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THESIS

EXPERIMENTAL HEAT EXCHANGER PERFORMANCE IN A THERMOACOUSTIC PRIME MOVER

by

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December, 1993

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by

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Lieutenant, United States Navy

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This thesis investigates the experimental heat exchanger performance in a neon filled thermoacoustic prime mover. The experimental approach is to measure the waveform and spectrum of the acoustic oscillations, as well as the relevant temperatures for heat exchangers 0.257, 0.569, and 0.82 cm in length. A temperature gradient is established across the stack by submerging the cold heat exchanger and cold end tube in liquid nitrogen and keeping the hot heat exchanger and hot end tube at ambient temperature. Measurements are made at various mean gas pressures ranging from 1.5 to 50 kPa and for various effective positions of the stack in the standing wave. Acoustic pressure amplitudes as high as 29% of mean gas pressure are generated by the prime mover. The primary experimental controls over heat exchanger performance are the various heat exchanger lengths mentioned above, and the control of the thermal penetration depth, which decreases with increasing mean gas pressure. Results indicated that the prime mover can generate peak-to-peak displacement amplitudes that are much longer than the heat exchanger lengths.

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LIST OF SYMBOLS

- c sound speed
- c_p isobaric heat capacity per unit mass
- d distance from heat exchanger or stack to plunger
- Q heat
- f frequency
- k wavenumber
- P_m mean gas pressure
- *P*_o peak pressure amplitude
- T temperature
- t_D dwell time in heat exchanger
- *u* longitudinal component of velocity
- *u_o* peak longitudinal component of velocity
- *x* longitudinal distance
- y transverse distance
- y_o plate half-gap
- β exponent for temperature dependence of dynamic viscosity
- γ ratio of isobaric to isochoric specific heats
- δ penetration depth
- Δx plate length
- κ thermal conductivity
- λ wavelength
- ξ displacement amplitude
- ξ_o peak displacement amplitude
- Π perimeter
- ρ density

- τ normalized dimensionless thermal equilibration time
- χ thermal diffusivity
- ω angular frequency

Subscrits and superscripts

- c cold
- h hot
- hx heat exchanger
- s stack
- κ thermal
- v viscous

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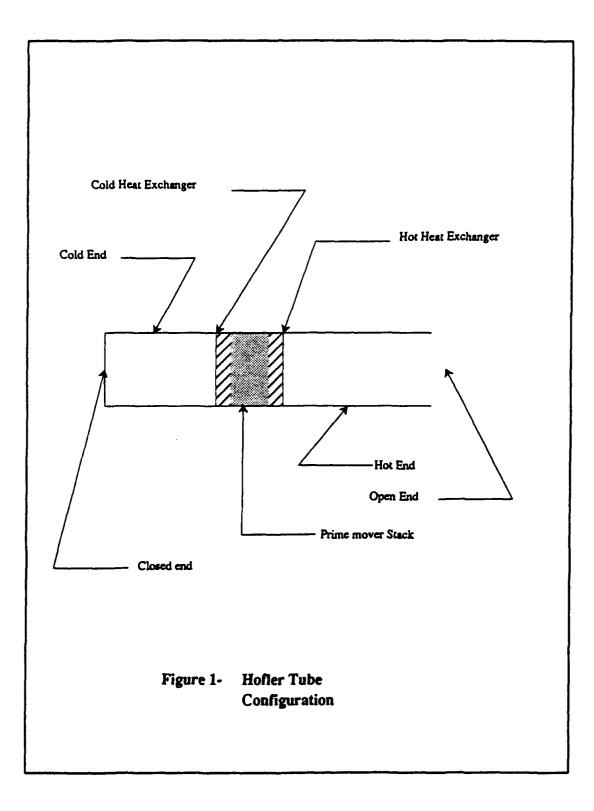
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I. INTRODUCTION

There are two types of thermoacoustic engines: heat pumps and prime movers [Refs 1-4]. A thermoacoustic heat pump uses a high amplitude acoustic standing wave to transport (or pump) heat along the boundary of a plate situated in the standing wave. The acoustically generated heat flow produces a thermal gradient across the plate. Acoustic energy is converted into stored thermal energy, which in turn can be used for other practical applications.

Thermoacoustic prime movers convert stored thermal energy into useful work in the form of sound. Such a device is the subject of this thesis. An example of a prime mover is shown in Figure 1. This particular configuration is referred to as the "Hofler Tube," described elsewhere [Ref 2]. The prime mover configuration consists of a stack of plates, called the prime mover stack (from here on referred as "the stack"), which is in thermal contact with two heat exchangers. In this thesis, the two heat exchangers are called the hot and cold heat exchangers. One end of the stack is held at ambient (or hot) temperature, and the other end of the stack is held near liquid nitrogen (or cold) temperature. The stack and heat exchangers are part of an acoustic resonator. The assembly consisting of a stack, heat exchangers, and a resonator is called a thermoacoustic prime mover.

Thermal energy is stored in the prime mover by imposing a temperature difference across the stack. In order for the prime mover to produce net positive work (i.e., to produce sound), the amount of stored energy converted to sound must exceed



the amount of acoustic energy dissipated by losses within the prime mover. The dominant loss mechanisms are thermal and viscous losses at the resonator walls, stack surface, and heat exchanger surfaces. The prime mover reaches "onset" when the temperature difference across the stack is sufficient for the prime mover to generate and sustain detectable levels of sound. Typically, the steady state acoustic pressure amplitudes are 1 to 10% of the mean gas pressure.

The main purpose of this thesis is an experimental investigation of heat exchanger performance in a thermoacoustic engine. A prime mover is chosen as a vehicle for these heat exchanger experiments partially because improving the performance of the thermoacoustic prime movers is valuable in its own right, but primarily because achieving the desired high level acoustic amplitudes and gas displacement amplitudes in a thermoacoustic refrigerator is a formidable engineering challenge. Achieving large gas displacement amplitudes with prime movers is generally much simpler, if sufficient amounts of heat can be sourced and sinked at the heat exchanger. It is hoped that the lessons learned in these experiments, will also be applicable to thermoacoustic refrigerators.

The prime mover is filled with low pressure neon gas to allow for low thermal heat flux at high pressure ratios (P_o/P_m). The neon gas in the prime mover is pressurized to approximately 1.5, 8.8, 10.9, 27.6, and 50 kPa. A temperature gradient is established across the stack by keeping the cold heat exchanger submerged in liquid nitrogen and the hot heat exchanger at ambient temperature. To determine the heat exchanger performance, the spectrum of acoustic waveform at steady state, the acoustic pressure amplitude, and the temperature at strategic locations in the prime move are measured for each heat exchanger length at the indicated mean pressures.

II. THEORY

This chapter contains a discussion of some theoretical aspects of heat exchanger performance. The first part is a summary of relevant parts of Swift's review article [Ref 1] and Figure 2 is taken from this article to briefly show Swift's arguments.

Heat exchangers are required in thermoacoustic engines to supply and extract heat to and from the ends of the stack. Figure 2.a shows one plate, cross-hatched, with heat exchangers at each end labeled T_h and T_c . Figure 2.a shows five parcels of fluid near the plate, and their extreme positions. Thus, (a) and (a') represent, respectively, the maximum expanded and maximum compressed positions of parcel (a), and similarly for (b) and (b'), (c) and (c'), (d) and (d'), and (e) and (e'). The articulated cycle for the typical "midplate" parcel (a) consists of a rapid adiabatic motion (1->2) from (a) to (a'), a wait (2->3) at (a') for thermal equilibrium with the plate, a rapid adiabatic motion (3->4) from (a') to (a), followed by a wait (4->1) again for thermal equilibrium.

One result of this cycle is a time-averaged hydrodynamic transport of heat along the direction of acoustic oscillation. For clarity in Figure 2.a, the parcels are drawn nonoverlapping, but in reality (a') and (b), or (d') and (a) occupy the same locations. It is easy to see that, at midplate locations, the time-averaged heat transfer from fluid to plate is zero; for instance, the dQ delivered to the plate by (a) is removed from the plate a half-cycle later by (d').

Now, observe what happens at the ends of the plate near the heat exchangers. For example, consider parcel (c), which makes thermal contact with the heat exchanger at point (c) of its cycle but is out of thermal contact at location (c'). After compression (1->2), the parcel is hotter at (c') than it was at (c), by the peak-to-peak adiabatic temperature oscillation, but it has no means at (c') to thermally relax. Moreover, nothing happens to it during the waiting period (2->3), so on decompression (3->4) it comes back at (c) to the same temperature that it had at (1). Parcel (c) just idles, not contributing to any time-averaged heat flow, even though it is hotter at (c') than at (c).

Hence, there is no compensating heat flow from (c) to match the dQ removed by (b') from the heat exchanger at T_h . In this simple picture, dQ is extracted from the reservoir at T_h by (b'), passed on down the plate in successive cycles by successive parcels, and then rejected to the reservoir at T_c by (d).

It is understood from these simple arguments that the heat exchanger should have a length of about twice the displacement amplitude, ξ_{hx} , of the gas parcel at the heat exchanger. The displacement amplitude is derived by using the Mach number relation,

$$Ma = \frac{u_o}{c} = \frac{1}{\gamma} \frac{P_o}{P_m}.$$

The relation between displacement and velocity amplitude for sinusoidal waves,

$$\xi_o=\frac{u_o}{\omega}.$$

The local peak displacement amplitude is given by,

$$\xi_{hx} = \xi_0 \sin(k x_{hx}).$$

For this experiment, it is assumed that $kx_{hx} \ll 1$ which allows for the approximation that $\xi_{hx} \simeq \xi_0 kx_{hx}$. Now, using the Mach number relation and the peak displacement amplitude relation, the local displacement amplitude equation is given by,

$$\xi_{hx} = \frac{1}{\gamma} \frac{P_o}{P_m} x_{hx}.$$
 [1]

Figure 2.b shows a heat exchanger that is longer than $2\xi_{hx}$. In that case, there are some parcels of fluid that contact only the heat exchanger, at both extremes of their positions, so that they perform no useful function. Their presence is purely dissipative, by virtue of both viscous and thermal processes.

On the other hand, Figure 2.c shows a heat exchanger that is shorter than $2\xi_{hx}$, so that there are parcels that "jump" from the plate to well past the heat exchanger, and contact the heat exchanger only for a very short time. Such parcels are, in Swift's estimation, ineffective in carrying heat to and from the heat exchangers.

This is because Swift assumes that a thermal boundary layer of gas exits between the parcel and the plate, and that this boundary layer represents too large of a thermal impedance for a significant amount of heat to diffuse through the boundary layer in such a short time. However, Hofler has proposed a simple heuristic model [Ref 5] that suggests that a heat exchanger can be both thermally effective and short in the longitudinal direction $\Delta x_{hx} \ll 2\xi_{hx}$, if the plate separation, in the transverse direction, is comparable to a thermal penetration depth $2y_o \approx \delta_{\kappa}$.

The model begins with a number of simplifying assumptions. Gas viscosity and compressibility are ignored as are the details of the thermal contact between the stack and the gas parcels. The latter issue is complicated by the fact that the stack has a longitudinal temperature distribution T(x). The thickness of the plates is assumed to be negligibly small. Referring to Figure 3, we will also assume that there is a discontinuity in the temperature T(x), between the end of the stack and the beginning of the heat exchanger plates.

Furthermore, the attention focuses on the second gas parcel from the top of Figure 3, which is the parcel with the shortest dwell time, t_D , in the heat exchanger. This is the shortest dwell time because the parcel has its peak velocity at the center of its excursion, or approximately $t_D \simeq \Delta x_{hx} / \omega \xi_{hx}$. This particular parcel is important because it has the poorest thermal contact with the heat exchanger.

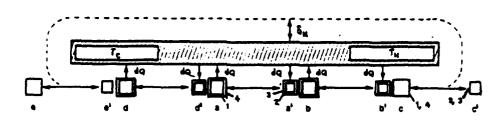


Figure 2.a- A set of oscillating fluid parcels about a thermal penetration depth from one plate of a thermoacoustic engine, with exchangers as its ends.

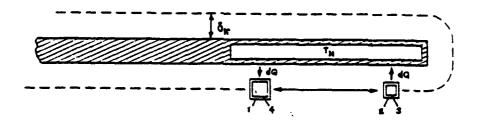


Figure 2.b- if the heat exchanger is too long, the parcel of fluid shown merely shuttles heat from one part of the heat exchanger to another, absorbing work.

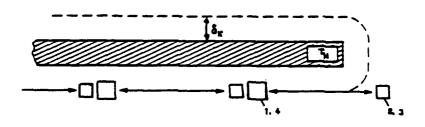
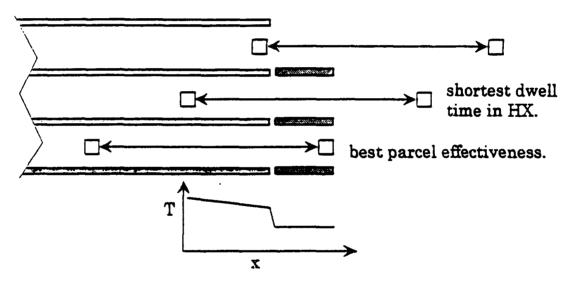


Figure 2.c- If the heat exchanger is too short, the chain of parcels shown has no source of heat.





The question of thermal effectiveness can be stated thus: How long does the parcel have to remain in the heat exchanger to thermally equilibrate with the exchanger plates? It is important to note that perfect thermal equilibration is probably not a sensible goal, because this will require long heat exchangers, such as those in Swift's analysis.

Long heat exchangers are not desirable because high acoustic dissipation is incurred. A more sensible goal might be 90% thermal equilibration. That is, the final temperature of the gas parcel after passing through the exchanger is different from that of the exchanger by an amount that is only 10% of the temperature discontinuity in T(x).

The problem can be recasted from one of a static temperature distribution in the plates and moving gas parcels, to one of a stationary slab of gas having a time dependent temperature at its boundaries. In Figure 4, it is assumed that the temperature of the slab boundaries has a step discontinuity in the temperature at time t=0 and lets the gas thermally equilibrate for a time equal to the dwell time, t_D .

The solution to the ratio of the gas temperature to the boundary temperature, averaged over the slab, can be found in a standard text [Ref 6]. This solution is indicated by the solid curve in Figure 5, where a y-axis value of one represents perfect equilibration, and the x-axis values are normalized time, $\tau = \chi t_D / y_o^2$. If the approximate value for the dwell time is substituted into the expression for τ , and the definition for the thermal penetration depth, $\delta_{\kappa} = \sqrt{2\chi/\omega}$, is used to obtain,

$$\tau \simeq \left[2\left(\frac{\xi_{hx}}{\Delta x_{hx}}\right)\left(\frac{y_o}{\delta_{\kappa}}\right)^2\right]^{-1}.$$

The solid curve represents the amount of equilibration after one pass through the exchanger and the long dash curve represents the amount of equilibration after two

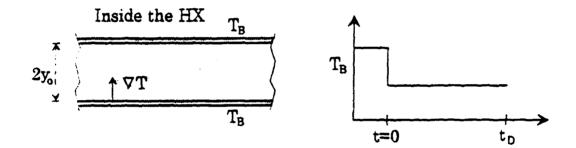


Figure 4- Temperature distribution in stationary plate.

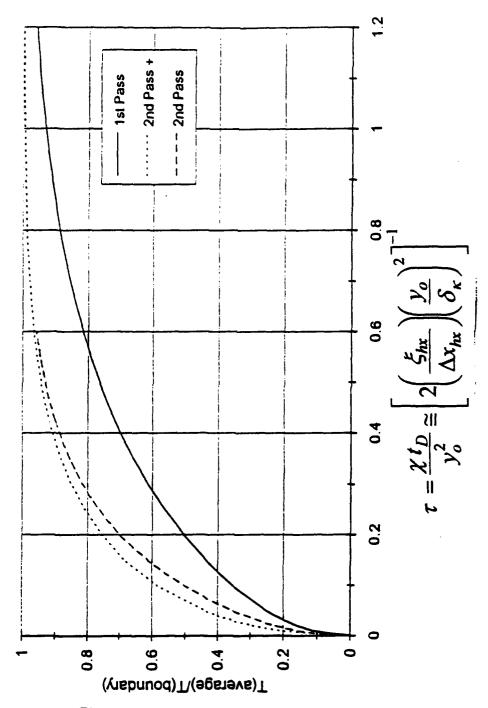


Figure 5- Averaged thermal equilibration.

passes, in one acoustic cycle. This assumes that the gas parcels are doing nothing when they are not in contact with either the stack or the heat exchanger. That is, they are thermally static.

In fact, the temperature of the gas parcels may be stratified in the transverse direction after one pass through the exchanger and they will be thermally equilibrating with their nearest neighbors. The short dash curve assumes perfect equilibration of the gas with itself, in between the first and second passes. The true solution will lie somewhere between the two dashed curves.

Note some values of equilibration versus τ . For $\tau = 0.25$, the equilibration is about 77%; for $\tau = 0.5$, the equilibration is about 93%; and for $\tau = 0.75$, the equilibration is about 98%. In Figure 6, the hatched area indicates the range of parameters y_o/δ_{κ} and $\xi_{hx}/\Delta x_{hx}$, for which the worst case of thermal effectiveness is between 77 and 93%. Short heat exchangers with $\xi_{hx}/\Delta x_{hx}$ in the range of 4 to 8 can be thermally effective if y_o/δ_{κ} is in the range of 0.75 to 0.5.

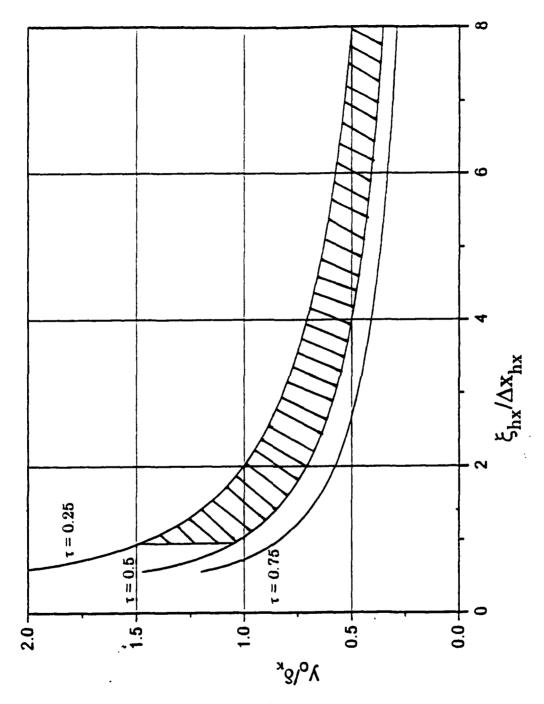


Figure 6- Curves of thermal effectiveness.

III. EXPERIMENTAL APPARATUS, INSTRUMENTATION, AND PROCEDURE

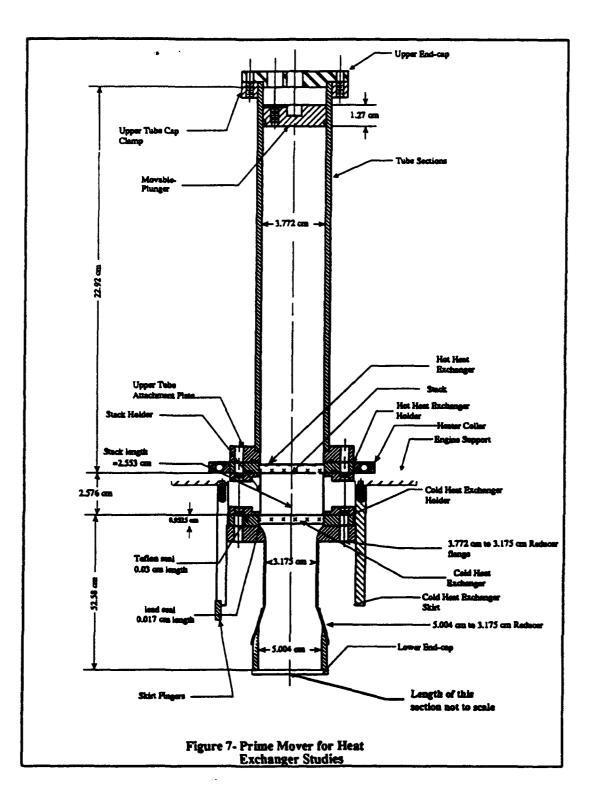
The primary goal of this thesis is to measure the pressure ratio, the steady state spectrum waveform, and temperatures at various locations in the prime mover. Figure 7 shows an overview of the main components of the prime mover. The prime mover is specifically designed to generate very high pressure ratios. Notice that the lower part of the drawing is not to scale. Detailed engineering drawings are found in Appendix A.

A. THE THERMOACOUSTIC PRIME MOVER

1. Hot End Tube Assembly

The hot end assembly consists of the tube, attachment plate, and an assembly consisting of an end-cap, end-cap clamp, and movable plunger. The upper copper tube (see A-1) is 3.81 cm I.D. and 22.018 cm long. The lower end of the tube is fitted with a 7.366 cm O.D., 3.81 cm I.D. and 0.9525 cm long copper attachment plate (see A-2) that has eight radial holes to be used as thermocouple feed-throughs to the inside of the prime mover. The attachment plate has a groove for an O-ring and eight holes used to bolt the assembly to the hot heat exchanger and stack holder.

The upper end of the tube is fitted with an assembly consisting of an endcap, end-cap clamp, and movable plunger. The copper end-cap (see A-3) fits into the copper tube and is held in place by a control rod threaded in a 6061-T6 Aluminum clamp (see A-4). The end-cap has three holes that allows free passage for the control rod, a transducer cable, and a 0.3175 *cm* copper gas line that attach to the movable plunger.



The copper movable plunger (see A-5) is 3.81 cm in diameter and 1.27 cm long. The holes for the transducer and the copper line are bored through the plunger to provide access to the interior of the prime mover. The plunger has a groove for an O-ring. The copper gas line is soldered into the hole. A copper-nickel capillary tube of 0.025 mm O.D. is inserted and soldered inside the 0.3175 cm copper gas line. The capillary provides a high impedance to the acoustic pressure oscillations.

The total internal length of the hot end tube assembly from the top most position of the plunger face to attachment plate bottom is 21.64 cm.

2. Hot Heat Exchanger Assembly

The assembly consists of hot heat exchanger, hot heat exchanger holder, hot heat exchanger heater collar, and spacers (if necessary). The purpose of the hot heat exchanger is to supply heat to the stack.

The copper heat exchangers (see A-6) are $3.924 \ cm$ in diameter and have lengths of 0.257, 0.569, and 0.82 cm. Each plate within the heat exchanger is 0.0292 cm thick. The plate gap is 0.0584 cm. The heat exchangers are soldered into the heat exchanger holder.

The heat exchanger holders (see A-7) are 7.366 cm O.D., 3.924 cm I.D., and 0.953 cm long. There is an O-ring groove in the side facing the stack. For the 0.257 and 0.569 cm heat exchangers, bronze spacers (see A-8) are used to support the heat exchangers flush with the side of the heat exchanger holder facing the stack.

The hot heat exchanger heater collar (see A-9) is made of copper and bolted onto the hot heat exchanger holder. It has a 7.37 cm I.D. and is 0.953 cm long. There are two 0.625 cm diameter holes fitted to hold two heater cartridges [Ref 7], used to maintain a constant hot heat exchanger temperature.

3. Stack Assembly

The prime mover stack assembly consists of the stack and the stack holder. The stack (see A-10) is made of polyester film of gauge 0.01 mm and 114.2 cm long. The film is rolled spirally around a phenolic core 0.635 cm in diameter and 2.529 cm long. The polyester film is separated by 0.851 mm diameter fishing line. The total spiral stack is 2.553 cm long, 3.772 cm O.D. and has a total surface area of $\Pi \Delta x = 583 \text{ cm}^2$.

The 304 stainless steel stack holder (see A-11) is 2.576 cm long and 3.772 cm I.D. The stack holder is 0.22 mm longer than the stack creating a 0.11 mm gap on each side of the stack inside the holder. The stack holder has eight 8-32 machine screw holes in each flange to attach the assembly consisting of the hot end tube and hot heat exchanger on one side and the assembly consisting of the cold end heat exchanger and cold end tube on the other.

4. Cold Heat Exchanger Assembly

In order to impose a temperature gradient across the stack, the cold heat exchanger is employed to remove heat from the stack and maintain one end of the stack close to liquid nitrogen temperature. The cold heat exchanger assembly is similar to the hot heat exchanger assembly, except that the cold heat exchanger (see A-12) holder has no O-ring groove. The cold heat exchanger is sealed by using a 0.03 *cm* thick teflon gasket in the stack side of the cold heat exchanger holder and a 0.17 *cm* lead wire diameter gasket in the cold end tube side of the cold heat exchanger holder.

5. Cold End Tube Assembly

The purpose of the cold end tube assembly is to provide a cold heat sink to the cold heat exchanger, determine the acoustic length, and to eliminate harmonics in the acoustic pressure waveform. It consists of a cold heat exchanger skirt, lower reducer flange, a 3.175 cm I.D. lower tube section, a cold end reducer, a 5.08 cm I.D. lower tube section, a lower tube end-cap, and an prime mover dewar support.

The purpose of the cold heat exchanger copper skirt (see A-13) is to provide a cold heat sink and to support the prime mover on the dewar. The skirt is 7.19 cm long. It has a 7.62 cm I.D. and 8.89 cm O.D. One end is bored to 8.64 cm I.D. and 4.763 cm deep to provide a lip to solder on to the lower reducer plate. The skirt has six 2.54 cm by 2.54 cm equally separated flat faces with 0.188 cm diameter holes to attach copper skirt fingers (see A-14) to increase the cold heat sink area. A 304 stainless steel dewar support plate (see A-15) is attached to the cold heat exchanger end of the skirt.

The purpose of the lower reducer copper flange (see A-16) is to provide the first change in diameter to eliminate the harmonics in the waveform and to conduct the liquid nitrogen temperature to the cold heat exchanger. The flange is 2.563 cm long and has a 7.72 cm O.D. The inside diameter starts at 3.797 cm on the face towards the cold heat exchanger and reduces to 3.23 cm. The plate is bored to 3.49 cm I.D. and 0.508 cm in depth to provide a lip to solder on to the 3.175 cm I.D. lower tube.

The 3.175 cm I.D. lower end copper tube (see A-17) is 26.99 cm long and has a 3.49 cm O.D. One end of the tube is soldered to the lower reducer flange and the other end to the cold end reducer. The purpose of the cold end copper reducer (see A-18) is to join the 3.175 cm I.D. lower tube to the 5.08 cm I.D. lower tube and provide the second diameter change to eliminate the harmonics. The cold end reducer is 2.95 cm long. The 5.08 cm I.D. lower end copper tube (see A-19) is 19.3 cm long and has a 5.595 cm O.D. The lower tube copper end-cap (see A-20) is soldered to the 5.08 cm I.D. lower tube. It is 1.016 cm long. It has a 5.395 cm I.D. and 5.76 cm O.D.

B. INSTRUMENTATION

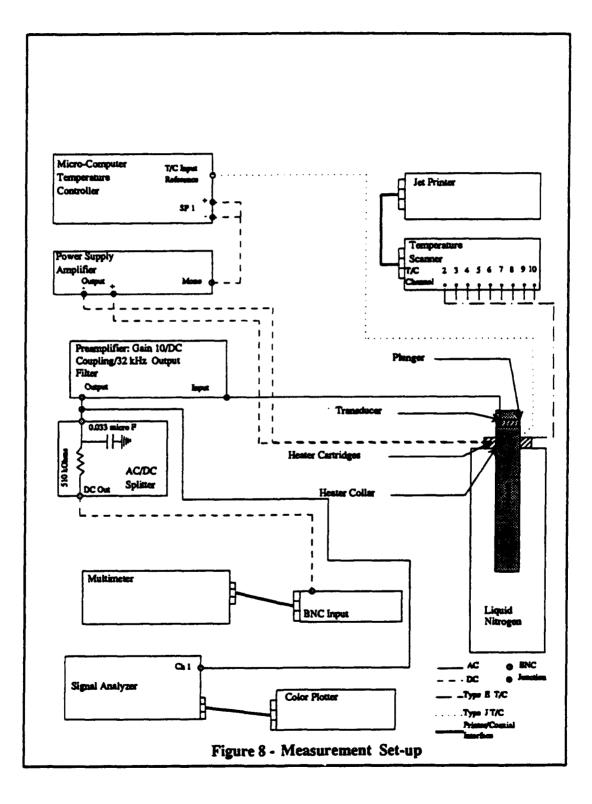
Figure 8 shows the overall measurement set-up utilized in this experiment.

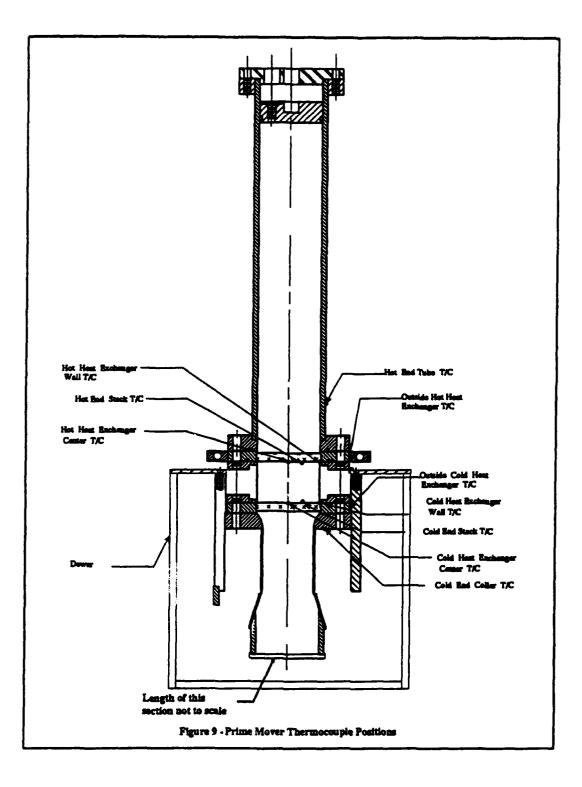
1. Temperature Measurement Equipment

Figure 9 shows the placement of the ten type E thermocouples (T/C) in the prime mover. These T/C are connected to the temperature scanner [Ref 8] which is connected to a printer [Ref 9] that provides printouts of all thermocouple readings. Each heat exchanger assembly has a T/C soldered at the center of the longest heat exchanger plate, the inside heat exchanger wall, and the outside heat exchanger holder wall. These points are chosen to provide a temperature map of the heat exchanger behavior in the radial direction. The T/C at the ends of the stack are placed there for two reasons: to provide an axial temperature measurement across the stack, and to provide an axial temperature measurement between the stack and center of the heat exchangers. The rest of the T/C are used to monitor resonator temperatures.

2. Temperature Control Equipment

Control of the hot heat exchanger is achieved by a digital temperature controller [Ref 10], heater cartridges, and a power supply amplifier [Ref 11]. The reference temperature for the temperature controller is taken from a type J T/C placed between the hot heat exchanger holder and the heater clamp. The signal is processed and is sent via the temperature controller's 10 volt DC analog output to the mono input in the power supply amplifier.





Power to the heater cartridges is provided by the output of the power supply amplifier. The heater cartridges are inserted into the copper heater collar. Once steady state thermal conditions are achieved in the experiment, the current and voltage delivered is approximately constant in time and can be measured in order to determine the heater power. We do not have a complete set of heater power measurements, because the heater controller circuitry was modified during the course of the experiment.

3. Pressure Measurement Equipment

The pressure amplitudes inside the prime mover are measured with a piezoresistive transducer [Ref 12] with a sensitivity of 16.38 mV/psi. The transducer is located on the prime mover plunger. The transducer senses absolute pressure and its zero offset was calibrated on a turbomolecular pump high vacuum system.

The transducer measures the mean pressure as a DC signal and the acoustic pressure amplitude as an AC signal, simultaneously. The transducer output is sent to a DC coupled preamplifier with a gain of 10 and a 32 kHz low-pass output filter. This output signal is split. One output is connected to channel 1 of the signal analyzer [Ref 13]. The signal analyzer displays the waveform and frequency spectrum of the AC component of the signal. These results are plotted on the color plotter [Ref 14]. The second output from the preamplifier is filtered to obtain a DC signal component by passing through a low-pass filter with a 0.033 μ F capacitor and 510 k Ω resistor. This DC voltage signal is then measured with the multimeter [Ref 15]. This allows the pressure ratio, P_o/P_m , to be accurately measured with which gas displacement amplitudes, ξ , can be calculated.

C. PROCEDURE

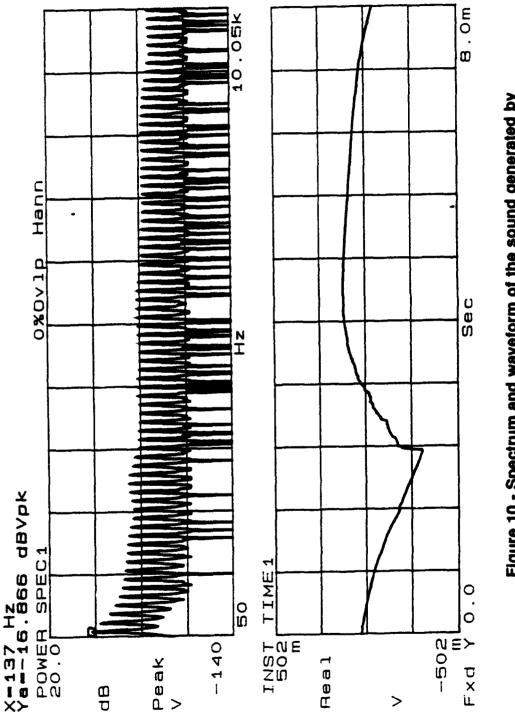
Two main modifications were made to the original straight bore prime mover design to achieve the desired experimental characteristics. The first modification was to reduce the harmonic distortion in the resonator. The second modification was to improve the heat flow within the prime mover parts.

As discussed at the beginning of this chapter, the primary goal is to measure the pressure ratio, temperatures, and steady state spectrum for three heat exchanger lengths. The measurement technique employed is to drive the prime mover by placing as big of a temperature gradient as possible across the stack at specific mean pressures. The design goal is to have the temperature gradient be nearly equal for all measurements. The three different heat exchanger lengths serve to vary $\xi/\Delta x$, and the differing mean pressures, change δ_{κ} and thus vary y_0/δ_{κ} . The measurements consisted of two phases: prime mover below onset and above onset.

1. Harmonic Distortion Suppression

The experiment was started using a straight bore resonator of 3.81 *cm* in diameter. However, the harmonic distortion was unacceptable as shown in Figure 10. There was a severe discontinuity in the waveform that produces the multiple harmonics observed in the frequency spectrum.

Large amounts of harmonic distortion of driven plane standing waves in constant bore tubes is common at higher amplitudes. This readily occurs when driving at the fundamental frequency because the higher modes of the tube have frequencies that are nearly integer multiples of the fundamental. Any distortion component contained in the drive waveform or arising from nonlinear propagation of the wave in the tube can coincide with one of the higher mode resonant frequencies and be "amplified" by the quality factor of that mode.





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More advanced theories for this type of waveform distortion suggest that there is an interaction between the first and second harmonics that generates the third harmonic, and so on; the acoustic energy cascades up to the higher frequencies. If this is true, then by modifying the interaction between the first and second harmonics, all of the higher harmonics may be suppressed.

This can be accomplished by shifting the frequency of the second mode relative to the first so that it is not exactly a factor of two higher. Changes in the cross section of the tube can produce this type of frequency shift, and can be very effective in suppressing harmonic distortion. This type of approach has been used for conventional piston driven tubes and thermoacoustically driven prime movers [Ref 16]. However, the thermoacoustic prime mover is unique in that it may spontaneously oscillate at a higher mode as well as at the fundamental. Whether or not this can occur, is an issue that is distinct from the issue of acoustic energy cascading to higher modes and suppression by frequency shifting.

A computer model for the thermoacoustic heat engines has been developed by Hofler based on the linear ordinary differentials equations of Rott [Ref 3]. This model has recently been improved to handle tubes with arbitrary cross-section correctly. This was used to propagate acoustic wave solutions from a starting point at the plunger, through the resonator having both a temperature distribution and a cross-section profile, and stopping at the cold termination. An arbitrary frequency is assumed at the beginning. While the acoustic boundary conditions are properly matched at the rigid plunger initially, the final boundary conditions at the cold termination are not generally satisfied because linear equations are being used to model a nonlinear system.

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At the rigid cold termination, the calculated acoustic work flow is generally not zero (as it must be in a real system) and this quantity can be interpreted as a thermoacoustic drive result. If it is positive, then the system will spontaneously oscillate at that modal frequency and if it is negative, then the mode will be damped. The frequencies of the modes can be determined by a peak in the magnitude of the acoustic impedance or by the imaginary part (i.e., the reactive part) of the impedance going to zero.

A frequency sweep of this type of computation is shown in Figure 11 for the prime mover with 3.81 *cm* and 3.17 *cm* bores. The first three modes can be seen from the solid curve representing the impedance magnitude. Note that the frequency of the second mode is somewhat higher than twice the first, although the frequency of the third mode is about three times the first. The height of the peaks is not important. The dashed curve shows that the work flow is positive for both the first and second modes, but negative for the third. This means that there will be substantial second harmonic distortion regardless of the frequency of the second mode.

To speed up the computational investigation, only the results at the peaks of the first three modes was generated. The frequencies of the first three peaks and work flows at those frequencies were generated for various resonator geometries having the original 3.81 cm and 3.17 cm diameters, but with a new larger third diameter at the cold end. The diameters investigated were 3.81, 4.45, and 5.08 cm. The length of the new tube section was a second variable.

These results are shown in Figure 12. The x-axis is the frequency normalized with respect to the fundamental. The two 4.45 cm diameter data points represent different lengths. Note that while both the 4.45 and 5.08 cm diameter tube segments shift the frequencies of the second and third modes substantially, only the 5.08 cm diameter tube segment causes the work flow to be negative at the second mode. If the second mode is damped and its frequency is substantially shifted from twice the fundamental, the harmonic distortion should be substantially suppressed.

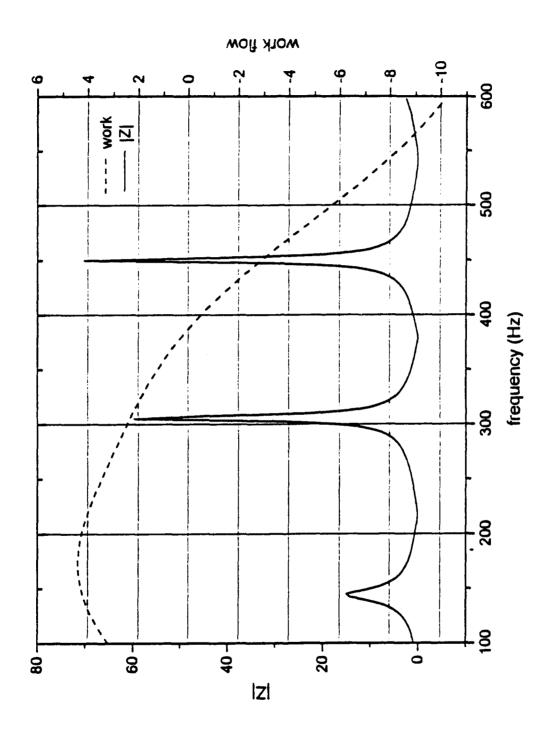


Figure 11 - Impedance and work flow at cold termination.

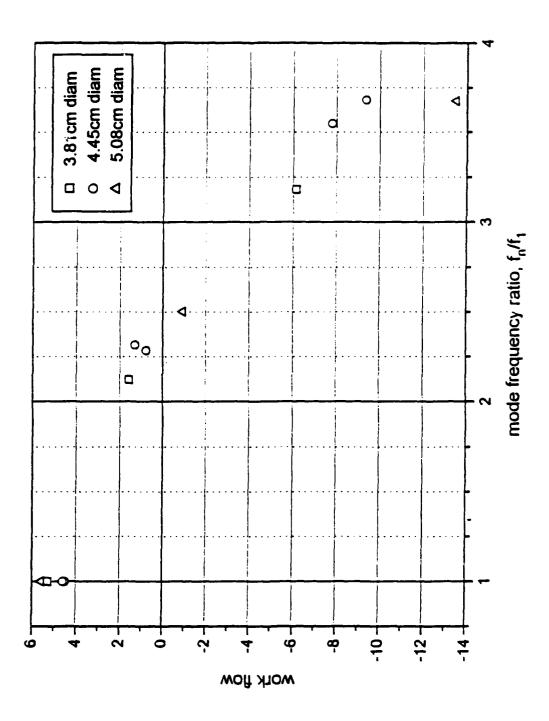


Figure 12 - Modes with three cold end diameters.

From the computer simulation results, geometric changes were done in the prime mover to reduce the harmonic distortion. First, the cold end resonator diameter was reduced from 3.81 cm to 3.17 cm. The waveform and frequency spectrum remained basically unchanged from the straight bore configuration.

The next step was to have two changes in the diameter of the cold end from 3.81 cm to 3.17 cm and back to 3.81 cm. This new geometry generated the sound waveform and spectrum shown in Figure 13. The discontinuity in the waveform was less dramatic in this prime mover configuration.

The final configuration used for this experiment is the one shown in Figure 3 of Chapter III, where the diameter of the resonator was decreased from $3.81 \ cm$ to $3.17 \ cm$ and then increased to $5.08 \ cm$. Appendix B shows the smooth sinusoidal waveforms and the spectra of the sound generated by the prime mover above onset at a $8.128 \ cm$ plunger position.

2. Heat Flow Modifications

The goal was to have the cold and hot heat exchanger temperatures to remain constant through this experiment. However, after the first experimental trial, it was observed that the heat exchanger temperatures were not constant. As a result, several modifications were performed. Again, drawings and specifications for the prime mover parts are provided in Appendix A.

Originally, heat was supplied by a nozzle heater clamped around the hot end tube just above the attachment plate. This heater was replaced with a heater collar clamped around the hot heat exchanger holder. This provides a constant supply of heat source going into the prime mover. The conductance of heat is improved by applying a silicone heat sink compound [Ref 17] at the side of the hot heat exchanger holder that faces the upper attachment plate. This modification worked fairly well in keeping the hot end of the prime mover at close to ambient temperature.

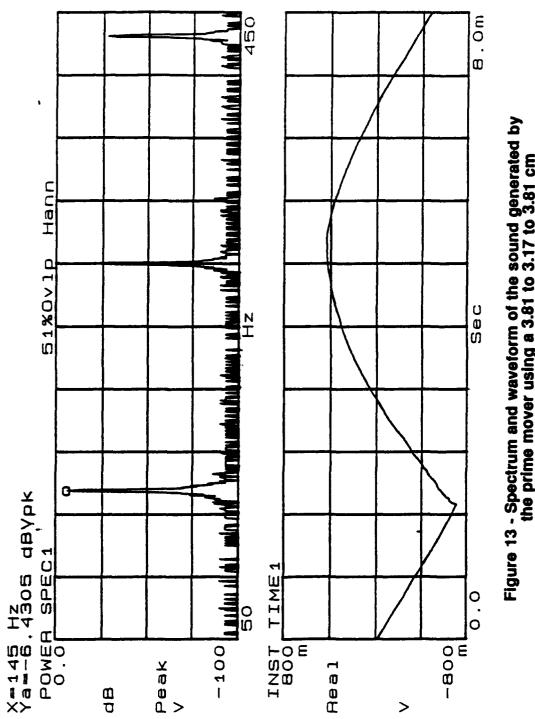


Figure 13 - Spectrum and waveform of the sound generated by the prime mover using a 3.81 to 3.17 to 3.81 cm diameter tube configuration.

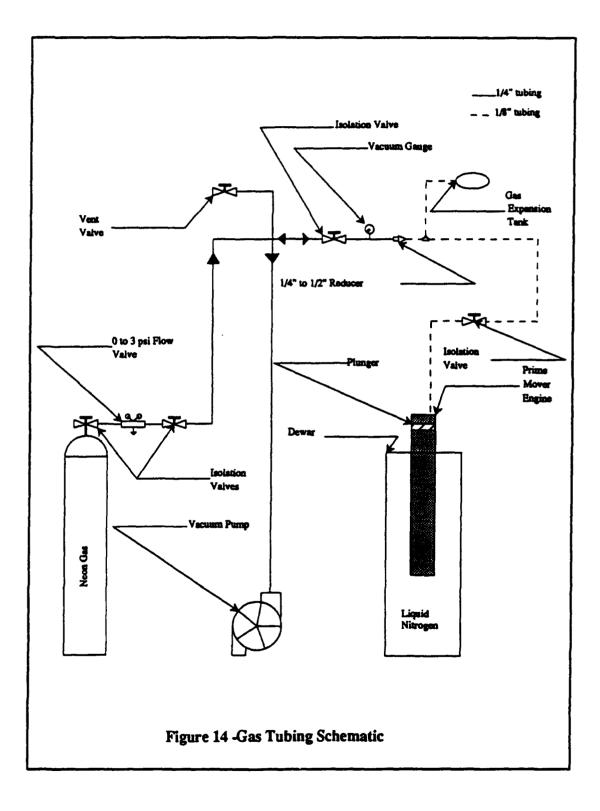
The heat flow into the cold heat exchanger was improved by increasing the surface area in contact with the liquid nitrogen. The cold heat exchanger copper skirt was modified to fit four copper skirt fingers. The conduction between the fingers and the skirt is improved by applying silicone heat sink compound between them. The thermal conductance between the cold heat exchanger holder and the lower reducer flange is improved by replacing the original teflon gasket with a conducting lead wire that is crushed between the two flanges.

Lastly, the stack holder is insulated from the ambient air and the hot end of the prime mover. The hot end tube and hot heat exchanger assemblies are insulated from the stack holder by wrapping them with fiberglass cloth. This increases the heat flux going into the hot heat exchanger, instead of the environment. Also, the cold end of the prime mover is insulated from the stack and the hot end of the prime mover by covering the dewar support and cold end skirt exposed to the environment with fiberglass cloth.

3. The Prime Mover Below Onset

The first phase of this experiment consists of measuring the thermal load of the prime mover below onset. Figure 14 shows the gas tubing schematic. The prime mover is pumped down using the vacuum pump to approximately 7 kPa and filled back with neon to approximately 110 kPa. This procedure is performed three times to ensure the gas inside the prime mover is substantially pure neon.

The prime mover is then pumped down to approximately 1.5 kPa. The prime mover is below onset at this pressure. The hot heat exhanger electric heater temperature is set at $25^{\circ}C$ and the dewar is filled with liquid nitrogen. The hot end tube assembly and the hot heat exchanger collar are insulated from the room environment by covering it with fiberglass cloth. After thermal equilibrium is reached,



the prime mover temperatures are obtained from the temperature scanner. This procedure is repeated for each heat exchanger length. Also, the electric power delivered to the heater can be measured with zero thermoacoustic heat transport. This power represents a nuisance heat leak to the colder parts below.

4. The Prime Mover Above Onset

The second phase consists of measuring the temperatures, the pressure ratio, and steady state waveform for two plunger positions at four mean pressures above onset. The hot heat exchanger electric heater temperature is set to $35 \degree C$, so at higher mean pressures the temperature in the wall of the hot heat exchanger does not drop below freezing. The plunger positions are measured from lower face of the plunger to the hot end of the stack. The two plunger positions used are 8.13 cm and 10.66 cm. The pressure is increased by using the low pressure flow valve to gradually fill the prime mover with neon to the selected pressures. The measurements are taken at mean pressures of approximately 8.8, 10.9, 27.6, and 50 kPa.

After the prime mover reaches steady state, the prime mover temperatures are obtained from the temperature scanner. The acoustic pressure amplitude and steady state frequency response are obtained from the signal analyzer. The mean pressure is obtained from the multimeter. This procedure is repeated for each heat exchanger length.

IV. RESULTS AND DISCUSSIONS

The results of the experiment are presented and discussed in this chapter. The results for the pressure ratios, the temperatures, the penetration depths, and the displacement amplitudes generated in the prime mover are presented for each heat exchanger lengt.

Figure 15 shows the ratio of the acoustic pressure amplitude of the fundamental component to the mean gas pressure in percent versus mean gas pressure for each heat exchanger length at two plunger positions. Table 1 shows the temperatures in the prime mover for each heat exchanger length at the selected experimental mean gas pressures for two plunger positions. Figure 16 shows the ratio of the center hot heat exchanger temperature to the center cold heat exchanger temperature versus mean gas pressure for each heat exchanger length. Figures 17 and 18 show the ratio of the heat exchanger and stack half plate separation to the thermal penetration depth versus mean gas pressure for each heat exchanger length. Figure 19 shows the ratio of the each heat exchanger length. Figure 19 shows the ratio of the each heat exchanger length versus mean gas pressure for each heat exchanger length. Figure 19 shows the ratio of the peak displacement amplitude at the midpoint of the element to the stack and heat exchanger length versus mean gas pressure for each heat exchanger for each heat exchanger length.

A. THE 0.257 cm HEAT EXCHANGER

The first mode acoustic pressure amplitudes are 1.082, 1.791, 5.366, and 8.84 kPa at mean gas pressures of 8.794, 10.92, 27.62, and 50.12 kPa, respectively. As shown in Figure 15, these amplitudes correspond to 12.3, 16.5, 19.4, and 17.6% of mean gas pressure. The first mode acoustic pressure amplitudes for a 10.688 *cm* plunger position are 0.749, 1.532, 5.021, and 8.05 kPa at mean gas pressures of 8.69, 10.81, 27.44, and 49.7 kPa, respectively. As shown in Figure 15, these amplitudes correspond to 8.62, 14.2, 18.3, and 16.2% of mean gas pressure.

From Table 1, the temperature differences across the stack vary from 172.7 to 77.8 $^{\circ}C$, from the lowest to the highest mean gas pressure. These temperatures correspond to temperature differences from the wall to the center of the cold heat exchanger of 2.4 to 9.2 $^{\circ}C$ and temperature differences from the wall to the center of the hot heat exchanger of 3.9 to 32.5 $^{\circ}C$.

As shown by Figure 16, the T_h / T_c ratios change from 3.51 to 2.02 for a 8.128 cm plunger position. This represents a T_h / T_c ratio decrease of 42% from the lowest to the highest mean gas pressure. The T_h / T_c ratios change from 3.53 to 1.97 for a 10.688 cm plunger position. This represents a T_h / T_c ratio decrease of 44% from the lowest to the highest mean gas pressure.

For purposes of reducing the experimental data, the thermal penetration depth is obtained for each pressure and frequency by using Equation [2],

$$\delta_{\kappa} = \sqrt{\frac{\kappa}{\pi f \rho c_p}}.$$
 [2]

Now, the relation $\rho c_p = \gamma P_m / T_m (\gamma - 1)$ and $T_m = 296 K$ are substituted into Equation [2] to obtain,

$$\delta_{\kappa}(T_m) = \sqrt{\frac{\kappa (\gamma - 1) T_m}{\pi f \gamma P_m}}.$$
[3]

For neon at 296 K, $\kappa = 0.0493 W/m$ K. Then the thermal penetration depth is obtained for each temperature, for its respective pressure and frequency, by using Equation [4],

$$\delta_{\kappa}(T_m) = \delta_{\kappa}(296 \ K) \ (T_m \ / \ 296)^{\frac{(1+\beta)}{2}}.$$
[4]

For neon, $\beta = 0.66$.

From Figure 17, the y_o/δ_{κ} ratios for the cold heat exchanger change from 0.638 to 1.066. The thermal penetration depth change is caused by the decreasing T_h / T_c ratio and the increasing temperature in the cold heat exchanger with increasing mean pressure. This represents a y_o/δ_{κ} ratio increase of 67% from the lowest to the highest

mean gas pressure. The y_o/δ_{κ} ratios for the hot heat exchanger change from 0.224 to 0.594. This represents a y_o/δ_{κ} ratio increase of 165% from the lowest to the highest mean gas pressure. From Figure 18, the y_o/δ_{κ} ratios for the stack change from 0.56 to 1.234. This represents a y_o/δ_{κ} ratio increase of 120% from the lowest to the highest mean gas pressure.

The displacement amplitude is calculated using Equation [2]. From Figure 19, the $\xi_{hx} / \Delta x_{hx}$ ratios for the cold heat exchanger change from 3.11 to 4.45, from the lowest to the highest mean gas pressure. The peak displacement amplitudes are as much as five times the heat exchanger length! For the hot heat exchanger, the $\xi_{hx} / \Delta x_{hx}$ ratios change from 2.25 to 3.21, from the lowest to the highest mean gas pressure. Note: the peak displacement amplitudes are as much as three and one-half times the heat exchanger length! For the stack, the $\xi_s / \Delta x_s$ ratios change from 0.27 to 0.39, from the lowest to the highest mean gas pressure. The peak-to-peak displacement amplitudes, except for the lowest pressure, are nearly as long as the stack length!

B. THE 0.569 cm HEAT EXCHANGER

The first mode acoustic pressure amplitudes are 0.398, 1.21, 4.65, and 9.0 kPa at mean gas pressures of 8.82, 10.91, 27.81, and 50.16 kPa, respectively. As shown in Figure 15, these amplitudes correspond to 4.5, 11.1, 16.7, and 17.9% of mean gas pressure. The first mode acoustic pressure amplitudes at a 10.688 *cm* plunger position are 0.885, 3.32, and 7.46 kPa at mean gas pressures of 10.77, 27.42, and 49.53 kPa, respectively. As shown in Figure 15, these amplitudes correspond to 8.2, 12.1, and 15.1% of mean gas pressure.

From Table 1, the temperature differences across the stack are from 186.3 to $65.4^{\circ}C$, from the lowest to the highest mean gas pressure. These temperatures correspond to temperature differences from the wall to the center of the cold heat exchanger of 0.6 to $28.5^{\circ}C$ and temperature differences from the wall to the center of the center of the hot heat exchanger of 0 to $12.4^{\circ}C$.

As shown by Figure 16, the T_h / T_c ratios change from 3.59 to 1.67 for a 8.128 *cm* plunger position. This represents a T_h / T_c ratio decrease of 53% from the lowest to the highest mean gas pressure. The T_h / T_c ratios change from 3.70 to 1.62 for a 10.688 *cm* plunger position. This represents a T_h / T_c ratio decrease of 56% from the lowest to the highest mean gas pressure.

From Figure 17, the y_0/δ_{κ} ratios for the cold heat exchanger change from 0.645 to 0.872. This represents a y_0/δ_{κ} ratio increase of 35% from the lowest to the highest mean gas pressure. The y_0/δ_{κ} ratios for the hot heat exchanger change from 0.223 to 0.574. This represents a y_0/δ_{κ} ratio increase of 157% from the lowest to the highest mean gas pressure. From Figure 18, the y_0/δ_{κ} ratios for the stack change from 0.559 to 1.154. This represents a y_0/δ_{κ} ratio increase of 106% from the lowest to the highest mean gas pressure.

From Figure 19, the $\xi_{hx} / \Delta x_{hx}$ ratios for the cold heat exchanger change from 0.523 to 2.08, from the lowest to the highest mean gas pressure. The peak displacement amplitudes are as much as two times the heat exchanger length. For the hot heat exchanger, the $\xi_{hx} / \Delta x_{hx}$ ratios change from 0.373 to 1.484, from the lowest to the highest mean gas pressure. The peak displacement amplitudes are as much as one and one-half times the heat exchanger length. For the stack, the $\xi_s / \Delta x_s$ ratios change from 0.099 to 0.396, from the lowest to the highest mean gas pressure. The peak to the highest mean gas pressure. The peak displacement amplitudes are as much as one and one-half times the heat exchanger length. For the stack, the $\xi_s / \Delta x_s$ ratios change from 0.099 to 0.396, from the lowest to the highest mean gas pressure. The peak-to-peak displacement amplitudes, except for the lowest pressure, are as long as the stack length.

C. THE 0.820 cm HEAT EXCHANGER

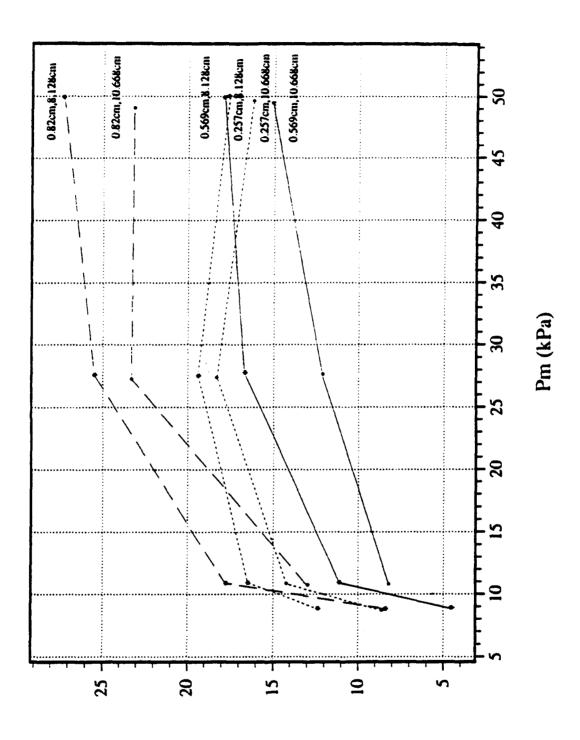
The first mode acoustic pressure amplitudes are 0.74, 1.944, 7.07, and 13.68 kPa at mean gas pressures of 8.814, 10.898, 27.62, and 49.95 kPa, respectively. As shown in Figure 15, these amplitudes correspond to 8.39, 17.8, 25.5, and 27.3% of mean gas pressure. The first mode acoustic pressure amplitudes at a 10.688 *cm* plunger position are 1.385, 6.355, and 11.36 kPa at mean gas pressures of 10.668, 27.287, and 48.966 kPa, respectively. As shown in Figure 15, these amplitudes correspond to 12.98, 23.3, and 23.2% of mean gas pressure.

From Table 1, the temperature differences across the stack are from 176 to 73.2 $^{\circ}C$, from the lowest to the highest mean gas pressure. These temperatures correspond to temperature differences from the wall to the center of the cold heat exchanger of 0.9 to 13.7 $^{\circ}C$ and temperature differences from the wall to the center of the hot heat exchanger of 0 to 24.9 $^{\circ}C$.

As shown by Figure 16, the T_h / T_c ratios change from 3.71 to 2.22 for a 8.128 cm plunger position. This represents a T_h / T_c ratio decrease of 40% from the lowest to the highest mean gas pressure. The T_h / T_c ratios change from 3.76 to 2.164 for a 10.688 cm plunger position. This represents a T_h / T_c ratio decrease of 42% from the lowest to the highest mean gas pressure.

From Figure 17, the y_o/δ_{κ} ratios for the cold heat exchanger change from 0.661 to 1.124. This represents a y_o/δ_{κ} ratio increase of 88% from the lowest to the highest mean gas pressure. The y_o/δ_{κ} ratios for the hot heat exchanger change from 0.222 to 0.579. This represents a y_o/δ_{κ} ratio increase of 161% from the lowest to the highest mean gas pressure. From Figure 18, the y_o/δ_{κ} ratios for the stack change from 0.563 to 1.252. This represents a y_o/δ_{κ} ratio increase of 122% from the lowest to the highest mean gas pressure.

From Figure 19, the $\xi_{kx} / \Delta x_{kx}$ ratios for the cold heat exchanger change from 0.684 to 2.224, from the lowest to the highest mean gas pressure. The peak displacement amplitudes are as much as two times the heat exchanger length. Also, the peak-to-peak displacement amplitude at the lowest pressure is as long as the heat exchanger length. For the hot heat exchanger, the $\xi_{kx} / \Delta x_{kx}$ ratios change from 0.473 to 1.541, from the lowest to the highest mean gas pressure. The peak displacement amplitudes are as much as one and one-half times the heat exchanger length. Again, the peak-to-peak displacement amplitude at the lowest pressure is as long as the heat exchanger length. For the stack, the $\xi_s / \Delta x_s$ ratios change from 0.186 to 0.604, from the lowest to the highest mean gas pressure. The peak displacement amplitudes, except for the lowest pressure, are as long as the stack length.

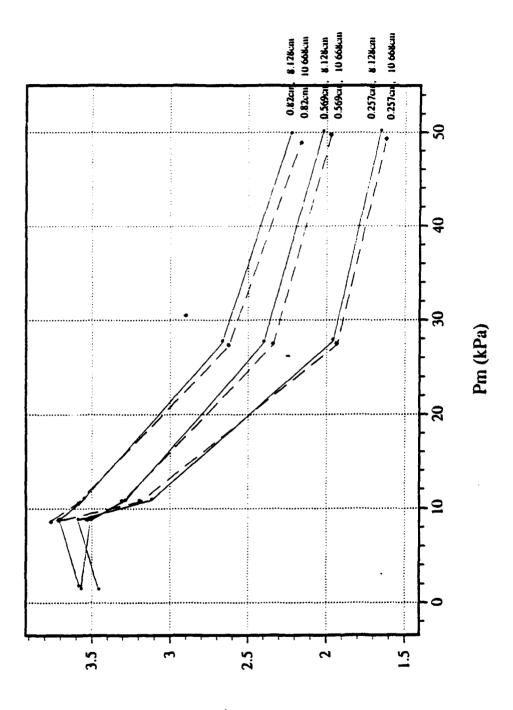


(%) WJ/OJ Figure 15 - Ratio of acoustic pressure amplitude to mean gas pressure versus mean gas pressure.

TABLE 1 - PRIME MOVER THERMAL LOAD FOR EACH HEAT EXCHANGER LENGTH AT DESIGNATED MEAN PRESSURES AND PLUNGER POSISTIONS.

Man Person	<u>Sa</u> M	1 6 40-		F	ľ						
										10.9 kPa	
			and the second								
								Management	Concentration of the local distribution of t	CONTRACTOR DEPENDING	Source and the second se
Cold He Length (cm)	(cm)	0.267	0.669	0.82	0.267	0.669	0.82		C 3 C 4		
Hot His Length (cm)	(cm)	0.263	0.560	0.62	0.263	0.644	0.02				
Mean presure (kPa)	(kPa)	1.428	1.42	1.763	8.794	0.02	A LA LA				
Acoustic P	resure (Ma)	•	0	0	1.082	0.298	0.74		1 707		
Frequency	(H2)	•	9	0	172.85	173.24	172 44				200 00 0
Prime Mover	Temperatures	(Centorado) 8								1.011	
Cold End	Collec	101.7	- 101 -								
Cold	Outside	101.7	101.	191.28	101.2						<u> 2.191-</u>
Cold He	3	101-	-108.4	-191.4							
Cold In	Center	-190.6	-106.2	190.9	10.001	A 7 A 1					
End	Stack	-106.7	-100.1	-150.0	141.	176.6	2 1 2 1				
Hot End	Stack	1.6	2.3	.1.	1.7		2 6				
T I	Center	20.4	20.1	21.3	30.0		34.0			0.9	
Ha H	M	21.1	20.1	21.78	34.7		2				
	Outside	21.3	20	21.48	35.7	1.95					
Hol End	Ithe	21.2	20.7	91.7%							
							33.7		19.3	9.55	36.1
Purse	Distance-	10.066 cm						CANCELORIE SHORE S			
			Non-				No. of Concession, Name	New York Concerning			
Cold He Lengel (cm	(cm)				0.257	0.520					
Hot Hu Length (cm	(cm)				0.243						28.9
Mean presure (NPa)	(InPa)					727			202.0	0.960	0.62
Acousto P	(Pa) (Pa)								10.01	10.77	10.01
Frequency						5	5		269.1	0.665	1.365
Prime Mover	Temestatures	(Certification)								166.99	199.4
Cold End	Collec				101-	7 101-					
	Outside				1.101-	6.101.			2.001		2.191.
	ł				-109.1	0.001-	101.				
	Center				-106.9	190	181-		5.001.	.177 6	
	Black				-170.7	-170.7	.104.6				
	Stack				3.6	9.6	12.1				
보	Center				C.16	34.3	35.4		2.45		
	Ţ				34.9	34.3			66	1.66	
Ŧ	Orelde				36.7	36			36	36	1.40
	3				36.4	34	34.63		36.36	34.9	9.40
											T T T

TABLE 1 - PRIME MOVER THERMAL LOAD FOR EACH HEAT EXCHANGER LENGTH AT DESIGNATED MEAN PRESSLIRES AND PLLINGER POSISTIONS. (CONTINUED)



JT / TT

Figure 16 - Ratio of center of hot heat exchanger temperature to center center cold heat exchanger temperature versus mean gas pressure.

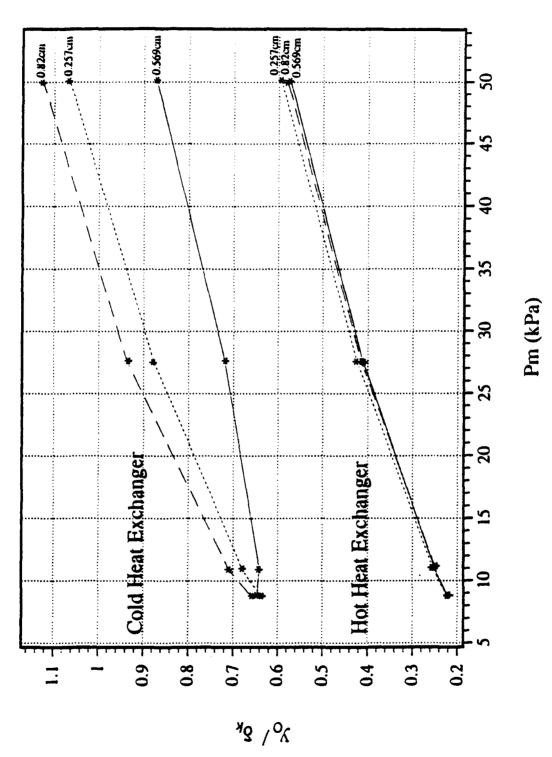


Figure 17- Ratio of heat exchanger plate half separation to thermal penetration depth versus mean gas pressure.

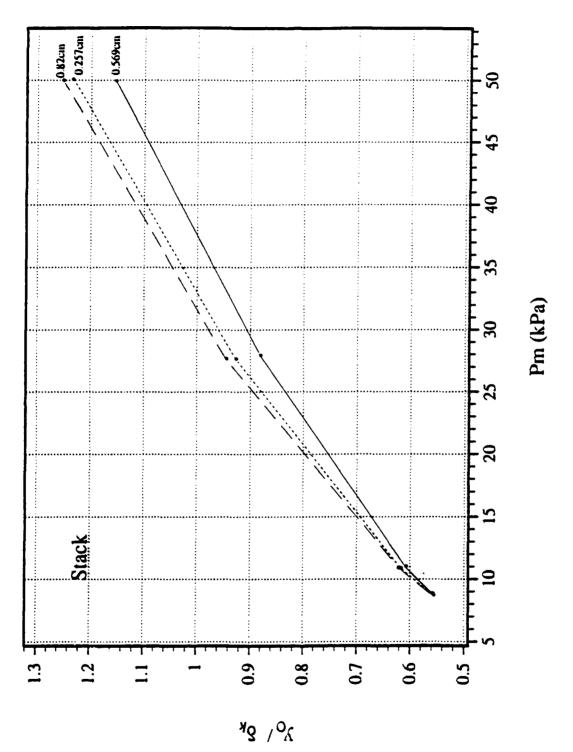
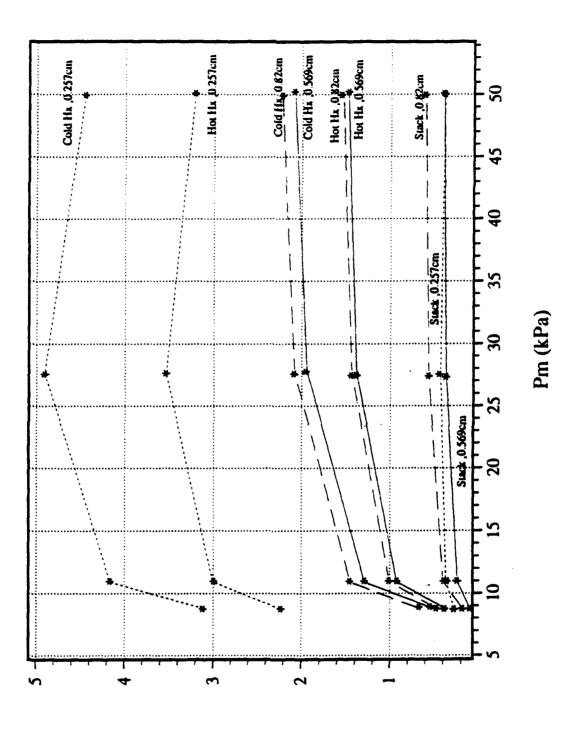


Figure 18 - Ratio of heat exchanger plate half separation to thermal penetration depth versus mean gas pressure.



x V / (x x) 3

Figure 19 - Ratio of peak displacement amplitude to stack or heat exchanger length versus mean gas pressure.

V. SUMMARY, CONCLUSIONS, AND RECOMMENDATIONS

A. SUMMARY

The goal of this thesis is to investigate the experimental heat exchanger performance in a thermoacoustic prime mover filled with neon gas. The experimental approach is to measure the pressure ratio, the steady state waveforms, and the temperatures for heat exchangers lengths of 0.257, 0.569, and 0.82 cm for a variety of mean gas pressures. A temperature gradient is established across the stack by keeping the cold heat exchanger submerged in liquid nitrogen and the hot heat exchanger near ambient temperature. The prime mover is designed to achieve very high acoustic displacement amplitudes with modest heat flows. Measurements are made at mean gas pressures of approximately 1.5, 8.8, 10.9, 27.6, and 50 kPa and plunger positions of 8.128 and 10.688 cm. The signal waveforms and spectra generated by the prime mover are also reported.

The overall results can be summarized as follows. Acoustic pressure amplitudes as high as 29% of mean gas pressure are generated by the prime mover. As the mean gas pressure increases, the pressure ratio increases for all heat exchanger lengths as the penetration depth approaches optimum value for this prime mover design. The shorter 0.257 *cm* long heat exchanger has higher acoustic pressure amplitudes at low mean gas pressures when compared to the longer heat exchangers.

There are three possible effects that may explain this behavior. First, the heat flows are small at low gas pressures, P_m , and so the limited thermal conductance of the shorter 0.257 cm heat exchanger is not a problem as it is at higher gas pressures and heat flows. Secondly, the acoustic losses for the 0.257 cm heat exchanger are the lowest of the three heat exchangers, which should increase the amplitude. Lastly, the

larger value of δ_{κ} at low gas pressures decreases the y_o/δ_{κ} , and thus the effective thermal diffusion distance from gas parcel to heat exchanger plate. This tends to corroborate Hofler's hypothesis that the short heat exchangers with $(\Delta x_{hx} / \xi_{hx}) < 1$, may be thermally effective if $(y_o / \delta_{\kappa}) < 1$.

In theory, heat exchangers are assumed to be isothermal. However, a pronounced radial temperature difference of as much a 9.2, 27.6, and $13.7 \degree C$ in the cold heat exchanger and 32.5, 11.9 and 24.9 $\degree C$ in the hot heat exchanger for lengths of 0.257, 0.569, and 0.82 cm, respectively, are generated at a mean gas pressure of 50 kPa. This difference results in radially varied thermal penetration depths.

The T_k / T_c ratio ideally is constant. However, experimentally, it is not. This ratio decreases by as much as 56%. This decrease is primarily due to the cold heat exchanger. Its temperature is not constant. On the other hand, the hot heat exchanger temperature is nearly constant.

The y_o/δ_{κ} ratios are similar in slope for the hot heat exchangers, for all the heat exchangers lengths, and increases approximately by 160% from the lowest to the highest mean gas pressure. However, the y_o/δ_{κ} ratio for different cold heat exchanger lengths increases between 35 to 88% and has different slopes with increasing mean gas pressure. Similarly, the y_o/δ_{κ} ratio for the stack increases by a much as 122%. The δ_{κ} is decreasing with increasing mean gas pressure which decreases the effective distance over which heat can diffuse transversely between a gas parcel and the heat exchanger plate.

The peak displacement amplitudes for the 0.257 *cm* long heat exchanger are as much as two times the heat exchanger length at the lowest mean gas pressure to as much as five times the heat exchanger length at the highest mean gas pressure.

Overall, the peak-to-peak displacement amplitudes are bigger than any of the heat exchanger lengths even at the lowest mean gas pressure. For the stack, the peak-to-peak displacement amplitudes are as long as the stack length, except at the lowest mean gas pressure. The displacement amplitudes increase with increasing mean gas pressure for all heat exchanger lengths.

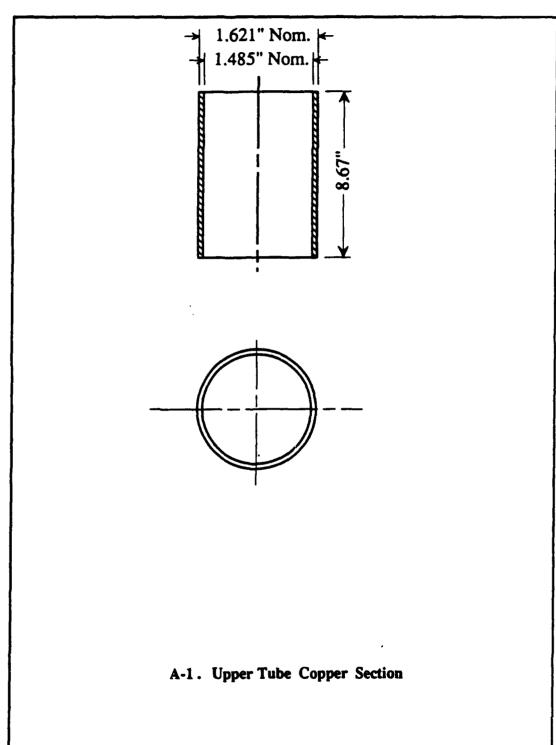
B. CONCLUSIONS

Several conclusions can be drawn from these results. The first conclusion is that the longest heat exchanger generated the largest pressure amplitudes. Secondly, the theoretical equations need to take into account the radial temperature gradient in the heat exchangers. The final conclusion is that the heat exchangers generate large pressure amplitudes even though the peak-to-peak displacement amplitudes are much larger than the heat exchanger lengths.

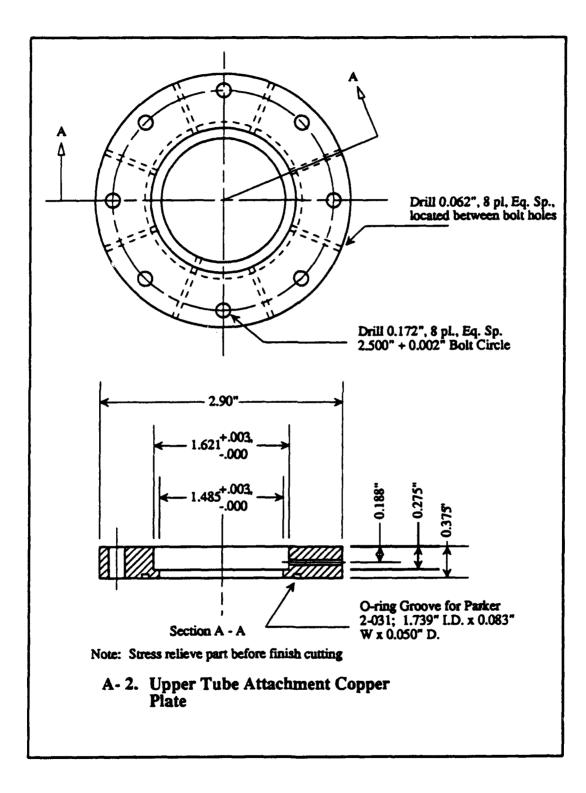
C. RECOMMENDATIONS

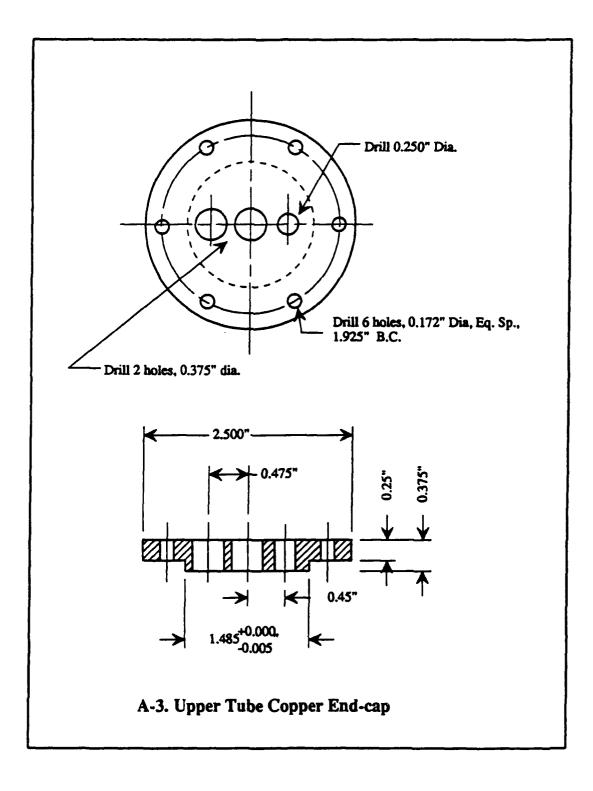
In order to better understand the thermoacoustic process, there is a need for more theoretical work to be performed to understand these experimental results. It is uncertain at this point what is the optimum heat exchanger length. However, several recommendations are proposed as follows:

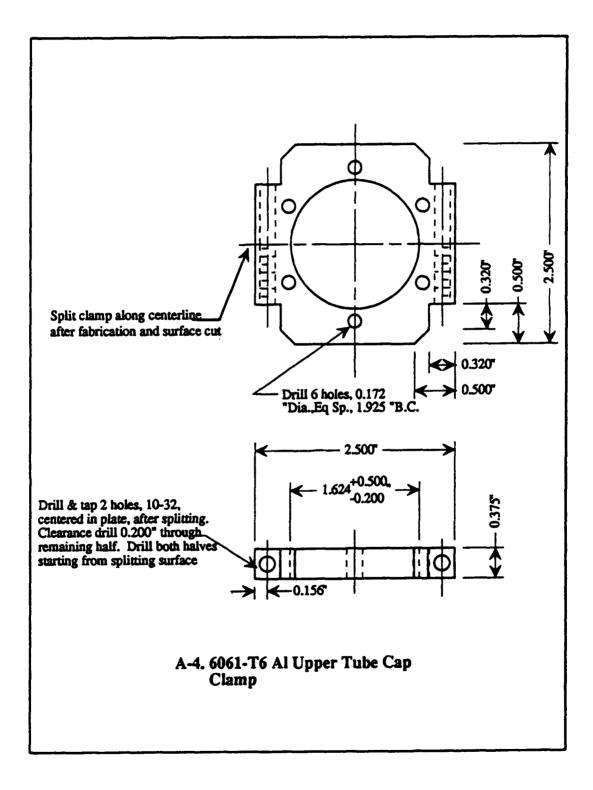
- Keep the T_c also constant and as low as possible.
- Measure heater powers carefully for all data points.
- Use cold heat exchangers of smaller y_o so that y_o/δ_{κ} is similar for both hot and cold heat exchangers.

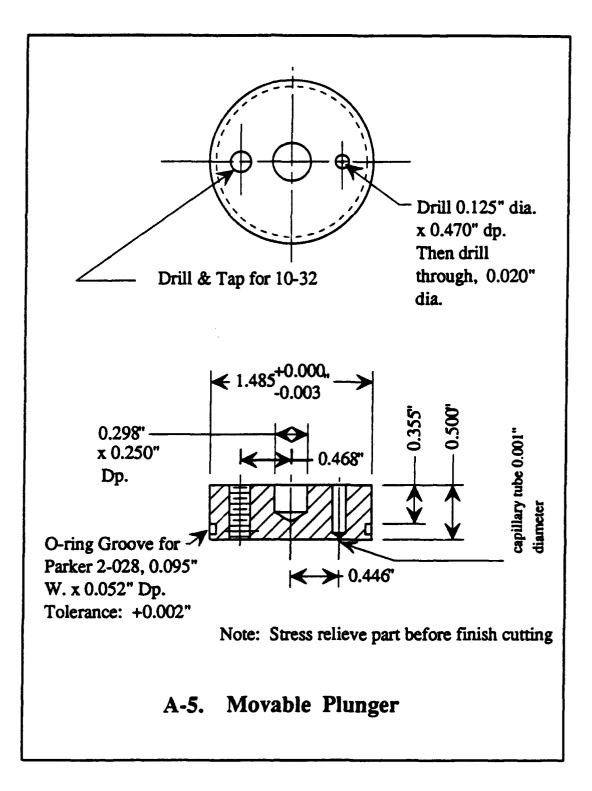


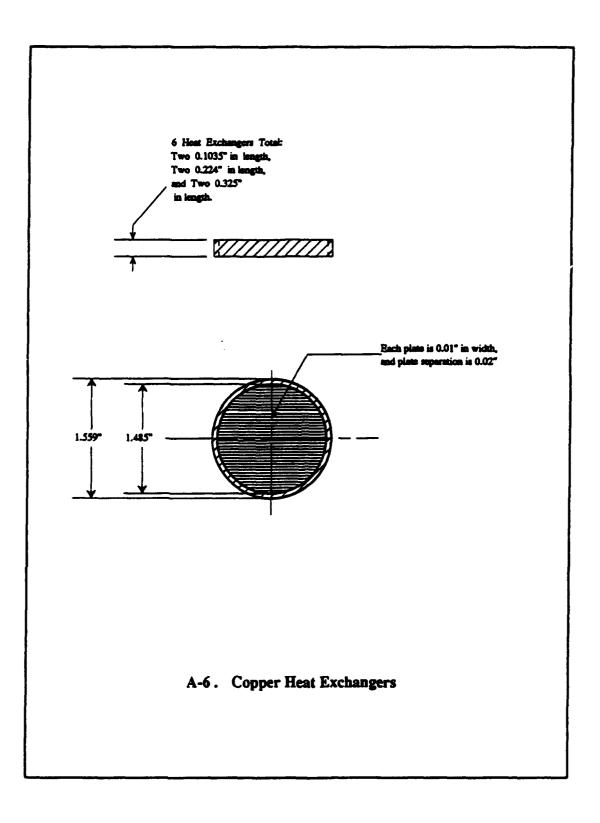
APPENDIX A - PRIME MOVER SPECIFICATIONS

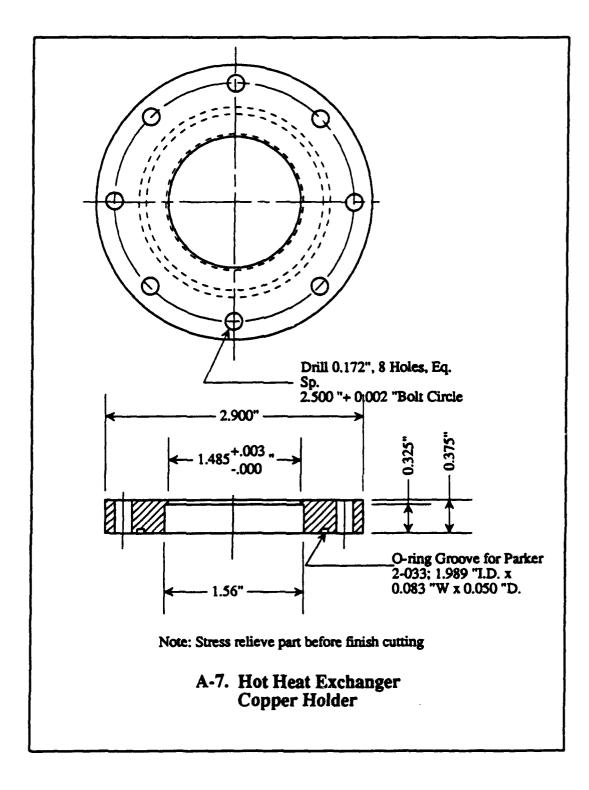


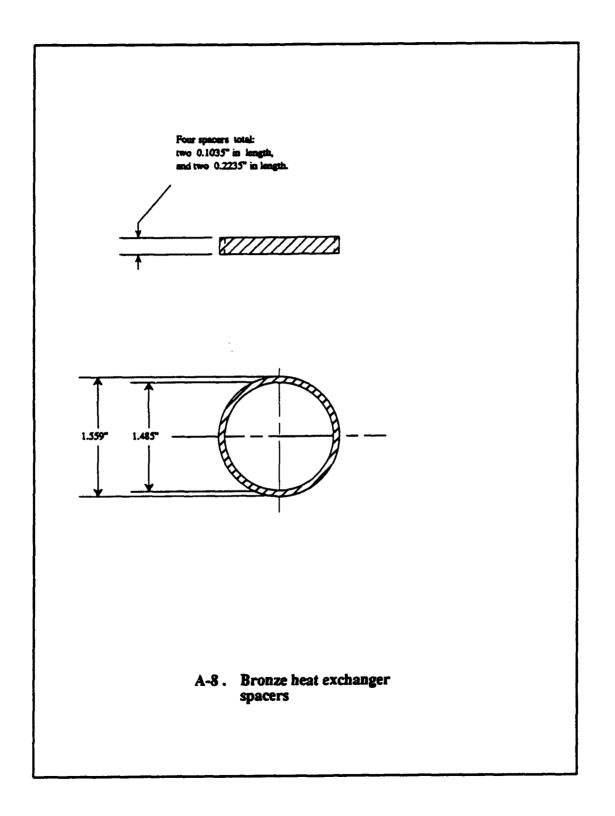


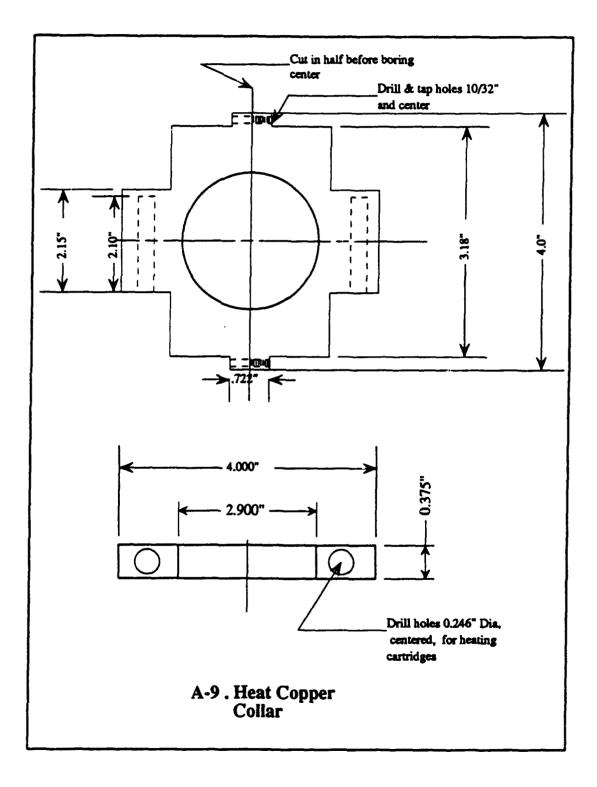


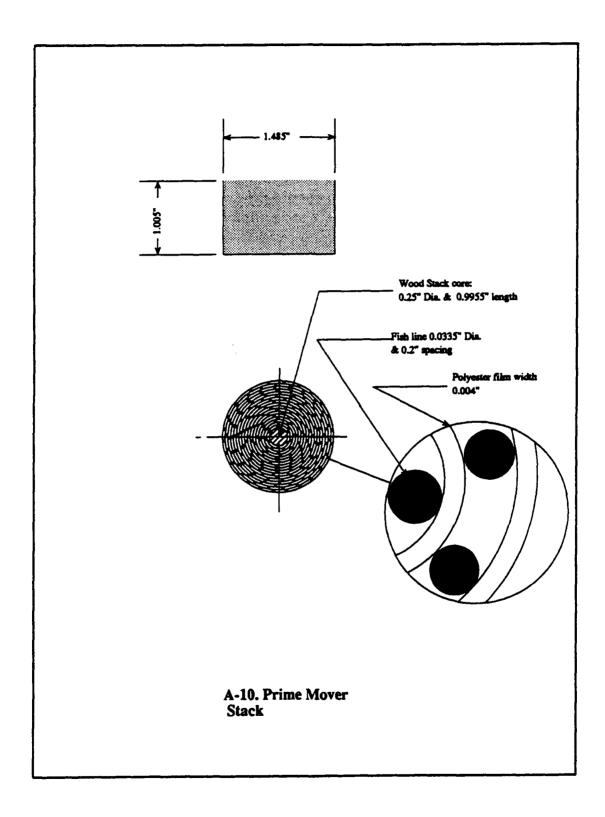


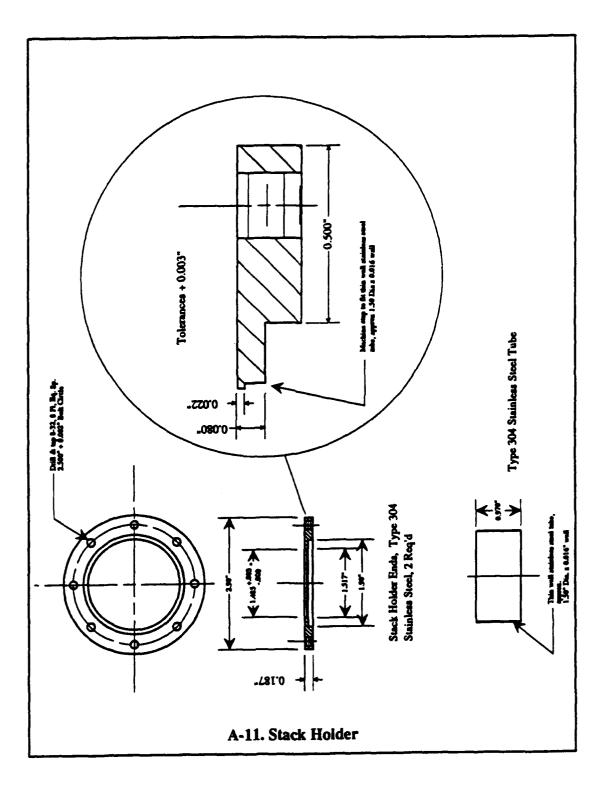


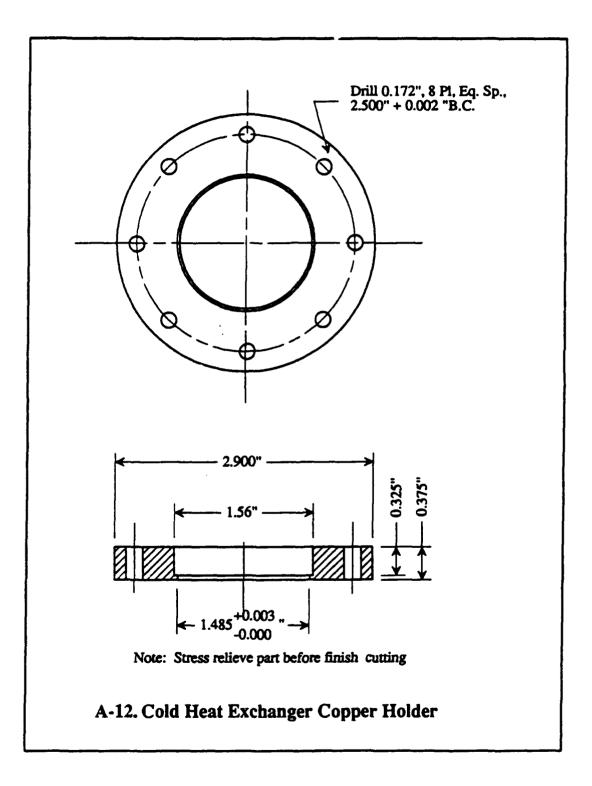


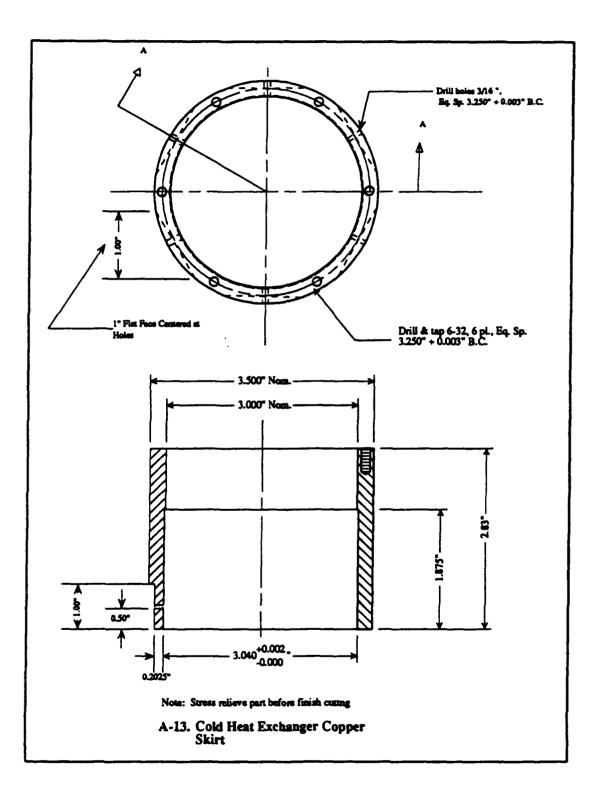


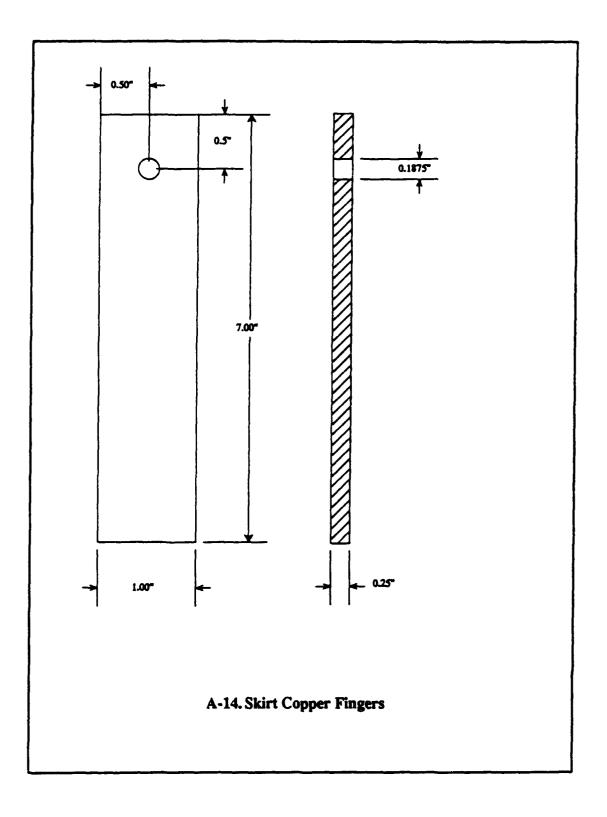


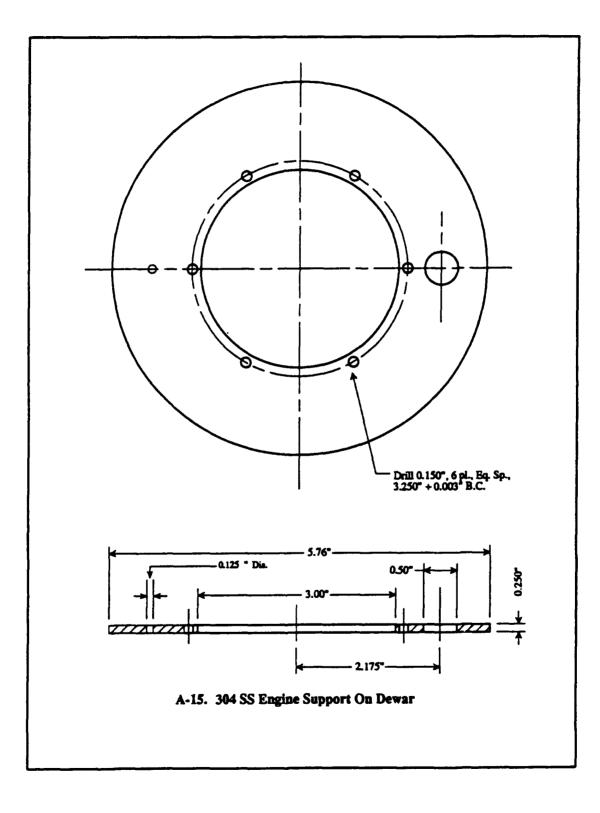


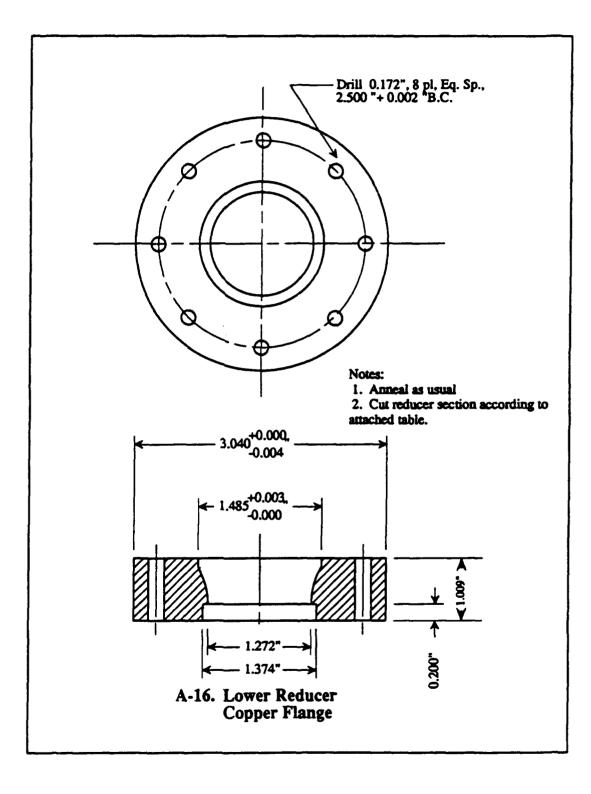


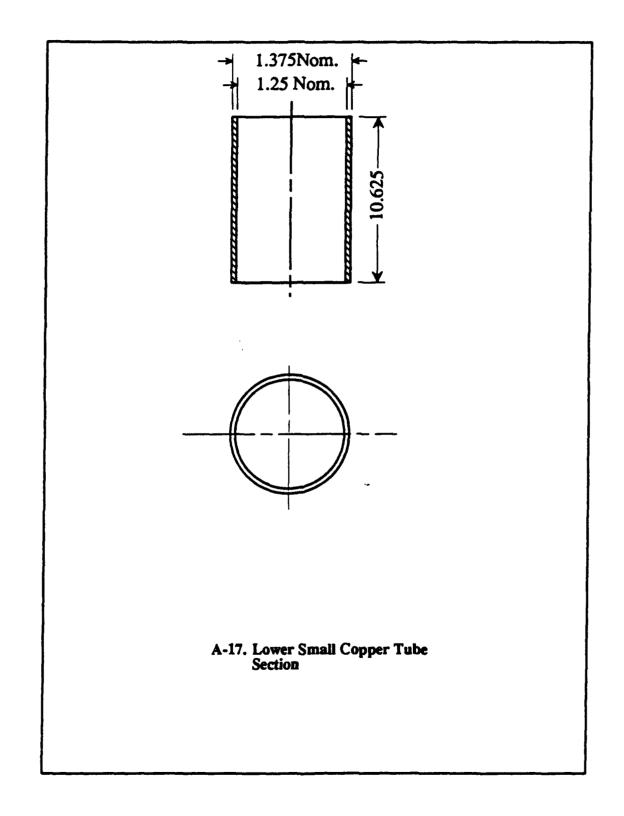


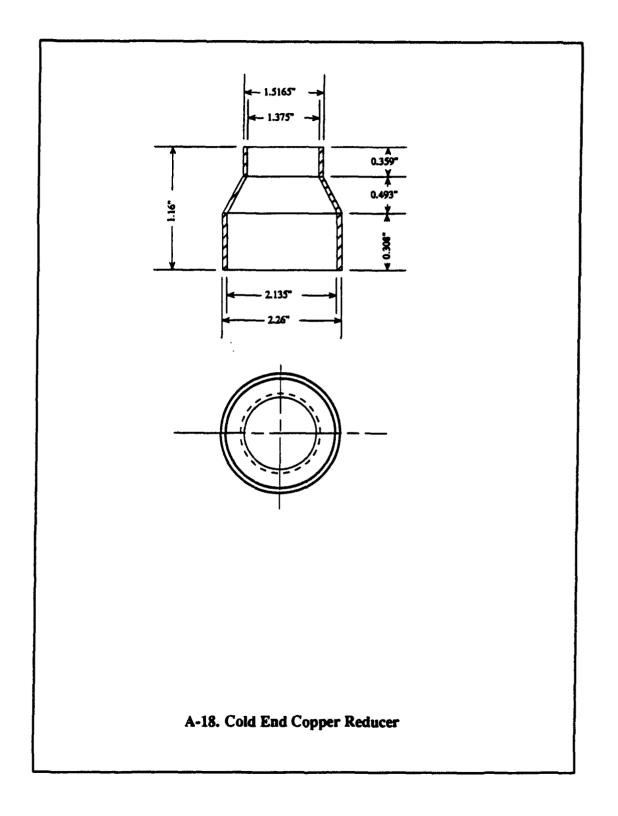


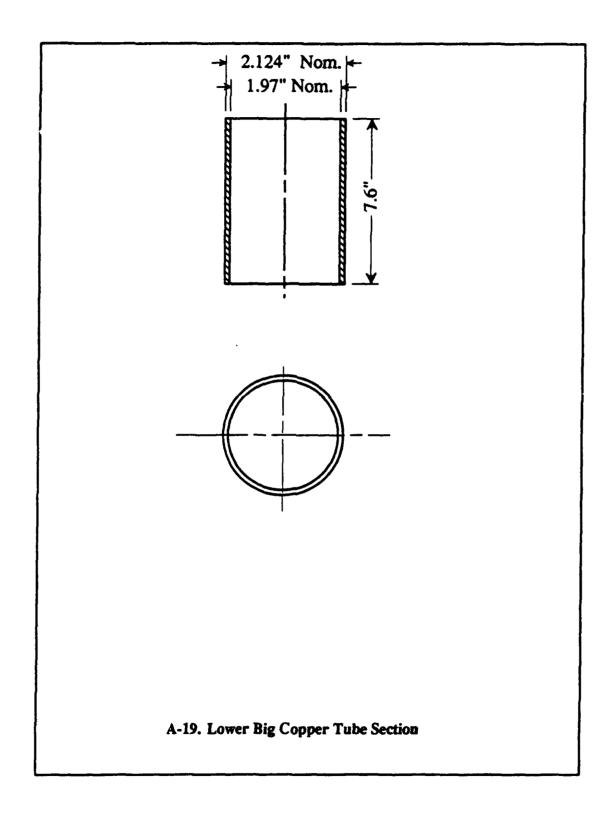


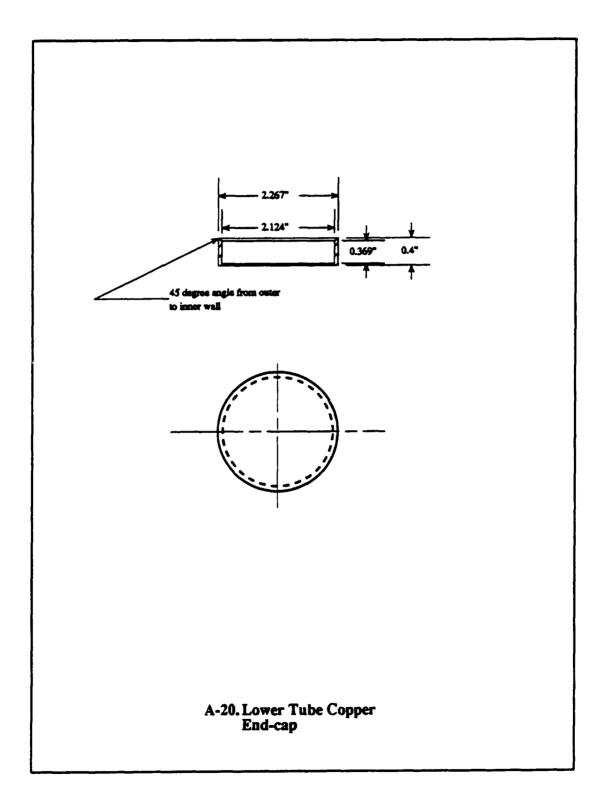


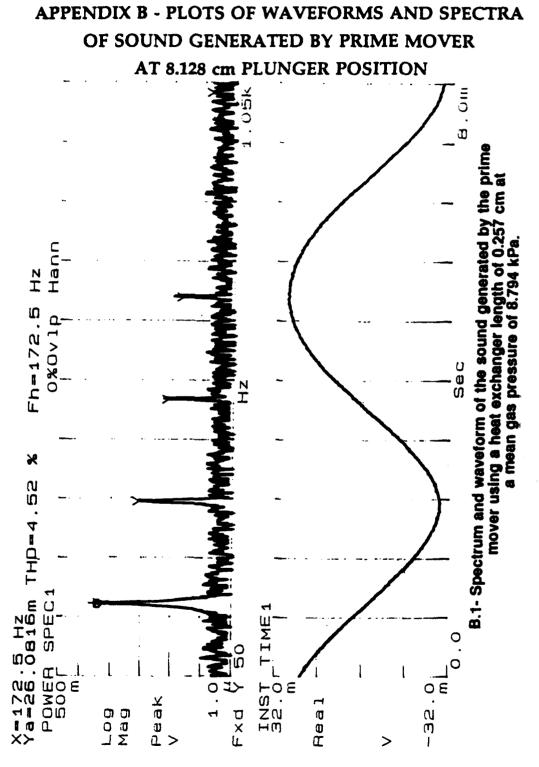




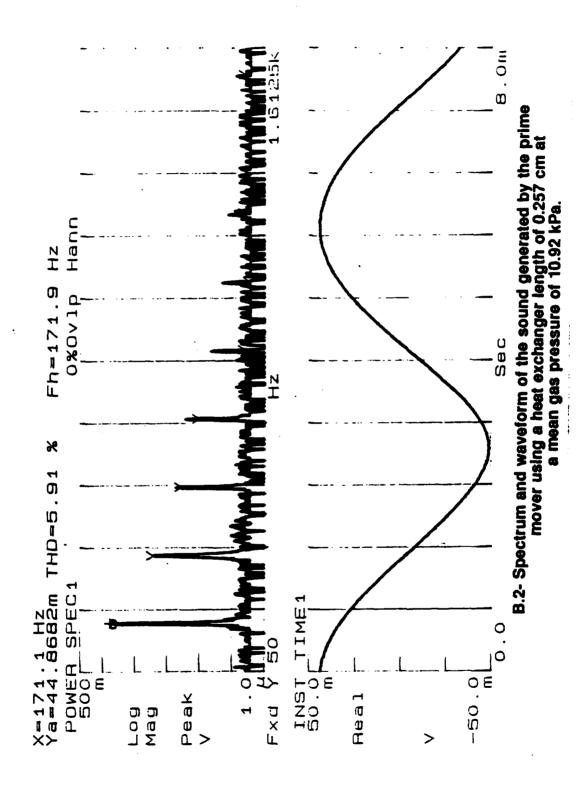


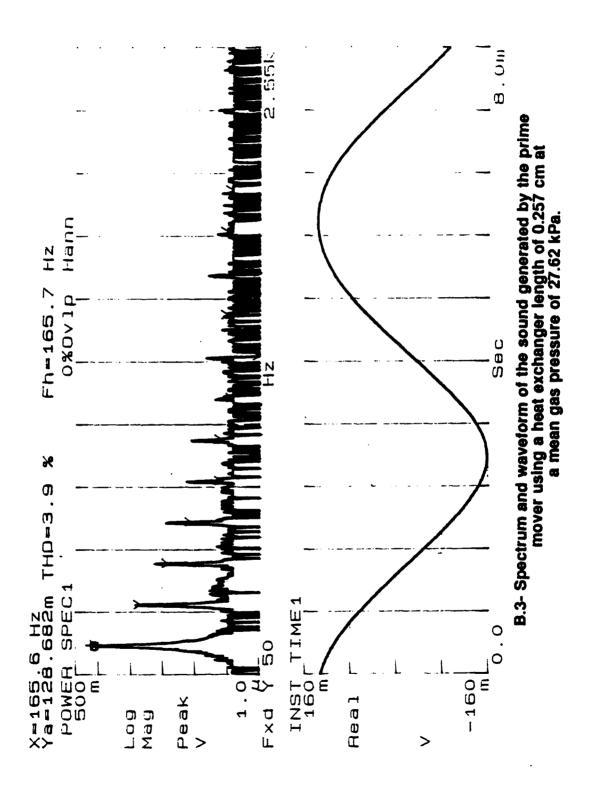


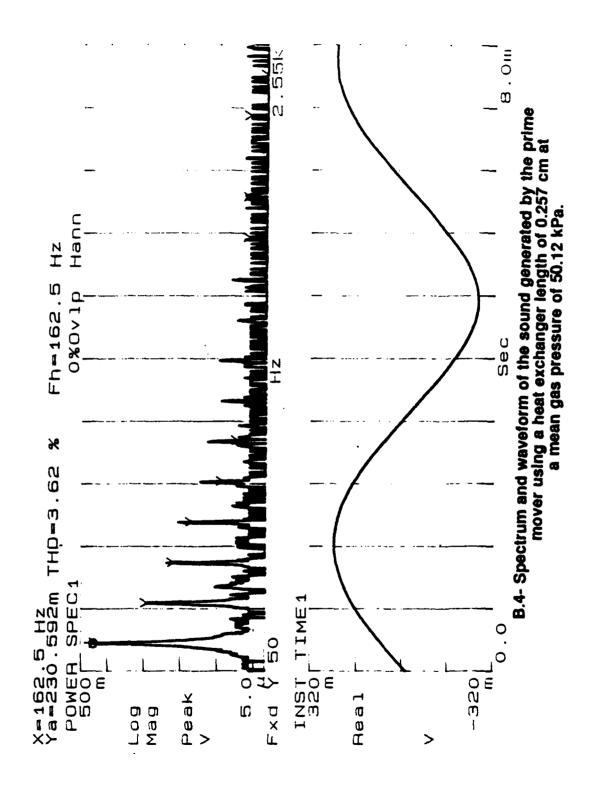


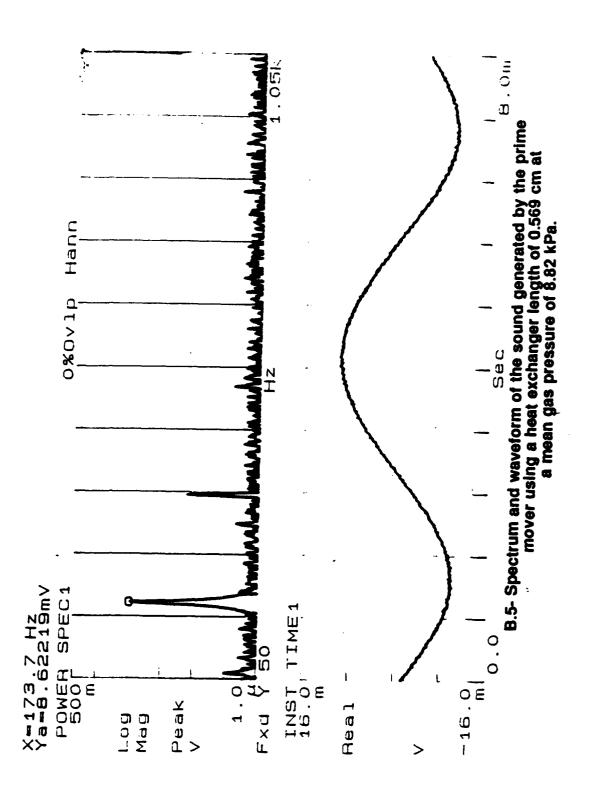


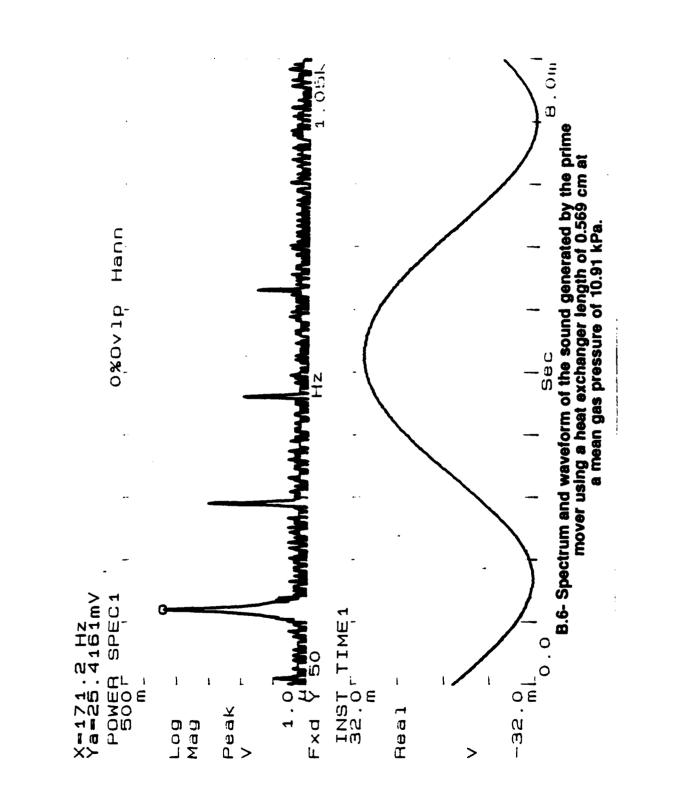
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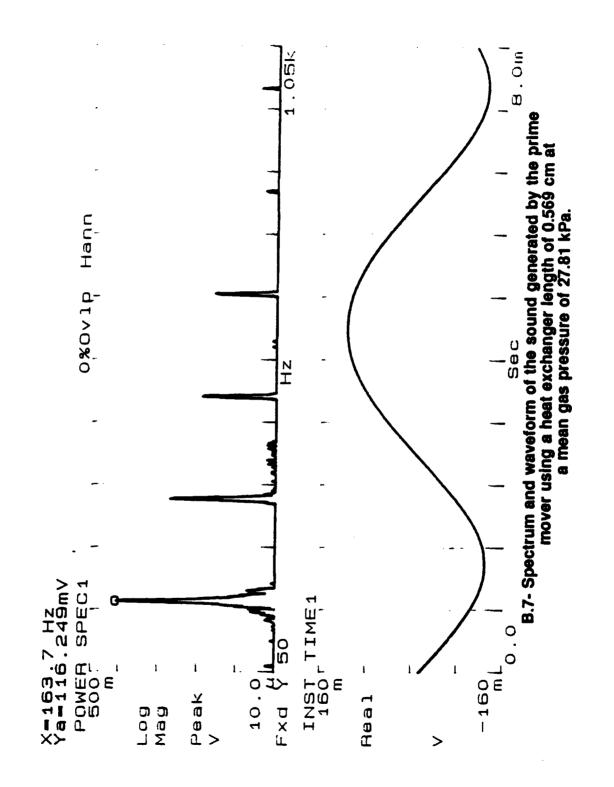


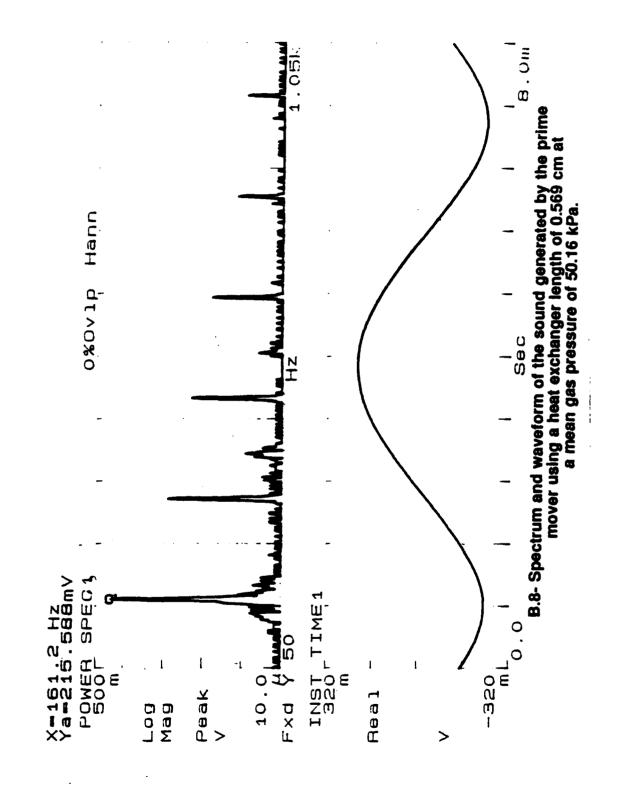


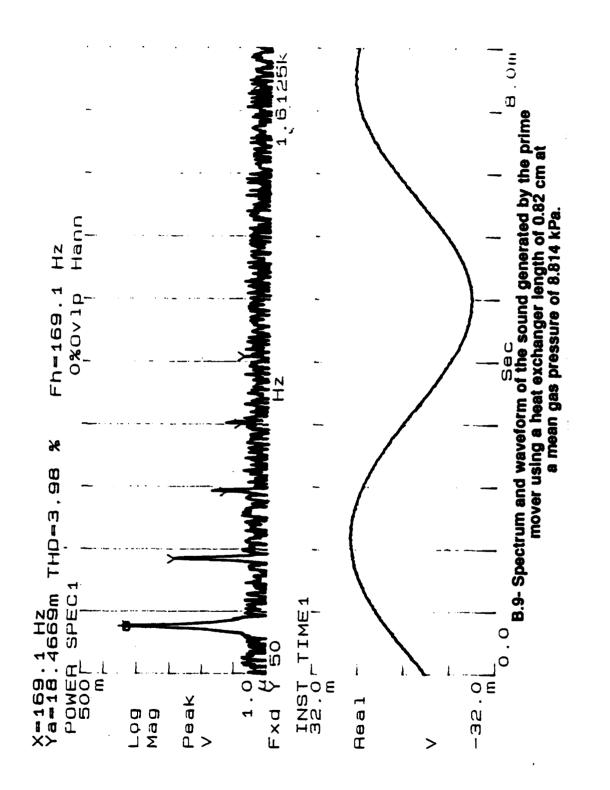


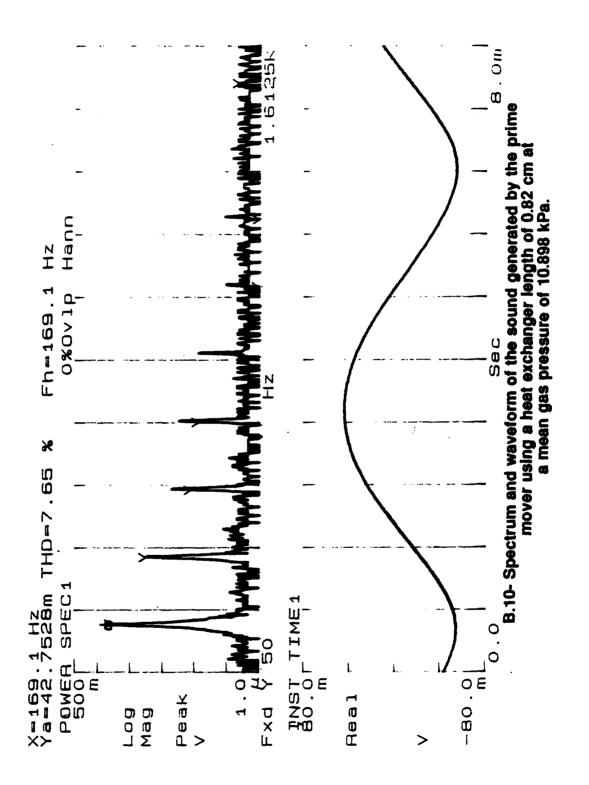


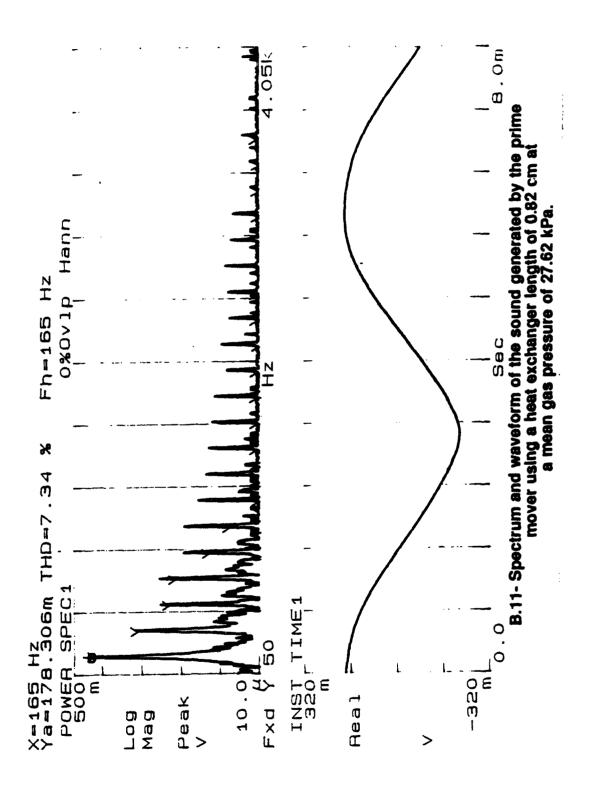


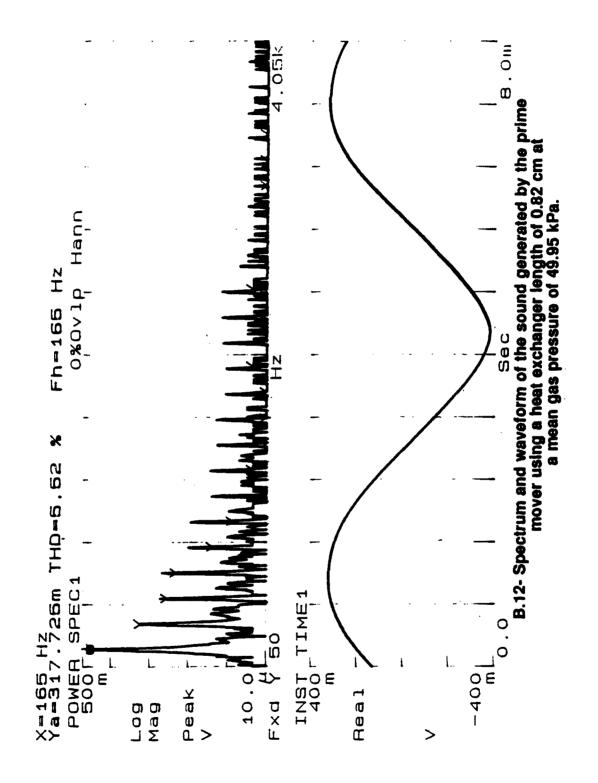












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