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**EFFECT OF BANDWIDTH ON WIDEBAND-STAP  
PERFORMANCE (PREPRINT)**

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# Effect of Bandwidth on Wideband-STAP Performance

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**Abstract**—A wideband signal occupies a finite bandwidth that is significant compared to its carrier frequency inasmuch as when transmitted, its returns cause bandwidth dispersion across the antenna. It is shown here that the effect of the finite bandwidth is to introduce a set of uncorrelated return signals for every physical scatter in the field. Further, each such uncorrelated return contains a set of coherent signals with different directional and Doppler components that result from a jittering effect both in angle and Doppler domain. As a result adaptive clutter cancellation using traditional processing schemes don't work well. Though in principle it is possible to correct these decorrelating effects by 3D space-time adaptive processing (STAP), the present day methods are quite costly and difficult to implement. In addition to the new wideband signal modeling framework mentioned above, we discuss a new hierarchical processing scheme which has the potential for dramatically reducing both processing and sample support burdens.

## I. INTRODUCTION

In the context of target detection in a wideband transmit/receive environment, consider an  $N$  element,  $M$  pulse array receiving signals from its field of view. In the narrowband setup with  $e^{j\omega_s t}$  representing the transmit signal, after down converting to base band the array output space-time data vector  $\mathbf{x}(t)$  from any range bin has the form

$$\mathbf{x}(t) = \sum_k \alpha_k \mathbf{s}(\theta_k, \omega_{d_k}) + \mathbf{n}(t) \quad (1)$$

where  $\alpha_k$  represents the random scatter return from the  $k^{\text{th}}$  bin in the cross-range (azimuth) domain. Here

$$\mathbf{s}_k = \mathbf{s}(\theta_k, \omega_{d_k}) = \underline{b}(\omega_{d_k}) \otimes \underline{a}(\theta_k) \quad (2)$$

represents the  $MN \times 1$  space-time steering vector with

$$\underline{a}(\theta_k) = \left[ 1 \quad e^{-j\omega_o \frac{d \sin \theta_k}{c}} \quad e^{-j\omega_o \frac{2d \sin \theta_k}{c}} \quad \dots \quad e^{-j\omega_o \frac{(N-1)d \sin \theta_k}{c}} \right]^T \quad (3)$$

representing the spatial steering vector along  $\theta_k$  and

$$\underline{b}(\omega_{d_k}) = \left[ 1 \quad e^{-j\pi\omega_{d_k}} \quad e^{-j2\pi\omega_{d_k}} \quad \dots \quad e^{-j(M-1)\pi\omega_{d_k}} \right]^T \quad (4)$$

representing the temporal steering vector along the Doppler frequency

$$\omega_{d_k} = \frac{2VT \sin \theta_k}{\lambda / 2}. \quad (5)$$

In (5),  $V$  represents the platform speed,  $T$  the pulse repetition interval and  $\lambda$  the operating wavelength.

In an uncorrelated clutter and noise scene, from (1), we get the space-time covariance matrix to be

$$\mathbf{R}_x = \sum P_k \mathbf{s}_k \mathbf{s}_k^* + \sigma^2 \mathbf{I} \quad (6)$$

and its properties are well documented [1]. Because of the angle-Doppler linear relationship in (5), in the narrowband case for a specific look direction, the adaptive space-time processor using (6) will null out clutter components induced along that direction thereby making target detection possible.

## II. WIDBAND STAP PERFORMANCE

To obtain a similar understanding for the structure of the data and array output covariance matrix in the wideband case, assume that matched filtering and adaptive processing is employed in that case as well. Consider a wideband signal  $f(t)$  with bandwidth  $B_o$  that is modulated by frequency  $\omega_o$  to generate  $f(t)e^{j\omega_o t}$  which acts as the transmit signal. When a wideband waveform is employed in the context of space time adaptive processing (STAP), the finite bandwidth dispersion across the antenna must be taken into account. To understand the bandwidth dispersion issues, let  $e^{j\omega_s t}$  and  $f(t)e^{j\omega_o t}$  refer to the transmit waveforms in the narrowband and wideband situation respectively across an  $N$  element array employing  $M$  pulses. Here (see Fig. 1)

$$f(t) \leftrightarrow F(\omega) = \begin{cases} F(\omega), & |\omega| \leq B_o / 2 \\ 0, & |\omega| > B_o / 2 \end{cases}. \quad (7)$$

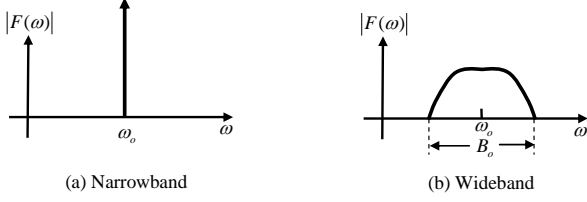


Fig. 1 Narrowband and wideband transmit signals.

Let  $y_1^{(1)}(t) = f(t)e^{j\omega_0 t}$  represent the scattered return at the reference sensor due to the first pulse from a ground location at an azimuth angle  $\theta_k$  for a specific range. In that case, the  $i^{\text{th}}$  sensor output (suppressing the noise component) is given by

$$y_1^{(i)}(t) = f(t - (i-1)\tau_1) e^{j\omega_0(t - (i-1)\tau_1)}, \quad i=1, 2, \dots, N \quad (8)$$

where

$$\tau_1 = \frac{d \sin \theta_k}{c} \quad (9)$$

represents the interelement time delay corresponding to azimuth angle  $\theta_k$ . This gives the first pulse output across the array to be

$$\underline{y}_1(t) = [f(t)e^{j\omega_0 t}, \dots, f(t - (N-1)\tau_1) e^{j\omega_0(t - (N-1)\tau_1)}]^T. \quad (10)$$

Similarly, the received signal at the reference sensor due to the  $n^{\text{th}}$  pulse is given by

$$\begin{aligned} y_n^{(1)}(t) &= f(t - (n-1)\tau_2) e^{j\omega_0(t - (n-1)\tau_2)} \\ &= f(t - (n-1)\tau_2) e^{-j\pi(n-1)\omega_{d_k}} e^{j\omega_0 t} \end{aligned} \quad (11)$$

where  $\omega_{d_k}$  is as defined in (5) and

$$\tau_2 = \frac{2VT \sin \theta_k}{c} \quad (12)$$

represents the interpulse delay for azimuth angle  $\theta_k$ . After demodulation, the output at the  $i^{\text{th}}$  sensor is given by

$$\begin{aligned} z_n^{(i)}(t) &= f(t - (n-1)\tau_2 - (i-1)\tau_1) e^{-j\pi(n-1)\omega_{d_k}} e^{-j(i-1)\tau_1}, \\ & \quad i=1, 2, \dots, N, \quad n=1, 2, \dots, M. \end{aligned} \quad (13)$$

Notice that the time-delayed version of the baseband waveform in the return signal is characteristic of the wideband situation. Using (11) - (12) with  $\underline{z}_n(t)$ ,  $n=1, 2, \dots, M$  representing the array output vector

for the  $1^{\text{st}}$ ,  $2^{\text{nd}}$ ,  $\dots$ ,  $M^{\text{th}}$  pulse, we get the space-time vector  $\mathbf{z}(t)$  due to a single scatter to be [2] (details omitted)

$$\mathbf{z}(t) = \begin{bmatrix} \underline{z}_1(t) \\ \underline{z}_2(t) \\ \vdots \\ \underline{z}_M(t) \end{bmatrix} \quad (14)$$

where

$$\underline{z}_n(t) = [z_n^{(1)}(t), z_n^{(2)}(t), \dots, z_n^{(N)}(t)]^T. \quad (15)$$

This gives

$$\mathbf{z}(t) = \mathbf{f}_k(t) \circ \mathbf{s}(\theta_k, \omega_{d_k}) \quad (16)$$

where  $\mathbf{s}(\theta_k, \omega_{d_k})$  is as defined in (2) - (5) and  $\circ$  represents the element-wise (Schur-Hadamard) multiplication. Here  $\mathbf{f}_k(t)$  represents an  $MN \times 1$  transmit signal dependent vector whose  $(iN+n)^{\text{th}}$  element is given by (see (13))

$$\begin{aligned} f(t - n\tau_2 - i\tau_1), \quad & i=0, 1, 2, \dots, N-1, \\ & n=0, 1, 2, \dots, M-1. \end{aligned} \quad (17)$$

Let

$$\mathbf{T}_k(\tau) = E \{ \mathbf{f}_k(t) \mathbf{f}_k^*(t + \tau) \} > 0 \quad (18)$$

represent the covariance matrix associated with the return wideband signal vector  $\mathbf{f}_k(t)$  in (16). We will assume that the actual return signal components in (17) resulting from random scatter modulations are stochastic in nature with wide sense stationary behavior, such that its power spectrum coincides with the original wideband transmit signal spectrum in Fig. 1 (b). This is a reasonable assumption, since the returns are inherently stochastic due to the physics of the problem, nevertheless their spectral content is assumed to be dictated by the transmit signal. As a result, the autocorrelation function  $\mathbf{T}_k(\tau)$  in (18) is dictated by the inverse Fourier transform of the baseband version of the wideband signal power spectrum.

Extending the summation over all scatterers in (16), the array output covariance matrix in the fully wideband case takes the form

$$\mathbf{R}_z = E \{ \mathbf{z}(t) \mathbf{z}^*(t) \} = \sum_k P_k \mathbf{T}_k \circ \mathbf{s}_k \mathbf{s}_k^* + \sigma^2 \mathbf{I} \quad (19)$$

where  $\mathbf{T}_k = \mathbf{T}_k(0)$ . For example, a low pass power spectrum for the transmit signal gives

$$\mathbf{T}_k(m, n) = \text{sinc} B_o ((n_1 - n_2)\tau_2(k) + (i_1 - i_2)\tau_1(k)) \quad (20)$$

where  $n_1, n_2, i_1$  and  $i_2$  satisfy

$$m = i_1 N + n_1, \quad n = i_2 N + n_2. \quad (21)$$

In the special case, such as when the azimuth spread generates only a small angular dispersion  $\Delta\theta$  due to a strong main beam pattern, the  $\mathbf{T}_k$  in (19) may be replaced by a positive definite matrix  $\mathbf{T}$  that is valid for all  $k$ . In that case (19) reduces to

$$\mathbf{R}_z = \mathbf{T} \circ \mathbf{R}_x \quad (22)$$

where  $\mathbf{R}_x$  represents the narrowband situation in (6). It is well known that even for moderate bandwidth the Schur-Hadamard operation in (22) results in spreading the eigenvalue spectrum of  $\mathbf{R}_x$  thereby increasing the clutter degrees of freedom [1]. This is shown in Fig. 2 for a 14 sensor 16 pulse array with  $CNR = 40\text{dB}$  and  $\mathbf{T}_k$  as in (20).

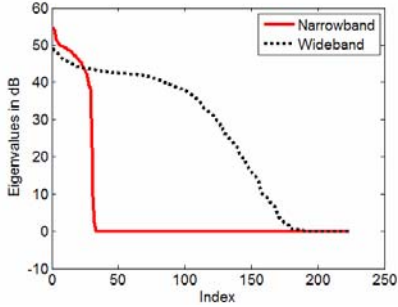


Fig. 2 Clutter eigen spectra for narrowband and wideband cases.

### III. WIDEBAND ANALYSIS

To fully understand the implication of the Schur-Hadamard product operation in (16) - (19), let us consider a simplified version of (19), where the corresponding narrowband scatter set corresponds to a single scatter from  $\theta_o$ ; i.e.,

$$\mathbf{R}_x = P_o \mathbf{s}(\theta_o, \omega_{d_o}) \mathbf{s}^*(\theta_o, \omega_{d_o}) \quad (23)$$

and hence from (19)

$$\mathbf{R}_z = \mathbf{T}_o \circ (P_o \mathbf{s}(\theta_o, \omega_{d_o}) \mathbf{s}^*(\theta_o, \omega_{d_o})). \quad (24)$$

To make further progress, let

$$\mathbf{T}_o = \sum_{i=1}^{MN} \lambda_i \mathbf{e}_i \mathbf{e}_i^* \geq 0 \quad (25)$$

represent the eigen decomposition of  $\mathbf{T}_o$  in (18). Substituting (25) into (24) we get

$$\mathbf{R}_z = \sum_{k=1}^{MN} P_k \tilde{\mathbf{s}}_k \tilde{\mathbf{s}}_k^* + \sigma^2 \mathbf{I} \quad (26)$$

where

$$P_k = P_o \lambda_k > 0 \quad (27)$$

and

$$\tilde{\mathbf{s}}_k = \mathbf{e}_k \circ \mathbf{s}(\theta_o, \omega_{d_o}). \quad (28)$$

Observe that (26) represents a bunch of  $MN$  uncorrelated returns, all of them associated with the single scatter located along  $\theta_o$ . Interestingly,  $\tilde{\mathbf{s}}_k$  in (26) represents a spatio-temporally amplitude modulated steering vector associated with the Doppler frequency  $\omega_{d_o}$  and location  $\theta_o$ .

Let  $\mathbf{d}_k$  represent the ordinary DFT vector associated with the eigenvector  $\mathbf{e}_k$ . The entries in  $\mathbf{e}_k$  correspond to a double sampling period of  $\tau_1$  followed by  $\tau_2$ . As a result, it can be shown that

$$\mathbf{e}_k = \sum_{n=1}^{MN} d_k(n) \left( \underline{b} \left( \frac{2j_n}{M} \right) \otimes \underline{a} \left( \frac{i_n c_o}{N} \right) \right) \quad (29)$$

where  $d_k(n)$  represents the  $n^{\text{th}}$  entry of  $\mathbf{d}_k$  and  $c_o$  is a normalization constant. Here  $i_n$  and  $j_n$  represent the unique solution to the equation  $n = (j_n - 1)N + i_n$ ,  $j_n = 1, 2, \dots, M$ ,  $i_n = 1, 2, \dots, N$ .

Using (29) in (28) we get the modified steering vectors in (26) - (28) to be

$$\begin{aligned} \tilde{\mathbf{s}}_k &= \sum_{n=1}^{MN} d_k(n) \left\{ \underline{b} \left( \frac{2j_n}{M} \right) \otimes \underline{a} \left( \frac{i_n c_o}{N} \right) \right\} \circ \left( \underline{b}(\omega_{d_o}) \otimes \underline{a}(\theta_o) \right) \\ &= \sum_{n=1}^{MN} d_k(n) \left\{ \underline{b} \left( \omega_{d_o} + \frac{2j_n}{M} \right) \otimes \underline{a} \left( \theta_o + \frac{i_n c_o}{N} \right) \right\} \\ &= \sum_{n=1}^{MN} d_k(n) \mathbf{s} \left( \theta_o + \frac{i_n c_o}{N}, \omega_{d_o} + \frac{2j_n}{M} \right) = \sum_{n=1}^{MN} d_k(n) \mathbf{s}(\theta_n, \omega_{d_n}) \end{aligned} \quad (30)$$

where

$$\theta_n = \theta_o + \frac{i_n c_o}{N}, \quad \omega_{d_n} = \omega_{d_o} + \frac{2j_n}{M}. \quad (31)$$

From (30) - (31), the amplitude modulated steering vector  $\tilde{\mathbf{s}}_k$  in (26)-(28) represents a coherent sum of steering vectors originating from locations  $\theta_n, n=1 \rightarrow MN$  with Doppler components  $\omega_{d_n}, n=1 \rightarrow MN$  as in (31). Referring back to (26), for each scatter return, the effect of finite bandwidth is

to introduce an apparent jittering effect both on angle and Doppler domains so that a bunch of uncorrelated returns are generated. Each such uncorrelated return contains a set of coherent returns with different directional and Doppler components as in (31). This is illustrated in Fig. 3. Observe that the vector  $\tilde{\mathbf{s}}_k$  in (30) contains several coherent returns, and it *does not* correspond to any physical steering vector. As a result, the adaptive processor based on (26) will not be able to null out either the original return from location  $\theta_o$  or any of the coherently combined vectors. From (26), the wideband data  $\mathbf{z}(t)$  in (16) can be equivalently expressed as

$$\mathbf{z}(t) = \sum_{k=1}^{MN} \alpha_k \sum_{n=1}^{MN} d_i(n) \mathbf{s}(\theta_n, \omega_{d_n}) = \sum_{k=1}^{MN} \alpha_k \tilde{\mathbf{s}}_k \quad (32)$$

and it represents  $MN$  uncorrelated bunches, each bunch containing  $MN$  coherent returns from various jittered locations around with  $\theta_o$ .

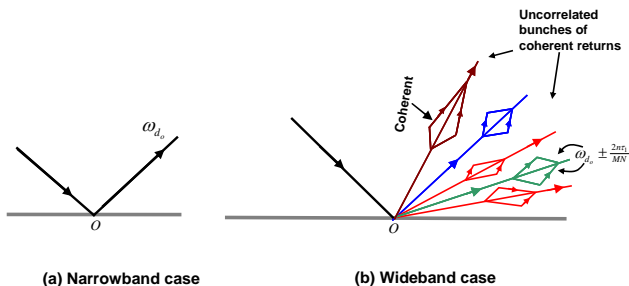


Fig. 3 Narrowband returns and effective wideband returns.

Fig. 4 shows the SINR loss of the adaptive processing with and without bandwidth dispersion. From there, the clutter null is wider due to the wideband operation nature of problem.

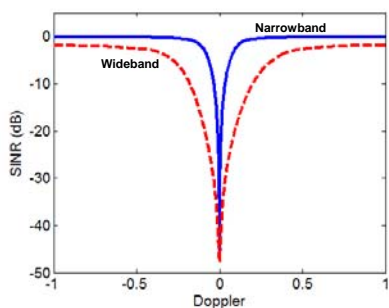


Fig. 4 Effect of wideband on clutter nulling performance.

### III. SUMMARY AND AREAS FOR FUTURE INVESTIGATION

In summary, when traditional processing is employed in the wideband case a single scatterer appears to generate several isolated coherent bundles, thereby making clutter cancellation impossible. Thus, alternate processing schemes must be considered to fully resolve the wideband situation.

Presently, there are only two basic classes of techniques that have been proposed to remedy the wideband STAP clutter cancellation problem: Subband based methods or 3D STAP [3]-[4]. The subbanding approach, as the name implies, attempts to transform the wideband processing problem into a parallel set of narrowband 2D (angle-Doppler) STAP filters. Since a single 2D STAP processing thread is already a quite substantial processing burden, a bank of say  $L$  subband processors would represent an  $L \times 100\%$  increase in processing burden!

The situation for 3D STAP is even worse. “3D” here refers to including the fast-time (i.e., reciprocal of receiver bandwidth) or equivalently instantaneous frequency domain. In other words, the dimensionality of the STAP problem is increased from  $NM$  to  $NML$ , where  $N$  is the number of spatial degrees-of-freedom (DOFs),  $M$  is the number of so-called “slow-time” or Doppler DOFs [1], and  $L$  is the number of fast-time or instantaneous DOFs. Since processing burdens are roughly on the order of  $O(k^2) \rightarrow O(k^3)$ , 3D STAP can result in an  $L^2 \rightarrow L^3$ -fold increase!

The above processing burden increases are in addition to the many other practical difficulties associated with implementing such adaptive filtering schemes, such as sufficient sample support for weights calculation [1] – a problem which is already acute for 2D STAP.

A practical framework for addressing the above issues is hierarchical processing in which an “energy preserving” lower resolution processing stage, to which 2D STAP could be applied, is followed by wideband processing only if a detection was obtained in the first stage. Even for the most stressing MTI applications, detections represent a very small fraction of all possible resolution cells—thus resulting in a substantially reduced wideband processing burden.

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