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Contract N00014-75-C-0694; NR-097-395

Interim Report

Heat and Momentum Transfer to Internal

Turbulent Flow of Helium-Argon Mixtures

in Circular Tubes

by Paul E. Pickett

CONVECTIVE HEAT TRANSFER FOR SHIP PROPULSION.

Dr. Donald M. McEligot Aerospace and Mechanical Engineering Department

3 January, 1978

Prepared for:

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## Interim Report

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HEAT AND MOMENTUM TRANSFER TO INTERNAL, TURBULENT FLOW OF HELIUM-ARGON MIXTURES IN CIRCULAR TUBES

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Research Sponsored by

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#### ABSTRACT

S.

The results of an experimental investigation of friction and heat transfer parameters for turbulent flow of helium-argon, mixtures in smooth, electrically heated, circular tubes are presented. Experimental results are compared to existing experimental correlations and to analytical results. Results of air and helium from the same experimental apparatus are included for comparison.

In this experiment helium-argon mixtures with molecular weights between 15.3 and 29.7 are used, This range resulted in Prandtl numbers between 0.42 and 0.49. Inlet Reynolds numbers range from 31200 to 102000, maximum wall temperatures from 392 to 828 K, maximum wall-to-bulk temperature ratios to 1.82, maximum wall heat flux values to 511 KW/m<sup>2</sup>, and pressures from 469 to 967KPa.(4.7 to 9.7 atmospheres).

Existing experimental correlations, developed using gases with Prandtl numbers of approximately 0.7, are compared to the measured friction and heat transfer results. Adiabatic friction factors and friction factors with heat addition are predicted within ±4 and ±10 percent, respectively. Nusselt numbers for fully developed, constant property conditions are predicted within ±5.0 percent. An empirical equation that correlates the helium-argon data within ±15 percent, and includes entrance and variable

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property effects is presented.

Using a recently developed technique that compares numerically calculated and measured constant property Nusselt numbers, turbulent Prandtl numbers in the wall region for helium-argon mixtures are determined. The validity of using these turbulent Prandtl numbers in a variable property numerical analysis is examined. The variation of turbulent Prandtl number with respect to Reynolds number and molecular Prandtl number is also inspected.

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### NOMENCLATURE

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a,	exponent used to account for temperature variation
	of viscosity;
<sup>a</sup> 1,	array of system parameters;
Α',	calibration constant for the laminar flow element;
A <sub>cs</sub> ,	cross sectional area of tube;
b,	exponent used to account for temperature variation
	of conductivity;
в',	calibration constant for the laminar flow element;
с,	velocity of sound;
° <sub>p</sub> ,	specific heat at constant pressure;
D,	inside diameter;
Ε,	voltage drop;
g,	gravitational constant;
s <sub>c</sub> ,	dimensional conversion factor;
G,	mass flow rate per unit area;
h,	heat transfer coefficient;
i,	enthalpy per unit mass;
k,	force constant in Lennard-Jones (6-12) potential;
K,	thermal conductivity;
L,	length between pressure taps in laminar flow
	element;
l,	mixing length;
m ,	mass flow rate;
ĥ,	molal mass;

vi

ΔP, pressure drop; Ρ, power; heat transfer rate; **q**, q', heat transfer rate per unit length; q", heat flux; volume flow rate; Q, radius; r, R, gas constant for a particular gas; R', resistance per unit length; R. universal gas constant; Τ, temperature; velocity in axial direction; u. axial distance from start of heating; x, radial distance from wall; у, array of measured values; Υ,, z, a calculated quantity. Greek symbols thermal diffusivity,  $K/c_p \rho$ ; α, ε, force constant in Lennard-Jones potential; eddy diffusivity for heat; ε<sub>н</sub>, eddy diffusivity for momentum; ε<sub>M</sub>, ratio of specific heats,  $c_p/c_v$ ; γ, von Karman constant κ, 0.4; μ, absolute viscosity; ν, kinematic viscosity;

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ρ,	density;
σ,	variance or standard deviation;
τ,	shear stress.
Non-dimens	sional parameters
f,	friction factor, $2g_c \rho \tau_w / G^2$ ;
Gr,	Grashof number based on wall heat flux,
	$gD^4 q_w'/(v^2 \mu c_D T)_i;$
м	Mach number, j/c
Nu, Pr,	Nusselt number, hD/K; Prandtl number, c µ/K;
q <sup>+</sup> ,	heat flux parameter, $q''_w/(Gc_{p,i}T_i)$ ;
Re,	Reynolds number, GD/µ;
y <sup>+</sup> ,	wall distance parameter, $y(g_c \tau_w \rho)^{\frac{1}{2}}/v$ ;
y <sup>+</sup>	empirical constant in van Driest mixing length
	model, 26.

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Subscripts

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b,	evaluated at bulk temperature;
cond,	heat conduction;
cp,	constant property condition;
DB,	Dittus-Boelter;
gen,	heat generation;
i,	inlet; an index;
Max,	maximum;
ref,	reference;
t,	turbulent;
VD,	van Driest;
w,	wall;
Xe,	xenon;
∞.	environment conditions.

viii

#### INTRODUCTION

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The closed Brayton cycle using inert gases as working fluids has been considered for use in many current applications. The Navy has investigated its use for undersea and surface ship propulsion. NASA has examined it for future space missions requiring relatively large amounts of electric power (100 ~ 500 Kw) [1,2,3,4]. Binary mixtures of helium and heavier inert gases, such as argon or xenon, have been considered as possible working fluids in these closed Brayton systems. The increase in density, due to the heavier inert gas, reduces the size of the compressor and turbine. The thermal conductivity of the binary mixture is lower than that of helium, thus causing an increase in the size of the heat exchangers. At an intermediate molecular weight an optimum can be attained.

Fig. 1 illustrates the relative heat transfer of helium-argon and helium-xenon mixtures compared to the pure gases and air. The relative heat transfer coefficients were calculated using the Dittus-Boelter type relation

$$h = 0.021 \ \text{Re}^{0.8} \ \text{Pr}^{0.4} \ (\text{K/D}) \tag{1}$$

and were normalized with respect to the lowest value. The geometry and mass flux were kept constant. This resulted in a relative heat transfer coefficient of the form

$$h/h_{Xe} = (c_p \mu_{Xe}/c_{pXe} \mu)^{0.4} (K/K_{Xe})^{0.6}.$$
 (2)



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All properties were taken at 860K and 101KPa. Vanco [4] performed a similar analysis, but kept the geometry and molal flow rate constant, which gave quite different curves. Examination of Fig. 1 shows why helium-xenon is the prime candidate for a working fluid in the closed Brayton cycle. Helium-argon has been investigated initially due to expense and convenience of experimental apparatus.

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The purpose of this research was to determine, for turbulent flow in tubes, the momentum and heat transfer characteristics of helium-argon mixtures. No basic momentum and heat transfer experimental work for fluids with Prandtl numbers between 0.1 and 0.67 presently exists in the literature. Until recently, it was thought that no fluids existed in this Prandtl number range [5,6]. The mixtures of helium and heavier inert gases fill this void, having Prandtl numbers between 0.25 and 0.67. Fig. 2 shows the variation of molecular Prandtl number, Pr, as a function of molecular weight and temperature for heliumargon and helium-xenon [7]. It can be seen that the Prandtl number varies little with temperature.

Experimental correlations, such as equation 1, were developed using air ( $\Pr \simeq 0.7$ ) and helium ( $\Pr = 0.67$ ). Extension of similar experimental correlations for calculating adiabatic friction factors, average friction factors with heat addition, Nusselt numbers at constant property



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Fig. 2. Variation of Prandtl number with respect to molecular weight for helium-argon and helium-xenon.

conditions, and Nusselt numbers with variable property and thermal entry effects included were examined in this study using helium-argon mixtures. Mixtures at molecular weights of approximately 15 (Pr = 0.42), 27 (Pr = 0.46), and 30 (Pr = 0.49) were used. For comparison, experiments with air and helium were also performed. Experimental studies similar to this one, except using air, helium, or hydrogen include those by Perkins and Worsøe-Schmidt [8], McEligot and Magee [9], Taylor [10], and Dalle Donne and Bowditch [11].

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In many analyses that predict turbulent heat transfer results, the value of the turbulent Prandtl number is needed [12]. The turbulent Prandtl number,  $Pr_t$ , is defined as the ratio of eddy diffusivity of momentum and eddy diffusivity of heat,  $\varepsilon_M / \varepsilon_H$ . The eddy diffusivities are defined by the transport relationships,

$$\tau/\rho = (v + \varepsilon_{\rm M}) \frac{\partial u}{\partial y}$$

$$q''/\rho c_{\rm p} = -(\alpha + \varepsilon_{\rm H}) \frac{\partial T}{\partial y}$$
(3)

and are used to account for the additional momentum and heat transport caused by turbulent mixing.

Much work has been done, both analytical and experimental, to develop methods to predict  $Pr_t$ . As of yet, no generally accepted method exists. Reynolds [13] examined more than thirty ways that have been developed to determine  $Pr_t$ . For more background information his review can be consulted. Quarmby and Quirk [14] demonstrated the wide

range of Pr<sub>t</sub> values that are predicted by different analyses and measured data. For air and other common gases, they showed that different methods predict Pr<sub>t</sub> near the wall from 0.5 to infinity.

Due to large uncertainties [15], experimental measurements haven't clarified the discrepancies. The measurements have indicated that  $Pr_t$  is a function of Pr, position in the flow, and turbulence intensity [13]. It has been generally observed that  $Pr_t$  increases as the wall is approached, and that the relationship between  $Pr_t$  and Pris [12,13]

 $\Pr_{+} \leq 1 \text{ for } \Pr \geq 1 \text{ unless } \Pr \leq 1.$  (4)

A recent technique, developed by McEligot, Pickett, and Taylor [16], determines  $Pr_t$  in the wall region by comparing the experimentally measured and numerically calculated axial variation of Nusselt number. The Nusselt numbers are calculated for the constant properties condition, and the measured Nusselt numbers are extrapolated to a constant properties condition. The technique was used in this investigation to determine  $Pr_t$  in the wall region for mixtures of helium-argon. By comparing  $Pr_t$ for helium-argon mixtures with results for air [16], the variation of  $Pr_t$  as a function of Pr was examined. The variation of  $Pr_t$  as a function of Reynolds number was also examined.

Relatively high heating rates could possibly occur

in the heater tubes of the closed Brayton cycle. These high heating rates cause significant variation of properties, and the constant properties idealization becomes invalid. To calculate bulk Nusselt numbers of heliumargon mixtures at these conditions, the  $Pr_t$  determined for constant properties was used in a numerical analysis in which the properties were allowed to vary. To validate using  $Pr_t$  determined for constant property conditions in a variable properties analysis, calculated and measured bulk Nusselt numbers were compared. By examining this comparison, the possibility that the helium and argon had separated, due to the Soret effect [17], was examined.

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#### GAS PROPERTIES

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The properties needed for this study were the compressibility, viscosity, thermal conductivity, specific heat, enthalpy, speed of sound, and gas constant. The properties of air have been studied extensively, and tables listing these properties are readily available. The <u>Tables of the Thermal Properties of Gases</u> [18] were used in this investigation. The properties of helium and helium-argon mixtures were calculated theoretically. For all of the gases, the viscosity and thermal conductivity were assumed to be independent of pressure.

The helium and helium-argon mixtures were assumed to be ideal gases, thus making the compressibility equal to a constant value of one. This is a reasonable assumption for the range of pressures (101.3 - 967.3KPa) and temperatures (294 -  $828^{\circ}$ K) used in this experiment. Since helium and argon are monatomic, and the temperatures used in this study were not too great, the equation [19]

$$c_{\rm p} = (5/2) R$$
 (5)

was used to calculate the specific heat. The specific heat was assumed to be constant, and the gas constant was calculated from the relation

$$\mathbf{R} = \mathcal{R}/\hat{\mathbf{M}}.$$
 (6)

Using the ideal gas and constant specific heat assumptions, simple equations for the enthalpy and speed of sound can be derived [20]

$$i = c_p (T - T_{ref})$$
(7)

 $c = \sqrt{\gamma R T} = \sqrt{5/3 R T}$ (8)

 $T_{ref}$  is an arbitrary reference temperature. From the assumptions already mentioned, the ratio of specific heats,  $\gamma$ , becomes a constant value of 5/3.

The viscosity and thermal conductivity of the helium and helium-argon mixtures were calculated using the Lennard-Jones (6-12) potential in the Chapman-Enskog kinetic theory [17]. The predicted properties were compared with experimental measurements.

The force constants,  $\varepsilon/k$  and  $\sigma$ , suggested by Hirschfelder, Curtiss and Bird [17] were tried originally. The predicted properties were compared with the experimental values only for the range of temperatures used in this study. The predicted helium viscosities were five percent below the experimental measurements of Dawe and Smith [21] and Kalelkar and Kestin [22]. The predicted thermal conductivities of helium agreed within one percent of the measurements by Saxena and Saxena [23], but were five percent below the values calculated from experimental viscosity measurements of Kalelkar and Kestin [22]. The predicted viscosities of helium-argon mixtures at 870°K were three to five percent below the measured values of Kalelkar and Kestin [22]. The predicted thermal conductivities of helium-argon mixtures at 790°K were five to

nine percent below the measured values of the Thermophysical Properties Research Center [24], and the measured values of von Ubisch repeated by Gandhi and Saxena [25].

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In an attempt to get better agreement between predicted and measured values, force constants suggested by DiPippo and Kestin [26] were tried. With these force constants, the predicted viscosities of both helium and helium-argon agreed within one percent of the measured values mentioned in the previous paragraph. The predicted thermal conductivities of helium agreed within one percent of the values of Kalelkar and Kestin [22], but were five percent above the measurements of Saxena and Saxena [23]. The predicted thermal conductivities of helium-argon were essentially unchanged. Since the agreement between the predicted and measured viscosities was improved, and the agreement between the predicted and measured thermal conductivities remained approximately the same, the force constants suggested by DiPippo and Kestin [26] were used. The calculated properties of helium and the helium-argon mixtures used in this investigation are listed in Appendix A.

The properties were inserted in tabular form in the numerical programs that reduced the experimental friction and heat transfer measurements. In the numerical program used to predict heat transfer results, the properties were inserted in equation form. The ideal gas law was used, the specific heat was assumed constant, and the variation

of viscosity and thermal conductivity with temperature was accounted for with the following relations.

$$\mu/\mu_{ref} = (T/T_{ref})^{a}$$
(9)
  
 $K/K_{ref} = (T/T_{ref})^{b}$ 
(10)

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As discussed by McEligot, Taylor, and Durst [7], the exponent "a" ranges from 0.7 to 0.8, and the exponent "b" ranges from 0.7 to 0.75 for the inert gases and their mixtures. The exponents, "a" and "b", of air for the range of temperatures in this study are 0.67 and 0.81, respectively. Thus, the viscosity and thermal conductivity of air, helium, and helium-argon vary with temperature in approximately the same manner.

For the present study the following values of the exponents were used for the mixtures:

at  $\hat{M} = 15.83$ , a = 0.745 and b = 0.718at  $\hat{M} = 27.53$ , a = 0.772 and b = 0.741.

#### EXPERIMENTAL APPARATUS AND PROCEDURE

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The experimental apparatus, arrangement, and procedure was similar to that used by Perkins, Schade, and McEligot [27]. Only differences in the two experiments will be noted here. Instead of a square duct, a circular tube made of Hastelloy-X was used as a test section. The tube had an inside diameter of 0.312 cm. and a wall thickness of 0.056 cm. The test section consisted of a heated section 98 diameters in length preceded by an unheated section 92 diameters in length. The unheated section ensured that the velocity profile was fully developed at the inlet of the heated section. For attachment of the a.c. power cables, stainless steel electrodes were brazed at the upper and lower ends of the heated section. Two pressure taps were used. One was located in the lower electrode and the other 8.0 diameters below the upper electrode. Sixteen premium grade chromel-alumel thermocouples (0.013 cm. diameter) were spot welded to the heated section of the tube using the parallel junction suggested by Moen [28].

In addition to the power supply used by Perkins et al. [27], an a.c. Lincoln welder was used in order to reach the high temperatures at the larger Reynolds numbers used in this experiment. To measure the higher flow rates, the positive displacement meter was replaced by a Meriam

laminar flow element. The latter was calibrated to measure the flow rate within  $^+$ 1.5 percent. Heise gages, inclined water manometers, and vertical mercury or water manometers were used to measure static pressure and pressure drop.

A vacuum external environment was not used in this experiment. The test section was completely enclosed with a heat shield that restricted the convective air currents and helped stabilize the heat loss from the tube to the environment.

The experimental procedure was slightly different than that used by Perkins et al. [27]. The "radiating thermocouple conduction error", discussed by Hess [29], was not exactly appropriate since the test section was surrounded by air at atmospheric pressure. Instead a correlation for natural convection from small wires was introduced, in addition to radiation, as detailed in Appendix E. The heat loss from the tube to the environment was determined using the method described by Campbell and Perkins [30].

To reduce the heat transfer data the same computer program that was used by Perkins et al. [27], was employed in this study, but was modified for use with a circular tube. The basics of this computer program are described

in other reports [30,31,32]. Table 1 summarizes the range of variables covered in this investigation. A more detailed discussion of the experiment is contained in Appendix B. A list of the experimental data is contained Appendix D.

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	Range of Variables in the Present Experiment			
	Air	Helium	Helium-Argon	
Experimental runs	25	4	28	
Molecular weight	28.97	4.003	15.30 - 29.70	
Inlet bulk Reynolds number	32,900 - 100,000	30,200	31,200 - 102,000	
Exit bulk Reynolds number	19900 - 89000	18,400 - 26,600	17,000 - 68,000	
Inlet bulk Prandtl number	0.719	0.667	0.419 - 0.486	
Exit bulk Prandtl number	0.682 - 0,708	0.667	0.426 - 0.495	
Maximum T <sub>w</sub> /T <sub>b</sub>	1.90	1.75	1.82	
Maximum T <sub>w</sub> ( <sup>0</sup> K)	817	789	828	
Maximum q <sup>+</sup>	0.0027	0.0027	0.0032	
Maximum Gr/Re <sub>i</sub> 2	8.90 x 10 <sup>-5</sup>	4.84 x 10 <sup>-5</sup>	$3.22 \times 10^{-3}$	
Maximum Mach number	0.26	0.25	0.33	
<pre>x/D for local bulk Nusselt     numbers     '</pre>	2.1 - 82.0	2.1 - 82.0	2.1 - 82.0	

TABLE 1 -

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#### EXPERIMENTAL RESULTS

#### Friction Results

Adiabatic friction factors were measured before each series of heated runs. These were compared to other researcher's results, and were also used as a check of the pressure, mixture molecular weight, and flow rate measurements. The method described by Shapiro [33] was used to calculate the adiabatic friction factors. The measured friction factors were compared to the experimental correlation of Drew, Koo, and McAdams [34],

 $f = 0.0014 + 0.125 \text{ Re}^{-0.32}.$  (11)

This correlation is for turbulent flow in tubes, and was used because of its simplicity and close agreement with the Kármán-Nikuradse relation. Fig. 3 shows the measured friction factor divided by that calculated from equation (11) plotted as a function of Reynolds number. Air and helium data points are included for comparison. All the measured friction factors are within ±4.0 percent of equation (11), and 76 percent are within ±2.0 percent.

Since only two pressure taps were used, local friction factors could not be determined for experiments with heat addition. Average friction factors were determined in the manner of Humble, Lowdermilk, and Desmon [35]. The average friction factors were compared to an experimental correlation suggested by Taylor [36]. This correlation is for

turbulent flow in tubes with heat addition.

 $f = (0.0014 + 0.125 \text{ Re}_w^{-0.32}) (T_w/T_b)^{-0.5}$  (12) This relation is similar to equation (11), but the bulk Reynolds number is replaced by the modified wall Reynolds number. The term  $(T_w/T_b)^{-0.5}$  is included to account for variation of properties with temperature. Equation (12) was used by Taylor to correlate average friction coefficients measured by several different people. It predicted most of the data within ±10 percent.

Fig. 4 shows the average friction factors with heat addition as measured in this investigation. The friction coefficients are divided by equation (12) and plotted as a function of modified wall Reynolds number. Again, helium and air are included for comparison. All of the data is predicted to within  $\pm 10$  percent by equation (12) and 84 percent is predicted to within  $\pm 4.0$  percent.

#### Heat Transfer Results

To determine the effects of the lower helium-argon Prandtl number on the heat transfer results, the variation of properties with temperature, and the entrance effects were minimized. The entrance effects were minimized by considering primarily the results at which fully developed conditions existed (x/D>20). A method described by Malina and Sparrow [37] was used to approach the constant properties idealization.

For the method described by Malina and Sparrow, a



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Figure 3. Comparison of Adiabatic Friction Factors to Drew, Koo and McAdams Correlation for Air, Helium and Helium-Argon Mixtures.



Figure 4 Comparison of Average Friction Factors to Taylor Variable Properties Correlation for Air, Helium and Helium-Argon Hixtures.

fixed inlet Reynolds number is maintained while the wallto-bulk temperature difference is varied. At a particular axial location, the ratio of experimentally determined bulk Nusselt number to a Dittus-Boelter type correlation is plotted as a function of the difference between wall and bulk temperature. Extrapolation to a difference of zero between the wall and bulk temperature gives a ratio that can be directly used to calculate a constant property Nusselt number, Nu<sub>cp</sub>. Since the ratio of bulk Nusselt number to a Dittus-Boelter type correlation partially eliminates any effects caused by small deviations of the Reynolds number, these deviations should be kept as small as possible. The procedure described in the previous paragraph is demonstrated in Fig. 5 for a helium-argon mixture with a molecular weight of 15.30, inlet Reynolds number of 55200, and inlet Prandtl number of 0.419. Extrapolation for four different axial locations is shown. For this investigation the Dittus-Boelter type correlation used was (equation 1 rearranged)

$$Nu_{DB} = 0.021 \ \text{Re}^{0.8} \text{Pr}^{0.4}$$
. (13)

For a sequence of runs at a nominal inlet Reynolds number, all individual runs had inlet Reynolds numbers within 1.8 percent of the nominal value.

The dashed lines in Fig. 5 show how the error in the constant property Nusselt number was estimated. This



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technique was used by Reynolds, Swearingen, and McEligot [38]. The error in this investigation varied from ±9 percent at small x/D to ±5.4 percent at large x/D. The dominant uncertainty in the Nusselt number is due to uncertainty in the wall-to-bulk temperature difference. This difference is small in the entrance region, thus causing a large error in the Nusselt number. At large x/D, the flow is fully developed and the wall-to-bulk temperature difference is relatively constant. The error in the Nusselt number becomes a minimum, and then increases with increasing x/D due to greater uncertainty in the bulk temperature. Appendix C describes the method used for calculating error in the measured Nusselt number.

For fully developed conditions (x/D>20), with air or helium as the experimental fluid, the ratio  $Nu_{cp}/Nu_{DB}$ varied from 0.94 to 1.00. No dependence on Reynolds number was noticed for the Reynolds number range used in this experiment. For fully developed conditions, with heliumargon mixtures as the experimental fluid, the ratio  $Nu_{cp}/Nu_{DB}$ varied from 0.83 to 0.93. For a helium-argon mixture at a molecular weight of 15.30, inlet Reynolds number of 55200, and x/D value of 56.9, Fig. 5 shows the ratio  $Nu_{cp}/Nu_{DB}$  to be approximately 0.84. From these results, it was determined that the Dittus-Boelter type equation (equation 13) did not predict correct Nusselt numbers for

the Prandtl number range between 0.42 and 0.50.

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A correlation suggested by Kays [39] predicted the constant property Nusselt numbers of the helium-argon mixtures within  $\pm 6.0$  percent.

$$Nu = 0.022 \ \text{Re}^{0.8} \text{Pr}^{0.6} \tag{14}$$

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This equation was recommended for fluids with Prandtl numbers between 0.5 and 1.0, constant properties, a constant heat flux boundary condition, and fully developed turbulent flow. If the coefficient of this equation is changed to 0.021, and the Prandtl number exponent adjusted so that approximately equivalent results are obtained, the resulting equation is

$$Nu = 0.021 \ Re^{0.8} Pr^{0.55}. \tag{15}$$

This equation shows that the exponent of the Prandtl number in equation (13) should be changed from 0.4 to 0.55 in order to accurately predict the constant property Nusselt numbers of the helium-argon mixtures. Fig. 6 shows the constant property Nusselt number divided by equation (15) plotted as a function of Prandtl number for the mixtures. Results are plotted for three Prandtl numbers, four Reynolds numbers, and axial positions at which the flow was fully developed. From the figure it can be seen that equation (15) predicts the constant property Nusselt numbers within ±5.0 percent. At a Reynolds number of 32000, a small effect of the Prandtl number varying between 0.419 and 0.486 can be noticed.

A dependence on Reynolds number was observed for the

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helium-argon mixtures. Since the effect on the constant property Nusselt number was about equivalent to the error in the constant property Nusselt number, only general trends can be discussed. Two trends were observed (Fig. 6). For a particular Prandtl number the ratio of Nu<sub>cp</sub> divided by equation (15) decreased as the Reynolds number increased, and this effect became more pronounced as the Prandtl number decreased.

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To account for the variation of properties and entrance effects in this investigation, the correction factors suggested by Magee [40] were used.

$$\left[\left(T_{w}/T_{b}\right)^{-0.4} + 0.6D/x\right]$$
(16)

The term  $(T_w/T_b)^{-0.4}$  accounts for the variation of properties, and the term 0.6D/x accounts for the entrance effects. If these correction factors are applied to equation (13), the resulting equation is

 $Nu_b = 0.021 \text{ Re}_b^{0.8} Pr_b^{0.4} [(T_w/T_b)^{-0.4} + 0.6D/x].$  (17) For x/D between 2.1 and 81.6 this equation predicted all of the present measured Nusselt numbers for air and helium within ±15 percent and 97 percent of the Nusselt numbers within ±10 percent.

If the correction factors (16) are applied to equation (15) the resulting relation is

 $Nu_{b} = 0.021 \text{ Re}_{b}^{0.8} Pr_{b}^{0.55} [(T_{w}/T_{b})^{-0.4} + 0.6D/x]. \quad (18)$ This equation predicted the helium-argon Nusselt numbers in the fully developed region within ±13 percent, but
underpredicted the Nusselt numbers in the entrance region by as much as 22 percent. To have the same type of accuracy with the helium-argon data that was obtained with the air and helium data changes to the correlation were necessary. 2

As previously discussed, the transport properties of helium-argon vary with temperature in approximately the same manner as those of air and helium. For this reason the term  $(T_{\mu}/T_{b})^{-0.4}$  was retained as a reasonably accurate correction factor for the variation of properties. Kays [39] discusses the effect of different Prandtl numbers in the thermal entrance region of circular tubes. He shows that as the Prandtl number decreases, the effect of the entrance region on the Nusselt number is more pronounced. Because of this, the coefficient in the term 0.6D/x of equation (18) was changed. Since helium-argon has a lower Prandtl number than air, one would expect the coefficient to have a larger value than 0.6. Different values for the coefficient of the entrance effects term were used in equation (18), and compared to the experimentally determined bulk Nusselt numbers of helium-argon. From this comparison it was determined that a value of 0.85 worked best for the coefficient of the entrance effects term. The complete correlation, accounting for entrance effects and variation of properties is

 $Nu_b = 0.021 \text{ Re}_b^{0.8} Pr_b^{0.55} [(T_w/T_b)^{-0.4} + 0.85D/x].$  (19) For x/D between 2.1 and 81.6 this equation predicted all

of the measured bulk Nusselt numbers for helium-argon within  $\pm 15$  percent and 92 percent within  $\pm 10$  percent.

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Measured bulk Nusselt numbers divided by equation (19) are plotted on Fig. 7 as a function of x/D. For clarity, only results from four helium-argon experimental runs were plotted. The data plotted are from experimental runs that include the complete range of experimental variables for the helium-argon mixtures. The greatest difference between the experimental data and equation (19) occurred at high heating rates in the x/D range between 4.0 and 16.0. In this range equation (19) underpredicted the measured Nusselt numbers by 5 to 15 percent.

Few correlations for gases with Prandtl numbers between 0.1 and 0.67 presently exist in the literature. Sleicher and Rouse [41] suggest a correlation for Prandtl numbers between 0.1 and  $10^5$ , and Reynolds numbers between  $10^4$  and  $10^6$ . The correlation is for fully developed conditions, and accounts for property variation.

> $Nu_{h} = 5 + 0.015 \text{ Re}_{f}^{m} \text{Pr}_{u}^{n}$  $m = 0.88 - 0.24/(4 + Pr_{tr})$ (20) $n = 1/3 + 0.5 \exp(-0.6 Pr_w)$

For the helium-argon mixtures, this equation predicted Nusselt numbers that were 15 to 40 percent lower than the Nusselt numbers measured in the fully developed region of this investigation. Equation (19) correlated the data more accurately.

 $(T_w/T_b)_{Max}$ 1.54 1.39 1.19 1.78 0.0019 0.0012 0.0005 0.0032 a<sup>+</sup> Max 2 -20 0.419 0.419 0.465 0.486 Prb Comparison of local bulk Nusselt numbers to equation 19 Ħ Outlet Reb 36200 60900 87500 17000 54800 80600 101000 31200 Inlet Reb 0.0 Molecular Weight 15.30 15.83 27.53 29.70 <u>∩</u>□⊲0 Fig. 7. Comparison of loc for helium-argon mixtures. 2018 **O** 15 TIN 27

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#### NUMERICAL ANALYSIS

## Procedure for Determining the Turbulent Prandtl Number

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To determine Pr<sub>t</sub> for helium-argon mixtures, the numerical method of Bankston and McEligot [42] was used in conjunction with the technique developed by McEligot, Pickett, and Taylor [16]. The numerical method uses finite control volume approximations. It was developed to solve the coupled, partial differential, axisymmetric, boundary layer equations; but can also be used for constant property conditions which uncouples the boundary layer equations. The boundary conditions are the no-slip and impermeable-wall conditions, the inlet conditions, and the wall heat flux.

The technique of McEligot et al. [16] uses the axial variation of the Nusselt number to determine Pr<sub>t</sub> in the wall region. By examination of the simplified energy equation,

$$u \frac{\partial T}{\partial x} = 1/r \frac{\partial}{\partial r} [r(\alpha + \epsilon_{M}/Pr_{t}) \frac{\partial T}{\partial r}]$$
(21)

they showed that the functional dependence of the Nusselt number is

$$Nu = Nu \{x, u(r), \varepsilon_{M}(r), Pr_{+}\}.$$
 (22)

The energy equation was simplified from the general form by using the following assumptions: the axisymmetric boundary layer approximations, hydrodynamic fully developed flow, steady flow at low velocities, and constant fluid properties. By using one of the semi-empirical relationships for  $\varepsilon_{M}(r)$  to determine the velocity profile, u(r), they obtained the result Nu = Nu {x,Pr<sub>t</sub>}. They inverted this relationship to obtain Pr<sub>t</sub> = Pr<sub>t</sub> {Nu(x)}. If Pr<sub>t</sub> is considered one-dimensional, comparison of experimental measurements of Nu(x) with calculated values of Nu(x) can be used to determine Pr<sub>t</sub>(r). McEligot et al. [16] pointed out that direct inversion would be difficult, and iterative use of the numerical procedure described in the previous paragraph was used. The radial variation of the turbulent Prandtl number was assumed to be

$$P_{r} = Pr_{t,w} + \frac{d(Pr_{t})}{d(y/r_{u})} (y/r_{w}).$$
(23)

The results of McEligot et al. [16] showed that a change of  $Pr_t$  in the wall region from 1 to  $\frac{1}{2}$  caused changes of 30 to 45 percent in Nu(x), whereas a change of  $Pr_t$  in the core only caused small changes. For air at a Reynolds number of 44500 and a Prandtl number of 0.72 they determined that

$$Pr_{t,w} = 0.9 \pm 0.1 \text{ and } \frac{d(Pr_t)}{d(y/r_w)} = 0.$$
 (24)

The typical errors of the experimentally measured Nusselt numbers did not allow calculation of  $\frac{d(Pr_t)}{d(y/r_w)}$ .

To determine  $\varepsilon_{M}(r)$  in this investigation, the van Driest mixing length model [43] was used in conjunction with the

Reichardt middle law [54].

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 $\ell_{\rm VD} = \kappa y [1 - \exp(-y^+/y_{\ell}^+)]$  $\epsilon_{\rm VD} = \ell \frac{2}{\rm VD} \frac{\partial u}{\partial y} \text{ and } \epsilon_{\rm M} = \epsilon_{\rm VD} \cdot (2 - \frac{y}{r_{\rm H}}) \cdot [1 + 2(\frac{r^2}{r_{\rm H}})]/6 \quad (25)$ 

The values of  $\kappa$  and  $y_{\ell}^+$  were 0.4 and 26, respectively. With these constants, the predicted friction factors agreed within one percent of equation (11) for the range of Reynolds numbers used in this study. In this study, as in the study by McEligot et al. [16], the errors in the experimentally measured Nusselt numbers did not allow calculation of  $\frac{d(Pr_t)}{d(y/r_t)}$ . The inlet Reynolds number, inlet Prandtl number, constant properties condition, wall heat flux variation, and different values of Pr, were used as input to the numerical procedure. For the first three diameters, the experimental axial wall heat flux variation resembled an exponential approach to a constant value as x increased. For the remaining length, the wall heat flux was constant within two percent. The same axial variation of wall heat flux was used for all of the constant property numerical calculations.

From the numerical analysis the axial variation of Nu<sub>cp</sub> was calculated. By comparing graphs of the experimentally measured Nu<sub>cp</sub> and the calculated Nu (examples in <sup>cp</sup> Fig. 8), Pr<sub>t,w</sub> for helium-argon mixtures was determined. The variation of Pr<sub>t,w</sub> with respect to Reynolds number was examined by comparing Nu<sub>cp</sub> at different Reynolds numbers,



but the same Prandtl number. The variation of Pr<sub>t,w</sub> with respect to Prandtl number was examined by comparing Nu<sub>cp</sub> at different Prandtl numbers, but the same Reynolds number.

### Procedure for Studying High Heating Rates

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As mentioned in the Introduction, relatively high heating rates could possibly occur in the heater tubes of the closed Brayton cycle. To calculate Nu, for heliumargon mixtures at these high heating rates, the numerical method of Bankston and McEligot [42], discussed in the previous section, was used. Properties were allowed to vary, and the relations (9,10) discussed in the Gas Properties section were used. The simple van Driest/Reichardt model eqn. (25), and  $Pr_{t,w}$  determined for helium-argon at constant property conditions were incorporated. No radial variation of  $Pr_t$  was included, thus,  $Pr_t = Pr_{t,w}$ . The axial variation of wall heat flux was similar to the one used for the constant property calculations, but was modified slightly for each experimental run to agree with the wall heat flux variation determined from the experimental measurements.

The axial variation of measured Nu<sub>b</sub> from two heliumargon experimental runs was compared to the calculated axial variation of Nu<sub>b</sub>. From a series of runs with approximately equivalent inlet Reynolds and Prandtl numbers, the runs with the highest and lowest heating rate were

chosen. Since the Pr<sub>t</sub> used was for constant properties, one would expect agreement of measured and calculated Nu<sub>b</sub> at the low heating rate. If the measured and calculated Nu<sub>b</sub> at the high heating rate agreed, this would validate the use of Pr<sub>t</sub> determined from constant property conditions for conditions in which properties varied significantly. If the results at the high heating rate did not agree, this might indicate that either, Pr<sub>t</sub> determined for constant properties couldnot be used for variable property conditions, or that some other phenomenon, such as the Soret effect, was acting.

## Turbulent Prandtl Number Results and Discussion

Fig. 8 illustrates examples of the comparisons between measured  $Nu_{cp}$  and calculated  $Nu_{cp}$  used to determine  $Pr_{t,w}$ . Examples for three Reynolds numbers and two Prandtl numbers are shown. Curves of the calculated  $Nu_{cp}$  are included at four different  $Pr_t$  (0.9,1.0,1.1,1.2). Brackets indicating the experimental error of the measured  $Nu_{cp}$  are also included. Because of the large error in the immediate thermal entry, only results for x/D greater than eight were used to determine  $Pr_{r,w}$ .

Fig. 8a shows the measured and calculated  $Nu_{cp}$  for a helium-argon mixture with a molecular weight of 15.30, Prandtl number of 0.419, and Reynolds number of 32000. By examining results of similar graphs,  $Pr_{t,w}$  was determined to be 1.1 ± 0.1 for helium-argon mixtures with molecular

weights of approximately 15, Prandtl numbers of 0.42, and Reynolds numbers between 32000 and 55200. The measured and calculated Nu<sub>cp</sub> are shown in Fig. 8b and 8c for a heliumargon mixture at a molecular weight of 29.70, Prandtl number of 0.486, and Reynolds numbers of 31600 and 82100. From results of similar graphs,  $Pr_{t,w}$  was determined to be 1.0 ± 0.1 for helium-argon mixtures with molecular weights between 27 and 30, Prandtl numbers between 0.46 and 0.49, and Reynolds numbers between 31600 and 102000.

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The effect of Reynolds number on  $Pr_{t,w}$  can be examined qualitatively using the results in Fig. 8b and 8c. These results are for the same Prandtl number (Pr = 0.486), but Reynolds numbers of 31600 and 82100. For x/D greater than eight, and at the low Reynolds number, the measured  $Nu_{cp}$ are slightly below the calculated  $Nu_{cp}$  for a  $Pr_{t,w}$  of 1.0. At the high Reynolds number and same axial length, the measured  $Nu_{cp}$  are slightly above the calculated  $Nu_{cp}$  for a  $Pr_{t,w}$  of 1.0. For the stated conditions, it appears that  $Pr_{t,w}$  has a weak dependence on Reynolds number, and decreases slightly as the Reynolds number increases.

The effect of molecular Prandtl number on  $Pr_{t,w}$  can be examined using the results from Fig. 8a and 8c summarized in Table 2. The results (24) of McEligot, Pickett and Taylor [16] for air may also be used since the mixture results appear to show that  $Pr_{t,w}$  varies only slightly with Reynolds number. Table 2. Variation of  $Pr_{t,w}$  with respect to Prandtl number.

Gas	Molecular Weight	Prandtl Number	Pr <sub>t,w</sub>	Reynolds Number
Helium-argon	15.30	0.419	$1.1 \pm 0.1$	32000
Helium-argon	29.70	0.486	$1.0 \pm 0.1$	31600
Air	28.97	0.72	$0.9 \pm 0.1$	44500

For the range of Prandtl numbers in Table 2, Pr<sub>t,w</sub> has a relatively strong dependence on Prandtl number and decreases as Prandtl number increases. This dependence agrees with that (equation 4) noted by Reynolds [13].

#### High Heating Rate Results and Discussion

Fig. 9 shows the results of the measured and calculated axial variation of Nu<sub>b</sub> for the two experimental runs that were investigated. The two runs were for a helium-argon mixture at a molecular weight of 29.70, inlet Prandtl number of 0.486, inlet Reynolds numbers of 32000 and 31200, and maximum heating rates of  $q^+ = 0.0006$  and  $q^+ = 0.0032$ . Since at a Pr, value of 1.0, the measured Nu in Fig. 8b were slightly below the calculated Nu<sub>cp</sub>, a Pr<sub>t</sub> value of 1.02 was used. The constants, "a" and "b" in equations (9) and (10) were 0.772 and 0.741, respectively. For both heating rates, the calculated axial variation of Nu, in Fig. 9 agreed with the measured axial variation of Nu<sub>b</sub>, within the accuracy of the measured values. From this example, it appears that  $\Pr_{t,w}$  determined from constant property results can be used to calculate Nu, for variable property conditions with heating rates up to,  $q^+ = 0.0032$ .

At high heating rates a large temperature gradient exists from the wall to the centerline of the tube. At

sufficiently high heating rates, the possibility of separation of the helium and argon due to the Soret effect arises. If separation did occur,  $Nu_b$  at a particular axial location would be expected to change since pure helium or argon have higher Prandtl numbers than helium-argon mixtures. For the high heating experimental run in Fig. 9, the largest wall-to-bulk temperature ratios occur in the axial range, 8.1<x/D<16.4. In this axial range, the measured  $Nu_b$  do fall slightly above the calculated  $Nu_b$ , but this can not necessarily be attributed to the Soret effect, since the calculated  $Nu_b$ .

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The effect of high heating on the axial variation of  $Nu_b$  can be examined by comparing the low and high heating rate results in Fig. 9. Since the thermal conductivity and viscosity of helium-argon increase as the temperature increases (equations 9,10), this causes  $Nu_b$  for high heat flux conditions to be lower than  $Nu_b$  for low heat flux conditions. In the immediate thermal entrance region (x/D<5), the small rise in bulk gas temperature has not caused significant bulk property variation, and the  $Nu_b$  for the two heating rates are approximately the same. In the fully developed region, the large rise in bulk gas temperature has caused large property variations, and  $Nu_b$  are quite different. At x/D = 57,  $Nu_b$  for  $q_{Max}^+ = 0.0032$  is 29 percent lower than  $Nu_b$  for  $q_{Max}^+ = 0.0006$ .



## CONCLUSIONS

The object of this investigation was to study the momentum and heat transfer characteristics for turbulent flow of helium-argon mixtures in tubes. Experimental results were compared to existing experimental correlations, and to results from a numerical analysis. From this investigation the following conclusions have been made:

Existing experimental correlations, such as the
 Drew, Koo, and McAdams relation [34],

 $f = 0.0014 + 0.125 \text{ Re}^{-0.32}$ 

predict the helium-argon adiabatic friction factors within  $\pm 4.0$  percent for turbulent flow in tubes with Reynolds numbers between 31200 and 102000.

2. A correlation suggested by Taylor [36],

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 $f = (0.0014 + 0.125 \text{ Re}_w^{-0.32}) (T_w/T_b)^{-0.5}$ predicts average friction factors within ±10 percent for heated turbulent flow of helium-argon mixtures in tubes with inlet Reynolds numbers between 31200 and 102000.

3. Dittus-Boelter type correlations developed from air and helium experimental data

 $Nu = 0.021 \text{ Re}^{0.8} \text{ Pr}^{0.4}$ 

overpredict helium-argon Nusselt numbers for constant property, fully developed conditions by as much as 17 percent. An equation of similar form,

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# $Nu = 0.021 \text{ Re}^{0.8} \text{ Pr}^{0.55}$

but with the exponent of the Prandtl number changed to 0.55 predicts constant property Nusselt numbers of helium-argon mixtures within  $\pm 5.0$  percent. The range of Prandtl numbers was between 0.419 and 0.486, and the range of Reynolds numbers was between 31200 and 102000. 4. For the same range of Reynolds numbers and Prandtl numbers, the entrance and properties variation effects can be accounted for by using the equation Nu<sub>b</sub> = 0.021 Re<sub>b</sub><sup>0.8</sup> Pr<sub>b</sub><sup>0.55</sup>  $[(T_w/T_b)^{-0.4} + 0.85 D/x]$ . This equation predicted the bulk Nusselt numbers of helium-argon mixtures within  $\pm 15$  percent for x/D between 2.1 and 81.6 and a maximum wall-to-bulk temperature ratio of 1.82.

5. For helium-argon mixtures with molecular weights between 14 and 20, Prandtl numbers of 0.42, Reynolds numbers between 32000 and 55000, and constant property conditions the turbulent Prandtl number in the wall region, Pr<sub>t,w</sub> was determined to be 1.1 ± 0.1.
6. For helium-argon mixtures with molecular weights between 27 and 30, Prandtl numbers between 0.46 and 0.49, Reynolds numbers between 32000 and 102000, and constant property conditions Pr<sub>t,w</sub> was determined to be 1.0 ± 0.1.
7. For Reynolds numbers between 30000 and 100000, Pr<sub>t,w</sub> is a weak function of Reynolds number. For the Prandtl number range between 0.42 and 0.72, Pr<sub>t,w</sub> is

a strong function of Prandtl number, and decreases as Prandtl number increases.

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At maximum wall heating rates of  $q^+ = 0.0032$ 8.  $((T_w/T_b)_{Max} = 1.78)$ , Pr<sub>t,w</sub> determined from constant property conditions can be used in a variable properties numerical analysis to calculate  $Nu_{b}(x)$ . For the particular experimental run studied (Re, = 31200, Pr, = 0.486,  $Pr_{t,w} = Pr_{t} = 1.02$ , calculated  $Nu_{b}(x)$  agreed with measured  $Nu_{b}(x)$ , within the accuracy of the measured values. No separation of the helium-argon mixture was apparent.

APPENDICES

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Molecular Weight = 4.0026

Specific Heat at Constant Pressure = 1.24036 BTU/LB-R

Temperature	Viscosity	Conductivity	Sound Velocity
(F)	(LB/FT-HR)	(BTU/HR-FT-F)	(FT/SEC)
$\begin{array}{c} 40.00000\\ 70.00000\\ 100.00000\\ 100.00000\\ 100.00000\\ 220.00000\\ 2250.00000\\ 2250.00000\\ 2350.00000\\ 310.00000\\ 310.00000\\ 370.00000\\ 370.00000\\ 400.00000\\ 520.00000\\ 520.00000\\ 530.00000\\ 5550.00000\\ 5550.00000\\ 5550.00000\\ 5550.00000\\ 5550.00000\\ 5550.00000\\ 5550.00000\\ 5550.00000\\ 5550.00000\\ 5550.00000\\ 5550.00000\\ 5550.00000\\ 5550.00000\\ 5550.00000\\ 550.0000\\ 550.$	$\begin{array}{c} 4.5339250 \pm -02\\ 4.7348801 \pm -02\\ 4.7348801 \pm -02\\ 5.3488072 \pm -02\\ 5.30565294 \pm -02\\ 7.3257294 \pm -02\\ 7.22529494 \pm -02\\ 7.24233957494 \pm -02\\ 7.571795643574 \pm -02\\ 7.57179943574 \pm -02\\ 7.57179943574 \pm -02\\ 7.5502004894 \pm -02\\ 2.42339244 \pm -02\\ 2.42357497898 \pm -02\\ 2.4235397978 \pm -02\\ 9.477569898 \pm -02\\ 9.4779898 \pm -02\\ 2.42897998 \pm -02\\ 2.42837998 \pm -02\\ 2.42839239 \pm -02\\ 2.42837998 \pm -02\\ 2.4283928 \pm -02\\ 2.4283928 \pm -02\\ 2.4283928 \pm -02\\ 2.4284798 \pm -02\\ 2.448488 \pm -02\\ 2.44848 \pm -0$	$\begin{array}{c} 8 \cdot 4 3533951 \\ 9 \cdot 5 3710 \\ 9 \cdot 5 26293770 \\ 1 \cdot 021 \\ 1 \cdot 0253555 \\ 1 \cdot 0011 \\ 1 \cdot 0053605 \\ 1 \cdot 0011 \\ 1 \cdot 1835913500 \\ 1 \cdot 1835917550 \\ 1 \cdot 18358007755 \\ 1 \cdot 185580 \\ 1 \cdot 185580 \\ 1 \cdot 195580 \\ 1 \cdot 105881 \\ 1 \cdot 10581 $	$\begin{array}{c} 3.3 \\ 2162 \\ 3.3 \\ 2162 \\ 3.3 \\ 2162 \\ 3.3 \\ 3.4 \\ 3.5 \\ 3.5 \\ 3.5 \\ 3.5 \\ 3.5 \\ 3.5 \\ 3.5 \\ 3.5 \\ 3.5 \\ 3.5 \\ 5.$

Helium cont.

Temperature	Enthalpy
(F)	(ETU/LB)
$\begin{array}{c} 40.00000\\ 70.00000\\ 100.00000\\ 130.00000\\ 160.00000\\ 220.00000\\ 220.00000\\ 250.00000\\ 250.00000\\ 310.00000\\ 340.00000\\ 340.00000\\ 430.00000\\ 430.00000\\ 430.00000\\ 490.00000\\ 490.00000\\ 560.00000\\ 100.00000\\ 100.00000\\ 100.00000\\ 100.00000\\ 100.00000\\ 100000\\ 100000\\ 1000000\\ 1000000\\ 100000\\ 120.0000$	

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# Helium-Argon

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 Molecular Weight = 15.30

Specific Heat at Constant Pressure = 0.32449 BTU/LB-R

Temperature	Viscosity	Conductivity	Sound Velocity
(F)	(LB/FT-HR)	(BTU/HR-FT-F)	(FT/SEC)
$\begin{array}{c} 49.30000\\ 76.00000\\ 150.00000\\ 150.00000\\ 150.00000\\ 150.00000\\ 2250.00000\\ 2250.00000\\ 2250.00000\\ 2250.00000\\ 310.000000\\ 340.000000\\ 430.000000\\ 430.000000\\ 430.000000\\ 5550.000000\\ 5550.00000\\ 5550.00000\\ 5550.00000\\ 5790.00000\\ 5790.00000\\ 5790.00000\\ 5790.00000\\ 5790.00000\\ 5790.00000\\ 5790.00000\\ 5790.00000\\ 5790.00000\\ 5790.00000\\ 5790.00000\\ 5790.00000\\ 5790.00000\\ 5790.00000\\ 5790.00000\\ 5790.00000\\ 5700.0000\\ 5700.0000\\ 5700.0000\\ 5700.0000\\ 5700.0000\\$	$\begin{array}{c} 5 \cdot 6 + 3 + 3 + 3 + 2 + 3 + 3 + 3 + 3 + 4 + 2 + 3 + 3 + 4 + 2 + 4 + 4$	$\begin{array}{l} \bullet \bullet$	$\begin{array}{c} 1.6449620E+03\\ 1.7409220E+03\\ 1.7409220E+03\\ 1.7809220E+03\\ 1.7809220E+03\\ 1.7809220E+03\\ 1.78018072E+03\\ 1.870013062E+03\\ 1.970013062E+03\\ 1.970013062E+03\\ 1.970013062E+03\\ 1.9700130627E+03\\ 1.9700130627E+03\\ 2.00809766E+03\\ 2.00809766E+03\\ 2.00809765E+03\\ 2.0080855E+03\\ 2.0080855E+03\\ 2.0080855E+03\\ 2.00808558E+03\\ 2.00808558E+00\\ 2.008085858E+00\\ 2.008085858E+0\\ 2.008085858585858585858585858585858585858$

# Helium-Argon, M = 15.30 cont.

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Temperature (F)	Enthalpy (BTU/LB)
$\begin{array}{c} 40.30000\\ 70.00000\\ 130.00000\\ 130.00000\\ 130.00000\\ 190.30000\\ 220.90000\\ 220.90000\\ 220.90000\\ 250.00000\\ 310.00000\\ 340.00000\\ 370.00000\\ 430.00000\\ 430.00000\\ 430.00000\\ 520.0000\\ 520.00000\\ 520.00000\\ 520.00000\\ 520.00000\\ 520.000$	$\begin{array}{l} 1 \cdot 623 + 292 - 222 - 2$
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# Helium-Argon

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Molecular Weight = 15.83

Specific Heat at Constant Pressure = 0.31362 BTU/LB-R

Temperature	Viscosity	Conductivity	Sound Velocity
(F)	(LB/FT-HR)	(BTU/HR-FT-F)	(FT/SEC)
$\begin{array}{c} 46.90630\\ 70.30930\\ 100.00000\\ 130.00000\\ 130.00000\\ 150.99600\\ 226.00000\\ 226.000000\\ 2250.000000\\ 2259.000000\\ 310.000000\\ 440.000000\\ 440.000000\\ 440.000000\\ 440.000000\\ 440.000000\\ 550.000000\\ 440.000000\\ 550.0000000\\ 550.0000000\\ 550.000000\\ 550.000000\\ 550.0000000\\ 550.0000000\\ 550.0000000\\ 550.0000000\\ 550.000000\\ 550.000000\\ 550.0000000\\ 550.0000000\\ 550.0000000\\ 550.000000\\ 550.000000\\ 550.000000\\ 550.00000$	$ \begin{array}{l} 5 \cdot 2384611764-002\\ 5 \cdot 492742907464-002\\ 5 \cdot 9907464-002\\ 5 \cdot 9907464-002\\ 5 \cdot 9907464-002\\ 5 \cdot 9912844790464-002\\ 6 \cdot 1337066624-002\\ 6 \cdot 1337057686264-002\\ 7 \cdot 3537057686264-002\\ 7 \cdot 35574407765571264-002\\ 7 \cdot 3557491567703664-002\\ 7 \cdot 35574407765571264-002\\ 8 \cdot 35711765677033644-002\\ 1 \cdot 35769775777755712765626-002\\ 2 \cdot 2275576446622-002\\ 2 \cdot 227557646622-002\\ 2 \cdot 227557646622-002\\ 1 \cdot 1365776667754-002\\ 1 \cdot 1267667752762-002\\ 1 \cdot 12676677266-001\\ 1 \cdot 12676677266-001\\ 1 \cdot 1267666662\\ 1 \cdot 226667622\\ 2 \cdot 26677226\\ 1 \cdot 22666762\\ 1 \cdot 22666762\\ 1 \cdot 2266666\\ 1 \cdot 2266666\\ 1 \cdot 226666\\ 1 \cdot 226666$	$\begin{array}{l} 3.939613548448222222222222222222222222222222222$	$\begin{array}{c} 1.6171903 \pm 0.03\\ 1.60171903 \pm 0.03\\ 1.601719303 \pm 0.03\\ 1.75090351 \pm 0.03\\ 1.75099351 \pm 0.03\\ 1.8009051 \pm 0.03\\ 1.8009051 \pm 0.03\\ 1.8009051 \pm 0.03\\ 1.9007593097 \pm 0.03\\ 1.9007593097 \pm 0.03\\ 2.004332097 \pm 0.03\\ 2.035677 \pm 0.03\\ 2.03577 \pm 0.0$

## Helium-Argon, M = 15.83 cont.

Temperature	Enthalpy
(F)	(BTU/LB)
$\begin{array}{c} 40.00000\\ 70.00000\\ 100.00000\\ 130.00000\\ 140.00000\\ 220.00000\\ 250.00000\\ 250.00000\\ 250.00000\\ 250.00000\\ 340.00000\\ 340.00000\\ 340.00000\\ 490.00000\\ 490.00000\\ 550.00000\\ 520.00000\\ 510.0000\\ 510.0000\\ 510.0000\\ 510.0000\\ 510.0000\\ 510.0000\\ 510.0000\\ 510.0000\\ 510.0000\\ 510.0000\\ 510.0000\\ 510.000\\ 510$	$\begin{array}{c} 1.671330E+022\\ 2.671330E+02222\\ 1.87553169E+02222\\ 1.87553169E+0222222222222222222222222222222222222$

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# Helium-Argon

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Molecular Weight = 27.53

Specific Heat at Constant Pressure = 0.18034 BTU/LB-R

Temperature (F)	Viscosity	Conductivity	Sound Velocity
	(LB/FT-HR)	(BTU/HR-FT-F)	(FT/SEC)
$\begin{array}{c} 40.00000\\ 70.00000\\ 130.00000\\ 130.00000\\ 130.00000\\ 220.00000\\ 2250.00000\\ 250.00000\\ 250.00000\\ 310.00000\\ 340.00000\\ 340.00000\\ 340.00000\\ 430.00000\\ 520.00000\\ 550.00000\\ 550.00000\\ 550.00000\\ 550.00000\\ 550.00000\\ 550.00000\\ 550.00000\\ 730.00000\\ 730.00000\\ 730.00000\\ 730.00000\\ 730.00000\\ 730.00000\\ 730.00000\\ 730.00000\\ 730.00000\\ 750.00000\\ 750.00000\\ 750.00000\\ 750.00000\\ 750.00000\\ 750.00000\\ 750.00000\\ 750.00000\\ 750.00000\\ 750.00000\\ 750.00000\\ 750.00000\\ 750.00000\\ 120.0000\\ 120.0000\\ 120.00000\\ 120.00000\\ 120.00000\\ 120.000$	$\begin{array}{c} 5 \cdot 2563842 \pm -02\\ 5 \cdot 7827260 \pm -02\\ 6 \cdot 0365858 \pm -02\\ 6 \cdot 2346258 \pm -02\\ 6 \cdot 2346258 \pm -02\\ 6 \cdot 5271322 \pm -02\\ 7 \cdot 241997 \pm -02\\ 7 \cdot 2241997 \pm -02\\ 7 \cdot 68170221 \pm -02\\ 8 \cdot 5033497 \pm -02\\ 8 \cdot 5033499207 \pm -02\\ 1 \cdot 0023607 \pm -02\\ 9 \cdot 6702492 \pm -02\\ 1 \cdot 0036871 \pm -01\\ 1 \cdot 03398698 \pm -01\\ 1 \cdot 0336871 \pm -01\\ 1 \cdot 053686091 \pm -01\\ 1 \cdot 0336871 \pm -01\\ 1 \cdot 053686091 \pm -01\\ 1 \cdot 12526719 \pm -01\\ 1 \cdot 125366359 \pm -01\\ 1 \cdot 22566359 \pm -01\\ 1 \cdot 2875466 \pm -01\\ 1 \cdot 333180649 \pm -01\\ 1 \cdot 3380649 \pm -01\\ 1 \cdot 338$	$\begin{array}{c} \bullet 0.4345322E + 0.02\\ \bullet 1.402132E + 0.02\\ \bullet 2.4473336E + 0.02\\ \bullet 2.447336E + 0.02\\ \bullet 2.4486 + 0.02$	$\begin{array}{c} 1 & 22625300 \pm +03\\ 1 & 29784232 \pm +03\\ 1 & 39504232 \pm +03\\ 1 & 39513742 \pm +03\\ 1 & 365632255 \pm +03\\ 1 & 365632255 \pm +03\\ 1 & 462293707255 \pm +03\\ 1 & 462293707255 \pm +03\\ 1 & 462293707255 \pm +03\\ 1 & 462293775 \pm +03\\ 1 & 462293775 \pm +03\\ 1 & 46329375 \pm +03\\ 1 & 463293775 \pm +03\\ 1 & 555013197775 \pm +03\\ 1 & 555013197775 \pm +03\\ 1 & 555013197775 \pm +03\\ 1 & 55503197775 \pm +03\\ 1 & 550363379799 \pm +03\\ 1 & 55036337775 \pm +03\\ 1 & 550363377799 \pm +03\\ 1 & 550363377799 \pm +03\\ 1 & 5036337775 \pm +03\\ 1 & 503645577799 \pm +03\\ 1 & 5036457775 \pm +03\\ 1 & 5036457775 \pm +03\\ 1 & 5037799 \pm +03\\ 1 & 50364557775 \pm +03\\ 1 & 5037799 \pm +03\\ 1 & 50364557775 \pm +03\\ 1 & 5037799 \pm +03\\ 1 & 50364557775 \pm +03\\ 1 & 5037799 \pm +03\\ 1 & 5036455775 \pm +03\\ 2 & 503024799 \pm +03\\ 2 & 53365665 \pm +03\\ 2 & 22415573975 \pm +03\\ 2 & 224155933765 \pm +03\\ 2 & 22415593275 \pm +03\\ 2 & 2 & 22415593275 \pm +03\\ 2 & 2 & 2 & 22415593275 \pm +03\\ 2 & 2 & 2 & 2 & 2 & 2 & 2 & 2 & 2 & 2 $

# Helium-Argon, M = 27.53 cont.

Temperature	Enthalpy
(F)	(BTU/LB)
$\begin{array}{c} 40.00000\\ 100.00000\\ 130.00000\\ 160.00000\\ 160.00000\\ 190.00000\\ 250.00000\\ 250.00000\\ 310.00000\\ 310.00000\\ 340.00000\\ 340.00000\\ 490.00000\\ 490.00000\\ 490.00000\\ 490.00000\\ 550.00000\\ 100000\\ 100000\\ 10$	9.5593245231 $\pm\pm\pm02$ 48 $\pm\pm022$ 48 $\pm5025$ 48 $\pm5025$ 48 $\pm5025$ 48 $\pm5025$ 48 $\pm5025$ 48 $\pm5025$ 48 $\pm5025$ 48 $\pm5027$ 48 $\pm5027$ 49 $\pm5027$ 49 $\pm5027$ 40 $\pm2027$ 40 $\pm2027$

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## Helium-Argon

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Molecular Weight = 29.70

Specific Heat at Constant Pressure = 0.16716 BTU/LB-R

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Temperature (F)	Viscosity (LB/FT-HR)	Conductivity (BTU/HR-FT-F)	Sound Velocity (FT/SEC)
$\begin{array}{c} 40.90000\\ 100.00000\\ 100.00000\\ 100.00000\\ 100.00000\\ 220.00000\\ 250.00000\\ 250.00000\\ 250.00000\\ 310.00000\\ 340.00000\\ 400.00000\\ 400.00000\\ 400.00000\\ 400.00000\\ 400.00000\\ 520.0000\\ 520.00000\\ 520.0000\\ 520.0000\\ 520.0000\\ 520.0000\\$	$5 \cdot 24186661E - 02$ $5 \cdot 5089766E - 02$ $5 \cdot 50897667E - 02$ $5 \cdot 7697317E - 02$ $6 \cdot 27328E - 022$ $5 \cdot 5154783097E - 022$ $6 \cdot 5154783097E - 022$ $7 \cdot 246622596E - 022$ $7 \cdot 246622596E - 022$ $7 \cdot 246622596E - 022$ $7 \cdot 24662596E - 022$ $7 \cdot 24662596E - 022$ $7 \cdot 24662596E - 022$ $8 \cdot 3099518E - 022$ $8 \cdot 3099508396E - 022$ $8 \cdot 309950839088E - 022$ $9 \cdot 246643E - 022$ $9 \cdot 246643E - 022$ $9 \cdot 24664503E - 022$ $9 \cdot 24664503E - 022$ $9 \cdot 2466503E - 022$ $1 \cdot 0239698578E - 022$ $1 \cdot 0239698578E - 022$ $1 \cdot 0239698578E - 011$ $1 \cdot 12559378E - 011$ $1 \cdot 12559378E - 011$ $1 \cdot 12559578E - 011$ $1 \cdot 12559578E - 011$ $1 \cdot 12559578E - 011$ $1 \cdot 12559578E - 011$ $1 \cdot 25728548 - 011$ $1 \cdot 28823992 - 011$ $1 \cdot 3035776 - 011$ $1 \cdot 30357776 - 011$ $1 \cdot 30577776 - 011$ $1 \cdot 30577776 - 011$ $1 \cdot 3057$	$ \begin{array}{c} 1.6071936 = -022\\ 1.6936 9202 = -002\\ 2.6022 = -0022\\ $	1.1000 1.1000
1496.00000	1.3339004E-01	4.48717396402	2.21777886403

Helium-Argon, M = 29.70 cont.

Temperature (F)	Enthalpy (BTU/LB)	
$\begin{array}{c} 40.00000\\ 70.00000\\ 100.0000\\ 100.0000\\ 130.00000\\ 220.00000\\ 220.00000\\ 230.00000\\ 230.00000\\ 310.00000\\ 400.00000\\ 400.00000\\ 400.00000\\ 400.00000\\ 400.00000\\ 400.00000\\ 400.00000\\ 400.00000\\ 610.00000\\ 580.00000\\ 580.00000\\ 610.00000\\ 100.00000\\ 100.00000\\ 100.00000\\ 100.00000\\ 1150.00000\\ 1150.00000\\ 1210.00000\\ 120.00000\\ 120.00000\\ 120.00000\\ 1300.00000\\ 10000000\\ 1000000\\ 1000000\\ 10000000\\ 1000000\\ 1000000\\ 1000000\\ 1000000$	$8 \cdot 3557257200000000000000000000000000000000$	
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#### APPENDIX B. EXPERIMENT

### Apparatus

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A schematic diagram of the experiment is shown in Fig. Bl. The helium and helium-argon mixtures were bought from manufacturers in high pressure bottles. The air was obtained from a large storage tank that was replenished by a compressor. Two regulators were used to reduce and stabilize the pressure. A Brooks rotometer was used to obtain a rough measurement of the flow rate. A Bourdon tube Heise gage measured the pressure just downstream of the rotometer. A small tank constructed to mix the gas, and instrumented with a thermocouple measured the inlet stagnation temperature. A sketch of the test section from the inlet to just below the outlet mixing tank, and displaying the location of thermocouples, pressure taps, electrodes, and voltage taps is shown in Fig. B2.

Power was measured using a Fluke differential voltmeter and Weston ammeter in the same manner as Perkins et al. [27]. Whenever possible the power supply described by Perkins et al. [27] was used. When it did not supply sufficient power an a.c. Lincoln welder (Model TM-500/500) was used. To determine a power factor when the welder was used, the power measured with a Weston watt meter was compared to that calculated from voltage (Fluke voltmeter)



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Fig. 32. Diagram of test section.

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and current (Weston ammeter) measurements. A power factor of 1.0 was used since the two measurements agreed within 3.3 percent.

Thermocouple output was measured with a Hewlett Packard digital voltmeter. An ice bath was used as a reference for all of the thermocouples. Thermocouples were selected for measurement using a manual switch. The numbered thermocouples on Fig. B2 were used in the computer program that reduced the experimental data for the heated runs. This computer program is described by a number of people [30,31,32,44]. The unnumbered thermocouples were used to determine the amount of preheating of the gas before it entered the heated section.

Pressure drop in the test section was measured with Meriam 60 inch vertical water or mercury manometers. Inlet static pressure was measured with a Bourdon tube Heise gage.

After the gas passed through the test section, it was cooled by a chilled water counterflow heat exchanger. The valve used to control the flow rate was located just downstream of the heat exchanger. The heat exchanger was necessary in order that a Meriam laminar flow element could be used. The laminar flow element was used to obtain an accurate measurement of the flow rate. A Meriam 60 inch inclined water manometer with a 10 inch range was used to measure pressure drop across the laminar flow

element. The temperature of the gas in the flow element was measured with a thermocouple, and the pressure was measured with a Meriam 60 inch vertical mercury manometer. The accuracies of the instruments used in this investigation are listed in appendix C (Table Cl).

#### Procedure

Before any experimental runs with gas flow were performed, the test section was heated without flow in order that the heat lost to the environment and the resistance of the test section could be determined. These items are discussed in detail in the following two sections.

The system was purged and all of the instruments zeroed before each set of experimental runs. The purging was done by pressurizing the system to approximately 100 psig with the gas to be used. The system was then allowed to blow down to approximately 10 psig. This sequence was performed four times.

The desired inlet Reynolds number was established by adjusting the pressure level and mass flow rate. Before power was supplied to the test section, measurements were taken so that calculation of the adiabatic friction coefficient was possible. The measured adiabatic friction coefficients were compared to the Drew, Koo, and McAdams correlation [34] (equation 11), and agreement ensured that the pressure measurements, mass flow rate measurements, and mixture molecular weights were correct.

The test section was then heated to the desired level. Since a small period of time elapsed while the thermocouples were manually recorded, pressure drop, pressure level, voltage, current, and mass flow rate measurements were taken before and after the thermocouple measurements. The average of the two measurements was used for data reduction. The inlet Reynolds number was maintained approximately constant while the test section wall temperature was varied by varying the power input. Measurements were taken for a number of different power inputs.

### Heat Loss Calibration

In order to calculate the heat transfer coefficient, the heat addition to the gas, q'<sub>gas</sub>, must be determined. If an energy balance for a small section of the tube is performed, the result is:

$$q'_{gas} = q'_{gen} - (q'_{cond} + q'_{loss}).$$
(B1)

The heat generated in the small section of the tube,  $q'_{gen}$ , is:

$$q'_{gen} = I^2 R'.$$
 (B2)

The current was measured, and the calculation of the resistance per unit length is discussed in the next section. The axial heat loss due to conduction is:

$$q'_{cond} = -K A_{cs} d^2 T_{w}/dx^2.$$
 (B3)

The second derivative was determined using a numerical parabolic fit described by McEligot [31]. The variation of thermal conductivity with temperature for Hastelloy-X was determined from data supplied by the manufacturer [46].

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K = [5.1 + (0.00622)(T)](Btu/hr-ft-F)(B4) (T in degrees Fahrenheit)

To determine the heat loss,  $q'_{loss}$ , the test section was heated at different levels without gas flow. A program described by Coon [32] was used to calculate the heat loss at each thermocouple. The heat loss was determined as a function of the tube wall and environment temperature difference. The environment temperature was measured with a thermocouple a few inches away from the test section. Figures B3, B4, B5, and B6 show the results for each thermocouple (thermocouple 4 is not included). Except for thermocouple three and four, the data for each thermocouple was fitted with an equation of the form:  $q'_{loss} = C_1(T_w - T_w) + C_2(T_w - T_w)^2 + C_3(T_w - T_w)^3$  (Btu/hr-ft) (B5)  $(T_w - T_w$  in degrees Fahrenheit)

The numerical values of  $C_1^{}$ ,  $C_2^{}$ , and  $C_3^{}$  are listed below.

Thermocouple	° <sub>1</sub>	с <sub>2</sub>	° <sub>3</sub>
2	2.43E+00	9.43E-03	-2.53E-05
5	1.22E-01	1.84E-04	-1.64E-08
6	9.83E-Ò2	2.30E-04	-6.57E-08
7	8.84E-02	1.66E-04	-3.79E-09
8	8.30E-02	1.35E-04	3.07E-08
9	8.16E-02	1.03E-04	6.04E-08
10	7.89E-02	1.12E-04	4.96E-08
11	7.70E-02	8.82E-05	7.13E-08
12	7.40E-02	1.03E-04	5.50E-08
13	7.54E-02	8.08E-05	7.59E-08
14	7.31E-02	9.77E-05	5.73E-08
15	8.09E-02	8.07E-05	7.46E-08
16	1.18E-01	1.59E-04	5.59E-08
17	2.76E-01	1.14E-03	1.19E-06

The data for thermocouple three was fitted with a straight line determined using the method of least squares.






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$$q'_{loss}(T.C. 3) = 0.91(T_w - T_w) - 2.20(Btu/hr-ft) \quad (B6)$$

$$(T_w - T_w \text{ in degrees Fahrenheit})$$

The data for thermocouple four was scattered and a representative curve could not be fitted. For this reason the heat loss at this thermocouple was neglected. This introduced only a small error since the heat loss for thermocouple four, even at the highest heating rates, was less than five percent of the heat added to the gas. The fitted curves are also shown on Figures B3, B4, B5, and B6.

#### Measurement of Test Section Resistance

The variation of resistance with temperature was measured by heating the test section without gas flow in the same manner as was done for the heat loss calibration. Using the thermocouple wires as voltage taps, a measurement of the voltage drop between thermocouple 14 and the lowermost thermocouple on the tube was taken. Another measurement of the voltage drop between thermocouple 12 and the lowermost thermocouple was taken. The difference between these two measurements gave the voltage drop between thermocouples 12 and 14. The section of the tube between thermocouples 12 and 14 was used since the wall temperature for this length was approximately constant.

For a particular power setting the voltage drop discussed in the previous paragraph, the current, and the average of the temperatures at thermocouples 12, 13, and

14 were recorded. Using these measurements the resistance per unit length was determined as a function of temperature. The results are shown on Fig. B7. Also shown on Fig. B7 is the line that was used to approximate the variation of resistance with temperature.

$$R' = [3.98 \times 10^{-4} (T) + 4.745] (m\Omega/in) (B7) (T in degrees Fahrenheit)$$

### Meriam Laminar Flow Element Calibration

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The laminar flow element was calibrated using a Parkinson-Cowan Type Dl positive displacement flow meter as a standard. It is specified to have 1/2 percent accuracy at ambient conditions and was calibrated by Tucson Gas and Electric before being used. Meriam [45] suggests the following equation for the laminar flow element:

$$\Delta P = A'Q\mu L/D^4 + B'\rho Q^2/D^4. \qquad (B8)$$

A' and B' are the constants to be determined by calibration. Since the length, L, and the hydraulic diameter, D, of the laminar flow element passages remain constant, they can be incorporated into new calibration constants, A and B.

$$A = A'L/D^4 \qquad B = B'/D^4$$

Equation B8 now becomes:

$$\Delta P = AQ\mu + B\rho Q^2. \tag{B9}$$

If this equation is solved for Q the result is:

$$Q = [-A\mu/B + \sqrt{(A\mu/B)^2 + 4\rho\Delta P/B}]/2\rho.$$
(B10)

If both sides of this equation are multiplied by the density,  $\rho$ , the result is:



$$\dot{m} = \rho Q = 1/2[-A\mu/B + \sqrt{(A\mu/B)^2 + 4\rho\Delta P/B}].$$
 (B11)

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Mass flow rate measurements were taken simultaneously with the positive displacement flow meter and the laminar flow element. Both air and helium were used. These gases bound the range of helium-argon mixture molecular weights that were used. From these measurements, A and B were determined so that the maximum difference between the positive displacement flow meter measurements and the laminar flow element measurements was 1.5 percent. The numerical values of A and B were 1375 and 672, respectively, if the parameters in equation Bll have the following units.

> m - lb/sec  $\mu$  - lb/hr-ft  $\rho$  - lb/ft<sup>3</sup>  $\Delta P$  - inches of water

### APPENDIX C Uncertainty Analysis

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An analysis to determine the uncertainty of the results calculated from the measured experimental data was performed. The uncertainties of the directly measured quantities were determined from manufacturers' specifications and experience. Table Cl lists the uncertainties of the instruments used in this investigation. The error propagated to the calculated results from the uncertainties in the measured quantities was determined using a method described by Bottaccini [47]. The general equation used was

$$\sigma_{\mathbf{Z}}^{2} = \left(\frac{\partial Z}{\partial \mathbf{Y}_{\mathbf{i}}}\right)^{2} \sigma_{\mathbf{Y}_{\mathbf{i}}}^{2} + \left(\frac{\partial Z}{\partial \mathbf{Y}_{2}}\right)^{2} \sigma_{\mathbf{Y}_{2}}^{2} + \dots + \left(\frac{\partial Z}{\partial \mathbf{a}_{\mathbf{i}}}\right)^{2} \sigma_{\mathbf{a}_{\mathbf{i}}}^{2} + \left(\frac{\partial Z}{\partial \mathbf{a}_{2}}\right)^{2} \sigma_{\mathbf{a}_{2}}^{2} + \dots$$
(C1)

 $\sigma_x$  is the variance or standard deviation of the xth quantity. Z is the calculated quantity, and Y<sub>i</sub> and a<sub>i</sub> are the measured values and system parameters used to calculate Z.

To illustrate the above technique, a simple example will be done. The power supplied to the tube can be determined using the relation (assuming a power factor of one)

$$P = E I.$$
 (C2)

If equation Cl is applied, the error or variance in the power caused by uncertainties in the voltage and current measurements is

$$\sigma_{\rm p} = \sqrt{({\rm I})^2 \sigma_{\rm E}^2 + ({\rm E})^2 \sigma_{\rm I}^2}.$$
 (C3)

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The result can be presented in the following form so that

the percent of uncertainty can easily be determined.

$$\sigma_{\rm P}/{\rm P} = \sqrt{(\sigma_{\rm E}/{\rm E})^2 + (\sigma_{\rm I}/{\rm I})^2}$$
 (C4)

# TABLE C1 Uncertainties of Measured Values

Measured quantity	Instrument	Uncertainty
Current	Weston 370 AC/DC ammeter	±0.25% of full scale
Voltage	Fluke 883AB differen- tial voltmeter	±0.1% of input
Mass flow rate	Meriam 50MH10-1 lam- inar flow element	±1.5% of flow rate
Wall and inlet bulk temperature	Premium grade chromel- alumel thermocouples	±2 <sup>°</sup> F, 3/8% of read- ing above 535°F
Thermocouple location Pressure tap location	Gaertner M911 Cathetometer	±0.1mm
Diameter	Manufacturers' specifications	±0.001 in.
Pressure	12 inch Heise gage	±0.15 psi
Pressure drop	60" Meriam 30EB25 vertica] H_O manometer 60" Meriam 30EB25 vertica] Hg manometer	±0.05 in-H <sub>2</sub> 0 L ±0.05 in-Hg

Table C2 lists the percentage uncertainty in the measured bulk Nusselt numbers for two representative helium-argon runs. The dominant uncertainty in the bulk Nusselt number is the bulk stagnation temperature. For convenience no uncertainty was included for the gas properties. The values used were assumed to be precise.

# TABLE C2 Percentage Uncertainties in the Measured Bulk Nusselt Numbers of Helium-Argon

Run Mclecular weight Re <sub>i</sub> (T <sub>w</sub> /T <sub>b</sub> ) <sub>MAX</sub>	126H 15.30 56100 1.17	131H 15.30 54700 1.77
x / D	Percentage	Uncertainty
1.17	13	10
2.07	8	6
4.14	6	4
8.13	5	3
24.52	4	3
40.79	4	3
56.88	4	3

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### APPENDIX D Helium-Argon Experimental Data

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The headings and their definitions used in the listing of the adiabatic friction data are below.

Heading	Definition						
Run	Experiment run number						
Date	Date on which experimental run was made						
Gas	Gas used in the experiment						
Molec. wt.	Molecular weight						
T <sub>i</sub>	Inlet mixer temperature						
• m	Gas flow rate						
Re <sub>i</sub>	Inlet Reynolds number						
P <sub>1</sub>	Static pressure at inlet pressure tap						
P <sub>2</sub>	Static pressure at outlet pressure tap						
Static Mach <sub>l</sub>	Static mach number at inlet pressure tap						
Static Mach <sub>2</sub>	Static mach number at outlet pressure tap						
fad	Adiabatic friction factor						

Table D-1 HELIUM-ARGON ADIABATIC FRICTION FACTOR DATA

Static f <sub>ad</sub> Mach <sub>2</sub>
Mach <sub>1</sub> Mac
P2 (psia)
P <sub>l</sub> (psia)
Re <sub>1</sub> .
m (1b/hr)
$\begin{pmatrix} \mathbf{T}_{1} \\ (0_{\mathbf{F}}) \end{pmatrix}$
Molec. Wt.
Gas
Date
Run

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The headings and their definitions used in the listing of the heated flow data are below. The headings that are self-explanatory, or that were used in the listing of the adiabatic friction data are not included.

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Heading	Definition
TIN	Inlet mixer temperature
TOUT	Outlet mixer temperature
I	Alternating current
Е	Voltage drop between voltage taps
TC	Thermocouple number
X/D	Corresponds to x/D in text
TW	Inside tube wall temperature
TW/TB	Wall-to-bulk temperature ratio
QGAS	Wall heat flux
Q <sup>+</sup>	Non-dimensional heat flux parameter. Corresponds to q in text.
PT	Pressure tap: 1-inlet, 2-outlet
ТВ	Bulk static temperature

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17	98.0	435.6	1.165 P1	40749. T X/D	STATIC PRESS. (PSIA)	92.24 TW/TB	TB PSESS	.001074
				90.1	80-3 67.8	1.10	291.4 .1996+0	1
	Y A.	ERAJE BULK 4812	AT REYNOLDS	VERAGZ PAR Avera	AMETERS BETWE GE WALL REYNO 36951.	EN PRESSUR LDS AV	E TAPS ERAGE PRICTION PA: .00554	ELDE
cin = PB,im	AV 72.8 P = 41	BUN 90H TOUT = 5 , Ga/Reso	RETNOLDS 9. 19. DATE 10 14.7.4 P. M. 2 = 166	VERAGZ PAR AVZRA 0/08/75, G ASS FLOW R E-03, MACE	AMETERS BETWE GE WALL REYNO 36951. AS HE-AR, MOL ATE = 25.1 LB (2) = .204,	EN PRESSUR LDS AV ECULAR WT. //HR, I = 1 MACH (16) =	E TAPS ERAGE PRICTION PA: .))554 42.7 AMPS, E = 3 .340, T,SURE =	- 523 Y2LI 79.5 P
TP C 234557890112	AV 72.8 7 = .41 X/D 1.12 2.1 4.1 16.4 232.9 3403.9 57.0	ERAJE BULK BULK BUN 90 TGB TGB TGB TGB TFP 195.00 384.92 5593.03 5593.03 5593.03 744.68	E DATE 10 REYNOLDS 9. 19. 10. 10. 10. 10. 10. 10. 10. 10	VERAGZ PAR $\lambda$ VERA $\lambda$ VERA $\lambda$ VERA $\lambda$ VERA $\lambda$ VERA $\lambda$	AM ETERS BETWE GE WALL ASYNO 36951. AS HE-AR, MOL ATE = 25.1 LB (2) = .204, HL/GGAS .073 .023 .025 .025 .025 .025 .029 .034	EN PRESSUR LDS AV ECULAR WT. /HR. (16) = SULK T 2335.332 101.328 73.080 59.925 50.925 50.51	E TAPS ERAGZ PRICTION PA: .JJ554 42.7 AMPS, E = 3 .J42.7 AMPS, E = 3 .J40, T, SURE = 03AS STU/HAPT2 116516.4 114231.3 122912.J 124446.7 125787.1 126740.6 126207.1 126315.1 126508.3	523 Y3L2 73.5 P 2+ .002255 .102455 .102455 .012455 .012465 .012465 .012465 .012465 .012465 .012465 .012465 .012465 .012465 .012465 .012465
IN I 23456789011234567	AV 72.8 41 X/D 12.11 164.465 34.65 75.4 5675.4 5675.3 50 8 30 8 30 8 30 8 30 8 30 8 30 8 30	ERAJE BULK BULK BULK BULK BULK BULK BULK BULK	RETNOLDS 8. 8. 19. 19. 10. 10. 10. 10. 10. 10. 10. 10	VERAGZ PAR AVERA VERAGZ PAR AVERA VERAGE 0/08/75, GR E-03. REC 8 DLLK 8 DLLK 8 DLLK 8 DLLK 8 DLLK 8 DLLK 8 DLLK 9 568271 55782 558271 55782 5	AM ETERS BETWE GE WALL REYNO 36951. AS HE-A8, MOL ATE = 25.1 LB (2) = .204, HL/GGAS .073. .023. .0223. .0225. .024. .025. .025. .025. .026. .036. .036. .036. .036. .042. .046. .437. STATIC	EN PRESSUR LDS AV ECULAR WT. HR.H [16] SUSSELT9 LK T 2322-233 101-328 65-38 556-38 556-38 556-55 552-55 552-55 556-55 552-55 555	E TAPS ERAGZ PRICTION PA: .JJ554 .JJ554 .J2554 .J2554 .J2554 .J2554 .J2556 .J25767. .J2516.4 .J2516.4 .J2516.4 .J25787.1 .J25797.1 .J257	- 523 Y2L - 523 Y2L - 79.5 P - 2.4 - 30224457 - 30224459 - 3022459 - 302459 - 30259 - 302

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TIY = Pa,Iy	72.3	RUN 938 F. TOUT = 3 19, GR/RESQ	6, DATE 10 48.8 P, MA = .187E	/13/75, G SS PLOW a -03, MACH	AS HE-AR, MOL ATE = 14.1 LE (2) = .107,	ECULAR #T. 3/88, 1 = MACH (16) =	= 15.33 80.0 AMPS, E = .134, T,SU2E	4.535 ¥2LIS
E 2345678901234567	L 1211145487881310 L 124564208643108 12344597891310	3121713593337596 3121713593337596 12270350757574192211 122223355702363911	T T B 1.0993 0993 1.12462 1.231266 1.322813 1.322813 1.322813 1.322813 1.322813 1.22256 1.22256 1.22256 1.22256 1.22256 1.22256 1.22256 1.22256 1.22256 1.22256 1.22256 1.22256 1.22256 1.22256 1.22256 1.22256 1.22256 1.22575 1.22575 1.22575 1.22575 1.22575 1.22575 1.22575 1.22575 1.22575 1.22575 1.22575 1.22575 1.22575 1.22575 1.22575 1.22575 1.22756 1.22756 1.22756 1.22756 1.22756 1.22756 1.22756 1.22756 1.22756 1.22756 1.22756 1.22756 1.22756 1.22756 1.22756 1.22756 1.22756 1.22756 1.22756 1.2077 1.2077 1.22756 1.20777 1.20777 1.20777 1.207777 1.207777 1.20777777777777777777777777777777777777	BY165231. BY16523. BY16523.	HL/QGAS 067 140 024 0224 0224 0225 0226 0226 0228 0228 0228 0228 0228 0228	N1 N1 N1 N1 N1 N1 N1 N1 N1 N1 N1 N1 N1 N	OGAS BTU/MBF12 346390.1 37309.7 38355.7 386705.0 197777.9 386705.0 197777.9 3867057.9 38674.3 386674.3 386599.9 393439.8 31339.4	2+ .))12311 .))12311 .))12357 .))13567 .))13567 .))13663 .))13663 .))13663 .))13663 .))13663 .))13663 .))13663 .))13663 .))13665 .))13665 .))13665 .))13665 .))13665 .))13557 .))13665 .))13665 .))13665 .))13655 .))13655 .))13655 .))13655 .))13655 .))13655 .))13655 .))13655 .))13655 .))13655 .))13655 .))13655 .))13655 .)]13655 .)]13655 .)]13655 .)]13655 .)]13655 .)]13655 .)]13655 .)]13655 .)]13655 .)]13655 .]]13655 .]]13655 .]]136555 .]]136555 .]]136555 .]]136555 .]]136555 .]]136555 .]]136555577 .]]136555577 .]]136555577 .]]136555777 .]]1365557777777777777777777777777777777777
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TIN = Pr,in	λ 72.8 1 Ξ .4	VERAGE BOLK 2750 RUN 9461 F. FOOT = 50 18, GR/RES2	RETNOLDS	13/75, GJ S PLOW RJ 03, MACH	AMETERS BETWE GE WALL REYNO 22429- AS HE-AR, MOL AE = 14.1 LB (2) = -107, 1	ELULAR T. /HR, I = 1 AACH(16) =	E TAPS ERAGE FRICTION .JO638 .J0638 .J0638 .J0638 .E TS.83 	7.395 VQLI3 73.5 P
r 2345679901234567	X 12344567899004544 12344567899004544	0996225993483277	F¥/TB 1.2349 1.5550 1.5550 1.5550 1.55561 1.55561 1.5460 1.55561 1.43055 1.3330 1.5130 1.515561 1.3330 1.515561 1.3330 1.515561 1.3330 1.515561 1.5155561 1.51555561 1.51555561 1.51555561 1.51555561 1.51555561 1.51555561 1.51555561 1.51555561 1.51555561 1.51555561 1.51555561 1.51555555561 1.51555561 1.515555561 1.515555561 1.515555561 1.515555561 1.51555555555555555555555555555555555	BULD499113 BULD499113 BU152595913 BU152595913 BU152595913 BU111109754426 CO149748420697 CO149748420697 CO149748420 CO149748420 CO149748420 CO149748420 CO1497484 CO1497484 CO1497484 CO149748 CO	HL/QGAS .093 .143 .037 .036 .037 .0341 .0443 .0443 .0557 .059	T4295787950721835 305776544621835 30577654432224460835	277519474 377519499 885596 885596 885596 885596 885596 885596 895517 895178 805	Q+ .)2775 .)22663 .)22984 .)22984 .)22984 .)22984 .)2292984 .)2292984 .)2292984 .)2292984 .)22928 .)20330222 .)2330222 .)2330222 .)2330222 .)2330225 .)233025 .)2325 .)233025 .)2355 .)23555 .)23555 .)23555 .)23555 .)235555 .)2355555 .)23555555555555555555555555555555555555
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	٨	VERAGE BULK 2495)	AVI REYNOLDS	AVERAGE PARI	A 27295 BETWE Se Wall Beyno 17318.	EN PRESSURI LDS AVI	E TAPS ERAGE PRICTION . J0655	PACIOR

BULK 8 27 NOLDS 59554. 68401. 67285. 65254. BULK NUSSELT 285.58 135.69 115.83 97.85 03AS BIŪ/HBPT2 150418.8 147299.2 157043.5 94938176666396807 155195405959410621 23549540505937906 23549540505937906 108 108 108 108 02417 80900044690m5997 0000000000003335997 ·211146599004644 124964208753108 12344567990 935084 558 1.744 45327098 49083. 47376. 45758. 44325. 42947. 41994. 39.3 99.3 84.5 11.6 1.465 x/d 90:4 STATIC PRESS. (PSIA) 109.4 88.7 PRESS DEFECT -.488E-02 .275E+01 PT T¥/TB TB (F) 68.2 537.9 1.20 12 URE TAPS AVERAGE FRICTION FACTOR .00517 AVERAGE PABAMETERS BETWEEN AVERAGE BULK BEYNOLDS · AVERAGE WALL BEYNOLDS 56295. 36396. PRESSURE RUN 97HP, DATE 10/20/75, GAS HE-AR, SOLECULAR WT. = 15.03 TIN = 72.3 P, FOUT = 290.2 P, MASS PLOW RATE = 35.3 LB/HR, I = 117.2 AMPS, Z = 5.793 VOLES PR.IN = .410, JR/BESQ = .336E-04, MACH(2) = .238, MACH(16) = .337, T,SURR = 75.3 F BULK REYNOLDS 80349. 90006. 79735. 90LK NUSSELT 333.86 156.22 132.51 112.68 Q3AS 0/H8772 9993.4 8204.2 HL/QGAS Q+ X/D ΓW TW/TB 8Ţ 5266303922192208 11914111608328687 119256925690356887 112222353334444444 001124 0011099 001160 001171 001171 98473239 9847323344677701 9001123344677701 9000111239 1211145487881310 186 185 187 187 187

RUN 98HF, DATE 10/20/75, GAS HE-AR, NOLECULAR WT. = 15.83 TIN = 73.2 F, FOUT = 570.4 F, MASS FLOW RATE = 30.8 LB/HR, I = 160.9 AMPS, E = 9.670 VOLTS PR,IN = .413, GR/RESQ = .213E-03, MACH(2) = .204, MACH(16) = .340, F,SUER = 79.0 F

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4496.5 4511.9 4428.2 3508.6 4059.2 10 PRESS DEFECT -.473E-02 .193E+01 STATIC PRESS. (PSIA) 106.9 98.5 TW/TB TB (F) 63.5 272.3 PT X/D 90:1 1.10 12 AVERAGE BULK BEYNOLDS AVERAGE FAR METERS BETWEEN PRESSURE TAPS AVERAGE BULK BEYNOLDS AVERAGE FALL REYNOLDS AVERAGE PRICTION FACTOR 71091. 53547. .00491

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67 890	16.4 16.4 124.05 124.59	573.2 663.5 726.4 776.3 920.5	1.832		76799. 72178. 68290. 64994. 61915.	- 033 - 034 - 037 - 040 - 043	111.67 96.47 88.45 83.42 79.88	100 100 100 100	727.5 240.4 391.2 492.6 492.6	.)0256) .002573 .)02577 .)0258) .00258)
1234567	45-53 104644 198	903.0 939.0 9939.0 993.0 993.0 786.6	1.599 1.599 1.5562 1.5562 1.205		56964. 54893. 52922. 51185. 49522. 48437.	- 002 - 002 - 00557 - 00958 - 00958 - 00958	74.035 74.035 722.78 724.48	100 100 100 100 99	**** 362.5 379.9 180.5 184.0 569.4 569.4	
				27	¥ ZD	STATIC	TW/TB	TB	PRESS	
				••	~/ •	D8255./95TI			072207	
TN = 9.IN	71.0 F = .43	RUN 10 TOUT = 5	3H DATE 471.8 F 5	1 2 0 1/2 9 3 E-	23/76, G	PHESS. (PSIA) 10.6 97.1 AS H2-AR, NOL ATE = 36.2 LB (2) = .173.	1.20 1.46 2CJLAB WT /HR.I.16	541.2 541.2 113.8 Au	D372CF - 470E-02 - 245E+01	.715 ∛OL 90-0 F
TN = 2,1%	71.0 g = .48	RUN 10 Tour = 6, ga/res	3H DATE 471.8 F. 2 = .5	1 2 0 1/ 5 32-		AS HE-AR, MOL ATE = 36.2 LE (2) = .173,	1.20 1.46 2CJLAR HT /HR I T MACH (16)	541.2 541.2 113.8 29.7 113.8 ME = .251,	DEFECT - 470E-02 245E+01 - 245E+01	.715 ¥01 90.5 ₽
23 29,1N	71.0 5 - 43 X/D	RUN 10 TOUT = 6, 32/RES IN IS 191.6	3H DATE 2 = 5 TW/TB 1-234	1 2 0 1/2 9 32-	23/76, G 35. FLOM R 03, MAGH 25YNOLDS 315023	PHESS. (PSIA) 10. (PSIA) 97.1 AS HE-AR, MOL ATE = 0.173, HL/QGAS - 109	1.20 1.46 2001AB b1 MR.1 1 MACH (16) NUSSELT 126.29	57.0 541.2 113.8 29.7 113.8 AME = 251, 8 TO	D 2 72CF - 4 70E-02 245E+01 - 245E+01 - 245E+01 - 245E+01 - 245E+01 - 245E+01 - 245E+01 - 245E+01 - 245E+01 - 245E+01 - 245E+01	-715 VQL 90-3 P 2+ -00184
N. I.N. I.N. I.N. I.N. I.N. I.N. I.N. I	71.0 F = .40 X/D 1.2 2.1 4.1	RUN 10 TOUT = 5, 32/RES 191.6 299.7 354.5 412.8	3H DATE 471.8 2 = .5 TW/TB 1.234 1.428 1.521 1.604 1.604 1.604		23,76, G 35,FLOW P 03, MAGH 301LK 25YNOLDS 31602. 30522. 77421. 77193.	PHESS. (PSIA) 10.6 97.1 AS HE-AR, MOL ATE = 36.2 LB (2) = .173, HL/QGAS .109 .146 .035 .029	1.20 1.46 2.201.AB WT MR.I. MACH (16) 9ULK NUS26.299 169.154 127.999	57.0 541.2 113.8 29.7 113.8 AMP 251, 8 TO 71 69 77	DEFECT - 475E+01 - 245E+01 - 245E+0000+000+000+000+000+000+000+000+000+	•715 VOL 90-0 F 00184 •00183 •00183 •00183 •00201
100 C 233567899	71.0 54 X/D 12.11 12.11 12.11 12.12 1.	RUNF 10 RUNF 25 TO32 TYP5 1999457 199957 199957 199957 199957 1997577 1997577 1997577 1997577 1997577 1997	3H DATE 471.855 1.23281 1.23281 1.56465 1.665224 1.665224	1 2 0 1/2 9 3 2 -		PHESS. (PSIA) 10.6 97.1 AS HE-AR, MOL ATE = 36.2 LB (2) = .173, HL/QGAS .109 .146 .035 .029 .021 .031 .032	1.20 1.46 2.01LAB + 1 HAC4 (16) 9ULSE F 90.799 1609.754 1214.9465 104.577 994.55	57.0 541.2 113.8 29.7 113.8 29.7 29.7 113.8 29.7 113.8 29.7 113.8 29.7 9 779 79 79 79	D = 700E-02 - 445E+01 - 245E+01 - 245E+00	-715 VOL 90-0 F 001993 -001993 -00203 -002203 -002203 -002203 -002203 -002203
N	7 = / D 1211114548890	ROUR #57 RUUR #57 RUUR #599422 199542 19955555555555555555555555555555	347 1.8 4.2 T 1.2 1.23220405234914 1.5665524914 1.5665524914 1.5665524914 1.556317 1.55677 1.55677 1.556777 1.556777 1.5567777 1.5577777 1.5577777777777777777777777777777777777	1 2 0 1/5 9 3 E -	23/76, G 90.4 90.4 23/76, G FLO4 B 31002 31002 31002 31002 30522 773587 77393 773587 77493 64701 64260 62360 20176 90376 90376 90522	PHESS. (PSIA) 110.6 97.1 AS HE-AR, MOL ATE = 36.2 LB (2) = .173, HL/QGAS .109 .046 .029 .032 .037 .037 .040	1.20 1.46 2.01.A B + T 4.40 2.01.A B + T 4.40 4.40 4.40 4.40 4.40 4.40 4.40 4.4	57.0 541.2 113.8 29.7 113.8 3.4 57.1 541.2 547.2 54	D = 70E + 0 1 D = 70E + 0 1 445 E + 0 1 245 E	715 VOL 90-0 F 0001990-0 F 0001990-0000000 00000000000000000000000
HR C 234567890123456	7= X 121111548899353	TOTAL STREET STR	347 = 0 A T E .5 0 A T E .5 1 . 23220 40552 4914 1 . 235533272 4 155533272 4 155572 4 391 155772 4 155772 4 157772	1 2 0 1/2 9 3 2 -	23 76, G 90.4 23 76, G 23 76, G 23 76, G 23 70, MAGH 25 10022. 30022. 79493. 647367. 70293. 647367. 70293. 64736. 64736. 5982404. 54743. 54137.	PHESS. (PSIA) 110.6 97.1 AS HE-AR, MOL ATE = 36.2 LB (2) = .173, HL/QGAS .109 .046 .035 .029 .037 .037 .037 .035 .040 .045 .055 .082	1.20 1.46 XAR AL LAR AL LAR AL LAR AL LSE. (16) SUS2690.79985907 10274.9454907 10274.9454907 10274.9454907 10274.9454907 102752 1	57.0 541.2 29.7 113.8 3.9 - 251. 8 71 696 779 799 799 799 799 799 799	D = 470E + 01 D = 47	715 VF 90-75 VF 001190144444 00000000000000000000000000
The C 2345678901234567	7 # / D 12111145488993532	ROUR #55 RUUF R RUUF R 1999422+067 1999422+067 1999422+06 1999422+06 1999422+06 1999422+06 1999422+06 199942 1999422+06 19994 19904 19904 19904	347 1.8 4.2 T W / 1.8 1.23220465524814 1.23553274 1.665524814 1.665524814 1.65533274 1.55774 1.43375 1.43375 1.43375	1 2 0 1/2 9 3 2 	23 76, G 90.4 23 76, G 23 76, G 23 76, G 23 8 24 76, G 24 76, G 25 76, G 26 76, G 27 77 76, G 27 76, G 27 76, G 27 76, G 27 76, G 27 77 76, G 27 77 76, G 27 77 76, G 27 77 77 77 77 77 77 77 77 77 77 77 77 7	PHESS. (PSIA) 110.6 97.1 AS HE-AR, MOL ATE = 36.2 LB (2) = .173, HL/QGAS .109 .046 .035 .029 .032 .037 .037 .037 .037 .037 .035 .035 .045 .045 .052 .082 .382	1.20 1.46 2.11,46 2.15,49	57.0 541.2 29.7 113.8 AMP 29.7 113.8 AMP 29.7 113.8 AMP 71 67 77 79 79 79 79 79 79 79 79 7	D = 702 + 01 D	715 VF 90-0749444444444444444444444444444444444

RUN 104H, DATE 01/28/76, GAS HE-AR, MOLECULAR WT. = 29.70 TIN = 72.3 F, TOUT = 305.7 F, MASS FLOW RATE = 36.2 LB/HR, I = 37.4 AMPS, E = PR,IN = .486, GB/RESQ = .3482-03, MACH(2) = .173, MACH(16) = .223, T,SURR = 5.020 VOLTS  $\begin{array}{c} & & & & \\ & & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & & \\ & & & \\ & & & & \\$ BULK NUSSELT 295.35 167.11 151.83 135.08 TC X/D T¥/T3 HL/QGAS 2+ . 001056 .001056 .001159 .001139 1.23047707481833398748183339874818333764333376483339748183337332074 11426m2m3367892605 3 1211145487381310 124864299643103 124864299643103 ม์ 5 001130 001136 001131 001131 001131 001131 001132 001132 001132 001132 001132 001132 001132 001132 001132 678901234567 11111157 STATIC PRESS. (PSIA) 110.1 101.1 PRESS DEFECT -.471E-02 .162E+01 TB (F) 288.4 X/D PT TW/TB 1.12 90:1 22 AVERAGE PARAMETERS BETWEEN PRESSURE TAPS AVERAGE BULK REYNOLDS AVERAGE FALL REYNOLDS AVERAGE FRICTION FACTOR 71562. 53128. 200484 RUN 105H, DATE 01/23/76, GAS HE-AR, MOLECULAR WT. = 29.70 TIN = 71.9 P, TOUT = 197.5 P, MASS FLOW HATE = 36.3 LE/HP, I = 64.9 AMPS, E = 3.535 VOLTS PR,IN = .486, GB/RESQ = .193E-03, MACH(2) = .173, MACH(16) = .204, T, SURA = 73.5 P BULK 8 1899. 316541. 316541. 31773. 791798. 76596. 75357. 753229. 72308. BULK VUSSELT 164.83 155.74 135.57 126.54 117.60 112.75 107.18 TC X/D 11531769279223971 11467802345679028... dL/QGAS TW/TB Q+ 2345678901234567 1211145477371209 107.18 105.46 103.81 104.13 101.97 102.70 108.99 182.22 73129. 72086. 71032. 70048. 69037. 68229. .000653 .000653 .000653 .000650 .000650 .000650 .000582 25421.3 25335.4 25328.5 24836.6 22686.8 STATIC PRESS. (PSIA) 110.5 103.3 197.7 PT X/D PRESS DEFECT -.471E-02 .1302+01 TW/TB 1:00 yō:5 12 

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TIN = PR,IN	72.3 F	RUN 107H TOUT = 70 , Ga/RESQ	9.6 F153	1/28/76, G ASS FLOW R E-02, MACH	AS HE-AR, MOL ATE = 13.9 LB (2) = .033,	ECULAR NT. 1/HR, I = MACÉ(16) =	= 29.70 90.0 AMPS, E .123, T,SU	= 5.370 VOLIS E2 = 79.5 F
IC 34567890 1111345 11111567	X D 1211146599004644 12164208755108	1920-0-7 C 1 65 87 4 77 6 0 2 3 5 9 9 9 9 9 9 9 9 9 9 9 9 9 9 9 9 9 9 9	T#/ 1941 245941 1.45941 1.77995396 1.775936 1.775936 1.775936 1.775936 1.775936 1.754394 1.54394 1.53344 1.533444 1.533444 1.533444 1.533444 1.533444 1.533444 1.53344444 1.53344444444444444444444444444444444444	BY105139 3105139 22105139 22105139 22105139 22105139 22105139 22105139 22105139 22105139 22105139 22105139 22105139 2102156 10992478 10992478 10992478 177003	HL/QGAS - 715 - 273 - 118 - 077 - 077 - 076 - 084 - 090 - 099 - 105 - 114 - 122 - 131 - 146 - 239 1.463	T170111 KE000111 US0641.09095714 US0654.090957144 US06741.07654448663	Sat 32,47,42,27,42,27,44,42,42,42,44,44,44,44,44,44,44,44,44,	2 2 2 2 2 2 2 2 2 2 2 2 2 2
			Þ.	T X/D	STATIC PRESS. (PSIA)	TW/TB		
				2 90.4	36.1	1.33	674.0 .29	ĴŹĒ+ĊĨ
rin = Pg,in	λVI 73.2 ₽ ≖ .48	RAGE BULK 24213 24213 5001 = 51 5, GR/RESQ	DATE 0 9.7 F. N = .106	1/28/76, G ASS 7LCW R 2-02, MACU	GE 4813 BE490 16879. AS HE-AR, NOL ATE = 14.0 LB (2) = .083,	ECJLAR WT. /Ha, I. /Ha, I.	29.70 75.1 AMPS, 5 113, 7,501	= 4.320 VOLIS R = 90.0 F
IC 234567 89011234567	X/D 1211145489993422 123345993422	6214587055793476 (F)340.001906095200176 16546001906095200176 16556019060952001500 19060952001500 190609520000 19060952000000000000000000000000000000000	TW 1213744 1.14455449495 1.1444375273 1.1444375273 1.1444375273 1.11111111111111111111111111111111111	BUNCOTA 3 SUNCOTA 3	HL/QGAS 180 265 1026 0056 0056 0056 0066 0076 0082 0082 0082 0082 0099 1051 1951 1759	T9124733484073972 US68445469532.9751969541 US68445445454.009751969541 US6953444449903	QGAS BTU/25644 319/25644 31169564 33337457 33337457 33337457 3332476 3332476 3332476 33307754 33307754 33307754 33207754 33207754 33207754 33207754 33207754 33207754 322536	2 0 0 0 0 0 0 0 0 0 0 0 0 0
			5	T X/D	STATIC PEESS. (PSIA) 32.9	ти/тв 1+14	TB PR (P) DE 73.55	BSS FBC7 90 E - 02 30 F + 01
	۸¥	ERAGE BULK 25502	AEYNOLDS	VERAGE PAR	ANETERS BETWE Ge Wall Betwo 15693.	EN PRESSUR	E TAPS ERAGE PRICTI .3062	UN PACTOR

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BULK BUNCLDS 31140. 30995. 30872. 30573. 30573. 30945. TH/TB rc X/D τw HL/QGAS 2+ 238353838C339620 238353838C339620 290016307773201044 112228773345792468003 112228773344445003 . 00 1222 . 00 1109 . 00 1349 . 00 1401 . 00 1411 . 00 1413 . 00 1413 . 00 1413 7:6037 671 48441059 23 1211145487381310 5 57 001413 001413 001414 001407 001407 0001407 0001407 0001300 0001300 0001312 0001312 28945. 27983. 27121. 26276. 25534. 24335. 24198. 000 1011213145167 48.71 47.04 46.37 46.55 48.55 48.53 21159.7 21051.3 20997.9 19730.9 15665.4 23573 23005 22452 22064 STATIC PRESS. (PSIA) 09.1 97.1 PRESS DE7ECT -.5912-02 .1912+01 PT X/D TW/TB TB (F) 73.9 341.5 1.10 90.1 12 AVERAGE PARAMETERS BETWEEN PRESSURE TAPS AVERAGE BULK REYNOLDS AVERAGE WALL REYNOLDS AVERAGE PRICTION FACTOR 26805. 21228. .00612 RUN 111H, DATE 01/29/76, GAS HE-AR, MOLECULAR WT. = 29.70 TIN = 67.0 P, TOUT = 247.1 P, MASS FLOW RATE = 14.2 LB/HR, I = 48.0 AMPS, E = 2.920 VOLTS PR,IN = .486, JR/RESQ = .446E-03, MACH(2) = .084, MACH(16) = .097, T,SURR = 72.0 P BULK REYNOLDS J2094. J1993. J1911. J1345. J0501. 292.56. 03AS 1257797. 1257797. 1257797. 134774. 135777. 135577. 13577. 13 HL/QGAS X/D TZ TW/IB TC <u>ୁ</u>+ 7007922431034170 [19091751670794826 [11111222222233333332 - 00833 - 100734 - 000852 - 000852 - 000852 - 000852 - 000852 - 000852 - 000852 - 000852 - 000852 - 000852 - 000833 - 100833 - 100833 - 100833 - 100734 - 000833 - 100734 - 000833 - 100734 - 000833 - 100734 - 000833 - 000833 - 000833 - 000833 - 000833 - 000833 - 000832 - 000852 - 000832 - 00082 - 00082 - 00082 - 00082 - 000822 - 000822 - 000822 - 00082 1.067 1.130 1.170 1.2233 1.2233 1.2213 1.2273 1.2273 2345.07 0901234567 1211145477771209 124864208643107 124864208643107 09 -2009 -0009 -0009 84. 29084. 27536. 27022. 26513. 26039. 25572. 25217. - 000695 - 200695 - 200890 - 200890 - 200890 - 200890 - 200890 - 200890 - 200890 - 200890 - 200895 STATIC PEESS. (PSIA) 89.4 87.7 TB (F) 234.1 PRESS DE7ECT -.587E-02 .158E+01 PT X/D TW/TB

RUN 109H, DATE 01/28/76, GAS HE-AR, MOLECULAR WT. = 29.70 FIN = 74.1 7, FOUT = 361.3 7, MASS FLOW RATE = 13.9 L3/H3, I = 60.0 AMPS, E = 3.300 VOLTS PR,IN = .486, GR/RESQ = .681E-03, MACH(2) = .083, MACH(16) = .103, T,SURR = 80.7 7

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C 2345677901234567	X/D 12111 1230 16420 16420 16420 16420 197 197 197 197 197	5465975157950951 11111778901234450	T # / T B 1.0591 1.114 1.144 1.1555 1.157 1.144 1.1557 1.144 1.157 1.144 1.157 1.144 1.157 1.147 1.127 1.127	BULK DS BUJLK DS BY 20038. J13459. J1520. J1520. J0571. 3097096. 2285200. 289520. 289520. 289520. 27432. 277432. 27749.	HL/QGAS 172 050 036 038 039 044 0449 0449 0449 0449 0449 0444 0449 0444 0449 0444 0449 0444 0449 0564 116 234	H07546883862005599 L30537660555544444 L315357665555544444 L31535766555557494897 20153766555555749459 20152668833852005599 20152668833852005599	2 T 2 S F 6 S F 7 S	2+ 00544546 -35005577 -3005577 -3005577 -3005577 -30055774 -30055774 -30055774 -30055772 -30055773 -30055774 -30055775 -30055775 -30055775 -30055775 -30055775 -30055775 -300557
			1	x/D 90:0	STATIC PBESS, (PSIA) 89•3 87•9	TW/TB 1.04 1.12	TB 28255 (E) DEFECT (7.15478-) 174.4 .1402+(	02
	¥ A	ERAGE BUL 297	AN REYNOLDS 20.	ZEBAGE PAL Avera	AMETERS BETWI GZ WALL REYNC 25898.	EZN PRESSU DLDS A	RE TAPS VERAGE FRICTION F: .00590	RCTOR
TIN = PR.IN	70.1 P = .465	BUN 116 TOUT = 1 GR/RESQ	H, DATE 02 54.2 P, MA = .908e	/21/76, GJ 55 Plov RJ -04, Mach	AS HE-AR, MOL VTE = 44.1 LB (2) = .247,	ECULAR WT. /HR, I = MACH(16) =	= 27.53 65.2 AMPS, E = 3 .307, T,SURR =	.76) VOLIS 73.0 P
TC 2345678901234567	x 1211145377870299 12111453778870299	0465506652795752 111111112222222222	TW/T3 1.069 1.151 1.159 1.159 1.190 1.190 1.192 1.188 1.188 1.188 1.188 1.188 1.188 1.180 1.	BY04316229 BY04316229 BY04316229 BY04316229 BY04316229 BY04316229 BY04316229 BY04316229 BY04316229 BY04316229 BY043162 BY04316 BY0400 BY0400 BY0400 BY0400 BY0400 BY04	HL/QGAS .057 .130 .027 .014 .014 .014 .014 .015 .016 .017 .018 .021 .025 .096	T3022606 KE23022606 L22022606 L2202260 L2202260 L220226 L220229 L22029 L22029 L22029 L22029 L22029 L22029 L22029 L22029 L22029 L22029 L22029 L22020 L22000 L22020 L22020 L22020 L22020 L22020 L22020 L22020 L22020 L22020 L22020 L22020 L22020 L22020 L22020 L22020 L22020 L22020 L22020 L2200 L2000 L200 L2000 L2000 L2000 L2000 L2000 L2000 L2000 L2000 L2000 L200	24 5 F 4 5 S F 4 5 S F 4 6 5 7 5 5 S F 4 5 S F 5 5 S F	2 • 4454744 00000555055555 • 00000555055555 • 00000055555555 • 00000055555555 • 00000005555555 • 00000005555555 • 0000000000 • 000000000 • 00000000 • 00000000 • 0000000 • 0000000 • 000000 • 000000 • 000000 • 000000 • 00000 • 0000 • 00000 • 000000 • 00000 • 000000 • 000000 • 000000 • 000000 • 000000 • 0000000000
			PT 1 2	X/D 90.0	STATIC PRESS. (PSIA) 97.1 84.2	T4/TB 1.36 1.17	TB PRESS (P) DEPECT 60.0450E-0 147.3 .131E+0	2
	A Y B	RAJE BULK 9440	AV REYNOLDS 2.	ERAGZ PARA Averag	METERS BETWE BE HALL REYNO 78875.	EN PRESSUR LDS AV	E TAPS ZRAGE PRICTION PA .JJ464	SIDE

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RUN 112H, DATE 01/29/76, GAS HE-AR, MOLECULAR WT. = 29.70 TIN = 67.9 F, TOUT = 182.9 F, MASS FLOW RATE = 14.2 LB/HR, I = 38.4 AMPS, E = 2.250 VOLTS PR,IN = .486, GR/BESQ = .284E-03, MACH(2) = .084, MACH(16) = .093, T,SURB = 72.3 F

2GAS 3TU/HBET2 58166.6 559203.2 61333.4 1607.4 BULK NUSSELT 451.90 206.43 179.63 160.09 7977141060108486 1912522169357923579915 122222357923579915 16743339395156407 1262633340085156407 111111111111556407 061470 00114889012224755 1249.45487881310 1249.487881310 1249.4878813108 60203.2 61333.4 61607.4 61761.4 61965.7 61922.7 61944.0 61992.2 61991.4 160 143 14 128 122 118 26 80 39 116 112 111 109 165 75 32 687 501 514 62031.4 61931.5 61885.8 60861.7 53799.3 X/D STATIC PRESS, (PSIA) 97.8 80.5 TB (F) 272.1 PRESS DEFECT -.450E-02 .177E+01 PT TV/TB 1:38 90:1 12 AVERAGE PARAMETERS BETWEEN PRESSURE TAPS AVERAGE BOLK RETNOLDS AVERAGE FALL REINOLDS AVERAGE FRICTION PACTOR 88371. 65677. .00453

TW/TB

RUN 117H, DATE 02/21/76, GAS HE-AR, MOLECULAR #T. = 27.53 TIN = 70.1 P, TOUT = 287.8 F, MASS PLOW RATE = 44.2 LE/HR, I = 100.7 AMPS, E = 5.750 VOLIS PE,IN = .465, GB/RESQ = .195E-03, MACH(2) = .246, MACH(16) = .352, T, SURR = 74.0 F

HL/QGAS

X/D

TC

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TIN = PR,IN	70.1 P = .465,	RUN 118 Tout = 4 , gr/bes(	H DATE 0 07.9 2, 4 2 = .354	2/21/76, G ASS FLOW R 2-03, MACH	AS HE-AR, NOL ATE = 44.9 LB (2) = .223,	ECULAR WT. /HR, I = MACH (16) =	= 27.53 123.4 ANPS, 2 = = .339, T,SURR =	7.220 VOLTS 78.5 F
r 23:75678901234567	X 1211145498993422 12.11145498993422	x y y y y y y z y y z y z y z y z y z y	B 106445901405519464997	BULK REYNOLDS 1006472. 1000431. 97016. 89109. 8587053. 775584. 73127. 71175. 67918.	HL/QGAS 084 111 036 025 0221 0221 0225 0225 0226 0229 0324 0324 057 283	BUSS204.21 15212.54866 20741.54866 12771.54866 12771.54866 9976.987.0286 9976.987 9977.208 977.208 977.2	2 3 A S 5 1 5 1 5 1 5 1 5 1 5 1 5 1 5 1	2+ .)01662 .)01758 .001758 .001796 .001807 .001807 .001807 .001813 .001813 .001812 .00182 .00184 .00082 .000184 .00082 .00
			2	T I/D 15 2 90.2	STATIC PRESS.(PSIA) 109.6 90.6	TW/TB 1.18 1.40	TB PRESS (F) DEFEC 63.0448E 384.6 .209E	02
	AVE.	RAGE BULA	REYNOLDS	VERAGE PAR Avera	AMETERS BETHE G2 HALL REYNO 57344.	EN PRESSUS LDS AV	TAPS VERAGE PRICTION 1 .00456	PACTOR

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	T 2345678901234557	X/D 121111 24454889935532 123445643108	F (950	T # 75099551 445405551 15666655631459 156666659631459 1542424 159451 15429 15549 1	BUN06019 9100019 9959290332 9959290332 9973833297 77205017 77205017 665258 655258	HL/QGAS DS 088 0116 022 0224 0224 0224 0224 0225 0224 0227 0227 0228 0227 0228 0231 0331 0332 035 039 061 329	r79965 LS9452 US97975559 N3175559 N3175559 N3175559 N3175559 N3175559 N3175559 N3175559 N3175559 N3175559 N3175559 N317559 N3175559 N3175559 N3175559 N3175559 N3175559 N3175559 N3175559 N3175559 N3175559 N317559 N317559 N3175559 N3175559 N3175559 N3175559 N3175559 N3175559 N3175559 N3175559 N3175559 N317559 N31755559 N3175559 N31755559 N31755559 N31755559 N31755559 N31755559 N31755559 N31755559 N317555555555555555555555555555555555555	0064244555555555555555555555555555555555	24H8612.52 5397374.52 162054.630 162054.630 162054.630 162054.09 1620555555555555555555555555555555555555	Q+ 0018827 0018985 001200210 000220037 0002200378 000220041 0002200437 0002200437 0002200437 0002200437 0002200437 0002200437 00021937 000219577
· · ·				1	PT X/D 2 90:	STATIC PRESS.(PSIA) 5 109.1 3 88.5	TW/TB 1.19 1.43	TB (F) 64-8 427-0	PRESS DEFECT 449E-32 .227E+01	
		AVES	AGE BULK 83880	REYNOLD: •	AVERAGE P S AVE	ARAMETERS BETWE RAGE WALL REYNC 54606.	EN PRESSU DLDS A	RE TAPS Verage pe •	ICTION PAC 39460	TOB
	TIN = PR,IX	71.9 P. = .465,	RUN 120H Tout = 20 , gb/resq	DATE ( 8.2 P, 8.12	02/21/76, MASS PLOW 12-03, MA	GAS KE-AR, MOI RATE = 43.7 LE CH(2) = .249,	LECULAR HT 3/HR, I = MACH(16)	. 3 27.5 ∂0.8 AMP = 330,	] S, E = 4. T, SURR *	750 VJLIS 30.7 P
	E 23456789011234567	X/D 12111 24111 8642 1230 1200	W) 1.5.5.1.97 1.1.5.5.4.7.2.2.2.9.2. 1.1.1.2.2.2.2.2.2.1.3.3.5.1 1.1.1.2.2.2.2.2.2.3.3.3.5.1	TW/TB 11935089521 1.193569521 1.22569521 1.22569521 1.2226951 1.2226951 1.2226952 1.22291 1.2226952 1.222695 1.22265 1.22265 1.22265 1.22265 1.22265 1.22265 1.22265 1.22265 1.22265 1.22265 1.22265 1.22265 1.22265 1.22265 1.22265 1.22265 1.22265 1.22265 1.22265 1.2265 1.2256 1.2256 1.	BY998376682 9998376682 999837450682 99753703982487 99753703982487 99753703982487 992088447 992088447 992088447 992088447 99208847 992088447 992088447 99208847 992088447 99208847 99208847 99208847 99208847 99208847 99208847 99208847 99208847 99208847 99208847 9920847 99208947 9920847 9970847 9970847 9970847 9970847 9970847 9970847 9970847 9970847 9970847 9970847 9970847 99700847 9970847 9970847 9970847 9970847 99708700	HL/QGAS DS 096 031 014 015 015 015 015 017 017 017 017 019 019 024 038 146	LKE 11 USSE 304000316955 N298678.0400031697 114807.04316397 111905.1397 11198055665 11108055665 11108055665 11118055665 11118055665	0006839999999999999999999999999999999999	2 5 5 5 5 5 5 5 5 5 5 5 5 5 5 5 5 5 5 5	2+ .000775 .000775 .000775 .0007781 .0007781 .000783 .000783 .000783 .000783 .000783 .000783 .0007785 .0007785 .00007785 .000007785 .00007785 .00007785 .00007785 .00007785 .00007785 .000000000000000000000000000000000000
				I	PT X/D 2 90:	STATIC PRESS.(PSIA) 5 95,5 0 90.6	T¥/TB 1.11 1.24	TB (2) 61.7 197.8	FRESS DEFECT 451E-02 .1512+01	
		¥ 4 61	RAGE BULK 90788	REYNOLD	AVERAGE P S AVE	ARAMETERS BETW RAGE HALL BETW 69459.	EEN PRESSU DLDS A	RT TAPS Verage 78	ICTION PAC	ECE
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RUN 119H, DATE 02/21/76, GAS HE-AR, MOLECULAR #T. = 27.53 TIN = 71.9 F, TOUT = 452.4 F, MASS PLOW RATE = 44.7 LD/HR, I = 130.3 AMPS, E = 7.720 VOLTS PR,IM = .465, GB/RESQ = .392E-03, MACH(2) = .224, MACH(16) = .355, T,SURR = 31.5 F

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RUN 124H, DATE 03/01/76, GAS HE-AR, MOLECULAR WT. = 15.30 TIN = 70.5 P, TOUT = 183.2 P, MASS FLOW RATE = 14.2 LB/H3, I = 52.7 AMPS, E = 3.070 VOLTS PR,IN = .419, GZ/RESQ = .682E-04, MACH(2) = .113, MACH(16) = .127, T, SURR = 74.0 P BULK REYNOLDS 32217. 32152. 32102. LS27753772302 US279015544599628484 N177765544459962884 N1777655444553 N17776554444553 TC X/D T¥ TW/TB HL/QGAS 2+ 7640932266992890 1124579999989719 1124567999989719 23 .000500 .000550 .000550 .000550 121145377970299 32102. 31984. 31758. 31301. 30867. 30465. 30051. 29668. 567 3 64 29294 29294 28939 28576 28236 27888 100562 100562 200562 200563 604 148 14866. STATIC PRESS. (PSIA) 92.7 90.0 TB (F) 68.8 174.5 PBESS DEPECT -.585E-02 .1402+01 PT X/D TW/TB 90:0 1.03 1 2 AVERAGE PARAMETERS BETWEEN PRESSURE TAPS AVERAGE BULK REYNOLDS AVERAGE WALL REYNOLDS AVERAGE PRICTION PACTOR 30057. 26548. .30534 RUN 125H, DATE 03/01/76, GAS HE-AR, MOLECULAR WT. = 15.30 TIN = 71.0 P, TOUT = 345.8 P, MASS PLOW RATE = 14.2 L3/HR, I = 31.5 AMPS, E = 4.305 VOLTS PR,IN = .419, G2/R25Q = .163E-03, MACH(2) = .113, MACH(16) = .142, T,SUBR = 75.5 P BULK BEYNOLDS 32051. 31902. 31504. 3097 23 ASS P2 37 434 • 2 357 77 79 • 1 99 9 • 1 400 5 4 • 9 400 15 4 • 9 400 15 4 • 9 400 15 7 • 1 400 15 7 • 1 400 16 5 • 9 BUSSE.094 NUSSE.094 700.133 531.525 45.74 45.74 45.74 4 7917375344427310 112222335325380524871021 112222335380524871021 HL/QGAS TC X/D TW/TB Q+ .))1267 .001209 .))1323 .))1343 .)01356 .)01356 .)01366 .)01366 .)01366 23 .083 1211145487881310 124954208643108 124954208643108 567 29961. 29044. 99011234567 111234567 45.79 44.45.29 42.22 40.92 40.25 39.84 40.95 39.85 40.91 81.51 001361 001359 001357 001357 001353 001315 001315 26689. 26007. 25381. 24765. 24203. 23635. 23210. 40217.4 40165.1 40176.9 40131.7 39978.7 38357.9 32487.7 TB (P) 69.9 325.8 STATIC PRESS. (PSIA) 92.7 89.0 PRESS DEFECT -.5872-02 .1872+01 27 X/D TW/T8 90:1 1.38 12 

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			PT	X/D	STATIC BRESS (DSTA)	TW/TB	TB PRESS	
			1	90.0	76.7 66.8	1.36	62.85142- 159.8 .142E	02 01
	A V 3	29 AGE BULK 5248	aernolds'	AVERA	GE 7411 AEYNO 44157.	LOS AV	ERAGE PRICTION P	ACTOR
tin = Pr,iy	73.2 P. = .419	29 AGE 801 K 5248 5248 5248 5248 5248 5248 5248 5248	H DATE 03	/01/76, G/ SS PLOW 8 -03, MACH	GE WALL BEYNO 44157. AS HE-AR, MOL ATE = 14, 1 LB, (2) = .112, 1	LJS AV BCJLAR 97. /HR, I = 10 /HR, I = 10 /HR, I = 10	= 15.30 02.2 AMPS, E = .156, T,SUBR =	ASTOR 5.930 ₹31.2: 78.5 P
rin = Pa, in TC	AVS 73.2 P, 2.419 X/D	29 AGE 801 K 5248 20 N 127 Tout = 5 , GB/RESQ T7 (7)	H DATE 03 10.4 P 256E TW/TB	AVERA AVERA SS PLOW R -03, MACH BULK REYNOLDS	GE WALL EENO 44157. ATE = 14.1 LB (2) = .112, 1 HL/QG AS	LJS AV BCJLAR WT. /HR, I = 10 JACH (16) = BOLK NUSSELT	= 15.30 02.2 AMPS, E = .156, T,SURR = 0GAS BTU/HRFT2	AJTOR 5.990 Volt: 78.5 P 2+
FIN = FR, IN TC 34	73.2 P. = .419 X/D 1.2 2.1	23 AGE 801K 5248 5248 120T = 5 , GB/RESQ T3 (2) 152.7 290-2	H. DATE 03 10.4 P 34 256E TW/TB 1.149 1.313 1.385	/01/76, G/ SS PLOW R/ -03, MACH/ BULK REYNOLDS 31735- 31316-	GE WALL HEYNO 44157. AS HE-AR, MOL ATE = 14.1 LB. (2) = .112, 1 HL/QGAS .074 .146 .046	LJS AV ECJLAR 91. /HR, I = 10 AACH (16) = BULK NUSSELT 184.51 H0.63 71.05	ERAGE PRICTION R .00504 .005	5.990 YOLE 78.5 P 2+ .002016 .002016 .0020191 .0020191
FIR C 2345578	XVS 73.2 P. = .419 X/D 1.2 2.1 3.1 3.1 16.4 24.5	23 AGE 801K 5248 5248 20 N 127 TOUT = 5 7, GB/RESQ 152.7 290.2 314.9 343.9 449.5	H DATE 03 10.4 Pi MA = .256E TW/TB 1.149 1.313 1.385 1.484 1.492 1.492 1.476	/01/76, G/ SS PLOW 82 -03, MACH BULK 82YNOLDS 31735- 31316- 30094- 28608- 27325,	GE WALL HEYNO 44157. AS HE-AR, MOL ATE = 14.1 LB (2) = .112, HL/QGAS .074 .146 .046 .046 .027 .027 .027	LJS AV ECULAR WT. /HR, I = 1( HACH (16) = BULK NUSSELT 194.51 A0.63 71.05 61.06 53.20 46.97 43.95	ERAGE PRICTION R .00504 02.2 ANPS, E = .156, T,SUBR = OGAS BTU/HRPT2 59526.6 51737.6 62919.0 63524.8 63524.8	5.990 VOLT 78.5 P 2+ .002016 .002016 .002011 .0020131 .002131 .002114 .002155 .032155
FIR C 234557890111	AVS 73.2 P. 73.2 419 X/D 1 1.2 1.4 1.5 32-4 3.1 2.1 1.4 5.4 5.4 5.4 5.4 5.4 5.4 5.4 5.4 5.4 5	EDN 127 EDN 127 EDN 127 COR = 5248 EDN 127 COR = 50 1528 1528 2994 1528 2994 1528 2994 1528 2994 1528 2994 1528 2994 1528 2994 1528 2994 1528 2994 1528 2994 1528 2994 1528 2994 1528 2994 1528 2994 1528 2994 1528 2994 1528 2994 1528 1538 1528 1538 15	H DATE 03 10.4 P. 03 10.4 P. 03 10.4 P. 03 10.4 P. 03 1.49 1.385 1.443 1.385 1.4423 1.492 1.476 1.451 1.492 1.451 1.399	/01/76, G/ SS PLOW 8/ -03, MACH BULK BULK BULK BULK BULK 31502. 31316. 30094. 2865. 30094. 27325. 26151. 24240.	GE WALL HEYNO 44157. AS HE-AR, MOL ATE = 14.1 LB (2) = 112, 1 HL/QGAS .074 .146 .046 .046 .027 .027 .027 .027 .027 .0235	LJS AV ECULAR RT. /HR, I = 1( HACH (16) = BULK NUSSELT 184.51 H0.63 71.05 53.20 46.97 43.95 42.31 41.16 40.21	2RAGE PRICTION R .00504 .005	5.990 VOLT 78.5 P 2+ .002016 .002016 .002131 .002131 .002132 .002153 .002155 .002156 .002156 .002156 .002156
TR C 2345578901112345 1112345 1112345	AVS 73.2419 X/D 1211142 1012114 10121114 10121114 10121114 10121114 10121114 10121114 101211114 101211114 1012111114 10121111111111	23 AGE 8011 80 M 127 5248 20 M 127 50 M 125 127 50 M 125 127 127 127 127 127 127 127 127	H DAT 03 10.4 P 34 10.4 P 34 10.4 P 34 10.4 P 34 1.315 1.345 1.485 1.4976 1.4976 1.4976 1.4976 1.4976 1.4976 1.4976 1.3991 1.3971 1.3971 1.3421	/01/76, G/ SS PLOW 8/ -03, NACH BULK REYNOLDS 317352. 31316. 300944. 287325. 262151. 224240. 23383. 225151. 224240. 23383. 22525. 224240. 22328.	GE WALL HEYNO 44157. AS HE-AR, NOL. ATE = 14.1 LB (2) = .112, HL/QGAS .074 .146 .046 .027 .027 .027 .029 .031 .043 .041 .043 .043 .043	LJS AV ECULAR #T. /HR, I = 1( HACH (16) = BULK NUSSELT 184.63 71.06 53.20 43.95 45.22 35.95 45.22 35.	PRASE PRICTION R .00504 .00504 .00504 .00504 .02.2 ANPS, E = .156, T.SUBR = .156, T.SUBR = .156, T.SUBR = .156, C.SUBR = .156, C.SU	5.990 V0LE 78.5 2+ .002016 .002091 .002091 .002131 .002153 .002153 .002153 .002153 .002153 .002153 .002153 .002154 .002154 .002154
TR C 23455789011234567 11123455789011234567	AVS 73.2419 X D 12111454889924 121114564310.8 1211145645710.8 1211145645710.8 1211145645710.8 1211145645710.8 1211145645710.8 1211145645710.8 1211145645710.8 1211145645710.8 1211145645710.8 1211145645710.8 1211145645710.8 1211145645710.8 1211145645710.8 12111145710.8 1211145710.8 121111145710.8 12111111111111111111111111	23 AGE 80112 8011 127 5248 2011 127 5001 127 5001 127 127 127 127 1240 1240 1240 1240 1240 1240 1240 1240 1240 1240 1240 127 1240 127 1256 127 1256 127 1248 127 127 127 127 127 127 127 127	H, DATE 03 10.4 P 34 10.4 P 34	/01/76, G/ SS PLOW R/ BULK DS BULK DS 31502. 31502. 31502. 300094. 2252. 20094. 2252. 2252. 224240. 2252. 2255. 25	GE WALL HEYNO 44157. AS HE-AR, NOL. ATE = 14.1 LB. (2) = .112, HL/QG AS .074 .046 .046 .046 .046 .049 .027 .029 .031 .049 .031 .041 .043 .043 .043 .043 .043 .045	LJS AV ECULAR #T. /HR, I = 10 HR, I = 10 BULK NUSSELT NUSSELT NUSSELT NUSSELT 184.51 A0.63 53.20 43.95 42.31 41.16 40.21 37.22 36.97 43.95 43.95 47.31 37.92 73.31	ERAGE PRICTION R .00504 .00564 .005	5.990 VOLE 78.5 78.5 2+ .002016 .002110 .002111 .002131 .002131 .002153 .002153 .002153 .002153 .002153 .002153 .002153 .002153 .002153 .002153
TPR C 234557890111234567	AVS 73. 2 4 19 X /D 121114 1641-548 1642-548 164	23 AGE 8012 80 M 127 5248 20 M 127 50 M 127 50 M 127 50 M 127 1528 15260 15260 15260 15260 15260 15260 1240 15260 1240 15260 1240 15260 1250 15260 12750 12750 15260 12750 15260 12750 15260 12750 12460 12750 12460 12750 12400 12750 12400 12750 12400 12750 12400 12500 12400 12500 12400 12500 12400 12500 12400 12400 12500 12400 12500 12400 12400 12500 124000 12400 12400 124000 12400 124000 12400	ABYNOLDS ABYNOLDS H DATE 03 10.4 P: 03 10.3 P: 03	/01/76, G/ SS PLOW 8/ -03, NACH BULK REYNOLDS 315302. 31316. 300994. 28635. 26151. 224240. 223383. 21905. 21905. 21905. 21905. 21905. 21905. 21999. 20199. 20199.	GE WALL HEYNO 44157. AS HE-AR, NOL. ATE = 14.1 LB (2) = .112, HL/QGAS .074 .146 .046 .027 .027 .027 .027 .027 .029 .031 .045 .041 .043 .041 .045 .041 .045 .041 .045 .041 .045 .041 .045 .045 .041 .045 .045 .045 .045 .045 .045 .045 .045	LJS AV ECULAR #T. /HR, I = 1( HACH (16) = BULK NUSSELT H0.63 71.06 53.20 43.95 45.20 45.2	PRASE PRICTION R .00504 .00504 .00504 .00504 .00504 .02.2 ANPS, E = = .156, T.SUBR =	5.990 VOLT: 78.5 P 2+ .002016 .002016 .002150 .002153 .002153 .002153 .002154 .002153 .002154 .002154 .002153 .002153 .002153 .002153 .001538

BUN 128H, DATE 03/01/76, GAS HE-AZ, HOLECULAR WT. = 15.30 TIN = 74.1 P, TOUT = 688.6 P, MASS FLOH RATE = 14.1 LB/HR, I = 121.2 AMPS, E = 7.280 VOLTS PR,IN = .419, GZ/RESQ = .355E-03, MACH(2) = .113, MACH(16) = .171, T,SURR = 31.5 F BULK YNOLDS 1682. 11357. 11095. 6315 BTU/HEFT2 83373.1 IC I/D HL/QGAS BULK r W TW/TB 2+ 30LK NOSSELI 165.34 77.77 67.10 56.17 [#) 198.0 198.0 196.4 195.0 195.0 195.0 195.0 195.0 RĘ 1.229 1.453 1.564 1.569 1.723 1.735 1.735 002816 23456799011234567 124364208004544 12437445673108 9470.5 30496 4544771.3392537 8844771.3392537 885457772773392537 8914504537 8914504537 1.500 1.660 1.611 1.525 1.405 1.403 1.403 1.403 1.349 1.39 060 061 074 117 20470. 19694. 18974. 18330. 82 14 38 57 x/D 90.4 STATIC PRESS. (PSIA) 92.7 87.0 PRESS DEFECT -.588E-02 .2902+01 9T TB (F) 74-4 650-4 TW/TB 1.18 12 AVERAGE PARAMETERS BETWEEN PRESSURE TAPS AVERAGE BULK REYNOLDS AVERAGE VALL BETNOLDS AVERAGE PRICTION PACTOR 25029. 17357. .00616 RUN 129H, DATE 03/01/76, GAS RE-AB, NOLECULAR WT. = 15.30 TIN = 75.4 P, TOUT = 306.3 P, MASS FLOW RATE = 24.4 LB/HR, I = 102.4 AMPS, Z = PR,IN = .419, GR/RESQ = .609E-04, MACH(2) = .233, MACH(16) = .335, T,SURR = i volis 5.925 BULK BUNCLDS 55262. 55020. 54830. 54398. 53582. 51992. BULK NUSSELT 204.06 113.76 101.31 88.79 79.42 71.72 TC X/D TZ TV/TB HL/QGAS 33AS 17HBF 12 9450 - 1 8758 - 9 2852 - 8 2000993867936430 402592479146766430 122222333334444454 зţ .079 .096 .026 .018 .015 .015 2345678901234567 148 121145437881310 124364238643108 124364238643108 1.2495 2495 1.2295 1.33676 3.3757 3.3566 3.3566 3.3566 3.3566 3.3566 3.3566 3. 4223.2 4189.6 4093.8 3112.5 4878.2 96 023 026 042 194 966 2019 340 631 5<sup>4</sup> : Śċ STATIC PRESS. (PSIA) 76.7 63.3 PRESS DEFECT -.515E+02 .194E+01 PT X/D 17/18 TB (F) 67.0 289.7 90:1 1.12 12

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AVERAGE PARAMETERS BETWEEN PRESSURE TAPS AVERAGE BULK REYNOLDS AVERAGE WALL REYNOLDS AVERAGE PRICTION PACTOR 48657. 364 11. .JOS 16

RUN 130R, DATE 03/01/76, GAS HE-AR, HOLECULAR 4T. = 15.30 TIN = 76.3 P, FOUT = 453.5 P. MASS FLON RATE = 24.3 LE/HR, I = 126.0 AMPS, E = 7.440 VOLFS PR,IN = -419, GE/RESQ = .143E-03, MACH(2) = .186, MACH(16) = .272, F,SURR = 04.7 P BULK BEYNQLDS 54594. 54232. GAS 1/HRFT2 675.1 X/D T7/TB HL/QGAS 3+ TW tc BULA NUSSELT 218-59 2345578901234567 124864208643108 124864208643108 12344567399 STATIC PRESS. (PSIA) 96.3 84.1 T8 (F) 71.9 426.8 PRESS D2PECT -.516E-02 .222E+01 27 X/D TW/TB 1.15 90.2 12 RUN 1314, DATE 03/01/76, GAS HE-AR, MOLECULAR WT. = 15.30 TIN = 76.8 P, TOUT = 623.6 P, MASS FLOW RATE = 24.2 LB/HR, I = 151.1 AMPS, E = 9.100 VOLTS PR,IN = .419, GR/RESQ = .206E-03, MACH(2) = .106, MACH(16) = .305, T,SURR = 37.3 F BULK YNOLDS 4345. 3822. TW/TB HL/QGAS 3+ 1/5 art f i TC 23 1.2 4567 1234456799904644 1021.4 1029.2 STATIC PEESS. (PSIA) 96.3 80.9 PRESS DEPECT -.517E-02 .28CE+01 TB (P) 73.1 588.0 9T X/D TW/TB 1.21 90.4 12 

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			5	T X/D	STATIC PRESS. (PSIA) 103.1 96.3	TW/TB 1.04 1.12	T9 PRESS (F) DEFECT 73.9 - 5962-0 166.8 - 1512+0	2
<b></b> .	<u>=</u> 77.7 <u>-</u>	8UN 134 10UT = 1	H DATE 0	3/10/76, G Ass Flc4_P	AS HE, MOL ATE = 11,6 La	ECULAE WT.	. = 4.0026 135.4 AMPS, 2 = 1 - 211. TSTAR = 1	7.390 VOLIS 78.0 P
23,I	N = .66 $K/D$ $1.2$ $2.1$ $4.1$ $3.1$ $16.4$ $24.5$ $32.4$ $40.8$	7, G8/E2S TW 142.6 211.5 233.1 2639.3 122.0 146.7 147.7 146.7	2 = .272 TW/T3 1.129 1.235 1.325 1.325 1.351 1.354 1.343 1.328 1.321	E-04, JACH BULK REINOLDS 309105. 299705. 299705. 299705. 277148. 26974. 26974.	(2) = .152, UL/QGAS .051 .063 .016 .011 .009 .009 .009 .010 .011 .011	ACA (10, 10) BUS93	QGAS BTÜ/HRFT2 106721.3 166024.3 111132.7 111915.1 112277.2 112565.3 112741.9 112966.1 113097.4 113097.4	Q+ .C01143 .C01135 .001196 .C01205 .C01205 .001208 .C01209 .C01209 .C01209 .C01211

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ZUN 135H, DATE 03/10/76, GAS HE, NOLECULAR WT. # 4.0026 TIN = 77.7 P, TOUT # 471.7 P, MASS FLOW RATE # 11.6 LB/HR, I = 171.2 AMPS, E = 10.170 VOLTS PR,IN = .667, GR/RESQ = .4352-04, MACH(2) = .163, MACH(16) = .235, T,SURR = 32.0 P BULK NUSSELT 73.556 41.555 QGAS BTU/HERT2 170678.3 "70966.9 "44.3 BULK REYNOLDS 30091. 29894. 29741. HL/QGAS 2+ 29/73 TC X/D TT .055 .061 .017 .012 .011 .011 (<sup>2</sup>) 190.3 294.1 1.215 1.398 1.471 1.523 1.557 1211145438993422 3 844. 275. C28. 30 35678901234567 iù d 143. 4555666677776 722386077776 905 965 947 551.8 537.3 493.5 019 028 183 STATIC PEESS. (PSIA) 103.1 91.6 PRESS DEFECT -.596E-02 .253E+01 TB (F) 74.7 442.3 X/D TW/TB 51 1:32 90:3 12 AVERAGE PARAMETERS BETWEEN PRESSURE TAPS AVERAGE BULK REYNOLDS AVERAGE FRICTION FACTOR 25428. 18139. .00652 RUN 136H, DATE 03/10/76, GAS HE, HOLECULAR WT. = 4.0026 TIN = 76.8 P, TOUT = 618.7 P, MASS FLOW RATE = 11.6 LB/HB, I = 199.6 AMPS, E = 12.015 VOLTS PR.IN = .667, GR/BESQ = .598E-04, MACH(2) = .162, MACH(16) = .258, T, SURR = 86.5 F BULK EYNOLDS 30055. 29786. 29578. QGAS 9TU/HBFT2 233859.2 233974.0 30 L K 0+ HL/QGAS TW T7/T8 I/D TC NUSSELT 179-54 94-72 .05C .019 .013 .013 .013 .013 RE (Z) 30. 289 2345 1 1.533 1.640 1.711 1.760 1.760 1.725 434. 493. 559. 642. 6789011234567 1.6 άÓ 560.5 496.9 770.0 45.11 45.11 48.57 57.04 021 024 035 260 382 25a STATIC SS. (PSIA) 103.2 89.2 PRESS DEFECT -.596E-02 .309E+01 TW/TB IB (P) 74.5 579.9 PT I/D PRE 90.4 1.22 12 

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BULK NUSSEET 24294949 2017 17 17 3ULX 9EYNGLDS 100225. 100225. 100057. 99947. TC x/0 TW/TS HL/AGAS ¢+ 451955566789850 .006417 .000392 .000436 234 767 89 01234 567 1241145477870299 124564235648197 124564235648197 99667. 99667. 99137. 99056. 970250. 95072. C 3C 44 19-0044 94151 62.00 92392 92392 91445 90529 39740 39740 39002 157.44 155.42 155.42 166.39 304.14 CCC444 COC444 CCO443 CCO443 OCO437 24915.0 24915.0 24904.3 24847.9 24527.5 23178.2 STATIC PRESS.(PSIA) 92.1 84.4 PRESS DEFECT -.45CE-02 .1175+01 T3 (F) 69.6 151.7 PT x/0 TW/T3 1.05 90.0 12 AVERAGE PARAMÉTERS BETWEEN PRESSURE TAPS AVERAGE BULK REYNOLDS AVERAGE WALL REYNOLDS AVERAGE FRICTION FACTOR 94992. 80621. .CO451 RUN 1534. DATE 04/J2/76, GAS AIR , MOLECULAR WT. = 29.97 TIN = 75.9 F, TOUT = 236.5 F, MASS FLOW RATE = 35.4 LB/HR, I = C9.1 AMPS, E = 5.720 VOLTS PR.IN = .72C, GR/RES3 = .386E-03, MACH(2) = .224. MACH(16) = .293, T,SURR = 79.5 F SY 97 8 8 9 6 4 2 2 2 4 7 7 8 8 9 6 4 2 2 2 4 7 7 8 8 9 6 4 2 2 2 4 7 7 8 9 9 6 4 2 2 2 4 7 7 8 9 9 7 8 9 7 6 4 2 2 2 5 7 7 8 9 7 6 4 2 2 2 5 7 7 8 9 7 6 4 2 2 2 5 7 7 8 9 7 7 7 8 9 7 7 7 8 9 7 7 8 2 TC X/0 TW/TB HL/QGAS 0+ .C009336 .C009336 .C01056 .C01056 .C01056 .C010571 271510002244607 93422202002220000 2345678901224567 90227. 932499. 94299. 922990. 922990. 92290. 92290. 92290. 92290. 9227. 92490. 9227. 9227. 9227. 9329. 9329. 9227. 9329. 9339. 939 CC1071 CC1C72 CC1072 :3 679

RUN 1494, DATE 04/02/76, GAS AIR , MOLECULAR WT. = 25.97 TIN = 74.5 F, TCUT = 153.5 F, MASS FLOW PATE = 35.6 LB/NR, I = 64.2 AMPS, F = PF-IN = .72C, GR/RESO = .1625-U3, MACH(2) = .225, FACH(16) = .250, T, SUPP =

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3.555 VOLTS

.001072 .001072 .001072 .001051 .001051 PRESS+(PSIA) 92+1 92+0 PRESCT DEFECT -.4511-02 .1555+01 T%/T3 13 (F) 71.5 273.7 PT X/D ł 90:1 1.12 AVERAGE BULK REYNOLDS AVERAGE WALL REYNOLDS AVERAGE FRICTION FACTOR 80442. 63447. 63447.

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C 2345678901234567	1314 4 545 6 67 7 7 7	9361274656744089 747732734556744089			/ 24506666555555442 7 3541565319742942		UN 92 089 425 6800 1855 9799764 19520354 209 8999764 19520354 209 8999764 19520354 209	KL6463734294811484	; H		G G442222222222223333304	n, 1 4 22 1 - 1 - 1 - 1 - 1 - 1 - 1 - 1 - 1 -		T 500 JA 004 400 040 04		10349448888888888888888888888888888888888	G197939377408164 6010	2		
						РТ 1	X	/0	PRE	STA SS. 91	TIC (PSIA	TW) 1.	/T2 19		TB (F) 73.	3	-	DEFEC	T C2 +01	
						-														
	AVÉRA	GË BI	ULK R 4755.	EY	NOL	D S D S	ERAGE	P 44 VERI	RAMET AGE W 5	ERS ALL 594	SETW Reyn 7.	EEN OLDS	PRE	SURE AVE	TAP Rage	S FR		T 10N 462	FACT	CR
N = 77.	∆V≟RA 7 F, T: .719,	RUN J	ULK Q 4755. 527, 527	. 5	NOL ATE F,		02/72		AGE W AGE S	ERS ALL 594 IR = 3	95TW REYN 7.		PR::	SURE AVE # 14:	TAP RAGE 3.2 .35	5 FR	7 5 7	E = ; ;URR ;	FACT 8.6	CR 35 VQ 7.0 F
N = 77. , IN = X/O	AVERA 7 F. T. .719, 1 2	RUN J GR/RE TJ (F) 47,4	ULK 9 4755. 527 527	EY D, •6	NOL F, .00 / TB 322	443 062-	02/76 3 FLC 03, 9921	VER VER	AGE 5	ERS ALL 594	9 # T W R 7 • 5 • 2 2 3 • 5 • 2 2 3 • 6 • 6 •		PR = 1	WT. 1 = 14 ] =	TAP Rage 3.2 .35	S FR. 3.4 9115	100 7	5 =	FACT 8.6	35 VQ 7-0 F 3+ • <u>6</u> 2 2 0
N = 77. , IN = X/O 1.	AVERA 7 F. T. 719, 12 4 1 5 1 5 1 5 1 5 1 5 1 5 1 5 1 5 1 5 1 5	GE 81 BI BI BI BI BI BI BI BI BI BI	ULK Q 4755.	EY D. .6 TW	NOL F. 0 7 2247320 7 2247320 7 2247320	A V D S 0 4 / 3 0 6 2 -	ERAGEA 02/70 03/70 01/10/10/10/10/10/10/10/10/10/10/10/10/1		AGE S AGE S ATE (2)	ERSL ALL4 594 IR 3: 1000000000000000000000000000000000000	9 4 5 Y Y Y Y Y Y Y Y Y Y Y Y Y Y Y Y Y Y		P 4 1 US030854	WT 14 T479521	TAP R AGE 3.2 3.2 .35		10 75T 6/102074	EUR 7744a2	8.6 5	35 VGF 35 VGF 00022222
x / 0 X / 0 120 140 140 140 140 140 140 140 14	AVERA 7 F. T. 7 719, 12 44 11 55 65 77 88 89 98	G III R 4835425719 RUG V(4185206159	ULK 9 4755. 527, 527 527	EY 0. TW	N L F. U 24732158700 TF. U 2473215870	04/3 962-	ER AGA 7L, UNC1:::00 2/F, UNC1:::00 2/F, UNC1::00 2/F, UNC		A S A A TE	ERS 494	7 5.2 G 4143992368		R L L L L L L L L L L L L L L L L L L L	SUPE SUPE Main T4795210570	TAP R AGE 3.235		LO 75T G/1021120221	T4 EU F7744a24252	8.6 - 6	CR 35 0 0000202020202 000000000000000000000
N = 77. IN = 12. 12. 12. 12. 12. 12. 12. 12. 12. 12.	AVERA F19, 2445677888999004664107	G RUG T(4185200159309109 C RUG T(4187394961459214590	ULK 9 4755. 527, 527 527	EY D6 T 1111111111111111111111111111111111	N L F. T 214491384077665536	04 96 96 96 96	R AGA 7L, UNX21017173 7F, UNX210170120 7F, UNX210170120 8Y8776152840742076020 97797979784077777766642	PE 514 11231729603333575143	A MET AGE 5 ATE (2)	RL9 RL9 I U 010000000000000000000000000000000000	7 5. G 944339923334445775		R         I         US63045747874           R         I         US63045747874           I         US63045747874         US630454787874	S UAV W=) T4795510570819059	TAP RAGE 3.2 .35		10 7ST G/10211202239974		8.6 • 6	CR 97 97 97 97 97 97 97 97 97 97

TIN = PR+IN	77.7 F, • .719	RUN 154 Tout = 1 Gr/res3	4. 04TE C4 71.2 F, 44 2395	/02/76, G/ SS #LJW R/ -03, MACH(	AS AIR + MGLI ATE = 29.3 LA (2) = .17G+ 1	ECULAR HT. /HR, I = MACH(16) =	# 23.97 59.5 AMPS, E # 188, T.SURR	3.260 VOLTS • 32.0 F
TC 2345 0785	×/) 1.2.1 12 1	T + 1071.07 1051.07 1051.07 1050.06 1770.01 1901.9 2013.9	TW/T3 1143 11452 11452 1182 1182 1182 1182 1182 1182	3ULU 4UL 3ULU 3ULU 3ULU 3ULU 3ULU 3ULU 3	HL/JGAS .083 .105 .022 .0117 .0117 .0117	TZ234 09.48 403407784 49 203407784 44 203407784 44 1154444 11544444	Z S C C C C C C C C C C C C C	9+ • 4444 • 20004444 • 2000444444 • 20004444444 • 20004444444 • 20004444444 • 20004444444 • 20004444444 • 200044444444 • 20004444444444 • 2000444444444 • 20004444444444444444444444444444444444
11234567 11234567	434.7 344.0 73.0 997.9 97.9	2137.38 22345.00 2553.00 2555.00 2555.00 2555.00 2555.00 2555.00 2555.00 2555.	1 • 180 1 • 177 1 • 177 1 • 174 1 • 172 1 • 168 1 • 153 1 • 157	75750. 75963. 751825. 744655. 72918. 72918. 72526.	.0149 .0140 .02212 .02224 .02429 .010	137.54 135.42 132.42 132.472 137.10 137.59 261.59	2765 2765 2765 2765 2765 2765 2765 2765	
			PT 1 2	×/D 90.0	STATIC PRESS.(PSIA) 99.9 95.2	TW/TB 1.05 1.15	TB DEFESS (F) DEFESS 75.0472F 163.0 .115F	T -C2 +01
TIN T PR.IN	4∨E 76.6 F, ■ .715	RUN 155 TOUT = 2 , GR/RESQ	45YNOLOS 5. 9. 94.5 F, Ma 94.5 F, Ma 94.5 F, Ma	VO2/75, GA S FLOW PA -03, MACH	AMETERS BETWE E WALL REYNOL 65365. AS AIR , MOL ATE = 29.1 L3 (2) = .170, 1	ECULAR WT. 199, 1 MACH(16) -	23.37 90.55 212, T.SUPR	FACTUR 5.220 VOLTS
T 2344 07 490 121111107	X 124111454878882310 124964203643108	4440m20-m7-104200 11120-m7719-m6- 12120-14-77719-m6- 1220-7771-57-000	1       1	S UN0504005173413755 UN0504099910000000000000000000000000000000	HL/9GAS .015194 .0155194 .000000000000000000000000000000000000	T976224596034 L524214649459545545446 L524214649457442144 BJ539864322211111144 BJ539864322211111144 L11111111114	2 5 5 5 5 5 5 5 5 5 5 5 5 5 5 5 5 5 5 5	2+ .CC210279 .CC210279 .CC210279 .CC21026822 .CC01026822 .CC01026822 .CC01026822 .CC01026823 .CC001000 .CC001000 .CC00100 .CC001000 .CC00000 .CC00000 .CC00000 .CC00000 .CC00000 .CC00000 .CC00000 .CC00000 .CC00000 .CC00000 .CC00000 .CC00000 .CC00000000
			PT 1 2	x/D c0.1	STATIC PRESS (PSIA) 99.9 93.7	TH/T8 1.11 1.30	TB PRESS (F) DEFEC 76.4472E 262.0 .151	T -02 +01
	<b>∆</b> V (	RAGE BULK	RETHCLDS	ERAGE PARI AVERA	AMETERS BETWE Ge Wall Peyng 54080.	EN PPESSUA LOS AV	E TAPS EFAGE FPICTION .CO475	FACTOR

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2 3 4	×/0 1.2 2.1	TW (F) 195.2 312.7 368.3	TW/TB 1+213 1+429 1-522	i	ALLK REYNOLDS 30768. 30297. 79523.	HL/JGAS .108 .161 .350	3ULK NUSSELT 413.54 200.43 130.37	9 310 70 67	GAS /HRFT2 496-4 935-5 253-7	0+ .C01529 .CC1471 .CO163
56789C1234567	1214548993422	3231160179493 749325702448 749325702448	1452851394641 5545435554894641 1111111111111111111111111111111111		7777565066531	.030704 307790336704 .02023336704 .02023336704 .0203336704 .02036 .0204 .0004		767 777 757 757 757 757 757 757 757 757	J10J3 725.4	
				PT	X/D	STATIC PRESS.(PSIA)	TW/T8	19 (F) 77.3	PPFSS	,
	A V Ē	RAGE 8ULK 6403	( ZEYNOLD 32.		90.3 RAGE PAP Avera	92.1 AMETERS BETWE Ge Wall Reynd 40754.	1:45 EN PRESSUE	343.3 Vêrage F <b>?</b> Vêrage F <b>?</b>	- 472F-02 - 146E+01	TCF
TIN = Pr,IN	AVE 79.5 F.	RUN 157 7 JUT = 5 7 GR/RES	( REYNOLD 32. 74. 0ATE 555.5 F, 0 = .11	12 4 VE 03 4 AS 03, 4 S 04 AS	90.3 RAGE PAP AVERA 02/75. G 5 FLJW R 02, TACH	40754. AMETERS BETWE GE WALL REYNO 40754. ATE = 25.9 LB (2) = .170.	1.45 EN PRESSUS LDS ECULAR WT MACH (16)	343.3 RETAPS VERAGE FR 132.2 <sup>26.97</sup> 132.2 <sup>349</sup>	- 4721-02 - 1465+01 IICTION FAC 00479 - 79 -	- 105 - 165 val -
TIN = PR·IN TC	4VE 79.5 F .713 X/D 1.1	RUN 157 RUN 157 TOUT 157 GR/RE30 (F) 243.2	( 2EYNOLO 22. 74. 0ATE 555.5 F. 11 14/TB 1.305	12 4 VE 0345-	90.3 RAGE PAP AVERA 02/75. GR 02/75. GR	92:1 AMETERS BETWE GE WALL REYND 40754. AS AIR 25.9 KL (2) = .170. HL/3GAS .109	1.45 EN PRESSU LDS AN ECULAR WT /HR. I = /HR. I = /HACH (16) BULK NUSSELT 409.60	343.3 S = TAPS VEFAGE FR 132.2 132.2 132.2 1 1 2 2 4 1 2 2 4 1 2 2 4 1 2 4 1 2 4 1 2 4 4 1 1 4 1 1 4 1 1 1 1 1 1 1 1 1 1 1 1 1	- • 72 - 02 • 146 - 01 ICTION FAC 00479 7. E - 7. 7. SUPR - 7. SUPR - 7. SUPR - 7. SUPR -	
TIN =	AVE 79.5713 X/D 1.2 2.1 2.1 2.1 2.1 2.4	RUN 150 RUN 150 TOUT 150 TOUT 150 TGR/RE30 T431-5 545-8 545-8 545-8 545-8 545-8	( REYNOLD 22. 74, DATE 555.5 F. 1.305 1.305 1.305 1.305 1.303 1.303 1.303 1.303 1.303 1.303	12 AVE 05 01.6	02 / 15 90.3 RAGE PAP AVERA 02 / 15 FL JWA 03 FL JWAC 303277 791621 759533 725954 759537 7259574	92.1 92.1 AMETERS BETWE GE WALL REYND 46754. AS AIR . MCL 46754. 101 HL/DGAS 101 037 033 033 039	ECULAR WT ECULAR WT MACH(16) BULKELT 40900407 136.27 136.27 1212.42	343.3 S = TAPS VEPAGE FR 132.2 132.2 4 132.2 4 1 1 1 1 1 1 1 1 1 1 1 1 1	- • • 72 02 • 1 - 6 E + 01 PICTION FAC • 00479 • 7 • 5 • E • 7 • 7 • 5 • E • 7 • 7 • 5 • E • 7 • 7 • 5 • • 7 • 7 • 7 • 5 • • 7 • 7 • 7 • 7 • 7 • 7 • 7 •	
TPR TC 23456765012	AVE F13 7. / 12465690 7. / 12465690 12111465690	R AG E 800 R AG E 800 N 1 4 50 N 1 4 50 R DUR 1 8 30 N 1 4 50 R DUR 1 8 30 R DU	<pre>X 2 E Y Y OLC Z 2 . Z 4 0 A T E 555 .5 F , 1 T 4 / TB 1 .305 1 .305 1 .305</pre>	1 2 4 5 5 4 4 5 6 5 4 4 5	90.3 90.4 90.4 90.4 90.4 90.4 90.4 90.4 90.4	92.1 92.1 AMETERS BETWE GE WALL REYNO 40754. AS AIR . MCL 40754. (2) = .170. HL/3GAS .109 .161 .037 .0333 .0367 .037 .0333 .037 .037 .037 .037 .037 .0	ECULAR WT MACHIGON 101 102 102 102 102 102 102 102 102 102	343.3 S = TAPS V = FAGE FA V = FAGE FA 2 = 2 2 = 2 2 = 2 1		TCF TCF TCF TCF TCF TCF TCF TCF TCF TCF
TPP T 2345676601234567	AVE F1 9 / 121114659004644 7 X 12344590004644	AGE 690 AGE 690 NT R TUR WE319964916335236 249744549243335236 249744549243335236 249743335236 24974937935236 249749335236 299966	<pre>X = Y + OLC Z = Z = Z = Z = Z = Z = Z = Z =</pre>	12 AVE 55 445 67.6	90         PAR           91         PAR           92         PAR           93         PAR           94         PAR           95         PAR           96         PAR           97         PAR           98         PAR           99         PAR           90         PAR           91         MC           92         PAR           91         NC           92         PAR           91         NC           92         PAR           93         PAR           94         PAR           95         PAR	92.1 AMETERS BETWE GE WALL REYND 40754. AS AIR 26.9 (2) = .170. HL/3GAS .109 .037 .0333 .0337 .0333 .0337 .043 .043 .045 .0568 .445	ECUL ALIO EN PRESSUA EN PRESSUA EN PRESSUA ELDS ELDS ELDS ECUL ALIO ECUL ALIO E	343. 343. 95 TAPS VEFAGE F4 VEFAGE F4 132.253 132.253 1331.1331.1331.1331.1331.1331.1331.13		

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TPR IN TC 23455789C11234567	7. 7. 7. 7. 7. 7. 7. 7. 7. 7.	RU17 / 1 3 4 3 4 5 6 6 5 3 2 8 2 1 2 8 3 4 3 1 4 3 3 4 5 7 5 6 5 3 2 8 2 1 2 8 3 9 8 0 1 1 1 1 2 2 2 2 2 2 2 2 2 2 2 2 2 2 2	1	3         3	S AIR , MOLI TE = 2C.1 L3. 2] = .157.1 HL/3GAS .104 .209 .0254 .02555 .02555 .02555 .025555 .025555555555		• 20.97 2.2 4 9 5. E • 6 • 174. T. SUD • • 3 100 3.7 1 574 3.3 1 574 7.2 1 574 7.2	2 • 33C 0 V F 73379 45 6 6 6 5 5 11 1 1 1 1 1 2 2 9 • 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0
	4 V E	RAGE BULK	PT 1 2 Reynolos	X/D 90.0 RAGE PARA AVERAG	STATIC PRESS.(PSIA) 71.2 METERS BETWEN 5 MALL REYNOL	TW/TB 1.05 1.15 En pressurg Los Ave	TAPS CAUE FRICTION FA	D2 D1 ACTOP
TIN . FR.IN	74.5 F.	RUN 1601 TCUT = 3 Gripesa	4. DATE 04 15.2 F. Mai	(08/75, GA 55 = L]# PA -33. Mach(	S AIR / MOL Tc = 20.1 La 2 = 157.1	ECULAR #T. /MR. 1 = 7 Mac4(16) =	■ 23.97 5.2 1495. E = 4	•• 425 VELTS
t 2345678901234567	X 121114542796231C8 121114542796231C8 123445973699		TW/TB 1234 123366 1366 1442229 1442229 1442229 1442229 1442229 1442095 133552 1335427 1335427 13246 13246 13246 13246 13246 13246 13246 13246 133546 13266 133546 13266 133546 13266 133546 1326 1	BULK BULK BS 500 50 50 50 50 50 50 50 50 50 50 50 50	HL/JGAS .119 .2051 .0740 .0332 .0332 .0335 .0359 .0349 .0444 .0448 .0448 .0448 .0448	Tiezoobsseestootoj Kiivootoj USSOteitoj USSOteitoj USSOteitoj USSOteitoj USSOteitoj USSOteitoj USSOteitoj USSOteitoj USSOteitoj	CGAR 9 - 12 GAR 9	G+ .C0101477 .C011178934 .C0011178934 .C0011178934 .C00111199464 .C00111199464 .C00111199464 .C00111199464 .C00111199464 .C0011114893 .C001114893 .C0011114893 .C0011114893 .C001114893 .C0011114893 .C0011114893 .C001114893 .C0011114893 .C0011114893 .C001114893 .C0011114893 .C001114893 .C001114893 .C001114893 .C00114893 .C00114893 .C00114883 .C001148
			рт 1 2	×/0 5 -0-1	STATIC PRESS.(PSIA) 74.5 70.2	TW/T3 1.11 1.29	T3 PP=25 (7) D57557 72.7 - 7.13E-( 72.7 - 6.625 16+5+	02
	<b>▲</b> VE	RAGE BULK 4978:	REYNOLDS	ERAGE PAPA Averag	METERS BETWE E WALL REYNJI 37493.	EN PRESSURE LOS AVE	TAPS RAGE PPICTION F .00513	ACTCP

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TIN = PR+IN	-75.4 F, • •720	TOUT = 4 , GR/RESQ	43.3 F, 	MASS FLJW R 3E-03, MACH	(2) + .136+	MACH(lè) =	-214. T.SURR :	5.725 VO • 30.3 F
TC 2345678901234	x 12111454eeeea		T 11006000 1 11055555542269 1 1.55555542269 1 1.55565226 1 1.555554226 1 1.555554226 1 1.555554226 1 1.555554226 1 1.5555554226 1 1.5555554226 1 1.55555555555555555555555555555555555	<pre></pre>	HL/7GAS .138 .225 .3439 .3439 .3449 .0449 .0441 .0445 .0445 .0557	L22:17.19.19.49.40 L22:17.19.19.49.4.9.40 L22:17.19.19.4.9.4.19.59.7 L23:17.19.14.9.4.19.59.7 L1111.19.14.9.4.19.59.7 L1111.19.14.9.4.7.777	2 5 5 7 5 7 5 7 5 7 7 7 7 7 7 7 7 7 7 7	
15 16 17	e1.4 90.2 98.2	753.a 753.5 567.0	1.433 1.380 1.134	39798. 33778. 38070.	. 362 . 103 . 392	73.76 75.79 155.80	57279.7 55121.4 43121.0	.CCI .COI .COI3
			1	PT X/D	STATIC	TW/TB	79 <b>29</b> 55	
	<b>∆</b> ve,	RAGE BULK 4735	REYNCLD 5.	2 50.3 AVERAGE PAR S AVERA	AMETERS BETWE GE WALL REYNG 32532.	1.17 1.33 En pressure Los al	(F) 05FEC 74.3514E 419.3 .203E 419.3 .203E 419.3 .203E 514E .00525	T C2 +C1 FACTCF
TIN = PR•IN	AV 57	R4GE BULK 4735 4735 1001 = 5 5 63/RES3	REYNOLD 5. 4. Date 3. F.	2 50.3 AVERAGE PAR S AVERA G4/00/75, G MASS FLOW R 3E-02, MACH	AMETERS BETWE GE WALL REYND 32532. AS AIR , MOL ATE = 20.0 L3 (2) = .157.	1.17 1.33 EN PRESSUR LOS AL ECULAR #T. /MR, I = 11 MACH(10) =	(F) 05525 74.3 -5145 419.8 -2035 TAPS RAGE FPICTION 1 .00525	T C2 +C1 FACTCF
TIN = IN FPR - IN TC 2345076961	AVE; 76.8719 X/D 124864659 124209	R 4GE 8ULK 4735 R 4GE 4735 R 164735 R 164757 R 1647577 R 1647577 R 16475777 R 1647577777777777777777777777777777777777	REYNCLD 7. 7. 7. 7. 7. 7. 7. 7. 7. 7.	1 - 5 2 - 0 - 3 A V ER A GE PAR A V ER A V ER A V ER A	AS A IR	1.17 1.33 EN PRESSUR LDS ELDS ELUS E	(F) 055126 (419.3 -55135 (419.3 -5	T = C2 +C1 F ▲ C T C = • • • • • • • • • • • • • • • • • •
TP T 1111111111111111111111111111111111	AV 57 76.8719 X/D 12111465 1234659 12486420075750 77500 775700 77600 1234659 124864200 123465	R 4 G E 4 7 3 5 R 4 G E 4 7 3 5 R 4 G E 4 7 3 5 R 1 6 5 7 R 1	R E YN CLD R E YN CLD A TE C A TE C	1 - 5 5 - 6 5 - 6 4 V ER A GE PAR A	A S A IR (Z) HL/QGAS 130757 HL/QGAS 130757 HL/QGAS 130757 00551 00609 00551 00609 00551 00609 00551 00609 00551 00609 00551 00609 00551 00609 00551 00609 00551 00609 00551 00609 00551 00609 00551 00609 00551 0060 00551 0060 00551 0	1.17 1.33 EN PRESSUR LOS 4.1 ELOS 4.1 ELO	(F) 055145 74.355135 419.355135 14.95 FPICTION 25.97, 52 25.97, 52 27.97, 52	T C 2 + C 1 F A C T C F C 2 F A C T C F C 4 C 4 C 4 C 4 C 4 C 4 C 4 C 4

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T 2345678901234567	x 121114546779823.44 1245642376673108 12456473108	00000000000000000000000000000000000000		2	S C L2 6 6 30 0 4 6 0 8 0 8 7 4 4 L2 6 0 7 1 8 10 1 7 6 0 8 1 7 4 4 D 7 8 7 8 1 8 10 9 8 7 6 1 7 7 7 7 3 7 7 8 7 8 7 8 7 7 7 7 7 7 7 7 7 7 7 7	HL/9G4S 201 2262 100400 0055000000	T3-14167C272680179 L347C977268755554222 U547C977264555554222 3U1 3U1		2 5 6 7 1 7 7 1 7 1 7 7 1 7 7 1 7 7 1 7 7 1 7 7 1 7 7 1 7 7 1 7 7 1 7 7 1 7 7 7 1 7 7 1 7 7 1 7 7 7 1 7 7 7 1 7 7 7 1 7 7 7 1 7 7 7 1 7 7 7 1 7 7 7 1 7 7 7 1 7 7 7 1 7 7 7 1 7 7 7 1 7 7 1 7 7 1 7 7 1 7 7 1 7 7 7 1 7 7 1 7 7 1 7 7 1 7 7 7 1 7 7 1 7 7 7 1 7 7 7 1 7 7 7 1 7 7 7 1 7 7 7 1 7 7 1 7 7 7 1 7 7 7 7 7 7 7 7 7 7 7 7 7	
				PT	×/0	STATIC PRESS.(PSIA)	TW/TB	TB (F) 75.3	PRESS DEFECT 5135-02	
	<b>A</b> V :	ERAGE BULA 2875	( 254NOL0 30.	AVER DS	AGE PAR Avera	AMETERS BETWE Ge Wall Peynd 21742.	EN PRESSUR LDS	TAPS VERAGE FRI	ICTION FAC	TCP
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APPENDIX E. Thermocouple Conduction Error D. M. McEligot, D.C. Tribolet and B. Bingham Aerospace and Mechanical Engineering University of Arizona

A thermocouple attached to the outside surface of a heated tube acts as a fin or extended surface which lowers the temperature at the point of attachment. Since this point also serves as the thermocouple junction, it measures a wall temperature lower than the value which would occur without its presence. The difference is known as "thermocouple conduction error." Consequently, in an experiment such as the present study, the deduced Nusselt number is systematically increased unless corrected for this effect.

## Analyses

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Based on extended surface analyses, the normalized thermocouple conduction error  $\theta$  has been shown by Schneider [48] to be approximately

$$\theta = \frac{t_w - t_{TC}}{t_w - t_\infty} = \left[1 + \frac{2\pi k_{tube} \delta_{tube} \lambda r_o}{h_{TC} A_s} \frac{K_1(\lambda r_o)}{K_o(\lambda r_o)}\right]^{-1}$$
(E1)

where  $t_w$  is the temperature of the undisturbed tube,  $h_{TC}A_s$ is the thermocouple conductance and  $r_o$  is the effective radius of the thermocouple attachment.  $K_o$  and  $K_1$  are modified Bessel functions of the second kind of order zero and one, respectively [49,50]. The quantity  $\lambda$  is defined as

 $\lambda = \frac{h_o + h_i}{k_{tube} \delta_{tube}}$ 

and the thermocouple conductance is defined by the equation

 $q_{TC} = h_{TC}A_s(t_{TC}-t_{\infty})$ . Thus, in this approximation  $\theta$  is a function of the non-dimensional parameters  $\lambda r_o$  and  $M = k_{tube} \delta tube / (h_{TC}A_s)$ .

Using approximations to the Bessel functions, valid at small values of the argument, one may reduce equation (E1) further so that it takes the form

$$\theta \approx - \ln(\lambda r_o)/2\pi M$$
 (E2)

when  $2\pi M >> 1$ . This form is useful for estimates of the maginitude of  $\theta$  when desiring to determine whether it is significant in a given case. It is presented as Figure E1.

In calibration for thermocouple conduction error data are normally obtained without flow so a probe can be used to measure the tube wall temperature in the vicinity of the thermocouple. In this case,  $h_1 = 0$ . Examining Figure El, one can see that the effect of flow (i.e., non zero value of  $h_1$ ) is to increase  $\lambda r_0$  and reduce  $\theta$  for the same thermocouple attachment and environment.

Hess [29] extended and improved Schneider's analysis for application to electrically heated tubes with internal flow such as the present experiment. His representation takes the form

$$\theta = \frac{\frac{(1-h_o/h_{TC})}{h_{TC}} + \frac{2k_w\delta}{h_{TC}r_o^2} + \frac{\lambda r_o K_1(\lambda r_o)}{K_o(\lambda r_o)}}$$
(E3)

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Evaluation of thermocouple conductance, h<sub>TC</sub>As

For fine wire thermocouples in an environment at atmospheric pressure, free convection dominates as the mechanism for heat loss from the thermocouple compared to radiation. As an approximation, one may consider the total heat transfer coefficient and properties to be constant for the thermocouple wire and derive

$$q_{TC} = \sqrt{(h_r + h_{NC}) P k_{TC} A_{CS}} (t_{TC} - t_{\infty}) \tanh m\ell$$
(E4)

as by Schneider [48]. For long wires, i.e.,  $m\ell = \sqrt{\ell hP/kA > 5}$ , tanh ml approaches unity; then  $h_{TC}A_S$  becomes  $\sqrt{(h_r + h_{NC})Pk_{TC}A_{CS}}$ .

The wire can be considered a small body in large surroundings so

$$h_{r} \approx \varepsilon \sigma \left[T_{TC}^{4} - T_{\infty}^{4}\right] / (T_{TC} - T_{\infty})$$
 (E5)

The heat transfer coefficient for natural convection can be determined from a correlation of the form  $Nu_f = fn(Gr_fPr)$ . The Grashof number is typically small for wires of the size of our thermocouples. For the range  $10^{-3}$  <GrPr<10<sup>-1</sup> the curve recommended by Kreith [51, Fig.7-3] can be represented as

$$Nu_f = \frac{h_{NC}d}{K_f} = 0.315 + 0.8 [Gr_f Pr_f]^{0.18}$$
 (E6)

In an unpublished note, Hess, Deardorff and McEligot [52, included herein as Appendix F] examined available calibration data for radiating thermocouples attached in the parallel junction form of Moen [28]. From comparisons between predictions and measurements, they concluded that the effective radius of the thermocouple attachment,  $r_o$ , was approximately equal to the actual radius of the thermocouple wire. Consequently,  $A_s = \pi r_o^2 \approx A_{CS} = \pi d^2/4$ . The calibration data of Campbell [53] for an atmospheric environment are also in approximate agreement with the choice of d/2 as  $r_o$ .

## Application to present experiment

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The heat transfer coefficient from the outside of the tube,  $h_o$ , may be deduced from the heat loss calibration equation (B5), to be

$$h_{o} = [C_{1} + C_{2}(t_{w} - t_{\infty}) + C_{3}(t_{w} - t_{\infty})^{2}]/\pi D \qquad (E7)$$

For  $h_i$  either a correlation such as equation (19) or tabular data from the initial data reduction can be employed. In this experiment  $h_i$  is typically of the order of 200 Btu/hrft<sup>2°</sup>F or more.

As an example of the magnitude of thermocouple conduction error to be expected, Figure E2 has been plotted for thermocouple 10. The value of  $h_1$  was taken as 200 Btu/hrft<sup>2°</sup>F and the environmental temperature was assumed to be about 70°F for this presentation. The reduction in  $\theta$  with flow is clear and it is also seen that in the heat loss runs (no flow)  $\theta$  decreases slightly as the wall temperature increases.

For the tabular results of Appendix D the thermocouple conduction error was calculated from equations (E3) through (E7) and correlations (19) and (B5). Material and fluid



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properties used in evaluating the thermocouple conductance were based on the temperature at the junction and its resulting film temperature,  $t_f = (t_{TC} + t_{\infty})/2$ , as appropriate. While a numerical solution could be applied to improve the analysis predicting the thermocouple conductance, such sophistication does not appear warranted due to the uncertain knowledge of several quantities and the small magnitude of  $\theta$ .

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APPENDIX F. Radiating Thermocouple Conduction Error W. G. Hess<sup>1</sup>, A. F. Deardorff<sup>2</sup> and D. M. McEligot Energy, Mass and Momentum Transfer Laboratory Aerospace and Mechanical Engineering Dep rtment The University of Arizona Tucson, Arizona

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In space environments and space simulation chambers, temperatures are often measured with thermocouples attached to exposed surfaces. The primary means of energy exchange then are conduction through the solid material and thermal radiation. In our Laboratory, we .'so often use a vacuum environment to minimize and/or localize the heat loss from thin-walled tubes in which we perform internal convective heat transfer measurements [38]. In these situations the thermocouple attachment usually acts as a radiating fin which reduces the local surface temperature near the point of measurement. This systematic effect may be called the radiating thermocouple conduction error.

Schneider [48] presents an analysis to predict the thermocouple conduction error in a convective environment by idealizing the thermocouple as a cylinder mounted perpendicular to the surface at a single point.

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Including energy generation in the wall by electrical resistive heating and energy transfer from the surface opposite the idealized thermocouple, one may extend Schneider's result to

$$\frac{T_{TC} - T_{w,u}}{T_{w,u} - T_{\infty}} = \frac{(h_o - h_{TC})r K_o(\lambda r)}{(h_i + h_{TC})r K_o(\lambda r) + 2\lambda k_w \delta K_1(\lambda r)}$$
(F1)

where the heat transfer coefficients may represent convective or radiative processes as appropriate. In the case of infinite radiating thermocouple leads, the effective heat transfer coefficient over the contact area of the thermocouple may be shown to be

$$h_{TC} = \frac{k_{TC} \left[ \frac{2\beta}{5} (T_{TC}^{5} - ST_{TC}T_{\infty}^{4} + 4T_{\infty}^{5}) \right]^{\frac{1}{2}}}{T_{TC} - T_{\infty}}$$
(62)

if its emissivity and thermal conductivity are constant. Thus, provided that the material properties are known, prediction of the radiating thermocouple conduction error reduces to the problem of determining the effective thermocouple radius, r.

For many applications the parallel type thermocouple junction, shown in the insert of Figure Fl, is more accurate than the more common cross type junction because the location of the measuring plane is effectively on the tube surface rather than being spread perpendicular to it [28].

Rather than satisfying the idealization of a single cylindrical interface between the thermocouple and the surface, the attachment region for the parallel junction consists of two roughly elliptical areas slightly separated from each other. Accordingly, the objectives of the present work were taken to be (1) to determine r for a parallel junction configuration and (2) to investigate the reproducibility of the conduction error when such thermocouples are produced by using normal laboratory standards for equipment construction.

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Measurements were conducted on three circular test sections of 0.010 inch thick Inconel 600, two feet long. Premium grade bare Chromel and Alumel thermocouple wires of 0.005 inch diameter were spot welded to the test section by the electrical discharge technique. Circumferential distance between the two wires was approximately 1/8 inch and the attached area of each covered approximately one to two wire diameters. Tests included about fourteen such thermocouples with all wires taken from the same spools.

These resistively heated test sections were mounted in glass vacuum chambers. With no intermal flow,  $h_i =$ quals zero and  $h_o$  can be determined from the tube emissivity which one also deduces from the tests. "Undisturbed"

tube wall temperatures, T<sub>w,u</sub>, were determined with a traveling internal thermocouple probe, also of premium grade Chromel-Alumel, which measured the wall temperature profile axially betwenn the thermocouples. Calculations show the maximum temperature drop through the wall to be less than 0.01°F so the thin wall idealization is valid. Readings were accepted without correction for deviation from standard N.B.S. emf tables since Hoskins Manufacturing Company certified the deviation as less than 1°F.

<u>Results</u> are demonstrated on Figure F1. The dashed curves are predictions based on equations (F1) and (F2) in conjunction with manufacturers' information for emissivities and thermal conductivities of the thermocouple wires and the tube. The solid curve represents predictions based on the measured emissibity of the Inconel tube used by Hess and on an effective thermoucouple radius equal to the actual wire diameter; otherwise the bases are the same.

Hess' data points are averages of the thermocouple readings for the central protion of the tube and they show the effective radius to be approximately equal to the wire diameter or slightly less. The measurements of Swearingen and of Reynolds and Deardorff are from test sections with different thermal histories, hence

emissivities, but of the same materials and dimensions. Their calibrations suggest that r is about one-half the wire diameter. Different welding jigs were used for each and, consequently, the region of attachment varied from test section to test section but would be approximately uniform for different thermocouples on the same test section. Accordingly, one would expect the level of the thermocouple conduction error to vary from test section to test section as it does in Figure F1.

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We <u>conclude</u> that the effective radius of the thermocouple attachment is approximately one-half to one wire diameter when constructed in the manner described. One may use this observation with manufacturers' information and equations (F1) and (F2) to determine whether the systematic error will be significant in his specific application. (For the results shown, in an internal convective heating experiment with  $T_w \simeq 1000^{\circ}F$  and  $T_b \simeq 900^{\circ}F$ , the resulting error in Nusselt number would be 5 to 10 per cent.) If such predictions indicate that the errors would be important, we recommend individual calibration since the values of a number of the pertinent input variables are not readily available.



## NOMENCALTURE

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h	Heat transfer coefficient; h <sub>i</sub> , inner surface: (thermocouple) surface.
k	thermal conductivity
к <sub>о</sub> , к <sub>1</sub>	Bessel functions
r	effective thermocouple attachment radius
ттс	temperature measured by thermocouple
T <sub>w,u</sub>	"undisturbed" wall temperature
β	thermocouple heat transfer constant, $2\sigma\epsilon/(k_{TC}r)$
δ	wall thickness
ε	emissivity
λ	wall heat transfer constant, $[h_0 + h_1)/(k_w r)]^2$
σ	Stefan-Boltzman constant

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