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MEMORANDUM REPORT NO. 1709

BASE PRESSURE MEASUREMENTS ON SHARP AND
BLUNT 9° CONES AT MACH NUMBERS FROM 3.50 TO 9.20

by

Neil A. Zarin

November 1965

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NOVEMBER 1965

1. Cones - Base pressure
2. Base pressure measurements
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AT MACH NUMBERS FROM 3.50 TO 9.20

Neil A. Zarin

Exterior Ballistics Laboratory

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ABERDEEN PROVING GROUND, MARYLAND

B A L L I S T I C R E S E A R C H L A B O R A T O R I E S

MEMORANDUM REPORT NO. 1709

NAZarin/gk
Aberdeen Proving Ground, Md.
November 1965

BASE PRESSURE MEASUREMENTS ON SHARP AND BLUNT 9° CONES
AT MACH NUMBERS FROM 3.50 TO 9.20

ABSTRACT

Base pressure measurements were made on sharp and hemispherically blunted 9° cones at Mach numbers from 3.50 to 9.20. The tests were carried out in the Ballistic Research Laboratories' Supersonic and Hypersonic Wind Tunnels at Aberdeen Proving Ground, Maryland. The data obtained are compared to experimental data and to data from semiempirical analyses from other sources. An empirical correlation for the base pressure data is presented. The relative contributions of base and form drag to total drag are compared.

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DEFINITION OF SYMBOLS

C_D	Drag coefficient = $\frac{D}{q S}$
D	Drag force
k	Constant used in base pressure coefficient correlation equation
ℓ	Model wetted length
M_∞	Free stream Mach number
p_b	Model base pressure (average of 4 taps)
P_b	Base pressure coefficient = $\frac{p_b - p_\infty}{q}$
p_c	Cone surface pressure (measured 0.60 in. ahead of base)
p_∞	Free stream static pressure
q	Dynamic pressure = $\frac{1}{2}\rho V^2$
r	Nose radius of blunted cone
R	Base radius of cones
Re_L	Reynolds number based on wetted length and local conditions at outer edge of boundary layer just ahead of base of body
Re_∞	Reynolds number based on wetted length and free stream conditions
S	Model reference area = 7.069 square inches
V	Free stream air velocity
ρ	Mass density

1. INTRODUCTION

While there has been much interest in the subject of base pressure in the past several years, there has been a noticeable lack of experimental data at hypersonic velocities. Many investigators have obtained a wealth of data at velocities up to Mach 5, and there have been several semiempirical theories advanced; however, little has been done above Mach 5.

The present limited investigation was initiated in order to begin to fill the void of hypersonic base pressure data, to evaluate our ability to accurately measure low pressures, to determine the effect of Mach number, Reynolds number, and boundary layer trip devices on base pressure, and to compare our test results with other experimental data, at least at the lower Mach numbers.

The present tests were considered a success, and further tests on bodies of different shapes are being planned.

2. MODELS AND APPARATUS

2.1 Wind Tunnels

The tests were conducted in Supersonic Wind Tunnel No. 1 and Hypersonic Wind Tunnel No. 4. The supersonic tunnel is of the continuous flow, closed circuit, variable density type and has a flexible nozzle for obtaining a range of Mach numbers from 1.50 to 5.00. The test section size is 13 inches wide by 15 inches high. The hypersonic tunnel is of the continuous flow, open jet, closed circuit, variable density variety. It has interchangeable axisymmetric nozzles for Mach numbers 6.0, 7.5, and 9.2 with exit diameters of 14.6, 15.6, and 18.7 inches, respectively. A combustion and an electric heater provide stagnation temperatures up to 1960° Rankine--sufficient to prevent air liquefaction. The specific humidity was maintained at a value less than 0.0002 lb of water vapor per pound of air for all tests. Further information on the tunnels may be found in Ref. 1.

2.2 Instrumentation

The base pressures were transmitted to four 0-1 psi Statham absolute pressure transducers which were located outside the test section. The sensitivity of the transducers was increased by using a supply voltage of 6.0 volts rather than the design voltage of 3.5 volts. The transducers then had a range

of 0-0.6 psia with an accuracy of better than ± 0.25 percent of their range. The cone pressure was measured on a 0-5 psi Statham absolute pressure transducer whose sensitivity had been increased to give it a 0-3 psi range. Its accuracy was also better than ± 0.25 percent of its range. The transducers were zero referenced to a vacuum system which measured less than 0.025 mm Hg at all times.

The electrical signals from the transducers were converted by the automatic data readout system to proportional digital readings which were typed on data sheets and punched in code on a tape. A schlieren system with camera provided continuous visual indication, as well as photographs, of the flow conditions in the test section. Spark shadowgraph photos were also taken.

2.3 Models

The models tested were sharp and hemispherically blunted 9° half angle cones. They were tested both with and without a square trip ring of height and width 0.050 inch, which had an inside diameter of 1.25 inches. The purpose of the trip ring was to artificially induce a turbulent boundary layer on the model. The sharp nose model was 9.471 inches long and had a 3.000-inch base diameter. The blunt nosed model used the same base and was 7.160 inches long. It had a nose radius of 0.429 inch. There were four base pressure taps located on a vertical diameter at varying distances from the model axis. One cone surface pressure tap was located on the top of the model, 0.600 inch from the base. The model physical characteristics are illustrated in the drawing, Figure 1.

3. TEST PROCEDURE

With the model installed on the tunnel centerline, flow was established at the desired Mach number, but at reduced stagnation pressure and temperature. Next, pressure and temperature were increased to their proper value for the test, and schlieren and shadowgraph photographs were taken. Then, pressure data were taken at one minute intervals until there was no noticeable change in pressure with time. After that, tunnel flow conditions were changed and the process repeated. Data were taken at zero angle of attack only.

4. DATA REDUCTION

4.1 Procedure

The raw numerical data from the typed data sheets were reduced to gage pressure by using the measured transducer calibration constants. The reference pressure was then added to these values to yield absolute pressures. The four base pressure readings were averaged arithmetically to yield an average base pressure, p_b . This pressure was reduced to coefficient form by subtracting from it the test section static pressure, p_∞ , and then dividing the difference by the dynamic pressure, q . The local Reynolds number, Re_L , was obtained by determining the local Mach number at the base of the model and then using Chart 25 in NACA Report 1135, using stagnation conditions in the case of the sharp cone and total conditions behind a normal shock in the case of a blunt cone.

4.2 Accuracy

The maximum deviation of any of the measured base pressures from the average values does not exceed ± 0.002 psi, which is ± 0.25 percent of the range of the transducers. The cone surface pressure readings were reduced in a similar manner and were also found to be consistent with transducer accuracy. The maximum error in base pressure coefficient due to transducer inaccuracy is ± 0.0025 . The range of p_b measured was from 0.0035 to 0.165 psia.

5. PRESENTATION OF DATA

The data from the wind tunnel tests are presented in several different ways in order to better illustrate certain trends, and to compare with the theoretical and experimental work of others.

In Figs. 2-5, the base pressure coefficient, P_b , is plotted as a function of free stream Reynolds number, Re_∞ , based on model wetted length. Each figure is for a different configuration, with a curve for each Mach number. In Figs. 6-12, we also plot P_b vs Re_∞ ; however, in these, each figure is for a different Mach number, with a curve for each configuration.

In Figure 13, P_b is plotted versus local Reynolds number, Re_L , based on conditions just outside the boundary layer at the model base and the model

wetted length. The points plotted are for cases where boundary layer transition has occurred before the model base. Each curve is for a different Mach number.

Fig. 14 shows the present data for sharp and blunt cones, without trip rings, plotted as a ratio of base pressure to cone surface pressure, P_b/P_c , vs Re_L . This is compared to data of Whitfield and Potter from Ref. 2, which is also for flow over sharp and blunt 9° cones. Fig. 15 is similar to Fig. 14, the only difference being that the present data shown are for models with the trip ring. It is again compared to Whitfield and Potter's data which are for models without boundary layer tripping devices.

Fig. 16 shows P_b plotted as a function of free stream Mach number, M_∞ . The points shown are for a sharp cone with a trip ring at a free stream Reynolds number of 6×10^6 . The data are compared to a compilation by Chapman, found in Ref. 3, for flow over axisymmetric models with cylindrical afterbodies and data for a $3/4$ power law body from Ref. 4 by Reller and Hamaker.

In Figs. 17 and 18, drag coefficient is plotted against Mach number for the sharp and blunt cones, respectively. The contributions of base and wave drag are compared. The contribution of friction drag is small and has been neglected.

Fig. 19 shows the ratio of base pressure to free stream static pressure plotted as a function of free stream Mach number. An "approximate estimate" from Ref. 2 for a 9° blunt cone with $r/R = 0.3$ and $Re_L \geq 40 \times 10^6$ is compared to present data where the cone angle is 9° , $r/R = 0.286$ and Re_L is sufficient to insure a turbulent boundary layer on the model.

6. DISCUSSION OF RESULTS

In Figs. 2-5, we see the effect of free stream Reynolds number and Mach number on the base pressure coefficient for the different configurations tested. We can see that, with increasing Reynolds number, the base pressure coefficient decreases (becomes more negative), and the base drag increases. This is due to the fact that at the higher Reynolds number the boundary layer becomes turbulent on the body and improves the mixing in the base region.

The base pressure coefficient increases with increasing Mach number, and the base drag therefore decreases.

The difference between the sharp and blunt cones and the effect of the boundary layer trip ring can best be seen in Figs. 6-12. Here we see that transition and the values of P_b characteristic of turbulent boundary layers occur at lower values of Re_∞ for the models with the trip ring. We may also note that once turbulent flow on the model is obtained, there is very little dependence on whether transition occurred naturally or by means of the trip ring. This is more evident in the case of the sharp cone since, at some of the Mach numbers tested, it was not possible to obtain turbulent flow on the blunt cone without the trip ring.

Fig. 13, which shows P_b plotted against local Reynolds number, shows that, for the turbulent data, an empirical correlation between the sharp and blunt cones is possible. For the present tests, the equation

$$-P_b = 0.00483 \log Re_L + k$$

best describes the data, where "k" varies with Mach number as follows:

M	k
3.50	0.0643
4.00	0.0445
4.50	0.0289
5.00	0.0158
6.00	0.0033
7.50	-0.0102
9.20	-0.0157

Fig. 14 shows the ratio of base pressure to cone pressure plotted as a function of local Reynolds number for sharp and blunt cones without trip rings. In addition to the present data, those of Whitfield and Potter, found in Reference 2, are also plotted. From the discussion of Figures 2-5 and from Figure 14, it is clear that local Reynolds numbers for the present tests were not high enough to produce a turbulent boundary layer over the model in the majority of cases. In the cases where there were turbulent boundary layers in both the present tests and those of Ref. 2, agreement between the two sets

of data is fairly good. The difference is primarily that of transition Reynolds number. This can be explained by the fact that model surface roughness and tunnel turbulence level, which are undoubtedly different for the two sets of data, have a significant effect.

In Fig. 15, which is similar to Fig. 14, the present data for the models with the trip ring are compared to those of Whitfield and Potter without a tripping device. The present data seem to level off to values fairly close to the higher Reynolds number data of Whitfield and Potter. Little more can be said about this data, due to the very large scatter. This scatter is believed to be due to the inability to obtain very high Reynolds numbers. Thus, it fixes transition well forward of the base on the model.

Fig. 16 shows base pressure coefficient as a function of Mach number. Present data on the sharp cone with trip ring, at Re_L of 6×10^6 , is compared to compilations by Chapman and Love, found in Refs. 3 and 5, respectively, for axisymmetric models with cylindrical afterbodies. Data from Ref. 4, by Reller and Hamaker on a model with a $3/4$ power law shape, is also shown. The present data compares quite well with the Chapman and Love curve at Mach 5.0 and above. Below Mach 5.0, there is good agreement between present tests and the work of Reller and Hamaker. The difference between the present data and the Chapman and Love compilation may be explained by the fact that all of their data were for models with cylindrical afterbodies. The base pressure on a cone would naturally be lower than that on a model with a cylindrical afterbody. The close agreement between the present data and the data of Reller and Hamaker on the $3/4$ power law body serves to illustrate this point.

Figs. 17 and 18 show the drag contributions for the sharp and blunt cones, respectively. Friction drag was found to be small and was considered negligible. The pressure, or "fore," drag contribution was calculated by inviscid cone theory for the sharp cone, and by modified Newtonian theory for the blunt cone. The base drag was obtained from the present data for turbulent flow, for both sharp and blunt configurations.

For the sharp cone at $M = 3.5$, base drag is about 57 percent of the total drag, while at $M = 9.2$ it is only 19 percent of the total drag. In the case

of the blunt cone, these percentages are 41 and 7 percent, respectively. Thus, we see that for the more slender, streamlined shapes, base pressure is extremely important.

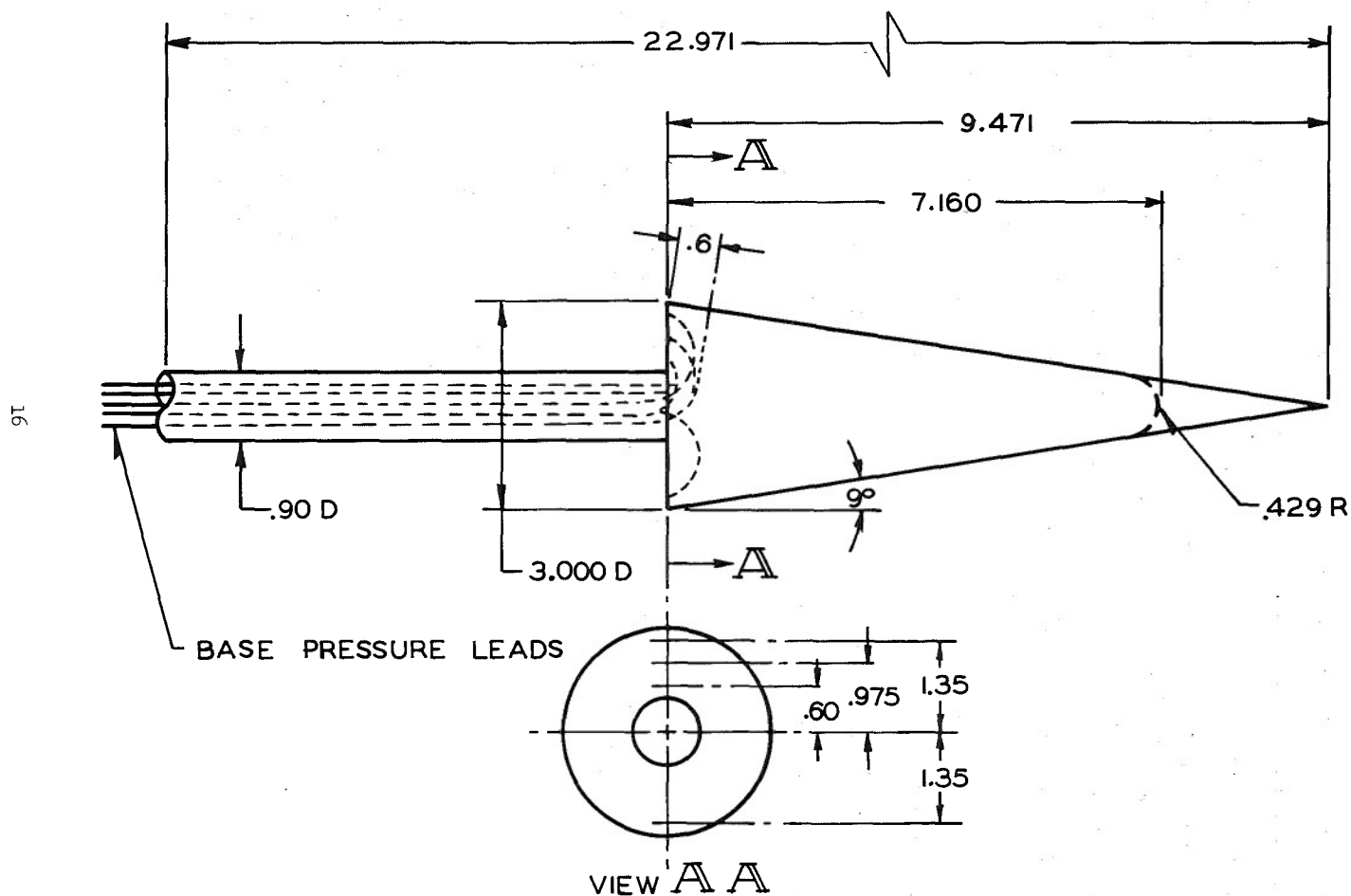
Figure 19 shows the ratio of base pressure to free stream static pressure plotted versus free stream Mach number. Present data for a blunt 9° cone with $r/R = 0.286$ and Re_L , sufficiently large for a turbulent boundary layer on the model, is compared to an "approximate estimate" made by Whitfield and Potter in Reference 2 for a blunt 9° cone with $r/R = 0.3$ and $Re_L \geq 40 \times 10^6$. The estimate of Whitfield and Potter was made to Mach 20 from data which went only to Mach 5.1, and the present data are in clear disagreement with it above Mach 4.5. At this Mach number and below, there is fairly good agreement between the present data and the estimate. It is felt that the estimate given in Reference 2 may lead to erroneous conclusions if used above Mach 4.5.

Further investigations on hypersonic base pressures are planned, and it is hoped that they will reinforce the present data and conclusions.

ACKNOWLEDGMENT

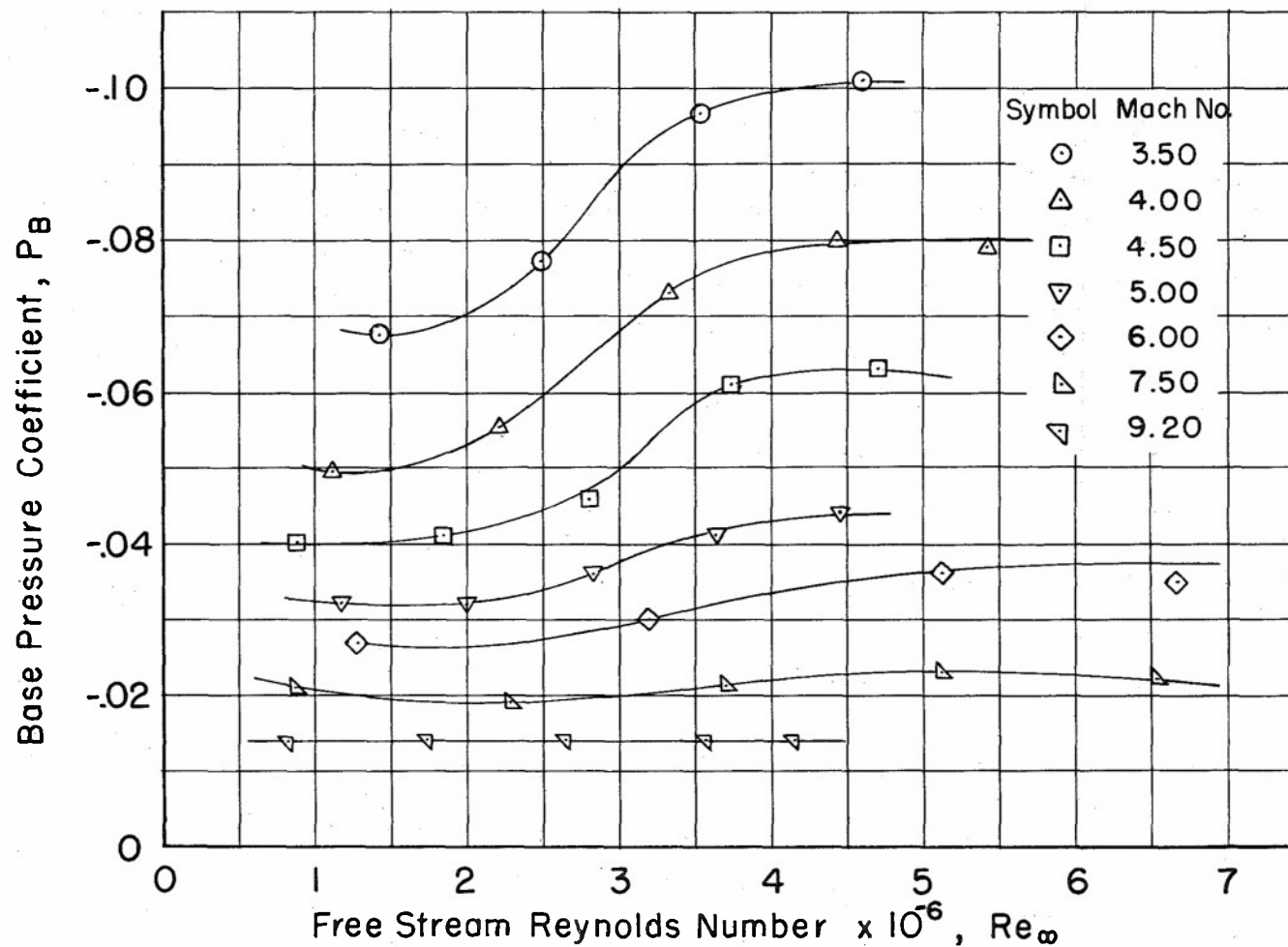
The author wishes to express his thanks to Mr. Robert H. Krieger, Chief, Wind Tunnel Testing Section, Supersonic Wind Tunnels Branch, for suggesting the program and for giving his invaluable help and support.

NEIL A. ZARIN



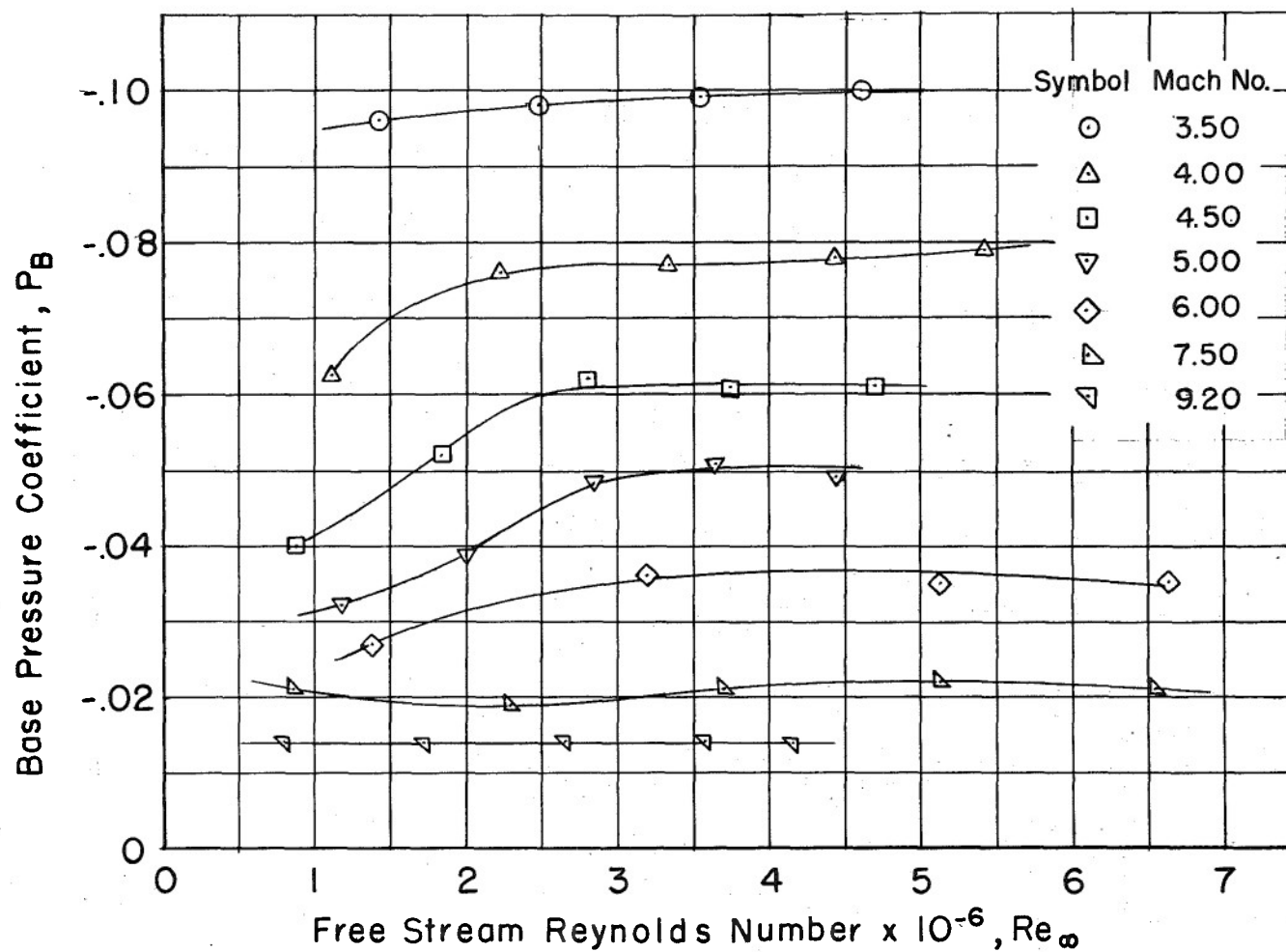
BASE PRESSURE MODELS

Fig. 1



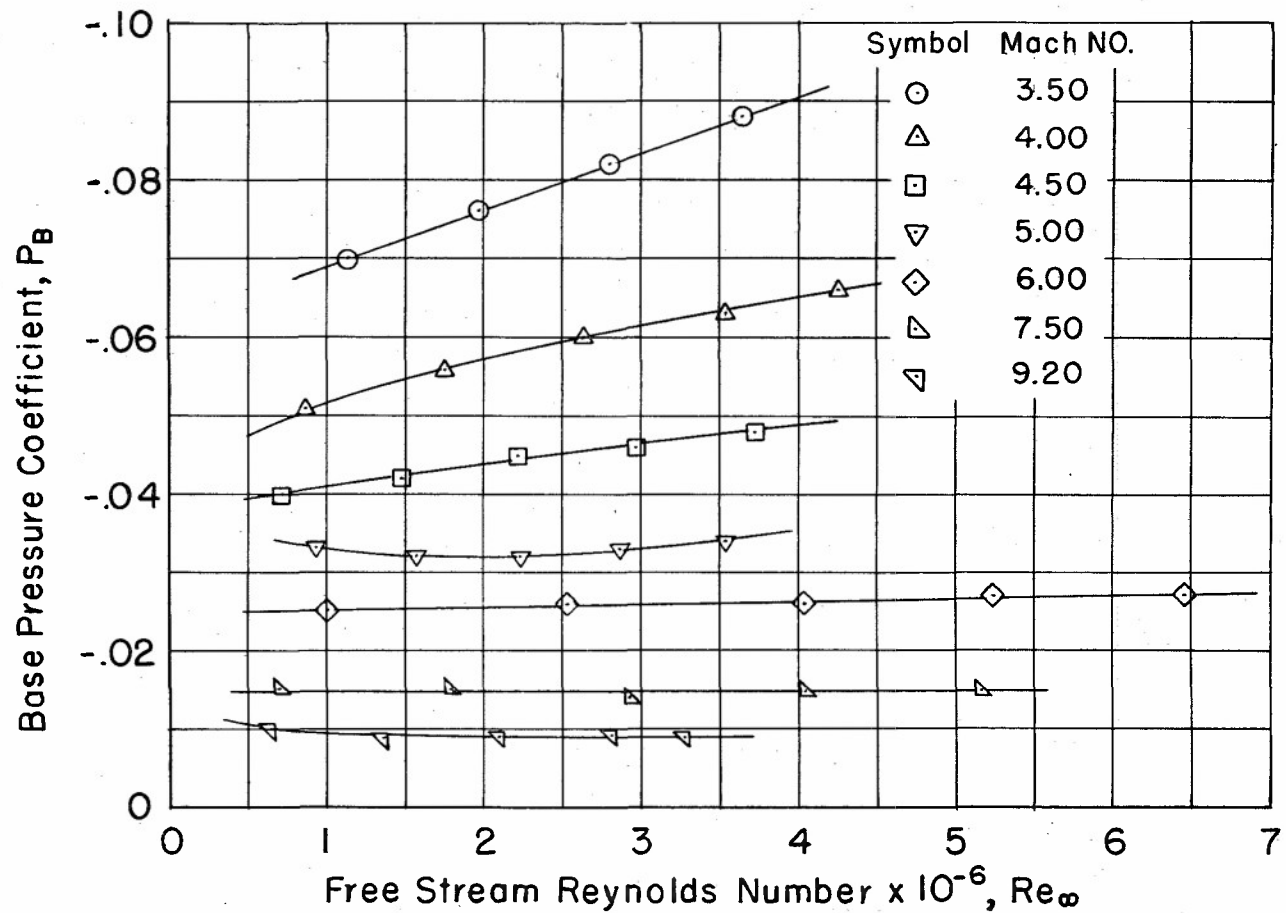
BASE PRESSURE COEFFICIENT VS. REYNOLDS NUMBER FOR
A SHARP 9° CONE WITHOUT TRIP RING

Fig.2



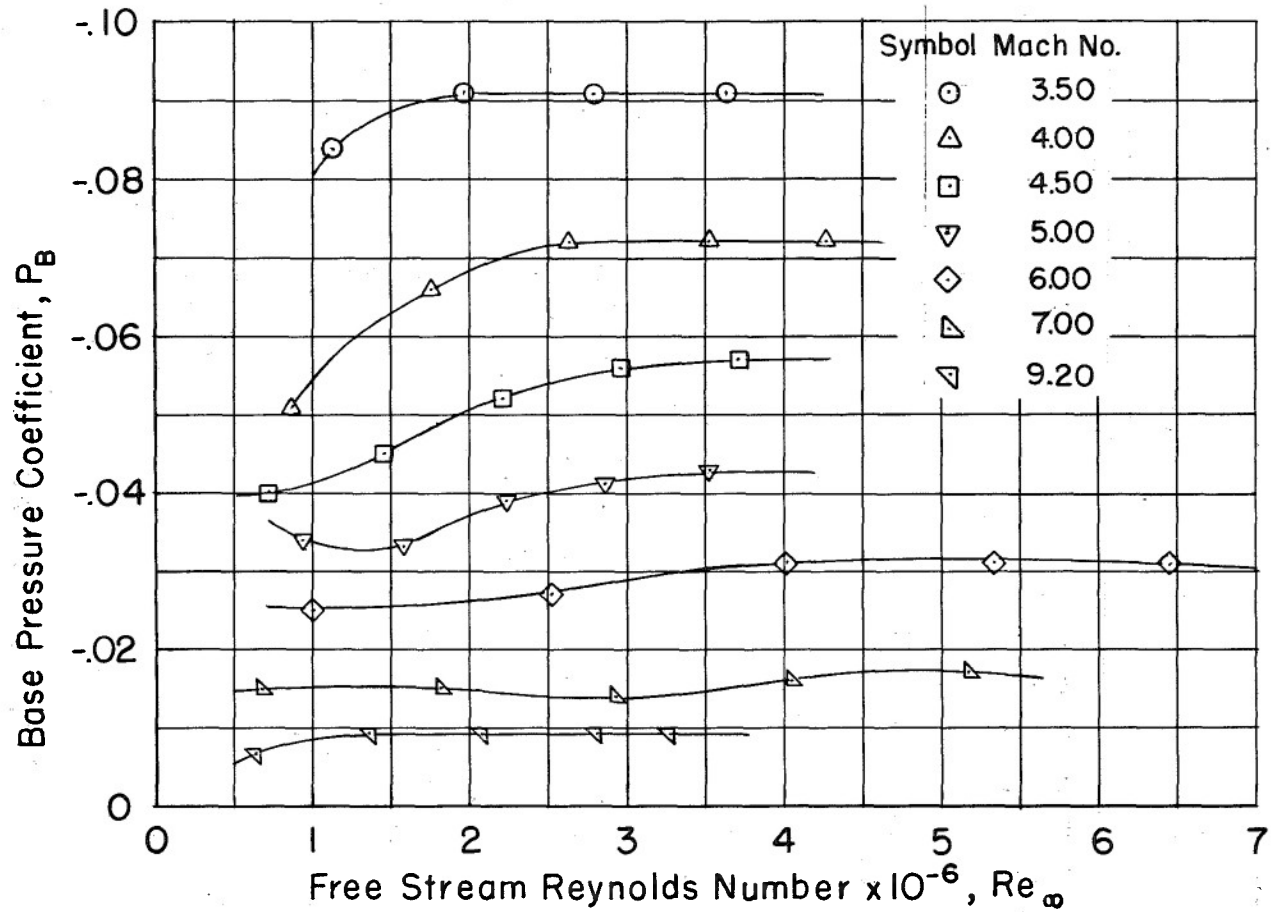
BASE PRESSURE COEFFICIENT VS REYNOLDS NUMBER FOR
SHARP 9° CONE WITH BOUNDARY LAYER TRIP RING

Fig. 3



BASE PRESSURE COEFFICIENT VS FREE STREAM REYNOLDS
NUMBER FOR A BLUNT 9° CONE WITHOUT TRIP RING

Fig.4



BASE PRESSURE COEFFICIENT VS REYNOLDS NUMBER FOR
BLUNT 9° CONE WITH BOUNDARY LAYER TRIP RING

Fig. 5

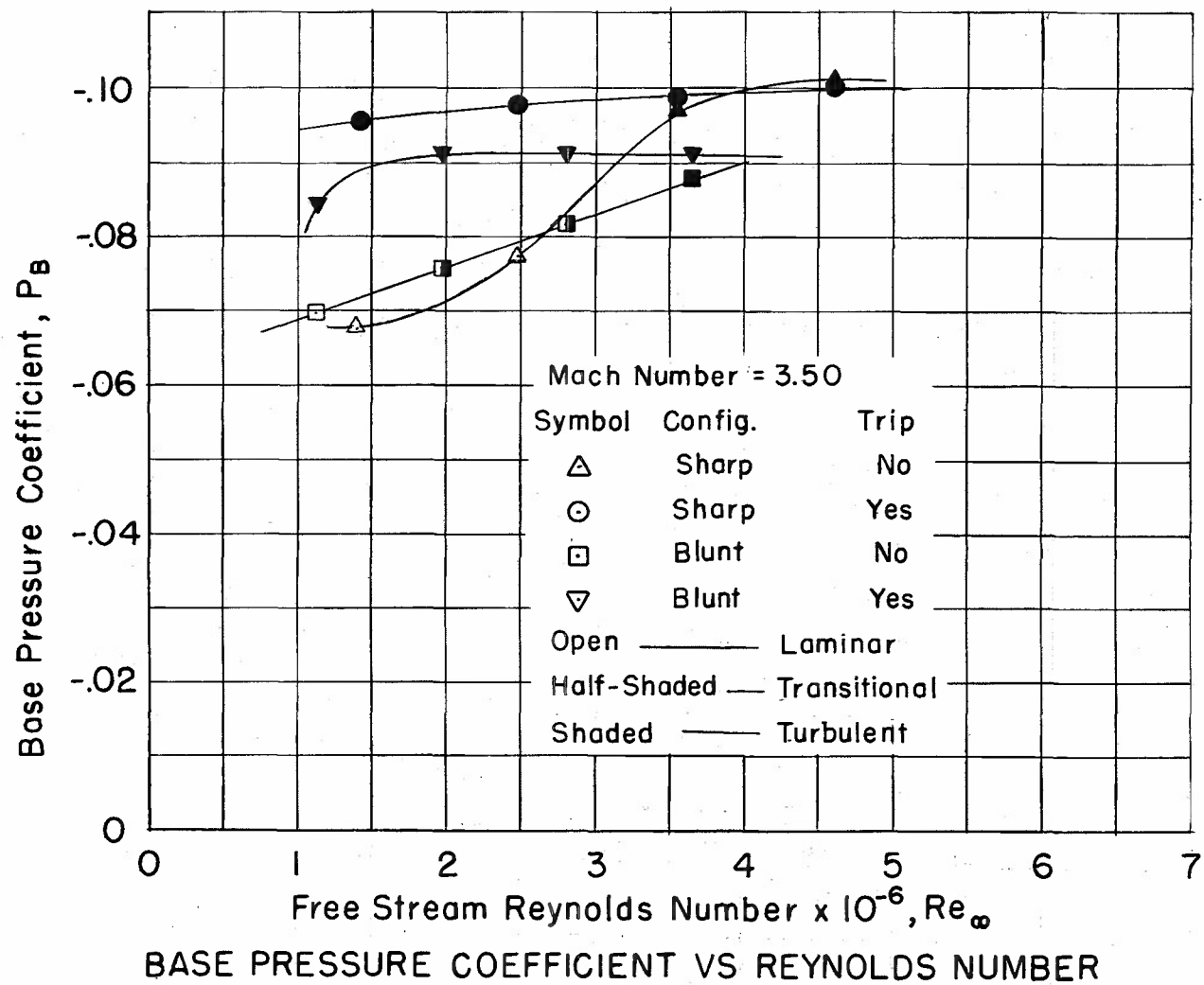
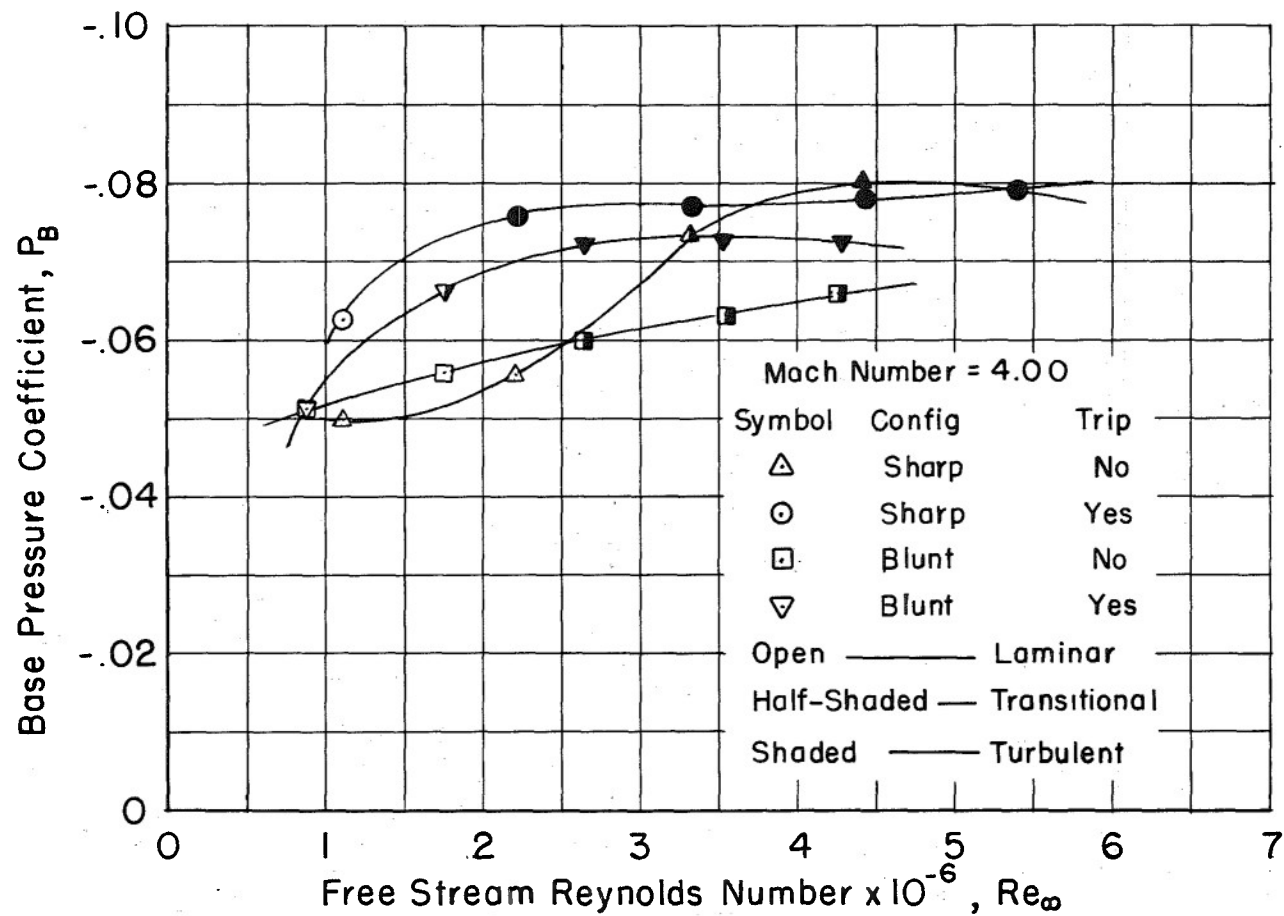
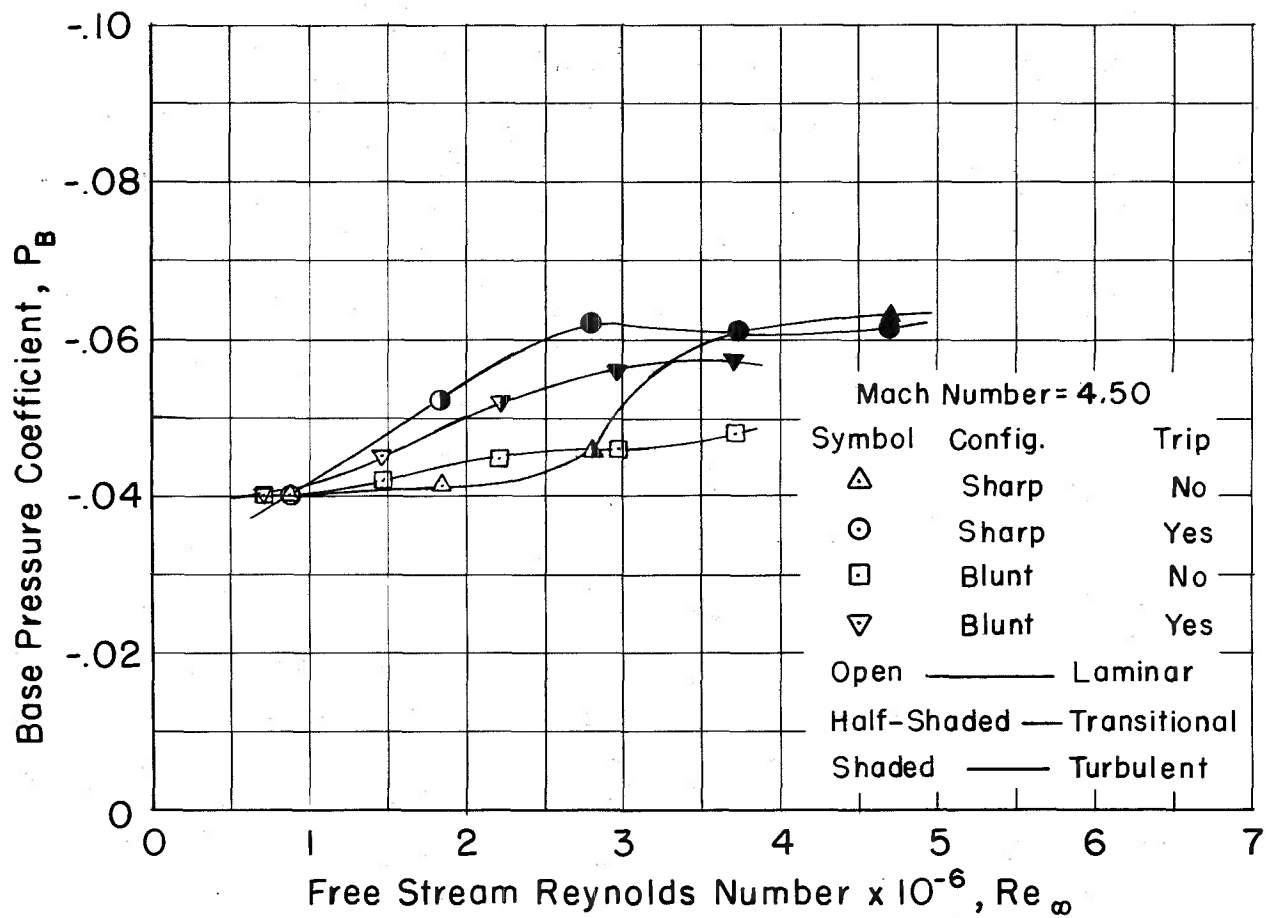


Fig. 6



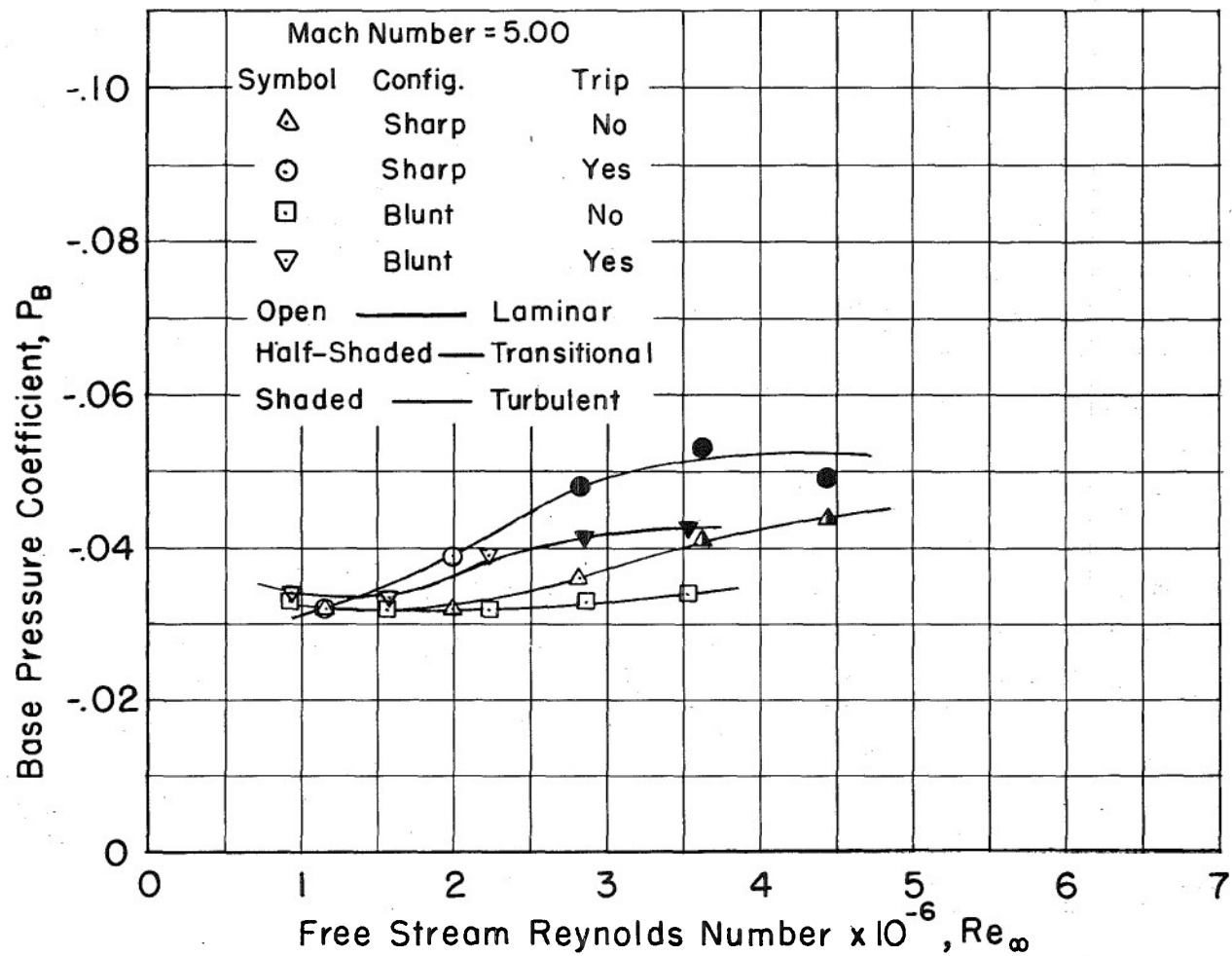
BASE PRESSURE COEFFICIENT VS REYNOLDS NUMBER

Fig. 7



BASE PRESSURE COEFFICIENT VS REYNOLDS NUMBER

Fig.8



BASE PRESSURE COEFFICIENT VS REYNOLDS NUMBER

Fig. 9

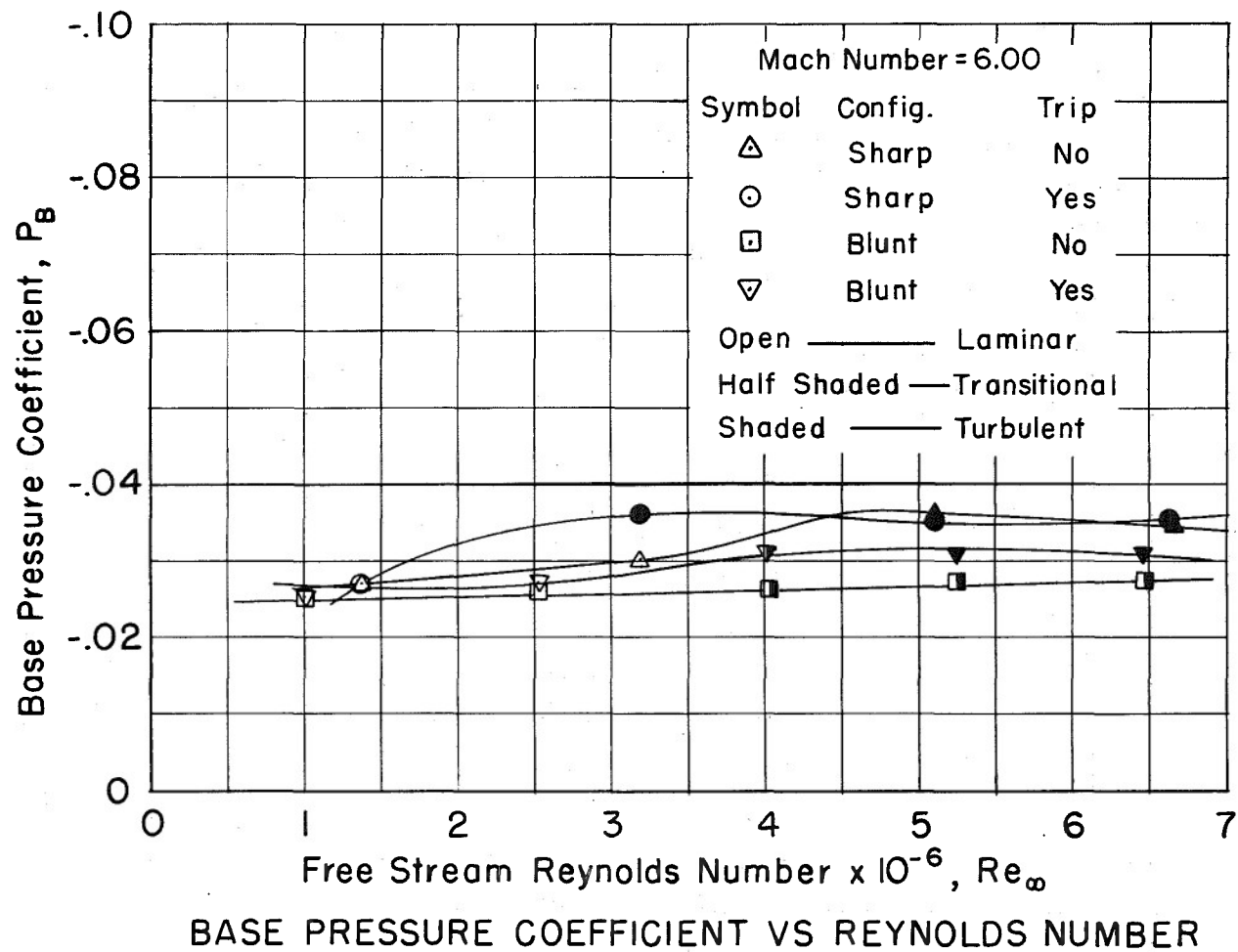


Fig.10

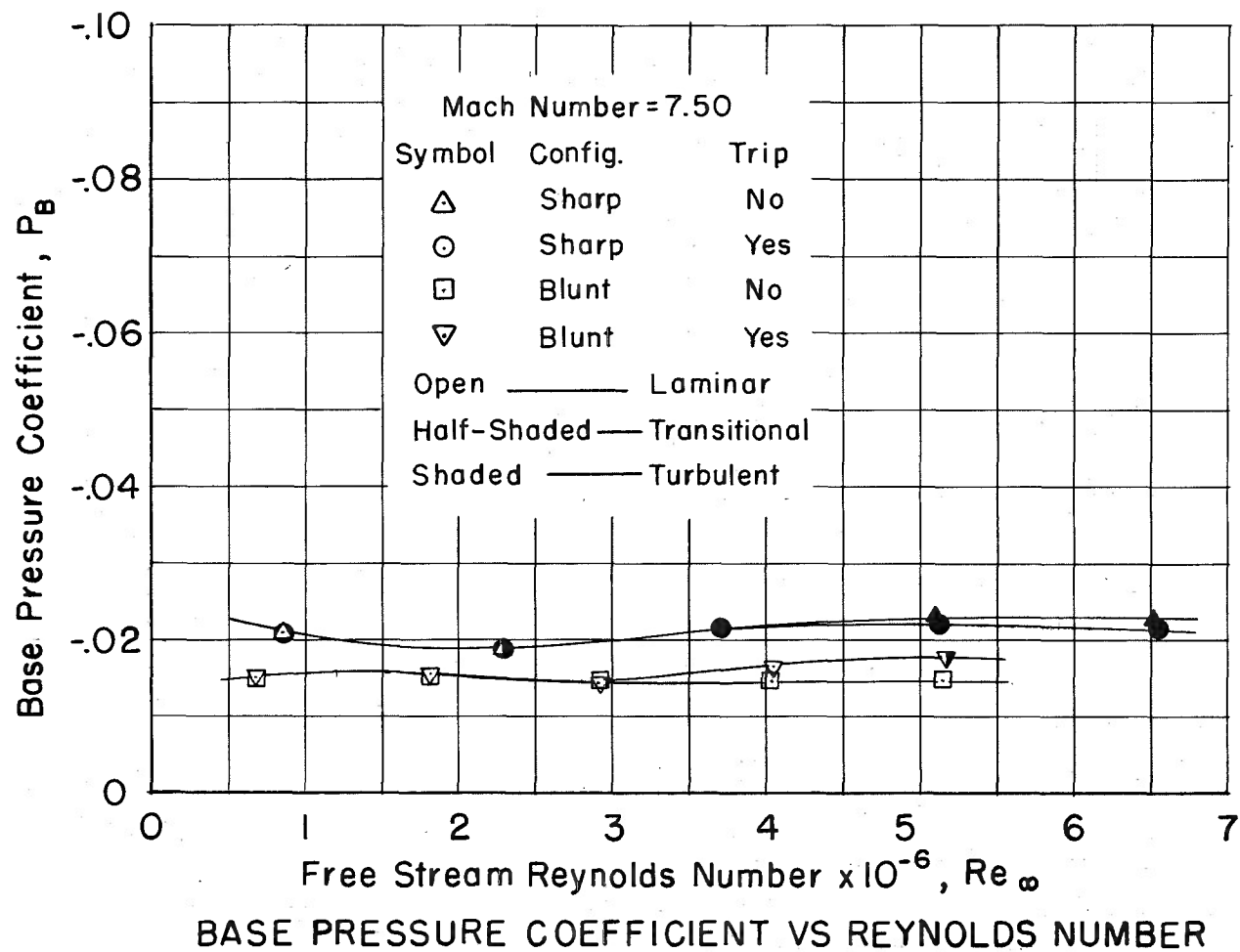
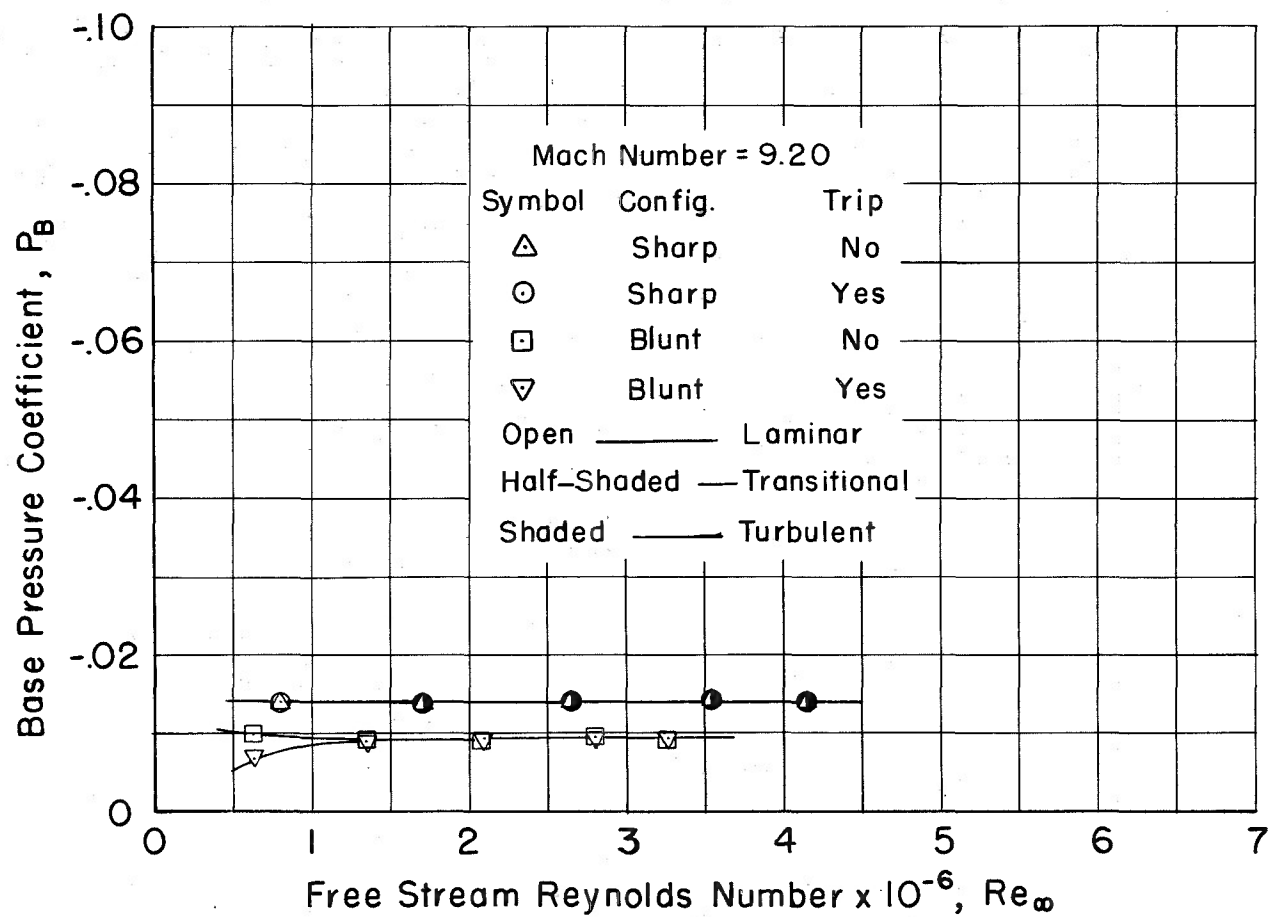
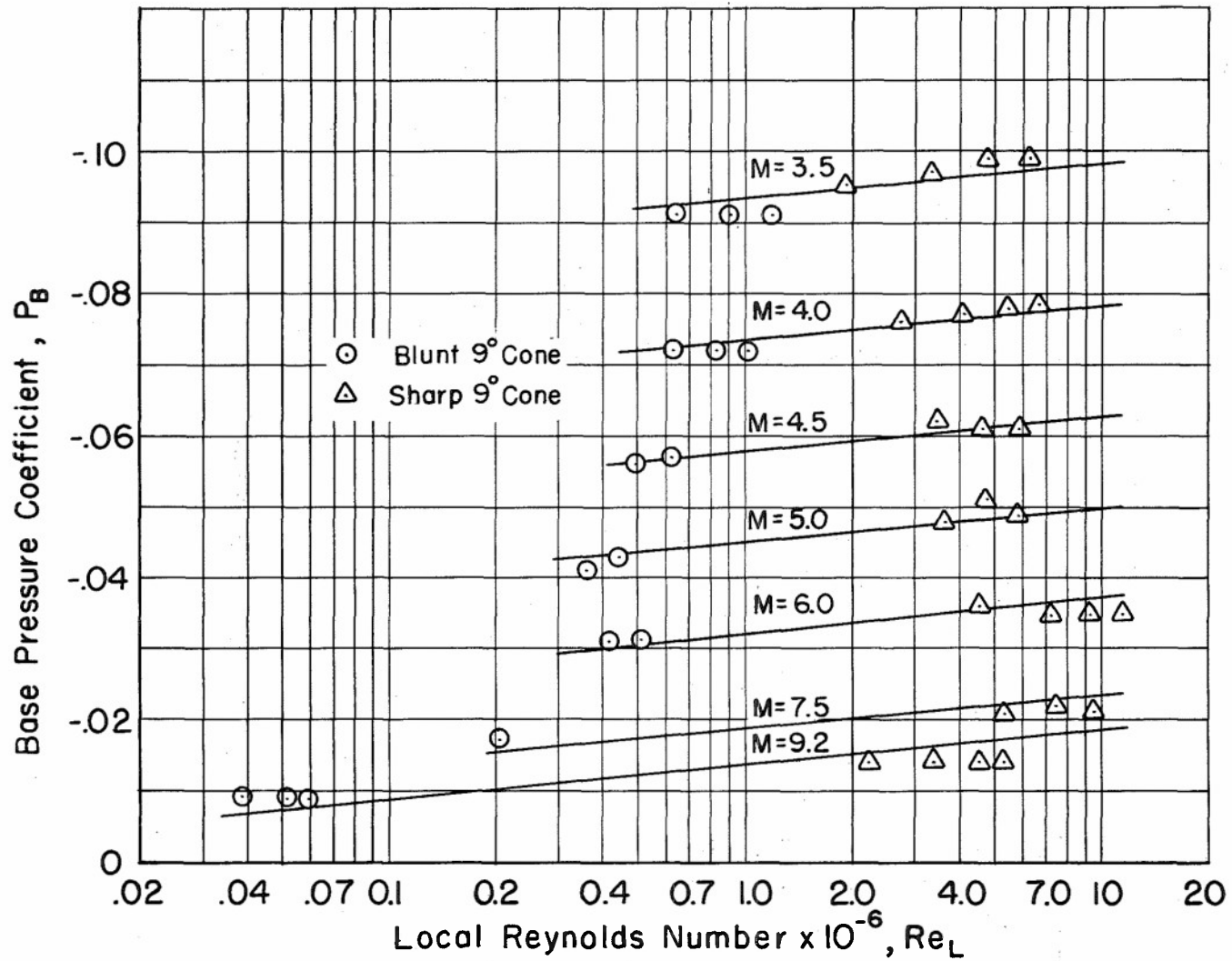


Fig. II



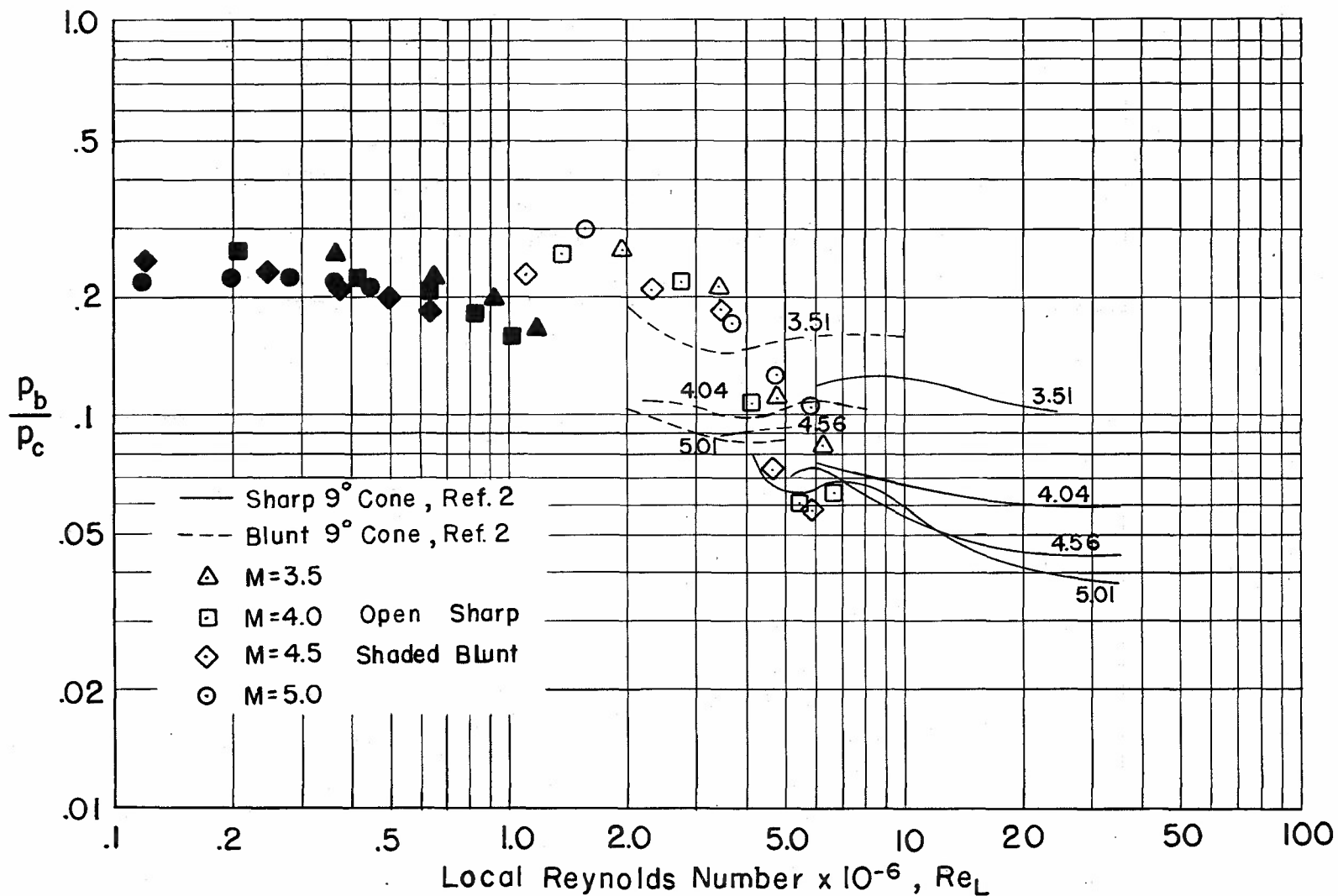
BASE PRESSURE COEFFICIENT VS REYNOLDS NUMBER

Fig.12



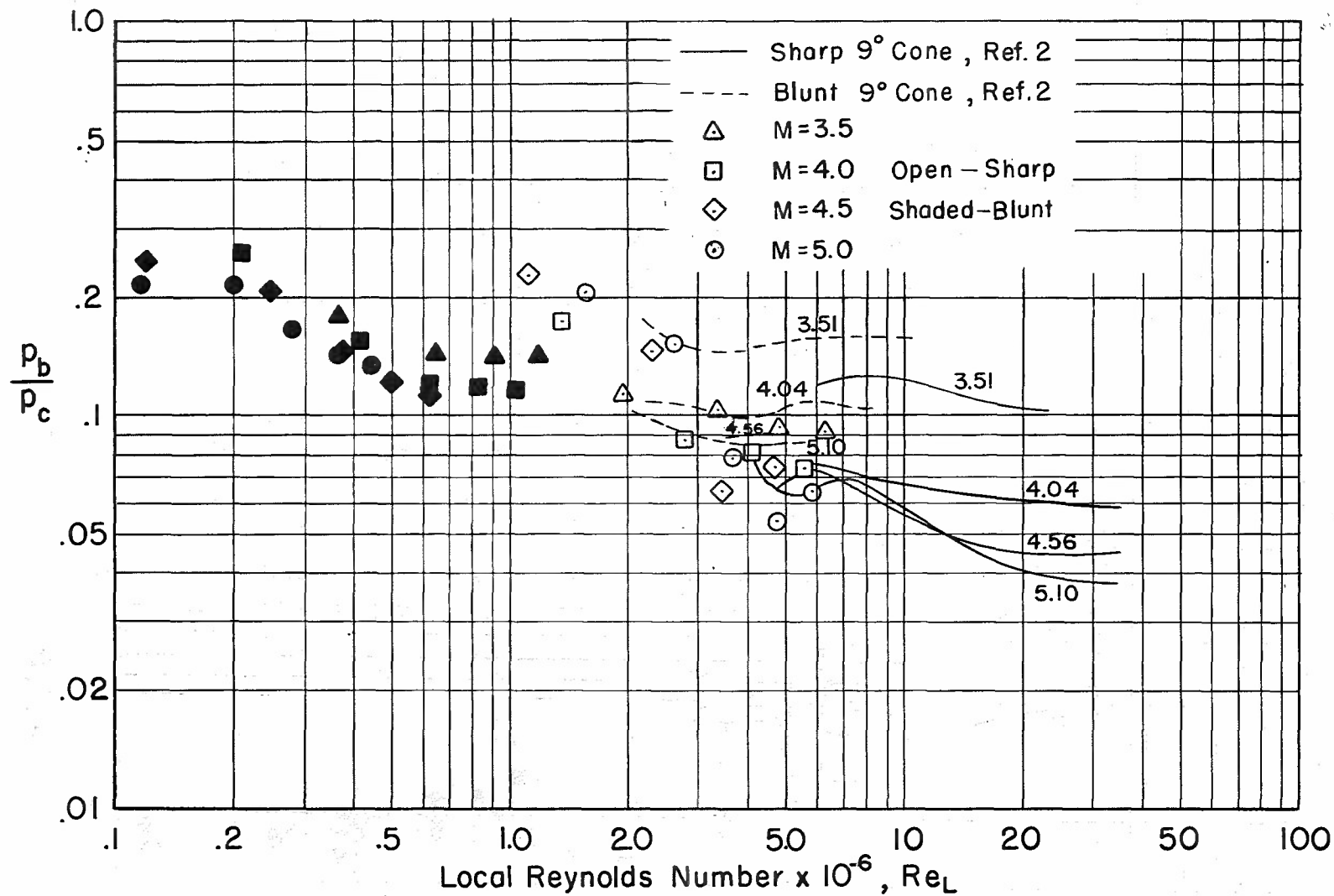
VARIATION OF BASE PRESSURE COEFFICIENT WITH LOCAL REYNOLDS NUMBER FOR RUNS WITH TURBULENT BOUNDARY LAYER ON MODEL

Fig.13



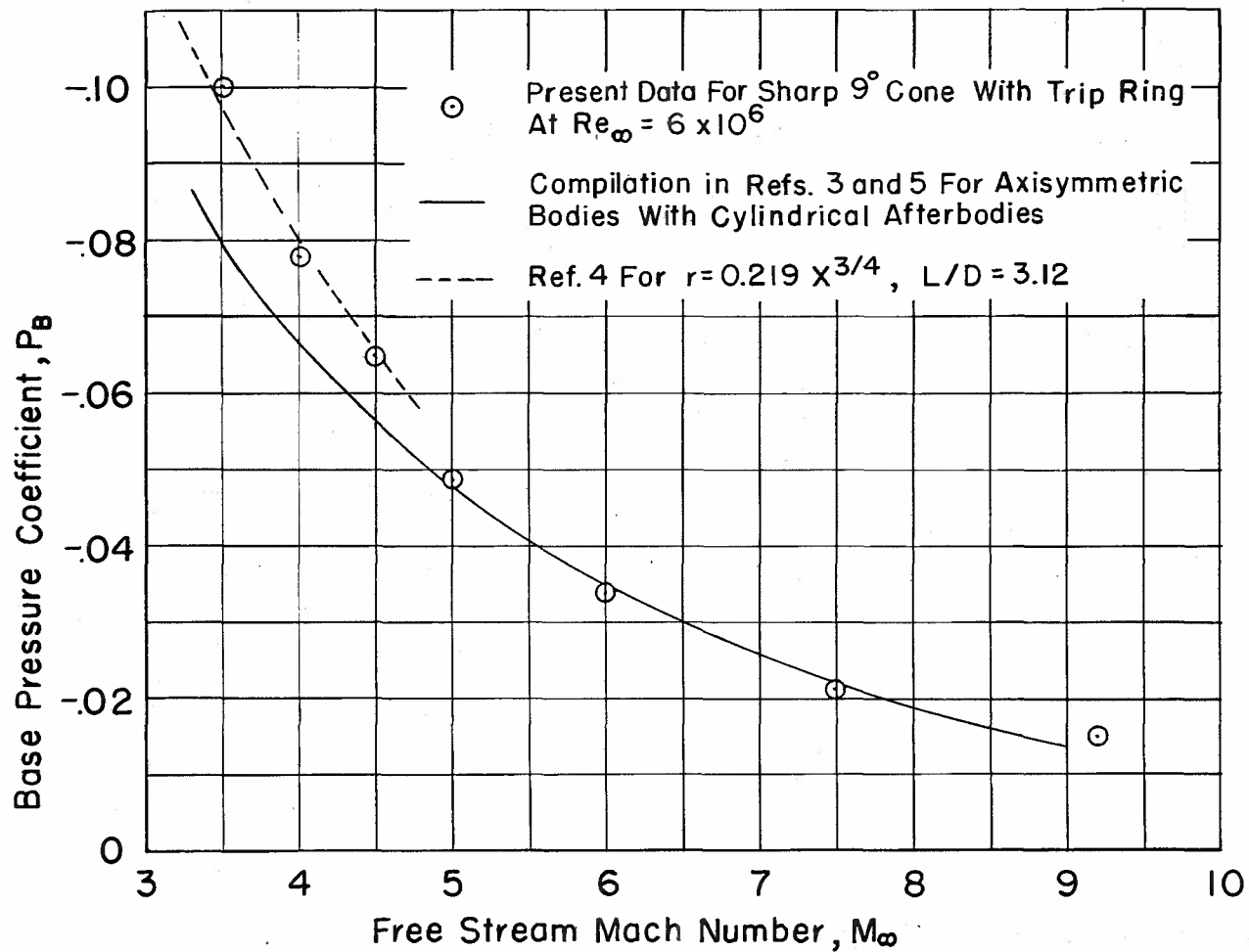
RATIO OF BASE PRESSURE TO CONE PRESSURE VS LOCAL REYNOLDS
NUMBER FOR SHARP AND BLUNT 9° CONES WITHOUT TRIP RING

Fig.14



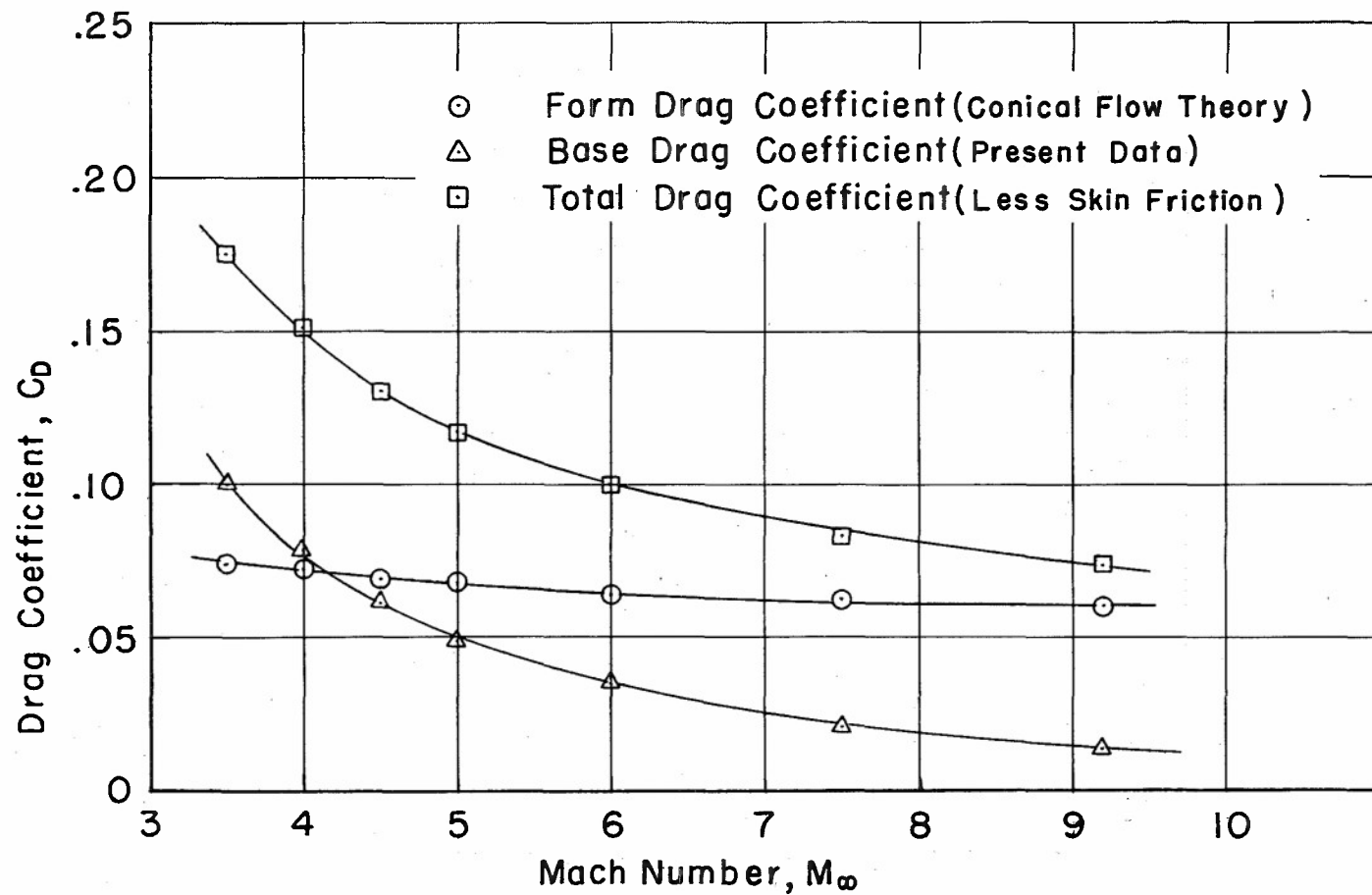
RATIO OF BASE PRESSURE TO CONE PRESSURE VS LOCAL REYNOLDS
NUMBER FOR SHARP AND BLUNT 9° CONES WITH TRIP RING

Fig.15



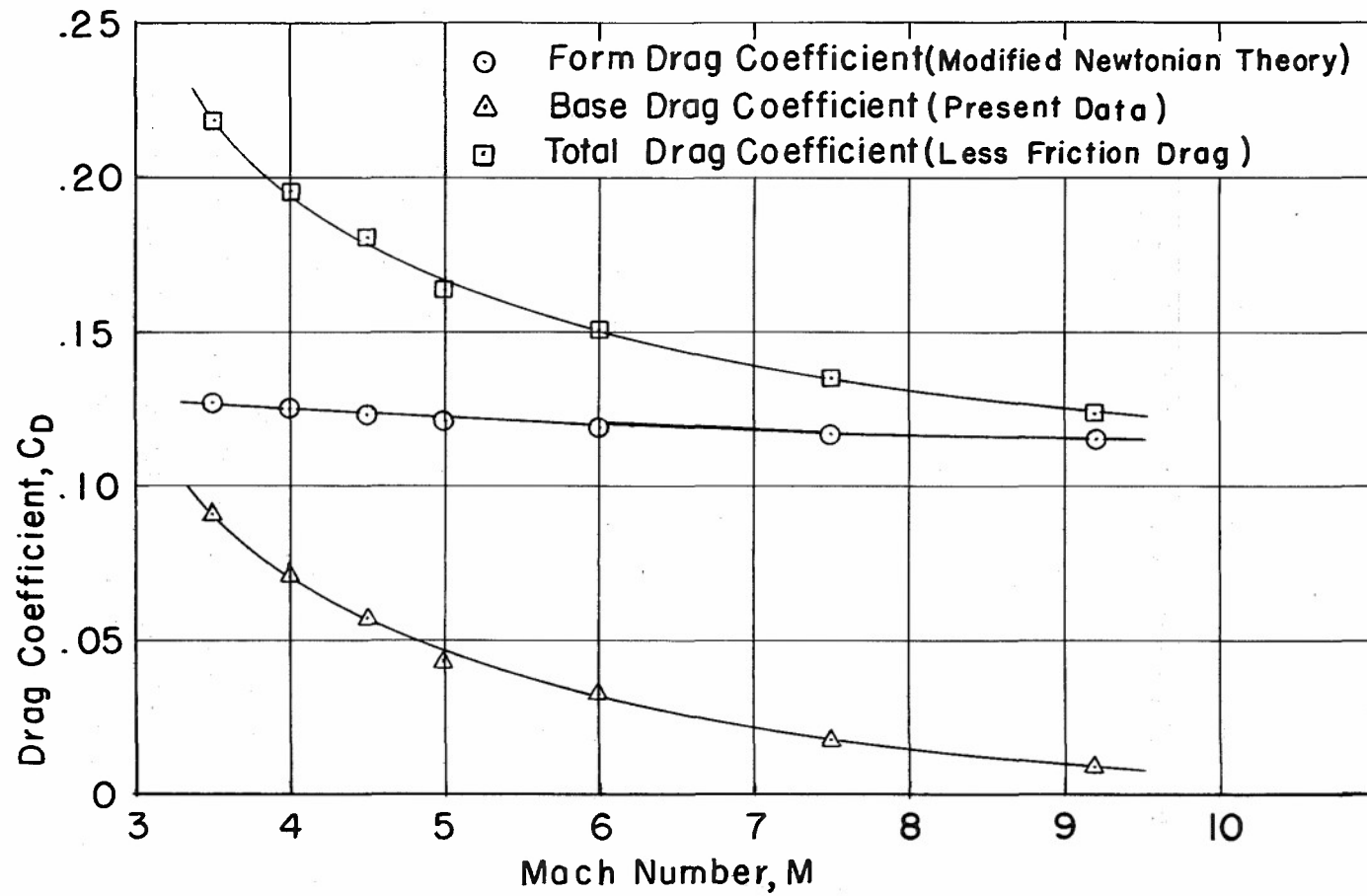
VARIATION OF BASE PRESSURE COEFFICIENT WITH FREE STREAM MACH NUMBER FOR SEVERAL NOSE SHAPES

Fig. 16



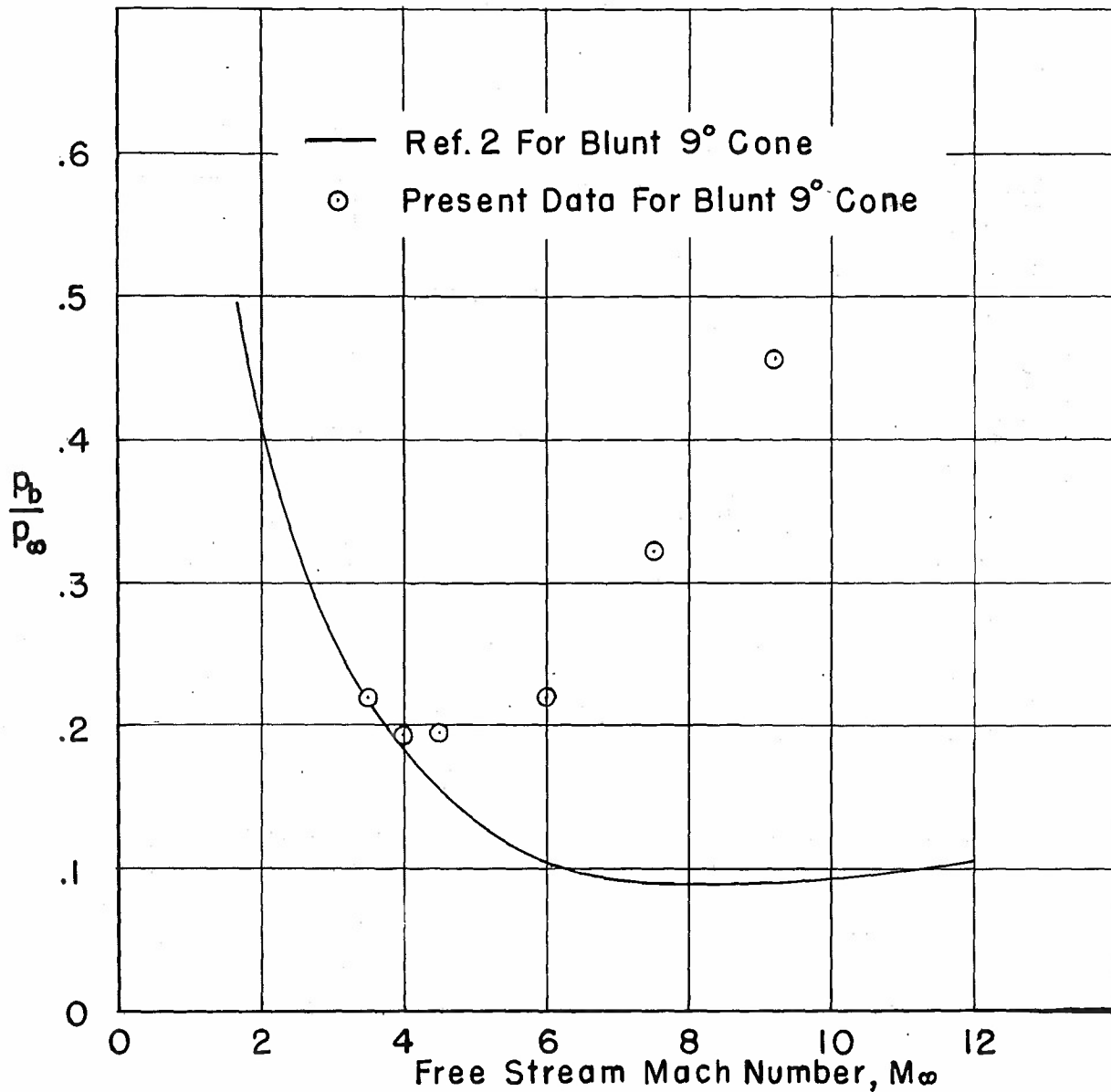
VARIATION OF DRAG COEFFICIENT WITH MACH NUMBER
FOR SHARP 9° CONE

Fig.17



VARIATION OF DRAG COEFFICIENT WITH MACH NUMBER FOR
BLUNT 9° CONE

Fig. 18



RATIO OF BASE PRESSURE TO FREE STREAM STATIC
PRESSURE VS. FREE STREAM MACH NUMBER

Fig. 19

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13. ABSTRACT Base pressure measurements were made on sharp and hemispherically blunted 9° cones at Mach numbers from 3.50 to 9.20. The tests were carried out in the Ballistic Research Laboratories' Supersonic and Hypersonic Wind Tunnels at Aberdeen Proving Ground, Maryland. The data obtained are compared to experimental data and to data from semiempirical analyses from other sources. An empirical correlation for the base pressure data is presented. The relative contributions of base and form drag to total drag are compared.			

14. KEY WORDS	LINK A		LINK B		LINK C	
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