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Case-Study Inverse Thermal Analyses of Friction Stir Welds Using Numerical-Analytical Basis Functions

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Introduction

Firction stir welding (FSW) ia a solid-state process where temperatures are below melting, and is significant for various applications because of unique weld characteristics [1-3]. FSW entails joining of a workpiece by means of a stirring tool of given geometry, which has specified translational and rotational speeds (see Fig.1) (see references [4-18]). The present study concerns inverse thermal analyses of FSW processes using a methodology formulated in terms of numerical-analytical basis functions, equivalent source distributions and temperaturess field constraints [19]. References [20-32] describe the general problem of inverse thermal analysis. This report describes inverse thermal analyses of AZ31-Mg-Alloy and Ti-6Al-4V FSWs, which can predict temperature histories within a workpiece for the range of process parameters considered [19, 33-40].

The subject areas presented are organized as follows. First, the procedure for inverse thermal analysis of FSWs using generalized numerical-analytical basis functions and equivalent source distributions is discussed. Second, case study inverse thermal analyses of AZ31-Mg-Alloy and Ti-6A1-4V FSWs are presented. Third, discussion is given concerning aspects of the inverse thermal analysis methodology. Finally, a conclusion is given.



Fig. 1 Schematic representation of friction stir welding and processing.

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Inverse Thermal Analysis Procedure

Following the procedure introduced in [40], the region consisting of stirring tool and stirred material (see Fig. 2) is segmented into a finite set of slices that are perpendicular to the *z* axis. Next, a circle is constructing within each slice , such that the circumference of each circle is defined by the interface between stirred and unstirred material (see Fig. 2). The lines joining centers of these circles and locations of discrete sources, having different strengths, are parallel to V. Accordingly, the *x* and *y* coordinates of each source depends on the width of the stirred-material cross section at its *z* coordinate. This procedure for assigning the locations of discrete sources, of given strengths and diiffusivity vectors $\hat{\kappa}$, permits modeling of asymmetric heat deposition associated with shape differences of FSW advancing and retreating sides.



Fig. 2 Schematic representation of FSW cross section consisting of stirring tool and stirred material, where equivalent source distribution consists of point sources located at centers of circles whose diameters are the SZ width as a function of z.

Next, a parametric representation of temperature fields for heat deposition during welding of plate structures is adopted, which is terms of an effectively complete set of basis functions. A numerical-analytical basis function whose formulation should be relatively optimal for parametric representation of FSWs is given by

$$T(\hat{x}) = T_A + \sum_{k=1}^{N_k} \frac{C(\hat{x}_k)}{r} \exp\left[-\frac{V}{2\kappa}(r+x-x_k)\right] \left(\sum_{n=1}^{N_t} G(z, z_k, n\Delta t, \kappa)\right)$$
(Eq 1)

and

$$T(\hat{x}_n^c, t_n^c) = T_n^c, \qquad (\text{Eq } 2)$$

where

$$G(z, z_k, t, \gamma \kappa) = \left\{ 1 + 2 \sum_{m=1}^{\infty} \exp\left[-\frac{(\gamma \kappa)m^2 \pi^2 t}{l^2}\right] \cos\left[\frac{m\pi z}{l}\right] \cos\left[\frac{m\pi z_k}{l}\right] \right\}$$
(Eq 3)

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and

$$r = \sqrt{(x - x_k)^2 + (y - y_k)^2}$$
 (Eq 4)

$$C(\hat{x}) = \sum_{k=1}^{N_k} C(\hat{x}_k) \delta(\hat{x} - \hat{x}_k)$$
(Eq 5)

and where $C(\hat{x}_k)$ is the value of the discrete source function at location \hat{x}_k . Equation (1) is the solution to the heat conduction equation for a point source located at position (x_k , y_k , z_k) within a region having non-conducting boundaries in *z*. Specifically, Eq.(1) is constructed using combinations of two general forms of the solution to the heat conduction equation. These general forms are the heat-kernel solution of the time-independent steaty-state heat conduction equation for an unbounded region, and the Fourier series solution of the time-dependent heat conduction equation for a region having non-conduction boundaries. Derivation of these solutions are given in Reference [41]. Equation (5) is the source term of the heat conduction equation associated with a spatial distribution of point sources, which results in the solution given by Eq, (1). The quantities κ , *V*, *l* and γ are the thermal diffusivity, welding speed and workpiece thickness and weight coefficient for modeling equivalent effective-diffusion (associated with equivalent source distributions [40]), respectively. The constraint conditions defined by Eq.(2), representing input quantities to the model, are imposed on the temperature field by minimization of the objective function defined by

$$Z_T = \sum_{n=1}^{N} w_n \left(T(\hat{x}_n^c, t_n^c) - T_n^c \right)^2$$
(Eq 6)

where T_n^c is the target temperature for position $\hat{x}_n^c = (x_n^c, y_n^c, z_n^c)$. The quantities w_n (n=1,...,N) are weight coefficients specifying relative levels of influence associated with constraint conditions T_n^c . The output quantity

of the parametric model defined by Eqs.(1)-(6) is the three-diemsional temperature field $T(\hat{x}, t)$ spanning the entire volume of the workpiece. Specifically, the quantities $C(\hat{x}_k)$ are adjusted such that temperature-field values calculated according to Eq.(1) are within a small error tolerances of target temperatures at specified positions (see Tables 1). For the present study, conditions on the objective function defined by Eq.(6) were $w_n = 1$ and $\sqrt{Z_T} < 1$ °C, for all *n*. The parameter γ , defined by Eq.(3), implies a diffusivity vector $\hat{\kappa} = (\kappa, \kappa, \gamma \kappa)$ for the equivalent effective-diffusion, and provides convenient adjustment of the temperature field within upstream regions of FSWs that satisfy boundary conditions defined by the SZ, as well as downstream conditions that are determined only by the diffusivity κ (not κ and $\gamma \kappa$). This follows from the mathematical property of heat diffusion within plate structures (see reference [40] for discussion).

The procedure for inverse thermal analysis defined by Eqs.(1)-(6) entails adjustment of parameters $C(\hat{x}_k)$, \hat{x}_k , Δt and γ . The parametric model combines numerical integration with optimization of linear combinations of numerical-analytical basis functions, which include fundamental solutions to the heat conduction equation and their Fourier-series representation [41]. Equation (1) defines a discrete numerical integration over time, where the time step Δt is specified according to the average energy per length deposited, for transition of the temperature field, at steady state, from uptream regions (close to the SZ) to downstream regions where there is no longer a *z*-coordinate dependence. It should be noted that the formulation of the inverse analysis methodology defined by Eq.(1)-(6) is equipped with a mathematical structure that satisfies all boundary conditions associated with welding of plate structures (see [19]).

Case Study Inverse Thermal Analysis of AZ31-Mg-Alloy FSWs

In this section results of inverse thermal analyses of AZ31-Mg-Alloy FSWs are described, which correspond to different weld process conditions and associated process-control parameters. AZ31-Mg-Alloy is commercially available in sheet form and offers good mechanical properties, but has limited ductility and tends to be brittle at

room temperature. It has been shown, however, that it is possible to form AZ31 sheets having improved ductility and conformability using FS processing, which modifies microstructure [42-50].

The significance of the inverse-problem approach for thermal analysis of FSWs, as for thermal analysis of different types of complex welding processes [19], is that the nature of the energy-source coupling to the workpiece, which is a function of tool geometry and process control parameters, is in principle difficult to specify relative to analysis based on the direct-problem approach. For this study, motivation for adopting SZ boundaries as constraint conditions is that for AZ31-Mg-Alloy FSWs one can associate (approximately) this boundary with an isothermal boundary of known temperature. In the case of the AZ31 magnesium alloy, reference [47] provides an empirical relationship for the estimated uniform SZ temperature as a function of FSW process parameters, which is

$$\frac{T}{T_m} = K \left(\frac{\Omega^2}{V \times 10^4}\right)^{\alpha}$$
(Eq. 7)

where $\alpha = 0.0442$, K = 0.8052 and $T_m = 610$ °C. The present study uses experimentally estimated SZ boundaries as measured in the laboratory for assigning volumetric constraints (see Eq. (2)) on the calculated temperature fields.

The analyses presented here entail calculation of the steady state temperature field for different shapes of the SZ, which are based on experimentally observed estimates of SZ boundaries. The shapes of these boundaries are determined experimentally by analysis of transverse FSW cross sections showing microstructure revealing estimated SZ boundaries, e.g., see references [46,49,50]. For calculations of the temperature field, which adopt SZ boundaries as constraints, the parameter values assumed are the SZ-edge temperature (T_{sz} determined by Eq.(7)), $\kappa = 4.858 \times 10^{-5} \text{ m}^2\text{s}^{-1}$ and $\Delta t = 0.5 \text{ s}$. The diffusivity weight-factor γ , for representation of advective influences (see [40]) is adjusted according to the location of the pseeudo-nonconducting boundary. As discussed previously [37-39], reasonable estimates of κ and isothermal surfaces adopted as field constraints, e.g., SZ-edge temperature T_{sz} , are sufficient for inverse analysis. This follows in that the parameters $C(\hat{x}_k)$, $k=1,...,N_k$, and κ , as well as Δt and γ , are in principle not uniquely determined by inverse analysis. Accordingly, different estimated values of κ , and assigned values of phenomenological parameters Δt and γ , require different values of $C(\hat{x}_k)$ in order to satisfy specified constraint conditions associated with a given isothermal surface.

One goal of the present analysis is determination of parameters that can serve as initial estimates for parameter adjustment with respect to AZ31-Mg-Alloy FSWs, whose process parameters are within similar regimes. Parameter adjustment with respect to other FSWs, which assume the results of this study as initial estimates, would adopt κ and T_{sz} as adjustable parameters, as well as the parameters $C(\hat{x}_k)$, Δt , γ and l_{nc} . Another goal of the present analysis of AZ31-Mg-Alloy FSWs is to provide prototype analyses for demonstrating extension of the methodology for application to FSW analysis in general.

Figures 3 through 8 show estimated transverse cross sections of SZ boundaries for AZ31-Mg-Alloy FSWs obtained from experiment [50] and different two-dimensional slices of three-dimensional temperature fields (°C) calculated using cross section information given in Table 1. Values of the workpiece thickness *l* and welding speed *V* for each FSW considered for analysis are given in these figures. The upstream boundary constraints on the temperature field, $T_c = T_{sz}$ for (y_{c}, z_c) defined in Eq. (2), are given in Table 1 for the SZ boundaries. These constraints are obtained using the estimated transverse weld cross sections of SZ boundaries shown in figures below for the corresponding FSWs, i.e., Welds 1 and 2. The FSW process parameters resulting in these cross sections are given in these figures.

Given in Tables 2 and 3 are values of the discrete source function that have been calculated according to the constraint conditions given in Table 1. Also indicated in Tables 2 and 3 are the assigned values of parameters Δt , γ and *l*. With respect to inverse thermal analysis, the constraint conditions given in Table 1 represent target values of temperature for objective function minimization, which were obtained by a distributed sampling of estimated SZ cross sections as measured in the laboratory (see Figs. 3 and 6). The relative location of each discrete source is specified following the procedure for constructing equivalent source distributions, consistent with FSW

processes, that is described above. Figures 4, 5, 7 and 8 show different planer slices of steady state temperature fields that have been calculated according to the constraint conditions given in Table 1 for estimated SZ-edge boundaries. Referring to the planar slices of the calculated temperature fields shown in these figures, it should be noted that all constraint and boundary conditions are satisfied, namely the condition $T(\hat{x},t) = T_{sz}$ at the SZ edge, and $\nabla T \cdot \hat{n} = 0$ at workpiece surface boundaries, where \hat{n} is normal to the surface. As shown in these figures, the calculated temperature fields have good agreement with experimentally measured cross sections for SZ-edge boundaries. This agreement does not represent model verification in the same sense as models based on first principles, but rather demonstrates parameter optimization with respect to a given upstream isothermal boundary and workpiece boundary conditions.

Table 1 Estimated SZ-edge boundaries on transverse cross sections of Welds 1 and 2.

WELDS 1 AND 2		
(zc mm, half width mm)		
(0.08, 4.48)		
(0.4, 3.4)		
(0.8, 2.64)		
(1.2, 2.0)		
(1.6, 1.56)		
(2.0, 1.24)		

Table 2 Volumetric source function $C(\hat{x}_k)$ calculated according to SZ-boundary constraint conditions given in Table 1, where $\gamma = 0.00605$, l = 2.0 mm, $\Delta l = (2.0/60)$ mm, $x_k = y_k = 0.0$ for k = 1 to 5 (Weld 1).

k	$C(\hat{x}_k)$ /(1.0x10 ⁻	$z_k \left(\Delta l ight)$
	⁴)	
1	6.2	1
2	0.5	25
3	0.5	30
4	0.5	35
5	0.5	40

Table 3 Volumetric source function $C(\hat{x}_k)$ calculated according to SZ-boundary constraint conditions given in Table 1, where $\gamma = 0.00605$, l = 2.0 mm, $\Delta l = (2.0/60)$ mm, $x_k = y_k = 0.0$ for k = 1 to 5 (Weld 2).

k	$C(\hat{x}_k)/(1 \times 10^{-4})$	$z_k \left(\Delta l ight)$
1	5.9	1
2	0.5	25
3	0.5	30
4	0.5	34
5	0.5	40



Fig. 3 Experimentally estimated transverse weld cross section of SZ boundary for AZ31-Mg-Alloy FSW [50] (Weld 1).



Fig. 4 Two-dimensional slices, at half workpiece top surface and longitudinal cross section at symmetry plane, of three-dimensional temperature field (°C) and isothermal boundary at SZ edge calculated using cross section information given in Table 1, where time = x/V and V = 5 mm/min (Weld 1).





Fig. 5 Temperature history (°C) of transverse cross section of weld calculated using SZ cross-section constraints given in Table 1, where $\Delta \tau = \Delta l/V$, $\Delta l = (3.0/60)$ mm and V = 5 mm/min (Weld 1). Temperature scale and time origin t = 0 are shown in Fig. 4.



Fig. 6 Experimentally estimated transverse weld cross section of SZ boundary for AZ31-Mg-Alloy FSW [50] (Weld 2).



Fig. 7 Two-dimensional slices, at half workpiece top surface and longitudinal cross section at symmetry plane, of three-dimensional temperature field (°C) and isothermal boundary at SZ edge calculated using cross section information given in Table 1, where time = x/V and V = 20 mm/min (Weld 2).





Fig. 8 Temperature history (°C) of transverse cross section of weld calculated using SZ cross-section constraints given in Table 1, where $\Delta \tau = \Delta l/V$, $\Delta l = (2.0/60)$ mm and V = 20 mm/min (Weld 2). Temperature scale and time origin t = 0 are shown in Fig. 7.

Case Study Inverse Thermal Analysis of Ti-6Al-4V FSWs

In this section results of inverse thermal analyses of Ti-6Al-4V FSWs are described, which correspond to different weld process conditions and associated process-control parameters. For this study, motivation for adopting $\alpha - \beta$ phase transformation boundaries as constraint conditions is that in practice, for welds of Ti and its alloys, one can associate (approximately) this boundary with the observed edge of the HAZ, and accordingly, specify an isothermal boundary of known temperature. The present study uses experimentally estimated HAZ-edge boundaries as measured in the laboratory for assigning volumetric constraints (see Eq.(2)) on the calculated temperature fields.

The analyses presented here entail calculation of the steady state temperature field for different shapes of SZ boundaries within the neighborhood of the stirring tool boundary, and experimentally observed estimates of the HAZ edge. The shapes of these boundaries are determined experimentally by analysis of transverse weld cross sections showing microstructure revealing estimated SZ and HAZ-edge boundaries (see reference [51]). For calculations of the temperature field, adopting HAZ-edge boundaries as constraints, parameter values assumed are $\kappa = 8.6 \times 10^{-6} \text{ m}^2 \text{s}^{-1}$, $T_{HAZ} = 995 \text{ °C}$, $T_{max} = \text{Cm} \times T_M$, where $T_M = 1604.85 \text{ °C}$ and $\Delta t = 0.5 \text{ s}$.

Figures 9 through 23 show estimated transverse cross sections of SZ and HAZ-edge boundaries for Ti-6Al-4V FSWs obtained from experiment (see references [40, 51]) and different two-dimensional slices of threedimensional temperature fields (°C) calculated using cross section information given in Tables 2, 3 and 4. Values of the workpiece thickness *l* and welding speed *V* for each FSW considered for analysis are given in these figures. The upstream boundary constraints on the temperature field, $T_c = T_{HAZ}$ for (y_c, z_c) defined in Eq. (2), are given in Tables 4, 5 and 6 for the HAZ-edge boundaries. Given in Tables 7-9 are values of the discrete source function that have been calculated according to the constraint conditions and weld specifications given in Tables 4-6. Specifically, given Tables 4-6 provide target values for objective function minimization, which were obtained by a distributed sampling of estimated HAZ cross sections as measured in the laboratory. The relative location of each discrete source is specified following the procedure for constructing equivalent source distributions, consistent with FSW processes, that is described above.

Shown in figures below are different planer slices of the steady state temperature field that have been calculated according to the constraint conditions given in Tables 4-6 for the estimated HAZ-edge boundary. Referring to the planar slices of the calculated temperature fields shown in these figures, it should be noted that all constraint and boundary conditions are satisfied, namely the condition $T(\hat{x}, t) = T_{HAZ} = 995^{\circ}$ C at the HAZ edge, and $\nabla T \cdot \hat{n} = 0$ at surface boundaries, where \hat{n} is normal to the surface. As shown in these figures, the calculated temperature fields have good agreement with experimentally measured cross sections for HAZ-edge boundaries.

Table 4 Estimated SZ-edge and HAZ-edge boundaries on transverse cross section	of Weld	3.
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ADVANCING	RETREATING		
SIDE	SIDE		
(y _c mm, z _c mm)	(y₀ mm, z₀ mm)		
(5.75, 0.25)	(5.625, 0.25)		
(5.25, 0.75)	(5.125, 0.75)		
(4.5, 1.375)	(4.25, 1.375)		
(2.625, 2.0)	(2.625, 2.0)		
(2.0, 2.625)	(2.0, 2.625)		
(2.0, 3.0)	(2.0, 3.0)		

HAZ			
ADVANCING	RETREATING		
SIDE	SIDE		
(y₀ mm, z₀ mm)	(y₀ mm, z₀ mm)		
(6.818, 0.1364)	(6.682, 0.1364)		
(6.682, 0.6818)	(6.682, 0.6818)		
(5.000, 2.0000)	(4.375, 2.0000)		

	SZ			
ADVANCING	RETREATING			
SIDE	SIDE			
$(y_c mm, z_c mm)$	(y₀ mm, z₀ mm)			
(6.25, 0.25)	(6.5, 0.25)			
(5.25, 0.75)	(5.875, 0.75)			
(3.75, 1.375)	(4.375, 1.375)			

(3.125, 2.0)

(2.625, 2.0)

Table 5 Estimated SZ-edge and HAZ-edge boundaries on transverse cross section of Weld 4.

(2.0, 2.625)	(2.25, 2.625)	
(1.375, 3.0)	(1.75, 3.0)	
H	AZ	
ADVANCING	RETREATING	
SIDE	SIDE	
(y _c mm, z _c mm)	(y _c mm, z _c mm)	
(7.363, 0.1364)	(7.363, 0.1364)	
(6.682, 0.6818)	(6.954, 0.6818)	
(6.136, 1.3636)	(6.000, 1.3636)	
(4.910, 2.0454)	(4.636, 2.0454)	
(3.545, 2.7272)	(3.136, 2.7272)	
(2.182, 3.0000)	(2.454, 3.0000)	

Table 6 Estimated SZ-edge and HAZ-edge boundaries on transverse cross section of Weld 5.

SZ			
ADVANCING	RETREATING		
SIDE	SIDE		
(y _c mm, z _c mm)	(y₀ mm, z₀ mm)		
(5.875, 0.25)	(6.125, 0.25)		
(5.25, 0.75)	(4.875, 0.75)		
(3.25, 1.375)	(3.375, 1.375)		
(2.375, 2.0)	(2.5, 2.0)		
(2.0, 2.625)	(2.0, 2.625)		
(1.625, 3.0)	(1.75, 3.0)		

HAZ

ADVANCING	RETREATING	
SIDE	SIDE	
(y₀ mm, z₀ mm)	(y₀ mm, z₀ mm)	
(7.363, 0.1364)	(6.954, 0.1364)	
(6.818, 0.6818)	(6.545, 0.6818)	
(5.454, 1.3636)	(5.863, 1.3636)	
(4.091, 2.0454)	(4.500, 2.0454)	
(2.727, 2.7272)	(3.136, 2.7272)	
(2.182, 3.0000)	(2.318, 3.0000)	

Table 7 Source function $C(\hat{x}_k)$ calculated according to HAZ-boundary constraint conditions given in Table 4, where $\gamma = 0.17$, $\Delta l = (3.0/60)$ mm, $x_k = y_k = 0.0$ for k = 1 to 4 (Weld 3).

k	$C(\hat{x}_k)/(1.95 \times 10^{-3})$	$z_k (\Delta l)$
1	0.3	10
2	0.57	15
3	0.57	20
4	0.4	25
5	0.4	30

Table 8 Source function $C(\hat{x}_k)$ calculated according to HAZ-boundary constraint conditions given in Table 5, where $\gamma = 0.14$, $\Delta l = (3.0/60)$ mm, $x_k = y_k = 0.0$ for k = 1 to 3 (Weld 4).

k	$C(\hat{x}_k)/(4.9 \mathrm{x} 10^{-3})$	$z_k\left(\Delta l ight)$
1	0.177	15
2	0.177	20
3	0.174	25

Table 9 Source function $C(\hat{x}_k)$ calculated according to HAZ-boundary constraint conditions given in Table 6, where $\Delta l = (3.0/60) \text{ mm}, x_k = y_k = 0.0 \text{ for } k = 1 \text{ to } 5 \text{ (Weld 5)}.$

k	$C(\hat{x}_k)/(2.4 \mathrm{x} 10^{-3})$	$z_k \left(\Delta l \right)$
1	0.2	10
2	0.57	15
3	0.57	20
4	0.4	25
5	0.4	30



Fig. 9 Experimentally estimated transverse weld cross sections of SZ and HAZ-edge boundaries for Ti-6Al-4V FSW as measured in laboratory [51] (Weld 3).



Fig. 10 Two-dimensional slices, at half workpiece top surface and longitudinal cross section at symmetry plane, of three-dimensional temperature field (°C) and isothermal boundary at HAZ edge calculated using cross section information given in Table 4, where time = x/V and V = 45 mm/min (Weld 3).





Fig. 11 Temperature history (°C) of transverse cross section of weld calculated using HAZ cross-section constraints given in Table 4, where $\Delta \tau = \Delta l/V$, $\Delta l = (3.0/60)$ mm and V = 45 mm/min (Weld 3). Temperature scale shown in Fig. 10.



Fig. 12 Two-dimensional slices, at half workpiece top surface and longitudinal cross section at symmetry plane, of three-dimensional temperature field (°C) and isothermal boundary at SZ edge calculated using cross section information given in Table 4, where time = x/V and V = 45 mm/min (Weld 3).





Fig. 13 Temperature history (°C) of transverse cross section of weld calculated using SZ cross-section constraints given in Table 4, where $\Delta \tau = \Delta l/V$, $\Delta l = (3.0/60)$ mm and V = 45 mm/min (Weld 2). Temperature scale shown in Fig. 12.



Fig. 14 Experimentally estimated transverse weld cross sections of SZ and HAZ-edge boundaries for Ti-6Al-4V FSW as measured in laboratory [51] (Weld 4).



Fig. 15 Two-dimensional slices, at half workpiece top surface and longitudinal cross section at symmetry plane, of three-dimensional temperature field (°C) and isothermal boundary at HAZ edge calculated using cross section information given in Table 5, where time = x/V and V = 55 mm/min (Weld 4).







Fig. 16 Temperature history (°C) of transverse cross section of weld calculated using HAZ cross-section constraints given in Table 5, where $\Delta \tau = \Delta l/V$, $\Delta l = (3.0/60)$ mm and V = 55 mm/min (Weld 4). Temperature scale shown in Fig. 15.



Fig. 17 Two-dimensional slices, at half workpiece top surface and longitudinal cross section at symmetry plane, of three-dimensional temperature field (°C) and isothermal boundary at SZ edge calculated using cross section information given in Table 5, where time = x/V and V = 55 mm/min (Weld 4).





Fig. 18 Temperature history (°C) of transverse cross section of weld calculated using SZ cross-section constraints given in Table 5, where $\Delta \tau = \Delta l/V$, $\Delta l = (3.0/60)$ mm and V = 55 mm/min (Weld 4). Temperature scale shown in Fig. 17.



Fig. 19 Experimentally estimated transverse weld cross sections of SZ and HAZ-edge boundaries for Ti-6Al-4V FSW as measured in laboratory [51] (Weld 5).



Fig. 20 Two-dimensional slices, at half workpiece top surface and longitudinal cross section at symmetry plane, of three-dimensional temperature field (°C) and isothermal boundary at HAZ edge calculated using cross section information given in Table 6, where time = x/V and V = 105 mm/min (Weld 5).







Fig. 21 Two-dimensional slices, at half workpiece top surface and longitudinal cross section at symmetry plane, of three-dimensional temperature field (°C) and isothermal boundary at HAZ edge calculated using cross section information given in Table 6, where time = x/V and V = 105 mm/min (Weld 5).



Fig. 22 Two-dimensional slices, at half workpiece top surface and longitudinal cross section at symmetry plane, of three-dimensional temperature field (°C) and isothermal boundary at SZ edge calculated using cross section information given in Table 6, where time = x/V and V = 105 mm/min (Weld 5).





Fig. 23 Two-dimensional slices, at half workpiece top surface and longitudinal cross section at symmetry plane, of three-dimensional temperature field (°C) and isothermal boundary at SZ edge calculated using cross section information given in Table 6, where time = x/V and V = 105 mm/min (Weld 5).

Discussion and Conclusion

The results of this study can be adopted as initial estimates for inverse thermal analysis of other FSWs, i.e., parameter optimization can be made more efficient using initial estimates of parameter values, requiring only fine adjustment with respect to constraint conditions (see [40] and references therein). As for inverse thermal analysis using numerical-analytical basis functions and equivalent source distributions, which have been applied to other types of welding processes [19, 33-36], the parametric temperature histories given here can contribute to a parameter space that contains parameters corresponding to different FSW processes, process conditions and different types of metals and their alloys. As discussed previously [40], adopting estimated SZ-edge or HAZ-edge boundaries as constraint conditions is formally equivalent to using thermocouple measurements for this purpose, i.e., thermocouple measurements can be associated with points on three-dimensional isothermal surfaces. In addition, as discussed previously [40], the parametric FSW temperature fields determined in this study may used for extrapolation of temperature histories from the regions close to the SZ-edge to those within the SZ, and thus provide a means for connecting results of inverse thermal analysis based on parameteric modeling, e.g., this study, with those of FSW modeling using basic theory.

Finally, this study demonstrates extension of a methodology for inverse thermal analysis of welds [19, 33-36, 40] with respect to its formulation, which is for application to FSWs. The extension is by inclusion of numerical-analytical basis functions equipped with an effective-duffusity parameterization. Prototype inverse thermal analyses of AZ31-Mg-Alloy and Ti-6Al-4V FSWs are presented that provide proof of concept for inverse thermal analysis using these extended basis functions. This proof of concept is with respect to parameter optimization for different types of SZ shape and HAZ-edge characteristics and boundary conditions.

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