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FINITE ELEMENT FORMULATION AND SOLUTION OF CONTACT-
IMPACT PROBLEMS IN CONTINUUM MECHANICS

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California University

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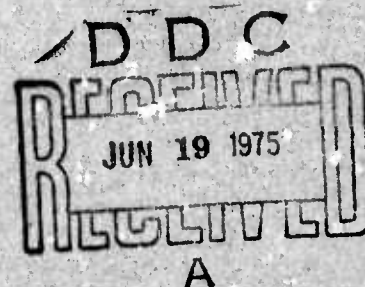
**FINITE ELEMENT FORMULATION AND SOLUTION
OF CONTACT-IMPACT PROBLEMS IN CONTINUUM
MECHANICS**

May 1974

An Investigation Conducted by
**STRUCTURAL ENGINEERING LABORATORY
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Introduction

In this report we consider the general problem of contact and impact between two bodies. The report is divided into three basic parts. These parts describe: (I) The general theory of contact-impact problems, (II) A numerical scheme for the analysis of contact-impact problems, and (III) The description of computer program FEAP 74 for the solution of contact-impact problems. In an appendix we include the program subroutines and general input description for FEAP 74.

In Sections 1 to 6, Part I, we deal with spatial aspects of the theory and in Section 7, Part I, we deal with temporal aspects. This splitting of the theory is motivated by the way we intend to numerically solve the equations, i.e., the finite element method spatially and a finite difference method temporally.

Part II considers a numerical implementation of the theory given in Part I. Section 9 deals with spatial notions of the numerical problem and Section 10 the temporal. The solution scheme for the resulting algebraic problem is discussed in Sections 11 and 12.

The computer program FEAP 74 was modified to incorporate the numerical contact-impact model. The program modifications and capabilities together with two numerical examples are contained in Part III.

Finally, in the appendix we give listings for the contact subroutines together with the data input instructions.

PART I
VARIATIONAL FORMULATION OF CONTACT-IMPACT
PROBLEMS IN CONTINUUM MECHANICS

1. Preliminaries

Our conventions on indices are as follows:

Superscripts indicate to which body an entity pertains. Summation is to take place only when explicitly indicated.

Latin subscripts range over 1,2,3, while Greek subscripts range over 1,2. The summation convention is assumed to hold for both.

A body B is a nice connected region of \mathbb{R}^3 with a piecewise smooth boundary ∂B . A contact^{*} problem is a boundary value problem, or an initial-boundary value problem, in which two bodies, B^1 and B^2 , interact according to the principles of mechanics. Thus the primary kinematic axiom of a contact problem is that configurations \mathcal{B}^1 and \mathcal{B}^2 , of B^1 and B^2 , respectively, do not penetrate each other, i.e.,

$$\begin{aligned} (\mathcal{B}^1)^\circ \cap \mathcal{B}^2 &= \emptyset, \\ \mathcal{B}^1 \cap (\mathcal{B}^2)^\circ &= \emptyset, \end{aligned} \tag{1}$$

where $(\mathcal{B}^\alpha)^\circ$ denotes the interior of \mathcal{B}^α , $\alpha = 1, 2$.

On the other hand the unique condition which characterizes contact problems is that material points on the boundaries of B^1 and B^2 may coalesce during the motion of the bodies. Thus we say B^1 and B^2 are in contact if $\partial \mathcal{B}^1 \cap \partial \mathcal{B}^2 \neq \emptyset$, and we define the contact surface e by

* It is usual for the term contact to have a static connotation while the term impact has a dynamic connotation. We shall use contact in the general sense to include static as well as dynamic phenomena.

$$\mathcal{C} = \partial \mathcal{B}^1 \cap \partial \mathcal{B}^2. \quad (2)$$

If \mathcal{B}^1 and \mathcal{B}^2 are never in contact then $\mathcal{C} = \emptyset$ for all configurations \mathcal{B}^1 and \mathcal{B}^2 , and in this case an initial-boundary value problem for \mathcal{B}^1 and \mathcal{B}^2 reduces to one in which \mathcal{B}^1 and \mathcal{B}^2 may be treated separately. Thus a non-trivial contact problem is one in which $\mathcal{C} \neq \emptyset$ for at least one instant during the motion of \mathcal{B}^1 and \mathcal{B}^2 . The picture (Fig. 1) illustrates these notions.

Equation (1) implies that \mathcal{C} is a material surface with respect to both bodies, i.e., one which is not crossed by material particles. From this we may deduce the interface conditions on \mathcal{C} .

Let \underline{x} be a persistent point of \mathcal{C} (one at which joining or releasing of the bodies is not instantaneously occurring) and \underline{v} be the velocity of \underline{x} ($\underline{v} = \dot{\underline{x}}$). Note that only the normal part of \underline{v} is independent of the parametrization of \mathcal{C} . Let \underline{v}^1 and \underline{v}^2 be the velocities of the material particles located at the points \underline{x}^1 and \underline{x}^2 , contained in $\partial \mathcal{B}^1$ and $\partial \mathcal{B}^2$, respectively, such that $\underline{x} = \underline{x}^1 = \underline{x}^2$ at the present instant. Then since \mathcal{C} is material and \underline{x} is persistent

$$\underline{v} \cdot \underline{n} = \underline{v}^1 \cdot \underline{n} = \underline{v}^2 \cdot \underline{n}, \quad (3)$$

where \underline{n} is a unit normal vector to \mathcal{C} at \underline{x} . From this it follows that a necessary condition for momentum to be balanced at \underline{x} is that

$$(\underline{t}^1 + \underline{t}^2) \cdot \underline{n} = 0, \quad (4)$$

where \underline{t}^{α} is the Cauchy traction vector with respect to $\partial \mathcal{B}^{\alpha}$.

In addition we assume that no tensile tractions can occur on \mathcal{C} ,

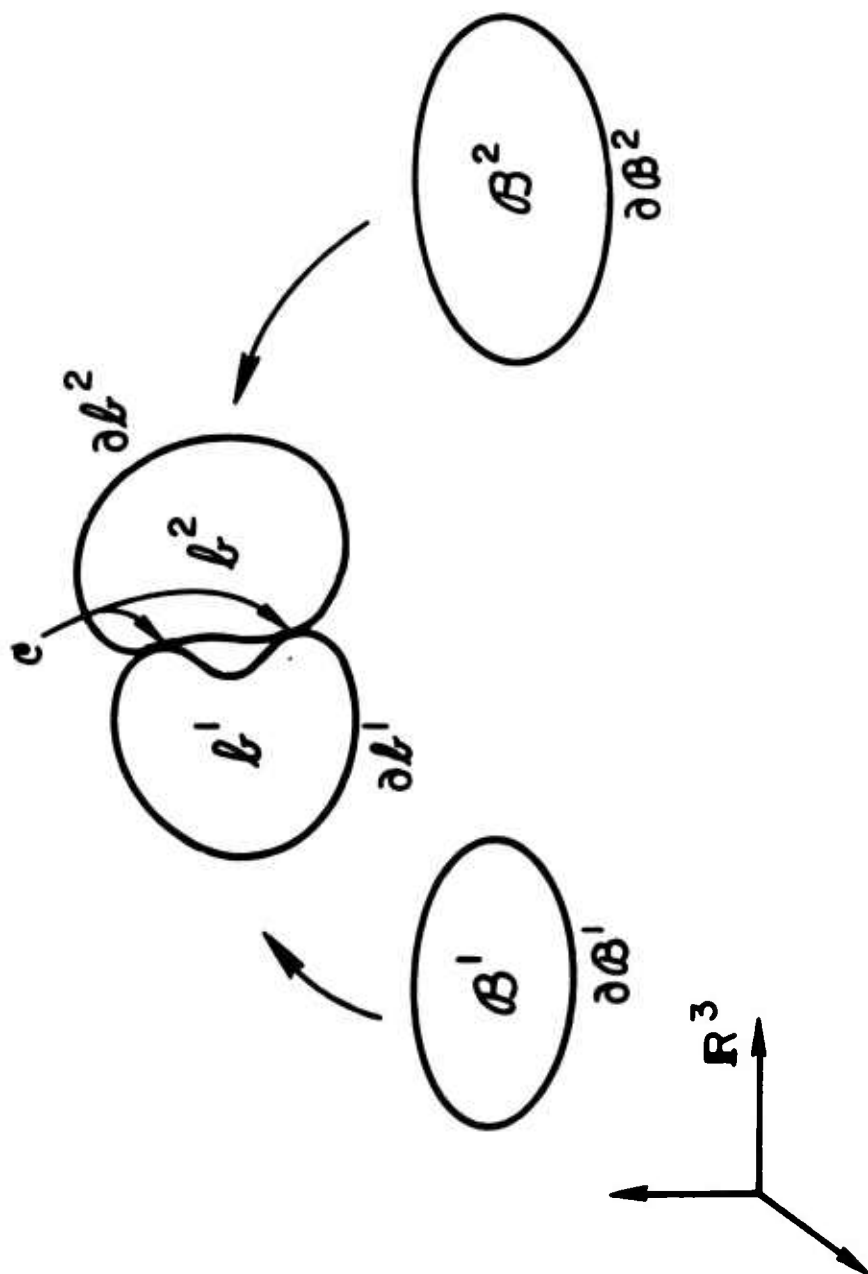


Figure 1

$$\underline{t}^{\alpha} \cdot \underline{n}^{\alpha} \leq 0, \quad (5)$$

where \underline{n}^{α} is the outward unit normal vector to $\partial \mathcal{B}^{\alpha}$. This condition excludes the possibility of the two bodies being glued together. Conditions (1-5) characterize our notion of a contact problem.

Note that thus far we have said nothing about the tangential parts of \underline{v}^{α} and \underline{t}^{α} . These remaining conditions are determined by the frictional nature of the contact. We shall study two simple cases.

Case I: If we assume that points, once in contact, move with \mathcal{C} until released, we have that

$$\underline{v}^1 = \underline{v}^2, \quad (6)$$

and therefore

$$\underline{t}^1 + \underline{t}^2 = \underline{0}. \quad (7)$$

For this model we say that a no-slip, or perfect friction, condition is achieved on \mathcal{C} . Thus condition (5) and equations (6) and (7) are the interface conditions for this case.

Case II: We may create the interface conditions for a frictionless, sliding contact by asserting that the tangential part of each \underline{t}^{α} is identically zero,

$$\underline{t}^{\alpha} - (\underline{t}^{\alpha} \cdot \underline{n}^{\alpha}) \underline{n}^{\alpha} = \underline{0}. \quad (8)$$

Eq. (8), along with (3-5), are the interface conditions for this case.

2. Variational Theorems

We will formulate a variational theorem for the contact problem of finite elastodynamics. We point out, however, that our treatment is entirely general and could be used in conjunction with any field theory, as the only unique feature of the formulation involves the handling of interface conditions. At the same time finite elastodynamics, though lending itself to a clean and simple variational statement, is a case of wide practical interest.

We shall first obtain a variational theorem for the usual initial-boundary value problem of finite elastodynamics by a trivial generalization of some work done by S. Nemat-Nasser [1].

For notational simplicity let \mathcal{C} denote $\partial\mathcal{B}$, and let $d\mathcal{A}$ and $d\mathcal{B}$ denote area and volume forms for \mathcal{B} and \mathcal{C} , respectively. Let $\mathcal{Q}_\tau \subset \mathcal{Q}$ be that part of \mathcal{C} where surface tractions are prescribed, and denote by $\bar{\mathbf{T}}$ the Piola - Kirchhoff traction vector representing these prescribed tractions. Call ρ_0 the density of \mathcal{B} in the initial configuration, $\underline{\mathbf{F}}$ the extrinsic body force vector and let $\underline{\mathbf{x}} = \underline{\mathbf{x}}_0(\underline{\mathbf{X}})$ represent the position at time t of the material particle located at $\underline{\mathbf{X}}$ in the initial configuration. For convenience we take \mathcal{B} to be the initial configuration. We denote by $\partial\underline{\mathbf{x}}/\partial\underline{\mathbf{X}}$ the deformation gradients and by $\Phi(\partial\underline{\mathbf{x}}/\partial\underline{\mathbf{X}})$ the strain energy density. Then if $\underline{\mathbf{x}}$ satisfies the kinematic boundary conditions

$$\underline{\mathbf{x}} = \bar{\underline{\mathbf{x}}} \quad (9)$$

on $\mathcal{Q}_\kappa \subset \mathcal{Q}$, where

$$\begin{aligned} \mathcal{Q}_\kappa \cup \mathcal{Q}_\tau &= \mathcal{Q}, \\ \mathcal{Q}_\kappa \cap \mathcal{Q}_\tau &= \emptyset, \end{aligned}$$

the functional Π defined by

$$\begin{aligned} \Pi(\underline{x}) = \int_0^t \left\{ \int_B \left(\Phi(\partial \underline{x} / \partial \underline{x}) - \rho \cdot \dot{\underline{x}} \cdot \dot{\underline{x}} / 2 \right. \right. \\ \left. \left. - \rho \cdot \underline{F} \cdot \underline{x} \right) d\mathcal{B} - \int_{A_r} \underline{x} \cdot \underline{T} d\mathcal{A} \right\} dt, \end{aligned} \quad (10)$$

is stationary, i.e., its first variation vanishes

$$\begin{aligned} 0 = \delta \Pi(\underline{x}, \delta \underline{x}) = \int_0^t \left\{ \int_B \left(\rho (\dot{\underline{x}} - \underline{F}) - \text{DIV } \underline{P} \right) \cdot \delta \underline{x} d\mathcal{B} \right. \\ \left. + \int_{A_r} (\underline{T} - \bar{\underline{T}}) \cdot \delta \underline{x} d\mathcal{A} \right\} dt, \end{aligned} \quad (11)$$

subject to the constraint on variations $\delta \underline{x}_t = \delta \underline{x}_0 = 0$, (12)

if and only if the equations of motion and traction boundary conditions are satisfied

$$\rho (\dot{\underline{x}} - \underline{F}) = \text{DIV } \underline{P}, \quad \text{in } \mathcal{B}, \quad (13)$$

$$\underline{T} = \bar{\underline{T}}, \quad \text{on } A_r, \quad (14)$$

where $\underline{P} = \partial \Phi / \partial (\partial \underline{x} / \partial \underline{x})$ is the first Piola - Kirchhoff stress tensor,

$\underline{T} = \underline{N} \cdot \underline{P}$ is the Piola - Kirchhoff traction vector, and \underline{N} is the outward unit normal vector to \mathcal{A} . The solution to the initial-boundary value problem must also satisfy the given initial conditions

$$\left. \begin{aligned} \underline{x} &= \underline{x}_0 \\ \dot{\underline{x}} &= \dot{\underline{x}}_0 \end{aligned} \right\} \quad \text{in } \mathcal{B}. \quad (15)$$

To interpret this variational theorem for two (non-interacting) bodies set

$$\mathcal{B} = \mathcal{B}^1 \cup \mathcal{B}^2,$$

$$\mathcal{A} = \mathcal{A}^1 \cup \mathcal{A}^2, \quad \text{etc.},$$

and write

$$\Pi(\underline{x}) = \Pi^1(\underline{x}^1) + \Pi^2(\underline{x}^2).$$

The next step is to add to Π terms manifesting the interface conditions on \mathcal{C} and to stipulate the constraints under which the vanishing of the first variation of the appended functional corresponds to a solution of the contact problem. To do this we must consider further the kinematics and geometry of \mathcal{C} .

Define two piecewise smooth, invertible maps $\underline{x}^1, \underline{x}^2$ by the condition

$$(\underline{x}^\alpha)^{-1}: \mathcal{C} \longrightarrow \mathcal{C}^\alpha \subset \mathcal{Q}^\alpha, \quad (16)$$

where each \underline{x}^α identifies points on the boundary of the initial configuration \mathcal{B}^α which map into the contact surface \mathcal{C} at each instant of time. If $\underline{x} \in \mathcal{C}$, then $\underline{x}^1 = (\underline{x}^1)^{-1}(\underline{x})$ and $\underline{x}^2 = (\underline{x}^2)^{-1}(\underline{x})$ are the positions of particles in \mathcal{Q}^1 and \mathcal{Q}^2 , respectively, which have coalesced at $\underline{x} \in \mathcal{C}$. It is clear what the \underline{x}^α 's really are, viz., if $\underline{x}^\alpha = \underline{x}_t^\alpha(\underline{X}^\alpha)$, for all $\underline{X}^\alpha \in \mathcal{B}^\alpha$, represents the motion of body \mathcal{B}^α from the original configuration \mathcal{B}^α to the present one \mathcal{B}^α , then \underline{x}^α is the restriction of \underline{x}^α to \mathcal{C}^α ,

$$\underline{x}^\alpha(\underline{X}^\alpha) = \underline{x}^\alpha(\underline{X}^\alpha), \quad (17)$$

for each $\underline{X}^\alpha \in \mathcal{C}^\alpha$, $\alpha = 1, 2$. For the time being we consider the \underline{x}^α 's as maps defined independently of the \underline{x}^α 's and consider (17) a constraint on possible motions.

We are interested in to what extent the relation

$$\underline{x} = \underline{x}^1(\underline{X}^1) = \underline{x}^2(\underline{X}^2) , \quad (18)$$

is smooth in time and analogously under what circumstances the variations of the \underline{x}^a 's are equal. In general the \underline{x}^a 's will not even be continuous in time since contact surfaces can be instantaneously created or destroyed. If we eliminate such exceptional instants and consider only persistent points, the bodies still may slide with respect to each other, as depicted in Fig. 2. Thus tangential velocities are seen to be unequal in general. However, when \underline{x} is persistent, the impenetrability condition (1) forces the normal velocity components to be equal, and concomitantly the normal components of variations of the \underline{x}^a 's are also equal

$$\delta \underline{x}^{1 \cdot n} = \delta \underline{x}^{2 \cdot n} \quad (19)$$

For sliding contact (Case II), Eq. (19) characterizes the constraint on variations of the \underline{x}^a 's equivalent to the velocity constraint (3).

For no-slip contact (Case I),

$$\delta \underline{x}^{1 \cdot t} = \delta \underline{x}^{2 \cdot t} , \quad (20)$$

is easily seen to be the condition on variations equivalent to Eq. (6). We shall see that Eqs. (19) and (20) lead to the proper interface conditions in the variational theorems.

Introduce vector valued Lagrange multipliers $\underline{\tau}^a$, and add

$$\mathcal{X} = - \sum_{a=1}^2 \int_0^t \int_{C^a} \underline{\tau}^a \cdot (\underline{x}^a - \underline{x}^a) dC^a dt , \quad (21)$$

to the functional Π (Eq. 10). Note that when $C \neq \emptyset$,

$$C = C_* \cup C_T \cup C^a ,$$

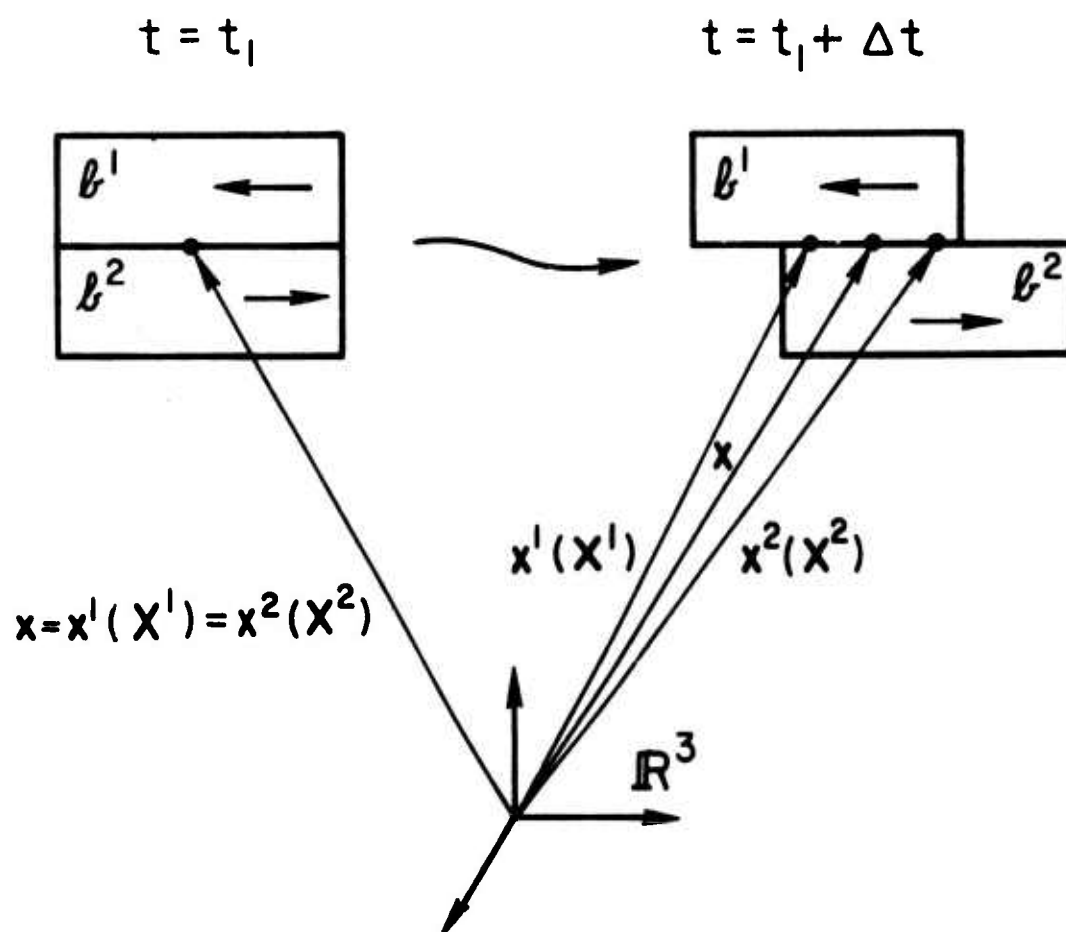


Figure 2

and assume for consistency's sake that

$$\begin{aligned} \mathcal{C}^{\alpha} &\subset \mathcal{Q}_T^{\alpha}, \\ \mathbf{T} &= \mathbf{0} \quad \text{on} \quad \mathcal{C}^{\alpha}. \end{aligned} \quad (22)$$

This condition will preclude the ambiguous circumstance of non-zero tractions being specified on the contact area. Upon taking variations of $\mathcal{J} = \mathbb{I} + \mathcal{K}$ we get Eqs. (13), (14) and,

$$\begin{aligned} 0 = & - \sum_{\alpha=1}^2 \int_0^t \int_{\mathcal{C}^{\alpha}} \left\{ \delta \mathcal{Q}^{\alpha} \cdot (\mathbf{x}^{\alpha} - \mathbf{x}^{\alpha}) + \right. \\ & + \delta \mathbf{x}^{\alpha} \cdot (\mathcal{Q}^{\alpha} - \mathbf{T}^{\alpha}) - \delta \mathbf{x}^{\alpha} \cdot \mathcal{Q}^{\alpha} \left. \right\} d\mathcal{C}^{\alpha} dt + \\ & + \text{transversality condition.} \end{aligned} \quad (23)$$

The transversality condition is the classical terminology for variations associated with the domain \mathcal{C}^{α} .

The first summand of (23) gives us (17) which insures that the \mathbf{x}^{α} 's map into \mathcal{C} properly. The second summand identifies \mathcal{Q}^{α} as the Piola - Kirchhoff traction vector \mathbf{T}^{α} on \mathcal{C}^{α} . Let us investigate the third summand.

Consider first Case I and define

$$\delta \mathbf{x}^{\alpha} = \delta \mathbf{x}^{\alpha}, \quad \alpha = 1, 2, \quad (24)$$

which makes sense because of Eq. (20). This condition is equivalent to insisting

$$\dot{\mathbf{x}}^{\alpha} = \dot{\mathbf{x}}^{\alpha}, \quad \alpha = 1, 2,$$

thus the first summand of (23) also implies (6) holds whenever we have a

persistent point. Let j^{α} denote the Jacobian determinant associated with χ^{α} ,

$$de = j^{\alpha} d\zeta^{\alpha}. \quad (25)$$

Notice then that since ζ^{α} is the Piola - Kirchhoff traction vector, $(1/j^{\alpha}) \zeta^{\alpha}$ is the corresponding Cauchy traction vector. With these we have for the third summand,

$$0 = \sum_{\alpha=1}^2 \int_{\zeta^{\alpha}} \delta \chi^{\alpha} \cdot \zeta^{\alpha} d\zeta^{\alpha} = \int_{\mathfrak{e}} \delta \chi \cdot ((1/j^1) \zeta^1 + (1/j^2) \zeta^2) de, \quad (26)$$

which in words means the Cauchy traction vectors are in equilibrium. Thus the momentum balance, Eq. (7), is satisfied on \mathfrak{e} .

In Case II we only have that (19) holds, so define

$$\delta \chi(\eta) = \delta \chi^{\alpha} \cdot \eta, \quad \alpha = 1, 2. \quad (27)$$

This requirement also insures that,

$$\dot{\chi}^1 \cdot \eta = \dot{\chi}^2 \cdot \eta, \quad$$

thus the first summand of (23) implies (3). For this case the third summand takes the form,

$$\begin{aligned} 0 = \sum_{\alpha=1}^2 \int_{\zeta^{\alpha}} \delta \chi^{\alpha} \cdot \zeta^{\alpha} d\zeta^{\alpha} &= \int_{\mathfrak{e}} \delta \chi(\eta) ((1/j^1) \zeta^1 \cdot \eta + (1/j^2) \zeta^2 \cdot \eta) de \\ &\quad + \sum_{\alpha=1}^2 \int_{\zeta^{\alpha}} (\delta \chi^{\alpha} - \delta \chi(\eta) \eta) \cdot \zeta^{\alpha} d\zeta^{\alpha}. \end{aligned} \quad (28)$$

The integral over \mathfrak{e} gives us Eq. (4). The significance of the second integral hinges on the observation that $(\delta \chi^{\alpha} - \delta \chi(\eta) \eta)$ is a tangent vector to ζ^{α} for each α . Thus the tangential part of each ζ^{α} is identically zero, which is equivalent to the shear free condition, Eq. (8),

which we require for Case II.

A standard calculation enables us to write the transversality condition as,

$$0 = \sum_{\alpha=1}^2 \int_{\partial C^{\alpha}} (\delta X^{\alpha} \cdot T^{\alpha} \cdot \underline{C}^{\alpha} \cdot (\underline{x}^{\alpha} - \underline{x}'^{\alpha})) d(\partial C^{\alpha}) , \quad (29)$$

where the transversal T^{α} is a unit vector field tangent to C^{α} , and perpendicular and pointing outward with respect to ∂C^{α} , Fig. 3. Thus (29) implies that

$$\underline{C}^{\alpha} \cdot (\underline{x}^{\alpha} - \underline{x}'^{\alpha}) = 0 \quad \text{on } \partial C^{\alpha}, \alpha=1,2 \quad (30)$$

Assuming continuity of the integrands of (21) on the closure of C^{α} , condition (30) is already implied by the first summand of (23). This assumption precludes \underline{C}^{α} taking the form of a δ -distribution on ∂C^{α} .

Although this assumption is warranted here it may not be true when one employs certain approximate theories in mechanics. For instance consider the case where a Bernoulli-Euler beam is uniformly loaded and sits on a rigid parabolic surface (Fig. 4). At the contact points a, a' , concentrated reactions must exist to balance shear forces. This example is actually from a completely different class of contact problems in that contact is made along a part of the interior rather than the boundary. Such problems as the contact of plates and shells also fall into this class. We could summarize such situations by the description -- m -dimensional contact of m -dimensional bodies, e.g., for the beam $m=1$, and for plates and shells $m=2$. The case under investigation in this paper ($m=3$) is an example of the $(m-1)$ -dimensional contact of m -dimensional bodies.

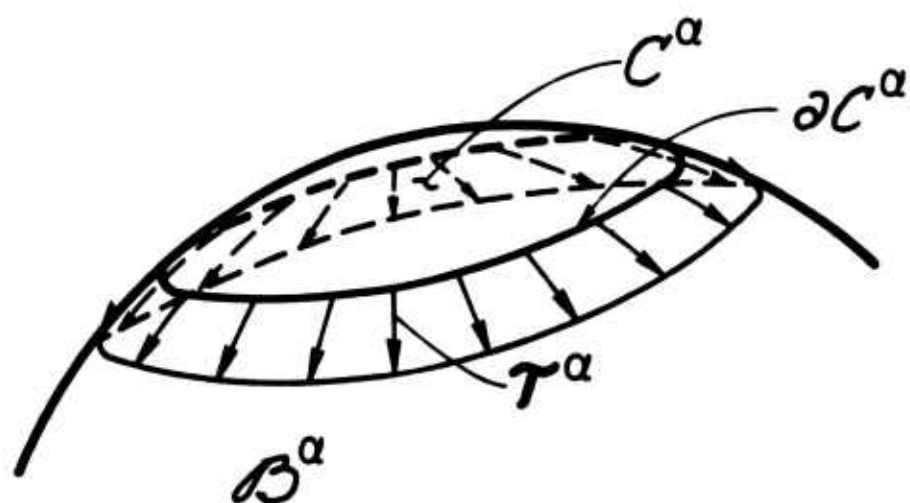


Figure 3

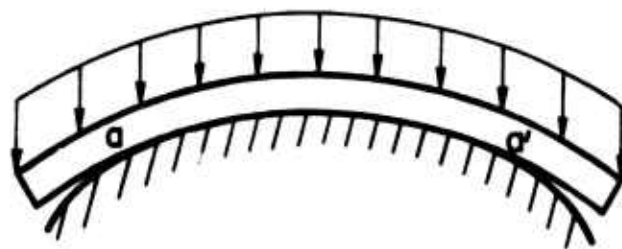


Figure 4

It is good to keep in mind cases such as that illustrated in Fig. 4 when considering specific boundary value problems.

A further point worth mentioning here is that the transversality condition will in general be an independent one in a numerical algorithm. For example, if the fields in the integrand of (21) are approximated by a family of trial functions, Eq. (23) only implies that some weighted integrals over the \tilde{C}^α 's vanish. The condition (29) requires that weighted integrals over the $\delta\tilde{C}^\alpha$'s also vanish.

We now summarize our results in the following theorems:

Theorem I: Let (1), (2), (5), (9), (12), (15), and (20) hold. Then \underline{x} is a solution to the no-slip contact problem (Case I), that is, (6), (7), (13), (14) and (17) also hold if, and only if, $\delta\mathcal{J}=0$ for arbitrary variations of \underline{x}^α , $\underline{\chi}^\alpha$ and \tilde{C}^α , $\alpha=1,2$.

Theorem II: Let (1), (2), (5), (9), (12), (15), and (19) hold. Then \underline{x} is a solution to the sliding contact problem (Case II), that is, (3), (4), (8), (13), (14) and (17) also hold if, and only if, $\delta\mathcal{J}=0$ for arbitrary variations of \underline{x}^α , $\underline{\chi}^\alpha$ and \tilde{C}^α , $\alpha=1,2$.

3. Consideration of Theorems I and II as Computational Tools

Theorems I and II may be employed to generate numerical algorithms for the solution of contact problems. The basic idea is to represent \mathbf{u}^* , \mathbf{v}^* and \mathbf{c}^* as the product of known functions on \mathbb{R}^3 with unknown parameters depending on time. Then Theorems I and II provide us with a method for generating an approximate system of equations (e.g., by the classical Ritz-Galerkin technique) in terms of these unknown parameters, which then can be solved incrementally and/or iteratively, subject to the side conditions of the theorems. The constraints (1) and (5) will both take the form of inequalities in actual computations, thus the ideas of optimization theory will probably be useful in the actual construction of a numerical algorithm.

The finite element method is a powerful technique for obtaining a system of approximate equations, and it is of interest to find out how amenable are Theorems I and II to a finite element formulation. Unfortunately the term χ would result in a terrible mess if the integrand was represented by typical finite element functions. This is because the boundaries of the \mathcal{C}^* 's are unknown and thus a parametric integration would bury the defining parameters of the \mathcal{C}^* 's in the arguments of Heaviside functions representing the supports of the elements. Note that a classical Ritz-Galerkin approximation would not be subject to this pitfall, since the associated trial functions could be chosen to be real analytic and thus easily integrated parametrically to a relatively simple form. However, such a formulation is restricted to a geometrically simpler class of problems. Thus it is desirable to seek a generalization that will lend itself cleanly to a finite element formulation.

4. Variational Theorems Without Transversality Conditions

Let \tilde{C}^n be a fixed part of \mathcal{A}_T^n such that

$$\tilde{C}^n \supset C^n, \quad (31)$$

and

$$\bar{I} = 0 \quad \text{on} \quad \tilde{C}^n \sim C^n. \quad (32)$$

Define a scalar valued function η^n on \tilde{C}^n such that

$$\eta^n(x^n) = 0 \quad \text{if} \quad x^n \in \tilde{C}^n \sim C^n. \quad (33)$$

Let $\tilde{e} \supset e$, and define the maps χ^n by the condition

$$(\chi^n)^{-1}: \tilde{e} \rightarrow \tilde{C}^n,$$

where, as before, χ^n represents x^n on C^n ; but on $\tilde{C}^n \sim C^n$ we place no physical interpretation on χ^n . Thus on \tilde{C}^n we will always have that,

$$\eta^n(x^n - \chi^n) = 0, \quad (34)$$

since $x^n = \chi^n$ on C^n and $\eta^n = 0$ on the relative complement $\tilde{C}^n \sim C^n$.

Introduce vector valued Lagrange multipliers σ^n and let $\mathcal{L} = \mathbb{I} + \mathcal{M}$ where

$$\mathcal{M} = - \sum_{i=1}^k \int_0^1 \left\{ \sigma^n \cdot \eta^n(x^n - \chi^n) \right\} d\tilde{C}^n dt. \quad (35)$$

We require that the variations of χ^n satisfy the same conditions as before, but now for all \tilde{C}^n :

$$\left. \begin{array}{l} \text{Case I:} \quad \delta \chi^n \equiv \delta \chi^n \\ \text{Case II:} \quad \delta \chi^n(n) \equiv \delta \chi^n \cdot n \end{array} \right\} \quad \text{on} \quad \tilde{C}^n \quad (36)$$

where \underline{n} is a unit normal vector to \tilde{E} . Computing the first variation of \mathcal{L} we have the usual conditions emanating from II and

$$\begin{aligned}
 0 = & - \sum_{\alpha=1}^2 \int_0^t \int_{\tilde{E}^-} \left\{ \delta \underline{\sigma}^\alpha (\underline{\eta}^\alpha (\underline{x}^\alpha - \underline{x}'^\alpha)) \right. \\
 & + \delta \underline{\eta}^\alpha (\underline{\sigma}^\alpha (\underline{x}^\alpha - \underline{x}'^\alpha)) \\
 & + \delta \underline{x}^\alpha \cdot (\underline{\eta}^\alpha \underline{\sigma}^\alpha - \underline{I}^\alpha) \\
 & \left. - \delta \underline{x}'^\alpha \cdot (\underline{\eta}^\alpha \underline{\sigma}^\alpha) \right\} d\tilde{E}^\alpha dt.
 \end{aligned} \tag{37}$$

The first summand gives us (34) and we define

$$\tilde{C}^\alpha = \{ \underline{x}^\alpha \in \tilde{E}^\alpha : \underline{x}^\alpha(\underline{x}^\alpha) = \underline{x}'^\alpha(\underline{x}^\alpha) \}. \tag{38}$$

The third summand defines $\underline{\eta}^\alpha \underline{\sigma}^\alpha$ as the Piola - Kirchhoff traction vector. Note that this insures that $\underline{I}^\alpha = 0$ on $\tilde{E}^\alpha \sim \tilde{C}^\alpha$ since $\underline{\eta}^\alpha = 0$ there. The fourth summand gives us the appropriate Cauchy traction condition across E for each case of (36). The second summand is identically satisfied on \tilde{C}^α since $\underline{x}^\alpha = \underline{x}'^\alpha$. On $\tilde{E}^\alpha \sim \tilde{C}^\alpha$ it tells us that $\underline{\sigma}^\alpha$ is orthogonol to $\underline{x}^\alpha - \underline{x}'^\alpha$, but this is of no physical interest.

Thus we can state the following theorems:

Theorem I': Let (1), (2), (5), (9), (12), (15), (31), (32) and (36)₁ hold. Then \underline{x} is a solution to the no-slip contact problem (Case I), that is, (6), (7), (13), (14) and (17) also hold where \tilde{C}^α is defined by (38), if $\delta \mathcal{L} = 0$ for arbitrary variations of \underline{x}^α , \underline{x}'^α , $\underline{\eta}^\alpha$ and $\underline{\sigma}^\alpha$, $\alpha = 1, 2$.

Theorem II': Let (1), (2), (5), (9), (12), (15), (31), (32) and (36)₂ hold. Then \underline{x} is a solution to the sliding contact problem (Case II), that is, (3), (4), (8), (13), (14) and (17) hold where \tilde{C}^α is defined by (38), if $\delta \mathcal{L} = 0$ for arbitrary variations of \underline{x}^α , \underline{x}'^α , $\underline{\eta}^\alpha$ and $\underline{\sigma}^\alpha$, $\alpha = 1, 2$.

The important feature of these theorems is that the regions \tilde{E}^{\pm} are fixed. Thus transversality conditions are absent, and the theorems may be applied to finite element formulations. In fact one would naturally take \tilde{E}^{\pm} to be a union of elements in \mathcal{Q}^{\pm} , large enough to contain \tilde{C}^{\pm} throughout the motion.

Thus far our considerations have been quite general and, in fact, more general than would be required for the solution of particular classes of contact problems. In the next section we illustrate the many simplifications which can be made in the application of the preceding theorems to a class of problems of wide practical interest.

5. Hertzian Contact Problems

We wish to characterize contact problems in which the contact surface is approximately planar and the bodies have undergone small deformations in the neighborhood of the contact surface.

Assume the following:

- (1) $\underline{n} \approx n_i \underline{e}_i \approx \underline{e}_3$ on \mathcal{C} , where the n_i indicate components with respect to the standard basis $\{\underline{e}_i\}_1^3$ for \mathbb{R}^3 ,
(see Fig. 5).

- (2) $j \approx 1$, $\alpha = 1, 2$, thus $\underline{t} \approx \underline{T}$ on \mathcal{C} .

Assumptions (1) and (2) together imply that,

$$\underline{t}_3 \approx \underline{t} \cdot \underline{n} \approx \underline{T} \cdot \underline{n} \approx T_3,$$

and that,

$$(\underline{t}_1, \underline{t}_2, 0) \approx \underline{t} - (\underline{t} \cdot \underline{n})\underline{n} \approx \underline{T} - (\underline{T} \cdot \underline{n})\underline{n} \approx (T_1, T_2, 0).$$

- (3) Material points which eventually contact have, to the first order, the same initial coordinates \underline{z}_1 and \underline{z}_2 . Explicitly we manifest this idea by requiring that the χ^α 's satisfy

$$\chi^1(\underline{z}_1, \underline{z}_2, X_3^1(\underline{z}_1, \underline{z}_2)) = \chi^2(\underline{z}_1, \underline{z}_2, X_3^2(\underline{z}_1, \underline{z}_2)). \quad (39)$$

This is depicted in Fig. 6. Since X_3^α are given functions which define the surfaces \mathcal{C}^α , it follows from (39) that,

$$\delta \chi^1 = \delta \chi^2.$$

We term problems for which these assumptions hold Hertzian, since these assumptions are implicit in Hertz' classical theory [2] (see

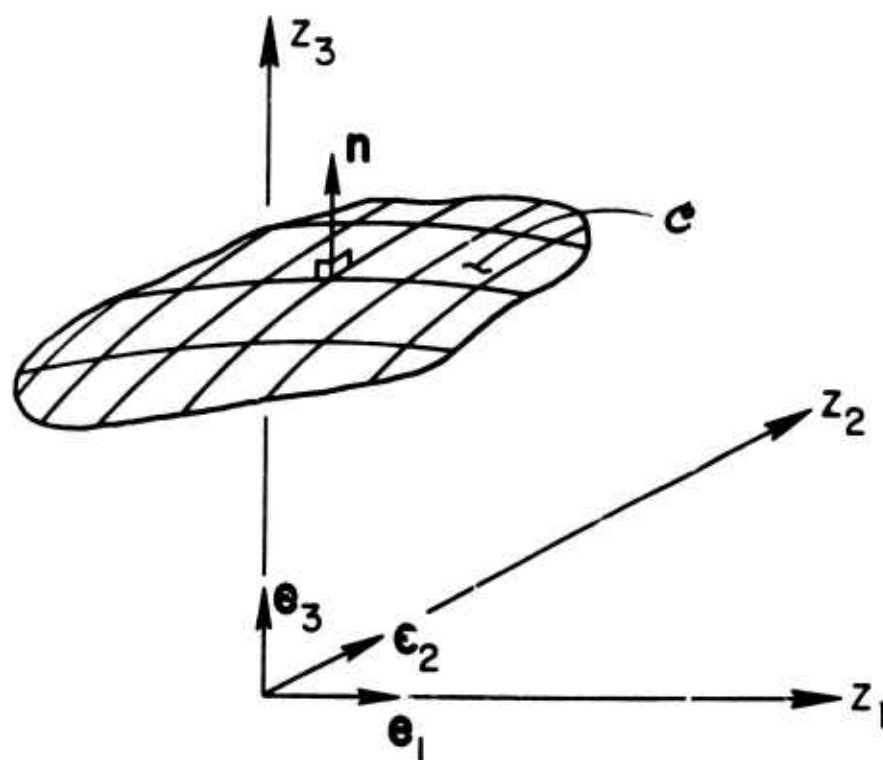


Figure 5

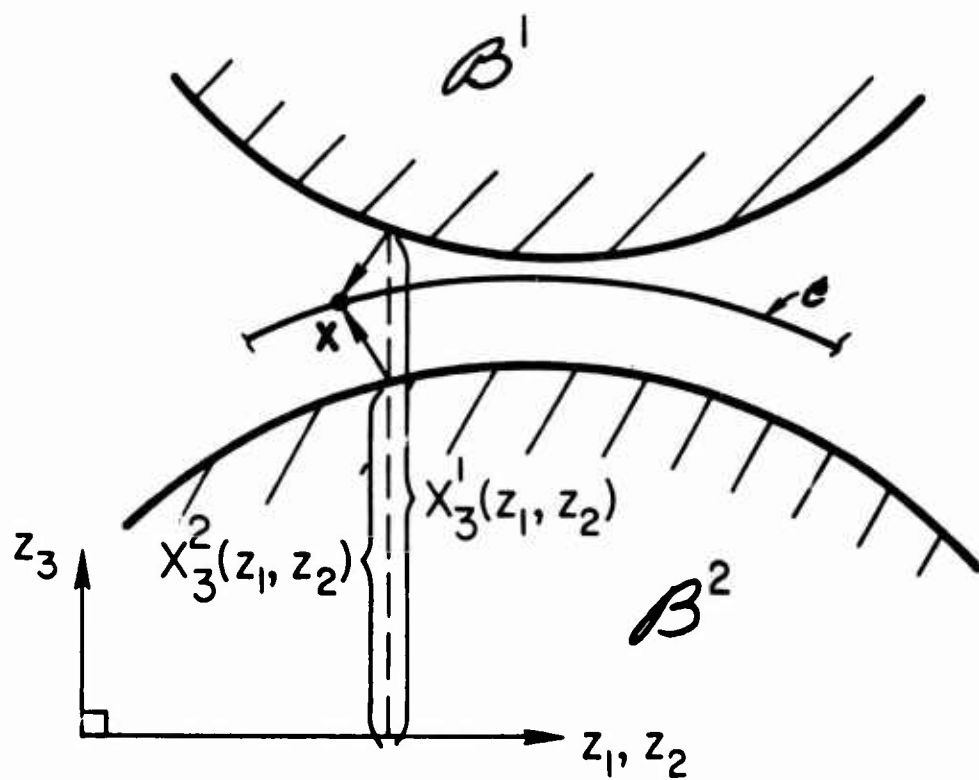


Figure 6

Goldsmith [3] for an excellent exposition of this work and also many applications of Hertz' theory to impact problems). It should be pointed out that the formulation we are about to give is still considerably more general than those to which Hertz' theory applies.

We now show how these assumptions allow us to make simplifications in the preceding theorems.

Theorems I and II:

Due to assumption (3) the term \mathcal{K} can be replaced by an integral over a region in the z_1, z_2 -plane. This region, say c , is the projection of \mathcal{C} onto the z_1, z_2 -plane, and due to assumption (2) it coincides, to the first order, with the projections of the \mathcal{C}^α 's. Thus \mathcal{K} can be written

$$\mathcal{K} = - \sum_{\alpha=1}^2 \int_0^t \int_c \underline{\tau}^\alpha \cdot (\underline{x}^\alpha - \underline{x}^\alpha) \, dc \, dt . \quad (40)$$

Since, for Case I, we know that the momentum balance on \mathcal{C} requires that

$$\underline{\tau}^1 + \underline{\tau}^2 = \underline{0} ,$$

we may make use of this relation immediately. Thus define

$$\underline{\tau} = \underline{\tau}^1 = -\underline{\tau}^2 ,$$

and substitute into (40). Employing (39), the integrand simplifies to

$$\underline{\tau} \cdot (\underline{x}^2 - \underline{x}^1) . \quad (41)$$

The analog of (23) becomes

$$\begin{aligned}
0 = & \int_0^t \int_C \left\{ \delta \tau \cdot (\underline{x}^2 - \underline{x}^1) + \right. \\
& \left. + \delta \underline{x}^1 \cdot (\underline{T}^1 - \underline{\tau}) + \delta \underline{x}^2 \cdot (\underline{T}^2 + \underline{\tau}) \right\} dc dt \\
& + \text{transversality condition.}
\end{aligned} \tag{42}$$

Thus the same conclusions of Theorem I can be drawn. However, from a numerical standpoint things are considerably different. First of all, since the \underline{x}^* 's are absent in this formulation, we do not get a uniquely defined $\underline{\tau}$; \underline{x}^1 and \underline{x}^2 will not in general be the same pointwise. If the graph of $\underline{\tau}$ is important it could be constructed by averaging \underline{x}^1 and \underline{x}^2 , which, if the solution is any good, should be reasonably close pointwise. On the other hand, the \underline{x}^* 's being absent engenders a considerable saving in the number of equations to be solved and in their complexity.

The analogous case for Theorem II is constructed simply by setting

$$\tau_1 = \tau_2 = 0, \quad \tau \stackrel{\text{def.}}{=} \tau_3.$$

Then the integrand of \mathcal{K} becomes

$$\tau (\underline{x}_3^2 - \underline{x}_3^1) \tag{43}$$

and (42) reduces to

$$\begin{aligned}
0 = & \int_0^t \int_C \left\{ \delta \tau (\underline{x}_3^2 - \underline{x}_3^1) + \right. \\
& + \delta \underline{x}_3^1 (\underline{T}_3^1 - \tau) + \delta \underline{x}_3^2 (\underline{T}_3^2 + \tau) \\
& + \delta \underline{x}_a^1 \underline{T}_a^1 + \delta \underline{x}_a^2 \underline{T}_a^2 \left. \right\} dc dt \\
& + \text{transversality condition.}
\end{aligned} \tag{44}$$

Hence the conclusions of Theorem II hold.

Thus in the case of Hertzian contact we can add the simplifications manifested in (41) and (43) to the conditions of Theorems I and II, respectively, and still garner the same conclusions.

Theorems I' and II':

For these cases \mathcal{M} can be written as an integral over $\tilde{\mathcal{C}}$, the projection of $\tilde{\mathcal{E}}$:

$$\mathcal{M} = - \int_{-\infty}^{\infty} \int_0^t \int_{\tilde{\mathcal{C}}} \underline{\sigma}^a \cdot \underline{n}^a (\underline{x}^a - \underline{x}^a) d\tilde{\mathcal{C}} dt.$$

Due to the present geometric situation, it is appropriate to take

$$\underline{n}^1 = \underline{n}^2,$$

and thus define

$$\underline{n} = \underline{n}^a, \quad a = 1, 2.$$

Analagous to the considerations for Theorems I and II, the momentum balance across $\tilde{\mathcal{E}}$ motivates the simplification

$$\underline{\sigma} \stackrel{\text{def}}{=} \underline{\sigma}^1 = -\underline{\sigma}^2.$$

With these and (39), the integrand of \mathcal{M} can be written

$$\underline{\sigma} \cdot \underline{n} (\underline{x}^2 - \underline{x}^1).$$

A further simplification can be made by setting^{*}

$$\sigma_3 = -\mathcal{M}.$$

This eliminates one unknown function and, as we shall see, has the effect of satisfying (5) naturally. Thus the integrand of \mathcal{M} becomes

* This is a standard ploy of optimization theory, see p. 82, [4].

$$\sigma_a \eta (x_a^2 - x_a^1) - (\eta)^2 (x_3^2 - x_3^1) , \quad (45)$$

and the analog of (23) is

$$\begin{aligned} 0 = \int_0^t \int_{\tilde{C}} \{ & \delta \eta (\sigma_a (x_a^2 - x_a^1) - 2\eta (x_3^2 - x_3^1)) \\ & + \delta \sigma_a (\eta (x_a^2 - x_a^1)) \\ & + \delta x_a^1 (T_a^1 - \eta \sigma_a) + \delta x_a^2 (T_a^2 + \eta \sigma_a) \\ & + \delta x_3^1 (T_3^1 + (\eta)^2) + \delta x_3^2 (T_3^2 - (\eta)^2) \} d\tilde{C} dt . \end{aligned} \quad (46)$$

Summand two tells us that either $\eta = 0$ or $x_a^1 = x_a^2$, on \tilde{C} .

Suppose $\eta \neq 0$, then $x_a^1 = x_a^2$, $a=1,2$. Summand one then gives us that $x_3^1 = x_3^2$ on \tilde{C} . Thus we have

$$\eta (x^2 - x^1) = 0 , \quad \text{on } \tilde{C} ,$$

as required, and \mathcal{C} is defined as the subset of \tilde{C} where $x^1 = x^2$.

The last four summands give the momentum balance conditions, as usual, and, in addition, the last two summands imply that the normal tractions are compressive (since $(\eta)^2 \geq 0$). Thus we have the conclusions of Theorem I' and condition (5).

The analogous set up for Theorem II' is accomplished by setting $\sigma_a = 0$ in (45) yielding

$$- (\eta)^2 (x_3^2 - x_3^1) \quad (47)$$

for the integrand of $\eta \eta$. With this Eq. (46) becomes

$$\begin{aligned}
0 = \int_0^t \int_{\bar{c}} \{ & -2 \delta \eta (\eta (x_3^2 - x_3^1)) + \\
& + \delta x_a^1 T_a^1 + \delta x_a^2 T_a^2 \\
& + \delta x_3^1 (T_3^1 + (\eta)^2) + \delta x_3^2 (T_3^2 - (\eta)^2) \} d\bar{c} dt
\end{aligned}$$

In this case we achieve the conclusion of Theorem II' and condition (5).

Thus to Theorems I' and II' we can delete condition (5), add the simplifications manifested in (45) and (47), and achieve the conclusions of Theorems I' and II', respectively, plus condition (5).

6. Contact Problems for One, Two and Three-dimensional Bodies

The previous work needs only trivial modification to be made applicable to contact problems involving bodies of different dimensions. There are many cases of considerable interest which fall into this category. For example, models consisting of a shell and a plate, or a solid and a plate, are useful for the study of head impact. The modifications necessary are essentially interpretative. An example illustrates this assertion.

Consider the frictionless Hertzian contact of a three-dimensional solid and a two-dimensional plate. Let B^1 represent the solid and B^2 the plate. In evaluating Π , the B^1 part is as before while the B^2 part would manifest the particular plate theory used. The contact term \mathcal{K} (or \mathcal{H}) would be exactly as before. However note that c (or \bar{c}) is, in this case, also identifiable with part of the two-dimensional "volume" of the plate, rather than its boundary. Taking variations, everything is as before except that the term $\tau \delta x_3^1$ (or $-(\tau_2)^2 \delta x_3^2$) contributes to the transverse momentum equation of the plate, rather than to its boundary conditions. The interpretation of τ (or $-(\tau_2)^2$) is thus two-fold, i.e., it is the normal component of the traction vector with respect to B^1 , as before, and it is also the equivalent normal "body force" with respect to B^2 , manifested by the interaction with B^1 (Fig. 7).

This interpretation is general, namely, for one and two-dimensional bodies the contact force is an equivalent "body force" which contributes to the momentum equations, rather than the boundary conditions. With this interpretation in mind, the construction of variational theorems, analogous to the ones constructed in Sections 2, 4 and 5, for the class of one, two and three-dimensional contact problems, is just a formal deductive exercise involving only appropriate definitions for Π .

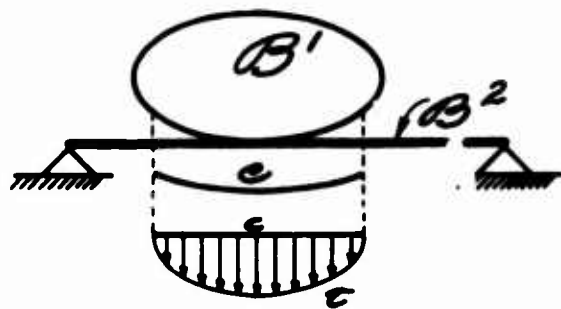


Figure 7

7. Impact

The previous sections deal with spatial aspects of contact problems. In this section we investigate temporal considerations, i.e., those phenomena which are unique to dynamic contact or impact. To manifest the problem encountered in such situations consider the following hypothetical situation. Assume that we are in the process of numerically solving some impact problem and suppose that it is discovered as we monitor the motion of the bodies that they impact somewhere in the time interval (t_1, t_2) . At time t_1 we know the states of both bodies and we know that somewhere between t_1 and t_2 they have coalesced over a portion of their boundaries. Assume for the moment we know the geometry of the contact surface \mathcal{C} . The question which arises then is what is the state of \mathcal{C} at time t_2 , i.e., what are the velocity and traction vectors on \mathcal{C} ? It is necessary to know this information to carry forth the step forward time integration. The question though seems improperly posed without specifying considerable data about the nature of the impact. To get a handle on things, we will initially formulate a simple one-dimensional problem involving the impact of two elastic rods. Although this problem is trivial, it provides considerable insight into the general nature of impact of continuum bodies. Since we are interested in the state of \mathcal{C} (in this case a point) immediately after impact, whether the rods are finite or semi-infinite is immaterial.

Assume that the pre-impact states of the two bodies are given by the following data:

$$\mathcal{B}^\alpha, \quad v_-^\alpha, \quad (\partial x / \partial X)_-^\alpha, \quad P_-^\alpha; \quad \alpha = 1, 2. \quad (48)$$

At impact the rods coalesce at e , and for some finite time interval thereafter (at least) $x \in e$ is persistent. At the moment of impact shock waves begin to propagate in each body. The space-time picture is depicted in Fig. 8. As discussed in section 1, since e is material and x is persistent, we have

$$\underline{v} \stackrel{\text{def}}{=} \underline{v}_+^1 = \underline{v}_+^2, \quad P \stackrel{\text{def}}{=} T_+^1 = -T_+^2, \quad (49)^*$$

for the post-impact state (t_+). In addition to (49), the well known shock conditions must hold across the wave fronts:

$$\begin{aligned} [\underline{v}^a] + U^a [(\partial x / \partial X)^a] &= 0, \\ \rho_0^a U^a [\underline{v}^a] &= [P^a], \end{aligned} \quad (50)$$

where U^a is the material velocity of the shock in B^a , and $[\]$ is the wave-front jump operator which assigns to a function the difference in its values behind and in front of the wave, i.e., $[f(X,t)] = f(X^-,t) - f(X^+,t)$ where X is a material point denoting the location of the wave-front. As can be deduced from Fig. 8, the states into which the shocks initially propagate are the pre-impact states given by (48), and the state at e , immediately after the shocks pass, is given by the post-impact state (49). These observations in conjunction with (50) yield,

$$\begin{aligned} \underline{v}_-^a - \underline{v} + U^a \{ (\partial x / \partial X)_-^a - (\partial x / \partial X)_+^a \} &= 0, \\ \rho_0^a U^a (\underline{v}_-^a - \underline{v}) + P_-^a - P &= 0. \end{aligned} \quad (51)^{**}$$

*For convenience we choose the initial state to be the pre-impact state, thus we need not distinguish between Cauchy and Piola tractions.

**A consistency condition for these equations is that $\underline{v}_-^1 - \underline{v}_-^2 > 0$. Otherwise the impact would not occur.

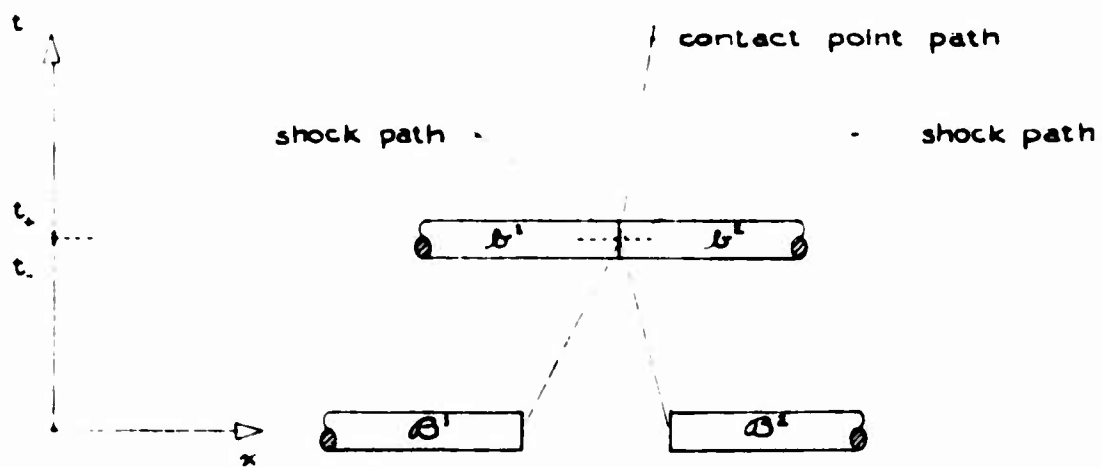


Figure 8

The four Eqs. (51) and constitutive equations relating P^α to $(\partial x / \partial X)^\alpha$ yield a formally deterministic system of six equations in the six unknowns $v, P, U^\alpha, (\partial x / \partial X)^\alpha$. Thus we see that the desired quantities v and P depend on the pre-impact states and material properties of both B^α . The precise form of this relationship depends upon the constitutive equations of the bodies. As a simple example, assume we have linear constitutive equations $P^\alpha = E^\alpha \{ (\partial x / \partial X)^\alpha - 1 \}$, E^α constant, and let the pre-impact state be given by

$$\begin{aligned} v_-^\alpha &= V^\alpha, \\ (\partial x / \partial X)_-^\alpha &= 1, \\ P_-^\alpha &= 0. \end{aligned} \quad (52)$$

These conditions, when inserted in Eqs. (51), lead to:

$$\begin{aligned} v &= \frac{\rho_0^2 U^2 V^2 - \rho_0^1 U^1 V^1}{\rho_0^2 U^2 - \rho_0^1 U^1}, \\ P &= \frac{V^2 - V^1}{\left(\frac{U^2}{E^2} - \frac{U^1}{E^1} \right)}, \\ (U^\alpha)^2 &= E^\alpha / \rho_0^\alpha. \end{aligned} \quad (53)$$

Note that the denominators in Eqs. (53)_{1,2} present no problems since $\rho_0^\alpha > 0$, $E^\alpha > 0$ and $+1 = \text{sgn } U^2 = -\text{sgn } U^1$.

This result is also appropriate whenever the intensity of the impact is small enough such that the non-linear constitutive equation can be replaced by its linear approximation about the pre-impact state. In this case E^α is a tangent modulus evaluated at the pre-impact strain

$\{ (\partial x / \partial X)_-^\alpha - 1 \} = 0$. To further simplify, consider the case when both rods have identical properties (i.e., $\rho_0 = \rho_0^\alpha$, $E = E^\alpha$, $\alpha = 1, 2$). Then

$$\begin{aligned}
 v &= \frac{V^1 + V^2}{2} , \\
 P &= \rho U (V^2 - V^1) / 2 , \\
 (U)^2 &= E / \rho .
 \end{aligned}
 \tag{54}$$

In Eqs. (54) U is positive, and since consistency requires $V^1 - V^2 > 0$, P is compressive.

Thus for the one-dimensional case at least the problem of computing the post-impact state is easily achieved. The solution of (51) for the fully non-linear case can be automated as part of a numerical algorithm. Although this problem is trivial, it serves to indicate that the post-impact problem, the solution of which is essential in a numerical algorithm, is one of wave propagation.

In the analysis of higher dimensional bodies the solution of the post-impact problem becomes greatly complicated due to the geometric variety of impact conditions. However, considerable simplifications can be taken advantage of if one keeps in mind the nature of the discrete problem. For instance, if a certain portion of the boundaries of two bodies have coalesced in e , each interior point of e , at which the tangent plane is well defined, may be treated, to the first order, as a point on the mating surface of two impacting half-spaces. As long as time steps are kept small enough, the local behavior is well represented. The post-impact problem for the general case, analogous to (51), can be automated as part of the numerical algorithm, and for many simple cases can be solved explicitly.

With these notions in mind, let us return to the case of main interest in this report, namely three-dimensional continuum bodies. We shall consider only the case of a frictionless contact surface (Case II), and leave the solution of the post-impact problem for the no-slip case (Case I), which is more difficult, for future work. With the proper interpretations, the one-dimensional rod formulation (Eqs. (48-54)) suffices to completely characterize this case. This is so because no tangential motions or stresses may be communicated across a frictionless surface, and thus we need only consider the configuration of normal incidence. In this case the requisite constitutive functionals in (51) would be those relating P^α , the normal Piola stress, to the normal component of strain, holding all other components of strain fixed at the pre-impact values. For example, in the linear isotropic case, E^α (Young's modulus) in Eqs. (53,54) would be replaced by $\lambda^\alpha + 2\mu^\alpha$ ($\lambda^\alpha, \mu^\alpha$ are the Lamé and shear moduli, respectively) and the propagation velocity would be that of dilatational waves.

PART II

A NUMERICAL SCHEME FOR ANALYSIS OF
CONTACT-IMPACT PROBLEMS8. Numerical Solution of Contact-Impact Problems

In performing numerical computations based on the above described variational formulation for contact-impact problems we have employed three distinct levels of approximation: (1) a spatial discretization of the bodies and contact surfaces, (2) a temporal discretization to determine the response of the discretized bodies, and (3) a numerical solution for the resulting system of nonlinear algebraic equations.

In the following sections we shall restrict our attention to the Hertzian contact problem described in Section 5. Significant numerical difficulties are encountered in the solution of impact problems; to complicate the problem further by introducing the additional steps necessary to determine the contact surface maps for the full kinematically nonlinear case is left for a future study. While this is a simple impact problem in terms of determining the contact surface and the full power of the preceding theory is neither necessary nor exploited in its solution, many of the features of the general problem are employed here.

9. Spatial Discretization of the Bodies and Contact Surface

The bodies \mathcal{B}^1 and \mathcal{B}^2 are discretized using standard finite element methods, (e.g., see [5]). In order to facilitate the computation of a discrete Hertzian contact surface the nodes of \mathcal{B}^1 are arranged so that they align with the nodes of \mathcal{B}^2 . This is consistent with the notions of condition 3 of Section 5 and ensures that during determination of the approximation to the contact surface contiguous nodes of the two bodies will meet. Thus, the simulation of the contact surface is trivial. The development of a numerical model for Hertzian contact problems is based upon the form of Theorem II' which uses (47) for the integrand of \mathcal{M} . For numerical computations we introduce the displacement vector \underline{u} such that

$$\underline{x} = \underline{X} + \underline{u} \quad (55)$$

For a compatible finite element displacement field the integrand of \mathcal{M} can be approximated by taking $\eta^2(\underline{x}, t)$ as the product of $\mathcal{E}^2(t)$ and $\delta(\underline{x} - \underline{x}_i)$ (i.e., Dirac delta functions in space). This corresponds to taking η^2 as "concentrated nodal loads" which are the generalized forces of the contact pressure. With this discretization we can describe pseudo contact elements between each pair of candidate contact nodes. Let these nodes be denoted as $(\cdot)_i$ and the generalized force as $(\mathcal{E}_i)^2$; then

$$\mathcal{M} = \int_0^t \sum_i (\mathcal{E}_i(t))^2 (u_{z_i}^2(t) - u_{z_i}^1(t) + X_{z_i}^2 - X_{z_i}^1) dt \quad (56)$$

where $\{i\}$ are the set of candidate contact nodes which span $\tilde{\mathcal{C}}$; $u_{z_i}^{\cdot}$ are the nodal displacements in the z_3 direction and $X_{z_i}^{\cdot}$ are the nodal coordinates of the candidate contact nodes.*

*We assume here that 3 is the direction nominally normal to the contact surfaces, e.g., see Fig. 5.

Use of the finite element method in Theorem II' with \mathcal{M} given by (56) produces a set of nonlinear second order ordinary differential equations which together with the impenetrability conditions define the discretized contact impact problem. These equations take the form:

$$\underline{\underline{M}} \ddot{\underline{u}} + \underline{\underline{K}}(\underline{u}) = \underline{\underline{R}}, \quad (57)$$

where $\underline{\underline{M}}$ is the usual finite element mass matrix, $\underline{\underline{K}}$ represents the elastic stiffness forces together with the contact terms, $\underline{\underline{R}}$ is the set of generalized forces resulting from boundary tractions and \underline{u} is the set of time dependent nodal displacements (which also include the $(\epsilon_i)^2$). For inelastic materials Theorem II' can be extended by treating the first variation as a Galerkin method (principle of virtual work) and replacing the elastic constitutive model by more general theories, e.g., viscoelastic, elastoplastic, viscoplastic, etc. In this case

$$\underline{\underline{K}}(\underline{u}) \rightarrow \underline{\underline{K}}(\underline{u}, \dot{\underline{u}}) \quad (58)$$

in (57).

10. Temporal Discretization

A temporal discretization of the second order ordinary differential equations which result from a finite element spatial discretization of the contact-impact problem is accomplished herein by using the Newmark family of methods [6]. The Newmark family of methods is a one-step integration method with two free parameters which can be used to control stability and numerical damping. The method is essentially a difference method in time. The behavior of the method for linear elasto-dynamics problems is discussed in [6,7]. The algorithm is given by

$$\underline{u}_{n+1} = \underline{u}_n + \Delta t \dot{\underline{u}}_n + \left(\frac{1}{2} - \beta\right) \Delta t^2 \ddot{\underline{u}}_n + \beta \Delta t^2 \ddot{\underline{u}}_{n+1}, \quad (59)$$

$$\text{and } \dot{\underline{u}}_{n+1} = \dot{\underline{u}}_n + (1-\gamma) \Delta t \ddot{\underline{u}}_n + \gamma \Delta t \ddot{\underline{u}}_{n+1},$$

where $\underline{u}_n = \underline{u}(t_n)$, $\Delta t = t_{n+1} - t_n$, and β, γ are the two parameters. For linear problems $\gamma = .5 + \delta = .5$ produces no artificial viscosity and $\beta \geq \frac{1}{4} (1 + \gamma)^2$ produces unconditional stability (i.e., the method is stiffly stable). Such generalization is not possible for nonlinear problems and during solution it may be necessary to monitor the solution for any signs of instability. In (59) $\beta = 0$ produces an explicit method for \underline{u}_{n+1} and if \underline{M} is diagonal (lumped mass) with \underline{K} and \underline{Q} independent of $\dot{\underline{u}}$ the solution can be advanced without solving a large set of simultaneous equations; for all other cases the method is implicit and equations must be solved. Solution of (59)₁ for $\ddot{\underline{u}}_{n+1}$ in terms of the solution at t_n and \underline{u}_{n+1} gives

$$\ddot{\underline{u}}_{n+1} = \frac{1}{\beta \Delta t^2} (\underline{u}_{n+1} - \underline{u}_n) - \frac{1}{\beta \Delta t} \dot{\underline{u}}_n - \left(\frac{1-2\beta}{2\beta}\right) \ddot{\underline{u}}_n \quad (60)$$

which can also be used in $(59)_2$ to express the velocity in terms of the solution at t_n and u_{n+1} . Since in this process we divide by β and Δt it is no longer possible to consider zero β or zero time steps.

11. Solution of the Nonlinear Algebraic Problem

Use of the Newmark method in (57) (including (58)) yields the set of nonlinear algebraic equations:

$$\frac{1}{\beta \Delta t^2} \underline{M} \underline{u}_{n+1} + \underline{K}(\underline{u}_{n+1}, \underline{u}_n, \dot{\underline{u}}_n, \ddot{\underline{u}}_n) = \underline{R}_{n+1} + \underline{M} \underline{A}_n, \quad (61)$$

where

$$\underline{A}_n = \frac{1}{\beta \Delta t^2} \underline{u}_n + \frac{1}{\beta \Delta t} \dot{\underline{u}}_n + \left(\frac{1-2\beta}{2\beta} \right) \ddot{\underline{u}}_n.$$

A Newton-Raphson iterative solution to this set of equations can formally be constructed, giving:

$$\left(\frac{1}{\beta \Delta t^2} \underline{M} + \partial_u \underline{K} - \partial_u \underline{R} \right) \Delta \underline{u}^{(i)} = \underline{R} - \underline{K}(\underline{u}_{n+1}^{(i)}, \underline{u}_n, \dot{\underline{u}}_n, \ddot{\underline{u}}_n) - \underline{M} \ddot{\underline{u}}_{n+1}^{(i)}, \quad (62)$$

where $\partial_u \underline{R}$ is the effect of loads varying with the deformation and

$$(\partial_u \underline{K})_{ij} = \partial K_i / \partial u_j, \quad (63)$$

is the tangent stiffness matrix. The coefficient to $\Delta \underline{u}^{(i)}$ is generally called the Jacobian matrix of the Newton-Raphson iteration. The solution is advanced by taking

$$\underline{u}_{n+1}^{(i+1)} = \underline{u}_{n+1}^{(i)} + \Delta \underline{u}^{(i)}, \quad (64)$$

and iterating until a norm of the solution satisfies

$$\|\Delta \underline{u}^{(i)}\| \leq \epsilon \|\underline{u}_{n+1}^{(i)}\|, \quad (65)$$

where ϵ is some small positive error tolerance. In the work reported here the norm $\| \cdot \|$ is taken as the Euclidian norm

$$\| \underline{x} \| = \left(\sum_i x_i^2 \right)^{1/2}, \quad (66)$$

and the load vector \underline{R} is assumed to be independent of \underline{u} . For stable elastic materials the resulting tangent stiffness is then symmetric and positive definite, consequently, standard direct solution methods normally employed in the solution of linear finite element problems can be used. For inelastic materials or deformation dependent loads the tangent stiffness may be asymmetric. In these cases some special methods may be necessary to effect a solution.

12. Discretized Impact Conditions

In the previous numerical development \tilde{c} has been defined by discrete points which correspond to nodes along the boundaries of \mathcal{B}^1 and \mathcal{B}^2 . When, during the course of advancing the solution in time, any one of these points violates the impenetrability condition a re-solution must be obtained in which the $(\epsilon_i)^2$ are now non-zero and the u_3^* satisfy the impenetrability condition. Some control and monitoring are required to effect this in a computer program. In addition to satisfying these conditions, the impact relations denoted in Section 7 must be invoked. In the present study these conditions are applied to the solution at the end of a time step in which points first go into contact. Accordingly we compute from (50)*

$$\dot{u}_+ = \frac{e_2^2 U^2 \dot{u}_-^2 - e_1^1 U^1 \dot{u}_-^1}{e_2^2 U^2 - e_1^1 U^1}, \quad (67)$$

and assign this value to the appropriate node of \mathcal{B}^1 and \mathcal{B}^2 .

To determine the solution vector \underline{u} at t_{n+1} , we have solved the set of equations (61). As described above the shock conditions are then used to determine the value of the velocity at time t_{n+1} for all points which have come into contact during the time interval. In order to get a consistent solution at these points we must modify the accelerations and contact force to reflect the shock conditions. This is accomplished by re-solving the equilibrium conditions of \mathcal{B}^1 and \mathcal{B}^2 at point i . The expanded forms of the appropriate equations are:

*The $()_-$ denotes a value which is computed before impact, whereas $()_+$ denotes the value after impact.

$$M^1 \ddot{u}_-^1 + K^1(\underline{u}) + (\mathcal{E}_i)_-^2 = R^1, \quad (68)$$

and

$$M^2 \ddot{u}_-^2 + K^2(\underline{u}) + (\mathcal{E}_i)_-^2 = R^2.$$

For nodes which have come into contact we must enforce the condition on acceleration

$$\ddot{u}_+^1 = \ddot{u}_+^2 = \ddot{u}_+, \quad (69)$$

and compute the contact force $(\mathcal{E}_i)_+^2$. The solution for these is obtained from

$$M^1 \ddot{u}_+ + K^1(\underline{u}) + (\mathcal{E}_i)_+^2 = R^1,$$

and

$$M^2 \ddot{u}_+ + K^2(\underline{u}) + (\mathcal{E}_i)_+^2 = R^2.$$

These are two equations in two unknowns which can be solved for the \ddot{u}_+ and $(\mathcal{E}_i)_+^2$. If $K^{\alpha}(\underline{u})$ is independent of velocity the stiffness forces and R^{α} will remain unchanged during the impact, hence we can solve the simpler problem

$$M^1 \ddot{u}_+ + (\mathcal{E}_i)_+^2 = M^1 \ddot{u}_-^1 + (\mathcal{E}_i)_-^2$$

$$M^2 \ddot{u}_+ - (\mathcal{E}_i)_-^2 = M^2 \ddot{u}_-^2 - (\mathcal{E}_i)_-^2$$

whose solution is

$$\ddot{u}_+ = \frac{M^1 \ddot{u}_-^1 + M^2 \ddot{u}_-^2}{M^1 + M^2},$$

and

$$\begin{aligned} 2(\mathcal{E}_i)_+^2 &= 2(\mathcal{E}_i)_-^2 + M^1 (\ddot{u}_-^1 - \ddot{u}_+) \\ &\quad - M^2 (\ddot{u}_-^2 - \ddot{u}_+). \end{aligned}$$

This completes the numerical specification of the solution at t_{n+1} ; this solution process is now repeated for each of the succeeding time steps.

At this point it is important to compare the solution procedure for impact of a continuum discretized by a finite element method with the solution procedure for a physically discrete body, i.e., a body composed of mass points joined by massless elastic springs. Both problems may be described by algebraic equations of the form of (57). The impenetrability condition is also identical. The impact conditions, however, are different. For the discretized continuum the procedure is described above. The study of the impact of mass points is considered in elementary mechanics books, e.g. [8]. The impact of two mass points is described by impulsive motion such that at t_- the velocities of the two mass points are V_-^1 and V_-^2 ; after impact at time t_+ , the two points have velocities V_+^1 and V_+^2 . The two points will not in general stay in contact (i.e., $V_+^1 \neq V_+^2$) but will rebound. The conditions used to compute the V_+^1 and V_+^2 are:

Balance of Momentum*

$$M^1 \{V^1\} + M^2 \{V^2\} = 0, \quad (71)$$

and use of an equation involving the "coefficient of restitution", e :

$$\frac{V_+^2 - V_+^1}{V_-^1 - V_-^2} = e. \quad (72)$$

For $e=1$ energy is conserved whereas for $e=0$ the points "stick" and energy is dissipated. We must comment in passing that (72) is the energy

* $\{f(t)\} = f(t_+) - f(t_-)$.

equation in disguise. To see this we can write the jump conditions for energy as

$$\frac{1}{2} M^1 \{(\dot{V}^1)^2\} + \frac{1}{2} M^2 \{(\dot{V}^2)^2\} = \{\dot{V}\} \quad (73)$$

The term $\{\dot{V}\}$ can exist only if other energies are dissipated during the jump. We rewrite (73) by using

$$\frac{1}{2} \{(\dot{V}^1)^2\} = [\dot{V}^1] \langle \dot{V}^1 \rangle ,$$

where

$$\langle \dot{V}^1 \rangle = \frac{1}{2} (\dot{V}_+^1 + \dot{V}_-^1) \quad (74)$$

Use of the momentum equation (71) then gives, after dividing by $M^1 \{\dot{V}^1\}$

$$\langle \dot{V}^1 \rangle - \langle \dot{V}^2 \rangle = \frac{\{\dot{V}\}}{M^1 \{\dot{V}^1\}} ,$$

or after recollecting terms and dividing by $(\dot{V}_-^1 - \dot{V}_-^2)$ we obtain:

$$\frac{\dot{V}_+^2 - \dot{V}_+^1}{\dot{V}_-^1 - \dot{V}_-^2} = 1 - \frac{\{\dot{V}\}}{M^1 \{\dot{V}^1\} (\dot{V}_-^1 - \dot{V}_-^2)} \quad (75)$$

The significance of the coefficient of restitution then is associated with the right hand side of (75).

It is clear from the above developments that the numerical simulation of the discretized continuum and the physically discrete system involve two distinct methods for treating the impact conditions. It is imperative then to associate the correct method for the problem at hand. In the

present study we are interested in the impact of continua, and in this case we shall employ the discrete shock condition to effect the solution. This a priori assumes that the response we are computing involves a time scale associated with wave propagation problems. Consequently, we cannot expect the computation procedure for advancing the solution in time to be accurate if we take time steps greatly in excess of transit times through each body. In this context it may be important to consider an "explicit" time integration procedure in future work. The stability restrictions may be too severe to make this feasible.

PART III

FEAP 74 - A COMPUTER PROGRAM FOR
SOLUTION OF CONTACT-IMPACT PROBLEMS13. Development of a Contact-Impact Model for FEAP

In order to incorporate an ability to compute solutions to contact-impact problems using a finite element method as described above it is necessary to have available a computer program which can solve the nonlinear equations of motion given by (61). The computer program FEAP is a general program to solve finite element problems. The program has a capability of solving both quasistatic and dynamic problems and can incorporate several types of elements simultaneously. The nonlinear capabilities required for the solution of contact-impact problems have been incorporated into FEAP and currently includes the user options (see Input Instructions):

- (1) Selection of quasistatic or dynamic option: The dynamic option will integrate the equations of motion using the one-step Newmark method to advance the solution in time. Quasistatic analysis is accommodated by any one-step algorithm. The algorithm employed is incorporated into each element routine and thus is defined by the developer of each element. Impact problems require description of the contact surface and wave speeds.
- (2) Selection of the nonlinear method to advance the solution: Options include:
 - (a) No iterations in each time step. Unbalanced forces at each time are added to the next time step.
 - (b) Iterations in each time step to achieve a balance of force within each time step. In this option the user can select to reform

the Jacobian matrix for each iteration or only at the first iteration in each time step.

In the impact problems solved to date it has been necessary to use the general form of the Newton-Raphson algorithm. This includes a complete forming and factoring of the Jacobian matrix for each iteration of each time step in the analysis. If the method described herein is to become computationally effective improvements in the computer program are paramount. Undoubtedly the most important aspect in reducing computer times is to introduce a substructuring system so that the highly nonlinear equations in the vicinity of the contact surface can be isolated from the remainder of the bodies. This will normally involve only a small number of equations in the total system of (62). The solution of a large finite element problem will generally concentrate the computer solution time in the forming and factoring of the tangent stiffness matrix. The fewer times that it is necessary to perform this costly step the more efficient the solution algorithm. Substructuring can be used then to restrict the part of the equations which must be formed and factored often, and thus greatly reduce the computer costs in analyzing impact problems.

The version of FEAP which can currently be used to analyze contact-impact problems includes, in addition to the nonlinear Newton-Raphson iterative algorithm, a new special contact-impact element and a new subroutine to describe impact surfaces and the discrete shock conditions described in Section 12. These are described in the following sections.

14. Contact Element for Hertzian Contact

The contact-impact element which has been developed is called ELMT05 and can be used along any coordinate direction. As developed it cannot be used along normals which are in non-coordinate directions. The development of the contact element assumes that within the framework of linear elasticity theory a node on ω^1 will impact on a node of ω^2 . In using this contact element we shall assume that the contact surface on ω^2 is located at larger coordinate values than the contact surface of ω^1 . The contact element is described by three nodes. Node 1 is associated with ω^1 , Node 3 is associated with ω^2 , and Node 2 is used as storage for the contact force $(E)^2$. The user can select the direction of contact motion by specifying the degree of freedom of the nodal unknowns to which the contact is to be measured; this must agree with the physical direction of the element (see Fig. 9). The degree of freedom for the contact element is specified during the MATERIAL data input and consists of a single card in I5 format. The element nodes are described along with all other elements according to Section 4 of the Input Instructions. The Node order as shown in Fig. 9 must be observed.

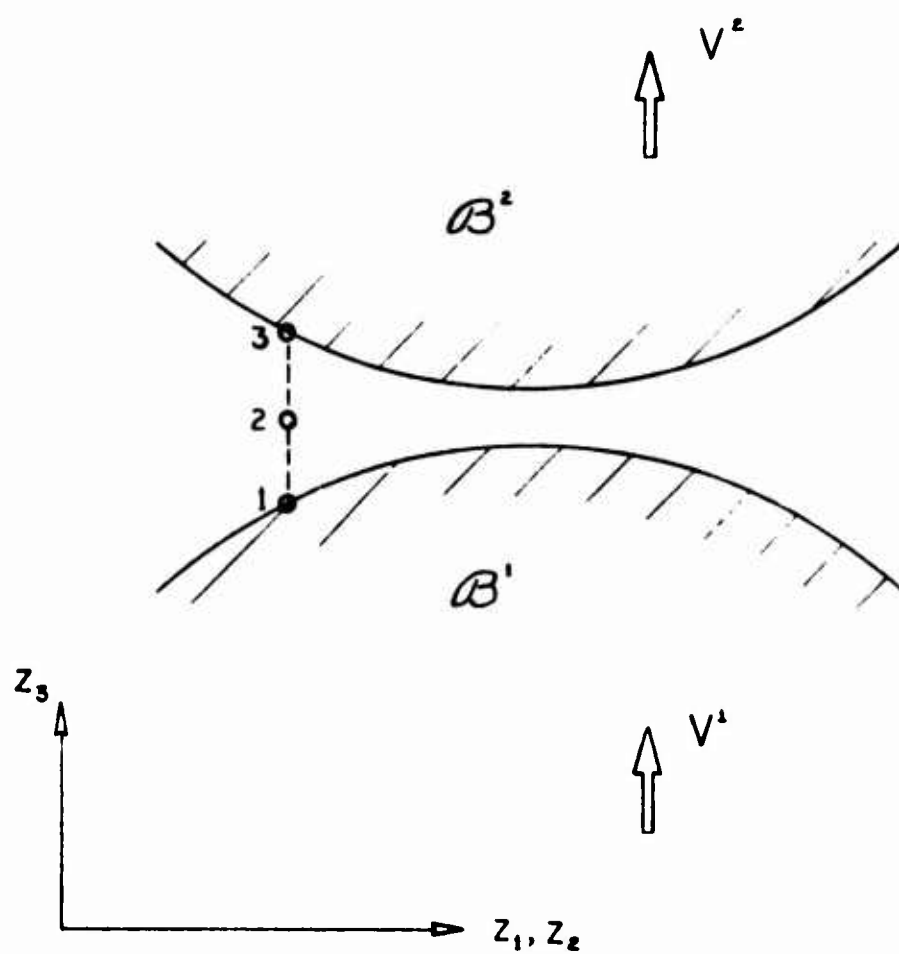


Figure 9

15. Impact Surface Description

The definition of the impact surface includes a list of all elements on the contact surface together with the degree of freedom describing the direction of contact motion (as described above). In addition, the product of mass density and wave velocity (always a positive number) for each body is input. This assumes, currently, that (1) each contact surface belongs to a linear material, and (2) the same material exists along all of the contact surface. This data need be prescribed only for impact problems, quasistatic contact problems do not require this data since no velocity or acceleration computations are performed in this class of problems. Data to be input for the impact surface is given in Table I.

Table I - Impact Surface Data

CARD 1) (6X,A6)

COL. 7 to 12

Must contain CONTAC

CARD 2) (2F10.0)

COL. 1 to 10

ρU of body 1

COL. 11 to 20

ρU of body 2

CARD 3) (I5)

COL. 1 to 5

NLIST, number of elements on contact surface

CARD 4) (2I5)

Repeat NLIST times

COL. 1 to 5

Contact element number

COL. 6 to 10

Degree of freedom of this contact element

16. Example Problems

Two example problems are included to illustrate the characteristics of the methodology and the associated computer program described above for Hertzian contact problems. The first problem is a quasistatic contact problem which is used to demonstrate the ability of the computer program to compute an evolving contact surface. The second problem will demonstrate the ability of the program to properly model the temporal response of an impact problem.

To model a problem in which a contact surface will change under different load levels we consider two beams with an initial parabolic curvature. A symmetric configuration is analyzed and the resulting finite element model is shown in Fig. 10. Each element is nominally one unit by one unit. The gap at the load end is initially 0.5 units. The material properties used are $E = 500$ and $\nu = 0$. The load P is applied as shown and allowed to increase linearly in time. The problem then is to determine the contact surface at various load levels. In order to eliminate a singularity in the system of equations it was necessary to permanently attach the two nodes at the symmetry axis of the contact surface. All other nodes along the boundaries between the two bodies are assumed to be possible contact points and contact elements are assigned between vertical nodal pairs. The load was varied from 0.2 to 0.8 in increments of 0.1 and the computed contact surface and forces were computed. These are given in Table II. The deformed shape at a load of 0.4 is also shown in Fig. 10 as a dotted form. The attached node at the center has influenced the solution at loads above 0.3 since the contact pressure there is tensile (negative). The force is small and should not greatly affect the actual contact region computed. As the load increases the

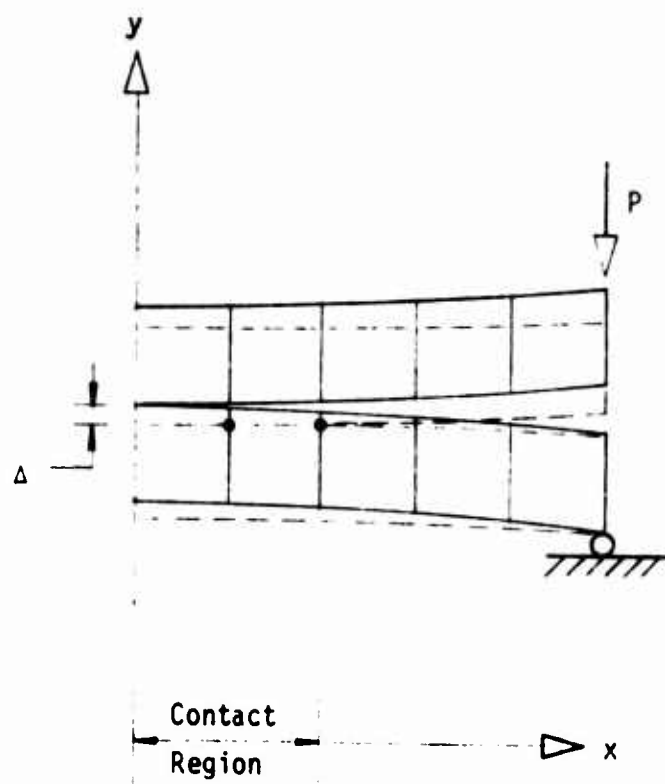


Figure 10

contact surface moves toward the load. This is conceptually correct since if the beams were modeled according to Euler-Bernoulli theory the contact force would be a point load which gradually moves from the center to the outer edge according to the relationship (using the above values for sizes and material properties)

$$X = 5 \left(1 - \frac{1}{6P} \right) .$$

This relation predicts that the contact point will be non-zero only after P exceeds 1/6. The finite element model is in qualitative agreement with this beam theory, but since shear deformations are included the finite element solution gives a distributed load on the contact surface. It is interesting to also note that the contact force over the center of the beams is zero, just as in the beam theory.

Table II - Contact Forces

LOAD	X-COORDINATE						BEAM THEORY-X
	0	1	2	3	4	5	
0.2	0.2	-	-	-	-	-	0.83
0.3	.07	.23	-	-	-	-	2.22
0.4	-.01	.07	.34	-	-	-	2.92
0.5	-0.00	.02	.23	.25	-	-	3.33
0.6	-.01	-	.09	.51	.01	-	3.61
0.7	-.01	-	.07	.45	.19	-	3.81
0.8	-.01	-	.05	.38	.38	-	3.96

This problem demonstrates that the computer program can model the evolution of a contact surface. Of particular importance is to note that as the load increases the program can both attach and detach a contact point. This is an essential requirement for the analysis of

impact problems as is shown in the next problem.

As a simple example we consider the impact against a rigid wall of a finite, linear elastic rod traveling at constant velocity. The rod has a modulus of elasticity E of 100, and a mass density ρ of 0.1. The arrival velocity is taken to be 0.1 (units may be assigned in any convenient system). The rod is taken to be 10 units long and is divided into 10 elements plus one contact element as shown in Fig. 11. At time zero the rod is just arriving at the wall. The exact solution predicts a contact duration of 0.2 time units. This corresponds to the time required for a wave to travel from the contact point to the left end and back to the contact point at which time the rod will part from the wall. The problem was analyzed using FEAP with time steps of 0.01 unit (transit time across an element) and the rod remains in contact until time 0.20 units and has rebounded at time 0.21. Thus the program can predict accurately the contact duration of the rod. The finite element solution obtained is compared with the exact solution in Fig. 12. The agreement of stresses and contact force is good. The largest discrepancy exists in defining the shock front, which is "smeared" by the finite element method and ordinary differential equation solution method used here. This is the same type of solutions which are commonly obtained with numerical solutions of this type even without impact. Solutions such as the impact shocks generated are probably one of the most difficult responses to accurately calculate by a finite element method.

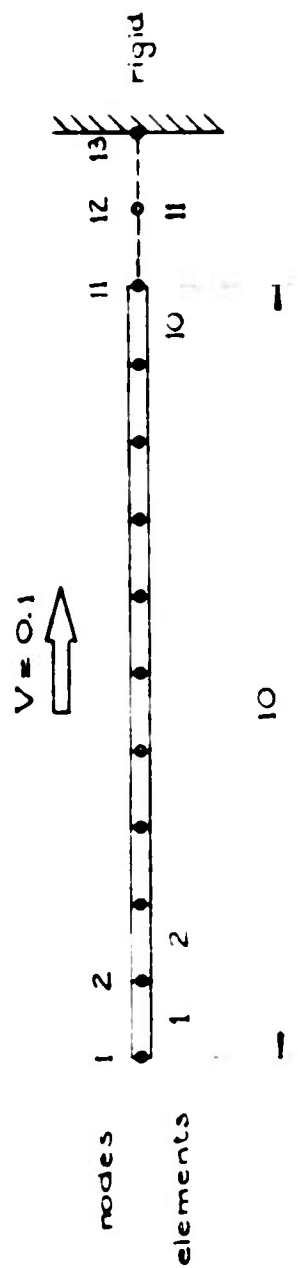


Figure 11

Comparison of Numerical Data with Exact Solution for Bar Impacting Rigid Wall

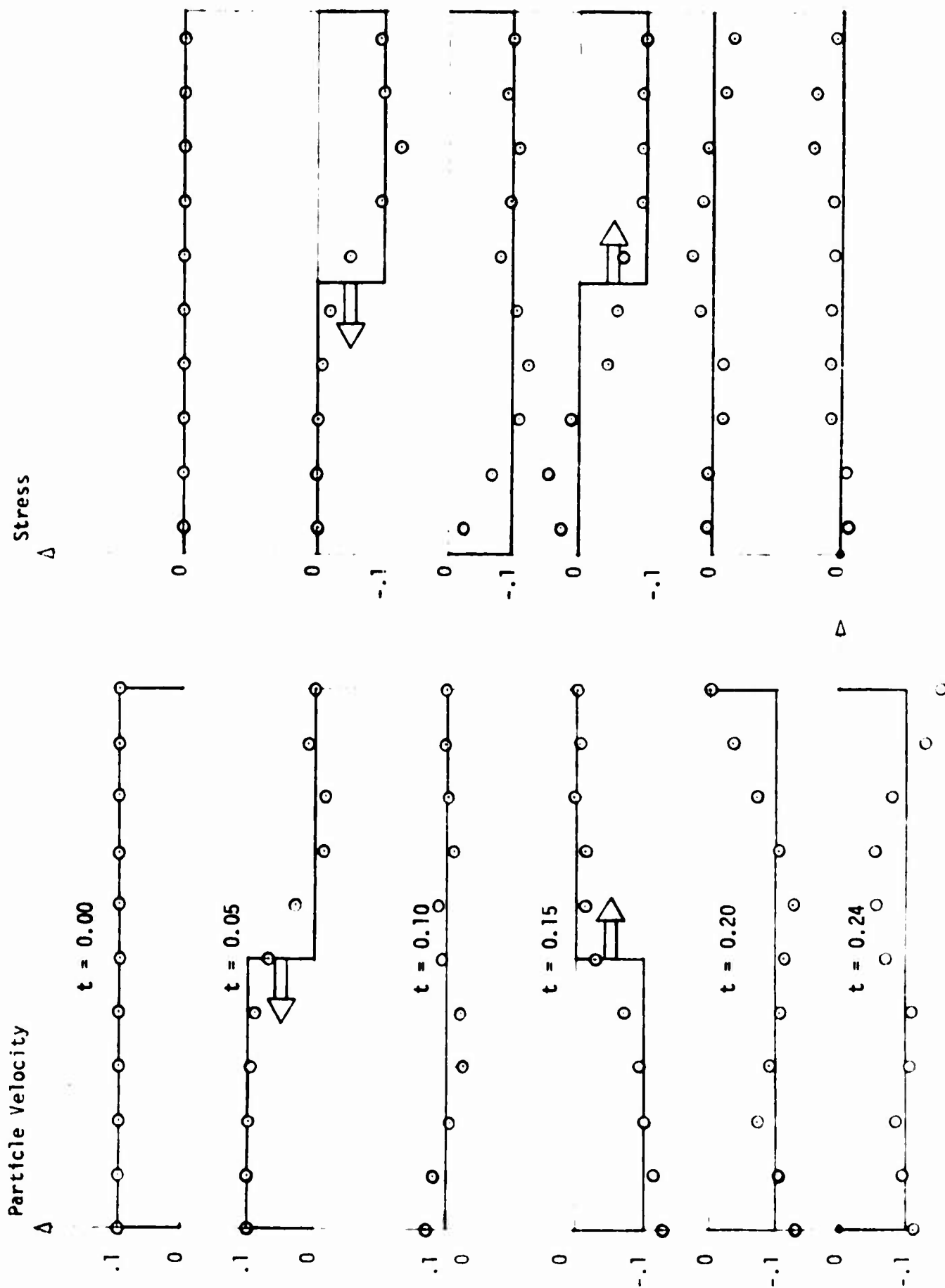


Figure 12

17. Closure and Recommendations for Future Work

In the preceding sections we have presented a theory for contact-impact problems together with the numerical development of a Hertzian contact-impact model. The computer program FEAP 74 has been modified to include the model and has successfully solved a contact problem and an impact problem. The work reported herein must be considered initiatory; the general theories and their numerical implementation have not been completed. The problem is of such a complicated nature and the literature existing prior to this study was so meager that we consider it fitting to document the work completed thus far.

We have attempted to qualify each stage of the development throughout the report, however, it may be fitting to reiterate future work which we consider to be essential for numerical models to be effective and efficient tools for predictive analyses.

- (1) The restriction of Hertzian type contact must be removed. This involves the non-trivial task of finding appropriate numerical methods to handle the χ^4 maps.
- (2) Improved methods for solving the set of nonlinear algebraic equations must be found. We have suggested two methods which should be considered: (a) Substructure the problem about the contact regime so that a more efficient forming and factoring of the tangent stiffness can be performed; and (b) Since the impact problem is a wave propagation problem an explicit time integration of the equations of motion should be explored. In complex situations the explicit integration method may have severe stability limitations which could make it unacceptable.

- (3) Methods of utilizing the shock conditions need to be explored further. We have noted some peculiar anomalies when the bodies separate. These appear to be caused by a shock like separation phenomena.
- (4) When the wave propagation property of the impact problem is ignored by taking time steps greatly in excess of the transit times in a body the computed response is meaningless. Under such situations the bodies rebound within a single time step. Currently the rebound velocity is much too large. When the shock conditions are used for a class of problems where the response desired is in the target instead of in the impactor, it may be expedient to take a large time step. Methods should be explored to accomplish this capability.

The above recommendations for future work should in no way minimize what has been accomplished by the present study. For the first time a contact-impact theory in the form of a variational problem has been presented in a general form. This formulation was motivated by the fact that numerical solutions would be obtained by a finite element method. In addition the necessary foundation for the numerical solution has been thought out and within this context a computer program has been developed for Hertzian contact-impact problems.

The implementations considered here have produced results which are hopeful signs for the eventual success of the more general impact problems.

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APPENDICES

A. Input Instructions for Contact/Impact Problems

In order to analyze contact/impact problems in FEAP, users must prepare the data for a time dependent analysis. This will include the following Data Type Identification Cards (see Section 1, Appendix B):

FEAP 74

MATERIAL

NODAL

ELEMENT

CONTACT (for impact problems only)

loadings

INITIAL CONDITIONS (if non-zero)

and

VISCOE (for quasi-static contact problems)

or

IMPLICIT (for impact problems)

In performing the necessary solution to (62) the full Newton-Raphson method must be employed; this is controlled by the data in col. 76-80 of the first card following the VISCOE or IMPLICIT card, and consists of a negative number (negative uses full Newton-Raphson iteration with the absolute value of the number giving the number of iterations to be performed before going to the next time step). As an example of the required input data Table A shows the input data used for the impact problem reported in Section 16.

14.19.14 OUTPUT FEAP74 * * ONE DIMENSIONAL BAR CONTACT PROBLEM

```

1 1 X
1 1 U
3 13 NODAL
1 1 0.
11 10.
13 10.
11 111111
11 ELEMENTS
1 1 1
1 2 2
11 12 13
11 3 INITIAL CONDITIONS
DISPLACEMENT
VELOCITY
ACCELERATIONS
1 1
11 0

```

```

2 MATERIALS
1 ELMO9
100. 1. .01
2 ELMO5
1 CONTACT
1 1.0 1.0E+30
11 1

```

```

1 IMPLICIT
1.0 5 1 1 13 1 11 0 0 0.25 -3
.01 25 1 1 13 1 11 0 0 0.25 -3
.01 STOP 10 1 1 13 3 11 0 0 0.25 -3

```

B. Input Instructions for FEAP 74

The input instructions for the description of a finite element mesh, together with the initial and boundary conditions, is described by subroutine MANUAL listed on the following pages. The input of the contact surface for impact problems is described in Table I of this report.

The description of material properties for the contact element is described in Section 14. For material properties for other elements in FEAP special input instructions must be supplied.

SUBROUTINE MANUAL

***** USER INSTRUCTIONS AND INPUT FORMATS FOR FEAP-74 *****

FINITE ELEMENT ANALYSIS PROGRAM

***** I N D E X *****

SECTION	WORD	C O N T E N T
1.0	FEAP74	DATA DIRECTORY WORDS
2.0	REMARK	PROBLEM INITIATION AND CONTROL CARDS
2.1	TITLE	REMARK CARD
2.2	STOP	TITLE CHANGE
2.3	MATERI	EXECUTION TERMINATION
3.0	NODAL	MATERIAL CHARACTERIZATION
4.0	GENERA	SEQUENTIAL NODAL GENERATION
4.1	BOUND	NON-SEQUENTIAL NODAL GENERATION
4.2	POLAR	BOUNDARY CODE INPUT
4.3	ELEMEN	POLAR TO CARTESIAN CONVERSION
5.0	BLOCK	ELEMENT GENERATION
5.1	VECTOR	BLOCK MESH GENERATION
6.0	INITIA	USER VECTOR INPUT
6.1	FORCE	INITIAL CONDITION
7.0	ELOADS	GENERALIZED NODAL FORCE
7.1	ELOADS	SURFACE LOADS
7.2	ELOADS	ELEMENT LOADS
7.3	ELOADS	PROPORTIONAL LOADS
8.0	SOLVE	INITIATION OF STATIC SOLUTION
8.1	EXPLICIT	STATIC SOLUTION
8.2	IMPLICIT	EXPLICIT DYNAMIC INTEGRATION
8.3	VISCOE	IMPLICIT INTEGRATION
8.4	MESH	IMPLICIT TIME INTEGRATION
9.0	PLOT	MESH CHECK
9.1	FOURIE	PLOT MESH
9.2	FOURIE	FOURIER SERIES HARMONICS
9.3	FOURIE	OUTPUT CONTROL

FEAP74 IS A GENERAL (F)INITE (E)LEMENT (A)NALYSIS (P)ROGRAM WHICH FURNISHES TO THE USER MESH INPUT/OUTPUT, ELEMENT ASSEMBLY AND SOLUTION OF EQUATIONS (LINEAR, IMPLICIT AND EXPLICIT TIME DEPENDENT, NONLINEAR), PRESCRIBED GENERALIZED NODAL FORCES, PRESCRIBED NODAL AND ELEMENT DATA, AND OUTPUT OF THE GENERALIZED DISPLACEMENTS AND FORCES. ELEMENT MATRICES FOR TWO AND THREE DIMENSIONAL LINEAR ELASTICITY, SHELLS, PLATES, AND FIELD (LAPLACE EQUATION) PROBLEMS ARE AVAILABLE. ALTERNATIVELY USERS MAY SUPPLY THEIR OWN ELEMENT LIBRARY BY PROVIDING A SUBROUTINE CALLED ELEMTH. WHERE NH IS A TWO DIGIT NUMBER (01-10). IDENTIFYING THE ELEMENT SUBROUTINE. EACH ELEMENT SUBROUTINE HAS AT LEAST FOUR BASIC FUNCTIONS WHICH ARE DELINEATED BY A SWITCHING PARAMETER. ISW. IN THE SUBROUTINE.

MAN 1C
MAN 2C
MAN 3C
MAN 4C
MAN 5C
MAN 6C
MAN 7C
MAN 8C
MAN 9C
MAN 10C
MAN 11C
MAN 12C
MAN 13C
MAN 14C
MAN 15C
MAN 16C
MAN 17C
MAN 18C
MAN 19C
MAN 20C
MAN 21C
MAN 22C
MAN 23C
MAN 24C
MAN 25C
MAN 26C
MAN 27C
MAN 28C
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ELMTN(N,MA,NDIM,NDF,NEL,NEL1,NSTF,NSIZV,NVEC,NCT,IM,D,XYZ,
IX,F,FORCE,ESTIF,U,VECT,ISW)

N IS ELEMENT NUMBER.
MA IS THE MATERIAL NUMBER.
NDIM IS SPATIAL DIMENSION, 1-2 OR 3.
NDF IS NUMBER OF DEGREES OF FREEDOM PER NODE.
NEL IS THE NUMBER OF EXTERNAL NODES PER ELEMENT
NEL1 IS DIMENSION OF ELEMENT PROPERTY ARRAY.
NSTF IS THE SIZE OF THE ELEMENT STIFFNESS.
NSIZV IS THE SIZE OF UTILITY VECTORS.
NVEC IS THE NUMBER OF UTILITY VECTORS.
NCT IS A PRINTER LINE COUNTER.
IM IS A PARAMETER FOR MATERIAL IDENTIFICATION.
D(L,1) IS MATERIAL PROPERTY MATRIX (63 CELLS).
XYZ(NDIM,1) ARE NODAL COORDINATES.
IX(NEL1,1) ARE ELEMENT PROPERTIES, NODES, ETC.
F(NDF,1) ARE NODAL GENERALIZED FORCES.
FORCE(NSTF,2) IS ELEMENT FORCE VECTOR TO BE
COMPUTED. COLUMN 2 IS LUMPED MASS
ESTIF(NSTF,NSTF) IS ELEMENT MATRIX TO BE
COMPUTED.
VECT(NSIZV,1) ARE PRESCRIBED NODAL OR ELEMENT
QUANTITIES, TEMPERATURES ETC.
U(NDF,1) IS SOLUTION VECTOR.
ISW IS SWITCHING PARAMETER.

ISW=1. ** MATERIAL CHARACTERIZATION**
ISW=2. ** CHECK ELEMENT FOR POSITIVE AREA *
ISW=3. ** ELEMENT STIFFNESS AND
LOAD COMPUTATION**
ISW=4. ** ELEMENT STRESSES AND PRINTOUT**
ISW=5. ** ELEMENT LOAD COMPUTATION ONLY**
ISW=6. ** NONLINEAR GENERALIZED FORCES**
OTHER ISW MAY BE USED FOR SPECIAL PURPOSES.

USERS CAN GENERATE SURFACE LOADINGS BY PROVIDING SLDNN
SUBROUTINES (WHERE NN IS A TWO DIGIT NUMBER BETWEEN 01 AND 05)
THAT SPECIFY THE LOAD ROUTINE. THE SUBROUTINE IS ACCESSED BY
THE CALL TO

SLDNN(NDIM,NDF,NPRES,IPRES,PP,XYZ,FS)

WHERE IN ADDITION TO QUANTITIES DEFINED ABOVE
FOR ELMTN.

NDF IS THE DIMENSION OF LOADED SURFACE

NPRES IS NUMBER OF LOADED NODES (MAX 8)

IPRES(8) ARE NODE NUMBERS OF LOADED NODES.

PP(8) ARE LOAD VALUES AT CORRESPONDING IPRES
NODES.

FS(8,8) ARE THE COMPUTED GENERALIZED (NODAL)
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MAN 1358C
MAN 1359C
MAN 1360C
MAN 1361C
MAN 1362C
MAN 1363C
MAN 1364C
MAN 1365C
MAN 1366C
MAN 1367C
MAN 1368C
MAN 1369C
MAN 1370C
MAN 1371C
MAN 1372C
MAN 1373C
MAN 1374C
MAN 1375C
MAN 1376C
MAN 1377C
MAN 1378C
MAN 1379C
MAN 1380C
MAN 1381C
MAN 1382C
MAN 1383C
MAN 1384C
MAN 1385C
MAN 1386C
MAN 1387C
MAN 1388C
MAN 1389C
MAN 1390C
MAN 1391C
MAN 1392C
MAN 1393C
MAN 1394C
MAN 1395C
MAN 1396C
MAN 1397C
MAN 1398C
MAN 1399C

SEE SECTION 7.1 FOR DATA INPUT DETAILS.

INTEGRATION TABLE IS ACCESSED BY THE CALL

CALL INTEGL(LIM,NCI,NDIM,LINT,STUW)

STUW(4,M) INTEGRATION POINTS AND WEIGHTS.

NOTE M MUST BE SET EXPLICITLY AND BE LARGER THAN OR EQUAL TO LINT.

LINT - RETURNS WITH NUMBER INTEGRATION POINTS.

NCI = 0 RETURNS GAUSS POINTS AND WEIGHTS IN

STUW.

LIM = 1 TO 5 IS NUMBER OF GAUSS POINTS DIP-
ECTION.

NCI = 1 RETURNS A SPECIAL 3-D GAUSS FORMULA.

SET LIM = 1 FOR 6 PT. CUBIC ACCURACY

SET LIM = 2 FOR 14 PT. QUINTIC ACCURACY.

NCI = 2 RETURNS TRIANGULAR INTEGRATION FORMULA

SET LIM = 1 FOR 1 PT. LINEAR ACCURACY.

SET LIM = 2 FOR 3 PT. QUADRATIC ACCURACY.

SET LIM = 3 FOR 7 PT. QUARTIC ACCURACY.

1.) DATA TYPE IDENTIFICATION CARDS (15,1X,12A6).

EACH DATA SEGMENT IS PRECEDED BY A CARD WHICH IDENTIFIES THE TYPE OF DATA AND LIMITS ON THE AMOUNT OF DATA WHICH IMMEDIATELY FOLLOWS THE CARD. EXCEPT AS NOTED THE DATA SEGMENTS MAY APPEAR IN ANY ORDER. THE IDENTITY CARDS MAY ALSO AID THE USER IN INTERPRETING THE INPUT DATA CARDS. AS SUPPLIED THERE ARE TWENTY-FIVE DIFFERENT DATA IDENTIFICATION CARDS. THESE ARE

COL 7 TO 12 IDENTITY(RESTRICTIONS)

FEAP74 START OF EACH PROBLEM (MUST PRECEDE ALL OTHER DATA).

TITLE CHANGE OUTPUT PAGE HEADINGS

REMARK COMMENTS ON OUTPUT

MATPE1 MATERIAL CHARACTERIZATION.

NODAL NODAL CARDS

POLAR POLAR CONVERSION. (PRECEDE BY NODAL, GENERA. OR BLOCK.)

ELEMEN ELEMENT CONNECTION CARDS.

GENERA GENERATE NODES IN A LINEAR PATH BY ANY INCREMENT

BLOCK GENERATE ALL MESH DATA (BOTH NODAL AND ELEMENT)

FOR 4, 2, OR 3 DIMENSIONAL REGION WHOSE BOUNDARY

MAY BE DEFINED BY 408 OR 8120 COLLOCATED POINTS

BOUNDARY CODE PRESCRIPTION (PRECEDE BY NODAL OR

GENERATE OF BLOCK)

BOUND1A PRECEDED NODAL OR ELEMENT DATA (PRECEDE BY

VECTOR NODAL, OF GENERA AND ELEMEN. OF BLOCK)

FORCE NODAL GENERALIZED FORCES (PRECEDE BY NODAL OR

GENERATE OF BLOCK)

MAN104C
MAN105C
MAN106C
MAN107C
MAN108C
MAN109C
MAN110C
MAN111C
MAN112C
MAN113C
MAN114C
MAN115C
MAN116C
MAN117C
MAN118C
MAN119C
MAN120C
MAN121C
MAN122C
MAN123C
MAN124C
MAN125C
MAN126C
MAN127C
MAN128C
MAN129C
MAN130C
MAN131C
MAN132C
MAN133C
MAN134C
MAN135C
MAN136C
MAN137C
MAN138C
MAN139C
MAN140C
MAN141C
MAN142C
MAN143C
MAN144C
MAN145C
MAN146C
MAN147C
MAN148C
MAN149C
MAN150C
MAN151C
MAN152C
MAN153C
MAN154C
MAN155C
MAN156C
MAN157C

C BLOADS SURFACE LOADINGS (SAME AS FORCE).
 C ELOADS ELEMENT LOADINGS (SAME AS FORCE).
 C MESH CHECK CONSISTENCY OF MESH ONLY (SAME AS SOLVE)
 C PLOT PLOT MESH (SAME AS SOLVE)
 C INITI INITIAL CONDITION PRESCRIPTION FOR DYNAMIC
 C SOLVE ANALYSIS (PRECEDS BY NODAL, GENERAL OR BLOCK)
 C COMPLETE COMPLETE FORMULATION AND SOLUTION FROM ELEMENTS
 C (PRECEDS BY MATERIAL, NODAL OR GENERAL, AND ELEMENT
 C OR PRECEDES BY MATERIAL AND BLOCK)
 C RESOLVE USE PREVIOUS PROBLEM DESCRIPTION WITH NEW LOAD
 C ONLY (PRECEDS BY SOLVE AND NEW LOADING CARDS).
 C EXPLICIT DYNAMIC SOLUTION BY EXPLICIT INTEGRATION. (SAME
 C AS SOLVE)
 C IMPLICIT IMPLICIT INTEGRATION OF DYNAMIC PROBLEMS
 C (PRECEDS BY SAME DATA AS FOR SOLVE)
 C VISCOE QUASI-STATIC LINEAR VISCOELASTIC INTEGRATION
 C (PRECEDS BY SAME DATA AS FOR SOLVE)
 C FOURIE FOURIER COMPOSITION (SAME AS SOLVE)
 C ADDUP ACCUMULATE FOURIER SOLUTION (AFTER FOURIE)
 C STOP NORMAL EXIT (MUST FOLLOW ALL DATA)

NOTE EACH IDENTIFIER IS PUNCHED STARTING IN COL 7 (LEFT
 JUSTIFIED).

EXCESS CARDS MAY EXIST BETWEEN EACH SECTION OF DATA. HOWEVER,
 THE DATA TO BE USED MUST IMMEDIATELY FOLLOW THE TYPE CARD AND
 MUST BE IN PROPER ORDER. NO PARTICULAR ORDER OF THE TYPE
 CARDS IS NECESSARY EXCEPT THAT THE FEAP74 CARD MUST ALWAYS BE
 THE FIRST CARD IN EACH SET OF DATA. AND RESTRICTIONS MUST BE
 OBSERVED.

2.) PROBLEM INITIATION AND CONTROL CARDS

CARD 1. (6X,12A6)

COL 7 TO 12 MUST CONTAIN WORD FEAP74
 COL 13 TO 78 OUTPUT PAGE HEADER

CARD 2. (15,1X,3A6)

COL 1 TO 5 NDIM - SPATIAL DIMENSION OF PROBLEM (1 TO 3)
 COL 7 TO 12 NAMES TO BE PRINTED AS OUTPUT HEADERS TO
 COL 13 TO 18 COORDINATES - IF BLANK SET TO 1.2.3 AS NEEDED.
 COL 19 TO 24

CARD 3. (15,1X,6A6)

COL 1 TO 5 NPF - NUMBER OF UNKNOWN PER NODE (1 TO 6)
 COL 7 TO 12 NAMES TO BE PRINTED AS OUTPUT HEADERS OF THE
 COL 13 TO 18 GENERALIZED DISPLACEMENTS AND FORCES - IF
 BLANK SET TO 1.2.3.4.5.6 AS NECESSARY
 COL 19 TO 24

MAN158C
 MAN159C
 MAN160C
 MAN161C
 MAN162C
 MAN163C
 MAN164C
 MAN165C
 MAN166C
 MAN167C
 MAN168C
 MAN169C
 MAN170C
 MAN171C
 MAN172C
 MAN173C
 MAN174C
 MAN175C
 MAN176C
 MAN177C
 MAN178C
 MAN179C
 MAN180C
 MAN181C
 MAN182C
 MAN183C
 MAN184C
 MAN185C
 MAN186C
 MAN187C
 MAN188C
 MAN189C
 MAN190C
 MAN191C
 MAN192C
 MAN193C
 MAN194C
 MAN195C
 MAN196C
 MAN197C
 MAN198C
 MAN199C
 MAN200C
 MAN201C
 MAN202C
 MAN203C
 MAN204C
 MAN205C
 MAN206C
 MAN207C
 MAN208C
 MAN209C
 MAN210C
 MAN211C

CARD 4. (615.5F10.0)

COL 1 TO 5 NEN - MAXIMUM NUMBER OF NOIES CONNECTED TO ANY ELEMENT (1 TO 20).
 COL 6 TO 10 NEXTPA - INCREASES ELEMENT MATRIX SIZE FROM NDF*NEN TO NDF*NEN + NEXTPA
 COL 11 TO 15 IREC - COMPUTE GENERALIZED FORCE CHECK IF NONZERO (FOR TIME INVARIANT ANALYSIS ONLY)
 COL 16 TO 20 MBAN - MAXIMUM EXPECTED BANDWIDTH, DEFAULT IS SET TO 100. USED AS AN ERROR CHECK TO PREVENT RUNNING WITH AN OBVIOUS ERROR.
 COL 21 TO 25 IBUF - BUFFER SIZE FOR STORAGE OF HISTORY EFFECTS IN TIME DEPENDENT ANALYSIS. DEFAULT IS IBUF = 1520720
 COL 26 TO 30 NC1 - USER INTEGER CONSTANT
 COL 31 TO 40 CON1 - USER DEFINED CONSTANT
 COL 41 TO 50 CON2 - USER DEFINED CONSTANT
 COL 51 TO 60 CON3(1) - USER DEFINED CONSTANT
 COL 61 TO 70 CON3(2) - USER DEFINED CONSTANT
 COL 71 TO 80 CON3(3) - USER DEFINED CONSTANT

2.1) REMARK - USER COMMENTS ON OUTPUT. (6X.12A6)

SUBSEQUENT CARDS

COL 7 TO 12 MUST CONTAIN REMARK
 COL 13 TO 78 STATEMENTS TO BE OUTPUT. USE AS MANY REMARK CARDS AS DESIRED. INSERT BEFORE ANY TYPE CARD.

2.2) TITLE CHANGE ON OUTPUT (6X.12A6)

COL 7 TO 12 MUST CONTAIN TITLE
 COL 13 TO 78 NEW TITLE DESCRIPTOR

2.3) EXECUTION TERMINATION (6X.64)

COL 7 TO 10 MUST CONTAIN STOP. INSERT AFTER LAST PROBLEM.

3.) MATERIAL CHARACTERIZATION (15.1X.12A6)

COL 1 TO 5 NUMMAT - NUMBER OF DIFFERENT MATERIAL CHARACTERIZATIONS TO FOLLOW.
 COL 7 TO 12 MUST CONTAIN WORD MATEP1

THE FOLLOWING CARDS ARE SUPPLIED FOR EACH MATERIAL TO BE CHARACTERIZED. MUST BE EXACTLY NUMMAT SETS OF CARDS)

CARD 1. ELEMENT SELECTOR CARD (15.1X.45.11A6)

COL 1 TO 5 MATERIAL NUMBER (1 TO NUMMAT)
 COL 7 TO 12 ELEMENT CLASS (01 TO 10) TO WHICH THE CHARACTERIZATION BELONGS.
 COL 13 TO 78 ELEMENT IDENTIFICATION INFORMATION TO BE OUTPUT.

MAN212C
 MAN213C
 MAN214C
 MAN215C
 MAN216C
 MAN217C
 MAN218C
 MAN219C
 MAN220C
 MAN221C
 MAN222C
 MAN223C
 MAN224C
 MAN225C
 MAN226C
 MAN227C
 MAN228C
 MAN229C
 MAN230C
 MAN231C
 MAN232C
 MAN233C
 MAN234C
 MAN235C
 MAN236C
 MAN237C
 MAN238C
 MAN239C
 MAN240C
 MAN241C
 MAN242C
 MAN243C
 MAN244C
 MAN245C
 MAN246C
 MAN247C
 MAN248C
 MAN249C
 MAN250C
 MAN251C
 MAN252C
 MAN253C
 MAN254C
 MAN255C
 MAN256C
 MAN257C
 MAN258C
 MAN259C
 MAN260C
 MAN261C
 MAN262C
 MAN263C
 MAN264C
 MAN265C

CARD 2, 1, ETC. ** USER DEFINED FOR EACH ELEMENT TYPE PROVIDED.

EXCESS BLANK CARDS ARE PERMISSIBLE BETWEEN EACH MATERIAL SET.

4.0) NODAL CARDS (15.1X.45)

COL 1 TO 5 NUNIP - NUMBER OF NODAL POINTS
COL 7 TO 12 MUST CONTAIN NODAL

SUBSEQUENT CARDS LAST NODAL CARD MUST NOT BE GENERATED.
+15.115.3F10.0)

COL 1 TO 5 NODE NUMBER
COL 15 1 IF 1 DISPLACEMENT IS SPECIFIED
COL 16 1 IF 2 DISPLACEMENT IS SPECIFIED
COL 17 1 IF 3 DISPLACEMENT IS SPECIFIED
COL 18 1 IF 4 DISPLACEMENT IS SPECIFIED
COL 19 1 IF 5 DISPLACEMENT IS SPECIFIED
COL 20 1 IF 6 DISPLACEMENT IS SPECIFIED
COL 21 TO 30 1 COORDINATE VALUE
COL 31 TO 40 2 COORDINATE VALUE * AS REQUIRED
COL 41 TO 50 3 COORDINATE VALUE

NODAL CARDS MUST BE IN ORDER. MISSING NODES ARE INTERPOLATED LINEARLY FROM INPUT NODES. IF SUCCEEDING CARDS HAVE IDENTICAL BOUNDARY CODES, THIS BOUNDARY CODE WILL BE ASSIGNED TO THE INTERVENING NODES. IN ALL OTHER CASES THE BOUNDARY CODE IS SET TO ZERO *TERMINATE ON NODE NUNIP OR A BLANK CARD*

4.1) NON SEQUENTIAL NODAL GENERATOR OPTION. (15.1X.46)

COL 1 TO 5 NUMBER OF NODAL POINTS
COL 7 TO 12 MUST CONTAIN GENERA

SUBSEQUENT CARDS (215.110.3F10.0)

COL 1 TO 5 NODE-NUMBER
COL 6 TO 10 NODE-NUMBER-INCREMENT WHICH WILL BE SUCCESSIVELY ADDED TO NODE-NUMBER UNTIL SUM IS GREATER THAN NODE-NUMBER ON FOLLOWING CARD (ALGEBRAIC).
COL 15 TO 20 BOUNDARY CODE. SAME AS INPUT FOR NODAL.
IF SUCCEEDING CARDS HAVE IDENTICAL BOUNDARY CODES, THIS BOUNDARY CODE WILL BE ASSIGNED TO THE INTERVENING NODES. IN ALL OTHER CASES THE BOUNDARY CODE IS SET TO ZERO.

COL 21 TO 30 1 COORDINATE VALUE *
COL 31 TO 40 2 COORDINATE VALUE * AS REQUIRED +
COL 41 TO 50 3 COORDINATE VALUE +

*TERMINATE WITH BLANK CARD *

4.2) BOUNDARY CODE CATCHING OPTION. (15.1X.47)

MAN2660
MAN2670
MAN2680
MAN2690
MAN2700
MAN2710
MAN2720
MAN2730
MAN2740
MAN2750
MAN2760
MAN2770
MAN2780
MAN2790
MAN2800
MAN2810
MAN2820
MAN2830
MAN2840
MAN2850
MAN2860
MAN2870
MAN2880
MAN2890
MAN2900
MAN2910
MAN2920
MAN2930
MAN2940
MAN2950
MAN2960
MAN2970
MAN2980
MAN2990
MAN3000
MAN3010
MAN3020
MAN3030
MAN3040
MAN3050
MAN3060
MAN3070
MAN3080
MAN3090
MAN3100
MAN3110
MAN3120
MAN3130
MAN3140
MAN3150
MAN3160
MAN3170

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C      COL 7 TO 12  MUST CONTAIN BOUNDA
C      SUBSEQUENT CARDS. (815)
C
C      COL 1 TO 5  N. NODE NUMBER TO HAVE REDEFINED BOUNDARY CODE.
C      COL 6 TO 10 N%. GENERATOR INCREMENT TO BE ADDED ALGEBRAICALLY
C      CALLY TO N. UNTIL SUM EXCEEDS (MAX OR MIN) THE
C      N OF THE FOLLOWING CARD.
C
C      COL 11 TO 15 IBC(I), (I=1,2,...,NDF) CODE FOR SPECIFYING FORCE
C      COL 16 TO 20 OR DISPLACEMENT BOUNDARY CONDITIONS.
C      COL ...
C      IBC(I) .EQ. 0. FORCE SPECIFIED.
C      IBC(I) .GT. 0. DISPLACEMENT SPECIFIED. NO
C      INTERVENING GENERATION.
C      IBC(I) .LT. 0. DISPLACEMENT SPECIFIED.
C      GENERATE BETWEEN MISSING NODES IN ALGEBRAIC
C      INCREMENTS OF NX.
C
C      * TERMINATE WITH A BLANK CARD. *
C
C      4.3) POLAR OR CYLINDRICAL COORDINATE CONVERSION TO CARTESIAN
C      COORDINATES (6X,A6)
C
C      COL 7 TO 12  MUST CONTAIN POLAR (LEFT JUSTIFIED)
C      CARD 1. (315.5%,2F10.0)
C
C      COL 1 TO 5  N1. FIRST NODE TO BE CONVERTED
C      COL 6 TO 10 N2. LAST NODE TO BE CONVERTED
C      COL 11 TO 15 N3. INCREMENT ADDED (ALGEBRAICALLY). N1 TO N2
C      COL 21 TO 30 X0. ORIGIN OF POLAR X-COORDINATE
C      COL 31 TO 40 Y0. ORIGIN OF POLAR Y-COORDINATE
C
C      5.) ELEMENT CARDS (15.1X,A6)
C
C      COL 1 TO 5  NUMEL - NUMBER OF ELEMENTS
C      COL 7 TO 12 MUST CONTAIN ELEMEN
C
C      SUBSEQUENT CARDS (415.2013/2014)
C      CARD 1.
C
C      COL 1 TO 5  ELEMENT NUMBER
C      COL 6 TO 10 MATERIAL NUMBER
C      COL 11 TO 15 NUMBER OF SUBSEQUENT ELEMENTS USING SAME
C      STIFFNESS MATRIX * SAVES RECOMPUTATION OF
C      SIMILAR MATRICES. ELEMENT MUST ALSO HAVE
C      SAME ELEMENT FORCE VECTOR * IF THESE ARE
C      IN THE STIFFNESS SUBROUTINE *
C      PRINT ELEMENT MATRIX IF NONZERO.
C      INDIC1 ELEMENT INCREMENT ARRAY ON NODE 1.
C      INDIC2 * IF NOT INPUT IS SET AUTOMATICALLY
C      UP TO
C      INDIC3 1 TO 60 FOR SEPERATING ELEMENTS * SEE REPORT

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MAN318C
MAN319C
MAN320C
MAN321C
MAN322C
MAN323C
MAN324C
MAN325C
MAN326C
MAN327C
MAN328C
MAN329C
MAN330C
MAN331C
MAN332C
MAN333C
MAN334C
MAN335C
MAN336C
MAN337C
MAN338C
MAN339C
MAN340C
MAN341C
MAN342C
MAN343C
MAN344C
MAN345C
MAN346C
MAN347C
MAN348C
MAN349C
MAN350C
MAN351C
MAN352C
MAN353C
MAN354C
MAN355C
MAN356C
MAN357C
MAN358C
MAN359C
MAN360C
MAN361C
MAN362C
MAN363C
MAN364C
MAN365C
MAN366C
MAN367C
MAN368C
MAN369C
MAN370C
MAN371C

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RUN FORTRAN COMPILER VERSION 2.3 B.3

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C      COL 21 TO 30  2-COORDINATE OF BOUNDARY-DEFINING-POINT.
C      COL 31 TO 40  3-COORDINATE OF BOUNDARY-DEFINING-POINT.
C
C      NOTES:  1.  BLOCK GENERATES ONLY 4 PT. QUADRIATERALS OR 8 PT.
C              BRICKS.
C              2.  INPUT OF CARDS 3. FOLLOW ORDER RULES FOR ELEMENT
C              INPUT (SEE SECTION 5.1).
C              3.  R-S-T ARE LOCAL COORDINATES.
C              I.E. (-1 .LE. R.S.T .LE. 1).  WHERE R IS DIRECTED
C              FROM NODE 1 TO 2. S IS IN PLANE OF FIRST THREE NODES
C              AND T IS NORMAL TO R-S PLANE.
C
C      6.) VECTOR CARDS, I.E. USER DEFINED INPUT (15.1X.A6)
C
C      COL 1 TO 5  NVEC. NUMBER OF DIFFERENT VECTORS (7 MAX)
C      COL 7 TO 12 MUST CONTAIN VECTOR
C
C      SUBSEQUENT CARDS
C
C      CARD 1.  (215)
C
C      COL 1 TO 5  NS12V. VECTOR LENGTH.COMMON TO ALL NVEC VECTORS
C      COL 6 TO 10 IPICK. CODED PARAMETER.
C
C              IPICK = 0. VECTORS ASSOCIATED WITH NODES
C              IPICK = 1. VECTORS ASSOCIATED WITH DEG.FREEDOM
C              IPICK = 2. VECTORS ASSOCIATED WITH ELEMENTS
C
C      CARD 2.  (6X. 2A6) REPEAT NVEC TIMES
C
C      COL 7 TO 13  DESCRIPTIVE TITLE FOR VECTOR
C
C      CARD 3.  (215.7F10.0)
C
C      COL 1 TO 5  POSITION NUMBER OF VECTOR ELEMENT. 1 TO NS12V
C      COL 6 TO 10 GENERATOR INCREMENT
C      COL 11 TO 20 VECTOR ELEMENT VALUE OF VECTOR 1
C      COL 21 TO 30 VECTOR ELEMENT VALUE OF VECTOR 2
C      COL ..... AS REQUIRED FOR NVEC VECTORS
C
C      LINEAR INTERPOLATION IS PERFORMED ON ALL VECTORS BETWEEN
C      NON-CONSECUTIVE POSITION NUMBERS SPECIFIED IN COL 1 TO 5
C      IF INCREMENT IS NONZERO.
C      IF DESCRIPTIVE TITLES OF ALL VECTORS ARE BLANK CARDS.
C      PRINTING OF THE VECTOR VALUES IS SUPPRESSED.
C
C      * TERMINATE ON BLANK CARD *
C
C      5.1. INITIAL CONDITIONS FOR TIME DEPENDENT ANALYSIS.
C
C      COL 1 TO 5  WITH NUMBER OF INITIAL CONDITION VECTORS
C      COL 7 TO 12 MUST CONTAIN INITIAL
C
C      15.1X.A6

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MAN427C
MAN428C
MAN429C
MAN430C
MAN431C
MAN432C
MAN433C
MAN433C
MAN433C
MAN433C
MAN435C
MAN436C
MAN437C
MAN438C
MAN439C
MAN440C
MAN441C
MAN442C
MAN443C
MAN444C
MAN445C
MAN446C
MAN447C
MAN448C
MAN449C
MAN450C
MAN451C
MAN452C
MAN453C
MAN454C
MAN455C
MAN456C
MAN457C
MAN458C
MAN459C
MAN460C
MAN461C
MAN462C
MAN463C
MAN464C
MAN465C
MAN466C
MAN467C
MAN468C
MAN469C
MAN470C
MAN471C
MAN472C
MAN473C
MAN474C
MAN475C
MAN476C
MAN477C

PROP = A8*(SIN(W1*T))**K + A2*(COS(W3*T))**K + W5

LOAD TYPE 3.

PPROP = USER DEFINED FUNCTION FROM SUBROUTINE
EXPRLD(PPROP,T,A) WHERE A IS AN ARRAY WITH
INFORMATION FOR COLUMNS 6-80 OF DATA CARDS.

***NOTE** PROPORTIONAL LOADS CAN BE ACCUMULATED FROM DIFFERENT
TYPES AT THE SAME TIME.

8.) INITIATION OF TIME INDEPENDENT SOLUTION (15.12.66)

COL 1 TO 5 IOUT. OUTPUT CONTROL CODE.

IOUT .EQ. 0. ALL STRESSES AND DISP. PRINTED
IOUT .NE. 0. SELECTED PRINTOUT. MORE DATA INPUT
SEE SECTION 9 FOR DATA PREPARATION.

COL 7 TO 12 MUST CONTAIN SOLVE *INDICATES ALL DATA INPUT*
COMPLETE FORMULATION AND SOLUTION OF EQUATIONS.
COL 7 TO 12 MUST CONTAIN RESOLV TO OBTAIN SUBSEQUENT
SOLUTIONS WHERE BOUNDARY CODES DO NOT CHANGE
AND ALL PRESCRIBED DISPLACEMENTS ARE ZERO.

8.1) INITIATION OF DYNAMIC SOLUTION BY EXPLICIT INTEGRATION.

COL 1 TO 5 IPR. OUTPUT CONTROL FOR DISPLACEMENT AND
STRESS PRINTOUT. SEE SECT. 9 FOR DATA INPUT.
IOUT = 1 - MIN(1,IPR)

IF IOUT .NE. 0. THE SPATIAL CONTROL DATA
COMES AT THE END OF THE DYNAMIC SEGMENT.
COL 7 TO 12 MUST CONTAIN EXPLIC

SUBSEQUENT CARDS (215.2F10.0.215)

COL 1 TO 5 NUMBER OF TIME STEPS

COL 6 TO 10 PRINT INTERVAL

COL 11 TO 20 TIME INCREMENT

COL 21 TO 30 NEWMARK DELTA-DAMPING TERM (GAMMA - .5)

COL 31 TO 35 NUMBER OF TIME EVOLUTION ELEMENT VARIABLE PLOTS

COL 36 TO 40 NPROP. NUMBER OF PROPORTIONAL LOADS TO BE INPUT

COL 41 TO 45 NFORC. LAST NODE ON WHICH A FORCE IS CHANGED
DURING EACH TIME STEP.

COL 46 TO 50 KKK. STABILITY CHECK OVERRIDE ** CAUTION USE
ONLY WHEN A BETTER ESTIMATE OF THE STABLE TIME

STEP IS AVAILABLE THAN CAN BE PERFORMED BY CODE

*** ZERO. USES INTERNAL STABILITY CHECK.

*** NONZERO. DISREGARDS STABILITY CHECK.

COL 51 TO 55 NOKL. 0 FOR LINEAR

1 FOR NON LINEAR

COL 56 TO 60 LPE. 0 PRINTS LOADS

1 SUPPRESS PRINT

MAN550C
MAN557C
MAN558C
MAN559C
MAN560C
MAN561C
MAN562C
MAN563C
MAN564C
MAN565C
MAN566C
MAN567C
MAN568C
MAN569C
MAN570C
MAN571C
MAN572C
MAN573C
MAN574C
MAN575C
MAN576C
MAN577C
MAN578C
MAN579C
MAN580C
MAN581C
MAN582C
MAN583C
MAN584C
MAN585C
MAN586C
MAN587C
MAN588C
MAN589C
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MAN592C
MAN593C
MAN594C
MAN595C
MAN596C
MAN597C
MAN598C
MAN599C
MAN600C
MAN601C
MAN602C
MAN603C
MAN604C
MAN605C
MAN606C
MAN607C
MAN608C
MAN609C
MAN610C
MAN611C
MAN612C
MAN613C
MAN614C
MAN615C
MAN616C
MAN617C
MAN618C
MAN619C
MAN620C
MAN621C
MAN622C
MAN623C
MAN624C
MAN625C
MAN626C
MAN627C
MAN628C
MAN629C
MAN630C
MAN631C
MAN632C
MAN633C
MAN634C
MAN635C
MAN636C
MAN637C
MAN638C
MAN639C
MAN640C

SUBSEQUENT CARDS (315) ONE FOR EACH STRESS PLOT.

COL 1 TO 5 ELEMENT NUMBER CONTAINING STRESS TO BE PLOTTED.
COL 6 TO 10 LOCAL COORDINATE POINT CODE, 1 TO 9, AS
PATTERNED AFTER, COL 11 TO COL 19, IN SECT. 9.
COL 11 TO 15 PLOT COMPONENT CODE, 1 TO 6 FOR SIGMA(1,J),
I.E., SIGMA(1,1)=1, SIGMA(1,2)=2, SIGMA(1,3)=3,
SIGMA(2,2) = 4, SIGMA(2,3) = 5, SIGMA(3,3) = 6.

IF(NPROP,NE,0) READ PROPORTIONAL LOAD CARDS, SEE SECT. 7.3

IF(NFORC,NE,0) READ FORCE CARDS AT EACH TIME STEP. IF OUTPUT IS
LIMITED BY IOUT NONZERO, THE FIRST FORCE CARD SET PRECEDES
OUTPUT CARDS AND THE REMAINDER FOLLOW THE OUTPUT CARDS NO BLANK
CARDS MAY BE USED BETWEEN SETS OF CARDS OTHER THAN THE USUAL
BLANK TERMINATOR CARD FOR FORCE INPUT CARDS.

IF(IOUT,NE,0) DATA FOR SPATIAL PRINTOUT CONTROL. SEE SECT.9.

SPECIAL COMMENTS FOR DYNAMIC OPTION

- (1) ONLY COLUMNS 1 TO 66 ARE AVAILABLE FOR PAGE HEADING.
- (2) MAXIMUM ADVANTAGE OF ELEMENT REUSE OPTION SHOULD BE TAKEN.
- (3) INITIAL CONDITIONS FOR DISPLACEMENT AND VELOCITY VECTORS,
AS WELL AS STORAGE FOR ACCELERATION VECTOR, MAY BE MADE
THROUGH INPUT OF AN INITIAL CONDITION CARD SET, WITHOUT
SPECIFIED INITIAL CONDITIONS THEY ARE AUTOMATICALLY SET ZERO.
- (4) SPATIAL LOADING IS INPUT THROUGH FORCE OR BOUNDARY
PRESSURE CARDS.
- (5) EXTREME CAUTION ON ORDER OF DATA CARDS MUST BE OBSERVED. NO
EXTRA CARDS ARE PERMITTED AND STRICT COUNTS ARE OBSERVED
EXCEPT FOR THE NUMBER OF FORCE CARDS USED IN EACH TIME STEP.

8.2) INITIATION OF IMPLICIT TIME INTEGRATIONS (15.1X.A6)

COL 1 TO 5 NSEQ, NUMBER OF TIME SEQUENCES
COL 7 TO 12 MUST CONTAIN VISCOE FOR LINEAR VISCOELASTIC
QUASI-STATIC PROBLEMS (ONE INITIAL CONDITION
ONLY MUST BE USED)
COL 7 TO 12 MUST CONTAIN IMPLIC FOR DYNAMIC IMPLICIT
INTEGRATIONS (THREE INITIAL CONDITIONS ARE
REQUIRED, MORE CAN BE SPECIFIED WITHOUT ERROR)

SUBSEQUENT CARDS. ONE SET FOR EACH TIME SEQUENCE

CARD 1. F10.0.815.2F10.0.215)

COL 1 TO 10 DT, TIME INCREMENT (NONZERO FOR IMPLIC)
COL 11 TO 15 NTS, NUMBER OF TIME STEPS IN SEQUENCE
COL 16 TO 20 INT, PRINT INTERVAL (DEFAULT 1)
COL 21 TO 25 NNT, FIRST NODE PRINTED
COL 26 TO 30 NNE, LAST NODE PRINTED
COL 31 TO 35 NEI, FIRST ELEMENT STRESS TO BE PRINTED

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MAN698C

C. Listing of the Contact/Impact Subroutines Added to FEAP 74

The listings for the subroutines which are added to FEAP 74 for the contact/impact theory described herein are given below.


```

451      UDD(IDEG,J) = CU
J454      300      CONTINUE

000457      RETURN

000457      1000      FORMAT(15.5X,2F10.0)
000457      1001      FORMAT(215)
000457      2000      FORMAT(A1,12A6,30X,4HPAGE,14//5X,25HCONTACT SHEET DESCRIPTION//
C          10X,13HBODY 1 RO*U =,E13.5//
C          10X,13HBODY 2 RO*U =,E13.5//
X          11X,7HELEMENT,1X,9HDIRECTION (15,110))
000457      2001      FORMAT(110,6E13.4)
000457      2002      FORMAT(10A9H ELEMENT,13H BODY 1 MASS,13H BODY 2 MASS)
000457      END
```



```

000256      PLOT(7) = UDL(IDEG,3)
000260      PLOT(8) = UDL(IDEG,3)
000262      DO 60 K = 1,NUMPLT
000263      KK = NPLT(K,2)
000265      IF (NPLT(K,1).GT.0) CALL PLDATA(NDIM,NPLT(K,1),THED(KK),X(I,2))
000306      C      PLOT(KK)
000311      C      CONTINUE
000312      C      RETURN
000312      C      FORMAT(215)
000312      C      FORMAT(5X,21HP0INT CONTACT ELEMENT
000312      C      5X,27HC0NTACT DEGREE OF FREEDOM = 13.5X,5HMP5 = 15)
000312      C      FORMAT(5X,7HELEMENT,15.5X,A5,5X,8HMATERIAL,13.5X,5HTRU =,E15.5,5X,
000312      X      SHETA =,F3.1)
000312      END

```



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ELM 590ELM 600
ELM 610
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ELM 750
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ELM 810
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ELM 830
ELM 840ELM 870
ELM 880

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000573 615 FORCE(KK+IU,1) = FORCE(IY+U,1) - AR*MOD(UEK,1) - BB*ST(UE,1)
000615 600 IU = IU + NDF
000621 400 CONTINUE
000624 5 RETURN
000635 1000 FORMAT(3F10.0)
000635 2000 FORMAT(14I12,4E13.5,17F4HPAGE,13F10H ELEMENT,10F20H)
000625 X,3X,5F20.0)
000625 2001 FORMAT(110,15,5X,45,5E12.4)
000625 2100 FORMAT(13HOLINEAR ELASTIC MATERIAL, BAR ELEMENT//
5,17HE =,E15.5,5X,6HAREH =,E15.5,5X,9HDENSITY =,E15.5,/)
000635 4000 FORMAT(45,15,1F3E12.5,24X,110)
000625 END

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ELM 310

ELM 320

ELM 360

ELM 370

ELM 380

ELM 390


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000017 SUBROUTINE HLNORM( ISW, MDEG, NN, U, DU, DF, F, PROP, NCT, NTT, IDEST, IBLK )
000017 C***** NLNORM ***** 12/14/73 *****
000017 DIMENSION U(1), DU(1), DF(1), F(1), IDEST(1)
000017 COMMON NORMS UNORM, DNORM, ANORM, CS, CSP, DP, DNP, IFL, XF, AG, EFF
000017 GO TO (100,200,200,100), ISW
000026 .00 CS=1.0
000027 CSP=1.0
000030 DP=0.75
000032 IFL=0
000032 XFLAG=.FALSE.
000033 ERR=1.0E-3
000035 RETURN
000036 UNORM = 0.
000037 DNORM = 0.
000037 ANORM=0.
000041 DO 500 N = 1, MDEG
000042 DUN = DU(N)
000044 UN = U(N)
000046 IF (ISW.EQ.3) UN = UN + DF(N) + DUN
000053 ANOR1=ANOR1+UN*DUN
000056 UNORM=UNORM+UN*DUN
000060 DNORM=DNORM+DUN*DUN
000065 UNORM=SQRT(UNORM)
000067 DNORM=SQRT(DNORM)
000070 IF (NN.EQ.1) DNP=UNORM
000077 CS=1.0
000100 IF (UNORM*DNORM.NE.0.) CS=ANORM/(UNORM*DNORM)
000105 IF (IFL.EQ.1) DP=0.75
000111 IF (IFL.EQ.1.AND.XFLAG) CSP=CS
000120 IFL=0
000121 IF (DNORM.LE.0.5*UNORM) GO TO 550
000124 IFL=1
000125 DP=0.5*UNORM/DNORM
000127 IF (DP*DNORM.GT.DNP) DP=[DNP/DNORM]
000134 IF (DF.EQ.0.0) DF = 1.0
000136 IF (CS*CSP.LT.0.) DP=DP/2.
000142 IF (CS*CSP.GT.0.) IF=1.25*IF
000146 DNP=DP*DNORM
000150 CSP=CS
000150 IF (ISW.EQ.2) RETURN
000153 DO 600 N = 1, MDEG
000155 DF(N) = DF(N) + DUN
000160 IF (NCT.EQ.0) DF(N) = 1.0/N
000164 DUN = DF(N)
000171 PERM=
000171 ENR

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 NLN 430
 NLN 440
 NLN 450


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000152 DTP = 0.
000153 TIME = 0.
000154 MB = MAXBAN + 1
000156 NEI = NEN + 1
C..... INITIALIZE INCREMENTAL FORCE VECTOR
000157 DO 6 N = 1, NSICD
000160 DF(N) = 0.0
000165 DO 10 N = 1, NUMEL
000166 IX(NEI,N) = 0
C..... PUT INITIAL DATA ON TAPE
000177 ISZH = IBUF
000201 ITRD = ITP13
000202 ITR = ITP14
000204 REWIND ITRD
000206 IF(IBLK.GT.0) WRITE(ITRD) ((U(I,J), I=1,NDF), J=1,KU)
000241 DO 11 N = 1, ISZH
000243 H(N) = 0.
000247 NTB = (NH+ISZH-1)/ISZH
000253 WRITE(ITRD) H
000264 IF(NTB.LE.1) GO TO 13
000272 NT = NTB*2
000273 DO 12 N = 1, NT
000275 WRITE(ITRD) H
000314 REWIND ITRD
000316 IF(IBLK.GT.0) READ(ITRD) DM
000331 NEP = NUMNP + NUMNP
000332 C6 = 0.0
000333 NST = 0
000334 NST = 0
000335 NUMPLT = 0
000336 PROP = 0.
000337 CFLAG = .FALSE.
000340 DO 900 M = 1, NSEQ
000341 READ(ITP5,1000) DT,NTS,INT,NNI,NNF,NEI,NEF,NPROP,NFORC,C0,DM,NT,NL
000376 IF(NICD.GT.1) C6 = DM
000406 IF(NL.NE.0) NST = NL
000410 NTT = IABS(NST)
000412 DFLAG = .FALSE.
000413 IF (NTT.NE.0) DFLAG = .TRUE.
000415 NSTEP = NSTEPAINTS
000417 IF(M.EQ.1) NUMPLT = NT
000422 IF(INT.LE.0) INT = 1
000425 WRITE(ITP6,2000) 0,HEAD,TIME,IPG,DT,NTS,INT,NNI,NNF,NEI,NEF
000457 IF(DFLAG) WRITE(ITP6,2032) NST
000476 IPG = IPG + 1
000500 IF(NICD.EQ.1) GO TO 1
000501 IF(DT.EQ.0.0) GO TO 901
000502 CALL UPDATE(3)
000504 IF(NPROP.GT.0) PROP = PROPLD(TIME,NPROP)
000517 IF(NFORC.GT.0) AND,NPROP,DT,0) WRITE(ITP5,4031)
C..... READ STRESS PLOT INFORMATION
000536 IF(NUMPLT.LE.0) OR,M.GT.1) GO TO 500
000541 REWIND 12
000542 WRITE(ITP6,2005)

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TSO 44C
TSO 45C
TSO 46C
TSO 47C
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TSO 94C
TSO 95C
TSO 96C

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555      DO 501 N=1,NUMPLT
556      READ(ITP5,1006) (NEDATA(N,I),I=1,3)
501      WRITE(ITP6,2006) N,(NEDATA(N,I),I=1,3)
500      CALL PLZERO
500      VFLAG = .FALSC.
500      DO 500 NT = 1,NTS
500      PPROP = PROP
500      PROP = 1.0
500      IF(CDFLAG)CALL NLNORM(4,MDEG,NT,U,DF,DU,F,PPROP,NCT,NTT,IDEST,IBLK)
500      IF(NPROP.GT.0) PROP = PROPLD(TIME+DT,0)
500      IF(NFORC.GT.0) CALL RESET(-NFORC,NUMP,NDF,F)
500      NCT = 0
500      REWIND ITLR
500      IF(1BLK.EQ.0) GO TO 15
500      IF(.NOT.VFLAG) GO TO 20
500      DO 21 I = 1,NDEG
500      DF(I) = 0.
500      GO TO 20
500      DO 19 N = 1,NUMP
500      DO 18 K = 1,NDF
500      J = IDEST(K,N)
500      IF(J.EQ.0) GO TO 18
500      IF(VFLAG) GO TO 17
500      DO 16 I = 1,MAXBAN
500      A(J,I) = 0.0
500      DF(J) = F(K,N)*PROP
500      CONTINUE
500      CONTINUE
500      NH = 1
500      IF(NTB.GT.1) READ(ITRD) H
500      IF(NCT.GT.0) GO TO 44
500      OUTPUT THE SOLUTION VECTOR FOR THE CURRENT TIME
500      IF(NH1.EQ.0.OR.(NT-1)/INT)*INT.NE.NT-1) GO TO 40
500      NCT = 0
500      DO 30 N = NNI,NNF
500      NCT = NCT - 1
500      IF(NCT.GT.0) GO TO 31
500      IF(NICD.EQ.1) WRITE(ITP6,2001) 0,HEAD,TIME,IPG,PPROP,M,NT,
500      X (XHED(I),XH,I=1,NDIM),(UHED(I),UH,I=1,NDF)
500      IF(NICD.NE.1) WRITE(ITP6,2001) 0,HEAD,TIME,IPG,PPROP,M,NT,
500      X (XHED(I),XH,I=1,NDIM),(UHED(I),UH,I=1,NDF),(UHED(I),UH,I=1,NDF)
500      X (UHED(I),UH,I=1,NDF)
500      IPG = IPG + 1
500      NCT = 50
500      IF(NICD.EQ.1)WRITE(ITP6,LABELN,(XYZ(I,N),I=1,NDIM),(U(I,N),
500      X ,I=1,NDF)
500      IF(NICD.NE.1)WRITE(ITP6,LABELN,(XYZ(I,N),I=1,NDIM),(U(I,N),I=1,NDF
500      X ),(U(I,NUMP+N),I=1,NDF),(U(I,NEP+N),I=1,NDF)
500      CONTINUE
500      UPDATE (1) * * FOR DYNAMIC SOLUTIONS ONLY
500      IF(NICD.GT.1) CALL UPDATE(I,MDEG,NICD,U,U(1,NUMP+1),U(1,NEP+1),
500      X DU,F,DF,IDEST,PROP)
500      IF(1BLK.GT.0) WRITE(ITMP,(U(I,J),I=1,NDF), J=1,40)
500      IF
5001235      IPG = IPG + 1
5001237      NCT = 50
5001240      IF(NICD.EQ.1)WRITE(ITP6,LABELN,(XYZ(I,N),I=1,NDIM),(U(I,N),
5001311      X ,I=1,NDF)
5001311      IF(NICD.NE.1)WRITE(ITP6,LABELN,(XYZ(I,N),I=1,NDIM),(U(I,N),I=1,NDF
5001420      X ),(U(I,NUMP+N),I=1,NDF),(U(I,NEP+N),I=1,NDF)
5001420      CONTINUE
5001423      UPDATE (1) * * FOR DYNAMIC SOLUTIONS ONLY
5001423      IF(NICD.GT.1) CALL UPDATE(I,MDEG,NICD,U,U(1,NUMP+1),U(1,NEP+1),
5001457      X DU,F,DF,IDEST,PROP)
5001457      IF(1BLK.GT.0) WRITE(ITMP,(U(I,J),I=1,NDF), J=1,40)
5001512      IF

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TS0 97C
TS0 98C
TS0 99C
TS0100C
TS0101C
TS0102C
TS0103C
TS0104C
TS0105C
TS0106C
TS0107C
TS0108C
TS0109C
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TS0115C
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TS0123C
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TS0125C
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TS0136C
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TS0141C
TS0142C
TS0143C

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001513      TEMP = 0.
001514      DO 400 N = 1,NUMEL
001515      NPR = .TRUE.
001516      DO 43 I = 1,9
001520      NPLT(I,1) = 0
001521      NPLT(I,2) = 0
001524      IF(NCT.GT.0) GO TO 46
001526      IF(N.GE.NEL.AND.N.LE.NEF) NPR = .FALSE.
001540      IF(NUMPLT.LE.0) GO TO 46
001542      DO 45 I = 1,NUMPLT
001543      IF(NEDATA(I,1).NE.N) GO TO 45
001546      J = NEDATA(I,2)
001547      NPLT(J,1) = 1
001551      NPLT(J,2) = NEDATA(I,3)
001553      CONTINUE
001556      MA = MOD(IX(NEL1,N),100)
001567      IF(MR.LE.0) MRR = IX(NEL1,N)/1000
001600      IF(MR.LE.0) MR = MRR
001603      DO 60 I = 1,NSTF
001605      FORCE(I,1) = 0.
001612      FORCE(I,2) = 0.
001616      LD(I) = 0
001620      IF(MR.NE.MRR.OR.VFLAG) GO TO 60
001625      DO 50 J = 1,NSTF
001627      ESTIF(I,J) = 0.0
001640      CONTINUE
001643      L = 0
001644      DO 110 I = 1,NEN
001645      K = IX(I,N)
001652      DO 90 J = 1,NDIM
001654      X(J,I) = 0.
001661      IF(K.EQ.0) GO TO 120
001663      NEL = 1
001665      DO 100 J = 1,NDIM
001667      X(J,I) = XYZ(J,K)
001701      DO 110 J = 1,NDF
001702      L = L + 1
001704      DUL(J,I) = DU(J,K)
001713      UL(J,I) = U(J,K)
001722      IF(NCT.GT.0) UL(J,I) = UL(J,I) + DUL(J,I)
001731      IF(NICD.EQ.1) GO TO 110
001734      UDL(J,I) = U(J,I) + NUMPLT
001744      UDDL(J,I) = U(J,I) + NEP
001754      IF(NCT.GT.0) UDDL(J,I) = UDDL(J,I) + CS*DUL(J,I)
001765      LD(I) = IDEST(J,K)
002001      CM = TYPE(MA)
002004      IF(MI.NE.TEMP) MCT = 0
002007      TEMP = TM
002010      CALL TICTOC(TIME,5)
C..... COMPUTE ELEMENT STRESSES AND UPDATE FORCES
C..... CALL ELIN18(NEN,NDIM,NDF,NEL,NEL1,NSTF,NEL1ZV,NVEC,MCT,DM,D,XYZ,
X IN-H.FORCE,ESTIF,UDIRECT,5)
C..... CALL TICTOC(TIME,5)
C..... FORM STIFFNESS IF NEEDED FOR THE NEXT TIME STEP

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TS0144C
TS0145C
TS0146C
TS0147C
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TS0185C
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002052      IF (.NOT. VFLAG.AND. NR.EQ. IPR).OR. (VFLAG.AND. IX.NE1.N.NE.EQ.1))
XCALL ELM18(N,MA,NDIM,NDF,NEL,NEL1,NSTF,NSIZV,NVEC,MCT,IM,D,X,YZ,
X IX,H,FORCE,ESTIF,U,VECT,3)
      IF (MOD(IX,NEL1.N),1000)/100.GT.0.AND..NOT.VFLAG)
X CALL PPTMAT(HEAD,IPG,N,NSTF,ESTIF,FORCE,LD,NSTF,0)
C..... MODIFY FOR THE DISPLACEMENT B.C.
      CALL MODIFY(NDF,NEL,NEL1,NEL,IBLK,NSTF,PROP,IX,ICOD,F,FORCE,ESTIF,
X N)
      IF (VFLAG) GO TO 300
      IF (IBLK.EQ.0) CALL GUEIN(A,DF,NEOB,ESTIF,FORCE,LD,NSTF)
      IF (IBLK.GT.0) WRITE(7) ESTIF,FORCE,LD
      IF (.NOT.FLAGC.OR.CFLAG) GO TO 400
      L = 1
      DO 130 I = 1,NEL
      K = IX(I,N)
      DO 130 J = 1,NDF
      DU(J,K) = DU(J,K) + FORCE(L,2)
      L = L + 1
      GO TO 400
300      CONTINUE
C..... ADD THE FORCE TO THE SOLUTION FOR A RESOLVE
      DO 310 K = 1,NSTF
      J = LD(K)
      IF (J.GT.0) DF(J) = DF(J) + FORCE(K,1)
      CONTINUE
310      CALL TICLOC(TYNE,3)
400      IF (FLAGC.AND..NOT.CFLAG) CALL CONTAC(2,IX,NEL1,NDF,NUMIP,DU)
      CFLAG = .TRUE.
      IF (NTB.GT.1) WRITE(ITUP,H)
      IF (IBLK.GT.0) WRITE(ITUP) (IX(NE1.N),N=1,NUMEL)
      IF (.NOT.VFLAG) CALL SOLVEQ(NUMIP,NUMEL,NDF,IDI,M8,MAXBAN,9,NSTF,
1 ISZA,NEOB,IBLK,A,DU,DF,IDEST,FORCE,ESTIF,LD,MAXB,NDEG)
      IF (VFLAG) CALL RESVEQ(NUMIP,NDF,M8,MAXBAN,ISZA,NEOB,IBLK,A,DF,DF,
1 IDEST,IDEST,MAXB,IFLG)
      CALL TICLOC(TYNE,4)
C..... UPDATE THE SOLUTION
      IF (NT.EQ.NTS.AND.M.EQ.NSEQ) GO TO 900
      IF (NCT.GT.0) GO TO 410
      DTP = DT
      I = ITRD
      ITRD = ITUP
      ITUP = I
410      IF (IBLK.GT.0) BACKSPACE ITRD
      IF (IBLK.GT.0) READ(ITRD) (IX(NE1.N),N=1,NUMEL)
      REWIND ITRD
      IF (IBLK.GT.0) READ(ITRD) (DU(I,J),I=1,NDF,J=1,KU)
      IF (VFLAG) CALL RESORP3(NVEC,NCT,J,DF,DU,F,PEOP,NCT,NTT,IDEST,IBLK)
      VFLAG = .TRUE.
      IF (NT.GT.0) VFLAG = .FALSE.
      NCT = NCT + 1
      IF (.NOT.IFLOG) GO TO 500
      WRITE(6,900) NCT,NUMEL,ITUP,IBLK,DO
      IF (NT.GT.0) IFLOG = .TRUE.
      IF (NT.GT.0) IFLOG = .FALSE.

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TS0190C
TS0198C
TS0200C
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TS0226C
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TS0229C
TS0230C
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TS0232C
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TS0235C
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TS0239C

TS0240C

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000016 SUBROUTINE UPDATE (ISU,NDEG,NICD,U,UD,UDD,DU,F,DF,IDEST,PROP)
000016 C***** UPDATE ***** 12/14/73 *****
000016 DIMENSION U(1),DU(1),F(1),DF(1),IDEST(1),UD(1),UDD(1)
000016 COMMON TAPES,ITP5,ITP6
000016 COMMON TIMITS,TIME,DT,DTP,NH,ISZH,C0,C1,C2,C3,C4,C5,C6,NCT
000016 GO TO (100,300,500),ISU
000024 UPDATE (1) * * PREUPDATE OF ACCELERATIONS FOR DYNAMIC SOLUTIONS.
000026 DO 200 N = 1,NDEG
000035 UDD(N) = C4*UDD(N) - C5*UD(N)
000035 C.....
000037 UPDATE (2) * * UPDATE THE SOLUTION FOR GENL. DISPL. VEL. ACCL.
000043 DO 400 N = 1,NDFG
000043 DU(N) = DF(N)
000043 IF (IDEST(N).EQ.0) DU(N) = F(N)*PROP - U(N)
000053 TEMP = DU(N)
000056 IF (NICD.EQ.1) GO TO 400
000060 P = UDD(N)
000061 UD(N) = C1*UD(N) + C2*P + C3*TEMP
000067 UDD(N) = P + C6*TEMP
000073 U(N) = U(N) + TEMP
000100 RETURN
000100 C.....
000102 UPDATE(3) * * SET INTEGRATION CONSTANTS
000102 C6 = C6 + 0.5
000104 IF (C0.EQ.0.0) C0 = 0.25
000104 WRITE (ITP6,2002) C0,C6
000114 C5 = 1./C0/DT
000116 C4 = 1. - 0.5/C0
000121 C3 = C5*C5
000127 C2 = 1. - C6/C0/2.)*DT/C4
000134 C1 = 1. - C6*C0 + C2+C5
000137 C6 = C5*DT
000137 RETURN
000137 2002 FORMAT(10X,30HNEWMARK INTEGRATION PARAMETERS/
C10X17HBETA VALUE = F6.3/
C10X17HGAUTTA VALUE = F6.3)
000137 END
000137
UPD 1C
UPD 2C
UPD 3C
UPD 4C
UPD 5C
UPD 6C
UPD 7C
UPD 8C
UPD 9C
UPD 10C
UPD 11C
UPD 12C
UPD 13C
UPD 14C
UPD 15C
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UPD 35C

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