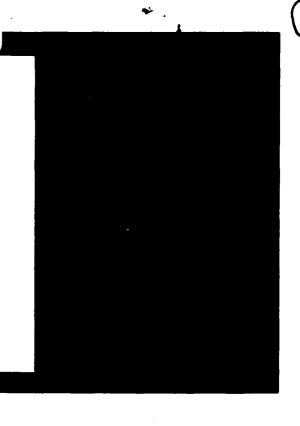


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LOWER BOUNDS FOR THE LIFE SPAN OF SOLUTIONS OF MONLINEAR WAVE EQUATIONS IN THREE DIMENSIONS

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ABSTRACT

The paper deals with strict solutions $u(x,t) = u(x_1,x_2,x_3,t)$ of an equation

$$u_{tt} - \sum_{i,k=1}^{3} a_{ik}(Du)u_{x_ix_k} = 0$$

where Du is the set of 4 first derivatives of u. For given initial values $u(x,0)=\epsilon F(x)$, $u_t(x,0)=\epsilon G(x)$ the life span $T(\epsilon)$ is defined as the supremum of all t to which the local solution can be extended for all x. Blow-up in finite time corresponds to $T(\epsilon)<\infty$. Examples show that this can occur for arbitrarily small ϵ . On the other hand $T(\epsilon)$ must at least be very large for small ϵ . Assuming that $a_{ik}^{}$, $F,G\in C$, that $a_{ik}^{}$, and that F,G have compact support, it is shown that $\lim_{\epsilon\to 0} \epsilon^{N} T(\epsilon) = \infty$ for every N. This result had been established previously only for N < 4.

AMS (MOS) Subject Classifications: 35L15, 35L67, 35L70, 35B40

Key Words: partial differential equations, hyperbolic equations, wave equations, second order nonlinear equations, shocks and singularities

Work Unit Number 1 (Applied Analysis)

This article represents work performed at the Courant Institute of Mathematical Sciences and supported by the National Science Foundation Grant No. MCS-79-00812 and the Office of Naval Research Grant No. N00014-76-C-0439. An outline is to appear in the Proceedings of the National Academy of Sciences, June 1982. This report was prepared at the Mathematics Research Center, sponsored by the United States Army under Contract No. DAAG 29-80-C-0041.

LOWER BOUNDS FOR THE LIFE SPAN OF SOLUTIONS OF NONLINEAR WAVE EQUATIONS IN THREE DIMENSIONS*

Fritz John

This paper deals with existence of solutions $u(x_1,x_2,x_3)=u(x,t)$ of a nonlinear wave equation of the form

$$u_{ee} - \sum_{i,k=1}^{3} a_{ik}(u^{i})u_{x_{i}x_{k}} = 0$$
 (1a)

for large times t. Here u' stands for the gradient vector

$$u' = (u_{x_1}, u_{x_2}, u_{x_3}, u_t) = (D_1 u, D_2 u, D_3 u, D_4 u) = Du$$
 (1b)

We assume that the $a_{ik}(U)$ are in C in a closed ball $|U| \le \delta$ in R^4 , and that

$$a_{ik}(0) = \delta_{ik} , \qquad (1c)$$

so that (1a) goes over into the classical linear wave equation

$$\Box u = u_{tt} - \Delta u = 0 \tag{1d}$$

for "infinitesimal" u. The solution u of (1a) is to be found from initial conditions for t=0. Here we use initial data of the form

$$u(x,0) = \varepsilon f(x), \quad u_{\varepsilon}(x,0) = \varepsilon g(x) \quad \text{for } x \in \mathbb{R}^3$$
 (1e)

where f,g are fixed functions in $C_0^{\infty}(\mathbb{R}^3)$ and ϵ is a parameter that serves to measure the <u>amplitude</u> of the initial values.

For a given choice of functions $f(x),g(x),a_{ik}(0)$ we define the <u>life span</u> $T=T(\varepsilon)$ as the supremum of all s such that a C-solution of (1a,e) with $|u'|<\delta$ exists for $x\in\mathbb{R}^3$ and $0\leq t\leq s$. One knows that $T(\varepsilon)>0$ for sufficiently small $|\varepsilon|$; ("local" solutions of the initial value problem exist). Existence of "global" solutions would correspond to $T(\varepsilon)=0$. One knows also (see [4]) that $T(\varepsilon)<0$ at least in some cases;

[&]quot;This is a continuation of the author's paper DELAYED SINGULARITY FORMATION OF SOLUTIONS OF MONLINEAR NAVE SQUATIONS IN HIGHER DIMENSIONS, Comm. Pure Appl. Nath. 29, (1976), 649-682, referred to as (*) in the sequel.

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(the local solution "blows up" in finite time). Thus for the equation*

$$u_{\text{tt}} - \frac{\Delta u}{1 - 2u_{\bullet}} = 0 \tag{2a}$$

we have $T(E) \leftarrow F$ for all sufficiently small positive E, when

$$\int_{\mathbb{R}^3} g(x) dx > 0 \tag{2b}$$

Realistic bounds for T(E) are difficult to obtain. In the example (2a,b) one can show that

$$T(\varepsilon) < \lambda \exp(B\varepsilon^{-2})$$
 (2c)

with certain constants A,B. In the present paper we show that T(E) increases with diminishing E faster than any reciprocal, power of E:

For any real positive N

$$\lim_{\varepsilon \to 0} |\varepsilon|^{H} T(\varepsilon) = 0 \tag{3}$$

Remarks

- (a) Statement (3) could be wrong for <u>plane wave</u> solutions of (1a,e). Indeed it does not hold for such solutions for M > 1, if equation (1a) is genuinely nonlinear. But plane wave solutions are excluded by our assumption that f and g have compact support.
- (b) Relation (3) had been proved in (*) for N restricted to the interval 0 < M < 4. The proof given here for general N closely follows the ideas developed in (*). The extension to M > 4 requires some not so obvious additional estimates, contained in the MAIN LEMMA below.
- (c) The arguments leading to (3) would also permit to derive more specific lower bounds for T(5) for fixed 5 for specific f,g,a_{ik}. These bounds would depend on assumptions on the <u>growth</u> of the derivatives of those functions with order. The methods used here do not yield (3) for general N, when only a <u>finite</u> number of derivatives of coefficients and data are available.

This contrasts with the situation in more than 5 space dimensions, where $T(\epsilon) = \infty$ for all sufficiently small ϵ , as shown by Klainerman [1]. See also [2], [3].

The proof of the THEOREM is broken up into a sequence of lemmas. Roughly one argues as follows. First a combination of a priori L_2 -estimates ("energy" estimates) and Sobolev inequalities leads to a <u>local</u> existence proof. To get existence for really large times to one has to establish pointwise <u>decay</u> of u^* with time. Decay cannot be inferred easily from L_2 -estimates. It is established here by approximating u^* in the L_2 -norm by vectors u^* that can be shown to decay. We select here for u^* the partial sums of the formal power series of u^* in terms of e^* :

$$u \approx \varepsilon u_1 + \varepsilon^2 u_2 + \varepsilon^3 u_3 + \dots \tag{4}$$

The $u_k^-(x,t)$ can be found explicitly by quadratures from linear recursion formulae. The THEOREM follows if one can prove that the derivatives of the u_k^- decay at least like $t^{-1}(\log t)^{2k-2}$ for large t. (The solutions of the <u>linear</u> wave equation (1d) decay like t^{-1} for initial data of compact support). In (*) this is shown for k=1,2,3. The crude technique used there, only estimating absolute values, does not work for k>3. One has to rely instead on <u>cancellation</u> of the worst contributions (exploited here by integration by parts). This comes about because the u_k^- satisfy certain <u>radiation</u> conditions, making them behave asymptotically like outgoing spherical waves. One also has to make use of the fact that the derivatives of the u_k^- decay more strongly, like $t^{-2}(\log t)^{2k-2}$, except for small |x|/t or small 1-|x|/t. This is the essential content of the <u>Main Lemma</u>.

Notation and assumptions.

For a vector $U \in \mathbb{R}^4$ we define |U| as its euclidean length.* For U = U(x,t) with $x \in \mathbb{R}^3$ and a non-negative integer n we set

$$|U(x,t)|_{n} = \sqrt{\sum_{|\alpha| \leq n} |D^{\alpha}U(x,t)|^{2}}$$
 (5a)

and introduce for fixed t the two norms

$$\left\{ U(t) \right\}_{n} = \sup_{x \in \mathbb{R}^{3}} \left\{ U(x,t) \right\}_{n}$$
 (5b)

[&]quot;This differs slightly from the definition of |U| in (*), (37).

$$\|u(t)\|_{n} = \sqrt{\int_{\mathbb{R}^{3}} (\|u(x,t)\|_{n})^{2} dx}$$
 (5c)

In what follows we shall use almost exclusively n=3 in (5b) and n=5 in (5c). By Sobolev's inequality there exists universal C such that

$$\{v(e)\}_{q} \leq civ(e)I_{q} \tag{5d}$$

We assume that the $a_{ik}(U)$ are of class C in the closed ball

and satisfy

$$\mathbf{a}_{ik}(\mathbf{U}) = \mathbf{a}_{ik}(\mathbf{U}) \tag{6b}$$

$$\mathbf{a_{ik}}(0) = \delta_{ik} \tag{6c}$$

Without restriction generality we can assume here that δ is so small that

and that

$$\frac{1}{2} |v|^2 \le v_4^2 + \sum_{i,k=1}^3 a_{ik}(v) v_i v_k \le 2|v|^2$$
 (6e)

for all U satisfying (6a) and all $V = (V_1, V_2, \overline{V}, ..., V_4) \in \mathbb{R}^4$. With a given vectorfield U = U(x,t) we associate the <u>linear</u> differential coerator

$$P(U) = D_4^2 - \sum_{i=k=1}^{3} a_{ik}(U)D_iD_k$$
 (7a)

The result of applying P(U) to a scalar function u(x,t) will be written P(U)[u]. In this notation (1a) becomes P(u')[u] = 0.

LEMMA I. (a priori estimate).

Let for a certain T > 0 the function u(x,t) be a C -solution of

$$P(u')[u] = 0$$
 for $x \in \mathbb{R}^3$, $0 \le t \le T$ (84)

for which

$$|u'(x,t)| < \delta \tag{3b}$$

Set

$$\lambda(t) = \max_{1 \le k \le 5} (\{u^{\epsilon}(t)\}_3)^k$$
(80)

There exists a constant C (depending only on the suprema of the $a_{ik}(U)$ and their derivatives of orders < 5 for |U| < δ), such that*

$$\|u'(t)\|_{5} \le 2\|u'(0)\|_{5} \exp(C \int_{0}^{t} \lambda(s) ds)$$
 (84)

<u>Proof.</u> See (*), p. 659 for N=5, N'=3. The extra factor 2 in (8d) is due to the difference in the definitions of |U|, and the use of (6e).

We assign initial data

$$u(x,0) = ef(x), u_{\mu}(x,0) = eg(x)$$
 (9a)

to the solution u(x,t) of (Sa), where $f,g\in C_0^\infty(\mathbb{R}^3)$. Without restriction of generality we can assume that

$$f(x) = g(x) = 0$$
 for $|x| > 1$. (9b)

Higher derivatives of u(x,t) for t=0 can be computed from f,g with the help of the differential equation (8a). In particular there exists a constant σ (depending on the f,g,a_{ik}) such that

$$8u^{\epsilon}(0)8_{5} \le c\epsilon$$
 for $|\epsilon| \le 1$ (9a)

LENGA II. (Local existence).

Let for given f,g,a,,

$$|\varepsilon| < \min(1, \frac{\delta}{4\sigma^2}) \tag{10a}$$

We can find a value $T_0 = T_0(\epsilon) > 0$ and a solution $u \in C^{\infty}$ of

$$P(u')\{u\} = 0 \text{ for } x \in \mathbb{R}^3, \quad 0 \le t \le T_0$$
 (19b)

$$u(x,0) = \epsilon f(x), \ u_{\xi}(x,0) = \epsilon g(x) \text{ for } x \in \mathbb{R}^3.$$
 (10c)

for which

$$\int_{0}^{T_0} \left\{ u^*(s) \right\}_3 ds = \frac{1}{C} \log 2$$
 (10d)

$$\|u^{\epsilon}(t)\|_{5} \le 4\|u^{\epsilon}(0)\|_{5} \le 4c\|\epsilon\| \le \frac{\delta}{C} \text{ for } 0 \le t \le T_{0}$$
 (10e)

[&]quot;We can use the same C in (5d) and (8d).

$$|u'(x,t)| < \delta < 1 \text{ for } x \in \mathbb{R}^3, 0 \le t \le T_0$$
 (10f)

$$u(x,t) = 0$$
 for $|x| > t + 1$, $0 \le t \le T_0$ (10g)

Proof. By (5b), (5d), (9c), (10a)

$$|u^*(x,0)| \le \{u^*(0)\}_3 \le c|u^*(0)|_5 \le cc|\varepsilon| \le \frac{\delta}{4} < \delta$$
. (11a)

Let T = T(E) be the life span, as defined on p. 1. Here T(E) > 0 because of (11a).

Then either

$$\lim_{t \to T} \int_{0}^{t} \lambda(s) ds = 0$$
 (11b)

or

$$\sup |u^{\varepsilon}(x,t)| = \delta \quad \text{for} \quad x \in \mathbb{R}^3, \quad 0 \le t \le T \tag{11c}$$

(See (*), pp. 660-661). We define $T_0 = T_0(\epsilon)$ by

$$\int_0^{T_0} \lambda(s) ds = \frac{\log 2}{C} \text{ when } \int_0^T \lambda(s) ds > \frac{\log 2}{C}$$
 (11d)

$$T_0 = T$$
 when $\int_0^T \lambda(s) ds < \frac{\log 2}{C}$ (11e)

In either case

$$\int_{0}^{T_0} \lambda(s) ds < \frac{\log 2}{C}$$
 (11f)

This implies by (8d), (9c), (10a), (5d) that for $0 \le t < T_0$

$$||u'(t)||_{S} \le 4||u'(0)||_{S} \le 4c||\varepsilon||$$
 (11f*)

$$\left\{ \mathbf{u}^{\mathbf{t}}(\mathbf{t}) \right\}_{3} \leq 4cC|\mathbf{c}| < \delta < 1 \tag{11g}$$

But then also by (10a)

$$|u^{\varepsilon}(x,t)| \le 4cC|\varepsilon| < \delta$$
 for $x \in \mathbb{R}^3$, $0 \le t \le T_0$ (11h)

This is incompatible with

$$\int_{0}^{T} \lambda(s) ds < \frac{\log 2}{C} < \infty$$

which would imply (11c) with $T = T_0$. Hence

$$\int_{0}^{T_0} \lambda(s) ds = \frac{\log 2}{C}, \quad T_0 < T$$
 (111)

It follows from (11g) and the definition (8c) of $\lambda(t)$ that

$$\lambda(t) = \{u^{t}(t)\}_{3} \text{ for } 0 \le t \le T_{0}. \tag{11j}$$

Thus (11i) implies (10d). Moreover (11f*), (11h) yield (10e), (10f). Finally (10g) is a consequence of assumption (9b); the effect of zero initial data is the same as for the linear wave equation; (see [4], p. 49).

Relation (10d) permits to derive lower bounds for T_0 and hence for T from upper bounds for $\{u^i(t)\}_{q}$. Trivially one has from (11g), (10d)

$$4cc^2T|\varepsilon| > 4cc^2T_n|\varepsilon| > \log 2$$
 (12)

Not much more can be extracted from an upper bound for $\{u'(t)\}_3$ as long as this bound does not show decay in t. Bounds showing decay cannot be obtained from bounds for $\{u'(t)\}_5$, which is not likely to decay; (it does not in the linear case). One way to find better estimates for $\{u'(t)\}_3$ is to compare u with an "approximation" \tilde{u} , for which $\{\tilde{u}'\}_3$ shows the appropriate decay.

Let $\varepsilon_{,u,T_0}$ be as in Lemma II. Let \widetilde{u} be an "approximation" of u, for which $P(u')[\widetilde{u}] = w(x,t) \text{ for } x \in \mathbb{R}^3, \ 0 \le t \le T_0 \tag{12a}$

$$\tilde{u}(x,0) = u(x,0); \quad \tilde{u}_{\epsilon}(x,0) = u_{\epsilon}(x,0)$$
 (12b)

There exists a constant C^* depending on the a_{ik} , but not on ϵ and \tilde{u} such that

$$I(u' - \tilde{u}')(t)I_5 \le 4M \int_0^{T_0} Iw(s)I_5 ds$$
 for $0 \le t < T_0$ (12c)

where

$$H = \exp(C^{*} \int_{0}^{T_{0}} \{\widetilde{u}^{*}(s)\}_{5} ds)$$
 (12d)

It follows from (5d) that

$$\{u'(t)\}_3 \le \{\widetilde{u}'(t)\}_3 + 4NC \int_0^{T_0} \|w(s)\|_5 ds$$
 (12e)

<u>Proof.</u> See (*), pp. 662, where (12a) is written as a symmetric hyperbolic system for the vector $\tilde{\mathbf{u}}$. The factor 4 in (12c) arises again from the difference in the definitions of

|U|. The norm $|w(s)|_{5}$ for the scalar w is defined exactly as that for vectors U in (5a,c).

We shall apply Lemma III to the case where \tilde{u} is one of the partial sums of the formal power series expansion for u with respect to ϵ_1

$$u = \varepsilon u_1 + \varepsilon^2 u_2 + \varepsilon^3 u_3 + \dots$$
 (13a)

To get recursion formulae for the u_k we expand the operator P(U) of (7a) formally with respect to U_2

$$P(0) = \Box + P_1(0) + P_2(0) + \dots$$
 (13b)

where $P_k(U)$ is a form of degree k in the components of U with coefficients, that are quadratic in D_1,D_2,D_3 . Substituting (13a,b) into the equation $P(u^*)[u] = 0$ and comparing terms with the same power of E we arrive at a sequence of equations

$$\Box_{\mathbf{u}_{\mathbf{k}}} = Q_{\mathbf{k}} \tag{13c}$$

for M=1,2,3,... Here Q_{M} is a polynomial in the first and second derivatives of the u_{k} . Each term of this polynomial is (except for a constant factor) of the form

$$\begin{array}{ccc}
\mathbf{r} & \mathbf{a}_{\mathbf{i}} \\
\mathbf{II} & (\mathbf{D}^{\mathbf{u}}_{\mathbf{k}_{\mathbf{i}}}) \\
\mathbf{i} = 1 & \mathbf{k}_{\mathbf{i}}
\end{array} \tag{13d}$$

with multi-indices α_i and integers k_i , where $|\alpha_i| = 1,2$ and

$$\sum_{i=1}^{T} k_{i} = H_{i} \quad k_{i} \ge 1, \quad r \ge 2$$
 (13e)

In particular

$$Q_1 = 0$$

$$Q_2 = -P_1(u_1^t)[u_1]$$
 (132)

$$Q_{3} = -P_{1}(u_{1}^{1})(u_{2}) - P_{1}(u_{2}^{1})(u_{1}) - P_{2}(u_{1}^{1})(u_{1})$$
 (13g)

$$Q_{4} = -P_{1}(u_{1}^{1})[u_{3}] - P_{1}(u_{2}^{1})[u_{2}^{1}] - P_{1}(u_{3}^{1})[u_{1}]$$

$$-P_{2}(u_{1}^{1})[u_{2}] - 2P_{2}(u_{1}^{1}, u_{2}^{1})[u_{1}] - P_{3}(u_{1}^{1})[u_{1}]$$
(13g)

(with $P_2(U,V)$ denoting the polar from of $P_2(U)$). The equations (13c) combined with the initial conditions

$$u_1 = f(x), D_2 u_1 = g(x)$$
 for $t = 0$ (13h)

$$u_k = 0$$
, $D_A u_k = 0$ for $t = 0$ when $k > 1$ (13i)

recursively determine the un. Setting

$$\widetilde{\mathbf{u}}_{\mathbf{N}} = \varepsilon \mathbf{u}_{1} + \varepsilon^{2} \mathbf{u}_{2} + \dots + \varepsilon^{N} \mathbf{u}_{N}$$
 (13j)

we shall have to estimate $\|P(u^*)\{\widetilde{u}_N^*\}\|_{S}$ and $\{\widetilde{u}_N^*\}_{S}$ and then apply Lemma III.

The required estimates for the \widetilde{u}_N involve the asymptotic behavior of the u_k for large t. To describe this behavior adequately we introduce the <u>radiation operators</u>

$$L_{\underline{i}} = D_{\underline{i}} - \sum_{k=1}^{3} x_{\underline{i}} x_{k} |x|^{-2} D_{\underline{k}}$$
 for $\underline{i} = 1,2,3$ (14a)

$$L_4 = D_4 + \sum_{k=1}^{3} x_k |x|^{-1} D_k$$
 (14b)

Finally we denote by $\Box^{-1}w(x,t)$ the solution v of the equation

$$\Box v = w(x,t) \tag{14c}$$

with vanishing initial data:

$$v(x,0) = v_{+}(x,0) = 0$$
 (14d)

We first note the asymptotic behavior of solutions of the linear wave equation.

LEDOGA IV.

Let v(x,t) be the solution of

$$\Box \forall (x,t) = 0 \text{ for } x \in \mathbb{R}^3, \ t \ge 0$$
 (15a)

$$v(x,0) = f(x), v_{\pm}(x,0) = g(x)$$
 (15b)

where $f \in C^3$, $g \in C^2$ and

$$f(x) = g(x) = 0$$
 for $|x| > 1$ (15c)

Then

The Arthur St.

$$v(x,t) = 0$$
 for $|t - |x|| > 1$ (15d)

$$v = 0(\frac{1}{t+2}); D_{\underline{1}}v = 0(\frac{1}{t+2}), L_{\underline{1}}v = 0(\frac{1}{(t+2)^2})$$
 (15e)

Here "0" stands for a constant depending on $\{f\}_2$ and $\{g\}_1$.

Proof. Classically

$$v = v_1(x,t) + D_4 v_2(x,t)$$
 (16a)

where

$$v_1 = \frac{1}{4\pi t} \iint_{|y-x|=t} g(y) ds_y; \quad v_2 = \frac{1}{4\pi t} \iint_{|y-x|=t} f(y) ds_y$$
 (16b)

(15d) is obvious from (15c) since for |t - |x|| > 1 the ball |y| < 1 and the sphere |y - x| = t do not intersect. Since the area of intersection of the ball and the sphere is at most equal to $\min(4\pi, 4\pi t^2)$ it follows immediately that

$$v_1(x,t) \leq \left(\min\left(\frac{1}{t}, t\right)\right) \sup |g| = 0\left(\frac{1}{t+2}\right)$$
 (16c)

The same argument shows that $D_i v_1 = 0(1/(t+2))$, where

$$D_i v_1 = \frac{1}{4\pi t} \iint_{|y-x|=t} D_i g(y) dS_y$$
 for $i = 1,2,3$ (16d)

$$D_{\underline{q}}v_{\underline{q}} = \frac{1}{4\pi t} \iiint_{|y-x|=t} \left(\frac{d\underline{q}}{d\underline{n}} + \frac{\underline{q}}{t}\right) d\underline{s}_{\underline{y}}$$
 (16e)

Here

$$\frac{dq}{dn} = \sum_{k=1}^{3} \xi_k D_k q(y)$$
 (162)

is the normal derivative of g on the sphere |y - x| = t with direction cosines

$$\xi_{k} = \frac{1}{k} \left(y_{k} - x_{k} \right) \tag{16g}$$

of the exterior normal.

The same rate of decay is then found for $D_4^{}v_2^{}$ and for

$$D_{i}D_{4}v_{2} = D_{4} \frac{1}{4\pi t} \iint_{|y-x|=t} D_{i}f dS_{y} \text{ for } i = 1,2,3$$

$$D_{4}^{2}v_{2} = \Delta v_{2} = \frac{1}{4\pi t} \iint_{|y-x|=t} \Delta f dS_{y}$$

completing the proof of

$$v = 0(\frac{1}{t+2}), D_1 v = 0(\frac{1}{t+2})$$
 (16h)

One easily verifies, by transforming the surface integral (16b) for v_{\uparrow} into a volume integral, that for i=1,2,3

$$D_{i}v_{1} = \frac{1}{4\pi\epsilon} \iiint_{|y-x|=\epsilon} \xi_{i}(\frac{2q}{\epsilon} + \frac{dq}{dn})ds_{y} = \frac{1}{4\pi\epsilon^{2}} \iiint_{|y-x|=\epsilon} (2\xi_{i}q + y_{i}\frac{dq}{dn} - x_{i}\frac{dq}{dn})ds_{y}$$
 (161)

It follows that

$$L_{i}v_{1} = \frac{1}{4\pi t^{2}} \iiint_{y-x|=t} (2\xi_{i}g + y_{i} \frac{dg}{dn})dS_{y}$$

$$- \sum_{k=1}^{3} \frac{1}{4\pi t^{2}} x_{i}x_{k}|x|^{-2} \iiint_{y-x|=t} (2\xi_{k}g + y_{k} \frac{dg}{dn})dS_{y} = 0((t+2)^{-2})$$

for 1 = 1,2,3, since

$$\iint_{|y-x|=t} ds_{y} \le 4\pi, |\xi_{1}| \le 1, |y_{1}| \le 1$$

$$|y|<1$$

Moreover by (16e,1), (15d)

$$\begin{split} \mathbf{L}_{4}\mathbf{v}_{1} &= \frac{1}{4\pi\epsilon^{2}} \iint_{|\mathbf{y}-\mathbf{x}|=\epsilon} (1+2\sum_{k} \mathbf{x}_{k} \mathbf{\xi}_{k} |\mathbf{x}|^{-1}) \mathbf{g} \ d\mathbf{s}_{\mathbf{y}} \\ &+ \frac{1}{4\pi\epsilon^{2}} \iint_{|\mathbf{y}-\mathbf{x}|=\epsilon} ((\epsilon-|\mathbf{x}|) + \sum_{k} \mathbf{y}_{k} \mathbf{x}_{k} |\mathbf{x}|^{-1}) \frac{d\mathbf{g}}{d\mathbf{n}} \ d\mathbf{s}_{\mathbf{y}} = 0((\epsilon+2)^{-2}) \end{split}$$

Then also

$$L_{i}D_{k}v_{2} = 0((t+2)^{-2})$$
 for $i = 1,2,3,4$; $k = 1,2,3$,

since $D_k v_2$ is obtained from v_1 by replacing g by $D_k f_*$. From the identities

$$L_{i}D_{4}v_{2} = L_{4}D_{i}v_{2} - \sum_{k=1}^{3} (x_{i}x_{k}|x|^{-2}L_{4} + x_{k}|x|^{-1}L_{i})D_{k}v_{2}$$

for i = 1, 2, 3, and

$$L_4D_4v_2 = \sum_{k=1}^{3} (L_k + x_k|x|^{-1}L_4)D_kv_2$$

it then follows that

$$L_i D_4 v_2 = (0((t+2)^{-2}) \text{ for } i = 1,2,3,4$$

This completes the proof of (15e).

LENGA V. (The MAIN LENGA).

Let $w(x,t) \in C^2$ for $x \in \mathbb{R}^3$, t > 0, and let

$$w(x,t) = 0$$
 for $|x| > t + 1$ (17a)

$$|w|, |D_1 w| \le H \frac{\log^k(t+2)}{(|x|+2)^2(t-|x|+2)^2}$$
 (17b)

$$|L_i w| \le H \frac{\log^k(t+2)}{(|x|+2)^2(t-|x|+2)(t+2)}$$
 (17c)

for i = 1,2,3,4 with a certain k. Then

$$\mathbf{u} = \mathbf{0}^{-1} \mathbf{v} \tag{17d}$$

satisfies

$$u(x,t) = 0$$
 for $|x| > t + 1$ (17e)

$$|D_{i}u| \le AH \frac{\log^{k+2}(t+2)}{(|x|+2)(t-|x|+2)}$$
 (17f)

$$|L_1 u| \le AM \frac{\log^{k+2}(t+2)}{(|x|+2)(t+2)}$$
 (17g)

for i = 1,2,3,4, where λ is a universal constant.

<u>Proof.</u> We postpone the lengthy proof of the MAIN LENGA which only deals with a property of the operator \Box , to the Appendix, in order not to interrupt the arguments leading to the proof of the THEOREM.

LEHOLA VI.

Let the u_N be defined recursively for N=1,2,3,... by the differential equations (13c) with initial conditions (13h,i) satisfying (9b). Then

$$u_{x}(x,t) = 0$$
 for $|x| > t + 1$ (18a)

$$D^{\alpha}u_{H} = 0\left(\frac{\log^{2H-2}(t+2)}{(|x|+2)(t-|x|+2)}\right) \text{ for } |\alpha| > 1$$
 (18b)

$$L_1 D^{\alpha} u_N = 0 \left(\frac{\log^{2N-2}(t+2)}{(|x|+2)(t+2)} \right) \text{ for } |\alpha| \ge 0$$
 (18c)

Here "0" stands for a constant depending on f,g,a_{ib},N,a. Relations (18a,b) imply that

$$\left\{u_{k}^{*}(t)\right\}_{k} = 0\left(\frac{\log^{2k-2}(t+2)}{t+2}\right)$$
 (184)

$$\|u_{H}^{*}(t)\|_{\dot{R}} = 0(\log^{2H-2}(t+2))$$
 (18e)

for all k > 0.

<u>Proof.</u> We use induction over N. Since $\Box D^{\alpha}u_{1} = 0$ and the initial data of $D^{\alpha}u_{1}$ have their support in the ball $|x| \le 1$, we find from Lemma IV that (18a,b,c) holds for M = 1. Let (18a,b,c) hold, when M is replaced by a smaller number. Then $\Box D^{\alpha}u_{1} = D^{\alpha}\Box u_{2}$ for $|\alpha| > 0$ is a linear combination of terms of the form

$$\begin{array}{ccc}
\mathbf{r} & \mathbf{a}_{\mathbf{i}} \\
\mathbf{I} & (\mathbf{D}^{\mathbf{i}} \mathbf{u}_{\mathbf{k}}) \\
\mathbf{i} = 1
\end{array}$$
(19a)

where

$$\sum_{i=1}^{r} k_{i} = M_{i} \quad k_{i} > 1, \quad r > 2, \quad |\alpha_{i}| > 1$$
 (19b)

(see (13d)). Since here $k_i \le N-1$ by (19b), the term (19a) can be estimated by induction assumption by

$$(|x| + 2)^{-r}(t - |x| + 2)^{-r}\log^{\mu}(t + 2)$$

with

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$$\mu = (2k_1 - 2) + (2k_2 - 2) + ... + (2k_r - 2) = 2u - 2r$$

Since $r \ge 2$ it follows that

$$\Box D^{\alpha} u_{\alpha} = 0((|\alpha| + 2)^{-2}(\epsilon - |\alpha| + 2)^{-2}\log^{2N-4}(\epsilon + 2))$$

Similarly for i = 1,2,3,4

is a linear combination of terms

$$(\mathbf{L}_{\mathbf{i}}^{\alpha_{1}}\mathbf{u}_{\mathbf{k}_{1}}^{\alpha_{1}})(\mathbf{D}^{\alpha_{2}}\mathbf{u}_{\mathbf{k}_{2}}) \dots (\mathbf{D}^{\alpha_{r}}\mathbf{u}_{\mathbf{k}_{r}})$$

satisfying (19b). Using the induction assumption we can estimate such a term by

$$(|x| + 2)^{-x}(t + 2)^{-1}(t - |x| + 2)^{1-x}\log^{y}(t + 2)$$

and find that

$$\mathbf{L}_{\mathbf{i}} \Box \mathbf{D}^{\mathbf{G}} \mathbf{u}_{\mathbf{N}} = \mathbf{0} ((|\mathbf{x}| + 2)^{-2} (\mathbf{t} - |\mathbf{x}| + 2)^{-1} (\mathbf{t} + 2)^{-1} \log^{2N-4} (\mathbf{t} + 2))$$

None

$$\mathbf{p}^{\alpha}\mathbf{u}_{\mathbf{N}} = \mathbf{v}_{\mathbf{N}} + \mathbf{p}^{-1} \mathbf{p}^{\alpha}\mathbf{u}_{\mathbf{N}}$$

where v_N is a solution of $v_N=0$ with the same initial values as $D^\alpha u_N$. Since all derivatives of u_N (including t-derivatives:) vanish for |x|>1, it follows from LEMMAS IV and V that u_N satisfies (18a,b,c).

We define \widetilde{u}_N by (13j). By (18d) there exists a constant Γ_N depending on f,g,a_{ik} such that

$$\left\{\widetilde{u}_{N}^{\prime}(t)\right\}_{5} \leq \Gamma_{N}|\varepsilon| \frac{\log^{2N-2}(t+2)}{t+2} \text{ for } |\varepsilon| \leq 1$$
 (20a)

and thus

$$\int_{0}^{T_{0}} \left\{ \widetilde{u}_{N}^{*}(s) \right\}_{5} ds \leq \Gamma_{N} |\varepsilon| \log^{2N-1}(T_{0} + 2)$$
 (20b)

We denote by En the set of & for which

$$C^{\circ}\Gamma_{n}|\varepsilon|\log^{2n-1}(\tau_{0}(\varepsilon)+2) < \log 2 \text{ for } n=1,2,...,H$$
 (20c)

$$|\varepsilon| < 1, \quad |\varepsilon| < \frac{\delta}{4cC} \tag{20d}$$

with C* as in LEMMA III and c as in (9c). We have

$$\exp\left(C^{4}\int_{0}^{\tilde{x}_{0}} \left\{\widetilde{u}_{N}^{1}(s)\right\}_{5}\right) ds < 2 \text{ for } n = 1, 2, ..., N$$
 (20e)

when & E Em. By definition

$$\mathbf{E}_{\mathbf{N}} \subset \mathbf{E}_{\mathbf{k}}$$
 for $\mathbf{k} \leq \mathbf{N}$ (20f)

LEDUCA VII.

Let ε, u, T_0 be as in LEMMA II. For given N there exists a constant γ (depending on f, g, a_{ik}, N) and an integer ν such that

$$\|(\mathbf{u}^{\mathsf{t}} - \widetilde{\mathbf{u}}_{\mathbf{u}}^{\mathsf{t}})(\mathsf{t})\|_{\mathsf{g}} \leq \gamma |\mathsf{g}|^{\mathsf{N}+1} \log^{\mathsf{V}}(\mathsf{t} + 2) \tag{21}$$

for $0 \le t \le T_0(\epsilon)$ and $\epsilon \in E_N$.

<u>Proof.</u> The proof proceeds by induction over M. (It is given for M=1,2,3 in (*), pp. 675-676). Using LEMMA III we have from (20e)

$$I(u^* - \tilde{u}_{H}^*)(t)I_{5} \le 8 \int_{0}^{T_{0}} I_{W_{H}}(s)I_{5}ds$$
 (22a)

where

$$w_{N} = P(u^{\epsilon}) \left[\widetilde{u}_{N}^{i} \right] = \sum_{j=1}^{M} P(u^{\epsilon}) \left[\epsilon^{j} u_{j} \right]$$
 (22b)

$$w_{N} = w_{N}^{+} + w_{N}^{++} + w_{N}^{+++}$$
 (22c)

with

$$w_{N}^{e} = \sum_{j=1}^{N} \left((P(u^{i}) - \Box - \sum_{k=1}^{N-j} P_{k}(u^{i})) [\epsilon^{j} u_{j}] \right)$$
 (22d)

$$w_{N}^{o+} = \sum_{j=1}^{N} \left(\Box + \sum_{k=1}^{N-j} P_{k}(\widetilde{u}_{N-j-k+1}^{*}) \right) \{ \varepsilon^{j} u_{j} \}$$
 (22e)

$$\mathbf{w}_{H}^{a+*} = \sum_{j=1}^{N} \sum_{k=1}^{N-j} \left(\mathbf{p}_{k}(\mathbf{u}^{*}) - \mathbf{p}_{k}(\widetilde{\mathbf{u}}_{H-j-k+1}^{*}) \right) [\epsilon^{j} \mathbf{u}_{j}]$$
 (222)

(This is the analogue of the decomposition for N=3 in (*), p. 676, formula (145)).

Setting

$$P(u^{*}) - \Box - \sum_{k=1}^{N-j} P_{k}(u^{*}) = R(u^{*}) = \sum_{i,k=1}^{3} b_{ik}(u^{*}) D_{i}D_{k}$$
 (22g)

we see that

$$D^{\alpha}(R(u^{i})[\epsilon^{j}u_{j}]) \qquad (22h)$$

is a linear combination of terms of the form

$$(\delta^{\beta}b_{ik}(u^{i}))(D^{\beta_{1}}u)(D^{\beta_{2}}u) \dots (D^{\beta_{2}}u)e^{j}D^{\gamma}u_{j}$$
 (221)

where δ stands for differentiation with respect to u^* and the multi-indices $\beta, \beta_1, \dots, \beta_x, Y$ satisfy

$$|\beta| = r > 0; \quad |\beta_1| + |\beta_2| + ... + |\beta_n| + |\gamma| = |\alpha| + r + 2$$
 (22j)

$$|\beta_1| > 2, |\beta_2| > 2, ..., |\beta_n| > 2, |\gamma| > 2$$
 (22k)

Here $b_{ik}(u^i)$ is the truncated Taylor expansion of $-a_{ik}(u^i)$ starting with terms of order W = j + 1. It follows then from the boundedness of the derivatives of the $a_{ik}(U)$ for $|U| < \delta$ that

$$\delta^{\beta}b_{jk}(u^{*}) = \begin{cases} 0(|u^{*}|^{M-j+1-r}) & \text{for } r \leq M-j+1 \\ 0(1) & \text{for } r > M-j+1 \end{cases}$$

By (22j,k) for $|\alpha| \le 5$ at most one of the $|\beta_k|$ can exceed 4. Hence

$$(D^{\beta_1}u)(D^{\beta_2}u) \dots (D^{\beta_r}u) = 0((|u^r|_3)^{r-1}|u^r|_5)$$
 for $r \ge 1$

It follows from (11g) that in all cases

$$D^{\alpha}(R(u^{*})\{\epsilon^{j}u_{j}\}) = 0(|\epsilon|^{N}|u^{*}|_{5}\{u_{j}^{*}\}_{6})$$

Hence by (112*), (18d)

$$\|u_{H}^{\bullet}\|_{5} = 0(|\epsilon|^{M+1} \frac{\log^{2M-2}(t+2)}{t+2})$$
 (221)

We turn to the estimate of $w_{H}^{a,a}$. By definition of the Q_{H} in (13c) we have the formal identities

$$0 = \sum_{M=1}^{n} \varepsilon^{M} Q_{M}(D^{\alpha} u_{1})$$

$$= \mathbb{P}(\varepsilon u_{1}^{i} + \varepsilon^{2} u_{2}^{i} + \dots) [\varepsilon u_{1} + \varepsilon^{2} u_{2} + \dots]$$

$$= \mathbb{Q}(\varepsilon u_{1} + \varepsilon^{2} u_{2} + \dots) + \sum_{k=1}^{n} \sum_{j=1}^{n} \varepsilon^{k+j} P_{k}(u_{1}^{i} + \varepsilon u_{2}^{i} + \varepsilon^{2} u_{3}^{i} + \dots) [u_{j}]$$

which yield actual finite identities when we collect the terms with equal powers of ε . Now the coefficients of ε^k with $k \le N$ will not be affected if we replace $u_1^s + \varepsilon u_2^s + \varepsilon^2 u_3^s + \dots$ in the last expression by the finite sum

$$u_1' + \varepsilon u_2' + \varepsilon^2 u_3' + \dots + \varepsilon^{N-k-j} u_{N-k-j+1}'$$

and restrict j to values < M and k to values < M - j. This means that

$$\begin{aligned} w_{H}^{\bullet,\bullet} &= \Box (\varepsilon u_{1} + \varepsilon^{2} u_{2} + \dots + \varepsilon^{M} u_{H}) \\ &+ \sum_{j=1}^{M} \sum_{k=1}^{M-j} \varepsilon^{j+k} p_{k} (u_{1}^{i} + \varepsilon u_{2}^{i} + \dots + \varepsilon^{M-k-j} u_{M-k-j+1}^{i}) \{u_{j}\} \end{aligned}$$

contains no terms ϵ^{m} with $m \le N$ and is equal to a linear combination of terms of the form

$$\epsilon^{N+1+i}(D^{\beta_1}u_{k_1})(D^{\beta_2}u_{k_2}) \cdots (D^{\beta_r}u_{k_r})$$

with

$$i > 0$$
, $r > 2$, $|\beta_1| > 1$,..., $|\beta_r| > 1$, $2 \le r \le H$.

The same holds for any $D^{\Omega} w_{\tilde{N}}^{e+}$. It follows from (18d,e) that

$$i w_{N}^{\bullet,\bullet}(t) i_{5} = 0 \left(\epsilon^{N+1} \sup_{i,j \leq N} \left(\left\{ u_{i}^{*}(t) \right\}_{5} \right)^{N-1} i u_{j}^{*}(t) i_{6} \right)$$

$$= 0 \left(\epsilon^{N+1} \frac{\log^{4N-4}(t+2)}{t+2} \right)$$
(22m)

Finally

$$P_{k}(u^{i}) - P_{k}(\widetilde{u}_{k-j-k+1}^{i})$$

is a form of degree k in $\widetilde{u}_{M-j-k+1}^t$ and $u^t-\widetilde{u}_{M-j-k+1}^t$, where each term contains at least one of the latter factors. Thus

$$\mathfrak{p}^{\alpha}((\mathfrak{P}_{k}(u^{*})-\mathfrak{P}_{k}(\widetilde{u}_{k-j-k+1}^{*}))[\epsilon^{j}u_{j}])$$

is a linear combination of terms of the form

$$\varepsilon^{j}(D^{\alpha_{1}}\widetilde{u}_{M-j-k+1})\cdots(D^{\alpha_{n}}\widetilde{u}_{M-j-k+1})(D^{\beta_{1}}(u-\widetilde{u}_{M-j-k+1}))\cdots(D^{\beta_{n}}(u-\widetilde{u}_{M-j-k+1}))(D^{\gamma_{n}}u_{j})$$

with

$$|\alpha_{\underline{1}}| > 1; \quad |\beta_{\underline{k}}| > 1; \quad |Y| > 2; \quad x + a = k; \quad |a| > 1; \quad H - j - k + 1 \le H - 1$$

$$|\alpha_{\underline{1}}| + \dots + |\alpha|_{\underline{r}} + |\beta_{\underline{1}}| + \dots + |\beta_{\underline{n}}| + |Y| = |\alpha| + 2 + x + a$$

For $|\alpha| \le 5$ it follows that at most one of the $|\beta_k|$ can exceed 3 and that none of them exceeds 6. Using (20a), (18d) and the induction assumption we find

$$\begin{split} & \mathbb{I}(\mathbb{P}_{k}(\mathbf{u}^{i}) - \mathbb{P}_{k}(\widetilde{\mathbf{u}}_{k-j-k+1}^{i}))[\epsilon^{j}\mathbf{u}_{j}](\epsilon)\mathbb{I}_{5} \\ & = 0(\epsilon^{x+j}(\{\mathbf{u}^{i} - \widetilde{\mathbf{u}}_{k-j-k+1}^{i})(\epsilon)\}_{3})^{s-1}\mathbb{I}(\mathbf{u}^{i} - \widetilde{\mathbf{u}}_{k-j-k+1}^{i})(\epsilon)\mathbb{I}_{5}(\mathbf{u}_{j}^{i}(\epsilon))_{6}) \\ & = 0(\epsilon^{x+j}(\mathbb{I}(\mathbf{u}^{i} - \widetilde{\mathbf{u}}_{k-j-k+1}^{i})(\epsilon)\mathbb{I}_{5})^{s}(\mathbf{u}_{j}^{i}(\epsilon))_{6}) \\ & = 0(\epsilon^{x+j+(N-j-k+2)s}(\epsilon + 2)^{-1}\log^{\mu}(\epsilon + 2)) \end{split}$$

with a certain μ . Here, since $s \ge 1$ and r + s = k,

$$x+j+(N-j-k+2)s = x+j+s+(N-j-k+1)s > x+j+s+(N-j-k+1) = N+1$$

It follows that

$$\|u_{N}^{***}(t)\|_{5} = 0(\epsilon^{N+1} \frac{\log^{\mu}(t+2)}{t+2})$$
.

Altogether then

$$\|\mathbf{w}_{N}(t)\|_{5} = 0(\epsilon^{N+1} \frac{\log^{N}(t+2)}{t+2})$$

with a certain μ . Hence by (22a)

$$\|u^{t} - \widetilde{u}_{h}^{t}\|(t)\|_{5} \approx 0(\epsilon^{H+1}\log^{\mu+1}(t+2))$$
.

This completes the proof of (21) by induction over N, provided we still verify the case N = 1. In that case w_1^{***} and $w_1^{***} = \Box \varepsilon u_1$ vanish. By (221)

$$||v_1^*||_5 = o(\epsilon^2(\epsilon + 2)^{-1})$$
,

and (22a) furnishes the desired relation

$$|u' - \widetilde{u}_1'|(t)|_{\xi} = 0(\epsilon^2 \log(t + 2))$$

LEDWA VIII

$$\lim_{\varepsilon \to 0} \left(\varepsilon \right)^{M} T_{0}(\varepsilon) = 0 \tag{23}$$

(This implies the THEOREM, since $T(\varepsilon) \ge T_0(\varepsilon)$).

<u>Proof.</u> Assume that (23) does not hold. There exists then a constant K and a sequence of
e, tending to zero for which

$$|\epsilon_j|^{N}T_0(\epsilon_j) \leq K$$
.

Then $\varepsilon = \varepsilon_j$ will satisfy (20c,d) for all sufficiently large j, and thus $\varepsilon_j \in E_{\mathbb{N}}$. By (21), (5d), (20a)

$$\{u^{*}(t)\}_{3} \leq \{\tilde{u}^{*}_{N}(t)\}_{3} + \gamma C|\varepsilon|^{N+1} \log^{V}(t+2)$$

$$\leq \Gamma_{N}|\varepsilon| \frac{\log^{2N-2}(t+2)}{t+2} + \gamma C|\varepsilon|^{N+1} \log^{V}(t+2)$$

for $\varepsilon = \varepsilon_j$ with j sufficiently large. It follows from (10d) that

$$\frac{\log 2}{C} = \int_{0}^{T_0} \left\{ u^{\tau}(\mathbf{s}) \right\}_{3} d\mathbf{s}$$

$$\leq \Gamma_{\mathbf{N}} |\varepsilon_{j}| \log^{2N-1} (\mathbf{T}_{0}(\varepsilon_{j}) + 2) + \gamma C |\varepsilon_{j}|^{N+1} \mathbf{T}_{0}(\varepsilon_{j}) \log^{V} (\mathbf{T}_{0}(\varepsilon_{j}) + 2)$$

$$\leq \Gamma_{\mathbf{N}} |\varepsilon_{j}| \log^{2N-1} (\mathbf{T}_{0}(\varepsilon_{j}) + 2) + \gamma C |\varepsilon_{j}| \log^{V} (\mathbf{T}_{0}(\varepsilon_{j}) + 2)$$
(24)

Since here

$$log(T_0(\varepsilon_j) + 2) \leq log(K|\varepsilon_j|^{-N} + 2);$$
 $lim_{j \leftrightarrow j} = 0$

(24) leads to a contradiction for j + . Thus (23) holds.

THE MATH LEMMA

Let $w(x,t) \in C^2$ for $x \in \mathbb{R}^3$, t > 0, and let

$$w(x,t) = 0$$
 for $|x| > t + 1$ (25a)

$$|w|, |D_1w| \le \frac{N \log^k(t+2)}{(|x|+2)^2(t-|x|+2)^2}$$
 for $|x| < t+1$ (25b)

$$|L_i w| \le \frac{H \log^k(t+2)}{(|x|+2)^2(t-|x|+2)(t+2)}$$
 for $|x| < t+1$ (25c)

with a certain k > 0 and L_i defined by (14a,b) for i = 1,2,3,4. Then

$$u = \Box^{-1} \mathbf{v} \tag{25d}$$

estisfies

$$u(x,t) = 0$$
 for $|x| > t + 1$ (25e)

$$|D_{i}u| \le \frac{AM \log^{k+2}(t+2)}{(|x|+2)(t-|x|+2)}$$
 for $|x| < t+1$ (25f)

$$|\mathbf{L}_{\mathbf{i}}\mathbf{u}| \le \frac{\mathbf{AM} \log^{k+2}(t+2)}{(|\mathbf{x}|+2)(t+2)} \text{ for } |\mathbf{x}| \le t+1$$
 (25g)

for i = 1,2,3,4 with a universal constant A.

Proof.

By Duhamel's principle

$$u(x,t) = \Box^{-1}w(x,t) = \int_{0}^{t} \omega(x,t,s)ds$$
 (26a)

where $\Theta(x,t,s)$ for $x \in \mathbb{R}^3$, t > s is the solution of

$$\square \omega(x,t,s) = 0 \text{ for } x \in \mathbb{R}^3, \quad 0 \le s \le t$$
 (26b)

$$w(x,t,s) = 0$$
, $w_{+}(x,t,s) = w(x,s)$ for $t = s$ (26c)

For w we have (see (16b)) the integral representation

$$w(x,t,s) = \frac{1}{4\pi(t-s)} \int_{|y-x|=t-s} w(y,s)ds_y$$
 (26d)

(25e) is an immediate consequence of (25e), (26d) since w(y,s) = 0 for |y| > s + 1, and

$$|y| > |x| - t + s > s + 1$$
.

for |y - x| = t - s, |x| > t + 1.

Since by (25a,b)

$$w(x,t) = 0\left(H \frac{\log^k(t+2)}{(t+2)^2}\right) = 0(H)$$
 (26e)

we have the trivial estimates

$$w = 0(M(t - s)), u = \Box^{-1}w = 0(Mt^2)$$
 (26f)

We observe that Diu is the solution of

with initial values

$$D_{i}u = 0$$
, $D_{d}D_{i}u = \delta_{id}w(x,0)$ for $t = 0$ (26g)

Thus

$$D_{i}u = \Box^{-1}D_{i}w + \delta_{i4}u^{0}$$
 (26g)

where u0 is the solution of

$$\exists u^0 = 0 \tag{26h}$$

$$u^0 = 0$$
, $u_t^0 = w(x,0)$ for $t = 0$ (261)

Since also $D_i w = O(M)$, we find in analogy to (26f)

$$D_{a}u = 0(Mt^{2} + Mt)$$
 (26j)

In particular $D_1 u = 0(M)$, if we prescribe a numerical upper bound for t, say t < 10. It follows that (25f,g) are satisfied for t < 10 since here |x| is restricted to values less than t + 1. In what follows we can assume that

Modified integral representation for u.

In (26e) introduce spherical coordinates θ, ϕ on the sphere $|y-x|\approx t-s$, with the polar axis pointing in the direction from x to 0, with θ = polar distance, ϕ = latitude. Then (see Figure 1)

$$\omega(x,t,s) = \frac{t-s}{4\pi} \int_{0}^{\pi} \sin \theta d\theta \int_{0}^{2\pi} w d\phi$$
 (28a)

Set

$$r = |x|; p = s - |y|; \theta = | (0xy); \psi = | (x0y); q = \cos \psi$$
 (29b)

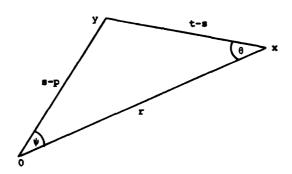


Figure 1

The ϕ -integration corresponds to y varying over the circle of intersection of the cone $y^*x = q|y||x|$ with the sphere |y| = s - p. Denote generally by j the average of w on that circle:

$$j(x,p,q,s) = \frac{1}{2\pi} \int_{\substack{y^*x=q|y||x|\\|y|=s-p}} w(y,s)d\phi$$
 (28c)

for $x \neq 0$, $|q| \leq 1$, $p \leq s$. Here ϕ is an angular measure on the circle (ϕ = arc length divided by radius). Then

$$\omega(x,t,s) = \frac{t-s}{2} \int_{0}^{\pi} j(x,p,q,s) \sin \theta d\theta \qquad (28d)$$

where p and q are the functions of θ defined by

$$p = s - \sqrt{(t - s)^2 + r^2 - 2(t - s)r \cos \theta}$$
 (28e)

$$q = Q(p,r,s,t) = \frac{(s-p)^2 + r^2 - (t-s)^2}{2r(s-p)}$$
 (28f)

Introducing p instead of θ as variable of integration in (28d) results in the expression

$$\omega(x,t,s) = \int_{\lambda}^{B} \frac{s-p}{2r} k(x,p,s,t) dp$$
 (28g)

where we define

$$k(x,p,s,t) = j(x,p,Q(p,|x|,s,t),s)$$
 (28h)

$$B = s - |s - t + r|$$
; $A = s - |t + r - s| = 2s - t - r$ (281)

We introduce new independent variables in w(y,s) which are better suited for describing <u>outgoing waves</u>. For that purpose we associate with w(y,s) the function w defined by

$$v(y,p,s) = w((s-p)\frac{y}{|y|}, s)$$
 for $y \neq 0, p \leq s, 0 \leq s$ (29a)

which is homogeneous of degree 0 in y. We have conversely

$$w(y,s) = v(y,s - |y|,s)$$
 for $y \neq 0$, $s > 0$ (29b)

Substituting this expression for w(y,s) into (28c) and replacing y by (s-p)y yields

$$j(x,p,q,s) = \frac{1}{2\pi} \int_{\substack{y \in x = q \mid x \mid \\ |y|=1}} v(y,p,s) d\phi$$
 (29c)

We observe that by assumption (25a)

$$v(y,p,s) = 0$$
 for $p < -1$ (29d)

and then also by (29c), (28h)

$$j(x,p,q,s) = 0$$
, $k(x,p,s,t) = 0$ for $p < -1$ (29e)

It follows from (28g,i) that

$$\omega(x,t,s) = \int_{-1}^{B} \frac{s-p}{2r} k(x,p,s,t) dp$$
 for $s < \frac{t+r-1}{2}$ (29f)

$$\omega(x,t,s)=0 \text{ for } s<\frac{t-r-1}{2} \tag{29g}$$

Derivatives of u. Direct estimates.

Let i = 1,2,3,4 be fixed. Set

$$W(x,t) = D_i w(x,t)$$
 (30a)

By (26g)

$$D_i u = \Box^{-1} w + \delta_{i4} u^0$$

Here (see (16c), (26e))

$$u^0 = 0(\frac{M}{t+2}) = 0(\frac{M}{(r+2)(t-r+2)})$$

since $u^0(x,t) \neq 0$ only for t-1 < r < t+1. Thus for the proof of (25f) we only have to show that

$$\Box^{-1}W = 0\left(M \frac{\log^{k+2}(t+2)}{(r+2)(t-r+2)}\right)$$
 (30b)

We represent $\Box^{-1}W$ in complete analogy to $u = \Box^{-1}W$, introducing

$$V(y,p,s) = W((s-p) \frac{y}{|y|}, s)$$
 (30c)

$$J(x,p,q,s) = \frac{1}{2\pi} \int_{\substack{y \in x = q \mid x \mid \\ |y \mid = 1}} V(y,p,s) d\phi$$
 (30d)

$$K(x,p,s,t) = J(x,p,Q(p,|x|,s,t),s)$$
 (30e)

$$\Omega(x,t,s) = \int_{A}^{B} \frac{s-p}{2r} \, K(x,p,s,t) \, dp \qquad (30f)$$

with A,B,Q as in (28f,i), so that

$$\Box^{-1}W(x,t) = \int_{0}^{t} \Omega(x,t,s)ds$$
 (30g)

For V we find from assumptions (25a,b) that

$$V(y,p,s) = 0$$
 for $p < 1$

$$V(y,p,s) = 0\left(\frac{M \log^{k}(s+2)}{(s-p+2)^{2}(p+2)^{2}}\right)$$
(30h)

which implies that

$$J(x,p,q,s), K(x,p,s,t) = 0 \text{ for } p < -1$$
 (301)

$$J(x,p,q,s), K(x,p,s,t) = 0\left(\frac{M \log^{k}(s+2)}{(s-p+2)^{2}(p+2)^{2}}\right)$$
(30j)

We can verify (30b) immediately for bounded t - |x|, say for

$$t - 3 < r < t + 1 \tag{31a}$$

For by (30j) with -1

$$\frac{s-p}{2r}K(x,p,s,t) = 0\left(\frac{M \log^{k}(s+2)}{r(s-p+4)}\right)$$

while by (281) B \leq s - (s - t + r) = t - r \leq 3. Thus, since here r $> \frac{1}{2}$ (t + 2) ,

$$\Omega(x,t,s) = 0 \left(\int_{-1}^{3} \frac{M \log^{k}(s+2)}{r(s-p+4)} dp \right) = 0 \left(\frac{M \log^{k}(t+2)}{(t+2)(s+1)} \right)$$

$$\Box^{-1}W = 0\left(\frac{M \cdot \log^{k+1}(t+2)}{t+2}\right) = 0\left(\frac{M \cdot \log^{k+2}(t+2)}{(x+2)(t-x+2)}\right)$$

for r satisfying (31a). Henceforth in the proof of (30b) we can assume that

$$0 < r < t - 3; t > 10$$
 (31b)

In this section we find the relevant estimate for $\Omega(x,t,s)$ in the case where

$$\frac{1}{2}$$
 (t + r - 1) \leq s \leq t (32a)

without making use of the radiation conditions (25c). (32a) implies that

$$-1 \le A = 2s - t - r \tag{32b}$$

By (30£,j)

$$\Omega(x,t,s) = 0 \left(\int_{A}^{B} \frac{M \log^{k}(t+2)}{r(s-p+2)(p+2)^{2}} dp \right)$$

$$= 0 \left(\frac{M \log^{k}(t+2)}{r(s+4)^{2}} \left(\frac{(B-A)(S+4)}{(A+2)(B+2)} + \log \frac{(s-A+2)(B+2)}{(s-B+2)(A+2)} \right) \right)$$
(32c)

Here

$$-1 \le \lambda \le B = s - |t - r - s| \le s$$

$$0 \le B - \lambda = 2r + (t - r - s) - |t - r - s| \le 2r$$

$$\frac{B - \lambda}{B + 2} \le \min(1, 2r) = 0 \left(\frac{r}{r + 2}\right)$$

$$\log \frac{(s - \lambda + 2)(B + 2)}{(s - B + 2)(\lambda + 2)} = \log\left(\left(1 + \frac{B - \lambda}{s - B + 2}\right)\left(1 + \frac{B - \lambda}{\lambda + 2}\right)\right)$$

$$\le 2r\left(\frac{1}{s - B + 2} + \frac{1}{\lambda + 2}\right)$$

$$\frac{1}{s + 4} \le \frac{2}{t + r + 7} = 0\left(\frac{1}{t + 2}\right) = 0\left(\frac{1}{r + 2}\right)$$

Hence

$$\Omega(x,t,s) = 0\left(\frac{M \log^{k}(t+2)}{(r+2)(t+2)}\left(\frac{1}{|t-r-s|+2} + \frac{1}{2s-t-r+2}\right)\right)$$

and thus

$$\int_{2}^{t} \Omega(x,t,s)ds = 0\left(\frac{M \log^{k+1}(t+2)}{(r+2)(t+2)}\right)$$
(32d)

Derivative of u. Integration by parts.

To complete the proof of (25f) we need in addition to (32d) that

$$I = \int_{\frac{\pi}{2}}^{\frac{\pi}{2}} \frac{(t+r-1)}{(t-r-1)} \Omega(x,t,s) ds = 0\left(\frac{M \log^{k+2}(t+2)}{(r+2)(t+2)}\right)$$
(33)

(Note that as in (29g)

$$\Omega(x,t,s) = 0$$
 for $s < \frac{1}{2}(t-r-1)$ (34a)

since then B < -1, and that

$$\frac{1}{2}(t-r-1) > 1$$
 (34b)

by (31b)). The considerations leading to (32d) are insufficient for the proof of (33). The estimate (32c) still holds when A is replaced by -1, but is not good enough to yield (33), and we have to have recourse to Kore complicated estimates involving the raditation conditions (25c). In this section we restrict ourselves to values s with

$$\frac{1}{2}(t-r-1) \le s \le \frac{1}{2}(t+r-1)$$
 (34c)

The relation $W=D_{\underline{i}} w$ between w and W yields relations between w and V, namely

$$V(y,p,a) = \frac{|y|}{a-p} v_{y_i}(y,p,a) - \frac{y_i}{|y|} v_p(y,p,a)$$
 when $i = 1,2,3$ (34d)

$$V(y,p,s) = v_{g}(y,p,s) + v_{p}(y,p,s)$$
 when i = 4 (34e)

In order to unify the arguments we introduce

$$V^{\pm}(y,p,s) = \frac{|y|}{s-p} v_{y_{1}}(y,p,s); \quad v^{\pm}(y,p,s) = -\frac{y_{1}}{|y|} v(y,p,s) \quad \text{when} \quad i = 1,2,3$$

$$V^{\pm}(y,p,s) = v_{g}(y,p,s); \quad v^{\pm}(y,p,s) = v(y,p,s) \quad \text{when} \quad i = 4$$

so that V° and V° are homogeneous of degree 0 in y and

$$V(y,p,s) = V^*(y,p,s) + V^*_p(y,p,s)$$
 (35a)

In analogy to previous notation we set again

$$J^{+}(x,p,q,s) = \frac{1}{2\pi} \int_{\substack{y^{+} x = q \mid x \mid \\ |y| = 1}} V^{+}(y,p,s) d\phi$$
 (35b)

$$j^{+}(x,p,q,s) = \frac{1}{2\pi} \int_{\substack{y^{+}x=q|x|\\|y|=1}} v^{+}(y,p,s)d\phi$$
 (35c)

$$K^{+}(x,p,s,t) = J^{+}(x,p,Q(p,|x|,s,t),s)$$
 (35d)

$$k^{*}(x,p,s,t) = j^{*}(x,p,Q(p,|x|,s,t),s)$$
 (35e)

Then by (35a), (30d), (30e)

$$J(x,p,q,s) = J^{+}(x,p,q,s) + j_{p}^{+}(x,p,q,s)$$
 (35f)

$$K(x,p,s,t) = K^*(x,p,s,t) + k_p^*(x,p,s,t)$$
 (35g)

$$-Q_{p}(p,r,s,t)j_{q}^{*}(x,p,Q(p,r,s,t),s)$$
 (35g)

Consequently (see (29f)), using that $k^*(x,-1,s,t) = 0$,

$$\begin{split} \Omega(\mathbf{x}, \mathbf{t}, \mathbf{s}) &= \int_{-1}^{B} \frac{\mathbf{s} - \mathbf{p}}{2\mathbf{r}} \left(\mathbf{K}^{a} + \mathbf{k}_{\mathbf{p}}^{b} - Q_{\mathbf{p}} \mathbf{j}_{\mathbf{q}}^{a} \right) d\mathbf{p} \\ &= \frac{\mathbf{s} - \mathbf{B}}{2\mathbf{r}} \mathbf{k}^{a} (\mathbf{x}, \mathbf{B}, \mathbf{s}, \mathbf{t}) + \int_{-1}^{B} \left(\frac{\mathbf{s} - \mathbf{p}}{2\mathbf{r}} \mathbf{K}^{a} (\mathbf{x}, \mathbf{p}, \mathbf{s}, \mathbf{t}) + \frac{1}{2\mathbf{r}} \mathbf{k}^{a} (\mathbf{x}, \mathbf{p}, \mathbf{s}, \mathbf{t}) \right) d\mathbf{p} \\ &- \int_{-1}^{B} \frac{\mathbf{s} - \mathbf{p}}{2\mathbf{r}} Q_{\mathbf{p}} (\mathbf{p}, \mathbf{r}, \mathbf{s}, \mathbf{t}) \mathbf{j}_{\mathbf{q}}^{a} (\mathbf{x}, \mathbf{p}, \mathbf{Q}(\mathbf{p}, \mathbf{r}, \mathbf{s}, \mathbf{t}), \mathbf{s}) d\mathbf{p} \end{split} \tag{35h}$$

One verifies easily from (29a) and the definition of V* that

$$V^{+}(y,p,a) = L_{\underline{i}}w((a-p)\frac{y}{|y|},a)$$
 (36a)

for i = 1, 2, 3, 4, and that

$$v_{p}(y,p,s) = -\sum_{k=1}^{3} \frac{y_{k}}{|y|} w_{x_{k}}((s-p) \frac{y}{|y|},s)$$
 (36b)

It follows from assumptions (25b,c) that

$$V^{+}(y,p,s) = 0\left(\frac{M \log^{k}(s+2)}{(s-p+2)^{2}(s+2)(p+2)}\right)$$
(36c)

$$v^*(y,p,s), v^*_p(y,p,s) = 0 \left(\frac{M \log^k(s+2)}{(s-p+2)^2(p+2)^2} \right)$$
 (36d)

Finally we obtain from

$$v_{y_k}(y,p,s) = \frac{s-p}{|y|} I_k w((s-p) \frac{y}{|y|},s)$$
 (36e)

and the definition of v* that

$$v_{y_{k}}^{+}(y,p,s) = 0\left(\frac{1}{|y|} |w| + \frac{s-p}{|y|} |L_{k}w|\right)$$

$$= 0\left(\frac{M \log^{k}(s+2)}{|y|(s-p+2)(p+2)} \left(\frac{1}{(s-p+2)(p+2)} + \frac{1}{s+2}\right)\right)$$

$$= 0\left(\frac{M \log^{k}(s+2)}{|y|(s-p+2)(s+2)(p+2)}\right)$$
(36f)

Relations (36c,d), (35b,c,d,e) imply that

$$J^{+}(x,p,q,s), K^{+}(x,p,s,t) = 0\left(\frac{H \log^{k}(s+2)}{(s-p+2)^{2}(s+2)(p+2)}\right)$$
(36g)

$$j^*(x,p,q,s), k^*(x,p,s,t) = 0\left(\frac{H \log^k(s+2)}{(s-p+2)^2(p+2)}\right)$$
 (36h)

We first consider the contribution to I of the term

$$\alpha = \frac{a - B}{2v} k^{\alpha}(x, B, a, t)$$
 (37a)

in (35h). By (36h)

$$\alpha = 0\left(\frac{H \log^{k}(t+2)}{r(s-B+2)(B+2)^{2}}\right) = 0\left(\frac{H \log^{k}(t+2)}{r(s+4)^{2}}\left(\frac{1}{s-B+2} + \frac{1}{B+2} + \frac{s+4}{(B+2)^{2}}\right)\right)$$
(37b)

where

$$-1 \le B = s - |s-t+r| \le s$$
; $\frac{1}{2}(t-r-1) \le s \le \frac{1}{2}(t+r-1)$; $t-r > 3$

Thus

$$\frac{1}{x+4} = 0(\frac{1}{x-x+2})$$

If here r < 1 we have from (37b)

$$\alpha = 0\left(\frac{M \log^k(t+2)}{r(t-r+2)}\right)$$

$$\int_{\frac{1}{2}}^{\frac{1}{2}} \frac{(t+r-1)}{\int_{\frac{1}{2}}^{\frac{1}{2}} (t-r-1)} ads = 0 \left(\frac{\frac{M \log^{k}(t+2)}{(t-r+2)}}{(t-r+2)} \right) = 0 \left(\frac{\frac{M \log^{k}(t+2)}{(r+2)(t-r+2)}}{(r+2)(t-r+2)} \right)$$
(37c)

If, on the other hand, r > 1 we have

$$\frac{1}{x(s+4)^2} \left(\frac{1}{s-3+2} + \frac{1}{3+2} + \frac{s+4}{(3+2)^2} \right)$$

$$= 0 \left(\frac{1}{x+2} \left(\frac{1}{(s+4)^2} + \frac{1}{(s+4)(t-x+2)} + \frac{1}{(t-x+2)(2s-t+x+2)} \right) \right)$$

and hence

$$\frac{\frac{1}{2}(t+r-1)}{\int_{\frac{1}{2}(t-r-1)}^{\frac{1}{2}(t-r-1)}} ads = 0\left(\frac{\frac{11\log^{k+1}(t+2)}{(r+2)(t-r+2)}}{(37d)}\right)$$

We next consider

$$\beta = \int_{-1}^{B} \left(\frac{s-p}{2r} \, K^{a}(x,p,s,t) + \frac{1}{2r} \, k^{a}(x,p,s,t) \right) dp \tag{38a}$$

Here by (36g,h)

$$\frac{s-p}{2r} g^{a} + \frac{1}{2r} k^{a} = 0 \left(\frac{H \log^{k}(t+2)}{r(s-p+2)(p+2)} \left(\frac{1}{s+2} + \frac{1}{(s-p+2)(p+2)} \right) \right)$$

$$= 0 \left(\frac{H \log^{k}(t+2)}{r(s-p+2)(p+2)(s+2)} \right)$$

$$= 0 \left(\frac{H \log^{k}(t+2)}{r(s+2)(s+4)} \left(\frac{1}{s-p+2} + \frac{1}{p+2} \right) \right)$$

since $p \le B \le s$. Hence by (34c)

$$\beta = 0 \left(\frac{M \log^{k+1}(t+2)}{r(s+2)(s+4)} \right)$$

$$\beta = 0 \left(\frac{M \log^{k+1}(t+2)}{r(t-r+2)} \right) \text{ for } r < 1$$

$$\beta = 0 \left(\frac{M \log^{k+1}(t+2)}{r(s+2)^2(r+2)} \right) \text{ for } r > 1$$

In either case

$$\frac{\frac{1}{2}(t+r-1)}{\int_{2}^{\infty}(t-r-1)} \beta ds = 0\left(\frac{M \log^{k+1}(t+2)}{(r+2)(t-r+2)}\right)$$
(38b)

This leaves

$$Y = \int_{-1}^{B} \frac{s - p}{2\pi} Q_{p}(p,r,s,t) j_{q}^{*}(x,p,Q(p,r,s,t),s) dp$$
 (39a)

For j* given by (35c) one easily derives the differentiation formula

$$j_{\mathbf{q}}^{*}(\mathbf{x}, \mathbf{p}, \mathbf{q}, \mathbf{s}) = \frac{1}{2\pi \mathbf{r}(1 - \mathbf{q}^{2})} \int_{\substack{\mathbf{y} \in \mathbf{q} \\ |\mathbf{y}| = 1}}^{3} (\mathbf{x}_{k} - \mathbf{r} \mathbf{q} \mathbf{y}_{k}) \mathbf{v}_{k}^{*}(\mathbf{y}, \mathbf{p}, \mathbf{s}) d\mathbf{p}$$

Here for $x \cdot y = qr$, |y| = 1

$$|x_k - rqy_k| \le |x - rqy| = r\sqrt{1 - q^2}$$
.

It follows from (36f) that

$$j_{q}^{+}(x,p,q,s) = 0\left(\frac{M \log^{k}(s+2)}{(s-p+2)(s+2)(p+2)} (1-q^{2})^{-1/2}\right)$$
(39b)

By definition (28f) of Q(p,r,s,t) we have

$$\sqrt{1-Q^2(p,r,s,t)} = \frac{R}{2r(s-p)}$$
 (39c)

where

$$R = \sqrt{(t+r-p)(t-r-p)(t+r-2s+p(2s-t+r-p))}$$
 (39d)

is Heron's expression for four times the area of the triangle with vertices 0,x,y; (see Figure 1). Moreover

$$Q_p(p,r,s,t) = \frac{8}{2r(s-p)^2}$$
 (39e)

where

$$s = -(s - p)^2 + r^2 - (t - s)^2$$
 (39f)

Consequently

$$\gamma = 0 \left(\int_{-1}^{B} \frac{N \log^{k}(t+2)}{r(s-p+2)(s+2)(p+2)} \frac{s}{R} dp \right)$$
 (39g)

Writing

$$S = -2(s - p)^{2} + (t - p)(2s-t+r-p)-r(t-r-p)$$
 (40a)

and using that

$$0 < s-B < s-p < t-r-p;$$
 $0 < B-p = 2s-t+r-p < t-r-p$ for $s < t-r$
 $0 < s-B < s-p < 2s-t+r-p;$ $0 < B-p = t-r-p < 2s-t+r-p$ for $s > t-r$

we see that

$$g = 0((t + 2)(t-r-p))$$
 for $s < t-r$ (40b)

$$s = 0((t + 2)(2s-t+r-p))$$
 for $s > t-r$ (40c)

We first take up the case

$$s > t-r$$
 (40d)

This case only occurs when

$$r > \frac{t+1}{3} \tag{40e}$$

because of (34c). Here

$$B = t-r; \frac{1}{t+r-p} = 0(\frac{1}{t+2}); s-p > \frac{1}{2}(2s-t+r-p)$$

Hence by (40c)

$$= 0\left(\frac{t+2}{s+4}\left(\frac{1}{s-p+2} + \frac{1}{p+2}\right) \sqrt{\frac{2s-t+r-p}{(t+r-p)(t-r-p)(t+r-2s+p)}}\right)$$

$$= 0\left(\frac{\sqrt{t+2}}{(s+4)\sqrt{t+r-2s-1}}\left(\frac{1}{\sqrt{(t-r-p)(2s-t+r-p+4)}} + \frac{\sqrt{2s-t+r+1}}{(p+2)\sqrt{t-r-p}}\right)\right)$$

Hence by (39g), (40e)

$$Y = 0\left(\frac{M\sqrt{t+2} \log^{k+1}(t+2)}{r(s+2)^2\sqrt{t+r-2s-1}} \left(1 + \sqrt{\frac{2s-t+r+1}{t-r+2}}\right)\right)$$
$$= 0\left(\frac{M \log^{k+1}(t+2)}{\sqrt{(t+2)(t-r+2)}} \frac{\sqrt{2s-t+r+1}}{(s+2)^2\sqrt{t+r-2s-1}}\right)$$

With the substitution

$$s + 2 = \frac{(t+r+3)(t-r+3)(1+v^2)}{2(t+r+3) + 2(t-r+3)v^2}$$

we find

$$\int_{t-r}^{1/2} \gamma ds = 0 \frac{M \cdot \log^{k+1}(t+2)}{\sqrt{(t+2)(t-r+2)}} r(t-r+3)^{-1/2} (t+r+3)^{-3/2} \int_{a}^{\infty} \frac{v^2 dv}{(1+v^2)^2}$$

$$= 0 \left(\frac{M \cdot \log^{k+1}(t+2)}{(r+2)(t-r+2)} \right)$$
(40f)

with

$$a = \sqrt{\frac{(t+r+3)(t-r+1)}{(t-r+3)(3r-t-1)}}$$

We come to the contribution of the values s with [see (34c)]

$$\frac{1}{2}(t-r-1) < a < D = Min(\frac{1}{2}(t+r-1), t-r)$$
 (41a)

Here

$$S = O((t + 2)(t-r-p)), S = 2s-t+r$$

$$\frac{t-r-p}{t+r-p} < \frac{t-r+1}{t+r+1} = 0(\frac{t-r+1}{t+2})$$

t+r-2s+p > t+r-2s-1

Hence by (39g)

$$Y = 0 \left(\int_{-1}^{B} \frac{H \log^{k}(t+2)}{r(s+2)(s-p+2)(p+2)} \sqrt{\frac{(t+2)(t-r+1)}{(t+r-2s-1)(2s-t+r-p)}} \, dp \right)$$

$$= 0 \left(\frac{H \log^{k}(t+2)}{r(s+2)(s+4)} \sqrt{\frac{(t+2)(t-r+1)}{t+r-2s-1}} \, g \right)$$
(41b)

with

$$H = \int_{-1}^{2s-t+r} \left(\frac{1}{s-p+2} + \frac{1}{p+2}\right) \sqrt{\frac{1}{2s-t+r-p}} \, dp$$

$$= \frac{1}{\sqrt{t-r-s+2}} \int_{0}^{a} \frac{d\theta}{(1+\theta)\sqrt{\theta}} + \frac{1}{\sqrt{2s-t+r+2}} \int_{0}^{b} \frac{d\theta}{(1-\theta)\sqrt{\theta}}$$

$$= 0\left(\sqrt{\frac{a}{(a+1)(t-r-s+2)}} + \sqrt{\frac{b}{2s-t+r+2}} \log \frac{2}{1-b}\right)$$
(41c)

$$a = \frac{2s-t+r+1}{t-r-s+2}$$
, $b = \frac{2s-t+r+1}{2s-t+r+2}$ (41d)

Here by (41a)

$$a = \frac{2r - (t+r-1-2s)}{2 + (t-r-s)} < r; \quad \frac{a}{a+1} < \frac{r}{r+1}$$

$$b = \frac{2r - (t+r-1-2s)}{2r+1 - (t+r-1-2s)} < \frac{2r}{2r+1}$$

$$b \log \frac{2}{1-b} < \frac{2r}{2r+1} \log(4r+2)$$

so that

$$\gamma = 0 \left(\frac{\text{M log}^{k+1}(t+2)}{(s+2)(s+4)} \sqrt{\frac{(t+2)(t-r+1)}{r(r+1)(t+r-2s-1)}} \left(\sqrt{\frac{1}{t-r-s+2}} + \sqrt{\frac{1}{2s-t+r+2}} \right) \right)$$

In the special case 0 < r < 1 we have by (41a)

 $D = \frac{1}{2}(t+r-1); s+2 > \frac{1}{2}(t+2); t-r-s+2 > \frac{1}{2}(t+2); 2s-t+r+2 > 1$

Hence

$$\gamma = 0\left(\frac{M \log^{k+1}(t+2)}{(t+2)\sqrt{r(t+r-2s-1})}\right)$$

and thus

$$\int_{\frac{1}{2}}^{\frac{1}{2}} \int_{(t-r-1)}^{(t+r-1)} \gamma ds = 0 \left(\frac{H \log^{k+1}(t+2)}{t+2} \right) = 0 \left(\frac{H \log^{k+1}(t+2)}{(r+2)(t+2)} \right)$$
(41e)

In the remaining case 1 < r < t-3 we have

$$\gamma = 0 \left(\frac{\text{M log}^{k+1}(t+2)}{(r+2)(s+2)} \sqrt{\frac{t+2}{(t-r+2)(t+r-2s-1)}} \left(\sqrt{\frac{1}{t-r-s+2}} + \sqrt{\frac{1}{2s-t+r+2}} \right) \right)$$

Here

$$\int_{\frac{1}{2}(t-r-1)}^{D} \frac{1}{s+2} \sqrt{\frac{1}{(t+r-2s-1)(2s-t+r+2)}} ds$$

$$< \int_{\frac{1}{2}(t-r-1)}^{\frac{1}{2}(t+r-1)} \frac{1}{s+2} \sqrt{\frac{1}{(t+r-2s-1)(2s-t+r+1)}} ds$$

$$= 0\left(\sqrt{\frac{1}{(t+r+3)(t-r+3)}}\right) = 0\left(\sqrt{\frac{1}{(t+2)(t-r+2)}}\right) \tag{442}$$

We need a similar estimate for

$$\int_{\frac{1}{2}(t-r-1)}^{D} \frac{1}{s+2} \sqrt{\frac{1}{(t+r-2s-1)(t-r-s+2)}} ds = G$$
 (44g)

Let at first 1 < r < t/2. Then $s+2 > \frac{1}{2}(t-r+3) > (t+2)/8$, and

$$G = \frac{1}{t+2} \int_{\frac{1}{2}(t-r-1)}^{D} \sqrt{\frac{1}{(t+r-2s-1)(t-r-s+2)}} ds = 0 \left(\frac{1}{t+2} \int_{a}^{b} \frac{d\theta}{\sqrt{\theta(1+\theta)}}\right)$$

-1+h

$$a = 0, b = \frac{2r}{t-3r+5} \text{ for } 1 < r < \frac{t+1}{3}$$

$$a = \frac{3r-t-1}{t-3r+5}, b = \frac{2r}{t-3r+5} \text{ for } \frac{t+1}{3} < r < \frac{t+5}{3}$$

$$a = \frac{4}{3r-t-5}, b = \frac{t-r+5}{3r-t-5} \text{ for } \frac{t+5}{3} < r < \frac{t}{2}$$

In each of the three sub-cases

$$G = 0\left(\frac{\log(t+2)}{t+2}\right)$$

as is easily verified. If instead t/2 < r < t-3 we have

$$D = t-r$$
; $t+r-2s-1 > t+r-2(t-r)-1 = 3r-t-1 > \frac{1}{4}(t+2)$

$$G = O\left(\frac{1}{\sqrt{t+2}} \int_{\frac{1}{2}(t-r-1)}^{t-r} \frac{ds}{(s+2)\sqrt{t-r-s+2}}\right) = O\left(\frac{1}{\sqrt{(t-r+4)(t+2)}} \int_{a}^{b} \frac{d\theta}{(1-\theta)\sqrt{\theta}}\right)$$

with

$$0 < a = \frac{2}{t-r+4} < b = \frac{t-r+5}{2(t-r+4)} < \frac{4}{7}$$

Here also

$$G = 0(\sqrt{\frac{1}{(t+2)(t-r+2)}})$$

Altogether then for 1 < r < t-3

$$\int_{V_2(t-r-1)}^{D} \gamma ds = 0 \left(\frac{M \log^{k+2}(t+2)}{(r+2)(t-r+2)} \right)$$

This, together with (41e), (40f), (38b), (37d), (37c), (32d) completes the proof of (25f).

Proof of the radiation conditions.

We notice that (25f) implies (25g) when $|x| = r < \frac{1}{2}t$. Thus in the proof of (25g) we can restrict ourselves to the case

$$\frac{1}{2}t < r < t+1$$
 (45a)

By (26d) and the assumption $w \in C^2$ we see that

$$w(x,ts) = \frac{t-s}{4\pi} \iint_{|\xi|=1} w(x + (t-s)\xi,s)dS_{\xi}$$

belongs to C^2 in x,t,s for $0 \le s \le t$. Because of (26a), (26c)

$$L_{\underline{i}}u = \int_{0}^{t} L_{\underline{i}}w(x,t,s)ds \tag{45b}$$

Since

$$L_{i}t = L_{i}r = \delta_{i4} \tag{45c}$$

we find from (28g), (28i) that

$$L_{\underline{i}}u = \delta_{\underline{i}4} \int_{0}^{t} \frac{s - h}{2r} k(x,h,s,t) ds - \delta_{\underline{i}4} \int_{0}^{t} ds \int_{\underline{A}}^{\underline{B}} \frac{s - p}{2r^{2}} k(x,p,s,t) dp$$

$$+ \int_{0}^{t} ds \int_{\underline{A}}^{\underline{B}} \frac{s - p}{2r} L_{\underline{i}}k(x,p,s,t) dp \qquad (45d)$$

Take first the case i = 1,2,3. By (28h), (45c), (14a)

$$L_{ik} = \sum_{m=1}^{3} (\delta_{im} - r^{-2}x_{i}x_{m})j_{x_{m}}(x,p,q,s)$$
 (45e)

Here by (29c), (36e), (25c)

$$j_{\mathbf{x}}(\mathbf{x}, \mathbf{p}, \mathbf{q}, \mathbf{s}) = \frac{1}{2\pi} \int_{\substack{\mathbf{y}^* \mathbf{x} = \mathbf{q}\mathbf{r} \\ |\mathbf{y}| = 1}} \left(\frac{\mathbf{q}}{\mathbf{r}} \mathbf{v}_{\mathbf{y}} - \sum_{n=1}^{3} \mathbf{r}^{-2} \mathbf{y}_{\mathbf{n}} \mathbf{x}_{n} \mathbf{v}_{\mathbf{y}_{n}} \right) d\phi \\
 = 0 \left(\frac{\mathbf{H} \cdot \log^{k}(\mathbf{s} + 2)}{\mathbf{r}(\mathbf{s} - \mathbf{p} + 2)(\mathbf{p} + 2)} \right)$$
(452)

It follows from (45d), (45a) that for i = 1,2,3

$$L_{i}u = 0\left(\int_{0}^{t} ds \int_{A}^{B} \frac{M \log^{k}(s+2)}{r^{2}(s+2)(p+2)} dp\right)$$

$$= 0\left(\frac{M \log^{k+2}(t+2)}{(r+2)(t+2)}\right) \tag{45g}$$

(More precisely the p-integration is taken over that portion of the interval (A,B) in which p > -1).

We turn to the case i = 4. By (28c,h), (25b)

$$\frac{s-p}{k} k(x,p,s,t) = 0 \left(\frac{M \log^{k}(s+2)}{r(s-p+2)(p+2^{2})} \right)$$

$$= 0 \left(\frac{M \log^{k}(t+2)}{r(s+4)^{2}} \left(\frac{1}{s-p+2} + \frac{1}{p+2} + \frac{s+4}{(p+2)^{2}} \right) \right)$$

It follows that

$$\int_{0}^{t} ds \int_{A}^{B} \frac{s-p}{2r^{2}} k(x,p,s,t) dp$$

$$= 0 \left(\frac{M \log^{k+1}(t+2)}{r^{2}} \right) = 0 \left(\frac{M \log^{k+1}(t+2)}{(r+2)(t+2)} \right)$$
(46a)

Similarly

$$\frac{s-h}{2r} k(x,h,s,t) = 0 \text{ for } h = 2s-t-r < -1$$

$$\frac{s-h}{2r} k(x,h,s,t) = 0 \left(\frac{H \log^{k}(t+2)}{r(t+r-s+2)(2s-t-r+2)^{2}} \right)$$

$$= 0 \left(\frac{H \log^{k}(t+2)}{r(r+1)(2s-t-r+2)^{2}} \right) \text{ for } h = 2s-t-r > -1$$

It follows that

$$\int_{0}^{t} \frac{s-\lambda}{2r} k(x,\lambda,s,t) ds = 0 \left(\frac{N \log^{k}(t+2)}{(r+2)(t+2)} \right)$$
 (46b)

This leaves the last term in (45d). By (14b), (28h)

$$L_{4}k = (Q_{x} + Q_{y})j_{q}(x,p,Q,z) + \sum_{m=1}^{3} x^{-1}x_{m}j_{x_{m}}$$
 (47a)

The contribution to $L_4 u$ of the terms with j_{R_2} is again of order

$$0(\frac{M \log^{k+2}(t+2)}{(r+2)(t+2)})$$

as follows from (45f), (45a). We estimate next

$$(Q_r + Q_t)j_q = -\frac{(2s-t+r-p)(t-r-p)}{2(s-p)r^2}j_q$$
 (47b)

We find from (39b) (which applies just as well to j_q as to j_{σ}^*), and (39c,d) that

$$\frac{s-p}{r} (Q_r + Q_t) j_q = 0 \left(\frac{M \log^k(s+2)}{r^2(s+2)(p+2)} \sqrt{\frac{(t-r-p)(2s-t+r-p)}{(t+r-p)(t+r-2s+p)}} \right)$$

Here

2s-t+r-p < t+r-p; t-r-p < t-r+1

If also $s < \frac{1}{2}(t+r-1)$ we have t+r-2+p > t+r-2s-1, $\lambda < -1$, and

$$\int_{-1}^{B} \frac{s-p}{r} (Q_{r} + Q_{t}) j_{q} dp = 0 \left(\frac{M \log^{k+1}(t+2)}{r^{2}(s+2)} \sqrt{\frac{t-r+1}{t+r-2s-1}} \right)$$

$$\int_{0}^{1/2} (t+r-1) ds \int_{-1}^{B} \frac{s-p}{r} (Q_{r} + Q_{t}) j_{q} dp = 0 \left(\frac{M \log^{k+2}(t+2)}{r^{2}} \sqrt{\frac{t-r+1}{t+r+3}} \right)$$

$$= 0 \left(\frac{M \log^{k+2}(t+2)}{(r+2)(t+2)} \right)$$

If instead $\frac{1}{2}(t+r-1) < s < t$ we have $s+2 > \frac{1}{2}(t+3)$, A > -1 and

$$\frac{s-p}{r} (Q_r + Q_t) j_q = 0 \left(\frac{M \log^k(t+2)}{r^2(t+2)(p+2)} \sqrt{\frac{t-r+1}{p-A}} \right)$$

$$\int_A^B \frac{s-p}{r} (Q_r + Q_p) j_q dp = 0 \left(\frac{M \log^k(t+2)}{r^2\sqrt{t+2}} \right) \int_A^\infty \frac{dp}{(p+2)\sqrt{p-A}}$$

$$= 0 \left(\frac{M \log^k(t+2)}{r^2\sqrt{(t+2)(A+2)}} \right)$$

$$\int_A^t ds \int_A^B \frac{s-p}{r} (Q_r + Q_t) j_q dp \approx 0 \left(\frac{M \log^k(t+2)}{r^2} \right)$$

$$= 0 \left(\frac{M \log^k(t+2)}{(r+2)(t+2)} \right)$$

This completes the proof of the radiation conditions (25g) and of the MAIN LEMMA.

REFERENCES

- [1] S. Klainerman. Global existence for nonlinear wave equations. Comm. Pure Appl.

 Mathematics 33 (1980) 43-101.
- [2] J. Shatah. Global existence of small solutions to nonlinear evolution equations.

 (Preprint)
- [3] S. Klainerman and G. Ponce. Global, small amplitude solutions to nonlinear evolution equations. (Preprint)
- [4] F. John. Blow-up for quasi-linear wave equations in three space dimensions. Comm.

 Pure Appl. Mathematics 34 (1981), 29-51.

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The paper deals with strict solutions $u(x,t) = u(x_1,x_2,x_3,t)$ of an equation

$$u_{tt} - \sum_{i,k=1}^{3} a_{ik}(Du)u_{x_ix_k} = 0$$

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where Du is the set of 4 first derivatives of u. For given initial values $u(x,0)=\epsilon F(x),\ u_t(x,0)=\epsilon G(x)$ the life span $T(\epsilon)$ is defined as the supremum of all t to which the local solution can be extended for all x. Blow-up in finite time corresponds to $T(\epsilon)<\infty$. Examples show that this can occur for arbitrarily small ϵ . On the other hand $T(\epsilon)$ must at least be very large for small ϵ . Assuming that a_{ik} , $F,G\in C$, that $a_{ik}(0)=\delta_{ik}$, and that F,G have compact support, it is shown that $\lim_{\epsilon\to 0} \epsilon^N T(\epsilon)=\infty$ for every N. This result had been established previously only for N < 4.

