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theorem (WFDPT) and establish a relationship between holography with wavelength diversity and inverse scattering. Coherent wave vector diversity radar techniques are used to access a finite volume of the 3-D Fourier space (also known as reciprocal space or \bar{p} -space) of the scattering object. The WFDPT permits the use of hybrid (digital/optical) computing that enables the retrieval and display of 3-D image information in parallel slices or cross-sectional outlines. The major attributes of this approach as compared to totally digital computing are its potential for displaying a true 3-D image in real-time. Wave-vector diversity methods appear suitable for the imaging of two classes of practical objects namely non-dispersive perfectly reflecting objects of the type often encountered in radar (and sonar) and semitransparent weakly scattering objects such as certain ultrasound and light scattering objects encountered in biology and medicine. The method has several unique characteristics. It furnishes true super-resolution, i.e. resolution exceeding the classical Rayliegh Limit of the available recording aperture, in this case a highly thinned (widely dispersed) broad-band coherent receiver array. True super-resolution is achieved because of an inherent aperture synthesis due to frequency diversity (frequency scanning, stepping or comb illumination) and conversion of spectral degrees of freedom into spatial image detail. Accordingly, when applied to the imaging of a dispersive object, a target signature rather than a geometrical image should be expected. Such a target signature could still be useful in object identification and classification since it contains information pertaining to the material composition of the object intermixed with geometrical image detail. The use of frequency diversity was found to lead to a unipolar impulse response. This is very useful in suppressing coherent noise (speckle) which is known to be the major drawback of coherent imaging.

Preliminary work on 3-D image display has yielded encouraging results on 3-D display from a series of weighted projection holograms of various slices of a test object. The projection holograms were viewed in rapid succession using the virtual Fourier transform.

To identify optimal and practical approaches to wave-vector diversity data acquisition a unique microwave measurement system has been assembled, , installed and tested in our anechoic chamber facility. The system was used also in the study of TDR (target derived reference) methods in which a reference for phase measurement can be furnished by the scattering object eliminating thus the need for costly local oscillator distribution networks and eliminating at the same time undesirable range phase ambiguities from the collected data. The measurement facility is computer aided furnishing thereby semi-automatic control . of object positioning or orientation, frequency stepping, data acquisition and storage, and final data correction and analysis. A high resolution CRT display enables the display of weighted projection holograms computed from the p-space data making use of the WFDPT. Preliminary results of this measurement system capabilities included in this report confirm its tremendous versatility. At this stage, the program strongly suggests the practical feasibility of a new generation of cost effective, real-time, super-resolving 3-D imaging radars that can because of their 3-D image slicing characteristic be appropriately referred to as Tomographic Radars (Tomos=slice in Greek).

Finally, a study of 3-D imaging using other forms of broadband radiation such as impulsive, random noise and particularly thermal emission for passive 3-D wavelength diversity imaging has been also initiated.

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The findings in this report are those of the authors and are not to be interpreted as the official position of the Air Force Office of Scientific Research or the U.S. Government.

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UNIVERSITY OF PENNSYLVANIA Philadelphia, Pennsylvania 19104

FINAL REPORT

SUPER RESOLUTION IMAGERY BY FREQUENCY SWEEPING

AIR FORCE OFFICE OF SCIENTIFIC RESEARCH/NE BUILDING 410 BOLLING AIR FORCE BASE WASHINGTON, D.C. 20332

August 15, 1980

GRANT NUMBER AFOSR-77-3256

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PREPARED BY

N.H. FARHAT AND C. WERNER

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HIGH RESOLUTION FREQUENCY SWEPT IMAGING

1. Introduction

The aim of the research work outlined in this final report was the analysis and investigation of methods by which frequency or wavelength diversity techniques can be employed to impart to a highly thinned, and therefore cost-effective, longwave (microwave or ultrasound) imaging aperture resolution capabilities better than its monochromatic classical (Rayligh) limit achieving thereby super-resolution by means of frequency synthesized apertures. This approach to longwave imaging gains practical significance when one considers the current highly developed state of the art of broadband microwave gear suitable for use in a new generation of cost-effective high resolution microwave imaging radars utilizing frequency diversity techniques.

It is well known that the development of longwave holographic imaging systems possessing resolution and image quality approaching those of optical systems is hampered by three factors: (a) prohibitive cost and size of longwave imaging apertures, (b) rapid deterioration of longitudinal resolution with range, (c) inability to view a 3-D image as with optical Fresnel holograms because of a wavelength scaling problem and (d) degradion of image quality by speckle or coherent noise because of the low numerical apertures attainable with present techniques. For example, a longwave imaging aperture operating at a wavelength of 3 cm should be about 3 km in size in order to achieve image resolution comparable to an ordinary photographic camera. In addition to inconvenient size, the cost of filling such a large aperture with suitable coherent sensors is clearly prohibitive. Furthermore, recall that in conventional longwave holography when optical image retrieval is utilized, it is necessary to store the longwave hologram data (fringe pattern) in an optical transparency suitable for processing on the optical bench using laser light. In order to avoid longitudinal distortion* of the reconstructed image, the size of the optical hologram replica must be m $(=\lambda \log/\lambda)$ times smaller than the longwave recording aperture. For the example cited earlier, this an optical hologram replica of less than a millimeter in size. It is means certainly not possible to view a virtual 3-D image through such a minute hologram even with optical aids since these tend to introduce their own longitudinal distortion. As a result, longwave holographers have long learned to forgo 3-D imagery and settled instead for 2-D imagery obtained by projecting the reconstructed real image on a screen. This permits lowering of the reduction factor m and consequently relaxing the resolution requirements of the photographic film which allows in turn the use of highly convenient Polaroid transparency film for preparation of the optical hologram replica. Because of the small size (measured in wavelength) of longwave apertures attainable in practice and the above methods of viewing the real image, speckle noise is always present leading to degration in image quality.

* Longitudinal distortion causes for example the image of a sphere to appear elongated in the range direction like a very long ellipsoid.

In this report we summarize the main results of our investigation under this grant. Our findings show that frequency diversity techniques not only circumvent the limitations discussed above but provide a means of viewing true 3-D images of distant objects such as satellites and aircraft. It is worthwhile to point out that our studies of wave-vector diversity imaging (or frequency swept imaging) were motivated to some extent by evidence of super-resolved "imaging" capabilities in the dolphin and the bat which are known to use frequency swept (chirp) signals in their "sonar" to discern small objects in their environment.

2. Summary of Important Results

The main findings of the study, details of which are given in the appendices and our publications (see list of publications), are outlined next.

(a) Wave-vector diversity (multifrequency and multiaspect) techniques can be used to enhance the amount of object information collected by a broadband coherent aperture deployed in the far field of the scattering object. Thus the data collected by a highly thinned array of coherent receivers intercepting the wavefield scattered from a distant 3-D reflecting object, as the frequency of its illumination (and/or its direction of incidence) are changed (see Fig. 1-a for example), can be stored as a 3-D data manifold in p-space (Fig. 2) from which an image of the object can be retrieved by means of a 3-D Fourier Transform. The size and shape of the 3-D data manifold, and therefore the resolution, depend on the relative positions of the object, the transmitter (illuminator), and the receiving array and on the spectral width of the illumination utilized.

(b) The data collected must be corrected for a quadratic phase factor F (caused by the unequal distances between the object and the receiving stations forming the widely dispersed imaging array) before it is stored in a 3-D manifold in p-space and an undistorted image of the 3-D reflecting object reconstructed through the 3-D Fourier transform operation. A bothersome range-azimuth ambiguity is also avoided through elimination of this quadratic phase term.

The most promising methods for data acquisition and correction is that which utilizes a target derived reference (TDR) at the synchronous detectors of the various receivers to correct for the unequal phase shifts or propagation time delays from the object to each receiver. In this approach the data furnished by the various receivers of the recording array is free of the undesired factor F. Therefore no additional processing by a computer will be necessary before filing the data at the appropriate locations in \tilde{p} -space. The TDR method has several advantages which include:

(i) Elimination of the need for a costly and unreliable central local oscillator distribution network.

(ii) Because TDR results in a recording configuration similar to that of a lensless Fourier Transform hologram, the resolution requirements from the recording device are greatly relaxed*.

*A. Macovski, "Hologram Information Capacity", J. Opt. Soc. Am., Vol. 60, Jan. 1970, pp. 21-29.

In longwave holography this fact is translated into a significant reduction of the number of receiving elements in the recording aperture. In addition the use of TDR allows us to place all the resolving power of the recording aperture on the target. This means that high resolution images of distant isolated targets should be feasible with array apertures consisting of tens of elements. The ability to synthesize a 2-D receiving aperture with a Wells arrayt consisting of two orthogonal linear arrays one of transmitters and the other of receivers provides further means of reducing the number of stations needed for data acquisition without sacrifice in resolution. A frequency swept Wells array of 10 transmitters and 10 receivers using a (2-4) GHz sweep should be able to easily furnish 10^4 3-D distinguishable resolution cells on the target which is more than sufficient for discerning the scattering centers on practical targets.

(iv) Greater immunity to phase fluctuations arising from turbulance and inhomogenieties in the propagation medium because both the reference and imaging signals arriving at each receiving element of the aperture travel roughly over the same path.

(v) TDR eliminates the range azimuth ambiguity and excessive bandwidth problems that arise in fast frequency swept imaging when the reference signal for the array aperture is distributed instead from the illumination source or a centrally located local oscillator phase locked to it.

Two TDR methods have been considered to some extent in our work to date. In one method which we term LFTDR (Low Frequency Target Derived Reference), the object is assumed to be illuminated simultaneously with a high frequency imaging signal and a low frequency signal that is a subharmonic of the illuminating frequency. The subharmonic reference frequency $\omega_{\rm c}$ is chosen such that k k<1, being the maximum linear dimension of the object and k = $\omega_{\rm c}/c$, c being the velocity of light. This places scattering from the object in the Rayliegh region where the object behaves as point scatterer with zero phase contribution. The far field phase of the reference signal at any receiver is therefore entirely due to propagation between a reference point formed at the object to the receiver. A method for measuring this reference signal phase and using it to correct the imaging signal phase due to propagation has been proposed by Porter* and analyzed for a one-dimensional object geometry. The reference signal phase and the imaging signal phase are measured separately at each receiving station with the aid of two receivers whose local oscillators (L.0)'s, one at the reference frequency and one at the imaging frequency, are phase-locked only to each other and not to a central local oscillator as would be the case were we to use a conventional receiver array. Phase locking of the two LO's can be accomplished by simply making the imaging L.O a harmonic of the

 +C.N. Nilsen and D.N. Swingler, "Quasi-Real-Time Inertialess Microwave Holography", Proc. IEEE (Letters), Vol. 65, March 1975, pp. 491-492.
 *R.P. Porter, "A Radar Imaging System Using the Object as Reference", Proc. IEEE (Letters), Vol. 59, Feb. 1971, pp. 307-308.

reference L.O. This would eliminate the difficulties encountered in the implementation of large or giant thinned coherent receiving arrays of the type required here, namely the distribution of a central local oscillator signal. A great reduction in cost and effort associated with installation of a central L.O. distribution network can thus be achieved. This cost reduction should be compared however with the cost of implementing a LFTDR. Because of the large difference between the high frequency imaging frequencies and the low frequency reference frequency required for the high resolution imaging of practical objects, the same microwave gear can not be used for both frequencies. This could increase system cost. In addition since the measured reference phase must be multiplied by a factor β equal to the ratio of the imaging to the reference frequency before being used as a reference phase in the imaging signal measurement, any errors in the reference phase measurement will also be amplified by this ratio. The precision of the reference phase measurement and phase error analysis are important and will have therefore to be examined further.

Another TDR methods which we call the Frequency Displaced Target Derived Reference (FDTDR) also shows promise. In this method, the analytical details of which are outlined in appendix I, the object is illuminated simultaneously during the sweep with two phase locked imaging frequencies ω_1 and $\omega_2 = \omega_1 + \Delta \omega$, $\Delta \omega$ being a small incremental frequency. This can be realized also by single side band modulation of the swept signal or by phase locking two sweep oscillators. Measurement of the differential phase between the signals scattered from the target at these frequencies yields $\frac{\Delta \omega}{c}$ ($R_{\rm T} + R_{\rm R}$), $R_{\rm T}$ being

the distance from the transmitter to the object and R_p being the distance from the object to the receiver. Multiplication of this phase by $\omega_1/\Delta\omega$ yields the phase factor F at frequency ω_1 which would be used to correct the phase measured at ω_1 . At first look this method would appear to still require a reference local oscillator. This however is not so since the procedure outlined above need not involve explicit phase measurements and multiplications. For example by mixing the two received signals at ω_1 and $\omega_1 + \Delta\omega$ in a square law detector at each receiver a beat signal at frequency $\Delta\omega$ is derived whose phase is equal to $\frac{\Delta\omega}{c}$ ($R_R + R_T$). The phase shift of

this signal due to the object is effectively zero because the wavelength at $\Delta \omega$ is much larger than the object extent making it behave effectively as a point scatterer. Harmonic mixing of the signal ω_1 received at each receiver with this beat signal should yield the corrected p-space data at ω_1 . Because of the small difference $\Delta \omega$ between the two frequencies ω_1 and ω_2 utilized, the effect of phase errors due to system and atmospheric propagation could be more completely cancelled in this method than in the low frequency TDR methods. The small difference $\Delta \omega$ means also that unlike the LFTDR case the same microwave gear (antennas, transmission lines and other microwave circuit components) can be utilized in the handling of the reference and imaging signals. A variation of the TDR technique involving double side-band modulation is also possible and appears to be more simple to implement than the single side-band method.

(c) Because in addition to being dependent on geometry, the dimensions of the 3-D data in p-space shown in Fig. 2 are dependent on the spectral range of the illumination, super-resolution (i.e. resolution beyond the classical limit of the available physical aperture) is achieved. This aperture synthesis by wave-vector or frequency diversity helps cut down array cost (since a thinned array can be used to frequency synthesize a large array with higher filling factor).

(d) Fourier Domain Projection Theorems (see appendixes II and III for details) enable the generation of two dimensional holograms from projections (or weighted projections) of the corrected \bar{p} -space 3-D data manifold of Fig. 2 permitting thereby optical image retrieval of the 3-D object in slices parallel to the projection plane one at a time. For example, Fig. 1-b shows the projection hologram for the \bar{p} -space data obtained in a computer simulation of the arrangement shown in Fig. 1-a. The central cross-sectional outline of the object (the two 1 m diameter reflecting spheres of Fig. 1-a) retrieved from this projection hologram by means of a 2-D Fourier transform carried out on the optical bench is shown in Fig. 1-c. A similar example is shown in Figs. 3 and 4. Figure 3 shows a second test object consisting of 3-D distribution of a set of 8 point scatterers with locations and spacings given in the Figure. Figure 4 shows the projection holograms corresponding to the three slices of the object containing the point scatterers and the image retrieved from each. The sweep width in this example, as in the previous example, was (2-4) GHz however the number of receivers in the recording array has been reduced from 50 to 16. These computer simulations demonstrate that a 3-D (lateral and longitudinal) resolution of the order of twenty centimeters*is easily achieved with a frequency sweep covering only (2-4) GHz using a broad-band array of 16 receivers and one transmitter. Wider-spectral windows should yield better resolution. It is worthwhile to note in this respect that commercial microwave sweepers and synthesizers are available with a spectral coverage of (.1-25) GHz indicating a potential for practical resolutions of the order of possibility few centimeters with cost-effective broad-band apertures consisting of tens of receivers operating with one central illuminator.

(e) The viewing or the display of a true 3-D image of the various slices or cross-sectional outlines should be possible by reconstruction of the various projection holograms in rapid succession while projecting the reconstructed real images of the corresponding slices on a rapidly moving projection screen. The screen would be displaced rapidly (together with the Fourier transforming lens) on the optical bench in the axial directions by small amounts proportional to the distances between the various slices. In another approach we have found that the 2-D virtual Fourier transform of a projection hologram can be carried out by simply viewing (with the unaided eye) a transparency containing an array of reduced replicas of the projection hologram arranged side-by-side with a point source. The image retrieved in this fashion would lie in the plane of the point source. This approach has the potential for 3-D display by viewing the virtual images retrieved from a series of projection holograms corresponding to different slices or cross-sectional outlines

*This means 10^3 distinguishable 3-D resolution cells in the $(2 \times 2 \times 2)m^3$ volume of the assumed object.

of the object passed in front of the eye in rapid succession while moving the reconstruction point source axially back and forth at a suitable rate of incremental axial displacements. A proposed electrooptical scheme that permits carrying out this procedure in real-time using a rapidly recyclable spatial light modulator (SLM) operating in a reflection mode is shown in Fig. 5. The computer, the high resolution CRT and the projection optics are used to project reduced noncoherent images of the various projection holograms in rapid succession on the SLM while the axial position of the reconstruction point sources is altered rapidly also under computer control. The point source need not be derived from a laser in order to yield an image but could also be a miniature "grain of wheat" light bulb. Details of this task are found in Appendix III.

(f) As seen in (e), unlike monochromatic longwave holographic imaging, there is no specific scaling requirements imposed on the projection holograms in order to avoid longitudinal distortion in the optical reconstruction circumventing thus the wavelength scaling problem.

(g) Because of the broad spectral extent of the illumination used and ability to display the reconstructed image in separate slices, speckle or coherent noise, which is known to plague coherent imaging systems, is suppressed making the system behave in as far as image noise is concerned like a noncoherent imaging system but at the same time enjoy the superior detection characteristics associated with synchroneous detection techniques.

(h) The broad-band nature of the imaging process also helps suppress undesirable image detail that could arise from object resonances which could seriously degrade image quality in a monochromatic imaging system.

(i) The data collected at every receiver, represents after correction, essentially the frequency response of the scattering object measured from a different aspect angle. Assuming the scattering process is linear, this frequency response is related to the impulse response of the object by a Fourier transform (see ref. 5 in List of publications). This suggests that impulse illumination can be utilized instead of frequency swept illumination. When this is done, the 3-D data manifold in p-space may be generated by Fourier transforming the impulse response at each receiver, correcting the data for the Factor F mentioned in (b), and storing the result in the appropriate p-space locations for each receiver. The resulting p-space volume accessed in this fashion can then be employed as described earlier to yield 3-D image information. Impulse illumination is desirable in certain instances of rapid target motion but may be more difficult to implement than frequency swept illumination. Since the impulse reponse of a time invariant linear system can also be deduced from white noise excitation and corellation of the output response with the input as described elsewhere in more detail (see 5 in list of publications), it follows that the techniques described in this report for coherent broadband radiation should be equally applicable with minor signal processing modification to noise-like broadband

radiation including passive black-body radiation.

(j) Experimental verification for both the principle of frequency diversity imaging and the TDR concept were obtained with the aid of a semi-automated network analyzer configured and installed in a recently refurbished anechoic chamber within the scope of this program(see Figs. 6 and 7). This versatile system is capable of vector (amplitude and phase) measurements of wavefields scattered from test objects situated in the anechoic chamber over any frequency range lying in the (.1-18)GHz range for a variety of polarizations. A test object consisting of two parallel cylinders 25 cm apart each 5 cm in diameter 50 cm long was mounted on a rotating styrofoam pedestal that is and under computer control and illuminated as shown in Fig. 8. The distance from the center of the object to the illuminating parabolic antenna to the left and the receiving horn feeding the network analyzer was 2.5 m. The complex frequency response of this object was measured in the (5-14)GHz range and the data stored for 128 object orientation covering 360°. The stored data was corrected for rangephase with a synthetic TDR generated in the computer and the corrected data displayed and photographed yielding the frequency swept hologram shown in Fig. 9 (c). The image retrieved from this hologram via an optical Fourier transform carried out on the optical bench is shown in Fig. 9 (d). For reasons of comparison a computer simulation of this experiment assuring a (2-18)GHz sweep was performed. The resultant range-phase corrected hologram and the image retrieved from it optically are shown in Fig. 9 (a) and (b). Further detail on this phase of the program were reported in an MSc. thesis made part of this report in Appendix IV. This part of the program is being continued with the aim of further enhancing measurement accuracy and demonstrating imaging of a nonsimple 3-D test object such as a model aircraft utilizing polarization diversity to further enhance image quality.

3. Conclusions.

The primarily analytical and numerical study of frequency diversity imaging performed under this grant demonstrates conclusively the feasibility of a new generation of coherent broadband imaging radars capable of furnishing 3-D image detail of distant target with cost effective giant apertures and efficient digital/optical signal processing.

Future work in this area will focus more on the analysis and identification of optimal methods for data acquisition, processing and 3-D display. The ultimate aim is the generation of design criteria for a prototype system and its assessment in the 3-D imaging of low flying aircraft passing within range of our facilities on route for landing at the Philadelphia Airport.

List of Publications

- N.H. Farhat, "Frequency Synthesized Imaging Apertures", Proc. 1976, International Optical Computing Conference, IEEE Cat. #76 CH 1100-7C, pp. 19-24.
- N.H. Farhat, M.S. Chang, J.D. Blackwell and C.K. Chan, "Frequency Swept Imaging of a Strip", Proc. 1976, Ultrasonics Symposium, IEEE Cat. #76 CH 1120-5SU.
- 3. J.D. Blackwell and N.H. Farhat, "Image Enhancement in Longwave Holography by Electronic Differentiation", Optics Communications, Vol. 20, Jan. 1977, pp. 76-80.
- 4. C.K. Chan, N.H. Farhat, M.S. Chang and J.D. Blackwell, "New Results in Computer Simulated Frequency Swept Imaging", Proc. IEEE (Letters), Vol. 65, pp. 1214-1215, Aug. 1977.
- 5. N.H. Farhat, "Principles of Broad-Band Coherent Imaging", J. Opt. Soc. Am., Vol. 67, pp. 1015-1020, Aug. 1977.
- 6. N.H. Farhat, "Comment on Computer Simulation of Frequency Swept Imaging", Proc. IEEE, Vol. 65, pp. 1223-1226, Aug. 1977.
- 7. N.H. Farhat, "Comment on a New Imaging Principle", Proc. IEEE (Letters), Vol. 66, pp. 609-700, May 1978.
- 8. N.H. Farhat, "Microwave Holographic Imaging Prospects For a Real-Time Camera", SPIE, Vol. 180, Real-Time Signal Processing II, (1979).
- 9. N.H. Farhat and C.K. Chan, "Three-Dimensional Imaging by Wave-Vector Diversity", Acoustical Imaging, Vol. 8, A. Metherell (ed.), Plenum Press, New York (1980), pp. 499-515.
- 10. C.K. Chan and N.H. Farhat, "Frequency Swept Imaging of Three Dimensional Perfectly Reflecting Objects", IEEE Trans. on Antennas and Propagation -Special Issue on Inverse Scattering. (Accepted for publication.)
- 11. C.K. Chan, "Analytical and Numberical Studies of Frequency Swept Imaging", University of Pennsylvania, Ph.D. Dissertation (1978).
- 12. N.H. Farhat, "Microwave Holography and Coherent Tomography", (Invited paper). Presented at 1980 IEEE/MTT's International Microwave Symposium, Electromagnetic Dosimetric Imaging. (To be published in special conference proceedings).

Related Publications

- 1. N.H. Farhat, "New Imaging Principle", Proc. IEEE (Letters), Vol. 64, pp. 379-380, March 1976.
- 2. N.H. Farhat, T. Dzekov and E. Ledet, "Computer Simulation of Frequency Swept Imaging", Proc. IEEE (Letters), Vol. 64, pp. 1453-1454, Jan. 1977.
- 3. G. Tricoles and N.H. Farhat, "Microwave Holography: Applications and Techniques", Invited paper, Proc. IEEE, Vol. 65, pp. 108-121, Jan. 1977.
- M.A. Kujoory and N.H. Farhat, "Microwave Holographic Substraction for Imaging of Buried Objects", Proc. IEEE (Letters), Vol. 66, pp. 94-96, Jan. 1978.
- 5. M.A. Kujoory and N.H. Farhat, "Format Generation For Double Circular Scanners For Use in Longwave Holography", Acoustical Imaging and Holography, Vol. 1, No. 2, pp. 133-141 (1979).
- 6. N.H. Farhat and J. Bordogna, "An Electro-Optics and Microwave-Optics Program In Electrical Engineering", IEEE Trans. on Education - Special Issue on Optics Education (accepted for publication).
- N.H. Farhat, "Holographically Steered Millimeter Wave Antennas", IEEE Trans. on Antennas and Propagation, Vol. AP-28, July 1980, pp. 476-480.



Fig. 1. Computer Simulation of Wave-vector Diversity Imaging, (a) Geometry, (b) Projection Hologram, (c) Retrieved Central Cross-sectional Image.

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Fig.2 Three dimensional p-space data generated by frequency sweeping and collected by a 2-D array of receivers.











Projection holograms and their optical reconstructions for the set of point scatterers in Fig.7.10 at different R' planes. (a) Hologram and reconstructed image of scatterers at $R_{z=-z_{0}}^{t=-z_{0}}$ plane. (b) Hologram and image at $R_{z=0}^{t=0}$ plane. (c) Hologram and image at $R_{z=z_{0}}^{t=2}$ plane. $x_{0}=y_{0}$ $=z_{0}=100$ cm.



True 3-D image reconstruction based on the virtual Fourier transform. ۍ. Fig.



i



View of Microwave Anechoic chamber showing illuminator antenna and a calibration sphere in Lackground.



(a)



Fig. 8. Two views of dual-cylinder test object in Anechoic chamber. (a) View showing illuminator to the left and the receiving horn on the right separated by absorbing barrier. (b) View showing test object mounted on rotating styrofoam pedastel. Cylinders are 5 cm in diameter, 50 cm long, 25 cm apart.



Fig. 9. Frequency swept holograms and retrieved images for a dual-cylinder test object. (a) Computed frequency swept hologram for a (2-18)GHz sweep and (b) retrieved image; (c) measured frequency swept hologram for a (5-,14)GHz sweep and (d) retrieved image.



APPENDIX I

The Frequency Displaced Target Derived Reference

A second TDR method which we refer to as a Frequency Displaced Target Derived Reference (FDTDR) method also shows promise. This method involves simultaneous illumination of the object with two phase locked imaging frequencies ω_1 and $\omega_2 = \omega_1 + \Delta \omega$ that differ by a small frequency increment $\Delta \omega$. 1 2 - 1 Referring to eq. (9) of ref. 10 (see list of publications) we can write for the far field at a given receiver location R_p ,

$$\psi_1(k_1,R_R) = \frac{jk_1}{2\pi R_R} e^{-jk_1(R_T+R_R)} \int U(\bar{r}) e^{-j\bar{p}_1\cdot\bar{r}} d\bar{r}$$
 (1)

$$\psi_{2}(k_{2},R_{R}) = \frac{j(k_{1}+\Delta k)}{2\pi R_{R}} e^{-jk_{1}(R_{T}+R_{R})} e^{-j\Delta k(R_{T}+R_{R})}$$

$$x \int U(\bar{r})e^{-j\bar{p}_{1}} (1+\frac{\Delta \omega}{\omega 1}) \cdot \bar{r}$$

$$d\bar{r} \qquad (2)$$

where $k_{1,2} = \omega_{1,2}/c$ and $\Delta k = \Delta \omega/c$.

By making $\Delta\omega/\omega \ll 1$ the integral in (2) will approach that in (1). The only difference between the far fields ψ_1 and ψ_2 at the receivers is then the phase term $\Delta k(R_T + R_p)$. Measurement of this phase difference yields $(R_T + R_p)$ since Δk is known. This information can be used to correct the phase of either the ψ_1 or ψ_2 signals to obtain the required \bar{p} -space information.

$$\Gamma(\bar{p}) = \int U(\bar{r}) e^{j \bar{p} \cdot \bar{r}} d\bar{r}$$
(3)

Appendix II

HOLOGRAPHY, WAVE-LENGTH DIVERSITY AND INVERSE SCATTERING

ABSTRACT

The use of wavelength diversity to enhance the performance of thinned coherent imaging aperatures is discussed. It is shown that wavelength diversity lensless Fourier transform recording arrangements that utilize a reference point source in the vicinity of the object can be used to access the three-dimensional Fourier space of nondispersive perfectly reflecting or weakly scattering objects. Hybrid (opto-digital) computing applied to the acquired 3-D Fourier space data is shown to yield tomographic reconstruction of 3-D image detail either in parallel or meridonal (central) slices. Because of an inherent ability of converting spectral degrees of freedom into spatial 3-D image detail true super-resolution is achieved together with suppression of coherent noise. The similarity of the key equations derived to those of inverse scattering theory is pointed out and the feasibility of using other forms of broadband radiation such as impulsive, noise and thermal is discussed. Finally, the potential of utilizing wavelength diversity imaging in microscopy and telescopy are discussed.

INTRODUCTION

A frequently encountered question in the science of image formation is how to make an available aperature collect more information about. the scene or object being imaged in order to enhance its resolving power beyond the classical Rayleigh limit. This process is known as super-resolution and is relevant to all imaging systems whether holographic or conventional. There are five known methods for achieving super-resolution. These include: weighting or apodization of the aperature data^{1,2}; analytic continuation of the wavefield measured over the aperature^{3,4}; use of evanescent wave illumination⁵; maximum entropy method⁶: and use of the time channel⁷. Weighting and analytical continuation techniques are known to become rapidly ineffective as the signal to noise ratio of the data collected decreases. Maximum entropy techniques are known to be more robust as far as noise is concerned but involve ususally extensive computation. Illumination with evanescent waves is practical in limited situations where full control of the recording arrangement exists as in microscopy for example. This leaves the time channel approach in which one can collect in time more information about the object through the available recording aperature by altering the object aspect relative to the aperature by

means of roation or linear motion^{8,9} or by altering the parameters of the illumation such as directions of incidence, wavelength and/or polarization. These later operations are known to increase the degrees of freedom of the wavefields impinging on the recording aperature enhancing thereby their ability to convey information about the nature of the scattering object. Sophisticated imaging systems endevour to convert the nonspatial degrees of freedom of the wavefield, e.g., angular, spectral and polarization to spatialimage detail enhancing thereby the resolution capability beyond the classical Rayleigh limit of the available physical aperature. Obviously such procedures involve more signal processing than that performed by conventional imaging with lens systems or holography.

In this paper we consider generalizing the holographic concept to include wavelength diversity as a means of enhancing resolution. A quick examination of the basic equations of holography reveals that the lensless Fourier transform hologram recording arrangement is amenable to this generalization. This conclusion is used then as a starting point for a Fourier optics formulation of wavelength diversity imaging of 3-D (three dimensional) nondispersive objects. The results show that measurement of the multiaspect or multistatic frequency (or wavelength) response of the 3-D object permits accessing its 3-D Fourier space. The resulting formulas are identical to those obtained from a multistatic generalization of inverse scattering^{10,11,12} establishing thus a clear connection between holography and the inverse scattering imaging problem. The inclusion of wavelength diversity in

scattering imaging problem. The inclusion of wavelength diversity in holography is shown to have several important features: (a) the availability of the 3-D Fourier space data permits 3-D image retrieval tomographically in parallel or meridonal (central) slices or crosssectional outlines by the application of Fourier domain projection theorems, (b) suppression of coherent noise and speckle in the retrieved image, (c) removal of several longstanding constraints on longwave (microwave and acoustical) holography such as the impractically high cost of the aperatures needed, the inability to view a true 3-D image as in optical holography because of a wavelength scaling problem, and minimization of the effects of resonances on the object.

WAVELENGTH DIVERSITY

We start by inquiring into the conditions under which the data from N holograms of the same nondispersive object recorded over the same aperture, each at a different wavelength, can be combined to yield a single image superior in quality to the image retrieved from any of the individual holograms.

One approach to answering the question posed above would be to determine the conditions under which the well known formulas 13 for the focusing condition, magnification and image location in holography can be made independent of wavelength. This quickly leads to the conclusion that wavelength independence can be met if a reference point source centered on the object is used and proper scaling of the individual holograms by the ratio of recording to the reconstruction wavelength is performed before super-position 15,24. The former condition is that for recording a lensless Fourier transform hologram 14 where the presence of the reference point source in the object plane leads to the recording of a Fraunhofer diffraction pattern of the object rather than its Fresnel diffraction pattern because of the elimination of a quadratic phase term in the object wavefield in the recorded hologram. This is known to result in a highly desirable reduction in the resolution required from the hologram recording medium and is therefore of . practical importance especially in nonoptical holography. More detail of the processing involved in combining the data in multi

wavelength hologram can be found $elsewhere^{15}$.

Additional insight into the process of attaining super-resolution by wavelength diversity is obtained by considering the concept of wavelength or frequency synthesized aperature 16-20The synthesis of a one dimensional aperature by wavelength diversity is based on the simple fact that the Fraunhofer or far field diffraction pattern of a nondispersive planar object changes its scale, i.e. it "breathes", but does not change its shape (functional dependence), as the wavelength is changed. A stationary array of broadband sensors capable of measuring the complex field variations deployed in this breathing diffraction pattern at suitably chosen locations would sense different parts of the diffraction pattern as the wavelength is altered collecting thereby more information on the nature of the diffraction pattern and therefore on the object that gave rise to it than if the wavelength was fixed (stationary diffraction pattern). Each stationary sensor in the array is thus able to collect as the wavelength is changed, and the breathing diffraction pattern sweeps over it, the same set of data or information collected by a movable sensor mechanically scanned over the appropriate part of the diffraction pattern when it is kept stationary by fixing the wavelength. Hence the term wavelength or frequency synthesized aperture.

The orientation and location of the wavelength synthesized aperture for any planar distribution of sensors deployed in the Fraunhofer diffraction pattern of a planar object and the retrieval of an image from the data collected has been treated earlier^{16,17}. It was clear, however, that extension of the wavelength diversity concept to the case of 3-D objects is necessary before its generality and practical use could be established.

For this purpose we considered²⁰ as shown in Fig. 1(a) an isolated planar object of finite extent with reflectivity $D(\bar{\rho}_0)$, where $\bar{\rho}_0$ is a two dimensional position vector in the object plane (x_0, y_0) . The object is illuminated by a coherent plane wave of unit-amplitude and of wave vector $\bar{k}_i = k \ 1_{k_i}$ produced for example by a distant source located at \bar{R}_T . The wavefield scattered by the object is monitored at a receiving point designated by position vector \bar{R}_R belonging to a recording aperture lying in the far field region of the object. The receiving point will henceforth be referred to as the receiver and the source point at the transmitter. The position vectors $\bar{\rho}_0$, \bar{R}_T and \bar{R}_R are measured from the origin of a cartesean coordinate system (x_0, y_0, z_0) centered in the object. The object is assumed to be nondispersive i.e., D is independent of k. However, when the object is dispersive such that $D(\bar{\rho}_0, k) = D_1(\bar{\rho}_1)D_2(k)$ and $D_2(k)$ is known, the analysis presented here can easily be modified to account for such object dispersive figures.

Referring to Figure 1(a) and ignoring polarization effects, the field amplitude at \bar{R}_R caused by the object scattered wavefield may be expressed as,

$$\psi(k,\bar{R}_{R}) = \frac{jk}{2\pi} \int D(\bar{\rho}_{0}) e^{-j\bar{k}_{1}} \cdot \bar{r}_{T} - \frac{-jk}{r_{R}} d\bar{\rho}_{0} \qquad (1)$$

where $d\bar{\rho}_0$ is an abbreviation for $dx_0 dy_0$ and the integration is carried out over the extent of the object. Noting that $\bar{r}_T = \bar{\rho}_0 - \bar{R}_T$, $\bar{R}_T = -R_T \bar{I}_{k_1}$ and using the usual approximations valid here: $r_R \simeq R_R + \rho_0^2/2R_R - \bar{I}_R \cdot \bar{\rho}_0$ for the exponential in (1) and $r_R \simeq R_R$ for the denominator in (1) where $\bar{I}_R = \bar{R}_R/R_R$ and $\bar{I}_{k_1} = \bar{k}_1/k$ are unit vectors in the \bar{R}_R and \bar{k}_1 directions respectively, one can write eq. (1) as,

$$\psi(k, \bar{R}_R) = \frac{jk}{2\pi R_R} e^{-jk(R_T + R_R)} \int D(\bar{\rho}_0) e^{-j\bar{p}} \cdot \bar{\rho}_0 d\bar{\rho}_0,$$
 (2)

where we have used the fact that the observation point is in the far field of the object so that $\exp(-jk\rho^2/2R_R)$ under the integral sign can be replaced by unity. In eq. (2), $\bar{p} = k(\bar{1}_{k_1} - \bar{1}_R)^{\Delta} p_x \bar{1}_x + p_y \bar{1}_y + p_z \bar{1}_z$ is a three dimensional vector whose length and orientation depend





Fig. 1. Geometries for wavelength diversity imaging. (a) Two dimensional object, (b) Three dimensional object with the n-th meridonal (central) slice and cross sectional outline c shown.

on the wavenumber k and the angular positions of the transmitter and the receiver. For each receiver and/or transmitter present, \bar{p} indicates the position vector for data storage. An array of receivers for example would yield therefore as k is changed (frequency diversity) or as $\bar{k} (=k\bar{l}_{k_1})$ is charged (wave-vector diversity) a 3-D data manifold. The projection of this 3-D data manifold on the object plane yields $\psi(k,R_T)$ because $\bar{p} \cdot \bar{\rho}_0 = \bar{p}_t \cdot \bar{\rho}_0 = p_{\chi\chi_0} + p_y y_0$ where $p_\chi = k(\bar{l}_{k_1} - \bar{l}_R)_{\chi}$ and $p_y = k(\bar{l}_{k_1} - \bar{l}_R)_y$ are the cartesian components of the projection \bar{p}_+ of \bar{p} on the object plane. Accordingly eq. (2) can be expressed as,

$$\psi(k, R_{R}) = \frac{jk}{2\pi R_{R}} e^{-jk(R_{T} + R_{R})} \int D(x_{0}, y_{0}) e^{-j(P_{X}x_{0} + P_{y}y_{0})} dx_{0} dy_{0}$$
(3)

Because of the finite extent of the object, the limits on the integral can be extended to infinity without altering the result. The integral in (3) is recognized then as the two dimensional Fourier transform $D(p_x, p_y)$ of $D(x_0, y_0)$. It is seen to be dependent on the object reflectivity function, the angular positions of the transmitter and the receiver and on the values assumed by the wavenumber k but is entirely independent of range. Information about D can thus be collected by varying these parameters. Note that the range information is contained solely in the factor $F = jk \exp \left[-jk(R_T + R_R)\right]/2\pi R_R$ preceeding the integral. The field observed at $\bar{\mathtt{R}}_{R}$ has thus been separated into two terms one of which, the integral \tilde{D} , contains the lateral object information and the other F, contains the range information. The presence of F in eq. (3) hinders the imaging process since it complicates data acquisition and if not removed, gives rise to image distortion because R_p is generally not the same for all receivers. To retrieve an image of the object via a 2-D Fourier transform of eq. (3), the factor F must first be eliminated. Holographic recording of the complex field amplitude given in (3) using a reference point source located at the center of the object will result in the elimination of the factor F and the recording of a Fourier transform hologram. This

operation yields D over a two dimensional region in the p_x, p_y plane.

The size of this region, which determines the resolution of the retrieved image depends on the angular positions of the transmitter and the receiver and on the values assumed by k, i.e. the extent of the spectral window used. The later dependence on k implies super-resolution imaging capability because of the frequency synthesized dimension of the 2-D data manifold generated. Because of the dependence of resolution on the relative positions of the object, the transmitter, and receiving aperture, the impulse response is clearly spatially variant. In fact a receiver point situated at \bar{R}_R for which \bar{p} is normal to the object

plane can not collect any lateral object information because for this condition (\bar{p} , $\bar{\rho}_0 = 0$) the integrals in (2) and (3) yield a constant. Such receiving point is located in the direction of specular reflection from the object where the diffraction pattern is stationary i.e. does not change with k. In this case the observed field is solely proportional to F containing thus range information only. Obviously this case can easily be avoided through the use of more than one receiver which is required anyway when 2-D or 3-D object resolution is sought^{20,21}.

The analysis presented above can be extended to three dimensional objects by viewing a 3-D object as a collection of thin merdional or central slices as depicted in Fig. 1(b) each of which representing a two dimensional object of the type analyzed above. With the n-th slice we associate a cartesean coordinate system x_0, y_0, z_0 that

differ from other slices by rotation about the common x_0 axis. Since the vectors \bar{p} , \bar{R}_T and \bar{R}_R are the same in all n-coordinate systems, eq. (3) holds. $\psi_n(k,\bar{R}_R)$ is then obtained from projection of the three dimensional data manifold collected for the 3-D object on the x_0, y_{on}

plane associated with the n-th slice. An image for each slice can then be obtained as described before. An inherent assumption in this argument is that all slices are illuminated by the same plane wave. This is a reasonable approximation when the 3-D object is weakly scattering and the Born approximation is applicable or when the 3-D object is perfectly reflecting and does not give rise to multiple reflections between its parts. In the later case the two dimensional meridonal slices $D_n(\bar{\rho}_0)$ deteriorate into contours, such as C in Fig. 1(b) defined by the intersection of the meridonal planes with

the illuminated portion of the surface of the object. Accordingly we can write for the n-th meridonal slice or contour,

$$\psi_{n}(k,R_{R}) = F \int D_{n}(\bar{\rho}_{0}) e^{-j\bar{p}\cdot\bar{\rho}_{0}} n d\bar{\rho}_{0}$$
(4)

We can regard $D_n(\bar{\rho}_0)$ as the n-th meridonal slice or contour of a three dimensional object of reflectivity $U(\bar{r})$ where \bar{r} is a three dimensional position vector in object space. This means that $D_n(\bar{\rho}_0) = U(\bar{r}) \quad \delta(z_0)$ where δ is the Dirac delta "function". Consequently eq. (4) becomes,

$$\psi_{n}(k,R_{R}) = F \int U(\vec{r}) \delta(z_{0}) e^{-j\vec{p}\cdot\vec{\rho}_{0}} n d\vec{\rho}_{0}$$
$$= F \int U(\vec{r}) \delta(z_{0}) e^{-j\vec{p}\cdot\vec{r}} d\vec{r}$$

where $d\bar{r}$ designated an element of volume in object space and where the last equation is obtained by virture of the sifting property of the delta function.

Summing up the data from all slices or contours of the object we obtain,

because

$$\sum_{n}^{\Sigma} U(\bar{r}) \delta(z_{0}) = U(\bar{r}).$$

Assuming that the Factor F in eq. (6) is eliminated as before, equation (6) reduces to

 $\psi(\bar{p}) = \int U(\bar{r}) e^{-j\bar{p}\cdot\bar{r}}$ (7)

which is the 3-D Fourier transform of the object reflectivity $U(\bar{r})$. Wavelength diversity permits therefore accessing the 3-D Fourier space of a nondispersive object providing thereby the basis for 3-D Lensless Fourier transform holography. An alternate formulation to that given above of super-resolved wave-vector diversity imaging of 3-D perfectly conducting objects is possible²² by extending the formulation of the inverse scattering imaging problem^{10,11} to the multistatic case, along lines that are similar but somewhat different than those given by Raz¹². The resulting scalarized formulas are identical to (7) establishing thus the connection between the holographic and the inverse scattering approaches to the imaging problem.

THREE DIMENSIONAL IMAGE RETRIEVAL

The above considerations of multiwavelength holography have lead us to determining a means by which the 3-D Fourier space of the object can be accessed employing synchroneous detection. It is clear that once the 3-D Fourier space data is available, 3-D image detail can be retrieved by means of an inverse 3-D Fourier transform which can be carried out digitally. Alternately, holographic techniques

(5)



Fig. 2. 3-D object consisting of a set of eight point scatterers shown in isometric and $R'_x-R'_y$ plane views at $R'_z=-z_0, 0, z_0, x_0=y_0=z_0$ =100 cm.


Fig. 3. Projection holograms and their optical reconstructions for the set of point scatterers in Fig. 2 at different R'_z planes. (a) Hologram and reconstructed image of scatterers at $R'_z=-z_0$ plane. (b) Hologram and image at $R'_z=0$ plane. (c) Hologram and image at $R'_z=z_0$ plane. $x_0=y_0=z_0=100$ cm.



Fig. 4. Arrangement used in computer simulation of wavelength diversity imaging.

can be invoked again. Fourier domain projection theorems²³ that are dual to the spatial domain projection theorem^{25,26} can be applied to the Fourier space data to produce a series of projection holograms from which 2-D images of meridonal or parallel slices of the object

can be retrieved on the optical $\operatorname{bench}^{20}$. This procedure does not involve any specific scaling of the size of the optical hologram transparency relative to the size of the original recording aperture by the ratio of the recording to the reconstruction wavelengths as in longwave holography where the scaling necessary for viewing a 3-D image free of longitudinal distortion ususally leads to an impractically minute equivalent hologram transparancy that cannot be readily viewed by an observer. The lateral and longitudinal resolutions in the retrieved image depend now on the dimensions of the volume in Fourier space accessed by wavelength diversity. This volume depends on the wavelength range and on the recording geometry. Thus the longitudinal resolution does not deteriorate now as rapidly with range as in conventional monochromatic imaging systems.

An example of computer simulations of frequency diversity holographic imaging of a 3-D object consisting of eight point scatterers distributed as shown in Fig. 2 is given in Fig. 3. Shown in Fig. 3 are three weighted Fourier domain projection holograms and the corresponding optically retrieved images for three equally spaced parallel slices of the object containing distinguishable 2-D distributions of scatterers. The simulated recording arrangement shown in Fig. 4 consisted of an array of 16 receivers equally distributed on an arc extending from $\phi =$

 40° to $\phi = 77.5^{\circ}$ surrounding a central transmitter capable of providing plane wave illumination of the object. The results shown were obtained with microwave imaging in mind assuming a frequency sweep of (2-4)GHz. They clearly indicate a leteral and longitudinal resolution capability of the order of 25 cm. Wider sweep widths yield better resolution. For example a (1-18)GHz sweep would yield a 3-D resolution of the order of 1.5 cm.

DISCUSSION AND CONCLUSIONS

Seeking means by which the information content in a hologram can be increased for example by wavelength diversity we have arrived at a formulation of 3-D Lensless Fourier transform holography capable of furnishing 3-D image detail tomographically. This ability of producing 3-D images in slices from coherently detected wavefields enable us to regard the method also as coherent tomography. The Fourier space accessed in the above fashion by wavelength diversity can be viewed as a generalized 3-D hologram in which one dimension has been synthesized by wavelength diversity. Such a generalized hologram contains not only spatial amplitude and phase data as in conventional holography but also spectral information and hence can yield better resolution than the classical Rayleigh limit of the available aperture operating at the shortest wavelength of the spectral window used. This super-resolving property is further enhanced through an inherent suppression of the effects of object resonances and coherent noise in the retrieved image, the latter being so because frequency diversity tends to make the impulse response of the system unipolar resembling that of a non-coherent imaging system that is free of speckle and coherent noise artifacts¹⁵. Futher enhancement of information content and resolution can be achieved by polarization diversity where the p space can be multiply accessed for different nonredundant polarizations of the illumination and the receivers and the resulting polarization diversity images added either coherently or non-coherently in order to achieve a degree of noise averaging as discussed elsewhere¹⁵.

The removal of several longstanding constraints on conventional longwave (microwave and acoustic) holography attained through the use of wavelength diversity as described here leads to a new class of imaging systems capable of converting spectral degrees of freedom into 3-D spatial image detail furnishing thereby true super-resolution. Wavelength diversity is applicable to the imaging of two classes of objects: perfectly reflecting objects of the type encountered in radar and sonar and weakly scattering objects of low or known dispersion of the type encountered in biology and medicine. The practical application of the concepts presented here to optical wavefield is presently under consideration. The availability of tunable dye lasers and electronic imaging devices suggest interesting possibilities of three dimensional wavelength diversity microscopy. Here one can conceive of an arrangement in which a minute semitransparent object with homogeneous or known dispersion is transilluminated by a collimated coherent light beam from a tunable dye laser which can also be made to provide a coherent reference point source in the immediate vicinity of the object. The resulting reference and the object scattered wavefields are intercepted by the photocathode of an electronic imaging device of known spectral response such as a vidicon. Because of the minute size of the object, the photocathode can easily be situated in the far field of the object. Thus nearly a lensless Fourier transform hologram recording arrangement results. The spatial frequency content in the resulting hologram is therefore expected to be sufficiently low to be resolved by a high resolution electronic imaging device. By recording and digitally storing the resulting detected hologram fringe pattern as a function of dye laser wavelength one can gain access to the 3-D Fourier space of the object since \bar{I}_{k_i} and \bar{I}_R

for the recording geometry are precisely known.

A similar recording arrangement can be envisioned in the active coherent imaging of a distant reflecting object (active telescopy) where the object can be made to furnish a reference point source situated on its surface like a wavelength independent stationary glint point or an intentionally placed retroreflector. Because in such an arrangement the reference and the object wavefields travel over the same path, atmospheric effects are expected to be minimized. The generation of an object derived reference geometry in longwave (microwave and acoustic) wavelength diversity imaging has been described elsewhere 20,27.

Finally it is worthwhile to note that since the scattering process is linear the multiaspect or multistatic frequency or wavelength response measurements referred to in this paper can be obtained also by measuring the multiaspect impulse response followed by Fourier

transformation of the individual impulse responses measured¹⁹. This means that impulsive illumination can also be utilized. Because the impulse response of a linear system can be measured by using random noise excitation and cross-correlating the output with the input¹⁹, a possibility of using random noise (white light) illumination and

cross-correlation detection techniques as a means for accessing the 3-D Fourier space of the object also emerges.

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APPENDIX III

THE VIRTUAL FOURIER TRANSFORM AND ITS APPLICATION IN THREE DIMENSIONAL DISPLAY

ABSTRACT

In contrast to the well known and widely used instantaneous Fourier transforming property of the convergent lens in coherent (laser) light, the "Virtual Fourier Transform" (VFT) capability of the divergent lens is less widely known or used despite many advantages. We will review the principle of the VFT and discuss its advantages in certain applications. In particular a method for viewing the virtual Fourier transform of a two dimensional function with the naked eye using an ordinary point source will be presented. A scheme for threedimensional image display based on a "Fourier domain projection theorem" utilizing varifocal VFT is described and a discussion of the properties of the displayed image given.

INTRODUCTION

Several sophisticated three dimensional (3-D) imaging techniques such as x-ray tomography¹, electron microscopy², crystallography², wave-vector diversity imaging and inverse scattering³, involve measurements that give access to a finite volume in the 3-D Fourier space of a 3-D object function. A 3-D image of the original object can then be reconstructed by computing the inverse 3-D Fourier transform. The retrieved iamge normally represents the spatial distribution of a relevant parameter of the object such as absorption, reflectivity, scattering potential, etc.

Obviously, the required inverse transform can be performed digitally. Digital techniques however often preclude real-time operation particularly when the object being imaged is not simple but contains considerable resolvable intricate detail. More importantly, because of the inherent two dimensionality of CRT computer displays, direct true 3-D image display is not possible. Present day computer graphic displays are capable of displaying 3-D image detail either in seperate crosssections or slices, or in a computed perspective (isometric) view of the object, or in some instances stereoscopically where an illusion of a 3-D scene is created in the mind of the observer who is required

usually to use special viewing glasses^{4,5}.

Hybrid (opto-digital) computing techniques offer an alternate approach to 3-D image retrieval from 3-D Fourier space data. They furnish as shown in this paper the ability to display true 3-D image detail. The approach is based on "Fourier Domain Projection Theorems"^{2,3} that are dual to "Spatial or Object Domain Projection Theorems" used in radio-astronomy⁷ and tomography¹. These theorems permit the reconstruction of 3-D image detail tomographically* i.e. in slices from 2-D projections of the 3-D Fourier space data^{2,3}. Although the required 2-D Fourier transform can be carried out digitally, the emphasis in this paper is on coherent optical techniques for performing the 2-D Fourier transform with particular attention to implementations that permit the execution of the necessary 2-D optical transforms of the various projection hologram sequentially in real-time. Specific attention is given to a technique that utilizes the virtual Fourier transform which permits the viewing of a virtual 3-D image in real-time.

FOURIER DOMAIN PROJECTION THEOREMS

There are two Fourier domain projection theorems. One leads to tomographic object reconstruction in parallel slcies and is called the "weighted Fourier domain projection theorem", the other leads to tomographic object reconstruction in meridonal or central slices and can therefore be called the "meridonal or central slice Fourier domain projection theorem".

We begin by considering a 3-D object function $f(\bar{r})$ with $\bar{r} = xI_x + yI_y + zI_z$ being a position vector in object space. Let $F(\bar{w})$ be the 3-D Fourier transform of $f(\bar{r})$ defined by,

$$F(\vec{w}) = \int f(\vec{r}) e^{-j\vec{w}\cdot\vec{r}} d\vec{r}$$
(1)

where $d\bar{r} = dx dy dz$ and $\bar{w} = w_x I_x + w_y I_y + w_z I_z$ is a position vector in the Fourier or spatial frequency domain.

Consider next the projection of $F(\bar{w})$ on the w_x , w_y plane defined by,

$$F_{p}(w_{x},w_{y}) = \oint_{W_{z}} F(\overline{w}) dw_{z} . \qquad (2)$$

and combining eq. (1) and (2),

$$F_{p}(w_{x},w_{y}) = \oint_{W_{z}} \{ \iint_{z} f(x,y,z) e^{-j(w_{x}x + w_{y}y + w_{z}z)} dxdydz \} dw_{z}$$
(3)

From the Greek work Tomos meaning slice.

Integrating with respect to w_z first and assuming that the volume in \bar{w} space occupied by $F(\bar{w})$ is sufficiently large we obtain,

$$F_{p}(w_{x},w_{y}) = \begin{cases} \int f(x,y,z) \delta(z) e^{-j(w_{x}x + w_{y}y)} \\ \int g(z) e^{-j(w_{x}x + w_{y}y)} \\ dxdydz \end{cases}$$
(4)

$$= \oint_{X} \int_{Y} f(x,y,o) e^{-j(w_{X}x + w_{y}y)} dxdy$$
 (5)

The 2-D Fourier domain projection $F_p(w_x, w_y)$ and the central slice f(x,y,o) through the object form thus a Fourier transform pair. This may be symbolically expressed as,

$$F_{p}(w_{x},w_{y}) \leftrightarrow f(x,y,o)$$
(6)

Other parallel slices through the object at $z = z_n$, z_n being a constant describing the z coordinate of the n-th parallel slice, can in a similar manner be related to "weighted" Fourier domain projections of $F(\bar{w})$ defined by,

$$F_{p,n}(w_x,w_y) = \oint_{w_z} F(w) e^{\int_{w_z}^{z} w_z} dw_z$$
(7)

Making use of eq. (1) and again performing the integration with respect to w, first we obtain,

$$F_{p,n}(w_x, w_y) \leftrightarrow f(x, y, z_n)$$
(8)

which indicates that the weighted projection $F_{p,n}(w_x, w_y)$ and the n-th parallel object slice $f(x, y, z_n)$ form a Fourier transform pair. Equation (6) is seen to be a special case of eq. (8) when $z_n = o$.

Given the 3-D Fourier space data manifold $F(\tilde{w})$ one can digitally compute and display a set of "weighted projection holograms" $F_{p,n}(w_x, w_y)$. A corresponding set of images of parallel slices or cross-sectional outlines of the 3-D object can then be retrieved via 2-D Fourier transform operations which can most conveniently be carried out optically from photographic transparency records of the weighted projection holograms displayed by the computer. Returning to eqs. (1) and (2) one can also show that projections of F(w) on arbitrarily oriented planes other than the w_x, w_y plane

chosen for eq. (2), yields "meridonal projection holograms" that are 2-D Fourier transforms of corresponding merdional (central) slices of the object. This is the "meridonal Fourier domain projection theorem. It furnishes the basis for angular multiplexing of the resulting meridonal projection holograms into a single composite hologram which can be used to form a 3-D iamge of the object in a

manner similar to that in integral holography⁸ which is increasingly being referred to as Cross holography*.

THE VIRTUAL FOURIER TRANSFORM

In contrast to the well known spatial Fourier transforming pro-⁹ of the convergent lens widely used in coherent optical computing, the complementary virtual Fourier transform capability of a divergent

lens¹⁰ is less widely known or used despite many attractive features. This is surprising since the power spectrum associated with the VFT is a phenomenon that is frequently observed in daily life when one happens to look at a distant point source such as a street light through a fine mesh screen or the fine fabric of transparent curtain material. The spectrum of the screen transmittance appears then as a virtual image in the plane of the point source.

The VFT concept of the divergent lens is easily derived from the Fourier transform expression of the convergent lens. Figure 1 illustrates the well known process of forming a real Fourier transform with a convergent lens. The object transparency, with complex transmittance t(x, y), is placed at a distance d in front of a convergent lens of focal length F and illuminated with a normally incidnet collimated laser beam. The complex field amplitude of the wavefield in the back focal plane, the transform plane, is given by the well known formula

$$T(x,y) = \frac{j}{\lambda F} e^{-j\frac{k}{2F} [(1 - \frac{d}{F})(x^{2} + y^{2})]}$$

$$\sum_{x \text{ ff } t(x_{0},y_{0})} \sum_{x \text{ ff } t(x_{0},y_{0})} \frac{j\frac{k}{F}(x x_{0} + y y_{0})}{dx_{0} dy_{0}}$$
(1)

in which the integral is recognized as the two dimensional Fourier transform of the object transmittance. T(x,y) becomes the exact Fourier transform of $t(x_0, y_0)$ when d = F that is when the object transparency is placed in the Front focal plane of the lens. The power spectrum associated with the transform is real and can be projected on a screen placed in the back focal plane. It is also well known that a scaled version of the transform can be obtained in the back focal plane by placing the object transparency in the converging laser beam to the right of the lens⁹.

Named after Lloyd Cross the originator of integral holography.









Noting that eq. (1) does not change when we replace d by -d, F by -F, x and y by -x and -y respectively, we can arrive at the complementary VFT arrangement illustrated in Fig. 2. An inverted transparency $t(x_0, y_0)$ is placed now in the divergent coherent beam to the right of the divergent lens (of focal length -F) and a VFT given by eq. (1) is observed in the virtual focal plane of the lens. The same VFT can be seen by removing the divergent lens and replacing the laser beam with a point source placed at the origin of the VFT plane as depicted in Fig. 3. Thus a simple way of viewing the power spectrum associated with the VFT of a given diffracting screen (which is usually a Fourier transform hologram or a projection hologram of the type described above) is to hold the screen close to the eye and look through it at a distant bright point source. The point source used need not be derived from a laser. In fact it is preferable for safety purposes to use an LED or a spectrally filtered minute white light source such as a "grain-of-wheat" subminiature incandescent lamp or a miniature Christmas tree decorating lamp covered by a color or interference filter. This has the added advantage of furnishing a measure of control over the coherence properties of the wavefield impinging on the screen providing thereby a means for reducing coherent noise in the observed VFT and also, as will be discussed below, a means for coherent or noncoherent superposition of VFT's. As the distance of the point source from the diffracting screen is decreased in order to make it compatible with typical laboratory or optical bench dimensions, the size of the observed VFT decreased because of the change in the curvature of the wavefield illuminating the diffracting screen. To compensate for this effect it is necessary to reduce the size of the diffracting screen or transparency often to such a scale where viewing the VFT throught the small available aperture becomes difficult. To overcome this limitation the displacement property of the Fourier transform can be utilized. A composite transparency containing an ordered or random array of reduced replicas of the transmittance function $t(x_0, y_0)$ arranged side

by side as illustrated in Fig. 4 is prepared. When such a composite transparency is viewed with the point source, the VFT's formed by the individual elements will overlap in the virtual Fourier plane. The VFT's are identical except for Linear phase dependence on x,y which depends in each VFT on the central position of each element in the composite transparency. This leads to a desirable noise averaging effect and the appearance of fine checkered texture in the image detail. All this leads to an enhancement of the quality of the observed power spectrum. Both coherent and noncoherent superposition of the over-lapping VFT's is possible using this scheme by varying the coherence area of the wavefield illuminating the composite transparency. When the coherence area is roughly equal to the size of the individual elements of the composite transparency noncoherent superposition results, while a coherence area equal or greater than the size of the composite transparency would yield coherent superposition.



Fig. 3. Arrangement for viewing a virtual Fourier transform with a point source



Fig. 4. A composite screen consisting of an ordered array of identical Fourier transform projection holograms.

THREE DIMENSIONAL DISPLAY

The VFT concept and the "weighted Fourier domain projection theorem" discussed above can be combined in an attractive scheme for the reconstruction and display of a 3-D image from a series of weighted projection holograms corresponding to different parallel slices through the object. The scheme is based on viewing a series of weighted projection holograms sequentially in the proper order of the occurance of their corresponding slices in the original object while displacing the point source axially for one hologram to next by an axial increment proportional to the spacings between adjacent object slices. In this fashion the reconstructed virtual images of the various slices are seen in depth at different VFT planes that are determined by the positions of the axially incremented point source. Repeated rapid execution of this procedure by displacing the point source back and forth leads the observer to see a virtual 3-D image tomographically in parallel slices or sections as he looks through the series of projection holograms passed rapidly, as in a motion picture film, infront of his eyes.

More specifically the scheme is based on preparing a series of N weighted Fourier domain projection holograms from the 3-D Fourier domain data $F(\bar{w})$ of a given object $f(\bar{r})$ as described in the preceeding sections. Each of the projection holograms would correspond to a different parallel slice through the object. A composite transparency similar to that shown in Fig. 4 is formed for each projection hologram. In fact Fig. 4 is an example of a computer generated composite hologram containing an array of identical weighted projection holograms corresponding to one slice of the test object shown in Fig. 5. The test object chosen consisted of eight point scatterers arranged as shown. The 3-D Fourier space of this test object was accessed in a computer simulation of wavelength diversity imaging as described in a companion paper in this volume*. The resulting computer generated Fourier space data manifold $F(\bar{w})$ was used to compute three weighted projection holograms corresponding to the three planes $R_2^{-1}=1m, 0, 1m$ of Fig. 5

containing the three different distributions of point scatterers. composite array such as that of Fig. 4 was formed and displayed by the computer for each of the three projection holograms, each was photographed yielding a set of three projection hologram composite transparencies. Copies of these were then mounted on a rotating wheel as shown in Fig. 6 (a) and viewed with an axially scanned point source. Four sets of transparency copies of these three composite projection holograms were mounted in the order 1,2,3,2,1,2 ... on the periphery of a rotating wheel as shown in Fig. 6 (a). The wheel is driven by a computer controlled stepper motor. The axially scanned point source was produced by scanning a focused laser beam back and forth on a length of fine nylon thread with the aid of a deflecting mirror mounted on the shaft of a second computer controlled stepper motor as shown in Fig. 6 (b). The laser and optical bench arrangement for forming the scanned focused beam appear in the background of Fig. 6 (a). The computer controlled steppers enable precise positioning of the secondary point source on the scattering thread in synchronism with the hologram

*See paper entitled "Holography, Wavelength Diversity Inverse Scattering" in this volume.



Fig. 5. A three-dimensional test object consisting of a set of eight point scatterers shown in isometric and $R'_x-R'_x$ plane views at $R'_z-z_0, 0, z_0$. $x_0=y_0=z_0=100$ cm.





(b)

Fig. 6. Quasi real-time three-dimensional image reconstruction and tomographic display in successive slices from a series of projection holograms mounted on rotating wheel seen in fore-front of (a); Detail of laser scanner used to produce linearly scanned point source is shown in (b).



being viewed so that the VFTs are formed in their proper planes. A viewer looking at the axially displaced point source through each transparency mounted on the wheel as it passes infront of his eye will see a 3-D virtual image. Photographs of the three virtual images seen by an observer in this fashion are shown in Fig. 7. An optodigital scheme for rapid real-time implementation of the procedure realized above is shown in Fig. 8. This scheme, presently under study, utilizes a rapid recyclable spatial light modulator (SLM) such as the Itek PROM in order to form VFT's of the projection holograms displayed by the computer in real-time.

CONCLUSIONS

We have presented the basic principles of tomographic 3-D image display based on Fourier domain projection theorems. One possible implementation of the principle using the virtual Fourier transform and a series Fourier domain projection holograms has been described. There are several advantages for using the VFT rather than the real Fourier transform (RFT), the most important of which is the ease with which the position of the VFT plane can be moved axially by simply moving the position of the reconstruction point source. The VFT approach was adopted in the present study because it is much easier to move a point source rapidly than to move the display screen needed in the RFT approach. Furthermore focusing in the VFT approach is carried out by the observer while in the RFT approach it must be performed by the system. Other attractive features of the VFT are:

(a) Simplicity - enables direct viewing of the power spectrum of a transparency or a hologram with a variety of simple point sources.

(b) The scale of the observed VFT can be easily altered by changing the distance between the projection hologram transparency and the reconstruction point source.

(c) Lower speckle noise and therefore higher reconstructed image quality can be attained by using nonlaser point sources in the reconstruction such as LED or miniature spectrally filtered incandescent lamps. Further reduction in speckle noise occurs when an array of the projection hologram rather than a single hologram is used and when the hologram is slightly vibrated or is in motion because of a noise averaging effect.

(d) Coherent and noncoherent superposition of VFT's is possible by altering the coherence area of the reconstruction wavelfield.

(e) Because of the Fourier transform nature of the projection holograms utilized, the resolution requirements from the storage medium (photographic film or the CRT/SLM system of Fig. 8) are much lower than would be needed in the recording of a Fresnel hologram of the object as a means of 3-D image display. The 3-D image detail contained in the single Fresnel hologram is now distributed over a series of lower resolution projection holograms which are used to form the 3-D image sequentially in time in individual slices.



(f) Because 3-D image reconstruction is tomographic (in separate slices) there is no interference between the wavefields forming the various slices.

(g) Permits other forms of 3-D image display involving spatial or angular multiplexing in a fashion similar to integral holography.

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APPENDIX IV

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AN AUTOMATED FREQUENCY RESPONSE AND RADAR CROSS-SECTION MEASUREMENT FACILITY FOR MICROWAVE IMAGING

UNIVERSITY OF PENNSYLVANIA

THE MOORE SCHOOL OF ELECTRICAL ENGINEERING

AN AUTOMATED FREQUENCY RESPONSE AND RADAR CROSS-SECTION MEASUREMENT FACILITY FOR MICROWAVE IMAGING

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Presented to the faculty of the Moore School of Electrical Engineering (Department of Electrical Engineering & Science) in partial fulfillment of the requirements for the degree of Master of Science in Engineering.

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May 1980

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THE MOORE SCHOOL OF ELECTRICAL ENGINEERING

AN AUTOMATED FREQUENCY RESPONSE AND RADAR CROSS-SECTION

MEASUREMENT FACILITY FOR MICROWAVE IMAGING

ABSTRACT

This thesis investigates the development of a broadband microwave holographic imaging facility. Different methods for the correction of microwave target scatter data are discussed and implemented. A minicomputer automates all system functions including data acquisition, storage, calculation, and graphic display. The effects of range phase shift on holographic frequency diversity imaging is considered and techniques for the removal of this phase shift in a laboratory environment. The frequency dependent backscatter of several test targets is derived analytically and simulations done of the corresponding holograms. These holograms are compared to those measured experimentally. Finally both simulated and experimental holograms are optically reconstructed to yield target images using optical Fourier transforms and shown to be in excellent agreement.

Degree: Master of Science in Engineering for graduate work in electrical Engineering and Science

Date: May 1980

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Author

ACKNOWLEDGEMENTS

I would like to extend my most sincere thanks to Dr. Nabil H. Farhat for making my studies at the University of Pennsylvania possible. His patience and wisdom were an inspiration for my research. I would also like to thank my parents for their love and understanding these many years. Finally I want to express appreciation to the Moore School for giving me the chance to study here.

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I INTRODUCTION

Frequency diversity imaging has been under study at the Electro-Optics and Microwave Optics laboratory of the Moore School Graduate Research Center.[1],[2],[3],[4] This study has established the theoretic feasibility of imaging objects by means of their multiaspect frequency response . For the purpose of experimentally studying frequency diversity imaging, an experimental measurement system has been assembled and installed in the Moore School anechoic chamber. In this thesis we will describe the automation of this measurement system and characterize its peroformance in the measurement of complex field amplitudes of scattered fields. A system block diagram fig.1.1, shows the major system components. The central element is the DEC MINC LSI 11/2 minicomputer. This computer performs several important functions. The MINC controls laboratory instrumentation via the IEEE-488 bus protocol standard. This allows the Hewlett-Packard 8620C microwave sweeper to be precisely tuned to any frequency in the 2.0 to 18.0 GHz range (fig.1.2). The computer collects data from the HP 8410B network analyzer through the four analog input channels available. These analog values are proportional to the amplitude in db and phase in the range $(-\pi)$ to $(+\pi)$ radians relative to the reference signal supplied to the network analyzer. The computer stores Experimental data on floppy The available storage capacity is large, over discs. 500,000 measurement pairs of complex field amplitude and



Fig. 1.1 System block diagram.



Fig. 1.2 HP 8620C microwave sweeper and HP 1410B network analyzer.

phase may be stored on a single disc. This data can be accessed for both processing and display on a Tektronix 606A high resolution CRT monitor. The disc system allows the MINC to operate under a sophisticated software system, DEC RT-11 V3.0B permitting programing in MINC BASIC,FORTRAN IV, and MACRO languages.

The processing capability of the system allows the removal of system response errors due to anechoic chamber clutter, antenna cross-coupling and receiver channel characteristics. The data may be processed for target range calculations and the removal of the phase shift due to the target range. In addition the collected data may be filtered to improve imaging and finally displayed on a high resolution Tektronix 606A X-Y CRT monitor using the MINC system D/A converter module.

The second part of this thesis will describe the simulation and actual operation of a holographic radar system verifying the theory of frequency diversity imaging. This is done using the system described in the first section. The scattering of various targets is derived and holograms using these results are generated for comparison with experimental data. Finally a system for the experimental measurement of the scattering for these targets is outlined and the results from this system compared to theory.

II SYSTEM OPERATION AND ERROR REMOVAL

This chapter will cover the operation of the microwave backscatter data aquisition system. This will include actual interfacing information and an analysis of the types of errors encountered when making microwave measurements. The error correction techniques developed are later used in the experimental verification of the frequency diversity imaging theory.[4]

2.1 Microwave sweeper operation

The first task in the development of an automated data acquisition system is the implementation of a data communications link between the intelligent controller and instrumentation. The IEEE~488 bus protocol is utilized in this application for the transfer of data to the Hewlett-Packard 8620C microwave sweeper from the MINC LSI 11/2 computer. This bus is a high speed 8 bit wide bidirectional data path with 5 additional lines dedicated to control. Data is transferred in ASCII format over the bus. For example the number 1 is transmitted as the ASCI code for the character '1'. Certain sequences of characters make the sweeper perform different functions or enter different modes of operation via its IEEE-488 interface.

In order to set the frequency of the sweeper a number must be sent to the IEEE-488 interface in the range 1-10000. Each sweeper frequency band has been split into 10000 frequency points. The frequency of operation is controlled

by an internal analog voltage that varies between 0.0 and 10.0 volts. A D/A converter on the 8620C IEEE-488 interface changes the data transmitted from the MINC into the frequency controlling voltage. The correspondence between voltage and output frequency is essentially linear. In order to obtain the interpolating function for frequency versus control voltage; a microwave frequency counter was the voltage-frequency used to measure characteristic function. Using the MINC-BASIC program CALAB.BAS a least squares fit for both linear and quadratic functions was made on the frequency vs voltage data. This type of program is used for determining the best polynomial fit to the frequency-voltage characteristic of the sweeper. Sample output and program listings are in appendix I along with an explanation of program operation. The results from this work indicate that the quadratic fit was statistically superior for all bands on the microwave sweeper. The FORTRAN subroutine SWEEP was written which utilizes the quadratic interpolation polynomial for each of the four bands of the 8620C. This subroutine automatically calculates the control voltage and band to generate any the 2-18 GHz frequency in range and transmits the appropriate commands over the IEEE-488 bus. The variance of the frequency setting using the quadratic fit for the three bands are as follows: .6 MHz in the 2.0-6.3 GHz band; 1.2 MHz in the 6.3-12.0 GHZ band; and 1.6 MHZ in the 12.0-18.0 GHz band. For higher accuracy the sweeper may be phase

locked to the reference in a locking frequency counter yielding very high accuracy as precise as the frequency reference itself. The EIP 371 locking counter may be used in this application to lock the HP 8620C sweep oscillator to the correct frequency once it is within 20 MHz of the desired frequency. The auxiliary output of the HP 8620C supplies a sample of the signal generated by the fundamental 2.0-6.3 GHz oscillator module within the sweeper to the locking counter. Sweeper output on higher bands is this fundamental multiplied by a factor of 2 or 3 for the 6.3-12.0 and 12.0-18.00 GHz bands respectively. For example the locking frequency for a 10.0 GHz sweeper output on band 2 would be 5.0 GHZ. In operation, subroutine SWEEP will set the frequency of the sweeper within 2.0 MHz of the desired frequency and subroutine SLOCK will be called to calculate the locking frequency; lock the HP 8620C and return to the main calling routine when lock will have occurred. Lock time varies from .1 to 3 seconds and resolution is 100 KHZ. These subroutines are called whenever the sweeper must be set to a particular frequency or the sweeper must be placed in or be released from computer control.

2.2 Network analyzer-computer interface

The Hewlett-Packard 8410B network analyzer is the focus of the measurement capabilities of the microwave measurement system. It can make vector(amplitude and phase) measurements in the (.1-18.0)GHz range. The range of amplitude measurement is 80 db and phase may be measured modulo (2π) . The system reference signal is fed from a 20db directional coupler to the HP8411A harmonic converter sampling head of the network analyzer. This reference is compared to the backscatter from the illuminated target; both in amplitude and phase. The reference signal amplitude is kept constant by leveling the sweeper with a feedback signal derived from its amplified output by means of a crystal detector. This allows the sweeper-TWT (Traveling Wave Tube) system to yield nearly constant output in the (2.0-16.5) GHz range; see fig.2.1. TWT power output is on the order of 1 watt over these frequencies.

The complex field amplitude measurements are available as analog voltages from the back panel of the 8410B. The outputs are proportional to amplitude and phase: 25 MV/db and 10 MV/DEGREE These values are digitized by the MINC using its built in analog to digital conversion channels. The MINC A/D converters digitize voltages lying in the range of -5.12 to +5.12 volts to the range of 0-4096 yielding 12 bit resolution. If the signal is corrupted by noise; the user has the option of employing signal averaging to cancel the effects of noise uncorrelated to the received signal.

A difficulty encountered when measuring phase angle modulo (2π) occurs when the phase is close to $(+\pi)$ or $(-\pi)$. At this point a small change in the signal phase may cause the phase to flip between these two equivalent extremes rapidly. If a data sample is taken close to $(+\pi)$ or $(-\pi)$ it may be in transition between them and therefore incorrect.

Since such points occur infrequently in a typical measurement they may be ignored. ; their presence does not seriously hinder any holographic imaging due to the inherent redundancy and therefore noise immunity of the holographic reconstruction process.[5] If it is desired that these points be identified and removed it is necessary to estimate the mean and variance of the samples taken at each frequency point. At frequencies where the phase is flipping between $(+\pi)$ and $(-\pi)$, the variance will be much larger than at other frequencies. The mean of the samples when this is occurring will be near zero. Hence to resolve the ambiguity problem it is necessary to decide if the variance exceeds a predetermined threshold when the mean in the neighborhood of Two FORTRAN subroutines PHAMP and PHAMP2 which zero. implement these algorithms are listed in appendix I.

Since such points occur infrequently in a typical measurement they may be ignored. ; their presence does not seriously hinder any holographic imaging due to the inherent redundancy and therefore noise immunity of the holographic reconstruction process.[5] If it is desired that these points be identified and removed it is necessary to estimate the mean and variance of the samples taken at each frequency point. At frequencies where the phase is flipping between $(+\pi)$ and $(-\pi)$, the variance will be much larger than at other frequencies. The mean of the samples when this is occurring will be near zero. Hence to resolve the ambiguity problem it is necessary to decide if the variance exceeds a predetermined threshold when the mean in the neighborhood of Two FORTRAN subroutines PHAMP and PHAMP2 which zero. implement these algorithms are listed in appendix I.
2.3 Implementation of the data aquisition system

In an experimental environment it is important to be aware of the various types of errors inherent in the equipment and the experimental procedure adopted. Conditions and equipment always vary from theoretical ideals. A clear understanding of the error removal process leads to the development of practical implementations of theoretical concepts and enhancement of measurement accuracy unattainable otherwise.

Errors in complex field amplitude measurement may be caused by several factors. These may be grouped into two categories. Errors caused by the instruments themselves fall into the first class. Such factors as measurement variations caused by electronic noise, ,inaccurate A/D and logarithmic conversions, and inaccuracy and instability in the microwave source make up this category. The second group of errors is caused by the test set, antennas , cables, connectors, amplifier, and room clutter. All these factors interact with each other in the microwave region and are the significant cause of error in microwave measurements. Little can be done about the first class of errors since they are inherent in the characteristics of the equipment used. The second group of errors can be removed through the use of automated measurement of system parameters in the frequency range of interest. These errors can be removed form any measurements of scattering objects by digital processing and the results stored for later recall. This

essentially provides an automated and improved version of the conventional two antenna radar cross-section measurement technique [6]; in which a microwave bridge is balanced in the absence of the target and the degree of imbalance is measured when the target is introduced into the microwave field.

Let us look at the first class of errors more closely since these errors will set the ultimate performance limits on the system. The characteristics of the signal source are important in this regard. In this case the signal source is a Hewlett-Packard 8620C microwave sweeper. Since the sweeper is not phase locked frequency and stability problems exist. The carrier also has significant FM noise which appears as phase noise in the scattered signal. The phase shift of the scattered scattered signal as a function of frequency for small frequency variations is given by:

$$\Delta \Theta = \Delta k \cdot R \tag{2.1}$$

where $k=(2\pi/\lambda)$ and R is the path length. For R greater than a few tens of meters ,the FM noise on the signal source causes measurable variations in the phase of the scattered signal. The stability of a synthesizer is required for the implementation of a holographic radar system when target ranges are in terms of kilometers.

Another limit on the ultimate accuracy of the system is the resolution of the A/D conversions and the accuracy of the network analyzer. Given the 2.2MV resolution of the MINC A/D converter and the network analyzer analog output of

50 MV/db, the system can resolve .048 db steps. This resolution limit restricts the minimum signal to system error ratio. If the error signal consisting of clutter, antenna coupling, system directivity and noise exceeds the scattered target signal the resolution of the target signal suffers. For example, if the system error signal and target scattered signal are of equal intensity, then a 1 db change in the scattered signal causes a .53 db change in the total received signal vector. When the system error is 13 db above the scattered signal it becomes impossible to resolve a 1 db change in the target signal given the resolution capabilities of the system. This difficulty is further compounded by errors in the network analyzer. These too are amplified when the clutter exceeds the target scatter In fig.2.1 is a plot of the minimum system signal. resolution in order to detect a 1 db change in the target return signal versus the system error signal to scattered signal level in db. In fig.2.2 is plotted the target signal resolution versus the noise to signal level in db. When the system error is 20 db below the target signal then the resolution of the target scatter signal is very close to the ultimate resolution, .048db. Clutter is the component of the received signal not scattered by the target but that signal that is the result of coupling between antennas and signal scattered by the anechoic chamber walls. As the clutter/signal ratio increases the target scatter signal resolution decreases exponentially. Clutter may be reduced



Fig. 2.1 The required resolution in db for a data acquisition system to detect a 1 db change in the desired signal vs. signal/error -signal ratio.



Fig. 2.2 Minimum change in desired signal level detectable versus signal/error signal level given .048 db system data acquisiton resolution.

by improving the isolation between the transmitter and receiving antennas with the introduction of absorbing foam panels such as Emerson and Cummings' Ecosorb panels between the antennas.

2.4 System error correction

Turning next to the second class of errors; those directly measurable and therefore removable; define the following quantities which are functions of frequency :

C(f) -- Corrected antenna clutter and coupling

R1(f)-- uncorrected reference target backscatter

R2(f) -- Corrected reference target data

(uncorrected for antenna system response)

R(f) -- Corrected reference target data Several possible techniques exist for the removal of system errors. The particular technique is dependent on the relative signal levels involved, the accuracy desired, ease of implementation, and computational speed. The first technique described here is similar to that used by Weir et al.[7]

The first step in the correction procedure measures the transfer function of an attenuator A(f) as a function of frequency. The equipment setup for this procedure is shown in fig. 2.3. The two ports of the reflection-transmission unit connected through a precision HP 11605A flexible coaxial arm. The MINC then steps the sweeper to a number of frequency points and stores the system response (log amplitude and phase vs. frequency) in memory. When this completed the attenuator is placed in series with the arm and another set of measurements is made at the same frequency points before. system as the response characteristic is subtracted from the combined attenuator plus system response measurement made on the second sweep. An example of this procedure is shown in fig 2.4. Computer subroutine PAD performs this operation. It is listed in appendix I along with all other computer program listings and output pertaining to system response measurement and removal.

The next step in the calibration process is measurement of the reference to test channel isolation I(f). This characteristic is dependent on the directivity of the network analyzer harmonic converter, and the reflection transmission unit. For the Hewlett-Packard 8411A harmonic converter the isolation is greater that 50 db. In this measurement the ports of the reflection-transmission unit are terminated in the cables used for the later target scatter measurement. These coaxial cables are terminated in



Fig. 2.3 Equipment setup for attenuator measurement.



Fig. 2.4 30 db attenuator characteristic for 8.0-8.5 GHz.

50 ohm resistive loads as shown in fig.2.5. The results of a typical run are shown in fig.2.6. Computer subroutine IST automates this stage in the correction procedure. As can be seen the coupled signal is well in the noise of the system and would not affect later scatter data. If the isolation effect is ignored later calculation would be simplified greatly; but is included here to be consistent with the procedure outlined in the literature.[7]

The next stage in generating the data for correction of scatter data is measurement of the system transfer function T(f). This consists of the characteristics of the traveling wave tube amplifier, system cables and connectors. This is done by connecting the cable from the transmitting antenna to the receiving antenna cable in fig.2.6 and placing the 30 db attenuator characterized previously in the line to avoid damage to the harmonic converter. The raw measurement MT(f) is a combination of several factors:

$$MT(f) = T(f) * A(f) + I(f)$$
(2.2)

Solving for T(f):

$$T(f) = (MT(f) - I(f)) / A(f)$$
 (2.2a)

The equipment setup for this procedure is shown in fig 2.7 and typical uncorrected and corrected transfer function data in fig.2.8 and 2.9.

Measurement of the antenna cross-coupling and room clutter is the next step in the correction process. In this procedure shown in fig.2.7b, the target is removed from the anechoic chamber and the antennas pointed to the target



Fig. 2.5 System configuration for measuring isolation between reference and unknown ports.









(Ь)

Fig. 2.7 a) System configuration for transfer characteristic measurement. b) Connection to antennas for scattering and clutter measurement. c) Transmitting antenna with spiral AEL antenna and parabolic dish. d) Receiving dual polarization horn; EMI A6100 with Norsal 90 hybrid.



(C)



Fig. 2.7 (contd.) a) system configuration for transfer characteristic measurement. b) Connection to antennas for scattering and clutter measurement c) Transmitting antenna with spiral AEL antenna and parabolic dish. d) Receiving dual polariztion horn; EMI A6100 with Norsal 90 hybrid.



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Fig. 2.8 Uncorrected transfer characteristic of system; 8.0-8.5 GHz.



Fig. 2.9 Transfer characteristic corrected for attenuator response; 8.0-8.5 GHz.

location. A high gain parabolic dish antenna is used for illumination of the target since the narrow beam pattern of the antenna places most of the radiated power on the target area. A smaller dual polarization horn is used for receiving in order to sample a small area of the scattered field. The uncorrected clutter C1(f) is given by:

Cl(f) = C2(f) * T(f) + I(f) (2.3)

Solving for the corrected clutter and coupling:

$$C2(f) = (C1(f) - I(f)) / T(f)$$
 (2.3a)

Subroutine ANTEN does the system clutter and coupling removal. Examples of the uncorrected and corrected clutter Cl(f) and C2(f) are shown in figs.2.10 and 2.11. The corrected clutter represents the actual signal reflected from the anechoic chamber walls and that signal coupled between the antennas with the system response removed.

These subroutines: PAD,IST,TRANS, and ANTEN, were combined into a program SYSRES. Data from each of these subroutines may be stored on disc for later recall or display. Theoretically if the system is not disturbed then the system response will remain constant. Then only the target data need be recorded in any run for a new corrected backscatter measurement.

The transfer function of the system and the range clutter- antenna cross coupling data are utilized the the next step of the error correction process. A reference object of known constant cross section is measured and the result stored. This data includes all the errors previously



Fig. 2.10 Uncorrected system clutter; 8.0-8.5 GHz.



Fig. 2.11 System clutter corrected for transfer function; 8.0-8.5 GHz.

described but also takes into account the antenna system variations as a function of frequency:

Rl(f) = C2(f) * T(f) + I(f) + R2(f) * T(f) (2.4)

Here C2(f) and R2(f) contain the antenna system response multiplying the actual values of corrected reference target and system clutter data.

$$R2(f) = R(f) * A1(f)$$
 (2.5)

C2(f) = C(f) * A1(f) (2.6)

This response Al(f) takes into account the varying amount of power received and transmitted as a function of frequency in the antenna system. If the reference target is chosen to have a constant cross section and linear phase over the frequency range of interest then R(f) is of constant amplitude and linear phase. Solving for R2(f):

R2(f) = (Rl(f) - I(f) - Cl(f)) / T(f)This leaves R2(f) proportional to Al(f) shifted by a linear phase corresponding to the reference target range.
(2.7)

When the actual target is measured; it is corrected for system errors as was the reference target data and this result is divided the corrected reference target data to yield the final target scatter data.

$$S2(f) = (S1(f) - I(f) - C1(f)) / T(f)$$
 (2.8)

S(f) = S2(f) / R2(f)

(2.9)

This technique may be simplified considerably in the laboratory environment given the signal to noise ratio is greater than 10 db for the scattered signals. In this case only multiplicative errors remain and the additive errors are masked by the high target scatter signal amplitude. Hence:

and therefore:

Rl(f) = R2(f) * T(f)

$$S(f) = S1(f) / R1(f)$$
 (2.11)

Weir and his group have reported that this technique has reduced equivalent range clutter to -45 db below 1 sq. meter.

There remains one source of error in the scattered signal measurement that cannot be removed by calculation. This error comes from multipath scattering from the object. The target scatters power in all directions. Some of this signal may be reflected off the walls or floor of the anechoic chamber. The signals reflected off the walls is over 48 db down from the incident wave amplitude in the 6.0-12.0 GHz range. However if the target is small in cross section; on the order of 100 sq. cm.; then it is possible for the walls (on the order of 100 000 sq. cm.) to contribute a significant component to the received signal.

In order to analyze the effect of the multipath scattering, let the directly received signal be written:

$$S_{p}(t) = A \cos(\omega t + \Psi) \qquad (2.12)$$

and the indirectly received signal:

$$S_{i}(t) = B \cos(\omega t + \Psi + \Theta) \qquad (2.13)$$

Then the total received signal is given by: $S_{\gamma}(t) = (A^{2} + 2AB\cos + B^{2})\cos(\omega t + \Psi + t_{exc}) \left(\frac{-B\sin \theta}{A + B\sin \theta}\right)$ (2.14) Since theta is a function of the indirect path length differences and frequency, it will lead to a periodic variation in the amplitude of the scattered signal. As an example; if the paths differ in length by 1 meter then the amplitude oscillations will occur every 300 MHz given all other factors remain constant.

A series of programs was written to test these various techniques of error removal. They differ in the only in the error removal technique employed; not in file storage or display formats nor in range calculation and removal to be described. These programs used together:

> 1)Measure and correct the reference target data with files of transfer function and clutter data generated by SYSRES.

> 2)Measure and correct the .target data and finally take the corrected reference target data and remove the antenna system response from the target data. 3)Alternately for high SNR; calculate the correct scattered target signal directly using the reference target data.

> 4)Calculate and remove phase shift due to range from the target signal after error correction

These programs are briefly described and differ in the mentioned categories:

SPHERE-Reads transfer function and clutter data files generated by SYSRES; takes the reference or object data and corrects for errors in the

equipment.

SPHER2-Measures clutter and reference target data. It then subtracts clutter and divides the target data by the reference target transfer function. SPHER3-Measures reference target and object signals ignoring clutter, and divides the object data by the reference target complex field amplitude data.

In order to implement the complete error correction process program SPHERE would be run twice; once for the reference target and then again for the test object. This data would be stored for later retrieval. Program SPHER3 would read these files and process then such that the test object complex field data would be divided by the reference target data. For high SNR cases; program SPHER3 alone would be run: first measuring the reference target and then the test target and finally dividing them yielding the final result. Listings and a more complete description of program operation is given in the program appendix I.

III RANGE PHASE SHIFT ANALYSIS

This chapter contains an analysis of the effects on the phase shift due to target range on coherent imaging of a target. It includes a discussion of some of the available techniques for the removal of this phase shift and the required perfor mance of these systems based on bandwidth and signal to noise ratio.

3.1 The effects of range phase on imaging

As previously described by Farhat and Chan ;[4]; the scattered field of the scattering target is given by:

$$\mathcal{W}(R, R_{\tau}, R_{\tau}) = \frac{j k}{2 \pi (R_{\tau} + R_{\tau})} e^{-j k (\vec{R}_{\tau} + \vec{R}_{\tau})} \int_{-\infty}^{\infty} \mathcal{U}(\vec{r}) e^{-j k (\vec{R}_{\tau} + \vec{R}_{\tau}) \cdot \vec{r}} d\vec{r} \qquad (3.1)$$

Where $\mathbf{\bar{R}}_r$ and $\mathbf{\bar{R}}_r$ are the vectors from the receiving and transmitting antennas to the target respectively and k is the wave vector of the illuminating wave. In this case the integration is over the extent of the object. The integral term is independent of the range of the target. While the term preceeding the integral is target independent and contains range information. The argument of the complex exponential is a linear phase function of frequency.

In the single receiver=transmitter pair arrangement; shown in fig.3.1, the spatial frequency domain (\vec{p} space) data is collected in a plane perpendicular to the axis of rotation. In this case the multiplication of the range phase factor leads to the convolution of the transforms of the the range and target functions in the spatial domain. The real part of the range frequency domain function is a





Fig. 2.12 a) Target and antenna placement relative to target considered in analysis. b) Experimental configuration in anechoic chamber; note target on rotating pedestal in forefront.

radially symmetric sinusoidal function:

$$R\left\{e^{-j \cdot h_{r}(R_{T}+R_{r})}\right\} = \cos -h_{r}(R_{r}+R_{r})$$
(3.2)

This transforms to a circular ring of radius $(R_r + R)$. This indicates that each point of the reconstructed image is convolved with this circular ring pattern; seriously degrading the target image. In fact any error in the removal of phase will distort the image in this manner. The effect of removing the range phase factor is equivalent to the focusing of the system on the target.

There are several other reasons for the removal of the factor:

$$\frac{jk}{2\pi(R_{\tau}+R_{\tau})} e^{-jk(R_{\tau}+R_{\tau})}$$
(3.3)

from the received data. If the data is discrete then the considerations of aliasing and sufficient data sampling rate are introduced. When the sampling rate in the frequency domain is (f) then the maximum target range before aliasing will occur is (c/4* f). However if the phase factor is removed and the target is of smaller dimension L than the range, then the sampling rate in the frequency domain need only be sufficient to prevent aliasing over the dimensions of the object :

$$\Delta f \stackrel{e}{=} \frac{C}{4L} \tag{3.4}$$

This allows the entire resolution capability of the system to be placed on the target itself greatly reducing the data

volume.

Another reason for the removal of this phase factor term can be seen in systems involving multiple receiver-transmitter pairs. For each pair (\vec{R}_r) and (\vec{R}_r) is different and in order for the data to be coherent all phase centers must be equal for an image to be formed.

Several methods have been suggested for the removal of the range phase factor. [4] In all cases it is required that the range removal technique be accurate to within $(\lambda/5)$ for there not to be serious image degradation. It is important to investigate the constraints on the ranging system parameters necessary to attain the required accuracy.

From the scaling theorem in Fourier analysis; a signal cannot be both of narrow bandwidth and short duration.[8]

$$af(at) \longrightarrow F(\frac{a}{2})$$
 (3.5)

We define the duration and the bandwidth of the signal in the following manner: []

$$(\Delta t)^{2} = \frac{1}{E} \int_{-\infty}^{\infty} t^{2} |f(t)|^{2} dt \quad (a) \qquad (\Delta \omega)^{2} = \frac{1}{2\pi E} \int_{-\infty}^{\infty} \omega^{2} |F(\omega)|^{2} d\omega \quad (b) \quad (3.6 \, e_{,b})$$

where:

$$E = \int_{-\infty}^{\infty} |f(t)|^{2} dt = \frac{1}{2\pi} \int_{-\infty}^{\infty} |F(\omega)|^{2} d\omega \qquad (3.6c)$$

is the signal energy. Then if:

$$\lim_{|t|\to\infty} \sqrt{t} f(t) = 0 \tag{3.6d}$$

When applyed to signals scattered by a target this may be translated to range uncertainty:

 $\Delta R \Delta f \ge \frac{c}{4\pi} \tag{3.7}$

This represents a limit when the durations of the signal pairs are as previously defined. However it may be possible to improve the range resolution given that the signal to noise ratio is greater than 10 db.

Consider a sinusoid in narrow band gaussian noise:

$$r(t) = A \cos \omega t + n_{c}(t) \cos \omega t - n_{s}(t) \sin \omega t$$

$$n(t) = n_{c}(t) \cos \omega t - n_{s}(t) \sin \omega t$$
(3.8)

where n(t) is the noise signal and ns(t) and nc(t) are the quadrature components of the noise signal. The amplitude and phase may be expressed as:[9]

$$|r(t)| = 1/(A + n_{c}(t))^{2} + n_{s}(t)^{2}$$
 (3.9a)

$$Arg(r(t)) = tan' \left\{ \frac{n_{\delta}(t)}{A + n_{c}(t)} \right\}$$
(3.96)

Since the noise is gaussian nc(t) and ns(t) are gaussian. Assuming that the SNR is high the phase may be approximated by:

$$Arg(r(t)) = \tan^{-1}\left\{\frac{n_s(t)}{A}\right\} \simeq \frac{n_s(t)}{A}$$
(3.10)

Also since the noise is white $n_{\xi}(t)$ may be related to the bandwidth of the receiving system.

$$\sigma_{n_{s}} = 2N_{s}B \qquad (3.1)$$

$$\int (\phi) = \sqrt{2N_0} \frac{A}{\pi B} e = N(0, \frac{\sigma_{n_s}^{L}}{A^{L}}) \qquad (3.12)$$

This leads to an interesting result; the variance of the 31.

phase is the inverse of the SNR:

$$\sigma_{0}^{1} = \frac{\sigma_{n_{3}}}{A^{1}} = \frac{1}{SNR}$$
(3.13)

Translating this result to an uncertainty relationship the phase uncertainty may be expressed:

$$\Delta \omega \Delta t \ge \sqrt{\frac{8NB}{A}}$$
 (3.14)

The uncertainty in phase is the product of the time and frequency uncertainties. This leads to an expression for the range resolution as a function of the SNR.

$$\Delta f \Delta R \ge \frac{C}{2\pi V SNR}$$
 (3.15)

Wide band ranging systems measure range by calculating the propagation delay of the signal. In one such system a high speed code is transmitted. The received signal is correlated with the original coded signal in a delay locked loop. The value of the control signal in the loop is proportional to the target range. Obviously the resolution is only as good as the the period of one of the code bits(chips). Other wide band systems use other signals for ranging(chirps,walsh functions) to obtain high range resolution.[/e],["],[12]

The system implemented here at the Graduate Research center of the Moore school is a coherent amplitude-phase measurement facility. Ranging techniques that may be integrated into this system are therefore of special interest. The phase of a point scatterer is a linear function the frequency. This directly corresponds to the

phase factor preceeding the scattering integral in eq (). If the target consists of multiple scattering centers; then each of these will be represented as a delta function in the in the reconstructed image. This suggests a technique using Fourier analysis to determine range. First place the reference target in the microwave field and measure the phase/amplitude response over as wide a frequency range as Then inverse transform this one dimensional possible. collection of data. This will transform to essentially a delta function occurring at the time corresponding to the propagation delay. This will occur when the target has a single scattering center such as a sphere. Even when the object is more complex the inverse transform will be centered around the transit time. In this Fourier technique for range determination the resolution (Δx) is inversely related to the frequency sweep width.

$$\Delta x = \frac{C}{2 \Delta f} \tag{3.16}$$

The factor of 1/2 results from the fact that the range is half the signal propagation path length.

A factor to be considered is that different scattering centers are visible from different receiver positions. It is imperitive that all when several receivers are used simultaneously that all choose the same scattering center at the phase center for range removal. If the various procenters do not coincide, the fringes of the hologram be skewed. The data sets from each of the receiver brought into alignment using an adaptive



Information is lost in the hologram if the phase centers for the scan lines are separated.

When this process is automated the sweep is done by measuring the response at discrete frequency points. The amplitude and phase are stored at N frequency points in the sweep range. The range at which aliasing will occur is given by:

$$R_{alias} = \frac{C}{4 \Delta f} \qquad (3.17)$$

This system exhibits processing gain ; a quality of all systems which spread a baseband signal into a wide spectrum. For this system the processing gain is a function of the number of measurements and the sweep width.

$$G_{ain} = \frac{\Delta f}{f_{sLepsize}} = n$$

$$G_{ain} db = 10 \log_{10} n$$
(3.18)

As an example, for 256 measurement points, this would give 24 db of processing gain.

A set of subroutines was written to test this range removal technique. In one of them, RANGE, the range of the strongest scattering center is calculated and in the other ,RANCOR, the phase factor is calculated and removed from the target data. These subroutines are used in programs SPHERE,SPHER2 , and, SPHER3.

The time domain equivalent of this technique is fitting a linear trend to the phase signal from the network analyzer. Since the phase signal is modulo (2w); this means estimating the frequency of a ramp waveform; either using a least squares approximation or implementing the equivalent of a phase locked loop. The slope of the ramp waveform is proportional to the range of the target. The accuracy to which the range may be determined depends on the sweep width, the target structure and the noise in the system. If the sweep is wide, then there will be more data with which to estimate the slope. Noise in the data will obviously interfere with the estimation process as will as any phase shifts due to the target structure.

Another type of system for generating a reference signal utilizes a Target Derived Reference. This system has been extensively studied at the Electro-Optics and Microwave Optics laboratory. [13] A brief review of the ideas developed to date in this regard are given below. The complex exponential in the integral term of scattered signal remains constant when the target is small relative to the illuminating signal wavelength.

$$\Psi(\vec{p}) = \int \mathcal{U}(\vec{r}) e^{-j\vec{p}\cdot\vec{r}} d\vec{r} \simeq e^{-j\vec{p}\cdot\vec{r}} \int \mathcal{U}(\vec{r}) d\vec{r} \qquad (3.19)$$

This is true if (L_{w_i}) is sufficiently small such that $(\not p \cdot \not r <<)$. In this case the integral value approaches a constant multiplied by a linear phase; i.e. a point scatterer. It will only occur when the target dimensions are less than a tenth wavelength of the illuminating signal, placing it in the Rayleigh scattering region. The TDR signal 15 mixed harmonically with a phase locked scattered

estimating the frequency of a ramp waveform; either using a least squares approximation or implementing the equivalent of a phase locked loop. The slope of the ramp waveform is proportional to the range of the target. The accuracy to which the range may be determined depends on the sweep width, the target structure and the noise in the system. If the sweep is wide, then there will be more data with which to estimate the slope. Noise in the data will obviously interfere with the estimation process as will as any phase shifts due to the target structure.

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$$\Psi(\vec{p}) = \int_{-\infty}^{\infty} \mathcal{U}(\vec{r}) e^{-j\vec{p}\cdot\vec{r}} d\vec{r} \simeq e^{-j\vec{p}\cdot\vec{r}} \int_{-\infty}^{\infty} \mathcal{U}(\vec{r}) d\vec{r} \qquad (3.19)$$

This is true if (L_{w_i}) is sufficiently small such that $(\not p \cdot \not r << 1)$. In this case the integral value approaches a constant multiplied by a linear phase; i.e. a point scatterer. It will only occur when the target dimensions are less than a tenth wavelength of the illuminating signal, placing it in the Rayleigh scattering region. The TDR signal is mixed harmonically with a phase locked scattered

signal at the imaging frequency. here the two signal sources are phase locked by a phase synchronizer. The high frequency sweeper acts as the slave signal source. These two signals simultaneously illuminate the target. Harmonic mixing of the suitably limited TDR and imaging signals will yield the desired phase corrected data.

Another technique which is useful for ranging is the frequency displaced reference method. Here the carrier is displaced (Δ f). It is important that the object not contribute to the phase shift; hence the displacement must satisfy the following condition:

$$\Delta f \leftarrow \frac{2c}{L}$$
(3.20)

In this case the target structure will not contribute to the net phase shift. The range to the target phase center is given by:

$$R = \frac{\Delta \emptyset c}{\Delta f^{2} \pi}$$
 (3.21)

c - speed of light

Where (44) is the change in phase for the carrier and displace carrier signals respectively. If the frequency shift is small then the phase shift will not be large. The resolution of the system then is directly related to the SNR in the receiver channel since the accuracy to which the phase may be measured is a function of the channel noise. If the displacement is large then the phase shift is increased and the SNR requirements on the signal for a specific resolution is decreased. This relationship may be expressed:

Resolution
$$\frac{r_{oC}}{4f^{2\pi}} = \frac{1/2N_{0B}}{4\pi A \cdot Af}$$
 3.22

The factor of 2 comes about due to the phase uncertainty existing in both the carrier and displaced carrier signals.

An alternate method for implementation of the displaced frequency ranging system is a swept frequency chirp system. The ranging signal frequency is given by :

$$f(t) = f_0 (1 + dt)$$
 where $d = \frac{f_1 - f_0}{T}$ 323

T- Sweep period

f- Initial frequency for sweep

f- final frequency

The scattered signal from the target is given by:

$$f_{g}(t) = f_{r} + \frac{d\phi}{dt} - \frac{f_{r} d}{c} (R_{T} + R_{r})$$
 (3.24)

where $(d \neq /d_t)$ is the change in frequency due to the target structure. The range of the target may then be simply calculated using a frequency counter and sweep time T and sweep width $(f_i - f_i)$. This method could be used in an analog imaging system.

Other Target Derived Reference systems simulate the low frequency reference carrier by measuring the change in phase of the imaging frequency over a narrow band. Over this small band the phase shift is assumed to be linear. In one system a series measurements is made for each frequency point. The first displaced down by a small amount, (Δf) ; the second at f; and the final measurement at $(f+\Delta f)$. Phase and amplitude are measured at each frequency and processed to obtain the target range. In a similar system the three signals are transmitted simultaneously by amplitude modulation of the carrier. These systems are presently under intense investigation by other workers at the Electro-Optics and Microwave Optics Laboratory of the Moore School.

3.2 Practical considerations for range phase removal

Several factors influence which of these systems would be of value in a long range imaging radar system versus a controlled laboratory environment. The distribution of the reference signal for complex field amplitude measurement makes implementation of system requiring a central reference difficult to implement. Techniques are being considered in the E.O. laboaraty for reference distribution using fiber optic that might remove this limitation. Reference distribution is accomplished in the lab readily since the distances are small.

transmitter-receiver pattern Typical arrangements might be the Wells array [14], an orthagonal pattern of receivers and transmitters, a circular array of receivers with central transmitter or a random array. Each combination of receiver and transmitter contributes another line in the frequency domain 3D data volume. For this reason it is advantageous that all combinations of the receivers and transmitters are utilized for data collection. Reference signal distribution difficulty therefore leaves the TDR systems as the only practical alternatives for long range imaging systems. The AM TDR system eliminates

reference distribution by transmitting the reference signal along with the imaging signal and automatically corrects for the target range. TDR systems have the additional advantage that they have immunity to turbulence and inhomogeneities in the propagation medium since both the reference and imaging signals follow the same path.

There are several considerations for determining the best TDR system. Narrow band systems yield only a weighted average of the range to the phase center of the object while wide band system can resolve individual scatters on the target body. A wide band system could adaptively choose one of the scattering centers for the phase reference of the system. A narrow band system could not do this, and any error would introduce image distortion. The narrow band systems have the advantage of automatically correcting the target data for the range phase factor.

For the laboratory imaging experiments a TDR system would not yield the correct phase factor since in this arrangement the object rotates about an axis and the correct phase factor would be a constant, representing the phase shift to the axis of rotation , not the target. The TDR system looks only at the range to the strongest specular reflector on the target surface.

If movement of the target is utilized for aperture synthesis then only one receiver-transmitter pair is required for 3-D imaging and the adaptive system is not required.

Knowledge of the placement of the data in the 3-D frequency domain volume is necessary for the reconstruction of the hologram. The azimuth and elevation angle of the target can be obtained from a conventional radar located at a central location where the data processing and reconstruction is taking place.

IV SYSTEM IMPLEMETATION OF SWEPT FREQUENCY IMAGING

This section of the thesis will describe the research that was done in order to obtain a clear understanding of system performance. Following this will be a section of the thesis devoted to the experimental verification of the swept frequency imaging theory.[] The theory is applied in the simulation of the experiments performed.

4.1 System repeatability

An important parameter of system performance is the repeatability of an experimental measurement. The frequency range over which this possible for the equipment used indicates the bandwidth for which imaging is possible. There are several feedback loops in the system which allow it to track variations in the transmitted signal level. The traveling wave tube amplifier characteristic is shown in fig.4.1. This plot is on a logarithmic scale, indicating that the TWT amplifier gain drops off exponentially below 7.0 GHz and above 15.0 GHz. This measurement was done by first measuring the system response (cables, connectors, attenuator) less the TWT amplifier and then subtracting this response from the TWT amplifier plus system data. A sample of the amplifier output is sampled using a 20 db directional coupler and this is fed to the RF input of the HP 8743A reflection-transmission (R-T) unit. A crystal detector at the 'unknown' port of the R-T unit rectifies a portion of this signal. The detector output is brought to the external signal leveling input of the HP 8620C sweeper.







Fig. 4.1 a) Traveling Wave Tube amplifier gain characteristic in db; 4.0-16.0 GHz. b) Varian VA 618G TWT amplifier (bottom) and HP 8743A reflection-transmission unit (top).
The leveling circuit of the sweeper can level the output over a 20 db range. In addition the AGC in the HP 8410B network analyzer can track the reference signal amplitude over a 40 db dynamic range. System performance may be seen in fig.4.2 for a cylindrical target 7 meters distant from the receiving and transmitting antennas. Two consecutive measurements of the target were made and the results divided. The ideal response would be 0 db flat amplitude and 0 degree phase difference over the entire frequency sweep. With few exceptions due to phase noise at the (+/-)transition point, the system has the desired repeatability in the 5.5 to 16.0 GHz range. Below 4.5 GHZ there are phase errors due to insufficient reference power whereas above 15 GHz errors come about due to the low amplitude of the received signal. Noise may be cancelled by taking multiple measurements and finding the mean. These results indicate the useful data can be recorded in the 5.5-16.0 GHz range

4.2 Computer control of target rotation

When implementing a frequency diversity system with just one pair of receiving and transmitting antennas, the target must then be rotated in the electromagnetic field and the scattering measured for different rotation angles. The target used in the experimental system rotates on a stepper motor driven pedestal. The column of the pedestal is $1 \frac{1}{2}$ meters in length and is made of styrofoam material with minimal cross section. A stepper motor controls table rotation precisely. In order for the pedestal to rotate one $\frac{43}{43}$

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Fig. 4.2 System repeatability for two consecutive scattering measurements of a 80 cm long cylinder 7 cm in diameter; 4.0-16.0 GHz.

revolution the motor must be stepped 10 000 times. The motor is under direct computer control using the digital output port the MINC. A FORTRAN subroutine STEP2 was written for control of the stepper motor. It calculates the number of steps required for a specified angular rotation and moves the table clockwise or counter clockwise based on the direction parameters passed in the subroutine call. The stepper and pedestal are shown in fig.4.3.

4.3 Sphere Simulation

For the calibration of a radar system a reference target is required. The most commonly used reference target is the conducting sphere since its high degree of symmetry does not favor any particular polarization for the incident illumination. Both the bistatic and monostatic scattering of a metallic sphere was simulated.

The general solution for the plane wave electro-magnetic scattering of the sphere was first done by Mie in 1908.[15],[16] In the far field approximation, the scattered field is given by:

$$E_{s}(A, r, o, \phi) = E \frac{e^{i A_{s} r}}{r_{s} r} \left[\cos \phi S_{i}(o) \hat{\Theta} - \sin \phi S_{2}(o) \hat{\phi} \right] \qquad (4.1)$$

where

$$S_{1}(\Theta) = \sum_{n=1}^{\infty} (-1)^{n+1} \left[A_{n} \frac{P_{n}^{1}(\cos \Theta)}{s \ln \Theta} + i B_{n} \frac{d}{d\Theta} \left\{ P_{n}^{1}(\cos \Theta) \right\} \right] \quad (4.2)$$

and

$$S_{L}(e) = \sum_{n=1}^{\infty} (-1)^{n+1} \left[A_{n} \frac{d}{de} \left\{ \frac{P_{n}(Lose)}{SLne} \right\} + i B_{n} \frac{P_{n}(Lose)}{SLne} \right]$$

$$45$$



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Fig. 4.3 Stepper motor and rotating pedestal.

 S_1 (Θ) and S_2 (Θ) are called the complex far field amplitudes for the ($\hat{\Theta}$) and ($\hat{\Phi}$) polarizations respectively. The quantity in the square brackets of eq. 4.1 is called the scattering function.

$$F(\phi,\phi) = \cos \phi \, S_i(\phi) \, \hat{\Theta} - \sin \phi \, S_i(\phi) \, \hat{\phi} \qquad (4.4)$$

The scattering cross section in any arbitrary polariztion ($\hat{\eta}$) for an incident wave polarized in the ($\hat{\gamma}$) direction may be written:

$$\sigma_{n}(\theta,\phi) = \frac{4\pi}{\hbar^{2}} \left| F(\theta,\phi) \right|^{2} \left| \hat{\tau} \cdot \hat{\eta} \right|^{2}$$
(4.5)

Where $(\hat{\gamma})$ is the polariztion of the incident wave and $(\hat{\gamma})$ is the polariztion vector of the receiving system. For the perfectly conducting sphere the coefficients A_n and B_n are: $J_n(A_n e)$ (4.6.)

$$A_{n} = -(-i)^{n} \frac{2n+1}{n(n+1)} \frac{J_{n}(J_{0}, \alpha)}{J_{n}'(J_{0}, \alpha)}$$
(4.6a)

$$B_{n} = (-i)^{n} \frac{2n+1}{n(nii)} \frac{[\frac{1}{2} a_{n} a_{n}]}{[\frac{1}{2} a_{n} b_{n}(\frac{1}{2} a_{n})]}$$
(4.6b)

j_(k_a) - Spherical Bessel function h'_(k_a) - Spherical Hankel function P'_(x) - Associated Legendre function k_ - Wave number of incident wave a - Sphere radius

The prime on the expression for B_{a} denotes differentiation with respect to (k, a).

Polynomial approximations exist for the Mie series exact solution.[15] Different polynomials are used for the three frequency regions for scattering. These are: low frequency or Rayleigh region (k_a) <.4 ; the resonance region .4<(k, a)<20 ; and finally the high frequency or physical optics region (k,a)>20. Two programs BISCAT and SPSCAT implement both bistatic and monostatic cases in the three frequency regions. Fig.4.4 shows the monostatic scattering of the sphere as calculated. This is exactly the same answer for the scattering of the sphere as the exact horizontal axis is solution. The in terms of the dimensionless quantity (k_a). Figure 4.5 a,b,c and d show the bistatic scattering of the metallic sphere of bistatic angles of 30,60,90, and 120 degrees. Note that the approximation are only valid in the range:

where

$S = O\left(\frac{1}{2m}\right)$

which leads to the discontinuities for small (k a) at large bistatic angles. BASIC programs BIDISP and SDISP generate the graphs for the sphere simulations. These programs are in appendix II.

An important result from these simulations is that at high frequencies the scattered signal is of constant amplitude and linear phase irrespective of the bistatic scattering angle. The only exception to this is the forward scattering case when (Θ) equals (π) , where the cross section grows without bound as k increases. This indicates that the only portion of the sphere that is scattering for large (k,a) is the front face closest to both the receiver and transmitter; and therefore a ray optics approximation



Fig. 4.4 Monostatic scattering for the perfectly conducting sphere. (k a) varies from .2 to 10.5 which corresponds to the scattering of a 20 cm. diameter sphere in the frequency range of .1 to 5.0 GHz. Log normalized cross section: log().

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Fig. 4.5 Bistatic scattering of a 20 cm. diameter conducting sphere in the frequency range .1 to 6.0 GHz; .2<(k, a)<12.6; polarization of the receiver equal to scattered wave polarization. a) Geometry for scattering expression. b) Bistatic angle 30° c) Bistatic angle 60° d) Bistatic angle 89° e) Bistatic angle 120°



(b)

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(C)

Fig. 4.5 (contd.) Bistatic scattering of a 20 cm. diameter conducting sphere in the freequency range .1 to 6.0 GHz; .2<(k_a)<12.6; polarization of the receiver equal to the scattered wave polariztion. a) Geometry for scattering expression. b) Bistatic angle 30° . c) Bistatic angle 60° . d) Bistatic angle 89° . e) Bistatic angle 120°.







(C)

Fig. 4.5 (contd.) Bistatic scattering of a 20 cm. diameter conducting sphere in the freequency range .1 to 6.0 GHz; .2<(k,a)<12.6; polarization of the receiver equal to the scattered wave polariztion. a) Geometry for scattering expression. b) Bistatic angle 30°. c) Bistatic angle 60°. d) Bistatic angle 89°. e) Bistatic angle 120°. may be applied to find the scattered field.

1!

$$S_{1}(0) = S_{2}(0) = -\frac{1}{2} k_{2} a(e^{-i 2 k_{2} a \cos \theta/2})$$
 (4.8)

for the bistatic case. In the monostatic case this reduces to:

$$F(0) \rightarrow -\frac{1}{2} + a = -\frac{1}{2} + -$$

For targets with features larger than a few wavelengths in size resonance effects become minimal.

Another possible reference target is the long cylinder (1 >> a). The scattering of the cylinder in the high frequency region when it is oriented vertically yields an answer similar to that of the sphere. For (k a) > 5 where a is the cylinder radius, resonance effects disappear and the copolarized scattered field is given by: [15]

$$E_{s} = E_{i} \left(\sqrt{\frac{3 \cdot \sigma_{s} \frac{6}{2}}{2r}} \right) \exp \left\{ i \cdot \frac{1}{2r} \cos \frac{6}{2} \right\}$$

$$r = (R_{T} + R_{r})$$
(4.10)

The equation for the scattering of the cylinder is used for the computer simulations of frequency swept holography.

4.4 Simulation of frequency swept imaging

A series of frequency swept hologram simulations were done of targets that would later be imaged experimentally. The basic arrangement consists of separate receiving and transmitting antennas which measure the scattering of a target that rotates about an axis. The center of rotation is chosen as the phase center of the imaging system. For this configuration the frequency domain data lies in a plane perpendicular to the axis of rotation. Therefore the transforms of the holograms will be slices in this plane. The first object hologram simulated was comprised of two cylinders equidistant from the rotational axis. This target is shown in fig.4.6. Approximations for the various distances were derived:

$$r_1 \simeq r - \frac{4}{2} \sin \left(\frac{9}{2} - \Theta \right) (4.11a) \qquad r_2 \simeq r + \frac{4}{2} \sin \left(\frac{9}{2} - \Theta \right) (4.11b)$$

$$x_{1} \simeq r + \frac{1}{2} \sin(\frac{1}{2} + \theta)$$
 (4.11c) $x_{2} \simeq r - \frac{1}{2} \sin(\frac{1}{2} + \theta)$ (4.11d)

the waves striking cylinders Cl and C2 are given by:

$$E_{c_{1}} = E_{o} e^{-i \cdot k_{o}(\vec{x}_{i})}$$
 (4.12a)

$$E_{e2} = E_{0} e^{-i \frac{1}{2} k_{0}(\vec{x}_{2})}$$
 (4.12b)

The scattered waves from the two cylinders including the cylinder response then follows:

$$E_{sc_1} = E_{c_1} \sqrt{\frac{a \cos \varphi_{r_1}}{c(R_{P})}} e^{-i \beta_{c_0}(r_1 - 2a \cos \frac{\varphi}{2})}$$
 (4.13a)

$$E_{SC_{2}} = E_{c_{1}} \sqrt{\frac{\alpha \cos \theta_{L}}{Z(R_{r})}} e^{-i \frac{1}{R_{0}} (r_{L} - 2\alpha \cos \theta_{L})} \qquad (4.13b)$$

This is further simplified by combining terms :

$$E_{s} = 2 E_{o} e^{i \frac{1}{2} \cdot (2 \alpha \cos \frac{4}{2})} \sqrt{\frac{\alpha \cos \frac{4}{2}}{2 (Rr)}} \left\{ \cos \frac{1}{2} \cos \frac{4}{2} \sin \frac{4}{2} \right\}$$
(4.14 b)

Finally take the real part of this function for display:

$$R(E_s) = C \cdot \cos(2k_e \cos\frac{\phi}{2}) \cos(k_e + \cos\frac{\phi}{2}) \sin \theta \qquad (4.14b)$$

This was done for a two cylinder target with cylinders 5 cm. 54





in radius ,separated by 25 cm. using program CYLIN. The results were displayed by program CDISP on a Tektronix 606A CRT display. CDISP gives the option of varying the gray scale compression of the hologram either logarithmically or by constant multiplication. These programs are listed in appendix II. The resultant hologram and the reconstructions obtained through Fourier transformation on the optical bench appear in figs.4.7 a,b The hologram simulated a sweep from 2.0 to 18.0 GHz in 64 frequency steps. The target in the simulation rotated 360 degrees in 128 steps.

Another target simulated which did not have the symmetry of the first target was comprised of two cylinders both mounted to one side of the rotational axis, as shown in Two simulations were done of this target with fig.4.8. varying diameter cylinders. In the first case 7 cm radius cylinders were used. The hologram for this case and the Fourier transform reconstructions are shown in figs.4.8 For the second simulation the target was two b.c.d. cylinders 3.5 cm in radius. In both cases the cylinders were located 10 cm from the center of rotation and the simulation was for a 2.0 to 18.0 GHz sweep. The hologram and the transformed images are shown in figs.4.9 a,b,c. The two cylinder off axis target was simulated by first calculating the copolarized scattered field for a single cylinder:

$$\vec{F}_{g} = 2\vec{E}_{g} e^{-i\frac{\pi}{2}2a} e^{-i\frac{$$



(a)



Fig. 4.7 a) Simulation Hologram of two cylinder target; 10 cm. diameter, 25 cm. apart. Frequency range: 2.0-18.0 GHz; 128 lines, 64 points/line. b) Optical Fourier transform of hologram.







(b)



(C**)**



Fig. 4.8 a) Geometry of two cylinder off-axis target. b) Hologram simulation two cylinder off-axis target; 2.0-18.0 GHz; 64 points/line; 128 lines. b) Transform with zero order term. c) Transform with zero order removed.



(a)



Fig. 4.9 a) Hologram simulation of off-axis two cylinder target with cylinders 7 cm. in diameter; 2.0-18.0 GHz b) Optical Fourier transform with zero order. c) Transform with zero order term removed. taking the real part:

R{E_} = 2E_ { cos-kza cos he sin 2 a by sin 2 a by sin 2 a by (4. 50) = 2E_ cos (to 2 a - to 1 sin 0)

For the two cylinder off axis target, one cylinder is at $(\ominus)=0$ and the other at $(\ominus)=90$, therefore the scattered field is given by:

In general for an arbitrary set of circular scatterers of radius a and distance 1 from the origin; the scattered field may be written:

$$R\{E_{s}\} = \sum_{n} \cos \left\{ \operatorname{le}_{n} - \operatorname{le}_{n} \operatorname{sin}_{(\Theta + \Theta_{n})} \right\} (4.17)$$

This gives the capability to simulate the scattering of any target given that it can be decomposed into N spherical scattering centers.

4.6 Experimental results

An experimental system for the implementation of swept frequency imaging was setup in the anechoic chamber at the Graduate Research Center in the Moore School. The frequency range for these experiments was from 6.3 to 16.0 GHz in 64 discrete steps. The targets were rotated 360 degrees in 128 steps. These holograms were then identical in form to the simulations previously done.

The system for error correction and range phase shift removal was that used for high signal to noise ratio signals. The reference target was a cylinder positioned so that its front face was located on the axis of rotation as in fig.4.10. A plot of system response is shown in fig.4.11. This data represents the combined characteristics of the antennas, amplifier , cables and clutter. In addition it contains the linear phase shift that is the range phase factor. As an example for the two cylinder target shown in 4.12 the raw data, magnitude and phase is shown in fig.4.13. Figure 4.14 shows how this data has been corrected for range phase and system response. This data was generated using the Fortran program SPHER3.

The experimental properties of the two cylinder target were studied extensively. Both the scattering as a function of frequency for a specific orientation of the target and the scattering as a function of angular rotation at specific frequencies was obtained. In figs. 4.15 a,b,c are shown the corrected frequency response of the targe for orientations of 45,90° and 135° degrees.

Another computer program ANTPAT was written to obtain the radiation pattern of an arbitrary target or antenna. In the two cylinder case the pattern was measured at 5.0,10.0and 15.0 GHz. Note that when the cylinders are collinear all that is seen is the front surface specular reflection of the one cylinder hence the pattern of a point scatterer in the vicinity of 0° degrees. These patterns are shown in figs.4.16 a,b,c. At high frequencies the lobe spacing is much closer than at low frequencies, consistent with the theoretical result for the pattern. To see this examine the





Fig. 4.10 Reference target on pe cylinder, 7 cm. in diameter.

pedestal; 80 cm. long













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Fig. 4.13 Uncorrected scatter data for symmetrical two cylinder target ; $(\Theta) = 0^\circ$, 6.3-16.0 GHz.

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Fig. 4.14 Corrected two cylinder target data using system response of Fig. 4.11.



(Q)



(b)



Fig. 4.15 a) Corrected two cylinder symmetrical target response; 6.3-16.0 GHz, $(\Theta) = 45^{\circ}$. b) $(\Theta) = 90^{\circ}$. c) $(\Theta) = 135^{\circ}$.



(Q)



(b)





Fig. 4.16 Scattering pattern of two cylinder target of Fig. 4.12 at different frequencies as a function target rotation angle. a) 5.0 GHz. b) 10.0 GHz. c) 15.0 GHz.

expression for the monostatic scattering of the symmetrical two cylinder target:

$$E_s = C \cdot \left\{ e^{i \frac{1}{2} e_s 2 a} \cdot c_s \frac{1}{2} e^{i \frac{1}{2} e_s 2 a} + c_s \frac{1}{2} e^{i \frac{1}{2} e_$$

The expression for the scattering of the target at a given frequency as a function of angular rotation may be written:

As for the swept frequency response; $l \sin(\Theta)$ remains constant and the hence the swept frequency response is sinusoidal with period dependent on (k) and l.

The final test for the system was the generation of actual holograms. The first target measured was the two cylinder target shown in fig.4.17a. The cylinders are of aluminum, 80 cm. in length and 7 cm. in diameter. The real part of the corrected swept frequency data in the range 6.3 to 16.0 GHz was stored and displayed on the CRT. The targets were rotated 360 degrees in 128 steps yielding a total of 8192 points in the hologram (64 points/line * 128 lines). The center of the hologram is at 0 HZ with radial distance directly proportional to frequency. An example is shown in fig.4.17b and the reconstructions in figs.4.17c and d. These Fourier transforms where done optically . [17]

This procedure was followed for other targets not having the symmetry of the first object used. The target type was the same as the simulations done previously. The



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(Q.)



Fig. 4.17 a) Two cylinder target in anechoic chamber on rotating pedestal. b) Hologram of target measured between 5.3 and 16 GHz corrected for range and system response; 128 lines, 64 points/line. c) Optical Fourier transform of hologram. c) Optical Fourier transform without zero order term. first of these was a single cylinder mounted off axis as shown in fig.4.18a. This cylinder was the same as the others used and the frequency range and angular sweep were identical to that of the two cylinder target. The hologram and reconstructions are shown in figs.4.18 b,c,d. The final target was the two off axis cylinder target pictured in fig.4.19a. The frequency diversity hologram and transforms are in figs.4.19 b,c,d.





(b)

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(C)

(4)

Fig. 4.18 a) Single cylinder off-axis target and position in anechoic chamber. b) Experimental hologram of target; 6.3-16 GHz; 128 lines, 64 points/line; corrected for range and system response. c) Optical Fourier transform of hologram. d) Optical Fourier transform without zero order term.







Fig. 4.19 a) Two cylinder off-axis target geometry and position in anechoic chamber. corrected for range and system response. c) Optical Fourier transform of hologram. d) Optical Fourier transform without zero order term.

V Conclusion

This thesis has described an automated swept frequency measuring system. This system may be used for radar cross section measurement, antenna pattern measurement and swept frequency holography. The system has a useful range of 5.0-17.0 GHz in which amplitude and phase of the scattered microwaves from targets in the anechoic chamber of the Graduate Research Center may be recorded and stored. A DEC MINC LSI-11/2 completely automates the data acquisition complete error process. Α correction algorithm was implemented using the data storage and processing capabilities of the minicomputer.

The effect of range phase shift on swept frequency holograms was investigated and various techniques for its removal were investigated. It is believed that a TDR system for range phase removal is required for implementation of a practical radar system. This system must have extremely high resolution for coherent imaging. The relationship between TDR system bandwidth, receiver channel bandwidth and resolution was derived:

Resolution =
$$\frac{\sigma_0 C}{\Delta f \ z \pi} = \frac{\sqrt{2} N_0 B C}{4 \pi A \ \Delta f}$$
 322

where (Δf) is the imaging bandwidth, (R) the range uncertainty, B the receiver channel bandwidth and N. the noise power spectral density.

Simulations were performed for the scattering of various radar targets which include the conducting sphere, 72

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 322

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Simulations were performed for the scattering of various radar targets which include the conducting sphere,

infinite cylinder and combinations of these. The holograms recorded from these analytical results reconstruct the targets extremely well and set a goal for practical system performance. An expression was derived for the scattering of N spherical/cylindrical cylinders in a plane passing through their centers:

$$R\{E_{s}\}=\sum_{n}\cos\left\{-kz\alpha_{n}-kl_{n}\sin\left(\Theta+\Theta_{n}\right)\right\}$$
(4.17)

where 1 is the distance from the axis of rotation (Θ_n) the angle relative to some reference for the target angular position and a the target radius.

Finally experimental swept frequency holograms were generated using a rotating pedestal under computer control to scan the target in one dimension. The experiments done indicate the feasibility of implementing a practical holographic radar system. The holograms obtained for various targets agree well with theory even though the error correction and range phase shift removal techniques used were robust in nature. It is believed that better images are possible given that the error correction techniques previously outlined are implemented.

Further work can be done in testing the TDR techniques for their suitability for an imaging system. The system may also be expanded to include scanning in the () direction to give true 3-D imaging capability. This may be implemented by adding a stepper motor controlled azimuthal scanner to the top of the rotating column.

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APPENDIX I

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	PHASE=FLO	WRITE (IUN	CONTINUE	ME LINKN	*******		FORMAT(/]	GENERAT 10		FURMAT (//	AMPLITUD	FORMATC//	FORMAT (/ /	SYSTEM ++	FORMAT(//	FORMAT(/'	FORMAT (/	FORMAT(/'	FORMAT (/	(FORMAT (/*	FORMAT(//	SYSTEM **	FORMAT (//		FOPMAT (10	FORMAT (/*	EDRMAT(/ 1	ENDING F	11 510d1 .	STOP	END	STORAG		OFFSET	000032	000034		/ PADA /.	OFFSET	000000	004002	
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004000 000000 400900 00400 000004 OFF3ET OFF3E1 OFF SET TYPE 1YPE DIMENSIONS (512) (512) (512) (512) (512) (512) (512) (512) (512) (512) (512) (512) (512) (512) (512) (512) (512) (512) (512) TYPE IFLA02 1+2 IFLAG3 1+2 **1**+2 IFLAG4 1+2 1+2 /, SIZE = 004004 (1026. WORDS) NOME **WAR** TPH4 **WVW** IPHS HORDS) HORDS) 512 TRANSA/, SIZE = 010004 (2050. 002000 ANTA /, SIZE = 010004 (2050. OFFSET 002000 002000 OFFSET 004004 OFF SET 004004 002000 002000 002000 002000 002000 002000 002000 002000 002000 002000 002000 002000 002000 002000 002000 012000 TYPE TYPE TYPE 1+2 1#2 1+2 1+2 I+2 000000 0014004 0014004 002000 0022000 0022000 0022000 0022000 0FF%E1 000000 1 AMP4 000000 1 AMP 6 W MAME MAME SH41 IPH2 EH41 SECTION O PAGA ISTA ANTA ANTA ANTA ISTA ANTA ANTA ANTA ANTA ANTA ANTA ANTA ION ARRAYS: K /1STA 00000 00000 FFSET FFSET 00000 04002 FISET 04002 04002

FUNCTIONS. STATEMENT AND PROCESSOR-DEFINED FUNCTIONS

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TIFE Mon TYPE W TYPE MAR TYPE NAME

7 **.** K TRANG -PAD 1+2 151 FLOAT R+4

FACE OOI

THU 09-AUG-79 00.46.59 V02.1-1

ERSION 1 0

HIS ROUTINE WILL GENERATE THE CHARACTERISTIC OF THE FAD USED IN THE DETERMINATION OF THE SYSTEM RESPONSE IN THE FRECUENCY DIVERSITY THADING SYSTEM. IT INVOLVES FLACING THE FLECTENC ARM WITH THE REFLECTION TRANSMISSION UNIT WITH THE ARM OMMEGIED FROM THE UNKNOWN TO THE TRANSMISSION RETURN FORTS DITH A 6 DB PAD IN SERIES WITH THE RAM. THE PRECISION AFC-7

004020 NGAPP 1+2 04016

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THIS SECTION WILL READ AN OLD PAD RESPONSE FORMAT(///%', ENTER NUMBER OF POINTS IN PAD NE-4 CNSE ') 2 FORMAT(//%', ENTER START FREUMENCY IN GIGAMERT2 ') FORMAT(///%', CENTER ENDING FREUMENCY IN GIGAMERT2 ') FORMAT(//1%', CONNECT THE FLEXIBLE ANN PLUS & DB FAD 'N'; FORMAT(1) . 'ENTER THE PAD RESPINSE FILE WHE PRESET LEEE BUS (IPHA(K) 01 720) IPHA(K)=IPHA(K)-1440 (IFHA(Y) LT -720) IPHA(K)=IPHA(K)+1440 CALL ASSIGN(12, --1, 'OLD', 'NC', 1) REAU(12, +) FSTART REAU(12, +) FEND READ(12. .) NPDINT STEF=(FEND-FSTART)/FLOAT (NPOINT) STEP= (FEND-FSTART) /FLOAT (NPOINT) CALL SWEEP(1, FSTART, 1, 4) 03 'ENDING FREQUENCY IN ONZ FREISEFSTART + FLOAT (K-1) + STEP CALL SWEFP (1, FREO, 1, 4) CALL PHAMP2(14, 1P, NSAMP) DO 204) K=1, (NPUINT) FREQ=FSTART+FLOAT(K=1)+STEP CALL SWEEP(1, FSTART, 1, 4) TYPE OOS CALL SHEEP(1, FREQ, 1, 4) CALL PHAMP2(1A, 1P, NSAMP) CALL SWEEP(10, FRED, 1.4) (AMPAD())=IA-IAMPAD(K) WRITE(12.*) IAMPAD(K) WRITE(12.*) IPHA(K) READ(12.+) IANPAD(K) READ(12.+) IPHA(K) WRITE(12, +) FSTART WRITE(12, +) FEND ARITE(12. .) NPOINT [PHA(V)=]P-]PHA(V) INICIAN TERLOWY ON [AMPAD(1') = [A-2048 DO 300 F-1, MPUINT [PHA(F)=]P-2048 ALL CLOSE(12) CALL CLOSE(12) ACCEPT +, FEND P= 1P-2048 8402-01=01 00 10 700 **BINI INUD** CONTINUE BUNINUE Ş **BUNITING** TYPE ODS NETTER PAUSE PAUSE FYPE u **6**03 u 902 TYPE FORTRAN IV 6 409 0062 403 0064 900 <u>888</u> 8 ç Ş 80 8 8 u 000 0055 0056 0056 0056 0056 0056 0030 0031 0032 0063 0063 00045 00045 00045 00045 00045 00045 00045 00045 1100 0013 0012) 3) 3 I AMPAD ARRAY WITH AMPLITUTE OF PAD RESPONSE IN INTEGER FORM TPHA ARRAY WITH PHASE STURED IN INTEGER FORMAT RECIEVED IFAGD OF OF ARRAY WITH PHASE STURED IN INTEGER FORMAT RECIEVED IFAED OF DID DIA FSTART STARTING FREQUENCY FOR FREQ RESPONSE FSTART STARTING FREQUENCY FOR FREQ RESPONSE FREQUENCY STARTING FREQUENCY FOR FREQ RESPONSE CREATE SERVENTIAL AFCESS ASCIT FILE USING SIMPLE LIST NOT FORTAM CARRATICE CONTROL. THE FILE WILL EXPAND AS FORTAM CARRATICE CONTROL. THE FILE WILL EXPAND AS MECESSARY THE RECORD SIZE IS & CHARAGETERS, THE WITT ANTHER IS LODICAL UNIT 12 THE MARE WILL BE ASSIGN USING THE ASSIGN SUBMOUTINE IN THE LIPRARY. THIS IS A NEW FILE ON THE DISC AND THE MICRYMANE SHEEFER DATA IS OBTINED IN THE FRECHER'S FROM REGION SPECIFIED AND IN THE INCREMENT SPECIFIED AT THIS POINT THE PAD MARSE CHARAGTERISTIC IS DESTRED IS FLACED INTO THE MCMITURE AND PHASE OF THE TWO MEASUMMENTS ARE MEASURED THEN AND THE RESULTANT VALUES ARE THE CHARACTERISTIC THE ONLY THE PAD RF DIRECTLY FROM SECTION TO DETAIN CHAPACTERISTICS OF ARM AND & DE PAD. NOT THE ACTUAL PAD TO BE MEASURED IN THE EVENT THAT THE FILE CONTAINING THE PAD RESPONSE ALREADY EXISTS THEN THE PROCHAM WILL FILL THE ARRAY PAD WITH THE PAD RESPONSE FACE 002 comicn/pada/144Pad. 1PMA, 1FLAG, FSTART, FEND, STEP, NPOINT, NSAP Integer 144Pad(512), 1PMA(512) THE CONVERSION FROM INTEGER TO REAL ANPLITUDE IS WINDER OF POINTS IN SYSTEM RESPONSE UNIT OBTAINS IF (IFLAD) 100, 100, 500 THEN OR OLD DATA? TYPE 900 PRINT FIRST INSTRUCTIONS USE ASSION SUBROUTINE TO DET NAME FOR FILE V02. 1-1 THU 09-AUG-79 00 46:59 MILL BE SAVED WHEN THE FILE IS CLOSED STARTING FREQUENCY IN CHI AMPLITUDE=(IAMPAD(I)-2048)+. 05 CALL ASSIGN(12...-1. MEN' MEN' 1.) 7-E 3-1 PHASE=([PHA-2048) + 25 AMPTER IN SERIES. COMPON DECLARATION ACCEPT - MPOINT TYPE 902 'STAR THE HICKNANCE AVERAGED **TYPE 901** υ 2 8 G U FORTRAN U, 0000 8003 8003 8003 * **S** 800 ٢ ¢ e ١

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PAGE 001		PETHEEN THE THO	LEN Y RALIDE ALD COMMON BLOCK NY 1574 THE	TIC INFEDENCE AUD	IN ARE REMOVED		s I (audu	ENATE NEH DATA	T EACH FREGUENU	5TEP. NPOINT, N33MP		PH2 (512)																		FAGE 001		
VO2 1-1 THU 09-6440-79 00 57.36	RESIGN 1 3	IS SUBROUTINE WILL MEASURE THE ISOLATION	ANNELS OF THE NETWORK ANALYZER THE FREOU E MANBER OF STEPS WILL BE PASSED THROUGH VIA THE OUTPUT WILL BE THROUGH COMMAN PLO	IBLES TO BE USEN IN THE ACTUAL MEASUREMENT STEM ARE TERMINATED IN THEIR CHARACTERIST	re magniture and physe of the lear age sign the same with of the n/a front panel as	ASURENTS	HIP2 ARRAY WITH AND ITUDE OF LEAKED S IPH2 ARRAY WITH PMASE OF LEAKED SI	LAGZ FLAGFOR CLD OK NEW LATA O=>GENE 1±> READ JLD DATA	SAMP2 NUMBER OF SAMPLE POINTS TAKEN AT	MHON VPADA / I.MHAD, IPHA, IFLAG, FSTART, FEND, 1	JHHON/ISTA/IAHP2, IPH2, IFLA02, NSAMP2	NIEGER 14400(512), IPMA(512), IA402(512), I	F(JFLAG2) 100.100.500	APPE VIN) MISE	ALL SWEEP(1, FSTART, 1, 4) 0.200 f=1, NF01NT	RE0-FS1ART+FLUAT(K+1)+STEP M. (EBE0.1.4)	ALL SWEET (JATTER JAT)	AMP2(K)=1A-2048 PW2(t)=1P-2048	ONTINJE ALL SWEEP(10,FSTART,1,4)	YPE 901 ALL 0551GN(12++-1+/NEM++*NC++1+)	RITE(12, 4) FSTART Dite(12, 4) EEND	RITE(12, •) NPOINT	O BUD K=1.NPOINT	RITE(12. +) IPH2(K)	(N1) N.E. 2014 - C.D.S.E.(32)	0 10 700	VPE 402	ALL_ASSION(12,,-1,'OLD','NC',1)	EAD(12,4) FSTART FAD(12,4) FEND	V02 1-1 THU 09-MUG-79 00:37:38	EAD(12, *) NPCINT STEP=(FEND-FSTART)/FLOAT(NPOINT)	0 600 K=1.MPUINT TEAD(12, +) TAMP2(K)
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ETURN', FOR	H THE ', 000	ŗ.		FF861	00032				FFACT	04000 0002	04012	4000	2000	0008 0007	8000 8000	0100	0011 0012 0012	TVPE 0013	Re4 001	88		002	000	0024	() 0025	002)	200 1	031	C FORTR	0034	003
D THE RETURN'S	IES WITH THE "	C	tos)	TYPE OFFSET	1*2 000032				TYPE OFFRET	1=2 004000 0003	1+4 004012 0003	40000 40000	5000	0009	8000 SM	0100	FUNCTIONS. 0011 0012	NAME TYPE 0013	SHEEP Reit 001			002	002	0024	(0025 0024	002)	200	033 033		0034	
NOUN TO THE RETURN', FOR	IN BERIES WITH THE	SE FILE NAME	18 MORDS)	AME TYPE OFFSET	P 1+2 000032			8	ANE TYPE OFFRET	FLAG 1+2 004000 0002	TEP R+4 004012 0003	9000	2000	0006	INENSIONS 0009 512) 0009	312) 0010	EFINED FUNCTIONS. 0011	IVPE NAME TYPE 0013	Re4 SMEEP Re4 001			002	002	0024	0025	002)	200	0.031	C FORTH	0034	003
THE UNDOON TO THE RETURN', FOR	ED PAD IN SERIES WITH THE	RESPONSE FILE NAVE ')	r PAD 14 (138 MORDS)	NAME TYPE OFFSET) IP 1+2 000032			(SOLON)	NAVE TYPE OFFRET	0002 IFLAG 1+2 004000	STEP R+4 004012 0003		5000	0000	DIMENSIONS 0009 112) (512) 0009	12.) (512) 0010	SSOR-DEFINED FUNCTIONS. 0012 0012	ANE TYPE NAME TYPE 0013	MANP2 Re4 SHEEP Re4 001			002	002	0024	0025	002)	2000	0031		0034	003
FROM THE UNBOROMAN TO THE RETURN', FOR N. VJFR')	DESIDED PAD IN SERIES WITH THE'.	LD PND RESPONSE FILE NAVE. ')	Militiand = Occo44 (18. NORDS)	OFFSET NAME TYPE OFFSET	000030 1P 1+2 000032			(1033. MDRDS)	OFFSET NAME TYPE OFFSET	002000 IFLA0 1+2 004000 0000	004006 STEP R+4 004012 0003	004020	2000		DIMENSIONS 0 (512) (512) (6009 0000	0 (512) (512) (010	PROCESSOR-DEFINED FUNCTIONS. 0012 0012	PE NAME TYPE NAME TYPE 0013	100 PHONP2 Re4 SHEEP Re4 001	100		002	002	0029	002)	2000	0.031	Locate	0034	003
CONECT FROM THE UNDODEN TO THE RETURN', FOR	CT THE DESIDED PAD IN SERIES WITH THE	THE OLD PAD RESPONSE FILE NAME ")	PROGRAM UNIT PAD Size = 000044 (18. MORDS)	TYPE OFFSET NAME TYPE OFFSET	1•2 000030 1P 1•2 000032			004022 (1033. MDMDS)	TYPE OFFSET NAME TYPE OFFSET	1+2 002000 1FLAG 1+2 004000 0002	R+4 004006 STEP R+4 004012 0003	1•2 004020 004020	2000	0005	002000 (512) (512) 0000 0000	002000 (512) (512) 0010	ENT AND PROCESSOR-DEFINED FUNCTIONS. 0012 0012	NE TYPE NAME TYPE NAME TYPE 0013	0017 Re4 PMMMP2 Re4 SMEEP Re4 0012			002	002	0024	002)	200 			0034	
AND COMECT FROM THE UNCOUND TO THE RETURN', FOR E NETWORK ANN YSER'S	COMMECT THE DESIMED PAD IN MERLES WITH THE',	(ENTER THE OLD PAD RESPONSE FILE NAME: ')	AP FOR PROGRAM UNIT PAD Subata, Size = Occo44 (18. Nords)	NNE TYPE OFISET NAME TYPE OFFSET	IA I+2 000030 IP I+2 000032			12E = 004022 (1033. MDRDS)	NAME TYPE OFFSET NAME TYPE OFFSET	1PHA 1+2 002000 1FLAG 1+2 004000 0002	FEND R+4 004006 STEP R+4 004012	NSAMP 1+2 004020 000020	2000 SA		200000 002000 (512) (512) (009		STATEMENT AND PROCESSOR-DEFINED FUNCTIONS. 0012 0012	E NAME TYPE NAME TYPE NAME TYPE 0013	FLOAT Res PHANP2 Res SHEEP Res 001			002	002	0029	002	002		200 			0034	
IES', / MAD COMMECT FROM THE UNCOUGH TO THE RETURN', FOR	T(//IX. COMMECT THE DESIRED PAD IN SERIES WITH THE'.	r(//\$.'.KWTER THE OLD PAD RESPONSE FILE NAME ')	MANGE MAP FOR PROGRAM UNIT PAD PSECT SDATA, Size = 000044 (18, MORDS)	T NAME TYPE OFFSET NAME TYPE OFFSET	4 IA 142 000030 IP 142 000032	2		· /· \$12E = 004022 (1033. MDRDS)	I NAVE TYPE OFISET NAVE TYPE OFISET	0 1PHA 1+2 002000 1FLAG 1+2 004000 0002	2 FEND Re4 004006 STEP Re4 004012 0003	6 NSAMP 1+2 004020	SUDDE		0000 00000 002000 (512) (512) (512) (5000 00000 00000 00000 00000 00000 00000 0000	UN UNZOUD 002000 (512.) (512) 0010	IONS. STATEMENT AND PROCESSOR-DEFINED FUNCTIONS. 0012 0012	TYPE NAME TYPE NAME TYPE NAME TYPE 0013	E Ret FLOAT Ret PHONE2 Ret SHEEP Ret 001			007 002	002	0029	0023			200 		Lostin Lostin	003	003
Y SERIES',/' AND COMMECT FROM THE UMCOURN TO THE RETURN', FOR P. PORT ON THE NETROPK ANNLYTER'S	FORMAT(//1% CCAMECT THE DESIMED PAD IN SERIES WITH THE',	FORMAT(//4/, 'ENTER THE OLD PAD RESPONSE FILE NAVE ') End	STORAGE MAP FOR PROGRAM UNIT PAD Les. Psect Buata, size = occo44 (18. mords)	OFFSET NAME TYPE OFFSET NAME TYPE OFFSET	000024 IA I+2 000030 IP I+2 000032	000022		: /PADA /. SIZE = 004022 (1033. MDRDS)	OFFSET NAME TYPE OFFSET NAME TYPE OFFSET	000000 IPHA 1+2 002000 IFLAG 1+2 004000 0002	004002 FEND R+4 004006 STEP R+4 004012 0003	004016 NSAMP 1+2 004020	COMPON ARRAYS		PADA 00000 002000 (512) (512) (000	PHUM UNZOUD 002000 (512.) (512) 0010	FUNCTIONS, STATEMENT AND PROCESSOR-DEFINED FUNCTIONS. 0011 0012	NAME TYPE NAME TYPE NAME TYPE NAME TYPE 0013	CLOSE Red FLOAT Red PHANP2 Red SHEEP Red 0013	100 100		002	002	0024	0022			200 			003	
1' SERIES', /' AND COMMECT FROM THE UMONOMM TO THE RETURN', FOR 2' PORT ON THE METHODIX ANNA YJER')	905 FORMAT(//1X, COMMECT THE DESIMED PAD IN SERIES WITH THE', 1, 6 DB PAD')	906 FORMAT(//4/, 'ENTER THE OLD PAD RESPONSE FILE NAME: ') End	W IV STORAGE MAP FOR PROGRAM UNIT PAD VAPIABLES. PSECT BLATA, SIZE = 000044 (18. MORDS)	TVPE OFFSET NAME TVPE OFFSET NAME TVPE OFFSET	R+4 000024 IA I+2 000030 IP I+2 000032	1•2 000022		1 BLOCK /PADA /, SIZE = 004022 (1033. MDRDS)	TYPE OFFSET NAME TYPE OFFSET NAME TYPE OFFSET	11+2 000000 1PHA 1+2 002000 1FLA0 1+2 004000 0002	R+4 004002 FEND R+4 004006 STEP R+4 004012 0003	1+2 004016 NSAMP 1+2 004020	AL AND COMPON ARRAYS		142 PADA 00000 002000 (512) (512) 0009 142 PADA 00000 002000 (512) (512) 0009	1.4 Filter W2000 002000 (512) (512) 0010	TINES, FUNCTIONS, STATENENT AND PROCESSOR-DEFINED FUNCTIONS. 0011 0012	TYPE NAME TYPE NAME TYPE NAME TYPE NAME TYPE 0013	Red CLOSE Red FLOAT Red PMANP2 Red SHEEP Red 0012	80 · · · · · · · · · · · · · · · · · · ·		002	002	0024	002 200	803 		200 		Loste	0034	003

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FORTRAN IV VO2 1-1 THU 0001 SUBROUTINE TRANS	191 MEASURES ANDLITUKE AND PHASE FROM . C VERSION 1 3 ER IN ORDER TO OPTAIN ITS DIRECTIVITY . C VERSION 1 3 SYSTEM CABLES TERMIMATED IN THEIR CHANNA. C THIS SUBROUTINE HILL F	EVELLE NAME FOR THE ISOLATION DATA *) C CHARATTERISTICS ANE FOR IST READS MADNITUDE AND PHASE DATA ', C CHARATTERISTICS ANE FOR DATA FILE '/') C C ATTANIS AND THE CONNECTED OLD FILE NAME FOR THE ISOLATION DATA. C IS ALREADY FUONT THE	RAM UNIT 1ST CULATED AND STORED C CULATED AND STORED C IAMP3 ARRAY CONTA C IAMP3 ARRAY CONTA	OFFSET NAME TYPE OFFSET C AMPLITUDE OFFSET NAME TYPE OFFSET C IPH3 ARAY CONTA	0000124 IP 142 000024 CONTA CO	C T/R= (C+P+G)/R	C T RECEIVER CHANNEL SI	E OFFSET NAME TYPE OFF3ET C R: REFERENCE CHANNEL S C C: NETWORK ANALYZER REI	002000 IFLAD I*2 004000 C ISOLATION C P PAD CHARACTERISTIC	004006 STEP Re4 004012 C 0 FREGUENCY CHARACTER	004/20 CORRECT THIS FOR THE F	I (1026. MORDS) C AND THE PAD FREQUENCY	OFFSET NAME TYPE OFFSET C C/R => M1 MEASURMENT C C/R => M1 MEASURMENT	002000 IFLA02 1+2 004000 C C OVER REFERE	C (C+P+0) => M2 Her C (C+P+0) => M2 ISOLATION C C		00 (512) (512) 00 (512) (512)	D PROCESSOR-DEFINED FUNCTIONS (11PH3(512), 15 (11) PROCESSOR-DEFINED FUNCTIONS)	YPE NAME TYPE NAME TYPE O 0006 IF (IFLAG3) 100, 100.50	Re4 PHANP2 Re4 SHEEF Re4 0007 100 TYPE 900 0008 TYPE 901
5(12)	K, 'SUBPROUTING ETHORY ANALY THE ANTENNU IMPEDENCE '/)	, 'ENTER THE A IX.' 'SUBROUTIN THE SPECIFIEI G', 'ENTER THE	T BDATA, SIZE	NAME TYPE	IA [•2	20100 - <u>2010</u> 0	20400 - 3116	NAME TYP	IPHA [+2	FEND R+4	NSAMP 1+2	512E = 00400	NAME TYPE	IPH2 1+2		I OFFSET	000000 000 000000 000 002000 0000 002000 0000	STATEMENT A	JINN 34	1+4 FLOAT

DO 2000 M-1, MOUNT REEDESTRATE +ELANDICU-19-STEP CALL SURFECT (FERE), 1-0 TOTAL STERIOL (FERE), 1-0 TOTAL STERIOL (FERE), 1-0 TOTAL STERIOL (FERE), 1-0 STATE THE UNCORRECTED DATA MITE(12, 0+15) CALL STERIOL (12, -1, WEW', WC'LL)) STORE THE UNCORRECTED DATA MITE(12, 0+15) CALL STERIOL CALL STATE MITE(12, 0+15) CALL STATE MITE(12, 0+15) CALL STATE WITE(12, 0+15) CALL CLOSE CONTING CALL CLOSE CALL CLOSE	0047 ATEMPY=ATRNX-AISTX 0048 ATEMPY=ATRNY-AISTY C	C CONVERT BACK TO DEGREES AND DB	0049 ATEMP=SGRT (ATEMPX++2+ATEMPX++2) 0050 TPHASE=ATAN2 (ATEMPY, ATEMPX)	C DIVIDE BY PAD RESPONSE TO REMOVE IT FROM TRANSFER FLUXCY	ONSI DDRTEM=10. *ALOG10(ATEMP)	0052 APAD=. 05+FLOAT (IAMPAD(K))	0005 HTMM547051744705 144054705 HT440547051514(2)	0055 PTRANS=TPHASE-(.25#FLOAT(IPHA(K))+DEG)	0056 [F (PIRANS (GI PI) PIRANS#PIRANS+2.4PI	$\frac{1}{1000} = \frac{1}{1000} + 1$	0061 WRITE(12. *) IAMP4(K)	0042 WRITE(12, *) IPH4(K)	0063 400 CONTINUE	0064 CALL CLUSE(12)	00/ 01 UU COURT AND COURT AND COURT AND COURT AND COURT AND COURT	0057 CALL ASSIGN(12, -1, OLD', 'NC', 1)	0068 READ(12, *) FSTART	0069 READ(12,*) FENU 0070 READ(12,*) NP01NT	0071 D0 600 K=1, NP01NT	0072 READ(12,*) IAMP3(K)	00/3 KEHULLZA ITTAIN 0074 600 CONTINUE	0075 CALL CLOSE(12)	0077 CALL ASSIGN(12, , -1, 'OLD', 'NC', 1)	0078 READ(12,*) FSTART	00/7 REHULLZ * READ	0081 D0 650 K=1, NP0INT	0082 HEAU(12, *) IMP4(K) 0083 READ(12, *) IPH4(K)	0084 650 CONTINUE	0085 CALL CLOSE(12) 0084 700 DETUDN	0087 900 FORMAT(//1X SUBRNUTINE TRANS DETERMINES THE TRANSFER 1 (FINCTION OF THE NETWORK ANALYZEK AND CABLES.)	FORTRAN IV VUZ. 1-1 THU VY-TRUET OF V V V V V V VOL TO VOZI V V V V V V V V V V V V V V V V V V V	1 PRECISION PAD () AMONG SOR EDDMATIZITY FUTER THE NAME OF THE NEW FILE TO WANTAIN	1, ***, *UNCORRECTED TRANSFER FUNCTION DATA: *)	0090 903 FURMAT(//1X, ENTER THE RIVER OF THE TARK THE TAR	0091 904 FORMAT(//IX, ENTER THE NAME OF THE OLD FILE CONTAINING 0091 904 FORMAT(//IX, ENTER TEANSTER FUNCTION DATA ()	0092 905 FORMAT(//1X, 'ENTER THE NAME OF THE OLD FILE CONTRINING	1, 4, 'UNCUMPELIEU INHINSEEN FUNCTION MITTE	0094 END	FORTRAN IV STORAGE MAP FOR PROCRAM UNIT TRANS
D0 ZOO K-1. MEDINT CALL FWERFLIFFERD. 1.4 CALL FWERFLIFFERD. 1.4 CALL FWERFLIFFERD. 1.4 CALL FWERFLIFA-ZOOB IMPROVINE CALL FWERFLIFA-ZOOB IMPROVINE CALL FWERFLIFA-ZOOB IMPROVINE CALL FWERFLIFA-ZOOB IMPROVINE CALL FWERFLIFA-ZOOB CALL FWERFLIFA-ZOOB CALL FWERFLIFA-ZOOB METECLIZ-61 FEAT. 1.4 WITECLIZ-61 FEA	(•		,			•		,	Ũ	Ū	•							- •			•				•			4))	· =))	1
D0 200 K=1.WOINT FREO-FSTARTFELONT(K=1)*SIEP CALL SWERLINGT(K=1)*SIEP CALL SWERLINGT(K=1)*SIEP CALL SWERLINGT(K=1)*SIEN PH3(K)*(1P-2048) DP3(K)*(1P-2048) DD3 FSTART LA-2048) DD3 FSTART LA-2048) DD3 FSTART LA-2048) DD3 FSTART LA-2048) DD3 FSTART LA-2048) DD3 FSTART LA-2048) DD3 FSTART LA-2048) DSTART SWERLING FSTART L1. FSTART SWERLING FSTART L1. FSTART SWERLIZ1. WEW'. 'NC'.1L) STORE THE UNCORRECTED DATA WITE(12.*) FSTART WITE(12.*) FSTART MATTE(12.*) FSTART MATTA MATTE(12.*) FSTART MATTA MATTE(12.*) FSTART MATTE(12.*) FSTART MATTE(12.*) FSTART MATTA MATTE(12.*) FSTART MATTA MATTE(12.*) FSTART MATTA	ŀ		-	-			•-					-		•.	•				·			•	•	•	•			•							FAOE 003		r		
··· 0	+ (K-1)*STEP	FREQ. 1.4)	-2048) -2048)			SN(12, , -1, 'NEW', 'NC', 1,)			2 1) ECTADY	2.4) FEND	2. #) NP0INT	r=1, woint	12. •) 1AMP3(K)	12.0) [PH3(K)	.c .05E(12)			·ILE WITH CORRECTED DATA		903 056160013 -1 /NEH/ /NC/ 1 /	E(12, +) FSTART	5(12.*) FEND 5(12.*) NUMAT		CREATE CORRECTED DATA AND STORE	1 1415924	100 K=1, NP01NT	MERT TO DEGREES AND DB		T= Z5+FLOAT(IPH2(K))+DEG IST= A5+FEAAT(IAMP2(K))	N= 25+FLOAT(IPH3(K))+DEG KN= 05+FLOAT(IAHP3(K))+DEG	ENT ID MAILS	=10 ==(ADBIST/10) =10 ==(ATRTEN/10)		KEAL AND IMAGINARY PARTS	2.1-1 THU 09-946-79 01:03:11	ATRN+COS (PTRAN)	ATRN+SIN(PTKAN) AIST+COS(PIST)	AIST*SIN(PIST)	ACT ISOLATION FROM TRANSFER+ISOLATION

	÷								·	PAGE 003
÷ .	•		CULATE THE SYSTEM CLUTTER		NEW', 'NC', 1.)	K))+DEG K)+DEO P2(K)) P5(K))		ROM ANTEWNA RESPONSE OR FORM	TEMPX) TEMPX(A) ATEMPY(A) ATEMPY(A)) ATEMPA(K)) TTEMPA(-4U0-79 01: 12: 37 0LD^, ^NC^, 1)
- 1		WRITE(12, +) NPOINT K=1,NPOINT WRITE(12, +) IMM5(K) WRITE(12, +) IPM5(K) CONTINUE CALL CLOSE(12)	THIS SECTION WILL CA	P]=3.1415926 DEG=P1/180. TYPE 902	CALL ASSIGN(12,,-1,' WRITE(12,*) FSTART WRITE(12,*) FEND WRITE(12,*) NPCINT	D0 400 K=1, NPULN PISI= 25+1, 0AT (1PH2 (1 PANT= 25+1, 0AT (1PH2 (1 ADBATT= 05+F1, 0AT (1AM ADBANT= 05+FL, 0AT (1AM	AIST=10. +* (AURIST/10 AANT=10. +* (AURANT/10 AISTZ=AIST+C(S(PIST) AISTY=AIST+C(S(PIST) AANTX-AANT+C(S(PANT) AANTY-AANT+SIN(PANT)	SUBTRACT ISOLATION F ATEMPX=AANTX-AISTY ATEMPY=AANTY-AISTY CONVERT BACK TO PHAS	DIENT CONTRACT CALLENY, A PTEMP-ATAN2 (ATEMPY, A ATEMP-SQRT (ATEMPY, A ALEKEM=1C. ALGGIO (A) ARES=ADBTEMP-(055FLO ARES=ATEMP-(255FLO FF (PRES GT P1)PRES IF (PRES GT P1)PRES IFMS(K)=IFIX(4, PPRES IPMS(K)=IFIX(4, PPRES I	VO2.1-1 THU 09 CALL CLOSE (12) 00 TO 700 TYPE 904 CALL ASSIGN(12.,-1,' READ(12) FSTART READ(12) FSTART READ(12) FEND READ(12) FOUNT DO 600 K-1, MPOINT
		. 90 30 9 8 9 9 9 8 9 9	00000	50 2 2 2	2020	86823	04040 0	000 000 22		72 72
	•	003 003 003 003 003 003 003 003 003 003		003 003 003	88888 88888 88888	000000 0000000	888888 888888 888888888888888888888888	000)
•	()) j j j
-		PAGE 00	D ROCM 1ENT 1EN HE 12 CM CHEATED .	ICTERISTIC			RENCE	itep, NPOLNT, NSANP 4, IPH4 IPH6	42 (512), IANF3 (512) 5 (512), IPH5 (512) 4	FAQE 002
		1-1 THU 09-AUG-79 01:12:37 E ANTEN 4	UUTINE WILL MEASURE THE ANTENNA AN 15TICS EXPLICITY. FIRST A MEASURER THE SYSTEM CHARACTERISTICS IS TAK AT EACH FREQUENCY IN ADDITION T 17FP AND ANTENNA CHARACTERISTICS A	CABLE RESPONSE AND ISULATION CHAR INCAR ANALYZER ARE REMOVED. THE TRANSFER CHARACTERISTIC OF T	TINE TRANS 1)/R BECETTER STEAM ERDM CONCE COLOR	CLUTTRA CLUTTRA TRANSFER FUNCTION OF N/A + CABLE REFERENCE CHANNEL SIGNAL RECEIVED REGAURMENT => M3 TECEIVED RECEIVED		JA/IAMPAD, IPHA, IFLAG, FSTART, FEND, S TA/IAMP2, IPH2, IFLAG2, NSAMP2 NSA/IAMP3, IPH3, IFLAG3, NSAMP3, IAMP (A/IAMP3, IPH5, IFLAG4, NSAMP4, IAMP6,	иней((512), ГРНА(512), ГАМР2(512), ГР Н3(512), ГАНРА(512), ГАМР2(512), ГАМР ИР6(512), ГРН6(512) 1 100, 100, 500 1 1 57АКТ, 1, 4) 1, МРОПИ 1, МРОПИ 1, КССЛТ(K-1)=STEP	(1-2048) 2048 2048 2048 10, FSTART, 1, 4)
		U VOZ SURROUTINE VERSION 1	THIS SUBRC CHARACTERI INCLUDING AND STORET RYMM CLUT	MEN THE (OF THE NET B/R	SUBROUT T/R=(C+0=) MMERE Y-	1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1	X/R= (M3 ALL MEA FHANHEL	COMMON/PAE COMMON/151 COMMON/TRA COMMON/TRA	INTEGER IF INTEGER IF INTEGER IF INTEGER IF INTEGER IF INTEGER IF PAUSE DO 200 Y-1 EREGEFSIME	I ANDE (K) = I ANDE (K) = I COLL INEE COLL SHEE VO2 1-1 VO2 1-1 VO2 1-1 CALL ASSIG
·		FORTRAW 01 0 0 0		0000	00000			0 01 20 20 20 20 20 20 20 20 20 20 20 20 20	252 20 20 20 20 00 00 00 00 00 00 00 00 00	23 23 23 24 24 24 24 24 24 24 24 24 24 24 24 24

? ?	NLT SLOSE	TYPE 903	READ(12, +)	TEAD(12. +)	D0 550 K=1	READ(12. +)	CONTINUE	CALL CLOSE RETURN		FORMAT (//1	1, 'CHARACTE 2///1X, 'CON	FORMAT(//	FORMAT (/	1 NAME. 1) FORMAT (7/2)	1 NAME. ') STOP	END	/ STORAGE			000110	001000	000104	000160	000140	000020	000070	¥91000		ck /Pada /.	E OFFSET	000000	004002	
IPHS(K)	(12)	M(121.	FSTART	FEND NPOINT	NPOINT	[PHA(K)		(12)		X ^ SUBROUT	RISTICS /	¢'.'ENTE	• · · · ENTER	\$'.'ENTER			MAP FOR			AANTX	ADBIST	AISTX	ATEMP	DEG	41	Ы	PTEMP		\$12E = (NAME	IPHA	FEND	
~						~				TINE AN	AND // 1 X	RTHEA	A THE A	₹ THE &			PROGRA	3710		R=4	R=4	R*4	R=4	R=4	1+2	R=4	R*4		004022	TYPE	1+2	R=4	
	•				•					TEN DETERM	. ANTENNA EM IN ITS	NTENNA SYS	NTENNA CLU	NTENNA SYS			M UNIT ANT	1 017000		000124	000074	000114	00150	090000	000052	000024	000144		(1033. MD	OFFSET	002000	004006	
	•									INES TH	FINNL FI	TEM RES	TTER DA	TEM RES			EN 21			AANTY	ADBTEM	AISTY	ATEMPX	FREQ	¥	PIST			RDS)	MAME	IFLAG	STEP	
										E ANTEN	UKM ()	PONSE 1	TA FIL	PONSE ((auau		ITPE	F*4	R#4	R#4	R=4	R#4	1+2	R.4				TYPE	1+2	R=4	
						•				, eni		1EN FILE	E NAME. *)	ארם גזרב					G-1961	000130	000154	021000	000134	000044	000042	600044				OFFSET	004000	004012	
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NIOdN	8	8	NAME	I AMP2	NSAMP		COMMO	NAME	IAMP3	MARN	COMMO	NAME	IAMPS	NSAMP	LOCAL	NAME	I AMP2	DAMPI	IAMES	IAMPI IPHA	IPH2 IPH3	e I	SHdI		FORTF SUBRC	NAME	AL 001	FL0A1	SHEEF				
T 1+2	1		TYPE	1+2	2 I#2	l 1	N BLOCK	TYPE	I#2	3 1+2	N BLOCK	TYPE	1+2	4 1+2	AND CO	TVP	1#2	2 1	142	1 1 1 1 1 1 1	1+2 1+2	14 1	1+2 1+2		IN IV	TYPE	0 R=4	R*4	R.4				
004016		OCK /ISTA	OFFSET	000000	004002		/TRANSA/	OFFSET	000000	004002	ANTA /	OFFSET	000000	004002	HINN ARRA	E SECTI	ISIO	TRANS	ANTA	E PADA	ISTA TRANS	2 TR	ANTA		STORAC FUNCTION	NAME	ASSION	IFIX					
NSAM		1 /	NAME	IPH2			', SIZE	NAME	ена	- I AMP	', SIZE	NAME	CH41	INHP	WS:	ON DEFS	0000	50 0000 50 0000	0000	0040	M 0020	ANSA 0	0020		JE MAP F 4S. STAT	TYPE	₽÷Я	I+2					
• I+2		2E = 00	TYPE	1+2			- 01000	TYPE	1+2	4 I+2	• 01000	TYPE	I+2	6 I#2		ET	200 005 00	00 002	00 002	00 00 00 00 00 00 00	800 800 800 800	06004	00 04 002 002		OR PROG	NAME	ATAN2	PHAMP2					
004020		004 1 10	OFFSET	002000			1 (2050.	OFFSET	002000	00400	1 (2050.	OFFSET	002000	004004				000 51	15 000	000 - 21 000 - 21	000 - 21 000 - 21	002000	000 C 21		RAM UNIT	TYPE NA	Re4 CL	Re4 SI					
		Z6. MUKI	MAN	IFL			HORDS)	MAN	IFLA	HdI	WORDS)	MAN	IFLA	IPH		HIG	2) (51)	2.) (51:	2.) (51	50 (21)	2) (3)	512.)	2) (31)		ANTEN SOR-DEF	ME TY	OSER	ά z					
	5	5	TYPE	02 1+2				TYPE	63 I+2	I+2		TYPE	04 I+2	1+2		SNOISNE			2	55	ລລ	(512)	ລລ		INED FU	¥.	5	50					
			5	004				CFF31	00400	00400		CFFSE	00100	00600											ACT I CHIS	ы. з	ě	41 R.					

<pre>IV vo2 1-1 PROJRAM SPHERE VERSION 1.2 VERSION 1.2 THIS PEPOGRAM WILL TAKE EXPERIME AND CORRECT IT FOR THE SYSTEM RE AND CORRECT IT FOR THE SYSTEM RE AND CORRECT IT FOR THE SYSTEM RE AND CORRECTED AND UNCORRECT AND THE HIGH RESOLUTION CRT AND THE SECTION AND THE POSITIAN ACCEPT 10001 AND THE FORMATION AND THE AND ACCEPT 10001 AND THE AND THE AND THE AND THE AND ACCEPT 10001 AND THE AND THE AND ACCEPT 10000 AND ACCEPT 10</pre>		(0035 HRITE(IUNIT, 906)	FACE 001 0037 IND1=K 0038 AMPLIT=FLOAT(OBUA(K))=.05	0039 PHASE=FLOAT (OBLP (K))+ 25 0040 FREQ=F START+FLOAT (INDI-1)+STEP	VIAL DATA FOR THE SPHERE 0041 MRITECIUNIT,907)K,FREG.AMPLIT,PHASE SPONSE. IT WILL FINT 0042 277 CONTINUE FED DATA AND FINNLY DISPLAY 0043 250 CONTINUE THIS SECTION WILL STORE THE CORRECTED DATA FOR	E IF DESIRED ANN.TITICAL 0044 TYPE 802 T IN A FILE IT WILL ALSO 0045 ACCEPT 1000.A FILE AND ALSO DISFLAY IT ON 0045 IF (A E0 'N') 60 TO 400	0049 TYPE 801 0049 CALL ASSIGN 201.'NEW'.'NC'.) 0050 WRITE (20F)FTAT 0051 WRITE (20F)FTAT 0052 WRITE (20F)FDD	0033 00 310 K=1, K=0 NT 40(512), NPOINT, NSAMF2 0054 WRITE(20, +)08.40(K) UTP(512), NPOINT, NSAMF2 0056 310 CONTINUE 0057 CALL CLOSE(20)	C THIS SECTION WILL CORRECT FOR RANGE	0058 400 TVPE 914 0059 ACCEPT +, IRFLAD	ECTED SPHERE DATA FROM COTED SPHERE DATA FROM Only Type 913 Dist	0062 CALL RANCOR 0063 TYPE 916	0064 ACCEPT 1000,A 0065 IF (A ED 'N') GO TO 470	0067 WRITE(IUWIT, 913)DIST 0068 WRITE(IUNIT, 911)	0069 WRITE(JUNIT, 905)FSTART, FEND, STEP, NPOINT 0070 WRITE(JUNIT, 906)	>, NPOINT . 0071 D0 410 K=1, NPOINT . 0072 INDI=K	0073 AMPLIT=FL0AT(0BUA(K))+. 05 0074 PHASE=FL0AT(0BUP(K))+. 25	0075 FREQUESTART+FLOAT(IND1-1)+STEP 0076 WRITE(IUNIT,907)K,FREQ,AWPLIT,PHASE	ANSE 0077 410 CONTINUE C THIS SECTION WILL STORE THE RANGE CORRECTED DA			0079 470 TYPE 901 0079 ACCET 1000,A	PAGE CO2 0082 TYPE 800 0083 CALL ASSIGN(20,-1, 'NEW', 'NC',) 0083 CALL ASSIGN(20,-1, 'NEW', 'NC',)
· · · · · · · · · · · · · · · · · · ·				PROGRAM SPARE	VERSION 1. 2	THIS PROGRAM WILL TAVE EXPERIMEN AND CORRECT IT FOR THE SYSTEM RE BOTH THE CORRECTED AND UNCORRECT IT ON THE HIGH RESOLUTION CRT.	ALSO IT DESIRED IT WILL GENERATE DATA FOR THE SPHERE AND STORE I PRINT THIS DATA IF DESIRED TO A THE HIGH RESOLUTION CRT MONITOR.	COMMON/PANGE1/DIST, IKFLAG, IUNIT COMMON/PANGE1/DIST, IKFLAG, IUNIT COMMON/OBJ/OBJA, OBJP, NSAMP, CLUTA, C	INTEGER OBJA(512), OBJP(512), TRA INTEGER TRANP(512), CLUTA(512), CI BYTE A	TYPE 400	TYPE 912 ACCEPT +, IUNIT	THIS SECTION WILL GENERATE CORRE EYPERIMENTAL DATA	TYPE 902	CHLL SYSIME STEP=(FEND-FSTART)/FLOAT(NPOINT) CALL SPHEAT	TYPE STATT	HE (10 11, 504 10 220 WRITE(1011, 504)	WPITE/IUNIT, 905) FSTART, FEND, STEP WRITE(TUNIT, 905)	D0 210 k=1.NP0INT AMPLIT=FLOAT(OB.46(K))+ O5	PHASE=FLOAT(0BJP(K))+ 25 FREG=FSTART+FLOAT(K-1)= 25	WEITE (IUNIT, 907)K, FREQ, AMPLIT, P. CONTINIE	MERE LORREC MERE LIST THE IMPORTION CONCOS I	TYPE 915	ACCEPT 1000. A	V VO2 1-1 IF (A.E0 'N') 00 TO 290 MRITE(IUNIT,909)

;-FORMAT(//1X, ***** SUBROUTINE SYSINP READS THE SYSTEM RESPONSE INTEGER TRANA(512), TRANP(512), CLUTA(512), CLUTP(512), NP0111T FORMAT(///%:/ENTER THE TRANSFER FUNCTION FILE NUME. ') FORMAT(///%', 'ENTER THE ANTENNA SYSTEM FUNCTION FILE NAME COMMON/SYSTA/TRANA, TRANP, CLUTA, CLUTP, FSTART, FEND. NPOINT OFFSET OFFSET 004000 1000010 THIS SUBROUTINE WILL READ IN THE SYSTEM RESPONSE FILES AND PLACE THEM IN A COMMON BLOCK TO BE PASSED TO OTHER RAVITINES. PAGE 001 TYPE TYPE B. HORDS) 1+2 * -----SIZE----- DIMENSIONS CLUTA NAME FEND MAN WORDS) FORTRAN IV STORAGE MAP FOR PROGRAM UNIT SYSINP LOCAL VARIABLES, PSECT \$DATA, SIZE = 000020 (ASSIGN(12. . -1. 'OLD', 'NC', 1) COMMON BLOCK /SYSTA /, SIZE = 010012 (2053. ASSIGN(12, , -1, 'OLD', 'NC', 1) TYPE OFFSET OFFSET 002000 0000010 VO2. 1-1 THU 09-AUG-79 01:42:40 SUBROUTINE SYSINP 1=2 FSTART R#4 READ(12, *) TRANA(K) READ(12, *) TRANP(K) READ(12, +) CLUTA(K) READ(12, +) CLUTP(K) SECTION OFFSET TRANP READ(12, *) FSTART READ(12, *) FEND READ(12. *) NPOINT **WAR** NAME READ(12. +) NP0INT READ(12. +) FSTART DO 100 K=1, NPUINT DO 200 K=1, NP01NT 2223 CL0SE(12) CALL CLOSE(12) 1 FILES ****') LOCAL AND COMPON ARRAYS: READ(12, +) READ(12. +) 900 CONTINUE ŝ 501 CONTINUE TYPE OFFSET 000012 OFFSET 000000 000900 010010 RETURN CALL TVPE GALL TΥΡΕ END TYPE TYPE 1+2 1+2 1#2 NPOINT 1+2 006 8 902 80 901 0000000 υu Q **FRANA** CLUTP 0003 0026 MAN W 0028 MARK 0002 0027 ORTRAN IV 000 ? 3) ; [COOPECTED DATA ') FORMAT(///, 'SENTER THE FILE NAME FOR THE COORECTED DATA. ') FORMAT(///, 'SDO YOU WANT TO STORE THE CORRECTED DATA ('Y OR N) FORMATI// THIS SECTION WILL GENERATE CORRECTED SPHERE DATA FROM // EXPERIMENTAL DATA) FROM // EXPERIMENTAL DATA) FROM // SPHERE DATA (Y OR N): ') FORMATI///// ****** UNCORRECTED DATA *****') FORMATI///// ****** UNCORRECTED SPHERE DATA *****') FORMATI///// ****** UNCORRECTED SPHERE DATA *****') FORMATI////// ****** UNCORRECTED SPHERE DATA *****') PAGE 004 FORMAT(///SPRINT SPHERE DATA CORRECTED FOR RANGE (" OR N) AAT ////// EXPERIMENTAL SPHERE DATA CORRECTED FORMAT(/// % ' ENTER LOGICAL UNIT NUMBER FOR OUTFUT (7=) FORMAT(///// CALCULATED RANGE TO THE TARGET = ", 1FG15.7. OFFSET 000046 000042 FORMAT(////.**** PROGRAM SPHERE ****') FORMAT(///.*6D0 YOU WANT TO STORE THE RANGE CORRECTED [DATA (Y OR N) ? *) FOPMAT(//1X, 75'**)/12X, POINT #'.5X, FREQUENCY, 7X, APPLITUGE DB '.5X, PHASE REGRES/1X, 75'*//) FOPMAT(///1, 17, 5X, 1P015, 7) 3X, 1P015, 7) FOPMAT(////// ***** CORRECTED EXPERIMENTAL SPHERE TYPE (SORDS) ***** 1.4 FORMAT(///, "SENTER THE FILE NAME FOR THE RANGE PHASE MAN FREQ FORTRAN IV STORAGE MAP FOR PROGRAM UNIT SPHERE LOCAL VARIABLES, PSECT SDATA, SIZE = 000066 (27) STOP ****** END OF PROGRAM ****** THIS SECTION WILL DISPLAY THE DATA OFFSET 960000 4E0000 TYPE 1•2 AMPLIT R44 WRITE (20, +) 08-14 (K) WRITE (20, +) 08-19 (K) CONTINUE CALL CLOSE (20) FOR RANGE ++++/). WRITE(20. +) NP01NT D0 460 K=1, NP01NT FORMAT STATEMENTS W HRITE(20. +) FEND DATA #####^) OR NU. FORMAT(A1) V02. 1-1 I METERS) (TERMINAL) CONT DAJE OFFSET 000026 000052 000000 Q 2 TYPE 8 3 VI NERTECT 5 1 619 10 516 Ş 8 8 932 ş 8 888 800 ŝ 110 912 ŝ ŝ 600 015 5000 6600 F 0102 1900 2600 6000 8000 510 233 901ŭ 0107 0103 oùli 0110 91112 6110 Tour ŝ \$600 8000 010 ¥ 97EP 1021 2002 9800 **C** £ ſ 0 0 0 C ۱

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FORI	RAN IV VO2 1-1 THU 09-RUG-79 01:43:34 FAGE 001	9999	L	- - 2	OINT: ')						
į			FORTRAN		STORAGE	MAP FOG	PROGR	AM UNIT S	PHDAT	P. MORE	ŝ
	THIS SUPPROUTINE WILL CALL OBUDAT AND OBTAIN ONE LINE OF DATA FUR THE SPHERE IF NEW DATA IS SPECIFIED OTHERNISE An OLD DATA FILE WILL BE READ. IN THE CASE OF NEW DATA THIS ROUTINE WILL WRIE THE NEW FILE WITH THE DESIRED NAME		Nore A	TYPE .	OFFSET X00012	N N N	TYPE I+2	0FFSET 000014	NAK	TME	OFFS
903 803 80	COMMON/OBJ/OBJA, OBJP, INSAMP2 COMMON/SYSTA/TRANA, TRANP, CLUTA, CLUTP, FSTART, FEND, 1201NT		COMMON	BLOCK	10BJ /1	size =	004002	(1025.	, JORDS)		
10 10 10 10 10 10 10 10 10 10 10 10 10 1	INTEGER 08.M(512), 08.P(512) INTEGER TRAMA(512), TRAMA(512), M INA/54.01, M INTEGER 18.MA(512), TRAMA(512), M		SHAME N	TYPE	DEFSET	AMEN DB. ID	TYPE 1#2	OFFSET	NAME	TYPE	OFFSE 00400
9000 1000			COMMON .	B OCK	SYSTA /.	SIZE #	010012	(2053. 1	ORDS)		
0007	TYPE 900		NAME	TYPE	DFFSET	NAME	TYPE	OFFSET	MAN	TYPE	OFFS
5000 6000	TYPE 901 ACCEPT 700. A		TRANG	1+2	000000	TRANP	I+2	002000	CLUTA	1+2	00400
812 812 812	IF(A EQ. 'N') GO TO 300 TYPE 905		CLUTP	1#2	000000	FSTART	R#4	010000	FEND	4.8	01000
0013 0014	ACCEPT *. NSAMP2 TYPE 902		NPOINT	1#2	010010						
88 1 2 1 2 1 2 1 2 1 2 1 2 1 2 1 2 1 2 1	PALISE		LOCAL A	ND COM	ION ARRAYS						
817 819	TYPE 903			1005	011030	1.0000	•		DIMEN.		
613 0020	HATTE(12, e) FSTART HATTE(12, e) FSTART HATTE(12, e) FEND			1+2 1+2	SYSTA SYSTA SYSTA	000000	00200	00 (512) (512)		
0021 0022	MRITE(12. +) NPOINT BD 200 F-1. NP0INT		08.UA 08.UP	I+2 I+2	080 080	000000	00200	00 (512 00 (512) (512)) (512)		
6023 00 24	WFITE(12, +) OB.MA(K) WRITE(12, +) OB.UP(K)		TRANA TRANP	1+2	SYSTA SYSTA	00000 002000	00 00 00 00 00 00 00 00 00 00 00 00 00	00 (512 00 (512) (512)) (512)		
5255 5255	00 CONTINUE FALL CLASE(12)		SUBROUT	INES, 1	-UNCTIONS.	STATEM	ENT AN	D PROCESSI	OR-DEFIN	ED FUNC	SNOT
0023 3	60 T0 500 0 TYPE 904		MANE	TYPE	NAME TY	JN 3d	L M	VPE NAM	E TYPE		TVI
6029	CALL ASSIGN(12.,-1, COLD', NC', 1,)							1			
6031 6031	READ(12.*) FSTART READ(12.*) FEND		ASSIGN	4 4	CLOSE			4			
25.00 (0.32	PEAD(12.*) NPOINT Po. 400 X+1 MEDINT										
100	READ(12, +) 08.4(K)										
883	READ(12, +) OBJP(K) A CONTINE										
0037	CALL CLOSE(12))									
883	0.) RETURN GAI FORMATCAI)										
0000	UN FORMATIONIX, ***** SUBRIJUTINE SPHDAT OBTAINS SWEPT FREQUENCY)									
(mai +	DI FORMAT(///%',/TAYE NEW DATA (Y OR N) ?') 02 FORMAT(///%',SET UP THE SFHERE, HIT RETURN TO CONTINUE')										
100100)		•							
- CAO?	IV VOZ 1-1 THU 09-RUI5-79 01:43:34 FAGE 002 03 FOPMAT(///s*, 'ENTER NEW SPHERE INTA FILE NAME' '))									
	04 FORMAT(///\$/, 'ENTER OLD SPHERE DATA FILE NAME: ')										

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-	FORTRAN IV VO2. 1-1 SUN 07-OCT-79 00: 11: 12 0001 SUBROUTIME CORREC	C THIS SUBROUTINE WILL TAKE THE DATA THAT WAS TAKEN I C SPHEAT AND CORRECT IT WITH THE DATA FROM SUBROUTIN C SVSINP, WHICH IS THE SVSTEM RESPONSE.	C VERSION 1. 1 7-007-79	COO2 COMMON/OBJ/OBJA, OBJP, NSAMP2 COMMON/SVST6/TERNA TRANE CLITA CLITA ESTART SEM #		0004 INTEGER 08-0(512), 08-0 (512), TRANA(512), TRANP(512:	0005 INTEGER CLUTA(512), CLUTP(512), NPOINT, NSAMP2 C	iningangangangangangangangangangangangangan	C MUST NOT !!! BE CORRECTED FOR THE SYSTEM TRANSFER FI		C OF THE SYSTEM IN ARRAYS TRANALAHPLITUDE) AND TRANG	,	C C Piers 141502A	0007 [FEG=P]/180		0010 CLPH=FL0AT(CLUTA(K))+. 05 0010 CLPH=FL0AT(CLUTF(K))+. 25+DEG	0011 TDPAMP=FL0AT(TRANA(K))=.05 0012 TPPH=FL0AT(TRANF(K))= 25=FFG	0013 00BAMP=FL0AT(0B.A(K))+ 05 0014 0B.PH=FL0AT(0B.P(K))+ 25+7E0		C CDBAMP=/CLUTTER IN DRM C CLPH=/CLUTTER PHASE IN RADIANS	C TUBAMP=>TRANSFER CHARACTERISTIC IN DRM C TRPH=>PHASF OF TRANSFER CHARACTERISTIC IN BADIANS	C ODBAMPENDLECT AMPLITUDE IN DBM	C OBJEMENT OF OF OFFICIAL DATE IN ALLIANTS C OBJEMPENOESECT AMPLITUUE IN ALLIANTS C OLAMPENCE AMPLITICE AMPLITICE IN ALLIANTS		0015 CLAMP=10. ++(CDRAMP/10) 0016 OB.AMP=10. ++(CDRAMP/10)		C CLCOS=>REAL PART OF CLUTTER	C CLESIN=PIERGINGHY PART OF CLUTTER C OBJOOS=PREAL PART OBJECT DATA	C OBJSIN=>IMARY PART OF OBJECT DATA C #************************************
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ς		OVER THE FREG		D. NPOINT		512)										hetter		000022				OFF3ET	004000		OFFSET	004000	010004	·)	
ζ		BUECT OVER THE FREG		RT, FEND. NPOINT		CLUTP(512)									, suenes			I=2 000022				TYPE OFFSET	2 1*2 004000		TYPE OFFSET	1+2 004000	R#4 010004	·)	- 2
ζ		OR AN OBJECT OVER THE FREG		TP, FSTART, FEND. NPOINT		A(512), CLUTP(512)									BUDAT 11 LUDEDS)			IP I=2 000022			DRDS)	NAME TYPE OFFSET	NSAMP2 1+2 004000	ORDS)	NAME TYPE OFFSET	CLUTA 1+2 004000	FEND Re4 010004	· · · · · · · · · · · · · · · · · · ·	-
ζ		IN DATA FOR AN OBJECT OVER THE FREQ	*2	CLUTA, CLUTP, FSTART, FEND, NPOINT	2) . MPOINT	512), CLUTA(512), CLUTP(512)	VPOINT >								MI UNIT OBUDAT OODODA (11 MODA)	TECCET LANE TURGET		000020 IP I=2 000022	00009		1025 MORDS)	RFSET NAME TYPE OFFSET	02000 NSAMP2 1+2 004000	2053. MORDS)	RFSET NAME TYPE OFFSET	02000 CLUTA 1+2 004000	10000 FEND Re4 010004	· · · · · · · · · · · · · · · · · · ·	
ζ		L OBTAIN DATA FOR AN OBJECT OVER THE FREQ	zanosu a	TRAMP. CLUTA. CLUTP. FSTART, FEND. NPOINT	IB.P.(512), NPOINT	TRAMP(512), CLUTA(512), CLUTP(512)	FLOAT (NPOINT)		-1) +STEP	, 4) Samp2)					PROGRAM UNIT OBUDAT SIZE = 000024 (11 MARTS)	TYDE DECCT LAME TYDE DECCT		1+2 000020 IP I+2 000022	R=4 000006		04002 (1025 MORDS)	TYPE OFFSET NAME TYPE OFFSET	1*2 002000 NSAMP2 1*2 004000	10012 (2053. MORDS)	TYPE OFFSET NAME TYPE OFFSET	1+2 002000 CLUTA 1+2 004000	R++ 010000 FEND R++ 010004	· · · · · · · · · · · · · · · · · · ·	
ζ		THE WILL OBTAIN DATA FOR AN OBJECT CVER THE FREG	B.IA. OB.P. NSAMP2	1/TRANG, TRAND, CLUTA, CLUTP, FBTART, FEND, NPOINT	(512), 08.PC(512), NPOINT	M(512), TRANP(512), CLUTA(512), CLUTP(512)	START)/FLOAT (NPOINT)		FLOAT (K-1) +STEP	. FREU. 1.4) 14. IP. NSAMP2)	040	EE100T 1 4)			MAP FOR PROGRAM UNIT OBLUAT BDATA, SIZE = 000024 (11 MARAS)	VAME TVDE DEFCET LANE TVDE DEFCET		IA I+2 000020 IP I+2 000022	STEP R=4 000006		12E = 004002 (1025 MORDS)	NAME TYPE OFFSET NAME TYPE OFFSET	08-JP I #2 002000 NSAMP2 I #2 004000	12E = 010012 (2053. MORDS)	WANE TYPE OFFSET NAME TYPE OFFSET	TRAMP 1+2 002000 CLUTA 1+2 004030	-57ART Re4 010000 FEND Re4 010004		
ζ		SUBPOUTINE WILL OBTAIN DATA FOR AN OBJECT CVER THE FREQ	W/OBJ/OBJA, OBJP, NSAMP2	N/SYSTA/TRANA. TRANP. CLUTA, CLUTP, FBTART, FEND, NPOINT	KR ORJA(512),OBJP(512),NPOINT	KR TRAMA(512), TRANP(512), CLUTA(512), CLUTP(512)	C FEND-FSTART)/FLOAT (NPOINT)		651A01+FLQAT(K−1)+STEP 6urcey • Eaco • • •	2455711.74544.1.4) Phamp2(1A,1P,NSAMP2)	F)=16-2048	NJE Suesdii Eetadt (1)			TORAGE MAP FOR PROGRAM UNIT OBJOAT PSECT BDATA, SIZE = 000024 (ET MANE TVDE DECCET LANE TVDE DECCET		14 IA I+2 000020 IP I+2 000022	12 STEP R+4 000006		/, SIZE = 004002 (1025 MORDS)	ET NAME TYPE OFFSET NAME TYPE OFFSET	00 08.JP 1+2 002000 NSAMP2 1+2 004000	TA /, SIZE = 010012 (2053. MORDS)	ET NAME TYPE OFFSET NAME TYPE OFFSET	00 TRAMP 1+2 002000 CLUTA 1+2 004000	00 FSTART R+4 010000 FEND R+4 010004		
ζ		THIS SUBMOUTINE WILL OBTAIN DATA FOR AN OBJECT OVER THE FREQ	COMPONYOBJYOBJA, OBJP, NSAMP2	COMPON/SYSTA/TRANN, TRAND, CLUTA, CLUTP, FSTART, FEND, NPOINT	INTEGER OBJA(512), OBJP(512), NPOINT	INTEGER TRANA(512), TRANP(512), CLUTA(512), CLUTP(512)	STEP=(FEND-FSTART)/FLOAT(NPOINT)	DO 100 Fat. NPOINT	FPED=FSTAP1+FLOAT(K-1)+STEP CAL SUFFAC: FAFC	CALL PHERFILLFREUIL4] CALL PHAMP2(1A, 1P, NSAMP2)	08.JA(Y)=1A-2048 06.JP(Y)=1P-2048	CONTINUE Cont succepti fotodi i At	RETURN		STORAGE MAP FOR PROGRAM UNIT OBLUDAT BLES PSECT SDATA, SIZE = 000024 ()1 LANDRS)	AFFACT MANNE TVDE AFFACT MANNE TVDE AFFACT		000014 IA I+2 000020 IP I+2 000022	000012 STEP R+4 000006		# /0B-) /, SIZE = 004002 (1023 MORDS)	OFFSET NAME TYPE OFFSET NAME TYPE OFFSET	000000 0B-JP 1+2 002000 NSAMP2 1+2 004000	k /svsta /, size = 010012 (2053. MORDS)	OFFSET NAME TYPE OFFSET NAME TYPE OFFSET	000000 TRAMP 1+2 002000 CLUTA 1+2 004000	006000 FSTART Re4 010000 FEND Re4 010004	010010	
ζ.		C THIS SUBPOUTINE WILL OBTAIN DATA FOR AN OBJECT OVER THE FREQ	C COMPONY OB JYOB. 14, OB. 14, NSAMP2	COMMON/SYSTA/TRANA, TRAND, CLUTA, CLUTP, FETART, FEND, NPOINT	C INTEGER ORJA(512), OBJP(512), NPOINT	INTEGER TRANA(512), TRANP(512), CLUTA(512), CLUTP(512) C	STEP=(FERD-FSTART)/FLOAT(NPOINT)	DO 100 Felverstration	FPE(n=KSTAPT+FL()AT(K=1)+STEP	CALL PHERFILLEREW.1.4) Call Phamp2(1A, 1P, NSAmp2)	08 JA (Y) = 1 A-2048 06.97 (Y) = 1 P-2048	100 CONTINUE CMI SUESPILEETADT 1 AN	RETURN	END	NAV IV STORAGE MAP FOR PROGRAM UNIT OBUDAT VARIABLES. PSECT BDATA. SIZE = 000024 (11 MADAS)	TVDE AFFGET MANE TVDE AFFGET LAME TUDE AFFGET		Re4 000014 1A 1=2 000020 1P 1=2 000022	I+2 000012 STEP R+4 000006		M BLOCK /0BJ /, SIZE = 004002 (1025 MORDS)	TYPE OFFSET NAME TYPE OFFSET NAME TYPE OFFSET	1+2 000000 08.JP 1+2 002000 NSAMP2 1+2 004000	M BLOCK /SVSTA /* SIZE = 010012 (2053. MORDS)	TVPE OFFSET NAME TVPE OFFSET NAME TYPE OFFSET	- 1+2 000000 TRANE 1+2 002000 CLUTA 1+2 004000	1=2 006000 FSTART R+4 010000 FEND R+4 010004	T 1+2 010010	AND COMMYN ASSAVE

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SCALE R+4 004090 SCALE R+4 004090 COMPON BLOCK /PANGELY. SIZE = 000010 (4 MORDS) NAME TYPE OFFSET NAME TYPE OFFSET NAME TYPE OFFSET NAME TYPE OFFSET NAME TYPE OFFSET NAME TYPE OFFSET NAME TYPE OFFSET NAME TYPE OFFSET NA	P	¥•4 ₩	004054	Id						r	υ	4 ÷ Č	900000	DEOR	••	ш 8	ELTA Re	1 000034
COMMON BLOCK / PANJELY, SITE = 000010 (4, MURLS) MANE TYPE OFFSET NAME TYPE OFFSET OFFSET NAME TYPE OFFSET NAME TYPE OFFSET OFFSET OFFSET OFFSET NAME TYPE OFFSET OF	SCA	F R.4	004050	•					211100		4	(•)	0000	KTFLTA R	++ 0000	30 K	BCO Re.	300024
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HOME TYPE OFFSET NAME TYPE OFFSET NAME TYPE OFFSET C RAD R+4 000044 THOPT R+4 000040 DIST R+4 0////////////////////////////////////	ŧ	ON BLOD	K /PANGE1/	= 321S	010000	•	HORDS))	PHASE	4.2	950000	PHNEW R	0000	44 44	Ł	100000
DIST R+4 0/0000 INFLAD I+2 000004 IUNIT I+2 000006 COMMON BLOCK /0BJ /, SIZE = 004002 (1025. MORDS) COMMON BLOCK /RANGEL/, SIZE = 000010 (4. MORDS)	NOW	TYPE	OFFSET	NAME	TYPE	OFFSET	MAME	. TYPE	OFFSET	J	RAD	R 4 4	000044	THOPI R	e4 0000	0		
COMMON BLOCK /OBJ /, SIZE = 004002 (1025. MORDS)	121Q	R.4	000000	IRFLAD	1+2	00000	TINUIT	1+2	00000	?								
	6	ION BLOCK	X /080 /.	\$12E =	004002	(1025.	(SORDS)			•	COMIC	N BLOCK	/RANGE1/.	812E = 00	0010 (4. WORD	â	
NAME TYPE OFFSET NAME TYPE OFFSET NAME TYPE OFFSET NAME TYPE OFFSET NAME	3 and	TYPE	OFFSET														1	

~	0.2 1-1 PROSPAM SPHERZ		038 040 040 040	- 4 - 3	YPE 503 CCEPT 1000.A F 14 EQ 11/10 00 10 220
	AEPSION 1 OF 100-10		640 440	33	RITE (1UNIT, 903)FSTART, FEND, STEP, NPOINT RITE (1UNIT, 906)
	THIS PPOGPAM WILL TAYE EXPERIMENTAL DATA FOR THE SPHERE AND CORRECT IT FOR THE SYSTEM RESPONSE. IT WILL PRINT BOTH THE CORPECTED AND UNCORRECTED DATA AND FINNALY DISPLAY IT ON THE HIGH RESALUTION CRT	600000	00460 0440 00446 00446 00446 00446 00446 00466 00000000		0 210 K=1.NPOINT WPLIT=ELOAT(OBJ0K(K))=.05 \\ASEFELOAT(OBJ0K)\S)=25 \\RECEFETART+FLOAT(K-1)=STEP \\RECEILINIT,907)K,FREQ.AMPLIT,P\\ASE
	ALSO IT DESIPED IT WILL GENERATE IF DESIRED AWALVITICAL Data for the structe and store it in a file. It will also PPINT this data if desired to a file and also display it on the high pestution of thomitor	• (00088 88	00 I	CONTINUE All Correc Here List the Corrected Sphere Data
	COMMON/PANGE1/D151, IRFLAQ, IUNIT COMMON/OBJ/OBJA, OBJP, NSAMP2 COMMON/SYSTA/TRANA, TRANP, CLUTA, CLUTP, FSTART, FEND, NPOINT	000000	052 25 055 055 055 057 057 055	0	TYPE 915 ACCEPT 1000.A ACCEPT 1000.A ARITE(1UNIT,905) ARITE(1UNIT,905)FSTART,FEND,STEP,NPOINT ARITE(1UNIT,906) ARITE(1UNIT,906)
	INTEGEP OBJA(512),OBJP(512),TRANA(512) INTEGER TRANP(512),CLUTA(512),CLUTP(512),NPOINT,NSANP2 BYTE A		062 062 062 062 062 062		00_277 K#1.WPUINI 1001=K Amelt=k=cot(08.JA(K))*.05 Paaset=f_cot(08.JP(K))*.25 Frec=fstart+float(1001-1)*\$tep
	TYPE 900 Type 912 Accept •. Lunit	000	0337 993 993 993 993	500	JRITE (ILWIT, 907)K, FREQ. AMPLIT, PHASE CONTINUE CONTINUE
	THIS SECTION WILL GENERATE CORRECTED SPHERE DATA FROM Evferimental data	, -	00 50		THIS SECTION WILL STORE THE CORRECTED DATA FOR SYSTEM NEW VDE 202
			869 968 969 969 969 969 969 969 969 969		ACCEPT 1000, A ACCEPT 1000, A ACCEPT 1000, A
	STEP*(FEND-FSTART)/FLOAT(NPOINT) TYFE + PRINT THE SYSTEM REPONSE DATA (Y OR N), '		072 072		ALL ASSIGN(20, -1, 'NEW', 'NC',) Write(20).*)FSTART
	ACCEPT 1000.A 15 (A En 'N') ED TO 100	00	074		WEITE(20, *)FEND WRITE(20, *)NPOINT
	44-116-110/11.912) 14-116-110/11.9055574RT, FEND, STEP, NPOINT		076 077		DO 310 K#1. NPOINT HRITE (20, +)0B.04(K)
	Control = 100000000000000000000000000000000000		078 079 080 31	0	HRITE(20, + JUBUPTK) Continue Call Close(20)
	Protecter T. (M. 1. Profile (F.) = 25 FFE(==551APT+ELOAT(K=1)=\$15 WPTE((UNIT, 907)K, FREG, AMPLIT, PMASE)	000	·	THIS SECTION WILL CORRECT FOR RANGE
	ÇîMTINJE Liditertinjî qis)		1800	8	TVPE 914
	WEITE(10411,905)FSTART, FEND, STEP, NPOINT WEITE(10411,905))	0 82 083		ACCEPT *. IRFLAG CALL RANGE
	С0 ЭС НЕ .НРОЛИТ Амецтеголт(сцита(к))е. 05 FPE0:#FSTAPT+CL0AT(к))е. 05 PMASE=FL0AT(CLUTP(K))е. 25	<u>```</u>	2 2 2 2 2 2 2 2 2 2 2 2 2 2 2 2 2 2 2		TYPE 913.DIST CALL PANCOR TYPE 915
	MRITE(IUNIT, 907)K, FREG. ANDLIT, PHASE Contime Call Shedat Call Sheepit	•	080 080 080 080 080		ACCEPT 1000.A Artestumit.913/D157 Writestumit.913/D157 Writestumit.911/051/55201,5550,001/1

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FORMAT(//// SPRINT SPHERE DATA CORRECTED FOR RANDE (Y OR N). ') DIMENSIONS 002000 (512) (512) 002000 (512) (512) 002000 (512) (512) 002000 (512) (512) 002000 (512) (512) 002000 (512) (512) CLUTA TVPE 4 1 PHASE **IUNIT** NAME FREG MAME WWW MAM FEND 33 **WORDS) HORDS (HORDS**) FLOAT VARIABLES. PSECT +DATA, SIZE + 000102 (WWW **SYS1** ¢ OFFSET OFFSET (1025 OFFSET 002000 (2053. OFFSET 002000 0000010 000046 000044 000000 TYPE 4.4 4 COMMON BLOCK /RANGEL/, SIZE = 000010 (SIZE = 010012 COMMON BLOCK /OBJ // SIZE = 004002 TYPE TYPE TYPE TYPE CORREC 1+2 1+2 ANPLIT R+4 1+2 FSTART R44 IRFLAG 1+2 SHEEP WWW 006000 002000 002000 002000 000000 000000 OFFSET TRANP MAN MAME MAN MAN 08.6 TYPE R.4 4 SECTION SYSTA ; LOCAL AND COMMON ARRAYS ~ ORJ SYSTA SYSTA L(V OR N): SPHDAT (,****** BLOCK /SYSTA 000062 OFFSET 000000 OFFSET 000000 006000 010010 CLOSE 000036 000000 OFFSET 000000 OFFSET MAN END TYPE SUBPOUT INES. TYPE TYPE TYPE TYPE 4 ***** 0001 618 1+2 **8** 8*4 **2***1 1#2 1•2 NPOINT 1+2 1+2 Ī 4 * 2 ¥ 1•2 5 ŝ 916 NÚMHOU ASSION OB. 19 OB. 19 TRANA TRANA CLUTP PANA RANGE LOCAL CLUTA DLUTP MAR 0133 0135 0136 0136 0138 NAME STEP MAME DIST MAR OB. M WWW NAME IQNI)) 3 3 COMPECTED DATA: ') FORMAT(///.'SENTER THE FILE NAME FOR THE COORECTED DATA: ') FORMAT(///.'SDO YOU WANT TO STORE THE CORRECTED DATA (Y OR N) FORMAT(//. THIS SECTION WILL GENERATE CORRECTED SPHERE DATA c0PMaT(////// ***** UNCORRECTED SPHERE DATA *****)
F0PMaT(/15/, ***** UNCORRECTED SPHERE DATA *****)
F0PMaT(/15/, *STATING FPECUENCY 0H2: *.1PG15 7/15%, *FREQUENCY STEP 0H2:
E1POING 715%, *UNPRER OF 57EPS: *.17)
F0PMAT(/17,75(**/)/12%, POINT **.5%, FREQUENCY 7%,
F0PMAT(/10,17), 75(**/)/12%, POINT **.5%, FREQUENCY 7%,
F0PMAT(10,117,25,24), PDISE DEGREES 7%, IPG15 7)
F0PMAT(10,117,25,24), PDISE DEGREES 7%, IPG15 7)
F0PMAT(10,117,25), ***** COPRECTED EXPERIMENTAL SPHERE FGPMATY///, ENTER THE RANGE CALCULATION FLAG/// 1=>DIRECT 1 MEASJREMENT// 2=>FOURIER ANALYSIS ///9INPUT FLAG. ') FGPMAT(///9PRINT SPHERE DATA CORRECTED FOR &YSTEM RESPONSE CIPMAT (/ / / / / / + + + + EXPERIMENTAL SPHERE DATA CORRECTED FORMATI /// 4. CHIER LOOICAL UNIT NUMBER FOR OUTPUT (7=> OPMAT (////, CALCULATED PANOE TO THE TARGET . ', 1PG15. 7. FORMAT(////, **** PROGRAM SPHERE ****') FORMAT(///, *9D0 YOU MANT TO STORE THE RANGE CORRECTED LGTA Y OP N) 2 *) ç FROM // EXPERIMENTAL DATA') FORMAT///.4FRINT THE UNCORRECTED DATA (Y OR N): THIS SECTION WILL STORE THE RANGE CORRECTED DATA FOPMAT(///. "SENTER THE FILE NAME FOR THE RANGE AMEL IT=FLOAT (08.44(K))+. 05 PMASE=FLOAT (08.94(K))+. 25 FFE:0=FIZAPT-FLOAT (1NDI-1)+5TEP WPITE(1UNIT-907)K, FREO, AMPLIT, PMASE THIS SECTION WILL DISPLAY THE DATA CALL ASSIGN(20...-1, 'NEW', 'NC',) JP1TE(20...)FSTART WP1TE(20...)FEND ACCEPT 1000. A WETE(20, +)NP0INT D0 460 Y=1, NP0INT WETE(20, +)08,14(K) WETE(20, +)08,14(K) WOITE([UNIT.905) Do 410 K=1. NPOINT Itid1=K (v***** BUNJA 404 FORMAT STATEMENTS CALL CLOSE (20) (, METERS() TREMINAL) 100 CCATT INME IVPE 800 CONTINUE TALL THE IS TVPE DATA ÷., 475 ŝ 613 515 ŝ 902 606 55 912 ¢9₹ 8⁷² ŝ 906 Į, e So 116 • 901 ē 00 55555 0112 0100 0113 ¢110 0121 0130 5933 1010 5102 6919 01:0 0109 2010 6010 0110 0112 0117 0113 0110 0110 0122 0123 0125 9127 0128 0129 1610 0132 0111 Poro C

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HORDS)

000056 000023 OFFSET TYPE R=4 R.4 OFFSET 004000 000000 0001000 OFFSET OFFSET 010004 TYPE TYPE TYPE 1+2 1+2 1.2 4=2 NSAMP2

FUNCTIONS. STATEMENT AND PROCESSOR-DEFINED FUNCTIONS

WWW

TYPE 44

RANCOR

FORMAT(//1X, **** SUBROUTINE SYSINP ORTAINS THE SYSTEM RESPONSE FORMAT(//1X, **** SUBROUTINE SYSINP ORTAINS THE SYSTEM RESPONSE 1 FILES ****) FORMAT(//**, FUTER THE ANTENNA SYSTEM FUNCTION FILE NAME: ') OFFSET 000032 004000 010004 OFF SET TVPE FUNCTIONS. STATEMENT AND PROCESSOR-DEFINED FUNCTIONS. W TYPE TYPE 1+2 (SORDS) **I**#2 4 • X -- DIMENSIONS CLUTA 512.) (512) 512.) (512) 512.) (512) 512.) (512) 512.) (512) TYPE MANE FEND 18. W COMMON BLOCK /SYSTA /, SIZE = 010012 (2053. WORDS) = 000044 (NATE CALL ASSIGN(12,,-1,'QLD','NC',1) READ(12,-9) FSTART READ(12,-8) FSTAD READ(12,-8) NPOINT DO 100 //EL, NPOINT DO 100 //EL, NPOINT READ(12,-8) TRANA(K) READ(12,-8) TRANA(K) CALL CLOSE(12) TYPE 902 CALL ASSIGN(12,.-1, 'QLD', 'NC',1) READ(12.*) FSTART READ(12.*) FSTART READ(12.*) NPOINT DO 200 (*1, NPOINT ---SIZE---OFFSET 002000 000010 000024 OFFSET 002000 (002000 (002000 (002000 (TYPE 442 TYPE TYPE LOCAL VARIABLES. . PSECT \$DATA. SIZE R#4 I#2 FSTART R+4 . SWDAT **WAR** SECTION OFFSET SYSTA 004000 SYSTA 004000 SYSTA 006000 SYSTA 000000 SYSTA 002000 READ(12. +) CLUTA(K) READ(12. +) CLUTP(K) TRANP NAME MAME 58 CALL CLOSE(12) RETURN 4+2 TYPE LOCAL AND COMMON ARRAYS: ORMAT (A1) CONTINUE CONTINUE CLOSE 000022 000900 010010 OFFSET OFFSET 000000 060000 **NAME** END TYPE TYPE SUBROUTINES, TYPE TYPE 4=2 Ξ 1+2 NSAMP2 1+2 1#2 NPOINT 1+2 14 1 222 901 902 ğ 00000 88 CLUTP ASSION RANA CLUTA TRAND MAME NAME MAME 0071 0072 0073 NAME))) ACCEPT - SEND TYPE - CEND TYPE - CHIEP THE NUMBER OF FREQ POINTS: ' ACCEPT - NEWER THE NUMBER OF SAMPLES AT EACH FREQ. ' ACCEPT - NISAME2 TYPE - SET UP FETLECTING PLANE FOR TRANSFER FUNCTION MEASURMENT' ACCEPT - NISAME2 TYPE - SET UP FOR CLUTTER MEASURMENT' AUSE 'A SET UP FOR CLUTTER MEASURMENT' FAUSE 'S TO P FOR CLUTTER MEASURMENT' FAUSE 'S STORE TO PROCEED + + + + + CALL SUBATICTUTA. CLUTP FETAD. NOINT. NSAMP2) TYPE - 'S STORE DATA ON DISC? (Y OR N)' INTEGER TPANA(512), TRANP(512), CLUTA(512), CLUTP(512), NPOINT COMMON/SYSTA/TRANA, TRANP, CLUTA, CLUTP, FSTART, FEND, NPOINT THIS SUBPOUTINE WILL READ IN THE SYSTEM RESPONSE FILES AND PLACE THEM IN A COMMON BLOCK TO BE PASSED TO OTHER POUTINES PAGE 001 TYPE +, "ENTER THE STARTING FREQ IN GH2: " CCCEPT +, FSTART TYPE +, "ENTER THE ENDING FREQ IN GH2: " VPE +, 'OLD OR NEW DATA (1=)NEW, O=)OLD)' CCEPT +, FLAG GALL ASSIGN(12,.-1,'NEW','NC',1,) Mette(12,e) FSTART Heite(12,e) FSTART Heite(12,e) NPOINT VIEW / VIC / 1) F IT EO (N.) 6010 300 FORTRAN IVVO2 1-1 2001 SJBROUTINE SVS1 C VERSION 1. 0 29-NOV-79 DO 30 K=1, NPOINT MPITE(12, +) TRANA(K) MPITE(12, +) TRANA(K) WPITE(12.*) CLUTA(K) WPITE(12.*) CLUTP(K) CONTINJE ASSIGN(12, -1, (FLAG) 50. 50. 10 WRITE(12, +) FSTART WRITE(12, +) FEND WPITE(12, +) NP0INT DO 40 K=1, NPOINT CALL CLOSE(12) TYPE 902 CALL CLOSE(12) ACCEPT BOOLC 00 TO 300 TYPE 901 17FE 900 **JUNI INUC** 100 3d . BVTE C 8 ŝ Ş 8 00000 C v 6000 22228 0035 5100 00.45 00.45 0.02 1000 9449 5455 5455 0100 1.2 \$1:00 **3**23 62.53 62.50 e con 1000 2012 0023 0.40 0.042 0018 ŝ 2222 1023 9024 n Second ŝ Serve (COU) 2002 5 4000 2 10.00 1710 1.144 ŝ ć `\$

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THIS SECTION WILL STORE THE CORRECTED DATA FOR SYSTEM RESPONSE WRITE(IUNIT,911) MRITE(IUNIT,975) MRITE(IUNIT,905) MRITE(IUNIT,906) D0 410 K=1, NPOINT IF (HEAL W) JU 10 200 WRITE(LWIT, 900) WRITE(LWIT, 905)FSTART, FEND, STEP, NPOINT WRITE(LWIT, 905)FSTART, FEND, STEP, NPOINT UD: 277 K#1, NPOINT INDI#Y DO 210 K=1, NPOINT AMPLIT#FLOAT(OBJAK(X))* (N5 AMPLE#FLOAT(OBJE(X))* (N5 FMEO#FSTART+FLOAT(K-1)*STEP MRITE(JUNIT,907)K,FREQ,AMPLIT,PMASE AMFLIT=FLOAT(OBJA(K))* 05 PHASE=FLOAT(OBJA(K))* 05 FREG=FSTART+FLOAT(IND)-1)*STEP WRITE(INUT,907)K,FREQ.AMPLIT,PHASE HERE LIST THE CORRECTED SPHERE DATA THIS SECTION WILL CORRECT FOR RANDE CALL ASSIGN(20...-1, 'NEW', 'NC',) WRITE(20.+)FSTART WRITE(20.+)FEND INDI=K AMPLIT=FLOAT(OBUA(K))+. 05 PHASE=FLOAT(OBUP(K))+. 25 FREQ=FSTART+FLOAT(INDI-1)+STEP • ACCEPT 1000. A ACCEPT 1000, A IF (A E0 'N') 60 TO 470 WRITE(LUNIT, 913)DIST ACCEPT 1000, A IF (A EQ 'N') 60 TO 280 D0 310 K=1.NP01NT WRITE(20, *)08JA(K) WRITE(20, *)08JP(K) CONTINUE WRITE (20. +) NPOINT TYPE 914 ACCEPT + IRFLAD CALL CLOSE(20) TYPE 913. DIST CALL RANCOR CONTINUE CALL CORDAT CALL RANGE 916 TYPE 915 CONTINUE CONTINUE TYPE SO2 **TYPE 801** TYPE 8 C C C 8.1 310 000 0072 0073 0074 0076 0077 0078 0078 0078))) 0089 0089 0089 0089 0082 6300 0083 0075 3) Э • (THIS PROCPUM WILL TAYE EXPERIMENTAL DATA FOR THE SPHERE BUT THE SYSTEM RESPONSE. IT WILL PRINT BUTH THE COPPECTED AND UNCORRECTED DATA AND FINNALY DISPLAY IT ON THE HIGH RESOLUTION CRT. ALSO IT DESIFED IT WILL GENERATE IF DESIRED ANALVITICAL DATA FOR THE SPHERE AND STORE IT IN A FILE. IT WILL ALSO PENT THIS DATA IF DESIRED TO A FILE AND ALSO DISPLAY IT ON THE HIGH RESOLUTION CRT MONITOR. COMMON/PANGE1/DIST. IRFLAG. IUNIT COMMON/PANGE1/DISJ. OBJP, NSAMP2 COMMON/SYSTA/TRAND, TRAND, CLUTA, CLUTP, FSTART, FEND, NPOINT INTEGER OB.M(512), OB.P(512), TRANA(512) INTEGER TRANP(512), CLUTA(512), CLUTP(512), NPOINT, NSAMP2 BYTE A THIS SECTION WILL GENERATE CORRECTED SPHERE DATA FROM EXPERIMENTAL DATA 1 ŝ PAGE 001 3TEP≞(FEND-FSTART)/FLOAT(NPOINT) TVFE +. 'FRINT THE SVSTEM REPONSE DATA (Y OR ACCEPT 1000.A IF (A E0 'N') G0 T0 220 WPTF6(IUNIT, 904) WPTF6(IUNIT, 905)F5TART, FEND, STEP, NPOINT WEITE FUNIT, SUS FSTART, FEND, STEP, NPOINT DO 20 K=1, NPOINT APPLIT=ELOAT(TRANA(K)) + 05 PHOEEELOAT(TRANA(K)) + 05 PHOEEEELOAT(TRANA(K)) + 25 HPLTETLUNIT, 007)V, FREO, AMPLIT, PHOSE FORTRAM [VV/2] 1-1 FR1 30-NOV-79 00: 25: 58 1001 PR05PAN SPHER3 IF (A EQ (N') 60 TO 100 UPTTE(100117, 417) VEPSION 1. 0 30-NOV-79 CALL SPHDAT CALL SUEEP(10....) UPITE(IUNIT, 205) HPITE (JUNIT, 906) ACCEPT +, IUNIT ACCEPT 1000.A 2373 ŝ 503 912 **LVPE** and CONT INUE 1 , PE TYPE 88 00 ωu 0000 0100 00100 00100 000 ₹úuij **909** 2229 2225 30 5.00**8** 5.005 5.005 Sec 5000 C.C.C. 1155 2100 1900 C . C C

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STOP ****** END OF PROGRAM ******* END

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LOCAL	VARIAB	LES. PSECT	BDATA,	SIZE	= 000074 (30	MORDS)	
NAME	TYPE	OFFSET	NAME	TYPE	OFFSET	NAME	TYPE	OFFSET
٩	3	000032	AMPLIT	R+4	000042	FREQ	R=4	000052
1 0 N1	1+2	000026	¥	1 •2	000040	PHAS	E R+4	000046
STEP	A . A	10000						

			•							
	OFFSET	900,490		067361	004000			0940-0	c10004	
	TYPE		•	ΙΥΡΕ	I+2		TVPF	14.2	5+4	
WORDS)	NAME	IUNIT	HORDS)	NAME	NSAMP2	WORDS)	NAME	CLUTA	FEND	
¥	OFFSET	000004	(1025.	OFFSET	002000	(2053.	OFFSET	00200	000010	
000010	TYPE	i 1+2	004002	TYPE	I+2	010012	TVPE,	1+2	R*4	
= 321S	NAME	IRFLAG	\$12E =	NAME	OBJP	\$12E •	NAME	TRANP	FSTARI	
~			~			:				
/RANGE1	OFFSET	000000	/GBJ	OFFSET	000000	/SYSTA	OFFSET	000000	000000	010010
BLOCK	TYPE	R + 4	BLOCK	TYPE	1+2	BLOCK	TYPE	1+2	I+2	1+2
COMMON	NAME	DIST	COMMON	NAME	OBJA	COMMON	NAME	TRANA	CLUTP	NPOINT

LOCAL AND COMMON ARRAYS:

SECTION DEFSETSIZE DIMENSIONS	SYSTA 004000 002000 (512) (512)	SYSTA 006000 002000 (512) (512)	nr.1 000000 002000 (512) (512)	08.1 002000 002000 (512.) (512)	SYSTA 000000 002000 (512) (512)	EVETA 002000 002000 (512) (512)
SPOTTON OFFSET	SYSTA 004000	SYSTA 006000	OB.1 000000	002000 UBJ	SYSTA 000000	CVETA 002000
TVPF	1#2	(+) (+)	1+2		•1	
NAME		C UTP			TRANA	

SUBPOUTINES. FUNCTIONS, STATEMENT AND PROCESSOR-DEFINED FUNCTIONS

TYPE Re4

NAME RANCOR

NAME FLOAT

R. 4

CORDAT SWEEP

TYPE R#4

ASSIGN

) **)**

NAME

SVS2

NAME CLOSE SPHDAT

4*2

TYPE Re4 Re4

TYPE

NAME

TYPE R#4 R#4

FORMAT(////PRINT SPHERE DATA CORRECTED FOR RANGE (Y OR N)) " THIS SECTION WILL GENERATE CORPECTED SPHERE DATA FORMATIVIN CENTER THE RANGE CALCULATION FLAGYNY 1=>DIRECT 1 MERCINEMENT// 2=>FOURIEP ANALYSIS ///*INPUT FLAG) FORMATI///*9PKINT SPHERE DATA CORRECTED FOR SYSTEM RESPONSE FORMAT (//////, ***** EXPERIMENTAL SPHERE DATA CORRECTED ENEMAT(////\$', 'ENTER LOGICAL UNIT NUMBER FOR OUTPUT (7=> FREND TILLIN, CALCULATED RANGE TO THE TARGET . . 19615. 7. FOWMAT(////,**** PROGRAM SPHERE ****) FORMAT(///.*5D0 YOU MANT TO STORE THE RANGE CORRECTED DATA (Y OP N) 2 *) FORMAT(//1Y.75////2X.POINT *.5X.FREQUENCY.7X. AMPLITUDE DB ~.5X.PHASE DEGREES/1X.75(*/)/) FORMAT(101.17.5X.PD015 7.3X.1P015 7.3X.1P015 7) FORMAT(////// ***** CORRECTED EXPERIMENTAL SPHERE FüRMAT(/////// ********* SYSTEM CLUTTER ********* THIS SECTION WILL STORE THE RANGE CORRECTED DATA FORMAT (/ / / SENTER THE FILE NAME FOR THE RANGE WPITE(IUNIT, 907)K, FREQ, AMPLIT, PHASE Continue THIS SECTION WILL DISPLAY THE DATA CALL ASSIGN(20., -1. "NEW", "NC",) WPITE 20. +) FSTAPT (A E0 'N') GO TO 500 elte(20. +)FEND 91te(20. +)NP0INT 9.460 K+1. NP0INT WPITE(20. *)08.JA(K) WPITE(20. *)08.JP(K) FORMAT STATEMENTS (CARARA BONDA AUG TYPE 901 ACCEPT 1000.A CL0SE(20) ()日本本本・シュリコン ç METERS ') FORMAT(A1) (TECM [NGL) 106 ŝ **JUNI 1MOD** IN AD IN TVPE **P**LL CALL ų, ş 410 170 0001 3 ŝ ŝ 802 800 110 503 8 513 513 610 E lo 5 908 600 608 012 410 \$10 0000 5000 \$11¢ 5003 5010 000 ŶŬĨŬ 2010 0112 6113 6113 9116 0113 0216 0121 0122 6210 2000 60.00 Ş 1010 0102 6010 0010 0111 2110 0124 0124 0110 0110 0125 8000 (

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MAME 3 > (SUBTRACT ANTENNA CLUTTER FROM OBJECT DATA. THIS CLUTTER DATA WIST NOT ... BE CORRECTED FOR THE SYSTEM TRANSFER FUCTION ****************** ********************************* PPPPPPPPP REPLACE DATA INTO INTEGER ARRAY **************** THE MEXT STEP IS TO DIVIDE BY THE TRANSFER CHARACTERISTIC OF THE SYSTEM IN ARRAYS TRANA(AMPLITUDE) AND TRANP(PHASE). CONTON/OBJ/OBJA, OBJP, NSAMP2 CONTON/SYSTA/TRAND, TRANP, CLUTA, CLUTP, FSTART, FEND, NPOINT OFFSET 000022 000012 000000 THIS SUBROUTINE WILL TAKE THE DATA THAT HAS TAKEN FROM SPHDAT AND CORRECT IT WITH THE DATA FROM SUBROUTINE SYSINP. WHICH IS THE SYSTEM RESPONSE. [NTEGER 08/4(512),08/9(512),TRANA(512),TRANP(512) [NTEGER CLUTA(512),CLUTP(512),NPOINT,NSAMP2 TYPE WORDS) 1+2 4 4 4 . K PAGE 001 WWW Hora 27 ï LOCAL VARIABLES. PSECT SDATA, SIZE = 000066 (FRI 30-NDV-79 00:21:16 OFFSET 910000 \$E0000 000024 ADGPES=/DBAHP-TDBAHP PHESs=(JDB4H-TDBAHP PHESs=(JD1)-PRES=PRES-2, #PI IF(PRES_LT, -PI)PPES=PRES-2, #PI IF(PRES_LT, -PI)PPES=PRES-2, #PI TDEAPP=FL0AT(TRANA(K))+ 05 TPPH=FL0AT(TRANP(K))+ 05 00BAPP=FL0AT(08.JA(K))+ 05 00JPH=FL0AT(08.JA(K))+ 05 00JPH=FL0AT(08.JP(K))+ 25+DEG 06.JA(K)=IFIX((20 +ADBRES) 08.JP(K)=IFIX((PRES/DE0)+4.) TYPE R.4 ODBAMP Re4 1 VERSION 1. 0 30-NOV-79 **FDBAMP** FOPTPAN IW02 1-1 FK1 *** 100 K=1. NP01NT MAN ŝ PI=3 1415926 DEG=P1/180 CONTINUE RETURN END 0000040 000020 OFFSET 000044 ĥ 8 TYPE 4.4 ADPRES Re4 1 8 uυ 000 0 **0 c** 000000000 Hango 8118 8818 0 MAME PRES 8002 8000 2008 2008 2008 111 2883 2883 8 Ċ) C

FORTRAN IVVO2. 1-1 FRI 30-NOV-79 10: 10: 36 PAGE 001 8UBROUTINE SWDAT (00JA, 00JP, FSTART, FEND, NPOINT, NSAMP2)

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THIS SUBROUTINE WILL OBTAIN DATA FOR AN OBJECT OVER THE FREG RANGE SET WITH THE NUMBER OF POINTS SPECIFIED.

INTEGER OBJA(\$12), OBJP(\$12), NPOINT, NSAMP2 INTEGER TRANA(\$12), TRANP(\$12), CLUTA(\$12), CLUTP(\$12) STEP=(FEND-FSTART)/FLOAT(NPOINT) DO 100 K=1, NPOINT FREQ=FSTART+FLQAT(K-1)+STEP CALL SWEEP(1, FREQ, 1, 4) CALL PHAMP2(1A, IP, NSAMP2) CALL SWEEP(1. FSTART, 1.4) CALL SWEEP(1, FSTART, 1.4) 08JA(K)=1A-2048 08JP(K)=1P-2048 CONTINUE RETURN END 8 000000000 0003 0005 0007 00010 00110 00111 00113 00113 00113 00113 00113 0004

LOCAL VARIABLES, PSECT SDATA, SIZE = \$10040 (2064, WORDS)

010024 OFFSET FSTART R+4 @ 000004 . 010020 TYPE I#2 R=4 MAME STEP ¥ OFFSET 010026 NSAMP2 1+2 @ 000012 010034 TYPE I#2 4.4 MAME FREQ e. LOCAL AND COMMON ARRAYS: TYPE OFFSET R+4 @ 000006 010032 NPOINT 1+2 @ 000010 1#2 MAME FEND Ā

512) (512) 512) (512) 512) (512) 512) (512) 512) (512) 512) (512) 512) (512) 002000 (002000 (002000 (002000 002000 002014 006014 000000 000002 •DATA •DATA •DATA •DATA •DATA •DATA 22222 * * * * * * CLUTA OBUA POR TRANA

DIMENSIONS

----SI 2E

SECTION OFFSET 004014

TYPE

1+2

002000

FUNCTIONS. STATEMENT AND PROCESSOR-DEFINED FUNCTIONS. SUBROUT INES,

TYPE MARE TYPE NAME TYPE NAME NAME TYPE TYPE MAN

8 *** *** SHEEP PHAMP2 Re4 4 FLOAT

IV VO2. 1-1 FR1 24-AUG-79 00: 18: 52 FAGE 00: SUBROUTINE SWEEP(MODE, FREQ, LOAND, IEEEND) MODE 13 THE SWEEPER MODE 1-9, MODE 1 15 DESCRETE MODE FOR SETT THE CH FREQUENCY OTHER MODES ARE OUTLINED IN THE SWEEFER IEEE MANUAL. ILE IBAND IS EQUAL TO 4 THEN TRAUSH IF IBAND-4 THEN FULL SWEEP 2-18 OH2 OTHERWISE THE WILL COMPUTE THE NECESSARY BAND. FAGE 001 ITIME DELAY TO CHAGE MD 15.3 XS THESE CONSTANTS WERE DERIVED USING PROGRAM CALAR BAS CN AUG-24-1979 6. 05)) IBANDI#1 12. 4)) IBANDI#2 18. 0)) IBANDI#3 IF MODE=10 THEN THE PROGRAM WILL RESET THE INTERFACE INITIALIZE IBAND2 AND MODEL TO ZERO IN THE MAIN LINE BYTE CHD(7), MD(3), BND(3) DATA CHD(7), TIS, MD(3), TIS, BND(3), "15/ DATA CHD(1), YV, C, CHD(6), TIS, BND(1), YY, **BND(1), Y3**, / IF (MDE_10) 3, 2, 3 FRI 24-AUG-79 00:18:52 IF ((FREG) GE, Z. 00), AND (FREG, LT. IF ((FREG) GE, 6, 05), AND (FREG, LT. IF ((FREG) GE, 6, 05), AND (FREG, LT. IF ((BANDZ-IBANDI) 36, 38, 36 IF (BANDZ-IBANDI) 36, 38, 36 FICODE(1, 1000, BND(2)) IBANDI CALL IPSEND(BND, IEEENO) CALL IFIFC 60 T0 2000 1F (MODE-MODE1) 5, 10, 5 ENCODE(1, 1000, MD(2)) MODE (40. 50. 60. 70), IBAND IBSEND(MD, , IEEENO) ITER=1, 500 . (IEAND-4) 25, 20, 25 60 T0 80 B0=12 0018 B1=5 98325E-04 B2=9 44848E-11 38009E-04 B1=4. 18734E-04 B2=1. 27261E-10 81=1. 59929E-03 82=6. 01267E-11 V02. 1-1 BAND2=I BANDI GO TO 80 BO=5. 99786 95966 40DE1=M0DE CALL IBIFC **JEREN** 00022 I BAND: PROCRAM DONT INCO EAND œ î 010% B2=2 39 F 29 F 010 Ŀ 8 2 FORTRAN I 0001 FORTRAN IV 89 8 \$ œ٩ 2 ۶ 3 88 2 ო 000 0044 0043 0044 00043 ્ર C 100 304 0 9) OFFSET 000022 000016 000012 TVPE FUNCTIONS. STATEMENT AND PROCESSOR-DEFINED FUNCTIONS. V VO2 1-1 FRI 24-AUG-79 00: 19: 19 Subpointine Mandriane, Iphase, NUM) This foutthe Mill tree NUM Reddings From A/D CHWNELS AND 1 Osphase 1-DAMPLITUDE **MAR** TVPE 4*A 1#2 20. WORDS) **1**#2 IAMPI TVPE NAME PHAS FORTRAN IV STORAGE MAP FOR PROGRAM UNIT PHAMP2 Local variables, psect sdata, size = 000050 (20 **MAR** TYPE OFFSET 1+2 8 00000 000024 I+2 € 000004 I AMP = I F I X (AMP / FL OAT (NUM)) I PHASE = I F I X (PHAS / FL OAT (NUM)) RE TLAN END TVPE 1#2 IPHASI 1+2 МР=4МР+FLΩAT(IAMP1) №65=РН45+FL0AT(IPHAS1) № 20. J=1,20 NAME IFIX MASI=1ADINP(0.0.) 10 F=1, NUM PP1=1ADINP(0, 1,) IAHP MAME Ę 1. 2 TYPE DO 5 J=1, 200 CONTINUE IADINP CONTINUE BUNI LNOU **ÚFFSET** 900000 00000 000020 W 0=501 Ú≡ di 2 • ¢ TYPE SUBROUT INES. SORTHAN 1 TYPE 8 **4** 8 e 4 PHASE 1+2 1+2 22 ųυ FLOAT 0000 0000 0000 1911 8883 8883 We we 888 888 \$913 8013 AME 0003 1160 9000 ţ C י: הי C C

OFFSET 000032 000030 000023 000046 TYPE 482 FUNCTIONS, STATEMENT AND PROCESSOR-DEFINED FUNCTIONS. IDELAY OF 5. 32 HS TO STEP |V1=|FTX((-B1+50RT(B1++2-4+B2+(B0-FREQ)))/(2+B2)) EKCODE(4,1001,CHD(2)) |V1 CALL IBSEND(CHD,,IEEEND) D0 40 |TER=1,200 !!DELAY OF 5.32 MS TO 5 MAN TYPE SORT W IV STORAGE MAP FOR PROGRAM UNIT SWEEP VAPIABLES: PSECT \$DATA: SIZE = 000070 (28. WORDS) 2.) (3) 4.) (7) 2.) (3) 2.) (3) R#4 I#2 IBANDI 1+2 MODE1 1+2 TYPE I#2 ITER W ß **WWW** IFIX TYPE OFFSET IEAND 1+2 & 00004 000042 IEEEN0 1+2 @ 000006 I+2 € 000000 SECTION OFFSET -----SIZ 8DATA 000022 000003 (9DATA 000010 000007 (9DATA 000017 000007 (TYPE IBSEND I+2 R+4 W N **Man** MODE 10 NAME TYPE I=2 LOCAL AND COMMON APRAYS: FORMAT(1)) FORMAT(14) CONTINUE IBREN OPESET Re4 000036 R+4 @ 000002 000034 000052 RETURN TYPE TYPE SUBROUT INES, 8<u>8</u>88 [BAVD2 1+2 TYPE 1•2 1+2 53 Ξ u 8 FORTRAL **FRE**Q **IBIFC** MAM MAR ¥2000 121 2 Ç

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STOP FREQ. OHZ. ", F2. " SATA FOINTS", N 230 GRAPH (* LINES * N.F. () A(),,,1) 240 LAREL (* UNDERLINE * * AMPLITUDE IN DB*,,1) 250 GRAPH (* LINES * N.F.F.(). B(),,,2] 250 LABEL (* UNDERLINE * * PHASE IN DEGREES *, 2) 270 LINPUT 78 280 DISFLAY + CLEAR 290 DISFLAY + CLEAR 300 DD 70 30 310 END 20 UN UN HICOU, B(200), F(300), F(300) 30 FRINT (NFUT DATE FILE NAME', 50 FRINT (DATE ', DATE 50 OPEN AS FOR INDUT AS FILE #1 100 INDUT #1, FREQ. OHZ. ', FLI' 110 FRINT 'START FREQ. OHZ. ', FLI' 120 FOR I = 0 TO N-1 130 DE (FZ-FI)/N-1 130 DE (FZ-F 10 DIM 4(500), B(500), F(500)

> , <mark>,</mark> 3

APPENDIX II

こうざい いまんせい FORMAT(/// FIELD IN THETA POLARIZATION DIRECTION'///// FORMAT(/// FIELD IN PSI POLARIZATION DIRECTION'///// FORMAT(//75(**)/// POINT #*// FREQUENCY // K#A APP DB // POINT #*// FREQUENCY // K#A APP DB // POINT #*// PHASE ///75(**/// 000000 OFFSET 000000 00)044 01:042 OFFSET OFF SET 011010 9010104 090010 010045 01-014 010100 ex charter to the second secon 7/* FSTART#*, F15. 7/ FEND#*, F15. 7/* ENC TYPE TYPE ž 4 • X TYPE (SORON) 4 4 = 2 8.0 R=4 1#2 4.4 4.2 THETA **NAM** MAR FREQ 154 VARIABLES, PSECT SDATA, SIZE = 010134 (2094. **WAR** (SONOM ğ \$C1 ÿ XdO KAS REA Var WORDS) (ANS. EQ. /N') GO TO 1000 DE 4, 'ENTER THE NAME FOR DATA FILE:' CALL ASSIGN(20.,-1.'NEW','NC'.1.) WRITE (20.+NUMPTS WRITE (20.+)FSTART WRITE (20.+)FEND WRITE (20.+)FEND WRITE (20.+)FHTA WRITE (20.+)PSI WRITE (20.+)FPOL *++++ END OF PROGRAM ++++* C 10. 010000 \$12E = 000050 (20. 000010 000040 OFFSET 010054 OFFSET RAD1US=', F15. 7/' 010010 010020 010050 OFFSET 010074 010000 010036 = 000024 START K#A=', F15. 7/' TYPE TYPE 1VPE 8. U 0 1 1 1 1#2 ****** 4.8 4.2 1 ŝ WRITE (20, #)MAG(N) WRITE (20, #)ANGLE(N) SLOWN STEPK Wan FPSI MAN MAN FEND /, SIZE DO 200 1=1, NUMPTS **803** Å ₹ ACCEPT 700, ANS LOCAL AND COMMON ARRAYS: : FORMAT (/// FORMAT (A1) 000000 OFFSET 000030 CONTINUE 000000 010072 TYPE OFFSET 010052 OFFSET 010026 010024 010004 010032 010064 COMMON BLOCK /SCT STOP END TYPE COMMON BLOCK / Ī TYPE 6 0 TYPE FTHETA C+8 0 * 0 1+2 1+2 4.84 2000 *** 8 Ē R.4 83 8 810 801 FSTART ¥₹ No TINUI STEPF LOC PL NAME MAN 0078 0079 0080 0082 <u> 502</u> 0072 0074 0075 0075 0076 0081 õ 0069 Ŷ 0071 0900)) 9) SI=PSI+3 1415926/180.0 VPE +. 1N MHICH POLARIZATION CALCULATE FIELD (O=>THETA, 1=>FSI). (CEPT +. IPOL URITE (IUNIT, 800)A. FSTART, KAS, FEND, KAE, NUMPTS, THETA+(60/3. 1415926 THETA=THETA=1 1415926/190. O 1776 - 1. CENTEP POLARIZATION SCATTERING ANGLE PS1 (DECREES). 2006 PT -, PSI COMMON /SCT/FTHETA, FPS1, PS1 COMMON 501, 502, SC1, SC2, KA, THETA COMMLEX 501, 502, SC1, SC2, FTHETA, FP31, F FE44-4 YA, MAG(5121, ANGLE(512), KAS, KAE, PS1, THETA PYTE ANS DATA C/2, 997925E+10/ 'ENTER SCATTERING ANOLE THETA (DEGREES); ' PACE 001 CEPT+, NUMPTS PPE +, CENTER LOOICAL UNIT NUMBER FOR OUTPUT' CEPT +, IUNIT HRITE (TUNIT, +)N. FREQ, KA. 10+AL0010(AMP), AND FREAEFREA-STEPF *A=ra+StEPr YPE +, 'STORE DATA FOR DISPLAY (Y OR N) :' IYPE =, 'ENTER SPHERE RADIUS (A) (CM.): " CENTER STARTING FREG (GH2): " MG=ATAN2(CPX, PEA)+(190/3, 1415926) TYPE +. 'ENTER ENDING FREQ (CH2): ' ACCEPT +. FEND IW02 1-1 SUN 21-0CT-79 02: 12: 32 PROGRAM BISCAT . 'ENTER NUMBER OF POINTS' PS1=190/3 1415926 IF (1P0L E0 0) WRITE(IUNET, 910) IF (1P0L E0 1) WRITE(IUNET, 811) (AMP ED 0 0) AMP=1. 0E-35 1 1415026+(1 0E+9/C) ART-WAV+A TEPF=/FEND-FSTART)/NUMPTS (IPOL EO O) F=FTHETA (IPOL En 1) F=FPSI STEPY=(YAE-YAS)/WMP1S IF (100L E0 0) HAL IF (100L E0 1) HAL HPITE(1UNIT, 801) CEPT +, THETA CEPT«. FSTART A-UNHUNATER LE (N) = ANG PEn=FSTART PY=AIMAG'F dWU= (N) SOL NCEPT .. A CALL PSCAT A=FEAL (F CUTINUE 540=100 1991100 L. 14 L 8 FORTRAN 1 (-001 Q υU $\circ \circ$ 828 કુટ્ટુટ્ટે 53666 83666 83666 26.92 0010 -004040888888 62.50 503 342 000 200 8000 5555 8888 8888 9044 Š ¢ C.

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<u>\$02</u> ₹ TYPE TYPE OFFSET 000030 000000 **MAR** END COMMON BLOCK / 0 • TYPE 4=2 0 0 0 000000 υu NAME MAR 201 23 g ?)) Э C THE SCATERED FIELD AMPLITUDE IN THE THETA AND PSI POLARIZATIONS THIS SUBPOUTINE WILL CALL SUBROUTINES BSCATI-BSCAT4 FOR THE CALCULATION OF THE BISTATIC SCATTERING OF A PERFECTUL CONDUCTING SHERE IT WILL RETURN THE SCATTERING COFFICIENT F AS A FUNCTION OF SCATERING ANOLE THETA AUD OFFSET 000020 OFFSET **TYPE** (SOROH) PAGE 001 MAN W Ś /. SIZE = 000050 (20. WORDS) 5 COMPLEX 501. 502. 9C1. 9C2. FTHETA. FP51 COMPLEX 81. 82. 83. 84 PEAL+4 FA. THETA. P51 VAPIABLES. PSECT STATA, SIZE = 000040 (FORTRAN IVVO2 1-1 SUN 21-OCT-79 01: 15: 59 0001 SUBROUTINE BSCAT COMMON/SCT/FTHETA, FPS1, PS1 COMMON S01, S02, SC1, SC2, KA, THETA OFFSET 010000 TYPE OFFSET 0 10 1000 F (ra. 07. (20. 0)) 00 70 300 ALL BSCAT3 F (14 67 (1. 0)) 00 TO 200 AL BSCAT2 heta=cos(PS1)+SC1 S1=-S1N(PS1)+SC2 F (KA 0T (4)) 00 T0 100 ALL BSCAT1 THETA=COS(PS1)+SC1 PS1=-SIN(PS1)+SC2 TYPE 8 FTHETA=COS(PSI)+SC1 FPSI=-SIN(PSI)+SC2 PETURN END APE FTHETA AND FPSI 82=502+552 FTHETA=505(PSI)+81 FP51=-51N(PSI)+82 50 T0_1000 MAN W 2 ALL BSCAT4 79 1000 NOLE PSI -201+50 **BUNITNO** KA>20. 0 TYPE OFFSET 000000 TYPE OFFSET 000000 / JOOTB NGALLOS L 8 • • 1000 6.0 8 3 õ 000 COCAL 88 Man ¥ 0003 0003 200 555 Š 5 * (c c O C C

PAGE 001 FORTRAN IVVO2: 1-1 SUN 21-OCT-79 01: 15: 39 0001 SUBROUTINE BSCAT1

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1117

THIS SUBROUTINE CALCULATES SI AND S2 FOR A PERFECTLY CONDUCTING SPHERE IN THE RANGE KA C. 4 AS A FUNCTION OF THETA THE BISTATIC SCATTERING ANOLE

COMMON SO1, SO2, SC1, SC2, KA, THETA COMPLEX SO1, SO2, SC1, SC2 REAL +4 KA, THETA

SC1=KA++3+(.5+COS(THETA)) SC2=KA++3+(.5+COS(THETA)+1.) RETURN

/. SIZE = 000050 (20. MORDS)

OFF SET 000000 00:044 TYPE 800 THETA Re4 MAME SCI TYPE OFFSET 000000 000010 R#4 8 • • MAME

TVFE SUBROUTINES, FUNCTIONS, STATEMENT AND PROCESSOR-DEFINED FUNCTIONS. MAN TYPE MAN TYPE MAN

CT-79 38: 24: 50 PAGE 001 FURTRAN 1VV02. 1-	CALCULATE THE SCATTERING CONSTANTS NUCCTING SPHERE IN THE RANGE UNCTION OF BISTATIC ANGLE THETA ECTION HANDBOOK GEORGE T RUCK ED. 0002 COMM	0004 CONT CLI SC2, KA, THETA SC1. SC2 9005 CONT 9005 CONT 9009 REAL 0009 CONT	0010 X=K 0011 X3=X 0012 X13=0012 X13=00012 X13=000000000000000000000000000000000000	5.0)+COS(THETA)+(1.0/12.0)+COS(2+THETA))+KA+*2 0013 013 XH2; (1.1807.0/70.0)-(2531.0/105.0)+COS(THETA)+ 0014 COS 3(2+THETA)+(1.0/2.0)+COS(3+THETA))+KA++4	/6_0)=(4=COS(THETA)=1)+((1_0/5_0)=(1+2=COS(THETA)) 0017 COS 0018 XT=(0019 XT=(0019 XT=)	-T2+T3) 0020 P1= 0021 PT= 0022 PT= 0022 PT=	0023 PI=- 0024 TI=-	a)+1 0055 7 52= 0025 7 52= 0055 7 52= 0055 7 52= 0056 72= 0026 72=	(-(1343.0/105.0)+(3769.0/280.0)+CDS(THETA)+ 0027 73= 0027 13= 0028.0/105.0)+(1.0/8.0)+CDS(14)+(1.0/8.0)+(1.0/8.0)+CDS(14)+KAM+A) + 0028 501+	/6. 0)+(4-COS(THETA))+. 2+(COS(THETA)+2)+(2)+(2)+(2)+(2)+(2)+(2)+(2)+(2)+(2)	0031 C T5= 0031 C 00	0032 C 502	SC21) C	υυ 8033 1 1	TA. SIZE = 000110 (36. WORDS) 0034 P1= E TYPE OFFSET NAME TYPE OFFSET 1 0036 B2=	Re4 000034 C3 Re4 000040 C CI=	I Re4 000024 SCIR Re4 001020
I-1 PACE 001 PROUTINE BSCAT3	IS SUBROUTINE FINDS THE SCATTERING COEFFICIENTS SCI-SCO- INO A FOLYNOMIAL APPROXIMATION IN THE RESOMANCE REGION 3 A PERFECTLY CONDUCTING SPHERE. RANGE 1X(AC20. K=113\)TE 7 ADDUS OF SPHERE THETA=BISTATIC SCATTERING ANGLE THON SOI, SO2. SCI, SC2, KA, THETA	HELEX T1, T2, T3, T4, T5, T6, B1, B2, B3, CON HELEX CON1, CON2, CON3, CON4, C1, C2, C3, C4, C5, C6, C7, C9 HELEX C11, C12, C13, C14, C15, XT, XTH HELEX CON5, CON5, CON5, CON6, CON9 HELEX CON5, CON6, CON7, CON6, CON9 HELEX C0, 1, C0,	KA =X**(1_0/3_0) =X*0/13_0	23=X++ (-2.0/3.0) 512=COS(THETA/2.0) 112=SIN(THETA/2.0)	51=COS(THETA) =CON=THETA=X H=-XT	=3.1415926 . =(P1+THETA) M=(P1-THETA)	=-2*KA+COST2 =(KA/2)+CEXP(CON+P1)	=-(1./(2*KA+(COST2*=3))) =1.0+CON+P2	=(7, 0/(4*(KA++2)))+((SINT2++2)/(COST2++6)). 1=11+(T2-T3)	=-CoST/(2*KA+(CoST2 *+3)) =(1.0+CoN*P4)	i=(1. 0/(4+KA++2))+((4. 0+COST)+(SINT2++2)/ 0ST2++6))	2=11+(T4+T5)	********* CALCULATE CREEPING WAVE TERMS *********	Ā	=P]=(X+(1,0/12,0)) =CEXr(CON+P1) =-CON+X3+SQRT(X/(2+P]+SINT))	=82=81	

14=P2+(C4-CON+C6)

P2=, 173966/(X*SINT) T5=P2+(C9-CON+C11)

CON9=(. 848747, 1. 470073) C15=1. 614208+XM23+CON9 T6=C15+((C3+C13)-(C5+C14))

U

0077

0076

800

C

3

(

6208 0080 0081

SC2=C1+(T4+T5+T6)

0083

o 00

0082

RETURN END

LOCAL VARIABLES. PSECT 6DATA, SIZE = 000730 (236. WORDS)

00100 001110 003270 001320 012100 £9£000 001200 001130 003210 111000 000000 001100 001424 001100 001020 OFFICET 010100 TYPE 0 0 8 • • 800 800 8 2 8 R 8 COST2 SINT2 CONS SNOC CONB H N Ê C12 (SQNOM C15 ő 68 2 ഇ 2 5 /, SIZE = 000050 (20. OFFSET 000260 000310 002000 000230 000460 000474 000450 010000 000000 000434 000340 000070 000120 000150 000370 000454 TYPE 8 1 1 8 0 1 1 1 4.8 8 8.0 8 ** 0.0 8.0 8 000 *** 4 • X 4 *** 8 XM23 MAME SINT Son CON7 COST 4NOC XTM 110 614 p 8 2 8 Z 80 000030 000000 000250 000420 000110 00410 000170 000470 000000 000330 OFFSET 090000 000140 096000 091000 000220 000504 COMMON BLOCK / TYPE 9 10 e t 8.0 0 * 0 0 1 1 0 *** 8 2 0 8 0 * 0 0.00 8 (0) * () MAME CONG CONS CONS ž 63 HE a 8 ដ ខ ĉ 2 4 ž)

OFFSET TYPE Wen С 2 Э

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SUBROUTINES, FUNCTIONS, STATEMENT AND PROCESSOR-DEFINED FUNCTIONS

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TVPE

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TYPE OFFSET 000010 000040

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THE RELEASED AND THE STATE STATE • PHON LOUIS CONT

******* CALCULATE SECOND CREEPING WAVE TERM ********* PI=(8750/(X+TAN(THETA))) C3=(1 0-C0N+PI) CON2=(700283.-. 404308) CON3=(141774..081853) B2=PTH=(XX3)=CON2 B3=ETH=(XM3)=CON3 C4=CEXP(XTH=B2+B3) CON4=(1 032306,1,788036) CON5=(2 232697,-1,289048) CON6=(,141492, 081684) r2=C7+((C3+C9)-(C5+C11)) CON1=(933508, 1. 616911) C2=2 715175+(XH23+CON1) CON7=(1 607133.-..927879) CON5=(00415. 057397) 71=C2+((C3+C4)-(C5+C6)) C7=1 391727+(XM23)+COM C12+ 201776/(X+SINT) T3=C12+(C13-CON+C14) B2=PTM+X3+C0N7 B3=PTM+XH3+C0N9 C13=CEXP(XTM-B2-B3) C2=CON+(1 0+CON+b1) (INIS+X)/16666 =24 B2=PTM=X3=CON5 B3=PTM=XM3=CON5 C9=CEXP(XTM=B2-B3) 83-PT+XM3+CON5 C11+CEXP(XT-82-83) B3=PT+XM3+CON9 C14=CEXP(XT-B2-B3) B2=PT+X3+C0N2 B3=PT+XM3+C0N3 C6=CEXP(XT-B2+B3) SCI=CI+(TI+T2+T3) 2=PT+X3+CON5 32=PT+13+CON7 10--10 υ c υ 9900 0.42 0043 9047 ŝ 1.800 88 88 1100 9990 2049 0.00 รีะ ŝ ŝ 328 22 8079 8073 8073 0073 ŝ 0052 ŝŝ 8 9074 527 5000 500 0072 Š (C C

AFTER NUMBER OF POINTS' AFTER NUMBER OF POINTS' ACCEPT +, ENTER LOOICAL UNIT NUMBER FOR OUTPUT' ACCEPT +, ENTER LOOICAL UNIT NUMBER FOR OUTPUT' ACCEPT +, IUNIT MRITE (IUNIT, 801) MRITE (IUNIT, 801) PACE 001 WRITE (IUNIT, +)N, FREQ, KA, 10+ALOO10(AMP), AND FREG#FRED+STEPF KANKA+STEPK EE *.'STORE DATA FOR DISPLAY (Y OR N) :' EET 700, ANS (ANS EQ 'N') GO TO 1000 FE *.'ENTER THE NAME FOR DATA FILE:' +, 'ENTER SPHERE RADIUS (A) (CH.): COMPLEX F REAL + KA, MAG(512), ANGLE(512), KAS, KAE REAL + 4 KA, MAG(512), ANGLE(512), KAS, KAE DATA C/2, 997925E+10/ VPE +, 'ENTER STARTING FREQ (GHZ):' PY=A1MAG(F) NG=ATAN2(CPX,REA)⇒(180/3. 1415926) IAG(N)=AMP TYPE +, Call Assion(20,,--1, 'NEW', 'NC',1,') URITE (20,+)NUMTS WRITE (20,+)FSTART WRITE (20,+)FEND TYFE +, 'ENTER ENDING FREG (GHZ):' Accept +, Fend COPTRAN 1W02 1-1 THU 18-OCT-79 00:04:05 0001 PROGRAM SPSCAT EPF= (FEND-FSTART) /NUMPTS AV=2+3 1415026+(1. 0E+9/C) TEPK= (KAE-KAS) /NUMPTS RE@#FSTART N#I WRITE (20, #)MAQ(N) WRITE (20, #)ANOLE(N) 00 100 I=1, NUMPTS 00 200 I=1, NUMPTS AS=FSTART+WAV+A NGLE (N) = ANO ALL SCAT MP=CABS(F) COMMON KA.F EA=REAL (F γ₽<u>Ε</u> ... ΥPE 8 υυ υu 82.00 0052 0052 0053 0036 0040 800 0025 0025 0025 0025 0043 8 0000 0032 0033 0033 2041 440 0028 029 0034 ŝ **6**33 ŝ) Э THIS SUBPRICAME WILL CALCULATE BISTATIC SCATTERIND Coefficients for the Perfectly conducting Sphere. In the Pange Va>20 OFFSET OFFSET 0000000 000044 JANT JANN SUPPOUTINES. FUNCTIONS. STATEMENT AND PROCESSOR-DEFINED FUNCTIONS. TYPE TYPE 1 (SORDS) 8 THETA ROG PAGE 001 TYPE į WWW 0 MAN **3**C1 /. SIZE = 000050 (20. WORDS) TVPE NAME LOCAL VARIABLES. PSECT 8DATA, SIZE = 000024 (SORT Елетрам IVVA2 1-1 S.N 21-0СТ-79 38:26:12 0001 SUBROUTINE BSCAT4 COMMON SOI, SO2, SCI, SC2, KA, THETA COMPLEX SOI, SO2, SC1, SC2, C1 REAL+4 KA, THETA TYPE OFFSET 010000 TYPE OFFSET 000010 000040 1 1 **1** MAN P1=2+KA+COS(THETA/2) C1=CMPLX(0, 0, -P1) SC1=- 5+KA+CEXP(C1) **5**1N 8 MAN W **2**03 2 Ä M . SC2=SC1 FETURN END TYPE DEFSET 000000 OFFSET 000000 060000 Ĭ 8 Xiona Mouno 6•3 C **J**AY 2 Ĕ 3 . . 0000 Way a 2000 W CEXP ŝ Ţ 23 ŝ 5))

3) 2 END K44-", F15. 7 • ビーン 71,75 7#U 010074 010072 010010 010035 010042 OFFSET 010022 OFF3E1 SUPPOUTINES. FUNCTIONS. STATEMENT AND PROCESSOR-DEFINED FUNCTIONS 7/* FSTART=*,F15.7; FEND=*,F15.7/* END `` NAME CABS TYPE CONTINUE FORMAT(///* RADIUS=',F15.7/* FSTART=',F; 'STARTKAA=',F15.7/* FEND=',F15.7/* '' NUMPTS=',17////) FORMAT(//75(*+')/* POINT *', FREQUENCY ', FORMAT(//75(*+')/* POINT *', FREQUENCY ', TYPE VARIABLES, PSECT ODATA, SIZE = 010112 (2085. WORDS) 8 8 8 4 844 4*2 4.8 NUMPTS 1+2 OFF3ET -----SIZE----- DIMENSIONS 004000 004000 (1024) (512) 000000 004000 (1024) (512) FSTART STEPK 4=2 TVPE MAR MAME Å ž â /. SIZE = 000014 (& WORDS) **ATAN2** MAME ***** END OF PROGRAM +++++ OFFSET OFFSET 010062 010000 010040 010060 010052 010046 00000 TVPE ASSIGN R+4 TYPE TYPE *** **4**•2 R=4 1*2 8+0 I#2 ¥•2 MAME SECTION OFFSET TINUI STEPF MARE MAN FREG FORMAT (AL) ŧ U z 4.2 NAME TYPE 4.5 LOCAL AND COMMON APRAYS: SDATA SDATA AL0010 TYPE OFFSET TYPE OFFSET 010015 010014 000000 010026 010055 010004 010066 010032 SCAT COMMON BLOCK / END. TYPE 1-4 ŝŝ TYPE . **t**•d 4 - d **4**•d 1.2 1+2 1 ž 22 4.0 ŝ LOCAL Sangle Burgle 7 F 8 8 8600 **GAMIA** 868 868 868 20.57 **JHVH** FEND Mar BHOH BHON PEAL £ 50.7 ₩ Be 2 2 2 3 Ą (C (C C C C

ALS FROM AS AND I SUMMER PARCELOND JF (kA GT (8)) 00 T0 200 P1=(? ∩/2 ∩)+(KA++3 0)+(1-(5 0/54 0)+(KA++2)+(17 0/700 0)+ 1(KA++4)-(6 551922€+06/11 907€+06)+(KA++2)) P2=(1 0/2 0)+(KA++6)+(1+(6 0/5 0)+(KA++2)) P2=CPL(101,P2) G0 T0 1000 THIS PROGRAM WILL CALCULATE THE SCATTERED FIELD FROT A PEFFECTLY CONDUCTING SHERE THIS IS A FAR FIELD APPROXIMATION AFTER RADAR CROSS SECTION HANDBOOK OCOMOLE T RUCK EDITOR PLENUM PRESS 1970. PAGES 146-155 PAGE 001 COMMON KA, F COMPLEX F. T1, T2, T3, C1, C2, C3, C4, C5, C6, C7 COMPLEX CON, CON1, CON2, CON3, CON4, B1, B2, F01, FC1 COMPLEX CON5, CON4, CON2, CON4, CON4 REAL KA BI=CON#(1/12#X)) B2=CON#(12#X) F01=(X/2)#CEYP(B2)#(1 0=B1) #******* (FEFFING MAVE TERM FC1 ******** IF (K4 GT. 20) C0 T0 300 ******* 0PTICAL TERM C1 ******** C1+-CON+/***(4 0/2 0)/*CEXP(B1) C2=1 257588+(CON1*XM23) FORTRAN IVVO2: 1-1 THU 18-OCT-79 00: 03: 33 0001 SUBROUTINE SCAT XM3244441-1.0/3 0) XM3244441-2.0/3 0) CON1++ 741194,1 253789 CON2+(2 200022 1 270172) CON2+(2 200022 1 270172) CON3+(445394, 1 570720) CON3+(445394, 1 570720) CONFE(7, 012224, -4, 049663) CONFE(444477, 25,419) CONTE(759521, 1, 255545) CONFE(5, 048954, -7, 315,114) CON=(0.0.1.00) FF (FA GT .4) G0 T0 100 FI=(3.0.2.0)+(KA)++3 F=CMELX(F1.0.0) G0 T0 1000 CON--- 312321, 180319) VERSION 1 3 0CT-17-79 (10 \$10 11-1)+[a=[d 13=K4++11 015 01 PLF? 1415026 PI-CONSCR-1 Y=Y.O ی ۲ ۵ ۵ ۵ $\omega \omega$ u 000000000000 0018 0018 0018 0019 0041 0041 0010 0010 0011 0013 0021 0022 0023 0023 2000 1100 0012

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THIS PROGRAM WILL MEASURE THE PATTERN OF AN ANTENNA USING A STEPPER MOTOR CONTROLLED TURNTABLE AND THE MEMETT-PACKARD 84200 MICHONAUS SWEEPER AND 84109 METWORK ANALYZER IN THE AND ANALYZER IN THE AND 84109 METWORK ANALYZER IN THE : TYPE *, '****** ANTEMNA PATTERN MEASUREMENT SYSTEM ******** Type *, IF(IFLAD E0.0)60 TO 500 TYPE 4. ENTER THE FREQUENCY OF OPERATION IN GM2: * Accept 4.Fort Type 4. 'Enter The Number of Data Points: * ACCEPT *.N ACCEPT *.N TYPE *. 'ENTER THE MUMBER OF SAMPLES/DATA POINT:' ACCEPT *.NSAMP ACCEPT *.NSAMP ACCEPT *.ANGLE ACCEPT *.ANGLE TYPE *.'ENTER THE FILENAME FOR THE DATA:' CALL ASSIGN(15, --L, NEW, 'NC', 1,) TYPE *, ***** IS THE ANTENNA POSITIONED ? ***** PAUSE '****** HI RETURN TO CONTINUE ***** THERAPHOLES. 1415926/190. O DTHETAATHETA/(N-1) PAGE 001 TYPE +, ^OLD OR NEW DATA? (O=>OLD 1=>NEW) ^ ACCEPT +, IFLAG 00 TO 700 TYPE +, 'ENTER FILENAME FOR DATA:' FORTRAN IVVO2. 1-1 WED 02-APR-80 00:02:41 0001 PROGRAM ANTPAT INTEGER AMP(512), PHA(512) BYTE STRING CALL FHAMP2(14, 1P, NSAMP) AMP(INDEX)=1A-2048 PHA(INDEX)=1P-2048 CALL STEP2(DTHETA, 0) CONTINUE CALL STEP2(THETA/2, 0, 1) AUTHOR C. WERNER 6-MAR-80 CALL SWEEP(1, FOPT, 1, 4) CALL STEP2(THETA/2, 0, 1) DO 100 1=1, N CALL SWEEP(10. FOPT. 1. 4) 00 200 J=1.N JRITE(15.+)AMP(J) RITE(15, +)PHA(J) WRITE(15. +) ANGLE WRITE(15. +) N WRITE(15. +) NSAMP HRITE(15. *) FOP1 ţ ACCEPT +. CONTINUE INDEX=I TYPE*, A=0. 0 ļ 8 800 8 υu 0000 0003 5400 1400 0028 0033 1200 6400 6400 0043 0044 0029 ÚÊ00 0038 0039 0042 0036 0037 1600 ? 2 0) OFFSET 001100 000130 003240 012100 0000200 001000 012(00 020(-00 1111 000000 0030200 TYPE OFFSET FLANCTIONS. STATEMENT AND PROCESSOR-DEFINED FUNCTIONS MAN TYPE VAPIATLES. PSECT STATA. SIZE = 000424 (139 WORDS) 8 * U 8 * U 0 1 1 8*0 8 8 4 8#0 ##¥ 2 TYPE キャッシュ キャンロフィード NAME CONG CONS CONS MAM NOU £ FOI 2 SIZE = 000014 (6. WOPDS) £ ŝ 2 NAME TYPE NAME TYPE OFFSET 002000 010000 TYPE OFFSET 0011000 000230 000260 000040 000220 002000 000000 000324 000000 (VEND3+EHX)+(2NU)+07-774A3-560 144 \$45864+(CONA#XM23) 154584044-Y34CON45)-(YM24CON5)) C\$= &(~104+(C0N7+XM23) (7=[E¥P+++7+C0N9)-(XM2+C0N9)) Beach and a state with a back F=-(1 0/2 0/+/A+CEXP(B1) 8 1 1 0+3 88 0 8 8. 2. 800 8.0 3 482 5 444 و: ورا •و\$1 ور: درا •و\$1 (+يت-يخ) NAME MAN CON2 VM23 CONSC SNU: 5 82 8 g 7 ā 347T L ŝ . 10404040 F1=12+15 73454407 <u>nerser</u> Buch 0211200 061000 091000 000250 **ΰ**ευύψυ ON CALCARD 0110000 **UEE**UUU 000000 015000 000000 12000 ABL 1 JAN 1 0-0-、 EN C 10012 Signitudes. TUPE アイトビ JYPE 0) * () 2 **-** 3 1000 01 10 10 10 * C 8+0 0.0 500 01 • • 10 20 7+3 4.4 ** 1 ¢. e ۰. ۰. NUMBER OF 900 77.5 SA02 5 5 \$ 1200 1.... ÷, n de la Swo? 1.1 ŀ Sec. 2 CEYP 1 . 6-2 ç С

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IF (ТАКР(К) 01. 720) ТАКР(К)=ТАКР(К)-1440 IF (ТАКР(К).LT -720) ТАКР(К)=ТАКР(К)+1440 RN06M=10 0++(((ТАКА(К)+05)-АМР)/10.0)+COS(TARP(К)+RADC) 001200 06:2040 00120 OFFSET 002074 002072 002024 TYPE CONTINUE CONTINUE AMP-AMP/NPTS TYPE *, AVERAGE AMPLITUDE IN DB: ', AMP TYPE *, 'ENTER DYNAMIC RANGE FOR DISPLAY IN DB: ' 566 WORDS) R#4 4 * X н. Н. ĩ 1+2 **#**# DIMENSIONS DO 500 J=1.NLINES CALL SWDAT(TARA,TARP,FBEOIN,FSTOP,NPTS,NSAMP) DBNORM (128) (128) (128) AMMAX FSTOP SCALE NSAMP MAME IF (ТАRР (К) GT 720) ТАRР (К)=ТАRР (К)-1440 IF (ТАRР (К) LT -720) ТАRР (К)=ТАRР (К)+1440 АМР=АМР+(ТАRА(К)*.05) <u>(</u>20 AMAX=10 0++(AMMAY/10.0) AM1X=10 0++(AMMAY/10.0) TYPE - XAMAX='AMAX', AM1N=',AM1N TYPE - XAMAX=AM1N) TYPE - SCALE FACTOR FOR DATA=',SCALE . 1288128 VARIABLES, PSECT \$DATA, SIZE = 002154 **** END OF PROGRAM ***** OFFSET 002114 002126 002066 002034 002044 002124 SIZE-000024 000400 000400 000400 000400 TYPE rarp(k)=tarp(k)-REFP(k) TARA(K)=TARA(K)-REFA(K) TARP(K)=TARP(K)-REFP(K) 4*2 R#4 R#4 FBEGIN R*4 1+2 1=2 COLLECT HOLOGRAM DATA CALL STEP2(DTHETA, 0) CONTINUE [Z=IFIX (RNORM+SCALE) 000000 001000 001000 OFFSET 002000 RNORM NAME AMIN NP1S ACCEPT +, DENORM AMMAX=(DENORM/2) D0 100 K=1. NPTS đ CLOSE (UNIT=15) STOP **** END 7 XUMMU-EN I MMU WRITE(15) IZ SECTION (AND COMMON ARRAYS: TYPE +, K, 12 DATA
 DATA
 DATA
 DATA OFFSET 002110 002062 002104 002056 002132 002052 002046 UN11NU Ñ TYPE TYPE 4.4 4 # H R#4 4 DTHETA R+4 I#2 1#2 ŝ Ŧ 8 8 20 000 NL INES LOCAL 1 LOCAL AMMIN THETA NAME RADC REFP TARA TARP 0038 0038 0040 0040 0040 0040 AMAX REFA NAME NAME ZI オード C Э)) t • Carvair's Libbor. THIS PROGRAM WILL FIND THE FREQUENCY SWEPT HOLOGRAM OF A TAPGET WITH VIRTICAL SYNHERY. THE FREQUENCY RANDE WILL BE BETWEEN AND FSTOCAL SYNHERY. THE FREQUENCY PTS. THE ANOLE PANGE WILL BE THETA WITH MLINES IN THAT ANGULAR SWEEP THE FEAL PART OF THE SCATTEED FILL BE SCALED FOR DISELAR BY PROGRAM CDISP. THE TANGET WILL BE CORRECTED FOR SYSTEM RESPOSE AND RANGE BY USE OF A REFERENCE TARGET TVPE *. **** PLACE REFERENCE TARGET IN THE FIELD **** Pause **** Hit return to continue **** Call Sudat(refa, refp, fregin, fstop, npts, nsamp) OPEN(UNIT=15, NAME=NAME, TYPE='NEW', ACCESS='SEQUENTIAL', [FOPM='UNFOPMATTED'] ACCEPT + FSTOP TYPE + CENTER THE NUMBER OF FREQUENCY POINTS (C128); C . 'ENTEP THE MUMBER OF SAMPLES/FREQUENCY POINT: PT +, NSAMP +, 'ENTER THE NAME OF THE FILE FOR DATA STORAGE: TYPE «, ***** PLACE IMAGE TARGET IN THE FIELD **** PAUSE ***** HIT REFURN TO CONTINUE **** Call sudat'tard, tarp, feboin, fstop, npts, nsamp) Padde= 23*3. 1413926/180. 05 INTEGER REFA(128), REFP(128), TARA(128), TARP(128) B/TE (IAME/20) PAGE 001 WPITE(15) FREGIN.FSTOP, MLINES, THETA, NPTS DTHETA=(THETA/NLINES)+(3, 141525/180,) . 'ENTEP ANGULAR SHEEP IN DEGREES: ' TYPE 0. THE THE NUMBER OF LINES: / ACCEPT 0. NLINES ACCEPT 0. NLINES ACCEPT 0. NLINES ACCEPT 0. YAMP ACCEPT 0. YAMP ACCEPT 0. YEAHP ACCEPT 0. YEAHP ACCEPT 0. YEAHP ACCEPT 0. YEAHPAN ACCEPT 0. FORTRAN IVVO2 1-1 . THU 21-FEB-80 05: 03: 15 0001 PROGRAM CVEXP GET SCALE FACTOR FOR DISPLAY AUTHOR C WERNER 21-AUG-80 [ARA(K)=TARA(K)-REFA(K) GTLIN(NAME, 221) GET SYSTEM PESPONSE +. THETA ACCEPT * NPTS TVPE * 'ENTEP DO 10 K=1.NPTS ACCEPT ACCEPT O=dub TVPE TYPE CALL 00 000 000 0000 0000 0023 0023 m21 m22 0032 (ſ C C

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IF (A.EQ. 'N') GO TO 30 Type +, 'Enter the Length of the Line (0-1000);

0059 0059 0059 0060 0061

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PAGE 001 FORTRAN IVVO2 1-1 SAT 23-FEB-80 01:31:06 2001 PROGRAM CDISP

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AUTHOR C. WERNER 19-FEB-BO

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TVPE 558

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CALL GTLIN(NAME, /?/) OPEN(UNIT=15, NAME=NAME, TYPE=/OLD', ACCESS#'SEQUENTIAL',

(^ CARM= ' UNFORMATTED')

PEQUISTSTAPT.FEUD.NLINES.ANGLE.NPTS TYPE +. FPEQUENCY GHZ (START): ', FSTART TYPE +. FREQUENCY GHZ (START): ', FSTART TYPE +. NUMER OF LINES: ', MLINES. 'POINTS/LINE: ', NPTS TYPE -. SWEEP ANGLE DEG. ', ANGLE NPOINT=NLINES*NPTS

DO 10 J=1, NPOINT PEAD'15)2(J) CONTINUE

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OCAL

0087

MAME

0081

0071 0073 0074 0074 0077 0077 0077 0080

TYPE *, "LOG COMPRESSION (Y OR N):" 0700 623

IFIA ED 'Y') 50 TO 25 TYPE -. ENTER CONTRAST SCALE FACTOR: " ACCEPT -. SCALE

SCALF=SCALF+

00 20 J=1.NP0INT Z(J)=IFIX(SCALF+Z(J)) IF (Z(J).5T 100)Z(J)=100 IF (Z(J).LT -100)Z(J)=-100

3

CUTINUE CUTINUE Ē

E + ENTEP LOG COMPRESSION SCALE FACTOR: * FPT + SCALL 28 J=1.NPOINT 3

F(1) LT ()2(J)=IFIX(-(100 0+AL0010(1, 0-SCALL+Z(J)/10, 0)))
F(Z(J) GE ()2(J)=IFIX(100 0+AL0010(1, 0+SCALL+Z(J)/10, 0))

CHIT THUE ន្តន

(PE +, 'ENTEP THE INITIAL Z VALUE (500-1000)' CEPT + IZINIT

+, "ENTER THE STARTING X POSITION (0, 0-1, 0)

CEPT + XSTAPT

INS#IFIX(NSTART+40%6) TVPE +, 'ENTER THE START Y POSITION (0.0-1.0): ACCEPT +, YSTART

IVS=IFIX(VSTART=4096) CALL IDISP(IXS, IVS, IZINIT)

0053

TVPE ... POSITION AND INTENSITY CORRECT (Y OR N): ACCEPT 1000.A

040064 OFFSET 040114 040042 040100 040052 040102 TYPE (8259. WORDS) R=4 444 I +2 NLINES 1+2 SCALF R+4 VSTART R+4 ****** END OF DISPLAY PROGRAM ******* FSTART DFREQ MAME SXI CALL TITSP(1XS, IYS, IZINIT) TYPE *. DISPLAY SAME DATA AGAIN (Y OR N) ACCEPT 1000, A ACCEPT +, ILEN PAUSE 'HIT RETURN TO CONTINUE : ' DTHETA=ANGLE+PI/(FLOAT(NLINES)+180) ILENS-IFIX(ILEN*(FSTART/DFREQ)) ILENS-IFIX(ILEN*(FSTART/DFREQ)) O 400 41).41.NETA INEFA:(J-1).00THETA IXEEG=IXS-IFIX(ILENS+COS(THETA)) IXEEG=IXS-IFIX(ILENS+SIN(THETA)) D 200 1NUE*1.NFTS D 200 1NUE*1.NFTS IZ(INDEX)=Z((J-1)*NFTS+INDEX) VARIABLES, PSECT \$DATA, SIZE = 040206 IF (A. EO 'Y') 60 TO 30 TYPE +. RESCALE DATA (Y OR N) / ACCEPT 1000, A 040046 040122 040062 040074 OFFSET 040054 040024 TYPE R#4 4#X F (A. EQ YY) 60 TO 15 4*2 4=2 1+2 1#2 XSTART INDEX ANGLE DFREQ=FEND-FSTART **MAR** FEND 2 CALL LINDSP CONTINUE CONTINUE OFFSET 040040 040110 040120 040060 040070 040106 STOP END TYPE 1000 3 R#4 I#2 **1**#2 4 # U NPOINT 1+2 ş 8 DTHETA ILENS

WORDS) BLOCK /DISPL /, SIZE = 002016 (519. COMMON

SCALL

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Sλ

002002 010200 OFFCEV TYPE 1+2 1#2 NARE ILEN NPTS OFFSET 00200 002006 TYPE IZINIT 1+2 IYBEG I+2 MAME OFFSET 000000 002004 002012 TYPE CUTHETA R#4 1+2 I#2 I X BEG MAME 71 3 Э

LOCAL AND COMMON ARRAYS:

SECTION OFFSET -----SIZE----- DIMENSIONS DISPL 000000 002000 (512.) (512) TYPE 1+2 NAME 21
	PAGE (101																				20. WORDS)	AME TYPE OFFSET	ON 1+2 0:0022	WCH 1+2 0:0010	CCH 1+2 0:0006	ADROT R+4 € 00000		CEFINED FUNCTIONS.
	RADROT, NDIR)	L-FEB-80	IN RADIANS _AG, O=>CW 1=>CCW							159					10FF)		ION)	ĺ	10-1-1		, SIZE = 000050 (TYPE OFFSET N	I*2 000024 I	I 1*2 000012 1	1+2 000032 M	R 000002 R		CHENT AND PROCESSOR-D
	2. 1-1 SUBROUTINE STEP2(F	AUTHOR C. WERNER 21	RADROT* ROTATION NDIR= DIRECTION FL	MCW=1024	MCCW=1536	IWCW=1280 IWCCW=1792		STEPS/RAD=10000/P		STRD=10000.0/3.14	ISTEFELFIXISTRUPR TE (NDTR) 10,10,20	ION=MEM	ICFF=IWCW Go TO 22	ION=MCCW	IGFFELW.CW CALL DOUT (., IERR.	D0 25 J=1, 200 CONTINUE	DO 500 K=1, ISTEP CALL DOUT(, , IERR,	CONTINUE	CALL DOUT (., IERK. DO 70 J=1,200	CONTINUE . CONTINUE . RETURN	END N FSCT \$DATA	OFFSET NAME	000026 IOFF	000020 INCCH	000030 K	000004 NDIR	000014	FUNCTIONS, STATE
	RTRAN IWO	000	0000	- , c S	50	4 K	0	00	0	909	200	04 10	010	012 20	113	015 22 016 25	017 30 018	020 50	021	023 224 500 225 500	026 Drai Variar	AME TYPE	ERR 1+2	STEP 1+2	1+2	CM 1+2	TRD R+4	UBROUT I NES,
(100	Ϊð			ð	č	ŏč	÷			Ō	õč	00	ĊĊ Ć	0	00	00	c 0 '		ос		o _	, z	1		י י י		Ŋ.)
PAGE		TA .														OFFSET	000020	000042	000024			OFFSET	002002	00200				TIONS.
	TA .	LEN. TH													WORDS)	TYPE	R*4	I#2	R#4			TYPE	1+2	I+2			SIONS	ed Func
3: 02: 43	LINE OF DA	EO, IYBEO, 1							•						LINDSP 5 (23.	NAME	5	IZVAL	×	c	MORDS)	NAME	NPTS	ILEN			DIMEN	SOR-DEFIN
-	_ ح	IXB									~	(0			00KAM UNI E = 000054	E OFFSET	10000	000040	000010		16 (519	E OFFSET	002000	002006			SIZE	AND PROCES
2-AUG-79 1	ISPLAY	, NPTS,													ΥN	Ā	4.5	1+2	R*4		020	μγ	ŝ	ŝ			Ιŝ	3 5
WED 22-AUG-79 1 UDSP	NE WILL DISPLAY NTS LONG.	(2. 121NIT, NPTS. 1)	23	IPTS)	Ś	•			INIT	17. IZVAL	EG, IVBEG,			ATA S	ų ų	. .		Z		ة • س	Ŵ	TINI	BEO 1			FSET	ATENED
V VO2 1-1 MED 22-AUG-79 1 SUBROUTINE LINDSP	up to subroutine will display	COMMON/DISPL/IZ, IZINIT, NPTS, INTEGER IZ(512)	COSVIETA) COSVIECOS(THETA)	DY#ILEN*SINV/NPTS	X#FLOAT (IXBEG)	Y≑FLΩAT(IYBEG) Dů 100 K=1,NPTS	xu+x=x	Y=Y+UY IX=IFIX(X)	IV*IFIX(V)	IZVAL=1Z(K)+1ZINIT	CONTINUE	CALL IDISP(IXBEG, IYBEG,	RETURN		LES, PSECT +DATA, S	OFFSET NAME 1	000004 DX	V100036	000034 SINV	06.0000	/DISPL /' SIZE = O	OFFSET NAME	TINIZI 000000	002004 IYBE0 1	002012	SABAR NOM	E SECTION OFFSET	FUNCTIONS, STATEMEN
OPTRAN IV VO2 1-1 WED 22-AUG-79 1 SUBROUTINE LINDSP	Č THIS SUBROUTINE WILL DISPLAY C UP TO 512 POINTS LONG.	COMMON/DISPL/IZ. IZINIT, NPTS. INTEGER IZ(512)	C COSV=COS(THETA) STNU=STN(THETA)	DY=ILEN+COSVINTS DY=ILEN+COSVINTS DY=ILEN+SINVINTS	AFFLOAT (IXEG)	V ♦FLΩAT(IVBEG) Dû 100 K=1,NPTS	x (1 + x = x	Y=Y+UY IX=IFIX(X)	[\\]]	IZVAL=IZ(K)+IZINIT	100 CONTINUE	CALL IDISP(IXBEG, IYBEG,	RETURN		L VARIABLES, PSECT +DATA, S	TYPE OFFSET NAME	R+4 000004 DX	I+2 000036 IV	I+2 000034 SINV	R+4 000030	GN BLOCK /DISPL /' SIZE = O	TYPE OFFSET NAME	1+2 000000 Isinit	0 I+2 002004 IYBE0 I	A R+4 002012	L AND COMPON ARRAYS:	TYPE SECTION OFFSET	DUTINES, FUNCTIONS, STATEMEN

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