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An Investigation of Si-SiO₂ Interface Charges in Thermally Oxidized (100), (110), (111), (511) Silicon

by

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An Investigation of Si-SiO₂ Interface Charges in Thermally Oxidized (100), (110), (111), and (511) Silicon

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ABSTRACT

Trends in the electronic properties of the Si-SiO₂ interface with various processing have been frequently reported. The present study focuses on silicon substrate orientation dependent trends in fixed oxide charge, Q_f , and interface trap charge, D_{it} , for four silicon orientations: (100), (110), (111), and (511), for oxidation temperatures in the 750-1100°C range, with and without hydrogen-containing post-metal anneals. and for processing within and without a cleanroom. It is found that the presence of mobile ionic charge in noncleanroom processing and the lack of post-metal annealing can either obscure or enhance some trends. Both Q_f and D_{it} increase for decreasing oxidation temperature for all silicon orientations. The orientational ordering of the charges varies with oxidation temperature and is dominated by the silicon atom areal density at the lowest temperatures with (110) Si having the highest charge, but a change to the (111) orientation is observed at higher oxidation temperatures. This orientational charge ordering parallels the orientational oxidation rate ordering but not the intrinsic stress. A model is proposed that considers the orientationally dominated oxidation rate, viscous relaxation, and strain accommodation across the interface as crucial processes.

INTRODUCTION

The control of charges at the Si-SiO₂ interface is of crucial importance for the fabrication of small-geometry devices. Thus, the relationship between the chemical and physical nature of the interfacial region and the dependence on various processing parameters has been the focus of much investigation. Early studies reported that the silicon surface could be electronically stabilized by a thermally grown oxide.¹ More recently, it was suggested that successful device fabrication was dependent upon the control of four types of charges that occur in thermally oxidized silicon, namely the interface trap charge (Q_{it}) , the fixed oxide charge (Q_f) , the mobile ionic charge (Q_m) , and the oxide trapped charge (Q_{ot}) .² In particular, most studies have focused on the two types of electrical charge that are closely related to the interfacial region, namely the interface trap charge and the oxide fixed charge. The interface trap charge is the charge on interface traps, also called surface states.

Both the fixed oxide charge and interface trap charge display strong dependencies on oxidation processing variables. In particular, Q_f depends on whether the oxidation ambient is dry oxygen or steam,³⁻⁵ the oxidation temperature with higher temperatures resulting in lower charge densities,³⁻⁷ and on the orientation of the silicon substrate with the (111) orientation having a higher Q_f value than (100).^{3,5-8} In addition, Q_f can be reduced by high temperature annealing in such inert ambients as nitrogen, argon, or helium.³⁻¹¹

There exist some discrepancies in the literature regarding the dependencies of the interface trap charge Q_{it} . Usually reported as the interface trap density distribution D_{it} (cm⁻²/eV), it displays the same orientation dependence as does Q_f . However, conflicting data exists on the effect of oxidation temperature^{4-6,13-15} and high-temperature annealing^{4-6,10,13-16} on Q_{it} . Though high-temperature annealing of interface traps is less well characterized, substantial data exists that shows a minimization of Q_{it} by low-temperature annealing in a hydrogen-containing ambient.^{5,15-17}.

Since fixed oxide charge and interface trap charge often display the same process dependencies, it is tempting to speculate that these charges have a similar origin. Recent oxidation studies have shown a possible relationship between the oxidation temperature and many oxide properties as is illustrated schematically in Fig. 1.¹⁸ It is seen that Q_f and D_{it} as well as the SiO₂ film properties of intrinsic stress, density and refractive index all increase with decreasing oxidation temperature and all anneal to lower values. Though speculation is tempting, the exact chemical and physical origin of these is still unknown. It should be noted that Q_f and D_{it} are usually measured after one or several processing steps, which must include oxidation or deposition of a dielectric. Hence, the measured D_{it} is a sum of initial number of states or charges associated with the silicon surface plus any changes due to processing. Therefore, if we assume in a first-order model that the initial number of states for a bare silicon surface, $D_{it}(I)$, is proportional to the silicon surface atom density, hence onstant, any changes in the measured D_{it} are wholly contained in the processing term, $D_{it}(P)$, which may be positive or negative. The measured D_{it} can be expressed as

$$D_{it} = D_{it}(I) + D_{it}(P) , \qquad (1)$$

and similarly for Q_{f} . The processing terms for thermally grown SiO₂ films contain the influence of temperature, ambient, rate, impurities and substrate orientation.

The purpose of this study is to attempt to separate the processing effects of temperature, oxidation rate, substrate orientation, and cleanliness. The implications of film stress, post-metal annealing, and sodium contamination on Q_f and D_{it} are directly addressed. The observed trends in Q_f and D_{it} will be compared and discussed in the light of current models and recent observations. A complex picture emerges, but important parameters are identified. Much of the data reported herein is new as four different orientations of silicon were processed simultaneously, but there is some overlap with previous studies with appropriate comparisons. Since controversies exist in this area, we believe that overlaping systematic studies and comparisons with similar studies are useful. This is particularly important, since we show the effects of processing cleanliness on the trends in the electrical data.

EXPERIMENTAL PROCEDURES

Two independent sets of metal-oxide-semiconductor (MOS) structures were prepared on single-crystal silicon substrates under similar processing conditions, in order to determine the effect of mobile-ion contamination, most likely Na ions, on the observed trends in Q_f and Q_{it} . The first set of samples was prepared in the non-cleanroom facility of the University of North Carolina. The second set was prepared in the cleanroom facilities of Duke University and the Microelectronics Center of North Carolina (MCNC). The results from these two sets of samples will be referred to hereafter as the cleanroom and noncleanroom results.

Processing steps for the non-clearoom samples:

Lightly boron-doped p-type silicon (100), (110), and (111) wafers with resistivities in the 2-5 Ω -cm range were used in the non-clearoom study. The samples were cleaned using a modification of the RCA method,¹⁹, which consists of mixtures of NH₄OH:H₂O₂:H₂O and HCl:H₂O₂:H₂O with both in the respective portions of 1:1:5, followed by a 10-sec concentrated HF dip with each step separated by thorough deionized (DI) water rinses and finally a clean-N₂ blow dry.

The wafers were loaded into a double-walled fused-silica tube of the oxidation furnace in an inert N_2 ambient, oxidized for a specified time, and pulled fast into the endcap in the oxidizing ambient with a pull time shorter than 5 sec. Some wafers received a postoxidation anneal (POA) in N_2 at the oxidation temperature. Dry oxygen of 99.99% purity, containing less than 6 ppm of water as measured at the exit of the oxidation tube, was used as the oxidizing ambient.

The wafers were oxidized at five temperatures in the 750°C-1100°C range, with the resulting oxide thicknesses varying from 25 nm on (100) wafers oxidized at 750°C to 90 nm on (110) wafers oxidized at 1100°C. Oxide thicknesses were measured with a high-precision manual ellipsometer.

Aluminum dots 0.6 μ m-thick were evaporated on the oxidized wafers in vacuum using a resistively heated source. Some of the wafers then received a 30 min post-metallization anneal (PMA) at 400°C in forming gas (FG), consisting of 10% H₂ and 90% N₂ at 400°C. Samples with and without PMA will be hereafter referred to as PMA and NPMA samples. respectively.

Processing steps for the cleanroom samples:

Oxidations of p-type (100), (110), (111), and (511) silicon wafers with resistivities in the 2-20 Ω -cm range were performed in the cleanroom facilities at Duke University. All wafers were cleaned using the modified RCA cleaning sequences prior to oxidation.¹⁹ The wafers were oxidized at the temperatures and for the times indicated in Table I. All samples were pulled fast into the endcap in the oxidizing ambient. Upon cooling, the wafers were stored in a dry box continuously purged with pure nitrogen until all oxidations were completed.

A blanket layer of aluminum was evaporated on the oxidized wafers using a resistively heated evaporator located in the class-1 cleanroom facility at MCNC. The (100), (111), and (511) 4-inch diameter wafers from all oxidations were metallized together, and the (110) 3-inch diameter wafers were metallized in a separate run. Capacitors with guard-ring sturctures were defined using lithography, and a second blanket layer of Al was evaporated on the backside of the wafers after the backside oxide had been etched for backside contact.

Upon completion of the metallization steps, the wafers were cleaved in half, with one half of each wafer receiving a 30 min PMA at 400°C in FG, and the other half receiving no anneal (NPMA).

DATA ANALYSIS

The oxide fixed charge density is calculated from the relationship²²

$$Q_{\rm f} = C_{\rm ox} \left(\phi_{\rm ms} - V_{\rm FB} \right), \qquad (2)$$

where C_{0x} is the oxide capacitance per unit area, V_{FB} the measured flatband voltage, and ϕ_{ms} the metal-semiconductor work-function difference. The oxide fixed charge density Q_f

is usually obtained from high-frequency (1 MHz) capacitance-voltage made on MOS structures that received a post-metallization anneal. This ensures minimum flatband-voltage variations due to interface traps.⁵ In this study, we will also report Q_f measurements from NPMA samples. Measurements of Q_f on NPMA samples also include the portion of Q_{it} that responds to the high-frequency measurement. For high concentrations of interface traps, this component is expected to be significant. While this charge reporting procedure could be misleading, we include the results in order to explore the contention that incomplete or ineffective PMA can confuse results and obscure or create artificial trends.

Interface trap density D_{it} was obtained combining the high-frequency capacitancevoltage measurements with quasi-static capacitance-voltage measurements using a voltage ramp rate of less than 100 mV/sec. The interface trap density is given by

$$D_{it} = \frac{1}{q A} \left[\frac{C_{LF}}{(C_{ox} - C_{LF})} - \frac{C_{HF}}{(C_{ox} - C_{HF})} \right], \qquad (3)$$

where C_{LF} and C_{HF} are the low-frequency (or quasi-static) and high-frequency capacitances at a given voltage, respectively, C_{OX} the oxide capacitance per unit area, A the area of the electrode, and q the elementary charge. The high-frequency C-V curve was corrected for the presence of a series resistance by measuring the capacitance and the equivalent parallel conductance in strong accumulation from a conductance-voltage measurement.²⁰ All data acquisition and analysis were computer-assisted. The surface potential was obtained by numerical integration of the low-frequency C-V curve.²¹ All of the measurements of fixed oxide charge and interface trap charge reported in this study were the average of eight or more measurements as obtained from four measurements per sample, with two samples of each orientation prepared for each oxidation, and the results are shown in Table II. The standard deviation of Q_f and D_{it} for the cleanroom samples was less than 5% across each wafer, and the two samples processed identically agreed to within 10% of each other. The non-clearoom (111) and (110) samples showed greater deviations of 15% across a single wafer and 20-30% between two wafers processed identically. However, the non-cleanroom (100) samples displayed even greater deviations of up to 100% across a wafer. These results are considered unacceptable and will not be reported. The \pm values reported in Tables III and IV are the maximum spread of the data.

EXPERIMENTAL RESULTS

Mobile-Ion Contamination. It will be demonstrated here that the presence of mobile ionic charges can alter the observation of trends in other charges. Bias-temperature stress high-frequency C-V measurements²⁰ on the non-cleanroom samples yielded flatband voltage shifts due to mobile-ion contamination Q_m that ranged from - 0.25 V to - 1.0 V, which corresponds to positive Q_m values ranging from 1 to 5×10^{11} cm⁻², respectively, and cleanroom samples had flatband voltage shifts of less than - 0.04 V which corresponds to a positive value of Q_m of about 1.5×10^{10} cm⁻² for both the 3-inch and 4-inch wafers. An additional limitation imposed by the non-cleanroom evaporator is that a maximum of only four samples could be prepared simultaneously. Therefore, when preparing a large matrix of samples such as used in this study, there may be significant differences in Q_m between samples oxidized together but metallized at different times.

Table II contains a sampling of Q_f and D_{it} data for PMA samples. Cleanroom samples generally have lower Q_f and D_{it} values by factor as large as $4\times$ when compared with non-cleanroom samples. Also the variations with oxidation temperature are more pronounced and the spread in the data is larger for non-cleanroom samples. With the larger absolute values of charges arising from ionic contamination, and the larger variations from sample to sample and on single samples, possible trends in the data from process variations may be different. An illustration of this is that no trend is observed in Q_f with oxidation temperature for (100) cleanroom samples, but a clear trend is seen for the correponding non-cleanroom samples for all orientations. Also a minimum is seen in D_{it} for non-cleanroom samples. It is therefore crucial to report Q_m values as well as the other charge levels since it is clear that complex cooperative effects occur. Other cleanliness effects are reported below. Effects of Low-Temperature Hydrogen Annealing. Post-metallization annealing (PMA) at 400°C for 30 min in forming gas (FG) profoundly affects the shape of the C-V curves as is most evident in the depletion region which indicates the presence of interface traps.²⁰ With PMA, there is good agreement between the high- and low-frequency C-V curves in the depletion region, ; ad a corresponding lower interface trap density.

The distribution of interface arap densities as a function of energy in the bandgap indicated that the midgap D_{it} has been reduced by an order of magnitude on the PMA sample. Additionally, the well reported peak^{5,15,23} in the interface trap density distribution at about 0.3 eV above the valence band edge for NPMA samples usually observed for damaged Si surfaces is no longer present in the (100) samples and considerably reduced in the (111) samples in the interface trap density distribution of PMA samples. This peak was found on all samples of different orientations that were rapidly removed from the oxidation furnace and cooled in the oxidizing ambient. No other distinct peaks were observed in our samples. Electron-spin resonance studies (ESR)²³ have related this peak to the P_b center which is a trivalent silicon defect or dangling bond. Thus, hydrogen effectively reduces this state on all orientations studied as well as reducing the concentration of all of the silicon bandgap interface traps.

For some NPMA samples, the interface trap densities were so high that a midgap D_{it} was not measurable. These high concentrations of interface traps effectively result in Fermi level pinning thereby rendering part of the silicon gap inaccessible. This was particularly true for (111) samples and occasionally true for (110) samples, which had the highest interface trap densities of the samples studied. For the purpose of comparing orientations, however, the D_{it} data was obtained at potentials in the silicon gap other than at midgap, as noted in the tables. Also from Fig. 2, it is observed that the PMA more dramatically affects the upper part of the silicon bandgap where the deep acceptor traps are located. The upper portion of the bandgap had unmeasurably large concentrations of interface traps - as indicated in the quasi-static C-V curve - prior to PMA which are reduced to below the other levels after PMA. It is interesting to note that for the (111) and (110) orientations. a minimum occurs at about 0.3 eV above the valence band edge for the NPMA samples.

which enables the use of this potential to compare D_{it} levels. By contrast, (100) samples and radiation-damaged silicon show a peak at this surface potential.

Substrate Orientation and Oxidation Temperature Effects. Fixed oxide charge for the cleanroom samples is plotted as a function of oxidation temperature for the four substrate orientations in Fig. 3 for PMA and NPMA samples and the data is shown in Table III along with the maximum spread in the data. The ambiguity in Q_f values for the NPMA samples as discussed above should be remembered and thus the usefullness of such data in the interpretation of fundamental phenomena is limited. However, these data are included to investigate the effect of the silicon orientation and oxidation temperature on trends in the charges of samples that have received either no anneals or incomplete anneals, since discrepancies exist in the literature. For all samples and temperatures greater than 750°C, the (111) orientation has the highest fixed oxide charge, with the order in Q_f being

$$Q_{f}(111) > Q_{f}(110) \ge Q_{f}(100) \ge Q_{f}(511)$$

At 750°C, there is an apparent change in the ordering, with the (110) orientation having the largest fixed oxide charge. However, while this effect appears to be real and is observed in both PMA and NPMA samples, the maximum spread in the data for NPMA samples precludes an unambiguous interpretation.

Higher oxidation temperatures generally yield lower values of Q_f and while the trend is discernible for PMA samples it is particularly evident for NPMA samples. It is also observed that (110), (511), and (100) PMA samples, oxidized at a temperature higher than 750°C, have approximately the same Q_f value of $1 - 2 \times 10^{11}$ cm⁻². The largest change between PMA and NPMA samples is seen with the (511) orientation which has a low value of 2×10^{10} cm⁻² at 800°C. Oxide fixed charges on NPMA samples are $4 \times$ larger and the spread in the data $2 \times$ larger than on PMA samples.

Fixed charge density was also obtained for the three orienations processed in nonclean room conditions. These results are plotted in Fig. 4 for PMA samples where it can seen that Q_f decreases with increasing temperature more strongly than in wafers processed under similar conditions in cleanroom conditions. However, the values of non-cleanroom Q_f are also significantly larger, the non-cleanroom results did not exhibit a crossover with different orientations. For example, Q_f at 750°C is 3× larger for the non-cleanroom (111) orientation than on cleanroom samples, and 6× larger on the (100) orientation. Similar results are found for NPMA non-cleanroom samples but with both larger values and greater deviations.

Interface trap charge density (D_{it}) displays an orientation dependence similar to that of $Q_{\rm f}$, as is shown in Fig. 5 for PMA and NPMA samples, and that data with maximum spreads is listed in Table IV. For temperatures above 750°C, the orientational ordering is clearly

$$Q_{f}(111) > Q_{f}(110) > Q_{f}(100) \ge Q_{f}(511)$$

with NPMA samples having D_{it} an order of magnitude higher than on PMA samples. As discussed above for Q_{f} , for both the NPMA and PMA samples grown at 750°C, there is an apparent change in the ordering of D_{it} , with the (110) orientation having the highest interface trap charge density. However, if we consider that for (110), D_{it} is $16 \pm 1 \times 10^{10}$ cm⁻²eV⁻¹ at midgap and for (111) it is $13 \pm 0.6 \times 10^{10}$ cm⁻²eV⁻¹, the maximum deviations make the values nearly the same within 10%. Without post-metallization annealing, D_{it} values for the (111) and (11C) orientations were too large to be determined at midgap. However, near a surface potential of 0.3 eV there was no crossover until the oxidation temperature was 1000°C, but here again the maximum error limits in the data render conclusions about the crossover tenuous. There was no crossover in D_{it} for the noncleanroom samples. Thus, there is a definite orientation dependence for both Q_{f} and D_{it} . However, whether the (111) or (110) orientation exhibits the larger D_{it} is a function of processing, and for NPMA samples, the (110) orientation has a higher D_{it} but the (111) orientation has a higher Q_{f} . It is shown that impurities and annealing greatly alter both the magnitude and orientational dependence of both fixed oxide and interface trap charges.

The temperature dependence of interface trap charges on cleanroom samples is also shown in Fig. 5. These are all midgap values and, as stated above, midgap values were not always measurable on NPMA samples. This explains why the much higher values of midgap D_{it} on (111) samples are not included in Fig. 5. All samples show an overall decreasing D_{it} with increasing oxidation temperature. However, the magnitude of the decrease depends on both the substrate orientation and the treatment. The midgap D_{it} on the (100) and (110) orientations of PMA samples decrease more sharply with increasing oxidation temperature than on the other orientations. The midgap D_{it} at 1000°C is 2.5 and 2.0 times smaller than those on samples oxidized at 750°C for the (100) and (110) orientations, respectively. The midgap D_{it} on the (511) samples shows only a slight decrease at the higher oxidation temperatures, and the (111) orientation shows little temperature dependence within experimental error. If the displayed temperature dependence is also considered to be dependent on the substrate orientation, then we find that the change in D_{it} with temperature (ΔD_{it}) over the temperature range in this study follows the order

$$\Delta D_{it}(100) \geq \Delta D_{it}(110) > \Delta D_{it}(511) > \Delta D_{it}(111) .$$

There is a stronger temperature dependence of D_{it} on NPMA samples, with D_{it} at 1000°C being three times smaller than that at 750°C for the (100) and (511) orientations.

This orientation-dependent temperature behavior was also observed in the non-cleanroom experiment, where an even larger temperature range was investigated. The interface trap density at 0.3 eV for (111), and at midgap for the (110) orientation of NPMA samples is plotted in Fig. 6. The interface trap density for the (111) orientation is fairly constant over this temperature range. Again, the non-cleanroom values are higher than the cleanroom values by approximately a factor of two.

Similar observations have already been reported. Razouk and Deal⁵ reported a temperature dependence of interface traps in the 900°C-1200°C range for NPMA cleanroom samples prepared in a manner almost identical to that used in this study. From their data.⁵ the slope of the interface trap density as a function of temperature for the (100) and (111) orientations was calculated, and the results are compared with the results of this study in Table V. The (100) orientation displays a stronger temperature dependence by nearly an order of magnitude, and good agreement is seen between the results of this study and those of Razouk and Deal.⁵ As the oxides studied by Razouk and Deal were much thicker than those investigated in this study, a quantitative comparison of D_{it} values with this study could only be done for the 8-hour 1000°C oxidation. As-oxidized midgap interface trap density reported by Razouk and Deal⁵ at 1000°C were 2.5×10^{12} cm⁻² and 9.0×10^{10} cm⁻² for the (111) and (100) orientations, respectively. The corresponding values obtained in this study were 2.2×10^{12} cm⁻² and 1.7×10^{11} cm⁻² for the (111) and (100) orientations. respectively.

Oxidation Rate Effects. In order to investigate the influence of the oxidation rate on Q_f and D_{it} , several experiments were carried out on SiO₂ films grown at 1000°C. This temperature was selected to minimize any effects that might arise due to the presence of intrinsic stress at the interface which is known to be smaller at higher oxidation temperatures.^{24,25} The thickness of the oxides in this experiment varied from 500-640Å for the 1-hour oxidation to 2050-2400Å for the 8-hour oxidation for the (100) and (111) orientations, respectively.

The dependence of fixed oxide charge results on oxidation time at 1000°C is plotted in Fig. 7 for PMA samples with the data listed in Table VI. For all but the (110) orientation, a slight decrease in Q_f is observed for longer oxidation times, i.e. for slower oxidation rates and thicker oxides. The results are similar for NPMA samples, whose data is listed in Table VI, but with approximately 2× larger charges. An increase in Q_f for thinner oxides has also been reported.¹¹

The dependence of interface trap charges of PMA samples on oxidation time at 1000°C is plotted in Fig. 8 with the data listed in Table VII. A more discernible decrease in interface trap charge with increasing oxidation time and decreasing oxidation rate is ovserved. Hung and Cheng¹⁵ have reported a similar dependence of D_{it} for oxide thicknesses less than 500Å, and their etch-back experiment confirmed that the observed thickness dependence was due to the oxidation process, and not due to any subsequent steps. Here again, the results are similar for NPMA samples, whose data is listed in Table VII, but with a $10 \times$ larger charge density and a smaller change with oxidation time.

The dependence of the overall SiO_2 growth rate on oxidation temperature for the four orientations is shown in Fig. 9. The overall oxidation rate was obtained by dividing the final oxide thickness by the total oxidation time. A crossover between the oxidation rates on the (110) and (111) orientations is seen at 750°C. Such a rate crossover has been previously reported,²⁶⁻²⁸ and studied in detail. Figures 3 and 5, showing the dependence of Q_f and D_{it} on oxidation temperature, display evidence for crossover or at least for the convergence of the charge levels with oxidation temperature. It should be noted that the overall oxidation rate is different from the oxidation rate at the end of the oxidation before the sample is removed from the furnace. In this study, the oxidation time at 750°C is pasthe reported interfacial oxidation-rate crossover point.²⁶⁻²⁸ Therefore, the final interfacial oxidation rate ordering is actually (111)>(110). The overall thicknesses of the oxides do not yet reflect this change, as the (111) must catch up with (110), i.e. the thickness crossover lags the rate crossover.²⁶⁻²⁸ The non-cleanroom samples were oxidized one hour longer than the cleanroom samples, and the change in interfacial oxidation rate is evident in the overall growth rate, i.e. the thickness of the oxides on the (111) orientation samples was larger than that on (110) samples after the additional hour. The non-cleanroom results did not show the same crossover in D_{it} and Q_{f} , with the (111) orientation always showing greater charge levels, but when one considers the actual oxidation rate at the interface. this is not necessarily inconsistent with the above results. Thus, the data strongly suggest that the charges scale with the fastest oxidizing surface.

For charges that exist at the Si-SiO₂ interface and possibly scale with the SiO₂ growth. a correlation should be attempted with the interfacial oxidation rate and not the overall oxidation rate. It is also possible that the differences in D_{it} at 750°C for the cleanroom samples are not statistically significant, with D_{it} of 1.55×10^{11} cm⁻² and 1.34×10^{11} cm⁻² on (110) and (111), respectively. However, the difference in Q_f is significant, with values of 1.1×10^{12} cm⁻² and 5.9×10^{11} cm⁻² on (110) and (111), respectively. This can be understood by considering the oxidation rate at the interface. Initially, the (110) orientation is oxidizing the fastest. We have shown that the substrate orientation is important in determining the ordering of Q_f and D_{it} which are both the highest on (110). Since new oxide is always being formed at the interface, any change in oxidation rate ordering – and consequently the orientational ordering of Q_f and D_{it} - will first be evident at the interface and in the interface trap charge. The fixed oxide charge, on the other hand, is thought to be displaced by about 50Å from the interface,²² and the change will not be detected until those 50Å have replaced by the new oxide. It appears that a slow interfacial oxidation rate at a constant oxidation temperature yields low values of Q_f and D_{it} for otherwise similar sample preparation procedures. While this is not yet fully understood, it is useful processing information. One need simply follow a rapid oxidation or other film formation procedure by a slower thermal oxidation which will effectively decrease the interfacial charges to acceptable levels. It is also well known that a zero oxidation rate or an anneal in an inert ambient decreases both Q_f and D_{it} values.⁷ However, we have no reason to believe that this anneal effect is the same as the low oxidation-rate effect.

DISCUSSION

Silicon Surface Orientation. The observed dependence of the interface traps and fixed oxide charges on the silicon substrate orientation can be compared with the density of silicon surface atoms (C_{Si}) which is known to follow the order²⁷

$$C_{Si}(110) > C_{Si}(111) > C_{Si}(511) > C_{Si}(100)$$
.

The (511) orientation is of particular interest because recent diffraction studies have identified the (511) plane as a vicinal surface consisting of (100) steps and (111) risers.²⁹ It has been establised that the silicon oxidation rate in the initial oxidation regime of oxides up to 300Å follows the order of C_{Si} over the 600°C-1000°C temperature range.^{27,28} For larger thicknesses, a crossover is observed with the (111) orientation yielding a higher oxidation rate than the (110) orientation. In the present study with 500Å-thick oxide films, the oxidation regime was above the crossover to where the (111) orientation is the fastest. The observed D_{it} and Q_f values track with the oxidation rates rather than the silicon atom areal density with the (111) yielding both the higher charge levels and interfacial oxidation rate. Near the lowest temperature investigated of 750°C, D_{it} values on the (111) and (110) orientations become nearly equal with a possibility of a crossover. This suggests that at lower temperatures, the dominance of C_{Si} is regained. Additional studies using very thin oxides grown at even lower temperatures will be required to further elucidate this issue.

A quantitative comparison of D_{it} to the surface atom density is made by comparing the ratio of surface atom densities for the same two orientations. This correlation is shown in Table VIII for the data at 750°C, with the quantitative agreement being quite good but diverging above 750°C for all orientations. This gives credence to the idea that at lower oxidation temperatures, the silicon atom density is important. It is noted that Q_f and D_{it} for the (511) orientation displayed no clear distinction from the (100) orientation, with values both greater than and smaller than Q_f and D_{it} on the (100) orientation. There was considerably greater variation in the measured values on (511), possibly owing to the more complex vicinal nature of the orientation.²⁹ It should also be noted that the orientation and oxidation temperature effects are somewhat obscured both by PMA in hydrogencontaining ambients and the lack of cleanroom conditions. Hydrogen atoms bonded to the silicon surface reduce all charges while other impurities usually tend to increase them. The result is that the strictly silicon-orientation-dictated amounts of charges are altered and perhaps masked by the influence of hydrogen and/or impurities on the charges. It has often been reported that chemical-vapor-deposited oxide (CVD SiO_2) films exhibit higher Q_f and D_{it} levels than thermally oxidized silicon. However, recently published studies show that with the use of very low deposition rates along with careful control of other processing variables and cleanroom conditions, the charge levels for CVD oxides can approach those of thermally oxidized silicon, suggesting that the improved registration obtainable with slower deposition rates across the substrate-film interface is important.^{30,31} Furthermore, SiO₂ grown at 1100°C on (100) silicon with PMA show decidedly smaller Q_f as the partial pressure of oxygen is decreased.³² The decrease in oxygen reduces the oxidation rate at the same oxidation temperature.

Intrinsic Stress. Intrinsic SiO₂ stress (σ_i) which arises from the molar volume change

from Si to SiO_2 has been observed to follow the silicon substrate orientation order as $^{25,33-35}$

$$\sigma_i(110) > \sigma_i(511) \ge \sigma_i(100) > \sigma_i(111)$$
.

As the oxide growth strain which is caused by the molar volume change during oxidation is thought to be the same regarless of the silicon orientation, the ordering for the resulting σ_i was predicted to follow Young's modulus. This would lead to the (111) orientation as having the highest stress value.²⁸ This prediction was not borne out in experiment and several similar models have proposed that if oxidation occurs laterally at surface steps then a smaller stress would result.^{36,37} It has been shown that the (111) orientation of silicon has a step density nearly twice that on the (100) orientation,³⁸ and a detailed mechanism has been proposed for the production and oxidation of steps on a silicon surface.³⁷ For all silicon orientations, σ_i decreases with increasing oxidation temperature and after hightemperature annealing.^{24,25,39,40}

With the exception of the (111) orientation, both Q_f and D_{it} follow the intrinsic stress related orientation and oxidation temperature dependence. It should be noted that the thermal expansion component of the total stress (σ_{th}) develops upon cooling from the oxidation temperature and has the opposite temperature dependence of both σ_i and the charges.²⁵ The anomalously low σ_i observed for the (111) orientation was attributed to lateral oxidation at steps or at rough surface regions, and may also be the cause for the highest charge level found on this orientation. Roughness is known to increase both interface traps and charges.³⁸

The present experimental results combined with published studies show that no single process paramater determines the variation of the interface charges. However, based on the experimental results now available, it is reasonable to propose that the measured values of Q_f and D_{it} can derive from the two primary sources. First, on the oxidized silicon surface, an initial charge level arises from the registry of the silicon surface with the oxide. At the interface, the large change in volume causes a misregistry of bonds with commensurate intrinsic stress leading to orientational and stress sensitive interface charges. It is then easy to imagine how the oxidation rate, the oxidation temperature, the oxygen partial

pressure, the silicon orientation, and the specific geometry across the interface can affect the registry and hence the charges. Second, there is a relaxation component that can reduce the misregistry and thereby lower the interfacial charge. Higher oxidation temperatures (reduced stress), reduced oxidation rate (stress relaxation and registry) and annealing are the primary relaxation procedures. Annealing (temperature and time) enhances the attainment of optimum registry, and annealing with hydrogen atoms can further reduce interface charges through hydrogen-bond formation at sites that for geometrical (steric) reasons cannot attain full bond registry across the interface.

SUMMARY

1. The effects of cleanliness on interface charges and traps were measured using Q_m , and it was found that both Q_f and D_{it} increased with an increase in Q_m . This observed trend can be both enhanced or obscured with high Q_m levels. This factor along with different anneals may explain why different trends were reported in the literature for the processing effects on charges.

2. Strong orientation effects are observed for both Q_f and D_{it} with the (111) orientation displaying the highest charge levels at all but the lowest oxidation temperature in this study where the (110) orientation dominates. This change in ordering strongly suggests that at higher oxidation temperatures, factors other than the silicon surface atom density is important for controlling charges.

3. Both Q_f and D_{it} decrease with higher oxidation temperatures. Intrinsic stress also decreases with increasing oxidation temperature and at first a direct correlation seems to exist between charge and stress. However, the (111) orientation has the largest charge at oxiation temperatures higher than 750°C and the lowest intrinsic stress. The anomalously low stress for the (111) orientation was attributed to surface steps or roughness. Perhaps this also causes the higher charge levels.

4. A crossover in both oxidation rate and charges leads to the notion of the dependence of charge on interfacial oxidation rate. The specific process steps leading to better interfacial

registry or the attainment of registry seem to provide a working model for predicting the final interfacial charge levels both for thermal oxides in this study and for CVD oxides reported in the literature. Thus, annealing, high oxidation temperatures, and low oxidation rates all reduce charge and stress, while the opposite processing increase charge.

5. Anneals in hydrogen-containing ambients can dominate both Q_f and D_{it} by reducing charge levels by orders of magnitude and leveling both oxidation temperature and silicon surface atom density trends in the charges. The marked effect on deep acceptor states was observed for the (111) and (110) silicon orientations. The notion that hydrogen bonds preferentially to certain energetically different sites may provide an analysis tool for the different sites.

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FIGURE CAPTIONS

Figure 1. Schematic representation of the effect of oxidation temperature on some oxide property F(T), where F(T) is the refractive index n, the oxide density ρ , the intrinsic stress σ_i , and the interface charges Q_f and D_{it} from Ref. 18.

Figure 2. Interface trap density D_{it} plots for a p-type (111) sample (a) without a postmetallization anneal (NPMA) and (b) with a post-metallization anneal (PMA).

Figure 3. The dependence of fixed oxide charge Q_f on oxidation temperature for the four silicon orientations studied in the cleanroom experiment, (a) with a post-metallization anneal (PMA), and (b) without a post-metallization anneal (NPMA).

Figure 4. The dependence of fixed oxide charge Q_f on oxidation temperature for the three silicon orientations studied in the non-cleanroom experiment including a post-metallization anneal (PMA).

Figure 5. The dependence of interface trap density D_{it} on oxidation temperature for the four silicon orientations studied in the cleanroom experiment, (a) with a post-metallization anneal (PMA), and (b) without a post-metallization anneal (NPMA).

Figure 6. The dependence of interface trap density D_{it} on oxidation temperature for the (111) and (110) silicon orientations studied in the non-cleanroom experiment without a post-metallization anneal (NPMA).

Figure 7. The dependence of oxide fixed charge Q_f on oxidation time for the four silicon orientations studied in the cleanroom experiment including a post-metallization anneal (PMA).

Figure 8. The dependence of interface trap density D_{it} on oxidation time for the four silicon orientations studied in the cleanroom experiment including a post-metallization anneal (PMA).

Figure 9. Overall oxiation rate at 750°C and 800°C showing a crossover in the rates of oxidation of the (110) and (111) orientations.



SiO₂ Film Properties, F(T), vs Oxidation Temperature









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(nim\Å) JTAA NOITADIXO

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Table I.	Oxidation	temperatures	and	corresponding	times	used	for	each	oxidation.	
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Oxidation Temperature	Oxidation Time
(°C)	(hr)
Non-Cleanroom Conditions	s: (100), (110), and (111).
750	24.0
800	18.0
900	4.0
1000	1.0
1100	0.5
Cleanroom Conditions: (100	0), (110), (111) and (511).
750	23.0
800	18.0
900	4.5
1000	8.0
1000	4.0
1000	2.5
1000	1.0

	Oxidation	Oxidation Fixed Oxide Charge Density N_f ($N_f (10^{11} { m cm^{-2}})$
	Temperature (°C)	(100)	(110)	(111)
Cleanroom	750	0.74	10.80	5.89
Conditions:	800	1.60	1.92	4.99
	900	1.27	1.71	4.86
	1000	1.73	1.08	4.33
Non-cleanroom	750	6.30	8.20	11.80
Conditions:	800	3.12	3.35	6.73
	900	2.87	3.45	6.78
	1000	2.43	2.90	6.60
	1100	1.42	2.02	3.92
Deal et al. ^{a)}	920	1.70	2.10	4.70
	1200	0.20	0.60	1.70
Razouk and Deal ^{b)}	1000	0.80	_	3.30

Table II. Comparison of fixed oxide charge and interface trap density for samples that were post-metallization annealed (PMA).

a) Ref. 3.

.

b) Ref. 4.

Oxidation	Interface Trap Charge Density D_{it} at Midgap			
Temperature (°C)	for (111) Si $(10^{11} \text{ cm}^{-2} \text{eV}^{-1})$			
	Cleanroom Conditions:	Non-cleanroom Conditions:		
750	1.30	3.20		
800	1.30	2.80		
900	1.00	1.60		
1000	1.20	4.90		

Oxidation	Cleanroom Fixed Oxide Charge Density N_f (10 ¹¹ cm ⁻²)					
Temperature	(Ma:	ximum Spread in	Data Indicated a	us ±)		
(°C)	(102)	(511)	(110)	(111)		
Post-metallization Annealed Samples (PMA)						
750	$0.7{\pm}0.4$	$1.6 {\pm} 0.4$	10.8 ± 1.2	$5.9{\pm}0.4$		
800	$1.6 {\pm} 0.5$	$0.2{\pm}0.2$	$1.9{\pm}0.3$	$5.0{\pm}0.6$		
900	$1.3 {\pm} 0.2$	$0.6{\pm}0.2$	$1.7{\pm}0.3$	$4.9{\pm}0.1$		
1000	$1.7{\pm}0.1$	1.8 ± 0.1	1.8 ± 0.1	$4.3 {\pm} 0.2$		
N	Ion-Post-metalliza	tion Annealed Sa	amples (NPMA)			
750	7.8 ± 1.6	6.2 ± 0.4	16.0 ± 1.3	15.2 ± 0.6		
800	$6.5 {\pm} 0.9$	6.2 ± 0.3	$8.5 {\pm} 0.8$	$13.3{\pm}0.4$		
900	$4.3 {\pm} 0.1$	$4.2{\pm}0.1$	6.5 ± 0.1	11.5 ± 0.3		
1000	3.3 '-0.3	$3.3 {\pm} 0.3$	$5.1 {\pm} 0.5$	9.6 ± 0.6		

Table III. Fixed oxide charge for post-metallization annealed (PMA) and non-post-metallization annealed (NPMA) samples.

Oxidation	Cleanroom Interface Trap Density D_{it} (10 ¹¹ cm ⁻² eV ⁻¹)					
Temperature	(Ma	ximum Spread in	Data Indicated a	us ±)		
(°C)	(100)	(511)	(110)	(111)		
Post-metailization Annealed Samples (PMA)						
750	12.0 ± 0.1	$7.5 {\pm} 0.4$	16.0±1.0	$13.0 {\pm} 0.6$		
800	$9.0 {\pm} 0.6$	$8.0{\pm}1.2$	$9.5 {\pm} 0.4$	13.0 ± 0.1		
900	$6.0 {\pm} 0.3$	$5.6 {\pm} 0.2$	$9.3{\pm}0.4$	$10.0 {\pm} 0.3$		
1000	$5.0{\pm}0.1$	$5.7{\pm}0.1$	$7.9{\pm}1.0$	12.0 ± 0.6		
N	Ion-Post-metalliza	ation Annealed Sa	amples (NPMA)			
750	8.2±0.2	$9.5{\pm}2.0$	35.0 ± 0.5^{a}	$22.0 \pm - a^{a}$		
800	$5.4{\pm}0.2$	$6.7{\pm}0.2$	22.0 ± 0.1^{a}	$16.0 \pm -a)$		
900	$4.1 {\pm} 0.2$	$6.8{\pm}1.4$	36.0 ± 14^{a}	$14.0 \pm - a$		
1000	$2.6 {\pm} 0.1$	3.2 ± 0.2	8.0 ± 1.2	$11.0 \pm -a$		

Table IV. Interfac trap density for cleanroom post-metallization annealed (PMA) and non-post-metallization annealed (NPMA) samples.

a) Values reported at 0.3 eV above the valence band where a minimum occurs.

Table V. The oxidation-temperature dependence of the slope of interface trap density (D_{it}/T) for non-post-metallization annealed (111) and (100) samples, in the 900-12-°C range from Ref. 4 and the 750-1000°C in this work.

	$\Delta D_{it} / \Delta T \ (10^{11} \ {\rm cm}^{-2} {\rm eV}^{-1} {\rm K}^{-1})$		
	(111)	(100)	
Razouk and Deal	-0.00433	-0.030	
as oxidized. ^{a)}			
This Work, as-oxidized,	-0.00488	-0.023	
cleanroom conditions.			
This Work, as-oxidized,	-0.00314	_	
non-cleanroom conditions.			

a) Ref. 5.

Oxidation	Cleanroom Fixed Oxide Charge Density $N_f (10^{11} \text{ cm}^{-2})$						
Time (hr)	(100)	(511)	(110)	(111)			
Post-metallization Annealed Samples (PMA)							
1.0	1.7	1.8	1.1	4.3			
2.5	1.7	1.2	1.0	4.4			
4.0	0.9	0.7	1.2	3.4			
8.0	1.3	1.1	1.0	3.8			
	Non-Post-metalli	zation Annealed S	Samples (NPMA)				
1.0	3.3	3.3	5.1	9.6			
2.5	2.8	2.5	4.3	8.9			
4.0	2.8	2.8	3.4	8.0			
8.0	3.0	3.0	3.9	8.4			

Table VI. Fixed oxide charge Q_f on samples oxidized at 1000°C with and without post-metallization annealing.

Oxidation	Interface Trap Charge Density D_{it} at Midgap (10 ¹⁰ cm ⁻² eV ⁻¹)							
Time (hr)	(100)	(511)	(110)	(111)				
Post-metallization Annealed Samples (PMA)								
1.0	4.9	5.7	7.9	12.0				
2.5	2.1	3.4	5.7	7.9				
4.0	2.6	1.5	2.8	5.0				
8.0	2.2	2.3	2.8	6.0				
	Non-Post-metal	lization Annealed	Samples (NPMA	.)				
1.0	26	32	80	11 a)				
2.5	17	22	89	73 a) .				
4.0	15	22	73	56 ^{a)}				
8.0	17	39	80	89 a)				

Table VII. Interface trap charge D_{it} on samples oxidized at 1000°C with and without post-metallization annealing.

a) Measured at 0.3 eV above the valence band edge and not at midgap.

Table VIII. Ratios of the measured interface trap density D_{it} and the silicon surface atom density C_{Si} for the various orrientations oxidized at 750°C and post-metallization annealed.

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	(110)/(100)	(111)/(100)	(511)/(100)	(111)/(110)
ratios of D_{it}	1.30	1.10	0.63	0.81
ratios of C_{Si}	1.40	1.15	0.63	0.86