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NO. 486

THE IMPACTION FORCE OF AIRBORNE PARTICLES
ON SPHERES AND CYLINDERS (U)

by

Stanley B. Mellisen

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ABSTRACT

The effect of dust on aerodynamic drag of spheres and cylinders was calculated by using a mathematical model developed for this purpose. The results were compared to experiments previously done by other workers. The calculated and experimental results agree favourably, showing that the mathematical model is satisfactory. Impaction efficiencies and drag coefficients due to dust alone were then obtained using the model for a wide range of the inertia parameter and the results are presented graphically. The model can also be used for calculating velocity distributions and points of impact for a stream of airborne particles flowing over a sphere or cylinder.

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NOTATION

A	frontal area of target, cm^2
A_d	far upstream cross sectional area of the envelope of particles which eventually hit the sphere or cylinder, cm^2
a	tube radius, cm
C_D	dimensionless drag coefficient for spheres or cylinders in air
C_{DC} or C_{DS}	drag coefficient of cylinder or sphere due to stream of particles alone
C_{De}	total effective drag coefficient for spheres or cylinders due to particles and air together
C_{DL}	drag coefficient for spheres in tubes adjusted to a sphere in free space by the Landenburg correction
C_{DpU}	drag coefficient for particles alone adjusted from v_{rms} to U
d	particle diameter, cm
E_m	impaction efficiency of particles on a target
F	force exerted on a target by the air and total stream of particles hitting it, dynes
F_a	force exerted on a target by the air alone, dynes
F_s or F_c	force exerted on the sphere or cylinder by the total stream of particles which hit it, dynes
F_p	force exerted on a target by the particles alone travelling in air, dynes
L	cylinder or sphere radius, cm
m'	total mass of particles which had its momentum changed in unit time, g s^{-1}
p_s or p_c	drag pressure on sphere or cylinder due to stream of particles alone, dynes cm^{-2}
r	distance from any point to the origin, cm
t	time, seconds
U	free-stream velocity, cm s^{-1}
u	local fluid velocity, cm s^{-1}

u_r	radial component of fluid velocity, cm s^{-1}
u_θ	circumferential component of fluid velocity, cm s^{-1}
v	local particle velocity, cm s^{-1}
v_{rms}	root-mean-square free stream particle speed upstream near target position, cm s^{-1}
v_x	component of particle velocity parallel to free field flow direction immediately before impact, cm s^{-1}
v_x'	component of particle velocity parallel to flow direction immediately after impact and reflection, cm s^{-1}
v_{x0}	far upstream particle velocity, cm s^{-1}
v_y	component of particle velocity perpendicular to free field flow direction immediately before impact, cm s^{-1}
Δv_x	change in the component of particle velocity parallel to the free stream flow direction caused by reflection from the sphere or cylinder, cm s^{-1}
x	co-ordinate (origin at centre of sphere or cylinder) of particle position parallel to free stream flow direction, cm
y	transverse co-ordinate of particle position, cm
y'	off-axis distance of a particle at point of impact with the sphere or cylinder, cm
y_c	far upstream transverse co-ordinate of the envelope of particles which eventually hit the sphere or cylinder, cm
α	angle of incidence of a particle against the sphere or cylinder, radians
β	angle between particle direction and free field flow direction just before the particle hits the sphere or cylinder, radians
γ	total angle change of a particle from its far upstream direction to its direction after reflection from the cylinder or sphere, radians
θ	polar angle between x axis and <u>radius vector</u> to particle position, radians
μ	absolute velocity of fluid, poise
ρ	fluid density, g cm^{-3}
ρ^*	total mass of particles and air per unit volume of space, g cm^{-3}
ρ_p	uniform particle density per unit volume of air far upstream, g cm^{-3}

σ particle density, $g\text{ cm}^{-3}$

The following were dimensionless:

$f(\bar{y})$	change in component of particle velocity due to reflection from sphere or cylinder as a function of off-axis distance far upstream
K	particle inertia parameter
Re	spherical particle Reynolds number in flow influenced by presence of sphere or cylinder
Re ₀	spherical particle Reynolds number in free stream
\bar{r}	$\frac{r}{L}$, distance from any point to the origin
T	time $\frac{tU}{L}$
\bar{u}	local fluid velocity $\frac{u}{U}$, \bar{u}_x and \bar{u}_y are x and y components
\bar{v}	local particle velocity $\frac{v}{U}$, \bar{v}_x and \bar{v}_y are x and y components
\bar{v}_x	$\frac{d\bar{x}}{dT}$, parallel component of particle velocity
\bar{v}_y	$\frac{d\bar{y}}{dT}$, transverse component of particle velocity
\bar{x}	parallel co-ordinate $\frac{x}{L}$
\bar{y}	transverse co-ordinate $\frac{y}{L}$
\bar{y}_d	far upstream transverse co-ordinate of the envelope of particles which eventually hit the cylinder
\bar{y}_i	co-ordinate of particle used in discretization for integration; y_i increases with i from $\bar{y}_1 = 0$ to $\bar{y}_n = \bar{y}_d$
ϕ	dimensionless group independent of drop size formed by combining Re ₀ and K

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1. INTRODUCTION

Extensive studies have been made of the drag exerted on spheres and transverse cylinders by moving air (Schlichting, 1960) and attempts have been made to determine the effect of introducing dust into the air stream (Gillespie and Gunter, 1957; 1959). Also, the trajectories of water drops moving in the neighborhood of a circular cylinder placed in a uniform stream of air have been calculated by numerical solution of the equations of motion (Glauert, 1940) and with a mechanical analog (Brun et al., 1953). The latter method has been used in conjunction with flight instruments used to study droplet size and distribution in icing clouds (Brun and Mergler, 1953). Measurements of aerodynamic drag on circular cylinders due to the blast wave from large TNT bursts have been done at DRES (Mellisen, 1974). The results of these tests have been used to provide blast loading input for structural analysis of lattice type masts also tested in the trials. On some of the trials there have been visible quantities of dust associated with the blast wave. In general the measured

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strains in the masts have agreed favorably in an individual field trial with the ones predicted from theoretical analysis based upon the blast loading input from the drag data averaged over several blast trials (Laidlaw, 1977). The influence of dust is difficult to determine and generally considered negligible.

Another area of interest in connection with moving particles in an air stream is the impingement of liquid drops on protective clothing and military equipment. The impaction efficiency and force, and the droplet distribution are all of interest. The study of these effects on transverse cylinders and spheres is useful in this application.

The purpose of this report is to describe numerical solutions to the equations of motion of particles in a moving air stream, with possible applications to both the effect on aerodynamic drag from the point of view of blast loading of structures and the interaction of droplets with spheres and cylinders from a chemical defence point of view. The mathematical techniques along with their associated computer programs are described herein. Also, the results are compared with experiments done at DRES (Gillespie and Gunter, 1957).

2. DEFINITION OF THE PROBLEM

A spherical particle flowing at free stream velocity far upstream of a spherical or transverse cylindrical solid target will not necessarily follow a streamline in the vicinity of the target where the radial and axial velocity components of the fluid may be changing markedly. A particle that just impinges on the outer circumference of the target was considered. All particles of the same diameter within an envelope generated by this limiting particle trajectory would hit the target, assuming that no particles collide with other particles reflected from the target.

All particles in the steady flow field far upstream were assumed to be travelling at the velocity of the fluid but the velocities of the particles will differ from the fluid velocity as the particles approach the target. This is due to inertial effects of the particles in

the flow field near the object. The velocity of a particle immediately before contact, and also the point of contact with the target, depend upon the position of the particle within the limiting envelope far upstream from which it travels to the target. These particles all have their momentum changed when they collide with the target and in the case of solid particles are deflected in a new direction. The total change in momentum is balanced by the drag force exerted by the particles on the target.

The basic problem was, therefore, to determine the envelope of particles hitting the target and the change in particle motion due to the presence of the target in the flow. Then the drag coefficient due to particles alone could be obtained.

3. EQUATIONS OF MOTION

The motion of an individual spherical particle has been shown (Batchelor, 1967) to be determined by the following ordinary differential equations:

$$\frac{d\bar{v}_y}{dT} = \frac{C_D \text{Re}(\bar{u}_y - \bar{v}_y)}{24K} \quad (\text{Eq. 1})$$

$$\frac{d\bar{v}_x}{dT} = \frac{C_D \text{Re}(\bar{u}_x - \bar{v}_x)}{24K} \quad (\text{Eq. 2})$$

where $\text{Re} = \text{Re}_0 [(\bar{u}_y - \bar{v}_y)^2 + (\bar{u}_x - \bar{v}_x)^2]^{1/2}$ (Eq. 3)

$$K = \frac{\sigma d^2 U}{18\mu L} \quad \text{particle inertia parameter} \quad (\text{Eq. 4})$$

$$\text{Re}_0 = \frac{U d \rho}{\mu} \quad \text{free stream Reynolds number} \quad (\text{Eq. 5})$$

The symbols are defined in the notation section near the front of this report and the basic geometry of the flow system is illustrated in Fig. 1.

Several assumptions were inherent in the development of Eq. 1 and Eq. 2 for calculating impaction efficiency and force due to a stream of particles, including,

- (a) uniform particle distribution,
- (b) no gravitational or electrostatic forces of consequence,
- (c) monodisperse spherical particles with diameter very small in relation to the target sphere or cylinder diameter, and
- (d) free stream flow that was steady, incompressible and irrotational.

The drag coefficient is a function of Reynolds number and was available in the form of definitive empirical equations (Davies, 1945). These equations are stated as follows:

$$Re = \frac{C_D Re^2}{24} - 2.3363 \times 10^{-4} (C_D Re^2)^2 + 2.0154 \times 10^{-6} (C_D Re^2)^3 - 6.9105 \times 10^{-9} (C_D Re^2)^4 \quad (\text{Eq. 6})$$

for $Re < 4$ or $C_D Re^2 < 140$

$$\log_{10} Re = -1.29536 + 9.86 \times 10^{-1} (\log_{10} C_D Re^2) - 4.6677 \times 10^{-2} (\log_{10} C_D Re^2)^3 + 1.1235 \times 10^{-3} (\log_{10} C_D Re^2)^3 \quad (\text{Eq. 7})$$

for $3 < Re < 10^4$ or $100 < C_D Re^2 < 4.5 \times 10^7$

4. AIR FLOW FIELD

The assumption was made that the air flow in front of the cylinder or sphere is given by classical hydrodynamics theory.

a) Cylinder

The equations of fluid velocity were derived from the stream function for ideal flow around a cylinder (Milne-Thomson, 1960). These were normalized to the free stream velocity and cylinder diameter and written as follows:

$$\bar{u}_x = 1 - \frac{\bar{x}^2 - \bar{y}^2}{\bar{r}^4} \quad (\text{Eq. 8})$$

$$\bar{u}_y = 1 - \frac{\bar{x}\bar{y}}{\bar{r}^4} \quad (\text{Eq. 9})$$

b) Sphere

For purposes of comparison of the results of the theoretical work reported herein to existing experimental data the flow field about a sphere on the axis of a circular tube was applied. The flow about a sphere in the absence of the tube was then readily available from this as a limiting case.

The flow velocity was derived (Mellsen et al., 1966) from the vector potential (Smythe, 1964). The velocity components are given in terms of a series solution. For a sphere diameter of 0.8 cm and tube diameter of 2.54 cm used in the experiments (Gillespie and Gunter, 1957) sufficient accuracy (Mellsen et al., 1966) was obtained by retaining three terms of the series.

The equations were

$$\frac{u_r}{U} = -\cos \theta \left[1 - \frac{2C_0}{3} \frac{C}{r} - \frac{2A_0}{a} + \dots \right] \quad (\text{Eq. 10})$$

$$\frac{u_\theta}{U} = \sin \theta \left[1 + \frac{C_0}{3} \frac{C}{r} - \frac{2A_0}{a} + \dots \right] \quad (\text{Eq. 11})$$

The constants C_0 , C and A_0 which are dependent upon the sphere and tube diameters (Smythe, 1964) are given as follows:

$$C_0 = 1.5401075$$

$$C = 1.0 \text{ (normalized sphere radius)}$$

$$A_0 = -C \left(\frac{C}{a} \right)^2 \frac{I(2)C_0}{9\pi}, \quad I(2) = 7.5098907$$

$$a = \frac{2.54}{0.8} C \text{ (normalized tube radius)}$$

u_θ and u_r are shown in Fig. 2.

Converting to rectangular co-ordinates gives

$$\bar{u}_x = -\frac{u_r}{U} \cos \theta + \frac{u_\theta}{U} \sin \theta \quad (\text{Eq. 12})$$

$$\bar{u}_y = \frac{u_r}{U} \sin \theta + \frac{u_\theta}{U} \cos \theta \quad (\text{Eq. 13})$$

The equation for flow around a sphere in the absence of the tube is given by letting the value of "a" approach infinity.

For $a \rightarrow \infty$, $C_D \rightarrow 1.5$ (Smythe, 1964)

$$\text{Eq. 10 becomes } \frac{u_r}{U} = -\cos \theta \left[1 - \left(\frac{L}{r} \right)^3 \right] \quad (\text{Eq. 14})$$

$$\text{Eq. 11 becomes } \frac{u_\theta}{U} = \sin \theta \left[1 + \frac{1}{2} \left(\frac{L}{r} \right)^3 \right] \quad (\text{Eq. 15})$$

Eqs. 14 and 15 are well known in ideal flow theory (Batchelor, 1967) and therefore serve as a partial check on the correctness of Eqs. 10 and 11.

5. VELOCITY CHANGE OF AN INDIVIDUAL PARTICLE

To calculate the drag force due to airborne particles, the velocity change of an individual particle due to the presence of the cylinder or sphere in the flow was first considered. The velocity change determined applied to both the cylindrical and spherical targets because the same geometry could be used for both (Fig. 3).

The motion considered concerned a spherical particle which starts far upstream within the envelope of particles which travel to the target, eventually hit it and are reflected by it in a new direction. The total change in direction of a particle which comes from a far upstream off-axis position y , and hits the target at a point defined by y' (Fig. 3), is represented by γ . For a spherical particle the angle of incidence, α , is equal to the angle of reflection. Then

$$\gamma = \left(\frac{\pi}{2} - \theta \right) + \alpha \quad (\text{Eq. 16})$$

where

$$\theta = \sin^{-1} \left(\frac{y'}{L} \right) = \tan^{-1} \left(\frac{y'}{\sqrt{L^2 - (y')^2}} \right) \quad (\text{Eq. 17})$$

and $\alpha + \beta = \frac{\pi}{2} - \theta$ (Eq. 18)

where $\beta = \tan^{-1} \left(\frac{v_y}{v_x} \right)$ (Eq. 19)

and v_x and v_y are the components of particle velocity just before impact with the target. Substituting the value of α obtained from Eq. 18 into Eq. 16 gives

$$\gamma = \pi - 2\theta - \beta \quad (\text{Eq. 20})$$

Substituting Eqs. 17 and 19 into Eq. 20 gives an equation of γ suitable for computation. Thus:

$$\gamma = \pi - 2 \tan^{-1} \left(\frac{y'}{\sqrt{1 - y'^2}} \right) - \tan^{-1} \left(\frac{v_y}{v_x} \right) \quad (\text{Eq. 21})$$

or in non-dimensional terms

$$\gamma = \pi - 2 \tan^{-1} \left(\frac{\bar{y}'}{\sqrt{1 - \bar{y}'^2}} \right) - \tan^{-1} \left(\frac{\bar{v}_y}{\bar{v}_x} \right) \quad (\text{Eq. 21a})$$

The component of velocity in the direction of free field flow immediately after particle reflection is

$$v_x' = v \cos \gamma \quad (\text{Eq. 22})$$

where $v = \sqrt{v_x^2 + v_y^2}$ (Eq. 23)

or in non-dimensional terms $\bar{v}_x' = \bar{v} \cos \gamma$ (Eq. 22a)

and $\bar{v} = \sqrt{\bar{v}_x^2 + \bar{v}_y^2}$ (Eq. 23a)

The total change of the component of velocity in the direction of free stream flow due to reflection from the target is then given by the following relationship:

$$\Delta v_x = v_x - v \cos \gamma \quad (\text{Eq. 24})$$

6. THE DRAG FORCE DUE TO PARTICLES ALONE

The force F_p of a stream of particles hitting a cylinder or sphere is related to the change in momentum which in general is given by the following relationship:

$$F_p = m' \Delta v_x \quad (\text{Eq. 25})$$

where m' is the total mass of particles having its momentum changed per unit time and Δv_x is its change of velocity in the free stream flow direction.

The quantity m' can be calculated using the following equation:

$$m' = \rho_p A_d v_{x0} \quad (\text{Eq. 26})$$

where ρ_p is the uniform particle density per unit volume of air far upstream, A_d is the far upstream cross sectional area of the envelope of particles which eventually hits the sphere or cylinder, and v_{x0} is the far upstream particle velocity which is assumed equal to the fluid velocity. Combining Eqs. 25 and 26 gives

$$F = \rho_p A_d v_{x0} \Delta v_x \quad (\text{Eq. 27})$$

Δv_x was not constant but was shown in the previous section to be a function of the off-axis, upstream starting location of an individual particle. This can be accounted for by integrating the velocity change over the range of starting locations across the particle envelope. Then the equation of force becomes

$$F = \rho_p v_{x0} \int_0^{A_d} \Delta v_x dA \quad (\text{Eq. 28})$$

where dA is an element of cross sectional area at location y .

Assuming that the particles do not interact with each other, the equation of drag force due to particles alone is then obtained in usable form by substituting Eq. 24 into Eq. 28 which gives the following equation:

$$F_p = \rho_p v_{x0} \int_0^{A_d} (v_x - v \cos \gamma) dA \quad (\text{Eq. 29})$$

7. THE DRAG COEFFICIENT

(a) Cylinder

For a transverse cylinder of unit length, Eq. 29 becomes

$$F_s = 2\rho_p v_{x0} \int_0^{y_c} (v_x - v \cos \gamma) dy \quad (\text{Eq. 30})$$

In non-dimensional terms this is written

$$\frac{F_s}{\rho_p U^2 L} = 2\bar{v}_{x0} \int_0^{\bar{y}_d} (\bar{v}_x - \bar{v} \cos \gamma) d\bar{y} \quad (\text{Eq. 31})$$

The drag coefficient is defined as follows:

$$C_D = \frac{p}{\frac{1}{2} \rho_p U^2} \quad (\text{Eq. 32})$$

where p is the average drag pressure (Schlichting, 1960).

The average drag pressure for a cylinder of unit length is given by the following equation:

$$p_c = \frac{F}{2L} \quad (\text{Eq. 33})$$

Substituting Eq. 33 into Eq. 32 and the resulting equation into Eq. 32 gives

$$C_{DC} = 2\bar{v}_{x0} \int_0^{\bar{y}_d} (\bar{v}_x - \bar{v} \cos \gamma) d\bar{y} \quad (\text{Eq. 34})$$

b) Sphere

For a sphere Eq. 29 becomes

$$F_s = 2\pi\rho_p v_{x0} \int_0^{y_c} (v_x - v \cos \gamma) y dy \quad (\text{Eq. 35})$$

In non-dimensional terms this is written

$$\frac{F_s}{\rho_p U^2 L^2} = 2\pi\bar{v}_{x0} \int_0^{\bar{y}_d} (\bar{v}_x - \bar{v} \cos \gamma) \bar{y} d\bar{y} \quad (\text{Eq. 36})$$

The drag pressure on the sphere is given by

$$p_s = \frac{F_s}{\pi L^2} \quad (\text{Eq. 37})$$

Substituting Eq. 36 into Eq. 37 and the resulting equation into Eq. 32 gives the drag coefficient for a sphere as follows:

$$C_{Ds} = 4\bar{v}_{x0} \int_0^{\bar{y}_d} (\bar{v}_x - \bar{v}' \cos \gamma) \bar{y} \, d\bar{y} \quad (\text{Eq. 38})$$

8. INTEGRATION METHOD FOR THE DRAG COEFFICIENT EQUATIONS

a) Cylinder

$$\text{Let } f(\bar{y}) = \bar{v}_x - \bar{v} \cos \gamma \quad (\text{Eq. 39})$$

Then Eq. 33 can be written

$$C_D = 2\bar{v}_{x0} \int_0^{\bar{y}_d} f(\bar{y}) \, d\bar{y} \quad (\text{Eq. 40})$$

$\int_0^{\bar{y}_d} f(\bar{y}) \, d\bar{y}$ was integrated numerically by dividing the upstream

cross sectional area into strips of equal width Δy , and calculating the contribution from each strip. These contributions were then summed as follows:

$$\int_0^{\bar{y}_d} f(\bar{y}) \, d\bar{y} = \frac{f(\bar{y}_1) + f(\bar{y}_n)}{2} + \sum_{i=2}^{i=n-1} f(\bar{y}_i) \Delta \bar{y} \quad (\text{Eq. 41})$$

where y_i increases with i from $\bar{y}' = 0$ to $\bar{y}_n = \bar{y}_d$.

b) Sphere

Using Eq. 39 in Eq. 38 gives the following equation for the drag coefficient of a sphere

$$C_{Ds} = 4\bar{v}_{x0} \int_0^{\bar{y}_d} f(\bar{y}) \bar{y} \, d\bar{y} \quad (\text{Eq. 42})$$

$\int_0^{\bar{y}_d} f(\bar{y}) \bar{y} d\bar{y}$ was integrated numerically by dividing the upstream

cross sectional area into concentric annuli of equal width $\Delta\bar{y}$ and calculating the contribution from each ring. These contributions were then summed as follows:

$$\int_0^{\bar{y}_d} f(\bar{y}) d\bar{y} = \sum_{i=1}^n f(\bar{y}_i) \bar{y}_i \Delta\bar{y} \quad (\text{Eq. 43})$$

where y_i increases with i from $\bar{y}_1 = \Delta\bar{y}$ to $\bar{y}_n = \bar{y}_d$

9. SOLUTION OF THE DRAG COEFFICIENT EQUATIONS

To solve Eq. 34 or Eq. 38 for the drag coefficient it was necessary to know the values of \bar{v}_x , \bar{v}' and γ in the range $0 \leq \bar{y} \leq \bar{y}_d$.

The first step was to find the value of \bar{y}_d which was the upper limit of the integral. This was done by an iterative procedure called the half interval method (Carnahan et al., 1969). The value of \bar{y} for the critical particle was estimated far upstream and the path followed to the target. The difference between the ordinate of the target path and the ordinate of the point on the actual path parallel to the tangent path was the miss criterion used. The direction of the tangent path was not known a priori but was assumed parallel to the actual path. The half interval method previously mentioned was applied to determine a better initial estimate. Then the path was followed to the target again for another calculation of miss distance. This process was repeated several times until sufficient accuracy was achieved. The plane of initial position of particles which was perpendicular to the flow direction was located far enough from the target so that free stream conditions prevailed. A distance of five target radii upstream of the target centre was considered adequate (Batchelor, 1967).

The path of an individual particle was determined step by step by applying a fourth order Runge-Kutta method (Carnahan et al., 1969) to

the equations of motion (Eqs. 1 and 2). The values of Re and K in these equations were easily determined for each new step by direct substitution of previously determined values into Eqs. 3, 4 and 5, but the value of $C_D Re$ in Eqs. 1 and 2 had to be calculated in each step by numerical solution of the definitive empirical equations (Eqs. 6 and 7). This was done using Newton's Method (Carnahan et al., 1969) for finding the zeros of a function. The values of \bar{u}_x and \bar{v}_y in Eqs. 1, 2 and 3 were calculated from Eqs. 8 and 9 for a cylindrical target and Eqs. 10, 11, 12 and 13 for a spherical target.

Once the value of \bar{y}_d had been determined the integration procedures described in Section 8 of this report were applied. The values of $f(\bar{y}_i)$ used in Eqs. 41 and 43 were obtained from Eq. 39. The quantities \bar{v} and γ used in the latter equation were calculated from Eqs. 23a and 21a respectively. The values of \bar{v}_x , \bar{v}_y and \bar{y}' used in the latter two equations were calculated by following a particle from its initial position to the cylindrical or spherical target using the previously described step by step procedure applied in determining the value of \bar{y}_d .

10. IMPACTION EFFICIENCY

The impaction efficiency was defined as the ratio of the cross sectional area of the far upstream envelope of particles which eventually hit the target to the cross sectional area of the target itself. For a transverse cylindrical target this is simply given by the value of \bar{y}_d . For a spherical target it is given by \bar{y}_d^2 .

11. COMPUTER PROGRAMS

Computer programs were written in Fortran IV for the DRES IBM 1130 computer to obtain solutions to the drag coefficient equations by the method described in Section 9 of this report. The complete programs, together with one set of results for each, are shown in Appendix A for the cylinder in free space and in Appendix B for the sphere in a tube or in free space (Appendix B). These programs are annotated so that the

functions of their various parts can be understood without further description. The application of the method described in Section 9, which was used in the programs, is straight forward with the exception of the step by step integration procedure near the target sphere or cylinder. This procedure is therefore explained as follows.

The particle motion was calculated in time steps, ΔT , until the particle was just one time step away from the target. Then the time increment size was decreased by a factor of ten and step by step integration continued until the particle was one new time step away from the target. At this point the time increment size was decreased by another factor of ten and the integration allowed to proceed until the target was reached. This method ensured that the position of the target and particle coincided within a maximum error given by the distance travelled during one percent of the original time step size, while allowing the particle to reach the proximity of the target in an adequate but small number of steps.

12. RESULTS

Drag coefficients and impaction efficiencies were calculated by means of the computer programs (Appendix A and B) from two main points of view. First, they were calculated using the particle sizes, fluid velocities and target configurations used in experiments done at Suffield many years ago (Gillespie and Gunter, 1957; 1959). This was done so that a comparison of the results from the theoretical calculation of drag due to particles in air could be made with existing information obtained from experiments. The results were obtained for a narrow range of particle sizes and fluid velocities and, correspondingly, a narrow range of inertia parameters and free stream particle Reynolds numbers. Next, calculations of impaction efficiencies and drag coefficients were done for a wide range of inertia parameters. The impaction efficiencies were compared with the results of other workers (Friedlander, 1977) but no results were available for comparison to the calculated drag coefficients over this

wide range of the inertia parameter. All the results previously mentioned are shown in Tables 1 and 2 and Fig. 4 - Fig. 9, along with plots of the calculated distribution of forces due to elastic reflection of particles from a cylinder (Fig. 10) and a sphere (Fig. 11). The latter two figures were plotted from the sample computer results shown with their corresponding programs in Appendix A and B. Further details explaining the significance of the tables and figures are described as follows.

The calculated results for the various particle sizes and fluid velocities used in the experiments done at Suffield (Gillespie and Gunter, 1957; 1959) are shown for a cylinder (Table 1) and for a sphere (Table 2). The sizes of these two targets also correspond to those in the experiments. The calculations were done for a 0.2 centimetre diameter cylinder in free space and a 0.8 centimetre sphere in a 2.54 centimetre tube. The tube which was used in all the experiments was included for the sphere since the velocity potential was readily available (Smythe, 1964). However, the effect of the presence of the tube on the potential flow and the final theoretical results was found to be negligible. Similarly, the effect of not including the tube in the potential flow for the cylinder is expected to be negligible.

A reason for tabulating the calculated results is to show clearly that the effect of particle size on the calculated drag coefficient for the sizes used in the experiments was negligibly small. This conclusion was also drawn from the experiments reported by Gillespie and Gunter (1957; 1959). The drag coefficients due to the middle size particles alone were plotted against target Reynolds number for the cylinder (Fig. 4) and for the sphere (Fig. 5). Also shown for the purpose of comparison of experiment to theory, described in the following section, are the drag coefficients due to air alone.

The impaction efficiencies for a wide range of the inertia parameter were plotted for cylinders (Fig. 6) and for spheres (Fig. 8). The cylinder results were plotted for $\phi = 1000$. The definition of ϕ

and the reason for its value choice are described as follows. The collection efficiency and also the drag coefficient are a function of two dimensionless groups, the inertia parameter, K , and the free stream particle Reynolds number, Re_0 . A new dimensionless group

$$\phi = \frac{Re_0^2}{K} \quad (\text{Eq. 44})$$

independent of particle size was introduced. According to the rules of dimensionless analysis this is permissible, but the efficiency is still determined by two groups which for convenience were chosen to be K and ϕ . The collection efficiencies were calculated by means of a mechanical analog and plotted for values of ϕ by other workers (Brun et al., 1953). The results for the middle value of ϕ from this work are also shown (Fig. 6) for comparison. The sphere results were plotted (Fig. 8) for an Re_0 value of 128. Again this choice was made for convenience of comparison to other work. The choice was the middle value of seven for which curves calculated from inviscid flow theory by Dorsch, Saper and Kadow are shown (Friedlander, 1977). The curve is also shown in Fig. 8.

The calculated drag coefficient due to particles alone were plotted for a wide range of inertia parameters for cylinders (Fig. 7) and for spheres (Fig. 9). In these two figures each curve was fitted by eye through several data points obtained by applying the computer programs (Appendix A and B) to various particle diameters for the cylinder case and various target diameters for the sphere case.

10. COMPARISON OF THEORETICAL RESULTS TO EXPERIMENTAL DATA

The experimental work previously done at Suffield (Gillespie and Gunter, 1957; 1959) indicated that the drag force of a sphere or cylinder is given by the following equation:

$$F = 1/2 \rho^* U^2 C_{D_e} A \quad (\text{Eq. 45})$$

where $\rho^* = \rho + \rho_p \quad (\text{Eq. 46})$

and generally ρ and ρ_p are unequal. For Eq. 45 to be true, the drag coefficients due to particles alone carried by air and due to air alone must be equal. This is shown as follows. The total drag force consisted of two parts expressed by the following equation:

$$F_p + F_a = 1/2\rho_p U^2 A C_{Dp} + 1/2\rho U^2 A C_D \quad (\text{Eq. 47})$$

Comparison of Eq. 45 and 46 to Eq. 47 shows that

$$(\rho + \rho_p) C_{De} = \rho_p C_{Dp} + \rho C_D \quad (\text{Eq. 48})$$

C_{Dp} and C_D must be equal for Eq. 48 to be true.

The results obtained herein (Tables 1 and 2) and illustrated in Fig. 4 and 5 show that the calculated drag coefficients due to particles carried by air are considerably greater than the established drag coefficients due to air alone (Schlichting, 1960). The discrepancies between the experimental and calculated drag coefficients were accounted for by showing that the velocity of the particles in the experiments were considerably lower than the mean air velocity.

The particles which were introduced by Gillespie and Gunter into the air stream with zero velocity in the direction of air flow were accelerated along a horizontal section of 2.54 cm tube, 17 cm long, and a vertical section, 78 cm long, giving a total tube length of 95 cm between the starting point and the test section. Calculations of velocity as a function of distance for particles accelerated from rest were done to simulate Gillespie's experiments (Mellisen, in draft). The results for a travelled distance of 95 cm for the various mean air velocities and number median diameter used in the experiments are shown in Table 3. Also, the curves of velocity versus distance for the three diameters and one air velocity are illustrated in Fig. 14.

The particle velocities were difficult to calculate because the behavior of the particles travelling in the air stream around the corner in the tube was very complex. The velocity of the air in a straight cross section of the tube was zero at the walls and increased to the centre to 1.22 to 1.25 times the mean velocity (Prandtl and Tietjens,

1957). Therefore, the particle velocities were affected by the location of the particles in the tube cross section. This could not be accounted for because the time history of the location of the particles was not known. As a simplifying approximation, the mean air velocity was assumed to act on the particles over their length of travel in the tube which was also assumed vertical over its entire length. The experimental results indicate that the variation in velocities was less than this method of calculation shows. A mean value of drag force due to the presence of all particle sizes was used to offset this.

The drag coefficients which were calculated for particles with a free stream velocity equal to the mean air velocity (Table 1 and 2) were then adjusted to simulate the special conditions of the experiments as follows. The drag coefficients in the experiments were based on the mean air velocities as measured by a rotameter. The drag coefficients due to the particles travelling at lower velocities were adjusted to the mean air velocity by the following relationship for flow over the cylinder.

$$C_{DpU} = C_{Dc} \left(\frac{v_{rms}}{U} \right)^2 \quad (\text{Eq. 49})$$

This method was justified because, for the purpose of comparison, the calculated drag coefficients over the range of velocities and particle sizes used were sufficiently near a constant value (Table 1). All three particle sizes were assumed present in equal quantities in the air flow. Then, since the drag forces vary as the square of the velocity, the root mean square particle velocity, v_{rms} , was calculated for each mean fluid velocity (Table 3) and used in Eq. 49 to obtain a value for the drag coefficient due to dust alone at each mean air velocity (Table 4). The overall drag coefficients due to air and particles were then calculated for the dust concentrations used in the experiments, namely 2 and 5 kilograms per cubic metre of space. This was done by solving Eq. 48 for C_{De} . That is

$$C_{De} = \frac{\rho_p C_{Dc} + \rho C_D}{\rho + \rho_p} \quad (\text{Eq. 50})$$

The results are shown for each mean air velocity (Table 4) and plotted along with the drag coefficient for air alone (Fig. 12) as was done for the experiments of Gillespie and Gunter (1957; 1959).

A similar method was used to obtain the overall drag coefficients for the sphere with the exception that, as for the experimental results, they were corrected for the presence of the tube walls by the Landenburg correction (Gillespie and Gunter, 1957). That is, the drag coefficients due to particles alone were divided by the factor $1 + 2.4 L/a$ to obtain the corrected drag coefficient, C_{DL} . For the sphere and tube radii used, the factor is 1.756. Then, the overall drag coefficient for the sphere was obtained from

$$C_{De} = \frac{\rho_p C_{DL} + \rho C_D}{\rho + \rho_p} \quad (\text{Eq. 51})$$

The results are shown for each mean air velocity in Table 4, and plotted along with the drag coefficient for air alone in Fig. 13.

14. DISCUSSION

The equations of motion (Eq. 1 and 2) have never been proven experimentally (Friedlander, 1977) although they have been used to calculate deposition efficiencies for cylinders (Glauert, 1940). The mathematical model described herein used these equations to calculate collection efficiencies for cylinders which compare favorably (Fig. 6) with the results obtained using a differential analyzer (Brun et al., 1953). Also, the impaction efficiencies for spheres (Fig. 8) agree with solutions to the equations obtained by other workers (Friedlander, 1977).

In the work reported herein, the mathematical model for calculating impaction efficiency was extended to obtain the impaction forces due to particles in air flowing over spheres and cylinders. These results agree favorably with experiments previously done at Suffield (Gillespie and Gunter, 1957) when adjusted to the special conditions of the experiment. The experiments do support the theory, but the theoretical results

show that the results of the experiments are not applicable to the free field case of drag due to dusty air flowing over a cylinder or sphere.

The model could also be extended to blast loading problems. Since the effect of compressibility becomes important in these problems it would have to be accounted for. This could be done in the mathematical model by changing the air flow field equations to include it. The model could also be used to calculate the actual distributions of impaction velocities on cylinders (Fig. 10) and spheres (Fig. 11) in special cases which may have applications in chemical and nuclear defence.

15. CONCLUSIONS

A mathematical model was developed for calculating the impaction forces due to airborne particles on spheres and cylinders. The model can also be used for calculating impaction efficiencies, velocity distributions, and points of impact for a stream of airborne particles flowing over a sphere or cylinder. Hence, the drag coefficient for the particles on the sphere or cylinder can be determined.

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TABLE 1

COEFFICIENT OF DRAG ON A 0.2 CM CYLINDER FOR VARIOUS PARTICLE SIZES

	PARTICLE SIZE			AIR VELOCITY	P	Re ₀ ² /K
	470μm	155μm	55μm	cm sec ⁻¹		
Coefficient	2.367	2.362	2.333	700	3.99	
of Drag due	2.367	2.361	2.327	560	3.19	
to Particles	2.367	2.360	2.318	435	2.48	
alone, C _{Dc}	2.367	2.357	2.300	294	1.67	

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TABLE 2

COEFFICIENT OF DRAG ON A 0.8 CM SPHERE IN A 2.54 CM TUBE
FOR VARIOUS PARTICLE SIZES USING POTENTIAL FLOW

	PARTICLE SIZE			AIR VELOCITY cm sec ⁻¹	Re ₀ For 155μm
	470μm	155μm	55μm		
Coefficient	2.048	2.036	1.970	700	73.6
of Drag due	2.047	2.033	1.955	560	58.9
to Particles	2.047	2.030	1.934	426	44.8
alone, C _{Ds}	2.046	2.023	1.887	272	28.6

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TABLE 3
VELOCITY OF PARTICLES AT TARGET POSITION
ACCELERATED VERTICALLY UPWARD IN A TUBE FROM REST
95 cm UPSTREAM

<u>PARTICLE</u> <u>DIAMETER</u> cm	<u>MEAN AIR</u> <u>VELOCITY IN</u> TUBE cm sec ⁻¹	<u>PARTICLE VELOCITY</u> <u>AT TARGET</u> POSITION cm sec ⁻¹	<u>RMS PARTICLE</u> <u>VELOCITY AT TARGET</u> POSITION cm sec ⁻¹
0.0470	700	299	
0.0155	700	572	534
0.0055	700	664*	
0.0470	560	191	
0.0155	560	445	398
0.0055	560	527*	
0.0470	435	75*	
0.0155	435	325*	302
0.0055	435	405*	
0.0470	426	66	
0.0155	426	317	295
0.0055	426	396*	
0.0470	294	-	
0.0155	294	187*	188
0.0055	294	267*	
0.0470	272	-	
0.0155	272	166*	171
0.0055	272	245*	

* Upward terminal velocity reached

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TABLE 4

DRAG COEFFICIENT OF A 0.2 CM CYLINDER IN A 2.54 CM TUBE DUE TO AIR
CONTAINING EQUAL MASSES OF 0.0470, 0.0155 AND 0.0055 CM PARTICLES
CORRECTED FOR PARTICLE VELOCITY DEFICIENCY

U cm sec ⁻¹	C _{Dc}	C _D	U _{rms}	C _{DpU}	C _{De}		Re
					2 kg m ⁻³ Dust	AIR + DUST 5 kg m ⁻³ Dust	
700	2.354	1.00	534	1.369	1.23	1.30	949
560	2.352	1.05	398	1.188	1.14	1.16	759
435	2.348	1.15	302	1.132	1.14	1.14	590
294	2.341	1.25	188	0.957	1.07	1.01	399

TABLE 5

DRAG COEFFICIENT OF A 0.8 cm SPHERE IN A 2.54 cm TUBE DUE TO AIR CONTAINING
EQUAL MASSES OF 0.0470, 0.155 AND 0.055 cm PARTICLES CORRECTED FOR
REAL FLUID VELOCITY AND PARTICLE VELOCITY DEFICIENCY

U cm sec ⁻¹	C _{Ds}	C _D	U _{rms}	C _{Dpu}	C _{DL}	C _{De}		Re
						AIR	AIR + DUST	
700	2.018	0.391	534	1.174	0.669	0.564	0.615	3796
560	2.012	0.396	398	1.016	0.579	0.510	0.543	3076
426	2.004	0.406	295	0.961	0.547	0.494	0.519	2310
272	1.985	0.433	171	0.785	0.447	0.442	0.444	1475

APPENDIX A

COMPUTER PROGRAM FOR A CYLINDER

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PAGE 1

// JOB 1

LOG DRIVE CART SPEC CART AVAIL PHY DRIVE
0000 0205 0205 0000

V2 '11 ACTUAL 16K CONFIG 16K

// EJCT

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PAGE 2

// FOR
*ONE WORD INTEGERS
*LIST ALL
SUBROUTINE SBM2RIG,TAU,DTAM,XK,REZ,UX,UY,XCDRE)

C THIS SUBROUTINE CALCULATES PARTICLE MOTION DURING THE FINAL
C INCREMENT OF TAU

C DIMENSION G(4),DG(4)
C M=0

C CALL ON RUNGE KUTTA SUBROUTINE

C R CONTINUE
C M=M+1

C CALL SBM22(4,G,DG,TAU,DTAM,IRUNG,M)

C IF(IRUNG-1)GO,9,10

9 PF=REZ*(UY-G(2))**2+(UX-G(1))**2)**0.5

C XCDRE=CDRE(RE)

C DG(1)=(XCDRE/(24.0*XK))*(UX-G(1))

C DG(2)=(XCDRE/(24.0*XK))*(UY-G(2))

C DG(3)=G(1)

C DG(4)=G(2)

C GO TO R

C 10 CONTINUE

C X=C

C CALCULATE FLUID VELOCITY AT PARTICLE POSITION

C RSQSQ=(G(3)**2+G(4)**2)**2

C UX=1.0-I(G(3)**2-G(4)**2)/RSQSQ

C UY=-(G(3)*G(4))/RSQSQ

C INTEGRATE ACROSS ANOTHER STEP IF REQUIRED

C IF(G(4) - 1.0)11,11,12

11 DELX = SORT(1.0 - G(4)**2)

C GO TO 13

12 DELX = 0.0

13 HITS=G(3)+G(1)*1.0*DTAM+DELX

C IF(HITS)8,8,18

C 18 CONTINUE

C RETURN

C END

VARIABLE ALLOCATIONS

DG(1) = 0006-0000

PE(1) = 0008

RSQSQ(1) = 000A

DELX(1) = 000C

HITS(1) = 000E

M(1) = 0012

PAGE 3

IRING(I)=0013

STATEMENT ALLOCATIONS

R =0055 9 =C06A 1G =00C0 11 =0104 12 =0115 13 =0119 18 =012E

FEATURES SUPPORTED
ONE WORD INTEGERS

CALLED SUBPROGRAMS

SRM22 CORE FSORT FAXB FADD FADDX FSUB FSUBX FMPY FMPYX FDIV FLD FLDX FSTD FSTOX
FSRR FDVR FAXI SNR SUBIN

REAL CONSTANTS

.500000F 00=0018 .240000E 02=001A .100000E 01=001C .000000E 00=001E .110000E 01=002D

INTEGER CONSTANTS

0=0072 1=0023 4=0024 2=0025

CORE REQUIREMENTS FOR SRM28
COMMON 0 VARIABLES 24 PROGRAM 280

RELATIVE ENTRY POINT ADDRESS IS 0026 (HEX)

END OF COMPILATION

// DUP

*STORE MS UA SRM28

CART ID 0205 DB ADDR 5F48 DB CNT 0015

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PAGE 4

// FOR
*ONE WORD INTEGERS
*LIST ALL

C FUNCTION CDRE(RE)

C THIS FUNCTION COMPUTES THE PRODUCT OF DRAG COEFFICIENT
C AND REYNOLDS NUMBER FOR A SPHERE AS A FUNCTION OF
C REYNOLDS NUMBER

C CONSTANT COEFFICIENTS

A1=1./24.
A2=-2.3363*1.E-04
A3=2.0154*1.E-06
A4=-6.9105*1.E-09
R0=-1.29536
P1=9.86*1.E-01
P2=-4.6677*1.E-02
P3=1.1235*1.E-03

C CHOOSE THE APPROPRIATE POLYNOMIAL

C IF(IRE-4.012*7.7

C INITIAL ESTIMATE

2 IF(IRE - 0.00113*4.4

3 CDRE = 24.0

GO TO 30

4 X=24.*RE

C BEGIN NEWTON METHOD ITERATION

C CONTINUE

DO 6 ITER=1,20

FX=A1*X+A2*X**2+A3*X**3+A4*X**4-RE

FX=A1*X+A2*X**2+A3*X**3+A4*X**4-RE

DELX=FX/FPX

X=X-DFLX

C CHECK FOR CONVERGENCE

C EPS=1.E-06

IF(ABS(DELX/X)-EPS)5,5,6

5 CDRE=X/RE

PAGE 5

GO TO 30
6 CONTINUE
GO TO 29

C INITIAL ESTIMATE

C 7 CD = 1.0
FLOG = 0.434294481903252
X=ALOG(CD*RE**2)*ELOG

C RFGIN NEWTON METHOD ITERATION

C DO 24 ITER=1,20
FX=RO+R1*X+R2*X**2+R3*X**3 - ALOG(RE)*ELOG
FPX=R1+2.*R2*X+3.*R3*X**2
DELX=FX/FPX
X=X-DELX

C CHECK FOR CONVERGENCE

C FPS=1.E-06
IF(ABS(DELX/X)-EPS)22,22,24

22 CORE=10.**X/RE
GO TO 30
24 CONTINUE
29 WRITE(3,702)
30 RETURN

C FORMATS FOR OUTPUT STATEMENTS

C 202 FORMAT(16HC NO CONVERGENCE)

C END

C VARIABLE ALLOCATIONS

CORE(I) I=0000
R1(I) I=000C
DELX(I) I=001R

A2(I) I=0004
R3(I) I=0010
CD(I) I=001C

A3(I) I=0006
X(I) I=0012
ELOG(I) I=001E

A4(I) I=0008
FX(I) I=0014
ITER(I) I=0028

B0(I) I=000A
FPX(I) I=0016

STATEMENT ALLOCATIONS
202 =0057 2 =00A6 3 =00AD 4 =00R3 5 =00R3 5
30 =01CF

=012B 6 =0133 7 =013E 22 =0186 24 =01C1 29 =01CA

FEATURES SUPPORTED
ONE WORD INTEGERS

CALLED SUBPROGRAMS
FARS FALOG FAXR

FADD FSUR FMPY FDIV FLD FSTO FSBR FAXI SHRT SCOMP SNR SUBIN

PAGE 6

REAL CONSTANTS

.100000E 01=002A .240000E 02=002C .233630E 01=002E .100000E-03=0030 .201540E 01=0032 .100000E-05=0034
.691050E 01=0036 .100000E-08=0038 .129536E 01=003A .986000E 01=003C .100000E 03=003E .466770E 01=0040
.100000E-01=0042 .112350E 01=0044 .100000E-02=0046 .400000E 01=0048 .200000E 01=004A
.434794E 00=004E .100000E 02=0050

INTEGER CONSTANTS

1=0052 20=0053 2=0054 3=0055 4=0056

CORE REQUIREMENTS FOR CDRE

COMMON 0 VARIABLES 42 PROGRAM 424

RELATIVE ENTRY POINT ADDRESS IS 0061 (HEX)

END OF COMPILATION

// DUP

*STOP %S UA CDRE

CART ID 0205 DB ADDR 5F60 DB CNT 0022

// FJCT

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PAGE 7
// FOR
#0MF WORD INTEGERS
*LIST ALL
SUBROUTINE SRM22(N,Y,F,X,H,IRUNG,M)
C
C FOURTH ORDER RUNGE KUTTA METHOD
C FOR N FIRST ORDER O.D.E.
C
C DIMENSION PHI(50),SAVY(50),Y(50),F(50)
C GO TO (2,3,4,5,6),M
C
C PASS 1
C
C 2 IRUNG=1
C RETURN
C
C PASS 2
C
C 3 DO 72 J=1,N
C SAVY(J)=Y(J)
C PHI(J)=F(J)
C 22 Y(J)=SAVY(J)+0.5*H*F(J)
C X=X+0.5*H
C IRUNG=1
C RETURN
C
C PASS 3
C
C 4 DO 23 J=1,N
C PHI(J)=PHI(J)+2.0*H*F(J)
C 33 Y(J)=SAVY(J)+0.5*H*F(J)
C IRUNG=1
C RETURN
C
C PASS 4
C
C 5 DO 44 J=1,N
C PHI(J)=PHI(J)+2.0*H*F(J)
C 44 Y(J)=SAVY(J)+H*F(J)
C X=X+0.5*H
C IRUNG=1
C RETURN
C
C PASS 5
C
C 6 DO 55 J=1,N
C 55 Y(J) = SAVY(J) + (PHI(J) + F(J))*H/6.0
C IRUNG=2
C RETURN
C
C END
```

PAGE 8

VARIABLE ALLOCATIONS
PHIR)=0062-0000 SAVYR)=00C6-0064 JII)=00C8

STATEMENT ALLOCATIONS
2 =0105 3 =0108 22 =011E 4 =0141 33 =0152 5 =016D 44 =017E 6 =019F 55 =01A3

FEATURES SUPPORTED
ONE WORD INTEGERS

CALLED SUBPROGRAMS
FADD FADDX FMPY FMPYX FDIV FLD FLDX FSTO FSTOX SUBSC SUBIN

REAL CONSTANTS
.500000F 00=00CA .200000E 01=00CC .600000E 01=00CE

INTEGER CONSTANTS
1=0000 2=0001

CORE REQUIREMENTS FOR SRM22
COMMON 0 VARIABLES 202 PROGRAM 250

RELATIVE ENTRY POINT ADDRESS IS 00D2 (HEX)

END OF COMPILATION

// DUP

*STORE W5 UA SBM22
CART ID 0205 DB ADDR 5F82 DR CNT 0012

// EJECT

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PAGE 9
// FOR
*ONE WORD INTEGERS
*LIST ALL
SUBROUTINE SBM29(G4LFT,G4RIT,SIGNL,DTAU,NIBP,NSBP,NX,XL,XK,REZ,
1G4ZFR)
C
C THIS SUBROUTINE CALCULATES THE
C IMPACTION EFFICIENCY OF A CIRCULAR CYLINDER
C
C DIMENSION G(4),DG(4)
C WRITE(3,200)
C WRITE(3,201)G4LFT,G4RIT,SIGNL,DTAU,NIBP,NSBP,NX
C HALF INTERVAL ITERATION FOR INITIAL G4 VALUE
C DO 47 ITER=1,NX
C
C SET AND PRINT INITIAL CONDITIONS
C
C N=0
C NSTFP=0
C TAU=0.0
C G(3)=XL
C G4ZFR=(G4LFT+G4RIT)/2.0
C G(4)=G4ZER
C RSQSO=(G(3)**2+G(4)**2)**2
C UX=1.0-(G(3)**2-G(4)**2)/RSQSO
C UY=-2.0*(G(3)*G(4))/RSQSO
C G(1)=UX
C G(2)=UY
C RE=REZ*((UY-G(2))**2+(UX-G(1))**2)**0.5
C XCDRE=CDRE/RE
C IP=ITER/NIRP*NIRP
C IF(IP-ITER)5,7,5
C
C 5 CONTINUE
C IF(ITER-1)6,7,6
C 6 CONTINUE
C IF(ITER-NX)8,7,8
C 7 CONTINUE
C WRITE(3,205)
C WRITE(3,203)ITER,G4LFT,G4ZER,G4RIT,TAU,G(1),G(2),G(3),G(4),UX,UY,
1XCDRE
C
C CALL ON RUNGE KUTTA SUBROUTINE
C
C A CONTINUE
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PAGE 10

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4**+1
CALL SRMZ2(4,G,DG,TAU,DTAU,IRUNG,M)
IF(IRUNG-1)10,9,10
9 RF=REZ*((UY-G(2))**2+(UX-G(1))**2)**0.5
  XCDRE=CDRE(RF)
  DG(1)=(XCDRE/(24.0**X))*(UX-G(1))
  DG(2)=(XCDRE/(24.0**X))*(UY-G(2))
  DG(3)=G(1)
  DG(4)=G(2)
GO TO 8
10 CONTINUE
**=0
C
C
C
CALCULATE FLUID VELOCITY AT PARTICLE POSITION
RSOSQ=(G(3)**2+G(4)**2)**2
UX=1.0-(G(3)**2-G(4)**2)/RSOSQ
UY=-2.0*(G(3)*G(4))/RSOSQ
C
C
C
PRINT SOLUTIONS
IS = ITER/NIRP*NIRP
IF(IS-ITER)11,13,11
11 CONTINUE
IF(ITER-1)12,13,12
12 CONTINUE
IF(ITER-NX)16,13,16
13 CONTINUE
NSTP=NSTEP+1
IF(NSTEP-NSRP)16,14,16
14 CONTINUE
NSTEP=0
TAY = TAU+0.0001
WRITE(3,204)TAY,G(1),G(2),G(3),G(4),UX,UY,XCDRE
C
C
C
INTEGRATE ACROSS ANOTHER STEP IF REQUIRED
16 CONTINUE
IF(G(4)-1.0)17,17,18
17 DELX=SORI(1.0-G(4))**2)
GO TO 19
18 DELX=0.0
19 WITS=G(3)+G(1)*I.1=DTAU+DELX
  IF(WITS)9,8,20
20 CONTINUE

```

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PAGE 11

C CHANGE INCREMENT SIZE NEAR CYLINDER AND INTEGRATE FURTHER

C DTAV=DTAU/16.0

C CALL SPW28(G,TAU,DTAV,XX,REZ,UJ,XCDRE)

C PRINT SOLUTIONS

C IF(I5-ITER)21,23,21

21 CONTINUE

C IF(ITER-1)22,23,22

22 CONTINUE

C IF(ITER-NX)24,23,24

23 CONTINUE

C WRITE(3,204)TAU,G(1),G(2),G(3),G(4),UJ,UY,XCDRE

24 CONTINUE

C DTAV=DTAU/160.0

C CALL SPW28(G,TAU,DTAV,XX,REZ,UJ,XCDRE)

C PRINT SOLUTIONS

C IF(I5-ITER)31,33,31

31 CONTINUE

C IF(ITER-1)32,33,32

32 CONTINUE

C IF(ITER-NX)34,33,34

33 CONTINUE

C WRITE(3,204)TAU,G(1),G(2),G(3),G(4),UJ,UY,XCDRE

34 CONTINUE

C CALCULATE ORDINATE AT TANGENT POINT OF TANGENT PATH

C OPD = G(1)/SORT(G(1)**2 + G(2)**2)

C FIND INTERVAL HALF WITH THE SIGN CHANGE

C IF((G(4)-ORD)*SIGNL-0.0)45,45,46

45 G=RT=GZER

GO TO 47

46 G=LEFT=GZER

47 CONTINUE

EM = GZER

WRITE(3,207) F⁴

C RETURN

C

QUALITY IS UNACCEPTABLE
 QUALITY IS UNACCEPTABLE

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C FORMATS FOR OUTPUT STATEMENTS
 C

200 FORMAT(1H0, 35X, 44H IMPACTION EFFICIENCY OF A CIRCULAR CYLINDER/
 1 1H0)

201 FORMAT(10HG4LEF = ,F10.6/ 10H G4RIT = ,F10.6/10H SIGML = ,
 1 F3.0/ 10H TAU = ,F10.6/ 10H NIBP = ,I3/ 10H NSRP = ,I3/
 2 10H NX = ,I3)

203 FORMAT(10H0ITER = ,I3/ 10H G4LEF = ,F10.6/ 10H G4ZER = ,
 1 F10.6/ 10H G4RIT = ,F10.6/ 7H0 TAU, 11X, 4HG(1), 12X,
 2 4HG(2), 12X, 4HG(3), 12X, 4HG(4), 14X, 2HUX, 14X, 2HUY,
 3 12X, 4HCDRE /
 4 1H0, F7.4, 4F16.6, 3F16.4)

204 FORMAT (1H, F7.4, 4F16.6, 3F16.4)
 205 FORMAT (46HITHE MOTION OF A CRITICAL PARTICLE IS GIVEN BY)
 207 FORMAT(30H0THE IMPACTION EFFICIENCY IS ,E10.4)
 END

VARIABLE ALLOCATIONS

GPR)=C036-0000 DGIR)=000E-0008 TAUUR)=0010 RSO50(R)=0012 UYIR)=0016
 REIR)=0018 XCDREIR)=001A TAWIR)=001C DELX(R)=001E DTAMIR)=0022
 ORD(R)=0024 EM(R)=0026 IYER(I)=002C IYER(I)=002E JPI(I)=002F
 IYUNG(I)=0030 IS(I)=0031

STATEMENT ALLOCATIONS

200	=004D	201	=006R	203	=00A3	204	=00E8	205	=00F3	207	=010C	5	=020A	6	=0210	7	=0216	8	=02AE
C	=0253	1C	=02A9	11	=02FA	12	=0300	13	=0306	14	=0312	16	=0328	17	=03A1	18	=0352	19	=0356
20	=036R	21	=0381	22	=0387	23	=038D	24	=03A9	31	=03BF	32	=03C5	33	=03C8	34	=03E7	45	=040F
46	=0415	47	=0419																

FEATURES SUPPORTED
 ONE WORD INTEGERS

CALLED SUBPROGRAMS

CDRE	SRX22	FSORT	SEM28	FAXB	FADD	FADDX	FSUB	FSUBX	FMPY	FMPYX	FDIV	FLD	FLDX	FSTC
FSTOX	FSR22	FDR	FDRX	FAXI	SWRT	SCOMP	SIOFX	SIOF	SIOI	SNR	SUBIN			

REAL CONSTANTS

.00000E 00=0036 .20000E 01=0038 .100000E 01=003A .500000E 00=003C .240000E 02=003E .100000E-03=0040
 .110000F 01=0042 .100000E G2=0044 .100000E 03=0046

INTEGER CONSTANTS

3=0048 1=0049 0=004A 2=004B 4=004C

CORE REQUIREMENTS FOR SEM29

COMMON C VARIABLES 54 PROGRAM 1016

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RELATIVE ENTRY POINT ADDRESS IS 011E (HEX)

END OF COMPILATION

// DUP

*STORE MS UA SBN29
CART ID 0205 DB ADDR 5F94 DB CNT 0047

// EJECT

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PAGE 14
// FOR
*ONE WORD INTFGERS
*LIST ALL
SUBROUTINE SBM26(XL,GAZER,DTAU,XK,REZ,PIX,NZER,NJBP)
C
C THIS SUBROUTINE CALCULATES THE MOTION OF PARTICLES
C IN A FLUID STREAM MOVING TOWARD A CIRCULAR CYLINDER AND
C CALCULATES THE FORCE OF PARTICLE IMPACT ON THE CYLINDER
C
C DIMENSION G(4),DG(4)
C DIMENSION YZER(500)
C
C SFT NUMBER OF INCREMENTS AT INITIAL POSITION
WRITE(3,200)
WRITE(3,201)NZER,NJBP
WRITE(3,202)
NCE=NZER+1
DELY=GAZER/FLOAT(NZER)
FSUM=0.0
C
C STEPWISE INTEGRATION FOLLOWING PARTICLE POSITION
DO 40 ITER=1,NCE
C
C SET AND PRINT INITIAL CONDITIONS
M=0
TAU=C*G
G(3)=XL
YZER(ITER)=FLOAT(ITER-1)*DELY
GAZER=YZER(ITER)
G(4)=GAZER
RSOSQ=(G(3)**2+G(4)**2)**2
UX=1.0-(G(3)**2-G(4)**2)/RSOSQ
UY=-2.0*(G(3)*G(4))/RSOSQ
UXZER=UX
UYZER=UY
G(1)=UX
G(2)=UY
C
C CALL ON RUNGE KUTTA SUBROUTINE
C
C CONTINUE
N=N+1
CALL SBM22(4,G,DG,TAU,DTAU,IRUNG,N)
```

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IF (IRUNG-1)10,9,10
9 9F=PEZ*((UY-G(2))**2+(UX-G(1))**2)**0.5
XCDPF=CPRE(IRE)

SG(1)=XCDRE/(24.0*XK);*(UX-G(1))

DG(2)=XCDRE/(24.0*XK);*(UY-G(2))

DG(3)=G(1)

DG(4)=G(2)

GO TO A

10 CONTINUE

M=0

C CALCULATE FLUID VELOCITY AT PARTICLE POSITION

PSOSO=(G(3)**2+G(4)**2)**2

UX=1.0-(G(3)**2-G(4)**2)/PSOSO

UY=-2.0*(G(3)*G(4))/PSOSO

C INTEGRATE ACROSS ANOTHER STEP IF REQUIRED

IF(G(4) - 1.0)11,11,12

11 DELX = SORT(1.0 - G(4)**2)

GO TO 13

12 DELX = 0.0

13 HITS=G(3)+G(1)*1.1*DTAU+DELX

IF(HITS)R,R,1R

1P CONTINUE

C CHANGE INCREMENT SIZE NEAR CYLINDER AND INTEGRATE FURTHER

DTAW=DTAU/10.0

CALL SPW2BIG,TAU,DTAW,XK,REZ,UX,UY,XCDRE

DTAW=DTAU/100.0

CALL SPW2BIG,TAU,DTAW,XK,REZ,UX,UY,XCDRE

C CO-ORDINATES, VELOCITY, AND PRESSURE DERIVATIVE AT CYLINDER

VX=G(1)

VY=G(2)

X=G(3)

Y=G(4)

V=SORT(VX**2+VY**2)

PI = 3.14159265358979323846264

GAMMA = PI - 2.0*ATAN(Y/SORT(1.0-Y**2)) - ATAN(VY/VX)

FY= UXZER*(VX - V*COS(GAMMA))

IF(ITER-1)24,25,24

24 CONTINUE

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PAGE 16

IF(IITER-NCE)26,25,26
25 FSUM=FSUM+0.5*FY
GO TO 29
26 FSUM=FSUM+FY
28 CONTINUE

C PRINT SOLUTIONS
C
C

IS = IITER/NJBP*NJBP
IF(IIS-IITER)31,33,31
31 IF(IITER-1)32,33,32
32 IF(IITER-NCE)34,33,34
33 CONTINUE

TAU = TAU+0.0001
WRITE(3,203)TAW,VX,VY,X,Y,UX,UY,FY
34 CONTINUE
4C CONTINUE

C DRAG COEFFICIENT OF CYLINDER DUE TO PARTICLE IMPACT
C
C

PX=2.*FSUM*DELY
WRITE(3,204)PX
RETURN

C FORMATS FOR OUTPUT STATEMENTS
C
C

200 FORMAT(65H,THE MOTION OF THE PARTICLES AT THE CYLINDER POSITION I
IS GIVEN BY)

201 FORMAT(10H,OMZER = ,I3,20X, 10H,NJRP = ,I3)

202 FORMAT(7H0 TAU, 11X, 4H VX , 12X, 4H VY , 11X, 4H X , 12X,
1 4H Y , 14X, 2H,UX, 14X, 2H,UY , 14X, 4H FY /1H)

203 FORMAT (F7.4, 4F16.6, 2F16.4, F16.5)

204 FORMAT(77H,THE DRAG COEFFICIENT OF THE CYLINDER DUE TO THE REFLEC
TION OF PARTICLES IS , F9.4)

END

VARIABLE ALLOCATIONS

GIR I=0006-0000	DGIR I=000E-0008	YZERIR I=03F6-0010	DELY(I) I=03FB	FSUM(I) I=03FA	TAU(I) I=03FC
RSQ50(I) I=03FE	UX(I) I=0400	UY(I) I=0402	UTZERIR I=0404	RE(I) I=0406	XCDRE(I) I=0408
DFLX(I) I=040A	HITS(I) I=040C	DTA(I) I=040E	VX(I) I=0410	VY(I) I=0412	X(I) I=0414
Y(I) I=0416	V(I) I=0418	PI(I) I=041A	GAMMA(I) I=041C	FY(I) I=041E	TAU(I) I=0420
NCE(I) I=042A	IITER(I) I=042B	M(I) I=042C	IRUNGI(I) I=042D	IS(I) I=042E	

STATEMENT ALLOCATIONS

200 =0440 201 =047C 202 =0480 203 =04A3 204 =04AA 8 =05A3 9 =0588 10 =060E 11 =0658 12 =06F9
13 =066D 18 =0682 24 =0707 25 =070D 26 =0717 29 =071D 31 =072D 32 =0733 33 =0739 34 =0753
4C =0753

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FEATURES SUPPORTED
ONE WORD INTEGERS

CALLER SUBPROGRAMS

SMZ2	CORE	FSORT	SMZ8	FATAN	FCGS	FAXB	FADD	FADDX	FSUB	FSUBX	FMPYX	FDIV	FLD
FLDY	FSTC	FSTOX	FSBR	FDVR	FAXI	FLOAT	SWRT	SCOMP	SIOF	SIOI	SNR	SUBIN	

REAL CONSTANTS

*100000E 00=0434	*100000E 01=0436	*200000E 01=0438	*500000E 00=043A	*2*0000E 02=043C	*110000E 01=043E
*10000E 02=0440	*100000E 03=0442	*31*159E 01=0444	*100000E-09=0446		

INTEGER CONSTANTS

3=044F	1=0449	0=044A	2=044B	4=044C
--------	--------	--------	--------	--------

CORE REQUIREMENTS FOR SMZ26

COMMON C VARIABLES 1076 PROGRAM 824

RELATIVE ENTRY POINT ADDRESS IS 0404 (HEX)

END OF COMPILATION

// DUP

*STORE 45 UA SMZ6
CART ID 0205 DR ADDR SFDR DB CNT 003B

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PAGE 18
// FOR
*ONF WORD INTEGERS
*IOCS(CARD)
*ICCS(1132 PRINTER)
*LIST ALL
C
C ESTABLISH PHYSICAL PROPERTIES FOR CALCULATING IMPACTION EFFICIENCY
C
C ESTABLISH PARAMETERS FOR STEP BY STEP INTEGRATION AND HALF
C INTERVAL METHOD
C
C GALT AND GART ARE LOWER AND UPPER ROUNDS RESPECTIVELY FOR
C STARTING LOCATION OF PARTICLE
C SIGNAL IS SIGN CORRECTION FOR HALF INTERVAL METHOD
C IN THIS PROBLEM IT IS -1.0
C DTAU IS THE TIME INCREMENT FOR STEP BY STEP INTEGRATION
C NSRP IS INTERVAL OF WRITTEN PARTICLE PATHS
C *X IS NUMBER OF ITERATIONS USED TO DETERMINE PARTICLE STARTING
C POINT
C 1 READ(2,100)GALT,GART,SIGNL,DTAU,NIRP,NSBP,*X
C
C ESTABLISH PHYSICAL PROPERTIES FOR CALCULATING IMPACTION EFFICIENCY
C
C DC IS CYLINDER DIAMETER, CM
C DP IS PARTICLE DIAMETER, CM
C RHO IS FLUID DENSITY, GM/CC
C SIGMA IS PARTICLE DENSITY, GM/CC
C XMU IS ABSOLUTE VISCOSITY OF FLUID, POISE
C UF IS FREE STEAM VELOCITY, CM/SEC
C
C REAR(2,101)DC,DP,RHO,SIGMA,XMU,UF
C
C XL IS UPSTREAM STARTING POSITION FOR CALCULATING PARTICLE MOTION
C
C READ(2,1,2)XL
C REZ=RHO*DP*UF/XMU
C XK=SIGMA*DP**2*UF/(9.*XMU*DC)
C P = 9.*RHO**2*UF*DC/(XMU*SIGMA)
C WRITE(3,200)
C WRITE(3,201)
C
C
C XK IS INERTIAL PARAMETER CALLED STOKES NUMBER
C REZ IS FREE STEAM REYNOLDS NUMBER OF SPHERICAL PARTICLE
C P IS DIMENSIONLESS GROUP INDEPENDENT OF DROP SIZE
C
C WRITE(3,202)XL,REZ,XK,P,DC,DP,RHO,SIGMA,XMU,UF
```

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C CALCULATE STARTING POSITION OF TANGENT PARTICLE

C CALL SRW29(G4LFT,G4RIT,SIGNL,DTAU,NIBP,N5BP,NX,XL,XK,REZ,G4ZER)

C ESTABLISH NUMBER OF INCREMENTS FOR INTEGRATING PARTICLE FORCE

C NZER IS NUMFR OF INCREMENTS

C NJRP IS INTERVAL OF WRITTEN INCREMENTS

C READ(2,103)NZER,NJRP

C CALCULATE FORCE PRODUCED BY PARTICLES

C CALL SRW26(XL,G4ZER,DTAU,XK,REZ,PX,NZER,NJBP)

C CDAIR IS DRAG COEFFICIENT OF CYLINDER IN FREE AIR

C READ(2,104) CDAIR

C RECYL=RHO*DC*UF/XMU

C WRITE(3,204)RECYL

C WRITE(3,206)CDAIR

C READ(2,111)NSTOP

C IF(NSTOP)1,30,30

C 30 CALL EXIT

C FORMATS FOR INPUT AND OUTPUT STATEMENTS

100 FORMAT(10X,F10.7,20X,F10.7,19X,F3.0/ 10X,F10.7,20X,13,27X,I2/
1 10X,I2)

101 FORMAT(F10.5,F10.7,F10.6, F10.6,F10.7,F10.1)

102 FORMAT (F10.7)

103 FORMAT(I5,I5)

104 FORMAT(F10.6)

111 FORMAT(I5)

200 FORMAT(1H1, 34X, 51HTHE DRAG ON CYLINDERS IN A STREAM OF DUST-LAD

1EN AIR / 140)

201 FORMAT(28H0THE PHYSICAL PARAMETERS ARE)

202 FORMAT(10HXL = ,F10.7/10H REZ = ,E12.6/10H XK = ,

1 F12.6/10H P = ,E12.6/

2 10H0DC = , F10.5/10H DP = ,F10.7/10H RHO = ,F10.6/

3 10H SIGMA = , F10.6/10H XMU = ,F10.7/10H UF = ,F10.1)

204 FORMAT(32H0THE TARGET REYNOLDS NUMBER IS ,E10.4)

206 FORMAT(61H0THE DRAG COEFFICIENT OF THE CYLINDER DUE TO THE AIR AL

10NE IS ,F9.3)

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END

VARIABLE ALLOCATIONS

G4LFT(R) = 0000 GARIT(R) = 0002 SIGNAL(R) = 0004 DTAU(R) = 0006 XL(R) = 0008 DC(R) = 0008 DPIR = 000A
RHO(R) = 000C SIGMA(R) = 000E XMR(R) = 0010 UF(R) = 0012 XL(R) = 0014 REZIR = 0016
XK(R) = 0018 GAZER(R) = 001C PX(R) = 001E CDATR(R) = 0020 CDATR(R) = 0020 RECYLIR = 0022
NIRP(I) = 0026 NSBP(I) = 0027 NX(I) = 0028 NZER(I) = 0029 NSBP(I) = 002A NSTOP(I) = 002B

STATEMENT ALLOCATIONS

100 = 0040 ICI = 0041 102 = 0048 103 = 004A 104 = 004D 111 = 004F 200 = 0051 201 = 0073 202 = 0083 204 = 0093
206 = 00F6 1 = -011F 30 = -01D8

FEATURES SUPPORTED

ONE WORD INTEGERS

IOCS

CALLED SUBPROGRAMS

SRW29 SBM26 FMPY FDIY FLD FSTO FDVR FAXI CARDZ PRNTZ SRED SWRT SCOMP SF10 STOF
SI01

REAL CONSTANTS

.990000E 01=002C

INTEGER CONSTANTS

2=002E 3=002F

CORE REQUIREMENTS FOR

COMMON 0 VARIABLES 44 PROGRAM 430

END OF COMPILATION

// XFO

THE DRAG ON CYLINDERS IN A STREAM OF DUST-LADEN AIR

THE PHYSICAL PARAMETERS ARE

XL = -5.0000009
 REZ = 0.735665E 02
 XK = 0.339462E 03
 P = 0.159429E 02
 DC = 0.80000
 DP = 0.015000
 RHO = 0.001213
 SIGMA = 2.600000
 XMU = 0.0001789
 UF = 700.0

IMPACTION EFFICIENCY OF A CIRCULAR CYLINDER

G4LFF = 0.000000
 G4RIT = 1.000000
 SIGML = -1.
 DTAU = 0.100000
 NIAP = 10
 NSFP = 2
 NX = 20

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THE MOTION OF A CRITICAL PARTICLE IS GIVEN BY

ITER = 1
 G4LEF = 0.000000
 G4ZER = 0.500000
 G4RIT = 1.000000

TAU	G(1)	G(2)	G(3)	G(4)	UX	UY	CDRE
0.0000	0.961180	0.007842	-5.000090	0.500000	0.9611	0.0078	24.0000
0.2000	0.961179	0.007842	-4.807763	0.501568	0.9581	0.0088	24.3728
0.4000	0.961177	0.007843	-4.615527	0.503136	0.9566	0.0099	25.2001
0.6000	0.961172	0.007844	-4.423291	0.504705	0.9508	0.0113	26.1397
0.8000	0.961164	0.007847	-4.231056	0.506274	0.9464	0.0129	27.1814
1.0000	0.961153	0.007851	-4.038823	0.507844	0.9415	0.0149	28.2857
1.2000	0.961138	0.007856	-3.846592	0.509415	0.9358	0.0172	29.3761
1.4000	0.961118	0.007864	-3.654366	0.510987	0.9293	0.0201	30.3616
1.6000	0.961093	0.007874	-3.462144	0.512560	0.9218	0.0236	31.2151
1.8000	0.961060	0.007887	-3.269928	0.514136	0.9131	0.0280	32.1351
2.0000	0.961019	0.007904	-3.077720	0.515715	0.9029	0.0334	34.0064
2.2000	0.960966	0.007928	-2.885521	0.517298	0.8908	0.0404	35.7737
2.4000	0.960898	0.007960	-2.693334	0.518887	0.8765	0.0493	37.7117
2.6000	0.960813	0.008002	-2.501165	0.520483	0.8595	0.0611	39.8650
2.8000	0.960704	0.008060	-2.309010	0.522089	0.8389	0.0767	42.2912
3.0000	0.960565	0.008141	-2.116882	0.523709	0.8139	0.0980	45.0651
3.2000	0.960386	0.008253	-1.924786	0.525348	0.7836	0.1276	48.2873
3.4000	0.960153	0.008415	-1.732731	0.527014	0.7467	0.1697	52.0961
3.6000	0.959848	0.008654	-1.540730	0.528720	0.7025	0.2314	56.6872
3.8000	0.959444	0.009018	-1.348799	0.530485	0.6515	0.3242	62.3423
4.0000	0.958911	0.009596	-1.156961	0.532343	0.5989	0.4682	69.4785
4.2000	0.958216	0.010551	-0.965246	0.534351	0.5638	0.6962	78.7198
4.3199	0.957729	0.011394	-0.850289	0.535669	0.5724	0.4465	73.4206
4.3250	0.957707	0.011418	-0.844543	0.535737	0.5739	0.4522	73.8589

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THE MOTION OF A CRITICAL PARTICLE IS GIVEN BY

ITER = 10
 G4LFF = 0.919921
 G4ZER = 0.920898
 G4RIT = 0.921875

TAU	G(1)	G(2)	G(3)	G(4)	UX	UY	CDRE
0.0000	0.963851	0.013783	-5.000000	0.920898	0.9638	0.0137	24.0000
0.2000	0.963850	0.013783	-4.807229	0.923655	0.9612	0.0154	24.3607
0.4000	0.963848	0.013784	-4.614459	0.926411	0.9583	0.0174	25.1591
0.6000	0.963844	0.013787	-4.421689	0.929169	0.9551	0.0197	26.0637
0.8000	0.963837	0.013791	-4.228919	0.931926	0.9516	0.0224	27.0664
1.0000	0.963828	0.013798	-4.036151	0.934685	0.9476	0.0256	28.1339
1.2000	0.963816	0.013807	-3.843386	0.937446	0.9432	0.0294	29.1998
1.4000	0.963800	0.013819	-3.650624	0.940208	0.9383	0.0339	30.1789
1.6000	0.963780	0.013835	-3.457865	0.942973	0.9329	0.0395	31.0271
1.8000	0.963755	0.013857	-3.265111	0.945743	0.9268	0.0462	31.8578
2.0000	0.963723	0.013884	-3.072362	0.948517	0.9201	0.0545	33.3259
2.2000	0.963684	0.013921	-2.879621	0.951297	0.9126	0.0647	35.1617
2.4000	0.963637	0.013969	-2.686889	0.954086	0.9045	0.0775	36.9541
2.6000	0.963579	0.014033	-2.494166	0.956886	0.8958	0.0937	38.9166
2.8000	0.963509	0.014117	-2.301457	0.959700	0.8868	0.1132	41.0873
3.0000	0.963426	0.014229	-2.108763	0.962535	0.8780	0.1405	43.5126
3.2000	0.963328	0.014379	-1.916086	0.965395	0.8707	0.1745	46.2488
3.4000	0.963216	0.014583	-1.723432	0.968290	0.8668	0.2185	49.3641
3.6000	0.963094	0.014860	-1.530800	0.971233	0.8703	0.2752	52.9396
3.8000	0.962971	0.015242	-1.338194	0.974242	0.8878	0.3473	57.0675
4.0000	0.962872	0.015768	-1.145610	0.977351	0.9305	0.4354	61.8417
4.2000	0.962849	0.016487	-0.953039	0.980564	1.0152	0.5346	67.3322
4.4000	0.963000	0.017445	-0.760457	0.983954	1.1630	0.6257	73.5249
4.6000	0.963486	0.018646	-0.567821	0.987560	1.3876	0.6659	80.2095
4.8000	0.964518	0.019978	-0.375021	0.991422	1.6672	0.5890	86.8126
5.0000	0.966240	0.021135	-0.181953	0.995598	1.9132	0.3453	92.2867
5.1099	0.967489	0.022292	-0.075599	0.997873	1.9871	0.0752	94.9705
5.1249	0.967668	0.0221300	-0.061086	0.998192	1.9924	0.0609	95.3191

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THE POSITION OF A CRITICAL PARTICLE IS GIVEN BY

ITER = 20
 G4LEF = 0.921587
 G4ZER = 0.921588
 G4RIT = 0.921589

TAU	G(1)	G(2)	G(3)	G(4)	UX	UY	CORE
0.0000	0.963856	0.013792	-5.000000	0.921588	0.9638	0.0137	24.0000
0.2000	0.963855	0.013792	-4.807229	0.924346	0.9612	0.0154	24.3607
0.4000	0.963853	0.013793	-4.614457	0.927104	0.9583	0.0174	25.1520
0.6000	0.963849	0.013796	-4.421685	0.929863	0.9551	0.0197	26.0635
0.8000	0.963843	0.013800	-4.228915	0.932623	0.9516	0.0224	27.0662
1.0000	0.963834	0.013807	-4.036146	0.935384	0.9476	0.0256	28.1336
1.2000	0.963822	0.013816	-3.843379	0.938146	0.9432	0.0294	29.1994
1.4000	0.963806	0.013828	-3.650616	0.940910	0.9384	0.0340	30.1785
1.6000	0.963785	0.013844	-3.457856	0.943677	0.9329	0.0395	31.0267
1.8000	0.963760	0.013866	-3.265101	0.946448	0.9268	0.0462	31.8572
2.0000	0.963729	0.013893	-3.072351	0.949224	0.9201	0.0545	33.3243
2.2000	0.963690	0.013930	-2.879609	0.952006	0.9127	0.0647	35.1604
2.4000	0.963642	0.013978	-2.686875	0.954797	0.9045	0.0776	36.9526
2.6000	0.963585	0.014042	-2.494152	0.957599	0.8958	0.0937	38.9147
2.8000	0.963515	0.014126	-2.301441	0.960415	0.8869	0.1142	41.0849
3.0000	0.963432	0.014238	-2.108746	0.963251	0.8781	0.1406	43.5096
3.2000	0.963334	0.014368	-1.916069	0.966114	0.8708	0.1746	46.2449
3.4000	0.963222	0.014592	-1.723412	0.969011	0.8670	0.2185	49.3590
3.6000	0.963100	0.014869	-1.530780	0.971956	0.8706	0.2752	52.9328
3.8000	0.962977	0.015251	-1.338172	0.974966	0.8801	0.3472	57.0582
4.0000	0.962879	0.015776	-1.145586	0.978067	0.8908	0.4352	61.8289
4.2000	0.962857	0.016495	-0.953014	0.981291	1.0156	0.5341	67.3143
4.4000	0.963008	0.017452	-0.760430	0.984683	1.1633	0.6250	73.5000
4.6000	0.963495	0.018651	-0.567785	0.988290	1.3877	0.6649	80.1751
4.8000	0.964526	0.019980	-0.374991	0.992153	1.6566	0.5879	86.7667
5.0000	0.966245	0.021134	-0.181922	0.996270	1.9120	0.3445	92.2295
5.1999	0.967850	0.021306	-0.046538	0.999243	1.9950	0.0464	95.3180
5.1969	0.968415	0.021308	-0.001035	1.000241	1.9995	0.0010	95.5894

THE IMPACTION EFFICIENCY IS 0.9215E 00

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THE MOTION OF THE PARTICLES AT THE CYLINDER POSITION IS GIVEN BY

NZER = 12C

NJBP = 4

TAU	VX	VY	X	Y	UX	UY	FY
4.1690	0.955357	0.000000	-1.000526	0.000000	0.0010	0.0000	1.83428
4.1690	0.955364	0.000336	-1.000514	0.023886	0.0027	0.0238	1.83324
4.1710	0.955378	0.000785	-0.998553	0.055736	0.0064	0.0556	1.82856
4.1730	0.955417	0.001232	-0.996544	0.087586	0.0160	0.0871	1.82018
4.1760	0.955474	0.001677	-0.993535	0.119438	0.0298	0.1183	1.80807
4.1810	0.955537	0.002121	-0.988577	0.151292	0.0459	0.1495	1.79221
4.1850	0.955628	0.002561	-0.983570	0.183147	0.0679	0.1797	1.77265
4.1920	0.955739	0.002995	-0.977564	0.215004	0.0939	0.2094	1.74935
4.2000	0.955858	0.003428	-0.969602	0.246867	0.1226	0.2388	1.72231
4.2090	0.955999	0.003855	-0.960640	0.278733	0.1557	0.2674	1.69153
4.2190	0.956184	0.004289	-0.950676	0.310616	0.1931	0.2951	1.65693
4.2300	0.956354	0.004719	-0.939714	0.342495	0.2347	0.3216	1.61864
4.2420	0.956547	0.005340	-0.927748	0.374378	0.2807	0.3467	1.57661
4.2550	0.956763	0.005751	-0.914779	0.406266	0.3307	0.3702	1.53084
4.2700	0.956996	0.006154	-0.899855	0.438165	0.3843	0.3929	1.48131
4.2870	0.957249	0.006548	-0.882964	0.470077	0.4420	0.4145	1.42801
4.3050	0.957529	0.006930	-0.865068	0.501995	0.5039	0.4339	1.37096
4.3240	0.957769	0.007698	-0.846172	0.533943	0.5700	0.4508	1.30967
4.3450	0.958093	0.008059	-0.825305	0.565890	0.6400	0.4657	1.24505
4.3680	0.958439	0.008408	-0.802470	0.597853	0.7142	0.4784	1.17663
4.3930	0.958810	0.008745	-0.777668	0.629832	0.7925	0.4883	1.10441
4.4210	0.959137	0.009586	-0.749939	0.661867	0.8757	0.4958	1.02771
4.4510	0.959565	0.009888	-0.720228	0.693903	0.9627	0.4995	0.94782
4.4830	0.960020	0.010179	-0.688544	0.725957	1.0528	0.4987	0.86410
4.5190	0.960474	0.011027	-0.652960	0.758082	1.1480	0.4939	0.77584
4.5590	0.961009	0.011281	-0.613464	0.790235	1.2477	0.4839	0.68431
4.6040	0.961593	0.011538	-0.569094	0.822441	1.3523	0.4677	0.58882
4.6540	0.962233	0.012317	-0.519841	0.854754	1.4595	0.4435	0.48878
4.7130	0.962973	0.013062	-0.461844	0.887174	1.5733	0.4094	0.38473
4.7830	0.963851	0.013253	-0.393145	0.919734	1.6906	0.3612	0.27699
4.8730	0.964955	0.013895	-0.305038	0.952602	1.8135	0.2902	0.16438
4.9010	0.965300	0.013940	-0.277681	0.960891	1.8434	0.2665	0.13555

THE DRAG COEFFICIENT OF THE CYLINDER DUE TO THE REFLECTION OF PARTICLES IS 2.3435

THE TARGET REYNOLDS NUMBER IS 0.3796E 04

THE DRAG COEFFICIENT OF THE CYLINDER DUE TO THE AIR ALONE IS 1.000

// PAUS

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APPENDIX B

COMPUTER PROGRAM FOR A SPHERE IN A

TUBE OR IN FREE SPACE

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PAGE 1

// JOB T

LOG DRIVE CART SPEC CART AVAIL PHY DRIVE
0000 0305 0305 0001

V2 411 ACTUAL 16K CNFIG 16K

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PAGE 2

// FOR
*ONE WORD INTEGERS
*LIST ALL

SUBROUTINE SPM34(G,UX,UY,ITURE)

C THIS SUBROUTINE CALCULATES THE FLOW ABOUT
C A SPHERE IN A CIRCULAR TUBE FOR A SPHERE TO TUBE DIAMETER RATIO
C OF 0.8/2.54 OR FOR A SPHERE IN FREE SPACE.
C

C DIMENSION G(4)

IF(ITURE)1,1,2

1 RSQ = G(3)**2+G(4)**2

UX = 1.0 - 1.0/RSQ**1.5 + 1.5*G(4)**2/RSQ**2.5

UY = 1.5*G(3)*G(4)/RSQ**2.5

GO TO 3

2 R = SORT(G(3)**2 + G(4)**2)

COST = -G(3)/R

SINT = G(4)/R

DC = 0.8

C = 1.0

A = (2.54/DC)*C

S2 = 7.5098907

CZERO = 1.5401075

PI = 3.14159265358979323846264

AZERO = -C*(C/A)**2*S2*CZERO/(9.0*PI)

FT1 = CZERO*C**3/(3.*R**3)

FT2 = 2.0*AZERO/A

UT = SINT*(1.0+FT1-FT2)

FRI = 2.0*CZERO*C**3/(3.*R**3)

FR2 = FT2

UR = -COST*(1.0-FRI-FR2)

UX = -UR*COST + UT*SINT

UY = UR*SINT + UT*COST

3 CONTINUE

RETURN

END

VARIABLE ALLOCATIONS

RSQ(R) = 0000

A(R) = 000C

FT2(R) = 001B

R(R) = 0002

S2(R) = 000E

UT(R) = 001A

COST(R) = 0004

CZERO(R) = 0010

FRI(R) = 001C

SINT(R) = 0006

PI(R) = 0012

FR2(R) = 001E

DC(R) = 0008

AZERO(R) = 0014

UR(R) = 0020

C(R) = 000A

FT1(R) = 0016

STATEMENT ALLOCATIONS

1 = 0059 2 = 00A3 3 = 016E

FEATURES SUPPORTED
ONE WORD INTEGERS

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PAGE 3
CALLED SURPROGRAMS
FSORT FAXB FADD FSUB FMPY FMPYX FDIW FLD FLDX FSTO FSBR FDVR FAXI SNR SUBIN
REAL CONSTANTS
.100000F 01=002R .150000E 01=002A .250000E 01=002C .800000E 00=002E .254000E 01=0030 .750989E 01=0032
.15401CF 01=0034 .314159E 01=0036 .900000E 01=0038 .300000E 01=003A .200000E 01=003C

INTEGER CONSTANTS
2=003E 3=009F

CORP REQUIREMENTS FOR SBM34
COMMON 0 VARIABLES 40 PROGRAM 328

RELATIVE ENTRY POINT ADDRESS IS 0040 (HEX)

END OF COMPILATION

// DUP

*STORE WS UA SBM34
CART ID 0305 DB ADDR 5A0A DB CNT 0018

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```
PAGE 4
// FOR
*OME WORD INTEGERS
*LTST ALL
C
FUNCTION CDRE(RE)
C
C THIS FUNCTION COMPUTES THE PRODUCT OF DRAG COEFFICIENT
C AND REYNOLDS NUMBER FOR A SPHERE AS A FUNCTION
C OF REYNOLDS NUMBER LESS THAN 10000.
C
A1=1.0/24.
A2=-2.3363*1.E-04
A3=2.0154*1.E-06
A4=-6.9105*1.E-09
R0=-1.29536
R1=9.86*1.E-01
R2=-4.6677*1.E-02
R3=1.1235*1.E-03
C
C CHOOSE THE APPROPRIATE POLYNOMIAL
C IF (RE-4.0)2,7,7
C
C INITIAL ESTIMATE
C
2 IF (RE-0.00001)3,4,4
3 CDRE = 24.0
GO TO 30
4 X=24.*RE
C
C BEGIN NEWTON METHOD ITERATION
C
CONTINUE
DO 6 ITER=1,20
FX=A1*X+A2*X**2+A3*X**3+A4*X**4-RE
FPX=A1+2.*A2*X+3.*A3*X**2+4.*A4*X**3
DELX=FPX/FPX
X=X-DELX
C
C CHECK FOR CONVERGENCE
C
EPS=1.E-06
IF (ABS(DELX/X)-EPS)5,5,6
5 CDRE=X/RE
GO TO 30
6 CONTINUE
```

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```
PAGE 5
GO TO 29
C
C INITIAL ESTIMATE
C
7 CD = 1.0
ELOG = 0.434294481903252
X=ALOG(CD*RE**2)*ELOG
C
C REGIN NEWTON METHOD ITERATION
C
DO 74 ITER=1,20
FX=R0+R1*X+R2**X**2+33**X**3 - ALOG(RE)*ELOG
FPX=B1+2.*B2*X+3.*B3*X**2
DELX=FX/FPX
X=X-DELX
C
C CHECK FOR CONVERGENCE
C
C
EPS=1.E-06
IF(ABS(DELX/X)-EPS)22,22,24
22 CDRE=10.**X/RE
GO TO 30
24 CONTINUE
29 WRITE(3,201)
30 RETURN
C
C FORMATS FOR OUTPUT STATEMENTS
C
202 FORMAT(16H0 NO CONVERGENCE)
C
END
VARIABLE ALLOCATIONS
CDRE(I) =0000 A1(I) =0002 A2(I) =0004 A3(I) =0006 A4(I) =0008 B0(I) =000A
B1(I) =000C B2(I) =000E B3(I) =0010 B4(I) =0012 B5(I) =0014 B6(I) =0016
DELX(I) =0018 EPS(I) =001A ELOG(I) =001C ELOG6(I) =001E ITER(I) =0028
STATEMENT ALLOCATIONS
207 =0059 2 =00AR 3 =00AF 4 =0085 5 =012D 6 =0135 7 =014C 22 =0188 24 =01C3 29 =C1CC
30 =D1D0
FEATURES SUPPORTED
ONE WORD INTEGERS
CALLED SUBPROGRAMS
FARS FALOG FAXB FADD FSUB FMPY FMPY FSTO F58R FAXI SWRT SCOMP SNR SUBIN
REAL CONSTANTS
.100000F 01=002A .240000E 02=002C .233630E 01=002E .100000E-03=0030 .201540E 01=0032 .100000E-05=0034
```

PAGE 6

.691050F 01=0036 .100000E-08=0038 .129536E 01=003A .986000E 01=003C .100000E 00=003E .466770E 01=0040
.100000E-01=0042 .112500E 01=0044 .100000E-02=0046 .400000E 01=0048 .100000E-04=004A .200000E 01=004C
.300000E 01=004E .634294E 00=0050 .100000E 02=0052

INTEGER CONSTANTS

1=0054 20=0055 2=0056 3=0057 4=0058

CORE REQUIREMENTS FOR CDRE

COMMON 0 VARIABLES 42 PROGRAM 426

RELATIVE ENTRY POINT ADDRESS IS 0063 (HEX)

END OF COMPILATION

// DUP

*STORE WS UA CDRE

CART ID 0305 DB ADDR 5A25 DB CNT 0022

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PAGE 7
// FOR
#ONE WORD INTEGERS
*LIST ALL
SURROUTINE SRM22(IN,Y,F,X,H,IRUNG,M)
C
C FOURTH ORDER RUNGE KUTTA METHOD
C FOR N FIRST ORDER O.D.E.
C DIMENSION PH(50),SAVY(50),Y(50),F(50)
C GO TO (2,3,4,5,6),M
C
C PASS 1
C
C 2 IRUNG=1
C RETURN
C
C PASS 2
C
C 3 DO 22 J=1,N
C SAVY(J)=Y(J)
C PH(J)=F(J)
C 22 Y(J)=SAVY(J)+0.5*H*F(J)
C X=X+0.5*H
C IRUNG=1
C RETURN
C
C PASS 3
C 4 DO 23 J=1,N
C PH(J)=PH(J)+2.0*F(J)
C 33 Y(J)=SAVY(J)+0.5*H*F(J)
C IRUNG=1
C RETURN
C
C PASS 4
C 5 DO 44 J=1,N
C PH(J)=PH(J)+2.0*F(J)
C 44 Y(J)=SAVY(J)+H*F(J)
C X=X+0.5*H
C IRUNG=1
C RETURN
C PASS 5
C 6 DO 55 J=1,N
C 55 Y(J) = SAVY(J) + (PH(J) + F(J))*H/6.0
C IRUNG=2
C RETURN
C
C END
```

PAGE 8

VARIABLE ALLOCATIONS
PPI(R) 0062-0000 SAVY(R) 00C6-0064 JII 00C6

STATEMENT ALLOCATIONS
2 0105 3 0108 22 011E 4 0141 33 0192 5 016D 44 017E 6 019F 55 01A3

FEATURES SUPPORTED
ONE WORD INTEGERS

CALLLED SUBPROGRAMS
FADD FADDX FMPY FMPYX FD:V FLD FLDX FSTO FSTOX SUBSC SUBIN

REAL CONSTANTS
0500000E 00=00CA 0200000E 01=00CC 0600000E 01=00CE

INTEGER CONSTANTS
1=0000 2=0001

CORE REQUIREMENTS FOR SRM22
COMMON 0 VARIABLES 202 PROGRAM 250

RELATIVE ENTRY POINT ADDRESS IS 00D2 (HEX)

END OF COMPILATION

// DUP

*STORE MS UA SRM22

CART ID 0305 DR ADDR 5A47 DB CNT 0C12

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```
// FOR
*ONE WORD INTEGERS
*LIST ALL
SURROUTINE SBM32(G,TAU,DTAW,XK,REZ,UX,UY,XCDRE)
C
C THIS SUBROUTINE CALCULATES PARTICLE MOTION DURING THE FINAL
C INCREMENT OF TAU
C
C DIMENSION G(4),DG(4)
C M=N
C
C CALL ON RUNGE KUTTA SUBROUTINE
C
C CONTINUE
C M=M+1
C CALL SBM22(4,G,DG,TAU,DTAW,IRUNG,M)
C IF(IRUNG-1)10,9,10
C 9 RE=REZ*(UY-G(2))**2+(UX-G(1))**2)**0.5
C XCDRE=CDRE(RE)
C DG(1)=(XCDRE/(24.0*XK))*(UX-G(1))
C DG(2)=(XCDRE/(24.0*XK))*(UY-G(2))
C DG(3)=G(1)
C DG(4)=G(2)
C GO TO 8
C 10 CONTINUE
C M=0
C
C CALCULATE FLUID VELOCITY AT PARTICLE POSITION
C
C CALL SBM34 (G,UX,UY,ITUBE)
C
C INTEGRATE ACROSS ANOTHER STEP IF REQUIRED
C
C IF(G(4) - 1.0)11,11,12
C 11 DELX = SORT(1.0 - G(4)**2)
C GO TO 13
C 12 DELX = 0.0
C 13 HITS=G(3)+G(1)*1.1*DTAW+DELX
C IF(HITS)8,8,12
C 16 CONTINUE
C RETURN
C END
VARIABLE ALLOCATIONS
DGIR I=0006-0000 REIR J=0008 DELXR J=000A HITSR J=000C MII J=0010 IRUNG(I)=0011
ITITURE(I)=0012
```

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STATEMENT ALLOCATIONS

8 =0050 9 =0065 10 =008B 11 =00CE 12 =00DF 13 =00E3 18 =00F8

FEATURES SUPPORTED
ONE WORD INTEGERS

CALLED SURPROGRAMS
SBM22 CORE
FDVR FA... SUBIN

SBM34
FSORT
FAXB

FADD

FADDD

FADXX

FASUB

FASUBX

FAPY

FLD

FLDX

FSTO

FSTOX

FSBR

REAL CONSTANTS

.50 000E 00=001B

.24000E 02=001A

.10000E 01=001C

.00000E 00=001E

.11000E 01=0020

INTEGER CONSTANTS

0=0022

4=0024

2=0025

CORE REQUIREMENTS FOR SRM32

COMMON 0 VARIABLES 24 PROGRAM 226

RELATIVE ENTRY POINT ADDRESS IS 0026 (HEX)

END OF COMPILATION

// DUP

*STORE WS UA SBM32

CART ID 0305 DB ADDR 5A59 DB CNT 0011

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```
// FOR
*ONE WORD INTEGERS
*LIST ALL
SURROUTINE SBM30(G4LFT,G4RIT,SIGNL,DTAU,NIBP,NSBP,NX,XL,XK,REZ,
IGAZER)

C
C THIS SUBROUTINE CALCULATES THE
C IMPACTION EFFICIENCY OF A SPHERE
C
C DIMENSION G(4),DG(4)
C WRITE(3,200)
C WRITE(9,201)G4LFT,G4RIT,SIGNL,DTAU,NIBP,NSBP,NX
C
C HALF INTERVAL ITERATION FOR INITIAL G4 VALUE
C DO 31 ITER=1,NX
C
C SET AND PRINT INITIAL CONDITIONS
C
C M=0
C NSTEP=0
C TAU=0.0
C G(1)=XL
C G4ZER=(G4LFT+G4RIT)/2.0
C G(4)=GAZER
C CALL SBM34 (G,UX,UY,ITUBE)
C G(1) = UX
C G(2)=UY
C RE=REZ*((UY-G(2))**2+(UX-G(1))**2)**0.5
C XCDRE=CDRE(IRE)
C IP=ITER/NIRP*NIBP
C IF(IP-ITER)5,7,5
C 5 IF(ITER-1)6,7,6
C 6 IF(ITER-NX)8,7,8
C 7 CONTINUE
C WRITE (3,205)
C WRITE(3,203)ITER,G4LFT,G4ZER,G4RIT,TAU,G(1),G(2),G(3),G(4),UX,UY,
C IXCDRE
C
C CALL ON RUNGE KUTTA SUBROUTINE
C
C 8 CONTINUE
C M=M+1
C CALL SBM22(4,G,DG,TAU,DTAU,IRUNG,M)
C IF(IRUNG-1)10,9,10
C 9 RE=REZ*((UY-G(2))**2+(UX-G(1))**2)**0.5
```

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XCDRE=CDRE(RE)
DG(1)=(XCDRE/(24.0#XK1))*(UX-G(1))
DG(2)=(XCDRE/(24.0#XK1))*(UY-G(2))
DG(3)=G(1)
DG(4)=G(2)

GO TO 8

10 CONTINUE

M=0

C CALCULATE FLUID VELOCITY AT PARTICLE POSITION

C CALL SBM34 (G,UX,UY,ITUBE)

C PRINT SOLUTIONS

IS = ITER/NIBP-NIBP

IF(15-ITER)11,13,11

11 IF(ITER-1)12,13,12

12 IF(ITER-NX)16,13,16

13 NSTEP=NSTEP+1

IF(NSTEP-NSBP)16,14,16

14 CONTINUE

NSTEP=0

TAU = TAU+0.0001

WRITE(9,204)TAU,G(1),G(2),G(3),G(4),UX,UY,XCDRE

C INTEGRATE ACROSS ANOTHER STEP IF REQUIRED

16 CONTINUE

IF(G(4)-1.0)17,17,18

17 DELX=SQRT(1.0-G(4)**2)

GO TO 19

18 DELX=0.0

19 HITS=G(3)+G(1)*1.1#DTAU+DELX

IF(HITS)18,8,20

20 CONTINUE

C CHANGE INCREMENT SIZE NEAR SPHERE AND INTEGRATE FURTHER

DTAU=DTAU/10.0

CALL SBM32(G,TAU,DTAU,XK,REZ,UX,UY,XCDRE)

C PRINT SOLUTIONS

IF(15-ITER)21,23,21

21 IF(ITER-1)22,23,22

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PAGE 13

22 IF (ITER-NX) 24, 23, 24
23 CONTINUE
WRITE(3, 204) TAU, G(1), G(2), G(3), G(4), UX, UY, XCDRE

24 CONTINUE
DTAU=DTAU/100.0
CALL SBMSZ(G, TAU, DTAW, XK, REZ, UX, UY, XCDRE)

C
C
C
PRINT SOLUTIONS

IF (IS-ITER) 25, 27, 25
25 IF (ITER-1) 26, 27, 26
26 IF (ITER-NX) 28, 27, 28
27 CONTINUE
WRITE(3, 204) TAU, G(1), G(2), G(3), G(4), UX, UY, XCDRE

28 CONTINUE
C
C
C
CALCULATE ORDINATE AT TANGENT POINT OF TANGENT PATH

ORD = G(1)/SQRT(G(1)**2 + G(2)**2)

FIND INTERVAL HALF WITH THE SIGN CHANGE

IF (G(4)-ORD)*SIGNL-O.0) 29, 29, 30

29 GARIT=GAZER
GO TO 31
30 G4LFT=GAZER
31 CONTINUE
EM = GAZER**2
WRITE(3, 207) EM

C
C
C
RETURN

FORMATS FOR OUTPUT STATEMENTS

200 FORMAT (1H0, 41X, 3HIMPACTION EFFICIENCY OF A SPHERE /

1 1H0)
201 FORMAT (10HOG4LEF = , F10.6/ 10H G4RIT = , F10.6/ 10H SIGNL = ,
1 F3.0/ 10H DTAU = , F10.6/ 10H NLRP = , F13/ 10H NSBP = , F13/
2 10H NX = , F13)

203 FORMAT (10HITER = , F13/ 10H G4LEF = , F10.6/ 10H GAZER = ,
1 F10.6/ 10H G4RIT = , F10.6/ 7H0 TAU, 11X, 4HG(1), 12X,
2 4HG(2), 12X, 4HG(3), 12X, 4HG(4), 14X, 2HUX, 14X, 2HUY ,
3 12X, 4HCORE /
4 1H0, F7.4, 4F16.6, 3F16.4)

204 FORMAT (1H, F7.4, 4F16.6, 3F16.4)
205 FORMAT (46HITHE MOTION OF A CRITICAL PARTICLE IS GIVEN BY)

ZOT FORMAT(30)THE IMPACTION EFFICIENCY IS (E10.6)

VARIABLE ALLOCATIONS

GIR J=0006-0000	DGIR J=000E-0008	TAUJR J=0010	UKJR J=0012	UYJR J=0014	REJR J=0016
XCDREJR J=0018	TAVJR J=001A	DELKJR J=001C	MITSJR J=001E	DTAMJR J=0020	ORDJR J=0022
FMJR J=0024	ITERJR J=002A	MJR J=002B	NSTEPJR J=002C	ITUBJR J=002D	IPJR J=002E
IRUNGJR J=002F	ISJR J=0030				

STATEMENT ALLOCATIONS

200 -004D 201 -0046	203 -009E 204 -00E6	205 -00EE 207 -0107	5 -01CE 6 -01D4	7 -01DA 8 -0202
9 -0217 10 -026D	11 -0287 12 -028D	13 -0293 14 -029F	16 -02C5 17 -02CE	18 -02DF 19 -02E3
20 -02F8 21 -030E	22 -0314 23 -031A	24 -0336 25 -034C	26 -0352 27 -0358	28 -0374 29 -039C
30 -03A2 31 -03AC				

FEATURES SUPPORTED
ONE WORD INTEGERS

CALLED SUBPROGRAMS

SRM34	CDRE	SRM2Z	FSORT	SRM3Z	FAZB	FADD	FADDY	FSUBX	FSUB	SIOP	SIOT	FMPY	FDIV	FLD	FLDX	FSTO
FS-DX	F5BR	FDVR	FDVFX	FAXI	SWRT	SCOMP	STOFX	STOI	STOF			SUBIN				

REAL CONSTANTS

.000000E 00=0036 .200000E 01=0038 .500000E 00=003A .240000E 02=003C .100000E -03=003E .100000E 01=0040
 .110000E 01=0042 .100000E 02=0044 .100000E 03=0046

INTEGER CONSTANTS

3=0048 1=0049 0=004A 2=004B 4=004C

CORE REQUIREMENTS FOR SRM30

COMMON 0 VARIABLES 54 PROGRAM 904

RELATIVE ENTRY POINT ADDRESS IS 0119 (HEX)

END OF COMPILATION

// DUP

*STORE WS UA SRM30

CART ID 0305 DR ADDR 5AGA DR CNT 003E

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// FOR  
#ONE WORD INTEGERS  
*LIST ALL  
SUBROUTINE SPM31(XL,G,ZER,DTAU,XK,REZ,PX,NZER,NJBP)
```

C THIS SUBROUTINE CALCULATES THE MOTION OF PARTICLES
C IN A FLUID STREAM MOVING TOWARD A SPHERE
C AND CALCULATES THE FORCE OF PARTICLE IMPACT ON THE SPHERE
C

C DIMENSION G(4),DG(4)
C DIMENSION YZER(500)

C SET NUMBER OF INCREMENTS AT INITIAL POSITION

```
WRITE(3,20C)  
WRITE(3,20I)NZER,NJBP  
WRITE(3,20Z)  
NCE=NZER+1  
DELY=G/ZER/FLOAT(NZER)  
FSUM=0.0
```

C STEPWISE INTEGRATION FOLLOWING PARTICLE POSITION

```
DO 30 ITR=1,NCE  
C  
C SET AND PRINT INITIAL CONDITIONS  
M=0  
TAU=0.0  
G(3)=XL  
YZER(ITR)=FLOAT(ITR-1)*DELY  
GAZER=YZER(ITR)  
G(4)=GAZER  
CALL SPM34 (G,U,X,U,Y,I,TURSE)  
UZER=UX  
G(1)=UX  
G(2)=UY  
C  
C CALL ON RUNGE KUTTA SUBROUTINE  
C  
C CONTINUE  
M=M+1  
CALL SPM22(4,G,DG,TAU,DTAU,IRUNG,M)  
IF(IRUNG-1)10,9,10  
9 RE=REZ*(UY-G(2))**2+(UX-G(1))**2)**0.5  
XCORE=CDRE(RE)
```

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DG(1)=(XCDRE/(24.0*WK))*(UX-G(1))
DG(2)=(XCDRE/(24.0*WK))*(UY-G(2))
DG(3)=G(1)
DG(4)=G(2)

GO TO 8

10 CONTINUE

M=0

C CALCULATE FLUID VELOCITY AT PARTICLE POSITION

C CALL SBM34 (G,UX,UY,ITUBE)

C INTEGRATE ACROSS ANOTHER STEP IF REQUIRED

IF(IG(4) - 1.0)11,11,12

11 DELX = SORT(1.0 - G(4)**2)

GO TO 13

12 DELX = 0.0

13 HITS=G(3)+G(1)*1.1*DTAU+DELX

IF(HITS)8,8,18

18 CONTINUE

C CHANGE INCREMENT SIZE NEAR SPHERE AND INTEGRATE FURTHER

DTAW=DTAU/10.0

CALL SBM32(G,TAU,DTAW,XK,REZ,UX,UY,XCDRE)

DTAW=DTAU/100.0

CALL SBM32(G,TAU,DTAW,XK,REZ,UX,UY,XCDRE)

C CO-ORDINATES, VELOCITY, AND PRESSURE DERIVATIVE AT SPHERE

VX=G(1)

VY=G(2)

X=G(3)

Y=G(4)

V=SQRT(VX**2+VY**2)

PI = 3.14159265358979323846264

GAMMA = PI - 2.0*ATAN(Y/SORT(1.0-Y**2)) - ATAN(VY/VX)

FY= UXZER*VX - V*COS(GAMMA)

FSUM = FSUM + FLOAT(2*ITER-1)*FY

CONTINUE

C PRINT SOLUTIONS

IS = ITER/NJRP+NJBP

IF(15-ITER)21,23,21

PAGE 17

21 IF(ITER-1)22,23,22
 22 IF(ITER-NCE)24,23,24
 23 CONTINUE
 WRITE(3,203)TAU,VX,VY,X,Y,UX,UY,FY
 24 CONTINUE
 30 CONTINUE

C DRAG COEFFICIENT OF SPHERE DUE TO PARTICLE IMPACT
 C
 C

PX = 2.*FSUM*DELY**2
 WRITE(3,204)PX
 RETURN

C FORMATS FOR OUTPUT STATEMENTS
 C
 C

200 FORMAT: 63H1THE MOTION OF THE PARTICLES AT THE SPHERE POSITION IS

1 GIVEN BY 1

201 FORMAT(10HONZER = ,I3,20X, 10H NJBP = ,I3)
 202 FORMAT(7H0 TAU, 11X, 4H VX , 12X, 4H VY , 11X, 4H X , 12X,
 1 4H Y , 14X, 2HUX, 14X, 2HUY , 14X, 4H FY /1H)
 203 FORMAT (F7.4, 4F16.6, 2F16.4, F16.5)
 204 FORMAT(75H0THE DRAG COEFFICIENT OF THE SPHERE DUE TO THE REFLECTI

ION OF PARTICLES IS , F9.4)
 END

VARIABLE ALLOCATIONS

GIR 1=0006-0000	DG(R 1=000E-0008	YZER(R 1=03F6-0010	DELY(R 1=03F8	FSUM(R)=03FA	TAU(R 1=03FC
UX(R 1=03FE	UY(R 1=0400	UXZER(R 1=0402	RE(R 1=0404	XCDRE(R)=0406	DELX(R 1=0408
HITS(R 1=040A	DTAW(R 1=040C	VX(R 1=040E	VY(R 1=0410	X(R 1=0412	Y(C 1=0414
VIR 1=0416	PI(R 1=0418	GAMMA(R 1=041A	FY(R 1=041C	NCE(I 1=0426	ITER(I 1=0427
MI 1=042R	ITUBE(I 1=0429	IRUNG(I 1=042A	IS(I 1=042B		

STATEMENT ALLOCATIONS																			
200	=0447	201	=0469	202	=0479	203	=049C	204	=04A3	8	=0557	9	=056C	10	=05C2	11	=05D5	12	=05E6
11	=05EA	18	=05FF	21	=06A0	22	=06A6	23	=06AC	24	=06C0	30	=06CO						

FEATURES SUPPORTED
 ONE WORD INTEGERS

CALLD SUPPROGRAMS

SRM34	SRM22	CDRE	FSORT	SRM32	FATAN	FCOS	FAXB	FADD	FADDX	FSUB	FSUBX	FMPY	FDIV	FLD
FLDX	FSTO	FSTOX	FSBR	FDVR	FAXI	FLOAT	SWRT	SCOMP	SI0F	SI0I	SUBSC	SUBIN		

REAL CONSTANTS

.000000E 00=0430	.500000E 00=0432	.240000E 02=0434	.100000E 01=0436	.110000E 01=0438	.1C3000E 02=043A
.100000E 03=043C	.314159E 01=043E	.200000E 01=0440			

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INTEGER CONSTANTS 3=0442 1=0443 0=0444 4=0445 2=0446

CORE REQUIREMENTS FOR SBM31
COMMON 0 VARIABLES 1072 PROGRAM 688

RELATIVE ENTRY POINT ADDRESS IS 04CC (HEX)

END OF COMPILATION

// DUP

#STORE MS UA SBM31
CART ID 0305 DB ADDR 5AAB DB CNT 0030

// EJECT

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// FOR
*ONE WORD INTEGERS
*IOCS(CARD)
*IOCS(1132 PRINTER)
*LIST ALL

C ESTABLISH PARAMETERS FOR STEP BY STEP INTEGRATION AND HALF
C INTERVAL METHOD
C GALFT AND GARIT ARE LOWER AND UPPER BOUNDS RESPECTIVELY FOR
C STARTING LOCATION OF PARTICLE

C SIGNAL IS SIGN CORRECTION FOR HALF INTERVAL METHOD

C IN THIS PROBLEM IT IS -1.0

C DTAU IS THE TIME INCREMENT FOR STEP BY STEP INTEGRATION

C NIRP IS INTERVAL OF WRITTEN PARTICLE PATHS

C NSRP IS INTERVAL OF WRITTEN INTEGRATION STEPS

C NX IS NUMBER OF ITERATIONS USED TO DETERMINE PARTICLE STARTING
C POINT

C ESTABLISH WHETHER A TUBE IS PRESENT

C IF ITURE IS POSITIVE A TUBE IS PRESENT

C IF ITURE IS ZERO OR NEGATIVE SPHERE IS IN FREE SPACE

C 1 READ(2,1001)GALFT,GARIT,SIGML,DTAU,NIRP,NSRP,NX,ITURE

C ESTABLISH PHYSICAL PROPERTIES FOR CALCULATING IMPACTION EFFICIENCY

C DC IS SPHERE DIAMETER, CM

C DP IS PARTICLE DIAMETER, CM

C RHO IS FLUID DENSITY, GM/CC

C SIGMA IS PARTICLE DENSITY, GM/CC

C XMU IS ABSOLUTE VISCOSITY OF FLUID, POISE

C UF IS FREE STEAM VELOCITY, CM/SEC

C READ(2,1011)DC,DP,RHO,SIGMA,XMU,UF

C XL IS UPSTREAM STARTING POSITION FOR CALCULATING PARTICLE MOTION

C READ(2,1021)XL

C REFZ=RHO*DP*UF/XMU

C XK=SIGMA*DP**2*UF/19.*XMU*DC)

C P = 9.*RHO**2*UF*DC/(XMU*SIGMA)

C WRITE(3,200)

C WRITE(3,201)

C C XL IS INERTIAL PARAMETER CALLED STOKES NUMBER

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C REZ IS FREE STREAM REYNOLDS NUMBER OF SPHERICAL PARTICLE
C P IS DIMENSIONLESS GROUP INDEPENDENT OF DROP SIZE
C

C WRITE(3,202)XL,REZ,XK,P,DC,DP,RHO,SIGMA,XNU,U,F,IITUBE

C CALCULATE STARTING POSITION OF TANGENT PARTICLE

C CALL SBM30(GALFT,GARIT,SIGNL,DIAU,NIBP,NSBP,NX,XL,XK,REZ,G4ZER)

C ESTABLISH NUMBER OF INCREMENTS FOR INTEGRATING PARTICLE FORCE

C NZER IS NUMBER OF INCREMENTS

C NJBP IS INTERVAL OF WRITTEN INCREMENTS

C READ(2,103)NZER,NJBP

C CALCULATE FORCE PRODUCED BY PARTICLES

C CALL SBM31(XL,G4ZER,DIAU,XK,REZ,PA,NZER,NJBP)

C RECYL=RHO*DC*U*F/XMU

C CDAIR = CDRE(RECYL)/RECYL

C OUTPUT RESULTS

C WRITE(3,204)RECYL

C IF(RECYL - 10000)10,11,11

10 CONTINUE

11 WRITE(3,206)CDAIR

READ(2,111)NSTOP

IF(NSTOP)1,30,30

30 CALL EXIT

C FORMATS FOR INPUT AND OUTPUT STATEMENTS

C 100 FORMAT(10X,F10.7,20X,F10.7,19X,F3.0Y 10X,F10.7,20X,I3,27X,I2/

1 10X,I2,10X,I2)

101 FORMAT(F10.5,F10.7,F10.6, F10.6,F10.7,F10.1)

102 FORMAT (F10.7)

103 FORMAT(I5,I5)

111 FORMAT(I5)

200 FORMAT(1H1, 35X, 49HTHE DRAG ON SPHERES IN A STREAM OF DUST-LADEN

1 AIR / 1H0)

201 FORMAT(28H0THE PHYSICAL PARAMETERS ARE)

202 FORMAT(10H0XL = F10.7/10H REZ = E12.6/10H XK = ,

1 E12.6/10H P = E12.6/

PAGE 21

2 10H0DC = , F10.5/10H DP = , F10.7/10H RMO = , F10.6/
3 10H SIGMA = , F10.6/10H XMU = , F10.7/10H UF = , F10.1/
4 10H I:TURE = , I3)
204 FORMAT(32H0THE TARGET REYNOLDS NUMBER IS ,E10.4)
206 FORMAT(59H0THE DRAG COEFFICIENT OF THE SPHERE DUE TO THE AIR ALON
IE IS,9.4)
END

VARIABLE ALLOCATIONS

G4LFT(R)=0090 G4RIT(P)=0002 SIGNL(R)=0004 DTAU(R)=0006 DC(R)=0008 DP(R)=000A
PHO(R)=000C SIGMA(R)=000E XMU(R)=0010 UF(R)=0012 XL(R)=0014 REZ(R)=0016
XK(R)=0018 P(R)=001A G4ZER(R)=001C PAIR)=0020 RECYL(R)=0020 CDAIR(R)=0022
NIRP(I)=0026 MSBP(I)=0027 NX(I)=0028 ITUBE(I)=002A NZER(I)=002B NJBP(I)=002B
NSTOP(I)=002C

STATEMENT ALLOCATIONS

100 =0013 101 =0046 102 =004D 103 =004F 111 =0052 200 =0054 201 =0075 202 =0085 204 =00DD 206 =00F0
1 =0178 10 =01E0 11 =01E6 30 =01EF

FEATURES SUPPORTED
ONE WORD INTEGERS
IOCS

CALLED SURPROGRAMS

SRM30 SRM31 CDRE FMPY FDIV FLD FSTO FSBR FDVR FAXI FLOAT CARDZ PRNTZ SRED SMKT
SCOMP SF10 SIOF SIOI

REAL CONSTANTS

.900000E 01=002E

INTEGER CONSTANTS

2=001C 3=0031 10000=0032

CORE REQUIREMENTS FOR

COMMON 0 VARIABLES 46 PROGRAM 450

END OF COMPILATION

// XEO

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THE DRAG ON SPHERES IN A STREAM OF DUST-LADEN AIR

THE PHYSICAL PARAMETERS ARE

XL = -5.0000009
REZ = 0.795665E 02
XK = 0.339462E 03
P = 0.159429E 02
DC = 0.80000
DP = 0.0155000
RHO = 0.001213
SIGMA = 2.600000
XMU = 0.0001789
UF = 700.0
ITURE = 1

IMPACTION EFFICIENCY OF A SPHERE

GALEF = 0.000000
GARIT = 1.000000
SIGML = -1.
OTAU = 0.100000
NIRP = 10
NSHP = 2
NX = 20

THE MOTION OF A CRITICAL PARTICLE IS GIVEN BY

ITER = 1
 G4LEF = 0.000000
 G4TER = 0.500000
 G4PIT = 1.000000

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0.0000	0.992235	0.001170	-5.000000	0.500000	0.9922	0.0011	24.0000
0.2000	0.992235	0.001170	-4.801552	0.500234	0.9912	0.0013	24.1144
0.4000	0.992234	0.001170	-4.603104	0.500468	0.9901	0.0016	24.3144
0.6000	0.992232	0.001171	-4.404656	0.500702	0.9887	0.0019	24.6838
0.8000	0.992230	0.001171	-4.206209	0.500936	0.9871	0.0023	25.0545
1.0000	0.992226	0.001172	-4.007762	0.501170	0.9851	0.0028	25.5006
1.2000	0.992221	0.001173	-3.809317	0.501405	0.9828	0.0034	26.0393
1.4000	0.992215	0.001175	-3.610873	0.501639	0.9799	0.0042	26.6888
1.6000	0.992206	0.001177	-3.412430	0.501874	0.9764	0.0052	27.4637
1.8000	0.992194	0.001180	-3.213989	0.502110	0.9719	0.0066	28.3636
2.0000	0.992179	0.001184	-3.015552	0.502346	0.9664	0.0085	29.3535
2.2000	0.992158	0.001190	-2.817118	0.502584	0.9592	0.0110	30.3524
2.4000	0.992131	0.001198	-2.618688	0.502822	0.9500	0.0146	31.3039
2.6000	0.992096	0.001210	-2.420265	0.503063	0.9379	0.0197	32.5793
2.8000	0.992046	0.001227	-2.221849	0.503307	0.9216	0.0273	35.1227
3.0000	0.991977	0.001253	-2.023447	0.503555	0.8993	0.0387	37.8687
3.2000	0.991878	0.001294	-1.825060	0.503809	0.8683	0.0566	41.2151
3.4000	0.991733	0.001361	-1.626698	0.504074	0.8241	0.0858	45.4350
3.6000	0.991511	0.001477	-1.428372	0.504357	0.7601	0.1354	50.8992
3.8000	0.991162	0.001689	-1.230102	0.504672	0.6665	0.2240	58.1905
4.0000	0.990588	0.002105	-1.031924	0.505048	0.5317	0.3305	68.2592
4.1599	0.989818	0.002824	-0.873486	0.505434	0.4024	0.6495	87.4018
4.1699	0.989753	0.002896	-0.863587	0.505462	0.3930	0.6701	89.2955

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ITER = 10
 GALEF = 0.986328
 G4ZFR = 0.987304
 G4RIT = 0.988281

TAU	G(1)	G(2)	G(3)	G(4)	UX	UY	CDRE
0.0000	0.992871	0.002153	-5.000000	0.987304	0.9928	0.0021	24.0000
0.2000	0.992871	0.002153	-4.601424	0.987735	0.9920	0.0025	24.1050
0.4000	0.992870	0.002153	-4.602849	0.988165	0.9910	0.0029	24.2421
0.6000	0.992869	0.002154	-4.404275	0.988596	0.9899	0.0034	24.6219
0.8000	0.992866	0.002155	-4.205701	0.989027	0.9885	0.0041	24.9539
1.0000	0.992863	0.002156	-4.007127	0.989458	0.9870	0.0049	25.3495
1.2000	0.992859	0.002158	-3.808554	0.989889	0.9851	0.0059	25.8227
1.4000	0.992854	0.002161	-3.609982	0.990321	0.9829	0.0072	26.3884
1.6000	0.992847	0.002164	-3.411412	0.990754	0.9802	0.0089	27.0605
1.8000	0.992838	0.002169	-3.212843	0.991187	0.9771	0.0111	27.8458
2.0000	0.992826	0.002176	-3.014276	0.991622	0.9732	0.0139	28.7317
2.2000	0.992811	0.002185	-2.815712	0.992058	0.9686	0.0176	29.6714
2.4000	0.992792	0.002198	-2.617151	0.992496	0.9629	0.0226	30.5895
2.6000	0.992767	0.002214	-2.418595	0.992937	0.9561	0.0294	31.4721
2.8000	0.992737	0.002238	-2.220044	0.993382	0.9478	0.0388	32.3164
3.0000	0.992696	0.002272	-2.021500	0.993833	0.9380	0.0519	35.1330
3.2000	0.992644	0.002320	-1.822965	0.994292	0.9267	0.0704	37.5175
3.4000	0.992577	0.002392	-1.624443	0.994763	0.9145	0.0966	40.3086
3.6000	0.992492	0.002500	-1.425935	0.995252	0.9032	0.1338	43.5836
3.8000	0.992388	0.002664	-1.227447	0.995767	0.8975	0.1858	47.4096
4.0000	0.992271	0.002912	-1.028981	0.996324	0.9066	0.2551	51.8058
4.2000	0.992165	0.003283	-0.830537	0.996941	0.9475	0.3376	56.6614
4.4000	0.992129	0.003809	-0.632109	0.997649	1.0426	0.4116	61.6017
4.6000	0.992268	0.004474	-0.433672	0.998476	1.2030	0.4248	65.8702
4.8000	0.992700	0.005142	-0.235180	0.999439	1.3892	0.3089	68.5349
5.0000	0.993434	0.005561	-0.036569	1.000514	1.4962	0.0545	69.3931
5.0299	0.993572	0.005570	-0.006764	1.000681	1.5277	0.0103	71.6809
5.0359	0.993600	0.005571	-0.000803	1.000714	1.5378	0.0012	71.6874

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ITER = ZC
 GALEF = 0.986612
 GAZFR = 0.986613
 GARIT = 0.986614

TAU	G(1)	G(2)	G(3)	G(4)	UX	UY	CDRE
0.0000	0.992870	0.002152	-5.000000	0.986613	0.9928	0.0021	24.0000
0.2000	0.992869	0.002152	-4.801426	0.987043	0.9920	0.0025	24.1050
0.4000	0.992869	0.002152	-4.602852	0.987474	0.9910	0.0029	24.3421
0.6000	0.992867	0.002153	-4.404278	0.987904	0.9899	0.0034	24.6220
0.8000	0.992865	0.002154	-4.205704	0.988335	0.9885	0.0041	24.9540
1.0000	0.992862	0.002155	-4.007130	0.988765	0.9870	0.0049	25.3498
1.2000	0.992858	0.002157	-3.808557	0.989197	0.9851	0.0059	25.8231
1.4000	0.992853	0.002159	-3.609985	0.989628	0.9829	0.0072	26.3889
1.6000	0.992845	0.002163	-3.411415	0.990061	0.9802	0.0089	27.0612
1.8000	0.992836	0.002168	-3.212846	0.990494	0.9771	0.0111	27.8467
2.0000	0.992825	0.002175	-3.014279	0.990928	0.9732	0.0139	28.7328
2.2000	0.992810	0.002184	-2.815715	0.991364	0.9686	0.0176	29.6726
2.4000	0.992790	0.002196	-2.617155	0.991802	0.9629	0.0226	30.5908
2.6000	0.992766	0.002213	-2.418598	0.992243	0.9561	0.0294	31.4736
2.8000	0.992735	0.002237	-2.220047	0.992687	0.9478	0.0388	32.4702
3.0000	0.992695	0.002270	-2.021504	0.993138	0.9380	0.0519	35.1374
3.2000	0.992643	0.002319	-1.822969	0.993597	0.9266	0.0704	37.5233
3.4000	0.992575	0.002391	-1.624447	0.994067	0.9143	0.0966	40.3164
3.6000	0.992490	0.002499	-1.425940	0.994556	0.9031	0.1339	43.5941
3.8000	0.992386	0.002663	-1.227452	0.995071	0.8973	0.1859	47.4243
4.0000	0.992268	0.002912	-1.028986	0.995627	0.9064	0.2553	51.8262
4.2000	0.992162	0.003283	-0.830543	0.996245	0.9472	0.3381	56.6896
4.4000	0.992125	0.003810	-0.632116	0.996952	1.0424	0.4123	61.6398
4.6000	0.992264	0.004476	-0.433679	0.997780	1.2031	0.4258	65.9190
4.8000	0.992097	0.005147	-0.235187	0.998744	1.3899	0.3098	68.5916
5.0000	0.993539	0.005525	-0.026640	0.999873	1.5374	0.0409	71.7265
5.0280	0.993629	0.005529	-0.007762	0.999977	1.5388	0.0119	71.7420

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TAU	VX	VY	X	Y	UX	UY	FY
3.9319	1.014362	0.000300	-1.001004	0.000000	0.0019	0.0000	2.06391
3.9329	1.014354	0.000144	-0.999986	0.024930	0.0006	0.0383	2.06261
3.9339	1.014364	0.000338	-0.998960	0.058170	0.0060	0.0892	2.05690
3.9359	1.014378	0.000531	-0.996900	0.091410	0.0149	0.1395	2.04662
3.9399	1.014386	0.000726	-0.992805	0.124651	0.0245	0.1900	2.03174
3.9449	1.014403	0.000922	-0.987683	0.157893	0.0378	0.2399	2.01229
3.9499	1.014435	0.001117	-0.982546	0.191134	0.0577	0.2878	1.98830
3.9569	1.014465	0.001314	-0.975374	0.224377	0.0786	0.3356	1.95971
3.9649	1.014502	0.001512	-0.967173	0.257621	0.1033	0.3820	1.92655
3.9749	1.014538	0.001714	-0.956930	0.290867	0.1295	0.4283	1.88880
3.9849	1.014594	0.001912	-0.944678	0.324113	0.1620	0.4710	1.84651
3.9969	1.014649	0.002114	-0.934380	0.357362	0.1963	0.5132	1.79862
4.0109	1.014813	0.002230	-0.920041	0.390612	0.2329	0.5547	1.74860
4.0239	1.014887	0.002425	-0.906702	0.423860	0.2770	0.5892	1.69238
4.0399	1.014956	0.002631	-0.890306	0.457115	0.3218	0.6242	1.63175
4.0579	1.015036	0.002838	-0.871872	0.490374	0.3695	0.6574	1.56653
4.0769	1.015130	0.003041	-0.852407	0.523634	0.4214	0.6860	1.49676
4.0979	1.015234	0.003245	-0.830898	0.556900	0.4766	0.7117	1.42239
4.1209	1.015395	0.003351	-0.807344	0.590168	0.5352	0.7336	1.34361
4.1449	1.015517	0.003554	-0.782760	0.623440	0.5974	0.7489	1.26008
4.1719	1.015653	0.003759	-0.755115	0.656722	0.6628	0.7609	1.17195
4.2019	1.015804	0.003964	-0.724405	0.690015	0.7318	0.7681	1.07922
4.2349	1.015970	0.004086	-0.690093	0.723312	0.8045	0.7690	0.98198
4.2699	1.016153	0.004285	-0.654810	0.756620	0.8796	0.7606	0.88008
4.3109	1.016330	0.004425	-0.612877	0.789990	0.9601	0.7463	0.77356
4.3549	1.016556	0.004625	-0.567873	0.823292	1.0424	0.7195	0.66241
4.4049	1.016812	0.004815	-0.516742	0.856660	1.1279	0.6802	0.54663
4.4639	1.017063	0.004970	-0.456439	0.890066	1.2180	0.6248	0.42618
4.5349	1.017361	0.005144	-0.383902	0.923532	1.3118	0.5456	0.30099
4.6269	1.017753	0.005334	-0.289958	0.957102	1.4093	0.4272	0.17098
4.7789	1.018412	0.005533	-0.134866	0.990992	1.5107	0.2057	0.03570
4.8979	1.018916	0.005627	-0.013958	0.999916	1.5386	0.0208	0.00019

THE DRAG COEFFICIENT OF THE SPHERE DUE TO THE REFLECTION OF PARTICLES IS 2.0356

THE TARGET REYNOLDS NUMBER IS 0.3796E 04

THE DRAG COEFFICIENT OF THE SPHERE DUE TO THE AIR ALONE IS 0.3909
 // PAUS

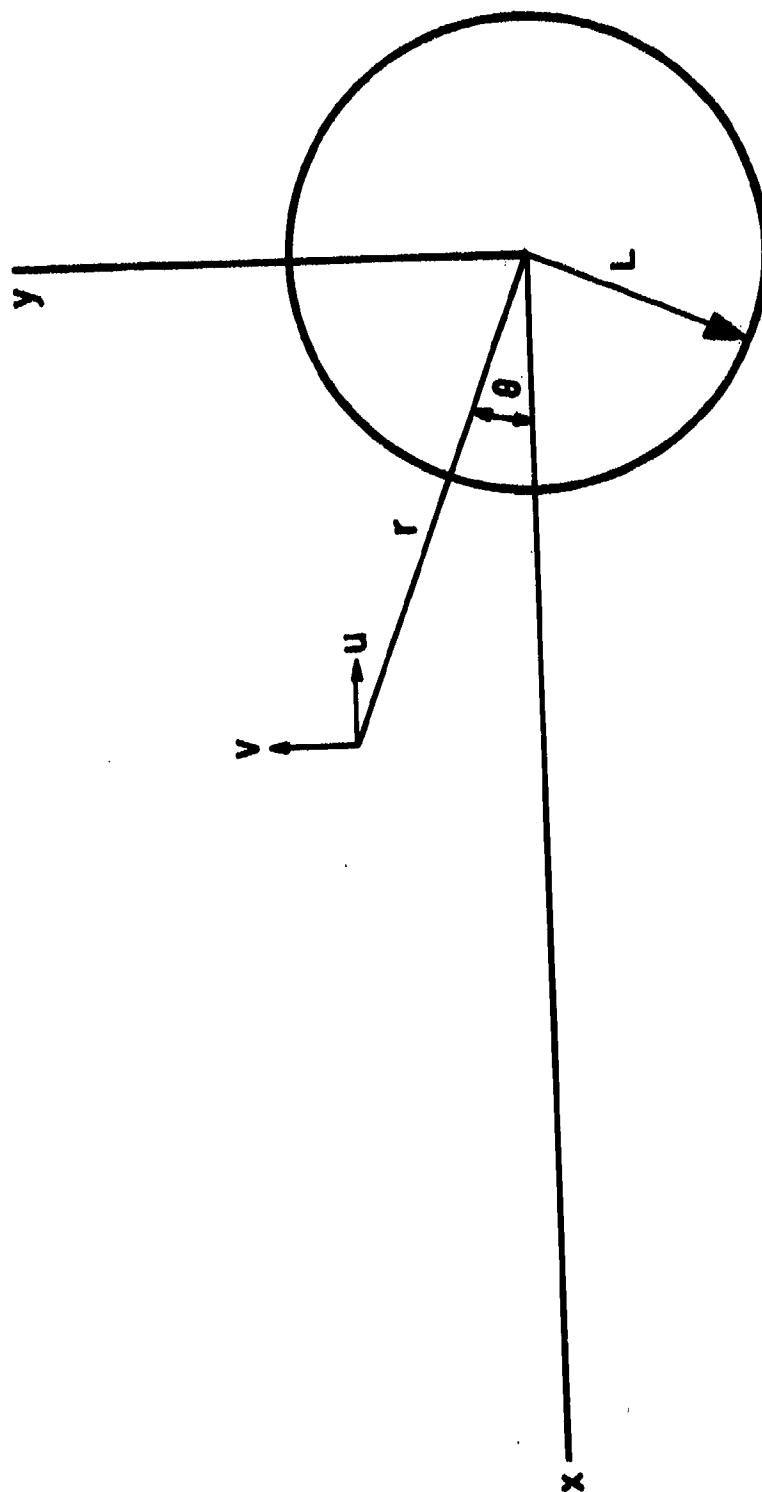


Fig. 1: Co-ordinate System for Transverse Flow with respect to a Circular Cylinder.

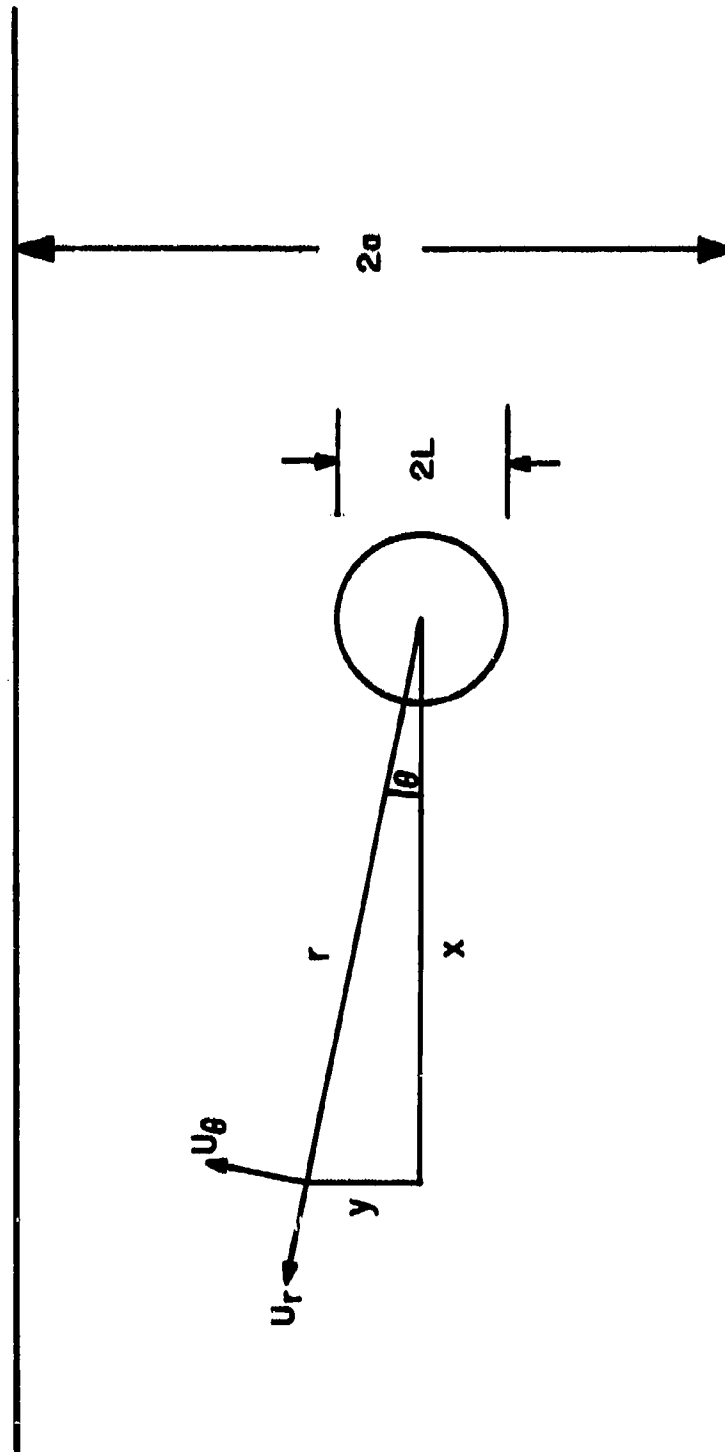


Fig. 2: Co-ordinate System for Flow around a Sphere in a Circular Cylinder.

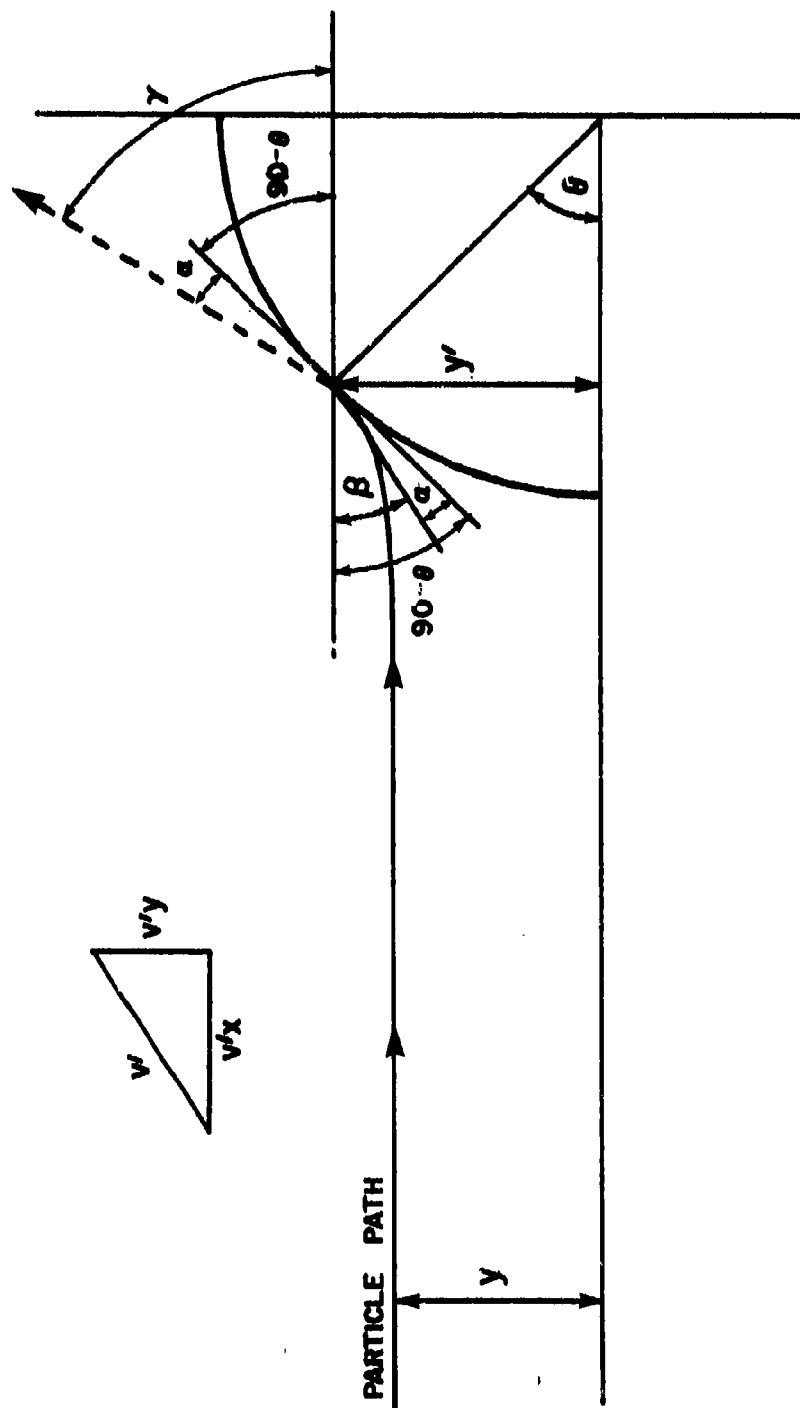


Fig. 3: Velocity Change of a Particle due to Interaction with a Sphere or Cylinder.

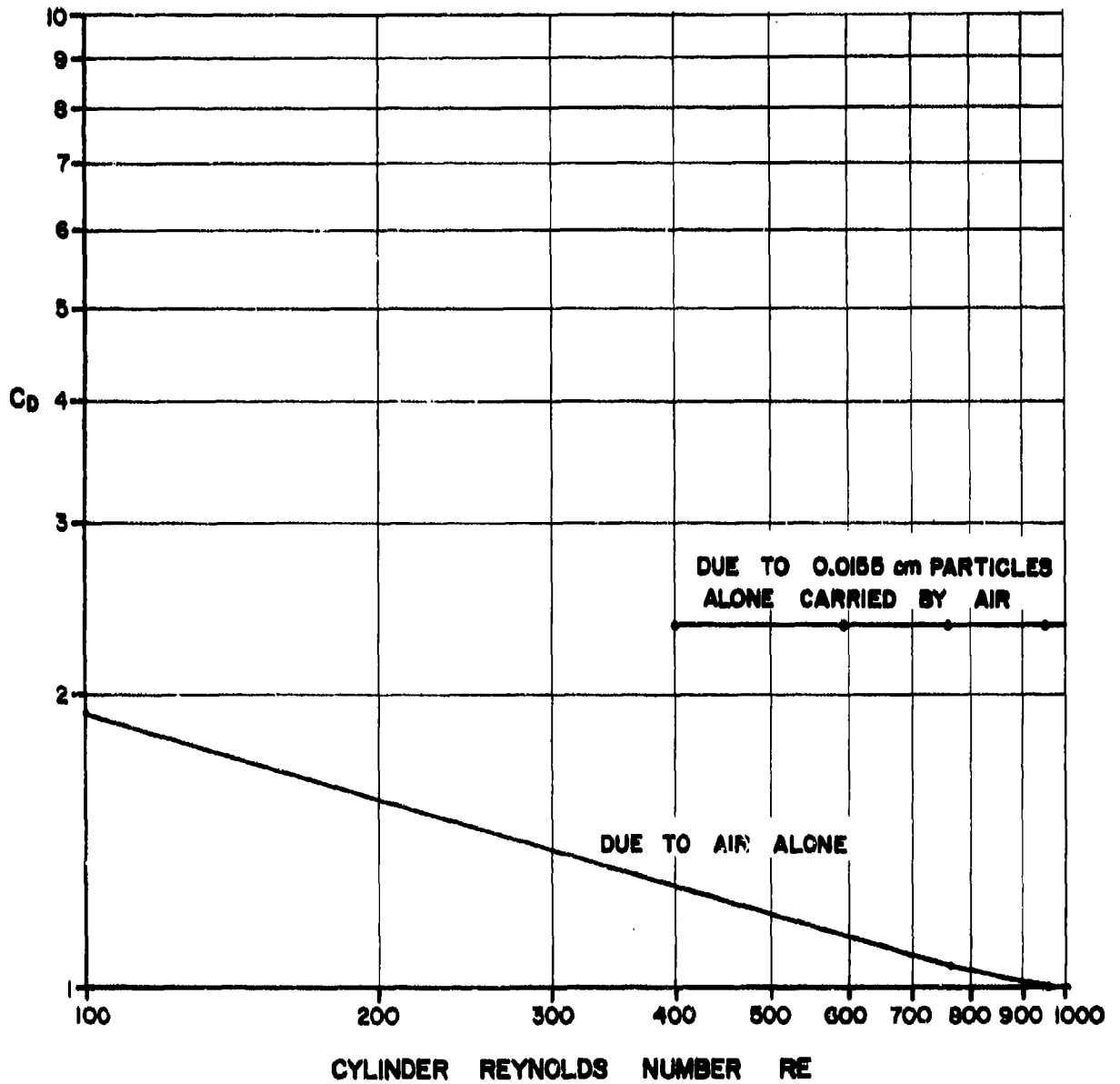


Fig. 4: Variation of Drag Coefficient for Cylinders.

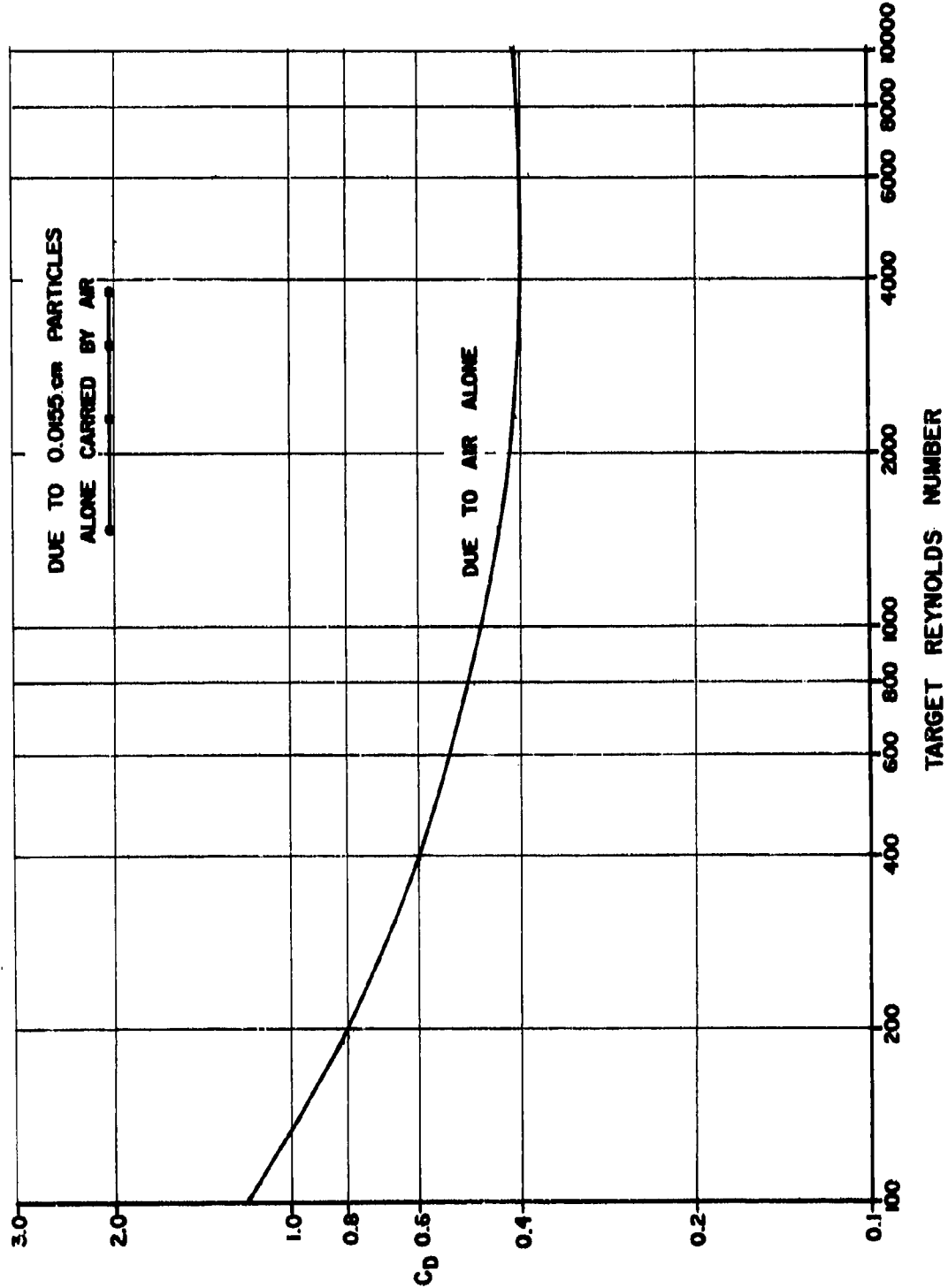
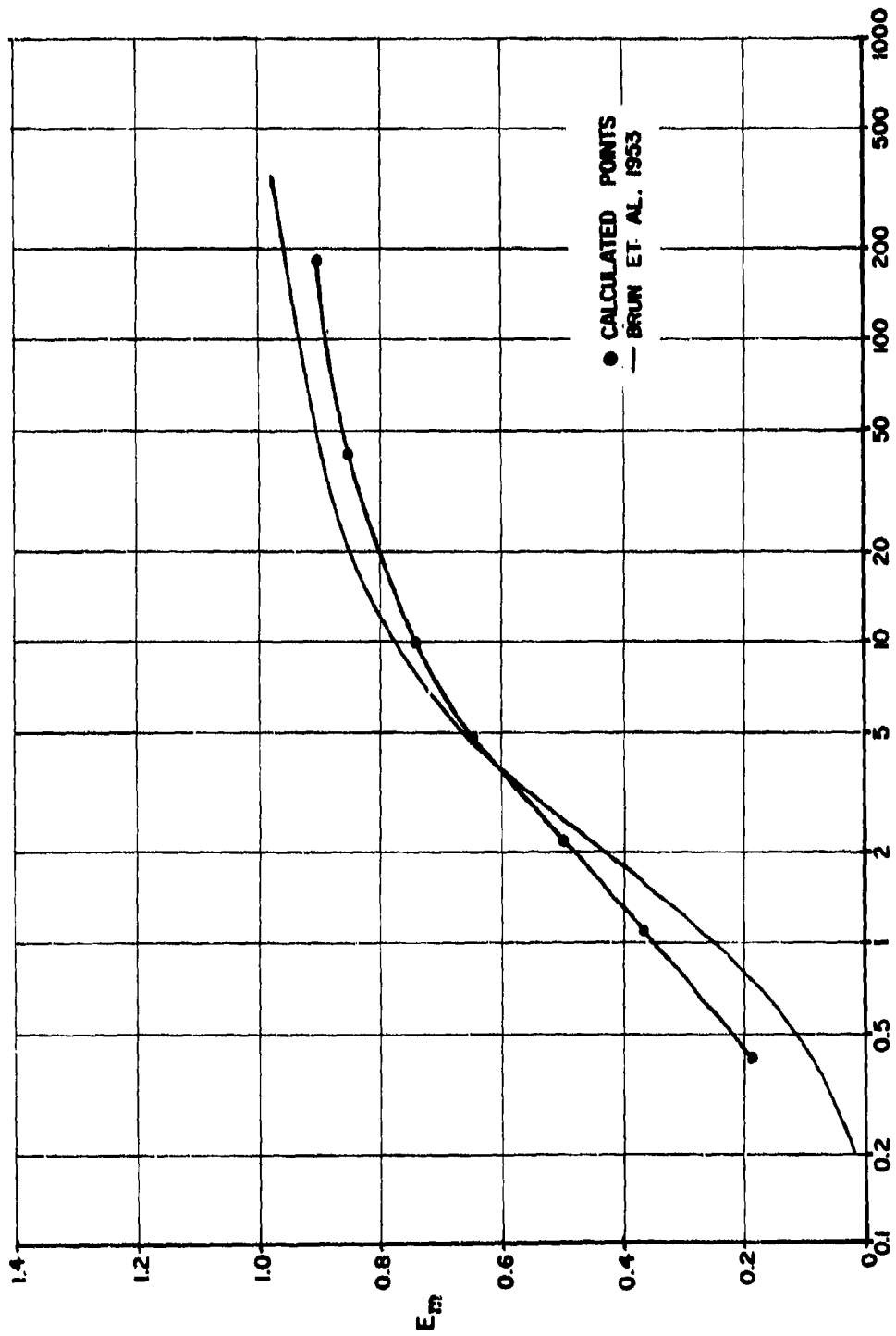


Fig. 5: Variation of Drag Coefficient for a Spherical Target.



INERTIA PARAMETER $K = \sigma^2 U^2 / 18 \mu L$

Fig. 6: Impaction Efficiency for Cylinders, $\phi = 1000$.

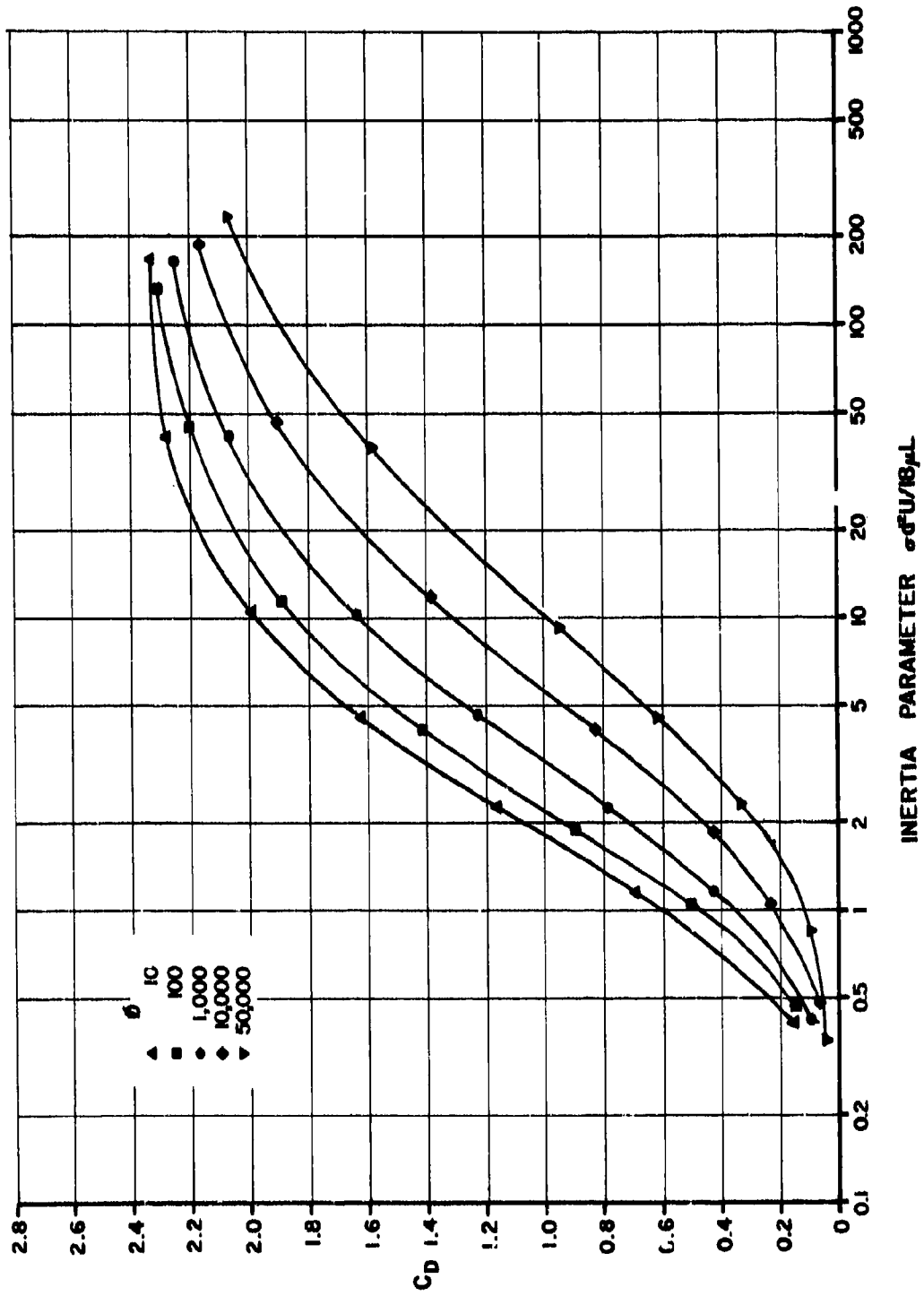


Fig. 7: Calculated Drag Coefficient for Cylinders due to Particles alone in an Inviscid Flow.

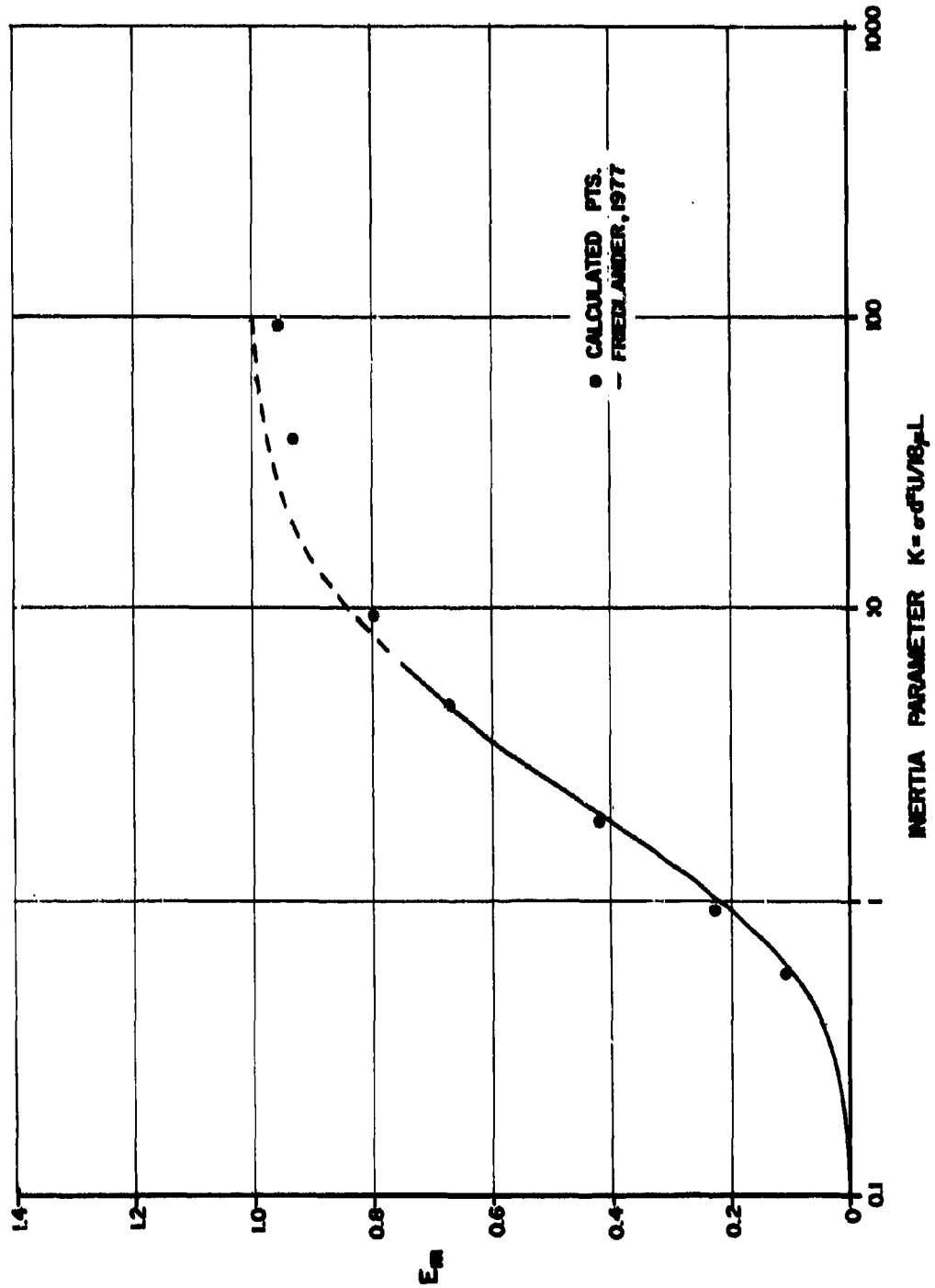
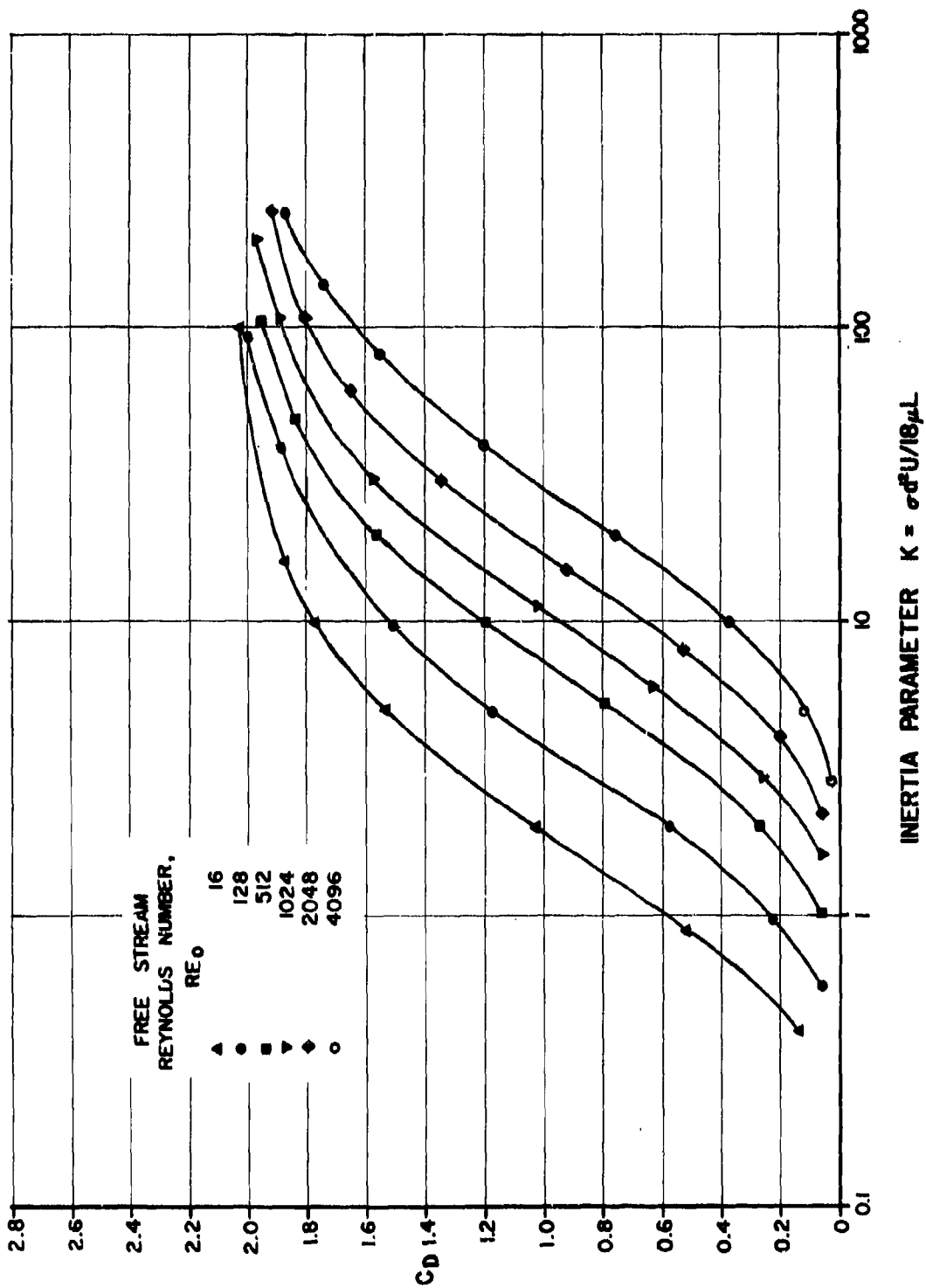


Fig. 8: Impaction Efficiency for Spheres for Free Stream Particle Reynolds Number of 128.



INERTIA PARAMETER $K = \sigma d^2 U / 18 \mu L$

Fig. 9: Calculated Drag Coefficient for Spheres due to Particles alone in an Inviscid Flow.

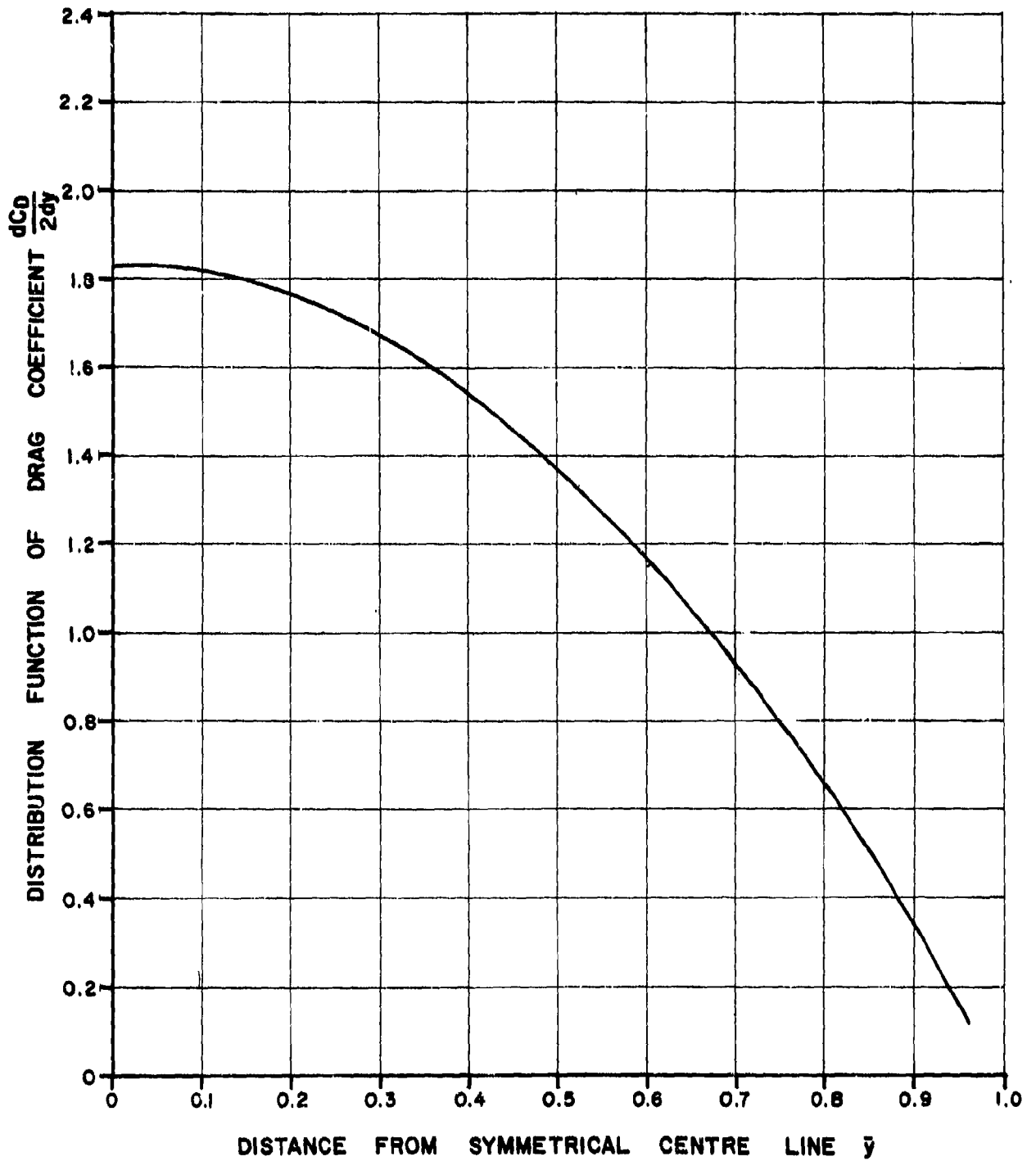


Fig. 10: Distribution of Forces due to Elastic Reflection of Particles from a Cylinder.

$Re_0 = 73.6$ and $K = 339$

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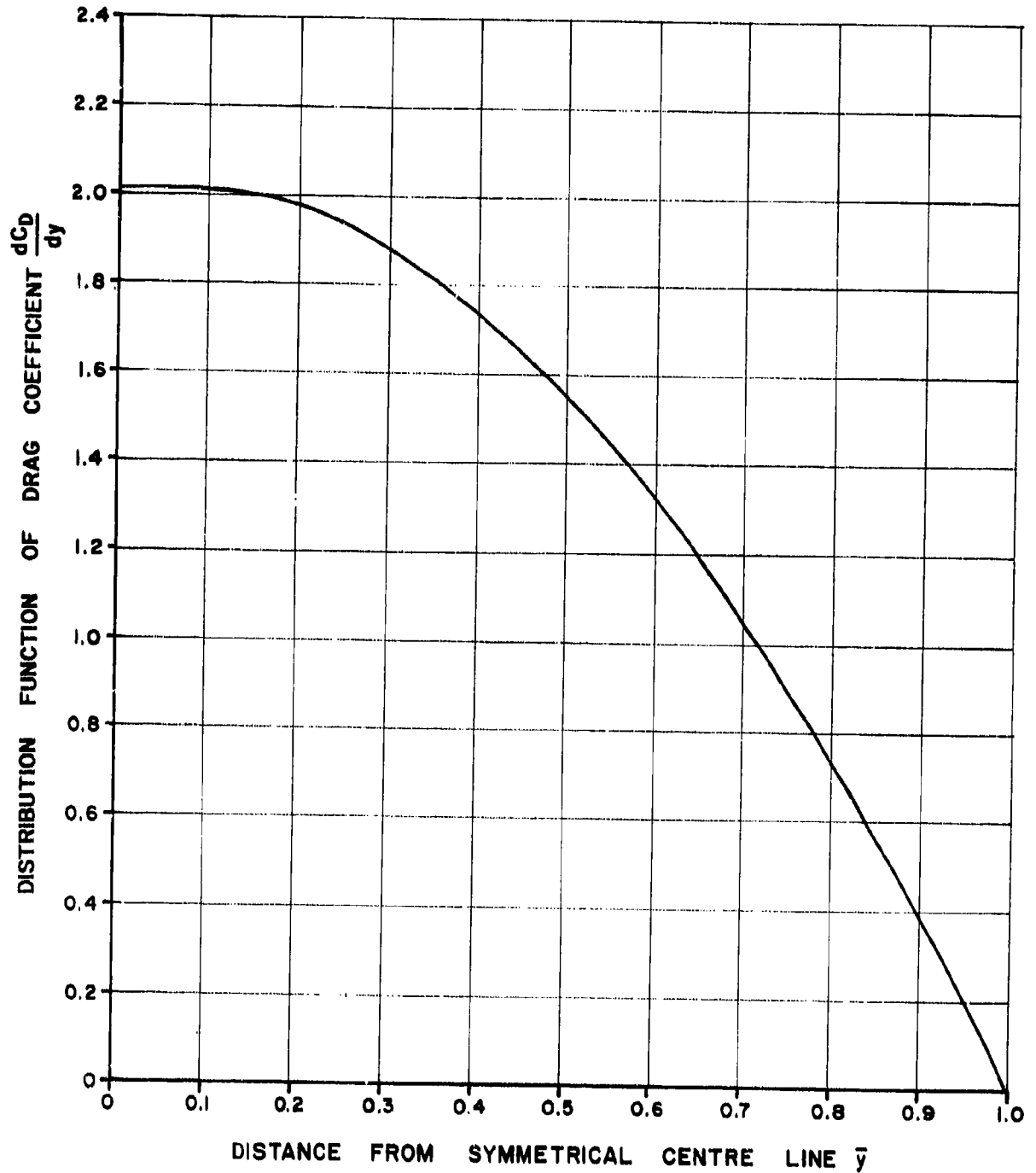


Fig. 11: Distribution of Forces due to Elastic Reflection of Particles from a Sphere in a Tube.

$Re_0 = 73.6$, $K = 339$ and $L = 0.3150 a$

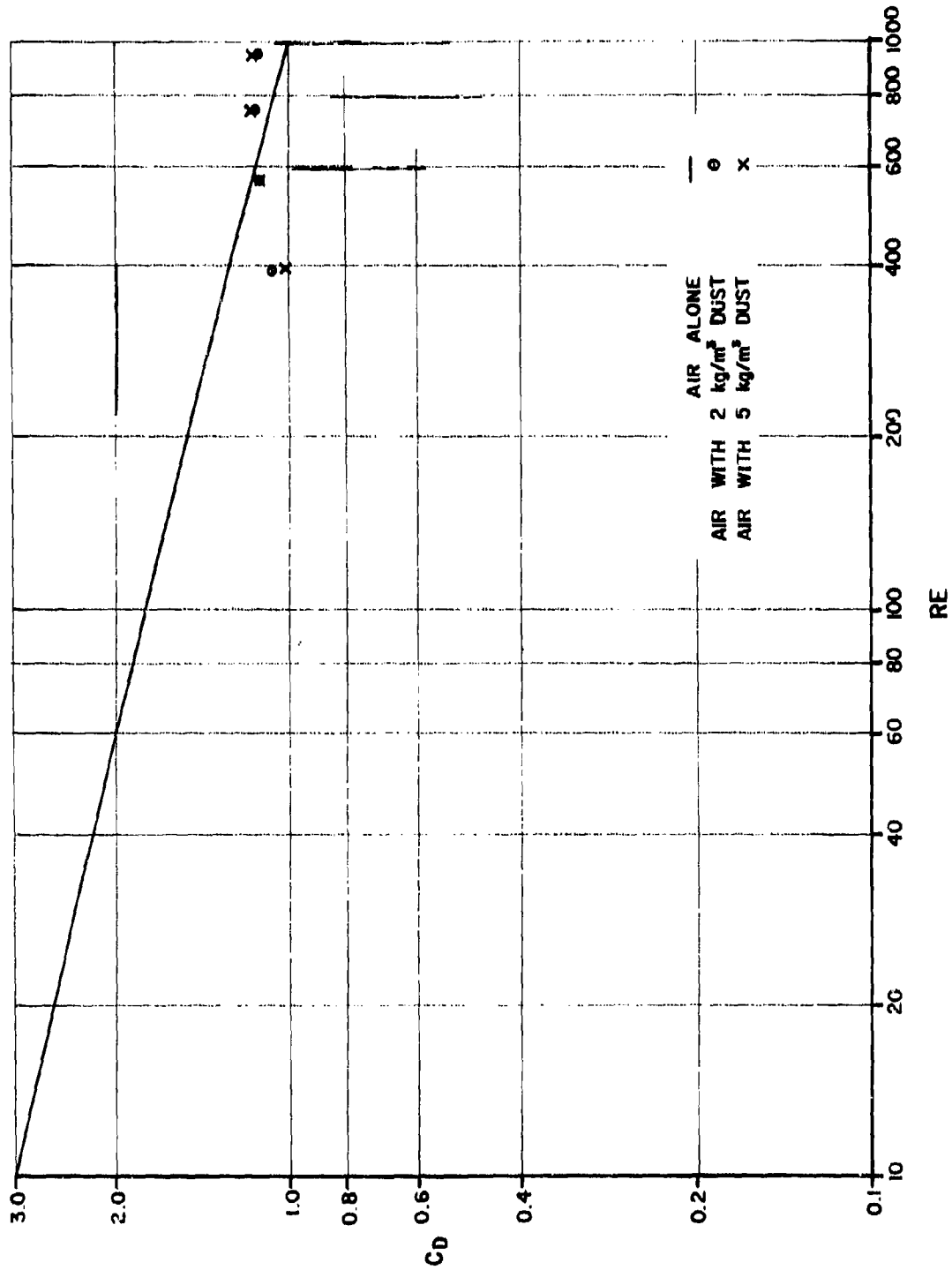


Fig. 12: Drag Coefficient of a 0.2 cm Cylinder in a 2.54 cm Tube due to Equal Masses of 0.0470, 0.0155 and 0.0055 cm Particles corrected for Particle Velocity Deficiency.

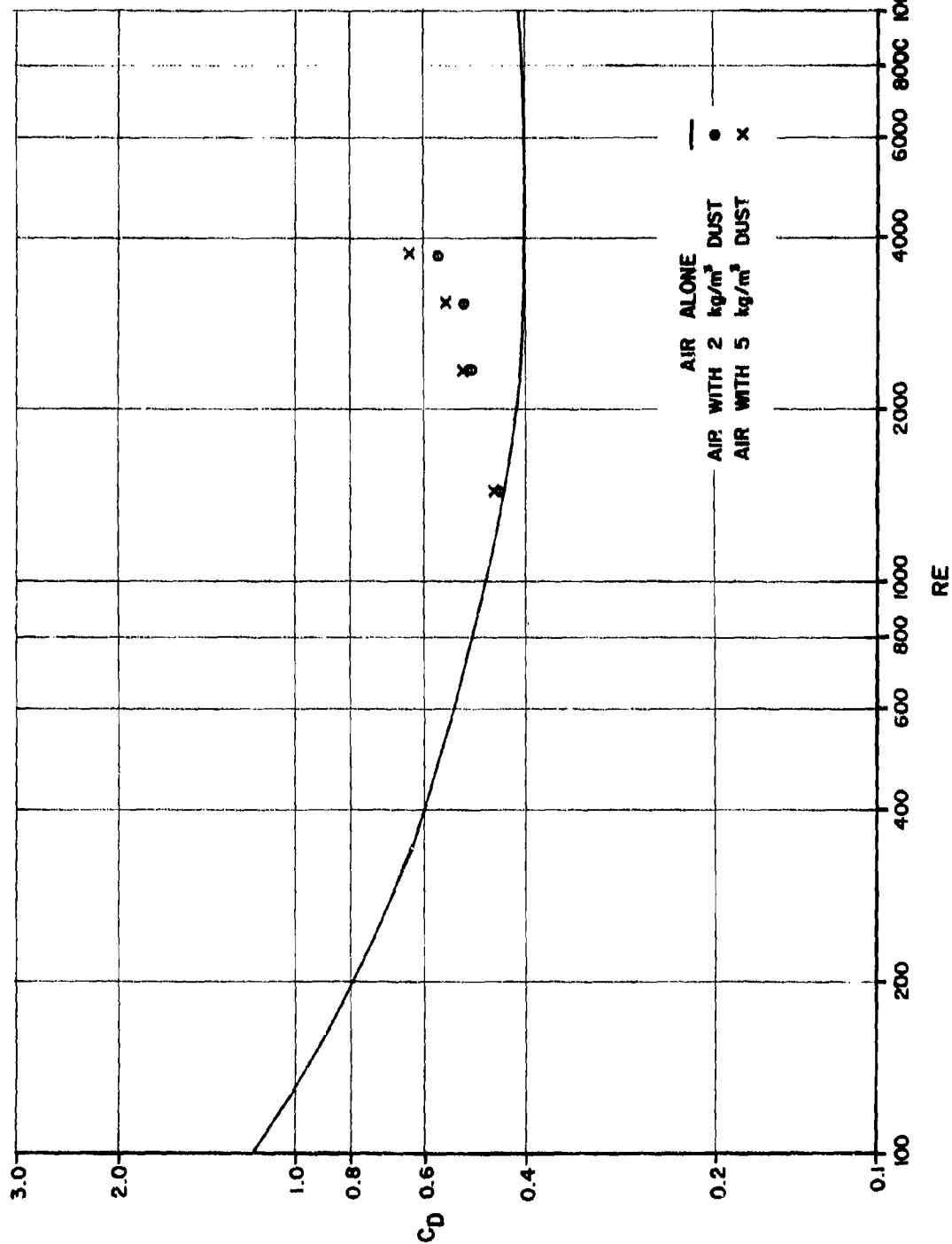


Fig. 13: Drag Coefficient of a 0.8 cm Sphere in a 2.54 cm Tube due to Equal Masses of 0.0470, 0.0155 and 0.0055 cm Particles corrected for real Fluid Velocity and Particle Velocity Deficiency.

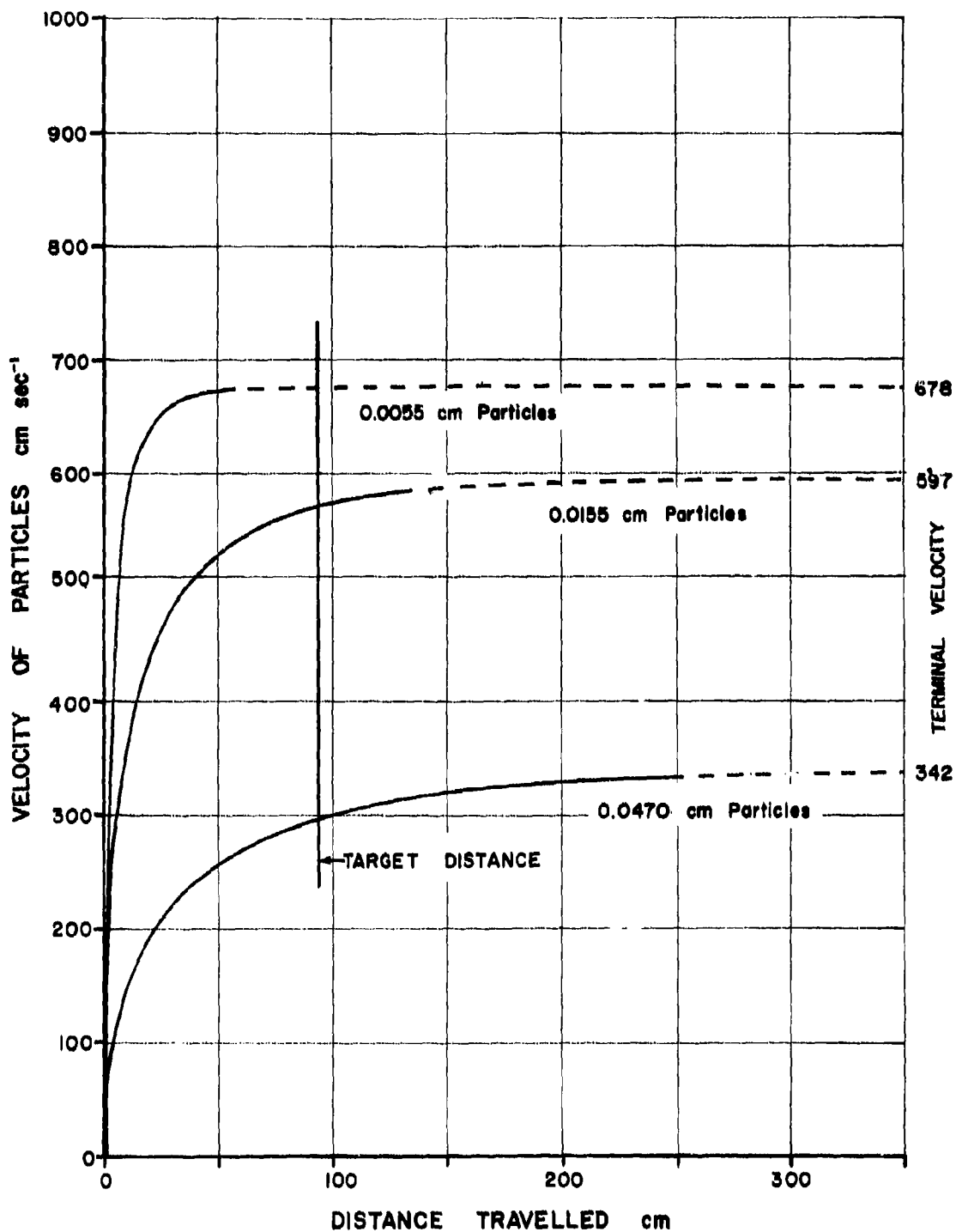


Fig. 14: Upward Velocity of Spherical Dust Particles accelerated from rest in a vertical Air Stream of 700 cm sec⁻¹ Velocity

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	2b. GROUP	
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4. DESCRIPTIVE NOTES (Type of report and inclusive dates) Technical Paper		
5. AUTHOR(S) (Last name, first name, middle initial) Mellsen, Stanley B.		
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13. ABSTRACT <p>The effect of dust on aerodynamic drag of spheres and cylinders was calculated by using a mathematical model developed for this purpose. The results were compared to experiments previously done by other workers. The calculated and experimental results agree favourably, showing that the mathematical model is satisfactory. Impaction efficiencies and drag coefficients due to dust alone were then obtained using the model for a wide range of the inertia parameter and the results are presented graphically. The model can also be used for calculating velocity distributions and points of impact for a stream of airborne particles flowing over a sphere or cylinder.</p> <p style="text-align: right;">(U)</p>		

KEY WORDS

Aerodynamic Drag
Cylindrical Bodies
Spheres
Blast Loads
Dust Particles
Drops (Liquid)
Impact
Aerosols - Penetration
Aerosols - Sampling

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