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Photoluminescent Response of Palladium-Cadmium Sulfide and Palladium-Graded Cadmium Sulfoselenide Schottky Diodes to Molecular Hydrogen

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INTRODUCTION

The bulk photoluminescence (PL) of semiconductors has proven to be a useful probe of Schottky barrier characteristics in both semiconductor/metal diodes¹ and photoelectrochemical cells (PEC's).² In particular, the electric field in the semiconductor can be estimated from PL intensity using a dead-layer model: Electron-hole (e⁻-h⁺) pairs formed within a distance from the interface on the order of the depletion width do not contribute to PL. While the importance of surface properties on bulk PL has been described^{3,4}, their interrelationship has not been examined experimentally in detail.

In choosing a system for such studies, our interest was drawn to Schottky diodes constructed with Pd because of the known sensitivity of their current-voltage (i-V) properties to gaseous H_2 .⁵⁻⁸ In general, Schottky barriers resulting from junctions of Pd with a variety of n-type semiconductors decreased upon exposure of the diode to H_2 , reflecting diminution of the Pd work function by the gas.⁹ The effect is sufficiently sensitive to H_2 concentration that Pd-CdS⁵ and Pd-TiO₂⁶ Schottky diodes have been proposed as H_2 detectors. Variation in Schottky barrier height with H_2 appears to also strongly influence the efficiency with which metal-coated semiconductor electrodes evolve H_2 in PEC's.¹⁰

We report herein that PL intensity from Pd-CdS Schottky diodes is substantially perturbed by exposure of the diode to H₂. Moreover, the PL properties are consistent with a dead-layer model, allowing calculation of the variation in electric field thickness with exposure to H₂. Schottky diodes constructed from Pd and a graded CdS_xSe_{1-x} ($0 \le x \le 1$; x = 1 at surface) semiconductor^{11,12} (Pd-CdS_xSe_{1-x}) evince a change in spectral distribution upon exposure to H₂; the color-coded nature of the light emitted from this semiconductor can be used to map the effective electric field in this solid and its perturbation by H₂. Besides illustrating the coupling of surface interactions to bulk PL, these studies demonstrate the feasibility of constructing optically-coupled chemical sensors.

EXPERIMENTAL SECTION

Samples of n-type, single-crystal, CdS and CdSe <u>c</u>-plates (1-mm thickness; ~ 2 ohm-cm resistivity) were purchased from Cleveland Crystals, Inc. and cut into pieces of ~ 0.25 -cm² area. Graded samples of n-CdS_xSe_{1-x} ($0 \leq x \leq 1$) where the graded zone had a thickness of $\sim 1.0 \mu$ m were prepared by vapor-phase diffusion of S into CdSe and characterized as described previously.¹² Prior to Pd deposition on its 0001 Cd-rich face, CdS was etched in Br₂/MeOH (1:10 v/v); the graded CdS_xSe_{1-x} samples were not etched owing to the thinness of the graded layer.

Deposition utilized Pd foil (50x50x0.1 mm; >99.997% metallic purity; Aesar Co., Seabrook, NH) and a SPI Super Sputter apparatus. Sputtering was conducted at $2x10^{-4}$ torr Ar pressure and $60-\mu A$ beam current for ~ 45 s; a parallel deposition onto Pt foil was used in conjunction with electrochemical stripping (0.75 V vs. SCE in 1M HCl aqueous electrolyte) to estimate the Pd layer thickness to be $\sim 100 \text{ Å}$. After deposition, electrical contact was made to the translucent Pd layer with Ag epoxy and to the back surface with Ga/In eutectic and Ag epoxy; Cu wires were connected to the Ag epoxy and current-voltage properties obtained with a PAR Model 173 potentiostat and Model 175 programmer.

PL spectra were recorded using 457.9- and 488.0-nm excitation from a CR-12 Ar⁺ laser and an Aminco-Bowman spectrometer equipped with a Hamamatsu R446S PMT. The sample was enclosed in a cell which permitted dry N₂, a 3:1 N₂:H₂ mixture (Air Products tank of H₂ mixed with in-house N₂), or air to bathe the sample; flow rates of ~ 0.5 l/min were employed. Optical properties of the Pd film were examined by depositing the metal on a microscope slide; transmission and reflectivity of 457.9-, 488.0-, and 514.5-nm light were measured in air and the N₂/H₂ medium.

RESULTS AND DISCUSSION

The PL spectrum from a Pd-CdS sample in air, illuminated through the metal with 457.9-nm ultraband gap (Eg \sim 2.4 eV) light, is characterized by band edge emission at $\sqrt{510}$ nm,¹³ as shown in Fig. 1. When a 3:1 mixture of N₂:H₂ is passed over the sample, Figure 1 reveals that the PL intensity is enhanced by approximately 70%; the enhancement occurs over $\sqrt{30}$ s with the $\sqrt{100-A}$ -thick Pd layer. PL enhancement requires the presence of Pd but does not appear optical transmission to derive from changes in the metal's optical properties: and reflectivity of the metal film, deposited on a microscope slide, were insensitive to the presence of H_2 . After flushing the sample cell with N_2 and then with air, the PL intensity returns to its original value. This effect is reversible over at least 10 cycles. If more penetrating 488.0-nm light is used for excitation, the PL intensity is augmented by v40% with exposure From current-voltage data, the Pd-CdS structure exhibits typical diode to H₂. behavior in air and in the N_2/H_2 medium.

These spectral changes are consistent with a reduction in Schottky barrier height resulting from the dissolution of H₂ in Pd, a well-studied phenomenon.⁵⁻¹⁰ Qualitatively, the PL intensity is expected to rise because the smaller electric field in the semiconductor allows a larger fraction of e^-h^+ pairs to radiatively recombine. By regarding the region supporting the electric field as being completely nonemissive, i.e., a dead layer, a quantitative expression for relative PL intensity can be obtained, eq. (1); this treatment assumes that the CdS surface recombination velocity in the diode is either very large or insensitive to H₂.²,³ In eq. (1), ϕ_{H_2} and ϕ_{air} are the radiative quantum yields in H₂ and

$$\frac{\phi_{\text{air}}}{\phi_{\text{H}_2}} = \exp(-\alpha'\Delta D) \tag{1}$$

in air; ΔD is the difference in dead-layer thickness between the two media; and $\alpha' = (\alpha + \beta)$ with α and β the solid's absorptivities for the exciting and emitted light, respectively. For CdS, α for $E \perp c$ polarized light is 6×10^4 and $\nu(9-10) \times 10^4$ cm⁻¹ for 488.0- and 457.9-nm light, respectively; β is 7×10^3 cm⁻¹ at 510 nm.¹⁴

The PL enhancements seen for the two excitation wavelengths employed give a consistent value for ΔD of $\sim 500-600$ Å. This value for the contraction of the electric field upon exposure to H₂ can be used to calculate a reduction in Schottky barrier height of the initial height is known. Literature estimates for the Pd-CdS barrier height in air range from $\sim 0.5-0.8 \cdot eV_3$; 5,7 our i-V curves yield an estimated height of 0.6 eV. The depletion width W is related to barrier height qV by eq. (2), where ε_0 is the permittivity of free space, q is the electronic

$$W = \sqrt{\frac{2\varepsilon\varepsilon_0 V}{q N_D}}$$
(2)

charge, and ε and N_D are the semiconductor's dielectric constant and charge carrier density; for our samples, ε and N_D are \sim 10 and 9×10^{15} cm⁻³, respectively.¹⁵ Substitution into eq. (2), equating D with W, leads to an estimated decline in barrier height of 0.2 eV. Literature values vary from \sim 0.5 eV⁵ to \sim 0.2 eV (extrapolated).⁷ The range of values likely reflects variations in sample preparation.¹⁶

A complementary system with H₂-sensitive PL is the Pd-CdS_xSe_{1-x} Schottky diode. The PL spectrum in air is shown in Fig. 2 and consists of edge emission from all of the CdS_xSe_{1-x} compositions which comprise the ν 1.0-µm-thick graded zone, from CdS at the surface to the CdSe (E_g ν 1.7 eV; $\lambda_{max} \nu$ 720 nm) substrate.¹² A linear correlation exists between composition x and the emission band maximum, eq. (3).¹³ In conjunction with AES/Ar⁺ sputter etch data, eq. (3) provides

$$\lambda_{\max}(nm) = 718 - 210 \times$$
 (3)

a map of radiative recombination in the solid: The PL is color-coded to indicate the depth from the surface at which e^--h^+ pair recombination occurs. Perturbation of PL by applied potential reflects changes in the effective electric field

11,12 (EEF) in the solid ; we use the term "EEF" to reflect the fact that the electric field of this solid is complex and contains contributions, e.g., from band-edge gradients, in addition to the field arising from the Schottky barrier.

Exposure of a Pd-CdS_xSe_{1-x} diode to H₂ results in asymmetric enhancement of as much as 50% in the blue end of the PL spectrum, Fig. 2; the effect corresponds to a modest color change. The material has diode i-V characteristics in both environments. Qualitatively, the PL enhancement at shorter wavelengths indicates a reduction in the EEF in the near-surface region of the semiconductor, since it is the near-surface, S-rich compositions which give rise to the emission. A more quantitative estimate of the affected region is afforded by the cessation of spectral changes for $\lambda \geq 600$ nm. From eq. (3) and AES/Ar⁺ sputter etch data, exposure to H₂ influences the EEF to a depth of $\sim 0.1 \ \mu m$ (1000 Å) from the surface.

In summary, bulk PL from Pd-CdS and Pd-CdS_xSe_{1-x} Schottky diodes provides a sensitive probe of changes in the electric field of the semiconductors resulting from a surface interaction with H₂. The ability to transform molecular surface interactions into a change in bulk PL intensity (Pd-CdS) or color (Pd-CdS_xSe_{1-x}) could have applications to the design of optically-coupled chemical sensors. Further experiments designed to couple analyte sensitivity to PL are underway in our laboratories.

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Figure Captions

Figure 1. Uncorrected PL spectra of a Pd-CdS Schottky diode in air (solid line) and in a $3:1, N_2: H_2$ atmosphere (dashed line). The sample was excited with the same intensity ($\sim40 \text{ mW/cm}^2$) of 457.9-nm light in both experiments, using an identical sample-detection optics geometry.

Figure 2. Uncorrected PL spectra of a Pd-CdS_xSe_{1-x} Schottky diode in air (solid line) and in a 3:1, N₂:H₂ atmosphere (dashed line). The sample was excited with the same intensity (\sim 40 mW/cm²) of 457.9-nm light in both experiments, using an identical sample-detection optics geometry.





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