two or three days from any epoch is very small.

The decision not to have an eighth parameter M<sub>3</sub>, which would have made the M polynomial a cubic, was justified by some test runs, early in the lifetime, which showed that no significant improvement in fit would result and that the value of M<sub>3</sub> itself would not be significant. For the two epochs during the week of missing data, however, M<sub>3</sub> was included in the model and a reasonable fit thereby obtained to data before and after the gap, covering a period of 10½ days. In retrospect, the decision to omit M<sub>3</sub> is open to question because drag increased fairly steadily through the 27½ months considered (as the values of M<sub>2</sub> in Table 2 show), the maximum effects being at the end of March 1969; repetition of two of the runs, with M<sub>3</sub> included, showed that at this stage significant improvement in fit would be obtained (see section 4).

Apart from the omission of  $M_3$  during periods of high drag, the main limitation of the orbital model is in the neglect of important tesseral-harmonic perturbations and, in particular, of the perturbation in inclination due to  $J_{2,2}$ . This perturbation has a period of just under 12 hours and an amplitude of about  $0^{\circ}.002$ , equivalent to a maximum position error of about  $\frac{1}{2}$  km.

### 4 RESULTS

### 4.1 Main results

The orbital parameters obtained from the orbit determinations, together with certain additional information, are listed in Table 2. Successive columns of the table provide the following quantities, zero suffixes being omitted from a etc:-

Epoch date (0<sup>h</sup> UTC understood).

Semi-major axis, a (km).

Eccentricity, e.

Perigee height, h (km).

Inclination, i (degrees).

Right ascension of the node,  $\Omega$  (degrees).

Argument of perigee,  $\omega$  (degrees).

Mean argument of latitutde,  $M_0 + \omega$  (degrees).

Mean motion, M, (degrees/day).

Half acceleration, M<sub>2</sub> (degrees/day<sup>2</sup>).

Number of observations used, N.

Number of observations rejected, K.

Extent of the observations, D (days).

Standard deviation of an observation of unit weight, &.

Modified Julian Day number of epoch date, MJD.

The orbital parameters are the seven quantities e, i,  $\Omega$ ,  $\omega$ ,  $M_0$ ,  $M_1$  and  $M_2$ , but  $M_0 + \omega$  is given instead of  $M_0$  because of the high correlation between  $M_0$  and  $\omega$ . This correlation arises directly from the fact that the orbit is so nearly circular, and Ref.3 may be consulted for further explanation. (The appropriate value of the control parameter JELTYP was used to give the variance of  $M_0 + \omega$  directly.)

The semi-major axis, a, is the mean element, as used by Merson, defined from M, by

$$a = (\mu/M_1)^{\frac{1}{3}} - \frac{1}{4} J_2 R^2 (\mu/M_1^2)^{-\frac{1}{3}} (2 - 3 \sin^2 i) (1 - e^2)^{-\frac{1}{2}}$$
,

where  $\mu$  is the earth's gravitational constant,  $J_2$  is its second zonal harmonic coefficient and R is its mean equatorial radius, the values from Ref.3 being used.

The perigee height is given by

$$h_p = r_p - R_p ,$$

where\*

$$r_{p} = a (1 - e) + \frac{J_{2} R^{2}}{4a(1 - e^{2})} \left\{ \sin^{2} i \cos 2\omega - (2 - 3 \sin^{2} i) \left( 1 + \frac{e}{1 + \sqrt{1 - e^{2}}} - \frac{\sqrt{1 - e^{2}}}{1 + e} \right) \right\}$$

and

$$R_{p} = R - 21.379 \sin^{2} i \sin^{2} \omega$$

The right ascension of the node is nominally referred to the standard PROP equinox<sup>3</sup>, but contains a small error due to the fact that the times are given in UTC and no correction to UTl has been made. To correct  $\Omega$  to the time PROP equinox (epoch date still understood to be  $0^h$  UTC) add  $0^\circ.004 \times (UTl - UTC)$ , where the time difference is in seconds.

The 'number of observations used' includes the number rejected; i.e. the parameters have been determined, in the end, from N - K observations.

After nine of the tabulated quanitites - the seven orbital parameters plus semi-major axis and perigee height - are given their computed standard

<sup>\*</sup>The difference between  $r_p$  and a (1-e) is important. Thus, though  $r_p$  (or rather  $h_p$ ) is the right parameter to use when correlating drag behaviour with perigee height, a (1-e) is the right parameter to work with when studying the effects of the earth's odd harmonics. For Ariel 3 the difference is approximately 1.54 cos  $2\omega$  km, and it was the use of  $r_p$  instead of a (1-e) which led to the apparent discrepancy mentioned in section 3.20 of Ref.13.

deviations, to one or two significant figures, the unit in each case being that of the final figure quoted for the main quantity. Every standard deviation includes  $\varepsilon$  as a factor, where  $\varepsilon$  is given by

$$\varepsilon = \{\Sigma (Res/0.00029)^2/(2N - 2K - 7)\}^{\frac{1}{2}};$$

here the summation is over all residuals, Res, in the N-K accepted east-west and north-south direction cosines, and 0.00029 is the a priori accuracy referred to in section 2. Since the actual accuracy is worse than this a priori figure  $^4$ , the values of  $\varepsilon$  in Table 2 are expected to be and are - larger than unity.

Table 2 was obtained as direct computer output from a program known as TOP (Tabulation of Orbital Parameters). This program takes, as part of its input data, the punched card output from PROP runs, so there should not be any errors in the table.

The secular rates of change  $e_1$ ,  $i_1$ ,  $\Omega_1$  and  $\omega_1$  are not given in Table 2, since they are computed internally by PROP as part of the model. It is remarked, however, that the computation of the  $J_2^2$  component of  $\omega_1$ contained an error until the PROP3 version of the program was introduced at the end of January, 1969. PROP2, which had this error in  $\omega_1$ , was able to compensate for it almost exactly, by fitting a slightly wrong value of M1, and this was one reason why the error was for a long time undiscovered. To correct the results from the PROP2 runs it was only necessary to correct M, by an amount equal to the error in  $\omega_1$ , and a special program was written to do this. The only reason for mentioning this point is that the values of M, in Table 2 are the corrected values, and so are different from the values provided in the first four provisional lists of Ariel 3 parameters to be issued. (To avoid having to ask AWRE Foulness to make a correction to the Ariel 3 telemetry data analysis program after PROP3 had been introduced, it was decided to continue sending incorrect values to AWRE, by adding the appropriate deliberate error to  $M_1$ .)

Fig.1 gives a plot of orbital inclination, each value being represented by a vertical line, two standard deviations in length, centred on the fitted value. Fig.2 gives a plot of eccentricity, but most of the time the scale is too small for standard deviations to be shown. Fig.3 shows a short section of the eccentricity curve (covering just over half a period of the perigee) with the scale expanded sufficiently for the standard deviations to be indicated as on the inclination plot. Figs.1 and 2 give, in addition to the definitive inclinations and eccentricities obtained at R.A.E., the SDC (NORAD) values published in Spacetrack bullatins.

# 4.2 Results involving the parameter M<sub>3</sub>

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For two of the runs covered by Table 2 the parameter  $M_3$  was included in the orbital model, namely, for those of epochs 1967 NOV 29 and 1967 DEC 2. These were the epochs which occurred during the week when no Minitrack data were supplied. Without  $M_3$  the fit was twice as bad ( $\epsilon$  5.7 for the first epoch instead of 2.8), due to the number of days spanned by the observations. The value of  $M_3$ , omitted from Table 2 to retain a regular format, was -0.00069, with standard deviation 0.00003, for both epochs; the same observations were used in both runs, so the second set of elements is really just the first set advanced three days, with small variations in the residuals (and hence a small change in  $\epsilon$ ) due to limitations of the orbital model.

It was stated in section 3 that, for epochs early in the satellite's lifetime, general introduction of M, would not have helped. To illustrate, the complete set of parameters, when  $M_3$  is included, for 1967 DEC 20 is as follows (with last-figure standard deviations in brackets):e = 0.007329 (15), i = 80.1802 (19),  $\Omega = 239.0276$  (19),  $\omega = 155.22$  (10),  $M_0 + \omega = 91.7171$  (15),  $M_1 = 5433.0021$  (22),  $M_2 = 0.0701$  (9) and  $M_2 = 0.0007$  (9); the value of  $\varepsilon$ , viz. 3.6, was actually <u>larger</u> (unrounded value, 3.553 as against 3.546) than for the run without  $M_2$ , due to the loss of a degree of freedom. For certain epochs later in the lifetime, however, introduction of Ma would have led to better fits. This may be illustrated by considering the two worst fits obtained, namely, for epochs 1969 MAR 20 ( $\epsilon$  of 5.0 in Table 2) and 1969 MAR 23 ( $\epsilon$  of 5.1); on repeating these runs, with M, included, the following results were obtained:- for 1969 MAR 20, e = 0.006902 (45), i = 80.1665 (20),  $\Omega = 10.3404$  (19),  $\omega = 124.14$  (17),  $M_0 + \omega = 189.4465$  (46),  $M_1 = 5475.7830$  (24),  $M_2 = 0.1282$  (9) and  $M_2 = 0.0077$  (9), with  $\varepsilon = 3.1$ ; for 1969 MAR 23,  $\varepsilon = 0.006936$  (49),  $i = 80.1677 (30), \Omega = 6.4259 (22), \omega = 116.38 (16), M_0 + \omega = 48.2682 (47),$  $M_1 = 5476.6144$  (32),  $M_2 = 0.1444$  (12) and  $M_3 = 0.0085$  (13), with  $\varepsilon = 3.6$ .

## 4.3 Results involving Hewitt camera observations

The runs at epochs 1969 APR 10 and 1969 JUN 21 were repeated, with (respectively 8 and 12) Hewitt camera observations included. The following results were obtained:- for 1969 APR 10, e = 0.006778 (7), i = 80.1671 (12),  $\Omega = 342.9333$  (13),  $\omega = 67.28$  (12),  $M_{\Omega} + \omega = 326.4188$  (10),  $M_{\eta} = 5480.4510$  (8) and M, = 0.0850 (7), with  $\varepsilon$  = 2.0 and only the same Minitrack observation rejected as was originally rejected; for 1969 JUN 21, e = 0.005403 (8),  $i = 80.1610 (10), \Omega = 248.6570 (13), \omega = 175.80 (6), M_0 + \omega = 120.8871 (10),$  $M_1$  = 5490.2793 (7) and  $M_2$  = 0.0584 (3), with  $\varepsilon$  = 2.4 and two of the Minitrack observations rejected that had previously been accepted. On comparison with corresponding entries in Table 2 it may be seen that for the first run there is little change - the maximum change in a parameter is for  $M_{\alpha}$  +  $\omega$ , the change being about twice the original standard deviation, and no standard deviation has decreased by a factor of more than 1; for the other run, however, there is a large change in eccentricity, nearly five times the original standard deviation, and the standard deviations for e,  $\omega$  and M +  $\omega$  have all been reduced by factors of more than 2. (It is worth remarking that the change in e was caused entirely by the introduction of the Hewitt camera observations, and not at all by the subsequent rejection of two Minitrack observations.)

A reasonable conclusion is that Hewitt camera observations, of high accuracy, are compatible with Minitrack observations of poorer accuracy. For a high-inclination satellite like Ariel 3 the effect is not very significant if the Hewitt camera observations all come from a single pass, but there is a great improvement in accuracy when observations from two or more passes are available.

## 5 ACCURACY OF POSITION COMPUTATION

As with Ariel 2 it was required, for correlation with on-board experiments, that the definitive orbital parameters should be good enough for position to be computable from them to better than 1 km. In the paper on the Ariel 2 orbit the accuracy of position computation was considered by reference to plots of  $\{\sigma^2(\mathbf{x}) + \sigma^2(\mathbf{y}) + \sigma^2(\mathbf{z})\}^{\frac{1}{2}}$ , where the variances  $\sigma^2(\mathbf{x})$ ,  $\sigma^2(\mathbf{y})$  and  $\sigma^2(\mathbf{z})$  are functions of time and the covariance matrix of the orbital parameters, and by comparison of such plots with plots, during overlap periods, of  $\{(\mathbf{x}_1 - \mathbf{x}_2)^2 + (\mathbf{y}_1 - \mathbf{y}_2)^2 + (\mathbf{z}_1 - \mathbf{z}_2)^2\}^{\frac{1}{2}}$ , where  $\mathbf{x}_1$ ,  $\mathbf{y}_1$ ,  $\mathbf{z}_1$  denote satellite

co-ordinates computed from orbital parameters at the epoch before the given overlap period and  $x_2$ ,  $y_2$ ,  $z_2$  denote co-ordinates computed from parameters at the following epoch. This approach would have been equally possible for Ariel 3, using the program PREP<sup>8</sup>, but it was decided that it would be adequate to consider the question by looking directly at standard deviations of orbital parameters and interpreting these as maximum position errors after  $1\frac{1}{2}$  days.

The justification for this approach is that, with  $M_0 + \omega$  rather than  $M_0$  taken as a parameter, large correlations between parameters did not occur. (Occasional correlations as large as  $\pm 0.4$  occurred, usually involving e and one other parameter.) For a satellite, like Ariel 3, in an orbit which is nearly circular and not too far from polar, the maximum effects, on position after  $1\frac{1}{2}$  days, of one-sigma errors in the parameters are approximately as follows:  $2 \text{ a } \sigma(e)$ , a  $\sigma(i)$ , a  $\sigma(\Omega)$ ,  $2 \text{ a } e \sigma(\omega)$ , a  $\sigma(M_0 + \omega)$ ,  $1\frac{1}{2}$  a  $\sigma(M_1)$  and  $2 \text{ a } \sigma(M_2)$ , where angle sigmas are now taken to be in radians. The main effects here of e,  $\omega$  and  $M_0 + \omega$  are the along-track errors which arise from the expression of argument of latitude in the form

$$u = (M + \omega) + 2 e \sin \{(M + \omega) - \omega\}$$
.

Let us consider 'maximum position effects' for two different sets of sigmas: first, the maximum value of each sigma that occurs anywhere in Table 2, and, second, maximum values during, roughly speaking, the best ninety per cent of the time. Denoting sigmas from the two sets by  $\sigma_1$  and  $\sigma_2$  respectively (with angles in degrees again), and the corresponding maximum position effects, in km, by MPE<sub>1</sub> and MPE<sub>2</sub>, we have the following table:-

	$\frac{\sigma_1}{2}$	MPE <sub>1</sub>	$\frac{\sigma_2}{}$	MPE <sub>2</sub>
e	0.000072	1.1	0.000020	0.3
i	0.0043	0.5	0.0030	0.4
Ω	0.0048	0.6	0.0025	0.3
ω	0.51	0.8	0.17	0.3
M <sub>0</sub> + ω	0.0074	0.9	0.0031	0.4
M <sub>1</sub>	0.0043	0.8	0.0023	0.4
M <sub>2</sub>	C.0019	0.5	0.0019	0.5

Since  $(\Sigma \text{ MPE}_2^2)^{\frac{1}{2}} = 1.0 \text{ km}$  it is reasonable to claim that the accuracy requirements are met most of the time. If we consider the accuracy of height only, i.e. of  $r = a \{1 - e \cos (M + \omega - \omega)\}$ , then only the maximum onesigma effects  $a \sigma(e)$  and  $a e \sigma(\omega)$  are significant; this gives  $a \Sigma^{\frac{1}{2}}$  of 0.4 km corresponding to the MPE<sub>2</sub> column in the table.

Some comments may be useful on the reason for some of the larger sigmas in Table 2. The large  $\sigma(M_1)$  (and hence  $\sigma(a)$ ) and  $\sigma(M_2)$  at the first epoch arose partly because this epoch was only 8 hours after launch and partly because of the complete absence of observations between 2<sup>h</sup> 43<sup>m</sup> on 6 May 1967 and 20<sup>h</sup> 18<sup>m</sup> on 7 May; there was a correlation of -0.988 between the computed values of  $M_1$  and  $M_2$ . Similarly, a large  $\sigma(M_1)$  arose for epoch 1967 DEC 5 because of the missing data for 3-4 December, which has already been mentioned. High sigmas for epochs from 1968 DEC 20 to 1969 JAN 16, inclusive, arose because of the paucity of observations during this period; with only 14 observations accepted, the run of 1969 JAN 1 gave the largest sigmas, for the parameters  $\Omega$  and  $\omega$ , of all the runs in Table 2. High sigmas for epochs from 1969 FEB 27 to 1969 APR 4 arose partly from paucity of observations and partly from the large values of & during that period; the epochs 1969 MAR 20 and 1969 MAR 23, for which the largest values of  $\epsilon$ (and of sigmas for e, i and  $M_0 + \omega$ ) of all the runs in Table 2 were obtained, have already been discussed in section 4.

#### 6 DISCUSSION

Although the perigee height of Ariel 3, during the period considered, was around 500 km, as opposed to about 300 km for Ariel 2, air drag was still important enough to be the chief limitation in the computation of orbital parameters by PROP. The value of  $M_2$ , equal to half the mean angular acceleration of the satellite, was only about  $0.02^{\circ}/d^2$  immediately after launch; when this parameter approached or exceeded  $0.1^{\circ}/d^2$ , as it did for a few days in October-November 1968 and for longer periods in 1969, or when it changed by more than about  $0.1^{\circ}/d^2$  from epoch to epoch, for example in late December 1967, the orbit does not fit the data so well, as indicated by higher values of  $\varepsilon$ . The period of validity of a set of orbital parameters is the same as the period spanned by the observations used in determining the parameters, i.e. between 3 and 4 days. If a set of orbital parameters is used to predict beyond the period of validity, then, when  $M_2$  is changing rapidly (and  $\varepsilon$  is large), error increases rapidly. Now since each set of orbital

parameters (after the first) was obtained by iteration from an initial set equivalent to the parameters at the preceding epoch, an immediate guide to the accuracy of three-day prediction - i.e. up to five days from a given epoch - is provided by the largest absolute value for the residuals in the first iteration of the orbit determination at the next epoch. This largest absolute value can change violently. As an extreme example, the figures for a series of successive epochs, starting at 1969 OCT 12, are:- 29, 45, 13, 7, 12, 67, 154, 189, 628, 52, 30, 59, 46, 17, 12, 74, 10; the value 628, equivalent to an angular error of about 10°, occurred for epoch 1968 NOV 5, and was obviously due to the unusually high value (at 1968 NOV 2) of 0.1144 for M<sub>2</sub>, which immediately afterwards fell to 0.0530. (A very large magnetic storm occurred on 1 November and had a devastating effect on the upper atmosphere 9.)

The behaviour of the orbital inclination, as evidenced by Fig.1, is worth discussing in detail. There are two distinct features. First, and very striking, is the secular behaviour: i remained essentially constant at 80.18° until the middle of December 1967, then dropped to about 80.163 in a period of about three months, and thereafter again remained essentially constant. Second, there are the superimposed oscillations, in which certain frequencies and amplitudes can fairly readily be seen. It is not entirely easy to explain either of these features.

Apart from resonances - and there should be no relevant resonance associated with the orbital parameters of Ariel 3 - the only known cause of secular variation in the orbital inclination of an earth satellite is the rotation of the atmosphere. Applying the formula of King-Hele and Scott  $^{10}$ , if the atmosphere at a height of 500 km is taken to rotate at twice the angular velocity of the earth ( $\Lambda$  = 2.0 in Ref.10), then i in Fig.1 should show a secular drop of about  $0.02^{\circ}$ , i.e. just about what it does show. The rate of drop should be proportional to  $M_1$ , however, whereas in Fig.1, as already remarked, the total drop is concentrated into a period of about three months, starting in December 1967.

The phenomenon is sufficiently remarkable for a sceptical reader to wonder whether the inclinations in Fig.1 really are right. Here the SDC values, though less accurate than the R.A.E. values, are useful; they are completely independent of the R.A.E. values, and confirm - not that there was a serious doubt - the secular behaviour indicated by the R.A.E. values.

Slightly more credible, though still unlikely, is the possibility that the sharp drop in i is not a purely secular effect but an oscillation superimposed on the change due to atmospheric rotation. Such an oscillation would require a period of 500 days or more, however, and even then the next cycle might have been expected to appear before the end of the graph. The amplitude of the oscillation would have to be nearly 0.01°. The author is unable to see whence such a term could arise.

There remain two possibilities: a single, complete discontinuity, due for example to meteoric impact, and a genuine (secular) perturbation over the (roughly) three-month period. The former seems very unlikely, though a discontinuity near the beginning of January 1968 cannot be completely ruled out; so we are left with the possibility of a real perturbation. Bearing in mind that a perturbing force, to produce an effect on i without affecting i, has to act in a direction perpendicular to the orbital plane, and that, to avoid cancellation, it has to act in opposite directions at the ascending and descending nodes of the orbit, it is difficult to see what the force can be, other than atmospheric rotation.

Attempts to produce an explanation must therefore degenerate into mere speculation. The three-month period of interest corresponded to a period of maximum solar activity (mean 107 mm solar radiation in excess of  $150 \times 10^{-22}$  W m<sup>-2</sup> Hz<sup>-1</sup>), and during this period M<sub>2</sub> was greater than before and after, though not enough greater to explain Fig.1 at once. Could it be that, at heights above 400 km, where no accurate measurements have been made, atmospheric rotation is significantly faster during periods of high solar activity, i.e. that King-Hele and Scott's  $\Lambda$  parameter is considerably less than 2.0 for most of Fig.1, but very much larger during the short period of maximum activity?

Apart from the correlation with solar activity, two other interesting (and unexplained) correlations should be mentioned. First, the direction of the spin axis of Ariel 3 has been monitored by RSRS, Slough<sup>1</sup>. At injection the spin axis pointed 69° south, i.e. to a point on the celestial sphere with declination -69°. During the first three months the axis looped towards the south, reaching declination -86° on 19 June 1967, but after this it moved north and the declination remained positive after August 1967, for as long as the satellite was still spinning. The axis looped towards the north, reaching a declination of almost 90° on 25 or 26 February 1968; from early

January to early April 1968, i.e. roughly the period of sharp fall in orbital inclination, was the period during which the declination of the spin axis exceeded 45°. (There was another period, starting about the middle of September 1968, when the declination again exceeded 45°, but this only lasted for about one month instead of three.) Second, the satellite's tape recorder, which contained two litres of air at a pressure of one atmosphere, was operating only intermittently during the period of interest. Until 28 November 1967 the tape recorder worked successfully. It then failed, but recovered and worked, apparently perfectly again, for three periods (of two, four and three weeks), until it failed for the last time on 14 April 1968. It is tempting to speculate that air was leaking, but, even if this was so, it is difficult to see how the right inclination-reducing force could result.

On turning to the oscillations, it is clear from Fig.1 that a number of components, of differing frequency, amplitude and phase, are present. Since cos i is small and e is very small, the oscillatory perturbations due to the earth's odd harmonics, the amplitude of which is proportional to e cos i, is completely negligible. The effect of the earth's tesseral harmonics, as remarked in section 3, is not negligible, and the amplitude of the  $J_{2,2}$ perturbation is more than twice as big as some of the values of  $\sigma(i)$  in Table 2; however, such effects should not appear in Fig.1, since they are averaged out during orbit determination. Hence the oscillatory components in Fig.1 may be thought of as being due solely to luni-solar perturbations, for which the various terms in di/dt are given in Ref.11 (equation (31)). The main term in the integral of the equation for di/dt is, for Ariel 3,  $0^{\circ}.0015$  cos  $2(u_{g} - \Omega)$ , where  $u_{g}$  is the argument of latitute of the sun; the period of this term is 80 days and a complete cycle may be seen, in Fig.1, for example between MJD 39673 and 39753 and between MJD 39993 and 40073. The next largest terms are combined terms for the sun and moon which, if we ignore the small non-zero value  $(\Omega_m)$  for the right ascension of the node of the moon's orbit, are given by  $0^{\circ}.0012 \cos \Omega$  and  $0^{\circ}.0007 \cos 2\Omega$ , of period 280 days and 140 days respectively. The fourth largest term is also the principal one in which  $u_m$ , the argument of latitude of the moon, appears; ignoring  $\Omega_{\rm m}$  again, it is 0°.0005 cos 2(u<sub>m</sub> -  $\Omega$ ), of period 12½ days approximately. Other terms are of smaller argument, but a combination of such terms could produce a detectable contribution to the graph of i.

In the absence of a spectral analysis or a complete analysis of all terms from Ref.11 it is difficult to be sure whether the oscillatory component of Fig.1 can be fully explained by luni-solar perturbations. It does appear, however, that Fig.1 contains a sinusoidal term with period about 30 days. A term in  $\cos (k\Omega \pm u_m)$ , for small integral k, would be appropriate here, but the only terms which arise have the eccentricity of the moon, i.e. 0.055, as a factor, and their amplitudes are too small.

One other known source of sinusoidal contributions to inclination variation should be mentioned. This is the precession and nutation of the earth's axis, which provides the reference with respect to which (the complement of) an orbital inclination is measured. The main contribution is from precession and may be taken from Ref.12 (in which the nutation terms have the wrong sign but the precession term is correct). For Ariel 3 this gives  $-0^{\circ}.0007 \cos \Omega$ , and so reduces (to about half) the amplitude of the direct luni-solar perturbation term of argument  $\Omega$ .

#### 7 CONCLUSIONS

Orbital parameters for the satellite Ariel 3, as for Ariel 2, have been determined at R.A.E., at three-day intervals, from Minitrack observations supplied by NASA. The accuracy of the computed parameters is about the same, in general, as was obtained for Ariel 2, i.e. better than 1 km except on very rare occasions, and should be adequate for experimenters' requirements. (Values of eccentricity were accurate enough to be used in determining the odd harmonics in the geopotential 13.) During periods of high drag, better accuracy could have been obtained by inclusion of an eighth parameter, M<sub>3</sub>, in the orbital model, as was done with Ariel 2.

Oribital inclination was determined rather less accurately than for Ariel 2, no doubt due to the fact that there are no Minitrack stations, in either hemisphere, at latitudes as high as 80°. The accuracy was good enough, however, for an anomalous secular behaviour in inclination to be clearly observable. This behaviour has been discussed but not explained.

#### Acknowledgement

The author wishes to thank Jennifer Davies for preparing the card decks and supervising the computer analysis.

Table 1

MINITRACK STATIONS OBSERVING ARIEL 3

Station name	Location	x (km)	y (km)	z (km)
Fort Myers	Fort Myers, Florida, U.S.A.	807.885	-5652.020	2833.549
Johannesburg	Hartebeshoek, South Africa	5084.798	2670.474	-2768.164
Lima	Lima, Peru	1388.818	-6088.429	-1293.207
Newfoundland	St. Johns, Newfoundland	2602.801	-3419.184	4697.694
Quito	Quito, Ecuador	1263.617	-6255.010	-68.856
Santiago	Santiago, Chile	1769.707	-5044.642	-3468.192
Winkfield	Winkfield, England	3983.130	-48.404	4964.711
Ulaska	Fairbanks, Alaska	-2282.332	-1452.667	5756.942
Madagascar	Tananarive, Malagasy	4091.903	4434.373	-2064.537
Orroral	Canberra, Australia	-4447.361	2677.215	-3695.209

Table 2. Orbital parameters of Ariel 3

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Table 2(cont'd.) Orbital parameters of Ariel 3

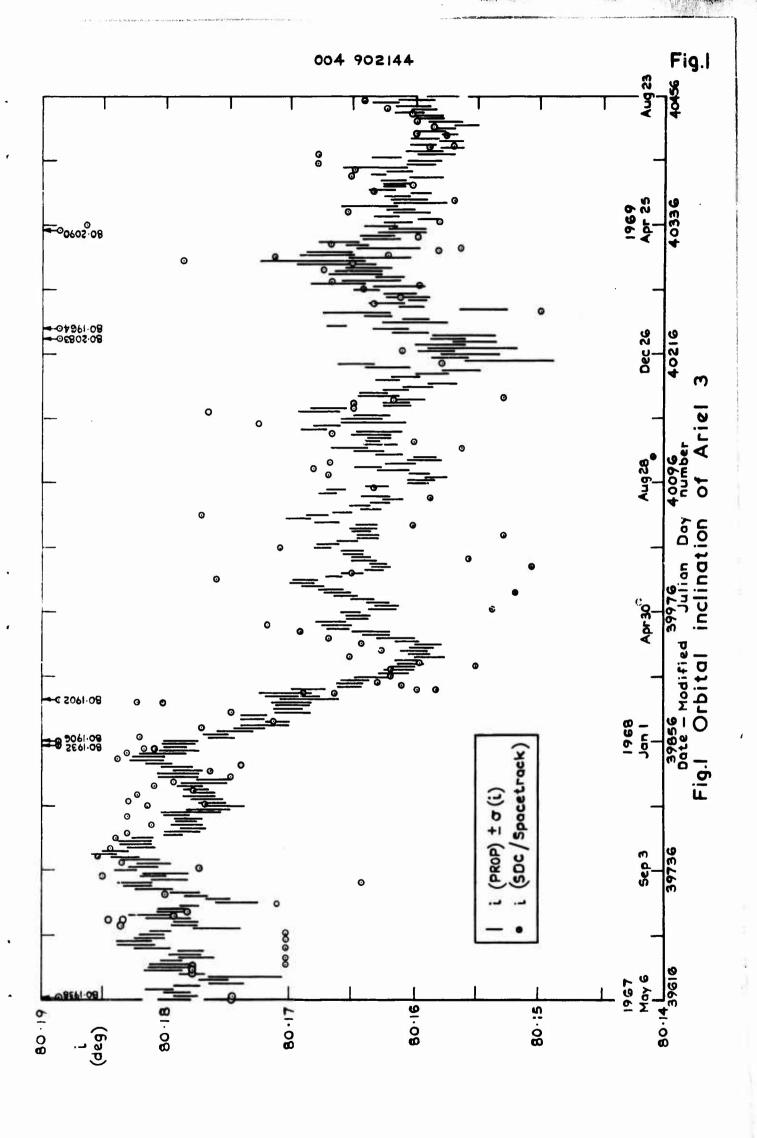
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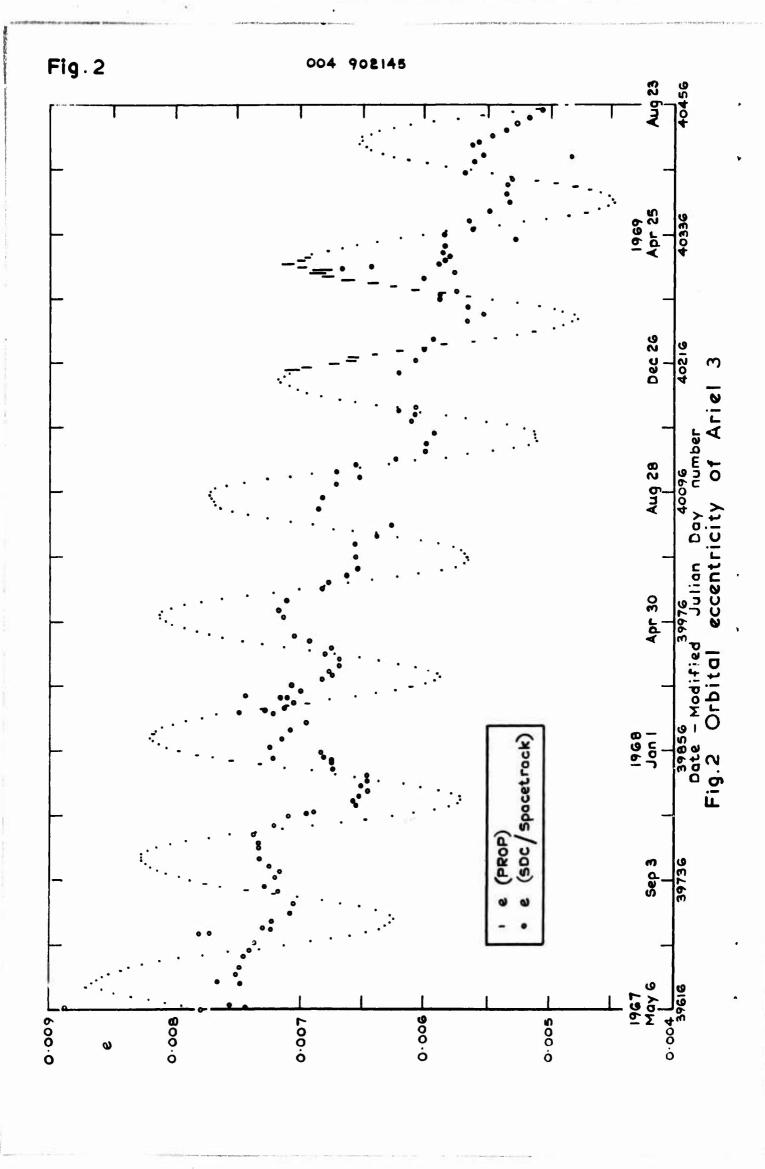
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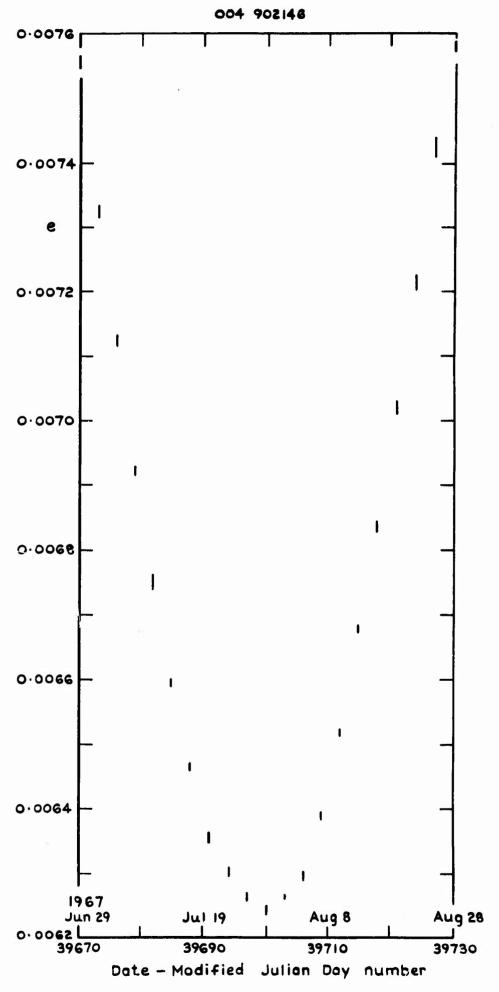
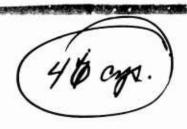


Fig. 3 Eccentricity of Ariel 3 during 60 days, on expanded scale



## ROYAL AIRCRAFT ESTABLISHMENT

Technical Report 69275

December 1969

THE ORBIT OF ARIEL 3 (1967-42A)

by

R. H. Gooding

#### ADDENDUM

The striking change of i (orbital inclination), from  $80.18^{\circ}$  in December 1967 to about  $80.163^{\circ}$  in March 1968, was remarked upon in section 6, but left unexplained. It is believed that the explanation is now known. During this period Ariel 3 was passing through a resonance associated with the earth's tesseral harmonics of order 15. In fact the mean motion of the satellite was exactly commensurate with the earth's rotation rate just before 0 hours on MJD 39889 (3 February 1968), and the variable M +  $\omega$  + 15 ( $\Omega$  - n t) of Ref.14 varied by no more that  $120^{\circ}$  from its resonance value, during the three-month period. The author, in section 6, dismissed resonance, thinking of luni-solar resonance rather than tesseral-harmonic resonance, and discussed alternative explanations which must now be considered irrelevant.

#### ERRATA

Page 7, line 3: for the first 'M<sub>1</sub>' read 'M<sub>1</sub>'.

Page 14, line 13: for the second 'i' read 'M<sub>1</sub>'.

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