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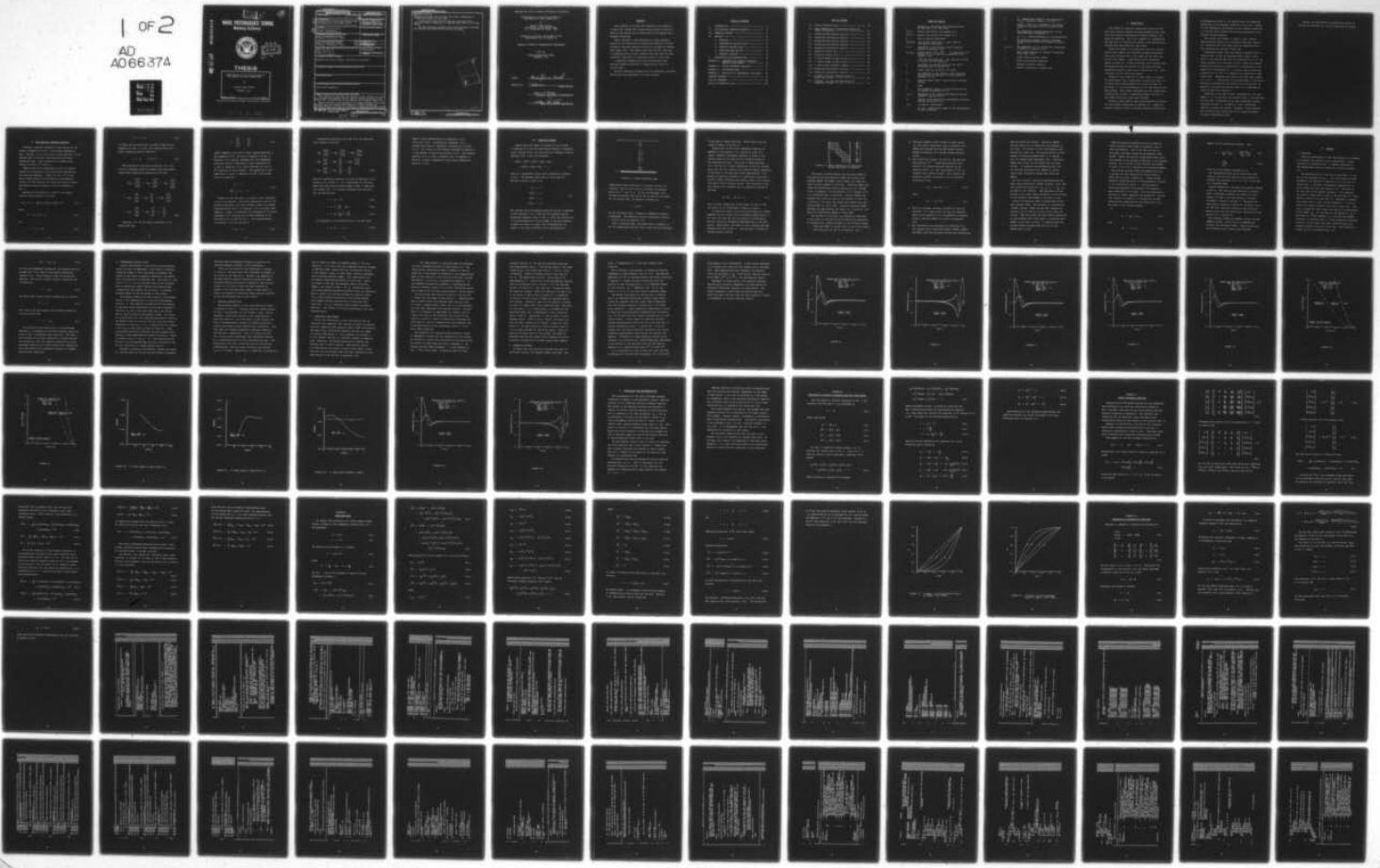
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## THESIS

INVESTIGATION OF PIPE FLOW INSTABILITY  
AND RESULTS FOR WAVE NUMBER ZERO

by

Michael James Arnold

December 1978

Thesis Advisor:

T. H. Gawain

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ABSTRACT (Cont'd)

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A nonuniform computational mesh was developed which provided dramatic reductions in computational time on a limited basis.

Two data reduction programs were also developed to process and display data generated by the main program.

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Investigation of Pipe Flow Instability  
and Results for Wave Number Zero

by

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Lieutenant, United States Navy  
B.S., University of Idaho, 1969

Submitted in partial fulfillment of the  
requirements for the degree of

MASTER OF SCIENCE IN AERONAUTICAL ENGINEERING

from the

NAVAL POSTGRADUATE SCHOOL  
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## ABSTRACT

Past research by Harrison and Johnston on the stability of pipe flow yielded only tenuous results owing to errors in setup of the problem and in formulation of the complex axis boundary conditions.

Recent advances in the formulation of these boundary conditions and application of generalized stability criteria allowed an accurate numerical solution to be made for angular wave number zero. The results show that flow for this case is characterized by certain instabilities that have not been previously identified in linearized studies of this type.

A nonuniform computational mesh was developed which provided dramatic reductions in computational time on a limited basis.

Two data reduction programs were also developed to process and display data generated by the main program.



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TABLE OF SYMBOLS

C	Constant in non-uniform mesh functions given by equations (C-32) and (C-40)
$D, D^2, \dots$	Partial derivatives with respect to $r$ .
$D^*, D^{*2}, \dots$	Partial derivatives with respect to $\eta$ .
e	Base of natural logarithms.
$\bar{e}_x, \bar{e}_r, \bar{e}_\theta$	Unit vectors along the $x$ , $r$ and $\theta$ axes in cylindrical coordinates.
F, G, H	Components of the velocity vector potential defined in equation (2-6).
$f_{11}, f_{22}, \dots$	Coefficients of $D^*Q$ , $D^{*2}Q$ , ... in equations (C-9) through (C-12) as defined in equations (C-13) through (C-22).
i	$+\sqrt{-1}$ , the imaginary unit. Also used as an index in Section III and Appendix D.
N	The number of interior points in the finite difference mesh of Section III.
O	Symbol denoting the phrase "of order".
Q	The component of the velocity vector potential derived from the component H by the change of variable, $H = rQ$ .
$R_e$	Reynolds number based on mean velocity and pipe radius.
t	Time.
U	The streamwise velocity in Pipe Poiseuille Flow as defined by equation (2-11).
u, v, w	Components of the complex perturbation velocity defined in equation (D-1).
$\bar{W}$	Complex vector potential of perturbation velocity defined in equation (D-2).
x, r, $\theta$	Cylindrical coordinates.
$\alpha$	$\alpha_R + i\alpha_I$ . Complex wave number of the perturbation in the $x$ -direction.



$\beta$	in. Complex wave number of the perturbation in the $\theta$ direction, where $n = 0, 1, 2, 3, \dots$
$\delta$	$1/(N+1)$ . The $r$ or $\eta$ increment in the finite difference approximations of the derivatives of $Q$ .
$\eta$	The independent variable replacing $r$ in the nonuniform mesh of Appendix C.
$\gamma$	$\gamma_R + i\gamma_I$ . Complex frequency of the perturbation.
$\bar{\Gamma}$	The vorticity transport equation expressed in abbreviated notation as defined in equation (2-7).
$\Gamma_x, \Gamma_r, \Gamma_\theta$	The components of $\bar{\Gamma}$ in cylindrical coordinates as defined in equation (2-7).
$\lambda$	Mesh offset parameter as defined in equations (C-32) and (C-40).
$\nabla$	Linear vector operator (nabla)
$\times$	Vector cross-product operator.
[ ]	Brackets enclosing a matrix.
{ }	Brackets enclosing a column vector.



## I. INTRODUCTION

The problem of finding an analytical solution to the pipe flow stability problem has been pursued actively ever since the classical experiments of Osborne Reynolds [10] about 100 years ago. Up to now, however, no investigation has been able to satisfactorily predict flow instabilities, although many approaches have been taken.

Salwen and Grosch [11] studied pipe flow with various angular wave numbers and sinusoidal streamwise perturbations and concluded that it was stable for all axial and angular wave numbers. Perturbations with exponential growth in space but a purely sinusoidal time variation were researched by Garg and Rouleau [2] and those with both exponential growth in space and in time by Gill [3]. Both concluded that the flows were stable.

Because of this inability of linear theory to account for experimental fact, explanations by Davey and Drazin [1] involving finite disturbances and by Huang and Chen [5] and Leite [7] involving conditions at the pipe entrance have been offered. While these investigations have indeed shown instabilities to exist, a completely general solution to the linear problem has never been achieved.

Recently a more general theory was presented by Harrison [4] and further investigated by Johnston [6]. These two studies, however, failed to produce conclusive results due

to mathematical errors in the problem setup and inadequate formulation of the boundary conditions at the axis. Gawain [9] has subsequently formulated the axis boundary conditions in a new way which corrects the previous discrepancies and promises further advances.

For angular wave number,  $n$ , equal to zero, radical simplifications result in the governing equations (Section II), indicating that this case should be approached first. This investigation centers on that case.

Preliminary checks using the computer program of Ref. 6 revealed that, of the two eigenfunctions,  $G$  and  $H$ , which occur in this problem and which are uncoupled for  $n = 0$ , the latter appeared to be the more critical. Hence the present research was arbitrarily restricted to investigation of the stability of eigenfunction  $H$ . A similar study of the other eigenfunction,  $G$ , for  $n = 0$  remains to be completed at some future time. Comparable calculations for other wave numbers ( $n = 1, 2, 3, \dots$ ) also remain to be accomplished in the future. Extensive and systematic calculations of this type will be essential to provide the factual basis for a comprehensive theory of pipe flow stability.

Reverting to the case at hand, eigenfunction  $H$  for wave number  $n = 0$ , we note that the program of Ref. 6 was rewritten for this case, incorporating the newly formulated boundary conditions of Ref. 9. In addition, a new, generalized stability criteria was adopted. Moreover, a new technique was introduced which allows the use of nonuniform meshes to reduce computational time.

Lastly, two data reduction programs were written to process data produced by the main investigative program.



## II. THE VORTICITY TRANSPORT EQUATION

Although a complete treatment of this subject is contained in Appendix A of Ref. 4 and further addressed in Ref. 6 and Ref. 9, it is felt that a brief overview is still required here to maintain continuity with previously referenced works. This discussion is an abbreviated version of Section II of Ref. 6.

Laminar flow of an incompressible fluid of constant viscosity is governed by the Navier-Stokes equation and the continuity equation. Taking the curl ( $\nabla \times$ ) of the Navier-Stokes equation and introducing a perturbation velocity ( $\bar{v}$ ) and vorticity ( $\bar{\omega}$ ) gives the vorticity transport equation which is equation (A-10) of Appendix A, Ref. 4.

Expressing this equation in terms of the complex velocity vector potential,  $\bar{W}$ , gives

$$W(x, r, \theta, t) = (\bar{e}_x F(r) + \bar{e}_r G(r) + \bar{e}_\theta H(r)) e^X \quad (2-1)$$

where

$$X = \alpha x + \beta \theta + \gamma t \quad (2-2)$$

and

$$\bar{v} = \nabla \times \bar{W} \quad (2-3)$$



$$\bar{\omega} = \nabla \times \bar{v} . \quad (2-4)$$

It should also be noted that, as shown in part one of Appendix G in Ref. 4,  $\alpha$  and  $\gamma$  are complex while  $\beta$  is a purely imaginary quantity defined by

$$\beta = i n \quad n = 0, 1, 2, \dots \quad (2-5)$$

When expressed in the form of equation (2-1), the vorticity transport equation becomes three simultaneous fourth-order differential equations of the form

$$\begin{aligned} & [M_4] \begin{Bmatrix} D^4 F \\ D^4 G \\ D^4 H \end{Bmatrix} + [M_3] \begin{Bmatrix} D^3 F \\ D^3 G \\ D^3 H \end{Bmatrix} + [M_2] \begin{Bmatrix} D^2 F \\ D^2 G \\ D^2 H \end{Bmatrix} \\ & + [M_1] \begin{Bmatrix} DF \\ DG \\ DH \end{Bmatrix} + [M_0] \begin{Bmatrix} F \\ G \\ H \end{Bmatrix} - \gamma ([N_2] \begin{Bmatrix} D^2 F \\ D^2 G \\ D^2 H \end{Bmatrix} \\ & + [N_1] \begin{Bmatrix} DF \\ DG \\ DH \end{Bmatrix} + [N_0] \begin{Bmatrix} F \\ G \\ H \end{Bmatrix}) = \begin{Bmatrix} 0 \\ 0 \\ 0 \end{Bmatrix} \end{aligned} \quad (2-6)$$

Equations (2-5) may be further expressed in the abbreviated form

$$\bar{\Gamma} = \begin{pmatrix} \bar{\Gamma}_x \\ \bar{\Gamma}_r \\ \bar{\Gamma}_\theta \end{pmatrix} = \begin{pmatrix} 0 \\ 0 \\ 0 \end{pmatrix} \quad (2-7)$$

where  $\bar{\Gamma}$  appears to be a set of three coupled equations in the components of  $\bar{W}$ . As given in Appendix B of Ref. 4, equations (2-7) actually represent only two independent conditions and by an appropriate linear combination of  $\Gamma_x$  and  $\Gamma_\theta$ , equations (2-6) can be expressed as a set of two equations in three unknowns. The appropriate linear combination is given in Appendix B of Ref. 4 and yields the set of equations

$$\begin{aligned} \Gamma_r &= 0 \\ -\frac{in}{r} \Gamma_x + \alpha \Gamma_\theta &= 0. \end{aligned} \quad (2-8)$$

Except for the case where  $n$  is equal to zero, equations (2-8) do not uncouple. The linear combination given by the second of equations (2-8) does, however, reduce the highest order derivative of  $G(r)$  in equations (2-6) to second order. Appendix C of Ref. 4 illustrates the redundancy of the three components of  $\bar{W}$ , allowing one of these components to be arbitrarily set to zero for all  $r$ . The maximum benefits of equations (2-8) are obtained if

$$F(r) = 0 \quad (2-9)$$

Incorporating equations (2-8) and (2-9) into equations (2-6) results in the form

$$\begin{aligned}
 & [M'_4] \begin{Bmatrix} D^4 G \\ D^4 H \end{Bmatrix} + [M'_3] \begin{Bmatrix} D^3 G \\ D^3 H \end{Bmatrix} + [M'_2] \begin{Bmatrix} D^2 G \\ D^2 H \end{Bmatrix} \\
 + & [M'_1] \begin{Bmatrix} DG \\ DH \end{Bmatrix} + [M'_0] \begin{Bmatrix} G \\ H \end{Bmatrix} - \gamma ([N'_2] \begin{Bmatrix} D^2 G \\ D^2 H \end{Bmatrix} \\
 + & [N'_1] \begin{Bmatrix} DG \\ DH \end{Bmatrix} + [N'_0] \begin{Bmatrix} G \\ H \end{Bmatrix} ) = \begin{Bmatrix} 0 \\ 0 \end{Bmatrix} \tag{2-10}
 \end{aligned}$$

where the coefficient matrices are given by equations (2-10) through (2-17) of Ref. 6. It is appropriate to note that these same coefficient matrices appear in Ref. 9, equations (A1) through (A9), in a slightly different form resulting from the substitutions

$$U = 2(1 - r^2) \tag{2-11}$$

$$t = \alpha^2 + \frac{\beta^2}{r^2} \quad \text{and} \tag{2-12}$$

$$T = \alpha U - \frac{1}{R_e} \left( \alpha^2 + \frac{\beta^2}{r^2} \right) . \tag{2-13}$$

As discussed in the previous section, the case where

$$\beta = in, \quad n = 0 \tag{2-14}$$



leads to great simplifications in equations (2-10), (2-12) and (2-13). In particular, equations (2-10) uncouple and allow an independent investigation of either H or G. As a result of the findings discussed in Section I, it was decided to explore the function H only. This reduced equation (2-10) to that of equation (A-6) of Appendix A, which is a linear, homogeneous fourth order differential equation in  $H(r)$ .



### III. NUMERICAL METHODS

Substituting the change of variable  $H = rQ$  as given in equation (A-1) and the coefficients defined in equations (A-11) through (A-18) into the vorticity transport relation, equation (A-6), gives the expression

$$\begin{aligned} M_4 D^4 Q + M_3 D^3 Q + M_2 D^2 Q + M_1 DQ + M_0 Q \\ - \gamma [N_2 D^2 Q + N_1 DQ + N_0 Q] = 0, \end{aligned} \quad (3-1)$$

which is a homogeneous fourth order differential equation in  $Q(r)$ . The boundary conditions for this case are derived in detail in Ref. 9 as

$$\begin{aligned} Q(1) &= 0 \\ DQ(1) &= 0 \\ DQ(0) &= 0 \\ D^3 Q(0) &= 0. \end{aligned} \quad (3-2)$$

The boundary finite difference equations derived in Appendix B from equations (3-2), along with the standard central difference equations given in Ref. 6, allow the function  $Q(r)$  to be approximated by a finite number of discrete unknowns. As shown by Figure 3-1 below, the non-dimensionalized radius of the pipe is divided into a one-dimensional

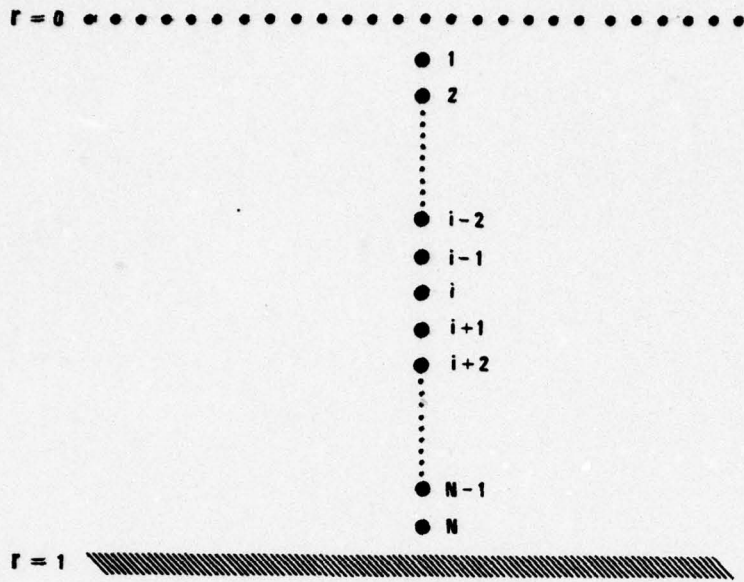


Figure 3-1 Finite Difference Mesh

computational mesh consisting of  $N$  interior points,  $N+1$  intervals, and  $N+2$  total points, including the boundary points at  $r = 1$  and  $r = 0$ . As will be discussed later, the spacing between these points may or may not be uniform. For the uniform case, the spacing is defined by

$$\delta = 1/(N+1) . \quad (3-3)$$

For the nonuniform case, a change of independent variable is performed. The spacing of the new independent variable,  $\eta$ , is still given by equation (3-3).

With a nonuniform mesh, the points shown in Figure 3-1 will be concentrated near the axis or near the wall according

to the type of offset specified. These effects are discussed in detail in Section IV.

Substitution of the finite difference equations of Appendix B into equation (3-1) results in a set of  $N$ , linear, algebraic difference equations in terms of the unknown value of  $Q$  at each of the  $N$  interior points of the computational mesh. Since each of these equations is of the form of a linear combination of the  $i$ th, central, point and the two, three or four adjacent points (depending on the order of the derivative being approximated), this system of equations consists of a coefficient array multiplying a vector containing the unknown value of the function  $Q$  at each of the  $N$  interior points. This technique allows the problem to be converted into an eigenvalue problem of the form

$$[X] \{Q\} - \gamma [Y] \{Q\} = 0 \quad (3-4)$$

with the basic composition of the arrays  $[X]$  and  $[Y]$  and the vector  $\{Q\}$  as illustrated in Figure 3-2 below.

It should be noted at this point that Figure 3-2 differs somewhat from the normal finite difference banded matrix in the first two rows and last row because of the method of deriving the finite difference approximations at the boundaries. Additionally, the order of the  $N$  unknowns has been reversed from that of Ref. 6. This was done to conform to standard matrix notation.



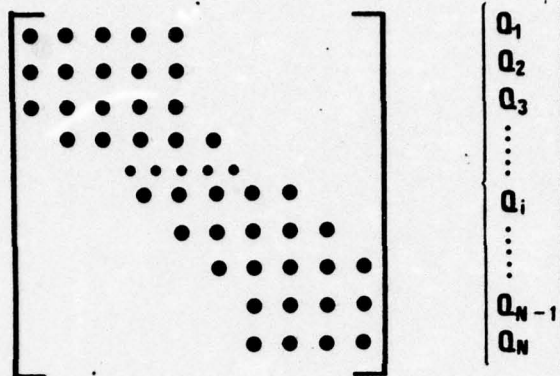


Figure 3-2 Basic Composition of Coefficient Arrays and Vector of Unknowns

This array is established by the subroutine MSET2 in conjunction with the subroutine MSET1 and function subprograms CQM1E1 and CQM2E1, which compute the numerical value for each element in the array. Subroutine MSET1 provides the coefficients given by equations (A-11) through (A-18) of Appendix A or by equations (C-24) through (C-31) if a nonuniform mesh is specified. Function CQM1E1 then computes the values for each of the elements of array [X] in equation (3-4) using the coefficients passed from subroutine MSET1 in vector CQM1. Function subprogram CQM2E1 performs the same function for matrix [Y] in equation (3-4) using the coefficients passed in vector CQM2.

The solution of the eigenvalue problem as formulated to this point is carried out by the controlling subroutine of program PIPE0, subroutine STAB, by the following steps:

- 1) Subroutine MSET2 is called twice to set up the coefficient matrices [X] and [Y] of equation (3-4).

- 2) Subroutine CDMTIN is then called to invert matrix [Y], the second coefficient array in equation (3-4). CDMTIN was obtained from the IBM Library routine CMTRIN by modifying it to accept double precision arrays.
- 3) Both coefficient arrays, [X] and [Y], are then pre-multiplied by  $[Y]^{-1}$ . Since multiplication of an array by its inverse invariably results in the identity matrix, [I], only the product  $[Y]^{-1}[X]$  is computed using subroutine MULM. This converts the eigenvalue problem of equation (3-4) to the more conventional form

$$([Z] - \gamma[I])\{Q\} = 0 \quad (3-5)$$

where

$$[Z] = [Y]^{-1}[X] \quad (3-6)$$

- 4) Since all programs currently available for solving equations (3-5) require that the real and imaginary parts of the elements of [Z] be presented in separate arrays, subroutine DSPLIT is called to accomplish this.
- 5) The eigenvalues and eigenvectors of equations (3-5) are computed using subroutines EBALAC, EHESSC, ELRH2C and EBBCKC which are available through the International

Math and Statistics Library. Subroutine EBALAC balances matrix [Z] by equalizing the exponents of all terms. The details of this transformation are retained for later use. The balanced matrix is then passed to subroutine EHESSC where it is reduced into the complex upper Hessenberg form. Subroutine ELRH2C then solves for the eigenvalues and eigenvectors. To transform the eigenvectors back into the original unbalanced form, EBBCKC is finally called using information passed from subroutine EBALAC.

For each solution, subroutine STAB determines the least stable eigenvalue (largest algebraic value) and then writes the values of  $N$ ,  $R_e$ ,  $\alpha_R$ ,  $\alpha_I$ ,  $\lambda$ ,  $\gamma_{RL}$ ,  $\gamma_{IL}$  and KSET to file FT02F001. The eigenvector corresponding to the least stable eigenvalue is also written to FILE FT02F001 when MODENO is set equal to one.

Control of subroutine STAB is accomplished by the main program, PIPE0. This program is a time-sharing (CP/CMS) program. Modes one and three compute the stability of the flow for a given set of input conditions. Mode one writes the least stable eigenvector to FILE FT02F001 while this output is inhibited when MODENO is set equal to three. To generate data for program EIGFCN, program PIPE0 must be run with MODENO equal to one.



Mode two operation generates a grid of stability values (stability map) based on parameters read in from FILE FT01F001. Due to the long run time in this mode, only small meshes can be generated under CP/CMS. Longer runs must be accomplished under batch, with changes to the program as specified in the comments section. Data is output to file FT03F001 when MODENO is equal to two and is compatible with program STBCONT.

The plotting programs EIGFCN and STBCONT were used to process the data generated by program PIPE0 in modes one and three, respectively. Program EIGFCN generates normalized plots of the perturbation velocity,  $u$ , as a function of radius,  $r$ . The perturbation velocities generated in accordance with Appendix D were normalized in two steps. First the perturbation velocity of largest magnitude was determined. Letting this velocity be termed  $u_c$ , a normalizing constant producing unit magnitude and zero phase angle in  $u_c$  was found in the following manner:

If

$$u_c = u_{RC} + iu_{IC} , \quad (3-7)$$

then

$$Cu_c = 1 + i(0) \quad (3-8)$$

where C is the normalizing constant. Thus,

$$C = \frac{1}{u_{RC} + iu_{iC}} = \frac{u_{RC} - u_{iC}}{(u_{RC}^2 - u_{iC}^2)} \quad (3-9)$$

$$= \frac{\bar{u}_C}{|u_C|^2} \quad (3-10)$$

where  $\bar{u}_C$  is the complex conjugate of  $u_C$ .

The nondimensionalized radius values were taken directly from the data cards for uniform meshes or computed from equations (C-32) or (C-40) in the case of a nonuniform mesh.

Program STBCONT plots the stability contours against  $\alpha_R$  and  $\alpha_I$ . The stability map generated by program PIPE0 is searched columnwise and rowwise for sign changes for each of the three stability criteria discussed in Section V and Ref. 9. The points are then plotted, producing contours of incipient, critical and fully developed instability and areas that denote stable flow and subcritical, supercritical and hypercritical instability.

Both programs, EIGFCN and STBCONT, utilize the NPS VERSATEC plotter, certain built-in VERSATEC subroutines, and subroutine PLOTG. These routines are only accessible when running under FORTCLGW.

#### IV. RESULTS

##### A. STABILITY

Since an understanding of the term stability is necessary to interpret the results of this investigation, a brief discussion is presented here. A complete discussion of the generalized criteria of stability is given by Gawain [9].

The characteristics of the flow for the case  $n = 0$  are set by the parameters  $R_e$  and  $\alpha$ . For fixed values of these parameters, the solution of equations (3-5) is a set of  $N$  eigenvalues,  $\gamma$ , and their corresponding eigenvectors,  $Q$ . As can readily be seen from equation (2-1), the value of the real part of the complex eigenvalue  $\gamma$  will determine the growth or decay rate in time of the perturbation. Since positive values of the real part of  $\gamma$  represent an exponential growth rate in time, the most important  $\gamma$  is the one having the largest algebraic value for its real part. This root is termed the least stable root and will be represented by the symbol  $\gamma_{RL}$ . As the stability represented by  $\gamma_{RL}$  is that seen by a fixed observer, it is not the most general criterion. As derived in Ref. 9, a more appropriate stability criterion is that based on an axis system moving at the average volumetric velocity of the flow. This criteria is termed  $\gamma_{RL}^*$  and is defined by Ref. 9 as



$$\gamma_{RL}^* = \gamma_{RL} + \alpha_R . \quad (4-1)$$

For this and subsequent discussions, the subscript will be dropped and  $\gamma^*$  will refer to the quantity defined by equation (4-1). Three stability cases arise from this equation. The first is termed incipient instability and is defined by

$$\gamma^* = -|\alpha_R| . \quad (4-2)$$

The second case, termed critical instability, is given by

$$\gamma^* = 0 \quad (4-3)$$

and, lastly, the case termed fully developed instability is said to exist when

$$\gamma^* = +|\alpha_R| . \quad (4-4)$$

The transition from stable flow to fully developed instability is progressive and several distinct stages are given in Ref. 9 to describe this transition. The region from incipient to critical instability is termed subcritical instability, that from critical instability to fully developed instability is called supercritical instability while that beyond fully developed instability is termed hypercritical instability.

## B. PERTURBATION VELOCITY PLOTS

Initial investigation of the function  $Q$  was centered around plotting its appearance in the region of interest. A Reynolds number of 1150 (2300 based on diameter) was chosen as this value is generally accepted as the nominal value for transition to turbulent flow. The value of  $\alpha$  was set at  $-0.5 + i 10.0$  for the major part of the investigation as preliminary checks revealed that supercritical instabilities were present for this value. A secondary Reynolds number of 4000 was chosen to show trends.

The quantity chosen as the most realistic and representative of the eigenfunction  $Q$  is the axial perturbation velocity,  $u$ . This quantity was derived from the elements of the least stable eigenvector as outlined in Appendix D. Initially,  $R_e$  and  $\alpha_I$  were held fixed and  $\alpha_R$  was varied over a range of positive and negative values. For values of  $\alpha_R$  below about two, the normalized perturbation velocity was found to have all activity near the axis with a decay essentially to zero by  $r = 0.3$ . A typical plot of  $u$  versus  $r$  for an  $\alpha_R$  in this range is shown in Figure 4-1. When  $\alpha_R$  was made sufficiently positive, the plot changed significantly in both appearance and region of activity. Figure 4-2 shows a plot of  $u$  for  $\alpha_R = 2.5$ . The activity can now be seen to be concentrated near the wall, with most of the activity occurring at  $r$  values greater than 0.7.

Although no particular relationship between the nature of  $u$  and the stability of the flow was evident or expected,

the plots were nevertheless valuable as indicators for various parameters involved in the investigation.

First, as can be seen by the differences in Figures 4-1 and 4-2, the plots were ideal indicators of changes in the nature of the function  $Q$ . Secondly, the adequacy of the mesh could be directly observed by noting the number of points defining the curves in regions of high activity. Figures 4-3, 4-4 and 4-5 show the same conditions as Figure 4-1 but with decreasing number of mesh points,  $N$ . Lastly, the effects of nonuniform meshes could be observed as will be discussed later in this section.

#### C. STABILITY CONTOUR PLOTS

The principal results of this investigation are shown in Figures 4-6 and 4-7. Although these two figures pertain to only a limited portion of the complex  $\alpha$  plane, they do represent a significant advance in the investigation of pipe flow stability. As can be seen in these figures, the flow is characterized by regions of differing stability, ranging from stable through supercritical instability. Note that these two figures correspond to Reynolds numbers of 1150 and 4000, respectively. This is a result that has not, to this writer's knowledge, been heretofore achieved by a linearized analysis of fully developed pipe flow. The figures also show that, as has been born out by previous investigations, flow for purely sinusoidal oscillations ( $\alpha_R = 0$ ) is stable. Additionally, a comparison of Figures 4-6



and 4-7 shows the effect of Reynolds number on the flow stability. It is clear from this comparison that an increase in Reynolds number reduces the size of the stable regions in the complex  $\alpha$  plane; in other words, stability decreases with increasing Reynolds number. This trend agrees with our general experience pertaining to fluid flow. Lastly, the effect of the real and imaginary parts of the wave number  $\alpha$  can readily be seen. For  $\alpha_R$ , increasingly negative values produce successively greater levels of instability. While a contour plot was not produced for positive values of  $\alpha_R$ , point checks of stability in this region suggest that somewhat similar contours exist in the right half-plane also. For  $\alpha_I$ , increasing values produce increasing stability. This effect is also more pronounced at the lower Reynolds number.

#### D. NONUNIFORM MESH EFFECTS

One of the difficulties in this investigation was the relatively long computing time required to obtain an accurate solution, especially when operating under CP/CMS (time-sharing). The major factor controlling computing time was the number of interior mesh points,  $N$ . As an example, an increase in  $N$  of 50 percent resulted in a fourfold increase in computing time. Therefore, the desired objectives of rapidity and accuracy were in direct conflict. Additionally, follow-on investigations for values of angular wave number  $n$  other than zero involve matrices twice the order required for this case because of the coupling of equations (2-8).

For these reasons, a nonuniform mesh was developed to obtain increased accuracy at lower values of  $N$ . The nature of the velocities as seen in Figures 4-1 and 4-2 shows that a high degree of resolution in the computational mesh is only required in the vicinity of the axis ( $\alpha_R$  less than about 2) or the wall ( $\alpha_R$  greater than about 2). It was therefore theoretically possible to redistribute the points at moderate values of  $N$  to attain resolutions equivalent to much finer (and more time-consuming) uniform meshes.

As can be seen from Figures 4-8 and 4-9, the value of  $\gamma^*$  varies with the number of mesh points,  $N$ . Theoretically, each of these curves would approach some limiting value if  $N$  were increased without bound, and it is this theoretical limit that represents the required solution. In practice, it is adequate to approximate the unknown limit by a point that lies on the relatively flat portion of the curve at a value of  $N$  which is practically attainable and which does not involve a prohibitively long computing time. It has been found in this investigation that  $N = 79$  fulfills these conditions.

The conversion to a nonuniform mesh involved a change of independent variable and the introduction of an analytical function to control the distribution of the mesh points. The details of these steps are given in Appendix C. By varying the mesh offset parameter,  $\lambda$ , it was possible to vary  $\gamma^*$  over a wide range. To determine when the high

accuracy solution ( $N = 79$ ) and the nonuniform solutions were approximately equal,  $\gamma^*$  was plotted versus  $\lambda$  for fixed values of  $R_e$ ,  $\alpha$  and  $N$  with the value of  $\gamma^*$  for  $N = 79$  as a reference. Figure 4-10 shows a plot of this type for  $N = 31$ . The appropriate value of  $\lambda$  can be seen to be approximately 1.1. Figure 4-11 is the perturbation velocity plot of the solution for  $N = 31$  and  $\lambda = 1.1$  for the same  $R_e$  and  $\alpha$  as Figure 4-1. Note that the  $\gamma^*$  values are equal for these two figures. While the resolution of Figure 4-11 is not quite as fine as that of Figure 4-1, a comparison of Figure 4-11 with Figure 4-5 makes the improved resolution obvious. Figures 4-2 and 4-12 are similar to Figures 4-1 and 4-11 except that a wall offset was used. Note that for this case  $\lambda = 1.2$ , which points to a drawback of the nonuniform mesh, that of dependence on input conditions. While a check of  $\lambda$  dependence on  $\alpha$  was not made, it most probably exists. There is also, however, the possibility that for small regions of the complex  $\alpha$  plane, the variations in  $\lambda$  are small enough to allow an average value of  $\lambda$  to be nearly optimum for the entire region. While not used for the main results of this study, the method as developed here may well prove to be of maximum utility in follow-on investigations of higher angular wave numbers.

#### E. NUMERICAL ACCURACY

To ensure that the solutions presented here were of sufficient accuracy, two separate checks were made. The



first,  $\gamma^*$  dependence on  $N$ , is the most commonly used criterion.

For a solution to be accurate, it should be virtually independent of mesh fineness, that is, of  $N$ . The required magnitude of  $N$  for an accurate solution was found by plotting  $\gamma^*$  against  $N$ . Figures 4-8 and 4-9 both show that the solution is well converged for  $N = 79$  at Reynolds numbers of 1150 and 4000, as  $\gamma^*$  changes by only .001 to .003 from  $N = 31$  to  $N = 79$  for both values of Reynolds number.

The second verification of the solution, so obvious that it is sometimes overlooked, involves simply substituting the numerical solution (least stable eigenvector) into the governing equation to ensure that it is indeed being satisfied. A short program was independently written to check the finite difference representations of equation (3-1) at the first and last interior stations and at a mid-radius station. Initial checks of numerical solutions yielded unsatisfactory results and led to the discovery of various programming errors. In particular, it was discovered that four double precision constants in the finite difference approximations were lacking the required "D0" exponent. Elimination of these seemingly trivial errors resulted in a surprising four order-of-magnitude improvement in the accuracy of the solution, with the left side of equation (3-1) improving from order  $10^{-4}$  to order  $10^{-8}$ .

It is instructive to note at this point that the order of magnitude of the left side of equation (3-1) is not the

true measure of its satisfaction. A more correct procedure is to compare this value with the largest term in the equation. When examined from this viewpoint, the relative error for solutions at  $R_e = 1150$  and  $R_e = 4000$  are found to be of order  $10^{-11}$  to  $10^{-12}$ , a very satisfactory result.

Therefore, by these results, the solutions presented here are both virtually independent of  $N$  and satisfy the governing differential equation to a high degree. The efforts expended to reach these conclusions were well worth the result and also point out that attention to detail is fundamental to accurate numerical results.

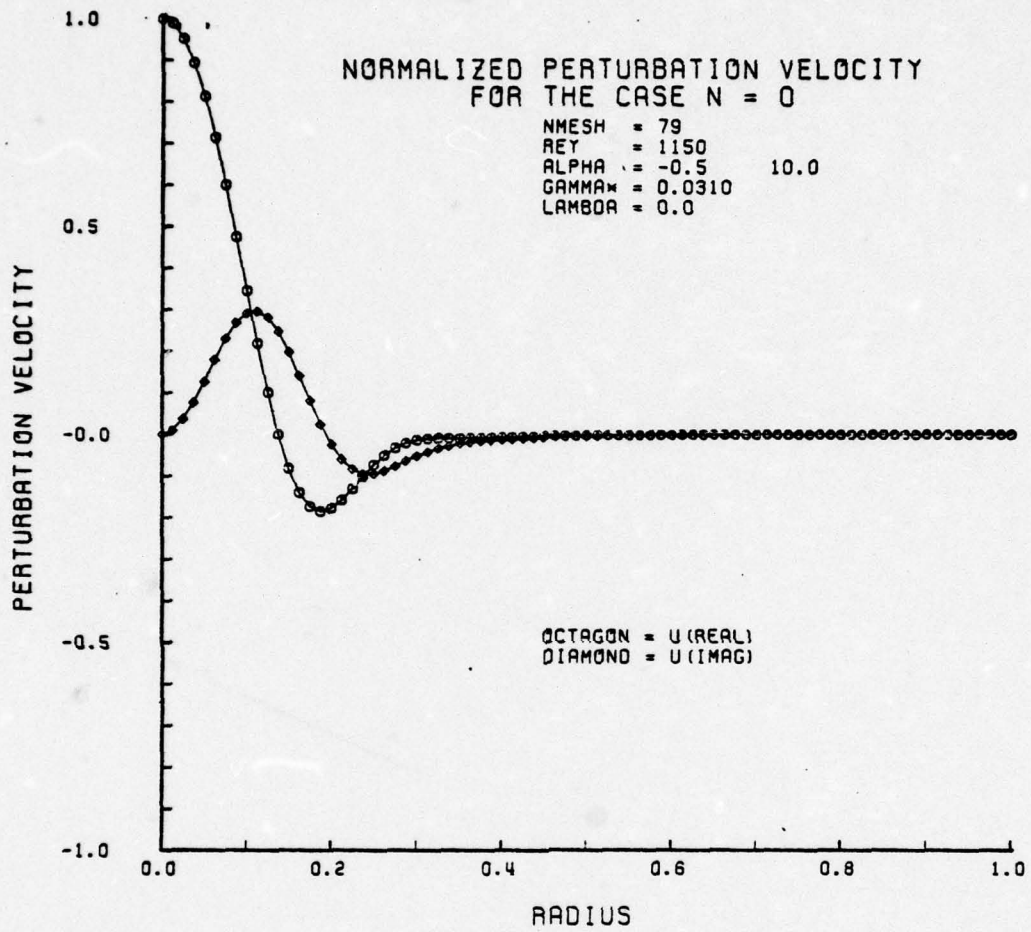


FIGURE 4-1



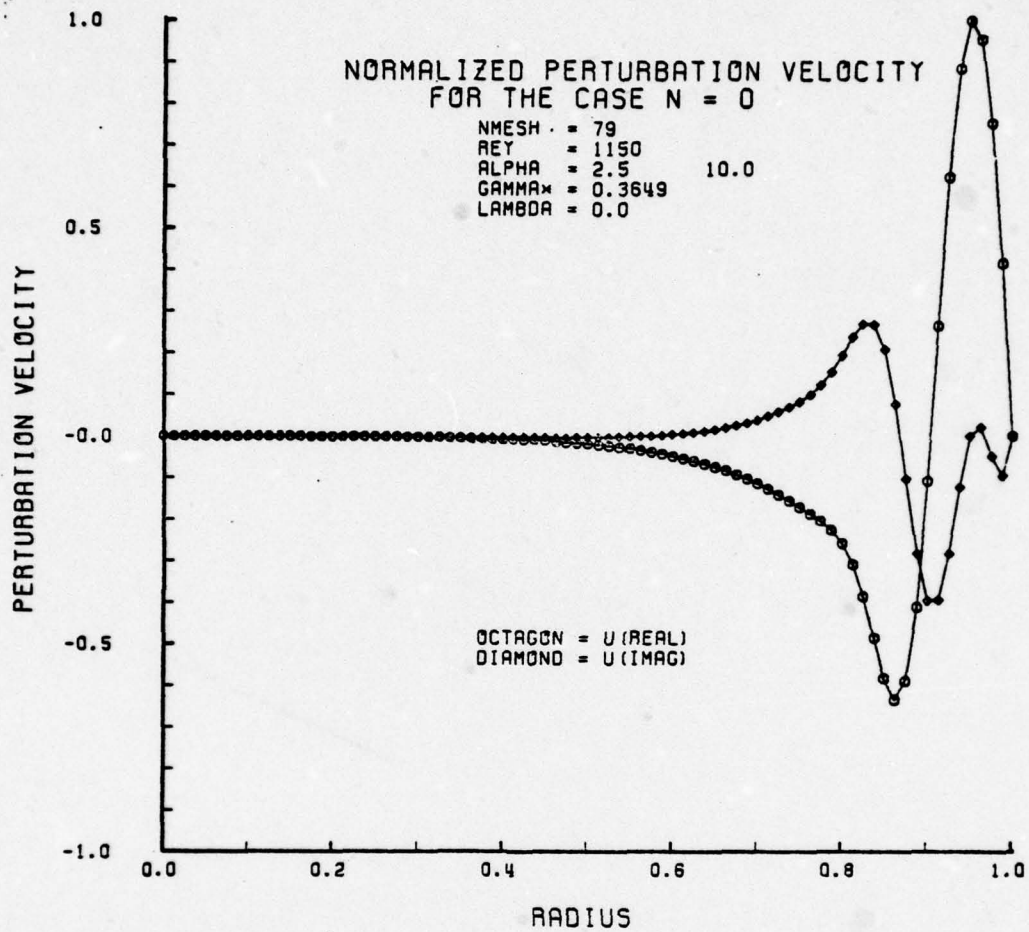


FIGURE 4-2

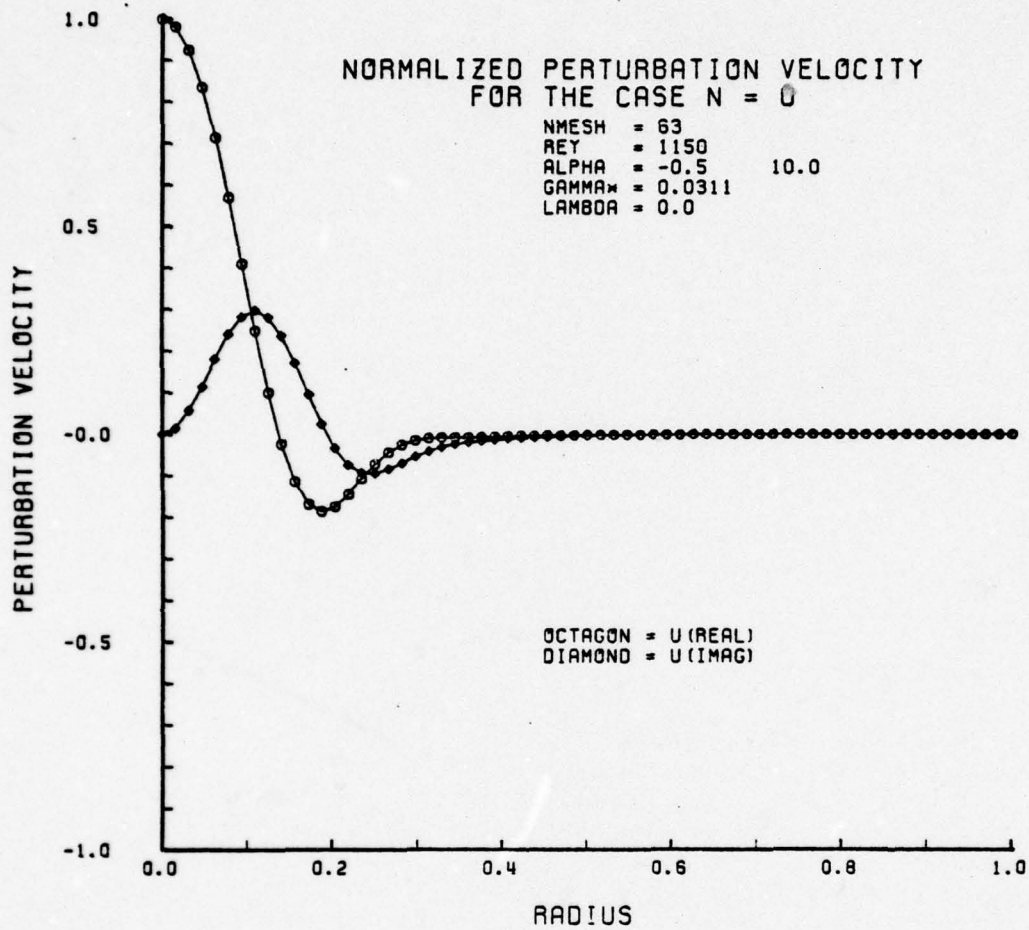


FIGURE 4-3

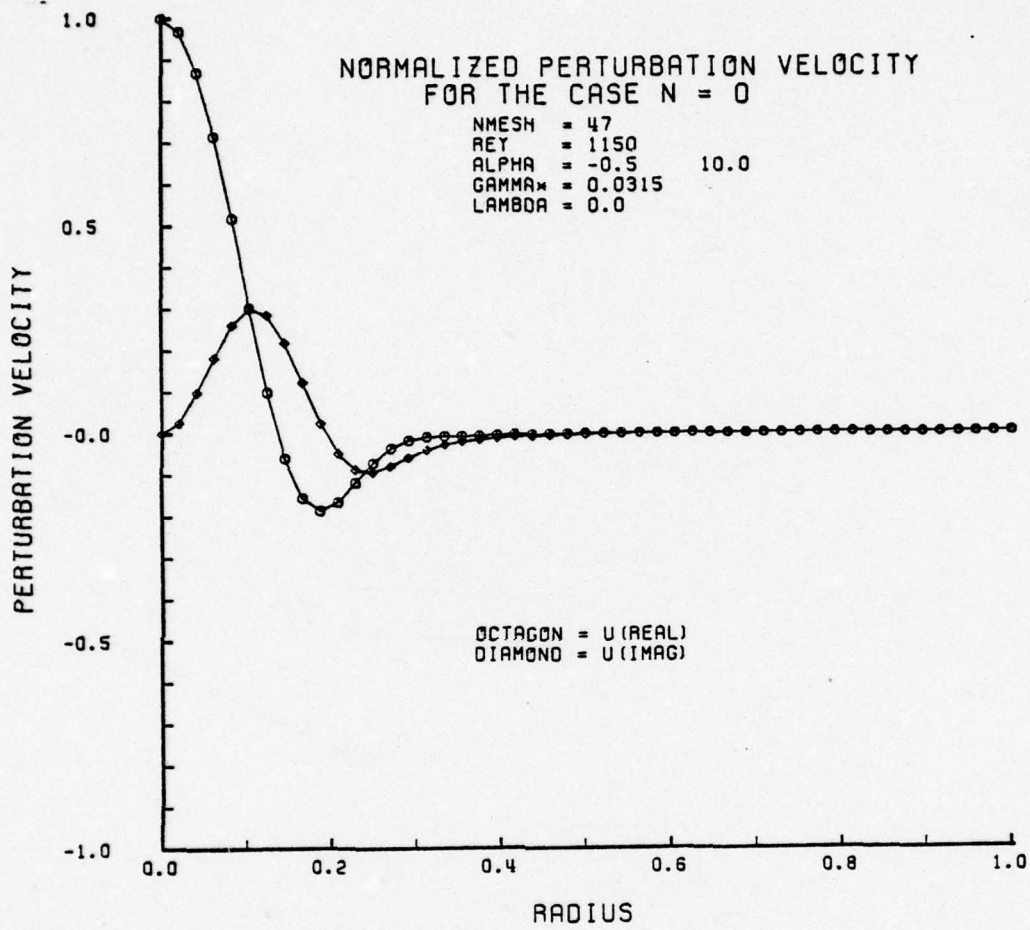


FIGURE 4-4



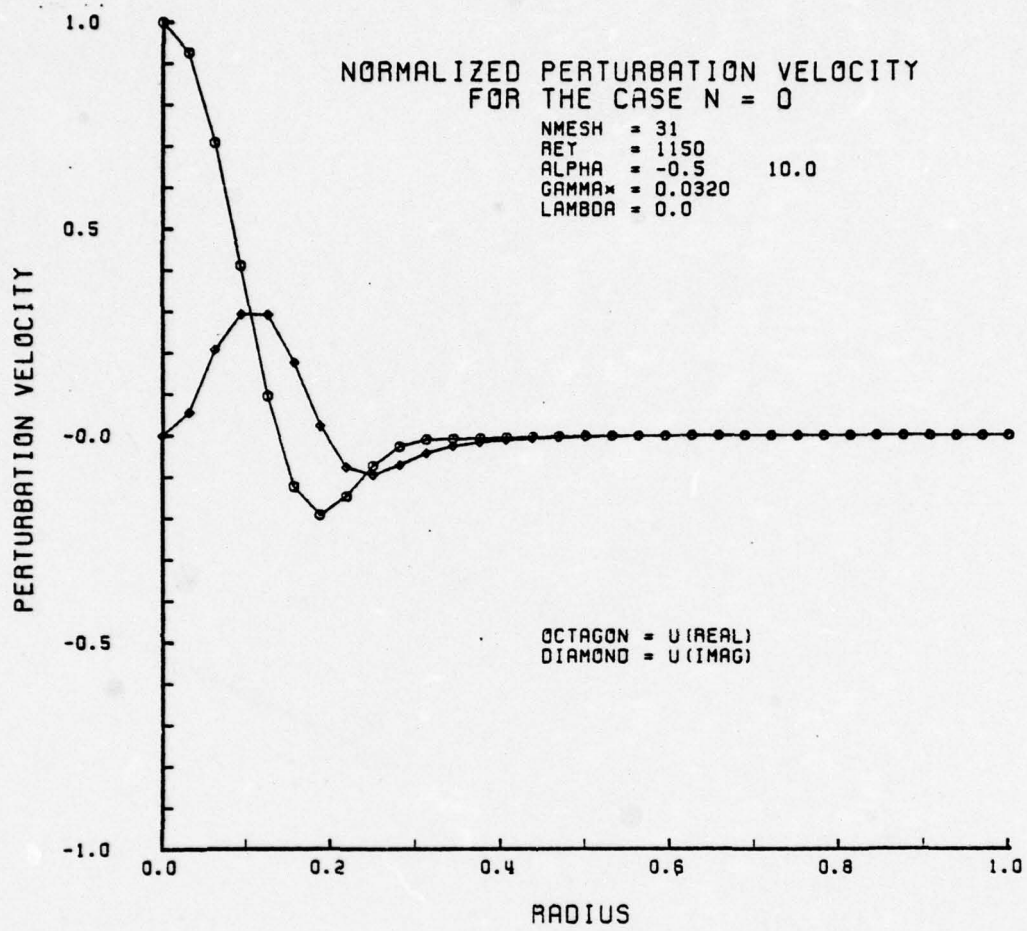


FIGURE 4-5

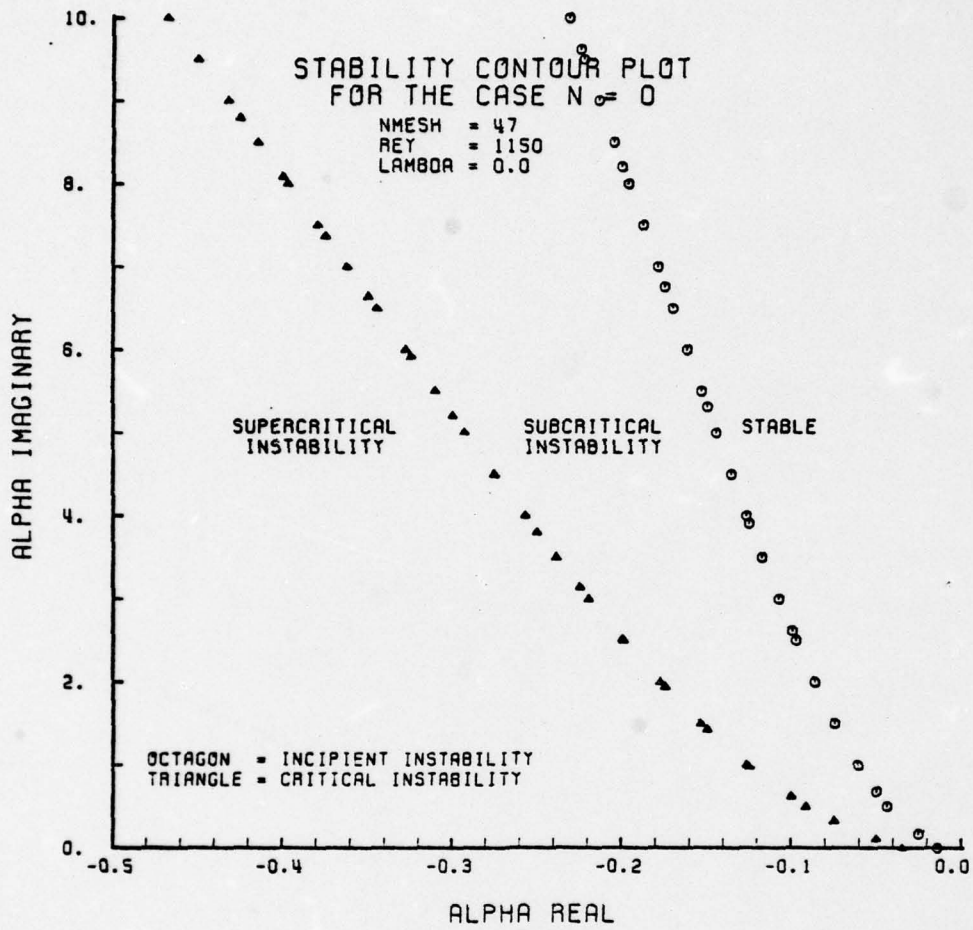


FIGURE 4-6

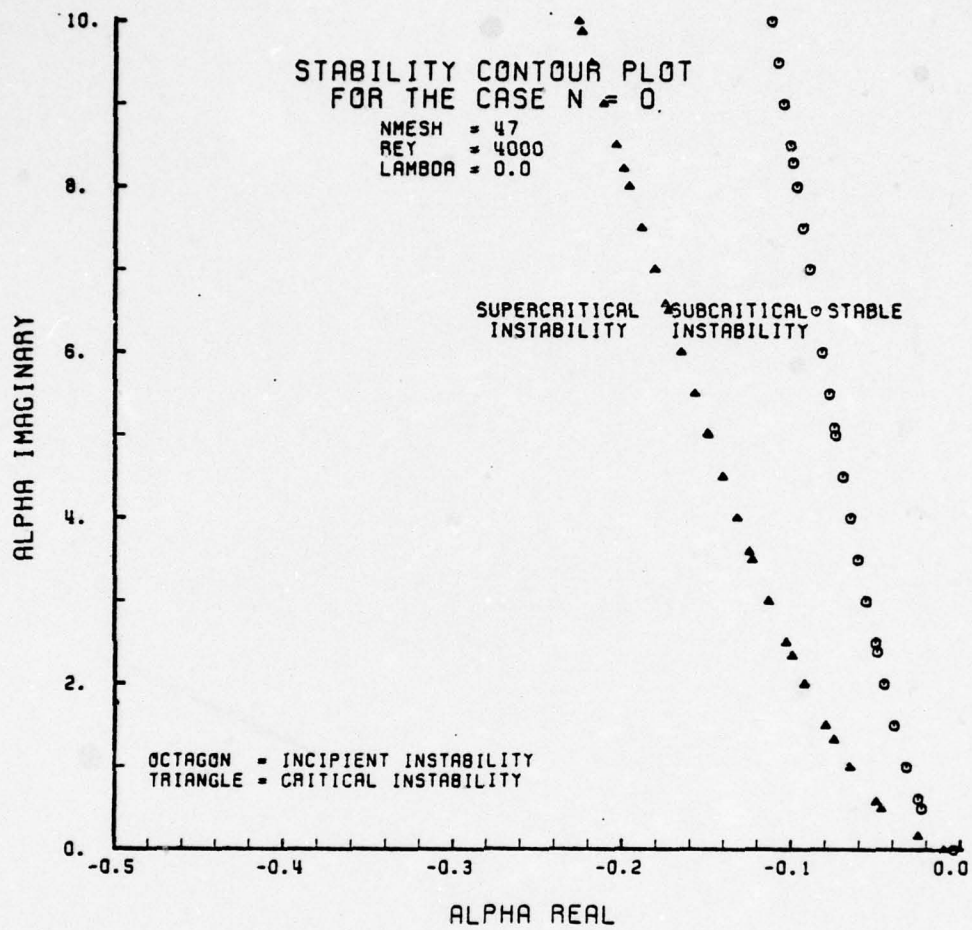


FIGURE 4-7



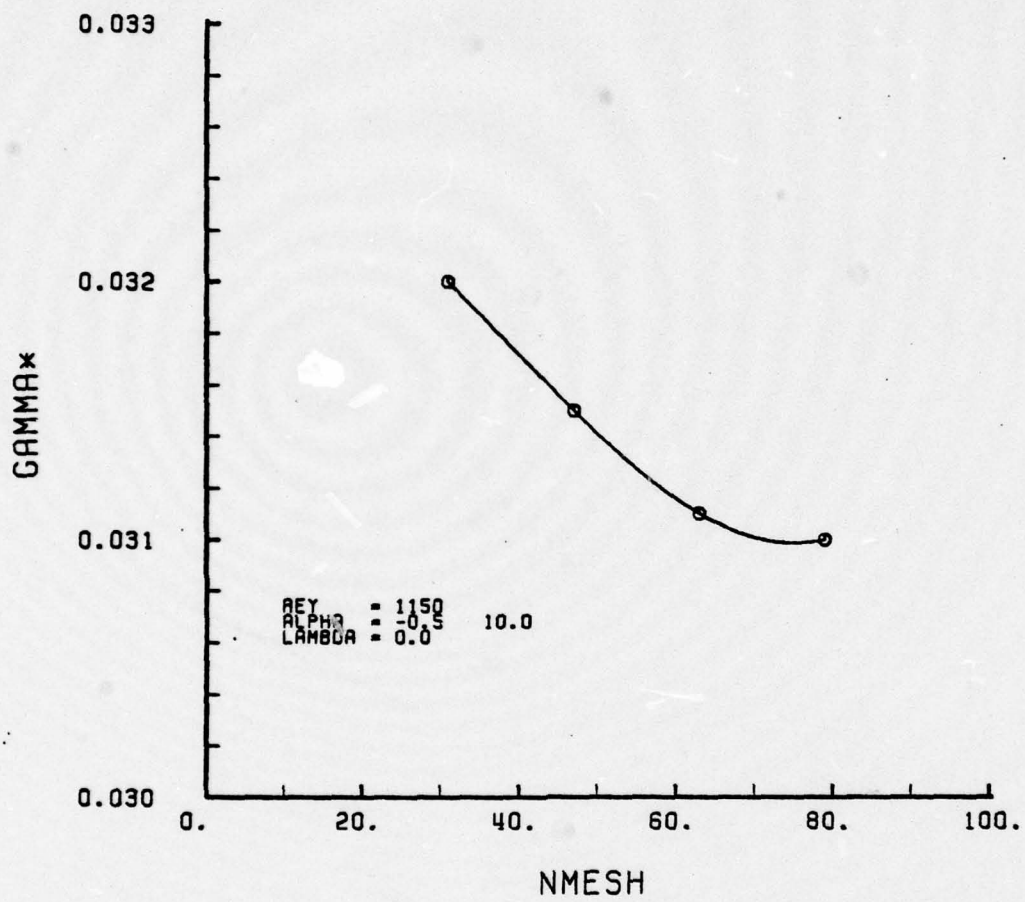


FIGURE 4-8.  $\gamma^*$  Versus Number of Mesh Points, N.

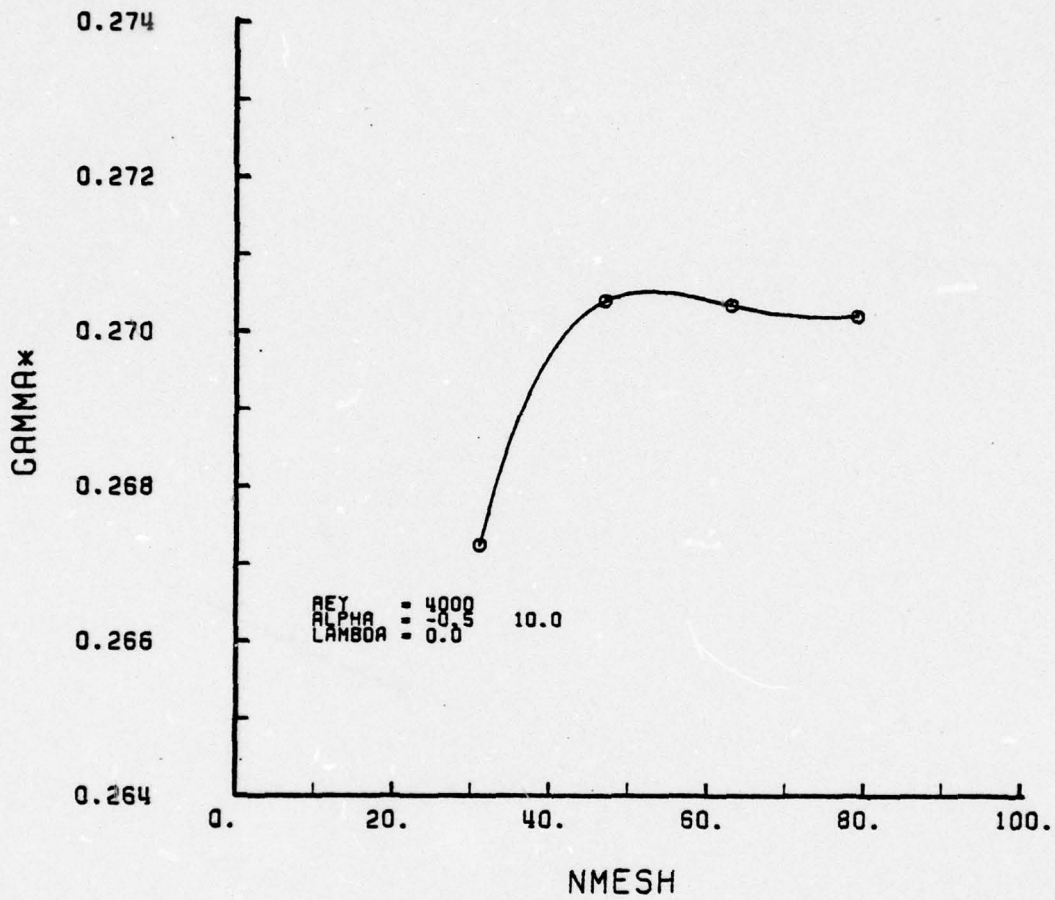


FIGURE 4-9.  $\gamma^*$  Versus Number of Mesh Points, N.

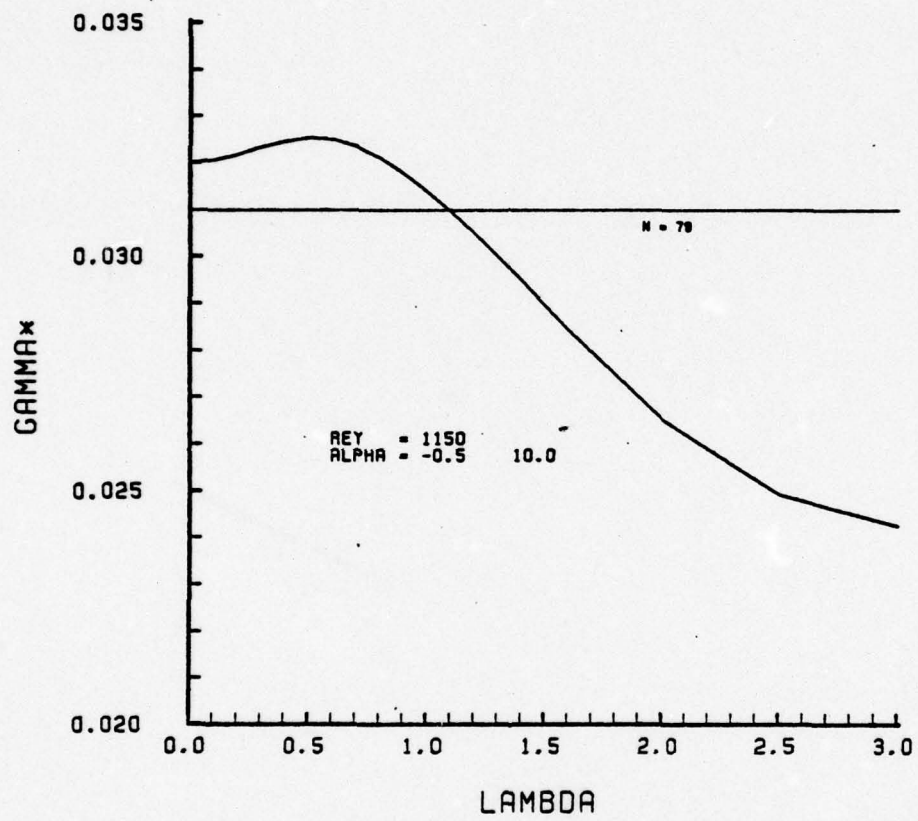


FIGURE 4-10.  $\gamma^*$  Versus Mesh Parameter, Lambda



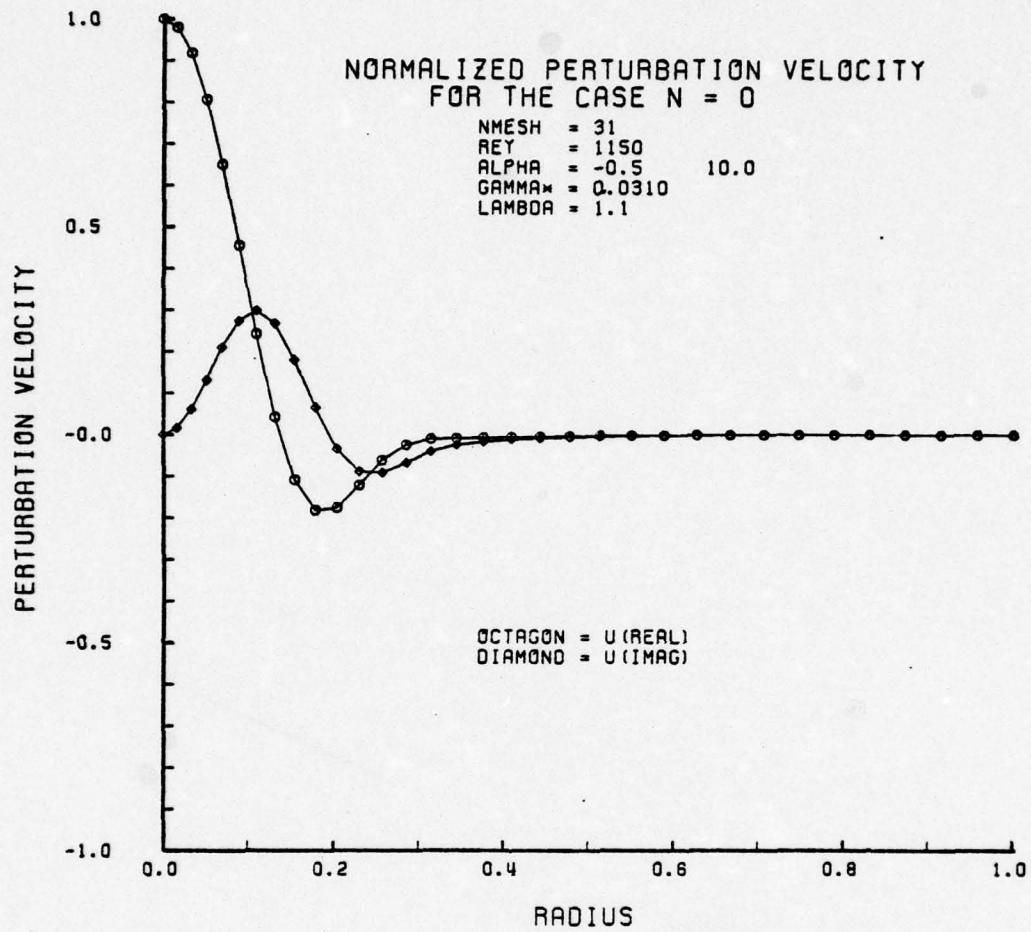


FIGURE 4-11

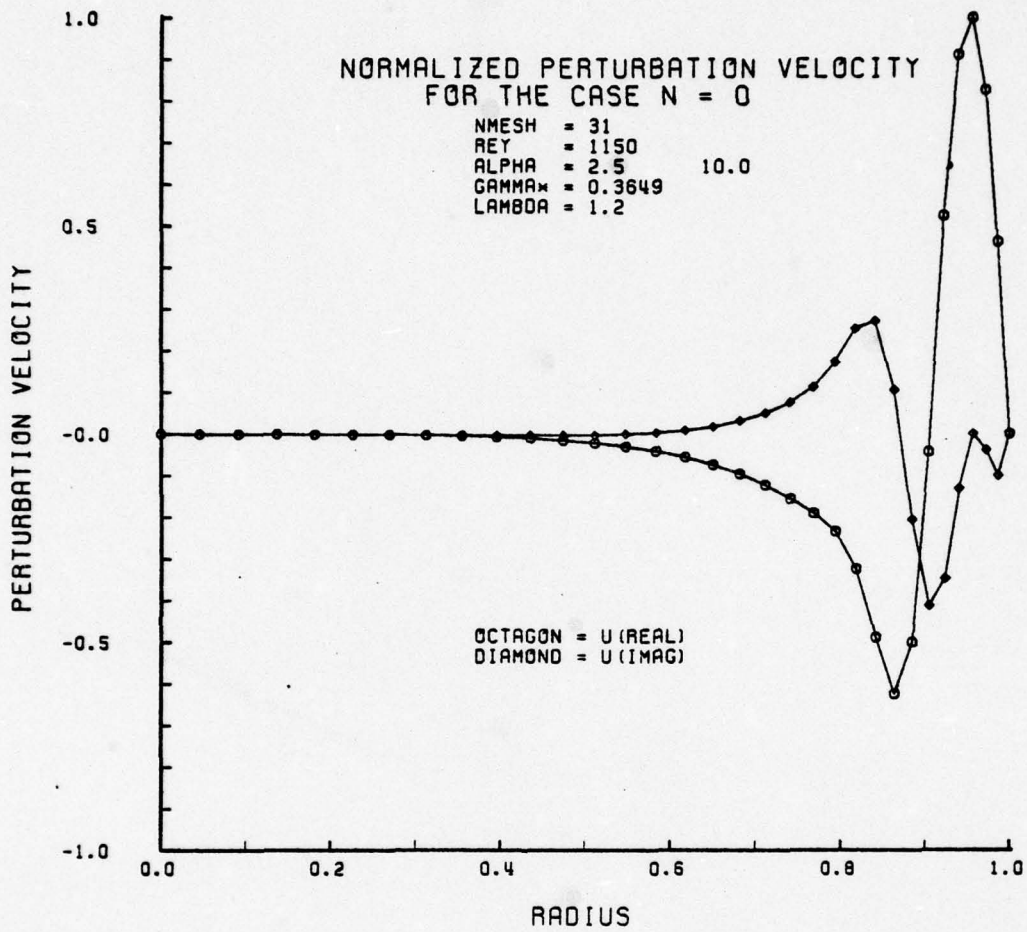


FIGURE 4-12

## V. CONCLUSIONS AND RECOMMENDATIONS

The implementation of the newly developed boundary conditions of Gawain [9] has permitted a stable, numerical solution to the linearized vorticity transport equation. The results of the numerical solution are presented in Section IV and show that the stability of pipe Poiseuille flow is governed by the three parameters,  $\alpha_R$ ,  $\alpha_I$  and  $R_e$ . In particular, both positive and negative values of  $\alpha_R$ , that is, streamwise growth and decay in space, if sufficiently large, produce unstable growth rates in time. This result is new and it is consistent with the known experimental fact that transition to turbulent flow depends not only on Reynolds number but also on the general character of the perturbations which exist in the flow.

The perturbation velocity plots of Section IV represent the first practical look at the function  $Q$ . These plots were valuable indicators for adequacy of mesh fineness, that is,  $N$ , changes in the nature of the function  $Q$  and effects of a nonuniform mesh.

No instabilities were discovered for purely sinusoidal perturbations ( $\alpha_R = 0$ ). This is consistent with the previous investigation of Ref. 11, but should not be assumed for investigations of other angular wave numbers, ( $n = 1, 2, 3, \dots$ ).



Adequate numerical accuracy was proven by demonstrating that the solution was virtually independent of the number of mesh points,  $N$ , and that it satisfied to a high degree an independent check of the governing differential equation. This procedure should also be carried out in future investigations prior to conducting full scale data runs.

This study suggests that similar, and perhaps even more rewarding results will be obtained for the higher angular wave numbers. Although lengthy, programming is straightforward if approached systematically. The general organization of the programs of Ref. 4 or Ref. 6 should be helpful in this task. It is recommended that the case for  $n = 1$  be undertaken as a follow-on to this study.

The nonuniform computational mesh was shown to be a powerful tool in the reduction of computational time. At the same time, however, the dependence of the mesh offset parameter,  $\lambda$ , on input conditions needs to be investigated further to realize the full potential of this technique.

APPENDIX A

DERIVATION OF VORTICITY TRANSPORT EQUATION COEFFICIENTS

From the change of variable introduced in Ref. 9, the function H for the case  $n = 0$  is expressed by

$$H = rQ \quad (A-1)$$

Taking derivatives

$$DH = rDQ + Q \quad (A-2)$$

$$D^2H = rD^2Q + 2DQ \quad (A-3)$$

$$D^3H = rD^3Q + 3D^2Q \quad (A-4)$$

$$D^4H = rD^4Q + 4D^3Q \quad (A-5)$$

Let the '\*' superscript denote element (2,2) of matrices (A1) through (A9) of Ref. 9. Since for  $n = 0$ , only the function H was investigated, equations (2-10) become

$$M_4^* D^4H + M_3^* D^3H + M_2^* D^2H + M_1^* DH + M_0^* H - \gamma [N_2^* D^2H + N_1^* DH + N_0^* H] = 0 \quad (A-6)$$

Substituting for H, equation (A-6) becomes

$$\begin{aligned}
& M_4^* \{rD^4Q + 4D^3Q\} + M_3^* \{rD^3Q + 3D^2Q\} + M_2^* \{rD^2Q + 2DQ\} \\
& + M_1^* \{rDQ + Q\} + M_0^* \{rQ\} - \gamma [N_2^* \{rD^2Q + 2DQ\} \\
& + N_1^* \{rDQ + Q\} + N_0^* \{rQ\}] = 0
\end{aligned} \tag{A-7}$$

Before proceeding further, it should be noted that the Ref. 9 matrices from which the coefficients for equation (A-7) were taken were obtained from matrices (2-10) through (2-17) of Ref. 6 by means of the following substitutions:

$$U = 2(1 - r^2) \tag{A-8}$$

$$t = \alpha^2 \frac{n_2}{r^2} \tag{A-9}$$

$$T = \alpha U - \frac{1}{R_e} \left( \alpha^2 - \frac{n_2}{r^2} \right) \tag{A-10}$$

Defining the new coefficients for equation (A-7) as  $M_0$  through  $M_4$  and  $N_0$  through  $N_2$

$$M_4 = rM_4^* = -\frac{r}{R_e} \tag{A-11}$$

$$M_3 = 4M_4^* + rM_3^* = -\frac{6}{R_e} \tag{A-12}$$

$$M_2 = 3M_3^* + rM_2^* = r\alpha U - \frac{1}{R_e} \left\{ \frac{3}{r} + 2\alpha^2 r \right\} \tag{A-13}$$

$$M_1 = 2M_2^* + rM_1^* = 3\alpha U + \frac{3}{R_e} \left\{ \frac{1}{r^2} - 2\alpha^2 \right\} \tag{A-14}$$

$$M_0 = M_1^* + rM_0^* = r\alpha^3 U - \frac{\alpha^4 r}{R_e} \tag{A-15}$$



$$N_2 = rN_2^* = -r \quad (\text{A-16})$$

$$N_1 = 2N_2^* + rN_1^* = -3 \quad (\text{A-17})$$

$$N_0 = N_1^* + rN_0^* = -\alpha^2 r \quad (\text{A-18})$$

Upon making use of the foregoing substitutions, the governing relation can finally be reduced to the form previously shown in equation (3-1).

APPENDIX B  
FINITE DIFFERENCE EQUATIONS

Improved finite difference equations for the boundaries were obtained by not using the virtual point method of Ref. 4 and Ref. 6 and deriving the forms directly from the boundary conditions of Appendix A. The equations thus formed are also of consistent order truncation error, significantly improving the accuracy of the solution [Ref. 8].

Because of a peculiarity in the form of the consistent second order truncation error equations at the axis, a singularity resulted for  $\alpha$  equal to zero. Consistent third order truncation error equations eliminated this problem.

From Appendix A, the axis boundary conditions are

$$DQ(0) = 0 \quad \text{and} \quad D^3Q(0) = 0 \quad (\text{B-1})$$

Representing  $Q$  by a power series and applying equations (B-1) yields

$$\begin{aligned} Q(r) = & Q(0) + D^2Q(0)\frac{r^2}{2!} + D^4Q(0)\frac{r^4}{4!} + D^5Q(0)\frac{r^5}{5!} \\ & + D^6Q(0)\frac{r^6}{6!} + \dots \end{aligned} \quad (\text{B-2})$$

Using five mesh points at  $r = \delta, 2\delta, 3\delta, 4\delta$  and  $5\delta$  results in the matrix

$$\begin{array}{l}
 \left. \begin{array}{l} Q_1 \\ Q_2 \\ Q_3 \\ Q_4 \\ Q_5 \end{array} \right\} = \begin{bmatrix} 1 & \frac{1}{2} & \frac{1}{24} & \frac{1}{120} & \frac{1}{720} \\ 1 & 2 & \frac{16}{24} & \frac{32}{120} & \frac{64}{720} \\ 1 & \frac{9}{2} & \frac{81}{24} & \frac{243}{120} & \frac{729}{720} \\ 1 & 8 & \frac{256}{24} & \frac{1024}{120} & \frac{4096}{720} \\ 1 & \frac{25}{2} & \frac{625}{24} & \frac{3125}{120} & \frac{15625}{720} \end{bmatrix} \begin{array}{l} Q(0) \\ \delta^2 D^2 Q(0) \\ \delta^4 D^4 Q(0) \\ \delta^5 D^5 Q(0) \\ \delta^6 D^6 Q(0) \end{array} + O\delta^7
 \end{array}
 \tag{B-3}$$

Differentiating equation (B-2) and substituting  $r = \delta$  gives  
(in matrix form)

$$\begin{array}{l}
 \left. \begin{array}{l} Q(\delta) \\ \delta D Q(\delta) \\ \delta^2 D^2 Q(\delta) \\ \delta^3 D^3 Q(\delta) \\ \delta^4 D^4 Q(\delta) \end{array} \right\} = \begin{bmatrix} 1 & \frac{1}{2!} & \frac{1}{4!} & \frac{1}{5!} & \frac{1}{6!} \\ 0 & 1 & \frac{1}{3!} & \frac{1}{4!} & \frac{1}{5!} \\ 0 & 1 & \frac{1}{2!} & \frac{1}{3!} & \frac{1}{4!} \\ 0 & 0 & 1 & \frac{1}{2!} & \frac{1}{3!} \\ 0 & 0 & 1 & 1 & \frac{1}{2!} \end{bmatrix} \begin{array}{l} Q(0) \\ \delta^2 D^2 Q(0) \\ \delta^4 D^4 Q(0) \\ \delta^5 D^5 Q(0) \\ \delta^6 D^6 Q(0) \end{array} + O\delta^7
 \end{array}
 \tag{B-4}$$

Let [A] and [B] denote the coefficient matrices of equations (B-3) and (B-4) respectively. The values of  $Q(0)$ ,  $\delta^2 D^2 Q(0)$ ,  $\delta^4 D^4 Q(0)$ ,  $\delta^5 D^5 Q(0)$  and  $\delta^6 D^6 Q(0)$  may be solved for by



$$\begin{bmatrix} Q(0) \\ \delta^2 D^2 Q(0) \\ \delta^4 D^4 Q(0) \\ \delta^5 D^5 Q(0) \\ \delta^6 D^6 Q(0) \end{bmatrix} = [A]^{-1} \begin{bmatrix} Q_1 \\ Q_2 \\ Q_3 \\ Q_4 \\ Q_5 \end{bmatrix} + O\delta^7 \quad (\text{B-5})$$

Putting equation (B-5) into equation (B-4),

$$\begin{bmatrix} Q(\delta) \\ \delta D Q(\delta) \\ \delta^2 D^2 Q(\delta) \\ \delta^3 D^3 Q(\delta) \\ \delta^4 D^4 Q(\delta) \end{bmatrix} = [B][A]^{-1} \begin{bmatrix} Q_1 \\ Q_2 \\ Q_3 \\ Q_4 \\ Q_5 \end{bmatrix} + O\delta^7 \quad (\text{B-6})$$

The last line of this set of equations gives

$$\begin{aligned}
 D^4 Q(\delta) &= \frac{1}{\delta^4} (-.911564626Q_1 + 2.750242955Q_2 - 3.043731779Q_3 \\
 &+ 1.42468416Q_4 - .219630709Q_5) + O\delta^3 \quad (\text{B-7})
 \end{aligned}$$

To solve for  $D^3 Q(\delta)$ , the rightmost column and bottom row are eliminated from matrices [A] and [B] then these new matrices are inserted into equations (B-5) and (B-6).

The bottom line of equation (B-6) will now give the expression for  $D^3Q(\delta)$  with a consistent third order truncation error.  $D^2Q(\delta)$  and  $DQ(\delta)$  were solved for in a similar manner.

$$D^3Q(\delta) = \frac{1}{\delta^3}(1.825165563Q_1 - 3.250331126Q_2 + 1.660927152Q_3 - .235761589Q_4) + O\delta^3 \quad (B-8)$$

$$D^2Q = \frac{1}{\delta^2}(-\frac{35}{60}Q_1 + \frac{8}{15}Q_2 + \frac{1}{20}Q_3) + O\delta^3 \quad (B-9)$$

$$DQ = \frac{1}{\delta}(-\frac{2}{3}Q_1 + \frac{2}{3}Q_2) + O\delta^3 \quad (B-10)$$

Due to the complexity of the boundary conditions, it was decided that consistent third order truncation error equations should also be used at  $r = 2\delta$ . For this the [B] matrix only need be changed as equation (B-2) is unchanged at this station. The new matrix [B] is formed by differentiating equation (B-2) and making the substitution  $r = 2\delta$ . Proceeding as for  $r = \delta$  gives the following finite difference approximations

$$D^4Q(2\delta) = \frac{1}{\delta^4}(-3.10340136Q_1 + 6.903012634Q_2 - 5.342274053Q_3 + 1.66083577Q_4 - 0.123420797Q_5) + O\delta^3 \quad (B-11)$$

$$D^3Q(2\delta) = \frac{1}{\delta^3}(.868874172Q_1 - .937748345Q_2 - .254304636Q_3 + .323178808Q_4) + O\delta^3 \quad (B-12)$$

$$D^2Q(2\delta) = \frac{1}{\delta^2} \left( \frac{11}{12}Q_1 - \frac{28}{15}Q_2 + \frac{19}{20}Q_3 \right) + O\delta^3 \quad (B-13)$$

$$DQ(2\delta) = \frac{1}{\delta} \left( -\frac{4}{3}Q_1 + \frac{4}{3}Q_2 \right) + O\delta^3 \quad (B-14)$$

It should also be noted that the value of  $Q$  at  $r = 0$  may be solved for from the top line of equations (B-5)

$$Q(0) = (1.795918367Q_1 - 1.24781341Q_2 + .606413994Q_3 - .177842566Q_4 + .023323615Q_5) + O\delta^3 \quad (B-15)$$

The central difference equations given by Ref. 6 were already consistent second order truncation error equations as confirmed by Ref. 8 and were retained.

For the wall, the clamped end, consistent second order equations (5) through (8) of Table II, Ref. 8 were modified for the "right boundary" using the procedure given in Section 5 of that reference.

$$D^4Q(1-\delta) = \frac{1}{\delta^4} \left( -\frac{1}{4}Q_{N-3} + \frac{8}{3}Q_{N-2} - 9Q_{N-1} + 16Q_N \right) + O\delta^2 \quad (B-16)$$

$$D^3Q(1-\delta) = \frac{1}{\delta^3} \left( -\frac{1}{3}Q_{N-2} + 3Q_N \right) + O\delta^2 \quad (B-17)$$

$$D^2Q(1-\delta) = \frac{1}{\delta^2} \left( Q_{N-1} - 2Q_N \right) + O\delta^2 \quad (B-18)$$

$$DQ(1-\delta) = \frac{1}{\delta} \left( -\frac{1}{2}Q_{N-1} \right) + O\delta^2 \quad (B-19)$$



Since the wall finite difference approximations were of only second order truncation error, the approximations for  $DQ$  through  $D^4Q$  at  $r = 1-2\delta$  were obtained directly from the central difference equations with  $Q(1) = 0$ .

$$D^4Q(1-2\delta) = \frac{1}{\delta^4}(Q_{N-3} - 4Q_{N-2} + 6Q_{N-1} - 4Q_N) + O\delta^2 \quad (B-20)$$

$$D^3Q(1-2\delta) = \frac{1}{\delta^3}(-\frac{1}{2}Q_{N-3} + Q_{N-2} - Q_N) + O\delta^2 \quad (B-21)$$

$$D^2Q(1-2\delta) = \frac{1}{\delta^2}(Q_{N-2} - 2Q_{N-1} + Q_N) + O\delta^2 \quad (B-22)$$

$$DQ(1-2\delta) = \frac{1}{\delta}(-\frac{1}{2}Q_{N-2} + \frac{1}{2}Q_N) + O\delta^2 \quad (B-23)$$

APPENDIX C

NONUNIFORM MESH

To control the distribution of a fixed number of mesh points, a change of the independent variable from  $r$  to  $\eta$  was performed.

$$Q = Q(\eta) \quad (C-1)$$

$$r = r(\eta) \quad (C-2)$$

The derivative with respect to  $r$  becomes

$$D = (D^*r)^{-1}D^* \quad (C-3)$$

where

$$D^* = \frac{d}{d\eta} \quad \text{and} \quad D = \frac{d}{dr} \quad (C-4)$$

$DQ, D^2Q \dots$  can now be expressed in terms of the new independent variable,  $\eta$ .

$$DQ = (D^*r)^{-1}D^*Q \quad (C-5)$$

$$\begin{aligned} D^2Q &= D(DQ) = (D^*r)^{-1}D^*(DQ) \\ &= (D^*r)^{-2}D^{*2}Q - (D^*r)^{-3}(D^{*2}r)D^*Q \end{aligned} \quad (C-6)$$

$$\begin{aligned}
D^3Q &= D(D^2Q) = (D^*R)^{-1}D^*(D^2Q) \\
&= (D^*r)^{-3}D^{*3}Q - 3(D^*r)^{-4}(D^{*2}r)D^{*2}Q \\
&\quad - [(D^*r)^{-4}(D^{*3}r) - 3(D^*r)^{-5}(D^{*2}r)^2]DQ \quad (C-7)
\end{aligned}$$

$$\begin{aligned}
D^4Q &= D(D^3Q) = (D^*r)^{-1}D^*(D^3Q) \\
&= (D^*r)^{-4}D^{*4}Q - 6(D^*r)^{-5}(D^{*2}r)D^{*3}Q \\
&\quad + [15(D^*r)^{-6}(D^{*2}r) - 4(D^*r)^{-5}(D^{*3}r)]D^{*2}Q \\
&\quad - [15(D^*r)^{-7}(D^{*2}r)^3 - 10(D^*r)^{-6}(D^{*2}r)(D^{*3}r) \\
&\quad\quad + (D^*r)^{-5}(D^{*4}r)]DQ \quad (C-8)
\end{aligned}$$

The derivatives of Q with respect to r can now be written

$$DQ = f_{11}D^*Q \quad (C-9)$$

$$D^2Q = f_{22}D^{*2}Q + f_{21}D^*Q \quad (C-10)$$

$$D^3Q = f_{33}D^{*3}Q + f_{32}D^{*2}Q + f_{31}D^*Q \quad (C-11)$$

$$D^4Q = f_{44}D^{*4}Q + f_{43}D^{*3}Q + f_{42}D^{*2}Q + f_{41}D^*Q \quad (C-12)$$

where

$$f_{11} = (D^*r)^{-1} \quad (C-13)$$



$$f_{22} = (D^* r)^{-2} \quad (C-14)$$

$$f_{21} = -(D^* r)^{-3} (D^{*2} r) \quad (C-15)$$

$$f_{33} = (D^* r)^{-3} \quad (C-16)$$

$$f_{32} = -3(D^* r)^{-4} (D^{*2} r) \quad (C-17)$$

$$f_{31} = 3(D^* r)^{-5} (D^{*2} r)^2 - (D^* r)^{-4} (D^{*3} r) \quad (C-18)$$

$$f_{44} = (D^* r)^{-4} \quad (C-19)$$

$$f_{43} = -6(D^* r)^{-5} (D^{*2} r) \quad (C-20)$$

$$f_{42} = 15(D^* r)^{-6} (D^{*2} r)^2 - 4(D^* r)^{-5} (D^{*3} r) \quad (C-21)$$

$$f_{41} = -15(D^* r)^{-7} (D^{*2} r)^3 + 10(D^* r)^{-6} (D^{*2} r) (D^{*3} r) \\ - (D^* r)^{-5} (D^{*4} r) \quad (C-22)$$

Substituting equations (C-9) through (C-12) into the vorticity transport equation (A-6) yields

$$M_4^* D^{*4} Q + M_3^* D^{*3} Q + M_2^* D^{*2} Q + M_1^* D^* Q + M_0^* Q \\ - \gamma [N_2^* D^{*2} Q + N_1^* D^* Q + N_0^* Q] = 0 \quad (C-23)$$

where

$$M_4^* = M_4 f_{44} \quad (C-24)$$

$$M_3^* = M_4 f_{43} + M_3 f_{33} \quad (C-25)$$

$$M_2^* = M_4 f_{42} + M_3 f_{32} + M_2 f_{22} \quad (C-26)$$

$$M_1^* = M_4 f_{41} + M_3 f_{31} + M_2 f_{21} \quad (C-27)$$

$$M_0^* = M_0 \quad (C-28)$$

$$N_2^* = N_2 f_{22} \quad (C-29)$$

$$N_1^* = N_2 f_{21} + N_1 f_{11} \quad (C-30)$$

$$N_0^* = N \quad (C-31)$$

In order to concentrate the mesh points at the axis, the function

$$r = 1 - C \tanh \lambda(1-\eta) \quad (C-32)$$

was chosen where  $\lambda$  is a parameter controlling the degree of concentration of mesh points near the axis. Equation (C-32) must satisfy the two conditions

$$r = 0 \quad \text{at} \quad \eta = 0 \quad (C-33)$$

and

$$r = 1 \quad \text{at} \quad \eta = 1 .$$

Substituting equation (C-33) into (C-32) gives

$$C = 1/\tanh \lambda . \quad (C-35)$$

Computing derivatives

$$D^* r = C\lambda/\cosh^2 \lambda(1-\eta) \quad (C-36)$$

$$D^{*2} r = 2C\lambda^2 [\tanh \lambda(1-\eta)/\cosh^2 \lambda(1-\eta)] \quad (C-37)$$

$$D^{*3} r = -2C\lambda^3 \{ [1-2\sinh^2 \lambda(1-\eta)]/\cosh^4 \lambda(1-\eta) \} \quad (C-38)$$

$$D^{*4} r = 8C\lambda^4 [\tanh^3 \lambda(1-\eta)/\cosh^2 \lambda(1-\eta)] \quad (C-39)$$

To shift the mesh point concentration to the wall, the function

$$r = C \tanh \lambda \eta \quad (C-40)$$

was selected. Satisfying equations (C-33) and (C-34) for this equation also gives equation (C-35). The derivatives



of (C-40) are given by equations (C-36) through (C-39) if  $\eta$  is substituted for all occurrences of  $(1-\eta)$  and the signs of equations (C-37) and (C-39) are reversed. Figures C-1 and C-2 show equations (C-32) and (C-40) for four selected values of the parameter  $\lambda$ .

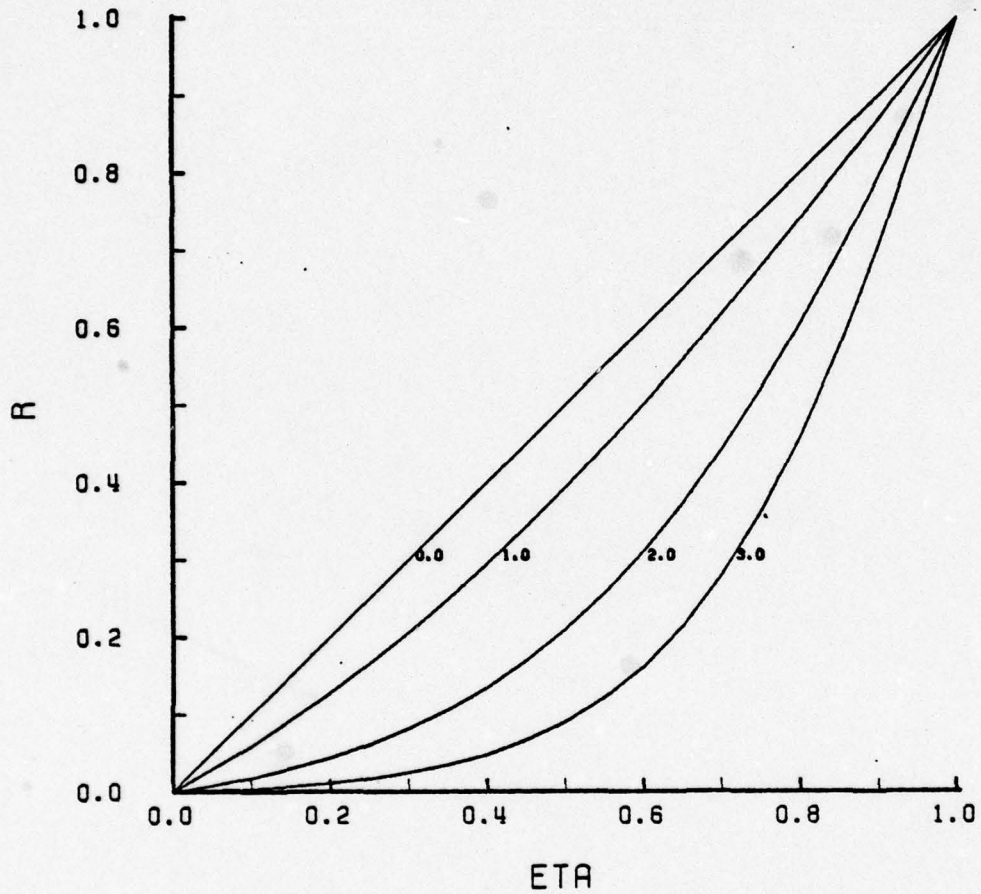


FIGURE C-1. R Versus  $\eta$  for Four Selected Values of Lambda - Axis Offset

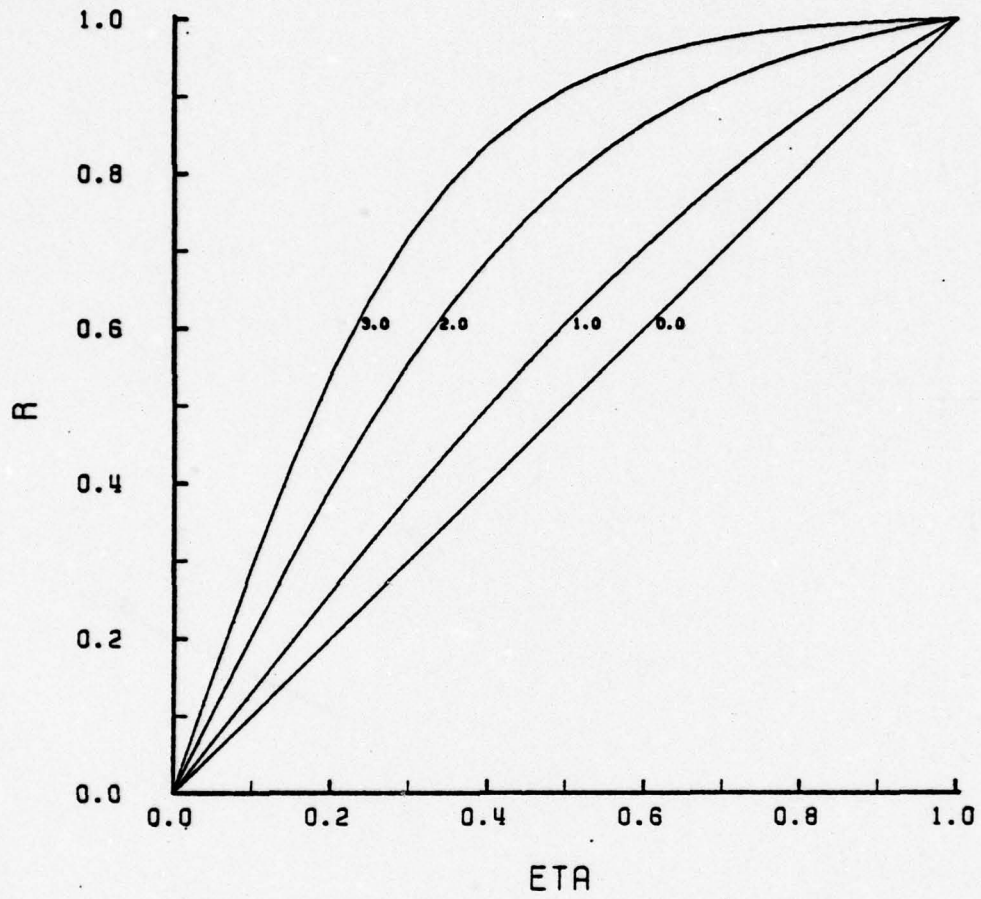


FIGURE C-2. R Versus  $\eta$  for Four Selected Values of Lambda - Wall Offset



APPENDIX D

DERIVATION OF PERTURBATION VELOCITIES

From Ref. 4, Appendix E, equations E-6 through E-8:

$$\begin{Bmatrix} u(r) \\ v(r) \\ w(r) \end{Bmatrix} = [A]\bar{W} + [B]D\bar{W} \quad (D-1)$$

$$= \begin{bmatrix} 0 & -\frac{\beta}{r} & \frac{1}{r} \\ \frac{\beta}{r} & 0 & -\alpha \\ 0 & \alpha & 0 \end{bmatrix} \begin{Bmatrix} F \\ G \\ H \end{Bmatrix} + \begin{bmatrix} 0 & 0 & 1 \\ 0 & 0 & 0 \\ -1 & 0 & 0 \end{bmatrix} \begin{Bmatrix} DF \\ DG \\ DG \end{Bmatrix} \quad (D-2)$$

For this case  $\beta = \eta_i = 0$  and  $F = DF = 0$ . Restricting the investigation to the function  $H$  for the reason expressed in Section I and solving for  $u(r)$  gives

$$u(r) = \frac{H}{r} + DH \quad (D-3)$$

Performing the change of variable

$$H = rQ \quad (D-4)$$

$$DH = Q + rDQ \quad (D-5)$$

$$u(r) = \frac{rQ}{r} + (Q + rDQ) = 2Q + rDQ \quad (D-6)$$

In order to implement this derivation in a numerical analysis, equation (D-6) was rewritten as

$$u_i = 2Q_i + r_i DQ_i \quad (D-7)$$

Performing the change of independent variable (Appendix C) to accommodate a nonuniform mesh

$$Q_i = Q(\eta_i) \quad (D-8)$$

$$r_i = r(\eta_i) \quad (D-9)$$

$$DQ_i = (D^* r_i)^{-1} D^* Q(\eta_i) \quad (D-10)$$

Substituting equations (D-8), (D-9) and (D-10) into equation (D-7) gives

$$u_i = 2Q(\eta_i) + r(\eta_i) (D^* r_i)^{-1} D^* Q(\eta_i) \quad (D-11)$$

For the axis offset nonuniform mesh,  $r(\eta)$  is given by equation (C-32) and  $(D^* r)$  by equation (C-36). Substituting into equation (D-11) using equation (C-35) results in

$$\begin{aligned}
u_i &= 2Q(\eta_i) + \left\{ 1 - \frac{\tanh[\lambda(1-\eta_i)]}{\tanh \lambda} \right\} \left\{ \frac{\cosh^2[\lambda(1-\eta_i)]}{C\lambda} \right\} D^* Q(\eta_i) \\
&= 2Q(\eta_i) + \left\{ 1 - \frac{\tanh[\lambda(1-\eta_i)]}{\tanh \lambda} \frac{\tanh \lambda \cosh^2[\lambda(1-\eta_i)]}{\lambda} \right\} D^* Q(\eta_i)
\end{aligned}
\tag{D-12}$$

For the wall offset mesh, equation (C-40) is substituted for equation (C-32) and all occurrences of the term  $1-\eta_i$  are replaced by the term  $\eta_i$ .

The value of  $u$  at the axis ( $u_0$ ) and at the wall ( $u_{N+1}$ ) were solved for by using the boundary conditions specified in Ref. 9, namely

$$Q(1) = 0 \tag{D-13}$$

$$DQ(1) = 0 \tag{D-14}$$

$$DQ(0) = 0 \tag{D-15}$$

$$D^3Q(0) = 0 \tag{D-16}$$

From equations (D-13) and (D-14), using equation (D-7) it is obvious that

$$u_{N+1} = 0 \tag{D-17}$$

and from equations (D-15) and (D-7), it is similarly found that



$$u_0 = 2Q(0) ,$$

(D-18)

where the finite difference approximation for  $Q(0)$  is given by equation (B-15).

```

10 PIP0
20 PIP0
30 PIP0
40 PIP0
50 PIP0
60 PIP0
70 PIP0
80 PIP0
90 PIP0
100 PIP0
110 PIP0
120 PIP0
130 PIP0
140 PIP0
150 PIP0
160 PIP0
170 PIP0
180 PIP0
190 PIP0
200 PIP0
210 PIP0
220 PIP0
230 PIP0
240 PIP0
250 PIP0
260 PIP0
270 PIP0
280 PIP0
290 PIP0
300 PIP0
310 PIP0
320 PIP0
330 PIP0
340 PIP0
350 PIP0
360 PIP0
370 PIP0
380 PIP0
390 PIP0
400 PIP0
410 PIP0
420 PIP0
430 PIP0
440 PIP0
450 PIP0
460 PIP0
470 PIP0
480 PIP0

.....
PROGRAM PIPEO (CP/CMS VERSION)
PROGRAM TO INVESTIGATE FLOW STABILITY AND CHARACTERISTICS
FOR THE 3-D CYLINDRICAL FLOW PROBLEM
NI=0; FOR THE FUNCTION Q
.....
TO OBTAIN A FLOW CHART OF THIS PROGRAM, CONSULT NAVAL
POSTGRADUATE SCHOOL TECHNICAL NOTE TN 0141-25, "USER'S
GUIDE TO THE PROGRAMMING AIDS LIBRARY", UNDER PROGRAM
FLOWCH.
.....
-----
IMPLICIT REAL*(A-H,O-Z)
DIMENSION A(PLT(20), REYPLT(20))
COMPLEX *16A,G
INTEGER *4CLOCK(6)
COMMON /COEFNT/ A,G,REY,DELR,AMDA
-----
INPUT DESIRED MODE NUMBER

WRITE (6,6) MODENO
READ (5,7) MODENO
CALL IXCLOCK (CLOCK)
WRITE (6,8) CLOCK(3),CLOCK(4)
IF (MODENO.EQ.2) GO TO 2

*****
CALCULATE STABILITY AT A POINT (MODENO = 1 & 3)

OUTPUTS DATA TO FILE FT02F001 COMPATIBLE WITH PROGRAM EIGFCN &
PRINTS THE VALUE OF THE LEAST STABLE EIGENVALUE (GAMMA*) AT THE
CONSOLE FOR EACH SET OF INPUT CONDITIONS. DATA IS ENTERED IN
'D' FORMAT WITH A ZERO EXPONENT (I.E., 1.D0, 1.5D0, ETC.).
RUN IS TERMINATED WHEN A NEGATIVE VALUE FOR REYNOLDS NUMBER IS
ENTERED. THE TYPE OF MESH OFFSET IS DETERMINED BY THE VALUE
OF LAMDA. IF LAMDA < -1.D-10, KSET IS SET EQUAL TO -1; IF
LAMDA > 1.D-10, KSET IS SET EQUAL TO 1; OTHERWISE KSET = 0.
IF MODENO = 1, THE EIGENVECTOR CORRESPONDING TO THE LEAST
STABLE EIGENVALUE (GAMMA*) WILL BE WRITTEN TO FILE FT02F001.
*****

```







```

10 FORMAT (1,5X,'INPUT ALPHA REAL')
11 FORMAT (D20.10)
12 FORMAT (1,5X,'INPUT ALPHA IMAG')
13 FORMAT (1,5X,'INPUT REYNOLDS NO')
14 FORMAT (1,5X,'STAB=',D20.10)
15 FORMAT (3I2,3D20.10)
16 FCRMAT (2D20.10)
17 FORMAT (1,5X,'REYNOLDS NUMBER = ',F8.1,/,10X,'AR',18X,
18 'AI',16X,
19 'STAB',//)
19 FORMAT (I2,3D20.10)
20 FORMAT (3E20.10)
21 FORMAT ('0','STOP TIME=',2A4)
END

```

```

.....SUBROUTINE STAB(AR,AI,GRMAX,KSET,MODENO).....
PURPOSE

```

RETURNS THE REAL PORTION OF THE LEAST STABLE EIGENVALUE  
FOR THE GIVEN INPUT CONDITIONS. THIS VALUE DETERMINES THE  
STABILITY OF THE FLOW.

USAGE

CALL STAB(AR,AI,GRMAX,KSET,MODENO)

DESCRIPTION OF PARAMETERS

AR - THE REAL PART OF ALPHA

AI - THE IMAGINARY PART OF ALPHA.

GRMAX - THE REAL PORTION OF THE LEAST STABLE EIGENVALUE.  
THIS VALUE IS RETURNED TO THE CALLING PROGRAM.

N - THE NUMBER OF INTERIOR MESH POINTS

KSET - AN INTEGER DENOTING THE TYPE OF MESH OFFSET  
USED.

KSET = -1

FOR WALL OFFSET

KSET = 0

FOR UNIFORM MESH

KSET = 1

FOR AXIS OFFSET

```

STAB 10
STAB 20
STAB 30
STAB 40
STAB 50
STAB 60
STAB 70
STAB 80
STAB 90
STAB 100
STAB 110
STAB 120
STAB 130
STAB 140
STAB 150
STAB 160
STAB 170
STAB 180
STAB 190
STAB 200
STAB 210
STAB 220
STAB 230
STAB 240
STAB 250
STAB 260
STAB 270
STAB 280

```

290 STAB  
 300 STAB  
 310 STAB  
 320 STAB  
 330 STAB  
 340 STAB  
 350 STAB  
 360 STAB  
 370 STAB  
 380 STAB  
 390 STAB  
 400 STAB  
 410 STAB  
 420 STAB  
 430 STAB  
 440 STAB  
 450 STAB  
 460 STAB  
 470 STAB  
 480 STAB  
 490 STAB  
 500 STAB  
 510 STAB  
 520 STAB  
 530 STAB  
 540 STAB  
 550 STAB  
 560 STAB  
 570 STAB  
 580 STAB  
 590 STAB  
 600 STAB  
 610 STAB  
 620 STAB  
 630 STAB  
 640 STAB  
 650 STAB  
 660 STAB  
 670 STAB  
 680 STAB  
 690 STAB  
 700 STAB  
 710 STAB  
 720 STAB  
 730 STAB  
 740 STAB  
 750 STAB  
 760 STAB

MODENO - AN INTEGER CONTROLLING THE OUTPUT OF  
 OF EIGENVECTORS TO FILE FT02FO01. IF MODENO  
 IS EQUAL TO +1, EIGENVECTORS ARE OUTPUT;  
 OTHERWISE OUTPUT IS INHIBITED.

OTHER ROUTINES NEEDED

MSET2,CDMTIN,MULM,DSPLIT,EHESCC,ELRH2C,EEALAC,EBBCKC

SUBROUTINE STAB (AR, AI, GRMAX, KSET, MODENO)  
 IMPLICIT REAL\*8 (A-H, O-Z)  
 COMPLEX \*16 A, G  
 COMPLEX \*16 CQM1E1, CQM2E1

NOTE --- CHANGE DIMENSIONS FROM HERE THROUGH 'N = ' FOR  
 NEW NMESH. CP/CMS MAX NMESH IS 79. LOGIN WITH 520K OF CORE.

REAL \*8 GR(79), GI(79), ZR(79,79), ZI(79,79), RADIUS(79), CVEC(79)  
 COMPLEX \*16 XMAT(79,79), YMAT(79,79), WV(79)  
 DIMENSION IVEC(79)  
 COMMON /COEFNT/ A, G, REY, DELR, AMDA  
 EXTERNAL CQM1E1, CQM2E1

A = DC MPLX(AR, AI)  
 MDIM = 79  
 N = 79

SET UP THE CENTRAL DIFFERENCE APPROXIMATION AT EACH POINT IN  
 THE MESH FOR THOSE TERMS IN THE VORTICITY TRANSPORT EQUATION  
 WHICH DO NOT CONTAIN GAMMA AS A FACTOR.

CALL MSET2 (XMAT, N, MDIM, CQM1E1, KSET)

SET UP THE CENTRAL DIFFERENCE APPROXIMATION AT EACH POINT IN  
 THE MESH FOR THOSE TERMS IN THE VORTICITY TRANSPORT EQUATION  
 WHICH CONTAIN GAMMA AS A FACTOR.

CALL MSET2 (YMAT, N, MDIM, CQM2E1, KSET)



```

CC      INVERT THE RESULTING ARRAY
CC      CALL CDMT IN (MDIM, YMAT, MDIM, IERR)
CC      PREMULTIPLY XMAT BY THE INVERSE OF YMAT TO CONVERT THE PROBLEM
CC      TO THE STANDARD FORM: (A - GAMMA*I)*X = 0
CC      CALL MULM (YMAT, XMAT, MDIM, MDIM, WV)
CC      SPLIT THE RESULTING ARRAY INTO REAL (XMAT) AND IMAGINARY (YMAT)
CC      PARTS.
CC      CALL DSPLIT (MDIM, MDIM, YMAT, XMAT, YMAT)
CC      CALCULATE THE EIGENVALUES (GAMMA) FOR THE EQUATION.
CC      CALL EBALAC (XMAT, YMAT, MDIM, MDIM, KBND, LBND, DVEC)
CC      CALL EHESSC (XMAT, YMAT, KBND, LBND, MDIM, MDIM, IVEC)
CC      CALL ELRH2C (XMAT, YMAT, KBND, LBND, MDIM, MDIM, GR, GI, ZR, ZI, IVEC, INERR,
1 IER)
CC      CALL EBBCKC (ZR, ZI, MDIM, MDIM, KBND, LBND, MDIM, DVEC)
CC      IF (INERR.NE.0) WRITE (6,4) INERR, IER
CC      DETERMINE LARGEST GAMMA REAL
CC      NEIG = 1
CC      GRMAX = -1.0D02
CC      DO 1 I=1, MDIM
CC      IF (GR(I).GT.GRMAX) NEIG=I
CC      IF (GR(I).GT.GRMAX) GRMAX=GR(I)
1 CONTINUE
CC      GRMAX = GRMAX+AR
CC      DEL10 = 1.00/DFLOAT(N+1)
CC      WRITE (2,5) MDIM, REY, AR, AI
CC      WRITE (2,7) AMDA, GR(NEIG), GI(NEIG)
CC      IF (MOD(ENO,NE.1)) GO TO 3
CC-----
770 STAB
780 STAB
790 STAB
800 STAB
810 STAB
820 STAB
830 STAB
840 STAB
850 STAB
860 STAB
870 STAB
880 STAB
890 STAB
900 STAB
910 STAB
920 STAB
930 STAB
940 STAB
950 STAB
960 STAB
970 STAB
980 STAB
990 STAB
1000 STAB
1010 STAB
1020 STAB
1030 STAB
1040 STAB
1050 STAB
1060 STAB
1070 STAB
1080 STAB
1090 STAB
1100 STAB
1110 STAB
1120 STAB
1130 STAB
1140 STAB
1150 STAB
1160 STAB
1170 STAB
1180 STAB
1190 STAB
1200 STAB
1210 STAB
1220 STAB
1230 STAB
1240 STAB

```

STABI250  
 STABI260  
 STABI270  
 STABI280  
 STABI290  
 STABI300  
 STABI310  
 STABI320  
 STABI330  
 STABI340  
 STABI350  
 STABI360  
 STABI370  
 STABI380  
 STABI390

```

C      DO 2 I=1,MDIM
      RADIUS(I)=DFLOAT(I)*DELL0
      WRITE (2,7) RADIUS(I),ZR(I,NEIG),ZI(I,NEIG)
C-----
C      2 CONTINUE
C      3 RETURN
C      4 FCRMAT ('0* * * ERROR NUMBER',I7,' ON EIGENVALUE',
      1 I7,' * * *'////)
      5 FORMAT (I2,3D20.10)
      6 FORMAT (I2)
      7 FORMAT (F15.7,2(1PD20.10))
      END
  
```

MSTI1  
 MSTI2  
 MSTI3  
 MSTI4  
 MSTI5  
 MSTI6  
 MSTI7  
 MSTI8  
 MSTI9  
 MSTI10  
 MSTI11  
 MSTI12  
 MSTI13  
 MSTI14  
 MSTI15  
 MSTI16  
 MSTI17  
 MSTI18  
 MSTI19  
 MSTI20  
 MSTI21  
 MSTI22  
 MSTI23  
 MSTI24  
 MSTI25  
 MSTI26  
 MSTI27  
 MSTI28  
 MSTI29  
 MSTI30  
 MSTI31

```

C-----
C      .....SUBROUTINE MSET1(ETA,CQM1,CQM2,KSET).....
C      PURPOSE
C      MSET1 GENERATES THE COEFFICIENTS FOR THE FINITE DIFFERENCE
C      APPROXIMATION OF THE COMPONENT Q.
C      CALL MSET1(ETA,CQM1,CQM2,KSET)
C      DESCRIPTION OF PARAMETERS
C      ETA -- INDEPENDENT VARIABLE REPLACING R
C            IN NONUNIFORM MESH.
C      CQM1 -- COEFFICIENTS OF Q AND ITS DERIVATIVES
C            IN THE FINITE DIFFERENCE APPROXIMATION OF
C            THE NON-GAMMA TERMS. CQM1(1) IS THE COEF-
C            FICIENT FOR Q, CQM1(2) IS THE COEF-
C            FICIENT FOR DQ, AND SO ON TO CQM1(5).
C      CQM2 -- SAME AS CQM1 EXCEPT GAMMA TERMS AND
C            DIMENSIONED 3 INSTEAD OF 5.
C      KSET -- MESH OFFSET PARAMETER AS DESCRIBED FOR
C            SUBROUTINE STAB.
C      OTHER ROUTINES NEEDED
C      NONE
  
```







```

MST2 100
MST2 110
MST2 120
MST2 130
MST2 140
MST2 150
MST2 160
MST2 170
MST2 180
MST2 190
MST2 200
MST2 210
MST2 220
MST2 230
MST2 240
MST2 250
MST2 260
MST2 270
MST2 280
MST2 290
MST2 300
MST2 310
MST2 320
MST2 330
MST2 340
MST2 350
MST2 360
MST2 370
MST2 380
MST2 390
MST2 400
MST2 410
MST2 420
MST2 430
MST2 440
MST2 450
MST2 460
MST2 470
MST2 480
MST2 490
MST2 500
MST2 510
MST2 520
MST2 530
MST2 540
MST2 550
MST2 560
MST2 570

```

```

CALL MSET2(X,N,MDIM,CFMAT,KSET)

```

```

DESCRIPTION OF PARAMETERS

```

```

X - THE NAME OF THE ARRAY BEING GENERATED. MUST BE DIMENSIONED
    IN THE CALLING PROGRAM

```

```

N- THE ROW DIMENSION OF THE MATRIX X. MUST BE .GE. N.

```

```

MDIM - THE COLUMN DIMENSION OF THE MATRIX X. MUST BE .GE. N.

```

```

CFMAT - THE NAME OF A FUNCTION SUBPROGRAM WITH 4 PARAMETERS,
        JSTA, K, CQM1 & CQM2. CFMAT MUST BE DECLARED
        EXTERNAL IN THE CALLING PROGRAM.

```

```

THE FOLLOWING IS OUTPUT BY MSET2

```

```

X - THE N BY N MATRIX INTO WHICH THE COEFFICIENTS OF THE CENTRAL
    DIFFERENCING ARE PUT.

```

```

OTHER ROUTINES NEEDED

```

```

FUNCTION SUBPROGRAM NAME PASSED IN THE CALLING PARAMETER 'CFMAT'
AND MSET1.

```

```

SUBROUTINE MSET2 (X,N,MDIM,CFMAT,KSET)
REAL *8REY,R,DEL,DFLOAT,AMDA,ETA
COMPLEX *16X(MDIM,MDIM),CQM1(5),CQM2(3)
COMPLEX *16A,G
COMPLEX *16CFMAT
COMMON /COEFNT/ A,G,REY,DEL,AMDA

```

```

DEFINE THE SPACING OF THE INTERIOR MESH POINTS.

```

```

DEL = 1D0/DFLOAT(N+1)

```

```

INITIALIZE ALL ELEMENTS IN THE ARRAY TO ZERO.

```

```

DO 1 I=1,N

```

```

DO 1 J=1,N

```

```

580 MST2
590 MST2
600 MST2
610 MST2
620 MST2
630 MST2
640 MST2
650 MST2
660 MST2
670 MST2
680 MST2
690 MST2
700 MST2
710 MST2
720 MST2
730 MST2
740 MST2
750 MST2
760 MST2
770 MST2
780 MST2
790 MST2
800 MST2
810 MST2
820 MST2
830 MST2
840 MST2
850 MST2
860 MST2
870 MST2
880 MST2
890 MST2
900 MST2
910 MST2
920 MST2
930 MST2
940 MST2
950 MST2
960 MST2
970 MST2
980 MST2
990 MST2
1000 MST2
1010 MST2
1020 MST2
1030 MST2
1040 MST2
1050 MST2

```

ESTABLISH THE CENTRAL DIFFERENCE APPROXIMATION AT EACH POINT IN THE MESH.

```
1 X(I,J) = (ODO,ODO)
```

CCCCC

```

ETA = DEL
CALL MSET1 (ETA, CQM1, CQM2, KSET)
X(1,1) = CFMAT(1,1, CQM1, CQM2)
X(1,2) = CFMAT(1,2, CQM1, CQM2)
X(1,3) = CFMAT(1,3, CQM1, CQM2)
X(1,4) = CFMAT(1,4, CQM1, CQM2)
X(1,5) = CFMAT(1,5, CQM1, CQM2)

```

C

```

ETA = 2DO*DEL
CALL MSET1 (ETA, CQM1, CQM2, KSET)
X(2,1) = CFMAT(2,1, CQM1, CQM2)
X(2,2) = CFMAT(2,2, CQM1, CQM2)
X(2,3) = CFMAT(2,3, CQM1, CQM2)
X(2,4) = CFMAT(2,4, CQM1, CQM2)
X(2,5) = CFMAT(2,5, CQM1, CQM2)

```

C

IL = N-2

C

DO 2 I=3, IL

K=1-3

ETA = DEL\*DFLOAT(I)

CALL MSET1 (ETA, CQM1, CQM2, KSET)

C

DC 2 J=1, 5

2 X(I, K+J) = CFMAT(3, J, CQM1, CQM2)

CC

```

ETA = 1DO*DEL
CALL MSET1 (ETA, CQM1, CQM2, KSET)
X(N-1, N-3) = CFMAT(4,1, CQM1, CQM2)
X(N-1, N-2) = CFMAT(4,2, CQM1, CQM2)
X(N-1, N-1) = CFMAT(4,3, CQM1, CQM2)
X(N-1, N) = CFMAT(4,4, CQM1, CQM2)

```

C

```

ETA = 1DO*DEL
CALL MSET1 (ETA, CQM1, CQM2, KSET)
X(N, N-3) = CFMAT(5,1, CQM1, CQM2)
X(N, N-2) = CFMAT(5,2, CQM1, CQM2)
X(N, N-1) = CFMAT(5,3, CQM1, CQM2)
X(N, N) = CFMAT(5,4, CQM1, CQM2)

```



RETURN  
END

CC

.....FUNCTION COM1E1(JSTA,K,CQM1,CQM2).....  
(POLAR COORDINATES)

PURPOSE  
RETURNS THE VALUES FOR THE COEFFICIENTS IN THE ARRAYS  
REPRESENTING THE CENTRAL DIFFERENCE APPROXIMATION OF THE  
VORTICITY TRANSPORT EQUATION USING THE COEFFICIENTS COMPUTED  
BY SUBROUTINE MSET1.

DESCRIPTION OF PARAMETERS  
JSTA - INDICATES WHICH DIFFERENCE EQUATION SET WILL BE USED.  
JSTA=1 - CONSISTENT 3RD ORDER TRUNCATION ERROR FINITE  
DIFFERENCE EQUATIONS FOR R=DEL WILL BE USED.  
JSTA=2 - SAME AS ABOVE BUT R=2DO\*DEL  
JSTA=3 - CENTRAL DIFFERENCE EQUATIONS WITH CONSISTENT 2ND  
ORDER TRUNCATION ERROR WILL BE USED.  
JSTA=4 - SAME AS JSTA=3 BUT FOR R=1DO-2DO\*DEL.  
JSTA=5 - SAME AS ABOVE BUT FOR R=1DO-DEL.

K - INDICATES THE ABSOLUTE POSITION OF THE POINT IN EACH ROW  
OF THE FINITE DIFFERENCE MESH. IF THE FIRST NON-ZERO ENTRY  
IN ROW J IS ELEMENT(J,3), THEN K=1 DENOTES ELEMENT(J,3),  
K=2 DENOTES ELEMENT(J,4), ETC.

CQM1, CQM2 - THE COEFFICIENT ARRAYS FOR THE FINITE DIFFERENCE  
APPROXIMATION OF THE FUNCTION Q. CQM1 CONTAINS THE  
COEFFICIENTS FOR THE NON-GAMMA TERMS, WHILE CQM2 CONTAINS  
THE COEFFICIENTS OF THE GAMMA TERMS. BOTH ARRAYS MUST BE  
DIMENSIONED COMPLEX\*16.

EXAMPLE OF THE CALLING ARGUMENT:  
CQ M(1,2) E1(JSTA,K,CQM1,CQM2)  
CQ - 0 COMPONENT OF THE VELOCITY VECTOR POTENTIAL.  
M(1,2) - 1 REFERS TO TERMS NOT CONTAINING GAMMA AS A  
FACTOR.  
2 REFERS TO TERMS CONTAINING GAMMA AS A FACTOR.

CQM1 10  
CQM1 20  
CQM1 30  
CQM1 40  
CQM1 50  
CQM1 60  
CQM1 70  
CQM1 80  
CQM1 90  
CQM1 100  
CQM1 110  
CQM1 120  
CQM1 130  
CQM1 140  
CQM1 150  
CQM1 160  
CQM1 170  
CQM1 180  
CQM1 190  
CQM1 200  
CQM1 210  
CQM1 220  
CQM1 230  
CQM1 240  
CQM1 250  
CQM1 260  
CQM1 270  
CQM1 280  
CQM1 290  
CQM1 300  
CQM1 310  
CQM1 320  
CQM1 330  
CQM1 340  
CQM1 350  
CQM1 360  
CQM1 370  
CQM1 380  
CQM1 390  
CQM1 400  
CQM1 410  
CQM1 420  
CQM1 430  
CQM1 440



```

12300*CQM1(3)/(1500*DEL**2)+400*CQM1(2)/(300*DEL)+CQM1(1)
GO TO 29
10 CCM1E1 = -5.34227405300*CQM1(5)/DEL**4-.25420463600*CQM1(4)/DEL**3
1+1900*CQM1(3)/(2000*DEL**2)
GO TO 29
11 CCM1E1 = 1.66608357700*CQM1(5)/DEL**4+.3231788000*CQM1(4)/DEL**3
GO TO 29
12 CCM1E1 = -0.12342079700*CQM1(5)/DEL**4
GO TO 29
CENTRAL DIFFERENCE APPROXIMATION FOR COMPONENT Q (NON GAMMA).
13 GO TO (14,15,16,17,18), K
14 CCM1E1 = CQM1(5)/DEL**4-CQM1(4)/(200*DEL**3)
GO TO 29
15 CCM1E1 = -400*CQM1(5)/DEL**4+CQM1(4)/DEL**3+CQM1(3)/DEL**2-CQM1(2)
1/(200*DEL)
GO TO 29
16 CCM1E1 = 600*CQM1(5)/DEL**4-200*CQM1(3)/DEL**2+CQM1(1)
GO TO 29
17 CCM1E1 = -400*CQM1(5)/DEL**4-CQM1(4)/DEL**3+CQM1(3)/DEL**2+CQM1(2)
1/(200*DEL)
GO TO 29
18 CCM1E1 = CQM1(5)/DEL**4+CQM1(4)/(200*DEL**3)
GO TO 29
FINITE DIFFERENCE EQUATIONS AT ETA=100-200*DEL (NCN-GAMMA).
19 GO TO (20,21,22,23), K
20 CCM1E1 = CQM1(5)/DEL**4-0.500*CQM1(4)/DEL**3
GO TO 29
21 CCM1E1 = -400*CQM1(5)/DEL**4+CQM1(4)/DEL**3+CQM1(3)/DEL**2-CQM1(2)
1/(200*DEL)
GO TO 29
22 CCM1E1 = 600*CQM1(5)/DEL**4-200*CQM1(3)/DEL**2+CQM1(1)
GO TO 29
23 CCM1E1 = -400*CQM1(5)/DEL**4-CQM1(4)/DEL**3+CQM1(3)/DEL**2+CQM1(2)
1/(200*DEL)
GO TO 29
FINITE DIFFERENCE EQUATIONS AT ETA=100-DEL (NON GAMMA).
24 GO TO (25,26,27,28), K
25 CCM1E1 = -0.2500*CQM1(5)/DEL**4
GO TO 29
26 CCM1E1 = 800*CQM1(5)/(300*DEL**4)-CQM1(4)/(300*DEL**3)
GO TO 29
27 CCM1E1 = -900*CQM1(5)/DEL**4+CQM1(3)/DEL**2-CQM1(2)/(200*DEL)
GO TO 29

```

C  
C

C  
C

C  
C



```

GO TO 29
28 CQMIE1 = 16D0*CQM1(5) / DEL**4+3D0*CQM1 (4) / DEL**3-2D0*CQM1 (3) / DEL**2
29 CQM1(1)
25 RETURN
C
ENTRY CQM2E1(JSTA,K,CQM1,CQM2)
C
GO TO (30,35,40,45,5C), JSTA
C
FINITE DIFFERENCE EQUATIONS AT ETA=DEL ( GAMMA).
C
30 GO TO (31,32,33,34,34), K
31 CQM2E1 = -35D0*CQM2(3) / (60D0*DEL**2)-2D0*CQM2(2) / (3D0*DEL)+CQM2(1)
GO TO 54
32 CQM2E1 = 8D0*CQM2(3) / (15D0*DEL**2)+2D0*CQM2(2) / (3D0*DEL)
GO TO 54
33 CQM2E1 = CQM2(3) / (20D0*DEL**2)
GO TO 54
34 CQM2E1 = (0D0,0D0)
C
FINITE DIFFERENCE EQUATIONS AT ETA=2D0*DEL ( GAMMA)
C
35 GO TO (36,37,38,39,39), K
36 CQM2E1 = 11D0*CQM2(3) / (12D0*DEL**2)-4D0*CQM2(2) / (3D0*DEL)
GO TO 54
37 CQM2E1 = -28D0*CQM2(3) / (15D0*DEL**2)+D0*CQM2(2) / (3D0*DEL)+CQM2(1)
GO TO 54
38 CQM2E1 = 19D0*CQM2(3) / (20D0*DEL**2)
GO TO 54
39 CQM2E1 = (0D0,0D0)
C
CENTRAL DIFFERENCE EQUATIONS FOR THE COMPONENT Q ( GAMMA).
C
40 GO TO (41,42,43,44,41), K
41 CQM2E1 = (0D0,0D0)
GO TO 54
42 CQM2E1 = CQM2(3) / DEL**2-CQM2(2) / (2D0*DEL)
GO TO 54
43 CQM2E1 = -2D0*CQM2(3) / DEL**2+CQM2(1)
GO TO 54
44 CQM2E1 = CQM2(3) / DEL**2+CQM2(2) / (2D0*DEL)
GO TO 54
C
FINITE DIFFERENCE EQUATIONS AT ETA=1D0-2D0*DEL ( GAMMA).
C
45 GO TO (49,46,47,48), K
46 CQM2E1 = CQM2(3) / DEL**2-CQM2(2) / (2D0*DEL)
CQM11410
CQM11420
CQM11430
CQM11440
CQM11450
CQM11460
CQM11470
CQM11480
CQM11490
CQM11500
CQM11510
CQM11520
CQM11530
CQM11540
CQM11550
CQM11560
CQM11570
CQM11580
CQM11590
CQM11600
CQM11610
CQM11620
CQM11630
CQM11640
CQM11650
CQM11660
CQM11670
CQM11680
CQM11690
CQM11700
CQM11710
CQM11720
CQM11730
CQM11740
CQM11750
CQM11760
CQM11770
CQM11780
CQM11790
CQM11800
CQM11810
CQM11820
CQM11830
CQM11840
CQM11850
CQM11860
CQM11870
CQM11880

```

```

47 CQM2E1 = -200*CQM2(3)/DEL**2+CQM2(1)
48 CQM2E1 = CQM2(3)/DEL**2+CQM2(2)/(200*DEL)
49 CQM2E1 = (000,000)
50 GO TO (53,52,51,52),K
51 CQM2E1 = CQM2(3)/DEL**2-CQM2(2)/(200*DEL)
52 CQM2E1 = -200*CQM2(3)/DEL**2+CQM2(1)
53 CQM2E1 = (000,000)
54 RETURN
END

```

FINITE DIFFERENCE EQUATIONS AT ETA=1D0-DEL ( GAMMA ).

```

CQM11890
CQM11900
CQM11910
CQM11920
CQM11930
CQM11940
CQM11950
CQM11960
CQM11970
CQM11980
CQM11990
CQM12000
CQM12010
CQM12020
CQM12030
CQM12040
CQM12050
CQM12060

```

```

C.....SUBROUTINE CDMTIN(N,A,NDIM,IERR).....
PURPOSE
  INVERT A COMPLEX*16 MATRIX
USAGE
  CALL CDMTIN(N,A,NDIM,DETERM)
DESCRIPTION OF PARAMETERS
  N - ORDER OF COMPLEX*16 MATRIX TO BE INVERTED
      (INTEGER) MAXIMUM 'N' IS 100
  A - COMPLEX*16 INPUT MATRIX (DESTROYED). THE
      INVERSE OF 'A' IS RETURNED IN ITS PLACE
  NDIM - THE SIZE TO WHICH 'A' IS DIMENSIONED
      (ROW DIMENSION OF 'A' ACTUALLY APPEARING
      IN THE DIMENSION STATEMENT OF USER'S
      CALLING PROGRAM)
  IERR - ERROR PARAMETER RETURNED BY CDMTIN. IERR = 0 INDICATES
      NORMAL INVERSION. IERR = 9999 INDICATES SINGULAR MATRIX.
REMARKS

```

```

CDMT 10
CDMT 20
CDMT 30
CDMT 40
CDMT 50
CDMT 60
CDMT 70
CDMT 80
CDMT 90
CDMT 100
CDMT 110
CDMT 120
CDMT 130
CDMT 140
CDMT 150
CDMT 160
CDMT 170
CDMT 180
CDMT 190
CDMT 200
CDMT 210
CDMT 220
CDMT 230
CDMT 240
CDMT 250
CDMT 260
CDMT 270
CDMT 280

```





CDMT 770  
 CDMT 780  
 CDMT 790  
 CDMT 800  
 CDMT 810  
 CDMT 820  
 CDMT 830  
 CDMT 840  
 CDMT 850  
 CDMT 860  
 CDMT 870  
 CDMT 880  
 CDMT 890  
 CDMT 900  
 CDMT 910  
 CDMT 920  
 CDMT 930  
 CDMT 940  
 CDMT 950  
 CDMT 960  
 CDMT 970  
 CDMT 980  
 CDMT 990  
 CDMT 1000  
 CDMT 1010  
 CDMT 1020  
 CDMT 1030  
 CDMT 1040  
 CDMT 1050  
 CDMT 1060  
 CDMT 1070  
 CDMT 1080  
 CDMT 1090  
 CDMT 1100  
 CDMT 1110  
 CDMT 1120  
 CDMT 1130  
 CDMT 1140  
 CDMT 1150  
 CDMT 1160  
 CDMT 1170  
 CDMT 1180  
 CDMT 1190  
 CDMT 1200  
 CDMT 1210  
 CDMT 1220  
 CDMT 1230  
 CDMT 1240

```

    ICOLUM = K
    AMAX = A(J,K)
  6 CONTINUE
  7 CCNTINUE
    IPIVOT(ICOLUM) = IPIVOT(ICOLUM)+1
    INTERCHANGE ROWS TO PUT PIVOT ELEMENT ON DIAGONAL
  8 CONTINUE
    IF (IROW-ICCLUM) 8,10,8
  9 A(ICOLUM,L) = SWAP
    DC 9 L=1,N
    SWAP = A(IROW,L)
    A(IROW,L) = A(ICOLUM,L)
    A(ICOLUM,L) = SWAP
  10 INDEX(I,1) = IROW
    INDEX(I,2) = ICOLUM
    PIVOT(I) = A(ICOLUM, ICOLUM)
    U = PIVOT(I)
    TEMP = PIVOT(I)*DCONJG(PIVOT(I))
    IF (TEMP) 11,20,11
    DIVIDE PIVOT ROW BY PIVOT ELEMENT
  11 A(ICOLUM,ICOLUM) = (1/CO,ODO)
    DO 12 L=1,N
    U = PIVOT(I)
  12 A(ICOLUM,L) = A(ICOLUM,L)/U
    REDUCE NON-PIVOT ROWS
    DO 15 L=1,N
    IF (L1-ICOLUM) 13,15,13
  13 T = A(L1,ICOLUM)
    A(L1,ICOLUM) = (OCO,ODO)
    DO 14 L=1,N
    U = A(ICOLUM,L)
  14 A(L1,L) = A(L1,L)-U*T
  
```

CDMT11250  
 CDMT11260  
 CDMT11270  
 CDMT11280  
 CDMT11290  
 CDMT11300  
 CDMT11310  
 CDMT11320  
 CDMT11330  
 CDMT11340  
 CDMT11350  
 CDMT11360  
 CDMT11370  
 CDMT11380  
 CDMT11390  
 CDMT11400  
 CDMT11410  
 CDMT11420  
 CDMT11430  
 CDMT11440  
 CDMT11450  
 CDMT11460  
 CDMT11470  
 CDMT11480  
 CDMT11490  
 CDMT11500  
 CDMT11510  
 CDMT11520  
 CDMT11530  
 CDMT11540

```

15 CONTINUE
16 CONTINUE
      INTERCHANGE COLUMNS
      DO 19 I=1,N
      L = N+1-I
      IF (INDEX(L,1)-INDEX(L,2)) 17,19,17
17 JROW = INDEX(L,1)
      JCOLUM = INDEX(L,2)
      DO 18 K=1,N
      SWAP = A(K,JROW)
      A(K,JROW) = A(K,JCOLUM)
      A(K,JCOLUM) = SWAP
18 CONTINUE
19 CONTINUE
20 RETURN (6,22)
21 IERR = 9999
      RETURN
22 FFORMAT (20H MATRIX IS SINGULAR)
      END
  
```

MULM 10  
 MULM 20  
 MULM 30  
 MULM 40  
 MULM 50  
 MULM 60  
 MULM 70  
 MULM 80  
 MULM 90  
 MULM 100  
 MULM 110  
 MULM 120  
 MULM 130  
 MULM 140  
 MULM 150  
 MULM 160

```

..... SUBROUTINE MULM(X1,X2,N,MDIM,TEMPV).....
PURPOSE
PERFORMS THE MATRIX MULTIPLICATION OF A SQUARE MATRIX BY A
SQUARE MATRIX. THE RESULT IS RETURNED IN MATRIX X1.
USAGE
CALL MULM(X1,X2,N,MDIM,TEMPV)
DESCRIPTION OF PARAMETERS
X1 - THE MULTIPLYING MATRIX ON INPUT AND THE RESULTANT PRODUCT
      ON OUTPUT.
  
```





```

C.....SUBROUTINE DSPLIT(N,MDIM,A,AR,AI).....
PURPOSE
DSPLIT TAKES A MATRIX OF COMPLEX*16 NUMBERS AND
SPLITS IT INTO TWO MATRICES, ONE CONTAINING THE REAL
PART OF THE ORIGINAL MATRIX, AND ONE CONTAINING THE
IMAGINARY PART.
USAGE
CALL DSPLIT(N,MDIM,A,AREAL,AIMAG)
DESCRIPTION OF PARAMETERS
N - THE SIZE OF THE MATRIX A, AN N BY N SQUARE
MATRIX.
MDIM - THE COLUMN DIMENSION OF MATRIX A
A - THE INPUT MATRIX. MUST BE DIMENSIONED MDIM BY
AT LEAST N IN THE CALLING PROGRAM (COMPLEX*16)
AREAL,AIMAG - THE OUTPUT MATRICES CONTAINING THE
REAL AND IMAGINARY PARTS, RESPECTIVELY, OF
MATRIX A. MUST BE DIMENSIONED (MDIM,MDIM) IN THE
CALLING PROGRAM.
NGTES...
MATRIX A AND MATRIX AREAL MAY OVERLAP IF THEY ARE
DIMENSIONED IN THE CALLING PROGRAM AS FOLLOWS...
COMPLEX*16 A(MDIM,MDIM)
REAL*8 AREAL(MDIM,MDIM),AIMAG(MDIM,MDIM)
EQUIVALENCE(A(1,1),AREAL(1,1))
OTHER ROUTINES NEEDED
NONE
SUBROUTINE DSPLIT(N,MDIM,A,AR,AI)
REAL*8A(2,MDIM,MDIM),AR(MDIM,MDIM),AI(MDIM,MDIM)
C

```

```

DSPL 10
DSPL 20
DSPL 30
DSPL 40
DSPL 50
DSPL 60
DSPL 70
DSPL 80
DSPL 90
DSPL 100
DSPL 110
DSPL 120
DSPL 130
DSPL 140
DSPL 150
DSPL 160
DSPL 170
DSPL 180
DSPL 190
DSPL 200
DSPL 210
DSPL 220
DSPL 230
DSPL 240
DSPL 250
DSPL 260
DSPL 270
DSPL 280
DSPL 290
DSPL 300
DSPL 310
DSPL 320
DSPL 330
DSPL 340
DSPL 350
DSPL 360
DSPL 370
DSPL 380
DSPL 390
DSPL 400
DSPL 410
DSPL 420
DSPL 430
DSPL 440
DSPL 450
DSPL 460
DSPL 470
DSPL 480

```

DSPL 490  
 DSPL 510  
 DSPL 520  
 DSPL 530  
 DSPL 540  
 DSPL 550  
 DSPL 560  
 DSPL 570  
 DSPL 580

DO 1 J=1,N  
 DO 1 I=1,N  
 AR(I,J) = A(1,I,J)  
 1 AI(I,J) = A(2,I,J)

RETURN  
 END

SUBROUTINE EBALAC (AR, AI, N, IA, K, L, D)

-----D-----LIBRARY I-----

FUNCTION

USAGE  
 PARAMETERS

AR  
 AI

N  
 IA  
 K  
 L

D

PRECISION  
 LANGUAGE

LATEST REVISION

SUBROUTINE EBALAC (AR, AI, N, IA, K, L, D)

DIMENSION

AR(IA, 1), AI(IA, 1), D(N)

EBALAC 0010  
 EBALAC 0020  
 EBALAC 0030  
 EBALAC 0040  
 EBALAC 0050  
 EBALAC 0060  
 EBALAC 0070  
 EBALAC 0080  
 EBALAC 0090  
 EBALAC 0100  
 EBALAC 0110  
 EBALAC 0120  
 EBALAC 0130  
 EBALAC 0140  
 EBALAC 0150  
 EBALAC 0160  
 EBALAC 0170  
 EBALAC 0180  
 EBALAC 0190  
 EBALAC 0200  
 EBALAC 0210  
 EBALAC 0220  
 EBALAC 0230  
 EBALAC 0240  
 EBALAC 0250  
 EBALAC 0260  
 EBALAC 0270  
 EBALAC 0280  
 EBALAC 0290  
 EBALAC 0300  
 EBALAC 0310  
 EBALAC 0320  
 EBALAC 0330  
 EBALAC 0340  
 EBALAC 0350  
 EBALAC 0360

-----  
 - BALANCES A COMPLEX GENERAL MATRIX AND ISOLATES  
 EIGENVALUES WHENEVER POSSIBLE.  
 - CALL EBALAC (AR, AI, N, IA, K, L, D)  
 - INPUT/OUTPUT MATRICES OF DIMENSION N BY N.  
 ON INPUT, AR AND AI CONTAIN THE REAL  
 AND IMAGINARY PARTS, RESPECTIVELY, OF  
 THE COMPLEX MATRIX, OF ORDER N TO BE  
 BALANCED. ON OUTPUT, AR AND AI CONTAIN THE  
 REAL AND IMAGINARY PARTS OF THE  
 TRANSFORMED MATRIX.  
 - INPUT VARIABLE CONTAINING THE ORDER  
 OF THE MATRIX A = (AR, AI) TO BE BALANCED.  
 - INPUT VARIABLE CONTAINING THE ROW DIMENSION OF  
 AR AND AI IN THE CALLING PROGRAM.  
 - OUTPUT INTEGERS CONTAINING THE BOUNDARY  
 INDICES FOR THE BALANCED MATRIX A = (AR, AI)  
 SUCH THAT  
 AR(I, J) = 0. AND AI(I, J) = 0. IF  
 (1) J IS GREATER THAN J AND  
 (2) J = 1, ..., K-1 OR  
 I = L+1, ..., N  
 - OUTPUT VECTOR OF LENGTH N CONTAINING  
 INFORMATION DETERMINING THE PERMUTATIONS  
 USED AND THE SCALING FACTORS.  
 - SINGLE/DOUBLE  
 - FORTRAN

-----  
 - MARCH 9, 1977

```

C          LOGICAL          NOCONV          RADIX IS A MACHINE DEPENDENT
C          DOUBLE PRECISION  AR, AI, 0, RADIX, ZERO, ONE, PT95, B2, F, C, G, R, S
C          DATA            RRADIX, RB2
C          DATA            RADIX/16.0D0/
C          DATA            ZERO, ONE, PT95/0.0D0, 1.0D0, 0.95D0/
C          B2 = RADIX*RRADIX
C          RRADIX = ONE/RADIX
C          RB2 = RRADIX*RRADIX
C          K = 1
C          L = N
C          GO TO 30
C
C          5 D(M) = J.EQ.M) GO TO 20
C          IF (J.I.EQ.1, L
C          DO 10 I = 1, L
C          F = AR(I, J)
C          AR(I, J) = AR(I, M)
C          AR(I, M) = F
C          F = AI(I, J)
C          AI(I, J) = AI(I, M)
C          AI(I, M) = F
C          10 CONTINUE
C          DO 15 I = K, N
C          F = AR(J, I)
C          AR(J, I) = AR(M, I)
C          AR(M, I) = F
C          F = AI(J, I)
C          AI(J, I) = AI(M, I)
C          AI(M, I) = F
C          15 CONTINUE
C          20 GO TO (25, 45), IEXC
C
C          25 IF (L.EQ. 1) GO TO 115
C          L = L-1
C
C          30 L1 = L+1
C          DO 40 JJ = 1, L
C          J = L1-JJ
C          DO 35 I = 1, L
C          IF (I.EQ. J) GO TO 35
C          IF (AR(J, I).NE. ZERO .OR. AI(J, I) .NE. ZERC) GO TO 40
C          35 CONTINUE
C
C          SEARCH FOR ROWS ISOLATING AN
C          EIGENVALUE AND PUSH THEM DOWN
C          DO J=L, 1, -1
C
C          IN-LINE PROCEDURE FOR ROW AND COLUMN
C          EXCHANGE
C
C          EBAC0370
C          EBAC0380
C          EBAC0390
C          EBAC0400
C          EBAC0410
C          EBAC0420
C          EBAC0430
C          EBAC0440
C          EBAC0450
C          EBAC0480
C          EBAC0490
C          EBAC0500
C          EBAC0510
C          EBAC0520
C          EBAC0530
C          EBAC0540
C          EBAC0550
C          EBAC0560
C          EBAC0570
C          EBAC0580
C          EBAC0590
C          EBAC0600
C          EBAC0610
C          EBAC0620
C          EBAC0630
C          EBAC0640
C          EBAC0650
C          EBAC0660
C          EBAC0670
C          EBAC0680
C          EBAC0690
C          EBAC0700
C          EBAC0710
C          EBAC0720
C          EBAC0730
C          EBAC0740
C          EBAC0750
C          EBAC0760
C          EBAC0770
C          EBAC0780
C          EBAC0790
C          EBAC0800
C          EBAC0810
C          EBAC0820
C          EBAC0830
C          EBAC0840
C          EBAC0850
C          EBAC0860

```



```

EBAC 0870
EBAC 0880
EBAC 0890
EBAC 0900
EBAC 0910
EBAC 0920
EBAC 0930
EBAC 0940
EBAC 0950
EBAC 0960
EBAC 0970
EBAC 0980
EBAC 0990
EBAC 1000
EBAC 1010
EBAC 1020
EBAC 1030
EBAC 1040
EBAC 1050
EBAC 1060
EBAC 1070
EBAC 1080
EBAC 1090
EBAC 1100
EBAC 1110
EBAC 1120
EBAC 1130
EBAC 1140
EBAC 1150
EBAC 1160
EBAC 1180
EBAC 1200
EBAC 1210
EBAC 1220
EBAC 1230
EBAC 1240
EBAC 1250
EBAC 1260
EBAC 1270
EBAC 1280
EBAC 1290
EBAC 1300
EBAC 1310
EBAC 1320
EBAC 1330
EBAC 1340
EBAC 1350
EBAC 1360

```

```

C
C
M = L
IEXC = 1
GO TO 5
40 CONTINUE
GO TO 50
C
45 K = K+1
50 DO 60 J = K,L
DO 55 I = K,L
IF (I.EQ. J) GC TO 55
IF (AR(I,J) .NE. ZERO .OR. AI(I,J) .NE. ZERO) GO TO 60
55 CONTINUE
M = K
IEXC = 2
GO TO 5
60 CONTINUE
C
DO 65 I = K,L
D(I) = ONE
65 CONTINUE
C
70 NOCONV = .FALSE.
DO 110 I = K,L
C = ZERO
DO 75 J = K,L
IF (J.EQ. I) GO TO 75
C = C+DABS(AR(J,I))+DABS(AI(J,I))
R = R+DABS(AR(I,J))+DABS(AI(I,J))
75 CONTINUE
G = R*RRADIX
F = ONE
S = C+R
IF (C .GE. G) GO TO 85
F = F*RADIX
C = C*B2
GO TO 80
85 G = R*RRADIX
90 IF (C .LT. G) GO TO 95
F = F*RRADIX
C = C*RB2
GO TO 90
C
95 IF ((C+R)/F .GE. PT95*S) GO TO 110
G = ONE/F
D(I) = D(I)*F

```

SEARCH FOR COLUMNS ISOLATING AN EIGENVALUE AND PUSH THEM LEFT

BALANCE THE SUBMATRIX IN ROWS  
K TO L

ITERATIVE LOOP FOR NORM REDUCTION

BALANCE

EBAC1370  
 EBAC1380  
 EBAC1390  
 EBAC1400  
 EBAC1410  
 EBAC1420  
 EBAC1430  
 EBAC1440  
 EBAC1450  
 EBAC1460  
 EBAC1470  
 EBAC1480  
 EBAC1490

EHEC0010  
 EHEC0020  
 EHEC0030  
 EHEC0040  
 EHEC0050  
 EHEC0060  
 EHEC0070  
 EHEC0080  
 EHEC0090  
 EHEC0100  
 EHEC0110  
 EHEC0120  
 EHEC0130  
 EHEC0140  
 EHEC0150  
 EHEC0160  
 EHEC0170  
 EHEC0180  
 EHEC0190  
 EHEC0200  
 EHEC0210  
 EHEC0220  
 EHEC0230  
 EHEC0240  
 EHEC0250  
 EHEC0260  
 EHEC0270  
 EHEC0280  
 EHEC0290  
 EHEC0300  
 EHEC0310  
 EHEC0320  
 EHEC0330  
 EHEC0340

```

NOCNV = .TRUE.
DO 100 J = K,N
  AR(I,J) = AR(I,J)*G
  AI(I,J) = AI(I,J)*G
CONTINUE
DO 105 J = 1,L
  AR(J,I)*F
  AI(J,I) = AI(J,I)*F
CONTINUE
100 CONTINUE
105 CONTINUE
110 IF (NOCNV) GO TO 70
115 RETURN
END
  
```

-----D-----LIBRARY 1-----  
 SUBROUTINE EHESSC (AR,AI,K,L,N,IA,ID)

FUNCTION  
 USAGE  
 PARAMETERS AR  
 AI  
 K  
 L  
 N  
 IA  
 ID  
 PRECISION

- REDUCTION OF A COMPLEX MATRIX TO COMPLEX UPPER HESSENBERG FORM.  
 - CALL EHESSC(AR,AI,K,L,N,IA,ID)  
 - INPUT/OUTPUT MATRIX OF DIMENSION N BY N. OF  
 ON INPUT CONTAINS THE REAL COMPONENTS OF  
 THE MATRIX TO BE REDUCED.  
 ON OUTPUT CONTAINS THE REAL COMPONENTS  
 OF THE REDUCED HESSENBERG FORM IN THE  
 UPPER TRIANGULAR PORTION (INCLUDING MAIN  
 AND SUB-DIAGONAL) AND THE DETAILS OF  
 THE REDUCTION IN THE LOWER TRIANGULAR  
 PORTION.  
 - INPUT/OUTPUT MATRIX OF DIMENSION N BY N  
 CONTAINING THE IMAGINARY COUNTERPARTS  
 TO AR ABOVE.  
 - INPUT SCALAR CONTAINING THE ROW AND COLUMN  
 INDEX OF THE STARTING ELEMENT TO BE  
 REDUCED BY ROW SCALING. FOR UNBALANCED  
 MATRICES SET K = 1.  
 - INPUT SCALAR CONTAINING THE ROW AND  
 COLUMN INDEX OF THE LAST ELEMENT TO BE  
 REDUCED BY ROW SCALING. FOR UNBALANCED  
 MATRICES SET L = N.  
 - INPUT SCALAR CONTAINING THE ORDER OF  
 THE MATRIX TO BE REDUCED.  
 - INPUT SCALAR CONTAINING ROW DIMENSION  
 OF AR AND AI IN THE CALLING PROGRAM.  
 - OUTPUT VECTOR OF LENGTH L CONTAINING  
 DETAILS OF THE TRANSFORMATIONS.  
 - SINGLE/DOUBLE

```

C-----
C LANGUAGE - FORTRAN
C-----
C LATEST REVISION - FEBRUARY 7, 1973
C
C SUBROUTINE EHESSC (AR, AI, K, L, N, IA, ID)
C
C DIMENSION
C DOUBLE PRECISION
C COMPLEX *16
C EQUIVALENCE
C
C DATA
C LA=L-1
C KPI=K+1
C IF (LA .LT. KPI) GO TO 45
C DC 40 M=KPI, LA
C I=M
C XR=ZERO
C XI=ZERO
C DO 5 J=M, L
C IF (DABS(AR(J, M-1))+DABS(AI(J, M-1)) .LE. DABS(XR)+DABS(XI))
C GO TO 5
C XR=AR(J, M-1)
C XI=AI(J, M-1)
C I=J
C CONTINUE
C ID(M)=I
C IF (I .EQ. M) GO TO 20
C
C MM1=M-1
C DO 10 J=MM1, N
C YR=AR(I, J)
C AR(I, J)=AR(M, J)
C AR(M, J)=YR
C YI=AI(I, J)
C AI(I, J)=AI(M, J)
C AI(M, J)=YI
C CONTINUE
C DO 15 J=1, L
C YR=AR(J, I)
C AR(J, I)=AR(J, M)
C AR(J, M)=YR
C YI=AI(J, I)
C AI(J, I)=AI(J, M)
C AI(J, M)=YI
C CONTINUE
C
C INTERCHANGE ROWS AND COLUMNS OF
C ARRAYS AR AND AI
C
C END INTERCHANGE

```

```

EHEC 0350
EHEC 0360
EHEC 0370
EHEC 0380
EHEC 0390
EHEC 0400
EHEC 0410
EHEC 0420
EHEC 0430
EHEC 0450
EHEC 0460
EHEC 0470
EHEC 0490
EHEC 0500
EHEC 0510
EHEC 0520
EHEC 0530
EHEC 0540
EHEC 0550
EHEC 0560
EHEC 0570
EHEC 0580
EHEC 0610
EHEC 0620
EHEC 0630
EHEC 0640
EHEC 0650
EHEC 0660
EHEC 0670
EHEC 0680
EHEC 0690
EHEC 0700
EHEC 0710
EHEC 0720
EHEC 0730
EHEC 0740
EHEC 0750
EHEC 0760
EHEC 0770
EHEC 0780
EHEC 0790
EHEC 0800
EHEC 0810
EHEC 0820
EHEC 0830
EHEC 0840
EHEC 0850
EHEC 0860

```



EHEC0870  
 EHEC0880  
 EHEC0890  
 EHEC0900  
 EHEC0910  
 EHEC0920  
 EHEC0930  
 EHEC0940  
 EHEC0950  
 EHEC0960  
 EHEC0970  
 EHEC0980  
 EHEC0990  
 EHEC1000  
 EHEC1010  
 EHEC1020  
 EHEC1030  
 EHEC1040  
 EHEC1050  
 EHEC1060  
 EHEC1070

```

20 IF (XR .EQ. ZERO .AND. XI .EQ. ZERO) GO TO 40
    MPI=M+1
    DO 35 I=MPI,L
      YI=AR(I,M-1)
      YI=AI(I,M-1)
      IF (YR .EQ. ZERO .AND. YI .EQ. ZERO) GO TO 35
      Y=Y/X
      AR(I,M-1)=YR
      AR(I,J)=AR(I,J)-YR*AR(M,J)+YI*AI(M,J)
      AI(I,M-1)=YI
      AI(I,J)=AI(I,J)-YR*AI(M,J)-YI*AR(M,J)
    DO 25 J=M,N
      CONTINUE
    DO 30 J=1,L
      AR(J,M)+YR*AR(J,I)-YI*AI(J,I)
      AI(J,M)=AI(J,I)+YI*AR(J,I)
    CONTINUE
  CONTINUE
  CONTINUE
  CONTINUE
  RETURN
  END
  
```

ELR20010  
 ELR20020  
 ELR20030  
 ELR20040  
 ELR20050  
 ELR20060  
 ELR20070  
 ELR20080  
 ELR20090  
 ELR20100  
 ELR20110  
 ELR20120  
 ELR20130  
 ELR20140  
 ELR20150  
 ELR20160  
 ELR20170  
 ELR20180  
 ELR20190  
 ELR20200  
 ELR20210  
 ELR20220  
 ELR20230  
 ELR20240  
 ELR20250

```

SUBROUTINE ELRH2C (HR,HI,K,L,N,IH,WR,WI,ZR,ZI,ID,INFER,IER)
-----LIBRARY 1-----
FUNCTION
  - COMPUTE THE EIGENVALUES AND EIGENVECTORS OF
  A COMPLEX UPPER HESSENBERG MATRIX AND
  BACK TRANSFORM THE EIGENVECTORS.
USAGE
  - CALL ELRH2C (FR,HI,K,L,N,IH,WR,WI,ZR,ZI,ID,
  INFER,IER)
PARAMETERS
  FR INPUT MATRIX OF DIMENSION N BY N CONTAINING
  THE REAL COMPONENTS OF THE COMPLEX
  HESSENBERG MATRIX. HR IS DESTROYED ON
  OUTPUT.
  HI INPUT MATRIX OF DIMENSION N BY N CONTAINING
  THE IMAGINARY COUNTERTPARTS TO HR, ABOVE.
  HI IS DESTROYED ON OUTPUT.
  K INPUT SCALAR CONTAINING THE LOWER BOUNDARY
  INDEX FOR THE INPUT MATRIX.
  L FOR UNBALANCED MATRICES SET K = 1.
  L INPUT SCALAR CONTAINING THE UPPER BOUNDARY
  INDEX FOR THE INPUT MATRIX.
  N FOR UNBALANCED MATRICES THE ORDER OF THE
  HESSENBERG MATRIX AND THE EIGENVECTOR
  MATRIX.
  
```

AD-A066 374

NAVAL POSTGRADUATE SCHOOL MONTEREY CALIF  
INVESTIGATION OF PIPE FLOW INSTABILITY AND RESULTS FOR WAVE NUM--ETC(U)  
DEC 78 M J ARNOLD

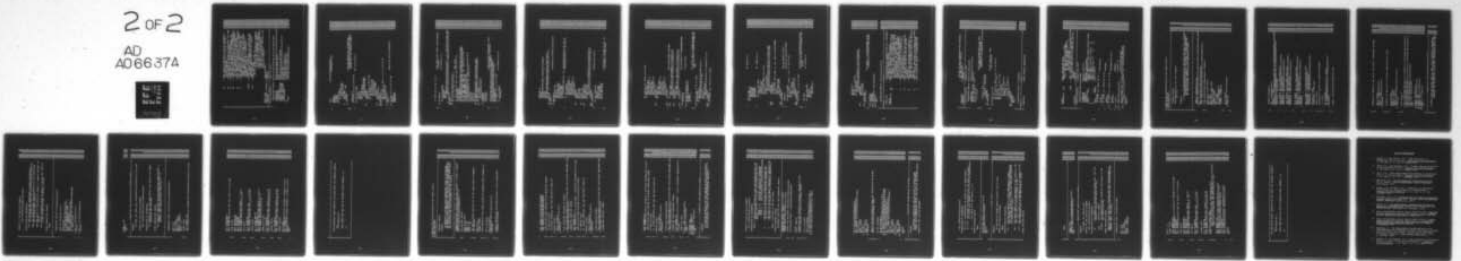
F/G 20/4

UNCLASSIFIED

NL

2 OF 2

AD  
A066374



END  
DATE  
FILMED

'5--79  
DDC

ELR20260  
 ELR20270  
 ELR20280  
 ELR20290  
 ELR20300  
 ELR20310  
 ELR20320  
 ELR20330  
 ELR20340  
 ELR20350  
 ELR20360  
 ELR20370  
 ELR20380  
 ELR20390  
 ELR20400  
 ELR20410  
 ELR20420  
 ELR20430  
 ELR20440  
 ELR20450  
 ELR20460  
 ELR20470  
 ELR20480  
 ELR20490  
 ELR20500  
 ELR20510  
 ELR20520  
 ELR20530  
 ELR20540  
 ELR20550  
 ELR20560  
 ELR20570  
 ELR20580  
 ELR20590  
 ELR20600  
 ELR20610  
 ELR20620  
 ELR20630  
 ELR20640  
 ELR20650  
 ELR20660  
 ELR20680  
 ELR20690  
 ELR20700  
 ELR20710  
 ELR20720  
 ELR20730  
 ELR20740

IH - INPUT SCALAR CONTAINING THE ROW DIMENSION  
 OF MATRICES HR, HI, ZR AND ZI IN THE  
 CALLING PROGRAM.  
 WR - OUTPUT VECTOR OF LENGTH N CONTAINING THE REAL  
 COMPONENTS OF THE EIGENVALUES.  
 WI - OUTPUT VECTOR OF LENGTH N CONTAINING THE  
 IMAGINARY COMPONENTS OF THE EIGENVALUES.  
 ZR - OUTPUT MATRIX OF DIMENSION N BY N CONTAINING  
 THE REAL COMPONENTS ARE NOT NORMALIZED.  
 ZI - OUTPUT MATRIX OF DIMENSION N BY N CONTAINING  
 THE IMAGINARY COMPONENTS TO ZR, ABOVE.  
 ID - INPUT VECTOR OF LENGTH N CONTAINING THE  
 INFORMATION GENERATED BY IMSL ROUTINE  
 HESSC IDENTIFYING THE ROWS AND COLUMNS  
 INTERCHANGED DURING THE REDUCTION TO  
 HESSENBERG FORM. ONLY COMPONENTS K  
 THROUGH L ARE USED.  
 INFER - OUTPUT SCALAR CONTAINING THE INDEX OF THE  
 EIGENVALUE WHICH GENERATED THE TERMINAL  
 ERROR (SEE DESCRIPTION OF IER, BELOW).  
 IER - ERROR PARAMETER  
 N = 1 INDICATES THE EIGENVALUE RECORDED  
 IN THE OUTPUT PARAMETER, INFER, 30  
 ITERATIONS. IF THE J-TH EIGENVALUE  
 COULD NOT BE SO DETERMINED,  
 THEN THE EIGENVALUES J+1, J+2, ..., N  
 SHOULD BE CORRECT.

PRECISION - SINGLE/DOUBLE  
 REQD. IMSL - UERTST  
 LANGUAGE - FORTRAN

LATEST REVISION - APRIL 5, 1977

SUBROUTINE ELRH2C (HR, HI, K, L, N, IH, WR, WI, ZR, ZI, ID, INFER, IER)  
 DIMENSION  
 COMPLEX\*16  
 DOUBLE PRECISION  
 DOUBLE PRECISION  
 DCUBLE PRECISION  
 EQUIVALENCE  
 1  
 2 DATA

CC



ELR2 0750  
 ELR2 0770  
 ELR2 0780  
 ELR2 0790  
 ELR2 0800  
 ELR2 0810  
 ELR2 0820  
 ELR2 0830  
 ELR2 0840  
 ELR2 0850  
 ELR2 0860  
 ELR2 0870  
 ELR2 0880  
 ELR2 0890  
 ELR2 0900  
 ELR2 0910  
 ELR2 0920  
 ELR2 0930  
 ELR2 0940  
 ELR2 0950  
 ELR2 0960  
 ELR2 0970  
 ELR2 0980  
 ELR2 0990  
 ELR2 1000  
 ELR2 1010  
 ELR2 1020  
 ELR2 1030  
 ELR2 1040  
 ELR2 1050  
 ELR2 1060  
 ELR2 1070  
 ELR2 1080  
 ELR2 1090  
 ELR2 1100  
 ELR2 1110  
 ELR2 1120  
 ELR2 1130  
 ELR2 1140  
 ELR2 1150  
 ELR2 1160  
 ELR2 1170  
 ELR2 1180  
 ELR2 1190  
 ELR2 1200  
 ELR2 1210  
 ELR2 1220  
 ELR2 1230

EPS/Z34100000000000000000/  
 INITIALIZE IER

C DATA

```
IER=0
INFER=0
TR=ZERO
TI=ZERO
DC 5 I=1, N
DO 3 J=1, N
  ZR(I, J)=ZERO
  ZI(I, J)=ZERO
CONTINUE
ZR(I, I) = ONE
5 CONTINUE
```

FORM THE MATRIX OF ACCUMULATED  
 TRANSFORMATIONS FROM THE INFOR-  
 MATION LEFT BY ROUTINE 'EHESSC'

C C

```
IEND=L-K-1
IF (IEND .LE. 0) GO TO 25
DO 20 I=1, IEND
  I=L-I
```

DO I=L-1, K+1, -1

C

```
IP1=I+1
IMI=I-1
DO 10 M=IP1, L
  ZR(M, I)=HR(M, IM1)
  ZI(M, I)=HI(M, IM1)
CONTINUE
J=I+1
IF (I .EQ. J) GO TO 20
DO 15 M=I, L
  ZR(I, M)=ZR(J, M)
  ZI(I, M)=ZI(J, M)
  ZR(J, M)=ZERO
  ZI(J, M)=ZERO
CONTINUE
ZR(J, I)=ONE
```

15

```
20 CONTINUE
25 DO 30 I=1, N
  IF (I .GE. K .AND. I .LE. L) GO TO 30
  WR(I)=HR(I, I)
  WI(I)=HI(I, I)
CONTINUE
NN=L
```

30

SEARCH FOR NEXT EIGENVALUE

C

```
35 IF (NN .LT. K) GO TO 150
ITS=0
NNM1=NN-1
NNM2=NN-2
```

```

C      IF (NN .EQ. K) GO TO 50
C      LOOK FOR SINGLE SMALL SUB-DIAGONAL
C      ELEMENT
C      DO M=NN,K+1,-1
C      DO 45 KK=K,NNM1
C      M=NPL-KK
C      NMI=M-1
C      IF (DABS(HR(M,MM1))+DABS(HI(M,MM1))) .LE. EPS*(DABS(HR(MM1,MM1))) GO TO 55
C      IF +DABS(HI(MM1,MM1))+DABS(HR(M,M))) GO TO 55
C      1 CONTINUE
C      45 M=K
C      55 IF (M .EQ. NN) GO TO 145
C      IF (ITS .EQ. 30) GO TO 205
C      IF (ITS .EQ. 10 .OR. ITS .EQ. 20) GO TO 60
C      FORM SHIFT
C      SR=HR(NN,NN)
C      SI=HI(NN,NN)
C      XR=HR(NNMI,NN)*HR(NN,NNMI)-HI(NNMI,NN)*HI(NN,NNMI)
C      XI=HR(NNMI,NN)*HI(NN,NNMI)+HI(NNMI,NN)*HR(NN,NNMI)
C      IF (XR .EQ. ZERO .AND. XI .EQ. ZERC) GO TO 65
C      YR=(HR(NNMI,NNMI)-SR)/TWO
C      YI=(HI(NNMI,NNMI)-SI)/TWO
C      Z=CDSQRT(DC MPLX(YR*#2-YI*#2+XR,TWO*YR*YI+XI))
C      IF (YR*ZZR+YI*ZZI .LT. ZERO) Z=-Z
C      X=X/(Y+Z)
C      SR=SR-XR
C      SI=SI-XI
C      GO TO 65
C      60 SR=DABS(HR(NN,NNMI))+DABS(HR(NNMI,NNM2))
C      65 SI=DABS(HI(NN,NNMI))+DABS(HI(NNMI,NNM2))
C      DO 70 I=K,NN
C      HR(I,I)=HR(I,I)-SR
C      HI(I,I)=HI(I,I)-SI
C      CONTINUE
C      TR=TR+SR
C      TI=TI+SI
C      ITS=ITS+1
C      XR=DABS(HR(NNMI,NNMI))+DABS(HI(NNMI,NNMI))
C      YR=DABS(HR(NN,NNMI))+DABS(HI(NNMI,NNMI))
C      ZZR=DABS(HR(NN,NN))+DABS(HI(NN,NNMI))
C      NNMJ=NNM1-M
C      IF (NNMJ .EQ. 0) GO TO 80
C      DO MM=NN-1,M+1,-1
C      DO 75 NM=1,NNMJ
C      PM=NN-NM

```

```

ELR21240
ELR21250
ELR21260
ELR21270
ELR21280
ELR21290
ELR21300
ELR21310
ELR21320
ELR21330
ELR21360
ELR21370
ELR21380
ELR21390
ELR21400
ELR21410
ELR21420
ELR21430
ELR21440
ELR21450
ELR21460
ELR21470
ELR21480
ELR21490
ELR21510
ELR21520
ELR21530
ELR21540
ELR21550
ELR21560
ELR21570
ELR21600
ELR21610
ELR21620
ELR21630
ELR21640
ELR21650
ELR21660
ELR21670
ELR21680
ELR21690
ELR21700
ELR21710
ELR21750
ELR21760
ELR21770
ELR21780
ELR21790

```



ELR21800  
 ELR21810  
 ELR21820  
 ELR21840  
 ELR21850  
 ELR21860  
 ELR21880  
 ELR21890  
 ELR21900  
 ELR21910  
 ELR21920  
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 ELR22120  
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 ELR22190  
 ELR22200  
 ELR22210  
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 ELR22230  
 ELR22240  
 ELR22250  
 ELR22260  
 ELR22270  
 ELR22280  
 ELR22290  
 ELR22300

```

MMM1=MM-1
YI=YR
YR=DABS(HR(MM,MMM1))+DABS(HI(MM,MMM1))
XI=ZZR
ZZR=XR
XR=DABS(HR(MMM1,MMM1))+DABS(HI(MMM1,MMM1))
IF(YR.LE.EPS*ZZR/YI*(ZZR+XR+XI)) GO TO 85
75 CONTINUE
80 MM=M
C
85 MP1=MM+1
DO 110 I=MP1,NN
  IM1=I-1
  XR=HR(IM1,IM1)
  XI=HI(IM1,IM1)
  YR=HR(I,IM1)
  YI=HI(I,IM1)
  IF(DABS(XR)+DABS(XI).GE.DABS(YR)+DABS(YI)) GO TO 95
  DO 90 J=IM1,N
    ZZR=HR(IM1,J)
    HR(IM1,J)=HR(I,J)
    HR(I,J)=ZZR
    ZZI=HI(IM1,J)
    HI(IM1,J)=HI(I,J)
    HI(I,J)=ZZI
  CONTINUE
  Z=X/Y
  WR(I)=ONE
  GO TO 100
  Z=Y/X
  WR(I)=-ONE
  HR(I,IM1)=ZZR
  HI(I,IM1)=ZZI
  DO 105 J=I,N
    HR(I,J)=HR(I,J)-ZZR*HR(IM1,J)+ZZI*HI(IM1,J)
    HI(I,J)=HI(I,J)-ZZR*HI(IM1,J)-ZZI*HR(IM1,J)
  CONTINUE
110 CONTINUE
C
DO 140 J=MP1,NN
  JM1=J-1
  XR=HR(J,JM1)
  XI=HI(J,JM1)
  HR(J,JM1)=ZERO
  HI(J,JM1)=ZERO
C
C
INTERCHANGE COLUMNS OF HR, HI,
ZR, AND ZI IF NECESSARY
COMPOSITION R*L=H
TRIANGULAR DECOMPOSITION H=L*R
INTERCHANGE COLUMNS OF HR, HI,
ZR, AND ZI IF NECESSARY

```



ELR22310  
 ELR22320  
 ELR22330  
 ELR22340  
 ELR22350  
 ELR22360  
 ELR22370  
 ELR22380  
 ELR22390  
 ELR22400  
 ELR22410  
 ELR22420  
 ELR22430  
 ELR22440  
 ELR22450  
 ELR22460  
 ELR22470  
 ELR22480  
 ELR22490  
 ELR22500  
 ELR22510  
 ELR22520  
 ELR22530  
 ELR22540  
 ELR22550  
 ELR22560  
 ELR22570  
 ELR22580  
 ELR22590  
 ELR22600  
 ELR22610  
 ELR22620  
 ELR22630  
 ELR22640  
 ELR22650  
 ELR22660  
 ELR22670  
 ELR22680  
 ELR22690  
 ELR22700  
 ELR22710  
 ELR22720  
 ELR22730  
 ELR22740  
 ELR22750  
 ELR22760  
 ELR22780  
 ELR22790  
 ELR22800

```

IF (WR(J) .LE. ZERC) GO TO 125
DO 115 I=1, J
  ZZR=HR(I, JMI)
  HR(I, J)=ZZR
  ZZI=HI(I, JMI)
  HI(I, J)=ZZI
CONTINUE
DO 120 I=K, L
  ZZR=ZR(I, JMI)
  ZR(I, J)=ZZR
  ZZI=ZI(I, JMI)
  ZI(I, J)=ZZI
CONTINUE
DO 130 I=1, J
  HR(I, JMI)=HR(I, JMI)+XR*HR(I, J)-XI*HI(I, J)
  HI(I, JMI)=HI(I, JMI)+XR*HI(I, J)+XI*HR(I, J)
CONTINUE
DO 135 I=K, L
  ZR(I, JMI)=ZR(I, JMI)+XR*ZR(I, J)-XI*ZI(I, J)
  ZI(I, JMI)=ZI(I, JMI)+XR*ZI(I, J)+XI*ZR(I, J)
CONTINUE
C 140 CONTINUE
GC TO 40
C 145 WR(NN)=HR(NN, NN)+TR
  WI(NN)=HI(NN, NN)+TI
  NN=NNMI
  GO TO 35
C 150 IF (N .EQ. 1) GO TO 9005
  FNORM=ZERO
  DC 160 I=1, N
    FNORM=FNORM+DABS(WR(I))+DABS(WI(I))
    IF (I .EQ. N) GO TO 160
    IP1=I+1
    DO 155 J=IP1, N
      FNORM=FNORM+DABS(HR(I, J))+DABS(HI(I, J))
    CONTINUE
  155 CONTINUE
  160 IF (FNORM .EQ. ZERO) GO TO 9005
  END INTERCHANGE COLUMNS
  END ACCUMULATE TRANSFORMATIONS
  A ROOT FOJND
  ALL ROOTS FOUND. BACKSUBSTITUTE TO
  FIND VECTORS OF UPPER TRIANGULAR
  FORM
  
```

ELR222810  
 ELR222820  
 ELR222830  
 ELR222840  
 ELR222850  
 ELR222860  
 ELR222870  
 ELR222880  
 ELR222890  
 ELR222900  
 ELR222910  
 ELR222920  
 ELR222930  
 ELR222940  
 ELR222950  
 ELR222960  
 ELR222970  
 ELR222980  
 ELR222990  
 ELR23000  
 ELR23010  
 ELR23020  
 ELR23030  
 ELR23040  
 ELR23050  
 ELR23060  
 ELR23070  
 ELR23080  
 ELR23090  
 ELR23100  
 ELR23110  
 ELR23120  
 ELR23130  
 ELR23140  
 ELR23150  
 ELR23160  
 ELR23170  
 ELR23180  
 ELR23190  
 ELR23200  
 ELR23210  
 ELR23220  
 ELR23230  
 ELR23240  
 ELR23250  
 ELR23260  
 ELR23270  
 ELR23280

```

C      DO NN=N,2,-1
      NP2 = N+2
      DO 180 NM=2,N
        NN=NP2-NM
        XR=WR(NN)
        XI=WI(NN)
        NNM1=NN-1
      DO 175 I=1,NNM1
        I=NN-I
        ZZR=HR(I,NN)
        ZZI=HI(I,NN)
        IF (I .EQ. NNM1) GO TO 170
        IP1=I+1
        DO 165 J=IP1,NNM1
          ZZR=ZZR+HR(I,J)*HR(J,NN)-HI(I,J)*HI(J,NN)
          ZZI=ZZI+HR(I,J)*HI(J,NN)+HI(I,J)*HR(J,NN)
        CONTINUE
        YR=XR-WR(I)
        YI=XI-WI(I)
        Z=Z/Y
        HR(I,NN)=T3(1)
        HI(I,NN)=T3(2)
      CONTINUE
      YR=EPS*FNORM
      END BACKSUBSTITUTION
      VECTOR OF ISOLATED ROOTS
      DO 190 I=1,NM1
        IF (I .GE. K .AND. I .LE. L) GO TO 190
        IP1=I+1
        DO 185 J=IP1,N
          ZR(I,J)=HR(I,J)
          ZI(I,J)=HI(I,J)
        CONTINUE
      CONTINUE
      IF (L .EQ. 0) GO TO 9005
      NPL=N+K
      DO 200 JJ=K,NM1
        J=NPL-JJ
        JM1=J-1
        DO 200 I=K,L
          ZZR=ZR(I,J)
      CONTINUE
      MULTIPLY BY TRANSFORMATION MATRIX
      TO GIVE VECTORS OF ORIGINAL FULL
      MATRIX
      DO J=N,K+1,-1
  
```



```
ELR233290
ELR233300
ELR233310
ELR233320
ELR233330
ELR233340
ELR233350
ELR233360
ELR233370
ELR233380
ELR233390
ELR233400
ELR233410
ELR233420
ELR233430
ELR233440
ELR233450
ELR233460
ELR233470
```

```

ZZI=ZI(I,J)
MM=JMI
IF (L.LT.J) MM=L
DO 195 M=K,MM
  ZZR=ZZR+ZR(I,M)*HR(M,J)-ZI(I,M)*HI(M,J)
  ZZI=ZZI+ZR(I,M)*HR(M,J)+ZI(I,M)*HR(M,J)
CONTINUE
ZR(I,J)=ZZR
ZI(I,J)=ZZI
195 CONTINUE
200 GO TO 9005
C
205 IER=129
  INFER=NN
9000 CONTINUE
  CALL UERTST (IER,6HELRH2C)
9005 RETURN
  END
```

SET ERROR - NO CONVERGENCE TO AN EIGENVALUE AFTER 30 ITERATIONS

```
EBBC0010
EBBC0020
EBBC0030
EBBC0040
EBBC0050
EBBC0060
EBBC0070
EBBC0080
EBBC0090
EBBC0100
EBBC0110
EBBC0120
EBBC0130
EBBC0140
EBBC0150
EBBC0160
EBBC0170
EBBC0180
EBBC0190
EBBC0200
EBBC0210
EBBC0220
EBBC0230
EBBC0240
EBBC0250
EBBC0260
EBBC0270
```

```

SUBROUTINE EBBCKC (ZR,ZI,N,IZ,K,L,M,D)
-----LIBRARY 1-----
FUNCTION
USAGE
PARAMETERS ZR ZI
N
IZ
K
L
M
D
-EBBCKC-----
- BACKTRANSFORM THE EIGENVECTORS OF A BALANCED
  COMPLEX GENERAL MATRIX.
- CALL EBBCKC (ZR,ZI,N,IZ,K,L,M,D)
- INPUT/OUTPUT MATRICES OF DIMENSION N BY M.
  ON INPUT THE FIRST M COLUMNS OF ZR AND
  ZI CONTAIN THE REAL AND IMAGINARY PARTS,
  RESPECTIVELY, OF THE EIGENVECTORS TO BE
  BACK TRANSFORMED. ON OUTPUT THESE M
  COLUMNS CONTAIN THE REAL AND IMAGINARY
  PARTS OF THE TRANSFORMED EIGENVECTORS.
- INPUT SCALAR CONTAINING THE NUMBER OF
  ROWS IN THE MATRIX Z = (ZR,ZI). N MUST
  NOT BE GREATER THAN IZ.
- INPUT SCALAR CONTAINING THE ROW DIMENSION
  OF MATRICES ZR AND ZI IN THE CALLING
  PROGRAM.
- INPUT SCALARS CONTAINING THE BOUNDARY
  INDICES FOR THE BALANCED MATRIX. K AND L
  ARE TWO OUTPUT PARAMETERS FROM ROUTINE
  EBALAC.
- INPUT SCALAR CONTAINING THE NUMBER OF COLUMNS
  OF Z = (ZR,ZI) TO BE BACK TRANSFORMED.
- INPUT VECTOR OF LENGTH N CONTAINING THE
```







```

//EIG$FCN JOB (1719,0947,AX74), 'SMC 1882', TIME=2
// EXEC FORTCLGM
// FORT.SYSIN DD *

```

```

.....
PROGRAM EIGFCN
PERTURBATION VELOCITY PLOT PROGRAM
NI = 0

```

```

PURPOSE

```

```

.....
TO PLOT THE NONDIMENSIONALIZED PERTURBATION VELOCITY U AGAINST
NONDIMENSIONALIZED RADIUS UTILIZING THE DATA GENERATED
BY PROGRAM PIPEO (MODE=1). PLOTTING IS PERFORMED ON
THE NPS VERSATEC PLOTTER USING SUBROUTINE PLOTG.
.....

```

```

.....
IMPLICIT REAL*8(A-H,O-Z)
COMPLEX *16OPRIM(85), DQPRIM(85), ALPHA, UP(85), UPRIM(85), CONST, GAMMA

```

```

1A REAL *BETA(85), U(85), UR(85), UI(85)
REAL *4URI(85), U1(85), RAD1(85), SREY, SLAMDA, SGAMMA, AR, AI

```

```

READ N, REY, ALPHA, LAMDA, GAMMA & QPRIME'S

```

```

READ (5,6) N, REY, ALPHA

```

```

NO = N+1

```

```

N1 = N+2

```

```

READ (5,7) AMDA, GAMMA

```

```

READ (5,8) KSET

```

```

SLAMDA = AMDA

```

```

SREY = REY

```

```

AR = ALPHA

```

```

AI = AIMAG(ALPHA)

```

```

SGAMMA = GAMMA

```

```

SGAMMA = SGAMMA + AR

```

```

DC 1 I=2, NO

```

```

READ (5,9) ETA(1), QPRIM(1)

```

```

1 CONTINUE

```

```

EIGF 10
EIGF 20
EIGF 30
EIGF 40
EIGF 50
EIGF 60
EIGF 70
EIGF 80
EIGF 90
EIGF 100
EIGF 110
EIGF 120
EIGF 130
EIGF 140
EIGF 150
EIGF 160
EIGF 170
EIGF 180
EIGF 190
EIGF 200
EIGF 210
EIGF 220
EIGF 230
EIGF 240
EIGF 250
EIGF 260
EIGF 270
EIGF 280
EIGF 290
EIGF 300
EIGF 310
EIGF 320
EIGF 330
EIGF 340
EIGF 350
EIGF 360
EIGF 370
EIGF 380
EIGF 390
EIGF 400
EIGF 410
EIGF 420
EIGF 430
EIGF 440
EIGF 450

```



```

C      COMPUTE UPPRIME'S
C      DEL = 1.00/DFLOAT (N+1)
C      ETA(1) = 0.000
C      QPRIM(1) = 1.79591836700*QPRIM(2) - 1.2478134100*QPRIM(3) + 0.60641399
C      UPPRIM(1) = 2.00*QPRIM(1)
C      CALL COEFNT (ETA(2), AMDA, COEF, KSET)
C      DQPRIM(2) = 2.00*(-QPRIM(2)+QPRIM(3))/(3.00*DEL)
C      DQPRIM(2) = COEF*DQPRIM(2)
C      UPPRIM(2) = 2.00*QPRIM(2)+ETA(2)*DQPRIM(2)
C      CALL COEFNT (ETA(3), AMDA, COEF, KSET)
C      DQPRIM(3) = 4.00*(-QPRIM(2)+QPRIM(3))/(3.00*DEL)
C      DQPRIM(3) = COEF*DQPRIM(3)
C      UPPRIM(3) = 2.00*QPRIM(3)+ETA(3)*DQPRIM(3)
C
C      DO 2 I=4,N
C      CALL COEFNT (ETA(I), AMDA, COEF, KSET)
C      DQPRIM(I) = (QPRIM(I+1)-QPRIM(I-1))/(2.00*DEL)
C      DQPRIM(I) = COEF*DQPRIM(I)
C      UPPRIM(I) = 2.00*QPRIM(I)+ETA(I)*DQPRIM(I)
C      2 CONTINUE
C
C      CALL COEFNT (ETA(NO), AMDA, COEF, KSET)
C      DQPRIM(NO) = -QPRIM(NO)/(2.00*DEL)
C      DQPRIM(NO) = COEF*DQPRIM(NO)
C      UPPRIM(NO) = 2.00*QPRIM(NO)+ETA(NO)*DQPRIM(NO)
C
C      ETA(N1) = 1.000
C      UPPRIM(N1) = (0.00,0.00)
C      WRITE (6,10)
C      DETERMINE U VECTOR OF LARGEST MAGNITUDE
C      C = 0.00
C      DO 3 I=1,N1
C      IF (CDABS(UPPRIM(I)).GT.C) INDEX=I
C      IF (CDABS(UPPRIM(I)).GT.C) C=CDABS(UPPRIM(I))
C      3 CONTINUE
C      CONST = DCONJG(UPPRIM(INDEX))/C**2

```

```

EIGF 460
EIGF 470
EIGF 480
EIGF 490
EIGF 500
EIGF 510
EIGF 520
EIGF 530
EIGF 540
EIGF 550
EIGF 560
EIGF 570
EIGF 580
EIGF 590
EIGF 600
EIGF 610
EIGF 620
EIGF 630
EIGF 640
EIGF 650
EIGF 660
EIGF 670
EIGF 680
EIGF 690
EIGF 700
EIGF 710
EIGF 720
EIGF 730
EIGF 740
EIGF 750
EIGF 760
EIGF 770
EIGF 780
EIGF 790
EIGF 800
EIGF 810
EIGF 820
EIGF 830
EIGF 840
EIGF 850
EIGF 860
EIGF 870
EIGF 880
EIGF 890
EIGF 900
EIGF 910
EIGF 920
EIGF 930

```

```

CCGC
NORMALIZE UPRIMES'S AND SPLIT INTO REAL & IMAGINARY VECTORS
DO 4 I=1,N1
UP(I) = CONST*UPPRIM(I)
UR(I) = UP(I)
UI(I) = (ODO,-1DO)*UP(I)
4 CONTINUE
CCGCC
CONVERT U'S AND ETA'S TO SINGLE PRECISION FCR PLOTG
DO 5 I=1,N1
RADI(I) = ETA(I)
URI(I) = UR(I)
UII(I) = UI(I)
WRITE (6,I1) RADI(I),URI(I),UII(I)
5 CONTINUE
CCGCC
PLOT RESULTS
CALL PLOTG(RADI,URI,N1,1,1,1,'RADIUS',6,'PERTURBATION VELOCITY',
$ 21,0,1,1,1,1,7,7,7)
CALL PLOTG(RADI,URI,N1,2,1,5,'RADIUS',6,'PERTURBATION VELOCITY',
$ 21,0,1,1,1,1,7,7,7)
CALL CHART (N,SREY,AR,AI,SGAMMA,SLAMDA)
CALL PLOT (0.0,0.0,999)
STOP
C
6 FORMAT (I2,3D20.10)
7 FORMAT (F15.7,2(I PD20.10))
8 FORMAT (I2)
9 FORMAT (F15.7,2(I PD20.10))
10 FORMAT ('I,')
11 FORMAT ('',3F15.7)
END
CCGCC
SUBROUTINE COEFNT(ETA,AMDA,COEF,KSET).....COEF
COEF 10
COEF 20
COEF 30
COEF 40
COEF 50
COEF 60
COEF 70
PURPOSE--WHEN AN OFFSET MESH IS USED, THIS SUBROUTINE GENERATES
THE COEFFICIENT REQUIRED TO CONVERT DQ/DETA TO DQ/DR AND
CONVERTS THE UNIFORM ETA VALUE INTO THE NONUNIFORM R VALUE

```

```

EIGF 940
EIGF 950
EIGF 960
EIGF 970
EIGF 980
EIGF 990
EIGF 1000
EIGF 1010
EIGF 1020
EIGF 1030
EIGF 1040
EIGF 1050
EIGF 1060
EIGF 1070
EIGF 1080
EIGF 1090
EIGF 1100
EIGF 1110
EIGF 1120
EIGF 1130
EIGF 1140
EIGF 1150
EIGF 1160
EIGF 1170
EIGF 1180
EIGF 1190
EIGF 1200
EIGF 1210
EIGF 1220
EIGF 1230
EIGF 1240
EIGF 1250
EIGF 1260
EIGF 1270
EIGF 1280
EIGF 1290
EIGF 1300
EIGF 1310
EIGF 1320

```

COEF 80  
 COEF 90  
 COEF 110  
 COEF 120  
 COEF 130  
 COEF 140  
 COEF 150  
 COEF 160  
 COEF 170  
 COEF 180  
 COEF 190  
 COEF 200  
 COEF 210  
 COEF 220  
 COEF 230  
 COEF 240  
 COEF 250  
 COEF 260  
 COEF 270  
 COEF 280  
 COEF 290  
 COEF 300  
 COEF 310  
 COEF 320  
 COEF 330  
 COEF 340  
 COEF 350  
 COEF 360  
 COEF 370  
 COEF 380  
 COEF 390  
 COEF 400  
 COEF 410  
 COEF 420  
 COEF 430  
 COEF 440  
 COEF 450  
 COEF 460  
 COEF 470  
 COEF 480  
 COEF 490  
 COEF 500  
 COEF 510  
 COEF 520  
 COEF 530  
 COEF 540  
 COEF 550

EXAMPLE OF THE CALLING ARGUMENT  
 CALL COEFNT(ETA,AMDA,COEF,KSET)

DESCRIPTION OF PARAMETERS  
 ETA --THE VALUE OF THE INDEPENDENT VARIABLE REPLACING R  
 IN THE NONUNIFORM MESH FOR THE STATION BEING COMPUTED.  
 AMDA--THE NONDIMENSIONAL MESH PARAMETER FOR THE PARTICULAR  
 SET OF DATA BEING PLOTTED.  
 COEF--THE VALUE OF (DR/DELTA)\*\*-1 NEEDED TO CONVERT DQ/DELTA TO  
 DQ/DR.  
 KSET-OFFSET PARAMETER WHICH IS EQUAL TO -1 FOR WALL OFFSET,  
 0 FOR UNIFORM MESH AND 1 FOR AXIS OFFSET.

OTHER SUBROUTINES NEEDED  
 NONE

.....

SUBROUTINE COEFNT (ETA,AMDA,COEF,KSET)

IMPLICIT REAL\*8(A-H,O-Z)

TETA = ETA  
 IF (AMDA.LE.1D-10) GO TO 3  
 IF (KSET.EQ.1) TETA=1D0-TETA  
 CNST = DTANH(AMDA)/AMDA  
 ETAP = AMDA\*TETA  
 COEF = CNST\*(DCOSH(ETAP))\*\*2  
 IF (AMDA.GE.1D-10) CONST=1D0/DTANH(AMDA)  
 IF (KSET) 1,3,2

1 ETA = CONST\*DTANH(ETAP)  
 RETURN

2 ETA = 1.D0-CONST\*DTANH(ETAP)  
 RETURN

C  
 C  
 C



COEF 560  
 COEF 570  
 COEF 580

CHAR 10  
 CHAR 20  
 CHAR 30  
 CHAR 40  
 CHAR 50  
 CHAR 60  
 CHAR 70  
 CHAR 80  
 CHAR 90  
 CHAR 100  
 CHAR 110  
 CHAR 120  
 CHAR 130  
 CHAR 140  
 CHAR 150  
 CHAR 160  
 CHAR 170  
 CHAR 180  
 CHAR 190  
 CHAR 200  
 CHAR 210  
 CHAR 220  
 CHAR 230  
 CHAR 240  
 CHAR 250  
 CHAR 260  
 CHAR 270  
 CHAR 280  
 CHAR 290  
 CHAR 300  
 CHAR 310  
 CHAR 320  
 CHAR 330  
 CHAR 340  
 CHAR 350  
 CHAR 360  
 CHAR 370  
 CHAR 380  
 CHAR 390  
 CHAR 400  
 CHAR 410  
 CHAR 420  
 CHAR 430

3 COEF = 100  
 RETURN  
 END

.....SUBROUTINE CHART(N,SREY,AR,AI,SGAMMA,SLAMDA).....  
 PURPOSE  
 TO LABEL THE GRAPH WITH INFORMATION PERTAINING TO THE PLOT  
 EXAMPLE OF THE CALLING ARGUMENT  
 CALL CHART(N,SREY,AR,AI,SGAMMA,SLAMDA)  
 DESCRIPTION OF PARAMETERS  
 THE PARAMETERS ARE SELF-EXPLANATORY AND MUST BE IN SINGLE  
 PRECISION FOR PLOTTING.  
 OTHER SUBROUTINES NEEDED  
 ONLY BUILT-IN VERSATEC PLOTTING FUNCTIONS NEWPEN,SYMBOL &  
 NUMBER. NOTE THAT THESE ROUTINES MAY ONLY BE  
 ACCESSED WHEN RUNNING UNDER 'FORTCLGW'.  
 .....  
 SUBROUTINE CHART (N,SREY,AR,AI,SGAMMA,SLAMDA)  
 X0 = 2.5  
 Y0 = 6.5  
 HT = 0.15  
 HT1 = 0.7\*HT  
 DELY1 = .08\*HT  
 DELY2 = .065\*HT1  
 DELX = .1  
 GRAPH TITLE  
 CALL NEWPEN (2)  
 CALL SYMBOL (X0,Y0,HT,'NORMALIZED PERTURBATION VELOCITY',0.,32)

CC  
 CCC  
 CCCCC

CHAR 440  
 CHAR 450  
 CHAR 460  
 CHAR 470  
 CHAR 480  
 CHAR 490  
 CHAR 500  
 CHAR 510  
 CHAR 520  
 CHAR 530  
 CHAR 540  
 CHAR 550  
 CHAR 560  
 CHAR 570  
 CHAR 580  
 CHAR 590  
 CHAR 600  
 CHAR 610  
 CHAR 620  
 CHAR 630  
 CHAR 640  
 CHAR 650  
 CHAR 660  
 CHAR 670  
 CHAR 680  
 CHAR 690  
 CHAR 700  
 CHAR 710  
 CHAR 720  
 CHAR 730  
 CHAR 740  
 CHAR 750  
 CHAR 760  
 CHAR 770  
 CHAR 780  
 CHAR 790  
 CHAR 800  
 CHAR 810  
 CHAR 820  
 CHAR 830  
 CHAR 840  
 CHAR 850  
 CHAR 860  
 CHAR 870  
 CHAR 880  
 CHAR 890  
 CHAR 900

```

XO = XO+7.*DELX
YO = YO-DELY1
CALL SYMBOL(XO, YO, HT, ' FOR THE CASE N = 0', 0., 0., 18)
MESH VALUE
CALL NEWPEN (1)
XO = XO+4.*DELX
YO = YO-DELY1
SN = FLOAT(N)
CALL SYMBOL(XO, YO, HT1, 'NMESH = ', 0., 9)
CALL NUMBER (999., 999., HT1, SN, 0., -1)
REY VALUE
YC = YO-DELY2
CALL SYMBOL(XO, YO, HT1, 'REY = ', 0., 9)
CALL NUMBER (999., 999., HT1, SREY, 0., -1)
ALPHA VALUE
X1 = XO+11.*DELY2
Y1 = YO-DELY2
CALL SYMBOL(XO, YO, HT1, 'ALPHA = ', 0., 9)
CALL NUMBER (999., 999., HT1, AR, 0., 1)
CALL NUMBER (X1, Y1, HT1, AI, 0., 1)
GAMMA RL* VALUE
YO = YO-DELY2
CALL SYMBOL(XO, YO, HT1, 'GAMMA* = ', 0., 9)
CALL NUMBER (999., 999., HT1, SGAMMA, 0., 4)
LAMBDA VALUE
YO = YO-DELY2
CALL SYMBOL(XO, YO, HT1, 'LAMBDA = ', 0., 9)
CALL NUMBER (999., 999., HT1, SLAMDA, 0., 1)
SYMBOL LEGEND
YO = 1.75
CALL SYMBOL(XO, YO, HT1, 'OCTAGON = U(REAL)', 0., 17)
YO = YO-DELY2
CALL SYMBOL(XO, YO, HT1, 'DIAMOND = U(IMAG)', 0., 17)
RETURN
END
  
```







```

C      CALL SEARCH (-1,X1,Y1,NPLT1,NDIM)
C      CALL SEARCH (0,X2,Y2,NPLT2,NDIM)
C      CALL SEARCH (1,X3,Y3,NPLT3,NDIM)
C      JUMP TO PLOT LABEL ROUTINE IF NO INCIPIENT POINTS
C      IF (NPLT1) 8,8,2
C      IF POINTS COMPUTED, NEW PAGE AND WRITE THEM OUT
C      2 WRITE (6,11)
C      DC 3 I=1,NPLT1
C      WRITE (6,12) X1(I),Y1(I)
C      3 CONTINUE
C      PLOT INCIPIENT INSTABILITY POINTS
C      CALL PLOTG(X1,Y1,NPLT1,1,0,1,'ALPHA REAL',10,'ALPHA IMAGINARY',15,
C      $ XMIN,XMAX,YMIN,YMAX,7,7)
C      LEGEND FOR INCIPIENT SYMBOL
C      CALL NEWPEN (1)
C      CALL SYMBOL (1.3,0.7,.1,'OCTAGON = INCIPIENT INSTABILITY',0.,32)
C      JUMP TO PLOT LABEL ROUTINE IF NO CRITICAL POINTS
C      IF (NPLT2) 8,8,4
C      IF POINTS COMPUTED, NEW PAGE AND PRINT THEM OUT
C      4 WRITE (6,11)
C      DO 5 I=1,NPLT2
C      WRITE (6,12) X2(I),Y2(I)
C      5 CONTINUE
C      PLOT CRITICAL POINTS
C      CALL PLOTG(X2,Y2,NPLT2,2,0,2,'ALPHA REAL',10,'ALPHA IMAGINARY',15,
C      $ XMIN,XMAX,YMIN,YMAX,7,7)
C      LEGEND FOR CRITICAL SYMBOL

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STBC 460
STBC 470
STBC 480
STBC 490
STBC 500
STBC 510
STBC 520
STBC 530
STBC 540
STBC 550
STBC 560
STBC 570
STBC 580
STBC 590
STBC 600
STBC 610
STBC 620
STBC 630
STBC 640
STBC 650
STBC 660
STBC 670
STBC 680
STBC 690
STBC 700
STBC 710
STBC 720
STBC 730
STBC 740
STBC 750
STBC 760
STBC 770
STBC 780
STBC 790
STBC 800
STBC 810
STBC 820
STBC 830
STBC 840
STBC 850
STBC 860
STBC 870
STBC 880
STBC 890
STBC 900
STBC 910
STBC 920
STBC 930

```



```

C C C CALL NEWPEN (1)
C C C CALL SYMBOL(1.3,.53,.1,'TRIANGLE = CRITICAL INSTABILITY',0.,31)
C C C JUMP TO PLOT LABEL ROUTINE IF NO FULLY DEVELOPED POINTS
C C C IF (NPLT3) 8,8,6
C C C IF POINTS COMPUTED, NEW PAGE AND PRINT THEM OUT
C C C € WRITE (6,11)
C C C DC 7 I=1,NPLT3
C C C WRITE (6,12) X3(1),Y3(1)
C C C 7 CONTINUE
C C C PLOT FULLY DEVELOPED POINTS
C C C CALL PLOTG(X3,Y3,NPLT3,3,0,5,'ALPHA REAL',10,'ALPHA IMAGINARY',15,
C C C $ XMIN,XMAX,YMIN,YMAX,7.,7.)
C C C LEGEND FOR FULLY DEVELOPED SYMBOL
C C C CALL NEWPEN (1)
C C C CALL SYMBOL(1.3,.36,.1,'DIAMOND = FULLY DEVELOPED INSTABILITY',
C C C $ C.,38)
C C C LABEL THE PLOT
C C C 8 CALL CHART (SN,SREY,SLAMDA)
C C C CALL PLOT (0.,0.,999)
C C C STOP
C C C 9 FORMAT (I2)
C C C 10 FORMAT (3E20.10)
C C C 11 FORMAT ('I')
C C C 12 FORMAT ('.',2E20.10)
C C C END
C C C .....SUBROUTINE SEARCH(NCASE, X, Y, NDIM).....
C C C PURPOSE
C C C TO SCAN THE STABILITY MAP FOR CHANGES OF SIGN WITH RESPECT
C C C TO A SPECIFIED STABILITY VALUE AND GENERATE AN ARRAY OF X, Y
C C C POINTS DEFINING A CONTOUR OF THE SPECIFIED STABILITY.
C C C SEAR 10
C C C SEAR 30
C C C SEAR 40
C C C SEAR 50
C C C SEAR 60
C C C SEAR 70
C C C SEAR 80

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90  
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SAMPLE OF THE CALLING ARGUMENT
CALL SEARCH(NCASE,X,Y,NDIM)
DESCRIPTION OF PARAMETERS
NCASE - DEFINES THE INSTABILITY CASE (INCIPIENT, CRITICAL OR
FULLY DEVELOPED) TO BE USED WHEN GENERATING THE X,Y
ARRAYS.
NCASE = -1 INCIPIENT INSTABILITY CRITERION
NCASE = 0 CRITICAL INSTABILITY CRITERION
NCASE = 1 FULLY DEVELOPED INSTABILITY CRITERION
X - THE ARRAY OF ALPHA REAL COORDINATES DEFINING THE
LOCATION OF THE POINTS OF SPECIFIED INSTABILITY.
Y - THE ARRAY OF ALPHA IMAGINARY COORDINATES DEFINING THE
LOCATION OF THE POINTS OF THE SPECIFIED INSTABILITY.
NDIM - THE ORDER OF THE MAP ARRAY.
OTHER ROUTINES NEEDED
STATEMENT FUNCTION CRIT AND SUBROUTINE INTERP.
.....
SUBROUTINE SEARCH (NCASE,X,Y,K,NDIM)
DIMENSION X(500), Y(500)
COMMON /ARRAY/ G(41,41),AR(41),AI(41)
DEFINE THE STATEMENT FUNCTION CRIT(NCASE,ALPHA)
CRIT(NCASE,ALPHA) = FLOAT(NCASE)*ABS(ALPHA)
K = 0
NDIM = NDIM-1
SEARCH FOR SIGN CHANGES BY COLUMN & INTERPOLATE FOR
ALPHA IMAGINARY AT WHICH SIGN CHANGE OCCURS
DO 5 I=1,NDIM
DO 5 J=1,NDIM
IF (G(I,J)-CRIT(NCASE,AR(I))) 2,4,1
1 IF (G(I,J+1)-CRIT(NCASE,AR(I))) 3,3,5
2 IF (G(I,J+1)-CRIT(NCASE,AR(I))) 5,3,3

```











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CHAR 380
CHAR 390
CHAR 400
CHAR 410
CHAR 420
CHAR 430
CHAR 440
CHAR 450
CHAR 460
CHAR 470
CHAR 480
CHAR 490
CHAR 500
CHAR 510
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CHAR 680
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CHAR 790
CHAR 800
CHAR 810
CHAR 820
CHAR 830

```

```

DELX = .1
GRAPH TITLE
CALL NEWPEN (2)
CALL SYMBOL (X0, Y0, HT, 'STABILITY CONTOUR PLOT', 0., 22)
X0 = X0 + 3. * DELX
YC = YO - DELY1
CALL SYMBOL (X0, Y0, HT, 'FOR THE CASE N = 0', 0., 18)
MESH VALUE
CALL NEWPEN (1)
X0 = X0 + 4. * DELX
Y0 = YO - DELY1
CALL SYMBOL (X0, Y0, HT1, 'NMESH = ', 0., 9)
CALL NUMBER (999., 999., HT1, SN, 0., -1)
REY VALUE
Y0 = YO - DELY2
CALL SYMBOL (X0, Y0, HT1, 'REY = ', 0., 9)
CALL NUMBER (999., 999., HT1, SREY, 0., -1)
LAMBDA VALUE
YC = YO - DELY2
CALL SYMBOL (X0, Y0, HT1, 'LAMBDA = ', 0., 9)
CALL NUMBER (999., 999., HT1, SLAMDA, 0., 1)
STABILITY AREA LABELS
NOTE - SINCE THE SHAPE OF THE CURVE VARIES WITH
      EACH SET OF INPUT DATA, THE COORDINATES OF THE FOLLOWING
      LABELS MUST BE ADJUSTED FOR EACH SPECIFIC PLOT.
CALL NEWPEN (2)
CALL SYMBOL (4.0, 4.5, HT1, 'SUPERCRITICAL', 0., 13)
CALL SYMBOL (5.6, 4.5, HT1, 'SUBCRITICAL', 0., 11)
CALL SYMBOL (6.9, 4.5, HT1, 'STABLE', 0., 6)
YC = 4.5 - DELY2
CALL SYMBOL (4.0, YO, HT1, 'INSTABILITY', 0., 12)
CALL SYMBOL (5.6, YO, HT1, 'INSTABILITY', 0., 11)
RETURN
END

```

C C

C C C

C C C

C C C C

C C C C C

C C C

.....  
THE FOLLOWING CARDS CCMPRISE THE CATA DECK FOR PROGRAM STBCONT.  
.....  
/\*  
//GO.SYSIN DD \*  
.....  
DATA DECK FROM ONE RUN OF PROGRAM PIPEO (MODEND = 2)  
.....  
/\*  
.....



### LIST OF REFERENCES

1. Davey, A., and Drazin, P.G., "The Stability of Poiseuille Flow in a Pipe," Journal of Fluid Mechanics, v. 36, part 2, p. 209, 22 August 1968.
2. Garg, V.K., and Rouleau, W.T., "Linear Spatial Stability of Pipe Poiseuille Flow," Journal of Fluid Mechanics, v. 54, part 1, p. 113, 6 January 1969.
3. Gill, A.E., "The Least-Damped Disturbance to Poiseuille Flow in a Circular Pipe," Journal of Fluid Mechanics, v. 61, part 1, p. 765, 3 December 1973.
4. Harrison, W.F., On the Stability of Poiseuille Flow, Ae. E. Thesis, Naval Postgraduate School, Monterey, California, 1975.
5. Huang, L.M. and Chen, T.S., "Stability of Developing Flow Subject to Non-axisymmetric Disturbances," Journal of Fluid Mechanics, v. 63, part 1, p. 183, 16 April 1973.
6. Johnston, R.H. III, A Program for the Stability Analysis of Pipe Poiseuille Flow, M.S. Thesis, Naval Postgraduate School, Monterey, California, 1976.
7. Leite, R.J., An Experimental Investigation of Axially Symmetric Poiseuille Flow, Report No. OSR-TR-56-2, Air Force Contract AF18(600)-350, November, 1956.
8. Naval Postgraduate School Report NPS-67Gn77051, Improved Finite Difference Formulas for Boundary Value Problems, by T.H. Gawain and R.E. Ball, 1 May 1977.
9. Naval Postgraduate School Report NPS67-78-006, A Basic Reformulation of the Pipe Flow Stability Problem and Some Preliminary Numerical Results, by T.H. Gawain, 1 September 1978.
10. Reynolds, O., "An Experimental Investigation of the Circumstances which Determine whether the Motion of Water Shall be Direct or Sinuous, and the Law of Resistance in Parallel Channels," Phil. Trans. Royal Soc., 174, p. 935-982, 1883.
11. Salwen, H., and Grosch, C.E., "The Stability of Poiseuille Flow in a Pipe of Circular Cross-section," Journal of Fluid Mechanics, v. 54, part 1, p. 93, 6 March 1972.

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