# ROUTING AND ACTION MEMORANDUM

#### **ROUTING**

TO:(1) Materials Science Division (Varanasi, Chakrapani)

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TITLE: Superstructures based on 2D crystals

#### **ACTION TAKEN BY DIVISION**

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- (x) Perfomance of the research effort was accomplished in a satisfactory manner and all other technical requirements have been fulfilled.
- (x) Based upon my knowledge of the research project, I agree with the patent information disclosed.

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All applications are limited by the properties of the materials that are available to us. Here we propose a new paradigm in material science: superstructures made by 2-dimensional (2D) crystals.

Single-layer crystals (such as graphene, monolayers of boron nitride, molybdenum disulphide, etc) have a number of exciting properties, often unique and very different from those of their bulk counterparts. However, it is the

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#### **Report Title**

### Superstructures based on 2D crystals

#### **ABSTRACT**

All applications are limited by the properties of the materials that are available to us. Here we propose a new paradigm in material science: superstructures made by 2-dimensional (2D) crystals.

Single-layer crystals (such as graphene, monolayers of boron nitride, molybdenum disulphide, etc) have a number of exciting properties, often unique and very different from those of their bulk counterparts. However, it is the combinations of those crystals to form 2D heterostructures that promise really endless possibilities. One can create such Van deer Waals (VdW) heterostructures by placing individual 2D crystals one on top of another. This way, 2D materials with very different properties can be combined into one 3D structure, thus producing novel, multifunctional materials: we can combine conductivity of one 2D crystal, mechanical strength of another, chemical reactivity of the third, while the optical properties would be given by the whole superstructure. By carefully choosing and arranging the individual layers in the stack one can tune the parameters of the heterostructure, creating materials with tailored properties, or "materials on demand".

In this project several type of heterostructures have been created, either to study fundamental properties of graphene and other 2D materials or to fabricate functional devices.

Enter List of papers submitted or published that acknowledge ARO support from the start of the project to the date of this printing. List the papers, including journal references, in the following categories:

(a) Papers published in peer-reviewed journals (N/A for none)

Received	<u>Paper</u>
10/18/2016	2 H. Yang, L. Britnell, A. P. Rooney, E. Lewis, A. Felten, C. R. Woods, V. Sanchez Romaguera, T. Georgiou, A. Eckmann, Y. J. Kim, S. G. Yeates, S. J. Haigh, A. K. Geim, K. S. Novoselov, and C. Casiraghi. Heterostructures Produced from Nanosheet-Based Inks, Nano Letters, (07 2014): 0. doi: 10.1021/nl501355j
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**Inventions (DD882)** 

**Scientific Progress** 

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**Technology Transfer** 

#### **Final Report**

Period covered: 01/01/2013 - 30/06/2016

#### INTRODUCTION

All applications are limited by the properties of the materials that are available to us: the bandgap in silicon dictates the voltages used in computers and the Young modulus of steel determines the size of the beams used in building construction. Here we propose a new paradigm in material science: superstructures made by 2-dimensional (2D) crystals.

Single-layer crystals (such as graphene, monolayers of boron nitride, molybdenum disulphide, *etc*) have a number of exciting properties, often unique and very different from those of their bulk counterparts. However, it is the combinations of those crystals to form 2D heterostructures that promise really endless possibilities. One can create such Van deer Waals (VdW) heterostructures by placing individual 2D crystals one on top of another. This way, 2D materials with very different properties can be combined into one 3D structure, thus producing novel, multifunctional materials: we can combine conductivity of one 2D crystal, mechanical strength of another, chemical reactivity of the third, while the optical properties would be given by the whole superstructure. By carefully choosing and arranging the individual layers in the stack one can tune the parameters of the heterostructure, creating materials with tailored properties, or "materials on demand". In this way, part of the functionality (and there might be several functionalities) is brought on the level of the design of the material itself (usually it is done on the level of creating a structure from a given material). As the range of available 2D materials broadens, so the possible functionality of the 2D-based heterostructures will cover larger and larger areas.

In this project several type of heterostructures have been created, either to study fundamental properties of graphene and other 2D materials or to fabricate functional devices. Different types of 2D crystals have been used, from single and few-layers made by micro-mechanical exfoliation to liquid-phase exfoliated 2D-materials, which are more suitable for industrial applications. In some cases graphene grown by Chemical Vapour Deposition (CVD) has been included in the heterostructure.

#### **SUMMARY OF THE RESULTS**

#### i) Heterostructures made with micro-chemically exfoliated flakes

We first investigated the best materials to create heterostructures, i.e. those giving sharp and smooth interfaces. While contamination (*ie* hydrocarbons trapped between the crystals) can never been completely removed [1], we observed that substrates such as hexagonal-Boron Nitride (hBN) or Transition Metal Dichalcogenides (TMDCs) are able of "self cleaning": surface contamination assembles into large pockets allowing the rest of the interface to become atomically clean [1]. Devices made by selecting area of the sample that are "bubbles-free" show very high mobilities (20000-500000 cm²/V·s). This is fingerprint of the high quality of the interface between the two crystals. However, not all the substrates have the ability of self cleaning: for example images taken by electron microscopy of graphene deposited on atomically

smooth surfaces such as mica, BSCOO, etc show that the contamination is uniformly distributed over all the sample, giving devices with very poor performance (mobility below 1000 cm<sup>2</sup>/V·s).

We further investigate the quality of VdW heterostructures by Raman spectroscopy [2]. Currently, the only method able to provide information about the interface is cross-sectional high-resolution transmission electron microscopy. However, this method is time-consuming and destructive. Therefore, it is essential to develop a fast, simple, and non-destructive technique that can be easily implemented into the device fabrication processes. We have shown that Raman spectroscopy can be used as a non-destructive and rapid technique for probing the VdW forces acting between two atomically thin crystals, where one is a TMDC. In particular, we studied a variety of heterostructures containing MoS<sub>2</sub> monolayer, which is used as a probe to extract information about the interface; graphene, WS2, and mica acted as the other component of the heterostructure. The Raman spectrum of MoS<sub>2</sub> shows two main Raman-active modes, E<sub>2g</sub><sup>1</sup> and  $A_{1g}$ . The  $E_{2g}^{-1}$  mode represents an in-plane vibration, and the  $A_{1g}$  mode corresponds to an out-ofplane lattice expansion. The A<sub>1g</sub> strongly depends on the interaction with its neighboring material: for example its position changes with the thickness of few-layers MoS<sub>2</sub>. Therefore, we used this Raman mode to extract information on the strength of the vdW interaction between the two adjacent layers. We observed that when the contact between graphene and single-layer MoS<sub>2</sub> is perfect, the  $A_{1g}$  position of the heterostructure corresponds to the position of 2-layers MoS<sub>2</sub>. The position does not change for increasing number of graphene layers, indicating that the shift is produced only by the interfacial contact between graphene and MoS<sub>2</sub>. However, when mica is used instead of graphene, the A<sub>1g</sub> position is not uniform, indicating that the interface is not perfect over the contact area. It is also noted that the blue shift of  $A_{1g}$  on mica is smaller than that shift observed in graphene/MoS<sub>2</sub> heterostructures. This indicates that the interfacial contact is not as good as in the graphene heterostructures, in agreement with transport results [1].

Following optimization of the fabrication process, we then demonstrated functional devices. In previous works, we already demonstrated tunneling transistors and photodetector VdW heterostructures [3-5]. Here we took the complexity and functionality of van der Waals heterostructures to the next level by introducing quantum wells (QWs) engineered with one atomic plane precision. We demonstrated light emitting diodes (LEDs) made by stacking up metallic graphene, insulating hBN and various semiconducting monolayers into complex but carefully designed sequences. Our devices exhibit extrinsic quantum efficiency of nearly 10% and the emission can be tuned over a wide range of frequencies by appropriately choosing and combining 2D semiconductors (monolayers of TMDCs) [6].

Until now, electroluminescence (EL) in TMDC devices has been reported only for lateral monolayer devices and attributed to thermally assisted processes arising from impact ionization across a Schottky barrier and formation of p-n junctions. The use of vertical heterostructures allows us to improve the performance of LED in many respects: reduced contact resistance, higher current densities allowing brighter LEDs, luminescence from the whole device area and wider choice of TMDC and their combinations allowed in designing such heterostructures. The same technology can be extended to create other QW-based devices such as indirect excitonic devices, LEDs based on several different QWs and lasers. Furthermore, we demonstrated that this technology is compatible with flexible substrates: the devices are stable under 1% uniaxial strain.

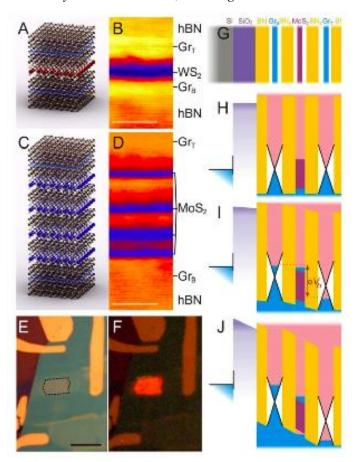


Fig. 1. Heterostructure devices with SQW and MQW. (A) Schematics of SQW heterostructure BN/GrB/2hBN/WS2/2hBN/GrT/hBN. (B) Cross-sectional bright field STEM image of the type of heterostructures presented in (A). Scale bar 5nm. (C, D) Schematics and STEM imaging of hBN/GrB/2hBN/MoS2/2hBN/MoS2/2hBN/MoS2/2hBN/MoS2/2hBN/GrT/hBN. Scale bar 5nm. (E) Optical image of an operational device (hBN/GrB/3hBN/MoS2/3hBN/GrT/hBN). The dashed curve outlines the heterostructure area. Scale bar  $10\mu$ m. (F) Optical image of EL from the same device.  $V_b$ =2.5V, T=300K. (G) Schematic of heterostructure consisting of Si/SiO<sub>2</sub>/ $hBN/GrB/3hBN/MoS_2/3hBN/GrT/hBN$ . (H-J) Band diagrams for the case of zero applied bias (H), intermediate applied bias (I) and high bias (J) for the heterostructure presented in (G).

Figure 1 schematically shows the architecture of single quantum well (SQW) and multiple quantum well (MQW) structures along with optical and STEM images of a typical device. As active layers we used single and multiple layers of TMDC flakes of  $MoS_2$ ,  $WS_2$  and  $WSe_2$ . Figure 2 shows the current-voltage (I-V<sub>b</sub>) characteristics, photoluminescence (PL) and EL spectra from symmetric devices based on  $MoS_2$ . At low  $V_b$ , the PL in Fig. 2A is dominated by the neutral A exciton,  $X_0$  peak at 1.93 eV, while the two weaker and broader peaks at 1.87 and 1.79 eV are attributed to bound excitons. At certain  $V_b$ , the PL spectrum changes abruptly with another peak emerging at 1.90 eV. This transition is correlated with an increase in the differential conductivity (Fig. 2A) and it is expected to be caused by the increase in the Fermi level in the bottom graphene, which rises above the conduction band in  $MoS_2$ , allowing injection of electrons into the QW. In contrast to PL, EL appears only at  $V_b$  above a certain threshold, Figs. 2B. We associate such behavior with the Fermi level of top graphene being brought below the

edge of the valence band so that holes can be injected to  $MoS_2$  from the top graphene (in addition to electrons already injected from the bottom graphene), as sketched in Fig. 1J. This creates conditions for exciton formation inside the QW and their radiative recombination. We find that the EL frequency is close to that of PL at  $V_b \approx 2.4V$  (Figs. 2A-C), which allows us to attribute the

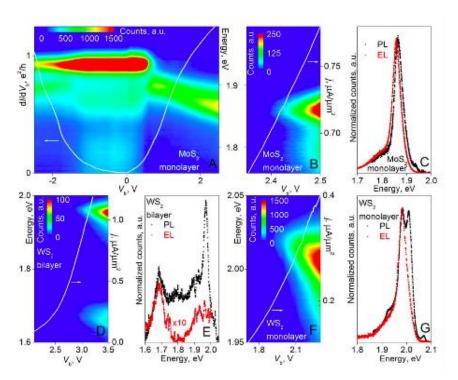


Fig. 2. Optical and transport characterisation of SQW devices, T=7K. (A) Colour map of the PL spectra as a function of  $V_b$  for a  $MoS_2$  –based SQW. The white curve is the device's dI/dVb. Excitation energy is 2.33eV. (B) EL spectra as a function of Vb for the same device as in (A). White curve: its j- $V_b$  characteristics (j is the current density). (C) Comparison of the PL and EL spectra for the same device. As PL and EL occur in the same spectral range, we measured them separately. (D,E) and (F,G) Same as (B,C) but for the bilayer and monolayer  $WS_2$  QWs, respectively. The PL curves were taken at  $V_b$ =2.4V (C), 2.5V (E) and 2V (G); EL – at  $V_b$ =2.5V (C), 2.5V (E) and 2.3V (G).

EL to radiative recombination of X-. The quantum emission QE (defined as 2Ne/I, where e is the electron charge, N is the number of the emitted photons and I is the current) is ~ 1%, which is ten times larger than that of planar p-n diodes and 100 time larger than EL from Schottky barrier devices.

We have demonstrated that the QΕ can be further increased by using **MQWs** in series, which stacked increases the overall thickness of the tunnel barrier and enhances the probability injected carriers recombine radiatively.

This technology also offers the possibility of combining various semiconductor QWs in one device. For example we demonstrated a LED made from WSe<sub>2</sub> and MoS<sub>2</sub> QWs, with the following sequence: Si/SiO<sub>2</sub>/hBN/GrB/3hBN/WS<sub>2</sub>/3hBN/MoS<sub>2</sub>/3hBN/GrT/hB. In comparison with SQW devices, this type of device exhibits more than an order of magnitude stronger both PL and EL, and ~ 5% QE. We

associate this with charge transfer between the  $MoS_2$  and  $WSe_2$  layers such that electron-hole pairs are created in both layers but transfer to and recombine in the material with the smaller band-gap. Such a process is expected to depend strongly on band alignment, which is controlled by bias and gate voltages, giving a complex, asymmetric  $V_b$  dependence of PL and EL in this device.

Furthermore, we used a heterostructure composed by graphene and hBN to study the effect of a superlattice on the properties of graphene [7]. When a crystal is subjected to a periodic potential,

under certain circumstances (such as when the period of the potential is close to the crystal periodicity; the potential is strong enough, etc.) it might adjust itself to follow the periodicity of the potential, resulting in a, so called, commensurate state. Such commensurate-incommensurate transitions are ubiquitous phenomena in many areas of condensed matter physics: from magnetism and dislocations in crystals, to vortices in superconductors, and atomic layers adsorbed on a crystalline surface. We observed a commensurate-incommensurate transition for graphene on top of hBN. Depending on the rotational angle between the two hexagonal lattices, graphene can either stretch to adjust to a slightly different hBN periodicity (the commensurate state found for small rotational angles) or exhibit little adjustment (the incommensurate state). In the commensurate state, areas with matching lattice constants are separated by domain walls of ~ 2 nm size, where the strain accumulates [7]. This effect might explain recent observations about changes in the electronic, optical, Raman and other properties of graphene-hBN heterostructures. In a more recent work [8], we demonstrated that even the weak interaction which exist between the Graphene and hBN is enough to rotate graphene mechanically by a few degrees, which might be used to create nanomotors operating on atom scale.

Finally, we studied tunnelling phenomena in Gr/hBN/Gr heterostructures. In case of the crystallographic lattices of the two graphene electrodes being almost aligned, the tunnelling current between the graphene layers is determined by the elastic processes, where both energy and the momentum of the electrons are conserved. The J-V<sub>b</sub> curve shows a strong peak, followed

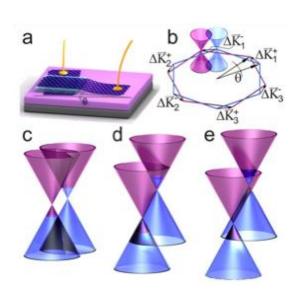


Fig. 3 a) Device schematics: two graphene layers (separated by a BN tunnel barrier shown in light blue); b) A rotation by  $\theta$  of the two graphene layers in real space corresponds to the momentum shift  $\Delta K \pm$  between two Dirac points. Diagrams c-e represent the relative alignment between top (left cones) and bottom (right cones) graphene Dirac points, for different bias conditions.

by a region of negative differential conductance (NDC) [9]. We attribute this peak to resonant tunnelling of carriers between the two graphene electrodes (with momentum conservation). The resonant conditions are demostrated in Fig. 3d, since the tunnelling current is maximized when the momentum difference is compensated by changing electrostatically the energy of the two Dirac cones by an amount  $\pm \hbar v_F \Delta K$  (valid only for small mismatch angle). If the bias voltage is too high, the in-plane momentum is conserved only for a small number of states, thus, leading to the reduction in the current. We also performed measurements under magnetic field, applied parallel to the graphene layers, i.e. perpendicular to the tunnel current. One expects the electron to acquire an additional in-plane momentum due to the action of the Lorentz force, which will modify the resonance. This was experimentally confirmed as the bias positions of the resonances are strongly shifted by the in-plane field. Finally, we demonstrated that the resonant tunneling device, due to its resonance peak negative differential conductance and characteristics, induces a tuneable radio-frequency oscillatory current which has potential for future high frequency technology.

In case of the two graphene electrodes being misoriented, the momentum conservation is satisfied by the phonon emission. Thus, this system presents a very convenient tool to study the phonon spectra of graphene, hBN and their interfaces [10].

In a recent work [11] we demonstrated that the tunnelling current is controlled not only by momentum and energy conservation, but also by pseudospin. Electrons in graphene are described by a vector, two-component, wave-function. Furthermore, they are chiral, which means that there is a very specific relation between the components of the wave-function and the momentum of the electrons. When electrons are tunnelling from one graphene layer to another the interference between the two components of the wave-function results in the effect that only electrons with particular momentum can participate in tunnelling. We can control the selection of the electrons with a particular momentum with the use of in-plane magnetic field – which also allows us to select electrons with particular pseudospin. Thus, we have got a tool that allows injection of electrons in a particular quantum and chiral state.

#### ii) Heterostructures made with chemically exfoliated flakes

Unfortunately micro-mechanical exfoliation is not suitable for commercial exploitation of graphene-based devices, as this fabrication technique is time consuming, expensive and not mass scalable. Therefore, we also investigated the use of chemically exfoliated 2D crystals for fabrication of heterostructure-based devices. Our approach is based on the use of liquid-phase exfoliation (LPE) to produce dispersions of various 2D crystals. One can then use such inks to deposit platelet layers of different materials sequentially by standard low-cost fabrication techniques (drop- and spray-coatings, inkjet printing, etc.), Figure 4. One of the most important advantages of LPE is that the same method can be used to create inks made of nanosheets of

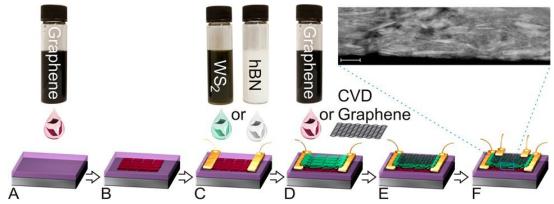


Figure 4 Schematic of a general heterostructure device fabrication process made by using 2D-crystal inks. Last panel shows the final device and the cross-sectional STEM HAADF image of the  $WS_2$  thin film. The scale bar is 35 nm.

different 2D crystals, covering a large variety of properties. Furthermore, this technique is compatible with low-cost and flexible substrates, so it can be combined with flexible electronics. We used different types of inks, from traditional inks based on organic solvents such as NMP, to water-based inks. The aqueous dispersions have been developed in the framework of another project (GOSPEL, funded by the European Science Foundation). The method is based on the idea of using small poly-aromatic molecules to exfoliate graphite in water and stabilize the dispersion [12]. Our approach has been used to produce inks of graphene, hBN and TMDCs (MoS<sub>2</sub>, WS<sub>2</sub>, MoSe<sub>2</sub> and MoTe<sub>2</sub>) [13]. In particular hBN inks have relatively high concentration

and are composed by several single-layers, while TMDCs inks are mostly composed by 3-6 layers. We have shown that the 2D crystals nanosheets are defects-free and are compatible with graphene-based devices. In particular, we used LPE hBN as gating dielectric in a planar graphene-based photodetector [13].

We fabricated and tested different types of devices with the following general structure: BGr/barrier/TGr, where TGr and BGr refer to top and bottom graphene electrodes, respectively (Figure 4). We demonstrated that such devices can act as [14]: (i) tunnelling transistors, where tunnelling between TGr and BGr through a barrier (typically made of TMDC) is controlled by a back gate; (ii) photovoltaic devices, where light absorbed in the barrier (TMDC) is converted into photocurrent through TGr and BGr; (iii) in-plane transistors, where TGr is used as a gate and the barrier as a gate dielectric to control the in-plane current in BGr.

The films were deposited by drop-casting and vacuum filtration and they were characterized by Atomic Force Microscopy, Electron Microscopy, Raman spectroscopy and X-Ray Photoelectron Spectroscopy. We did not observe strong changes in the structural properties of the film with the fabrication method. In order to avoid the shortcut between the two electrodes, it is essential to use a pin-hole free barrier in vertical heterostructures: we did observe that filtration is a perfect technique to obtain pin-holes free TMDC films. By placing the membrane in water, we observed that only the first 20 nm of the membrane get removed and floats in the water. By repeating the fishing process several times, ie by depositing these 20 layers one on top of each other, one can strongly minimize the pin-holes density. We observed that already after 2 depositions (*ie* using a TMDC film with a total thickness of ~40 nm), the pin-holes are almost completely eliminated. Therefore, our devices are characterized by TMDCs film of 40-60 nm thickness, which are pin-hole free.

The tunnelling transistors with structure:  $Si/SiO_2/BGr/WS_2/TGr$  are characterized by current-bias voltage (I–V<sub>b</sub>) characteristics which are strongly nonlinear - fingerprint of those type of devices. However, the device shows characteristics very different from the one obtained by micromechanical exfoliation [3]. For example the zero-bias conductivity decreases dramatically with decreasing temperature. Such a strong temperature dependence suggests an excitation mechanism for charge carrier generation, either from the graphene electrodes (in this case the tunnelling barrier is the Schottky barrier between graphene and  $WS_2$ ) or from the impurity band in  $WS_2$  (a strong impurity band is expected due to the large fraction of edges in our nanocrystals of  $WS_2$ ).

The same device, used as photovoltaic, shows a linear I-V curve under illumination, which demonstrates that in this regime the current is dominated by the photo-excited carriers. Also, finite photocurrent has been observed even at zero bias voltage, demonstrating that such structures can be indeed used as photovoltaic devices. The device is also sensitive to the gate voltage, although again we find that the overall behaviour of the device is very different from the one reported for the same device produced by micro-mechanically exfoliated flakes [5]. We achieved photoresponsivity values of 0.1 mA/W. We also fabricated the same device on plastic and investigate the photo-current stability under bending conditions: the device is stable at least up to 1.5% strain, making this technology compatible with flexible electronics.

In addition we have tested some of the properties of LPE hBN: first, we looked at its dielectric properties [13-14]. Its excellent chemical, mechanical and thermal properties make h-BN thin films a promising dielectric alternative in the next generation of nanodevices. We tested

Si/SiO<sub>2</sub>/BGr/hBN/Au devices, where LPE hBN (prepared through filtering of a hBN suspension, with subsequent transfer of the hBN paper from the filter to the device) served as transparent dielectric between the channel (BGr, CVD graphene) and the gate (Au). From the resistivity of the BGr channel as a function of top gate voltage we can estimate the effective dielectric constant of LPE hBN, which is ~ 1.5. The significant deviation from the bulk value (4, as established in recent tunnelling experiments) is due to lose packing of hBN laminates. This low value of the dielectric constant of the LPE hBN could be an advantageous property when considering its incorporation in densely packed electronic elements, where loss needs to be kept to a minimum. Knowing the capacitance to the top gate allows us to estimate the mobility of the BGr to be of the order of  $3 \times 10^3$  cm<sup>2</sup>/V·s, which is typical of CVD graphene. This clearly indicates that deposition of LPE hBN does not deteriorate the properties of graphene. Finally, we have also tested the breakdown voltage for our LPE h-BN, which turned out to be 0.25 V/nm [14]. This is comparable (or better) than traditionally used inkjet printed dielectric or the dielectric strength of spattered films. This demonstrates that LPE h-BN can be used as a dielectric for transparent, flexible transistor applications. More recently, we have looked also at its thermal properties, we found that hBN laminates have high thermal conductivity of around 30 W/mK, which is very significant for insulating materials [15].

Finally, in the framework of another project, we have carefully engineered our 2D-crystal formulations to be compatible also with ink-jet printing - a simple, versatile, and low-cost technique for large-scale manufacturing of functional devices.

Available inkjet printable formulations, produced by LPE, are still far from ideal as they are either based on toxic and expensive solvents, or require time-consuming and expensive formulation processing, substrate functionalization or need relatively high temperature to dry, which limits the range of substrates that can be used. In addition, none of those formulations are suitable for thin-films heterostructure fabrication, which requires multi-stack formation with well defined interfaces. Fully printing a multi-layer stack is a very well known challenge in printing technology: the different materials in the stack tend to re-disperse at the interface, producing uncontrolled interfaces, resulting in poor device reproducibility and performance. Better control of interfacial effects and processing conditions allowed significant improvements in the performance of organic field effect transistors. Making new printable formulations of functional materials is very challenging: inkjet printing requires the ink to have specific physical properties; water, for example, is unsuitable for both LPE and inkjet printing. We developed a simple method for production of highly concentrated, stable, inkjet printable, water-based inks that can be formulated for a range of 2D materials [16]. The ink composition has been optimized to achieve optimal film formation for multi-stack formation, allowing fabrication of all-inkjet printed heterostructures, such as arrays of all-printed photosensors over cm<sup>2</sup> areas on plastic and paper and programmable logic devices completely made of 2D-crystals. Devices fully printed on paper show responsivity well above 1 mA/W – several orders of magnitude above typical responsivity found with chemically exfoliated materials. Because the inks are expected to find utility in several applications as different consumer products, we also investigated possible adverse effects from exposure and determined their safety limitations by performing doseescalation cytotoxicity assays in vitro using lung and skin cells. Overall, no significant cytotoxic responses compared to untreated cells were observed for all doses. Therefore, our inks could find important applications also in drug delivery and biomedical applications.

#### iii) Other works

- 1) Graphene can be used as tunable plasmonic material. Here we utilised it in quantum cascade lasers, where graphene tunable plasmonic properties allowed us to discriminate between different lasing modes.
- S. Chakraborty, O. P. Marshall, T. G. Folland, Y. J. Kim, A. N. Grigorenko, and K. S. Novoselov, "Gain modulation by graphene plasmons in aperiodic lattice lasers", Science 351(6270), 246-248 (2016).
- 2) CVD graphene can be used for a number of applications where electrical conductivity and optical transparency are required. In this project we created adaptable contact lenses by creating liquid crystals lenses. The use of graphene is essential here, since no other transparent materials can guarantee conductivity on flexible substrate
- S. Kaur, Y. J. Kim, H. Milton, D. Mistry, I. M. Syed, J. Bailey, K. S. Novoselov, J. C. Jones, P. B. Morgan, J. Clamp, and H. F. Gleeson, "Graphene electrodes for adaptive liquid crystal contact lenses", Optics Express 24(8), 8782-8787 (2016).
- 3) The ability of functionalize graphene with several methods, such as radical reactions, cyclo-additions, hydrogenation and oxidations, allows this material to be used in a large range of applications. In this framework, it is essential to be able to control the efficiency and stability of the functionalization process this requires understanding how the graphene reactivity is affected by the environment, including the substrate. In this work we provided an insight on the substrate dependence of graphene reactivity towards hydrogenation by comparing three different substrates: silicon, hexagonal boron nitride (h-BN) and molybdenum disulfide (MoS<sub>2</sub>). Although MoS<sub>2</sub> and h-BN have flatter surfaces than silicon, we found that the H coverage of graphene on h-BN is about half of the H coverage on graphene on both silicon and MoS<sub>2</sub>. Therefore, graphene shows strongly reduced reactivity towards hydrogenation when placed on h-BN.
- S. Son, C. Holroyd, J. Clough, A. Horn, S. Koehler, and C. Casiraghi, "Effect of the substrate on hydrogenation and hydrogenation mechanism", Appl. Phys. Lett, under review
- 4) The hydrogenation mechanism of graphene is not completely understood. From a theoretical point of view, hydrogenation is expected to happen by clustering, i.e. starting from a nucleation point, which corresponds to the formation of the first C-H bond. However, this has never completely verified experimentally. Here we used the roughening dynamic theory to get further insights on the H coverage mechanism.
- S. Son, C. Casiraghi, "Insights on the hydrogenation of graphene by using the roughnening dynamics theory", in preparation
- 5) The superposition of graphene on hexagonal—Boron Nitride (hBN) leads to the creation of a superlattice, which causes a strong reconstruction of the graphene's electronic spectrum with new superlattice Dirac points (SDP) emerging at sub-eV. Changes in the electronic structure are expected to reflect in the Raman spectrum, in particular when shifting the Fermi energy close to

the new Dirac points. In this work we study the Raman spectrum of graphene (Gr) on hBN in both aligned (i.e. Gr/hBN superlattice) and misaligned configurations. We observed strong changes in position and width of G- and 2D-Raman modes in case of gated graphene/hBN superlattices. Those are direct fingerprint of the presence of a period potential distribution.

S. Son, Y. Shin, C. Casiraghi, "Raman spectroscopy of gated graphene/hexagonal-boron nitride superlattices", in preparation

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