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TABLES AND GRAPHS OF THE THEORETICAL PEAK PRESSURES, ENERGIES, AND IMPULSES OF SHOCK WAVES FROM EXPLOSIVE SOURCES IN SEA WATER

by

Stuart R. Brinkley, Jr., and John G. Kirkwood Cornell University

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E. Bright Wilson, Jr., Chief

Division 2

Effects of Impact and Explosion



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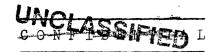
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### Preface

The work described in this report is pertinent to War Department Project OD-03 and to Navy Department Project NO-224. This work was carried out and reported by Cornell University as part of its performance under Contract OEMsr-121.

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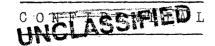
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# TABLES AND GRAPHS OF THE THEORETICAL PEAK PRESSURES, ENERGIES, AND INPULSES OF SHOCK WAVES FROM EXPLOSIVE SOURCES IN SEA WATER

### Abstract

A theory of the propagation of shock waves from explosive sources was presented in NDRC Report A-318 (OSRD-4814). In that report, a pair of ordinary differential equations for peak pressure and shock-wave energy as functions of the distance from the source were formulated from the equations of hydrodynamics by imposing a similarity restraint on the shape of the energy-time curve of the shock wave. Two-parameter families of peak pressure-distance curves are obtained by the solution of these propagation equations. The parameters are conveniently chosen as the initial values of the pressure and shock-wave energy.

In the present report, tables and graphs of the twoparameter families of curves for shock waves from explosive sources in sea water are presented. A method is outlined for the determination of the parameters from experimental values of the peak pressure and impulse over a limited range of distances from the source.

I. DESCRIPTION OF THE TABLES AND GRAPHS OF THE PROPERTIES OF SHOCK WAVES FROM EXPLOSIVE SOURCES IN WATER

### 1. Introduction

A theory of propagation of shock waves from explosive sources has been described in an earlier report. A pair of ordinary differential equations for peak pressure and shock-wave energy as functions of distance from the source were obtained from the equations of hydrodynamics and the Hugoniot relations by imposing a similarity restraint on the shape of the energy-time curve and by utilizing the second law of thermodynamics to determine, at an arbitrary distance, the distribution of the initial energy input from the explosive between energy available for further propagation and dissipated energy residual in the fluid already traversed by the shock wave. The theory

Theory of the propagation of shock waves from explosive sources in air and water, by J. G. Kirkwood and S. R. Brinkley, Jr., NDRC Report A-318 (OSRD-4814).

takes proper account of the finite entropy increment in the fluid resulting from the passage of the shock wave and permits the use of the exact Hugoniot curve for the fluid.

The basic equations of the theory are

$$\frac{dp}{dR} + \frac{p}{R} n(p) = -\frac{R^2 p^4}{K} F(p),$$

$$\frac{dK}{dR} = -p^3 R^2 f(p),$$
(1)

where

$$f(p) = \frac{12 \sigma^{3}}{\sigma + 1} \frac{h}{c_{o}^{2} p^{3}},$$

$$F(p) = \frac{12G}{p} \frac{\sigma}{\sigma + 1} \left(\frac{c_{o}}{U}\right)^{2} \frac{1}{2(1 + g) - G},$$

$$n(p) = \frac{4\rho_{o}/\rho + 2(1 - \rho_{o}/\rho)G}{2(1 + g) - G},$$

$$G = 1 - \left(\frac{\rho_{o}U}{\rho c}\right)^{2},$$

$$g = 1 - \frac{p}{U} \frac{dU}{dp},$$

and  $\underline{p}$  is the excess peak pressure measured in units of a characteristic pressure  $\underline{B}$ , at a distance  $\underline{Ra}$  from the center of a spherical charge of explosive of radius  $\underline{a}$ .  $\underline{B}$  is the characteristic pressure of the Tait equation of state,

$$p = B[(\rho/\rho_0)^{\tilde{g}} - 1]$$
 (2)

and  $\underline{\tilde{g}}$  is the exponent of  $\rho/\rho_0$  in the Tait equation. The quantities  $\underline{\rho}$  and  $\rho_0$  are the densities of the water at distances  $\underline{R}$  and infinity, respectively;  $\underline{c}$  and  $\underline{c}_0$  denote the sound velocity in water at densities  $\underline{\rho}$  and  $\rho_0$ , respectively;  $\underline{U}$  is the velocity of the shock front;  $\underline{h}$  is the specific enthalpy increment of an element of fluid, traversed by a shock wave of peak pressure  $\underline{p}$ , after re-

turn to pressure p along its new adiabatic. The variable  $\underline{K}$  is related to the shock-wave energy  $\underline{\epsilon}$  per gram of the initial explosive charge by

$$\varepsilon = \frac{3 + 1}{143^2} \frac{B}{\rho_e} K, \tag{3}$$

where  $\rho_{\rm e}$  is the density of loading of the explosive. The positive impulse is given by

$$I = p\Theta B, \qquad (4)$$

where

$$-\frac{\mathbf{a_o}}{\Theta} = \frac{\mathbf{U}}{\mathbf{R}} \left\{ \frac{2}{\mathbf{G}} + \left[ \left( \frac{1+\mathbf{g}}{\mathbf{G}} \right) \frac{\mathbf{p}}{\mathbf{p_o}} + \left( 1 - \frac{\mathbf{p}}{\mathbf{p_o}} \right) \right] \frac{\mathbf{R}}{\mathbf{p}} \frac{\mathbf{dp}}{\mathbf{dR}} \right\}.$$

Tables of the functions defined in Eqs. (1) and  $(l_4)$  are presented in an appendix.

# 2. Description of the tables and graphs

The two-parameter family of peak pressure-distance curves and the family of shock-wave energy-distance curves for selected values of the constants of integration were obtained by numerical integration of Eqs. (1), employing the methods of Part III of the previous report. Lach member of the two families is characterized by a value of  $p_1$ , the excess peak pressure at the initial instant of time measured in kilobars and by a value of  $Q_1$ , related to the initial shock-wave energy by

$$Q_{1} = \frac{\mu \eta^{2}}{\eta + 1} \frac{\rho_{e} \ell_{1}}{Bp_{1}} . \tag{5}$$

The family of peak pressure-distance curves is presented in tabular form in Table I. In this and following tables, a reduced distance variable equal to  $\underline{\mathbb{R}}(\rho_{\mathrm{e}}/\mathbb{W})^{1/3}$  is employed, where  $\underline{\mathbb{R}}$  is expressed in feet,  $\rho_{\mathrm{e}}$  in grams per cubic

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centimeter, and  $\underline{W}$  in pounds. In order to facilitate interpolation with respect to distance, the entries of the table are of the function  $10^{-3} \mathrm{Rp} (\rho_{\mathrm{e}}/\mathrm{W})^{1/3}$  with p expressed in pounds per square inch. 2/

Portions of the peak pressure-distance curves for moderately large distances are presented in a series of graphs, Figs. 1 to 6. Each graph gives a family of peak pressure-distance curves for a fixed value of the parameter  $Q_1$ , plotted as  $\log \operatorname{Rp}(\rho_e/W)^{1/3}$  versus  $\log \operatorname{R}(\rho_e/W)^{1/3}$ . These graphs are designed to aid in the determination of the parameters from an experimental peak pressure-distance curve in a manner to be described in the following.

The family of shock-wave energy-distance curves is tabulated in Table II, the entries of which are  $\rho_{\rm e} \ell$  with the reduced distance defined in the foregoing as argument. The units of  $\underline{\ell}$  are calories per gram of explosive.

The family of positive impulse-distance curves was constructed with the aid of Eqs. (3), and it is tabulated in Table III. Here, the function  $10^3 \mathrm{RI}(\rho_\mathrm{c}/\mathrm{W})^{2/3}$  is listed with the reduced distance as argument. The units of I are pound-second per square inch.

# 3. Determination of the initial parameters from experimental data

The graphs, Figs. 1 to 6, are designed to aid in the determination of the parameters of those members of the families of peak-pressure curves that give the best fit to a particular set of experimental data. For this purpose, the experimental data may be plotted to the correct scale on tracing cloth and compared with the theoretical values by superimposing the tracing on the published graphs. Unless the data are of high accuracy, this procedure will not lead to a unique determination of the parameters but to the selection of several pairs of values of the parameters between which a choice could not easily be made by visual methods. The method of least squares may then be used to narrow the choice of parameters. One computes the mean of the squares of the weighted deviations of the individual experimental values of the peak pres-

 $<sup>\</sup>frac{2}{N}$  Note that in the tables p is expressed in pounds per square inch, while in the equations p was a dimensionless quantity expressed in units of B.

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sure and positive impulse from the theoretical values given by the curves selected by visual matching, employing the reciprocal of the root mean square deviations of the experimental points at a given distance from their mean as the weight factor at that distance. One then selects the parameter set yielding the smallest value of this quantity. 3

II. TABLES AND GRAPHS OF THE PEAK PRESSURE, ENERGY,
AND IMPULSE OF SHOCK WAVES FROM EXPLOSIVE
SOURCES IN SEA WATER

## Symbols and Units

Symbol	Unit	<u>Definition</u>
R	ft	Distance from the center of spherical charge of explosive.
$ ho_{ m e}$	gm/cm <sup>3</sup>	Density of explosive.
W	lb	Weight of charge of explosive.
р	lb/in <sup>2</sup>	Excess peak pressure of shock wave.
P1	kilobars	Initial excess peak pressure.
Qı	inche Marie Santa	Initial value of reduced energy variable.
٤	cal/gm	Shock-wave energy.
I	lb-sec/in?	Impulse of shock wave.

Note: The first tabulated value of the distance argument  $(\rho_e/\mathbb{V})^{1/3} = 0.156$ , corresponds to the surface of the sphere of intact explosive.

<sup>3/</sup> See E. T. Whittaker and G. Robinson, The calculus of observations, 3rd (Blackie and Son, 1940), pp. 182, 221.

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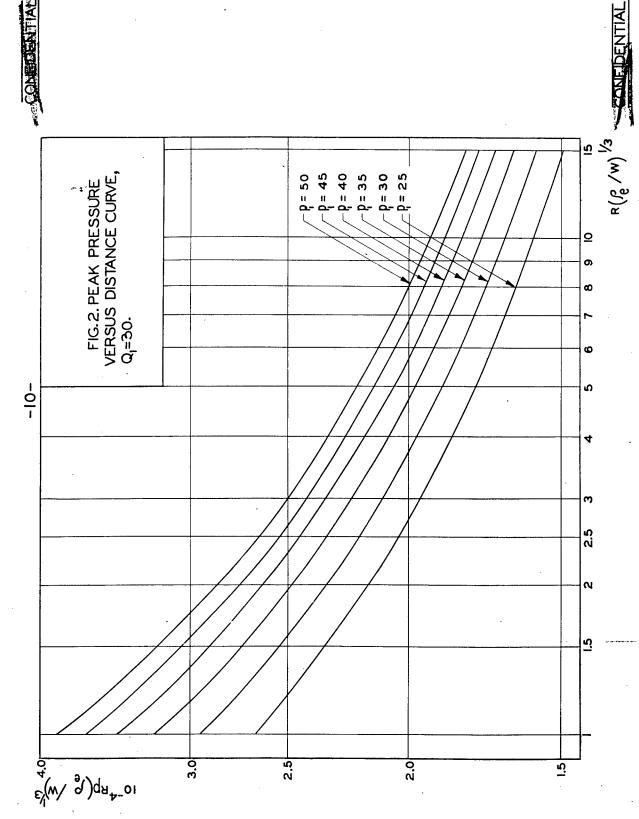
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		3.0	18.0 19.0 20.0 22.2 22.2	22.22 22.22 23.22 24.22	20.7 22.3 23.7 26.0 26.0	23.23.33.33.33.33.33.33.33.33.33.33.33.3	22.7 24.6 26.2 27.7 29.0	67 1-80 C-
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		0•3	120 20 20 20 20 20 20 20 20 20 20 20 20 2	13.5 56.5 72.5 72.5 72.5 73.5 74.5 75.5 75.5 75.5 75.5 75.5 75.5 75	44.2 55.4 74.3 74.3 81.3	441.8 52.6 60.4 68.0 75.7 83.2	45.2 53.2 61.1 69.0 76.7 84.4	45.6 63.7 69.8 77.7 85.4
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		0.156	56.0 7.02 1.02 1.3	58.0 79.4 102 113	56.7 68.0 79.4 90.7 102	56.7 66.0 79.4 90.7 102	56.7 68.0 79.4 90.7 102 113	56,7 68.0 79.4 90.7 112
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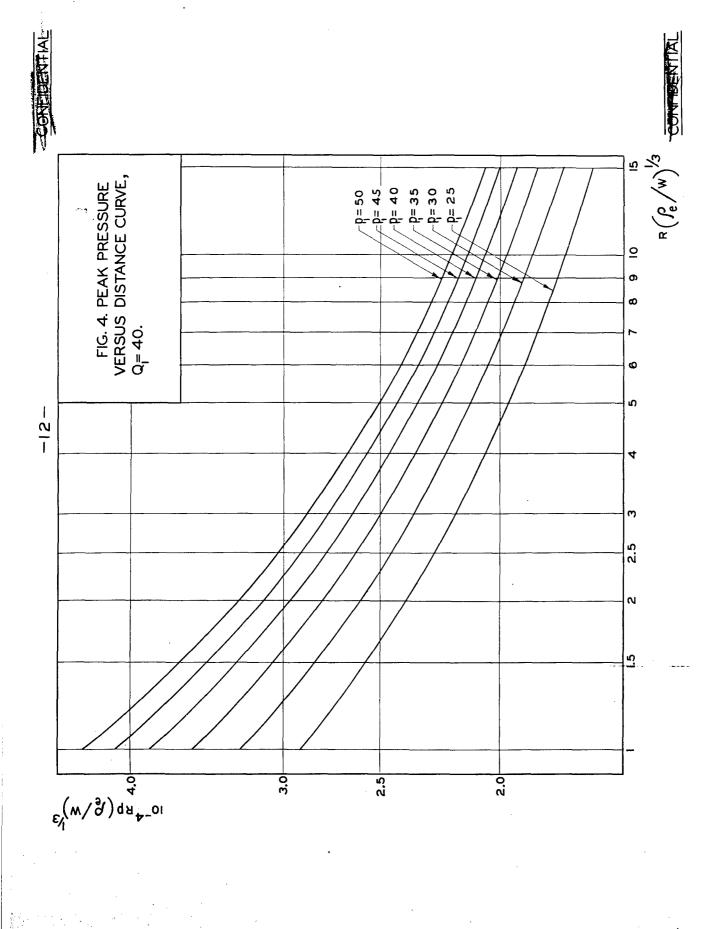
Table II. The shock-wave energy.

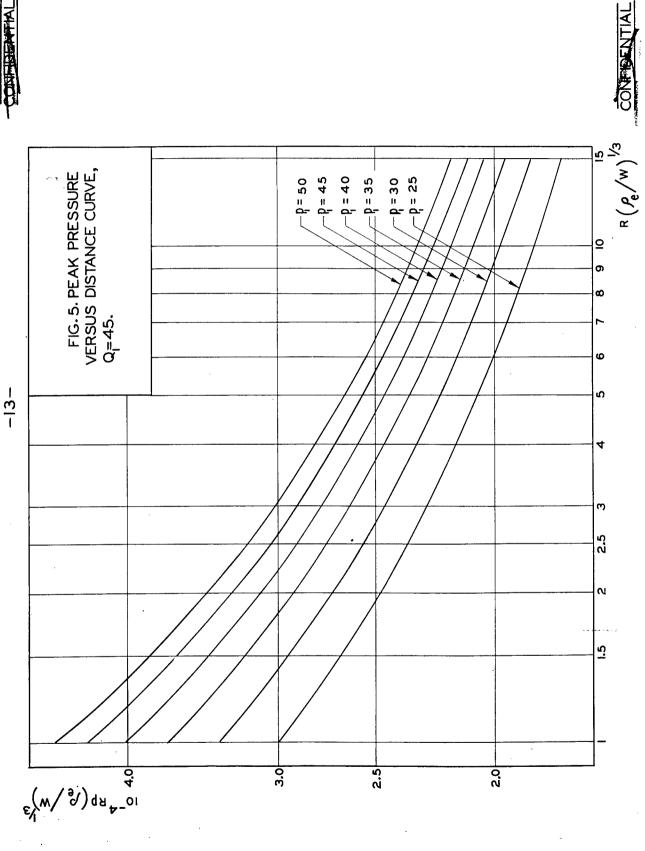
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		٠ <u>۲</u>	25	%	35.	01	1,5	55	

Table III. The impulse.

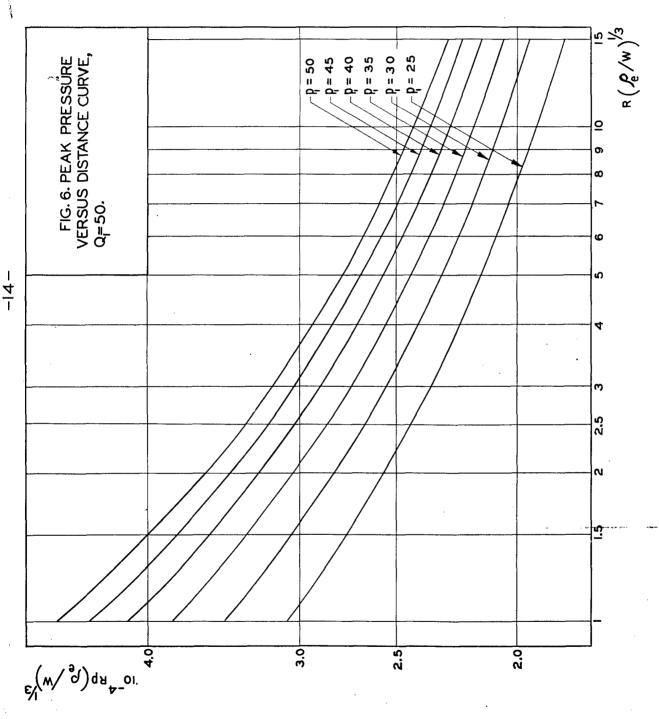
0.2         0.3         0.5         0.7         1.0           242         279         321         378         433           242         301         326         392         449           257         303         326         392         449           257         303         326         405         449           267         333         338         405         449           280         355         364         422         488           315         363         424         499           284         376         389         465         536           317         383         407         475         552           318         403         484         564         536           317         383         407         475         552           318         417         507         585           318         417         507         585           323         415         489         564           324         413         522         606           344         450         540         650           344         480 <td< th=""></td<>
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504 630 534 654 567 673 569 682 634 687 669 705











#### APPENDIX

# Methods of Numerical Integration

In the present section, methods that are appropriate for the numerical integration of the basic differential equations, Eqs. (1), are described and tables are given of the auxiliary functions required for this integration.

For shock waves in water, it is convenient to employ the following variables:

In these variables, Eqs. (1) become

$$-\frac{dY}{dX} = \ell(p) + \frac{P^2}{Q} m(p),$$

$$\frac{dZ}{dX} = \ell(p) + \frac{P^2}{Q} \eta(p),$$
(A2)

where

$$\ell(p) = n(p)-1,$$
 $m(p) = F(p),$ 
 $\eta(p) = F(p)-f(p).$ 

Tables of the functions  $\ell(p)$ , m(p), and  $\eta(p)$  are presented in Table IV. Table V lists certain auxiliary quantities necessary for the calculation of the impulse by Eq. (4).

A procedure for the numerical integration of Eqs. (A2) has been outlined in an earlier report.—

It is suggested that the integration be started with the Runge-Kutta method, and an iterative procedure for continuing the integration is outlined.

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Table IV. Functions for the integration of the propagation equation in sea water.

equation in sea water.					
p (Units of B)	ℓ(p)	m(p)	n (p)		
0	0	1.5000	0.5000		
<b>0.01</b>	0.001495	1.4843	•4950		
•02	.002921	1.4670	.4882		
•03	.004355	1.4524	.4839		
.04	.005725	1.4363	•4778		
•05	.007075	1.4207	.4721		
•06	•008424	1.4054	.4665		
•07	.009774	1.3871	•4600		
•08	•01112	1.3759	.4560		
•09	.01247	1.3617	•4510		
0.1	•01380	1.3479	•4462		
•2	•02622	1.2175	•3984		
•3	.03745	1.1058	<b>.</b> 3576		
•4	.04763	1.0095	•3215		
•5	•05676	0.9305	•2939		
•6	•06581	•8555	•2673		
• 7	.07388	•7923	•2450		
•8	•08156	•7378	•2268		
•9	•08867	<b>.</b> 6886	•2102		
1.0	•09517	.61439	<b>.</b> 1945		
1,5	.1235	.4793	.1404		
2.0	.1416	•3844	•1133		
2.5	.1563	•3123	•0889		
3.0	.1696	.2524	•0690		
3.5	.1815	•2048	<b>.</b> 05 <b>36</b>		
4.0	.1919	<b>.</b> 1695	.0428		
4.5	•2009	.1465	•0366		
5.0	•2085	<b>.</b> 1348	•0345		
10	•2506	<b>.</b> 05140	•01276		
15	•2678	•02770	.00694		
20 .	•2848	.01731	•00452		
25	•3143	•01242	•00359		

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Table V. Auxiliary table for blast waves in sea water.

p (Units of B)	U/c <sub>o</sub>	. g	G	Alpo
0.01 .02 .03 .04 .05 .06 .07 .08 .09	1.003 1.006 1.008 1.011 1.014 1.017 1.020 1.023 1.025	0.9971 .9944 .9917 .9890 .9863 .9838 .9812 .9787	0.005654 .01121 .01669 .02211 .02737 .03258 .03773 .04278	1.001393 1.002774 1.002774 1.005502 1.006846 1.008180 1.009499 1.01082 1.01213
0.1 .2 .3 .4 .5 .6 .7 .8	1.027 1.054 1.072 1.103 1.126 1.148 1.170 1.191	.9737 .9512 .9316 .9147 .8995 .8859 .8738 .8627	.05266 .09779 .1371 .1706 .2011 .2283 .2525 .2748 .2949	1.01342 1.02584 1.03738 1.04818 1.05835 1.06797 1.07705 1.08568
1.0 1.5 2.0 2.5 3.0 3.5 4.0	1.231 1.321 1.407 1.486 1.555 1.619 1.682 1.742	.8435 .8069 .7654 .7408 .7298 .7241 .7110	•3130 •3862 •4500 •4975 •5192 •5296 •5515 •5682	1.1018 1.1368 1.1673 1.1935 1.2136 1.2302 1.2478 1.2630
5 10 15 20 25	1.797 2.272 2.462 2.961 3.252	.6906 .6381 .6149 .5943 .5638	•5799 •6593 •6944 •7142 •7255	1.2762 1.3720 1.4294 1.4687 1.4945