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# ERASABLE OPTICAL INFORMATION STORAGE IN THE POLYANILINES

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## ABSTRACT

The polyaniline family of polymers exhibits properties that allow for the storage of optical information. Photoexcitation of several forms of polyaniline results in long-lived changes in the absorption spectra, associated with trapped charged defect states. These new absorptions exhibit long lifetimes at low temperatures (below  $\sim 250$  K), with erasure occurring upon warming the sample. Use in erasable optical information storage devices and in related technologies is discussed.

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Efforts in recent years to develop erasable optical storage devices have focused on various optical phenomena in a number of materials. The most successful efforts have made use of magneto-optical effects in rare earth-transition metal alloys and amorphous-to-crystalline phase transitions in chalcogenide glasses [1,2]. Numerous other approaches have been proposed, including photochromic effects in organic materials such as TCNQ salts [3] and frequency domain optical storage based on persistent spectral hole burning in a variety of organic and inorganic materials [4,5]. Recently the polyaniline family of conducting polymers has been shown to exhibit a number of interesting optical and photoinduced optical properties [6-8] that indicate that this family of materials holds promise for use as an erasable optical storage medium.

Polyaniline is a large bandgap ( $E_g \sim 3.8$  eV) polymer that consists of alternating  $C_6H_4$  (phenyl) rings and amine or imine nitrogens, the exact composition depending on the oxidation state. Oxidation of the fully reduced leucoemeraldine base toward the half-oxidized emeraldine base and the fully oxidized pernigraniline base leads to the appearance of an absorption band near 2 eV, which is associated with charge transfer from benzenoid to quinoid groups [9-11]. Photoexcitation of polyaniline via pumping into the 2-eV absorption band or into the 3.8-eV bandgap transition results in the formation of long-lived charged defect states, trapped states, and excitons [8]. Various derivatives of polyaniline, including poly(*o*-toluidine) and poly(2-ethoxyaniline), show similar results [12].

The dominant signature of the long-lived photoinduced defects in polyaniline is an increased absorption that peaks near 1.4-1.5 eV [7,8,12]. The origin of this phenomenon has been proposed [8,13] to result from the sensitivity of the electronic states of polyaniline to rotations of the phenyl rings in and out of the plane formed by the nitrogens in the backbone. The new absorption near 1.4 eV results from the presence of a polaron, described in terms of ring-angle distortions, near a planar quinoid moiety [8,9]. Sufficient  $\pi$ -electron delocalization occurs so as to render this configuration as a bound state [8]. Relaxation to the ground state is hindered because of the positive binding energy of this configuration,

but may also be retarded by interchain packing due to a lack of free volume in the material. Previous experiments have shown that the magnitude of the 1.4-eV photoinduced absorption peak in emeraldine base remains nearly constant for several hours when the sample is held at 80 K [7,8], with shorter lifetimes as the temperature is raised; complete erasure occurs upon heating the sample to room temperature. In addition, the magnitude of this peak varies as the square root of the pump intensity, i.e.,  $I^{\frac{1}{2}}$ .

On the basis of these results we have investigated polyaniline for use as an erasable optical storage medium [14]. Samples of numerous forms of polyaniline and its derivatives have been prepared as thin films or as powders dispersed in KBr pellets. Samples are mounted in a cryostat and cooled in the dark to low temperatures. The fractional change in transmission of each sample,  $-\Delta T/T$ , near 1.4–1.5 eV is determined using a grating monochromator or a FTIR spectrometer to measure the transmission before and after photoexcitation by an argon-ion laser or a dye laser.

Table I lists several of the sample configurations studied for photoexcitation using the 488-nm or 514-nm lines of an argon-ion laser at power densities of  $\sim 400$  mW/cm<sup>2</sup>. Also included are the lifetimes of  $-\Delta T/T$  determined for exponential decay at different temperatures. Magnitudes of  $-\Delta T/T$  range from a few percent to greater than 30%, depending on the particular sample studied, the concentration of the sample, and the temperature of the sample. Most notable are those materials with lifetimes of at least several days including pernigraniline base at 77 K, emeraldine poly(*o*-toluidine) at 77 K, and, in particular, pernigraniline poly(*o*-toluidine) at 250 K. The photoinduced absorption of each of these samples showed some initial decay within 1–2 h after photoexcitation, with very little further decay during the measurement time of 18–24 h. Experiments involving longer measurement times are in progress. The results at 250 K for pernigraniline poly(*o*-toluidine) indicate that through molecular design a material may be synthesized that possesses long-lived effects at room temperature, which is necessary for practical devices.

As an example of the optical effects observed in these materials, Fig. 1 shows a sample of pernigraniline base in KBr mounted at the end of a liquid nitrogen cold finger. The

sample was exposed to a focused laser beam (488-nm, 100 mW,  $\sim 600 \mu\text{m}$  dia. beam size) whose position was changed to form the pattern "OSU". The photograph was taken using a 35mm camera with infrared film; the illuminating source was the output of a tungsten lamp passed through a long-wavelength-pass filter ( $\lambda > 750 \text{ nm}$ ) in order to screen out visible light. The increased absorption of the sample is clearly seen as the written pattern "OSU". Changes in intensity of the pattern and the background are due to differences in exposure time of the write beam and due to nonuniformity of the material in the KBr pellet. This written pattern continued to be observed as long as the sample was maintained at low temperature (77 K) and disappeared upon warming the sample to room temperature.

Figure 2 shows a contour plot of the transmission of an emeraldine base film of thickness  $\lesssim 0.5 \mu\text{m}$  on a quartz substrate held at 77 K after being exposed to a pump laser beam at 580 nm of size  $\sim 70 \mu\text{m}$  dia. The transmission was measured using a laser beam at 827 nm (1.5 eV) of size  $\sim 30 \mu\text{m}$  dia. The sample was scanned over a  $200 \mu\text{m}$  square. Small irregularities in the contour reflect variations in sample uniformity in the area probed. Although the maximum photoinduced effect is concentrated within the combined write beam and read beam diameters, the region showing a nonzero photoinduced effect is somewhat larger. This larger size may be due to diffusion of the photoexcited species, slightly unfocused optics, competition with photoinduced bleaching (because the probe energy is very close to the crossover energy between photoinduced absorption and photoinduced bleaching [8]), and competition with heating effects [15]. Work is currently in progress to determine minimum detectable spot sizes in order to obtain maximum bit storage density.

In addition to studies of lifetimes and bit sizes, other parameters that are important for practical optical information storage have been investigated. Our experiments have shown that there is no degradation with thermal cycling between room temperature and the desired low temperature. Preliminary experiments also indicate that 1.4-eV photoinduced absorption in emeraldine base occurs within a few picoseconds following photoexcitation, so that writing times may be fast. It is estimated that in order to write a  $1 \mu\text{m}^2$  bit where

$-\Delta T/T \sim 0.1$ , an energy of  $\lesssim 10$  nJ is needed, assuming a square root dependence on write beam energy.

The use of polyaniline or its derivatives for optical information storage would involve the writing of information on a thin film, composite, or blend of polyaniline with a visible or uv laser beam, and subsequently, reading the information with a diode laser (and an appropriate detector) tuned near the peak in the photoinduced absorption spectrum (1.4–1.5 eV). In principle, a sufficiently intense laser operating near 1.4–1.5 eV can be used to heat the exposed portions of the sample to room temperature or higher for erasure of bit information, or the entire sample can be warmed for bulk erasure. The laser-intensity dependence of the photoinduced absorption spectrum may allow for analog data storage at any bit. Other possible uses based on the photoinduced absorption phenomena include photorefractive devices and photoinduced gratings.

In summary, we have found that the polyaniline family of polymers exhibits photoinduced optical absorptions associated with long-lived defect states that make them candidates as media for erasable optical information storage. A variety of forms of polyaniline have been determined to store information for long times with erasure occurring upon heating the sample. Other studies in progress include a survey of additional materials that exhibit similar long-lived optical effects, determination of threshold powers for erasure, maximum bit density, the maximum number of write-erase cycles, time dynamics for writing and erasing, and signal-to-noise ratios for fast reading times. The availability of a broad range of polyanilines, derivatized polyanilines, and similar ring-containing polymers provides a large number of materials options for optical information storage within desired temperature ranges.

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## FIGURE CAPTIONS

Fig. 1: Optical pattern written with the 488-nm argon-ion laser line on a composite pellet of pernigraniline base in KBr mounted at the end of a liquid nitrogen cold finger. The circular hole in the sample holder is 11.5 mm in diameter. The double-ring halo is infrared light reflected from the cryostat.

Fig. 2: Contour plot of the transmission of an emeraldine base film on a quartz substrate held at 77 K. The circular region corresponds to increased absorption in the area where the film was exposed to a  $\sim 70\text{-}\mu\text{m}$  size laser beam at  $\lambda = 580\text{ nm}$  and probed with a  $\sim 30\text{-}\mu\text{m}$  size laser beam at  $\lambda = 827\text{ nm}$ .

Table I. Materials exhibiting long-lived photoinduced effects.

Material	Temperature (K)	Lifetime
Emeraldine base film	77	110 min
Emeraldine base in KBr	77	24 hr
Pernigraniline base in KBr	77	≥24 hr
Pernigraniline base in KBr	200	>24 hr
Emeraldine poly( <i>o</i> -toluidine) in KBr	77	≥24 hr
Emeraldine poly( <i>o</i> -toluidine) in KBr	250	~2.5 hr
Pernigraniline poly( <i>o</i> -toluidine) in KBr	77	≥24 hr
Pernigraniline poly( <i>o</i> -toluidine) in KBr	250	≥24 hr
Pernigraniline poly(2-ethoxyaniline) in KBr	77	140 min



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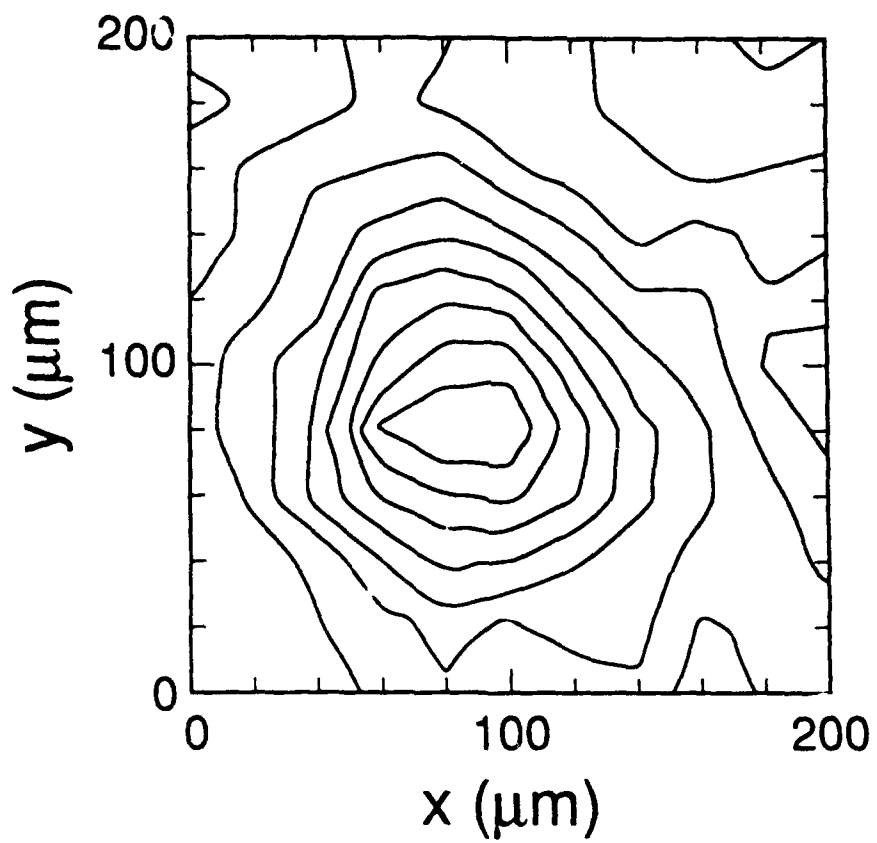


Fig. 2 R.P. McCall et al.