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Yung K. Choo
Lewis Research Center
Cleveland, Ohio

Seokkwan Yoon
Sverdrup Technology, Inc.
Lewis Research Center
Cleveland, Ohio

and

Kestutis C. Civinskas
Propulsion Directorate
U.S. Army Aviation Research and Technology Activity—AVSCOM
Lewis Research Center
Cleveland, Ohio

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COMPOSITE GRID AND FINITE-VOLUME LU IMPLICIT SCHEME

FOR TURBINE FLOW ANALYSIS

Yung K. Choo*
National Aeronautics and Space Administration
Lewis Research Center
Cleveland, Ohio 44135

Seokkwan Yoon**
Sverdrup Technology, Inc.
Lewis Research Center
Cleveland, Ohio 44135

and

Kestutis C. Civinskas†
Propulsion Directorate
U.S. Army Aviation Research and Technology Activity - AVSCOM
Lewis Research Center
Cleveland, Ohio 44135

Abstract

A composite grid was generated in an attempt to improve grid quality for a typical turbine blade with large camber in terms of mesh control, smoothness, and orthogonality. This composite grid consists of the C-grid (or O-grid) in the immediate vicinity of the blade and the H-grid in the upstream region and in the middle of the blade passage between the C-grids. It provides a good boundary layer resolution around the leading-edge region for viscous calculation, has orthogonality at the blade surface and slope continuity at the C-H (or O-H) interface, and has flexibility in controlling the mesh distribution in the upstream region without using excessive grid points. This composite grid eliminates the undesirable qualities of a single grid when generated for a typical turbine geometry.

A finite-volume lower-upper (LU) implicit scheme can be used in solving for the turbine flows on the composite grid. This grid has a special grid node that is connected to more than four neighboring nodes in two dimensions and to more than six nodes in three dimensions. But the finite-volume approach poses no problem at the special point because each interior cell has only four neighboring cells in two dimensions and only six cells in three dimensions. The finite-volume LU implicit scheme was demonstrated to be robust and efficient for both external and internal flows in a broad flow regime.

Introduction

Turbine blades are often designed to have substantial thickness and camber and rounded leading and trailing edges. For a typical turbine blade the H-grid does not provide a good boundary layer resolution around the leading edge. If very fine meshes are used to resolve this, the number of grid points in the upstream region becomes excessive. Thus mesh points would be wasted. A turbine flow

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*Research Scientist, Computational Applications Branch.
**Senior Research Engineer; Member AIAA.
†Aerospace Research Engineer, Computational Applications Branch.

analysis on an H-grid indicated that the standard H-grid needs cusps at the leading and trailing edges to reach a convergent solution and that it causes somewhat large entropy changes around the leading edge, where the H-grid has a large grid skewness and slope discontinuity.[1] These numerical errors around leading edges would be reduced if a C-grid were used because it provides a better resolution around the leading edge. A standard C-grid was generated for a core turbine stator vane.[2] It reveals large grid skewing on the suction side of the flow passage, which should increase numerical errors.[3,4] For a typical turbine blade, no single grid offers satisfactory grid properties in the entire turbine stator or rotor passage.

A two-dimensional O/H patched grid used in a turbine cascade computation[5] shows the slope discontinuity at the O-H grid interface and has extremely large-aspect-ratio meshes near the leading edge. The numerical scheme used in that work is a cell-centered scheme based on the Beam and Warming approximate factorization.[6] Although the flexibility of the scheme is indicated in two dimensions, its likely limitations in three dimensions suggest an alternative approach.

Two- and three-dimensional composite grids were generated in an attempt to improve the grid quality in terms of mesh control, smoothness, and orthogonality. A numerical scheme that will run on this composite grid is discussed.

Composite Grid

A stator vane ring and the vane geometry at mean section are shown in Figs. 1(a) and (b), respectively. This vane ring was used for an annular cascade experiment reported in Ref. 2. To alleviate numerical errors associated with the H-grid skewness for the blunt leading edge, a standard C-grid (Fig. 2(a)) was generated to examine any possible advantage over the H-grid for this turbine vane geometry. But the grid became very skewed on the suction side of the passage because the blade was highly cambered (Fig. 2(b)). In addition, the control of the mesh distribution in the blade region was limited with the C-grid, because it affected the mesh distribution in the upstream region. Results of numerical simulations

using this grid are likely to be affected by the grid skewing.[3]

Two- and three-dimensional composite grids were generated as shown in Figs. 3(a) and (b), respectively. The blades had a constant profile from hub to tip and were stacked at the trailing edge (Fig. 3(c)). Only a selected number of grid lines are shown for illustration. This composite grid consists of the C-grid (or O-grid) in the immediate vicinity of the blades and the H-grid in the upstream region and in the middle of the blade passage between the C-grids. The C-grid (or O-grid) portion can be generated by using either the elliptic method[7] or the algebraic method.[8] At the C-grid and H-grid interface (or O-H interface) the slope continuity was preserved so that no special numerical approximations are needed for the derivatives at the interface. The C-grid (or O-grid) is orthogonal to the blade surface and provides a good boundary layer resolution near the leading edge. This composite grid has better smoothness and orthogonality and provides more flexibility in controlling the meshes than any single grid for a typical turbine blade with large camber and a rounded leading edge.

The C-grid portion of the composite grid was generated by the elliptic grid generation code.[9] Two subroutines of the code were modified.[10] The outer boundary subroutine was modified, to improve generation of the periodic boundaries, by using the mean camber line for high-solidity blade rows. The inner boundary subroutine was modified to allow a more general clustering of points about the leading and trailing edges of highly cambered turbine blades. Figure 4 illustrates the portion of the C-grid between $\eta = 0$ and $\eta = \eta_c$ that was retained for the composite grid. The grid lines at $\eta = 0$ and $\eta = \eta_{max}$ are inner and outer boundaries, respectively. The choice of η_c is arbitrary and may depend on the extent of the shear flow region. The η' coordinate is a translation of the η coordinate by one pitch in circumferential direction. For the H-grid in the middle between the blades, the C-grid points at $\eta = \eta_c$ and $\eta = \eta_{c+1}$ on the suction side and at $\eta' = \eta_c'$ and $\eta' = \eta_{c+1}'$ on the pressure side are used in the cubic spline to preserve slope continuity at the C-H interface and to effect a smooth change in mesh size. The three-dimensional grid in Fig. 3(b) was constructed algebraically from the composite grids of the hub and shroud surfaces by using the cubic spline and exponential stretching.

A two-dimensional composite grid was also generated for the first-stage stator and rotor of the space shuttle main engine (SSME) fuel-turbopump turbine. Figure 5 shows the construction of the C-H composite grid for the rotor blade. Figure 6 shows the composite grid generated for the stator vane. One merit of the composite grid is the flexibility in constructing the H-grid in the upstream region. One can choose the number and distribution of meshes, and the H grid can be extended in the upstream direction independently of the C-grid part of the composite grid around the blade. With this flexibility the downstream end of the grid for the stator can easily be matched with the upstream end of the grid for the rotor in the vaneless space between the stator and rotor. Figure 7 illustrates that the composite grid can be easily constructed for a stator/rotor interaction study.

Numerical Scheme

Semidiscrete Finite-Volume Scheme

The two-dimensional C-H composite grid has an unusual grid point that is connected to more than four neighboring nodes (Fig. 8(a)). At this point the usual differencing techniques cannot be applied. If a standard finite-difference scheme were used to solve a flow problem on this composite grid, it would be difficult to treat this special point. A finite-volume scheme[11] to be used with this composite grid presents no problem and requires no special treatment because each interior cell of this composite grid has only four neighboring cells in two dimensions and six neighboring cells in three dimensions (Fig. 8(b)). The finite-volume scheme is described briefly here.

The Euler equations in integral form can be written as

$$\frac{\partial}{\partial t} \iiint_{\Omega} w \, d\Omega + \iint_{\partial\Omega} \underline{F} \cdot d\underline{S} = 0 \qquad (1)$$

for a fixed region Ω with boundary $\partial\Omega$. Here w represents the conserved quantity, $\underline{F}$ is the corresponding flux term, and t is time.

A convenient way to ensure a steady-state solution independent of the time step is to separate the space and time discretization procedures. In the semidiscrete finite-volume scheme one begins by applying a semidiscretization in which only the spatial derivatives are approximated. To derive a semidiscrete model that can be used to treat complex geometric domains, the computational domain is divided into quadrilateral cells. Assuming that the dependent variables are known at the center of each cell, a system of ordinary differential equations is obtained by applying equation (1) separately to each cell. These equations have the form

$$\frac{d}{dt} (S_{ij} w_{ij}) + Q_{ij} = 0 \qquad (2)$$

where S_{ij} is the cell area and Q_{ij} is the net flux out of the cell. This can be evaluated as

$$\sum_{k=1}^{4} (\Delta y_k \, f_k - \Delta x_k \, g_k) \qquad (3)$$

where f_k and g_k denote values of the flux vectors $\underline{f}$ and $\underline{g}$ on the kth edge, Δx_k and Δy_k are the increments of x and y along the edge with appropriate signs, and the sum is over the four sides of the cell. The flux vectors are evaluated, for example, by averaging the values in the cells on either side of the edge:

$$f_1 = \frac{1}{2} (f_{i+1,j} + f_{i,j})$$

The scheme constructed in this manner reduces to a central-difference scheme on a Cartesian grid and is second-order accurate in space provided that the mesh is smooth enough. It also has the property that uniform flow is an exact solution of the difference equations.

LU Implicit Scheme

The composite grid is generated to be used for both inviscid and viscous flow calculations. For

viscous calculations the C- or O-meshes must be very fine to resolve the boundary layer. And it is likely that the time step imposed by an explicit stability bound will be much less than that imposed by the accuracy bound of an implicit scheme. Since an obvious way to accelerate convergence to a steady state is to increase the size of the time step, an implicit scheme is expected to have a faster convergence. Although the alternating direction implicit (ADI) scheme has been valuable in two-dimensional problems, its inherent limitations in three dimensions suggest an alternative approach. An LU implicit scheme was demonstrated to be efficient and robust for both external and internal flows in a broad flow regime.[12-14] The scheme was extended for three-dimensional flows on an H-grid.[15] This scheme will be used for the numerical simulation of turbine flows on the composite grid.

The conservation law form of the Euler equations in Cartesian coordinates for two-dimensional flow is

$$\frac{\partial \underline{W}}{\partial t} + \frac{\partial \underline{F}}{\partial x} + \frac{\partial \underline{G}}{\partial y} = 0 \qquad (4)$$

where $\underline{W}$ is the vector of dependent variables and $\underline{F}$ and $\underline{G}$ are convective flux vectors:

$$\underline{W} = (\rho, \rho u, \rho v, \rho E)^T$$

$$\underline{F} = [\rho u, \rho u^2 + p, \rho v u, u(\rho E + p)]^T \qquad (5)$$

$$\underline{G} = [\rho v, \rho u v, \rho v^2 + p, v(\rho E + p)]^T$$

where ρ, u, v, E, and p are density, velocity components, total energy, and pressure. The pressure is obtained from the equation of state

$$p = \rho(\gamma - 1)\left[E - \frac{1}{2}(u^2 + v^2)\right] \qquad (6)$$

where γ is the ratio of specific heats. These equations are to be solved for a steady state $\partial \underline{W}/\partial t = 0$, where t denotes time.

Let the Jacobian matrices be

$$A = \frac{\partial \underline{F}}{\partial \underline{W}} \quad , \quad B = \frac{\partial \underline{G}}{\partial \underline{W}}$$

and let the correction be

$$\delta \underline{W} = \underline{W}^{n+1} - \underline{W}^n$$

where n denotes the time level.

The linearized implicit scheme for a system of nonlinear hyperbolic equations such as the Euler equations can be formulated as

$$\left[I + \beta \Delta t (D_x A + D_y B)\right] \delta \underline{W} + \Delta t \underline{R} = 0 \qquad (7)$$

where I is the identity matrix and $\underline{R}$ is the residual

$$\underline{R} = D_x \underline{F}(\underline{W}^n) + D_y \underline{G}(\underline{W}^n)$$

Here D_x and D_y are central difference operators that approximate $\partial/\partial x$ and $\partial/\partial y$.

If $\beta = 1/2$, the scheme remains second order accurate in time; for other values of β the time

accuracy drops to first order. The unfactored implicit scheme (eq. (7)) produces a large block-banded matrix, which is very costly to invert and requires huge storage. An unconditionally stable implicit scheme that has error terms at most of order $(\Delta t)^2$ in any number of space dimensions can be derived by LU factorization

$$\left[I + \beta \Delta t (D_x^- A^+ + D_y^- B^+)\right] \left[I + \beta \Delta t (D_x^+ A^- + D_y^+ B^-)\right]$$

$$\times \delta \underline{W} + \Delta t \underline{R} = 0 \qquad (8)$$

where D_x^- and D_y^- are backward difference operators and D_x^+ and D_y^+ are forward difference operators. The reason for splitting is to ensure the diagonal dominance of lower and upper factors as well as to make use of the built-in implicit dissipation.

Here A^+, A^-, B^+, and B^- are constructed so that the eigenvalues of "+" matrices are nonnegative and those of "-" matrices are nonpositive.

$$A^+ = \frac{1}{2}(A + r_A I), \qquad A^- = \frac{1}{2}(A - r_A I)$$

$$B^+ = \frac{1}{2}(B + r_B I), \qquad B^- = \frac{1}{2}(B - r_B I)$$

where

$$r_A \geq \max(|\lambda_A|), \qquad r_B \geq \max(|\lambda_B|)$$

Here λ_A and λ_B represent eigenvalues of Jacobian matrices. Equation (8) can be inverted in two steps. The LU implicit scheme needs the inversion of sparse triangular matrices, which can be done efficiently without using large storage. This scheme has only two factors in three dimensions.

Concluding Remarks

Two- and three-dimensional composite grids were generated to improve grid quality for an annular turbine cascade in terms of smoothness, resolution, and orthogonality. This composite grid, which has a C-type (or O-type) grid in the immediate vicinity of the turbine blade, provides a good boundary layer resolution around the leading-edge region for viscous calculation, has orthogonality at the blade surface and slope continuity at the C-H (or O-H) interface, and controls the mesh distribution in the upstream region without using excessive grid points. This composite grid eliminates the undesirable qualities of a single grid when generated for a typical turbine geometry.

The C-H composite grid has flexibility in constructing the H-grid in the upstream region so that the grids for the stator vanes and rotor blades of the SSME fuel-turbopump turbine can easily be matched in the vaneless space and can be used for the stator-rotor interaction study.

A finite-volume lower-upper (LU) implicit scheme is to be used in solving the turbine flows on the composite grid. This grid has a special grid node that is connected to more than four neighboring nodes in two dimensions and to more than six nodes in three dimensions. But the finite-volume approach poses no problem at the special point because each interior cell has only four neighboring cells in two dimensions and only

six cells in three dimensions. The finite-volume LU implicit scheme was proven to be robust and efficient in a broad flow regime and is expected to yield accurate solutions on the improved composite grid.

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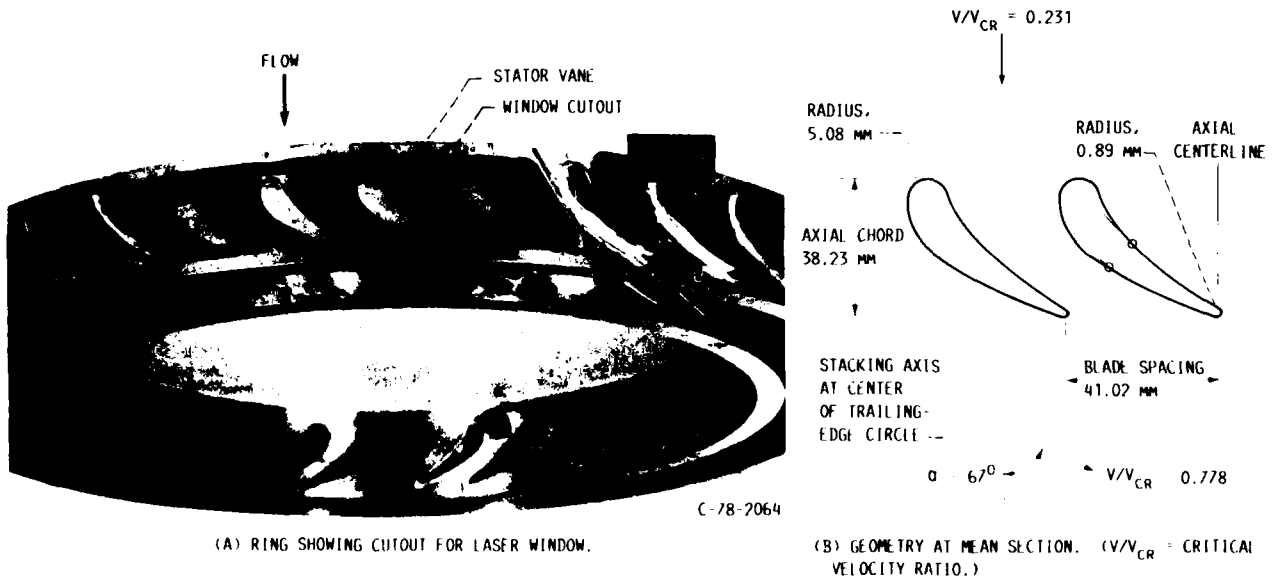
FLOW

STATOR VANE

WINDOW CUTOUT

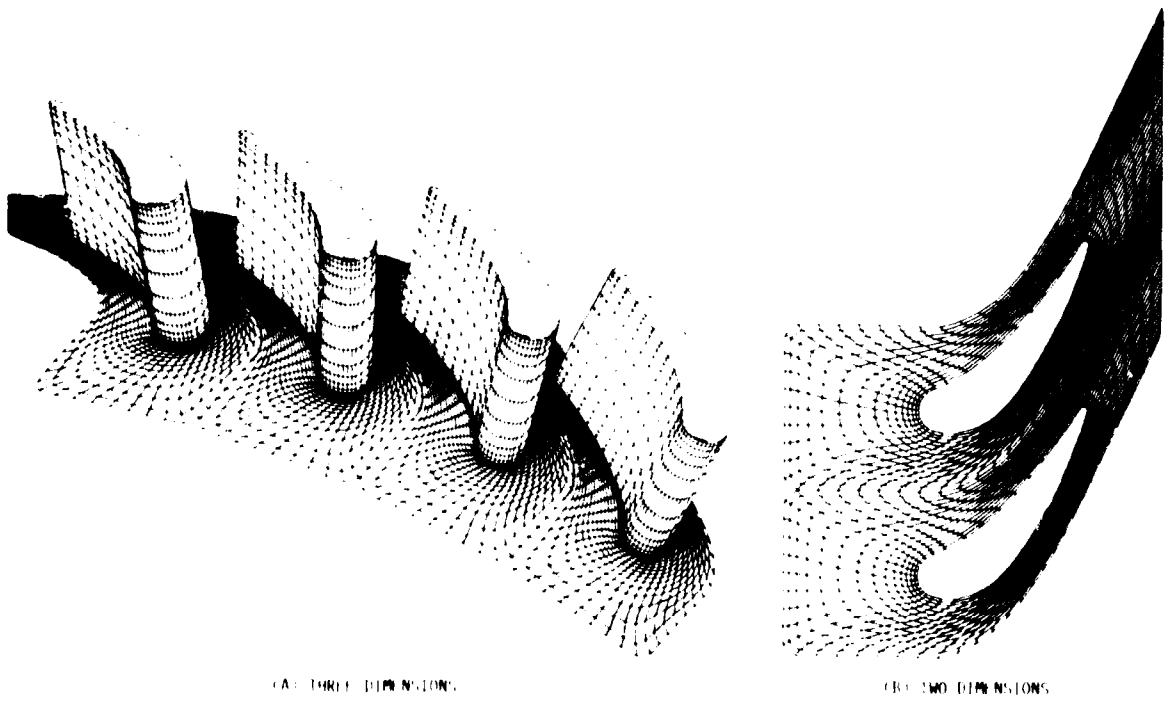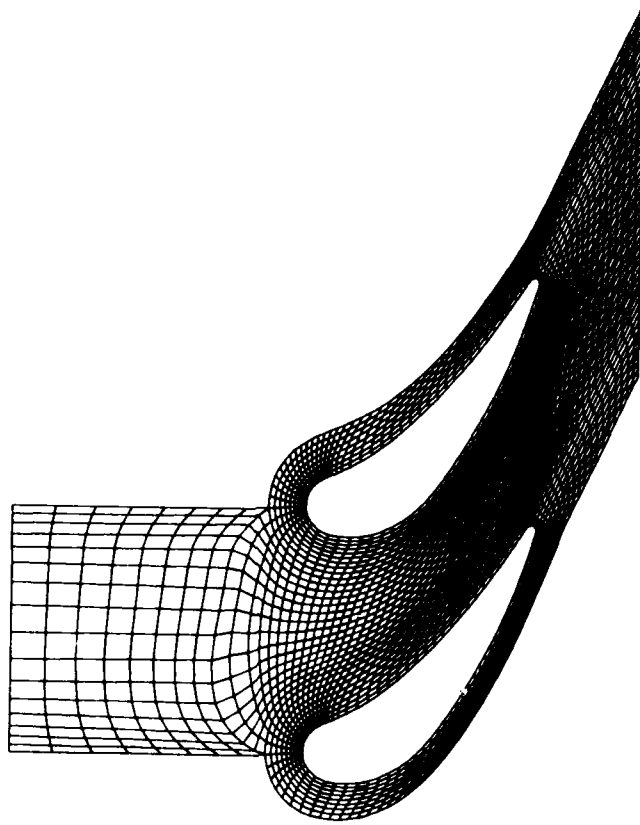C-78-2064

(A) RING SHOWING CUTOUT FOR LASER WINDOW.

$V/V_{CR} = 0.231$

RADIUS, 5.08 MM

RADIUS, 0.89 MM

AXIAL CENTERLINE

AXIAL CHORD 38.23 MM

STACKING AXIS AT CENTER OF TRAILING- EDGE CIRCLE

BLADE SPACING 41.02 MM

$\alpha = 67^0$

$V/V_{CR} = 0.778$

(B) GEOMETRY AT MEAN SECTION. (V/V_{CR} = CRITICAL VELOCITY RATIO.)

FIGURE 1. CORE TURBINE STATOR VANE.

(A) THREE DIMENSIONS.

(B) TWO DIMENSIONS

FIGURE 2. STANDARD C GRID FOR CORE TURBINE VANES.

(A) TWO DIMENSIONS.
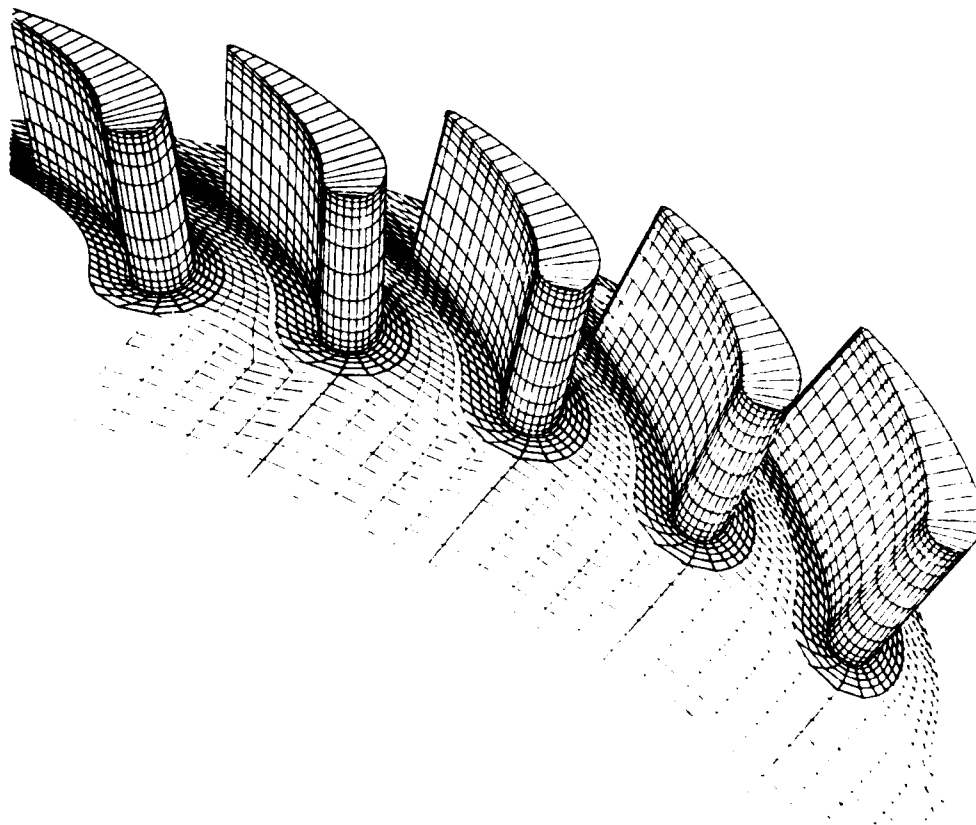
(B) THREE DIMENSIONS.

FIGURE 3. - COMPOSITE GRID FOR CORE TURBINE STATOR VANES.
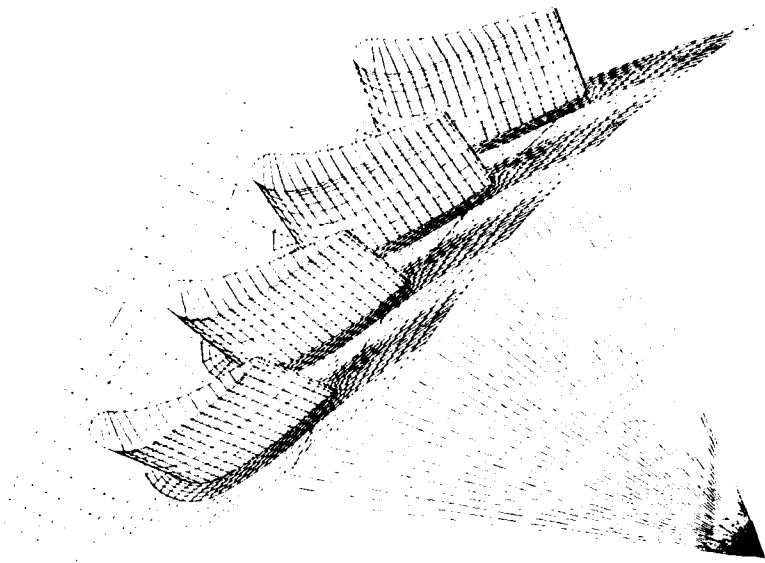
(C) STACKED AT TRAILING EDGES.
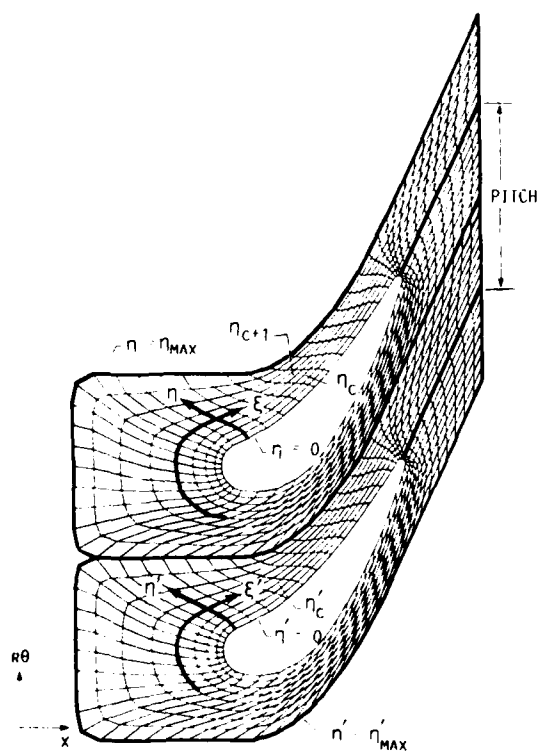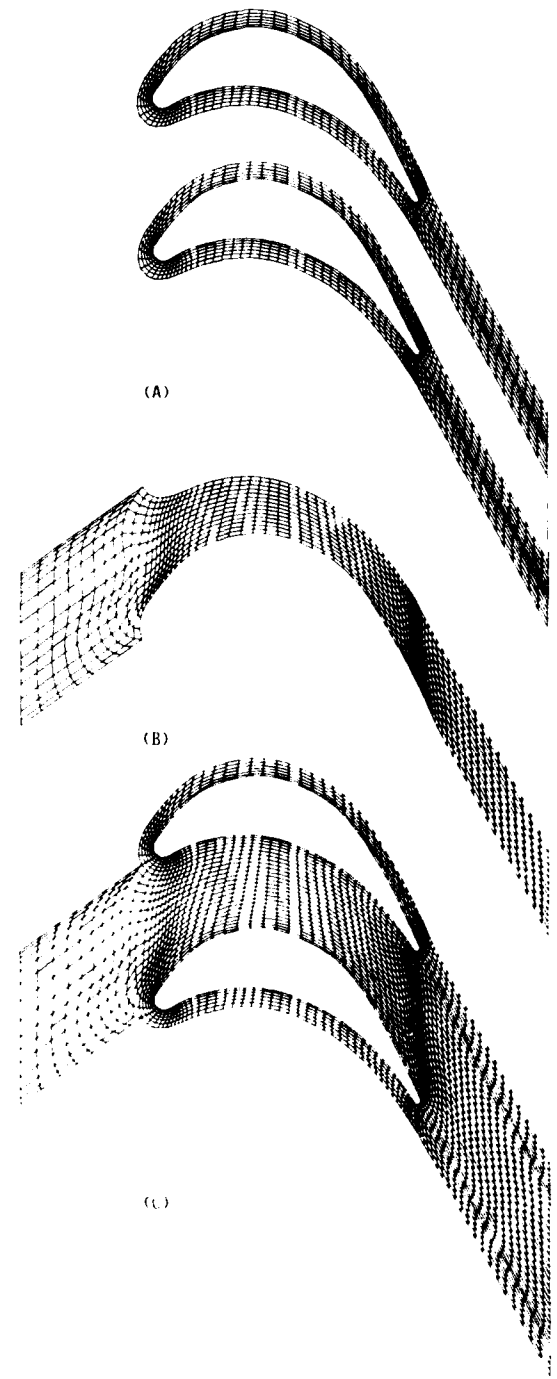
FIGURE 3. - CONCLUDED.

(A)

(B)

(C)

FIGURE 4. - STANDARD TWO-DIMENSIONAL C-GRID FOR CORE TURBINE VANES WITH BODY-FITTED COORDINATES SHOWN.

(A) C-GRID IN IMMEDIATE VICINITY OF ROTOR BLADE.
(B) PATCHING OF GRID IN BLADE REGION AND EXTENSION OF H-GRID INTO UPSTREAM REGION.
(C) COMPOSITE GRID.

FIGURE 5. - CONSTRUCTION OF C-H COMPOSITE GRID FOR FIRST-STAGE ROTOR OF SSME FUEL TURBOPUMP TURBINE.

7

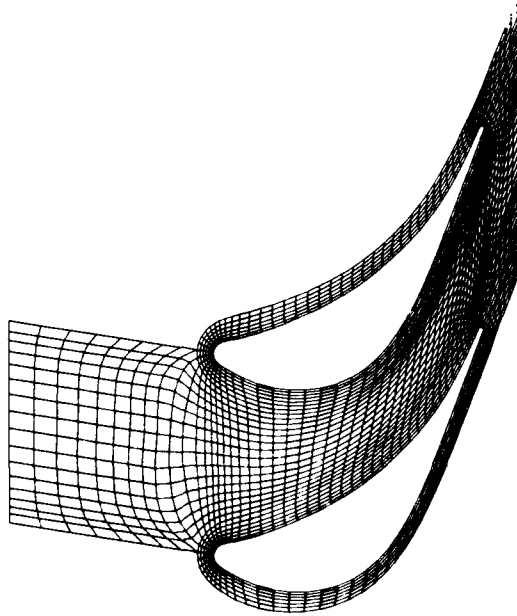FIGURE 6. - COMPOSITE GRID FOR FIRST-STAGE STATOR OF SSME
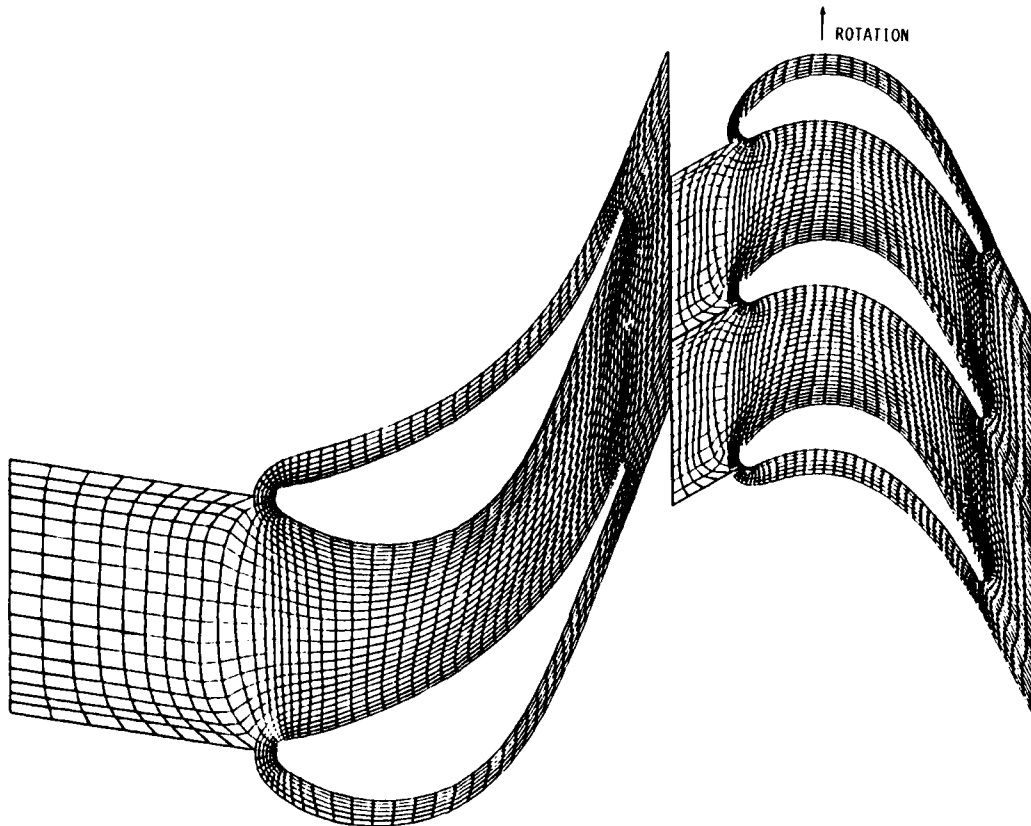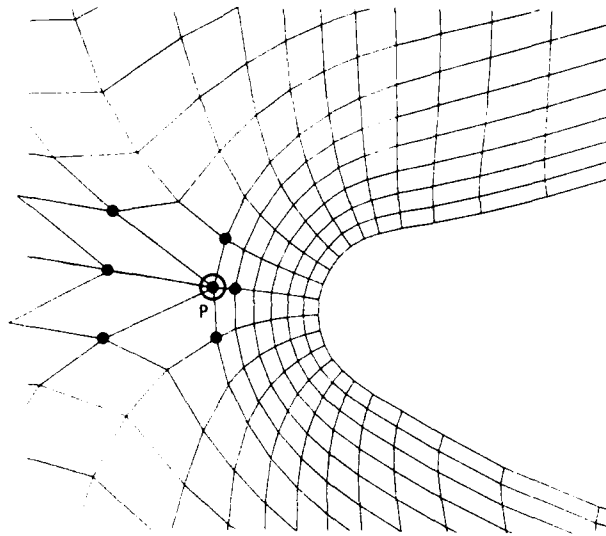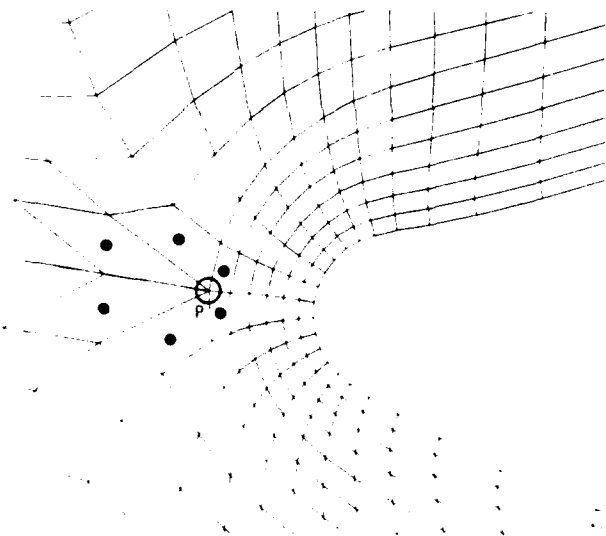FUEL-TURBOPUMP TURBINE.

↑ ROTATION

FIGURE 7. - COMPOSITE GRID FOR FIRST STAGE OF SSME FUEL-TURBOPUMP TURBINE.

(A) GRID NODES TO BE USED IN FINITE-DIFFERENCE SCHEME.

(B) CELL CENTERS TO BE USED IN FINITE-VOLUME SCHEME.

FIGURE 8. - SPECIAL POINT (P) OF COMPOSITE GRID.

9

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16. Abstract

A composite grid was generated in an attempt to improve grid quality for a typical turbine blade with large camber in terms of mesh control, smoothness, and orthogonality. This composite grid consists of the C grid (or O grid) in the immediate vicinity of the blade and the H-grid in the upstream region and in the middle of the blade passage between the C-grids. It provides a good boundary layer resolution around the leading-edge region for viscous calculation, has orthogonality at the blade surface and slope continuity at the C-H (or O-H) interface, and has flexibility in controlling the mesh distribution in the upstream region without using excessive grid points. This composite grid eliminates the undesirable qualities of a single grid when generated for a typical turbine geometry. A finite-volume lower-upper (LU) implicit scheme can be used in solving for the turbine flows on the composite grid. This grid has a special grid node that is connected to more than four neighboring nodes in two dimensions and to more than six nodes in three dimensions. But the finite-volume approach poses no problem at the special point because each interior cell has only four neighboring cells in two dimensions and only six cells in three dimensions. The finite-volume LU implicit scheme was demonstrated to be robust and efficient for both external and internal flows in a broad flow regime.

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6 — 81

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