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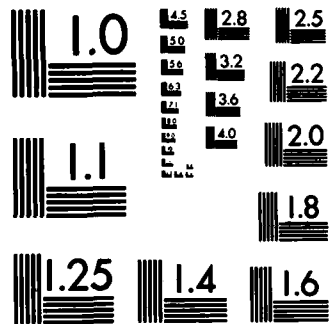
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UNDERSTANDING DIELECTRIC CURE MONITORING

By

Stephen D. Senturia

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An elementary description of the dielectric properties of curing systems, how they are measured, and how they can be interpreted.		

UNDERSTANDING DIELECTRIC CURE MONITORING

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Introduction

An ideal process monitoring method is one which does not perturb the process, yet somehow provides complete documentation of what is going on. Like all ideals, it is rarely approached in practice, but nevertheless serves to guide the development of suitable methods. The monitoring of the cure of thermosets is a good example. The continuous transition from viscous liquid to rubbery gel to vitreous solid characteristic of thermosets produces such large changes in properties that almost any method of physical measurement can be expected to yield a useful result. In practice, however, techniques such as calorimetry, infrared spectroscopy, and dynamical mechanical measurements, all of which yield highly useful information, tend to be confined to carefully controlled laboratory situations because of the difficulties of in-situ use. For cure monitoring in the manufacturing environment, alternative techniques are needed.

Electrical measurements offer an attractive possibility. During the past decade, various in-situ methods for examining changes in the electrical conductivity of resins during cure have been developed, as have methods for studying changes in dielectric properties. It is useful to group "electrical conductivity" and "dielectric properties" under the common heading "dielectric properties", because, as will be shown below, measurements of the frequency dependence of dielectric properties can yield information both on the conductivity and on the dielectric properties usually associated with hindered dipole rotation.

Parallel Plate Measurements

The most common form of dielectric measurement involves a pair of parallel plates or foils either placed in contact with or embedded into the specimen. The electrode pair form a circuit element which can be represented as a capacitor in parallel with a conductor. The capacitance of the capacitor depends on the area of the plates, the spacing between the plates, and on the dielectric permittivity ϵ' of the medium between the plates. If the permittivity is dispersive (i.e., frequency dependent), then the measured capacitance would also depend on frequency. The conductance of the conductor in parallel with the capacitor depends similarly on the plate area and spacing, and also depends on the conductivity of the medium σ . If σ depends on frequency, then the conductance would also. A conventional way of relating the permittivity and conductivity contributions to one another is to define a quantity with the same dimensions as ϵ' , but proportional to conductivity. From basic

theory, the correct corresponding quantity, called the loss factor and denoted by ϵ'' , and is equal to σ/ω , where ω is the angular frequency at which both ϵ' and ϵ'' are measured. Both ϵ' and ϵ'' are usually reported in units of the permittivity of free space ϵ_0 , which has the value 8.85×10^{-14} Farads/cm. In this case, the ϵ' value corresponds to the conventionally defined dielectric constant, and a material with an ϵ'' value of unity at a frequency of 1 Hz has a conductivity equal to 5.56×10^{-13} (Ohm cm)⁻¹.

Because both the capacitance and conductance measured from parallel plates depend directly on plate spacing, and because this spacing typically changes during a cure cycle either by shrinkage of the curing material or through the application of pressure part-way through the cure cycle, one usually finds reports of the ratio of ϵ'' to ϵ' , called the loss tangent (or, conventionally, $\tan \delta$). Other authors report the dissipation factor D of the capacitor, which is experimentally equivalent to $\tan \delta$. In either case, the reason for reporting $\tan \delta$ or D is that the plate spacing and area both cancel out in the ratio. This has the benefit of yielding a result that is relatively insensitive to geometric changes, but has the disadvantage of preventing simultaneous separate measurements of ϵ' or ϵ'' . When $\tan \delta$ is observed to change, one does not know a priori whether to attribute that change to ϵ' , to ϵ'' , or to both.

Measurements on parallel plate electrodes involve some form of capacitance bridge. Since the measurement apparatus is separated from the plates by cables, the very tiny currents that charge and discharge the cable capacitances coupled with electrostatic noise pickup on the cabling ultimately limit the lowest frequency at which satisfactory measurements can be performed. Typical values of this low-frequency limit are on the order of 100 Hz, although some extensions to lower frequencies are possible with extremely careful design. At the high frequency end, cable inductance eventually limits measurement accuracy, typically in the 100 kHz - 1 Mhz range.

Microdielectrometry

An alternative to the parallel plate approach to dielectric property measurements has recently been developed in our laboratories at MIT. The electrodes form an interdigitated pair, or comb pattern, on the surface of a 2x4 mm integrated circuit (or "chip"). The sample is placed directly on these electrodes; alternatively, the chip is embedded in the material undergoing cure. One of the electrodes is driven with a sinusoidal signal, and field-effect-transistor amplifiers within the chip detect the amount of charge reaching the other electrode through the sample medium. For a given chip geometry, calibration tables which relate the magnitude and relative phase of this charge to the ϵ' and ϵ'' values for the medium have been developed. The electrode geometry does not depend on temperature, pressure, or on shrinkage during cure. Furthermore, given the precision of modern microelectronic fabrication methods, different chips are sufficiently identical to one another so that individual calibration is not required. Because of the built-in amplification, there is no low-frequency limitation to the measurement other than the intrinsic reaction rate of the medium being studied, which represents a significant advantage over parallel

plate methods. A semiconductor diode temperature sensor is also built into the chip, so that extremely local temperature changes during cure can be monitored. The chip is connected through fine wires to electronic instrumentation which, under control of a desk-top calculator, sequences the measurement frequency, performs the measurement, interprets the data to yield ϵ' and ϵ'' (and, if desired, temperature), logs the data onto magnetic cassette, and generates real-time display of the data.

Using the Microdielectrometry system, we have studied the cure of a number of thermosetting systems as functions of time, temperature, and frequency. We consistently find that the frequency dependence of both ϵ' and ϵ'' must be examined in order to provide meaningful interpretation of the data. A general model for how ϵ' and ϵ'' change is described below.

Understanding the Data

In our most elementary model, the permittivity is represented as consisting of a constant ϵ_∞ (the equivalent high-frequency dielectric constant) plus a Debye-like dipole orientation term,

$$\epsilon' = \epsilon_\infty + \frac{(\epsilon_s - \epsilon_\infty)}{1 + (\omega\tau)^2}$$

where ϵ_s is the permittivity one measures at low frequency, and τ is the dielectric relaxation time. Note that for $\omega\tau=1$, ϵ' falls at the midpoint between ϵ_s and ϵ_∞ . Corresponding to the dipole term in ϵ' there is a frequency dependent contribution to ϵ'' which exhibits a peak at $\omega\tau=1$. We also add a bulk conductivity term, σ/ω , where in the elementary model, σ is frequency-independent, yielding a total ϵ'' of the form,

$$\epsilon'' = \frac{\sigma}{\omega} + \frac{(\epsilon_s - \epsilon_\infty)\omega\tau}{1 + (\omega\tau)^2}$$

The parameters ϵ_∞ , ϵ_s , and τ characterize the permittivity, and σ plus the other parameters characterize the loss factor. In particular, the frequency dependence of ϵ'' can provide immediate insight into whether one's measurements are dominated by the conductivity term or by the dipole term. An example will be helpful.

If we assume that the bulk conductivity of a curing system falls exponentially with time due to an increase in viscosity produced by crosslinking, and if we also assume that the dipole relaxation time grows exponentially with time for the same reason, then for typical initial values of σ and τ , the "Theory" curves of ϵ' and ϵ'' versus time shown in Figure 1 would be expected. Note the characteristic decrease in ϵ' at successively lower frequencies as τ grows with time (a classic relaxation), and note the corresponding loss peaks in ϵ'' . At 1000 Hz in the example, the loss factor drops, then increases, then drops again during cure. Further, note that early in cure, ϵ'' decreases monotonically with increasing frequency, but when the point is reached where the dipole loss peaks show up, the frequency dependence becomes complex. Since as one approaches the loss peaks, ϵ' is also changing with frequency, if one can only

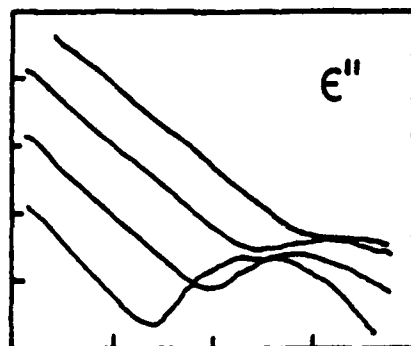
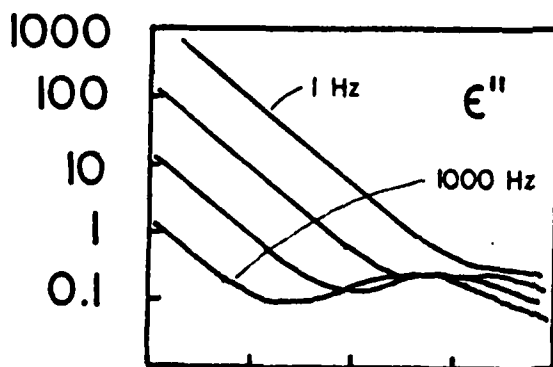
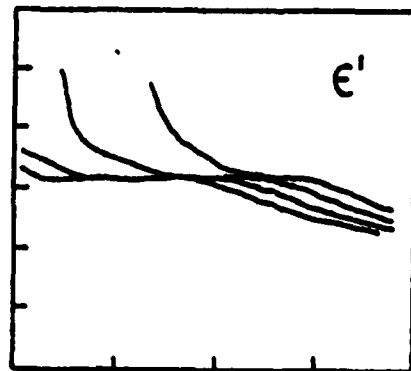
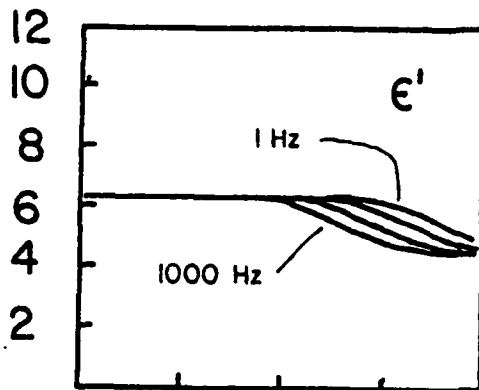
measure the ratio of ϵ'' to ϵ' , one cannot easily assign an origin to observed frequency variation. There is a distinct advantage, therefore, in obtaining separate measurements of both ϵ' and ϵ'' .

To illustrate the reasonableness of the elementary model, the "Experiment" side of Figure 1 shows actual data from an isothermal cure of an amine-epoxy system. The only major differences between the Experiment and the Theory in Figure 1 are the dispersion early in cure in ϵ' , which is attributable to ionic conduction with blocked electrodes, and the persistence of relatively high loss factors at the end of cure, which is attributable to the drop in reaction rate due to vitrification of the sample.

The model outlined above represents a useful starting point for interpreting dielectric data during cure. Rate of change of loss factor can be related to reaction rate. The detailed shape of the dipole orientation peaks can be related to the distribution of relaxation times within the sample, which can, in principle at least, be related to properties of the network being formed by crosslinking. In every case, however, measurements of both ϵ' and ϵ'' over a wide frequency range are required to permit meaningful interpretation of the results.

THEORY

EXPERIMENT



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