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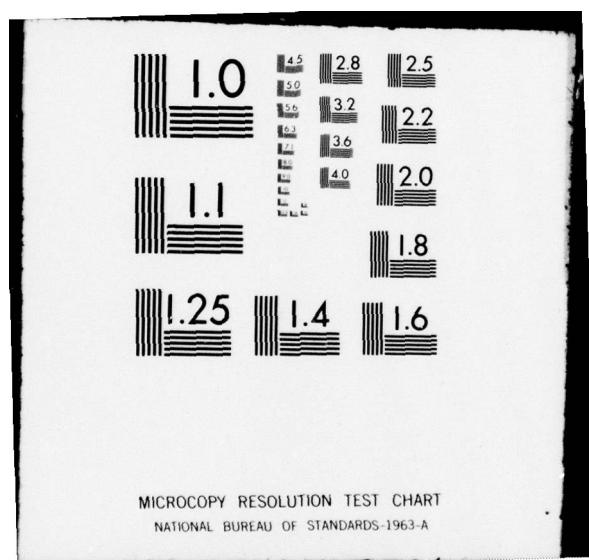
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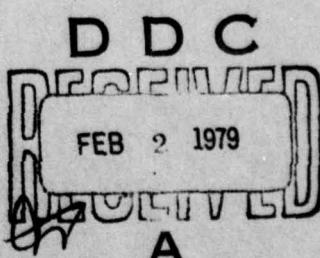
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A METHOD FOR CALCULATING WAVE PACKET  
TRAJECTORIES AND WAVE HEIGHTS. PART II.

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## ABSTRACT

In Part I a method for calculating wave packet trajectories and wave heights is based on the assumption that the water depth contours are locally parallel in the vicinity of each ray point. This method is extended in order to predict the modification to surface gravity water waves in shoaling water when the water depth contours are not parallel. The calculations are greatly simplified by choosing a coordinate system at each ray point in which one axis is aligned parallel with the direction of the gradient of the water depth. Example printouts and plots are presented to illustrate the wave prediction method. It is discovered that when waves initially approach sinuous water depth contours symmetrically with respect to the beach there can be more energy in the bays than at the headlands.

## ACKNOWLEDGMENTS

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## CHAPTER I INTRODUCTION

This is a supplement to the report by Breeding, J. Ernest, Jr., K.C. Matson, and Nourollah Riahi, "A Method for Calculating Wave Packet Trajectories and Wave Heights," Department of Oceanography, Florida State University, Tallahassee, Report No. JEB-1 (1978) hereafter referred to as Part I. The program for predicting the modification to waves in shoaling water described in Part I is developed on the assumption that the water depth contours are locally parallel in the vicinity of each ray point. This method results in accurate wave predictions for many examples of bottom topography. However, the wave height predictions are not accurate when the water depth contours are highly non-parallel in the region of a ray point.

The objective in this report is to extend Part I in order to make accurate wave predictions when the water depth contours are not parallel. The theory is developed in Chapter II. The calculations are greatly simplified by making them in a  $x'y'$ -coordinate system which is chosen so that at each ray point the positive  $x'$ -axis is in the direction of the gradient of the water depth. As a result, the first partial derivatives of the water depth and wave velocities with respect to  $y'$  vanish. Further, there is a simplification in the second partial derivatives involving  $y'$ . In Chapter III modifications to the computer program are described. To illustrate the wave prediction program two examples of nonparallel water depth contours are presented in Chapter IV. An interesting result is that when waves initially approach sinuous water depth contours symmetrically with respect to the beach (headon) it is possible for there to be more energy in the bays than at the headlands. This is in contrast to the result for monochromatic trajectories where under similar conditions there is always more energy predicted at the headlands than in the bays.

CHAPTER II THEORY FOR NONPARALLEL WATER DEPTH CONTOURS

2.1 Spatial Derivatives of the Water Depth. At each point of a wave packet trajectory the calculations are made in a  $x'y'$ -coordinate system where the  $x'$ -axis is taken in the direction of the gradient of the water depth. The particulars of a wave packet trajectory are tabulated in a  $xy$ -coordinate system which retains a fixed orientation with respect to the water depth grid. The relationships between these coordinate systems and a specific ray point for a set of nonparallel water depth contours are shown in Figure (2-1). Equations relating these coordinate systems are given by

$$x' = x \cos \alpha + y \sin \alpha \quad (2-1)$$

$$y' = -x \sin \alpha + y \cos \alpha \quad (2-2)$$

$$\tan \alpha = \frac{\partial h}{\partial y} / \frac{\partial h}{\partial x} \quad (2-3)$$

where  $\alpha$  is the angle by which the  $x'$ -axis is rotated with respect to the  $x$ -axis and  $h$  is the water depth.

The partial derivatives of  $h$  in the  $x'y'$ -coordinate system with respect to the partial derivatives of  $h$  in the  $xy$ -coordinate system are given by

$$\frac{\partial h}{\partial x'} = \frac{\partial h}{\partial x} \cos \alpha + \frac{\partial h}{\partial y} \sin \alpha \quad (2-4)$$

$$\frac{\partial h}{\partial y'} = -\frac{\partial h}{\partial x} \sin \alpha + \frac{\partial h}{\partial y} \cos \alpha = 0 \quad (2-5)$$

$$\frac{\partial^2 h}{(\partial x')^2} = \frac{\partial^2 h}{\partial x^2} \cos^2 \alpha + 2 \frac{\partial^2 h}{\partial x \partial y} \sin \alpha \cos \alpha + \frac{\partial^2 h}{\partial y^2} \sin^2 \alpha \quad (2-6)$$

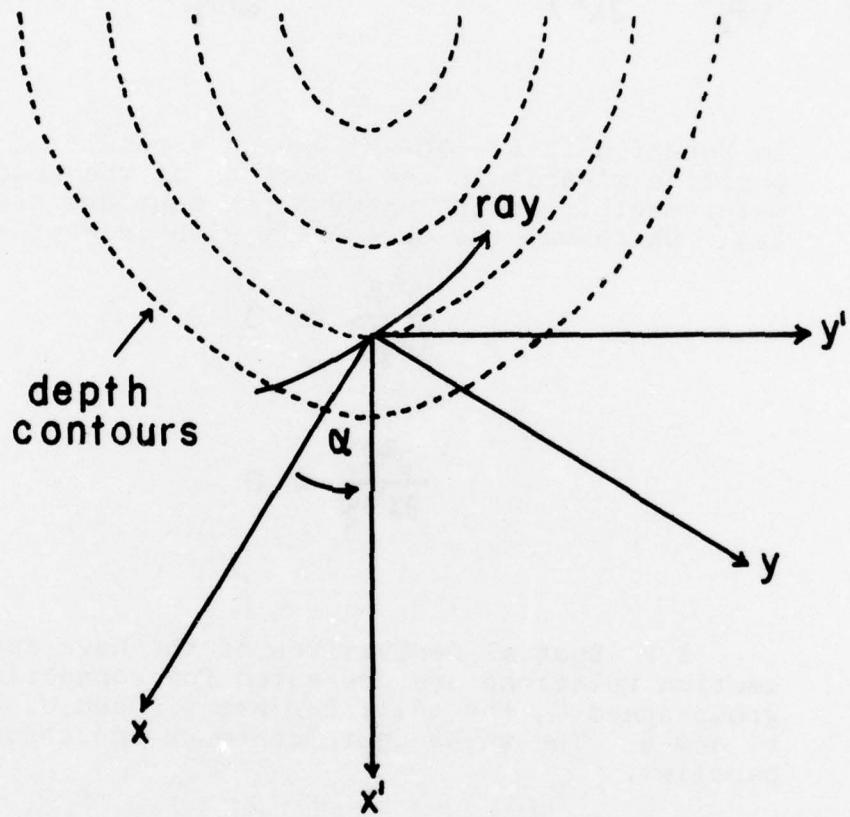


Figure (2-1). Relationships between the coordinate systems and the water depth contours.

$$\frac{\partial^2 h}{(\partial y')^2} = \frac{\partial^2 h}{\partial x^2} \sin^2 \alpha - 2 \frac{\partial^2 h}{\partial x \partial y} \sin \alpha \cos \alpha + \frac{\partial^2 h}{\partial y^2} \cos^2 \alpha \quad (2-7)$$

$$\frac{\partial^2 h}{\partial x' \partial y'} = \left( \frac{\partial^2 h}{\partial y^2} - \frac{\partial^2 h}{\partial x^2} \right) \sin \alpha \cos \alpha + \frac{\partial^2 h}{\partial x \partial y} (\cos^2 \alpha - \sin^2 \alpha) \quad (2-8)$$

In Equation (2-5),  $\partial h / \partial y' = 0$  as a result of choosing the positive  $x'$ -axis in the direction of the gradient of the water depth. If the water depth contours are locally parallel, which was the case dealt with in Part I, then

$$\frac{\partial^2 h}{(\partial y')^2} = 0 \quad (2-9)$$

$$\frac{\partial^2 h}{\partial x' \partial y'} = 0 \quad (2-10)$$

2.2 Spatial Derivatives of the Wave Speeds. In this section relations are presented for connecting the geometric group speed  $G$ , the collinear group speed  $U$ , the phase speed  $v$ , and  $h$ . The water depth contours are assumed to be non-parallel.

#### a. Derivatives of $v$

The phase speed of a surface gravity water wave is defined (Part I)

$$v = \frac{1}{a} \tanh \frac{1}{2} \quad (2-11)$$

in which

$$a = \frac{2\pi}{gT} \quad (2-12)$$

$$b = \frac{T}{4\pi} \quad (2-13)$$

$$I = \frac{h}{bv} \quad (2-14)$$

where  $T$  is the wave period and  $g$  is the acceleration due to gravity.

The first partial derivatives of  $v$  in the  $x'y'$ -coordinate system are given by

$$\frac{\partial v}{\partial x'} = w \frac{\partial h}{\partial x'} \quad (2-15)$$

$$\frac{\partial v}{\partial y'} = w \frac{\partial h}{\partial y'} = 0 \quad (2-16)$$

where

$$w = \frac{v(1 - \alpha^2 v^2)}{[2\alpha b v^2 + h(1 - \alpha^2 v^2)]} \quad (2-17)$$

The second partial derivatives of  $v$  are defined by

$$\frac{\partial^2 v}{(\partial x')^2} = w \left[ \frac{\partial^2 h}{(\partial x')^2} + \gamma \left( \frac{\partial h}{\partial x'} \right)^2 \right] \quad (2-18)$$

$$\frac{\partial^2 v}{(\partial y')^2} = w \frac{\partial^2 h}{(\partial y')^2} \quad (2-19)$$

$$\frac{\partial^2 v}{\partial x' \partial y'} = w \frac{\partial^2 h}{\partial x' \partial y'} \quad (2-20)$$

where

$$\gamma = - \frac{4abv^2}{[2abv^2 + h(1-a^2v^2)]^2} \quad (2-21)$$

b. Derivatives of U

The collinear group speed of a surface gravity water wave is defined (Part I)

$$U = \frac{1}{2} \left( 1 + \frac{I}{\sinh I} \right) v \quad (2-22)$$

The first partial derivatives of U in the  $x'y'$ -coordinate system are given by

$$\frac{\partial U}{\partial x'} = \frac{\cosh I}{2} \left[ (\sinh I + I) \frac{\partial v}{\partial x'} + v \frac{\partial I}{\partial x'} (1 - I \coth I) \right] \quad (2-23)$$

$$\frac{\partial U}{\partial y'} = \frac{\cosh I}{2} \left[ (\sinh I + I) \frac{\partial v}{\partial y'} + v \frac{\partial I}{\partial y'} (1 - I \coth I) \right] = 0 \quad (2-24)$$

where

$$\frac{\partial I}{\partial x'} = I \left( \frac{1}{h} \frac{\partial h}{\partial x'} - \frac{1}{v} \frac{\partial v}{\partial x'} \right) \quad (2-25)$$

$$\frac{\partial I}{\partial y'} = I \left( \frac{1}{h} \frac{\partial h}{\partial y'} - \frac{1}{v} \frac{\partial v}{\partial y'} \right) = 0 \quad (2-26)$$

The second partial derivatives of U are given by

$$\begin{aligned} \frac{\partial^2 U}{(\partial x')^2} = & \frac{\cosh I}{2} \left[ (\sinh I + I) \frac{\partial^2 V}{(\partial x')^2} + V \frac{\partial^2 I}{(\partial x')^2} (1 - I \coth I) \right] + \\ & + \frac{\partial I}{\partial x'} \left\{ - \frac{\partial U}{\partial x'} \coth I + \frac{\cosh I}{2} \left[ \frac{\partial V}{\partial x'} (2 + \cosh I - I \coth I) + \right. \right. \\ & \left. \left. + V \frac{\partial I}{\partial x'} \cosh I (I \cosh I - \cosh I) \right] \right\} \end{aligned} \quad (2-27)$$

$$\frac{\partial^2 U}{(\partial y')^2} = \frac{\cosh I}{2} \left[ (\sinh I + I) \frac{\partial^2 V}{(\partial y')^2} + V \frac{\partial^2 I}{(\partial y')^2} (1 - I \coth I) \right] \quad (2-28)$$

$$\frac{\partial^2 U}{\partial x' \partial y'} = \frac{\cosh I}{2} \left[ (\sinh I + I) \frac{\partial^2 V}{\partial x' \partial y'} + V \frac{\partial^2 I}{\partial x' \partial y'} (1 - I \coth I) \right] \quad (2-29)$$

where

$$\begin{aligned} \frac{\partial^2 I}{(\partial x')^2} = & \frac{1}{I} \left( \frac{\partial I}{\partial x'} \right)^2 + I \left\{ \frac{1}{h} \left[ \frac{\partial^2 h}{(\partial x')^2} - \frac{1}{k} \left( \frac{\partial h}{\partial x'} \right)^2 \right] - \right. \\ & \left. - \frac{1}{V} \left[ \frac{\partial^2 V}{(\partial x')^2} - \frac{1}{V} \left( \frac{\partial V}{\partial x'} \right)^2 \right] \right\} \end{aligned} \quad (2-30)$$

$$\frac{\partial^2 I}{(\partial y')^2} = I \left( \frac{1}{h} \frac{\partial^2 h}{(\partial y')^2} - \frac{1}{V} \frac{\partial^2 V}{(\partial y')^2} \right) \quad (2-31)$$

$$\frac{\partial^2 I}{\partial x' \partial y'} = I \left( \frac{1}{h} \frac{\partial^2 h}{\partial x' \partial y'} - \frac{1}{V} \frac{\partial^2 V}{\partial x' \partial y'} \right) \quad (2-32)$$

### c. Derivatives of G

The geometric group speed is defined by

$$G = U \cos \phi \quad (2-33)$$

where

$$\phi = \theta' - \gamma'$$

(2-34)

in which  $\theta'$  is the direction of the wave packet,  $\gamma'$  is the direction of the wavelets, and both directions are measured with respect to the positive  $x'$ -axis. The first spatial derivatives of  $G$  are

$$\frac{\partial G}{\partial x'} = \frac{\partial U}{\partial x'} \cos \phi - U \sin \phi \frac{\partial \phi}{\partial x'} \quad (2-35)$$

$$\frac{\partial G}{\partial y'} = \frac{\partial U}{\partial y'} \cos \phi - U \sin \phi \frac{\partial \phi}{\partial y'} = 0 \quad (2-36)$$

The ray curvature for a wave packet is given by (Part I)

$$\frac{\partial \theta'}{\partial s_G} = \cos \theta' \frac{\partial \theta'}{\partial x'} + \sin \theta' \frac{\partial \theta'}{\partial y'} = \frac{1}{G} \left( \sin \theta' \frac{\partial G}{\partial x'} - \cos \theta' \frac{\partial G}{\partial y'} \right) \quad (2-37)$$

where  $ds_G$  is an element of arc length along the packet trajectory. Since the first derivatives in  $y'$  vanish it is found that

$$\frac{\partial \theta'}{\partial x'} = \frac{\tan \theta'}{G} \frac{\partial G}{\partial x'} \quad (2-38)$$

$$\frac{\partial \theta'}{\partial y'} = \frac{-\cot \theta'}{G} \frac{\partial G}{\partial y'} = 0 \quad (2-39)$$

In like manner, when considering the ray curvature of the wavelets it is found that

$$\frac{\partial \gamma'}{\partial x'} = \frac{\tan \gamma'}{v} \frac{\partial v}{\partial x'} \quad (2-40)$$

$$\frac{\partial \gamma'}{\partial y'} = -\frac{\cot \gamma'}{v} \frac{\partial v}{\partial y'} = 0 \quad (2-41)$$

Therefore

$$\frac{\partial \phi}{\partial x'} = \frac{\tan \theta'}{G} \frac{\partial G}{\partial x'} - \frac{\tan \gamma'}{v} \frac{\partial v}{\partial x'} \quad (2-42)$$

$$\frac{\partial \phi}{\partial y'} = -\frac{\cot \theta'}{G} \frac{\partial G}{\partial y'} + \frac{\cot \gamma'}{v} \frac{\partial v}{\partial y'} = 0 \quad (2-43)$$

When Equation (2-42) is substituted into (2-35) and the result is simplified it is found that

$$\frac{\partial G}{\partial x'} = \rho \left( \frac{\partial U}{\partial x'} \cos \phi + \sigma \frac{\partial V}{\partial x'} \right) \quad (2-44)$$

where

$$\rho = (1 + \tan \phi \tan \theta')^{-1} \quad (2-45)$$

$$\sigma = \frac{U}{V} \sin \phi \tan \gamma' \quad (2-46)$$

The second partial derivative of  $G$  with respect to  $x'$  is given by

$$\begin{aligned} \frac{\partial^2 G}{(\partial x')^2} &= \rho \left( \frac{\partial^2 U}{(\partial x')^2} \cos \phi + \sigma \frac{\partial^2 V}{(\partial x')^2} \right) + \left( \cos \phi \frac{\partial \rho}{\partial x'} - \right. \\ &\quad \left. - \rho \sin \phi \frac{\partial \phi}{\partial x'} \right) \frac{\partial U}{\partial x'} + \left( \rho \frac{\partial \sigma}{\partial x'} + \sigma \frac{\partial \rho}{\partial x'} \right) \frac{\partial V}{\partial x'} \end{aligned} \quad (2-47)$$

where

$$\frac{\partial \rho}{\partial x'} = -\rho^2 \left( \tan \theta' \sec^2 \phi \frac{\partial \phi}{\partial x'} + \tan \phi \sec^2 \theta' \frac{\partial \theta'}{\partial x'} \right) \quad (2-48)$$

$$\frac{\partial G}{\partial x'} = \sigma \left( \frac{1}{U} \frac{\partial U}{\partial x'} - \frac{1}{v} \frac{\partial v}{\partial x'} \right) + \frac{U}{v} \left( \cos \phi \tan \gamma' \frac{\partial \phi}{\partial x'} + \sin \phi \sec^2 \gamma' \frac{\partial \gamma'}{\partial x'} \right) \quad (2-49)$$

The remaining second order spatial derivatives of  $G$  are obtained by differentiating Equation (2-36). Since the first derivatives in  $y'$  vanish it is found that

$$\frac{\partial^2 G}{(\partial y')^2} = \frac{\partial^2 U}{(\partial y')^2} \cos \phi - U \sin \phi \frac{\partial^2 \phi}{(\partial y')^2} \quad (2-50)$$

$$\frac{\partial^2 G}{\partial x' \partial y'} = \frac{\partial^2 U}{\partial x' \partial y'} \cos \phi - U \sin \phi \frac{\partial^2 \phi}{\partial x' \partial y'} \quad (2-51)$$

The second partial derivative of  $\phi$  with respect to  $y'$  is found by differentiating Equation (2-43). The result is

$$\frac{\partial^2 \phi}{(\partial y')^2} = - \frac{\cot \theta'}{G} \frac{\partial^2 G}{(\partial y')^2} + \frac{\cot \gamma'}{v} \frac{\partial^2 v}{(\partial y')^2} \quad (2-52)$$

When Equation (2-52) is substituted into (2-50) and the terms rearranged it is found that

$$\frac{\partial^2 G}{(\partial y')^2} = \zeta \left( \frac{\partial^2 U}{(\partial y')^2} \cos \phi - \xi \frac{\partial^2 v}{(\partial y')^2} \right) \quad (2-53)$$

where

$$\zeta = (1 - \tan \phi \cot \theta')^{-1} \quad (2-54)$$

$$\xi = \frac{U}{v} \sin \phi \cot \gamma' \quad (2-55)$$

After differentiating Equation (2-43) with respect to  $x'$  it is found that

$$\frac{\partial^2 \phi}{\partial x' \partial y'} = \frac{\tan \theta'}{G} \frac{\partial^2 G}{\partial x' \partial y'} - \frac{\tan \gamma'}{v} \frac{\partial^2 v}{\partial x' \partial y'} \quad (2-56)$$

When Equation (2-56) is substituted into (2-51) the simplified result becomes

$$\frac{\partial^2 G}{\partial x' \partial y'} = \rho \left( \frac{\partial^2 u}{\partial x' \partial y'} \cos \phi + v \frac{\partial^2 v}{\partial x' \partial y'} \right) \quad (2-57)$$

2.3 Ray Curvature for Nonparallel Water Depth Contours. The ray curvature  $\kappa_G$  depends only on derivatives of the first order. By making the calculations in the  $x'y'$ -coordinate system the first derivatives in  $y'$  vanish. This results in a simplified ray curvature expression which is formally the same as presented in Part I for locally parallel water depth contours.

$$\kappa_G = \frac{\sin \theta'}{G} \frac{\partial G}{\partial x'} \quad (2-58)$$

2.4 Ray Separation Equation for Nonparallel Water Depth Contours. The ray separation equation is defined (Part I)

$$\frac{d^2 \beta}{dt^2} + p \frac{d\beta}{dt} + q \beta = 0 \quad (2-59)$$

where  $\beta$  is the ray separation factor,  $t$  is time, and

$$p = -2 \cos \theta' \frac{\partial G}{\partial x'} \quad (2-60)$$

$$q = G \left( \sin^2 \theta' \frac{\partial^2 G}{(\partial x')^2} - 2 \sin \theta' \cos \theta' \frac{\partial^2 G}{\partial x' \partial y'} + \cos^2 \theta' \frac{\partial^2 G}{(\partial y')^2} \right) \quad (2-61)$$

In the  $x'y'$ -coordinate system  $p$  is simplified as is  $\kappa_G$ . In  $q$  there is also a reduction in the expressions for the second order derivatives involving  $y'$ .

### CHAPTER III SUBROUTINE SURFCE

3.1 Modifications to Subroutine SURFCE. Modifications have been made to the program subroutine SURFCE in Part I in order to remove the restriction that the water depth contours be locally parallel about each ray point and to simplify some of the computations. When the program is run with these modifications minor changes will occur in the printed output for the sample input data presented in Part I.

Referring to the program listing for SURFCE on pages 52-54 of Part I the following modifications have been made. In place of SURFACE 58 an alternative expression is used to compute  $\partial h / \partial x'$  (DHDX). Between SURFACE 87 and SURFACE 88  $\partial I / \partial x'$  (DIDX) is computed as are the hyperbolic sine, cosine, and tangent of I. Statements SURFACE 91 through SURFACE 93 for computing  $\partial U / \partial x'$  (DUDX) have been replaced by a simplified expression. Between SURFACE 107 and SURFACE 108  $\partial^2 h / (\partial y')^2$  (DHDYY) and  $\partial^2 h / \partial x' \partial y'$  (DHDXY) are computed. SURFACE 111 has been replaced by expressions to compute  $\partial^2 v / (\partial y')^2$  (DVDYY) and  $\partial^2 v / \partial x' \partial y'$  (DVDXY). Statements SURFACE 114 through SURFACE 120 have been replaced by expressions for calculating  $\partial^2 I / (\partial y')^2$  (DIDYY) and  $\partial^2 I / \partial x' \partial y'$  (DIDXY). A simplified expression has been substituted for  $\partial^2 U / (\partial x')^2$  (DUDXX). Also, the values of  $\partial^2 U / (\partial y')^2$  (DUDYY) and  $\partial^2 U / \partial x' \partial y'$  (DUDXY) are computed. Between SURFACE 122 and SURFACE 123  $\zeta$  (ZETA),  $\xi$  (XI),  $\partial^2 G / (\partial y')^2$  (DGDYY), and  $\partial^2 G / \partial x' \partial y'$  (DGDXY) are determined. Simplified expressions are used to compute  $\partial^2 G / (\partial y')^2$  and  $\partial^2 G / \partial x' \partial y'$  when  $|\tan \theta'| \leq \tan 5^\circ$  and  $|\tan \gamma'| \leq \tan 5^\circ$ . Statement SURFACE 124 has been replaced by the complete expression for computing q (QOT).

3.2 Description of Subroutine SURFCE. SURFCE is called by RAYN and MOVE to calculate  $h$ ,  $\alpha$ ,  $\gamma$ ,  $G$ ,  $p$ ,  $q$ ,  $k_G$ , and other ray particulars. At the first ray point twelve values of  $h$  from CMAT are selected about the point as shown in Figure (2-2) of Part I. A quadratic surface is fit to the set of water depths. At successive ray points the quadratic surface is determined only if there is a change in the set of twelve water depths. The water depth and its partial derivatives in the fixed xy-system,  $\partial h / \partial x$ ,  $\partial h / \partial y$ ,  $\partial^2 h / \partial x^2$ ,  $\partial^2 h / \partial y^2$ , and  $\partial^2 h / \partial x \partial y$ , are determined at the ray point by interpolating on the quadratic surface.

If  $h \leq 0$ , NDP = 2 and there is a RETURN. If  $h > 0$  the ratio of the water depth to the deep water wavelength is computed. If  $h/\lambda_d > 0.64$ , which defines deep water, NFK = 1. If  $h/\lambda_d \leq 0.64$ , NFK = 2. VELCTY is called, and after the return if NFK = 1, W = 0. If NFK = 2, CONDER is called to

compute W. The values of  $\partial v / \partial x$  and  $\partial v / \partial y$  are calculated using W.

At each ray point calculations are made in a  $x'y'$ - coordinate system which is chosen with the positive  $x'$ -axis in the direction of the gradient of the water depth. The value of  $\partial h / \partial x'$  is computed, and if it exceeds 0.00001 the angle  $\alpha$  by which the  $x'$ -axis is rotated with respect to the  $x$ -axis is computed. If  $|\partial h / \partial x'| \leq 0.00001$  the water depth is assumed to be constant and  $\alpha$  remains constant.

If FLAG1 = 0,  $\gamma'$  is computed, and if necessary it is placed within the range  $|\gamma'| \leq 360^\circ$ . A check is made to determine if there is total reflection. If there is, FLAG2 = 1 and there is a RETURN. Otherwise, FLAG2 = 0 and the new  $\gamma'$  is computed using Snell's law with phase velocity following a set of rules. Using the values of  $\gamma'$ ,  $\gamma$  is computed. When FLAG 1  $\neq 0$  these steps for computing the new wavelet direction and the test for total reflection are omitted.

The values of  $\phi$ , G, and  $\partial v / \partial x'$  are calculated. If NFK = 2,  $\partial U / \partial x'$  is determined using its unsimplified expression. If NFK  $\neq 2$ , the deep water formula is used to calculate  $\partial U / \partial x'$ . The value of  $\partial U / \partial x'$  is used in computing  $\partial G / \partial x'$ .

If NFK  $\neq 2$ , the coefficients of the ray separation equation and the ray curvature are set equal to zero. Then there is a RETURN. If NFK = 2, p,  $\partial^2 h / (\partial x')^2$ ,  $\partial^2 h / (\partial y')^2$ ,  $\partial^2 h / \partial x' \partial y'$ ,  $\partial^2 v / (\partial x')^2$ ,  $\partial^2 v / (\partial y')^2$ ,  $\partial^2 v / \partial x' \partial y'$ ,  $\partial^2 U / (\partial x')^2$ ,  $\partial^2 U / (\partial y')^2$ ,  $\partial^2 U / \partial x' \partial y'$ ,  $\partial^2 G / (\partial x')^2$ ,  $\partial^2 G / (\partial y')^2$ ,  $\partial^2 G / \partial x' \partial y'$ , q, and  $\kappa_G$  are computed. This is followed by a return.

3.3 Listing of Subroutine SURFCE

```

SUBROUTINE SURFCF(X,Y,A,FK,NFK,NDP,AV)          SURFACE 1
DIMENSION S(6,F),EM(6,12),C(12),YVH(6),E(6)    SURFACE 2
$,CMAT(100,100),AX(2000),AY(2000),CONTUR(9)    SURFACE 3
RFL,KR,KF,KS,KRTOL,KFC                         SURFACE 4
INTEGER FLAG1,FLAG2                           SURFACE 5
COMMON S,EM,E,YVH,CMAT,C,AX,AY,CONTUR,PROJCT,GRID,DCON,FAN,DATE1 SURFACE 6
$,DATE2,CIN,DIR,POP,TT,WBCOP,MOE,DY,DELTAT,SOLTAT,O,HGT,HGTZ,SVX SURFACE 7
$,SVY,SOEP,W,DEF,NL,V,SAVV,FREV,SPPEV,U,SAVU,GZERO,G,SG,SVG,DUD,KS SURFACE 8
$,DGDX,SVA,TPI,SAV,SAV,V,PHI,ALFA,SVALFA,SSALFA,CNVRSA,DELA,UHDX SURFACE 9
$,SVFKB,SAVFK,FKEAR,MAXQ,SK,FLC,NUMT,INU1,IFLC,RCOUNT,AMM,ANN SURFACE 10
$,PFLCT,RFLBUM,PFLPCT,FFRUM,BRK,FLAG1,FLAG2,FLAG3,FLAGR,KFC,CF,BZ SURFACE 11
$,SBZ,RNZ,SBNZ,KRTOL,KR,POT,P1,P2,P3,P4,P5,QOT,Q1,Q2,Q3,Q4,Q5 SURFACE 12
C IN THIS SUBROUTINE THE WATER DEPTH, ROTATION ANGLE, WAVELET SURFACE 13
C DIRECTION, GEOMETRIC GROUP VELOCITY, COEFFICIENTS OF THE RAY SURFACE 14
C SEPARATION EQUATION, AND THE PACKET RAY CURVATURE ARE COMPUTED. SURFACE 15
C SUBROUTINES VELCTY AND CONDER ARE CALLED. SURFACE 16
C
I=X $ J=Y $ FI=I $ FJ=J $ XL=X+1.-FI $ YL=Y+1.-FJ SURFACE 17
IF (MAXQ .LE. 1) GO TO 1 SURFACE 18
IF (ZI .NE. FI) GU TO 1 SURFACE 19
IF (ZJ .EQ. FJ) GO TO 3 SURFACE 20
1 ZI=FI & ZJ=FJ SURFACE 21
C SELECT 12 WATER DEPTHS ABOUT RAY POINT SURFACE 22
C(1)=CMAT(J+1,I) $ C(2)=CMAT(J+2,I) $ C(3)=CMAT(J,I+1) SURFACE 23
C(4)=CMAT(J+1,I+1) $ C(5)=CMAT(J+2,I+1) $ C(6)=CMAT(J+3,I+1) SURFACE 24
C(7)=CMAT(J,I+2) $ C(8)=CMAT(J+1,I+2) $ C(9)=CMAT(J+2,I+2) SURFACE 25
C(10)=CMAT(J+3,I+2) $ C(11)=CMAT(J+1,I+3) $ C(12)=CMAT(J+2,I+3) SURFACE 26
C FIT QUADRATIC SURFACE TO 12 WATER DEPTHS SURFACE 27
DO 318 II=1,6 SURFACE 28
YVW(II)=0. SURFACE 29
DO 319 L=1,12 SURFACE 30
YVW(II)=YVW(II)+C(L)*EM(II,L) SURFACE 31
318 CONTINUE SURFACE 32
DO 319 II=1,6 SURFACE 33
E(II)=0. SURFACE 34
DO 319 JJ=1,6 SURFACE 35
E(II)=E(II)+S(JJ,II)*YVW(JJ) SURFACE 36
319 CONTINUE SURFACE 37
C COMPUTE INTERPOLATED WATER DEPTH SURFACE 38
3 DEP=(E(1)+E(2)*XL+E(3)*YL+E(4)*XL**2+E(5)*XL*YL+E(6)*YL**2)*DCON SURFACE 39
C COMPUTE PARTIAL DERIVATIVES OF WATER DEPTH IN FIXED XY-SYSTEM SURFACE 40
HX=(E(2)+2.*E(4)*XL+E(5)*YL)*DCON SURFACE 41
HY=(E(3)+E(5)*XL+E(6)*YL)*DCON SURFACE 42
HXX=2.*E(4)*DCON $ HYX=2.*E(6)*DCON $ HXY=E(5)*DCON SURFACE 43
IF (DEP .GT. 0.) GO TO 324 SURFACE 44
NDF=2 SURFACE 45
GO TO 403 SURFACE 46
324 IF (DEP/WL .GT. .64) GO TO 322 SURFACE 47
NFK=2 SURFACE 48
GO TO 323 SURFACE 49
322 NFK=1 SURFACE 50
323 CALL VELCTY(V,TT,MAXQ,DEP,NFK,U) SURFACE 51
IF (NFK .EQ. 2) GO TO 402 SURFACE 52
W=0. SURFACE 53
GO TO 10 SURFACE 54
402 CN=1. SURFACE 55
CALL CONDER(DN,TT,V,MAXQ,NFK) SURFACE 56
W=DN SURFACE 57

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10  UX=W*HX $ VY=W*HY $ DHDX=HX*COS(ALFA)+HY*SIN(ALFA) SURFACE 58
  IF (ABS(DHDX/GRID) .GT. 0.00001) GO TO 8 SURFACE 59
  GO TO 9 SURFACE 60
C COMPUTE ROTATION ANGLE SURFACE 61
  ALFA=ATAN2(HY,HX) SURFACE 62
  IF (FLAG1 .NE. 0) GO TO 12 SURFACE 63
C COMPUTE WAVELET DIRECTION IN ROTATED XY-SYSTEM USING SNELLS LAW SURFACE 64
C WITH V AND TEST FOR TOTAL REFLECTION DUE TO THE WAVELETS SURFACE 65
  GP=SAV-ALFA SURFACE 66
  14  IF (ABS(GP) .LE. 6.2831853) GO TO 13 SURFACE 67
  IF (GP) 16,13,17 SURFACE 68
  16  GP=GP+6.2831853 SURFACE 69
  GO TO 14 SURFACE 70
  17  GP=GP-6.2831853 SURFACE 71
  GO TO 14 SURFACE 72
  13  ARG1=V*SIN(GP)/SAVV SURFACE 73
  IF (ABS(ARG1) .LE. 1.) GO TO 18 SURFACE 74
  FLAG2=1 SURFACE 75
  GO TO 403 SURFACE 76
  18  FLAG2=0. $ GPT=ASIN(ARG1) SURFACE 77
  IF (ABS(GP) .LE. 4.7123889) GO TO 20 SURFACE 78
  AVP=6.2831853+GPT SURFACE 79
  GO TO 22 SURFACE 80
  20  IF (ABS(GP) .LE. 1.5707963) GO TO 23 SURFACE 81
  AVP=3.1415927-GPT SURFACE 82
  GO TO 22 SURFACE 83
  AVP=GPT SURFACE 84
  22  AV=AVP+ALFA SURFACE 85
  12  PHI=A-AV $ G=U+COS(PHI) SURFACE 86
  DVDX=X=W*DHDX $ PAR3=12.5663708/T7 $ BAR4=PAR3*DEP/V SURFACE 87
  DIDX=BAR4*(DHDX/DEP-[DVDX/V]) $ SINHI=(EXP(BAR4)-EXP(-BAR4))/2. SURFACE 88
  COSHI=(EXP(BAR4)+EXP(-BAR4))/2. $ TANHI=SINHI/COSHI SURFACE 89
  IF (INFK .EQ. 2) GO TO 25 SURFACE 90
  DVDX=DVDX/2. SURFACE 91
  GO TO 27 SURFACE 92
  25  DUDX=(.5/SINHI)*((SINHI+BAR4)*DVDX+V*DIDX*(1.-BAR4/TANHI)) SURFACE 93
  RHO=1./((1.+TAN(PHI))+TAN(A-ALFA)) $ SIGMA=U*SIN(PHI)*TAN(AV-ALFA)/V SURFACE 94
  DVDX=RHO*(CUDX*COS(PHI)+SIGMA*DVOX) SURFACE 95
  IF (INFK .EQ. 2) GO TO 28 SURFACE 96
  POT=0. $ QOT=1. $ FK=0. SURFACE 97
  GO TO 403 SURFACE 98
C COMPUTE P IN ROTATED XY-SYSTEM SURFACE 99
  28  POT=-2.*COS(A-ALFA)*DGDX/GRID $ DAVDX=TAN(AV-ALFA)*DVOX/V SURFACE 100
  DVDX=TAN(A-ALFA)*DGDX/G $ DPDX=DAVDX-DAVDX SURFACE 101
  DRHDX=-(RHO**2)*(TAN(A-ALFA))*DPDX/(COS(PHI)**2)+ SURFACE 102
  $TAN(PHI)*DAVDX/(COS(AV-ALFA)**2) SURFACE 103
  $SIGDX=SIGMA*(DUDX/U-DVDX/V)+U*(COS(PHI)*TAN(AV-ALFA)*DPDX+ SURFACE 104
  $SIN(PHI)*DAVDX/(COS(AV-ALFA)**2))/V SURFACE 105
  DHDX=ECOS(ALFA)**2)*HXX+2.*SIN(ALFA)*COS(ALFA)*HXY+(SIN(ALFA)**2)*HYY SURFACE 106
  DHDX=-(SIN(ALFA)**2)*HXX-2.*SIN(ALFA)*COS(ALFA)*HXY SURFACE 107
  $+(COS(ALFA)**2)*HYY SURFACE 108
  DHDX=(SIN(ALFA)*COS(ALFA))*(HYY**2)-(HXX**2)+ SURFACE 109
  $((COS(ALFA)**2)-(SIN(ALFA)**2))*HXY SURFACE 110
  SMA=6.2831854/(32.2*T) $ SMAB=1./64.4 SURFACE 111
  DVDX=V*(DHDX+(DHDX**2)*(-4.*SMAB*(V**2)/(12.*SMAB*(V**2)+DEP* SURFACE 112
  $(1.-(SMAB*V)**2)**2))) SURFACE 113
  DVDX=V*(DHDX+(DHDX**2)*(-4.*SMAB*(V**2)/(12.*SMAB*(V**2)+DEP* SURFACE 114
  DVDX=V*(DHDX+(DHDX**2)*(-4.*SMAB*(V**2)/(12.*SMAB*(V**2)+DEP* SURFACE 115

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	DIDXX=(DIOUX**2)/BAR4+BAR4*((DHDXX-(DHDX**2)/DEP)/DEP-(DVDXX \$-(UVUX**2)/V)/V)	SURFACE	116
	DIDYY=BAR4*(DHDXY/DEP-DVDXY/V) \$ DIDXY=BAR4*(DHDXY/DEP-DVDXY/V)	SURFACE	117
30	DUDXX=(.5/SINH1)*(-1*(SINH1+BAR4)+DVDX*V*DIOX*(1.-BAR4/TANH1))*DIOX \$/TANH1*(SINH1+BAR4)*DVDXX+DVDX*OIDX*(2.*COSHI-BAR4/TANH1)+V*(CIDXX \$*(1.-BAR4/TANH1)+(CIDX**2)*(BAR4/SINH1-COSHI)/SINH1)	SURFACE	118
	DUDYY=(.5/SINH1)*((SINH1+BAR4)*DVDYY+V*DIDYY*(1.-BAR4/TANH1))	SURFACE	119
	DUDXY=(.5/SINH1)*((SINH1+BAR4)*DVDXY+V*DIDXY*(1.-BAR4/TANH1))	SURFACE	120
32	DGDXX=RHO*(COS(PHI)*CUDXX+SIGMA*DVDXX)+(CUS(PHI)*DRHO*DX-RHO*SIN \$(PHI)*OPHI*DX)*CUDX+(RHO*DSIG*DX+SIGMA*DRHO*DX)*DVDX	SURFACE	121
	IF (ABS(TAN(A-ALFA)) .GT. U.0.0748866 .AND. ABS(TAN(AV-ALFA)) .GT. \$U.0.0748866) GO TO 35	SURFACE	122
	ZETA=1. \$ XI=0.	SURFACE	123
	GO TO 36	SURFACE	124
35	ZETA=1.0/(1.0-TAN(PHI)/TAN(A-ALFA)), \$ XI=U*SIN(PHI)/V*TAN(AV-ALFA)	SURFACE	125
36	DGDYY=ZETA*(DUCYY*COS(PHI)-XI*DVCYY), \$ CGDXY=RHO*(DUDXY*LCS(PHI)+SIGMA*DVDXY)	SURFACE	126
C	COMPUTE Q IN ROTATED XY-SYSTEM	SURFACE	127
	QOT=(G*((SIN(A-ALFA)**2)*DGDXX-2.*SIN(A-ALFA)*COS(A-ALFA)*DGDYX+ \$ (COS(A-ALFA)**2)*DGDYY))/(GRIC**2)	SURFACE	128
C	COMPUTE PACKET RAY CURVATURE IN ROTATED XY-SYSTEM	SURFACE	129
	FK=SIN(A-ALFA)*UGDX/G	SURFACE	130
403	RETURN	SURFACE	131
	END	SURFACE	132
		SURFACE	133
		SURFACE	134
		SURFACE	135
		SURFACE	136
		SURFACE	137
		SURFACE	138
		SURFACE	139

## CHAPTER IV WAVE PREDICTION EXAMPLES FOR NONPARALLEL WATER DEPTH CONTOURS

Two examples of nonparallel water depth contours are used to illustrate the wave prediction program.

4.1 Sinusoidal to Parallel Water Depth Contours. Figure (4-1) shows a set of wave packet trajectories for a wave period of 7.0 seconds. The rays begin in deep water (initial water depth equal to or greater than one half the wavelength). The water depth contours are sinusoidal at the shoreline and gradually become parallel at a water depth of 38.6 meters. The amplitude of the water depth contour at the shoreline is 2.5 kilometers, and the contour wavelength is 10 kilometers. GRID, the distance between grid points, is 156.25 meters.

From the figure it is seen that the energy is fairly evenly distributed along the coastline with slightly more energy at the headland than in the bay. Figures (4-2) through (4-4) show the printed output for rays number 4, 8, and 16, respectively, of Figure (4-1). The computed refraction coefficients are in good agreement with values estimated from the plot.

4.2 Sinusoidal Water Depth Contours. In Figure (4-5) the period of the wave packet trajectories is 14.0 seconds and they begin in deep water. The water depth contours are sinusoidal with an amplitude of 5 kilometers and a wavelength of 20 kilometers. GRID has a value of 312.5 meters.

This example is quite interesting since there is decidedly more energy in the bay than at the headland. The opposite result would be expected for monochromatic trajectories. The refraction of wave packets could explain why there is more erosion in bays than at headlands for some coastlines. Figures (4-6) through (4-8) show the printed output for rays number 4, 8, and 16, respectively, of Figure (4-5). The computed refraction coefficients agree favorably with values estimated from the plot.

PROJ. NO. SIP 2, 78/11/13  
 PLOT = 1/98425 CIN = 0  
 SPLT NO. 1, DIR. = WAVPAK

SOUNDINGS IN METERS

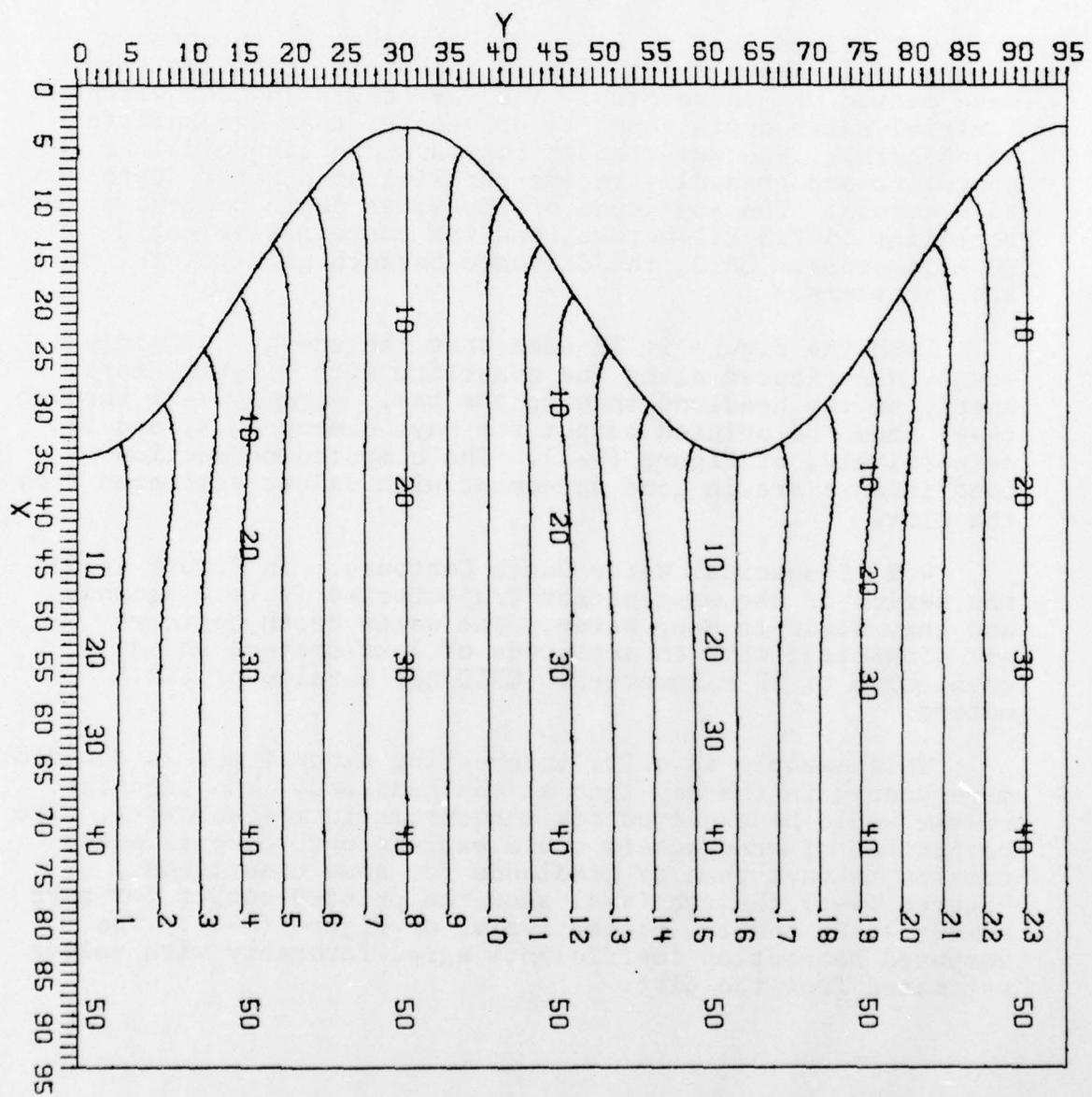


Figure (4-1). Plot for sinusoidal to parallel water depth contours.

Figure (4-2). Printed output for Ray No. 4 in Figure (4-1).

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PROJECT NO. SIP 2, 78/11/13 , PLOT NO. 1, PERIOD= 7.0SEC., RAY NO. 16, ELLTAT= 10.00, CF=0.000000, KRTOL= .001000

THE OUTPUT IS IN METRIC UNITS. H=HGT(METER). G,U,V(METER/SECOND).

MAX X	Y	Z	H	PACK	WAVE D	FK	ALFA	G	U	V	HGT	KS	KF	NC KR	PCTOFF
80.00	64.00	45.10	180.00	180.00	3.53E-01	5.70E-07	10.92	1.0000	1.0000	1.0000	1.0000	1.0000	1.0000	1.0000	.01
78.59	64.00	44.25	180.00	180.00	3.53E-01	6.66E-01	10.92	0.9995	0.9995	0.9995	0.9995	0.9995	0.9995	0.9995	.01
76.82	64.00	43.28	180.00	180.00	3.53E-01	7.60E-01	10.92	0.9989	0.9989	0.9989	0.9989	0.9989	0.9989	0.9989	.01
75.06	64.00	42.20	180.00	180.00	3.53E-01	8.69E-01	10.92	0.9982	0.9982	0.9982	0.9982	0.9982	0.9982	0.9982	.01
73.24	64.00	41.14	180.00	180.00	3.53E-01	9.79E-01	10.92	0.9975	0.9975	0.9975	0.9975	0.9975	0.9975	0.9975	.01
71.51	64.00	40.00	180.00	180.00	3.53E-01	1.08E-01	10.92	0.9963	0.9963	0.9963	0.9963	0.9963	0.9963	0.9963	.01
69.73	64.00	39.00	180.00	180.00	3.53E-01	2.07E-01	10.92	0.9953	0.9953	0.9953	0.9953	0.9953	0.9953	0.9953	.01
68.00	64.00	37.93	180.00	180.00	3.53E-01	3.06E-01	10.92	0.9941	0.9941	0.9941	0.9941	0.9941	0.9941	0.9941	.01
66.27	64.00	35.97	180.00	180.00	3.53E-01	4.05E-01	10.92	0.9931	0.9931	0.9931	0.9931	0.9931	0.9931	0.9931	.01
64.54	64.00	33.92	180.00	180.00	3.53E-01	5.04E-01	10.92	0.9921	0.9921	0.9921	0.9921	0.9921	0.9921	0.9921	.01
62.82	64.00	31.87	180.00	180.00	3.53E-01	6.03E-01	10.92	0.9911	0.9911	0.9911	0.9911	0.9911	0.9911	0.9911	.01
61.10	64.00	30.00	180.00	180.00	3.53E-01	7.02E-01	10.92	0.9901	0.9901	0.9901	0.9901	0.9901	0.9901	0.9901	.01
59.38	64.00	28.06	180.00	180.00	3.53E-01	8.01E-01	10.92	0.9891	0.9891	0.9891	0.9891	0.9891	0.9891	0.9891	.01
57.65	64.00	26.14	180.00	180.00	3.53E-01	9.00E-01	10.92	0.9881	0.9881	0.9881	0.9881	0.9881	0.9881	0.9881	.01
55.93	64.00	24.23	180.00	180.00	3.53E-01	1.00E-01	10.92	0.9871	0.9871	0.9871	0.9871	0.9871	0.9871	0.9871	.01
54.20	64.00	22.32	180.00	180.00	3.53E-01	2.00E-01	10.92	0.9861	0.9861	0.9861	0.9861	0.9861	0.9861	0.9861	.01
52.47	64.00	20.40	180.00	180.00	3.53E-01	3.00E-01	10.92	0.9851	0.9851	0.9851	0.9851	0.9851	0.9851	0.9851	.01
50.74	64.00	18.50	180.00	180.00	3.53E-01	4.00E-01	10.92	0.9841	0.9841	0.9841	0.9841	0.9841	0.9841	0.9841	.01
49.01	64.00	16.59	180.00	180.00	3.53E-01	5.00E-01	10.92	0.9831	0.9831	0.9831	0.9831	0.9831	0.9831	0.9831	.01
47.28	64.00	14.68	180.00	180.00	3.53E-01	6.00E-01	10.92	0.9821	0.9821	0.9821	0.9821	0.9821	0.9821	0.9821	.01
45.55	64.00	12.77	180.00	180.00	3.53E-01	7.00E-01	10.92	0.9811	0.9811	0.9811	0.9811	0.9811	0.9811	0.9811	.01
43.82	64.00	10.86	180.00	180.00	3.53E-01	8.00E-01	10.92	0.9801	0.9801	0.9801	0.9801	0.9801	0.9801	0.9801	.01
42.09	64.00	8.95	180.00	180.00	3.53E-01	9.00E-01	10.92	0.9791	0.9791	0.9791	0.9791	0.9791	0.9791	0.9791	.01
40.36	64.00	7.04	180.00	180.00	3.53E-01	1.00E-01	10.92	0.9781	0.9781	0.9781	0.9781	0.9781	0.9781	0.9781	.01
38.63	64.00	5.13	180.00	180.00	3.53E-01	2.00E-01	10.92	0.9771	0.9771	0.9771	0.9771	0.9771	0.9771	0.9771	.01
36.80	64.00	3.22	180.00	180.00	3.53E-01	3.00E-01	10.92	0.9761	0.9761	0.9761	0.9761	0.9761	0.9761	0.9761	.01
35.07	64.00	1.31	180.00	180.00	3.53E-01	4.00E-01	10.92	0.9751	0.9751	0.9751	0.9751	0.9751	0.9751	0.9751	.01
33.34	64.00	-0.90	180.00	180.00	3.53E-01	5.00E-01	10.92	0.9741	0.9741	0.9741	0.9741	0.9741	0.9741	0.9741	.01
31.61	64.00	-3.80	180.00	180.00	3.53E-01	6.00E-01	10.92	0.9731	0.9731	0.9731	0.9731	0.9731	0.9731	0.9731	.01
29.88	64.00	-5.70	180.00	180.00	3.53E-01	7.00E-01	10.92	0.9721	0.9721	0.9721	0.9721	0.9721	0.9721	0.9721	.01
28.15	64.00	-7.59	180.00	180.00	3.53E-01	8.00E-01	10.92	0.9711	0.9711	0.9711	0.9711	0.9711	0.9711	0.9711	.01
26.42	64.00	-9.48	180.00	180.00	3.53E-01	9.00E-01	10.92	0.9701	0.9701	0.9701	0.9701	0.9701	0.9701	0.9701	.01
24.69	64.00	-11.37	180.00	180.00	3.53E-01	1.00E-01	10.92	0.9691	0.9691	0.9691	0.9691	0.9691	0.9691	0.9691	.01
22.96	64.00	-13.26	180.00	180.00	3.53E-01	2.00E-01	10.92	0.9681	0.9681	0.9681	0.9681	0.9681	0.9681	0.9681	.01
21.23	64.00	-15.15	180.00	180.00	3.53E-01	3.00E-01	10.92	0.9671	0.9671	0.9671	0.9671	0.9671	0.9671	0.9671	.01
19.50	64.00	-17.04	180.00	180.00	3.53E-01	4.00E-01	10.92	0.9661	0.9661	0.9661	0.9661	0.9661	0.9661	0.9661	.01
17.77	64.00	-18.93	180.00	180.00	3.53E-01	5.00E-01	10.92	0.9651	0.9651	0.9651	0.9651	0.9651	0.9651	0.9651	.01
16.04	64.00	-20.82	180.00	180.00	3.53E-01	6.00E-01	10.92	0.9641	0.9641	0.9641	0.9641	0.9641	0.9641	0.9641	.01
14.31	64.00	-22.71	180.00	180.00	3.53E-01	7.00E-01	10.92	0.9631	0.9631	0.9631	0.9631	0.9631	0.9631	0.9631	.01
12.58	64.00	-24.60	180.00	180.00	3.53E-01	8.00E-01	10.92	0.9621	0.9621	0.9621	0.9621	0.9621	0.9621	0.9621	.01
10.85	64.00	-26.49	180.00	180.00	3.53E-01	9.00E-01	10.92	0.9611	0.9611	0.9611	0.9611	0.9611	0.9611	0.9611	.01
9.12	64.00	-28.38	180.00	180.00	3.53E-01	1.00E-01	10.92	0.9601	0.9601	0.9601	0.9601	0.9601	0.9601	0.9601	.01
7.39	64.00	-30.27	180.00	180.00	3.53E-01	2.00E-01	10.92	0.9591	0.9591	0.9591	0.9591	0.9591	0.9591	0.9591	.01
5.66	64.00	-32.16	180.00	180.00	3.53E-01	3.00E-01	10.92	0.9581	0.9581	0.9581	0.9581	0.9581	0.9581	0.9581	.01
3.93	64.00	-34.05	180.00	180.00	3.53E-01	4.00E-01	10.92	0.9571	0.9571	0.9571	0.9571	0.9571	0.9571	0.9571	.01
2.20	64.00	-35.94	180.00	180.00	3.53E-01	5.00E-01	10.92	0.9561	0.9561	0.9561	0.9561	0.9561	0.9561	0.9561	.01
0.47	64.00	-37.83	180.00	180.00	3.53E-01	6.00E-01	10.92	0.9551	0.9551	0.9551	0.9551	0.9551	0.9551	0.9551	.01
-1.50	64.00	-39.72	180.00	180.00	3.53E-01	7.00E-01	10.92	0.9541	0.9541	0.9541	0.9541	0.9541	0.9541	0.9541	.01
-3.23	64.00	-41.61	180.00	180.00	3.53E-01	8.00E-01	10.92	0.9531	0.9531	0.9531	0.9531	0.9531	0.9531	0.9531	.01
-4.99	64.00	-43.50	180.00	180.00	3.53E-01	9.00E-01	10.92	0.9521	0.9521	0.9521	0.9521	0.9521	0.9521	0.9521	.01
-6.76	64.00	-45.39	180.00	180.00	3.53E-01	1.00E-01	10.92	0.9511	0.9511	0.9511	0.9511	0.9511	0.9511	0.9511	.01
-8.53	64.00	-47.28	180.00	180.00	3.53E-01	2.00E-01	10.92	0.9501	0.9501	0.9501	0.9501	0.9501	0.9501	0.9501	.01
-10.30	64.00	-49.17	180.00	180.00	3.53E-01	3.00E-01	10.92	0.9491	0.9491	0.9491	0.9491	0.9491	0.9491	0.9491	.01
-12.07	64.00	-50.96	180.00	180.00	3.53E-01	4.00E-01	10.92	0.9481	0.9481	0.9481	0.9481	0.9481	0.9481	0.9481	.01
-13.84	64.00	-52.85	180.00	180.00	3.53E-01	5.00E-01	10.92	0.9471	0.9471	0.9471	0.9471	0.9471	0.9471	0.9471	.01
-15.61	64.00	-54.74	180.00	180.00	3.53E-01	6.00E-01	10.92	0.9461	0.9461	0.9461	0.9461	0.9461	0.9461	0.9461	.01
-17.38	64.00	-56.63	180.00	180.00	3.53E-01	7.00E-01	10.92	0.9451	0.9451	0.9451	0.9451	0.9451	0.9451	0.9451	.01
-19.15	64.00	-58.52	180.00	180.00	3.53E-01	8.00E-01	10.92	0.9441	0.9441	0.9441	0.9441	0.9441	0.9441	0.9441	.01
-20.92	64.00	-60.41	180.00	180.00	3.53E-01	9.00E-01	10.92	0.9431	0.9431	0.9431	0.9431	0.9431	0.9431	0.9431	.01
-22.69	64.00	-62.30	180.00	180.00	3.53E-01	1.00E-01	10.92	0.9421	0.9421	0.9421	0.9421	0.9421	0.9421	0.9421	.01
-24.56	64.00	-64.19	180.00	180.00	3.53E-01	2.00E-01	10.92	0.9411	0.9411	0.9411	0.9411	0.9411	0.9411	0.9411	.01
-26.33	64.00	-66.08	180.00	180.00	3.53E-01	3.00E-01									

PRJ. NO. SIN 3, 78/08/31  
PLOT NO. 19685 CIN = 30  
PLT NO. 1, DIR. = WAVPAK

SOUNDINGS IN METERS

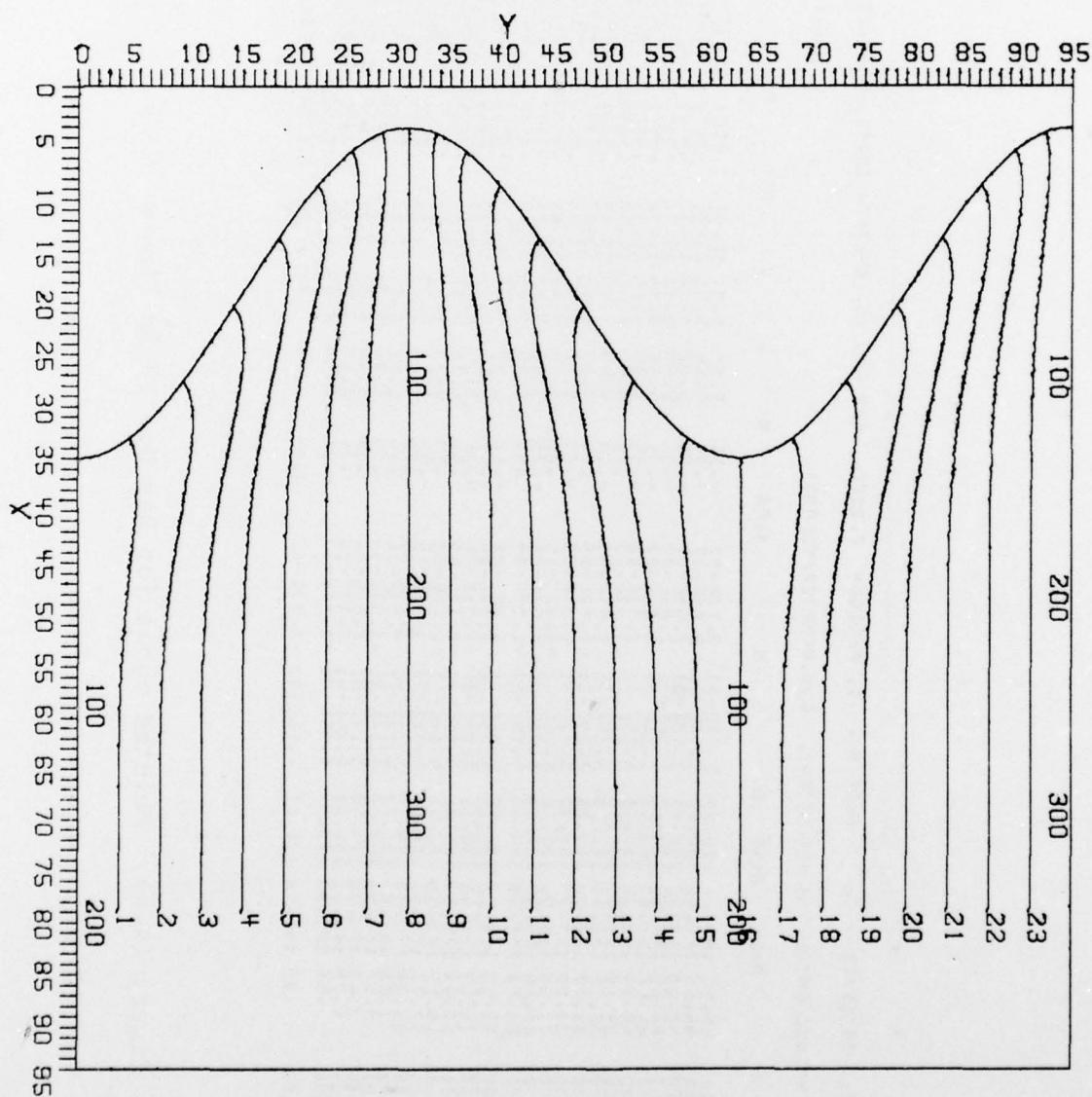


Figure (4-5). Plot for sinusoidal water depth contours.

Figure (4-6). Printed Output for Ray No. 4 in Figure (4-5).

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PROJECT NO. SIN 3, 78/06/31 PLOT NO. 1, PERIOD = 14.0 SEC., RAY NO. 0. ELEVAT= 10.00. KTR10L= 001000.

THE OUTLINE IS AS FOLLOWS:—

Figure (4-7). Printed output for Ray No. 8 in Figure (4-5).

PROJECT NO. SIN 3. 78/08/31 \* PLOT NO. 1, PERIOD= 14.0 SEC., RAY NO. 16, ELEVAT= 10.00, CF=0.000000, KRTOL= .001000  
 THE OUTPUT IS IN METRIC UNITS. H, HGT(METER). G, U,V(METER/SECOND).  
 MAX X Y H PACK HAVE D FK ALFA G U V HGT KS KF NO KR PCTDIF

X	Y	H	PACK	HAVE	D	FK	ALFA	G	U	V	HGT	KS	KF	NO KR	PCTDIF
0.00	0.00	180.00	3.50E-01								10.93	1.0000	1.0000		0.2
64.00	0.00	180.00	3.52E-01								10.93	1.0000	1.0000		0.2
128.00	0.00	180.00	3.54E-01								10.93	1.0000	1.0000		0.2
192.00	0.00	180.00	3.56E-01								10.93	1.0000	1.0000		0.2
256.00	0.00	180.00	3.58E-01								10.93	1.0000	1.0000		0.2
320.00	0.00	180.00	3.60E-01								10.93	1.0000	1.0000		0.2
384.00	0.00	180.00	3.62E-01								10.93	1.0000	1.0000		0.2
448.00	0.00	180.00	3.64E-01								10.93	1.0000	1.0000		0.2
512.00	0.00	180.00	3.66E-01								10.93	1.0000	1.0000		0.2
576.00	0.00	180.00	3.68E-01								10.93	1.0000	1.0000		0.2
640.00	0.00	180.00	3.70E-01								10.93	1.0000	1.0000		0.2
704.00	0.00	180.00	3.72E-01								10.93	1.0000	1.0000		0.2
768.00	0.00	180.00	3.74E-01								10.93	1.0000	1.0000		0.2
832.00	0.00	180.00	3.76E-01								10.93	1.0000	1.0000		0.2
896.00	0.00	180.00	3.78E-01								10.93	1.0000	1.0000		0.2
960.00	0.00	180.00	3.80E-01								10.93	1.0000	1.0000		0.2
1024.00	0.00	180.00	3.82E-01								10.93	1.0000	1.0000		0.2
1088.00	0.00	180.00	3.84E-01								10.93	1.0000	1.0000		0.2
1152.00	0.00	180.00	3.86E-01								10.93	1.0000	1.0000		0.2
1216.00	0.00	180.00	3.88E-01								10.93	1.0000	1.0000		0.2
1280.00	0.00	180.00	3.90E-01								10.93	1.0000	1.0000		0.2
1344.00	0.00	180.00	3.92E-01								10.93	1.0000	1.0000		0.2
1408.00	0.00	180.00	3.94E-01								10.93	1.0000	1.0000		0.2
1472.00	0.00	180.00	3.96E-01								10.93	1.0000	1.0000		0.2
1536.00	0.00	180.00	3.98E-01								10.93	1.0000	1.0000		0.2
1600.00	0.00	180.00	4.00E-01								10.93	1.0000	1.0000		0.2
1664.00	0.00	180.00	4.02E-01								10.93	1.0000	1.0000		0.2
1728.00	0.00	180.00	4.04E-01								10.93	1.0000	1.0000		0.2
1792.00	0.00	180.00	4.06E-01								10.93	1.0000	1.0000		0.2
1856.00	0.00	180.00	4.08E-01								10.93	1.0000	1.0000		0.2
1920.00	0.00	180.00	4.10E-01								10.93	1.0000	1.0000		0.2
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2496.00	0.00	180.00	4.28E-01								10.93	1.0000	1.0000		0.2
2560.00	0.00	180.00	4.30E-01								10.93	1.0000	1.0000		0.2
2624.00	0.00	180.00	4.32E-01								10.93	1.0000	1.0000		0.2
2688.00	0.00	180.00	4.34E-01								10.93	1.0000	1.0000		0.2
2752.00	0.00	180.00	4.36E-01								10.93	1.0000	1.0000		0.2
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3072.00	0.00	180.00	4.46E-01								10.93	1.0000	1.0000		0.2
3136.00	0.00	180.00	4.48E-01								10.93	1.0000	1.0000		0.2
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In Part I a method for calculating wave packet trajectories and wave heights is based on the assumption that the water depth contours are locally parallel in the vicinity of each ray point. This method is extended in order to predict the modification to surface gravity water waves in shoaling water when the water depth contours are not parallel. The calculations are greatly simplified by choosing a coordinate system at each ray point in which one axis is aligned parallel with the direction of the gradient of the water depth. Example printouts and plots are presented to illustrate the wave prediction method. It is discovered that when waves initially approach sinuous water depth contours symmetrically with respect to the beach there can be more energy in the bays than at the headlands.