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Experimental Investigation of Turbulent Heat Transfer in Straight and Curved Rectangular Ducts

by

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ABSTRACT

An experimental apparatus has been constructed and tested to examine the convective heat transfer in straight and curved ducts of rectangular cross-section. The channel can be configured with both walls independently heated at a constant heat flux, or one wall at a constant heat flux and the opposite wall adiabatic. Local and average Nusselt numbers can be calculated, and used to evaluate the effect of curvature on heat transfer.

Experiments were conducted at steady state for turbulent flow with one wall heated at a constant heat flux and the opposite wall adiabatic. The heat transfer characteristics of the straight and curved sections, on both the inner and outer walls, were compared. The heat transfer rate of the concave curved wall proved to be enhanced over both the convex curved wall and both straight sections.

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LIST OF SYMBOLS

Symbol	Meaning	Units
A	cross-sectional area of the orifice	m ²
A _c	cross-sectional area of the channel	m ²
A pipe	cross-sectional area of the pipe	m ²
A _{PL}	area of the heated wall	m ²
C _{pair}	specific heat of air at constant pressure	J/KgK
Dc	channel height	m
De	Dean number	
D _{hd}	hydraulic diameter	m
D _{orf}	diameter of the orifice	m
D pipe	diameter of the pipe	m
F _{wo-wi}	radiation shape factor	
а ^с		Kg m N sec 2
ħ	average heat transfer coefficient	W/m ² C
к	flow coefficient	
^K air	thermal conductivity of air	W/mC
K ins	thermal conductivity of insulation	W/mC
• m	mass flow rate of air	Kg/sec
Nu	local Nusselt number	
Nu	average Nusselt number	
Pr	Prandtl number	
Patm	atmospheric pressure	N/m ²

Symbol	Meaning	Units
P ₁	pressure upstream of orifice	N/m ²
Pwet	wetted perimeter of channel	m
Q _{air}	heat convected to air	W
Q _{li}	heat lost through inner wall	W
Q _{lo}	heat lost through outer wall	W
° _p	power supplied	W
°r,	heat transferred by radiation	W
R	gas constant for air	Nm/Kg K
Red	Reynolds number based on channel height	
Re hd	Reynolds number based on hydraulic diameter	
^{Re} pipe	Reynolds number based on pipe diameter	
Ri	radius of curvature of inner convex wall	m
RPR	electrical resistance of precision resistor	Ω
R _R	total radiation resistance	m ⁻²
Та	Taylor number	
T _{blk}	bulk temperature of flow	с
T _{in}	average flow inlet temperature	с
^T ins,il	temperature between Plexiglas and first layer of inner insulation	С
^T ins,i2	temperature between first and second layers of inner insulation	С
^T ins,i3	temperature between second and third layers of inner insulation	С
^T ins,ol	temperature between Lexan and first layer of outer insulation	С
^T ins,o2	temperature between first and second layers of outer insulation	с

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Symbol	Meaning	Units
^T ins,03	temperature between second and third layers of outer insulation	С
Torf	temperature of air flowing downstream of orifice	С
Tout	average flow outlet temperature	С
T _{wi}	average temperature of inner wall	С
T _{wo}	average temperature of outer wall	С
V _{PR}	voltage across precision resistor	v
v _H	voltage across the wall heater	v
W _{id}	width of channel	m
Y	expansion factor	
β	ratio or orifice diameter to pipe diameter	
εcu	emissivity of copper plate	
Y	ratio of specific heats of air	
^µ air	dynamic viscosity of air	Kg/m sec
^ρ air	density of air	Kg∕m ³
σ	Stefan-Boltzmann constant	W/m^2K^4
ΔP	pressure drop across the orifice	N/m^2
ΔT	mean temperature difference	с
۵X	thickness of insulation layer	m

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I. INTRODUCTION

A. TAYLOR-GORTLER VORTICES

Since the early part of the twentieth century, considerable research has shown that the fully developed laminar flow along a concave wall does not remain twodimensional [Refs. 1, 2, 3]. The flow instead forms a system of spiral vortices, of counter rotating pairs, whose axes are aligned in the principle flow direction. This phenomenon is the result of the variations in the centrifugal forces acting on the fluid particles, and is known as Taylor-Gortler vortices. Figure 1 illustrates the type of fluid motion just described.

In a channel that is curved in the streamwise direction, those fluid particles located in the center of the flow cross-section are subject to higher centrifugal forces than those fluid particles traveling along the channel's outer boundary wall. As a result, the fluid in the center of the channel moves outwardly toward the concave boundary. As the process continues, the fluid particles near the boundary wall, move in a spanwise direction, and finally radially inward replacing the outwardly moving particles. These particles then come under the same centrifugal force and the process repeats itself. The resulting cyclic motion causes the formation of the counter rotating Taylor-Gortler



Figure 1. Schematic of Taylor-Gortler Vortices in a Curved Channel.

vortices, considered to be primarily a laminar flow phenomenon.

It has been observed [Ref. 4] that the heat transfer rate from flow along a concave curved wall is greater than that for flow along a straight wall, a phenomenon attributed to the additional mixing provided by the secondary motion of the Taylor-Gortler vortices.

There are many possible applications that could result from a more thorough understanding of Taylor-Gortler vortices and their effect on heat transfer and fluid flow characteristics. Two such applications could lead to improved heat exchanger designs and improved turbine blade cooling [Refs. 5, 6, 7].

B. HISTORY

The stability of an inviscid fluid flowing past a curved boundary, was first considered by Lord Rayleigh in 1916 [Ref. 8]. By assuming that the fluid was non-viscous, he determined that for the motion to remain stable, its circulation must increase with increasing radius. G. I. Taylor, in 1923, [Refs. 1, 9], continued this study with an extensive analytical and experimental study of viscous fluids. His investigations focused on the flow between two cylinders, in which the inner cylinder rotated while the outer cylinder remained stationary. Taylor ascertained that such flows become unstable when the value of a dimensionless parameter exceeded a critical value of 41.3.

The parameter, known as the Taylor number is defined as:

where 'd' is the width of the gap, assumed small when compared to 'Ri', the radius of the inner cylinder, and 'Re' is the Reynolds number based on the peripheral velocity of the inner cylinder. Taylor determined that for those cases in which the value of the Taylor number exceeded the critical value, a secondary motion developed and the Taylor vortices formed. Figure 2 illustrates this fluid motion.

Instability of a similar nature is also observed when a viscous fluid flows in a curved channel due to a pressure gradient acting along the channel wall. This problem was first considered analytically by W. R. Dean [Ref. 10] in 1928, for a channel formed by two concentric cylinders, where the radius of the inner cylinder was large in comparison to the spacing between the inner and outer cylinder walls. Dean concluded that there would be an initiation of flow instability, and the formation of vortices, when a dimensionless parameter, the Dean number, exceeded a value of 36. The Dean number is defined as:

$$De = Re \sqrt{\frac{d}{Ri}}$$

where 'd' represents the channel half-width, 'Ri' is the inner cylinder radius, and 'Re' is the Reynolds number based



on the mean velocity of the undisturbed flow and the channel half-width 'd'. The analytical work of Dean was later verified by W. H. Reid [Ref. 11], using an approximate solution.

In 1940, H. Gortler [Ref. 2] studied the stability of laminar boundary layer profiles on curved walls under the influence of small disturbances. He found that these disturbances were similar to the vortices studied by G. I. Taylor. Using approximate numerical calculations, Gortler concluded that the disturbances of vortices were produced only on the concave boundary walls and that the overall flow profile appeared to remain laminar in nature. These results were verified, with an exact solution, by G. Hammerlin, and reported by H. Schlichting [Ref. 12] in 1955. A. M. O. Smith completed an even more extensive numerical analysis that further substantiated these findings [Ref. 3]. The results of these numerical solutions have recently been demonstrated with the use of hot wire anemometry, laser doppler systems, and flow visualizations techniques, [Refs. 13, 14]; and in 1976, Y. Aihara [Ref. 15] conducted a nonlinear analysis of Gortler Vortices.

As interest began to develop concerning the effects of the secondary flows associated with the Taylor-Gortler vortices, many studies were published concerning the influences of these vortices on heat transfer in the laminar

and turbulent flow regimes. In 1955, F. Kreith [Ref. 16], studied the influence of heat transfer with respect to the curvature of the boundary wall for fully turbulent flows. He concluded that the heat transfer from the heated concave boundary wall was considerably higher than that transferred from the convex boundary wall of the same curvature and under similar turbulent conditions.

In 1965, L. Persen [Ref. 17], considered the special cases of very high and very low Prandtl number fluids and related the increase in heat transfer from a curved wall to the presence of the Taylor-Gortler vortices.

There has been only a limited amount of published literature dealing with the flow and heat transfer in curved channels of rectangular cross-section and large aspect ratios, where the aspect ratio is equal to the spanwise distance of the channel divided by the channel height. Much of what has been published involves the development of numerical approximations and solutions for heat and mass transfer in curved ducts of various geometries. K. Cheng and M. Akiyama [Ref. 18] developed a numerical solution for forced convection heat transfer with laminar flows in curved channels of rectangular cross-section, but only for small aspect ratios.

In 1976, A. A. Shibani and M. N. Ozisik [Ref. 19], using matched asymptotic expansion techniques for a wide range of

Prandtl numbers, solved the heat transfer problem between straight parallel plates with turbulent flow, for the case of uniform wall temperature.

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Other experimental and analytical studies that are worth citing with regard to this present study are Y. Mori [Ref. 20], who obtained results for hydrodynamically fully developed flows with constant wall heat flux in curved channels of square cross-section. W. M. Kays and E. Y. Leung [Ref. 21] reported solutions for turbulent flow heat transfer in a concentric circular tube annulus with a fully developed velocity profile and constant heat rate per unit length, for a fluid of Prandtl number 0.7. P. F. Brinich and R. W. Graham [Ref. 22] reported results for turbulent flows in a rectangular curved channel with an aspect ratio of 6, for the inner wall heated, the outer wall heated, and both walls heated.

In 1973, at the Naval Postgraduate School, R. J. Mckee [Ref. 23] confirmed the presence of Taylor-Gortler vortices in a curved rectangular channel of aspect ratio 40. M. Durao [Ref. 24] modified R. J. Mckee's channel and modeled it as infinite parallel plates with the outer wall heated and inner wall adiabatic. M. Durao and J. Ballard [Refs. 24, 25] investigated the effects of Taylor-Gortler vortices on the heat transfer in a straight and curved test section for laminar flows. R. Holihan, Jr. [Ref. 26] reported

results for laminar and transition flows. S. Daughety [Ref. 27] reported results for turbulent flows and J. Wilson [Ref. 28] reported results for transition and turbulent flows. Each time Mckee's apparatus was used it was modified slightly using prior recommendations to improve on the accuracy of the results.

II. INTENT OF STUDY

The purpose of this investigation was to build an improved curved rectangular duct of large aspect ratio which could be configured with a variety of boundary conditions, and yield more accurate data then previously attained. By combining new ideas and experimental equipment with past experiences and recommendations of R. Mckee and others previously mentioned [Refs. 23, 24, 25, 26, 27, 28], the effect of streamwise curvature on the heat transfer rate in a curved rectangular channel could be studied in greater detail while providing a smaller uncertainty in the results.

Some of the improvements incorporated in the new channel include: the ability to heat one or both walls over a wide range of constant heat fluxes; the ability to heat both walls at different constant heat fluxes; the use of more calibrated thermocouples which will allow greater precision in reading surface and air stream temperatures, combined with a better arrangement which will allow the calculation of the local heat transfer coefficient and Nusselt number; a longer transition region at the outlet which will allow for more mixing and a more accurate outlet temperature measurement. The use of an outer and inner section allowed both wall surfaces to be specially constructed with smooth

transitions between wall and heater surfaces in both straight and curved test sections, preventing unexpected tripping from laminar to turbulent flow. Finally, by constructing the channel with materials having low thermal conductivities and avoiding any heat sinks along the channel walls, heat losses could be minimized.

Taylor-Gortler vortices are known to enhance heat transfer in curved channels for transition and turbulent flows. This enhancement in the transfer of heat has been attributed to the secondary flow velocity components of the Taylor-Gortler vortices. These secondary components transport the heated fluid from the outer concave wall, inward toward the opposite wall of the channel, displacing the cooler fluid particles and causing them to move toward the heated concave wall. It was expected that similar results would be observed in this study.

This investigation was conducted using a single channel with a rectangular cross-section and constant aspect ratio, again defining aspect ratio as the spanwise distance divided by the channel height. The channel incorporates a straight and a curved test section. Each test section consists of opposing walls which can be independently heated. The results obtained in each test section at approximately the same flow rates were compared in an effort to determine the effects of the Taylor-Gortler vortices on the heat transfer

rate. Also, the results of both the straight and curved outer sections were compared to the results of Daughety [Ref. 27], and Wilson [Ref. 28] to compare the new channel to the old one.

The straight section results of this study were compared to the Dittus-Boelter correlation using hydraulic diameter [Ref. 29], and the results of Kays and Leung for turbulent and transition flow in annular passages [Ref. 21].

The curved section results of this study were compared with the results of Brinich and Graham [Ref. 22] for turbulent flow in a rectangular curved channel, and with Kreith [Ref. 16] for turbulent flows in concave and convex curved channels.

III. EXPERIMENTAL SETUP

A. DESCRIPTION OF THE APPARATUS

To meet the objectives of this investigation, a channel of rectangular cross-section and constant aspect ratio was constructed. The rectangular channel has four electrically heated wall sections, composed of a rubber heater bonded to a thin copper plate. All four heated wall sections are 30.48 centimeters long and 25.4 centimeters wide, with a heated area of 774.2 square centimeters. The rectangular channel is 0.635 centimeters high and 25.4 centimeters wide resulting in an aspect ratio of 40. The cross-sectional area of the channel is 16.13 square centimeters. The wetted perimeter is 52.07 centimeters, and the hydraulic diameter is 1.239 centimeters. Figure 3 shows a cross-sectional view presenting both the straight and curved sections of the channel.

The channel, shown in Figure 4, has a straight entrance region of 76.20 centimeters, which insures hydrodynamically fully developed flow prior to entering the straight heated test sections. The entrance region is followed by two parallel 30.48 centimeter straight heated test sections, which oppose each other, one in each wall. Following the straight heated test sections is a 15.32 centimeter straight



Figure 3. Schematic Cross-sectional View of Channel.



exit region, that completes the straight test section. The straight test section is immediately followed by the curved test section which subtends a 180 degree arc, with the inner wall having a radius of curvature of 29.7 centimeters. The curved test section has a 52.63 centimeter unheated convex curved section prior to the curved heated sections. The two concentric curved heated sections, inner one convex and outer one concave, are 30.48 centimeters long and the outer heated wall subtends an arc of 58.8 degrees. The curved heated sections are followed by a 10.16 centimeter convex curved exit region which completes the curved test section. Finally, the channel is completed by a 92.0 centimeter straight exit region which allows sufficient mixing prior to reading the outlet temperature. The entire channel was then thermally insulated with a plastic foam material.

An entrance bell constructed of Plexiglas is connected to the inlet of the straight entrance section of the channel. It was designed and manufactured according to ASME nozzle standards, with an elliptical curved base on a major axis equal to 25.4 centimeters and a minor axis of 2.54 centimeters. A fiberglass filter was attached to the entrance nozzle to prevent foreign matter from entering the channel. An aluminum exhaust nozzle was attached to the exit of the channel and directed the flow from the channel into a two inch diameter pvc pipe. The piping contained a

mass flow rate measuring section which could be fitted with standard ASME concentric orifices.

A total of one hundred and eighty-nine copper-constantan, 30 gauge thermocouples are located throughout the test apparatus. Only ninety-nine thermocouples are read at anyone time, through seventy-five data channels, to record the desired temperatures. Each heated wall section contains thirty thermocouples which are read individually and then averaged to give the wall section surface temperature. Four thermocouples span the width of the channel at the entrance and exit to give the average entrance and exit air temperatures. Three groups of five thermocouples are connected in parallel and inserted between each layer of insulation at each of the heated test sections. Each group of five thermocouples only uses one data channel, and these extra thermocouples comprise the difference between the number of thermocouples read and the data channels available. Finally, there is one thermcouple to measure the air temperature at the orifice.

B. DETAILS OF THE CONSTRUCTION

Each of the four heated wall sections was composed of a copper plate laminated to a wire wound silicone rubber heater. The rubber heater was made by Watlow Industries. There are thirty copper-constantan thermocouples placed between the copper plate and rubber heater. Figure 5 shows



Legend: numbers above location marks are outer and inner thermocouple numbers, numbers in parentheses below location mark apply to inner plates only.



the placement of the thermocouples on the back of the copper plate. The copper plate is 30.48 centimeters long, 26.72 centimeters wide, and 0.155 centimeters thick, but only the middle 25.4 centimeters of copper plate width is heated and centered in the channel. The back of the copper plate has a pattern of slots milled 0.076 centimeters deep. The pattern of slots for the two inner plates is different from the pattern of slots for the two outer plates. See Figures 6, 7, 8, and 9. The calibrated thermocouples were numbered and expoxied into the slots, as shown in Figures 10 and 11. The rubber heater has the same length and width dimensions as the copper plate, but only heats the middle 25.4 centimeters of width. This results in a heated wall area of 774.2 square centimeters. The excess copper plate is sandwiched between the wall and the side spacer thereby helping to support the copper plate in the channel. Dow Corning 3140 RTV was used to adhere the rubber heater to the back of the copper plate, and complete the heated section. Figure 12 and 13 represent a completed test wall.

The channel was built on a carefully constructed wood frame with steel rails, that also was used to support the entire channel. The inner wall was composed of two 0.343 centimeter thick and 29.85 centimeter wide pieces of Plexiglas laminated together, that had spaces milled out for the heated copper sections. The inner wall was constructed



Figure 6. Photograph of Milled Thermocouple Slots.





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Figure 8. Photograph of Both Inner Plates Showing Thermocouple Slot Pattern.



Figure 9. Photograph of Both Outer Plates Showing Thermocouple Slot Pattern.


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Figure 10. Photograph of Epoxied Thermocouples in Outer Curved Plate.



Photograph of Epoxied Thermocouples in Outer Straight Plate. Figure 11.



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Figure 12. Photograph of Convex Surface of Copper Plate.





by heating and formi: the Plexiglas around the wood frame. The Plexiglas wall was secured to the frame by screws tapped into the steel rails. The heated copper sections were then fitted into the machined spaces and two 0.635 centimeter thick and 2.223 centimeter wide Plexiglas spacers, that served as the sides of the channel and support for the copper plates, were attached by screws to the inner wall and steel rails. The two spacers were 25.4 centimeters apart and formed the width of the channel. Additionally, five flat head screws spanning the width of the channel at the joints where the copper and Plexiglas met were used to secure the copper plates and maintain an even joint. The seams were then worked to a smooth transition by applying "Bondo" to the joints and sanding. See Figures 14, 15, 16, and 17.

The outer wall was constructed in a similar manner, but using 0.318 centimeter thick and 29.85 centimeter wide pieces of Lexan. Lexan was used on the outer wall because of its flexibility and ease of bending without heat, which would be required when attaching the outer wall to the inner wall and frame. Also, the outer wall was formed in an outer wood frame so that the transition joints of the outer wall could be worked smooth like the inner wall. See Figures 18 and 19. Both walls were constructed so that no seams or butt joints existed on the inside of the channel, except



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Figure 14. Photograph of Channel Frame.



Figure 15. Photograph of Bottom View of Inner Straight Test Section.



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Figure 16. Photograph of Inner Stral & Wall at Copper-Plexiglas Joint.



Figure 17. Photograph of Inner Curved Wall at Copper-Plexiglas Joint.

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Figure 18. Photograph of Outer Frame and Lexan Wall Showing Straight Test Section.



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Figure 19. Photograph of Outer Curved Test Section.

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where the heated copper section joined the Plexiglas and Lexan walls, and all eight of those joints had a smooth transition.

Once both walls were completed, the outer wall and support frame was positioned onto the inner wall and main frame. The outer wall was aligned onto the inner wall and the outer wall support frame removed. The outer wall was then bolted to the inner wall and frame with steel bars running the length of the channel used as a spacer to insure even compression. Once this was completed the inlet and outlet nozzles were attached to the channel.

The sides of the channel and all connections between the channel, the pvc piping, and the orifice were sealed with Macklanburg-Duncan Silicone Rubber Sealant Caulk to ensure there was no leakage of air into the channel or the piping which could effect the recorded temperatures or mass flow rate. The entire channel was thermally insulated with three layers of 1/2 inch Armstrong Armaflex 22 Sheet Insulation, a flexible foamed plastic material. By positioning thermocouples between these insulation layers, the heat lost through the insulation could be computed at each test section. The heat loss through insulation surrounding the rest of the channel due to the heated air flowing through the channel was a small fraction of the total power supplied to the heater. The insulation was held in place by the use of silver duct tape and velcro strips attached to the frame.

C. INSTRUMENTATION

The data acquisition system used for this experiment was a Hewlett Packard 3054A Automatic Data Acquisition/Control System consisting of a 3456A Digital Voltmeter, a 3497A Data Acquisition/Control Unit, and four input panels capable of reading any mix of eight voltages or copper-constantan thermocouples. Also used in conjunction with this data acquisition system, were a Hewlett Packard 9826 Computer and 2671A Printer. The program used to reduce the data is listed in Appendix A.

The copper-constantan thermocouples were calibrated against a platinum resistance thermometer using a ROSEMONT Commutating Bridge model 920A, and a ROSEMONT Constant Temperature Bath model 913A. All thermocouples were calibrated between 20 and 85 degrees Celsius, and a second order polynomial was curve fitted to correct each thermocouple.

The mass flow rate in the channel was measured with a thin plate orifice measuring device which was placed in the 2 inch pvc piping between the channel and the compressor. In this study an orifice with a diameter of 2.7318 centimeters was used to measure hydraulic Reynold's numbers above 5000. The orifice device was installed in accordance with ASME Power Test Code [Ref. 30], with 3/8 inch pressure taps placed on either side of the orifice at 1 and 1/2

diameters, and one thermocouple inserted in the two inch pvc piping down stream of the orifice to record the temperature of the air flowing through the orifice. Pressure measurements were taken with Meriam vertical manometers, one whose fluid was water with a 0 to 60 inch range, another using water with a 0 to 30 inch range, and a third whose fluid was mercury, calibrated to read inches of water with a range of 0 to 415 inches. Atmospheric pressure was measured with a Princo precision barometer. The pressure readings were used to calculate the mass flow rate of the working fluid as outlined in [Refs. 30, 31].

Power was supplied to the heated test section by using two Lambda Model LK345A FM 60 volt power supplies in series. The power supplies are rated at 6.0 amps. maximum at 40 degrees Celsius. Since the electrical resistance of the rubber heater is not constant, but varies slowly with temperature, a precision resistor was calibrated and connected in series with the heater to allow the calculation of the instantaneous power being supplied. Both the outer and inner heated test sections have their own circuitry and precision resistor. The value of the precision resistor for the outer test sections is 2.0173 ohms, and the value for the inner test sections is 2.0084 ohms. See Figure 20 for diagram of circuit used to measure instantaneous power.



The working fluid was air at room temperature, which was drawn through the channel by an electrically driven Spencer Turbo Compressor, rated at 30 horsepower at 3500 rpm, and 550 cubic feet per minute at 70 degrees fahrenheit and one atmosphere. Photographs of the channel and associated test equipment are shown in Figures 21, 22, 23, and 24.





Figure 22.



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 Figure 23. Photograph of Spencer Turbo Compressor.



Figure 24. Photograph of Test Channel.

IV. EXPERIMENTAL PROCEDURES

The experimental procedures followed were the same for the straight and curved test sections. The procedure for starting up the experimental apparatus is as follows:

- (1) turn on the computer and data acquisition system
- (2) close globe valve at channel exit
- (3) close valve to manometer reading Pl
- (4) set compressor blast gate to 1/8 open
- (5) start turbo compressor
- (6) open intake valve on compressor 1/4 to 1 turn
- (7) open globe valve at channel
- (8) open manometer valve slowly to ensure Pl does not exceed manometer scale
- (9) turn on and adjust power supply to the heater
- (10) use intake valve on compressor to control mass flow rate through the channel

It was necessary to determine the approximate voltage setting required to bring the heated wall to approximately 50 degrees Celsius. This heated wall temperature would ensure a large temperature difference between the inlet and outlet temperatures, and the heated and unheated wall temperatures, while not damaging the apparatus. The larger the temperature difference the smaller the uncertainty in the results. The required power setting was found by

monitoring the heated plate temperature locally from the data acquisition system while the voltage to the heater was adjusted.

To compute the instantaneous power supplied to the heater the following relationship was used:

$$Q_{\mathbf{P}} = \frac{V_{\mathbf{H}} V_{\mathbf{PR}}}{R_{\mathbf{PR}}}$$

The precision resistor voltage and the heater voltage were read by the data acquisition system and the resistance of the precision resistor was a known constant. A diagram of the circuit used was shown in Figure 16.

The mass flow rate was monitored by the two manometers at the orifice. For each run, the atmospheric pressure, the pressure difference across the orifice, and the pressure upstream of the orifice were entered manually into the computer for the mass flow rate calculation. All other data was acquired by the data acquisition system directly from the experimental apparatus.

Once the apparatus was started properly, sufficient time was required to reach steady state prior to gathering data. Preliminary runs were performed to determine the time required for the test apparatus to reach steady state. For this experiment at least two hours were allowed for a cold test apparatus to reach steady state, and at least one hour

if the apparatus was already running. After waiting the required time, data was taken at ten minute intervals to ensure that a steady state condition had been achieved. The criteria for steady state was based on three variables:

(1) the difference between inlet and outlet temperature

(2) the heated test section wall temperature

(3) the unaccounted heat loss in the system When the first two of these quantities varied by less than two percent over a ten minute interval and the third was less than ten percent of the input power, it was considered that a steady state condition had been achieved.

Experiments were conducted at specific mass flow rates of air corresponding to Reynolds numbers between 7000 and 22000. The Reynolds number was based on the hydraulic diameter of the channel. For each mass flow rate, temperature and voltage data were taken automatically by the data acquisition system and immediately reduced and printed. The variables that were measured and used in this investigation were:

- (1) the temperature of the air entering the channel (T_{in}) (2) the temperature of the air leaving the channel (T_{out})
- (3) the temperature of the heated wall in the test section $(T_{wo} \text{ or } T_{wi})$
- (4) the temperature of the unheated wall in the test section $(T_{wi} \text{ or } T_{wo})$
- (5) the temperature between each layer of insulation at the test section (T_{ins})

- (6) the temperature of the air at the orifice (T_{orf})
- (7) the voltage across the precision resistor (V_{PR})
- (8) the voltage across the rubber heater $(V_{\rm H})$
- (9) the pressure upstream of the orifice (P1)
- (10) the difference in the upstream and downstream pressure across the orifice (ΔP)
- (11) the atmospheric pressure (P_{ATM})

The results obtained from the straight test sections served as a baseline for the comparison of the curved section results. An effort was made to match the mass flow rates for each test section as closely as possible.

V. PRESENTATION OF DATA

A. DATA REDUCTION

The constant heat flux surface for the heated wall of each test section was provided by the uniform electrical resistivity of the rubber heater. The insulated unheated wall was considered adiabatic since the heat losses through that wall were negligible. The straight portion of the channel upstream of the straight heated test section was of sufficient length to ensure that the air flow was hydrodynamically fully developed for the flow velocities of this experiment. The straight portion of the channel downstream of the curved section proved to be of sufficient length to ensure that the air flow exiting the channel was thoroughly mixed and that the measured outlet air temperature (T_{out}) was an average bulk temperature, as long as the hydraulic Reynolds number of the channel was less than 17000.

The analysis of the data required several quantities to be defined. The heat convected to the air was calculated using the equation:

 $Q_{air} = {}^{\bullet}C_{pair} (T_{out} - T_{in})$

where C_{pair} is the specific heat of the air at constant pressure and \dot{m} is the mass flow rate of the air.

The average heat transfer coefficient between the heated wall and the flow of air in the channel was defined by the equation:

$$\overline{h} = \frac{Q_{air}}{A_{PL} \Delta T}$$

where ' Q_{air} ' is defined above, ' A_{PL} ' is the area of the heated copper plate in the test section, and ' Δ T' is the difference between the arithmetic average heated wall temperature (T_w) and the average bulk air temperature (T_{blk}). The average bulk air temperature is defined as the arithmetic mean of the entrance and exit air temperatures (T_{in} , T_{out}).

The average Nusselt number was then calculated using:

$$\overline{Nu}_{hd} = \frac{\overline{h} D_{hd}}{K_{air}}$$

In this equation $'D_{hd}'$ is the hydraulic diameter and $'K_{air}'$ is the thermal conductivity of the air.

The hydraulic Reynolds number was calculated for each test run as follows:

$$Re_{hd} = \frac{\overset{m}{m} D_{hd}}{A_c \mu_{air}}$$

where again 'm' and $'D_{hd}'$ are the mass flow rate and hydraulic diameter of the channel, $'A_c'$ is the cross-

sectional area of the channel, and $'^{\mu}_{air}'$ is the dynamic viscosity of the air.

For the curved sections the Dean number was defined as:

$$D_{e} = \frac{\dot{m} D_{c}/2}{A_{c} \mu_{air}} \sqrt{\frac{D_{c}/2}{Ri}} = \frac{Re_{d}}{2} \sqrt{\frac{D_{c}/2}{Ri}}$$

where ' Re_d ' is the Reynolds number based on the channel height, ' D_c ' is the channel height, and 'Ri' is the radius of curvature of the inner convex wall.

A sketch of the control volume and a set of sample calculations for one test run of the curved section are given in Appendix B. A sample calculation for the uncertainty analysis is given in Appendix C.

B. RESULTS

The data obtained from each experimental run was reduced by the computer program listed in Appendix A, utilizing the expressions described in Part A above. The major parameters resulting from this process are shown in Tables I through IV. Tables I and II contain the straight sections results and Tables III and IV contain the results of the curved sections. A plot of the average Nusselt number versus hydraulic Reynolds number is given in Figures 25, 26, 27, 28, and 29 for the comparison of the results. "Error" bands have been indicated as a result of the uncertainty analysis.

TABLE I

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SUMMARY OF STRAIGHT OUTER WALL TEST RESULTS

Re _{hd}	Qair (W)	h (W/mC)	Δ Τ (C)	^{Nu} hd
7709	108.50	41.99	28.52	20.07
11953	135.26	61.57	25.49	29.51
15646	157.92	74.43	24.45	35.65
21455	217.53	87.42	27.36	41.95

Reynolds number and Nusselt number are based on hydraulic diameter. Results for Reynolds number of 21455 based on orifice air temperature instead of outlet air temperature.

TABLE II

SUMMARY OF STRAIGHT INNER WALL TEST RESULTS

Re _h đ	Q (Ŵ) ^r	h (₩/mC)	ΔT (C)	Nuhđ
7628	105.57	41.52	28.90	19.85
12169	138.36	59.50	27.42	28.50
15788	159.84	70.24	26.37	33.61
21312	204.42	81.82	27.84	39.17

Reynolds number and Nusselt number are based on hydraulic diameter. Results for Reynolds number of 21312 based on orifice air temperature instead of outlet air temperature.

TABLE III

SUMMARY OF CURVED CONCAVE WALL TEST RESULTS

Re _{hd}	De	Qair (W)	h (₩∕mC)	Т (С)	Nuhđ
7782	206	115.88	52.43	25.51	25.00
11726	311	157.45	68.16	29.91	32.53
15813	419	158.17	83.13	22.34	39.78
21246	563	231.31	93.02	27.20	44.52

Reynolds number and Nusselt number are based on hydraulic diameter. Results for Reynolds number of 21246 based on orifice air temperature instead of outlet air temperature.

TABLE IV

SUMMARY OF CURVED CONVEX WALL TEST RESULTS

Re _{hd}	De	Q (Øjr	h (₩/mC)	∆T (C)	Nuhd
7757	206	82.37	38.33	24.47	18.37
12021	318	115.50	57.62	23.92	27.65
15908	421	142.54	68.57	24.34	32.94
21603	572	164.08	79.28	23.43	38.11

Reynolds number and Nusselt number are based on hydraulic diameter. Results for Reynolds number of 21603 based on orifice air temperature instead of outlet air temperature.

Figure 25. Present Straight vs. Curved Section Results.



 LEGEND
BRESENT CONCAVE DATA
PRESENT CONVEX DATA Present Concave vs. Convex Section Results. HYDRAULIC REYNOLDS NUMBER Ö Figure 26.

AVERAGE NUSSELT NUMBER

Present Straight Outer vs. Inner Section Results. Figure 27.



30000 D = PRESENT CONCAVE DATA D = PRESENT STRAIGHT OUTER DATA Straight Outer Section Results. HYDRAULIC REYNOLDS NUMBER <u>а</u> о D 0 Present Concave vs. ۵ 10000 Figure 28. ۵ 4 7000 <u>0</u> F AVERAGE NUSSELT NUMBER

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Present Convex vs. Straight Inner Section Results. Figure 29.



As would be expected the results indicate an increase in the rate of heat transfer with increasing Reynolds number for both the straight and curved test sections. In addition, the heat transfer rate was highest in the concave curved section for each Reynolds number investigated, and lowest in the convex curved section.

In the low end of the turbulent regime, at a hydraulic Reynolds number of 7780, the rate of heat transfer was about twenty-two percent higher in the concave curved section than the outer straight section. As the hydraulic Reynolds number continued to increase, between 11700 and 15800, the rate of heat transfer was approximately eleven percent higher between the concave curved section and the outer straight section. At a hydraulic Reynolds number of 21000 the difference in heat transfer rate decreased to about seven percent.

The results obtained at a hydraulic Reynolds number of 21000 have a much larger uncertainty, because the orifice air temperature was higher than the average outlet air temperature. This signifies that the air was not fully mixed at the outlet, because no heat is added to the air after the outlet air temperature is read, yet the air temperature increased. For hydraulic Reynolds numbers greater than about 17000 the channel outlet section is not long enough to allow the air to mix properly.

The values obtained for the hydraulic Reynolds number of approximately 21000 are based on the orifice air temperature vice the outlet air temperature. This resulted in a Nusselt number higher than obtained by using the the outlet air temperature, but still lower than the actual Nusselt number, because of the small heat loss between the outlet nozzle and the thermocouple at the orifice. This heat loss cannot be accounted for and the temperature used as an outlet air temperature was still lower than the actual outlet air temperature.

In comparing the concave curved wall to the convex curved wall the rate of heat transfer was about thirty-six percent higher in the concave section for a hydraulic Reynolds number of 7780. For hydraulic Reynolds numbers between 11700 and 15800 the rate of heat transfer was twenty percent higher in the concave section, and at a hydraulic Reynolds number of 21200 the rate of heat transfer was seventeen percent higher.

VI. DISCUSSION AND CONCLUSIONS

J. Wilson [Ref. 28], in his study, reported that the actual presence of Taylor-Gortler vortices were verified by the observation of the liquid crystal isotherm distribution in the curved test section. He further reported, that at low values of hydraulic Reynolds number and through most of the transition regime, liquid crystal bands were oriented in the streamwise direction. As the Reynolds number increased turbulent forces began to dominate in the curved section and less evidence could be detected for the existence of the vortices.

The high aspect ratio of the channel provided the basis for the assumption that the experimental apparatus should behave like infinite parallel plates. Holihan's data [Ref. 26] tended to verify this assumption. His experimental data in the laminar flow region, approached the theoretical limit for average Nusselt number for parallel plates with one wall heated at a constant heat flux and the opposite wall adiabatic [Ref. 32]. Based on the assumptions mentioned above, comparisons were made with other analytical and experimental results that were of the same or similar problems.

A comparison of the experimental results of this study and the experimental results of Daughety [Ref. 27] and
Wilson [Ref. 28] are shown in Figure 30 for the curved section, and Figure 31 for the straight section. The channel aspect ratio was 40 for each of these studies and the experimental procedures were similar, with the exception of a slight modification in Wilson's procedures. Wilson [Ref. 28] made a correction to determine the temperature drop across the liquid crystal sheet covering the heated wall. By applying this temperature drop to the measured temperature, the true temperature of the heated wall would be known. After reviewing Wilson's data and experimental procedure it was detected that Wilson was over estimating the temperature drop across the liquid crystal by approximately five to seven percent, and this lead to an over estimation of T of ten to twelve percent. This resulted in an over estimation of the average heat transfer coefficient and the average Nusselt number of between ten and twelve percent. When Wilson's data is lowered by ten percent, both his data and Daughety's data fall within the uncertainty bands of the present study as shown in Figures 32 and 33. Differences were attributed to the inability to acccount for all of the heat transfer processes and the limitations inherent in any experimental work.

Ballard [Ref. 25] reported an increase in the rate of heat transfer between a concave wall and a straight outer wall of eleven percent for laminar flows. Holihan [Ref. 26] reported an increase of fifteen percent for laminar flows

Comparison of Present Data with Daughety and Wilson, Curved Section. 30000 CAVE DATA DATA EGEN ž HYDRAULIC REYNOLDS NUMBER 0 ۵ 4 □ = PRESENT
 ○ = PRESENT
 ○ = PRESENT
 + = DAUGHET 0 0 ۵ 10000 ۵ 0 ⊲ Figure 30. 7000 01 AVERAGE NUSSELT NUMBER

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Comparison of Present Data with Daughety and Wilson, Straight Section. Figure 31.



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Comparison of Present Data with Daughety and Corrected Wilson, Curved Section. Figure 32.

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and thirty percent for transition flows. Daughety [Ref. 27], reported an increase of twenty percent in the heat transfer rate for turbulent flows. Wilson [Ref. 28], reported an increase of forty-five percent for laminar flows, twenty percent for transition flows, thirteen percent for hydraulic Reynolds numbers between 5400 and 12000, and about twenty-six percent at a hydraulic Reynolds number of 24000. This compares very favorably with the present results, except for the data at a hydraulic Reynolds number of 21000. No conclusion can be made about the present results at a hydraulic Reynolds number of 21000 because of the high uncertainty of the data as previously mentioned. This area needs further experimentation before accurate conclusions can be made about the heat transfer rate at high hydraulic Reynolds numbers.

Kreith [Ref. 16] reported an increase in the rate of heat transfer along a concave wall as compared to a convex wall of twenty-five to sixty percent for hydraulic Reynolds numbers between 10^4 to 10^6 . This also compares favorably with the present results which indicate an increase of seventeen to thirty-eight percent when comparing concave wall to convex wall.

A comparison of analytical and experimental results with those of the present study for the straight test section is shown graphically in Figure 34. The Dittus-Boelter equation





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[Ref. 29] for heat transfer in a straight tube for constant wall temperature is plotted from:

 $Nu = 0.02 Re^{0.8}$

For a Prandtl number equal to 0.71. Additionally, the results of Kays and Leung [Ref. 21] for heat transfer in annular passages are shown for $r^* = 1.0$, which equates to parallel plates. When examined these results compare favorably with the present data.

The curved section data is plotted in Figure 35, and is compared to the experimental work of Brinich and Graham [Ref. 22]. The accuracy of the data points used from Brinich and Graham are subject to some error, in that the actual values are not given in their study and had to be taken from a plot of Stanton number versus Reynolds number.

After carefully examining the past and present data the trends and results are consistent. Heat transfer rates increase with increasing Reynolds numbers. Also, the concave curved section enhances the rate of heat transfer over the convex curved section and the straight test sections for laminar, transition, and turbulent flows. The presence of Taylor-Gortler vortices in the laminar and transition regime seems to be the major contributing factor to this enhancement. Whether the Taylor-Gortler vortices continue to exist in the turbulent flow regime remains to be seen.



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VII. RECOMMENDATIONS

The amount of experimental work remaining in the study of heat transfer in curved channels is still significant. The present experimental apparatus has been tested in the turbulent regime, and the test data proves that accurate results can be obtained for hydraulic Reynolds numbers below 17000, slightly higher when heating the straight test section. Additional experiments into the laminar, transition, and turbulent region should be conducted to better understand the heat transfer processes occurring in the straight and curved sections. By remaining below a hydraulic Reynolds number of 17000 accurate data can be obtained and significant conclusions can be drawn in the comparison of the curved sections to the straight section, and the concave curved wall to the convex curved wall.

To improve the accuracy of the data and results some slight modifications to the present channel should be done. First, the straight outlet section of the channel should be lengthened to ensure proper mixing of the heated outlet air. Second, some sort of honey combing, screening, or other type of device should be placed in the outlet section to enhance the mixing of the heated outlet air. Finally, extra thermocouples should be placed at the outlet and read more

then once during a data run to decrease the uncertainty of the outlet air temperature (T_{out}) and the uncertainty of the difference between the inlet and outlet air temperature. By applying these recommendations the uncertainty of the Nusselt number can be reduced to four percent or less.

Further experiments at hydraulic Reynolds numbers greater than 17000 also need to be examined. This can be accomplished if the above recommendations are taken along with one other. The channel walls need outer support to prevent the center of the channel from collapsing at the high hydraulic Reynolds numbers. Slight collapsing was noticed at hydraulic Reynolds numbers greater than 20000. This collapsing was caused by the laminating of two 0.3175 centimeter thick walls which were not as rigid as the single 0.635 centimeter thick walls used previously. The actual hydraulic Reynolds numbers that can be attained and still yield accurate results will be limited by the power of the compressor and the rigidity of the channel's walls.

Finally, once a strong data base can be obtained for each test section further experimentation can be conducted with different boundary conditions. For example, both curved walls heated at the same constant heat flux. Once this can be done a true comparison of the heat transfer rate of the total curved test section can be compared to the total straight test section, i.e., both curved walls versus both straight walls.

APPENDIX A: COMPUTER PROGRAM

FILE NAME: DRP2A LAST REVISION DATE: 21 OCTOBER 1985 AUTHOR: LT. GEORGE G. GALYO 1000 1 1010 1 1020 1 1030! PROGRAM FOR GATHERING AND REDUCING DATA WHEN HEATING ONE COPPER PLATE, 1040 1 WITH 10 MINUTES BETHEEN RUNS. PROGRAM IS MODIFIED TO ACCOUNT FOR BROKEN THERMOCDUPLE CHANNELS 5, 16, 21, AND 26 WHICH EFFECT OUTER CURVED HEATER, AND THERMOCDUPLE CHANNEL 47 WHICH EFFECTS INNER CURVED 1050 1060 1070 1080 HEATER. 1 1090! CDM /Co/ D(7),Aa(76,2) DIM Emf(79),T(79) 1100 1110 11201 1130 ! CORRELATION FACTORS FOR EMF TO DEGREES CELSIUS FOR CU-CO T/C. DATA 0.10086091,25727.94369,-767345.8295,78025595.81 DATA -9247486589,6.97688E+11,-2.66192E+13,3.94078E+14 1150 1160 READ D(+) 1170 1180 PRINTER IS 701 1190 BEEP INPUT "ENTER MONTH, DAY, AND TIME (MM:DD:HH:MM:SS)", Time\$ 1200 OUTPUT 709;"TD";Time\$ 1210 1220 1230 BEEP INPUT "ENTER INPUT MODE (0=3054A.1=FILE)".Im 1240 IF Im=1 THEN BEEP 1250 1260 1270 INPUT "ENTER NAME OF EXISTING DATA FILE",Oldfile\$ PRINT USING "20X,""These Results Are From Data File: "",10A";Oldfile\$ 1280 PRINT 1290 PRINT ELSE 1300 1310 INPUT "ENTER NAME OF NEW DATA FILE?",Newfile\$ PRINT USING "20X,""Data For This Run Is Stored In Data File: "",10A";Newf 1320 1330 ile\$ 1340 1350 PRINT PRINT END IF 1360 1370! ! ENTER THE CHANNEL CONFIGURATION. 1380 1390! PRINT USING "9X, "THE FOLLOWING CHANNEL CONFIGURATION WAS SELECTED: "" 1400 PRINT 1410 1420 BEEP INPUT "SELECT HEATER TYPE (0=STRAIGHT, 1=CURVED)",Itype IF Itype=0 THEN PRINT USING "12X,""Heating STRAIGHT Test Section.""" IF Itype=1 THEN PRINT USING "12X,""Heating CURVED Test Section.""" 1430 1440 1450 BEEP 1460 INPUT "SELECT HEATER POSITION (0-OUTER, 1-INNER)", Ipos IF Ipos=0 THEN PRINT USING "12X, "Heating OUTER Plate."" IF Ipos=1 THEN PRINT USING "12X, "Heating INNER Plate."" 1470 1480 1490 1500 PRINT 1510! 1520 1530! ! ENTERS FILES WITH T/C CALIBRATION COEFFICIENTS BASED ON HEATER TYPE. IF Itype=0 THEN ASSIGN @Filec TO "STRA_COEF" IF Itype=1 THEN ASSIGN @Filec TO "CURV_COEF" FOR I=0 TO 76 ENTER @Filec:Aa(I,0),Aa(I,1),Aa(I,2) 1540 1550 1560 1570 1580 NEXT I

1590 ASSIGN @Filec TO * 1600! 1610 ! INSERT DATA STORAGE DISK. 1620! BEEP 1630 INPUT "CHANGE DISK AND HIT ENTER",Ok IF Im=1 THEN ASSIGN @File TO Didfile\$ IF Im=0 THEN 1640 1650 1660 1670 CREATE BDAT Newfile\$,40 1680 ASSIGN @File TO Newfile\$ END IF ENTER 709;Time\$ CLEAR 709 PRINT USING "20X,""Month, Day, and Time: "".14A";Time\$ PRINT 1690 1700 1710 1720 1730 1740! ! FOR DATA RUN READ AND STORE ALL RAW EMF VALUES. 1750 1760! 1770 IF Im=0 THEN OUTPUT 709:"AR AF00 AL79" OUTPUT 722:"F1 R1 T1 Z1 FL1" 1780 1790 FOR 1=0 TO 79 OUTPUT 709; "AS" 1800 1810 ENTER 722:Emf(I) NEXT I 1820 1830 OUTPUT @File;Emf(+) 1840 1850 ELSE 1860 ENTER @File:Emf(*) 1870 END IF 1880 OUTPUT 709: "TD" 18901 ! CONVERT AND CALIBRATE ALL T/C READINGS TO DEGREES CELSIUS. 1900 1910! 1920 PRINT USING "9X," THE FOLLOWING DATA WERE RECORDED:""" 1930 PRINT 1940 FOR I=0 TO 79 IF I-38 OR I-39 OR I-77 OR I-78 THEN IF (Ipos-0 AND I-38) OR (Ipos-1 AND I-77) THEN 1950 1960 1970 Iadd=0 IF Ipos=1 THEN Iadd=39 PRINT USING "12X,""Heater Voltage (Vh) 1980 1990 = "",2D.3D,"" (V)""";Emf(38+Ia dd) 2000 PRINT USING *12X, **Resistor Voltage (Vr) = "",2D.3D,"" (V)""";Emf(39+Ia 2010 IF Ipos=0 THEN PRINT USING "12X,""Precision Resistor (Rpr) = 2.01729 (Ohm s)""" 2020 IF Ipos=1 THEN PRINT USING "12X, ""Precision Resistor (Rpr) = 2.00839 (Ohm s)"" 2030 PRINT END IF 2040 2050 2060! ! IF T/C DUT OF RANGE PRINT A QUE. 2070 2080! Aemf=Emf(I) IF Aemf<3.E-4 OR Aemf>3.E-3 THEN 2090 2100 2110 T(I)=999.99 2120 ELSE CALL Tvsv(Emf(I),Tt) IF I=79 THEN 2130 2140

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T(1)=Tt GOTO 2210 2150 2160 2170 END IF 2180 T(I)=FNTcorr(Tt,I) END IF 2190 2200 NEXT I 2210 2220! ! PRINT ALL T/C READINGS. 2230 2240! PRINT USING "12X,""OUTER PLATE:" PRINT USING "14X,""T/C Number: 2250 2260 2 3 5 1 2270 PRINT USING "14X,""Temp (C): **,6(3D.DD,2X)*;T(0),T(1),T(2),T(3),T(4),T (5) 2280 PRINT USING "14X,""T/C Number: 12"" 9 10 7 8 11 2290 PRINT USING "14X,""Temp (C): **,6(3D.DD,2X)*;T(6),T(7),T(8),T(9),T(10), T(11) 2300 PRINT USING "14X,""T/C Number: 18"" 13 14 15 16 17 "",6(3D.DD,2X)";T(12),T(13),T(14),T(15),T(2310 PRINT USING "14X,""Temp (C): 16),T(17) 2320 PRINT USING "14X,""T/C Number: 24"" 21 22 23 19 20 "",6(3D.DD,2X)";T(18),T(19),T(20),T(21),T(2330 PRINT USING "14X.""Temp (C): 22),T(23) 340 PRINT USING "14X,""T/C Number: 25 26 27 28 29 2350 PRINT USING "14X, "Temp (C): **,6(3D.DD,2X)*;T(24),T(25),T(26),T(27),T(28),T(29) 2360 PRINT 2370! ! CALCULATE AVERAGE TEMPERATURE OF OUTER HEATED WALL. 2380 23**90!** 2400 Sum=0. IF Itype=0 THEN FOR I=0 TO 29 410 2420 2430 Sum=Sum+T(I) NEXT I 2440 2450 Two=Sum/30 2460! ! CORRECTION FOR BROKEN T/C CHANNELS 5, 16, 21, AND 26 IN OUTER CURVED HE 2470 ATER. 2480! 2490 ELSE FOR I=0 TO 29 IF I=5 OR I=16 OR I=21 OR I=26 THEN 2530 2500 2510 Sum*Sum+T(I) NEXT I 2530 Two=Sum/26 2540 END IF 2550 PRINT USING "12X.""OUTER INSULATION:""" PRINT USING "14X.""I/C Number: 31 (T 2560 32 (Tinso2) 33 (Tinso3)" 2570 31 (Tinsol) "",3(3D.DD,8X)";T(30),T(31),T(32) 2580 PRINT USING "14X, ""Temp (C): 2590 Tinso1=T(30) 2600 Tinso2=T(31) 2610 Tinso3=T(32) PRINT 2620

PRINT USING "12X, "'ORIFICE TEMP (Torf) = "", 3D.DD, "" (C) """; T(33) 2630 2640 PRINT PRINT USING "12X."INLET TEMPERATURE:"" PRINT USING "14X."T/C Number: 35 PRINT USING "14X."Temp (C): ",4(3D.[2650 38*** 37 36 2660 **,4(3D.DD,2X)*;T(34),T(35),T(36),T(37) 2670 2680 PRINT 2690! 2700 ! CALCULATE AVERAGE INLET TEMPERATURE. 2710! 2720 Sum=0. FOR I=34 TO 37 Sum=Sum+T(I) 2730 2740 NEXTI 2750 Tin-Sum/4 PRINT USING "12X,""INNER PLATE:" PRINT USING "14X,""T/C Number: 2760 2770 2780 41 42 43 44 45 46** 2790 PRINT USING #14X, **Temp (C): "",6(3D.DD,2X)";T(40),T(41),T(42),T(43),T(44),T(45 2800 PRINT USING "14X,""T/C Number: 47 48 49 50 51 "",6(3D.DD,2X)";T(46),T(47),T(48),T(49),T(2810 PRINT USING "14X,""Temp (C): 2820 PRINT USING "14X,""T/C Number: 58"" 55 56 57 53 54 2830 PRINT USING "14X,""Temp (C): **,6(3D.DD,2X)*;T(52),T(53),T(54),T(55),T(2840 PRINT USING "14X,""T/C Number: 64"" 59 60 61 62 63 2850 PRINT USING "14X.""Temp (C): "",6(3D.DD.2X)";T(58),T(59),T(60),T(61),T(2860 PRINT USING "14X,""T/C Number: 70""" 65 66 67 68 69 · 2870 PRINT USING "14X, "Temp (C): **,6(3D.DD,2X)*;T(64),T(65),T(66),T(67),T(68),T(69) 2880 PRINT 2890! ! CALCULATE AVERAGE TEMPERATURE OF INNER HEATED WALL. 2900 2910! 2920 Sum=0. IF Itype=0 THEN FOR I=40 TO 69 2930 2940 2950 Sum=Sum+T(I) NEXT I 2960 2970 Twi=Sum/30 2980! 2990 ! CORRECTION FOR BROKEN T/C CHANNEL 47 IN INNER CURVED HEATER. 3000! 3010 ELSE 3020 FOR 1=40 TO 69 3030 IF 1=47 THEN 3050 3040 Sum=Sum+T(I) 3050 NEXT I 3060 Twi=Sum/29 3070 END IF 3080 PRINT USING "12X,""INNER INSULATION:" 3090 PRINT USING "14X,""T/C Number: 71 71 (Tinsi1) 72 (Tinsi2) 73 (Tinsi3)" 3100 PRINT USING "14X, ""Temp (C): "",3(3D.DD,8X)";T(70),T(71),T(7)) 3110 Tinsi1=T(70)

Tinsi2=T(71) 3120 3130 Tinsi3=T(72) 3140 PRINT PRINT USING "12X,""OUTLET TEMPERATURE:"" PRINT USING "14X,""T/C Number: 74 PRINT USING "14X,"Temp (C): ",4(3D. PRINT USING "0," 3150 74 75 76 77*** **,4(3D.DD,2X)*;T(73),T(74),T(75),T(76) 3160 3170 3180 31901 ! CALCULATE AVERAGE OUTLET TEMPERATURE. 3200 3210! 3220 Sum=0. FOR 1=73 TO 76 3230 Sum=Sum+T(I) NEXT I 3240 3250 3260 Tout=Sum/4 3270! ! CALCULATE DIFFERENCE BETWEEN CHANNEL INLET AND OUTLET TEMPERATURES. 3280 32901 Tdiff=Tout-Tin 3300 3310! 3320 ! CALCULATE CHANNEL BULK TEMPERATURE. 3330! Tblk=(Tin+Tout)/2 3340 3350! **! PRINT CALCULATED TEMPERATURES.** 3360 3370! PRINT USING "9X." THE FOLLOWING TEMPERATURES WERE CALCULATED: "" 3380 PRINT 3390 3400 PRINT USING "12X,""Average Outer Wall Temperature (Two) DD,"" (C)""";Two "",3D. 3410 PRINT USING "12X,""Average Inner Wall Temperature (Twi) DD,"" (C)""";Twi - "",3D. 3420 PRINT USING "12X,""Average Dutlet Temperature (Tout) DD,"" (C)""":Tout ••.3D. 3430 PRINT USING "12X,""Average Inlet Temperature (Tin) DD,"" (C)""":Tin --.3D. 3440 PRINT USING "12X,""Channel Inlet and Outlet Temp Diff (Tdiff) - "".3D. DD,"" (C)""";Tdiff 3450 PRINT USING "12X,""Average Bulk Temperature (Tblk) DD,"" (C)""";Tblk - **.3D. 3460! ! CALCULATE DELTA T (Tdel) FOR USE IN CALCULATING H LATER. 3470 3480! 3490 Tdelo=Two-Tblk IGELI-IWI-IDIK IF IPOS-O THEN Tdel-Tdelo IF IPOS-1 THEN Tdel-Tdeli PRINT USING "12X,""Mean Temperature Difference (Tdel) (C)""";Tdel 3500 3510 3520 35**30** = "",3D. 3540! 3550 ! CALCULATE DIFFERENCE BETWEEN OUTER AND INNER WALL. 3560! 3570 Twdiff=ABS(Two-Twi) 3580 PRINT USING "12X,"" DD,"" (C)""";Twdiff "Outer and Inner Wall Temp Difference (Twdiff) = "",3D. DD,"" 3590! ! CALCULATE LOCAL CHANNEL TEMPERATURES AT INLET. MIDDLE AND DUTLET. 3600 36101 3620 Tcin=Tin+(.5+Tdiff/12) Tcmid=Tin+(4.5+Tdiff/12) 3630

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3650 PRINT USING "12X,""Average Local Channel Temp at Inlet (Tcin) DD,"" (C)""";Tcin "".3D. 3660 PRINT USING "12X,""Average Local Channel Temp at Middle (Tcmid) DD,"" (C)""":Tcmid 3670 PRINT USING "12X,""Average Local Channel Temp at Outlet (Tcout) DD,"" (C)""":Tcout - **,3D. - **.3D. 3680! 3690 3700! ! CALCULATE AVERAGE LOCAL HEATED WALL TEMPERATURE AT INLET. 3710 Sum=0. Iadd=0 3720 IF Ipos=1 THEN Iadd=40 3730 3740! CORRECTION FOR BROKEN T/C CHANNELS 5, AND 47 CURVED HEATER ONLY. 3750 3760! IF Itype=1 THEN FOR I=0 TO 6 3770 3780 IF (Ipos=0 AND I=5) OR (Ipos=1 AND I=6) THEN 3810 3790 Sum=Sum+T(I+Iadd) 3800 3810 NEXT I 3820 Twin=Sum/6 3830 ELSE FOR I=0 TO 7 3840 3850 IF (Ipos=0 AND I=7) OR (Ipos=1 AND I=6) THEN 3870 3860 Sum=Sum+T(I+Iadd) 3870 NEXT I Twin=Sum/7 END IF 3880 3890 PRINT USING "12X,""Avg Local Heated Wall Temp at Inlet (Twin) (C)""";Twin - "",3D. 3900 DD, "" 3910! 3920 ! CALCULATE AVERAGE LOCAL HEATED WALL TEMPERATURE NEAR MIDDLE. 3930! 3940 Sum=0. 3950! ! CORRECTION FOR BROKEN T/C CHANNEL 16 CURVED HEATER ONLY. 3960 3970! IF Itype=1 AND Ipos=0 THEN FOR I=9 TO 15 IF I=10 THEN 4020 3980 3990 4000 4010 Sum=Sum+T(I) 4020 NEXT I 4030 Twmid=Sum/6 ELSE 4040 FOR I-9 TO 16 IF (Ipos-0 AND I-10) OR (Ipos-1 AND I-9) THEN 4080 4050 4060 4070 Sum=Sum+T(I+Iadd) NEXT I 4080 4090 Twmid=Sum/7 4110 PRINT USING "12X,""Avg Local Heated Wall Temp Near Middle (Twmid) = "",3D. DD,"" (C)""";Twmid 4120! 4100 END IF ! CALCULATE AVERAGE LOCAL HEATED WALL TEMPERATURE AT OUTLET. 4130 4140! 4150 Sum=0. 4160! ! CORRECTION FOR BROKEN T/C CHANNEL 26 CURVED HEATER ONLY. 4170 4180!

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IF Itype=1 AND Ipos=0 THEN FOR I=23 TO 29 IF I=26 THEN 4230 4190 4200 4210 4220 Sum=Sum+T(I) 4230 NEXT I Twout=Sum/6 ELSE FOR I=23 TO 29 4250 4260 Sum=Sum+T(I+Iadd) 4270 4280 4290 NEXT I Twout=Sum/7 PRINT USING "12X,""Avg Local Heated Hall Temp at Outlet (Twout) - "",3D, (C)""";Twout 4300 4310 DD, 4320 PRINT PRINT 4330 4340! ! ENTER ORIFICE PRESSURE DATA. 4350 4360! BEEP 4370 INPUT "ENTER PATM (inHg), DPM(inH20), P1M(inH20), RE1",Patm.Dpm,P1m.Re1 PRINT USING "9X,""THE FOLLOWING ORIFICE DATA WERE ENTERED:"" 4380 4390 4400 PRINT RE1*** PRINT USING "12X,""Patm(inHg) DPM(inH2O) P1H(inH2O) PRINT USING "14X,2(2D.DD,10X),(2D.DD,7X),5D.D";Patm,Dpm,P1m.Re1 DPM(inH2O) P1M(1nH20) 4410 4420 PRINT 4430 44401 ! CONVERT PRESSURE READINGS TO SI UNITS AND PIM TO ABSOLUTE. 4450 4460! 4470 Patm=Patm*3376.8 Pde1=Dpm*248.84 P1=Patm-(P1m*248.84) 4480 4490 4500 Torf=T(33) 4510! 4520 4530! ! PHYSICAL PROPERTIES AND CONSTANTS. 4540 R=286.987 !FOR AIR. Rho-P1/(R*(Torf+273.15)) !DENSITY OF AIR AT ORIF PLATE. Gamma=1.40 !FOR AIR. 4550 4560 Cp=1006 !Cp GOOD FOR TELK BETWEEN 12 AND 33 (DEG C). 4570 4580 Gc=1.0 Sigma=5.669E-8 Ecu=.12 !EMISSIVITY OF COPPER. 4590 4600 4610! ! DIMENSIONS OF CHANNEL. 4620 4630! 4640 Dc=.006350 4650 Wid=.2540 Ri=.297 4660 4670 Pwet=2*(Dc+Wid) 4680 Ac=Dc+Hid 4690 4700 Dhd=4+Ac/Pwet Ap1=.0774192 4710! ! PROPERTIES OF INSULATION. 4720 4730! 4740 4750 Xins=.0127 Kins=3.8E-2 4760! ! ENTER ORIFICE CONFIGURATION. 4770

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4780! 4790 BEEP 4800 INPUT "SELECT DIAMETER OF ORIFICE (0=.5334, 1=1.0755 (inches))", Size IF Size=0 THEN Dorf=.013548 4810 IF Size=1 THEN Dorf=.027318 4820 4830! 4840 ! DIMENSIONS OF PIPE AND ORIFICE PLATE. 4850! 4860 $A=(PI*Dorf^2)/4$ Dpipe=.051772 Apipe=PI*Dpipe*2/4 4870 4880 Beta=Dorf/Dpipe PRINT USING "12X,""Dorf(m) 4890 Beta"" 4900 A(m*2) PRINT USING *12X, (Z.6D,6X), (Z.3DE,6X), Z.4D*; Dorf, A.Beta 4910 4920 PRINT 4930 PRINT 4940! 4950 ! CORRELATION FOR EXPANSION FACTOR BASED ON 1D AND 1/2D TAPS. 4960! 4970 Y=1-(.41+.35*Beta*4)*(Pdel/(Gamma*P1)) 4980! ! CALCULATION OF FLOW COEFFICIENT (K) IN SI UNITS. 4990 5000! B=.0002+2.794E-5/Dpipe+(.0038+1.016E-5/Dpipe)*(Beta*2+(16.5+1.968504E+2*Dp 5010 ipe)*Beta 16) 5020! 5030 1 K1 AND K2 ARE DUMMY VARIABLES TO CALCULATE Ko. 5040! 5050 K1+.6014-5.3974064E-3/Dpipe*.25 5060 K2=(.376+2.8971138E-2/Dpipe*.25)*(1.6129E-7/(Dpipe*2*Beta*2+6.35E-5*Dpipe) +Beta*4+1.5*Beta*16) 5050 5060 5070 Ko=K1+K2 K=Ko+1000+B/Re1 .5 5080 5090! Mu CORRELATION GOOD FOR Torf BETWEEN 17 AND 41 DEGREES CELSIUS. 5100 5110! 5120 Mu=4.6971429E-8*Torf+1.7194722E-5 5130! 5140 ! CALCULATE Mdot AND REpipe AND COMPARE TO PREDICTED RE1. 5150! 5160 Mdot=K*A*Y*(2*Gc*Rho*Pdel)*.5 5170 Repipe=(Mdot+Dpipe)/(Mu+Apipe) Diff=(ABS(Re1-Repipe)/Re1)+100 5180 5190 Re1=Repipe IF Diff>.001 THEN 5080 PRINT USING "9X,"THE FOLLOWING DATA WERE CALCULATED:"" 5200 5210 5220 PRINT PRINT USING "12X,""Orifice Expansion Factor (Y) PRINT USING "12X,""Orifice Flow Coefficient (K) PRINT USING "12X,""Density Based on Torf (Rho) .Z.4D. .Z.4D. .Z.4D. 5230 : Y 5240 5250 3) (ka/m' :Rho 5260 PRINT USING "12X, ""Viscosity Based on Torf (Mu) = "".2Z.3DE."" (kg/m.s) ";Mu 5270! 5280 ! My CORRELATION GOOD FOR TELK BETHEEN 17 AND 41 DEGREES CELSIUS. 5290! 5300 Mu=4.6971429E-8+Tb1k+1.7194722E-5 Red=(Mdot+Dc)/(Mu+Ac) 5310 5320 Rehd=(Mdot+Dhd)/(Mu+Ac) 5330!

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! CALCULATION OF POWER INTO HEATER PLATE.
5340
5350!
5360
        Iadd=0
        IF Ipos=1 THEN Iadd-39
IF Ipos=0 THEN Rpr=2.01729
IF Ipos=1 THEN Rpr=2.00839
5370
5380
5390
5400
        Qp=(Emf(38+Iadd)*Emf(39+Iadd))/Rpr
5410!
5420
        ! CALCULATION OF HEAT INPUT TO AIR.
5430!
5440
        Dair=Mdot+Cp+(Tout-Tin)
5450!
5460
          CALCULATE LOCAL HEAT TRANSFER COEF. AND NUSSELT NUMBER.
5470!
5480
        Kair=7.7257143E-5*Tcin+.024165836
5490
        Hin=Qair/(Apl*(Twin-Tcin))
5500
        Nuin=(Hin*Dhd)/Kair
        Kair=7.7257143E-5*Tcmid+.024165836
5510
        Hmid=Qair/(Apl*(Twmid-Tcmid))
5520
        Numid=(Hmid+Dhd)/Kair
5530
        Kair=7.7257143E-5*Tcout+.024165836
Hout=Qair/(Apl*(Twout-Tcout))
5540
5550
5560
        Nuout=(Hout*Dhd)/Kair
5570!
5580
           Kair CORRELATION GOOD FOR THIK BETHEEN 17 AND 41 DEGREES CELSIUS.
5590!
5600
        Kair=7.7257143E-5*Tblk+.024165836
5610!
5620
        ! CALCULATE AVERAGE HEAT TRANSFER COEF. AND NUSSELT NUMBER.
5630!
5640
        Havg=Qair/(Apl+Tdel)
5650
        Nuavg=(Havg*Dhd)/Kair
5660!
        ! CALCULATE HEAT LOSSES.
5670
5680!
5690
        Qlo1=(Kins*Apl*(Tinso1-Tinso2))/Xins
5700
        Qlo2=(Kins+Apl+(Tinso2-Tinso3))/Xins
        Qlo=(Qlo1+Qlo2)/2
5710
        Qli1=(Kins*Apl*(Tinsi1-Tinsi2))/Xins
 5720
5730
        Qli2=(Kins*Apl*(Tinsi2-Tinsi3))/Xins
        Q11=(Q11+Q112)/2
Rr=2.0*((1-Ecu)/(Ap1*Ecu))+(1.0/Ap1)
Qr=(Sigma*(ABS((Two+273.15)*4-(Twi+273.15)*4)))/Rr
5740
 5750
 5760
        Odel=Op-Oair-Olo1-Or-Oli1
 5770
 5780!
        ! CALCULATE DEAN NUMBER.
 5790
 5800!
        De=.5*Red+((Dc/2)/Ri)*.5
PRINT USING "12X,""Viscosity Based on Tblk (Mu)
5810
5820 |
                                                                              = "".2Z.3DE,"" (kg/m.s)
 5830
        PRINT USING "12X,""Therm Cond Based on Tblk (Kair) = "",Z.3DE,"" (W/m.K)"
    ;Kaii
 5840
        PRINT
        PRINT

PRINT USING "12X, "Patm = "",6D.DD, ""(N/m<sup>2</sup>)"";Patm

PRINT USING "12X, "Pdel = ",4D.DD," (N/m<sup>2</sup>)"";Pdel

PRINT USING "12X, "P1 = ",6D.DD," (N/m<sup>2</sup>)"";Pdel

PRINT USING "12X, "Mdot = ",7.4D," (kg/s)"";Pd

PRINT USING "12X, "Repipe = ",6D.D";Repipe

PRINT USING "12X, "Red = "',6D.D";Red

PRINT USING "12X, "Rehd = ",6D.D";Rehd
 5850
 5860
 5870
 5880
 5890
5900
 5910
```

5920! ! ONLY PRINT DEAN NUMBER WHEN HEATING CURVED SECTION. 5930 59401 IF Itype=1 THEN PRINT USING "12X,""De PRINT - ",4D.D";De 5**950** 5960 PRINT 5970 PRINT USING "12X,""Qp -atts)"";Qp,Qair 5980 PRINT USING "12X,""Qlo1 -(Watts)"";Qlo1,Qli1 5990 PRINT USING "12X,""Qlo2 -(Watts)"";Qlo2,Qli2 6000 PRINT USING "12X,""Qlo -(Watts)"";Qlo,Qli 6010 PRINT USING "12X,""Qr -(Watts)"";Qr,Qdel 6020 PRINT 5960 - "",3D.3D,"" (Watts)"",7X,""Qair = "",3D.3D,"" (W ",2D.3D,"" (Watts)"",7X,""Qli1 = "",2D.3D,"" ",2D.3D,"" (Watts)",7X,""Qli2 = **,2D.3D,** "",2D.3D,"" (Watts)"",7X,""Qli = ".2D.3D." "",2D.3D,"" (Watts)"",7X,""Qdel = ".2D.3D." 6020 PRINT 6030 PRINT USING "12X, ""Hin = "", 3D.DD, "" (W/m 2C)"", 7X, ""Nuin = "", 3D.DD, ";H in-Nuin 6040 PRINT USING "12X,""Hmid = "",3D.DD,"" (W/m^2C)"",7X,""Numid = "",3D.DD,";H mid,Numid 6050 PRINT USING "12X,""Hout = "",3D.DD,"" (W/m'2C)"",7X,""Nuout = "",3D.DD,";H out,Nuout PRINT USING "12X,""Havg = "",3D.DD,"" (W/m°2C)"",7X,""Nuavg = "",3D.DD,";H 60**60** avg, Nuavg 6070 PRINT 6080 PRINT 6090 BEEP INPUT "WILL THERE BE ANDTHER RUN? (1-Y,0-N)",Go_on 6100 IF Go_on<>0 THEN PRINT USING "@,#" 6110 6120 IF Im=0 THEN WAIT 600 6130 6140 GOTO 1700 6150 END IF PRINT USING "BX.""END OF RUN""" PRINT USING "@.#" 6160 6170 PRINTER IS 1 6180 6190 ASSIGN @File TO * 6200 END 62101 6220 ! CONVERTS EMF TO DEGREES CELSIUS. 6230! 6240 SUB Tvsv(V,T) COM /Co/ D(7),Aa(76,2) 6250 5260 Sum=0 6270 FOR I=0 TO 7 6280 6290 Sum=Sum+D(I)+V*I NEXT I 6300 T=Sum 6310 SUBEND 6320! ! CALIBRATES T/C READINGS. 6330 6340! 6350 DEF FNTcorr(T,I) COM /Co/ D(7), Aa(76,2) 6360 Tc=Aa(I,0) FOR J=1 TO 2 6370 6380 6390 6400 Tc=Tc+Aa(I.J)*T'J NEXT J RETURN Tc 6410 6420 FNEND

APPENDIX B: SAMPLE CALCULATIONS

Figure 36 below, shows the major heat transfer components for each of the test sections. The sample calculations that follow, demonstrate the methods used by the computer to calculate these components as well as the Reynolds number, average heat transfer coefficient, and average Nusselt number for each set of data. Part A of Appendix B contains a sample printout of a data run, followed by a summation of the data. The sample calculations are for the concave curved section, but those for the straight section are similar.



Figure 36. Energy Balance in Test Section.

1. SAMPLE CALCULATION DATA

THE FOLLOWING CHANNEL CONFIGURATION WAS SELECTED:

Heating CURVED Test Section. Heating OUTER Plate.

Month, Day, and Time: 10:17:13:47:11

THE FOLLOWING DATA WERE RECORDED:

Heater Voltage (Vh) = 61.337 (V) Resistor Voltage (Vr) = 5.195 (V) Precision Resistor (Rpr) = 2.01729 (Ohms)						
DUTER PLATE: I/C Number: I/C Number: Temp (C): I/C Number: Iemp (C): I/C Number: Iemp (C): I/C Number: Temp (C): I/C Number: Temp (C):	1 42.75 7 46.54 13 43.40 19 45.79 25 45.30	2 41.57 8 42.61 14 44.18 20 48.45 26 46.47	3 44.39 9 43.51 15 44.10 21 49.18 27 30.31	4 41.23 10 43.47 16 41.97 22 33.83 28 45.80	5 44.55 11 44.77 17 38.06 23 45.32 29 45.98	6 40.59 12 43.86 18 45.73 24 45.51 30 47.28
DUTER INSULATIO T/C Number: Temp (C):)N: 31 (Ti 49.	nso1) 58	32 (Tins 37.29	io2) 33	(Tinso3 29.36)
DRIFICE TEMP (1	Torf) =	24.16 (()			
INLET TEMPERATU T/C Number: Temp (C):	JRE: 35 20.55	36 20.48	37 20.47	38 20.43		
INNER PLATE: T/C Number: Temp (C): T/C Number: Temp (C): T/C Number: Temp (C): T/C Number: Temp (C): T/C Number: Temp (C):	41 21.43 47 21.03 53 21.23 59 21.25 65 22.05	42 21.31 48 22.39 54 21.27 60 21.56 66 21.91	43 21.16 49 21.10 55 21.17 61 21.47 67 21.85	44 21.14 50 21.05 56 21.19 62 21.74 68 21.78	45 21.05 51 21.38 57 21.15 63 21.66 69 21.74	46 21.92 52 21.46 58 21.35 64 22.18 70 21.73
INNER INSULATI T/C Number: Temp (C):	ON: 71 (T) 21.	insi1) 53	72 (Tine 21.49	si2) 73	(Tinsi3 21.32))
OUTLET TEMPERA T/C Number: Temp (C):	TURE: 74 24.29	75 24.38	76 24.38	77 24.25		

THE FOLLOWING TEMPERATURES WERE CALCULATED:

Average Outer Wall Temperature (Two)	=	44.76	(C)
Average Inner Wall Temperature (Twi)	*	21.46	(C)
Average Outlet Temperature (Tout)	=	24.32	(C)
Average Inlet Temperature (Tin)	*	20.48	(C)
Channel Inlet and Dutlet Temp Diff (Tdiff)	*	3.84	(C)
Average Bulk Temperature (Tblk)	*	22.40	(C)
Mean Temperature Difference (Tdel)	=	22.36	(C)
Outer and Inner Wall Temp Difference (Twdiff)	*	23.30	(C)
Average Local Channel Temp at Inlet (Tcin)	*	20.64	(C)
Average Local Channel Temp at Middle (Tcmid)	*	21.92	(C)
Average Local Channel Temp at Outlet (Tcout)	¥	24.16	(C)
Avg Local Heated Wall Temp at Inlet (Twin)	*	43.51	(C)
Avg Local Heated Wall Temp Near Middle (Twmid)	×	43.50	(C)
Avg Local Heated Wall Temp at Dutlet (Twout)	×	46.06	(C)

THE FOLLOWING ORIFICE DATA WERE ENTERED:

Patm(inHg)	DPM(inH20)	P1M(inH20)	RE1
29.79	18.58	12.55	15500.0
Dorf(m)	A(m²2)	Beta	
0.027318	5.861E-04	0.5277	

THE FOLLOWING DATA WERE CALCULATED:

Drifice Expansion Factor (Y)	= 0.9852
Drifice Flow Coefficient (K)	= 0.6323
Density Based on Torf (Rho)	= 1.1424 (kg/m ³)
Viscosity Based on Torf (Mu)	= 18.330E-06 (kg/m.s)
Viscosity Based on Tblk (Mu)	= 18.247E-06 (kg/m.s)
Therm Cond Based on Tblk (Kai	(r) = 2.590E-02 (W/m.K)
Patm = 100594.87 (N/m ²) Pdel = 4623.45 (N/m ²) P1 = 97 ^{471.93} (N/m ²) Mdot = 0.0375 (kg/s) Repipe = 50352.2 Red = 8097.1 Rehd = 15799.1 De = 418.6	
Qp = 157.971 (Watts)	Qair = 144.950 (Watts)
Qlo1 = 2.847 (Watts)	Qli1 = .009 (Watts)
Qlo2 = 1.838 (Watts)	Qli2 = .038 (Watts)
Qlo = 2.342 (Watts)	Qli = .023 (Watts)
Qr = .751 (Watts)	Qdel = 9.415 (Watts)
Hin = 81.90 (W/m*2C)	Nuin = 39.39
Hmid = 86.78 (W/m*2C)	Numid = 41.58
Hout = 85.51 (W/m*2C)	Nuout = 40.70
Havg = 83.75 (W/m*2C)	Nuavg = 40.07

v _H	= 61.337 V
V _{PR}	= 5.195 V
R _{PR}	= 2.0173 A
Patm	= 29.79 in. Hg
۵P	= 18.58 in. H ₂ 0
P ₁	= 12.55 in. H ₂ 0
^T in	= 20.48 C
Tout	= 24.32 C
T _{wo}	= 44.76 C
T _{wi}	= 21.49 C
^T ins,il	= 21.53 C
T _{ins,i2}	= 21.49 C
T _{ins,01}	= 49.58 C
^T ins,02	= 37.29 C
^T orf	= 24.16
K _{air}	$= 25.90 \times 10^{-3} $ W/mC
C _{pair}	= 1006 J/Kg K
^µ air	= 4.6971 x 10^{-8} x T _{air} + 1.71947 x 10^{-5}
^K ins	$3.80 \times 10^{-2} $ W/mC
۵X _{ins}	= 0.0127 m
[€] cu	= 0.12
σ	= 5.669 x 10^{-8} W/m ² K ⁴
ß	= .5277
Ŷ	= 1.40
a ^c	= 1 Kg m/Ns ²

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R	= 286.98 Nm/K
F _{wo-wi}	= 1.0
A _{PL}	$= .07742 \text{ m}^2$
A _c	$= .0016 m^2$
A pipe	$= .00211 m^2$
D _c	= .00635 m
D pipe	= .0518 m
D _{orf}	= .027318 m
Pwet	= .5207 m
Ri	= .297 m

2. TEMPERATURE CALCULATIONS

a. Bulk Temperature (T_{blk})

$$T_{blk} = \frac{T_{in} + T_{out}}{2} = \frac{20.48 + 24.32}{2} = 22.40 C$$

К

b. Mean Temperature Difference (ΔT)

 $\Delta T = T_{wo} - T_{blk} = 44.76 - 22.40 = 22.36$

c. Temperature Difference Between Outer and Inner Wall (T_{wdiff})

 $T_{wdiff} = T_{wo} - T_{wi} = 44.76 - 21.49 = 23.27 C$



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MICROCOPY RESOLUTION TEST CHART

3. POWER CALCULATIONS

a. Power Supplied (Q_p)

$$Q_{p} = \frac{V_{H} V_{PR}}{R_{PR}} = \frac{(61.337)(5.195)}{(2.0173)} = 157.97 W$$

b. Heat Lost Through the Back of the Outer Plate (Q_{lo})

$$Q_{10} = \frac{T_{\text{ins,ol}} - T_{\text{ins,o2}}}{(\Delta X_{\text{ins}}) / (K_{\text{ins}}) (A_{\text{PL}})} = \frac{(49.58 - 37.29)}{(0.0127) / (3.80 \times 10^{-2}) (.07742)}$$

= 2.85 W

c. Heat Lost Through the Back of the Inner Plate (\textbf{Q}_i)

$$Q_{1i} = \frac{T_{ins,i1} - T_{ins,i2}}{(\Delta X_{ins})/(K_{ins})(A_{PL})} = \frac{(21.53 - 21.49)}{(0.0127)/(3.80 \times 10^{-2})(.07742)}$$

= .01 W

d. Heat Radiated (Q_r)

(1) Radiation Resistance (R_R)

$$R_{R} = 2 \times \left(\frac{1 - \epsilon_{Cu}}{A_{PL} \epsilon_{Cu}}\right) + \frac{1}{A_{PL} F_{wo-wi}}$$
$$= \frac{1}{A_{PL}} \left[\frac{2}{\epsilon_{Cu}} - 1\right] = \frac{1}{0.07742} \left[\frac{2}{.12} - 1\right] = 202.36 \text{ m}^{-2}$$

(2) Heat Radiated (Q_r)

$$Q_r = \frac{\sigma (T_{wo}^4 - T_{wi}^4)}{R_R} \qquad T [=] K$$

$$=\frac{(5.669 \times 10^{-8}) [(44.76 + 273)^4 - (21.49 + 273)^4]}{202.36}$$

= 0.75 W

4. MASS FLOW RATE CALCULATIONS

a. Pressure Conversions

(1)
$$P_{atm} = 29.79$$
 in Hg x 3376.8 = 100594.87 N/m²
(2) $\Delta P = 18.58$ in H₂O x 248.84 = 4623.5 N/m²
(3) $P_1 = (12.55$ in H₂O x 248.84) - 100594.87
= 97471.93 N/m²

b. Density of Air
$$(\rho_{air})$$

$$\rho_{air} = \frac{P_1}{R T_{orf}} = \frac{(97471.93)}{(286.98)(24.16 + 273)} = 1.1424 \text{ Kg/m}^3$$

c. Expansion Factor (Y)

$$Y = 1 - [.41 + .35\beta^{4}] \frac{\Delta P}{\gamma P_{1}}$$

= 1 - [.41 + .35(.5277)^{4}] $\frac{4623.5}{(1.40)(97471.93)}$
= .9852

d. Area of Orifice (A)

$$A = \frac{\pi (D_{orf})^2}{4} = \frac{\pi (.027318)^2}{4} = .005861 \text{ m}^2$$

e. Mass Flow Rate (m)

$$m = YKA 2g_{c} \rho_{air} \Delta P = (.9852)K(.0005861)$$

x 2(1)(1.1424)(4623.5)

m = .05935 K

Iterating:

Assume a Pipe Reynolds number = 50000 Obtain a value for K, the flow coefficient, from reference 30. K = .6325 Solve for \dot{m} . \dot{m} = .0375 Kg/s Solve for new Repipe = $\frac{\dot{m} D_{pipe}}{A_{pipe}\mu_{air}} = \frac{(.0375)(.0518)}{(.00211)(18.33\times10^{-6})}$ Repipe = 550225 Check convergence and repeat process if necessary. (Convergence if difference less than .001)

m = .0375 Kg/s

5. REYNOLDS NUMBER CALCULATIONS

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a.
$$\operatorname{Re}_{\operatorname{pipe}} = \frac{\stackrel{\circ}{\operatorname{m}} \stackrel{D}{\operatorname{pipe}}}{\operatorname{A}_{\operatorname{pipe}} \stackrel{\mu}{\operatorname{air}}} = \frac{(.0375)(.0518)}{(.00211)(18.33\times10^{-6})} = 50225$$

b.
$$\operatorname{Re}_{d} = \frac{\prod_{c}^{m} D_{c}}{A_{c} \mu_{air}} = \frac{(.0375)(.00635)}{(.0016)(18.25 \times 10^{-6})} = 8097.1$$

$$Re_{hd} = \frac{\overset{\bullet}{m} \overset{D}{D}_{hd}}{A_c \overset{\mu}{air}}$$
(1) $D_{hd} = \frac{4 \times A_c}{P_{wet}} = \frac{(4)(.0016)}{.5207} = .01229 \text{ m}$

$$Re_{hd} = \frac{\overset{\bullet}{m} \overset{D}{D}_{hd}}{A_c \overset{\mu}{air}} = \frac{(.0375)(.01229)}{(.0016)(18.25\times10^{-6})} = 15799.1$$

5. HEAT CONVECTED TO AIR CALCULATION
a.
$$Q_{air} = \frac{m}{m}C_{p}(T_{out} - T_{in}) = (.0375)(1006)(24.32 - 20.48) = 144.95 W$$

7. AVERAGE HEAT TRANSFER COEFFICIENT CALCULATION

$$\overline{h} = \frac{Q_{air}}{A_{PL} \Delta T} = \frac{144.95}{(.07742)(22.36)} = 83.72 \text{ W/m}^2\text{C}$$

8. AVERAGE NUSSELT NUMBER CALCULATION

$$\overline{Nu}_{hd} = \frac{\overline{h} \ D_{hd}}{K_{air}} = \frac{(83.73)(.01229)}{25.90 \ x \ 10^{-3}} = 39.73$$

9. DEAN NUMBER CALCULATION

$$De = \frac{Re_d}{2} \sqrt{\frac{Dc/2}{Ri}} = \frac{8097.1}{2} \sqrt{\frac{.00635/2}{.297}} = 418.6$$

APPENDIX C: EXPERIMENTAL UNCERTAINTY

The uncertainties for the major variable in the experiments, were calculated in accordance with the method described by S. Kline and F. McClintoch [Ref. 33]. The estimates of the uncertainty in the measured quantities were made conservatively. As a result, there is considerable confidence in the uncertainties as calculated.

The following equations were used to calculate the uncertainties:

$$(1) \quad \frac{d\rho}{\rho} = \left(\left(\frac{dR}{R}\right)^{2} + \left(\frac{dP_{1}}{P_{1}}\right)^{2} + \left(\frac{dT_{orf}}{T_{orf}}\right)^{2}\right)^{1/2}$$

$$(2) \quad \frac{d\tilde{m}}{\tilde{m}} = \left[\left(\left(\frac{dY}{Y}\right)^{2} + \left(\frac{dK}{K}\right)^{2} + \left(\frac{dA}{A}\right)^{2}\right) + \left(\frac{1/4}{\left(\left(\frac{d\rho_{air}}{\rho_{air}}\right)^{2} + \left(\frac{dA}{A}\right)^{2}\right)\right)^{1/2}$$

$$(3) \quad \frac{dQ_{air}}{Q_{air}} = \left(\left(\frac{d\tilde{m}}{\tilde{m}}\right)^{2} + \left(\frac{dC_{pair}}{C_{pair}}\right)^{2} + \left(\frac{d(T_{out}-T_{in})}{T_{out}-T_{in}}\right)^{2}\right)^{1/2}$$

$$(4) \quad \frac{d\tilde{m}}{\tilde{m}} = \left(\left(\frac{dQ_{air}}{Q_{air}}\right)^{2} + \left(\frac{dA_{PL}}{A_{PL}}\right)^{2} + \left(\frac{dAT}{AT}\right)^{2}\right)^{1/2}$$

(5)
$$\frac{dNu_{hd}}{Nu_{hd}} = \left(\left(\frac{dh}{h}\right)^2 + \left(\frac{dD_{hd}}{D_{hd}}\right)^2 + \left(\frac{dK_{air}}{K_{air}}\right)^2\right)$$

(6)
$$\frac{dRe_{hd}}{Re_{hd}} = \left(\left(\frac{dm}{m}\right)^2 + \left(\frac{dD_{hd}}{D_{hd}}\right)^2 + \left(\frac{d\mu_{air}}{\mu_{air}}\right)^2 + \left(\frac{dA_c}{A_c}\right)^2\right)$$

The following quantities had their uncertainty calculated by dividing the estimated error in the quantity by the value of the quantity, and these uncertainties are considered constant for the range of values in this experiment:

Quantity	Uncertainty
A	.0093
^A c	.0408
A _{PL}	.0080
C _{pair}	.0020
Dc	.0400
D _{hd}	.0396
К	.0030
K _{air}	.0038
Patm	.0010
R	.0003
W _{id}	.0080
Y	.0020
^µ air	.0054

The following quantities had the worst case error for hydraulic Reynolds numbers below 17000:

Quantity	Error
T _{blk}	<u>+</u> .1°C
T _{in}	<u>+</u> .1°C
T _{orf}	<u>+</u> .1°C
Tout	<u>+</u> .1°C
T _{wo}	<u>+</u> .1°C
T _{out} -T _{in}	<u>+</u> .2°c
ΔΤ	<u>+</u> .2 ⁰ C

For hydraulic Reynolds numbers greater than 17000 the error in the outlet temperature was much higher, as mentioned previously, and the following errors apply:

Quantity	Error
^T blk	<u>+</u> .2°c
^T in	<u>+</u> .1°C
^T orf	<u>+</u> .1°C
Tout	<u>+</u> .3°C
T _{wo}	<u>+</u> .1°C
T _{out} -T _{in}	<u>+</u> .4 [°] C
ΔΤ	+.3 [°] C

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The uncertainties for the following quantities were calculated using the sample data for the outer curved test section at a hydraulic Reynolds number of 15800:
Quantity	Uncertainty
P ₁	.0159
ΔP	.0108
^T orf	.0003
T _{out} -T _{in}	.0521
ΔΤ	.0089
^p air	.0159
• m	.0139
Q _{air}	.0540
h	.0553
Nu _{hd}	.0681
Rehd	.0588

The uncertainty for the following quantities were calculated for the outer curved test section at a hydraulic Reynolds number of 21200:

Quantity	Uncertainty
P ₁	.0043
ΔP	.0053
Torf	.0003
^T out ^{-T} in	.1039
ΔΤ	.0111
^p air	.0043
• M	.0105
Q _{air}	.1044

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h	.1053
Nuhd	.1126
Rehd	.0581

By multiplying the quantity by its uncertainty the possible error can be calculated. The previous uncertainty calculations for a hydraulic Reynolds number of 15800 yielded $\overline{\text{Nu}}_{\text{hd}}$ = 40.07 ±2.73 and a Re_{hd} = 15800 ± 928. For the run at a hydraulic Reynolds number of 21200 the $\overline{\text{Nu}}_{\text{hd}}$ = 44.71 + 5.03 and a Re_{hd} = 21200 ± 1232.

The major source of uncertainty in the average Nusselt number is the uncertainty in the difference between the outlet and inlet air temperature. If this uncertainty can be brought below four percent, then the uncertainty of the average Nusselt number will be dictated by the uncertainty of the measured dimensions of the channel. In particular, the channel height (D_c) which has an uncertainty $dD_c/D_c = .0400$. The uncertainty of the channel height (an only be reduced by increasing the channel height, and/or using materials with better tolerances.

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