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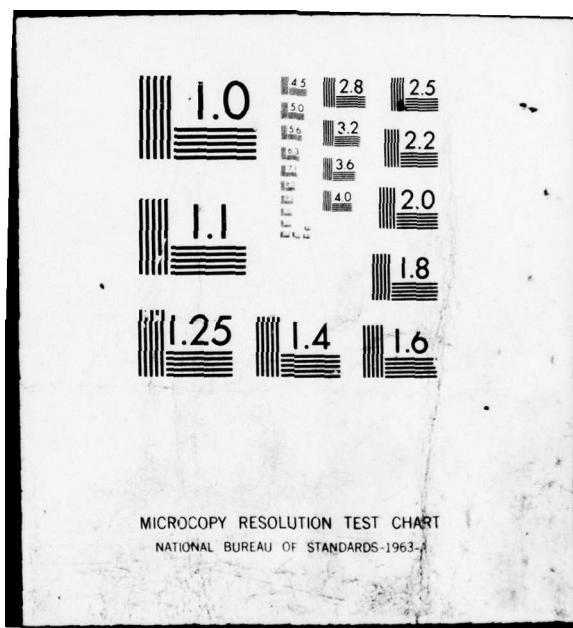
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THEORETICAL ANALYSIS OF  
TRANSONIC FLOW PAST  
UNSTAGGERED OSCILLATING CASCADES,

by

10 Peter Carlton/Clsen

11 Sept 1978

Thesis Advisor:

M.F. Platzer

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## (20. ABSTRACT Continued)

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Theoretical Analysis of  
Transonic Flow Past  
Unstaggered Oscillating Cascades  
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requirements for the degree of

AERONAUTICAL ENGINEER

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September 1978

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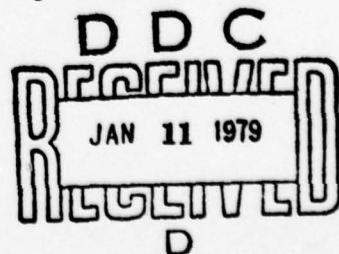
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ABSTRACT

This paper presents an independent verification of the collocation method as a technique for calculating the lift on an oscillating airfoil in an unstaggered cascade immersed in transonic flow. This method was originally proposed by Gorelov. Results presented here differ somewhat from those presented by him. Two formulations are shown; one is purely numerical, the second employs an analytic expansion for small frequency.

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LIST OF SYMBOLS

a	= local speed of sound	II
$a_0$	= speed of sound in the uniform flow	III
c	= blade semichord	III
$f_j$	= elementary function used in collocation solution	V
F	= specific energy ("head")	II
G	= function describing the surface of the airfoil as a function of time	III
H	= function specifying location of blade surface in the vertical axis	III
i	= $\sqrt{-1}$	II, III, IV, V
k	= Strouhal number, nondimensional frequency	III, IV, V
m	= $\sqrt{(\gamma+1)w}$	IV, V
M	= Mach number = $\frac{ \vec{V} }{a}$	III
n	= number of collocation points - 1, order of highest spanning function	V
p	= pressure	II
	= nondimensional interblade distance	V
R	= universal gas constant	II
R.P.	= "real part of"	III, IV, V
T	= temperature	II
t	= time, nondimensional time	II, III, IV, V
$U_0$	= uniform velocity from infinity	II, III, VI
u	= x-component of velocity	II, III
$u^0, u^1$	= interference vertical velocities due to reference and adjacent blades respectively, solved so as to satisfy the tangential flow conditions	V, VI

$u'$	= small disturbance velocity	III
$v$	= y-component of velocity	II, III, IV, V
$\vec{v}$	= general velocity vector	II
$v'$	= small disturbance velocity	III
$v^o, v^l$	= vertical velocities due to the reference and adjacent blades respectively, determined from the tangential flow condition	V
$w$	= $\tilde{\phi}_x$ , a constant used in Gorelov's approximation of the transonic flow potential	IV, V
$x$	= horizontal coordinate, may be non-dimensional	
$x_*$	= $m_p$ (transformed interblade distance in Gorelov's approximation)	V
$x_\ell$	= blade leading edge	IV
$x_o$	= center of pitch of the unstaggered cascade	IV, V
$y$	= vertical coordinate, may be non-dimensional	
$y, y_1$	= vertical coordinates attached to the reference and adjacent blades respectively, may be non-dimensional	IV, V
$z, z_1$	= transformed vertical coordinates used in Gorelov's approximation, attached to the reference and adjacent blades respectively. $z = my$ , $z_1 = my_1$	IV, V
$\frac{D}{Dt}$	= substantial derivative w.r.t. time $= \frac{\partial}{\partial t} + \frac{\partial x}{\partial t} \frac{\partial}{\partial x} + \frac{\partial y}{\partial t} \frac{\partial}{\partial y}$ $= \frac{\partial}{\partial t} + u \frac{\partial}{\partial x} + v \frac{\partial}{\partial y}$	II, III
$O(\omega^2)$	= "of the order of magnitude of $\omega^2$ "	V
$\alpha$	= angle of attack	II, III, IV, V
$\alpha_o$	= maximum amplitude of pitch oscillations	IV

$\gamma$	= ratio of specific heats, $c_p/c_v$	
$\delta_{io}$	= Dirac Delta function	V
	= 1 when $i = 0$	
	= 0 when $i \neq 0$	
$\eta$	= $\cos^{-1}(-x)$	V
$\eta_*$	= $\cos^{-1}(1-x_*)$	V
$\bar{\eta}$	= $\cos^{-1}(-s)$	V
$\hat{\eta}$	= $\cos^{-1}(x_* - x)$	V
$\theta_j^0, \theta_j^1$	= interference potential coefficients for reference and adjacent blades respectively	V
$\lambda$	= $k/m^2$	IV, V
$\mu_j^0, \mu_j^1$	= Fourier coefficients describing the motion of the reference and adjacent blades respectively	V
$\nu$	= angular frequency of oscillation	IV, V
$\rho$	= density	II
$\sigma$	= phase angle	V
$\tau$	= cascade solidity, $\frac{2}{p}$	V
$\Phi$	= general velocity potential	II, III, VI
$\Phi_0$	= uniform flow velocity potential	
$\tilde{\phi}$	= steady flow perturbation potential	III, IV
$\phi^0, \phi^1$	= perturbation potential in collocation solution due to reference and adjacent blades respectively	V, VI
$\Phi^0, \Phi^1$	= transformed potentials	V
$\psi$	= oscillatory flow potential	III, IV, V
$\Phi$	= transformed oscillatory potential in Gorenflo's coordinates corresponding to $\psi$	V
$\psi^0, \psi^1$	= interference potentials due to reference and adjacent blades respectively	V, VI

$\psi^0, \psi^1$  = transformed potentials in Gorelov's coordinates, corresponding to  $\psi^0$  and  $\psi^1$

$$\omega = \frac{k(1-m)^2}{m}$$

$\nabla = \vec{i} \frac{\partial}{\partial x} + \vec{j} \frac{\partial}{\partial y}$ , Gradient operator,  $\vec{i}$  and  $\vec{j}$  are unit vectors in the x and y direction respectively

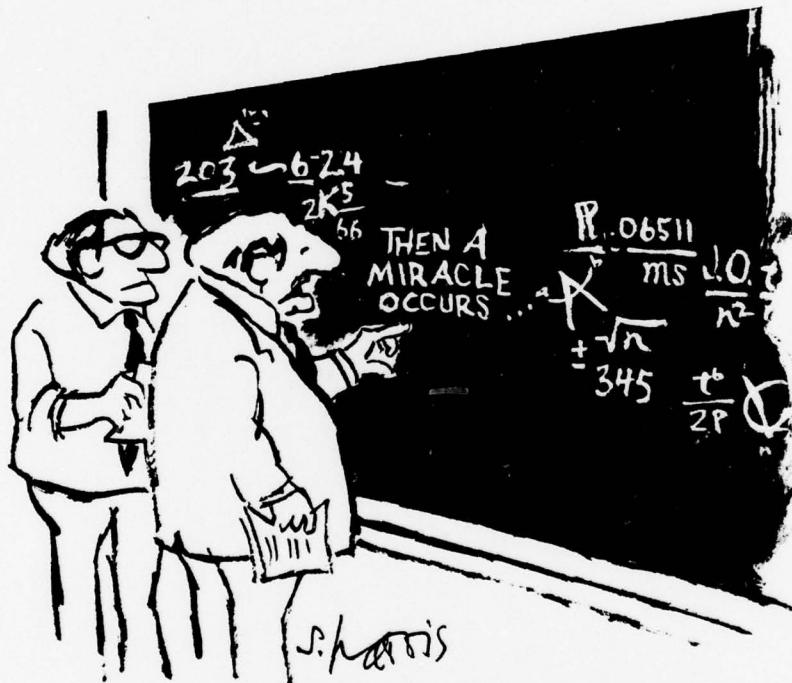
Computer Variables

DK = Reduced Frequency, k  
DLAMDA =  $\lambda$   
DM2 =  $m^2$   
DR = mp =  $x_*$   
ETA = n  
ETASTR =  $n_*$   
IPT = Print Parameter  
N = n  
NF = not used  
OFFSET = r  
OMEGA =  $\omega$   
QALPHA =  $0 + i(\lambda - k)$   
QCONST =  $e^{i\sigma}$   
QDCL =  $C_l_\alpha$   
QDCM =  $C_m_\alpha$   
QDK =  $0 + ik$   
QEXP =  $0 + i\lambda$   
QINTAP, QINTRP = Variables used to transmit boundary condition integrals  
Q1ABCF, Q1RBCF = Interference coefficients for adjacent and reference blades  
Q1COF = Right hand side vector in collocation solution  
Q1INT = Known matrix of integrals in collocation solution  
Q1RBP, Q1ABP = Not used  
Q2CP = Not used  
Q2EXP =  $e^{-i\lambda x}$

Q2PT	= Not used
RHO	= vertical displacement
RHP	= local input variable for rho
SIGMA	= $\sigma$
TAU	= $\tau$
XASTN	= Current x station in adjacent blade coordinates
XSTN	= Current x station in reference blade coordinates

ACKNOWLEDGEMENT

I would like to thank Professor Max F. Platzer without whose infinite patience this thesis would not have been written.



"I think you should be more explicit here in step two."

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## I. INTRODUCTION

The analysis of unsteady transonic flows in aircraft turbopropulsion is an area of intense current interest. Rising fuel prices and increasing thrust requirements both point toward the need of turbomachinery capable of performing well with transonic or supersonic internal flow. But, increased flow has increased both the costs and uncertainties of engine designs. Flutter problems have already become a major consideration in engine development. Problems unforeseen in earlier days of turbine engine production have caused long development delays, or forced acceptance of engines producing less than their initial design thrust. These uncertainties cannot be avoided when an attempt is made to extend the state of the art, but they can be reduced by extending the range of analytical modeling.

Such extension must now be done piecemeal. The three-dimensional flows in turbomachinery, including the simultaneous effects of boundary layers, rotation, finite blade thickness, spanwise Mach distributions, and shocks, are well beyond present capability. Perhaps one day complete analysis will be practical, but it is not today. The best that can be done now is to approach the problem from one aspect at a time. Flow through a two dimensional cascade has been a useful tool in this partial analysis.

This thesis was originally to have been an extension of the work of Elder [1] and Schlein [2] to the case of a staggered cascade. Their work, based on Teipel's [3] linearization of the unsteady transonic small perturbation equation, analyzed transonic flow through oscillating unstaggered cascades by use of the collocation method. While the problem was easy to state, it was difficult to solve. Both Elder and Schlein had encountered difficulty in employing the collocation method. Therefore, it was decided that verification of the basic collocation solution presented by Gorelov [4] using a different linearization would be a worthwhile goal in itself.

The following investigation presents a verification of the development in [4], along with numerical results and suggestions for further work.

### II. UNSTEADY TRANSONIC FLOW THEORY

Considering inviscid flow only, the following four equations govern the aerodynamic flow problem at hand:

The equation of state

$$p = \rho RT \quad (\text{II-1})$$

and the equations for the conservation of

$$1. \text{ Mass: } \operatorname{div}(\rho \vec{v}) + \frac{\partial \rho}{\partial t} = 0 \quad (\text{II-2})$$

$$2. \text{ Momentum: } \frac{D\vec{v}}{Dt} + \frac{1}{\rho} \nabla p = 0 \quad (\text{II-3})$$

$$3. \text{ Energy: } \frac{DS}{Dt} = 0 \quad (\text{II-4})$$

where

$\vec{v}$  = velocity

p = pressure

S = entropy

R = universal gas constant

T = temperature

t = time

$\rho$  = density

$$\frac{D}{Dt} = \frac{\partial}{\partial t} + \frac{\partial}{\partial x} \frac{\partial x}{\partial t} + \frac{\partial}{\partial y} \frac{\partial y}{\partial t}$$

The analysis starts with a uniform flow from infinity.

This flow has velocity  $U_0$  parallel to the x-axis. This

formulation can be simplified by working with the total velocity potential,  $\phi$ , where

$$u = \frac{\partial \phi}{\partial x} = \phi_x = x \text{ component of velocity} = \frac{\partial x}{\partial t} \quad (\text{II-5})$$

$$v = \frac{\partial \phi}{\partial y} = \phi_y = y \text{ component of velocity} = \frac{\partial y}{\partial t} \quad (\text{II-6})$$

Thus, the initial uniform flow is represented by the uniform flow potential

$$\phi_0 = U_0 x \quad (\text{II-7})$$

This notation may be applied to the conservation equations for mass and momentum. The equation for conservation of mass

$$\text{div}(\vec{\rho v}) + \frac{\partial \rho}{\partial t} = 0 \quad (\text{II-8})$$

becomes for two-dimensional unsteady flow

$$\frac{\partial(\rho u)}{\partial x} + \frac{\partial(\rho v)}{\partial y} + \frac{\partial \rho}{\partial t} = 0$$

$$[\frac{\partial \rho}{\partial t} + u \frac{\partial \rho}{\partial x} + v \frac{\partial \rho}{\partial y}] + \rho (\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y}) = 0 \quad (\text{II-9})$$

but

$$\frac{\partial \rho}{\partial t} + u \frac{\partial \rho}{\partial x} + v \frac{\partial \rho}{\partial y} = \frac{D\rho}{Dt}$$

and

$$u = \phi_x \quad \text{and} \quad v = \phi_y$$

Thus

$$\frac{D\rho}{Dt} + \rho(\phi_{xx} + \phi_{yy}) = 0$$

and

$$\phi_{xx} + \phi_{yy} = -\frac{1}{\rho} \frac{D\rho}{Dt} \quad (\text{II-10})$$

The speed of sound is given by

$$a^2 = \frac{dp}{d\rho}$$

Thus

$$\frac{D\rho}{Dt} = \frac{dp}{dp} \cdot \frac{Dp}{Dt} = \frac{1}{a^2} \frac{Dp}{Dt}$$

Applying this to equation (II-10) yields

$$\phi_{xx} + \phi_{yy} = -\frac{1}{\rho a^2} \frac{Dp}{Dt} \quad (\text{II-11a})$$

$$\phi_{xx} + \phi_{yy} = -\frac{1}{\rho a^2} (up_x + vp_y + p_t) \quad (\text{II-11b})$$

$$= -\frac{1}{\rho a^2} [(\nabla \phi) \cdot (\nabla p) + p_t] \quad (\text{II-11c})$$

where  $\nabla$  is the gradient operator,  $P_x = \frac{\partial P}{\partial x}$ ,  $P_y = \frac{\partial P}{\partial y}$ ,  $P_t = \frac{\partial P}{\partial t}$ .

Laying this aside for the moment, consider the momentum equation (II-3)

$$\frac{D\vec{v}}{Dt} + \frac{1}{\rho} \nabla p = 0$$

$$\frac{D\vec{v}}{Dt} = \frac{\partial \vec{v}}{\partial t} + (\vec{v} \cdot \nabla) \vec{v}$$

Thus

$$\rho \left[ \frac{\partial \vec{v}}{\partial t} + (\vec{v} \cdot \nabla) \vec{v} \right] = \nabla p \quad (\text{II-12})$$

$$\vec{v} = \nabla \Phi$$

so

$$\frac{\partial \vec{v}}{\partial t} = \frac{\partial}{\partial t}(\nabla \Phi) = \nabla \frac{\partial \Phi}{\partial t} \quad (\text{II-13})$$

and

$$(\vec{v} \cdot \nabla) \vec{v} = \nabla \frac{v^2}{2} - \vec{v} \times (\nabla \times \vec{v})$$

where

$$v^2 = u^2 + v^2$$

$$= \vec{v} \cdot \vec{v}$$

For irrotational flow

$$\nabla \times \vec{v} = 0$$

thus

$$(\vec{v} \cdot \nabla) \vec{v} = \frac{\nabla v^2}{2} \quad (\text{II-14})$$

Thus

$$\frac{\nabla p}{\rho} + \nabla [\phi_t + \frac{v^2}{2}] = 0 \quad (\text{II-15})$$

which after integration along a streamline becomes

$$\int \frac{dp}{\rho} + \phi_t + \frac{v^2}{2} = F(t) \quad (\text{II-16})$$

For uniform flow from infinity  $F(t) = \frac{1}{2} U_\infty^2$  and thus

the final result is

$$\int \frac{dp}{\rho} + \phi_t + \frac{v^2}{2} = \frac{1}{2} U_\infty^2 \quad (\text{II-17})$$

Differentiation with respect to t gives

$$p_t = -\rho(\phi_{tt} + \frac{1}{2} \frac{\partial v^2}{\partial t}) \quad (\text{II-18})$$

From (II-3)

$$-\nabla p = \rho \frac{\vec{D}\vec{v}}{Dt} \quad (\text{II-19})$$

Substitute (II-18) and (II-19) into (II-11c) to obtain

$$\phi_{xx} + \phi_{yy} = \frac{1}{a^2} [(\nabla\phi) \cdot \frac{\vec{D}\vec{v}}{Dt} + \phi_{tt} + \frac{1}{2} \frac{\partial v^2}{\partial t}] \quad (\text{II-20})$$

This may be further simplified

$$\begin{aligned} \frac{\vec{D}\vec{v}}{Dt} &= \frac{\partial \vec{v}}{\partial t} + (\vec{v} \cdot \nabla) \vec{v} \\ &= \frac{\partial \vec{v}}{\partial t} + \frac{1}{2} \nabla v^2 \end{aligned}$$

Hence

$$\phi_{xx} + \phi_{yy} = \frac{1}{a^2} [(\nabla\phi) \cdot (\frac{\partial \vec{v}}{\partial t} + \frac{1}{2} \nabla v^2) + \phi_{tt} + \frac{1}{2} \frac{\partial v^2}{\partial t}] \quad (\text{II-21})$$

Expanding terms

$$\begin{aligned} (\nabla\phi) \cdot (\frac{\partial \vec{v}}{\partial t}) &= \nabla\phi \cdot (\frac{\partial}{\partial t} \nabla\phi) = \phi_x \phi_{xt} + \phi_y \phi_{yt} \\ \nabla\phi \cdot \frac{\nabla v^2}{2} &= \frac{\phi_x^2 \phi_{xx}}{2} + \frac{\phi_y^2 \phi_{yy}}{2} + \frac{\phi_x \phi_y \phi_{xy}}{2} + \frac{\phi_y \phi_x \phi_{yx}}{2} \\ &= \frac{\phi_x^2 \phi_{xx}}{2} + \frac{\phi_y^2 \phi_{yy}}{2} + \frac{2 \phi_x \phi_y \phi_{xy}}{2} . \end{aligned}$$

$$\begin{aligned}
 \frac{1}{2} \frac{\partial V^2}{\partial t} &= \frac{1}{2} \frac{\partial}{\partial t} [(\nabla \Phi) \cdot (\nabla \Phi)] \\
 &= \frac{1}{2} [\Phi_x \Phi_{xt} + \Phi_{xt} \Phi_x + \Phi_y \Phi_{yt} + \Phi_{yt} \Phi_y] \\
 &= \Phi_x \Phi_{xt} + \Phi_y \Phi_{yt}
 \end{aligned}$$

The final result obtained by combining terms is

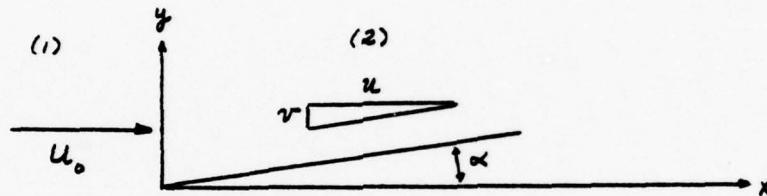
$$\begin{aligned}
 (1 - \frac{\Phi_x^2}{a^2})_{xx} + (1 - \frac{\Phi_y^2}{a^2})_{yy} - \frac{2\Phi_x \Phi_y \Phi_{xy}}{a^2} \\
 - \frac{2\Phi_x}{a^2} \Phi_{xt} - \frac{2\Phi_y}{a^2} \Phi_{yt} - \frac{1}{a^2} \Phi_{tt} = 0 \quad (\text{III-22})
 \end{aligned}$$

This result is valid for irrotational, inviscid, two-dimensional, unsteady, compressible flows where gravity has been neglected.

### III. SMALL PERTURBATION THEORY OF TRANSONIC FLOW

#### A. GENERAL CASE

A thin body at a small angle of attack will cause only a slight disturbance in the fluid. A flat plate is an example. Consider flow past a flat plate at angle of attack,  $\alpha$ .



The flow at (2) must be parallel to the plate. To achieve this, small disturbance velocities  $u'$  and  $v'$  must be added to the free stream velocity yielding

$$u = u_0 + u'$$

$$v = v'$$

The potential of the disturbed flow may be considered as the sum of the uniform flow potential,  $\Phi_0 = U_0 x$ , and the disturbance potential,  $\phi$

$$\Phi = \Phi_0 + \phi \quad (\text{III-1})$$

Thus

$$\phi_x = u_o + u' \quad (\text{III-2a})$$

$$\phi_y = v \quad (\text{III-2b})$$

If  $\phi$  is a function of time, then

$$\phi_t = \phi_t$$

This result may be substituted into (II-22) leading to

$$\begin{aligned} & [1 - \frac{(u_o + u')^2}{a^2}] \phi_{xx} + [1 - \frac{v^2}{a^2}] \phi_{yy} - 2 \frac{(u_o + u')v}{a^2} \phi_{xy} \\ & - 2 \frac{(u_o + u')}{a^2} \phi_{xt} - \frac{2v}{a^2} \phi_{yt} - \frac{1}{a^2} \phi_{tt} = 0 \quad (\text{III-3}) \end{aligned}$$

This expands to yield

$$\begin{aligned} & [1 - \frac{u_o^2 + 2u_o u' + u'^2}{a^2}] \phi_{xx} + [1 - \frac{v^2}{a^2}] \phi_{yy} - 2 \frac{(u_o v + u' v)}{a^2} \phi_{xy} \\ & - 2 \frac{(u_o + u')}{a^2} \phi_{xt} - \frac{2v}{a^2} \phi_{yt} - \frac{1}{a^2} \phi_{tt} = 0 \quad (\text{III-4}) \end{aligned}$$

Equation (III-4) may be further simplified as shown by Landahl [3]. All non-linear terms except the  $\phi_x \phi_{xx}$  product

term can be neglected yielding the following transonic small disturbance equation

$$\begin{aligned} & [(M^2 - 1) + (\gamma + 1)M^2 \frac{\phi_x}{U_0}] \phi_{xx} \\ & - \phi_{yy} + \frac{2M_0}{a_0} \phi_{xt} + \frac{1}{a_0} \phi_{tt} = 0 \end{aligned} \quad (\text{III-5})$$

where  $a_0$  is the velocity of sound in the free-stream, and  $\gamma$  is the ratio of specific heats, and  $M$  = Mach number.

#### B. BOUNDARY CONDITION

The tangential flow condition requires that the flow be tangent to the airfoil surface at each instant of time. This means that no fluid may flow through the surface of the airfoil and is expressed by the condition

$$\frac{DG}{Dt} = 0 \quad \text{on } G(x, y, t) \quad (\text{III-6})$$

where

$G(x, y, t)$  describes the surface of the body as a function of time.

For a thin airfoil restricted to small oscillations, this may be written as

$$G = y - H(x, t) \quad (\text{III-7})$$

where

$H(x,t)$  is the function describing the position of the airfoil.

$H(x,t)$  can be written for harmonic pitch oscillations as

$$H(x,t) = R.P. [\alpha_0 (x - x_0) e^{ivt}] \quad (\text{III-8})$$

where the time-varying angle of attack  $\alpha(t)$  is given by

$$\alpha(t) = R.P. [\alpha_0 e^{ivt}]$$

and  $\alpha_0$  = maximum amplitude of pitch oscillation

$x_0$  is the pitch axis

$v$  is the angular frequency of oscillation

$$i = \sqrt{-1}$$

R.P. = "real part of"

Inserting (III-8) into the flow tangency condition (III-6) gives, after linearization,

$$\phi_y(x,0) = v(x,0) = \alpha_0 [U_0 + iv(x - x_0)] e^{ivt} \quad (\text{III-9})$$

on  $y = 0$

This is a condition for the normal velocity to be prescribed at the airfoil's mean position  $y = 0$ .

#### C. NONDIMENSIONALIZATION

The terms in equations (III-5) and (III-9) are dimensional. For the following calculations it is convenient to use non-dimensional quantities. Define non-dimensional time and length to be

$$\bar{x} = \frac{x}{c}$$

$$\bar{y} = \frac{y}{c}$$

(III-10)

$$\bar{t} = \frac{tU_o}{c}$$

where

$U_o$  = uniform velocity from infinity

$c$  = reference length (blade semichord).

The velocity potential in equation (III-5) may be non-dimensionalized as follows. Let

$$\bar{\phi} = \frac{\phi}{U_o c}$$

Hence:

$$\phi = U_o c \bar{\phi}$$

$$\phi_x = U_o c \bar{\phi}_x \left( \frac{1}{c} \right)$$

$$= U_o \bar{\phi}_x \quad \text{(III-11)}$$

and similarly for the other derivatives in (III-5), yielding

$$[(M^2 - 1) + (\gamma + 1) M^2 \bar{\phi}_{\bar{x} \bar{x}}] \bar{\phi}_{\bar{y} \bar{y}} - \bar{\phi}_{\bar{y} \bar{y}} + 2M^2 \bar{\phi}_{\bar{x} \bar{t}} + M^2 \bar{\phi}_{\bar{t} \bar{t}} = 0$$

(III-12)

This equation is non-dimensional.

The boundary condition given in equation (III-9) may be non-dimensionalized in a similar fashion

$$\bar{\phi}_y(x, 0) = \alpha_o [U_o + i\nu(x - x_o)] e^{i\nu t} \quad (III-9)$$

Thus

$$U_o c \bar{\phi}_y = \alpha_o [U_o + ik \frac{U_o}{c} (c\bar{x} - c\bar{x}_o)] e^{i\nu t} \quad (III-13)$$

where

$$k = \frac{\nu c}{U_o} = \text{Strouhal number or reduced frequency}$$

$$U_o c \bar{\phi}_y \cdot \frac{1}{c} = \alpha_o U_o [1 + ik(\bar{x} - \bar{x}_o)] e^{ik\bar{t}} \quad (III-14)$$

Thus

$$\bar{\phi}_y = \bar{v}(\bar{x}, 0) = \alpha_o [1 + ik(\bar{x} - \bar{x}_o)] e^{ik\bar{t}} \quad (III-15)$$

Because the final operations are linear in  $\alpha_o$ , set  $\alpha_o = 1$ , yielding

$$\bar{\phi}_y = \bar{v}(\bar{x}, 0) = [1 + ik(\bar{x} - \bar{x}_o)] e^{ik\bar{t}} \quad (III-16)$$

The overbars denoting nondimensional quantities will be dropped from the remainder of the paper. All further quantities shall be assumed appropriately non-dimensional. This yields the following final equations

$$[(M^2 - 1) + (\gamma + 1) M^2 \phi_x] \phi_{xx} - \phi_{yy} + 2M^2 \phi_{xt} + M^2 \phi_{tt} = 0 \quad (III-17)$$

and

$$\phi_y(x, 0) = v(x, 0) = [1 + ik(x - x_0)] e^{ikt} \quad (III-18)$$

where

all quantities are nondimensional and

$$\alpha_0 = 1$$

#### D. HARMONIC OSCILLATIONS

In the case of harmonic oscillations, equation (III-17) may be simplified still further.

Let

$$\phi = \tilde{\phi} + R.P. [\psi e^{ikt}]$$

where

$\tilde{\phi}$  = non-dimensional steady flow potential

$\psi$  = non-dimensional oscillatory flow potential

R.P. = "real part of"

Equation (III-17) then becomes

$$(1-M^2)\psi_{xx} + \psi_{yy} - M^2(\gamma+1)[\tilde{\phi}_x\psi_{xx} + \tilde{\phi}_{xx}\psi_x] + M^2k^2\psi - 2imMk^2\psi_x = 0 \quad (\text{III-19})$$

For  $M$  close to 1, this is a nonlinear mixed elliptic-hyperbolic partial differential equation with variable coefficients, the exact type depending on  $\tilde{\phi}_x$  and  $\tilde{\phi}_{xx}$ . However, because flutter analysis is primarily concerned with the stability of small perturbations about a steady flow, the oscillatory component may be assumed small compared to the steady flow potential and therefore the product term  $\psi_x\psi_{xx}$  may be neglected, yielding,

$$(1-M^2)\psi_{xx} + \psi_{yy} - M^2(\gamma+1)[\tilde{\phi}_x\psi_{xx} + \tilde{\phi}_{xx}\psi_x] = 2imMk^2\psi_x + M^2k^2\psi = 0 \quad (\text{III-20})$$

#### IV. LINEARIZATION OF THE GOVERNING EQUATION

The basic flutter equation, (III-20), is still a non-linear, mixed elliptic-hyperbolic partial differential equation with variable coefficients and difficult to solve. It may yet be further simplified.

##### A. BASIC SOLUTION

For  $M = 1$ , equation (III-20) becomes

$$\psi_{yy} - (\gamma+1) [\tilde{\phi}_x \psi_{xx} + \tilde{\phi}_{xx} \psi_x] - 2ik\psi_x + k^2\psi = 0 \quad (IV-1)$$

Now assume

$$\tilde{\phi}_x \approx w = \text{constant} \quad (IV-2)$$

$$\tilde{\phi}_{xx} \approx 0$$

throughout the interblade channel. Setting

$$\tilde{\phi}_x(\gamma+1) = w(\gamma+1) = m^2 \quad (IV-3)$$

yields

$$m^2\psi_{xx} - \psi_{yy} + 2ik\psi_x - k^2\psi = 0 \quad (IV-4)$$

The solution to this equation is found in Garrick and Rubinow [5]

$$\psi(x,y) = -\frac{1}{m} \int_{x_\lambda}^{x-my} v(s) J_0 [\omega \sqrt{(x-s)^2 - (my)^2}] e^{i\lambda(s-x)} ds \quad (IV-5a)$$

for  $y > 0$

and

$$\psi(x,y) = \frac{1}{m} \int_{x_\lambda}^{x+my} v(s) J_0 [\omega \sqrt{(x-s)^2 - (my)^2}] e^{i\lambda(s-x)} dx \quad (IV-5b)$$

for  $y < 0$

where

$$v(x) = \lim_{y \rightarrow 0} \frac{\partial}{\partial y} \psi(x,y) .$$

$v(x)$  may be obtained directly from the tangential flow boundary condition, and

$$\omega = \frac{k^2(1-m^2)}{4m}$$

$$\lambda = \frac{k}{m^2} \sqrt{1+m^2} \approx \frac{k}{m^2} \quad (\text{where this paper employs the approximation used by Gorelov [4]})$$

$x_\lambda$  = blade leading edge,

Gorelov [4], has proposed a further simplification.

Set

$$z = my \quad (\text{IV-6a})$$

$$\Psi(x, z) = \psi(x, y) e^{i\lambda x}. \quad (\text{IV-6b})$$

Equation (IV-4) then becomes

$$\Psi_{xx} - \Psi_{zz} + \omega^2 \Psi = 0$$

with solution

$$\Psi(x, z) = -\frac{1}{m} \int_{x_\ell}^{x-z} v(s) J_0[\omega \sqrt{(x-s)^2 - z^2}] e^{i\lambda s} ds \quad (\text{IV-7a})$$

$x_\ell$

$$z > 0$$

$$\Psi(x, z) = \frac{1}{m} \int_{x_\ell}^{x-z} v(s) J_0[\omega \sqrt{(x-s)^2 - z^2}] e^{i\lambda s} ds \quad (\text{IV-7b})$$

$x_\ell$

$$z < 0$$

where

$$v(x) = m e^{-i\lambda x} \lim_{z \rightarrow 0} \Psi_z(x, z)$$

$v(x)$  is obtained from the tangential flow boundary condition.

For a thin body immersed in the flow, the solutions for  $y > 0$ ,  $z > 0$ , and  $y < 0$ ,  $z < 0$  apply above the body along left-running Mach lines, or below along right running Mach lines respectively.

#### B. BOUNDARY CONDITIONS

##### 1. Flow Tangency Condition

The boundary condition comes from the tangential flow condition, (III-18)

$$v(x) = [1 + ik(x - x_0)] \quad (\text{IV-8})$$

##### 2. Upstream Condition

The final linearized equation is a hyperbolic differential equation with boundary condition

$$\psi(x, y) = 0 \quad (\text{IV-9})$$

when

$$x - x_l < |my|$$

for the solution shown in equations (IV-5) or

$$\Psi(x, z) = 0$$

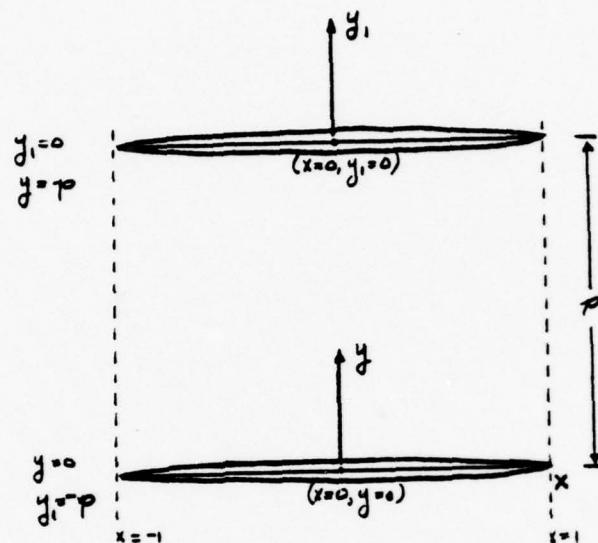
**when**

$$x - x_l < |z|$$

**for the solution shown in equations (IV-7).**

## V. PROBLEM FORMULATION

### A. CO-ORDINATE SYSTEM



Assume the geometry shown above. Both blades are thin airfoils of semichord  $c$ . All measurements are non-dimensional, normalized to  $c$ . The  $(x, y)$  co-ordinate system has its origin at the center of the reference (lower) blade. The  $(x, y_1)$  system is centered at the middle of the adjacent (upper) blade. The origin of the  $(x, y_1)$  system is located at  $(0, p)$  in the reference system. Generalizing this convention, the same symbols shall be used for the same quantities on both blades. Where discrimination is required, the quantity associated with the adjacent blade will be marked with

superscript <sup>1</sup>, the quantity associated with the reference blade will be either unsuperscripted or marked with a superscript <sup>0</sup>.

Each blade is assumed to perform a small amplitude harmonic oscillation about its mid-chord point. Both blades are assumed to have identical reduced frequencies,  $k$ , and the motion of the adjacent blade lags that of the reference blade by a phase angle  $\sigma$ .

The blades are immersed in a uniform flow from the left at  $M = 1$ . The objective is to determine the oscillatory pressure distributions and aerodynamic forces generated by the blades' oscillations. Cascade solidity,  $\tau = 2/p$ .

#### B. BOUNDARY CONDITIONS

##### 1. Upstream Condition

$$\psi = 0 \quad \text{whenever}$$

$$x + 1 < |my| \quad (\text{V-1})$$

and

$$x + 1 < |my_1|, \text{ simultaneously}$$

##### 2. Flow Tangency Condition

Along the reference blade

$$\lim_{y \rightarrow 0} \psi_y(x, y) = (1 + ikx) \quad (\text{V-2a})$$

Along the adjacent blade

$$\lim_{y_1 \rightarrow 0} \frac{\psi_{y_1}(x, y_1)}{y_1} = (1 + ikx)e^{i\sigma} \quad (V-2b)$$

where

$\sigma$  is the phase angle between the blades oscillations

### C. BASIC SOLUTION TECHNIQUE

Assume that the unsteady potential,  $\psi$ , may be written as the sum of four components

$$\psi(x, y) = \phi^0(x, y) + \psi^0(x, y) + \phi^1(x, y_1) + \psi^1(x, y_1) \quad (V-3)$$

where:

$\phi^0$  = potential due to the reference blade alone, known from equation (IV-7)

$\phi^1$  = potential due to the adjacent blade alone, known from equation (IV-7)

$\psi^0$  = interference potential required to satisfy tangential flow condition along reference blade, unknown

$\psi^1$  = interference potential required to satisfy tangential flow condition along adjacent blade, unknown.

This total potential must satisfy the tangential flow condition at the plane of both the reference and adjacent blades. Thus

$$\phi_y^0(x, y=0) + \phi_{y_1}^1(x, y_1=-p) + \psi_y^0(x, y=0) + \psi_{y_1}^1(x, y_1=-p) \quad (V-4a)$$

$$= (1 + ikx)$$

at the reference blade, and

$$\begin{aligned}\phi_y^o(x, y=p) + \phi_{y_1}^l(x, y_1=0) + \psi_y^o(x, y=p) + \psi_{y_1}^l(x, y_1=0) \\ = (1 + ikx) e^{i\sigma} \quad (V-4b)\end{aligned}$$

at the adjacent blade.

But from the unsteady potential solution for a single oscillating blade one has

$$\phi_y^o(x, y=0) = 1 + ikx \quad (V-5a)$$

and

$$\phi_{y_1}^l(x, y_1=0) = (1 + ikx) e^{i\sigma} \quad (V-5b)$$

Thus

$$\phi_{y_1}^l(x, y_1=-p) + \psi_{y_1}^l(x, y_1=-p) + \psi_y^o(x, y=0) = 0 \quad (V-6a)$$

along the reference blade, and

$$\phi_y^o(x, y=p) + \psi_y^o(x, y=p) + \psi_{y_1}^l(x, y_1=0) = 0 \quad (V-6b)$$

along the adjacent blade.

From equation (IV-7)

$$\phi^O(x, y) = -\frac{1}{m} \int_{-1}^{x-my} v^O(s) J_0[\omega \sqrt{(x-s)^2 - (my)^2}] e^{i\lambda(s-x)} ds$$

$$y > 0 \quad (V-7a)$$

$$= \frac{1}{m} \int_{-1}^{x+my} v^O(s) J_0[\omega \sqrt{(x-s)^2 - (my)^2}] e^{i\lambda(s-x)} ds$$

$$y < 0 \quad (V-7b)$$

where

$$v^O(s) = 1 + iks$$

$$\lambda = k/m^2$$

$$m = (\gamma+1)w$$

$$\omega = \frac{k^2(1-m^2)}{m^4}$$

w = mean value of  $\phi_x$  in the channel

$$\phi^1(x, y_1) = -\frac{1}{m} \int_{-1}^{x-my_1} v^1(s) J_0[\omega \sqrt{(x-s)^2 - (my_1)^2}] e^{i\lambda(s-x)} ds$$

$$y_1 > 0 \quad (IV-8a)$$

$$= \frac{1}{m} \int_{-1}^{x+my_1} v^1(s) J_0[\omega \sqrt{(x-s)^2 - (my_1)^2}] e^{i\lambda(s-x)} ds$$

$$y_1 < 0 \quad (IV-8b)$$

where

$$v^1(s) = (1 + iks)e^{i\sigma}$$

Henceforth attention will be restricted to the flow within the channel,  $0 \leq y \leq p$ ,  $-p \leq y_1 \leq 0$  leaving (IV-7a) and (IV-8b) as the governing equations of interest.

The two interference potentials are assumed to have forms identical to (IV-7a) and (IV-8b).

Set

$$\psi^0(x,y) = -\frac{1}{m} \int_{-1}^{x-my} u^0(s) J_0[\omega \sqrt{(x-s)^2 - (my)^2}] e^{i\lambda(s-x)} ds$$
$$y > 0 \quad (V-9a)$$

$$\psi^1(x,y_1) = \frac{1}{m} \int_{-1}^{x+m} u^1(s) J_0[\omega \sqrt{(x-s)^2 - (my_1)^2}] e^{i\lambda(s-x)} ds$$
$$y_1 < 0 \quad (V-9b)$$

where

$u^0(s)$  and  $u^1(s)$  are unknown functions to be determined so as to satisfy equations (V-6)

Substitution of (V-7a), (V-8b) and (V-9) into (V-6) yields

$$u^0(x) + \psi^1_{y_1}(x, y_1 = -p) + \phi^0_{y_1}(x, y_1 = -p) = 0 \quad (V-10a)$$

$$u^1(x) + \psi_Y^0(x, y=p) + \phi_Y^0(x, y=p) = 0 \quad (V-10b)$$

Recalling Gorelov's transformation discussed in [4] and shown in equations (IV-6) above, set

$$\phi^0(x, z) = \phi^0(x, y) e^{i\lambda x}$$

$$\phi^1(x, z_1) = \phi^1(x, y_1) e^{i\lambda x}$$

$$\psi^0(x, z) = \psi^0(x, y) e^{i\lambda x}$$

$$\psi^1(x, z_1) = \psi^1(x, y_1) e^{i\lambda x}$$

where  $z = my$

$$z_1 = my_1$$

Then

$$\frac{e^{i\lambda x}}{m} u^0(x) + \phi_{z_1}^1(x, z_1 = -x_*) + \psi_{z_1}^1(x, z_1 = -x_*) = 0 \quad (V-11a)$$

and

$$\frac{e^{i\lambda x}}{m} u^1(x) + \phi_z^0(x, z = x_*) + \psi_z^0(x, z = x_*) = 0 \quad (V-11b)$$

where

$$x_* = mp .$$

To employ the collocation method, assume that  $u^1(x)$  and  $u^0(x)$  can be approximated as the sum of a set of elementary functions  $f_j$  so that

$$u^0(x) \approx \sum_{j=0}^n \theta_j^0 f_j(x) \quad (V-12a)$$

$$u^1(x) \approx \sum_{j=0}^n \theta_j^1 f_j(x) \quad (V-12b)$$

where  $f_j(x) = 0$  when  $x \leq x_* - 1$

Note that here both  $u^0$  and  $u^1$  are expressed in terms of the same elementary functions,  $f_j$ .

Because of the slightly supersonic nature of the problem observe that  $u^0(x) = 0$  and  $u^1(x) = 0$  when  $x \leq x_* - 1$ .

Equations (V-12) may now be rewritten as

$$e^{i\lambda x} \sum_j \theta_j^0 f_j(x) + \frac{\partial}{\partial z_1} \int_{x_* - 1}^{x+z_1} \sum_j \theta_j^1 f_j(s) J_0[\omega \sqrt{(x-s)^2 - z_1^2}] e^{i\lambda s} ds$$

$$= -e^{i\sigma} \frac{\partial}{\partial z_1} \int_{-1}^{x+z_1} (1+iks) J_0[\omega \sqrt{(x-s)^2 - z_1^2}] e^{i\lambda s} ds$$

$$\text{at } z_1 = -x_* \quad (V-13a)$$

$$\begin{aligned}
 & e^{i\lambda x} \sum_j \theta_j^1 f_j(x) - \frac{\partial}{\partial z} \int_{x_*-1}^{x-z} \sum_j \theta_j^0 f_j(s) J_0[\omega \sqrt{(x-s)^2 - z^2}] e^{i\lambda s} ds \\
 & = \frac{\partial}{\partial z} \int_{-1}^{x+z} (1+iks) J_0[\omega \sqrt{(x-s)^2 - z^2}] e^{i\lambda s} ds \\
 & \text{at } z = x_* \tag{V-13b}
 \end{aligned}$$

where

$$f_j(x) = 0 \text{ for } x \leq x_* - 1$$

This simplifies to

$$\begin{aligned}
 & e^{i\lambda x} \sum_j \theta_j^0 f_j(x) + \sum_j \theta_j^1 \left\{ \frac{\partial}{\partial z_1} \int_{x_*-1}^{x+z_1} f_j(s) J_0[\omega \sqrt{(x-s)^2 - z_1^2}] e^{i\lambda s} ds \right\} \\
 & = -e^{i\sigma} \frac{\partial}{\partial z_1} \int_{-1}^{x+z_1} (1+iks) J_0[\omega \sqrt{(x-s)^2 - z_1^2}] e^{i\lambda s} ds \\
 & \text{at } z_1 = -x_* \tag{V-14a}
 \end{aligned}$$

and

$$\begin{aligned}
 & e^{i\lambda x} \sum_j \theta_j^1 f_j(x) - \sum_j \theta_j^0 \left\{ \frac{\partial}{\partial z} \int_{x_*-1}^{x-z} f_j(s) J_0[\omega \sqrt{(x-s)^2 - z^2}] e^{i\lambda s} ds \right\} \\
 & = \frac{\partial}{\partial z} \int_{-1}^{x-z} (1+iks) J_0[\omega \sqrt{(x-s)^2 - z^2}] e^{i\lambda s} ds \tag{V-14b}
 \end{aligned}$$

at  $z = x_*$  where

$$f_j(x) = 0 \text{ for } x \leq x_* - 1.$$

Performing the indicated differentiation yields

$$e^{i\lambda x} \sum_j \theta_j^1 f_j(x) + \sum_j \theta_j^1 \left\{ \int_{x_* - 1}^{x - x_*} f_j(s) \frac{J_1 [\omega \sqrt{(x-s)^2 - x_*^2}] \omega x_* e^{i\lambda s} ds}{\sqrt{(x-s)^2 - x_*^2}} \right.$$

$$\left. + f_j(x - x_*) e^{i\lambda(x - x_*)} \right\}$$

$$= e^{i\sigma} \left\{ - \int_{-1}^{x - x_*} (1 + iks) \frac{\omega x_* J_1 [\omega \sqrt{(x-s)^2 - x_*^2}] e^{i\lambda s} ds}{\sqrt{(x-s)^2 - x_*^2}} \right.$$

$$\left. - [1 + ik(x - x_*)] e^{i\lambda(x - x_*)} \right\} \quad (V-15a)$$

and

$$e^{i\lambda x} \sum_j \theta_j^1 f_j(x) + \sum_j \theta_j^0 \left\{ \int_{x_* - 1}^{x - x_*} f_j(s) \frac{J_1 [\omega \sqrt{(x-s)^2 - x_*^2}] \omega x_* e^{i\lambda s} ds}{\sqrt{(x-s)^2 - x_*^2}} \right.$$

$$\left. + f_j(x - x_*) e^{i\lambda(x - x_*)} \right\}$$

$$= - \int_{-1}^{x - x_*} (1 + iks) \frac{\omega x_* J_1 [\omega \sqrt{(x-s)^2 - x_*^2}] e^{i\lambda s} ds}{\sqrt{(x-s)^2 - x_*^2}} - [1 + ik(x - x_*)] e^{i\lambda(x - x_*)} \quad (V-15b)$$

Gorelov's formulation, equations [2.6, 2.7, 2.8, and 2.9] of [4], can be obtained directly from equations (V-15) by substituting

$$f_j(x) = \cos j\eta - \cos j\eta_*$$

$$v^0 = \sum_{j=0}^n \mu_j^0 \cos j\eta$$

$$v^1 = \sum_{j=0}^n \mu_j^1 \cos j\eta$$

where:

$$\eta = \cos(-x)$$

$$\eta_* = \cos(1-x_*)$$

In comparing the two systems care must be taken to note the differing symbols and coordinate systems. The corresponding quantities are:

Here

in [4]

$$\theta_j^0, \theta_j^1$$

$$v_{o\sigma}, v_{i\sigma}$$

$$\mu_j^0, \mu_j^1$$

$$\theta_{o\sigma}, \theta_{i\sigma}$$

$$\eta = \cos^{-1}(-x)$$

$$\eta = \cos^{-1}(1-x)$$

$$\eta_* = \cos^{-1}(1-x_*)$$

$$\eta_* = \cos^{-1}(1-x_*)$$

$$z, z_1$$

$$y, y_1$$

$$j$$

$$\sigma$$

$$\sigma$$

$$\psi$$

Here  $-1 \leq x \leq 1$ ; in [4]  $0 \leq x \leq 2$ . This transformation accounts for the differing definitions of  $\eta$ . Making the substitutions results in the system

$$\sum_{j=0}^n \{\theta_j^0 [\cos j\eta - (1-\delta_{01}) \cos j\eta_*]\}$$

$$+ \theta_j^1 \int_{x_*-1}^{x-x_*} \frac{\partial}{\partial z_1} J_0 [\omega \sqrt{(x-s)^2 - z_1^2}] [\cos j\hat{\eta} - (1-\delta_{10}) \cos j\eta_*] e^{i\lambda s} ds$$

$$+ \theta_i^1 [\cos j\bar{\eta} - (1-\delta_{i0}) \cos j\eta_*] e^{i\lambda(x-x_*)}\}$$

$$= - \sum_{j=0}^n \{-\mu_j^1 \int_{-1}^{x-x_*} \frac{\partial}{\partial z_1} J_0 [\sqrt{(x-s)^2 - z_1^2}] \cos j\hat{\eta} e^{i\lambda s} ds$$

$$- \mu_j^1 \cos j\bar{\eta} e^{i\lambda(x-x_*)}\} \quad (V-16a)$$

$$\text{at } z_1 = -x_*$$

$$x > x_* - 1$$

and

$$\begin{aligned}
& \sum \left\{ \theta_j^1 [\cos j\eta - (1-\delta_{0j}) \cos j\eta_*] \right. \\
& + \theta_j^0 \int_{x_*-1}^{x-x_*} \frac{\partial}{\partial z} J_0 [\omega \sqrt{(x-s)^2 - z^2}] [\cos j\eta - (1-\delta_{0j}) \cos j\eta_*] e^{i\lambda s} ds \\
& \left. + \theta_j^0 [\cos j\bar{\eta} - (1-\delta_{0j}) \cos j\eta_*] e^{i\lambda(x-x_*)} \right\} \\
= & \sum_{j=0}^n \left\{ \mu_j^0 \int_{-1}^{x-x_*} \frac{\partial}{\partial z} J_0 [\omega \sqrt{(x-s)^2 - z^2}] \cos j\hat{\eta} e^{i\lambda s} ds \right. \\
& \left. - \mu_j^0 \cos j\bar{\eta} e^{i\lambda(x-x_*)} \right\} \tag{V-16b}
\end{aligned}$$

where:

$$\eta = \text{arc cos } (-x)$$

$$\eta_* = \text{arc cos } (1-x_*)$$

$$\bar{\eta} = \text{arc cos } (-s)$$

$$\hat{\eta} = \text{arc cos } (-x+x_*)$$

$$\delta_{0j} = \text{Dirac } \delta \text{ function} = \begin{cases} 1 & \text{for } j = 0 \\ 0 & \text{for } j \neq 0 \end{cases}$$

Given the change in coordinates and notation, this system  
is equivalent to that shown in [4].

This was the system programmed for computer solution. Because the function,  $f_j(x)$ , is unaffected by the differentiation with respect to  $y$  (or  $z$ ) the exact form used need not be specified, so that the system shown in (V-15) may be programmed with  $f$  undetermined. A subroutine may be written to return the function desired and the remaining program left perfectly general. In the program developed with this thesis both the Gorelov functions shown above and the Legendre polynomials were employed. All the integrals may now be evaluated at  $n+1$  points,  $x_i$ , on both blades in  $mp-1 < x_i < 1$  and the resulting linear system solved for  $\theta_j^0$  and  $\theta_j^1$ ,  $j=1,2,\dots,n+1$ .

The interference potentials may be constructed by taking

$$\psi^0(x,z) = \frac{-1}{m} \int_1^{x-z} [\sum_j \theta_j^0 f_j(x)] J_0[\omega \sqrt{(x-s)^2 - z^2}] e^{i\lambda s} ds$$

(V-17a)

$$\psi^1(x,z_1) = \frac{1}{m} \int_{-1}^{x+z_1} [\sum_j \theta_j^1 f_j(x)] J_0[\omega \sqrt{(x-s)^2 - z_1^2}] e^{i\lambda s} ds$$

(V-17b)

where

$$f_j(x) = 0, \quad \text{for all } x \leq x_* - 1$$

Once the potentials have been calculated as outlined above, the surface pressure may be calculated using the relationship

$$C_p = -2(\psi_x + ik\psi) \quad (V-18a)$$

$$= -2[\psi_x + i(k-\lambda)\psi]e^{-i\lambda x} \quad (V-18b)$$

Because all the plates are assumed to be in steady oscillation with uniform phase shift,  $\sigma$ , between neighboring plates, then

$$v^1(x) = v^0(x)e^{i\sigma}, \quad u^1(x) = u^0(x)e^{i\sigma}, \quad \psi(x,y) = -\psi(x,-y)$$

in this case

$$C_{k\alpha} = 2 \int_{-1}^1 [\psi_x(x,+0) + i(k-\lambda)\psi(x,+0)]e^{-i\lambda x} dx \quad (V-19)$$

where

$$\begin{aligned} \psi(x_1,+0) &= \frac{-1}{m} \int_{-1}^x [(1+iks)+u^0(s)]J_0[\omega(x-s)]e^{i\lambda(s)} ds \\ &+ \frac{e^{i\sigma}}{m} \int_{-1}^{x-x_*} [(1+iks)+u^0(s)]J_0[\omega\sqrt{(x-s)^2-(x_*)^2}]e^{i\lambda(s)} ds \end{aligned}$$

where  $x_* = mp.$  (V-20)

Results from this approach, in the form of values of  $C_{k\alpha}$  for  $k = 0.1$ , at various values of  $w$  are presented in the results for

approximations based both on Gorelov's formulation, and on the Legendre polynomials.

D. COLLOCATION SOLUTION OF THE POTENTIAL EQUATION EXPANDED FOR SMALL  $k$ ,

In order to provide a partially independent check of the results of the main program, the Gorelov function representation of the collocation solution was expanded for small  $k$ , and solved at two collocation points,  $n = 2$ . The resulting potentials, and partial derivatives with respect to  $x$  and  $y$  were then used to replace the corresponding numerical routines in the main program. The output resulting from the approximations were compared with the purely numerical results obtained from the computer program.

1. Solution For The Unknown Potential Coefficients

The basic system of linear equations used to determine the unknown coefficients is

$$\frac{1}{m} e^{i\lambda x} u^0 + \phi_{z_1}^1 + \psi_{z_1}^1 = 0, \quad z = 0, \quad z_1 = -x_* = -mp$$

(V-21a)

$$\frac{1}{m} e^{i\lambda x} u^1 + \phi_z^0 + \psi_z^0 = 0, \quad z_1 = 0, \quad z = x_* = mp \quad (V-21b)$$

where

$$0 = u^0(x) = u^1(x) \quad \text{when } x \leq -1+x_*$$

otherwise

$$u^o = \sum_{j=1}^n \theta_j^o (\cos j\eta - \cos j\eta_*) + \theta_o^o$$

$$u^l = \sum_{j=1}^n \theta_j^l (\cos j\eta - \cos j\eta_*) + \theta_o^l$$

where

$$\eta = \arccos(-x)$$

$$\eta_* = \arccos(1-x_*) .$$

Thus, for  $n = 2$ , the system becomes

$$e^{i\lambda x} \{ \theta_o^o + \theta_1^o (\cos \eta - \cos \eta_*) \}$$

$$+ \frac{\partial}{\partial z_1} \int_{-1}^{x+z_1} v^l(s) J_o [\omega \sqrt{(x-s)^2 - z_1^2}] e^{i\lambda s} ds$$

$$+ \frac{\partial}{\partial z_1} \int_{x_*-1}^{x+z_1} u^l(s) J_o [\omega \sqrt{(x-s)^2 - z_1^2}] e^{i\lambda s} ds = 0$$

$$z_1 = -x_*$$

(V-22a)

along the reference blade and

$$e^{i\lambda x} \{ \theta_0^1 + \theta_1^1 (\cos n - \cos n_*) \} - \frac{\partial}{\partial z} \int_{-1}^{x-z} v^0(s) J_0 [\omega \sqrt{(x-s)^2 - z^2}] e^{i\lambda s} ds$$

$$- \frac{\partial}{\partial z} \int_{x_*-1}^{x-z} u^0(s) J_0 [\omega \sqrt{(x-s)^2 - z^2}] e^{i\lambda s} ds . \quad (V-22b)$$

along the adjacent blade

$$\text{where } u^0(s) = \theta_0^0 + \theta_1^0 (\cos n - \cos n_*)$$

$$u^1(s) = \theta_0^1 + \theta_1^1 (\cos n - \cos n_*)$$

$$v^0(s) = 1 + iks$$

$$v^1(s) = (1 + iks)e^{i\sigma}$$

For  $k$  sufficiently small, this system may be further simplified by the following approximations

$$J_0 [\omega \sqrt{(x-s)^2 - z^2}] \approx 1 - O(\omega^2) \approx 1 \quad (V-23a)$$

$$J_0 [\omega \sqrt{(x-s)^2 - z^2}] \approx 1 - O(\omega^2) \approx 1 \quad (V-23b)$$

$$e^{i\lambda x} \approx 1 + i\lambda x - O(\lambda^2 x^2) \approx 1 + i\lambda x$$

$$e^{i\lambda s} \approx 1 + i\lambda s - O(\lambda^2 s^2) \approx 1 + i\lambda s$$

where  $O(\omega^2)$  means "of the order of magnitude of  $\omega^2$ "

$$-1 \leq x \leq 1, \quad -1 \leq s \leq x-x_*$$

$$\lambda = k/m^2, \quad \omega^2 = \frac{k^2(1-m^2)}{m^4}$$

The interference source distributions may be replaced by

$$u^0(s) = \theta_0^0 + \theta_1^0(-s + x_* - 1)$$

$$u^1(s) = \theta_0^1 + \theta_1^1(-s + x_* - 1) .$$

If higher order terms are neglected, the result is a system linear in  $k$  and  $\lambda$

$$(1+i\lambda x)[\theta_0^0 + \theta_1^0(-x-1+x_*)] + \frac{\partial e^{i\sigma}}{\partial z_1} \int_{-1}^{x+z_1} (1+iks)(1+i\lambda s) ds$$

$$+ \frac{\partial}{\partial z_1} \int_{x_*-1}^{x+z_1} [\theta_0^1 + \theta_1^1(-s-1+x_*)](1+i\lambda s) ds = 0$$

$z_1 = -x_*$  (V-24a)

$$(1+i\lambda x)[\theta_0^1 + \theta_1^1(-x-1+x_*)] - \frac{\partial}{\partial z} \int_{-1}^{x-z} (1-iks)(1+i\lambda s) ds$$

$$- \frac{\partial}{\partial z} \int_{x_*-1}^{x-z} [\theta_0^0 + \theta_1^0(-s-1+x_*)](1+i\lambda s) ds = 0$$

$z = x_*$  (V-24b)

Product terms containing  $(k\lambda) = \frac{k^2}{m}$  may be neglected as of higher order in  $k$ , yielding

$$(1+i\lambda x) [\theta_0^0 + \theta_1^0(-x-1+x_*)] + \frac{\partial e^{i\sigma}}{\partial z} \int_{-1}^{x+z_1} [1+i(\lambda+k)s] ds$$

$$+ \frac{\partial}{\partial z} \int_{x_*-1}^{x+z_1} \{ [\theta_0^1 + \theta_1^1(-s-1+x_*)]$$

$$+ i\lambda s [\theta_0^1 + \theta_1^1(-s-1+x_*)] \} ds = 0$$

$$z_1 = -x_* \quad (V-25a)$$

$$(1+i\lambda x) [\theta_0^1 + \theta_1^1(-x-1+x_*)] - \frac{\partial}{\partial z} \int_{-1}^{x-z} [1+i(\lambda+k)s] ds$$

$$- \frac{\partial}{\partial z} \int_{x_*-1}^{x-z} \{ [\theta_0^0 + \theta_1^0(-s-1+x_*)] + i\lambda s [\theta_0^0 + \theta_1^0(-s-1+x_*)] \} ds = 0$$

$$z = x_* \quad (V-25b)$$

Evaluating the indicated derivatives yields

$$(1+i\lambda x) [\theta_0^0 + \theta_1^0(-x-1+x_*)] + [1+i(k+\lambda)(x-x_*)] e^{i\sigma}$$

$$+ [\theta_0^1 + \theta_1^1(2x_*-x-1)] + i\lambda(x-x_*) [\theta_0^1 + \theta_1^1(2x_*-x-1)] \} = 0$$

$$(V-26a)$$

$$(1+i\lambda x) [\theta_0^1 + \theta_1^1 (-x-1+x_*)] + [1+i(\lambda+k)(x-x_*)]$$

$$+ \{ [\theta_0^0 + \theta_1^0 (2x_* - x - 1)] + i\lambda(x-x_*) [\theta_0^0 + \theta_1^0 (2x_* - x - 1)] \} = 0$$

(V-26b)

Thus:

$$\theta_0^0 (1+i\lambda x) + \theta_1^0 [(-x-1+x_*) + i\lambda x (-x-1+x_*)]$$

$$+ \theta_0^1 [1 + i\lambda(x-x_*)] + \theta_1^1 [(2x_* - x - 1) + i\lambda(x-x_*) (2x_* - x - 1)]$$

$$= -e^{i\sigma} [1 + i(k+\lambda)(x-x_*)] \quad (V-27a)$$

$$\theta_0^1 [1 + i\lambda x] + \theta_1^1 [(-x+1-x_*) + i\lambda x (-x+1-x_*)]$$

$$+ \theta_0^0 [1+i\lambda x(x-x_*)] + \theta_1^0 [(2x_* - x - 1) + i\lambda(x-x_*) (2x_* - x - 1)]$$

$$= -[1 + i(\lambda+k)(x-x_*)] \quad (V-27b)$$

This system may be solved at two points,  $x_1$  and  $x_2$ , for  $\theta_0^0$ ,  $\theta_1^0$ ,  $\theta_0^1$ , and  $\theta_1^1$ .

2. Calculation Of The Potential

The potential is given by

$$\psi(x,y) = \psi(x,z)e^{-i\lambda x} \quad (V-28)$$

where

$$\begin{aligned} \psi(x,z) &= -\frac{1}{m} \int_{-1}^x [v^0(s) + u^0(s)] J_0[\omega \sqrt{(x-s)^2}] e^{i\lambda s} ds \\ &\quad + \frac{1}{m} \int_{-1}^{x-x_*} [v^1(s) + u^1(s)] J_1[\omega \sqrt{(x-s)^2 - x_*^2}] e^{i\lambda s} ds. \end{aligned}$$

$$u^0(s) = u^1(s) = 0 \quad \text{for all } s \leq x_* - 1.$$

Thus

$$\begin{aligned} \psi(x,z) &= -\frac{1}{m} \int_{-1}^x v^0(s) J_0[\omega \sqrt{(x-s)^2}] e^{i\lambda s} ds \\ &\quad - \frac{1}{m} \int_{x_*-1}^x u^0(s) J_0[\omega \sqrt{(x-s)^2}] e^{i\lambda s} ds \\ &\quad + \frac{1}{m} \int_{-1}^{x-x_*} v^1(s) J_1[\omega \sqrt{(x-s)^2 - x_*^2}] e^{i\lambda s} ds \\ &\quad + \frac{1}{m} \int_{x_*-1}^{x-x_*} u^1(s) J_1[\omega \sqrt{(x-s)^2 - x_*^2}] e^{i\lambda s} ds \quad (V-29) \end{aligned}$$

Making the same small frequency approximations as  
in the previous section yields

$$\Psi(x, z) = -\frac{1}{m} \int_{-1}^x v^0(s) (1+i\lambda s) ds$$

$$- \frac{1}{m} \int_{x_*-1}^x u^0(s) (1+i\lambda s) ds$$

$$+ \frac{1}{m} \int_{-1}^{x-x_*} v^1(s) (1+i\lambda s) ds$$

$$+ \frac{1}{m} \int_{x_*-1}^{x-x_*} u^1(s) (1+i\lambda s) ds \quad (V-30)$$

From the general formulation

$$\Psi = \phi^1 + \phi^0 + \psi^1 + \psi^0 \quad (V-31)$$

Thus, along the reference blade

$$-m\phi^0(x, z=0) = \int_{-1}^x v^0(s) (1+i\lambda s) ds = \int_{-1}^x (1+iks) (1+i\lambda s) ds$$

$$= \int_{-1}^x [1+i(k+\lambda)s] ds = [s+i(k+\lambda)\frac{s^2}{2}]_{-1}^x \quad (V-32)$$

$$= x + i\frac{(k+\lambda)}{2}x + 1 - i\frac{(k+\lambda)}{2}$$

$$\phi^O(x, z=0) = -\frac{1}{m} \left[ (1+x) + i \left( \frac{k+\lambda}{2} \right) (x^2 - 1) \right] \quad (V-33)$$

By inspection

$$\phi^1(x, z_1 = -x_*) = \frac{e^{i\sigma}}{m} \left\{ (1+x-x_*) + i \left( \frac{k+\lambda}{2} \right) [(x-x_*)^2 - 1] \right\} \quad (V-34)$$

$$-m\Psi^O(x, z=0) = \int_{x_*-1}^x u^O(s) (1+i\lambda s) ds \quad (V-35)$$

$$= \int_{x_*-1}^x [\theta_0^O + \theta_1^O (-s+x_*-1)] (1+i\lambda s) ds$$

$$= \int_{x_*-1}^x \theta_0^O (1+i\lambda s) + \theta_1^O (-s+x_*-1) (1+i\lambda s) ds$$

$$= \theta_0^O \left[ (x+1-x_*) + \frac{i\lambda}{2} (x^2 - x_*^2 + 2x_* - 1) \right]$$

$$+ \int_{x_*-1}^x \theta_1^O [-x+x_*-1+i\lambda(-s^2+sx_*-s)] ds$$

$$= \theta_0^O \left[ (x+1-x_*) + \frac{i\lambda}{2} (x^2 - x_*^2 + 2x_* - 1) \right]$$

$$+ \theta_1^O \left\{ \left( -\frac{s^2}{2} + sx_* - s \right) + i\lambda \left( -\frac{s^3}{3} + \frac{s^2 x_*}{2} - \frac{s^2}{2} \right) \right\} \Big|_{x_*-1}^x \quad (V-36)$$

$$\begin{aligned}
&= \theta_0^O [ (x+1-x_*) + \frac{i\lambda}{2} (x^2 - x_*^2 + 2x_* - 1) ] \\
&\quad + \theta_1^O \{ (-\frac{x^2}{2} + x_* - x) - [\frac{(x_*-1)^2}{2} + x_* (x_*-1) - (x_*-1)] \\
&\quad + i\lambda [ (-\frac{x^3}{3} + \frac{x^2 x_* - x^2}{2}) + \frac{(x_*-1)^3}{3} x_* \frac{(x_*-1)^2}{2} + \frac{(x_*-1)^2}{2} ] \} \\
&= \theta_0^O [ (x+1-x_*) + \frac{i\lambda}{2} (x^2 - x_*^2 + 2x_* - 1) ] \\
&\quad + \theta_1^O - \frac{x^2}{2} + x(x_*-1) + \frac{(x_*-1)^2}{2} - (x_*-1)^2 \\
&\quad + i\lambda [ -\frac{x^3}{3} + x^2 \frac{(x_*-1)}{2} + \frac{x_*^3}{3} - x_*^2 + x_* - \frac{1}{3} \\
&\quad - \frac{x_*^3}{2} + x_*^2 - \frac{x_*}{2} + \frac{x_*^2}{2} - x_* + \frac{1}{2} ] \} \\
&= \theta_0^O [ (x+1-x_*) + \frac{i\lambda}{2} (x^2 - x_*^2 + 2x_* - 1) ] \\
&\quad + \theta_1^O \{ -\frac{x^2}{2} + x(x_*-1) - \frac{(x_*-1)^2}{2} \\
&\quad + i\lambda [ -\frac{x^3}{3} + \frac{x^2(x_*-1)}{2} - \frac{x_*^3}{6} + \frac{x_*^2}{2} - \frac{x_*}{2} + \frac{1}{6} ] \}
\end{aligned}$$

$$\begin{aligned}
\psi^0 = & -\frac{1}{m}\{\theta_0^0[(x+1-x_*) + \frac{i\lambda}{2}(x^2 + x_*^2 + 2x_* - 1)] \\
& + \theta_1^0\{-\frac{x^2}{2} + x(x_* - 1) - \frac{(x_* - 1)^2}{2} \\
& + i\lambda[-\frac{x^3}{3} + \frac{x^2(x_* - 1)}{2} - \frac{(x_* - 1)^3}{6}]\}\} \quad (V-37)
\end{aligned}$$

$m\psi^1(x, z_1 = -x_*)$  may be evaluated by substituting  $x - x_*$  for  $x$  in the expression for  $m\psi^0$  and exchanging  $\theta_0^1$  and  $\theta_1^1$  for  $\theta_0^0$  and  $\theta_1^0$

$$\begin{aligned}
m\psi^1(x, z_1 = -x_*) = & \theta_0^1\{[(x - x_*) + 1 - x_*] + \frac{i\lambda}{2}[(x - x_*)^2 - x_*^2 + 2x_* - 1]\} \\
& + \theta_1^1\{-\frac{(x - x_*)^2}{2} + (x - x_*)(x_* - 1) - \frac{(x_* - 1)^2}{2} \\
& + i\lambda[-\frac{(x - x_*)^3}{3} + \frac{x^2(x_* - 1)}{2} - \frac{(x_* - 1)^3}{6}]\} \quad (V-38)
\end{aligned}$$

$$\begin{aligned}
= & \theta_0^1(x+1-2x_*) + \frac{i\lambda}{2}[(x^2 - 2xx_* + x_*^2) - x_*^2 + 2x_* - 1] \\
& + \theta_1^1\{-\frac{(x^2 - 2xx_* + x_*^2)}{2} + (x - \frac{3}{2}x_* + \frac{1}{2})(x_* - 1) \\
& + i\lambda[-\frac{x^3}{3} + x^2x_* - xx_*^2 + \frac{x_*^3}{3} + (x^2 - 2xx_* + x_*^2)\frac{(x_* - 1)}{2} \\
& - \frac{(x_* - 1)^3}{6}]\}
\end{aligned}$$

$$\begin{aligned}
&= \theta_0^1 \{ (x+1-2x_*) + \frac{i\lambda}{2} [x^2 - 2xx_* + 2x_* - 1] \\
&\quad + \theta_1^1 \{ -\frac{x^2}{2} + xx_* - \frac{x_*^2}{2} + x(x-1) + \frac{1}{2}(1-3x_*)(x_*-1) \\
&\quad + i\lambda [-\frac{x^3}{3} + x^2 \frac{(3x_*-1)}{2} + x(-\frac{x_*^2-2xx_*^2-1}{2}) \\
&\quad + \frac{x_*^3}{3} + \frac{x_*^3}{2} - \frac{1}{2}] \}
\end{aligned}$$

$$\begin{aligned}
&\psi^1(x, z_1 = -x_*) \\
&= \frac{1}{m} \{ \theta_0^1 \{ (x+1-2x_*) + \frac{i\lambda}{2} [x^2 - 2xx_* + 2x_* - 1] \\
&\quad + \theta_1^1 \{ [-\frac{x^2}{2} + x(2x_*-1) - xx_*^2 + 2x_* - \frac{1}{2}] \\
&\quad + i\lambda [-\frac{x^3}{3} + x^2 \frac{(3x_*-1)}{2} + x \frac{(-3x_*^2-1)}{2} + \frac{5x_*^3}{6} - \frac{1}{2}] \} \}
\end{aligned}$$

(V-39)

A comparison of the results for the full program and the approximation is given below for  $k = 0.01$ ,  $w = 0.05$ ,  $\sigma = \rho$ ,  $n = 2$ , yielding  $\omega^2 \approx 6.1 \times 10^{-3}$ ,  $\lambda \approx 0.083$

	Full program	Approx
$x = .1285$		
$\phi$	-5.363, .327i	-5.379, .549i
$\phi_x$	-8.634, .361i	-8.66, .379i
$x = .5643$		
$\phi$	-9.754, .809i	-9.804, 1.000i
$\phi_x$	-12.234, .692i	-12.272, .7608i
$C_{\ell_\alpha}$	+31.658, -3.1032i	$C_{\ell_\alpha}$ 31.7019, -3.2165i

## VI. RESULTS

The collocation method was used to solve the partial differential equation resulting from the Gorelov approximation of transonic potential flow in an unstaggered cascade. The system was solved using both the spanning functions proposed by Gorelov in [4], resulting in the equations (V-16); and the Legendre polynomials, resulting in equations (V-15) with  $f_j$  replaced by the Legendre polynomial,  $P_j$ . The resulting values of  $C_{\ell\alpha}$  for  $k = .1$ ,  $\tau = 1$ ,  $\sigma = 1$  and seven collocation points on each blade are presented in figures VI-2, VI-3, and VI-4.

Figure VI -1 presents a diagram which is useful in commenting on the other results. This shows the location of the collocation points and first three interference reflections as a function of  $w$  expressed as a percentage of that portion of the chord subject to reflection. The collocation points are equally spaced throughout this interval, 12.5% from the leading edge of the interference zone, 12.5% between each pair of points and 12.5% from the blade trailing edge. The independent variable,  $w$ , is plotted vertically so that the dependent variable, percent of chord subject to interference, may be more conveniently visualized along the blade. (The curves are not precisely linear, but are very nearly so in the range shown.)

Figure VI -2 shows the  $C_{\ell\alpha}$  calculated with  $k = 0.0$  in comparison with the results obtained from Ackeret theory.

Agreement is good where there is no reflection and the portion of the blade subject to interference is affected by a constant interference potential,  $w \geq 0.11$ . Throughout the rest of the curve the results calculated here oscillate above and below the theoretical values. This appears to be due to the discrete nature of the approximation used in the collocation method. Rarely is the fraction of the chord subject to interference reflection equal to the fraction of the collocation points which feel it. Where the collocation point fraction lags, as near  $w = 0.6$ , the collocation results are lower than those due to Ackeret theory. When the collocation point fraction leads, as it does for  $w \leq .04$  and briefly for  $w \approx .08$ , the the collocation results are higher than those due to Ackeret theory. The fault appears to be an intrinsic feature of the small number of points sampled. This results in a set of coefficients similar to those which would be obtained from a generalized Fourier series based on the integration of the Taylor series expansion about each point. This obviously cannot be a good approximation when both the function and its derivative are discontinuous at the reflections.

Figures VI -3 and VI -4 show the results of using Legendre polynomials and Gorelov's functions as spanning functions. The results for both formulations are identical. Gorelov's results are presented for comparison. Agreement is good for  $w > 0.05$  except for an anomalous point, marked A .

It is believed that this anomaly is due to the location of the first reflection just ahead of the last collocation point (cf. "A" on Figure VI -1). This will yield a very small contribution from the reflection potential to the linear system from which the collocation points are determined. The resulting system will have a large dynamic range and may be ill-conditioned.

The discrepancy between these results, and those in [3] for  $w < 0.5$  is still unexplained, as is the outlying value for  $w = 0.5$ .

The discontinuities in the imaginary results are believed to be due primarily to the reflection/collocation interaction explained above.

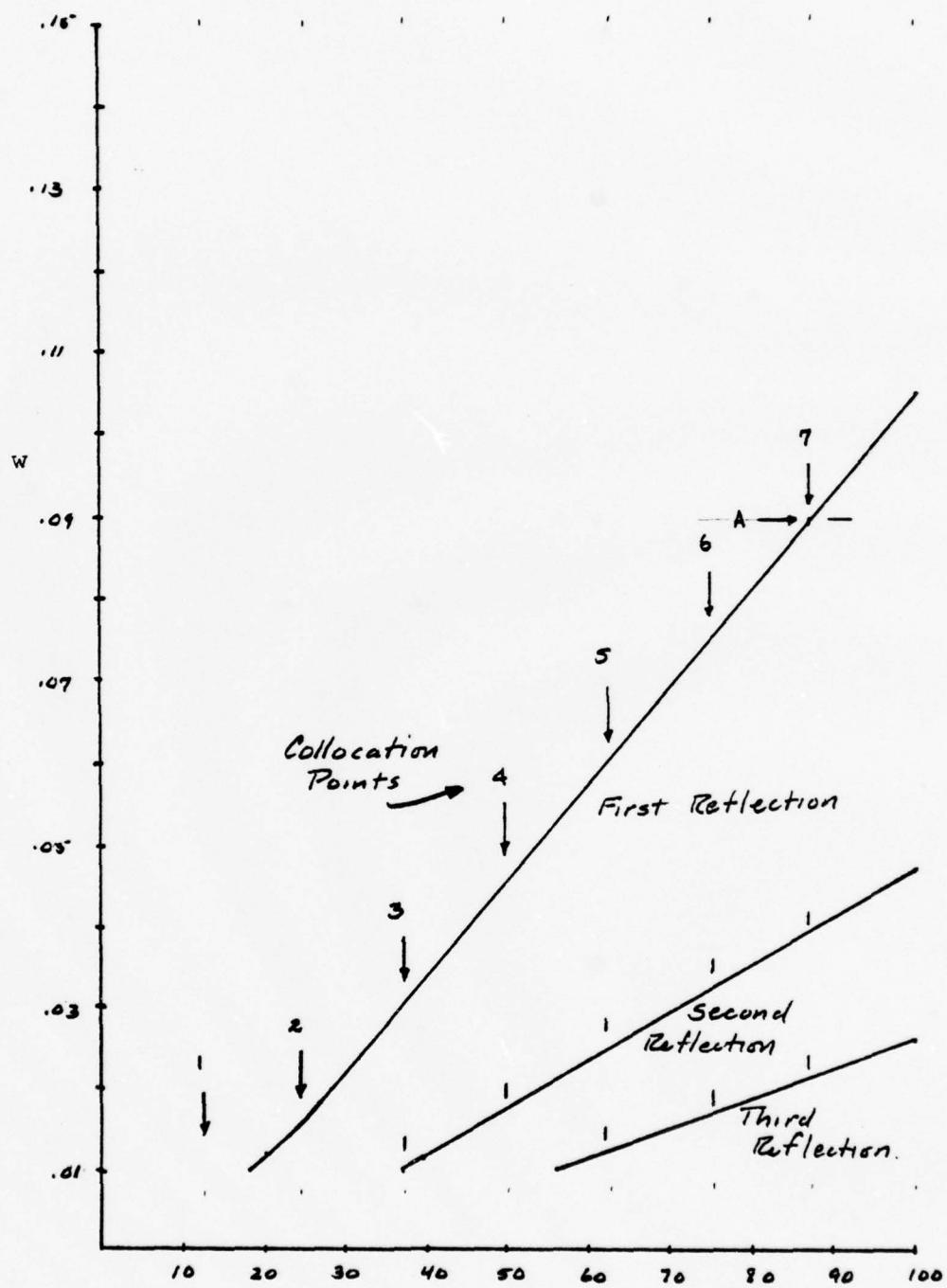


FIGURE VI-1. Location of Reflections and Collocation Points Shown as Percent of Chord Subject to Interference

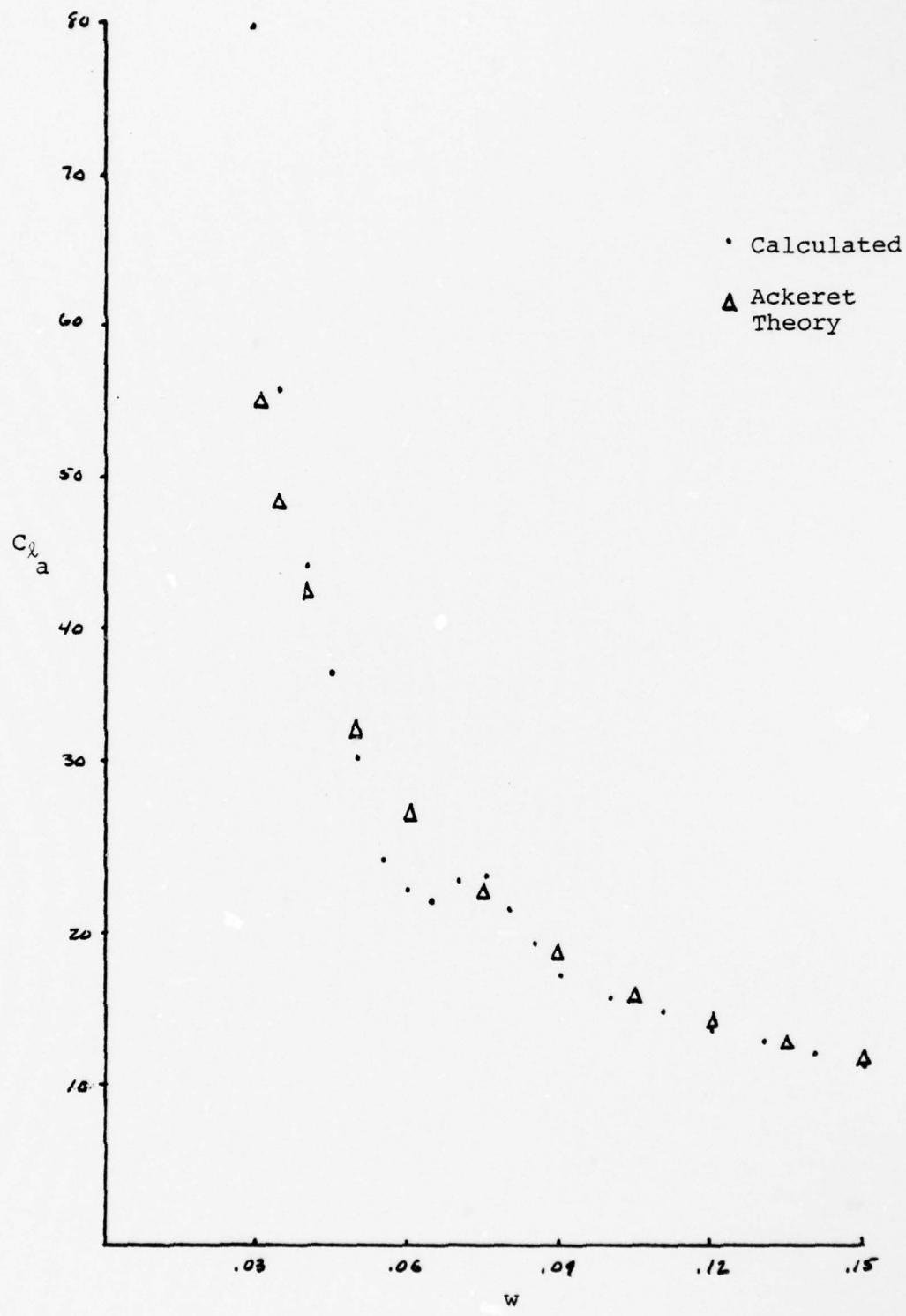


FIGURE VI-2. Comparison of  $C_{l\alpha}$ -vs- $w$  to that Obtained from Ackeret Theory

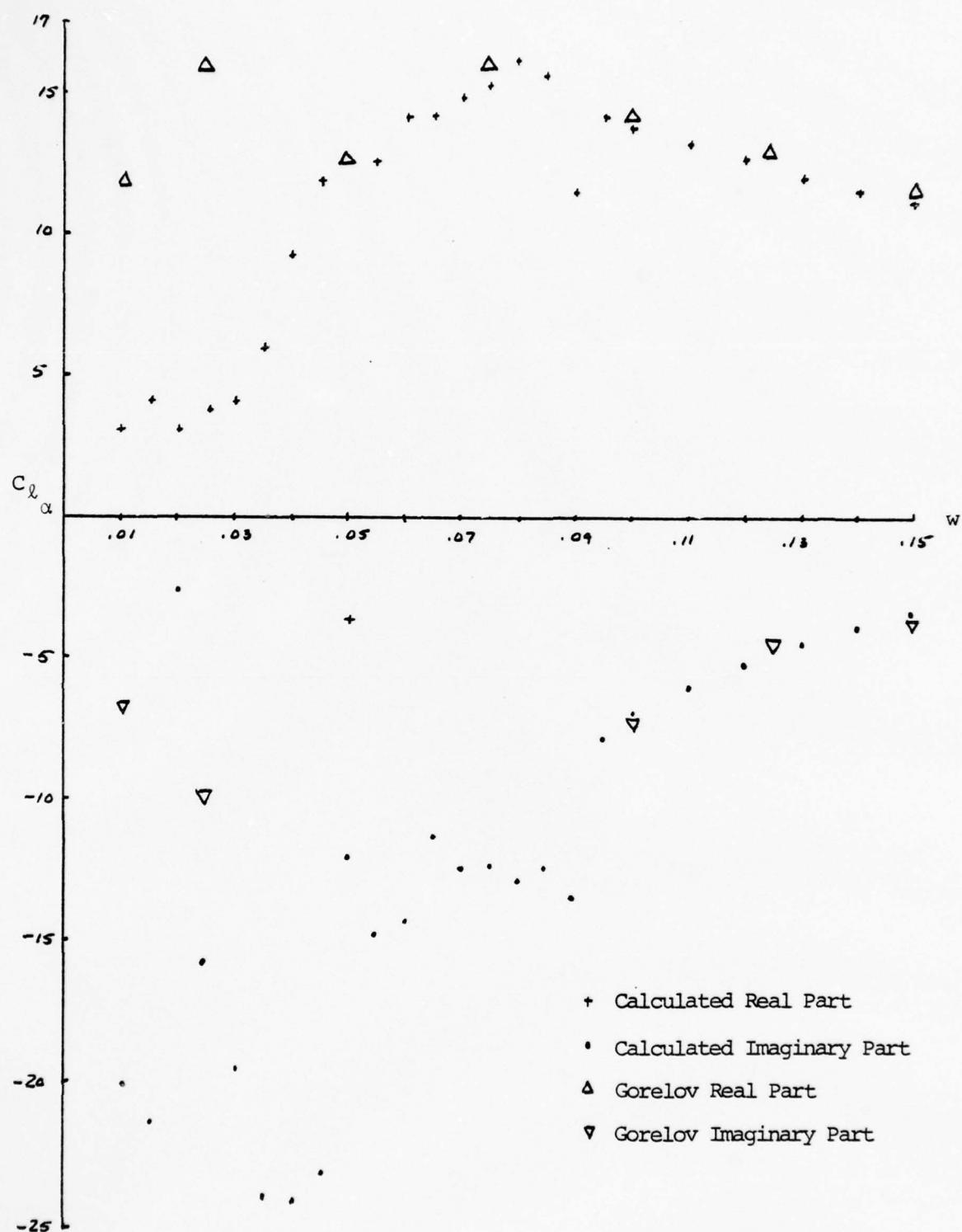


FIGURE VI-3. Plot of  $C_{l\alpha}$ -vs-w, Legendre Polynomials  
 $k = 0.1, \tau = 1.0, \sigma = \pi, n = 7$   
 compared with Gorelov's results

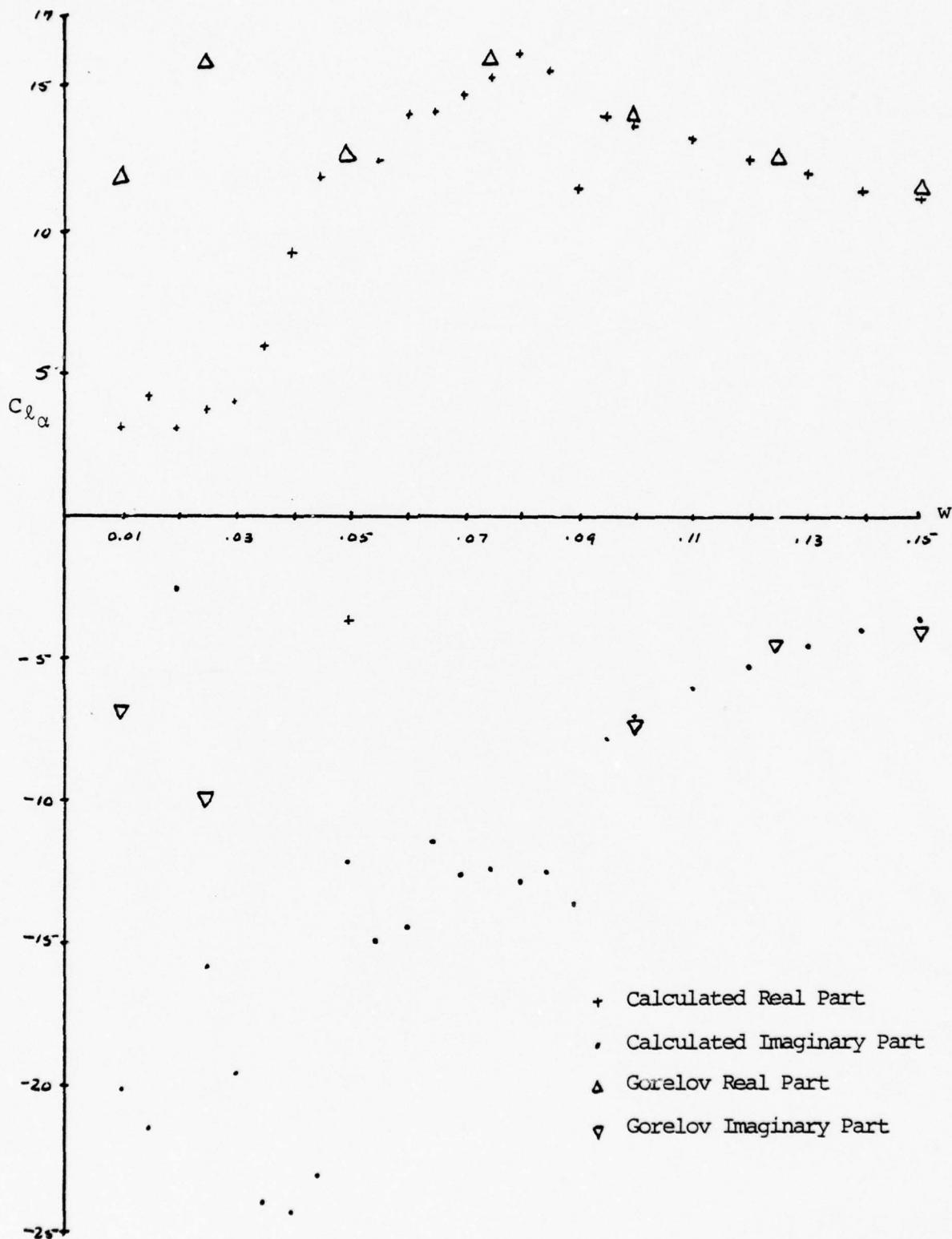


FIGURE VI-4. Plot of  $C_{l\alpha}$ -vs-w, Gorelov Spanning Function  
 $k = 0.1, \tau = 1.0, \sigma = \pi, n = 7$   
 compared with Gorelov's results

## VII . RECOMMENDATIONS

There are two recommendations to be made about the techniques used in the collocation method, and a new area in which it might be employed.

The program developed in the course of writing this thesis employs adaptive Simpson's integration to calculate the elements of a completely determined system. This system is solved to provide the coefficients of the spanning functions.

Two improvements may be made:

1. The Simpson's integration scheme may be replaced by a Gaussian integrator. Experience has shown that several thousand function evaluations are required by the Simpson's integration routine when  $C_{t_\alpha}$  is to be evaluated for small w. This entails large amounts of computer time and leads to increased accumulations of numerical error. Use of Gaussian integration would probably improve both of these characteristics with little loss of accuracy.

2. The present program treats a completely determined system of dimension  $2n+1$  by  $2n+1$ , and then solves that system to find the collocation coefficients. This procedure has worked satisfactorily in this thesis, but may not work as well at higher frequencies where the final linear system of equations may be ill-conditioned. As an alternative, it is recommended that the boundary conditions be applied at more than n points, say m points, where m is twice or three times as many points, and that the least squares technique be used to determine the

the collocation coefficients which give the minimum square error over-all. This may be thought of as "sampling more data" in order to get more information about the unknown function. The present program could be easily modified in this regard by replacing the spanning function matrix, Q1ZINT, by a new matrix of the form

$$Q1ZINT' = X^T X$$

where  $X$  is the new  $m$  by  $n+1$  ( $m > n+1$ ) matrix, and replacing the present right-hand-side vector, Q1COF with

$$Q1COF' = X^T Y$$

where  $Y$  is the new  $m$  by 1 right-hand-side vector. An alternative would be to employ a prepackaged statistical linear regression routine after either modifying the routine to accept complex data, or transforming the present system into a larger system of real numbers only.

The new area in which the collocation method might be employed is the calculation of the potential flow about a staggered cascade. The method could be employed to calculate both the potential in the channel and above the upper blade. The program presented has been designed to enable the

calculation of flow within the channel of a staggered cascade. Unfortunately, there was not enough time to extend the study to this case.

APPENDIX A  
PROGRAM DESCRIPTION

This section describes the computer program used to calculate the interference solution to the Gorelov linearization for unsteady transonic flow in a channel. The program written in IBM Fortran IV with the basic structure outlined by Stevens [5]. The basic points are:

1. Organization of the program into small subroutines, each of which performs a specific task.
2. Transmission data to and from subroutines via a formal parameter argument list. No common statements are used.

The end objective is code which is both easy to modify and maintain.

Each subroutine is designed with optional diagnostic printing of its input and output. This is controlled by the parameter IPT. The diagnostic output is printed (only) if  $IPT > 0$ . Each routine accepts IPT, sets IOT = IPT - 1, and then passes IOT as the print parameter to routines it calls. By this method, diagnostic output can be "cascaded" to any desired level. Large initial values of IPT should be avoided because of the spectacular amount of output which can be generated by the double integrals within Q1DCOF.

1. Main Program; including subroutines READ and ABSA.

The basic structure is given in Figure A-1. MAIN calls READ to read input data and then ABSA to calculate the

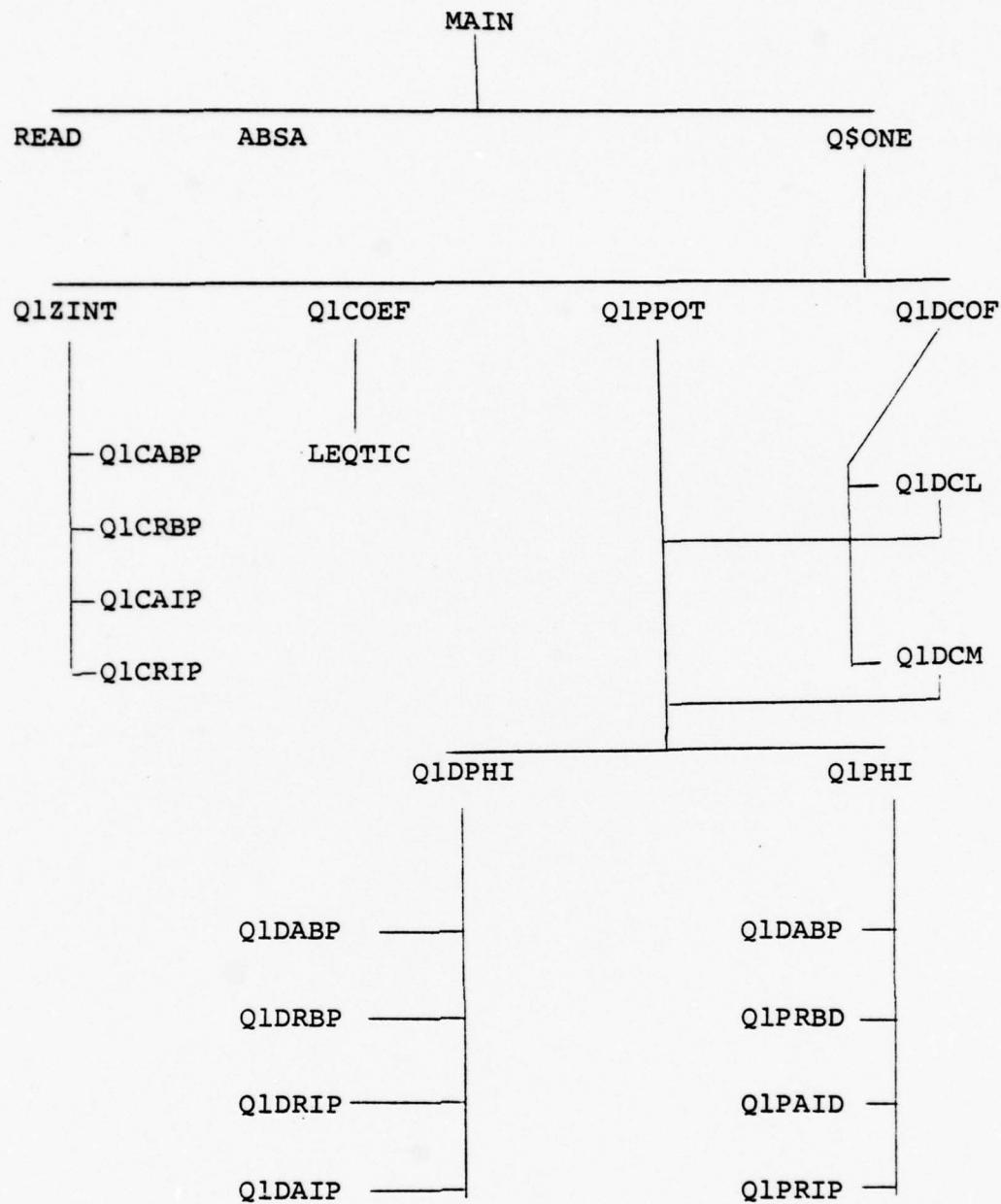


Figure A-1. Program Hierarchy

collocation points. The version of ABSA shown evenly spaces the collocation points across that portion of the blade subject to interference. ABSA may be easily replaced if different point spacing is desired, or if additional points are to be added for an overdetermined system and least squares approximation.

2. Q\$ONE This subroutine controls the actual potential calculation. It performs no calculation itself, but calls subordinate subroutines where the calculations are actually performed. The calling hierarchy is shown in Figure A-1.

3. Q1ZINT This subroutine calculates the linear equation system arising from the boundary conditions. Hierarchy is shown in Figure A-2.

The matrix output is carried through Q1INT. Q1ZINT calls the following subprograms

- a. Q1CRBP returns the value of  $\phi_z^0$
- b. Q1CABP returns the value of  $\phi_{z_1}^1$
- c. Q1CRIP returns the values of  $\psi_z^0$

$$\frac{\partial}{\partial z} \int_{-1+x_*}^{x-z} f_j(s) J_0 [\omega \sqrt{(x-s)^2 - z^2}] ds$$

where  $f_i$  is one of the set of i elementary functions,  $j=1,n$

- d. Q1CAIP returns the values of  $\psi_{z_1}^1$

**Q1ZINT**

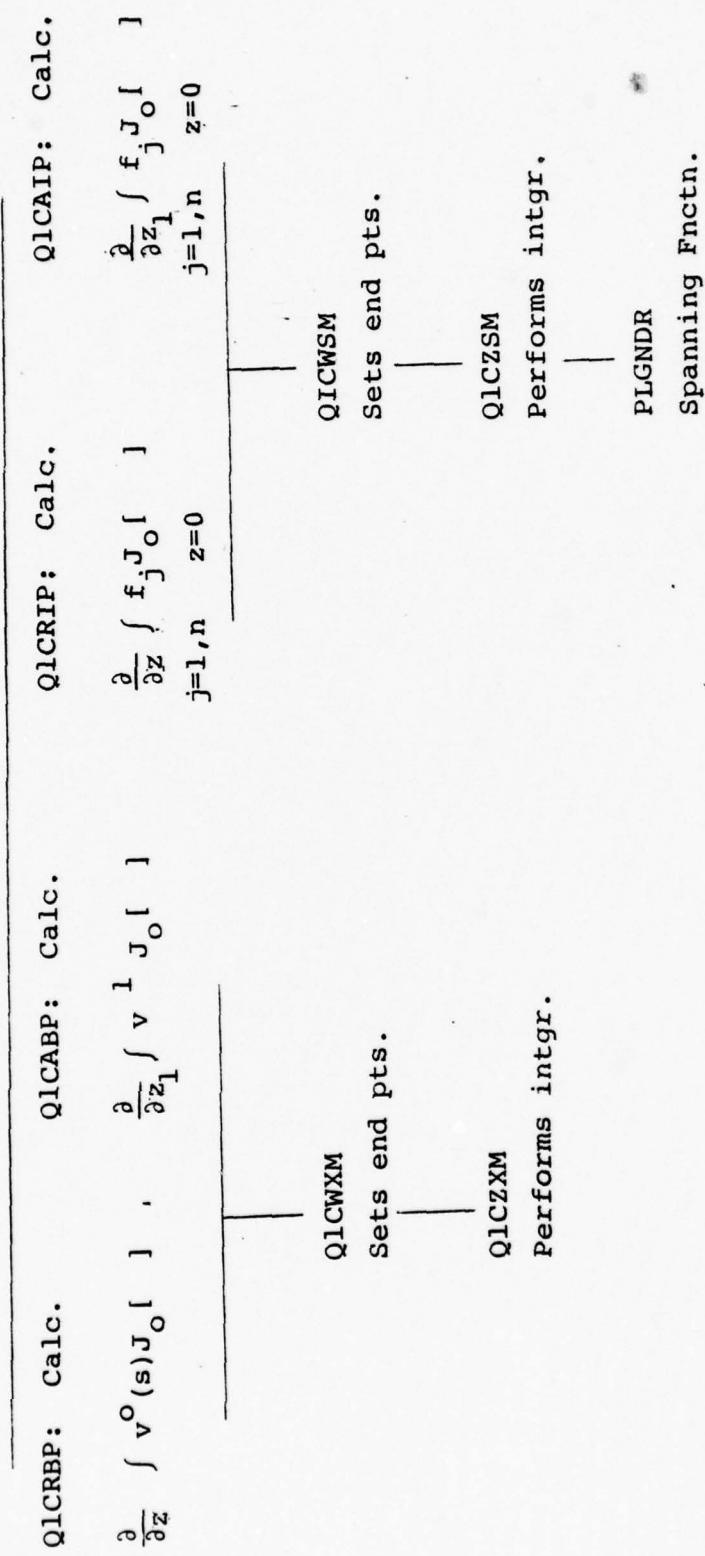


Figure A-2

$$\frac{\partial}{\partial z} \int_{z_1 - 1 + \text{OFFSET} + x_*}^{x+z_1} f_j(s) J_0 [\omega \sqrt{(x-s)^2 - y_1^2}] ds$$

where  $f_j$  is one of the set of elementary functions. OFFSET is a parameter included to facilitate program conversion to a staggered cascade

e. PLGNDR is the subprogram which returns  $f_j(x)$ , the elementary spanning function. No other routine contains explicit reference to the spanning function. This facilitates easy replacement of the spanning functions should this be desired.

Q1CRBP and Q1CABP in turn call Q1CWXM and Q1CZXM. Q1CWXM computes end-points and then calls Q1CZXM, a complex integration routine based on SIMP by Shampine and Allen [6]. Q1CAIP and Q1CRIP call Q1CWSM and Q1CZSM to perform the integration. Q1ZINT passes the constant matrix to Q1COEF in the array Q1INT and the right-hand-side vector in the array Q1COF.

4. Q1COEF This subroutine employs the IMSL routine LEQT1C to solve the linear system received from Q1ZINT. The resulting coefficients are Q1ABCF for the adjacent blade and Q1RBCF for the reference blade. LEQT2C, the high precision complex IMSL routine may be directly substituted for LEQT1C. Q1COEF may be rewritten to employ the generalized inverse required for least squares approximation

$$A = (x^T x)^{-1} x^T y \text{ where } x = Q1INT$$

$$y = Q1COF$$

after first performing the multiplication necessary to replace Q1INT and Q1COF with the proper matrix products in the call to LEQTIC.

5. Q1PPOT This subroutine calculates the potential,  $\phi$ , and  $\phi_x$ , at each collocation point along the reference blade but only if Q1PPOT receives a value of IPT > 0, requiring IPT  $\geq 2$  on input to the main program. If IPT  $\leq 0$ , then the subroutine is exited before any calculations are performed. This subroutine is most useful for debugging Q1ZINT and Q1COEF. Q1PPOT calls Q1PABP, Q1PRBP, Q1PRIP, Q1PAIP, Q1DABP, Q1DRBP, Q1DRIP, and Q1DAIP, all of which will be described in the next section.

6. Q1DCOF This subroutine calculates the dimensionless coefficients of lift and moment;  $C_{l\alpha}$ ,  $C_{m\alpha}$ . Its internal hierarchy is shown in Figure A-3.

- a. Q1DCL calculates the nondimensional complex coefficient  $C_{l\alpha}$
- b. Q1DCM calculates the nondimensional complex coefficient  $C_{m\alpha}$ .
- c. Q1PRBP and Q1PABP calculate the potentials due to the reference and adjacent blades respectively. Q1PWXM and Q1PZXM are called to perform the actual integration.

- d. Q1PRIP and Q1PAIP calculate the interference potentials along the reference and adjacent blades. Q1PZSM is called to perform the integration.
- e. Q1DRBP, Q1DABP, Q1DRIP, and Q1DAIP correspond exactly to subroutines above except that the value returned is the partial derivative of the potential with respect to X. Q1DWXM, Q1DZSM, and Q1DXSM perform the co-reponding integrals.

6. Program Listing The program listing shown below incorporates the Legendre functions as spanning functions. Listings for a subroutine employing Gorelov's spanning function and the linear approximation program follow.

```

      IMPLICIT REAL*8(A-H,C-P,R-Y),COMPLEX*16(Q,Z)
      DIMENSION X(13),Q2PT(13),Q2P(13)
      BLOCK ONE READ AND EDIT DATA
      CALL ERSET(208,256,-1,1)

      C WRITE(6,91C)
      C CALL IPT READ(DK,DR,DW,OFFSET,SIGMA,N,OFFSET) RHC,N,NF
      C 1 CALL IPT("I","5X","CR","ELCV",CASCACE,"PROGRAM")
      C 91C FCRMAT("0","0","OK","13X","CR","13X",DW,VALUES,TRANS)
      C 1 " " ,10X,"5CE12.53X","15X","NFCN"
      C 2 " " ,10X,"10X","10X","12X","3X","12X","4X",DW
      C 3 " " ,10X,"10X","10X","12X","3X","12X","4X",DW
      C CALL ABSA(N,CFFSET,X,RHO,DW)

      C BLOCK TWO CALLS CALCULATION ROUTINES FOR ZONE ONE AND ZONE TWO
      C CALL Q$CNE(CK,DR,EW,RHO,OFFSET,SIGMA,N,NF,IPT,X)
      C END
      C END
      C SUBROUTINE READ(CK,CR,DW,RHO,OFFSET,SIGMA,N,NF,IPT)
      C      SUBROUTINE READ*8(A-I,C-P,R-Y)
      C      READ(S,905) DK,DR,CW,RHP,OFFSET,SIGMA,N,NF,IPT
      C      IF(DK.LT.0.0D0) STOP
      C      IF(IPI.GT.C) WRITE(6,S10)
      C      GAMMA=1.4D0
      C      CFC=RHP*ESQRT((GAMMA + 1.0D0) * DW)
      C      RFC=DR
      C      IF(N.LE.13) GO TO 1
      C      N=13
      C      WRITE(6,935)
      C      1 WRITE(6,940) LE.1.9DC) GO TO 2
      C      OFFSET=CFFSET-1.0D0
      C      GO TO 1
      C      C CONTINUE
      C      IF(IPT.GT.0) WRITE(6,925) DK,DR,DW,SIGMA,CFFSET,RFC,N,NF,IPT
      C      RETURN(6,F10.4,312)
      C 905 FCRMAT("0","10X","STAGGERED SUPERSCALIC CASCADE PROGRAM")
      C 910 FCRMAT("0","10X","DK","13X","DR","13X","DW","SIGMA","10X","OFFSET")
      C 915 FCRMAT("0","10X","5CE12X","12X","15X","NFCN","12X","IPT")
      C 920 FCRMAT("0","10X","10X","10X","12X","12X","12X","I2")
      C 925 FCRMAT("0","10X","10X","10X","12X","12X","12X","I2")
      C 930 FCRMAT("0","EX","ORIGINAL TOTAL LARGE (.GT.1.3) - RESET TON = 13")
      C 935 FCRMAT("0","EX","OFFSET TOO LARGE (.GT.1.9), RESET AS CFFSET = 1")
      C 940 FCRMAT("0","EX","OFFSET - 1.CDD0")
      C END

```



```

XSTN = XSTN - DR
CEXP = CDEXP(DCNPLX(0.0D0, DLANDA*XSTN))
CALL Q1CAIP(DK, DR, DW, RHO, OFFSET, SIGMA, XSTN, QINTRP, N, IOT)
DC 20 J=1,N
      J=N+
J1 = J-1
Q1INT(I,N,J) = PLGNDR(XSTN,DR,J1) * CEXP
Q1INT(I,N,JN) = QINTAP(J)
2 C Q1NTINUE(I,N) = PLGENR(XSTN-OFFSET,DR,J1) * CEXP
Q1COF(I,N) = -Q1CREF(DK,DR,DW,RHO,XSTN,IOT)
Q1COF(I) = -Q1CABP(DK,DR,DW,RHO,OFFSET,SIGMA,XSTN,IOT)
50 CCNTINUE
      RETURN
END
COMPLEX FUNCTION Q1CRBP*16(DK,DR,[,RHO,XSTN,IPT])
COMPLICIT REAL * 8 (A-F,C,P,R-Y), COMPLEX * 16 (Q,Z)
DIMENSION QINP(2)
IF(IPT<GT) WRITE(6,990) DK,DF,DW,RHO,XSTN,IPT
990 FORMAT(0:X,Q1CREP ENTERED WITH: /,1X,DR: /,1X,DW:,11X,RHO:,10X,XSTN:,SX,' IPT'
12 IF(XSTN.LE.RHO-1.0E-20) GOTO 20
IICK = IPT - 1
QCK = DCMPLEX(0.0D0,DK)
CALL Q1CWXM(DK,DR,DW,FHO,XSTN,QINP,1)
Q1CRBP = -Q1K*Q1NP(2) - QINP(1)
1F(IPT LE.0) RETURN
GCTO 30
20 Q1CRBP = DCFLX(0.0D0)
1F(IPT LE.0) RETURN
30 WRITE(6,995) Q1CREP
995 FURMAT(0.0,10X,Q1CREP = ' , E14.7, *, E14.7)
      RETURN
END
COMPLEX FUNCTION Q1CABP*16(DK,DR,DW,RHO,OFFSET,SIGMA,XSTN,IPT)
COMPLICIT REAL * 8 (A-F,O,P,R-Y), COMPLEX * 16 (G,Z)
DIMENSION QINP(2)
IF(IPT<GT) WRITE(6,990) DK,DR,DW,RHO,OFFSET,SIGMA,XSTN,IPT
990 FORMAT(0:X,Q1CABP ENTERED WITH: /,1X,DR: /,1X,DW:,11X,RHO:,10X,OFFSET:,7X,
12 *SIGMA:,8X,XSTN-OFFSET
      XASTN=XSTN - OFFSET
      IF(XASTN.LE.RHO-1.CD0) GOTO 20

```

```

IF(XASTN.GT.2.0D0) GOTO 20
10T = IPT - 1
DK = DCMPLEX(0.0D0,0D0)
QCONST = CDEXP(DCMPLEX(0.0D0,SIGMA))
CALL Q1CXM(DK,DR,DWRHO,XASTN,1,0T)
CALLABP = -(DK*QINP(2)+QINP(1)) * QCONST
IF(IPT.LE.0) RETURN
10 TO 30
20 CALLABP = DCMPLEX(0.0DC,0.0DD0)
30 WRITE(6,995) Q1CABP
995 FORMAT(1X,10X,'Q1CABP = ',E14.7,I14.7)
RETURN
20
SUBROUTINE Q1CWM(XM(DK,DR,DW,RHO,XSTN,QINP,IPT)
C
C N IS THE MAXIMUM DEGREE OF THE TERM U IN THE INTEGRALS
C
IMPLICIT REAL*8 (A-H,O,P,R-Y), COMPLEX*16 (Z,Q)
CIMENSION QINP(2)
1 IF(IPT.GT.0) WRITE(6,990) DK,DR,RHC,XSTM(IP1)
990 FORMAT('0',10X,'Q1CWM ENTERED WITH ARGUMENTS:',IP1,
1 '10X,'10X,'DK','16X,'ER','16X,'RHO,',15X,'XSTM',
2 '10X,'4(13.6,5X),12,5X,13)
10T = IPT - 1
B = XSTN - RHO - 1.0D-10
A = -1.0D0
CALL Q1CZX(DK,DR,Eh,RHO,XSTN,A,B,1,QANS,10T)
CALLP(1)=QANS
CALL CZX(DK,DR,Eh,RHO,XSTN,A,B,2,QANS,10T)
CALLP(2)=QANS
1 IF(IPT.LE.0) RETURN
10T = IPT - 1
995 FORMAT(1X,10X,'Q1CWM RESULTS:',/,1,
1 DC1001=1,1=1,2
1001 DC1001,I=1,3,5XE12.6,/,E12.6)
1001 RETURN
1001 WRITE(6,996) I, QINP(1)
1001
SUBROUTINE Q1CZX(DK,DR,DW,RHC,XSTM,A,B,J,QANS,IPT)
IMPLICIT REAL*8(A-E,G,H,M,O,P,R-Y),COMPLEX*16(F1,F2)
DIMENSION FV(5),LORR(30),F1(30),F2T(30),F3T(30),QPSUM(30),
1 AREST(30),QEST(30),EPST(30),CAT(30)
1 F(X) = (X**J)*CDEXP(QEXP(X))
1 F(X) = (OMEGA*RHO/(DSQRT((XSTM-X)*(XSTM-X))-Y))**J
2 MMEJS1(OMEGA*DSQRT((XSTM-X)*(XSTM-X))-Y),IER)
2 GAMMA = 1.4C0

```

```

DN2 = ( GAMMA + 1.00 ) * DW
CNEGA = DSQRT( DK*EK* ( 1.00 - DM2 ) / ( DM2 * DM2 ) )
YY = RHO * RHC
CLANDA = DK / DM2
CEXP = DCMLX( 0.00, CLANDA )
JACC = J - 1.00 - 6
ACC = 1.00 - 13
IF ( IPT * GT * C ) WRITE( 6, 990 ) DK, DR, D, RHO, XSTN, A, B, J, IFT
FCRMAT( 1, 15X, Q1CZ, X, ENT, RED, W, AR, GUMEN, TS, J, IFT
1, '15X, '13X, 'CR, '13X, 'DW, '10X, 'RH, '12X, 'XSTN, '11X,
2, 'A, '14X, 'B, '14X, 'J, '2X, '1PT, '12, '2X, '13,
3, '15X, '7, 'E, '14, '7, 'J, '12, '2X, '13,
EFFOURU = 4.0*U
IFLAG = 1
EPS = ACC
CERROR = DCMLX( 0.00, 0.00 )
LVLFR( LVL ) = 1
LVLFSUM( LVL ) = 0.0
ALPHA = A
CA = B - A
AREA = 0.0
AREST = 0.0
FV( 1 ) = F( ALPHA )
FV( 3 ) = F( ALPHA + 0.5*DA )
FV( 5 ) = F( ALPHA + DA )
KCNT = 3
WT = DA / 6.0
WEST = WT * ( FV( 1 ) + 4.0*FV( 3 ) + FV( 5 ) )
CX = 0.5*DA
FV( 2 ) = F( ALPHA + 0.5*DX )
FV( 4 ) = F( ALPHA + 1.5*DX )
KCNT = KOUNT + 2
WESTL = WT * ( FV( 1 ) + 4.0*FV( 2 ) + FV( 3 ) )
WESTR = WT * ( FV( 3 ) + 4.0*FV( 4 ) + FV( 5 ) )
QSUM = QESTL + QESTR
ARESTL = WT * ( CDABS( FV( 1 ) + CDABS( FV( 2 ) ) + CDABS( FV( 3 ) ) )
ARESTR = WT * ( CDABS( FV( 3 ) + CDABS( FV( 4 ) ) + CDABS( FV( 5 ) ) )
AREFA = AREA + (( ARESTL + ARESTR ) - ARREST )
QDIFF = QSLM - ( CDABS( QDIFF ) * LE * EPS*DABS( AREA ) ) GO TO 5
1F( CDABS( DX ) * LE * EFCURU*DABS( ALPHA ) ) GO TO 5
1F( LVL * GE * 30 ) GO TO 5
1F( LVL * GE * 30 ) GO TO 6
LCRR( LVL ) = 0

```

```
SUBROUTINE CCAIP(CK,CR,DW,RTO,GFST,SGMA,XSTN,CCAIP,N,IPT)
IMPLICIT REAL*8 (A-H,C,P,R-Y), COMPLEX*16 (C,Z)
```

```

      DIMENSION CINP(13) QCAIP(13) CK, DR, DW, RHO, OFST, SGMA, XSTN, IPT
      IF(IPT .EQ. 0) WRITE(6,990) CK, DR, DW, RHO, OFST, SGMA, XSTN, IPT
      990 FORMAT(0,10X,10X,10X,10X,10X,10X,10X,10X,10X,10X,10X,10X)
      1   Q1CAIP ENTERED WITH 1/1X, 1/RHC, 1CX, 1C, 1/2, 1/2, 1/2, 1/2, 1/2, 1/2, 1/2, 1/2, 1/2, 1/2
      2   XASSTN = XSTN - OFST
      2   IF(XASSTN .LE. DR+RHO-1.000) GOTC 20
      1CT = IPT - 1
      1CT = CDEXP( DCMLX(0.000,SGMA) )
      CALL Q1CWSM(DK,DR,DW,RHO,XSTN,N,CINP,1OT)
      QCAIP(I) = CINP(I)
      CNTINUE = 1
      10 IF(IPT .NE. 0) RETURN
      10 GCTO 30
      20 ZERC = DCMLX(0.000,0.000)
      20 DC25 I = 1
      20 QCAIP(I) = ZERO
      25 CNTINUE
      25 CMLX(I) = 0 RETURN
      30 WRITE(6,995)
      30 995 FORMAT(10,10X,'Q1CAIP RESULTS: J      QCAIF(J)')
      30 995 QCAIF(J)
      30 995 CCRIP(1,10X,J=1,N)
      30 995 CCRIP(1,6,996) J=QCAIP(J)
      30 995 CCRIP(1,6,996) J=26X,12,3X,E14.7,1,1,E14.7)
      30 995 CNTINUE
      30 995 RETURN
      END
      SUBROUTINE Q1CRIP (CK, DR, DW, RHC, CFFSET, XSTN, CCRIP, N, IPT)
      DIMENSION CCRIP(13) REAL*8(A-H,O,P,R-Y), COMPLEX *16(C,2)
      IF(IPT .EQ. 0) WRITE(6,990) CK, DR, DW, RHO, XSTN, IPT
      990 FORMAT(10X,10X,10X,10X,10X,10X,10X,10X,10X,10X,10X,10X)
      1   CCRIP ENTERED WITH 1/1X, 1/RHC, 1CX, 1C, 1/2, 1/2, 1/2, 1/2, 1/2, 1/2, 1/2, 1/2, 1/2, 1/2
      2   XSTN = XSTN - OFSET
      2   IF(XSTN .LE. DR+RHC+CFFSET-1.000) GOTC 20
      1CT = IPT - 1
      1CT = Q1CWSM(DK,DR,DW,RHO,XSTN,N,CCRIP,IPT)
      1CT = IPT - 1
      1CT = CMLX(I) RETURN
      1CT = CMLX(I)
      20 ZERC = DCMLX(0.000,0.000)
      20 DC25 I = 1
      20 QCRIP(I) = ZERO
      25 CCRIP(I) = ZERO
      25 CCRIP(I) = ZERO
      30 WRITE(6,995)
      30 995 FORMAT(10,10X,'Q1CRIP RESULTS: J      QCRIF(J)')
      30 995 QCRIF(J)

```

```

DC 40 J = 11^ J,QCRIP(J)
996 FORMAT(1.26X,12,3X,E14.7,' ',E14.7)
940 CONTINUE
      RETURN
ENDC
      SUBROUTINE QICWSM (CK, DR, DW, RHO, XSTN, N, QINP, IFT)
      C
      C   A IS THE MAXIMUM DEGREE OF THE TERM U IN THE INTEGRALS
      C
      C   IMPLICIT REAL * 8 (A-H,O,P,R-Y), COMPLEX * 16 (Z,C)
      C
      C   DIMENSION QINP(13)
      C   IF(IPT.GT.0) WRITE(6,990) DK,DR,FHC,XSTN,N,IPT
      990 FORMAT(1.0X,Q1CH,16X,DR,16X,RHO,.15X,XSTN,14X,N!,6X,'IPT',/,,
      1   ,10X,1.0X,4(E13.6,5X),12.5X,13)
      2   IPT = IPT - 1
      B = XSTN - RHO - 1.0C-10
      A = DR-1.0D0
      DC 30 J=1N
      CALL QICZSM(DK,DR,DW,RHO,XSTN,A,B,J,QANS,IOT)
      CINP(J)=QANS
      C
      C   CONTINUE
      IF(IPT.LE.0) RETURN
      994 WRITE(6,995)
      995 FORMAT(1.10X,QICWSM,RESULTS:',/',,
      996 FORMAT(1.11X,'J 1,6X,QINP(I),F3.1
      DC 1001 I = 1,1
      1001 WRITE(6,996) I, CINP(I),E12.6, ',',E12.6)
      1001 RETURN
ENDC
      SUBROUTINE QICZSM(DK,DR,DW,RHO,XSTN,A,B,J,QANS,IPT)
      C
      C   IMPLICIT REAL*8(A-E,H,M,O,P,R-Y),COMPLEX*16(F,G,Z)
      C
      C   DIMENSION FV(5),LCRR(20),F1T(30),F2T(30),F3T(30),
      1   AREST(30),QESTT(30),QPSUM(30)
      1   F(X) = PLGCR(X,DR,J)*CDEXP(QEXP(X))
      1   *(CMSEG*RH0/(TSQRT((XSTN-X)*(XSTN-X)-YY)))*
      2   MMBSJ1(OMEGA*DSQRT((XSTN-X)*(XSTN-X)-YY)),IER)
      2   GAMMA = 1.4C0
      EN2 = (GAMMA + 1.0D0) * DW
      YY = RH0 * RHO
      CLAMDA = DK/DM2
      QEXP = DMPLX(0.0D0, LA MDA )
      J1 = J-1
      ACC = 1.0D-6
      U = 9.0D-13

```

```

1 IF (IIFT(6,0) WRITE(6,990) DKDR, XSTN, A, E, J, IPT
2   FORMAT(15X, Q1CZS, ENTERED WITH ARGUMENTS: A, E, J,
3   15X, CR, 13X, CR, 14X, J, 2, IPI, 13,
4   A, 15X, 7(E 14, 7, , 1, 12, 2X, 13),
5   EFCURU = 4. C*U
6   IFLAG = 1
7   EERROR = ACC
8   EPLX(0.0D0, 0.0D0)
990  LLRR(LLVLL) = 1
10  CFSUM(LLVLL) = 0.0
11  ALPHAA = A - A
12  AREA = C*0
13  AREEST = 0.0
14  FV(1) = F(ALPHA)
15  FV(3) = F(ALPHA + 0.5*DA)
16  FV(5) = F(ALPHA + DA)
17  KOUNT = 3
18  KDIFF = DA/6.0
19  KEST = WT*(FV(1) + 4.0*FV(3) + FV(5))
20  DX = 0.5*DA
21  FV(2) = F(ALPHA + 0.5*DX)
22  FV(4) = F(ALPHA + 1.5*DX)
23  KOUNT = KOUNT + 2
24  KESTL = DX/6.0
25  KESTR = WT*(FV(1) + 4.0*FV(3) + FV(5))
26  QESTW = QESTL + QESTR
27  QESTWL = WT*(CDABS(FV(1)) + CDABS(FV(3)) + CDABS(FV(5)))
28  QESTWR = WT*(CDABS(FV(1)) + CDABS(FV(3)) + CDABS(FV(5)))
29  AREA = AREA + (AREESTL + ARESTR) - AREEST
30  QDIFF = QEST - QSUM
31  IF (CDABS(QDIFF) .LE. EP * DABS(AREA)) GO TO 5
32  IF (LVL * GE * 3C) GO TO 5
33  IF (KOUNT .GE. 2000) CC TC 6
34  LVL = LVL + 1
35  LLRR(LLVLL) = 0
36  F1T(LLVLL) = FV(3)
37  F2T(LLVLL) = FV(4)
38  F3T(LLVLL) = FV(5)
39  DA = DX
40  AREEST(LVL) = ARESTR
41  AREESTT(LVL) = ARESTR
42  QEST = QESTL

```

```

QESTT(LVL) = QESTR
EPS = EPS/1.4
FV(5) = FV(3)
FV(3) = FV(2)

3  GERROR = QERROR + CDIFF/15.0
IF(LLCRR(LVL)EQ.0) GO TO 4
QSUM = QPSUM(LVL) + QSUM
LVL = LVL - 1
IF(LVL.GT.1) GO TO 3 * CDEXP(QEXP*B) - QSUM
IF(IPT.GT.C) GO TO 11 RETURN
IF(IER.EQ.129) GO TO 11
IF(IFLAG.EQ.1) WRITE(6,995) DKDRDW,RHCBXSTNALE,J,IPT
995 1 FORMAT(15X,"RESULTS: QANS=ERR",E14.7,".",E14.7,"/",15X,
      1 "IFLAG",5X,"IER:",E14.7,".",E14.7,".",E14.7)
2  RETURN

4  QPSUM(LVL) = QSUM
LCRR(LVL) = 1
ALPHA = ALPH + DA
DA = DAT(LVL)
FV(1) = F1T(LVL)
FV(3) = F3T(LVL)
FV(5) = F5T(LVL)
AREST = AREST(LVL)
CEPS = QEST(LVL)
CEPS = EPS
5  IFLAG = 2
6  IFLAG = 3
CEAD READ FUNCTION PLGNDR*B(X,DR,N)
IF(PLICIT REAL*B(A-T,0-Z) EQ.0) GOTO 10C
X2 = X*X
GCTO(101,1C2103,104,105,106,107,108,109,110,111,112), N
100 PLCNDR = 1.0D0
101 PLCNDR = X
102 PLCNDR = ((3.0DC)*X2-1.0D0)/2.0D0
RETURN

```



```

      IA = 26
      IJCB=0
      CALL LEQT1C(Q1CINT,N2,IA,Q1CCF,M,IE,IJOE,ZWA,IER)
      IFF(IER.EQ.0) GO TO 30
      IF(IER.EQ.129) GO TO 10
      WRITE(6,93)
      FORMAT('0.10X','Q1CCF - ITERATIVE IMPROVEMENT FAILED, MATRIX TOO
      93 1 FULL-CONDITIONED. USE RESULTS WITH CAUTION.')
      GC TO 30
      WRITE(6,95)
      95 FORMAT('0.10X','Q1CCF - MATRIX ALGORITHMICALLY SINGULAR. CCEFFIC
      1ENTS SET TO ZERO')
      ZERO = DCMPLX(0.0DC,0.0DO)
      DC 20 I = 1,N
      Q1COF(I) = ZERO
      Q1COF(I2) = ZERO
      CONTINUE
      20
      Q1ABC(I) = Q1COF(IN)
      Q1RBCF(I) = Q1COF(IN)
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      Q1CF(IPT*1E+0) RETURN
      DC 40 I = 1,N
      WRITE(6,99) I, IM1, Q1RBCF(I), Q1ABC(I)
      CONTINUE
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      CCNTINUE
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      FORMAT('0.5X12.10X,2(E14.7,5XE14.10X)', 'COMPLEX POWER SERIES COEFFIC
      SC FORMATS')
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      FORMAT('0.5X','INDEX',7X,'DEG FCFLY',4X,'REFERENCE FFLADE TERMS Q1C
      3CUT(INDEX)',7X,'ADJACENT BLADE TERMS Q1COF(2*INDEX),')
      RETURN
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QDRIPP = Q1DRIP(DK,DR,DW,RHO,OFFSET,XSTN,C1RBCF,N,IOT)
QDAIPO = Q1DAIP(DK,DR,DW,RHO,OFFSET,SIGMA,XSTN,Q1ABC,F,N,IOT)
QDAIPP = Q1DAIP(DK,DR,DW,RHO,OFFSET,SIGMA,XSTN,Q1AECF,N,IOT)
QCR = QDRO + QDAF + QCRIPQ + QCA1PP
QDA = QDRO + QDAF + QCRIPP + QDAIFC
QLPHI(I) = Q1CFPI(DK,DR,DW,RHO,OFFSET,SIGMA,N,C1ABC,F,C1RBCF,
1 WRITE(XSTN,10T)
1 WRITE(6,95) 1,XSTN
1 WRITE(6,91) QCR,QCRO,QDAP,QDAO,QCRIPP,QCAIPO
1 WRITE(6,92) QCA,QCRP,QDAF,QCRIPP,QCAIPO
1 QDRR(1) = QDR
1 QCAAC(I) = QDA
1 FFORMAT(0,0,X-STATION NUMBER,112,0,XSTN)=1F6.4,/,/
1 1 'BL TCTAL D(POT)/DX',14X,'REF BL D(POT)/DX',1F8X,
1 2 'ADJ BL D(FCT)/DX','REF BL INT D(POT)/DX',1F5X,
1 3 'ADJ BL INT D(POT)/DX,'
10 CONTINUE(6,994)
95 1 WRITE(6,994) '10X' 'SUMMARY LISTING'
1 0 '10X' 'XSTN' 'SINGLE BLADE TOTAL POTENTIAL';EX,
1 1 'REF BLADE POTENTIAL',15X,'ADJ BLADE POTENTIAL';EX,
1 DC 20 I = 1,N
XSTN = X(1)
94 1 WRITE(6,94) XSTN,CPHI(I),QRR(I);QAA(I)
1 0 '10X',15X,E14.7,;E14.7,)
20 CONTINUE(6,996)
956 1 WRITE(6,996) '10X' 'XSTN' 'SINGLE ELADE TOTAL DIRECT/DX',6X,
1 0 'REF ELADE D(POT)/DX',15X,'ADJ ELADE D(POT)/DX,'
1 DC 50 I = 1,N
XSTN = X(1)
94 1 WRITE(6,94) XSTN,CPHI(I),QRR(I),QAA(I)
20 CONTINUE
50 RETURN
ENDCUTINE Q1DCOF(CK,DR,DW,RHC,CFFSET,SIGMA,N,Q1ABC,F,Q1RBCF,IFT)
INPLICIT REAL*8 (A-H,O,P,R-Y) COMPLEX*16 (Q,Z)
DIMENSION Q1ABCF(13),Q1RBCF(13)
DIMPNT(G,O)WRITE(6,990) DKDR,E'RHO' OFFSET SIGMA,N,IPT
990 FCFORMAT(1:10X,Q1DCOF - CALCULATION OF COMPLEX DIMENSIONLESS AERO
1 1 DYNAMIC COEFFICIENTS',13X,'DR',13X,'IP',13X,'DW',13X,'RHO',12X,'OFFSET',SX,'SIGMA'
3 0 '10X',15X,'CK',13X,'IP',13X,'DW',10X,'12X,12,2X,13'
4 56 '1E12.5'3X),12,2X,13)
50 ICT = 1 PT - 3
995 QCCL = C1CCM(DK,DR,DW,RHO,OFFSET,SIGMA,N,Q1ABC,F,C1RBCF,IOT)
QCCM = Q1DCM(DK,DR,DW,RHC,OFFSET,SIGMA,N,Q1ABC,F,Q1RBCF,IOT)
CAMA = 1.4E0

```

```

TAL = ((2.0DC*DSQRT((GAMMA + 1.0DO)*DW))/DR
WRITE (6.90) DK TAU, DW2N, SIGMA2CL, GDCM
FCRMA(6.0, EX, DK = , F6.3, , TAU = , F7.4, , DW = , F6.3, , CM = , , N =
      12, SIGMA = , F6.3, , CL = , F9.4, , F9.4)
      2 RETURN
END

COMPLEX FUNCN Q1DCL*16(DK,DR,DW,RHO,OFFSET,SIGMA,N,Q1ABC,F,G1RBCG)
1 IMPLICIT REAL*8(A-E,G1B0FIR-V), COMPLEX*16(F,Q,Z)
2 DIMENSION Q1ABC(F(13),Q1RBCF(13))
3 DIMENSION FV(5) F1T(60) F2T(60), F3T(60), QESIT(60), QPSLM(60)
4 DIMENSION DAT(60), ARSTT(60)
5 DIMENSION LCR(60)
6 FIX1 = (C1DEFI(DK,DR,DW,RHO,OFFSET,SIGMA,N,Q1ABC,F,G1RBCF,X,ICT))
7 IF (IPT*GT*CN) WRITE (6,990) DK,ER,DW,RHO,CFFSET,SIGNA,N,IP1
8 FCRMAT(0,10X,10X,Q1DCL ENTERED WITH: /, RHC, CFFSET, SIGNA,N,IP1
9 4 10X,13X,DK,13X,DR,13X,DK,13X,IP1,12X,13X,IP1,12X,13X,IP1,12X,13X,IP1
10 A = -1.0DO
11 E = 1.0DO
12 ICT = IPT - 1
13 GAMMA = 1.4DO
14 CLAMDA = CK/((GAMMA + 1.0DO) * Dk)
15 QALPHA = DCMPLX(0.0DO,ELAMDA-Dk)
16 = 9.0E-13
17 ACC = 1.0D-5
18 EFCURU = 4.0*u
19 IFLAG = 1
20 IFS = ACC
21 QERROR = DCMPLX(0.0DO,0.000)
22 QYL = 1
23 LCCR(LVYL) = 1
24 QPSUN(LVYL) = 0.0
25 ALPHA = A
26 DA = B - A
27 AREA = 0.0
28 AREST = 0.0
29 FV(1) = F(ALPHA)
30 FV(3) = F(ALPHA + 0.5*DA)
31 FV(5) = F(ALPHA + DA)
32 KCLNT = 3
33 KTEST = DA/9.0
34 CTEST = HT*(FV(1) + 4.0*FV(3) + FV(5))
35 DA = 0.5*DA
36 F(2) = F(ALPHA + 0.5*DX)
37 DR = 3375
38 DR = 3380
39 DR = 3395
40 DR = 3400
41 DR = 3415
42 DR = 3420
43 DR = 3430
44 DR = 3435
45 DR = 3440
46 DR = 3445
47 DR = 3450
48 DR = 3455
49 DR = 3460
50 DR = 3465
51 DR = 3470
52 DR = 3475
53 DR = 3480
54 DR = 3490
55 DR = 3495
56 DR = 3500
57 DR = 3505
58 DR = 3510
59 DR = 3515
60 DR = 3520
61 DR = 3525
62 DR = 3530
63 DR = 3535
64 DR = 3540
65 DR = 3545
66 DR = 3550
67 DR = 3555
68 DR = 3560
69 DR = 3565
70 DR = 3570
71 DR = 3575
72 DR = 3580
73 DR = 3585
74 DR = 3590
75 DR = 3595
76 DR = 3600

```



AD-A063 083

NAVAL POSTGRADUATE SCHOOL MONTEREY CALIF  
THEORETICAL ANALYSIS OF TRANSONIC FLOW PAST UNSTAGGERED OSCILLA--ETC(U)  
SEP 78 P C OLSEN

F/G 20/4

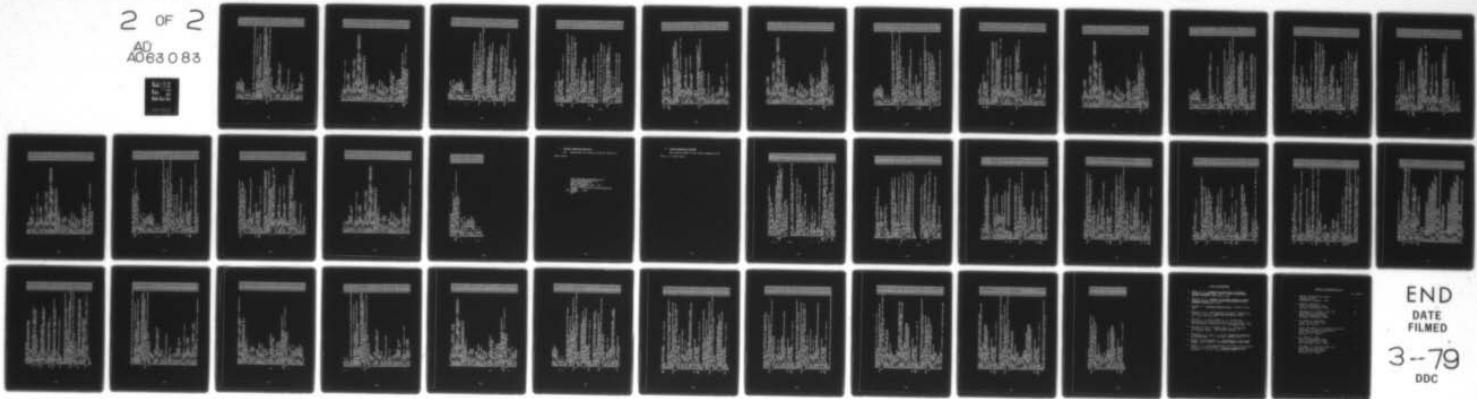
UNCLASSIFIED

2 OF 2

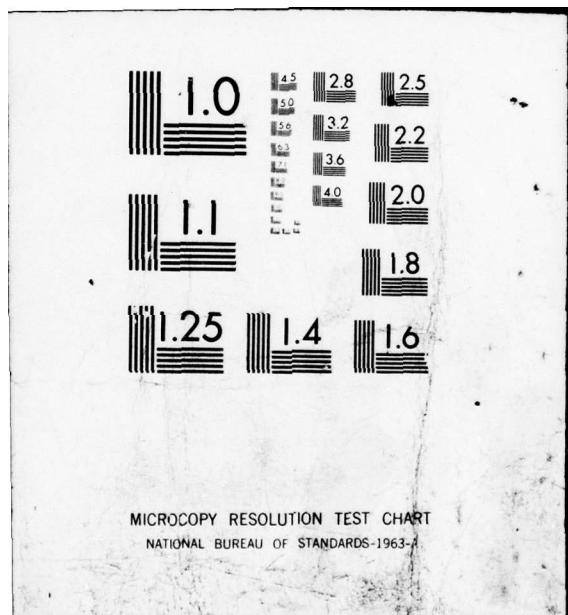
AD  
A063 083

FF

NL



END  
DATE  
FILMED  
3-79  
DDC





```

KCINT = 3
WT = DA / 6.0
CEST = WT*(FV(1) + 4.0*FV(3) + FV(5))
DX = 0.5*DA
FX(1) = F(ALPHA + C5*DX)
FX(2) = F(ALPHA + C5*DX)
FX(3) = F(ALPHA + C5*DX)
FX(4) = F(ALPHA + C5*DX)
KOUNT =

```

1

```

ARESTL = WT*(CDABS(FV(1)) + CDAES(4.0*FV(2)) + CDABS(FV(3)))
ARESTR = WT*(CDABS(FV(3)) + CDABS(4.0*FV(4)) + CDAES(FV(5)))
AREA = AREA + ((AFESTL + AR ESTR) - AREST)
CDIFF = QUEST - Q SUP
IF(CDAB(SQDIFF) <= EPS*DABS(AREA)) GO TO 2
IF(CDAB(SQDIFF) >= EPS*DABS(ALPHA)) GO TO 5
IF(LVL.GE.60) GO TO 6
IF(KOUNT.GE.14000) GO TO 6

```

```

L>L = L>L
L>R = L>L
L>T = L>L
F>T = F>F
F>3 = F>F

```

```

DATA = DX
DATE(LVL) = CX
AREST(L) = AREL
ARESTT(LVL) = ARESTR
ARESTL(LVL) = QESTR
CESTT(LVL) = QESTA
CEST(LVL) = QESTA
CEPST(LVL) = EPS/1.4
EFST(LVL) = EPS
EFST(5) = FV(2)
FV(3) = FV(2)

```

**GERROR** = CEFDRR + QCIFF15.0  
**GERROR** = CEFDRR + GO TO 4  
**GCFILM** = QSUM(LVL) + QSUM

115

23





```

995  WRITE(6,995)
      FORMAT(6,995)
      1  FORMAT(1X,'J',1X,'RESULTS : ',/,
      2  '001 1 = 1 2 001 1 13 5 X E12.6, 0, E12.61
      3  FORMAT(1X,'MM ESJO (6,996) 1, (INP(1),
      4  RETURN
END
SUBROUTINE QIPZXM (DK,DR,DW,RHO,XSTN,A,B,J,CMPLEX,IPT)
 1  SUBPLICION REAL*8 (A-E,F,G,I,O,P,R-V,T,U)
 2  DIMENSION FY(5),LCFR(60),F1(60),F2T(60),F3(60),DAT(60),
 3  1  REST(60),QEST(60),EPS(60),FSUM(60)
 4  F(X) = (X**J)*CDEXP(QEXP(X))
 5  1  * MM ESJO ((OMEGA * DSQRT ((XSTN-X)*(XSTN-X) - YY)), IER)
 6  GAPMA = 1.4D0
 7  YY = RHO * RHO
 8  D12 = (GAMMA + 1.0E0) * DW
 9  CNEGA = DSQRT((DK*CK*(1.0D0-DM2)/(CM2*DM2)))
10  CLANDA = DK/CM2
11  C2EXP = CDEXP(DCMPLX(0.0D0,CLAMDA))
12  Q1 = DCMPLX(0.0D0,-DLANDA*XSTN))
13  J = 1
14  ACC = 9.00-13
15  IF (IPT .EQ. 0) WRITE(6,990) DK,DR,DW,RHO,XSTN,A,B,J,1FT
16  1  FORMAT(1X,'0',15X,'1',13X,'CR',13X,'DW',10X,'RHO',12X,'XSTN',11X,
17  2  'A',14X,'B',14X,'J',2,X,'IPT',1,
18  3  'EFCURU = 4.0*U
19  1  ERROR = ACC
20  1  CERROR = 1
21  LCLR(LVL) = 1
22  CFSLM(LVL) = 0.0
23  AREA = B-A
24  AREAST = 0.0
25  ARREST = Q(0)
26  FY(1) = F(ALPHA)
27  FY(5) = F(ALPHA + Q(5)*DA)
28  KCOUNT = 3
29  KEST = DA/6*0
30  CX = 0.5*DA
31  CEST = WT*(FV(1) + 4.0*FV(3) + FV(5))

```

$FV(2) = F(\text{ALPHA} + 0.5*DX)$   
 $FV(4) = F(\text{ALPHA} + 1.5*DX)$   
 $KOUNT = KCNT + 2$   
 $DX/6^0 = WT*(FV(3) + 4.0*FV(2) + FV(3))$   
 $QESTL = WT*(FV(1) + 4.0*FV(4) + FV(5))$   
 $QESTR = CE\text{STL} + QESTR$   
 $ARESSTL = WT*(CDABS(FV(1)) + CDAES(4.0*FV(2)) + CDABS(FV(3)) + CCABS(FV(5)))$   
 $ARESTR = WT*(CDABS(FV(3)) + CDAES(4.0*FV(4)) + CCABS(FV(5)))$   
 $AREAST = AREAST - (QSUN - QESTL + ARESTR) - \text{AREST}$   
 $AREAFF = QESTL - QESTR$   
 $ACIICABSS(X) = LE*EPS*DABS(AREA)) GO TO 2$   
 $ACIICDABS(X) = LE*EFOUR*DABS(ALPHA)) GO TO 5$   
 $IF(COUNT*CE + 2000) GO TO 5$   
 $LVR(LVL) = LVL + 0$   
 $LFT(LVL) = FV(3)$   
 $LFT(LVL) = FV(4)$   
 $LFT(LVL) = FV(5)$   
 $DATA(LVL) = DX$   
 $AREST = AREST$   
 $ARESTT(LVL) = ARESTR$   
 $CESTT = CESTL$   
 $CESTT(LVL) = QESTR$   
 $CESTT(LVL) = EPS/1.4$   
 $CESTT(LVL) = FV(3)$   
 $CESTT(LVL) = FV(2)$   
 $QERROR = QERROR + CDIFF/15.0$   
 $IF(CORR(LVL) EQ 0) GO TO 4$   
 $QSUM(LVL) = QSUM(LVL) + QSUM$   
 $LVL = LVL - 1$   
 $IF(LVL = 1) GO TO 3$   
 $QANS = QSUM * Q2EXP / DM$   
 $IF(IPTAG = EC.1) RETURN$   
 $IF(IFIT(6,990) QANS IFLAG = EC.1) RETURN$   
 $IF(IFIT(6,995) QANS RESULTS: QANS = E14.7, J, IPT$   
 $FCRIMAT(6,15X, 15X, 1FLAG, 1ER, QERR(14.7, 1, E14.7),$   
 $15X, 13, 5, E14.7, E14.7)$   
 $RETURN$   
 $QSUM(LVL) = QSUN$   
 $LCRR(LVL) = ALPH$   
 $ALPHA = ALPH + DA$

$FV(1)$	$= CAT(LYL)$
$FV(3)$	$= F1T(LYL)$
$FV(5)$	$= F2T(LYL)$
$AREST$	$= F3TEST(LYL)$
$EPS$	$= QUEST(LYL)$
$GCTD1$	$= EPST(LYL)$
$GCTD2$	$= 2$
$GCTD3$	$= 3$
$GCTD4$	$= 2$
$GCTD5$	$= 2$

```

      DAT(LYL)
      FV(1) = F1T(LYL)
      FV(2) = F2T(LYL)
      FV(3) = F3T(LYL)
      AREST = AREST(LYL)
      EPS = EPST(LYL)
      GCUTD = 1
      GFLAG = 2
      GCTO = 2
      GFLAG = 3
      GCTO = 2
      END
      COMPLEX FUNCTION CIPAIPI*16(CK, CR, DR, RHO, OFST, SGMA, XSTN, C1CF, N, IPT)
      COMPLEX REAL*8 (A-H, C, P, R-Y), COMPLEX *16 {Q,Z}
      DIMENSION Q1CF(13)
      IF(IPT*GT*0) WRITE(*,990) DKDF1D!RHO, OFST, SGMA, XSTN, IPT
      990 FCRMAT('0', '0X', '0', 'Q1FAIP ENTERED', 'W', '1', '1', '1', '1', '1', '1', '1', '1', '1', '1', '1')
      1/2 'XSTN', '9X', 'CK', '11X', 'CR', '11X', 'DW', '11X', 'IP', '10X', '7(E12.5, ., ), 13)
      1/2 'XSTN = XSTN - OFST'
      IF(XASTA*LE*DR+RHC-1.0D0) GO TO 2C
      A = CR-1.0D0
      B = XASTN - RHO - 1.0D-8
      ICCNST = CDEXP( CCMLX(0.0D0, SGMA))
      CALL Q1PZSM(CK, DR, DW, FHO, XASTN, A, E, N, QANS, Q1CF, ICT)
      C1FAIP = QANS
      C1FAIP*LE.0) RETURN
      C1FAIP = CCMLX(0.0D0, 0.0D0)
      IF(IPT*LE.0) RETURN
      995 FORMAT('0.', '0X', 'Q1FAIP = ', E14.7, '., E14.7)
      RETURN
      END
      COMPLEX FUNCTION CIPRIP*16(CK, DR, DR, RHO, OFST, XSTN, C1CF, N, IPT)
      COMPLEX REAL*8 (A-H, C, P, R-Y), COMPLEX *16 {Q,Z}
      DIMENSION C1CF(13)
      IF(IPT*GT*0) WRITE(*,990) DKDF1D!RHO, XSTN, IPT
      990 FCRMAT('C', '0X', '0', 'Q1PRIP ENTERED', 'W', '1', '1', '1', '1', '1', '1', '1', '1', '1', '1', '1')
      1/2 'XSTN', '9X', 'IP', '10X', '5(E12.5, ., ), 13)
      2 IF(XSTN*LE*CR+RHO+CFFSET-1.0D0) GO TO 20
      A = CR-1.0C-8
      B = XSTN - RHO - 1.0D-8
      ICT = IPT - 1
      CALL Q1PZSM(CK, DR, CR, RHO, XSTN, A, B, N, QANS, C1CF, IOT)
      END

```

```

C1FRIP = -QANS
1F(1PT*LE,0) RETURN
CC10 60
CC1PR1P = DCMPLX(0.0D0,0.0D0)
2C Q1PR1P = DCMPLX(0.0D0,0.0D0)
60 WRITE(69,95) Q1PRIP
995 FORMAT(0.0,10X,C1FRIP = ' , E14.7,0,0,E14.7)
      RETURN
END
SUBROUTINE C1PZSN (DK'DR,DW,RHO,XSTN,A,B,J'CANS*Q1CF,IPT)
      REAL*8 (A-E,G,H,M,O,P,R)
      COMPLEX*16(F,Q,Z)
      DIMENSION Q1CF(13),LORR(60),F1T(60),F2T(60),CAT(C0),
     1 FEST(60),QEST(60),EPST(60),QPSUM(60)
      F(X) = GLGNCR(X,CR,J,Q1CF)*CDEXP((XEXP*(X))
      1 * HMBSC((CMEGA*DSQR((XSTN-X)*(XSTN-X))-YY)), IER)
      GAMMA = 1.4E0
      YY = RHC*RHO
      DM2 = (GAMMA + 1.0D0) * DW
      DMEGA = DSQRT(DK*DK*(1.0D0-DM2)/(DM2*DM2))
      CLAMDA = DK/DM2
      DK = DS_CRT(DM2)
      Q2EXP = CDEXP(DCMFLX(0.0D0,-CLAMDA*XSTN))
      QEXP = DCMPLX(0.0D0,CLAMDA)
      DU = 9.0D-13
      AC = 1.0D-6
      IF (1PT*GT-.C) WRITE(690) DK'DR,DH'RHO,XSTN,A,B,J,IPT
      990 FC_FMAT(.0,.15X*Q1PZSM ENTERED '1' ARGUMEN'S:;/,
     1 '15X*DK-.13X*DR*.13X*DW*.10X,RHC*,12X,'XSTN',11X,
     2 'A*'14X*B*14X*J*.2X*IP*.13',
     3 '15X*7(E14.7,0,12,2X,1,3),
     E FOURU=4.C*U
      EFLAG = 1
      EERROR = ACC
      LCLVL = 1
      LCCR(LV,L) = 1
      QFCSLM(LL,V) = 0.0
      ALPHAI = A
      DA = B-A
      AREA = 0.0
      AFEST = 0.0
      FV(1) = F(ALPHA)
      FV(3) = F(ALPHA + C*5*DA)
      FV(5) = F(ALPHA + EA)
      KCOUNT = 3
      WT = DA/6.0
      CEST = WT*(FV(1) + 4.0*FV(3) + FV(5))

```





```

SSC FCFORMAT('0','10X','Q1CRBP ENTERED WITH ','DW','1IX','RHO','1CX','XSTN','9X','IPT'.
2 '/XSTN'LE.RHO-1.0D0) GOT0 20
1CT = IPT - 1
QDK = DCMPXL(0,0DO,DK)
CALL QIDWM(DK,DR,DW,RHO,XSTN,QINP,IOT)
Q1DRBP = QCK*QINP(2) - QINP(1)
IF(IPT.LE.0) RETURN
GCI 60
Q1DRBP = DCMPXL(0,0DO,0.0D0)
20 IF(IPT.LE.0) RETURN
60 WRITE(6,995) Q1DRBP
995 FCFORMAT(0,10X,Q1CRBP = 'E14.7,' ,E14.7)
RETURN

C COMPLEX FUNCTION Q1DABP*1.6(DK,DR,DW,RHO,CFFSET,SIGMA,XSTN,IPT)
C IMPLICIT REAL*8 (A-H,O,P,R-Y), COMPLEX*16 {Q,Z}
CIMPLICIT QINP(2) WRITE(6,990) DKER,EW,RHO,OFFSET, SIGMA,XSTN,IPT
C IF(IPT*'C',10X,'Q1CABP ENTERED WITH ',EW,'10X,7(E12.5,',',I3)',7X,
C 1 'SIGMA',8X,'XSTN',SX,IPT),EW,1IX,'RHO',10X,7(E12.5,',',I3)',7X,
C 2 XASTN=XSTN - OFFSET
C IF(XASTN.LE.RHO-1.CD0) GO TO 20
C CT = IPT - 1
C CK = DCMPXL(0,0DO,DK)
CCCNST = CDEXP( CCMPXL(0,0DO,SIGMA))
CALL Q1DWXM(DK,DR,CW,RHO,XASTN,QINP,IOT)
C1DABP = (QCK*QINP(2) + QINP(1)) * QCONST
C IF(IPT.LE.0) RETURN
CCTO 60
20 Q1DABP = DCMPXL(0,0DO,0.0D0)
60 WRITE(6,995) Q1DAEP
995 FCFORMAT(0,10X,Q1DABP = 'E14.7,' ,E14.7)
RETURN

SUBROUTINE C1DWXM(CK,DR,DW,RHO, XSTN,QINP,IPT)
END

C C A IS THE MAXIMUM DEGREE OF THE TERM U IN THE INTEGRALS
C IMPLICIT REAL*8 (A-H,C,P,R-Y), COMPLEX * 16 (Z,C)
CIMPLICIT QINP(2) WRITE(6,990) DKER,RTC,XSTN,IPT
C IF(IPT*'C',10X,'Q1DWX ENTERED WITH ARGUMENTS:',IPT,',,
C 1 '10X,'16X,'DR,'16X,'RHO,'1IX,'XSTN',14X,'IPT',,
C 2 ',10X,4(E13.6,5X),12,5X,I3),

```

```

1CT = IPT - 1 - RHO - 1.0D-10
B = -1.0D0
CALL Q1DZXM(LK, DR, DW, RHO, XSTN, A, E, 1, QANS, 107)
CALL Q1CZXM(LK, DR, CW, RHC, XSTN, A, E, 2, QANS, 107)
CINF(2) = QANS
IF(IPT .LE. 0) RETURN
WFORMAT('6.995')
995 F1 = 10X,'J',6X,'QINP(I)',/
1 DC 1001,1 = 12 * 10X 12 EXP E12-6, ' , E12-6)
996 FFORMAT('6.996',1,CINP(I))
1001 WRITE(6,996)
      RETURN
END
SUBROUTINE C1DZXM (DK, DR, DW, RHO, XSTN, A, B, C, MFLX, IFL)
DIMENSION REA(8)(A-E2,G1-H1,D1-P1), F1(60), F2T(60), F3T(60),
          FV(5), LCRR(60), F1T(60), F2T(60), F3T(60), QPSUM(60)
1  AREST(60) = QESTT(60), QEPST(60), QPSUM(60)
1  F(X) = (XXXJ1)*CDEXP(QEXP*(X))
1  * MMESJ1((COMEGA*(XSTN-X)*(XSTN-X)-(YY))), IER)
2  * COMEGA*(XSTN-X)/(DSQRT((XSTN-X)*(XSTN-X)-(YY)))
2  GAMMA = 1.4E0
Y = RHO * RHO
CN2 = (GAMMA + 1.0E0) * DW
CNEGA = DSQRT(DK*DK*(1.0D0-DM2)/(CM2*DM2))
DLAMDA = DK/DM2
DN = DSQRT(CM2)
QEXP = CDEXF(DCMPLX(0.0D0, -CLAMDA*XSTN))
QEXP = CCMPLX(0.0D0, DLAMDA)
J1 = J - 1
J = 9.0D-13
ACC = 1.0D-6
1 IF(IPT .LT. 0) WRITE(6,990) DK, DR, RHO, XSTN, A, B, J, IFT
990 FCRPAT('0', '15X', 'Q1DZXM ENTERED WITH ARGUMENTS: /',
           '1', '15X', '0', 'DK', '13X', 'CR', '13X', 'CW', '10X, 'RHC', '12X', 'XSTN', '11X,
           '2', '14X', '0', 'BK', '14X', 'J', '2', 'X', 'IP', '1',
           '3', '15X', '7', 'E14.7', '1', '12', '2X', '13', '),
           EFCURU = 4.C*U
IFLAG = 1
IERS = ACC
QERROR = DCMPLX(CCDC, 0.0D0)
LVL = 1
LCRR(LVL) = 1
QFSLM(LVL) = 0.0
ALPHA = A - A
DA = B - A
QLDR6490
QLDR6495
QLDR6500
QLDR6505
QLDR6510
QLDR6520
QLDR6530
QLDR6535
QLDR6540
QLDR6545
QLDR6550
QLDR6555
QLDR6560
QLDR6565
QLDR6570
QLDR6575
QLDR6580
QLDR6585
QLDR6590
QLDR6595
QLDR6600
QLDR6610
QLDR6620
QLDR6630
QLDR6640
QLDR6650
QLDR6660
QLDR6670
QLDR6680
QLDR6690
QLDR6700
QLDR6710
QLDR6715
QLDR6720

```

```

AREA = 0.0
AREST = 0.0
FV(1) = F(ALPHA)
FV(3) = F(ALPHA + DA)
FV(5) = F(ALPHA + DA)
KCOUNT = 3
WT = DA/6.0
QEST = WT*(FV(1) + 4.0*FV(3) + FV(5))
DX = 0.5*DA
FV(2) = F(ALPHA + 0.5*DX)
FV(4) = F(ALPHA + 1.5*DX)
KCOUNT = KCOUNT + 2
WT = DX/6.0
QESTL = WT*(FV(1) + 4.0*FV(2) + FV(3))
QESTM = QESTL + QESTR
ARESTR = WT*(CDABS(FV(1)) + CDABS(FV(2)) + CDABS(FV(3)))
AREA = AREA + ((ARESTL + ARESTR) - AREST)
QDIFF = CDABS(QDIFF)
IF((DABS(DX) <= EPS*DABS(AREA))) GO TO 5
IF((LVLN*GE.4000) GO TO 6
IF((KOLNT*GE.6C) GO TO 5
LVL = LVL + 1
F1T(LVL) = FV(3)
F2T(LVL) = FV(4)
F3T(LVL) = FV(5)
LDT(LVL) = DX
AREST = ARESTL
ARESTR = ARESTR
QEST = QESTL
QESTR = QESTL
EEEST = EPS/1.4
EEEST(5) = EPS
FV(3) = FV(2)
FV(4) = FV(3)
FV(5) = FV(4)
GERROR = GERROR + QDIFF/15.0
GFSUM = GFSUM(LVL) EC.0 ) GO TO 4
LVL = LVL - 1
IF((LVL*GT.1) GO TO 3
GANS = ((B*J1)*C*EXP(QEXP*B) - CSLM)*Q2EXP/DM
IF(IFLAG.EQ.1) RETRN

```

1

2

3





```

ALPHA = A
DAE = B - A
AREA = 0.0
AREST = 0.0
FV(1) = F(ALPHA)
FV(3) = F(ALPHA + 0.5*DAI)
FV(5) = F(ALPHA + DAI)
KCLNT = 3
KWT = DA / C_0
KEST = WT*(FV(1) + 4.0*FV(3) + FV(5))
DX = 0.5*DA
FV(2) = F(ALPHA + C_5*DX)
FV(4) = F(ALPHA + 1.5*DX)
KCLNT = COUNT + 2
KWT = DX / KCLNT
QESTR = WT*(FV(1) + 4.0*FV(2) + FV(3))
CESTR = WT*(FV(3) + 4.0*FV(4) + FV(5))
QSUM = QESTR + CESTR
QAFFESTR = WT*(CDABS(FV(1)) + CDABS(FV(3)) + CDABS(FV(5)))
QAFFESTR = AREST + ((ARESTL + ARESTR) - AR)
QAFF = AREST - (QSLV
QAFF = QESTR - EPS*DABS(AREA) ) GO TO 1
QIF(CDABS(SQDIFF) .EQ. 4000) GO TO 5
LVL = KCLNT * GE * 4000) CC 10 6
LCCR(LVL) = 0
F1T(LVL) = FV(3)
F2T(LVL) = FV(4)
F3T(LVL) = FV(5)
DA = DX
DATE(LVL) = CX
ARESTT(LVL) = ARESTL
ARESTT = CESTL
CESTR = QESTR
CEFFEST(LVL) = QESTR
CEFFEST(LVL) = 1.0^4
CEFF(LVL) = FV(3)
F2(LVL) = FV(2)
GERRCR = GERROR + CCIFF/15.0
F1L = QSUM
LVL = QSUM(LVL) + GSUM
LVL = LVL - 1
IF(LVL .GT. 1) GO TO 2
CANS = (CLGNDR(B, CR, J, Q1CF)*CDEXF(CEX

```



7. Gorelov Spanning Function.

The subprogram for function used by Gorelov is shown below.

```
REAL FUNCTION PLGNDR*8(X,DR,N)
IMPLICIT REAL*8(A-F,C-Z)
IF(N.EQ.0) GOTO 100
ETA = CARCCS(-X)
ETASTR = CARCOS(1.000 - DR)
FN = CFLCAT(N)
PLGNDR = DCCS(FN*ETA)-DCOS(FN*ETASTR)
RETURN
100 PLGNDR = 1.000
RETURN
END
```

8. Linear Expansion Program

The program based on the linear expansion for  
small k is shown below.





```

IN = I + N
XSTN = XSTN - DR
XC = CDEXP(DCMPLX(C,0.0,DLAMDA*XSTN))
CALL Q1CAIP(DK,DR,EW,RHO,OFFSET,XSTN,SIGMA,XINTRP,K,ICF)N,10T)
CALL Q1CRIP(DK,DR,DW,RHC,OFFSET,XSTN,Q1NTRP,K,ICF)N,10T)
DO 20 J = 1,N
   J1 = J-1
   Q1INT(I,J1) = DCMPLX(1.0D0 ,DLAMDA*XSTN)
   Q1EW = LR-XSTN-1.0D0
   Q1INT(I,J2) = DCMPLX(TEMP,CLANCA*XSTN*TEMP)
   Q1INT(I,J3) = Q1NTRP(J)
   Q1INT(I,J4) = Q1NTAP(J)
   Q1XSTN = XSTN - OFFSET
   Q1INT(IN,3) = DCMPLX(1.0D0 ,DLAMDA*XASTN)
   TEMPINT(IN,4) = DCMPLX(TEMP ,DLAMDA*XASTN*TEMP)
2C   Q1CCF(IN) = -Q1CRBP(DK,DR,DW,RHO,XSTN,10T)
      Q1CCF(I) = -Q1CAEP(DK,DR,DW,RHC,CFFSET,SIGMA,XSTN,10T)
90  RETURN
ENDC
COMPLEX FUNCTION Q1CRBP*16(DK',DR,EW,RHO,XSTN,16 IPT)
COMPLEX FUNCTION Q1CAEP*16(A-H,0,P,R-Y), COMPLEX XSTN 16 (Q,Z)
COMPLEX QINP(2)
IF (IPT.GT.0) WRITE(6,990) DK, CR, EW, RHO, XSTN, IPT
990 FCRMAT(0,10X,1DK,1IX,ER,11X,DR,1IX,1RHO,10X,'XSTN',9X,'IPT'
12 IF (XSTN.LE.RHO-1.0D0) GOT 20
12 IF (XSTN.GT.-1.0D0) GOT 20
12 QICK = DCMPLX(0.0D0 ,DK)
GAMMA = 1.4D0
CLAMDA = DK/((GAMMA+1.0D0)*DW)
Q1CRBP = DCMPLX(1.0D0 ,(DK+DLAMDA)*(XSTN-DR))
12 IF (IPT.LE.0) RETURN
12 GCTOR30 = DCMPLX(0.0D0 ,0.0D0 )
2C   Q1CRBP = DCMPLX(0.0D0 ,0.0D0 )
2C   IF (IPT.LE.0) RETURN
3C   WRITE(6,995) Q1CRBP
995 FORMAT('C',10X,'Q1CRBP = ',E14.7,'.',E14.7)
3C   RETURN
ENDC
COMPLEX FUNCTION C1CABP*16(DK',DR,EW,RHO,OFFSET,SIGMA,XSTN,IPT)
1AFLICIT REAL * 8 (A-H,0,P,R-Y), COMPLEX * 16 (Q,Z)

```



```

DC 40 J = 11N J CCAIP (J)
956 FCNFORMAT (6, 996) N1 26X, 12, 3X, E14.7, ' , ' , E14.7)
4C CONTINUE
READ
SUBROUTINE CICRIP (OK, DR, Dh, RHO, CPLEX, XSTN, QCRIP, IPT)
DIMENSION QCRIP(13)
IF (IPT .GT. 0) WRQ1CRIP(6, 990) DKDR1H, RHO, XSTN, IPT
990 1 FCRNAT (0, !OK, !10X, !Q1CRIP ENTERED W1H, !/HC, !CX,
2 !XSTN, !9X, !IP, !1X, !DR, !11X, !11X, !RHC, !CX,
1F (XSTN, !E, !DR+RHO+CFFSET-1.0D0) GOTO 20
1F (XSTN, !G, !2D0) GOTO 20
IPT = IPT - 1
GAMMA = 1.4LO
DLAMDA = DK / ((GAMMA + 1.0D0)*DW)
QCRIP(1) = CCMLX(1, 0D0, DLAMDA*(XSTN-RHO))
TERP = DR+RHO-XSTN-1.0D0
QCRIP(2) = DCMLX(TEMP, DLAMDA*(XSTN-RHO)*TENF)
IF (IPT.LE.0) RETURN
GOTO 30
20 ZERO = CCMLX(0.0D0, 0.0D0)
DC 25 I = 11N
25 QCRIP(I) = ZERO
IF (IPT.LE.0) RETURN
30 WRITE (6, 995)
55 DC 40 J = 11N
DC 40 WRITE (6, 996) J, CCRIP (J)
996 FORMAT (6, 26X, 12, 3X, E14.7, ' , ' , E14.7)
40 CONTINUE
RETURN
END
SUBROUTINE C1COEF (Q1COOF, Q1INT, N1FT, Q1ABCF, Q1RBCF)
DIMENSION REAL*8(A-H, P-R-Y), COMPLEX*(Z, Q)
DIMENSION Q1COOF(26), Q1INT(26, 26), ZWA(300)
1F (IPT.GE.0) WRITE (6, 90)
1B = 26
M = 1
N2 = 2*N
1F (IPT.LE.0) GO TO 5
55 WRITE (6, 98) N1 N2
98 1 FCRNAT (0, !EX, !Q1COEF ENTERED WITH ' , 12, ' DEG PWR SERIES (' , 12,
1 CC 2 !SQUARE MATRIX)
1 CC 2 ! = 1, N2

```









```

ARESTR = WT*(CDABS(FV(3)) + CDABS(FV(4))) + CDABS(FV(5))
AREAAF = AREAT + (ARESTR - QSUM)
QDF(CDABS(QCUFF) * LE * EPS*DABS(AREA)) GO TO 2
IF(DABS(IDX) * LE * FCUR*DABS(ALPHA)) GO TO 5
IF(KCUNT * GE * 4000) GO TO 6
LCFR(LVL) = 0
F1T(LVL) = FV(3)
F2T(LVL) = FV(4)
F3T(LVL) = FV(5)
LCFR(LVL) = CX
DAT(LVL) = DX
ARESTT(LVL) = ARESTL
ARESTT(LVL) = ARESTR
QESTT(LVL) = QESTL
QESTT(LVL) = QESTR
EPSST = EPS/1.4
EPSST(LVL) = EPS
FV(3) = FV(2)
FV(5) = FV(2)
CURROR = QERROR + QDIFF/15.0
QSUM = QPSUM(LVL) + CSUM
LVL = LVL - 1
IF(LVL * GT * 160 TO 3
IQIDCL = QSUM * 20E0
IF(IPT * GT * 0) GO TO 1
IF(IFLAG * EQ * 1) RETURN
WRITE(6, 990) DKDR, DRHO, OFFSET, SIGMA, N, IPT
990 FORMAT(15X, 15X, IFLAG, RESULTS: QIDCL = ', E14.7, ', E14.7, ', 1,
2, 15X, 13, 2X, E14.7, ', E14.7)
RETURN
QSUM(LVL) = QSUM
LCRR(LVL) = 1
ALPHA = DAT(LVL) + DA
DA = DAT(LVL)
FV(1) = F1T(LVL)
FV(3) = F2T(LVL)
FV(5) = F3T(LVL)
AREST = QESTT(LVL)
EPS = EPST(LVL)
GC TO 1

```

```

5   IFLAG = 2
6   GC TO 2
      COMPLEX FUNCTION Q1DCM*16(DK,DR,CH,RHO,OFFSET,SIGMA,N,Q1BCF,C1RBC
      IF IPT1CIT REAL*8(A-E,G1H0,P13)-Y), COMPLEX*16(F,C,Z)
      IF PLENSION Q1ABCFC15)F1(60)F2(60)AREST(60),CEST(60),CPSUM(60)
      DCNTRN DAT(60),EFS(60)
      C(X)=X*(Q1DPHI(DK,DR,DW,RHO,OFFSET,SIGMA,N,Q1ABCFC1RBCF,X,ICT))
      IF (IPT10)WRITE(6,990)DK,ER1N,RHO,CFSET,SIGMA,N,IP1
      950 FORMAT(0.10X,Q1DCM ENTERED WITH:/'10X,13X,DR,13X,DK,12X,RHO,12X,OFFSET,9X,SIGMA'
      456 10X,13X,1PT1,13)
      A=1E12.5,3X,12,2X,13)
      IPT=1.0P-1
      GAMMA=1.4C0
      CLAMDA=LK/(GAMMA+1.0D0)*Lh)
      QALPHA=DCNPLX(0.0D0,DLAMDA-DK)
      =9.0E-13
      ACC=1.0D-5
      EFCURU=4.0CU
      IF FLAG=1
      CERROR=DCNPLX(0.0DC,0.0D0)
      LVR=1
      LCR(LVL)=1
      ALPFA=A
      CA=B-A
      AREA=0.0
      AREST=0.0
      FV1=F(ALPHA)
      FV2=F(ALPHA+0.5*DA)
      FCNT=3
      KFCNT=DA/6.0
      KEST=KFCNT*(FV1)+4.0*FV(3)+FV(5))
      DX=0.5*DA
      FV(2)=F(ALPHA+0.5*DX)
      FV(4)=F(ALPHA+1.5*DX)
      KCNT=KCNT+2

```

1

```

 $WT = DX/6.0$ 
 $CESTL = WT*(FV(3) + 4.0*FV(4) + FV(5))$ 
 $CSESTR = QESTL + QESTR$ 
 $ARESSTL = WT*(CDABS(FV(1)) + CDABS(FV(3)) + CDABS(FV(4)) + CDABS(FV(5)))$ 
 $AREAA = AREAT + (ARESTL + ARESTR) - ARESTR$ 
 $AREAFF = QUEST - CSCLW$ 
 $ARETDAB = SQDIFF(DX) * EPS * DABS(ALPHA) GO TO 5$ 
 $ARETDAB(GE .60) GO TO 5$ 
 $ARETDAB(GE .4000) GO TO 6$ 
 $QDIFF(LVL) = 0$ 
 $LORR(LVL) = FV(3)$ 
 $F1*(LVL) = FV(4)$ 
 $F2*(LVL) = FV(5)$ 
 $DATA(LVL) = DX$ 
 $ARESTT(LVL) = ARESTL$ 
 $QESTT(LVL) = QESTR$ 
 $QEFPT(LVL) = QESTR$ 
 $QEFPT(LVL) = EPS$ 
 $FV(3) = FV(2)$ 
 $FV(2) = 1$ 
 $QERROR = QERROR + QCIFF/15.0$ 
 $QSUM = QSUM(LVL) + QSUM$ 
 $LVL = LVL - 1$ 
 $QSUM = QSUM * 2.000$ 
 $IF(IFLAG .EQ. 1) RETRN$ 
 $IF(IFLAG .EQ. 1) DKDR, DW RHO, OFFSET, SIGMA, N, IPT$ 
 $WRITET(6,9901) IFLAG, QICCM, QERRR$ 
 $FORMAT(15X,15X,RESULTS: QICCM = , E14.7, , , E14.7, , , E14.7)$ 
 $12 FETURN$ 
 $QSUM(LVL) = QSUM$ 
 $LCRR(LVL) = 1$ 
 $ALPHA = ALPHA + CA$ 
 $D2 = CAT(LVL)$ 
 $FV(1) = F1T(LVL)$ 
 $FV(3) = F2T(LVL)$ 

```





```

CCPLEX FUNCTION C1CREP*16(CK'DR, CW'RHO, XSTN, IPT)
CCPLEX ICINP(2) WRITE(6,990) CK'DR, CW'RHO, XSTN, IPT
IF(IPT .EQ. 0) ENTERED
FCRMAT(.C.,10X,1CK,11X,1CR,11X,1RH,11X,1CX,1RC,1CY,1STN,9X,'IPT'
12 IF(XSTN .NE. RHO-1.0E0) GOTO 20
ICMMA = 1.4E0
CM2 = (GAMMA + 1.000) * DW
CLAMDA = DK/DM2
CLAMCA = (GAMMA + 1.CC0) * DW
PIMAG = (DK+DLAMDA)*(XSTN-RHO)
Q1CRBP = -DCMPLX(1.000,PIMAG-CLAMCA*XSTN)/CSQRT(DM2)
Q1F(IPT,LE,0) RETURN
G1CRBP = DCMPPLX(0.000,0.000)
20 Q1F(IPT,LE,0) RETURN
60 WRITE(6,995) Q1FBP
55 FCRMAT(.C.,10X,1Q1DRBP = ' , E14.7, ' , E14.7)
RECF
CCPLEX FUNCTION C1DAPP*16(DK'DR, CW'RHO, OFFSET, SIGMA, XSTN, IPT)
CCPLEX ICINP(2) WRITE(6,990) DK'DR, CW'RHO, OFFSET, SIGMA, XSTN, IPT
IF(IPT .EQ. 0) ENTERED
FCRMAT(.C.,10X,1CK,11X,1CR,11X,1RH,11X,1RC,1CY,1STN,9X,'IPT'
12 SIGMA = 8.2E-19X,1IPTR/1.10X,7(E12.5,1.12)
X4STN = XSTN - OFFSET
11CT = IPT-1
ICMNST = CCEXP( DCMPPLX(0.000,SIGMA))
CM2 = 1.4E0
CLAMDA = DK/DM2
CLAMAG = (DK+DLAMDA)*(XSTN-RHO)
Q1CAPP = DCMPLX(1.000,PIMAG-CLAMCA*XSTN)*CCCNST/CSQRT(DM2)
Q1F(IPT,LE,0) RETURN
11CT=TOP60
C1CAEP = DCMPPLX(0.0DC,0.0D0)
20 C1F(IPT,LE,C) RETURN
60 WRITE(6,995) Q1DABP
55 FCRMAT(.C.,10X,Q1CAEP = ' , E14.7, ' , E14.7)
RECF
CCPLEX FUNCTION C1FAIP*16(CK,DR,CW,RHO,CFST,SGMA,XSTN,C1CF,N,IPT) INC72CC

```

```

16 FUNCTION REAL*8 (A-H,C,P,R-Y), CCPLEX * 16 (C,2)
17 DIMENSION QJCF(13)
18 IF (IPT .GT. 0) WRITE(6,990) CK, DR1CW, RHO, OFST1, SCMA, XSTN, IPT
19 FCRMAT=1.0X-1.0X-Q1FA1*ENTERED, WITH: /RHO/ 10X, OFST 1, 9X, *SCMA*, SX
20 XSTN = XSTN - OFST
21 XSTN = XSTN - LE*DR+RHC-1.0D0) GO TO 20
22 IFCNSTA = CCEXP( DCMPLX(0.0D0, SGMA1)
23 DM2 = (GAMMA1+1.0D0)*DW
24 CLAMDAD = DK/DM2
25 XX = XASTTN-R1Q
26 Q1 = Q1ICF(1)
27 Q2 = Q1ICF(2)
28 PR1 = Q1XX+1.0D0-DR
29 PR2 = (LR-1.0D0-X/2.0D0)*XX - (CR-1.0D0)*(DR-1.0D0)/2.0D0
30 FIMAG1 = (XX*XX+(2.0D0-DR)*DR-1.0D0)*DLAMDA/2.0D0
31 PIMAG2 = (XX*XX*((DR-1.0D0)/2.0D0-XX)*(DR-1.0D0)*3)/6.0D0
32 PIMAG1 = PIMAG1*DLAMDA
33 PIMAG2 = PIMAG2*(PR1, FIMAG1-DLAMDA+PR1*XASTN)
34 QP1 = DCMPLX(PR1, PIMAG2-DLAMDA*FR2*XASTN)
35 QP2 = DCMPLX(PR2, PIMAG2-QP1+Q2*QP2)/CSQRT(LN2)
36 QFCAIP = (Q1*QP1+Q2*QP2)/CSQRT(LN2)
37 QFCAIP = LE.01 RETURN
38 CCCTCAIP = DCMPLX(0.0D0,0.0D0)
39 IF (IPT .LT. 0) RETURN
40 WRITE(6,991) QIPAIP
41 FORMAT(-0.,10X,01PAIP = , E14.7)
42 RETURN
43 END

44 COMPLEX FUNCTION QIPRIP*16(CK, DR1CW, RHO, OFFSET, XSTN, C1CF, N, IPT)
45 DIMPLEXIC1CF(13)
46 IF (IPT .GT. 0) WRITE(6,990) CK, DR1CW, RHO, XSTN, IPT
47 FCRMAT=C10X-Q1PRIP ENTERED, WITH: /RHO/ 11X, 10X, *RHO/ 10X,
48 XSTN = XSTN - LE*IP-1.0D0*FFSET-1.0D0) GO TO 20
49 GAMMA = 1.0D0
50 CN2 = (GAMMA+1.0D0)*DW
51 CLAMDAD = DK/DM2
52 XX = XSTN - RHO
53 C2 = Q1ICF(2)
54 PR1 = XXX+1.0D0-DR
55 PR2 = (CR-1.0D0-XX/2.0D0)*XX - (DR-1.0D0)*(DR-1.0D0)/2.0D0

```

```

N07690
N07700
N07710
N07720
N07730
N07740
N07750
N07760
N07770
N07780
N07790
N07800
N07810
N07820
N07830
N07840
N07850
N07860
N07870
N07880
N07890
N07900
N07910
N07920
N07930
N07940
N07950
N07960
N07970
N07980
N07990
N08000
N08010
N08020
N08030
N08040
N08050
N08060
N08070
N08080
N08090
N08100
N08110
N08120
N08130
N08140
N08150
N08160

PIMAG1 = (XX*XX + (2.0D0-DR)*DR-1.0D0)*DLAMDA/2.0D0
PIMAG2 = XX*(DR-1.0D0)/2.0D0-XX)-(DR-1.0D0)**3)/6.0D0
PIMAG1 = PIMAG1*DLAMDA
PIMAG2 = PIMAG2*DLAMDA
P1M1 = DCMPLX(PR1-PIMAG1-DLAMDA*PR1*XSTN)
P1M2 = DCMPLX(PR2-PIMAG2-DLAMDA*PR2*XSTN)
Q1P1P = -(Q1*QP1+C2*CP2)/DSQRT(LM2)
IF(IPT.LE.0) RETURN
CUTRIP = DCMPLX(0.0D0,0.0D0)
20 CIF(IPT.LE.0) RETURN
30 WRITE(6,995) Q1P1P
995 FCRNAT(6,995,10X,Q1P1P = ' ,E14.7,*,*,E14.7)
FETURN
EAD
CCMPLEX FUNCTION Q1DAIP*16(DK,DF,CH,RHO,OFST,SGMA,XSTN,C1CF,N,IPT)
10 IF(CPLXIT REAL#8(A-H,0,P,R-Y),CCMPLEX * 16 {C,Z)
11 IF(CPLXIT C1CF(13) WRITE(6,990) DK,DR1CH,RHO,CFS1,SGMA,XSTN,IPT
12 FCRNAT(10X,Q1DAIP*10X,1Q1X,DR,ENTERED WITH /'1/RHO'/CX,OFST*9X,'SGMA',SX
13 XSTN,9X,1PT/,*,10X,71E12.5,*,13},13
14 XSTN-OFST
15 FCRNAT(10X,LE,DR+RHC-1.0D0) GC,TO,20
16 GMM1=1.4D0
17 GMM2=DK/DM2
18 X1=XA$TN-RHO
19 Q1=C1CF(1)
20 Q2=C1CF(2)
21 PR1=1.0D0
22 PR2=-DR-1.0D0
23 P1MAG1=DLAMDA*XX
24 P1MAG2=DLAMDA*(DR-1.0D0-XX)*XX
25 CFS1=DCMPLX(PR1-PIMAG1-DLAMDA*PR1*XASTN)
26 CFS2=DCMPLX(PR2-PIMAG2-DLAMDA*PR2*XASTN)
27 C1CAIP=(Q1*QP1+Q2*CP2)/DSQRT(LM2)
28 IF(IPT.LE.0) RETURN
29 C1DAIP = DCMPLX(0.0D0,0.0D0)
30 CIF(IPT.LE.0) RETURN
31 FCRNAT(6,995,Q1DAIP = ' ,E14.7,*,*,E14.7)
32 FETURN
END
CCMPLEX FUNCTION C1DRIP*16(DK,CR,CH,RHO,OFFSET,XSTN,C1CF,N,IPT)
10 IF(CPLXIT REAL#8(A-H,0,P,R-Y),CCMPLEX * 16 {Q,Z)

```

```

      DIMENSION Q1CF(13) WRITE(6,990) OKIDRIF, RHO, XSTN, IPT
      IF (IPT .GT. 0) WRITE(6,990) OKIDRIF, ENTERED, ' / '
      FORMAT(10X,'10X','10X','10X','10X','10X','10X','10X',
     12   'XSTN','9X','10X','10X','10X','10X','10X','10X,
     12   'IF (XSTN .LE. LR+RHC+OFFSET-1.0) 20
      GAMMA = 1.4*CO
      CLAMDA = (GAMMA+1.*DD0)*Dh
      Q1 = Q1CF(1)
      Q2 = Q1CF(2)
      XX = XSTN - RHO
      PR1 = 1.0*DD0
      PR2 = DR-1.*DD0-XX
      PIAG2 = CLAMDA*(DR-1.*DD0-XX)*XX*XX
      PIAG1 = DLAMDAXX
      QP1=DCMPLX(FR1*PIAG1-DLAMDAA*PR1*XSTN)
      QP2=DCMPLX(PR2*PIAG2-DLAMDAA*PR2*XSTN)
      Q1ERIP=(C1*CP1+Q2*QP2)/CSQRT(DW2)
      IF (IPT.LE.0) RETURN
      CTC30 = DCMPLX(C*DD0,0.000)
      CIDRIP = DCMPLX(C*DD0,0.000)
      IF (IPT.LE.0) RETURN
      WRITE(6,995) Q1DRIP
      FORMAT(0.,10X,'Q1DRIP = ',E14.7,' , E14.7)
      RETURN
      END

```

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