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|--|------------------------|------------------------|
| NAME OF CONTRACTOR:                              | Paul M. Raccah         | DTIC                   |
| CONTRACT NUMBER:                                 | DAAK 70-83-K-0047      | ELECTE                 |
| EFFECTIVE DATE OF CONTRACT:                      | 2/1/83                 | AUG 0 6 1990           |
| EXPIRATION DATE OF CONTRACT:                     | 1/31/86                |                        |
| REPORTING PERIOD: / Title                        | FIFTH QUARTER for DAAN | (2 U<br>(70-83-K 04/10 |
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#### FIFTH QUARTER

#### RESEARCH PLAN

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This quarter has been devoted to the study of the physical significance of the parameter  $\Gamma$  which can be obtained by Electrolyte Electroreflectance (EER) or by spectroscopic ellipsometry.

### RESULTS

The generalized theory of electroreflectance (which has been carried out under separate funding) has permitted us to reconcile the  $\Gamma$  obtained by electroreflectance to the  $\Gamma$  obtained by ellipsometry.

However, it does remain that the value of  $\Gamma$  obtained either way can be as high as 120 meV as compared to the expected 70 meV. We have determined t at this is due to the breakdown of the virtual crustal assumption and is in fact a quantitative measure of it. It turns out that the measured  $\Gamma_{\rm m} = \Gamma_{\rm O}$  + kT +  $\Delta\Gamma$  where  $\Gamma_{\rm m}$  is the measured linewidth,  $\Gamma_{\rm O}$  is the expected 70 meV and  $\Delta\Gamma$  measures the departure from the virtual crystal approximation. (See Attached Expose)

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## SIMPLE APPROXIMATION FOR EFFECT OF ALLOYING ON THE PHENOMENOLOGICAL LINEWIDTH Γ

Consider that  $E_{CV}(\vec{k})$  is not a unique, sharply defined energy, but that for the absorption of a photon "locally", it depends on the local concentration  $c_{loc}$  averaged over a cluster of N atoms on the Hg-Cd sublattice.

$$E_{cv}^{loc}(\vec{k}) \approx E_{cv}^{loc}(\vec{k}_{cr}) + \frac{\aleph^2}{2\mu}(\vec{k} - \vec{k}_{cr})^2;$$

i.e., the dominant source of variation in  $E_{CV}^{loc}(\vec{k})$  is the variation in the local critical point energy

 $E_{o} (c_{loc}) \equiv E_{cv}^{loc} (\vec{k}_{cr}).$ 

Then, for N > > 1,

$$P(E_0) \approx (\sigma \sqrt{2\pi})^{-1} \exp \{-[E_0 - E_0(c)]^2/2\sigma^2\}$$

with  $\sigma^2 = E_1 c(1-c)/N$ ,

where  $E(c_{loc}) \approx E(c) + (c_{loc} - c) E_1$ .

This leads to a replacement of the lineshape

$$L(E,\vec{k},\Gamma_{O}) = -[E - E_{CV}(\vec{k}) + i\Gamma_{O}]^{-1}$$
  
by  $f_{O}(E,\vec{k},\Gamma_{O}) = \int_{-\infty}^{\infty} \{E - E_{CV}(\vec{k}) - [E_{O} - E_{O}(c)] + i\Gamma_{O}\}^{-1} P(E_{O}) dE_{O}$ 

The only simple analytic result is obtained by replacing the Gaussian probability  $P(E_0)$  by a Lorentzian probability. If one does this and chooses the Lorentzian probability to have width

 $\Gamma^1 = \sqrt{2\sigma}$ 

which follows from an expansion of  $e^{-u^2}$  as  $[1 + u^2 + ...]^{-1}$ , one finds that  $\Gamma_0$  is replaced by

$$\Gamma_{\rm m} = \Gamma_{\rm O} + \sqrt{2\sigma} = \Gamma_{\rm O} + E_1 \sqrt{2c(1-c)/N} + kT$$

This gives the following table:

| Γ-(Γ <sub>0</sub> +kT) | .04eV | .06eV | .08eV | .10eV |
|------------------------|-------|-------|-------|-------|
| N                      | 200   | 88    | 50    | 32    |

A better numerical approximation leads to values of N approx 40% larger.