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A Numerical Solution for Rocket Ascent Trajectory

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August 16, 1979



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INTRODUCTION

In a previous report [1] the author presented a theoretical framework for analyzing the trajectory of rocketsatellite systems, both in the powered rocket ascent or launch phase and in the subsequent orbital motion or coast phase. Techniques for solving the problem of launching space systems into proper earth orbit were elucidated, and subsequently these techniques have been tried on particular cases. In the process, a numerical solution of the equations of motion of rocket ascent has been programmed on a digital computer and found to be useful. It is reported here to make it more generally available.

This is obviously not the first time a computer program has been developed to handle the rocket ascent problem. It is possible to get one of these programs from an outside source, and adapt it to the specific tasks at hand. Very often, however, the information accompanying the program is insufficient, or the requirements of the program are not matched to the data base for the problem or the computer at hand. It is frequently easier, faster, and more reliable to develop a program ab initio. This was done, and the resulting program is presented and explained below.

It was shown in Reference [1] that the rocket ascent problem could be solved as if the earth (and atmosphere) were stationary in an inertial frame of reference, and that the corrections from earth's rotation could be added in later. The computer program developed here solves the launch problem for the stationary earth case. It will be shown later in a specific application how earth's rotation effects are adjoined to the computer program results. For the stationary earth model the simplest case of planar rocket motion is shown in Fig. 1. Here r is the radius vec-Note: Manuscript submitted March 26, 1979. tor from the center of the earth to the center of mass of the rocket, y is the rocket altitude, ϕ is the angular displacement from launch, which determines the ground range x when the radius of the earth R is factored in, ψ is the flight path angle of the center of mass (or "heading") of the rocket with respect to the local horizontal, and α is known as angle of attack of the rocket axis with respect to the flight direction.

The equations of motion for the rocket are deduced [1] straightforwardly from the force diagram shown in Fig. 2, which also contains a brief description of the indicated symbols. The forces indicated are: the gravitational force W, the thrust of the escaping exhaust gases T, which is canted at the angle β to the vehicle centerline, and the aerodynamic forces of drag D and lift L due to rocket motion through the atmosphere. Drag and lift forces are directed antiparallel and perpendicular to the flight path direction, respectively. In deducing the equations of motion, we make the assumptions found in many references [1-4]:

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- the aerodynamic forces act through the center of pressure,
- (2) the force of gravity acts through the center of gravity, and
- (3) the thrust force is applied through the center of combustion.

The equations of motion for the center of mass motion are determined [1] to be:

$$\Psi = -(K/r^2) \sin \psi - (D/M) + (T/M)\cos(\alpha + \beta)$$
(1)

$$\mathbf{v} \quad \mathbf{\dot{\psi}} = - \frac{\mathbf{K}}{\mathbf{r}^2} - \frac{\mathbf{v}^2}{\mathbf{r}} \cos \psi + \frac{\mathbf{T}}{\mathbf{M}} \sin (\alpha + \beta) + \frac{\mathbf{L}}{\mathbf{M}}$$
(2)

The speed of the rocket center of mass is here denoted by v, and the mass of the rocket is denoted by M. The symbol K is the product of the universal gravitational constant

-2-

and the mass of the earth, and has the value

 $\kappa = 3.986012 \times 10^5 \text{ km.}^3/\text{sec.}^2$

The other symbols have already been defined. The angle of thrust β is normally very small for large launch vehicles (e.g., $\beta/\alpha < 1$), and it is customary [2,3] to set β to zero in the equations of motion before solving for the launch trajectory. Hence, the good approximation is made

(3)

₿ ≈ 0

in Eq.'s (1) and (2), thus avoiding the complications of additional moment equations for the steering of the rocket. The variation of β is preprogrammed in the guidance and control system of the rocket, which causes the relatively much larger variations in α involved in the shaping of the launch trajectory. The approach to be used in the numerical solution of Eq's (1) and (2) is in this spirit. An $\alpha(t)$ profile will be selected to cause the relatively much larger variations in $\psi(t)$, and it will be done in such

a way as to obtain a correct value of ψ at rocket burnout

and a smooth variation in $\psi(t)$. The program user typically knows this burnout value of ψ from consideration of the orbit into which he desires to inject the payload. The procedure will be: (1) selection of an $\alpha(t)$ profile, (2) numerical solution of v(t) and $\psi(t)$ from the equations of motion, and then from orbit injection or other criteria, (3) selection of a new $\alpha(t)$ profile and a repetition of step (2). This process is iterated on the computer until a "correct" launch trajectory is obtained.

Other approximations in the solutions of Eq's (1) - (3)have been experimented with during the evolution of this program. One approximation is the choice of specific, analytic profile for $\psi(t)$ with free parameters adjusted so that $\psi(t)$ takes on specific end values. The known

-3-

end values are those at rocket launch and at orbit injection. The problem is that all the other values of $\psi(t)$ in between are forced by the analytic form; this will be elaborated on shortly. The advantage of assuming an analytic $\psi(t)$ profile is that only Eq. (1) needs to be solved if some other information about a(t) is known - e.g., that a is small enough so that cos and (recall Eq. (3)). The criterion of small a values is presumably desirable from the standpoint of launch efficiency; the rocket engine then does its work along the direction of motion. Ehricke [3] uses this approach, but points out that α could be calculated from $\psi(t)$ and the solution for v from Eq. (2). In particular, the assumption about the smallness of α could be checked. Ehricke uses a particular analytic form $\psi(t)$ for each stage of a multistage rocket. Experience with this approach, however, points out certain flaws in it. In the first place, the end points of $\psi(t)$, which nail down its free parameters, are typically unknown ahead of time. Hence, there is much time wasted experimenting with end values (i.e., stage ignition and burnout values of ψ) which will keep the range of α values small. A major problem is that the analytic form of $\psi(t)$ is apparently not very realistic, because it forces some rather wild variations of $\alpha(t)$. Neither is the criterion of smallness of $\alpha(t)$ necessarily realistic, particularly in the higher rocket stages when substantial tilts may be necessary to inject payloads at sufficiently high altitudes to avoid aerodynamic heating or drag effects. Ehricke's method was thus found to be too slow and inefficient in the trial and error aspect, not sufficiently unambiguous for inclusion of an automatic self-correcting procedure based on a small a criterion, and potentially too inaccurate for our purposes. Ruppe [3] similarly employs analytic forms for sin ψ and cos ψ which are polynomial expansions in t designed to give correct launch and orbit injection end point conditions. With

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certain other approximations this enables one to integrate Eq. (1) analytically. Unfortunately, the large disagreements in the ψ values inferred from the separate polynomial expansions for sin ψ and cos ψ and the inaccuracies in the method seemed unacceptable for our purposes.

Other approximations, often made in the solution of Eq's (1) - (3), involve the evaluation of the force terms in these equations. The evaluation of forces in the program will be detailed in the next section, but for present purposes it will suffice to just mention the qualitative reasoning behind the approximations. Due to ambient pressure decrease with altitude it is found that the effective rocket thrust force T in Eq's (1) and (2) increases from its launch value T, at sea level to its vacuum value T, at the orbit injection altitude. This increase of thrust effect is at least partially offset by the drag force in Eq. (1) which acts to slow the rocket down. This has prompted certain authors [3,4] to eliminate the drag and thrust increase effects from Eq. (1), since they tend to counterbalance each other, and use a constant thrust term instead. This is a nice simplification, particularly since the drag and lift coefficients for the rocket, which enter in the evaluation of D and L in Eq's (1) and (2), are usually not known very well. Similarly, the force L is often dropped from Eq. (2) on the grounds that it is not

very important for large rockets. It has been the author's experience that for large rockets, e.g., with liftoff weights of hundreds of thousands of lbs. or more, the thrust increase effect outweighs the drag effect. As a result, the constant thrust approximations give injection velocities smaller than the more complete evaluations by

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as much as 2-3%. This kind of discrepancy is serious for the orbit determination, as shown in Fig. 3, which gives the dependence of orbit apogee altitude on perigee injection velocity. A two percent variation injection speed can make a difference of 500 nmi. in apogee altitude. In questions of orbit determination (or even whether or not a stable orbit has been achieved) this kind of uncertainty may well be unacceptable. As Eq. (1) indicates, the thrust term also varies when the angle of attack a becomes large, which is often the case in the upper rocket stages. The constant thrust approximation was therefore abandoned. The decision was also made to not discard the lift term L in Eq. (2), since it was found [3] to significantly enhance the leverage exerted by the assumed $\alpha(t)$ profile on the calculated $\psi(t)$ solution at ascent altitudes where aerodynamic effects attain their largest values.

For the above reasons the approach adopted in the computer program is to assume an $\alpha(t)$ profile, solve the equations of motion, and then refine the original $\alpha(t)$ in an iterative approach. This proved to be easy to automate, and the results did not appear to be sensitive to the ambiguities in choosing the $\alpha(t)$ profile. With an evaluation of thrust increase and aerodynamic effects, and with a nonreliance on the constant thrust approximation, the program has reached a level of approximation where it should be useful for most purposes, including orbit determination. More on this later, but the next task is to supply details about the program and its use.

The Evaluation of Terms in the Equations of Motion

For a rocket with swivel control motors the total thrust T in Eq's (1) and (2) includes pressure forces, and is written [3-5] as

$$T = Mv_e + A [p_e - p(y)] ,$$

where M is the mass loss rate of exhaust gases, v_e is the actual average axial speed of these gases relative to the rocket, A is the exit aperture area, p_e is the pressure of the exhaust gases at the exit aperture, and p(y)is the ambient pressure of the atmosphere at altitude y. This pressure varies from p(o) to zero during the first part of rocket flight, usually by the time of first stage separation. Consequently, T varies from its sea level value T_{s1} to its vacuum value T_{vac} . Both T_{s1} and T_{vac} are normally specified for the first stage of the rocket, possibly also for the second stage, but usually only T_{vac} is given for higher stages. Hence, for the first stage, possibly also the second, the ratio

$$x = T_{vac}/T_{sl}$$
(5)

(4)

is known, and Eq. (4) can be rewritten as

$$\mathbf{T} = \mathbf{T}_{vac} \begin{bmatrix} 1 - \frac{\mathbf{x} - 1}{\mathbf{x}} & \frac{\mathbf{p}(\mathbf{y})}{\mathbf{p}(\mathbf{o})} \end{bmatrix}$$
(6)

The computer program evaluates thrust from this equation, which exhibits the increase of thrust with altitude.

The aerodynamic forces of lift L and drag D in Eq's (1) and (2) can be evaluated from [3]:

$$D = [C_{D0} + C_{DL} \alpha^2] (\rho v^2/2) S$$
(7)

$$L = [(\partial q, /\partial \alpha) \alpha] (\rho v^2 / 2) S$$
(8)

where ρ is the atmospheric mass density, v is the rocket speed and S is a reference cross-sectional area (e.g., that for the lowest rocket stage) to which the aerodynamic coefficients are referred. As seems plausible, the aerodynamic forces are seen to be proportional to the dynamic

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pressure of the atmosphere (i.e., $\rho v^2/2$) and the reference cross-sectional area of the rocket vehicle. The proportionality constants in square brackets in Eq's (7) and (8) are the drag and lift coefficients which depend on both the angle of attack α and the speed v. The dependence on α is to be expected, since the tilting of the rocket centerline away from the flight direction obviously exposes more of the rocket surface to the pressure forces of the atmosphere, thus increasing the aerodynamic forces. The lift force vanishes for $\alpha = 0$. The representation of the α dependence of the drag and lift coefficients in Eq's (7) and (8) is supposed to be valid, according to Ehricke [3], for moderate angles of attack (e.g., $\alpha < 10^{\circ}$). Here C_{DO}, C_{DL}, and $\partial C_{r}/\partial \alpha$ are constants which depend only on the speed of the rocket. Actually, the dependence is found to be on the mach number M, defined as

 $M = v/c , \qquad (9)$

where c is the speed of sound at the altitude of the rocket. We have used the dependences given by Ehricke (cf. his Fig. 5-8 in [3]) for a two-stage rocket, a close analytic fit to which has been found to be:

 $C_{DO} = \begin{cases} 0.25 + 0.3773 \exp \left[-16.6426 \left(M-1.14\right)^{2}\right] & (0 \le M \le 0.88889) \\ 0.6273 - 30.902 \left|M-1.14\right|^{3.5} & (0.88889 \le M \le 1.28) \\ 0.17913 + 0.5318/M & (1.28 \le M) & (10) \end{cases}$ $C_{DL} = \begin{cases} 1.822 - 0.298 \cos \left[M \pi/1.086\right] & (0 \le 1.5) \\ 0.2818 + 3.2727/M & (1.5 \le M) & (11) \\ (0 \le M \le 1.5) & (11) \\ (0 \le M \le 1.5) & (12) \end{cases}$ $DC_{L} / \partial \alpha = \begin{cases} 0.8 + 6/M & (M > 1.5) & (12) \\ M > 1.5) & (12) \end{cases}$ The graphs of Eq's (10) - (12) are shown in Fig. 4, along with a sketch of the type of rocket to which these coefficients apply. In this case, the reference value S in

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Eq's (7) and (8) is the cross-sectional area of the bottom

rocket stage. The dynamic pressure factor in Eq's (7) and (8) is seen to involve a competition between an atmospheric density ρ which is decreasing exponentially with altitude and a quadratic dependence on velocity which is increasing with altitude. The result of this in calculations is that aerodynamic forces are negligible above about 60 km, and peak somewhere around 10 km.

For the evaluation of thrust and aerodynamic forces an atmospheric model is needed, i.e., a specification of pressure ρ and density ρ vs altitude y in the region in which thrust variation and aerodynamic forces play a role (e.g., for $0 \le y < 100$ km). The temperature is related to p and ρ by the ideal gas law;

$$p = \rho RT/M$$
(13)

where M is the mean molecular weight of one mole of atmospheric gas, and \hat{R} is the gas constant given by $\hat{R} = 8.3143$ joules mole⁻¹deg⁻¹= 8.3143X 10⁷ergs mole⁻¹deg⁻¹. The speed of sound in Eq. (9) is calculated from $c^2 = 1.4 p/\rho$ (14)

where the factor 1.4 is the ratio of specific heat at constant pressure to the specific heat at constant volume for diatomic molecules, such as air. For use in the computer program it would be helpful to have an analytic expression of the atmospheric density and pressure altitude profiles. This is developed in Appendix A. Using the approach developed there and the numerical tabulation of Jastrow and Kyle [6], the following analytic forms for the atmospheric profile are obtained:

$$p = p_{0} (H/H_{0})^{-1/\beta} o$$

$$\rho = \rho_{0} [1-2(y-y_{0})/(R+y_{0})]^{-1} (H/H_{0})^{-(1+\beta_{0})/\beta} o,$$

where
$$H = H_0 + \beta_0 (y - y_0)$$
 $(y_i < y \le y_f)$ (15)

Different parameters $(\rho_0, p_0, H_0, \beta_0, y_0)$ are used for different altitude intervals $[Y_i, Y_f]$, and these are given in Table 1. Use of these expressions is found to give a good fit of the numerical tabulation of Jastrow and Kyle [6].

TABLE I

Parameters for Analytic Atmospheric Profiles (See Eq. (15) of text)

		-		_						_
(km) Y_{f} (km) Y_{O} (km) P_{O} (kg./m. ³) P_{O} (km) H_{O} (km)1010104.176-012.614+046.408725104.176-012.614+046.408725254.048-022.534+036.438855555.650-044.283+017.870785859.193-065.020-015.723590859.193-065.020-015.7235100903.842-065.020-015.7235120100903.642-074.005-026.34871401201203.613-084.324-0310.7986	Bo	-2.0378-01	2.0064 . 03 6.9940-02	2.4469-03	-7.2257-02	1.7711-03	6.1446-02	2.6826-01	2.3350-01	
(km) Y_{f} (km) Y_{o} (km) ρ_{o} (kg./m. ³) P_{o} (N/m ²)1010104.176-012.614+0425104.176-012.614+0425254.048-022.534+0355555.650-044.283+0185859.193-065.020-0190859.193-065.020-0190859.193-065.020-01100903.842-062.098-01120100903.642-074.005-021401203.613-084.324-03	H ₀ (km)	6.4087	6.4087 6.4388	7.8707	5.7235	5.7235	5.7323	6.3487	10.7986	
(km) Y_f (km) Y_o (km) ρ_o (kg./m. 3)1010104.176-0125104.176-0125254.048-0255555.650-0485859.193-0690859.193-0690859.193-061009085120100901201203.642-071401203.613-08	Po (N/m ²)	2.614+04	2.614+04 2.534+03	4.283+01	5.020-01	5.020-01	2.098-01	4.005-02	4.324-03	0
Yr Yr Yo Yo Km) 10 10 10 10 25 10 25 25 45 25 55 55 85 85 85 90 85 85 1100 90 100 120 120 100 120 120 120	ρ _o (kg./m. ³)	4.176-01	4.176-01 4.048-02	5.650-04	9.193-06	9.193-06	3.842-06	6.642-07	3.613-08	0
(km) Y _f (km) 10 25 45 55 85 90 100 120 120	Y _o (km)	10	10 25	55	85	85	06	100	120	
(Km)	Y _f (km)	10	25 45	55	85	06	100	120	140	D
Yi 0 10 25 25 25 25 85 90 100 120	y ₁ (km)	0	10 25	45	55	85	06	100	120	140

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Integration of the Equations of Motions

Before going on to the actual integration of the equations of motion, further approximations used should be indicated. These have to do with the magnitude of the angle of attack α during the flight of the rocket.

As Eq. (2) indicates, while the rocket speed is relatively low during first stage motion, a particular value of α has a larger effect on the turning rate $\dot{\psi}$ than when the rocket speed is much greater in second or higher stage motion. Values of α in first stage motion needed to a cause a particular orbit injection value of ψ (e.g., $\psi=0$) have been found not to exceed a few degrees. On the other hand, the maximum value of α in higher stages has been found to be as high as tens of degrees in some cases, depending on payload weight and the required orbit injection altitude. At least while aerodynamic effects are important, however, it appears to be a good approximation to replace sin α by α , and the approximations of Eq's (7) and (8) are valid. Hence to a good approximation, Eq's (1) - (3) become

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$$\dot{\mathbf{v}} = -(\mathbf{K}/\mathbf{r}^2)\sin\psi - (\mathbf{D}/\mathbf{M}) + (\mathbf{T}/\mathbf{M})\cos\alpha \qquad (16)$$

$$\mathbf{v}\psi = -\left[\frac{\mathbf{K}}{\mathbf{r}^2} - \frac{\mathbf{v}^2}{\mathbf{r}}\right]\cos\psi + \frac{1}{\mathbf{M}}\left[\mathbf{T} + \frac{\partial C_L}{\partial \alpha}\frac{\rho \mathbf{v}^2}{2}\mathbf{S}\right]\sin\alpha \qquad (17)$$

These are the equations of motion integrated in the computer program. They are considered valid for all rocket stages. Here D is given by Eq's (7), (10), and (11), and T is given by Eq. (6).

The procedure is to employ a particular profile for $\alpha(t)$ in the integration of Eq's (16) and (17). It has the form

$$\alpha(t) = \begin{array}{c} C (t-t_{o}) \left[1 - \frac{(t-t_{o})}{t_{d}}\right] & (t_{o} < t < t_{o} + t_{d}) \\ 0 & (t \le t_{o} \text{ or } t \ge t_{o} + t_{d}) \end{array}$$
(18)

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It is shown in Fig. 5 as an inverted parabola where it is non-vanishing, with an extremum point at $t = t_0 + t_d/2$ and associated value $\alpha_m = Ct_d/4$. It is a convenient profile to manipulate, and the results for orbit injection conditions have been found to be insensitive to variation of parameters in the profile. For example, a particular orbit injection altitude at $\psi = 0$ can be obtained for various C values by adjusting the t_d value, and the value of injection velocity is insensitive to this procedure (e.g., within 0.2%). The assumption of a particular form in Eq. (8) would be a major concern if this were not the case.

Beside having solutions for v(t) and $\psi(t)$ in Eq's (16) and (17), it would also be of interest to find the altitude and ground range (cf. Fig. 1), given as solutions of

$$\dot{\mathbf{y}} = \mathbf{v} \sin \psi \tag{19}$$

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$$\mathbf{x} = \mathbf{v} \, \mathbf{R} \, \cos \, \psi / \, (\mathbf{R} + \mathbf{y}) \tag{20}$$

The Eq's (16), (17), and (19), and (20) are a coupled set of ordinary first order differential equations in the time variable, and are therefore solvable by the Runge-Kutta method [7]. In explaining this method it will be convenient to rewrite these equations as

 $\dot{\mathbf{x}}^{(\alpha)} = f^{(\alpha)}(t, \{\mathbf{x}^{(\beta)}\})(\alpha, \beta=1, 2, 3, 4), (21)$

where the superscript α in this case runs over the four variables v, ψ , y, and x, which are respectively represented by $x^{(1)}$, $x^{(2)}$, $x^{(3)}$, and $x^{(4)}$. Eq. (21) states that the time derivative of one of these variables is a particular function of the time (e.g., through the given explicit time dependences of $\alpha(t)$ in Eq. (18) and mass of the rocket M(t)) and of the four variables $\{x^{(\beta)}\}$ at that time. Actually, inspection of Eq's (16), (17), (19), and

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(20) reveals no dependence of $f^{(\alpha)}$ on $x^{(4)} = x$ here. In the Runge-Kutta numerical method, the time scale is broken up into mesh points separated by a fixed interval width h. The integration is consecutively forward on the time scale, with the variable at the mesh point n determined from the known values at the mesh point n-1, as follows:

$$x_{n}^{(\alpha)} = x_{n-1}^{(\alpha)} + \frac{1}{6} \left(\Delta x_{1}^{(\alpha)} + 2\Delta x_{2}^{(\alpha)} + 2\Delta x_{3}^{(\alpha)} + \Delta x_{4}^{(\alpha)} \right) \quad (22)$$

where

$$\Delta x_{1}^{(\alpha)} = hf^{(\alpha)} (t_{n-1}, \{x_{n-1}^{(\beta)}\})$$

$$\Delta x_{2}^{(\alpha)} = hf^{(\alpha)} (t_{n-1} + h/2, \{x_{n-1}^{(\beta)} + \Delta x_{1}^{(\beta)}/2\})$$

$$\Delta x_{3}^{(\alpha)} = hf^{(\alpha)} (t_{n-1} + h/2, \{x_{n-1}^{(\beta)} + \Delta x_{2}^{(\beta)}/2\})$$

$$\Delta x_{4}^{(\alpha)} = hf^{(\alpha)} (t_{n-1}^{+} + h/2, \{x_{n-1}^{(\beta)} + \Delta x_{3}^{(\beta)}\}$$
(23)

The procedure, according to these equations, is to first evaluate the four values of $\Delta x_1^{(\alpha)}$, which then enables the evaluation of the four values of $\Delta x_2^{(\alpha)}$. Next comes the evaluation of $\Delta x_3^{(\alpha)}$, and then $\Delta x_4^{(\alpha)}$. Finally, $x_n^{(\alpha)}$ is evaluated from Eq. (22). This is also the procedure for any number of variables (α , β , = 1,2, ----), not just the four of interest here, and the method is easily programmable.

The computer program for the rocket trajectory consists of a main program and ten subroutines. The general approach of the coordinating main program RALLY 6 is to take a trial $\alpha(t)$ function from Eq. (18), and then perform the numerical integration of Eq's (16), (17), (19), and (20) by the Runge-Kutta method outlined in Eq's (21)-(23). Based on the final rocket burnout values of $\psi(t)$ and y(t)- i.e., the orbit injection values of these variables -

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the tilt parameters in Eq. (18) are varied, and the process is repeated iteratively until the desired orbit injection conditions are obtained. Instead of varying tilt parameters the program also provides for the variation of coast times after powered stage burnouts, in order to obtain orbit injection conditions. The program is interactive with the user in that corrections are made by him during the run, but at some point the user can choose an option which makes the program self-correcting, i.e., the program iterates on the corrections by a Newton-Raphson method [7] until convergence to a specified (by the user) stage burnout heading is obtained within a specified tolerance. The input and output options are very flexible, so that the program will handle a wide variety of rocket problems with ease. The Subroutines are: (1) MASS which evaluates the mass of the rocket in Eq's (16) and (17) from $M=M_{O}$ -Mt, where the initial stage mass M_O and mass ejection rate \dot{M} are known parameters for the rocket stage, and t is the burn duration in that stage; (2) TILT which evaluates angle of attack from Eq. (18); (3) DEN which evaluates density and pressure according to Eq. (15) and Table 1 and sound speed according to Eq. (14); (4) PSINCR which evaluates ψ in Eq. (17), which is needed for the Runge-Kutta integration in the main program; (5) VINCR which evalues v in Eq. (16); (6) ALT which evaluates y in Eq. (19); (7) GRRG which evaluates x in Eq. (20); (8) CNVG which computes the parameter corrections to tilt amplitude (C in Eq. (18)) or stage coast time in the self-correcting mode of operation of the program, and then tests for convergence; (9) STORE which either stores burnout values just computed for the present rocket stage to be used in calculations for the next stage, or recalls values stored from calculations on the preceding stage to be used in a repetition of the calculations for the present stage; and (10) FORCES which com-

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putes thrust and aerodynamic forces in the equations of motion according to Eq's (5)-(12).

Detailed Description of the Computer Program and Its Use

In Appendix B is the listing of the program for the DEC-10 system, and in Appendix C is the almost identical listing of the program for the PDP-11/70 computer system. One of the few changes made arose from the way the two different computers evaluated $\cos \psi$ in Eq's (17) and (20). There was no problem with the DEC-10, but the PDP-11/70 made small errors in the evaluation of $\cos \psi$ near $\psi=\pi/2$ rad. As small as these errors were, they were compounded in the integration procedure, so that final errors in the burnout ground range and altitude amounted to a few km. or more. This was remedied in the PDP-11/70 program version by replacing $\cos \psi$, wherever it occurred, by sin $(\pi/2-\psi)$. With respect to the DEC-10 listing it is seen that statements #12400 and 14900 were changed in the

PDP-11/70 version in this fashion. There is also a difference in the treatment of the vacuum-to-sea level thrust ratio which is given a value VSLR(I) for each rocket stage in the DEC-10 version (cf. statements #120, 2010, and 4255 in Appendix B). This ratio is given as x in Eq's (5) and (6). Very often the effect of VSLR or x is negligible, except for the first stage, because rocket altitude in the upper stages is high enough that the atmosphere is essentially a vacuum, i.e., T takes on the value T_{vac} in Eq. (6). This is why VSLR is treated as a single parameter for the first stage in the PDP-11/70 version in Appendix C. The only other change was necessitated by the difference in printing output between the two systems. While the output is typed at the DEC-10 terminal on a paper roll as TYPE statements are encountered in the program, it is printed on a CR-tube display in the

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PDP-11/70 system (which periodically erases itself) and subsequently printed on paper at a line printer terminal upon successful completion of the job. In the PDP-11/70 version of the program the IPR printing option variable in the DEC-10 version is replaced by ITY, and IPR becomes a printing option variable for the line printer. Otherwise, the programs are identical. Henceforth, the discussion will be confined to the DEC-10 version.

At this point the statements of the program listed in Appendix B will be itemized, and to help the reader follow the logic of the program, a flow chart of the main program (RALLY 6) is included in Fig. 6(a), (b), and (c). The first two DATA statements in Appendix B (#300 and 400) list the yf and y values in Table 1, while the next one lists the values of ρ_0 in g./cm.³. The next data statement (#700) lists the values of $-2(R+y_0)^{-1}$ in km⁻¹ (cf. Eq. (15)). Next come the values of H_o (in #1000), β_o (in #2000), and P_o in dynes/cm² (in #1350) from Table 1. This is followed by values for constants used in the program; sequentially, these are K in $km^{3/sec}^{2}$ (cf Eq's (1) and (2)), the radius of the earth R in km, pi, and the pressure at sea level (p(o) in Eq. (6)) in dynes/cm². After zeroing a few execution option parameters, the user reads in the number of rocket stages NS up to 6 (which can be changed to a higher number, if necessary, by changing statement #120) and SMZ (I), the initial mass for each rocket stage (in lbs.), which is defined as the mass of the rocket at the beginning of the stage. Periods of coasting can be included as separate stages. Next SMD (I) is read-in which is the rate of mass loss of each stage due to fuel consumption (in lbs/sec). This would be set to zero for a coasting stage. Then the vacuum thrust for each stage TH(I) is read (in lbs), which is the T of Eq. (6). Following this, the program calls for the vacuum-to-sea level thrust ratio VSLR for

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each stage (i.e., x in Eq. (5)), the duration time (in sec) for each stage TBD(I), the stage diameter DI(I) (in meters), and the number of time intervals NI (I) that each rocket stage is to be broken up into in the specification of the time mesh for the numerical integration. In statements #2605 - 2630, TBS(I) is computed, which is the time measured from launch in sec. at which the I'th rocket stage commences. Next the altitude YI, time TI, speed VI, heading PSI, altitude interval ID (associated with Table 1) and angle of attack AL are given values appropriate to launch conditions, and the program is then ready to enter the numerical integration procedure.

After testing a pair of parameters, which will be discussed later, the program calls for the specification of the angle of attack parameters in statements #3460 - 3550. Here CA TLZ, and TLD refer to C (in rad.), t_0 (in sec), and t_d (in sec) in Eq. (18), and thus refer to the tilt amplitude, the time measured from launch at which tilt begins, and the duration of non-vanishing tilt, respectively. Now the user supplies the parameters which determine the various execution options in statement #3575:

- NT specifies the number of rocket stages to be integrated as in Eq's (21) - (23).
- (2) IOP specifies what part of the program that control is returned to after the NT rocket stages have been integrated. IOP = 1 terminates the program, IOP=2 returns control to program statement 3 (PS3) where variables are given their launch values. IOP=3 returns control to PS69 which redefines the time mesh for the numerical integration. IOP = 4 returns control to PS25 which initiates the integration of rocket stages other than the first. IOP=5 returns control to PS2 which defines a new rocket problem.
- (3) ICE specifies the number of rocket stages which

have previously been integrated.

- (4) IPR specifies the extent of the printed output from the numerical integration of the NT rocket stages. IPR = 0 causes printout of the variables at all the mesh points. IPR \geq 1 causes printout of only the final or burnout values of the variables. IPR = 0,1 signals the program that corrections to tilt or coast time parameters will be supplied by the user. IPR=2 signals that tilt amplitude or coast time corrections will be computed automatically by the program in its self-correcting mode until convergence has been obtained, i.e., when the burnout value of the heading (ψ) takes on a specified value within a specified tolerance.
- (5) IBG specifies what the value of IOP will be after convergence has been obtained in the self-correcting mode of the program.
- (6) JCV specifies what is to be corrected in the selfcorrecting mode of the program. JCV = 2 specifies that TBD(NS) is to be adjusted, i.e., the coast time of stage NS. JCV * 2 specifies that CA is to be adjusted, i.e., the amplitude parameter for the angle of attack profile of a particular rocket stage.

- (7) IMA specifies whether or not the payload weight is to be adjusted. This control is convenient for obtaining the dependence of orbit injection conditions on payload weight.IMA = 1 specifies that a payload weight increment is to be read-in after the program has obtained covergence in its self-correcting mode of operation.
- (8) JNS specifies whether or not a new NS is to be read in for the integration procedure. JNS = 0 says not, but JNS ≠ 0 specifies that a new NS will be read in. A temporary NS value is used for the self-correcting mode of operation on parameters of an intermediate rocket stage when the NT stages being integrated do not include the final rocket stage.

If IPR = 2 the program enters its self-correcting mode of operation in the integration of NT rocket stages. The first part of this is the specification of the desired final result PSSS for the heading in the integration and the convergence within EPS of the parameter that is being corrected. These values are read in statements #3882 and 3885. The program then carries out the numerical integration of the Eq's of motion (16), (17), (19), and (20) as indicated in Eq's (21) - (23), and with the help of the subroutines discussed at the end of the preceding section. The reader will have little difficulty in following this. While some of the parameters were read in CGS or English units, everything is converted to MKS units in the numerical integration with the exception that distance units are km. in the final results.

Beginning with statement #8010 the program enters a series of tests which are designed to correctly set the parameter ICD for the subroutine STORE, which is called in statement #8050. If ICD = 1, STORE will store the trajectory values just computed, to be used as initial values for integration of the next stage. If ICD = 2, STORE will recall the initial values for the integration just computed, so that this integration can be repeated to get a heading which is closer to that desired. The rest of the main program, which is indicated by the flow chart, is mostly concerned with obtaining convergence in the self-correcting mode of operation. Convergence is on the tilt amplitude CA or on the coast time of the stage NS (i.e., TBD(NS)), as

determined by JCV. As previously stated, the sub-routine CNVG computes the corrections to these parameters in the self-correcting mode (see statement #8655). The parameter ICVG signals whether or not convergence has been obtained. ICVG = 0 means it has not; ICVG = 1,2 means it has.

Perhaps the best way to understand how the program works and how to use it is to trace through the flow chart in Fig. 6 with data on sample problems. For this purpose, consideration is directed to a two-stage rocket with four stages of motion. The author would like to make the distinction between rocket stages and stages of motion. The program is solely concerned with the latter usage, and the subsequent discussion here is too. The first stage of motion is powered, the second stage is defined as a

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coasting stage, the third stage is powered, and the fourth stage is a coasting stage for the payload. It is assumed that all the rocket datahave been read in, so that the point in the program has been reached where the angle of attack dataare to be read-in, i.e., at statement #3500 in the program. The solutions of four problems for this rocket are of interest, and these are given below.

Problem 1: The right amount of tilt must be put in the first stage so that the rocket will be injected into orbit at zero heading (i.e., $\psi = 0$ as the final solution of the equations of motion). Hence, the second through fourth stages fly in a gravity-turn ($\alpha = 0$ in Fig. 1) trajectory. The user must supply the data shown in Table 2, which also shows the statement number that demands the data, along with occasional comments. The reader should trace the path of logic through the flow chart which is dictated by the data. The initial reading of the tilt or angle of attack data will result in a heading which is non-zero in the printed output. This suggests a new CA value which is read-in (see Table 2), and the change of IPR from 1 to 2 in the next data input will cause the program to go into its selfcorrecting mode of operation. From then on the program iterates on CA until corrections to CA become less than 0.000002. This will typically take only a few seconds of actual runtime on the DEC-10 system. Then the complete launch trajectory is printed out at every time mesh point as the converged solution. Note that the user supplies two trial values of CA before the correction procedure is made automatic. Two such values are required in the calculation of the correction in Subroutine CNVG. Program control then returns to PS3, where now ICVG = 0 and IPR =0, and variables are initialized at their launch values. The program is now ready to solve the second problem.

TABLE II

Data Required for Problem 1 in Text

DATA SUPPLIED	STATEMENT NO. (#)	COMMENTS	
CA, TLZ, TLD	3500, 3505, 3510	CA<0	
42012	3626	JCV=IMA-JNS=0	
CA, TLZ, TLD	3500, 3505, 3510	A new CA	
42022	3626		
0.	3882	Final ψ	
0.000002	3885	Sample toler-	

TABLE III

Data Required for Problem 2 in the Text

DATA SUPPLIED	STATEMENT NO. (#)	COMMENTS
CA, TLZ, TLD	3500, 3505, 3510	First stage values
24002	3626	
CA, TLZ, TLD	3500, 3505, 3510	Third stage values
24212	3626	
CA, TLZ, TLD	3500, 3505, 3510	A new CA - third stage
24222	3626	
0.	3882	
0.000002	3885	

Problem 2: With a specified tilt in the first stage, the tilt in the third stage must be adjusted so that the final orbit injection heading is zero. Table 3 specifies the data for this problem. The first stage tilt value is read-in, the first two stages are integrated, with the trajectory values being printed out, and the final results for the variables are stored for use as initial values in the integration of subsequent stages. Then trial tilt parameters for the third stage are read-in, and the injection conditions are calculated and printed out as a result of the integration of the third and fourth stages. Based on this, a new CA for the third stage is entered by the user along with the other tilt parameters. Then the program is put into its self-correcting mode, iteration on CA for the third stage ensues, and finally the converged launch trajectory is printed out. Control is then returned to PS3, where the variables are initialized at their launch values. The program is now ready to solve the third problem.

<u>Problem 3</u>: With specified tilt functions in the first and third rocket stages and a specified coast time for the second stage (TBD(2)), it is desired to adjust the coast time of the fourth stage to obtain $\psi = 0$ as the orbit injection heading. The data supplied by the user is shown in Table 4. By putting JCV = 2, the correction procedure is now on TBD(4). It proceeds analagously to that on CA. After converging to the correct coast time for the fourth stage, program control reverts back to PS3, where variables are initialized at their launch values again, and the fourth problem is now ready to be solved.

<u>Problem 4:</u> With specified tilt functions in the first and third rocket stages, it is desired to adjust the coast

TABLE IV

Data Required for Problem 3 in Text

DATA SUPPLIED	STATEMENT NO. (#)	COMMENTS
CA, TLZ, TLD	3500, 3505, 3510	First stage
24012	3626	
CA, TLZ, TLD	3500, 3505, 3510	Third stage
14212	3626	
0.	3500	These tilt
0.	3505	parameters play
0.	3510	no role (no thrust)
143122	3626	
TLS	8687	New coast time fourth stage
0., 0., 0.	3500, 3505, 3510	
143222	3626	
0.	3882	
.000002	3885	

TABLE V

Data Supplied by User for Problem 4 in Text

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DATA SUPPLIED	STATEMENT NO. (#)	COMMENTS
CA, TLZ, TLD	3500, 3505, 3510	First stage
14001	3626	
0., 0., 0.	3500, 3505, 3510	Second stage
14114201	3626	Coast time iteration
2	3845	Temporary NS
TLS	8687	New coast time
0.,0.,0.	3500, 3505, 3510	Irrelevant tilt again
141242	3626	Self-correct coast time
PSSS	3882	42
.000002	3885	Tolerance
Ca, TLZ, TLD	3500, 3505, 3510	Third stage
14204001	3626	
4	3845	Restore ori- ginal NS
0., 0., 0.	3500, 3505, 3510	Fourth stage
143122	3626	
TLS	8687	New coast time
0., 0., 0.	3500, 3505, 3510	Irrelevant tilt again
143222	3626	
0.	3882	Orbit injec- tion heading
.000002	3885	

time of the second stage to obtain a particular heading ψ_2 at the end of the second stage, and then to adjust the coast time of the fourth stage to obtain an orbit injection heading $\psi=0$. The data supplied by the user is shown in Table 5. The rocket launch trajectory is printed out for all four stages in the course of the program.

As demonstrated by the examples discussed, the program is flexible enough to handle a wide variety of rocket problems. Whatever rocket problem the user is dealing with, it seems advisable to check the data to be supplied to the program by tracing its path through the flow chart.

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Example: A Hypothetical 2-stage Launch Vehicle

For typical input data the numerical output from a specific example is obtained next, and its use in computations is demonstrated. As in the preceding section, the rocket in this example has four stages of motion, the first and third being powered and the second and fourth being coasting stages. One might typically be given the rocket data of Table 6. It is a simple matter to calculate the input parameters for the program from Table 6 with the help of the following equations:

 $SMZ = M_{o}$ $SMD = T_{vac}/I_{sp}$ $TH = T_{vac}$ VSLR = x $(TBD)(SMD) = M_{p}$

The left-hand sides of these equations are written in the notation of the computer program (Appendix B), while the right-hand sides are written in the notation of Table 6. In Appendix D the solutions of <u>Problem 1</u> and <u>Problem 2</u> of the preceding section are obtained on a remote time-sharing DEC-10 terminal for the hypothetical rocket of

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Rocket Data for a Hypothetical 2-Stage Launch Vehicle

PARAMETER	STAGE	FIRST	SECOND
Ignition Wt.≡Mo	(1bs)	400000	120000
Stage Wt. =Ms	(lbs)	274000	110000
Total Propellant Wt. ≡ M	(lbs)	264000	105000
Vac. Spec. Impulse	(sec)	302.3	314.3
Vac, Thrust = T vac	(lbs)	665000	220000
Vac.Sea level Thrust Ratio = x		1.15	
Stage diameter	(ft)	10	10

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Table 6. The user enters data for the rocket problem when the program prompts him to do so with an ENTER---- message. The value 10 seconds is arbitrarily used for the coast times of the second and fourth stages of the rocket motion. This might, for example, allow enough time between burnout of one powered stage and ignition of the next powered stage to take care of jettisoning launch vehicle vestiges. In this particular exercise the user is trying to obtain the orbit injection velocity at zero heading (in this case, the perigee of the orbit) and at an altitude of 100 nmi. = 185.2 km. In order to find this, the user plots the three output values altitude vs. velocity, as in Fig. 7 (x - marks), and then interpolates (to o-mark) to find an injection velocity $v' \approx 7.833$ Km/sec. Incidentally, the total computer runtime for the output in Appendix D was 26 seconds.

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The value just found for the orbit injection velocity at 100 nmi. altitude is in a stationary earth approximation, as discussed in the Introduction. The earth's rotation is easy to account for [1] in separate computations. To do this, one notes first that the total ground range covered in this example is about 670km., which is small enough that a flat earth approximation[1] is in order. Then the actual orbit injection conditions for the rocket, including earth's rotation, are given by

$$v = [v'^{2} + 2v'v_{o}\sin a_{o} + v_{o}^{2}]^{\frac{1}{2}}$$

$$\psi = 0^{0}$$

$$\tan a = \frac{v'\sin a_{o} + v_{o}}{v'\cos a_{o}}$$
(25)

where

 $v_o \equiv \omega_E^R [1 + (y_1/R] \cos L_o]$

Here, v' is the injection velocity in the stationary earth

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approximation, v_0 is the effective initial velocity in an easterly direction imparted by earth's rotation, a is the launch azimuth angle as seen by an earth observer, which is established shortly after launch by suitable pitch and yaw rotations, and a is the actual azimuthal angle at orbit injection as seen by an inertial observer. These azimuthal angles are measured clockwise from the local north direction to the direction of motion projected on the local horizontal plane. In the expression for v_0 , w_E is the angular rotation rate of the earth, R is its radius, y1 is the rocket altitude after first stage burnout, and L is the latitude of the launch site. The orbit inclination angle i, which is the angle made by the normal to the orbital plane (determined from the right-hand rule by curling one's fingers in the direction of motion of the satellite) with the earth's rotation axis, is determined from the preceding parameters by

(26)

where

 $o \le a \le 2\pi$ and $|L_0| \le i \le \pi - |L_0|$

 $\cos i = \sin a \cos L_0$

in radian measure. By way of illustration it is supposed that the rocket in the example is launched from Cape Kennedy, for which $L_0 = 28.5^{\circ}$. One has that v' = 7.833Km/sec, $w_E =$ 7.292116 rad/sec., R = 6378.145km., $y_1 \approx 64$ km., and a_0 is variable in the example. Fig. 8 is a plot of a and i vs. a_0 and in Fig. 9 is a plot of v vs.i. This is as much information as is needed for the launch problem example. It will be seen from these figures and Fig. 3 that, as expected, the apogee altitude will be maximum when the launch direction is due east $(a_0 = 90^{\circ})$; then it is about 1075nmi. On the other hand, for launch directions in the range $190^{\circ} \le a_0 \le$ 350° there is a question of the stability of the parking orbit, since in this range orbit altitudes dip below 100nmi.

DISCUSSION

The program developed here was employed in calculations on large launch vehicles. Curves of the type of Fig. 7 were found for the injection conditions in the stationary earth approximation. Earth's rotation effects were subsequently included in the manner described previously. The results were compared to corresponding calculations which leave out explicit consideration of thrust increase and drag effects, but rather assume that they cancel in an approximation delineated by Ruppe [4]. A similar approximation was used by Ehricke [3]. The curves based on this approximation yielded injection speeds at a given altitude 2-3 % less than those calculated with the present program. For very large rockets the thrust increase effect outweighs the drag effect by about this amount. Because of the sensitive dependence of apogee altitude on perigee injection speed for eccentric orbits, it was concluded that explicit inclusion of thrust increase and drag effects was necessary. It was also possible to calculate results with the present program for a large launch vehicle which had been independently calculated elsewhere. These were results of the nature of Fig. 9, and the agreement between the two sets of calculations was within 0.05%. While this kind of agreement is probably fortuitous, it is also somewhat encouraging.

Realistically, one would expect the present program to yield information about large, unfinned rocket systems, consistent with the data base to within 1% for the injection speed. The neglect of angle of thrust (β in Fig. 2) is supposed to be a good approximation for large rockets [3, 4], and a particular form for the angle of attack profile is apparently unimportant for these systems. A variation of parameters in the profile used here produced $\stackrel{<}{\sim}$ 0.2% errors. Beyond this, one can argue that small errors

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in the aerodynamic coefficients or atmospheric model are not important, because the total contribution of aerodynamic effects is relatively small for large rockets [4]. Another consideration is that significant improvements on the present program are probably not possible for many rocket problems. Data for such a problem are often comparable with that of Table 6. Further refinements on the approximations would require more data than is given. One would need, for example, drag and lift coefficients specific to the rocket system, details of the angle of attack or thrust angle profiles, the time history of thrust and mass variations, and data on the atmospheric conditions. Furthermore, uncertainties in rocket parameters in the data base can and often do amount to a few percent, and this is therefore an inherent limitation to accuracy.

The present program is certainly flexible enough to handle a large variety of rocket problems. From the preceding discussion it appears to be entirely satisfactory for most launch system analyses, particularly those which involve large, unfinned rockets.

ACKNOWLEDGEMENT

The author would like to thank Mr. Frank D. Clarke who entered the program and adapted it to the DEC-10 and PDP-11/70 systems. In addition to his help in the debugging and in numerous revisions of the program, he also assisted in obtaining results for some rocket problems. The author also appreciates assistance from Dr. John N. Hayes on the manuscript.

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1.2 - VECTOR FROM CENTER OF GRAVITY TO CENTER OF PRESSURE

- ß - ANGLE OF THRUST D
- DRAG

Fig. 2 - Force diagram for the rocket with symbol legend



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Fig. 6(a) - Flow chart for RALLY 6 (Dec-10 version)

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Fig. 7 - Altitude vs. speed at orbit injection for the example rocket (see text) in stationary earth model

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Fig. 8 - Injection azimuth angle and orbit inclination angle dependence on launch azimuthal angle for the example rocket (see text)

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APPENDIX A

Analytic Atmospheric Density and Pressure Profiles

In hydrostatic equilibrium we have

$$dp/dy = -\rho g$$
, (A1)

where p is the pressure, y is the altitude, ρ is the density, and g is the acceleration due to gravity. The ideal gas law, Eq. (13) of the text, is reproduced here as

$$p = \rho \hat{R} T / M , \qquad (A2)$$

from which Eq. (Al) is rewritten as

$$dp/p = -dy/H$$
(A3)

where

$$H = RT/Mg$$
(A4)

is the "scale height". From Eq's (A2) and (A3), we get

$$d\rho/\rho = dp/p - dT/T + dM/M = dp/p - dH/H - dq/q$$
 (A5)

There are several numerical tabulations of standard atmosphere density and pressure profiles, and if we choose one of them [6], then for an altitude y near a particular tabular value y_0 , we have from Eq. (A4)

$$H = H_{0} + \beta_{0} (y-y_{0})$$
(A6)

where H and β are constants evaluated at y, and Eq. (A6) is simply a first order expansion of H about y₀. Direct evaluation yields

$$\beta_{0} = \left\{ \frac{dH}{dy} \right\}_{Y=Y_{0}} = \left\{ \frac{A}{RT} \right\}_{Mg} \left[\frac{1}{T} \frac{dT}{dZ} - \frac{1}{M} \frac{dM}{dZ} + \frac{2}{R+Y} \right] \left\{ \begin{array}{c} (A7) \\ Y=Y_{0} \end{array} \right\}_{Y=Y_{0}} \right\}$$

where R is the radius of the earth, and we have used the known variation of g with altitude. The case of constant gradient of scale height is a well known simplification in the integration of Eq's (A4) and (A5), which become

$$\frac{d\mathbf{p}}{\mathbf{p}} = \frac{-1}{\beta_0} \frac{d\mathbf{H}}{\mathbf{H}} \Longrightarrow \frac{\mathbf{p}}{\mathbf{p}_0} = \left(\frac{\mathbf{H}}{\mathbf{H}_0}\right)^{-1/\beta_0}$$
(A8)

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$$\frac{d\rho}{\rho} + \frac{dg}{g} = \frac{d(\rho g)}{\rho g} = -\left(\frac{\beta_0 + 1}{\beta_0}\right) \frac{dH}{H}, \quad \frac{\rho g}{\rho_0 g_0} = \left(\frac{H}{H_0}\right) \frac{-(1 + \beta_0)/\beta_0}{(A9)}$$

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These equations are thus interpolation formulae, which when used in conjunction with some of the values in the numerical tabulation [6], become the desired analytic expression of atmospheric density and pressure (see text).

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APPENDIX B

DEC-10 Version of RALLY 6

. TYPE S	RALLYS.FO	R
00100		COMMON/91/5C, R, SS, VSR, RHO, PRZ, P, VT
00110		COMMON/B1/YZ (9) + YD (9) + RHZ (9) + ALZ (9) + HZ (9) + BT (9) + PR (9)
00120		COMMON/C1/SMZ (5) ; SMD (5) ; TBS (5) ; TBD (5) ; DI (5) ; TH (5) ; NI (5)
00300		DATA (YD (1) + I=1,9)/10. +25. +45. +55. +35. +90. +100. +120. +140.
1		
00400		DATA (YZ (1) + 1=1.9) / 10. + 10. + 25. + 55. + 35. + 35. + 90. + 100. + 120. /
00500		DATA (RHZ (1) , 1=1,9) /4. 175E-04, 4. 175E-04, 4. 043E-05, 5. 550E-
07.		
00500		19, 1935-09, 9, 1935-09, 3, 3425-09, 5, 5425-10, 3, 5135-11/
00700		DATA (4, Z(1), 1=1, 4)/-3, 13085-04, -3, 13085-04, -3, 12355-04,
00900		1-3.10395-043.09455-043.09455-043.09215-043.0323
5-04.		
00900		3-3 07795-04/
01000		D9T9 (47 (1) . Ist. 9) /5. 4037.5. 4097.5. 4399.7. 9707.5. 7935.5.7
235.		
01100		15, 7323.4 3497.10, 7994/
01200		NATA (BT (1. 1 = 1. 0) /= 2 02705-01.2 00445-02.4 9945-02.2 444
95-02.		
01200		1-7 00575-00.1 77115-00.4 14455-00.0 20055-01.0 00505-01
01300		1-1.223/2-02/1.1/112-03/3.14452-02/2.55252-01/2.353502-01
01250		NOTO (00(1), 1+1.9) /2 \$14505, 2 \$14505, 2 \$24504, 4 202502
01330		15 03.5 03.2 030.4 0055-01.4 2345-037
01375		20+0 202010EAS
01400		
01500		X=53(5,140 01-0 (4)500(54
01550		
01500		78241.013255205
01525		100-0
01537		THE 130
01550		
015/5	120	PERMATIC ENTER INST
01700		READ (5,100) AS
01725		1142=142
01/50		TYPE 130.45
01300		8E4D (5,200) (5M2(1),1=1,MS)
01350		TYPE 140.115
01375	130	FORMAT (ENTER , 11, VALUES FOR SM2)
01900		READ (5,200) (SMU(1),1=1,45)
01950		TYPE 150.05
01975	140	FORMAT C ENTER STILL VALUES FOR SMD S
03000		READ (5,200) (14(1),1=1,45)
02005		TYPE 152
02005	152	FORMAT (ENTER VSLR)
02010		READ (5,200) VSLR
02050		TYPE 150+MS
02075	150	FORMAT (" ENTER ", II, " VALUES FOR VAC. TH")
02100		READ (5,200) (TBD(1),I=1,HS)
02150		TYPE 170, HS
02175	150	FORMAT (ENTER , II, VALUES FOR TBD/)
02200		READ (5,200) (DI(I),I=1,HS)
02250	59	TYPE 130+MS
02275	170	FORMAT (ENTER (11) VALUES FOR DI)
02300		READ (5,300) (NI(I),I=1,NS)
02375	130	FORMAT (SHTER (11) YALVES FOR HIT)
02400	100	FORMAT (11)
02500	200	FORMAT (F10.0)
02500	300	FORMAT (15)
02505		X=0.
02510		TB3(1)=0.
02515		IF (HS.E0.1) 50 TO 3
02520		00 5 I=2,45
02625		X=X+TBD(I-1)

02550	-	103117-0	
02700	3	YI=0.	
02750		TI=0.	
02775		VI=0.	
02737		XI=0.	
02900		PSI=PI/2.	
03200		ID=1	
03300		9 12VS (1-9 12VS 0	
03425		9 = 1).	
03450	25	TE (TOVE ED 1. TOVE = P	
02452			
03452		15 (109 GT 1, GD TO 210	
03455		TYPE 100	
03450			
03475	190	FURMAT (C ENTER VALUES FUR UNITEZITED)	
03500		REHD (5,200)CH	
03505		READ (5,200) TLZ	
03510		READ (5,200) TLD	
03550		TYPE 225	
03575	225	FORMAT (ENTER VALUES FOR HT, IOP, ICE, IPR, IBG, JC	V. IMA.
03575		1.0957)	
03525		READ (5,500) HT, IOP, ICE, IPR, IBS, JCV, IMA, JHS	
03300	500	FORMAT (SII)	
03315		IF (UHS.EQ. 0) 50 TO 303	
03930		TYPE 120	
03345		854D (5,100) NS	
03350	700	FORMAT (INT. TIME (SEC.) SPEED (K/S) HEADING (R)	9LT. (KM
03370		1.7 SR. 8NS (KM) ()	
03373	303	15 (199.1 T. 2) 50 TO 210	
03375		TYPE 195	
03379	195	CODMAT (" ENTED WALLES END DESS. COS",	
03392		0CAR (5.000) 0222	
03335		9540 (5.200) 500	
033399		2220=1200 VII	
03394	210	CONTINUE	
03900		DO 10 TIRLANT	
04000		THICATI	
04050			
04100			
04200			
04250		UT-TH/1.44 4400	
04250		S-DIADI/I. ADI/I	
04250		5=910D1(1)0D1(1)/4.	
04300			
04505		HH=5H2(1)/2.20452	
04352			
04352		15 (198.51.1) 50 10 412	3
04357		TYPE 700	9
04372	412	CONTINUE	9
04400		DD 20 J=1+!!!	2
04402		III=J	ö
04404		IF (VI) 22,21,22	2
04405	21	T=VT/VSLR	A
04403		FD=0.	E A
04410		₽J1=0.	30
04412		YL1=0.	34
04414		X:41=0.	äR
04415		CALL VINCE (PSI, T.AM. YI, FD, AL, VK)	HH
04413		VK1=H+VK	I'S I
04420		50 TO 23	e 3
04422	22	CALL DEN(ID, YI)	L'S
04424		CALL FORCES(VI,AL,S,T,FD,FL)	40
04425		CALL VINCE (PSI, T, AM, YI, FD, AL, VK)	0
04440		CALL PSINCR (VI, PSI, T, AM, YI, FL, AL, PSS)	4.8
			SIN
			E 02
			48

04445		CALL ALT (VI.951.YY)
04447		CALL 5885 (VI.921.VI.X)
04443		VK1=HeVK
04450		2290H=11.9
04455		YL1=H+YY
4477		X!11=4+X
04500	23	TB=TI+H/2.
04500		V1=V1+VK1/2.
4700		C 11 0+120=120
04300		Y1=Y1+YL1/2.
04900		CALL DENCID.YL
05000		CALL TILT (CA.TLZ.T.D.TR.AL)
05100		CALL MASS (TB. I. AM)
5102		CALL FORCES (VI.AL.S.T.FD.FL.
05200		CALL VINCE (PSI.T. AM. YI. FD. AL. VK)
05300		CALL PSINCR (VI. PSI. T. 84. VI. F 8 . PSS)
05400		CALL ALT (VI. PSI. YY)
05450		CALL 5885 (V1.951.V1.X)
05500		VK2=H+VK
05600		2290H=SL9
05700		YL2=H+YY
05750		X12=H•X
05300		V2=V1+VK2/2.
05900		PS2=PS1+PJ2/2.
05000		Y2=Y1+YL2/2.
66160		CALL DEN(ID.Y2)
5102		CALL FORCES (V2. AL. S. T. FD. FL)
05200		CALL VINCE (PS2, T. AM. Y2, FD. AL. VK)
06300		CALL PSINCR (V2, PS2, T. AM. Y2. FL. AL. PSS)
5400		CALL ALT (V2, PS2, YY)
05450		CALL GRRG (V2. PS2. Y2. X)
05500		VK3=H+VK
05500		2290H=ELG
06700		YL3=H+YY
05750		XM3=H•X
06890		V3=VI+VK3
05900		PS3=PS1+PJ3
07000		Y3=YI+YL3
07100		TB=TI+4
07200		CALL DEN(ID, Y3)
07250		CALL TILT (CA.TLZ,TLD,TB,AL)
07275		CALL MASS(TB.I.AM)
07277		CALL FORCES(V3+AL,S,T,FD,FL)
07300		CALL VINCR (PS3, T, AM, Y3, FD, AL, VK)
07400		CALL PSINCR (V3. PS3. T. AM. Y3. FL. AL. PSS)
07500		CALL ALT (V3.PS3.YY)
07502		CALL GRRG (V3, PS3, Y3, X)
07505		VK4=H+VK
07510		PJ4=H+PSS
07515		YL4=H+YY
07557		X!14=H•X
07500		VF=VI+(VK1+2. + (VK2+VK3)+VK4)/5.
07700		PSF=PSI+(PJI+2. +(PJ2+PJ3)+PJ4)/5.
07300		YF=YI+(YL1+2. +(YL2+YL3)+YL4)/5.
7325		XF=XI+(XM1+2. +(XM2+XM3)+XM4)/5.
7350		IF (IPR.ST.0) 50 TO 501
7900		TYPE 500. J. TB. VF. PSF. YF. XF
7950	500	FORMAT (1X, 15, 5(1X, 510, 4))
17935	501	II=IB
7950		AI=A-
7955		21 40 2 F
1.970		Υ Ι= Υ;•

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07935		XI=XF
03000	20	CONTINUE
03010		IF (109.85.4) 50 TO 33
05050		IF (11.95.9T) 50 TO 10
08030		ICD=1
09040		IF (1.59.MS) ICD=2
03045		15 (ICV5.59.2) ICD=1
09050		CALL STORE (VI, PSI, YI, XI, TI, VVI, PSSI, YYI, XXI, TTI, ID, IID, I
CD>		
09050		50 TO 32
06180	33	IF (1.55.HS) 60 TO 29
03500	10	CONTINUE
03505	29	ID=1
03505	32	IC=ICV5
03507		ICV5=0
03503		IF (IC.5E.1) 60 TO 31
03510		IF (IPR.50.0/ 35 TO 31
03620		TYPE 500, III, TB, VF, PSF, YF, XF
03525		IF (IPR.EQ.1) 50 TO 31
03530		Aid=b22-b222
03535		15 (JCV.E9.2) 50 10 215
03540		XNECH
03545		50 TO 220
03550	215	VUELER(UZ)
03500	220	
03505		FIDMAT // VM EQ /.FLA 9. / . VE EQ /.ELA 4.
03537	105	TE (INU ED 2) 20 TO 225
000000		17 (JOA'SA'S) 27 17 522
03550		
000175	225	
03575	200	50 TO 917
03477	21	15 (10 50 2) 50 TO 215
09479		IE (JCV.59.2) 50 TO 310
03532		X0=09
03533		50 TO 317
03534	31.0	XO=TBD(LJK)
03535		TYPE 154
03535	154	FORMAT (ENTER TLS)
03537		READ (5+200) TLS
03539	315	IF (JCV.HE.2) 60 TO 316
03590		IF (IJK.EQ.NMS) 50 TO 317
03591		TTT=TLS-TBD(IJK)
03592		HAB=IJK+1
03594		DO 316 I=NAB, MMS
03595	315	TBS(I)=TBS(I)+TTT
03693	317	TBD (IJK) =TLS
03700	315	IF (IC.LT.2) 60 TO 317
03710	315	IOP=IB6
03711		IF (IMA.MS.1) 50 TO 317
03712		TYPE 415
03713	415	PORMAT (* ENTER SMIM')
03714		READ (5,200) SAIN
03715	417	24771-19172 24771-24771-45414
03715	317	372(1)=372(1)=371(1) an in (30.3.20.35.3) inp
00720	20	CONTINUE
03300	50	2ND
09000		SURPOUTINE MASS (TR. L. 94)
091.00		COMMON/CI/SM7(S).SMD(S).TRS(S).TRD(S).DI(S).TH(S).MI(S)
09200		SM=SMZ(I) - SMB(I) + (TB-TBS(I))
09300		9/1=5/1/2.20452

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09325		RETURN		
09500		END		
09500		SUBROUTINE TILT (CA.TLZ, TLD, TB.AL)		
09700		AL=0.		
09300		IF (TB.ST.TLZ) 50 TO 1		
09900		50 TO 5		
09950	1	TT=TLZ+TLD		
09975		15 (TB.ST.TT) 50 TO 5		
10000	3	T=TR-T! 7		
10100		9L=C9+T+(1,-T/T!D)		
10200	5	QCTIIDIN		
10300	-	END		
10400		SUBORITINE DENCID.Y.		
10402		COMMON/01/60.0.92.02.040.007.0.07		
10500		COMMON 21/27 (9) . VD (2) . DH7 (9) . AL 7 (9) . H7 (9)		
10550		IE (ID GT 9) GD TD 10	· BI (9) ·	PR(3)
10500				
10300	•			
10900				
10900	1.0	PHO-0		
10900	10	R=0		
110950		50 TO 7		
11100				
11100	3	19 (10.21.2) 50 10 4		
11110		YY=YD(ID-1)		
11120		15 (7.51.77) 50 10 4		
11130		ID=ID-1		
11140		50 10 3		
11150	4	X=Y-YZ(ID)		
11200		T=BT (ID)		
11300		H=HZ(ID)+T+X		
11302		C=HZ(ID)/H		
11400		A=RHZ(ID)/(1.+ALZ(IB)+X)		
11402		AP=PR(ID)		
11500		B=C++((1.+T)/T)		
11502		BP=C++(1./T)		
11500		RH3=A+B+1000.		
11502		P=9P+BP/10.		
11504		SS=S0RT(1.4+P/RHD)		
11505		\$\$=\$\$/1000.		
11700	7	RETURN		
11300		E!YD		
11900		SUBROUTINE PSINCR (V.PS.T.AM.Y.FL.AL.PSS)		
12000		C04404/91/50.8.35.VSR.RHD.PRZ.P.VT		
12100		X=R+Y		
12200		B=X+X		
12300		PSS=50/(V+B)-V/X		
12400		PSS=-PSS+CDS(PS)		
13020		CD=FL+T		
13030		PSS=PSS+CD+SIN(9L)/(9M+V+1000.)		
13400		RETURN		
13500		END		
13500		SUBROUTINE VINCE (PS. T. AM. Y. FD. AL. VD)		
13700		COMMON/91/50.8.SS.VS8.840.987.9.VT		
13300		X=R+Y		
13900		X=X+X		
13905		1F (9L.E9.0.) 50 TO 5		
13910		TT=T+COS (9L) -FD		
13915		50 TO 7		
13920	5	TT=T-FD		
14000	7	VD=TT/ (94+1000,) - (50/X) + 210/025		
14100		RETURN		
14200		SHD		

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14300		SUBROUTINE ALT (V,PS,YD)
14400		AB#A+21(4(52)
14500		SETURA
14500		
14700		SOBREDITAE GREG (V.F.S.T.X)
14300		Common / A 1/50, x, SS, VSX, XHO, PK2, P, V1
14900		
15000		
15100		
15200		
15200		SIDONITING SNUCTON, VE. VN. COS. TOD. TOVES
15400		305/30/11/2 (3/15/3/2/3/14/14/14/14/2/3/14/2/3/16/3/
15500		
15500		
15700		YD=Y!4
15300		×H=×H+D×
15900		
15000		IF (9.5E.EPS) 50 TO 5
15100		158=0
15101		ICV5=1
15200	5	RETURN
15300		END
15310		SUBROUTINE STORE (VI, PSI, YI, XI, TI, VVI, PSSI, YYI, XXI, TTI, ID
, IID, IC.	D	
15315		D
16320		59 T9 (3,5) ICD
16330	3	VVI=VI
15340		PSSI=PSI
16350		YYI=YI
16360		××I=×I
15355		TTI=TI
15370		IID=ID
15330		50 10 7
15390	5	VI=VVI
15400		PS1=9S51
15410		Y1 = YY1
15420		
15425		
16440	-	
16450		
15550		
15550		COMMON/41/30.00.00.000.000.000.000
15750		T=VT+(1,-VS2+P/997)
15300		IF (8H0.59.0.) 50 TO 5
15350		XH=V/SS
16950		∃=×!1~1.14
17050		IF (XM.ST.0.33333) 50 TO 2
17150		B=-15.5425+A+A
17250		CD=.25+.3773+EXP(B)
17350		50 T 3 S
17450	2	IF (XM.ST.1.23) 50 TO 4
17550		B=ABS (A)
17650		B=B++3.5
17750		CD=.5273-30.902+B 53
17350		
17950	4	CD=. 17913+. 5313/XM
13050	3	15 (XM. ST. 1.5) 50 TO 5
13150		CL=4.5-0.5♦CBS(2.355♦XM)
19250		CDL=1.522-0.293+005(2.39231+XM)
19450	-	
19400	2	01-0.575.78M
		4.5
		20
		-51-
		. 6

 13550
 CDL =-. 2818+3.2727/XM

 13550
 5

 13750
 C=C+S+1000000.

 18350
 FD=CD+CDL+AL

 19350
 FD=FD+C

 19050
 FL=CL+C

 19150
 RETURN

 19250
 END

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APPENDIX C PDP-11/70 Version of RALLY 6

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COMMON/A1/GC, P, SS, VSP, PHO, PRZ, P, VT COMMOM/R1/Y7(9), YD(9), RHZ(9), AL7(9), HZ(9), BT(9), PR(9) COMMON/C1/SM2(61, SMD(6), TRS(6), TRD(6), DI(6), TH(6), NI(6) DATA YD/10.,25.,45.,55.,85.,90.,100.,120,,140./ DATA YZ/10.,10.,25.,55.,85.,85.,90.,100.,120./ DATA RH7/4.176E-04,4.176E-04,4.04RF-05,5.650E-07,9.193E-09, 19.193F-09, 3. R42F-09, 6. 642E-10, 3. 613E-11/ DATA ALZ/-3.1308F-04,-3.1308E-04,-3.1235E-04,-3.1099E-04,-3.0945E-04, 1-3.0945F-04,-3.0921F-04,-3.0873F-04,-3.0778E-04/ DATA HZ/6.4087,6.4087,6.4388,7.8707,5.7235,5.7235,5.7323, 16.3487,10.7986/ DATA HT/-2.037HF-01,2.0064F-03,6.994E-02,2.4469E-03,-7.2257E-02, 11.7711E-03.6.1446E-02.2.6826F-01.2.3350E-01/ DATA PP/7.614F05,7.614F05,2.534F04,4.283E02,5.02,5.02,2.098, 14.005F-01.4. 174F-02/ GC=3.986012F05 R=6378.145 PT=3.141502654 ITY=0 IPR=0 ICVG=0 PR7=1.013255F05 TYPE 110 FORMAT (' OUTPHT TO LINE PRINTER? (1=YES, 0=NO)') 110 PEAD (5,100) TPR TYPE 120 2 FORMAT (FATER NS') 120 PEAD (5,100) *5 FORMAT (11) 100 NNSENS TYPE 130, "S FORMAT (' FNTER ', TI, ' VALUES FOR SMZ') 130 RFAD (5,200) (547(T), T=1,NS) FOPMAT (FIO.0) 200 TYPE 140.45 FORMAT (' ENTER ', II, ' VALUES FOR SMD') 140 READ (5,200) (S"P(T), T=1, NS) TYPE 150, MS FORMAT (' FRITER ', 11, ' VALUES FOR VAC. TH.') 150 READ (5,200) (TH(1), I=1, NS) TYPE 152 FORMAT (' FRITER VSI.P') 152 READ (5,200) VSI.P TYPE 160,45 FORMAT (' FUTFU ',11,' VALUES FOR TRD') 150 PEAD (5,200) (TPD(T), J=1, NS) TYPE 170, NS FORMAT (' ENTER ', T1, ' VALUES FOR DI') 170 181-2 Cart and the state of the RFAD (5,200) (DT(T), T=1, 45) 69 TYPE IND, MS 190 FORMAT (' ENTER ', 11, ' VALUES FOR NT') RFAD (5,300) (NT(T), T=1, MS) FOPMAT (15) 300 X=O. TES(1)=0. JF (MS.FO.1) CO TO 3 nn 5 1=7, "S X=X+TRD(I-1) THSITIES 5 (Real OFT 2 YI=O. TT=0. VT=0.

	ALL UTUCE/DET T AN UT ED AT VEN
	CALL FUNCES(VI,AL,S,T,FD,FL)
	CALL FORCECTUT AL & T FO FLY
22	CALL DEN (TD, YT)
	GO TO 23
	VK1=H+VK
	CALL VINCH(PST, T, AM, YT, FD, AL, VK)
	XM1=0.
	11,1=0.
	P.11=0
	FD=0.
11	TEVT/VSLR
	IF (VI) 22.21.22
	ttt=J
	DO 20
412	
412	CONTINUE
	IF (IPP.FO.1) PPINT 700
	TYPE 700
	(1 FT. (ST. T) (G) T() 412
	AM=SM7(T)/2,20462
	NN=NJ(T)
	Sar, tu((); ")(())",
	S=PT+DT(T)+D((T)/4
	VT=TH(T)+4.44R7
	HETHD(I)/DDT
	DDI=NT(T)
	IJK=T
	I=ICE+II
	DO 10 II=1.NT
210	CONTINUE
	10=252-2588
	READ (5.200) EPS
	PFAD (5,200) PSSS
195	FORMAT (' ENTEP VALUES FOR PSSS, FPS')
	1177
803	IF (ITY.LT.2) GO TO 210
	110HHFADI%G(P),2X,9HALT.(KM.),1X,10HGR.RNG(KM))
100	The set of
700	FORMAT (27 ANINT 17 IGHTINE(SEC) 17 IGHEDEBO(K/C) 17
	RFAD (5,100) NS
	PRINT 170
	1166 120
1.000	15 (JNS 50.0) CO TO 803
500	FORMAT (911)
	READ (5,500) NT, IOP, ICE, ITY, JRG, JCV, IMA, JNS
225	FOWMAT (ENTER VALUES FOR NT, IOP, ICE, ITT, IBG, JCV, IMA, JNS')
	FORMA (CUTTO UNLING FOR NO TOP TOP TOP TOP IN THE INC.)
	TYDE 225
	READ (5,200) TLD
	FRAD (3,700) TUZ
	RFAD (5.200) CA
190	FORMAT (' ENTER VALUES FOR CA, TLZ, TLD')
	TYPE 190
	TE (TTY CT 1) CO TO 210
	IF (JCVG.GE.1) GO TO 210
25	IF (TCVG.EQ.1) TCVG=2
	ALEV.
	VSR=(VSI.R-1)/VSI.R
	TD=1
	POIERI//.
	¥7=0

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V1=VT+VK1/2. PS1=PSI+PJ1/2. Y1=YT+VL1/2. CALL DEN(ID, Y1) CALL TILT (CA, TLZ, TLD, TR, AL) CALL MASS(TR, I.AM) CALL FORCES(V1.AL.S.T.FD.FL) CALL VINCR(PS1.T.AM, Y1.FD, AL, VK) CALL PSINCP(V1, PS1, T, AM, Y1, FL, AL, PSS) CALL ALT (V1, PS1, YY) CALL GRPG (V1, PS1, Y1, X) VKJ=H+VK PJ2=H+PSS Y1.7=H+YY X42=4+X V2=VT+V#7/2. PS7=PS1+P.17/2. Y2=YI+Y1.7/2. CALL DEN(ID, Y2) CALL FORCES(V2, AL, S, T, FD, FL) CALL VINCR(PS7, T, AM, Y2, FD, AL, VK) CALL PSINCR(V7, PS2, T, AM, Y7, FL, AL, PSS) CALL ALT (V2, PS2, YY) CALL GRPG (V2, PS2, Y2, X) VK 3=H+VF PJ3=H+PSS Y1, 3=H*YY XM3=H#X V3=VI+VK3 PS3=PST+PJ3 ¥3=YT+Y1,3 TR=TI+H CALL DEN(TO, Y3) CALL TILT (CA, TIZ, TLD, TB, AL) CALL MASS(TB, T, AM) CALL FOPCES(V3, AL, S, T, FD, FL) CALL VINCE(PS3, T, AM, Y3, FD, AL, VK) CALL PSTUCE(V3, PS3, T, AM, Y3, FL, AL, PSS) CALL ALT (V3, FS3, YY) CALL GRPG (V3, PS3, Y3, X) VK4=H+VK PJ4=H+PSS YL4=H+YY XM4=4+Y VF=VI+(V×1+2.*(VK2+VK3)+VK4)/6. PSF=PSI+(P,11+2,*(P.12+P.13)+P.14)/6. YF=YI+(YL1+2.*(YL2+YL3)+YL4)/6. XF=XT+(XM1+7.*(XM7+XM3)+XM4)/6. IF (ITY.GT.0) GD TO 501 TYPE 600, 1, TH, VE, PSF, YF, XF IF (TPP.FO.1) PRTMT 610,1, TB, VF, PSF, YF, XF FORMAT (11, 15, 5(11, F10.4)) TTETR VT=VF PSI=PSF YT=YF XI=XF CONTINUE JF (TOP.NF.4) GO TO 33 IF (TT.NF.NT) CC TO 10 ICD=1 TF (T.FO."S) 100=2 IF (TCVC.FO.2) TCP=1 CALL STORE (VI, PST, VI, XT, TI, VVI, PSSI, 1YYT, YYT, TTT, TD, ITD, ICD) G1) TO 37

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	33	TF (T.GF.NS) GO TO 29
	10	CONTINUE
	29	ID=1
	12	
		TYPE 600.III.TB.VE.PSE.YE.YE
		TE (TPP.FC.1) PRINT 600.TTT.TB.VF.PSF.YF.XF
		IF (ITY.FC.1) GO TO 31
		YN=PSF-PSSS
		IF (JCV.FO.2) GD TO 215
		XH=CA
		GO TO 220
	215	XM=TRD(NS)
	220	CALL CHVG(XN, XD, YD, YN, FPS, ITY, ICVG)
		TYPE 156, XN, YN
		IF (IPP.FQ.1) PPINT 156, XN, YN
	156	FOPMAT (' XN FO. ',F14.P,' , YO FO. ',F10.4)
		TE (JCV.E0.2) GO TO 235
		CA=XN
		GO TO 315
	235	TLS=XM
		GD TO MIT
	31	
		IF (JCV.F0.2) GD TO 310
		AUELA CO TO 317
	210	
	310	
	154	FORMAT (1 FATED TISI)
	134	DEAD (5.200) TIS
	815	TE (JCV-NE-2) CC TO 316
		TE (1.14 . FO. MAS) GD TO 917
		TTT=TI.S-TSD(IJK)
		NAH=TJK+1
		DO RIG TENAR, UNS
	816	TBS(T)=TKS(T)+TTT
	817	TBD(J,I*)=TIS
	315	IF (IC.1.T.2) GO TO 317
	316	TOPETRC,
		IF (IMA.AF.1) GO TO 317
		TYPE 415
	415	FOPMAT (' FUTFR SMIN')
		READ (5,200) SHIN
		DO 417 T=1, VS
	417	5#2(T)=5*2(T)+5#TN
	317	GN TO (30, 3, 69, 25, 2) TOP
	30	CUNTINIE
		FND
		SUBROUTINE MASS(T4, T, 4M)
		(0,0,0,0), (0,0,0),
		SHESH/2 20442
	0	END
		SUBPOUTINE THET (CA.TLZ.TLD.TB.AL)
		IF (TB.GT.TUZ) CO TO 1
		CO TO 5
	1	TT=TLZ+TID
		IF (TR.GT.TT) GO TO 5
	3	T=TP-TLZ
		AL=CA+T+(1T/TLD)
	5	RETIRN
		FND
	1 4	A A A A A A A A A A A A A A A A A A A
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SUBPOUTTVE DEN (ID,Y) COMMON/A1/GC, R, SS, VSP, RHO, PRZ, P, VT COMMON/H1/YZ(9),YD(9),RHZ(9),ALZ(9),HZ(9),PT(9),PR(9) IF (ID.GT.9) GO TO 10 TE (Y.LE. YP(TO)) GO TO 3 ID=ID+1 TF (Th. LE. 9) CO TO 5 10 PHOIS . P=0. GO TO 7 IF (ID.1.T.2) GO TO 4 YY=YD(ID-1) IF (Y.GT.YY) GO TO 4 TP=ID-1 GO TO 3 X=Y-YZ(10) T=PT(TO) H=HZ(TD)+T+X C=42(17)/H A=RH7(10)/(1.+AL7(T0)*X) AP=PR(TO) A=C**((1.+T)/T) PP=C**(1./T) PHO=A+++1000. P=AP+RP/10. SS=SORT(1,4*P/840) SS=SS/1000. RETIRM FND SUBROUTINF PSINCP (V, PS, T, AM, Y, FL, AL, PSS) COMMINU/A1/GC, R, SS, VSR, RHO, PPZ, P, VT ANG=1.570796327 X=P+Y A=Y+X PSS=GC/(V+=)-V/X PSS=-PSS#SIN(AMG-PS) CD=FL+T PSS=PSS+CD#S1%(AL)/(A##V#1000.) RETHRN END SUBROUTTHE VINCE (PS.T.AM.Y.FD.AL., VD) COMMON/SI/GC, P, SS, VSP, RHO, PPZ, P, VT 4"G=1.570796327 X=R+Y X=X=X IF (AL.FO.0.) GO TO 5 TT=T*SIN(ANG-AL)-FD CO TO 7 FRIS PAGE IS BEST QUALITY PRACTICALLY TT=T=FP VD=TT/(A**1000.)-(GC/X)*SIN(PS) RETURN FND SURPOUTTME ALT (V.PS. YD) YD=V*STN(PS) PETIRN FND SUBPONTIME GPRG (V, PS, Y, X) COMMON/A1/GC, P, SS, VSP, RHO, PRZ, P, VT ANC=1.570796327 X=V*SI'(ANG-PS) X=X*P/(R+Y) RETURN FND SUBROUTTYF CNVG(XN, XO, YO, YN, FPS, ITY, ICVG) DY=(YN-YD)/(XN-XD) DX=-YN/DY

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XOSXN YPEYN XN=XI+DX A=ARS(TO=XN) IF (A.GF. FPS) GD TD 5 ITY=0 ICVG=1 5 RETURN FND SUBROUTINE STORE(VI, PST, YI, XT, TI, VVI, PSSI, 1YYT,XXI.TTI,IO,ITD,ICD) GC TO (3,5) 100 VVI=VT 3 PSSI=PSI YYT=YI XXT=YT TTTTTF ITP=TP GO TO 7 5 VI=VVI PSI=PSSI YJ=YYT TXX=TX TJETTI TD=TID 7 RETURN FND SUBROUTINE FORCES(V, AL.S.T.FD.FL) COMMON/A1/GC, P, SS, VSR, PHO, PRZ, P, VT ANG=1.570796327 T=VT+(1.-VSR+P/PRZ) TF (PHO.FO.0.) GO TO 5 XM=V/SS A=XM-1.14 IF (XM.GT.O. APARA) GD TO 2 B=-16.6476+4+4 CD=.25+.3773+FXP(R) GO TO 3 TE (XM.GT.1.28) GD TO 4 2 BEARSTAT R=R##3.5 CD=.6273-30.902*P GD TD 3 CD=.17913+.5318/X* 4 IF (XM.GT.1.5) GO TO 6 3 CL=4.6-0.5*STN(ANG-2.856*XM) CDL=1.822-0.298*SIN(ANG-2.89281*XM) GO TO 5 CL=0.8+6./X* 6 CDL=-. 2919+3.2727/XM C=RH0+V+V/2. 5 C=C+S+1000000. FD=CD+CDI. +AL +AL FD=FD+C FL=CL+C RETURN END

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.EXECUTE RALLYS.FOR LINK: LOADING [LNKXCT RALLYS EXECUTION] ENTER HS ENTER 4 VALUES FOR SHZ 4000000. 135000. 120000. 10000. ENTER 4 VALUES FOR SMD 2200. 0. 700. ð. ENTER 4 VALUES FOR VAC. TH 665000. 0. î. 220000. Û. ENTER 4 VALUES FOR VELR 1.15 1.15 1. 1. ENTER 4 VALUES FOR TBD 120. 10. 150. 10. ENTER 4 VALUES FOR DI 3. 3. 3. 3. ENTER 4 VALUES FOR HI 40 10 10 10 ENTER VALUES FOR CA.TLZ.TLD -. 001 10. 25. ENTER VALUES FOR NT, IOP, ICE, IPR, IBG, JCV, IMA, JHS 42012 THIS PAGE IS BEST QUALITY FRAGTICADIN INT. TIME (SEC.) SPEED (K/S) HEADING (R) ALT. (KH.) SR. RHS (KH) INT. TIME (SEC.) SPEED (K/S) HEADING (R) INT. TIME (SEC.) SPEED (K/S) HEADING (R) INT. TIME (SEC.) SPEED (K/S) HEADING (R) ALT. (KM.) SR. RNS (KM) ALT. (KH.) GR. RHG (KH) ALT. (KH.) SR. RHS (KH) .12615+01 .53155+03 .15375+03 .2900E+03 .6822E+01 10 ENTER VALUES FOR CA.TLZ.TLD -. 005 10. 25. ENTER VALUES FOR NT, IOP, ICE, IPR, IBG, JCV, IMA, JHS 42022 ENTER VALUES FOR PSSS, EPS ð. .000002 10 .2900E+03 .7729E+01 .3612E+00 .3379E+03 .6333E+03 XN EQ. -. 66051209E-02 . YD EQ. .3612E+00

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	10	.29005+03	.31175+01	.14945+00	.2029E+03	.73315+03
24	EU.	77358819	107 · 20-3	149424	.00	
	10	.29005+03	.93505+01	.31005-01	.1143E+03	.79246+03
XH	ES.	90332990)E-02 , YO S	EQ3100E-	•01	
	10	.29005+03	.34175+01	.31595-02	.92075+02	.79285+03
XH	50.	90669290	05-05 · 10 8	EQ3159E-	-02	
	10	.2900E+03	.94235+01	.69405-04	.39585+02	.79405+03
XH	59.	30676834	E-02 , YD 9	EQ 6940E-	-04	
I	IT.	TIME (SEC.)	SPEED (K/S)	HEADING (R)	ALT. (KM.)	SR. SHS (KH)
	1	.3000E+01	.13495-01	.15715+01	.20065-01	.00002+00
	2	.50005+01	.27745-01	.1571E+01	.81725-01	. 00005+00
	3	90005+01	49795-01	15715+01	19735+00	.00005+00
	4	12005+02	59295-01	15446+01	33935+00	20775-02
	-	15005+02	75505-01	15405+01	54025400	24995-09
	2	10002+02	02005-01	15075101	70005+00	154252-02
	-	.15002402	. 75252-01	14545+01	10035401	10402-01
	5	.21002402	.11222+00	.14542401	.10772+01	.43342-01
	3	.24002+02	.13232+00	.13952+01	.14522+01	.9/195-01
	3	.2700E+02	.15402+00	.13352+01	.1832E+01	.1946E+00
	10	.30005+02	.17735+00	.1275E+01	.23515+01	.3151E+00
	11	.3300E+02	.20245+00	.1218E+04	.2900E+01	.49695+00
	15	.35002+02	.2294E+00	.11565+01	.35015+01	.73685+00
	13	.39005+02	.25935+00	.1115E+01	.4165E+01	.10425+01
	14	.42005+02	.23905+00	.10665+01	.4892E+01	.14215+01
	15	.45005+02	.32095+00	.10175+01	.56815+01	.19935+01
	15	4900E+02	.35445+00	.96995+00	.65305+01	.24355+01
	17	51005+02	.39015+00	.90395+00	.74355+01	.30975+01
	19	54005+02	42955+00	37955+00	.93975+01	39475+01
	19	57005+02	44965+00	93705+00	94155+01	47975+01
	50	20005102	51255+00	79696444	LAAGEAAD	57056401
	20	- 30002+02	52055400	75745100	11205102	20015401
	22	.53002702		70005+00	10005100	. 55512+01
	22	.55002702	.010/2700	.72052400	.12502402	.51772+01
	23	.59002+02	.55412+00	.53505+00	.14042+02	.95322+01
	24	.72005+02	.72032+00	.55305+00	.15325+02	.11255+02
	25	.75005+02	.73095+00	.5220E+00	.1555E+02	.13065+02
	26	.78005+02	.8444E+00	.59275+00	.18055+02	.1506E+02
	27	.81005+02	.9115E+00	.56525+00	.19495+02	.1726E+02
	29	.34005+02	.9824E+00	.5393E+00	.3099E+03	.19675+02
	29	.37005+02	.10575+01	.5150E+00	.2252E+02	.22305+02
	30	.90002+02	.1135E+01	.49215+00	.24115+02	.2517E+02
	31	.93005+02	.1213E+01	.4705E+00	.2574E+02	.29295+02
	32	.9500E+02	.13055+01	.45045+00	.27428+02	.3166E+02
	33	.9900E+02	.13965+01	.4314E+00	.29155+02	.35315+02
	34	.10205+03	.14925+01	.4135E+00	.30926+02	. 39245+02
	35	.10505+03	.15935+01	.3970E+00	. 32746+02	43475+02
	34	10905+03	16995+01	39135+00	34425+02	49015+09
	37	11105+03	19106+01	24445+00	34545+00	SOCOFANO
	20	11405+02	19975+01	252964444	20516100	SOLOCEADO
	20	11705+03	20505101	24005+00	40526402	- 35102+02 - 2327EA09
	37	.11/02+03	.20502+01	- 3400E+00	40002402	.00072772
	40	.12002403	-21312701	· JEOUETUU	- +2012+02	-07032702
1:		ITME (SEC.)	SPEED(K/S/	DOALE LOA	HLT. (KH.)	BR. RIG (KP)
	1	.12102+03	21772+01	.32412+00	443912+02	.71532+02
	-	.1220E+03	.21742+01	.3202E+00	.44082+02	·/3/3E+02
	3	.12305+03	.2171E+01	.31635+00	.44685+02	·7578E+02
	4	.1240E+03	.2168E+01	.3124E+00	.45358+02	.7783E+02
	5	.1250E+03	.2165E+01	.3035E+00	.45015+02	.7987E+02
	6	.1260E+03	.2162E+01	.3045E+00	.46665+02	.\$192E+02
	7	.1270E+03	.2159E+01	.3006E+00	.4731E+02	.9397E+02
	3	.12905+03	.21565+01	.2966E+00	.4794E+02	.3602E+02
	9	.1290E+03	.2153E+01	.2927E+00	.48575+02	.3806E+02
	10	.13005+03	.2150E+01	.2887E+00	.49198+02	.9011E+02
1:	IT.	TIME (SEC.)	SPEED (K/S)	HEADING (R)	ALT. (KH.)	GR. RHG .KM.
	1	.1450E+03	.2394E+01	.2322E+00	.5732E+02	.12275+03

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.1600E+03 .2675E+01 .1922E+00 .6569E+02 .1595E+03 2 .2997E+01 .1385E+00 .1750E+03 .72435+02 .2010E+03 3 4 .1900E+03 .3369E+01 .1010E+00 .78095+02 .2478E+03 .3903E+01 .3007E+03 5 .2050E+03 .5931E-01 .3261E+02 .3599E+02 .4316E+01 .4337E-01 .36055+03 .2200E+03 6 .49395+01 .2306E-01 .4288E+03 7 .2350E+03 . 3824E+02 .5074E+03 .25005+03 .5725E+01 .9446E-02 .89445+02 9 .3977E+02 .59955+03 9 .26505+03 .5737E+01 -.1304E-03 .89595+02 .29005+03 .9424E+01 -.1712E-02 10 .7109E+03 INT. TIME (SEC.) SPEED (K/S) HEADING (R) ALT. (KM.) SR. RHG (KM. .29105+03 .94245+01 -.15415-02 .3959E+02 .7192E+03 .8424E+01 -.1370E-02 .72755+03 .3957E+02 .29205+03 2 .9424E+01 -.1199E-02 .99555+02 .29305+03 .73595+03 3 .29405+03 .39555+02 .74425+03 4 .94245+01 -.10275-02 .7525E+03 .29505+03 .94245+01 -.95595-03 . 39545+02 5 5 .28505+03 .9424E+01 -.5346E-03 .3953E+02 .75095+03 .9424E+01 -.5134E-03 .39532+02 .76915+03 .29705+03 . 3953E+02 .77745+03 .94245+01 -.34225-03 9 .29905+03 .94245+01 -.17105-03 0 .29905+03 . 99525+02 .78575+03 .34245+01 .23145-05 . 99525+02 .7940E+03 .29005+03 10 ENTER VALUES FOR CA.TLZ.TLD -. 005 10. 25. ENTER VALUES FOR HT. IOP. ICE. IPR. IBS. JCV. IMA. JHS 24012 INT. TIME (SEC.) SPEED (K/S) HEADING (R) ALT. (KM.) GR. RNG (KM) INT. TIME (SEC.) SPEED (K/S) HEADING (R) ALT. (KM.) SR. RHS (KH) .13005+03 .20335+01 .52095+00 .64985+02 .74925+02 10 ENTER VALUES FOR CA.TLZ.TLD -. 01 135. 140. ENTER VALUES FOR NT. 10P. 1CE. 1PR. IBG. JCV. IMA. JNS 24212 INT. TIME (SEC.) SPEED (K/S) MEADING (R) ALT. (KM.) SP. SHE (KM) INT. TIME (SEC.) SPEED (K/S) MEADING (R) ALT. (KM.) SR. RHS (KM) .90555+01 -.94225-01 .11415+03 .71995+03 .29005+03 10 ENTER VALUES FOR CA.TLZ.TLD -. 003 135. 140. ENTER VALUES FOR NT. IOP. ICE. IPR. IBG. JCV. IMA. JNS 24222 ENTER VALUES FOR PSSS.EPS ù. .000002 10 .29005+03 .30535+01 -.30535-01 .14205+03 .71995+03 XM E9. -.703695245-02 . YO E9. -.30635-01 10 .29005+03 .30635+01 -.70055-04 . .15565+03 .71955+03 XH EQ. -. 703464435-02 . YD EQ. -. 70055-04 10 .2900E+03 .9053E+01 -.4971E-07 .15565+03 .71955+03 XH EQ. -. 70346427E-02 . YD ED. -. 4971E-07 INT. TIME (SEC.) SPEED (K/S) HEADING (R) ALT. (KM.) SR. RNS (KM) .14505+03 .22475+01 .90125+02 .10295+03 .46315+00 1 .39925+00 .9496E+02 .1347E+03 3 .15005+03 .24965+01 .2795E+01 3 .17505+03 .33175+00 .10915+03 .17115+03 .31245+01 .26355+00 .12205+03 .19005+03 ·2126E+03 .3524E+01 .19725+00 .1333E+03 5 .20505+03 .25012+03 .31455+03 .22005+03 .40075+01 .13555+00 .14265+03 5 .14955+03 .23505+03 .46045+01 .37715+03 .30895-01 3 .25005+03 .53735+01 .3663E-01 .1537E+03 .44995+03

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- 9 .26505+03 .15555+03 .54275+01 .71175-02 .53595+03 .30535+01 -.75975-03 .29005+03 .15565+03 10 .54075+03 INT. TIME (SEC.) SPEED (K/S) HEADING (R) ALT. (KH.) SR. RHS (KH) .2910E+03 .9063E+01 -.6939E-03 .1556E+03 .6496E+03 1 .15565+03 .29205+03 .80535+01 -.50795-03 2 . 55555+03 .29305+03 .80632+01 -.53182-03 .15565+03 .55445+03 3 .15565+03 .29405+03 .90535+01 -.45585-03 ·6722E+03 4 .2950E+03 .1556E+03 .58015+03 .30635+01 -.37995-03 5 .2950E+03 .30532+01 -.30395-03 .1556E+03 .59905+03 ÷ .23705+03 .90535+01 -.22795-03 .15562+03 .69592+03 .90535+01 -.15195-03 .15565+03 3 .29905+03 .70375+03 .30532+01 -.75972-04 .29905+03 .15565+03 .7116E+03 10 .29005+03 .80635+01 .22955-08 .15565+03 .71955+03 ENTER VALUES FOR CA.TLZ.TLD -.005 10. 35. ENTER VALUES FOR NT. 10P. ICE. IPR. IBG. JCV. IMA. JHS 24012 INT. TIME (SEC.) SPEED (K/S) HEADING (R) ALT. (KM.) GR. RNS (KM) INT. TIME (SEC.) SPEED (K/S) HEADING (R) ALT. (KH.) SR. RHS (KH) 10 .13005+03 .19735+01 .65595+00 .72275+02 .65555+02 ENTER VALUES FOR CA, TLZ, TLD -. 01 135. 140. ENTER VALUES FOR NT. IDP. ICE. IPR. IBG. JCV. IMA. JHS 24212 ALT. (KH.) SR. RHG (KH) INT. TIME (SEC.) SPEED (K/S) HEADING (R) INT. TIME (SEC.) SPEED (K/S) HEADING (R) ALT. (KM.) SR. RNS (KM) 10 .29005+03 .79095+01 .27145-01 .20005+03 .56995+03 ENTER VALUES FOR CA.TLZ.TLD -. 013 135. 140. ENTER VALUES FOR NT, IOP, ICE, IPR, IBG, JCV, IMA, JNS 24000 ENTER VALUES FOR PSSS.EPS ù. .000002 10 .2900E+03 .7775E+01 -.7398E-01 .1574E+03 .5590E+03 XM E0. -. 10805297E-01 , YD E0. -. 7398E-01 10 .29005+03 .78015+01 .59805-04 .19965+03 .67015+03 XM EQ. -. 10807070E-01 . VD EQ. .5980E-04 INT. TIME (SEC.) SPEED (K/S) HEADING (R) ALT. (KM.) GR. RNS (KM) .14505+03 .21705+01 .59915+00 .90505+02 .90345+02 1 .10935+03 .1500E+03 .2400E+01 .52965+00 2 .11995+03 .26675+01 3 .17505+03 .44915+00 .12665+03 .1517E+03 .19005+03 .29785+01 .35245+00 .1432E+03 .18985+03 4 .15915+03 .23375+03 5 .20505+03 .33505+01 .27575+00 .22005+03 .33035+01 .1924E+00 .1704E+03 .28455+03 S .4373E+01 .17975+03 .11595+00 .23505+03 .34335+03 .25005+03 .1356E+03 3 .51235+01 .54375-01 .41205+03 .25505+03 .5157E+01 .1199E-01 .13925+03 .49375+03 .78015+01 -.31365-04 10 .29005+03 .18855+03 .59445+03 INT. TIME (SEC.) SPEED (K/S) HEADING (R) ALT. (KM.) SR. RHS (KM) .29105+03 .73015+01 -.29215-04 .19955+03 .50195+03 .13955+03 .23205+03 .73015+01 -.25065-04 .50955+03 2 .19955+03 .51715+03 .29305+03 .78015+01 -.21915-04 3 .29405+03 .13955+03 .62475+03 .78015+01 -.18765-04 .78015+01 -.15615-04 .23505+03 .13955+03 . 53235+03 5 ÷ .23505+03 .7901E+01 -.1246E-04 .13355+03 .63995+03

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Barry reaction of the San State of the state of a

7 .29705+03 .79015+01 -.9312E-05 .1995E+03 .6474E+03 3 .2990E+03 .7901E+01 -.5152E-05 .1995E+03 .6550E+03 9 .2990E+03 .7901E+01 -.3013E+05 .1985E+03 .6526E+03 10 .2900E+03 .7901E+01 .1361E-06 .1595E+03 .6701E+03 ENTER VALUES FOR CA+TLZ+TLD ~C the state of the s

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