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Investigation of DF Kinetics in the Presence of BF<sub>3</sub>

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#### I. INTRODUCTION

The technology for helium-free DF lasers is being examined to determine if the operational advantages afforded by solid-grain reactants and closedsystem operation can be exploited. All of the required reactants can be produced from storable solid-grain compounds. One solid-grain formulation designed to produce fluorine utilizes  $NF_4BF_4$ , which results in the appearance of  $BF_3$  by product in the combustion products. The development of solid-grain compounds and auxiliaries to remove all  $BF_3$  would add cost, complexity, and bulk to the solid-grain system. The object of the present study is to assess the effect of  $BF_3$  on laser performance, with emphasis on the deactivation of DF by  $BF_3$ . The study includes a review of the deactivation of DF(v) by  $BF_3$ , calculation of combustor products, and DESALE calculations to determine laser power outputs with and without  $BF_3$ .

#### II. DEACTIVATION OF DF(v) by $BF_3$

The rate coefficient for the deactivation process

$$DF(v = 1) + BF_3 \xrightarrow{K(1)} DF(v = 0) + BF_3(1)$$

has been measured at T = 295 K by the technique of laser-induced fluorescence.<sup>1</sup> Although the temperature dependence is not known, the deactivation rate coefficients for the deactivation of DF(v = 1) by other molecules depend rather weakly on temperature in the 300- to 500-K range.<sup>2</sup> We have chosen a temperature-independent rate for the modeling calculations given by

 $K(1) = 1.25 \times 10^{11} \text{ cm}^3/\text{mol sec}$ 

The rate coefficients for the high vibrational levels have been estimated on the basis of the  $v^2$  scaling that holds for DF(v) deactivation by diatomic molecules.<sup>3</sup> The remainder of the kinetics package used in the laser modeling calculations are reactions and rate coefficients recommended by Cohen of The Aerospace Corporation.<sup>4</sup> These reactions are periodically reviewed and updated for consistency with the latest literature.<sup>5</sup> Cohen has estimated rate coefficients for the upper vibrational levels where no data are available.

#### III. CALCULATIONS

#### A. PRIMARY- AND SECONDARY-FLOW CALCULATIONS

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Typical gas-flow conditions are shown in Table 1 (taken from statement of work). Adiabatic flame-temperature calculations were performed for the primary flow at a total pressure of 2 atm. The results of this calculation are shown in Table 2. The combustion produces a temperature of 2291 K, with the only species of any concentration being fluorine, HF,  $BF_3$ , and  $N_2$ . The large bond strengths of HF,  $BF_3$ , and  $N_2$  make this result quite understand-able. Although an adiabatic flame temperature of 2291 K is predicted, the actual temperature of the gas at the entrance to the expansion nozzles is probably closer to 1800 K after heat loss to the walls of the combustor.

The combustor gas was assumed to expand adiabatically from 2 atm at 1800 K to a velocity of  $1.75 \times 10^5$  cm/sec (Mach No. = 5.0) with T = 282 K, and P = 2.86 Torr. These conditions were used as input conditions for the laser modeling calculations. The mole fractions leaving the combustor are shown in Table 2 to be fluorine, 0.161; HF, 0.356; N<sub>2</sub>, 0.379; and BF<sub>3</sub>, 0.104.

The secondary flow of  $D_2$  is assumed to undergo a small expansion to a velocity of 1.1 ×  $10^5$  cm/sec corresponding to a Mach number of 1.25. If the  $D_2$  is initially at room temperature, this expansion leads to a final static temperature of 225 K.

#### B. DESALE CALCULATIONS OF LASER PERFORMANCE

DESALE<sup>6</sup> calculations were performed for a laser with the following configuration: nozzle heights = 0.5 cm; primary nozzle width = 0.08936 cm; secondary nozzle width = 0.01064 cm; total length of nozzle bank = 100 cm; and an unstable cavity with a magnification M = 2. The DESALE calculations take into account finite mixing rates of the primary and secondary flows in the laser cavity, reaction rates, and laser radiation. The mirrors have one edge at x = 0, the nozzle exit, and the other edge at position x downstream. The optimum

Primary Flows	
Gas	Molar Amount
F <sub>2</sub>	1.0
NF <sub>3</sub>	1.0
BF3	0-1.0
N <sub>2</sub>	3.16
H <sub>2</sub>	1.72
Secondary Flow	
D <sub>2</sub>	5.5
Static Pressure = 2-10 Torr Initial Static Temperature = 200-400 Initial Total Temperature = 1800 K	ĸ

# Table 1. Typical Gas Flow Conditions

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Calculations

Flow

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sized mirrors are those that produce the maximum power (Fig. 1). The calculations were performed for the conditions listed in Table 3, and the results of output power versus mirror width are shown in Figs. 1 and 2.

A maximum power of 4700 W was calculated for Run 1, which contained the full amount of  $BF_3$  and the deactivation kinetics for  $DF(v)-BF_3$ . The calculation for Run 2 was performed for exactly the same conditions except the deactivation kinetics for  $DF(v)-BF_3$  were deleted. The output power in Run 2 increased to 4875 W from 4700 W in Run 1, a change of only about 4%. The  $BF_3$  was removed from the primary flow for Run 3 with a consequent 10.4% decrease in the static pressure. In order to match the pressures in the primary and secondary flows, the  $D_2$  flow rate was reduced by 10.4%. The resulting power of 4750 W falls between the powers of Runs 1 and 2. The lower pressure allows mixing to occur sooner, but the lower  $D_2$  concentration tends to stretch out the reaction zone.

In order to check the sensitivity of the results to uncertainties in the deactivation rate coefficient, a calculation was performed in which the rate coefficients for DF(v) deactivation by BF<sub>3</sub> were arbitrarily increased by a factor of 2. The calculated peak power in this case was 4515 W, about 4% below the power for Run 1 with the most probable rates and about 8% below that of Run 2 where the deactivation rates are set equal to zero. The decrease in laser power is roughly proportional to the increase in the deactivation rate. The temperature in the laser cavity increased from about 280 to 380 K at the peak power in all cases so that the average temperature was 330 K, only about 12% above the temperature of the DF(1)-BF<sub>3</sub> rate measurements. A calculation performed with the standard rate coefficients, but with a v<sup>3</sup> scaling for the DF(v)-BF<sub>3</sub> rates gave 4370 W, compared to 4700 W for the v<sup>2</sup> scaled calculation.

The most important deactivators in the system are HF and  $D_2$ . Instead of taking BF<sub>3</sub> out of the primary flow, an alternative is to reduce heat losses in the combustor and use 10% less H<sub>2</sub>, which makes HF. Run 4 shows that the power is increased about 5% by a 10% reduction in HF.

Table 3. Calculated Laser Performance

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		Run No.		
	1	2	я	4
Pressure, atm	0.003776	0.003776	0.0033835	0.003248
Primary Flow				
F, %	16.1	16.1	18.0	18.75
НF, %	35.6	35.6	39.7	37.19
BF3, %	10.4	10.4	0	0
N2, %	37.9	37.9	42.3	44.06
Comments	With deactivation by BF <sub>3</sub>	Without deactivation by BF3	No BF3, HF	No BF3, reduced by 10%
Max Power, W	4700	4875	4750	5000

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#### IV. CONCLUSIONS

Laser modeling calculations have been performed for a helium-free DF laser based on solid-grain reactants. The presence of  $BF_3$ , a combustion byproduct, in the laser cavity, degrades the laser performance by about 4% for reasonable estimates of the temperature dependence of the deactivation rate coefficient and for its scaling to higher vibrational levels. It is concluded, on the basis of our DESALE calculations, that laser performance degradation is not severe enough to warrant the removal of  $BF_3$  from a combustor flow that utilizes an  $NF_4BF_4$  solid-grain formulation.

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#### LABORATORY OPERATIONS

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