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### NON-LINEAR AXISYMMETRIC POTENTIAL FLOW BOUNDARY MODEL FOR PARTIALLY CAVITATING HIGH SPEED BODIES

TO ALL WHOM IT MAY CONCERN:

BE IT KNOWN THAT (1) ABRAHAM N. VARGHESE and (2) JAMES S. UHLMAN, Citizens of the United States of America, Employees of the United States Government and residents of (1) Wakefield, County of Washington, State of Rhode Island; and (2) Newport, County of Newport, State of Rhode Island have invented certain new and useful improvements entitled as set forth above of which the following is a specification:

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> \*23523\* 23523 PATENT TRADEMARK OFFICE

1 Attorney Docket No. 79866 2 3 NON-LINEAR AXISYMMETRIC POTENTIAL FLOW BOUNDARY MODEL FOR 4 PARTIALLY CAVITATING HIGH SPEED BODIES 5 STATEMENT OF GOVERNMENT INTEREST 6 7 The invention described herein may be manufactured and used by or for the Government of the United States of America for 8 9 governmental purposes without the payment of any royalties 10 thereon or therefore. 11 12 CROSS-REFERENCE TO RELATED PATENT APPLICATIONS 13 Not applicable. 14 15 BACKGROUND OF THE INVENTION 16 (1) Field of the Invention 17 The present invention relates to computer model of 18 hydrodynamic flows and more particularly, relates to modeling 19 partially cavitating flows over a supercavitating axisymmetric 20 body. 21 (2) Description of the Prior Art 22 Modeling of boundary flows about objects subject to laminar 23 and turbulent flows is well known in the art. High speed 24 underwater vehicles, however, cause cavitation of the surrounding 25 fluid. Cavitation reduces pressure in the fluid below its vapor

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pressure causing the fluid to vaporize, allowing the undersea
 vehicle to travel with lower friction when the vehicle is
 completely surrounded by the cavity.

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4 Partial cavitation is an unsteady phenomenon that occurs 5 when part of the supercavitating vehicle is traveling in the cavity. Specifically, this phenomenon occurs during launch of 6 7 the vehicle. A steady, partial cavitation allows development of 8 vehicle designs which take advantage of drag reduction through 9 cavitation. It may also be possible to take advantage of drag 10 reduction with partial cavitation by properly directing the re-11 entrant jet that forms in the cavity closure region. Partial 12 cavitation often occurs during maneuvering of the supercavitating 13 vehicle.

A slender body theory has been developed to solve axisymmetric supercavitating flows. Using the slender body method, sources are defined along the body-cavity axis and control points along the body-cavity surface. A nonlinear differential equation is formed by imposing dynamic boundary conditions on the cavity. A conical cavity closure is assumed in order to solve the developed nonlinear differential equation.

A non-linear boundary element method for determining a cavity shape has been developed. Source and dipole strengths along the body-cavity surface are determined using kinematic boundary conditions on the wetted body surface and dynamic boundary conditions on the assumed cavity shape. The kinematic

boundary condition is then used to update the cavity shape. The
 process is then iterated to solve for the unknown cavity shape.

3 Two numerical hydrodynamics models have been developed by 4 the Naval Undersea Warfare Center for axisymmetric super 5 cavitating high speed bodies. These models are the slender body 6 theory (SBT) model and the boundary element (BE) model. Both of 7 these models have been proven to predict cavity shape and 8 parameters with good accuracy.

9 These models, however, do not account for the transition case when the vehicle is subjected to only partial cavitation. 10 In the SBT model, total drag is predicted by adding the 11 12 pressure drag obtained from the model solution and the viscous drag obtained by applying the Thwaites and Falkner-Skan 13 approximations along the wetted portions of the cavitator. 14 This 15 method is extended to subsonic compressible flows using the 16 compressible Green's function. In the BE model, sources and 17 dipoles are defined on the body-cavity shape and are solved using 18 Green's formula. This yields a Fredholm integral equation of the 19 second kind which gives the supercavitating cavity shape.

Partial cavitation modeling has been done by Uhlman, J.S.
(1987), The Surface Singularity Method Applied to Partially
Cavitating Hydrofoils, Journal of Ship Research, Vol. 31, No. 2,
pp. 107-24; Uhlman, J.S. (1989), The Surface Singularity or
Boundary Integral Method Applied to Supercavitating Hydrofoils,
Journal of Ship Research, Vol. 33, No. 1, pp. 16-20; Kinnas,

S.A., and Fine, N.E. (1990), Non-Linear Analysis of the Flow 1 Around Partially and Super-Cavitating Hydrofoils by a Potential 2 Based Panel Method, Proceedings of the IABEM-90 Symposium, 3 International Association for Boundary Element Methods, Rome, 4 Italy, and Kinnas, S.A., and Fine, N.E. (1993), A Numerical 5 Nonlinear Analysis of the Flow Around Two- and Three-Dimensional 6 Partially Cavitating Hydrofoils, Journal of Fluid Mechanics, Vol. 7 8 254. However, these methods are explicitly adapted for 9 hydrofoils, and the theories presented therein are not readily 10 adapted to supercavitating vehicles.

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#### SUMMARY OF THE INVENTION

One object of the present invention is a method for modeling partial cavitation.

Another object is that such method model partial cavitation about an axisymmetric vehicle.

17 Accordingly, the present invention provides a method for 18 calculating cavity shape for partial cavities about an 19 axisymmetric body having a cavitator located at the foremost end. 20 The method includes receiving system parameter data including 21 geometric data describing the axisymmetric body, a cavity length, 22 and a convergence tolerance. Boundary element panels are 23 distributed along the body-cavity surface and matrices are 24 initialized for each boundary element panel using the unit 25 dipole, unit source functions and known boundary values.

Disturbance potential matrices are formulated for each boundary 1 element panel using disturbance potentials, normal derivatives of 2 disturbance potentials, and no net flux boundary conditions. The 3 initialized matrices and the formulated matrices are solved for 4 each boundary panel to obtain unknown disturbance potentials 5 along the wetted body-cavity surfaces, and normal derivatives of 6 disturbance potentials along the cavity surface. The cavity 7 position is then updated by moving each panel to satisfy the 8 kinematic boundary condition, no flux across the cavity. 9 The method then tests for convergence against a tolerance, and steps 10 11 are iterated until convergence is achieved. The method then 12 provides parameters of interest and the location of the cavity as output. Another aspect of this invention allows the calculation 13 14 of cavity shape and cavity length for an input cavitation number. 15 This is accomplished by an outer loop adjusting cavity length until the model converges to the input cavitation number. 16

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#### BRIEF DESCRIPTION OF THE DRAWINGS

19 These and other features and advantages of the present 20 invention will be better understood in view of the following 21 description of the invention taken together with the drawings 22 wherein:

FIG. 1 is a diagram of a partially cavitating axisymmetric body related to the method of the current invention; and

FIG. 2 is a flow chart of the method of the current
 invention; and

3 FIG. 3 is a flow chart of another method of the current 4 invention.

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DESCRIPTION OF THE PREFERRED EMBODIMENT 6 FIG. 1 shows a diagram of the physical problem of partial 7 cavitation. FIG. 1 shows a radial cross section of an 8 axisymmetric body 10. Axis r represents the radius from the axis 9 of body 10. Axis x represents the length along the body 10 10 measured from a cavitator disk 12. Although a cavitator disk is 11 shown, the model can calculate cavities for cavitator cones as 12 well as cavitator disks. Flow,  $U_{\scriptscriptstyle\!\!\infty},$  is in the direction of arrow 13 14. A cavity 16 is shown extending from the edge of the 14 cavitator along the length of body 10. The length of the cavity, 15  $\ell_{c}$ , is shown by dimension arrows. Likewise, the length of the 16 17 body,  $\ell_{h}$ , is also shown by dimension arrows.

Body 10 extends beyond a cavity closure 18. Cavity 16 is closed to the body 10 with a modified Riabouchinsky cavity termination wall. Cavity closure 18 can be positioned in either body conical section 22 or body cylindrical section 24. The plane of cavity closure 18 is referenced in the following disclosure as an endplate.

Body 10 has a flat front area 20 followed by a conical section 22 and a cylindrical section 24. The diameter of flat

1 front area 20 should be less than or equal to the diameter of the 2 cavitator disk 12 base.

The mathematical formulations in of this algorithm are based 3 on using the cavitator diameter to remove dimensionality for all 4 lengths and using the free stream velocity,  $U_{\scriptscriptstyle\!\! \infty},$  to remove 5 dimensionality for all velocities. Alternate formulations using 6 standard units can also be developed. 7 The flow field is governed by Laplace's equation, 8  $\nabla^2 \Phi = 0$ (1)9 where  $\Phi$  is the total potential which is the sum of free 10 stream potential,  $\phi_{\infty}$ , and disturbance potential,  $\phi$ , giving: 11  $\Phi = \phi_{\infty} + \phi$ (2)12 The free stream potential is the product of the velocity and 13 the distance, x. Because the equation has been non-14 dimensionalized, the velocity is 1, and the free stream 15 potential,  $\varphi_{\scriptscriptstyle\!\!\infty}\,,$  is x. The disturbance potential,  $\phi,$  also obeys 16 Laplace's equation, giving: 17

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$$\nabla^2 \Phi = 0 \tag{3}$$

The disturbance potential satisfies Green's third identity, yielding a Fredholm integral equation of the second kind along the cavitator, cavity, endplate and body. Thus, at any point, x, on the body-cavity surface, the disturbance potential can be computed from:

1 
$$2\pi\phi(x) = \oint_{S} \left[ \phi(x) \frac{\partial}{\partial n} G(x; x') - \frac{\partial}{\partial n} \phi(x) G(x; x') \right] dS$$
(4)

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2 where x' are the points where the sources and dipoles are 3 distributed under the boundary element model;

4 S is the body-cavity surface; and

5 G(x, x') is the Green function.

6 The Green function is further identified as:

7 
$$G(x, x') = \frac{1}{|x - x'|}$$
 (5)

8 The dynamic condition on the cavity boundary is derived from 9 Bernoulli's equation. Along the cavity surface, this can be 10 written as:

11 
$$p_{\infty} + \frac{1}{2}\rho U_{\infty}^{2} = p_{c} + \frac{1}{2}\rho U_{S}^{2}$$
(6)

12 where  $p_{\infty}$  is the free stream ambient pressure;

14  $p_c$  is the pressure inside the cavity; and

15  $U_s$  is the flow velocity at the cavity surface.

16 The flow velocity at the cavity surface can be obtained from

17 equation (6) giving:

18 
$$U_s = \sqrt{1 + \sigma} \tag{7}$$

19 where  $\sigma$  is the cavitation number which is defined as:

20 
$$\sigma = \frac{p_{\infty} - p_c}{\frac{1}{2}\rho U_{\infty}^2}$$
(8)

21 The kinetic boundary condition is that no flow crosses the body-

1 cavity boundary,

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$$\frac{\partial \Phi}{\partial n} = -n_x \tag{9}$$

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3 where  $n_x$  is the axisymmetric body free-stream velocity power. 4 The no net flux condition,

$$\oint_{S} \frac{\partial \phi(x)}{\partial n} dS = 0 \tag{10}$$

6 is also required to make the problem a determinate system.

7 Total drag is calculated by adding the drag coefficients. 8 The pressure drag coefficient,  $C_p$ , at  $\bar{x}$  is calculated as 9 follows:

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$$C_p = 1 - U(\bar{x})^2 \tag{11}$$

11 The pressure contribution to the drag coefficient may then be 12 computed as:

$$C_{dp} = \frac{4}{\pi} \oint_{S} C_{p} n_{x} dS \tag{12}$$

The viscous contribution to the drag coefficient along the wetted portions of the conical and cylindrical body areas is calculated using the International Towing Tank Conference equation given by Newman, <u>Marine Hydrodynamics</u>, MIT Press, Cambridge, Mass. 1980, for the friction coefficient,  $C_f$ , at  $\bar{x}$  is as follows:

19 
$$C_{f} = \frac{0.075}{\left(\log_{10}(R(\bar{x}) - 2)\right)^{2}}$$
(13)

20 where  $R(\overline{x})$  is the local Reynolds number.

21 The total viscous drag coefficient,  $C_{dv}$ , is:

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$$C_{dv} = \frac{4}{\pi} \oint_{S} C_{f} s_{x} dS \tag{14}$$

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2 The base drag coefficient,  $C_{db}$ , which is the component of 3 pressure drag associated with the base of the body is:

 $C_{db} = \frac{0.029(2b_{base})^3}{\sqrt{C_{dv}}}, \qquad (15)$ 

5 where  $b_{\scriptscriptstyle base}$  is the body radius at the base.

6 The total drag coefficient is then given by

$$C_{d} = C_{dv} + C_{dv} + C_{db} . (16)$$

8 The panels are distributed along the cavitator, cavity, endplate, 9 and cylindrical body section aft of the cavity, according to the 10 partial floor method, known in the art. The partial floor method 11 optimizes the number of panels in accordance with requirements 12 for getting good convergence. Non-uniform panel spacing is used 13 in many locations, in order to reduce the number of panels 14 without reducing the accuracy of the solution.

15 During iteration, the end plate height is determined by integrating the cavity surface back from its detachment point on 16 the cavitator, and the number and distribution of panels along 17 the endplate changes according to the changes in the endplate 18 height. Smaller panels are required at highly non-linear flow 19 locations, such as the region near the cavitator. 20 Panel 21 distribution in the wetted body area after cavity closure 18 22 changes to keep the aspect ratio of the neighboring panels

1 between 0.5 and 2.0, in order to ensure good accuracy of the 2 results.

In following the method of the current invention, first an 3 initial cavity is defined. An arbitrary initial cavity can be 4 chosen as a cone extending from the cavitator edge to an assumed 5 endplate height of 0.2 or 0.3 is sufficient for most cases. 6 In this discussion, the endplate height is measured as the radial 7 offset from the body surface to the last point of the cavity. By 8 9 applying equation (4) on all panels along the cavity body 10 surface, S, a system of equations is obtained. This system is 11 solved for the disturbance potentials,  $\phi$ , along the wetted 12 portions of the boundary and on the Riabouchinsky endplate; the 13 normal derivative of the disturbance potential along the cavity 14boundary; and the cavitation number.

15 The kinetic boundary condition given in equation (9) is 16 applied along cavitator, endplate, and aft body to update the 17 cavity shape. In order to update the cavity, the program 18 calculates how much each panel has to be rotated to satisfy the 19 no flow condition. The program starts with the first panel at 20 the cavitator and shifts the aft most point of the panel in the 21 radial direction which satisfies the calculated rotation. The 22 panel is rotated with the aft most point. The foremost point of 23 the next panel is then shifted to the same radius as the previous 24 aft most point. This process is continued until the panel 25 adjacent to the endplate is updated. The endplate height is

adjusted to the aft most point of the aft cavity panel. The
 iteration continues until the kinetic boundary condition
 converges to within a tolerance, giving the cavity shape.
 From the converged disturbance potential along S, the

5 disturbance velocity components can be calculated:

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$$u_x = \frac{\partial \phi}{\partial x}$$
 and  $u_r = \frac{\partial \phi}{\partial r}$ . (17)

Referring now to FIG. 2, there is shown a flowchart of the 7 current invention. In the input step 30, geometric and other 8 9 system parameter data including the estimated cavitation number, 10 the estimated cavity length and the convergence criteria is read. 11 The routine then distributes boundary element panels along the 12 cavitator, cavity, endplate, body extension in the conical 13 section, body extension in the horizontal section, and the aft body. The panels are distributed in order to reduce the number 14 15 of panels and get an accurate result. In the initialize step 32, 16 the algorithm calculates the unit dipole and unit source 17 functions and initializes matrices for the influence functions with known boundary values wherever applicable. The formulate 18 19 equations step 34 formulates matrices for each panel using the 20 disturbance potential equation (4) and no net flux condition 21 given in equation (9). The solve equations step 36 solves the 22 matrices created in the formulate equations step 34 in order to 23 obtain the unknown disturbance potential along wetted body 24 sources, normal distributions of disturbance potentials along 25 cavity surfaces, and the cavitation number. The compute forces

step 38 computes velocity components such as those in equation 1 2 (17) and drag coefficients: including pressure drag, equation (12); viscous drag, equation (14); and base drag, equation (15) 3 from the solved equations. In the update cavity step 40, the 4 cavity is updated from the computed forces using the kinetic 5 boundary condition of equation (9). Convergence on cavity shape 6 7 is checked in the converges decision step 42. If the cavity is 8 not converged, the initialize step 32 is executed to calculate 9 influence functions for the updated cavity and next iteration 10 thus begins. Once the cavity has converged, the compute 11 parameters step 44 computes various output parameters of the 12 converged solution which include pressure drag, viscous drag, 13 base drag, total drag, cavitation number, cavity length, maximum 14 cavity radius, length of cavity to maximum radius location. The 15 output results step 46 then provides the location of the cavity 16 written as coordinates and the cavity's disturbance potential, 17 disturbance potential gradient, and pressure coefficient.

18 The basic algorithm enumerated above provides cavity shape 19 and cavitation number based on an input cavity length. In order 20 to obtain cavity shape and cavity length for an input cavitation 21 number, the embodiment of FIG. 3 adds an additional series of 22 iterations. The user inputs a cavitation number and an assumed 23 cavity length. This embodiment follows the previous embodiment 24 in converging on a new cavitation number,  $\sigma$ , for the assumed 25 cavity length. In step 48, if the new cavitation number is

within a tolerance of the given cavitation number, parameters are 1 computed, step 44, and the results are provided, step 46. 2 Otherwise the embodiment proceeds to step 50 wherein the 3 algorithm determines the relationship between the new cavitation 4 5 number,  $\sigma$ , and the given cavitation number. In step 52, cavity length is increased by a predetermined amount if the calculated 6 7 cavitation number is lower than the initial cavitation number, and in step 54 the cavity length is decreased by a predetermined 8 9 amount if the calculated cavitation number is greater than the 10 initial cavitation number. The routine loops back to the 11 initialize step 32 and recalculates the cavitation number for the 12 new cavity length. Operation continues until the calculated 13 cavitation number falls within a tolerance of the initial 14 cavitation number, the cavity length has converged, as tested in 15 step 48.

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16 Using this invention, partial cavitation for high-speed 17 underwater bodies can be analyzed. As disclosed, the invention 18 can analyze axisymmetrical bodies using two cavitator shapes, a 19 disk and a cone; however, the invention can easily be modified to 20 analyze other axisymmetric cavitator shapes. As disclosed the 21 inventive method can converge on cavity length or cavitation 22 number. Total drag is calculated by adding the pressure drag, 23 viscous drag and base drag. The invention can also be utilized 24 for studying the effects of body aft radius, body cone angle and 25 body cone angle starting at the cavity closure if the closure is

1 on conical section 22. This method provides new information 2 concerning the physics of cavitation which can be used in the 3 design of cavitating vehicles.

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## NON-LINEAR AXISYMMETRIC POTENTIAL FLOW BOUNDARY MODEL FOR PARTIALLY CAVITATING HIGH SPEED BODIES

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#### ABSTRACT OF THE DISCLOSURE

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A method for calculating parameters about an axisymmetric 7 body in a cavity is provided. The user provides data describing 8 the body, a cavity estimate, and convergence tolerances. 9 Boundary element panels are distributed along the body and the 10 estimated cavity. Matrices are initialized for each panel using 11 12 disturbance potentials and boundary values. Disturbance 13 potential matrices are formulated for each panel using 14 disturbance potential equations and boundary conditions. The 15 initialized matrices and the formulated matrices are solved for each boundary panel to obtain panel sources, dipoles and 16 17 cavitation numbers. Forces and velocities are computed giving 18 velocity and drag components. The cavity shape is updated by 19 moving each panel in accordance with the calculated values. The 20 method then tests for convergence against a tolerance, and 21 iterates until convergence is achieved. Upon completion, 22 parameters of interest and the cavity shape are provided. This 23 invention also allows determiniation of cavity shape for a 24 cavitation number.



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FIG. 3