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COY, Ro TYPE OF EMORANDU SUPPLEME FIELD ABSTRACT he calib park photo or a fir aw dampi on-linea ransonic	bert L. REPORT M INTARY NOTA COSATI GROUP 01 (Continue on per .50 Ba stography re control ing rates ar behavic and subs TION / AVAILAL SIFIED/UNLIMI F RESPONSIBL	I 3b TIME (C FROM	OVERED TO TO TO TO TO TO TO TO TERMS (Caliber .50 Bu Aerodynamic Cha Aerodynamic Dr. and identify by block r M8 and APIT, M termine the com odynamic drag, g ned at supersonic sus moment coeff RPT DTIC USERS	14 DATE OF REPO 1989 Decem (Continue on reverse llets aracteristics ag iumber) 20 munitions plete aerobal yroscopic sta c, transonic icient predic 21 ABSTRACT SE UNCLASSIFII 22b TELEPHONE (RT (Year, Month, iber 5 e if necessary and Gyros b Dynam Yaw L were recent listic char ibility, dyn and subsoni its a slow a curry classific ED	Dey) 15 indentify b copic S ic Stab imit-Cy ly fire acteris amic st c speed irm limit ATION e) 22c OF	PAGE COUNT 75 by block number) Stability vility vile ed in the BRL stics required ability and is. The observ it cycle yaw at
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Table of Contents

	List of Figures
	List of Tables
I.	Introduction
II.	Test Procedure and Material
III.	Results
	1. Drag Coefficient
	2. Overturning Moment Coefficient
	3. Gyroscopic Stability
	4. Lift Force Coefficient
	5. Magnus Moment Coefficient and Pitch Damping Moment Coefficient
	6. Damping Rates
IV.	Conclusions
v.	Recommendations
	References
	List of Symbols
	Distribution List



٠.

Page



iii

Ŋ

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List of Figures

Figure	<u><u>P</u></u>	age
1	Photograph of the Caliber .50 Projectiles	8
2	Photograph of the BRL Free Flight Aerodynamics Range	9
3	Coordinate System for the BRL Aerodynamics Range	10
4	Sketch of the Caliber .50 Projectiles	11
5	Shadowgraph of Ball, M33 Projectile at Mach 2.66	12
6	Shadowgraph of API, M8 Projectile at Mach 2.60.	13
7 .	Shadowgraph of API, M8 Projectile at Mach 2.60, Angle of Attack = 15.4 Degrees	14
8	Shadowgraph of APIT, M20 Projectile at Mach 2.33	15
· 9	Shadowgraph of Ball, M33 Projectile at Mach 1.99	16
10	Shadowgraph of API, M8 Projectile at Mach 2.04	17
11	Shadowgraph of APIT, M20 Projectile at Mach 2.01	18
12	Shadowgraph of Ball, M33 Projectile at Mach 1.53	19
13	Shadowgraph of API, M3 Projectile at Mach 1.51.	20
14	Shadowgraph of APIT, M20 Projectile at Mach 1.45	21
15	Shadowgraph of Ball, M33 Projectile at Mach 1.00	22
16	Shadowgraph of Ball, M33 Projectile at Mach 0.92	23
17	Shadowgraph of Ball, M33 Projectile at Mach 0.89	24
18	Shadowgraph of API, M8 Projectile at Mach 0.90.	25
19	Zero-Yaw Drag Force Coefficient versus Mach Number, Ball, M33	26
20	Zero-Yaw Drag Force Coefficient versus Mach Number, API, M8	27
21	Zero-Yaw Drag Force Coefficient versus Mach Number, APIT, M20	28
22	Yaw Drag Force Coefficient versus Mach Number	29
23	Comparison of Old and New Drag Coefficients for the API, M8 Projectile	3 0
24	Zero-Yaw Overturning Moment Coefficient versus Mach Number, Bail, M33.	31
25	Zero-Yaw Overturning Moment Coefficient versus Mach Number, API, M8.	32
26	Zero-Yaw Overturning Moment Coefficient versus Mach Number, APIT, M20.	33

v

List of Figures (Continued)

gure		Page
27	Cubic Overturning Moment Coefficient versus Mach Number	34
2 8	Zero-Yaw Lift Force Coefficient versus Mach Number, Ball, M33	35
29	Zero-Yaw Lift Force Coefficient versus Mach Number, API, M8	36
30	Zero-Yaw Lift Force Coefficient versus Mach Number, APIT, M20	37
31	Cubic Lift Force Coefficient versus Mach Number	38
32	Magnus Moment Coefficient versus Effective Squared Yaw	39
33	Magnus Moment Coefficient versus Effective Squared Yaw	40
34	Magnus Moment Coefficient versus Effective Squared Yaw	41
35	Zere-Yaw Magnus Moment Coefficient versus Mach Number, Ball, M33	42
36	Zero-Yaw Magnus Moment Coefficient versus Mach Number, API, M8	43
37	Zero-Yaw Magnus Moment Coefficient versus Mach Number, APIT, M20.	4 4
38	Cubic Magnus Moment Coefficient versus Mach Number	45
39	Zero-Yaw Pitch Damping Moment Coefficient versus Mach Number, Ball, M33	46
40	Zero-Yaw Pitch Damping Moment Coefficient versus Mach Number, API, M8	47
41	Zero-Yaw Pitch Damping Moment Coefficient versus Mach Number, APIT, M20	48
42	Fast Arm Damping Rate versus Effective Squared Yaw	49
43	Slow Arm Damping Rate versus Effective Squared Yaw	50
44	Fast Arm Damping Rate versus Effective Squared Yaw	51
45	Slow Arm Damping Rate versus Effective Squared Yaw	52
46	Fast Arm Damping Kate versus Effective Squared Yaw	53
47	Slow Arm Damping Rate versus Effective Squared Yaw	54

Figure

vi

List of Tables

 \mathcal{O}

<u>Table</u>		Page
1	Average Physical Characteristics of Caliber .50 Projectiles	55
2	Aerodynamic Characteristics of the Ball, M33 Projectile	56
3	Aerodynamic Characteristics of the API, M8 Projectile	57
4	Aerodynamic Characteristics of the APIT, M20 (Burnt Tracer)	58
5	Flight Motion Parameters of the Ball, M33 Projectile	59
6	Flight Motion Parameters of the API, M8 Projectile	60
7	Flight Motion Parameters of the APIT, M20 Projectile (Burnt Tracer) :	61
8	Tracer-On Drag Measurements for the APIT, M20 Projectile	62

vii

viii

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I. Introduction

The caliber .50 Armor Piercing Incendiary (API, M8) and Armor Piercing Incendiary Tracer (APIT, M20) munitions were developed in 1943-44, for wartime service use in various versions of the caliber .50, M2, Browning Machine Gun. Reference 1 contains a summary of the limited aerodynamic data obtained during development testing of the new munitions. The drag coefficient was determined from resistance firings over solenoid velocity screens, and the stability was measured using yaw card techniques.

The caliber .50 Ball, M33 round was developed in 1961, as a companion ball munition to the API, M8, and was intended to be a ballistic match of the M8. Apparently, no aerodynamic tests were ever conducted for the M33 projectile. Some unpublished aerodynamic data for the APIT, M20 were obtained by M. J. Piddington in 1979, in support of the M1 Abrams tank development program. Piddington's spark photography range data are included in this report.

In November 1987, the Fire Control division of the U.S. Army Armament Research, Development and Engineering Center (ARDEC) requested that the Ballistic Research Laboratory (BRL) provide trajectory data for a fire control study involving current 7.62mm and caliber .50 infantry weapons. The BRL advised ARDEC that the existing aeroballistic data base for the caliber .50 munitions was insufficient to permit accurate trajectory predictions, and recommended that testing be conducted in the BRL spark photography ranges.²³ In early April 1988, test material and funding for the BRL spark range firings of caliber .50 munitions were received from ARDEC.

Final plans for the spark photography range tests were nearing completion when the Air Force Armament Laboratory at Eglin Air Force Base, Florida, requested that the BRL conduct large-yaw firings of the caliber .50, API, M8 projectile in the Free Flight Aerodynamics Range² to provide aeroballistic data for side-fire from high speed aircraft. By mutual agreement between the BRL, ARDEC, and the Air Force; the two aeroballistic tests were combined, and the Air Force supplied additional funding to the BRL for the large-yaw firings. This report presents all the modern aeroballistic data collected in the two BRL spark photography ranges, for the caliber .50 Ball, M33, API, M8 and APIT, M20 munitions.

II. Test Procedure and Material

Figure 1 is a photograph of the three caliber .50 projectiles. Figure 2 is a photograph of the BRL Aerodynamics Range (circa 1958), and Figure 3 illustrates the local and master coordinate systems for the range.

Physical measurements were taken on a sample of five projectiles of each type. The average physical properties of the test projectiles are listed in Table 1. The Ball, M33. API, M8. and APIT, M20 designs all have the same nominal external dimensions, and differ only in minor surface details, such as rolled versus machined cannelures. Figure 4 is a sketch of the exterior contour, common to all three projectiles.

All test rounds were fired from a 114.3 cm (45 inch) caliber .50 Mann barrel, with a uniform rifling twist rate of one turn in 38.1 cm (15 inches). Suitable propellants and charges were selected to achieve test velocities varying from 915 metres/second down to 240 metres/second. For the large-yaw firings of the API, M8, a half-muzzle type yaw inducer was used, with a lip length ranging from 6.35mm(1/4 inch) to 12.7mm(1/2 inch). Average yaw levels exceeding 12 degrees were obtained for several test rounds with the 12.7mm lip yaw inducer.

Live tracer firings of small caliber projectiles present a problem for the BRL Aerodynamics Range, because the light emitted from the tracer tends to fog the film. The APIT, M20 live tracer firings were conducted in the BRL Transonic Range, ³ with the gun backed off approximately 100 metres from the range entrance, to insure a fully burning tracer over the instrumentation. Transonic Range shadowgraphs of typical small caliber projectiles do not permit accurate measurement of yawing or swerving motion, sr only drag is obtained from the live tracer firings.

Tracer-off firings of the APIT, M20 in the BRL Aerodynamics Range were conducted by pulling the projectiles, burning out the tracer mix, then firing the burned-out tracer round. All the Ball, M33 and API, MS test firings were conducted in the Aerodynamics Range.

An interesting and useful by-product of spark photography range testing is the high quality flowfield visualization provided by the spark shadowgraphs. Figures 5 through 18 show the flowfields around the three caliber .50 bullets at various supersonic, transonic, and subsonic speeds. Most of the shadowgraphs were selected from range stations where the angle of attack was less than one degree; Figure 7 illustrates the effect of large angle of attack on the flow past the API, M8 projectile.

The round-by-round aerodynamic data obtained for the three caliber .50 bullets are listed in Tables 2, 3 and 4. Free flight motion parameters for the three bullets are listed in Tables 5, 6 and 7. The tracer-on drag data obtained in the BRL Transonic Range for the APIT, M20 projectile is listed in Table 8.

III. Results

The free flight spark range data were fitted to solutions of the linearized equations of motion and the resulting flight motion $p_{\rm F}$ ameters were used to infer linearized aerodynamic coefficients, using the methods of Reference 4. Preliminary analysis of the aerodynamic data showed distinct variation of several coefficients with yaw level. In BRL Report 974, Murphy ⁵ has shown that aerodynamic coefficients derived from the linearized data reduction can be used to infer the coefficients in a nonlinear force and moment expansion, if sufficient data are available. For the caliber .50 bullets, sufficient data were obtained to permit determination of several nonlinear coefficients. A more detailed analysis of nonlinear effects is presented in the subtopics of this section, which discuss individual aerodynamic coefficients.

1. Drag Coefficient

The drag coefficient, C_D , is determined by fitting the time-distance measurements from the range flight. C_D is distinctly nonlinear with yaw level, and the value determined from an individual flight reflects both the zero-yaw drag coefficient, C_{D_0} , and the induced drag due to the average yaw level of the flight. The drag coefficient variation is expressed as an even power series in yaw amplitude:

$$C_D = C_{D_0} + C_{D_{\delta^2}} \delta^2 + \dots \tag{1}$$

where C_{D_0} is the zero-yaw drag coefficient, $C_{D_{\delta^2}}$ is the quadratic yaw-drag coefficient, and δ^2 is the total angle of attack squared.

Preliminary analysis of the drag coefficient data for the three caliber .50 projectiles showed that the zero-yaw drag coefficients were, for practical purposes, identical. The drag data for all three round types were therefore combined, and a single yaw-drag curve was determined. No significant variation of the yaw-drag coefficient with projectile type could be found, and the yaw-drag curve shown in Figure 22 was used to correct all the range values of C_D to zero-yaw conditions.

Figures 19 through 21 illustrate the variation of C_{D_0} with Mach number for the three caliber .50 bullets. The zero-yaw drag coefficients for the Ball, M33; API, M8; and APIT, M20 (Tracer off) are essentially identical at all speeds tested. The round-to-round standard deviation in zero-yaw drag coefficient is 1.3 percent at supersonic speeds, for all bullet types.

Figure 21 also illustrates the effect of the burning tracer on the zero-yaw drag coefficient of the M20 projectile. The tracer adds heat and mass flux into the wake, which raises the base pressure and lowers the base drag. For the APIT, M20 projectile, the tracer reduces the total zero-yaw drag coefficient by approximately 7 percent, at all speeds tested.

Figure 23 is a comparison of the API, M8 drag coefficient obtained by H. P. Hitchcock¹ with the current Aerodynamics Range test results for the same projectile. Hitchcock's curve was converted from the old K_D to the modern C_D nomenclature [$C_D = (8/\pi) K_D$], and was also corrected for the difference in reference diameter (Hitchcock used 0.50 inch, and the present tests use 0.51 inch). Hitchcock's drag coefficient averages about 4 percent lower than the spark range curve at supersonic speeds, and about 10 percent higher at transonic and subsonic speeds. Considering the relatively crude instrumentation used in the 1943 resistance firings, the agreement is satisfactory.

2. Oversurning Moment Coefficient

The range values of the overturning moment coefficient, $C_{M_{\alpha}}$, were fitted using the appropriate squared-yaw parameters from Reference 5. A weak dependence of $C_{M_{\alpha}}$ on yaw level was observed for the caliber .50 projectiles. The overturning moment is assumed to be cubic in yaw level, and the coefficient variation is given by:

$$C_{\mathcal{M}_{\alpha}} = C_{\mathcal{M}_{\alpha 0}} + C_2 \,\delta^2 + \dots \tag{2}$$

where $C_{M_{ab}}$ is the zero-yaw overturning moment coefficient, and C_2 is the cubic coefficient.

Figure 27 illustrates the observed variation of C_2 with Mach number, and this curve was used to correct all the range values of $C_{M_{\alpha}}$ to zero-yaw conditions. Figures 24 through 26 show the variation of $C_{M_{\alpha_0}}$ with Mach number for the three caliber .50 projectiles. The Ball, M33 has the highest overturning moment coefficient of the three bullets; $C_{M_{\alpha_0}}$ for the API, M8 averages about 2 percent lower than that of the Ball, M33 and the APIT, M20 curve is approximately 10 percent lower than the Ball, M33 curve.

3. Gyroscopic Stability

The launch gyroscopic stability factors of the three caliber .50 bullets, fired from a barrel with 15 inch twist of rifling, at a muzzle velocity of 2950 feet/second, into a sea-level ICAO standard atmosphere, are as follows:

Projectile	Launch Gyroscopic St	ability Factor
Ball, M33	. 1.8	
API, M8	1.9	·
APIT, M20	2 2	

A gyroscopic stability factor above 1.5 is usually considered adequate, so all the caliber .50 projectiles tested have sufficient gyroscopic stability to permit satisfactory flight in all expected conditions, including extreme cold weather (high air density) conditions. Since the caliber .50 ammunition is never fired at reduced muzzle velocities, the lower values of S_q observed in Tables 5 through 7 will never occur in field firings.

4. Lift Force Coefficient

The range values of the lift force coefficient, $C_{L_{\alpha}}$, were also analyzed using the methods of Reference 5. A weak dependence of $C_{L_{\alpha}}$ on yaw level was observed for the caliber .50 projectiles. The variation of $C_{L_{\alpha}}$ with yaw level is also assumed to be cubic:

$$C_{L_{\alpha}} = C_{L_{\alpha 0}} + a_2 \delta^2 + \dots$$
(3)

where $C_{L_{22}}$ is the zero-yaw lift force coefficient, and a_2 is the cubic coefficient.

Figure 31 illustrates the variation of the cubic lift force coefficient with Mach number.¹ The subsonic and supersonic regions showed distinctly different levels of cubic behavior, and the curve of Figure 31 was used to correct all range values of $C_{L_{\alpha}}$ to zero-yaw conditions.

The variation of $C_{L_{\alpha0}}$ with Mach number for the three caliber .50 bullets is shown in Figures 28 through 30. The zero-yaw lift force coefficients of the three projectiles are essentially identical.

5. Magnus Moment Coefficient and Pitch Damping Moment Coefficient

The Magnus moment coefficient, $C_{M_{p_{\alpha}}}$, and the pitch damping moment coefficient sum $(C_{M_q} + C_{M_{\alpha}})$, are discussed together, since if either coefficient is nonlinear with yaw level, both coefficients exhibit nonlinear coupling in the data reduction process.⁵ Due to mutual reaction, the analysis of $C_{M_{p_{\alpha}}}$ and $(C_{M_q} + C_{M_{\alpha}})$ must be performed simultaneously, even though the aerodynamic moments are not, in themselves, directly physically related.

If the dependence of the Magnus moment and the pitch damping moment are cubic in yaw level, the nonlinear variation of the two moment coefficients is of the general form:

$$C_{M_{p_{\alpha}}} = C_{M_{p_{\alpha_{0}}}} + \hat{C}_{2} \delta^{2}$$
 (4)

$$\left(C_{M_q} + C_{M_{\dot{\alpha}}}\right) = \left(C_{M_q} + C_{M_{\dot{\alpha}}}\right)_0 + d_2 \,\delta^2 \tag{5}$$

where $C_{M_{p_{\alpha_0}}}$ and $(C_{M_q} + C_{M_{\dot{\alpha}}})_0$ are the zero-yaw values of Magnus and pitch damping moment coefficients, respectively, and \hat{C}_2 and d_2 are the associated cubic coefficients.

In Reference 5, it is shown that the nonlinear coupling introduced through the least squares fitting process yields the following expressions for range values [*R*-subscript] of $C_{M_{p_{\alpha}}}$ and $(C_{M_{q}} + C_{M_{\alpha}})$:

5

$$\left[C_{M_{p_{\alpha}}}\right]_{R} = C_{M_{p_{\alpha_{0}}}} + C_{\varepsilon} \delta_{\varepsilon}^{2} + d_{2} \delta_{\varepsilon_{TH}}^{2}$$
(6)

$$\left[\left(C_{M_q} + C_{M_{\alpha}}\right)\right]_R = \left(C_{M_q} + C_M\right)_0 + \hat{C}_2 \,\delta^2_{\ell_{HH}} + d_2 \,\delta^2_{\ell_{HT}} \tag{7}$$

where the above effective squared yaws are defined as:

$$\delta_{c}^{2} = K_{F}^{2} + K_{S}^{2} + \frac{(\phi_{F}^{\prime} K_{F}^{2} - \phi_{S}^{\prime} K_{S}^{2})}{(\phi_{F}^{\prime} - \phi_{S}^{\prime})}$$
(8)

$$\delta_{e_{TH}}^{2} = \left(\frac{I_{x}}{I_{y}}\right) \left[\frac{(K_{F}^{2} \phi_{F}^{\prime 2} - K_{S}^{2} \phi_{S}^{\prime 2})}{(\phi_{F}^{\prime 2} - \phi_{S}^{\prime 2})}\right]$$
(9)

$$\delta_{e_{HH}}^{2} = \left(\frac{I_{y}}{I_{x}}\right) \left[\frac{(\phi_{F}' + \phi_{S}') (K_{S}^{2} - K_{F}^{2})}{(\phi_{F}' - \phi_{S}')}\right]$$
(10)

$$\delta_{e_{HT}}^{2} = \frac{(\phi_{F}' K_{S}^{2} - \phi_{S}' K_{F}^{2})}{(\phi_{F}' - \phi_{S}')}$$
(11)

The remaining symbols are defined in the List of Symbols in this report.

Preliminary analysis of the caliber .50 data indicated strong nonlinearity in the range values of $C_{M_{pa}}$ and $(C_{M_q} + C_{M_a})$ at angles of attack less than 3 degrees, but essentially no variation of either coefficient was observed at larger yaw levels. The data rounds were separated into Mach number groups, and an analysis was performed to determine the cubic coefficients at both small and large yaw levels. No significant values of the cubic pitch damping moment coefficient, d_2 , could be found.

Figures 32 through 34 illustrate the variation of the range values of $C_{M_{pa}}$ with the appropriate squared yaw parameter from Reference 5. The general characteristic of these plots is bi-cubic behavior, with strong nonlinearity at small yaw levels, followed by no significant variation of $C_{M_{pa}}$ with y_k level at larger yaws. The small-yaw cubic Magnus moment coefficient varies significantly with Mach number, but the large-yaw $C_{M_{pa}}$ appears to be essentially independent of Mach number at supersonic speeds.

Least squares fitting of the Magnus moment coefficient data yielded the curve of Figure 38 for \hat{C}_2 , which was then used to correct the range values of $C_{M_{po}}$ and $C_{M_q} + C_{M_o}$ to zero-yaw conditions. Figures 35 through 37 show the variation of $C_{M_{poo}}$ with Mach number for the three caliber .50 bullets, and Figures 39 through 41 illustrate the variation of $(C_{M_q} + C_{M_o})_o$ with Mach number.

It should be noted that the analysis of nonlinear Magnus and pitch damping moment data from free flight spark ranges is a delicate process at best, and the results are highly sensitive to small errors in determination of the damping exponents on the two modal arms. The uncertainties in damping rate determinations are reflected in the larger roundto-round data scatter in Magnus and pitch damping moment coefficients, compared with the smaller scatter observed in the overturning moment coefficients.

6. Damping Rates

The damping rates, λ_F and λ_S , of the fast and slow yaw modes indicate the dynamic stability of a projectile. Negative λ 's indicate damping; a positive λ means that its associated modal arm will grow with increasing distance along the trajectory.

For a projectile whose Magnus or pitch damping moments are nonlinear with yaw level, the damping rates also show a nonlinear dependence on yaw.⁶ Figures 42 through 47 illustrate the variations in damping rates with yaw level for the three caliber .50 bullets at supersonic and subsonic speeds. Figures 42 and 43 show that both modal arms are

damped at all yaw levels tested, for high supersonic speeds. At low supersonic speeds. Figures 44 and 45 shows a damped fast arm at all yaw levels tested, but the slow arm is unstable at small yaw and stable at larger yaw levels. This behavior is described by the aeroballistician as limit-cycle yaw, and Figure 45 predicts a slow arm limit-cycle yaw of about 3 degrees at low supersonic and transonic speeds.

Figures 46 and 47 show the variation of the fast and slow arm damping rates with yaw level at subsonic speeds, for the three caliber .50 bullets. The fast arm shows generally satisfactory damping for the Ball, M33 and API, M8 bullets at subsonic speeds, with very weak undamping of the fast arm observed for the APIT, M20 design. Figure 47 indicates a slow arm limit-cycle for all bullet types at subsonic speeds. The magnitude of the expected slow arm limit-cycle yaw at subsonic speeds is slightly greater than 4 degrees, for all three caliber .50 bullets.

IV. Conclusions

The drag coefficients of the caliber .50 Ball, M33, API, M8 and tracer-off APIT, M20 projectiles are essentially identical at all speeds tested. The effect of the burning tracer on the APIT, M20 is to reduce the tracer-off drag approximately 7 percent. The observed round-to-round standard deviation in drag coefficient is 1.3 percent at supersonic speeds. for all three bullet types.

The three caliber .50 bullets tested have sufficient gyroscopic stability to permit satisfactory flight at all conditions, including extreme cold weather (high air density), when fired at service velocity from a barrel with standard 381mm (15 inch) twist of rifling.

The non-linear Magnus moment characteristics of the three caliber .50 bullets predict a slow arm limit cycle yaw of approximately 3 degrees at low supersonic and transonic speeds. The low speed Magnus moment behavior predicts a slow arm limit cycle yaw slightly exceeding 4 degrees at subsonic speeds.

V. Recommendations

A long range limit cycle yaw test should be conducted with the caliber .50 service ammunition to verify the flight dynamic predictions made in this report.

A radar doppler velocimeter test should be conducted for the caliber .50 projectiles to verify the drag coefficient results out to long ranges, and very low subsonic speeds.



























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Figure 32. Magnus Moment Coefficient versus Effective Squared Yaw.

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Figure 46. Fast Arm Damping Rate versus Effective Squared Yaw.



Projectile	Reference Diameter	Weight	Center of Gravity	Axial Moment of Inertia	Transverse Moment of Inertia
	(mm)	(grams)	(cal - base)	(gm-¢m ²)	(gm-cm ²)
Ball, M33	12.95	42.02	1.78	7.85	74.5
API, M8	12.95	41.98	1.79	7.84	73.9
APIT, M20 (Live Tracer)	12.95	39.77	1.84	7.79	68.5
APIT, M20 (Burnt Tracer)	12.95	38.95	1.88	7.77.	66.7

 Table 1. Average Physical Characteristics of Caliber .50 Projectiles.

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Round Number	Mach Number	α_t (degrees)	C_{D}	С _{Ма}	$C_{L_{\alpha}}$	$C_{M_{ra}}$	$(C_{M_q} + C_{M_{\dot{\alpha}}})$	CP_N (cal - base)
18892	2.653	1.63	.2813	3.01	2.21	.15	-7.4	2.99
18924	2.589	2.54	.2971	3.01	2.41	.15	-6.4	2.90
18891	2.570	1.98	. 2923	3.07	2.20	.15	-6.2	3.02
18895	1.972	1.67	.3171	3.40	2.29	.12	-8.8	3.09
18896	1.953	3.54	.3410	3.39	2.43	. 20	-8.4	3.01
18899	1.516	1.59	. 3489	3.69		10	-6.6	
18898	1.475	2.79	.3637	3.70	2.06	.10	-7.1	3.31
18901	1.222	1.84	.3725	3.82	1.89	40	-3.9	3.47
18902	1.158	2.26	.3757	3.85	1.79	10	-4.6	3.56
18907	1.041	1.59	.3569	4.04	1.46	73	0.8	4.01
18908	1.003	2.97	.3461	4.19	1.65	10	-7.3	3.88
18906	. 989	3.08	.3215	4.24	1.45	41	1.0	4.17
18909	.918	3.15	.1505	4.39		.74	-12.6	
18936	.888	2.87	.1372	4.42	1.44	52	3.7	4.58
18911	.606	2.87	.1230	3.71		62	3.0	
18912	.551	3.30	.1324	3.54	1.67	23	-7.9	3.74

Table 2. Aerodynamic Characteristics of the Ball, M33 Projectile.

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Round Number	Mach Number	a _t (degrees)	C _D	С _{Мо}	C _L	$C_{M_{pa}}$	$\begin{array}{c} (C_{M_q} \\ + C_{M_{\dot{\alpha}}}) \end{array}$	CP _N (cal - base)
18857	2.686	1.39	. 2938	2.88	2.69	.05	-5.5	2.75
18922	2.669	4.62	.3200	2.81	2.48	.26	-7.3	2.79
18856	2.639	2.54	. 2991	2.85	2.45	.24	-7.8	2.82
18918	2.628	12.81	.4985	2.65	3.26	.13	-6.2	2.49
18923	2.605	2.08	.2950	2.87	2.38	. 22	-7.5	2.86
18920	2.600	1.75	. 2956	2.92	2.43	. 20	-8.0	2.86
18917	2.511	12.58	.5189	2.76	3.30	. 28	-7.9	2.51
18915	2.508	13.57	. 5683	2.82	3.40	.41	-7.4	2.50
18859	2.038	.60	.3108					
18860	1.926	1.63	.3260	3.33	2.33	02	-8.4	3.04
.18866	1.500	1.69	. 3507	3.60	2.35	14	-5.9	3.12
18867	1.496	1.86	.3525	3.65	2.10	07	-6.1	3.28
18929	1.198	7.41	.4542	3.69	1.97	. 27	-7.8	3.31
18930	1.197	7.34	.4507	3.67	1.98	. 26	-7.5	3.30
18926	1.178	5.00	.4100	3.69	1.79	. 24	-7.7	3.46
18875	1.158	4.48	. 3984	3.78	1.65	.31	-9.0	3.63
18874	1.109	1.84	.3763	3.84	1.63	68	2.7	3.70
18878	.976	2.74	. 2388	4.77	1.55	.61	-18.2	4.45
18879	. 959	2.55	.1624	4.88		.13	-8.2	
18932	.939	9.63	. 2331	4.00	1.50	.02	2.7	4.10
18933	.897	4.39	. 1454	4.16	1.30	.41	-8.5	4.67
18935	. 892	2.87	.1328	4.28	1.62	.20	-7.7	4.23
18881	.799	4.07	. 1407.	3.86	1.40	. 02	-4.0	4.29
18934	.692	1.75	.1232	3.80				

Table 3. Aerodynamic Characteristics of the API, M8 Projectile.

Table 4. Aerodynamic Characteristics of the APIT, M20 (Burnt Tracer).

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Round	Mach	àr	C_D	$C_{M_{\alpha}}$	$C_{L_{\alpha}}$	$C_{M_{p'a}}$	$(C_{M_{\bullet}})$	CP_N
Number	Number	(degrees)				·	$+C_{M_{\dot{\alpha}}})$	(cal - base)
13550 ·	2.309	1.59	.3013	2.86	2.23	.01	-8.8	3.02
13551	1.965	1.33	.3250	3.00	2.52	33	-8.0	2.94
13552	1.958	1.90	.3261	2.98	2.25	11	-4.8	3.04
13549	1.855	1.04	.3182	3.04	2.68	25	-6.5	2.90
13553	1.420	1.29	.3589	3.34	1.67			3.53
13554	1.362	2.13	.3729	3.34	2.10	07	-6.3	3.23
13555	1.134	1.78	.3743	3.47	1.54			3.70
13556	1.075	2.37	.3790	3.47	1.55	46	.7	3.68
13557	.941	1.50	.1715	4.11		'		
13559	.819	2.55	.1260	3.87	1.45	37	4.2	4.34
13560	.748	2.07	.1298		 .			
13561	.710	2.41	.1294	3.66		62	7.8	

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Table 5.	Flight Motion Parameters of the Ball, M33 Projectile.	
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Round	Mach	S_{g}	Sd	$\lambda_F imes 10^3$	$\lambda_S imes 10^3$	K_F	K_{S}	ϕ'_F	ϕ'_{S}	Spin
Number	Number	•	•	(1/cal)	(1/cal)			(r/cal)	(r/cal)	(r/cal)
18892	2.653	1.81	.8	147	075	.0191	.0204	.0187	.0037	. 213
. 18924	2.589	1.83	.9	120	089	.0297	.0318	.0188	.0037	.213
18891	2.570	1.80	.9	114	084	.0234	.0245	.0187	.0038	.213
18895	1.972	1.61	.6	210	054	.0177	.0222	.0182	.0043	. 214
18896	1.953	1.61	.8	165	092	.0403	.0445	.0182	.0044	. 214
18899	1.516	1.49	. 2	231	.037	.0132	.0238	.0178	.0048	.214
18898	1.475	1.48	.7	173	053	.0310	.0361	.0178	.0049	.214
18901	1.222	1.42	6	271	.126	.0127	.0287	.0174	.0051	.214
· 18902	1.158	1.42	. 3	178	.019	.0214	.0327	.0174 [.]	.0052	. 214
18907	1.041	1.34	-18.1	281	. 250	.0054	.0257	.0169	.0056	.213
18908	1.003	1.34	. 2	279	.073	.0263	.0430	.0172	.0057	.217
18906	.989	1.31	-10.0	134	.124	.0266	.0458	.0169	.0059	.216
13909	.918	1.18	1.2	087	232	.0498	.0203	.0153	.0067	. 209
18936	.888	1.26		099	. 140	.0209	.0449	.0165	.0062	.215
. 18911	.60€	1.43		124	. 183	.0243	.0434	.0173	.0050	.211
18912	.551	1.47	1	315	.077	.0352	.0445	.0172	.0048	.209

Round Number	Mach Numbe	S _g r	S_d	$\lambda_F imes 10^3 \ (1/{ m cal})$	$\lambda_S imes 10^3 \ (1/ ext{cal})$	K_F	K _s	ϕ_F^{i} (r/cal)	ϕ_S' (r/cal)	.Spin (r/cal)
18857	2.686	1.92	.8	128	071	.0157	.0179	.0191	.0035	.213
18922	2.669	1.93	1.0	107	121	.0576	.0538	.0191	.0034	.212
18856	2.639	1.93	.9	131	109	.0308	.0304	.0191	.0035	.212
18918	2.628	2.08	1.0	112	113	.1626	.1468	.0196	.0032	.214
18923	2.605	1.89	. 9	129	102	.0239	.0263	.0190	.0035	.212
18920	2.600	1.85	.8	150	096	.0200	.0222	.0188	.0036	.212
18917	2.511	2.04	1.1	113	148	.1590	.1442	.0197	.0033	. 217
18915	2.508	1.99	1.4	108	124	.1729	.1564	.0196	0034	.217
18859	2.038					.0038	.0095			
18860	1.926	1.67	.4	244	007	.0136	.0244	.0185	.0042	.213
18866	1.500	1.50	.3	215	.018	.0156	.0243	.0178	. 2047 -	.212
18867	1.496	1.49	.4	201	.003	.0i81	.0261	.0178	iu∩48	.213
18929	1.198	1.55	. 9	126	105	.0882	.0902	.0183	.0647	.217
18930	1.197	1.54	. 9	124	103	.0863	.0907	.0183	.047	. 216
18926	1.178	1.50	.9	138	088	.0576	.0629	.0179	.0048	.214
18875	1.158	1.45	.9	153	101	.0517	. 0559	.0177	.0050	.214
18874	1.109	1.42		168	.184	.0108	.0294	.0175	.0051	.213
18878	.976	1.20	.7	352	037	.0219	.0404	.0162	.0069	.218
18879	. 959	1.07	. 5	318	.063	.0192	.0389	.0140	.0082	. 209
18932	. 939	1.33		.119	- 115	. 1499	.0703	.0166	.0056	. 209
18933	.897	1.29	1.1	094	134	.0595	.0457	.0164	.0059	.210
18935	.692	1.24	. 8	163	038	.0370	.0318	.0159	.0062	. 209
18881	.799	1.39	.6	110	015	.0454	.0537	.0171	.0053	. 211
18934	. 692	1.44				.0079	.0277	.0175	.0050	.212

Table 6. Flight Motion Parameters of the API, M8 Projectile.

Table 7. Flight Motion Parameters of the APIT, M20 Projectile (Burnt Tracer).

Round	Mach	S_{g}	S_d	$\lambda_F imes 10^3$	$\lambda_S imes 10^3$	K_F	Ks	ϕ'_F	ϕ'_{S}	Spin
Number	Number	r -		(1/cal)	(1/cal)			(r/cal)	(r/cal)	(r/cal)
13550	2.309	2.09	.4	252	- 033	.0139	.0226	.0213	0034	213
13551	1.965	2.02	1	330	.062	.0066	.0210	.0213	.0036	.214
13552	1.958	2.01	.4	174	015	.0168	.0272	.0212	.0036	.213
13549	1.855	1.94	.1	270	.024	.0073	.0156	.0208	. 0037	.211
13553	1.420	1.83	-1.1	351	. 172	. 0039	.0205	.0207	0040	. 213
13554	1.362	1.81	.4	210	012 ·	.0168	.0315	.0208	.0041	. 214
13555	1.134	1.74				.0044	.0278	.0206	.0043	. 214
13556	1.075	1.75	-5.9	124	.096	.0166	.0365	.0207	.0043	.215
13557	.941	1.36				.0241	.0049	.0182	.0058	. 206
13559	.819	1.48		.019	.048	.0254	.0362	.0190	.0052	. 209
13560	.748				÷-	.0037	.0335			
13561	.710	1.57		.050	. 130	.0179	.0372	.0195	.0049	. 209

	Round Number	Mach Number	α_t (degrees)	$C_{D_{(R)}}$	C_{D_0}	
	30461	2.502	.31	. 2748	.274	
	30460	2.497	. 27	. 2656	. 265	
•	30467	2.478	.63	.2741	.274	
	17089	2.430	*	.2759	.275	

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1.79

1.00

1.05

2.30

1.63

1.28

.3029

.3254

.3293

.3193

.3031

.3028

.2263

.1596

.1603

.1439

.302

.325

.329

.319

.304

.302

.225

.152

.158

.142

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Round Number	Mach Number	a _t (degrees)	$C_{D_{(R)}}$	C_{D_0}	

1.882

1.533

1.528

1.525

1.015

1.007

.983

.973

.967

.966

Table 8. Tracer-On Drag Measurements for the APIT, M20 Projectile.

Notes:	* Very small yaw ($\alpha_t < .5$ degree)
	() Denotes split reduction

30462

17159

17158

17160

-3

30475(a)

30474(a)

30474(Ъ)

30475(Ъ)

30474(c)

30473

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List of Symbols

a ₂	=	cubic lift force coefficient	
<i>C</i> ²	=	cubic static moment coefficient	
Ĉ2	=	cubic Magnus moment coeffi- cient	
CD	=	$\frac{\text{Drag Force}}{[(1/2)\rho V^2 S]}$	· · ·
C_{D_0}	=	zero-yaw drag coefficient	
$C_{D_{\delta}}$,	=	quadratic yaw drag coefficient	
C _{La}	=	$\frac{\text{Lift Force}}{[(1/2)\rho V^2 S \delta]}$	Positive coefficient: Force in plane of total angle of attack, α_t , \perp to trajectory in direc- tion of α_t . (α_t directed from trajectory to missile axis.) $\delta = \sin \alpha_t$.
$C_{N_{\alpha}}$	=	$\frac{\text{Normal Force}}{[(1/2)\rho V^2 S \delta]}$	Positive coefficient: Force in plane of total angle of attack, α_t , \perp to missile axis in direc- tion of α_t . $C_{N_{\alpha}} \cong C_{L_{\alpha}} + C_D$
С _{Ма}	=	$\frac{\text{Static Moment}}{[(1/2)\rho V^2 S d\delta]}$	Positive coefficient: Moment increases angle of attack α_t .
(' _{Mpa}	E	$\frac{\text{Magnus Moment}}{[(1/2)\rho V^2 S d (p d/V)]}$	Positive coefficient: Moment rotates nose \perp to plane of α_t in direction of spin.
C _{Npa}	-	$\frac{\text{Magnus Force}}{[(1/2)\rho V^2 S (p d/V) \delta]}$	Negative coefficient: Force acts in direction of 90° rota- tion of the positive lift force against spin.

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List of Symbols (Continued)

For most exterior ballistic uses, where $\dot{\alpha} \approx q$, $\dot{\beta} \approx -r$, the definition of the damping moment sum is equivalent to:

$(C_{M_q}+C_{M_\alpha})$	=	$\frac{\text{Damping Moment}}{\left[\left(1/2\right)\rho V^2 S d \left(q_t d/V\right)\right]}$	Positive coefficient: Moment increases angular velocity.
C_{l_p}	=	$\frac{\text{Roll Damping Moment}}{[(1/2)\rho V^2 S d (p d/V)]}$	Negative coefficient: Moment decreases rotational velocity.
CPN		center of pressure of the nor- mal force, positive from base	
		to nose	· · ·
α, β	=	angle of attack, side slip	•
α_t	-	$(\alpha^2 + \beta^2)^{\frac{1}{2}} = \sin^{-1} \delta,$ total angle of attack	
λ_F	=	fast mode damping rate	negative λ indicates damping
λ_{S}	. =	slow mode damping rate	negative λ indicates damping
ρ	=	air density	
ϕ'_F	=	fast mode frequency	
ϕ'_{S}	=	slow mode frequency	•
<i>c.m</i> .	=	center of mass	
d	=	body diameter of projectile,	
		reference length	
d2	=	reference length cubic pitch damping moment coefficient	

List of Symbols (Continued)

Iy	=	transverse moment of inertia
K _F	=	magnitude of the fast yaw mode
Ks	=	magnitude of the slow yaw mode
l	=	length of projectile
m	=	mass of projectile
Μ	=	Mach number
p	=	roll rate
q , r	Ξ	transverse angular velocities
q.	=	$(q^2 + r^2)^{\frac{1}{2}}$
R	=	subscript denotes range value
\$	=	dimensionless arc length along the trajectory
S	=	$(\pi c'^2/4)$, reference area
Sd	=	dynamic stability factor
Sg	= '	gyroscopic stability factor
V	=	velocity of projectile

List of Symbols (Continued)

Effective Squared Yaw Parameters

$$\begin{split} \bar{\delta} &\cong K_{F}^{2} + K_{S}^{2} \\ \delta_{e}^{2} &= K_{F}^{2} + K_{S}^{2} + \frac{(\phi_{F}' K_{F}^{2} - \phi_{S}' K_{S}^{2})}{(\phi_{F}' - \phi_{S}')} \\ \delta_{e_{F}}^{2} &= K_{F}^{2} + 2K_{S}^{2} \\ \delta_{e_{S}}^{2} &= 2K_{F}^{2} + K_{S}^{2} \\ \delta_{e_{HH}}^{2} &= \left(\frac{I_{y}}{I_{x}}\right) \left[\frac{(\phi_{F}' + \phi_{S}') (K_{S}^{2} - K_{F}^{2})}{(\phi_{F}' - \phi_{S}')}\right] \\ \delta_{e_{HH}}^{2} &= \left(\frac{I_{x}}{I_{y}}\right) \left[\frac{(K_{F}^{2} \phi_{F}'^{2} - K_{S}^{2} \phi_{S}'^{2})}{(\phi_{F}'^{2} - \phi_{S}')}\right] \\ \delta_{e_{HT}}^{2} &= \frac{(\phi_{F}' K_{S}^{2} - \phi_{S}' K_{F}^{2})}{(\phi_{F}' - \phi_{S}')} \end{split}$$

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