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Technical Report 680

COHERENCE VARIABILITY OF ARRAYS DURING BEARING STAKE (U)

JA Neubert

August 1981


Prepared for Naval Electronic Systems Command

*Unauthorized Disclosure Subject to Criminal Sanctions*
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   (C) By the use of acoustic synoptic factors, this study investigates how the variable and random sound field structure influences the performance of wide-aperture, horizontal line array systems for low-frequency, narrowband sources in the northwestern Indian Ocean using data from the Bearing Stake exercise. Several acoustic synoptic factors are defined and illustrated. The patterned and coherence behavior of arrays is studied in detail. The degrading effects of an array turn are considered graphically and synoptically. The interaction of several of the acoustic synoptic factors is shown. The stability of time averages is investigated synoptically. It was found that the more stable the time average of the phase is, the better the coherence is. On the other hand, the stability of the time average of the amplitude had little impact on the coherence.

Finally, it was shown that a normal mode model was suitable for representing the Bearing Stake results.
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<td>$A_j$</td>
<td>instantaneous amplitude measured at sensor $j$</td>
</tr>
<tr>
<td>$\beta_j$</td>
<td>instantaneous residual phase of Eq (2)</td>
</tr>
<tr>
<td>$\Sigma_A$</td>
<td>normalized amplitude standard deviation of Eq (4)</td>
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<tr>
<td>$S+N$</td>
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<td>$N$</td>
<td>noise power of Eq (10)</td>
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<td>$C_P$</td>
<td>array phase coherence coefficient of Eq (12)</td>
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<tr>
<td>$\text{ASG}$</td>
<td>array signal gain of Eq (13)</td>
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<tr>
<td>$CA_j$</td>
<td>amplitude nonhomogeneity factor of Eq (15)</td>
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<td>$C^j_i$</td>
<td>percent relative amplitude standard deviation of Eq (16)</td>
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<td>array phase coherence matrix of Eq (21)</td>
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<td>$(\gamma^c_{jk})^2$</td>
<td>classical coherence matrix of Eq (22)</td>
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<tr>
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<td>plot notation for $\text{CNA}_j\phi$; see Eq (23)</td>
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<tr>
<td>$\text{CF}$</td>
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<td>structure factor of Eq (33)</td>
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<td>plot notation for $\text{CM}_j\phi$ of Eq (37)</td>
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<td>$\text{NAC}$</td>
<td>amount of negative amplitude correlation as defined by Eq (39)</td>
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<tr>
<td>$\text{ADS}_i$</td>
<td>plot notation of $A\Delta_i$ of Eq (43)</td>
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I. (U) INTRODUCTION

(C) This report investigates (via acoustic synoptic factors) how the variable and random sound field structure influences the performance of wide-aperture, horizontal line array systems for low-frequency, narrowband sources in the northwestern Indian Ocean using data from the Bearing Stake exercise.

(U) When operating an array system in the ocean, a complex situation results. There are environmental and system variables that result in a variable system behavior. Environmental variables are oceanic conditions (e.g., multipath), geographic location, and measurement conditions (e.g., array depth and surveillance range). Examples of system variables are averaging time and processing algorithms. The variable system behavior manifests itself by its impact on array gain, sidelobe degradation, etc.

(U) Obviously these variables interact. A convenient and representative summary is given below.

A. Environmental variables:
   1. Oceanic conditions, e.g., propagation loss, bathymetry, bottom interaction and losses, sea state, sound-speed channels, multipaths, internal waves, diffraction, dispersion, scattering, structure of the sound-speed profile, fronts and eddies, noise, fluctuations
   2. Geographic location
   3. Season
   4. Array deformation
   5. Measurement conditions, e.g., source level and tow speed, frequency, bandwidth, surveillance range, target bearing and aspect, array depth and tilt, array aperture length and configuration, type of target signal, track of maneuvering submarine, doppler.

B. System variables:
   1. Averaging time and method
   2. Bandwidth
   3. Beamformer weighting (array shading)
   4. Degree of signal clipping
   5. Processing algorithms and criteria
   6. Array operation procedure (e.g., tow speed and turns).

C. Variable system behavior can be characterized by
   1. Array gain
   2. Coherence
   3. Sidelobe degradation
   4. Quality of the main beam (e.g., effective beamwidth, beam splitting, beam jumping) (Ref 1)
   5. Quality of the CMAP passive ambiguity surface (Ref 1)
   6. Target detectability
   7. False alarms

(U) To permit consideration of as many of these factors as feasible, extensive experiments such as Bearing Stake (discussed below) have accumulated massive data bases. In order to use these large quantities of data to systematically appraise system performance, synoptic factors have been computed and then tabulated and/or plotted versus the variables of interest. For example, the array signal gain behavior of an array system can be summarized for each data set by the coefficient ASG of Eq (13). Then ASG can be tabulated versus, say, range and integration time. Examining these tabulations gives a quantitative appraisal of such things as the importance of array deformation and the limitations of integration time and range. These tabulations can be further simplified by taking the mean and standard deviation of ASG. These statistical values can then be compared as a function of, say, array depth and tilt, target bearing, geographic location, and season to analyze array system behavior and its models.

(U) What is a synoptic factor? It is not an abstract concept, but is a computed synoptic quantity determined by the actual data themselves. Hence synoptic factors can consider the complex effects of the environmental and systems variables on the behavior of specific systems. They use actual ocean data to concisely characterize the pertinent system behavior. Each synoptic factor is specifically classified by the conditions that occur during the ocean measurement and by the method of system application. Synoptic factors quantify the impact of complex environmental and system variables on the amplitude and phase and (and their interaction) across the array and relate these to the behavior of array systems. Actually it is easier to think in terms of two types of synoptic factors: acoustic synoptic factors and system synoptic factors. Acoustic synoptic factors are coefficients that summarize quantitatively the acoustic behavior across arrays as a function of the environmental variables. System synoptic factors are coefficients that summarize quantitatively the impact of the environmental and system variables on array system behavior. This report is primarily interested in the application of acoustic synoptic factors to array systems. However, $C_{ai}$ of Ref 1 is an example of an effective system synoptic factor for some cases (see Fig 6 and 8 of Ref 1).

(U) By themselves synoptic factors cannot predict array system behavior if the underlying pertinent conditions change. They merely synopsize this behavior in a useful form that is sensitive to the changes in the underlying pertinent conditions. Therefore, they are useful for quantitatively comparing changing oceanic situations and often for discerning which changing conditions affect the system behavior and by how much. Examples of this are found in Sections III and IV. However, synoptic factors can be used with deterministic theoretical models such as normal modes (Ref 2) and Raywave (Ref 3). Deterministic models can consider or approximate some of the wide range of interacting oceanic variables and then be input into the synoptic factor algorithms. The theoretical results produced thereby can then be compared with the synoptic factor algorithms using real oceanic data inputs. The results can be used to appraise the suitability of the theoretical models for the situation under consideration. By iterative use of this approach, the theoretical models can be improved and the pertinent oceanic conditions can be better understood (Ref 4). In short, synoptic factors can succinctly evaluate the accuracy and variability of deterministic theoretical approaches as a function of environmental and system variables in order to suggest improvements in these theoretical models.

2. M. A. Pedersen and G. S. Yee, "Propagation Loss Assessment of the Bearing Stake Exercise (U)," NOSC TR 467, 28 September 1979, CONF.
4. D. F. Gordon, "Array Simulation at the Bearing Stake Sites (U)," NOSC TR 664, April 1981, CONF.
(U) The synoptic factor technique discussed below is universal (i.e., it is not site specific), while the data from any site may be unique. Once developed for a sufficiently complex region (like the northwestern Indian Ocean), the synoptic factors can be applied with equal validity to other ocean regions (such as the Mediterranean Sea, the Norwegian Sea and the Sea of Japan) as the data are made available. By comparing archival data and using appropriate models, projections of array systems behavior to other regions, where data are not yet available, can be made from regions for which data presently exist. This makes synoptic factors useful for planning experiments and for projecting present and future array system performance to other regions.

(C) Figure 1 shows the northwestern Indian Ocean surveyed during Bearing Stake. It has been shown (Ref 5 through 10) that the sound field was highly structured and variable in this strategically important ocean region and that the sound field structure changed from site to site. For wide-aperture, horizontal line arrays in this multipath environment, the sound field structure can be described by

1. the amplitude nonhomogeneity across the array,
2. the wavefront corrugation across the array,
3. the fluctuations,
and
4. the interaction of these as discussed below.

(U) Three arrays were employed for several projector tows for five sites shown in Fig 1 during Bearing Stake. These three arrays were:

1. Ocean Acoustic Measurement System (OAMS) array,
2. Long Acoustic Towed Array (LATA) (formally called Lambda I), and

(U) In Section II several acoustic synoptic factors are given and explained. Section III extends the sound field structure study of Ref 7 by introducing several new acoustic synoptic factors. It also illustrates the impact of an array turn on the performance of an array. It concludes by showing some interactions between acoustic synoptic parameters. Section IV discusses the stability of time averages for the purposes of appraising array performance and introduces an appropriate class of acoustic synoptic factors. Section V discusses the merits of a normal mode array performance model using Bearing Stake data.

(U) In this report the first 11 signal array arrivals for projector tow 2P3A and the first 12 signal array arrivals for projector 4P1 are considered as representative of the Bearing Stake data. These data are shown in Fig 2 through 24 and pertain to the relations in Sections II and IV. Acoustic synoptic factors from Sections II through IV for all of projectors tows 2P3A and 4P1 are found in Tables 1 and 2.

II. (U) ACOUSTIC SYNOPTIC FACTORS

(U) Represent a narrowband signal arriving at the \( j \)th sensor along a wide-aperture, horizontal line array at time \( t \) by

\[
F_j = A_j \exp \left[ i \left( \phi_j - \omega t \right) \right]
\]  

(1)

where \( A_j \) is the amplitude, \( \phi_j \) is the phase and \( \omega \) is the angular frequency. Given the target bearing \( \Phi_0 \) (defined as zero for a forward endfire arrival), define the residual phase as

\[
\theta_j = \phi_j - k_0 d_j \cos \Phi_0,
\]  

(2)

where \( k_0 \) is the wavenumber and \( d_j \) is the distance from the first sensor to the \( j \)th sensor \((d_1 = 0)\). The amplitude fluctuations are given by

\[
\delta A_j = A_j - \langle A_j \rangle.
\]  

(3)

The total number of array sensors is denoted by \( J \).

(U) For an array of \( J \) sensors, \( \theta_j = \theta_0 \) for all \( j \neq \ell \) represents a plane arrival for a linear horizontal array. Likewise, \( A_j = A_0 \) for all \( j \neq \ell \) represents a homogeneous amplitude arrival along the array. Wavefront corrugation occurs when \( \theta_j \neq \theta_0 \) for any \( j \neq \ell \), i.e., it represents a nonlinear phase variation along the array. Amplitude nonhomogeneity (i.e., amplitude variation along the array) occurs when \( A_j \neq A_0 \) for any \( j \neq \ell \). Amplitude nonhomogeneity and wavefront corrugation represent the sound field structure as experienced by the array in a multipath ocean.

(U) In order to quantify, analyze, and model the sound field structure across the array, the following acoustic synoptic factors were devised. Two classes of acoustic synoptic factors will be considered below: total array synoptic factors and intra-array synoptic factors.

TOTAL ARRAY SYNOPTIC FACTORS (U)

\( \Sigma_A \) Factor (U)

(U) The signal plus noise amplitude fluctuation behavior is given by the normalized standard deviation

\[
\Sigma_A \equiv \sigma_A / \bar{A}_a > 0.
\]  

(4)

where

\[
\bar{A}_a = \frac{1}{J} \sum_{j=1}^{J} \langle A_j \rangle.
\]  

(5)

\[
\sigma_A^2 = \frac{1}{J} \sum_{j=1}^{J} \sigma_j^2.
\]  

(6)
and

\[ \sigma_j^2 \equiv \langle (A_j - \langle A_j \rangle)^2 \rangle = \langle (\delta A_j)^2 \rangle \]  

(7)

with the operator \( \langle \cdot \rangle \) representing a time average. Normalizing by \( \bar{A}_d \) avoids changes in the propagation conditions from biasing the measure \( \Sigma_A \).

\( \Sigma_A \) has been tabulated on a site-by-site basis and gives the following overview of the Bearing Stake exercise (Ref 9, 10):

1. \( \Sigma_A \) is sensitive to the propagation conditions.
   a. It discerned the bottom-limited propagation for Sites 3 and 1 from the non-bottom-limited propagation for Site 4.
   b. It sensed the propagation interference due to the presence of a seamount, i.e., it discerned the difference between projector tows 5P1 (no seamount) and 2P3A (seamount between source and OAMS array).

2. \( \Sigma_A \) is sensitive to array behavior (e.g., array tilt and/or tow depth).
   a. It discerned that the LATA (towed at 300 ft) was not sensitive to the propagation condition differences between Site 4 (not bottom-limited) and Sites 5 and 2 (bottom-limited) while the OAMS array (towed at 200 ft) was.
   b. It showed that the BMA and the OAMS array performed essentially the same at Sites 3 and 1B. [This was confirmed by the \( C_p \) and \( A_b \)G analyses (Ref 9, 10).]

3. \( \Sigma_A \) is sensitive to array sensor calibration and/or damage [e.g., it discerned the OAMS array "crush event" during Site 4 projector tows (Ref 5)].

SNR Factor (U)

\( \text{SNR} \equiv (S + N) - N \) \quad (8)

where

\[ (S + N) \equiv 10 \log_{10} \bar{A}^2 \quad \text{(for signal plus noise)} \]  

(9)

\[ N \equiv 10 \log_{10} \bar{A}^2 \quad \text{(for noise only)} \]  

(10)

and

\[ \bar{A}^2 \equiv \frac{1}{J} \sum_{j=1}^{J} \langle A_j^2 \rangle \]  

(11)

\( C_p \) Factor (U)

\( \text{The wavefront coherence is measured by the array phase coherence coefficient} \)
\[
C_p = \frac{1}{J^2} \sum_{j=1}^{J} \sum_{k=1}^{J} \langle \cos (\theta_j - \theta_k) \rangle \leq 1 ,
\]

where \( \langle \cos (\theta_j - \theta_k) \rangle \) measures the phase incoherence due to wavefront corrugation and array deformation.

\( \text{(U)} \) Plots of \( C_p \) versus range \( R \) measure the phase incoherence due to wavefront corrugation and array deformation, i.e., high \( C_p \) values occurring intermittently at all ranges show that the array deformation is effectively negligible for a radial projector tow. Then the existence of intermittent low values of \( C_p \) rangewise shows the occurrence of beam jumping during this projector tow due to variable wavefront corrugation. \( C_p \) is also used to evaluate the array modeling ability and limitations of normal mode theory (see Section V).

\text{ASG Factor (U)}

\( \text{(U)} \) By conveniently normalizing (to compensate for the use of various array shading schemes) the beam array signal gain [see Eq (26) of Ref 1], the array signal gain can be given by (Ref 8)

\[
\text{ASG} = 10 \log_{10} \text{asg} \leq 0 ,
\]

where

\[
\text{asg} = \frac{\sum_{j=1}^{J} \sum_{k=1}^{J} W_j W_k \langle A_j A_k \cos (\theta_j - \theta_k) \rangle}{\sum_{j=1}^{J} \sum_{k=1}^{J} W_j W_k A_j A_k}
\]

and the weighting factors are \( 0 \leq W_j, W_k \leq 1 \). The mean and standard deviation of \( \text{ASG} \) are usually of the most interest.

\text{INTRA-ARRAY SYNOPSIS FACTORS (U)}

\( \text{(U)} \) Two types of intra-array synoptic factors will be considered below: one-point structure factors and two-point structure factors.

\text{One-Point Structure Factors Versus Sensor Number } j \ (U)

\( \text{(U)} \) \( CA_j \) \text{ FACTOR. The amplitude nonhomogeneity factor}

\[
CA_j = 10 \log_{10} (A_j)
\]

measures the amplitude nonhomogeneity that impacts the degree of beam splitting and sidelobe suppression (Ref 1). In situ \( CA_j \) can check for sensor calibration and damage as was shown in Ref 7.
(U) \( C_j^\sigma \) FACTOR. The percent relative amplitude standard deviation

\[ C_j^\sigma \equiv 100 \frac{\sigma_j}{\langle A_j \rangle} \]  

(16)

is very sensitive to signal multipaths but not to noise multipaths and amplitude nonhomogeneity. See Ref 7.

Two-Point Structure Factors Versus Sensor Separation (U)

(U) \( C_{NJ} \) FACTOR. The normalized amplitude covariance

\[ C_{NJ} \equiv \frac{\langle A_j A_k \rangle}{\max (\langle A_j \rangle^2, \langle A_k \rangle^2)} \]  

(20)

gives the structure of the amplitude field across an array in a multipath ocean. Note that \( C_{NJ} \) varies about its value for zero covariance, which is \( 1/\max (\langle A_j A_k \rangle/\langle A_j \rangle \langle A_k \rangle) \). This zero covariance level is indicated in the \( C_{NJ} \) plots of Fig 2 through 24 by the solid horizontal line with arrows on either end.

(U) \( C_{JR}^p \) FACTOR. The array phase coherence matrix

\[ C_{JR}^p \equiv \langle \cos (\theta_j - \theta_k) \rangle \]  

(21)

gives the structure of the nonlinear phase field across an array in a multipath ocean.

(U) \( (\gamma_{JR}^c)^2 \) FACTOR. The classical coherence matrix

\[ (\gamma_{JR}^c)^2 \equiv \frac{(A_j A_k \cos (\theta_j - \theta_k))^2 + (A_j A_k \sin (\theta_j - \theta_k))^2}{\langle A_j^2 \rangle \langle A_k^2 \rangle} \]  

(22)

gives the interaction of the structure of the amplitude field and the nonlinear residual phase field across an array in a multipath ocean.
III. (U) SOUND FIELD STRUCTURE STUDY

(U) The plots of CNAj, Cj and (yj) (see Fig 2 through 24) show the nature of sound field structure at Sites 2 and 4 for Bearing Stake. These figures are somewhat difficult to follow and understand. In this section we introduce some additional synoptic factors that clarify the interpretation of these figures; they are included with each figure. Table 3 then aids the reader in prioritizing the figures in order of their coherence via the coherence factor CF of Eq(32). In addition, Table 3 shows that the impact of off-broadside incidence \Delta \Phi_o of the received signal has no real influence on the coherence. Also it is seen that Cj is a poor prioritizing factor compared to CF. Table 3 also shows that the phase time stability, as synopsized by ADS3 of Eq(43), is closely related to the quality of coherence, while the amplitude time stability, as synopsized by ADS1 of Eq(43), is not. On page 12 a narrative guide through Fig 2 through 24 is given in terms of the coherence structure. As a further aid to understanding Fig 2 through 24, Eq (34) through (36) relate some of the pertinent stochastic factors that appear on the figures.

(U) The sound field structure of Fig 2 through 24 is synopsized as follows. Cj is plotted as a function of sensor separation d = |d| - |d|; that is, Cj = Cj, i = 1, 2, 3, where

\[ C_1 = C_1 = C_{11} = C\ldots, \]
\[ C_2 = C_2 = C_{22} = C\ldots, \]
\[ C_3 = C_3 = C(y)^2 = C\ldots. \]

At each separation d there exists a number n of values C(k), 1 \leq k \leq n. That is, n is a function of d. There is a total N sensor separations. The first separation d = 0 and the maximum separation d is D. With each separation d there is an associated set of n values C(k). This can be represented by relating the vector

(d) = (d_1, ..., d_N = D)

to the vector

(n_d) = (n_1, ..., n_D = 1).

(U) At each value of d, compute

\[ \mu_d = \frac{1}{n_d} \sum_{k=1}^{n_d} C_d(k) \]
\( (\sigma_d^i)^2 \equiv \frac{1}{n_d} \sum_{k=1}^{n_d} \left[ C_d^i(k) - \mu_d^i \right]^2 \geq 0, \quad \sigma_d^i = SD_i. \) \hspace{1cm} (27)

\( \mu_d^i \) and \( \sigma_d^i \) represent the two essential properties of the behavior of the sound field structure as manifested in Fig 2 through 24. The mean structure \( \mu_d^i \) represents the structured behavior of \( C_d^i \) versus \( d \). The standard deviation \( \sigma_d^i \) represents the vertical spread about the mean structure \( \mu_d^i \) versus \( d \). Both \( \mu_d^i \) and \( \sigma_d^i \) come from normalized functions, and it has been found that the behavior of \( \mu_d^i \) and \( \sigma_d^i \) versus \( d \) have largely independent structure. The appropriate acoustic synoptic factors are

\[ \Sigma_g^i \equiv \frac{1}{N} \sum_{d=d_1}^{N} \sigma_d^i = SND_i, \] \hspace{1cm} (28)

\[ C_j^i \equiv \frac{1}{J} \sum_{j=1}^{J} \sum_{k=1}^{J} C_{j_k}^i = CI_i, \] \hspace{1cm} (29)

where \( C_j^2 = C_\psi \) of Eq (12),

\[ \mu_N^i \equiv \frac{1}{N} \sum_{d=d_1}^{N} \mu_d^i = MND_i, \] \hspace{1cm} (30)

and

\[ (\mu_g^i)^2 \equiv \frac{1}{N} \sum_{d=d_1}^{N} (\mu_d^i - \mu_N^i)^2 \geq 0, \quad \mu_g^i = MS_i. \] \hspace{1cm} (31)

Note that in general \( \mu_d^i \) will not equal \( C_j^i \). \( \mu_d^i \) is of special interest since it becomes smaller as the vertical variation of \( \mu_d^i \) versus \( d \) decreases in vertical extent. \( \mu_d^i \) is drawn as a solid curve through the plots of \( CD_i \) in Fig 2 through 24; \( CI_i, MND_i, SND_i, \) and \( MS_i \) are also given on each of these plots.

(U) Several characteristics of the sound field structure are evident in Fig 2 through 24. The quality of the classical coherence \( \left( \gamma_{12}^i \right)^2 \) is better the higher \( \mu_d^i \) remains and the smaller \( \sigma_d^i \) is. Thence, the Coherence Factor is defined by the convenient measure

\[ CF \equiv 1 - \frac{\Sigma_g^3}{C_j^2} \ll 1. \] \hspace{1cm} (32)
The closer this factor is to unity, the higher the quality of the coherence observed. In terms of CF, Table 1 shows the plots of $(\gamma_{jk}^C)^2$ rated in order of decreasing coherence. The amount of structure in $(\gamma_{12}^C)^2$ is determined by $\mu_3^3$ and $\sigma_3^3$. To synopsize this, define the Structure Factor by

$$SF = \sum_{\sigma} \mu_3^3.$$  

(33)

The smaller SF is, the less structure that is occurring in $(\gamma_{jk}^C)^2$. SF is shown in the plots (Fig 2 through 24) of $(\gamma_{jk}^C)^2$, where it is seen that SF correlates well with CF. However, visual inspection of the plots of $(\gamma_{jk}^C)^2$ shows that CF is the better coherence prioritizing factor. Table 3 shows that CF is also a far better prioritizing factor than $C_p$.

(U) The amplitude, phase, and coherence behavior for the first 12 signal periods for projector tows 2P3A and 4P1 is shown in Fig 2 through 24. The Tow 4P1 behavior is discussed first. Tow 4P1 has one turn at about time 1111 ZULU. Comparing Fig 16c (at the time of the turn) with Fig 14c (24 min before the turn) shows that during the turn, the coherence fluctuation increases greatly and the coherence becomes structured. This is also evident in the relative amplitude behavior (see Fig 14a and 16a) and in the relative phase behavior (cf. Fig 14b and 16b). In Fig 17 (7 min after the turn) the patterned amplitude and phase behavior persists but is not as well correlated. In Fig 18 (20 min after the turn) the coherence fluctuation has decreased while it has become much more patterned, and the amplitude and phase are patterned and in correlation. In Fig 19c (36 min after the turn) and in Fig 20c (50 min after the turn), the coherence shows its greatest pattern magnitudes. In Fig 21c (1 hr 45 min after the turn), the coherence pattern magnitude has decreased significantly (but shows more cycles of pattern). In Fig 22c (1 hr, 50 min after the turn) the coherence pattern magnitude has decreased greatly and again shows fewer cycles of pattern.

(U) Tow 2P3A has no turns. Figure 2c shows only fair coherence, but Fig 3c (24 min later) shows excellent coherence, and Fig 4c (5 min later) shows good coherence. Seventeen minutes later (Fig 5c) the coherence is again only fair, but the coherence is again excellent seventeen minutes later in Fig 6c. The coherence is good 29 min later (Fig 7c) and remains good in Fig 8c (9 min later), and in Fig 9c (45 min later). Sixty-one minutes later (Fig 10c) the coherence declines somewhat but remains good. However, in Fig 11c (16 min later) the coherence is only moderate and continues to be moderate in Fig 12c (11 min later). Thus, the coherence varied considerably during the straight-line tow 2P3A. (Table 3 also indicates this behavior.)

(U) Some observations can be made about Fig 2 through 24 using the synoptic factors that appear above each figure. In general, the mean vertical structure of CNA_{jk} times the mean structure of $C_{jk}^P$ nearly equals the mean structure of $(\gamma_{jk}^C)^2$, i.e.,

$$\frac{\mu_d^1 \cdot \mu_d^2}{\mu_d^3} \approx 1.$$  

(34)

(The only case in which this does not occur closely is for Fig 12, where the value is 0.764.) This relation between the mean structure is to be expected because of the definitions of CNA_{jk}, $C_{jk}^P$ and $(\gamma_{jk}^C)^2$ in Eq (20, 21) and (22), respectively.
The quantity \( (\gamma^2) \) is a measure of the correlation of the amplitude and phase structure. The smaller \( \Sigma^2 \) is relative to the sum of \( \Sigma^1 \) and \( \Sigma^2 \), the better the phase-amplitude correlation. This is measured by

\[
\Sigma_r \equiv \frac{\Sigma^1 + \Sigma^2}{\Sigma^3} .
\]

In general, \( \Sigma_r \) is greater than unity (the only exceptions are Fig 24, 21 and 17, where \( \Sigma_r \) equals 0.744, 0.9744 and 0.875, respectively). In general, \( \Sigma_r \) decreases as CF decreases. The relation of \( \Sigma^2 \) to \( \Sigma^1 \) is also of interest. In all cases, the phase coherence fluctuations exceed the amplitude fluctuations, ie,

\[
\frac{\Sigma^2}{\Sigma^1} > 1 ,
\]

and this ratio roughly declines as CF declines.

For signals (out less so for noise), \( \sigma^1 \) generally declines with d as the number of sample sets decreases. This effect is less pronounced when the structure of \( C_{ij}^P \) is more patterned, but it is very definite when the patterned structure of \( C_{ij}^P \) abates. For signals (but less so for noise), \( \sigma^2 \) tends to peak in mid-range for d. This represents a compensating trend between the occurrence of many samples for small separations d and reduced phase coherence for greater separations d.

The behavior of

\[
CM_{jk} \equiv \frac{\langle A_j A_p \rangle}{\langle A_j \rangle \langle A_p \rangle} = CM_{jk}
\]

plotted in Fig 2 through 24 is interesting. It correlates with the observation that the higher CF is and the less structured \( C_{ij}^Q \) is, the more likely \( C_{ij}^Q \) is to remain above its zero correlation line given by \( 1/\max(\langle A_j A_p \rangle/\langle A_j \rangle \langle A_p \rangle) \). Therefore, the relation of \( C_{ij}^Q \) to its zero correlation line and the amount of \( C_{ij}^Q \) with negative correlation will be determined as follows. \( C_{ij}^Q \) has a negative correlation when \( \langle A_j A_p \rangle/\langle A_j \rangle \langle A_p \rangle < 1 \). Therefore, define

\[
CA^3_d \equiv \frac{\langle A_j A_p \rangle}{\langle A_j \rangle \langle A_p \rangle}
\]

when \( \langle A_j A_p \rangle < \langle A_j \rangle \langle A_p \rangle \) and set it to zero otherwise. \( CA^3_d \) can be treated in the same manner that \( CA^1_d \) was. Briefly, at each separation d there may exist a number \( n_d \) of values for \( CA^3_d \) and \( CA^3_d(k) < 1, 1 \leq k \leq n_d \). with \( n_d \) a nonempty subset of \( n_d \). Denote the total

\[
\sum_{k=1}^{n_d} CA^3_d(k) = C_d^3
\]
number of separations when $n_d^-$ is nonempty by $N^-$, which makes $N^-$ a nonempty subset of $N$. Therefore, analogously to $\mu_N$, define the amount of Negative Amplitude Correlation by

$$\text{NAC} \equiv \frac{1}{N} \sum_{d=d_1}^{N^-} \frac{1}{n_d^-} \sum_{k=1}^{n_d^-} |CA_d^3(k)| .$$  \hspace{1cm} (39)
IV. (U) STABILITY OF THE TIME AVERAGES

(U) The question arises as to the stability of the synoptic factors as the averaging time $\Delta T$ is changed. This issue is addressed by considering the stability of some representative quantities [denoted by $S_i(\Delta T)$] of the sound field structure as $\Delta T$ is increased. In Ref 5 through 10, an averaging time of about 4 min was employed. (This was chosen after a preliminary study of CA and C for various values of $\Delta T$.) To see how sensitive the quantities $S_i(\Delta T)$ are to changes in $\Delta T$, assume a reference averaging time $\Delta T_r$ of about 4 min and consider how those quantities $S_i(\Delta T)$ vary relative to their values $S_i(\Delta T_r)$ as a function of $\Delta T$. In particular, consider

$$\Delta_i(\Delta T) = \frac{S_i(\Delta T)}{S_i(\Delta T_r)} = DS_i.$$  \hspace{1cm} (40)

(U) A suitable set of representative sound field structure quantities $S_i(\Delta T)$ was chosen for the acoustic synoptic factors of Section II. From these it was determined that the behavior of the quantities

$$S_1 = \bar{A}_a = \frac{1}{J} \sum_{j=1}^{J} \langle A_j \rangle$$  \hspace{1cm} (41)

and

$$S_3 = C_p = \frac{1}{J^2} \sum_{j=1}^{J} \sum_{k=1}^{J} \cos(\theta_j - \theta_k)$$  \hspace{1cm} (42)

were adequate for the purpose of this study.

(U) Let the samples $s(m)$ be indexed by $m$ and summed to $M$, which corresponds to $\Delta T$. Then

$$S_i(\Delta T) = S_i(M) = \frac{1}{M} \sum_{m=1}^{M} s_i(m), \ i = 1, 3.$$  

Since we are dealing with indexed samples, say $s_{ij}(m)$, the above result is obtained as follows (note that the function index $i$ has been suppressed):

$$S(\Delta T) = S(M) = \frac{1}{J^2} \sum_{j=1}^{J} \sum_{k=1}^{J} \frac{1}{M} \sum_{m=1}^{M} s_{jk}(m)$$  

$$= \frac{1}{M} \sum_{m=1}^{M} \left( \frac{1}{J^2} \sum_{j=1}^{J} \sum_{k=1}^{J} s_{jk}(m) \right)$$

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\[ s(m) = \frac{1}{M} \sum_{m=1}^{M} s(m) = \frac{s(1) + \ldots + s(M)}{M} \]

where \( s(1, \ldots, M) \) replaces \( s(1) + \ldots + s(M) \) for brevity. Therefore,

\[ \Delta_1(\Delta T) = \frac{\Delta_1(M)}{\Delta_1(M) + \Delta_2(M) + \Delta_3(M)} = \frac{M_r}{s(1, \ldots, M_r) + s(1, \ldots, M_r) + s(1, \ldots, M_r)} \]

where \( M_r \) corresponds to \( \Delta T_r \).

(U) To quantify the total area between the curve \( \Delta_1(\Delta T) \) and the line \( \Delta_1 = 1 \), the synoptic factor

\[ A\Delta_1 \equiv \sum_{M=1}^{M_r} |\Delta_1(M) - 1| = A\Delta_1 \]

is utilized and appears in Table 3 and Fig 2 through 24.

(U) The behavior of five functions \( \Delta_1(\Delta T) \) for the first 12 signal periods for projector tows 2P3A and 4P1 was plotted. It was found that \( \Delta_1(\Delta T) \) and \( \Delta_3(\Delta T) \) were related to the other three functions in a logical manner, so only these two functions are included in this report for each of the 23 signal periods (Fig 2 through 24). The synoptic factors \( A\Delta_1 \) and \( A\Delta_3 \) are shown at the top of each of these figures and in Table 3. Table 3 shows that \( A\Delta_3 \) is related to CF (except for those corresponding to Fig 10, 19 and 23) in that \( A\Delta_3 \) generally increases as CF decreases. On the other hand, Table 1 shows that there is no correlation between \( A\Delta_1 \) and CF. Therefore, it is concluded that the more stable the time average of \( C_p \) is, the better the coherence \( (\gamma_{12})^2 \) is. On the other hand, the stability of the time average of the amplitude has little impact on the coherence \( (\gamma_{12})^2 \) for the cases investigated in Fig 2 through 24.
V. (U) NORMAL MODE MODEL FOR BEARING STAKE

(U) A normal mode model has been applied to Bearing Stake. This section discusses the suitability of this model for representing Bearing Stake results (Ref 4, 9, 11). The model found in Ref 9 that the Bearing Stake sound field structure is at least three times as complex as that of the open Atlantic and Pacific Oceans. In its Pacific and Atlantic results, Ref 11 found that, with respect to sound field complex\(^{1+\infty}\) across arrays, vertical projections are about 15 times as important as radial projections. With regard to wavefront corrugations, both vertical and radial (ie, rangewise or endfire) projections should be more important than transverse (ie, broadside) projections. (Being essentially two-dimensional, the model cannot treat behavior transverse to the plane of sound projection.) Wavefront corrugations (a sound field propagation effect) and array deformations are indiscernible near broadside. Wavefront corrugation near endfire is mainly caused by the radial field along the array, which is averaged over 4 min of mainly relative radial motion of the array (the aspect angle determines the relative amount of radial and transverse source motion).

(U) The normal mode model can treat radial and vertical projections of the sound field structure but not the transverse sound field structure. Therefore the model was employed to examine the 50-ft difference in vertical projections between the OAMS array (925.4 m or 3036 ft long times 1 deg of tilt equals 53 ft vertical projection) and the LATA (1200 m or 3936 ft long times 1 deg of tilt equals 103 ft vertical projection). For Site 4, the OAMS array subset of the LATA (Ref 8) reduced the LATA vertical projection by about 1/4 to 79.48 ft (ie, 3036 ft long times sin 1\(\frac{1}{2}\) deg) due to the approximately 1/4 length reduction of the LATA. This reduces the residual phase variability by about 1/3 near broadside as shown in Fig 25. Therefore, reducing the vertical projection does reduce the wavefront corrugation significantly. It may even be true that the transverse wavefront corrugation is essentially negligible. This is important for modeling purposes when one considers how difficult it is to remove all array tilt.

(U) The broadside OAMS array behavior with respect to \(C_p\) is observed (see Fig 26) to be about the same as the OAMS array subset of the LATA (Fig 25a). Thus, it appears that the LATA and the OAMS array are comparable at broadside with respect to \(C_p\). For the LATA for Site 2 and for the OAMS array for Site 5, the LATA tilt was probably no worse (if not better) due to “array trimming” (Ref 8) than the OAMS array tilt. The OAMS array on Tow 5P1 (Fig 27) was comparable to the LATA on Tow 2P3A (Fig 29) and on Tow 4P1 (Fig 31) near broadside with respect to \(C_p\). However, the OAMS array on Tow 5P1 (Fig 27) was much better than the LATA on Tow 2P3 (Fig 33) near broadside with respect to \(C_p\). Note that although Fig 27 shows \(C_p\) behavior that is superior to Fig 28 and 30, when only the broadside values of Fig 29 and 31 are compared to Fig 27, it appears that the OAMS array and LATA performed similarly near broadside.

(U) Figure 34 illustrates a plausible explanation of how the angle-variable behavior for the LATA agrees with the normal mode model. The solid line in Fig 34 represents the mean of \(C_p\) and the dashed line represents the standard deviation about this mean. The normal mode model assumes that there is no transverse variation and assumes that the offset below unity of the mean \(C_p\) at broadside is due solely to the remaining vertical projection of the LATA array. The normal mode model generally agrees with the behavior shown in Fig 34 in that in moving from broadside toward endfire, the radial projection of the sound field increases, causing the mean of \(C_p\) to decrease and the standard deviation about this mean to increase.

VI. (U) CONCLUSIONS

(U) In Section I an overview of synoptic factors was given. The usefulness of synoptic factors for improving the utilization of theoretical models was emphasized, and Ref 4 gives examples of this. Section II defined several types of acoustic synoptic factors. Some of these are shown in Tables 1 and 2 for the Bearing Stake data addressed in this report. Two types of acoustic synoptic factors are given. These are total array acoustic synoptic factors and intra-array acoustic synoptic factors. The latter class consists of one-point structure factors and two-point structure factors. The latter type is plotted in Fig 2 through 24, and its structure received much attention in Sections III and IV.

(U) A detailed sound field structure study was developed in Section III. This consisted of developing some further acoustic synoptic factors for the structure of the sound field and finding meaningful relations between these factors. It was shown that the coherence factor CF, as determined by the structured behavior of the classical coherence \((\gamma_{jk}^c)^2\), is very useful for categorizing array coherence behavior. In fact, Table 3 shows that CF is a more reliable parameter than the usual coherence measure \(C_p\), the array phase coherence coefficient. In a discussion of projector Tow 4P1, it was shown that an array turn can introduce a patterned disturbance into \((\gamma_{jk}^c)^2\) and that this behavior can persist for at least 1.75 hr after the turn event. Examining Tow 2P3A showed that even a straight radial tow can produce significant changes in the coherence behavior.

(U) The stability of the time averages used in this coherence study is given in Section IV. The results are shown in Fig 2 through 24 and in Table 3. It is concluded that the more stable the time average of the phase coherence is, the better the classical coherence is, while the stability of time average of the amplitude has far less significance.

(U) Section V discussed normal mode modeling for Bearing Stake. It can model the impact of array tilt and bearing angle. The OAMS array and LATA differences discussed in Ref 7 were resolved, and it was shown that the LATA behavior shown in Fig 34 can be explained by the normal mode model.

(U) The question arises as to how the acoustic synoptic factors are used in practice. Reference 4 gives excellent examples as to the use of some of the acoustic synoptic factors discussed in this report. In Ref 4 the array configuration and tilt are modeled in their software along with the appropriate propagation conditions. This produces the appropriate values for the amplitude and phase. These latter values are then fed into software models of the acoustic synoptic factors (for example, ASG of Eq (13) and \((\gamma_{jk}^c)^2\) of Eq (22)) and then compared with the results obtained when actual Bearing Stake data are applied. These results were instructive, and the phenomenological ambiguities of Ref 7 were resolved. This procedure is the one recommended for predicting the behavior of the acoustic synoptic factors in other environments and/or for arrays other than those used in Bearing Stake. With the understanding gained from such an analysis, insight is gained into the sound field structure for the other environments and/or for the other arrays. With this insight, system relations such as appear in Ref 12 can be applied with confidence.

REFERENCES (U)


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(CONFIDENTIAL)

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$\Delta \phi_o$ gives the off-broadside angle.

(UNCLASSIFIED)

(U) Table 3. Study of CF.
(CONFIDENTIAL)

(U) Figure 1. Bearing Stake Site Locations.
(U) Figure 2a. Normalized amplitude covariance.
Figure 25. Array phase coherence.
Figure 2c. Classical coherence.
Figure 2d. Array amplitude covariance.
Source Bearing 102.70 deg
LATA Tow 2P3A, 27 April 1977 1635:15 to 1639:15 Zulu, Signal, Freq 25.10 Hz
RDSI = 8.626

Figure 2e. Amplitude time stability.
(CONFIDENTIAL)

(U) Figure 2f. Phase time stability.
Figure 3a. Normalized amplitude covariance.
Figure 3b. Array phase coherence.
Figure 3c. Classical coherence.

HYDROPHONE PAIR SEPARATION (FT)
Figure 3d. Array amplitude covariance.
Figure 3e. Amplitude time stability.
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SOURCE BEARING 101.45 DEG
LATA Tow 2P3A, 27 APRIL 1977 1859-47 TO 1903-47 ZULU, SIGNAL, FREQ 20.08 HZ
ALOS- .555

AVERAGING-TIME (SEC)

(U) Figure 3f. Phase time stability.
Figure 4a. Normalized amplitude covariance.
Figure 4b. Array phase coherence.
HYDROPHONE PAIR SEPARATION (FT)

(CONFIDENTIAL)

(U) Figure 4c. Classical coherence.
Figure 4d. Array amplitude covariance.
Figure 4e. Amplitude time stability.
(U) Figure 4f. Phase time stability.
(CONFIDENTIAL)

(U) Figure 5a. Normalized amplitude covariance.
Figure 5b. Array phase coherence.
(U) Figure 5c. Classical coherence.
(CONFIDENTIAL)

(U) Figure 5d. Array amplitude covariance.
Figure 5e. Amplitude time stability.
(U) Figure 5f. Phase time stability.
(U) Figure 6a. Normalized amplitude covariance.
(U) Figure 6b. Array phase coherence.
HYDROPHONE PAIR SEPARATION (FT)

(U) Figure 6c. Classical coherence.

CONFIDENTIAL
(U) Figure 6d. Array amplitude covariance.
Figure 6e. Amplitude time stability.
CONFIDENTIAL

LATA Tow 2P3A, 27 APRIL 1977 1939:16 TO 1943:16 ZULU, SIGNAL, FREQ 25.10 HZ
ADS3= .495

(U) Figure 6f. Phase time stability.

CONFIDENTIAL
Figure 7a. Normalized amplitude covariance.
CONFIDENTIAL

Figure 7b. Array phase coherence.
(U) Figure 7c. Classical coherence.
(U) Figure 7d. Array amplitude covariance.
Figure 7e. Amplitude time stability.
(U) Figure 7f. Phase time stability.
(U) Figure 8a. Normalized amplitude covariance.
(U) Figure 8b. Array phase coherence.
GJ3- .9645 MND3-.847 SN03-.025 HS5-.106 CC-.9721 SF-.0027 NAC-.0000

HYDROPHONE PAIR SEPARATION (FT)

(U) Figure 8c. Classical coherence.
(U) Figure 8d. Array amplitude covariance.
Figure 8e. Amplitude time stability.
(U) Figure 8f. Phase time stability.
Figure 9a. Normalized amplitude covariance.
Figure 9b. Array phase coherence.
Figure 9c. Classical coherence.

(U) Figure 9c. Classical coherence.
(U) Figure 9d. Array amplitude covariance.
Figure 9e. Amplitude time stability.
(U) Figure 9f. Phase time stability.
Figure 10a. Normalized amplitude covariance.
LATA Tow 2P3A, 27 April 1977 2205:20 to 2209:20 Zulu, Signal, Freq 25.10 Hz
G..-.0450 MND-.827 SND-.116 MS2-.068 CF-.9828 SF-.0039 NAO-.4816

Figure 10b. Array phase coherence.

(CONFIDENTIAL)
Figure 10c. Classical coherence.
(U) Figure 10d. Array amplitude covariance.
Figure 10e. Amplitude time stability.
Figure IV. Phase time stability.

(U) Figure 10f. Phase time stability.
HYDROPHONE PAIR SEPARATION (FT)

(CONFIDENTIAL)

(U) Figure 11a. Normalized amplitude covariance.
(U) Figure 11b. Array phase coherence.
Figure 11c. Classical coherence.
LATA Tow 2P3A, 27 April 1977 2231:40 to 2235:40 Zulu, signal, freq 25.10 Hz.

*Figure 11d. Array amplitude covariance.*
(U) Figure 1le. Amplitude time stability.
(U) Figure 11f. Phase time stability.
(U) Figure 12a. Normalized amplitude covariance.
Figure 12b. Array phase coherence.
LATA Tow 2P3A, 27 APRIL 1977 2333:49 TO 2337:49 ZULU, SIGNAL, FREQ 25.12 Hz
CJ= .7238 MND= .609 SNDS= .040 MS3= .211 CF= .9447 SF= .0084 MNO= .4087

CONFIDENTIAL

HYDROPHONE PAIR SEPARATION (FT)

(CONFIDENTIAL)

(U) Figure 12c. Classical coherence.

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CONFIDENTIAL
(U) Figure 12d. Array amplitude covariance.
Figure 12e. Amplitude time stability.
(CONFIDENTIAL)

(U) Figure 12f. Phase time stability.

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CONFIDENTIAL
(CONFIDENTIAL)

Figure 13a. Normalized amplitude covariance.
(U) Figure 3b. Array phase coherence.
(U) Figure 13c. Classical coherence.
Figure 13d. Array amplitude covariance.
Figure 13e. Amplitude time stability.
Figure 13f. Phase time stability.
Figure 14a. Normalized amplitude covariance.
(CONFIDENTIAL)

(U) Figure 14b. Array phase coherence.
Hydrophone pair separation (ft)

Figure 14c. Classical coherence.
(U) Figure 14d. Array amplitude covariance.
Figure 14e. Amplitude time stability.
(U) Figure 14f. Phase time stability.
(U) Figure 15a. Normalized amplitude covariance.
(U) Figure 15b. Array phase coherence.
Figure 15c. Classical coherence.
(CONFIDENTIAL)

(U) Figure 15d. Array amplitude covariance.

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CONFIDENTIAL
(U) Figure 15e. Amplitude time stability.
CONFIDENTIAL

SOURCE BEARING 107.70 DEG
LATA Tow 4P1, 13 March 1977 1056:11 TO 1100:11 ZULU, SIGNAL, FREQ 20.04 HZ

![Phase time stability graph]

AVERAGING-TIME (SEC)

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CONFIDENTIAL

(U) Figure 15f. Phase time stability.
Figure 16a. Normalized amplitude covariance.
(U) Figure 16b. Array phase coherence.
Figure 16c. Classical coherence.
(U) Figure 16d. Array amplitude covariance.
(CONFIDENTIAL)

(U) Figure 16e. Amplitude time stability.
(U) Figure 16f. Phase time stability.
Figure 17a. Normalized amplitude covariance.
(U) Figure 17b. Array phase coherence.
(U) Figure 17c. Classical coherence.
Figure 17d. Array amplitude covariance.
(U) Figure 17e. Amplitude time stability.
Figure 17f. Phase time stability.
Figure 18a. Normalized amplitude covariance.
(CONFIDENTIAL)

(U) Figure 18b. Array phase coherence.
Figure 18c. Classical coherence.

(U) Figure 18c. Classical coherence.
Figure 18d. Array amplitude covariance.
Figure 18e. Amplitude time stability.
Figure 18f. Phase time stability.

(U) Figure 18f. Phase time stability.
(U) Figure 19a. Normalized amplitude covariance.
Figure 19b. Array phase coherence.
CONFIDENTIAL

(U) Figure 19c. Classical coherence.

(CONFIDENTIAL)
Figure 19d. Array amplitude covariance.
Figure 19e. Amplitude time stability.
LATA Tow 4P1, 13 March 1977 1146+57 to 1150+57 Zulu, Signal, Freq 25.05 Hz
ADSS- 5.830

(CONFIDENTIAL)

(U) Figure 19f. Phase time stability.
(U) Figure 20a. Normalized amplitude covariance.
(CONFIDENTIAL)

(U) Figure 20b. Array phase coherence.
(U) Figure 20c Classical coherence.
(U) Figure 20d. Array amplitude covariance.
Figure 20e. Amplitude time stability.
(CONFIDENTIAL)

(U) Figure 20f. Phase time stability.
Figure 21a. Normalized amplitude covariance.
(U) Figure 21b. Array phase coherence.
CONFIDENTIAL

SOURCE BEARING 144.90 DEG
LAT A Tow 4PI, 13 MARCH 1977 1255+51 TO 1259+51 ZULU, SIGNAL, FREQ 20.04 Hz
C(3)-.6237 MND3-.613 SND3-.039 MS3-.202 OF-.8381 SF-.0078 NAC-.5149

HYDROPHONE PAIR SEPARATION (FT)

(CONFIDENTIAL)

(U) Figure 21c. Classical coherence.
Figure 21d. Array amplitude covariance.
(U) Figure 21e. Amplitude time stability.
(CONFIDENTIAL)

(U) Figure 21f. Phase time stability.
HYDROPHONE PAIR SEPARATION (FT)

(CONFIDENTIAL)

(U) Figure 22a. Normalized amplitude covariance.
Figure 22b. Array phase coherence.

(C) Figure 22b. Array phase coherence.

(C) CONFIDENTIAL
Figure 22c. Classical coherence.
(U) Figure 22d. Array amplitude covariance.
(U) Figure 22e. Amplitude time stability.
(U) Figure 22f. Phase time stability.
Figure 23a. Normalized amplitude covariance.
(U) Figure 23b. Array phase coherence.
Figure 23c. Classical coherence.
Figure 23d. Array amplitude covariance.
(U) Figure 23e. Amplitude time stability.
SOURCE BEARING 141.90 DEG
LATA Tow 4P1, 13 March 1877 1319:05 to 1323:05 Zulu, SIGNAL, FREQ 25.05 Hz
ADS3-2.496

(U) Figure 23f. Phase time stability.
HYDROPHONE PAIR SEPARATION (FT)

(CONFIDENTIAL)

(U) Figure 24a. Normalized amplitude covariance.
(U) Figure 24b  Array phase coherence.
CONFIDENTIAL

Figure 2c. Classical coherence.
(U) Figure 24d. Array amplitude covariance.
(U) Figure 24e. Amplitude time stability.
Figure 24f. Phase time stability.
(U) Figure 25. Comparison of $C_p$ versus source bearing for Tow 4P1 for the LATA Subset (upper plot) and the ordinary LATA (lower plot).
(CONFIDENTIAL)

(C) Figure 26. Phase coherence as a function of range; OAMS array; Tow 4P1, 25 Hz; 13 March 1977.
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