

TASK REPORT 97-03

Second Report for Research and Modeling of Water Particles in Adverse Weather Simulation Facilities

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1.0 INTRODUCTION TO THE RESEARCH

This report describes a continuation of research into the modeling of water particle freezing for application to adverse weather simulation facilities. The research was initiated in FY1996 to investigate the physics of freezing of submillimeter supercooled water particles or droplets in both natural and artificial or simulated adverse weather environments. The first phase of the research was reported and discussed in a report [1] and a paper [2]. The work has continued into FY1997 and has been expanded to include work done to model three-dimensional ice accretions on surfaces, as well as modeling the near field of water spray clouds produced by air-atomized water spray nozzles. Because of the increased scope of the work, a single report cannot cover all of the work phases. Therefore, the present report covers only the continued research and development of water particulate freezing models and their application in a one-dimensional multiphase flow code to predict water spray freezing in ducted air flows.

The outline of the report is as follows:

Section 2 provides an overview of applications of classical homogeneous nucleation theory to predict supercooled water particle freezing. Specifically, it covers and discusses the J(T) and I(T) functions (which are identical except for nomenclature used by various authors) that predict the rate-of-formation of freezing centers or ice-germ nuclei that initiated the freezing process.

Section 3 describes how J(T) and/or I(T) were evaluated in the present study and compared to their evaluations in the literature.

Section 4 describes how J(T) and/or I(T) were modified to apply to the heterogeneous freezing of nominal purity supercooled water droplets.

Section 5 then provides calculated results from incorporating the modified heterogeneous droplet freezing model into the AEDC 1DMP code and evaluating some of the same flow cases

as described in references [1,2] for a ten-droplet water spray. The results from a previously developed simple heterogeneous freezing model and the new modified homogeneous nucleation freezing function model are compared and evaluated.

Section 6 then summarizes the current status of research and development under this phase of the research task investigation, provides conclusions about accomplished and needed future research, and makes recommendations about how to pursue future development of adverse weather simulation test capabilities in ground test facilities.

2.0 OVERVIEW OF CLASSICAL HOMOGENEOUS NUCLEATION THEORY

2.1 Nucleation Probability Functions

This section very briefly reviews the homogeneous nucleation probability functions of the literature and describes the selection process used to select one of the functions for application to development of a heterogeneous freezing model for supercooled water droplets.

In this context, homogeneous nucleation refers to the creation of a condensed phase of a pure substance by the thermodynamically controlled formation of condensed phase precursors. These precursors are metastable clusters of molecules that potentially can serve to trigger the formation of the condensed phase, that is, act as condensation sites. Specifically, the precursors can be thought of as predroplets, or nandroplets, when a liquid condenses from a supersaturated vapor, or as ice germs when ice crystals form within a liquid. Heterogeneous nucleation follows the same thermodynamic process of molecular cluster formation as occurs in homogeneous nucleation, except that some foreign material, such as a particle or a surface, reduces the energy required to form the precursors. Thus, heterogeneous nucleation occurs at higher phase temperatures, in general, than homogeneous nucleation.

The literature on nucleation and condensation processes, that start with the formation of metastable molecular clusters in the non-condensed phase of a substance, is reasonably large in both papers [3-38], and books [39-43]. The "condensed" phase may be a liquid droplet condensing from a (non-condensed) supersaturated vapor, or it may represent a proto-crystal or ice germ, in the case where pure supercooled liquid water is beginning its spontaneous freezing process, at some temperature below 273.16 K. The freezing process that begins to occur in sprays of supercooled liquid metal droplets is essentially the same process as that in which pure supercooled water spray droplets freeze, except that liquid metals can be alloys and can freeze in different solid phases, depending on drop size, composition, and temperature at the time of solidification nucleation. However, pure liquid water droplets, supercooled and freezing, turn into ice I-type solid crystalline particles under normal "atmospheric" conditions.

As stated, the purpose of this section of the report is to briefly overview the probability functions, denoted in the literature as J(T) or I(T), which predict the production rate of nucleation sites per unit volume per unit time. In the present research effort, the theory for homogeneous nucleation of "pure" liquid water droplets, or homogeneous freezing, has been modified to account for heterogeneous (or mote-induced) freezing of water droplets. Heterogeneous freezing is the usual, typical, or most probable freezing process for liquid water of average purity, such as either unfiltered, or even filtered, and distilled drinking water. The modified homogeneous nucleation freezing theory can be applied in a straight-forward way to general numerical spray computation codes (numerical spray CFD codes) that track the position, lifetime, kinetic, and thermal states of particles or particle group packets that are injected and convected in gas flows. Moreover, the modified theory can incorporate probabilistic functions that extend the theory to stochastic processes of freezing. To explain the modified theory, it is necessary to review or overview classical homogeneous nucleation theory for particulate freezing. The theory begins in

Germany with the work of Volmer, et al. [3], Becker, et al. [4], and others with further development by Turnbull, and co-workers, e.g. [7,8]. Its exposition, however, is begun with Mason [14], for convenience.

2.2 <u>Mason (1952)</u>

In his paper, Mason [14] provides J(T) in the form

$$J(T) = \frac{nkT}{h} \exp\left\{-\frac{U}{kT} - \frac{W_c}{kT}\right\}$$
(1)

Mason then evaluates $\log J(T)^{1}$ as

$$\log J(T) = 32.84 + \log T - \frac{U}{2.30kT} - \frac{760\sigma_{sL}^3}{(T_o - T)^2 T}$$
(2)

In these expressions, the parameters are:

 $n \equiv$ number of molecules per cubic centimeter of the condensed phase,

 $k \equiv \text{Boltzmann's constant},$

 $h \equiv$ Planck's constant,

 $T \equiv$ the absolute temperature, Kelvins,

 $T_o \equiv$ the bulk water freezing temperature, Kelvins (273.16K)

 $U \equiv$ the "activation energy for self-diffusion of a molecule in the" non-condensed phase,

 $W_c \equiv$ "the work of nucleus formation".

Mason provides W_c as

$$W_c = \frac{1}{3}\sigma_{sL} A = \frac{1}{3}\sigma_{sL}\omega r_c^2$$
(3)

3) '

attributing this to Frenkel [39]. These parameters are:

¹ Footnote: ln () refers to the natural log, whereas log () refers to log base 10.

- $\sigma_{sL} \equiv$ the "specific surface energy of the crystal-liquid interface".
- $A \equiv$ the surface area of the nucleus (the molecular cluster, or ice germ precursor).
- $\omega \equiv$ shape factor for nucleus, relating the surface to volume ratio for the shape of the cluster (ω is approximately 21 23 in value).
- $r_c \equiv$ the radius of a sphere inscribed in a nucleus of critical size, i.e., that sized nucleus that can form a nucleation or crystallization site.

Mason evaluates ω by assuming that the nucleus forms a hexagonal prism with height equal to twice the radius to a prismatic side (see sketch below):







(7)

Top or Bottom View



The volume of this hexagonal prism is $V = 4\sqrt{3} r_c^3 = 6.9282 r_c^3$ (4)

The surface area is
$$A = 12\sqrt{3} r_c^2 = \omega r_c^2 = 20.785 r_c^2$$
 (5)

The surface to volume ratio is
$$\frac{A}{V} = \frac{12\sqrt{3} r_c^2}{4\sqrt{3} r_c^3} = \frac{3}{r_c}$$
 (6)

Note also that A can be written in terms of V as $A = 3^{\frac{7}{6}} 4^{\frac{1}{3}} V^{\frac{2}{3}}$ So that the derivative dA/dV is

$$\frac{dA}{dV} = 3^{\frac{7}{6}} 4^{\frac{1}{3}} \left(\frac{2}{3}\right) V^{-\frac{1}{3}} = 2 \cdot 3^{\frac{1}{6}} \cdot 4^{\frac{1}{3}} V^{-\frac{1}{3}}$$
(8)

Substituting for V

$$\frac{dA}{dV} = 2 \cdot 3^{\frac{1}{6}} \cdot 4^{\frac{1}{3}} / \left(4\sqrt{3} r_c^3\right)^{\frac{1}{3}}$$
(9)

or

$$\frac{dA}{dV} = \frac{2}{3^{\frac{1}{6}}} \frac{3^{\frac{1}{6}}}{r_c} = \frac{2}{r_c}$$
(10)

This parameter, dA/dV, is important because Mason, in a later paper [25(1960)], shows that the change in Gibbs Free Energy for a collection containing g molecules that have passed from the liquid phase to form a cluster in the solid phase is

$$\Delta G = (\mu_s - \mu_L)g + A\sigma_{sL} \tag{11}$$

where

 $\mu_{L} \equiv \text{Gibbs Free Energy of molecule in liquid phase,}$

 $\mu_s \equiv$ Gibbs Free Energy of molecule in an ice-like cluster,

 $g \equiv$ number of molecules in the cluster.

Mason [25(1960)] defines the critical cluster of molecules as the metastable cluster reached when ΔG achieves a maximum, thus the two conditions

$$\frac{d\Delta G}{dg} = 0 \tag{12}$$

and

$$\Delta G = \Delta G_{\max} \tag{13}$$

define the cluster size when a metastable ice germ (cluster) has formed.

Therefore, for this condition,

$$\frac{d\Delta G}{dg} = 0 = (\mu_s - \mu_L) + \frac{dA}{dg}\sigma_{sL}.$$
(14)

But,

and

where $\lambda \equiv$ the effective volume, per molecule, in the cluster.

Thus,

$$\frac{dA}{dg} = \frac{dA}{dV} \lambda \; .$$

 $\frac{dA}{dg} = \frac{dA}{dV} \frac{dV}{dg}$

 $V = \lambda g$

An evaluation of λ gives

 $\lambda = \frac{W_{mol}}{N_A \rho_s}$

where

 $W_{mol} \equiv$ the molecular weight of ice,

 $N_A \equiv$ Avogadro's number,

 $\rho_s \equiv$ the density of ice.

Thus,

$$\frac{dA}{dg} = \frac{dA}{dV} \frac{W_{mol}}{N_A \rho_s}.$$
(19)

Consequently, the condition for metastability of the cluster is

$$(\mu_L - \mu_S) = \frac{dA}{dg} \sigma_{SL} = \frac{\sigma_{SL} W_{mol}}{N_A \rho_S} \frac{dA}{dV}.$$
(20)

(17)

(18)

(15)

(16)

However, Mason, e.g. [1960], shows that

$$\frac{d}{dT}(\mu_{L} - \mu_{S}) = -(S_{L} - S_{S}) = \frac{L_{F}}{T}$$
(21)

where

 $L_F^* \equiv$ the heat of fusion, per molecule (treated as a negative quantity),

 $S_L \equiv$ the entropy of a molecule in the liquid phase,

 $S_s \equiv$ the entropy of a molecule in the ice-like cluster or solid phase.

Thus, by integration

$$(\mu_{L} - \mu_{S}) = \int_{T_{a}}^{T} \frac{L_{F}^{*} dT}{T}$$
(22)

and, by substitution, therefore, based on equation (20),

$$(\mu_L - \mu_S) = \frac{\sigma_{SL} W_{mol}}{N_A \rho_S} \frac{dA}{dV} = \int_{T_a}^T \frac{L_F^* dT}{T}.$$
(23)

Substituting for *dA/dV*

$$\frac{\sigma_{sL}W_{mol}}{N_A\rho_s}\frac{2}{r_c} = \int_{T_o}^T \frac{L_F^* dT}{T}.$$
(24)

Solving for r_c , and rearranging terms gives r_c as

$$r_{c} = \left(\frac{2\sigma_{SL}}{\rho_{s}}\right) / \int_{T_{o}}^{T} \frac{N_{A}L_{F}^{*}}{W_{mol}T} dT .$$
(25)

But, the heat of fusion per unit mass of the liquid is

$$L_F = \frac{N_A}{W_{mal}} L_F^*.$$
 (26)

Therefore, the equation for the critical radius of the metastable molecular cluster is

$$r_c = \frac{2\sigma_{SL}}{\rho_s \int\limits_{T_c}^{T} \frac{L_F dT}{T}}.$$

The work of formation of the critical metastable cluster which was given by equation (3)

$$W_c = \frac{1}{3}\sigma_{sL}A = \frac{1}{3}\sigma_{sL}\omega r_c^2$$
⁽²⁸⁾

can now be written substituting for r_c , as

Therefore, W_c can be written as

For the hexagonal prism defined by Mason, we had

$$A = 12\sqrt{3} r_c^2 = \omega r_c^2$$
 (5)

hence,

However, in his 1952 paper [14], Mason simply introduces an equation for r_c that Mason calls "Thompson's equation," without reference, as

$$r_c = \frac{2\sigma_{SL}}{\rho_s L_F} \frac{T_o}{T_o - T}.$$
(32)

Mason later provides a derivation of this equation in his 1957 text, The Physics of Clouds, [44]. Clearly, equation (32) is also obtainable from equation (27) when L_F is treated as a constant, by expanding and linearizing the resulting log function. In the previous equation, ρ_s is the density

$$W_{c} = \frac{4}{3}\omega \frac{\sigma_{SL}^{3}}{\rho_{s}^{2} \left[\int_{T_{o}}^{T} \frac{L_{F}dT}{T}\right]^{2}}$$
(29)

$$\omega = 12\sqrt{3}.$$
 (30)

$$W_{c} = \frac{16\sqrt{3} \sigma_{SL}^{3}}{\rho_{s}^{2} \left[\int_{T_{c}}^{T} \frac{L_{F} dT}{T}\right]^{2}}.$$
(31)

(27)

of the condensed phase, here assumed to be Ice-I, normal ice. By using equations (3), (30), and (32) in equation (1), we can write the equation for J(T) as

$$J(T) = \frac{nkT}{h} \exp\left\{-\frac{U}{kT} - \frac{\binom{1}{3}\sigma_{sL}\omega}{kT} \left(\frac{2\sigma_{sL}}{\rho_s L_F}\right)^2 \left(\frac{T_o}{T_o - T}\right)^2\right\}.$$
(33)

By using $T_o = 273.16$ K, and simplifying, we get

$$J(T) = \frac{nkT}{h} \exp\left\{-\frac{U}{kT} - \frac{4}{3} \frac{\sigma_{sL}^{3} \omega}{kT \rho_{s}^{2} L_{F}^{2}} \left(\frac{273.16}{273.16-T}\right)^{2}\right\}.$$
 (34)

Equation (2) is equation (34), after some of the parameters have been numerically evaluated. Mason provides values, in his 1952 paper, for the following parameters in J(T), as follows.

The interfacial surface tension, σ_{sL} , was evaluated by Mason as the difference between the surface tension of ice at -40°C and of liquid water at -40°C, that is,

$$\sigma_{sL} = \sigma_s - \sigma_L = 102 \frac{\text{erg}}{\text{cm}^2} - 80 \frac{\text{erg}}{\text{cm}^2}$$

or

$$\sigma_{sL} = 22 \frac{\text{erg}}{\text{cm}^2} = 0.022 \frac{\text{J}}{\text{m}^2}.$$

 σ_s was estimated from breaking hydrogen bonds normal to a selected face of an ice crystal.

The value for U was estimated at $U = 3.3 \times 10^{-13}$ erg, or $3.3 \times 10^{-20} J$, for $0 \le T \le -10^{\circ}$ C, and this same value taken at $T = -40^{\circ}$ C. The density of ice and its latent heat of fusion were taken as

$$\rho_{\rm c} = 0.92 \text{ gm/cm}^2 = 920 \text{ kg/m}^3$$

and

$$L_F = -3.33 \ x \ 10^9 \ \frac{\text{erg}}{\text{gm}} = -3.33 \ x \ 10^5 \ \frac{\text{J}}{\text{kg}}$$

The shape factor, ω , was taken by Mason to be $\omega = 23$. (However, it was computed in the present study as about 21 instead.)

The next parameter is the number n of molecules of liquid water per unit volume. This parameter is given by $n = \frac{\rho_L \cdot N_A}{W_{max}}$ where ρ_L is the liquid density, N_A is Avogadro's number, and

 W_{mol} is the molecular weight of water. Thus, for

$$\rho_L = 1 \frac{\text{gm}}{\text{cm}^3},$$

$$N_A = 6.022169 \text{ x } 10^{23} \frac{\text{molecules}}{\text{gm} - \text{mole}},$$

$$W_{\text{mol}} = 18 \text{ gm/gm} - \text{mole},$$

n is computed as $n = 3.3456 \times 10^{22} \frac{\text{molecules}}{\text{cm}^3} = 3.3456 \times 10^{28} \frac{\text{molecules}}{\text{m}^3}$.

The remaining parameters are Boltzmann's constant, k, and Planck's constant, h, given by

$$k = 1.380622 \ x \ 10^{-16} \ \frac{\text{erg}}{\text{K}} = 1.380622 \ x \ 10^{-23} \ \frac{\text{J}}{\text{K}}$$

and

$$h = 6.626196 \times 10^{-27} erg \cdot s = 6.626196 \times 10^{-34} J \cdot s$$
.

Evaluation of $\log J(T)$

Using these values for the parameters in the J(T) equation Mason evaluated log J(T) as $\log J(T) = \log(3.3456 \ x \ 10^{22}) + \log(1.380622 \ x \ 10^{-16}) + \log T - \log(6.626196 \ x \ 10^{-27})$ (35) $\frac{U}{2.303 kT} = \frac{4}{3} \cdot \frac{1}{2.303} \cdot \frac{\sigma_{SL}^3 (23)(273.16)^2}{(1.380622 x 10^{-16}) T (0.92)^2 (3.33 x 10^9)^2} \cdot \frac{1}{(T_a - T)^2}.$

Evaluating this, we get

$$\log J(T) = 32.843 + \log T - \frac{U}{2.303kT} - \frac{766.8\sigma_{sL}^3}{T(T_g - T)^2}$$
(36)

where σ_{sL} is in erg/cm² units and J(T) would be in units of nuclei per cubic centimeter per second. With minor numerical differences, this is Mason's equation, equation (2).

In the present study, using the values of the parameters as given in SI units, and using natural logarithms, this equation has been put into the form

$$\ln J(T) = 89.440 + \ln T - \frac{2390.2}{T} - \frac{252}{T} \left(\frac{273.16}{273.16 - T}\right)^2.$$
 (37)

Alternately, this equation gives J(T) as

$$J(T) = 6.9708 \ x \ 10^{38} T \exp\left\{-\frac{2390.2}{T} - \frac{252}{T} \left(\frac{273.16}{273.16-T}\right)^2\right\}$$
(38)

with units of nuclei per meter cubed per second.

In the following subsections of the report, alternate forms for J(T) are described which are given in the literature. J(T) has been also given the symbol I(T). In this report, J(T) and I(T)are exactly the same functions.

2.3 McDonald (1953)

In this paper, McDonald [18] provides a critical review of homogeneous nucleation theory. In particular, with reference to the application of the theory to predict the spontaneous freezing of supercooled submillimeter water droplets by homogeneous nucleation, McDonald reviews and assesses the values of the physical parameters in the argument of the function, J(T). McDonald writes the J(T) function as (with a sign correction).

$$J(T) \approx \frac{nkT}{h} \exp\left\{-\left(\frac{A+Fc}{kT}\right)\right\}$$
(39)

A is the activation energy for self-diffusion of liquid molecules near the water-ice germ interface. F_c is the work of formation of the ice germ or molecular cluster. *n* is the number of molecules per unit volume in the liquid phase, *k* is Boltzmann's constant, and *h* is Planck's constant.

McDonald gives F_c as

$$F_c = \sigma_s g r_c^2 / 3 \tag{40}$$

where r_c is the critical radius of the molecular cluster, given by

$$r_c = \frac{2\sigma_s T_o}{\left(\rho_s L_F (T_o - T)\right)}.$$
(41)

In these equations, σ_s is the specific surface free energy of the water-ice germ interface, g is a geometric factor such that $g r_c^2$ is the total surface area of the ice germ or critical embryo, ρ_s is the density of ice, L_F is the latent heat of fusion of ice, and T_o is the melting temperature of ice, 273.16 K. Thus, the expression for J(T) is essentially the same as that presented earlier by Mason and co-workers.

The main thrust of McDonald's investigation is an examination of the parameters A, σ_s , and L_r . McDonald reviews the work of others, including Mason [14], in an attempt to obtain thermodynamically correct values of these parameters, in terms of the best theoretical model of an ice germ structure. McDonald notes that L_r , the heat of fusion of ice is not constant, but decreases with decreasing temperature. (This was discussed by the present author in previous work [1, 2].) McDonald also finds both A and σ_s to vary with temperature.

McDonald's rigorous analysis is well worth examination by those with expertise in molecular physics because the levels of uncertainty in the estimations or calculations of A, σ_s , and L_r , that McDonald demonstrates, implies that serious further theoretical and experimental

research is required to put homogeneous nucleation theory on a rigorous basis. This is one of McDonald's main conclusions in 1953 and it appears to be valid even at present. In actual practice, therefore, the application of the freezing function J(T), by Mason and others, is a correlative theory, rather than a predictive theory, due to the uncertainties in the physical parameters of the theory.

2.4 Day (1958)

J. A. Day [21] presented a functional form for J(T) that he attributed to J. B. McDonald [18]. This equation is

$$J(T) = \frac{nkT}{h} \exp\left\{-\frac{A}{kT} - \frac{30.7\sigma_{SL}^{3}}{k\,\rho_{S}^{2}\,L_{F}^{2}\,T\left(\ln\frac{T_{o}}{T}\right)^{2}}\right\}$$
(42)

where T_o and the other parameters have been defined in the previous section. Obviously, this equation can be developed from equation (1) with W_c evaluated from equation (31) and with L_F treated as constant.

2.5 <u>Mason (1958)</u>

Mason presented another version of the J(T) equation in this 1958 paper [22]. In this paper, Mason uses I(T) to represent the nuclei density rate equation. The equation provided was

$$\log I(T) = 32.84 + \log T - \frac{U}{2.303kT} - \frac{1.11 \times 10^{17} \sigma_{SL}^3}{L_F^2 T \left(\ln \frac{T_o}{T} \right)^2}.$$
 (43)

The same comments apply here to this equation as for the equation Day presented.

2.6 Langham and Mason (1958)

In 1958, Langham and Mason reviewed the heterogeneous and homogeneous nucleation of supercooled water [23] and presented a functional equation for I(T) that is the same as that of Mason (1958) [22].

2.7 Mason (1960)

In a 1960 paper, Mason [25] applied nucleation theory to predict both the formation of water aerosol droplets from supersaturated water vapor, and also, homogeneous nucleation or freezing of supercooled water droplets. Mason also reviews the formation of water droplets by condensation of vapor on foreign nuclei such as ions, hygroscopic and non-hygroscopic motes, and on "mixed" nuclei.

Mason attributes the theoretical development of condensation theory to Becker and Doring [4], Zeldovich [45], and Turnbull and Fisher [7]. Mason discusses the validity of the theory, as well as the validity of the experiments that provided data on droplet condensation and freezing.

Mason discusses the freezing of submillimeter supercooled droplets by both heterogeneous and homogeneous nucleation processes and the applicable theories for both. Mason also reviews the effectiveness or activity of various mineral substances to act as ice nucleating motes. The most active substance Mason found was silver iodide which Mason says produces one ice crystal per 10000 AgI particles. This statement has relevant implications, therefore, in the interpretation of J(T). It implies that the population number densities for nucleation, as correlated by J(T) (or I(T)) must reflect only the active or effective nuclei, or the active/effective motes, and not the actual or physical nuclei/mote population number densities or their rates of creation. Especially this will be true if J(T) is used to predict heterogeneous freezing by correlating the J(T) function (actually σ_{st}) to a given set of experimental data.

A derivation of I(T) was provided by Mason in this paper [25] that leads to the equation

$$\log I(T) = 32.84 + \log T - \frac{U}{2.303kT} - \frac{16\sqrt{3}\sigma_{SL}^3}{2.303\rho_S^2 kT \left[\int_T^{T_s} \frac{L_F}{T} dT\right]^2}.$$
 (44)

It is clear that this equation can be derived from the approach described in section 2.2 of the present study. In discussing this equation, Mason notes that σ_{st} probably cannot be calculated with sufficient accuracy to permit using I(T) to predict freezing nucleation rates. However, he also notes, that by using experimental data for the freezing of pure supercooled water droplets, σ_{st} values can be obtained by correlating I(T) with appropriate sized drops freezing at measured temperatures. If the presence of motes in ordinary water acts to change the σ_{st} values, then it should also be possible to correlate I(T) to the heterogeneous freezing of supercooled droplets and obtain σ_{st} versus drop size for each set of experimental data. This has been done in the present research investigation to obtain a J(T) function applicable to heterogeneous nucleation freezing of supercooled water droplets and will be presented in section 4.

2.8 Fletcher (1960)

In this paper [24], Fletcher writes the nucleation site population density function as

$$J(T) = K \exp\{-\Delta G^* / kT\}$$
(45)

where "K is a kinetic constant typically of order 10^{25} cm⁻³ sec⁻¹ for the cases which we shall consider," k is Boltzmann's constant, and T is absolute temperature, as utilized previously. Fletcher discusses the extension or application of nucleation/condensation theory to account for the effect of foreign surfaces or particles in modifying (reducing) the value of the work of formation of the molecular cluster or embryo nuclei, denoted ΔG^* . Thus, if a theoretical model can be developed of how substances, either as motes or surfaces, can change the value of ΔG^* , then this model can be incorporated into J(T) to create a heterogeneous nucleation site population rate equation. Theoretically, therefore, as Becker, Turnbull, Mason, and others have also discussed, heterogeneous nucleation freezing of supercooled water droplets can be predicted with the same theoretical tools used to predict homogeneous freezing of submillimeter supercooled water droplets.

Fletcher discusses the effects of foreign surface geometry and size on ΔG^* for heterogeneous nucleation or the formation of the ice germ or embryo. He presents a function, f(m,x), that modifies the free energy of formation for homogeneous nucleation, i.e., Fletcher gives

$$\Delta G^* = \Delta G_o^* f(m, x) \tag{46}$$

where ΔG_o^* is the free energy of formation of a nucleus of critical radius under homogeneous nucleation conditions, such as given by equation (31) and where

$$m = \cos\theta \,. \tag{47}$$

Cosine θ is the cosine of the contact angle between ice germ (embryo) and the foreign particle and x is defined as

$$x = R/r^*. \tag{48}$$

R is the radius of a spherical foreign particle and r^* is the radius of curvature of the surface of the critical embryo.

Fletcher [24] provides distributions of f(m,x) in a figure (Fig. 1 of [24]) obtained from solving the "nucleation equations for a spherical cap embryo nucleating upon a perfect spherical particle." Details of this model are provided by Fletcher in another paper.

Fletcher shows that this theoretical model for accounting for the effects of foreign particles on the freezing of supercooled water predicts the correct trends when applied to silver iodide particles inducing higher spontaneous freezing temperatures in supercooled droplets of water. Thus, the use of modified homogeneous freezing theory to predict heterogeneous freezing can be justified by more than one theoretical approach.

2.9 Kuhns and Mason (1968)

This paper [29] continues the work of Mason on the assessment and theoretical evaluation of the freezing of submillimeter supercooled droplets of pure water. Kuhns and Mason review and describe the work of Mason (1957, 1958, 1960) and others who have studied the freezing of small water droplets. Kuhns and Mason report experimental data they obtained on freezing of freely falling water drops, noting that much of the existing data to 1968 consisted of freezing droplets that were supported or in contact with surfaces, wires, hypodermic needles, etc. Much of this paper describes the apparatus used by Kuhns and Mason to perform their experiments and to analyze the freely falling droplets to estimate their temperature time history, as well as, size. They also investigated the effect of different ambient gases surrounding the particles on the particle freezing process.

These authors review, and apparently expand, the heterogeneous freezing theory of Bigg [16], as well as, the homogeneous theory developed by Mason (1952, 1958, 1960) from the Turnbull-Fisher work [7]. They report the same J(T) function that Mason presented in his 1960 paper (see section 2.7 of the present report). Based on the probability theory of Bigg [16], and using the concept of a "median" freezing temperature for a group of drops all of the same size, but freezing stochastically about this median temperature, the authors define the "median" freezing event as occurring when the product of $J(T_F)$, particle volume, V, and time, t, reach a critical value. That is, when the freezing probability, P_F , reaches $P_F = \frac{1}{2}$, half of the drops are assumed frozen, [16]. To illustrate this theory, let the probability of a drop remaining liquid be given by

$$P_L = 1 - e^{-J(T)Vt}$$

Then, the probability of freezing is

$$P_F = 1 - P_L = e^{-J(T)Vt}$$

Taking the logs of both sides

$$\ln P_F = \ln(1 - P_L) = -J(T)Vt$$

If the probability of freezing of half of the droplets in a given size group is set at $\frac{1}{2}$, then by definition, the drop size group has reached its "median freezing temperature, T_{r} . Hence,

$$-\ln P_F = J(T_F)Vt = -\ln(1 - P_L) = -\ln(\frac{1}{2}) \approx 0.7.$$
(49)

The present author utilizes equation (49) differently. Equation (49) can be interpreted or utilized in several ways. With σ_{sL} determined, equation (49) can be used to predict the freezing temperature, T_F , of different sized droplets. Alternately, if the "median" freezing temperature, T_F , of a given sample of drops of all the same size is known, then equation (49) can be used to evaluate a "median" value for σ_{sL} . However, the present author interprets equation (49) somewhat differently.

There are two reasons for the present author using a different interpretation to the product, JVt. From a physical interpretation, the number, N_c , of (active) freezing nuclei present in a single sample or droplet of supercooled water at time, t, must be given by the following integral

$$N_{c}(t) = \int_{t_{0}}^{t_{r}} J(T, D, t) V(D, t) dt$$
(50)

where t_r is the time at which the particle begins to freeze, and t_o is the time at which the particle first began to supercool (the time when its temperature reached 273.16K). In this integral, the function, J, which gives the population net rate-of-creation, has been generalized to be a function not only of particle temperature, T, but also of particle size, given by its diameter, D, and time, t. Thus, equation (50) can apply equally well to heterogeneous nucleation events in supercooled water.

With this interpretation, the minimum or critical value for N_c is unity, that is, for a single drop undergoing freezing,

$$N_{c,critical} \equiv 1.0 \tag{51}$$

Hence, by the present author's interpretation at least one active freezing site must come into existence in each sample to initiate freezing.

The second reason for this interpretation is that the freezing model must be implemented, in the present study, into a code wherein all droplets of the same size freeze at the same time. At present, the limitations of the code preclude modeling wherein half of the droplets of a given size freeze at the median freezing temperature and half remain liquid. Hence, when a defined measure of the product JVt reaches a critical value for the droplets of that size class, assumed in this study to be unity, all of these droplets begin to freeze. The measure of JVt used in the present study is an integral of the product JV dt, as defined below.

In the present study, a criterion or a measure of JVt for both homogeneous and heterogeneous nucleation freezing was adopted, given by the integral equality

$$\int_{t_{g}}^{t_{F}} J(T, D, t) V(D, t) dt = 1.0$$
(52)

When this integral reaches unity, all drops of size D are assumed to begin freezing. This criterion (equation (52)) was adopted and has been applied in the present research study to predict heterogeneous nucleation freezing of water spray droplets in ducted air flow, by being incorporated into a numerical model for one-dimensional, multiphase flow described in a previous report [1], and paper [2]. Details of this model are presented later in the present report.

Returning to a discussion of the Kuhns and Mason paper, Kuhns and Mason determine the value of σ_{sL} required in J(T) by requiring 20-micron drops to freeze at a median freezing temperature of -37C after a period of 1 second. This resulted in a σ_{sL} value of

$$\sigma_{sL} = 19.7 \frac{\text{erg}}{\text{cm}^2} = 0.0197 \frac{\text{J}}{\text{m}^2}$$
(53)

where σ_{sL} is the same free energy of the ice-water interface or surface as described previously. With this set of conditions, with σ_{sL} of 19.7 erg/cm², the other freezing temperatures of droplets of different sizes can be determined when

$$J(T_{F_{i}}) V_{i}(1 \sec) = 0.7$$
(54)

where

 $T_{F_i} \equiv$ spontaneous (median) freezing temperature of droplets of diameter, d_i

$$V_i \equiv \frac{\pi}{6} d_i^3$$
, the droplet volume for a drop of diameter, d_r

They presented the resulting T_{F_i} vs d_i curve and compared experimental data from various sources against it, including their own data. The results of the comparison were good.

Kuhns and Mason also estimate the size of the critical embryo, or nucleus or ice germ that reaches a metestable equilibrium, and thus has an opportunity to nucleate the freezing process. The critical radius of the cluster was given by the same equation as equation (27), namely

$$r_c = \frac{2\sigma_{SL}}{\rho_s \int_T^{T_c} \frac{L_f \, dT}{T}}$$
(55)

where, as before,

 $\rho \equiv$ the density of ice,

 $L_f \equiv$ the heat of fusion (or melting) of ice, (now, however, treated as positive). If L_f is constant, r_c is given by

$$r_{c} = \frac{2\sigma_{sL}}{\rho_{s}L_{f}\ln\left\{\frac{T_{o}}{T}\right\}}$$
(56)

By assuming different values of σ_{sl} , the corresponding values of the critical radius r_c were computed, and the number of water molecules that form the ice precursor or embryos were estimated. The authors obtained 150-300 water molecules depending on σ_{sl} in the temperature range investigated. The authors also showed, albeit with some approximation, that the results of a statistical thermodynamics model of water molecule aggregates, termed the "flickering cluster" model, explored by Nemethy and Scheraga [46], can be interpreted to predict the same number of molecules required for a meta-molecular cluster and for a σ_{sl} value of about 20 erg/cm². Thus, it is implied that the methods of statistical thermodynamics, applied to a special model of molecular aggregates, could be used to predict the spontaneous freezing of pure supercooled water.

2.10 Anderson, Miller, Kassner, Jr., and Hagen (1980)

This paper [33] reviews condensation-freezing nucleation of small water droplets in an expansion cloud chamber. The authors report on their experimental cloud chamber test results where submillimeter liquid droplets first nucleate from a supersaturated vapor phase then undergo spontaneous homogeneous (nucleation) freezing. The freezing process is reported to occur where chamber temperature is "near -40°C." The authors also note minor effects of an electric field applied to the cloud chamber which reduces the presence of ions on which droplets can nucleate and thus freeze.

The authors investigate whether the occurrence of frozen nuclei in cold supersaturated vapors is due to the direct formation of ice germs or ice nuclei by vapor to solid nucleation processes or whether it occurs due to a two-step process of nucleation of vapor first to form liquid droplets, then, second, for the liquid droplets to freeze by liquid to solid nucleation. An understanding of which of these processes can account for the presence of atmospheric ice particles would improve our general knowledge about the mechanisms of formation of arctic precipitation clouds.

The authors discuss the literature on water nucleation, including vapor to liquid, vapor to solid, and liquid to solid nucleation. They provide expressions for two nucleation site creation rate functions, $J_{v_L}(T)$ and $J_{v_S}(T)$. $J_{v_L}(T)$ is the homogeneous nucleation site creation rate of liquid water droplets forming by nucleation from supersaturated water vapor. $J_{v_S}(T)$ is the homogeneous nucleation of solid ice particles directly from supersaturated water vapor by homogeneous nucleation. These functions are not presented herein, since it is the J(T) function for nucleation of ice from liquid water (i.e. $J_{LS}(T)$) that is of interest herein. This function was reported in a following paper.

This important paper presents experimental data on nucleation rates of droplets condensing from supersaturated cloud chamber environments, and also, confirms that small submicron and micron sized liquid droplets spontaneously freeze, in about 0.01 seconds, when suddenly formed and environmentally exposed to temperatures at or below -41°C. Other aspects, results, and observations of this paper will not be reviewed herein.

2.11 Hagen, Anderson, and Kassner, Jr. (1981)

This paper [34] is a continuation of the work presented in Anderson, et al., (1980), described in 2.10. In their 1981 paper, Hagen, et al., continue to analyze experimental data on ice nucleation to gain a better understanding of ice nucleation rates, as well as, better estimates of

the free energy of formation of ice germ embryos, under conditions of homogeneous nucleation, i.e., the nucleation of ice from pure supercooled water.

The authors discuss the important role of homogeneous nucleation in the ultimate objective of understanding heterogeneous nucleation. The authors present the J(T) function for the nucleation rate of ice germs, or ice precursors, or ice embryos, from pure supercooled liquid as $J_{LS}(T)$, in the form

$$J_{LS}(T) = n' \nu \left(\frac{4\sigma_{LS}}{kT}\right)^{\frac{1}{2}} \left(\frac{n_L kT}{h}\right) \exp\left\{-\frac{\Delta g}{kT} - \frac{\Delta G_{LS}^*}{kT}\right\}.$$
(57)

The parameters of this equation are defined as:

- $n' \equiv$ the number of molecules of water in contact with a unit area of the ice germ surface,
- $v \equiv$ the volume of a water molecule in ice,
- $\sigma_{LS} \equiv$ the interfacial free surface energy of an ice and water interface,
- $n_L \equiv$ the number of liquid water molecules per unit volume,
- $\Delta g \equiv$ the activation energy for the transfer of a water molecule across the water-ice boundary,
- $\Delta G_{LS}^* \equiv$ the increase in free energy of the system (of molecules) of a critical-sized embryo, i.e. an embryo that is metastable enough to become an ice germ or ice precursor.

The parameters k and h are the Boltzmann constant and the Planck constant, respectively. The expression for ΔG_{LS}^* was given as

$$\Delta G_{LS}^* = \frac{16\pi\sigma_{LS}^3}{\left[3\left(n_s kT \ln\left\{\frac{P_L}{P_S}\right\}\right)^2\right]}$$

(58)

where

 $n_s \equiv$ the number of ice molecules per unit volume,

 $P_{L} \equiv$ the saturated vapor pressure of water over a plane surface of liquid water,

 $P_s \equiv$ the saturated vapor pressure of water over a plane surface of ice

This form for $J_{LS}(T)$ was ascribed to the theoretical developments of Turnbull and Fisher [7], Dufour and Defay [47], and Hobbs [48].

The authors note that small, 1-20 μ m diameter drops, were found to freeze in time scales of order 0.01 seconds, when exposed to ambient temperatures at or below about -40°C.

By applying nucleation theory to fit the freezing populations in their cloud/expansion chamber experiments, Hagen, et al., determine the "energy barrier for freezing" as a set of empirical fits for $\Delta g + \Delta G_{LS}^*$, based on the minimum temperature achieved in each experiment. Then, an averaged fit for $\Delta g + \Delta G_{LS}^*$ was obtained by least squares, over all the experiments, as

$$\Delta g + \Delta G_{rs}^* = -1.739234 \ x \ 10^{-8} + 8.1157 \ x \ 10^{-21} T \quad J \tag{59}$$

in units of joules, for temperature in Kelvins. The authors present a figure (Fig. 6, [34]) wherein Δg has been separately estimated as a function of T, given by

$$\Delta g = -71.8 \ x \ 10^{-20} + 0.3400 \ x \ 10^{-20} \ T \quad J \tag{60}$$

where Δg has units of joules, and T is in Kelvin.

Note the Δg is the same as the parameter U in the equations for J(T) or I(T) that Mason and co-workers have published, or the parameter A of McDonald [18]. Values for Δg , U or A range from 3.3 x 10⁻²⁰ J (Mason [14]), to about 9.4 x 10⁻²⁰ J (Hagen, et al., [34]). Note that McDonald also discussed a range of values for Δg (or, A) that he estimated [18] as being from about 2.35 x 10⁻²⁰ to about 5.5 x 10⁻²⁰ J. Hagen, et al., discussed the differences in the various values of Δg and the differences of interpretation of Δg . The authors note that Δg , determined from experiments in homogeneous nucleation, can be used without modification for application to predicting heterogeneous nucleation events.

The authors incorporate their empirical fit of $\Delta g + \Delta G_{LS}^*$ into the expression for $J_{LS}(T)$, the homogeneous nucleation freezing nuclei rate equation, to get

$$J_{LS}(T) = J(T) = 5.92 \ x \ 10^{32} \ T \exp\left\{\frac{1.260 \ x \ 10^5}{T} - 588.1\right\}$$
(61)

which has units of nuclei per cubic centimeter per second. The authors show a plot [their Fig. 7] of $\log_{10} J_{LS}(T)$ versus ambient temperature, covering the range of their experiments, and even extrapolating the curve to greater ambient temperatures. Also plotted was a J(T) function attributed to McDonald [18], which does fall within part of their data in the lower temperature range. However, it is not clear to the present author that Hagen, et al., have interpreted the experimental data appropriately. The interpretation, and the evaluation of the J(T) function in the present authors opinion, should be based on the product J(T)Vt, or its integral, which reflects not only the thermodynamic requirements of nucleation, but also reflects the size of the droplet through its volume, as well as the time spent before nucleation, t. Furthermore, in the present author's opinion, when fit to experimental data, the product, J(T)Vt, must reflect the active or effective nucleation sites, i.e., those that lead to actual droplet freezing, not to just theoretical population rates. That is, the nucleation population-rate density function when correlated to experiment, must subsequently predict actual, effective freezing site population rates, not theoretical population rates.

In summary of this paper, it's importance includes the fact that the authors analyze the energy requirements for homogeneous nucleation freezing of small (generally less than 1 micron) water droplets formed by condensation from a supersaturated vapor cloud produced in a cloud/expansion chamber. Their analysis shows that homogeneous as well as heterogeneous nucleation freezing of water droplets can be fit by empirical functions for Δg (or U or A) and ΔG_{LS}^* (or W_c). Thus, the present author's opinion is that when such correlations fit the product J(T)Vt to actual freezing data, the fit must reflect the production rate of effective or active freezing nuclei, not just the total population rate of production.

2.12 Jensen, Toon, and Hamill (1991)

This paper [35] represents another instance wherein homogeneous nucleation freezing theory is modified to predict the reduction in the freezing temperature of atmospheric icing clouds by dissolved sulfuric and nitric acids and/or their hydrates in the supercooled droplets. The work also relates directly to understanding polar atmospheric cloud formation.

The authors review classic homogeneous nucleation theory, presenting a form of J(T) derived in Pruppacher and Klett's classic text <u>Microphysics of Clouds and Precipitation</u> [49], namely

$$J(T) = 2 N_c \left(\frac{\rho_w}{\rho_i h}\right) (\sigma_{iw} kT)^{\frac{1}{2}} \exp\left\{-\frac{\Delta F^+}{kT} - \frac{\Delta F_g}{kT}\right\}$$
(62)

in units of nuclei per cubic centimeter per second. The variables are:

- $N_c \equiv$ number of water molecules in contact with a unit surface area of the ice germ, embryo, or ice precursor molecular cluster.
- $\rho_{w} \equiv$ density of liquid water.
- $\rho_i \equiv \text{density of (normal) ice.}$
- $\sigma_{iw} \equiv$ the surface energy of the ice-water interface.

 $\Delta F^* \equiv$ the energy that must be overcome for the free water molecules to become bound to the ice crystal precursor;

 ΔF^{+} is equal to the difference between the equilibrium energy of the water molecule in the liquid state, and the energy of the water molecule in the ice phase. This quantity has been called the phase change activation energy or just the activation energy, and has been approximated by the energy of activation for the displacement of water molecules in bulk water. ΔF^{+} has been given the symbol Δg by Hagen, et al.. [34], U by Mason [14], and A by McDonald [18].

 $\Delta Fg \equiv$ the work against surface forces required to form a critical ice germ of a metastable size, that can become an ice crystal or ice germ, or embryo.

Jensen, et al., give ΔFg as

$$\Delta Fg = \frac{4}{3}\pi \sigma_{iw} a_g^2 \tag{63}$$

where

 $a_g \equiv$ the radius of the critical nucleus or ice germ. (the same as r_c , before). The expression for a_g given was

$$a_{g} = \frac{2M_{\omega} \sigma_{i\omega}}{L_{m}\rho_{i} \ln\left\{\frac{T_{o}}{T_{e}}\right\}}$$
(64)

where

 $M_{\omega} \equiv$ the molecular weight of water

 $L_m \equiv$ the latent heat of fusion of water (treated as positive)

 $T_a \equiv$ the melting temperature of ice (273.16 K)

 $T_{\star} \equiv$ the ambient (supercooled) temperature

These two expressions for ΔFg and a_g are the same as that provided by, for example, Mason [14] (1960), for W_c and r_c , respectively.

Jensen, et al., report that a_g is modified by acids or acid hydrates present or in solution with the water, so that, based on the theory of Pruppacher and Klett [49], for this case

$$a_{g} = \frac{2 M_{\omega} \sigma_{i/s}}{L_{m} \rho_{i} \ln\left\{\frac{T_{o}}{T_{e}}\right\} + \rho_{i} RT \ln a_{\omega}}$$
(65)

where

 $\sigma_{ii} \equiv$ the "ice-solute surface energy,"

 $a_{\omega} \equiv$ the activity of water in the solution of water and acid.

Because a_{ω} was reported, experimentally, over a fairly wide temperature range, Jensen, et al., could determine the variation of J(T) over the range of polar cloud temperatures for typical droplet sizes. The spontaneous freezing temperature of droplets now varies with the concentration of acid (such as H₂SO₄) present in solution. Freezing temperatures for solution droplets are predicted as low as 195K or -78C.

While this paper has other important results and conclusions concerning the applicability of modified classical homogeneous nucleation theory to predict the freezing of the atmospheric aerosol, these will not be reviewed here. The main point of the present review is merely to document the form for J(T) used by the authors and the method used to modify the function to account for the effects of dissolved material in water to modify the spontaneous freezing temperature of very small, -1 micron, aerosol droplets.

2.13 Stoyanova, Kashchiev, and Kupenova (1994)

This paper [36] develops and tests a method for analyzing and predicting the kinetics of supercooled droplet freezing that encompasses both homogeneous and heterogeneous (seeded)

nucleation freezing. Classical homogeneous nucleation theory is the starting point for the kinetic freezing model. The aim of the paper is "to propose a method for experimental determination of the nucleation rate in freezing droplets, to employ the method for obtaining the ice nucleation rate, to characterize quantitatively the nucleation activity of the aerosols, and to verify the theoretically expected linear dependence of the nucleation rate on the aerosol concentration."

The theoretical considerations begin with the assumption that for each resulting frozen drop, one ice germ or embryo was formed that lead to the frozen drop. Under this hypothesis, if the population of a sample of initially N_o identical drops is studied while being uniformly supercooled at a rate, q, then the time rate of increase in frozen droplets should be functionally related to J(T) (for both homogeneous as well as heterogeneous freezing processes). Stoyanova, et al., show that time can be eliminated as a variable of the problem as follows. If dN(t)/dt is the time-rate-of-change of the number of frozen droplets, then, from the definition of J(t),

$$\frac{dN(t)}{dt} = v[N_o - N(t)]J(t)$$
(66)

where v is the volume of each identical droplet. For a constant cooling rate q, by definition

$$\frac{dT}{dt} = -q. \tag{67}$$

Next, the time-rate-of-change of the number of drops freezing can be written

$$\frac{dN(t)}{dt} = \frac{dN(t)}{dT} \frac{dT}{dt} = v[N_o - N(t)]J(t)$$
(68)

or, substituting for $\frac{dT}{dt}$ and rearranging

3

$$\frac{dN(T)}{dT} = -\frac{v}{q} [N_o - N(T)] J(T).$$
(69)

Integrating this expression from the initial temperature, $T_o = 273.16$ K, to some supercooled temperature, T, gives

$$N(T) = N_o \left\{ 1 - \exp\left\{ -\frac{\nu}{q} \int_{T}^{T_o} J(T^*) dT^* \right\} \right\}.$$
 (70)

Note that $N(T_o) \equiv 0$, i.e. there were no frozen droplets at $T_o = 273.16$ K.

Stoyanova, et al., then solve this equation for J(T) by first rearranging the terms, and then differentiating with respect to T to get

$$J(T) = \frac{q}{v} \frac{d}{dT} \ln\left\{1 - \frac{N(T)}{N_o}\right\}$$
(71)

This equation "shows that the experimental determination of the temperature dependence of the nucleation rate J at known droplet volume and constant cooling rate q reduces to finding the temperature derivative of the experimentally obtainable quantity $ln(1 - N(T)/N_{o})$." Also, this procedure applies whether the droplets contain motes or foreign particles or not. Thus, J(T) for heterogeneous nucleation/droplet freezing events can also be experimentally determined. The authors also remark that this method of determining J(T) can also be applied to non-steady-state nucleation processes.

The authors next review classical nucleation theory and present J(T) in the form

$$J(T) = A(T) \exp\{-W^* / kT\}$$
(72)

where

 $W^* \equiv$ "the nucleation work,"

 $A(T) \equiv$ "a kinetic factor whose temperature dependence is usually weaker than that of the exponential term."

k is the Boltzmann constant as defined in a previous subsection. The authors note that J(T) is really J(T(t)) or J(t), since the phase temperature usually depends on time. Application of the classical theory is valid for relatively slow temperature transients ("...sufficiently small cooling rate, q"...), so that J(T) represents an approximation valid for slow cooling rates. That is, in other words, J(T) in the form due to Volmer [3], and others, was derived for steady-state isothermal processes. This does raise questions about whether this theory can be used directly to predict supercooled droplet freezing in rapidly accelerating and cooling ducted spray flow fields. Clearly, future research on the freezing of very rapidly cooled water droplets is needed.

Stoyanova, et al., present A(T) as

$$A(T) = Z f N_a \tag{73}$$

where

 $Z \equiv$ the Zeldovich factor

 $f \equiv$ the frequency of attachment of water molecules to the molecular cluster or ice germ

 $N_a \equiv$ "the concentration of active centers on which nuclei can be formed"

The Zeldovich factor, Z, was then given as

$$Z = \left(W * / 3\pi \ k \ T \ n *^2 \right)^{1/2} \tag{74}$$

where

 $n^* \equiv$ "the number of molecules in the nucleus"

Stoyanova, et al., give f as

$$f = n_s^* \beta k T / \eta v_m \tag{75}$$

where

 $n_{\star}^* \equiv$ "the number of attachment sites of molecules on the nucleus surface,"

 $\eta \equiv$ "the viscosity of the liquid around the nucleus,"

 $v_{\pm} \equiv$ "the molecular volume," i.e. the volume of a single molecule,
β = is a factor that accounts for the change in liquid viscosity near the surface of the

molecular cluster ($\beta \le 1$).

The expression for A(T) that then results is

$$A(T) = \left(\frac{Z n_s^* \beta k T}{\eta v_m}\right) N_a.$$
(76)

Stoyanova, et al., report that classical nucleation theory gives

$$W^* = \alpha \,\sigma_{ef}^3 \,v_m^2 \,/\,\Delta\mu^2 \tag{77}$$

and

$$n^* = 2W^* / \Delta \mu \tag{78}$$

where

- $\alpha \equiv$ "a numerical shape factor for the molecular cluster, e.g., $\alpha = 16\pi/3$ for spherical nuclei,"
- $\sigma_{cf} \equiv$ the specific (effective) surface free energy of the liquid/nucleus interface (equivalent to σ_{sL} or σ_{LS} used in other J(T) or I(T) expressions presented herein).
- $\Delta \mu \equiv$ the difference between the chemical potentials of the molecules in the liquid and in the ice germ or molecular cluster.

The authors note that, for heterogeneous nucleation, the value of σ_{e} is less than that for homogeneous nucleation. So that, in general, one can write

$$\sigma_{ef}^3 = \Phi \sigma^3 \tag{79}$$

where

- $\sigma \equiv$ the specific surface free energy for homogeneous nucleation (same as σ_{LS}).
- $\Phi \equiv 0 \le \Phi \le 1$; a parameter used for quantitative characterization of the nucleation activity of aerosol particles or "active centers."

Stoyanova, et al., note that "different theoretical models give different expressions for Φ ." Thus, theoretical models such as that discussed by, for example, Fletcher [24] can be used to estimate Φ for the application of J(T) to predict heterogeneous nucleation freezing of supercooled water drops. Stoyanova, et al., report a functional form for Φ in the case where a hemispherical cap shaped nucleus, or ice germ, begins to form on a flat surface:

$$\Phi(\theta) = \frac{(2 + \cos\theta)(1 - \cos\theta)^2}{4}.$$
(80)

This expression is attributed to Volmer [50]. In this expression for Φ , θ is the angle of "wetting" of the surface by the nucleus. Values of $\Phi(\theta)$ would be

 $\Phi(180) = 1$

for either non-wetting or homogeneous nucleation;

$$\Phi(90) = 0.5$$

for "half-wetting" and;

$$\Phi(0) = 0$$

at "full wetting." Full wetting would seem to imply particle freezing at the bulk water freezing temperature, T = 273.16K.

Stoyanova, et al., also provide thermodynamic expressions for the evaluation of $\Delta \mu$. For the freezing of water, they provide

$$\Delta \mu = \frac{\Delta S_m T \Delta T}{T_m} \tag{81}$$

where

 $\Delta S_m \equiv$ the entropy change of melting (see equation (21), for example),

 $\Delta T \equiv T_m - T$, the degree of supercooling experienced by the water sample,

 $T_m \equiv$ the melting temperature of ice, $T_m = 273.16$ K.

By combining the various expressions given so far, Stoyanova, et al., write J(T) as

$$J(T) = Z n_s^* \beta \left(\frac{kT}{\eta v_m}\right) N_a \exp\left\{\frac{-\alpha \sigma_{ef}^3 v_m^2 T_m^2}{\Delta S_m^2 k T^3 \Delta T^2}\right\}$$
(82)

where, with $\sigma_{ef}^3 = \Phi \sigma^3$, J(T) will apply to both heterogeneous and homogeneous nucleation freezing events. For values of some of the parameters in the equation for J(T), the authors provide:

 $\alpha = 16\pi/3$

 $\sigma = 0.02 \text{ J/m}^2 (\text{or } 20 \text{ erg/cm}^2)$

 $\Delta S_m = 2.65 \ k \ (k \text{ is Boltzmann's constant})$

$$v_{-} = 3 \times 10^{-23} \text{ cm}^{3}$$

An expression for the viscosity of liquid water was reported as

$$\eta(T) = 0.139 \left(\frac{T}{225} - 1\right)^{-1.64}$$
 poise (83)

but, in the temperature range of their experiments, the authors used

 $\eta = 0.005$ poise (constant η).

The authors note that analysis of the heterogeneous water sample freezing data that they had obtained then required a value of β of about 10⁻⁶. Recall that, for Stoyanova, et al., β is the factor that accounts for "the change in liquid viscosity very near to the nucleus or ice germ surface." The small value for β may "actually reflect the commonly observed failure of the classical nucleation theory to give absolute magnitudes correctly." These aspects of the formulation of J(T), due to Volmer [50], Walton [51], and others, as presented by Stoyanova, et

al., [36], as well as the extension of the theory to higher cooling rates of particles, need further review and clarification.

An application of the nucleation theory was made by Stoyanova, et al., to both seeded and unseed supercooled water freezing. The purpose was to evaluate key parameters of the theory, such as $\sigma_{e^{p}} \Phi$, θ , $\Delta\mu$, and W* for both "seeded" and "unseeded" water samples. The "unseeded" water was not pure enough to reflect homogeneous nucleation, being referred to as "distilled water." Stoyanova, et al., provided tabulated values for the various parameters listed above that resulted from analysis of supercooled water sample freezing due to different types and amounts or concentrations of "atmospheric aerosol particles." As expected, $\sigma_{e^{p}} \Phi$, θ , $\Delta\mu$, and W* all decrease as the concentrations of seed particles in the water samples increased. There were also changes in these variables for different types of seed particles or "active centers" as they were denoted.

In summary, Stoyanova, et al., like Fletcher [24], describe a general approach for the application of homogeneous nucleation theory to heterogeneous nucleation freezing processes. It is clear that atmospheric samples of water collected from flight through icing conditions will be needed in the future to begin the process of modeling freezing of atmospheric supercooled water, as well as for the simulation of icing tests in ground test facilities. The particle concentrations in these samples, as well as classification of particulate types will be needed for future computer-based modeling studies of the freezing of supercooled water. Water droplets that impinge on aircraft surfaces can form thin run off sheets whose small thickness may permit the use of modified homogeneous nucleation theory, or heterogeneous nucleation theory, to identify where in the run back process freezing begins. Other effects on the heterogeneous nucleation freezing process in supercooled water, such as the effect of different types of seed/mote particles simultaneously present in the water samples, as well as their concentrations, were also discussed

Stoyanova, et al., which makes their paper an excellent starting point for future researchers to review the theory.

2.14 Summary of the Overview

This section of the report was intended to provide a brief overview of some aspects of homogeneous nucleation theory applied to predict the nucleation or crystallization of supercooled liquid water droplets. This overview was not intended to be an inclusive survey of nucleation theory and it is not. It was also the purpose of this section to indicate how homogeneous nucleation theory has been extended to apply to heterogeneous nucleation events, that is, to instances where supercooled liquid water droplets, of ordinary purity, freeze due to contact with solid surfaces or to the presence of entrained foreign particles ("motes") in the water.

Because heterogeneous nucleation is the predominate mechanism by which supercooled water drops freeze in natural and artificial icing environments, the use of modified homogeneous nucleation theory to predict droplet freezing represents a physics-based freezing theory, compared to probability-based freezing theories developed earlier by Levine [11], Bigg [16,17], and applied by the present author to predict spray cloud freezing in ground test simulations of icing conditions [1,2]. The first-generation freezing model was based on the Levine-Bigg probability arguments for activities of the "freezing nuclei" in mote-induced freezing. The second-generation freezing model is based on modified homogeneous nucleation theory (MHNT) which can, in principle, account for the effects of mote-types on water particle freezing by accounting for the "wetting" characteristics of the mote-types, as well as for chemical additives that change either vapor pressures or interfacial surface free energies, and for surface geometries.

In the following sections, the development of a water particle crystallization model for heterogeneous freezing is described. The incorporation of this model into a numerical, one-dimensional, multiphase flow code is discussed and results are presented which were obtained from predictions made of water particle freezing in ducted flows using this code. Comparisons of these predictions are made to the same cases and similar results obtained with the first-generation water particle freezing model (the Levine-Bigg Model) [1,2].

Conclusions are drawn from the comparison of results and recommendations are made for further research and development of water droplet crystallization models for multidimensional numerical flow codes. The other implications of modified homogeneous nucleation theory are discussed, which include other aspects of adverse weather phenomena, as well as for weather simulation in ground test facilities.

3.0 BEHAVIOR OF THE FREEZING NUCLEI RATE OF CREATION FUNCTION, J(T)

3.1 J(T) Utilized in Present Study

The first step taken by the present author to understand the behavior of the function J(T) in predicting the onset of freezing of submillimeter, pure, supercooled water droplets was to reproduce the freezing temperature versus droplet diameter curve shown in Figure 54, page 496, of the paper of Langham and Mason [23]. This curve was based on their equation (4) of their paper which is provided as equation (43) of the present report (which corrects the typographical error in their equation (4)). The author presented the freezing temperature curve thus obtained, in Figure 1, of a previous report [1], as described in [1].

In the present study, the author found it convenient to use a simplified version of equation (43), which is the form first provided by Mason [14], shown in the present report as equation (36), page 11. The reason for using the simpler format for the equation is that by numerical

experimentation, it is obvious that the parameters having the greatest effect on the value of J(T)are σ_{sl} , L_{p} and T. For given values of L_{f} and T, the function J(T) is very sensitive to values of σ_{sl} , because its cube enters the exponential term of J(T). Therefore, adopting a constant value for L_{f} , the heat of melting of ice as

$$L_f = 3.33 \times 10^5 \text{ J/kg},$$

and a shape factor ω evaluated at

 $\omega = 23$

the form of the J(T) equation used in the present study is given by

$$J(T) = 6.9708 \ x \ 10^{38} T \exp\left\{-\frac{2390.2}{T} - \frac{2.367 \ x \ 10^7}{T} \sigma_{SL}^3 \left(\frac{273.16}{273.16-T}\right)^2\right\}.$$
 (84)

3.2 Properties and Behavior of J(T)

The form of J(T) given in equation (84) explicitly includes σ_{sL} as a free parameter meaning that σ_{sL} must be determined for both homogeneous as well as heterogeneous (mote-induced) freezing processes which are to be modeled in the present study. In the form given by equation (84), assuming, as did Langham and Mason [23], that 1 micron diameter pure water drops freeze in 0.6 seconds at -41°C, the value for σ_{sL} was found to be required at

$$\sigma_{SL}^* = 0.0212 \, \text{J/M}^2 \tag{85}$$

Figure 1 shows a plot based on J(T) of the calculated number of critical-sized active freezing nuclei for a one-micron diameter droplet of pure water as a function of the droplet temperature (in degrees centigrade). It was assumed, in the present study, that as a minimum, at least one freezing nuclei had to be present in the drop to initiate freezing at -41°C. The number of critical-sized nuclei present was evaluated from equation (51) and (52) of the present study, using J(T) given by equation (84), under the assumption that each value of N_c obtained was obtained for a

fixed particle size (1μ) and at each fixed temperature, *T*. Thus, under these conditions, equation (52) reduced to

$$N_{c} = J(T) * \frac{\pi}{6} D^{3} * t$$
(86)

where D = 1 micron, and t = 0.6 seconds. Note, that, the actual production rate term for nuclei generation, that is, J(T), has a distribution with temperature as shown in Figure 2. Thus, the use of the thermodynamic theory of spontaneous nucleation to predict particle freezing cannot be simply based on values taken by J(T), but rather, must be obtained with J(T) using some criteria such as that expressed by equations (51) and (52) in the present study. As a second point of interest, Figure 3 is provided to show the freezing nuclei production rate term J(T) normalized by the number of molecules per cubic meter present in the water sample, N. Thus, Figure 3 shows a plot of J(T)/N versus the temperature of the supercooled water. As this figure shows, the formation of molecular clusters containing between 20-300 molecules, that form the ice germs or ice precursors, is still a fairly rare molecular process in the water sample, even at -41°C.

3.3 <u>Review of Dorsch and Hacker Data for Heterogeneous Freezing of Submillimeter</u> <u>Drops</u>

In [1], the present author selected the data provided by Dorsch and Hacker [10] to obtain a Levine-Bigg type heterogeneous freezing function for use in a one-dimensional multiphase flow code. The Levine-Bigg freezing function fit to the Dorsch and Hacker data was

$$T_{\rm F} = 422.2 + 5.592 \ln (D) {\rm R} \tag{87}$$

where T_r is assumed to be the median freezing temperature of drops of diameter, D, in microns. In degrees centigrade, this function was

$$T_{\rm F} = -38.6 + 3.11 \ln (D)C.$$
(88)

The Dorsch and Hacker data set was also selected in the present study for consistency with the previous study [1] and for the reasons it was selected in the first study, as well as because this data set falls in the range typical of heterogeneous freezing experiments. The first step taken in the present study to develop a more physics-based particle freezing function for heterogeneous freezing, based on modified homogeneous nucleation theory, was to review the Dorsch-Hacker data set. Figure 4 shows a set of data plotted that came from Fig. 8 and part of Table I of the Dorsch-Hacker report [10]. The figure shows both the original, hemispherical "drop" sizes, as well as the equivalent spherical droplet sizes, D_{re} , where

$$D_{eq} = 2^{-\frac{1}{3}} D$$
 (89)

In the present study, additional data points were assembled from Table I, Fig. 8, Fig. 10, and Fig. 11 of the Dorsch-Hacker report which were then combined into a single figure, Fig. 5, of the present report. Also on Fig. 5 is a curve representing the data utilized in [1] to obtain the Levine-Bigg curve fit, equations (87) and (88). Thus, the previously used representation of the Dorsch-Hacker data in [1] was deemed adequate, for defining a Levine-Bigg freezing function. Figure 6 shows the data of Fig. 5 corrected to spherical diameters. In the present study, using all of the data of Fig. 6 to obtain a new representation of a Levine-Bigg freezing function, a new curvefit equation was obtained for particle heterogeneous freezing temperature, given by

$$T_F = -38.3 + 3.19 \ln (D)$$
 in degrees C (90)

This function is plotted in Fig. 7, and is essentially the same as that obtained previously [1], i.e., it compares favorably to equation (88).

Therefore, either of the Levine-Bigg type representations of the Dorsch-Hacker data are assumed accurate enough for expressing freezing behavior of the Dorsch-Hacker experiments on heterogeneous particle freezing.

The Levine-Bigg type of freezing function for heterogeneous freezing of water particles is considered to be a first-order or lowest-level freezing model from the standpoint of the physics of the process.

3.4 Freezing Function Based on Modified Homogeneous Nucleation Theory (MHNT)

A second order, or higher-level freezing model, was developed next in the present study based on the following criterion and equations:

(1) Criterion for initiation of freezing of a water particle;

with the number of active freezing nuclei being given by

$$N_{c}(T,D,t) = \int_{t_{o}}^{t} J(T,D,t) V(D,t) dt$$
(91)

the freezing criterion is

$$N_c(T, D, t_F) = 1.0$$
 (92)

That is, in the present study, it is assumed that at least one active freezing nuclei must be created in the drop of volume, V, in the time period τ , where $\tau = t_F - t_o$, to initiate freezing. The times are defined herein as

 t_o = the time when the droplet temperature first reaches the bulk water freezing temperature, i.e. 273.16 K

 $t_{\rm r}$ = the time when the drop begins to freeze, i.e., by definition, when $N_{\rm c} = 1.0$

(2) It is assumed that the homogeneous nucleation rate function, J(T), can be used, as modified, for heterogeneous freezing predictions. That is, the freezing nuclei rate equation,

J(T,D,t), is the same function as the homogeneous nucleation rate equations given by equation (84), except, that σ_{sL} is now a function of drop diameter; that is, $\sigma_{sL} = \sigma_{sL}(D)$. Thus, J(T,D,t) for heterogeneous freezing is assumed given by

$$J(T, D, t) = 6.9708 \times 10^{38} T \exp\left\{-\frac{2390.2}{T} - \frac{2.367 \times 10^7}{T} [\sigma_{sL}(D)]^3 \left(\frac{273.16}{273.16 - T}\right)^2\right\}$$
(93)

(3) It is assumed that the specific surface free energy for heterogeneous freezing is a function of drop volume, hence a function of the number and type of active freezing nuclei present in the droplet volume. For simplicity, at present, therefore, it was assumed that σ_{sL} could be represented by a smooth function of the droplet diameter, D, or,

$\sigma_{sL} = \sigma_{sL} (D)$ (heterogeneous freezing)

From a practical standpoint, therefore, the specific surface free energy, $\sigma_{sl}(D)$, for the heterogeneous nuclei (water plus mote) responsible for the initiation of droplet freezing in the Dorsch-Hacker experiments can be determined for representative drop sizes by matching the criterion given by equation (92) for each drop size, at its median freezing temperature. To implement this process, and to develop a smoothly varying σ_{sl} function of D, the median freezing temperature of each size drop was assumed to be given by equation (90). Thus, $\sigma_{sl}(D)$ was determined by requiring N_c (T_{FP} D_p t_{Fl}) to be unity for a selected set of freezing temperatures predicted by equation (90). Thus, if T_{Fi} is the median freezing temperature of droplet with diameter, D_{ij} then

$$N_{c}(T_{F_{i}}, D_{i}, t_{F_{i}}) = 1.0 = \int_{t_{o_{i}}}^{t_{F_{i}}} J(T_{F_{i}}, D_{i}, t) V(D_{i}, t) dt$$
(94)

 $T_{F_{i}}$ and $V(D_{p}, t)$ were assumed constant during an assumed freezing interval given by

$$t_i = t_{Fi} - t_{ai} = 1.0$$
 second.

Then $N_c(D_p, T_{Fi})$ could be approximated by

$$N_{c}(T_{F_{i}}, D_{i}) = 1.0 = J(T_{F_{i}}, D_{i})V(D_{i})(1.0)$$
(95)

With

$$V(D_i) = \frac{\pi}{6} D_i^3 \tag{96}$$

the freezing criterion actually applied to determine $\sigma_{sl}(D_i)$ was

$$N_{c}(T_{F_{i}}, D_{i}) = 1.0 = J(T_{F_{i}}, D_{i})\frac{\pi}{6}D_{i}^{3}$$
(97)

Figure 8 shows $N_c(T_{FF}, D_i)$ plotted as a function of T for drop sizes ranging from 1-10000 microns. The line $N_c = 1.0$ cuts each N_c curve at the median freezing temperature of that respective droplet size. For example, for the 5 micron droplet, the line $N_c = 1.0$ cuts the 5 micron N_c curve where $T = -33^{\circ}$ C. By inspection of Fig. 7, it can be seen that this corresponds to the point on the curve fit for the heterogeneous, spontaneous (median) freezing temperature where 5 micron droplets freeze at $T_F = -33^{\circ}$ C (the ordinate value of T_F for ln (5) = 1.609).

Each curve plotted on Fig. 8 was obtained by trial and error evaluation of σ_{sl} , changing the value of σ_{sl} for each drop size, until the N_c curve generated for that drop size passed through the appropriate freezing temperature when $N_c = 1.0$.

With the determination of each successful value of σ_{sL} , for each a'priori chosen drop size, the data in Fig. 9 was constructed. Shown in Fig. 9 are two curve fits for σ_{sL} that were obtained by the same method to deduce the variation of σ_{sL} for heterogeneous freezing. The first attempt produced the distribution denoted $\sigma_{sLI}(D)$. A second attempt to create a σ_{sL} distribution in the present study corresponds, in Fig. 8, to the curve $\sigma_{sL2}(D)$. The differences between σ_{sLI} and σ_{sL2} result from a later, more thorough reconstruction and analysis of the Dorsch-Hacker data by the author. The forms of the curve fit equations for σ_{sLI} and σ_{sL2} are

$$\sigma_{SL1}(D) = 0.02045 - 0.00044528 \ x \left(\ln(D)\right) - 7.2527 \ x \ 10^{-5} \left(\ln(D)\right)^2 \tag{98}$$

and

$$\sigma_{SL2} = 0.020228 - 0.00044755 \ x \left(\ln(D)\right) - 7.9081 \ x \ 10^{-5} \ x \left(\ln(D)\right)^2 \tag{99}$$

These two representations for σ_{sL} were kept in the present study to investigate the sensitivity of the predicted freezing behavior to variations in the values of σ_{sL} , for the same water sample data set. σ_{sL2} is the recommended curvefit because it is believed more accurate.

The differences in predicted submillimeter water particle freezing behavior in ducted flow when computed using the Levine-Bigg theory, σ_{sl} , and σ_{sl2} , will be shown in the next section.

4.0 ANALYSIS AND EVALUATION OF THE WATER PARTICLE FREEZING MODEL BASED ON MODIFIED HOMOGENEOUS NUCLEATION THEORY

4.1 Introduction to the Approach

Because of a lack of detailed data on the freezing of convected, submillimeter, supercooled water droplets in ducted flows, the analysis and evaluation of the new model for heterogeneous freezing must be based on comparing model results with model results. Therefore, the predicted results of a calculation of water particle freezing in a ducted flow, made based on the Levine-Bigg freezing model [1], will be compared to the corresponding results obtained based on the modified homogeneous nucleation theory (MHNT) presented in the last section. The results obtained with the Levine-Bigg model were made with a one-dimensional, multiphase flow code described in references 52, 53 and modified and updated in the study reported in [1]. The flow code solves the fully-coupled mass, energy, and momentum conservation equations for a dilute, multiphase flow of water particles entrained in a ducted

airflow [52, 53]. The solution procedure is a Runge-Kutta fifth-order numerical integration scheme [54, 55] that requires a set of input duct flow initial conditions as well as the duct geometry.

The code was denoted, in the author's first report on water particle modeling [1], as the "AEDC1DMP" code (pg. 34, [1]). This code designation will be retained in the present study, meaning AEDC <u>one-dimensional</u>, <u>multiphase</u> flow code. It will be made clear when the results of this code refer to predictions made based on the Levine-Bigg (LB) type freezing model or the modified homogenous nucleation theory (MHNT) model for particulate freezing.

4.2 Implementation of the Particle Freezing Model

4.2.1 Freezing Criterion

In implementing the MHNT model in AEDC1DMP, the transformation of the calculation of the time integral for $N_c(T, D, t)$ to an integral along the flow with respect to the axial flow coordinate, was done as follows. The time integral was rewritten as

$$N_{c}(T,D,x) = \int_{x_{o}}^{x} \frac{J(T,D,x)V(D,x)\,dx}{U(x)}$$
(100)

where

$$dx = U(x) dt \tag{101}$$

and U(x) is the particle velocity at x. The lower limit on the integral, x_o , is the axial location in the flow where the given sized droplet first reaches (or cools down to) a temperature of 32°F (0 C or 273.16 K). Thus, droplet temperatures are monitored and when a given size drop temperature reaches 32°F, the evaluation of the integral is begun. However, the integral was actually implemented in the code by a running, finite summation

$$N_{c}(T,D,x_{k}) = \sum_{i}^{k} \frac{J^{*} V^{*} dx_{i}}{U^{*}}$$
(102)

where J, V, and U are mean or representative values of J, V, and U over the spatial integration step interval, dx_i , at the axial location, x_i , given by

$$x = x_{k} = x_{o} + \sum_{i}^{k} dx_{i}$$
(103)

The integration step intervals, dx_i , are controlled by the numerical integration scheme to control numerical (truncation) errors. Note, that when x reaches a limiting, input preset value, x_{max} , then the AEDC1DMP code terminates the calculation (integration) process for the entire flow.

The implementation of the Levine-Bigg type freezing function in AEDC1DMP was discussed previously [1,2]. However, to review the process, the implementation was as follows. In the course of computing the flow, the temperatures of the liquid drops are monitored. When a drop of diameter D cools to a temperature at or below the corresponding median freezing temperature T_{L_r} , for that drop size, the droplet is considered to have begun the freezing process, that is, the freezing procedure is initiated.

4.2.2 Model of the Freezing Process

Whether the initiation of freezing is triggered by a criterion given by the Levine-Bigg model or by the modified homogeneous nucleation theory (MHNT) model, once freezing has been initiated, the computation of the drop freezing is the same. The steps are

- (1) the supercooled liquid drop is assumed to become, instantaneously, a slushball or mixture particle containing both liquid water and ice at the normal melting temperature of ice (492°R, 32°F, or 0°C).
- (2) The amount of the mixture particle that is ice is determined from an energy balance for the particle based on conditions just before freezing was initiated and just after freezing was initiated. Since the mixture temperature is assumed at 492°R, the only

free variable that can be determined from the energy balance is the fraction of liquid that remains after freezing is triggered, given by

$$\alpha_{F} = 1 - \frac{C_{L}}{H_{F}} \left(492 - T_{L_{F}} \right) \tag{104}$$

where

 $\alpha_F \equiv \text{mass fraction of drop or mixture particle that remains liquid at 492°R. (hence,$ $(1 - <math>\alpha_F$) of the particle has become ice at 492°R).

 $C_i \equiv$ specific heat of liquid water

 $H_{\rm F} \equiv$ heat of fusion of liquid water

 $T_{L_r} \equiv$ freezing temperature of the supercooled liquid water droplet.

Details of the energy balance are provided in [1].

(3) The density of the mixture particle is now the mean density of a two-phase system wherein the phases cannot occupy the same volume. Hence, the mean density of the mixture particle is given by

$$\overline{\rho} = \frac{\rho_L \rho_I}{\alpha_F \rho_I + (1 - \alpha_F) \rho_L} \tag{105}$$

where

 $\rho_L \equiv$ density of liquid water

 $\rho_{i} \equiv \text{density of ice}$

The specific heat of the particle is also redefined as a mass-average specific heat given by

$$\overline{C}_{P} = \alpha_{F} C_{P_{L}} + (1 - \alpha_{F}) C_{P_{I}}$$
(106)

where

- $C_{P_{l}}$ = specific heat of liquid water
- C_{P} = specific heat of ice
- (4) The size of the mixture particle is then given by requiring that the total mass of the particle remains constant during the (assumed instantaneous) transition of the particle from supercooled liquid drop to a mixture particle of water and ice at 492°R. Thus,

$$\overline{\rho} \, \frac{\pi}{6} \, \overline{D^3} = \rho_L \, \frac{\pi}{6} \, D_{l_F}^3 \tag{107}$$

or

$$\overline{D} = \left(\frac{\rho_L}{\overline{\rho}}\right)^{1/3} D_{l_F} \tag{108}$$

where

 \overline{D} = the new diameter of the mixture particle

 $D_{L_F} \equiv$ the diameter of the supercooled liquid drop just before freezing is initiated.

(5) Thereafter, the mixture "drop" is treated as any other droplet in the flow. Heat, mass, and momentum exchanges with the air flow are computed in the same way as that for the all droplets. However, as each amount of heat is extracted from the particle, the amount of liquid remaining in the droplet is reduced (α_r is reduced), and a new mean density, and particle diameter are re-calculated. In the recalculation, the mass loss or gain of the particle by mass transfer is taken into account, as well as energy lost or gained by the mass transfer process. The amount of new ice formed is also calculated. The details of the mass and energy balance during freezing of the evaporating, two-phase particle are given in Appendix B.

When the entire particle has become ice, α_r reaches zero, and the particle is now completely frozen.

During the freezeout process for the mixture particle, the particle temperature is held constant at 492°R.

(6) Once the particle completely freezes, its density becomes that of ice, as well as its specific heat. From this instant or location in the duct flow, the particle undergoes mass, energy, and momentum exchanges with the air flow as an ice particle. It can gain or loose mass by sublimation based on the partial pressure of ice and the air flow specific humidity.

5.0 EVALUATION OF AEDC1DMP CODE ON REPRESENTATIVE DUCT FLOW CASES

5.1 Baseline Case of Ducted, Two-Phase Flow

The base case computed for a comparison purpose was a ducted flow with a "single" water spray station at its inlet.

In the code AEDC1DMP, the water is assumed to enter the airflow at up to ten different injection stations that are separated, axially in the duct. Each injection station is capable of putting in a given amount of water, in a given sized water particle with its specified velocity and temperature. When all of the injection stations are bunched close together, axially, in the duct, they can be used to model a single spray station with a spray droplet size distribution characterized by ten discrete drop sizes. This was the approach taken in the previous study [1] to examine the freezing of two-phase, ducted flows representative of both full-scale and research-scale icing test facilities. Note, that to illustrate and emphasize the ducted-flow freezing processes of supercooled, submillimeter water particles, the air flow inlet total temperature utilized was significantly lower, at 460°R, than is typical of many icing tests, where an airflow inlet total temperature of 486°R is more likely. In any case, for consistency with previously

obtained results, the same airflow and water inlet conditions have been specified and are

provided below.

 Table 1. Duct Air Flow Inlet and Water Spray Conditions for Baseline Test Case 1

Air Inlet Conditio	ns
Total Pressure, psfa	1123.48
Total Temperature, °R	460°R
Velocity, ft/s	20
Mach Number	0.0189
Inlet Relative Humidity, percent	15.78

Water (Particle) Input Conditions				
Injection Station No.	Particle Diameter	Particle Temperature	Particle Velocity	Load Factor
(ft)	(microns)	(°R)	(ft/s)	(FL)
1.00	5	530	46	6.85 E-05
2,0,01	10	530	46	2.06 E-04
2,0.01	15	530	46	4.80 E-04
3, 0.02	15	530	46	3.62 E-04
4, 0.03	20	530	46	2.06 E-04
5, 0.04	30	530	46	3 43 E-05
6, 0.05	40	530	40	1 37 E-05
7, 0.06	50	530	40	1.37 E-05
8, 0.07	60	530	40	1.57 E-00
9, 0.08	80	530	46	1.37 E-07
10, 0.09	100	530	46	1.4 E-08
Total FL = 1.37 E-03				

FL = lbm of water/lbm of dry air injected at a given injection station

The duct geometry used for this reference or baseline case is shown in Fig. 10. As is clear from inspection of Fig. 10, and as discussed in [1], this duct geometry is not representative of full-scale icing test facilities. Rather, it was chosen in the earlier study [1] because it was related to a similar geometry of interest. For consistency in comparing results, therefore, this geometry was also retained in the present study.

5.2 Comparison of Predicted Results of Baseline Case for Three Particulate Freezing

<u>Models</u>

Figure 11 shows predicted temperatures of 10 micron water droplets, as they vary with axial distance along the flow, for three heterogeneous freezing models. The models utilized were the Levine-Bigg model, given specifically by equation (87), and the modified homogeneous nucleation theory (MHNT) using σ_{SLI} and σ_{SL2} for the specific surface free energy specifications.

As can be seen, the comparison is extremely good. The results obtained with the σ_{sL} function σ_{sL2} gives a closer agreement with the Levine-Bigg model than does σ_{sL2} . Since it is thought that σ_{sL2} was obtained from a more thorough data analysis than σ_{sL2} , the better agreement of the predicted results with σ_{sL2} is expected, although such good agreement in both quality and quantity displayed in Fig. 11 is remarkable.

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On the basis of the results shown in Fig. 11, and found to be consistently the same for all droplet sizes, it was decided to utilize the function $\sigma_{su}(D)$ for the heterogeneous freezing model for all further calculations made in the present study.

Figure 12 shows a comparison of predicted droplet temperatures made with two versions of AEDC1DMP, one version with the Levine-Bigg (LB) freezing model implemented, and the other version with the modified homogeneous nucleation theory (MHNT) model implemented, based on σ_{sL2} . The agreement between the first (LB) and second order (MHNT) models is good. An interesting difference in the predicted results occurs for the 5 micron droplets. The Levine-Bigg model results in the 5 micron droplets freezing at the location in the duct flow where their temperature has reached about 427°R. (The freezing occurs so rapidly, it occurred between print step sizes in the program and is not shown, therefore, in the plot, Fig. 12.) The drops have shrunken, by evaporation, to about 2 micron in diameter at this location. On the other hand, the modified homogeneous nucleation theory model predicts that the 5 micron drops, shrunken to 2 micron, do not freeze within the domain of the calculation, because, even though the drops have reached a freezing temperature, their volume is too small to generate at least one active freezing nuclei. In other words, the criterion

$$N_c = \int JV \, dt = 1 \tag{109}$$

was not satisfied for these small drops throughout the calculated length of flow. Thus, significant potential differences could arise in predicted freezing behaviors made with a Levine-Bigg model and with a modified homogeneous nucleation theory model, for certain circumstances.

The predicted temperature distributions for all of the other drop sizes appear similar or comparable, based on the two different freezing models.

(The calculated temperature distributions shown in Fig. 12 also indicate that the particles took longer (more distance) to freeze for the MHNT model predictions, than in the Levine-Bigg model predictions, than in the Levine-Bigg model predictions. This was not due to any differences caused by the freezing models, but due to improvements made in AEDC1DMP to integrate the heat and mass transfer processes of the freezing particles in a more recent version of AEDC1DMP in which the MHNT model was implemented.)

Figure 13 shows a comparison of the predicted droplet sizes from AEDC1DMP with the MHNT model implemented, for the ten different drop sizes as they move along the flow. Drop freezing shows up in these curves by sudden small increases in drop size (a kink) in the curves. Clearly, the 5 micron particle curve does not show a sudden increase, indicating that the drop stays liquid along the entire flow.

5.3 Calculations of Flow in a Representative Icing Test Facility

5.3.1 Introduction

The duct geometry of baseline case calculated, with results shown in Figures 10-13 was not representative of full scale icing test facility duct geometries. Therefore, a second test case was prepared and calculated wherein the duct or wind tunnel geometry and scale represents typical icing research or icing test wind tunnel geometry.

The results from computations of test case 2 are shown below.

5.3.2 Nominal Test Conditions in a Full-Scale Icing Facility

The duct or wind tunnel geometry assumed for test case 2 is shown in Figure 14. This wind tunnel like configuration has a large inlet and a steep contraction section, to minimize flow disturbances. Icing spray nozzles are assumed to be put in the plane of the large inlet. The test section is located at x = 46 ft, so that the state of the water particles injected at the inlet of the wind tunnel <u>should be</u> in kinetic and thermal equilibrium with the air flow by the time they reach the test section. The tunnel has a large contraction ratio (inlet flow area divided by test section flow area), hence, the range of permissible inlet air flow velocities is small, ranging from near 0 to about 45 ft per second. At the higher inlet air velocities, the test section has reached near-choking, or sonic flow conditions. Generally, icing tests are conducted at air flow or flight-simulating speeds in the range of a few hundred miles per hour. Therefore, a nominal set of inlet test conditions was defined for the representative calculation made with AEDC1DMP. The nominal test conditions are listed below.

Table 2. Nominal Tunnel Inlet Test Conditions for Icing Wind Tunnel-Test Case 2

Air Inlet Conditions

Total Pressure, psfa	2074
Total Temperature, °R	500
Velocity, ft/s	20.8
Mach Number	0.0188
Relative Humidity, %	54.7

Water (Particle) Input Conditions

Injection Station No. (ft)	Particle Dia. (microns)	Particle Temperature (°R)	Particle Velocity (ft/s)	Load Factor (FL)
(a) 1 (0.0)	50	634	15.6	0.00018
(b) 1 (0.0)	500	634	15.6	0.00018

Note: The case was computed first (case 2a) with a single particle size of 50 microns to represent the spray cloud, then computed (case 2b) with a single particle size of 500 microns. In

each computation, the 50 or 500 micron particle represented the mean volumetric diameter of the spray cloud particle size distribution at the inlet.

Some selected results from computing these flow cases are shown in Figures 15-17. Figure 15 shows the gas velocity together with the velocities of a 50 micron particle and a 500 micron particle. The calculation shows that the larger particle did not reach velocity (kinetic) equilibrium with the air flow at the test section. On the other hand, the 50 micron particle did achieve a velocity equal to the gas velocity at the test section location.

Figure 16 shows similar results for the gas temperature and the temperatures of the two different sized particles. In the case of the 500 micron particle, its temperature is above the gas temperature by about 10 degrees, R, at the test section location due to thermal lag, however, the 50 micron particle temperature has fallen quickly, through evaporation and cooling, to its wetbulb temperature about 5 degrees, R, below the gas temperature. The effects of air flow inlet humidity on droplet cooling and thermal non-equilibrium have been well-documented elsewhere, for example [53], and will not be discussed herein. However, it is possible to increase the inlet air flow humidity to increase the drop wet bulb temperature at the test section, to be equal to the gas temperature. Large droplet thermal lag effects could also be reduced, by reducing inlet humidity. Thus, opposite humidity modifications are required for large and small drops to account for thermal non-equilibrium effects at the test section.

Finally, Figure 17 shows the changes in the droplet diameters as the water particles pass down the duct. The changes in droplet diameter are greatest where the rate of evaporation is greatest, in the duct inlet throat region. (Note that kinks or slope discontinuities in the plotted curves are due to the data output frequency or axial spacing, not to discontinuities in computed results.)

Test Case 2, therefore, provided some nominal results which show how AEDC1DMP can be used for assessment of flow variations and flow quality typical of full-scale icing test or research wind tunnels. In this test case, neither the 50 micron drop nor the 500 micron drop were near their spontaneous (median) droplet freezing temperatures. The next test case computed will be one that represents a full-scale icing test with supercooled droplet freezing.

5.3.3 Test Conditions in a Full-Scale Icing Facility with Supercooled Droplet Freezing

By trial, computation, and output analysis, a set of wind tunnel inlet conditions was found that produced supercooled droplets, slushball or mixture particles, and ice particles in the test section located of the flow. The set of inlet conditions found is listed below in Table 3.

Table 3. Tunnel Inlet Test Conditions for Icing Wind Tunnel Flow withWater Particle Freezeout – Test Case 3

Air Inlet Conditions

Fotal Pressure, psfa	2076
Total Temperature, °R	475
Velocity, ft/s	40
Mach Number	0.037
Relative Humidity, %	26

Injection Station No.	Particle Dia.	Particle Temperature	Particle Velocity	Load Factor
(ft)	(microns)	(°R)	(ft/s)	(FL)
1 (0.0)	5	495	16	9.06 E-06
2 (0.02)	10	495	16	2.73 E-05
3 (0.04)	15	495	16	6.35 E-05
4 (0.06)	20	495	16	4.79 E-05
5 (0.08)	30	495	16	2.73 E-05
6 (0.10)	40	495	16	4.54 E-06
7 (0.12)	50	495	16	1.31 E-06
8 (0.14)	60	495	16	1.81 E-07
9 (0.16)	80	495	16	1.8 E-08
10 (0.18)	100	495	16	2.0 E-09
		• • • • • • • • • • • • • • • • • • •	Total FL	= 1.82 E-04

Water (Particle) Input Conditions

As the input data in Table 3 indicate, an inlet spray cloud with a 10 drop size distribution was input at the wind tunnel entrance station. The inlet air temperature and the water spray temperatures are rather low, the water being only slightly above freezing and the air at below freezing temperature (15°F). This combination of inlet conditions lead to predicted freezeout of some of the drop sizes in the wind tunnel flow at the test section station.

Figures 18 show some of the calculated results obtained with AEDC1DMP utilizing the modified homogeneous nucleation freezing model for predicting the freezing of supercooled water particles. Figure 18 shows the air flow Mach number distribution from duct entrance to the test section. Note that this case's inlet conditions lead to rather high test section Mach number, about 0.69. This is about twice the nominal or typical value for icing tests in full-scale facilities. The higher value of air flow Mach number ensured that air flow static temperatures would be low enough to induce particle freezeout, with a freezing model calibrated to the Dorsch-Hacker data set [10].

Figure 19 shows the predicted variations in drop or particle sizes from entrance to test section. Note that the initially 5 micron particle has disappeared, by evaporation, by about 4.5-5.0 feet from the duct entrance. This indicates that, for similar inlet conditions, a real water spray cloud with similar drop size distributions would experience all of the drops smaller than, say, 10 microns evaporating quickly in the inlet. The smallest surviving droplet at the test section would probably be in the range from 6-8 microns. These results, of course, depend on the inlet conditions, especially water injection temperature and inlet air flow humidity.

Figure 19 also shows size-jumps in the 15 to 80 micron particles which indicates that these water particles have begun to freezeout in the air flow where the size-jumps occur. The 10 and 100 micron droplets remain completely liquid at the test section.

Figure 20, which shows the predicted water particle or droplet temperatures, confirms that the 15-80 micron particles have begun to freeze in the flow. The (initially) 10 micron particle does not experience freezing, nor does the (initially) 100 micron particle, up to their arrival at the test section and, they have very different temperatures. The other particles have either become ice particles (15, 20, 30 micron particles) or are slushball or mixture particles (40, 50, 60, 80 micron particles) at the test section. The results of the calculation made with the MHNT freezing model indicate the potentially complex results that could be obtained in actual icing or adverse weather test simulations when either inlet flow, water spray nozzle, or duct geometry result in thermal and kinetic non-equilibrium flow conditions at the test station or test article.

Concluding the results of test case 3, Figure 21 shows the predicted water particle velocities. As Figure 21 shows, for this case, all of the particles, even the 100 micron droplet, are in, or very near to, kinetic or velocity equilibrium with the air flow at the test section (x = 46 feet).

Briefly summarizing this section, the AEDC1DMP code is capable of predicting results of two-phase, dilute, air and entrained water droplet flows in research scale and full-scale icing and adverse weather test facilities. The incorporation of a second generation or second order (in the hierarchy of physics of freezing modeling) modified homogeneous nucleation theory (MHNT) for particle freezing in AEDC1DMP gives the code the capability of accounting for particulate freezing over a wide range of flow conditions. Appendices C-E describe the data input audits format in detail, for user convenience. Included are the input data for test case 3. Appendix F lists a sample output from AEDC1DMP for test case 3.

6.0 DISCUSSION OF RESULTS OF STUDY, SUMMARY, CONCLUSIONS, AND

RECOMMENDATIONS

6.1 Summary of the Present Study

The study and the report had three parts. The first part was an examination of homogeneous nucleation theory and its modifications and applications to the prediction of the freezing of supercooled pure water particles and heterogeneous freezing of water particles containing motes or foreign substances. The second part addressed the behavior of the homogenous nucleation rate function, J(T), as a function of water temperature. This function was then calibrated to a set of heterogeneous particle freezing data, provided by Dorsch and Hacker [10], by assuming that the surface specific free energy for an ice-water interface could be represented as a function of the volume or diameter of the particles containing the motes. With this functional form for the surface free energy in hand, the function J(T) became, in essence, calibrated to the Dorsch-Hacker data set.

The third part of the study comprised some numerical computations of two-phase, dilute, air and entrained water particle flows, using a new version of the AEDC one-dimensional, multiphase flow code, AEDC1DMP. The results of these calculations demonstrated that modified homogeneous nucleation theory [MHNT] could be satisfactorily used in numerical flow models to identify the initiation of heterogeneous freezing events in submillimeter, supercooled water particles in the computed flows. The freezing model described in the report accounts for the effects of time, particle size, and particle temperature on the initiation of freezing.

6.2 Conclusions

The conclusions of the research reported herein are as follows:

- (a) Classical homogeneous nucleation theory can be easily modified to predict heterogeneous freezing of submillimeter, supercooled water particles.
- (b) Both empirical and theoretical methods can be used to account for the effect of foreign particles, surface wetting and dissolved chemicals in the initiation and formation of freezing events.
- (c) A particulate freezing model for heterogeneous freezing of submillimeter, supercooled water particles based on modified homogeneous nucleation theory has been incorporated in a one-dimensional, multiphase flow code (AEDC1DMP) with excellent results, and can easily be incorporated into multidimensional, multiphase flow codes.

6.3 Recommendations

The recommendations of the present study are as follows:

- (d) A multidimensional, multiphase flow analysis of water spray nozzle discharge plumes should be pursued with the highest priority. The analysis should include water particle freezing in the spray plumes to investigate the possibility of creating ice crystals in the spray plumes. This topic has been raised in the first part of this research program [1,2], and it is also seen as an important continuation of the present study.
- (e) The physics of water particle freezing should continue to be studied under heterogeneous freezing conditions. The further development of models for the specific surface free energy of the molecular clusters that trigger freezing events should be pursued. Models of σ_{sl} , such as that described by Volmer [50], Fletcher [24], Stoyanova [36], and others, should be developed so that the effects on particle

freezing of different kinds of motes in the atmospheric water particles, as well as in ground test facility water supplies, can begin to be addressed.

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- (f) In view of the previous recommendation, it is further recommended that additional atmospheric data be obtained on the composition of water particles in nominal icing conditions. The data should include the chemical and physical composition of supercooled water particles or water samples. This includes dissolved and/or entrained chemicals, as well as the kinds, shapes, and numbers of solid particles in the atmospheric icing environment. Similar effort should be made to catalogue the same information for water used in ground test simulations of icing and adverse weather simulations for aircraft engines and/or their components.
- (g) Consideration of the freezing process for supercooled water should be extended to thin films of water on aircraft surfaces, including those on ice accretions. In other words, it should be investigated whether there are water run-back modes wherein thin films of supercooled water can be formed in the ice accretion process. If so, then the initiation of freezeout of run back water, under such conditions, might be predicted by a modified homogeneous nucleation theory.

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RATE OF PRODUCTION OF FREEZING NUCLEI J (T) , NUCLEI/M'3*SEC







PLOT OF "AVERAGE" OR MEDIAN, HEMISHERICAL, SUPERCOOLED DROP SPONTANEOUS FREEZING TEMPERATURE AS A FUNCTION OF DROP DIAMETER, FROM THE DORSCH AND HACKER REPORT PLOT OF "AVERAGE" OR MEDIAN, 4 FIG.

DEG , ERUTARET



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REEXAMINATION OF THE DORSCH & HACKER "AVERAGE" SPONTANEOUS FREEZING TEMEERATURE DATA TO OBTAIN AN ACCEPTABLE LEVEL OF CONFIDENCE FOR MATCHING IT WITH MODIFIED HOMOEREOUS FREEZING THEORY FIG. 5

TEMPERATURE, DEG C





TEMPERATURE, DEG C



FIG. 7 PLOT SHOWING LINEAR FIT TO AVERAGED SPONTANEOUS FREEZING TEMPERATURES, $T_{\rm F}$, OF THE DORSCH AND HACKER DATA[10], BASED ON CORRECTED DROP DIAMETERS

TEMPERATURE, DEG C



1

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NUMBER OF ACTIVE FREEZING NUCLEI

10000 0006 σ_{SIII} (D) =0.02045-0.00044528* (IN(D)) -7.2527E-05* (IN(D)) ^2, J/M^2 og G BASED ON THE CURVES IN FIG. 8 LED TO THE SECOND CURVE FIT BELOW J/M^2 8000 A SECOND ANALYSIS TO IMPROVE AND REFINE THE VALUES FOR $\sigma_{\rm SL}$ TWO DIFFERENT ATTEMPTS TO OBTAIN THE BEST CURVE FITS OR EXPRESSIONS OF σ_{SL2} (D) =0.020228-0.00044755*IN(D) -7.9081E-05*(IN(D)) ^2, 7000 œ σ_{SL} values corresponding to the curves of Fig. DROPLET DIAMETER, MICRONS 6000 PRELIMINARY ANALYSIS GAVE THE FIT 5000 4000 3000 i 2000 0 + 100 C 22×10⁻³ S ø œ 5 ຊ **1**6 ភ 4 2 σ ~ 5 18 1 44 5 = TEW ' QT 2.W/C '

CURVE FITS OF THE SPECIFIC SURFACE FREE ENERGY, Og., OF AN ICE-WATER INTERFACE FOR HEITEROGENEOUS NUCLEATION FREEZING, EVALUATED BASED ON THE DORSCH-HACKER DATA [10] თ FIG.

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TT, RADIUS, R, FT

r.



TEMPERATURE, DEG R

FIG. 11 COMPARISON OF CALCULATED DROPLET TEMPERATURE DISTRIBUTIONS IN DUCTED FLOW WITH FREEZING CONDITIONS FREDICTED



FIG. 12 COMPARISON OF PREDICTED PARTICLE TEMPERATURES DURING FREEZING: LEVINE-BIGG THEORY VS MODIFIED HOMOGENEOUS NUCLEARTION THEORY (MENU)



FIG. 12 (CONTINUED) COMPARISON OF PREDICTED PARTICLE TEMPERATURES DURING FREEZING: LEVINE-BIGG THEORY VS MODIFIED HOMOGENEOUS NUCLEATION THEORY (MANT)

TEMPERATURE, DEG R





TEMPERATURE, DEC R

-



FIG. 12 (CONTINUED) COMPARISON OF PREDICTED PARTICLE TEMPERATURES DURING FREEZING: LEVINE-BIGG THEORY VS MODIFIED HOMOGENEOUS NUCLEATION THEORY (MENU)

> 95 24

ļ,



FIG. 12 (CONFINIED) COMPARISON OF PREDICTED PARTICLE TEMPERATURES DURING FREEZING: LEVINE-BIGG THEORY VS MODIFIED HONGENEOUS NUCLEATION THEORY (MANT)

TEMPERATURE, DEG R

2



TEMPERATURE, DEG R

FIG. 12 (CONCLUDED) COMPARISON OF PREDICTED PARTICLE TEMPERATURES DURING FREEZING: LEVINE-BIGG THEORY VS MODIFIED HOMOGENEOUS NUCLEATION THEORY (MHNT)

엮 Ф ព ф 쇡 Φ 97 ф Φ q 抣 ф Ŋ Œ 1 ф IJ Φ C 14 ф ወ D 扣 Ф 枊 ۵ D ង 巾 ⊕ 们 AXIAL DISTANCE FROM DUCT ENIRANCE, FT ф Ф D 枊 T ⋪ Ľ Φ 9 ф ወ 乜 D ф 乜 80 MICRON INITIAL DIAMETER Œ INITIAL DIAMETER 5 MICRON INITIAL DIAMETER INITIAL DIAMETER 15 MICRON INITIAL DIAMETER 20 MICRON INITIAL DIAMETER 30 MICRON INITIAL DIAMETER INITIAL DIAMETER INITIAL DIAMETER 100 MICRON INITIAL DIAMETER Φ Ш D IJ ω ⋪ Þ ſŤ Þ 巾 抣 Q -----Ð D 60 MICRON LO MICRON 40 MICRON 50 MICRON ф 쇡 þ Œ ф Q U Ф Þ Φ Ф 1 ф Φ 1 巾 Φ 2 Φ P m 1 C m 1 ጠ m 0 D ទ କ୍ଷ 8 ജ ą R 10 8 8 2

PREDICTED VARIATIONS IN DROP SIZES ALONG THE DUCTED FLOW FREEZING PROCESS FIG. 13

DEOPLET DIAMETER, MICRONS





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VELOCITY, FT/S



FIG. 16 PREDICTED AIR AND WATER PARTICLE TEMPERATURE VARIATIONS ALONG THE DUCTED FLOW, REFERENCE CASE 2

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DIAMETER, MICROWS

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FIG. 17 (CONCLUDED) PREDICTED VARIATIONS OF 50 AND 500 MICRON DROP SIZES ALONG THE DUCT FLOW, REFERENCE CASE 2



AXIAL DISTANCE FROM DUCT ENTRANCE, FT

DIAMETER, MICRONS

\$ ф 4 Ø 8 R ጽ Ř 8 AXIAL DISTANCE FROM DUCT ENTRANCE, FT ജ 8 ଞ୍ଚ 8 ผ କ୍ଷ 8 16 14 ង 9 ω 9 4 2 0 0.05 0.10 0.0 0.85 0.80 0.75 0.60 0.55 0.50 0.45 0.40 0.35 0.30 0.25 0.20 0.15 0.95 0.90 0.70 0.65 1.00

ATER FLOW MACH NUMBER

FIG. 18 PREDICTED AXIAL DISTRIBUTION OF AIR FLOW MACH NUMBER ALONG DUCTED FLOW



FIG. 19 PREDICTED AXIAL VARIATIONS OF WATER PARTICLE SIZES ALONG WIND TUNNEL FLOW

DIAMETER, MICRONS

-



FIG. 19 (CONCLUDED) PREDICTED AXIAL VARIATIONS OF WATER PARTICLE SIZES ALONG WIND TUNNEL FLOW



TEMPERATURE, DEG R





TEMPERATURE, DEG R

1



PREDICTED VARIATIONS OF WATER PARTICLE VELOCITIES ALONG DUCT FLOW

FIG. 21

VELOCITY, FT/S



FIG. 21 (CONCLUDED) PREDICTED VARIATIONS OF WATER PARTICLE VELOCITIES ALONG DUCT FLOW

<u>APPENDIX B</u>

MASS AND ENERGY BALANCES FOR TWO-PHASE PARTICLE DURING FREEZING

B.1 The particle initially starts out as a two-phase mixture of water and ice in equilibrium. At the end of an integration step (or time, δt), the particle is still two-phase with a different ratio of ice mass to liquid mass because of freezing and evaporation that occurred during the period, δt . Heat transfer and mass transfer effects have changed the particle phase composition. Let the subscripts "T" refer to the solid or ice phase, and "L" refer to the liquid phase.

B.1.2 Particle Energy Change

The initial particle energy, E, at time t, is

$$E = M_{I}h_{I} + M_{I}h_{I} \tag{1}$$

The final particle energy at time $t + \delta t$ is

$$E' = M_I \dot{h}_I + M_L \dot{h}_L$$
(2)

The energy change is given by

$$\Delta E = E' - E = M_{I} \dot{h}_{I} + M_{L} \dot{h}_{L} - M_{I} h_{I} + M_{L} h_{L}$$
(3)

Because the process is isothermal,

 $h_{i} = h_{i} \tag{4}$

$$h_L = h_L \tag{5}$$

(6)

Hence,

$$\Delta E = \left(M_{I}' - M_{I} \right) h_{I} + \left(M_{L}' - M_{L} \right) h_{L}$$

or

$$\Delta E = \Delta M_{I} h_{I} + \Delta M_{L} h_{L}$$

B.1.3 Liquid Mass Change by Evaporation

Evaporation of liquid mass occurs during the period, δt , in the amount

$$\delta M_{L} = \dot{M}_{e} \dot{A} \delta t \tag{8}$$

where A is the surface area of the particle.

For a moving particle

$$\delta t = \frac{dx}{V} \tag{9}$$

where V is the particle velocity. Hence,

$$\delta M_{L_{\epsilon}} = \frac{\dot{M}_{\epsilon} A dx}{V}$$
(10)

The mass flux of evaporation is given by

$$\dot{M}_{e}^{"} = h_{m} \left(\frac{Y_{s} - Y_{\infty}}{1 - Y_{s}} \right) \tag{11}$$

where

$$h_m = \frac{h}{c_{p_\ell} L_e^m} \tag{12}$$

and

h = average convective heat transfer coefficient for the particle

 L_{e} = the Lewis number for the particle "film" conditions

 c_{p_I} = the specific heat of the "film" fluid

Y = mass fraction of vapor phase

 $Y_{\rm r}$ = mass fraction at particle-gas "film" interface

 Y_{∞} = free stream mass fraction of vapor.

(7)

Thus,

$$\delta M_{L_{\epsilon}} = \frac{h A dx}{c_{p_f} L_{\epsilon}^m V} \left(\frac{Y_s - Y_{\infty}}{1 - Y_s} \right)$$
(13)

B.1.4 Total Liquid Mass Change

The total liquid mass change, ΔM_L , is thus in two parts: a part due to freezing of the liquid, with subscript L_F and a part due to evaporation, with subscript L_F . Hence,

$$-\Delta M_{L} = \delta M_{L_{E}} + \delta M_{L} \tag{14}$$

Clearly, however,

$$\delta M_{I_{\star}} = \Delta M_{I_{\star}} \tag{15}$$

Thus,

$$-\Delta M_{L} = \Delta M_{I} + \delta M_{L} \tag{16}$$

Solving,

$$\Delta M_{L} = -\Delta M_{I} - \delta M_{L} \tag{17}$$

B.1.5 Energy Balance Revisited

The energy balance can now be written

$$\Delta E = \Delta M_{I} h_{I} + \Delta M_{L} h_{L} \tag{18}$$

or, by substitution

$$\Delta E = \Delta M_{I} h_{I} + \left(-\Delta M_{I} - \delta M_{L_{c}} \right) h_{L}$$
⁽¹⁹⁾

$$\Delta E = \Delta M_{I} (h_{I} - h_{L}) - \delta M_{L} h_{L}$$
⁽²⁰⁾

B.1.6 Heat Balance with Freestream

The energy balance for the particle must equal the energy lost or gained by the particle to the freestream. Thus,

$$\Delta E = -\dot{q}^{*}A\delta t - \dot{M}_{e}^{*}Ah_{v}\delta t \qquad (21)$$

where

$$\dot{q}'' = h(T_s - T_{\omega}) \tag{22}$$

and

h_v = the enthalpy of the vapor leaving the droplet.

Using previously defined parameters in equations (10), (11), (13) and (21), the energy flow from the particle is given by

$$\Delta E = -\frac{hA\,dx}{V} \left(T_s - T_{\infty}\right) - \frac{hA\,dx\,h_v}{c_{p_s} L_e^m V} \left(\frac{Y_s - Y_{\infty}}{1 - Y_s}\right) \tag{23}$$

B.1.7 Energy Equation Equality

The two energy equations for the particle can thus be combined by setting equation (20) equal to equation (21)

$$\Delta E = \Delta M_{I} (h_{I} - h_{L}) - \delta M_{L_{e}} h_{L} = -\dot{q}^{"} A \,\delta t - \dot{M}_{e}^{"} A \,h_{v} \,\delta t \tag{24}$$

Using equation (8) in equation (24)

$$\Delta E = \Delta M_{I} (h_{I} - h_{L}) - \delta M_{L_{e}} h_{L} = -\dot{q}^{*} A \delta t - \delta M_{L_{e}} h_{\nu}$$
⁽²⁵⁾

Rearranging this equation

$$\Delta M_{I}(h_{I}-h_{L}) = -\dot{q}^{"}A\delta t - \delta M_{L}(h_{v}-h_{L})$$
⁽²⁶⁾

with

$$h_{\nu} - h_{L} = h_{fe} \tag{27}$$

where

 h_{fg} = the enthalpy of evaporation of water

and

 $h_I - h_L = -H_F$

where

 $H_{\rm F}$ = the enthalpy of melting of ice

the energy balance becomes

$$\Delta M_{I} = \frac{\dot{q}^{\prime\prime} A \delta t}{H_{F}} + \frac{\delta M_{L_{e}} h_{fg}}{H_{F}}$$
(29)

Using equations (9), (22) and (13) for δt , $\dot{q}^{"}$, and δM_{L} ,

$$\Delta M_{I} = \frac{hAdx}{VH_{F}} \left(T_{s} - T_{\infty}\right) + \frac{hAh_{fg}dx}{Vc_{p_{f}}L_{e}^{m}H_{F}} \left(\frac{Y_{s} - Y_{\infty}}{1 - Y_{s}}\right)$$
(30)

Thus, the amount of ice produced is given by this expression. In words, the amount of ice produced equals the heat lost by convection and evaporation, divided by the heat of fusion of water.

B.2 Method for Adjusting Phase Masses During Freezing Process

Define the initial amount of mass of the particle as M_{ϕ} , where

$$M_{\phi} = \frac{\pi}{6} \rho_{\phi} D_{\phi}^{3} = M_{L} + M_{I}$$
(31)

and the subscript " ϕ " refers to the initial values of the parameters before the integration step. After computing ΔM_1 and δM_{L_e} during the integration step, then determine the new values of the particle mass M_n as follows.

 $M_n = M_I + M_L$

$$M_{n} = M_{\phi} - \delta M_{L_{e}} \tag{32}$$

(28)

(33)

$$M_I = M_I + \Delta M_I \tag{34}$$

thus,

and

$$M_{L} = M_{n} - M_{I} = M_{L} - \delta M_{L_{e}} - \Delta M_{I}$$
 (35)

The mass fraction of the liquid phase is then given by

$$\alpha = M_{L}^{\prime} / M_{n} \tag{36}$$

B.3 Particle Properties

The particle properties for the two-phase mixture particle of ice and liquid water are given in terms of α as follows:

average density

$$\overline{\rho} = \frac{\rho_L \rho_I}{\alpha \rho_I + (1 - \alpha)\rho_I} \tag{37}$$

average specific heat

$$\overline{c}_{p} = \alpha c_{PL} + (1 - \alpha) c_{Pl} \tag{38}$$

average temperature

$$\overline{T} = T_L = T_I = 492^{\circ} R$$
 (isothermal assumption) (39)

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<u>APPENDIX C</u>

FORTRAN DATA INPUT FORMATS ("CARD IMAGES")

The following data input statements identify the data parameters that are input, their input order, and the ("card image") format that is required for their input. The term LCR is an acronym from the "old days" meaning line card reader. It is taken care of in the code compiling process, and is of no concern to users of the code.

```
READ(LCR,7) NXY,NUNITS,NCASES
READ(LCR,707)NFREZ
READ(LCR,8) (XS(I),YS(I),I=1,NXY)
READ(LCR,9) NSTA,XPRINT,DX0
READ(LCR,10) (ALP(I),I=1,10)
NS1=NSTA+1
READ(LCR,12) (STA(I),I=1,NS1)
READ(LCR,12) (STA(I),I=1,NS1)
READ(LCR,12) (FL(I),I=1,NSTA)
READ(LCR,12) (VL(I),I=1,NSTA)
READ(LCR,12) (TS(I),I=1,NSTA)
READ(LCR,12) (DMICRON(I),I=1,NSTA)
```

The actual input formats for the data are given below. They are extremely important because the code AEDC1DMP expects to read in the data in the defined format. Future releases of AEDC1DMP will use user friendly data input formats, but at present, this set of input formats must be used.

C-1

FORMAT(3I2)
FORMAT(1I2)
FORMAT(2E12.0)
FORMAT(112,2E12.0)
FORMAT(10A7)
FORMAT(4E10.0)
FORMAT(8E10.0)
APPENDIX D

SAMPLE DATA INPUT CASE

The data below are for a single data input case. Multiple cases can be run on AEDC1DMP by starting each new data set with card NSTA, XPRINT, DX0 and those that follow, through (DMICRON(I), I=1, NSTA); no spaces between data input sets.

800001	
01	
.0000	1.5788
.0250	1.42511
.0500	1.29351
.0750	1.18135
.1000	1.08625
.1250	1.00599
.1500	0.938571
.1750	0.882179
.2000	0.835174
.2250	0.796079
.2500	0.763576
.2750	0.736492
.3000	0.713789
.3250	0.694557
.3500	0.678003
.3750	0.663443
.4000	0.650297
.4250	0.638074
.4500	0.62637
.4750	0.614858
.5000	0.603283
.5250	0.59145
.5500	0.579224
.5750	0.566517
.6000	0.55329
.6250	0.539538
.6500	0.525291
.6750	0.510605
.7000	0.495562
.7250	0.48026
.7500	0.46481
.7750	0.449333
.8000	0.433958
.8250	0.418814
.8500	0.40403
.8750	0.38973

.90 00	0.376034
.9250	0.36305
.9500	0.350879
.9750	0.339605
1.0000	0.3293
1.0250	0.320021
1.0500	0.311807
1.0750	0.304681
1,1000	0.298648
1 1250	0 293696
1 1500	0.22000
1.1500	0.209790
1.1750	0.2009
1.2000	0.204740
1.2250	0.203037
1.2300	0.203341
1.2750	0.203341
1.3000	0.283541
1.3250	0.283541
1.3500	0.283541
1.3750	0.283541
1.4000	0.283541
1.4250	0.283541
1.4500	0.283541
1.4750	0.283541
1.5000	0.283541
1.5250	0.283541
1.5500	0.283541
1.5750	0.283541
1.6000	0.283541
1.6250	0.283541
1.6500	0.283541
1.6750	0.283541
1.7000	0.283541
1.7250	0.283541
1.7500	0.283541
1.7750	0.283541
1.8000	0.283541
1.8250	0.283541
1.8500	0.283541
2.0000	0.283541
2.2000	0.283541
2.4000	0.283541
2.6000	0.283541
2.8000	0.283541
10 0.025	

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GENERIC NOZZLE GEOMETRY FOR PARAMETRIC FREEZING STUDY 0.0 0.02 0.04 0.06 0.08 0.10 0.12 0.14 2.70 0.16 0.18 0.0002857 18.0 460.0 2116.0 6.85E-05 2.06E-04 4.80E-04 3.62E-04 2.06E-04 3.43E-05 1.37E-05 1.37E-06 1.37E-07 1.37E-08 46.0 46. 46. 46. 46. 46. 46. 46. 46. 46. 543. 543. 543. 543. 543. 543. 543. 543. 543. 543. 60.0 5.0 10.0 20.0 30.0 40.0 50.0 15.0 80.0 100.0

<u>APPENDIX E</u>

DEFINITIONS OF THE DATA INPUT PARAMETERS

The definitions of the data input parameters for AEDC1DMP are the same, for the most part, as those in the original version of the code reported in Reference 53. Their specific definitions, and units, where necessary, are provided below.

NXY= THE NUMBER OF DUCT WALL COORDINATES INPUT

NUNITS=0 OR1. IF NUNITS=0, ALL X AND Y COORDINATES, XPRINT, DX0, AND (STA(I), I=1,NS1) ARE INPUT IN UNITS OF FEET. IF NUNITS IS INPUT AS 1, THEN ALL SPACE COORDINATES ARE INPUT IN UNITS OF INCHES.

NCASES= THE NUMBER OF CASES INPUT TO BE COMPUTED.

NFREZ= 0 OR 1. IF NFREZ=0, THE WATER PARTICLE FREEZING MODEL IS NOT ACTIVATED; ALTERNATELY, FOR NFREZ=1, THE MODEL WILL BE UTILIZED IN THE DUCT FLOW COMPUTATIONS.

(XS(I), I=1,NXY)= THE AXIAL COORDINATES OF THE DUCT WALL CURVE. (YS(I), I=1,NXY)=THE CORRESPONDING RADIAL COORDINATES OF THE DUCT WALL CURVE.

NSTA= THE NUMBER OF WATER INJECTION STATIONS, WITH A MAXIMUM OF TEN(10).

XPRINT= THE DISTANCE BETWEEN AXIAL LOCATIONS WHERE THE FLOW SOLUTION IS PRINTED OUT, AND, CONSEQUENTLY, STORED FOR INPUT TO DATA PLOTTING PROGRAMS SUCH AS WAVEMETRICS(R)' IGOR™ PRO(http://www.wavemetrics.com/).

DX0= THE INITIAL NUMERICAL INTEGRATION STEP SIZE THAT THE USER DESIRES. BY SETTING DX0=0 OR BLANK, THE CODE SELECTS A DEFAULT STEP SIZE BASED ON THE DUCT LENGTH.

ALP= AN 80 COLUMN ALPANUMERIC FILE IDENTIFIER.

(STA(I), I=1, NS1) = THE AXIAL LOCATIONS, OR STATIONS, WHERE WATER IS INJECTED INTO THE FLOW, PLUS, THE LAST VALUE, STA(NS1), WHICH INDICATES THE TERMINAL AXIAL LOCATION FOR FLOW SOLUTION. STA(NS1) USUALLY EQUALS XS(NXY) OR SLIGHTLY LESS.

CV= FLOW INITIAL SPECIFIC HUMIDITY, LBM OF WATER VAPOR/LBM OF DRY AIR.

VG= THE FLOW INITIAL VELOCITY, FT/SECOND.

TG= THE FLOW INITIAL STATIC TEMPERATURE, DEGREES RANKINE.

P= THE FLOW INITIAL STATIC PRESSURE, LBFORCE/SQUARE FOOT.

(FL(I), I=1, NSTA) = THE WATER LOADING INJECTED AT EACH STATION, LBM OF WATER/LBM OF DRY AIR.

(VL(I), I=1,NSTA) = THE VELOCITIES OF THE WATER PARTICLES INJECTED AT THE WATER INJECTION STATIONS, FEET/SECOND.

(TS(I), I=1,NSTA) = THE TEMPERATURES OF THE WATER PARTICLES INJECTED AT EACH STATION, DEGREES RANKINE.

(DMICRONS(I), I=1,NSTA) = THE DIAMETERS OF THE WATER PARTICLES INJECTED AT EACH STATION, IN MICRONS (1 METER=1,000,000 MICRONS)

AS A POINT OF CLARIFICATION, FL(I), VL(I), TS(I), AND DMICRON(I) ALL CORRESPOND TO THE WATER PARTICLE CONDITIONS AT THE AXIAL INJECTION STATION LOCATED AT X=STA(I).

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APPENDIX F

SAMPLE OUTPUT FROM CASE 3

Hard Disk 2150:MPW: AEDC ADVERSE WEATHER 3/98:DEVELEMENTS OF AEDC MFF CODE: AEDCIDMP ...: AEDCIDMP. CUT 3/13/98 5:15 BM

THIS VERSION CONTAINS A DROPLET FREEZEOUT MODEL BASED ON MODIFIED HOMOGENEOUS NUCLEATION THEORY (DESCRIBED IN AEDC TR 73-144) UNDER AEDC-UTSI TASK ORDER 97-03; TOM TIBBALS SVERDRUP TECHNOLOGY, INC., PROJECT MANAGER THIS VERSION DEVELOPED BY R.J. SCHULZ, UTSI, TULLAHOMA, TN 37388: DECEMBER, 1997 A.E.D.C. ONE-DIMENSIONAL, MULTI-PHASE FLOW PROGRAM, AEDCIDMP CALIBRATED TO THE DATA OF DORSCH AND HACKER, NACA TN 2142 DEVELOPED FROM THE CODE OF C.E. WILLBANKS AND R.J. SCHULZ

INITIAL FLOW CONDITIONS AND THEIR UNITS:

TEMPERATURE, TG AND TS, DEG R VELOCITY, VG AND VL, FT/S ; WATER LOAD FACTOR, FL, LBM- LIQUID WATER/LBM-DRY AIR ; SPECIFIC HUMIDITY, CV, LBMS-WATER VAPOR/LBM-DRY AIR ; PRESSURE, P, PSFA ; DROP DIA., D, MICRONS ; YDUCT=DUCT RADIUS, FT

IRT CASE 3 RADIAL: V=40 FT/SEC, T=15 DEG F, 10 DROP SIZES IN THE INLET

475.00 .000181611 2073.60 INLET STATIC TEMPERATURE= 25.997580 INITIAL FLtot= .000500000 INLET STATIC PRESSURE= INLET RELATIVE HUMIDITY %= 40.000000 INLET SPECIFIC HUMIDITY= AIR INLET VELOCITY=

INITIAL X= .000000 INITIAL INTEGRATION STEP SIZE, DX= .000700000

60.00000 80.00000 20.00000 30.00000 40.00000 50.00000 5.00000 10.00000 15.00000 100.00000 占 占 Å ۵ 占 凸 ᆸ 4 ᆸ 495.000000 495.000000 495.00000 495.000000 495.000000 495.00000 195.000000 495.000000 495.000000 495.000000 **TS**= TS= TS= TS= TS= TS= TS= TS= TS= TS= 16.000000 16.000000 16.000000 16.000000 16.000000 16.000000 16.000000 16.000000 16.000000 16.000000 Ч" ۲<u>-</u> ۲<u>-</u> <u>-</u>72 ЧС= ۲<u>-</u> <u>"</u> ₹T= чц. Ч<u>г</u>= 000009060 000063500 .000047900 .000004540 .000001810 00000018 000027300 .000027300 00000000000 .000000181 ARE =<u></u>[] FL= =75 FL= -12 EL= <u>"</u> 믭 -12 = L EACH INJECTION STATION .06000000 .08000000 .180000000 .00000000 .020000000 .04000000 .100000000 .120000000 .14000000 .16000000 $=(1) \Omega X$ =(I) INIX =(I) MIX $=(1) \Omega XIX$ =(I) INIX =(I) INIX $=(1) \Omega X$ =(I) INIX =(1) NIX =(I) MIXSPRAY CONDITIONS AT STATION10 INJECTION STATION 1 STATION 6 STATION 8 STATION 9 STATION 3 STATION 4 STATION 5 STATION 7 STATION INUECTION INJECTION INJECTION INJECTION INJECTION INUECTION INUECTION INJECTION INJECTION

INITIAL AREA= 763.166377

475.133999	5	.000009060 0.144222E-12	475.125494 3	.000175752	0.119113E-12	0.109667E-11	0.379184E-11	0.906085E-11	0.307763E-10	0.731673E-10	0.143160E-09	0.247706E-09	0.588109E-09	0.114985E-08
TO= 01_00	80166.07	FLtot= DMASS=	TO= 26.31155	FLtot=	DMASS=	DMASS=	DMASS=	DMASS=	DMASS=	DMASS=	DMASS=	DMASS=	DMASS=	DMASS=
2075.6258	UM. 8=	2073.60 5.00000	2075.6254 IUM. %=	2073.60	4.69115	9.83222	14.86766	19.87691	29.87890	39.87783	49.87695	59.87860	79.88130	99.88631
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.037269 P	1,00000	475.000000 495.000000	.037296 P	474.991307	469.748029	469.748086	469.853291	470.676583	474.090880	477.535286	480.267645	482.556906	485.638121	487.735676
= DW	A(X)/A0=	=ST	MG= A (X) /A0=	=DT	TS=	TS=	TS=	=ST	=ST	=ST	TS=	TS=	TS=	=ST
15.586000	763.1664	40.000000	15.580455 762.6234	40.028167	40.028233	40.028426	40.028634	40.027198	39.852714	39.094691	37.926089	36.500310	33.853338	31.467418
YDUCT=	A (X) =		YDUC'T= A (X) =	= <u>0</u> /	СЦ =	- IV	-1N	=IV	-1N	="N	-17 	ΔT=	NL=	νL=
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3/13/98 5:15 PM Hard Disk 2150:MPM:AEDC ADVERSE WEATHER 3/98:DEVELEMENTS OF AEDC MFF CODE:AEDCIDMP ...: AEDCIDMP.OUT

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	475.106282)	.000171396 0.972066E-13 0.104946E-11 0.371986E-11 0.896241E-11 0.306022E-10 0.728614E-10 0.142657E-09 0.246928E-09 0.586558E-09 0.586558E-09 0.114719E-08	475.087051 1	.000167154 0.766273E-13 0.100244E-11 0.364765E-11 0.3846501E-11 0.304482E-10 0.726296E-10 0.142302E-09 0.246399E-09 0.246399E-09 0.114544E-08 0.114544E-08	.000163126 0.580697E-13 0.956984E-12 0.357724E-11 0.876996E-11 0.303023E-10 0.724218E-10 0.142005E-09 0.245978E-09 0.584743E-09 0.114412E-08	475.050919 5 .000159231 0.413439E-13 0.912201E-12 0.350702E-11
	TO= 26.55970(Fl.tot= DMASS= DMASS= DMASS= DMASS= DMASS= DMASS= DMASS= DMASS= DMASS=	TO= 26.80234	FLtot= DMASS= DMASS= DMASS= DMASS= DMASS= DMASS= DMASS= DMASS= DMASS= DMASS= DMASS= DMASS= DMASS= DMASS=	FLtot= DMASS= DMASS= DMASS= DMASS= DMASS= DMASS= DMASS= DMASS= DMASS= DMASS=	TO= 27.25809 FLLtot= DMASS= DMASS= DMASS=
	2075.6254 UM. %=	2073.60 4.38387 9.68905 14.77298 19.82246 29.82246 39.82248 49.81580 79.81103 99.80941	2075.6255 UM. %=	2073.60 4.04968 9.54214 14.67676 19.77234 29.77991 49.777991 49.77797 59.77310 79.76444 99.75864 99.75864	2073.60 3.69212 9.39566 14.58171 19.66188 29.72471 39.74192 59.73900 79.72861 99.72003	2075.6257 LUM. %= 2073.59 3.29683 9.24676 14.48567
	C= SEL H		O= REL H			
•	.037268 P	474.972296 469.750619 469.750736 469.751736 469.798763 469.859602 473.042335 477.450015 480.746702 483.188061	.037241 P	474.953266 469.752553 469.752738 469.752738 469.755634 469.755634 470.039948 471.271519 471.271519 471.271519 471.971382 477.971382 477.971382 480.471680 037268 P	474.934728 469.754228 469.754228 469.755426 469.755426 469.831073 470.471381 471.725164 473.242179 476.148065 478.604354	.037377 P .99705 474.916164 469.755338 469.755334 469.757301
	MG= A (X) /A0=	13 13 14 15 15 15 15 15 15 15 15 15 15 15 15 15	MG= A(X)/A0=	IG= TS= TS= TS= TS= TS= TS= TS= TS= TS= TS	= ST = ST = ST = ST = ST = ST = ST = ST	MG= A(X)/A0= TG= TS= TS= TS=
	15.586023 763.1686	39.998143 39.998347 39.998347 40.000448 40.002236 40.004206 39.935985 39.157493 37.809281 36.306763	15.591582 763.7132	39.968286 39.968332 39.968567 39.968910 39.973143 39.973143 39.973143 39.913617 39.913617 39.013419 39.013419 38.027579 38.027579 38.027579	39.995972 39.995851 39.994442 39.9944442 39.989520 39.989520 39.989520 39.964559 39.895277 39.895277 39.895277 38.833277	15.563031 760.9187 40.113036 40.112750 40.108290 40.101820
	YDUCT= A (X) =	VL= VL= VL= VL= VL= VL= VL= VL= VL= VL=	YDUCT= A(X) =	VG= VL= VL= VL= VL= VL= VL= VL= VL= VL= VL		YDUCT= A(X) = VG= VL= VL=
	1.002403 .007000000	.000510205 .000006106 .000024832 .000060666 .0000246510 .0000026818 .000001480 .000001790 .000000179 .00000018	1.506403	.000514443 .000004814 .000023719 .000059483 .000046005 .0000046683 .000001465 .000000179 .000000179 .000000018 .000000028 .003403	.000001782 .000003648 .000003648 .0000022644 .0000022644 .0000025555 .000001782 .0000001782 .000000179 .000000179 .00000018	2.501850 .005638855 .000522358 .000002597 .000021584
	X= DX=	CV= FI= FI= FI= FI= FI=	X= DX=	A C C C C C C C C C C C C C C C C C C C	- ADE EEEEEEEEEEEEEEEEEEEEEEEEEEEEEEEEEE	X= DX= FI= FI=

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3/13/98 5:15 PM Hard Disk 2150:MEW:AEDC ADVERSE WEATHER 3/98:DEVELEMENTS OF AEDC MFF CODE: AEDCIDMP ...: AEDCIDMP.GUT

0.867455E-11 0.301556E-10 0.722169E-10 0.141726E-09 0.245599E-09 0.584068E-09 0.584068E-09 0.114300E-08	475.033819 4	.000155486 0.2670508-13 0.868218E-12 0.343709E-11 0.85877E-11 0.30068E-10 0.720100E-10 0.141449E-09 0.245232E-09 0.583447E-09 0.583447E-09 0.583447E-09	475.017663 6	.000151949 0.145737E-13 0.825545E-12 0.336821E-11 0.848361E-11 0.288568E-10 0.718003E-10 0.141169E-09 0.582842E-09 0.582842E-09 0.114105E-08	475.002683 0	.000148669 0.538430E-14 0.784468E-12 0.330096E-11 0.838998E-11 0.297074E-10 0.715896E-10 0.140887E-09 0.244496E-09 0.582240E-09 0.582240E-09 0.114012E-08
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3/13/98 5:15 FM Hard Disk 2150:MPW: AEDC ADVERSE WEATHER 3/98:DEVELIMENTS OF AEDC MFF CODE: AEDCIDMP ...: AEDCIDMP. CUT

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474.989305 2	.000145741 0.153680E-15 0.745423E-12 0.3236078-11 0.829891E-11 0.295693E-10 0.14066E-09 0.14129E-09 0.244129E-09 0.581643E-09 0.581643E-09 0.113920E-08	474.977925 3	.000143248 0.708557E-12 0.317395E-11 0.821116E-11 0.294171E-10 0.711744E-10 0.11744E-09 0.243768E-09 0.281055E-09 0.113830E-08 0.113830E-08	.000140881 0.673733E-12 0.311456E-11 0.812686E-11 0.292786E-10 0.709743E-10 0.140058E-09	0.243415E-09 0.580484E-09 0.113742E-08 474.957027 0	.000138672 0.641544E-12 0.305898E-11 0.804760E-11 0.291477E-10 0.707845E-10 0.139801E-09
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3/13/98 5:15 FM Hard Disk 2150:MFW: AEDC ADVERSE WEATHER 3/98: DEVELEMENTS OF AEDC MFF CODE: AEDCITMP ...: AEDCITMP. CUT

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-	5= 0.243079E-09 5= 0.579945E-09 5= 0.113660E-08	474.947359)749	 .000136554 0.610983E-12 0.300557E-11 0.797106E-11 0.290206E-10 	<pre>>= 0.705995E-10 >= 0.705995E-10 >= 0.139550E-09 >= 0.242751E-09 >= 0.579420E-09 >= 0.579420E-09</pre>	474.938351)572	<pre>0001345810.582781E-12 - 0.295566E-11 - 0.28907E-11 - 0.28907E-10 - 0.78923E-10 - 0.139311E-09 - 0.139311E-09 - 0.278924E-09 - 0.13505E-08</pre>	474.929727 9491	 .000132692 0.556040E-12 0.290774E-11 0.280774E-11 0.78293E-11 0.78293E-11 0.78293E-10 0.287843E-10 0.287843E-10 0.287843E-10 0.281745 0.139078E-09 0.13432E-09 0.113432E-08 	474.921703 4525 = .000130934 S= 0.5314048-12 S= 0.2863048-11 S= 0.7764978-11
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	48.475664 47.560113 46.712606	13.554534 577.1904	52.903465 52.844621 52.736694 52.592599 52.218604	51.766881 51.273484 50.763666 49.7556195 48.814411	13.230057 549.8868	55.536401 55.472082 55.350648 55.189073 54.772454 54.273834 53.177653 53.177653 51.053341	12.901450 522.9099	58.409065 58.338447 58.201307 58.019540 57.554110 57.002091 57.002091 56.408501 55.801943 55.801943 53.495855 53.495855	12.577403 496.9719 61.466625 61.388770 61.233472 61.028521
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Hard Disk 2150: MPW: AEDC ADVERSE WEATHER 3/98: DEVELEMENTS OF AEDC MEF CODE: AEDCIDAP ...: AEDCIDAP. CUT 0.696425E-10 0.138435E-09 0.241292E-09 0.577104E-09 0.466669E-12 0.274265E-11 0.758827E-11 0.283723E-10 0.138237E-09 0.241032E-09 0.576691E-09 0.241845E-09 0.285696E-10 0.699364E-10 0.577538E-09 0.113297E-08 0.486646E-12 0.278030E-11 0.284680E-10 0.697853E-10 0.113171E-08 0.286746E-10 0.700917E-10 0.138856E-09 0.577981E-09 0.138643E-09 0.241566E-09 0.764383E-11 0.113232E-08 0.113363E-08 0.508398E-12 0.282077E-11 0.770324E-11 .000124819 .000126221 .000129276 .000127691 474.906909 474.900205 474.893815 474.914136 29.490125 29.630342 29.218605 29.354661 DMASS= FLtot= FLtot= FLtot= FLtot= å å Ę å 59.40254 79.42007 29.18265 39.31105 2075.6260 2075.6260 39.25370 2075.6260 18.73585 39.22689 79.36098 99.35832 7.49946 18.78146 49.29855 59.33587 49.37205 2075.6260 7.60957 13.47148 18.83000 29.14702 39.28200 49.34675 59.37964 13.40673 29.11241 49.32203 79.37990 7.39541 13.34594 29.07977 59.35721 99.37627 79.39977 2068.20 99.41451 99.3951 2070.31 2069.70 2069.01 HUM. %= HUM. S= REL HUM. %= HUM. %= <u>ا</u> م 占 占 占 씸 씸 444 占 REL 出出 REL 씸 占 REL ii A ᇤ 塭 å å å ቘ 469.690675 469.711943 469.757095 469.539345 469.552296 469.566219 469.594495 469.678190 469.585604 469.598013 469.647344 469.670345 469.704798 470.507689 469.645841 469.875054 470.383313 469.648875 469.670610 469.692388 469.730914 470.002904 470.653685 469.623362 469.934194 469.670840 470.082954 469.604705 469.614963 469.574101 469.621061 470.822451 469.626065 474.402152 474.562732 474.514906 474.462261 .58527 .52286 .55385 .61803 .063776 .071433 .060380 .067413 A(X)/A0= A(X)/A0= A(X)/A0= A (X) /A0= TS= TS= ΞST TS= TS=TS= TS= TS= TS# TS= ₽ ъ В **"** ង្ក ľ "B ¦¦ = DW ۳ ۳ Ë Ш Ш 70.093559 69.219370 68.415869 68.319864 65.647418 63.328825 72.206573 71.976138 71.676709 70.935815 68.350304 76.620454 63.696868 67.854735 67.196563 66.439895 64.854147 66.690963 59.896116 59.243963 58.581957 57.290776 64.776279 64.690031 64.513432 64.281456 63.017484 60.173987 11.923750 446.6585 65.167229 60.507740 56.085222 62.299511 61.575751 68.118231 61.919974 72.313731 58.871783 11.270093 12.252924 11.599271 422.6796 399.0294 471.6604 YDUCT-YDUCI-YDUCT= -LODOR ۲<u>-</u> A.(X)= =**T** ۲<u>–</u> Ч**г** A (X) = A (X) = ۲<u>-</u> ЧL= =75 g ۲<u>1</u>2 ۲<u>5</u> =75 =75 <u>=</u>Д -77= SL= **"** 5 55 ۲<u>-</u> 5 ۳<u>۲</u> 5 "B Ľ =7N <u>=</u>77 =75 27= A (X) = **"** ЧЧ-<u>-</u>2 <u>-</u>17 15 25= 8 24 .00700000 .000000000 000011515 000000175 .000039379 .000004282 .000556731 000000000 .000553863 000045339 000024948 00000018 000000000000 .000011042 .000044725 .000024864 000001735 000000175 00000018 000000000 .000004309 .000000176 00000018 0000000000 000552280 000045999 000025037 000004300 000001740 00000018 000039667 000001737 000025129 .000001743 000000000000 000012029 000039976 000000175 000000000000 000004291 .000555331 10.005321 9.501321 9.004321 8.500321 3/13/98 5:15 PM 밀밀 =]L 빌 빌 ЫL X= DX= Ē 믭 믭 22 =XC S =14 밀 EL. 빌 11--12 믭 = X = X = X =<u></u>____ =15 FL= Ш. -14 =15 **E**L= 빌딩 = X0 2 = ŝ Ш. Ē E. 믭 -15 믭 튑 ×

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... AEDCIDMP. OUT 0.280300E-10 0.691249E-10 0.137515E-09 0.240075E-09 0.575166E-09 0.399494E-12 0.261201E-11 0.112945E-08 0.137684E-09 0.240299E-09 0.739292E-11 0.575525E-09 Hard Disk 2150: MFW: AEDC ADVERSE WEATHER 3/98: DEVELEMENTS OF AEDC MFF CODE: AEDCIDMP 0.414324E-12 0.281088E-10 0.692451E-10 0.112998E-08 0.743736E-11 0.281928E-10 0.693724E-10 0.137862E-09 0.240536E-09 0.575902E-09 0.113054E-08 0.264146E-11 .000121146 0.282798E-10 0.695036E-10 0.138044E-09 0.240777E-09 0.576287E-09 0.430562E-12 0.267328E-11 0.748510E-11 0.447825E-12 0.753491E-11 0.113111E-08 0.270667E-11 .000122286 474.872242 .000123521 474.877091 474.882279 474.887900 30.083472 29.924874 29.773230 DMASS= FLtot= DMASS= DMASS= DMASS= DMASS= DMASS= DMASS= DMASS= DMASS= DMASS= å FLtot= ģ FLtot= å ģ 59.25726 79.29096 13.17974 18.61082 28.98945 39.15213 13.13058 18.57368 79.30745 99.30771 2075.6258 49.21255 13.23246 18.65055 28.96233 39.12948 79.32476 99.32407 49.23266 59.27570 2075.6259 7.02203 .29213 2075.6259 7.10786 29.01830 39.17612 2075.6260 49.25387 59.29512 18.69183 29.04812 39.20080 49.27562 59.31499 79.34243 7.19953 2064.92 7.29449 13.28733 99.34077 2066.18 2067.27 REL HUM. 8= REL HUM. &= REL HUM. 8= Å ᆸᆸᆸ 占 占 씸 ۵ 出 88888 씸 Å 盀 出 占 ᇤ å 붪 å å å 469.333190 469.355663 469.477544 469.515030 469.548520 469.582319 469.416136 469.436702 469.672389 469.542525 470.003185 469.468585 469.506233 469.770792 470.175393 469.379371 469.426077 469.396746 469.721410 470.085576 469.529057 469.560828 469.486647 469.522519 469.555719 469.617924 469.468726 469.585697 469.617583 469.649278 469.820869 470.272708 469.451911 469.513305 469.590462 474.168279 474.256790 469.498594 .43595 474.334833 .46388 .49314 .091523 .085789 .080582 .075767 A (X) /A0= A (X) /A0= A(X)/A0= "TS= TS= **T**S= TS= TS= TS= TS= ¶3= TS= TS= ±ST =ST TS= TS= "ST TS= **TS**= TS= TS= TS= TS= TS= TS= TS= TS= TS₌ TS= ¦¦ TS= ۳ ۲ HOM M ¦p ÿ HGH HG= 88.350422 87.033757 83.400509 81.282528 89.669969 85.761720 91.410615 90.892863 83.206129 91.821077 73.147254 72.192784 70.383401 68.731778 85.462050 84.384600 82.018637 78.704480 91.997002 76.499938 76.234897 75.892717 75.054942 74.114243 76.311758 74.338012 86.420950 86.266514 85.912471 80.863395 76.757092 9.966362 81.128130 80.822498 80.430656 79.482034 78.430577 547057 81.264172 77.360421 10.290903 10.615443 10.945054 332.7031 354.0186 376.3446 YDUCT= A(X)= YDUCT-**5** =<u></u>7 YDUCT= <u>"</u> 5 "T <u>"</u>П 25 5 ЧL= 24 4 "B <u>-</u>12 A (X) = g =**1**2 5 **"** 5 ч**г**g g 25 <u>-</u>Д Ч**Г**= =75 ЧЧ-17 = 5 =75 чL= Ľ, ЧГ= A (X)= <u>المع</u> =75 217= чL= 4 <u>-17</u> 000000174 .00000018 .000000002 11.503321 .000038365 .000024564 000004250 000001726 00000018 000009453 .000042594 11.006321 .007000000 .000043075 .000038596 .000024633 000001728 .000000175 000000000 000560401 .000004257 .000044138 .000024783 000000000 000559262 000009803 .000004265 000001730 00000018 .007000000 000010188 000038844 .000024707 000000175 .000001732 .000000175 .00000018 000000000 000558028 .000043594 000010596 12.000321 10.502321 Ma 5:15 3/13/98 =<u>П</u> =11 FL= EL= 믭 = X= 믭 빌 밀 믭 Ē 팊 =<u>1</u>= Ë Ľ 빌 Ē 8 DX= EL= =12 × EL= FL= 립 불 FL= 티= = XC 밀밀 造 믭 밀답 ЕĽ EL= 믭 FL= P B ×

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3/13/98 5:15 PM Hard Disk 2150:MPM:AEDC ADVERSE WEATHER 3/98:DEVELEMENTS OF AEDC MFF CODE:AEDCITMP ...: AEDCITMP. CUT

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1	.000120079 0.385775E-12 0.258440E-11 0.735101E-11 0.279551E-10 0.690101E-10 0.137353E-09 0.239859E-09 0.574820E-09 0.574820E-09 0.112894E-08	474.867657 1	.000119068 0.372938E-12 0.255821E-11 0.731098E-11 0.278829E-10 0.688989E-10 0.688989E-10 0.137196E-09 0.239649E-09 0.574481E-09 0.574481E-09 0.574481E-09	474.863452 5	.000118139 0.361285E-12 0.253409E-11 0.727388E-11 0.278154E-10 0.687942E-10 0.687942E-10 0.137047E-09 0.233449E-09 0.574158E-09 0.574158E-09 0.574158E-09	474.859497 4 .000117262 0.350433E-12 0.251132E-11 0.723858E-11 0.277505E-10 0.686931E-10 0.686931E-10 0.136903E-09 0.239255E-09 0.573843E-09
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A (X) /A0=	1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1	MG= A(X)/A0=	16 	MG= A(X)/A0=	1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1	MG= A(X)/A0= TG= TS= TS= TS= TS= TS= TS= TS= TS= TS= TS
312.0493	98.135393 97.933485 97.454776 96.856745 95.464082 93.983409 92.521148 91.118424 91.118424 88.532009 86.224699	9.636871 291.7575	105.024232 104.790253 104.227513 103.532143 101.937660 100.268955 98.638723 97.086275 94.243285 91.721116	9.311708 272.4009	112.567177 112.294775 111.631739 110.822154 108.995527 107.114033 105.295364 103.575766 100.447451 97.687045	8.981587 253.4289 121.098865 120.778548 119.991115 119.931145 119.939564 114.8099569 112.772781 110.860668 107.404828
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.00700000	.0000561467 .000009128 .0000042144 .0000038148 .0000024498 .000001724 .000000174 .000000174	12.504321 .007000000	.000562477 .00008824 .000081717 .000037940 .0000037940 .000001722 .0000001722 .000000174 .000000174	13.001321 .007000000	.000563405 .000008548 .0000041324 .000037748 .000024376 .0000024376 .000001720 .000000174 .000000174	13.505321 .00700000 .000564280 .000008292 .0000037564 .0000037564 .0000024319 .00000024319 .0000001718 .000000174
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3/13/98 5:15 PM Hard Disk 2150:MFW: AEDC ADVERSE WEATHER 3/98:DEVELPMENTS OF AEDC MFF CODE: AEDCIDMP ...: AEDCIDMP.OUT

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3= 0.112749E-08	474.855893 5416	 .000116460 0.340638E-12 0.249045E-11 0.220598E-11 0.276900E-10 0.276900E-10 0.136767E-09 0.136767E-09 0.137764E-09 0.112704E-08 474.852532 	1470	 .000115708 .000115708 .3315755-12 .3315755-12 .231575112-11 .7175115-11 .7175115-11 .7175115-12 .2763215-10 .1366355-09 .1366355-09 .1366355-09 .1366355-09 .126615-09 .1126615-09 	474.849502 4972	 000115023 000115023 000115023 0.323453E-11 0.245299E-11 114675E-11 114675E-11 275783E-10 275783E-10 275783E-10 13611E-09 238724E-09 238724E-09 238724E-09 238724E-09 112619E-08 112619E-08 	474.846753 2914	= .000114394 S= 0.316103E-12 S= 0.243655E-11 S= 0.712039E-11 S= 0.275277E-10
DMASS	TO= 31.166	FLtot= DMASS DMASS DMASS DMASS DMASS DMASS DMASS DMASS	31.491	FLtot= DMASS DMASS DMASS DMASS DMASS DMASS DMASS	TO= 31.87/	FLLOT DMASC DMASC DMASC DMASC DMASC DMASC DMASC	TO= 32.34	FLtot DMAS DMAS DMAS DMAS
99.23457	2075.6253 IUM. %=	2054.12 6.65871 12.92364 18.41579 28.84475 39.02982 49.12314 59.17454 79.21631 99.22149	TUM. 8=	2050.46 6.59912 12.88964 18.38945 28.82463 39.01246 49.10737 59.15981 79.20286 99.20870	2075.6247 HUM. &=	2046.04 6.54480 12.85852 18.36519 28.99621 49.09253 59.145911 79.19011 99.19655	2075.6243 HUM. %=*	2040.59 6.49485 12.82973 18.34258 28.78829
4	O= REL H		RELI	_╣ ╘╏╘┇╘┇	PQ=	_╬ ╝╝┙┙┙	PO= REL	^{ਸ਼} ਜ਼ਜ਼ਜ਼ਜ਼ ੶
469.689052	.121799 I	473.429259 468.785363 468.8840361 468.895703 468.9985703 469.087155 469.087155 469.22794 469.605009 469.605009		473.182016 468.599483 468.667495 468.735001 468.858572 468.963721 469.052408 469.128387 469.276570 469.512788	.143026 .26355	472.884430 468.375367 468.460391 468.543389 468.543389 468.923859 468.923059 469.011941 469.176075 469.412258	.155785 .24252	472.517168 468.098601 468.206515 468.309819 468.492449
=SI	MG= A(X)/A0=	1G= 17S= 17S= 17S= 17S= 17S= 17S= 17S= 17S	A(X) /A0=	TG TG TG TG TG TG TG TG TG TG TG TG TG T	MG= A (X) /A0=	10 10 10 10 10 10 10 10 10 10 10 10 10 1	MG= A(X)/A0=	=91 =S1 =S1 =S1 =S1
104.371604	8.656054 235.3910	130.512927 130.134563 129.197538 128.085122 128.085122 123.251695 120.969544 118.019278 111.681330 111.681330	217.8086	141.227373 140.775951 139.653201 138.342204 135.543628 135.543628 132.803421 130.237467 127.860529 123.616413 129.927944	8.001397 201.1322	153.171608 152.627915 151.274671 149.723709 149.723709 143.364645 143.364645 143.364645 143.364645 137.816848 133.096509 129.013222	7.675488 185.0810	166.771847 166.109239 164.464064 162.617366 158.843766
±	YDUCT= A (X) =	- EIN EIN EIN SALE EIN A EIN A EIN A EIN A EIN A A A A A A A A A A A A A A A A A A A	A(X) =	=10 =10 =10 =10 =10	YDUCT= A (X) =	=20 =20 =20 =20 =20 =20 =20 =20 =20 =20	YDUCT= A (X) =	=11 =11 M
.000000002	14.002321 .007000000	.000565081 .00008060 .00008060 .000037395 .000037395 .0000037395 .0000037395 .0000001716 .000000174 .000000174 .000000018	.007000000	.000565833 .00007846 .000040292 .000037235 .000024215 .000004212 .000004212 .000001715 .000000174 .000000174	15.003321 .00700000	00056517 00007653 00004001 000037088 0000024168 000004207 00000173 000000173	15.500321 .007000000	.000567145 .000007479 .000039733 .000036951 .000024124
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3/13/98 5:15 PM Hard Disk 2150:MEW: AEDC ADVERSE WEATHER 3/98: DEVELAMENTS OF AEDC MFF CODE: AEDCIDMP ...: AEDCIDMP.CUT

0.136393E-09 0.238563E-09 0.572709E-09 0.112579E-08	474.844258 57	.000113809 0.309404E-12 0.242129E-11 0.709567E-11 0.274796E-11 0.274796E-10 0.238408E-09 0.136280E-09 0.272450E-09 0.572450E-09 0.112540E-08	474.842102 14	.000113285 0.3035128-12 0.2407598-11 0.7073198-11 0.2743538-10 0.6819008-10 0.6819008-10 0.1361748-09 0.2382618-09 0.2382618-09 0.5722058-09 0.5722058-09	474.840250 51	.000112806 0.298241E-12 0.239507E-11 0.705240E-11 0.273936E-10 0.681213E-10 0.136072E-09 0.381213E-09 0.38121E-09 0.571969E-09 0.571968E-08	474.838774 38 .000112385 0.293694E-12	0.238405E-11
DMASS= DMASS= DMASS= DMASS= DMASS=	TO= 32.93375	Fl.tot= DMASS= DMASS= DMASS= DMASS= DMASS= DMASS= DMASS= DMASS= DMASS=	TO= 33.68711	Fltot= DMASS= DMASS= DMASS= DMASS= DMASS= DMASS= DMASS= DMASS= DMASS= DMASS=	TO= 34.64976	FLtot= DMASS= DMASS= DMASS= DMASS= DMASS= DMASS= DMASS= DMASS= DMASS=	TO= 35.8178; FLtot= DMASS=	DMASS=
49 .07840 59.13261 79.17786 99.18485	2075.6238 JM. %=	2033.69 6.44864 12.80290 18.32133 28.97152 38.96602 49.06477 59.11975 59.116593 99.17343	2075.6230 UM. %=	2024.95 6.40744 12.77869 18.30196 28.75604 38.95226 49.05201 59.10765 79.15466	2075.6220 UM. %=	2014.00 6.37013 6.37013 12.75650 18.28401 28.74461 28.75461 28.75919 49.39319 49.15206 99.15206	2075.6208 UM. %= 2001.11 6.33759	12.73691
┇╻┇	C REL H	_╬ ╩╬╬╬	C= REL H	[≝] ª₽₽₽₽₽₽₽₽₽	C= REL H	^r a a a a a a a a a a a a a a a a a a a	а Бала Н Н Н Н Н Н Н Н Н Н Н Н Н Н Н Н Н Н	4
468.768138 468.872985 469.058062 469.298945	.170628 P	472.052238 467.748268 467.887591 468.018124 468.244297 468.428011 468.578289 468.703333 469.166907	.187825 P .20248	471.462934 467.305311 467.489592 467.657278 468.166585 468.166585 468.349016 468.349016 468.746810 468.746810 468.746810	.207484 P .18414	470.722869 466.985496 466.985496 467.201479 467.559911 467.840293 468.064036 468.246571 468.246571 468.239233	.228617 E .16804 469.848785 466.078372	466.383876
TS= TS= TS=	MG= A (X) /A0=	IG 19 19 19 19 19 19 19 19 19 19 19 19 19	MG= A(X) /A0=	nd= 201 201 201 201 201 201 201 201 201 201	<u>MG=</u> A(X)/A0=	10 	MG= A(X)/A0= TG= TS=	=ST
152.014994 149.028242 143.756302 139.217263	7.344689 169.4715	182.572581 181.752410 179.735052 177.522391 173.105261 169.011614 165.296641 161.921286 155.996841 150.919927	7.013268 154.5222	200.850106 199.820947 197.304943 194.613782 189.383626 189.381404 180.381404 176.540362 169.840584 164.129202	6.688209 140.5302	221.700256 220.460186 217.443227 214.268247 202.721692 197.853740 197.853740 197.853740 197.853740 197.853740 197.853740	6.389047 128.2396 244.057328 242.623047	239.132672
	YDUCT= A(X) =	=11 =11 0 11 0 11 0 11 0 11 0 11 0 11 0	YDUCTI \ A(X)=	= 20 = 20 = 20 = 20 = 20 = 20 = 20 = 20	YDUCT= A(X)=		YDUCT= A(X) = VG= VL=	NL= N
.000001712 .000000173 .000000018 .000000018	16.004321 .007000000	.000567729 .000007321 .000039484 .000036823 .000024082 .000001710 .000000173 .000000173	16.501321 .007000000	.000568253 .000007182 .000039261 .000036706 .0000024043 .000004193 .000001709 .0000001709 .000000173	17.005321 .00700000	.000568731 .000007057 .000039057 .000036598 .000024006 .000001708 .000000173	17.502321 .007000000 .000569152 .000066949	.000038877
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3/13/98 5:15 PM Hard Disk 2150:MPM:AEDC ADVERSE WEATHER 3/98:DEVELEMENTS OF AEDC MFF CODE:AEDCIDMP ...: AEDCIDMP.CUT

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0.703388E-11 0.273559E-10 0.680586E-10 0.135978E-09 0.237991E-09 0.571748E-09 0.571748E-09	474.837634 8	.000112006 0.2896858-12 0.2374168-11 0.7017088-11 0.2732118-10 0.6800028-10 0.1358918-09 0.2378688-09 0.2378688-09 0.2378688-09 0.21124028-08	474.836798 5	.000111679 0.286254E-12 0.236562E-11 0.700245E-11 0.272904E-10 0.679481E-10 0.679481E-10 0.135812E-09 0.237757E-09 0.571346E-09 0.571346E-09	474.836211 0	.000111389 0.283237E-12 0.235807E-11 0.698946E-11 0.679011E-10 0.135740E-09 0.135740E-09 0.237655E-09 0.237655E-09 0.1123455E-09	474.835847 6
DWASS= DWASS= DWASS= DWASS= DWASS= DWASS= DWASS= DWASS=	TO= 37.20655	FLtot= DMASS= DMASS= DMASS= DMASS= DMASS= DMASS= DMASS= DMASS= DMASS=	TO= 38.61336	FLtot= DMASS= DMASS= DMASS= DMASS= DMASS= DMASS= DMASS= DMASS= DMASS=	TO= 40.08288	Fl.t.ot= DMASS= DMASS= DMASS= DMASS= DMASS= DMASS= DMASS= DMASS= DMASS= DMASS=	TO= 41.66236
18.26800 28.72827 38.92723 49.02856 59.08525 79.13355 99.14217	2075.6193 WM. %=	1986.38 6.30862 6.30862 12.71929 18.25344 28.71609 38.91609 49.01801 59.07508 79.12386 99.13271	2075.6178 IUM. %≕	1972.14 6.28362 12.70402 18.24074 28.70534 38.90615 49.00852 59.06588 79.11500 99.12403	2075.6163 IUM. %=	1957.91 6.26146 12.69049 18.22945 28.69571 38.89718 48.99991 59.05748 79.10686 99.11599	2075.6147 NM. %=
▙▙▙▙₽₽	Ó= REL H	_╣ ӒӒӒӒӒӒӒӒӒ	C= REL H	_╣ ӓӓӓӓӓӓӓ	REL F	[╨] ┱┱┱┱┱┱┱┱┱	REL E
466.656802 467.105287 467.451686 467.725345 467.946751 468.294334 468.294334	.250773 P	468.845780 465.302396 465.676464 465.676464 466.012931 466.989802 467.322746 467.322746 467.322746 468.003980 468.366641	.270671 P .14370	467.870214 464.523751 464.939485 465.329650 465.329650 465.882299 466.882280 467.199477 467.684611 468.092100	.289337 F	466.891512 463.743707 464.194523 464.628757 465.373377 465.373377 465.313377 465.313377 465.313377 465.782811 467.798430	.307476 F
75 75 75 75 75 75 75 75 75 75 75 75 75 7	MG= A(X)/A0=	16 17 18 18 18 19 19 19 19 19 19 19 19 19 19 19 19 19	MG= A(X)/A0=	IG= 17S= 17S= 17S= 17S= 17S= 17S= 17S= 17S	MG= A(X)/A0=	10 10 10 10 10 10 10 10 10 10 10 10 10 1	MG= A(X)/A0=
235.499305 228.600203 222.442613 216.966321 212.050872 203.519530 196.280650	6.119458 117.6456	267.428136 265.883795 262.048391 258.053838 259.053838 259.053838 259.053838 253.053838 253.053838 246750 243.65750 233.259416 232.788904 214.762531	5.908304 109.6669	288.350870 286.866538 283.120252 279.122789 271.340879 271.340879 264.258092 257.880231 257.880231 255.103812 253.320455 233.320455	5.732180 103.2261	307.917831 306.440636 302.680767 298.667275 299.819140 290.819140 283.605305 283.605305 271.068647 271.068647 251.374935 251.374935	5.578275 97.7574
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3/13/98 5:15 PM Hard Disk 2150:MPM:AEDC ADVERSE WEATHER 3/98:DEVELEMENTS OF AEDC MFF CODE:AEDCIDMP ...: AEDCIDMP.OUT

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.000111127 0.280528E-12 0.235125E-11 0.697767E-11 0.272378E-10 0.678578E-10 0.135674E-09 0.13561E-09 0.237561E-09 0.571004E-09 0.571004E-09	474.835691 8	.000110895 0.278153E-12 0.234523E-11 0.696721E-11	0.578189E-10 0.135614E-09 0.237475E-09 0.570853E-09 0.112296E-08	474.835746 2	.000110684 0.276006E-12 0.233974E-11 0.695764E-11 0.271946E-10 0.677827E-10 0.135558E-09 0.237395E-09 0.570711E-09 0.570711E-09 0.112273E-08	474.836000 [5 .000110497 0.274120E-12 0.233488E-11 0.694910E-11 0.677498E-10 0.677498E-10 0.135507E-09 0.135507E-09 0.237321E-09 0.570579E-09 0.570579E-09
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465.882261 462.938409 463.427030 463.902375 464.734291 465.394954 465.338709 466.338709 466.976975 467.484343	.324540 PC .12213 H	464.881197 462.133872 462.658555 463.172624	464.086415 464.822416 465.840902 465.882321 466.599871 467.160498	.341579 P	463.832264 461.292844 461.860413 462.415129 463.410879 464.874212 464.874212 465.401590 466.201613 466.818481	.358035 P .11220 462.773049 460.437798 461.050166 461.647165 461.647165 463.611494 463.611494 464.909537 465.793127 466.467697
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3/13/98 5:15 PM Hard Disk 2150:MFW: AEDC ADVERSE WEATHER 3/98:DEVELEMENTS OF AEDC MFF CODE: AEDCITMP ...: AEDCITMP. CUT

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2075.6078 IUM. &=	1883.03 6.18069 12.64072 18.18745 28.65905 38.86240 48.96600 59.02600 59.08298 99.08298	2075.6058 IUM. %=	1866.98 6.16941 12.63363 18.18134 28.65355 38.85705 48.96068 59.01875 79.06845 99.07754	2075.6036 IUM. \$ =	1850.36 6.15941 12.62727 12.62727 18.17582 28.64768 38.85210 48.95572 59.01372 79.06337 99.07237	2075.6012 HUM. %= 1833.05 6.1505(12.62155 18.17075 28.6438 38.84746 38.84746
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.374708 F	461.654303 459.534921 450.199223 460.842717 460.842717 462.970247 463.753349 465.362501 466.097939	.391063 I	460.513047 458.608176 459.327477 460.020540 461.272392 462.315114 463.165814 463.861357 465.7189100 465.718910	.407486 I	459.323867 457.640377 458.420101 459.166799 460.512217 461.636044 462.556339 463.32514 465.32514	.424115 .09761 .09761 458.076477 456.618367 457.462472 457.462472 459.712825 450.922019 461.915167
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5.121473 82.4024	396.559489 394.906419 390.676353 386.248386 377.711926 369.893140 362.746637 362.746637 362.746637 364.330495 333.886206	5.031572 79.5348	413.362434 411.656402 407.293095 402.738064 393.982919 385.984684 378.684031 371.961090 359.880976 349.207206	4.947965 76.9136	430.173288 428.401364 423.875246 419.175853 410.181129 401.986896 394.517887 387.644931 375.298951 375.298951 364.386013	4.869482 74.4930 447.128125 445.300955 440.632784 435.795347 426.556897 418.156635 410.508220
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21.506321 .007000000	.000571209 .00006446 .000038003 .000038003 .000023800 .0000023800 .000001700 .000000172 .000000172	22.003321 .007000000	.000571360 .0000037939 .000035985 .0000035985 .0000035985 .000001699 .000000172 .000000172	22.500321 .007000000	.000571496 .00006379 .000037882 .000035952 .000023774 .000004169 .000000172 .000000172	23.004321 .007000000 .0000571619 .000006352 .000037830 .000035923 .000023763 .0000023763
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3/13/98 5:15 FM Hard Disk 2150:MEW: AEDC ADVERSE WEATHER 3/98: DEVELEMENTS OF AEDC MFF CODE: AEDCITMP ...: AEDCITMP. OUT

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0.570123E-09 0.112180E-08	474.840588 2	.000109808 0.257441E-12 0.231708E-11 0.691712E-11 0.571035E-10 0.571035E-10 0.135299E-09 0.135299E-09 0.237017E-09 0.570025E-09 0.570025E-09	474.842206 5	.000109712 0.2665668-12 0.2314638-11 0.6912548-11 0.2709258-10 0.6759868-09 0.1352658-09 0.1352658-09 0.1121498-09 0.1121498-09 0.1121498-09 0.2312538-09 0.2312538-11 0.6908288-11 0.2708288-11 0.2568328-12 0.2312538-10 0.1352358-09 0.2568448-09 0.1352358-09 0.5568448-09 0.5568448-09 0.5568448-09	474.846040 36 .000109558	0.265208E-12 0.231070E-11 0.690502E-11 0.270739E-10
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79.05852 99.06740	2075.5987 UM. %=	1815.51 6.14283 12.61653 18.166534 28.63964 38.84325 48.94676 59.00461 79.05399 99.06274	2075.5960 UM. %=	1797.43 6.13613 12.61207 18.16234 28.63573 28.83933 48.94272 59.00047 79.04965 99.05824 99.05824 1779.52 1779.52 6.13049 12.60826 13.60826 12.60826 1	2075.5901 NUM. %= 1761.32	6.12568 12.60493 18.15575 28.62922
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463.978953 464.910324	.440513 PC .09475 1	456.804494 455.571195 456.481686 457.346909 458.896231 460.192779 461.260083 461.260083 464.486024	.457001 P	455.484255 454.475365 456.381787 456.381787 456.381787 456.381787 456.381787 461.517733 461.517733 461.517733 461.517733 461.517733 462.965591 462.965591 462.965591 464.040590 462.965591 454.166208 453.371847 454.412926 455.403942 457.172821 457.172821 457.172821 457.172821 457.172821 457.172821 457.172821 457.172821 457.172821 457.172821 457.172821 457.172821 457.172821 457.172821 457.172821 457.172821 457.27797	.488871 E .08761 452.817777	452.235723 453.337700 454.391031 456.272441
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390.845058 379.680091	4.797625 72.3107	463.778974 461.897403 457.088044 452.117406 442.643536 434.040492 426.212874 419.017513 406.100810 394.681395	4.730433 70.2994	480.451022 473.617495 473.617495 458.534033 458.534033 450.048515 442.045515 442.045515 442.045515 421.482585 421.482585 409.804932 421.482585 409.804932 421.482585 409.804932 426.507914 457.486774 457.486774 455.645348 457.486774 455.645348	4.613415 66.8644 512.469992	510.519128 505.491459 500.273626 490.289334
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Hard Disk 2150: MEW: AEDC ADVERSE WEATHER 3/98: DEVELEMENTS OF AEDC MFF CODE: AEDCIDMP ...: AEDCIDMP. CUT 0.270529E-10 0.675207E-10 0.135134E-09 0.236766E-09 0.569541E-09 0.112084E-08 .000109367 0.263722E-12 0.263963E-12 0.270591E-10 0.675335E-10 0.135156E-09 0.230681E-11 0.675475E-10 0.135180E-09 0.236838E-09 0.569683E-09 0.264289E-12 0.230788E-11 0.236801E-09 0.569610E-09 0.112095E-08 0.689716E-11 0.675629E-10 0.135206E-09 0.236878E-09 0.569760E-09 0.112121E-08 0.264699E-12 0.230917E-11 0.690199E-11 0.270661E-10 0.112108E-08 0.689939E-11 .000109402 .000109445 .000109497 474.853069 474.855602 474.850590 474.848238 82.184253 86.602750 73.711062 77.854202 DMASS= FLtot= FLtot= å FLtot= FLtot= å į å 18.15081 28.62401 38.82686 48.92955 1691.98 6.11422 18.14885 28.62182 38.82441 2075.5766 2075.5835 6.11860 58.98666 79.03480 2075.5801 6.11609 12.59786 48.92687 79.03159 99.03911 48.93558 79.04176 2075.5869 18.15309 28.62646 12.59980 99.04257 58.98377 58.99304 99.04997 12.60214 38.82954 48.93244 58.98973 79.03817 99.04617 38.83249 6.12177 1708.44 1725.69 1743.37 REL HUM. %= HUM. %= REL HUM. &= HUM. %= 4 ᆸᆸ ۵ 出出 씸 占占 4 占 씸 Ц 4 卢卢 ۵ 占 占 占 ᆸᆸ ۵ Ь طٰ ۵ 4 出 占 REL 盀 REL 비 塭 å å å å 2 450.077484 451.301480 460.229748 461.887910 447.586719 447.728893 457.848018 459.151232 455.369435 456.914818 458.199398 460.182976 452.345548 457.677135 460.765920 453.509958 457.037190 458.420685 459.567276 449.961114 451.174105 454.447817 456.212917 458.892349 448.824951 461.648088 452.260668 455.366898 462.637649 462.150181 463.115041 451.100223 453.374331 461.332551 450.148939 448.842442 451.477932 .08104 .08396 .08241 .08569 .546950 .533468 .519150 .504256 A(X)/A0= A (X) /A0= A(X)/A0= A (X) /A0= TS= TS= **TS**= TS= ±S= TS= "IS= TS= TS" TS= TS= TS= TS= TS= TS= TS= TS= **TS**= TS= ±ST TS= TS= TS= ±S= TS= TS= **TS**= TS= Ë TSL ₽ La ¦ ¦ ģ ۳ ۳ = W = W HOH 570.074388 568.239364 **451.540799 439.3859**07 556.791448 554.914564 534.619594 516.752194 508.846682 549.971273 525.301257 494.564148 487.824472 540.688006 535.636028 520.206058 510.916602 494.578244 480.442148 481.854610 481.193185 472.899942 465.265806 527.824901 525.864809 520.794466 505.438928 496.231318 480.077145 453.766310 542.625527 530.357230 502.419527 544.747911 515.528131 467.888907 466.131601 4.436968 4.516049 4.474166 4.562385 62.8889 61.8475 64.0718 65.3934 YDUCI-YDUCT-YDUCT-YDUCTч**г**ч A (X)= g ~I/ ۳<u>-</u> A (X) = <u>ح</u>12 =12 5 5 5 5 41-۲<u>-</u> Ľ 35 <u>-17</u> 45 g чц. 5 5 A (X)= ¦۳ ЧL= ЧL= ۲<u>ا</u>= ЧL= ۲<u>-</u> 2L= 25 <u>-1</u> 25 12 ۳**5** 21= 25 A (X) = 55 g .000572166 .000023713 .007000000 .000037618 .000001696 .000000172 .00000017 .007000000 .000006253 .000037635 .000035804 000001696 00000017 .007000000 .000006246 .000035793 .000023708 .000004151 .000000172 000572131 .000000000 .007000000 .000006263 .000037656 .000035818 .000023719 .000004153 .000001697 000000172 000000017 .000572088 .000000000 .000000000 .000004154 .000001697 .000000172 .00000017 .00000000 000572037 27.001321 26.000321 26.504321 25.503321 5:15 PM 3/13/98 25 티 =1 님 =12 ЕL= =XQ =X= -12 **F**L= =15 -25 = E = H = =XC 5 Ľ Ē 凒 빌 븝 EL= 2 -E X= DX= EL= EL= **"** 믭 = [] =]님 Ē -E Ë E EL= 12 ŝ × ä

....AEDCIDMP.OUT Hard Disk 2150: MEW: AEDC ADVERSE WEATHER 3/98: DEVELEMENTS OF AEDC MFF CODE: AEDCLIMP 0.270353E-10 0.674820E-10 0.135064E-09 0.236653E-09 0.569306E-09 0.270388E-10 0.674901E-10 0.263392E-12 0.112043E-08 0.569415E-09 0.135079E-09 0.236678E-09 0.569359E-09 0.112052E-08 0.230448E-11 0.689162E-11 0.230597E-11 0.689532E-11 0.270476E-10 0.675094E-10 0.135114E-09 0.569476E-09 0.112073E-08 0.263549E-12 0.230531E-11 0.689379E-11 0.270428E-10 0.674992E-10 0.135096E-09 0.236705E-09 0.112062E-08 0.263442E-12 0.230482E-11 0.689258E-11 0.236735E-09 .000109303 .000109339 .000109318 474.863738 474.866560 474.858263 474.860958 106.048075 96.011516 100.958962 91.262593 DMASS= FLtot= FLtot= FLtot= å å å ģ 2075.5658 2075.5621 6.11288 PO= 2075.5695 REL HUM. %= 6.11206 48.92022 79.02320 18.14400 28.61562 38.81698 48.91838 99.02714 79.02862 28.61826 58.97868 79.02580 12.59423 18.14484 28.61684 38.81854 58.97645 99.02990 6.11167 58.97437 79.02073 28.61993 48.92445 99.03588 2075.5731 18.14590 38.82027 48.92222 99.03278 18.14724 12.59361 12.59632 58.98114 12.59511 38.82224 1675.73 1660.20 1645.01 HUM. %= REL HUM. 8= REL HUM. %= ᆲᆸ 占 占 씸 REL ۳ ۳ <u></u> å å **445.577887 446.907613** 453.869709 455.419690 457.830656 454.635066 459.614614 450.279555 452.585983 454.535542 456.159422 457.511560 459.602939 446.636168 447.943788 451.652344 453.690119 455.391867 456.811481 459.011186 448.251915 450.736791 452.858366 456.120087 458.425435 460.129895 444.529184 445.876084 449.816234 452.019154 449.010571 449.252991 447.251251 461.147772 460.636604 443.955236 446.338130 445.137298 .572538 .07868 .07675 .07767 .07979 .596120 .584580 .560105 A(X)/A0= A(X)/A0= A(X)/A0= A (X) /A0= =ST TS= TS= TS= TS= TS= TS= TS= TS= "ST TS= TS= TS= TS= TS≔ TS= TS= ТS= TS= TS= 1S1 TS= TS= TS= **TS**⁼ TS= TS= TS= TS= TS= TS=TS= **T**S= TS" TS= ¦¦ ¦۳ ¦P HOH Ë ₩C¦ ₩Q# 563.390503 558.243454 548.195732 538.897920 583.661353 573.806430 605.129592 600.538437 576.752456 561.366018 535.624200 521.196562 593.366224 588.673128 564.613154 556.096862 548.164286 533.717398 595.609033 585.870973 568.282470 560.375759 545.938513 532.951974 530.337239 522.401193 508.025819 581.177574 576.399914 571.319862 543.567117 508.289822 595.129444 520.758031 **4**.343645 59.2732 606.849747 4.317875 58.5720 582.981110552.113851 195.201273 4.371915 60.0473 4.402649 60.8945 YDUCT-YDUCT-YDUCI-YDUCT-A(X) =32 Ч**Г**= = 5 5 ="I Ч. ЧL= A (X)= A (X) = <u>-</u>П <u>"</u> Ľ ЧL= 2 || <u>-</u>7 ЧL= **"** 5 25 ЧС= 2ª <u>-</u>2 ۳<u>1</u> ZL= <u>=</u>75 <u>"</u> <u>-П</u> <u>"</u> ۲<u>-</u> A (X) = ÿ **"** <u>-</u>П =75 **5** 44 25 ₹¶ **"** g ۲L= 29.003321 .007000000 .000006233 .000037585 .000035769 .000001695 .000000172 .007000000 .000572230 .000006232 .000037579 .000035764 .000023692 000004149 .000001695 000000172 000006236 .000035775 .000004150 000572215 .000023695 .000004149 .000000017 .000000017 .007000000 000037593 .000023699 000001695 000000172 000000017 0000000000 000000000 000000000000 .000035783 .000001696.000000017 .000572194 .000037604 .000023703 .000004151 000000172 .000000000 28.002321 28.506321 27.505321 3/13/98 5:15 PM = E = E <u>"</u>" X= DX= =<u></u>]_ =]E =1 -14 믭 밀 = 1 X= DX= E E -12 빌 믭 = [] **EL**= Ę =XC 5 ---Ē 14 =15 2 밀 -12 EL. =11 FL= =<u></u>]E **"** 님 EL. °, Ē 뿗

3/13/98 5:15 PM Hard Disk 2150:MFW: AEDC ADVERSE WEATHER 3/98: DEVELEMENTS OF AEDC MFF CODE: AEDCITMP ...: AEDCITMP. OUT

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	.000109293 0.263394E-12 0.230429E-11 0.689093E-11 0.270325E-10 0.674749E-10 0.135050E-09 0.236630E-09 0.569257E-09 0.112034E-09 0.112034E-09	474.869539 1	.000109288 0.263444E-12 0.230422E-11 0.689046E-11 0.270301E-10 0.574687E-10 0.135038E-09 0.135038E-09 0.2569211E-09 0.569211E-09 0.112026E-08	474.872680 6	.000109287 0.263540E-12 0.230426E-11 0.689019E-11 0.270281E-10 0.674631E-10 0.135026E-09 0.236590E-09 0.236590E-09 0.569167E-09 0.569167E-09	474.875871 1	.000109291 0.263676E-12 0.230442E-11 0.689012E-11 0.270265E-10 0.270265E-10 0.135016E-09 0.135016E-09 0.236572E-09 0.569126E-09 0.112011E-08
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	1630.34 6.11169 6.11169 12.59325 18.14339 28.61460 38.81563 48.91673 58.97248 79.01846 99.02456	2075.5582 JM. %=	1615.51 6.51313 12.59313 18.14298 18.14298 38.61375 38.81443 48.91524 58.97074 58.97074 59.02212	2075.5541 UM. %=	1600.56 6.11282 12.59321 18.14274 18.14274 28.61305 38.81337 38.91388 58.96912 79.01429 99.01978	2075.5500 UM. %=	1586.01 6.11387 6.11387 12.59349 18.14268 18.14268 28.61251 38.81246 48.91267 58.96765 79.01241 99.01758
	╉ ┛┓╗╗ ┙	C= REL H	_ॣ ॖ ॖ ॖ ॖ ॖ ॖ ॖ ॖ ॖ ॖ ॖ ॖ ॖ ॖ ॖ ॖ ॖ ॖ ॖ	O= REL H	[╨] ╺╗╗╗╗ ┺╗╗ ┺	o= Reil H	[╨] ╙ ╙ ╙ ╙ ╙ ╙ ╙ ╙ ╙ ╙ ╙ ╙ ╙ ╙ ╙ ╙ ╙ ╙ ╙
-	442.805327 443.511788 444.875702 446.278216 448.916689 451.196478 453.117804 454.730505 457.244095 459.105738	.607690 P	441.635759 442.486930 443.879251 445.310386 445.310386 445.375452 452.366379 452.366379 456.656373 458.595347	.619272 P .07505	440.448742 441.436963 442.865491 444.328375 447.111310 449.541649 451.602606 453.339719 456.057935 458.075298	.630462 F .07430	439.286677 440.400745 441.864720 443.360281 448.718365 450.847762 450.847762 455.465426 457.560062 457.560062
	퉒 ᇺ ᅂᅂᇊ ᅂ ᅂ ᅂ ᅂ ᅂ ᅂ ᅂ ᅂ ᅂ ᅂ ᅂ ᅂ ᅂ ᅂ ᅂ () ()	MG= A (X) /A0=	TG= TS= TS= TS= TS= TS= TS= TS= TS= TS= TS	MG= A(X)/A0=	10 12 12 12 12 12 12 12 12 12 12 12 12 12	MG= A(X)/A0=	=51 =51 = 21 = 21 = 21 = 21 = 21 = 21 = 21 = 2
	618.037327 616.326093 611.784429 606.917009 597.285049 588.242881 579.825593 571.953403 557.545646 544.552197	4.293277 57.9065	629.211406 627.465450 622.867494 617.993595 608.400210 599.404810 591.026117 581.82120 588.802566 555.807795	4 .269846 57.2762	640.352138 638.604045 633.983458 629.084191 629.479423 610.497244 602.136128 594.307100 579.942498 566.942406	4 .248294 56.6995	651.073370 649.334689 644.725843 639.819815 630.200846 621.218058 612.862384 605.039315 590.679325 577.670307
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	.000572240 .00006232 .000037576 .000035760 .000023690 .000004149 .0000001695 .000000172 .000000172	29.500321 .007000000	.000572245 .00006233 .000037575 .000035758 .000023688 .000023688 .000001448 .000000172 .000000172	30.004321 .00700000	.000572246 .00000572246 .000037576 .0000035756 .000003148 .000001485 .000000172 .000000172	30.501321 .007000000	.000572242 .000037578 .000035756 .000035756 .000003685 .000004148 .000001694 .000000172 .000000172
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3/13/98 5:15 FM Hard Disk 2150:MEW: AEDC ADVERSE WEATHER 3/98:DEVELEMENTS OF AEDC MFF CODE: AEDCIDMP ...: AEDCIDMP. OUT

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474.879208 1	.000109299 0.263853E-12 0.230467E-11 0.689022E-11 0.674543E-10 0.674543E-10 0.135007E-09 0.135007E-09 0.2569087E-09 0.569087E-09 0.112004E-08	474.882650 4	.000109310 0.264064E-12 0.230502E-11 0.689049E-11 0.270246E-11 0.270246E-10 0.134999E-09 0.134999E-09 0.236541E-09 0.569052E-09 0.111997E-08	474.886139 6	.000109324 0.264311E-12 0.230546E-11 0.689092E-11 0.270242E-10 0.134992E-09 0.134992E-09 0.236528E-09 0.569018E-09 0.1111990E-08	474.889406 1 .000109342 0.264584E-12 0.230597E-11 0.689149E-11 0.674458E-10 0.674458E-10 0.134986E-09
TO= 130.020404	FL.tot= DMASS= DMASS= DMASS= DMASS= DMASS= DMASS= DMASS= DMASS= DMASS= DMASS=	TO= 136.924224	Fl.tot= DMASS= DMASS= DMASS= DMASS= DMASS= DMASS= DMASS= DMASS= DMASS= DMASS=	TO= 144.08135	Fltot= DMASS= DMASS= DMASS= DMASS= DMASS= DMASS= DMASS= DMASS=	TO= 150.89689 FLtot= DMASS= DMASS= DMASS= DMASS= DMASS= DMASS=
075.5456 M. %=	1571.42 6.11524 12.59396 18.14277 28.61210 38.81167 48.91157 58.96629 58.96629 79.01063	2075.5411 JM. %=	1556.94 6.11686 6.11686 12.59459 18.14301 28.61183 38.81102 48.91061 58.96506 79.00898 99.01349	2075.5365 UM. %=	1542.86 6.11877 6.11877 12.59540 18.14338 18.14338 38.81048 48.90974 58.96393 79.00741 99.01159	UM. %= UM. %= 1530.24 6.12088 12.59633 18.14388 28.61164 38.81005 48.90900
EL HC	_╫ ╘╘╘╘╘	L L L L L L L L L L	_╨ ┨┇┇┇┇	REL H	_╢ ӒӒӒӒӒӒӒӒӒ	
.641624 PC .07358 1	438.113004 439.351409 440.850182 442.379580 447.882863 447.882863 451.940953 451.940953 454.862390 457.035317	.652643 P	4 36.940215 438.306391 439.844795 441.409751 444.407399 447.056858 449.322369 451.242943 451.242943 456.515201	.663305 P	435.792367 437.256172 438.827004 440.427279 446.218508 448.551732 450.533334 455.985534 455.985534	.672823 P .07178 434.757133 436.275209 437.849954 437.849954 437.476698 445.399106 445.399106
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4.227821 56.1543	661.724706 659.985981 655.385438 650.481290 640.854824 631.867891 623.512197 615.690225 601.328504 588.308004	4.208571 55.6441	672.198726 670.444588 665.796525 660.865891 651.208704 642.204088 633.837988 633.837988 633.837988 633.637040 598.597070	4.190821 55.1757	682.293765 680.617356 676.062864 671.155089 661.497365 652.482517 644.107180 636.270395 621.880996 608.824469	4.175681 54.7778 691.271653 689.751387 685.473163 680.713102 671.182255 662.218051 653.868448
YDUCT= A (X) =	=17 =17 =17 =17 =17 =17 =17 =17 =17 =17	YDUCT= A(X)=		YDUCT= A(X) =	NL= NL= NL= NL= NL= NL= NL= NL= NL= NL=	YDDCT= A(X) = VG= VL= VL= VL= VL= VL=
31.005321 .007000000	.000572235 .00006243 .000037583 .000035757 .000023684 .000001694 .000000172 .000000172 .000000172	31.502321 .007000000	.000572224 .00006248 .000035758 .000035758 .000023683 .000001694 .000001694 .0000001694 .000000172	32.006321 .00700000	.000572209 .00006254 .000037595 .000035760 .000023683 .000004147 .000001694 .000000172 .000000172	32.503321 .00700000 .000572192 .000006260 .000037604 .000035763 .000035763 .0000035882 .000004147
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3/13/98 5:15 PM Hard Disk 2150:MPM:AEDC ADVERSE WEATHER 3/98:DEVELAMENTS OF AEDC MFF CODE:AEDCITMP ...: AEDCITMP. CUT

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0.236515E-09 0.568987E-09 0.111984E-08	474.892388 3	.000109362 0.264881E-12 0.230655E-11 0.689217E-11 0.570243E-10 0.574442E-10 0.134981E-09 0.236505E-09 0.568958E-09 0.568958E-09 0.111979E-08	474.894957 5	.000109384 0.265202E-12 0.230719E-11 0.689298E-11 0.270247E-10 0.270247E-10 0.134976E-09 0.134976E-09 0.236495E-09 0.568931E-09 0.568931E-09	474.896927 7	.000109408 0.2655328-12 0.2307878-11 0.6893858-11 0.2702548-11 0.2702548-10 0.2744248-10 0.1349738-09 0.2364978-09 0.5689078-09 0.1119688-08	474.898149 4	.000109433 0.265874E-12 0.230857E-11 0.689479E-11
DMASS= DMASS= DMASS=	TO= 157.18755	Fltot= DMASS= DMASS= DMASS= DMASS= DMASS= DMASS= DMASS= DMASS=	TO= 162.59033	Fl.tot= DMASS= DMASS= DMASS= DMASS= DMASS= DMASS= DMASS= DMASS= DMASS=	TO= 166.62504	Fltot= DMASS= DMASS= DMASS= DMASS= DMASS= DMASS= DMASS= DMASS= DMASS=	TO= 168.84466	FLtot= DMASS= DMASS= DMASS= DMASS=
58.96292 79.00597 99.00982	2075.5279 UM. %=	1519.20 6.12317 12.59738 18.14448 28.61171 38.90874 48.90837 58.90837 58.90837 79.00815	2075.5242 UM. %=	1510.15 6.12564 12.59855 12.59855 12.14519 28.61187 38.9951 48.90784 58.90784 58.90784 59.00540 99.00657	2075.5212 UM. %=	1503.63 6.12818 12.59977 18.14595 28.61211 38.80939 48.90741 58.96053 79.00229 99.00513	2075.5191 UM. 8 =	1500.12 6.13081 12.60105 18.14678
╏┛╏	0= REL H	_┨ ╘╘╘ [┶] ╘╘╘	O= REL H	_ॣ ॖॖॖॖ॑॑॑॑॑॑॑॑॑॑॑॑॑॑॑॑॑॑॑॑॑॑	O= REL H	[₽] ₽₽₽₽₽₽₽₽	O= REL H	<u> </u>
449.836958 453.059011 455.463879	.681117 P .07134	433.846968 435.384559 436.926322 438.564549 441.747903 444.597494 447.055113 449.150976 452.467772 454.946926	.687901 P .07100	433.097055 434.609558 436.072761 437.701065 440.314812 443.814949 446.326649 446.326649 446.326649 441.881348 454.432434	.692781 P .07076	432.554726 434.002968 435.340610 435.340610 436.932803 445.642333 445.642333 447.834860 451.322935 453.940185	.695406 P	432.262115 433.589626 434.741555 436.264587
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646.046505 631.670135 618.612563	4.163013 54.4459	699.069554 697.748833 693.870302 689.382301 689.382301 689.382301 680.145158 671.328224 663.066592 653.0066592 653.303160 640.997971 627.977419	4.153004 54.1844	705.429941 704.379989 701.067875 696.991854 688.245747 689.700571 671.612120 649.822369 636.897424	4.145995 54.0017	709.994735 709.254261 709.254261 703.097914 695.037561 686.894858 679.073259 671.623259 671.735914 657.735914 657.735914	4.142291 53.9053	712.446682 712.092135 710.400697 707.595902
ענ≓ ענ=	YDUCT= A(X) =	=10 =10 =10 =10 =10 =10 =10 =10 =10 =10	YDUCT= A (X) =	=27 =17 =17 =17 =17 2 2 2 2 2 2 2 2 2 2 2 2 2 2 2 2 2 2 2	YDUCT= A(X)=	=71 =71 201 201 201 201 201 201 201 201 201 20	YDUCT= A(X) =	NL= VL= VL=
.000000172 .000000017 .000000002	33.000321 .00700000	000572172 00006267 000037613 000035767 000023683 000004147 00000172 000000172	33.504321 .007000000	.000572149 .00006275 .000037624 .000035771 .0000023683 .0000004147 .0000001694 .000000172	34.001321 .007000000	000572125 000006283 000037635 000035775 000023684 0000023684 0000004147 0000001694 000000172	34.505321	.000572100 .00006291 .000037646 .000035780
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Hard, Disk 2150:MPM: AEDC ADVERSE WEATHER 3/98: DEVELEMENTS OF AEDC MFF CODE: AEDCIDMP ...: AEDCIDMP. CUT 0.236479E-09 0.270263E-10 0.674421E-10 0.134970E-09 0.568885E-09 0.111964E-08 0.270273E-10 0.674423E-10 0.134968E-09 0.236474E-09 0.568866E-09 0.230995E-11 0.689670E-11 0.270284E-10 0.674427E-10 0.134966E-09 0.236469E-09 0.568849E-09 0.266890E-12 0.689759E-11 0.270296E-10 0.674433E-10 0.134965E-09 0.236464E-09 0.568834E-09 0.111953E-08 0.266211E-12 0.111960E-08 0.266550E-12 0.111956E-08 0.231063E-11 0.230926E-11 0.689574E-11 .000109508 .000109458 .000109364 .000109483 474.898919 474.900330 474.898767 474.898801 169.300693 168.377039 167.432146 169.627466 DMASS= FLtot= FLtot= FLtot= FLtot= å å ខ្ពុ å 28.61318 38.80944 28.61358 38.80956 48.90650 28.61243 38.80934 2075.5179 2075.5172 12.60356 2075.5145 48.90706 99.00378 6.13340 28.61279 38.80936 18.14845 48.90662 12.60480 58.95993 79.00127 79.00038 99.00256 12.60232 6.13600 58.95902 78.99957 99.00144 2075.5171 18.14923 58.95868 99.00044 18.14761 48.90681 58.95944 6.13861 78.99887 1500.80 1502.33 1498.88 1499.37 REL HUM. %= REL HUM. &= REL HUM. 8= PO= 2075.51 REL HUM. %= ۵ 4444 8888 44 44 ۵ ۵ 씸 씸 占 占 ᆸᆸ 出出 ۵ 씸 塭 ᇤ ۳ ۳ 塭 å å å 442.406292 444.995656 438.337685 441.268716 **433.828129 435.011677** 447.225460 444.414086 450.289048 435.310196 443.884619 446.160096 449.823960 433.360950 440.810139 450.785024 453.462865 434.304917 453.019006 437.916704 443.420926 445.705580 449.402813 439.452174 433.374607 435.729353 438.852854 441.805023 433.312997 434.009767 446.671311 452.598647 452.213821 432.320146 432.449312 432.200702 432.159283 .693752 .07071 .07059 .07066 .07061 .695960 .696329 .694893 A (X) /A0= A (X) /A0= A(X)/A0= A (X) /A0= TS= TS= TS= TS= TS= TS= ¶3= ₽ LS= TS= TS= TS= TS= TS= TS= TS= **TS** ≡ST TS= TS= TS= TS= TS= TS= TS= TS= TS= **TS**= IS= å Ë ľ ¦ WG= WC= E E E 700.473695 692.894514 685.454715 690.494370 694.413515 710.903867 678.284421 710.366129 704.350944 697.458916 683.678649 670.682994 658.527330 713.081749 712.952415 711.758638 700.829729 688.013968 675.620436 711.968545 712.552073 712.080574 697.264917 .305403 664.786181 652.287180 713.307912 713.221453 712.365575 712.964492 706.945811 663.886142 712.179722 708.437747 703.111101 679.561484 668.292353 4.143008 4.141512 4.140999 4.144623 53.8850 53.9660 53.8716 53.9239 691 YDUCT= YDUCT--LODOL YDUCT= A (X) = ۲L= Ę **1**2 ۲<u>-</u> **"** =12 =15 ۲<u>-</u> ЧС= 5 <u>-</u>7 <u>"</u>" 5 5 5 ="F =15 <u>-</u>12 <u>=</u>]5 **"** <u>-17</u> 15 **5** ۲<u>-</u> 5 ₹**1**1 A(X) =<u>"</u> <u>-</u>75 =75 A(X) =¦۳ <u>5</u> 5 5 A (X) = <u>-</u>22 5 35.002321 .007000000 .000572050 000037669 000035790 .007000000 000006315 000037680 .007000000 .000572169 .000023686 000000172 000572025 000035795 .000004147 .000000172 .000000017 .000006299 .000037657 .000023685 .000001694 000006307 000004147 000001694 .000004147 000001694 .000001694 .000572075 .000035785 .000000172 00000017 000000017 0000000000000 .000023687 000000172 .000023684 .000000000 .000004147 .000000000 000000017 2000000000 35.506321 36.500321 36.003321 3/13/98 5:15 PM н Ц Е. Е EL-Ë =XC =15 Ш_= 믭 E E =1E **FL**= ۳L= 팊 Ē =XC 2 믭 -1 1 ۲<u>-</u> Ē 밀 ŝ 믭 EL= 뿝 밀 닖 12 -12 ŝ EL-빌 =]L EL= Ŗ 뿝 튑 × ₿ Ŀ 붪 뿠 × 1

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Hard Disk 2150: MPM: AEDC ADVERSE WEATHER 3/98: DEVELEMENTS OF AEDC MFF CODE: AEDCIDMP ...: AEDCIDMP. CUT 0.269942E-10 0.674185E-10 0.134965E-09 0.568787E-09 0.111942E-08 0.674450E-10 0.270029E-10 0.674336E-10 0.236456E-09 0.268235E-12 0.236454E-09 0.270209E-10 0.674441E-10 0.270115E-10 0.134964E-09 0.267904E-12 0.134964E-09 .568797E-09 0.231350E-11 0.681442E-11 0.267231E-12 0.231135E-11 0.689246E-11 0.134965E-09 0.568820E-09 0.267574E-12 0.231209E-11 0.236458E-09 0.568808E-09 0.231280E-11 0.688106E-11 0.236461E-09 0.111949E-08 0.688670E-11 0.111947E-08 0.111944E - 08.000108662 .000109204 .000109047 474.907283 474.904059 474.905894 474.902081 166.568518 166.781018 166.643317 DMASS= DMASS= DMASS= DMASS= DMASS= DMASS= DMASS= DMASS= DMASS= **DMASS=** DMASS= DMASS= DMASS= DMASS= DMASS= DMASS= DMASS= DMASS= =SSAMC DMASS= DMASS= DMASS= DMASS= DMASS= DMASS= DMASS= DMASS= =SSAMC =SSAMC DMASS= DMASS= DMASS= DMASS= DMASS= DMASS= DMASS= FLtot= FLtot= FLtot= å ¦۳ Į å . 29.21153 18.65670 29.24004 39.51378 18.60569 29.26873 39.54103 18.56372 29.18043 18.61089 48.90639 58.95818 78.99769 98.99869 58.95800 78.99719 98.99792 58.95785 78.99673 98.99720 38.80971 48.90643 58.95841 78.99825 6.14122 12.60611 48.90638 48.90640 38.80988 2075.5087 12.60875 2075.5059 6.14891 12.61002 2075.5038 2075.5117 6.14385 12.60745 6.14637 1503.62 503.78 1503.60 HUM. %= PO= 2075.51 REL HUM. %= REL HUM. %= ᆲᆸ щ. Д =WO =WC ₽W= =MO =WQ =WO =IO 占님 卢님 щ 444 占 卢占 4 占 4 PQ= REL 塭 å å **4**32.559215 **4**33.572388 **4**32.561175 **4**33.651218 447.999920450450 **4**33.462180 **4**33.727839 441.984310 444.576628 475.261419 492.000000 444.250202 449.015989 440.063329 451.518608 433.636318 433.679486 492.000000 492.000000 492.000000 448.320475 451.204719 492.000000 492.000000 492.000000 440.413440 443.009400 445.295824 433.689324 492.000000 492.000000 442.636804 444.919172 442.301697 433.673068 448.654061 451.856450 432.574819 .07076 .07076 .07077 .692660 .692797 692801 692783 A (X) /A0= A(X)/A0= A (X) /A0= =ST =ST =ST TG= TS= **TSI=** =NST =ST TS= TS= TS= **TS**= TS= TS= =WST =WST TS= TS= TS= TS= TS= =WST TSM= TS= TS= TS= =WS1 TS= TS= TS= TS= ۳a ۳ ¦p ¦g Ë = 2 W "U =OM 711.105286 711.663720 711.720251 710.000484 710.147072 697.338793 687.543584 709.987269 710.003625 710.194488 703.033115 710.688986 711.014084 709.550546 695.778449 675.082785 709.886670 709.888497 710.123388 710.443707 709.639990 706.352437 709.716611 698.696002 709.247766 699.335050 .818235 682.741590 671.961265 705.609773 700.864203 685.365647 710.016227 706.978261 689.477453 702.032771 677.734367 704.618111 580.120811 4.146000 4.146000 4.146003 4.146187 54.0018 54.0019 54.0067 693. YDUCT= A(X)= KDOCT-YDUCI-YDUCT= g 5 A (X) = =IJI <u>-</u> A (X) = 21чL= =WIN =WTIN 25 5 **8** 25 -WIN =WLIV 44 ="IL ЧL= ЧL= 44 \$5 чL= =WTIN =WTIA 5 =<u></u>7 25 35 25 ЧL= .000006323 .000004139 000572486 000006339 .000037715 .000023576 000000172 .007000000 .000035025 .000023540 37.004321 .007000000 0000000000 .000035389 .000004143 000001694 000000017 000000000 000006347 000037727 000001694 000000172 .000035526 .000023612 .000004147 000001694.000000172 .000000017 000000002 000572871 000000017 0000000000 .000035668 .000023651 .000004147 .00000172 .00000017 000572329 .000006331 .000037704 .000001694 .000000000 38.502321 37.501321 38.005321 3/13/98 5:15 PM FLM= FLM= FLM= FLM= FL= FL= FL= FL= CV= FL= FL= FL= FL= FL= FL= FL= FL= X= DX= =]딘 =X= -12 =]F =]F =]F = X0 25 25 밀 FILM= FLM= =<u>1</u>-FLM= FLM= ×

0.269856E-10 0.134965E-09 0.236453E-09 0.568778E-09 0.673757E-10 0.268561E-12 0.231417E-11 0.674040E-10 0.673896E-10 0.134950E-09 0.236452E-09 0.269686E-10 0.568763E-09 0.269602E-10 0.673621E-10 0.134908E-09 0.236450E-09 0.568756E-09 0.679711E-11 0.111939E-08 0.268890E-12 0.231485E-11 0.269770E-10 0.568770E-09 0.678938E-11 0.134929E-09 0.236451E-09 0.269542E-12 0.229370E-11 0.678923E-11 0.679126E-11 0.269216E-12 0.231237E-11 0.111935E-08 0.111937E-06 .000108552 .000108500 .000108256 .000107754 474.908515 474.911780 474.915581 166.804816 166.812549 166.987665 166.863990 DMASS= FLtot= FLtot= FLtot= FLtot= å å å 12.61247 18.58459 39.56726 39.59319 49.78911 58.95749 78.99527 18.58992 29.32525 49.74205 39.61819 39.64268 29.29697 12.92658 18.58287 29.35287 49.76587 78.99559 98.99533 18.58273 1503.58 6.15391 12.94241 29.38027 2075.5020 6.15139 48.90644 78.99594 2075.4970 12.61123 58.95772 78.99632 98.99654 6.15639 58.95755 2075.4912 58.95762 6.15887 98.99591 1503.59 1503.57 1503.55 REL HUM. %= 90 11 HUM. %= HUM. %= REL HUM. ᇪᄲᄸ =IQ =WO =MQ ۵ =10 × EMC 씸 =IO =IQ =10 =WO =WQ ᆸᆸ HNG ₩ 出 44 占 REL Å REL ᇤ 盀 å å å 443.946339 447.699286 432.562773 433.653164 460.929042 492.000000 492.000000 433.654795492.000000 456.725282 492.000000 **492.000000 492.000000 443.655263 447.409067** 433.652079 433.667365 466.635660 492.000000 492.000000 441.690457 433.663494 443.383436 447.136039 481.096710 450.615642 450.338372 433.657924 443.125333 450.076304 453.412512 492.000000 492.000000 492.000000 446.874976 132.562016 432.564754 432.567007 .07076 .07076 .07076 .07076 .692804 .692813 .692824 A (X) /A0= A (X) /A0= A(X)/A0= A (X) /A0: TS= =IST =ISL =WST TS= TS= TS≓ =WST =WST "LS" =WST =IST =WST **TS**= TSM= TSM= TS= TS= =ISI =WST =MST TS= TS= 맓 TS" ^E ¦ ¦۳ ¦¦ "U ۳ M H U H 710.066651 709.849308 710.046205 709.930259 705.905220 702.538347 695.193171 710.046373 710.040018 710.019874 710.014585 710.108359 709.791484 .490671 703.879425 .163458 710.024158 710.020825 707.914045 692.676548 710.033289 710.049115 709.893695 708.256994 699.858889 682.225918 710.024898 704.658976 710.035790 710.028366 705.330263 701.763142 710.048966 708.541968 710.020877 700.882051 .004302 684.138603 4.146000 54.0018 4.146000 54.0018 585.836604 4.146000 54.0018 54.0018 694. 691. 707. YDUCT-ЧГ= A (X) = A (X) = ЧL= <u>=</u>12 =IJV чL= ۲<u>-</u> =IJV 25 ЧГ= 5 A (X) = A (X)= =WLIV -MLIV 21= <u>-</u>2 <u>=</u>15 ₿ B =WIN =WIN =WIN ₿ " |5 -WIN =IJIV =WTIA =WIV -WIN =15 <u>=</u>12 21-**" "** =IJN =IJV -WIN -WLIV =WIV ÿ .007000000 000000172 .007000000 000006355 .000037749 .000034906 .000006370 .000023433 000000000 000006378 000037078 .000572981 .000034936 .000004135 .000001694 000000172 00000017 .007000000 .000573033 .000006362 .000023468 .000573276 .000037547 .000034896 000573778 000034896 000023399 .000037737 .000023504 000000000 .000004131 .000001693 .000000172 000000017 000000000 .000004127 .000001692 .000000172 000000017 .000000000 000004123 000001690 40.000321 39.503321 39.006321 -12 =Xa =XQ FLI= FLM= FLM= =<u></u>___ EL= ЕL= X= DX= 믭 <u>"</u> ELLE EMER FILME Ē FLI= FLM= FLM= FI.M= -12 FILM= 빌딩 Ë =XC Ē =I'ILI =111 -12 25 8 8 12 FLM= FLM= FLM= รื Ë ່ຮື × ×

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Hard Disk 2150: MEW: AEDC ADVERSE WEATHER 3/98: DEVELPMENTS OF AEDC MEF CODE: AEDCIDMP ...: AEDCIDMP. OUT 3/13/98 5:15 PM

Hard Disk 2150: MPM: AEDC ADVERSE MEATHER 3/98: DEVELPMENTS OF AEDC MFF CODE: AEDC1DMP ...: AEDC1DMP. CUT 0.267259E-10 0.134849E-09 0.568738E-09 0.270859E-12 0.228084E-11 0.679481E-11 0.266739E-10 0.673094E-10 0.134869E-09 0.568744E-09 0.673222E-10 0.236391E-09 0.269873E-12 0.228037E-11 0.269518E-10 0.673485E-10 0.134888E-09 0.236445E-09 0.568749E-09 0.111932E-08 0.270200E-12 0.227909E-11 0.679129E-11 0.268360E-10 0.673354E-10 0.236418E-09 0.111930E-08 0.270532E-12 0.227967E-11 0.679295E-11 0.111928E-08 0.111933E-08 0.678999E-11 .000107309 .000107295 .000107510 .000107384 474.923269 474.920345 474.921968 474.918309 167.075000 167.066847 167.078832 167.042370 DMASS= MASS= DMASS= DMASS= DMASS= DMASS= DMASS= DMASS= DMASS= =SSAMC DMASS= DMASS= FLtot= FLtot= FLtot= FLtot= å å å å 18.58782 29.32233 39.73711 49.81216 59.90716 **49.83441** 59.92873 12.91816 59.95015 39.66706 39.69067 12.91597 29.34137 39.71423 18.58613 49.85653 12.91728 18.58343 29.40787 12.91487 29.38161 18.58461 6.16889 2075.4840 2075.4815 6.16640 8.99447 2075.4796 98.99478 2075.4870 6.16140 78.99498 98.99426 6.16388 98.99378 98.99332 78.99471 1503.52 1503.51 1503.53 1503.54 REL HUM. %= REL HUM. %= HUM. %= REL HUM. &= =IQ =10 =10 **DI=** =IQ ≡IQ =WC =WO =WO DI= Ē =10 =IO ₩ D =WQ ₩d =10 HA HA =WO =WQ 占占 ۵ ᆸ 4 出법 Ä 씸 REL ᇤ 盀 Ц 塭 å å -B å **433.666898 445.218521 444**.842182 **473**.598760 **492**.000000 **433.665735 449.003101** 492.000000 449.824616 485.917090 492.000000 446.381715 446.449778 478.379477 492.000000 492.000000 433.661644 450.652106 492.000000 492.000000 492.000000 492.000000 **4**32.56989**4 4**33.663970 454.556689 448.382057 492.000000 492.000000 449.346094 446.148151 449.118147 463.030942 449.579182 446.621481 432.570889 432.571698 432.568653 .07076 .07076 .07076 .07076 .692837 .692845 .692842 .692832 A(X)/A0= A(X)/A0= A(X)/A0= A (X) /A0= =IST =UST =NST TS= TS= =IST =ISL =IST =IST =WST =WST TSM= =WST TS= =IST =WST =WST **TS**= TS= =IST TSI= =WST =WST TSM= =WST **T**S≓ TS= =IST =IST TS= ц Ц ¦ ¦g ¦۳ ¦۳ = 2 = W Ē E E E 710.066723 710.021266 709.291269 698.124376 710.073592 710.070952 710.061545 710.005269 709.150976 710.056619 710.063660 709.985329 708.981849 703.901028 710.070200 710.069262 710.065951 707.210226 591.237420 710.074350 687.373442 710.051209 710.050128 709.960999 708.783055 706.406871 703.241922 696.276525 710.064893 710.059453 710.055731 706.834234 597.240015 690.059482 704.490881 710.058161 588.788427 4.146000 4.146000 4.146000 4.146000 54.0018 54.0018 54.0018 54.0018 YDUCT-YDUCT= A (X) = A (X) = 25 =IJN =WTIA =171A =171A -ILIV =171 =IJI -П =17 =WTIA =171A =WIN =WIN ₿ 8 2L= =WTIV =WTIN <u>"</u> =15 ۲<u>ا</u>= 5 A (X) = 25 =5 =WIN g A (X) = 5 -111 =IJV =WID =WIV -WIN 25 g ¦g 000036870 00000000000 .000034924 000023109 .007000000 .000034915 000574236 000006409 .007000000 000006393 .000004115 .000001688 00000017 000004107 .000000002 .007000000 000006386 .000036842 .000034906 .000023249 .000000172 000574222 .000006401 .000023154 .000004111 .000001687 .000000172 000000000000 .000034899 .000023364 .000004119.000001689 000574147 000000017 0000000000 .000036851 .000036862 .000000172 .000000017 000000000 .000574021 42.002321 41.001321 41.505321 40.504321 3/13/98 5:15 PM X= DX= FLI= DX= =I'IJ ş FLI= =11F FLI= FLI= FLM= =WTL 믭 8 =15 FLI= FLI= =W-L =WIE FLM= 븝 -T-W= Ē FLM= =1.15 FILM= FLM= EL= =XC FLI= Ē -12 =12 =111 =WTL =W11 Ē Ë =XC 8 ່ຮື × × ×

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....AEDC1DMP.OUT 2150:MPW: AEOC ADVERSE WEATHER 3/98: DEVELEMENTS OF AEDC MFF CODE: AEDCIDAP 0.236364E-09 0.568734E-09 0.672966E-10 0.266433E-10 0.134811E-09 0.568729E-09 0.134793E-09 0.266119E-10 0.672718E-10 0.134775E-09 0.568721E-09 0.134830E-09 0.111927E-08 0.271190E-12 0.228226E-11 0.679685E-11 0.236338E-09 0.271517E-12 0.266244E-10 0.672842E-10 0.568725E-09 0.236288E-09 0.111923E-08 0.272176E-12 0.228706E-11 0.111926E-08 0.228380E-11 0.679896E-11 0.236313E-09 0.111924E-08 0.271844E-12 0.228539E-11 0.680116E-11 .000107305 .000107327 .000107391 .000107357 474.926118 474.924373 474.925305 474.926850 167.065382 167.053266 167.039767 67.025267 DMASS= **DMASS**= DMASS= FLtot= FLtot= FLtot= FLtot= å å å ¦ 18.58968 29.31111 39.75999 49.89926 **49.87793** 59.97084 12.92375 18.59161 29.30415 39.78226 49.91994 18.59361 29.29957 39.80427 49.94030 1503.50 6.17887 12.92990 12.92086 59.99142 12.92676 60.03092 78.99424 98.99290 60.01134 78.99385 98.99211 2075.4780 6.17140 78.99368 78.99404 98.99249 2075.4766 6.17388 2075.4754 6.17636 98.99174 2075.4743 1503.51 1503.50 1503.50 HUM. %= REL HUM. &= HUM. %= HUM. %= <u>н</u> L =IQ =10 =WQ =WQ =1**U** =IQ =IQ =WO ۳ <mark>۳</mark> **D1=** =10 ۳ ۵ ۵ -WO 씸 ₩G =WQ =IQ -WO × ٣ =WC 씸 REL 씸 REL REL ll L Å å å å å 492.000000 492.000000 445.926598 448.901107 443.460998 469.918543 492.000000 492.000000 492.000000 445.710199 438.125888 433.667703 442.494262 439.160726 492.000000 433.669019 448.688343 440.562893 442.303454 466.955664 492.000000 492.000000 492.000000 448.485310 433.668693 441.316072 492.000000 492.000000 464.433611 448.288597 445.504421 445.305721 433.668271 432.572390 432.573492 432.572977 432.573957 .07076 .07076 .07076 .07076 692848 .692850 .692852 .692854 A (X) /A0= A(X)/A0= A(X)/A0= A (X) /A0= =WST TS= =ISL =WST =MST =MST =ISI **TSI**= =WST TS= "ISI TSI= ±S± =ST "ISI =WST =WST TS= =ISL =IST =IST TS= ±ST =ISI =WST =WST =WST TS= TS= ₽ E =2F ឌ្ឋ ¦ HQH <u>"9</u> HQ= HG= 710.085465 710.085049 710.083683 710.053721 709.596200 707.533085 705.006203 698.916283 692.302069 707.819185 710.034446 709.411269 705.470890 710.078362 710.075099 710.045036 700.910510 710.071268 699.647770 693.294384 710.080719 710.080180 709.511318 708.066488 700.306667 710.082757 710.077821 710.077191 594.196074 4.146000 54.0018 710.083226 710.081194 710.078382 708.284097 706.245863 595.029431 4.146000 54.0018 710.075031 705.879711 4.146000 4.146000 54.0018 54.0018 Hard Disk YDUCT= A(X)= YDUCTI= YDUCT-A (X) = A (X) = A (X) = ЧГ= ЧL= ~IJN ~IJV =IJN =111 =WIN <u>4</u>1-<u>-1</u>2 =IJV =WIN =WIN =IJV 5 =IJV =IJV 55 =IIN **"** =WIV 212 3 ₿ 35 -WIN -WIN -WIN <u>5</u> ₿ =IJV =WIN -WIN ۳ B =WIN =WTIN 17 .000574140 .000006440 .000036970 .007000000 .000001686 .000000171 000000017 .007000000 .000574226 .000036918 .000034946 .000023066 007000000 .000006432 .000004096 0000000000 00000000000 .000006417 .000034935 .000023082 000001685 000000017 .000004100 000574174 000001683 .000036893 000004104 000000171 000000000000 000574204 000006424 .000001684 .000000171 000000017 00000000000 .000036944 .000034957 000023055 000000000 000000171 000000017 43.003321 43.500321 44.004321 42.506321 5:15 BM 3/13/98 Ш Ш =XQ FLM= EL= 믭 =X= FL= =I'IJ FLI= FLJ= FLM= Ē = E =I'ILI FLI= ELL-FLM= 님 믭 =XC Ŗ = FLI= FLI= =IJT ELM= FLM= ş <u>_</u>___ FLI= FLM ŝ FILM= FLM= 2 FLM= FI.M= FLM= 밀 DX= Ë ¥ × ×

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3/13/98 5:15 PM Hard Disk 2150:MPW: AEDC ADVERSE WEATHER 3/98:DEVELPMENTS OF AEDC MFF CODE: AEDCIDMP ...: AEDCIDMP.OUT

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=11 =11 =W1A =W1A =W1A =W1A =17A	YDUCT= A(X)=	=WIN =WIN =ITA =ITA =IA =DV	VLM= VL= VL= A(X) =	=11 =W10 =W10 =110 =110 =110 =110 =110 =	YDUCT= A(X)=	=DV =LI =LI =LI =LI =LI =LI =LI =LI =LI = =LI = =LI = =LI = = = =
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