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## THESIS

A COMPUTER PROGRAM TO CALCULATE  
THE SUPERSONIC FLOW OVER A  
SOLID CONE IN AIR OR WATER

by

Patrick William Hughes  
June 1984

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Prepared for: Mr. Donald Phillips  
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Naval Surface Weapons Center, White Oak  
Silver Spring, Maryland 20910

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Generally, pressures involved in water flow are much larger than for air flow, and the cone semi-vertex angles for water flow are smaller than for air flow.



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A Computer Program to Calculate  
the Supersonic Flow Over a  
Solid Cone in Air or Water

by

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Lieutenant, United States Navy  
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Submitted in partial fulfillment of the  
requirements for the degree of

MASTER OF SCIENCE IN ENGINEERING SCIENCE

from the

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## ABSTRACT

The computer program calculates the supersonic flow over a cone in air or water. The main objective is to calculate the cone semi-vertex angle given prescribed initial conditions. The program is written in structured FORTRAN and implements Busemann's graphical integration technique. Supersonic flow over a cone in water is useful as a good first approximation to the motion of the metal jet from an explosive shaped-charge fired underwater.

A typical result for supersonic flow over a cone in water is as follows: given an upstream temperature, 323.16 Kelvin; upstream pressure, 1 bar; shock angle, 20.0 degrees; and pressure behind the shock front, 5 kilobars, the cone semi-vertex angle is calculated to be 7.23 degrees.

Generally, pressures involved in water flow are much larger than for air flow, and the cone semi-vertex angles for water flow are smaller than for air flow.

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## I. INTRODUCTION

### A. DESCRIPTION OF THE PROBLEM

The solution of the hydrodynamic equations describing supersonic flow over a cone in air has been well known since the 1930's. Until recently, the problem of describing the flow over a cone in water has been limited to solutions of the subsonic case. Primarily, calculations were limited to subsonic flow because researchers believed that supersonic flow in water was not feasible for normal vessels (such as a ship). For ordinary vessels in water, it is certainly true that supersonic flow past that vessel is highly impractical. However, the motion of the metal jet from an explosive shaped-charge fired underwater is supersonic.

This thesis presents a computer program which calculates the hydrodynamic flow past a cone in either water or air under supersonic conditions. The program utilizes the methods developed by previous researchers for calculating the supersonic flow in air and which have been suitably modified to describe the conditions in the water. Such modifications include utilizing the modified Tait equation, which is the "thermal" or "thermodynamic" equation of state for water, to describe the physical state of the water rather than the perfect gas law used for air.

In actuality, the cone liner in the jet from an explosive shaped-charge is blunt-nosed rather than conical. However, solution of the conical case is a preliminary requirement to solution of the actual blunt-nosed problem. The solution to the conical flow case will serve as an excellent test program for the solution to the blunt-nosed problem. This thesis presents a solution to the conical

case and it is hoped that the program will assist the continuing research into the problems of utilizing explosive shaped-charges in an underwater environment.

#### E. METEORIC TECHNOLOGY

The computer program presented in this work was originally developed in the BASIC computer language using a Hewlett-Packard HP-67 computer. That program is the basis for this thesis. It was desired to translate the program into a higher-order computer language for execution on a large, mainframe computer system. This translation was desired in order to make the program more readily accessible to a wider body of researchers and in order to speed the execution time of the program. In this thesis, the following goals have been accomplished:

- Successfully translate the program from BASIC into a higher-order language. This goal was met by utilizing FORTRAN as the high-level language of choice. While FORTRAN has many drawbacks as a high-order language, it is still widely used in the scientific community. FORTRAN was used, therefore, so that the program will be useful to as wide an area of researchers as possible.
- Follow modern programming practices in the design and implementation of the program. As before, the choice of FORTRAN as the high-level language makes this goal somewhat more difficult. However, many computer scientists have demonstrated that structured programming practices can be achieved using FORTRAN. To the largest extent possible structured programming practices have been utilized.
- Present a "user-friendly", well-documented program. In this regard, liberal use of comments occur in the

program itself, meaningful variable names are used, and detailed flowcharts which demonstrate the logic of the program are included. In addition, due to limited interaction with the user, the user's responses are verified before the program executes.

Similarly, in contrast to air, the thermodynamic changes which occur as a result of the shock process in water cannot be easily delineated by simple equations as in the air case. However, a simplification can be made in the water case because, unlike in the air case, the pressure jump across the shock in water is so very large. In air, pressure changes across the shock front on the order of 1 to 2 bars are considered large (at least for chemical explosives). In contrast, as pointed out by Richardson, et. al., [Ref. 3], the pressure jump across a shock in water is on the order of kilobars to tens of kilobars. Therefore, the calculations can be simplified by specifying the pressure on the downstream side of the shock front. This is valid since the upstream pressure is so small in comparison to the upstream dynamic pressure  $\rho_1 v_1^2 / 2$  and in comparison to both the downstream pressure and dynamic pressure. The specification of the downstream pressure is accomplished, in the program of this thesis, by allowing the user to input a "pressure multiplication factor (MFACT)", which converts the pressure upstream to a pressure downstream at point 2 which is given by:

$$P_2 = P_1 \times \text{MFACT} \times 1000.0 \quad (2.24)$$

The factor 1000.0 in equation 2.24 converts the right-hand side of the equation from a pressure in bars to a pressure in kilobars. As an example, if  $P_1$  is 1 bar and MFACT is 5.0, the pressure at point 2 downstream will be 5 kilobars.

It can be shown that the simplifying assumption made above is entirely valid by considering the momentum equation for steady frictionless flow along a streamline. The

## I. EQUATIONS SPECIFIC TO WATER

In contrast to air, which has a relatively elegant and simple state equation, the equation of state for water is rather more complicated. The most commonly used state equation for water is known as the modified Tait equation, which may be written as:

$$p = B(S) \left[ \left( \frac{\rho}{\bar{\rho}} \right)^n - 1 \right] \quad (2.21)$$

where  $B(S)$  is a slowly varying function of entropy alone,  $n$  is approximately a constant equal to 7.15, and  $\bar{\rho}$  is the value of the density for zero pressure. The above equation is from Eolt [Ref. 6], but, in different forms, it is also described quite extensively in Cole [Ref. 2], Richardson, et. al., [Ref. 3] and Rowlinson [Ref. 7]. As mentioned in the Introduction to this work, the modified Tait equation is the "thermal" equation of state for water. Cole [Ref. 2] shows that the modified Tait equation is of the form

$$p = B(S) \left[ \left( \frac{v(T,0)}{v(T,p)} \right)^n - 1 \right] \quad (2.22)$$

or, in simpler terms,

$$f = p(v, T) \quad (2.23)$$

Since the full modified Tait equation relates the three thermodynamic quantities of  $p$ ,  $v$ , and  $T$ , it is called the "thermodynamic" or "thermal" equation of state.

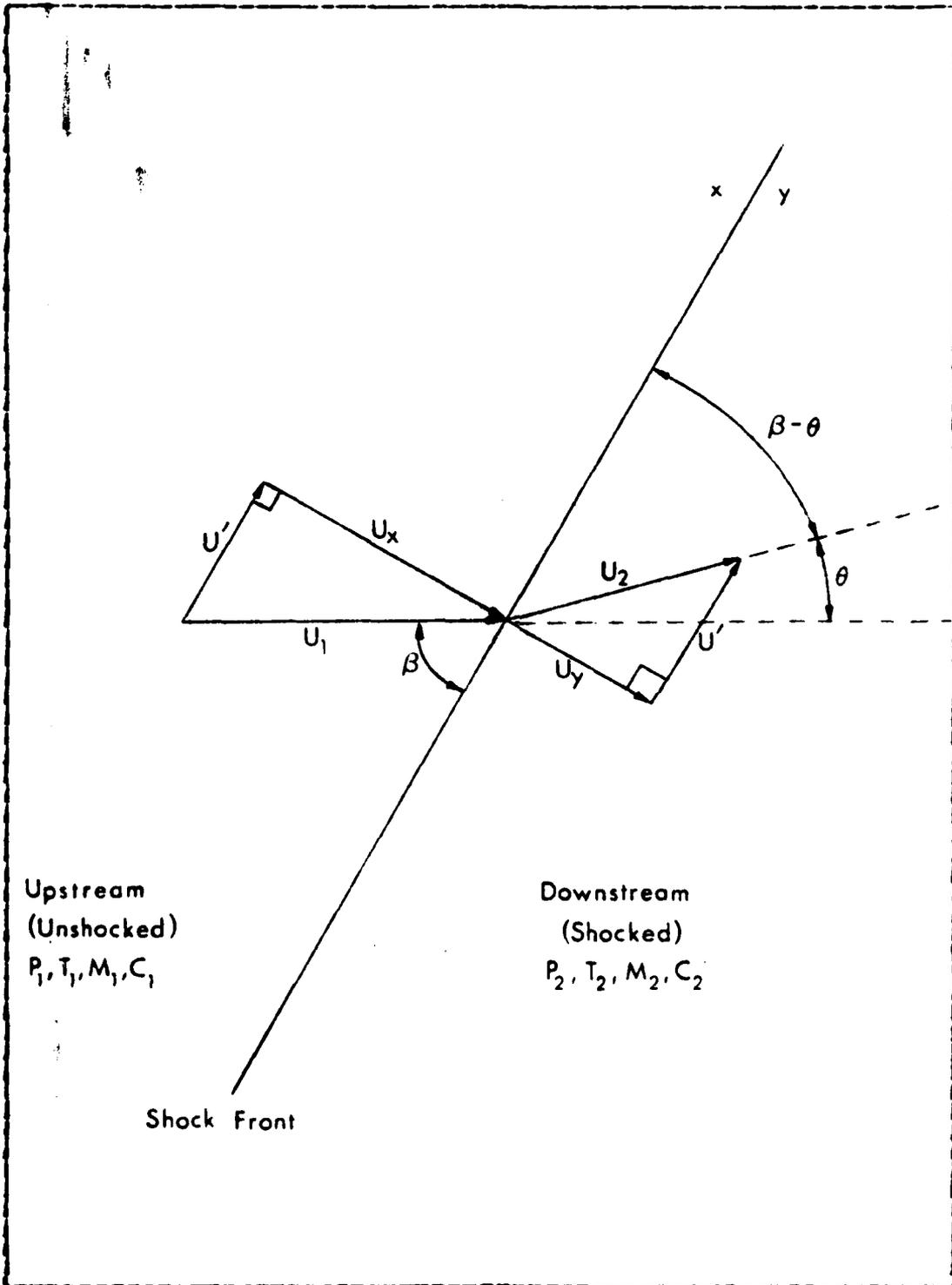


Figure 2.4 Geometry of the Oblique Shock Front.

$$\frac{p_2}{p_1} = \frac{kM_1^2 \sin^2 \beta - \left(\frac{k-1}{2}\right)}{\left(\frac{k+1}{2}\right)} \quad (2.17)$$

Equation 2.17 is used to determine the pressure downstream of the shock front.

$$\frac{T_2}{T_1} = \left(\frac{c_2}{c_1}\right)^2 = \frac{1 + \left(\frac{k-1}{2}\right)(M_1^2 \sin^2 \beta)(kM_1^2 \sin^2 \beta - \left[\frac{k-1}{2}\right])}{\left(\frac{k+1}{2}\right)^2 M_1^2 \sin^2 \beta} \quad (2.18)$$

Equation 2.18 is used to determine the temperature, and more importantly, the speed of sound,  $c$ , downstream of the shock front.

$$M_x = \frac{u_x}{c_1} = \frac{u_1 \sin \beta}{c_1} = M_1 \sin \beta \quad (2.19)$$

$$M_y = \frac{u_y}{c_2} = \frac{u_2 \sin(\beta - \theta)}{c_2} = M_2 \sin(\beta - \theta) \quad (2.20)$$

Finally, equations 2.19 and 2.20 are used to determine the velocity components of the flow across the shock front.

The geometry of the flow conditions across the shock front is illustrated by figure 2.4. Note that in the equations above and in figure 2.4, the '1' subscripts refer to conditions in the unshocked (upstream) fluid, the 'x' subscripts refer to the normal components of flow in the unshocked fluid, the '2' subscripts refer to conditions in the shocked (downstream) fluid and the 'y' subscripts refer to the normal components of flow in the shocked fluid. Note also that 'k' in the equations above is the designator for the ratio of the heat capacities  $c_p/c_v$ .

where  $R$  is the specific gas constant and is related to the universal gas constant,  $A$ , by:

$$R = \frac{A}{M_a} \quad (2.14)$$

In equation 2.14,  $M_a$  is the molecular weight of the air.

In air, the change in the thermodynamic properties of the gas as it crosses the shock front are easily calculated, as shown in Kinney and Graham [Ref. 5]. Since most fluid dynamics textbooks illustrate the development of the equations which follow, it is not necessary to derive them here. As mentioned previously, Kinney and Graham [Ref. 5] provide exceptionally lucid explanations and derivations. The principle equations used to calculate the thermodynamic changes which occur across the shock front in air are as follows:

$$\frac{\tan(\beta - \theta)}{\tan\beta} = \frac{2 + (k - 1) M_1^2 \sin^2\beta}{(k + 1) M_1^2 \sin^2\beta} \quad (2.15)$$

Equation 2.15 is used to iteratively determine the deflection angle  $\theta$ . All other quantities in this equation are known (i.e.  $\beta$  is the shock angle and  $M_1$  is the freestream Mach number, both of which are input parameters to the program for the air calculations).

$$[M_2 \sin(\beta - \theta)]^2 = \frac{2 + (k - 1) M_1^2 \sin^2\beta}{2kM_1^2 \sin^2\beta - (k - 1)} \quad (2.16)$$

Having determined  $\theta$  from equation 2.15, equation 2.16 is used to determine the Mach number on the downstream side of the shock front.

By combining Eqn 2.9 with Eqn 2.8, one arrives at:

$$R = \frac{\frac{\sin\theta}{V \sin\omega}}{1 - \frac{2V^2 \sin^2(\omega - \theta)}{(k-1)(V_{\max}^2 - V^2)}} \quad (2.10)$$

But, the energy equation asserts that:

$$c^2 = \left(\frac{k-1}{2}\right) (V_{\max}^2 - V^2) \quad (2.11)$$

therefore

$$R = \frac{\frac{\sin\theta}{V \sin\omega}}{1 - \frac{V^2 \sin^2(\omega - \theta)}{c^2}} \quad (2.12)$$

where  $c$  is the local speed of sound in the fluid.

Equation 2.12 is the basis for the calculation of the supersonic flow over the solid cone in either air or water. The graphical integration method invented by Busemann is adequately explained in Shapiro [Ref. 8] and need not be repeated here. Essentially, the computer program given in this work automates the Busemann graphical integration method for calculating the cone semi-vertex angle.

### C. EQUATIONS SPECIFIC TO AIR

The equation of state for air is specified by the perfect gas law (under the assumption, that is, that the air behaves as a perfect gas). This law is quite elegant and simple and allows easy manipulation to obtain various quantities. The form of the perfect gas law most often used in the calculations of this thesis is:

$$p = \rho RT \quad (2.13)$$

But, from Eqn 2.1 and Eqn 2.2, it can be seen that:

$$\frac{dV_r}{d\omega} = V_\omega = -V \sin(\omega - \theta) \quad (2.5)$$

Substitution of this result into Eqn 2.3 yields:

$$\frac{d\theta}{d\omega} = \frac{-\frac{dV}{d\omega}}{V \tan(\omega - \theta)} \quad (2.6)$$

Shapiro [Ref. 8] shows that the equation governing the flow in the critical region is:

$$\begin{aligned} \left(\frac{k-1}{2}\right)(2V_r + V_\omega \cot\omega + \frac{dV_\omega}{d\omega})(V_{\max}^2 - V_\omega^2 - V_r^2) \\ = (V_r \frac{dV_r}{d\omega} + V_\omega \frac{dV_\omega}{d\omega})V_\omega \end{aligned} \quad (2.7)$$

Eliminating  $\frac{d\theta}{d\omega}$  from Eqn 2.4 and substituting the expressions for  $V_r$ ,  $V_\omega$ ,  $\frac{dV_r}{d\omega}$  and  $\frac{dV_\omega}{d\omega}$  given by Eqns 2.1, 2.4 and 2.5, into Eqn 2.7 gives:

$$\frac{dV}{d\omega} = \frac{V \sin(\omega - \theta) \frac{\sin\theta}{\sin\omega}}{1 - \frac{2V^2 \sin^2(\omega - \theta)}{(k-1)(V_{\max}^2 - V^2)}} \quad (2.8)$$

Designating 'R' as the radius of curvature of the hodograph streamline, one obtains:

$$R = \frac{dV}{\sin(\omega - \theta)d\omega} \quad (2.9)$$

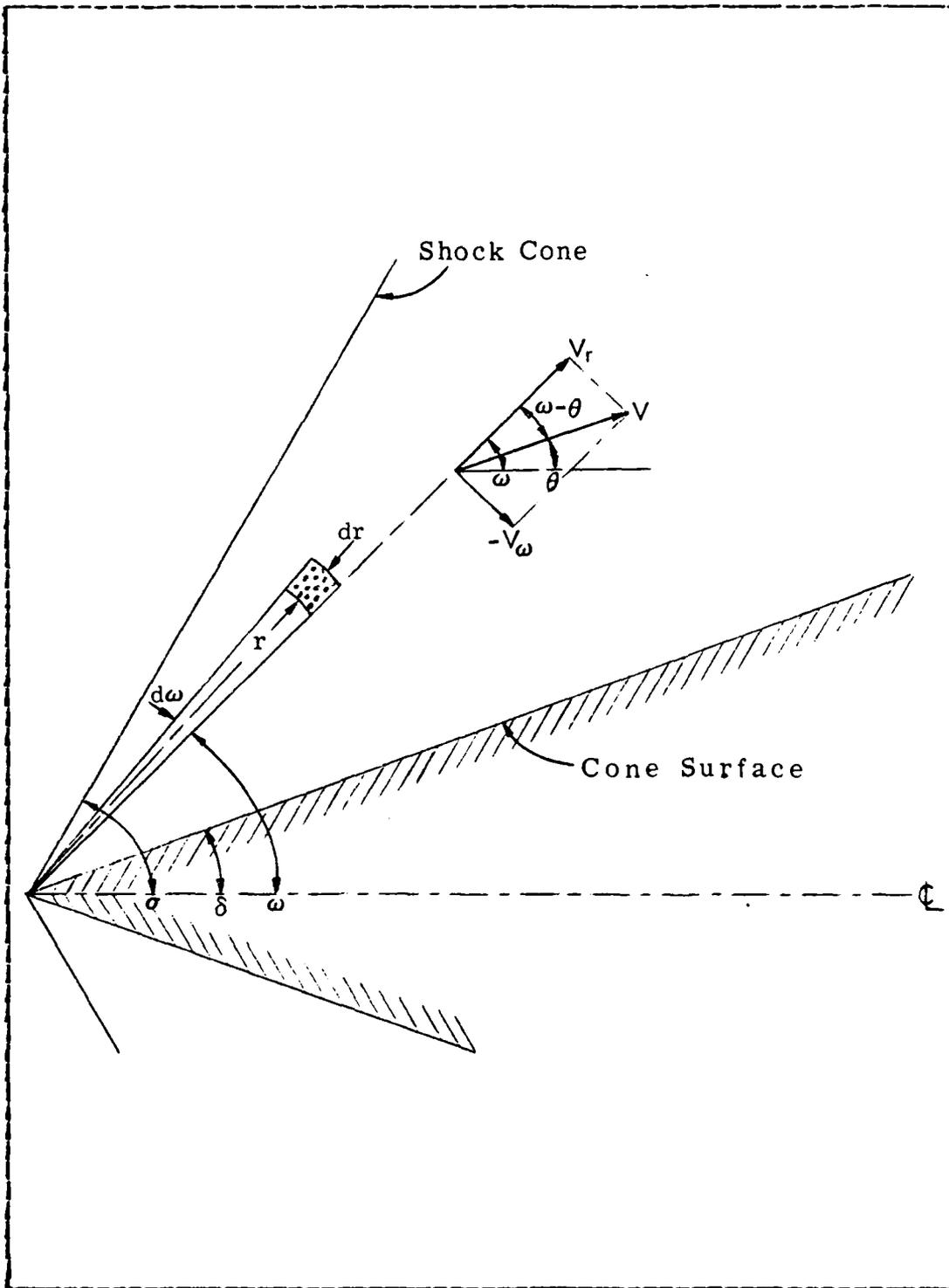


Figure 2.3 Cone and Flow Geometry.

In the following derivations, as per Shapiro [Ref. 8], the spherical coordinates  $r$  and  $\omega$  have been used with the corresponding velocity components  $V_r$  and  $V_\omega$  (see figure 2.3). Only the primary equations which are used in the computer program are presented in this thesis. A detailed derivation of the equations can be found in Shapiro [Ref. 8], and need not be repeated here. The nomenclature used in the equations developed in this chapter is detailed in Table I. In keeping with modern thought, the M-K-S (meter-kilogram-second) unit system has been used throughout this thesis except for occasional lapses during the water calculations when pressures are referred to in units of kilbars.

From the geometry of figure 2.3, it can be seen that:

$$V_r = V \cos(\omega - \theta) \quad \text{and} \quad V_\omega = -V \sin(\omega - \theta) \quad (2.1)$$

In the development of the actual second-order differential equation, Shapiro [Ref. 8] shows that, due to the condition of irrotationality, the following relation must be true:

$$V_\omega = \frac{dV_r}{d\omega} \quad (2.2)$$

Differentiating Eqn. 2.1 with respect to  $\omega$ , one obtains:

$$\frac{dV_r}{d\omega} = -V(1 - \frac{d\theta}{d\omega}) \sin(\omega - \theta) + \frac{dV}{d\omega} \cos(\omega - \theta) \quad (2.3)$$

and

$$\frac{dV_\omega}{d\omega} = -V(1 - \frac{d\theta}{d\omega}) \cos(\omega - \theta) - \frac{dV}{d\omega} \sin(\omega - \theta) \quad (2.4)$$

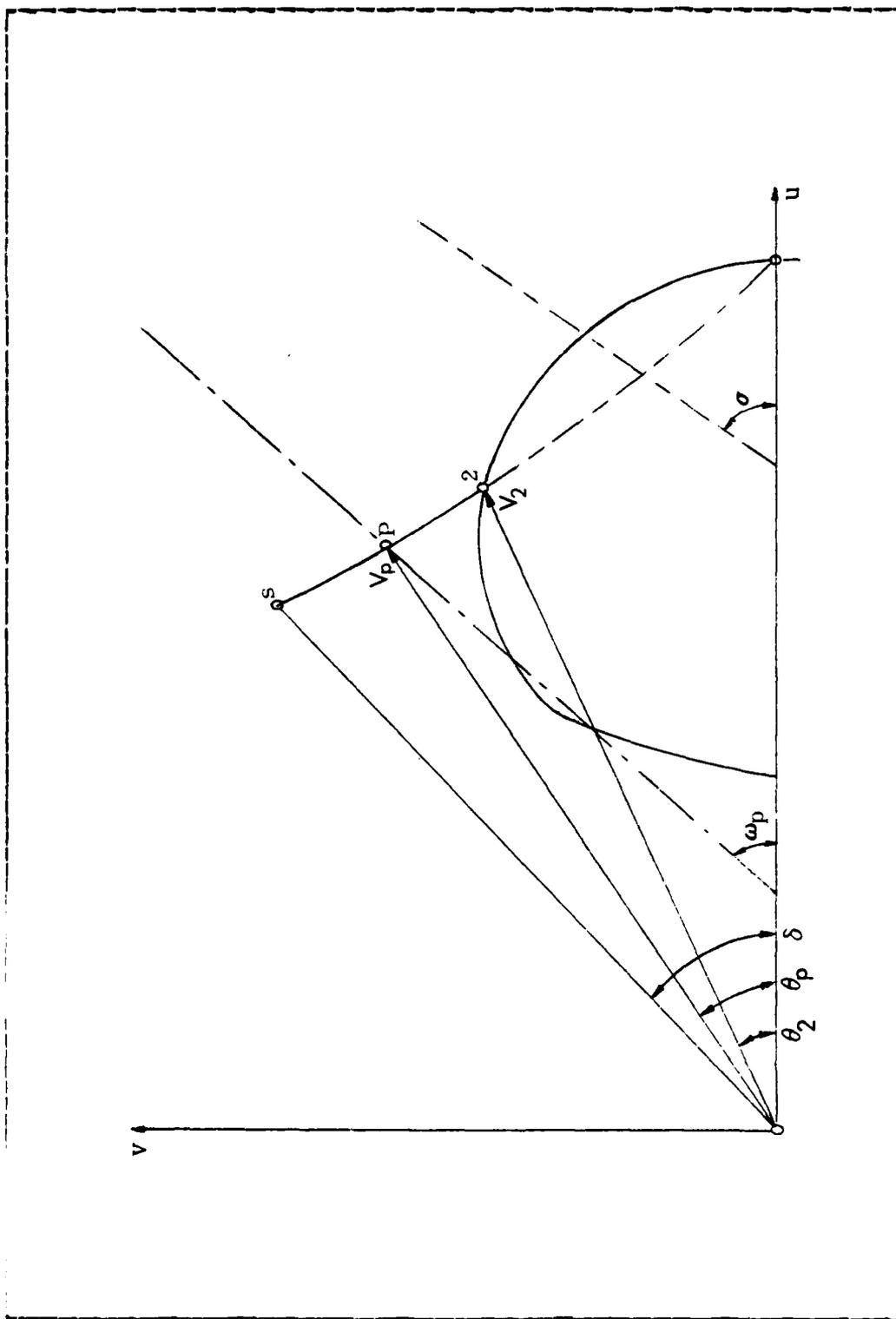


Figure 2.2 Hodograph Image of Typical Streamline.

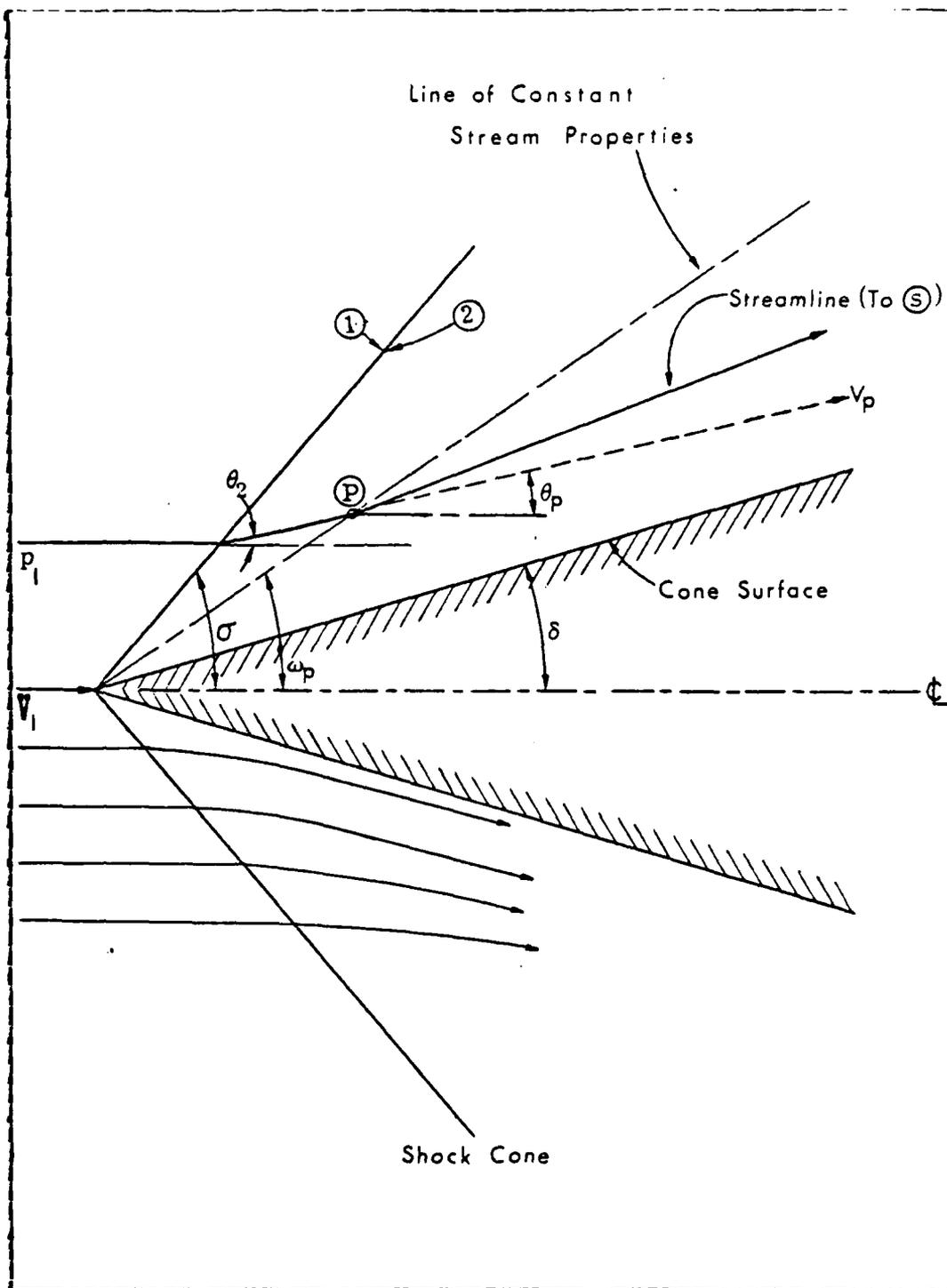


Figure 2.1 Shock Cone and Typical Streamline.

After calculating the thermodynamic changes which occur as a result of the shock front, it is assumed that the fluid properties will remain constant on imaginary cones having a common vertex. By this assumption, the flow past the cone can be calculated. The flow geometry is illustrated by figure 2.1 and figure 2.2, which show a typical streamline and its image in the hodograph plane.

As discussed thoroughly by Shapiro [Ref. 8], there is a discontinuous change in both direction and velocity across the oblique shock front. Points 1 and 2 (see figure 2.1 and figure 2.2) lie, therefore, on a common shock plane which originates at point 1. Between point 2 and the point 's', which is on the cone surface only at infinite distance, there is a region of conical flow where the stream properties vary continuously. The velocity vector to the point 's' in the hodograph plane defines what is called the cone semi-vertex angle with the centerline axis. In the methods which follow, the cone semi-vertex angle is the variable which is ultimately determined. Further, as pointed out by Shapiro [Ref. 8], since all streamlines in the flow experience the same entropy jump across the shock front, the flow between the shock front and the cone surface is both isentropic and irrotational.

The second-order differential equation which actually describes the flow of the fluid past the cone in air is fully developed by Shapiro [Ref. 8]. Shapiro notes that there are two methods commonly used to solve this equation. One method, developed first by Taylor and MacColl [Ref. 8], performs a numerical integration of the equation. The second method, which integrates the equation using a graphical construction method, was first developed by Busemann. The program developed in this thesis utilizes Busemann's method, modified for performance on a modern high-speed computer, to perform the integration of the full second-order differential equation.

The fundamental equation used to describe the thermodynamic state of air is the perfect gas law. In water, the modified Tait equation is the equation most often used to describe the thermodynamic state of the water. The modified Tait equation can be used to describe either pure water or sea water. The form of the modified Tait equation used in this thesis was taken from Helt [Ref. 6], who has continued to perform research in underwater explosion phenomena. An excellent discussion of the modified Tait equation and how it can be utilized is contained in Rowlinson [Ref. 7].

After the thermodynamic properties of the water (or air as the case may be) on the downstream side of the shock front have been calculated, an iteration method, utilizing an automated graphical construction first developed by Eusemann, [Ref. 8], is used to progress from the shock front to the cone surface. The equations needed for use by this iteration method are fully described, for air, by Shapiro [Ref. 8], who describes their development and use. The methods which apply to the flow past a cone in air can, with the necessary changes made for the differences in the thermodynamics of the two fluids, be used to calculate the conical supersonic flow in water. These calculations form the basis for the main part of the FORTRAN program which follows. It should be noted here that Shapiro also points out the pioneering work of Taylor and Maccoll in the 1930's and 1940's on the methods of solution of the flow over a cone in air problem [Ref. 8].

## E. GENERAL EQUATIONS

In the development of the equations which follow, it is assumed that these equations can be validly used to calculate the flow over the cone in either air or water. The equations were primarily developed by Shapiro [Ref. 8] for flow in air.

## II. FUNDAMENTAL EQUATIONS

### A. BACKGROUND

The first step necessary to describe supersonic flow over a cone is to calculate the thermodynamic properties across a shock wave. The basic research into the change of thermodynamic properties across a shock wave in water was extensively conducted and reported upon in Underwater Explosive Research [Ref. 1] during and just after World War II. The best summarization of these works can be found in Cole [Ref. 2].

The primary source used as reference for the calculation of the hydrodynamic properties of sea water at the front of a shock wave is the work of Richardson, Arons, and Halverson [Ref. 3]. They utilized graphical techniques, which were rather crude and tedious, to calculate the thermodynamic data needed to describe the conditions of the sea water. Fuhs [Ref. 4] used the work of Richardson, et al., [Ref. 3], to develop a computer program for the HP41CV hand-held calculator which efficiently calculated the same thermodynamic properties. Fuhs' [Ref. 4] programs provided the basis for the FORTRAN subroutines, used in this work, which calculate the same thermodynamic properties such as pressure, temperature, and density.

The fundamental equations used to calculate the thermodynamic changes which occur across an oblique shock front in air have been well known for decades. These equations are exceptionally well described in Kinney and Graham [Ref. 5] and form the basis for the initial calculations for the supersonic flow past a cone in air.

one-dimensional momentum equation states:

$$p_1 + \rho_1 v_1^2/2 = p_2 + \rho_2 v_2^2/2 \quad (2.25)$$

Now note that typical values for  $\rho_1$  and  $v_1$  for the problem under consideration are as follows:

$$\rho_1 = 1000 \text{ kg/m}^3 \quad \text{and} \quad v_1 = 1500 \text{ m/s}$$

Therefore,  $\rho_1 v_1^2/2 = 1.12 \times 10^9$  Pascals or about 11.1 kilobars. Typical values for  $p_1$ , the upstream pressure, are on the order of 1 to perhaps a few tens of bars. Thus, it can be safely assumed that  $p_1$  can be ignored compared to the upstream dynamic pressure.

Having specified this pressure, the program of this thesis utilizes the subroutines developed by Fuhs [Ref. 4] to calculate the density and velocity components at a point just downstream of the shock front. Since the geometry of the velocity flow across the shock front is the same in either air or water (see figure 2.4), the velocity components determined in the first step can be used to "jump back across" the shock front to the upstream side where the free-stream Mach number  $M_1$  can be determined.

After determining the conditions on the downstream side of the shock front, the program iterates along the streamlines to determine the cone semi-vertex angle. This iteration is performed using the Busemann graphical construction as for the air case. The major difference between the water and air cases, after the initial shocked conditions have been determined, is that the local speed of sound on a streamline must be determined iteratively, for the water case, by using Fuhs' [Ref. 4] subroutines to determine the density and pressure at each point. Knowing the pressure

and density, the speed of sound in the water can be calculated from:

$$c^2 = \left(\frac{nB}{\rho}\right) \left(\frac{c}{\bar{\rho}}\right)^n \quad (2.26)$$

which is derived from the modified Tait equation using the definition of the speed of sound as:

$$c^2 = \left. \frac{\partial p}{\partial \rho} \right|_s \quad (2.27)$$

These equations, combined with the equations developed in section 2 of this chapter, constitute all the equations needed to completely solve the supersonic flow over a cone.

TABLE I  
Nomenclature

c	Speed of sound
k	Ratio of specific heats ( $c_p/c_v$ )
M	Mach number
$M_a$	Molecular weight of air
P	Pressure
r	Radius in spherical coordinates
R	Radius of curvature of the hodograph streamline
T	Temperature
V	Velocity
$V_{max}$	Maximum velocity for adiabatic flow
$\beta$	Angle the streamline makes with the shock plane (same as the angle $\sigma$ ) in air calculations
$\delta$	Cone semi-vertex angle
$\theta$	Flow direction
$\rho$	Mass density
$\sigma$	Shock angle
$\omega$	Angle in spherical coordinates
$\Lambda$	Universal gas constant
( ) <sub>1</sub>	Signifies a condition upstream of conical shock or a shock front
( ) <sub>2</sub>	Signifies a condition downstream of conical shock or a shock front
( ) <sub>s</sub>	Signifies conditions at cone surface
( ) <sub>r</sub>	Signifies a component in r-direction
( ) <sub><math>\omega</math></sub>	Signifies a component in $\omega$ -direction

### III. DESCRIPTION OF THE COMPUTER PROGRAM

#### A. FBCGFAP LOGIC

As discussed in Chapter 1, the computer program developed for this thesis was written in the FORTRAN programming language. Structured programming practices have been followed throughout in that the program is divided into blocks (or modules) of code, each of which is designed to perform a single calculation sequence. Many of the modules have been placed inline rather than being written as separate subroutines or functions. This was done primarily for ease of use and understanding. In addition, however, most modules were placed inline because the parameter list exchanged between the module and the main program would have been excessively long otherwise. Regardless of whether the code is included inline or as a separate module, the program was written in a manner which ensured that the "side-effects" problem discussed by MacLennan [Ref. 9] did not occur. In addition, the use of FORTRAN's COMMON construct has been studiously avoided to prevent the aliasing problem discussed by MacLennan [Ref. 9]. A module hierarchy chart, which shows the major modules of the program and their interconnections, is given as figure 3.1.

The flow logic of the main computer program is illustrated by figures A.1 through A.7 of Appendix A. These flowcharts show that the initial part of the program (up to line 60) is used to initialize certain key variables and to gather input values from the program user. Note that, as stated in the objectives for this program, all user input is verified to ensure that the values entered fall within prescribed limits. In order to ensure that input values are

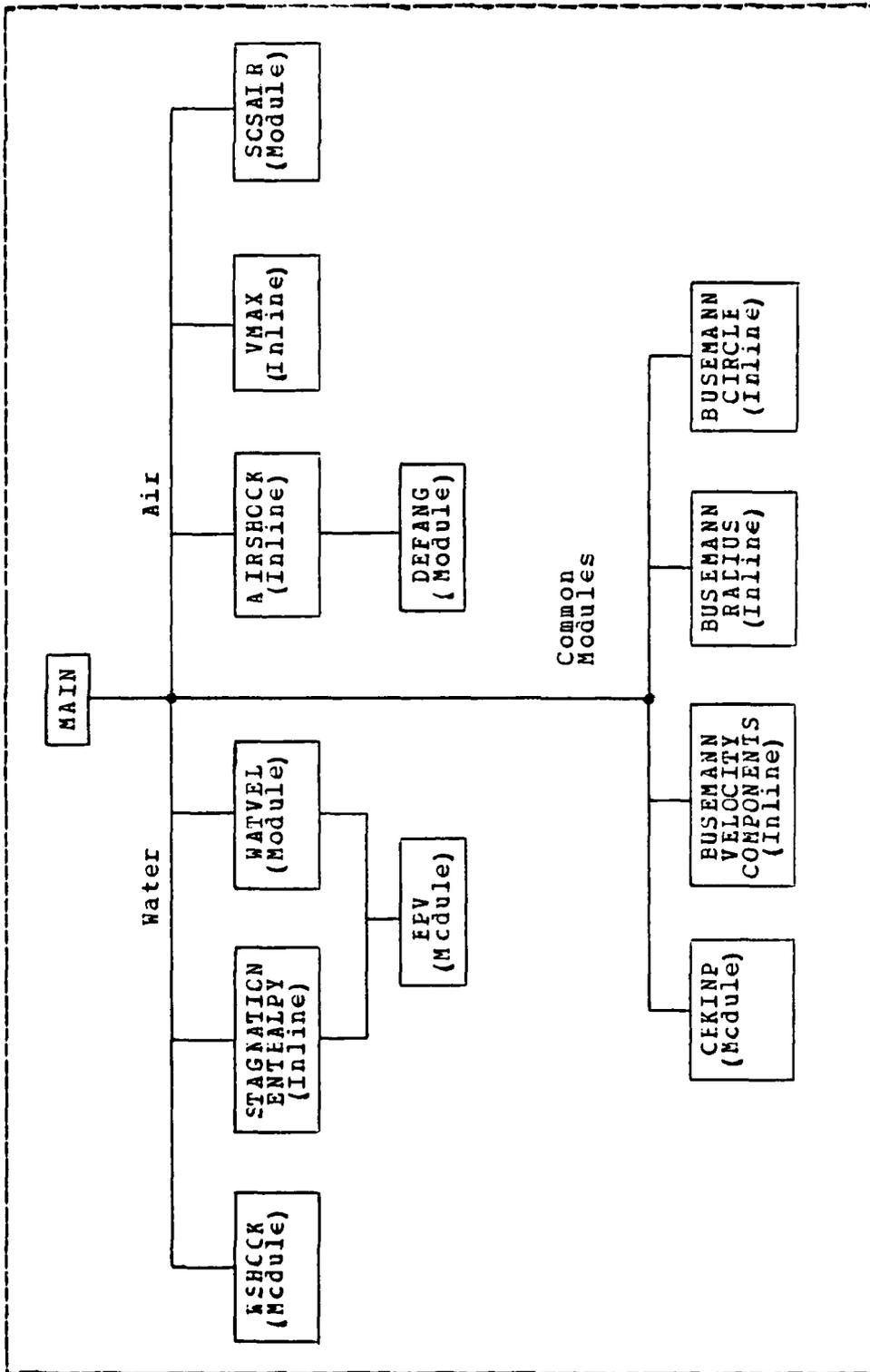


Figure 3.1 Module Hierarchy.

within limits, the subroutine CHKINP is invoked and the input values are passed as parameters. The flow logic of subroutine CHKINP, which is self-explanatory, is included as figures A.8 through A.10 of Appendix A.

After receiving and verifying the user's input, the program begins the calculations required to determine the cone semi-vertex angle. From line 60 to line 400, the program calculates the initial thermodynamic properties of the water or air, as the case may be, and calculates the changes which occur in the fluid properties as a result of the passage of the shock front. As mentioned previously, these calculations for the air case are rather straightforward and, because of this simplicity, the computer code for the air calculation has been placed inline. For the water case, the subroutines WSHOCK and EPV are needed to calculate the required thermodynamic properties. Subroutine WSHOCK calculates the thermodynamic properties of the water at point 1 upstream of the shock front and at point 2, just across the shock front on the downstream or shocked side. Appendix A, figures A.12 and A.13, contains the flowchart which describes the logic of this subroutine. Subroutine EPV was copied, with the author's permission, from Fuhs [Ref. 4] and is fully described in that work. However, the FORTRAN translation of EPV is included in the program listing contained in Appendix C.

Having calculated the initial thermodynamic fluid conditions on both sides of the shock front, the program begins the iteration process required to determine the cone semi-vertex angle. These calculations begin at line 400 of the computer program. This iteration process makes use of the Busemann graphical integration technique discussed by Shapiro [Ref. 8]. First, the radius of the Busemann curve for a point J is calculated using the velocity at that point, the streamline angle at that point, and the value of

the angle  $\omega$  at that point (see figure 3.2). Next, the center of the circle for the Eusemann curve at the point J is calculated using the velocity components at point J (which are designated U and V), the radius calculated in the previous step, and the value of the angle  $\omega$ . This circle center provides the point from which the program "draws" the arc used to calculate the next point on the Eusemann curve. Finally, the velocity and the Eusemann velocity components, at the next point, J+1, are calculated by "swinging an arc" from the circle center calculated in the previous step. In addition, the streamline angle at the next point is calculated.

The program next tests to see if the cone surface has been reached (line 475). This test is conducted by determining if the absolute value of the difference between the streamline angle and the angle  $\omega$  is less than a specified test value (which is set to  $1 \times 10^{-6}$ ). If this difference is less than the test value, the cone surface has been reached. In this case, the program then calculates the thermodynamic fluid conditions at the cone surface and the cone semi-vertex angle and displays the final results of the program. If the difference between the streamline angle and the angle  $\omega$  is greater than the test value, the cone surface has not been reached. In this case, the program calculates the thermodynamic fluid properties at the next point, J+1, then loops back to line 400 to begin the Eusemann graphical integration process for the new point. In the case of a water run, the subroutine WAIVEL is invoked just prior to looping back to line 400. This subroutine calculates the thermodynamic properties of the water at any given point. Appendix A, figures A.14 and A.15, gives a flowchart which demonstrates the logic flow within this subroutine.



The program includes a feature which allows the user to make repeated executions of the program without the requirement of "reloading" the program between each execution. This feature was included by soliciting a response from the user at line 1500 as to whether he/she wishes to make another execution of the program. An affirmative response causes the program to loop back to statement 1 at the beginning of the initialization section of the main program. A negative response causes program termination.

Note that certain separately compiled functions included with this program have not been flowcharted because their logic is so straightforward and because the subprograms in question usually consist of only one or two lines of executable code. The functions fitting into this category are: LTOR, BTCL, and SOSAIF.

#### E. USER INSTRUCTIONS

As discussed in the objectives listed in chapter 1, one goal of this thesis was to ensure that the program presented was easy to use and maintain. Ease of use has been facilitated by including code which verifies that the input provided by the user is within specified ranges. For example, in the water calculations, the pressure at point 2 can be no greater than 100.0 kilobars. This is due to the fact that Fuhs' [Ref. 4] subroutines, which are used to calculate the thermodynamic properties of the water, are based on the work of Richardson, et. al., [Ref. 3], which only gives results up to pressures of 100.0 kilobars. Therefore, while the subroutines could compute the water conditions at pressures greater than 100.0 kilobars, it is uncertain whether the values so calculated would be entirely correct. For this reason, the range of the input values has been restricted. Ease of use is further facilitated by

providing clear, meaningful output. Appendix B contains sample outputs generated by the program for the summary and complete print options of the program for both water and air.

One of the major criticisms leveled by computer scientists against the FCFTAN language is for its lack of a requirement for formally defining all variables used within a program. Similarly, FORTRAN has an implicit declaration policy whereby any variable whose first letter is I, J, K, L, M or N is implicitly declared to be of type Integer. In this program, these two unfortunate characteristics of FORTRAN are avoided by requiring explicit declaration of all program variables.

Ease of maintenance and ease of understanding of the program have been achieved through the use of a liberally commented program and through the use of relatively meaningful variable names. In this regard, the FCFTAN language is less than desirable since it limits variable names to six characters. Appendix C contains a fully documented listing of the computer program and all subroutines or functions used by the main program (other than standard library functions such as sin). Tables VI, VII, VIII and IX, located in Appendix A, contain lists of all variable names used in the main program, their meaning, and their MKS (meter-kilogram-second) units, if appropriate. Table X contains a similar variable list for the subroutine WSHOCK and Table XI contains a variable list for the subroutine WATVEL. These two tables are also located in Appendix A.

Finally, this program was developed using the FORTRAN IV language supplied by the IEM Corporation as part of their IBM 370/3033AP computer system which is the main computer system available at the Naval Postgraduate School. As far as is known, no implementation specific features have been included in this program. Therefore, the program should

execute on any system which has a FORTRAN compiler. In order to execute the program on the IBM 370 computer system available at the Naval Postgraduate School, the following steps must be accomplished in the order given:

- (1) The program must be compiled using the command:

```
FORTHX CCNEFLOW
```

This command invokes the FORTRAN H Extended compiler, which is an optimizing, production run compiler supplied as part of the IBM computer system. The program could also be compiled using another FORTRAN compiler such as the FORTRAN compiler. Note that the above assumes the program supplied by this thesis has been entered into a file which has the filename CCNEFLOW and which is of filetype FORTRAN. Note further that this step need only be performed one time provided errors do not occur.

- (2) Next, the libraries of standard subroutines and functions supplied by the computer center must be attached to the file. This is accomplished by issuing the following command:

```
GIORAI TXLIE FCRTMOD2 MOD2EEH
```

Note that this step need only be performed once provided abnormal terminations do not occur.

- (3) In the FORTRAN language, a formatted input/output (I/O) statement requires the definition of an I/O device. For example, at most installations, by default, I/O device 5 is used for input and I/O device 6 is used for output. To utilize the program of this thesis properly, the following commands must be issued to define the I/O devices to be used:

FILEDEF 06 DISK CONEFLOW OUTPUT  
FILEDEF 07 TERM (TERM)

As a result of the above commands, all output from the program will be written into a file on the user's disk named CONEFLOW with a filetype of OUTPUT. It is not strictly necessary to issue the FILEDEF 06 command shown above, but, if not, all output will be written at the terminal and will not be available for printing. The second file definition command (FILEDEF 07) is absolutely required or the program will not operate. This file definition tells the computer that I/O unit 07 is the computer terminal. In the program of this thesis, all input from the user is requested from I/O unit 07. Therefore, if unit 07 is not defined, the program cannot receive any input.

(4) Finally, the program can be executed by issuing the command:

ICAD CONEFLOW (START

When program execution is completed, a cryptic message of the form R; T=0.01,C.01 16:45:00 will be displayed at the terminal. This is simply a message from the computer's operating system indicating that the task just requested has been completed. The user can now utilize the system's operating commands to review the output from the execution of the program.

If the user desires to execute the program again, return to step (3) of the above procedure. If the output from the previous execution is no longer needed, the same FILEDEF 06 command can be issued. However, if the previous execution's output is needed, the filetype of the FILEDEF 06 command should be changed (e.g. change OUTPUT in the above command to CBTFT2).

#### IV. PROGRAM CALCULATION RESULTS

##### A. PROGRAM RESULTS FOR THE CALCULATION OF CONICAL FLOW IN AIR

The program accurately calculates the flow over a cone in air. Numerous program executions were made for the air case at various upstream Mach numbers and for various shock angles. The results obtained from these program executions are summarized in Tables II and III. The numbers produced by the program were compared for the variables shown in the tables to the tables given in Kinney and Graham [Ref. 5], against the graphs given by Shapiro [Ref. 8], and to the tables produced by Kopal [Ref. 10]. The values given in these sources are included in Tables II and III, where appropriate, along with the calculated results. As can be seen from these two tables, the calculations made by the program give quite accurate results (less than a 1% error in most cases). Certainly, the program is much more accurate than one's ability to read the graphs presented in Shapiro [Ref. 8].

Based on these comparisons, it is believed that the part of the program which calculates supersonic flow over a cone in air is accurate. The importance of this fact is that, by feeling confident that the procedure followed for the air case is valid and that this procedure has been correctly implemented in the programming language, it is safe to assume that the Eulerian calculation procedure utilized for the air case can be accurately applied to the water case provided the water conditions at each point are calculated correctly.

TABLE II  
Program Results and Comparisons

Shock Angle	Cone Angle		Mach Number at Cone Surface		Drag Coefficient	
	Prog.	(1)	Prog.	(2)	Prog.	(2)
15.719	5.7098	5.0	2.86	2.9	0.036	0.04
16.458	7.6728	7.5	2.71	2.75	0.089	0.08
17.196	10.0422	10.0	2.50	2.5	0.174	0.18
18.034	12.5096	12.5	2.29	2.3	0.285	0.29
18.872	15.0107	15.0	2.06	2.05	0.420	0.44
19.710	17.5024	17.5	1.83	1.85	0.576	0.575
20.548	20.0084	20.0	1.59	1.65	0.751	0.75
21.386	22.5080	22.5	1.35	1.38	0.941	0.94
22.224	25.0091	25.0	1.08	1.2		
23.062	27.5183	27.5				
23.900	30.0276	30.0				
24.738	32.5361	32.5				
25.576	35.0463	35.0				
26.414		40.0				
27.252		45.0				

Notes: (1) The shock angle values were taken directly from Kofal. In Kofal's work, the entering parameters are the cone semi-vertex angle and upstream Mach number. This is the reverse of the way the program operates. However, using Kofal's values allows for comparisons to be made more easily.

(2) The numbers in these columns were read from the graphs presented in Shapiro. Thus, the numbers are subject to interpretation depending upon how well one reads graphs. Regardless, the numbers presented were read from the graphs at the whole number cone angles given in the same row (e.g. the first row of the table was read at the cone angle of 5.0 degrees). Thus, there is some room for error here since the program values are calculated at slightly different cone angles. As one can see, however, the values match relatively well.

(3) The following variables were held constant at the values given for all program executions:  $T_1 = 298.16$  Kelvin;  $P_1 = 101300.0$  Pascals; and  $M_1 = 3.0$ .

## V. CONCLUSIONS

As discussed in the Introduction, this thesis had certain goals which it was desired to achieve. It is believed that these goals have all been successfully met. The program of this thesis has been successfully translated from Basic into FORTRAN as desired. In the translation of the original program, modern structured programming practices have been followed to the greatest extent possible. Finally, the program presented is quite "user-friendly" and is very well-documented.

As was demonstrated in the results chapter, this program calculates correct results for each of the air or water cases. It is mentioned here that, as pointed out frequently by computer scientists, it is virtually impossible to test a program for all possible cases. Therefore, no program is entirely error-free. It is believed that the program of this thesis is as free of errors as is possible without exhaustive, time-consuming and extremely expensive testing.

As a result of having met the three goals mentioned in the first paragraph, it is believed that the program of this thesis will provide an excellent working tool for researchers in the field of shaped-charge jet penetration in water. The fact that the program is well-documented will make any modifications or extensions to the program, should that be required or desired, easy to perform. Further, since the program is easy to use, only a cursory knowledge of computers is needed in order to utilize the program. These features are requirements for any computer program which is to be used by scientists as a tool. Too often, computers and their programs require that the people who simply desire to use their capabilities must learn a great deal of detail

PLOT OF DRAG COEFFICIENT  
VERSUS FREESTREAM MACH NUMBER  
WITH CONE ANGLE AS PARAMETER

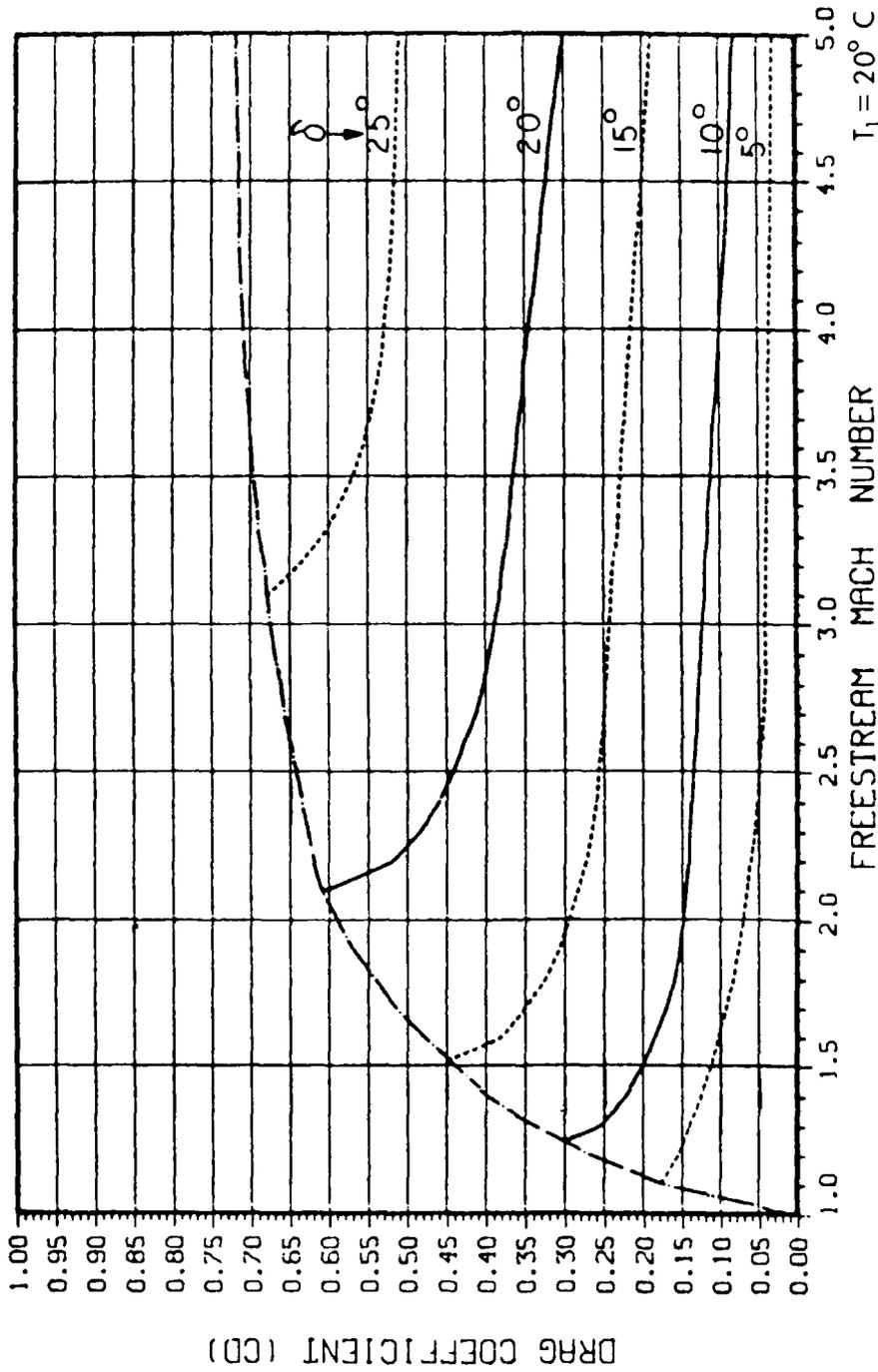


Figure 4.3 Freestream Mach Number vs Drag Coefficient for Supersonic Flow in Water.

PLOT OF CONE SURFACE MACH NUMBER  
VERSUS FREESTREAM MACH NUMBER  
WITH CONE ANGLE AS PARAMETER

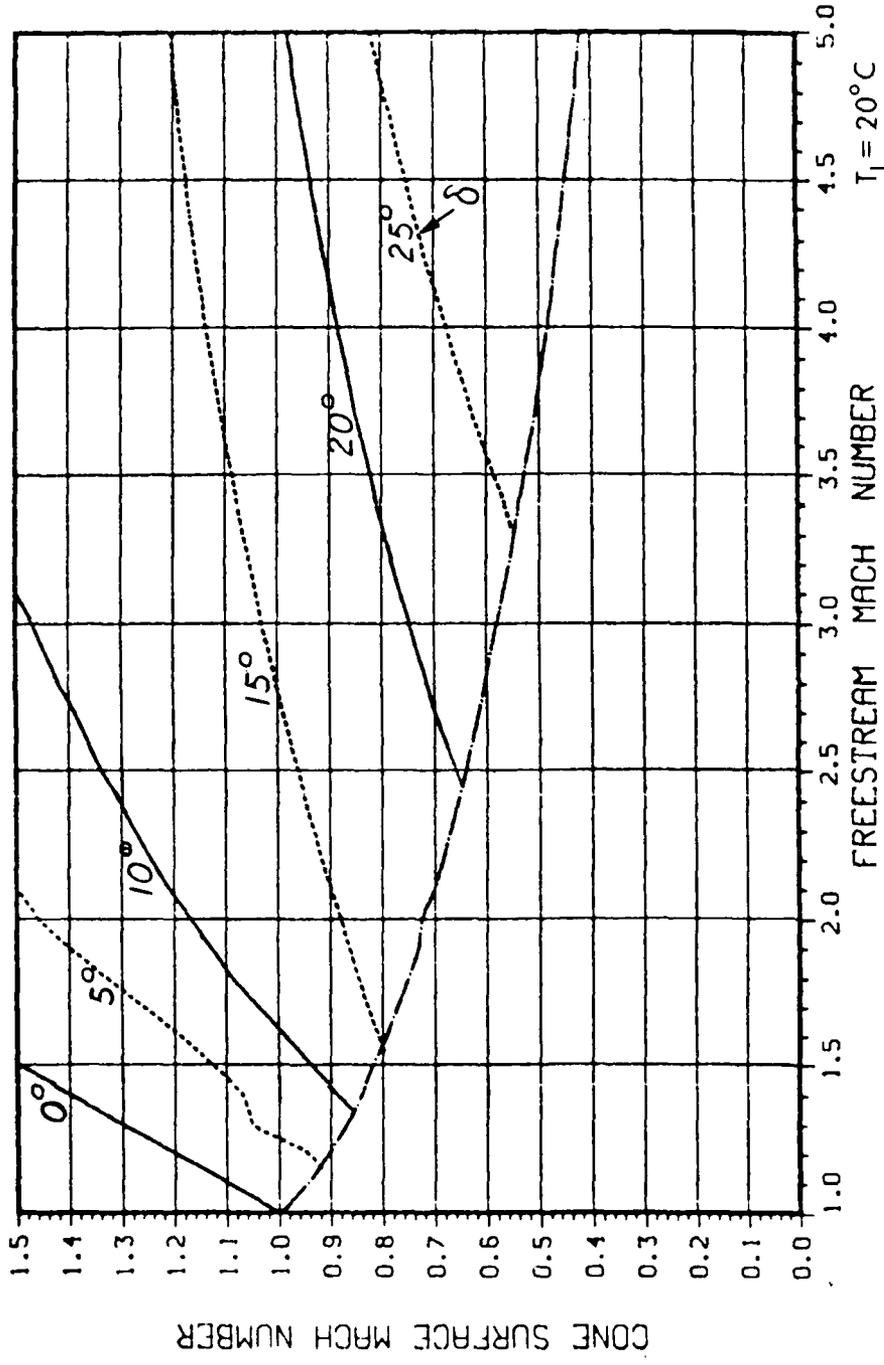


Figure 4.2 Freestream Mach Number vs Surface Mach Number for Supersonic Flow in Water.

PLOT OF SHOCK ANGLE  
 VERSUS FREESTREAM MACH NUMBER  
 WITH CONE ANGLE AS PARAMETER

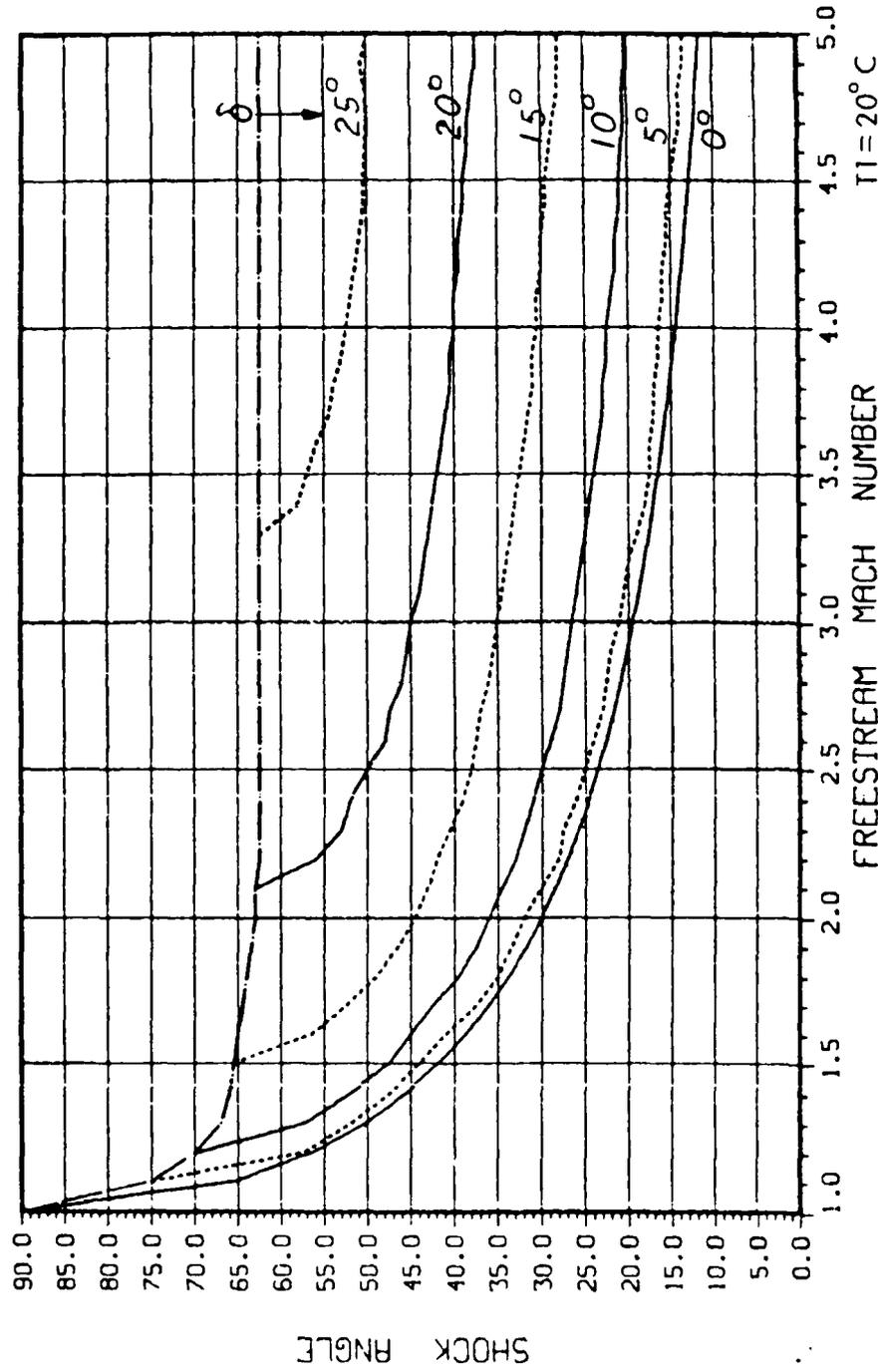


Figure 4.1 Freestream Mach Number vs Shock Angle for Supersonic Flow in Water.

computer system is being utilized at near its capacity, the same response time can increase up to about 2 to 3 minutes. This result, however, is symptomatic of any computer system being utilized at near capacity and is not a result or a reflection of the program's structure or design.

greater than 5.0 (a convenience for ease of presentation only). The pre-processed data points were passed to an interpolation and smoothing program developed by Frake [Ref. 12] which used a quadratic Shepard's method to smooth the data points generated. After smoothing, the data was post-processed to remove erroneous interpolation points which were generated because the original data was irregular. Finally, the processed data was passed to a FORTRAN program developed by the author of this thesis which utilized the Contouring feature of the DISSPLA graphing routines of the ISSCC company. The results of this extensive processing are the graphs presented in this thesis. Because of the irregularity of the original data (which causes difficulties in the quadratic Shepard's interpolation method), and because of the smoothing required by the contouring routines of DISSPLA, the graphs presented are somewhat rough and should, therefore, only be used for "back of the envelope" calculations. Accurate results are provided by the CONEFLOW program which should be utilized for more precise work.

Finally, one of the primary reasons for translating the original Basic programs into a higher-order language (in addition to the desire to make the programs more accessible) was to speed the execution of the program. In this regard, the work of this thesis has more than accomplished this result. It must be mentioned here that the program will execute extremely fast, considering all the iterations needed, provided the computer system in use is not heavily in use at the time. For example, at the Naval Postgraduate School, when the computer system is relatively free, the program operates so quickly that one execution requires less than 5 seconds (from the time of the last user input to the time of the request by the program for an indication of whether another execution is desired). Conversely, when the

This conclusion is true provided the upstream Mach number is held constant by varying the pressure behind the shock front through the mechanism of the pressure multiplication factor. Having determined that the calculations can be conducted independently of the temperature, plots of the variables of interest in water were made in a manner similar to the graphs for air shown in Shapiro [Ref. 8].

Finally, having determined that the calculations could be made independently of the upstream temperature and that these calculations would be accurate (within 1% for cone angle and 13% for pressure), it was decided to plot various parameters of interest for the flow over a cone in water. These graphs are presented as follows:

- (1) Figure 4.1 is a plot of the freestream Mach number versus the shock angle with the cone semi-vertex angle as an entry parameter. This graph is the water analog to Figure 17.7.(a) of Shapiro [Ref. 8].
- (2) Figure 4.2 is a plot of the freestream Mach number versus the Mach number at the cone surface with the cone semi-vertex angle as an entry parameter. This graph is the water analog to Figure 17.7.(c) of Shapiro [Ref. 8].
- (3) Finally, Figure 4.3 is a plot of the freestream Mach number versus the drag coefficient ( $C_D$ ) with the cone semi-vertex angle as an entry parameter. This graph is the water analog to Figure 17.7.(f) of Shapiro [Ref. 8].

It should be noted here that the graphs developed and presented in this thesis were obtained through repeated executions of the CONEFLOW program, which generated approximately 3600 data points per graph. These data points were pre-processed to remove entries with a Mach number of

TABLE V  
Cone Semi-vertex Angle Variation with Temperature

Upstream Temperature (Kelvin)	Upstream Mach Number	Pressure Multiplication Factor	Cone Semi-vertex Angle (Degrees)
273.16	3.197999	1.000	4.5866
283.16	3.198033	1.051	4.8844
293.16	3.198458	1.090	4.8159
303.16	3.198541	1.115	4.7786
313.16	3.198482	1.128	4.7565
323.16	3.198257	1.130	4.7513
333.16	3.198505	1.125	4.7653
343.16	3.198392	1.111	4.7922
353.16	3.198447	1.090	4.8322
363.16	3.198063	1.062	4.8786
373.16	3.198013	1.031	4.9321

Note: In the execution of the program, the upstream pressure was held constant at 101300.0 Pascals and the shock angle was held constant at 20 degrees.

parameter, does not exist. Since water is not a thermally and calorically perfect fluid, it is not possible to plot universal graphs of the variables of interest.

To gain insight into the sensitivity of the supersonic conical flow in water to the upstream conditions, a study was conducted in which calculations were performed to determine whether the cone semi-vertex angle is relatively independent of the upstream temperature or whether various program executions are required to show the variation of cone semi-vertex angle with water temperature.

The determination discussed above was made by holding the upstream pressure, the shock angle, and the upstream Mach number constant. Then the upstream temperature was varied to determine the variation in the cone semi-vertex angle. In order to hold  $M_1$  constant, the pressure multiplication factor, described in Chapter II.D, was varied until the Mach number upstream for the new temperature matched (within reasonable accuracy) the upstream Mach number for the original temperature. Note that "reasonable accuracy" for matching of the upstream Mach numbers meant matching the numbers to the fourth decimal place. Table V provides the results from these calculations. As can be seen from this table, the cone semi-vertex angles calculated for the various temperatures are all of approximately the same value (within 1%). The pressure downstream of the shock wave varies by 13% as can be seen in Table V. While there is small variation between the values for the cone semi-vertex angles, it must be remembered that, for practical purposes, these variations are very small and can be safely ignored in examining the flow over a cone in water. Thus, it was concluded that the calculation of the cone semi-vertex angle in water can be conducted independently of the water temperature (at least for the accuracy required for the calculation of the movement of the metal jet from a shaped-charge).

Liepmann and Roshko [Ref. 11] demonstrate that the thermal and caloric equations of state are thermodynamically related by:

$$\left. \frac{\partial h}{\partial p} \right|_T = v - T \left. \frac{\partial v}{\partial T} \right|_p \quad (4.6)$$

Introducing equation 4.2 into equation 4.6 results in the right-hand side of equation 4.6 becoming equal to zero. Thus, for a thermally perfect gas, the enthalpy is independent of the pressure, and, hence, enthalpy is a function only of temperature. Therefore, a necessary condition for equation 4.5 to be valid is to have a thermally perfect gas.

For water, Richardson, et. al., [Ref. 3], demonstrate that the heat capacities are functions of the temperature of the water. Thus, equations 4.4 and 4.5 are not valid for water, and water is not a calorically perfect fluid.

If a fluid is both thermally and calorically perfect, the conditions across a normal shock front depend only on  $k = c_p/c_v$  and on the freestream Mach number. Conversely, if the fluid is not thermally and calorically perfect, additional variables must be specified in order to define the shock conditions. Extending the argument, one can state that the supersonic flow of a thermally and calorically perfect fluid over a cone is a function only of the heat capacity ratio,  $k$ ; the freestream Mach number,  $M_1$ ; and the shock front angle,  $\sigma$ . As a result, a single graph is sufficient to represent all flows over the cone. This is the condition for a perfect gas.

In contrast, when the fluid is not thermally and calorically perfect, the solution for supersonic conical flow is dependent upon variables other than  $k$ ,  $M_1$ , and  $\sigma$ . Consequently, a universal graph of the shock angle versus the freestream Mach number, with the cone angle as an entry

which describe the cone flow under all conditions as is shown by Shapiro [Ref. 8]. Whether universal results for the water case could be obtained was not so clear; this point is now discussed.

Liepmann and Rosko [Ref. 11] discuss the general thermodynamics of fluids and introduce the concept of a "thermal" equation of state. In general, the thermal equation of state for any fluid is given by:

$$f(p, \rho, T) = 0 \quad (4.1)$$

For a perfect gas, the thermal equation of state is given by:

$$pv = RT \quad (4.2)$$

where  $R = \Lambda / M_a$ . Another equation which relates the thermodynamic variables  $e$ ,  $v$ , and  $T$  is the "caloric" equation of state, which, in a general form, is given by:

$$f(e, v, T) = 0 \quad (4.3)$$

A calorically perfect gas is defined by:

$$e = c_v T \quad (4.4)$$

or by:

$$h = c_p T \quad (4.5)$$

where  $e$  is internal energy and  $h$  is enthalpy. Thus, the heat capacity (either  $c_v$  or  $c_p$ ) is a constant for a calorically perfect gas.

TABLE IV  
Properties of Sea Water at a Shock Front

kbar	Water Velocity at Point 2 (m/s)		Shock velocity (m/s)		Speed of sound at Point 2 (m/s)		Specific Volume at Point 2 (cm <sup>3</sup> /gm)					
	F10G	Fuhs	RAH	P10G	Fuhs	RAH	E10G	Fuhs	E10G	Fuhs	RAH	
5	256	258	257	1934	1939	1930	2193	2201	2190	.8581	.8572	.8593
10	430	434	433	2298	2306	2290	2724	2736	2720	.8036	.8026	.8040
20	653	698	658	2856	2868	2845	3515	3533	3510	.7490	.7480	.7483
30	895	907	905	3299	3315	3285	4130	4152	4125	.7194	.7185	.7186
40	1076	1084	1080	3677	3695	3665	4646	4671	4640	.6995	.6987	.6989
50	1233	1243	1240	4010	4031	4000	5097	5125	5095	.6848	.6840	.6842
60	1376	1386	1385	4313	4336	4300	5503	5533	5495	.6735	.6727	.6728
70	1507	1518	1515	4595	4620	4585	5876	5909	5870	.6646	.6639	.6641
80	1626	1633	1635	4864	4892	4855	6229	6265	6225	.6582	.6576	.6579
90	173	1747	1740	5128	5159	5120	6569	6609	6570	.6542	.6539	.6542

Notes: (1) The conditions upstream of the shock front were established to be as follows: temperature = 0° C; salinity = 0.7 M NaCl; acoustic velocity,  $c_0$ , = 1443 m/s.

(2) The values for the two columns on the right of each major division identified as Fuhs and RAH, were taken from Table II of Fuhs [Ref. 4].

## E. PROGRAM RESULTS FOR THE CALCULATION OF CONICAL FLOW IN WATER

As mentioned in the previous section, the calculation procedure was verified for the air case by comparison with known results. This comparison showed the method used in the calculations was correct. The next step taken was to determine if the subroutines used to calculate the thermodynamic properties of the water gave results comparable to those of Richardson, et. al., [Ref. 3], and Fuhs [Ref. 4]. Calculations were performed which gave results which could be compared to Table II of Fuhs [Ref. 4]. Table IV presents a comparison between the numeric results calculated by the program and those given by Fuhs. As can be seen, these results are, within reasonable error, remarkably similar. This is not unexpected since the subroutines used by the main program to calculate these variables are essentially direct language translations of Fuhs' programs. Therefore, the results should be similar.

Having verified that the procedure utilized in the calculation of the cone semi-vertex angle was correct (through the air results) and having verified that the program correctly calculated the thermodynamic properties of the water at any point, it was believed that the program could be executed for the water case, for various initial conditions, with certainty that the results so calculated would be accurate. However, a question arose as to whether the calculations performed in the water case were independent of the upstream thermodynamic properties of the water. In air, the only quantities needed to calculate the cone semi-vertex angle are the upstream Mach number ( $M_1$ ) and the shock angle ( $\sigma$ ) (e.g. see figure 17.7(a) of Shapiro [Ref. 8]). Thus, for air, the calculation of the cone semi-vertex angle is independent of the upstream temperature or pressure, and, therefore, universal curves can be drawn

TABLE III  
Program Results and Comparisons (cont'd.)

Shock Angle	Pressure Ratio (P2/P1)		Sound Velocity Ratio (C2/C1)		Temperature Ratio (T2/T1)	
	Prog.	(2)	Prog.	(2)	Prog.	(2)
15-719	1-0287	1-0283	1-0040	1-0039	1-0081	1-0083
20-418	1-1161	1-1163	1-0158	1-0157	1-0319	1-0321
21-360	1-2708	1-2710	1-0350	1-0350	1-0713	1-0710
22-274	1-4841	1-4838	1-0590	1-0589	1-1214	1-1217
23-275	1-7473	1-7475	1-0857	1-0853	1-1788	1-1785
24-335	2-0533	2-0532	1-1142	1-1138	1-2415	1-2416
25-618	2-4006	2-4006	1-1442	1-1440	1-3092	1-3094
26-514	2-7849	2-7851	1-1754	1-1752	1-3815	1-3814
27-822	3-2043	3-2043	1-2076	1-2079	1-4582	1-4579
28-555	4-1356	4-1395	1-2746	1-2744	1-6245	1-6249
29-860	5-1350	5-1888	1-3439	1-3425	1-8061	1-8076
30-863	6-5285	6-3314	1-4149	1-4179	1-8020	1-8036
31-263	7-5905	7-5935	1-4888	1-4914	2-2164	2-2186

- Notes: (1) The shock angle values were taken directly from Kofal. In Kofal's work, the entering parameters are the cone semi-vertex angle and upstream Mach number. This is the reverse of the way the program operates. However, using Kofal's values allows for comparisons to be made more easily.
- (2) The numbers in these columns were read from Table IX of Kinney and Graham. A linear interpolation between the values listed in Table IX of Kinney and Graham was made to arrive at the numbers presented in these columns. This may account for some of the variation seen between the values calculated by the program and those of the reference material. As can be seen from the table, even these variations are quite small.
- (3) The following variables were held constant at the values given for all program executions:  $T_1 = 298.16$  Kelvin;  $P_1 = 101300.0$  Pascals; and  $M_1 = 3.0$ .

about the computer in order to use the tool. In this program, an understanding of FORTRAN is not required in order to utilize the program or its design. Unfortunately, due to the operating system of the IBM 370 computer system, the same cannot be said about the steps required in order to actually use the program. Some familiarity with the computer system in use at the users location will be a necessity in order to operate the program correctly.

Finally, as mentioned in the Introduction to this thesis, the program will provide an excellent test case and comparison model for a computer program which models the actual flow of the metal jet from an explosive shaped-charge fired through the water. The actual situation is a blunt-nosed, rather than a sharp-pointed conical, flow problem which is much more difficult to solve. Therefore, a known solution and methodology of solution for the easier problem is a necessary first step to the solution of the larger problem. It is believed that the program of this thesis will serve as this necessary first step. Further, it is believed that, when the equations solving the actual flow problem are developed in their final form, the program of this thesis, due to its ease of modification, will serve as the programming model for the program which calculates the blunt-nosed flow problem.

APPENDIX A  
PROGRAM FLOWCHARTS

This appendix contains the logic flowcharts of the main program and its subroutines (i.e. those which have not been described elsewhere or which are not part of standard computer center libraries). The language is kept rather general so the overall logic of the program can be demonstrated. If the reader desires to know how a particular logic sequence is implemented, he need only refer directly to the program segment which the logical flowchart is describing. The flowcharts included in this appendix are listed below:

- (1) The Main Program Flowchart consists of Figures A.1 through A.7.
- (2) Function CHKINP consists of Figures A.8 through A.10.
- (3) Subroutine DEFANG consists of Figure A.11.
- (4) Subroutine WSEOCK consists of Figures A.12 and A.13.
- (5) Subroutine WATVEL consists of Figures A.14 and A.15.

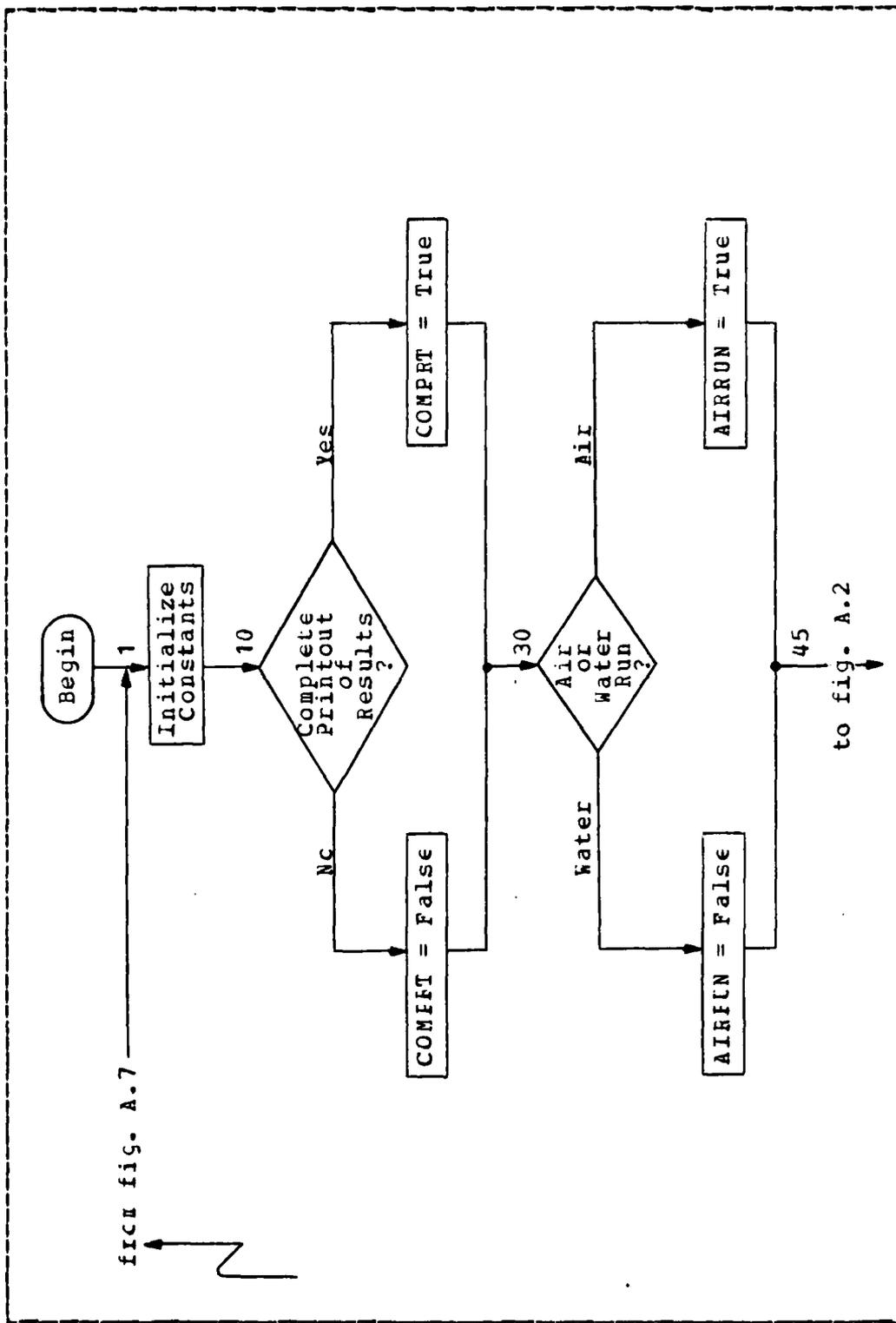


Figure A.1 Main Program Flowchart.

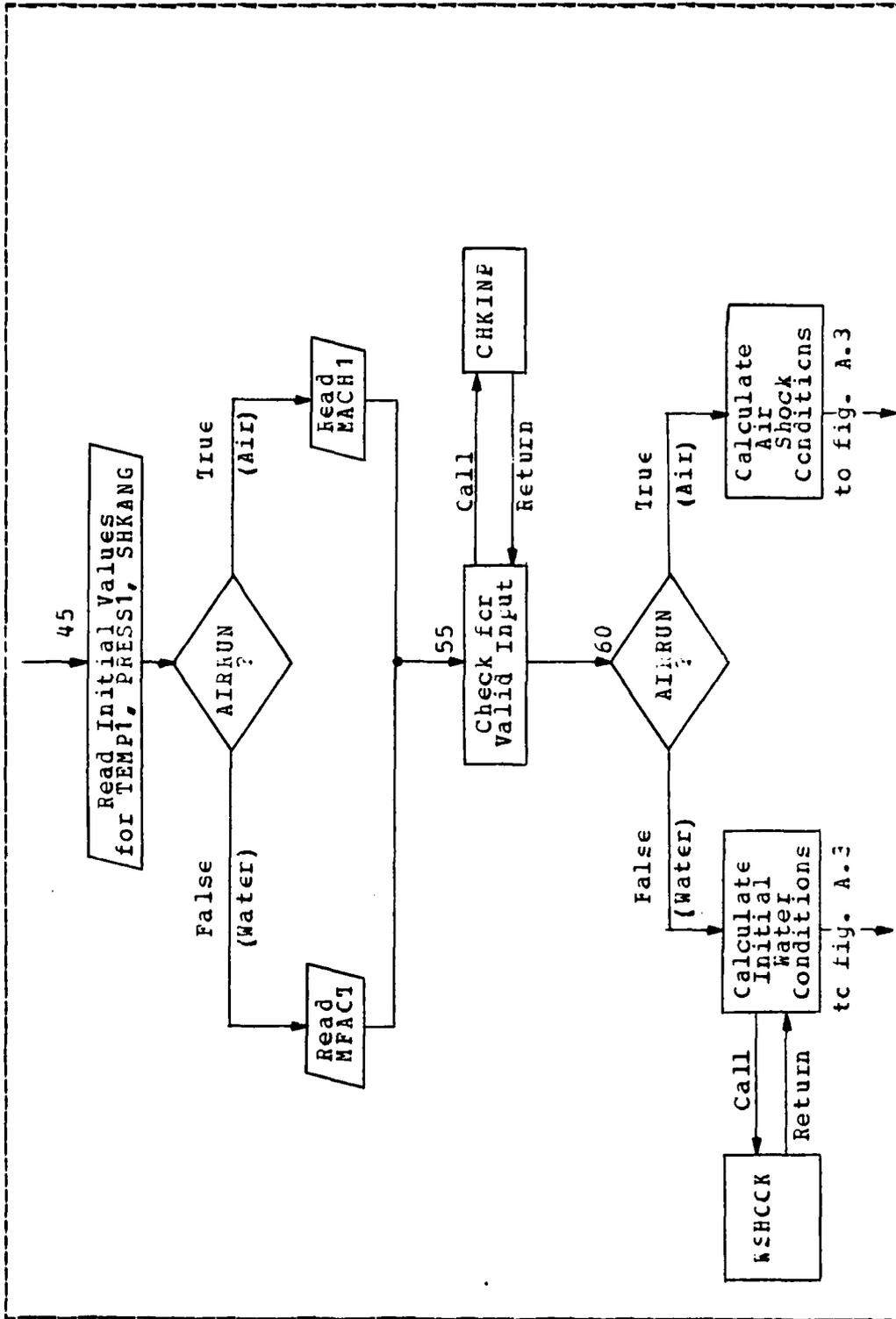


Figure A.2 Main Program Flowchart (ccnt'd.).

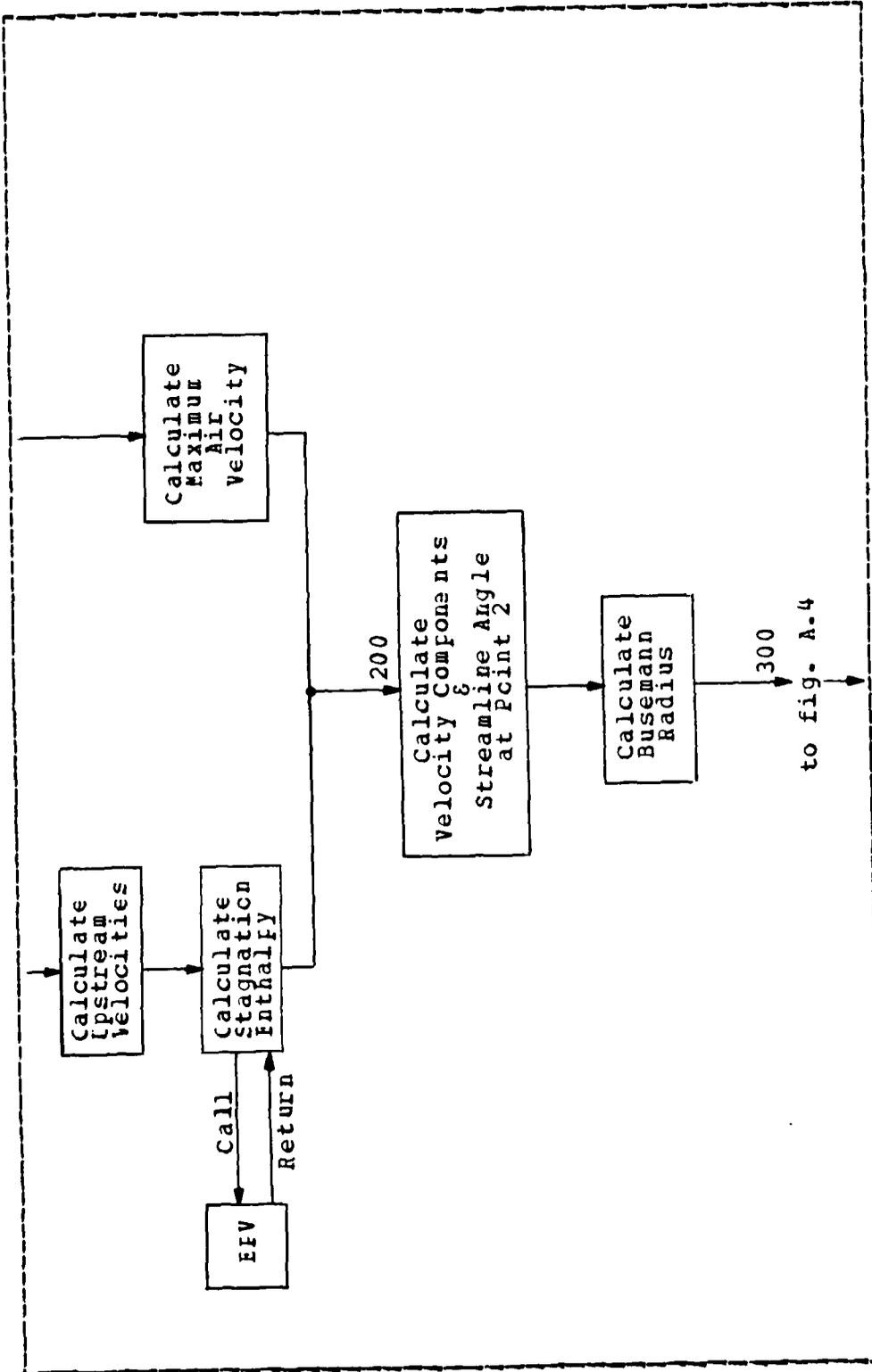


Figure A.3 Main Program Flowchart (ccnt'd.).

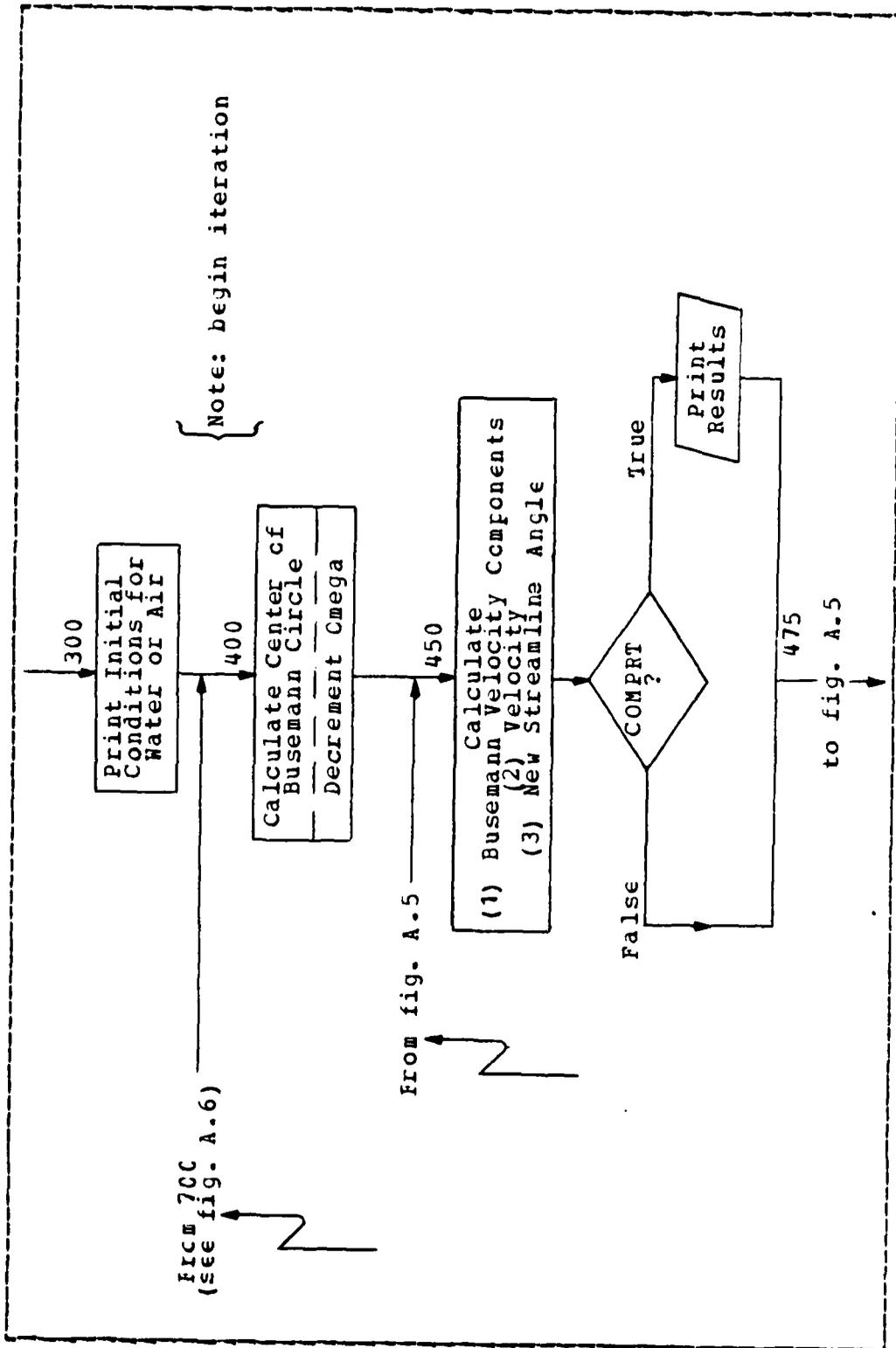


Figure A-4 Main Program Flowchart (ccnt'd.).

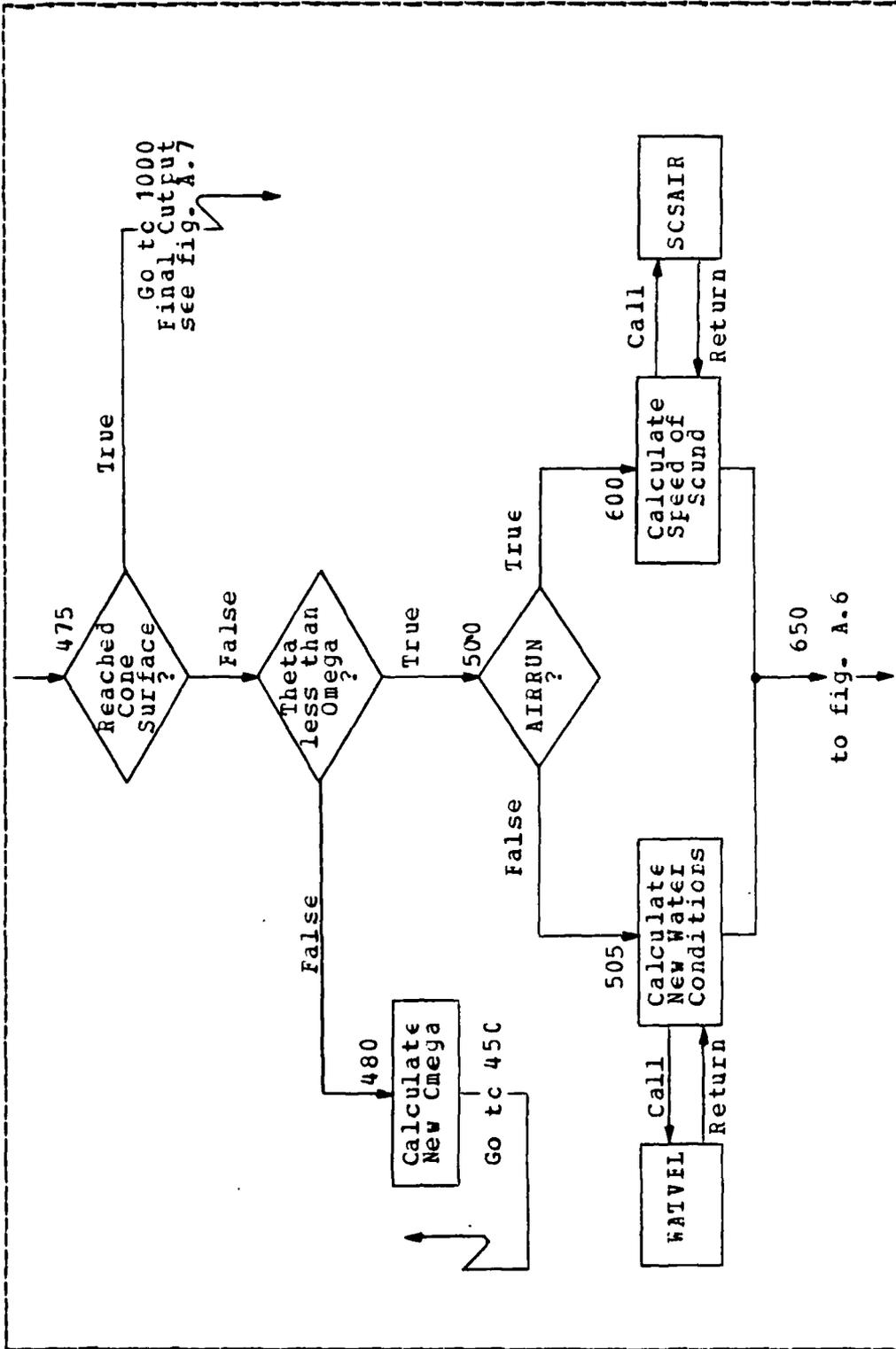


Figure A.5 Main Program Flowchart (cont'd.).

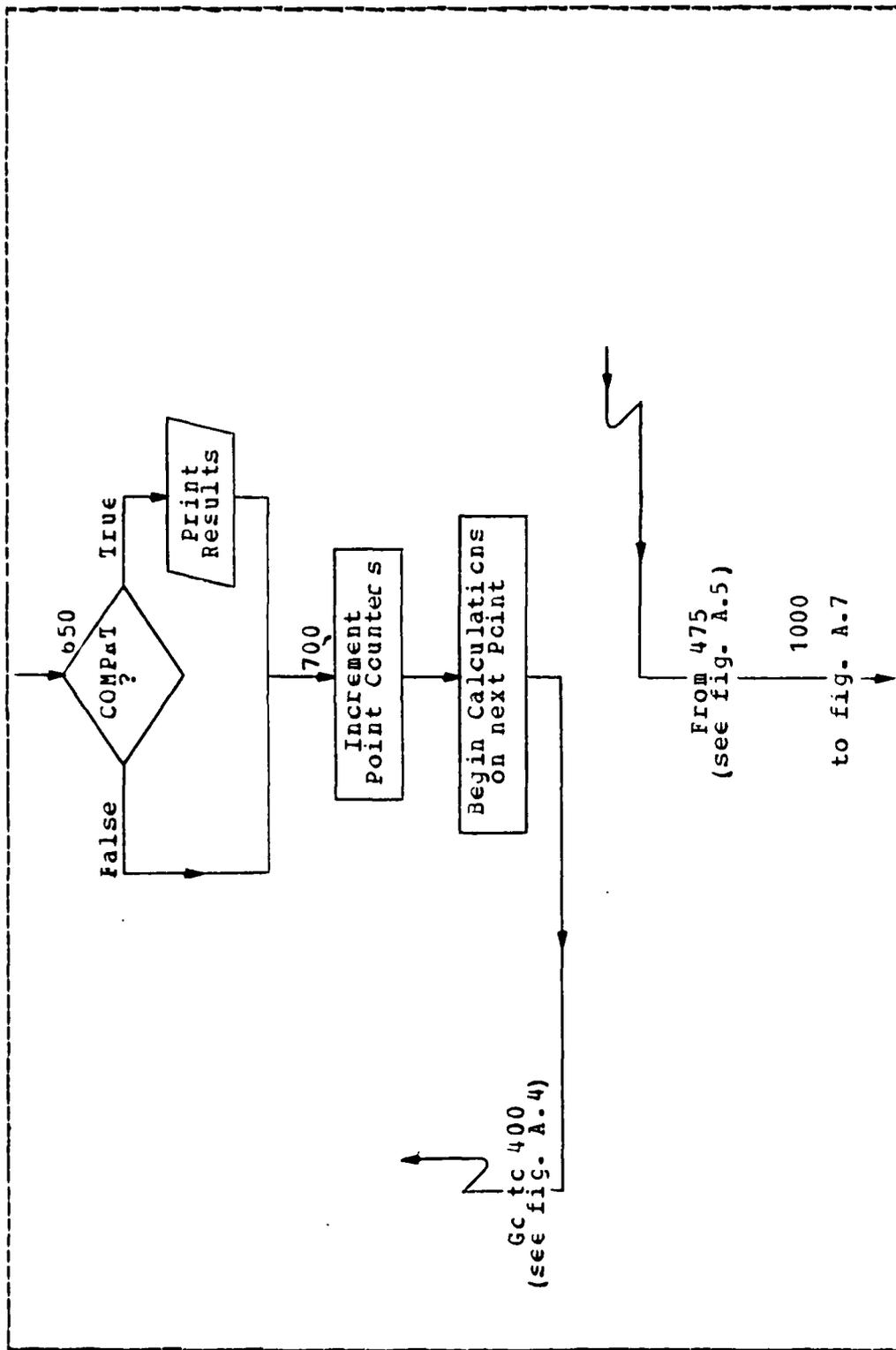


Figure A.6 Main Program Flowchart (ccnt'd.).

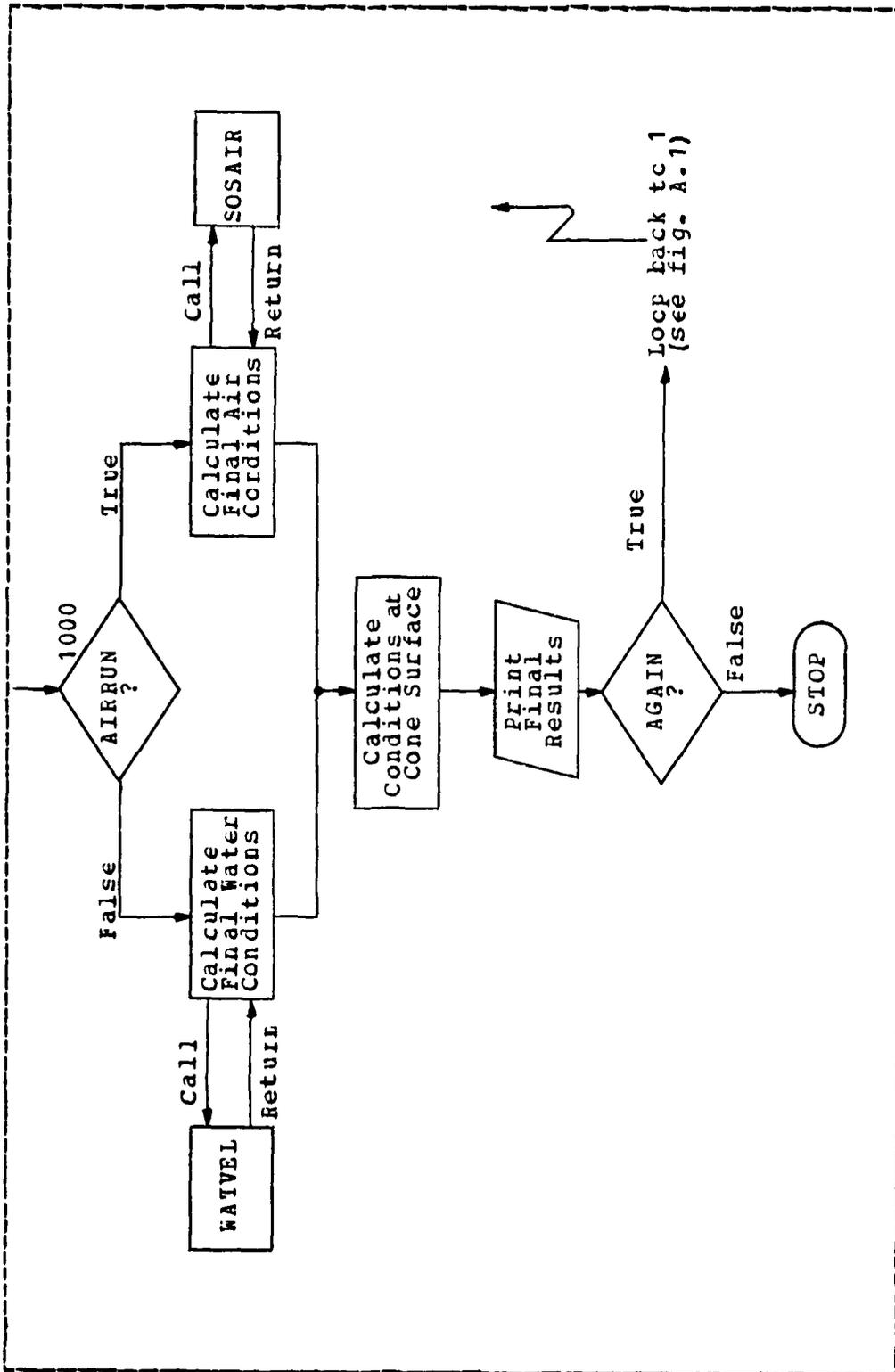


Figure A.7 Main Program Flowchart (ccnt'd.).

TABLE VI  
Main Program Variables

Variable	Variable's Meaning or Use	Type	Module	Units
AGAIN	Used to determine if another run of the program is desired	C	F.C.	-
AIRM	Molecular Weight of Air	R	A.S.	-
ALFB	Universal Gas Constant	R	A.S.	J/mol-K
AIRRUN	Indicates if run is for air or water	I	M.	-
ANGLEF	Deflection Angle after crossing the shock front	R	A.S.	Radians
AFISF	Test value to determine if run is for air or water	C	I.	-
ARGUM1	Holds argument from an air shock formula passed to SQRT function	R	A.S.	-
ARGUM2	Holds argument from an air shock formula passed to SQRT function	R	A.S.	-
AWBESP	User response to type of run option - an input variable	C	I.	-
C	An Array which holds the speed of sound at each point J	LP	M.	m/s
CCMFEI	Indicates if a complete or summary print is desired	I	M.	-
CCNANG	Cone Semi-vertex Angle	R	F.C.	Degrees
CSFF	Reference Speed of Sound on upstream side of shock front	LP	M.S.	m/s
LEN	Value of denominator in formula for Busemann Radius	LP	B.F.	-
LENSCS	Density at the Ccne Surface	R	F.C.	kg/m <sup>3</sup>
LENSY	Array which holds value of the Density at each point J	LP	M.	kg/m <sup>3</sup>
DENSY1	Density at Point 1 upstream	LP	M.	kg/m <sup>3</sup>
DENSY2	Density at Point 2 downstream	LP	M.	kg/m <sup>3</sup>
DRAGCC	Drag Coefficient for the Cone	R	F.C.	J/kg
ENERG	Energy term returned by EPV	LP	S.E.	J/kg
ENTH2	Enthalpy at Point 2 downstream	LP	S.E.	J/kg
ENTHCS	Enthalpy at Ccne surface - water case	R	F.C.	J/kg
ENTHLP	Array holding values of Enthalpy at each point J	LP	M.V.	J/kg

TABLE VII

Main Program variables (cont'd.)

Variable	Variable's Meaning or Use	Type	Module	Units
GAMMA	Ratio of Specific Heats (c/c)	R	A.S.	-
GAMMA1	Value of an expression of Gamma	R	A.S.	-
GAMMA2	Value of an expression of Gamma	R	A.S.	-
GAMMA3	Value of an expression of Gamma	R	A.S.	-
J	Counter used to indicate point J	I	M.	-
JAND1	Used to indicate point J+1	I	M.	-
MACH1	Mach number at Pcnt 1 upstream	R	M.	-
MACH2	Mach number at Pcnt 2 downstream	R	M.	-
MACHCS	Mach number on Cone Surface	R	F.C.	-
MAXVEL	Maximum Air Velocity	R	S.A.	m/s
MFACT	Kilobar Multiplication Factor	R	I.	-
NC	Input value for a Water Run	C	I.	-
NUM	Test value to determine if a complete print and/or another run is desired	LP	P.F.	m/s
CLEG	Value of the numerator in the calculation of Busemann Radius	R	M.	Degrees
OMEGA	Used to output value of angle	LP	M.	Radians
KREAF	Omega including values of angle	R	M.	Kilobars
PFESCS	Pressure at any point in kilobars	R	M.F.C.	Pascals
PFESF	Pressure at Cone Surface	C	I.	-
PFESS	User response to printout option - an input value	LP	M.	Pascals
PFESS1	Array including values of the Pressure at each Pcnt J	LP	M.	Pascals
PFESS2	Pressure at Point 1 upstream - an input value	LP	M.	Pascals
RADIUS	Pressure at Point 2 downstream	LP	M.	Pascals
RIFFI	Array including values of the Busemann radius at each point J	LP	M.	m/s
RIFVCL	Used to output the Deflection Angle Reference Specific Volume upstream	R	A.S.	Degrees
		LP	W.S.	m <sup>3</sup> /kg
			W.V.	

TABLE VIII  
Main Program Variables (cont'd.)

Variable	Variable's Meaning or Use	Type	Module	Units
SEKANG	Value of the Shock Angle	R	M.	Degrees
SHRAL	- an input value	R	A. S.	Radians
SFVCI2	Value of Shock Angle in Radians at Point 2	LP	W. S.	m <sup>3</sup> /kg
STEF	Specific Volume at Point 2	R	M.	Radians
STEG	Step size by which Omega is decreased	R	M.	Degrees
TEMP1	Used for output of streamline angle	LP	M.	Kelvin
TEMP2	Temperature at Point 1 upstream	LP	M.	Kelvin
TEMP3	- an input value	LP	M.	Kelvin
TEMP4	Temperature at Point 2 downstream	R	F. C.	Kelvin
TEMP5	Temperature at Cone Surface	LP	M.	Radians
TEMP6	Array holding values of the Streamline angles	LP	S. F.	J/kg
TCENT	Total Enthalpy	LP	W. V.	-
TVALUE	Test value used to determine if cone surface has been reached	F	M.	-
U	Array holding values of Busemann velocity components in x-direction at point J	LP	B. C.	m/s
V	Array holding values of Busemann velocity components in y-direction at point J	LP	B. C.	m/s
VEL	Array holding values of the velocity at each point J	LP	M.	m/s
VELFS	Free-stream Velocity	LP	M.	m/s
VNCFM1	Component of velocity Normal to the shock front upstream	LP	M.	m/s
VNCFM2	Component of velocity Normal to the shock front downstream	LP	M.	m/s
VTANG	tangential component of velocity at the shock front	LP	M.	m/s
WFSEF	test value for determining if the input is for air or water	C	I.	-
X	Array holding values of the x-coord. of the Busemann circle center	DP	B. C.	m/s

TABLE IX  
Main Program Variables (cont'd.)

Variable	Variable's Meaning or Use	Type	Module	Units
Y	Array holding values of the y-coord. of the Eusemann circle center	DP	B.C.	m/s
YES	Test value used to determine if complete print and/or another run is desired	C	I.	-

**LEGEND**

(1) TYPE is the type of data represented with designations as follows:

- R = Real Single Precision
- DP = Real Double Precision
- I = Integer (used for counters)
- L = Logical (i.e. True or False)
- C = Character (used to store character information - is an integer type in FORTRAN)

(2) MODULE is the primary, but not exclusive, location of use of the variable. The abbreviations used mean:

- A.S. = Airshock Calculation Section of Main Program
- B.C. = Eusemann Circle Center Calculation Section of Main Program
- B.R. = Eusemann Radius Calculation Section of Main Program
- I. = Input Section of Main Program
- F.C. = Final Output Section of Main Program
- M. = Variable Used Throughcut Main Program
- S.A. = Parameter to SOSAIR Module
- S.E. = Stagnation Enthalpy Calculation of Main Program
- W.S. = Parameter to WSHOCK Module
- W.V. = Parameter to WATVEL Module

**Note:** this legend applies to all Tables listing program variables.

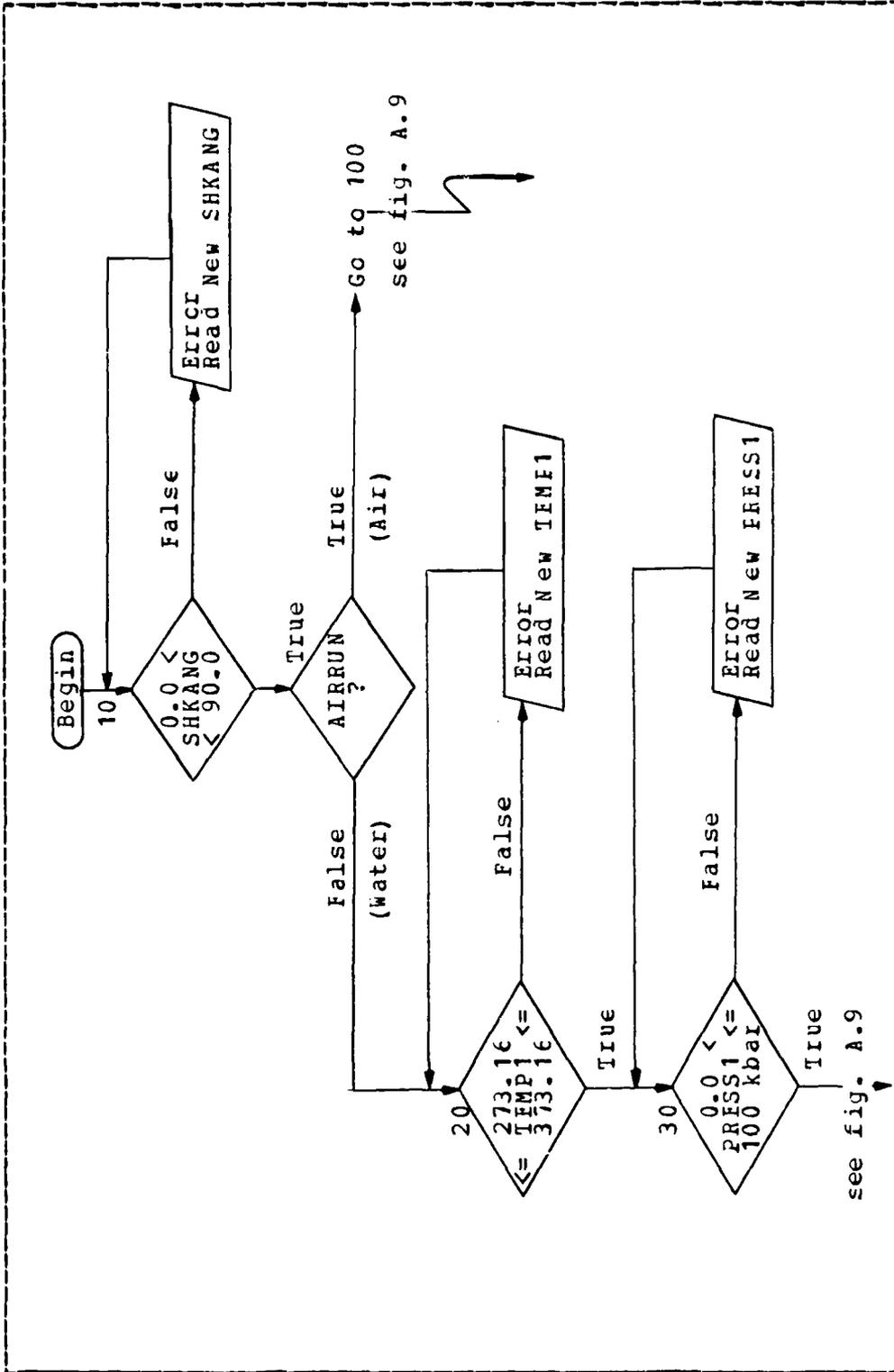


Figure A.8 Subroutine CHKINP Flowchart.

SAMPLE 2 - SUMMARY PRINTOUT (WATER CASE)

\*\*\* THIS RUN IS FOR - WATER - \*\*\*

INPUT VALUES WERE AS FOLLOWS:

UPSTREAM TEMPERATURE = 298.16 KELVIN  
 UPSTREAM PRESSURE = 101300.0 PASCALS  
 STICK ANGLE = 2.0 DEGREES  
 THE PRESSURE MULTIPLICATION FACTOR (MFACT) = 10.0

CONDITIONS CALCULATED FOR POINT 1 UPSTREAM:

REFERENCE SPECIFIC VOLUME = 0.00099599 M<sup>3</sup>/KG  
 DENSITY AT POINT 1 = 1005.8843 KG/M<sup>3</sup>  
 NORMAL COMPONENT OF VELOCITY = 2358.1095 M/S  
 TANGENTIAL COMPONENT OF VELOCITY = 5056.9817 M/S  
 UPSTREAM VELOCITY = 5579.7621 M/S  
 REFERENCE SPEED OF SOUND = 1494.8674 M/S  
 SPEED OF SOUND AT POINT 1 = 1485.0679 M/S  
 UPSTREAM MACH NUMBER = 3.7572

CONDITIONS CALCULATED FOR POINT 2 DOWNSTREAM

PRESSURE AT POINT 2 = 101300000.0 PASCALS  
 THIS PRESSURE IN KILOBARS = 10.13 KILOBARS  
 SPECIFIC VOLUME AT POINT 2 = 0.00081410 M<sup>3</sup>/KG  
 DENSITY AT POINT 2 = 1228.3457 KG/M<sup>3</sup>  
 ENTHALPY (FROM REF) AT POINT 2 = 91194.8388 J/KG  
 ENTHALPY AT POINT 2 = 915881.2343 J/KG  
 STAGNATION ENTHALPY = 15566873.3346 J/KG  
 NORMAL COMPONENT OF VELOCITY = 1931.0411 M/S  
 VELOCITY AT POINT 2 = 5413.1307 M/S  
 WATER VELOCITY AT POINT 2 = 427.0684 M/S  
 SPEED OF SOUND AT POINT 2 = 2779.0381 M/S  
 MACH NUMBER AT POINT 2 = 1.9478  
 X-COMPONENT OF EUSEMANN VELOCITY = 5399.2752 M/S  
 Y-COMPONENT OF EUSEMANN VELOCITY = 387.0554 M/S

KG/M3

1243.1082  
0.0630

DENSITY AT CONE SURFACE =  
LAG COEFFICIENT (CD) =

\* \* \* \* \*

OMEGA AT POINT 2 = 20.0 DEGREES  
 STREAMLINE ANGLE AT POINT 2 = 2.3018 DEGREES  
 EUSEMANN RADIUS AT POINT 2 = 1743.5342 M/S  
 X-COORDINATE OF EUSEMANN CENTER = 7452.2403 M/S  
 Y-COORDINATE OF EUSEMANN CENTER = 830.0185 M/S

\*\*\*\*\*

PCINT = 3

X-COORDINATE OF VELOCITY AT POINT 3 = 5803.6963 M/S  
 Y-COORDINATE OF VELOCITY AT POINT 3 = 262.3793 M/S  
 VELOCITY AT POINT 3 = 5809.6243 M/S

OMEGA AT POINT 3 = 19.0 DEGREES  
 STREAMLINE ANGLE AT POINT 3 = 2.5885 DEGREES

ENTHALPY AT POINT 3 = 522713.5817 J/KG  
 PRESSURE AT POINT 3 = 684748052.6934 PASCAL  
 TEMPERATURE IN KILOBARS = 6.8475 KILOBARS  
 DENSITY AT POINT 3 = 1183.3863 KG/M3  
 SPEED OF SOUND AT POINT 3 = 2567.7326 M/S

EUSEMANN RADIUS AT POINT 3 = 1362.8023 M/S  
 X-COORDINATE OF EUSEMANN CENTER = 7092.2512 M/S  
 Y-COORDINATE OF EUSEMANN CENTER = 706.0643 M/S

\*\*\*\*\*

CCNE HAS BEEN REACHED - FOLLOWING IS A SUMMARY OF RESULTS:

CCNE SEMI-VERTEX ANGLE = 7.2299 DEGREES

SECCOR ANGLE = 20.0 DEGREES  
 FLOWSTREAM MACH NUMBER = 3.9052  
 FLOWSTREAM VELOCITY = 5898.9120 M/S

VELOCITY AT CONE SURFACE = 5753.2645 M/S  
 SPEED OF SOUND AT CONE SURFACE = 2987.4511 M/S  
 MACH NUMBER AT CONE SURFACE = 1.9258

PRESSURE AT CONE SURFACE = 1092461570.0 PASCALS  
 PRESSURE AT CCNE SURFACE = 10.9246 KILOBARS  
 ENTHALPY AT CONE SURFACE = 848554.2500 J/KG

SAMPLE 1 - COMPLETE PRINTOUT (WATER CASE)

\*\*\* THIS RUN IS FOR - WATER - \*\*\*

INLET VALUES WERE AS FOLLOWS:

UPSTREAM TEMPERATURE = 323.16 KELVIN  
 UPSTREAM PRESSURE = 100000.0 PASCALS  
 SECCOR ANGLE = 20.0 DEGREES  
 THE PRESSURE MULTIPLICATION FACTOR (MFACT) = 5.0

CONDITIONS CALCULATED FOR POINT 1 UPSTREAM:

REFERENCE SPECIFIC VOLUME = 0.00100411 M<sup>3</sup>/KG  
 DENSITY AT PCINT 1 = 996.5139 KG/M<sup>3</sup>  
 NORMAL COMPONENT OF VELOCITY = 2017.5467 M/S  
 TANGENTIAL COMPONENT OF VELOCITY = 5543.1641 M/S  
 FREESTREAM VELOCITY = 5898.9120 M/S  
 REFERENCE SPEED OF SOUND = 1510.7939 M/S  
 SPEED OF SOUND AT PCINT 1 = 1510.5239 M/S  
 FREESTREAM MACH NUMBER = 3.9052

CONDITIONS CALCULATED FOR POINT 2 DOWNSTREAM

PRESSURE AT POINT 2 = 50000000.0 PASCALS  
 TOTAL PRESSURE IN KILCBARS = 5.0 KILCBARS  
 SPECIFIC VOLUME AT POINT 2 = 0.00087980 M<sup>3</sup>/KG  
 DENSITY AT POINT 2 = 1136.6191 KG/M<sup>3</sup>  
 HEAT (FROM EV) AT PCINT 2 = 30923.2949 J/KG  
 ENTHALPY AT POINT 2 = 470824.3162 J/KG  
 STAGNATION ENTHALPY = 17358580.6611 J/KG  
 NORMAL COMPONENT OF VELOCITY = 1768.8540 M/S  
 VELOCITY AT POINT 2 = 5818.5490 M/S  
 WATER VELOCITY AT PCINT 2 = 248.6927 M/S  
 SPEED OF SOUND AT PCINT 2 = 2268.3064 M/S  
 MACH NUMBER AT PCINT 2 = 2.5652  
 X-COMPONENT OF EUSEMANN VELOCITY = 5813.8541 M/S  
 Y-COMPONENT OF EUSEMANN VELOCITY = 233.6947 M/S

APPENDIX B  
SAMPLE PRINTOUTS

This appendix presents copies of the output from various runs of the computer program listed in Appendix C. The various samples illustrate the output from the following printouts:

- (1) A Complete Print of a Water Run
- (2) A Summary Print of a Water Run
- (3) A Complete Print of an Air Run, and
- (4) A Summary Print of an Air Run

(Note: in order to reduce the volume of the thesis, only a portion of the output from the complete printouts is included.)

TABLE XI  
Subroutine WATVEL Variables

Variable	Variable's Meaning or Use	Type	Units
E	Constant used in Tait equation of EPV	DP	Pascals
COUNT	Program counter	I	-
LENS3	An iteration density term	DP	kg/m <sup>3</sup>
LENS4	An iteration density term	DP	kg/m <sup>3</sup>
LENS5	An iteration density term	DP	kg/m <sup>3</sup>
LENSJ	Density at Point J	DP	kg/m <sup>3</sup>
LENSJ1	Density at Point J+1	DP	kg/m <sup>3</sup>
LENTHAL	Enthalpy iteration term	DP	J/kg
ENTHE1	An enthalpy iteration term	DP	J/kg
ENTHE2	An enthalpy iteration term	DP	J/kg
ENTHE3	Constant used in the Tait equation	DP	J/kg
N 1	An expression involving N	R	-
PFPSJ1	Pressure at Point J+1	R	-
PFPSJ2	Pressure at Point J	DP	Pascals
RVCI	Speed of Sound at Point 1	DP	m/S
SCCSJ1	Reference Specific Volume at Point 1	DP	m <sup>3</sup> /kg
VELJ	Stagnation Enthalpy Point J+1	DP	J/kg
VELJ1	Speed of Sound at Point J+1	DP	m/S
WLENS	Velocity at Point J	DP	m/S
WLENERG	Density at Point J+1	DP	kg/m <sup>3</sup>
WLENTF	Energy term returned by EPV	DP	J/kg
WLENTF1	Enthalpy at Point J+1	DP	J/kg
WLENTF2	Enthalpy at Point J	DP	J/kg
WLENTF3	Pressure at Point J+1	DP	Pascals
WLENTF4	Specific Volume at Point J+1	DP	m <sup>3</sup> /kg
WLENTF5	Temperature at Point J+1	DP	Kelvin
WLENTF6	Temperature at Point J	DP	-
WLENTF7	Test value	R	-

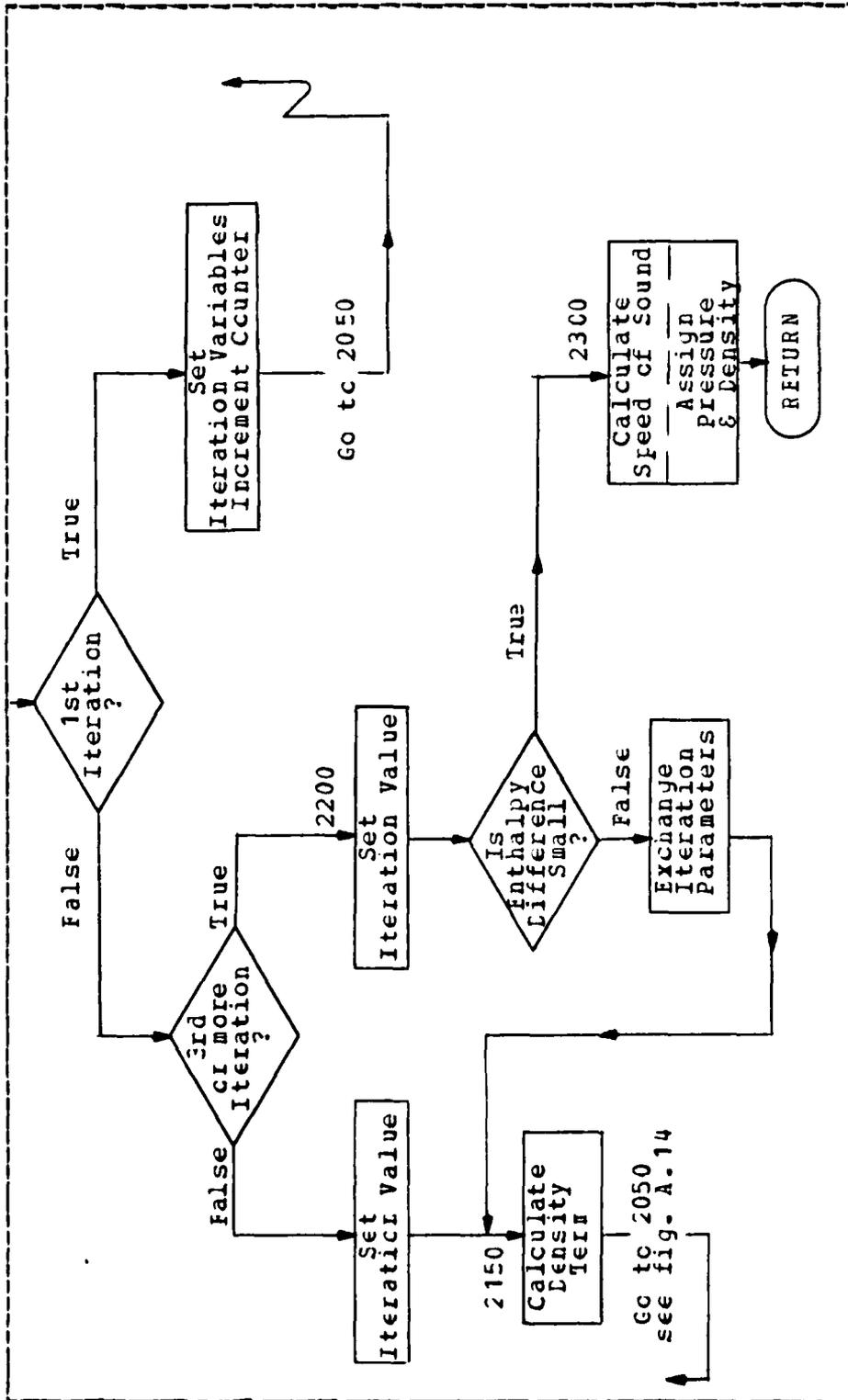


Figure A-15 Subroutine WATVEL Flowchart (cont'd.).

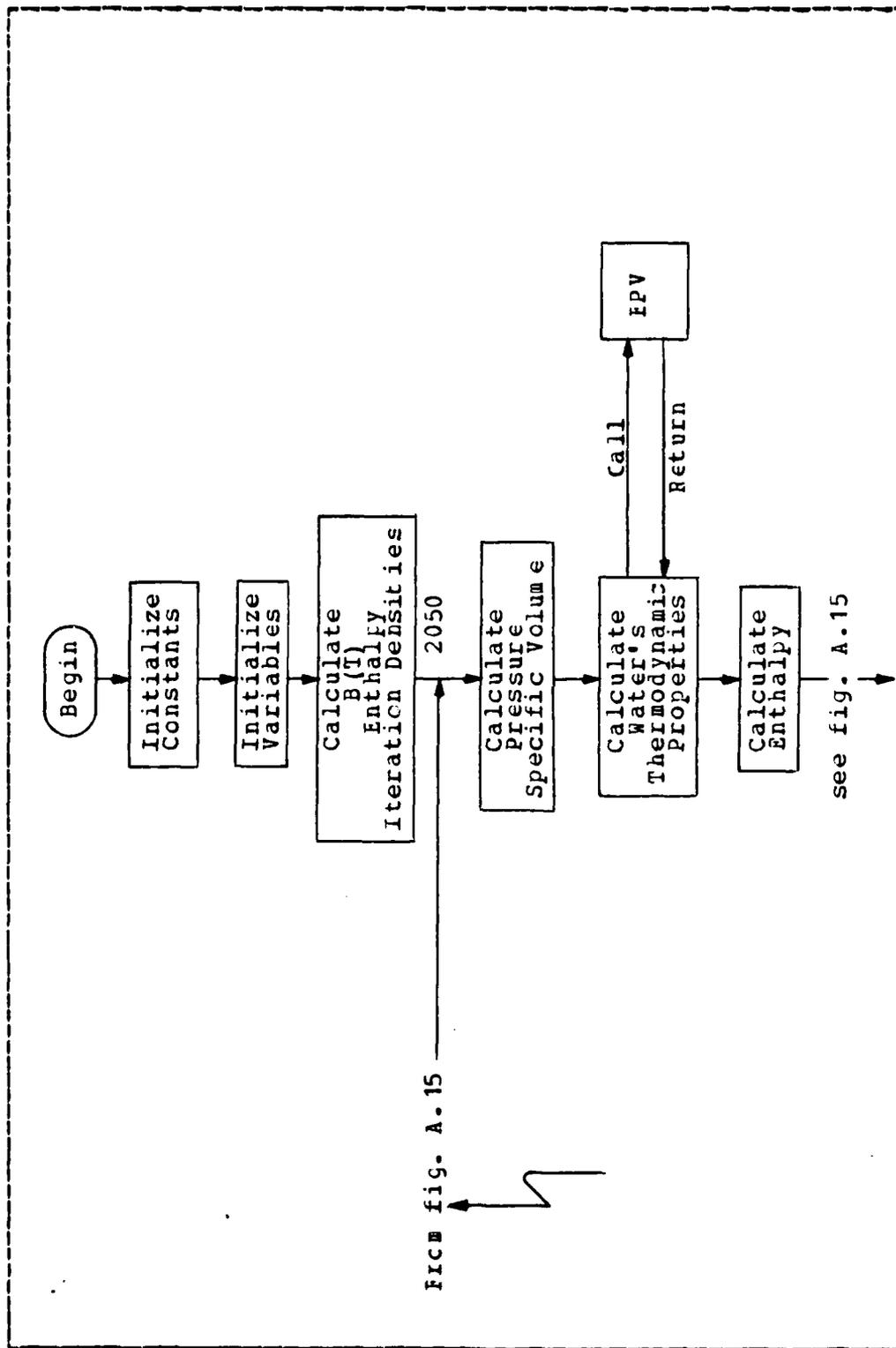


Figure A.14 Subroutine WATVEL Flowchart.

TABLE X  
Subroutine WSHOCK Variables

Variable	Variable's Meaning or Use	Type	Units
EF7	A constant used in the Tait equation	DP	Pascals
EFV	A constant used in the Tait equation	DP	Pascals
CCOUNT1	Program counter	I	m2/s2
CCOUNT2	Program counter	I	-
DENS1	Density at Point 1	DP	kg/m3
DENS2	Density at Point 2	DP	kg/m3
ISEVCI	Iteration Specific Volume	DP	kg/m3
ITEMF1	Iteration temperature	DP	m3/kg
ITEMF2	Iteration temperature	DP	Kelvin
ITEMF3	Iteration temperature	DP	Kelvin
N1	A constant used in the Tait equation	R	-
N2	An expression involving N	R	-
N3	An expression involving N	R	-
N4	An expression involving N	R	-
EFSCS	Reference Speed of Sound	R	-
SICFF1	Specific Volume vs. temperature curve	DP	m/s
SICFF2	Specific Volume vs. temperature curve	DP	J/Kelvin
SCCS1	Speed of Sound at Point 1	DP	m/s
SCCS2	Speed of Sound at Point 2	DP	m/s
SFVCI0	Specific Volume at Point 1	DP	m3/kg
SFVCI1	Specific Volume at Point 2	DP	m3/kg
SVRAT	Ratio of specific volumes	DP	-
VHI1	Velocity at Point 1	DP	m/s
VHI2	Velocity at Point 2	DP	m/s
WSEN13	Iteration enthalpy term	DP	J/kg
WSEN14	Iteration enthalpy term	DP	J/kg
WSEN15	Iteration enthalpy term	DP	J/kg
WSEN16	Iteration enthalpy term	DP	J/kg
WSEN17	Iteration enthalpy term	DP	J/kg
WSEN18	Iteration enthalpy term	DP	J/kg
WSEN19	Iteration enthalpy term	DP	J/kg
WSEFES1	Pressure at Pcnt	DP	Pascals
WTEFES1	A test value	R	-
WTEFES2	Temperature at Pcnt 1	R	Kelvin
WVVEL1	Velocity of the Shock Front	DP	m/s

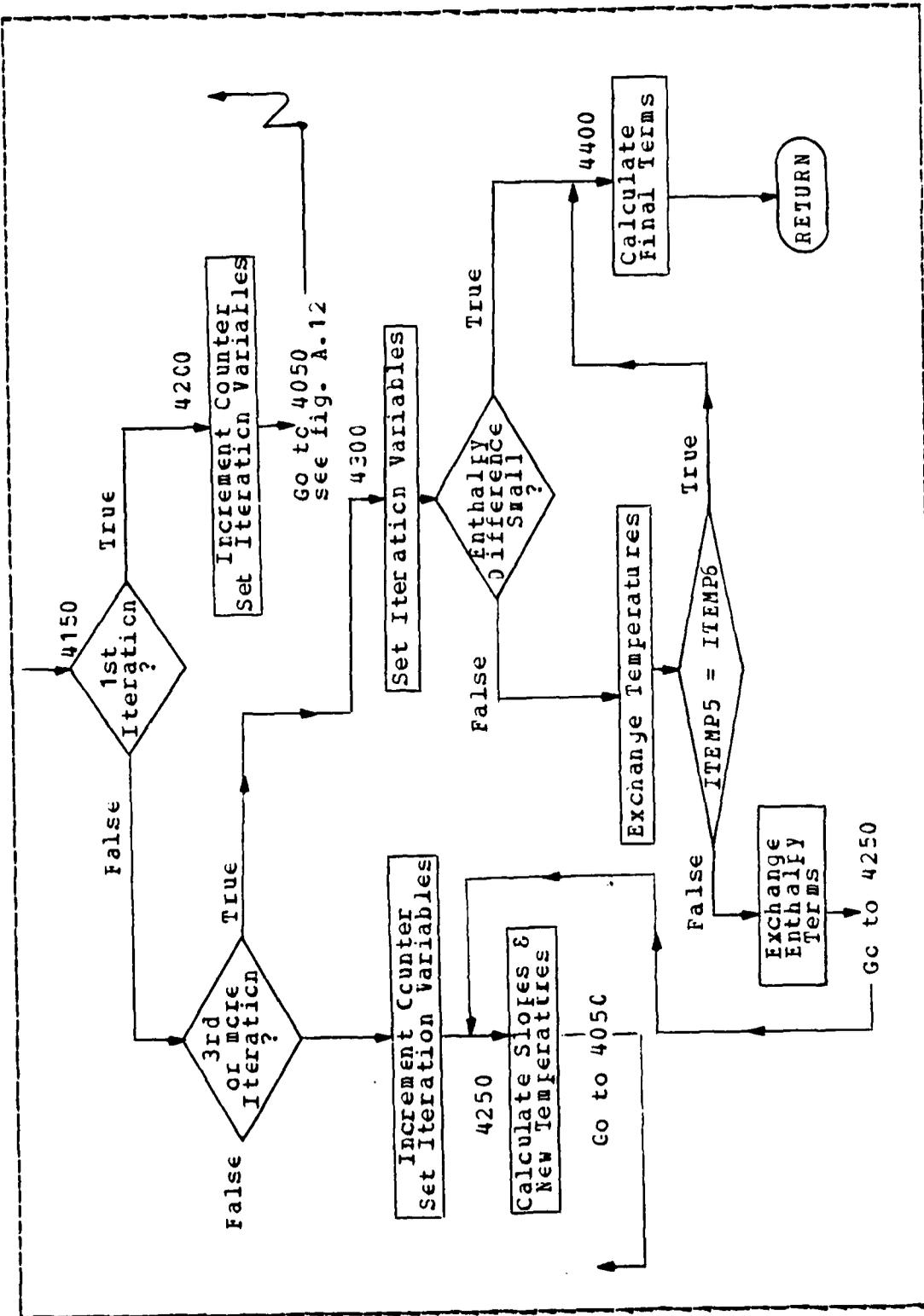


Figure A.13 Subroutine WSHOCK Flowchart (cont'd.).

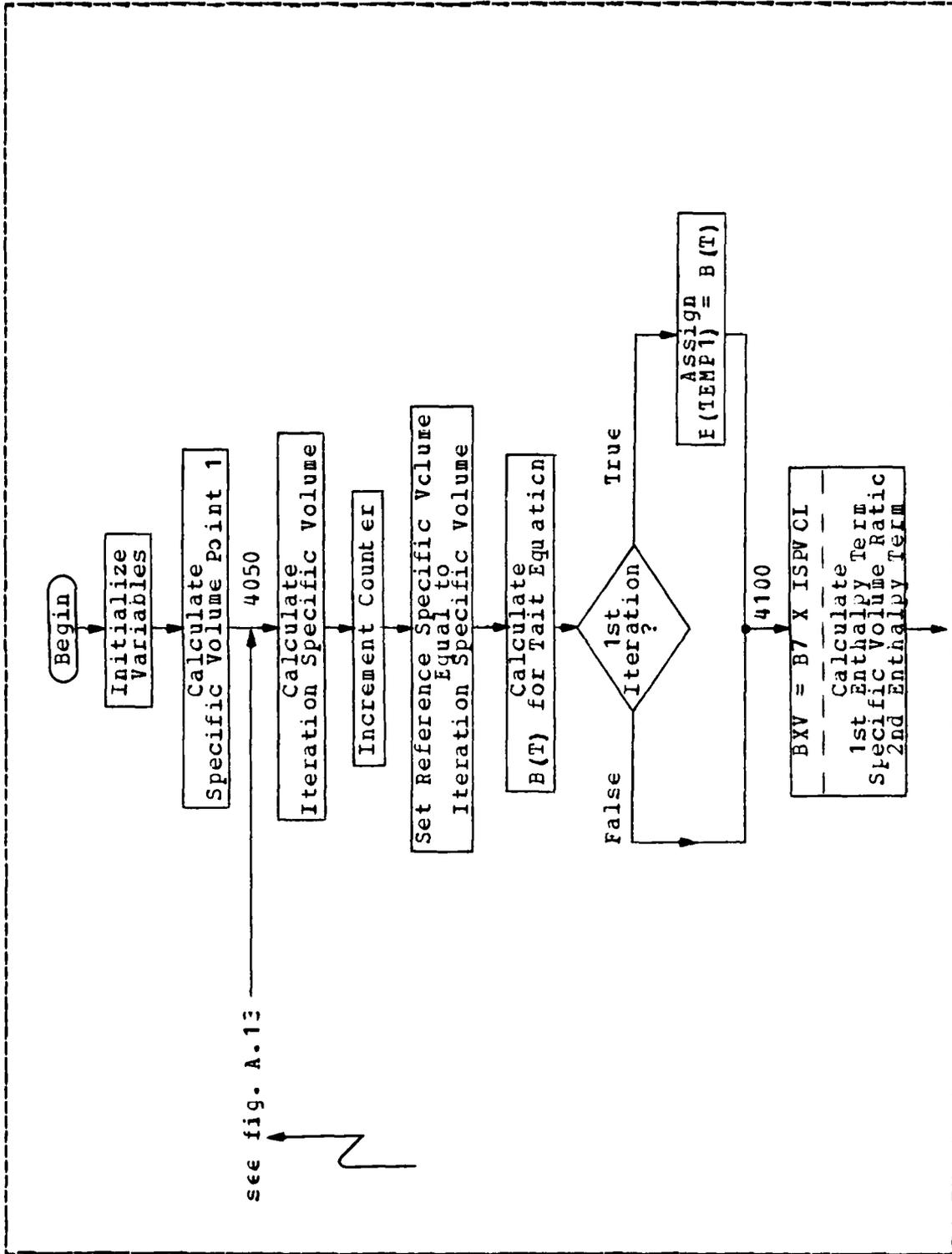


Figure A.12 Subroutine WSHOCK Flowchart.

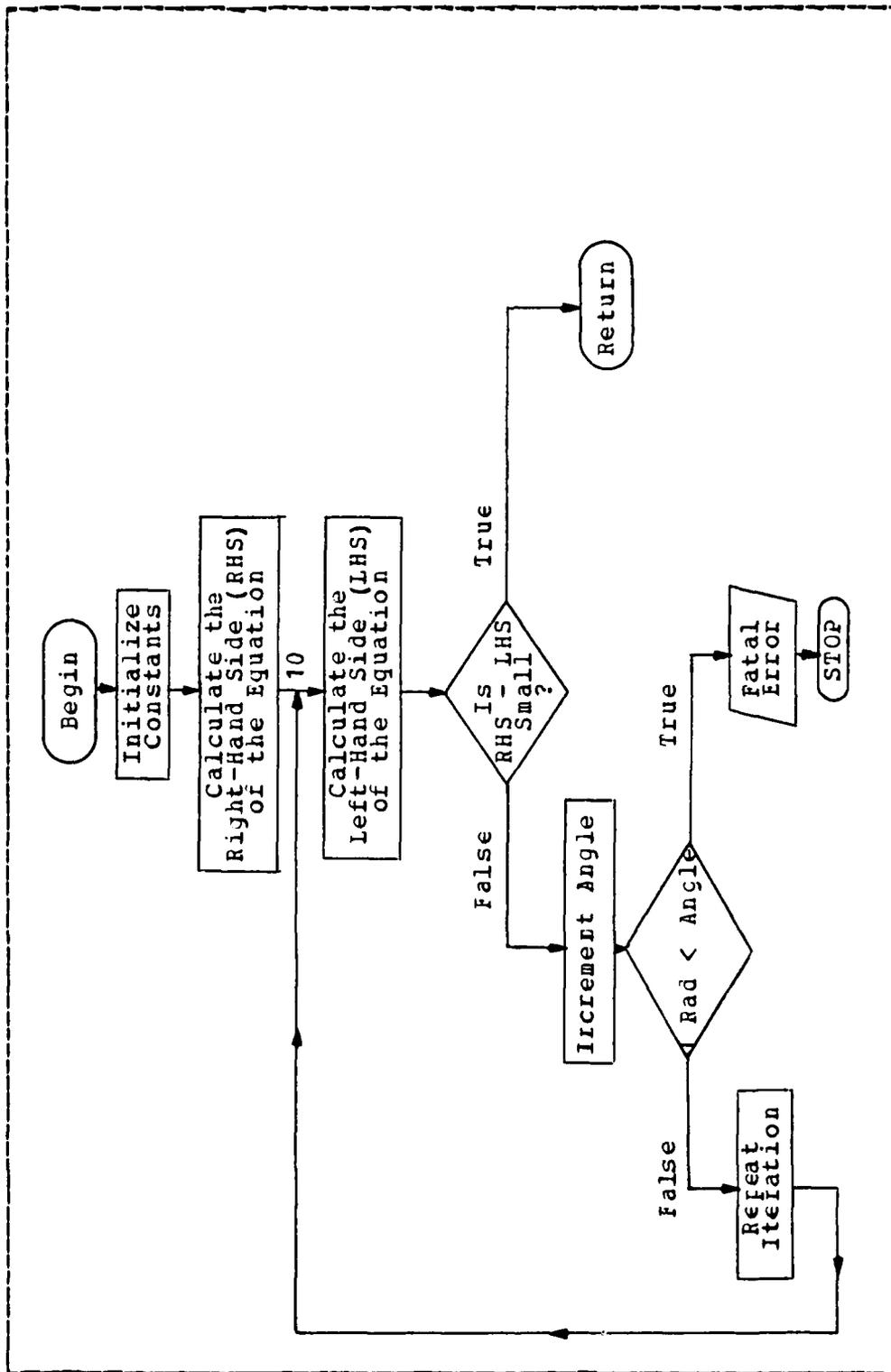
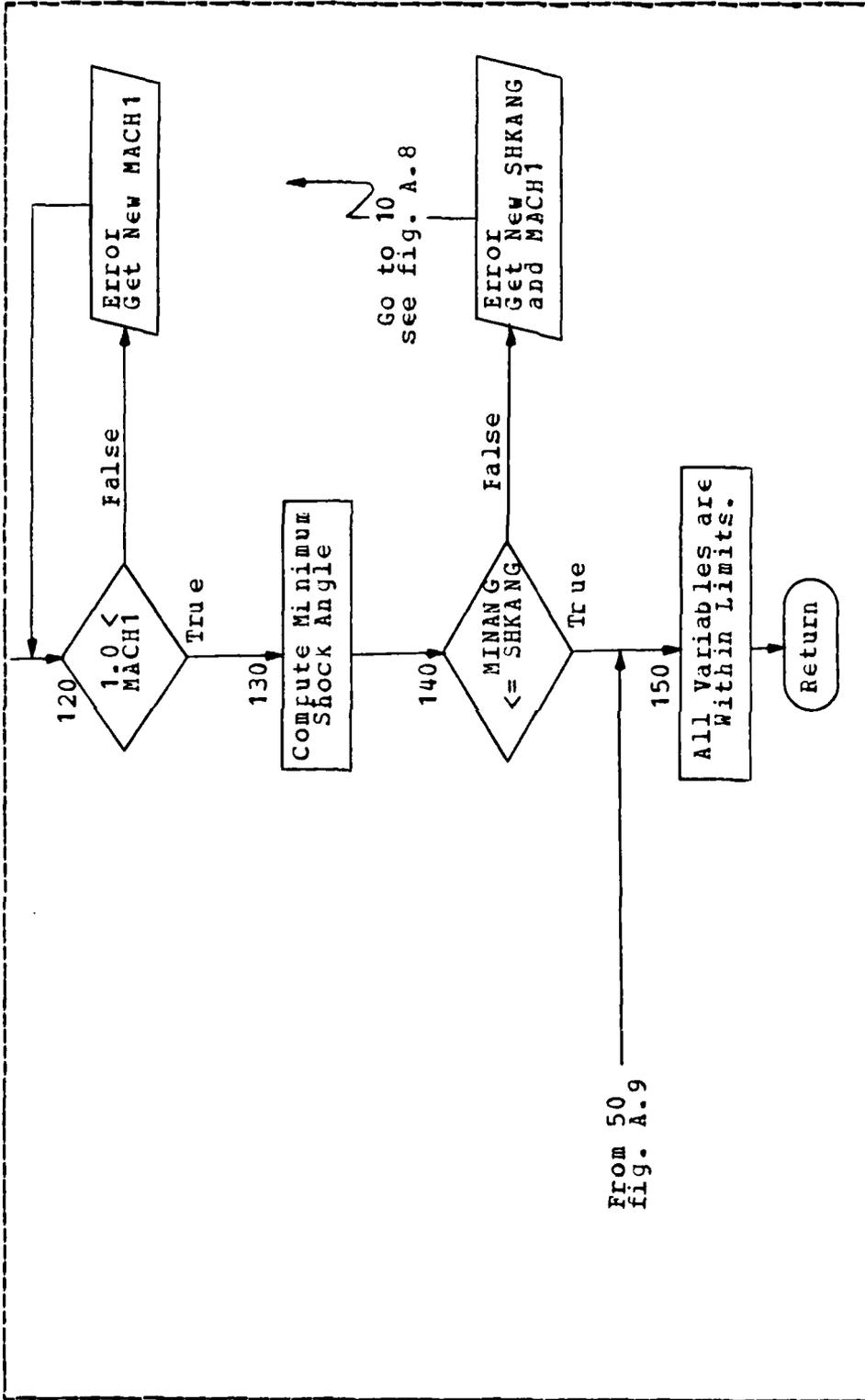


Figure A.11 Subroutine DEFANG Flowchart.



From 50  
fig. A.9

Figure A.10 Subroutine CHKINP Flowchart (cont'd.).

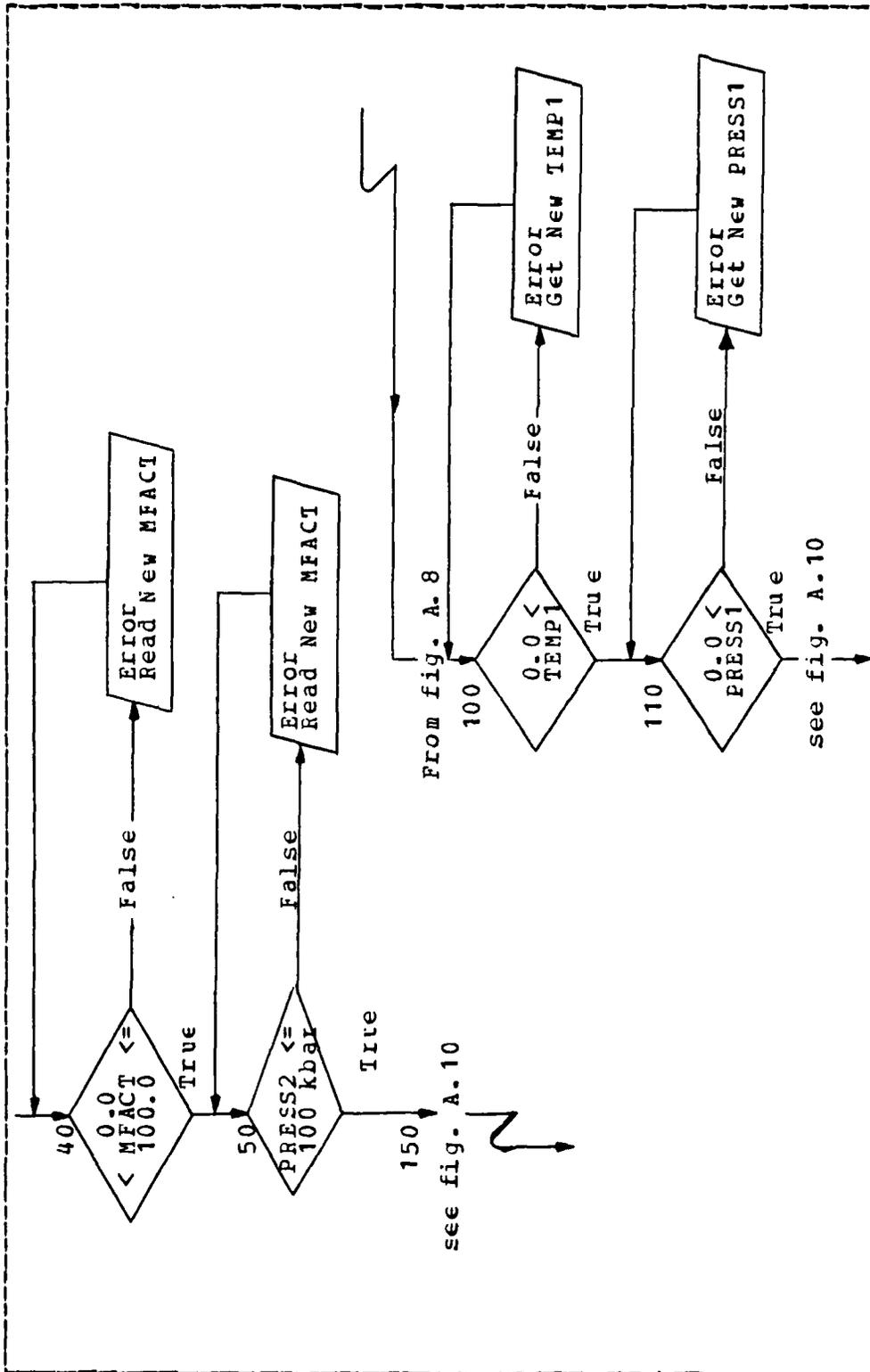


Figure A.9 Subroutine CHKINP Flowchart (cont'd.).

ANGLE AT POINT 2 = 25.0000 DEGREES  
STREAMLINE ANGLE AT POINT 2 = 4.7003 DEGREES

\*\*\*\*\*

CCONE HAS BEEN REACHED - FOLLOWING IS A SUMMARY OF RESULTS:

CCONE SEMI-VERTEX ANGLE = 10.5470 DEGREES

SPOCK ANGLE = 25.0 DEGREES

FREESTREAM MACH NUMBER = 3.7572

FREESTREAM VELOCITY = 5579.7621 M/S

VELOCITY AT CONE SURFACE = 5316.8138 M/S

SPEED OF SOUND AT CONE SURFACE = 3534.3125 M/S

MACH NUMBER AT CONE SURFACE = 1.5041

PRESSURE AT CONE SURFACE = 1806928000.0 PASCALS

PRESSURE AT CCONE SURFACE = 18.0093 KILOBARS

ENTHALPY AT CONE SURFACE = 1432619.0 J/KG

DENSITY AT CONE SURFACE = 1328.2969 KG/M3

DRAG COEFFICIENT (CD) = 0.1150

\*\*\*\*\*

SAMPLE 3 - COMPLETE PRINTOUT (AIR CASE)

\*\*\* THIS RUN IS FOR - AIR - \*\*\*

INLET VALUES WERE AS FOLLOWS:

UPSTREAM TEMPERATURE = 300.16 KELVIN  
 UPSTREAM PRESSURE = 100000.0 PASCALS  
 SWIRL ANGLE = 30.0 DEGREES  
 UPSTREAM MACH NUMBER = 3.0

CALCULATED INITIAL CONDITIONS WERE:

DEFLECTION ANGLE = 12.7735 DEGREES  
 DENSITY AT PCINT 1 = 1.1610 KG/M3  
 FREESTREAM VELOCITY = 1041.7729 M/S  
 MAXIMUM AIR VELOCITY = 1299.3167 M/S  
 MACH NUMBER AT PCINT 2 = 2.3673  
 VELOCITY AT POINT 2 = 944.5709 M/S 902.1986 M/S  
 TANGENTIAL COMPONENT OF VELOCITY =  
 X-COMPOONENT OF EISEMANN VELOCITY = 921.1945 M/S  
 Y-COMPOONENT OF EISEMANN VELOCITY = 208.8416 M/S

OMEGA AT POINT 2 = 30.0 DEGREES  
 STREAMLINE ANGLE AT POINT 2 = 12.7735 DEGREES  
 EISEMANN RADIUS AT PCINT 2 = 821.4434 M/S  
 X-COORDINATE OF EISEMANN CENTER = 1632.5854 M/S  
 Y-COORDINATE OF EISEMANN CENTER = 619.5633 M/S

\*\*\*\*\*

PCINT = 3

X-COMPOONENT OF VELOCITY AT PCINT 3 = 914.1348 M/S  
 Y-COMPOONENT OF VELOCITY AT PCINT 3 = 221.3196 M/S  
 VELOCITY AT POINT 3 = 940.5449 M/S

OMEGA AT POINT 3 = 25.0 DEGREES  
 STREAMLINE ANGLE AT POINT 3 = 13.6099 DEGREES  
 SPEED OF SOUND AT PCINT 3 = 400.8987 M/S

ELSEMANN RADIUS AT PCINT 3 = 745.5216 M/S  
X-COORDINATE OF ELSEMANN CENTER = 1566.1827 M/S  
Y-COORDINATE OF ELSEMANN CENTER = 582.7557 M/S

\* \* \* \* \*

CCNE HAS BEEN REACHED - FOLLOWING IS A SUMMARY OF RESULTS:

CCNE SEMI-VERTEX ANGLE = 20.3996 DEGREES

SHOCK ANGLE = 50.0 DEGREES

FREESTREAM MACH NUMBER = 3.0

FREESTREAM VELOCITY = 1041.7729 M/S

VELOCITY AT CONE SURFACE = 925.8313 M/S

WALL CF SOUND AT CCNE SURFACE = 407.6909 M/S

MACH NUMBER AT CCNE SURFACE = 2.2709

PRESSURE AT CONE SURFACE = 285849.8120 FASCALS

TEMPERATURE AT CCNE SURFACE = 413.7244 KELVIN

DENSITY AT CONE SURFACE = 2.4077 KG/M3

DRAG COEFFICIENT (CD) = 0.2950

\* \* \* \* \*

SAMPLE 4 - SUMMARY PRINTOUT (AIR CASE)

\*\*\* THIS RUN IS FOR - AIR - \*\*\*

INPUT VALUES WERE AS FOLLOWS:

UPSTREAM TEMPERATURE = 373.16 KELVIN  
 UPSTREAM PRESSURE = 202600.0 PASCALS  
 SHOCK ANGLE = 45.0 DEGREES  
 FREESTREAM MACH NUMBER = 3.25

CALCULATED INITIAL CONDITIONS WERE:

DEFLECTION ANGLE = 27.0239 DEGREES  
 DENSITY AT POINT 1 = 1.8920 KG/M3  
 FREESTREAM VELOCITY = 1258.3641 M/S  
 MAXIMUM AIR VELOCITY = 1527.4334 M/S  
 MACH NUMBER AT POINT 2 = 1.7324  
 VELOCITY AT POINT 2 = 935.4593 M/S  
 TANGENTIAL COMPONENT OF VELOCITY = 889.7951 M/S  
 X-COMPONENT OF EISEMANN VELOCITY = 833.3229 M/S  
 Y-COMPONENT OF EISEMANN VELOCITY = 425.0375 M/S  
 ANGLE AT POINT 2 = 45.0 DEGREES  
 STREAMLINE ANGLE AT POINT 2 = 27.0239 DEGREES

\*\*\*\*\*

CCNE HAS BEEN REACHED - FOLLOWING IS A SUMMARY OF RESULTS:

CCNE SEMI-VERTEX ANGLE = 35.3971 DEGREES  
 SHOCK ANGLE = 45.0 DEGREES  
 FREESTREAM MACH NUMBER = 3.25  
 FREESTREAM VELOCITY = 1258.3641 M/S  
 VELOCITY AT CONE SURFACE = 914.1494 M/S  
 SPEED OF SOUND AT CONE SURFACE = 547.2446 M/S  
 MACH NUMBER AT CCNE SURFACE = 1.6705  
 PRESSURE AT CONE SURFACE = 1333402.0 PASCALS  
 TEMPERATURE AT CCNE SURFACE = 745.4390 KELVIN

DENSITY AT CONE SURFACE = 6.2334 KG/M3  
DRAG COEFFICIENT (CD) = 0.7549

\* \* \* \* \*

APPENDIX C  
PROGRAM LISTING

This appendix contains a complete listing of the fully documented main program named CONEFLOW and all of its functions and subroutines. The functions and subroutines included in this appendix are as follows:

- (1) Function DTOR
- (2) Function RTOI
- (3) Function DEFANG
- (4) Function SOSAIR
- (5) Subroutine CEMINP
- (6) Subroutine WSECCK
- (7) Subroutine WAITVEL
- (8) Subroutine EFV



THE PROGRAM REQUIRES THAT THE USER SUPPLY THE FOLLOWING INPUT:

- (1) AN INDICATION OF WHETHER A COMETE CR A SUMMARY PRINTOUT IS DESIRED. THIS IS RECEIVED BY THE PROGRAM THROUGH VARIABLE -PFESP--.
- (2) AN INDICATION OF WHETHER THIS RUN IS FOR AIR CR FOR WATER. THIS IS RECEIVED BY THE PROGRAM THROUGH VARIABLE -AMRESP--.
- (3) A VALUE FOR THE TEMPERATURE UPSTREAM OF THE SHOCK FRONT. THIS IS RECEIVED BY THE PROGRAM THROUGH VARIABLE -TEMP1--.
- (4) A VALUE FOR THE PRESSURE DOWNSTREAM OF THE SHOCK FRONT. THIS IS RECEIVED BY THE PROGRAM THROUGH VARIABLE -PRESS1--.
- (5) A VALUE FOR THE ANGLE THE SHOCK FRONT MAKES WITH THE HORIZONTAL. THIS IS RECEIVED BY THE PROGRAM THROUGH VARIABLE -SHKANG--.
- (6) IF THE RUN IS FOR AIR, A VALUE FOR THE MACH NUMBER UPSTREAM OF THE SHOCK FRONT. THIS IS RECEIVED BY THE PROGRAM THROUGH VARIABLE -MACH1--.
- (7) IF THE RUN IS FOR WATER, A VALUE FOR THE PRESSURE MULTIPLICATION FACTOR USED TO CONVERT TO A PRESSURE AT POINT 2 DOWNSTREAM OF THE SHOCK FRONT, WHICH IS PRESS1 TIMES MFACT KILOBARS. THIS IS RECEIVED BY THE PROGRAM THROUGH VARIABLE -MFACT--.

LIST OF PROCEDURES: IN ADDITION TO STANDARD FUNCTIONS AND/OR PROCEDURES (SUCH AS SIN AND COS), FCUNT IN LIBRARIES AND THIS PROGRAM UTILIZES THE FOLLOWING FUNCTIONS AND SUBROUTINES:

DICE - A FUNCTION  
R1CD - A FUNCTION  
DEFANG - A FUNCTION  
SOSAIR - A FUNCTION  
CHKINF - A SUBROUTINE  
WSHCCK - A SUBROUTINE  
WATVEL - A SUBROUTINE

EPV - A SUBROUTINE

EACH OF THESE FUNCTIONS OR SUBROUTINES, EXCEPT EPV, IS EXPLAINED IN DETAIL IN THE SUBROUTINE EPV IS INCLUDED TO SHOW HOW IT WAS TRANSLATED TO THE FORTRAN LANGUAGE AND FOR USE BY THIS PROGRAM. IT IS COPIED BY PERMISSION OF THE AUTHOR FROM NAVAL POSTGRADUATE SCHOOL REPORT NPS 67-001 BY A. E. FUHS. THE SUBROUTINE EPV IS ADEQUATELY DESCRIBED AND FLOWCHARTED IN THAT WORK AND NEED NOT BE REPEATED HERE.

IMPLEMENTATION NOTES: THERE ARE NO IMPLEMENTATION SPECIFIC FEATURES OF THE NAVAL POSTGRADUATE SCHOOL COMPUTER SYSTEM INCLUDED IN THIS PROGRAM.

\*\*\*\*\*

MAIN PROGRAM CONEFLOW

DECLARE MAIN PROGRAM VARIABLES

REAL AIRM, AIRR, ANGLEF, ARGUM1, ARGUM2, CCNANG, DENSCS, DRAGCO,  
 ENTHCS, GAMMA, GAMMA1, GAMMA2, GAMMA3, MACH1, MACH2, MACHCS,  
 MAXVEL, MFACI, CDEG, PKBAR, PRESCS, RDEFI, SHKANG, SHRAD,  
 STEP, IDEG, TEMCS, TVALUE, WATSPD

DOUBLE PRECISION C, CREF, DEN, DENSITY, DENSY1, DENSY2, ENERG,  
 ENTH2, ENTHL, ENTHLP, NUM, OMEGA, PRESS, PRESS1, PRESS2,  
 RADIUS, REFVOL, SP, VOL2, TEMF1, TEMF2, THETA,  
 TCTENT, U, V, VEL, VELS, VNCM1, VNORE2, VIANG,  
 X, Y

LOGICAL AIRFUN, CCMPFT

THE FOLLOWING INTEGER VARIABLES ARE USED AS COUNTERS

INTEGER J, JAND1

THE FOLLOWING INTEGER VARIABLES ARE USED TO STORE CHARACTERS

INTEGER AGAIN, ARESF, AHRESP, NO, PRESP, WRESF, YES

ESTABLISH ARRAY SIZES

```

C      DIMENSION C(90), DENSITY(90), ENTHLP(90), U(90), V(90), WESP(90), X(90),
C      *          RADIUS(90), THETA(90), OMEGA(90), VEL(90), PRESS(90),
C      *          Y(90)
C
C      INITIALIZE ARRAYS TO PREVENT UNWANTED RESULTS
C
C      DATA C/90*C.00000000/DENSITY/90*0.00000000/ENTHLP/90*C.00000000/
C      *      DATA OMEGA/50*0.00000000/PRESS/90*0.00000000/
C      *      DATA RADIUS/90*0.00000000/THETA/90*0.00000000/U/90*0.00000000/
C      *      DATA V/90*C.00000000/VEL/90*0.00000000/X/90*C.00000000/
C      *      DATA Y/90*C.00000000/
C
C      DECLARE CONSTANTS
C
C      DATA GAMMA/1.4/AIRM/28.974/AIRR/8314.41/
C      *      DATA ARESP/AIR/WRESP/WATE/YES/YES/NC/NO '/'
C
C      INITIALIZE KEY VARIABLES
C
C      BEGIN (PROGRAM CONEFLOW)
C
C      J = 2
C      JAND1 = J + 1
C      GAMMA1 = (GAMMA - 1.0)/2.0
C      GAMMA2 = (GAMMA + 1.0)/(GAMMA * 2.0)
C      GAMMA3 = GAMMA/(GAMMA - 1.0)
C      STEP = DTOF(1.00000000)
C      TVALUE = 0.000001
C
C      DETERMINE IF A COMPLETE OR A SUMMARY PRINTOUT IS DESIRED
C
C      *      WRITE(7,9915)
C      *      READ(7,8902) PRESE
C      *      IF ((PRESE.EQ. YES) .OR. (PRESE.EQ. NO)) GO TO 20
C      *      WRITE(7,9916)
C      *      GO TO 10
C      *      IF (PRESE.EQ. YES) GO TO 25
C      *      IF (PRESE.EQ. FALSE)
C      *          GO TO 30
C      *      CCMPRT = .TRUE.
C
C      DETERMINE IF THIS IS AN -AIR- OR A -WAIT- RUN
C
C      *      WRITE(7,9920)
C      *      READ(7,8902) AMRESE

```

```

35 IF (AWRESE .EQ. ARESE) .OR. (AWRESP .EQ. WRESP)) GC TO 35
   WRITE(7,9921)
   GO TO 30
36 IF (AWRESP .EQ. ARESE) GO TO 40
   AIRRUN = .FALSE.
   GO TO 45
37 AIRRUN = .TRUE.

C
C GET VALUES FOR USER SPECIFIED VARIABLES
C
38 WRITE(7,99CC) TEMP1
   READ(7,890C) TEMP1
C
39 WRITE(7,99C5) PRESS1
   READ(7,890C) PRESS1
C
40 WRITE(7,991C) SHKANG
   READ(7,8901) SHKANG
C
41 IF (AIRRUN) GC TO 50
   WRITE(7,9930)
   READ(7,89C1) MFACT
   GO TO 55
42 WRITE(7,9925) MACH1
   READ(7,8901) MACH1
C
43 VERIFY USER INEUT IS VALID BY CALLING SUBROUTINE CHKINE
C
44 CALL CHKINE(TEMP1,PRESS1,SHKANG,AIRRUN,MACH1,MFACT)
C
45 CMEGA(2) = LIGR(SHKANG)
   IF (AIRRUN) GO TO 100
C
46 CALCULATE INITIAL WATER CONDITIONS AT POINTS 1 AND 2
C
47 IF (PRESS1 .GT. 0.0) GC TO 65
   PRESS2 = MFACT * 10000000.0
   GO TO 70
48 PRESS2 = PRESS1 * MFACT * 1000.0
   CALL WSHOCK(TEMP1,PRESS2,VNCRM1,VNORM2,C(1),SEVCL2,CREF,C(2),
   *
   * VLFPS = VNCFM1/DENST(2) REFVOL, DENSY1, WATSPD)
   VTANG = VEIIFS * DCCS(OMEGA(J))
   MACH1 = VEIIFS/C(1)
   VEL(2) = DSCRI(VTANG**2 + VNORM2**2)

```

```

MACH2 = VEL(2)/C(2)
CALCULATE STAGNATION ENTHALPY FOR CALCULATED WATER CONDITIONS
CALL EPV(TEMP1,SPVCL2,PRESS2,ENERG)
ENTH2 = ENFG + PRESS2 * SPVOL2
TCTENT = ENTH2 + VEL(2)**2/2.0
GC TO 200
END - CALCULATION OF INITIAL WATER CONDITIONS
CALCULATE INITIAL AIR CCNDITIONS AT POINTS 1 AND 2
100  ANGDEF = DEFANG(SHKANG,MACH1)
      KEFL = RTCL(ANGDEF)
      SHRAD = DTCE(SHKANG)
      CALCULATIONS FOR FCINT 1
      DENS1 = PRESS1 * AIRM/(AIRR * TEMP1)
      C(1) = DSQRT(GAMMA * AIRR * TEMP1/AIRM)
      VELFS = C(1) * MACH1
      CALCULATIONS FOR FCINT 2
      ARGUM1 = (5.+MACH1**2*SIN(SHRAD)**2)/(7.*MACH1**2*SIN(SHRAD)**2
      *-1.0)
      ARGUM2 = (5.+MACH1**2*SIN(SHRAD)**2) * (7.*MACH1**2*SIN(SHRAD)**2
      *-1.)/(36.* MACH1**2*SIN(SHRAD)**2)
      MACH2 = SQRT(ARGUM1)/SIN(SHRAL-ANGDEF)
      PRESS2 = PRESS1*(7.*MACH1**2*SIN(SHRAD)**2-1.)/6.)
      TEMP2 = TEMP1 * ARGUM2
      C(2) = C(1) * SQRT(ARGUM2)
      VEL(2) = C(2) * MACH2
      DENS2 = GAMMA2 * PRESS2/C(2)**2
      VNORM2 = MACH2 * C(2) * SIN(SHRAD-ANGDEF)
      VTANG = VNORM2/TAN(SHRAD-ANGDEF)
      CALCULATE MAXIMUM AIR VELOCITY
      MAXVEL = DSQRT(C(2)**2/GAMMA1 + VEL(2)**2)
      END - CALCULATION OF INITIAL AIR CONDITIONS
      CALCULATE BUSEMANN VELOCITY COMPONENTS AT FCINT 2
200  U(J) = VNORM2 * DSIN(CMEGA(J)) + VTANG * DCOS(OMEGA(J))

```

```

C
C      V(J) = -(VNCRM2 * DCCS(OMEGA(J)) + VTANG * ISIN(OMEGA(J))
C
C      CALCULATE THE STREAMLINE ANGLE AT POINT 2
C      THETA(J) = LATAN(V(J)/U(J))
C
C      PRINT INITIAL CONDITIONS AS CALCULATED ABOVE FOR AIR OR WATER
C300  IF (AIRRUN) GO TO 310
C      WRITE(6,9003)
C      GO TC 32C
C10
C20  WRITE(6,9005) TEMF1,PRESS1,SHKANG
C
C      IF (AIRRUN) GO TO 330
C
C      PRINT INITIAL WATER CONDITIONS
C
C      PKBAR = PRESS2/10000000.0
C      WRITE(6,9007) MFACT
C      WRITE(6,9013) REFVCL, DENSITY
C      WRITE(6,9015) VNCRM1, VTANG, VELFS
C      WRITE(6,9017) CREF, C(1)
C      WRITE(6,9020) MACH1
C      WRITE(6,9025) PRESS2, PKBAR, SPVOI2, DENSITY(2)
C      WRITE(6,9027) ENRG, ENTH2, TOTENT
C      WRITE(6,9030) VNCRM2, VEL(2), MATSPD
C      WRITE(6,9032) C(2), MACH2
C      WRITE(6,9034)
C      GO
C
C      PRINT INITIAL CONDITIONS FOR AIR
C330  WRITE(6,9042) MACH1
C      WRITE(6,9044) RDEF1, DENSITY1
C      WRITE(6,9046) VELFS, MAXVEL
C      WRITE(6,9048) MACH2, VEL(2), VTANG
C
C      PRINT COMMON INITIAL OUTPUT VARIABLES
C340  CLEG = RTOL(OMEGA(2))
C      TLEG = RIOT(THETA(2))
C      WRITE(6,9044) U(2), V(2)
C      WRITE(6,9044) CDEG, IDEG
C

```

```

C BEGIN - ITERATION SEQUENCE TO DETERMINE CONE SEMI-VERTEX ANGLE
C
C CALCULATE RADII OF BUSEMANN APPLE CURVE FOR AIR OF WATER
C
C 400 NUM = VEL(J) * DSIN(THETA(J)) / DSIN(OMEGA(J))
C DEN = 1.0 - (VEL(J)/C(J)) **2 * DSIN(OMEGA(J)) - THETA(J) **2
C RADIUS(J) = NUM/DEN
C
C CALCULATE CENTER OF EUSEMANN CIRCLE
C
C X(J) = U(J) + RADIUS(J) * DCOS(OMEGA(J))
C Y(J) = V(J) + RADIUS(J) * DSIN(OMEGA(J))
C
C DECREMENT THE ANGLE OMEGA
C
C OMEGA(J+1) = OMEGA(J) - STEP
C
C CALCULATE BUSEMANN VELOCITY COMPONENTS AT THE NEXT POINT
C
C 450 U(J+1) = X(J) - RADIUS(J) * DCOS(OMEGA(J+1))
C V(J+1) = Y(J) - RADIUS(J) * DSIN(OMEGA(J+1))
C
C CALCULATE THE VELOCITY AT THE NEXT POINT
C
C VEL(J+1) = ISQRT(V(J+1)**2 + U(J+1)**2)
C
C CALCULATE THE STREAMLINE ANGLE AT THE NEXT POINT
C
C THETA(J+1) = DATAN(V(J+1)/U(J+1))
C
C PRINT RESULTS IF THIS IS A COMPLETE PRINTOUT
C
C IF (.NOT. CCHEAT) GO TO 475
C WRITE(6,9C45) J, RADIUS(J)
C WRITE(6,9C50) X(J), Y(J)
C WRITE(6,9C75) JAND1, U(J+1), V(J+1)
C WRITE(6,9C80) JAND1, ODEG, JAND1, IDEG
C ODEG = RTCD(OMEGA(J+1))
C IDEG = RTCD(THETA(J+1))
C WRITE(6,9C70) JAND1, ODEG, JAND1, IDEG
C
C TEST TO DETERMINE IF CONE SURFACE HAS BEEN REACHED
C
C 475 IF (IVALUE -GT. DABS(THETA(J+1) - OMEGA(J+1))) GO TO 1000
C
C ENSURE STREAMLINE ANGLE IS LESS THAN THE ANGLE OMEGA

```

AD-A151 102

A COMPUTER PROGRAM TO CALCULATE THE SUPERSONIC FLOW  
OVER A SOLID CONE IN AIR OR WATER(U) NAVAL POSTGRADUATE  
SCHOOL MONTEREY CA P W HUGHES JUN 84 NPS-67-84-007

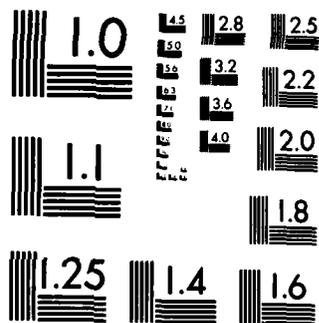
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MICROCOPY RESOLUTION TEST CHART  
NATIONAL BUREAU OF STANDARDS-1963-A

```

C
C
C
480 IF (THETA(J+1) - II. OMEGA (J+1)) GO TO 500
C OMEGA (J+1) = LAIAN (Y(J)/X(J))
C GO TO 450
C
C
500 IF (AIRRUN) GC TO 600
C
C CALCULATE THE WATER CCNLTIONS AT THE NEXT PCINT
C
505 * CALL WATVEL (J,TEMP1,CREF,TOTENT,REFVOL,DENSTY(J),VEL(J+1),C(J+1),
C ENTHLF(J+1),PRESS(J+1),DENSTY(J+1))
C
510 GC TO 650
C
C CALCULATE THE SPEED CF SCUND IN AIR AT THE NEXT FOINT
C
600 C(J+1) = SCCAIR(MAXVEL,VEL(J+1))
C
C EFFINT RESULTS IF A COMPIETE PRINTOUT IS DESIRED
C
650 IF (-NCT. CCMPRT) GO TO 700
C IF (AIRRUN) GO TO 675
C PKBAR = PRESS (J+1)/100000000.0
C WRITE (6,9145) JAND1,ENTHLP(J+1)
C WRITE (6,9150) JAND1,PRESS(J+1),PKBAR
C WRITE (6,9155) JAND1,DENSTY(J+1)
C WRITE (6,9160) JAND1,C(J+1)
C
C INCREMENT COUNTERS AND REPEAT CALCULATIONS
C
700 J = J + 1
C JAND1 = J + 1
C GC TO 400
C
C
C THE CONE SURFACE HAS BEEN REACHEL - CALCULATE FINAL CCNDITIONS
C
1000 IF (AIRRUN) GC TO 1100
C
C CALCULATE FINAL WATER CCNDITIONS AT THE CONE SURFACE
C
* CALL WATVEL (J,TEMP1,CREF,TOTENT,REFVOL,DENSTY(J),VEL(J+1),C(J+1),
C ENTHLF(J+1),PRESS(J+1),DENSTY(J+1))
C
PKBAR = PRESS (J+1)
PKBAR = PKBAR/100000000.0
DINSCS = DENSTY(J+1)
ENTHCS = ENTHLP(J+1)
GC TO 1200

```

```

C
C
C 1100 CALCULATE FINAL AIR CCNDITIONS AT THE CONE SURFACE
      C(J+1) = SCSAIR(MAXVEL,VEL(J+1))
      TEMPCS = C(J+1)**2 * AIRM/(GAMMA * AIRR)
      PRPSCS = PRESS2 * (TEMPCS/TEMP2)**GAMMA3
      DENSCS = GAMMA * PRESCS/C(J+1)**2
C
C 1200 CALCULATE CCNE CONDITIGNS COMMON TO AIR CR WATER CASES
      DRAGCO = 2.C * (PRESCS - PRESS1)/(DENSY1 * VEIIFS**2)
      MACHCS = VEL(J+1)/C(J+1)
      CCNANG = RICI(OMEGA{J+1})
C
C DISPLAY FINAL RESULTS
      WRITE(6,9075)
      WRITE(6,9085)
      WRITE(6,9100)
      WRITE(6,9085)
      WRITE(6,9090)
      WRITE(6,9095)
      WRITE(6,9100)
      WRITE(6,9110)
      WRITE(6,9120)
      WRITE(6,9125)
      WRITE(6,9135)
      WRITE(6,9075)
      CONANG
      SHKANG
      MACH1
      VELFS
      VEL(J+1)
      C(J+1)
      MACHCS
      PRESCS
C
C IF (AIRRUN) GO TO 1225
      WRITE(6,9106) PKRAF
      WRITE(6,9111) ENTHCS
      GO TO 1250
C 1225 WRITE(6,9110) TEMFCS
C 1250 WRITE(6,9115) DENSCS
      WRITE(6,9125) DRAGCC
      WRITE(6,9075)
C
C DETERMINE IF ANOTHER RUN IS DESIRED
C 1275 WRITE(7,9940)
      READ(7,9902) AGAIN
C
C VERIFY CORRECT RESPONSE
C IF ((AGAIN.EQ.YES) .OR. (AGAIN.EQ.NO)) GC TC 1300
      WRITE(7,9916)
      GO TO 1275

```

```

C RESECNSE IS CCONNECT, TAKE ACTION DIRECTED
C
1300 IF (AGAIN EQ NO) GC TO 1400
    WRITE(6,5185)
    WRITE(7,5945)
    GO TO 1
C
1400 WRITE(7,9935)
    STOP
C
C END PROGRAM CONEFLCW
C
FCMAT STATEMENTS
C
THE FOLLOWING FORMAT STATEMENTS ARE FOR READING DATA
C
8500 FCRMAT (F12.5)
8501 FCRMAT (F8.5)
8502 FCRMAT (A4)
C
C THE FOLLOWING FORMAT STATEMENTS ARE FOR WRITING TO THE OUTPUT FILE
C
9000 FCRMAT (*** THIS RUN IS FOR - WATER - ***/) INPUT VALUES WERE
    AS FOLLOWS:
9003 FCRMAT (*** THIS RUN IS FOR - AIR - ***/) INPUT VALUES WERE A
    S FOLLOWS:
9005 FCRMAT (UPSTREAM TEMPERATURE = F10.5, KELVIN/UPSTREAM PRESSU
    RE = F20.1, SHOCK ANGLE = F8.4, DEGREES)
9007 FCRMAT (PRESSURE MULTIPLICATION FACTOR (MFA, CT) = F8.5/)
9013 FCRMAT (CONDITIONS CALCULATED FOR POINT 1 UPSTREAM:
9015 FCRMAT (REFERENCE SPECIFIC VOLUME = F20.10, M3/KG/ DENSITY AT
    POINT 1 = F20.10, KG/M3)
9017 FCRMAT (NORMAL COMPONENT OF VELOCITY = F20.10, M/S/ TANGENTIA
    L COMPONENT OF VELOCITY = F20.10, M/S/ FREESTREAM VELOCITY =
    F20.10, M/S/)
9020 FCRMAT (REFERENCE SPEED OF SOUND = F20.10, M/S/ SPEED OF SOUN
    D AT POINT 1 = F20.10, M/S/)
9023 FCRMAT (FREESTREAM MACH NUMBER = F20.10/)
9025 FCRMAT (FREESTREAM MACH CALCULATED FOR PCIN 2 DOWNSTREAM/)
9027 FCRMAT (PRESSURE AT FCIN 2 = F25.10, PASCALS/ THIS PRESSUR
    E IN KILOBARS = F10.5, KILOBARS/ SPECIFIC VOLUME AT PCIN 2 =
    F15.10, M3/KG/ DENSITY AT POINT 2 = F20.10, KG/M3)
9030 FCRMAT (EBAR (FROM FEV) AT POINT 2 = F20.10, J/KG/ ENTHALPY
    AT POINT 2 = F20.10, J/KG/ STAGNATION ENTHALPY = F20.10,
    J/KG)
9032 FCRMAT (NORMAL COMPONENT OF VELOCITY = F20.10, M/S/ VELOCITY
    AT POINT 2 = F20.10, M/S/ WATER VELOCITY AT POINT 2 = F10.4

```



```

***
THE FOLLOWING FORMAT STATEMENTS ARE FOR WRITING TO THE TERMINAL
FCRMT(ENTER THE FREESTREAM TEMPERATURE IN DEGREES KEVIN: ')
FCRMT(ENTER THE FREESTREAM PRESSURE IN PASCALS: ')
FCRMT(ENTER THE SHOCK ANGLE IN DEGREES (E.G. 30.0): ')
FCRMT(DO YOU WANT A COMPLETE PRINTOUT? (YES OR NO): ')
** YES OR NO ** PLEASE TRY AGAIN.
FCRMT(IS THIS RUN FOR AIR OR WATER? (ANSWER = AIR OR WATER): ')
FCRMT(ENTER YOUR ANSWER TO THE QUESTION ABOVE MUST BE EITHER
** AIR ** OR ** WATER ** PLEASE TRY AGAIN.
FCRMT(ENTER THE FREESTREAM MACH NUMBER (E.G. 3.0): ')
FCRMT(ENTER THE KIICBAR MULTIPLICATION FACTOR: (NOTE: MFACT MUST
BE GREATER THAN 0.0 BUT LESS THAN 100.1): ')
FCRMT(PROGRAM EXECUTION IS COMPLETED. YOU CAN REVIEW THE OUTPUT
FROM THIS FUNCTION BY USING THE * BRCMSE * OR * XELIT * COMMANDS.
FCRMT(WOULD YOU LIKE TO MAKE ANOTHER RUN OF THIS PROGRAM? (YES OR
NO): ')
FCRMT(VERY WELL. EXECUTION OF NEW RUN BEGINS.)/
END

```

```

***
FUNCTION DTOR
THIS FUNCTION CONVERTS AN ANGULAR MEASUREMENT, RECEIVED AS AN INPUT
PARAMETER TO THE FUNCTION, FROM DEGREES TO THE CORRESPONDING MEASURE
IN RADIANS.
** GLOBAL VARIABLES ARE AFFECTED BY THIS FUNCTION.
**
FUNCTION DTOR(ANGLED)
DOUBLE PRECISION VARIABLES
BEGIN (FUNCTION DTOR)
DTOR = ANGLED * 3.1415926535/180.0

```

```

C C C      RETURN
C C C      END (FUNCTION DTOB)
C C C      END
C * * * * *
C * * * * *      FUNCTION RTOD
C * * * * *      THIS FUNCTION CONVERTS AN ANGULAR MEASUREMENT RECEIVED AS AN INPUT
C * * * * *      PARAMETER TO THE FUNCTION, FROM RADIAN TO THE CORRESPONDING MEASURE
C * * * * *      IN DEGREES.
C * * * * *      NO GLOBAL VARIABLES ARE AFFECTED BY THIS FUNCTION.
C * * * * *
C * * * * *      FUNCTION RTDI(ANGLE)
C * * * * *      DECLARE FUNCTION VARIABLES
C * * * * *      DOUBLE PRECISION ANGLE
C * * * * *      BEGIN (FUNCTION RTOD)
C * * * * *          RTOD = ANGLE * 180.0/3.1415926535
C * * * * *          RETURN
C * * * * *      END (FUNCTION RTOD)
C * * * * *      END
C * * * * *
C * * * * *      FUNCTION DEFANG
C * * * * *      THIS FUNCTION CALCULATES THE DEFLECTION ANGLE WHICH RESULTS WHEN A
C * * * * *      STREAMLINE PASSES ACROSS AN OBLIQUE SHOCK FRONT IN AIR. THE FUNCTION
C * * * * *      RECEIVES THE FOLLOWING MAIN PROGRAM VARIABLES AS INPUT PARAMETERS:
C * * * * *      SHKANG AND MACH1
C * * * * *      THE DEFLECTION ANGLE IS COMPUTED AND IS RETURNED IN RADIAN MEASURE.

```







```

30 IF ((PANS -GE. 0.0) -AND. (PANS -LE. 1C0CC0C0000.0)) GC TC 40
   WRITE(7,9020)
   WRITE(7,9025)
   READ(7,8955) FANS
   GO TC 30
40 IF ((MFANS -GT. 0.0) -AND. (MFANS -LE. 100.0)) GO TC 50
   WRITE(7,9030)
   WRITE(7,9035)
   READ(7,8950) MFANS
   GO TC 40
50 ENSURE PRESSURE AT PCINT 2 IS LESS THAN 100.0 KEARS
   IF ((PANS*MFANS) -LE. 10J0000000.0) GO TC 60
   WRITE(7,9040)
   WRITE(7,9035)
   READ(7,8950) MFANS
   GO TC 40
60 ALL INPUT VALUES FOR A -WATER- RUN ARE WITHIN LIMITS
   WRITE(7,9045)
   GO TO 260
70 FOLLOWING TESTS ARE CONDUCTED FOR AN -AIR- RUN
100 IF (TANS -GT. 0.0) GC TO 110
   WRITE(7,9050)
   WRITE(7,9055)
   READ(7,8955) TANS
   GO TO 110
110 IF (FANS -GT. 0.0) GC TO 120
   WRITE(7,9060)
   WRITE(7,9065)
   READ(7,8955) FANS
   GO TO 110
120 IF (MANS -GT. 1.0) GO TO 130
   WRITE(7,9070)
   WRITE(7,9075)
   READ(7,8950) MANS
130 ENSURE SHOCK ANGLE IS ABOVE THE MINIMUM FOR THE MACH NUMBER GIVEN
   MNMANG = ARSIN(1.0/MANS)

```

```

RSANS = IJOR(SANS)
IF (RSANS .GE. HINANG) GO TO 140
WRITE(7,9085)
WRITE(7,9005) SANS
READ(7,9050) SANS
WRITE(7,9075) MANS
WRITE(7,8950) MANS
GO TO 16

```

```

C ALL VALUES ARE WITHIN LIMITS FOR AN -AIR- RUN

```

```

C 140 WRITE(7,9080)

```

```

C 200 RETURN

```

```

C END (SUBROUTINE CHKINP)

```

```

C FCMAT STATEMENTS

```

```

C HEAD FORMAT STATEMENTS

```

```

C 8550 FCRMAT(F8.5)

```

```

C 8555 FCRMAT(F12.5)

```

```

C WRITE FORMAT STATEMENTS

```

```

C 9000 FCRMAT(/,ERROR - THE SHOCK ANGLE MUST BE GREATER THAN 0.0 BUT LESS

```

```

* THAN 90.0),PLEASE RE-ENTER THE SHOCK ANGLE IN DEGREES (E.G. 30.0):')

```

```

FCRMAT(/,ERROR - THE TEMPERATURE MUST BE GREATER THAN 273.16 KELVIN

```

```

* N BUT LESS THAN /,373.16 KELVIN.))

```

```

FCRMAT(/,PLEASE RE-ENTER THE FREESTREAM TEMPERATURE IN DEGREES KELV

```

```

* IN:))

```

```

FCRMAT(/,ERROR - THE PRESSURE MUST BE GREATER THAN 0.0 PASCALS /,

```

```

* BUT LESS THAN 101300000.0 PASCALS.

```

```

FCRMAT(/,PLEASE RE-ENTER THE FREE-STREAM PRESSURE IN PASCALS:'))

```

```

FCRMAT(/,ERROR - THE KILOBAR MUST BE GREATER

```

```

* THAN 0.0 /,BUT LESS THAN 100.1,

```

```

FCRMAT(/,PLEASE RE-ENTER THE KILOBAR MULTIPLICATION FACTOR:'))

```

```

FCRMAT(/,THE PRESSURE AT POINT 2 WILL BE GREATER THAN THE MULTI

```

```

* S PROGRAM. /,THE COMBINATION OF FREESTREAM PRESSURE TIMES THE MULTI

```

```

* PPLICATION FACTOR MUST BE LESS THAN 100.1 KILOBARS. YOU MUST EN

```

```

* TER A MULTIPLE FACTOR WHICH WILL BE GREATER THAN THE PRESSURE AT POI

```

```

* NT 2 WITHIN LIMITS. /))

```

```

FCRMAT(/,ALL INPUT VALUES FOR THE CALCULATIONS IN * WATER * ARE WI

```

```

* THIN LIMITS. /,EXECUTION OF YOUR PROGRAM WILL NOW BEGIN:'))

```

```

FCRMAT(/,ERROR - THE FREESTREAM AIR TEMPERATURE MUST BE GREATER TH

```



```

* * VFL2, WSENT2, WSENT3, WSENT4, WSENT5, WSENT6, WSENT16
* * WSENT7, WSENT8, WSENT9, WSERIES, WSTEP, WTEMP1
C
C INTEGER CCUNT1, CCUNT2
C
C DECLARE CCONSTANTS
C DATA N/7.15/, WSIEST/0.00001/
C
C INITIALIZE KEY VARIABLES
C
C ITEMF2 = 200.0
C N1 = (N + 1.0) / (N - 1.0)
C N2 = (N - 1.0) / N
C N3 = -(1 - C/N)
C N4 = (N - 1.0) / (2.0 * N)
C CCUNT1 = 0
C CCUNT2 = C
C
C BEGIN (SUBROUTINE WSHOCK)
C
C ITEMF1 = WTEMP1 - 273.16
C WTEMP = ITEMF1
C SPVOL1 = (0.99415 + 0.0002929 * (WTEMP1 - 298.16) + 0.000003241
* * (WTEMP1 - 258.16)**2) / 1000.0
C
C ISPVOL = (0.99415 + 0.0002929 * (WTEMP - 25.0) + 0.000003241 *
* * (WTEMP - 25.0)**2) / 1000.0
C COUNT2 = COUNT2 + 1
C SPVOL0 = ISPVOL
C B7 = 101.00000.0 * (3.134 - 0.00165 * (WTEMP - 55.0) -
* * (WTEMP - 55.0)**2 + 0.000000532 * (WTEMP -
* * 55.0)**3)
C
C IF (COUNT2 .GT. 1) GC TO 4100
C B6 = B7
C
C 4100 BXV = B7 * ISPVCL
C WSENT9 = 3.9644 * (WTEMP - WTEMP1 + 273.16) + 0.000312 * (WTEMP
* * 2 - (WTEMP1 - 273.16)**2)
C SVRAT = (SPRES/E7 + 1.0)
C WSENT8 = EXV/2.0 * (SVRAT - N1 * (SVRAT**2 - 1.0) - SVRAT**3 -
* * (ISPVOL - SPVOL1) * (SVRAT - 1.0) / ISPVOL) / 1000.0
C
C IF (COUNT1 .GT. 0) GC TO 4200
C WSENT2 = WSENT8
C WSENT3 = WSENT9
C WSTEMP = ITEMF2

```





```

C CALCULATE PRESSURE USING THE MODIFIED TAIT EQUATION
2050 WPRESS = I * ((WDENS * RVOL)**N - 1.0)
WSPVCL = 1.0/WDENS
C
C CALCULATE THE WATER'S THERMODYNAMIC PROPERTIES
CALL EPV (TEM1,WSPVCL,WPRESS,WENERG)
WENTH = WENERG + WRESS * WSPVCL
C
IF (COUNT .GT. 0) GO TO 2100
ENTHP1 = WENTH
WDENS = DENS4
COUNT = CCOUNT + 1
GO TO 2050
C
2100 IF (COUNT .GT. 1) GO TO 2200
ENTHP2 = WENTH
COUNT = CCOUNT + 1
2150 * DENS5 = DENS3 + (ENTHAL - ENTHP1) * (DENS4 - DENS3)/(ENTHE2 -
WDENS = DENS5
GO TO 2050
C
2200 ENTHE3 = WENTH
IF (WTEST .GT. DABS(ENTHAL - ENTHE3)) GO TO 2300
DENS3 = DENS4
DENS4 = DENS5
ENTHP1 = ENTHE2
ENTHP2 = ENTHE3
GO TO 2150
C
2300 PRESJ1 = WPRESS
DENSJ1 = WDENS
SOJJ1 = HFSPD * (DENS5 * RVOL)**N1
RETURN
C
C END (SUBROUTINE WATVFI)
C
C
C
C * * * * *
C
SUBROUTINE EPV

```

C THIS SUBROUTINE IS INCLUDED FOR THE INFORMATIONAL PURPOSE OF SHOWING  
 C THE USE OF THE PROGRAM WRITTEN BY A.E. FUHS WAS CONVERTED TO THE  
 C FORTRAN LANGUAGE. THE FOLLOWING IS A DIRECT TRANSLATION OF THAT WORK.  
 C  
 C THE SUBROUTINE IS DESCRIBED AND FLOWCHARTED IN NAVAL POSTGRADUATE  
 C SCHOLARSHIP REPORT NPS67-82-001 BY A.E. FUHS. IT IS USED IN THIS WORK BY  
 C PERMISSION OF THE AUTHCF.

C \* \* \* \* \*

C SUBROUTINE EPV(TIEPV,VEPV,PEPV,E3)

C DECLARE PROGRAM VARIABLES

C REAL LOC7,N,N5,N6,N7

C \* DOUBLE PRECISION X1,X2,X3,X4,Y1,Y2,Y3,B,V8,E9,Z1,CC,E1,E2,M6,  
 C YLIF,XDIF,TIEPV,TIEPV,VIEPV,EEFV,E3

C INTEGER J2,K1

C INITIALIZE KEY VARIABLES

N = 7.15  
 LOOP7 = C.000001  
 X1 = 0.0  
 X2 = 144.C  
 J2 = 0  
 K1 = 0  
 TIEPV = X1

C BEGIN (SUBROUTINE EPV)

C X4 = TIEPV - 273.16  
 C B = 1013CC000.0 \* (3.134 - 0.00165 \* (TEFV - 55.0) - 0.0001181 \*  
 C (TIEPV - 55.0)) \* 2 + 0.00000532 \* (TEFV - 55.0) \*\*3)  
 C V8 = (0.55415 + 0.0002929 \* (TEPV - 25.0) \*\*2) / 1000.0  
 C P9 = B \* ((V8/VEPV) \*\* N - 1.0)  
 C K1 = K1 +

C IF (J2.EC.0) GO TO 3100  
 C IF (J2.GT.1) GO TO 3200  
 C Y2 = P9  
 C J2 = J2 + 1  
 C GO TO 3150

```

C100      Y1 = P9
          IEPV = X2
          J2 = J2 + 1
          GO TO 30CC

C150      M6 = (Y2 - Y1) / (X2 - X1)
          X3 = X1 + (PEPV - Y1) / M6
          IEPV = X3
          GO TO 30CC

C200      Y3 = P9
          IF (LOOP7 .GT. DABS((Y3 - PEPV) / PEPV)) GC IC 3300
          X1 = Y2
          X2 = X3
          Y1 = Y3
          GO TO C150

C      CALCULATE OMEGA OR THE Z-TERM

C300      Z1 = V8 / IEPV
          CO = DSQRT(N * V8 * B)
          N5 = N - 1.0
          N6 = 1.0 / (N * N5)
          N7 = -(1.0 / N5)
          E2 = CO**2 * (N6 * Z1**N5 + 1.0 / (N * Z1) + N7)

C      CALCULATE H10 TERM

          E1 = (3.5644 * (X3 - X4) + 0.000312 * (X3**2 - X4**2)) * 1000.0
          E3 = E1 + E2
          RETURN

C      END (SUBROUTINE EEV)
          END

```

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