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THEORETICAL ANALYSIS OF MULTIMODE FIBER STRUCTURES.(U)  
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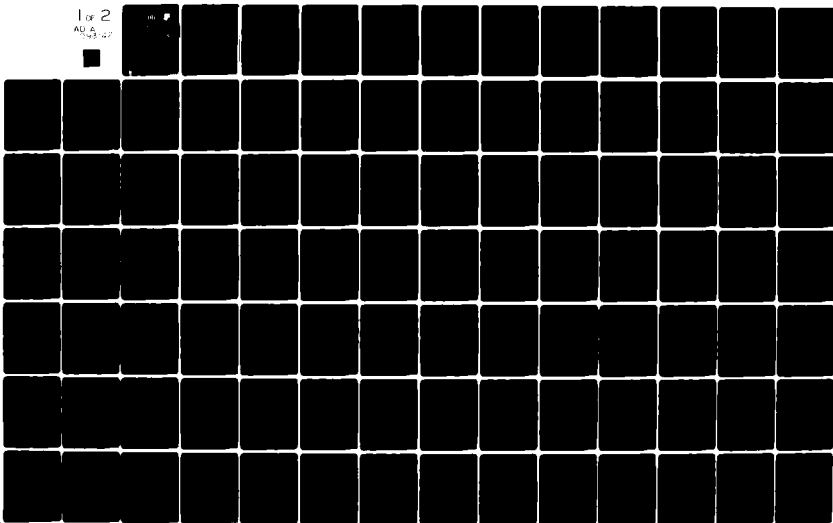
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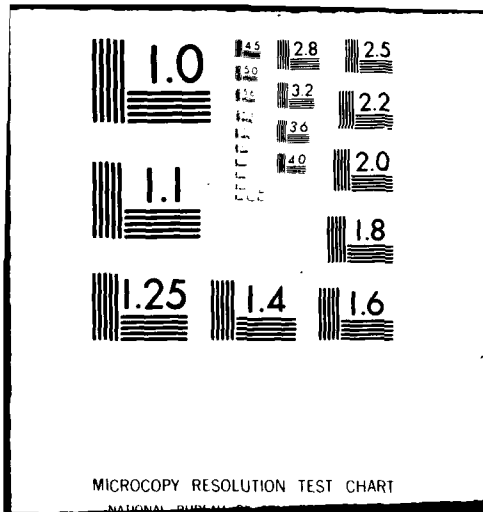
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Final Technical Report  
October 1980

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# THEORETICAL ANALYSIS OF MULTIMODE FIBER STRUCTURES

EMTEC Engineering Incorporated

C. Yeh

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Some referenced figures do not appear in this document. The references are to computer calculations too voluminous to publish. The resulting analysis of the computer data is adequately presented in this report. Therefore, the missing data is considered irrelevant to the conclusions presented herein. The missing data may be obtained by contacting RADC (ESO) Hanscom AFB MA 01731.

This report has been reviewed and is approved for publication.

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20. ABSTRACT (Continue on reverse side if necessary and identify by block number) This is a final report on the study of propagation characteristics of a light beam in multimode fiber structures. Realistic fiber structures made with commercially available fibers such as those provided by Corning or ITT were studied. The resultant computer programs may be used readily to generate design data for structures made with realistic fibers with step or parabolic index profiles. It is believed that our unique approach based on the scalar-wave FFT method may be extended to			

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treat problems dealing with nonlinear fibers or fibers with frozen-in statistically varying index profiles.

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The work reported herein was supported by the Electronic Systems Division of the Air Force Systems Command, USAD, Hanscom AFB. The author wishes to express his special thanks to Dr. L. Eyles for his sincere interest in this project and for his suggestion in the application of our technique to study the large-size single mode fiber. Technical discussion with Dr. P. Gianino was also appreciated. Continuing support of this project by Dr. A Yang is gratefully acknowledged.

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## EVALUATION

The effort as summarized in the report provides accurate theoretical analyses of the optical power transmission properties of a number of devices important to RADC efforts under TPO #4/D - Solid State Devices, Subthrust #3 - Electro-Optical Components. The devices include couplers, tapers, horns, and branches. The programs and techniques provided by the contractor permit in effect computer experiments to be done for a very large variety of design parameters. To do the actual experiments in the laboratory with this range of parameters would be enormously more expensive and time consuming. The report is then a crucial element in simplifying and accelerating the design process and in leading to final design specs in the shortest time.

  
LEONARD J. EYGES  
Project Engineer

## I. INTRODUCTION

This final report summarizes the work performed under Contract F19628-80-C-0053 which the Electronic Systems Division of the Air Force Systems Command granted to the EMtec Engineering, Los Angeles, California. The work was begun in January, 1980 and completed in August, 1980.

The principal thrusts of this R & D study in performing numerical analysis of multimode fiber components were two fold. Firstly, we wish to learn the limitation (and possible improvement) of our numerical scheme<sup>1</sup> and secondly, we wish to obtain numerical data for realistic multimode fiber structures.

Specifically, the following tasks were carried out:

- a) Study the effect of step index gradient and of tight beam confinement by an adaptive coordinate scheme.
- b) Study the effect of the presence of absorber at the edge of the mesh on the beam propagation characteristics of multimode fiber structures.
- c) Compute the coupling characteristics of tapered multimode fiber couplers and unequal size fiber couplers .
- d) Obtain data for reflection coefficients and beam waist changes for multimode fiber tapers, horns and branches.

In section II we shall present the implementation of the adaptive coordinates in our numerical solution of the scalar wave equation. Then, the scheme to include an absorber at the edge of the mesh will be described. Finally, an approximate approach to obtain the reflection coefficients for complex fiber structures will be shown. Detailed results of our study on the proposed tasks are given in Section III. Concluding remarks and recommendations for future work are included in Section IV.

## II. ANALYTICAL APPROACH

The basic approach taken to find the solution of wave propagation along complex fiber structures is to solve the reduced scalar wave equation via the fast Fourier transform (FFT) technique.<sup>2</sup> In this section we shall first indicate the conditions under which the exact vector wave equation may be simplified to yield the reduced scalar wave equation. Then we shall introduce the concept of adaptive coordinates<sup>3</sup> and incorporate this concept in the solution of the reduced scalar wave equation via the fast Fourier transform technique.

A. Formulation of the Scalar Wave Approach. Starting with the vector wave equation for the electric field vector  $\underline{E}$  in the fiber structure,

$$\nabla \times \nabla \times \underline{E} - \omega^2 \mu_0 \epsilon \underline{E} = 0 \quad (1)$$

where  $\omega$  is the frequency of the wave,  $\mu_0$  the permeability and  $\epsilon = \epsilon(r)$ , the inhomogeneous permittivity of the structure, and making use of the vector identity

$$\nabla \times \nabla \times \underline{E} = \nabla(\nabla \cdot \underline{E}) - \nabla^2 \underline{E} \quad (2)$$

and the relation

$$\nabla \cdot \underline{E} = -\frac{1}{\epsilon} \nabla \epsilon \cdot \underline{E}, \quad (3)$$

one has

$$\nabla^2 \underline{E} + \omega^2 \mu_0 \epsilon \underline{E} - \nabla \left( \frac{1}{\epsilon} \nabla \epsilon \cdot \underline{E} \right) = 0 \quad (4)$$

Rewriting Eq. (4) gives

$$\nabla^2 \underline{E} + \omega^2 \mu_0 \epsilon_0 \left\{ \frac{\epsilon}{\epsilon_0} \underline{E} - \left[ \frac{1}{\omega^2 \mu_0 \epsilon_0} \nabla \left( \frac{1}{\epsilon} \nabla \epsilon \cdot \underline{E} \right) \right] \right\} = 0$$

The relative importance of the terms within the curly brackets can be determined from the following

$$\frac{\epsilon}{\epsilon_0} \underline{E} = \mathcal{O}\left(\frac{\epsilon}{\epsilon_0} \underline{E}\right) \quad (6)$$

$$\frac{1}{\omega^2 \mu \epsilon_0} \nabla\left(\frac{1}{\epsilon} \nabla \epsilon \cdot \underline{E}\right) = \frac{1}{k_0^2} \mathcal{O}\left(\frac{\nabla \epsilon}{\epsilon} \cdot \nabla \underline{E}\right) = \mathcal{O}\left(\frac{\epsilon/\epsilon_0}{k_0 l} \underline{E}\right) \quad (7)$$

where the symbol  $\mathcal{O}$  means the "order of magnitude," and  $l$  is the smaller of the distance over which  $\epsilon/\epsilon_0$  and  $\underline{E}$  change appreciably. For single-mode fiber structures, the values of  $\epsilon/\epsilon_0$  and  $k_0 \lambda$  are typically in the range

$$\epsilon/\epsilon_0 = \mathcal{O}(2) \quad (8)$$

$$k_0 l = \frac{2\pi}{\lambda} l = \mathcal{O}(10^2 \text{ or } 10^3), \quad l = \mathcal{O}(10\mu \text{ to } 100) \quad (9)$$

$$\lambda = \mathcal{O}(1\mu)$$

It follows that the second term within the curly brackets in Eq. (5) is several orders of magnitude smaller than the first term  $\epsilon/\epsilon_0 \underline{E}$ . It is therefore justifiable to neglect the second term and write Eq. (5) in the form

$$\nabla^2 \underline{E} + k_0^2 \frac{\epsilon}{\epsilon_0} \underline{E} = 0$$

The physical significance of replacing Eq. (5) by Eq. (10) is this. By discarding the term  $\nabla \frac{1}{\epsilon} \nabla \epsilon \cdot \underline{E}$ , we are neglecting any depolarization effects that may occur. This means that the wave retains the polarization it has at the source, which is evidenced by the fact that Eq. (10) can be reduced to a scalar equation by writing  $\underline{E}(x)$  in the form

$$\underline{E}(\underline{x}) = \underline{e}_p u(\underline{x}) \quad (11)$$

where  $\underline{e}_p$  is a unit vector in the direction of the initial polarization of the wave.<sup>4</sup> Substituting Eq. (11) in Eq. (10), we find that  $u(\underline{x})$  satisfies the scalar wave equation,

$$\nabla^2 u + k_0^2 \frac{\epsilon}{\epsilon_0} u = 0 \quad (12)$$

This equation with the boundary condition on the initial surface, and the radiation condition at infinity, completely specifies  $u(\underline{x})$ , from which we can then obtain the electromagnetic field vectors  $\underline{E}$  and  $\underline{H}$ .

If we write  $u$  as the product of a factor  $e^{ik_0 z}$  that accounts for the rapid change in the phase of  $u$  along the direction of propagation and a complex amplitude  $A(\underline{x}, z)$ , a further simplification of the problem results

$$\left[ 2ik_0 \frac{\partial}{\partial z} + \nabla_T^2 + k^2 (n^2(\underline{x}, z) - n_0^2) \right] A(\underline{x}, z) = -\frac{\partial^2 A(\underline{x}, z)}{\partial z^2} \quad (13)$$

where  $\nabla_T^2$  is the transverse Laplacian  $\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2}$ , and  $n_0$  is a given constant which represents the refractive index of some uniform medium. At laser wavelengths the complex amplitude  $A(\underline{x})$  varies much more rapidly transverse to the direction of propagation than it does along the direction of propagation. This enables us to make the paraxial approximation wherein the term on the right side of Eq. (13) is neglected (in the Russian literature this is called the parabolic approximation). So, the complex amplitude now satisfies

$$\left[ i2kn_0 \frac{\partial}{\partial z} + \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + k^2 (n^2(\underline{x}, z) - n_0^2) \right] A(\underline{x}, z) = 0 \quad (14)$$

For given initial data, i.e., values of the field at points on the initial surface, the propagation simulator must generate the corresponding field values at the terminal aperture such that Eq. (14) is satisfied. To do this we divide the medium into slabs defined by planes on which  $z$  is constant. In going from one slab to the next, we write  $A(\underline{x}, z)$  in the form

$$A(\underline{x}, z) = e^{\Gamma(\underline{x}, z)} w(\underline{x}, z) \quad (15)$$

where  $\Gamma(\underline{x}, z)$  is a phase function associated with the medium inhomogeneities

$$\Gamma(\underline{x}, z) = \frac{1k}{2} \int_{z_0}^z \left[ n^2(\underline{x}, y, z') - n_0^2 \right] dz' \quad (16)$$

The modified complex amplitude  $w(\underline{x}, z)$  then satisfies the equation

$$\left[ i2kn_0 \frac{\partial}{\partial z} + e^{-\Gamma} \nabla_T^2 e^{\Gamma} \right] w(\underline{x}, z) = 0 \quad (17)$$

with the initial condition

$$w(\underline{x}, y, 0) = u(\underline{x}, y, 0) \quad (18)$$

Physically, these equations approximate the propagation in the inhomogeneous medium by a two-step process at each  $z$  increment. First, we propagate the field  $u(\underline{x})$  at  $z - \Delta z/2$  to  $z + \Delta z/2$ , assuming that the intervening space is homogeneous. The effect of the inhomogeneities between  $z - \Delta z/2$  and  $z + \Delta z/2$  is then accounted for by multiplying this solution by the phase factor  $\exp(\Gamma)$ .

#### B. Adaptive Coordinates

To reduce the size of the mesh required to solve Eq. (17) numerically,

let us introduce an adaptive coordinate system defined by the transformation<sup>5</sup>

$$\zeta_1 = \frac{x/\rho_0}{N(z)} \quad (19)$$

$$\zeta_2 = \frac{y/\rho_0}{N(z)} \quad (20)$$

$$N(z) = \alpha^{-1/2} \left[ \left(1 - \frac{z}{f}\right)^2 + \alpha^2 \left(\frac{z}{k\rho_0}\right)^2 \right]^{1/2} \quad (21)$$

$$\xi = \tan^{-1} \left[ \frac{(1+\beta) \frac{z}{f}}{\beta^{1/2}} - 1 \right] \quad (22)$$

$$\beta^{1/2} = \alpha \frac{f}{k\rho_0} \quad (23)$$

where  $\rho_0$  is a characteristic dimension of the beam at the initial surface (e.g., the e-folding radius of a gaussian beam),  $f$  is the distance to the focus, and  $\alpha$  is a constant determined by the requirement that the solution be confined within the boundaries of the mesh at the focal plane. The choice  $\alpha = 1$  yields a coordinate system that converges at a rate determined by the free-space diffraction of a gaussian beam having an e-folding radius  $\rho_0$ .

When written in terms of the converging coordinate variables defined above, Eqs. (15) and (17) for the complex amplitude are replaced by the relations

$$w(x, y, z) = \hat{w}(\zeta, \xi) \exp(\tilde{\Gamma}) v(\zeta, \xi) \quad (24)$$

$$\hat{w}(\zeta, \xi) = \left( \alpha^{1/2} N(z) \right)^{-1} \exp \left[ \frac{i}{2} (\zeta_1^2 + \zeta_2^2) \tan \xi \right] \quad (25)$$

$$\tilde{\Gamma} = \frac{ik}{2} \int_{z-\Delta z/2}^{z+\Delta z/2} dz' \left( n^2(x, y, z') - 1 \right) - \frac{i}{2} (\zeta_1^2 + \zeta_2^2) \Delta \xi \quad (26)$$

$$\left[ \frac{\partial}{\partial \xi} - \frac{i}{2} \exp(-\tilde{\Gamma}) \left( \frac{\partial^2}{\partial \zeta_1^2} + \frac{\partial^2}{\partial \zeta_2^2} \right) \exp(\tilde{\Gamma}) \right] v = 0 \quad (27)$$

where  $\Delta\xi$  is the increment in  $\xi$  in going from  $z-\Delta z/2$  to  $z+\Delta z/2$ . The initial condition for  $v$  is

$$v(\underline{z}, \xi) = w(x, y, z) / \tilde{w}(\underline{z}, \xi) \quad (28)$$

To solve Eq. (27) we utilize the fact that for sufficiently small values of  $\Delta\xi$  (i.e.,  $\Delta z$ ) the effect of the exponential factors  $\exp(\pm \tilde{\Gamma})$  in this equation is small. Hence, we solve the simpler equation obtained when these factors are equated to unity

$$\left[ \frac{\partial}{\partial \xi} - \frac{i}{2} \left( \frac{\partial^2}{\partial \zeta_1^2} + \frac{\partial^2}{\partial \zeta_2^2} \right) \right] v = 0 \quad (29)$$

We use a fast Fourier transform technique to solve Eq.(29). The basis of this approach is the fact that the solution of Eq. (29) can be expressed in the form of a discrete Fourier series

$$v(\underline{z}, \xi) = \sum_{m=0}^{N-1} \sum_{n=0}^{N-1} v_{mn}(\xi, t) \exp[i(p_m \zeta_1 + q_n \zeta_2)] \quad (30)$$

where the Fourier coefficients  $v_{mn}$  are determined from the initial data and Eq. (29) as follows. The initial values of  $v_{mn}$  are obtained by taking the discrete Fourier transform of the initial values of  $v(\underline{z}, \xi_1)$  over a mesh of points  $\zeta_1 = [l - (N/2)] \Delta\zeta$ ,  $\zeta_2 = [j - (N/2)] \Delta\zeta$  ( $l, j = 0, 1, \dots, N-1$ )

$$v_{mn}(\xi_i) = \frac{(-1)^{m+n}}{N^2} \sum_{l=0}^{N-1} \sum_{j=0}^{N-1} v\left(\left(l - \frac{N}{2}\right) \Delta\zeta, \left(j - \frac{N}{2}\right) \Delta\zeta, \xi_i\right) \exp\left[-\frac{i2\pi}{N} (ml + nj)\right] \quad (31)$$



The dependence of  $V_{mn}$  is then determined by substituting Eq. (30) in Eq. (31), which yields

$$\frac{\partial V_{mn}}{\partial \xi} + \frac{i}{2} (p_m^2 + q_n^2) V_{mn} = 0 \quad (32)$$

from which it follows that

$$V_{mn}(\xi) = V_{mn}(\xi_j) \exp \left[ - \frac{i(p_m^2 + q_n^2) \Delta \xi}{2} \right] \quad (33)$$

Finally, it can be shown that in order for the discrete Fourier series representation of  $v$  given in Eq. (30) to be real when  $v$  is real, the coefficients  $p_m$  and  $q_n$  must have the form

$$p_m = \frac{2\pi}{N\Delta\zeta} \left( m - \frac{N}{2} \right) \quad (34)$$

$$q_n = \frac{2\pi}{N\Delta\zeta} \left( n - \frac{N}{2} \right) \quad (35)$$

Hence, for discrete points  $\zeta_1 = (\ell - N/2)\Delta\zeta$ ,  $\zeta_2 = (j - N/2)\Delta\zeta$  ( $\ell, j = 0, 1, \dots, N-1$ )

$$\begin{aligned} & v \left( \left( \ell - \frac{N}{2} \right) \Delta \zeta, \left( j - \frac{N}{2} \right) \Delta \zeta, \xi \right) \\ &= (-1)^{\ell+j} \sum_{m=0}^{N-1} \sum_{n=0}^{N-1} (-1)^{m+n} V_{mn}(\xi_j) \\ & \exp \left[ - i \beta \left( \left( \frac{m - \frac{N}{2}}{N} \right)^2 + \left( \frac{n - \frac{N}{2}}{N} \right)^2 \right) + i \frac{2\pi}{N} (\ell m + j n) \right] \quad (36) \end{aligned}$$

where  $\hat{\beta} = 2\pi^2 \Delta\xi / (\Delta\xi)^2$ . Note that  $v$  is simply  $(-1)^{l+j}$  times the discrete Fourier transform of  $(-1)^{m+n} v_{mn}(\xi)$ .

The effect of the medium and the factor  $\exp\{-i/2(\xi_1^2 + \xi_2^2)\Delta\xi\}$  introduced by the coordinate transformation is taken into account at each  $\xi$  step in the calculation by multiplying the value of  $v$  obtained in the previous step by the quantity  $\exp(\tilde{\Gamma})$  defined in Eq. (26), i.e., the initial value inserted in Eq. (31) is  $\exp(\tilde{\Gamma})$  times the value of  $v$  determined from the previous steps.

Using this adaptive coordinate algorithm we have been successful in our treatment of various realistic multimode fiber structures. Results are summarized in Section III.

### C. Implementation of Lossy Outer Boundary.

It is believed that the field touching the outer boundary of the cladding region of a fiber structure will be attenuated due to radiation or absorption. To accommodate this situation in order to further improve our computer simulation, we have incorporated the presence of a lossy dielectric layer outside the cladding region in our computer program. An example of the index profile of a fiber is shown in Fig. 1:

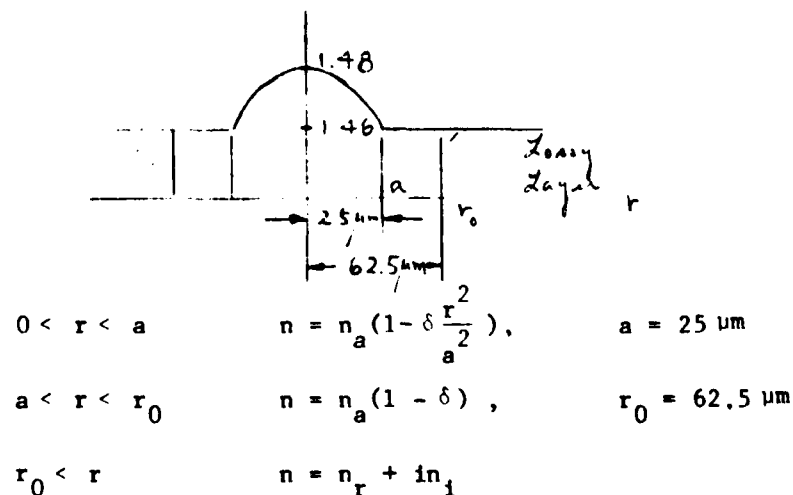


Figure 1: A Typical Index Profile with Lossy Outer Layer

Typical intensity patterns of beams propagating in a fiber with the index profile given by Fig. 1 are shown in Fig. 2. It was found that the propagation characteristics of the guided beams are not significantly affected by the presence of a lossy outer layer, except when the spot size of the beam is larger than the core diameter, as expected.

D. A Heuristic Approach in Obtaining the Reflection Coefficient.

One of the a'prior assumption in the development of the FFT scalar wave approach is that only paraxial rays are allowed and no reflection is permitted. This assumption enables us to develop an algorithm, there-by, we may obtain the propagating field by a forward stepping process as described earlier. As the field evolves from one  $z$  plane to the next  $z + \Delta z$  plane, the averaged value of the refractive index as seen by the field may be different as illustrated in Fig. 3:

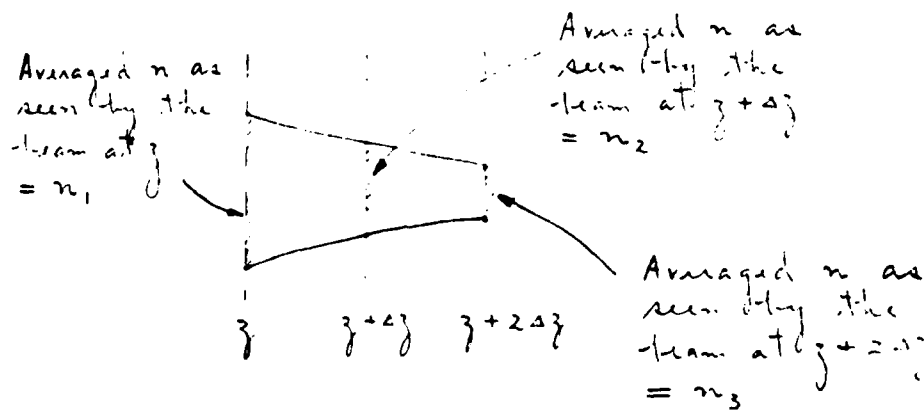
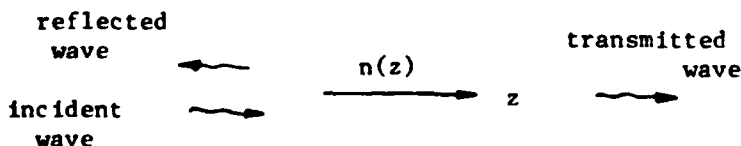


Figure 3: Illustration of the averaged  $n$  as seen by the beam.

In affect one may postulate that the wave is experiencing reflection in a medium with longitudinally slowly varying refractive index as shown below:



$n(z)$  is given by Fig. 3.

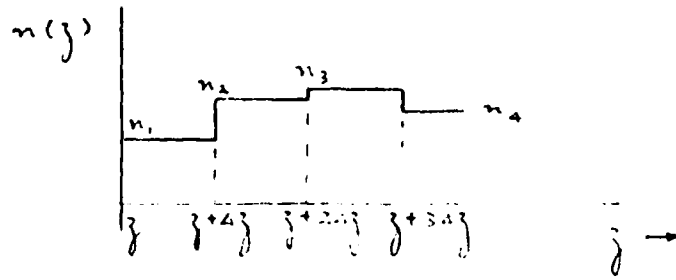


Figure 4: Equivalent Index Profile

The reflection coefficient for a plane wave propagating in this longitudinally non-uniform medium may be obtained according to a formula derived for the case of plane wave propagation in stratified layered medium:<sup>6</sup>

$$R(z) = - \exp[-is(z)] \int_z^{\infty} \gamma(z) \exp[is(z)] dz$$

$$s(z) = 2 \int^z \beta(z) dz \quad \beta(z) = k_0 n(z)$$

$$\gamma(z) = \frac{dn}{dz} / 2\beta$$

This is the heuristic approach that we shall use to calculate the reflection coefficient for waves in our multimode fiber structures.

### III. Results

The algorithms detailed above have been implemented in our computer programs. Results for the proposed tasks are given in the following:

#### (a) Effects of Step Index Gradient on the Propagation Characteristics.

The purpose of this study is to learn the effects of step index gradient on the propagation characteristics of waves in a multimode fiber guide. Let us introduce the following index profile:

$$n(r) = n_0 - \delta \left(\frac{r}{a}\right)^{2m} \quad (0 < r < a)$$

$$n(r) = n_0 - \delta \quad (a < r)$$

where  $n_0$ ,  $m$  and  $\delta$  are given constants and  $a$  is the core radius of the fiber. For a typical parabolic index profile fiber, one has

$$n_0 = 1.48, \quad \delta = 0.02, \quad a = 25\mu\text{m}, \quad m = 1.$$

The constant  $\delta$  must necessarily be small so that the depolarization effect may be ignored and the scalar wave approach may be justified. By varying  $m$ , the steepness of the index gradient may be varied as shown in Fig. 5. It should be kept in mind that even when the FFT technique is capable of handling steep index variations, the slope of the index profile must still be gentle enough so that the gradient term in the exact wave equation (Eq. (5)) may be ignored. We have carried out propagation calculation for the following specific cases:  $n_0 = 1.48$ ,  $\delta = 0.02$ ,  $a = 25\mu\text{m}$ ,  $m = 1, 4, 6, 10$ . Higher  $m$  values means steeper index gradient. As shown in Figs. 6, no computational difficulties were encountered for even the steepest case ( $m = 10$ ) in which the index changes from  $n = 1.48$  to  $n = 1.46$  in  $3\mu\text{m}$  distance for  $50\mu\text{m}$  core diameter fiber. However, one should be aware that we are pushing the limit of validity for the scalar wave approach.

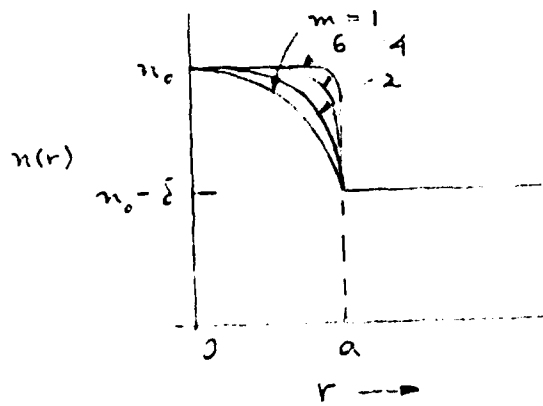


Fig. 5. Plot of  $n(r) = n_0 - \delta (r/a)^{2m}$ .

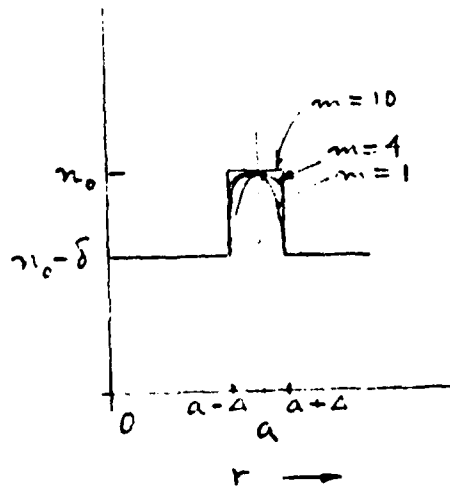


Fig 7 Plot of  $n = n_0 - \delta ((r-a)/\Delta)^{2m}$

From the results, it is of interest to note that as  $m$  increases from 1, i.e., as the index profile deviates from the parabolic profile, the beam profiles no longer remain to be of gaussian shapes, but take on ring-type structures. This implies that the phase front of the multimode beam is no longer a monotonic function of the radial distance but has become an oscillatory one.

We may conclude from these calculations that our program is capable of handling problems with steep index gradient. The index transition may occur in a distance as small as  $4\lambda$  where  $\lambda$  is the free-space wavelength. The limiting factor apparently is the justification for the elimination of the depolarization term in Eq. (5).

(b) Beam Propagation in a Ring Fiber

A typical single-mode fiber has a core diameter of the order of  $10\mu\text{m}$ . Consequently it is more difficult to handle than a multi-mode fiber. An idea to enlarge the single-mode fiber has recently been put forth by Dr. L. Eyges of RADC. He suggested that perhaps a ring-type structure may support a single mode and yet possesses larger dimension than the usual solid-core fiber. This task was undertaken to investigate this possibility. Let us postulate that the index profile of a ring fiber takes the following

$$\text{form} \quad n = n_0 - \delta \left(\frac{r-a}{\Delta}\right)^{2m} \quad \text{for } a-\Delta < r < a+\Delta$$

$$n = n_0 - \delta \quad \text{for } a-\Delta > r, r > a+\Delta$$

where  $n_0$ ,  $\delta$ ,  $a$ ,  $\Delta$ , and  $m$  are given constants. By increasing the  $m$  value, one may adjust the steepness of the index gradient as shown in Fig. 7. Two types of initial beam shapes will be studied: (1) a solid centered gaussian beam and (2) a hollow-centered donut beam. We wish to learn how well the ring fiber will confine these two types of beams. The field expression for a solid gaussian beam takes the form

$$u = e^{-\frac{\alpha}{2} \left(\frac{r}{a}\right)^2}$$

while the hollow-centered donut beam takes the form

$$u = e^{-\frac{\alpha}{2} \left[ \left( \frac{r}{a} \right)^2 - 1 \right]}$$

where  $\alpha$  and  $a$  are given constants. Results of our computation are shown in Figs. 8 and 9. By following the evolution of the beam intensities, one may determine how well the ring fiber is guiding the beam. It can be seen from these figures that the solid beams appear to be better confined than the hollow beams, although the spreading of the solid beam energy is quite noticeable. It also appears that simple insertion of beam energy in the high index region of the ring fiber does not insure good guidance of the beam energy. One may conclude from this preliminary study that neither solid Gaussian beams nor hollow-centered Gaussian beams correspond to the mode energy distribution of a single-mode in a ring fiber. One should first perform the classical modal analysis to obtain the mode pattern of the single mode and then use this mode pattern as the initial beam pattern for propagation down the ring fiber. It is believed that the use of ring-index fiber as large core single-mode fiber definitely possesses merit and should be studied further. What we have demonstrated with our present study is that our program is capable of handling this type of fibers.

#### (c) Fiber Couplers

One of the simplest type of light couplers is the fiber coupler. By placing two or more fibers in close proximity of each other light energy may transform from one to the other through the coupling effect. This coupling process is rather involved. The well-known coupled mode theory may be adequate for simple, single-mode structures such as slabs



with reasonable separations. But, when multi-mode complex structures such as the fiber couplers are involved, the coupled mode theory becomes grossly inadequate.\* On the other hand, our FFT-scalar wave approach is uniquely qualified to deal with this fiber coupler problem. This is because this technique provides the evolution of beam field as it propagates down a complex multimode inhomogeneous fiber structure. Four types of fiber couplers have been studied:

Case 1 Coupling between two equal parabolic index fibers.

Two graded-index fibers are fused together longitudinally with separation  $d$  between their centers. The index profile for each fiber is given by

$$n(r_{1,2}) = n_0 - \delta \left( \frac{r_{1,2}}{a} \right)^2$$

where  $n_0$ ,  $\delta$ , and  $a$  are given constants, and 1 or 2 refers, respectively, to #1 or #2 fiber. Typical values for a Corning or ITT graded index fibers are used:

$$n_0 = 1.48$$

$$\delta = 0.02$$

$$a = 25\mu\text{m}$$

---

\*Recent advances by L. Eyges and P. Gianino of RADC using the extended boundary condition technique have shown that single mode couplers involving arbitrarily shaped uniform core guides can be successfully and accurately treated.

Various separation  $d$  were used. A gaussian beam represented by

$$u(x,y) = u_0 \exp \left\{ \left[ - \left( x + \frac{d}{2} \right)^2 - y^2 \right] / w^2 \right\}$$

where  $u(x,y)$  is the scalar wave function of the beam, and  $u_0$ ,  $w$  are given constants, is incident on one of the fibers. Results have been obtained for

$$w = 2.5\mu\text{m}, 5\mu\text{m}, 10\mu\text{m}$$

$$d = 8\mu\text{m}, 12\mu\text{m}, 16\mu\text{m}, 20\mu\text{m}.$$

The evolution of the beam along this structure is shown in Figs. 10-11. Displayed in Figs. 12 is the %-power in one fiber as a function of longitudinal distance for various separations and initial beam sizes. For the situations considered above, many modes are excited. The coupling process is very involved as displayed in Fig. 10-11. It still appears that back and forth power exchange among the guides prevails. The complexity of the power exchange phenomenon for the multimode coupler re-emphasizes the importance of obtaining design data through analysis before the actual construction of fiber coupler.

Case 2 Coupling between two equal step-index fibers.

This coupler is identical to the previous one except step-index fibers were used. We shall approximate the index profile of a step index fiber by the following expression:

$a(0, -4), (0, 4)$   
 Separations  $b(0, -6), (0, 6)$   
 $c(0, -8), (0, 8)$

Two Fiber Core  
 Parabolic Index Profile  
 Initial Beam Spot Size =  $10\mu\text{m}$

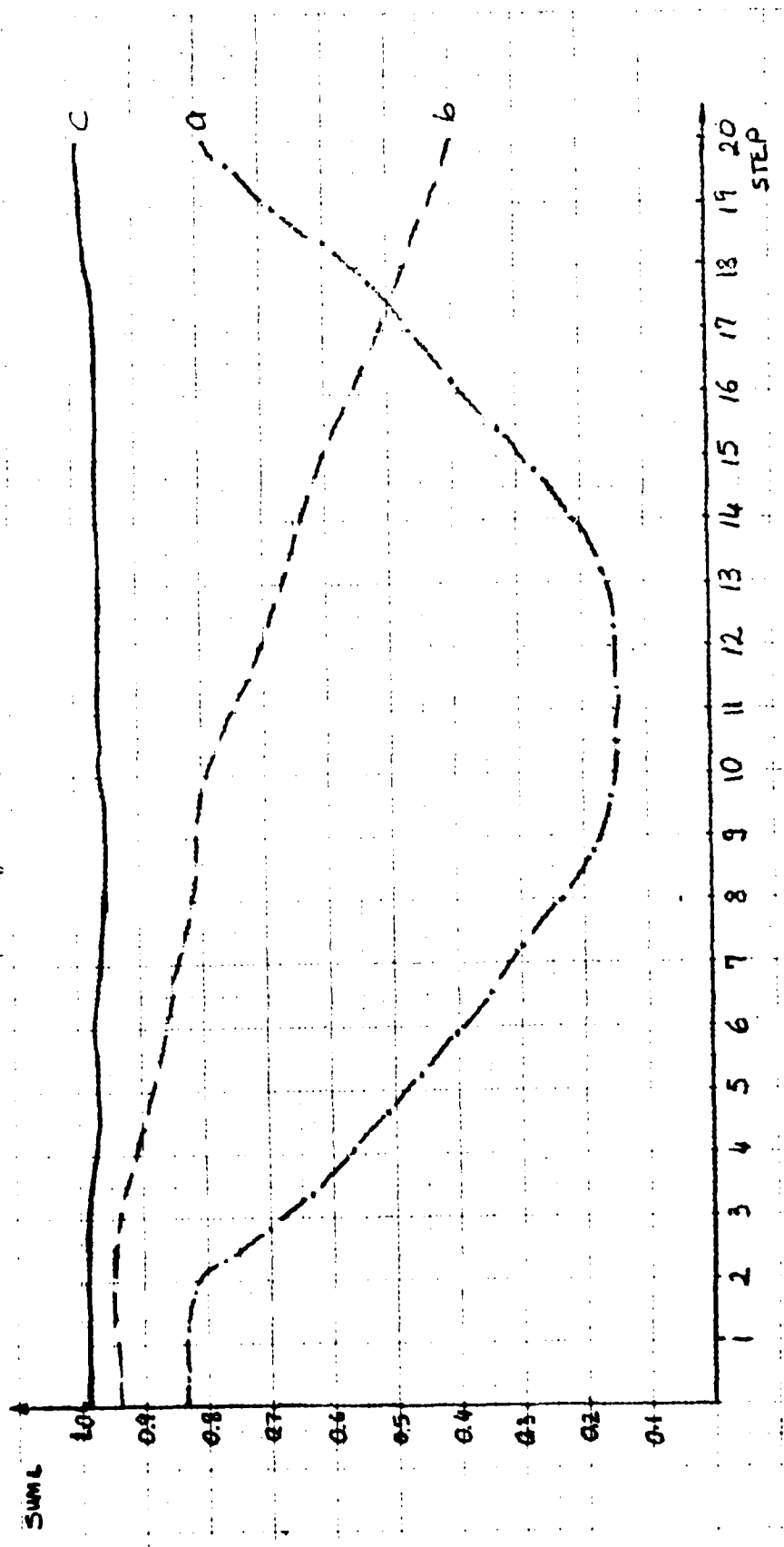


FIG 12 (a)

Two Fiber Composites  
 Parabolic Index Profiles  
 Initial Beam Spot Size = 20  $\mu$ m

Pressure - 1000x  
 Pressure - 200x  
 a(0, -4), (0, 4)  
 b(0, -6), (0, 6)  
 c(0, -8), (0, 8)  
 d(0, -11), (0, 11)

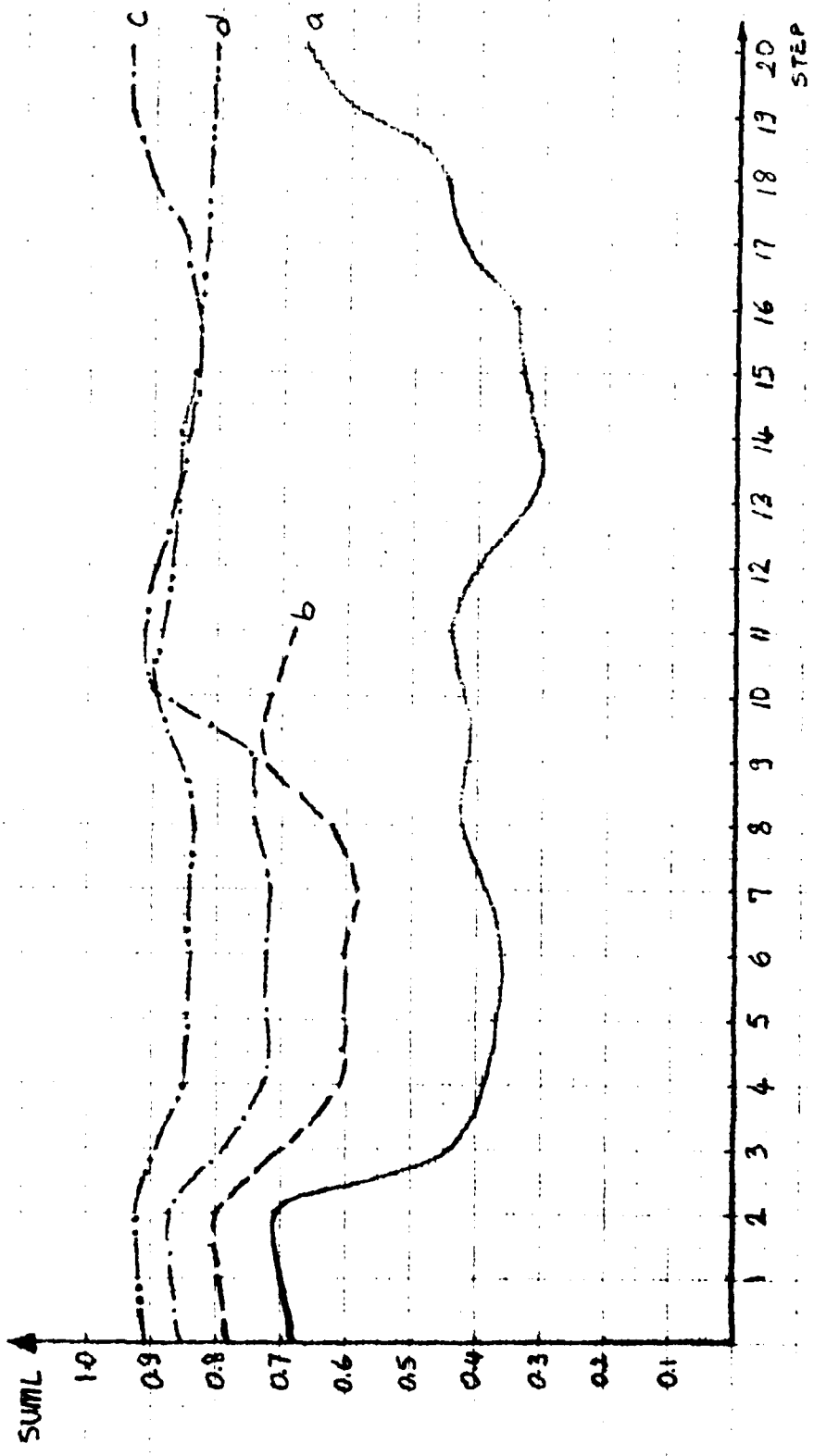


FIG 12 (b)

$$n(r_{1,2}) = n_0 - \delta \left( \frac{r_{1,2}}{a} \right)^{2m}$$

with  $m = 4$ . Again typical values for a Corning or ITT graded index fiber are used; i.e.,  $n_0 = 1.48$ ,  $\delta = 0.02$ ,  $a = 25\mu\text{m}$ . Results have been obtained for

$$w = 2.5\mu\text{m}, 5\mu\text{m}$$

$$d = 8\mu\text{m}, 12\mu\text{m}, 16\mu\text{m}$$

where  $d$  is the separation distance and  $w$  is the beam waist radius. Specific results are given in Figures 13. Displayed in Figs. 14 is the % power in one fiber as a function of longitudinal distance for various separations.

Case 3 Coupling between two unequal fibers.

It is of interest to learn, when two unequal size fibers are placed side by side, whether transfer of power would occur for the multimode case. The following fibers were used:

$$n(r_{1,2}) = n_0 - \delta \left( \frac{r_{1,2}}{a_{1,2}} \right)^{2m}$$

$$n_0 = 1.48 \quad m = 1, 4$$

$$\delta = 0.02$$

$$a_1 = 25\mu\text{m}$$

$$a_2 = 12.5\mu\text{m}$$

$$\text{Separation distance, } d = 8\mu\text{m}, 12\mu\text{m}$$

$$\text{Initial Beam Radius, } w = 5\mu\text{m}, 10\mu\text{m}$$

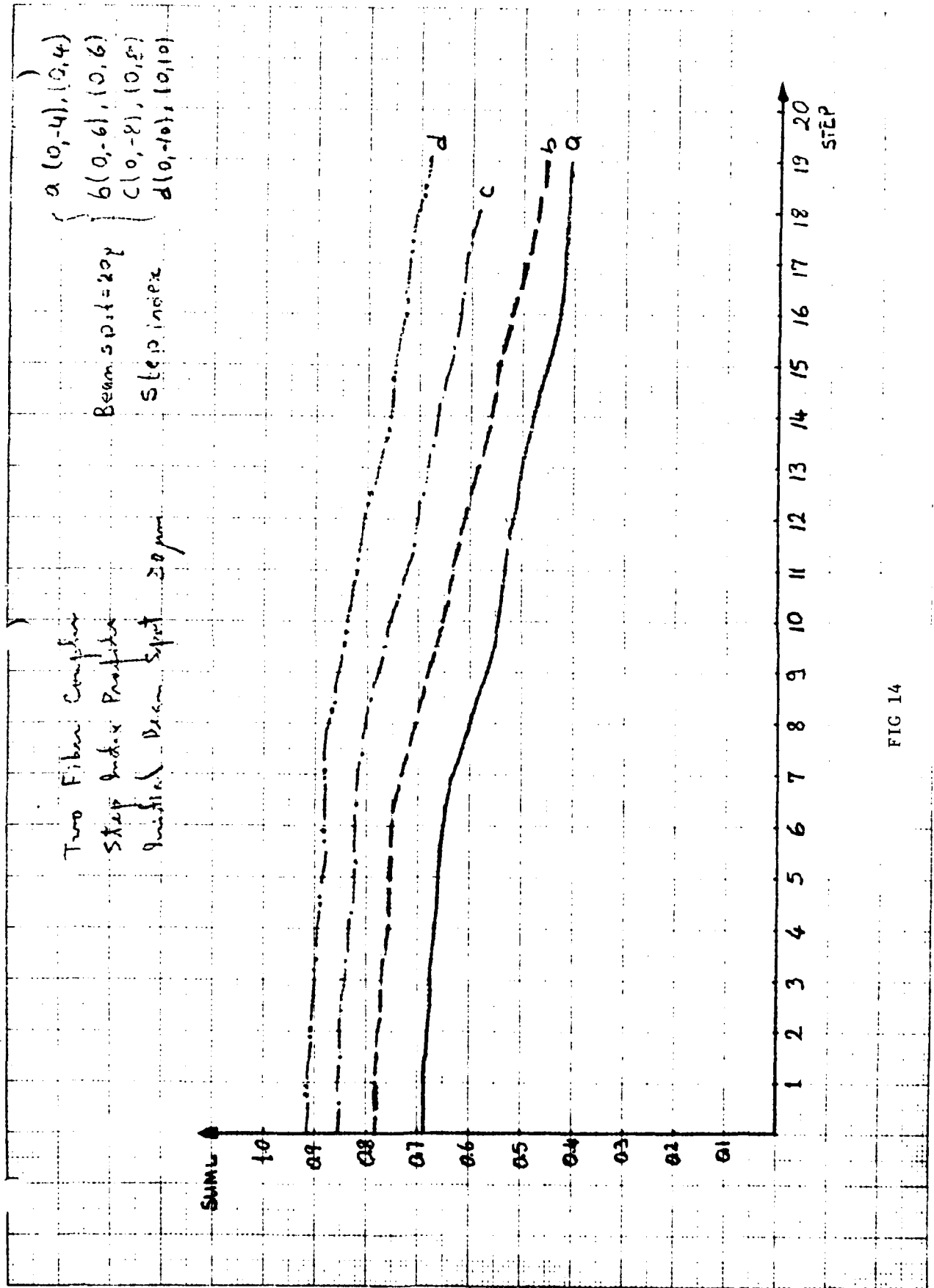


FIG 14

For the above chosen parameters, one may observe from Figs. 15-16 that only nominal coupling occurs between two unequal size fibers. In other words these structures are not efficient couplers.

Case 4 Coupling between more than two fibers.

As illustrations, we have considered two types of structures: Three equal size fibers located equal-distance from each other and three unequal size fibers arranged in a triangle shape. One fiber is initially illuminated, we wish to learn the power exchange characteristics in these two couplers. The following parameters were used:

$$n(r_{1,2,3}) = n_0 - \delta \left( \frac{r_{1,2,3}}{a_{1,2,3}} \right)^2$$

$$n_0 = 1.48$$

$$\delta = 0.02$$

$$a_1 = 25\mu\text{m}$$

$$a_2 = 25\mu\text{m}, 20\mu\text{m}$$

$$a_3 = 25\mu\text{m}, 15\mu\text{m}$$

$$\text{Separation distances, } d_1 = d_2 = d_3 = 12\mu\text{m}$$

$$\text{Initial beam radius, } w = 5\mu\text{m}$$

It can be seen from Fig.17 that when one of the fiber of 3 identical fibers, separated an equal distance from each other, is illuminated with a gaussian beam, power exchange takes place between two unilluminated fibers with the one illuminated fiber in a synchronous manner. In other words, each of the originally unilluminated fibers is receiving identical amount of power transfer from the illuminated one, as expected. On the other hand, the power exchange phenomenon for 3 non-identical fibers coupler is much more complicated as seen from Fig.18

Larger power transfer tends to take place among fibers with similar diameters.

Fiber coupler is one of the crucial components in any fiber optics system. As such, one must understand the detailed wave interaction phenomenon in this type of coupler so that correct design can be made. We have demonstrated the capability of our technique in dealing with this type of fiber structures. Systematic studies of different varieties of fiber couplers may now be undertaken.

(d) Reflection Coefficient Calculations.

We have implemented the heuristic approach discussed earlier in the computer program to yield reflection coefficients for waves propagating in the various multimode fiber structures. Although it is difficult to justify the accuracy of the absolute values for the reflection coefficients obtained according to this algorithm, nevertheless, we feel that their relative values for different structures can be believed. This is because our heuristic approach took into consideration the fundamental characteristic of wave reflection: i.e., reflection occurs when discontinuity of the propagation medium or structure is experienced by the wave. The larger is the discontinuity, the larger will be the reflected energy.

Results for sample calculations for the reflection coefficients for various fiber structures are shown in Figs.19-20



#### IV. Conclusions and Recommendations

Support of this program has enabled the contractor to develop and perfect a computer program based on the scalar wave - FFT algorithm to study the propagation characteristics of guided waves in several important, practical fiber structures such as fibers with general index profiles (step index, parabolic index, ring index, etc), multi-channel fiber couplers, and fiber horns or tapers. These fiber structures may be made with commercially available fibers whose index variation may be as large as 1 - 2%.

We have implemented the adaptive coordinates and lossy outer mesh boundary schemes in our computer program. However, for most practical situations of interest in which the fiber core radius is about  $25\mu\text{m}$ , the cladding index is about 1 - 2% less than the core index, the spot size of the beam is less than  $20\mu\text{m}$  and the free-space wavelength of the beam is larger than  $0.6\mu\text{m}$ , it is not necessary to implement the adaptive coordinates and lossy outer mesh boundary schemes. We also discovered that steep index gradient is not a hindrance for the program to produce accurate results as long as the scalar wave approach is still justifiable.

It is not unreasonable to ask the following question:

"Now that we have completed a beautiful program capable of producing propagation results for a variety of practical fiber structures, what can we do with it?"

The answer is "May be a lot!" Listed below are only a few of the problems that we can solve with this program:

- (1) Any single-mode or multimode weakly guiding fiber with arbitrary refractive index profile.

Our program provides the means to obtain the propagation characteristics of guided waves supported by this structure. The core of the structure may be circular, elliptical, rectangular, triangular or dumb-bell

shape with general index profile.<sup>8</sup>

- (2) Any fiber couplers composed of parallel strands of two or more of the above fibers. This is the only program which can provide the detailed coupling characteristics of this type of structure. Prior knowledge of coupling characteristics of a coupler is the key to successfully design and construct fiber couplers.
- (3) Any transition elements derived from the above fibers. Transition elements such as tapers, horns, or mode converters or branches can all be analyzed by our program.

#### Recommended Future Research

In addition to the important practical problems mentioned above that can be solved by our approach, it is worthwhile to look into the future and seek out problems of potential importance and interest. For example:

- (1) Nonlinear Fiber.

Very high intensity is achievable in fibers. One may wish to learn the propagation characteristics of waves in a fiber in which the induced nonlinearity of the medium plays a significant role. This problem may be solved by the scalar wave - FFT approach.

- (2) Mode Conversion in a Fiber Due to its statistically Varying Medium.

This problem may also be approached from the scalar wave - FFT point of view. It is known that the presence of frozen-in statistically varying medium contributes to the mode conversion phenomenon in a fiber.<sup>9</sup> A systematic study of this problem will reveal the severity of this effect in changing the dispersion characteristics of beam in this fiber.

- (3) Large-size Single-Mode Fiber

The advantages of having large-size single-mode fiber are well-known. Ring-index fiber as proposed by L Eyges of RADC appears

to be a promising one. Other type of fiber whose index variation may be radially unsymmetrical, such as, a layered fiber as shown in Fig. 21 may also be promising. Research should be encouraged on this type of fibers.

(4) Polarization Preserving Fibers

One of the main features of a weakly guiding fiber which can be analyzed by the scalar wave approach is that the guiding structure is polarization insensitive. However, for several important application areas, the polarization preserving characteristic of fiber is essential. Stress induced birefringent fiber or deformed core fiber may satisfy the needed requirement. However, to achieve the targeted isolation for the two orthogonal dominant modes, the required stress is excessively high for stress induced birefringent fiber and the loss is excessive for deformed-core fiber. We propose the use of layered dielectrics as shown in Fig. 21 to produce an equivalent birefringent effect to enable proper isolation for the two orthogonal dominant modes.<sup>10</sup> Initial indication is very promising. It would therefore be very worthwhile to pursue this research.

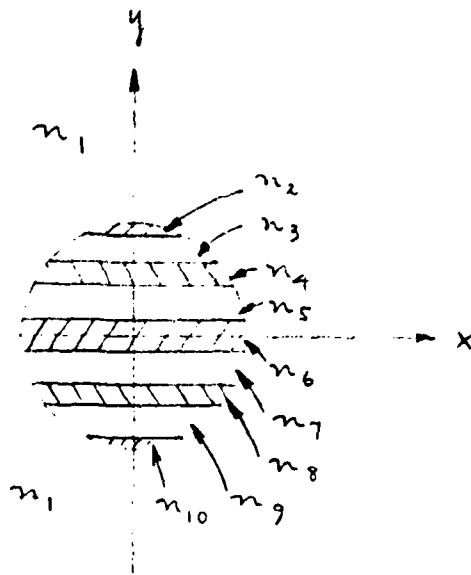


FIG. 21 Polarization Preserving Layered Fiber

**Personnel:**

The principal contributors of this contract have been:

C. Yeh                      Senior Research Engineer

P. Barber                  Research Engineer

F. Manshadi              Research Engineer

#### REFERENCES

1. C. Yeh, "Analysis of Multimode Fiber Couplers, Tapers, and Mode Converters", RADC-TR-79-341, Rome Air Development Center, Griffiss Air Force Base, New York 13441 (Jan. 1980). (A081669)
2. C. Yeh, L. Casperson, and B. Szejn, J. Opt. Soc. Am. 68, 989 (1978).
3. L.C. Bradley and J. Herrmann, Appl. Opt. 13, 331 (1974).
4. C. Yeh, L. Casperson, and W.P. Brown, Appl. Phys. Lett. 34, 460 (1979).
5. M. Born and E. Wolf, "Principles of Optics", 2nd Edition, p. 155. Pergamon Press, New York (1964).
6. L.M. Brekhovskikh, "Waves in Layered Media", Academic Press, New York (1960).
7. L. Eyges and P. Gianino, "Single Mode Coupler, Scalar Wave - Extended Boundary Condition Approach" to appear in Applied Optics.
8. C. Yeh, K. Ha, S.B. Dong and W.P. Brown, Appl. Optics 18, 1490 (1979).
9. M. Imai, T. Asakura and Y. Kinoshita, Opt. and Quan. Elec. 7, 95 (1975).
10. Patent Pending.

SAMPLE PROGRAM LISTINGS FOR THE CASE OF BEAM  
PROPAGATION IN A STEP INDEX FIBER

CORE SIZE (DIAMETER) = 50 $\mu$ m  
CORE INDEX = 1.48  
CLADDING INDEX = 1.46  
INITIAL GAUSSIAN BEAM SPOT SIZE (DIAMETER) = 40 $\mu$ m  
WAVELENGTH = 0.8 $\mu$ m

(STEP INDEX CASE, m = 6)

LEVEL 2.3 (JUNE 78) 05/360 04TRAY H EXTENDED DATE 80.19' 1.57.04

REQUESTED OPTIONS:

OPTIONS IN EFFECT: NAME(MAIN) OPTIMIZE(2) LINECOUNT(50) SIZE(MAX) AUTODBL(NONE) SOURCE EBCDIC NOLIST NODECK OBJECT VJ4AS NOFORMAT GOSTMT NUXRE= ALC NOANSF NOTICRM IDM FLA

```

C THIS PROGRAM CONTAINS: CORRECTED = IRLJN=M(JPTFIB,PRINTER,PEAK,SIZE,
C HARM,GREYSC).
C
C PROGRAM JPTFIB( INPUT,OUTPUT,=I-E).FA*E5=INPUT,TAPE6=OUTPUT,
C TAPE7=FILED)
C
C INPUT PARAMETERS
C CARD 1 (IIS FORMAT)
C NCASES NUMBER OF CASES TO BE READ
C CARD 2 (NAMELIST FORMAT-$DEFAULT $
C LAMBDA WAVE LENGTH (MICRONS)
C RO 1/E POINT IN IRRADIANCE (MIC*JNS)
C FR FIBER RADIUS (MICRONS)
C NO REFRACTIVE INDEX
C PCORP PERCENT DROP AT R=FR OF V/UJ
C OUTRAD OUTER RADIUS (MICRONS)
C DX MESH SPACING (MICRONS)
C NSTEPS NUMBER OF Z- STEPS
C NZINC LENGTH OF Z- STEP = ZMIN/NZINC
C IDOUT DEVICE NUMBER FOR OUTPUT
C IGREY IF TRUE PRINT IRRADIANCE PROFILE AT EACH STEP
C PWALST IF TRUE WRITE 2ND MOMENTS AT EACH STEP
C PLYMAX IF TRUE PLOT 2ND MOMENTS VS DISTANCE
C PLYFLD IF TRUE PLOT FIELD INTENSITY VS DISTANCE
C PLYFILE IF TRUE DO ABOVE PLOT AT EVERY STEP
C MESH GRID SIZE (32, 64, OR 128)
C
C COMMON /LCM2/REFNDX(16384),SV(128),CS(128),ZTSQ(128),PQ53(128)
C AMPARY(16384),RADARY(16384)
C LEVEL 2,REFNDX,SM,CS,ZTSQ,PQ53,AMPARY,RADARY
C DIMENSION WQR(256),MM(3),INV(32),S(32)
C DIMENSION X(3),Y(3),PCDRPA(3),=NA(3),REFCA(3)
C COMMON /ARRAYS/V(132768)
C COMMON /PARAM/ZINC,MESH,LAMBDA,RU,FR,NO,PCORP,OUTRAD,DX,NSTEPS,
C NZINC,MESH50,WSHS02,PI,WAVENR,AKS1,XS,NS,MF,YF,MSHPIS
C COMMON /PRINT/PGREY,PWALST,PLYST,PLYMAX,LAST,IDUT,IGREY
C PLYFLD,PLYFILE
C
C REAL LAMBDA,NO,MDSQ,NZ
C LOGICAL PGREY,PWALST,PLYST,PLYMAX,PLYFLD,PLYFILE,PLYMAX,PLYLAST,CALLPM
C DATA ICNCS/1/,ICNT/0/
C
C NAMELIST /DEFAULT/LAMBDA,RO,NO,PCORP,JJTRAD,DX,NSTEPS,NZINC,
C IDUT,IGREY,PGREY,PWALST,PLYST,PLYMAX,PLYFLD,PLYFILE,MESH
  
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ISN 0002 COMMON /LCM2/REFNDX(16384),SV(128),CS(128),ZTSQ(128),PQ53(128)  
 ISN 0003 AMPARY(16384),RADARY(16384)  
 ISN 0004 LEVEL 2,REFNDX,SM,CS,ZTSQ,PQ53,AMPARY,RADARY  
 ISN 0005 DIMENSION WQR(256),MM(3),INV(32),S(32)  
 ISN 0006 DIMENSION X(3),Y(3),PCDRPA(3),=NA(3),REFCA(3)  
 ISN 0007 COMMON /ARRAYS/V(132768)  
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 ISM 0010 COMMON /PRINT/PGREY,PWALST,PLYST,PLYMAX,LAST,IDUT,IGREY  
 ISM 0011 PLYFLD,PLYFILE  
 REAL LAMBDA,NO,MDSQ,NZ  
 LOGICAL PGREY,PWALST,PLYST,PLYMAX,PLYFLD,PLYFILE,PLYMAX,PLYLAST,CALLPM  
 DATA ICNCS/1/,ICNT/0/  
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 IDUT,IGREY,PGREY,PWALST,PLYST,PLYMAX,PLYFLD,PLYFILE,MESH



```

15N 0012 READ(5,1000) NCASES
15N 0013 WRITE(6,1000) NCASES
15N 0014 READ(5,11)XB,YB
15N 0015 WRITE(6,11)XB,YB
15N 0016 READ(5,1000) NFIR
15N 0017 WRITE(6,1000) NFIR
15N 0018 READ(5,12)(X0(I),Y0(I),I=1,NFIR)
15N 0019 WRITE(6,12)(X0(I),Y0(I),I=1,NFIR)
15N 0020 READ(5,13)(PCORPA(K),K=1,NFIR)
15N 0021 WRITE(6,13)(PCORPA(K),K=1,NFIR)
15N 0022 READ(5,11)(FRACK(K),K=1,NFIR)
15N 0023 WRITE(6,11)(FRACK(K),K=1,NFIR)
15N 0024 1000 FORMAT(2X,12)
15N 0025 11 FORMAT(1X,F4.1,1X,F4.1)
15N 0026 12 FORMAT(1X,F4.1,1X,F4.1,1X,F4.1,1X,F4.1)
15N 0027 13 FORMAT(1X,F12.5,1X,F12.5)
C1
C READ(5,DEFAULT)
MM(1)= 7
MM(2)= 7
MM(3)= 0
I(CN)=0
FLAG=NJ/ABS(ND)
ND=ABS(ND)
2050 IF(FLAG.LT.0.) WRITE(10,2050)
FORMAT(/,47M THE REFRACTIVE INDEX IS A CONSTANT EQUAL TO NJ)
PCORP=PCORPA(1)
FR=FR(1)
WRITE(10,1,DEFAULT)
C
C CALCULATE CONSTANTS
MESR=2*MESH
MESHSQ=MESH**2
MHSO2=2*MESHSQ
RN2ND=.02*PCORP/FR**2
ZMIN=PI/12.*SORT(RN2ND)
DXSI=ZMIN/NDZINC
DXSIH=DXSI/2.
DZET=DX/RO
WAVENM=2.*PI/LAMBDA
BETAHT=(2.*ZMIN*DXSI/(WAVEN*ND))*(PI/(MESH*DZET*RO))**2
FITCNT=(1.-1./MESH)*PI
XYUMESH/2
RDNMI=10J/HAD/RO)**2
NDNR=1.*MESHSQ
NZNJ*RN2ND
REPCF=N2*RU**2/2.
ALPHA=2.*ZMIN/(PI*WAVEN*ND**4)**2
NOSQ=NJ**2
XSIMUL=DXSI
LAST=.FALSE.
IF(FLAG.LT.0.) REFCF=0.

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ISN 0063          CALLPR=PGREY,OR.PWAIST,OR.P_T=SF.34.2,_TRAX,OR.PLTFLD,OR.PLT=LE    OPTFIB
C          WRITE THE IMPORTANT CALCULATED PARAMETERS    OPTFIB
C          WRITE(IJOUT,2000) ZMIN,DZINC,MN2YJ,ALPHA    OPTFIB
ISN 0064          2000  * 7MICRONS,/.9H ZMIN = .F10+.4,IX,74(-30VS,/.9H DZINC = .F10+.4,IX,    OPTFIB
ISN 0065          * 9H ALPHA = .F10.5/)    OPTFIB
C          * 9H ALPHA = .F10.5/)    OPTFIB
C          CALCULATE NECESSARY ARRAYS    OPTFIB
C          MP=6    TEMP
DO 800 K=1,NFIB    TEMP
X0(K)=X0(K)*XYD    TEMP
Y0(K)=Y0(K)*XYD    TEMP
MX=2*MP    TEMP
REFCFA(K)=YD*.02*PCORPA(K)*(30/3A(K))*MX/2.    TEMP
DO 100 K=1,MESH    OPTFIB
IK=K    OPTFIB
ARG=FIXNST*PK    OPTFIB
CS(K)=COS(ARG)    OPTFIB
SN(K)=SIN(ARG)    OPTFIB
ZFSQ(K)=((RK-XYD)*DZET)**2    OPTFIB
POSQ(K)=((RK-XYD+.5)**2    OPTFIB
CONTINUE    OPTFIB
C          SET UP REFRACTIVE INDEX ARRAY    OPTFIB
C          M=0    TEMP
DO 120 J=1,MESH    TEMP
DC 120 I=1,MESH    TEMP
N=M+1    TEMP
TMPNDX=0.    TEMP
DO 110 K=1,NFIB    TEMP
Z1=((J-1.-Y0(K))*DZET)**2    TEMP
Z2=((I-1.-X0(K))*DZET)**2    TEMP
RAD=Z1+Z2    TEMP
TMPNDA=A*MAXI(TMPNDX,(NO-REFCFA(K))*((NAJ**MPI)))    TEMP
REF=NUS(1.-D.01*PCORPA(K))    TEMP
REFN(X)=A*MAXI(TMPNDX,CREF)    TEMP
CONTINUE    TEMP
C          CREF=.52.03.(.EQ.58) WRITE(IJOUT,*) REFNDX(M)    TEMP
C          CONTINUE,N=128) GO TO 10    OPTFIB
IF(MESH.NE.128) GO TO 10    OPTFIB
MS=MESH/4+1    OPTFIB
MF=MS+MESH/2-1    OPTFIB
NS=NS    OPTFIB
NF=NF    OPTFIB
GO TO 40    OPTFIB
CONTINUE    OPTFIB
MS=1    OPTFIB
MF=MESH    OPTFIB
ISN 0066          ISN 0061    OPTFIB
ISN 0067          ISN 0062    OPTFIB
ISN 0068          ISN 0063    OPTFIB
ISN 0069          ISN 0064    OPTFIB
ISN 0070          ISN 0065    OPTFIB
ISN 0071          ISN 0066    OPTFIB
ISN 0072          ISN 0067    OPTFIB
ISN 0073          ISN 0068    OPTFIB
ISN 0074          ISN 0069    OPTFIB
ISN 0075          ISN 0070    OPTFIB
ISN 0076          ISN 0071    OPTFIB
ISN 0077          ISN 0072    OPTFIB
ISN 0078          ISN 0073    OPTFIB
ISN 0079          ISN 0074    OPTFIB
ISN 0080          ISN 0075    OPTFIB
ISN 0081          ISN 0076    OPTFIB
ISN 0082          ISN 0077    OPTFIB
ISN 0083          ISN 0078    OPTFIB
ISN 0084          ISN 0079    OPTFIB
ISN 0085          ISN 0080    OPTFIB
ISN 0086          ISN 0081    OPTFIB
ISN 0087          ISN 0082    OPTFIB
ISN 0088          ISN 0083    OPTFIB
ISN 0089          ISN 0084    OPTFIB
ISN 0090          ISN 0085    OPTFIB
ISN 0091          ISN 0086    OPTFIB
ISN 0092          ISN 0087    OPTFIB
ISN 0093          ISN 0088    OPTFIB
ISN 0094          ISN 0089    OPTFIB
ISN 0095          ISN 0090    OPTFIB
ISN 0096          ISN 0091    OPTFIB
ISN 0097          ISN 0092    OPTFIB
ISN 0098          ISN 0093    OPTFIB
ISN 0099          ISN 0094    OPTFIB
ISN 0100          ISN 0095    OPTFIB
ISN 0101          ISN 0096    OPTFIB
ISN 0102          ISN 0097    OPTFIB
ISN 0103          ISN 0098    OPTFIB
ISN 0104          ISN 0099    OPTFIB
ISN 0105          ISN 0100    OPTFIB

```



LEVEL 2.3.	( JUNE 78 )	MAIN	05/360	JATRAN H EXTENDED	DATE 80.194	1.57.04
ISN 0150		V(K)=VR*AR-VI*AI				OPTFIB
ISN 0151		V(KPI)=VI*AI+VR*AI				OPTFIB
ISN 0152	160	CONTINUE				OPTFIB
	C					OPTFIB
	C	DO TRANS=DRN				OPTFIB
ISN 0153		CALL HARM(V,MM,INV,S,1,IFERR)				OPTFIB
	C					OPTFIB
	C	SOLVE FIRST ORDER ODE				OPTFIB
ISN 0154		K=-1				OPTFIB
ISN 0155		DO 190 J=1,MESH				OPTFIB
ISN 0156		PHI1=SETAMT*POSQ(J)				OPTFIB
ISN 0157		DO 190 I=1,MESH				OPTFIB
ISN 0158		K=K+2				OPTFIB
ISN 0159		KPI=K+1				OPTFIB
ISN 0160		PHI2=BETAMT*POSQ(I)				OPTFIB
ISN 0161		VR=V(I)				OPTFIB
ISN 0162		VI=V(I*PI)				OPTFIB
ISN 0163		ANG2=(PHI1*PHI2)				OPTFIB
ISN 0164		CANG=CDS(ANG)				OPTFIB
ISN 0165		SANG=SLY(ANG)				OPTFIB
ISN 0166		V(K)=IVR*CANG-VI*SANG)				OPTFIB
ISN 0167		V(KPI)=(VR*SANG+VI*CANG)				OPTFIB
ISN 0168	180	CONTINUE				OPTFIB
	C					OPTFIB
	C	DO INVERSE TRANSFORM				OPTFIB
ISN 0169		CALL HARM(V,MM,INV,S,-1,IFERR)				OPTFIB
	C					OPTFIB
	C	RECONDITION V BECAUSE OF TRANSFORM				OPTFIB
ISN 0170		K=-1				OPTFIB
ISN 0171		DO 200 J=1,MESH				OPTFIB
ISN 0172		SAJ=SM(J)				OPTFIB
ISN 0173		CSJ=CS(J)				OPTFIB
ISN 0174		RZ1=(1-J-K*YO-1)*DZET)*2				OPTFIB
ISN 0175		CLOSSE=2				OPTFIB
ISN 0176		DO 200 I=1,MESH				OPTFIB
ISN 0177		BZ2=(1-I*YO-1)*DZET)*2				OPTFIB
ISN 0178		PRAD=BZ1*BZ2				OPTFIB
ISN 0179		CM=1				OPTFIB
ISN 0180		IF(BZ1*AD.GT.HADNRM) CM=EXP(-C.USS*(HAD-RADNRM))				OPTFIB
ISN 0181		K=K+2				OPTFIB
ISN 0182		KPI=K+1				OPTFIB
ISN 0183		SI=SM(I)				OPTFIB
ISN 0184		CS1=CS(I)				OPTFIB
ISN 0185		AR=CSJ*CS1-SNJ*SI				OPTFIB
ISN 0186		AI=CSJ*SI+SNJ*CS1				OPTFIB
ISN 0187		VR=V(K)*CM				OPTFIB
ISN 0188		VI=V(KPI)*CM				OPTFIB
ISN 0189		V(K)=VR*AR-VI*AI				OPTFIB
ISN 0190		V(KPI)=VR*AI+VI*AR				OPTFIB
ISN 0191		CONTINUE				OPTFIB
ISN 0192	200					OPTFIB



LEVEL 2.3. (JUNE 79) OS/360 J3RTRAN H EXTENDED DATE 80.195 1.57.07

REQUESTED OPTIONS:

OPTIONS IN EFFECT: NAME(MAIN) OPTIMIZE(2) LINECOUNT(56) SIZE(MAX) AUTODBL(NONE) SOURCE EBCDIC NLIST NOBECK OBJECT NJ4AP NOFORMAT GOSTMT NOXREF ALC NOANSF NOTERM IBM FLA

```

ISM 0002 SUBROUTINE PEAK(VMAX)
ISM 0003 COMMON /LCM2/REFNDX(16384),N(123),CS(129),ZTSO(128),POS(128)
      * .ARRAY(16384),RADARY(16384)
      * .LEVEL 2,REFNDX,SM,CS,ZTSO,PJSU,APADY,RADARY
      * COMMON /ARRAYS/V(32768)
      * COMMON /ARRAY/DZINC,MSH,LANJDA,Q,FR,N),PCOMP,OUTRAD,DX,NSTEPS.
      * NOZINC,NEHSO,MSHSO2,PI,WAVERN,DXS(L,MS,NS,MF,NF,MSHPTS
      * L=0
      * SUML=0.
      * SUMR=0.
      * VMAX=0.
      * DO 10 K=1,MSHSO2*2
      *   VR=V(K)
      *   L=L+1
      *   KPI=K+1
      *   VI=V(KPI)
      *   VRAD=VI**2+VR**2
      *   IF(K.LE.NEHSO) SUML=SUML+VRAD
      *   IF(K.LE.NEHSO) SUMR=SUMR+VRAD
      *   VMAX=MAX(VMAX,VRAD)
      * RADARY(L)=VRAD
      * CONTINUE
      * TOT=SUML+SUMR
      * SUML=SUML/TOT
      * SUMR=SUMR/TOT
      * WRITE(6,2000) SUML,SUMR
      * 2000 FORMAT(1X,7H$UML = ,E14.7,X,7H$JMR = ,E14.7,/)
      * RETURN
      * END
  10
  2000
  
```

OPTIONS IN EFFECT: NAME(MAIN) OPTIMIZE(2) LINECOUNT(56) SIZE(MAX) AUTODBL(NONE)  
 OPTIONS IN EFFECT: SOURCE EBCDIC NLIST NOBECK OBJECT NJ4AP NOFORMAT GOSTMT NOXREF ALC NOANSF NOTERM IBM FLA  
 \*STATISTICS\* SOURCE STATEMENTS = 28, PROGRAM SIZE = 540, SUBPROGRAM NAME = PEAK  
 \*STATISTICS\* NO DIAGNOSTICS GENERATED  
 \*\*\*\*\* END OF COMPILATION \*\*\*\*\*

112K BYTES OF CORE NOT USED

LEVEL 2.3. (JUNE 78) OS/360 JATRAM 1 EXTENDED DATE 80.199 .57.08

REQUESTED OPTIONS:

OPTIONS IN EFFECT: NAME(MAIN) OPTIMIZE(2) LINECOUNT(55) SIZE(MAX) AUTODBL(NONE) SOURCE EBCDIC NDLIST NOCHECK OBJECT VJMAP NOFORMAT GOSTMT NOXHEF ALC NOANSF NOTERM IRM FL/

```

ISN 0002 SURROUTINE SIZE(AARRAY,MESH,MESH,XO,YO,X2,Y2) SIZE
ISN 0003 DIMENSION AARRAY(MESH,MESH) SIZE
ISN 0004 MID=MESH/2+1 SIZE
ISN 0005 SUM1=0. SIZE
ISN 0006 SUM2=0. SIZE
ISN 0007 SUM3=0. SIZE
ISN 0008 SUM4=0. SIZE
ISN 0009 SUM5=0. SIZE
ISN 0010 DO 10 N=1,MESH SIZE
ISN 0011 RN=N-MID SIZE
ISN 0012 RMSQ=RN**2 SIZE
ISN 0013 DO 10 M=1,MESH SIZE
ISN 0014 RM=M-MID SIZE
ISN 0015 RMSO=RM**2 SIZE
ISN 0016 ARMN = ARRAY(M,N) SIZE
ISN 0017 SUM1=SUM1+RM*ARMN SIZE
ISN 0018 SUM2=SUM2+RN*ARMN SIZE
ISN 0019 SUM3=SUM3+RMSO*ARMN SIZE
ISN 0020 SUM4=SUM4+RMSO*ARMN SIZE
ISN 0021 SUM5=SUM5+ARMN SIZE
ISN 0022 CONTINUE SIZE
ISN 0023 SNDRM=1./SUM5 SIZE
ISN 0024 XO=SNDRM*SUM2 SIZE
ISN 0025 YO=SNDRM*SUM1 SIZE
ISN 0026 X2=SNDRM*SUM4-XO**2 SIZE
ISN 0027 Y2=SNDRM*SUM3-YO**2 SIZE
ISN 0028 X2=SORT(X2) SIZE
ISN 0029 Y2=SORT(Y2) SIZE
ISN 0030 RETURN SIZE
ISN 0031 END SIZE

```

OPTIONS IN EFFECT: NAME(MAIN) OPTIMIZE(2) LINECOUNT(55) SIZE(MAX) AUTODBL(NONE)

OPTIONS IN EFFECT: SOURCE EBCDIC NDLIST NOCHECK OBJECT VJMAP NOFORMAT GOSTMT NOXHEF ALC NOANSF NOTERM IRM FLA

\*STATISTICS\* SOURCE STATEMENTS = 30. PROGRAM SIZE = 786. SUBPROGRAM NAME = SIZE

\*STATISTICS\* NO DIAGNOSTICS GENERATED

\*\*\*\*\* END OF COMPILATION \*\*\*\*\*

112K BYTES OF CORE NOT USED

LEVEL 2.3.3 (JUNE 78)

OS/350 UNTRAN H EXTENDED

DATE 00.19 1.57.09

REQUESTED OPTIONS:

OPTIONS IN EFFECT: NAME(MAIN) OPTIMIZE(2) LINECOUNT(55) SIZE(MAX) AUTODBL(NONE)  
SOURCE EBCDIC NOLIST NODECK OBJECT VJ413 NOFORMAT COSTMT NOXREF ALC NOANSF NOTFRM (DM FL/

```

ISN 0002      BLOCK DATA
ISN 0003      COMMON /ARRAYS/V(32768)
ISN 0004      COMMON /PARAM/DZINC,MESH,LAMBDA,JD,FJ,NO,PCDRP,OUTRAD,DX,NSTEPS,
ISN 0005      COMMON /MESH/MESH02,PI,WA,VE,NO,DX,DI,MS,NS,MF,MF,MSHPTS,
ISN 0006      COMMON /PRINT/PGREY,PWAIST,PLTWST,PLTMAX,LAST,IGREY
ISN 0007      COMMON /PLTFLO,PLTFLE
ISN 0008      LOGICAL SGREY,PWAIST,PLTWST,PLTMAX,LAST
ISN 0009      LOGICAL PLTFLO,PLTFLE
ISN 0010      REAL LAMBDA,NO
ISN 0011      DATA PGRV,PWAIST,PLTWST,PLTMAX,AST,IGREY/2*,TRUE,2*,FALSE,
ISN 0012      DATA LAMBDA,NO,NO,PCDRP,OUTRAD,DX,NSTEPS,NOZINC
ISN 0013      DATA PLTFLO,PLTFLE/,FALSE,0.,1E+01,40.5/
            DATA MESH/128/,PI/3.14159265/
            END
  
```

```

OPTIONS IN EFFECT:NAME(MAIN) OPTIMIZE(2) LINECOUNT(55) SIZE(MAX) AUTODBL(NONE)
OPTIONS IN EFFECT:SOURCE EBCDIC NOLIST NODECK OBJECT VJ413 NOFORMAT COSTMT NOXREF ALC NOANSF NOTERM IRM FL
STATISTICS#  SOURCE STATEMENTS = 12, PROGRAM SIZE = 0, SUBPROGRAM NAME #ARRAYS
STATISTICS#  NO DIAGNOSTICS GENERATED
***** END OF COMPILATION *****
112K BYTES OF CORE NOT USED
  
```



REQUESTED OPTIONS:  
 OPTIONS IN EFFECT:

NAME(MAIN) OPTIMIZE(2) LINECOUNT(50) SIZE(MAX) AUTODRL(NONE)  
 SOURCE EBCDIC NOLIST NODECK OBJECT VJ4AP NDFORMAT COSTNT NOKREF ALC NOANSF NOTERM (IBM FL)

```

SUBROUTINE HARM
PURPOSE
PERFORMS DISCRETE COMPLEX FOURIER TRANSFORMS ON A COMPLEX
THREE DIMENSIONAL ARRAY
USAGE
CALL HARM (A,M,INV,S,IFSET,IFERR)
DESCRIPTION OF PARAMETERS
A - AS INPUT, A COMPLEX, 3-DIMENSIONAL
ARRAY TO BE TRANSFORMED, THE REAL PART OF
ALL (1,2,3) IS STORED IN VECTOR FASHION IN A CELL
WITH INDEX 2*(130N1+12*N2+11)+1 WHERE
N1 = 2*(M(I)-1)+2,3 AND N2 = 0,1,...,M1-1 ETC.
THE IMAGINARY PART IS IN THE CELL IMMEDIATELY
FOLLOWING. NOTE THAT THE SUBSCRIPT II INCREASES
MOST RAPIDLY AND I INCREASES LEAST RAPIDLY.
AS OUTPUT, A CONTAINS THE COMPLEX FOURIER
TRANSFORM. THE NUMBER OF CORE LOCATIONS OF
ARRAY A IS 2*(130N1+12*N2+11)
M - A THREE CELL VECTOR WHICH DETERMINES THE SIZES
OF THE 3 DIMENSIONS OF THE ARRAY. THE SIZE
OF THE I DIMENSION OF A IS 2*(M(I)-1)+2,3
INV - A VECTOR WORK AREA FOR FIT AND INDEX MANIPULATION
OF DIMENSION ONE-JUITY OF THE QUANTITY
MAX(130N1,12*N2)
S - A VECTOR WORK AREA FOR SINE TABLES WITH DIMENSION
THE SAME AS INV
IFSET - AN OPTION PARAMETER WITH THE FOLLOWING SETTINGS
0 SET UP SINE AND INV TABLES ONLY
1 CALCULATE FOURIER TRANSFORM ONLY AND
INV TABLES ONLY
-1 CALCULATE INVERSE FOURIER TRANSFORM ONLY AND
INV TABLES ONLY
2 CALCULATE INVERSE FOURIER TRANSFORM (FOR
UNDER 4E+13) BELOW
SINE AND INV TABLES EXIST)
-2 CALCULATE INVERSE FOURIER TRANSFORM ONLY
(CASSUME SINE AND INV TABLES EXIST)
IFERR - ERROR INDICATOR.
IFERR = 1 MEANS THE MAXIMUM M(I) IS GREATER THAN
20, I=1,2,3. WHEN IFSET IS 2,-2, IFERR = 1
MEANS THAT THE SINE AND INV TABLES ARE NOT LARGE
ENOUGH OR HAVE NOT BEEN COMPUTED.
IF ON RETURN IFERR = 0 THEN NONE OF THE ABOVE
CONDITIONS ARE PRESENT
HARM 40
HARM 50
HARM 60
HARM 70
HARM 80
HARM 90
HARM 100
HARM 110
HARM 120
HARM 130
HARM 140
HARM 150
HARM 160
HARM 170
HARM 180
HARM 190
HARM 200
HARM 210
HARM 220
HARM 230
HARM 240
HARM 250
HARM 260
HARM 270
HARM 280
HARM 290
HARM 300
HARM 310
HARM 320
HARM 330
HARM 340
HARM 350
HARM 360
HARM 370
HARM 380
HARM 390
HARM 400
HARM 410
HARM 420
HARM 430
HARM 440
HARM 450
HARM 460
HARM 470
HARM 480
HARM 490
HARM 500
HARM 510
HARM 520
  
```

REMARKS  
 THIS SUBROUTINE IS TO BE USED FOR COMPLEX, 3-DIMENSIONAL  
 ARRAYS IN WHICH EACH DIMENSION IS A POWER OF 2. THE  
 MAXIMUM N(1) MUST NOT BE LESS THAN 3 OR GREATER THAN 20.  
 I = 1,2,3

SUBROUTINES AND FUNCTION  $\rightarrow$  J8R10GNA4S REQUIRED

NONE

METHOD

FOR IFFSET = +1, OR +2, THE FOURIER TRANSFORM OF COMPLEX  
 ARRAY A IS OBTAINED.

$$X(J1,J2,J3) = \sum_{K1=0}^{N1-1} \sum_{K2=0}^{N2-1} \sum_{K3=0}^{N3-1} A(K1,K2,K3) * W1^{L1} * W2^{L2} * W3^{L3}$$

WHERE W1 IS THE N(1) ROOT OF UNITY AND  $L1=K1*J1$ ,  
 $L2=K2*J2$ ,  $L3=K3*J3$

FOR IFFSET = -1, OR -2, THE INVERSE FOURIER TRANSFORM A OF  
 COMPLEX ARRAY X IS OBTAINED.

$$A(K1,K2,K3) = \frac{1}{N1*N2*N3} \sum_{J1=0}^{N1-1} \sum_{J2=0}^{N2-1} \sum_{J3=0}^{N3-1} X(J1,J2,J3) * W1^{-L1} * W2^{-L2} * W3^{-L3}$$

SEE J.W. COOLEY AND J.W. TUCKER, AN ALGORITHM FOR THE  
 MACHINE CALCULATION OF COMPLEX FOURIER SERIES,  
 MATHEMATICS OF COMPUTATIONS, VOL. 19 (APR. 1965), P. 297.

ISN 0002  
 ISN 0003  
 ISN 0004  
 ISN 0005  
 ISN 0006  
 ISN 0007  
 ISN 0008  
 ISN 0009  
 ISN 0010  
 ISN 0011  
 ISN 0012  
 ISN 0013  
 ISN 0014  
 ISN 0015  
 ISN 0016  
 ISN 0017  
 ISN 0018  
 ISN 0019

.....  
 SUBROUTINE HARM(A,M,INVA5,IFFSET, IFFERR)  
 DIMENSION A(1),INV(1),S(1),M(3),NP(3),W(2),W2(2),W3(2)  
 EQUIVALENCE (N1,N(1)),(N2,N(2)),(N3,N(3))  
 IF( IABS( IFFSET ) - 1) 900,900,12  
 10 MTT=MAX0(M(1),M(2),M(3))-2  
 12 ROOT2 = SQRT(2.)  
 IF (MT-MT) 14,14,13  
 13 IFFERR=1  
 14 RETURN  
 IFFERR=0  
 W1=M(1)  
 W2=M(2)  
 W3=M(3)  
 N1=2\*\*M1  
 N2=2\*\*M2  
 N3=2\*\*M3  
 16 IF( IFFSET ) 16,16,20  
 18 NX= N1\*N2\*N3

HARM 531  
 HARM 580  
 HARM 580  
 HARM 580  
 HARM 570  
 HARM 580  
 HARM 590  
 HARM 600  
 HARM 610  
 HARM 620  
 HARM 630  
 HARM 640  
 HARM 650  
 HARM 660  
 HARM 670  
 HARM 680  
 HARM 690  
 HARM 700  
 HARM 710  
 HARM 720  
 HARM 730  
 HARM 740  
 HARM 750  
 HARM 760  
 HARM 770  
 HARM 780  
 HARM 790  
 HARM 800  
 HARM 810  
 HARM 820  
 HARM 830  
 HARM 840  
 HARM 850  
 HARM 860  
 HARM 870  
 HARM 880  
 HARM 890  
 HARM 900  
 HARM 910  
 HARM 920  
 HARM 930  
 HARM 940  
 HARM 950  
 HARM 960  
 HARM 970  
 HARM 980  
 HARM 990  
 HARM 1000  
 HARM 1010  
 HARM 1020  
 HARM 1030  
 HARM 1040  
 HARM 1050  
 HARM 1060

LEVEL 2. J (JUNE 78) HARM OS/360 J3TRAN H EXTENDED DATE 80.19 21.57.09

```

ISN 0020 FN = NX
ISN 0021 DD 19 I = I,NX
ISN 0022 A(2*1) = A(2*(I-1))/FN
ISN 0023 A(2*1) = -A(2*I)/FN
ISN 0024 NP(1) = N1*2
ISN 0025 NP(2) = NP(1)*N2
ISN 0026 NP(3) = NP(2)*N3
ISN 0027 DD 250 I = I+3
ISN 0028 IL = NP(3) - NP(1)
ISN 0029 I1 = I1+1
ISN 0030 MI = M(I)
ISN 0031 IF (N1)250.250.30
ISN 0032 IDIF = NP(1)
ISN 0033 KRIT = NP(1)
ISN 0034 MEV = 2*(M1/2)
ISN 0035 IF (N1 - MEV)60.60.40

C M IS 000: DD L=1 CASE
C 40 KRIT = 8(I/2)
C KL = KBIT-2
C DD 50 I = I, I1, IDIF
C KLAST = KL+1
C DD 50 K = (.KLAST*.2
C KD = K*KBIT

C DD ONE STEP WITH L=1,J=0
C A(K) = A(K)*A(KD)
C A(KD) = A(K) - A(KD)

T=A(D)
A(KD) = A(K) - T
A(K) = A(K)+T
T = A(D+1)
A(KD+1) = A(K+1) - T
A(K+1) = A(K)+T
IF (N1 - 1)250.250.52
52 LFIRST = 3

C DEF - JLAST = 2*(L-2) - 1
C JLAST = 1
C GO TO 70

C M IS EVEN
C LFIRST = 2
C JLAST = 0
C DD 200 L = LFIRST, MI, 2
C JJOIF = KBIT
C KRIT = KBIT/A
C KL = KBIT-2

C DD EJM J=0
C DD 80 I = I, I1, IDIF
C KLAST = I+KL
C DD 80 K = I, KLAST.2

```

HARM1070  
HARM1080  
HARM1090  
HARM1100  
HARM1110  
HARM1120  
HARM1130  
HARM1140  
HARM1150  
HARM1160  
HARM1170  
HARM1180  
HARM1190  
HARM1200  
HARM1210  
HARM1220  
HARM1230  
HARM1240  
HARM1250  
HARM1260  
HARM1270  
HARM1280  
HARM1290  
HARM1300  
HARM1310  
HARM1320  
HARM1330  
HARM1340  
HARM1350  
HARM1360  
HARM1370  
HARM1380  
HARM1390  
HARM1400  
HARM1410  
HARM1420  
HARM1430  
HARM1440  
HARM1450  
HARM1460  
HARM1470  
HARM1480  
HARM1490  
HARM1500  
HARM1510  
HARM1520  
HARM1530  
HARM1540  
HARM1550  
HARM1560  
HARM1570  
HARM1580  
HARM1590  
HARM1600

ISM 0061 K1K+KBIT  
 ISM 0062 K2=K1+KBIT  
 ISM 0063 K3=K2+KBIT  
  
 DO TWO STEPS WITH J=0  
 A(K)=A(K)+A(K2)  
 A(K2)=A(K)-A(K2)  
 A(K1)=A(K1)+A(K3)  
 A(K3)=A(K1)-A(K3)  
  
 A(K)=A(K)+A(K1)  
 A(K1)=A(K)-A(K1)  
 A(K2)=A(K2)+A(K3)+1  
 A(K3)=A(K2)-A(K3)+1  
  
 T=A(K2)  
 A(K2)=A(K)-T  
 A(K)=A(K)+T  
 T=A(K2+1)  
 A(K2+1)=A(K1)+T  
 A(K+1)=A(K+1)+T  
  
 T=A(K3)  
 A(K3)=A(K1)-T  
 A(K1)=A(K1)+T  
 T=A(K3+1)  
 A(K3+1)=A(K1)+T  
 A(K1+1)=A(K1)+T  
  
 T=A(K1)  
 A(K1)=A(K1)-T  
 A(K)=A(K)+T  
 T=A(K1+1)  
 A(K1+1)=A(K+1)+T  
 A(K+1)=A(K+1)+T  
  
 R=-A(K3+1)  
 T=A(K3)  
 A(K3)=A(K2)-R  
 A(K2)=A(K2)+R  
 A(K3+1)=A(K2+1)-T  
 A(K2+1)=A(K2+1)+T  
 IF (JLAST) 235,235,02  
 60 JJ=JJ+1  
 62 JJ=JJ+1  
  
 DO FOR J=1  
 ILAST = IL + JJ  
 DO 65 I = JJ, ILAST, IDIF  
 KLAST = KL + I  
 DO 65 K = I, KLAST, 2  
 K1 = K+KBIT  
 K2 = K1+KBIT  
 K3 = K2+KBIT  
  
 ISM 0064  
 ISM 0065  
 ISM 0066  
 ISM 0067  
 ISM 0068  
 ISM 0069  
  
 ISM 0070  
 ISM 0071  
 ISM 0072  
 ISM 0073  
 ISM 0074  
 ISM 0075  
  
 ISM 0076  
 ISM 0077  
 ISM 0078  
 ISM 0079  
 ISM 0080  
 ISM 0081  
  
 ISM 0082  
 ISM 0083  
 ISM 0084  
 ISM 0085  
 ISM 0086  
 ISM 0087  
 ISM 0088  
 ISM 0089  
  
 ISM 0090  
 ISM 0091  
 ISM 0092  
 ISM 0093  
 ISM 0094  
 ISM 0095  
 ISM 0096

LEVEL 2.3. - ( JUNE 78) HARM DS/350 JATKAN H EXTENDED DATE 80.199 1.57.09

```

C LETTING W=(1+I)/ROOT2, W3=(-1+I)/ROOT2, W2=I,
C A(K1)=A(K)+A(K2)*I
C A(K2)=A(K)-A(K2)*I
C A(K3)=A(K1)+W*A(K3)*W3
C A(K3)=A(K1)+W-A(K3)*W3
C A(K3)=A(K)+A(K1)
C A(K1)=A(K)-A(K1)
C A(K2)=A(K2)+A(K3)*I
C A(K3)=A(K2)-A(K3)*I
C R =-A(K2+I)
C T =A(K2)
C A(K2) = A(K)-R
C A(K) = A(K)+R
C A(K2+I)=A(K+I)-T
C A(K+I)=A(K+I)+T
C
C AMR=A(K1)-A(K1+I)
C AMI = A(1+I)+A(K1)
C R =-A(K3)-A(K3+I)
C T=A(K3)-A(K3+I)
C A(K3)=(AMR-H)/ROOT2
C A(K3+I)=(AMI-T)/ROOT2
C A(K1)=(AMR+H)/ROOT2
C A(K1+I)=(AMI+T)/ROOT2
C T = A(K1)
C A(K1)=A(K1)-T
C A(K) =A(K)+T
C T=A(K1+I)
C A(K1+I)=A(K+I)-T
C A(K+I)=A(K+I)+T
C R=-A(K3+I)
C T=A(K3)
C A(K3)=A(K2)+R
C A(K2)=A(K2)+R
C A(K3+I)=A(K2+I)-T
C A(K2+I)=A(K2+I)+T
85 IFL JLAST-I) 235+235.90
90 JJE JJ + JJDI
C
C NOW DO THE REMAINING J'S
C DO 230 J=2, JLAST
C
C FEIC W'S
C DEF W=W**INV(J), W2=W**2, WJ=W**3
C 96 IC=NI-I
C 98 W(1)=S(IC)
C W(2)=S(I)
C I2=2*I
C IFC=NI-I2
C IFL I2C) I20.110.100
C

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ISM 0097  
ISM 0098  
ISM 0099  
ISM 0100  
ISM 0101  
ISM 0102  
ISM 0103  
ISM 0104  
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ISM 0127  
ISM 0128  
ISM 0129  
ISM 0130  
ISM 0131  
ISM 0132

HARM2150  
HARM2160  
HARM2170  
HARM2180  
HARM2190  
HARM2200  
HARM2210  
HARM2220  
HARM2230  
HARM2240  
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HARM2500  
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HARM2640  
HARM2650  
HARM2660  
HARM2670  
HARM2680

LEVEL 2	(JUNE 78)	HARM	OS/350	ORTRAN H EXTENDED	DATE	80.19	21.57.09
ISM 0133	C	291 IS IN FIRST QUADRANT					
ISM 0134	100	W2(1)=S(12C)					HARM2690
ISM 0135		W2(2)=S(12)					HARM2710
ISM 0136		GO TO 130					HARM2720
ISM 0137	110	W2(1)=0.					HARM2730
ISM 0138		W2(2)=1.					HARM2740
		GO TO 130					HARM2750
ISM 0139	C	291 IS IN SECOND QUADRANT					HARM2760
ISM 0140	120	I2CC = I2CNT					HARM2770
ISM 0141		I2C = I2C					HARM2780
ISM 0142		W3(1)=S(12C)					HARM2790
ISM 0143	130	W3(2)=S(12CC)					HARM2800
ISM 0144		I2CNT = I3					HARM2810
ISM 0145		IF(I3C)160.150.140					HARM2820
ISM 0146	C	I3 IN FIRST QUADRANT					HARM2830
ISM 0147	140	W3(1)=S(13C)					HARM2840
ISM 0148		W3(2)=S(13)					HARM2850
ISM 0149	150	GC TO 200					HARM2860
ISM 0150		W3(1)=0.					HARM2870
ISM 0151		W3(2)=1.					HARM2880
		GO TO 200					HARM2890
ISM 0152	C	I3CC=I3CNT					HARM2900
ISM 0153	160	IF(I3CC)190.180.170					HARM2910
ISM 0154	C	I3 IN SECOND QUADRANT					HARM2920
ISM 0155	170	I3C = I3C					HARM2930
ISM 0156		W3(1)=S(13C)					HARM2940
ISM 0157		W3(2)=S(13CC)					HARM2950
ISM 0158	180	GC TO 200					HARM2960
ISM 0159		W3(1)=1.					HARM2970
ISM 0160		W3(2)=0.					HARM2980
		GO TO 200					HARM2990
ISM 0161	C	396 IN THIRD QUADRANT					HARM3000
ISM 0162	190	I3CC=I3CNT					HARM3010
ISM 0163		I3CC = I3CC					HARM3020
ISM 0164		W3(1)=S(13CC)					HARM3030
ISM 0165	200	W3(2)=S(13CC)					HARM3040
ISM 0166		LAST=I+JJ					HARM3050
ISM 0167		DD 220 I=JJ.LAST.IDIF					HARM3060
ISM 0168		KLAST=K+I					HARM3070
ISM 0169		DD 220 K=I.KLAST.2					HARM3080
ISM 0170		K=K+KBIT					HARM3090
ISM 0171		K2=K+KBIT					HARM3100
		K3=K2+KBIT					HARM3110
	C	DD TWO STEPS WITH J NOT 0					HARM3120
	C	A(KJ)=A(KJ)+A(K2)*M2					HARM3130
	C	A(K2)=A(K)-A(K2)*M2					HARM3140
	C	A(K1)=A(K1)+M*A(K3)*M3					HARM3150
							HARM3160
							HARM3170
							HARM3180
							HARM3190
							HARM3200
							HARM3210
							HARM3220

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LEVEL 2.3.0 ( JUNE 78)          HARM          DS/360          J3TRAN H EXTENDED          DATE 80.195 .1.57.09
C          A(K3)=A(K1)*W-A(K3)*W3
C          A(K1)=A(K1)+A(K1)
C          A(K1)=A(K1)-A(K1)
C          A(K2)=A(K2)+A(K3)*I
C          A(K3)=A(K2)-A(K3)*I
C
ISN 0172      R=A(K2)*W2(1)-A(K2+1)*W2(2)
ISN 0173      T=A(K2)*W2(2)+A(K2+1)*W2(1)
ISN 0174      A(K2)=A(K)-R
ISN 0175      A(K)=A(K)+R
ISN 0176      A(K2+1)=A(K+1)-T
ISN 0177      A(K+1)=A(K+1)+T
C
ISN 0178      R=A(K3)*W3(1)-A(K3+1)*W3(2)
ISN 0179      T=A(K3)*W3(2)+A(K3+1)*W3(1)
ISN 0180      WR=A(K1)*W(1)-A(K1+1)*W(2)
ISN 0181      WI=A(K1)*W(2)+A(K1+1)*W(1)
ISN 0182      A(K3)=A(K2)-R
ISN 0183      A(K3+1)=A(K1)-T
ISN 0184      A(K1)=A(K1)+T
ISN 0185      A(K1)=A(K1)-T
ISN 0186      A(K)=A(K)+T
ISN 0187      T=A(K1)+T
ISN 0188      A(K1+1)=A(K+1)+T
ISN 0189      A(K+1)=A(K+1)+T
ISN 0190      R=-A(K3+1)
ISN 0191      T=A(K3)
ISN 0192      A(K3)=A(K2)-R
ISN 0193      A(K2)=A(K2)+R
ISN 0194      A(K3+1)=A(K2+1)-T
ISN 0195      A(K2+1)=A(K2+1)+T
ISN 0196      A(K2+1)=A(K2+1)+T
ISN 0197      END JF I AND K LOOPS
C
ISN 0198      JJ=JJDIF+JJ
C          END JF J-LDOP
C
ISN 0199      JLAST=JLAST+3
ISN 0200      CONTINUE L LOOP
C          END DF L LOOP
C
ISN 0201      C 250 CONTINUE ID LOOP
C          END DF ID LOOP
C
ISN 0202      WE NOW HAVE THE COMPLEX FOURIER SUMS BUT THEIR ADDRESSES ARE
ISN 0203      BIT-REVERSED. THE FOLLOWING ROUTINE PUTS THEM IN ORDER
ISN 0204      N3QNT=N3-4I
C          N3=INT(N3/4I)
C          N3 GR. ON EQ. MT
C          I G03=I

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HARM3230
HARM3240
HARM3250
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HARM3270
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HARM3740
HARM3750
HARM3760

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ISM 0206      N3VNT=N3/NT
ISM 0207      MINN3=NT
ISM 0208      GO TO 380

ISM 0209      C 370 M3 LESS THAN MT
ISM 0210      N3VNT=1
ISM 0211      NTVN3=NT/N3
ISM 0212      MINN3=N3
ISM 0213      JJD3 = NY50/N3
ISM 0214      *2M=M2-MT
ISM 0215      450 IF (M2MT)470,460,460

ISM 0216      C 460 M2 GR. OR EQ. MT
ISM 0217      IG02=1
ISM 0218      N2VNT=N2/NT
ISM 0219      MINN2=NT
ISM 0219      GO TO 480

ISM 0220      C 470 M2 LESS THAN MT
ISM 0221      IG02 = 2
ISM 0222      N2VNT=1
ISM 0223      NTVN2=NT/N2
ISM 0224      MINN2=N2
ISM 0225      JJD2=NTS0/M2
ISM 0226      M1M=M1-MT

ISM 0227      C 560 M1 GR. OR EQ. MT
ISM 0228      IG01=1
ISM 0229      N1VNT=N1/NT
ISM 0230      MINN1=NT
ISM 0230      GO TO 580

ISM 0231      C 570 M1 LESS THAN MT
ISM 0232      IG01=2
ISM 0233      N1VNT=1
ISM 0234      NTVN1=NT/N1
ISM 0235      MINN1=NT
ISM 0236      JJD1=NTS0/M1
ISM 0237      JJ3=1
ISM 0238      DD 890 JPP3=1,N3VNT
ISM 0239      (PP3=INV(JJ3),MINN3
ISM 0240      DD 870 JPP3=1,MINN3
ISM 0241      GO TO (610,620),IG03
ISM 0242      IP3=INV(JP3),N3VNT
ISM 0243      GO TO 630
ISM 0244      IP3=INV(JP3),NTVN3
ISM 0245      JPP3=1,PP3=IP3,M2
ISM 0246      JJ2=1
ISM 0247      DD 870 JPP2=1,N2VNT
ISM 0248      (PP2=INV(JJ2),I3
ISM 0249      DD 850 JPP2=1,MINN2

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15N 0250          GO T3 (710,720),IGD2
15N 0251          IP2=INV(JP2)*N2VNT
15N 0252          GO TO 730
15N 0253          IP2=INV(JP2)/N1VNT
15N 0254          GO TO 730
15N 0255          IP2=INV(JP2)*N1
15N 0256          GO TO 730
15N 0257          J1=1
15N 0258          DO 850 J2=1,MINI
15N 0259          GO TO (810,820),IGD1
15N 0260          IP1=INV(JP1)*N1VNT
15N 0261          GO TO 830
15N 0262          IP1=INV(JP1)/N1VNT
15N 0263          GO TO (810,820),IGD1
15N 0264          I2=IPI*IPI+1
15N 0265          IF (J2-1) 840,850,850
15N 0266          T=A(I)
15N 0267          A(I)=T
15N 0268          T=A(I+1)
15N 0269          A(I+1)=A(J+1)
15N 0270          A(J+1)=T
15N 0271          J2=J2+1
15N 0272          END OF J2 AND JP1 AND JP2
C
C
C 870 J2=J2+1
C 880 END OF J2 AND JP2 AND JP3 LOOPS
C
C 880 J3 = J3+1
C 890 END OF J3 LOOP
C
C 890 IF (IFSET) 891, 895, 895
C 891 DO 892 I = 1, NX
C 892 A(2*I) = -A(2*I)
C 895 RETURN
C
C THE FOLLOWING PROGRAM COMPUTES THE SIN AND INV TABLES.
C
C 900 MT=MAXO(M(1),M(2),M(3)) -2
C 901 MT = MAXO(2,MT)
C 902 IF (MT-19) 906,906,13
C 903 IFEQ=0
C 904 NT=2*MT
C 905 NTV2=MT/2
C
C SET UP SIN TABLE
C 910 THETA=PI/2*(L+1) FOR L=1
C 911 THETA=.7853981634
C
C JSTEP=2*(NT-L+1) FOR L=1
C JSTEP=NT
C
C JDIF=2*(NT-L) FOR L=1
C JDIF=NTV2

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HARM4310  
HARM4320  
HARM4330  
HARM4340  
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HARM4370  
HARM4380  
HARM4390  
HARM4400  
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HARM4690  
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HARM4860  
HARM4870  
HARM4880

LEVEL 2.3. ( JUNE 78) HARM OS/360 JSTRAN H EXTENDED DATE 80.199 .57.04

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15N 0286 S(JDIF)=SIN(THETA)
15N 0287 DO 950 L=2,MT
15N 0288 THETA=THETA/2.
15N 0289 JSTEP2=JSTEP.
15N 0290 JSTEP=JDIF
15N 0291 JDIF=JSTEP/2
15N 0292 S(JDIF)=SIN(THETA)
15N 0293 JCI=NT-JDIF
15N 0294 S(JCI)=COS(THETA)
15N 0295 JLAST=VT-JSTEP2
15N 0296 IF(JLAST - JSTEP) 950,920,920
15N 0297 DO 940 J=JSTEP,JLAST,JSTEP
15N 0298 JC=NT-J
15N 0299 JD=J-JDIF
15N 0300 S(JD)=S(J)*S(JCI)+S(JDIF)*S(JC)
15N 0301 950 CONTINUE
15N 0302 C
15N 0303 C
15N 0304 SET UP INV(J) TABLE
15N 0305 C
15N 0306 MLEXP=2*(MT-L). FOR L=1
15N 0307 LMLEXP=1
15N 0308 C
15N 0309 LMLEXP=2*(L-1). FOR L=1
15N 0310 INV(L)=0
15N 0311 DO 970 L=1,MT
15N 0312 INV(L-MLEXP+1) = MLEXP
15N 0313 DO 970 J=2,L,VMLEXP
15N 0314 J=J+VMLEXP
15N 0315 INV(J)=INV(J)+MLEXP
15N 0316 MLEXP=VMLEXP/2
15N 0317 980 LMLEXP=LMLEXP*2
15N 0318 982 IF(IFSET)12,895,12
15N 0319 END

```

\*OPTIONS IN EFFECT\*NAME(MAIN) OPTIMIZE(2) LINECOUNT(55) SIZE(MAX) AUTOOBL(NONE)  
 \*OPTIONS IN EFFECT\*SOURCE EBCDIC NOLIST NODECK OBJECT YUNAP NJFORMAT GOSTMT NORREF ALC NDANSF NOTERN IOM FL/  
 \*STATISTICS\* SOURCE STATEMENTS = 314. PROGRAM SIZE = 5206. SUBPROGRAM NAME = HARM  
 \*STATISTICS\* NO DIAGNOSTICS GENERATED  
 \*\*\*\*\* END OF COMPILATION \*\*\*\*\*

56K BYTES OF CORE NOT USED

REQUESTED OPTIONS:

OPTIONS IN EFFECT: NAME(MAIN) OPTIMIZE(2) LINECOUNT(50) SIZE(MAX) AUTOOBL(NONE) SOURCE EBCDIC NDLIST NOCHECK OBJECT VJMAP NOFORMAT GOSTNT NOXRF= ALC NOANSF NOTERM (BM FL)

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    ISM 0002      SUBROUTINE GREYSC(IFILE1,NLEVEL,AMAT,IDI4,JDIM,IMIN,IMAX,IOEL,
                  JMIN,JMAX,JDEL,AMIN,AMX,TITLE,NCHAR)
    C
    ISM 0003      LEVEL=2,AMAT
    ISM 0004      DIMENSION ICHARS(10,3),LINE(1,32,3),LEVEL(10),AMAT(IDIM,JDIM)
    ISM 0005      DIMENSION TITLE(20)
    ISM 0006      DOUBLE PRECISION TITLE
    ISM 0007      DATA ICHARS/1H *,1H *,1H *,1H *,1H *,1H *,1H *,1H *,1H *,1H *,1H *,1H *,1H *,1H *,1H *,
    ISM 0008      DATA IBORDR/1H *,1H *,2H *,9H *,1H *,1H *,
    ISM 0009      DATA LEVEL/6H *,3H *,2H *,3H *,
    ISM 0010      DATA NCP/10,
                  DATA NLMAX/59,
                  IFILEX=-IFILE1
    C
    C      SCALE, IF NECESSARY
    C
    ISM 0011      AMX=AMAX
    ISM 0012      IFILE=I43S(IFILEX)
    ISM 0013      AMN=AMIN
    ISM 0014      IF(AMX-GT-AMIN) GO TO 100
    ISM 0016      AMX=-1,E20
    ISM 0017      AMN=-1,E20
    ISM 0018      DD 50 H=IMIN,IMAX,IOEL
    ISM 0019      DD 50 J=JMIN,JMAX,JDEL
    ISM 0020      AMN=AMINI(AMN,AMAT(I,J))
    ISM 0021      AMX=AMAXI(AMX,AMAT(I,J))
    ISM 0022      CONTINUE
    C
    C      PRINT HEADER
    C
    ISM 0023      NWORDS=(NCHAR-1)/NCP+1
    ISM 0024      IF(NCP.EQ.6) WRITE(IFILE,1030) (TITLE(I),I=1,NWORDS)
    ISM 0026      IF(NCP.EQ.4) WRITE(IFILE,1041) (TITLE(I),I=1,NWORDS)
    ISM 0028      IF(NCP.EQ.10) WRITE(IFILE,1042) (TITLE(I),I=1,NWORDS)
    ISM 0030      IF(NCP.EQ.5) WRITE(IFILE,1043) (TITLE(I),I=1,NWORDS)
    ISM 0032      IF(NK.EQ.AMN) RETURN
    C
    C      PREPARE A LINE FOR PRINTING
    C
    ISM 0034      LINE(I,1)=IBORDR
    ISM 0035      JLOC=(JMAX-JMIN)/JDEL+3
    ISM 0036      LINE(JLOC+2,1)=IRORDR
    ISM 0037      LINE(I,2)=IBLANK
    ISM 0038      LINE(I,3)=IBLANK
    ISM 0039      LINE(JLOC+2,2)=IBLANK
    ISM 0040      LINE(JLOC+2,3)=IBLANK
    ISM 0041      LASTL=JLOC+2
    ISM 0042      MLASTL=LASTL-1
    ISM 0043      WRITE(IFILE,1045) (IBORDR,I=1,MLASTL-1)

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LEVEL 2.3.0 (JUNE 78) GREYSC OS/360 TRAN H EXTENDED DATE 80.199, .57.13

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15N 0044 DO 400 I=IMIN,IMAX,1DEL
15N 0045 DO 250 J=1,3
15N 0046 DO 250 J=2,NLASTL
15N 0047 LINE(J,1)=IBLANK
15N 0048 LEV=1
15N 0049 J=2
15N 0050 DO 300 JJ=JMIN,JMAX,1DEL
15N 0051 J=JJ+1
15N 0052 L=(AMAT(I,JJ)-ANN)/(AMX-ANN)*FLJAT(NLEVEL)+1
15N 0053 L=MAX(1,L)
15N 0054 L=MIN(0,NLEVEL,L)
15N 0055 LEV=MAX(0,LEV,LEVEL(L))
15N 0056 LEVNOW=LEVEL(L)
15N 0057 DO 300 K=1,LEVNOW
15N 0058 LINE(J,K)=ICHARS(L,K)
15N 0059 CONTINUE
300 CONTINUE
C
C FIND LAST PRINT POSITION
C
15N 0060 DO 400 KL=1,LEV
15N 0061 DO 350 K=1,LAJLI
15N 0062 KK=LASTL-K+1
15N 0063 IF(LINE(KK,KL).NE.IBLANK) G3 TJ 375
15N 0064 CONTINUE
15N 0065 CONTINUE
15N 0066 IF(I=KK)
15N 0067 IF(I=KK) WRITE(IFILE,1050) (LINE(I,KL),
15N 0069 IF(I=KK)
15N 0070 IF(I=KK) WRITE(IFILE,1050) (LINE(I,KL),
15N 0071 IF(I=KK)
15N 0072 IF(I=KK) WRITE(IFILE,1060) (LINE(I,KL),
15N 0073 IF(I=KK)
15N 0074 IF(I=KK) WRITE(IFILE,1060) (LINE(I,KL),
15N 0075 CONTINUE
400 CONTINUE
C
C WRITE LAST LINE (BORDER)
C
15N 0076 WRITE(IFILE,1070) (IBORDR,I=1,LAJLI)
C
C PRINT TRAILING SCALE INFORMATION
C
15N 0077 IF(NCPW.EQ.6) WRITE(IFILE,1030) (TITLE(I),I=1,NWORDS)
15N 0078 IF(NCPW.EQ.4) WRITE(IFILE,1001) (TITLE(I),I=1,NWORDS)
15N 0079 IF(NCPW.EQ.10) WRITE(IFILE,1002) (TITLE(I),I=1,NWORDS)
15N 0080 IF(NCPW.EQ.5) WRITE(IFILE,1033) (TITLE(I),I=1,NWORDS)
15N 0081 WRITE(IFILE,1010)
15N 0082 DELTA=(AMX-ANN)/FLOAT(NLEVEL)
15N 0083 DO 200 I=1,NLEVEL
15N 0084 XMIN=FLJAT(I-1)/FLOAT(NLEVEL)*(AMX-ANN)+ANN
15N 0085 XMAX=XMIN+DELTA
15N 0086 LF=LEVEL(I)
15N 0087 DO 175 J=1,LEV
15N 0088 IF(I=J)
15N 0089 IF(I=J)
15N 0090 IF(I=J)
15N 0091 IF(I=J)
15N 0092 IF(I=J)

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LEVEL 2.3. ( JUNE 78)      JREYSC      05/360      JAFRAN M EXTENDED      DATE 80.199  1.57.13

ISN 0094      IF(IFI,EX,LT,0,AND,J,EQ,1) WRITE(ILE,1020) ICHARS(I,J),KMIN,
ISN 0096      I, KMAX
ISN 0098      IF(ILEX,GT,0,AND,J,NE,LEV) WRITE(IFILE,1030) ICHARS(I,J)
ISN 0100      IF(ILEX,LT,0,AND,J,NE,1) WRITE(ILE,1030) ICHARS(I,J)
ISN 0101      CONTINUE
ISN 0102      CONTINUE
ISN 0103      RETURN
ISN 0104      FORMAT(I10,20A5)
ISN 0105      FORMAT(I10,20A4)
ISN 0106      FORMAT(I10,12A10)
ISN 0107      FORMAT(I10,2A5)
ISN 0108      FORMAT(I10,32HGREY-SCALE CHARACTERS AND RANGES/IX)
ISN 0109      FORMAT(I10,5X,E15.6,5X,E15.6)
ISN 0110      FORMAT(I10,4X,A1)
ISN 0111      FORMAT(I11)
ISN 0112      FORMAT(I11/IX,20A6)
ISN 0113      FORMAT(I11/IX,20A4)
ISN 0114      FORMAT(I11/IX,12A10)
ISN 0115      FORMAT(I10,2A5)
ISN 0116      FORMAT(I10,132A1)
ISN 0117      FORMAT(I10,132A1)
ISN 0118      FORMAT(I10,132A1)
ISN 0119      END

*OPTIONS IN EFFECT*NAME(MAIN) OPTIMIZE(2) LINECOUNT(55) SIZE(MAX) AUTODBLINDONE)
*OPTIONS IN EFFECT*SOURCE EBCDIC NULIST MODECK OBJECT NUNAP NOFORMAT GOSTMT NOXREF ALC NOANSF NOTERM IRM FIA
*STATISTICS* SOURCE STATEMENTS = 118, PROGRAM SIZE = 4840, SUBPROGRAM NAME =GREYSC
*STATISTICS* NO DIAGNOSTICS GENERATED
***** END OF COMPILATION *****

80K BYTES OF CORE NOT USED

```

LEVEL 2.3. ( JUNE 78) 05/360 JATMAN H EXTENDED DATE 00.199 1.57.15

REQUESTED OPTIONS:

OPTIONS IN EFFECT:

NAME(MAIN) OPTIMIZE(2) LINECOUNT(56) SIZE(MAX) AUTOOBL(NONE)

SOURCE EBCDIC NOLIST NODECK OBJECT VJMA3 NOFORMAT GOSTMT NOXREF ALC NOANSF NOTERM IOM FLA

```

ISM 0002 SUBROUTINE PRINTER(I,STEP)
ISM 0003 COMMON /CM2/REFNDX(16384),SV(129),CS(128),ZTSZ(128),POSQ(128)
C      *AMPARY(16384),RADARY(16384)
ISM 0004 LEVEL 2,REFNDX,SN,CS,ZTSQ,PJSQ,AMPARY,RADARY
ISM 0005 COMMON /ARRAYS/V(32768)
ISM 0006 *NOZIVC,RESHSO,MSHSO2,PI,WAVENM,XXSI,MS,NS,MF,NF,MSHPTS
C      *PLTFLD,PLTFLE
C      DIMENSION X(130),KWAST(100),YWAIST(100),VRAD(100),WSTID(2),
C      *MAXI(2),MAXY(2),MSYX(2),MSYK(2),MSYI(2),XT(100)
C      DIMENSION FIELDS(130),XID(2),FLDID(2),FLOTTL(2)
C      EQUIVALENCE (MAXX,MSYX)
C      LOGICAL PGREY,PWAIST,PLTST,PLTMAX,PLTFLD,PLTFLE,LAST
REAL -AMBDA
DATA 5STID/19H2ND MOMENTS VS DIST/
DATA VAXID/18H2PEAK IRRAD VS DIST/
DATA MAXY/18HAXIAL DIST/
DATA MSY/15HPEAK IRRADIANCE/
DATA XSY/17HBEAM WAIST RATIO5/
DATA XID/18HAXIAL DIST (MICR)/
DATA FLDID/10HIRRADIANCE/
DATA FLOTTL/20HIRRAD VS RADIAL DIST/
CSTEP=1STEP+1
IF (PGREY .OR. PWAIST) WRITE(10JT,230J) I,STEP
IF (PGREY .OR. PWAIST) PLTMAX,OM,PLTST) CALL PEAK(VMAX)
IF (.NOT. PWAIST .OR. PLTST) GO TO 10
CALL SIZE(RADARY,MESH,XO,YJ,X2,Y2)
X2=1.41421356*X2*DX
Y2=1.41421356*Y2*DY
IF (PWAIST) WRITE(10UT,201D) X2,Y2
2000 FORMAT(14H THIS IS STEP ,(2//)
2010 =OMATI//,38H BEAM WAIST IN X AND Y DIRECTIONS IS ,2(F10.5,2X),
+ 7MICRONS)
10 CONTINUE
IF (PGREY) CALL GREYSC(IGREY,I0,RADARY,MESH,MESH,MS,MF,1,NS,NF,1,
+ 0,,0,,10-IRRADIANCE,10)
RETURN
END

```

\*\*\*\*\* END OF COMPILATION \*\*\*\*\*

112K BYTES OF CORE NOT USED

DATE 80.19.21.57.15

OS/360 F33TRAN H EXTENDED

LEVEL 2.3.0 (JUNE 78)

\*STATISTICS\* NO DIAGNOSTICS THIS STEP

```

IEF1421 - S..P EXECUTED - COND CODE 0000
IEF2851 SYS1.FORTRAN PASSED
IEF2851 VOL SER NOS= ACS002. PASSED
IEF2851 SYS80199.T215654.RV000.SBI404F2.LOADSET PASSED
IEF2851 VOL SER NOS= SYSDA1. DELETED
IEF2851 SYS80199.T215654.RV000.SBI404F2.R0000001 DELETED
IEF2851 VOL SER NOS= SYSDA3. DELETED
IEF2851 SYS80199.T215654.RV000.SBI404F2.R0000002 DELETED
IEF2851 VOL SER NOS= SYSDA2. SYSJUT
IEF2851 SYS80199.T215654.SV000.SBI404F2.R0000003 SYSIN
IEF2851 VOL SER NOS= SYSDA3. DELETED
IEF2851 SYS80199.T215654.RV000.SBI404F2.S0000004
IEF2851 VOL SER NOS= SYSDA2. DELETED
IEF2851 SYS80199.T215654.RV000.SBI404F2.S0000004
IEF2851 VOL SER NOS= SYSDA2.

```

CCN3011 SYSRINT PRINTED 24 TOTAL PAGES. REQUIRED 10 TRACKS.

REGION ALLOCATION CORE USED  
250K

STEP	CPU TIME	JOB CPU TIME	STEP I/O CJVT	JOB I/O COUNT
	3.14S	3.14S	306	306
EXEC	PGM=6LOADER.COND=(S,LT,FORT).TIME=6TG.REGION=6RG			00000320
IEF6531	SUBSTITUTION JCL - PGM=LLADER.COND=(S,LT,FORT).TIME=1439.REGION=450K			00000340
XXFT05F001	DD DDNAME=SYSIN			00000360
XXSTEPL1B	DD DISP=(SHR,PASS),DSN=SYS&LEVEL.*FORTRAN			00000390
IEF6531	SUBSTITUTION JCL - DISP=(SHR,PASS),DSN=SYSI.FORTRAN			00000400
XXSYSL1B	DD DISP=(SHR,PASS),DSN=SYS&LEVEL.*FORT-1B			00000420
IEF6531	SUBSTITUTION JCL - DISP=(SHR,PASS),DSN=SYSI.FORT-1B			00000440
XXSYSLIN	DD DSN=66LOADSET,DISP=(OLD,DELETE)			00000460
XXSYSROUT	DD SYSOUT=A,SPACE=(133,(200,75),RLSE)			
XXFT06F001	DD SYSOUT=A,UNIT=SYSOUT,SPACE=(TRK,6PG,RLSE).			
IEF6531	SUBSTITUTION JCL - SYSOUT=A,UNIT=SYSOUT,SPACE=(TRK,25,RLSE).			
XX	DCB=(RECFM=FBA,-RECL=133,BLKSIZE=4123)			
XX	GO.SYSIN DD *			

```

//
IEF2361 ALL OC. FOR SBI404F2 GO
IEF2371 153 ALLOCATED TO FT05F001
IEF2371 822 ALLOCATED TO STEPL1B
IEF2371 822 ALLOCATED TO SYS-1B
IEF2371 153 ALLOCATED TO SYSLIN
IEF2371 713 ALLOCATED TO SYSROUT
IEF2371 265 ALLOCATED TO FT05F001

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300 LOADER

OPTIONS USED - PRINT,MAP,NOLET,CALL,RES,NOTERM,SIZE=0322,NAME=00GD

NAME	TYPE	ADDR	NAME	TYPE	ADDR	SIZE	NAME	TYPE	ADDR	NAME	TYPE	ADDR	NAME	TYPE	ADDR	NAME	TYPE	ADDR
MAIN	SD	116010	PEAK	SD	1174D8		JME7JC71*	SD	1176F8	ARRAYS	SD	117410	PARAM	SD	117410			
PRINTP	SD	137A68	MAXM	SD	137A90		18081971*	SD	138E58	PRINTER	SD	13A1D0	IMDEFKPI*	LR	13A604			
FIXPFI*	LR	13A488	IMJECQMH*	SD	13A5D8		AP31971*	LR	13A604	ICUM*	LR	13A604	F310CS#*	SD	13A9A0			
IMDSMTC-1*	LR	13B3C0	FIQA#*	SD	13B488		PCVJUTP*	LR	13B522	IMDCM42*	SD	13B9A0	SC00ASD*	LR	13C682			
IMDFCVTH*	SD	13C5A8	ADCON#*	LR	13C5A8		FCVJUTP*	LR	13C5E2	FCVLOUTP*	LR	13C682	FCV7OUTP*	LR	13C682			
FCV7OUTP*	LR	13C8BA	FCVEDJTP*	LR	13C8AA		IMDEFIUS2*	LR	13C8A8	INT65C-1*	LR	13C682	IMDEFVTH*	LR	13C682			
ARITH#*	LR	13C790	ADJSTFC1*	LR	13D524		IMDEFIUS2*	SD	13C8A8	FILE2#*	LR	13D790	FILE2#*	LR	13D790			
IMOFIUS2*	SD	13F730	IMDFCON1*	SD	13E5F8		IMDEFIUS2*	LR	13E778	IMDEFIUS2*	SD	13D790	FILE2#*	LR	13D790			
IMOERRN*	SD	13F730	ERRN#*	LR	13E5F8		IMDEFIUS2*	LR	13E778	IMDEFIUS2*	SD	13D790	FILE2#*	LR	13D790			
IMONTRCM*	SD	13F730	IMDFCON1*	SD	13E5F8		IMDEFIUS2*	LR	13E778	IMDEFIUS2*	SD	13D790	FILE2#*	LR	13D790			
IMONAMEL*	SD	140C10	IMDFCON1*	SD	13E5F8		IMDEFIUS2*	LR	13E778	IMDEFIUS2*	SD	13D790	FILE2#*	LR	13D790			
IMOSSURT*	SD	141A28	FRDY#*	LR	140C10		IMDEFIUS2*	LR	13E778	IMDEFIUS2*	SD	13D790	FILE2#*	LR	13D790			
IM8ACOS*	LR	141B90	FRDY#*	LR	140C10		IMDEFIUS2*	LR	13E778	IMDEFIUS2*	SD	13D790	FILE2#*	LR	13D790			
EXP	LR	141D98	SIN	LR	141A28		IMDEFIUS2*	LR	13E778	IMDEFIUS2*	SD	13D790	FILE2#*	LR	13D790			
			LC42	LR	141BB2		IMDEFIUS2*	LR	13E778	IMDEFIUS2*	SD	13D790	FILE2#*	LR	13D790			
				CM	141F48		IMDEFIUS2*	LR	13E778	IMDEFIUS2*	SD	13D790	FILE2#*	LR	13D790			

TOTAL LENGTH 5C738  
ENTRY ADDRESS 116010

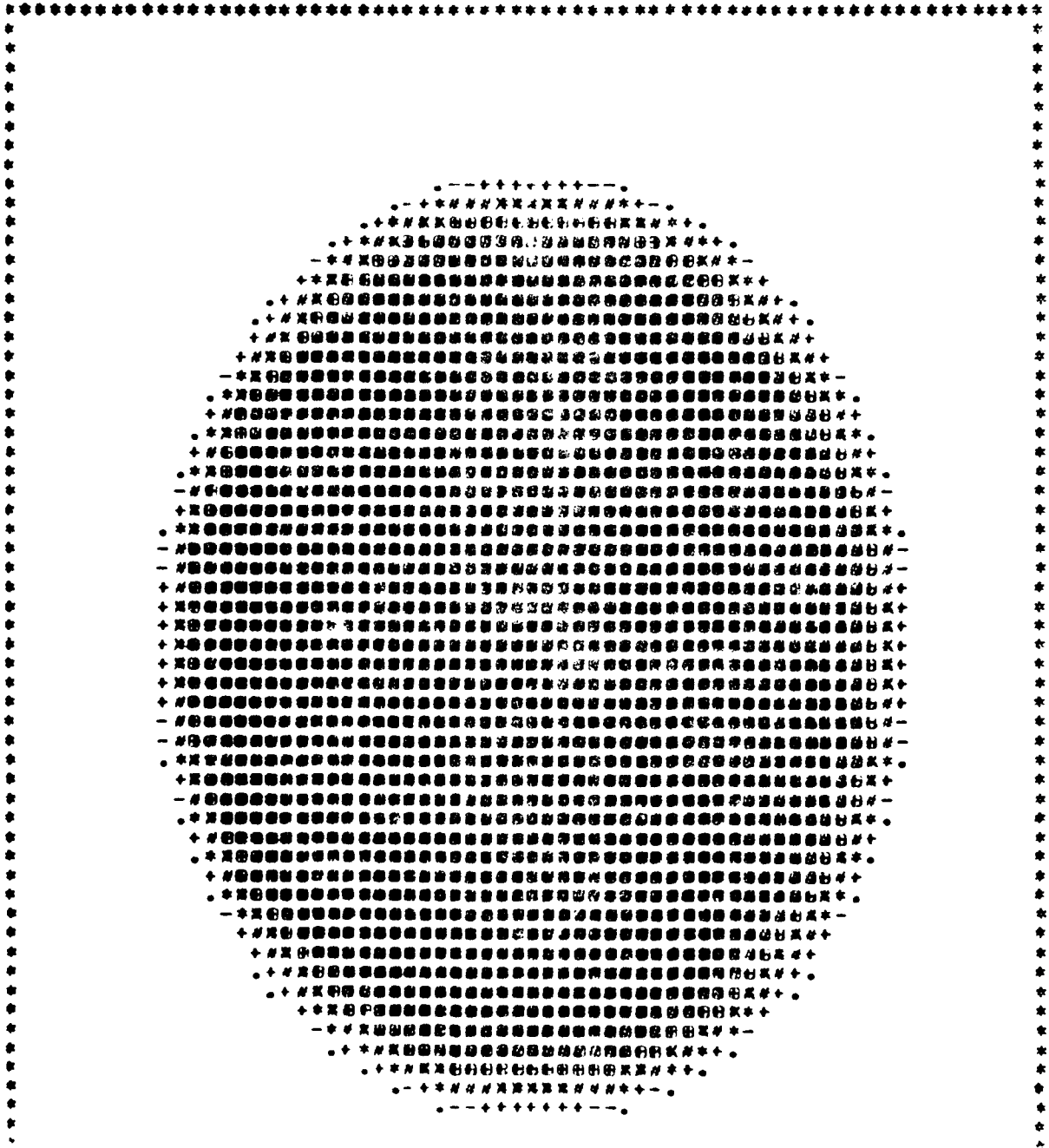
```

1 0.0 0.0
1 0.0 0.0
0.0 1.40000
25.0
CDEFAULT -0.00000012 .RO= 20.0000000 .NO= 1.47349954 .PCOR= 1.39999952 .OUTRAD= 62.0000000
LAMBDA= 40. .MDZINC= 5.10UT= 5.1.REV= 6. .PGREV=1.2 .VAISY=1. .PLWST=1. .PLTMAX=F.
N STEPS= 126 = MESH DX = 1.0
CEND
ZMIN = 234.6827 MICRONS
ZINC = 46.9365 MICRONS
ZEND = 0.448E-04 MICRONS*(-2)
ALPHA = 0.03213

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REFNOX

P-10 Study



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REFNDX

GREY-SCALE CHARACTERS AND RANGES

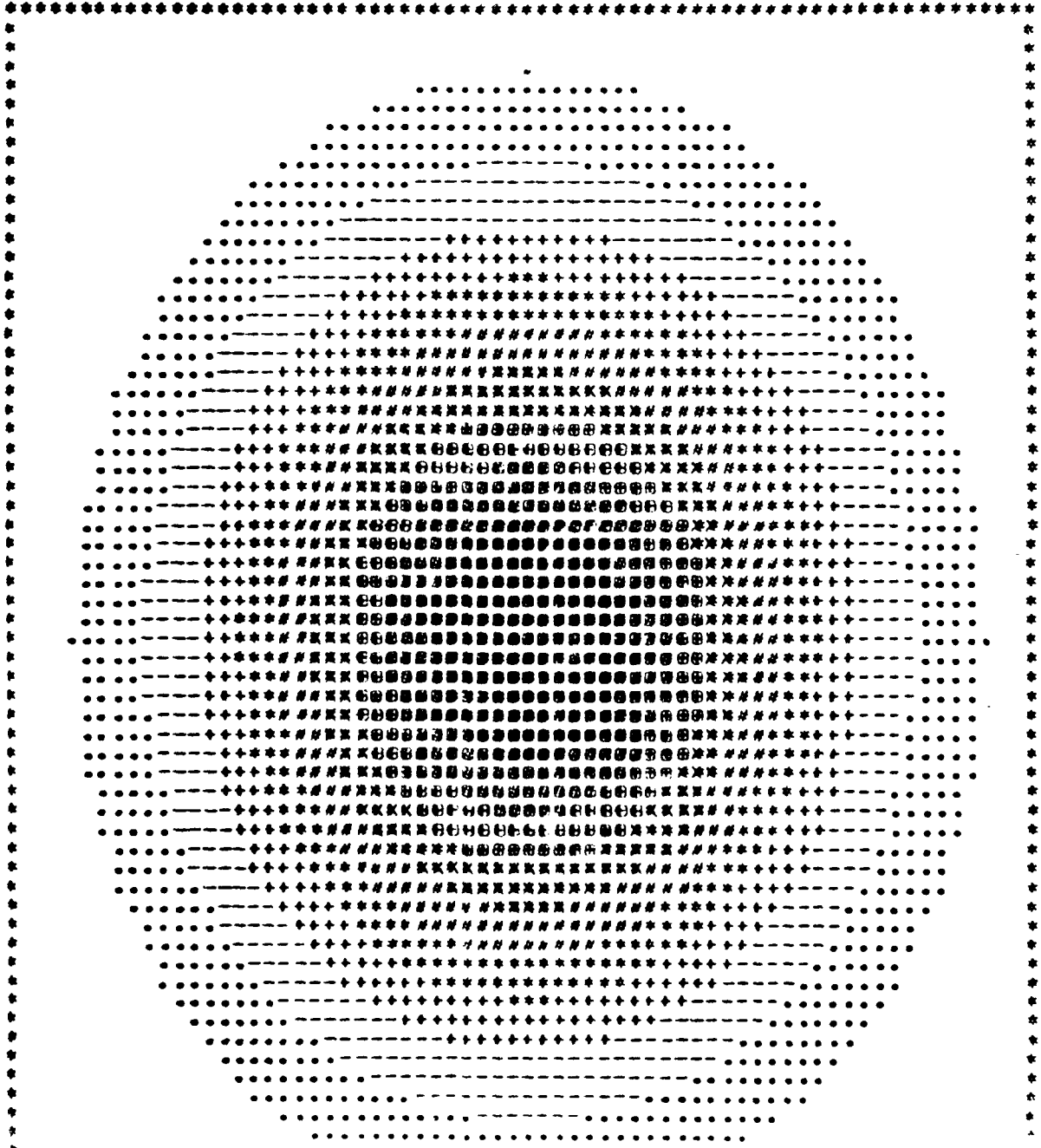
.	0.145928E+01	0.146135E+01
-	0.146135E+01	0.146342E+01
+	0.146342E+01	0.146549E+01
*	0.146550E+01	0.146757E+01
#	0.146757E+01	0.146964E+01
/	0.146964E+01	0.147171E+01
K	0.147171E+01	0.147378E+01
■	0.147378E+01	0.147585E+01
■	0.147585E+01	0.147793E+01
■	0.147793E+01	0.148000E+01

THIS IS STEP 0

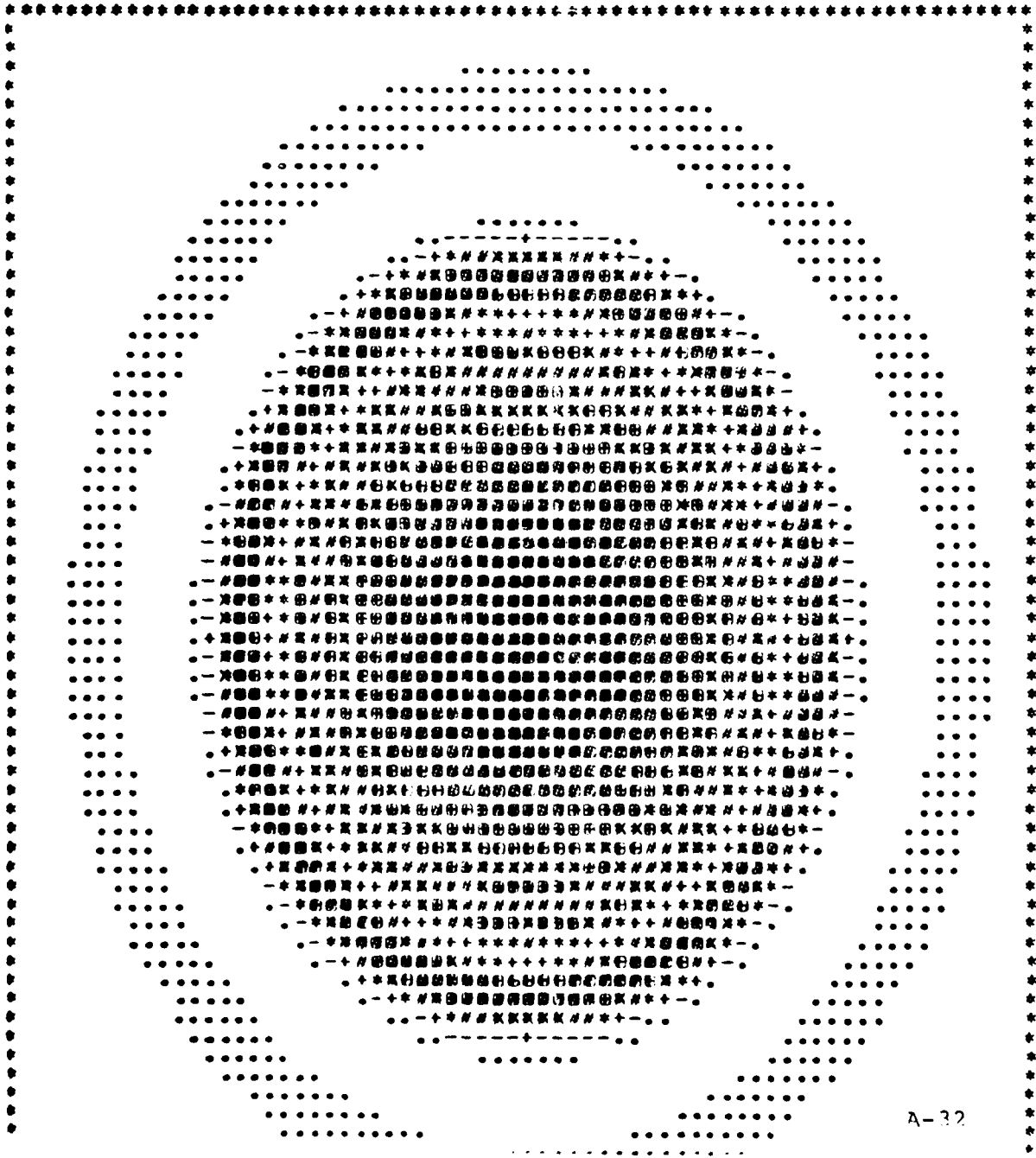
SUML = 0.4860770E+00 SU4R = 0.5139230E+00

BEAM WAIST IN X AND Y DIRECTIONS IS 19.99481 19.99631 MICRONS

IRRADIAN





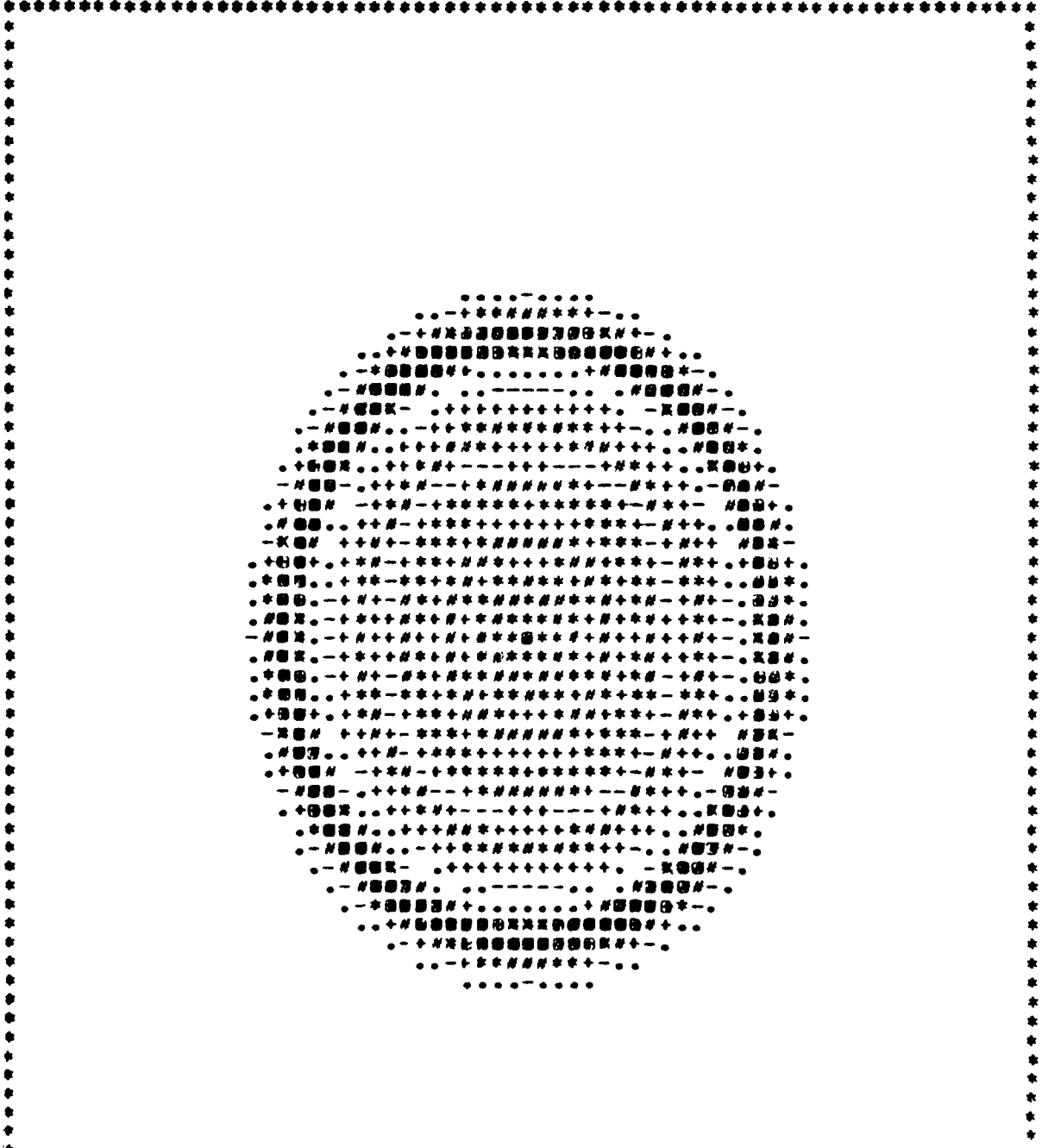






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IRRADIAN



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IRRADIAN

GREY-SCALE CHARACTERS AND RANGES

.	0.323949E-02	0.211144E+00
-	0.211144E+00	0.419048E+00
+	0.419048E+00	0.626953E+00
*	0.626953E+00	0.834857E+00
#	0.834857E+00	0.104276E+01
%	0.104276E+01	0.125066E+01
^	0.125067E+01	0.145857E+01
^	0.145857E+01	0.166647E+01
^	0.166647E+01	0.187438E+01
^	0.187438E+01	0.208228E+01

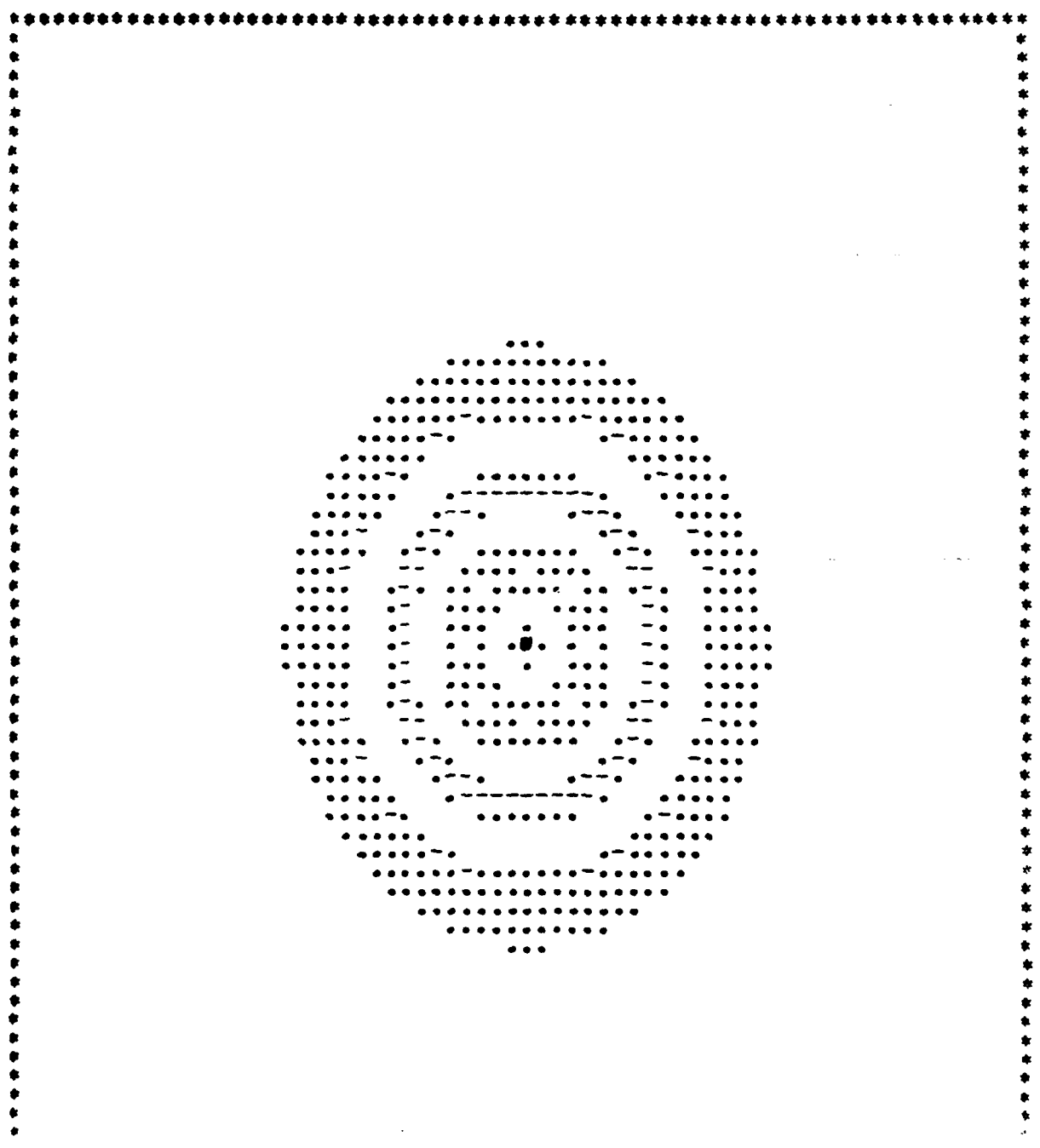
THIS IS STEP 3

SUNL = 0.4817133E+00    SUMR = 0.5182866E+00

BEAM WAIST IN X AND Y DIRECTIONS IS    18.07376    18.07498    MICRONS

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IRRADIAN



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IRRADIAN

GREY-SCALE CHARACTERS AND RANGES

	0.759699E-05	0.842067E+00
.	0.842067E+00	0.168413E+01
-	0.168412E+01	0.252618E+01
+	0.252618E+01	0.336824E+01
*	0.336824E+01	0.421030E+01
#	0.421030E+01	0.505236E+01
R	0.505236E+01	0.589442E+01
●	0.589442E+01	0.673648E+01
●	0.673648E+01	0.757854E+01
●	0.757854E+01	0.842060E+01

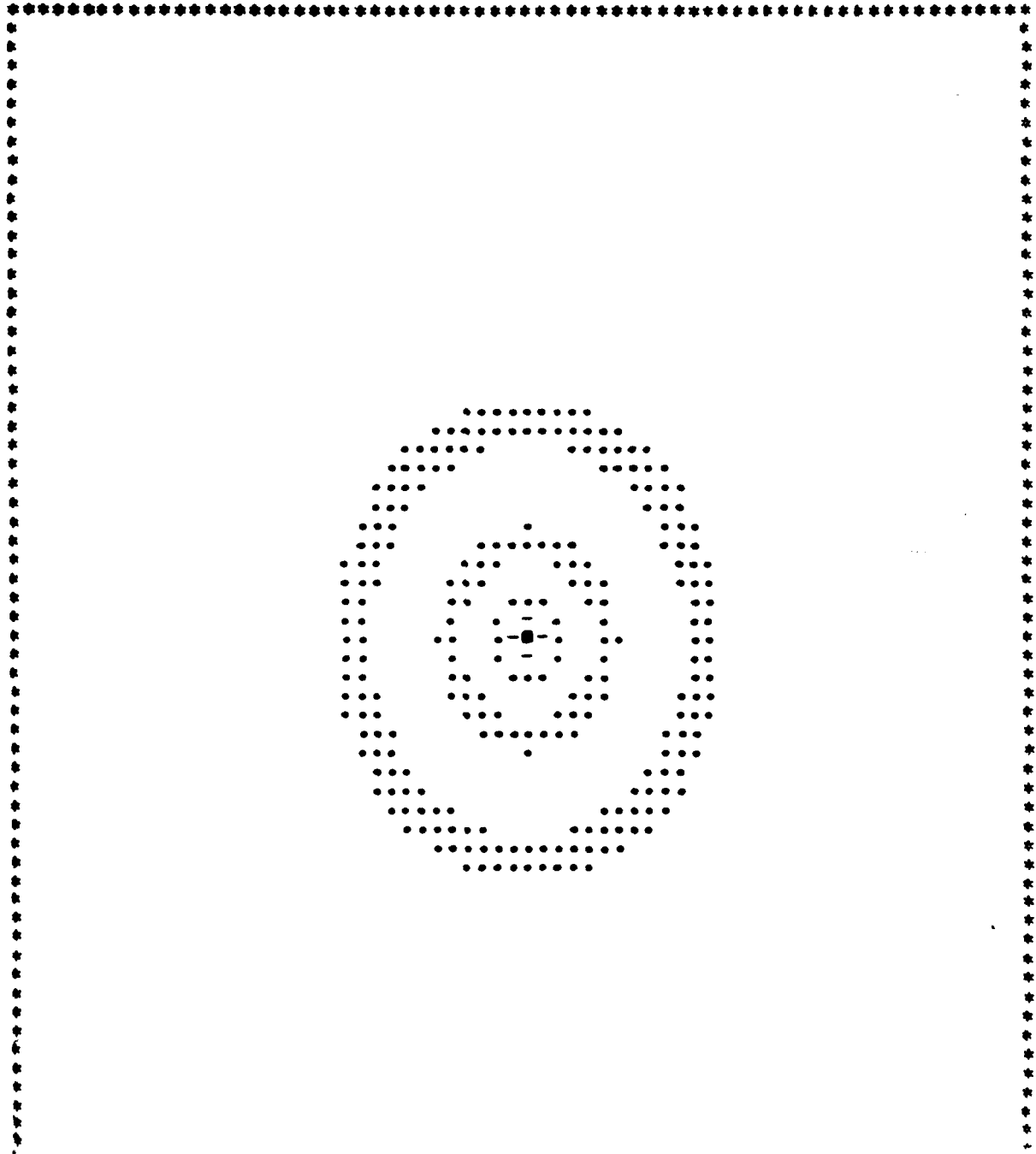
THIS IS STEP 4

SUML = 0.4754750E+00    SUMR = 0.5245249E+00

BEAM WAIST IN X AND Y DIRECTIONS IS    17.64449    17.64572    MICRONS

P-10

IRRADIAN



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IRRADIAN

GREY-SCALE CHARACTERS AND RANGES

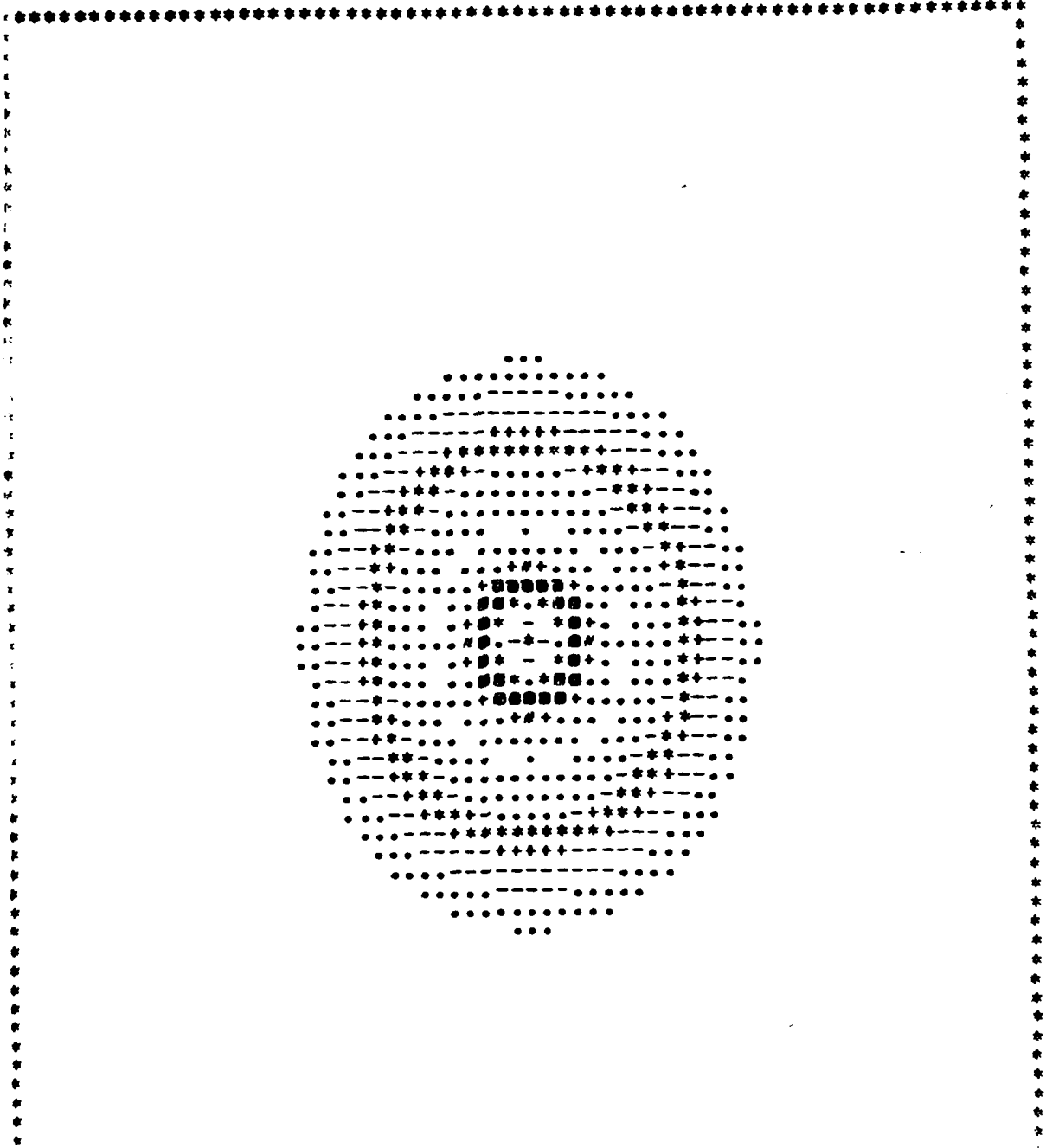
.	0.247009E-03	0.153054E+01
-	0.153054E+01	0.306083E+01
+	0.306083E+01	0.459112E+01
*	0.459112E+01	0.612142E+01
#	0.612142E+01	0.765171E+01
%	0.765171E+01	0.918200E+01
^	0.918200E+01	0.107123E+02
^	0.107123E+02	0.122426E+02
^	0.122426E+02	0.137729E+02
^	0.137729E+02	0.153032E+02

THIS IS STEP 5

SUML = 0.4803946E+00    SUMR = 0.5196053E+00

BEAM WAIST IN X AND Y DIRECTIONS IS    17.43616    17.43765 MICRONS

IRRADIAN



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IRRADIAN

GREY-SCALE CHARACTERS AND RANGES

	0.348578E-05	0.546719E+00
.	0.546719E+00	0.109343E+01
-	0.109343E+01	0.154015E+01
+	0.164015E+01	0.218686E+01
*	0.218696E+01	0.273358E+01
#	0.273358E+01	0.328029E+01
£	0.328029E+01	0.382701E+01
®	0.382701E+01	0.437372E+01
©	0.437373E+01	0.492044E+01
®	0.492044E+01	0.546716E+01

THIS IS STEP 6

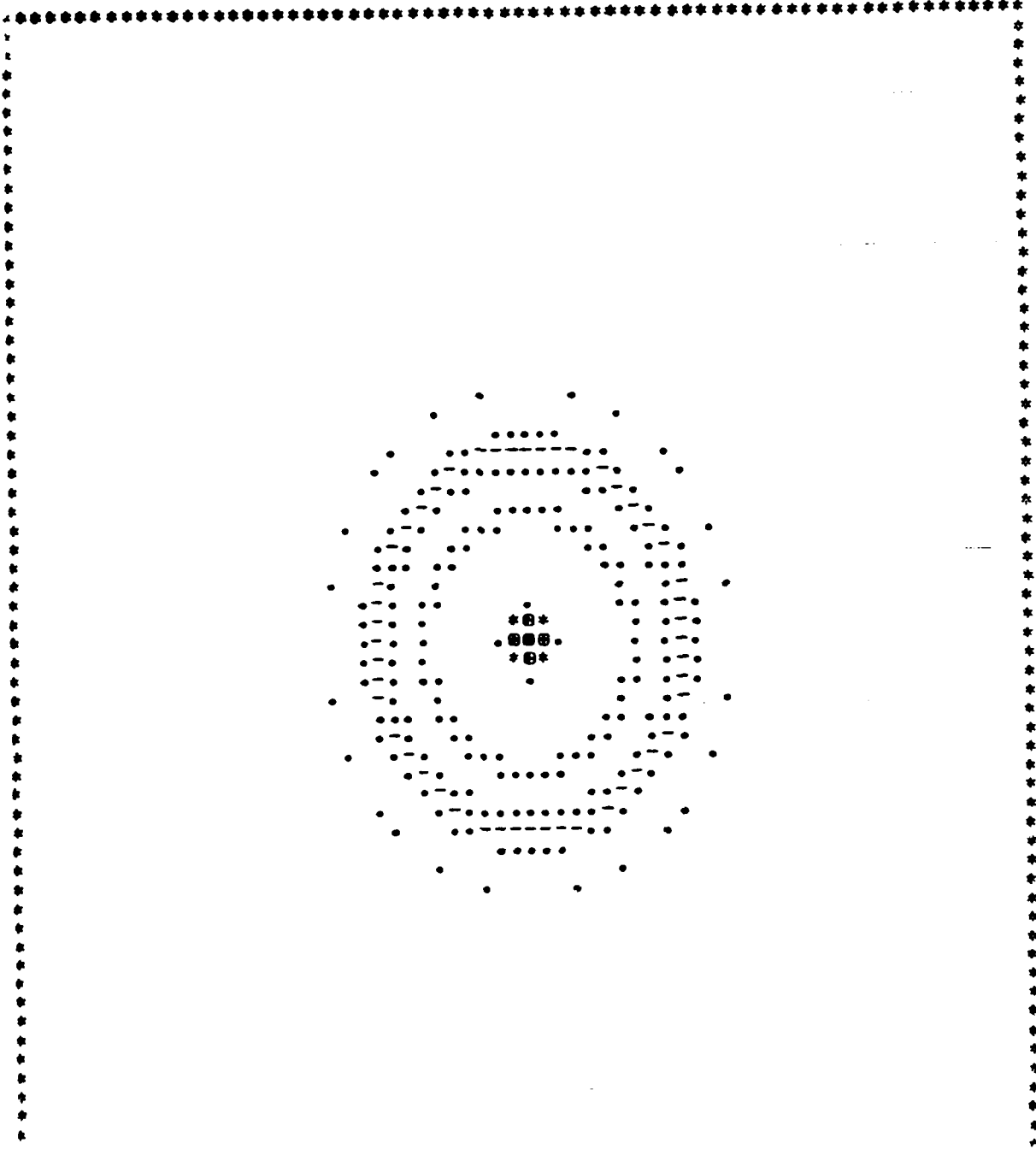
SUML = 0.4704280E+00      SUMR = 0.5295719E+00

BEAM WAIST IN X AND Y DIRECTIONS IS      17.47115      17.47259      MICRONS



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IRRADIATION



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IRRADIAN

GREY-SCALE CHARACTERS AND RANGES

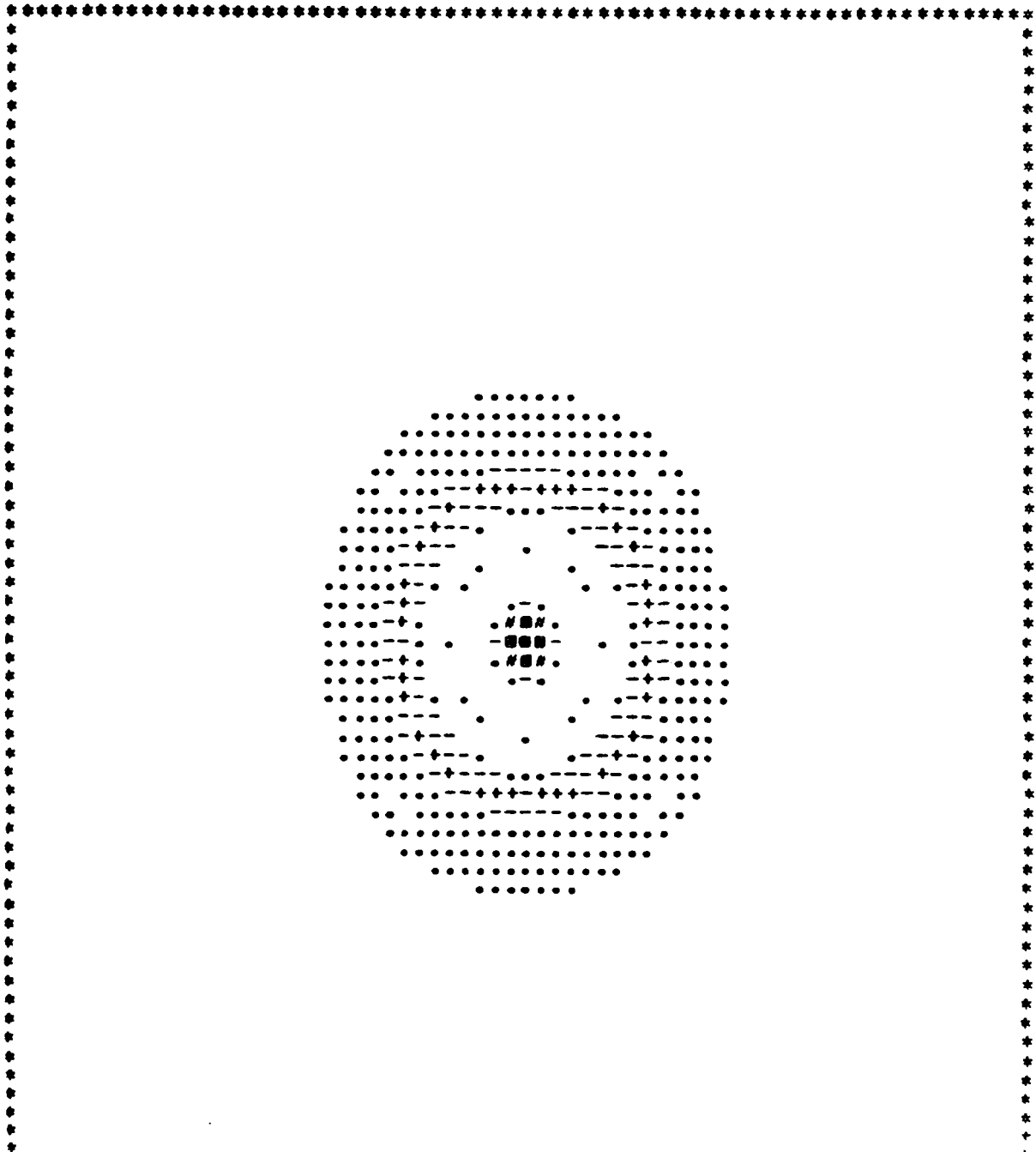
.	0.341598E-04	0.146743E+01
-	0.146743E+01	0.293483E+01
+	0.293483E+01	0.440223E+01
*	0.440223E+01	0.586963E+01
#	0.586963E+01	0.733703E+01
@	0.733703E+01	0.880443E+01
A	0.880443E+01	0.102718E+02
B	0.102718E+02	0.117392E+02
C	0.117392E+02	0.132066E+02
D	0.132066E+02	0.146740E+02

THIS IS STEP 7

SUML = 0.4736281E+00      SUMR = 0.5263718E+00

BEAM WAIST IN X AND Y DIRECTIONS IS      17.75745      17.75934      MICRONS

IRRADIAN



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IRRADIAN

GREY-SCALE CHARACTERS AND RANGES

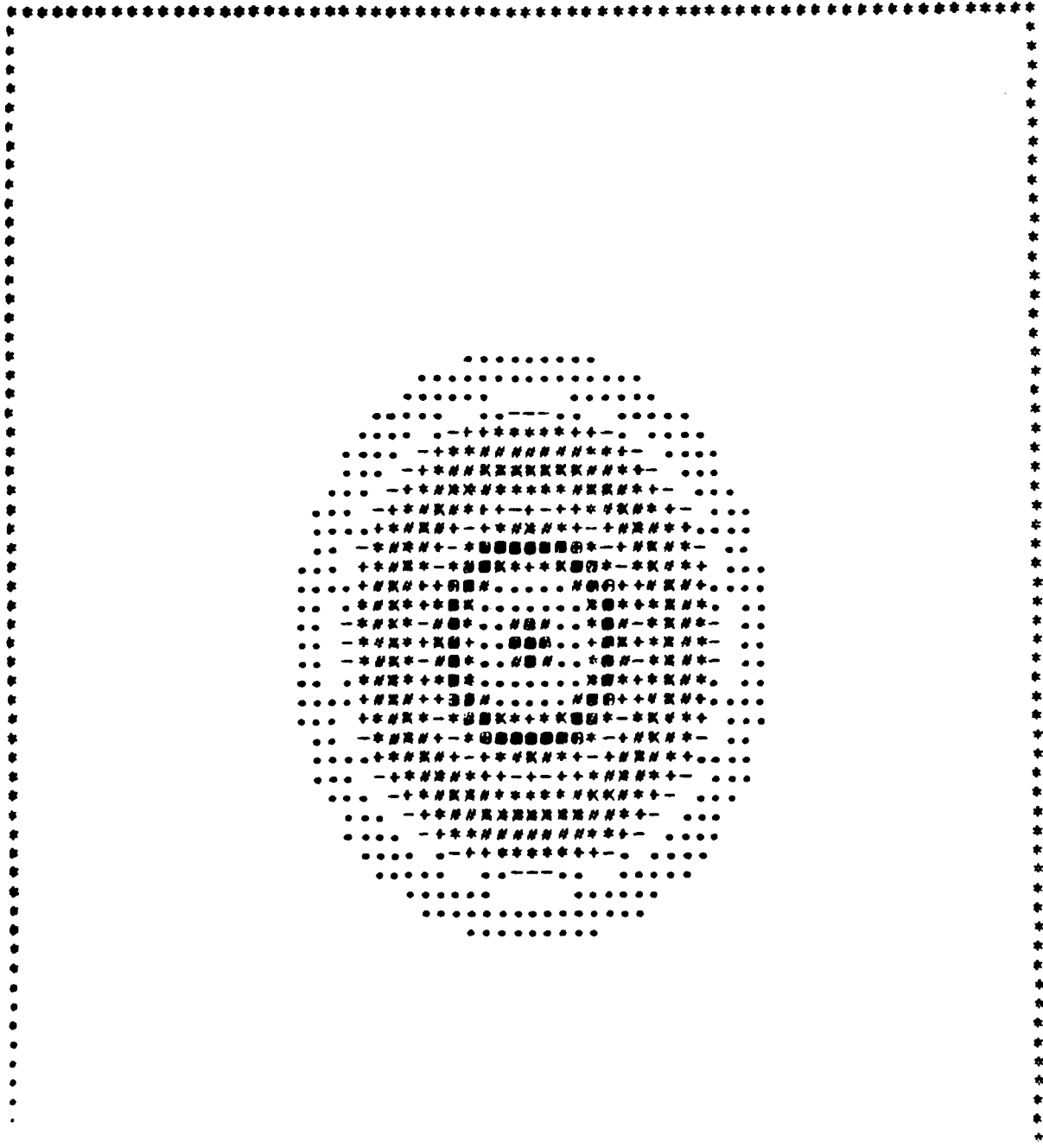
.	0.349578E-03	0.996981E+00
-	0.996981E+00	0.199361E+01
+	0.199361E+01	0.299024E+01
*	0.299024E+01	0.398687E+01
#	0.398688E+01	0.498351E+01
\$	0.498351E+01	0.598014E+01
%	0.598014E+01	0.697677E+01
@	0.697677E+01	0.797340E+01
^	0.797340E+01	0.897003E+01
^	0.897003E+01	0.996667E+01

THIS IS STEP 8

SUML = 0.4800833E+00      SUMR = 0.5199167E+00

BEAM WAIST IN X AND Y DIRECTIONS IS      18.20370      18.20607      MICRONS

IRRADIAN



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IRRADIAN

GREY-SCALE CHARACTERS AND RANGES

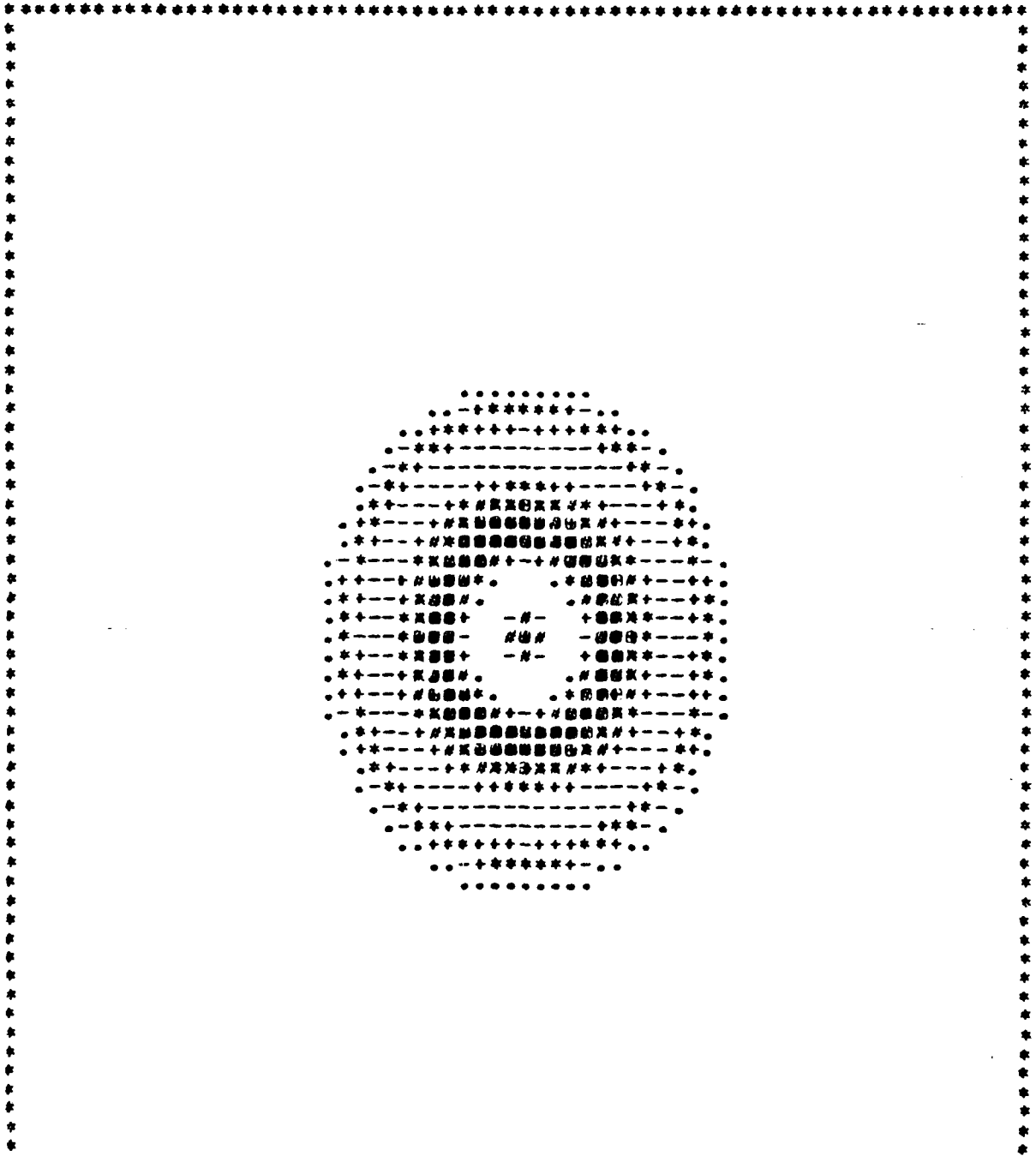
.	0.213262E-03	0.358114E+00
-	0.358113E+00	0.716014E+00
+	0.716014E+00	0.107391E+01
*	0.107391E+01	0.143181E+01
#	0.143181E+01	0.178971E+01
R	0.178971E+01	0.214761E+01
D	0.214761E+01	0.250551E+01
D	0.250551E+01	0.286341E+01
D	0.286341E+01	0.322131E+01
D	0.322132E+01	0.357922E+01

THIS IS STEP 9

SUML = 0.4812048E+00    SUMR = 0.5187951E+00

BEAM WAIST IN X AND Y DIRECTIONS IS    18.54669    19.54912    MICRONS

IRRADIAN



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IRRADIAN

GREY-SCALE CHARACTERS AND RANGES

.	0.753595E-04	0.379830E+00
-	0.379830E+00	0.759584E+00
+	0.759584E+00	0.113934E+01
*	0.113934E+01	0.151909E+01
#	0.151909E+01	0.189885E+01
@	0.189885E+01	0.227860E+01
^	0.227860E+01	0.265836E+01
^	0.265836E+01	0.303811E+01
^	0.303811E+01	0.341786E+01
^	0.341786E+01	0.379762E+01

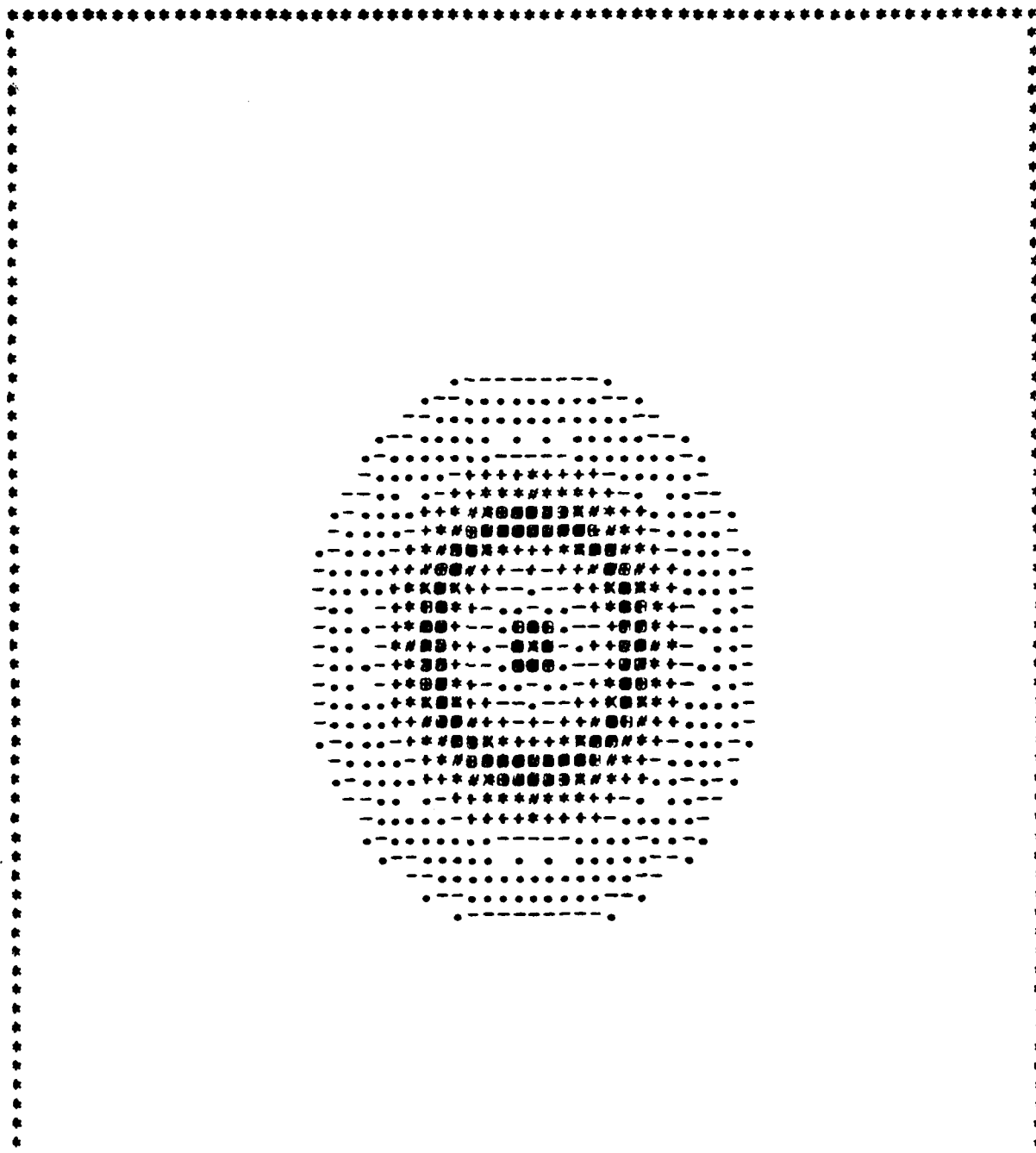
THIS IS STEP 10

SUML = 0.4789906E+00    SUMR = 0.5210093E+00

BEAM WAIST IN X AND Y DIRECTIONS IS    18.63527    18.63782    MICRONS



IRRADIAN



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IRRADIAN

GREY-SCALE CHARACTERS AND RANGES

.	0.103213E-03	0.410842E+00
-	0.410842E+00	0.821581E+00
+	0.821582E+00	0.123232E+01
*	0.123232E+01	0.164306E+01
#	0.164306E+01	0.205380E+01
@	0.205380E+01	0.246454E+01
R	0.246454E+01	0.287528E+01
Q	0.287528E+01	0.328602E+01
W	0.328602E+01	0.369676E+01
0	0.369676E+01	0.410749E+01

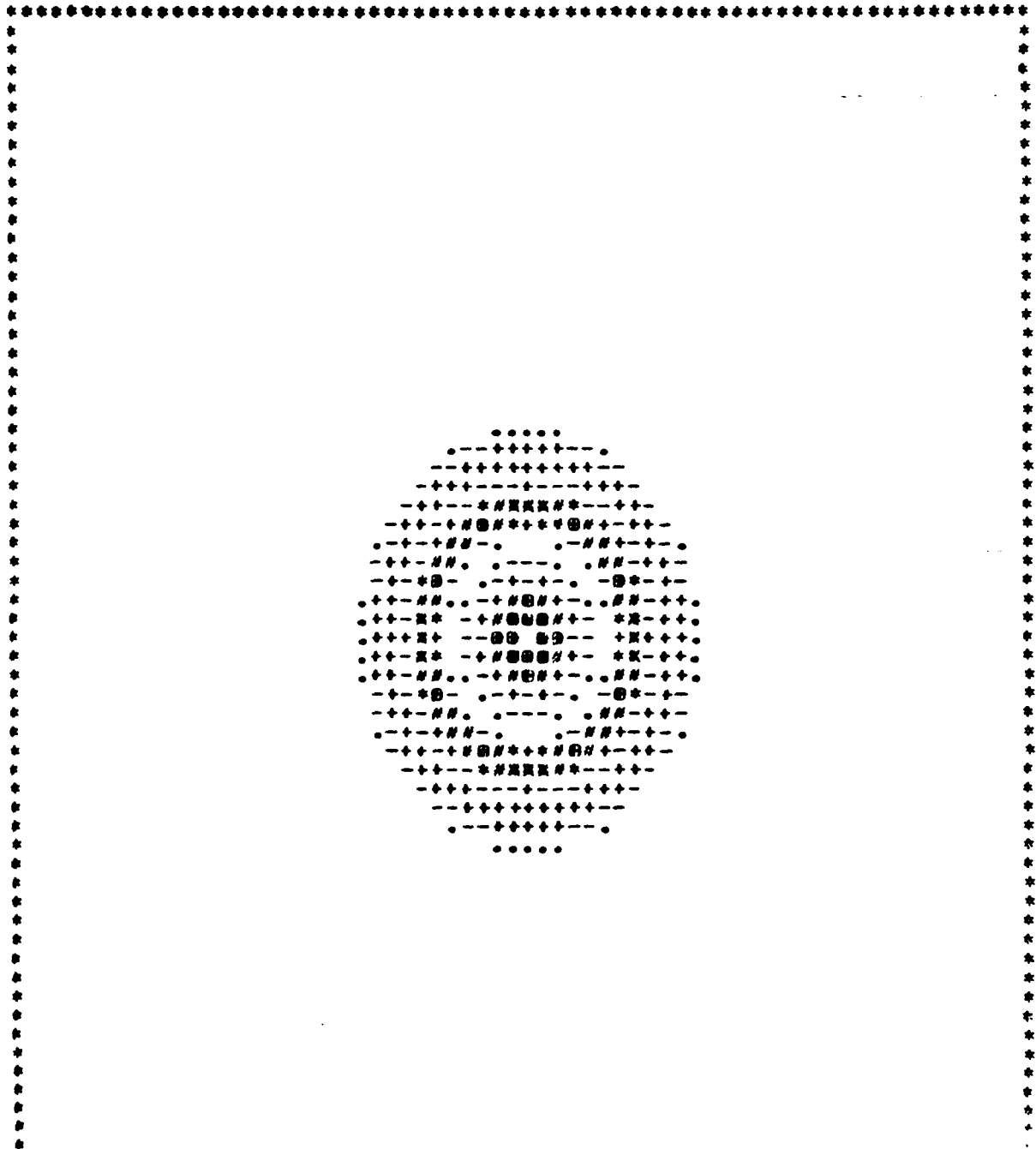
THIS IS STEP 11

SUML = 0.4775044E+00    SUMR = 0.5224955E+00

BEAM WAIST IN X AND Y DIRECTIONS IS    18.77403    18.77573    MICRONS

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IRRADIAN

GREY-SCALE CHARACTERS AND RANGES

.	0.373785E-04	0.594157E+00
-	0.594157E+00	0.118828E+01
0	0.118828E+01	0.178239E+01
+	0.178239E+01	0.237651E+01
*	0.237651E+01	0.297063E+01
#	0.297063E+01	0.356475E+01
\$	0.356475E+01	0.415887E+01
@	0.415887E+01	0.475299E+01
■	0.475299E+01	0.534711E+01
■	0.534711E+01	0.594123E+01

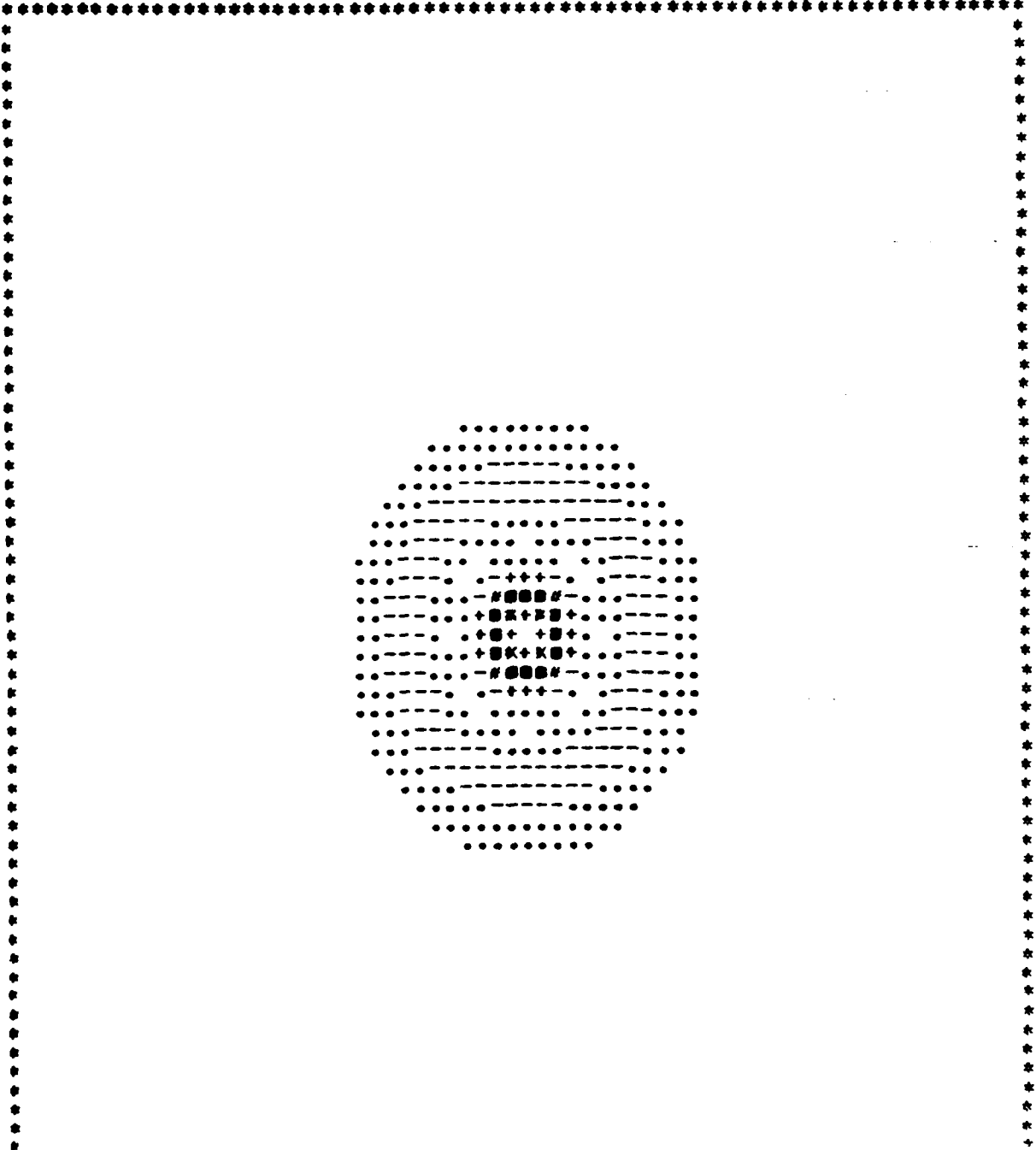
THIS IS STEP 12

SUML = 0.4750217E+00    SUMR = 0.5249783E+00

BEAM WAIST IN X AND Y DIRECTIONS IS    18.92630    18.92923    MICRONS

171000

IRRADIAN



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IRRADIAN

GREY-SCALE CHARACTERS AND RANGES

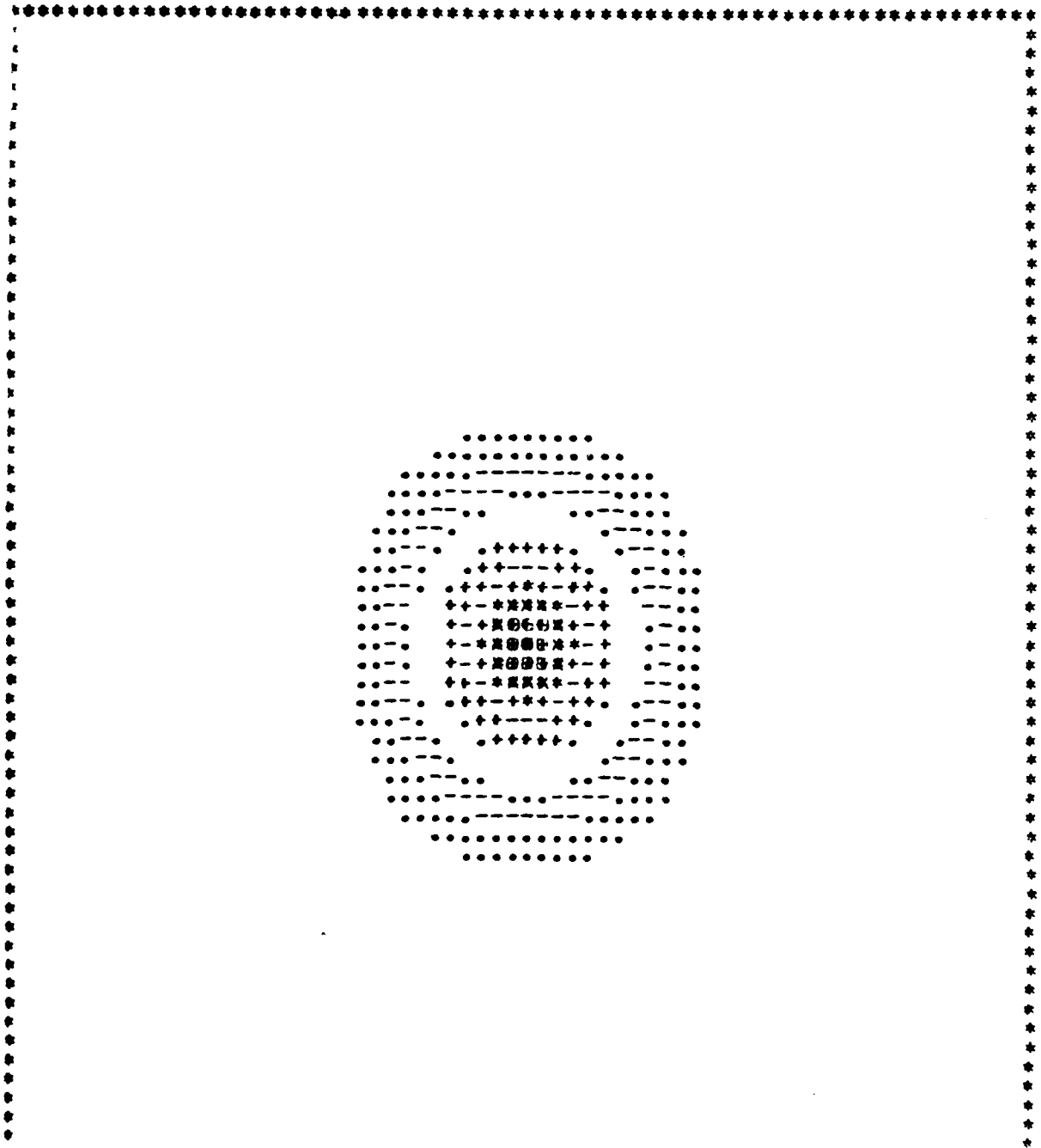
.	0.126287E-04	0.898516E+00
-	0.898516E+00	0.179702E+01
+	0.179702E+01	0.269552E+01
*	0.269552E+01	0.359403E+01
#	0.359403E+01	0.449253E+01
@	0.449253E+01	0.539103E+01
^	0.539103E+01	0.628954E+01
~	0.628954E+01	0.718804E+01
▧	0.718804E+01	0.808654E+01
▨	0.808654E+01	0.898505E+01

THIS IS STEP 13

SUML = 0.4702227E+00    SUMR = 0.5297772E+00

BEAM WAIST IN X AND Y DIRECTIONS IS    18.73090    13.73331    MICRONS

IRRADIAN



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IRRADIAN

GREY-SCALE CHARACTERS AND RANGES

.	0.703742E-04	0.848021E+00
-	0.848021E+00	0.169597E+01
+	0.169597E+01	0.254392E+01
*	0.254392E+01	0.339187E+01
#	0.339187E+01	0.423982E+01
\$	0.423982E+01	0.508778E+01
%	0.508778E+01	0.593573E+01
^	0.593573E+01	0.678368E+01
^	0.678368E+01	0.763163E+01
^	0.763163E+01	0.847958E+01

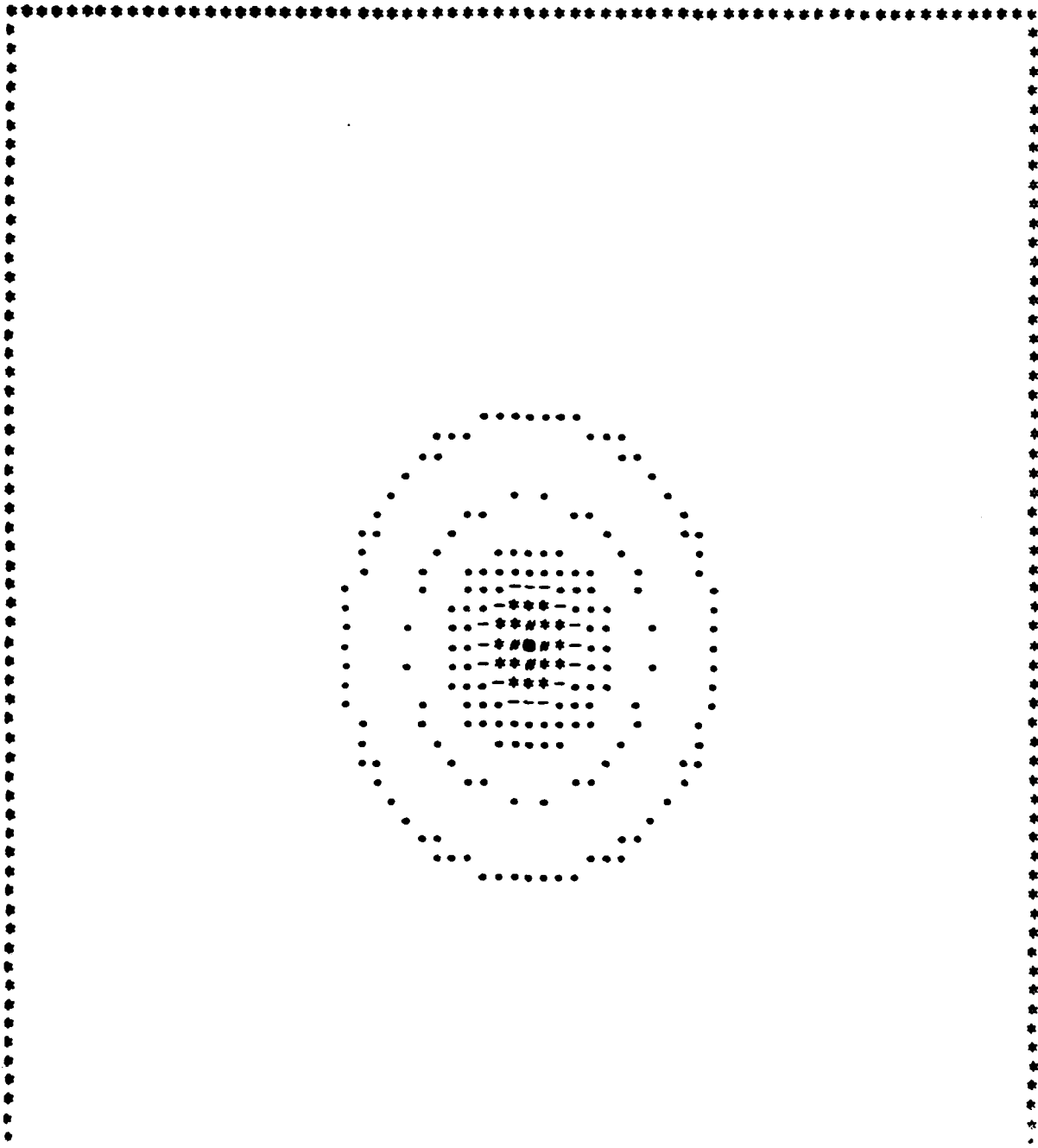
THIS IS STEP 14

SUML = 0.4651438E+00    SUMR = 0.5348561E+00

BEAM WAIST IN X AND Y DIRECTIONS IS    18.77950    18.78253    MICRONS



ERRADIAN



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IRRADIAN

GREY-SCALE CHARACTERS AND RANGES

.	0.244114E-03	0.155652E+01
-	0.155652E+01	0.311279E+01
+	0.311279E+01	0.466906E+01
*	0.466906E+01	0.622533E+01
#	0.622534E+01	0.778161E+01
R	0.778161E+01	0.933788E+01
□	0.933788E+01	0.108942E+02
■	0.108942E+02	0.124504E+02
■	0.124504E+02	0.140067E+02
■	0.140067E+02	0.155630E+02

THIS IS STEP 15

SUML = 0.4690364E+00      SUMR = 0.5309635E+00

BEAM WAIST IN X AND Y DIRECTIONS IS      18.87981      18.88269      MICRONS

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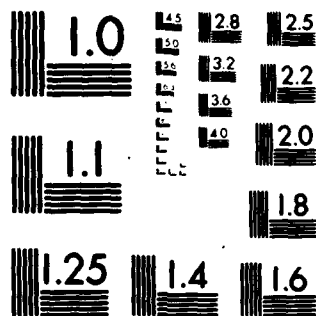
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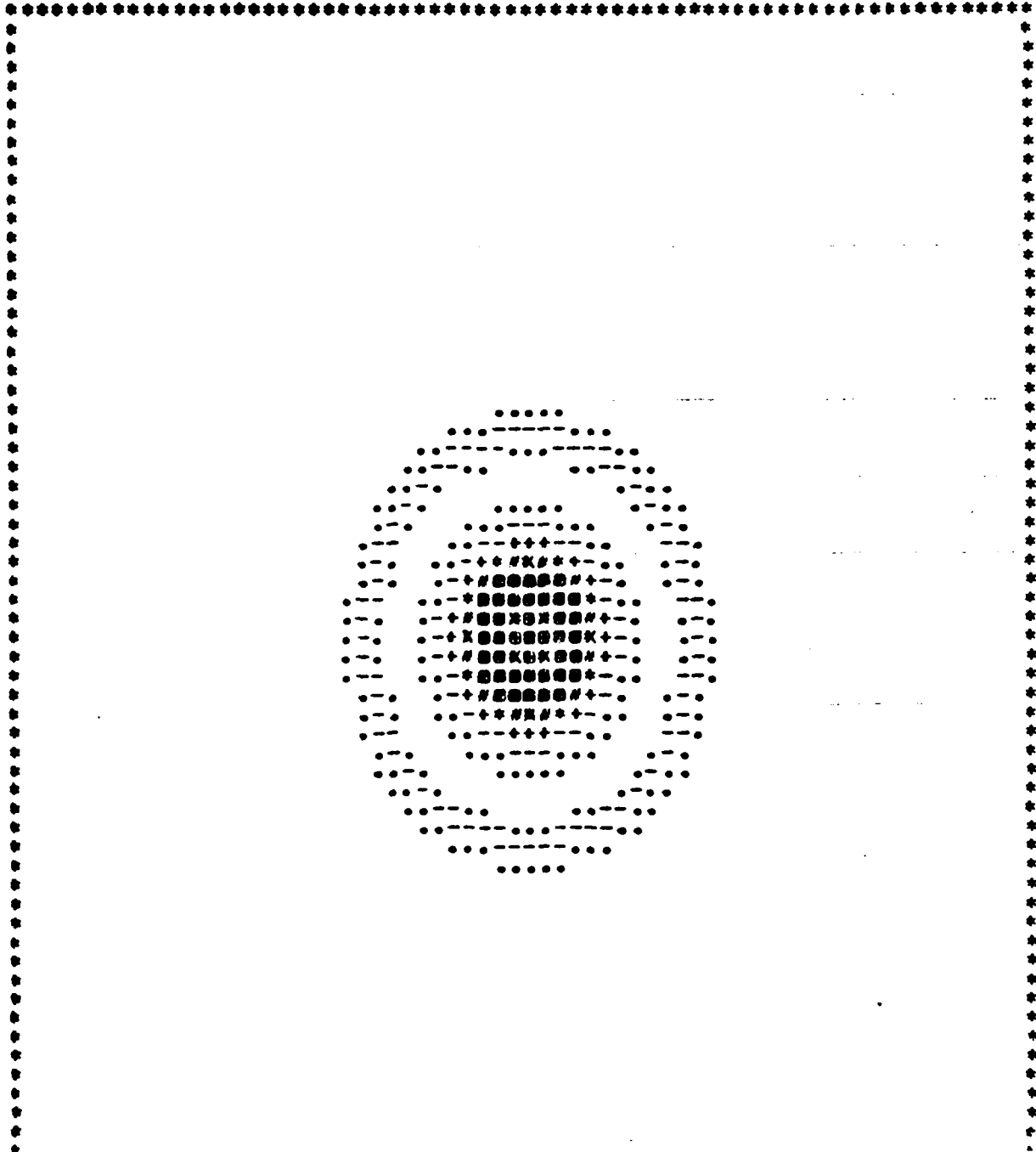
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IRRADIAN





IRRADIAN

GREY-SCALE CHARACTERS AND RANGES

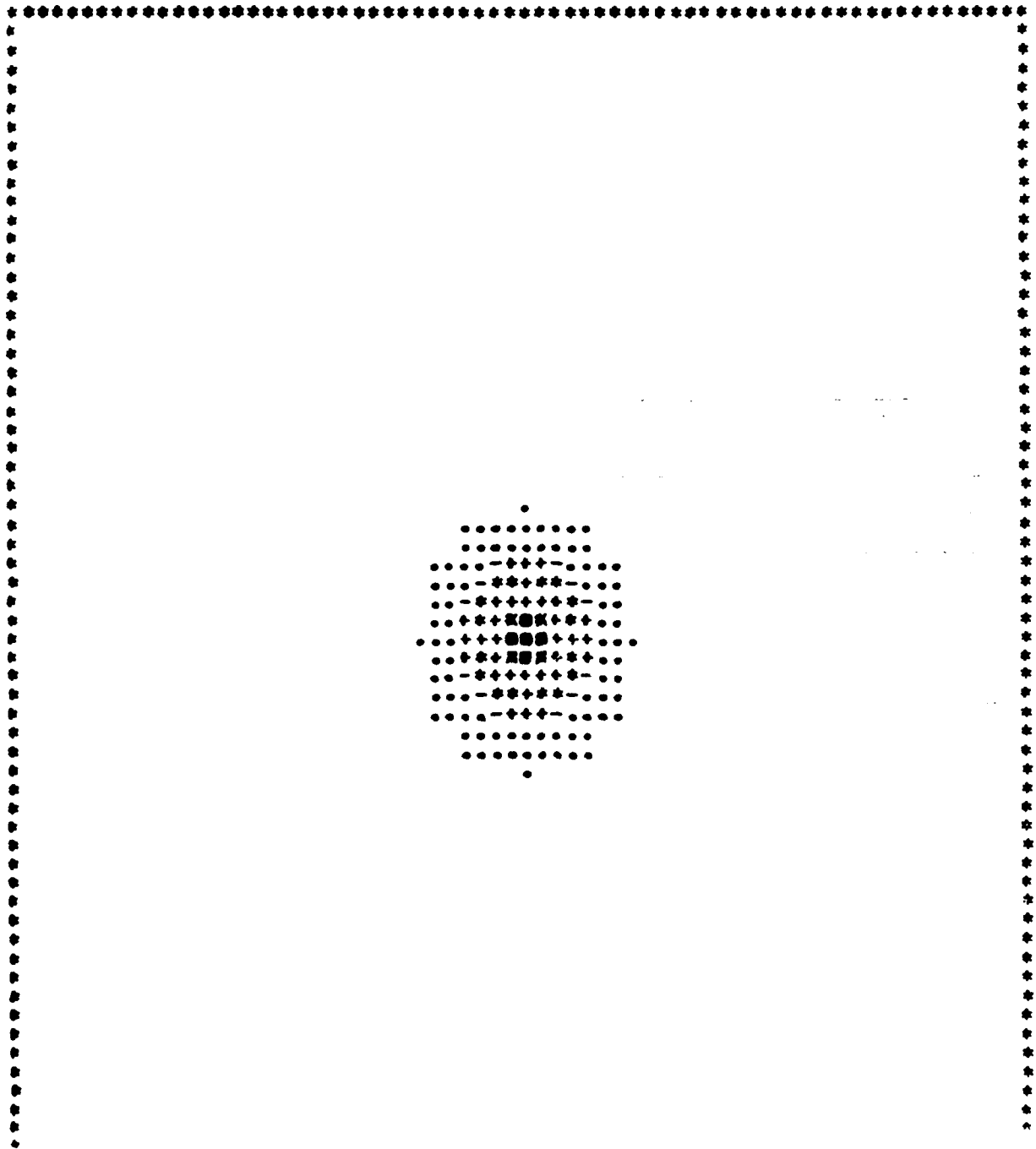
.	0.246905E-03	0.697666E+00
-	0.697666E+00	0.139509E+01
+	0.139508E+01	0.209250E+01
*	0.209250E+01	0.278992E+01
#	0.278992E+01	0.348734E+01
@	0.348734E+01	0.418476E+01
1	0.418476E+01	0.488218E+01
2	0.488218E+01	0.557960E+01
3	0.557960E+01	0.627702E+01
4	0.627702E+01	0.697444E+01

THIS IS STEP 16

SUML = 0.4632549E+00    SUMR = 0.5367450E+00

BEAM WAIST IN X AND Y DIRECTIONS IS    19.10669    19.10956    MICRONS

IRRADIAN



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ERRADIAN

KEY-SCALE CHARACTERS AND RANGES

.	0.911394E-03	0.131676E+01
-	0.131676E+01	0.263261E+01
+	0.263261E+01	0.394845E+01
*	0.394845E+01	0.526430E+01
#	0.526430E+01	0.658015E+01
@	0.658015E+01	0.789600E+01
R	0.789600E+01	0.921185E+01
Q	0.921185E+01	0.105277E+02
W	0.105277E+02	0.118435E+02
0	0.118435E+02	0.131594E+02

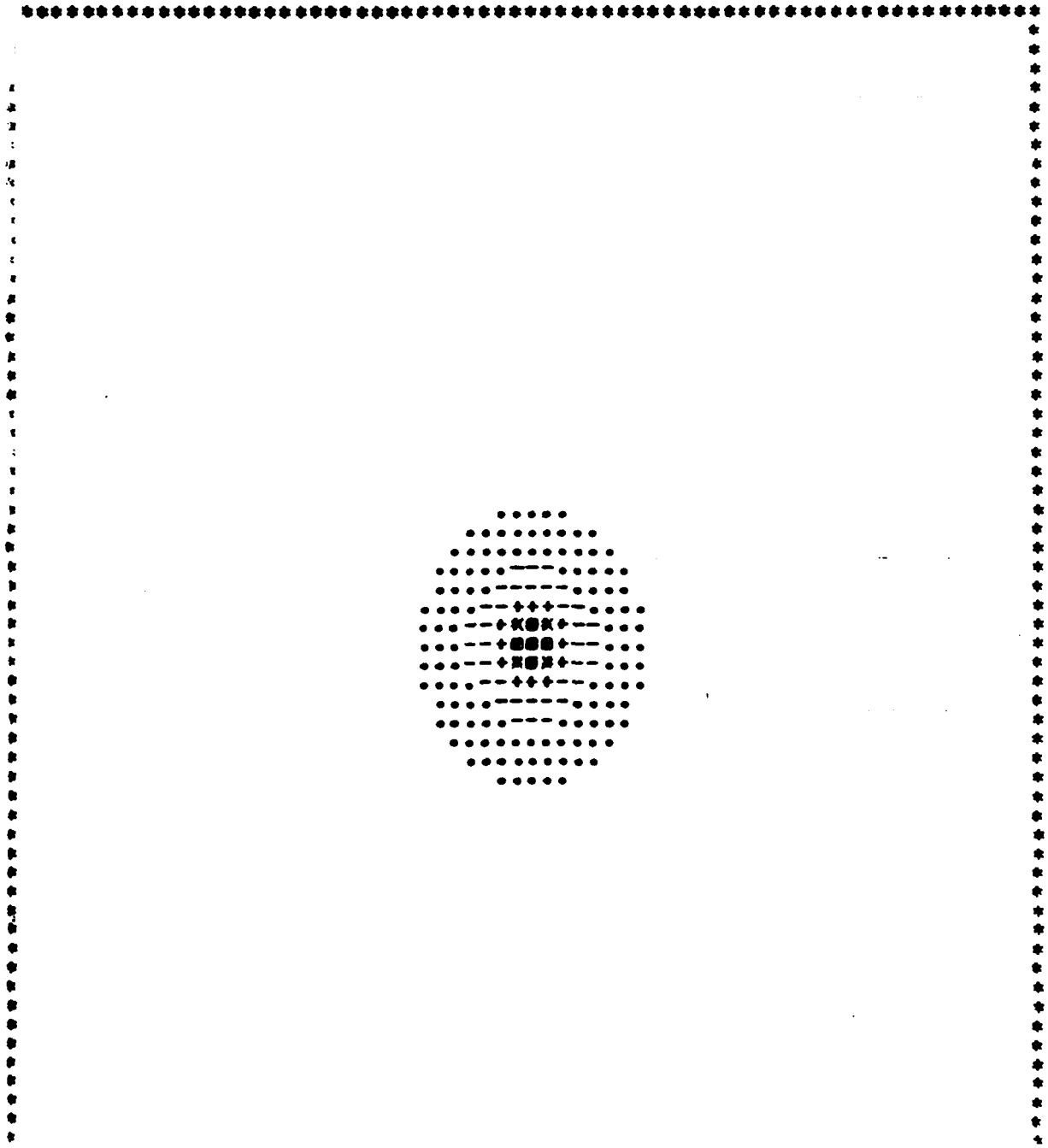
THIS IS STEP 17

SUML = 0.4618073E+00 SUQR = 0.5381926E+00

BEAM WAIST IN X AND Y DIRECTIONS IS 19.16985 19.17290 MICRONS



ERRADIAN



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IRRADIAN

GREY-SCALE CHARACTERS AND RANGES

	0.192583E-04	0.159542E+01
.	0.159542E+01	0.319081E+01
-	0.319081E+01	0.478621E+01
+	0.478621E+01	0.638161E+01
*	0.638161E+01	0.797700E+01
#	0.797700E+01	0.957240E+01
K	0.957240E+01	0.111678E+02
0	0.111678E+02	0.127632E+02
0	0.127632E+02	0.143586E+02
0	0.143586E+02	0.159540E+02

THIS IS STEP 18

SUML = 0.4730968E+00      SUMR = 0.5269032E+00

BEAM WAIST IN X AND Y DIRECTIONS IS      19.34006      19.34309      MICRONS

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**Rome Air Development Center**

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