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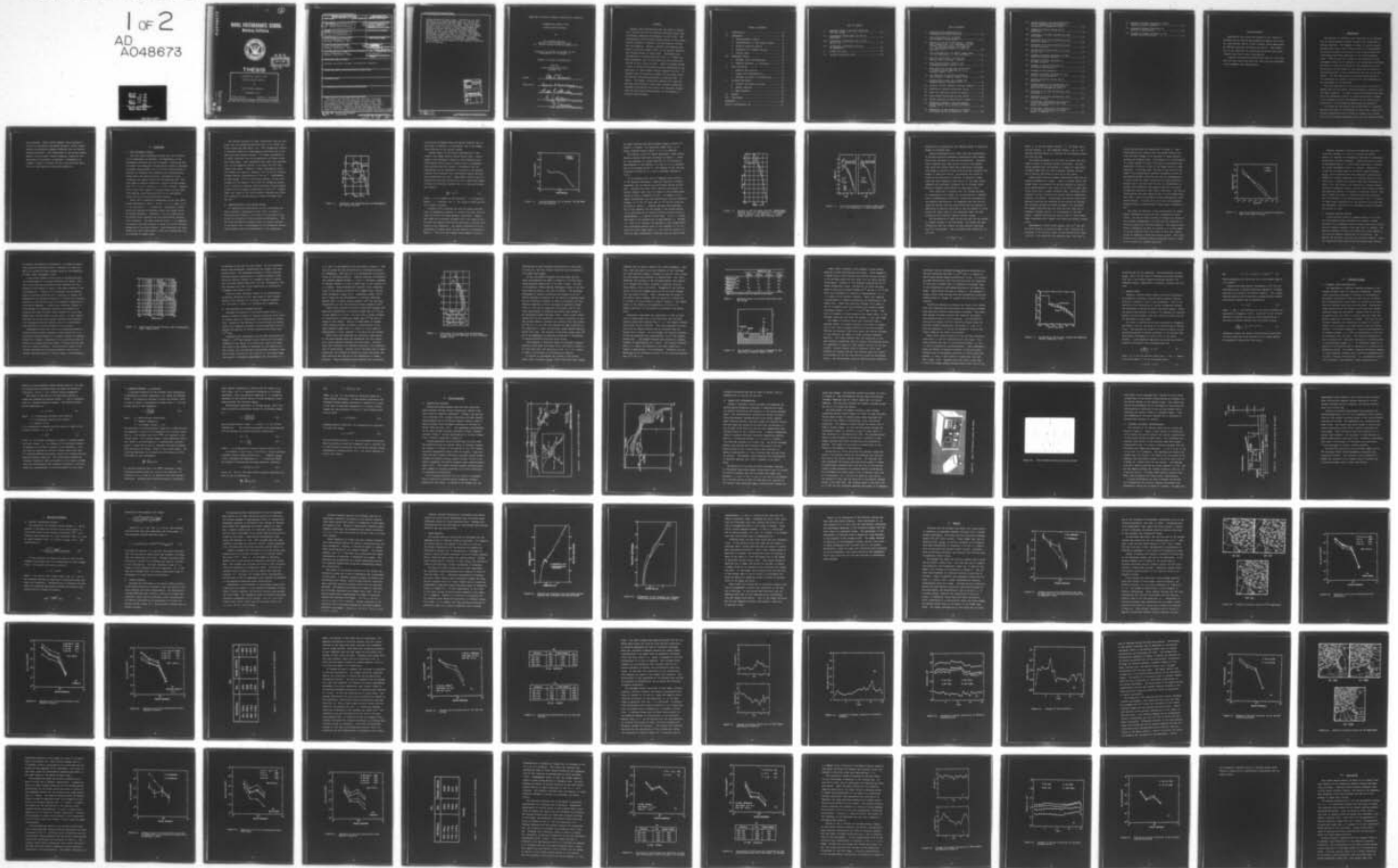
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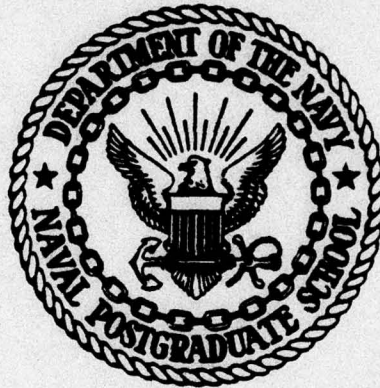


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Monterey, California



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## THESIS

A COMPARATIVE STUDY OF THE  
COASTAL MARINE AEROSOL

by

Alan Anthony Simoncic

December 1977

Thesis Advisor:

Kenneth L. Davidson

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A Comparative Study of the  
Coastal Marine Aerosol

by

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Captain, United States Air Force  
B.S., United States Air Force Academy, 1970

Submitted in partial fulfillment of the  
requirements for the degree of

MASTER OF SCIENCE IN METEOROLOGY

from the  
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## ABSTRACT

Aerosol size distributions near the coast of Panama City, Florida and off the Southern California coast near the Channel Islands are investigated in this study. The relationships of the coastal marine aerosol to wind speed, relative humidity, stability, and sub-synoptic circulation are examined. Relative humidity and stability are shown to have the largest effect on the aerosol distribution during periods of light winds. Coalescence and sedimentation of droplets greater than  $1.5 \mu$  radius are most pronounced when the wind speed and sea surface production of salt nuclei are weak. When wind speeds exceed  $7 \text{ m/sec}$ , a state of equilibrium between sedimentation and production of these larger droplets appears to exist. An apparent zone of transition between the two bubble bursting sea-salt producing mechanisms is observed near  $.5 \mu$  radius. The highest correlation between wind speed and particle concentration occurs under unstable conditions. Secondary circulations are shown to be important determinants of the coastal marine aerosol in the absence of synoptic scale forcing.

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## I. INTRODUCTION

The military is currently very interested in the performance of electro-optical weapons systems in an atmosphere of varying turbidity. For example, a number of electro-optical systems which utilize the visible as well as IR wavelengths are being developed by the Navy for use in surveillance and intelligence gathering operations in the marine boundary layer. These systems are limited by the extinction of the propagated energy due to absorption and scattering by aerosols. The effect of absorption depends on the composition of the particulates and wavelength of the energy and the effect of scattering depends on the concentration and size of the scatterers. For most applications the scattering processes in the atmosphere are caused by particles of size comparable to the wavelength of the radiation.

The size distribution of the marine aerosol is known to depend upon the wind speed, relative humidity, stability, and air mass trajectory. In order to evaluate accurately and predict the atmospheric effects on these electro-optic systems, it is necessary to know the dependence of the aerosol size distribution on the foregoing meteorological parameters.

The nature of the aerosol size distribution in a coastal marine environment is investigated in this study. Data from aerosol observations off the coast of Panama City, Florida and off the Southern California coast near the Channel Islands

were analyzed. These coastal regimes, which represent a mixture of continental and marine aerosols, should contain aerosol distributions somewhat different from the typical marine environment. The relationship of the coastal marine aerosol to wind speed, relative humidity, stability, and sub-synoptic circulation is examined. Furthermore, an attempt is made to evaluate the use of the friction velocity as a valid aerosol distribution predictor.

## II. BACKGROUND

### A. THE ATMOSPHERIC AEROSOL

With the recent increasing concern over the pollution of our atmospheric environment, the examination of the tropospheric aerosols has also increased. Particulate matter enters the atmosphere through either natural or man-made processes; approximately 10% of the total concentration is believed to originate from combustion and industrial processes while the natural sources, including soil dust, volcanoes, and oceans account for the remaining 90%. The size range of aerosols observed by current methods extends from  $10^{-3}$   $\mu$  to  $10^3$   $\mu$  radius ( $1 \mu = 10^{-6} \text{m} = \text{micron}$ ). Depending on their size, amount of soluble matter, and the relative humidity, these particles may act as condensation nuclei and aid in the precipitation process.

Mason (1975) classified condensation nuclei into three groups according to radius: Aitken ( $< 0.1 \mu$ ), Large ( $0.1 \mu - 1 \mu$ ), and Giant ( $> 1 \mu$ ) particles. Essentially, Aitken nuclei are produced by man-made sources and larger nuclei by natural processes. Therefore, it is not surprising to see Aitken nuclei dominate the size distribution spectrum over continents. The marine aerosol above  $.1 \mu$  is composed of sea-salt particles produced by spray and bubble bursting mechanisms on the water surface. These mechanisms are quite complex and their contribution to the size distribution will be discussed in detail later.

The vertical profiles of trace constituents that are produced over the continents have been shown to be rather uniform in space and time above 5 km. This "background" aerosol is affected slightly by anthropogenic activities and is far away from local natural sources. Its number concentration is almost identical with the concentration of Aitken nuclei over ocean areas. Past experiments have shown the concentration of the background aerosol to be about  $300 \text{ cm}^{-3}$  over remote ocean areas. However, a recent experimental cruise (R/V Meteor) by Junge and Jaenicke, 1971 in the mid Atlantic yielded observed concentrations of  $600 \text{ cm}^{-3}$ . Measurements by Hidy, et al. (1973) on San Nicolas Island, 130 km, west-southwest of Los Angeles, have shown the background aerosol to be a mixture of material from both marine and continental sources with an average concentration of  $2400 \text{ cm}^{-3}$ . Samples taken over oceans of the South Atlantic (Meszaros and Vissy, 1974) resulted in Aitken particle counts of between  $300\text{-}450 \text{ cm}^{-3}$ .

#### B. CHARACTERISTICS OF THE MARINE AEROSOL

An idealized size distribution of the continental and marine aerosol is supplied by Junge (1972) in Figure 1. The significant feature is the shift of the maximum of particles as a function of total particle concentration. Over the ocean the sea-salt aerosol, which is usually confined to the lower 2 km, is superimposed on the background aerosol. Junge reasoned that the concentration of the background



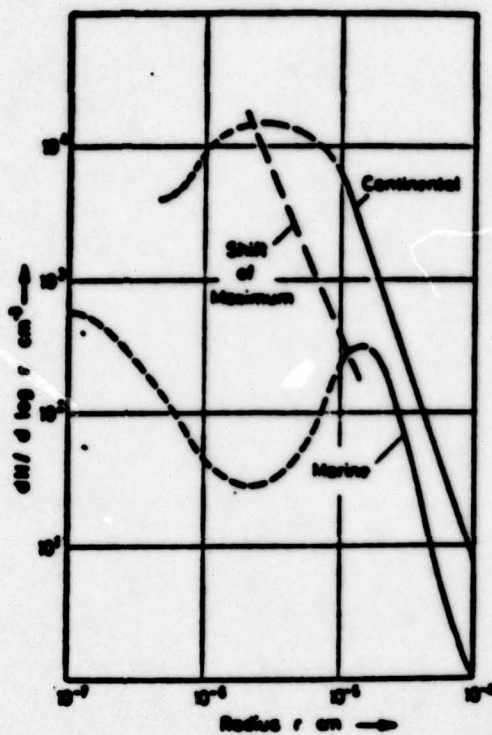


Figure 1. Idealized Size Distributions of Continental and Marine Aerosols

cloud nuclei decreases below the marine inversion due to the effect of washout, or coalescence, due to the larger water droplets in the cumulus clouds.

Another aspect of the aerosol distribution is the slope of the number density versus radius curve. Friedlander (1961) proposed a theory of self-preserving size distributions which helps to explain why all atmospheric size distributions are similar. He proposed that the similarities can be explained by solutions to the kinetic equation which describe the relationship between particle size distribution and time. Experimental results have indicated that the size distribution over a particular range of sizes of continental aerosol has a -4 slope and follows the relation

$$\frac{dN}{dr} = C\phi r^{-4} \quad (1)$$

where  $N$  is the number of particles/cm<sup>3</sup>,  $r$  the particle radius,  $C$  a constant, and  $\phi$  the volume of particles per unit volume of aerosol.

Blifford (1970) measured the size and number distributions of atmospheric aerosols at various altitudes over the ocean 250 km west of Santa Barbara, California. Samples were taken by an aircraft equipped with a jet impactor and the data was obtained from direct microscopic counting techniques in the laboratory. The aerosol distribution at approximately 15 meters above the sea surface is presented in Figure 2. The curve has a rather steep negative slope at

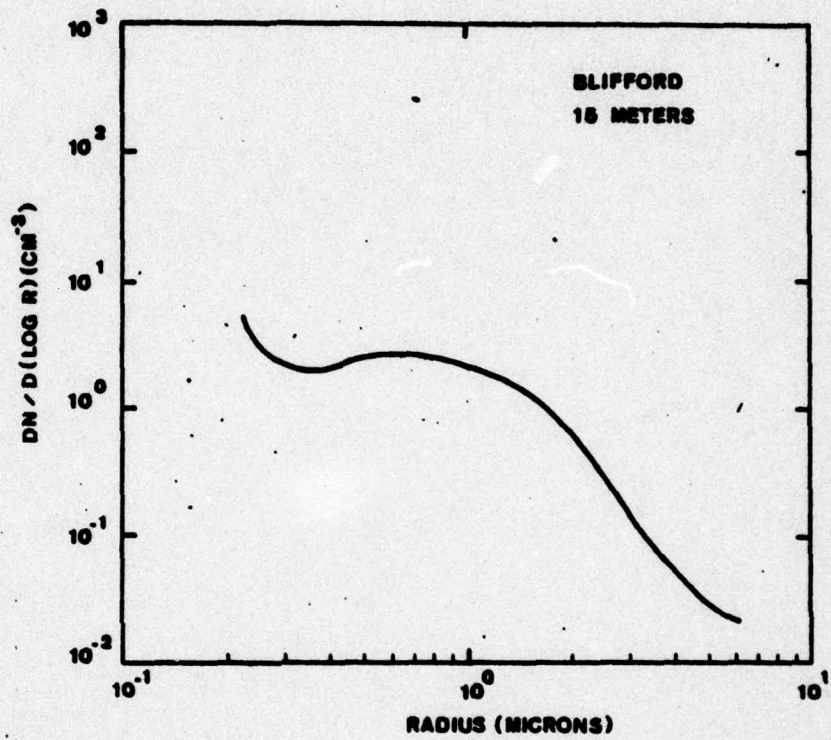


Figure 2. Size Distribution at 15 Meters, 250 km West of Santa Barbara

the small particle end which becomes slightly positive at around  $.4 \mu$  radius. For particles larger than  $.8 \mu$ , a fairly constant slope of about -2 to -3 is observed.

The results of the R/V Meteor experiment, where several aerosol counters were used, are shown in Figure 3. Above  $10 \mu$  the exponent of a power function fit to the data is approximately -6 and between  $0.3 \mu$  and  $10 \mu$  it is variable but on the average around -3. The maximum of the size distribution occurred at  $0.3 \mu$  with a secondary maximum at  $0.03 \mu$ .

It is possible that, due to increased human activity in the Northern Hemisphere, Junge and Jaenicke's Atlantic experiment did not explore the undisturbed marine environment. Meszaros and Vissy (1974) describe the results of aerosol samples taken over the oceans of the Southern Hemisphere by means of membrane filters. An example of the number concentration and size distribution over the Atlantic between (a)  $0^\circ$  and  $20^\circ$  South and (b)  $40^\circ$  and  $60^\circ$  South can be found in Figure 4. Chemical analyses were performed and it was observed that the maxima in the concentrations of all particles and of sodium chloride particles occur at approximately  $.1 \mu$  radius in both cases. Up to  $.5 \mu$  radius the slope of the distribution is approximately -5. Between  $0.5 \mu$  and  $1.5 \mu$ , however, the decrease of the concentration with increasing particle size is very moderate (-1 to -2), while for radii larger than  $1.5 \mu$  the slope is close to -3. This has been interpreted to indicate that the form of the

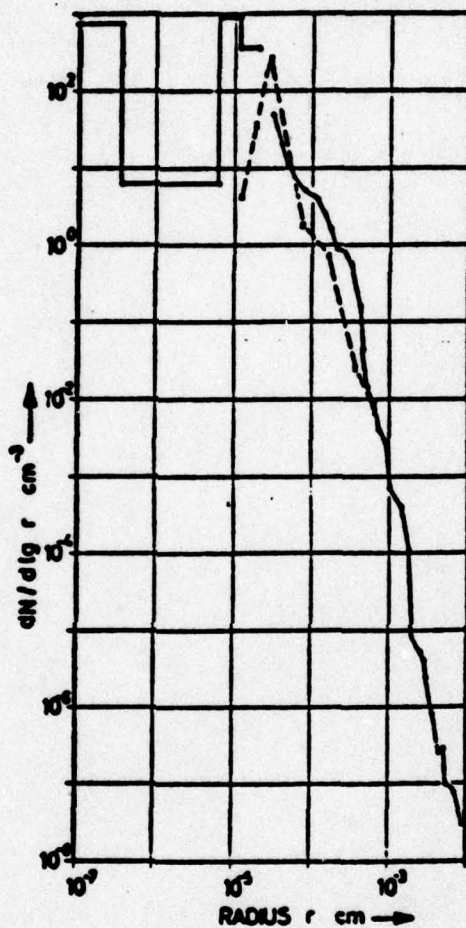


Figure 3. Results of the R/V Meter Cruise. Measurement Systems Used:  $\circ$ — $\circ$  Combination of CCN Counter, Optical Counter, and Impactors: --- Double Stage Impactor (Junge and Jaenicke, 1971)

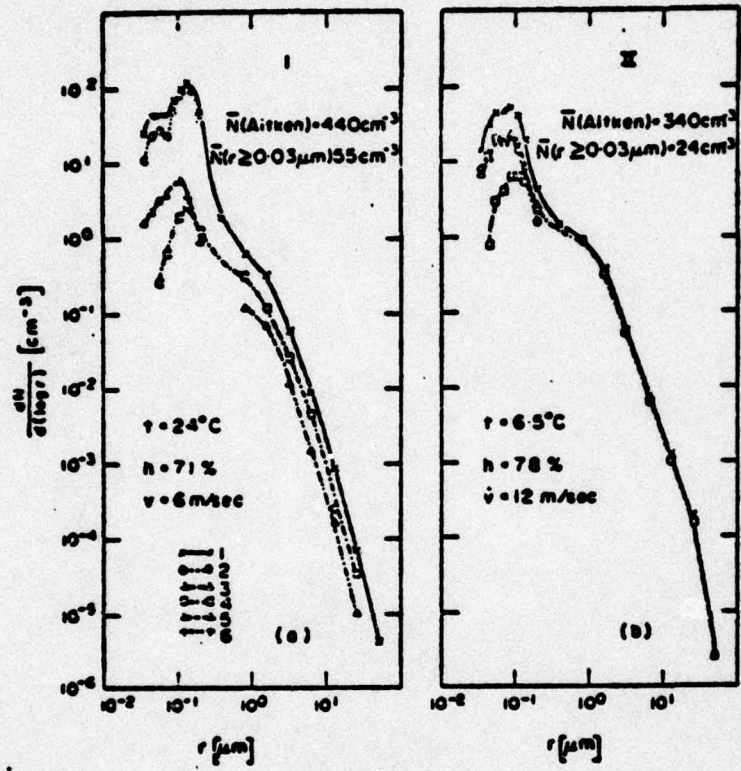


Figure 4. Size Distributions over Remote Ocean Areas  
 (— All Particles; - - NaCl Particles)

distribution is produced by the combined effect of particles formed in different ways.

Oceanic measurements have shown that the concentration of sea-salt particles decreases exponentially with height, with little variation of the size distribution. Ericksson (1959) reported that there exists a level a few hundred feet above the surface where the concentration decreases with height in periods of high wind force and increases with height in lower wind forces. He reasoned that there is little or no production of sea-salt in regions of light winds and that coagulation and fallout in the lower levels combined with horizontal transport due to vertical shear produce a maximum concentration at some upper level.

Toba (1965a and b) proposed that the average decrease in concentration with height can be explained by a combination of sedimentation, diffusion, convective processes and the humidity distribution. He suggested that the line between the aerosol vertical distribution and the process of production of sea-salt particles at the sea surface is found within the lowest layer of the atmosphere where the eddy diffusivity and relative humidity sharply change.

The distribution of eddy diffusivity near the sea surface is closely related to the wind speed. The larger the eddy diffusivity near the surface the more sea-salt particles that will be supplied. Toba considered eddy diffusivity in the form

$$D = kU_* (Z + Z_0) \quad (2)$$

where  $k$  is the von Karman constant,  $Z$  the height above the sea surface,  $Z_0$  the roughness length, and  $U_*$  the friction velocity which is a function of the momentum transfer over the sea.

The relative humidity in the first few meters over the ocean is known to decrease rapidly with height. The particles produced near the surface in a region of high humidity grow larger and thus have a greater terminal velocity due to gravity than those at the top of this layer.

During light winds the number concentration near the sea surface increases with height. Since it results from a non-steady state, an inversion of vertical gradient of the particle concentration is most likely to be found in small particles which have a longer residence time. Ericksson (1959) computed the fall velocities for given relative humidities, salinity, and radius. During high wind periods, giant size sea-salt particles are produced at the surface and through the diffusion process are mixed throughout the atmosphere. The largest particles may fall back into the ocean due to excessive terminal velocity or be entrained in the wave crests. Smaller particles are free to rise to cloud height where coalescence with larger cloud drops and washout usually occur.

Measurements of salt nuclei greater than  $10^{-14}$  gm over the North Atlantic by Moore and Mason (1954) revealed the existence of two distinct types of size distributions (Type I and II). The curves for the observed Type I and Type II



nuclei distributions are reproduced in Figure 5. Type I distributions were observed for wind speeds between 6-15 m/sec and were thought to be residuals of spray droplets produced by breaking waves. The presence of a discontinuity or a sharp change of slope in the Type I distribution was explained in terms of a loss of the larger nuclei by sedimentation. In strong winds, the part of the curve to the right of the discontinuity probably represents a state of equilibrium between production and loss by sedimentation. In light winds and stable conditions the slope should be steeper due to the fact that the loss by sedimentation is greater than production and larger nuclei are not easily transported vertically under stable conditions. The Type II distributions were only observed when the wind speeds were less than 7 m/sec and resembled a high concentration continental aerosol. In winds of up to 15 m/sec the measured concentrations of large sea-salt nuclei rarely exceeded  $10 \text{ cm}^{-3}$ .

The effect of stability on the concentration of atmospheric condensation nuclei was well documented by Moore (1952). He used an Aitken counter to measure the relationship between concentration of nuclei and the intensity of vertical mixing over the North Atlantic. The results indicated a decrease by as much as a factor of 4 in the number of Aitken particles near the surface on days with cumulus clouds as compared to days with stratus clouds. This would indicate that convection plays an important role, at least in the transport of smaller particles.

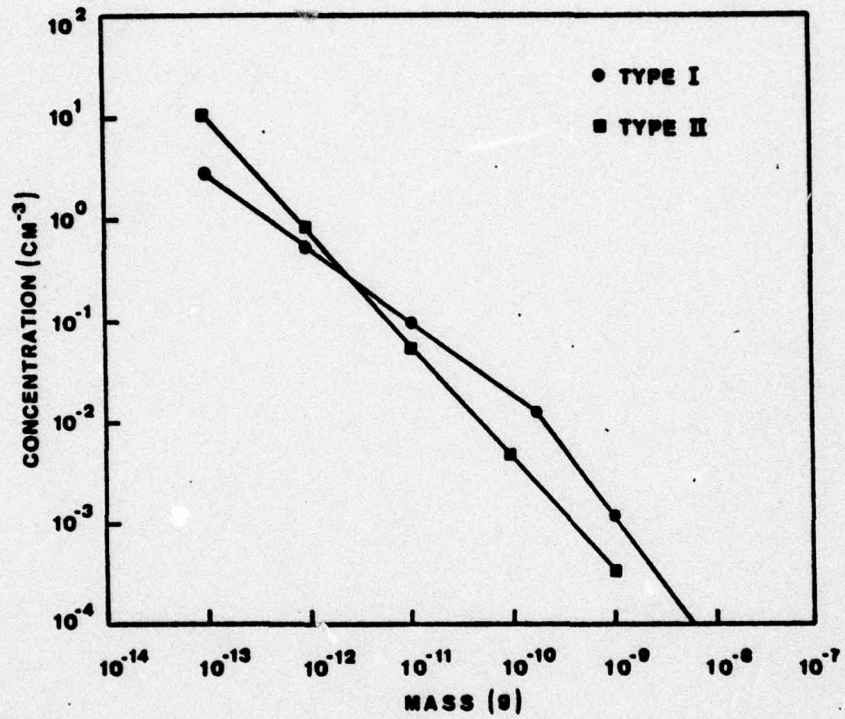


Figure 5. The Size Distribution of Moore and Mason's Type I and Type II Nuclei

Chemical analyses by various investigators have indicated that between 0.1  $\mu$  and 1.0  $\mu$  radius the marine aerosol is composed of a background component of continental origin and a sea-salt component. Sodium chloride was found to predominate above 1  $\mu$  radius while particles of continental origin predominate below 0.1  $\mu$  radius. Results of a cruise off the Grand Banks in the North Atlantic (Ruskin, et al., 1976) indicated that the continental particles are composed of sulfate compounds and a smaller amount of sulfuric acid. Aerosols over remote ocean areas (Meszaros and Vissy, 1974) were shown to be comprised of variable concentrations of ammonium sulfate, sulfuric acid, sodium chloride, and particles similar in structure to ammonium sulfate. The sum of these four types of identified particles accounted for 75-95 percent of all particles greater than .3  $\mu$  radius. In other words, practically all the particles in a pure marine atmosphere, undisturbed by continental particle sources, are soluble in water.

#### C. RELATIVE HUMIDITY EFFECTS

A solid particle which is composed wholly, or in part, of a pure water-soluble substance will undergo a sudden transition to a saturated solution droplet when some critical value of relative humidity, less than 100%, is reached. The relative humidity at which this transition occurs depends on the size and chemical composition of the particle. The smaller the particle, the lower the critical humidity. Below the transition point, solid particles acquire small amounts

of water by the process of adsorption. At relative humidities above the transition point, a particle (or, more properly, an aqueous solution droplet) grows by the absorption of water vapor (Fitzgerald, 1975).

A pure water droplet is said to be in equilibrium with its surroundings if it neither evaporates nor grows. This only occurs when the equilibrium vapor pressure over the surface of the droplet is equal to the vapor pressure of the surrounding air. Winkler (1973) describes the equilibrium growth of aerosol particles due to humidity as complex and depending on the relative proportion of soluble and insoluble material in the particles and on the chemical composition of the soluble component. Complex ionic mixtures, similar to those present in atmospheric aerosols, show material influences and lower the water vapor pressure to a much less degree than the same amount of pure salts. In such complex mixtures the various salts become dissolved only gradually with increasing relative humidity until at a sufficiently high humidity all soluble material is in solution.

Measurements have shown that with increasing humidity a sodium chloride crystal undergoes a phase transition to a saturated solution droplet at a relative humidity of approximately 78%. Figure 6 describes how the equilibrium radii of droplets containing specified masses of sodium chloride vary with the relative humidity. The equilibrium radius of the droplet increases with increasing humidity until the air becomes supersaturated by a critical amount, corresponding to

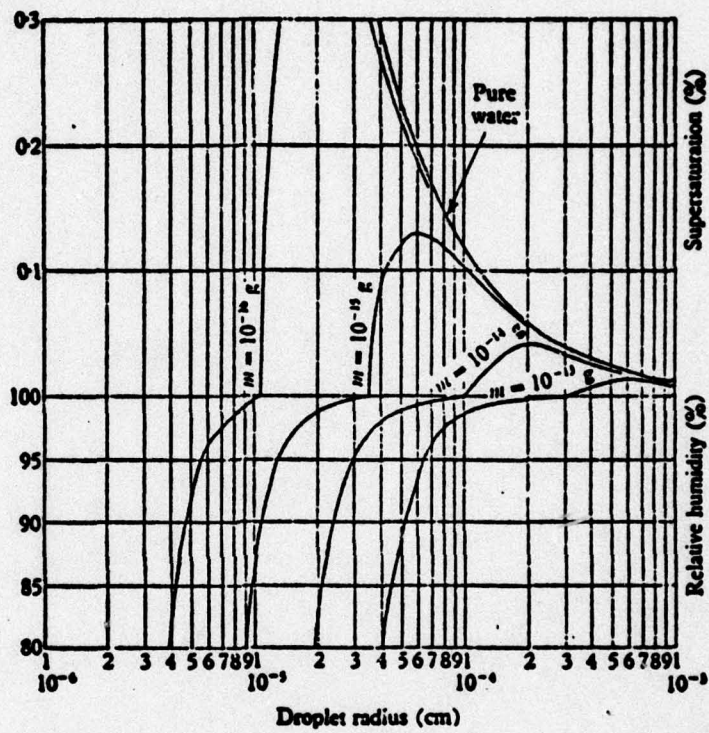


Figure 6. Equilibrium Relative Humidity and Corresponding Radii (Mason, 1975)

the maximum of the curve in this figure. If this supersaturation were maintained, theoretically the droplet will grow without bound. With decreasing humidity a sodium chloride solution droplet crystallizes at a humidity between 35-45%. Since the relative humidity at a height of about 15 meters over the ocean surface goes below 40% very infrequently, sea-salt droplets will have little opportunity to crystallize (Fitzgerald and Ruskin, 1977).

Since the later discussion refers to the distribution of sea-salt particles by bolt mass weight of salt (grams), radius of dry crystals ( $\mu$ ) or radius at ambient humidity ( $\mu$ ), the scale in figure 7 is furnished as a reference.

#### D. THE PRODUCTION OF AIRBORNE SEA-SALT

Although the spectrum of the marine aerosol above .1  $\mu$  radius is known to consist of sea-salt particles, very little is certain about the concentration and mechanisms of production. Because of the smallness of the particles and limitations of the sampling equipment, earlier experiments did not measure the quantity of sea-salt particles much less than  $10^{-12}$  gm.

Woodcock (1953) determined that the mass distribution of "giant" ( $> 10^{-12}$  gm) sea-salt nuclei varies with wind speed. Increases in the amount of air-borne salt near cloud bases were shown to be related to increases in wind speed at the sea surface, with the greatest proportionate increase in particle number occurring at the large end of the weight range. The results of Woodcock's measurements for wind forces of 1,

3, 5, and 7 on the Beaufort scale are shown in Figure 7. The line (a) gives the size distribution of continental aerosol for comparison. The line (b) is an extrapolated size distribution of the marine aerosol. Chemical analysis of Woodcock's bulk aerosol samples between  $.1 \mu$  and  $1 \mu$  indicated a maximum of sea-salt around  $0.3 \mu$  and a lower limit in the vicinity of  $.1 \mu$  radius. These distributions indicate total concentrations of all sea-salt particles of no higher than a few per cubic centimeter (Junge, 1972). According to Mason (1975), over a rough sea the concentration of sea-salt particles greater than  $2 \mu$  radius rarely exceeds  $1 \text{ cm}^{-3}$  and the total concentration of all salt particles rarely exceeds  $10 \text{ cm}^{-3}$ .

Moore (1952) observed a distinct correlation between wind speed and concentration of sea-salt larger than  $10^{-11}$  gm up to wind speeds of 15 m/sec. He also found a linear increase in concentration of particles larger than  $10^{-9}$  gm with increase in wave height. Results of experiments by Monahan (1968) reveal an abrupt increase in concentration of sea water droplets larger than  $45 \mu$  radius at a wind speed of approximately 9 m/sec, measured 47 cm above the sea surface.

Moore (1952) also analyzed the visibility observations at two ocean weather ships and determined that the opacity for a given humidity increases with wind speed. He attributed this increase to an observed increase in the concentration of large nuclei. Another result indicated that at lower humidities, the increase in opacity was more pronounced, and Moore believed this was due to the dehydration of larger droplets. These conclusions would indicate that the aerosol

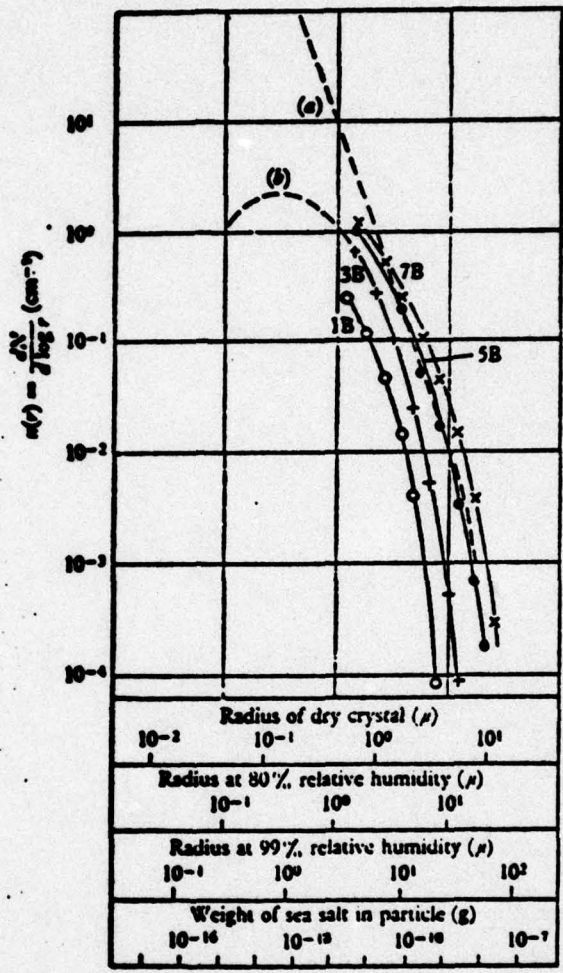


Figure 7. Size Scale and Average Size Distribution of Sea-Salt Nuclei Measured by A.H. Woodcock (Mason, 1975)



distribution is more variable and sensitive to wind speed in drier air, and this feature should be most noticeable in the larger size ranges.

As the wind speed increases over the ocean, gravity waves are generated and begin to break at a critical wind speed generally agreed upon to be near 7 m/sec. Air that is entrained by these breaking "whitecaps" rises to the surface sometime later in the form of bubbles. The principal mechanisms of sea-salt production are thought to be the direct spraying of droplets off the crests of breaking waves and the bursting of bubbles in areas of whitecaps and foam. Droplets produced by direct spraying are generally larger than  $45 \mu$  and, due to large fall velocities, are not airborne long enough to evaporate and become light enough to be transported upward (Monahan, 1968). Toba's model (1965b) showed that the net production of sea-salt particles at the sea surface seems to increase with particle mass even beyond  $10^{-8}$  gm ( $20 \mu$ ), but that the transport by eddy diffusion is not sufficient to carry the particles upward against gravity beyond this size. The presence of particles larger than  $10^{-8}$  gm in the atmosphere is generally attributed to coalescence of sea-salt droplets within and below clouds.

Some examples of residence times for different sea-salt particle sizes taken from Junge (1972) are found in Table I. It would seem then that particles in the  $.1 \mu - 20 \mu$  range, at least, are produced by the bursting of bubbles.

In efforts to photograph the rupture of the surface bubble film, Kientzler, et al. (1954) found that their camera

exposure was too long to capture this rapid phenomenon. However, they were able to see the formation of the "Rayleigh" jet which projects upward, continues to rise as a thin column, and then breaks into droplets of varying sizes. Day (1964) describes this process in the following manner. Each bubble, as it reaches the surface, develops a spherical film-cap which drains, thins, and bursts. Fragments of the film are thrown out and are dragged upward by the air which escapes from the bubble orifice. Water, rushing down the sides of the bubble cavity, emerges from the center as a narrow jet. A schematic of this process is shown in Figure 8. The larger drops (L) are formed by disintegration of the jet (J). Smaller particles (S) are formed by bursting of the bubble film.

Kientzler's experiment was significant in that no droplets of large enough size to be resolved by the film and optical system were observed from .2 - 1.8 mm diameter bubbles until after the jet formation. This was interpreted to indicate that the larger droplets are not produced when the bubble film is broken. On the average, the droplets produced by the jet mechanism were approximately 1/10 of the original bubble size. 1 mm diameter bubbles were observed to produce droplets of approximately 50  $\mu$  radius. The smallest observed were of 2  $\mu$  radius and deduced to have been formed by a bubble of approximately .04 mm diameter. Therefore, the jet mechanism can be considered a source of salt particles greater than  $10^{-12}$  gm (1  $\mu$ ).

	Residence Time $\tau$ , days			
	$M = 10^{20}$ grams	$M = 10^{11}$ grams	$M = 10^{10}$ grams	$M = 10^{-9}$ grams
Toba's value $\tau_1$ for 80% relative humidity	86	17	2.9	0.32
Toba's value $\tau_1$ for 91.4% relative humidity	39	11	3.1	0.23
Eriksson's estimate from sedimentation	82	16	2.6	0.4
Eriksson's estimate from production	3.5	1.0	0.6	0.5
Our estimate*	1.9	1.6	1.0	0.26

Table is taken from Toba (1963a, Table 2). The variable  $M$  is the mass of particles.

Table I. Residence Times of Sea-Salt Particles over the Oceans

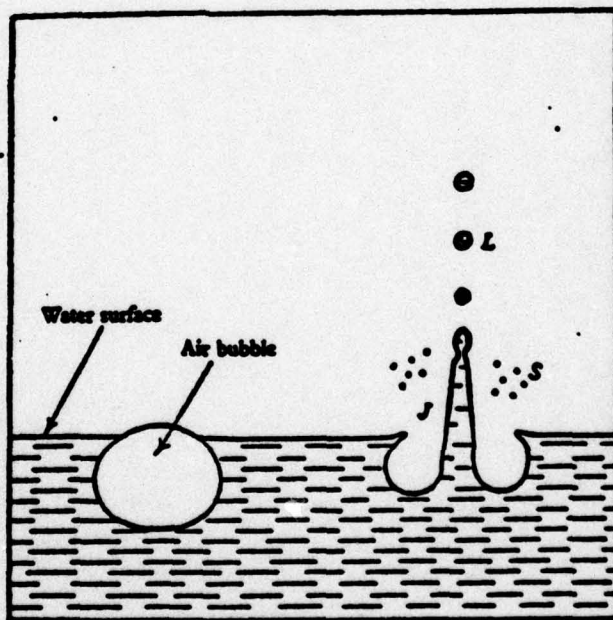


Figure 8. The Formation of Sea-Salt Droplets by the Bursting of Bubbles (Mason, 1975)

Mason (1954) utilized a cloud chamber to study bubble behavior in both distilled and salt water. After expansion, a dense cloud of tiny droplets was observed rising vertically in the space above the salt water, but not above the distilled water. Bubbles of 3 mm diameter produced 100-200 of these condensation nuclei, the majority of which are estimated to have salt contents between  $10^{-15}$  gm and  $2 \times 10^{-14}$  gm. This would correspond to droplets of approximately .1  $\mu$  to .3  $\mu$  radius at 80% relative humidity. Mason also observed a second group of droplets produced by the shattering of the bubble film. These were projected sideways at an angle of ten to 15 degrees above the horizontal and slightly larger, containing between  $2 \times 10^{-12}$  -  $5 \times 10^{-10}$  gms of salt. However, the numbers of these droplets were always small, on the average, there was only about one droplet in this size range.

The number of droplets which rise vertically from a bursting bubble is strongly related to the state of compression of the film of organic material on the water surface. Paterson and Spillane (1969) have shown that with an increase of film pressure the number of nuclei produced decreases markedly. This would indicate that the production of sea-salt droplets originating from the bubble film mechanism would be suppressed in regions of high organic activity on the sea surface. Aerosol samples taken by Woodcock (1972) over Hawaiian and Alaskan seas may help explain where the transition between the jet and film sea-salt production mechanisms occurs. His observations, using an improved slide collection

technique, show an increased average particle production for sea salt particles less than  $2 \times 10^{-14}$  gm ( $.3 \mu$  radius) in Hawaii where marine organic productivity is low. In contrast, the mean distribution curve for particles over the organically rich Gulf of Alaska fails to indicate an increased slope of the concentration curve among particles of the same size range. These curves are shown in Figure 9. The presence of surface active films arising from the biologically productive Alaskan waters is thought to suppress the production of film droplets.

Statistical analysis by Meszaros and Vissy (1974) showed that with increasing particle radius the correlation between wind speed and chloride concentration increased. This meant that smaller chloride particles are formed by the bubble film mechanism than by direct spraying. The distribution curve gives evidence that the transition between these two chloride formation mechanisms lies between  $.2 \mu$  and  $.4 \mu$ . Thus the maximum at  $.1 \mu$  gives the maximum of chloride particles formed by the bubble film process.

Moore (1952) found evidence that the particle concentrations below  $1 \mu$  are not correlated with wind speed. This would indicate that most of the particles between  $.1 \mu$  and  $1 \mu$  are not produced by bursting bubbles. Other experiments using the effects of relative humidity on particle growth indicate that a considerable proportion of marine particles between  $.1 \mu$  and  $1 \mu$  must differ in composition from sea-salt (Junge, 1972). Meszaros and Vissy (1974) found that, in this size range, sodium chloride varied from 4-50% of the

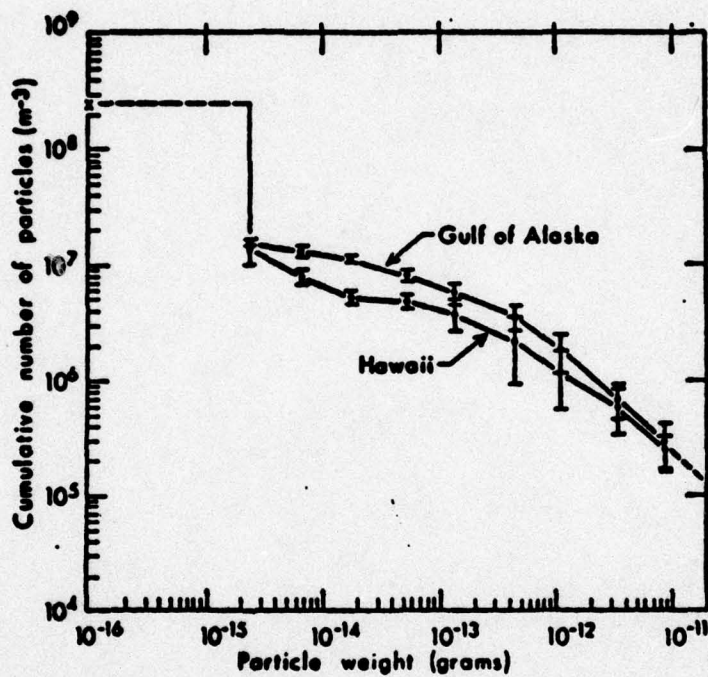


Figure 9. Distribution Curves over Alaskan and Hawaiian Waters (Woodcock, 1972)

concentration for all particles. The observations by Hidy et al. (1974) off the coast of Southern California revealed that 11% of the aerosol sampled contained sea-salt, the remainder being a combination of sulfates, nitrates and soil dust.

#### E. AEROSOL MODEL

Recently, various aerosol models have been developed in an attempt to accurately describe marine aerosol distributions as a function of one or more parameters. This is essential for the calculation of optical propagation through the atmosphere as aerosols scatter and absorb energy. Since the aerosol distribution is known to be dependent on relative humidity and wind speed, these two variables usually are the key parameters of each model.

One model in particular has been developed by Fitzgerald and Ruskin (1977) on the basis of the North Atlantic observations. They applied the effects of relative humidity on the equilibrium growth of aerosol particles to the sea-salt mass distribution determined by Lovett (1975) in the North Atlantic. Lovett presents empirical log radius mass distributions in the form of the following power law:

$$\frac{dN}{d \log r_d} = C r_d^{-\nu} \quad (3)$$

where  $r_d$  is the dry particle radius and  $C$  and  $\nu$  depend on the wind speed  $V$  in the following manner:

$$\nu = 3.317 - .03 V \quad (4)$$

and 
$$C = 0.2 - 0.0196 V + 0.0121 V^2 \quad (5)$$

These expressions are valid only over a wind speed range of 3-17 m/sec<sup>-1</sup>.

Formulae have been derived (Fitzgerald, 1975) for the equilibrium size of aerosol particles composed of a single pure salt as a function of relative humidity. For a sodium chloride particle the relationship between particle radius and relative humidity may be expressed as

$$r = \alpha r_d^\beta \quad (6)$$

where  $\alpha$  and  $\beta$  are functions of the relative humidity as described by Fitzgerald (1975). Equations (3) and (6) are combined to describe the aerosol size distribution as a function of relative humidity and wind speed, giving

$$\frac{dN}{d \log r} = \frac{C}{\beta} (\alpha^{v/\beta}) (r^{-v/\beta}) \quad (7)$$

Comparison between the aerosol distributions derived from the above model and those observed in two coastal marine environments is made within this study.



### III. TURBULENCE THEORY

#### A. BOUNDARY LAYER CONSIDERATIONS

The importance of turbulent exchange processes in the surface boundary layer has long been recognized. Panofsky (1969) describes atmospheric turbulence as consisting of horizontal and vertical eddies by which the air is mixed. The two mechanisms by which eddies are formed in the atmosphere are heating from below and wind shear. Heating produces convection and the change in wind speed with height produces mechanical turbulence. Because there is no wind at ground level, and there is usually some wind above the ground, mechanical turbulence is common. This type of turbulence increases with increasing wind speed (at a given height) and is greater over rough terrain than over smooth terrain. The terrain roughness is usually characterized by a roughness length,  $Z_0$ , which is proportional to the size of the eddies that can exist. The relative importance of heat convection and mechanical turbulence is characterized by the Richardson number,  $R_i$ . The Richardson number is a measure of the relative rate of conversion of convective to mechanical energy. For example, negative Richardson numbers of large magnitude indicate that convection predominates resulting in strong vertical motion. As the mechanical turbulence increases, the Richardson number approaches zero.

Finally, as the Richardson number becomes positive, the thermal stratification becomes stable and damps the mechanical turbulence. For  $R_i > 0.25$ , vertical mixing disappears.

The effect of the wind on the underlying surface is termed the shearing or Reynolds stress,  $\tau$ , and is characterized by a downward momentum transfer. The Reynolds stress may be represented by

$$\tau = -\rho \langle u'w' \rangle \quad (8)$$

where  $u'$  = fluctuating horizontal wind velocity

$w'$  = fluctuating vertical wind velocity

$\rho$  = density of air

It is convenient to express Reynolds stress in terms of the friction velocity  $U_*$  so that

$$\tau = \rho U_*^2 \quad (9)$$

where  $U_*$  is constant throughout a region of constant momentum flux. Hence,  $U_*$  is a measure of the downward transfer of momentum in the lower 50 meters of the atmosphere. Over the ocean an increase in the near surface winds would lead to a greater momentum and energy transfer for surface wave and sea-salt aerosol production. The relationship between the turbulent transfer of heat and moisture in the marine boundary layer and the generation and transfer of aerosols is not well known and, unfortunately, is not investigated in this study.

## B. MOMENTUM TRANSFER, $U_*$ , RELATIONS

A thorough discussion of the boundary layer expressions is presented in several references, e.g. Lumley and Panofsky (1964). The similarity approach of Monin and Obukhov (1954) is used to define a representative length scale,  $L$ , for the surface layer of the atmosphere,

$$L = \frac{-U_*^3 T_o}{kg \overline{w'T'}} \quad (10)$$

where  $g$  = gravitational acceleration

$T$  = ambient temperature

$k$  = von Karman constant = 0.35

The selection of the Monin-Obukhov length as a stability scaling parameter is based on the assumption that friction velocity,  $U_*$ , and vertical heat flux ( $\overline{w'T'}$ ) are constant in the surface layer. This scaling length, using dimensional analysis, leads to the development of a dimensionless function,  $\phi_m(Z/L)$ , which can be used to represent the mean horizontal wind variation with height,  $d\bar{u}/dZ$ , in the surface layer. The following expression is the empirical relationship for the wind shear in this development,

$$\frac{d\bar{u}}{dZ} = \frac{U_*}{kZ} \phi_m(Z/L) \quad (11)$$

As vertical turbulent heat flux ( $\overline{w'T'}$ ) decreases to zero, indicating neutral stability,  $\phi_m(Z/L)$  must approach 1 if Equation (11) is to take on its expected form under neutral conditions. Assuming that convective mixing is negligible

under neutral conditions it follows that for values of  $\phi_m$  ( $Z/L$ ) near 1 or  $Z \ll 1$  mechanical turbulence is of primary importance. Thus, the absolute magnitude of  $L$  becomes an indicator of the vertical extent to which mechanical turbulence controls the turbulent regime.

Observational experiments by Businger et al. (1971) produced a definite relationship between the Richardson number,  $R_i$ ,

$$R_i = \frac{g(\partial\theta_v/\partial Z)}{\bar{\theta}(\partial u/\partial Z)^2} \quad (12)$$

and the Monin-Obukhov length,  $L$ , where  $\theta_v$  is the virtual temperature. The following expressions are approximations for the unstable and stable conditions respectively,

$$Z/L = R_i \quad (13)$$

$$Z/L = \frac{R_i}{1-\alpha R_i} \quad (14)$$

where  $\alpha$  is an empirically derived constant equal to 0.5.

Of interest in this study is the rate of viscous molecular turbulent kinetic energy dissipation,  $\epsilon$ . Wyngaard, et al. (1971) considered the dependence of  $\epsilon$  on momentum fluxes and height in deriving the following empirical expression

$$\epsilon = U_*^3/kZ \phi_\epsilon (Z/L) \quad (15)$$

Since  $Z/L$  and  $R_i$  are functionally related, equations (11) and (15) can be rewritten as

$$\frac{d\bar{u}}{dZ} = \frac{U_*}{kZ} f_m (R_i) \quad (16)$$

and 
$$\epsilon = U_*^3 / kZ f_\epsilon (R_i) \quad (17)$$

where  $f_m$  and  $f_\epsilon$  are stability corrections equal to 1 under neutral conditions. In near neutral conditions, the turbulent kinetic energy production is assumed to be equal to the rate of molecular dissipation of turbulent kinetic energy and from equations (16) and (17) the following relation is valid

$$\epsilon = U_*^2 (\partial \bar{u} / \partial Z) \quad (18)$$

Assuming neutral conditions, the combinations of equations (16) and (18) yields

$$U_* = (\epsilon k Z)^{1/3} \quad (19)$$

Now the friction velocity  $U_*$  can be estimated from either mean wind profiles using the integrated form of equation (16) or from velocity fluctuation data involving turbulent energy dissipation by using equation (19). The latter approach is used in this study.

#### IV. DATA COLLECTION

##### A. DURATION AND LOCATION

Aerosol and meteorological data for this study were made available through Calspan Corporation, Buffalo, New York, from two separate experiments. During a ten day period in February 1977, Calspan Corporation provided limited meteorological and cloud physics support during a study of marine boundary layer phenomena conducted on the Gulf of Mexico (Mack and Katz, 1977). The experiment was performed on the Naval Coastal Systems Laboratory's (NCSL) offshore platform "Stage I" located approximately 20 km SW of Panama City, Florida as depicted in Figure 10.

A second experiment which provided data for this study was conducted along the coastal waters of Southern California (Figure 11) during a 12 day period in July 1977 aboard the Naval Postgraduate School (NPS) R/V Acania. Under contract from NPS, Calspan Corporation provided limited meteorological and aerosol physics support during a study of air quality parameters and marine boundary layer characteristics (Mack, 1977). This region contains primary shipping lanes and a number of drilling platforms all of which contribute to atmospheric contamination.

The following discussion will be limited to equipment used to measure the meteorological parameters actually analyzed in this study. A listing of the Panama City and

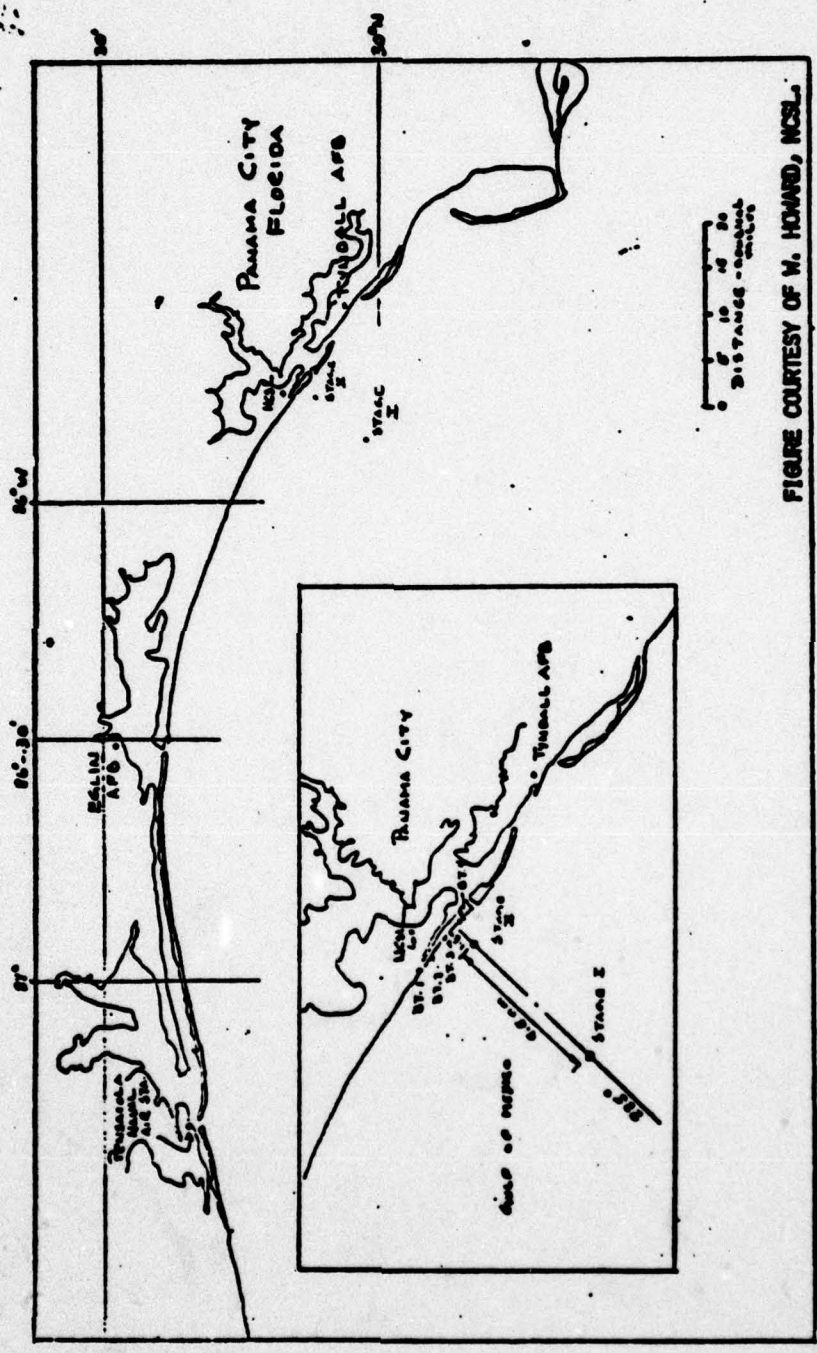


Figure 10. Location of NCSL Offshore Platform "Stage I"

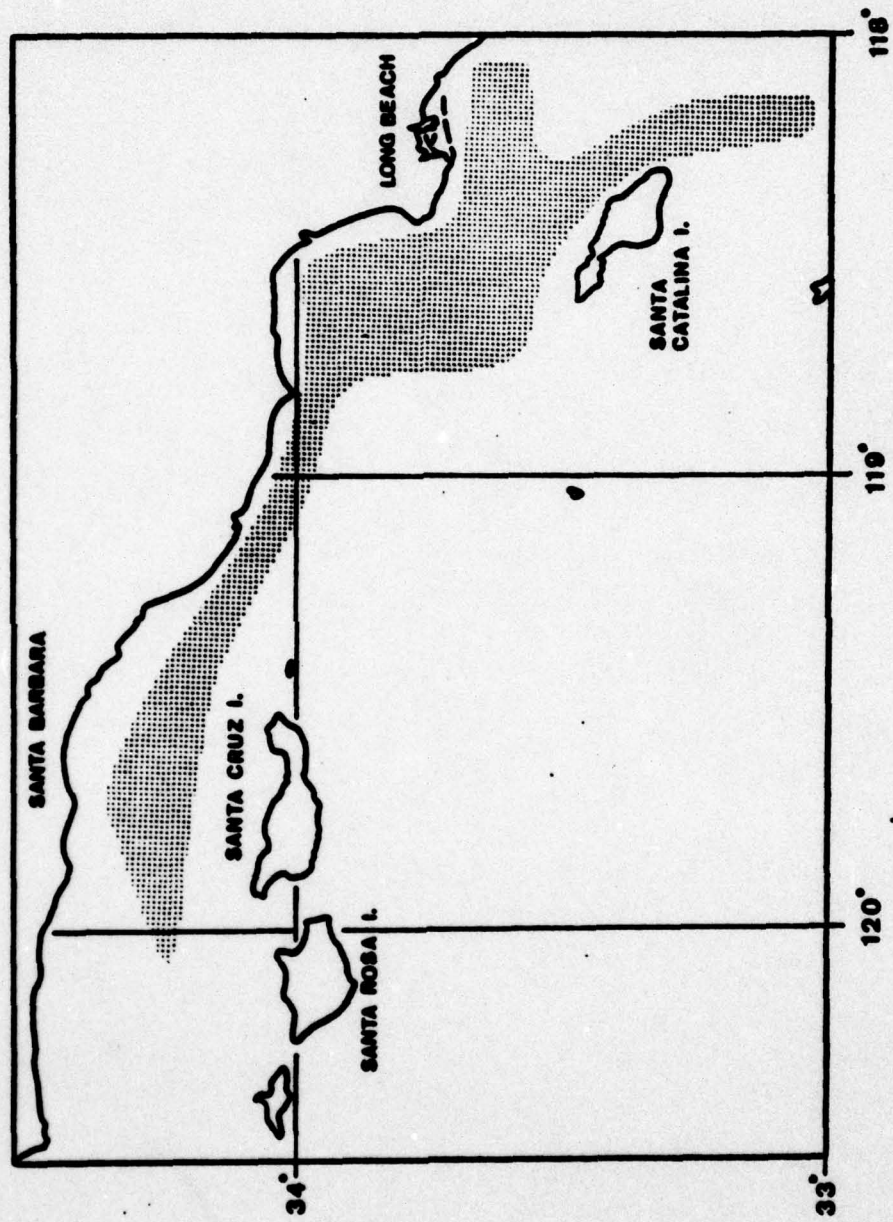


Figure 11. Location of Southern California Cruise



Southern California data may be found in Tables V and VI, respectively, at the end of the text.

#### B. PANAMA CITY INSTRUMENTATION

"Stage I" provided a stable platform for measuring the meteorological parameters necessary to describe and study the aerosol distribution and behavior in the marine boundary layer. The instrumentation installed by Calspan included a Sling Psychrometer, Bechman-Whitley wind system, Gardner small particle detector, and Royco Model 225 Particle Counter. The wind speed and direction was monitored continuously at the 20 meter level while wet and dry bulb temperatures were obtained hourly at the 17 meter level. A Foxboro temperature system (4 sensors) provided continuous temperature measurements at 4 levels; sea surface, 4.5, 9.0 and 24.5 meters. This data was recorded in an hourly log. Ten minute averaged aerosol size spectra were obtained continuously with the Royce counter at the 17 meter level, and a printout of aerosol concentration in 5 size intervals was provided every ten minutes. The Gardner Counter measured the concentration of particles greater than  $.0025 \mu$  diameter on an hourly basis.

The majority of the time the Royco instrument operated in "threshold" mode where number concentration (per 2.8 liters) of particles greater than the following size ranges were measured:  $0.5 \mu$ ,  $0.7 \mu$ ,  $1.4 \mu$ ,  $3.0 \mu$ , and  $5.0 \mu$  diameter. For a shorter period of time the instrument was operated in the "window" mode producing number concentrations between the

above size ranges. The particle counter and sensor are shown in Figure 12. The environmental air was drawn continuously through a sampling line of 3 meter length and 5 cm inside diameter. The flow rate through the counter's sensing volume was set at 2.8 liters per minute.

The Royco Model 225 sampler utilizes a near forward scattering optical system (Figure 13) which is ideal for monitoring large volumes of ambient gases where suspended particles can vary widely in composition, size, and optical properties. The aerosol is drawn through the sensor into a beam of focused light. As each particle passes through the illuminated volume, it scatters a pulse of light which is then detected by a photomultiplier tube. The photomultiplier output is then processed electronically to produce a pulse height spectrum from which the particle size spectrum is deduced. The height of each pulse is proportional to the square of the diameter of the particle.

Whitby and Liu (1973) note that the important characteristics of an optical counter are the sampling flow rate and the size of the optical viewing volume. The sampling flow rate determines the minimum counting period needed to obtain a statistically accurate count, and the size of the optical viewing volume determines the maximum aerosol concentration the instrument can accept without loss of particle count due to "coincidence", i.e., the loss of particle count due to the presence of more than one particle in the optical viewing volume at the same time. The viewing volume of the Royco 225 is  $4.0 \text{ mm}^3$  and the collection aperture half angle is 25 degrees.

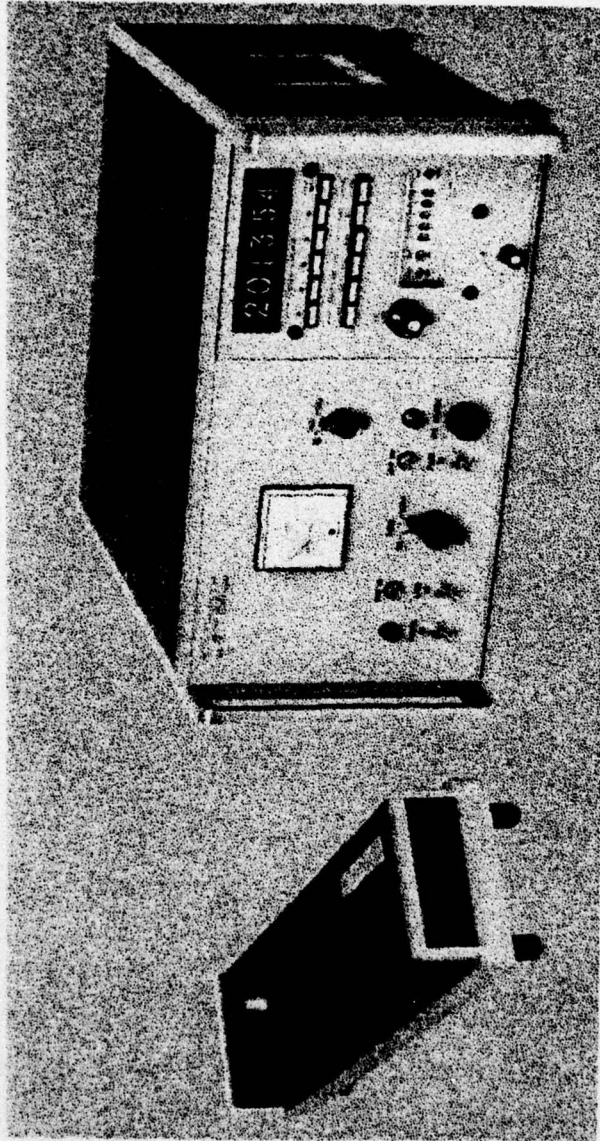


Figure 12. Royco 225 Particle Counter and Sensor

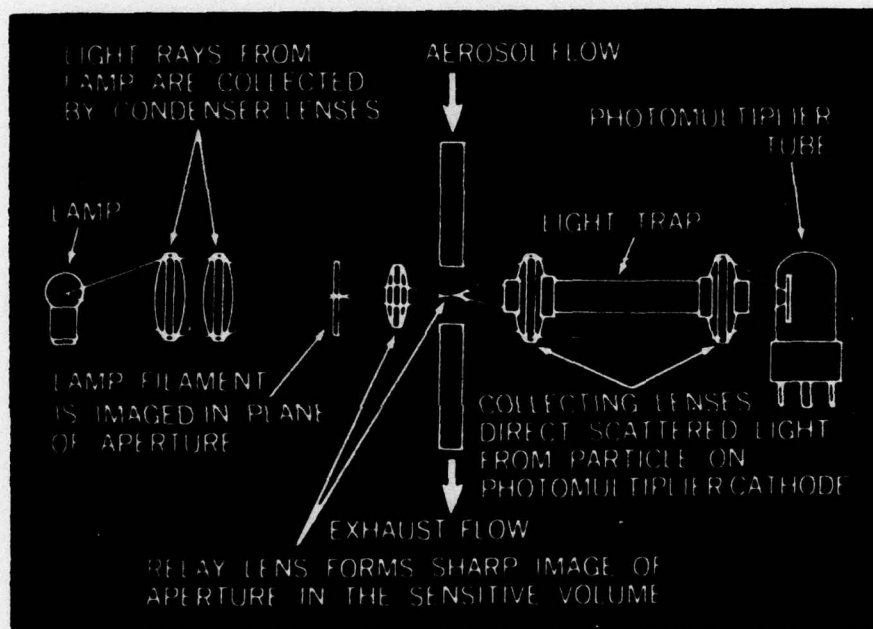


Figure 13. Near Forward Scattering Optical System

This model is also equipped with a sheath air inlet which diverts part of the aerosol stream through an external filter before reentry to the viewing volume. This sheath improves the performance of the instrument by preventing the recirculation of particles in the optical chamber and by confining the aerosol stream to a narrower region. Thus, the broadening of the pulse spectrum due to variation in illuminating intensity is reduced.

#### C. SOUTHERN CALIFORNIA INSTRUMENTATION

The location of the sensors aboard the R/V Acania are shown in Figure 14. Again, a Royco Model 225 Optical Particle Counter was used to measure the aerosol concentration of the coastal marine boundary layer. This instrument was operated continuously in the threshold mode where number concentration (per .28 liters) of aerosols greater than the following size ranges were measured: 0.3  $\mu\text{m}$ , 0.6  $\mu\text{m}$ , 1.2  $\mu\text{m}$ , 3.0  $\mu\text{m}$ , and 5.0  $\mu\text{m}$  diameter. The mainframe and sensor were located near the bridge of the Acania with the origin of the sampling line positioned forward of the pilot house roof at a height of 7 meters above the sea surface. The sampling line was 6 meters long with an inside diameter of 5 cm. The air was sampled through the viewing volume at a rate of .28 liters per minute. A Gardner small particle detector was again used to measure the Aitken nuclei concentration.

A sling psychrometer was used to measure the wet/dry bulb temperatures and relative humidity determined from psychrometric tables for a height of 5 meters. The mean wind

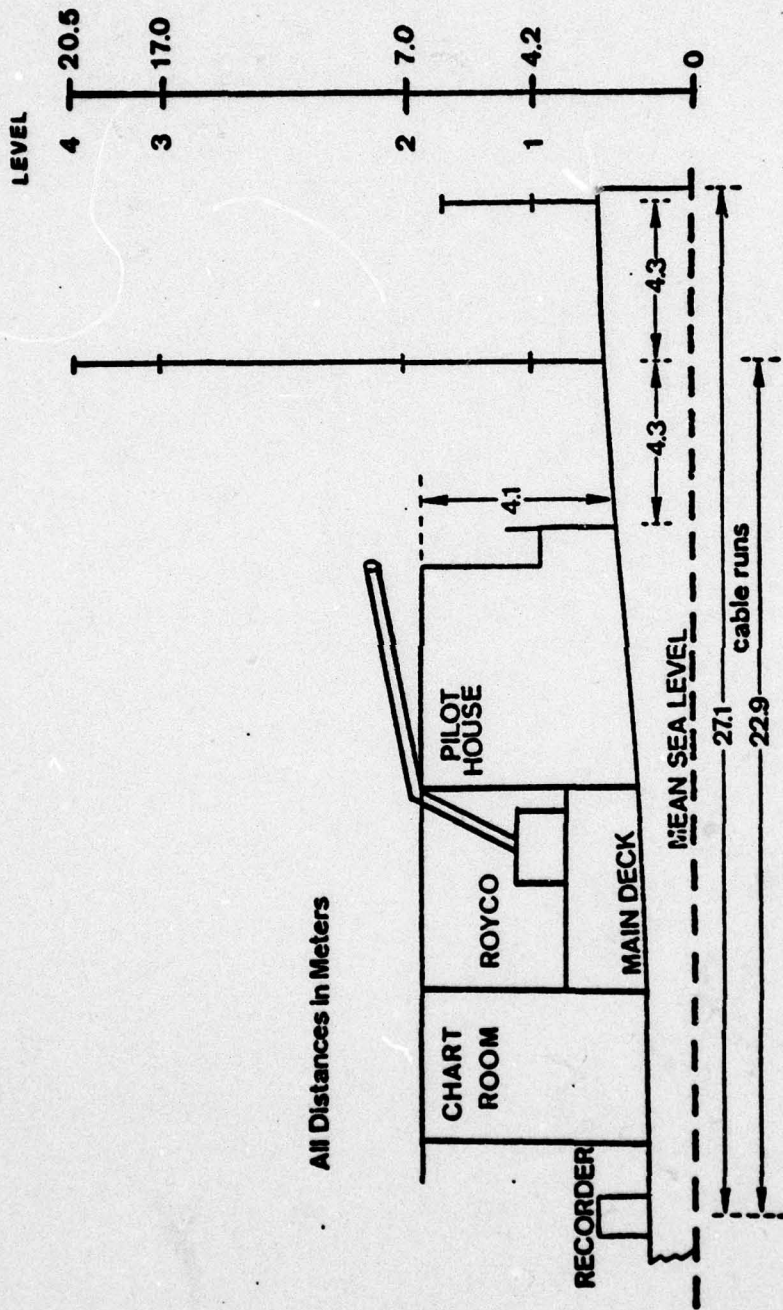


Figure 14. Sensor Locations on Board the R/V Acania

measurements were obtained at four levels using cup anemometer wind profile register systems supplied by the NPS. Calspan recorded the wind, humidity, and aerosol measurements in an hourly log.

Velocity fluctuation measurements were obtained with Thermo-Systems Model 1210 hot wire anemometer probes mounted with hot film sensors (platinum coated, 60 mil quartz fibers) installed by the NPS. The anemometer was a Thermo-Systems Model 1054B. The sensors were small enough to resolve the viscous dissipation scale without making corrections for wire length. Wind fluctuation data were recorded on a 14 channel tape recorder. The placement of these sensors required exceptionally long cable runs. Therefore, adjustments were made in the bridges for resistance and capacitance of the wirelength to insure a correct response.

The mean and fluctuation wind data were logged by the NPS developed MIDAS (Microprogrammable Integrated Data Acquisition System). This system is fully automated to sample the tailored list of sensors every 30 seconds and 20 minute averaged output values were printed.

## V. ANALYSES PROCEDURES

### A. VELOCITY FLUCTUATION ANALYSIS

The dissipation of turbulent kinetic energy,  $\epsilon$ , can be related to the mean wind velocity at any given level,  $\bar{u}$ , and the RMS value of the velocity fluctuation,  $\overline{u'^2}$ , in a frequency band specified by a lower frequency limit,  $f_l$ , and an upper frequency limit,  $f_u$  (Fairall, et al., 1977). The relationship is

$$\epsilon = \frac{(4/3)^{3/2} (u'_{\text{RMS}})^3}{(\bar{u}/2\pi)[f_l^{-2/3} - f_u^{-2/3}]^{3/2}} \quad (20)$$

In this procedure recordings were made of both the cup anemometer wind speed and the corresponding hot wire voltage output. The sensor wind speed is given by

$$v = V_o^2 + B(\bar{u})^{1/2} \quad (21)$$

where  $v$  is the hot wire voltage output, and  $V_o^2$  and  $B$  are constants obtained by laboratory calibration using a TSI Model 1125 Calibrator. Differentiation of equation (21) produces the following relationship between the velocity fluctuation and the voltage fluctuation:

$$u'_{\text{RMS}} = \frac{4v(\bar{u})^{1/2}}{B} v'_{\text{RMS}} \quad (22)$$



Substitution into equation (20) yields

$$\epsilon = \frac{(4/3)^{3/2} [4v(\bar{u})^{1/2}/B]^3 (v'_{RMS})^3}{(\bar{u}/2\pi) [f_l^{-2/3} - f_u^{-2/3}]^{3/2}} \quad (23)$$

Values of  $f_l = 5$  Hz and  $f_u = 200$  Hz were selected for the cruise and since amplifiers with known gains,  $G$ , were required, further reduction leads to

$$\epsilon = (3.53 \times 10^3) [V_o^2 + B(\bar{u})^{1/2}]^{3/2} (\bar{u})^{1/2} [v'_{RMS}/BG]^3 \quad (24)$$

The friction velocity,  $U_*$ , was then calculated from equation (19) for each of three levels and averaged to produce over 400 values from 19-27 July. Voltage fluctuation data from level 3 proved to be erroneous and were not included in the calculations. Obviously erroneous values of  $U_*$  owing to erratic behavior were also neglected.  $U_*$  values were then averaged about the aerosol observation times to correspond to a given aerosol distribution.

#### B. AEROSOL ANALYSIS

Analyses were performed on 215 aerosol samples during the SC cruise which were confined to the time period of the valid velocity fluctuation measurements. The observations included date and time, humidity, relative wind speed and direction, ship's speed and heading, Aitken concentration, and aerosol concentration as determined by the Royco 225 optical counter (Table VI). Wind and ship's speeds were recorded in knots.

The analyzed aerosol observations for the PC experiment were limited to 137 cases during the period 18-23 February. Cold frontal passage at approximately 0000Z, 24 February and subsequent advection of continental dust through 25 February were reasons for neglecting the aerosol samples for these days. Aerosol counts prior to 18 February were determined with the Royco instrument in the window mode and were not included in this study. Observations were generally made hourly and recorded in a log. They included date and time, humidity, wind speed and direction (knots), Aitken concentration, and data from the optical particle counter (Table V).

Computer programs were developed to plot the aerosol size distribution as a function of radius (R) in microns versus  $dN/d \log R \text{ (cm}^{-3}\text{)}$  where N is the number of particles greater than a given radius as measured by the Royco instrument. The program also included provisions to plot size distributions predicted by Fitzgerald's model. For this the observed relative humidity and wind speeds were used with equation (7). Initially the average aerosol distributions for both the SC and PC experiments were computed and compared to the respective predicted model distributions.

Subsequently, the variations in the average aerosol distributions with respect to four different categories of wind speed, relative humidity, and friction velocity were plotted for the SC data. The categories chosen for each of the above respective parameters are as follows: 0-2, 2-5, 5-8, 8-12 m/sec; 90-99, 80-90, 70-80, and 60-70 percent; and 0-.15, .15-.25, .25-.35, .35-.70 m/sec.

Friction velocity data was not available from the PC experiment; therefore, variations in the aerosol distributions were plotted with respect to categories of wind speed and humidity only. Because of essentially different meteorological conditions, the categories were chosen as follows: 0-3, 3-7, 7-10, and 10-15 m/sec; and 85-99, 70-85, 55-70 and 40-55 percent.

Visual inspection of these plots may indicate satisfactory relationships between the aerosol concentration and the above parameters. However, a statistical means of viewing these relationships was also deemed necessary. Wind speed, humidity, and  $U_*$  values were cross correlated with number concentration of particles in graduated size ranges. This procedure was accomplished by a Biomed Regression/Correlation computer program which produced corresponding correlation coefficients.

The nature of the diurnal variation of the aerosol concentration during the SC and PC experiments was investigated in this study. A computer program averaged the aerosol concentrations, wind speeds, humidities, and friction velocities about each hour and plots showing variations with time are produced. The aerosol plots depict the number of particles per  $\text{cm}^3$  within specified size ranges versus time. The SC data produced curves representing the number of particles between the following size ranges: .15-.30  $\mu$ , .30-.60  $\mu$ , .60-1.5  $\mu$ , and 1.5-2.5  $\mu$  radius. Diurnal variation of concentration for the PC data utilized the following slightly different size ranges: .25-.35  $\mu$ , .35-.70  $\mu$ , .70-1.5  $\mu$ , and 1.5-2.5  $\mu$  radius.

Finally, diurnal variations of the aerosol size distribution for the SC and PC experiments were calculated using techniques similar to those described above. Average size distributions for the following two time periods were plotted: 0000-1200 hrs and 1200-2400 hrs.

### C. ERROR ANALYSIS

The optical particle counter has an advantage over the membrane filter or impactor sampling techniques. For example, the latter require the samples to be taken to a lab for microscopic inspections and the aerosols may possibly be disturbed or altered due to contamination. Although the optical counter provides continuous "in situ" aerosol measurements, there are ample causes for counting errors. Because light scattering is a function of size, shape, and refractive index of the particles, careful calibration is necessary.

The Royco 225 model counter used in these experiments was calibrated using monodisperse latex spheres of known refractive index (1.6). Laboratory experiments by Lieberman and Allen (1969) showed a good correlation between the theoretical response curve for a near forward optical system and measurements using latex sphere and glass beads of refractive index 1.6 (Figure 15). Of most significance is the "fold" in the curve or zone of multi-valued response in the region of 1  $\mu$  diameter. Figure 16 is provided to illustrate how the response curve varies with particles of different refractive index. It is evident that when measuring particles of refractive index 1.6, a zone of ambiguity exists between

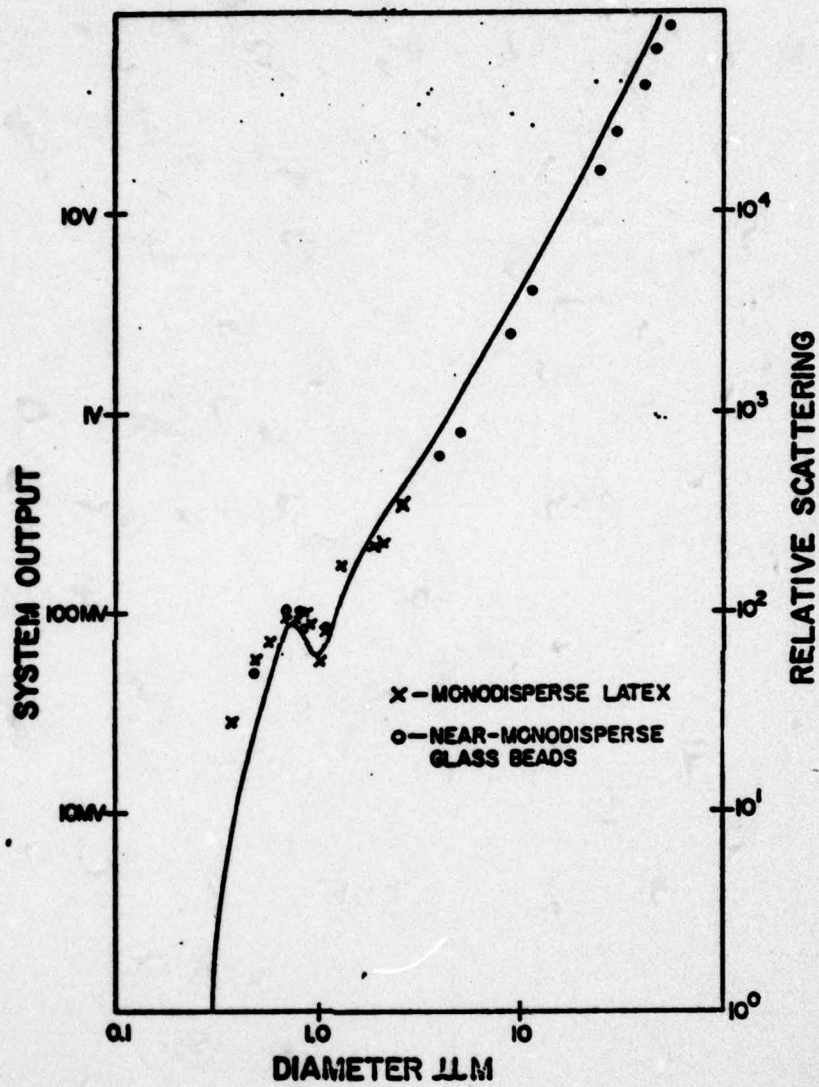


Figure 15. Theoretical Response Curve and Experimental Results for a Forward Scattering System

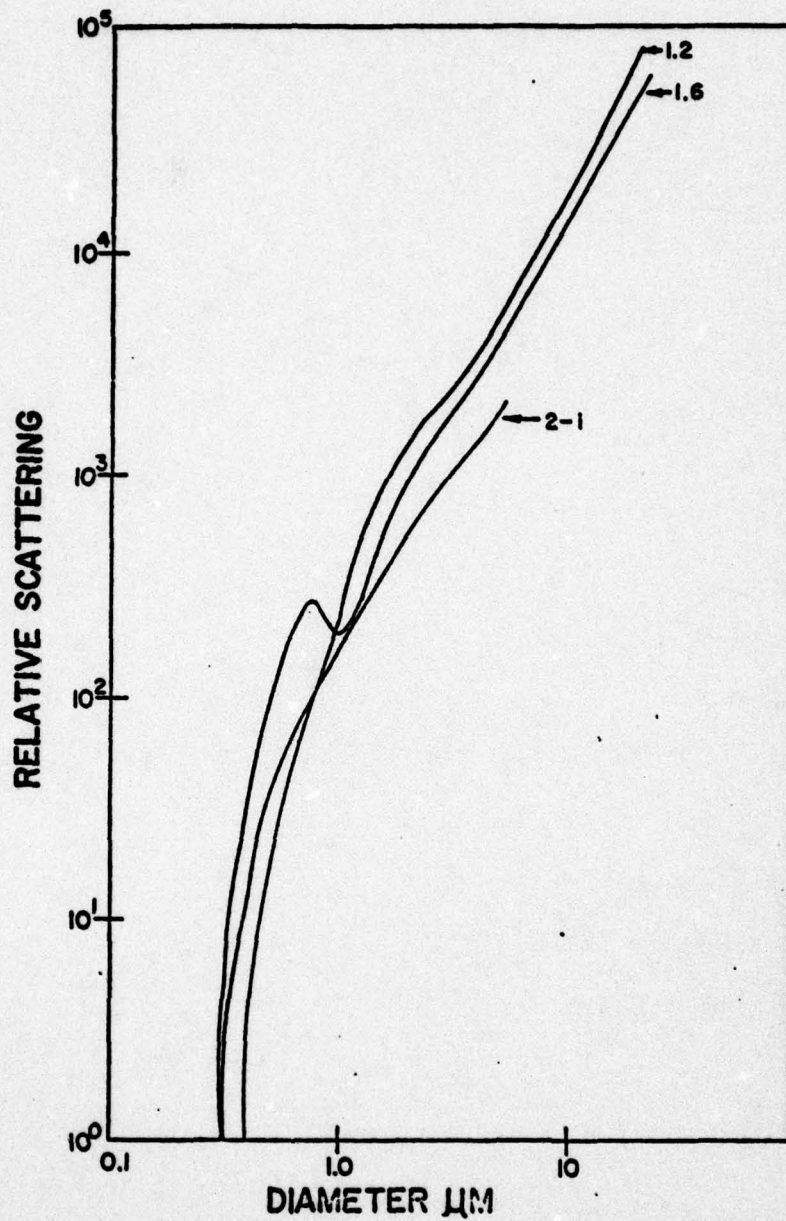


Figure 16. Dependency of the Response of a Forward Scattering System on Refractive Index

approximately  $.3 \mu$  and  $.6 \mu$  radius and may vary with the aerosol refractive index. Lieberman and Allen (1969) state that the instrument will still produce valid data if the zone is encompassed within a size range or channel. Since the SC counter measured between  $.3 \mu$  and  $.6 \mu$  radius and the PC counter between  $.35 \mu$  and  $.70 \mu$  radius, it is assumed that this multi-valued zone is compensated for.

Counting errors can also arise from flow rate considerations. If the particle sizes are large and the number of particles small, enough particles must be counted to obtain good statistical resolution. When a small random number of particles is counted, the statistical error in counting is equal to the ratio of 1 over the square root of the number of particles counted (Zinky, 1962). The counter should be operated over a longer time period (10 minutes) to sample a larger volume or an increase in the flow rate will reduce the error. It then seems quite possible that the flow rate of the counter used in the SC cruise ( $.28$  liter/min) provided too small of a sampling volume to obtain an accurate count of the larger particles.

Zinky (1962) also states that a vertically aligned inlet tube is recommended to prevent any deposition in the line due to settling. It has already been mentioned that the sampling lines used in each experiment were considerably long and aligned horizontally. Many of the larger particles may not have remained airborne long enough to reach the illuminated volume.

Errors in the calculation of the friction velocity may have come from various sources. Since calculation of  $U_*$  from Equation (19) is only valid for near-neutral conditions, any substantial departure of the Richardson number from zero would result in inaccurate values. The measurement of the dissipation of turbulent kinetic energy was large dependent on the accuracy of the voltage output. The signal response is sensitive to electromagnetic energy, and any local radio or radar transmission may introduce noise to the system. Additionally, under the light wind conditions which prevailed on the SC experiment, the lateral motion of the anemometers due to ship pitch and roll may have resulted in erroneously high readings.



## VI. RESULTS

The data from the Southern California (SC) cruise proved to represent an atmosphere somewhat different from a typical marine environment. The Aitken particle population averaged almost  $8500 \text{ cm}^{-3}$  which is about 4 times higher than that observed by Hidy, et al. 130 km west southwest of Los Angeles. This high concentration is suspected to be due to a combination of pollution from merchant ships' exhaust, combustion from the drilling platforms, and offshore flow from the nearby populated coastal cities.

The average wind speed and relative humidity were 3 m/sec and 86 percent, respectively. This data was used to compute the prediction from Fitzgerald's model (Eq. 7) which is compared to the average SC distribution in Figure 17. The vertical bars represent one standard deviation either side of the mean. There is generally good agreement between the two below  $.4 \mu$  radius, with a larger experimentally observed concentration above this range. Although sea-salt production should have been minimal during this time period because of low wind speeds, the characteristic hump at around  $1 \mu$ , to a certain extent, reflects the contribution by sea salt nuclei. A similarity exists here with Moore and Mason's (1954) observation of a discontinuity where the slope changes and becomes rather steep in the region of the larger size range. The larger concentration in this range may be solely

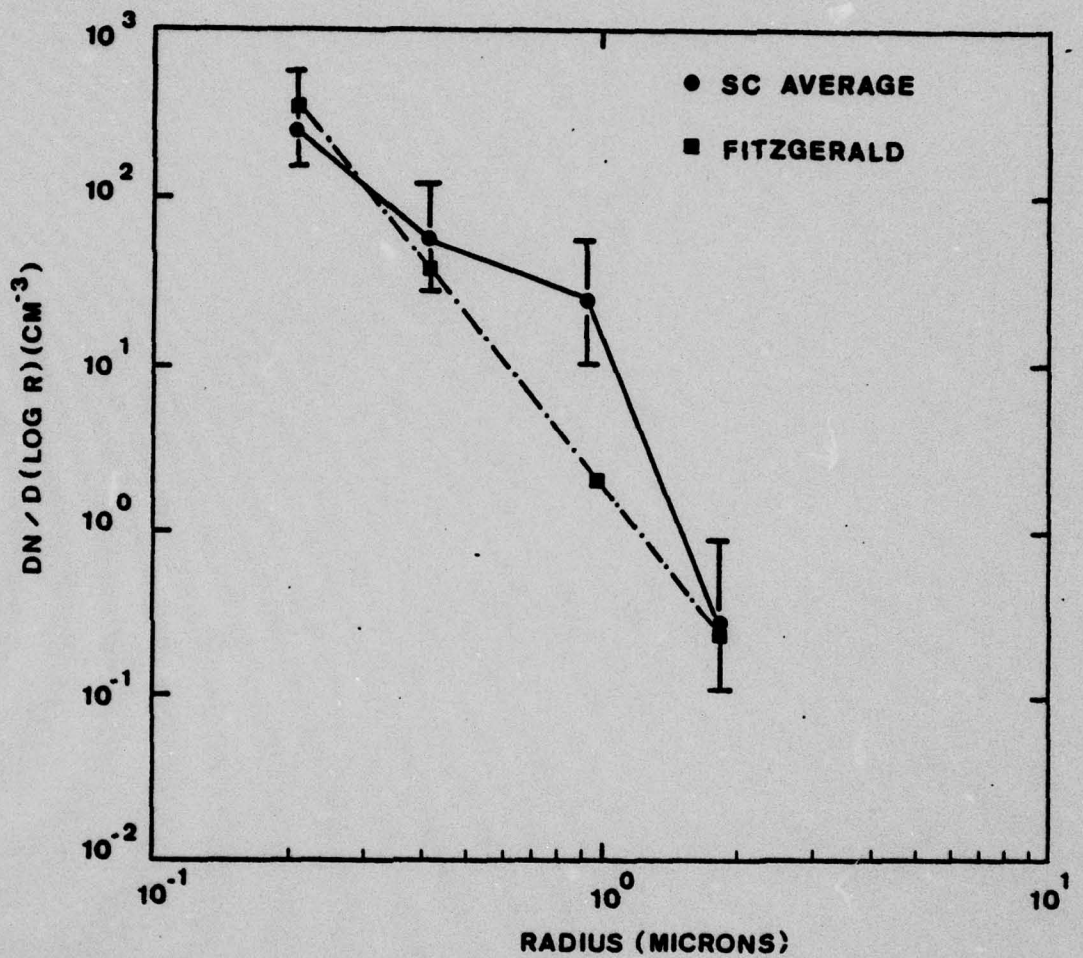


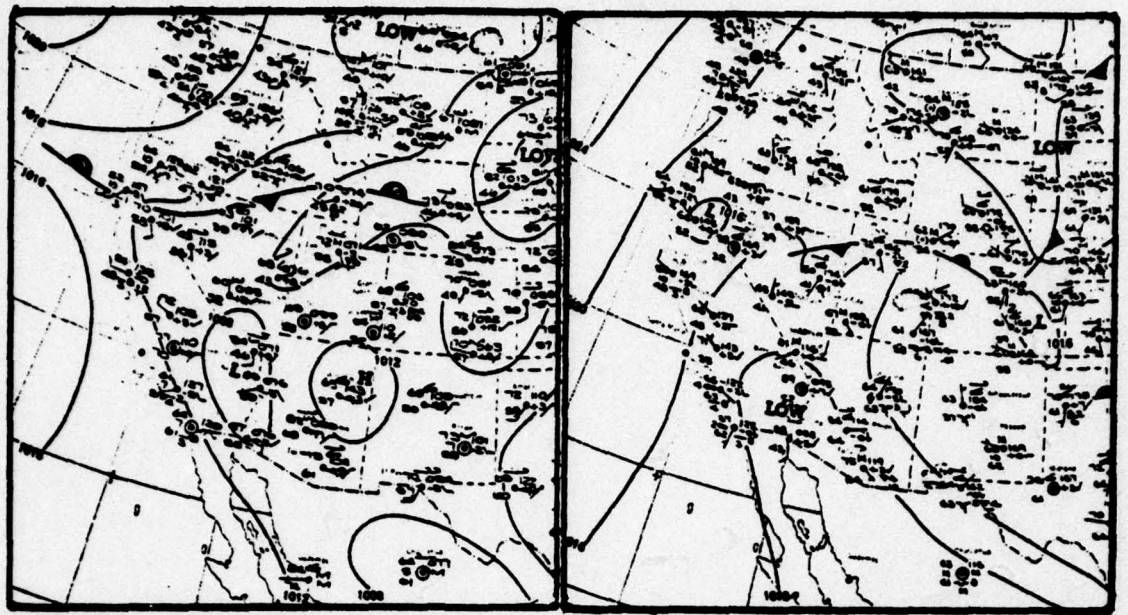
Figure 17. Average Aerosol Size Distribution for the SC Experiment and Distribution Predicted by Fitzgerald's Model

due to the influence of atmospheric contaminants such as combustion by-products, soil dust, or smoke. Considering previous experiments, this range does indeed contain a mixture of both continental and marine aerosols possibly resulting in the increase over Fitzgerald's model.

As previously mentioned, the low flow rate of the optical counter may account for the low concentrations at  $2 \mu$ . However, since the wind speed reached  $8 \text{ m sec}^{-1}$  only 6 times, this may have been a truly representative concentration of droplets as agreement is also shown with Fitzgerald's curve.

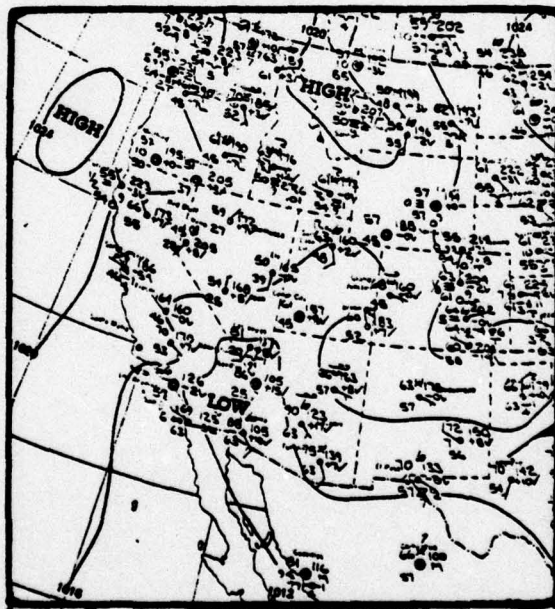
Figure 18 presents the synoptic situation for three days at the beginning, middle, and end of the experiment. A persistent thermal low is located in the desert area of Southern California and the isobaric pattern off the coast reflects a rather weak gradient. Therefore, smaller scale circulations should prevail in this area of little or no synoptic forcing.

Plots showing the variations of the average distributions with respect to wind speed, relative humidity, and friction velocity ( $U_*$ ), are shown separately in Figures 19, 20, 21. The number of observations in each category is placed in parentheses. These figures indicate that the size distribution has a better relationship with the relative humidity than to the wind speed and  $U_*$ . Correlation coefficients between these parameters and the number concentration of particles in a given size interval are produced in Table II. Since diurnal variations tend to reflect a negative relationship between relative humidity and wind



19 JUL

23 JUL



26 JUL

Figure 18. Synoptic Situation during the SC Experiment

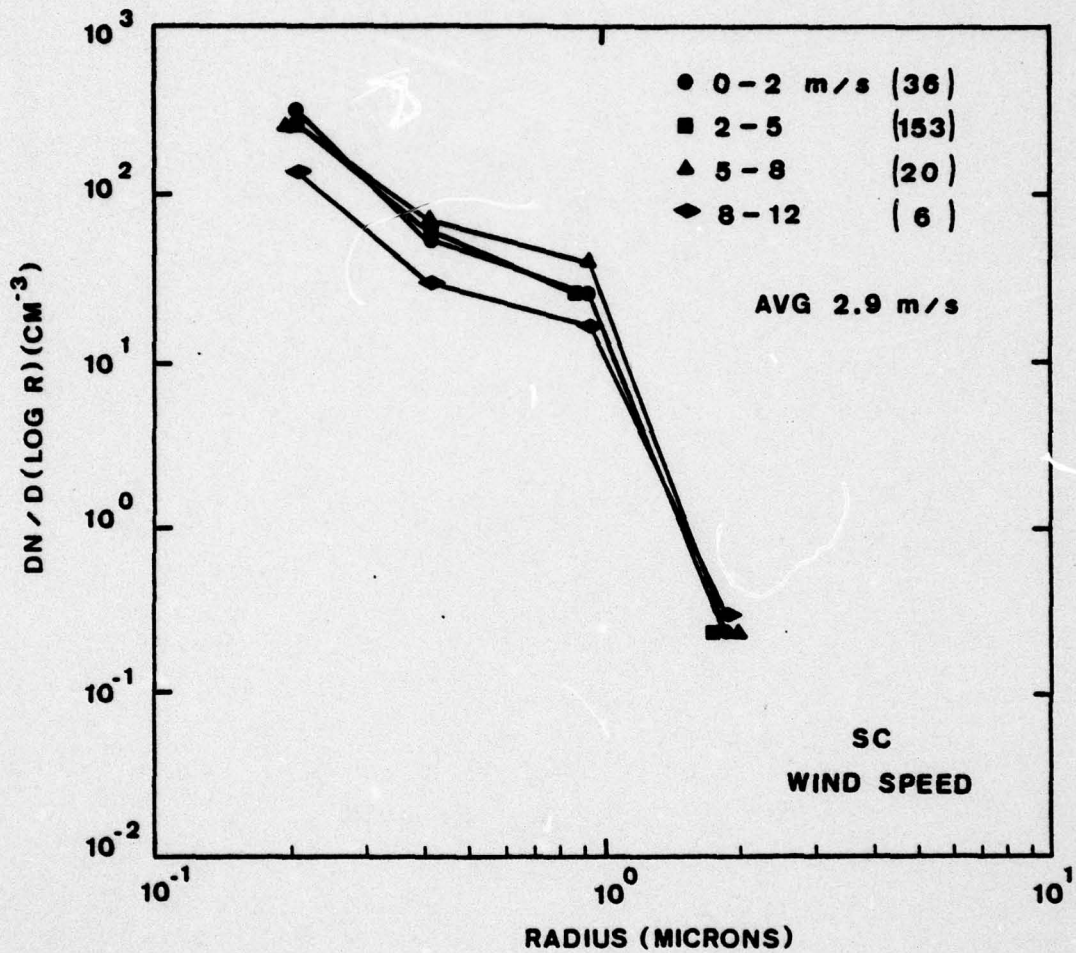


Figure 19. Variation of SC Size Distribution with Wind Speed

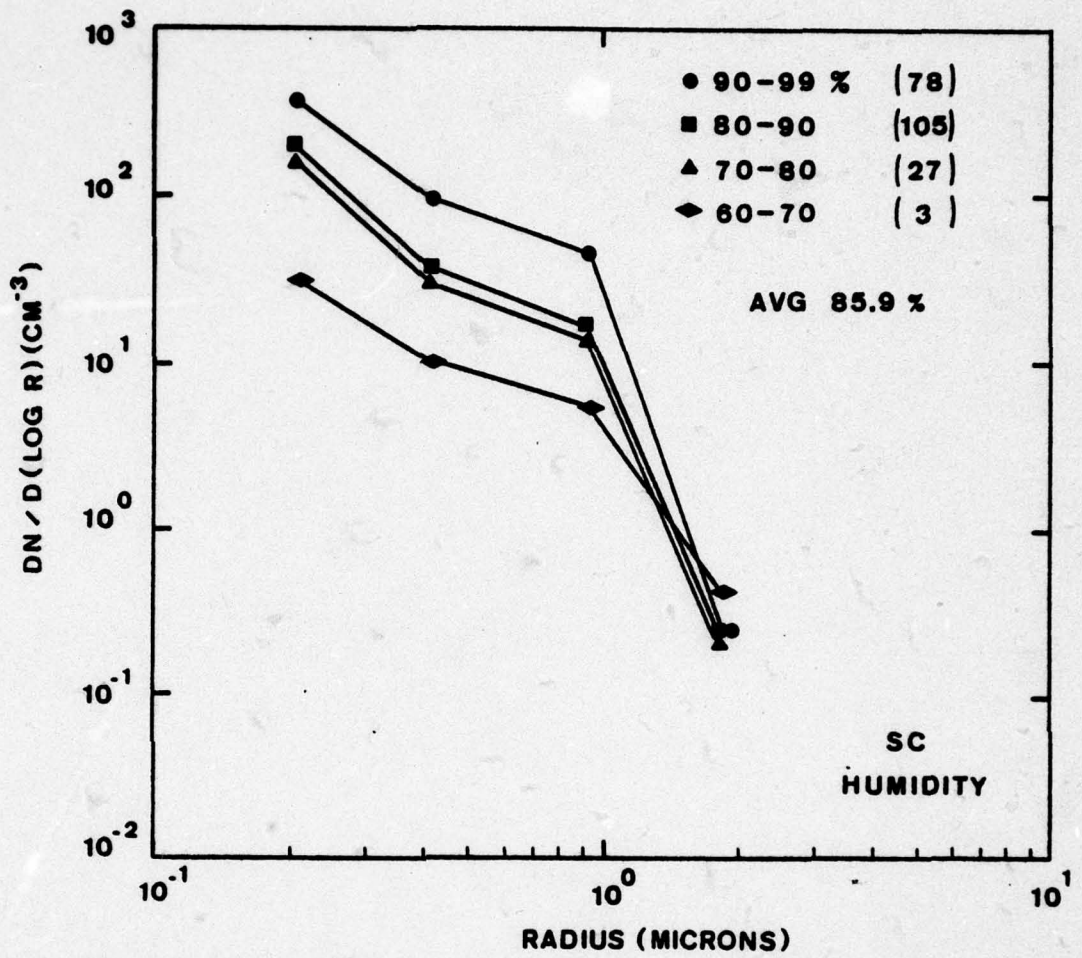


Figure 20. Variation of SC Size Distribution with Relative Humidity

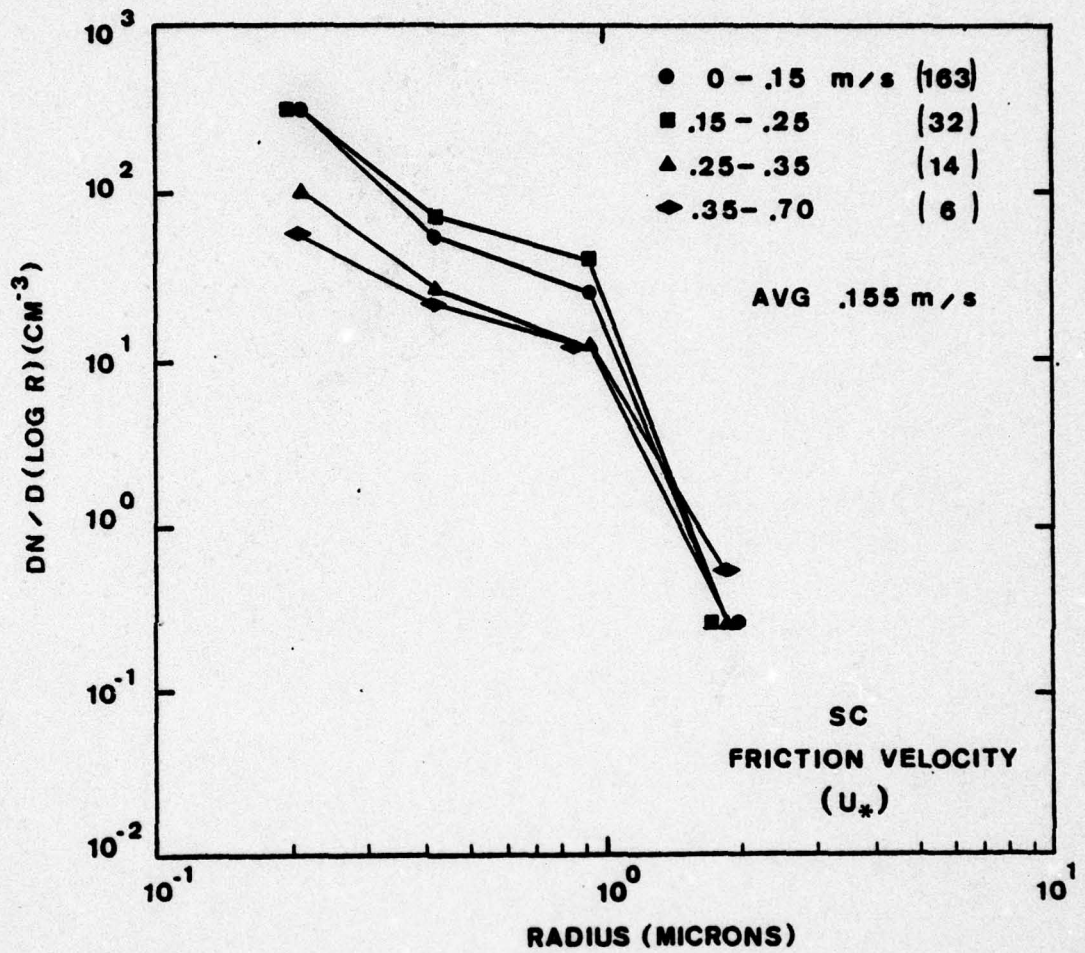


Figure 21. Variation of SC Size Distribution with Friction Velocity

SC

INTERVAL	RH	WIND SPEED	U*
.15 - .30 $\mu$	.375	-.127	-.162
.30 - .60 $\mu$	.383	-.056	-.107
.60 - 1.5 $\mu$	.376	-.017	-.077
1.5 - 2.5 $\mu$	-.037	-.022	.132

TOTAL

Table II. Correlation Coefficients for the SC Experiment



speed, the results in this table are not surprising. The negative correlation of relative humidity with the concentrations in the large size range indicates that sedimentation of large droplets, which grow with increasing humidity, is most important when the wind speed and sea surface production of salt nuclei are weak. Although these larger droplets also exhibit a small positive correlation with  $U_*$  while the wind speed correlation remains negative, this result does not appear to be significant.

An attempt was made to examine the influence of stability on the size distribution. The summer months are characterized by the occurrence of stratus and fog off-shore below the marine inversion. Two days are compared with the assumption that they represent the unstable and stable atmospheres. According to the daily observation log, stratus clouds in the morning becoming stratocumulus by afternoon were observed on 19 July. 26 July was characterized by clear skies. The average distribution for both days is presented in Figure 22. The correlation coefficients between concentration and wind speed and  $U_*$  show a trend toward positive values from the stable to the unstable day with  $U_*$  eventually becoming positively correlated in the unstable day (Table III). The increase in the size distribution on 26 July in the size range greater than  $.3 \mu$  seems to be due to increase in the average wind speed and occasional gustiness as whitecaps were reported during the afternoon. The stable stratification assumed in this case allows generated sea-salt nuclei to accumulate and the concentration to increase at the 7 meter

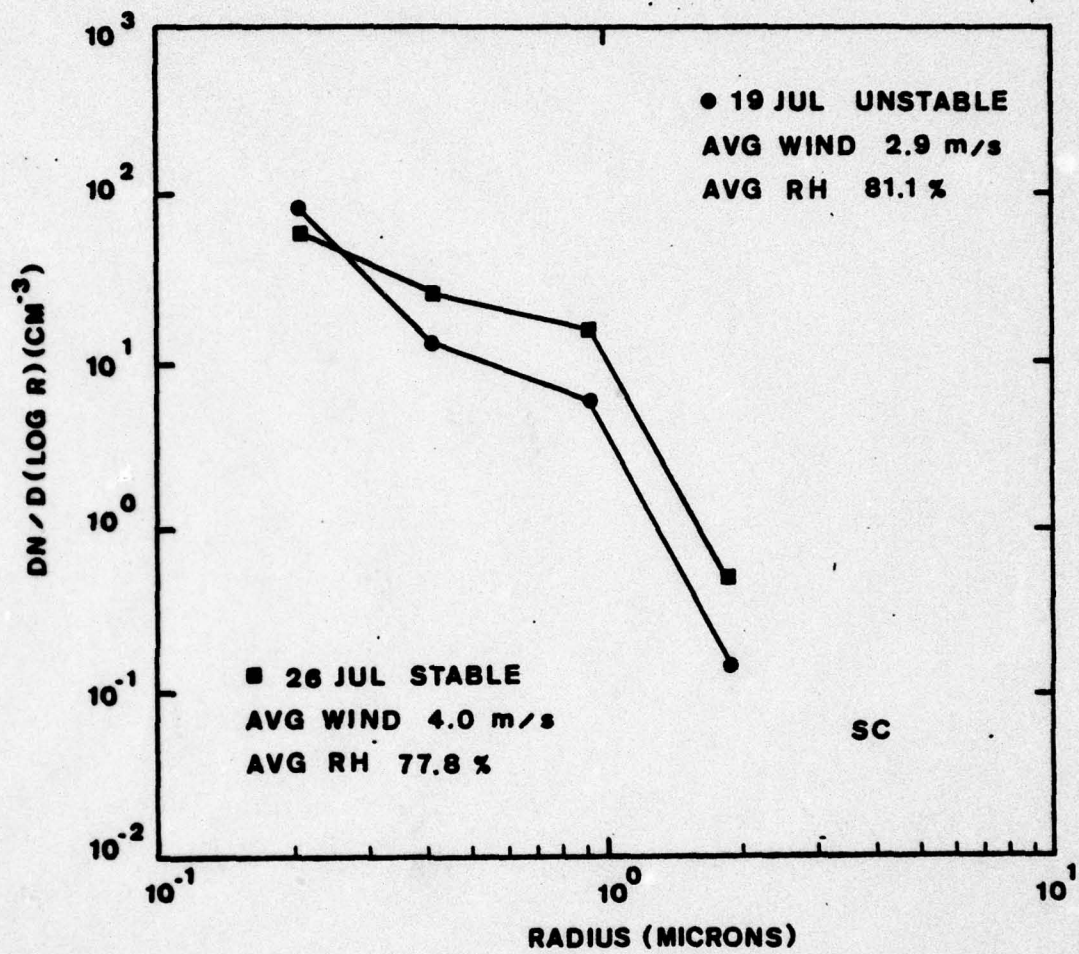


Figure 22. Average Size Distributions on 19 July and 26 July

SC

INTERVAL	RH	WIND SPEED	U*
.15 - .30 $\mu$	.814	-.205	.129
.30 - .60 $\mu$	.837	-.184	.165
.60 - 1.5 $\mu$	.792	-.110	.147
1.5 - 2.5 $\mu$	-.104	-.049	.075

19 JUL UNSTABLE

SC

INTERVAL	RH	WIND SPEED	U*
.15 - .30 $\mu$	.610	-.395	-.372
.30 - .60 $\mu$	.754	-.499	-.535
.60 - 1.5 $\mu$	.835	-.615	-.634
1.5 - 2.5 $\mu$	.355	-.413	-.273

26 JUL STABLE

Table III. Correlation Coefficients for 19 July and 26 July

level. The lower average wind speed associated with the unstable period does not allow for much sea-salt production. An unstable atmosphere can lead to convective processes which may vertically transport aerosols and create higher concentrations at an upper level as proposed by Ericksson (1959) and Toba (1965a & b). Hence, a decrease in the size distribution on 19 July is observed. This evidence gives credence to the possibility that friction velocity is a better indication of aerosol size distribution than wind speed. On both days the correlation of the concentration with humidity is lowest in the largest size interval. This relationship is most pronounced on the unstable days and may be explained by sedimentation due to mixing and resulting increased coalescence.

The averaged diurnal variations of wind speed, relative humidity, friction velocity, and aerosol concentration are shown in Figures 23, 24, and 25. Again the negative relationship of aerosols to wind speed and  $U_*$  in the size range of generally less than  $1 \mu$  is indicated. A satisfactory relationship with relative humidity is not evident and this is probably due to transport by a secondary circulation. A land-sea breeze type of effect could account for the observed decrease in concentration of the particles smaller than  $1.5 \mu$ . As the heating over the land generates an on-shore flow along the coast, the wind increases and persists through the afternoon. The average wind direction derived from the observations of five random days during the experiment is shown in Figure 26. A westerly wind is

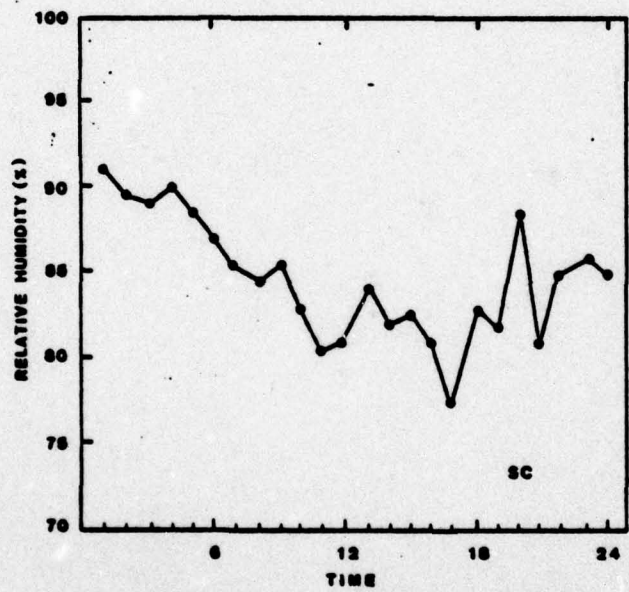
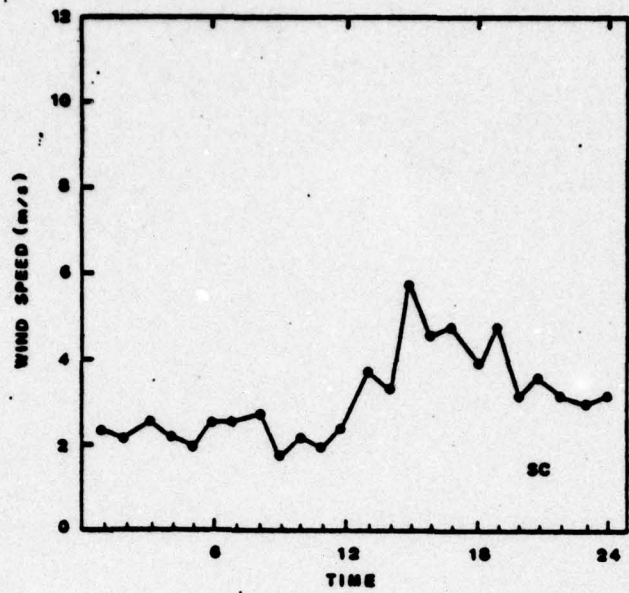


Figure 23. Average SC Diurnal Variations of Wind Speed and Relative Humidity

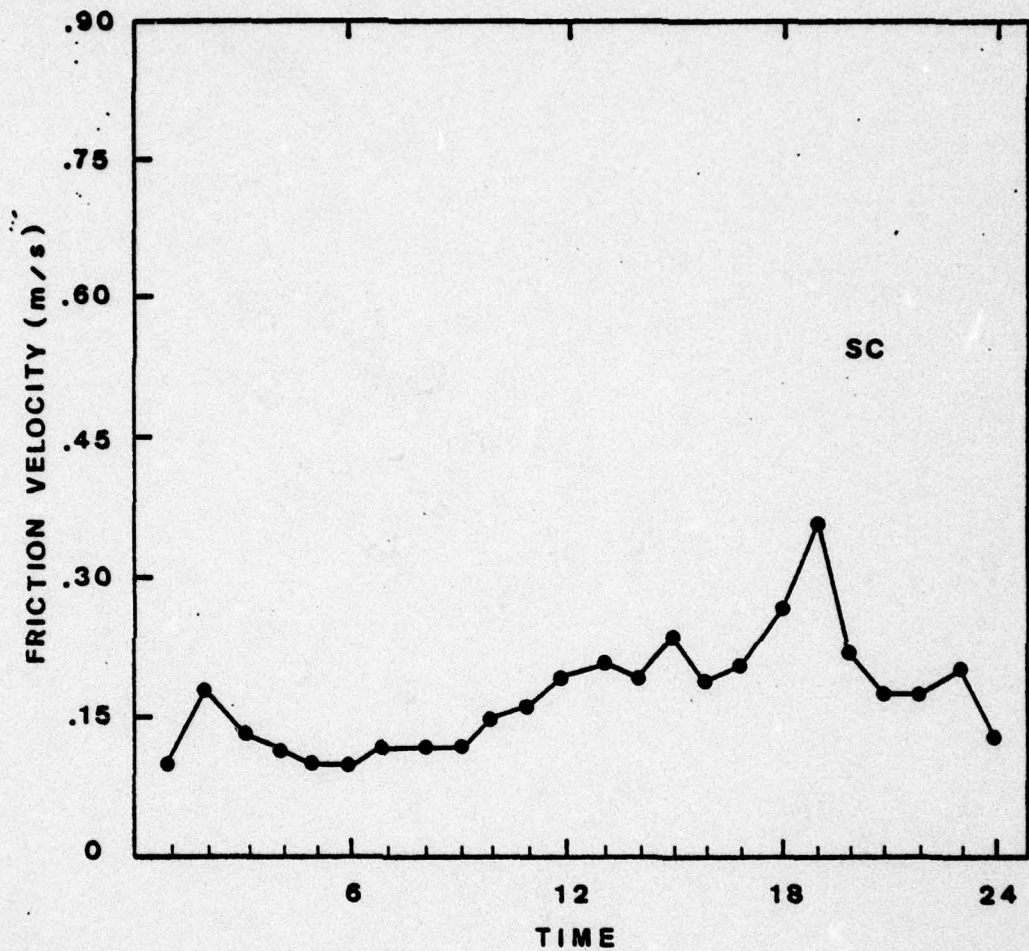


Figure 24. Average SC Diurnal Variation of Friction Velocity

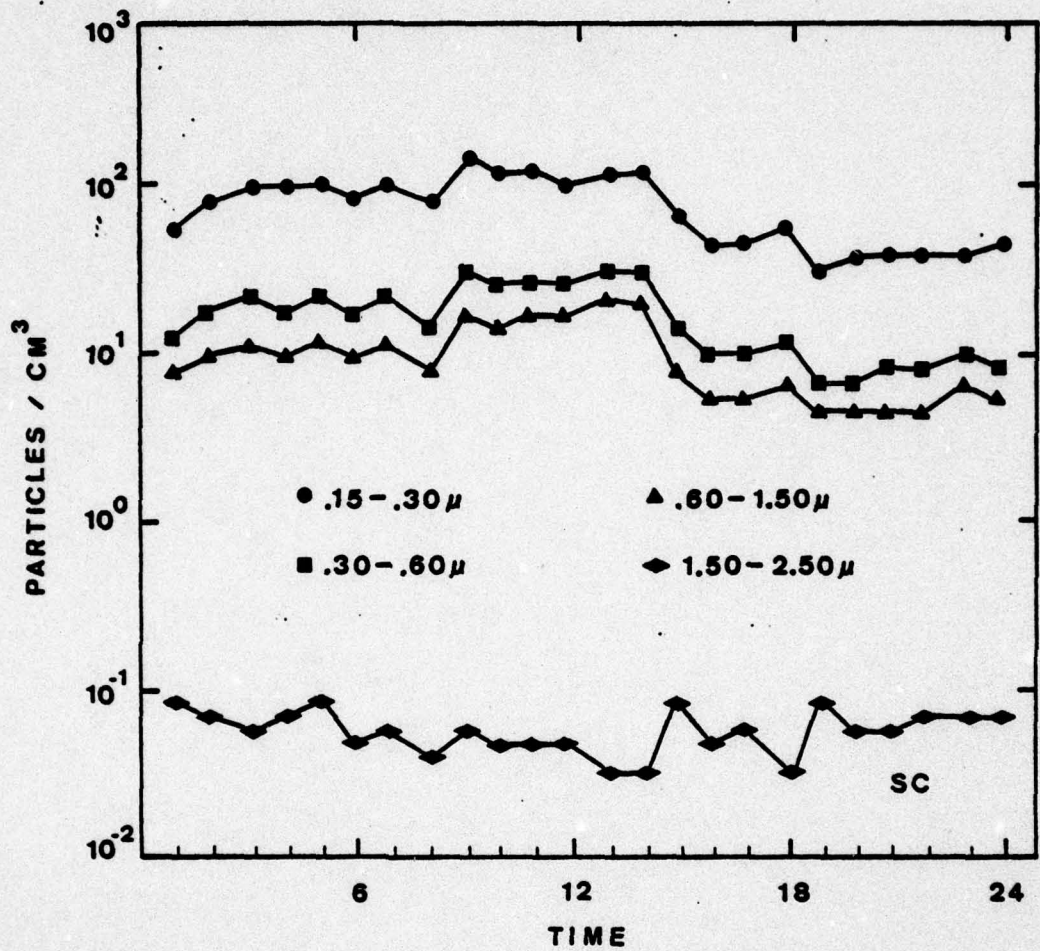


Figure 25. Average SC Diurnal Variations of Particle Concentrations

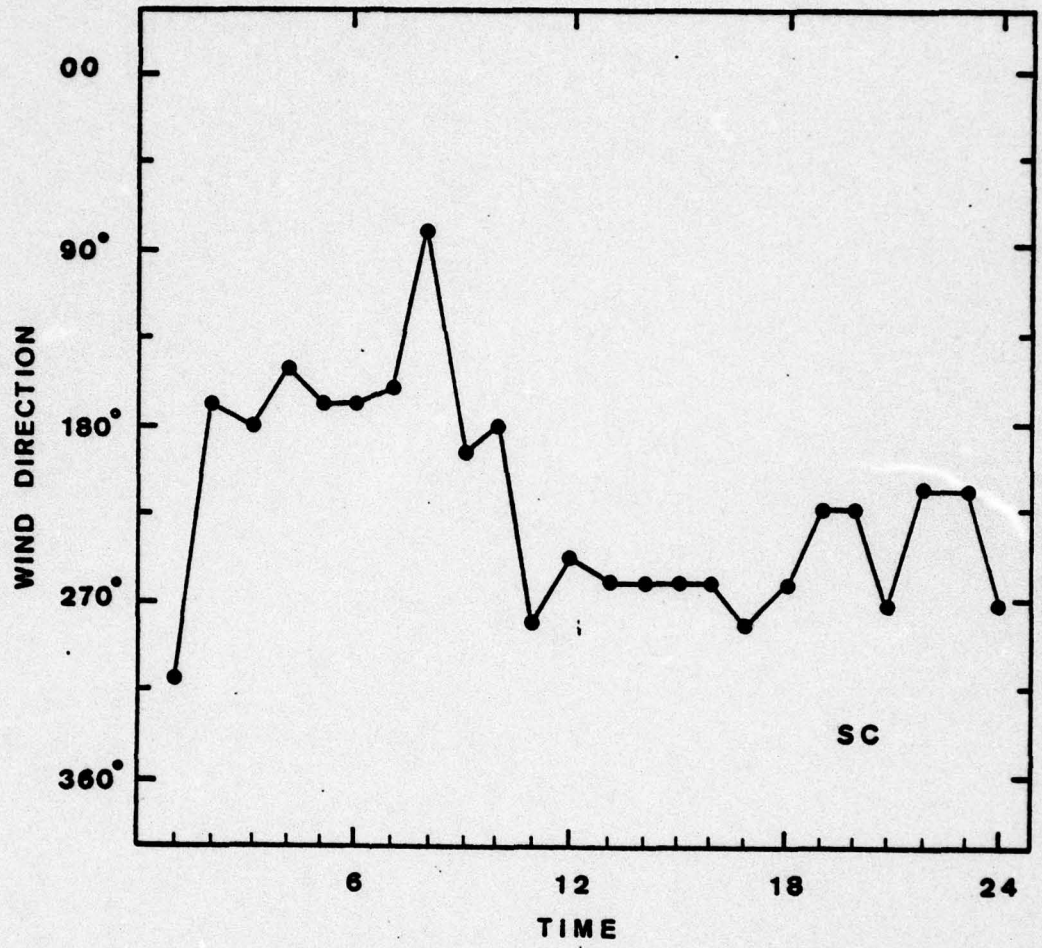


Figure 26. Average SC Wind Direction



seen to dominate during the peak wind periods. The decrease in the aerosol population may be explained by a horizontal divergence effect in the marine boundary layer as largest accelerations are found near the coast. The average size distributions displayed in Figure 27 reflect the decrease in the aerosol population due to this sub-synoptic circulation. Although the relative humidity increases slightly in the early evening hours, the smaller nuclei show a stronger relationship with the wind speed. This again implies that a large part of the coastal marine aerosol is of continental origin. The outflow of circulation aloft is probably responsible for the introduction of continental particulates to the marine environment. The minor peaks in the small particle concentration and also the somewhat greater increase in the large particles during the afternoon should be attributed to sea-salt production.

The Panama City (PC) observations more closely resembled a marine environment. The Aitken particle count was lower and averaged  $2600 \text{ cm}^{-3}$  while the distribution curve showed a marked change from the Southern California data. Winter time synoptic scale features predominate in this region of the Gulf Coast. Cold frontal passages and an accompanying influx of continental air into the Gulf of Mexico are frequent occurrences. Subsequent movement of the high pressure ridge into Florida and off its eastern seaboard provides the circulation which reestablishes moist southerly flow and return of the marine aerosol. Figure 28 provides the synoptic analyses for the period of the experiment. Stable

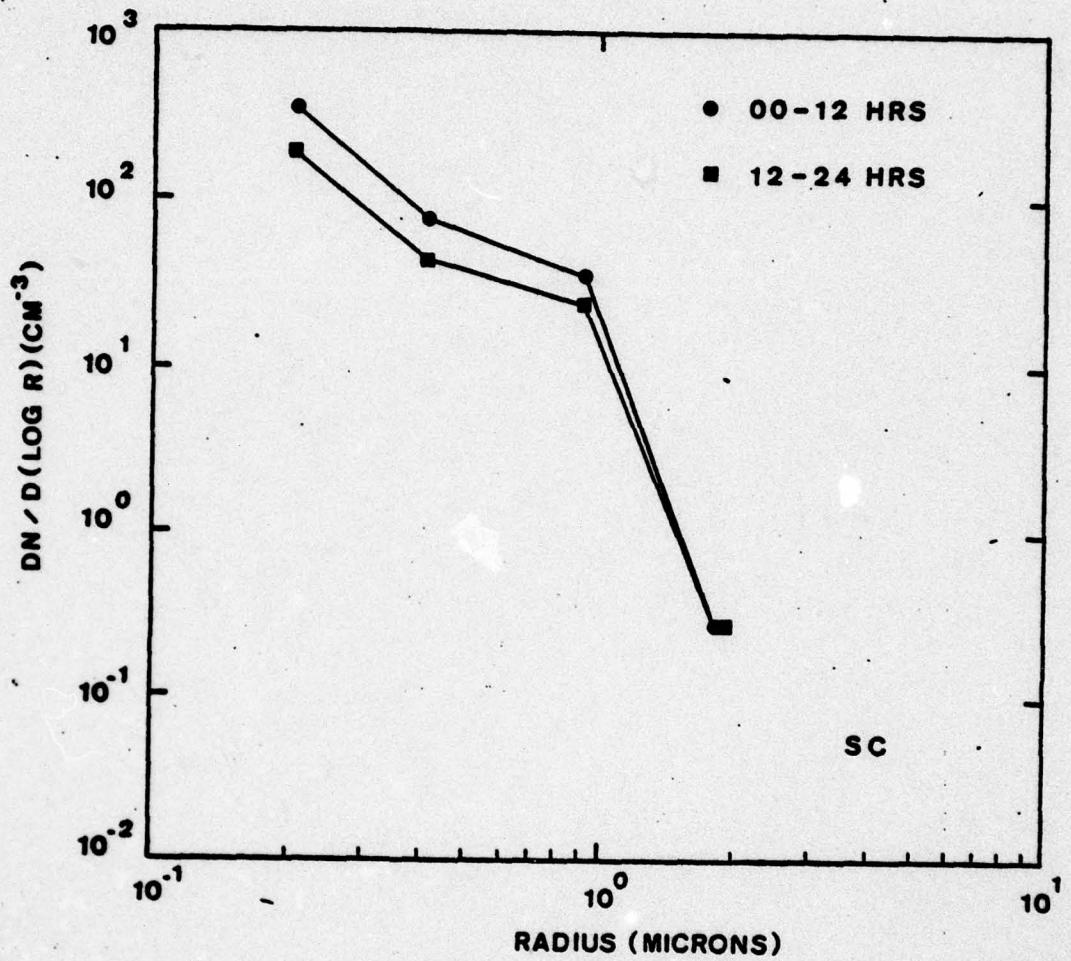
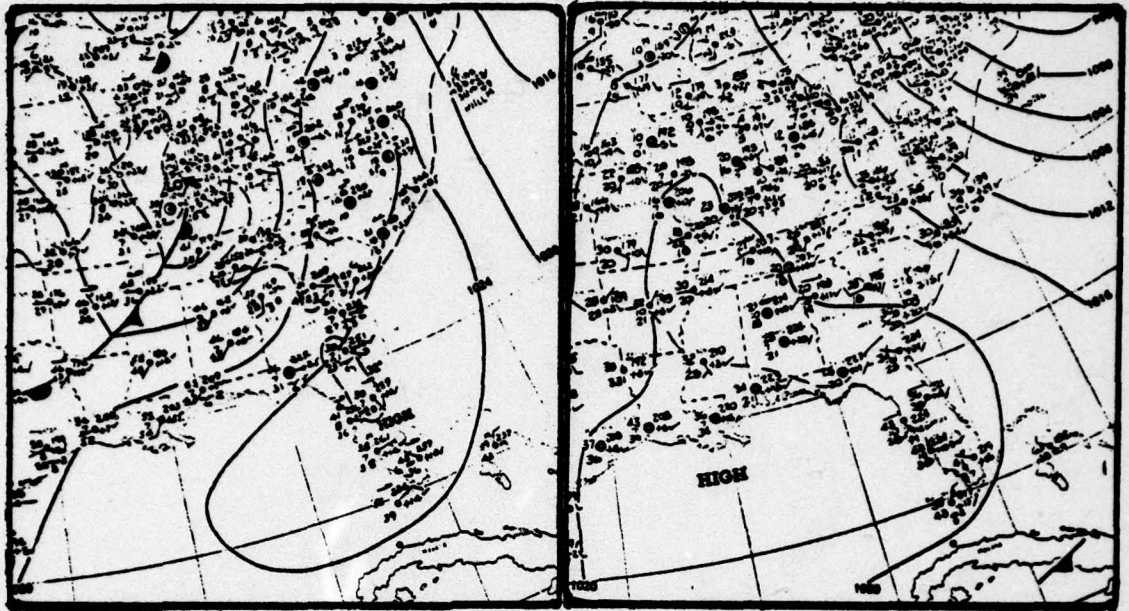
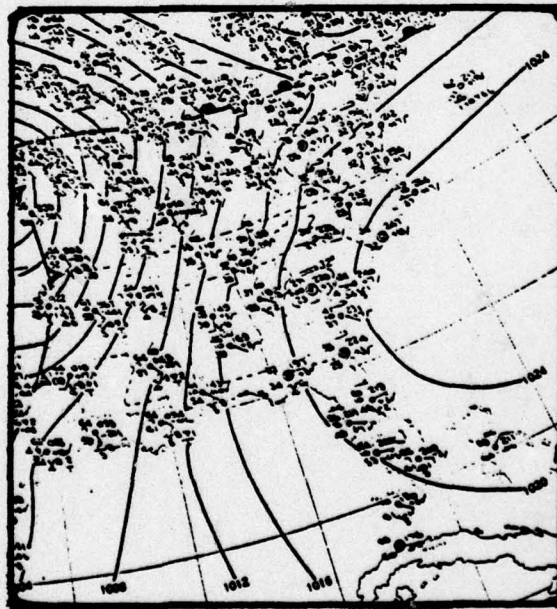


Figure 27. Average SC Diurnal Variation of the Aerosol Size Distribution



18 FEB

21 FEB



23 FEB

Figure 28. Synoptic Situation during the PC Experiment

conditions prevailed in the Panama City area at the beginning of the period; but, after frontal passage early on 20 February, south to southeasterly flow developed and persisted for the remainder of the experiment. The influx of this warm, moist air contributed a destabilizing effect in the lower levels of the marine boundary layer.

The average wind speed and relative humidity for PC were 8.4 m/sec and 71 percent, respectively. Fitzgerald's curve for these average conditions and the average aerosol distribution for the entire period are shown in Figure 29. Good agreement exists only for particle size range greater than  $.9 \mu$  radius. The observed concentrations are approximately an order of magnitude lower than Fitzgerald's prediction for aerosols smaller than  $.5 \mu$  radius. A significant aspect of the distribution is the positive slope observed between approximately  $.5 \mu$  and  $1 \mu$  radius which appears to be the result of sea-salt production. Actually, good agreement is shown with Blifford's (1970) observation off the Pacific coast with respect to both slope and number concentration.

Plots showing the effect on the average distributions due to wind speed and relative humidity separately are shown in Figures 30 and 31. Relatively good correlations seem to exist between these parameters and the aerosol distributions. Correlation coefficients are presented in Table IV. The synoptic scale effects predominate over diurnal variations and wind speed and relative humidity are both positively correlated to the concentration. The highest correlation of

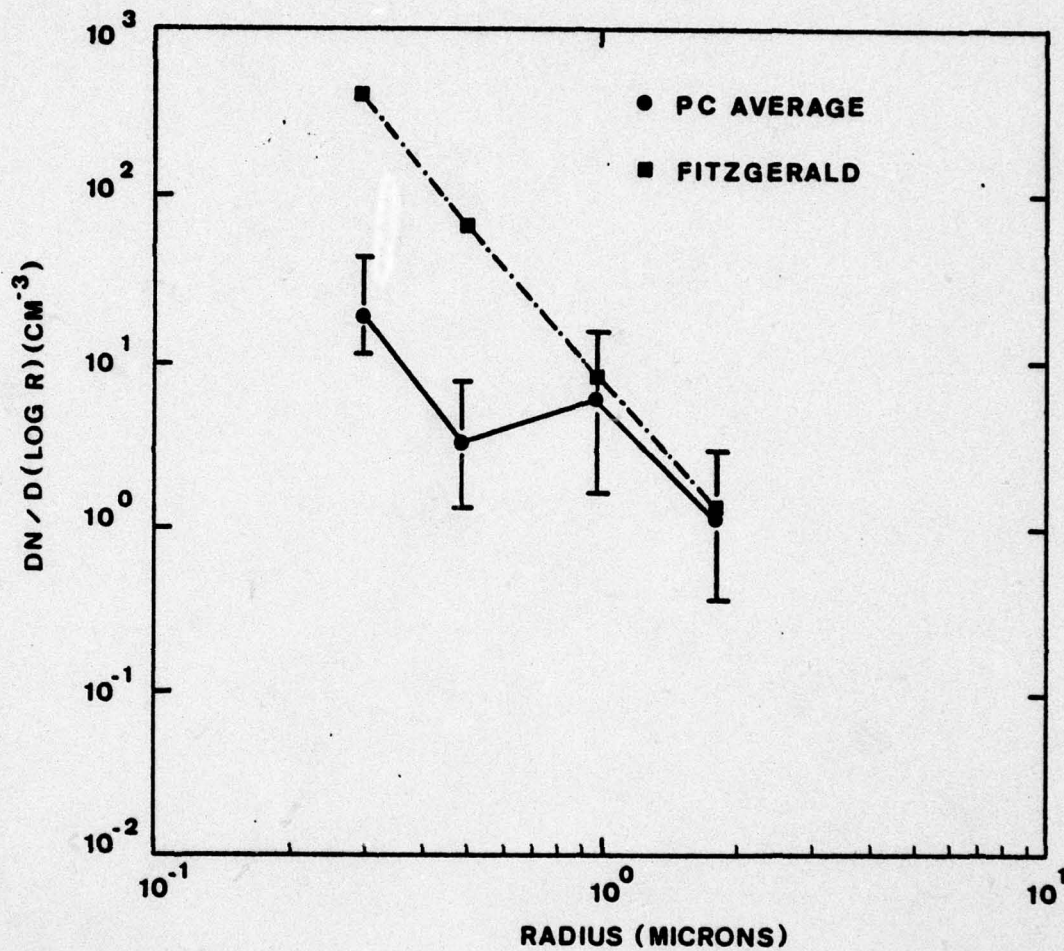


Figure 29. Average Aerosol Size Distribution for the PC Experiment and Distribution Predicted by Fitzgerald's Model

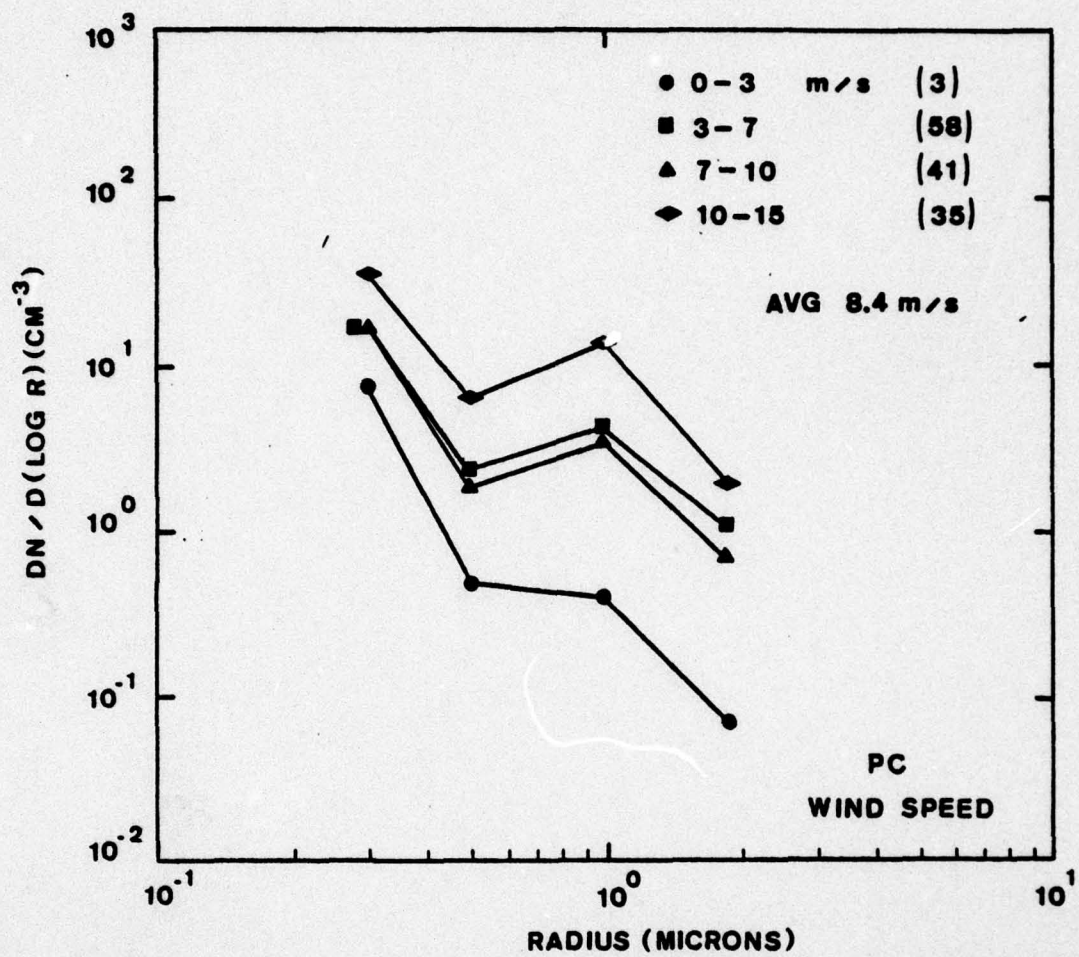


Figure 30. Variation of PC Size Distribution with Wind Speed

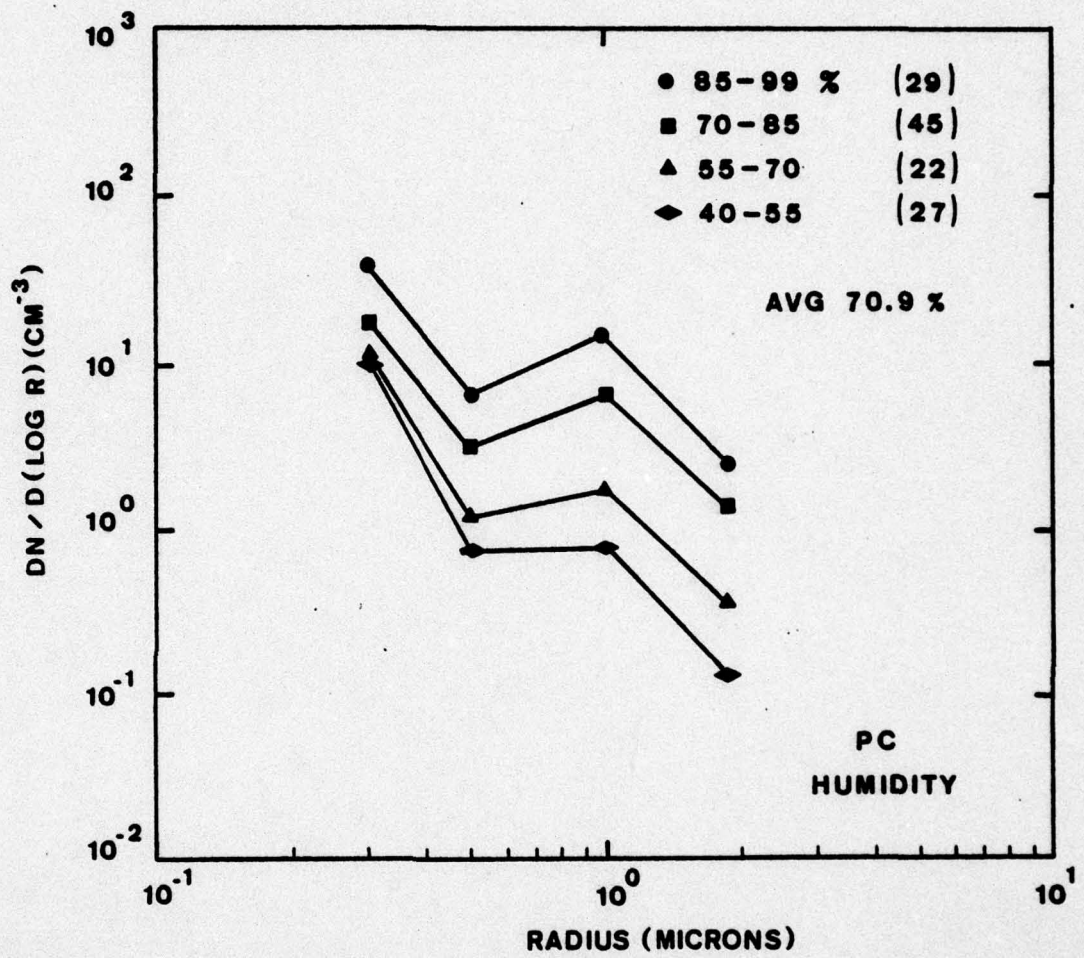


Figure 31. Variation of PC Size Distribution with Relative Humidity

PC

INTERVAL	RH	WIND SPEED
.25 - .35 $\mu$	.620	.559
.35 - .70 $\mu$	.726	.554
.70 - 1.5 $\mu$	.654	.511
1.5 - 2.5 $\mu$	.729	.442

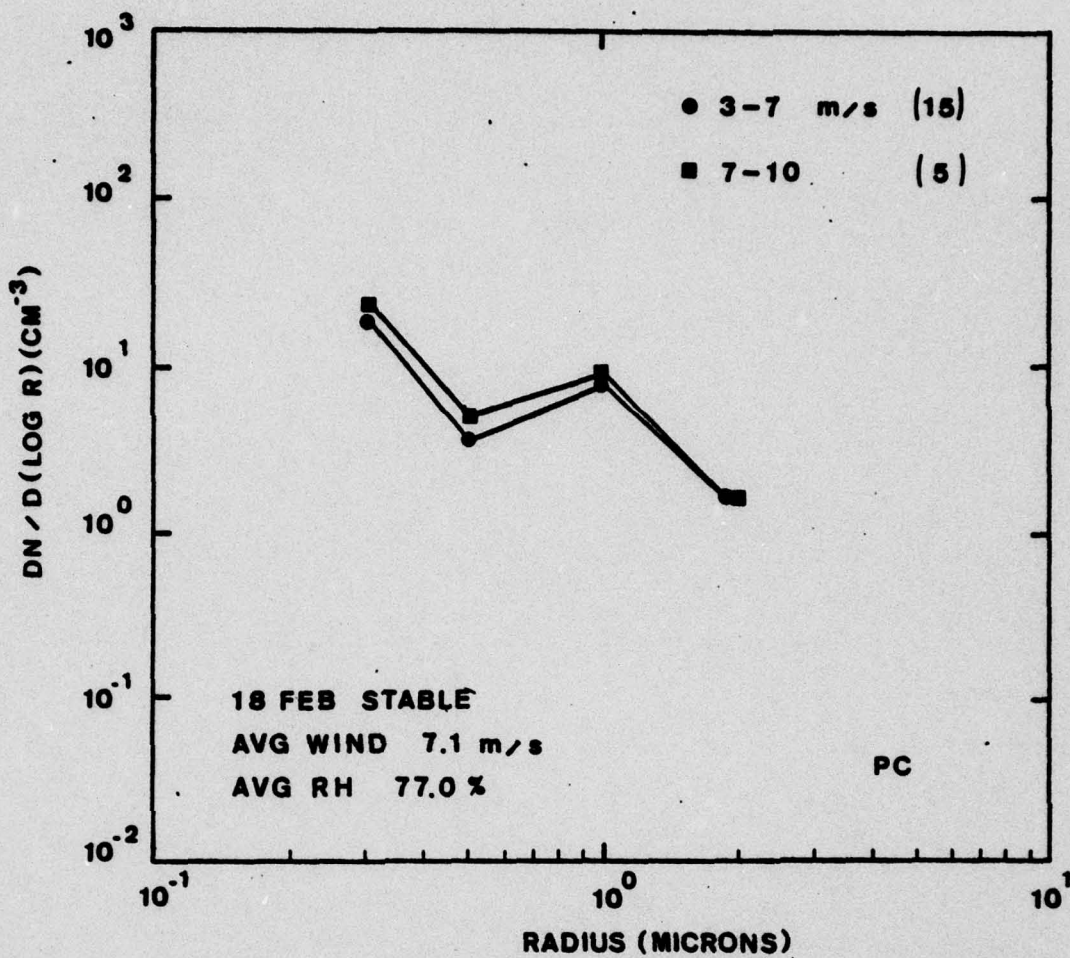
TOTAL

Table IV. Correlation Coefficients for the PC Experiment



concentration to humidity at Panama City is witnessed in the 1.5  $\mu$  to 2.5  $\mu$  interval. This result may indicate that equilibrium tends to exist between production and sedimentation in this interval as hypothesized by Moore and Mason (1954). Disagreement exists in that the steeper negative slope is found during periods of strongest wind. The plot showing the effect of relative humidity on the size distributions results in small variations in the .25  $\mu$  - .35  $\mu$  interval. This probably indicates that the majority of these particles represent a mixture of continental and marine nuclei.

The stability influence was investigated by comparing observations on 18 February and 21 February. Temperature measurements at various levels on the platform made it possible to examine the lapse rates and determine the stability. The average distributions for these days grouped according to wind speed, and respective correlation coefficients are shown in Figures 32 and 33. The low humidities on 21 February resulted from the earlier intrusion of continental air, but southeast to southwest flow persisted most of the day. Although this trajectory helped to advect in warmer air, production of sea-salt dropped off as the wind decreased considerably below 7 m/sec. A much larger decrease is observed in the distribution curve on 21 February as compared to 18 February when the wind speed decreased below 7 m/sec. This agrees well with Moore's (1952) finding that the change in opacity is well marked during periods of low humidity. Also the decrease in the slope of the curve between .5  $\mu$  and

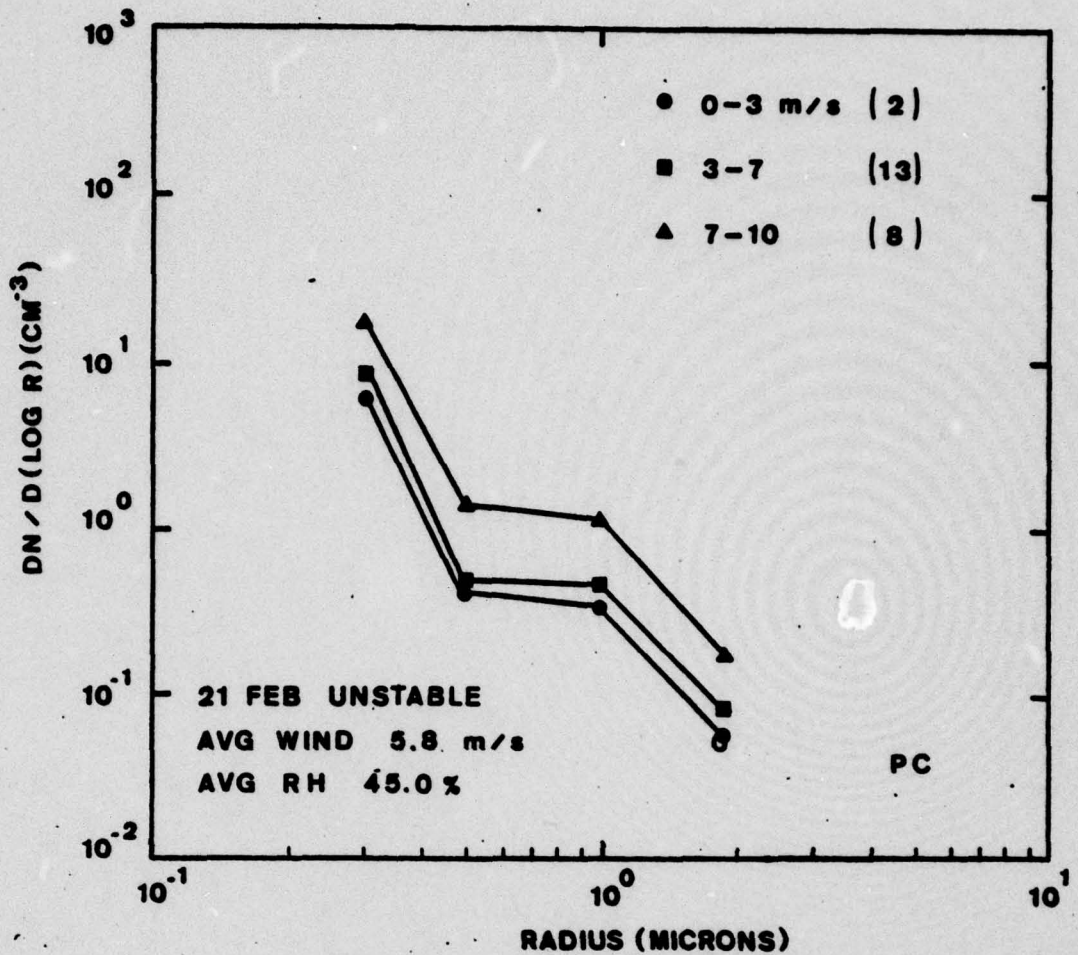


PC

INTERVAL	RH	WIND SPEED
.25 - .35 $\mu$	.169	.065
.35 - .70 $\mu$	.104	.334
.70 - 1.5 $\mu$	.279	.605
1.5 - 2.5 $\mu$	.385	.442

18 FEB STABLE

Figure 32. Correlation Coefficients and Variation of the Size Distribution with Wind Speed, 18 February



PC

INTERVAL	RH	WIND SPEED
.25 - .35 $\mu$	.746	.615
.35 - .70 $\mu$	.692	.651
.70 - 1.5 $\mu$	.558	.737
1.5 - 2.5 $\mu$	.253	.571

21 FEB UNSTABLE

Figure 33. Correlation Coefficients and Variation of the Size Distribution with Wind Speed, 21 February

1  $\mu$  appears to be a function of decreased relative humidity. This again reinforces the premise that sea-salt nuclei predominate in the size range above approximately .7  $\mu$ .

The correlation between concentration and wind speed for all size ranges is greatest on the unstable day. As with the SC data, this is consistent with momentum and diffusion theory. Again the small correlation with relative humidity exhibited by the larger nuclei is probable caused by growth and sedimentation in the absence of significant sea-salt production. When generation was occurring on 18 February, the large particles exhibited the largest correlations with humidity and wind speed. This stable stratification evidently was produced by a previous frontal passage and northerly flow of cold air and accompanying continental particulates. Therefore, a large portion of the aerosol at the beginning of the experiment may have been composed of non-hygroscopic material.

Figures 34 and 35 display the average diurnal changes in wind speed, relative humidity, and aerosol concentration. Again positive correlations are noted as relative humidity and wind speed, although containing quite a bit of scatter, tend to vary accordingly. Of most significance would be the obviously high concentration of droplets in the .7  $\mu$  - 1.5  $\mu$  range. Noting that the average wind seldom went below 7 m/sec, this would indicate that sea-salt nuclei production is greatest in this size range. A diurnal representation of the average aerosol distribution is presented in Figure 36.

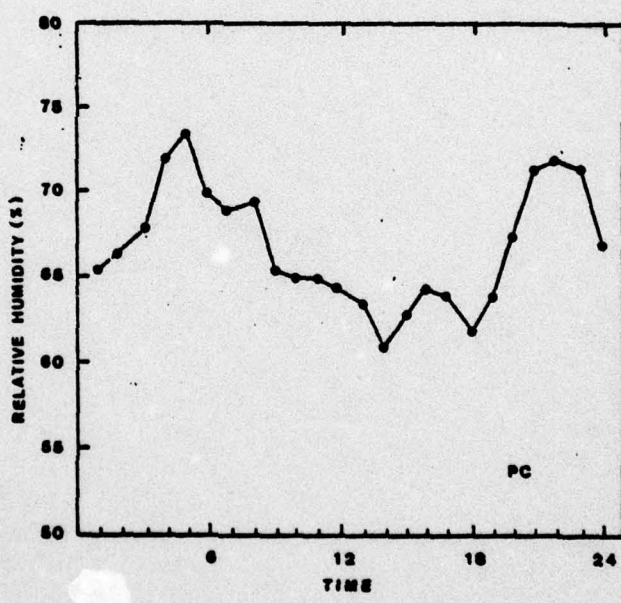
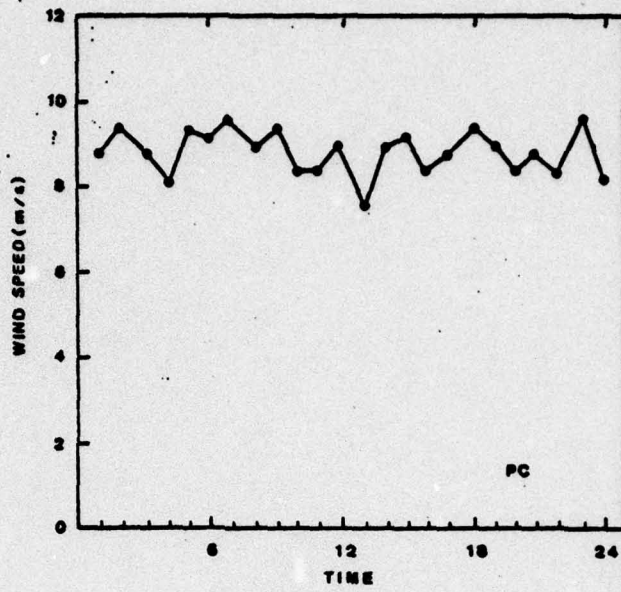


Figure 34. Average PC Diurnal Variations of Wind Speed and Relative Humidity

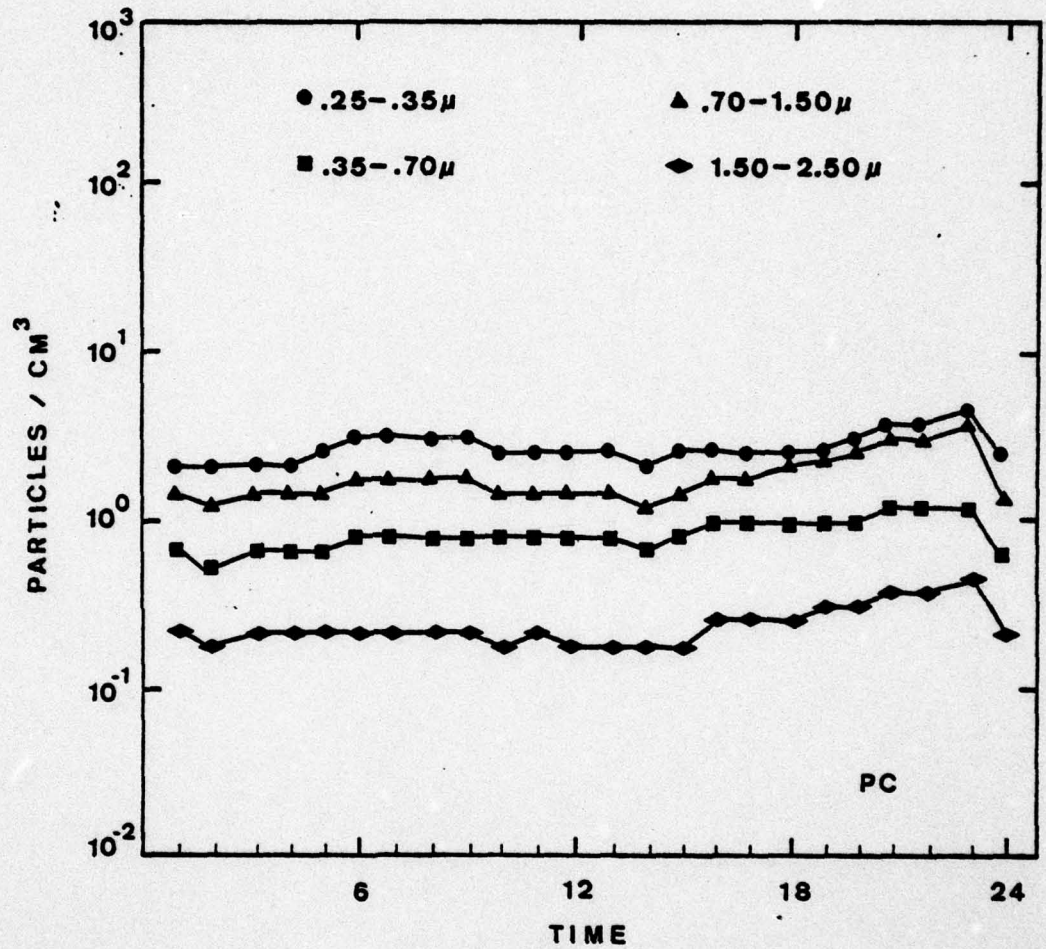


Figure 35. Average PC Diurnal Variations of Particle Concentrations

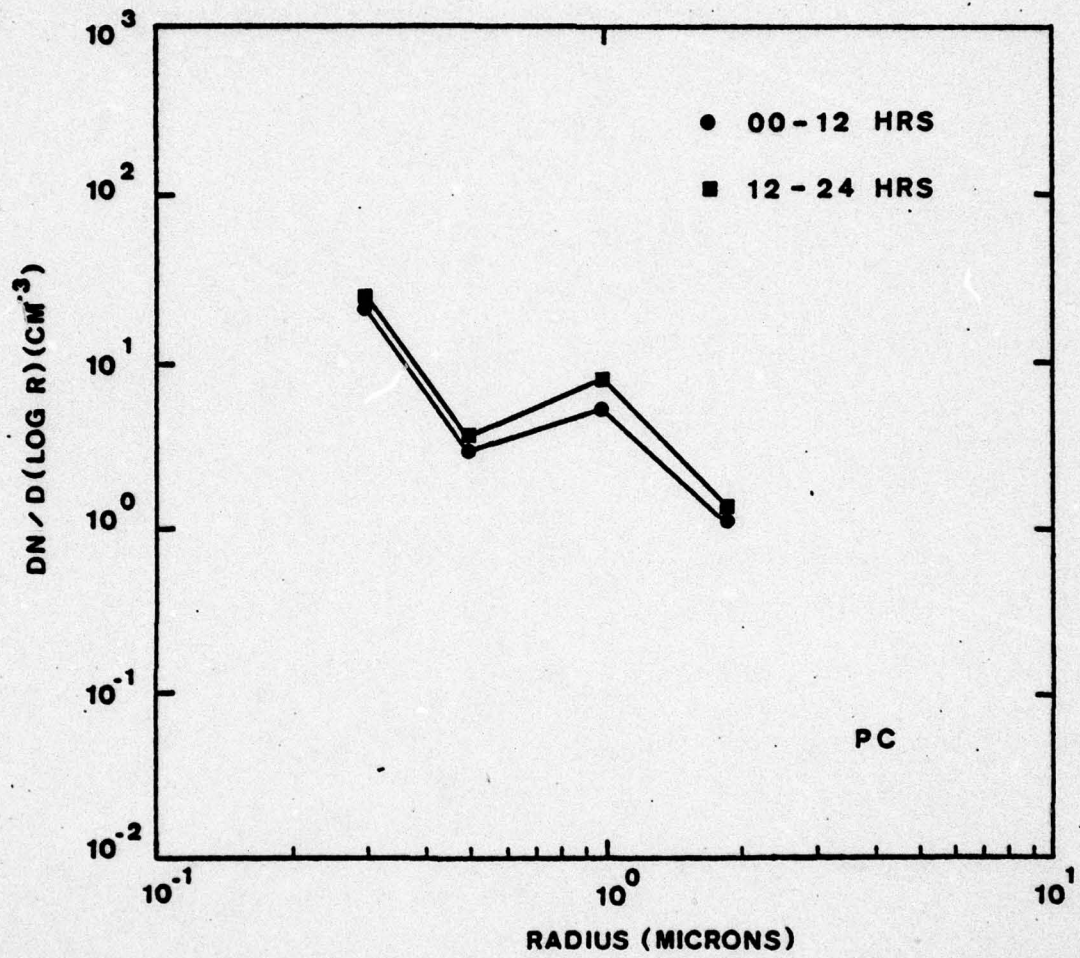


Figure 36. Average PC Diurnal Variation of the Aerosol Size Distribution

Any transport of aerosols due to a land-sea breeze effect should be ruled out as a satisfactory relationship does not seem to exist.



## VII. CONCLUSIONS

The coastal marine aerosol is shown to be a highly variable function of the interaction between synoptic and meso-scale processes. Important meteorological parameters such as wind speed, relative humidity, and stability are dependent upon secondary circulations between land and sea in the absence of large scale forcing.

The minimum concentration in the size distribution curves at  $.4 \mu - .5 \mu$  radius may indicate that this size is indeed the transition zone between the two bubble bursting sea-salt producing mechanisms. Since the slope on either side of this zone is steeper during the Panama City experiment, wind speeds of greater than 7 m/sec result in the generation of sea-salt particles larger than  $.25 \mu$  radius. Sedimentation of particles larger than  $1.5 \mu$  appears to be most significant during periods of low wind speed. During strong winds a state of equilibrium between sedimentation and production exists for these larger particles.

Relative humidity variations have the largest effect on the aerosol size distribution in the absence of sea-salt production. The concentration of the coastal marine aerosol is most sensitive to wind speed effects at low relative humidity. Friction velocity seems to be a better indication of the aerosol size distribution than wind speed under unstable atmospheric conditions. Also during light wind

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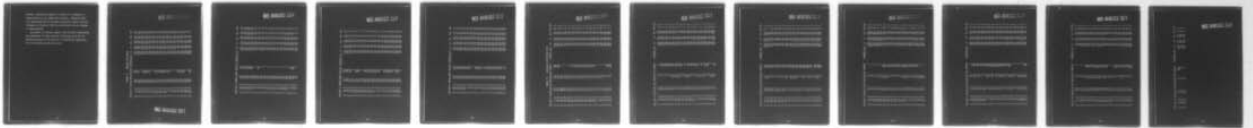
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A COMPARATIVE STUDY OF THE COASTAL MARINE AEROSOL.(U)  
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periods, instability appears to result in a decrease in concentration at the observation height. Enhanced diffusion during periods of sea-salt production causes vertical transport of sea-salt from the sea surface and an increase in concentration.

Any effect of surface organic film possibly suppressing the production of small sea-salt particles could not be examined because of the absence of significant generation off the Southern California coast.

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Table V, Panama City Data

DATE & TIME	PH	WIND DIR	SPEED	AITKEN CON	# PARTICLES / 2.0 L	> 0.5	> 0.7	> 1.4	> 3.0	> 5.0
16	C	77	28015	3000	21075	11206	7935	1287	200	
16	100	68	28016	2300	21791	12419	8553	1324	216	
16	200	72	27517	2000	18261	10588	7682	1147	179	
16	300	72	28014	0	22634	13357	9792	1598	308	
16	400	72	28514	0	23305	13514	10351	1644	265	
16	500	83	29015	1500	23266	13930	10350	1588	270	
16	600	82	28014	1600	24627	14956	11146	1701	294	
16	700	76	28516	1800	25487	15546	11471	1570	228	
16	800	78	28515	1800	24018	14461	10607	1450	220	
16	900	73	28017	1600	23508	14357	10342	1491	225	
16	1000	70	27515	0	22752	13351	5447	1317	190	
16	1100	72	25015	0	21002	11973	6285	1161	174	
16	1200	73	26514	1700	15682	10473	7014	1067	144	
16	1300	72	26016	1600	18450	8917	5471	879	109	
16	1400	74	25012	2100	17104	8210	5055	722	76	
16	1600	64	28011	1900	15002	10300	6521	1088	127	
16	1700	78	28514	2000	18102	10248	7048	1117	158	
16	2100	66	27014	1800	23647	13933	10385	1685	250	
16	2200	84	27514	1700	23440	13872	10274	1655	299	
16	2300	64	27015	1700	22644	13509	5625	1447	240	
15	0	81	26516	1600	22648	13432	5744	1480	257	
15	100	62	28018	1500	20363	11832	8320	1236	209	
15	200	63	28017	0	18567	10948	7658	1163	180	
15	300	64	28015	0	20337	11916	8551	1314	203	
15	400	64	28013	0	20082	11917	8735	1486	220	
15	500	85	28013	0	17003	5872	7170	1245	212	
15	600	66	28012	1600	15437	8631	6196	1047	174	
15	700	66	27512	1500	15227	8438	6035	1014	194	
15	800	87	27013	1500	13565	7784	5624	1002	190	
15	900	87	26512	3200	12605	6862	4832	804	126	
15	1000	86	26512	1350	12267	6484	4655	781	133	
15	1100	85	26510	0	12084	6339	4314	716	92	
15	1200	63	25011	0	10293	5183	3402	567	84	
15	1300	62	26511	1300	10085	4854	3057	485	65	
15	1400	77	25512	1500	11280	5547	3588	592	73	

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DATE & TIME	RF	WIND DIR	SPEED	AITKEN CON	# PARTICLES / 2.0 L	>.5	>.7	>1.4	>3.0	>5.0
15	1500	77	27516	1400		11553	6074	3956	621	75
15	1600	74	25015	1500		13707	7262	4518	818	119
15	1700	61	24014	1600		13113	6551	4752	822	115
15	1800	83	24014	1500		14767	8355	6040	1100	191
15	1900	66	24514	2400		15356	8762	6422	1134	156
15	2000	65	25014	3900		16216	8251	5921	1060	184
15	2100	81	27015	5000		22721	5815	6807	1206	199
15	2200	66	28512	5200		27510	12120	8448	1448	240
20	2000	65	35028	0		9030	2810	1640	201	37
20	3000	66	35025	0		8603	2198	1196	160	23
20	4000	62	35523	0		6051	1542	822	82	11
20	5000	56	35021	8000		8187	1558	1067	139	19
20	6000	51	35319	0		5283	2094	1125	142	21
20	7000	51	35319	0		10768	2421	1153	138	16
20	8000	51	35016	0		8003	1788	876	93	19
20	9000	58	33516	0		7142	1227	686	59	12
20	1000	45	34017	0		6425	1375	681	57	8
20	1100	31	32019	0		5225	1248	754	125	36
20	1200	35	31521	0		4356	572	566	77	25
20	1300	35	31015	0		3606	1027	637	135	35
20	1400	30	31516	0		3155	671	536	85	18
20	1500	31	30020	0		2932	680	416	60	23
20	1600	32	30517	0		2671	712	457	74	16
20	1700	26	30515	0		4768	934	534	68	17
20	1800	24	32013	0		4334	1005	565	78	14
20	1900	25	35012	0		3243	598	684	167	54
20	2000	33	32010	0		3482	853	530	92	22
20	2100	45	25017	0		7045	2528	1586	260	51
20	2200	65	28518	0		6583	2076	1238	169	23
20	2300	45	33026	0		5442	2659	1636	254	52
21	0000	44	33021	0		5284	2823	1667	267	74
21	1000	42	32516	2700		7582	2674	1691	304	79
21	2000	35	35016	2700		5155	1323	724	64	11
21	3000	44	318	0		5417	715	278	5	1
21	4000	45	1518	0		8835	1702	821	49	7

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DATE & TIME	RP	WIND DIR & SPEED	AITKEN CON	0 PARTICLES / 2.0 L	>.5	>.7	>1.4	>3.0	>5.0
21	500	57	5017	2400	16227	3181	1482	135	29
21	600	51	6516	2200	14456	3144	1565	171	42
21	700	51	8516	2200	15791	3463	1508	104	11
21	600	52	10511	3300	10451	2086	963	57	0
21	500	47	135 5	5000	6777	1305	688	64	9
21	1600	42	130 8	0	5754	1212	632	67	12
21	1100	45	135 4	0	4743	508	511	56	6
21	1200	42	175 6	8000	3550	510	454	62	10
21	1300	37	215 7	9000	2872	567	300	35	7
21	1400	40	265 8	7000	2650	546	324	40	4
21	1500	40	240 8	0	2465	502	266	33	2
21	1600	35	240 5	0	2643	597	310	35	7
21	1700	36	24510	4500	2585	678	371	41	11
21	1800	43	24010	4000	3040	702	402	55	7
21	1500	43	22510	2900	3282	808	451	64	10
21	2600	45	23511	3100	3133	780	435	65	10
21	2100	47	22510	2800	3124	757	467	76	8
21	2300	47	21510	3000	3321	888	520	71	10
21	2300	48	20010	3000	3356	800	492	73	11
22	100	55	200 5	2800	3404	856	485	71	11
22	200	55	18011	2700	3260	903	520	75	11
22	300	55	16511	2200	2502	846	510	89	6
22	400	55	14513	0	3556	1218	764	112	7
22	500	64	14014	0	3735	1478	971	135	18
22	600	71	13517	2000	5366	2287	1517	256	32
22	600	66	14015	2000	6182	2747	1508	278	42
22	700	66	14015	2000	6688	3031	2045	341	61
22	800	66	14021	2000	7297	3365	2262	311	54
22	900	61	13822	1800	7255	3345	2223	298	30
22	1000	62	13522	2100	6784	3125	2050	278	28
22	1100	62	13522	0	8511	4341	3010	402	61
22	1200	65	13522	2300	8645	3906	2565	336	34
22	1300	66	13522	2500	9368	4203	2663	375	45
22	1400	66	13822	3600	8305	3451	2225	333	34
22	1500	66	13322	0	8253	3500	2207	335	38

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DATE & TIME	RH	WIND DIR & SPEED	AIKEN COM	# PARTICLES / 2.0 L	> 0.5	> 0.7	21-4	23-0	> 5.0
22	1600	65	14518	2400	8153	3511	2283	326	34
22	1700	71	13025	3200	9455	4276	2846	443	49
22	1800	70	12026	3000	9013	3528	2631	332	44
22	1900	70	11926	3000	10166	4562	2972	421	56
22	2000	76	14021	2500	10704	4593	3058	450	60
22	2100	73	15022	3300	8087	3483	2251	273	21
22	2200	73	13817	2000	8353	4419	3067	508	82
22	2300	82	12815	1700	10158	5447	3882	678	111
23	0000	85	12018	2000	11035	6411	4566	678	87
23	0100	80	11821	1700	12617	7048	5002	686	86
23	0200	82	12922	2000	14337	8001	5818	835	129
23	0300	84	13018	0	17227	5785	7192	1125	154
23	0400	86	11518	0	18069	10235	7555	1076	191
23	0500	88	11027	3500	18055	5643	7041	901	116
23	0600	82	11527	2800	24655	14118	5970	1180	158
23	0700	82	11028	3200	27584	15405	10982	1297	179
23	0800	83	11024	3100	34453	18582	13313	1685	209
23	0900	88	11028	3300	34011	18410	13031	1563	170
23	1000	86	11525	0	31144	16629	11121	1223	102
23	1100	86	11027	2700	36246	15615	13208	1332	107
23	1200	86	11521	2600	36875	20600	13926	1504	132
23	1300	85	13026	2400	37036	21113	14036	1696	188
23	1400	51	12525	2400	37751	21301	14553	1712	145
23	1500	91	11528	1800	41245	22854	15325	1493	134
23	1600	52	12825	1700	48270	28627	20153	2165	234
23	1700	51	11526	1600	48641	25137	20620	2117	208
23	1800	91	12528	1800	52317	31873	23165	2510	275
23	1900	51	14025	1700	60162	37378	28007	3312	505
23	2000	54	13325	1500	68055	42388	32516	4167	690
23	2100	54	14026	1300	70670	44713	34658	3822	518
23	2200	54	14528	0	75264	46227	36507	3174	281
23	2300	54	15528	1200	55524	56409	45316	5071	516

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Table VI. Southern California Data

DATE & TIME	RA	REL WIND DIR & SPD	SHIPS HEAD & SPD	AIRKEN CON	# PARTICLES / .20 L	>.3	>.6	>1.2	>3.0	>5.0
15	C	85	287 3	270 0	3000	12344	2531	933	10	0
15	1C0	52	290 1	270 0	1500	14239	2526	896	10	0
15	20C	87	26C 1	270 0	1700	15632	2575	1039	5	0
15	30C	85	192 1	270 0	3500	12690	2660	995	11	0
15	4CC	85	18C10	180 9	0	12914	2475	797	10	0
15	5CC	85	18C10	183 9	0	14444	2485	841	6	0
15	60C	80	34C 7	140 9	0	6461	1455	578	7	0
15	7CC	82	15 8	140 9	0	6028	1371	550	10	0
15	800	81	340 9	205 9	0	5612	1673	624	9	0
15	160C	77	215 1	353 2	0	9261	1505	545	6	0
15	11CC	76	172 9	175 9	3800	6753	1305	479	6	0
15	1123	76	172 9	175 9	0	7046	1425	565	11	0
15	113C	76	172 9	175 9	8500	6785	1321	525	10	0
15	114C	76	172 2	155 2	2200	6530	1290	495	12	0
15	1150	76	172 2	155 2	3000	6374	1321	489	10	0
15	12CC	74	168 3	168 2	1300	6215	1357	579	14	0
15	1300	75	168 3	168 2	1200	6173	1291	503	12	0
15	14CC	78	148 3	168 2	1500	6215	1304	458	5	0
15	143C	78	234 7	254 0	2900	6760	1466	583	22	0
15	160C	78	305 6	317 0	2100	8352	1618	585	7	0
15	17CC	8C	284 2	274 0	1300	6632	134C	504	12	0
15	1800	82	190 1	133 9	1200	7095	1491	575	14	0
15	20CC	66	23C 8	167 9	1100	8226	1658	610	10	0
15	21CC	8C	175 5	147 9	1800	11798	2088	728	5	0
15	220C	83	166 1	147 0	1400	11950	2139	723	16	0
15	23CC	8C	315 C	147 0	0	1162C	2C33	689	12	0
2C	C	82	183 1	147 0	0	5226	1780	643	7	0
2C	1CC	82	67 2	147 0	0	5584	1747	637	11	0
2C	200	84	290 2	147 0	0	11695	1512	662	10	0
2C	300	82	90 4	147 0	0	15871	4111	1473	8	0
2C	4CC	84	28C 5	0 0	0	2300C	4668	1636	8	0
2C	5CC	84	280 4	0 0	0	22881	4185	1431	5	0
20	555	84	27C 3	0 0	5100	23114	4664	1628	8	0
2C	61C	84	270 3	0 0	8000	24115	4634	1784	5	0
20	625	84	270 3	0 0	4800	25023	5273	1922	8	0



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DATE & TIME	RM	REL WIND DIR & SPD	SNIPS HEAD & SPD	AITKEN CON	# PARTICLES / .20 L	>.3	>.6	>1.2	>3.0	>5.0
20	65C 84	270 3	0 0	0	23370	4875	1707	6	0	
20	70C 82	285 5	0 0	2403	23550	5020	1761	8	0	
20	71C 82	285 5	0 0	1950	22370	4650	1654	8	0	
20	72C 82	305 2	0 0	2700	21453	4445	1564	3	0	
20	74C 82	310 1	0 0	6000	22713	4723	1725	7	0	
20	75C 82	240 1	0 0	2400	21528	4536	1665	5	0	
20	66C 82	240 1	0 0	2200	22292	4675	1690	7	0	
20	62C 82	230 3	0 0	3000	24133	4141	1525	10	0	
20	635 82	315 5	345 9	3900	22244	4640	1637	9	0	
20	565 85	340 7	342 9	2900	22825	4953	1835	12	0	
20	100C 82	310 8	350 9	1650	23006	5055	1865	14	0	
20	110C 74	20 4	343 9	5500	13184	2310	793	7	0	
20	120C 76	325 3	340 9	3600	12645	2228	733	5	0	
20	130C 76	310 6	332 9	2500	14505	2066	629	6	0	
20	140C 80	280 2	0 0	1600	15433	1917	550	3	0	
20	130C 76	255 3	0 0	0	15117	3566	1392	7	0	
20	170C 81	275 10	350 0	9200	20615	4238	1448	8	0	
20	1615 81	275 10	350 0	0	15965	4032	1371	1	0	
20	150C 64	270 14	272 4	5800	21875	4718	1653	4	0	
20	2005 87	275 10	270 5	7000	26296	5768	2040	10	0	
20	265C 85	280 5	0 0	5100	25681	6755	2436	7	0	
20	220C 85	280 2	0 0	1900	27050	6880	2614	16	0	
20	230C 88	300 3	0 0	3100	67724	20710	8710	16	0	
21	53	330 7	360 9	2900	32513	6312	3119	14	0	
21	100 94	305 2	0 0	4100	45855	12409	4761	5	0	
21	200 55	225 2	265 9	0	51710	14480	5750	5	0	
21	300 97	135 4	0 0	4400	55620	17660	7350	3	0	
21	345 53	100 5	100 2	0	46050	12860	5264	5	0	
21	500 50	125 2	0 0	5000	57565	16320	6641	19	0	
21	545 50	160 2	0 0	5000	30950	6370	3548	17	0	
21	645 87	170 1	0 0	2600	25423	5767	2250	16	0	
21	800 87	170 1	0 0	0	34360	7496	2685	10	0	
21	500 65	30 6	360 7	0	37857	7156	2730	15	0	
21	1000 67	20 5	360 9	0	87734	25760	10098	15	0	
21	1115 50	270 5	0 0	8200	98190	32021	13122	15	0	

# BEST AVAILABLE COPY

DATE & TIME		PH	REL WIND DIR & SPD	SHIPS HEAD & SPD	AITKEN COM	# PARTICLES / .20 L	>.3	>.6	>1.2	>3.0	>5.0
21	1215	5C	26C 0	0 0	7600	107112	34552	15930	11	0	
21	1300	5C	26C16	265 5	9000	11054C	35236	16729	9	0	
21	1400	5C	27416	254 6	7300	116775	37650	16250	1	0	
21	1445	67	28C16	0 0	21000	65554	16527	6258	5	0	
21	1600	85	26C12	0 0	4200	31756	7017	2509	13	0	
21	1700	8C	27C 5	0 0	4300	11364	3261	1257	14	0	
21	1800	85	265 9	0 0	9500	6360	2645	1179	17	0	
21	1925	62	250 6	0 0	6400	5863	2685	1230	30	0	
21	2010	65	225 7	0 0	5000	6255	2865	1307	21	0	
21	2115	84	255 1	0 0	5300	6557	3069	1502	28	0	
21	2200	8E	8C 9	53 9	4700	6587	2543	1464	25	0	
21	2300	5C	100 1	0 0	4800	8985	3256	1491	21	0	
21	2315	54	100 1	0 0	4800	8736	3106	1416	16	0	
21	2325	54	15 8	360 9	5200	11636	3648	1570	24	0	
22	3C 53		10 1	0 0	5600	10065	3569	1474	21	0	
22	10C 51		2C 1	0 0	5200	5586	3214	1418	18	0	
22	13C 51		145 9	165 9	9000	12750	3885	1703	35	0	
22	14C 51		145 5	165 9	18500	34751	7800	2560	18	0	
22	15C 5C		145 9	160 9	14500	58358	13055	4711	11	0	
22	20C 50		145 5	16C 9	19000	72729	17835	6583	17	0	
22	215 5C		18C 1	0 0	7300	50505	11661	4030	16	0	
22	23C 52		180 1	0 0	10000	61164	14380	5195	16	0	
22	245 52		55 6	0 0	14000	60500	16452	5500	14	0	
22	30C 54		55 6	25 9	11000	6586C	1540C	5500	14	0	
22	315 54		65 8	25 9	28000	72405	17855	6525	12	0	
22	33C 53		122 3	0 0	27000	7080C	1438C	5523	16	0	
22	345 53		122 3	0 0	27000	57280	12460	4100	10	0	
22	355 52		300 5	267 9	19000	42050	8785	2637	27	0	
22	415 52		15 1	0 0	16000	30585	7654	2250	19	0	
22	445 52		15 1	0 0	0	73900	18800	7060	14	0	
22	515 55		12C 3	0 0	21000	74348	15837	7487	18	0	
22	60C 55		13C 5	138 9	5500	64051	16194	6190	15	1	
22	645 55		1C 7	36C 9	9000	78135	20548	7623	10	0	
22	715 55		20 1	0 0	16000	100065	25740	11800	11	0	
22	80C 53		20 1	0 0	9600	81079	21612	7799	12	0	

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DATE & TIME	RF	REL WIND DIR & SPD	SHIPS HEAD & SPD	AITKEN COM	# PARTICLES / .28 L	>.3	>.6	>1.2	>3.0	>5.0
22	635	53	20 1	0 0	0	104527	30760	11584	15	0
22	650	93	20 1	0 0	0	104660	31147	12317	15	0
22	510	53	20 1	0 0	0	120771	36450	14250	20	0
22	545	53	20 1	0 0	0	95714	22441	8413	6	0
22	1015	53	20 1	0 0	0	96659	27513	10500	10	0
22	1040	53	20 1	0 0	0	100800	26583	10641	8	0
22	1130	53	20 1	0 0	0	104548	31062	12386	9	0
22	1200	52	178 2	52 0	0	94874	24122	5211	3	0
22	1300	52	45 3	104 9	0	90750	25340	9591	8	0
22	1400	85	100 5	155 9	6400	114549	33457	12966	9	0
22	1500	84	225 7	155 9	79000	37164	6476	2055	36	0
22	1600	85	155 2	153 9	6200	14650	3333	1137	5	0
22	1700	84	215 6	179 9	3100	26285	5464	1546	9	0
22	1755	65	275 5	0 0	10500	17685	3778	1328	5	0
22	1810	65	260 3	0 0	4000	5248	2549	1102	9	0
22	1900	51	140 3	0 0	3800	10716	2870	1127	13	0
22	1915	91	150 3	0 0	8100	11925	3114	1242	20	0
22	1930	51	150 3	0 0	6000	12038	3151	1266	23	0
22	1945	51	150 3	0 0	9300	12081	2564	1257	20	0
22	1950	50	145 6	0 0	14000	12468	3145	1220	15	0
22	2020	50	255 2	0 0	4400	15438	3562	1453	12	0
22	2115	52	120 9	100 9	4300	14500	3471	1295	14	0
22	2200	57	11515	100 9	10000	14504	3245	1251	10	0
22	2300	85	13512	100 9	11500	11553	2657	1093	7	0
23	0000	87	13512	100 9	5100	11140	2751	1065	8	0
23	1000	67	13010	95 9	4200	13020	2710	518	6	0
23	2000	52	135 8	90 9	3800	33570	6770	2222	8	0
23	2355	52	160 2	0 0	5800	63264	14772	5139	10	0
23	3155	52	330 3	325 9	2900	63550	14540	5240	17	0
23	4000	52	340 4	324 9	4700	46733	10112	3350	5	0
23	5300	52	340 4	324 9	3600	38071	7235	2265	3	0
23	6000	52	325 3	300 5	0	25975	5441	1575	0	0
23	7000	52	310 2	300 9	0	35975	6897	2170	5	0
23	8000	85	315 3	290 9	0	32127	5870	1860	5	0
23	9000	80	345 2	315 9	0	22542	3620	1070	5	0

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DATE & TIME	RA	REL WIND DIR & SPD	SHIPS HEAD & SPD	AITKEN CON	# PARTICLES / .20 L	>.3	>.6	>1.2	>3.0	>5.0
23	1600	330 4	300 9	0	25258	4480	1255	5	0	0
23	1400	250 5	0 0	7300	18968	4242	1570	17	0	0
23	1500	215 8	215 5	2700	13656	3353	1327	21	0	0
23	1400	300 7	360 9	5100	16043	3675	1411	27	0	0
23	1600	265 8	0 0	8000	40053	9837	3524	17	0	0
23	1600	245 9	262 9	6400	76557	18645	6462	12	0	0
23	2000	250 1	0 0	3000	18165	4097	1571	18	0	0
23	2100	300 1	0 0	2800	20170	4600	1805	16	0	0
23	2200	330 1	0 0	3500	31142	4760	1852	21	0	0
23	2300	200 1	0 0	4100	22894	5540	2212	23	0	0
24	8 53	5 3	160 4	6000	26665	6232	2352	20	0	0
24	30 53	305 6	318 4	13800	34045	8019	3021	22	0	0
24	100 50	270 1	0 0	15500	75830	15535	7371	28	0	0
24	130 50	135 6	100 5	18000	66626	16280	6169	20	0	0
24	150 85	180 2	0 0	17500	31474	2770	2160	23	0	0
24	230 85	150 2	0 0	19500	36695	6711	2382	23	0	0
24	300 86	225 7	245 9	16000	35310	6150	2032	11	0	0
24	330 86	131 2	0 0	17500	45800	8344	2665	23	0	0
24	400 86	230 6	245 9	15500	45516	5750	3090	11	0	0
24	430 86	240 9	246 9	10000	34785	7334	2560	20	0	0
24	500 87	60 5	50 9	5000	43400	5248	3200	24	0	0
24	530 87	25 4	38 9	6900	46142	5792	3456	19	0	0
24	550 89	310 2	0 0	6200	62980	14200	5069	19	0	0
24	630 85	170 4	147 9	12000	55820	13200	4600	15	0	0
24	700 84	160 8	167 9	14000	45738	5260	3000	15	0	0
24	600 80	140 8	150 9	0	54537	16550	3285	14	0	0
24	630 80	140 8	150 9	0	37951	6447	2111	10	0	0
24	500 76	75 2	0 0	0	30717	5212	1866	10	0	0
24	530 76	75 2	0 0	0	32532	5742	1557	10	0	0
24	1000 77	320 8	325 9	0	41046	6516	2050	15	0	0
24	1030 77	320 8	325 9	0	63024	11604	3446	15	0	0
24	1100 77	320 11	330 9	0	75771	16170	5225	15	0	0
24	1200 71	135 1	100 9	0	92097	24256	10170	30	0	0
25	2100 83	270 11	245 1	3800	10824	2634	712	7	0	0
25	2230 83	285 6	305 1	3600	8284	3261	1161	18	0	0

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DATE & TIME	PH	REL WIND DIR & SPD	SHIPS HEAD & SPD	AIKEN CON	# PARTICLES / .20 L	>.3	>.6	>1.2	>3.0	>5.0
25	23CC	E3	28C10	0	0	9450	343C	1374	20	0
26	C	83	290 6	0	0	9317	3784	1531	40	0
26	1CC	87	45 1	0	0	9486	4147	1890	50	0
26	2CC	86	225 1	0	0	13584	5317	2390	40	0
26	3CC	50	240 1	0	0	10170	4661	2259	40	0
26	4CC	5C	36C 9	0	0	5960	460C	2300	40	0
26	5CC	86	300 2	0	10000	11768	5540	2663	51	0
26	6CC	51	270 6	5600	5600	11050	5090	2650	47	0
26	7CC	66	275 5	3000	3000	1C505	4825	2394	55	1
26	8CC	86	315 9	4500	4500	9229	4167	2088	26	0
26	5CC	66	3CC1C	6000	6000	14334	5851	268C	36	0
26	1CC	85	30012	7300	7300	14455	6C19	2791	36	0
26	11CC	81	30C12	12200	12200	11144	4776	2256	24	0
26	14CC	62	25017	10500	10500	1C552	4365	1575	19	0
26	15CC	81	29016	9600	9600	10857	4545	2055	12	0
26	14CC	77	27C17	5600	5600	10953	4605	2077	15	0
26	15CC	77	24514	5400	5400	12162	528C	2464	21	0
26	16CC	71	23512	4400	4400	11940	5086	2280	15	0
26	17CC	52	29021	20000	20000	7640	2442	1205	26	0
26	15CC	60	3C514	7800	7800	5229	1517	850	55	0
26	21CC	51	32512	0	0	5391	132C	520	25	0
26	22CC	62	325 8	0	0	4700	1446	573	25	0
26	23CC	62	325 4	0	0	3186	1067	452	15	0
27	C	73	150 3	7300	7300	5065	2654	1097	26	0
27	104	55	13C 4	10000	10000	7280	3094	1506	45	0
27	115	55	115 3	8400	8400	6134	2455	1215	23	0
27	145	54	200 2	7600	7600	6321	2358	1047	24	0
27	2CC	53	210 1	7000	7000	7654	3C97	1384	23	0
27	212	52	235 8	5602	5602	5576	385C	1682	24	0
27	222	62	24C 1	6800	6800	10586	4300	1849	25	0
27	25C	82	24C 1	7100	7100	1C122	4147	182C	25	0
27	3C6	84	315 1	6700	6700	10260	4205	1890	23	0
27	224	65	31C 1	16000	16000	10808	4253	1851	26	0
27	335	50	310 1	52000	52000	10212	3996	1650	15	0
27	356	50	31C 1	23000	23000	10383	4C25	1807	40	0

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DATE & TIME	RM	REL WIND DIR & SPD	SHIPS HEAD & SPD	AITKEN CON	# PARTICLES / .28 L	>.3	>.6	>1.2	>3.0	>5.0
27	417	SC	30 1	35000	10864	4320	1530	35	0	0
27	44C	SC	80 1	23500	12414	4915	2272	34	0	0
27	5CC	65	275 4	0	5864	4105	1508	25	0	0
27	6CC	80	270 2	0	5685	4014	1558	15	0	0
27	7CC	75	255 2	0	12690	4586	2504	25	0	0

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