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**A COMPUTATIONAL METHOD FOR EXACT,
DIRECT, AND UNIFIED SOLUTIONS FOR
AXISYMMETRIC FLOW OVER BLUNT BODIES
OF ARBITRARY SHAPE (PROGRAM BLUNT)**

**Ronald H. Aungier
Capt USAF**

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**AIR FORCE WEAPONS LABORATORY
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A COMPUTATIONAL METHOD FOR EXACT, DIRECT, AND UNIFIED SOLUTIONS FOR
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FOREWORD

This report was performed under Program Element 62601F, Project 5791, Task 27.

Inclusive dates of research were May 1969 through December 1969. The report was submitted 11 May 1970 by the Air Force Weapons Laboratory Project Officer, Captain Ronald H. Aungier (WLEE).

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This technical report has been reviewed and is approved.

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ABSTRACT

A time-dependent numerical method is presented that provides direct, exact, and unified solutions for axisymmetric flows about blunt nosed bodies of essentially arbitrary shape. The differencing scheme used ensures that the required stabilizing terms can be specified arbitrarily small and completely independent of the finite difference mesh sizes used. The method is shown to be more accurate than other reported time-dependent techniques. Computational procedures are introduced to enhance the computer efficiency obtainable with the time-dependent method. Extensive comparison with standard computational methods shows that the present method is comparable in both numerical accuracy and computer efficiency. A FORTRAN IV computer code and instructions for its use are provided.

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LIST OF SYMBOLS

A	weighting term defined by equation (45)
a	sound speed
B	weighting term defined by equation (46)
C_1	empirical constant used in equation (63)
C_2	empirical constant used in equation (63)
C_v	specific heat at constant volume
f	general function
G	metric of the coordinate system
g	general function
K	body curvature-also a constant used in the stability analysis
L	streamwise length of a body segment
M	Mach number
\dot{m}	mass flux across a wave
n	coordinate normal to the body
P	pressure
q	a general velocity component
r	body radius defined by equation (1)
S	entropy
s	coordinate parallel to the body
t	time
U	velocity component parallel to the shock wave
u	velocity component along s
V	velocity component normal to the shock wave
v	velocity component along n
W_s	the shock velocity
x	a general coordinate-also the coordinate normal to the shock wave
y	the coordinate parallel to the shock wave
z	distance along the symmetry axis, defined by equation (2)
α	a constant used in the stability analysis
β	a constant used in the stability analysis-also the shock angle
γ	the ratio of specific heats
δ	a constant defined by equation (88)

ζ	the ratio of the pressure in front of a wave to the pressure behind the wave
η	a nondimensional coordinate defined by equation (7)
θ	the body angle
κ	$1 - Kn$
μ	coefficients of the stabilizing terms
μ_0	an empirical constant used in equation (62)
ν	a constant used in the stability analysis
ξ	a coordinate defined by equation (6)
ρ	density
r	transformed time variable defined by equation (8)
Φ	stabilizing term
ϕ	a constant used in the stability analysis

SUPERSCRIPTS

(ξ)	relative to the ξ coordinate
(η)	relative to the η coordinate
(x)	relative to the x coordinate

SUBSCRIPTS

m	maximum value of a quantity
s	value of a quantity at a point on the shock wave
t2	value behind a normal shock
o	value of a quantity when $s = 0$
∞	free stream value of a quantity
2	behind a shock wave

SECTION I

INTRODUCTION

1. STATEMENT OF THE PROBLEM

A direct, exact, and unified numerical method for predicting the inviscid flow about high performance blunt bodies of essentially arbitrary shape is presented. This work was motivated by a requirement for extremely accurate definition of the edge conditions used in detailed boundary layer analyses. Accurate boundary layer analyses for slender high performance reentry vehicles must include consideration of entropy layers, bluntness induced pressure gradients, boundary layer displacement effects and vehicle shape deformation due to ablation mass loss. This requires an accurate inviscid flow solution capable of considering a general class of body shapes. Blunt body solutions of sufficient accuracy to be employed for these analyses are generally limited to body shapes defined by specific conic sections or power law profiles. Also, a description of the entire flow field usually requires the use of matched solutions from two or more numerical methods. A numerical method capable of considering the entire flow field is desirable, i.e., a unified solution procedure. An additional objective of this study was to develop a numerical method which can be easily generalized to three-dimensions to treat the angle-of-attack problem. An extensive development study considering three-dimensional problems is impractical due to the long computation times required to generate these solutions. Also, a numerical method for these problems should first be validated for axisymmetric flow, where standard solutions and experimental data are readily available. Consequently, it is believed that significant advances in developing techniques for three-dimensional problems must be based on axisymmetric flow studies.

2. THE SCOPE OF THIS STUDY

The time-dependent method suggested by the present author in reference 1 is generalized to realize the objectives of this study. Stability arguments are used to establish a differencing scheme that can be considered to be exact for practical purposes. A general body oriented coordinate system is used that is applicable to the desired class of body geometries. Computational procedures are introduced that greatly improve the computer efficiency of the time-dependent method. Extensive comparison with available exact solutions and experimental data is accomplished to demonstrate the accuracy of this method. A FORTRAN IV computer code and instructions for its use are included.

SECTION II

THE BLUNT BODY PROBLEM

1. DESCRIPTION OF THE PROBLEM

The blunt body problem is described by a set of nonlinear partial differential equations. These equations are elliptic in form in the subsonic region of the flow field and hyperbolic in the supersonic region. The mixed elliptic-hyperbolic form of these equations presents a formidable mathematical problem. In the supersonic region, the hyperbolic equations can be solved using the method of characteristics, once a starting line is established by a different numerical method. The method of characteristics has been generalized to three-dimensional flows, but with limited success. Solution of the elliptic equations in the subsonic region requires imposing boundary conditions at an unknown boundary, namely, the detached bow shock. The mixed elliptic-hyperbolic form of the governing equations has usually resulted in the use of two or more numerical techniques to describe the flow field. While a unified numerical solution is preferable to this matching procedure, a practical unified computational method for treating reentry vehicles of current interest has not been available. These other features can be identified as desirable characteristics for a blunt body solution. The solution should be direct, i.e., the body shape should be specified exactly. The governing equations should be treated in their entirety without simplifying assumptions, i.e., the numerical method should be exact. Finally, a practical computational method must allow efficient utilization of the computer to avoid excessive cost in generating solutions.

2. INVERSE TECHNIQUES

The inverse technique (Ref. 2 and 3) circumvents the problem of the unknown bow shock location by solving the flow field behind a specified bow shock. The body shape associated with this shock shape is computed during the solution. When solutions for specific body shapes are required, the shock shape is changed after each solution to achieve an approximation of the desired body shape through an iteration process. Clearly, the exact body shape cannot be achieved. Also, the body shape is extremely sensitive to the shock shape used. Small changes in shock shape can induce large changes in body shape or even preclude the existence of a solution. Consequently, this method is generally limited to body shapes for which the shock shape is well-known. Usually, only the subsonic portion of the flow field is treated due to the difficulty of obtaining convergence when the supersonic afterbody region is included.

3. THE METHOD OF INTEGRAL RELATIONS

The method of integral relations (Ref. 4) provides a direct solution to the blunt body problem. This method employs polynomial approximations to the flow field profile to reduce the governing equations to ordinary differential equations. The flow field is divided into strips to establish these ordinary differential equations. The number of equations that must be solved increases as more strips are used to improve the approximation. Further complication is introduced by the existence of movable saddle point singularities at the sonic line. The large number of equations that must be integrated through these singularities has generally limited the method to 1 or 2 strip approximations. An auxiliary assumption is required, typically that the maximum entropy streamline wets the body. This assumption is of questionable validity when three-dimensional effects are considered. Despite its limitations, the method of integral relations has provided some of the most accurate blunt body solutions reported to date.

4. THE TIME-DEPENDENT TECHNIQUE

The unsteady form of the governing equations are hyperbolic in form throughout the entire flow field. Consequently, the unsteady blunt body problem is a well-posed initial value problem. Several investigators (Refs. 1, 5-17) have employed the unsteady equations to advance the solution in time, asymptotically approaching the steady flow solution. Consequently, the unsteady equations become a means to an end. The additional complication of considering transient flows is justified by the relative simplicity of this problem as compared to the steady flow case. Additional terms must be introduced into the governing equations to stabilize an explicit finite difference scheme. If the influence of these terms on the solution is negligible, the method can be considered exact. The method is direct and, in principle, can provide a unified solution. In addition, it is easily generalized to treat complex body shapes, nonequilibrium flows and three-dimensional flow problems.

SECTION III

THE TIME-DEPENDENT METHOD

1. BACKGROUND

The time-dependent method was suggested by von Neumann and Richtmyer (Ref. 5) for treating flow problems involving shock waves. Their numerical method was motivated by the fact that shock waves tend to thicken in the presence of dissipative effects. Artificial dissipative terms are introduced into the governing equations to spread the shock to a finite thickness. Then, a differencing scheme can be established for the entire flow field which requires no special shock point computations. Lax and Wendroff (Refs. 6-8) developed more sophisticated differencing schemes using the conservation form of the governing equations. While these differencing schemes were motivated by mathematical rather than physical reasoning, the stabilizing terms used are, in effect, dissipative terms. Godunov, Zabrodin and Prokopov (Ref.9) applied the time-dependent method to the blunt body problem. Their method uses the one-dimensional Riemann problem to model the wave interaction between computational cells and at a discontinuous shock boundary. Burstein (Ref. 10) applied various differencing schemes to unsteady multidimensional flows. Bohachevsky, Rubin and Mates (Refs. 11, 12) used this technique to consider blunt body flows including bodies with corners, nonequilibrium flows and angles-of-attack. Moretti, Abbett and Bleich (Refs. 13, 14) introduced a one-dimensional unsteady characteristics technique to treat shock and body points for flows in two and three dimensions. This technique can be considered to be a generalization of Godunov's method where multidimensional effects are included through constant weighting terms. In reference 1, the present author suggested a differencing scheme which partially decouples the magnitude of the stabilizing terms from the finite difference mesh sizes used. This permits selecting the mesh sizes based on numerical accuracy requirements. Subsequently, the magnitude of the stabilizing terms can be established to properly model the inviscid flow problem.

2. ADVANTAGES

The governing equations for unsteady flow are hyperbolic in form. This results in a well-posed initial value problem. The unknown bow shock boundary is allowed to move in time until the proper steady-state profile is achieved. Since the procedure is applicable to the entire flow field, a unified solution procedure can be established. The method is direct and can be considered exact if the in-

fluence of the stabilizing terms is negligible. The method is easily programed for the computer, even when complex reentry phenomena are considered. While the time-dependent method is still in the early stages of development, no inherent limitations have been identified that would preclude its use for a practical computational method for generating exact, direct, and unified blunt body solutions.

3. COMPUTATIONAL ACCURACY

Time-dependent solutions reported in the literature often predict the correct surface pressure distribution. However, enthalpy and density distributions are usually in serious error. The errors are caused by the influence of the stabilizing terms and by the methods used to impose surface boundary conditions. It is believed that an exact solution can be achieved if a differencing scheme is developed that allows the magnitudes of the stabilizing terms to be specified arbitrarily small. In addition, a technique is required to achieve exact computations at points on the body surface. Since a solution to the blunt body problem requires the use of different mesh sizes throughout the flow field, the magnitudes of the stabilizing terms should be independent of the mesh sizes used. While a partial decoupling was achieved in reference 1, it is apparent that complete decoupling must be realized to obtain a practical computational method.

4. COMPUTER EFFICIENCY

Time-dependent methods reported in the literature have required extremely long computation times and extensive computer storage. As a consequence, no solutions for slender high performance vehicles of practical length have been reported. Even the relatively efficient numerical solution reported in reference 14 exhausted the storage of the IBM 7094 computer while treating the flow about an extremely short sphere-cone vehicle at angle of attack. The authors of reference 14 suggest that the consideration of more nodes on a larger computer would result in computation times too long for practical computation. The reduced computational field for axisymmetric problems will alleviate the computer storage problem. However, the computation time required to achieve a steady state increases with the body length. Consequently, an attempt to treat practical high performance vehicles with available time-dependent schemes must result in impractically long computation times. A significant reduction in the computer storage and computation times required by available time-dependent methods must be achieved to realize the full potential of this powerful technique.

SECTION IV

DESCRIPTION OF THE TECHNIQUE

1. THE COORDINATE SYSTEM

A body oriented coordinate system (Ref. 17) was selected for this numerical method. Body oriented coordinates allow a better finite difference nodal distribution than more conventional coordinate systems. As an example, the cylindrical coordinate system of reference 13 may be considered. The nodal spacing across the shock layer is reasonably good near the stagnation point of a blunt body, but the afterbody region for slender bodies must be treated with a very coarse mesh size in the axial direction. It is also apparent that a cylindrical afterbody cannot be treated with the transformed coordinate system of reference 13. The body oriented system is generally preferable for a numerical solution intended to treat a broad class of body shapes. The details of the coordinate system used in this study are illustrated in figure 1. The s coordinate is measured along the body surface from the stagnation point. The n coordinate is measured normal to the body surface with $n = 0$ on the body surface.

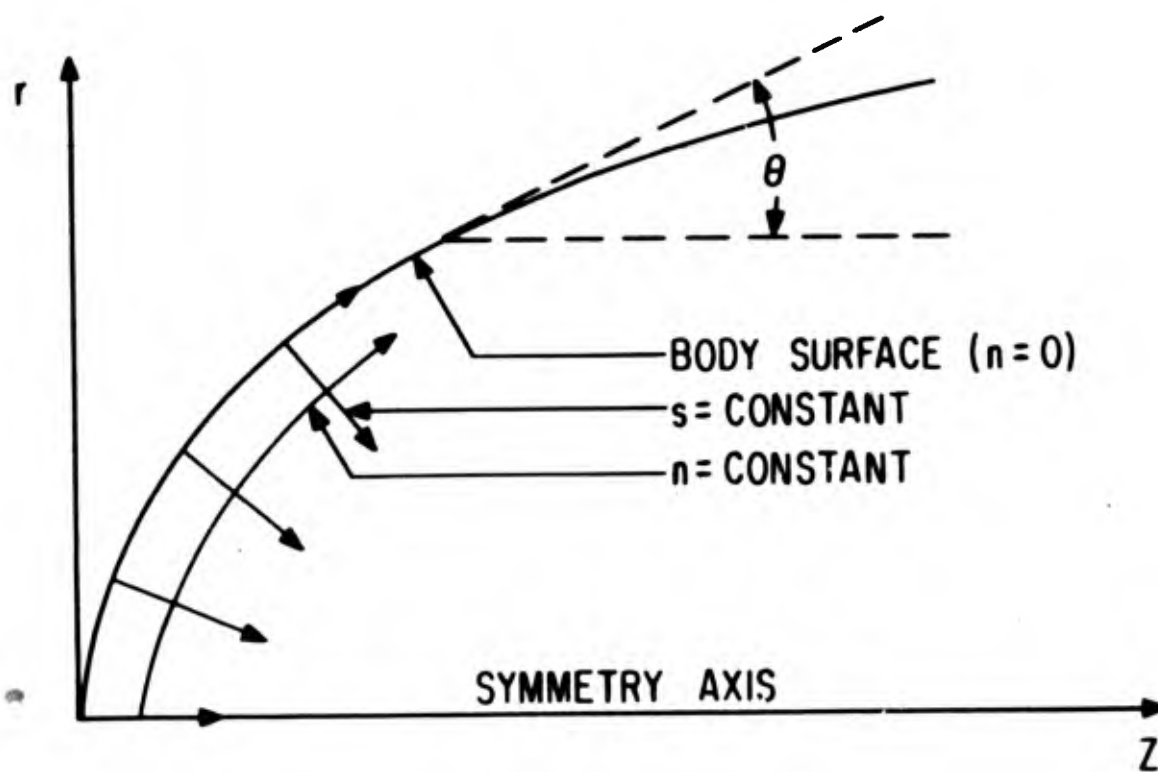


Figure 1. Body Oriented Coordinate System

It is easily shown that

$$r(s,n) = \int_0^s \sin \theta ds' - n \cos \theta \quad (1)$$

$$z(s,n) = \int_0^s \cos \theta ds' + n \sin \theta \quad (2)$$

The curvature of the body is given by

$$K(s) = - \frac{d\theta}{ds} \quad (3)$$

Defining

$$\kappa = 1 - Kn \quad (4)$$

the curvature of any line of constant n is given by K/κ . The metric tensor for this system can be expressed as

$$G = \begin{bmatrix} \kappa^2 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & r^2 \end{bmatrix} \quad (5)$$

These coordinates can be applied to smooth bodies, only, and are restricted to bodies whose curvature is continuous. It will be shown later that the latter restriction can be relaxed by properly staging the solution. Consequently, bodies with discontinuous curvature, such as sphere-cones, can be treated with the present technique. It is convenient to introduce the following transformations

$$\xi = s \quad (6)$$

$$\eta = n/n_s(s) \quad (7)$$

$$\tau = t \quad (8)$$

where $n_s(s)$ is the value of n at the shock and t is the time. Transformations of this type were applied in references 1 and 13 with considerable success. The principal merit is the simplification of the finite difference solution. It was shown in reference 1 that this type transformation can lead to serious errors unless applied with extreme care. The domain of solution is shown in figure 2.

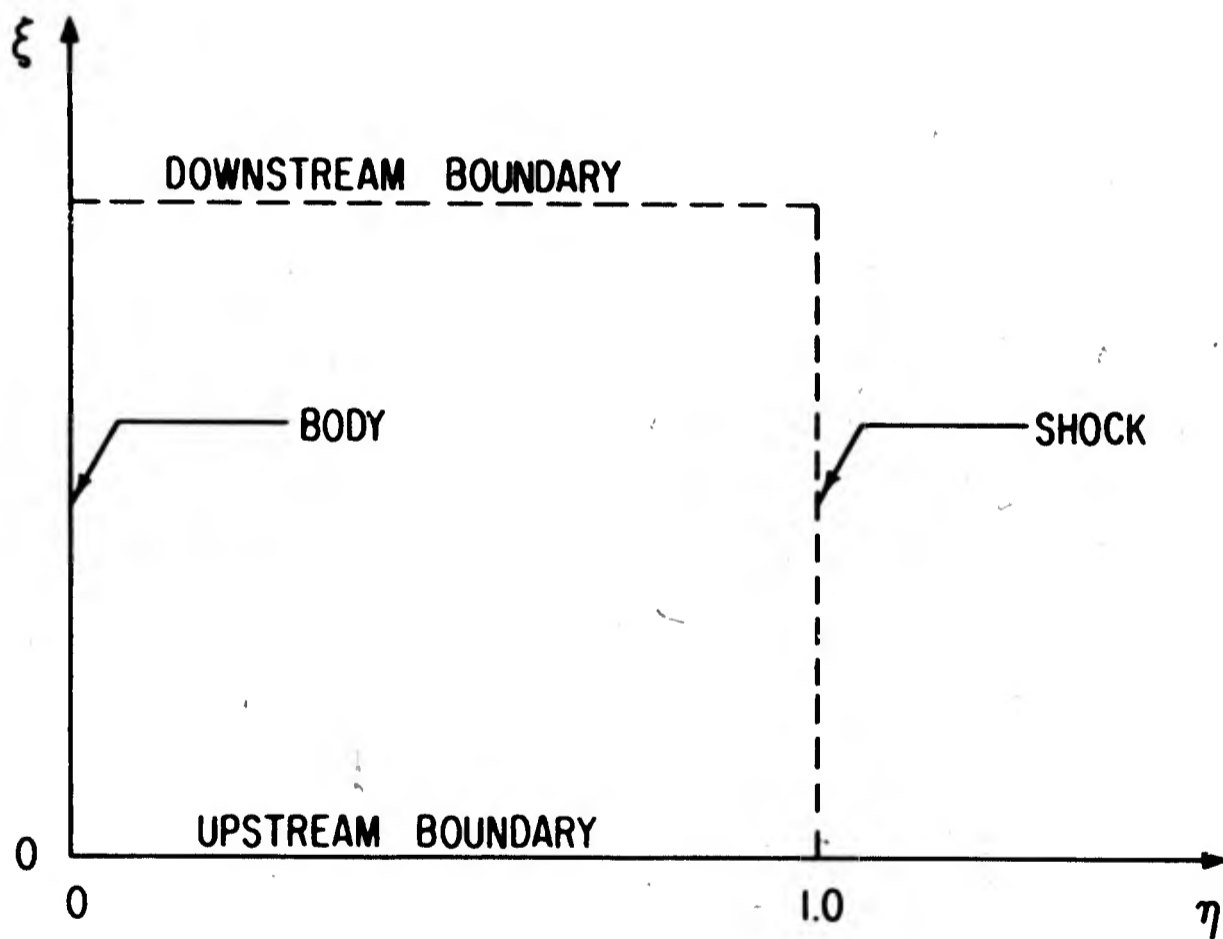


Figure 2. Transformed Coordinate System

It is generally possible to use a constant number of nodes in the η direction in the finite difference solution while maintaining a reasonably good node distribution. This greatly simplifies the numerical logic. Using equations (6) through (8) it is easily shown that

$$\frac{\partial}{\partial s} = \frac{\partial}{\partial \xi} - \frac{\eta}{n_s} \frac{\partial n_s}{\partial s} \frac{\partial}{\partial \eta} \quad (9)$$

$$\frac{\partial}{\partial n} = \frac{1}{n_s} \frac{\partial}{\partial \eta} \quad (10)$$

$$\frac{\partial}{\partial t} = \frac{\partial}{\partial \tau} - \frac{\eta}{n_s} \frac{\partial n_s}{\partial t} \frac{\partial}{\partial \eta} \quad (11)$$

where $\frac{\partial n_s}{\partial t}$ is the shock velocity component in the n direction. Equations (9) through (11) allow solution of the governing equations for the (s,n) coordinate space in the (ξ,η) coordinate space.

2. THE GOVERNING EQUATIONS

The governing equations for adiabatic inviscid flow in the (s,n) coordinate space are

$$\kappa r \rho_t + (r \rho u)_s + (\kappa r \rho v)_n = 0 \quad (12)$$

$$\kappa u_t + u u_s + \kappa v u_n + \frac{1}{\rho} P_s = \kappa u v \quad (13)$$

$$\kappa v_t + u v_s + \kappa v v_n + \frac{\kappa}{\rho} P_n = -\kappa u^2 \quad (14)$$

$$\kappa S_t + u S_s + \kappa v S_n = 0 \quad (15)$$

where the subscripts imply partial differentiation, ρ is the density, P is the pressure, S is the entropy and u and v are the s and n velocity components, respectively. Equation (12) is singular at $s = 0$ and is replaced by

$$\kappa^2 \rho_t + 2\kappa(\rho u)_s + (\kappa^2 \rho v)_n = 0 \quad (16)$$

Equation (16) is obtained from equation (12) by evaluating the limit as s approaches zero using L'Hospital's rule.

For convenience, nondimensional variables will be employed. Defining

$$\kappa' = \kappa / \kappa_0 \quad (17)$$

$$\rho' = \rho / \rho_\infty \quad (18)$$

$$P' = P/P_0 \quad (19)$$

$$S' = (S - S_0)/C_v \quad (20)$$

$$x' = K_0 x \quad (21)$$

$$q' = q/(P_0/\rho_0)^{1/2} \quad (22)$$

$$t' = K_0 (P_0/\rho_0)^{1/2} t \quad (23)$$

where x refers to a general spatial coordinate and q refers to a general velocity component. It is easily shown that equations (12) through (16) remain valid for these nondimensional variables. The prime notation will be omitted in the remainder of this report, and all variables will be assumed to be nondimensional. To complete the set of governing equations, an appropriate state solution is needed. For a perfect gas, the appropriate form is

$$P = \rho^\gamma \text{EXP}(S) \quad (24)$$

3. FINITE DIFFERENCES FOR INTERIOR POINTS.

For interior points, i.e., points not on the shock or body surface, central difference approximations are employed for the spatial partial derivatives

$$f_\xi \rightarrow [f(\xi+\Delta\xi, \eta, r) - f(\xi-\Delta\xi, \eta, r)]/2\Delta\xi \quad (25)$$

$$f_\eta \rightarrow [f(\xi, \eta+\Delta\eta, r) - f(\xi, \eta-\Delta\eta, r)]/2\Delta\eta \quad (26)$$

and a forward difference approximation is employed for the partial derivatives in r .

$$f_r \rightarrow [f(\xi, \eta, r+\Delta r) - f(\xi, \eta, r)]/\Delta r \quad (27)$$

It is known that equations (25) through (27) result in an unconditionally unstable finite difference scheme when applied to equations (12) through (15). Von Neumann (Ref. 5) showed that stability could be achieved by including stabilizing terms in the governing equations. Denoting equations (12) through (15) by

$$g_t = f(s, n, \rho, u, v, P) \quad (28)$$

where g is ρ , u , v or S , the stabilizing terms are added in the form

$$g_t = f - \Phi \quad (29)$$

where

$$\Phi = \mu^{(\xi)} g_{\xi\xi} + \mu^{(\eta)} g_{\eta\eta} \quad (30)$$

In reference 1, $\mu^{(\xi)}$ and $\mu^{(\eta)}$ are equal and are given by

$$\mu^{(x)} = \frac{\nu (\Delta x)^2}{\Delta t} \quad (31)$$

where ν is a constant and x refers to ξ or η . Using equation (31), it was shown in reference 1 that the appropriate stability criterion is

$$\Delta t \leq \frac{\sqrt{\nu} \Delta x}{(q+a)} \quad (32)$$

where Δx is the smallest spatial mesh size, q is the corresponding velocity component and a is the sound speed. An outline of these calculations from reference 1 has been included in appendix I. Although this stability analysis is extremely simplified, equation (32) proved to be extremely accurate in predicting the time step allowable for stable computations (Ref. 1). The impressive success of equation (32) in the numerical method of reference 1 prompted an attempt to exploit the simple stability analysis further to obtain a more sophisticated numerical method. First, equation (32) is inverted to predict the minimum value of ν allowed for stability,

$$\nu \geq [(q+a) \Delta t / \Delta x]^2 \quad (33)$$

Then equation (31) and (33) are combined to yield

$$\mu^{(x)} = (q+a)^2 \Delta t \quad (34)$$

Clearly, this decouples the magnitude of the stabilizing terms from the spatial mesh size. Equation (34) is now applied in the ξ and η directions to yield

$$\Phi = \Delta t [(u+a)^2 g_{\xi\xi} + (v+a)^2 g_{\eta\eta}] \quad (35)$$

which results in the magnitude of Φ being completely defined and minimized for any specified value of Δt . Equation (35) is applied at each node, independently in each coordinate direction, on each time iteration. As a result, the magnitudes

of the stabilizing terms are minimized for all points and are proportional to Δt . The Courant-Friedricks-Lewy (CFL) stability criterion (Ref. 18) must be satisfied

$$\Delta t \leq \frac{\Delta x}{q+a} \quad (36)$$

Consequently, any value of Δt between zero and CFL limiting value can be used. The stabilizing term effects can be reduced by decreasing Δt , but the computation time required to reach a steady state will increase accordingly.

4. FINITE DIFFERENCES FOR BODY POINTS.

For nodes on the body surface, the velocity component normal to the body must vanish. A popular technique for satisfying this requirement is known as the reflection principle (Refs. 1,11,12). Basically, this involves assigning values for the flow parameters at a dummy set of nodes inside the body. The values are identical to those at nodes above the body surface except for the normal velocity component which is assigned the negative of its value above the surface. Then the differencing scheme used at interior points is applied at the body point. It is difficult to assess the accuracy of this scheme. A more appealing approach would be to develop a stable differencing scheme for the body points which does not rely on empirical schemes. This would allow an assessment of the error in terms of the finite difference errors of the various derivative approximations. It is shown in appendix I that a backward difference at the body surface should be stable when $\mu^{(\eta)}$ vanishes. This proved to be true in practice and is the basis for the body point solution presented here. Equation (25) is used to approximate the partial derivative in ξ . The partial derivatives in η were represented by forward differences (Note: this is equivalent to a backward difference in n).

$$f_{\eta} \rightarrow [f(\xi, \eta + \Delta\eta, \tau) - f(\xi, \eta, \tau)] / \Delta\eta \quad (37)$$

At body points, $\mu^{(\eta)}$ was allowed to vanish, i.e.,

$$\Phi = \Delta t (u+a)^2 g_{\xi\xi} \quad (38)$$

and the velocity component normal to the body was required to vanish on the body surface. With these exceptions, the differencing scheme for interior points was applied for the body points. Comparison between this technique and the reflected boundary condition logic is given in the presentation of results. It will become apparent at that time that this approach is considerably more accurate.

5. FINITE DIFFERENCES FOR THE STAGNATION STREAMLINE.

For axisymmetric flow, points on the stagnation streamline, like the body points, are characterized by the vanishing of one of the velocity components. The stability arguments of appendix I should also apply to these points. Consequently, the procedure for interior points are modified by allowing $\mu^{(\xi)}$ to vanish such that

$$\Phi = \Delta t (v+a)^2 g_{\eta\eta} \quad (39)$$

The partial derivatives in ξ all vanish at $s = 0$ except for u_ξ . Since a forward, backward and central difference approximation for u_ξ are all identical at $S = 0$, it is of little concern as to which is applied. The velocity component, u , is required to vanish. The differencing scheme for interior points is modified to include equation (39) and applied at these points. Special attention is required at the stagnation point where the velocity vanishes. Equations (15) and (16) must be solved for entropy and density. The stability arguments of appendix I now indicate that both $\mu^{(\xi)}$ and $\mu^{(\eta)}$ can vanish, i.e., $\Phi = 0$. Then, equation (16) is solved using a forward difference approximation for the partial derivative in η . Equation (15) must be treated with care. The procedure used for equation (16) cannot be applied since equation (15) would then be degenerate. The proper value of entropy behind a steady normal shock could be imposed and held constant with advancing time. This would preclude the consideration of transient phenomena such as time varying free-stream conditions. An averaging procedure using entropy values at adjacent nodes was attempted, but resulted in longer computation times to achieve the steady-state solution. For steady-state solutions, the assignment of the value of entropy at the node upstream of the stagnation point at each time step proved to be a successful approach. However, for transient solutions, this proved to be unreliable, usually resulting in the generation of indefinite values in actual computation. A procedure that allows transient problems while providing rapid convergence to the steady-state solution is as follows:

$$S(o, o, r + \Delta t) = \mu S(o, o, r) + (1 - \mu) S(o, \Delta\eta, r) \quad (40)$$

where

$$\mu = \left[\frac{S(o, o, r) - S(o, \Delta\eta, r)}{S(o, o, r) + S(o, \Delta\eta, r)} \right]^2 \quad (41)$$

When the values of entropy at $\eta = 0$ and $\eta = \Delta\eta$ are essentially equal, the stagnation point entropy is assigned a value approximately equal to its upstream value. When these values of entropy are quite different, the value is held approximately constant with a weak dependence on the upstream value. Since viscous effects predominate in this region in the physical problem, this procedure appears to be a reasonable way to obtain realistic transient solutions.

6. COMPUTATIONS FOR THE SHOCK POINTS.

The shock points were treated with a quasi-one-dimensional unsteady characteristics technique suggested by Moretti and Abbett (Ref. 13). This technique assumes that the variation in the flow variables normal to the shock are the relevant quantities for determining the shock behavior, while variations in the flow variables parallel to the shock can be assigned constant values for use as weighting functions. The method is quite similar to Godunov's method (Ref. 9) as applied in reference 16 except for the use of these constant weighting functions. The equations of motion are written for the shock fixed Cartesian coordinate system shown in figure 3, where β is the local angle between the shock and the free stream velocity vector. Denoting the x and y velocity components by V and U, respectively, the governing equations are

$$\rho_t + V\rho_x + \rho V_x = -A \quad (42)$$

$$V_t + VV_x + \frac{1}{\rho}P_x = -B \quad (43)$$

$$S_t = \text{constant} \quad (44)$$

where the y momentum equation has been omitted since the value of U is known from M_∞ and β . A and B are defined by

$$A = (\rho U)_y \quad (45)$$

$$B = UV_y \quad (46)$$

and will be treated as constants. In appendix II it is shown that the relevant characteristics can be defined by

$$\frac{dx}{dt} = V + a \quad (47)$$

$$\frac{dx}{dt} = V - a \quad (48)$$

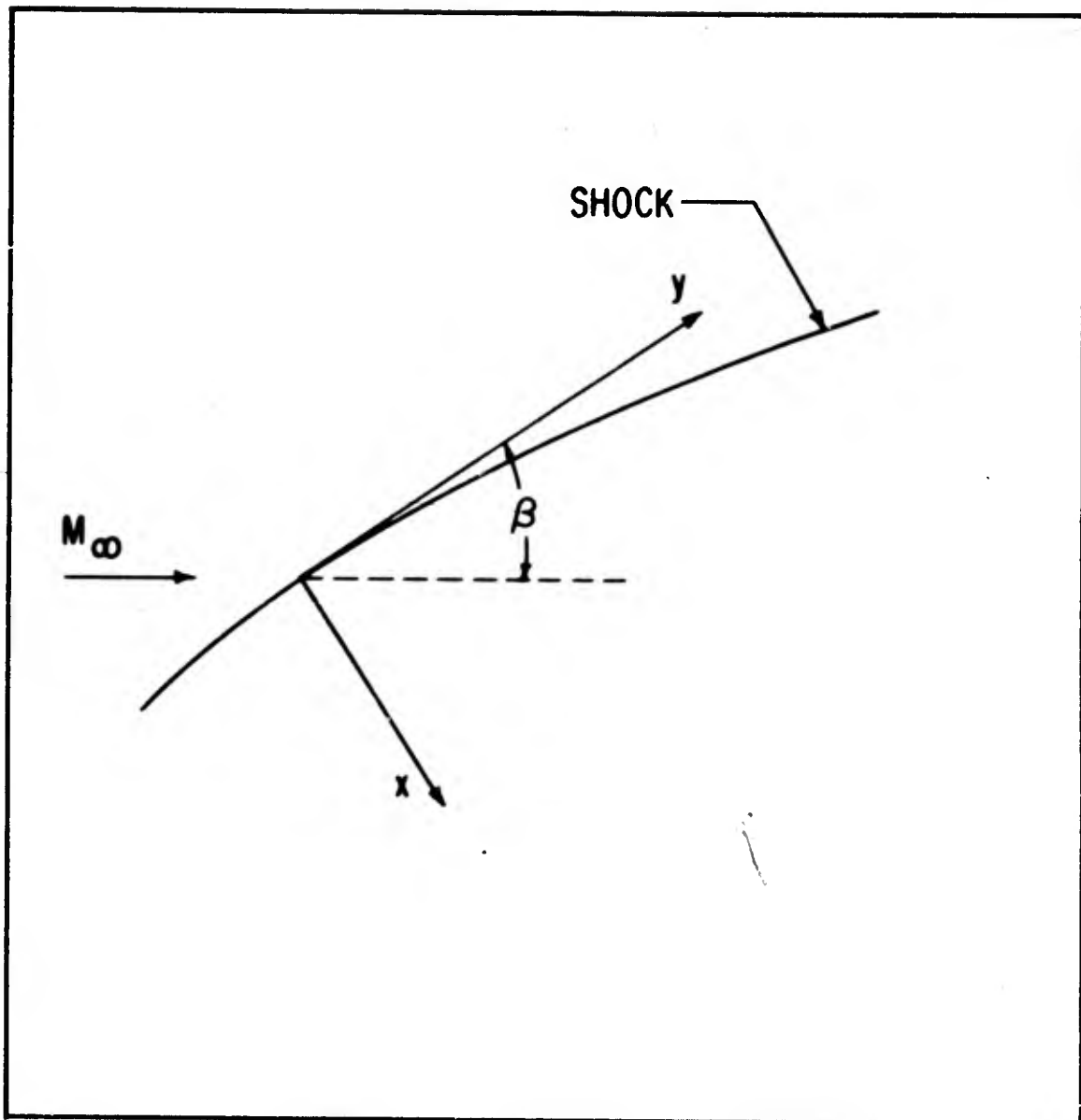


Figure 3. Shock Fixed Coordinates

and the compatibility equation associated with equation (48) is

$$\frac{dP}{dt} = \left[\rho a \frac{dV}{dt} + (B-A) \frac{a}{\rho} \right] \quad (49)$$

It is shown in appendix II that application of this method on the symmetry axis requires that A be multiplied by 2. The problem can be visualized as a one-dimensional Riemann problem with equation (49) replacing the usual Riemann invariants. The physical phenomena is illustrated in figure 4 where W_s refers to the shock velocity in the x direction. The problem is solved by an iteration procedure. The shock parameters are established from the Rankine-Hugoniot relations

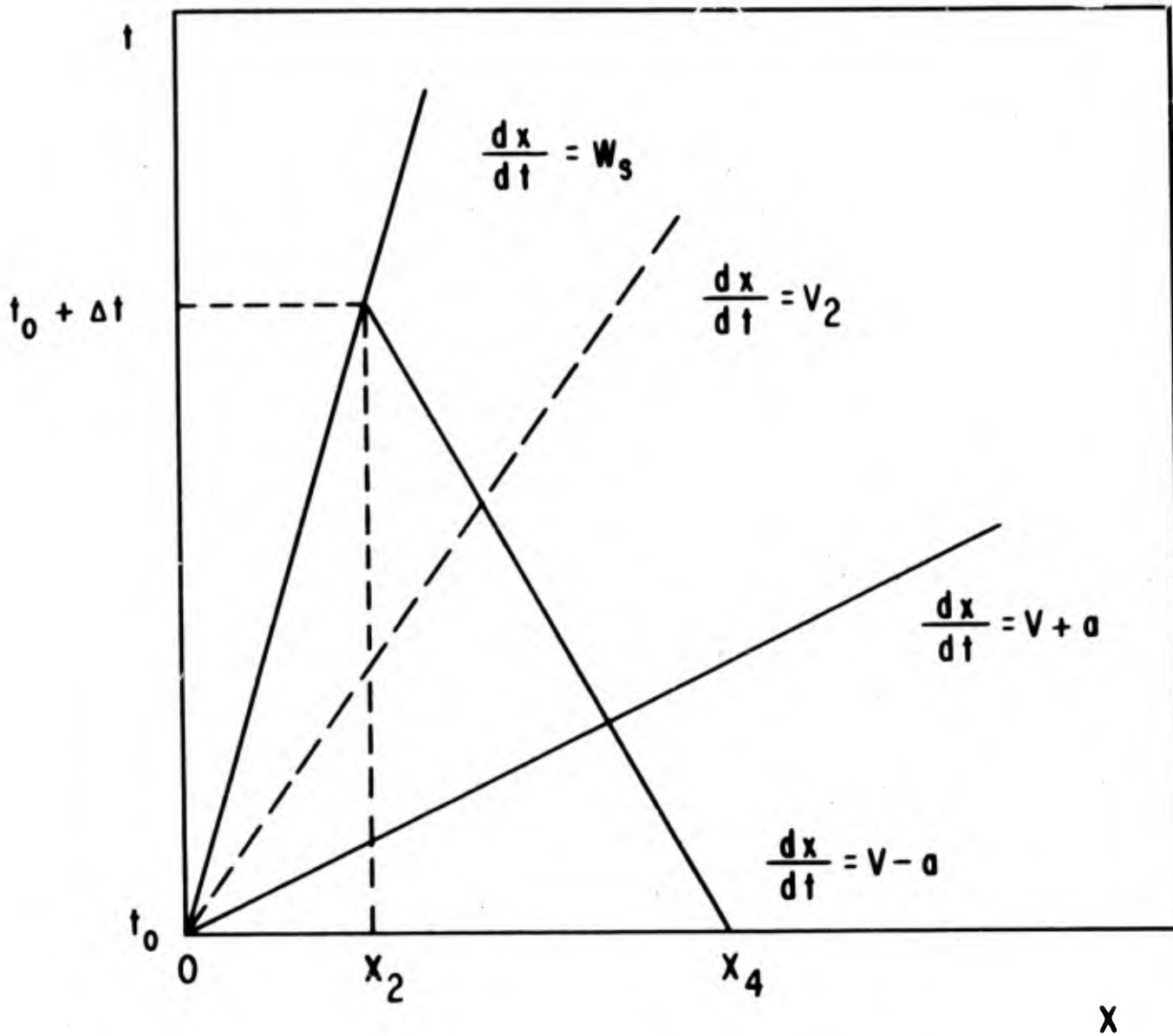


Figure 4. Shock Wave x-t Diagram

once satisfactory agreement between the Rankine-Hugoniot pressure prediction and the pressure prediction of equation (40) has been achieved. The Rankine-Hugoniot relations are applied in the form

$$\dot{m} = \left[\frac{\gamma+1}{2} P_s + \frac{\gamma-1}{2} \right]^{1/2} \quad (50)$$

$$\dot{m} V_s + P_s = \dot{m} V_\infty + 1 \quad (51)$$

$$W_s = V_\infty - \dot{m} \quad (52)$$

$$\rho_s = \frac{\dot{m}}{V_s - W_s} \quad (53)$$

Where a perfect gas has been assumed and the subscript, s, refers to a quantity behind the shock. The velocity component in the y direction is unchanged across the shock

$$U_s = U_\infty \quad (54)$$

where the subscript ∞ refers to a free stream value. The entropy behind the shock is computed from the equation of state

$$S_s = \ln(P_s) - \gamma \ln(\rho_s) \quad (55)$$

The iteration procedure is diagramed in figure 5. Starting with the shock parameters at point 0 (Fig. 4) as an initial guess, point 2 is located by

$$X_2 = W_s \Delta t \quad (56)$$

and the flow variables at 2 are assumed identical to their values at 0, initially. Point 4 is located from equation (48) in the form

$$X_4 = X_2 + (V_2 - a_2) \Delta t \quad (57)$$

Then, the point 4 is redetermined from

$$X_4 = X_2 + \frac{1}{2} (V_4 - a_4 + V_2 - a_2) \quad (58)$$

and an iteration on equation (58) is performed until successive values of X_4 agree to an acceptable degree. Using the known solution at time t_0 , the values of A and B are determined at 4, averaged with their values at point 0 and used to solve equation (49). The pressure at 2 is compared with the pressure predicted by equation (49). If they agree within an acceptable error, the shock parameters are computed from equations (50) through (55) and the iteration loop is terminated. If the agreement is not satisfactory, a new pressure at point 2 is computed as the average of the resident value and the value predicted by equation (49). Other shock parameters required are obtained from equation (50) through (53) and the entire iteration loop is repeated. The value of the shock angle, β , is held constant throughout the iteration procedure, equal to its value at t_0 .

$$\beta = \theta - \tan^{-1} \left[\frac{1}{\kappa_s} \frac{\partial n_s}{\partial \xi} \right] \quad (59)$$

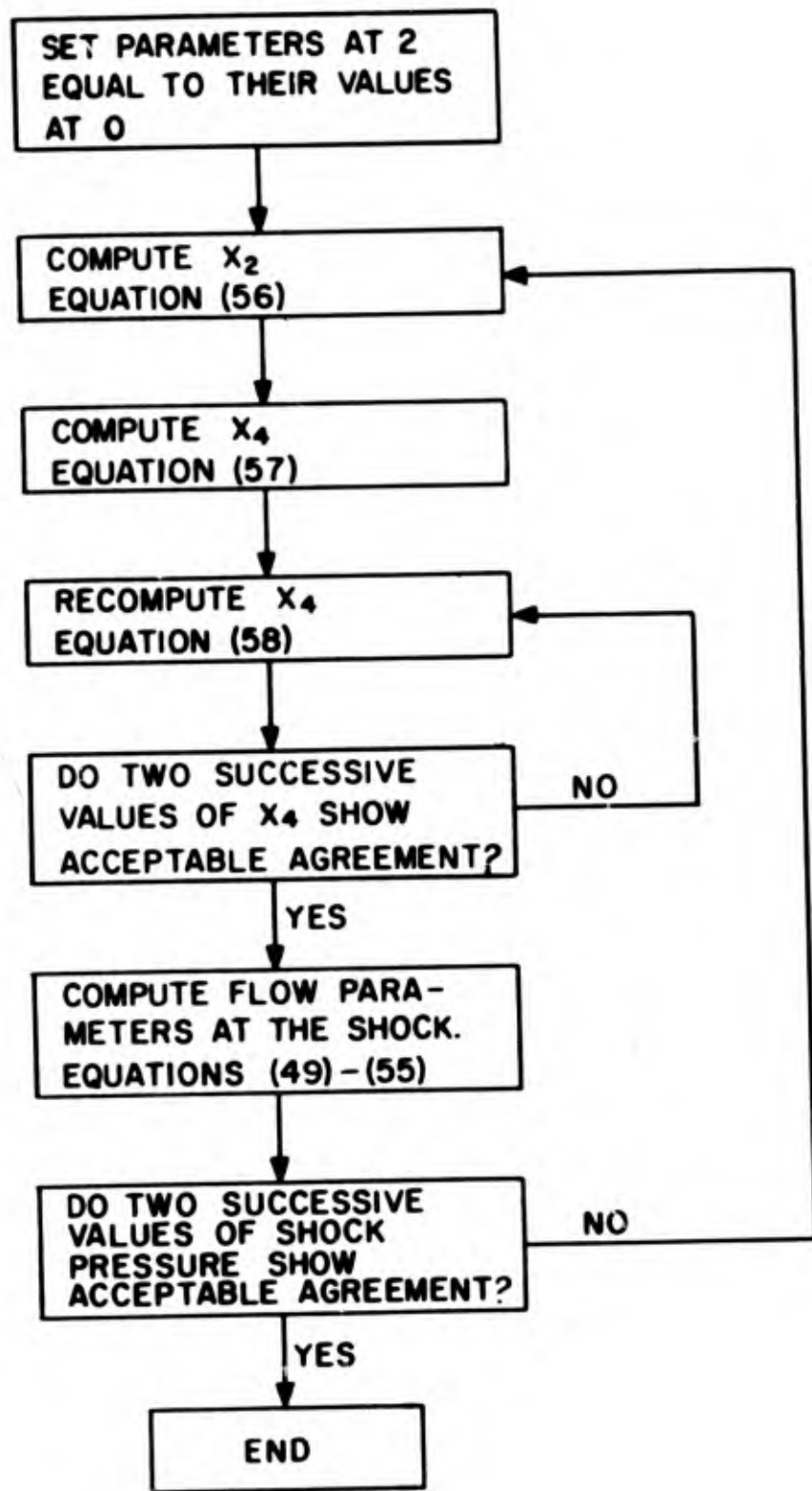


Figure 5. The Shock Point Iteration Scheme

The evaluation of the shock slope through the term containing the derivative of n_s with respect to ξ can be a source of serious error. When a central difference approximation is employed, the slope at any point is evaluated independently of the value of n_s at that point. If the initial guess for the shock shape is not carefully specified, a discontinuity in the predicted shock shape will often result. The present numerical scheme proved to be sufficiently reliable to avoid fatal errors in actual computer computations when this occurred, but computation times were increased by factors of 20 to 30. The necessity of requiring extremely accurate initial guesses for shock shape would greatly complicate the program operation. The problem arises at regions where the shock curvature is large such as the supersonic region on a hemisphere. The actual difference approximation provides a reasonable value for the shock slope despite the fact that the shock is discontinuous at the point under consideration. The use of a backward difference approximation eliminated this problem, but proved to be inaccurate. To obtain sufficient accuracy while eliminating discontinuities in the shock, an average of two backward differences and one forward difference is employed.

$$\frac{dn_s}{d\xi} = \left[n_s(\xi + \Delta\xi) + n_s(\xi) - 2n_s(\xi - \Delta\xi) \right] / 3\Delta\xi \quad (60)$$

Satisfactory results are obtained at all points with this approximation while discontinuities in the shock are avoided. Introducing a dependence on the local value of n_s in the finite difference approximations acts as a damping influence on the computation. Moretti and Bleich (Ref. 19) experienced this problem for very blunt elliptical bodies. The shock shapes predicted in the reference are very similar to the type predicted with this method using the central difference approximation. It is probable that equation (60) could be used in their solution to avoid this difficulty.

7. POINTS ON THE UPSTREAM BOUNDARY.

Solution for nodes on the upstream boundary are treated by extending the solution linearly in ξ one extra mesh width. The interior point logic is then applied in its entirety. It is noted in appendix I that this is equivalent to employing backward differences for the partial derivatives in ξ and allowing $\mu(\xi)$ to vanish. Since the linear extension logic was numerically faster and simpler, it has been adopted for this numerical method.

8. SELECTION OF THE TIME STEP.

The proper choice of a value of Δt for numerical computations requires that the quantity $\mu^{(x)}$ in equation (34) be small relative to the dynamic terms in the equations of motion, i.e., the effect of the stabilizing terms must be negligible. A simple estimate for the magnitude of the dynamic terms is given by the free-stream velocity. It is reasonable to require

$$\mu_0 = \frac{\mu^{(x)}}{q_\infty} \quad (61)$$

where μ_0 is a constant whose magnitude is small relative to unity. Using free-stream values for velocity and sound speed in equation (34) and combining that equation with equation (61)

$$\Delta t = \frac{\mu_0 q_\infty}{(q_\infty + a_\infty)^2} \quad (62)$$

Equation (62) has been used to predict Δt for a range of free-stream Mach numbers from 1.5 to 500. The percent error resulting from the stabilizing terms was essentially equal in all cases, indicating that equation (62) is a proper expression for computing Δt . The computer code makes an additional check to ensure that the Courant-Friedricks-Lewy stability criterion, equation (36), is satisfied. A typical value for the constant, μ_0 , is 0.04.

9. SEGMENTED SOLUTIONS

A computational procedure has been developed that avoids the problems of long computation times and excessive computer storage requirements. When the steady-state flow over typical high performance blunt bodies with relatively long afterbody regions is considered, e.g., sphere-cones with axial lengths in excess of 50 nose radii; the procedure used will be referred to as the segmented approach. The flow field is computed in a series of discrete segments marching back along the body. The steady-state solution is obtained for each segment before proceeding to the next. The solution at the end of each segment is used as a constant starting line for the next segment. The computations on each segment are terminated at the downstream boundary with the linear extension procedure described previously. The lengths of the segments is arbitrary except for the first, which must include the subsonic region and terminate in supersonic flow. The nodal spacing along the body and the number of nodes across the shock layer can be

changed for each segment before starting the solution at that segment. Typical slender blunt bodies may require 5 nodes across the shock in the stagnation region while as many as 40 may be required far back on the afterbody. The number of nodes considered along the body can vary. Generally, far downstream on the afterbody a coarse nodal spacing using one computational node in the streamwise direction is the optimum choice for minimum computational time. An empirical expression to predict the physical time required for essentially steady-state solutions, r_f , is

$$r_f = (C_1 + C_2 L) / (u_o + a_o) \quad (63)$$

where C_1 and C_2 are constants, L is the streamwise length of the segment and u_o and a_o are the velocity and sound speed at the body on the upstream boundary of the segment. This expression has been used for segments of widely varying values of L and has always predicted proper values for r_f . Equation (63) permits the segmented approach to be used in the solution without complicated input instructions. In practice, computations proceed until r exceeds r_f . Then the shock velocity at all points is checked to insure that all values are less than 10^{-3} . If any shock velocity is greater than that value, further iterations are performed. As a result, the input instructions to the computer code are minimal. Since the exact number of iterations to be performed for a given problem is not known in advance, an accurate estimate of computer time is difficult. To avoid wasting large amounts of computer time due to a time limit termination, flow data for each segment is stored in sequential files on a magnetic tape. A simple restart option allows renewing computations at the first incomplete segment with little waste of computer time. A solution for a sphere-cone vehicle with a 10° cone half-angle and axial length of 50 nose radii was performed in under 6 minutes on the AFWL CDC 6600 computer using the segmented approach. To resolve the afterbody flow, 17 nodes across the shock layer were required. A mesh size of 0.15 was required to resolve the flow in the stagnation region. Without the segmented approach these two quantities would have had to remain constant throughout the flow field. Using the maximum value of the denominator in equation (63) to ensure the estimate of computation time would be a minimum, equation (63) and the known computational efficiency of the computer code resulted in a predicted computational time of 72 hours for the same computer code, without the segmented

approach, on the same computer. The estimate is probably less than the time which would actually be required due to the use of the maximum value of the denominator in equation (63). It is easy to understand why time-dependent solutions for this type of vehicle have not been reported previously.

SECTION V

PRESENTATION OF RESULTS

1. INVESTIGATION OF THE STABILIZING TERM EFFECTS

A value of the parameter μ_0 must be selected empirically. A comparison of the predicted stagnation point pressure, P_0 , with the known stagnation pressure behind a normal shock, P_{t2} , was accomplished for this purpose. The error in P_0 and the number of iterations required for an approximate steady state are shown in figure 6 for a range of values of μ_0 . The case considered was Mach 4 flow over a sphere with 5 nodal points across the shock layer. The errors shown in figure 6 include both stabilizing term effects and finite difference errors. Since the finite difference errors should be approximately constant for all values of μ_0 , the relative values of the errors indicate the effects of the stabilizing terms. For values of $\mu_0 > 0.05$, the Courant-Friedricks-Lewy (CFL) stability limit was exceeded, causing the computer code to ignore the specified values of μ_0 . Based on the results presented in figure 6, a value of $\mu_0 = 0.04$ was selected and used for all other cases presented in this report. The error in stagnation point pressure is 0.53 percent for this value of μ_0 . This should be adequate for most applications. Since the computation time required varies approximately as the inverse of μ_0 , requiring greater accuracy will greatly increase the computation time while accepting larger errors will not greatly reduce it. The importance of the stabilizing term effects is further demonstrated in figures 7 and 8, where the entropy and total enthalpy distributions on the body predicted by the present technique and the method of reference 1 are compared. The method of reference 1 was selected for this comparison since it was shown (Ref.1) that accurate pressure distributions can be predicted with that method. It is seen that the present method predicts constant entropy on the body while the method of reference 1 shows considerable variation in that quantity. The error in total enthalpy on the body is less than 0.5 percent for the present method, while reference 1 predicts errors greater than 8 percent for this quantity. Clearly, the validation of a time-dependent technique cannot be accomplished by examining pressure profiles only.

2. UNSTEADY FLOW

By assuming an initial shock shape extremely close to the body surface, an approximate description of the flow over a body started impulsively from rest can

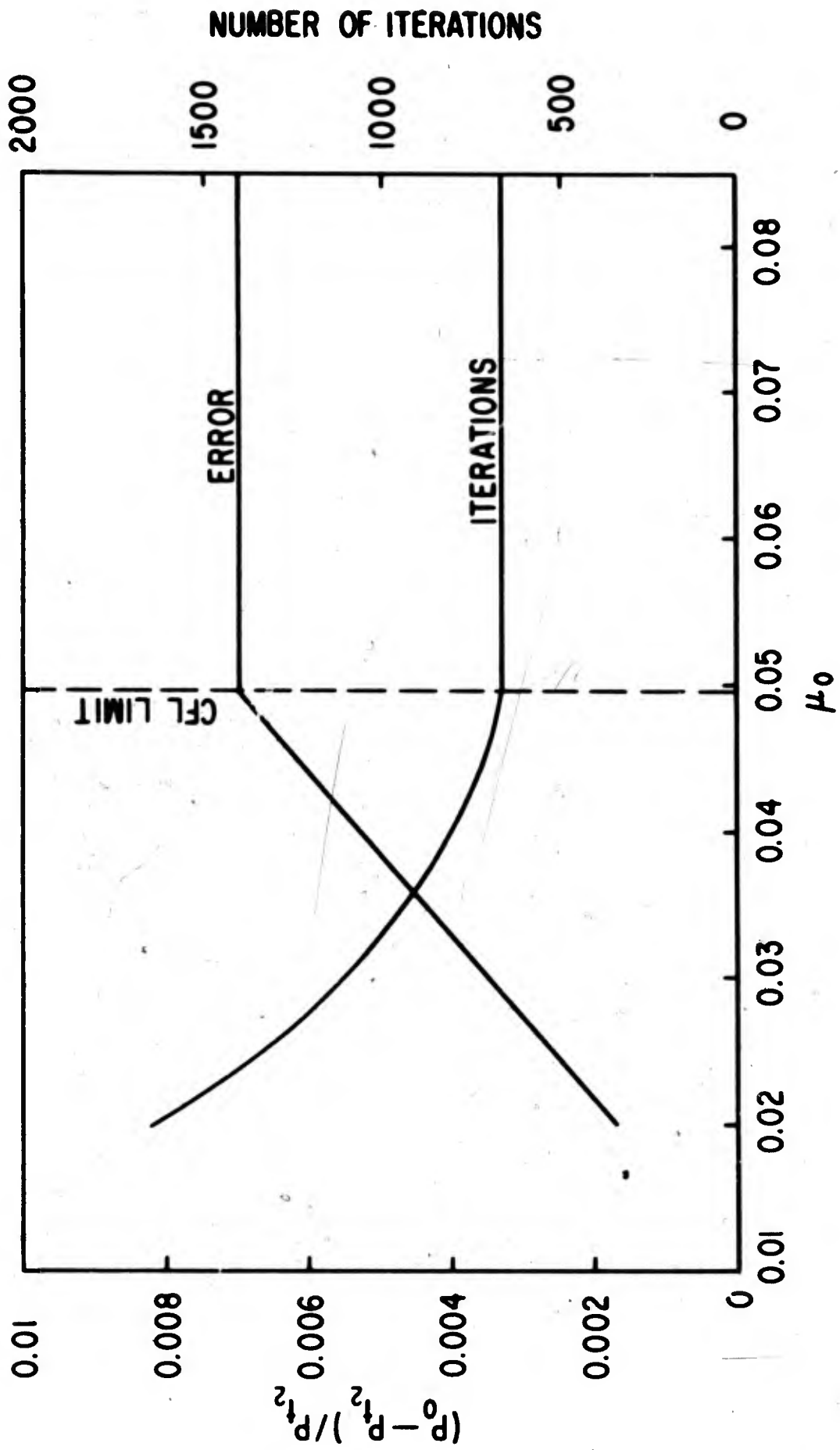


Figure 6. Investigation Of The Stabilizing Term Effects

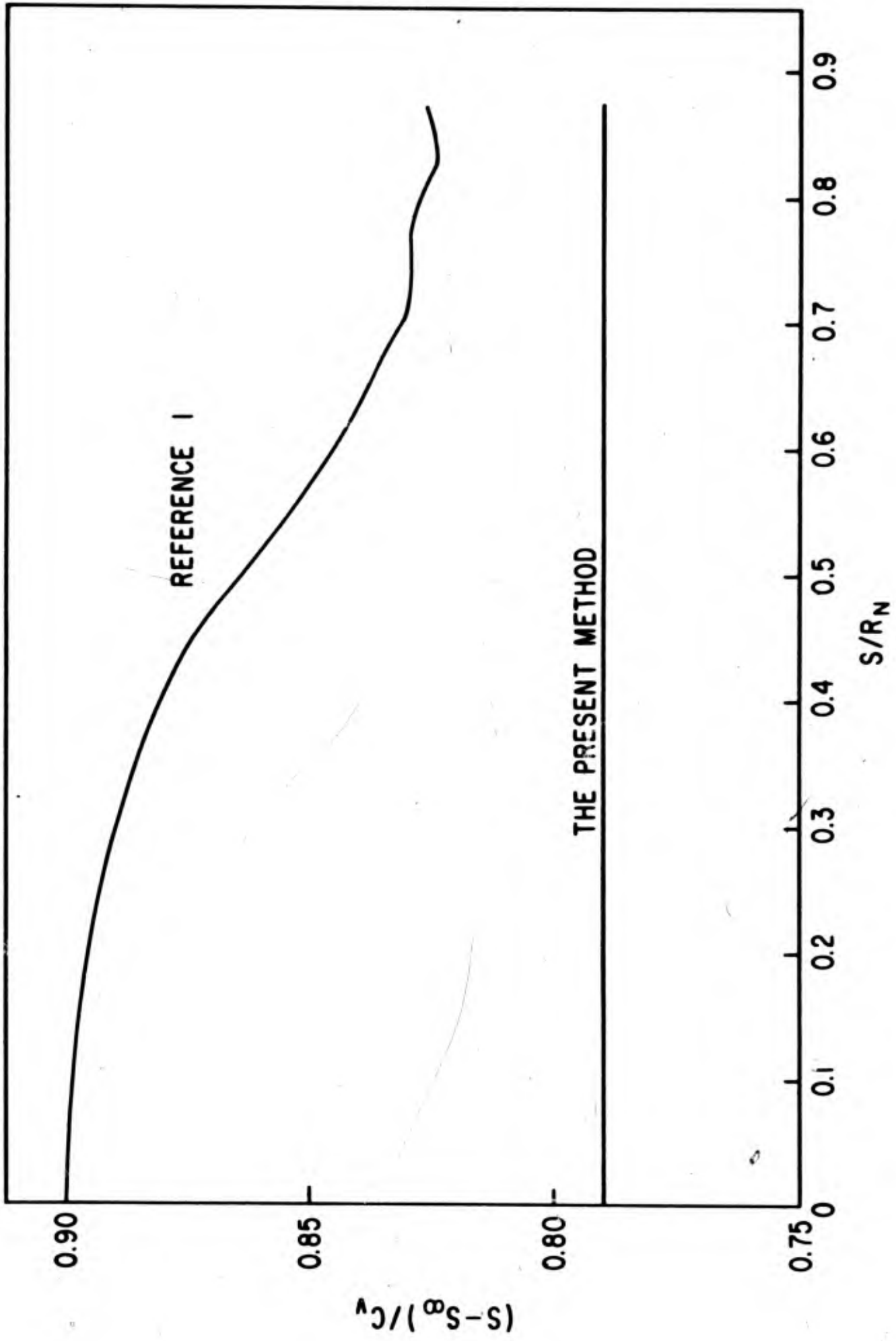


Figure 7. Surface Entropy Profiles For A Sphere In Mach 4 Flow

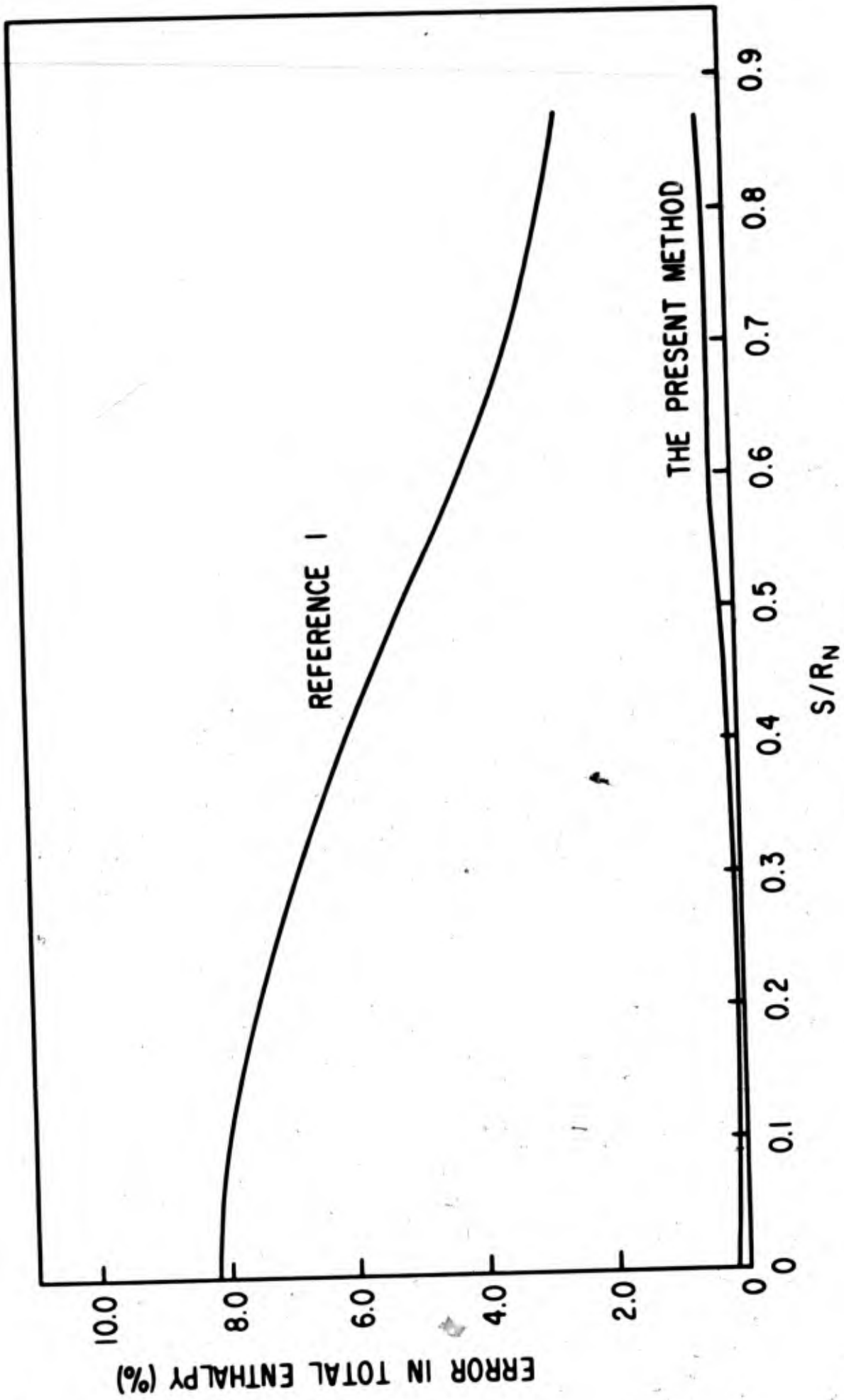


Figure 8. Surface Total Enthalpy Profiles For A Sphere In Mach 4 Flow

be achieved. Figure 9 illustrates the results of this procedure for a sphere in a Mach 30 flow. It is seen that the shock propagates out from the body and stabilizes to its steady-state profile after about 500 iterations. The stabilization of the shock shape has been used as a criterion for specifying an approximate steady state. This case demonstrated that this criterion is not sufficient since the flow field parameters had not stabilized after 500 iterations. Figure 10 illustrates the shock velocity at various times for this case. A reliable specification for achieving steady state flow was found to be the requirement that all shock velocities be less than 10^{-3} .

3. BLUNT BODY FLOW AT VARIOUS MACH NUMBERS

Many available solution procedures are extremely reliable for a limited Mach number range while experiencing difficulty at low and high Mach numbers.

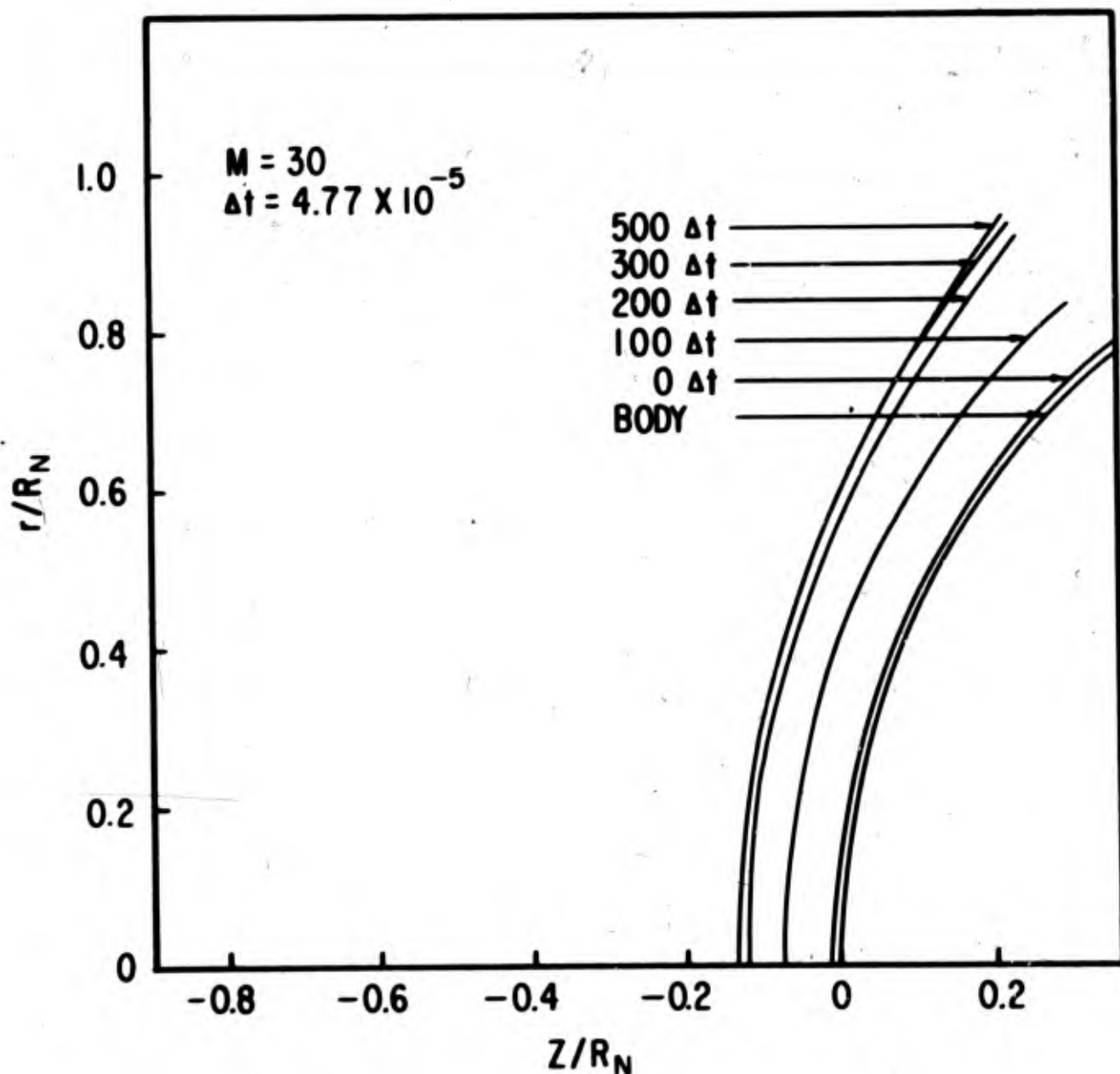


Figure 9. Shock Shapes At Various Times For Unsteady Flow Over A Sphere

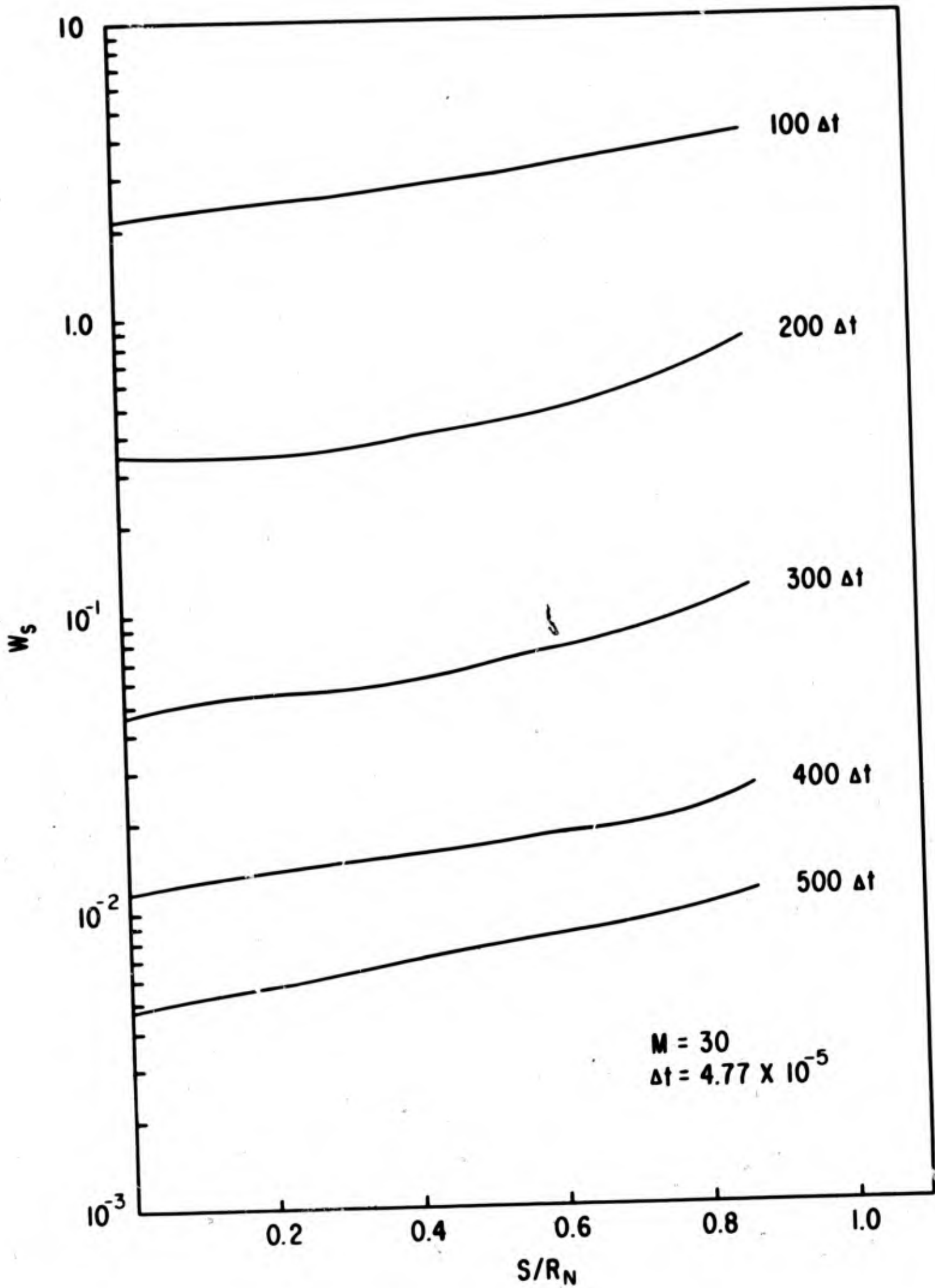


Figure 10. Shock Velocity At Various Times For Unsteady Flow Over A Sphere

Figure 11 illustrates shock shapes predicted for flow over a hemisphere at Mach numbers ranging from 1.5 to 500. No difficulties were encountered for any of these cases. Shock shapes predicted by the method of reference 13 are also shown. Since the shock point calculations for the present method are based on the method of reference 13, it is useful to ensure that both give the same results.

4. VALIDATION OF THE METHOD FOR BLUNT BODY FLOWS

To establish the accuracy of the present method for blunt body flows, the 3 strip integral relations solution of Belotserkovskii (data taken from reference 16) was used as a standard. Figure 12 illustrates the surface pressure distribution for a sphere in a Mach 4 flow. It is seen that the present method and Belotserkovskii's solution are in excellent agreement. Also shown are several time-dependent solutions reported in the literature. The solution by Moretti and Abbett (Ref. 13) is in excellent agreement with the present method in the subsonic region but predicts higher surface pressures in the supersonic region than either the present method or the integral relations solution. Solutions by the Godunov method reported by Godunov, Zabrodin and Prokopov (Ref. 9) and Masson (Ref. 16) are also shown. The solution by Masson is seen to agree very well with Belotserkovskii's results and the present method. It is concluded that the present method and Masson's method provide the best predictions of the time-dependent methods presented. Figure 13 illustrates the surface density distribution predicted by the present method, Belotserkovskii and Masson. Agreement between the present method and Belotserkovskii's solution is excellent. It is seen that Masson's solution does not predict the correct density distribution. It is concluded that the accuracy of the present method is superior to other reported time-dependent methods and equivalent to the method of integral relations.

5. AFTERBODY FLOWS

To validate the present method for afterbody flows, comparison with the NASA inverse scheme-method of characteristics program (Ref. 3) was accomplished. Since the surface entropy has been shown to be constant and correct, comparison of the predicted surface pressure distributions is sufficient to establish validity. Figures 14, 15 and 16 show the comparison of predicted surface pressures for sphere-cone vehicles with various cone half-angles and Mach numbers of 4, 6 and 10. Agreement between these methods is seen to be excellent.

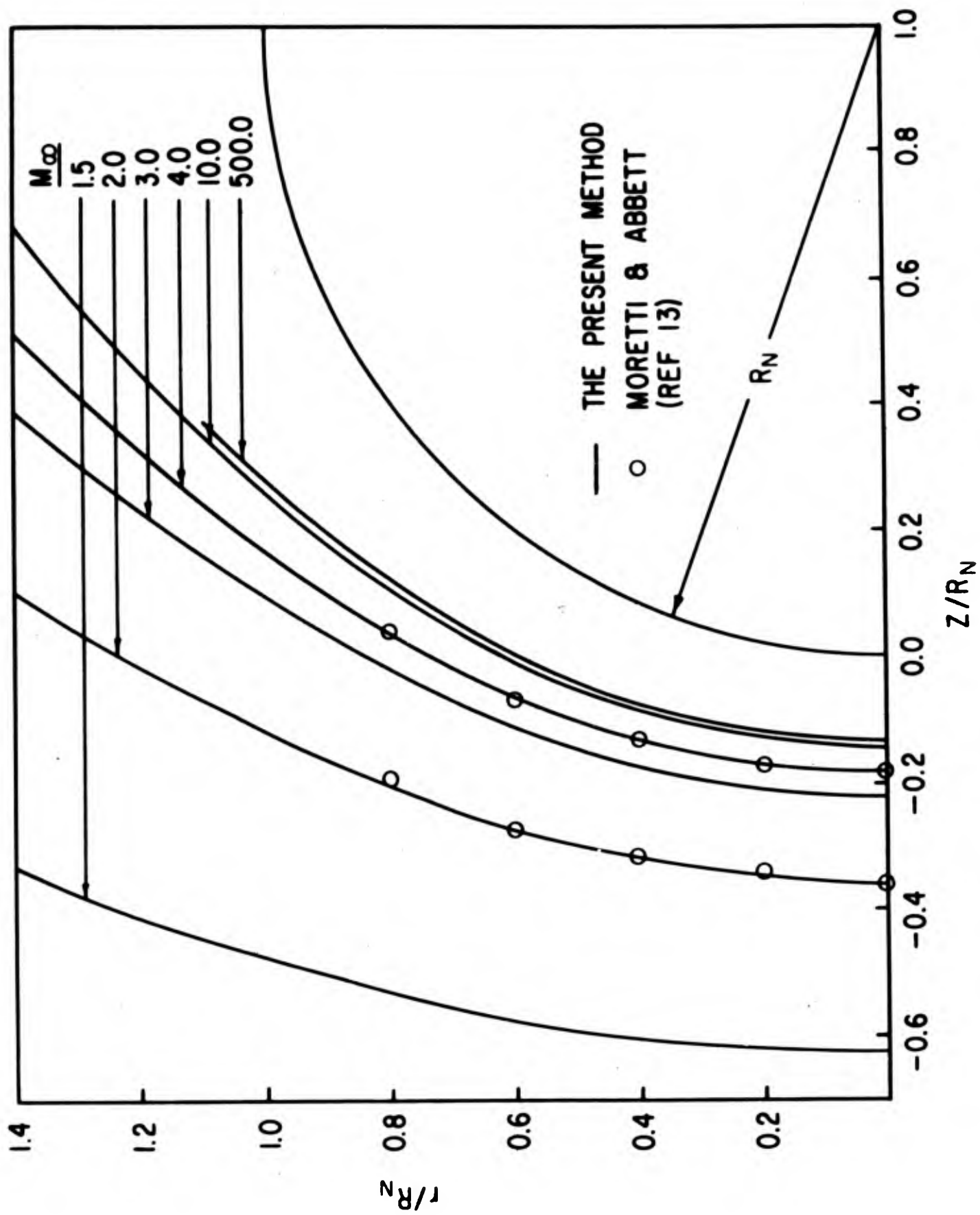


Figure 11. Shock Shapes For Flow Over A Hemisphere At Several Mach Numbers

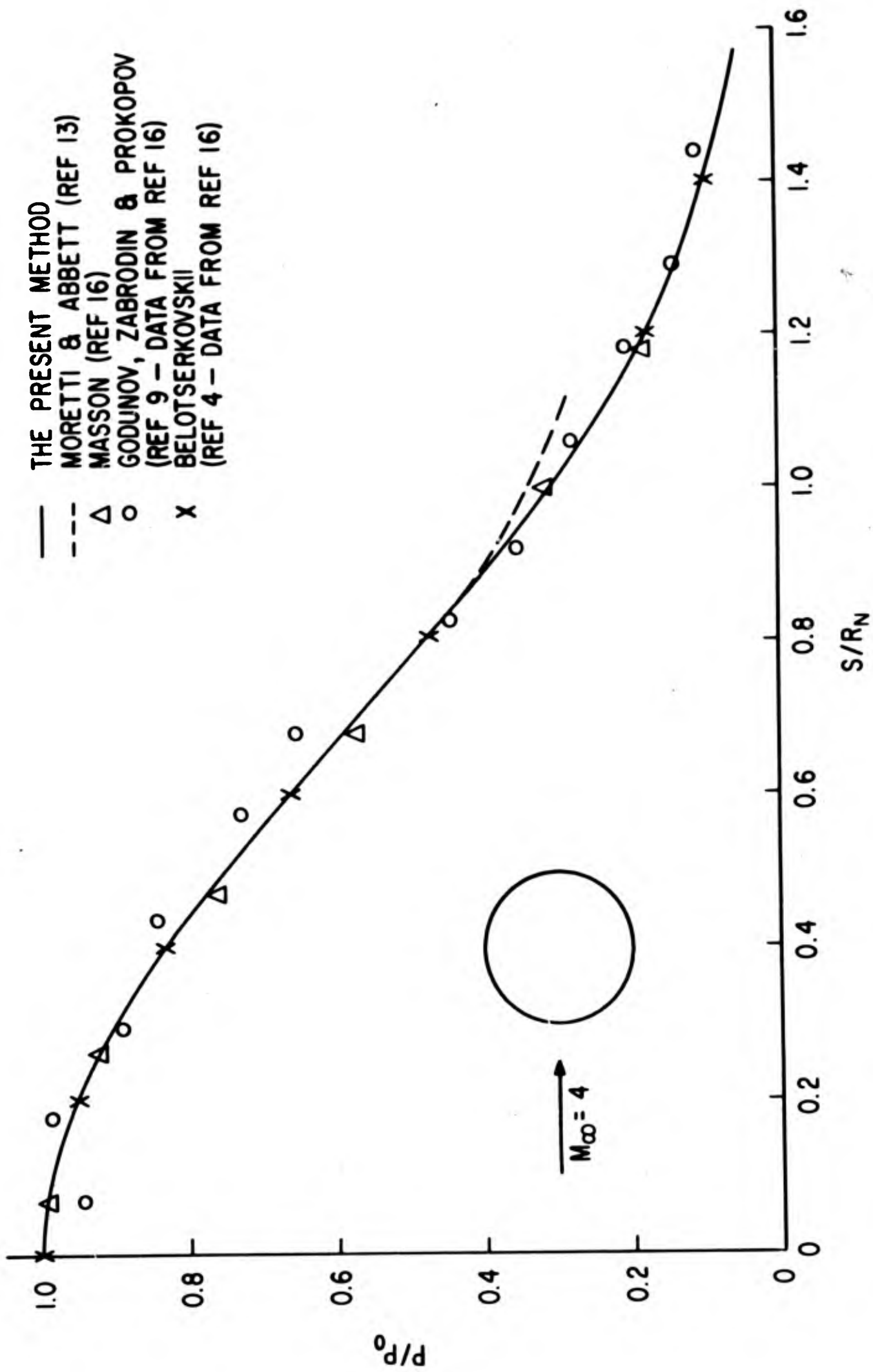


Figure 12. Comparison Of Time-Dependent Methods With An Integral Relations Solution

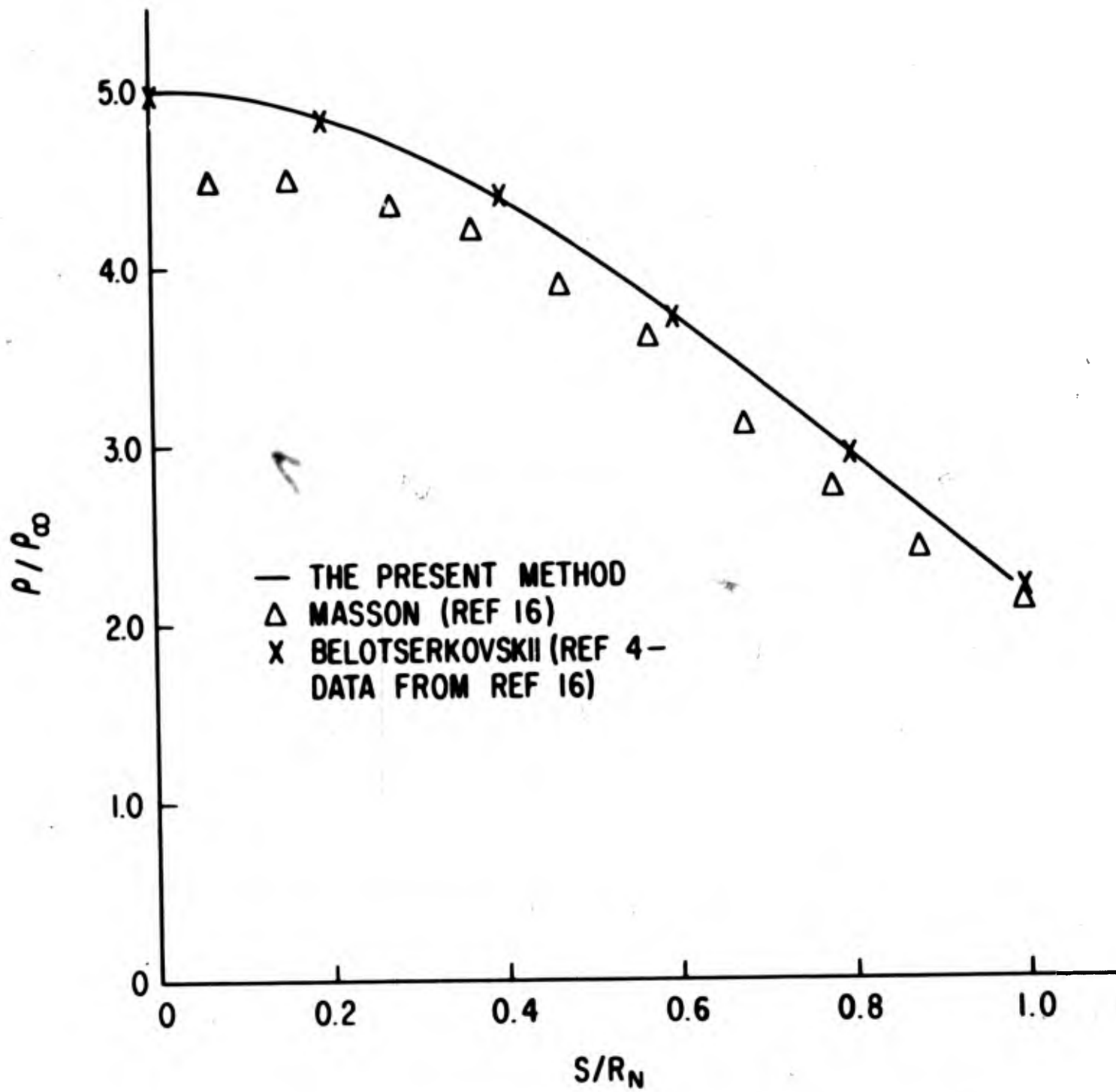


Figure 13. Surface Density Distributions For A Sphere In Mach 4 Flow

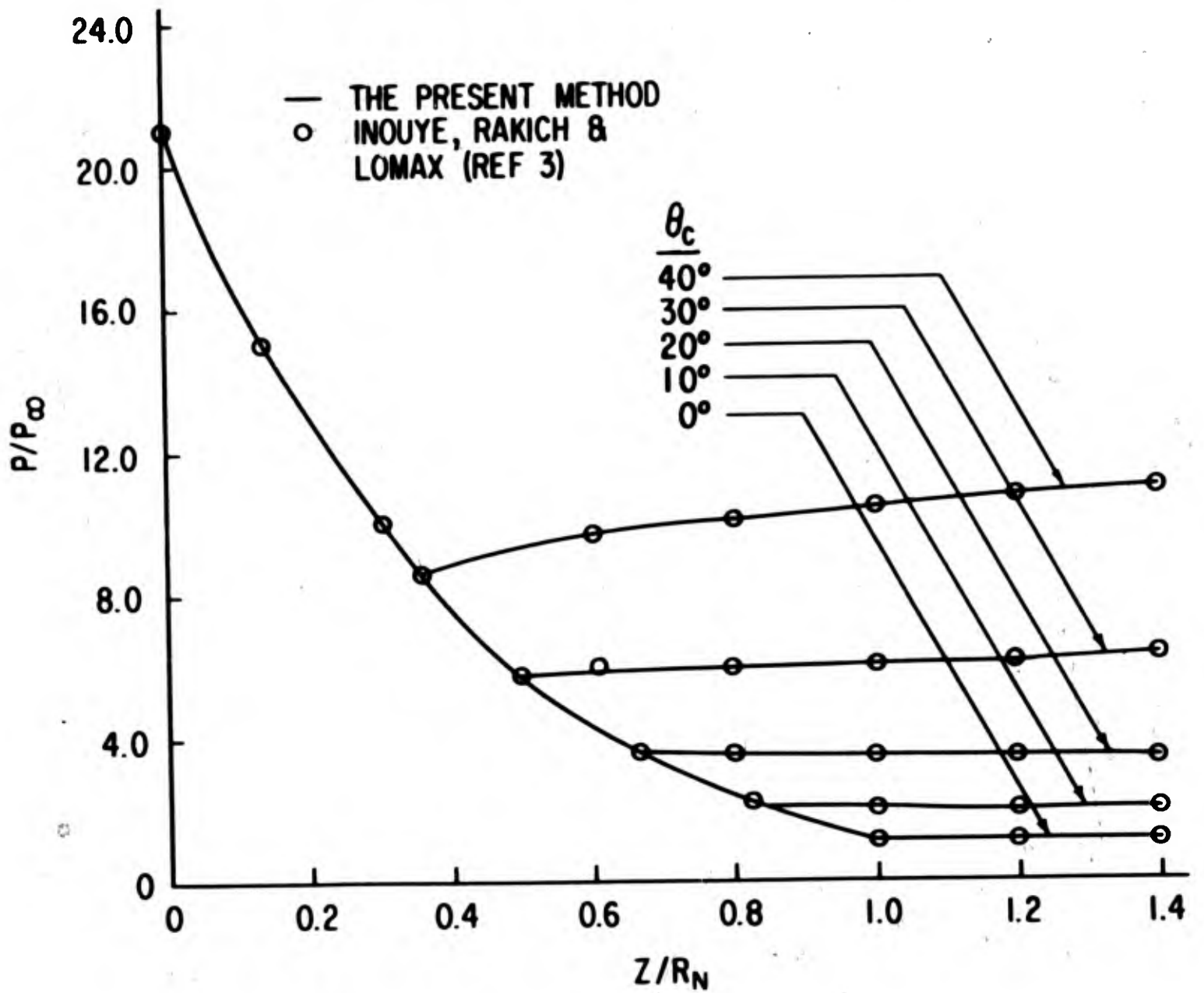


Figure 14. Surface Pressure Distributions For Sphere-Cones In Mach 4 Flow

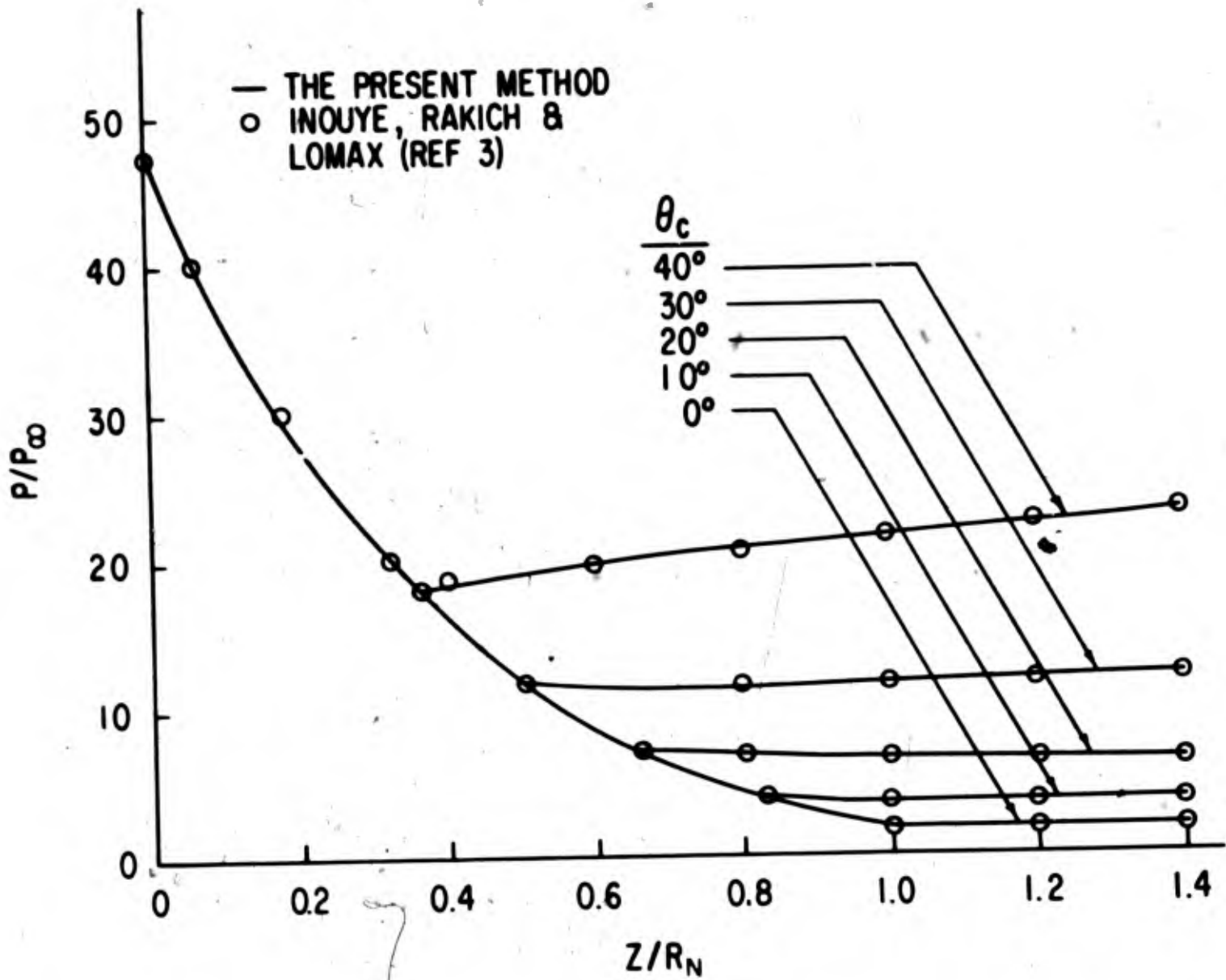


Figure 15. Surface Pressure Distribution For Sphere-Cones In Mach 6 Flow

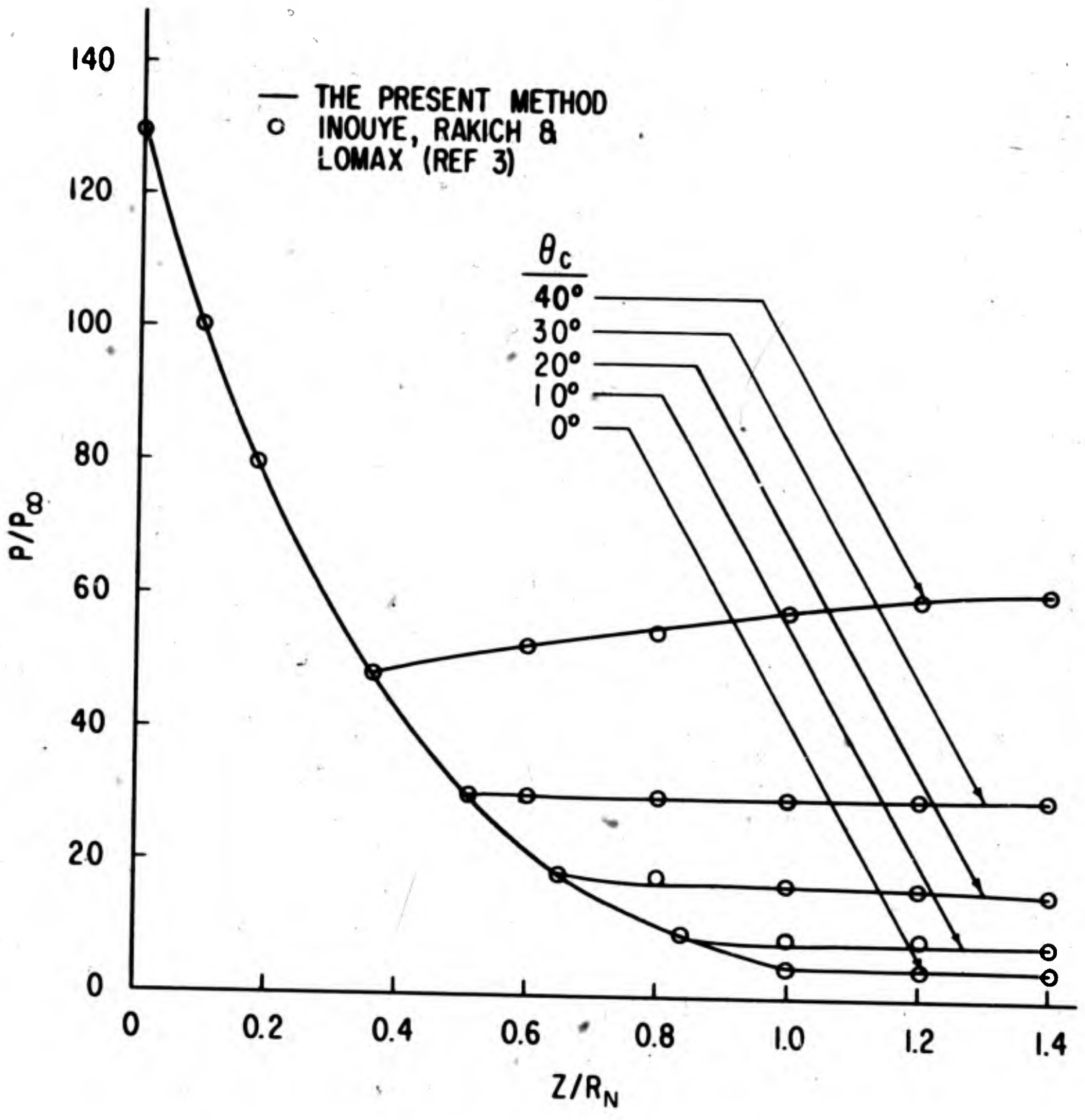


Figure 16. Surface Pressure Distribution For Sphere-Cones In Mach 10 Flow

6. ARBITRARY BODY SHAPES

The previous solutions for sphere-cone vehicles employed internal relations in the computer code to define the body geometry. To validate the numerical computations used to define an arbitrary body shape from input data, ellipsoid-cylinder vehicles were considered for comparison with reference 3. Figure 17 shows a comparison of surface pressures between these methods. It is concluded that the arbitrary body shape logic provides accuracy equivalent to reference 3 for those cases where comparison is possible.

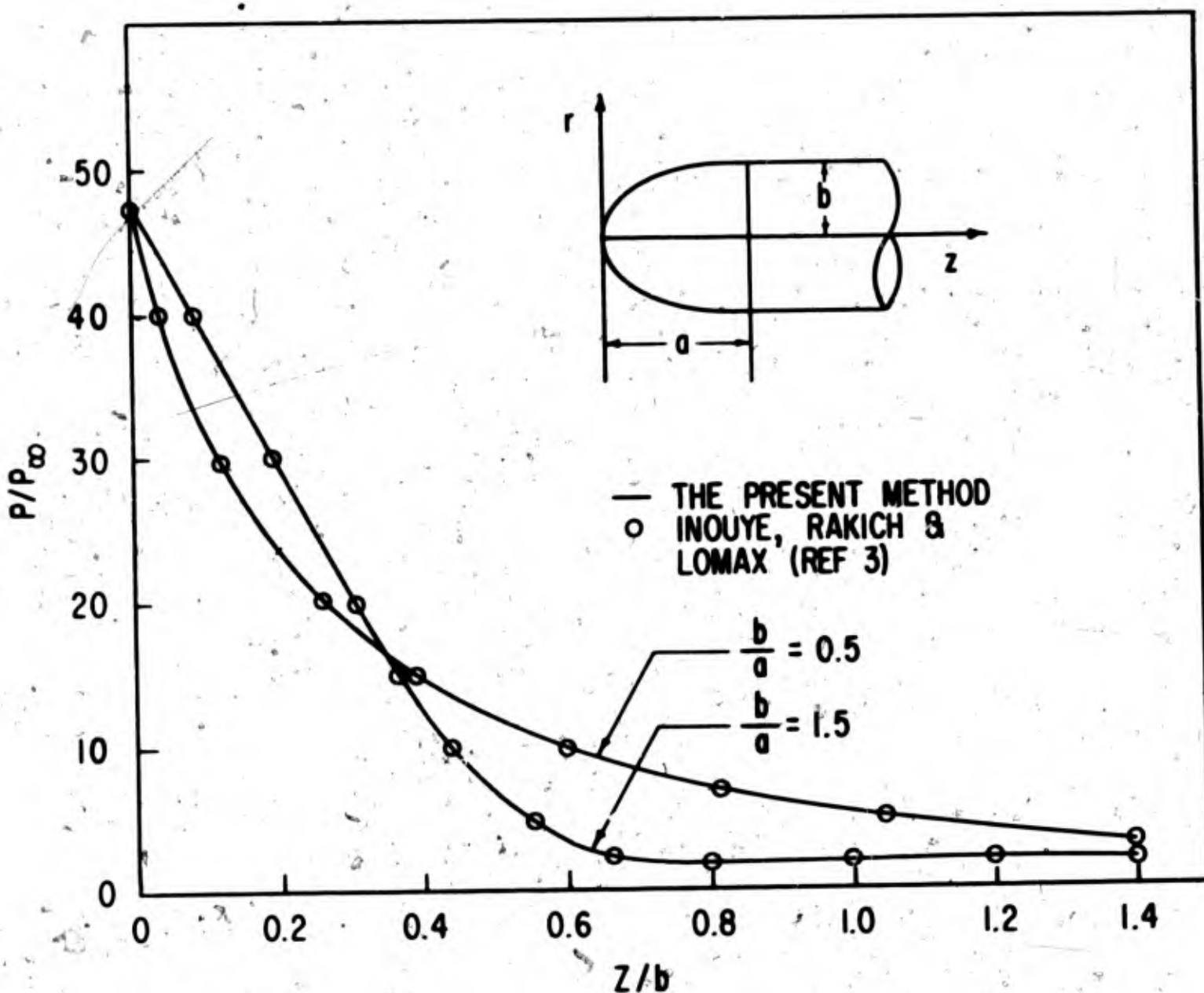


Figure 17. Surface Pressure Distributions For Ellipsoid-Cylinders In Mach 6 Flow

7. APPLICATION TO TYPICAL REENTRY VEHICLES

Since the accuracy of the present method has been established for the relatively simple problems normally used to demonstrate the capability of a computational method, it remains only to attempt a practical reentry problem. The AFWL version of the NASA computer code of reference 3 and the present method were used to generate a solution for a sphere-cone vehicle with a 10° cone half-angle in a flow with a Mach number of 20. The solution was carried downstream 50 nose radii measured along the symmetry axis. The surface pressure distribution comparison is shown in figure 18. Some disagreement is seen in the strong expansion region. Additional reduction of the stabilizing terms and a solution not using the segmented approach were attempted but had no significant effect on the predictions by the present method. It was concluded that this minor disagreement could not readily be explained. A comparison of the surface Mach number predictions is shown in figure 19. Except for the local disagreement in the strong expansion region, the solutions appear to be equivalent. To further investigate the local disagreement, a comparison of surface pressure distribution predicted by the NASA program and the present method are compared to experimental data reported in reference 19. Figure 20 shows that the present method agrees better with the experimental data than the NASA solution. This closer agreement is undoubtedly fortuitous, but does indicate that the local disagreement between the two computational methods is insignificant when compared to experimental data. The solution shown in Figures 17 and 18 represents the only known solution for a practical high performance reentry vehicle generated with a time-dependent method. The solution was generated in less than 6 minutes on the CDC 6600 computer and required only 40,000 octal locations for storage. As previously noted, the computation time required to treat this problem without the segmented approach is estimated to be in excess of 70 hours on the same computer.

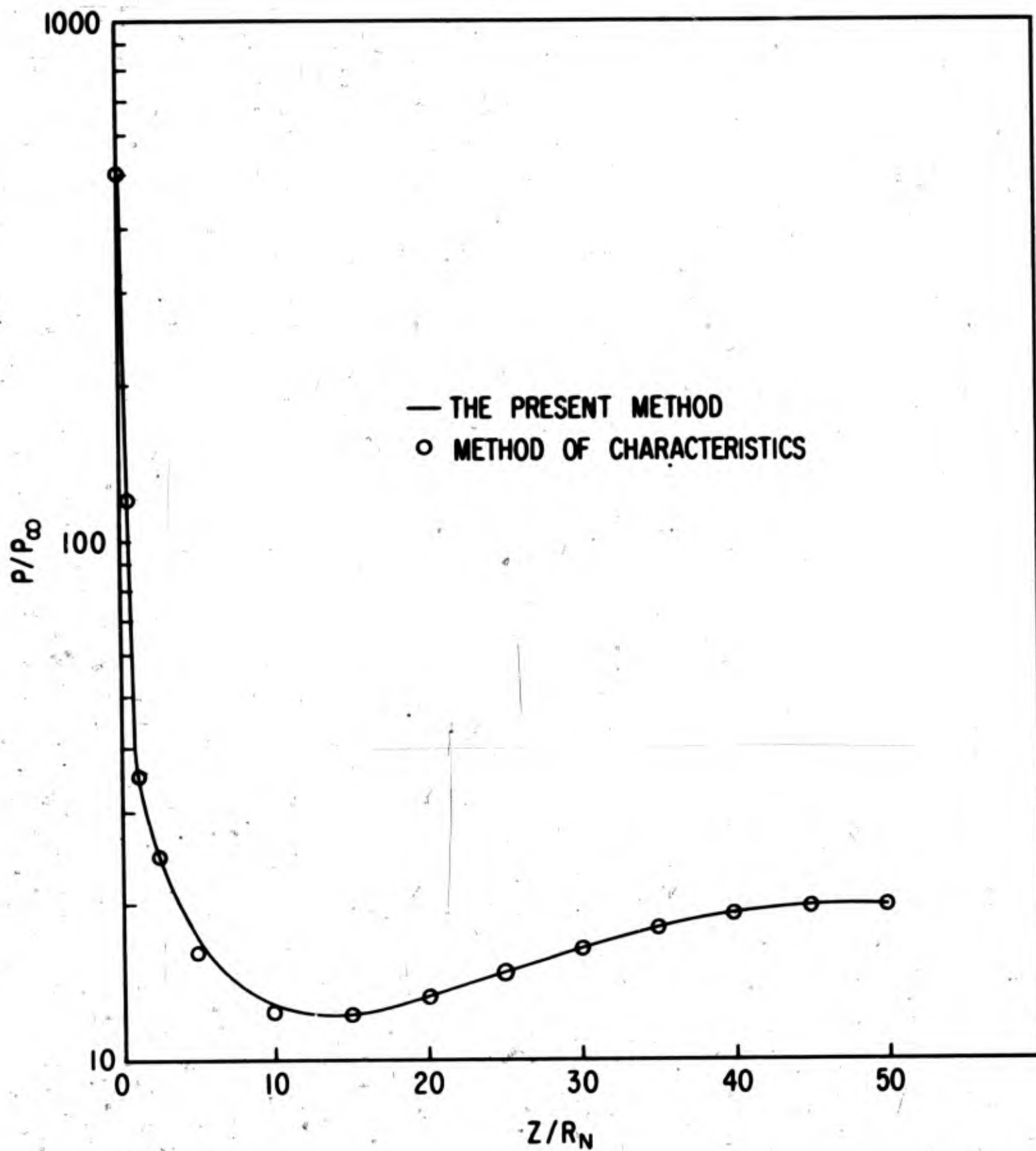


Figure 18. Surface Pressure Distribution For A 10° Sphere-Cone In Mach 20 Flow

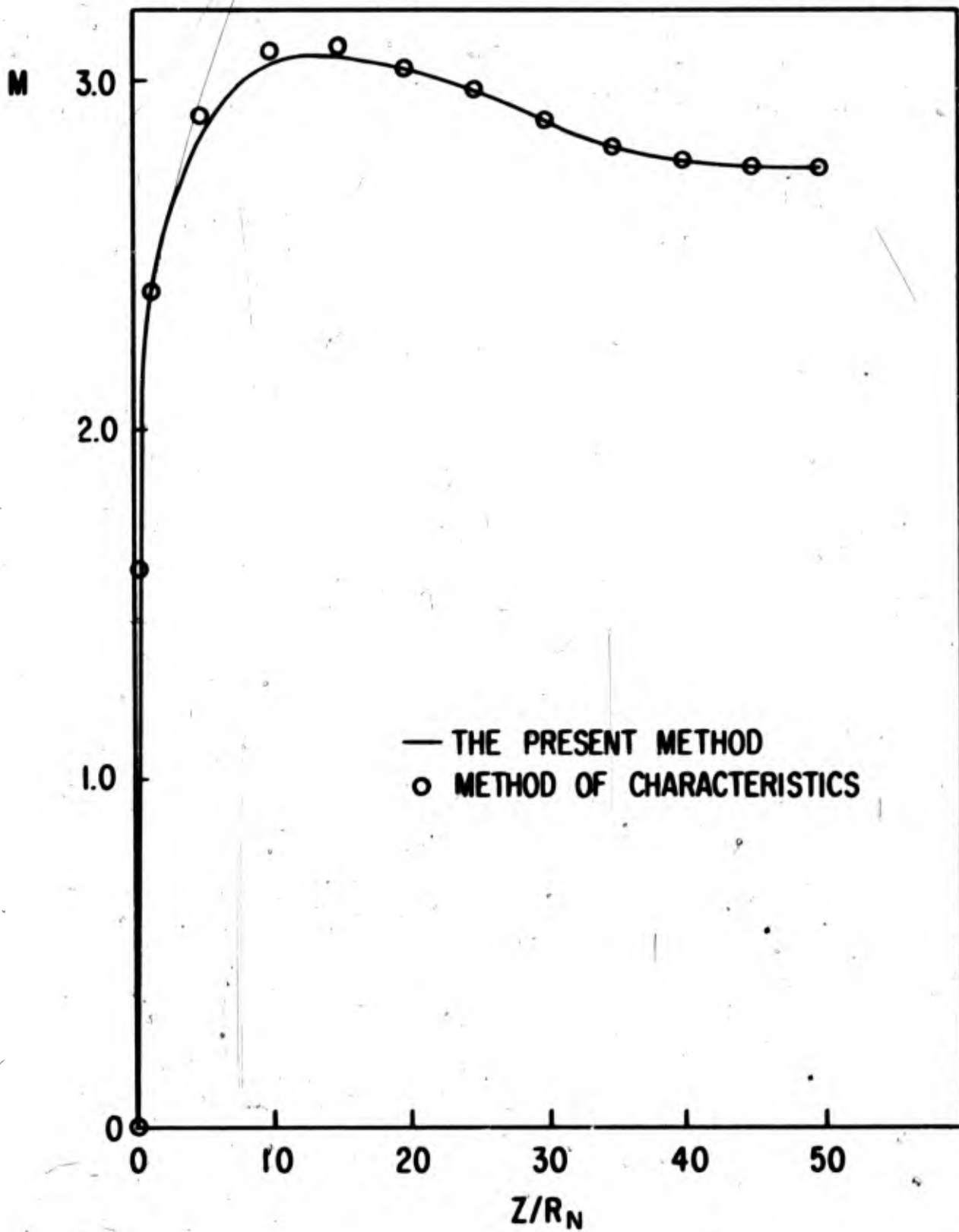


Figure 19. Surface Mach Number Distribution For A 10° Sphere-Cone In Mach 20 Flow

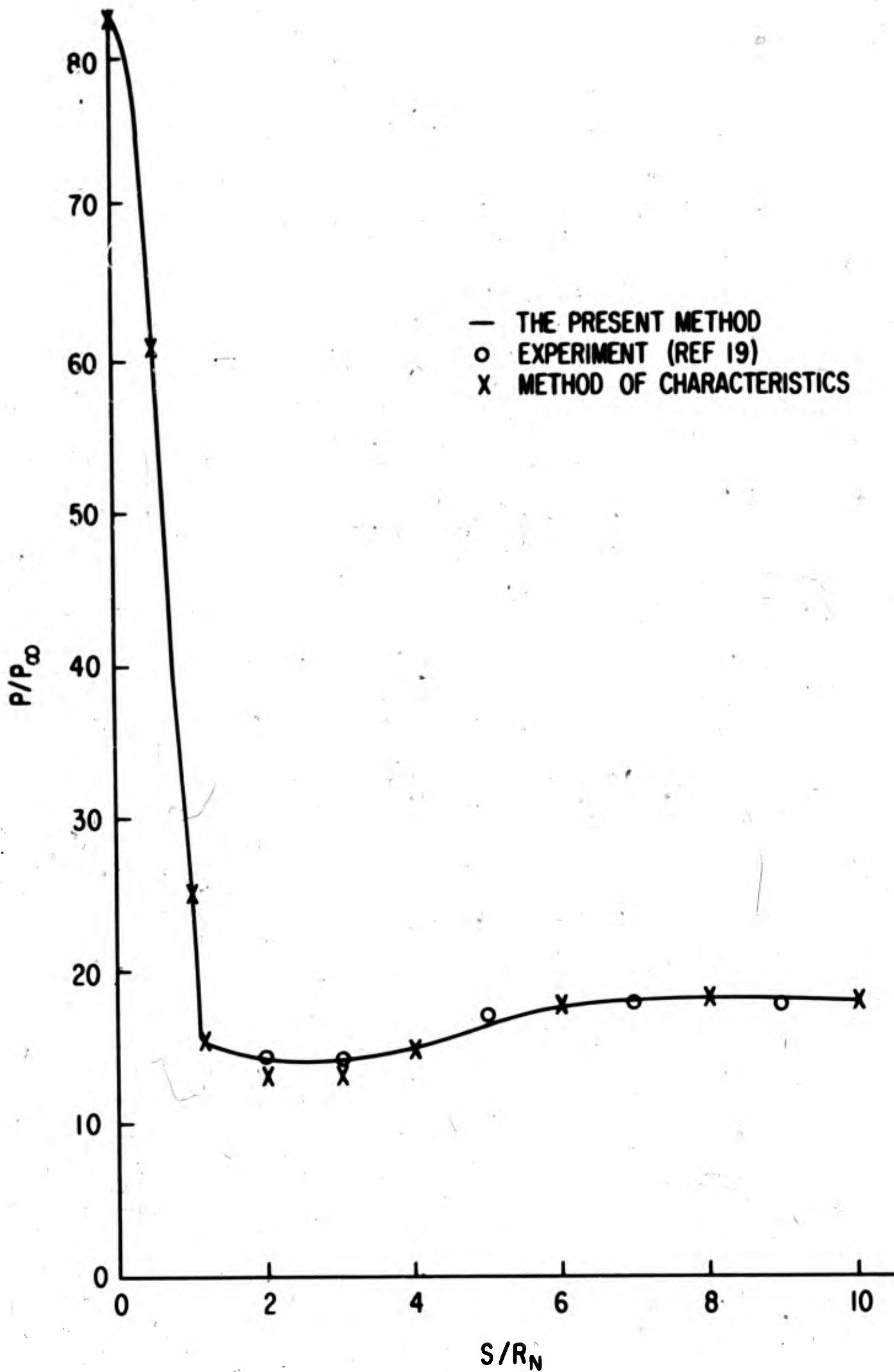


Figure 20. Surface Pressure Distribution For A 25° Sphere-Cone In Mach 8 Flow

SECTION VI

THE FORTRAN IV COMPUTER CODE

1. MACHINE ROUTINES

The computer code employs 3 machine routines available on the CDC 6600 computer. When a different computer is used, these routines may not be available. Substitute routines should be provided, as required, for the following routines:

INT (X) - a FUNCTION routine which converts the floating point variable, X, to integer form.

FLOAT (I) - a FUNCTION routine which converts the integer variable, I, to floating point form.

SKFILE (FILE, N) - a SUBROUTINE which skips N files on a magnetic tape unit specified by the Hollerith argument FILE. To simplify the replacing of this routine, calls to it are contained only in the dummy routine, SUBROUTINE SKIP.

2. DIMENSIONED VARIABLES

The variable dimensions currently contained in the computer code allow up to 15 nodes along the body and 40 nodes across the shock layer. To change these dimensions, new values for IMAX and NMAX should first be supplied in a DATA statement in the main program, where

IMAX - specifies the maximum number of nodes along the body permitted by the dimension sizes.

NMAX - specifies the maximum number of nodes across the shock layer permitted by the dimension sizes.

Then, the following COMMON statements should be modified consistent with the values of IMAX and NMAX:

```
COMMON/1/ V(IMAX,NMAX,4), D(NMAX,4), P(3,NMAX)
```

```
COMMON/2/ RB(IMAX), ZB(IMAX), TH(IMAX), CTH(IMAX), K(IMAX),  
NS(IMAX), WS(IMAX)
```

```
COMMON/11/ MAXIT, ISTOP, ITD, TIME, SI(IMAX), NST(IMAX)
```

```
COMMON/SPLINE/ C(4,IMAX)
```

SUBROUTINES INTERP and INTEG contain a dimensional variable which should be dimensioned X(IMAX).

3. THE MAIN PROGRAM

A. PURPOSE

The main program is designated as PROGRAM BLUNT. It's major functions are:

- a) Initialize constants and some of the case data variables.
- b) Define values for the control variables to select the proper options.
- c) Accomplish checks to determine if any errors were sensed while reading the input data. When errors are sensed, an error message is printed, the case terminated, and the next case started.
- d) Performs all flow field computations except for points on the shock and on the body surface.
- e) Schedules the flow field data output as requested by the user.
- f) Schedules the storing of the flow field data on magnetic tape.

B. VARIABLE LIST

- AMACH1 - M_∞
- AMUO - μ_0 , equation (62)
- AVN - parameter used to introduce the stabilizing terms relative to the n coordinate.
- AVS - parameter used to introduce the stabilizing terms relative to the s coordinate.
- Al - A, equation (49), evaluated at the shock.
- ENOSE - nose bluntness parameter
- B1 - B, equation (49), evaluated at the shock.
- CDEL - $\cos(\beta - \theta)$
- CN - $(\kappa \rho V)_n$
- CS - $2(\rho u)_s$
- CTH - array of $\cos \theta$ values
- C1 - $u/(2\Delta\xi)$
- C2 - $-\frac{\eta_s}{\partial s} \frac{u}{\chi(v-w_s \eta)}$
- C3 - $\chi(v-w_s \eta)/(2n_s \Delta\eta)$
- C30 - $\frac{1}{2n_s \Delta\eta}$
- D - array used for temporary storage of flow field parameters at $r + \Delta r$.
- DBET - $\beta - \theta$
- DELS - length of a body segment measured along the surface.
- DET - $\Delta\eta$
- DE2 - $2\Delta\eta$
- DKAP - $\Delta\kappa = K\Delta\eta$
- DN - Δn
- DR - $\Delta r = \Delta n \cos \theta$

DRSM - r at $\xi - \Delta\xi$
 DRSP - r at $\xi + \Delta\xi$
 DT - Δr
 DTO - $0.9 / (q_{\bullet} + a_{\bullet})$. Used for CFL stability check.
 DT1 - the value of Δr evaluated from equation (62).
 DUM - dummy variable used for intermediate computations.
 DX1 - $\Delta\xi$
 DX2 - $2\Delta\xi$
 DX3 - $3\Delta\xi$
 ERR - array specifying acceptable errors in convergence on n and P for shock iteration procedure.
 ETA - η
 GAMA - $(\gamma+1)/2$
 GAMB - $(\gamma-1)/2$
 GAML - γ
 HT1 - total enthalpy.
 I - index specifying body station in DO loops.
 IGEOM - control variable. Specifies method for selecting the number of streamwise nodes to be used.
 IM - $I - 1$
 IMAX - maximum number of streamwise nodes allowed by the dimension sizes.
 IMIN - value of I at which computations should start on the current segment.
 INTAP - the magnetic tape unit number to be used for flow-field input (when required).
 IP - $I + 1$
 IPR - array specifying iteration numbers or times when flow field output is to be provided.
 IPRINT - iteration number when next flow field output is to be provided.
 ISEG - segment number.
 ISTOP - error indicator.
 IT - index on major iteration DO loop.
 ITD - total number of iterations performed for the current segment.
 JPR - array used for temporary storage of IPR.
 K - array containing values of K .

KAP - κ
KBAR - $(K(\xi+\Delta\xi) + 2K(\xi) + K(\xi-\Delta\xi))/4$
KGEOM - control variable defining the option to be used for geometry calculations.
KSEG - index used to suppress the check on NSEG (see NSEG).
KZERO - body curvature at $S = 0$. Used to nondimensionalize all length parameters.
L - specifies first index in the P array corresponding to I.
LM - specifies first index in the P array corresponding to IM.
LP - specifies first index in the P array corresponding to IP.
LPLUS - array used to specify LP. $LP = LPLUS(L)$.
M - dummy index.
MAXI - total number of streamwise nodes employed to treat current segment.
MAXIT - total number of iterations to be performed.
MAXN - the total number of nodes across the shock layer.
MI - $MAXI - 1$
MN - $MAXN - 1$
M1 - used for temporary storage of M_{∞} .
N - index specifying nodes across shock layer in DO loops.
NCASE - case number (see section VII).
NFILE - number of the first segment to be treated (see section VII).
NGEOM - control variable specifying method for selecting MAXN.
NM - $N - 1$
NMAX - total number of nodes across the shock layer allowed by dimension sizes.
NMIN - $MAXN - 2$
NOPT - control variable specifying method to be used in selecting MAXIT.
NP - $N + 1$
NPR - index used to select proper value of IPRINT from the IPR array.
NS - array containing values of n_s .
NSEG - total number of segments to be treated.
NTAP - control variable specifying the method to be used for assigning initial conditions.
OTAP - magnetic tape unit number on which flow field data should be stored.

P - array containing values of P at the 3 body stations required for local computations.

PI - π

PI2 - $\pi/2$

PN - $\frac{\partial P}{\partial n}$

PS - $\frac{\partial P}{\partial s}$

PT2 - total pressure behind a normal shock.

P1 - P_{∞}

P90 - empirical constant used to assign initial conditions.

Q1 - q_{∞}

R - r

RB - array containing values of r at the body stations.

RGAS - gas constant.

RH01 - ρ_{∞}

RN - scale factor used for dimensional output (see section VII).

RSM - r at $(\xi - \Delta\xi)$

RSP - r at $(\xi + \Delta\xi)$

SDEL - $\sin(\beta - \theta)$

TCHECK - estimated value of r for steady state, equation (63).

TCRIT - array containing constants used in equation (63).

TDEL - $\tan(\beta - \theta)$

TH - array containing values of θ at all body stations.

THETC - cone half-angle.

TIME - r

TPR - array containing values of r when output should be provided.

T1 - T_{∞}

V - array containing values of ρ , u, v and S at all nodes.

VA - array containing the values of averages used to introduce the stabilizing terms.

VIS - 1.0 - AVS - AVN. Used to compute averages stored in VA.

VK - Kv

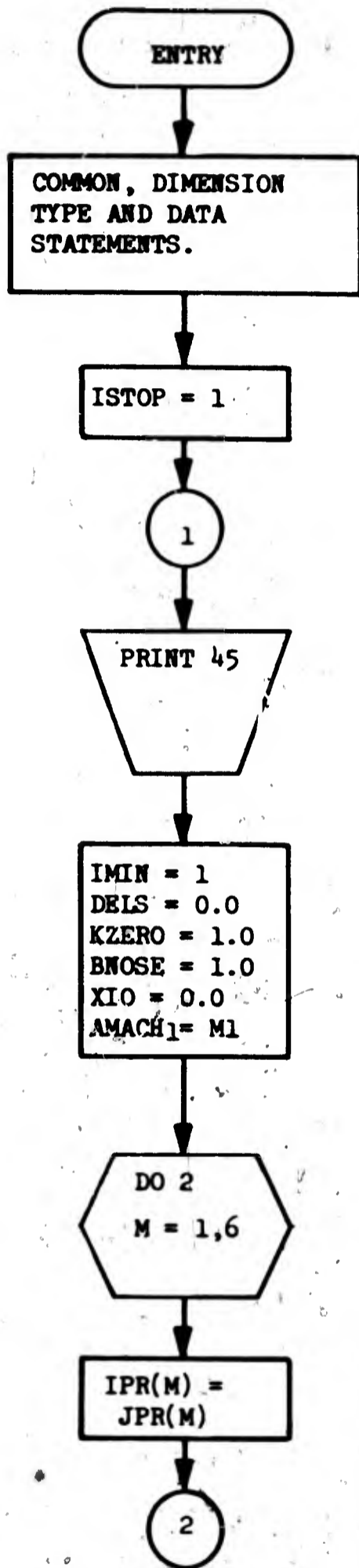
VMIN - array used to store streamwise derivatives required for shock computations.

VN - array containing values of derivatives with respect to n.

VS - array containing values of streamwise derivatives.

VS1 - a
VT - array containing values of derivatives with respect to r
WS - array containing values of W_s for all streamwise stations.
WSH - component of $W_s(\xi)$ normal to the body.
XIO - value of ξ at start of segment.
ZB - array containing values of z at all body stations.
ZRN - axial length (Z/R_N) for sphere-cones.
ZZ - dummy variable used for intermediate computations.

C. FLOW CHART



ENTER MAIN PROGRAM.

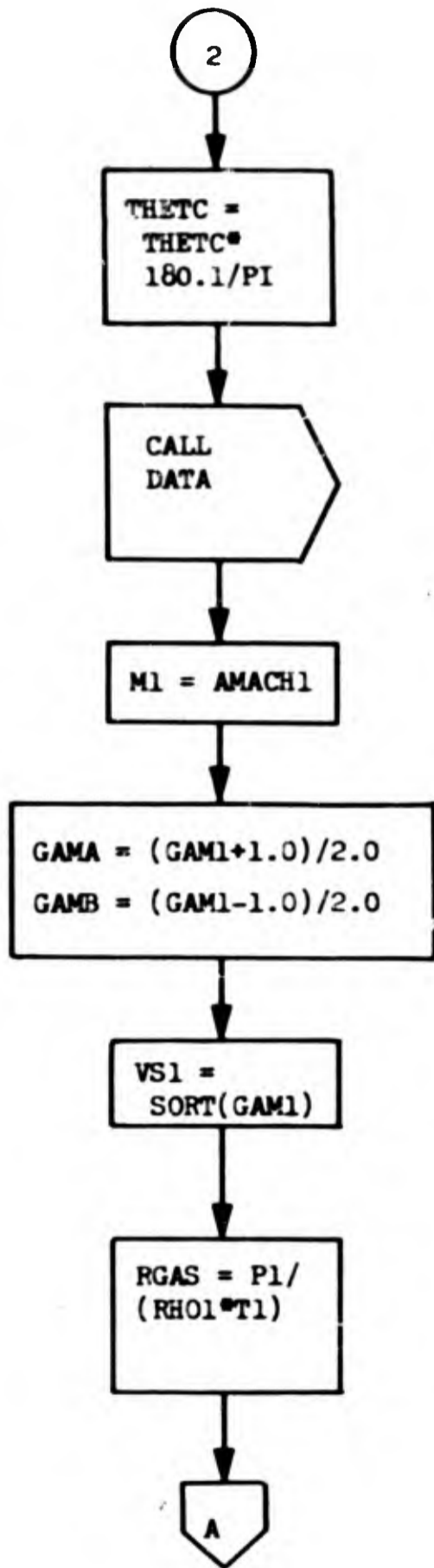
SET ERROR INDICATOR.

START POINT FOR ALL CASES.

PRINT TITLE.

INITIALIZE BODY GEOMETRY
CONSTANTS AND RECOVER THE
MACH NUMBER READ ON THE PREVIOUS
CASE.

RECOVER PRINTOUT ARRAY
READ ON PREVIOUS CASE.



CHANGE UNITS ON θ_c
TO DEGREES.

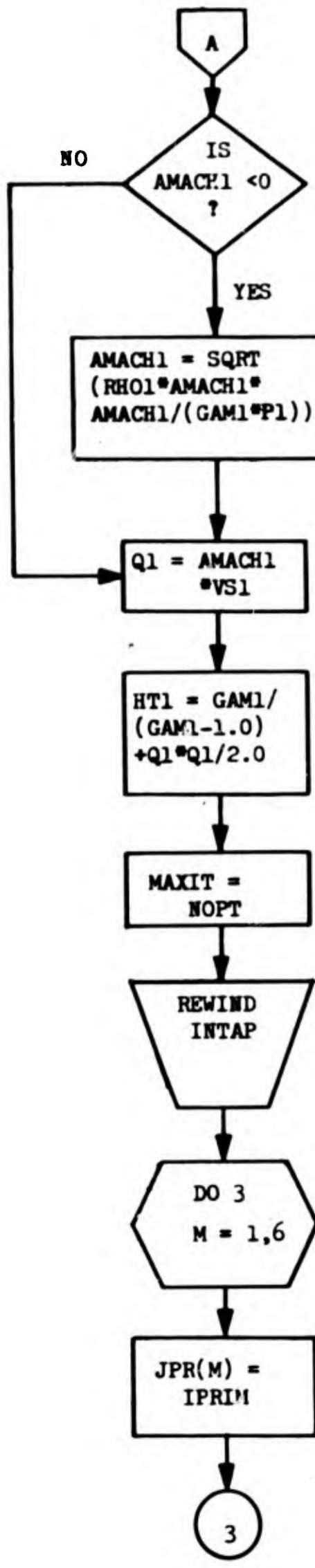
LOAD CASE DATA.

SAVE MACH NUMBER.

STORE $\frac{\gamma+1}{2}$ AND $\frac{\gamma-1}{2}$
FOR SHOCK COMPUTATIONS.

COMPUTE NONDIMENSIONAL
FREE STREAM SOUND SPEED.

COMPUTE GAS CONSTANT
FOR USE IN DATA PRINTOUT.



WAS A FREE STREAM VELOCITY
READ RATHER THAN A MACH
NUMBER ?

COMPUTE FREE STREAM MACH
NUMBER.

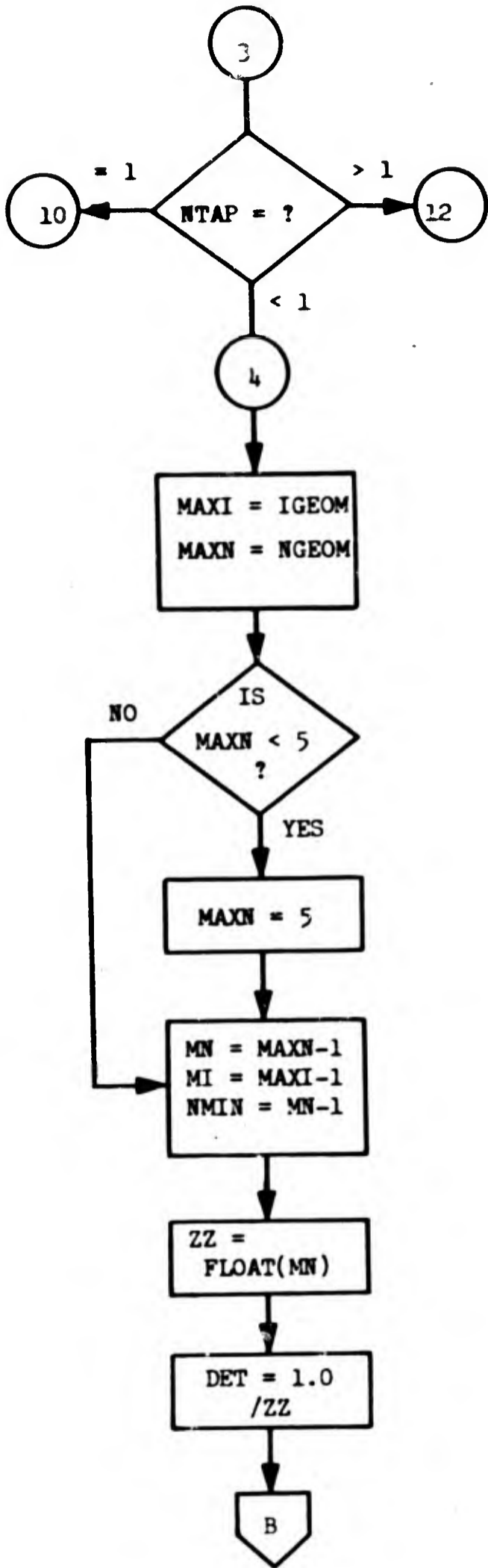
COMPUTE FREE STREAM
VELOCITY.

COMPUTE TOTAL ENTHALPY.

STORE VALUE OF MAXIT
READ IN CASE DATA.

REWIND INPUT MAGNETIC
TAPE UNIT.

STORE PRINTOUT ARRAY.



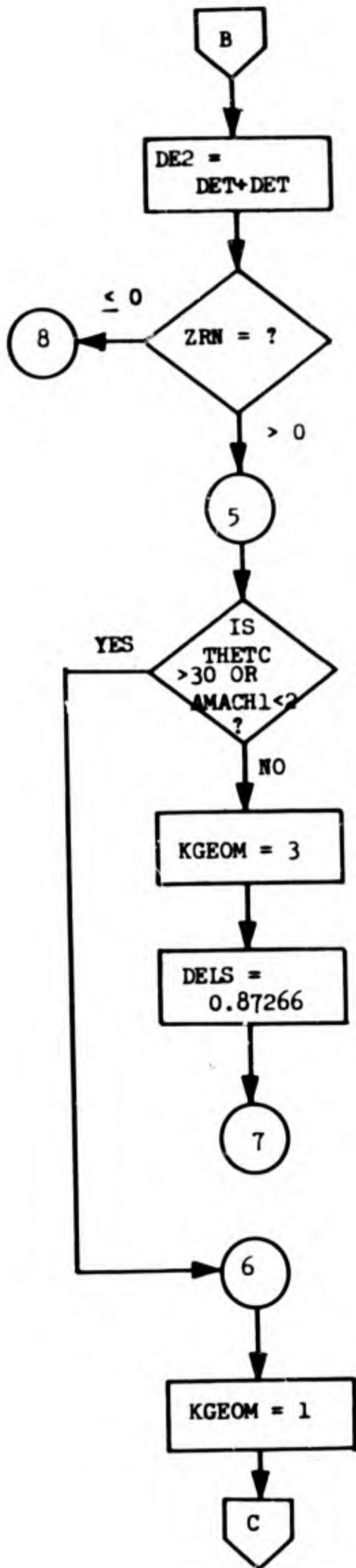
NEW CASE, RESTART OR
ADDITIONAL ITERATIONS ?

LOAD THE SPECIFIED
NUMBER OF NODES IN
THE ξ AND η DIRECTION.

MAXN < 5 ?

STORE MAXN-1, MAXI-1
AND MAXN-2.

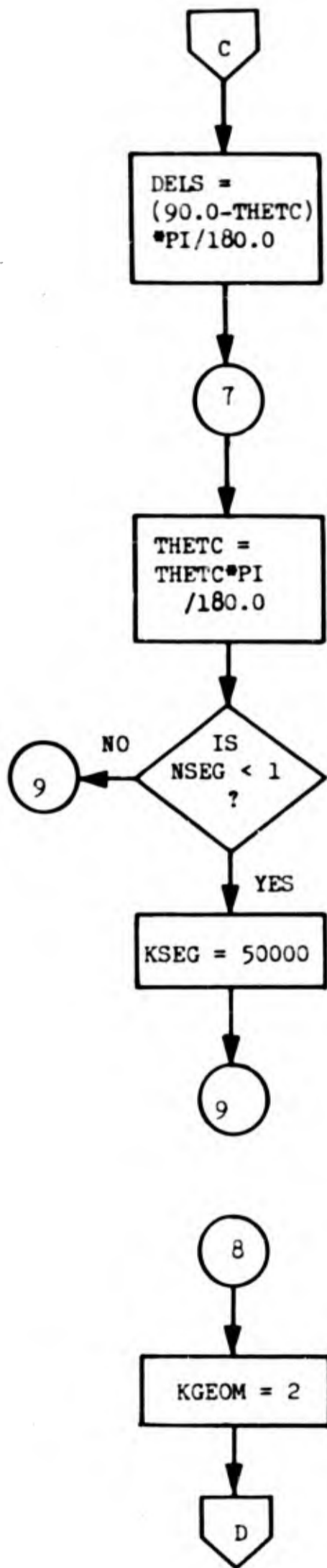
COMPUTE $\Delta\eta$.



SHOULD SPHERE NOSE BE TREATED WITH ONE OR TWO SEGMENTS ?

TWO SEGMENTS WILL BE USED.

ONE SEGMENT WILL BE USED.

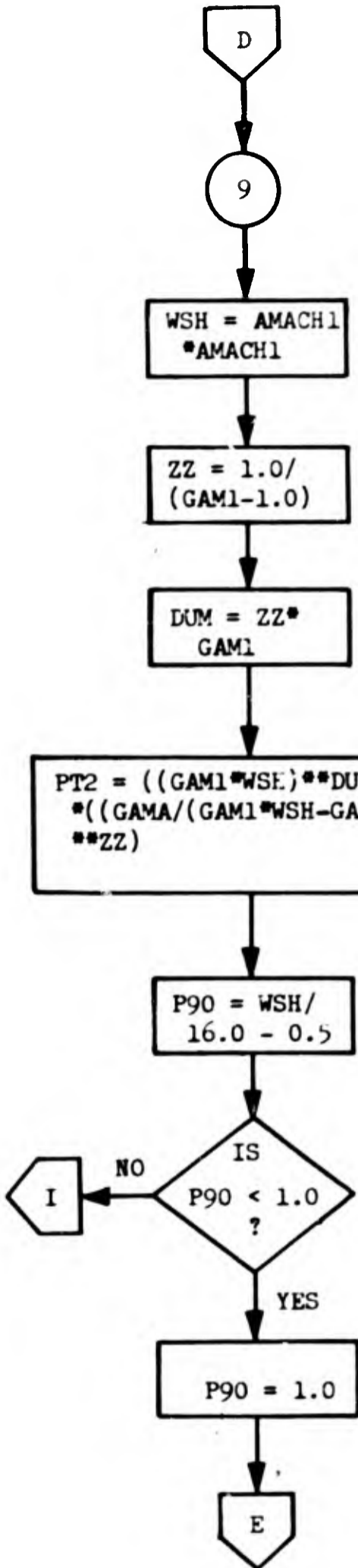


CONVERT θ_c TO RADIANS.

WAS THE NUMBER OF SEGMENTS TO BE TREATED SPECIFIED ?

NO. COMPUTATIONS WILL STOP WHEN THE TOTAL LENGTH, ZRN, HAS BEEN TREATED.

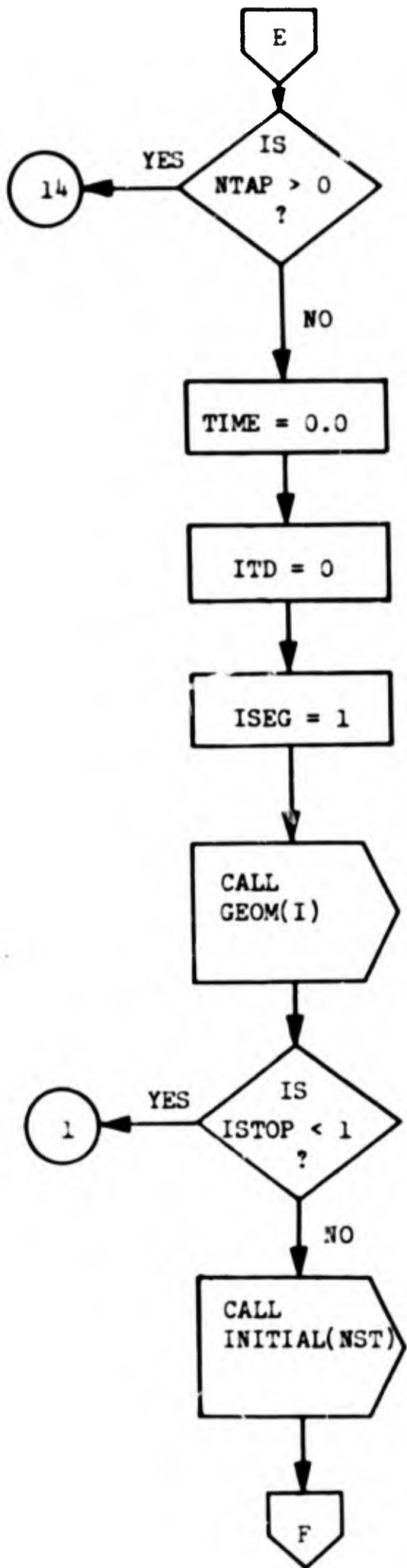
ARBITRARY BODY SHAPE OPTION.



COMPUTE CONSTANTS FOR STAGNATION PRESSURE CALCULATION.

COMPUTE STAGNATION PRESSURE BEHIND A NORMAL SHOCK. PERFECT GAS.

EMPIRICAL CONSTANT USED TO ESTIMATE INITIAL BODY PRESSURE.



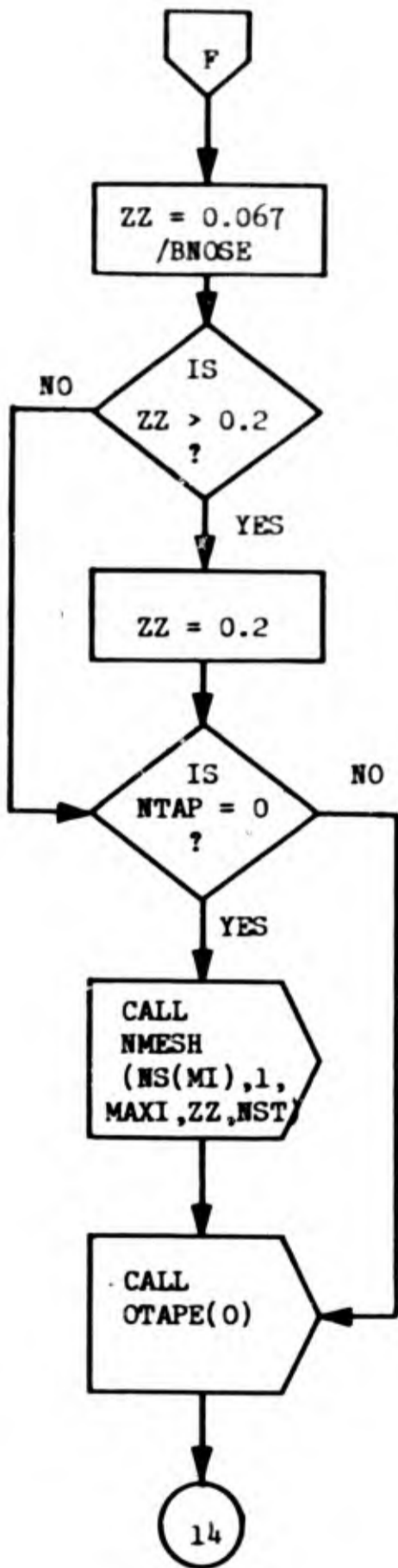
RESTART OR ADDITIONAL ITERATIONS ?

NO. NEW CASE. INITIALIZE CONSTANTS FOR FIRST SEGMENT.

SET UP BODY GEOMETRY FOR FIRST SEGMENT.

WAS DATA ERROR SENSED IN SUBROUTINE GEOM ?

INITIALIZE FLOW FIELD PARAMETERS FOR FIRST SEGMENT.



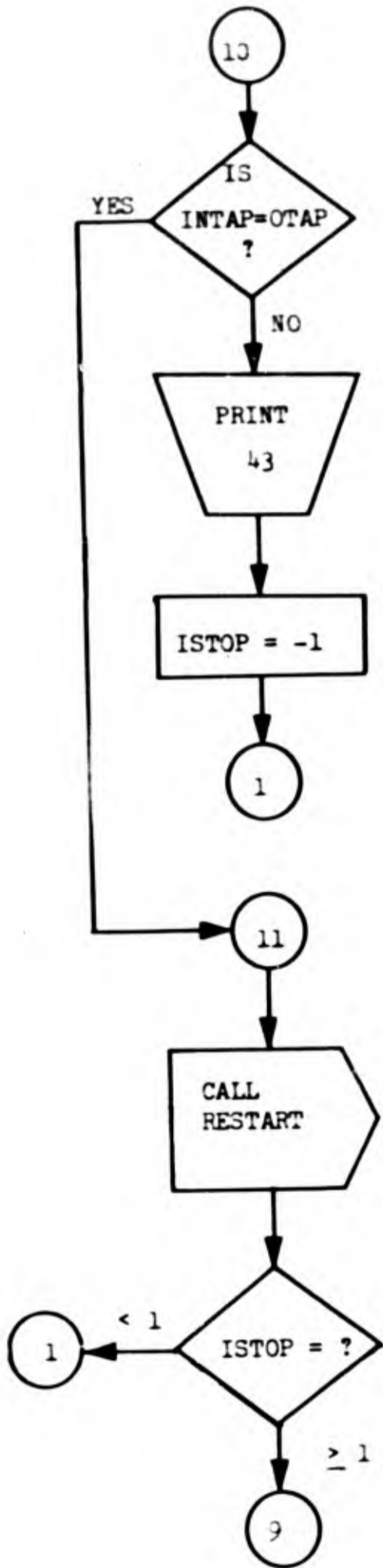
SET $\Delta\eta$ AT $S = 0$.

$\Delta\eta$ REQUIRED TO BE ≤ 0.2

NTAP < 0, UNSTEADY FLOW CASE,
DOES NOT EMPLOY AUTOMATIC
NODAL SPACING OPTION.

SET NODAL SPACING ACROSS
SHOCK LAYER.

STORE INPUT ARRAY ON
MAGNETIC TAPE.



RESTART OPTION.

THIS OPTION REQUIRES
INTAP = OTAP.

PRINT ERROR MESSAGE.

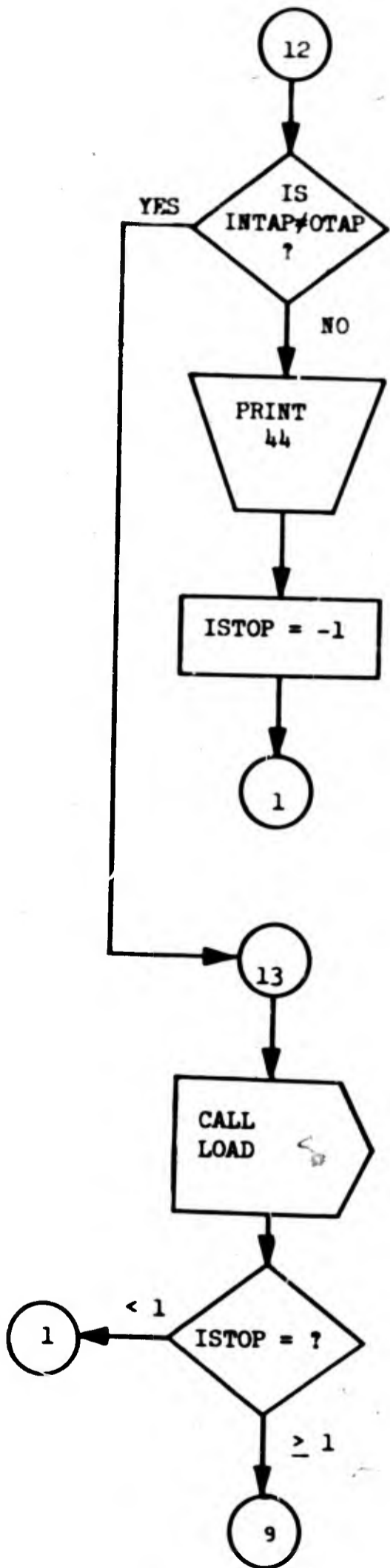
SET ERROR INDICATOR.

READ NEXT CASE.

LOAD DATA FOR RESTART.

WAS DATA ERROR SENSED
IN RESTART ?

GO TO 9 TO COMPUTE PT2.



ADDITIONAL ITERATIONS TO BE PERFORMED ON A CASE STORED ON MAGNETIC TAPE.

OPTION REQUIRES INTAP ≠ OTAP.

PRINT ERROR MESSAGE.

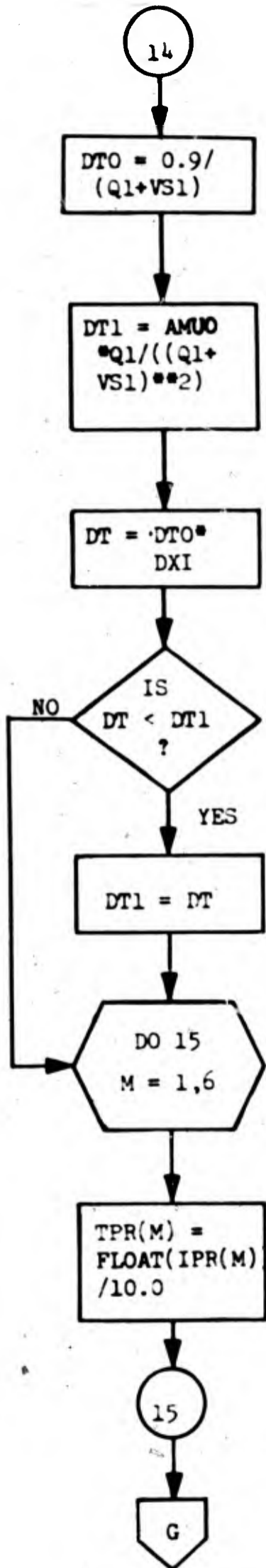
SET ERROR INDICATOR.

READ NEXT CASE.

LOAD DATA FROM MAGNETIC TAPE.

ERROR INDICATOR < 1 ? IN THIS CASE, IT SIGNIFIES THE END OF THE CASE.

GO TO 9 TO COMPUTE PT2.



COMPUTE CONSTANT, DTO, FOR USE IN CFL STABILITY CRITERION, EQUATION (36).

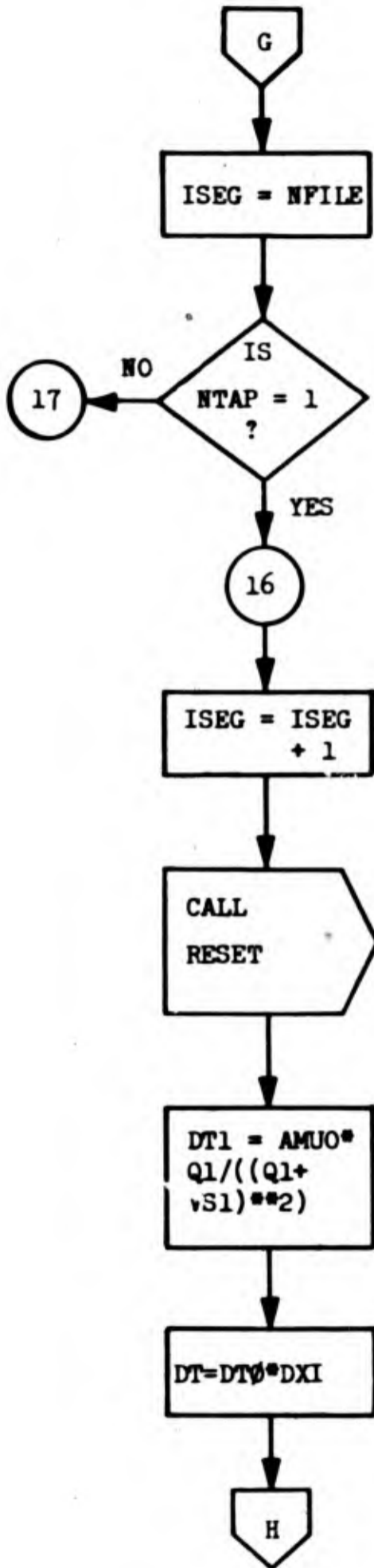
EQUATION (62).

COMPUTE CFL STABILITY LIMIT BASED ON $\Delta\xi$.

DOES DT1 EXCEED THIS LIMIT ?

YES. SET DT1 = DT.

SET TPR ARRAY FOR USE IN STEADY STATE FLOW OPTION.



INITIALIZE SEGMENT NUMBER.

RESTART OPTION ?

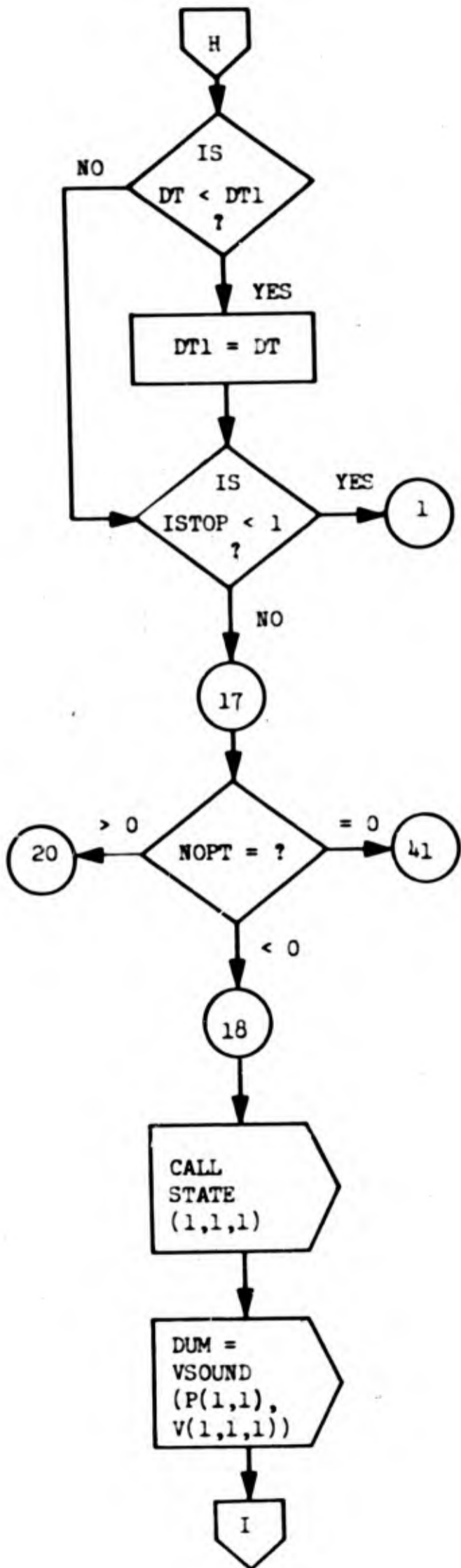
START OF SEGMENT COMPUTATIONS.

INCREMENT SEGMENT NUMBER.

SET UP INITIAL DATA FOR
PRESENT SEGMENT.

EQUATION (62).

COMPUTE CFL STABILITY LIMIT
BASED ON $\Delta\xi$.



DOES DT1 EXCEED THIS LIMIT ?

YES. SET DT1 = DT.

ERROR INDICATOR < 1 (THIS MEANS AN INPUT ERROR WAS SENSED OR THE CASE WAS COMPLETED).

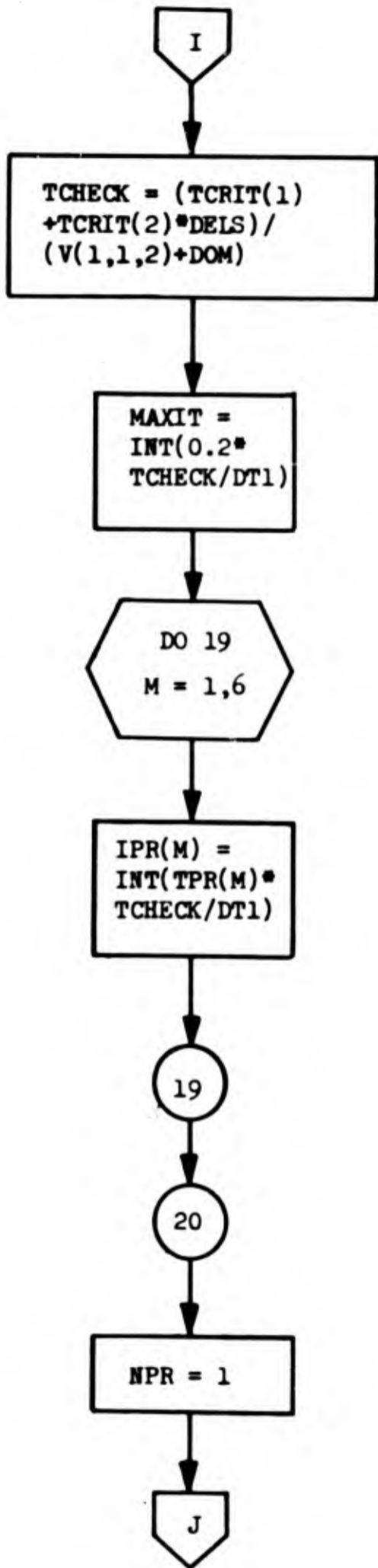
NOPT > 0 : NUMBER OF ITERATIONS SPECIFIED.

NOPT = 0 : NO ITERATIONS. JUST PRINT FLOW DATA.

NOPT < 0 : STEADY STATE FLOW SOLUTION.

COMPUTE BODY PRESSURE AT START OF SEGMENT.

COMPUTE BODY SOUND SPEED AT START OF SEGMENT.

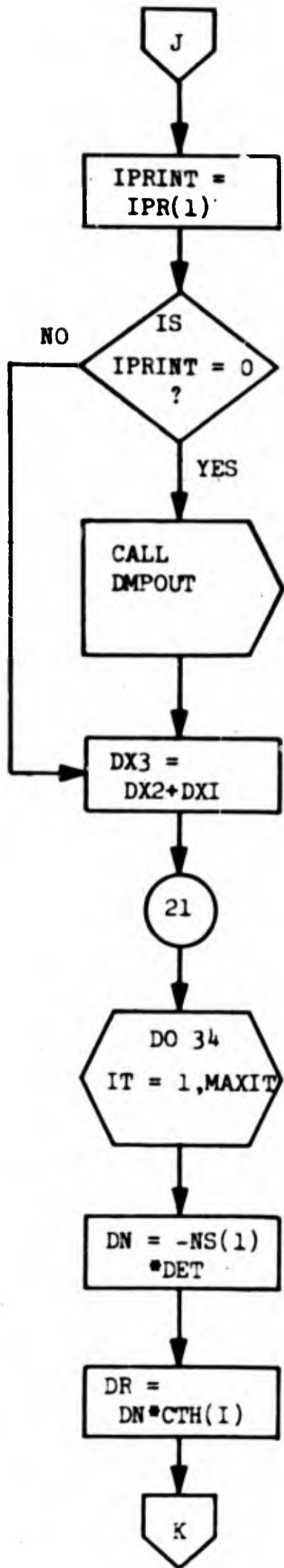


ESTIMATE TIME REQUIRED FOR STEADY STATE SOLUTION, EQUATION (63).

SET MAXIT TO ALLOW SUFFICIENT ITERATIONS FOR A MEANINGFUL CFL STABILITY CALCULATION BASED ON $\Delta\tau$.

ESTIMATE ITERATION NUMBERS FOR WHICH PRINTOUT SHOULD BE PROVIDED.

INITIALIZE PRINTOUT INDEX.



SET CONTROL VARIABLE FOR PRINTOUT.

ARE INITIAL CONDITIONS TO BE PRINTED ?

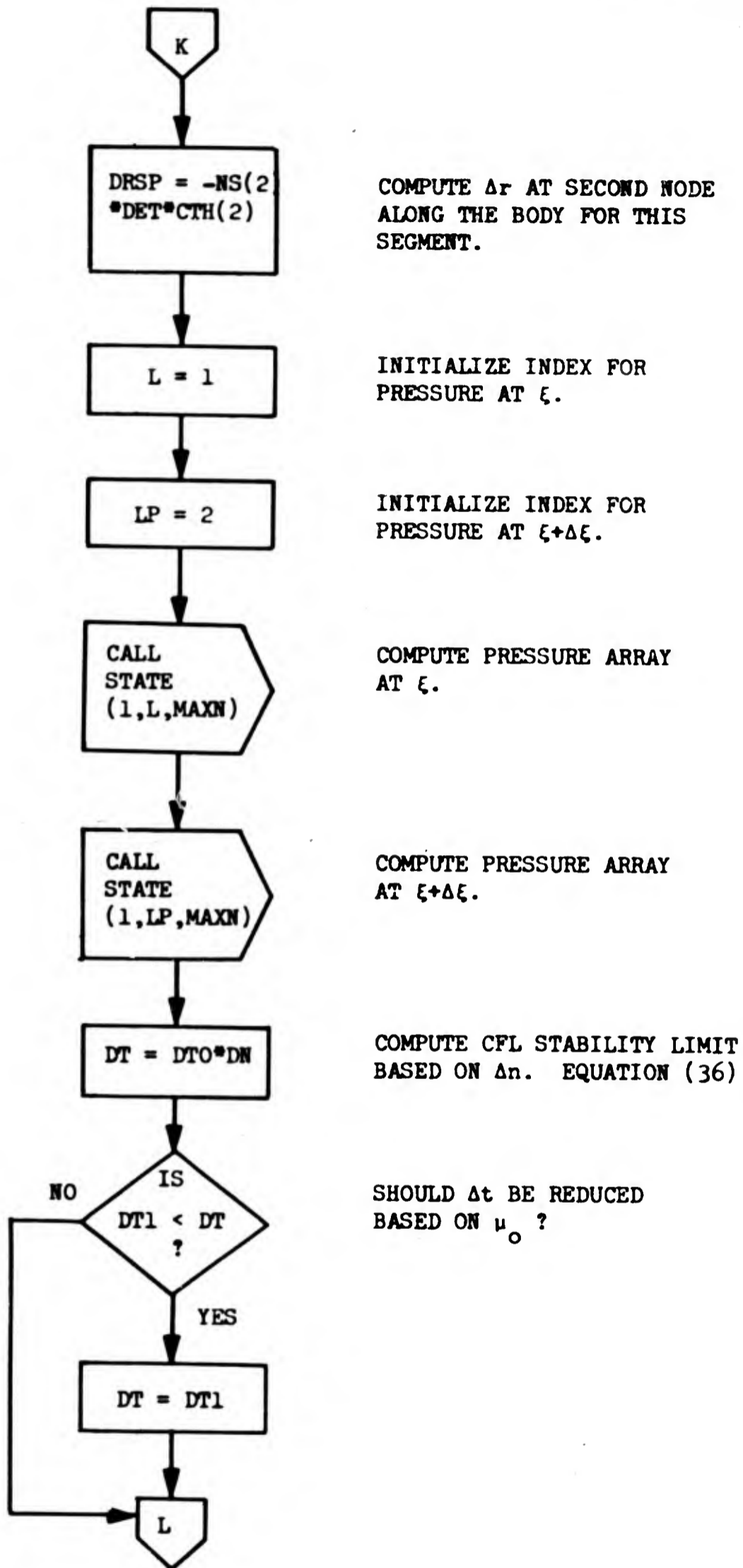
PRINT FLOW FIELD DATA.

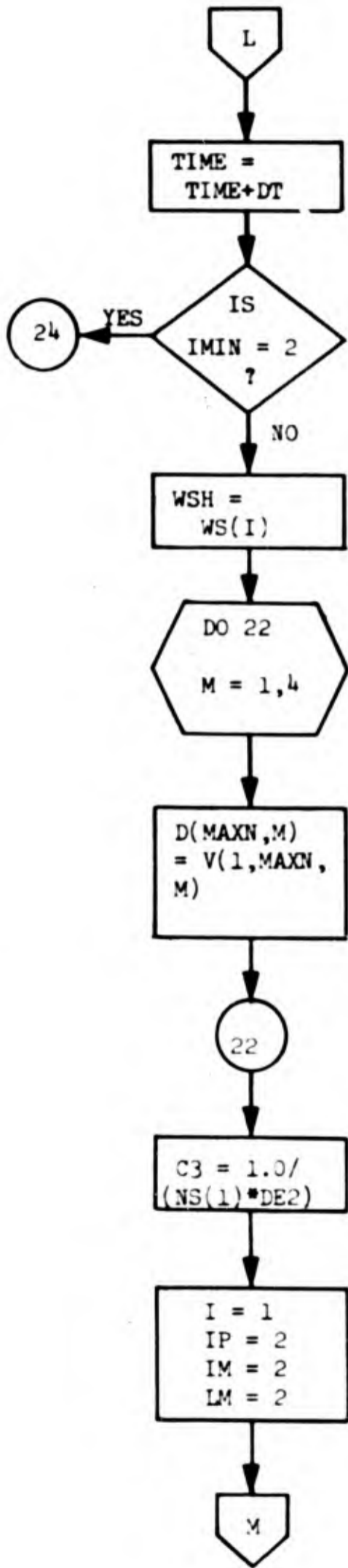
STORE $3\Delta\xi$.

START OF ITERATION LOOP.

COMPUTE Δn AT START OF SEGMENT; $\Delta n = -n_s \Delta\eta$.

COMPUTE Δr AT START OF SEGMENT; $\Delta r = \Delta n \cos(\theta)$.





INCREMENT TIME.

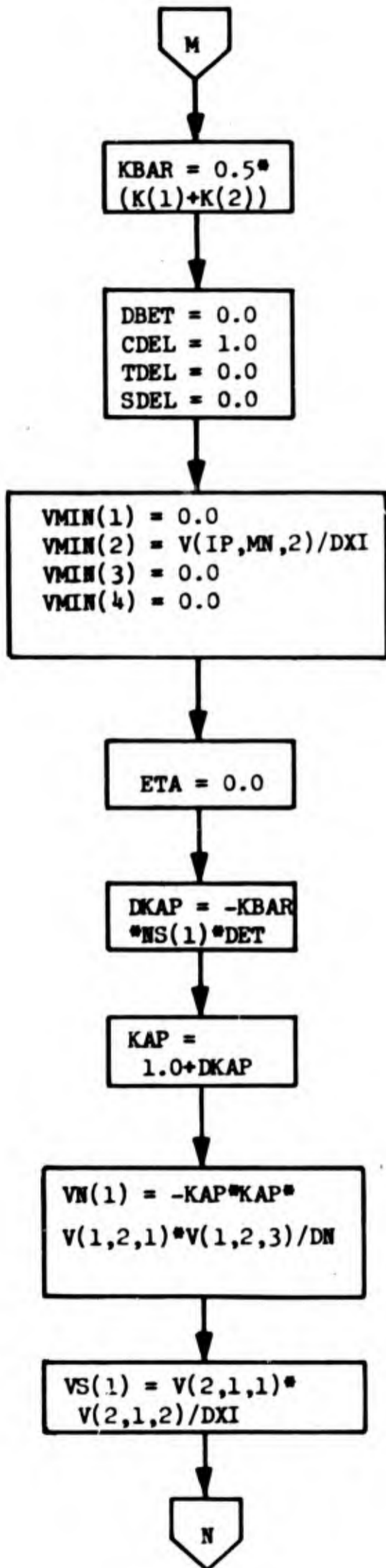
IMIN = 2 WHEN THE SEGMENT UNDER CONSIDERATION DOES NOT INCLUDE THE STAGNATION STREAMLINE.

COMPUTE SHOCK VELOCITY ALONG THE NORMAL COORDINATE.

STORE FLOW VARIABLES AT THE SHOCK.

COMPUTE CONSTANT FOR USE IN EVALUATING PARTIAL DERIVATIVES WITH RESPECT TO n .

SET NODAL INDICES FOR STAGNATION STREAMLINE COMPUTATIONS.



COMPUTE AVERAGE BODY CURVATURE FOR STAGNATION STREAMLINE DIFFERENCES.

SET β , $\cos(\beta)$, $\tan(\beta)$, AND $\sin(\beta)$ FOR NORMAL SHOCK COMPUTATIONS.

STORE $\frac{\partial \rho}{\partial \xi}$, $\frac{\partial u}{\partial \xi}$, $\frac{\partial v}{\partial \xi}$ AND $\frac{\partial s}{\partial \xi}$ AT THE FIRST NODE INSIDE THE SHOCK.

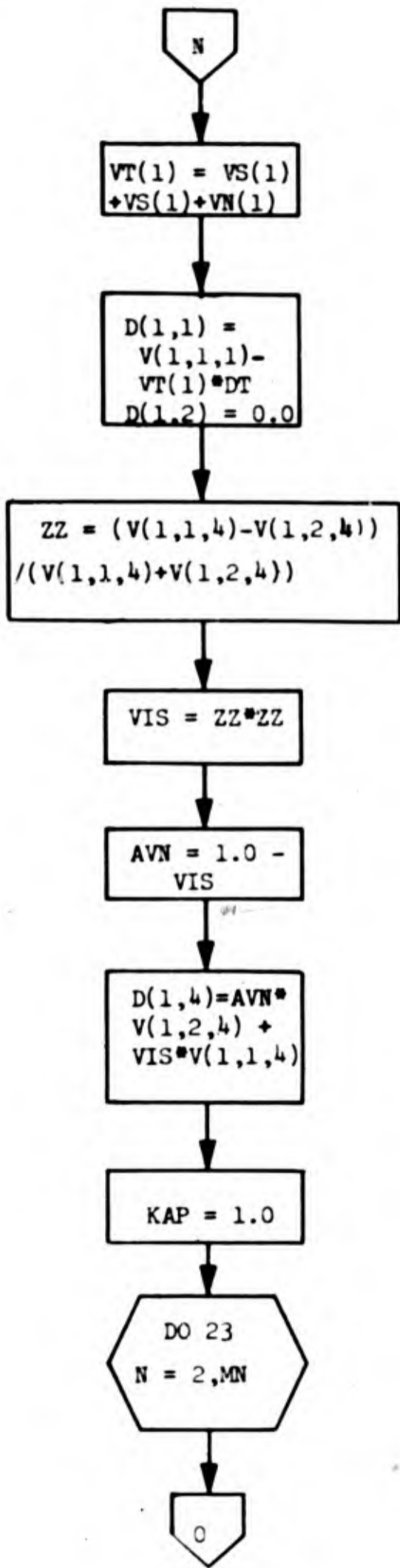
INITIALIZE η .

COMPUTE $\Delta \kappa = K \Delta \eta$.

COMPUTE κ AT $\eta = \Delta \eta$.

COMPUTE $\frac{\partial \kappa^2 \rho v}{\partial \eta}$ FOR USE IN EQUATION (16).

COMPUTE $\frac{\partial \rho u}{\partial s}$ FOR USE IN EQUATION (16).



COMPUTE $\frac{\partial \rho}{\partial t}$ AT THE STAGNATION POINT.

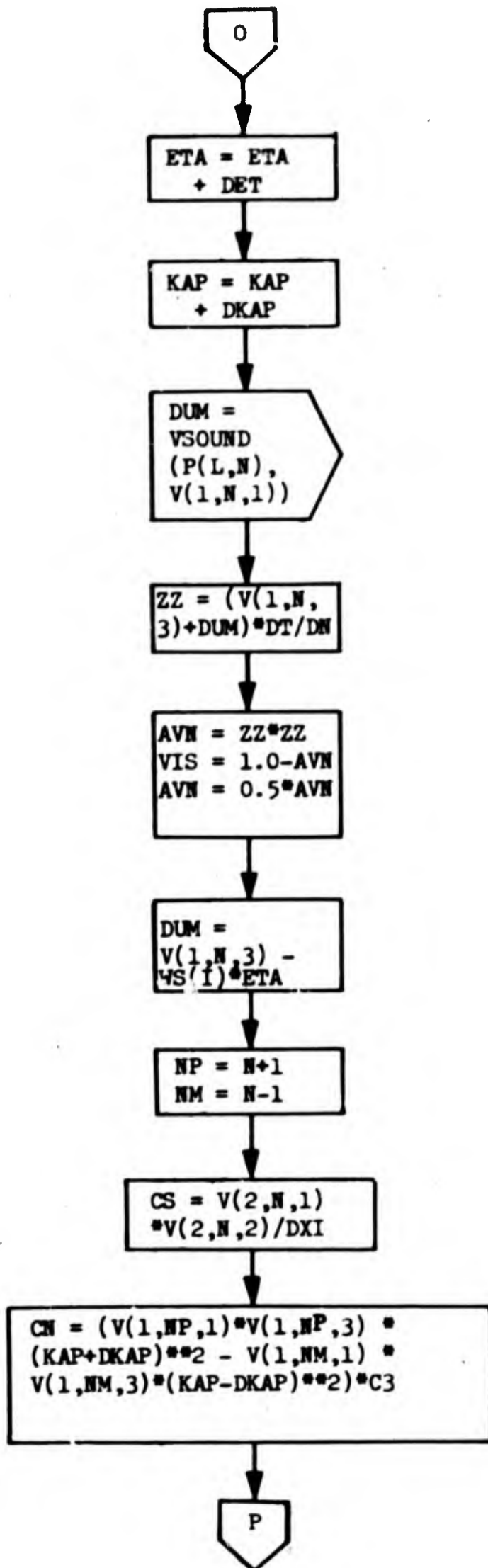
COMPUTE $\rho(\tau + \Delta\tau)$, EQUATION (16). SET $u=0$.

EQUATION (41).

COMPUTE $S(\tau + \Delta\tau)$, EQUATION (40).

INITIALIZE κ .

START LOOP ON N FOR POINTS ON THE STAGNATION STREAMLINE.



INCREMENT η .

INCREMENT κ .

COMPUTE LOCAL SOUND SPEED.

COMPUTE MINIMIZED STABILIZING TERMS.

STORE A USEFUL PARAMETER.

STORE INDICIES FOR THE NODES AT $\eta + \Delta\eta$ AND $\eta - \Delta\eta$.

COMPUTE $\frac{\partial \rho u}{\partial \xi}$

COMPUTE $\frac{\partial \kappa^2 \rho v}{\partial \eta}$



$$\begin{aligned} VN(1) &= C3 * (V(1, NP, 1) - V(1, NM, 1)) \\ VN(3) &= C3 * (V(1, NP, 1) - V(1, NM, 1)) \\ VN(4) &= C3 * (V(1, NP, 1) - V(1, NM, 1)) \end{aligned}$$

COMPUTE $\frac{\partial \rho}{\partial n}$, $\frac{\partial v}{\partial n}$, AND

$\frac{\partial s}{\partial n}$.

$$\begin{aligned} VA(1) &= AVN * (V(1, NP, 1) + V(1, NM, 1)) \\ &\quad + VIS * V(1, N, 1) \\ VA(3) &= AVN * (V(1, NP, 3) + V(1, NM, 3)) \\ &\quad + VIS * V(1, N, 3) \\ VA(4) &= AVN * (V(1, NP, 4) + V(1, NM, 4)) \\ &\quad + VIS * V(1, N, 4) \end{aligned}$$

COMPUTE AVERAGE VALUES OF ρ , v AND s TO INTRODUCE STABILIZING TERMS.

$$VT(1) = (CS + CS + CN / KAP) / KAP - WS(1) * ETA * VN(1)$$

COMPUTE $\frac{\partial \rho}{\partial \tau}$.

$$VT(3) = DUM * VN(3) + C3 * (P(L, NP) - P(L, NM)) / V(1, N, 1)$$

COMPUTE $\frac{\partial v}{\partial \tau}$.

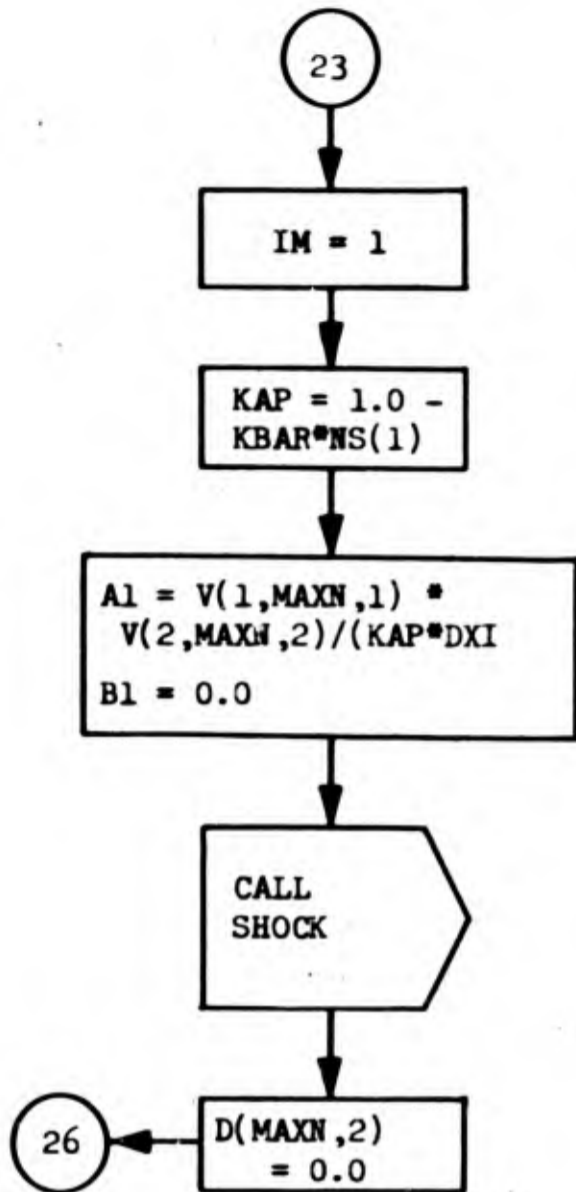
$$VT(4) = DUM * VN(4)$$

COMPUTE $\frac{\partial s}{\partial \tau}$.

$$\begin{aligned} D(N, 1) &= VA(1) - VT(1) * DT \\ D(N, 2) &= 0.0 \\ D(N, 3) &= VA(3) - VT(3) * DT \\ D(N, 4) &= VA(4) - VT(4) * DT \end{aligned}$$

SOLVE GOVERNING EQUATIONS.





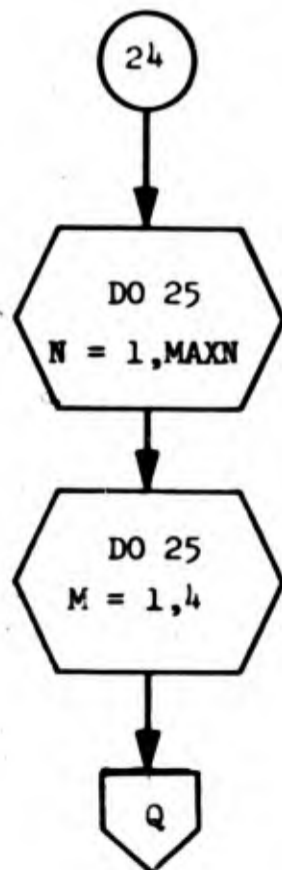
SET IM TO INSURE PROPER FLOW VARIABLE STORAGE IN SUBROUTINE SHOCK.

$$\kappa = 1 - Kn_s$$

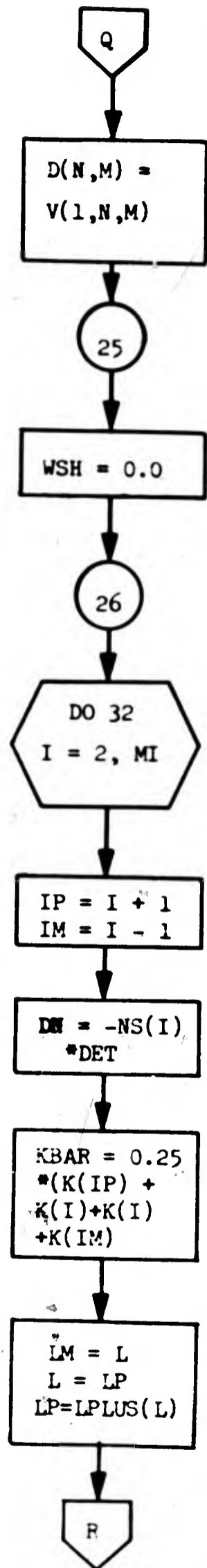
COMPUTE CONSTANTS FOR SHOCK CALCULATIONS.

COMPUTE SHOCK PARAMETERS.

SET U = 0 AT THE SHOCK.



COMPUTATIONAL STARTING POINT WHEN IMIN = 2



STORE STARTING LINE
PARAMETERS.

SET SHOCK VELOCITY COMPONENT
NORMAL TO THE BODY TO ZERO FOR
THE CONSTANT STARTING LINE.

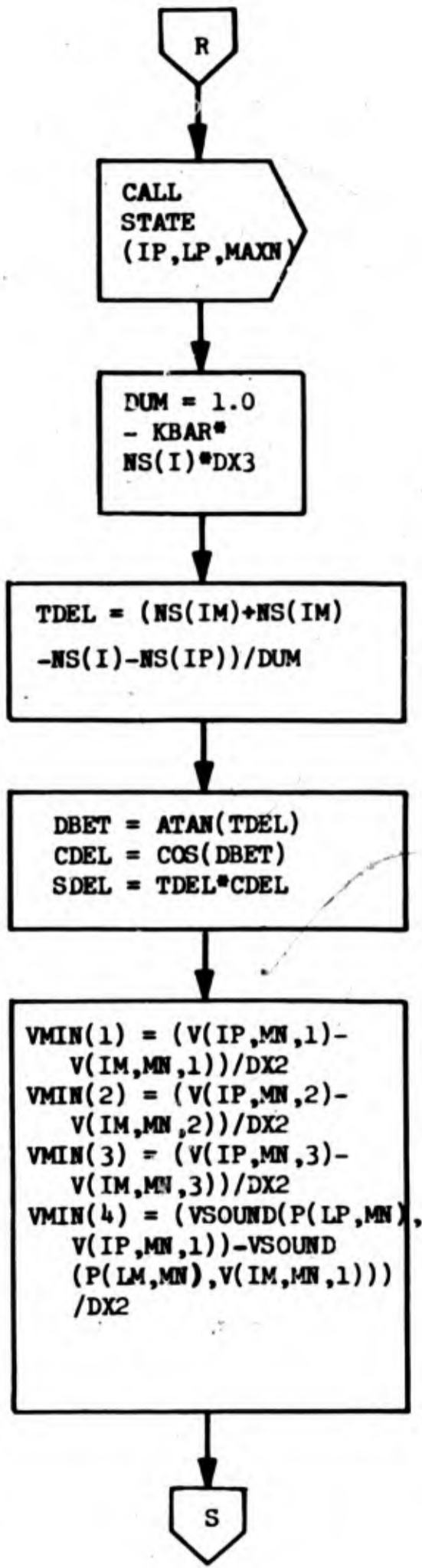
START OF LOOP TO TREAT POINTS
NOT ON THE STAGNATION STREAMLINE.

COMPUTE INDICIES FOR NODES AT
 $\xi + \Delta\xi$ AND $\xi - \Delta\xi$.

COMPUTE Δn .

COMPUTE AVERAGE CURVATURE FOR
FINITE DIFFERENCES.

SET INDICIES FOR THE PRESSURE
ARRAY.



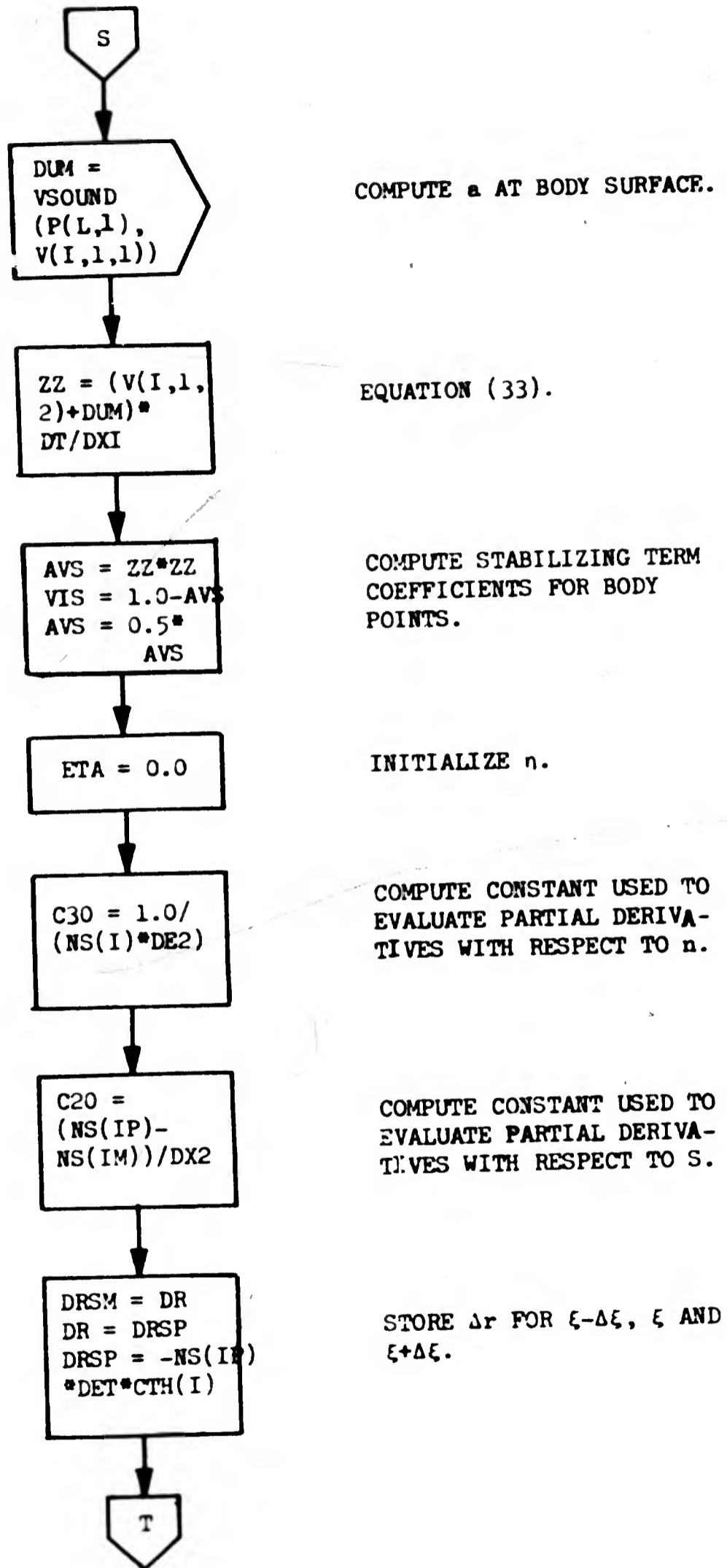
COMPUTE PRESSURE FOR
NODES AT $\xi + \Delta\xi$.

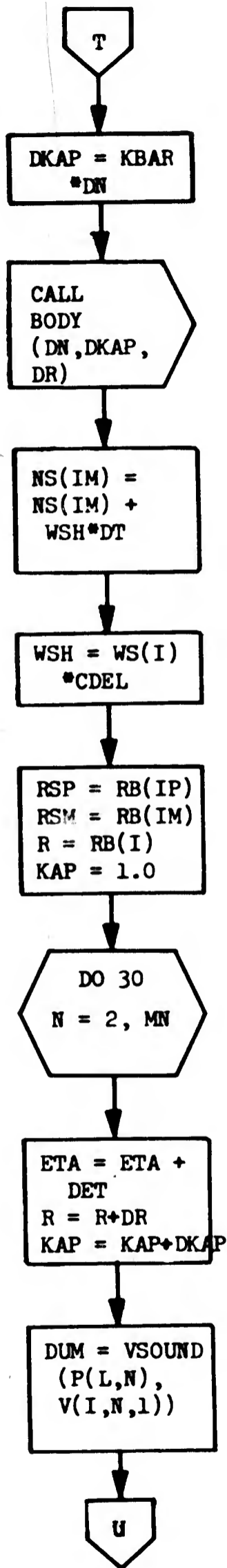
COMPUTE $3\kappa\Delta\xi$ AT THE
SHOCK.

SHOCK SLOPE, EQUATION (60).

COMPUTE CONSTANTS USED IN
SHOCK TRANSFORMATIONS.

COMPUTE $\frac{\partial p}{\partial \xi}$, $\frac{\partial u}{\partial \xi}$, $\frac{\partial v}{\partial \xi}$ AND $\frac{\partial a}{\partial \xi}$
AT THE FIRST NODE INSIDE
THE SHOCK FOR USE IN
SUBROUTINE SHOCK.





COMPUTE $\Delta\kappa = K\Delta n$.

BODY POINT COMPUTATIONS.

COMPUTE NEW SHOCK POINT LOCATION FOR $\xi - \Delta\xi$.

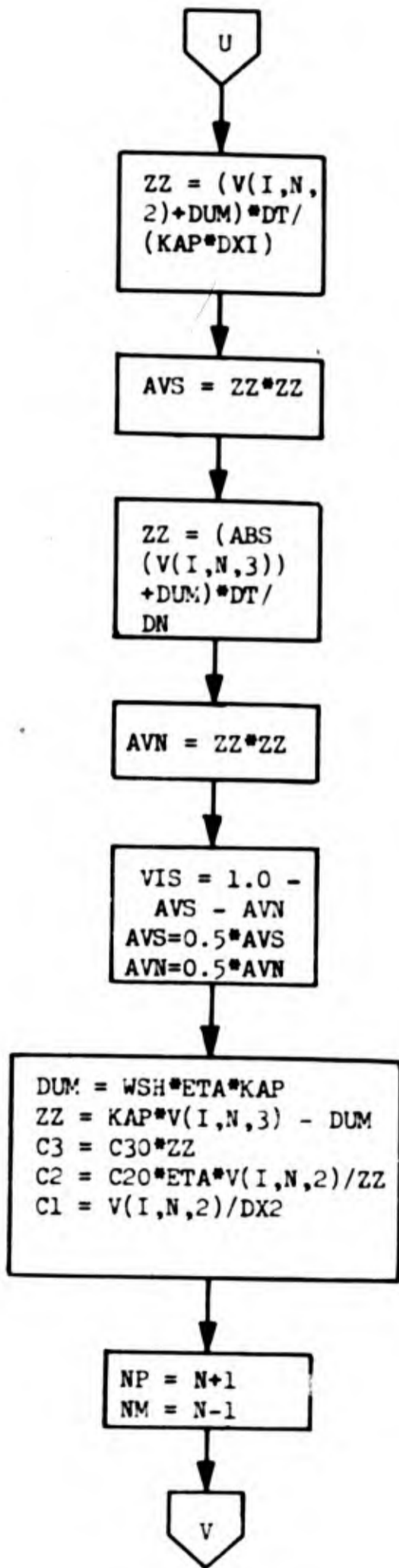
COMPUTE SHOCK VELOCITY COMPONENT NORMAL TO THE BODY.

INITIALIZE r AT $\xi + \Delta\xi$, $\xi - \Delta\xi$ AND ξ , AND κ .

START OF LOOP TO TREAT POINTS ALONG THE NORMAL COORDINATE AT ξ .

INCREMENT η , r AND κ .

COMPUTE SOUND SPEED.



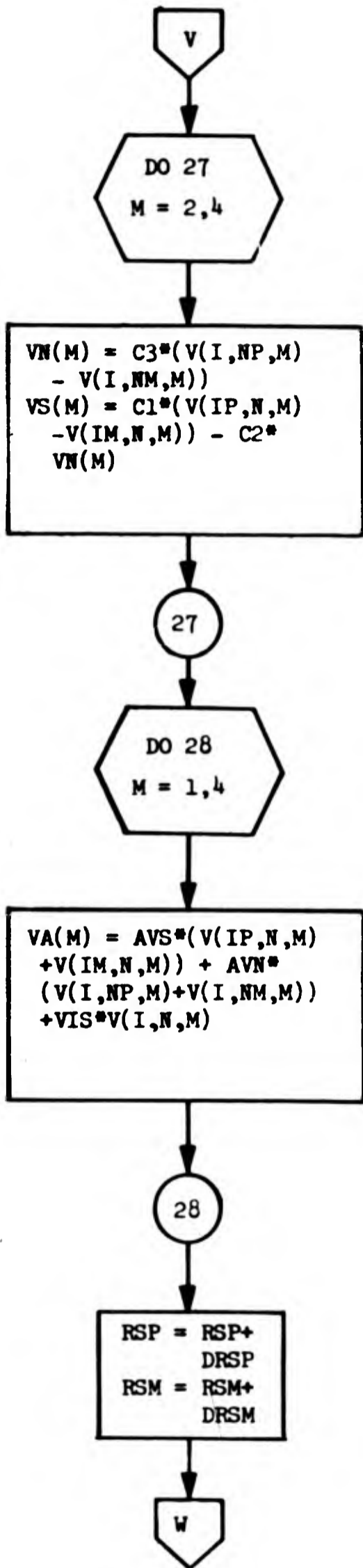
EQUATION (33) FOR ξ COORDINATE.

EQUATION (33) FOR η COORDINATE.

COEFFICIENTS FOR THE STABILIZING TERMS.

COMPUTE USEFUL PARAMETERS.

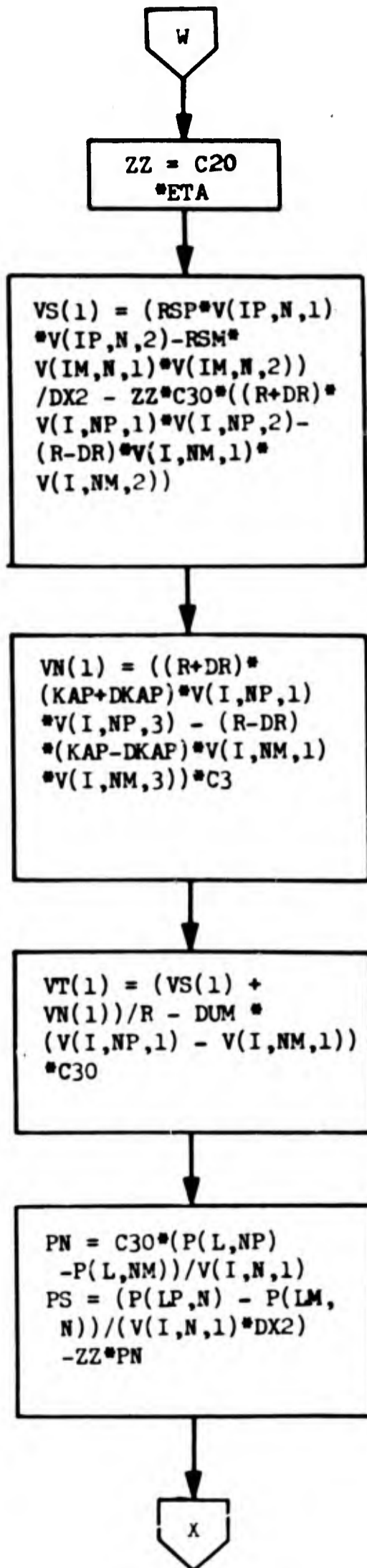
COMPUTE INDICES OF NODES AT $n+\Delta\eta$ AND $n-\Delta\eta$.



COMPUTE PARTIAL DERIVATIVES
FOR U, V AND S.

COMPUTE AVERAGES TO
INTRODUCE STABILIZING TERMS
FOR ρ , U, V AND S.

INCREMENT r AT $\xi + \Delta\xi$
AND $\xi - \Delta\xi$.

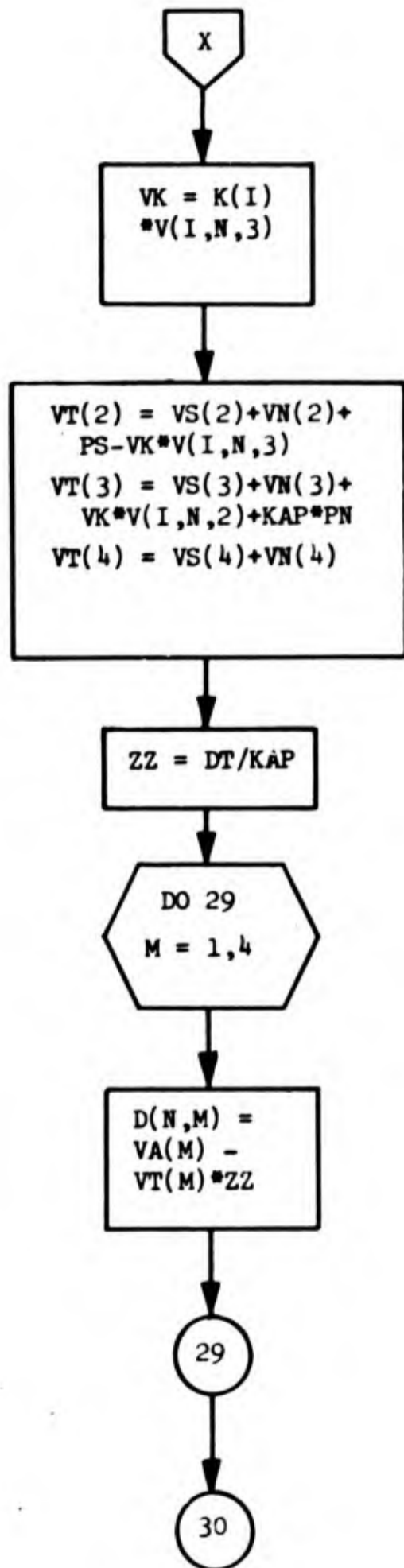


$$\frac{\partial \rho v}{\partial S}$$

$$\frac{\partial \kappa \rho v}{\partial n}$$

$$- \frac{\partial \rho}{\partial \tau}$$

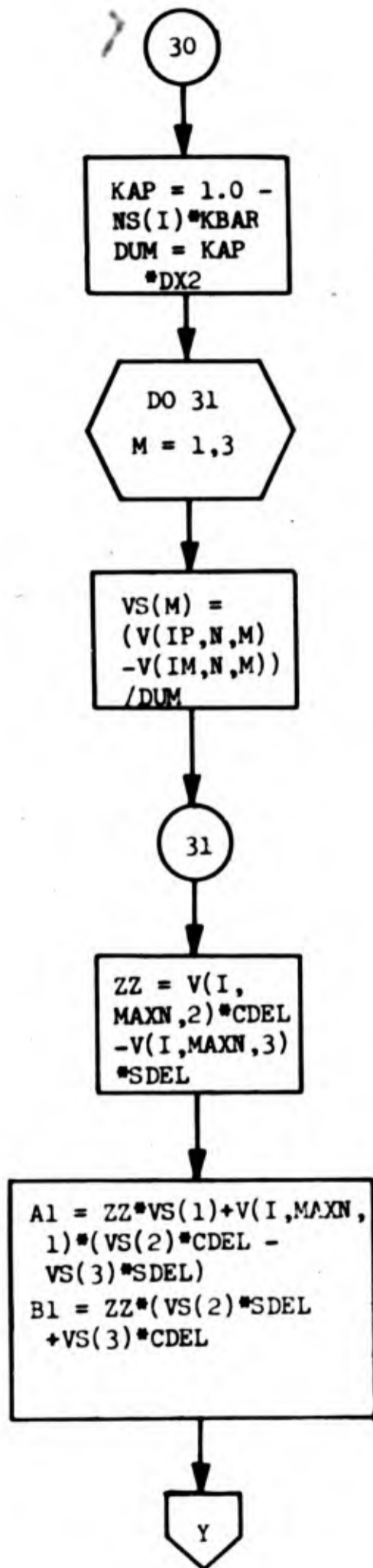
$$\frac{1}{\rho} \frac{\partial P}{\partial n} \quad \text{AND} \quad \frac{1}{\rho} \frac{\partial P}{\partial S}$$



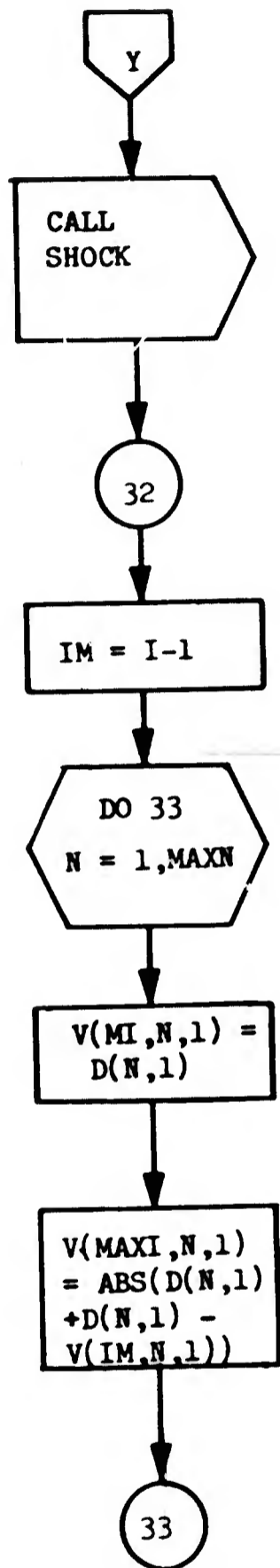
$$-\frac{\partial u}{\partial t}, -\frac{\partial v}{\partial t} \text{ AND } -\frac{\partial S}{\partial t}.$$

COMPUTE FLOW PARAMETERS
AT $\tau + \Delta\tau$.

END OF LOOP ON N.



COMPUTE CONSTANTS FOR
SHOCK CALCULATIONS.

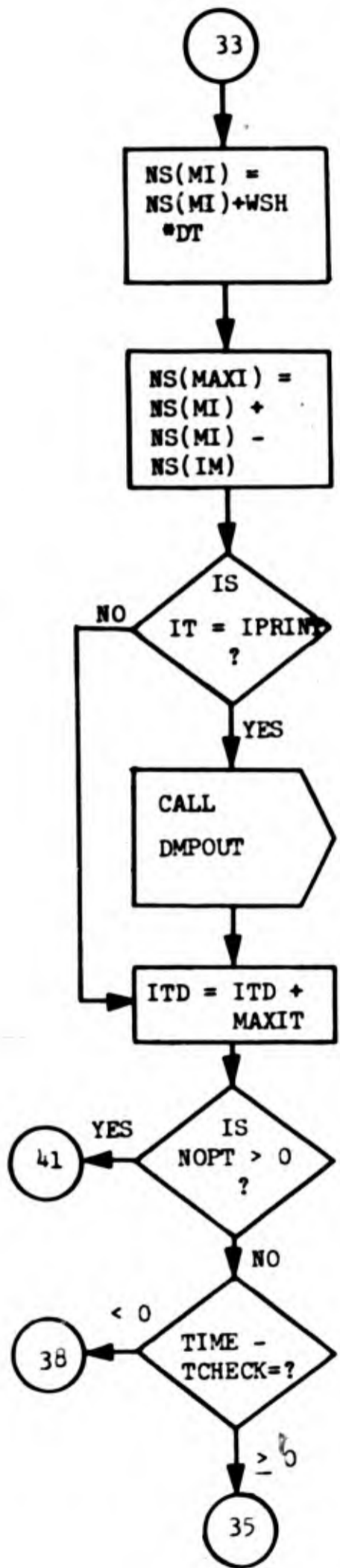


COMPUTE SHOCK PARAMETERS.

END OF LOOP ON I.

TRANSFER VALUES OF ρ FOR THE DOWNSTREAM SEGMENT BOUNDARY FROM TEMPORARY STORAGE.

LINEAR EXTENSION FOR ρ . ABSOLUTE VALUES USED TO INSURE INITIAL GUESS WILL NOT INTRODUCE NEGATIVE VALUES.



COMPUTE NEW SHOCK POINT
LOCATION AT DOWNSTREAM
SEGMENT BOUNDARY.

LINEAR EXTENSION OF SHOCK.

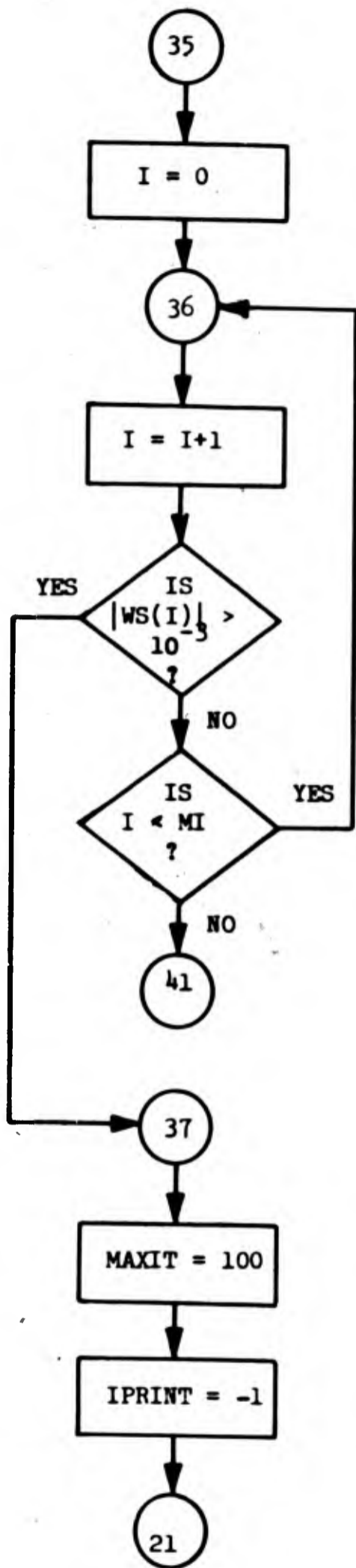
SHOULD PRINTOUT BE PROVIDED ?

PRINT FLOW FIELD VARIABLES.

COMPUTE TOTAL NUMBER OF
ITERATIONS PERFORMED ON
THIS SEGMENT.

WAS THE NUMBER OF ITERATIONS
TO BE PERFORMED SPECIFIED ?

HAS THE ESTIMATED TIME FOR
STEADY STATE BEEN EXCEEDED ?



INITIALIZE I.

INCREMENT I.

IS THE MAGNITUDE OF THE SHOCK VELOCITY GREATER THAN 10^{-3} ?

ARE THERE MORE POINTS IN THIS SEGMENT ?

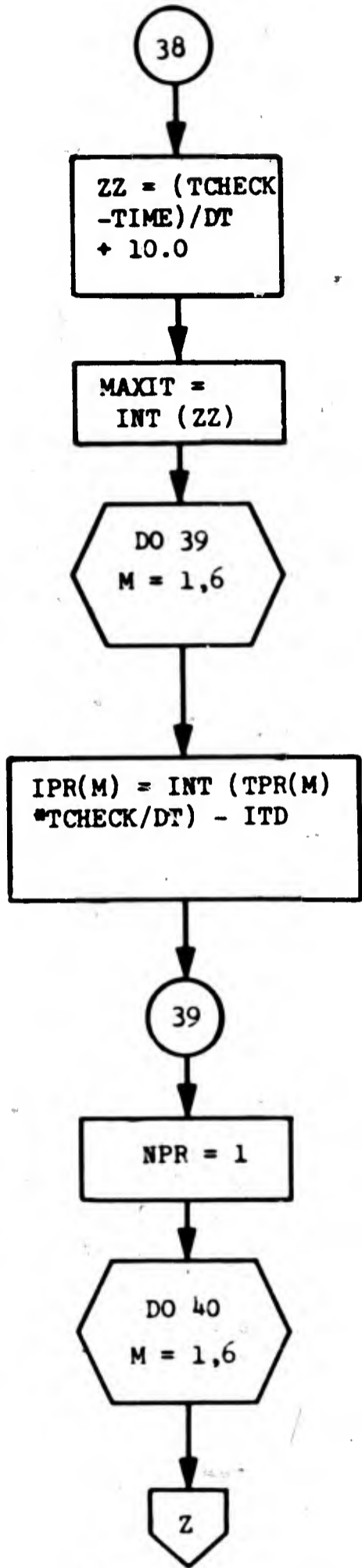
STEADY STATE ACHIEVED.

STEADY STATE NOT ACHIEVED.

SET MAXIT FOR 100 MORE ITERATIONS.

SUPPRESS EXTRA PRINTOUT.

START ITERATION LOOP.

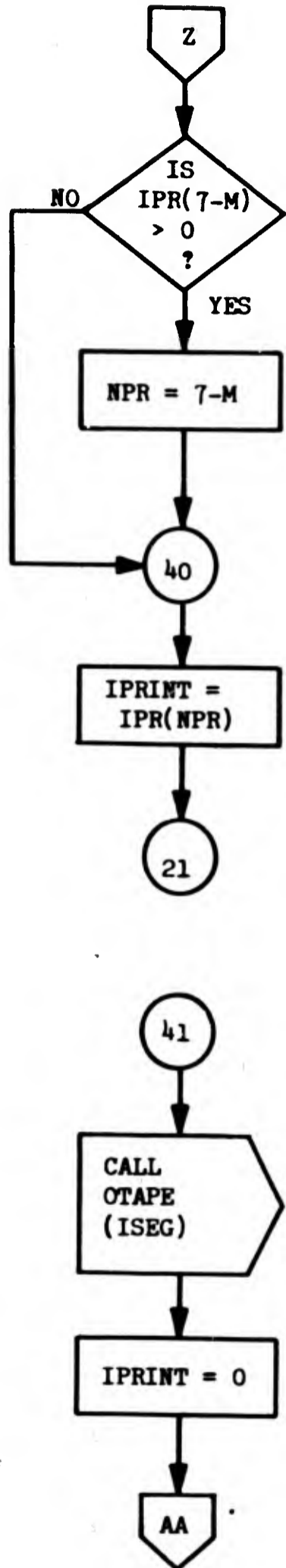


ESTIMATE THE NUMBER OF
ADDITIONAL ITERATIONS
REQUIRED FOR STEADY STATE.

CONVERT TO INTEGER FORM.

EVALUATE NEW PRINTOUT
ARRAY.

INITIALIZE PRINTOUT
CONTROL INDEX.



IS IPR(7-M) > 0 ?

RESET NPR.

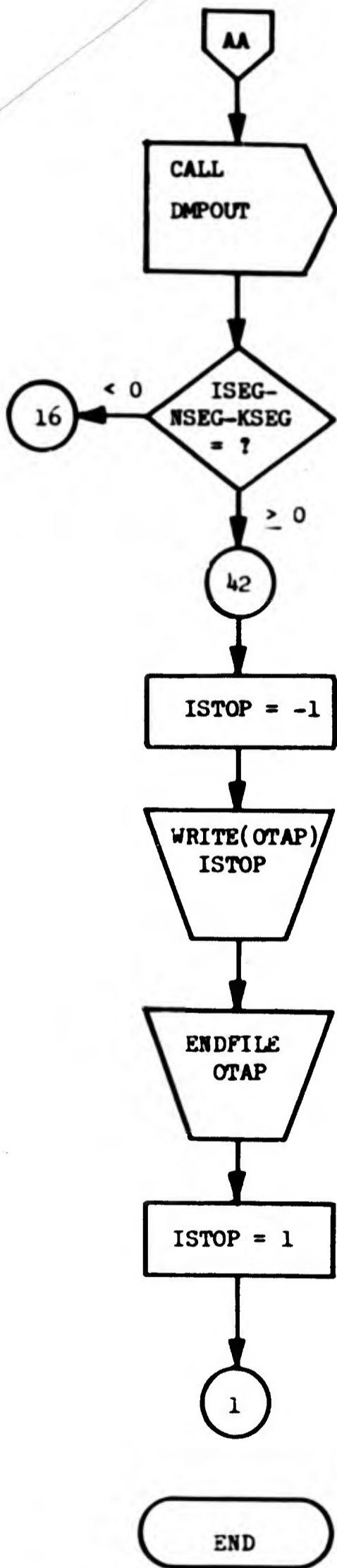
SET PRINT CONTROL VARIABLE.

START ITERATION LOOP.

SEGMENT COMPLETED.

WRITE FLOW FIELD DATA ON MAGNETIC TAPE.

SET IPRINT FOR FINAL PRINTOUT.



FINAL PRINTOUT FOR THIS SEGMENT.

HAVE ALL SEGMENTS BEEN TREATED ?

WRITE INDICATOR ON TAPE TO SIGNIFY END OF CASE.

ENDFILE TAPE.

SET ERROR INDICATOR.

READ NEXT CASE.

END.


```

PROGRAM BLUNT(INPUT,OUTPUT,TAPE1,TAPE2,TAPE5=INPUT)
BLUN 1
BLUN 2
BLUN 3
BLUN 4
BLUN 5
BLUN 6
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BLUN 8
BLUN 9
BLUN 10
BLUN 11
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BLUN 49
BLUN 50
BLUN 51
BLUN 52

      COMMON /1/ V(15,4),J(14),F(3,40)
      COMMON /2/ R3(15),Z3(15),TH(15),CTH(15),K(15),NS(15),WS(15)
      COMMON /3/ MAX1,MI,IHIN,MAXN,MN,MIN
      COMMON /4/ IPPINT,NPR,KGEOM,TPR(5),KSEG
      COMMON /5/ OXI,DET,OX2,DE2,DT
      COMMON /6/ O1,VS1,MT1,A1,H1
      COMMON /7/ I,IP,IL,L,LP,LH
      COMMON /8/ VMIN(4),KAP,KPAR,KZRO,KNOSE
      COMMON /9/ ISEG,XIG,JELS,KZERO,KNOSE
      COMMON /10/ VS(4),VN(4),VA(4),VT(4)
      COMMON /11/ MAXIT,ISTOP,ITJ,TIME,SI(15),NST(15)
      COMMON /12/ ORET,ODEL,SOEL,TOEL
      COMMON /13/ AVS,AVN,VIS
      COMMON /INPUT/ GAM1,ALPHA1,THETA,ZRN,TCRIT(2),P1,RHO1,R1,AMUJ,T1,=
      P2(2),IPR(5),INTAP,OTAP,NOPT,NCASE,NFILE,NSEG,IGEOM,NGEOM,NTAP
      COMMON /CONST/ PI,PI2,GAMA,GAMP,RGAS
      COMMON /MESH/ IMAX,NMAX
      COMMON /PRES/ PT2,PJ1
      DIMENSION LPLUS(3),JPR(5)
      INTEGER OTAP
      REAL NS,K,KAP,KPAR,KZERO,H1,NST
      DATA PI,PI2/3.14159265359,1.570796326795/
      DATA (LPLUS(M),M=1,3)/2,3,1/
      DATA (JPR(M),M=1,5)/-5,-5,-5,-5,-5/
      DATA
      AMUJ/ 0.14/, TCRIT(1)/ 2.0/,
      TCRIT(2)/ 1.0/, GAM1/ 1.4/,
      NSEG/ 0/, T1/ 1.0/,
      P1/ 1.0/, RHO1/ 1.0/,
      R1/ 1.0/, ZRN/ 0.0/,
      ERR(1)/ 1.0E-05/, ERR(2)/ 1.0E-08/,
      INTAP/ 2/, OTAP/ 1/,
      NOPT/ -1/, NFILE/ 1/,
      NTAP/ 0/, IGEOM/ 0/,
      NGEOM/ 0/, IIMAX/ 15/,
      NMAX/ 40/, NCASE/ 1/

      ISTOP=5
      PRINT 45
      IHIN=1
      DELS=0.0
      KZERO=1.0
      KNOSE=1.0
      XIG=0.0
      KSEG=0
      AMUJ=11
      O1=1,5
      TPR(2)=JPR(4)

```

NOT REPRODUCIBLE

```

CALL DATA (I,ICR)
M1=AMACH1
GAM1=(GAM1+1.0)/2.0
GAM1=(GAM1-1.0)/2.0
M1=INT(GAM1)
ICR=PI/(PHO1*PI)
IF (MACH1.LT.1.0) AMACH1=SQRT(PHO1*AMACH1*AMACH1/(GAM1*PI))
G1=AMACH1*VSI
HT1=GAM1/(GAM1-1.0)+G1*PI/2.0
MAXIT=NCOT
DO WHILE (I)
  I=I+1
  J=J+1
  IF (I) 4,12,12
  MAXIT=ICRGM
  MAXN=NGFOM
  IF (MAXN.LT.1) MAXN=5
  MN=MAXN-1
  MI=MAXI-1
  MM=MM-1
  Z=FLCAT(MN)
  ZT=1./Z
  ZC=ZT+ZT
  IF (ZC) 8,8,8
  IF (THEIC.GT.7).1.02.AMACH1.LT.2.0) GO TO 6
  KGFOM=7
  ZC=1.87266
  GO TO 7
  ZC=(91.1-THEIC)*PI/191.1
  KGFOM=1
  THEIC=THEIC*PI/191.1
  IF (MS.G.LT.1) KSEG=50111
  GO TO 8
  ZC=ZC
  KSH=AMACH1*AMACH1
  Z=1.0*(GAM1-1.0)
  ZC=Z*GAM1
  ZC=((GAM1*KSH)**ZC)*((GAM1/(GAM1*KSH-GAM1))**Z7)
  ZC=(ZC/16.0)-1.5
  IF (ZC.LT.1.0) ZC=1.0
  IF (MACH.GT.2) GO TO 14
  ZT=1.0
  ZC=1.0
  CALL GCOM (I)
  IF (MACH.LT.1) GO TO 1
  CALL INITIAL (N)
  Z=1.0/Z/NOSE
  IF (MACH.GT.2) Z=1.0
  IF (MACH.GT.2) CALL AMESH (NS(MI),1,MAXI,ZZ,N)
  CALL START (I)
  GO TO 14
  IF (MACH.GT.2) GO TO 11
  ZT=1.0

```

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PLUN 54
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PLUN 104
PLUN 105
PLUN 106
PLUN 107

```

	ISTOP=-1	PLUN 109
	GO TO 1	PLUN 110
11	CALL DEFSADT	PLUN 111
	IF (ISTOP) 1,1,9	PLUN 111
12	IF (INT(NT*CTAP) GO TO 13	PLUN 112
	NDT=44	PLUN 113
	ISTOP=-1	PLUN 114
	GO TO 1	PLUN 115
13	CALL LGR1	PLUN 116
	IF (ISTOP) 1,1,9	PLUN 117
14	DT1=1.0/(C1+VS1)	PLUN 119
	DT1=AMU1*C1/((C1+VS1)**2)	PLUN 119
	DT=DT1*DX1	PLUN 120
	IF (DT.LT.DT1) DT1=DT	PLUN 121
	DO 15 M=1,6	PLUN 122
15	ITP(M)=FLOOR(ITP(M)/11.)	PLUN 123
	ISEG=ITP(M)	PLUN 124
	IF (ITAP-1) 17,16,17	PLUN 125
16	ISEG=ISEG+1	PLUN 126
	CALL RESET	PLUN 127
	DT1=AMU1*C1/((C1+VS1)**2)	PLUN 128
	DT=DT1*DX1	PLUN 129
	IF (DT.LT.DT1) DT1=DT	PLUN 130
	IF (ISTOP) 1,1,17	PLUN 131
17	IF (MGT) 13,41,21	PLUN 132
18	CALL STATE (1,1,1)	PLUN 133
	DUM=VSOUNM(P(1,1),V(1,1,1))	PLUN 134
	TCHECK=(TDRIT(1)+TDRIT(2)*DELS)/(V(1,1,2)+DUM)	PLUN 135
	MAXIT=INT(9.2*TCHECK/DT1)	PLUN 136
	DO 19 M=1,5	PLUN 137
19	ITP(M)=INT(ITP(M)*TCHECK/DT1)	PLUN 138
20	MDS=1	PLUN 139
	ITPNT=ITP(1)	PLUN 140
	IF (IPRINT.FG.1) CALL OMFOUT	PLUN 141
	DX2=DX2+DX1	PLUN 142
21	DO 24 IT=1,MAXIT	PLUN 143
	NY=-NS(1)*DE1	PLUN 144
	NY=NY*CTH(1)	PLUN 145
	NYSP=-NS(2)*DE1*CTH(2)	PLUN 146
	L=1	PLUN 147
	LP=?	PLUN 148
	CALL STATE (1,L,MAXN)	PLUN 149
	CALL STATE (2,LP,MAXN)	PLUN 150
	DT=DT1*NY	PLUN 151
	IF (DT1.LT.DT) DT=DT1	PLUN 152
	TIME=TIME+DT	PLUN 153
	IF (IMIN.FG.2) GO TO 24	PLUN 154
	NSH=NS(1)	PLUN 155
	DO 22 M=1,4	PLUN 156
22	P(MAXN,M)=V(1,MAXN,M)	PLUN 157
	DT=1.0/(NS(1)*DE2)	PLUN 158
	I=1	PLUN 159
	IP=2	PLUN 160
	IM=2	PLUN 161
	LM=2	PLUN 162

<KAP=1.5*(K(I)+K(J))	PLUN 163
DT=1.0	PLUN 164
DT=L=1.0	PLUN 165
DTI=1.0	PLUN 166
DTL=1.0	PLUN 167
V(1)=1.0	PLUN 168
V(2)=1.0	PLUN 169
V(3)=1.0	PLUN 170
V(4)=1.0	PLUN 171
ETA=1.0	PLUN 172
KAP=-KAP*V(1)*DT	PLUN 173
KAP=1.0+KAP	PLUN 174
V(1)=-KAP*KAP*V(1,2,1)*V(1,2,3)/DN	PLUN 175
V(2)=V(2,1,1)*V(2,1,2)/DXT	PLUN 176
V(3)=V(3,1)+V(3,2)+V(3,3)	PLUN 177
V(4)=V(4,1,1)-V(4,1,2)*DT	PLUN 178
V(1,2)=1.0	PLUN 179
ZZ=(V(1,1,4)-V(1,2,4))/(V(1,1,4)+V(1,2,4))	PLUN 180
VTS=ZZ*Z	PLUN 181
V(1,3)=1.0-VTS	PLUN 182
V(1,4)=AVN*V(1,2,4)+VTS*V(1,1,4)	PLUN 183
KAP=1.0	PLUN 184
DO 23 N=2,M	PLUN 185
ETA=ETA+DT	PLUN 186
KAP=KAP+KAP	PLUN 187
V(I)=V(I,N,1)*V(I,N,2)	PLUN 188
ZZ=(V(1,1,3)+V(1,2,3))*DT/DN	PLUN 189
AVN=ZZ*ZZ	PLUN 190
VTS=1.0-AVN	PLUN 191
V(1,3)=1.0*AVN	PLUN 192
V(I)=V(I,1,3)-W(I)*ETA	PLUN 193
M=M+1	PLUN 194
N=N-1	PLUN 195
DT=V(2,1,1)*V(2,N,2)/DXT	PLUN 196
DN=(V(1,1,2)*V(1,N,2)*(KAP+DKAP))*V(1,1,1)*V(1,N,2)*(KAP-DKAP)	PLUN 197
*DT**2)*DT	PLUN 198
V(1)=DT*(V(I,1,1)-V(I,N,1))	PLUN 199
V(2)=DT*(V(I,1,2)-V(I,N,2))	PLUN 200
V(3)=DT*(V(I,1,3)-V(I,N,3))	PLUN 201
V(4)=DT*(V(I,1,4)-V(I,N,4))	PLUN 202
V(1,1)=AVN*(V(I,1,1)+V(I,N,1))+VTS*V(I,1,1)	PLUN 203
V(1,2)=AVN*(V(I,1,2)+V(I,N,2))+VTS*V(I,1,2)	PLUN 204
V(1,3)=AVN*(V(I,1,3)+V(I,N,3))+VTS*V(I,1,3)	PLUN 205
V(1,4)=AVN*(V(I,1,4)+V(I,N,4))+VTS*V(I,1,4)	PLUN 206
V(2,1)=(DT+DT+DT)/KAP-KAP*(ETA*V(1,1))	PLUN 207
V(2,2)=V(2,1)*V(2,1)	PLUN 208
V(2,3)=V(2,1)*V(2,1)	PLUN 209
V(2,4)=V(2,1)*V(2,1)	PLUN 210
V(2,5)=V(2,1)*V(2,1)	PLUN 211
V(2,6)=V(2,1)*V(2,1)	PLUN 212
V(2,7)=V(2,1)*V(2,1)	PLUN 213
V(2,8)=V(2,1)*V(2,1)	PLUN 214
V(2,9)=V(2,1)*V(2,1)	PLUN 215
V(2,10)=V(2,1)*V(2,1)	PLUN 216
V(2,11)=V(2,1)*V(2,1)	PLUN 217

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20 20 26
21 20 25 M=1,MAXM
22 20 25 M=1,4
23 20 25 M=1,4
24 20 25 M=1,4
25 20 25 M=1,4
26 20 25 M=1,4
27 20 25 M=1,4
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97 20 25 M=1,4
98 20 25 M=1,4
99 20 25 M=1,4
100 20 25 M=1,4

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BLUN 271
BLUN 272

```

1000	V(I,N)=V(I,N)*K1	273	PLUN
1001	V(I,N)=V(I,N)-DUM	274	PLUN
1002	V(I,N)=V(I,N)	275	PLUN
1003	V(I,N)=V(I,N)/Z	276	PLUN
1004	V(I,N)=V(I,N)/DX	277	PLUN
1005	V(I,N)=V(I,N)	278	PLUN
1006	V(I,N)=V(I,N)	279	PLUN
1007	V(I,N)=V(I,N)	280	PLUN
1008	V(I,N)=V(I,N)	281	PLUN
1009	V(I,N)=V(I,N)	282	PLUN
1010	V(I,N)=V(I,N)	283	PLUN
1011	V(I,N)=V(I,N)	284	PLUN
1012	V(I,N)=V(I,N)	285	PLUN
1013	V(I,N)=V(I,N)	286	PLUN
1014	V(I,N)=V(I,N)	287	PLUN
1015	V(I,N)=V(I,N)	288	PLUN
1016	V(I,N)=V(I,N)	289	PLUN
1017	V(I,N)=V(I,N)	290	PLUN
1018	V(I,N)=V(I,N)	291	PLUN
1019	V(I,N)=V(I,N)	292	PLUN
1020	V(I,N)=V(I,N)	293	PLUN
1021	V(I,N)=V(I,N)	294	PLUN
1022	V(I,N)=V(I,N)	295	PLUN
1023	V(I,N)=V(I,N)	296	PLUN
1024	V(I,N)=V(I,N)	297	PLUN
1025	V(I,N)=V(I,N)	298	PLUN
1026	V(I,N)=V(I,N)	299	PLUN
1027	V(I,N)=V(I,N)	300	PLUN
1028	V(I,N)=V(I,N)	301	PLUN
1029	V(I,N)=V(I,N)	302	PLUN
1030	V(I,N)=V(I,N)	303	PLUN
1031	V(I,N)=V(I,N)	304	PLUN
1032	V(I,N)=V(I,N)	305	PLUN
1033	V(I,N)=V(I,N)	306	PLUN
1034	V(I,N)=V(I,N)	307	PLUN
1035	V(I,N)=V(I,N)	308	PLUN
1036	V(I,N)=V(I,N)	309	PLUN
1037	V(I,N)=V(I,N)	310	PLUN
1038	V(I,N)=V(I,N)	311	PLUN
1039	V(I,N)=V(I,N)	312	PLUN
1040	V(I,N)=V(I,N)	313	PLUN
1041	V(I,N)=V(I,N)	314	PLUN
1042	V(I,N)=V(I,N)	315	PLUN
1043	V(I,N)=V(I,N)	316	PLUN
1044	V(I,N)=V(I,N)	317	PLUN
1045	V(I,N)=V(I,N)	318	PLUN
1046	V(I,N)=V(I,N)	319	PLUN
1047	V(I,N)=V(I,N)	320	PLUN
1048	V(I,N)=V(I,N)	321	PLUN
1049	V(I,N)=V(I,N)	322	PLUN
1050	V(I,N)=V(I,N)	323	PLUN
1051	V(I,N)=V(I,N)	324	PLUN
1052	V(I,N)=V(I,N)	325	PLUN
1053	V(I,N)=V(I,N)	326	PLUN
1054	V(I,N)=V(I,N)	327	PLUN

```

37 (1-10)*40(1),J1.1.17-10) GO TO 37
38 (1-10)*40(1),J1.1.17-10) GO TO 38
39 (1-10)*40(1),J1.1.17-10) GO TO 39
40 (1-10)*40(1),J1.1.17-10) GO TO 40
41 (1-10)*40(1),J1.1.17-10) GO TO 41
42 (1-10)*40(1),J1.1.17-10) GO TO 42
43 (1-10)*40(1),J1.1.17-10) GO TO 43
44 (1-10)*40(1),J1.1.17-10) GO TO 44
45 (1-10)*40(1),J1.1.17-10) GO TO 45
46 (1-10)*40(1),J1.1.17-10) GO TO 46
47 (1-10)*40(1),J1.1.17-10) GO TO 47
48 (1-10)*40(1),J1.1.17-10) GO TO 48
49 (1-10)*40(1),J1.1.17-10) GO TO 49
50 (1-10)*40(1),J1.1.17-10) GO TO 50
51 (1-10)*40(1),J1.1.17-10) GO TO 51
52 (1-10)*40(1),J1.1.17-10) GO TO 52
53 (1-10)*40(1),J1.1.17-10) GO TO 53
54 (1-10)*40(1),J1.1.17-10) GO TO 54
55 (1-10)*40(1),J1.1.17-10) GO TO 55
56 (1-10)*40(1),J1.1.17-10) GO TO 56
57 (1-10)*40(1),J1.1.17-10) GO TO 57
58 (1-10)*40(1),J1.1.17-10) GO TO 58
59 (1-10)*40(1),J1.1.17-10) GO TO 59
60 (1-10)*40(1),J1.1.17-10) GO TO 60
61 (1-10)*40(1),J1.1.17-10) GO TO 61
62 (1-10)*40(1),J1.1.17-10) GO TO 62
63 (1-10)*40(1),J1.1.17-10) GO TO 63
64 (1-10)*40(1),J1.1.17-10) GO TO 64
65 (1-10)*40(1),J1.1.17-10) GO TO 65
66 (1-10)*40(1),J1.1.17-10) GO TO 66
67 (1-10)*40(1),J1.1.17-10) GO TO 67
68 (1-10)*40(1),J1.1.17-10) GO TO 68
69 (1-10)*40(1),J1.1.17-10) GO TO 69
70 (1-10)*40(1),J1.1.17-10) GO TO 70
71 (1-10)*40(1),J1.1.17-10) GO TO 71
72 (1-10)*40(1),J1.1.17-10) GO TO 72
73 (1-10)*40(1),J1.1.17-10) GO TO 73
74 (1-10)*40(1),J1.1.17-10) GO TO 74
75 (1-10)*40(1),J1.1.17-10) GO TO 75
76 (1-10)*40(1),J1.1.17-10) GO TO 76
77 (1-10)*40(1),J1.1.17-10) GO TO 77
78 (1-10)*40(1),J1.1.17-10) GO TO 78
79 (1-10)*40(1),J1.1.17-10) GO TO 79
80 (1-10)*40(1),J1.1.17-10) GO TO 80
81 (1-10)*40(1),J1.1.17-10) GO TO 81
82 (1-10)*40(1),J1.1.17-10) GO TO 82
83 (1-10)*40(1),J1.1.17-10) GO TO 83
84 (1-10)*40(1),J1.1.17-10) GO TO 84
85 (1-10)*40(1),J1.1.17-10) GO TO 85
86 (1-10)*40(1),J1.1.17-10) GO TO 86
87 (1-10)*40(1),J1.1.17-10) GO TO 87
88 (1-10)*40(1),J1.1.17-10) GO TO 88
89 (1-10)*40(1),J1.1.17-10) GO TO 89
90 (1-10)*40(1),J1.1.17-10) GO TO 90
91 (1-10)*40(1),J1.1.17-10) GO TO 91
92 (1-10)*40(1),J1.1.17-10) GO TO 92
93 (1-10)*40(1),J1.1.17-10) GO TO 93
94 (1-10)*40(1),J1.1.17-10) GO TO 94
95 (1-10)*40(1),J1.1.17-10) GO TO 95
96 (1-10)*40(1),J1.1.17-10) GO TO 96
97 (1-10)*40(1),J1.1.17-10) GO TO 97
98 (1-10)*40(1),J1.1.17-10) GO TO 98
99 (1-10)*40(1),J1.1.17-10) GO TO 99
100 (1-10)*40(1),J1.1.17-10) GO TO 100

```

NOT REPRODUCIBLE

```

BLUN 328
BLUN 329
BLUN 330
BLUN 331
BLUN 332
BLUN 333
BLUN 334
BLUN 335
BLUN 336
BLUN 337
BLUN 338
BLUN 339
BLUN 340
BLUN 341
BLUN 342
BLUN 343
BLUN 344
BLUN 345
BLUN 346
BLUN 347
BLUN 348
BLUN 349
BLUN 350
BLUN 351
BLUN 352
BLUN 353
BLUN 354
BLUN 355
BLUN 356
BLUN 357
BLUN 358
BLUN 359

```

4. SUBROUTINE BODY

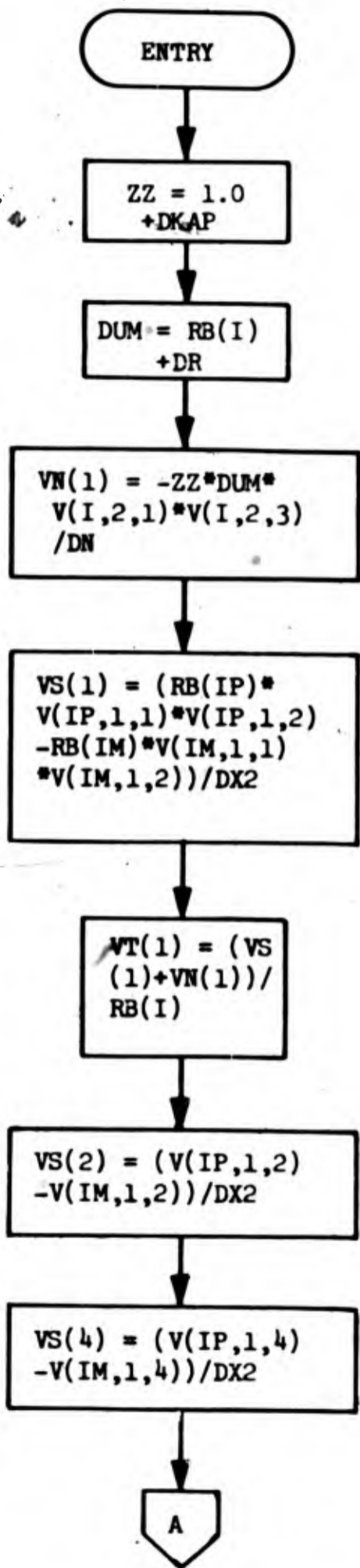
A. PURPOSE

This subroutine solves the governing equations at nodes on the body surface, except for the stagnation point. The special body point logic developed in section IV is employed.

B. VARIABLE LIST

All variables used in this subroutine have been defined in the description of the main program.

C. FLOW CHART



ENTER SUBROUTINE BODY.

COMPUTE κ AT $\eta = \Delta\eta$.

COMPUTE r AT $\eta = \Delta\eta$.

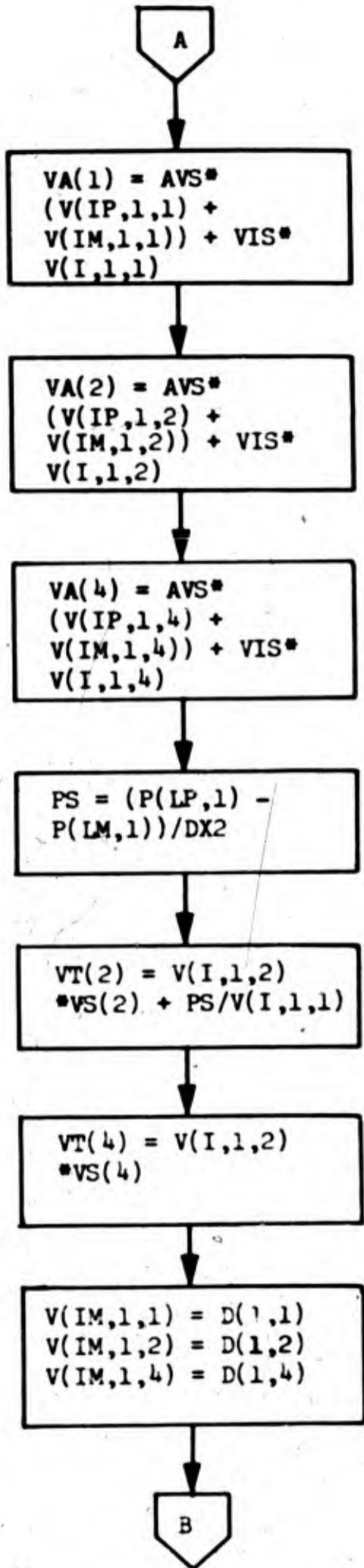
COMPUTE $(\kappa r \rho V)_n$.

COMPUTE $(r \rho u)_3$.

COMPUTE ρ_τ .

COMPUTE u_s .

COMPUTE S_s .



COMPUTE AVERAGE VALUE OF ρ
TO INTRODUCE THE STABILIZING
TERM.

COMPUTE AVERAGE VALUE OF u
TO INTRODUCE THE STABILIZING
TERM.

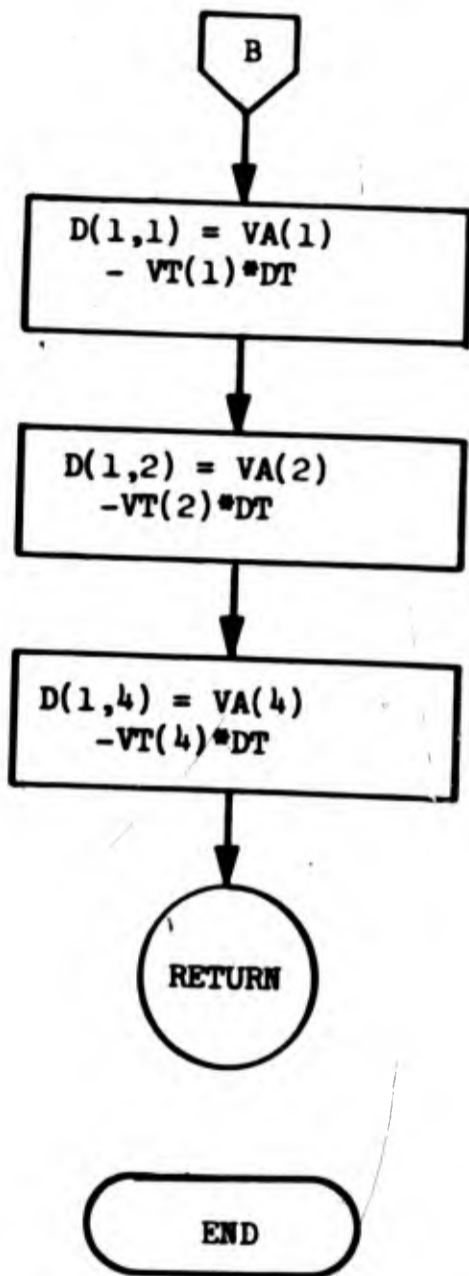
COMPUTE AVERAGE VALUE OF S
TO INTRODUCE THE STABILIZING
TERM.

COMPUTE P_s .

COMPUTE U_t .

COMPUTE S_t .

TRANSFER ρ , U AND S AT
 $S-\Delta S$ AND $t+\Delta t$ FROM TEMPORARY
TO PERMANENT STORAGE.



SOLVE FOR ρ AT $\tau + \Delta\tau$.

SOLVE FOR u AT $\tau + \Delta\tau$.

SOLVE FOR S AT $\tau + \Delta\tau$.

RETURN.

END STATEMENT.

0. LISTING

```

SUBROUTINE BODY (CM, CKAP, JN)
C
C THIS ROUTINE TREATS THE BODY ON THE BODY SURFACE EXCLUSIVE
C OF THE STAGNATION POINT.
C
COMMON /1/ V(15,4), J(4), P(3,4)
COMMON /2/ PR(15), ZR(15), TH(15), CTH(15), K(15), NS(15), WC(15)
COMMON /5/ DXI, DXT, DX2, DT2, DT
COMMON /7/ I, IE, IM, L, LB, LM
COMMON /17/ VS(4), VN(4), VA(4), VT(4)
COMMON /17/ AVS, AVN, AVT
C
REAL NS, K, KAP
ZZ=1.0+KAP
DUM=JN(I)+JN
VN(1)=-Z*SUM*V(I,2,1)*V(I,2,1)/DN
VS(1)=(P(I,IP)+V(IP,1,1)+V(IP,1,2)-PR(ITM)+V(IM,1,1)+V(IM,1,2))/DX2
VT(1)=(VS(1)+VN(1))/ZR(I)
VS(2)=(V(IP,1,2)-V(IM,1,2))/DX2
VS(4)=(V(IP,1,4)-V(IM,1,4))/DX2
VA(1)=AVS*(V(IP,1,1)+V(IM,1,1))+VIS*V(I,1,1)
VA(2)=AVS*(V(IE,1,2)+V(IM,1,2))+VTS*V(I,1,2)
VA(4)=AVS*(V(IP,1,4)+V(IM,1,4))+VIS*V(I,1,4)
PS=(P(LC,1)-P(LM,1))/DX2
VT(2)=V(T,1,2)+VS(2)+PS/V(I,1,1)
VT(4)=V(I,1,2)+VS(4)
V(TM,1,1)=J(1,1)
V(TM,1,2)=J(1,2)
V(IM,1,4)=J(1,4)
J(1,1)=VA(1)-VT(1)+DT
J(1,2)=VA(2)-VT(2)+DT
J(1,4)=VA(4)-VT(4)+DT
RETURN
END

```

```

BODY 1
BODY 2
BODY 3
BODY 4
BODY 5
BODY 6
BODY 7
BODY 8
BODY 9
BODY 17
BODY 11
BODY 12
BODY 13
BODY 14
BODY 15
BODY 16
BODY 17
BODY 18
BODY 19
BODY 20
BODY 21
BODY 22
BODY 23
BODY 24
BODY 25
BODY 26
BODY 27
BODY 28
BODY 29
BODY 30
BODY 31
BODY 32
BODY 33

```

5. SUBROUTINE COEFF

A. PURPOSE

This subroutine generates coefficients for a cubic fit of the data supplied under the arbitrary body shape option. Conventional curve fit techniques such as least squares or spline fit techniques were found to be inadequate for this application. The present curve fit procedure requires that the user supply a fairly complete description of the body geometry at all data points supplied (r , z , θ , and K). It is recognized that approximate methods may be required to generate this data for many body shapes of interest. However, this data is often known explicitly, but is not easily specified at equal intervals along the body, e.g., an ellipse. It was decided that an accurate technique should be incorporated into the program, leaving the choice of more approximate schemes to the user when that becomes necessary. The cubic fit is accomplished using arrays which specify the values of x , y and $\frac{dy}{dx}$ at a number of points. The relevant constraint relations for a cubic fit can be expressed as

$$y_j = \sum_{i=1}^4 c_{ij} x_j^{(i-1)} \quad (64)$$

$$\left(\frac{dy}{dx}\right)_j = \sum_{i=2}^4 (i-2) c_{ij} x_j^{(i-2)} \quad (65)$$

At end points of the array of X_y values, the values of y_j and $\left(\frac{dy}{dx}\right)_j$ at 2 adjacent points are used to generate the coefficient c_{ij} . Defining the matrixes

$$[B_{ij}] = \begin{bmatrix} 1 & x_1 & x_1^2 & x_1^3 \\ 1 & x_2 & x_2^2 & x_2^3 \\ 0 & 1 & 2x_1 & 3x_1^2 \\ 0 & 1 & 2x_2 & 3x_2^2 \end{bmatrix} \quad (66)$$

$$[Y_j] = \begin{bmatrix} y_1 \\ y_2 \\ (\frac{dy}{dx})_1 \\ (\frac{dy}{dx})_2 \end{bmatrix} \quad (67)$$

a matrix, $[A]$, is generated from $[B]$ by replacing the i^{th} column with $[Y_j]$ and the determinants $|A|$ and $|B|$ are used to generate C_{ij}

$$C_{ij} = |A_{ij}| / |B_{ij}| \quad (68)$$

For interior points, the values of y at 2 adjacent points are used with the local value and slope to define $[B_{ij}]$ by

$$[B_{ij}] = \begin{bmatrix} 1 & x_1 & x_1^2 & x_1^3 \\ 1 & x_2 & x_2^2 & x_2^3 \\ 1 & x_3 & x_3^2 & x_3^3 \\ 0 & 1 & 2x_2 & 3x_1^2 \end{bmatrix} \quad (69)$$

where the point under consideration is presumed to be x_2 . Using

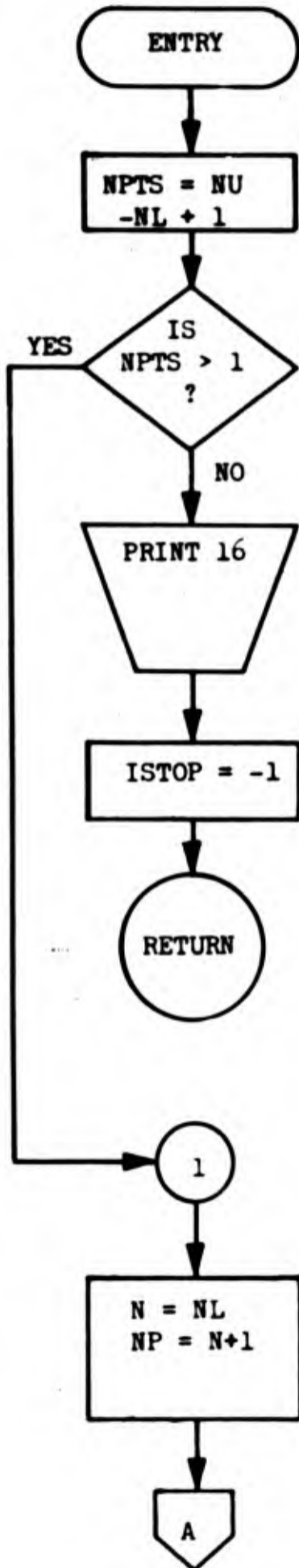
$$[Y_j] = \begin{bmatrix} y_1 \\ y_2 \\ y_3 \\ (\frac{dy}{dx})_2 \end{bmatrix} \quad (70)$$

Equation (68) supplies the desired cubic coefficients for the interior points.

B. VARIABLE LIST

- A - $[A_{ij}]$
- B - $[B_{ij}]$
- C - $[C_{ij}]$
- DET - $|B|$
- DETC - $|A|$
- DYDX - array containing values of $\frac{dy}{dx}$.
- I - index for DO loops.
- J - index for DO loops.
- K - index for DO loops.
- MN - index of last interior point.
- N - index specifying the location of the point under consideration in the input arrays.
- NEX - integer exponent used in loading arrays.
- NL - index of first point to be considered.
- NM - $N - 1$
- NN - index of first interior point.
- NP - $N + 1$
- NPTS - number of points to be considered.
- NU - index of the last point.
- X - array containing values of X.
- Y - array containing values of Y.
- YC - $[Y_j]$

C. FLOW CHART



ENTER SUBROUTINE COEFF.

COMPUTE TOTAL NUMBER OF POINTS
TO BE CONSIDERED.

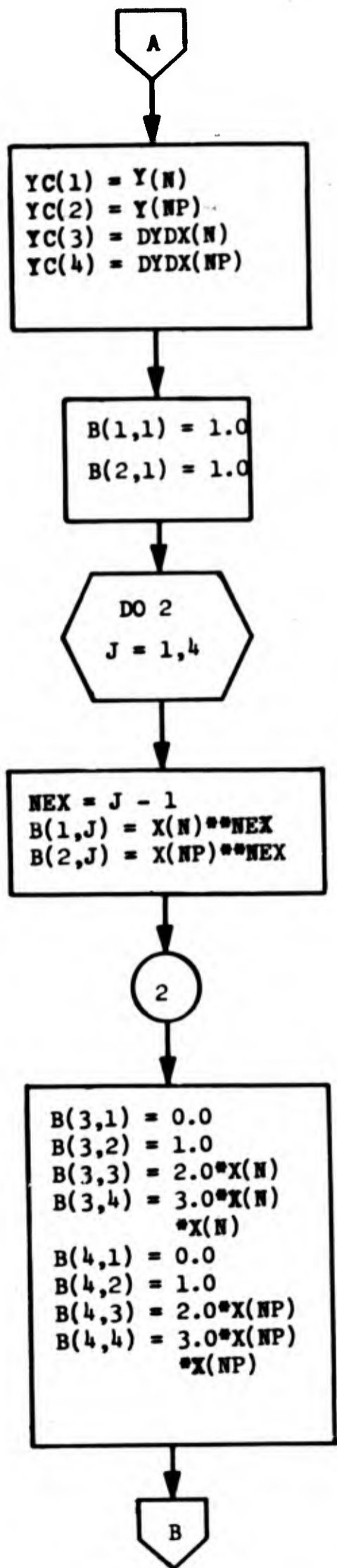
WERE AT LEAST 2 POINTS
PROVIDED ?

PRINT ERROR MESSAGE.

SET ERROR INDICATOR.

RETURN.

SET INDICES FOR FIRST
POINT.

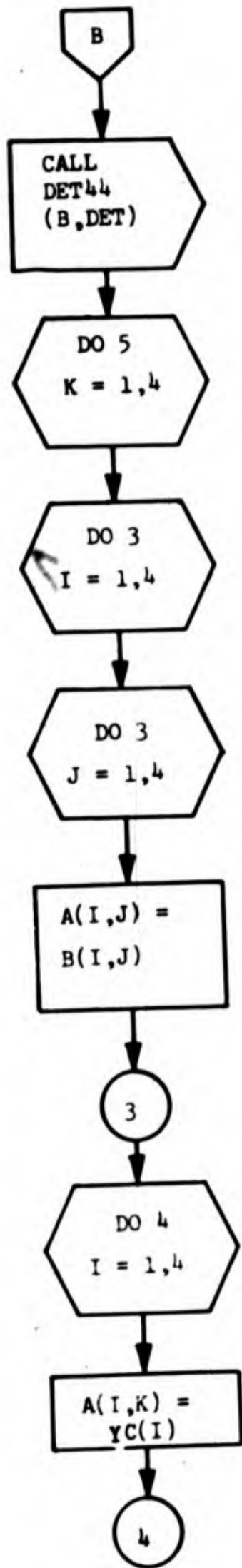


LOAD THE YC ARRAY FOR THE FIRST POINT.

SET b_{11} AND b_{21} .

LOAD COLUMNS 2 TO 3 OF ROWS 1 AND 2 OF THE B ARRAY FOR THE FIRST POINT.

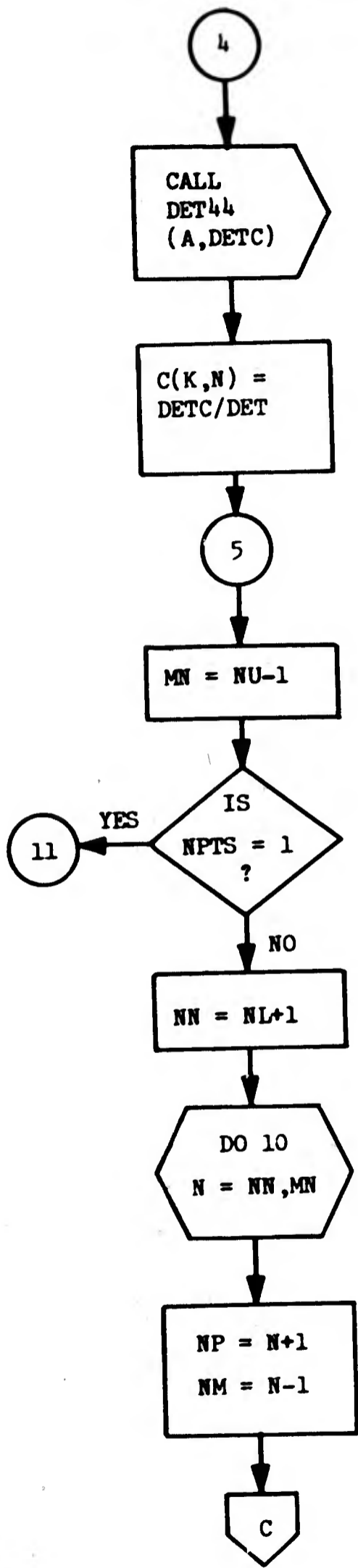
LOAD REMAINING TERMS IN THE B ARRAY FOR THE FIRST POINT.



COMPUTE DETERMINANT OF B.

LOAD A ARRAY FOR THE FIRST POINT.

LOAD YC ARRAY INTO THE A ARRAY FOR COMPUTATION OF C_k

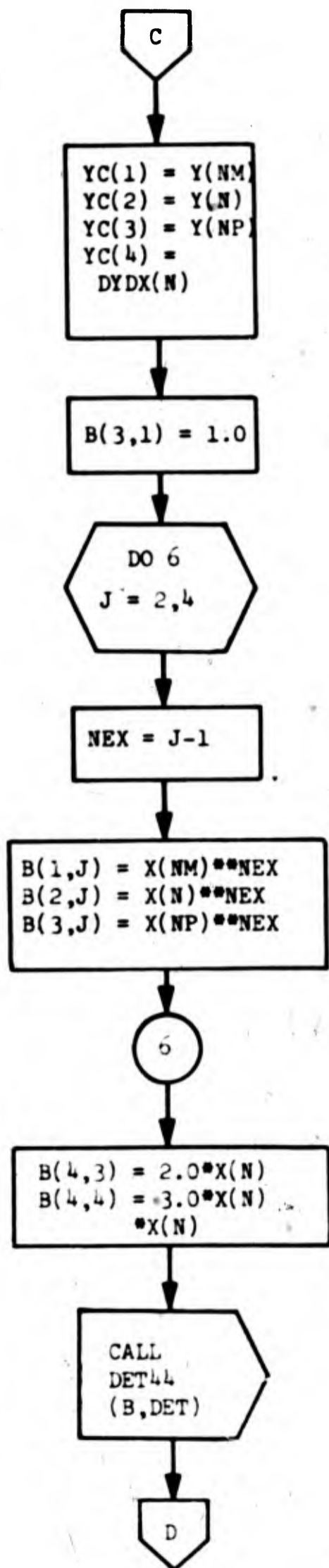


COMPUTE DETERMINANT OF A.

COMPUTE C_k 's FOR THE FIRST POINT.

WERE ONLY 2 POINTS PROVIDED ?

START LOOP FOR INTERIOR POINTS.



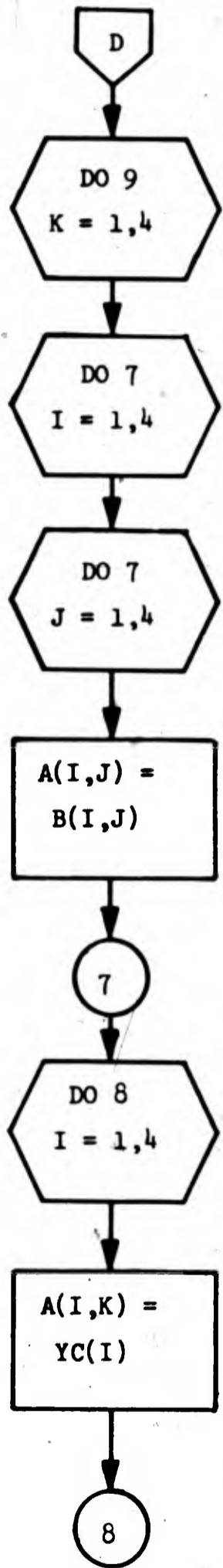
LOAD YC ARRAY.

COMPUTE EXPONENT.

LOAD COLUMNS 2, 3 AND 4
OF ROWS 1, 2 AND 3.

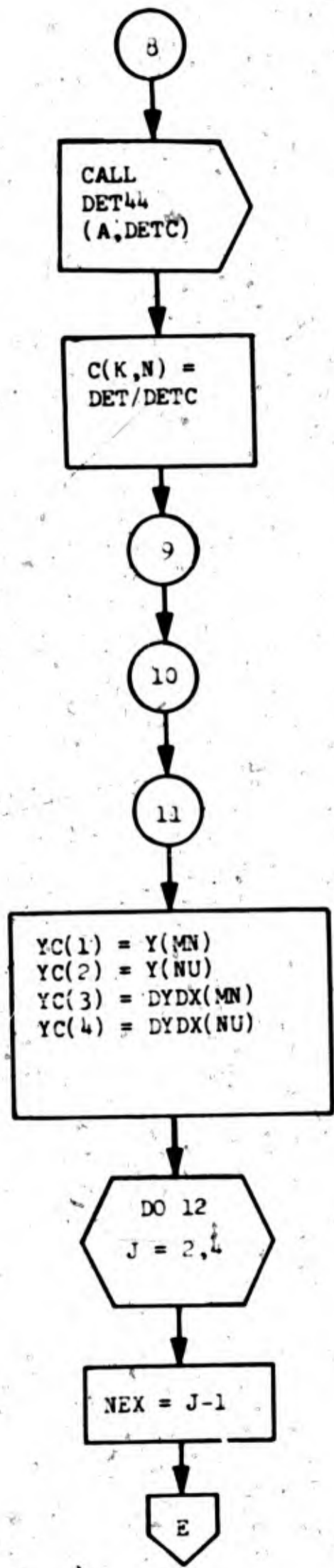
LOAD REMAINING TERMS.

COMPUTE DETERMINATION OF B.



LOAD A ARRAY FROM B ARRAY.

LOAD YC VALUES INTO THE KTH
COLUMN OF THE A ARRAY.



COMPUTE DETERMINANT OF A.

COMPUTE C_k 's.

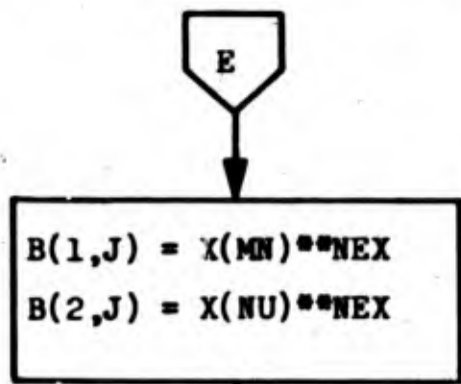
END OF DO LOOP ON K.

END OF DO LOOP OVER ALL INTERIOR POINTS.

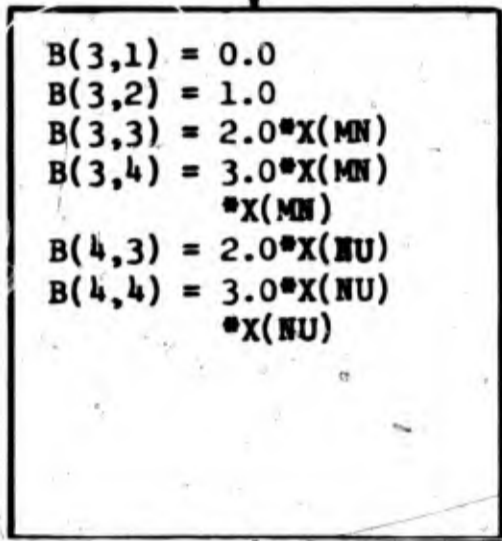
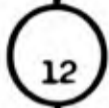
START COMPUTATIONS FOR THE LAST POINT.

LOAD YC FOR LAST POINT.

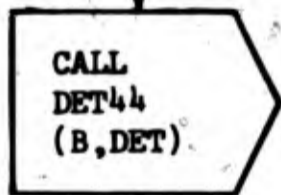
COMPUTE EXPONENT FOR THE J^{TH} COLUMN.



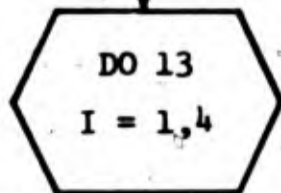
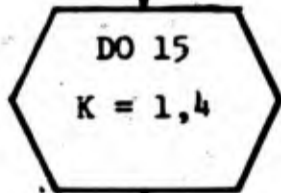
LOAD COLUMNS 2 TO 4 OF
ROWS 1 AND 2 OF THE B
ARRAY.

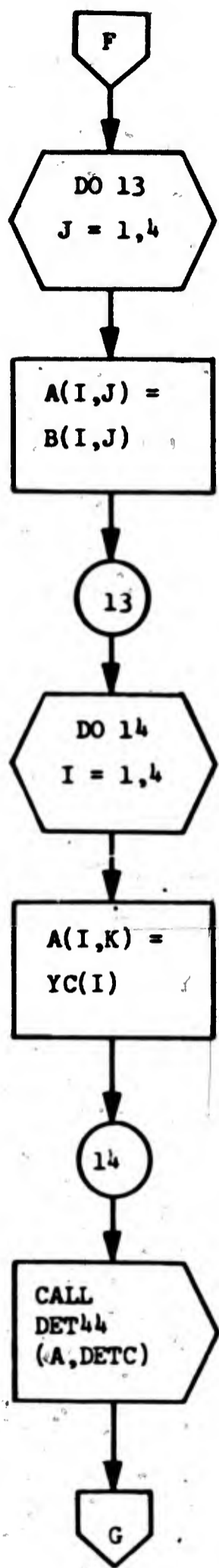


LOAD REMAINING TERMS
INTO THE B ARRAY.



COMPUTE DETERMINANT OF B.

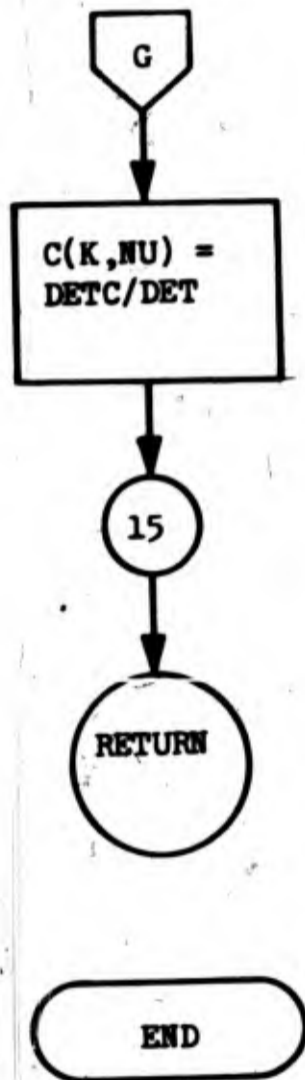




LOAD A ARRAY FROM THE B ARRAY.

LOAD YC VALUES INTO KTH COLUMN OF A ARRAY.

COMPUTE DETERMINANT OF A.



COMPUTE C_k FOR THE LAST POINT.

RETURN.

END

SUBROUTINE COEFF (NL,NU,X,Y,DYDX)

COEFF 1
COEFF 2
COEFF 3
COEFF 4
COEFF 5
COEFF 6
COEFF 7
COEFF 8
COEFF 9
COEFF 10
COEFF 11
COEFF 12
COEFF 13
COEFF 14
COEFF 15
COEFF 16
COEFF 17
COEFF 18
COEFF 19
COEFF 20
COEFF 21
COEFF 22
COEFF 23
COEFF 24
COEFF 25
COEFF 26
COEFF 27
COEFF 28
COEFF 29
COEFF 30
COEFF 31
COEFF 32
COEFF 33
COEFF 34
COEFF 35
COEFF 36
COEFF 37
COEFF 38
COEFF 39
COEFF 40
COEFF 41
COEFF 42
COEFF 43
COEFF 44
COEFF 45
COEFF 46
COEFF 47
COEFF 48
COEFF 49
COEFF 50
COEFF 51
COEFF 52

THIS ROUTINE COMPUTES THE COEFFICIENTS FOR A CUBIC FIT WHERE THE VALUE OF THE FUNCTION AND ITS SLOPE AT THE POINT UNDER CONSIDERATION AND ITS VALUES AT THE TWO ADJACENT POINTS ARE USED.

```
COMMON /11/ TOLMMY,ISTOP
COMMON /20LINE/ C(4,15)
DIMENSION X(15), Y(15), DYDX(15), VC(4), R(4,4), B(4,4)
NPTS=NU-1L+1
IF (NPTS,01,1) GO TO 1
POINT 16
ISTOP=-1
RETURN
N=NL
NP=N+1
VC(1)=Y(N)
VC(2)=Y(NP)
VC(3)=DYDX(N)
VC(4)=DYDX(NP)
R(1,1)=1.0
R(2,1)=1.0
DO 2 J=2,4
RTY=J-1
R(1,J)=X(N)**RTY
R(2,J)=Y(NP)**RTY
R(3,1)=1.0
R(3,2)=1.0
R(3,3)=2.0*X(N)
R(3,4)=2.0*X(N)*X(N)
R(4,1)=1.0
R(4,2)=1.0
R(4,3)=2.0*X(NP)
R(4,4)=3.0*X(NP)*X(NP)
CALL DET44 (R,DET)
DO 5 K=1,4
DO 3 I=1,4
DO 3 J=1,4
A(I,J)=R(I,J)
DO 4 I=1,4
A(I,K)=Y(VC(I))
CALL DET44 (A,DETC)
C(K,N)=DETC/DET
NN=NU-1
IF (NPTS,02,2) GO TO 11
NN=NL+1
DO 11 N=NN,NA
NP=N+1
NM=N-1
VC(1)=Y(NM)
VC(2)=Y(N)
VC(3)=Y(NP)
```

NOT REPRODUCIBLE

	V(4)=DYX(N)	COEFF 53
	R(1,1)=1.0	COEFF 54
	DO 5 J=2,4	COEFF 55
	IFY=J-1	COEFF 56
	R(1,J)=X(M)*IFY	COEFF 57
	R(2,J)=X(N)*IFY	COEFF 58
6	R(3,J)=X(M)*IFY	COEFF 59
	R(4,3)=2.0*X(N)	COEFF 61
	R(4,4)=2.0*X(N)*X(N)	COEFF 61
	CALL DET44 (8,DET)	COEFF 62
	DO 9 K=1,4	COEFF 63
	DO 7 I=1,4	COEFF 64
	DO 7 J=1,4	COEFF 65
7	A(I,J)=R(I,J)	COEFF 66
	DO 8 T=1,4	COEFF 67
8	A(I,K)=Y(I)	COEFF 68
	CALL DET44 (4,DETC)	COEFF 69
9	C(K,N)=DETC/DET	COEFF 71
11	CONTINUE	COEFF 71
11	Y(1)=Y(M)	COEFF 72
	Y(2)=Y(N)	COEFF 73
	Y(3)=DYX(M)	COEFF 74
	Y(4)=DYX(N)	COEFF 75
	DO 12 J=2,4	COEFF 76
	IFY=J-1	COEFF 77
	R(1,J)=X(M)*IFY	COEFF 74
12	R(2,J)=X(N)*IFY	COEFF 79
	R(3,1)=1.0	COEFF 81
	R(3,2)=1.0	COEFF 82
	R(3,3)=2.0*X(M)	COEFF 83
	R(3,4)=2.0*X(M)*X(M)	COEFF 83
	R(4,3)=2.0*X(N)	COEFF 84
	R(4,4)=2.0*X(N)*X(N)	COEFF 85
	CALL DET44 (8,DET)	COEFF 86
	DO 15 K=1,4	COEFF 87
	DO 13 I=1,4	COEFF 88
	DO 13 J=1,4	COEFF 89
13	A(I,J)=R(I,J)	COEFF 91
	DO 14 T=1,4	COEFF 91
14	A(I,K)=Y(I)	COEFF 92
	CALL DET44 (4,DETC)	COEFF 93
15	C(K,NU)=DETC/DET	COEFF 94
	DETC=N	COEFF 95
16	FORMAT (//3IX,25HERROR IN SUBROUTINE COEFF/1IX,33HNUMBER OF POINTS	COEFF 96
	* IS LESS THAN TWO/1IX,15HCASE TERMINATED)	COEFF 97
	END	COEFF 99

6. SUBROUTINE DATA

A. PURPOSE

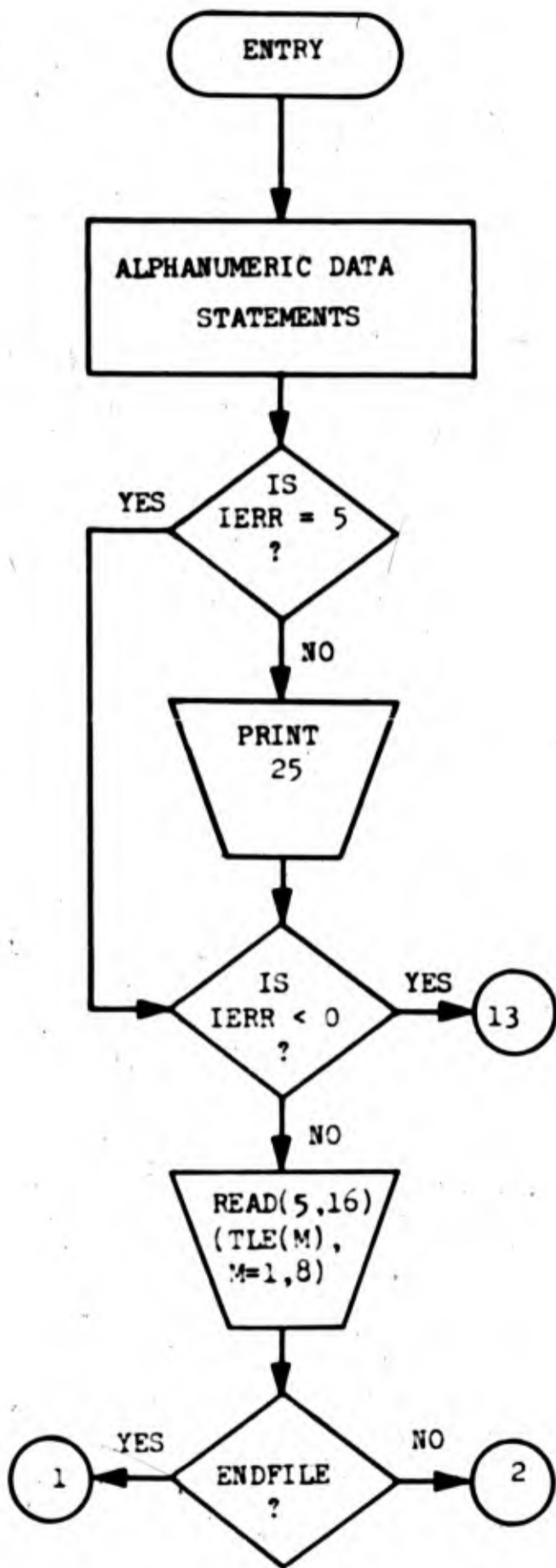
This subroutine loads all of the input data except for the arbitrary body shape data. It sorts the input data by matching alphanumeric specifications on the data cards with known variable names and with special alphanumeric flags signaling the start and end of a data packet. Consequently, the number of cards provided and the order of these cards are arbitrary. Since the computer code does not alter the values of the input variables, only variables whose resident values are to be changed need be input by the user. This simplifies the input data required for multicase runs. The principal reason for including a sorting routine of this type was to greatly simplify the input data required without imposing unnecessary restrictions on the numerical method. In addition to the empirical constants required for the numerical method, several option control variables had to be included. For example, to apply this method as a practical computational technique, a restart option had to be included since the exact number of iterations that will be required cannot be specified in advance. Also, if a steady state is not realized for a problem, an option to recover the data and perform additional iterations should be available. Other input controls considered desirable are specifications for additional flow field outputs, dimensional output, and nodal geometry. As a result, 23 input variables were required for a completely general computational method. Most of these variables have been supplied initial resident values in the main program which will be adequate for the majority of the problems encountered. Consequently, while the capability of inputting values for all 23 variables is available, the input required is usually quite brief, typically 3 to 5 cards. It is easily seen that without a routine of this type, a general input capability would be very complicated.

B. VARIABLE LIST

- ENFILE - alphanumeric word LAST in A10 format. Used to check for the end of the data packet.
- I - index for an implied DO loop in a print statement.
- IERR - error indicator. Called ISTOP in the main program.
- IV - array containing all integer variables in COMMON block /INPUT/.
- K - array used for temporary storage of the values of input integer variables.

- M - index used for various applications.
- T - array containing the alphanumeric names of all allowed input variables.
- TITLE - alphanumeric word TITLE in A10 format. Used to identify the first card in a data packet.
- TLE - array used to read and print an alphanumeric title.
- V - array containing all floating point variables in COMMON block /INPUT/.
- VAR - variable used to read the input variable name from data cards.
- ZZ - variable used for temporary storage of the values of input floating point variables.

C. FLOW CHART



ENTER SUBROUTINE DATA.

SET ALPHANUMERIC TITLES OF THE CASE DATA VARIABLES.

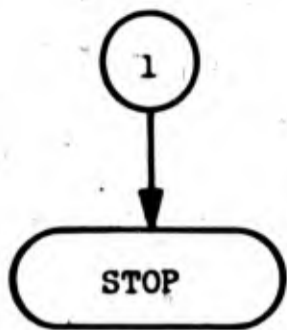
FIRST CASE ?

START A NEW PAGE.

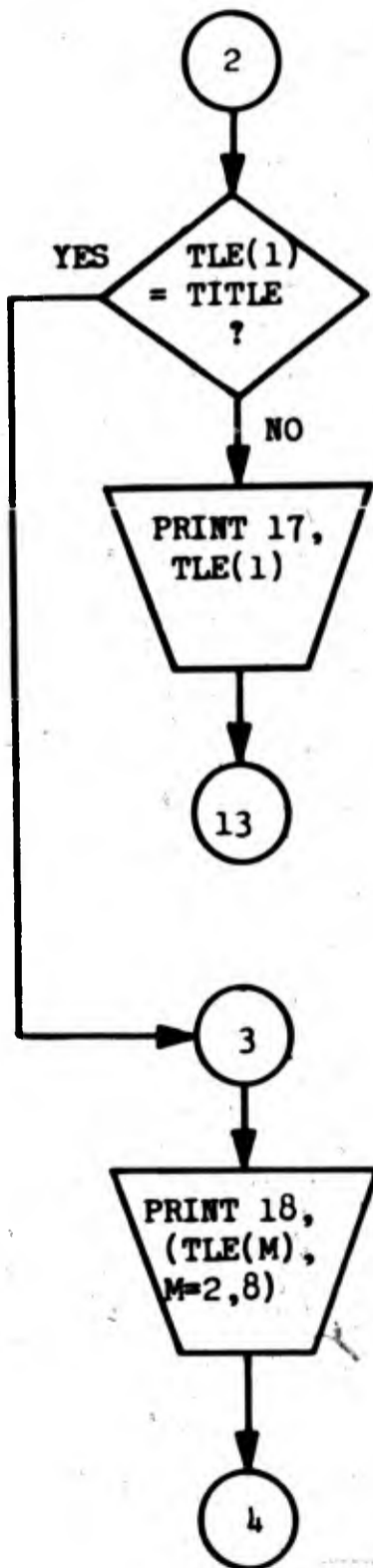
ERROR SENSED IN PREVIOUS CASE ?

READ TITLE CARD.

ENDFILE ENCOUNTERED ?



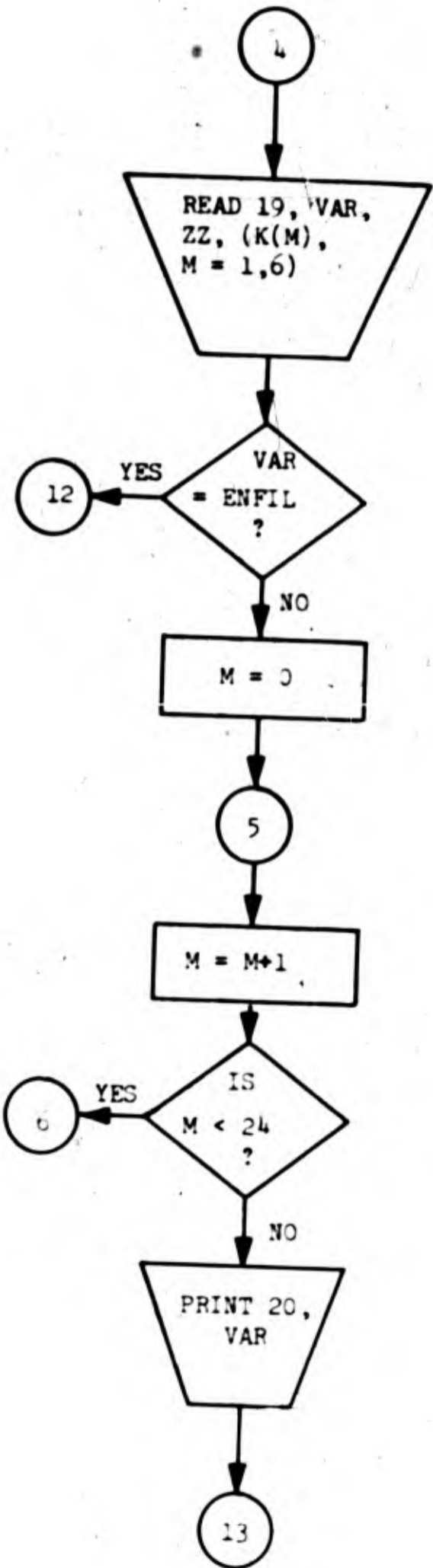
TERMINATE JOB.



WAS TITLE INDICATOR
CORRECT ?

PRINT ERROR MESSAGE.

PRINT TITLE.



READ CASE DATA CARD.

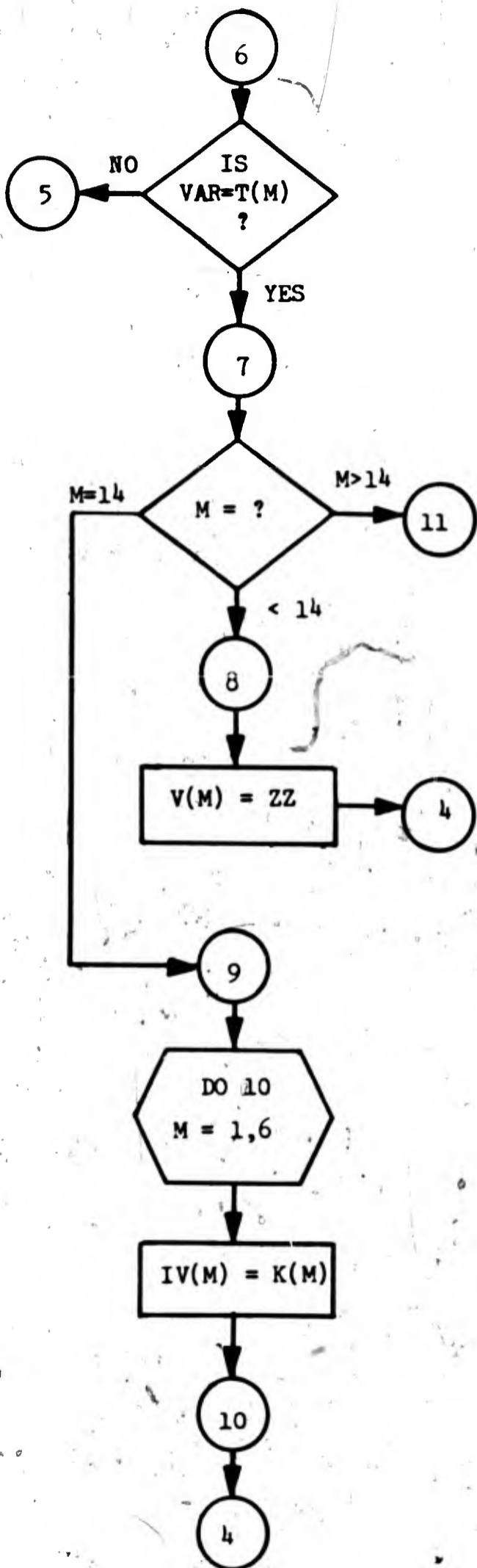
LAST CARD IN THE CASE ?

INITIALIZE INDEX, M, TO SEARCH TABLE OF VARIABLE NAMES FOR THE PRESENT VARIABLE.

INCREMENT M.

IS M WITHIN THE ALLOWABLE NUMBER OF CASE VARIABLES?

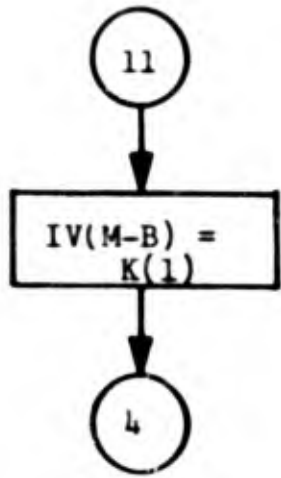
VARIABLE NAME COULD NOT BE MATCHED WITH AN INPUT VARIABLE. PRINT ERROR MESSAGE.



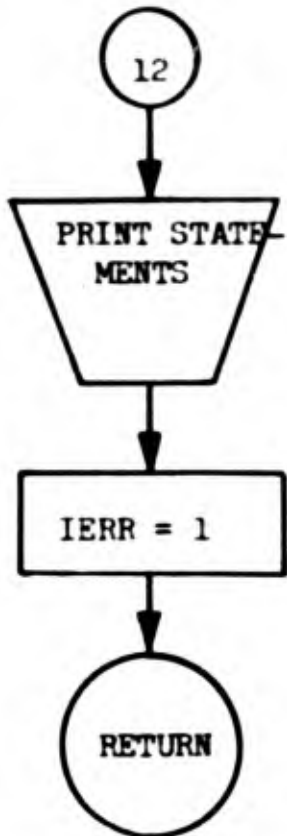
DOES VARIABLE NAME READ
MATCH THE NAME STORED IN
T(M) ?

FLOATING POINT VARIABLE.
STORE IN PROPER COMMON
LOCATION.

IPR(6) ARRAY WAS READ.
STORE IN PROPER COMMON
LOCATION.



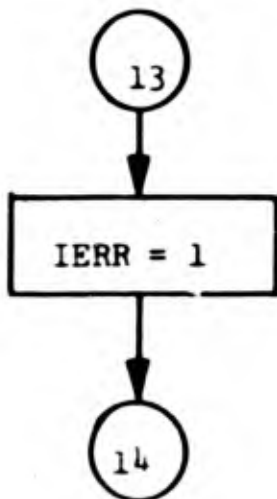
INTEGER VARIABLE. STORE
IN PROPER COMMON LOCATION.



PRINT INPUT DATA.

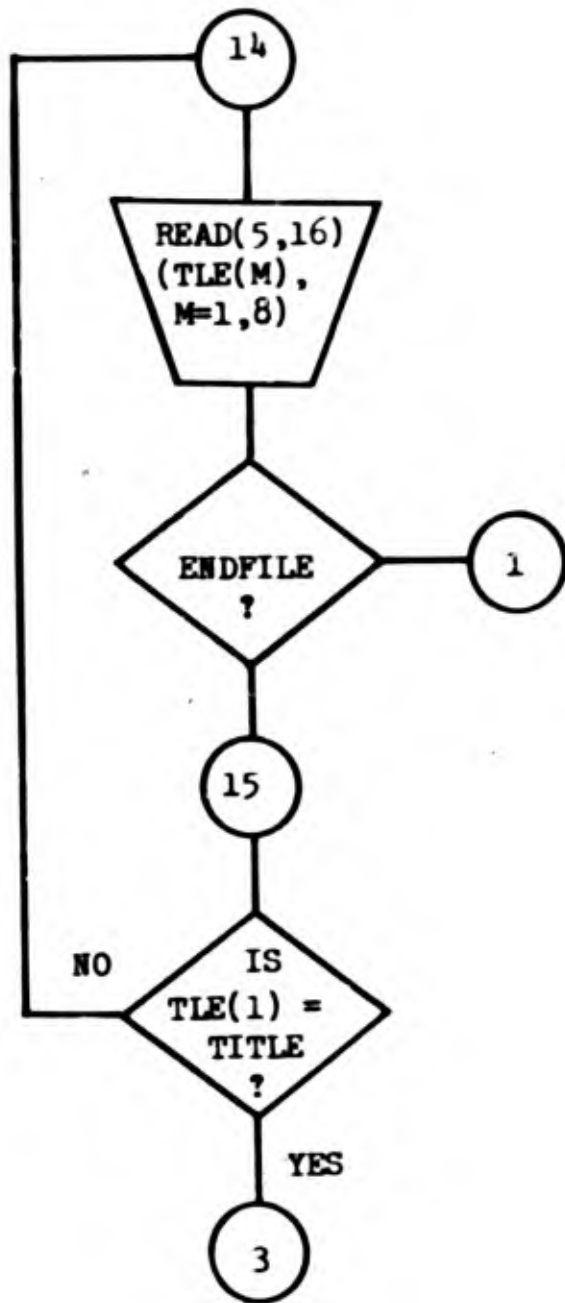
SET ERROR INDICATOR.

RETURN.



START PROCEDURES TO
BYPASS ERRONEOUS DATA.

SET ERROR INDICATOR.



READ DATA CARD.

WAS AN ENDFILE
ENCOUNTERED ?

WAS A TITLE CARD
READ ?

3. LISTING

```

CIRCUIT DATA (IFP)
    THIS ROUTINE LOADS THE CASE DATA.

    DO 10 M=1, NFILE(3), VAP, VAP
    DO 10 M=1, NFILE(4), IV(15)
    DO 10 M=1, NFILE(5), K(5)
    DATA (T(M), M=1, 23) /
    *I10GAMMA  ,I10MM           ,I10THETA      ,I10WZ/PK      *
    *I10TODITI  ,I10TODITD     ,I10MP1       ,I10D0401     *
    *I10MU      ,I10MU         ,I10MT1       ,I10E00000M   *
    *I10E00000  ,I10IPR       ,I10INTAD     ,I10CTAD      *
    *I10MAXIT   ,I10MCASE     ,I10NFITL    ,I10AC0G     *
    *I10MAXN    ,I10MAXN      ,I10NTAD     /
    DATA TITLE, ENFILE/I10TITLE
    IF (IFP2,LT,3) GO TO 17
    IF (IFP2,GE,3) PRINT 25
    DATA (5,16) (TLE(M), M=1, 8)
    IF (ENDFILE 5) 1, 2
    STOP
    IF (TLE(1),EQ, TITLE) GO TO 3
    PRINT 17, TLE(1)
    GO TO 17
    PRINT 19, (TLE(M), M=2, 8)
    DATA 13, VAP, 22, (K(M), M=1, 5)
    IF (VAP,EQ, ENFILE) GO TO 12
    M=1
    M=M+1
    IF (M,LT,24) GO TO 5
    PRINT 27, VAP
    GO TO 17
    IF (VAP-T(M)) 3, 7, 5
    IF (1-14) 3, 8, 11
    IV(M)=K(M)
    GO TO 4
    IV(M-1)=K(1)
    GO TO 4
    PRINT 21
    PRINT 22, ((T:I+9), TV(I)), I=7, 15)
    PRINT 23, ((M, IV(M)), M=1, 5)
    PRINT 24, ((T(M), V(M)), M=1, 17)
    IFP2=1
    ENFILE
    IFP2=1
    DATA (5,16) (TLE(M), M=1, 8)
    IF (ENDFILE 5) 1, 15
    IF (TLE(1),EQ, TITLE) GO TO 3
    GO TO 14
  
```

```

DATA 1
DATA 2
DATA 3
DATA 4
DATA 5
DATA 6
DATA 7
DATA 8
DATA 9
DATA 10
DATA 11
DATA 12
DATA 13
DATA 14
DATA 15
DATA 16
DATA 17
DATA 18
DATA 19
DATA 20
DATA 21
DATA 22
DATA 23
DATA 24
DATA 25
DATA 26
DATA 27
DATA 28
DATA 29
DATA 30
DATA 31
DATA 32
DATA 33
DATA 34
DATA 35
DATA 36
DATA 37
DATA 38
DATA 39
DATA 40
DATA 41
DATA 42
DATA 43
DATA 44
DATA 45
DATA 46
DATA 47
DATA 48
DATA 49
DATA 50
DATA 51
DATA 52
  
```

NOT REPRODUCIBLE

16	FORMAT (A10)		
17	FORMAT (1H1,7X,19HEP00 IN INPUT DATA/7X,21HTITLE CARD NOT SENT	DATA	53
	*/7X,4HDEAD,1X,7H10H,A10,1X,75HOTHER THAN 1HTITLE	DATA	54
	*CASE TERMINATED)	DATA	55
14	FORMAT (//25X,13HCASE TITLE - ,7A10)	DATA	56
18	FORMAT (A10,F10.4,6I10)	DATA	57
20	FORMAT (/70X,19HEP00 IN INPUT DATA/70X,42HAN INCORRECT VARIABLE	DATA	58
	*NAME WAS ENCOUNTERED/70X,12HTHE NAME WAS,1X,A11/70X,15HCASE TERMIN	DATA	59
	*TED)	DATA	60
21	FORMAT (//25X,80H***** INPUT DATA *****	DATA	61
	*****)	DATA	62
22	FORMAT (/15X,?(10X,A10,I10))	DATA	63
23	FORMAT (/15X,?(10X,4HTPO(,I1,1H),4X,I10))	DATA	64
24	FORMAT (/15X,?(10X,A1),F10.3))	DATA	65
25	FORMAT (1H1////////)	DATA	66
	END	DATA	67
		DATA	68
		DATA	69

7. DETERMINANT SUBROUTINES

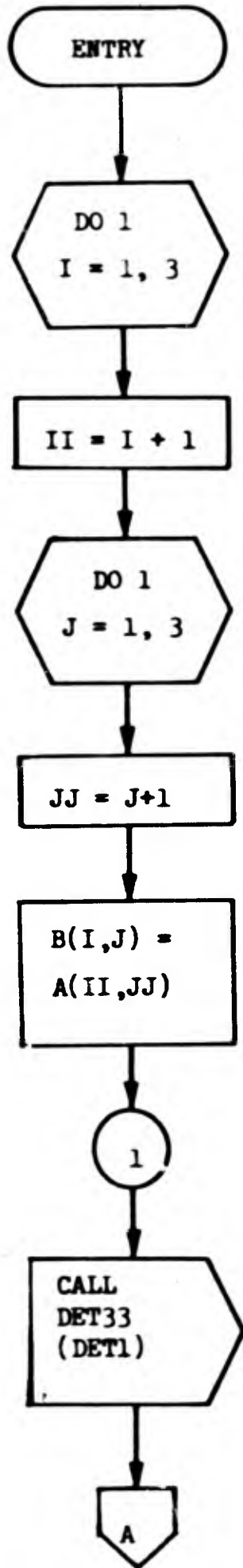
A. PURPOSE

SUBROUTINE DET44 computes the determinant of a 4 X 4 array. It accomplishes a standard expansion by minors using the explicit determinant solution for a 3 X 3 array contained in SUBROUTINE DET33.

B. VARIABLE LIST

- A - 4 X 4 array for which the determinant is to be evaluated.
- B - 3 X 3 array containing various minors of the A array.
- DET - value of the determinant evaluated in both routines.
- DET1 - determinant of the first minor.
- DET2 - determinant of the second minor.
- DET3 - determinant of the third minor.
- DET4 - determinant of the fourth minor.
- I - index used to load arrays.
- II - index used to load arrays.
- J - index used to load arrays.
- JJ - index used to load arrays.

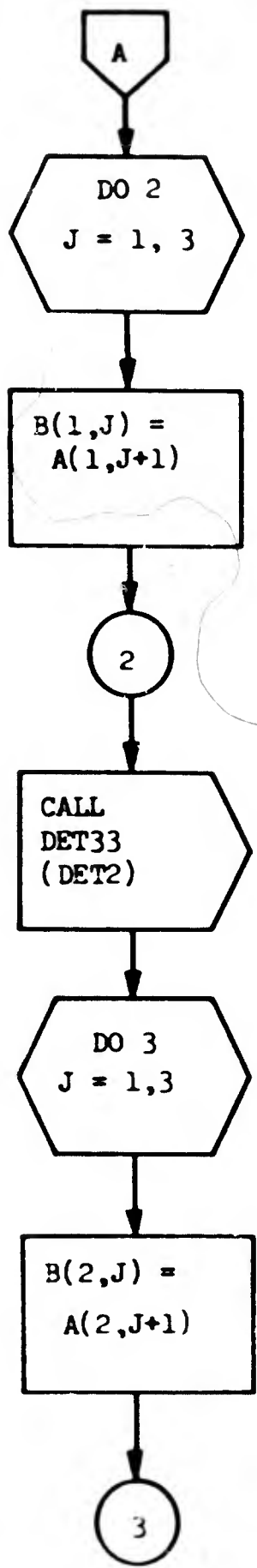
C. FLOW CHART



ENTER SUBROUTINE DET44.

LOAD 3X3 ARRAY FOR THE
FIRST MINOR.

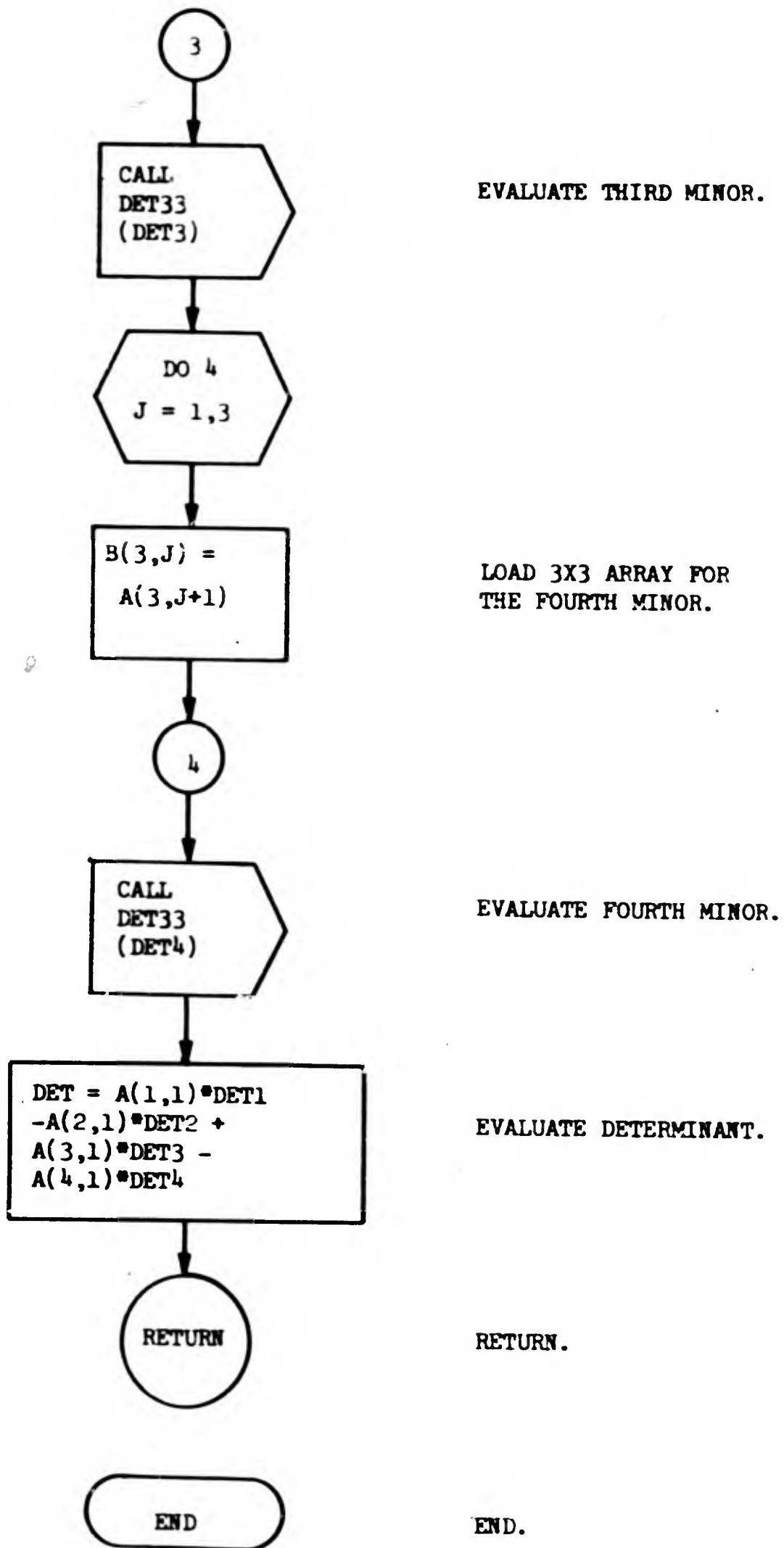
EVALUATE FIRST MINOR.

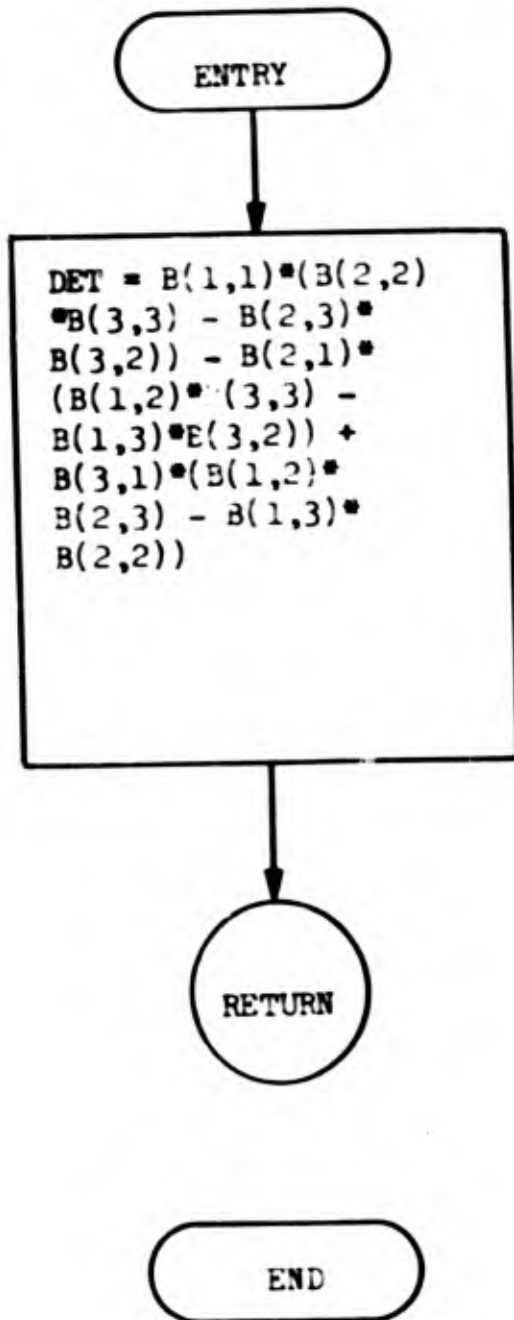


LOAD 3X3 ARRAY FOR
SECOND MINOR.

EVALUATE SECOND MINOR.

LOAD 3X3 ARRAY FOR
THIRD MINOR.





ENTER SUBROUTINE DET33.

EXPLICIT SOLUTION FOR
A 3X3 DETERMINANT.

RETURN.

END.

D. LISTING

```

C      SUBROUTINE DET44 (A, DFT)
C
C      THIS ROUTINE SOLVES A 4X4 DETERMINANT.
C
COMMON /DETERM/ B(3,3)
DIMENSION A(4,4)
DO 1 I=1,3
  II=I+1
  DO 1 J=1,3
    JJ=J+1
    R(I,J)=A(II,JJ)
    CALL DET33 (DET1)
  DO 2 J=1,3
    R(1,J)=A(1,J+1)
    CALL DET33 (DET2)
  DO 3 J=1,3
    R(2,J)=A(2,J+1)
    CALL DET33 (DET3)
  DO 4 J=1,3
    R(3,J)=A(3,J+1)
    CALL DET33 (DET4)
  DET=A(1,1)*DFT1-A(2,1)*DET2+A(3,1)*DET3-A(4,1)*DET4
  RETURN
END

```

DET44 1
 DET44 2
 DET44 3
 DET44 4
 DET44 5
 DET44 6
 DET44 7
 DET44 8
 DET44 9
 DET44 10
 DET44 11
 DET44 12
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 DET44 14
 DET44 15
 DET44 16
 DET44 17
 DET44 18
 DET44 19
 DET44 20
 DET44 21
 DET44 22
 DET44 23
 DET44 24

```

C      SUBROUTINE DET33 (DET)
C
C      THIS ROUTINE SOLVES A 3X3 DETERMINANT.
C
COMMON /DETERM/ B(3,3)
DET=B(1,1)*(B(2,2)*B(3,3)-B(2,3)*B(3,2))-B(1,2)*(B(1,3)*B(2,1)-B(1,1)*B(2,3))+B(1,3)*(B(1,2)*B(2,1)-B(1,1)*B(2,2))
RETURN
END

```

DET33 1
 DET33 2
 DET33 3
 DET33 4
 DET33 5
 DET33 6
 DET33 7
 DET33 8
 DET33 9

8. SUBROUTINE DMPDOUT

A. PURPOSE

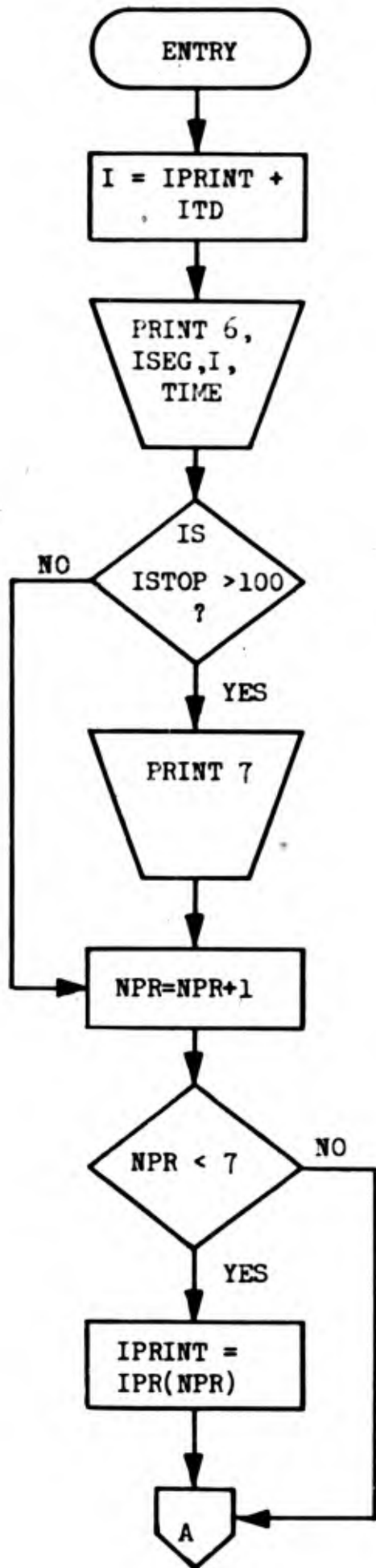
This subroutine prints the flow-field data for each segment. Printed outputs are provided as requested by the user through the input variable IPR, and after the last iteration for each segment. The output is printed in dimensional form (except for entropy) based on the values of P1, T1, RH01 and RN supplied in the input data. When all of these variables are unity, nondimensional output is provided. The normal velocity component, the coordinate A and the shock velocity are defined positive outward from the body in the printed output. The printed output is set up to include as many streamwise stations per page as possible, based on a maximum of 60 lines per page.

B. VARIABLE LIST

All variables previously defined in the description of the main program have been omitted from this list.

- DAT - array containing the values of various body parameters for each station.
- DS - ΔS
- IND - control variable specifying when a value of DAT should be included on a given printed line.
- IPAGE - number of streamwise stations which have been printed on the present page.
- M - index used in print statements.
- NPAGE - number of streamwise stations which can be printed on each page.
- NST - array used to store the pressure coefficient at each body station for later output.
- PREFIX - array containing the alphanumeric names of body parameters.
- Q - $[P_{\infty} \rho_{\infty}]^{1/2}$
- QSQ - $u^2 + v^2$
- SI - array used to store the values of the ratio of local pressure to total pressure on the body for later output.
- VS - array used to store dimensional values of output values as they are evaluated for later output.
- ZZ - dummy variable used for intermediate computations.

C. FLOW CHART



ENTER SUBROUTINE DMPOUT.

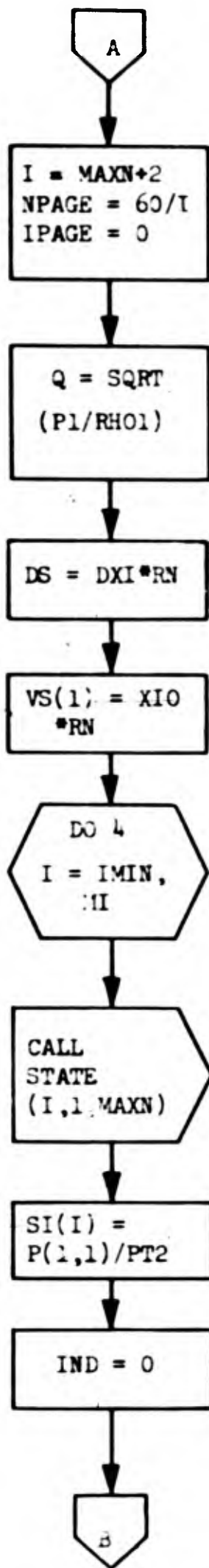
COMPUTE TOTAL NUMBER OF ITERATIONS PERFORMED ON THIS SEGMENT.

WAS THE MAXIMUM NUMBER OF NODES ACROSS THE SHOCK LAYER EXCEEDED IN SUBROUTINE GEOM ?

INCREMENT NPR.

ARE ANY VALUES OF IPR UNUSED ?

LOAD ITERATION NUMBER FOR NEXT PRINT.



COMPUTE THE NUMBER OF STATIONS WHICH CAN BE PRINTED ON EACH PAGE. INITIALIZE IPAGE.

COMPUTE SCALE FACTOR FOR DIMENSIONAL VELOCITIES.

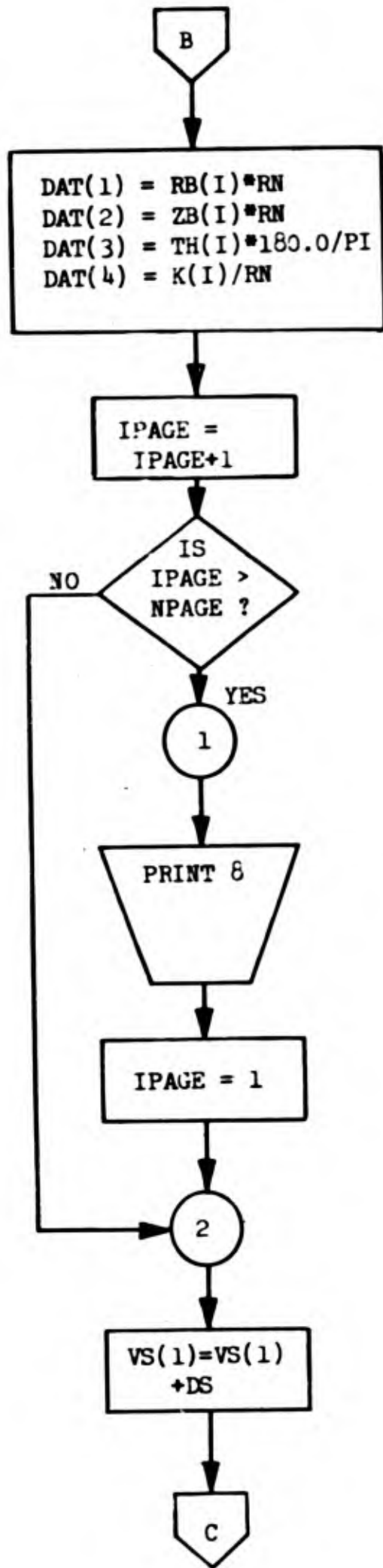
COMPUTE DIMENSIONAL ΔS.

INITIALIZE S.

COMPUTE PRESSURES.

STORE P/P₀.

INITIALIZE IND.



COMPUTE LOCAL BODY
PARAMETERS.

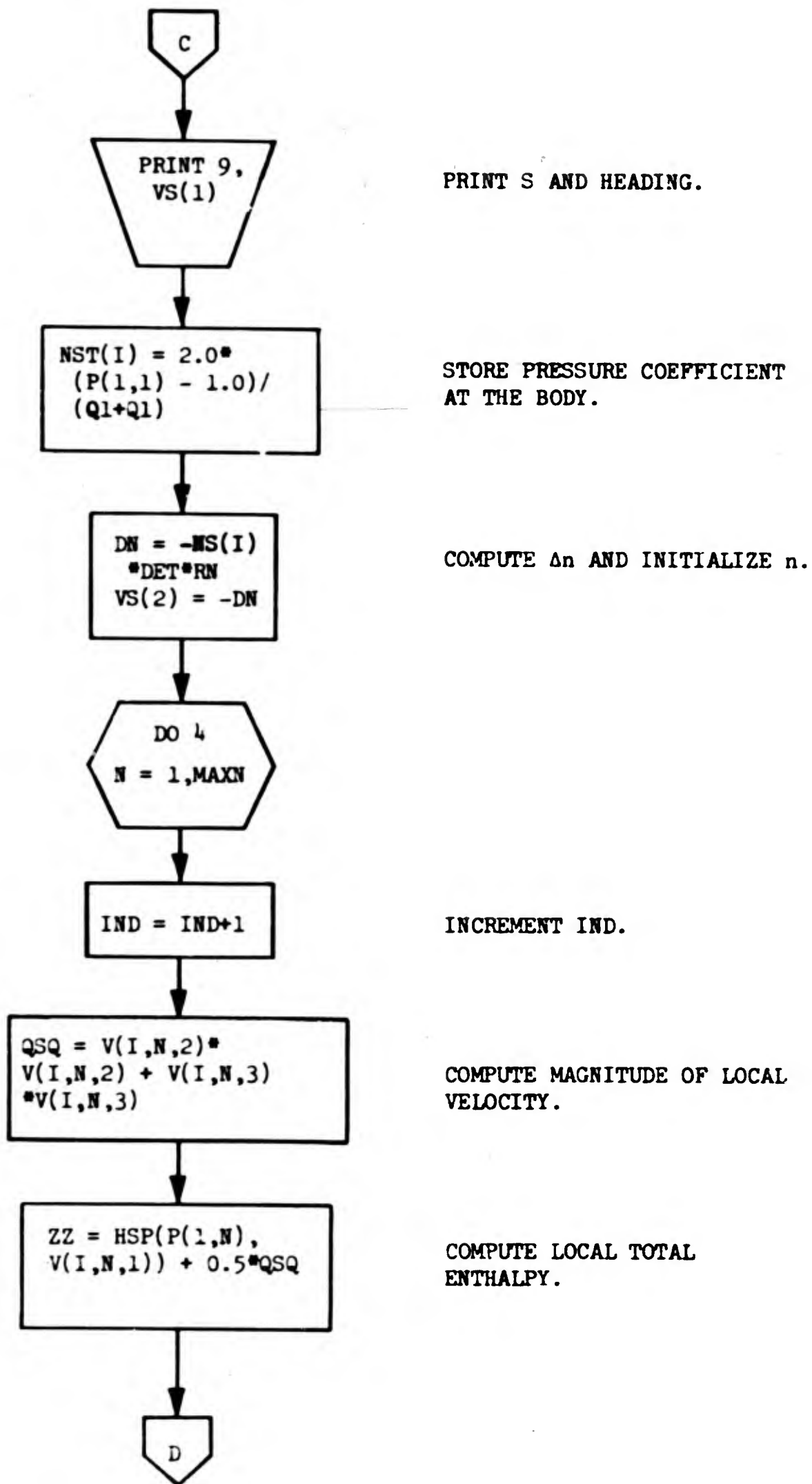
INCREMENT IPAGE.

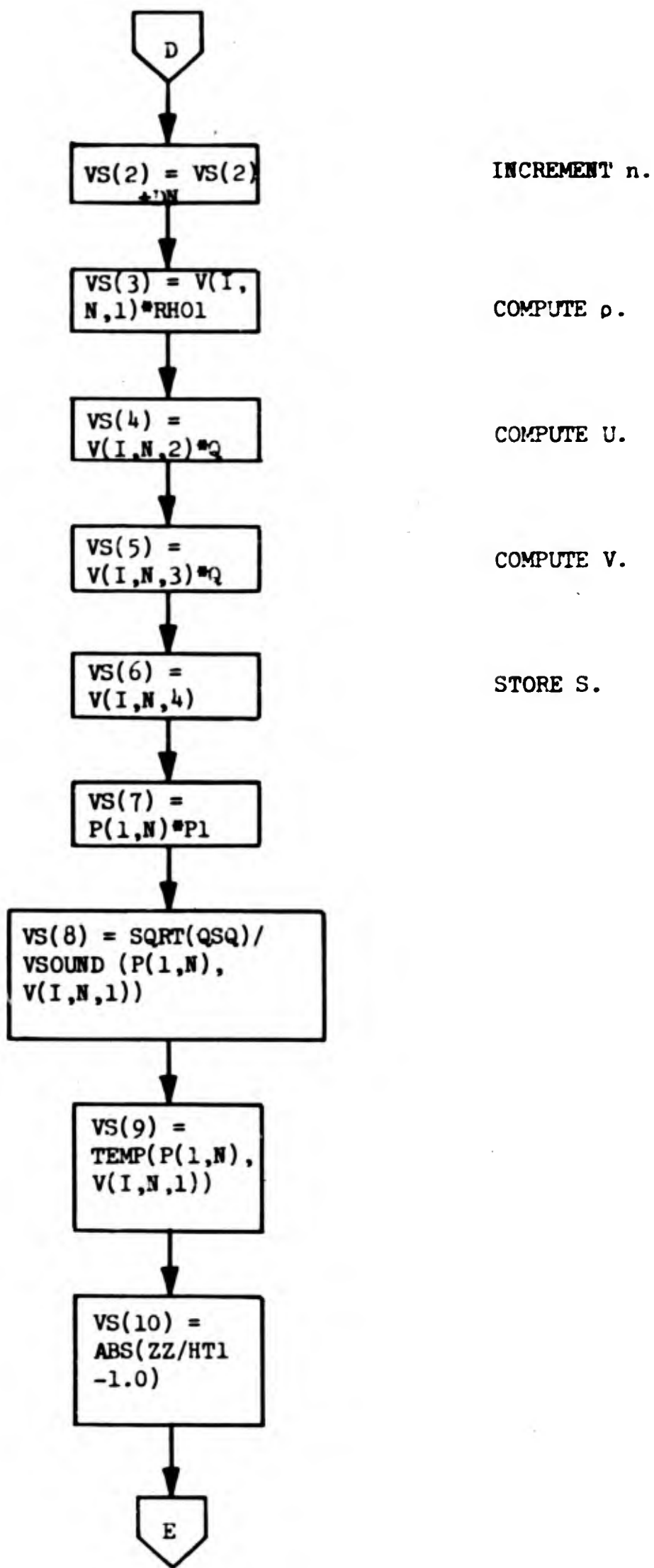
START A NEW PAGE ?

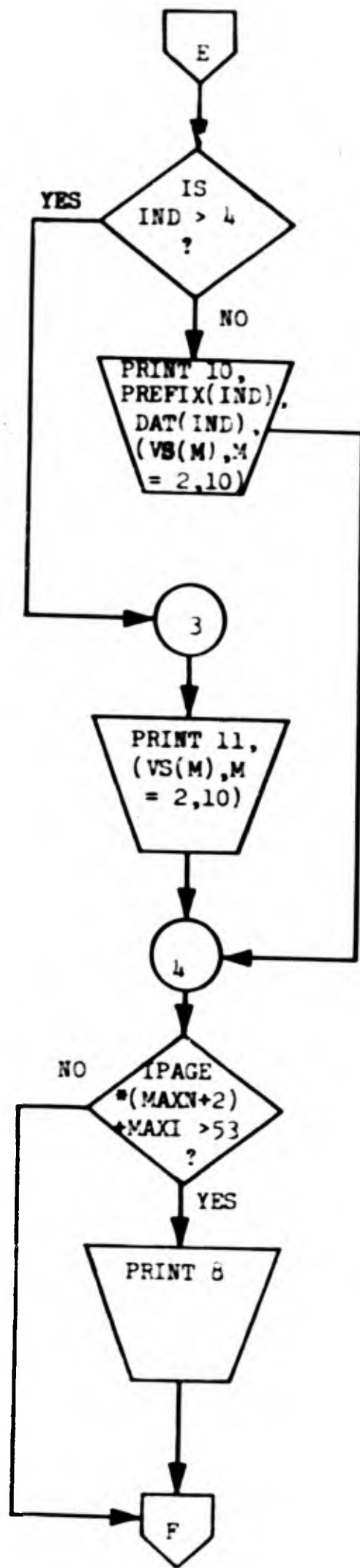
NEW PAGE.

INITIALIZE IPAGE.

INCREMENT S.



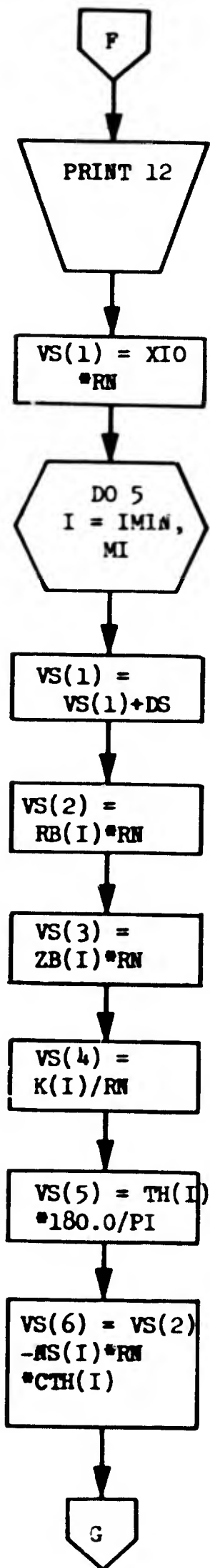




PRINT BODY PARAMETERS
AND FLOW FIELD DATA.

PRINT FLOW FIELD DATA.

NEW PAGE FOR SHOCK AND
BODY PARAMETERS ?



PRINT HEADING FOR SHOCK
AND BODY PARAMETERS.

INITIALIZE S.

INCREMENT S.

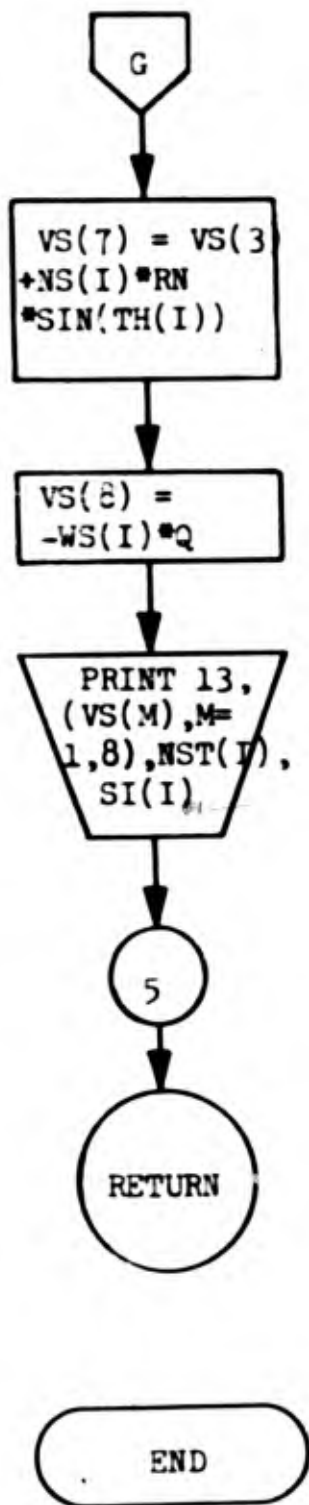
COMPUTE γ AT THE BODY.

COMPUTE Z AT THE BODY.

COMPUTE K AT THE BODY.

COMPUTE θ .

COMPUTE γ AT THE SHOCK.



COMPUTE Z AT THE SHOCK.

COMPUTE SHOCK VELOCITY.

PRINT SHOCK AND BODY PARAMETERS.

RETURN.

END OF SUBROUTINE DMPOUT.

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52

```

THIS ROUTINE PRINTS THE FLOW FIELD DATA.

COMMON /1/ V(15,4),U(4),P(3,4)

COMMON /2/ PR(15),ZR(15),TH(15),CYH(15),K(15),NS(15),WS(15)

COMMON /3/ MAXI,MI,IMIN,MAXN,MN,NMIN

COMMON /4/ IPRINT,NPR,KGEOM

COMMON /5/ OX1,DET,OX2,OF2,OT

COMMON /6/ O1,VS1,HT1,A1,B1

COMMON /9/ ISEG,XI2,DEL5

COMMON /10/ VS(10),DS,DN,C,CSC,ZZ

COMMON /11/ MAXIT,ISTOP,ITD,TIME,SI(15),NST(15)

COMMON /INPUT/ GAM1,AMACH1,THETA,ZRN,ICPT(2),P1,RHC1,RN,AMU0,T1,PC(2),IPR(6),IATAP,OTAP,NCPT,NCASE,NFILE,NSEG,ICEOM,NCFOM,NTAP

COMMON /CONST/ PI,PI2,GAMA,CAMB

COMMON /PRES/ PT2,PG

DIMENSION P(4), DAT(4)

DATA (P(I),I=1,4)/5H 00=,6H 70=,8HTHETA=,9H <=

DEL NS,K,NST

I=IPRINT+1

PRINT 6, ISEG,I,TIME

IF (ISTOP.GT.10) PRINT 7

NDD=NDD+1

IF (MOD(I,7) .EQ. 0) IPRINT=IPR(NPR)

I=MAXN+2

NPAGE=61/I

IPAGE=0

C=SQRT(P1/RHC1)

DS=OXT*DN

VS(1)=XIT*RN

DO 4 I=IMIN,MN

CALL STATE (I,1,MAXN)

SI(I)=P(1,1)/PI2

ITD=0

DAT(1)=P(1,1)*CA

DAT(2)=P(2,1)*CA

DAT(3)=TH(I)*180./PI

DAT(4)=K(I)/CA

IPAGE=IPAGE+1

IF (IPAGE=NPAGE) 2,2,1

PRINT 8

IPAGE=1

VS(1)=VS(1)+DS

PRINT 9, VS(1)

NST(I)=2.7*(P(1,1)-1.0)/(O1*O1)

DN=-NS(I)*DET*DN

VS(2)=-DN

DO 4 N=1,MAXN

ITD=ITD+1

DSQ=V(I,N,2)*V(I,N,2)+V(I,N,3)*V(I,N,3)

ZZ=HSP(P(1,N),V(I,N,1))+3.5*CSC

	VS(2)=VS(2)+CA	00000000 57
	VS(3)=V(I,N,1)*PH01	00000000 58
	VS(4)=V(I,N,2)*P2	00000000 59
	VS(5)=-V(I,N,3)*P3	00000000 60
	VS(6)=V(I,N,4)	00000000 61
	VS(7)=P(I,N)*P4	00000000 62
	VS(8)=SQRT(0.50)/VSOUND(P(I,N),V(I,N,1))	00000000 63
	VS(9)=TEMP(VS(7),VS(8))	00000000 64
	VS(10)=135(ZZ/HT1-1.1)	00000000 65
	IF (IND.GT.4) GO TO 3	00000000 66
	PRINT 10, PREFIX(IND), DAT(IND), (VS(M),M=2,10)	00000000 67
	GO TO 4	00000000 68
7	PRINT 11, (VS(M),M=1,10)	00000000 69
	CONTINUE	00000000 70
	IF (IPACT*(M3VX+2)+M3VT.GT.53) PRINT *	00000000 71
	PRINT 12	00000000 72
	VS(1)=X(I)*P2V	00000000 73
	DO 5 I=1,N,M1	00000000 74
	VS(1)=VS(1)+DS	00000000 75
	VS(2)=P(I)*P2A	00000000 76
	VS(3)=P(I)*P2B	00000000 77
	VS(4)=K(I)/PA	00000000 78
	VS(5)=TH(I)*P2C/P2	00000000 79
	VS(6)=VS(2)-VS(1)*P2A*STH(I)	00000000 80
	VS(7)=VS(3)+VS(1)*P2B*STH(I)	00000000 81
	VS(8)=-VS(1)*P2C	00000000 82
7	PRINT 13, (VS(M),M=1,8),NST(I),ST(I)	00000000 83
	CONTINUE	00000000 84
2		00000000 85
6	FORMAT (1H1,20X,16HSEGMENT NUMBER,13,10X,16HITERATION NUMBER,13,17X,	00000000 86
	*Y,8HTIME=,2I2,4)	00000000 87
7	FORMAT (31X,52HEXCEEDED MAXIMUM NUMBER OF NODES IN NORMAL DIRECTION	00000000 88
	*M/70X,27HMAXIMUM NUMBER WAS USED)	00000000 89
8	FORMAT (1H1)	00000000 90
9	FORMAT (74X,2F5=,F3,2,7X,1HN,11X,3H2FC,9X,1HD,11X,1HV,3X,3H3C-51)/	00000000 91
	*OV,6Y,1HD,11X,1HM,11X,1HT,9X,FHDH/M1)	00000000 92
10	FORMAT (1X,4E,97.2,1X,4E12.4,F8.5)	00000000 93
11	FORMAT (15X,4E12.4,F3.5)	00000000 94
12	FORMAT (777.7X,25HSHOCK AND BODY PARAMETERS//4X,1FS,11X,2HCP,11X,	00000000 95
	*RHTB,10X,1HK,9X,5HHTETA,7X,2HFS,10X,2H2S,11X,2HWS,11X,2HCP,9X,	00000000 96
	*LHD/31)	00000000 97
17	FORMAT (10F12.4)	00000000 98
	END	00000000 99

NOT REPRODUCIBLE

9. SUBROUTINE GEOM

A. PURPOSE

This subroutine computes the body geometry parameters and the initial guess for the shock shape. An option for generating these parameters for a sphere-cone vehicle is available. Currently, other body shapes must be specified with data cards. These geometry data cards are read in this routine. A computed GO TO statement selects the proper option based on the value of a variable KGEOM. Explicit computations for body shapes other than sphere-cones could be incorporated by using additional values of this variable (currently, values of 1 to 3 are used). When the body geometry is read from data cards, several options are available. The data can be read in at specified equal intervals along the body or at a series of representative points. In the latter case, curve fit routines accomplish the necessary integrations and interpolations to specify the data at equal intervals along the body. The initial shock shape can be input or computed using empirical relations. The number of nodes to be used along the body and across the shock layer can be specified for each segment. These specifications will be ignored if these quantities were defined in the input read in SUBROUTINE DATA. The body curvature can be changed between segments at the common segment boundary. If no specification for the number of nodes to be used along the body is provided, the number of data cards provided is used as the value for this quantity. An error check is accomplished to ensure that the number of points provided for a segment has not exceeded the dimension sizes. The segment number for which a current set of data cards has been provided is checked with the number of the segment under consideration. Data are skipped until the correct segment number is located. This allows the complete geometry data packet to be included in a restart, even though the first few segments will not be required. If a segment number greater than the value of the segment under consideration is used, it is considered as a data error. All data errors result in an error message and termination of the current case. An subsequent case will then be treated.

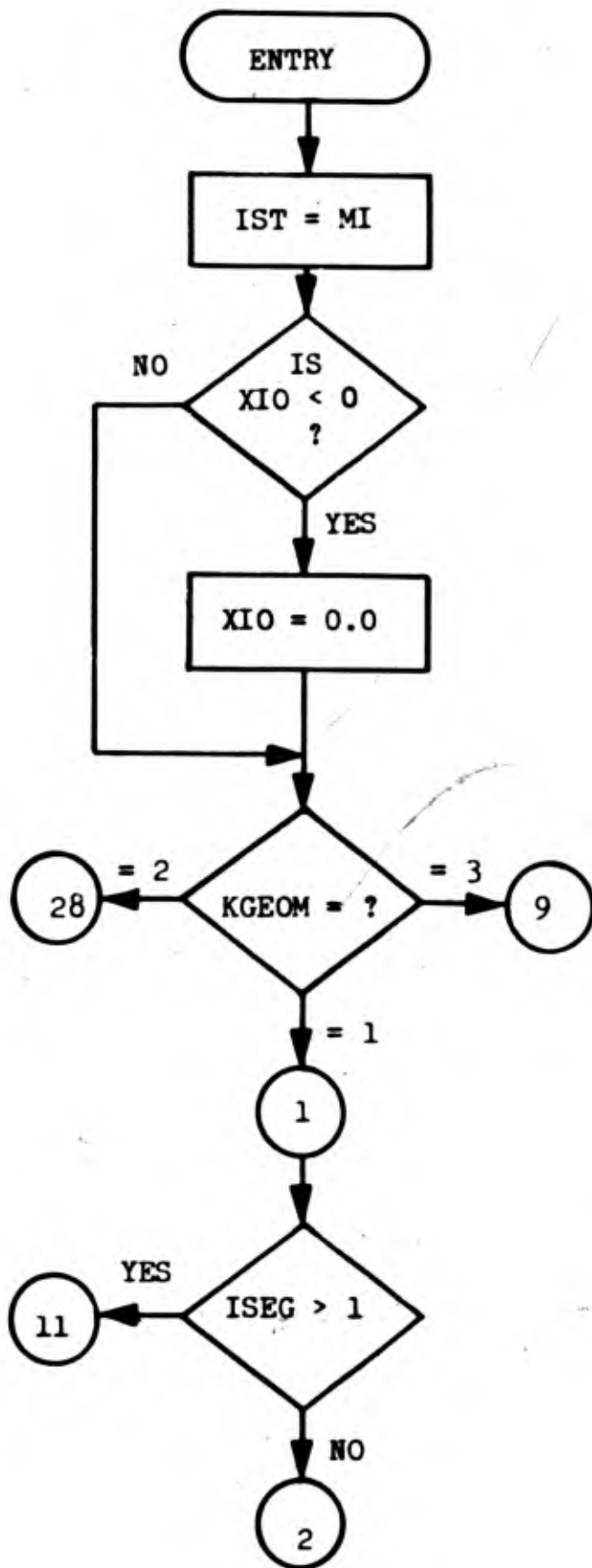
B. VARIABLE LIST

All variables previously defined in the description of the main program have been omitted from this list.

- A - array used for body shape computations under the arbitrary body shape option.
- ANS - variable used for temporary storage of the value of n_s at the end of a segment previously completed.

- BLANK** - alphanumeric blank under A10 format. Used to check for discontinuities in body curvature.
- C** - array used to store cubic fit coefficients for each data point under the arbitrary body shape option.
- DNS** - change in n_s between body stations.
- DR** - change in r at the body between body stations for a sphere-cone.
- DUM** - dummy variable used for intermediate computations.
- DZ** - change in Z between body stations for a sphere-cone vehicle.
- IC** - body station index at which integration with respect to r should be replaced with integration with respect to Z .
- IP** - dummy index used for various applications.
- IST** - index of the downstream boundary body station for the previous segment.
- KSH** - control variable specifying whether the initial shock shape has been specified on the data cards.
- M** - dummy index used for various applications.
- NPTS** - number of data cards to be read for the current segment.
- NST** - array used for temporary storage of an initial shock shape specification as provided on the data cards.
- SI** - array used for body shape computations under the arbitrary body shape option.
- SUM** - result of a numerical integration.
- ZZ** - dummy variable used for intermediate computations.

C. FLOW CHART



ENTER SUBROUTINE GEOM.

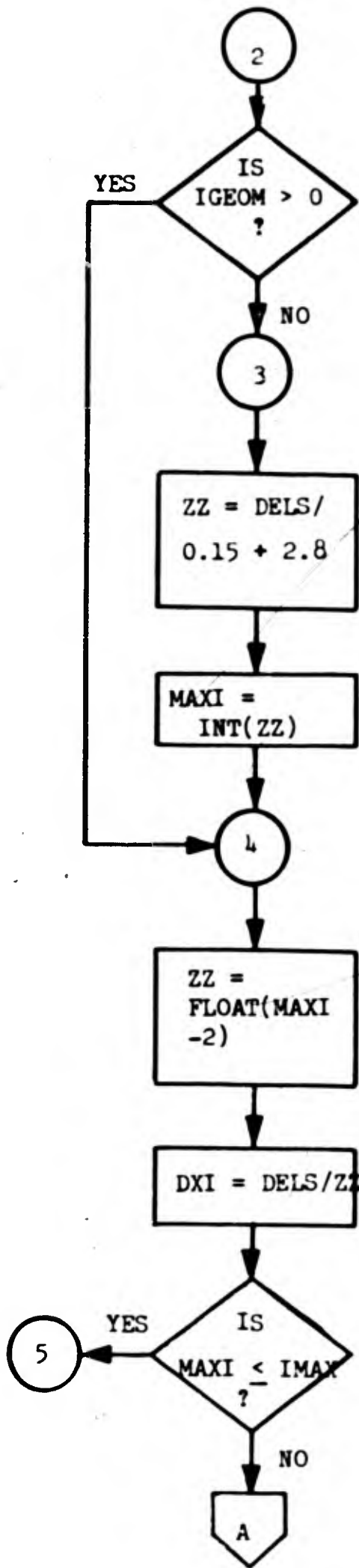
SAVE INDEX OF THE DOWNSTREAM BOUNDARY OF THE PREVIOUS SEGMENT.

IS THE VALUE OF ξ AT UPSTREAM BOUNDARY OF THE LAST SEGMENT NEGATIVE ?

WHAT GEOMETRY OPTION ?

SPHERE-CONE; SPHERE TREATED WITH A SINGLE SEGMENT.

HAS THE FIRST SEGMENT BEEN COMPLETED ?

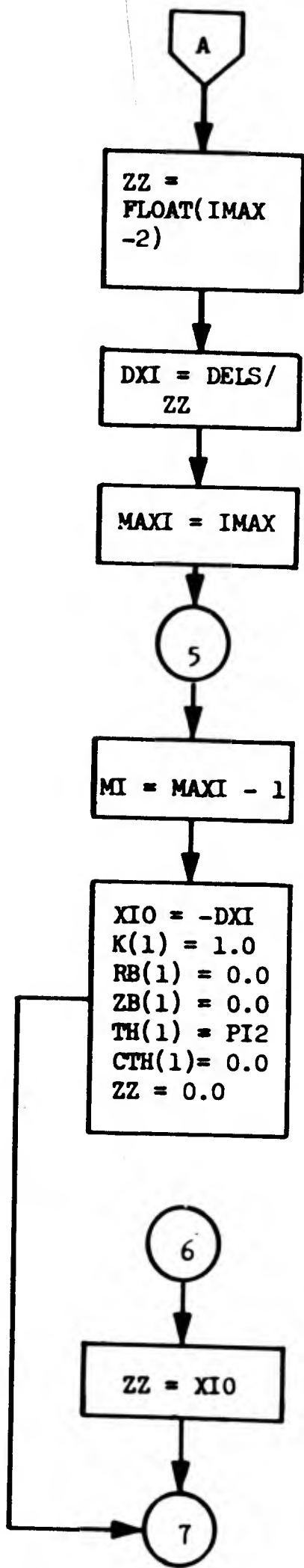


WAS THE NUMBER OF NODES ALONG THE BODY PER SEGMENT SPECIFIED IN THE CASE DATA ?

COMPUTE MAXI.

COMPUTE $\Delta\xi$.

ARE DIMENSIONED VARIABLES LARGE ENOUGH FOR THE VALUE OF MAXI ?



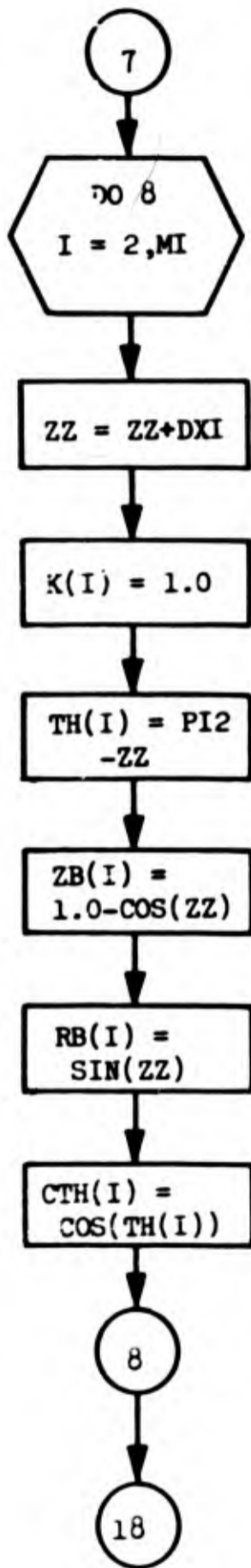
COMPUTE MINIMUM ALLOWABLE
VALUE OF $\Delta\xi$.

SET MAXI TO MAXIMUM VALUE.

COMPUTE INDEX OF DOWNSTREAM
BOUNDARY NODE.

INITIALIZE BODY PARAMETERS
FOR STAGNATION POINT.

INITIALIZE ξ FOR SEGMENTS
OTHER THAN THE FIRST.



START DO LOOP TO COMPUTE BODY
GEOMETRY FOR A SPHERE.

INCREMENT ξ .

CURVATURE.

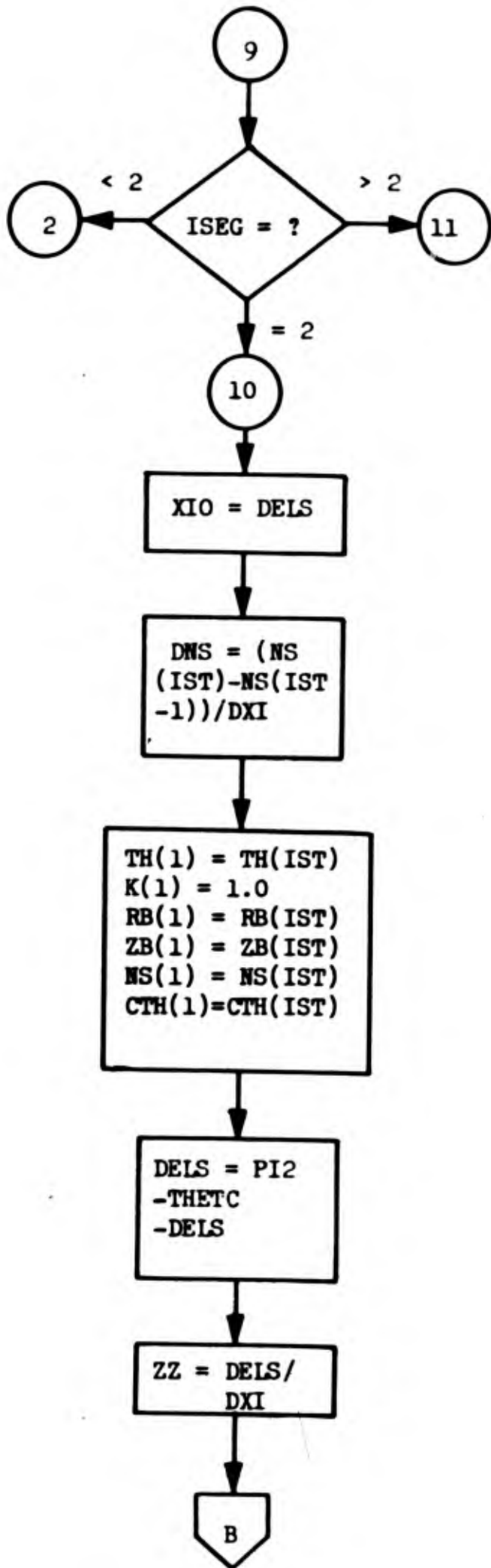
θ .

Z_b .

r_b .

$\cos(\theta)$.

END OF LOOP.



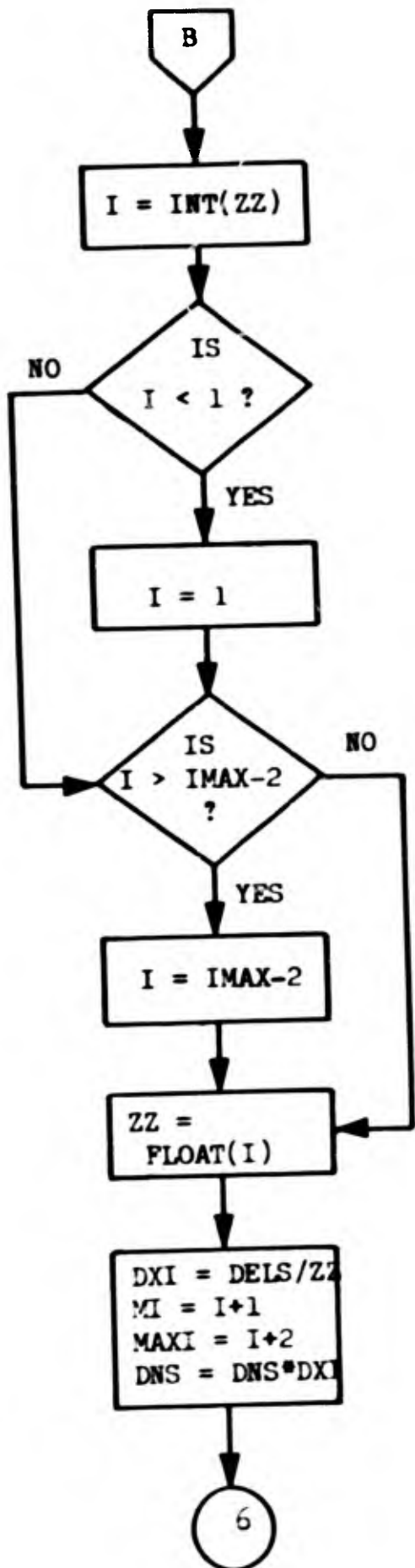
SPHERE-CONE VEHICLE WITH
SPHERE TREATED WITH 2
SEGMENTS.

INITIALIZE XIO FOR
SEGMENT 2.

SAVE SHOCK SLOPE AS COMPUTED
AT THE DOWNSTREAM BOUNDARY OF
SEGMENT 1.

INITIALIZE BODY GEOMETRY
PARAMETERS FOR UPSTREAM
BOUNDARY OF SEGMENT 2.

COMPUTE SEGMENT STREAMWISE
LENGTH.

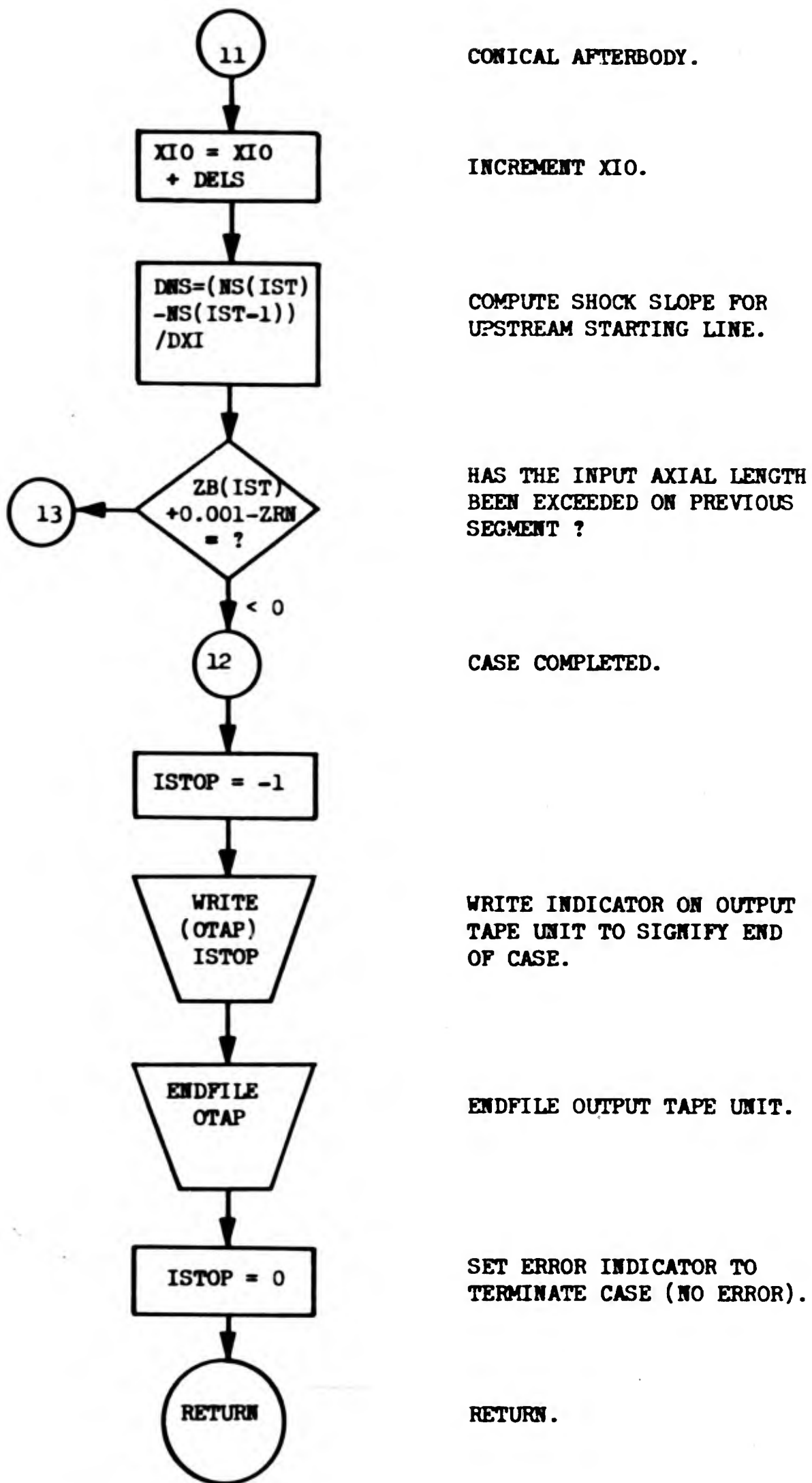


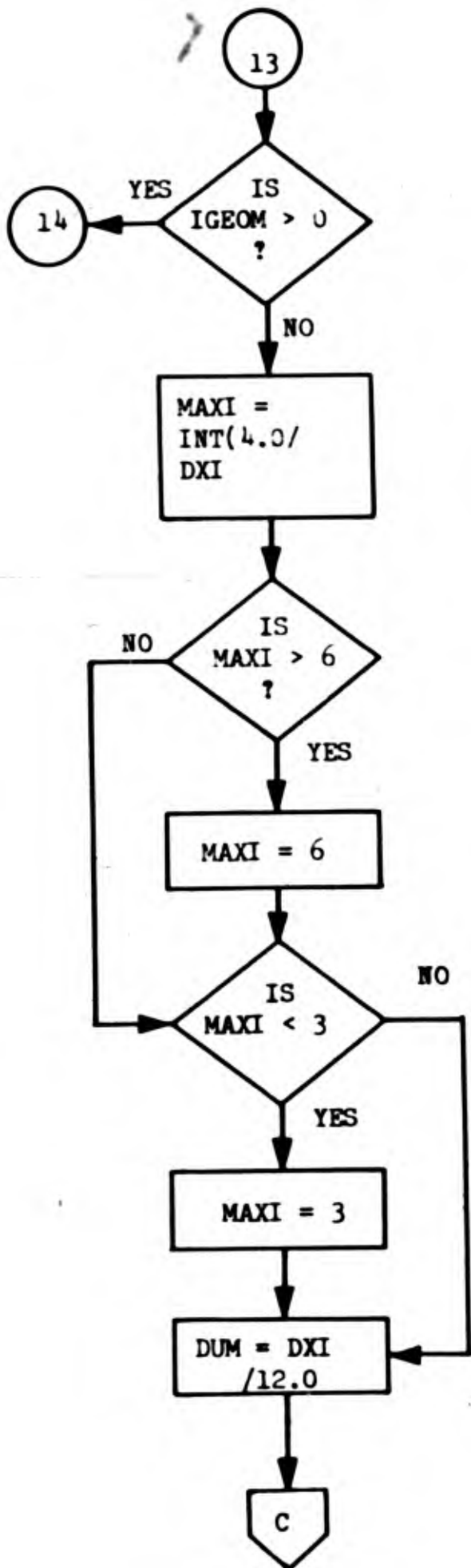
COMPUTE THE NUMBER OF NODES
REQUIRED TO GIVE THE SAME
VALUE OF $\Delta\xi$ AS USED FOR
SEGMENT 1.

EXCEEDED MAXIMUM NUMBER OF
NODES ALLOWED ?

USE MAXIMUM NUMBER.

COMPUTE $\Delta\xi$, MI, MAXI AND THE
CHANGE IN NS PER EACH $\Delta\xi$
BASED ON LINEAR EXTRAPOLATION.





CASE INCOMPLETE.

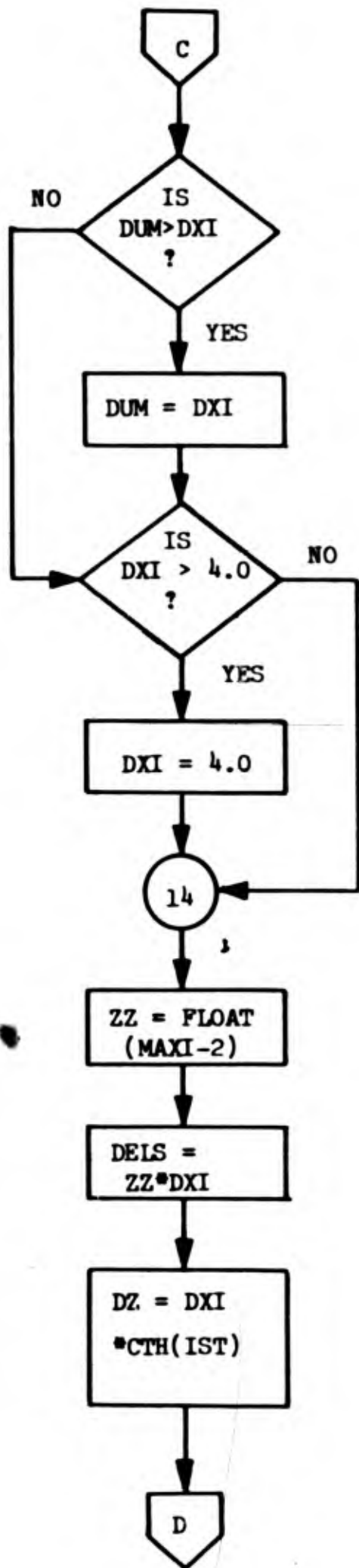
WAS MAXI SPECIFIED IN THE CASE DATA ?

ESTIMATE THE OPTIMUM NUMBER OF NODES FOR THIS SEGMENT.

USE MAXIMUM OF 6 NODES.

MUST HAVE AT LEAST 3 NODES.

ESTIMATE OPTIMUM VALUE OF $\Delta\xi$.

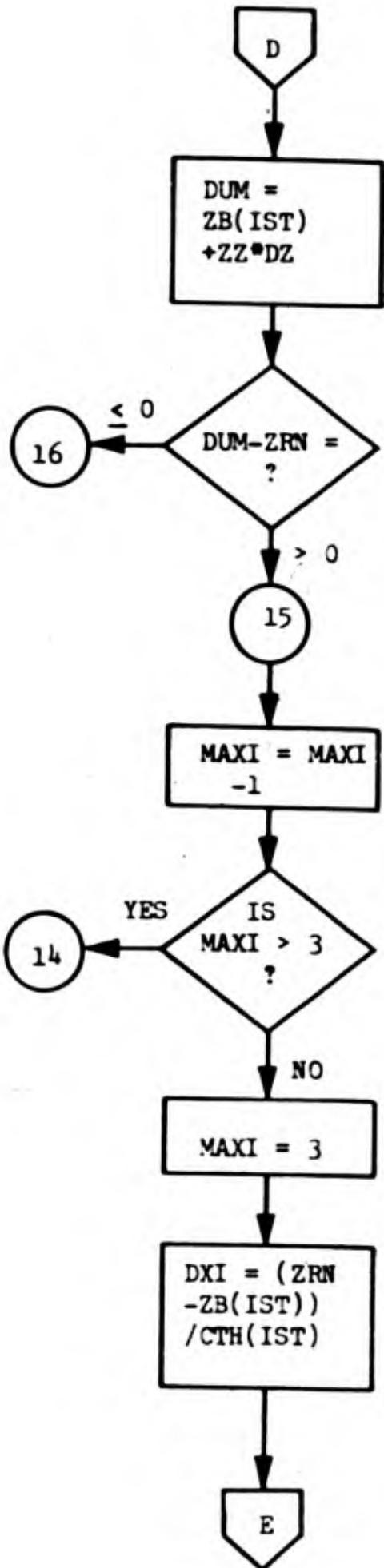


$\Delta\xi \geq$ PREVIOUS VALUE

$\Delta\xi \leq 4.0.$

COMPUTE SEGMENT LENGTH.

$\Delta Z = \Delta\xi \cos \theta.$



COMPUTE Z/RN AT THE END OF THE CURRENT SEGMENT.

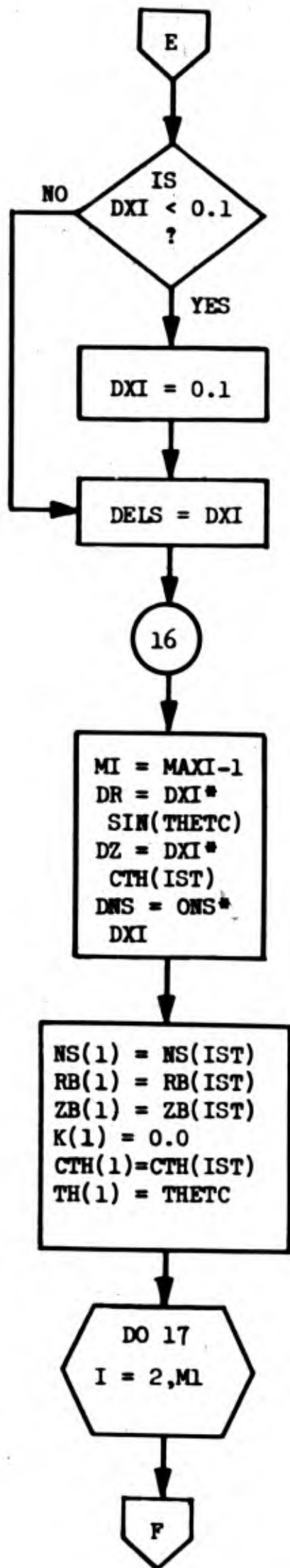
EXCEEDED SPECIFIED VEHICLE LENGTH ?

REDUCE MAXI.

ALLOWABLE VALUE ?

USE 3 NODES TO FINISH CASE.

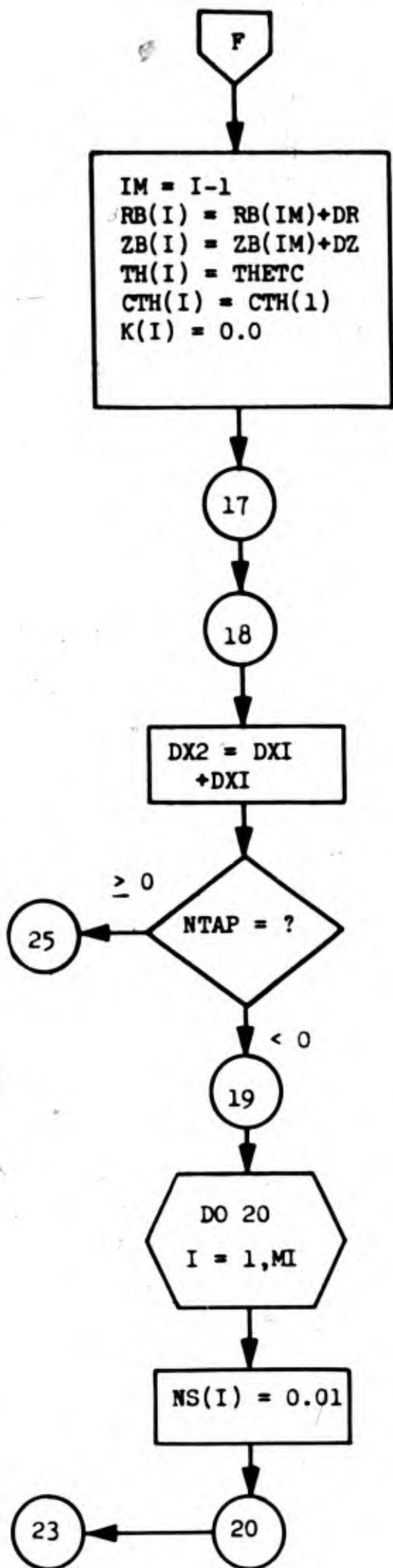
COMPUTE $\Delta \xi$ REQUIRED TO COMPLETE CASE.



REQUIRE $\Delta\xi \geq 0.1$.

COMPUTE INDEX OF NODE AT
DOWNSTREAM BOUNDARY; Δr , ΔZ
AND Δn , BETWEEN NODES.

ASSIGN BODY GEOMETRY
PARAMETERS AT UPSTREAM
STARTING LINE.



COMPUTE BODY GEOMETRY
PARAMETERS AT REMAINING
POINTS.

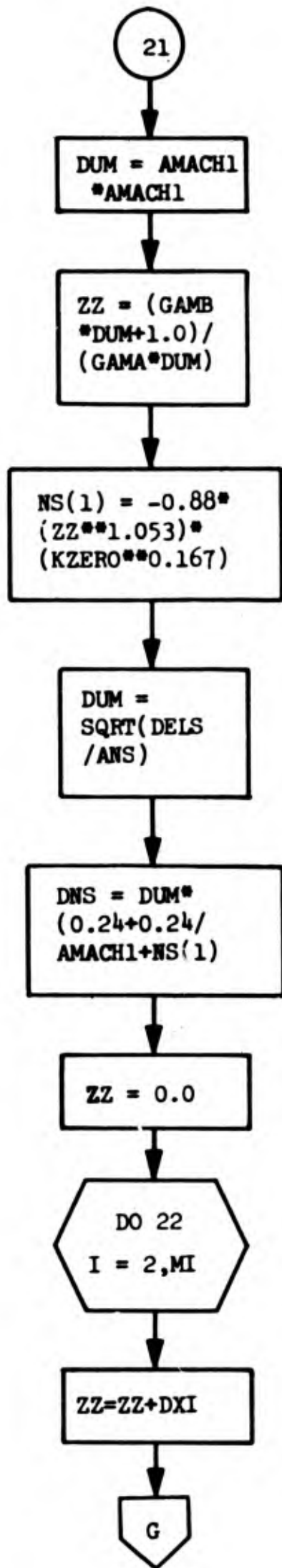
START SHOCK GEOMETRY
COMPUTATIONS (ALL OPTIONS).

COMPUTE $2\Delta\xi$.

UNSTEADY FLOW OPTION.

SHOCK SPECIFIED AS A
CONSTANT DISTANCE FROM
THE BODY OF 0.01.

ESTIMATE SHOCK SHAPE FOR FIRST
SEGMENT, STEADY STATE OPTION.

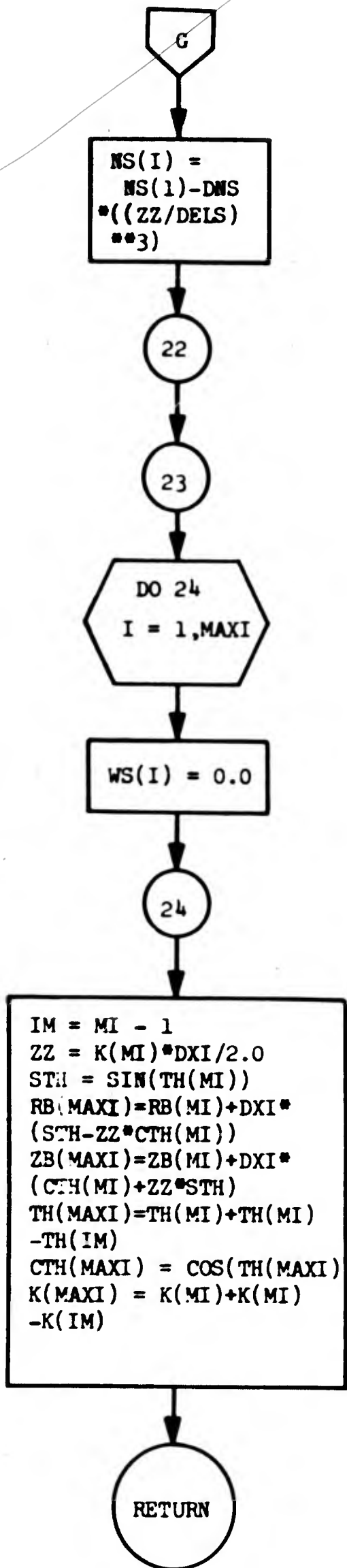


ESTIMATE SHOCK STANDOFF
DISTANCE.

ESTIMATE Δn_s BETWEEN NODES.

INITIALIZE ξ .

INCREMENT ξ .



ESTIMATE SHOCK SHAPE
FOR SEGMENT 1.

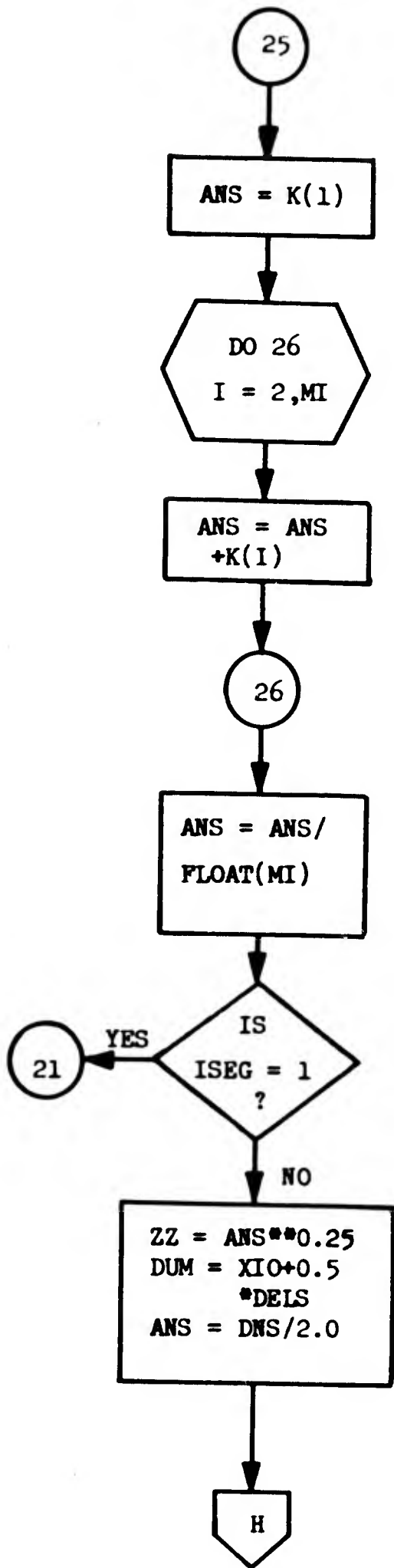
FINAL SERIES OF
COMPUTATIONS FOR ALL OPTIONS.

SET SHOCK VELOCITIES
TO ZERO.

EXTEND BODY GEOMETRY PARAMETERS
TO ADDITIONAL DOWNSTREAM NODE.

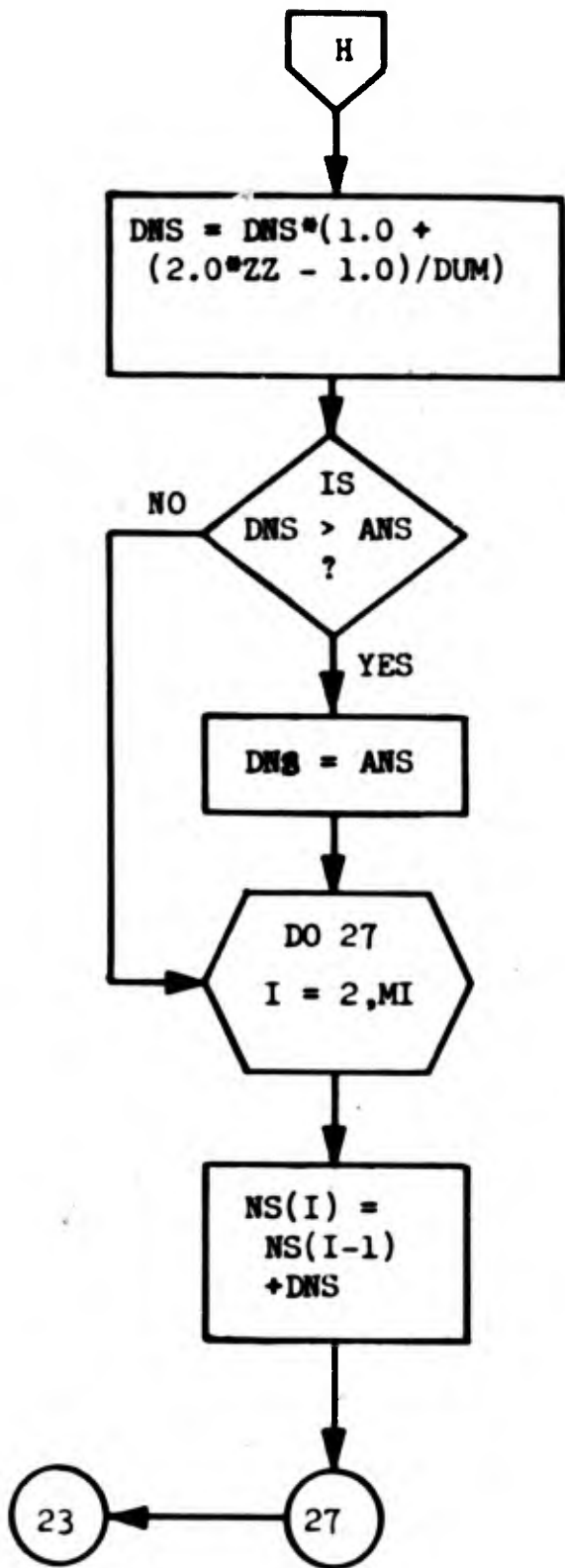
RETURN.

ENTRY POINT FOR SHOCK COMPUTATIONS
WHEN STEADY STATE INITIAL SHOCK
SHAPE IS TO BE ESTIMATED.



COMPUTE AVERAGE BODY CURVATURE
OF THE SEGMENT.

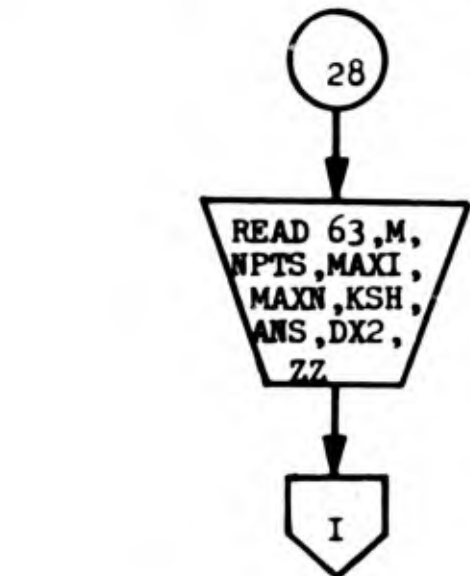
FIRST SEGMENT ?



ESTIMATE Δn_s BETWEEN NODES.

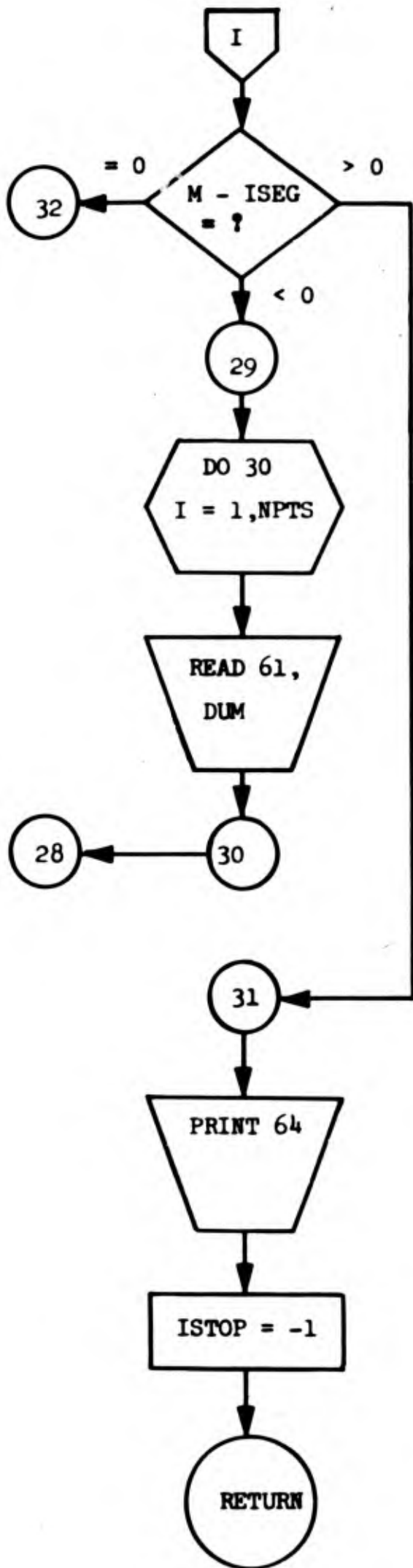
IS ESTIMATED VALUE SMALLER IN MAGNITUDE THAN HALF OF THE VALUE PREDICTED BY THE SHOCK SLOPE ?

ESTIMATE SHOCK SHAPE FOR OTHER POINTS.



ARBITRARY BODY SHAPE OPTION.

READ FIRST DATA CARD FOR THIS SEGMENT.



CORRECT SEGMENT NUMBER ?

SEGMENT NUMBER READ, M,
WAS LESS THAN THE CURRENT
SEGMENT NUMBER.

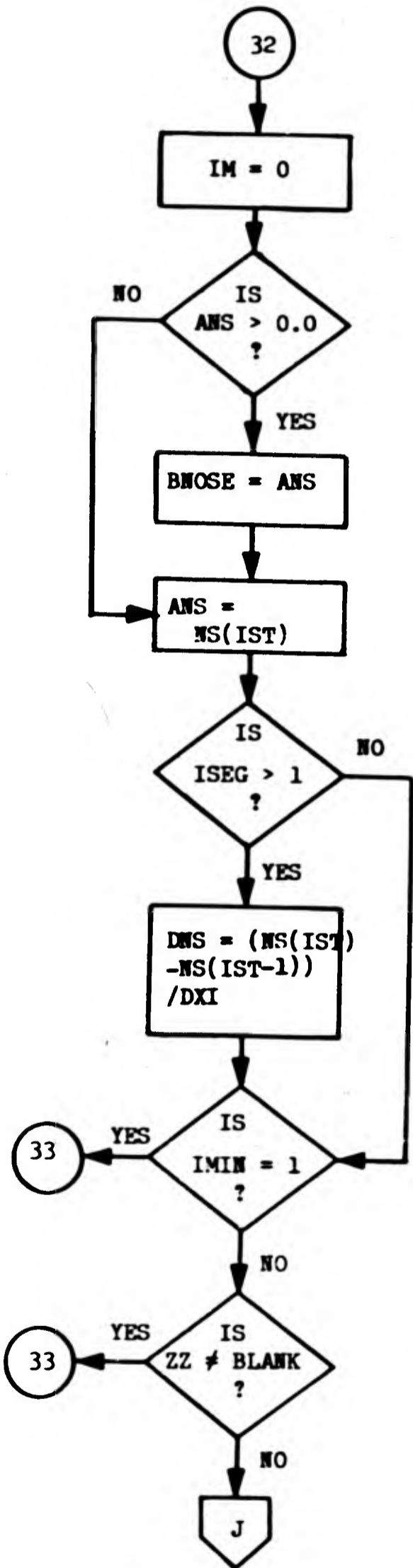
CLEAR UNNEEDED DATA CARDS
AND RETURN TO 28 TO READ
NEXT SET OF DATA CARDS.

FATAL ERROR. SEGMENT NUMBER
READ WAS GREATER THAN THE
CURRENT SEGMENT NUMBER.

PRINT ERROR MESSAGE.

SET ERROR INDICATOR.

RETURN.



CORRECT SEGMENT NUMBER READ.

SET IM TO READ GEOMETRY DATA FOR I = 1, MI.

WAS A NOSE BLUNTNES PARAMETER PROVIDED ?

LOAD NOSE BLUNTNES PARAMETER.

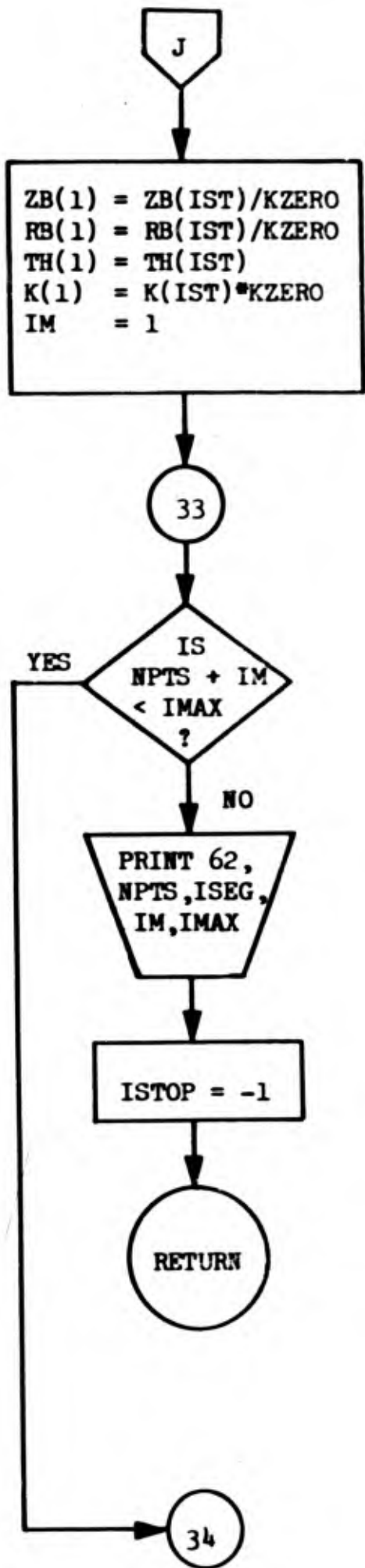
SAVE SHOCK LOCATION FOR UPSTREAM STARTING LINE.

SEGMENT NUMBER ≠ 1.

SAVE SHOCK SLOPE AT DOWNSTREAM BOUNDARY OF PREVIOUS SEGMENT.

NEED DATA FOR I = 1 ?

CHANGE CURVATURE AT I = 1 ?



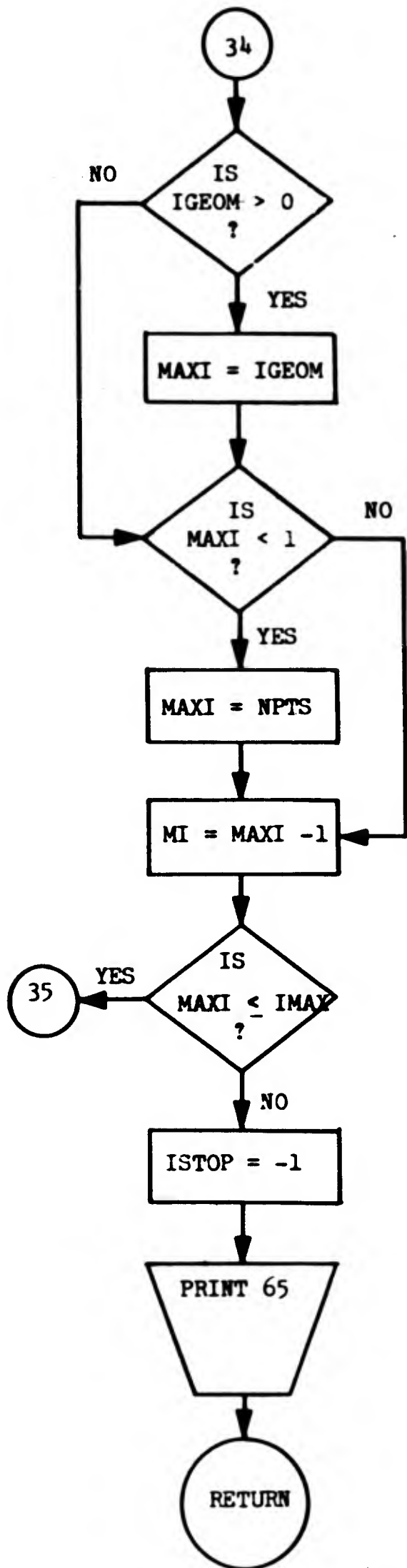
ASSIGN BODY GEOMETRY PARAMETERS FROM DOWNSTREAM BOUNDARY OF THE PREVIOUS SEGMENT TO THE UPSTREAM BOUNDARY OF THE CURRENT SEGMENT, SET IM TO READ DATA FOR I = 2, MI.

ARE DIMENSIONED VARIABLES LARGE ENOUGH TO HOLD ALL THE POINTS REQUIRED ?

PRINT ERROR MESSAGE.

SET ERROR INDICATOR.

RETURN.



WAS MAXI SPECIFIED IN CASE DATA ?

USE VALUE PROVIDED IN CASE DATA.

IS MAXI UNSPECIFIED ?

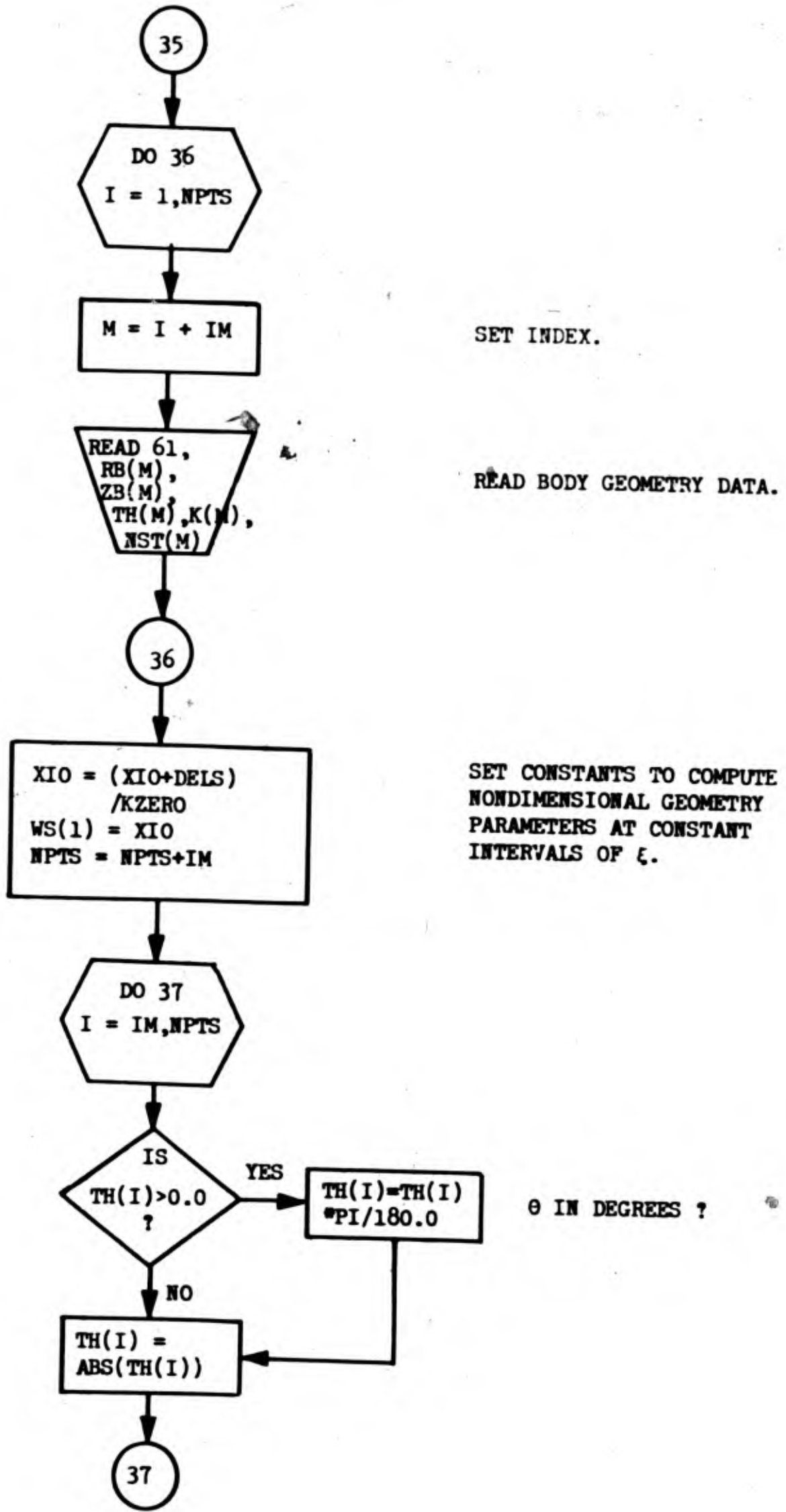
SET MAXI = NPTS.

DIMENSIONS LARGE ENOUGH FOR THIS VALUE OF IMAX ?

SET ERROR INDICATOR.

PRINT ERROR MESSAGE.

RETURN.

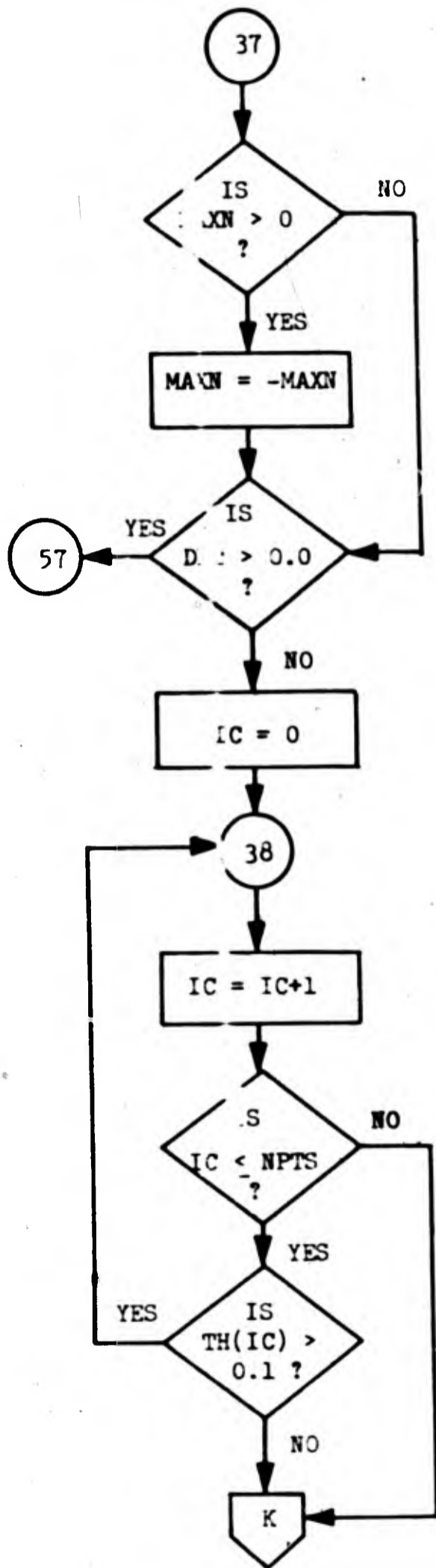


SET INDEX.

READ BODY GEOMETRY DATA.

SET CONSTANTS TO COMPUTE
NONDIMENSIONAL GEOMETRY
PARAMETERS AT CONSTANT
INTERVALS OF ξ .

θ IN DEGREES ?



WAS MAXN SPECIFIED ON
FIRST DATA CARD ?

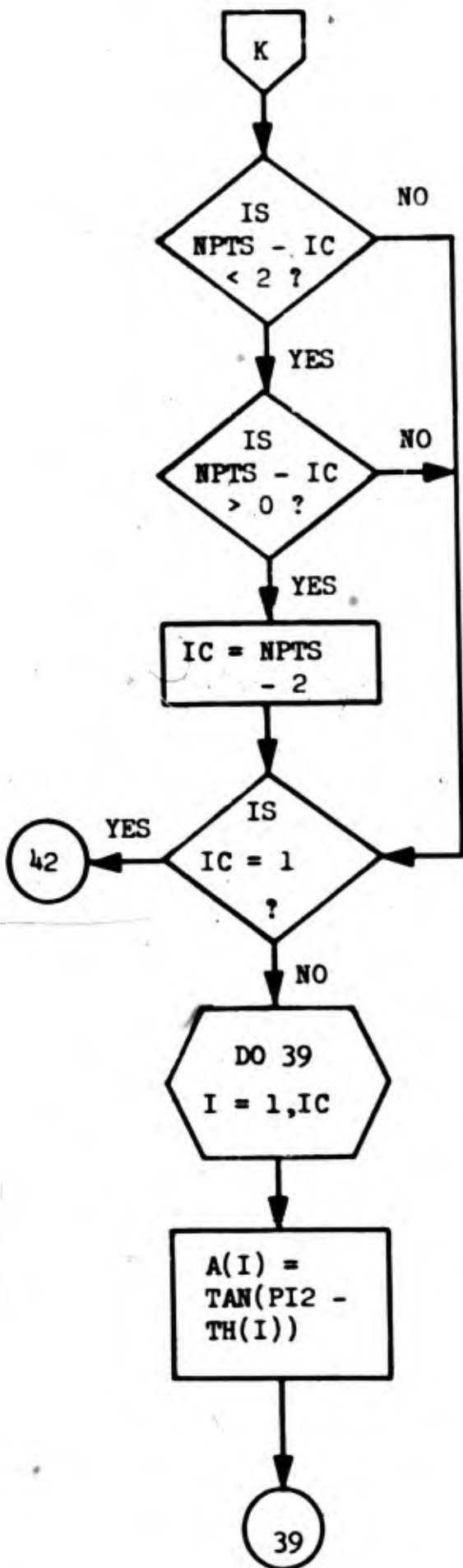
SET MAXN NEGATIVE TO
INDICATE THIS VALUE
SHOULD BE USED.

HAS DATA BEEN PROVIDED
AT EQUAL INTERVALS OF
 ξ ?

SET IC TO AVOID
SINGULAR DERIVATIVES.

IS IC WITHIN THE
SPECIFIED NUMBER OF POINTS ?

IS $\theta > 0.1$?



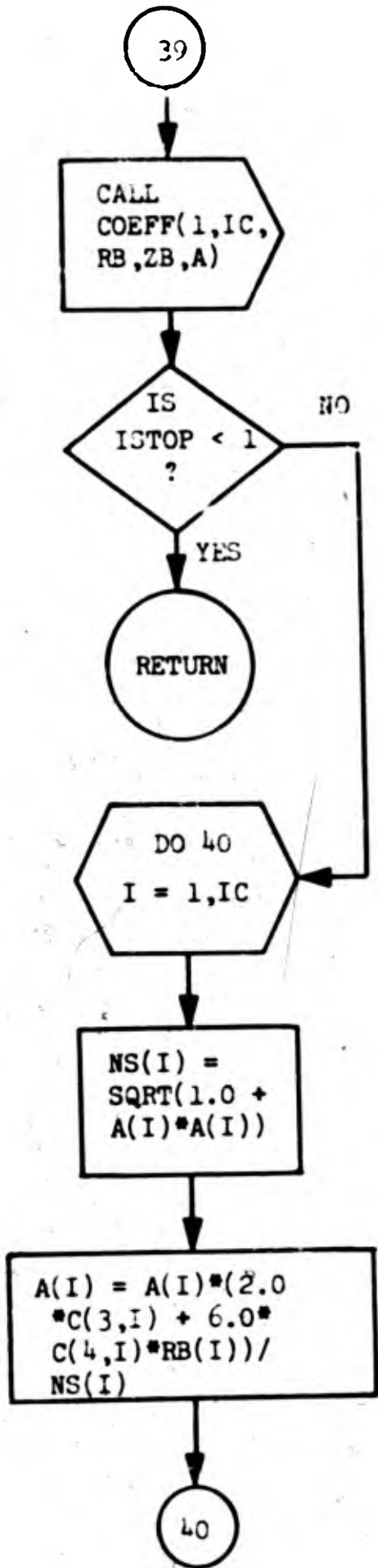
HAVE LESS THAN 2 POINTS BEEN ALLOCATED FOR INTEGRATION WITH RESPECT TO Z ?

HAVE ANY POINTS BEEN ALLOCATED ?

REQUIRE AT LEAST 2 IF ANY WERE USED.

ARE ALL INTEGRATIONS TO BE PERFORMED WITH RESPECT TO Z ?

$$\text{LOAD } \tan\left(\frac{\pi}{2} - \theta\right)$$



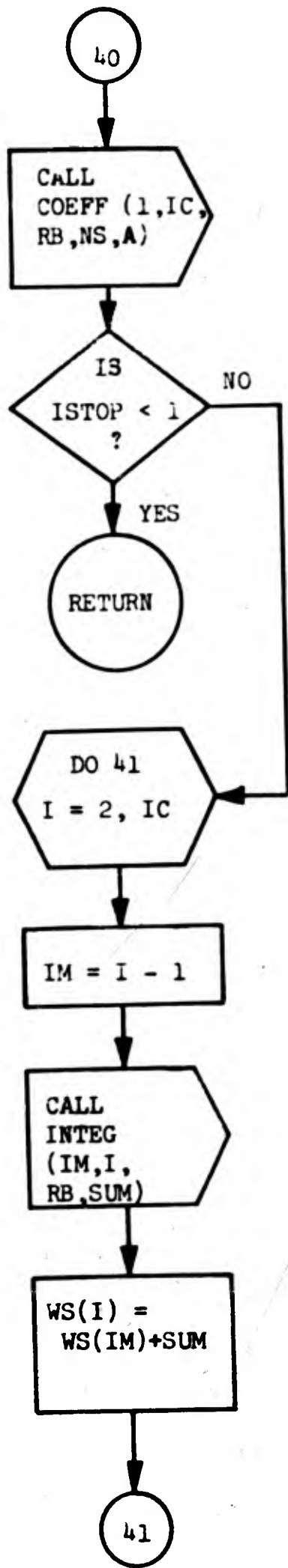
COMPUTE CUBIC COEFFICIENTS
FOR INTEGRATION WITH
RESPECT TO r.

ERROR IN SUBROUTINE COEFF ?

RETURN.

COMPUTE $\sqrt{1 + \left(\frac{dz}{dr}\right)^2}$

COMPUTE $\frac{d}{dr} \sqrt{1 + \left(\frac{dz}{dr}\right)^2}$

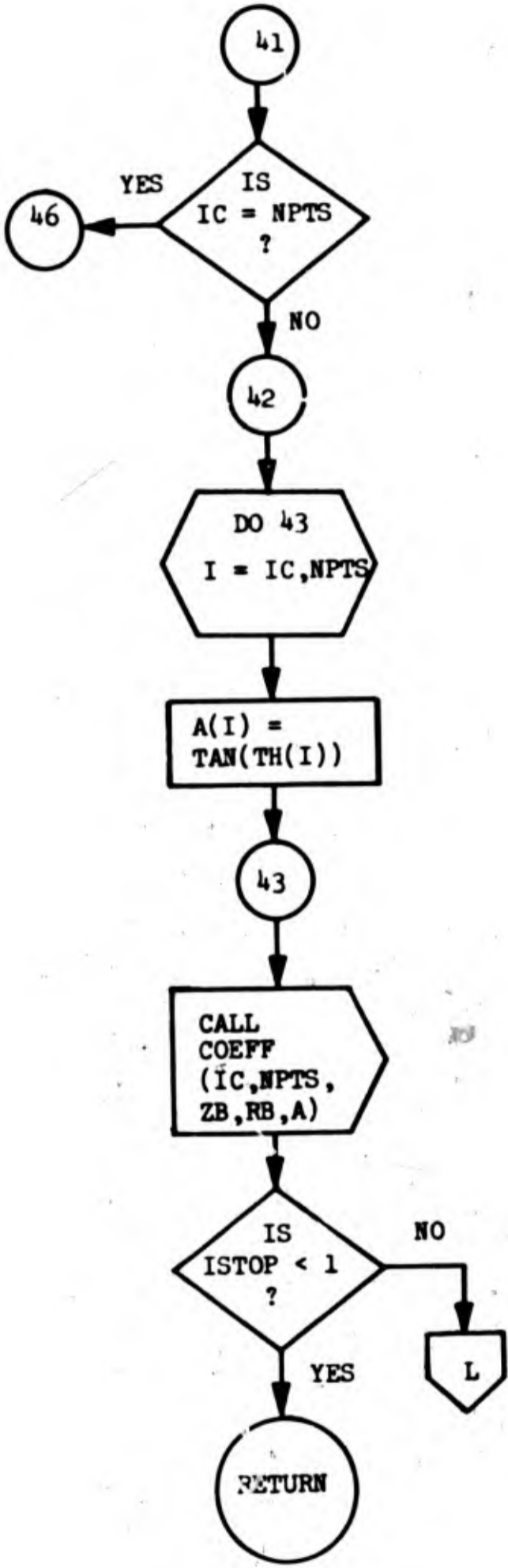


COMPUTE CUBIC COEFFICIENTS
FOR INTEGRATION WITH
RESPECT TO r .

ERROR IN SUBROUTINE COEFF ?

INTEGRATE FOR ξ AT I .

STORE ξ .



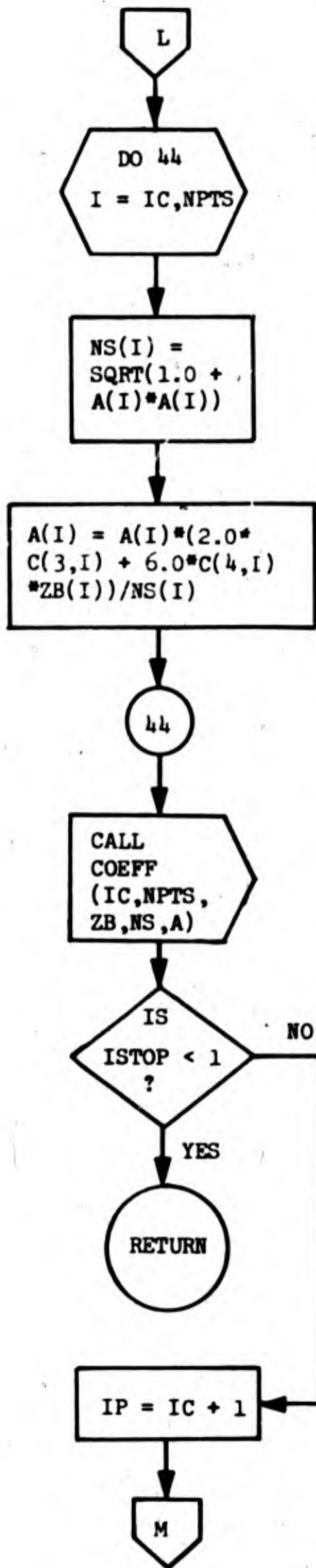
ARE ALL INTEGRATIONS TO
BE PERFORMED WITH RESPECT
TO r ?

COMPUTE $\frac{dr}{dz}$

COMPUTE COEFFICIENTS
FOR A CUBIC FIT OF
Z VS r.

ERROR IN SUBROUTINE COEFF ?

RETURN.



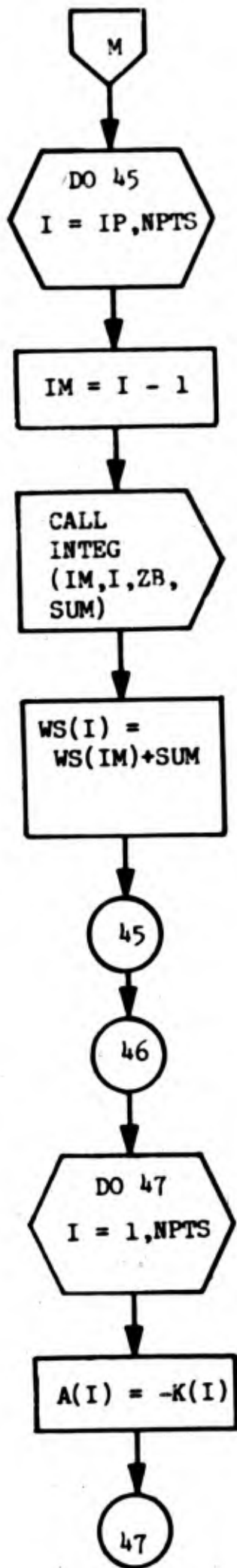
COMPUTE $\sqrt{1 + \left(\frac{dr}{dz}\right)^2}$

COMPUTE $\frac{d}{dz} \sqrt{1 + \left(\frac{dr}{dz}\right)^2}$

COMPUTE CUBIC COEFFICIENTS
FOR INTEGRATION WITH
RESPECT TO Z.

ERROR IN SUBROUTINE COEFF ?

RETURN.

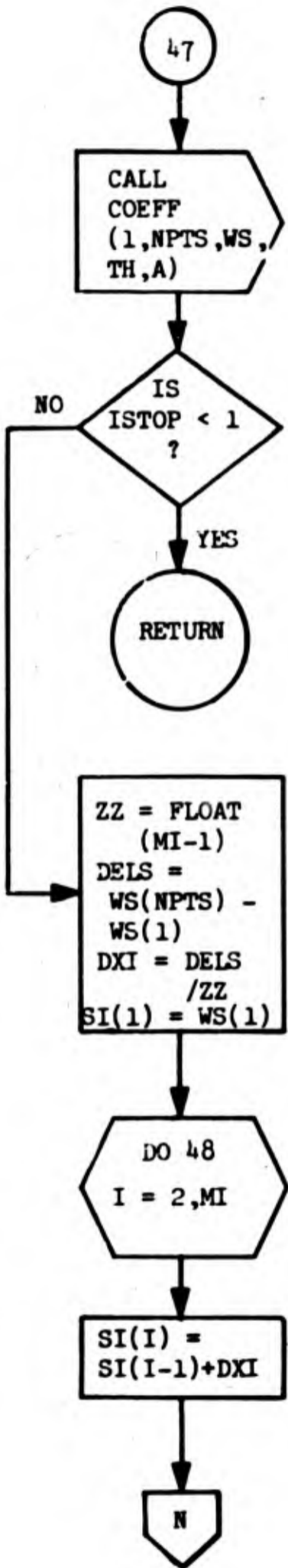


INTEGRATE FOR ξ AT I.

STORE ξ .

START INTERPOLATION.

LOAD $\frac{d\theta}{d\xi}$.

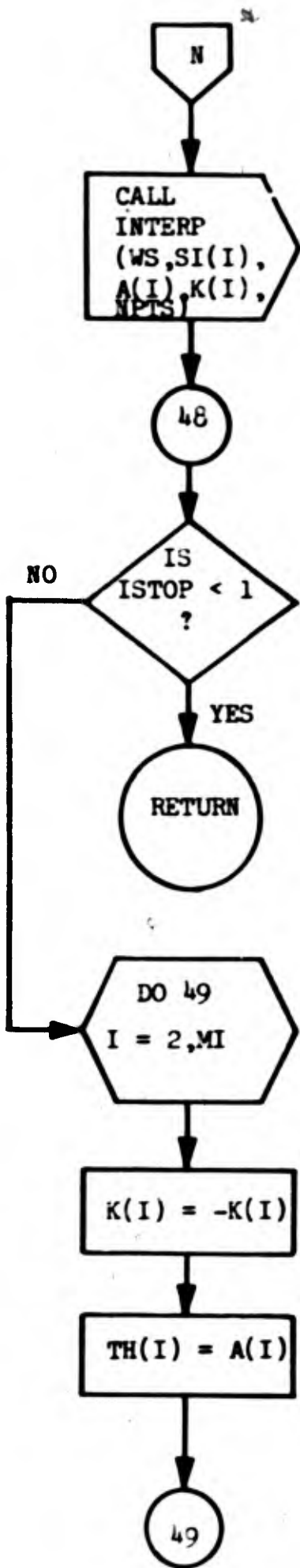


COMPUTE CUBIC COEFFICIENTS
FOR θ VS ξ .

ERROR IN SUBROUTINE COEFF ?

RETURN.

COMPUTE DELS AND $\Delta\xi$.



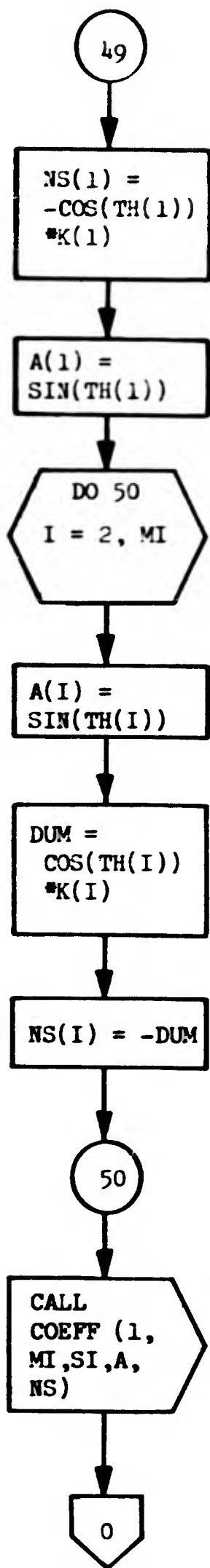
INTERPOLATE FOR θ AT ξ .

ERROR I: SUBROUTINE INTERP ?

RETURN.

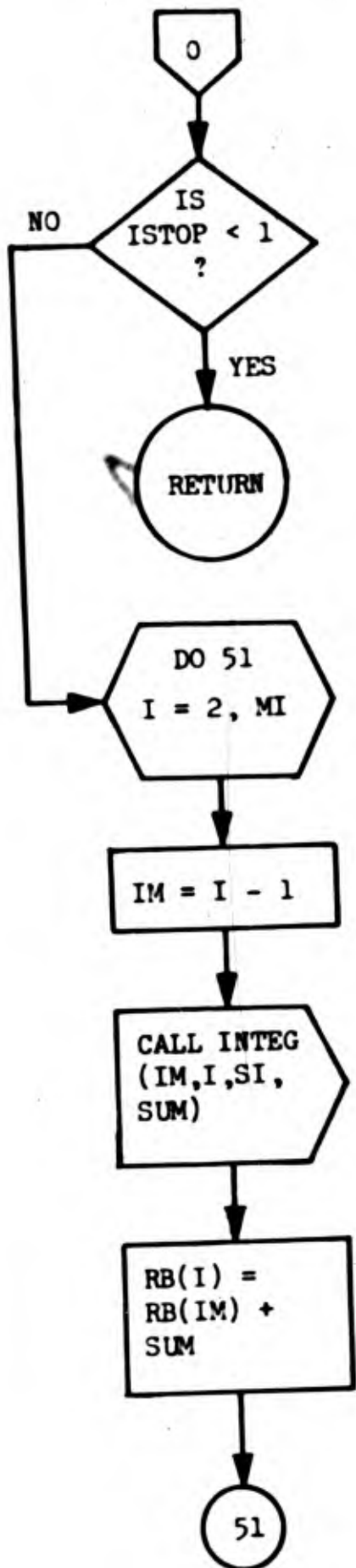
LOAD K AT ξ .

LOAD θ AT ξ .



LOOP TO LOAD PARAMETERS
IN EQUATION (1) TO
INTEGRATE FOR RB(I).

COMPUTE CUBIC COEFFICIENTS
TO INTEGRATE FOR RB(I).

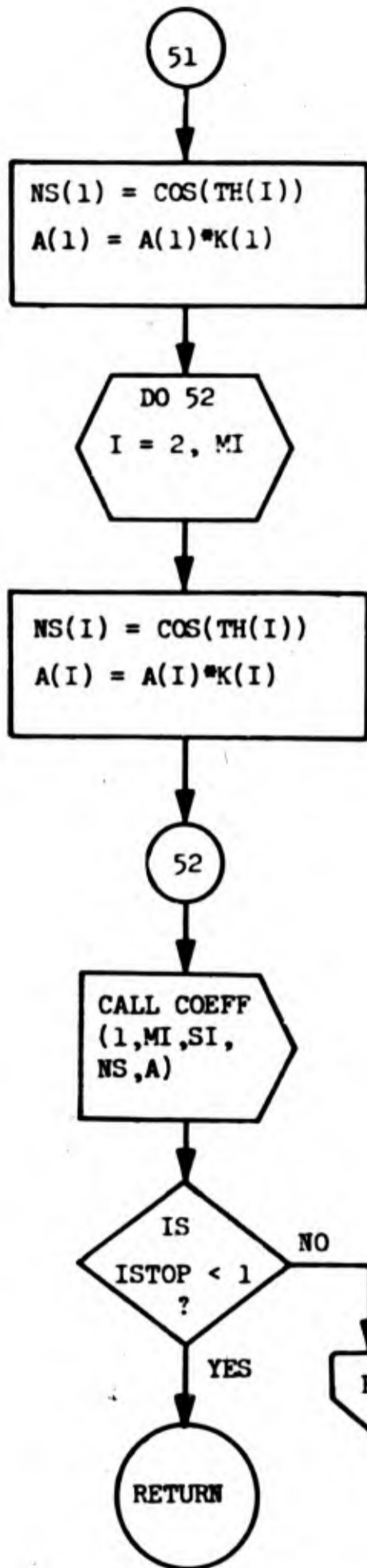


ERROR IN SUBROUTINE COEFF ?

RETURN.

INTEGRATE EQUATION (1).

STORE RB(I).

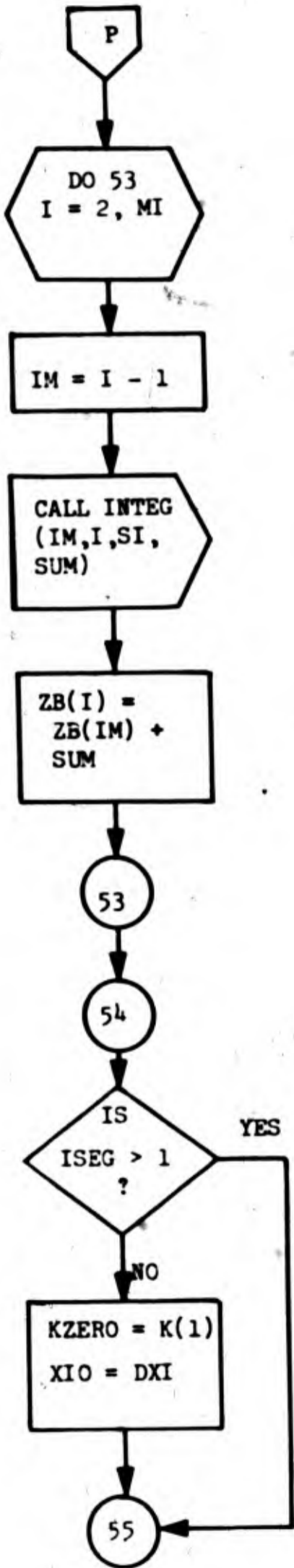


LOOP TO LOAD PARAMETERS FOR SOLVING EQUATION (2) FOR ZB(I).

COMPUTE CUBIC COEFFICIENTS FOR SOLVING EQUATION (2).

ERROR IN SUBROUTINE COEFF ?

RETURN.



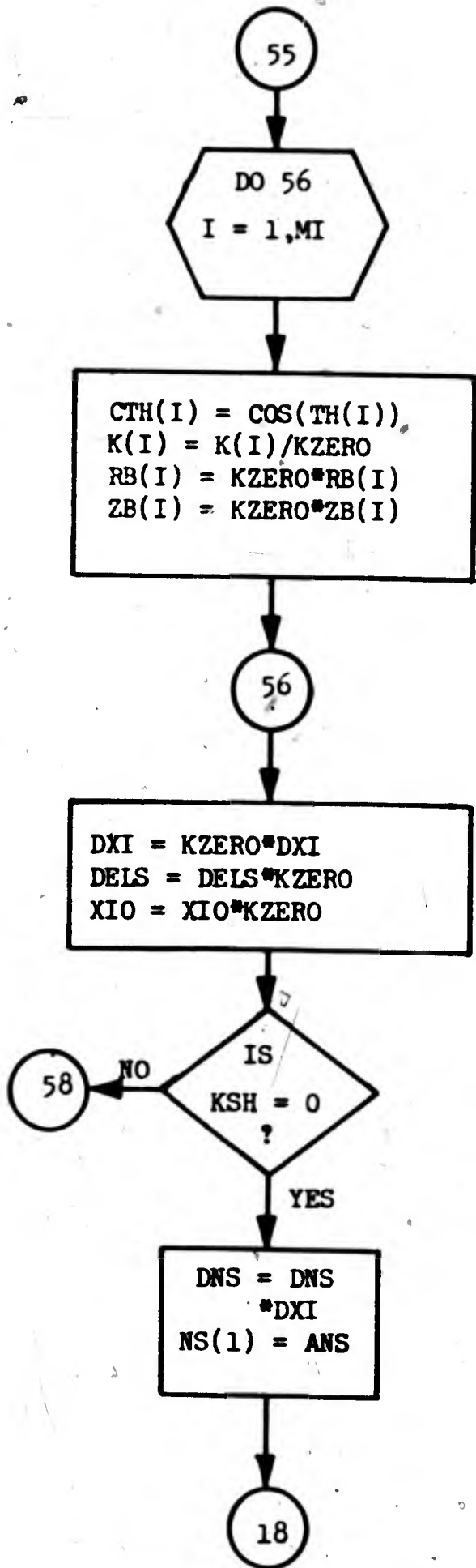
SOLVE EQUATION (2).

STORE ZB(I).

ENTRY POINT TO NON-DIMENSIONALIZE THE BODY GEOMETRY.

SEGMENT NUMBER > 1 ?

SAVE K(1) AND SET XIO FOR FIRST SEGMENT.



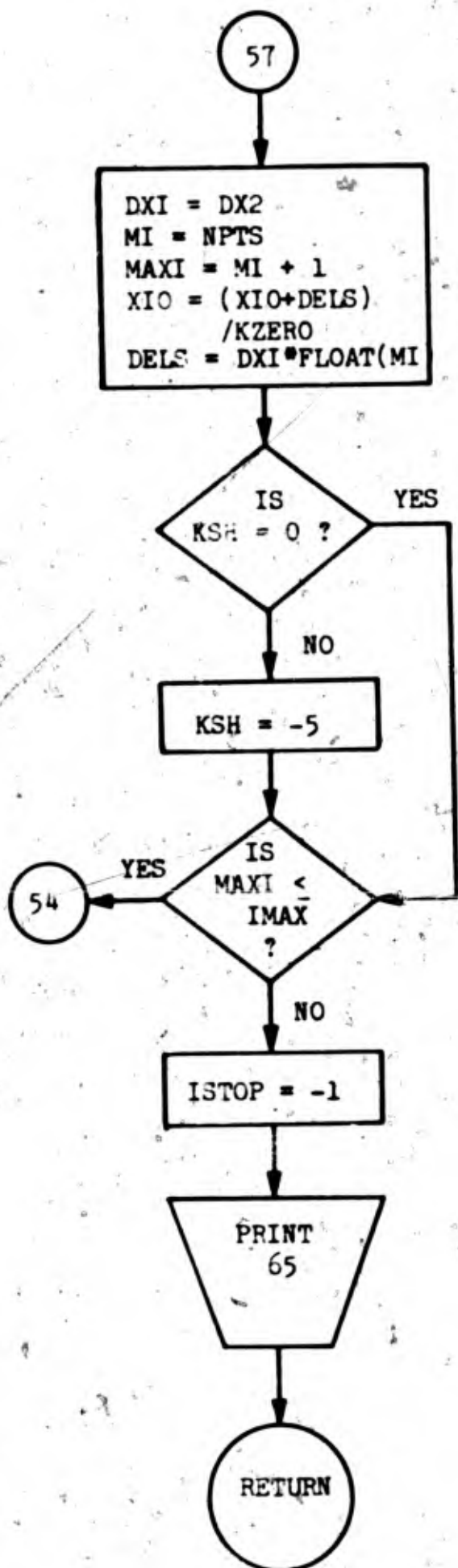
COMPUTE COS (θ) AND NON-DIMENSIONALIZE K, RB AND ZB.

NONDIMENSIONALIZE DXI, DELS AND XIO.

WAS THE SHOCK SHAPE READ IN ?

STORE SHOCK SLOPE AND NS AT THE CONSTANT STARTING LINE.

BODY GEOMETRY DATA WAS INPUT
IN EQUAL INCREMENTS OF S.



SET UP VALUES OF BODY GEOMETRY
PARAMETERS.

IS THE SHOCK SHAPE TO BE
ESTIMATED ?

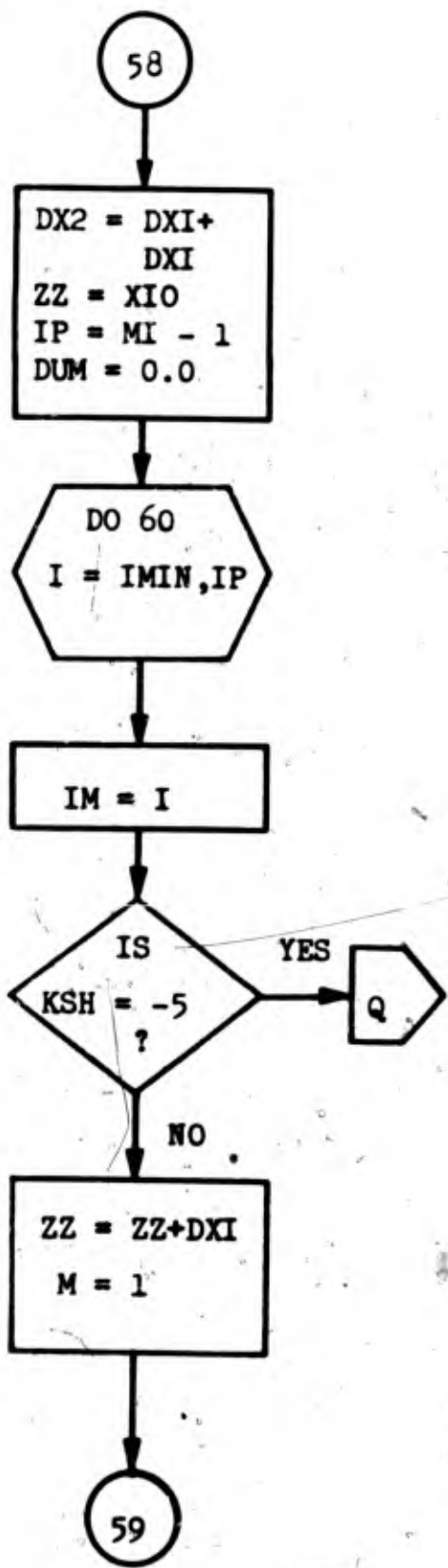
NO. PROVIDED IN INPUT DATA.
SET INDICATOR.

ARE DIMENSION SIZES LARGE ENOUGH
FOR THE NUMBER OF POINTS SUPPLIED ?

NO. SET ERROR FLAG.

PRINT ERROR MESSAGE.

RETURN.

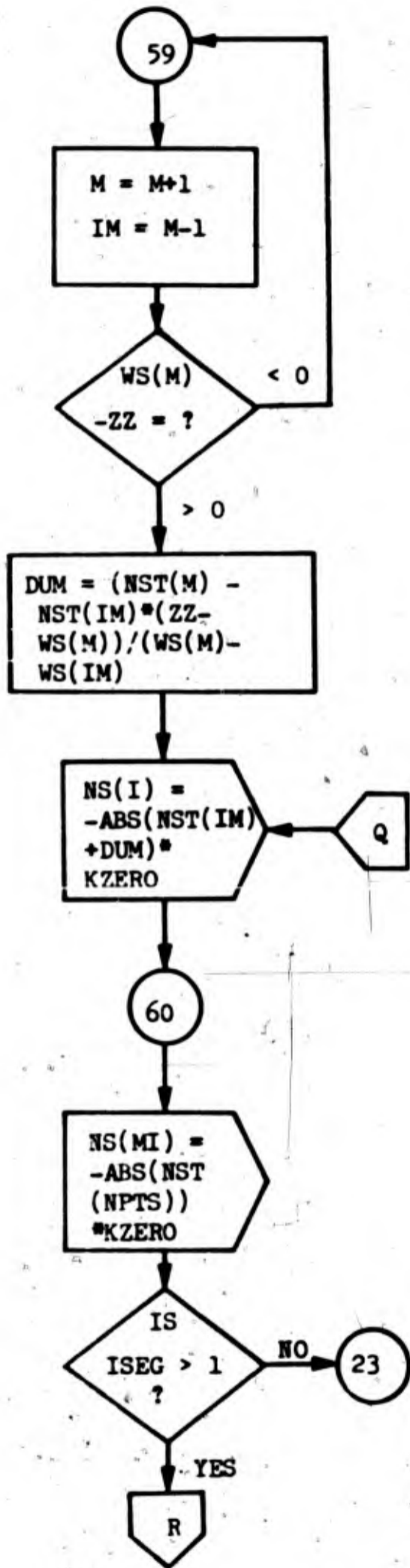


SHOCK SHAPE HAS BEEN INPUT.

SET GEOMETRY PARAMETERS.

START LOOP TO SET n_s .

WAS DATA INPUT IN EQUAL INCREMENTS OF S ?



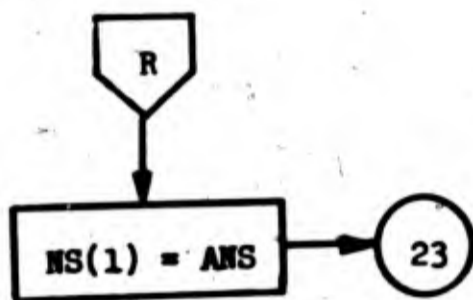
SEARCH FOR CORRECT
INPUT SHOCK VALUE.

INTERPOLATE FOR
POINT OF INTEREST.

SET n_s .

SET DOWNSTREAM
VALUE OF n_s .

SEGMENT NUMBER > 1 ?



RECOVER UPSTREAM
VALUE OF n_s .

END

7. LISTING

	SUBROUTINE GEOM (IST)	GEOM	1
C		GEOM	2
C	THIS ROUTINE PERFORMS THE BODY AND SHOCK GEOMETRY CALCULATIONS	GEOM	3
C	FOR EACH BODY SEGMENT.	GEOM	4
C		GEOM	5
	COMMON /2/ R2(15),Z0(15),TH(15),CTH(15),K(15),NS(15),MS(15)	GEOM	6
	COMMON /3/ MAXI,MI,IMIN,MAXN,MN,KMIN	GEOM	7
	COMMON /4/ IFEINT,NPR,KGECM	GEOM	8
	COMMON /5/ DXI,DET,DX2,DF2,DT	GEOM	9
	COMMON /6/ C1,VS1,HT1,A1,B1	GEOM	10
	COMMON /9/ ISEG,XI0,DELTA,KZERO,ROUSE	GEOM	11
	COMMON /11/ MAXIT,ISTCP,ITD,TIME,SI(15),NST(15)	GEOM	12
	COMMON /INPUT/ GAM1,AMACH1,THETA,ZRN,TCRIT(2),P1,RHO1,RN,AMU0,T1,EGEOM	GEOM	13
	IP2(2),IP2(6),INTAP,CTAP,NOPT,KCASE,NFTLE,NSEF,TGECM,NGEOM,NTAP	GEOM	14
	COMMON /CONST/ PI,PI2,GAMA,GAMB	GEOM	15
	COMMON /MESH/ IMAX,NMAX	GEOM	16
	COMMON /SPLINE/ C(4,15)	GEOM	17
	DIMENSION A(1)	GEOM	18
	DATA BLANK(17)	GEOM	19
	INTEGER DTAP	GEOM	20
	REAL NS,K,KZERO,NST	GEOM	21
	EQUIVALENCE (A,CTH)	GEOM	22
	IST=MI	GEOM	23
	IF (XI).LT.0.) XI0=0.0	GEOM	24
	GO TO (1,20,9), KGECM	GEOM	25
1	IF (ISEG.GT.1) GO TO 11	GEOM	26
2	IF (IGECM) 3,3,4	GEOM	27
3	ZZ=DELTA/1.15+2.0	GEOM	28
	MAXI=INT(ZZ)	GEOM	29
4	ZZ=FLOAT(MAXI-2)	GEOM	30
	DXI=DELTA/ZZ	GEOM	31
	IF (MAXI.LE.1MAX) GO TO 5	GEOM	32
	ZZ=FLOAT(1MAX-2)	GEOM	33
	DXI=DELTA/ZZ	GEOM	34
	MAXI=1MAX	GEOM	35
5	MI=MAXI-1	GEOM	36
	XI0=-DXI	GEOM	37
	K(1)=1.0	GEOM	38
	RO(1)=1.0	GEOM	39
	Z0(1)=0.0	GEOM	40
	TH(1)=PI2	GEOM	41
	CTH(1)=0.0	GEOM	42
	ZZ=0.0	GEOM	43
	GO TO 7	GEOM	44
6	ZZ=XI0	GEOM	45
7	DO 8 I=2,MI	GEOM	46
	ZZ=ZZ+DXI	GEOM	47
	K(I)=1.0	GEOM	48
	TH(I)=PI2-ZZ	GEOM	49
	Z0(I)=1.0-COS(ZZ)	GEOM	50
	RO(I)=SIN(ZZ)	GEOM	51
8	CTH(I)=COS(TH(I))	GEOM	52


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11 11 11,11
12 12 11
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15 15 11
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GECM 53
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GECM 105
GECM 106
GECM 107

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10  X(I)=Y(I)*Z(I)
11  Y(I)=Y(I)*Z(I)
12  Z(I)=Y(I)*Z(I)
13  Y(I)=Y(I)*Z(I)
14  Z(I)=Y(I)*Z(I)
15  Y(I)=Y(I)*Z(I)
16  Z(I)=Y(I)*Z(I)
17  K(I)=Z(I)
18  DX2=DX1+DXI
19  IF (MAYT) 10,25,25
20  DO 21 I=1,MY
21  NS(I)=-1.01
22  GO TO 23
23  DUM=AMACH1*AMACH1
24  Z7=(GAMMA*DUM+1.01)/(GAMMA*DUM)
25  NS(I)=-1.01*(Z7**1.157)*(KZE**0.157)
26  DUM=SQRT(DELTA/ANS)
27  DMS=(1.24+1.24/AMACH1+NS(I))
28  Z7=1.0
29  DO 22 I=1,MY
30  Z7=Z7+DXI
31  NS(I)=NS(I)-DMS*(Z7/DELTA)**7)
32  DO 24 I=1,MAXI
33  NS(I)=1.0
34  TH=M[-1]
35  Z7=K(MI)*DXI/2.0
36  CTH=STH/TH(MI)
37  D9(MAYT)=D9(MI)+DXI*(STH-Z7*CTH(MI))
38  Z7(MAYT)=Z7(MI)+DXI*(CTH(MI)+Z7*STH)
39  TH(MAYT)=TH(MI)+TH(MI)-TH(IM)
40  NS(MAYT)=NS(MI)+NS(MI)-NS(IM)
41  CTH(MAYT)=DMS/TH(MAYT)
42  K(MAYT)=K(MI)+K(MI)-K(IM)
43  D9TUM
44  DO 24 I=1,MY
45  DMS=K(I)
46  ANS=ANS+K(I)
47  ANS=ANS/FLOAT(MI)
48  IF (TSEG.50.1) GO TO 21
49  Z7=ANS**1.25
50  DUM=YT)+1.5*DELS
51  ANS=DMS/2.0
52  DMS=ANS*(1.0+(2.0*Z7-1.0)/DUM)
53  IF (ANS.GT.ANS) ANS=ANS
54  DO 22 I=1,MY
55  NS(I)=NS(I-1)+DMS
56  GO TO 23
57  READ 67, M,NEIS,MAXT,IP,KSH,ANS,DX2,Z7
58  IF (M.TSEG) 20,32,31
59  DO 31 M=1,NPTS
60  READ 61, DUM
61  GO TO 23
62  WRITE 66
63  STOP=-1

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GECM 109
GECM 110
GECM 111
GECM 112
GECM 113
GECM 114
GECM 115
GECM 116
GECM 117
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GECM 160
GECM 161
GECM 162

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RETURN
72 IM=0
IF (ANS.GT.0) ANOSE=ANS
ANS=NS(IST)
IF (ISEG.GT.1) DNS=(NS(IST)-NS(IST-1))/DXI
IF (IMIN.EQ.1.C7.ZZ.NF.BLANK) GO TO 33
Z9(1)=Z9(IST)/KZERO
P9(1)=P9(IST)/KZERO
TH(1)=TH(IST)
K(1)=K(IST)*KZERO
IM=1
73 IF (NPTS+IM.LE.IMAX) GO TO 34
PRINT 62,NPTS,ISEG,IM,IMAX
ISTOP=-1
RETURN
34 IF (IGEOM.GT.0) MAXI=IGEOM
IF (MAXI.LT.1) MAXI=NPTS
NI=MAXI-1
IF (MAXI.LE.IMAX) GO TO 35
ISTOP=-1
PRINT 65
RETURN
75 DO 36 I=1,NPTS
M=I+IM
76 READ 61, P9(M),Z9(M),TH(M),K(M),NST(M)
XIO=(XIO+DELS)/KZERO
NS(I)=YI?
NPTS=NPTS+IM
IM=IM+1
DO 37 I=IM,NPTS
IF (TH(I).GT.0.0) TH(I)=TH(I)*PI/180.0
77 TH(I)=ABS(TH(I))
IF (IP.GT.0) MAXN=-IP
IF (IX2.GT.0.0) GO TO 57
IC=0
38 IC=IC+1
IF (IC.LT.NPTS.AND.TH(IC).GT.0.1) GO TO 38
IF (NPTS-IC.LT.2.AND.NPTS-IC.NE.0) IC=NPTS-2
IF (IC.EQ.1) GO TO 42
DO 39 I=1,IC
79 A(I)=TAN(PI?-TH(I))
CALL COEFF (1,IC,RB,Z9,A)
IF (ISTOP.LT.1) RETURN
DO 40 I=1,IC
NS(I)=SQRT(1.)+A(I)*A(I)
40 A(I)=A(I)*(2.0*C(3,I)+6.0*C(4,I)*RB(I))/NS(I)
CALL COEFF (1,IC,RB,NS,A)
IF (ISTOP.LT.1) RETURN
DO 41 I=2,IC
IM=I-1
CALL INTEG (IM,I,RB,SUM)
41 WS(I)=WS(IM)+SUM
IF (IC.EQ.NPTS) GO TO 46
42 DO 43 I=IC,NPTS
43 A(I)=TAN(TH(I))

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GECM 163
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GECM 217

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CALL COEFF (IC,NPTS,ZB,PA,A)
IF (ISTOP.LT.1) RETURN
DO 44 I=IC,NPTS
NS(I)=SQRT(1.0+A(I)*A(I))
X(I)=A(I)*(2.)*C(3,I)+6.)*C(4,I)*ZB(I))/NS(I)
CALL COEFF (IC,NPTS,ZB,NS,A)
IF (ISTOP.LT.1) RETURN
IP=IC+1
DO 45 I=IP,NPTS
IM=I-1
CALL INTEG (IM,I,ZB,SUM)
45 WS(I)=WS(IM)+SUM
46 DO 47 I=1,NPTS
47 A(I)=-K(I)
CALL COEFF (1,NPTS,WS,TH,A)
IF (ISTOP.LT.1) RETURN
ZZ=FLOAT(MI-1)
DELS=WS(NPTS)-WS(1)
DXT=7FELS/ZZ
SI(1)=WS(1)
DO 48 I=2,MI
48 SI(I)=SI(I-1)+DXT
CALL INTERP (WS,SI(I),A(I),K(I),NPTS)
IF (ISTOP.LT.1) RETURN
DO 49 I=2,MI
49 K(I)=-K(I)
50 TH(I)=A(I)
NS(I)=-COS(TH(I))*K(I)
A(I)=SIN(TH(I))
DO 50 I=2,MI
50 A(I)=SIN(TH(I))
DUM=COS(TH(I))*K(I)
NS(I)=-DUM
CALL COEFF (1,MI,SI,A,NS)
IF (ISTOP.LT.1) RETURN
DO 51 I=2,MI
51 IM=I-1
CALL INTEG (IM,I,SI,SUM)
NS(I)=NS(IM)+SUM
A(I)=A(I)*K(I)
DO 52 I=2,MI
52 NS(I)=COS(TH(I))
A(I)=A(I)*K(I)
CALL COEFF (1,MI,SI,NS,A)
IF (ISTOP.LT.1) RETURN
DO 53 I=2,MI
53 IM=I-1
CALL INTEG (IM,I,SI,SUM)
ZB(I)=ZB(IM)+SUM
54 IF (ISEG.GT.1) GO TO 55
KZER0=K(1)
KI)=-DXT
55 DO 55 I=1,MI
OTH(I)=COS(TH(I))

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GECM 218
GECM 219
GECM 220
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GECM 272

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      X1)=K(X1)/K7500
      X2)=K7500*(X1)
      X3)=K7500*(X1)
      X4)=K7500*(X1)
      X5)=K7500*(X1)
      X6)=K7500*(X1)
      X7)=K7500*(X1)
      X8)=K7500*(X1)
      X9)=K7500*(X1)
      X10)=K7500*(X1)
      X11)=K7500*(X1)
      X12)=K7500*(X1)
      X13)=K7500*(X1)
      X14)=K7500*(X1)
      X15)=K7500*(X1)
      X16)=K7500*(X1)
      X17)=K7500*(X1)
      X18)=K7500*(X1)
      X19)=K7500*(X1)
      X20)=K7500*(X1)
      X21)=K7500*(X1)
      X22)=K7500*(X1)
      X23)=K7500*(X1)
      X24)=K7500*(X1)
      X25)=K7500*(X1)
      X26)=K7500*(X1)
      X27)=K7500*(X1)
      X28)=K7500*(X1)
      X29)=K7500*(X1)
      X30)=K7500*(X1)
      X31)=K7500*(X1)
      X32)=K7500*(X1)
      X33)=K7500*(X1)
      X34)=K7500*(X1)
      X35)=K7500*(X1)
      X36)=K7500*(X1)
      X37)=K7500*(X1)
      X38)=K7500*(X1)
      X39)=K7500*(X1)
      X40)=K7500*(X1)
      X41)=K7500*(X1)
      X42)=K7500*(X1)
      X43)=K7500*(X1)
      X44)=K7500*(X1)
      X45)=K7500*(X1)
      X46)=K7500*(X1)
      X47)=K7500*(X1)
      X48)=K7500*(X1)
      X49)=K7500*(X1)
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      X64)=K7500*(X1)
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      X67)=K7500*(X1)
      X68)=K7500*(X1)
      X69)=K7500*(X1)
      X70)=K7500*(X1)
      X71)=K7500*(X1)
      X72)=K7500*(X1)
      X73)=K7500*(X1)
      X74)=K7500*(X1)
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      X77)=K7500*(X1)
      X78)=K7500*(X1)
      X79)=K7500*(X1)
      X80)=K7500*(X1)
      X81)=K7500*(X1)
      X82)=K7500*(X1)
      X83)=K7500*(X1)
      X84)=K7500*(X1)
      X85)=K7500*(X1)
      X86)=K7500*(X1)
      X87)=K7500*(X1)
      X88)=K7500*(X1)
      X89)=K7500*(X1)
      X90)=K7500*(X1)
      X91)=K7500*(X1)
      X92)=K7500*(X1)
      X93)=K7500*(X1)
      X94)=K7500*(X1)
      X95)=K7500*(X1)
      X96)=K7500*(X1)
      X97)=K7500*(X1)
      X98)=K7500*(X1)
      X99)=K7500*(X1)
      X100)=K7500*(X1)

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NOT REPRODUCIBLE

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GEOM 323

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10. SUBROUTINE INITIAL

A. PURPOSE

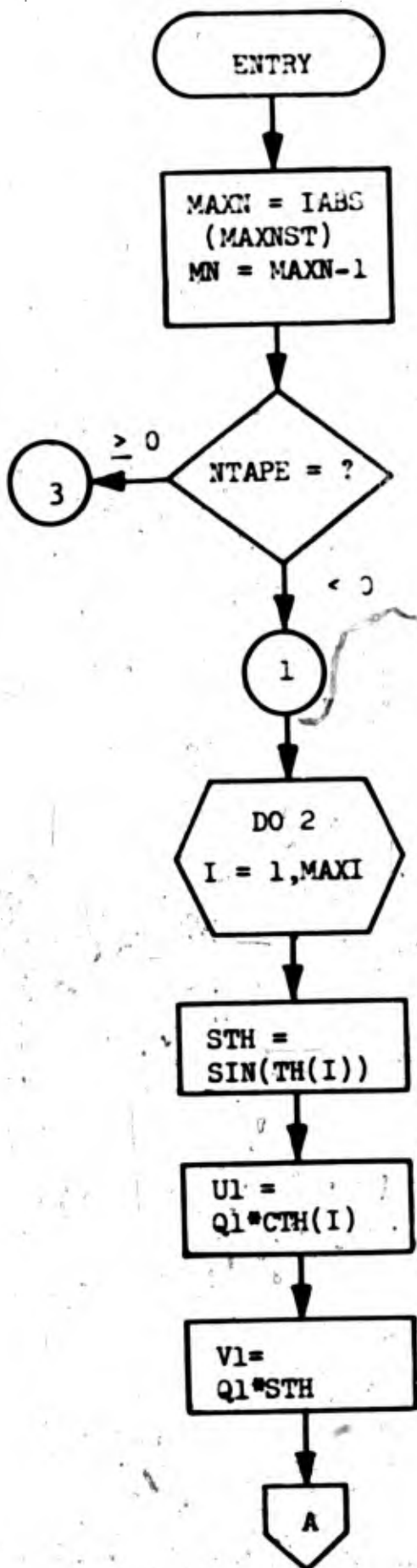
This subroutine provides initial conditions at all nodes to start the solution. The shock jump conditions are applied to the shock shape point parameters. At body points, the surface boundary condition, an empirical pressure distribution (similar to modified Newtonian theory), conservation of total enthalpy and the known body streamline entropy supply the parameters. Intermediate points are treated by interpolation between the shock and body based on a variation proportional to η^2 .

B. VARIABLE LIST

All variables previously defined in the description of the main program have been omitted from this list.

- AMN - $(M \sin \beta)^2$
- CBET - $\cos \beta$
- DUM - dummy variable used for intermediate computations.
- M - dummy index.
- MAXNST - specifies the number of nodes to be used across the shock layer (may be negative).
- PB - pressure on the body.
- SBET - $\sin \beta$
- U1 - component of the free-stream velocity parallel to the shock.
- V1 - component of the free-stream velocity normal to the shock.
- V2 - component of the velocity at a shock point normal to the shock.
- ZZ - dummy variable used for intermediate computations.

C. FLOW CHART



ENTER SUBROUTINE INITIAL

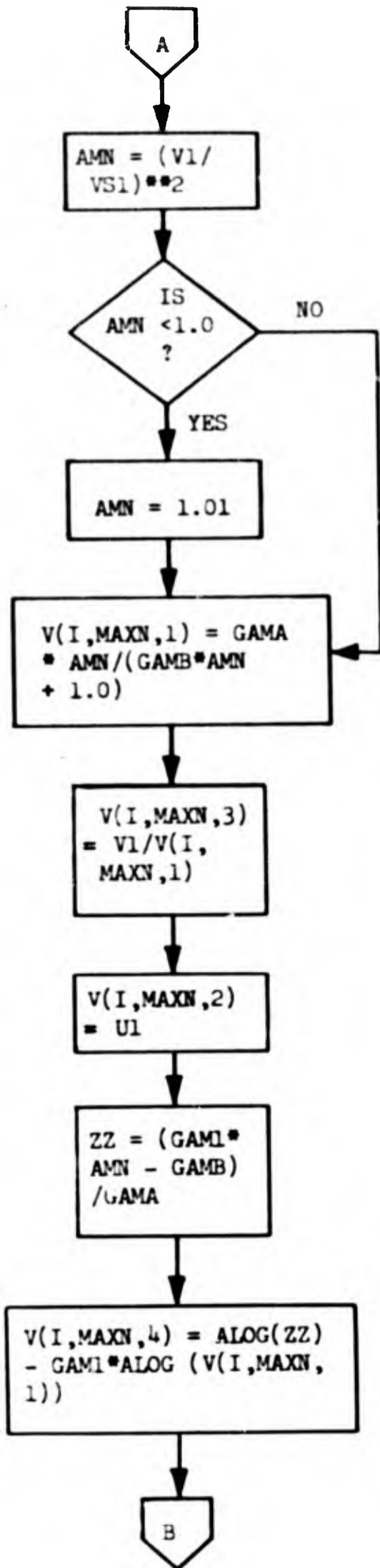
INSURE THAT THE INDEX OF THE SHOCK NODES IS POSITIVE.

UNSTEADY FLOW INITIAL CONDITIONS.

COMPUTE SIN θ .

COMPUTE U .

COMPUTE V .



COMPUTE THE SQUARE OF THE MACH NUMBER NORMAL TO THE SHOCK.

IS $AMN < 1.0$?

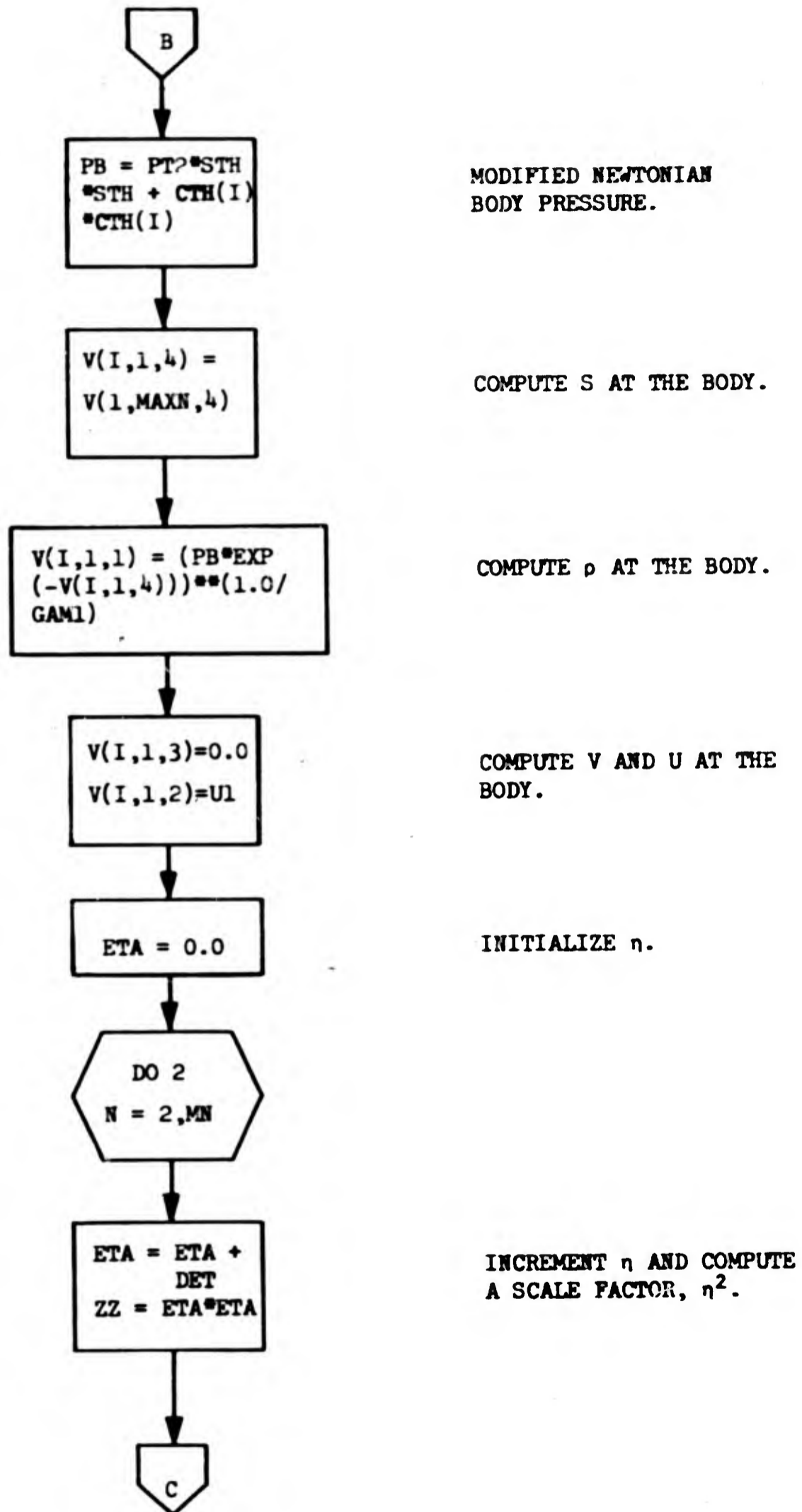
SET $AMN = 1.01$.

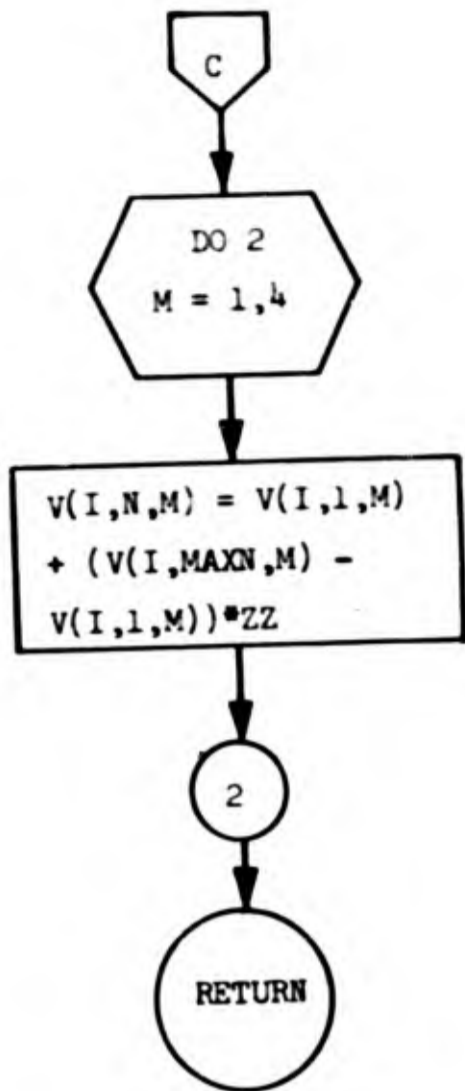
COMPUTE ρ AT THE SHOCK.

COMPUTE V AT THE SHOCK.

COMPUTE U AT THE SHOCK.

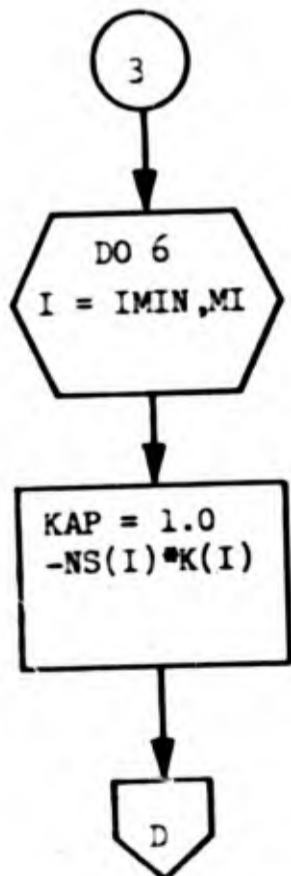
COMPUTE S AT THE SHOCK.





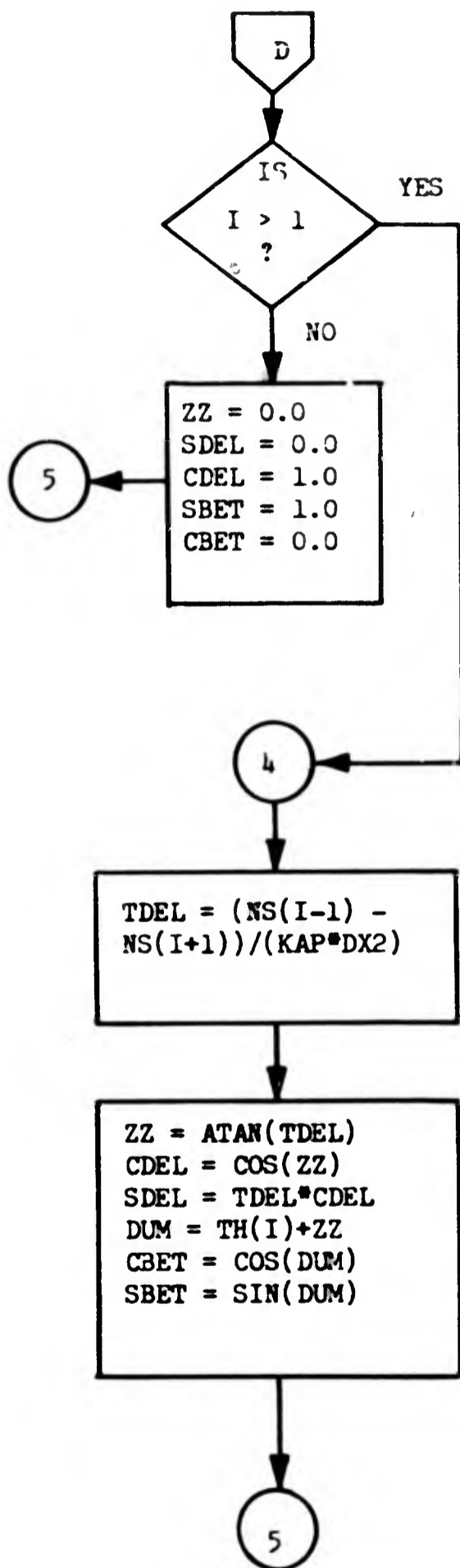
COMPUTE ρ , U, V AND S AT
REMAINING NODES ACROSS
THE SHOCK LAYER.

RETURN



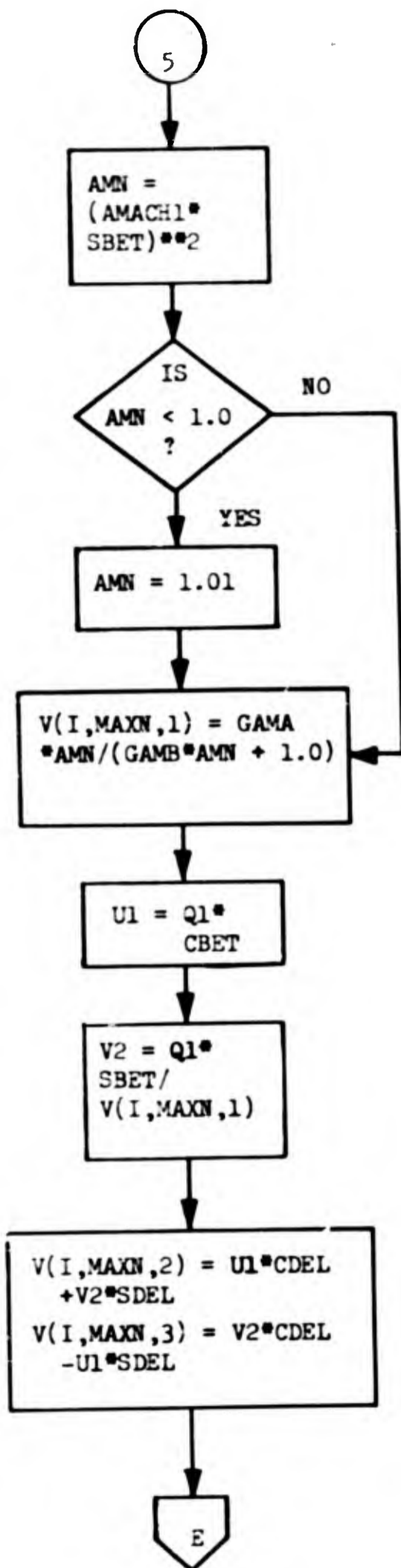
STEADY STATE INITIAL
CONDITIONS

COMPUTE κ AT THE SHOCK.



COMPUTE SHOCK GEOMETRY
FOR STAGNATION POINT.

COMPUTE SHOCK GEOMETRY
FOR OTHER POINTS.



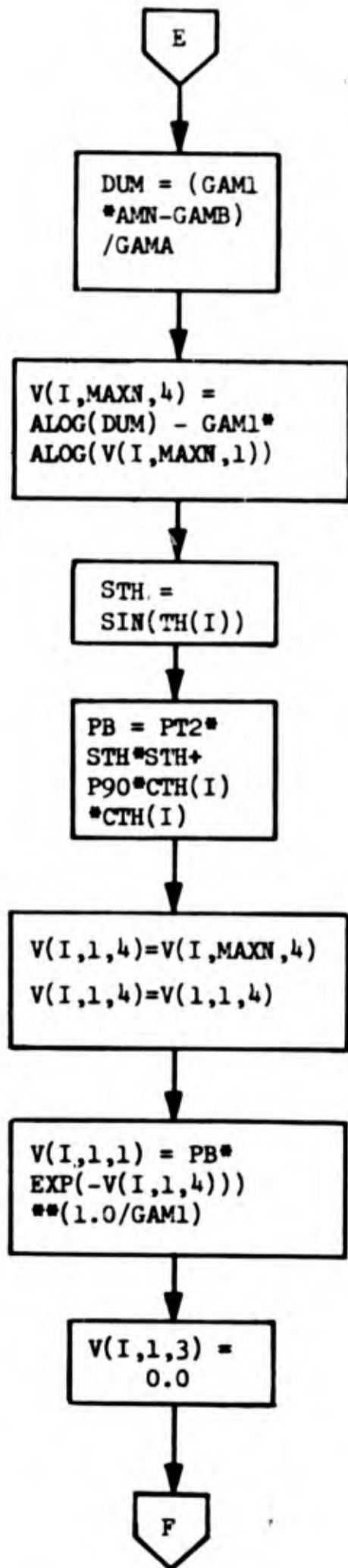
COMPUTE THE SQUARE OF THE MACH NUMBER NORMAL TO THE SHOCK.

AMN < 1.0 ?

COMPUTE ρ AT THE SHOCK.

COMPUTE VELOCITY COMPONENT TANGENT TO THE SHOCK.

COMPUTE VELOCITY COMPONENT NORMAL TO THE SHOCK.



COMPUTE S AT THE SHOCK.

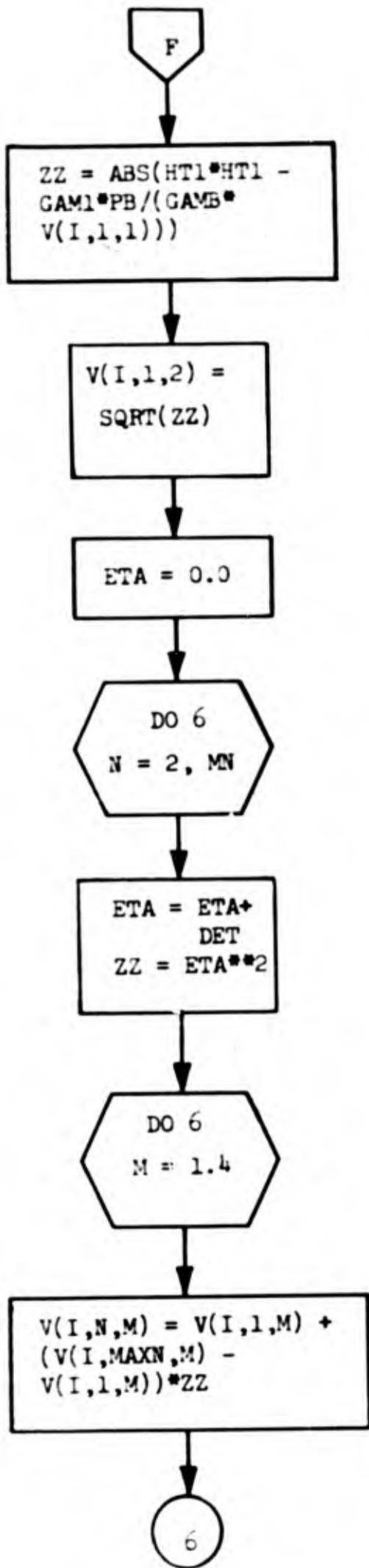
COMPUTE SIN θ .

ESTIMATE BODY PRESSURE.

CHOOSE MAXIMUM ENTROPY FOR BODY POINTS.

COMPUTE ρ AT THE BODY.

SET V = 0 AT THE BODY.

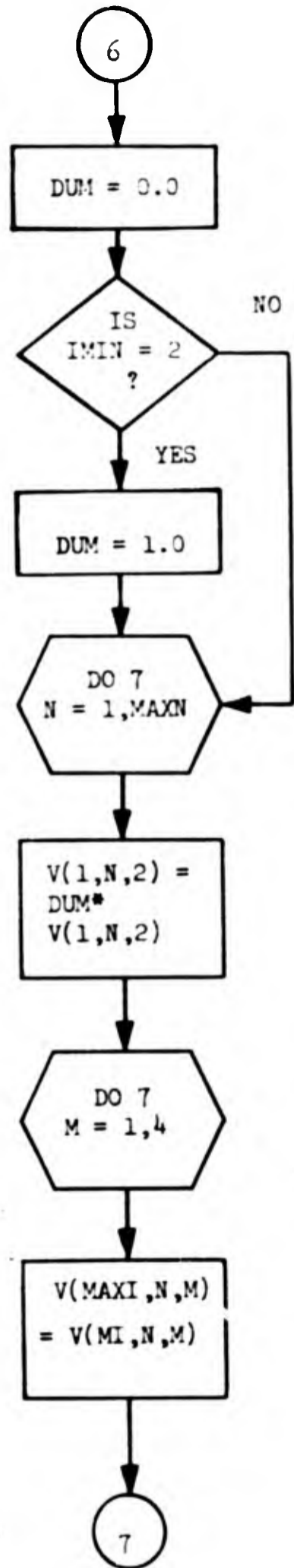


SET U AT THE BODY TO CONSERVE TOTAL ENTHALPY.

INITIALIZE η .

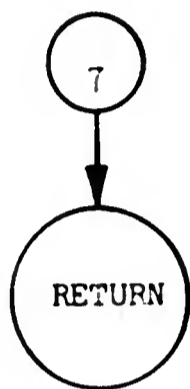
INCREMENT η AND COMPUTE SCALE FACTOR, η^2 .

COMPUTE ρ , U, V AND S AT REMAINING POINTS ACROSS THE SHOCK LAYER.

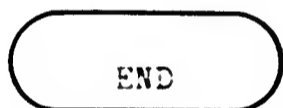


ZERO U AT I = 1 IF IT IS
THE STAGNATION STREAMLINE.

SET VALUES OF ρ , U, V AND
S AT LAST BODY STATION
NODES.



RETURN.



END OF SUBROUTINE INITIAL.


```

COMMON /INITIAL (MAXN)
      COMMON /INITIAL PROVIDES THE INITIAL GUESSES FOR FLOW FIELD
      CALCULATIONS.
COMMON /1/ (X1,4), (Y1,4), (Z1,4), (P1,4)
COMMON /2/ ZR(15), ZS(15), TH(15), CTH(15), K(15), NS(15), WS(15)
COMMON /3/ MAXI, MI, IMIN, MAXACT, MN, NMIN
COMMON /4/ DXT, DFT, DYT, DFT, DT
COMMON /5/ Q1, VS1, HT1, S1, P1
COMMON /INPUT/ GAM1, GAM2, THETA, ZRN, TPRIT(2), S1, SHC1, RN, AMN, Z1,
      (DZ1,1), (DZ1,2), (DZ1,3), (DZ1,4), (DZ1,5), (DZ1,6), (DZ1,7), (DZ1,8), (DZ1,9), (DZ1,10), (DZ1,11), (DZ1,12),
      (DZ1,13), (DZ1,14), (DZ1,15), (DZ1,16), (DZ1,17), (DZ1,18), (DZ1,19), (DZ1,20), (DZ1,21), (DZ1,22), (DZ1,23), (DZ1,24), (DZ1,25),
      (DZ1,26), (DZ1,27), (DZ1,28), (DZ1,29), (DZ1,30), (DZ1,31), (DZ1,32), (DZ1,33), (DZ1,34), (DZ1,35), (DZ1,36), (DZ1,37), (DZ1,38), (DZ1,39), (DZ1,40),
      (DZ1,41), (DZ1,42), (DZ1,43), (DZ1,44), (DZ1,45), (DZ1,46), (DZ1,47), (DZ1,48), (DZ1,49), (DZ1,50), (DZ1,51), (DZ1,52)
COMMON /PRES/ PT2, P3
      REAL NS, K, KAP
      MAXACT=100 (MAXACT)
      MN=MAXN-1
      IF (INTD) 1, 2, 3
      DO 2 I=1, MAXI
      CTH=CTH(I)
      S1=Q1*CTH(I)
      V1=Q1*CTH
      AMN=(V1/VS1)**2
      IF (AMN.LT.1.0) AMN=1.01
      V(I, MAXN, 1)=GAMA*AMN/(GAMA*AMN+1.0)
      V(I, MAXN, 2)=V1/V(I, MAXN, 1)
      V(I, MAXN, 3)=L1
      ZZ=(GAM1*AMN-GAMA)/GAMA
      V(I, MAXN, 4)=ALOG(ZZ)-GAM1*ALOG(V(I, MAXN, 1))
      ZR=PT2*STH*STH+CTH(I)*CTH(I)
      V(I, 1, 4)=V(I, MAXN, 4)
      V(I, 1, 1)=(ZRN*V(I, 1, 4))**(.1)/GAMA
      V(I, 1, 2)=1.1
      V(I, 1, 3)=Q1
      CTH=1.1
      DO 2 M=2, MN
      CTH=CFT+DFT
      ZR=CFT*DFT
      DO 2 N=1, 4
      V(I, N, M)=V(I, 1, M)+(V(I, MAXN, M)-V(I, 1, M))*ZR
      CFT=CFT+DFT
      DO 6 I=[MIN, MI]
      KAP=1.1-NS(I)*K(I)
      IF (I.GT.1) GO TO 4
      ZR=1.0
      CTH=1.1
      CFT=1.1
      CFT=1.1
      GO TO 5
      CTH=(NS(I-1)-NS(I+1))/(KAP*DZ(I))
      ZZ=ATAN(CTH)

```

```

INITI 1
INITI 2
INITI 3
INITI 4
INITI 5
INITI 6
INITI 7
INITI 8
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INITI 11
INITI 12
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INITI 51
INITI 52

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NOT REPRODUCIBLE

```

COFL=COS(ZZ)
SDFL=TH(I)*COFL
DUM=TH(I)+Z
DRT=COS(DUM)
SDFL=SIN(DUM)
AM=(1-MACH1*SEET)**2
IF (AMN.LT.1.0) AMN=1.0
V(I,MAXN,1)=GAM1*AMN/(GAM1*AMN+1.0)
U1=01*CRST
V2=01*SBET/V(I,MAXN,1)
V(I,MAXN,2)=U1*COFL+V2*SDFL
V(I,MAXN,3)=V2*COFL-U1*SDFL
DUM=(GAM1*AMN-GAM1)/GAM1
V(I,MAXN,4)=ALOG(DUM)-GAM1*ALOG(V(I,MAXN,1))
STH=SIN(TH(I))
DR=DT2*STH*STH+D9)*DTH(I)*DTH(I)
V(I,1,4)=V(I,MAXN,4)
V(I,1,4)=V(I,1,4)
V(I,1,1)=(D1*EXP(-V(I,1,4)))** (1.0/GAM1)
V(I,1,3)=0.0
Z2=ABS(HT1+HT1-GAM1*DR/(GAM1*V(I,1,1)))
V(I,1,2)=SQRT(Z2)
ETA=0.0
DO 6 N=2,MN
ETA=ETA+DFT
Z2=ETA**2
DO 6 M=1,4
5 V(I,N,M)=V(I,1,M)+(V(I,MAXN,M)-V(I,1,M))*Z2
DUM=1.0
IF (IMIN.EQ.2) DUM=1.0
DO 7 N=1,MAXN
V(I,N,2)=DUM*V(I,N,2)
DO 7 M=1,4
7 V(MAYI,N,M)=V(I,N,M)
RETURN
END

```

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INITI 53
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INITI 87
INITI 88

```

11. SUBROUTINE INTAPE

A. PURPOSE

This subroutine loads data and initial conditions from a magnetic tape unit. This type of initialization is used for restarting an unfinished case or for performing additional iterations on a completed case.

B. VARIABLE LIST

Most of the variables in this routine are read from magnetic tape, only. Their definitions can be determined by matching their locations in the various common blocks with the corresponding variables in the main program.

IRD - control variable specifying the type of read operation.

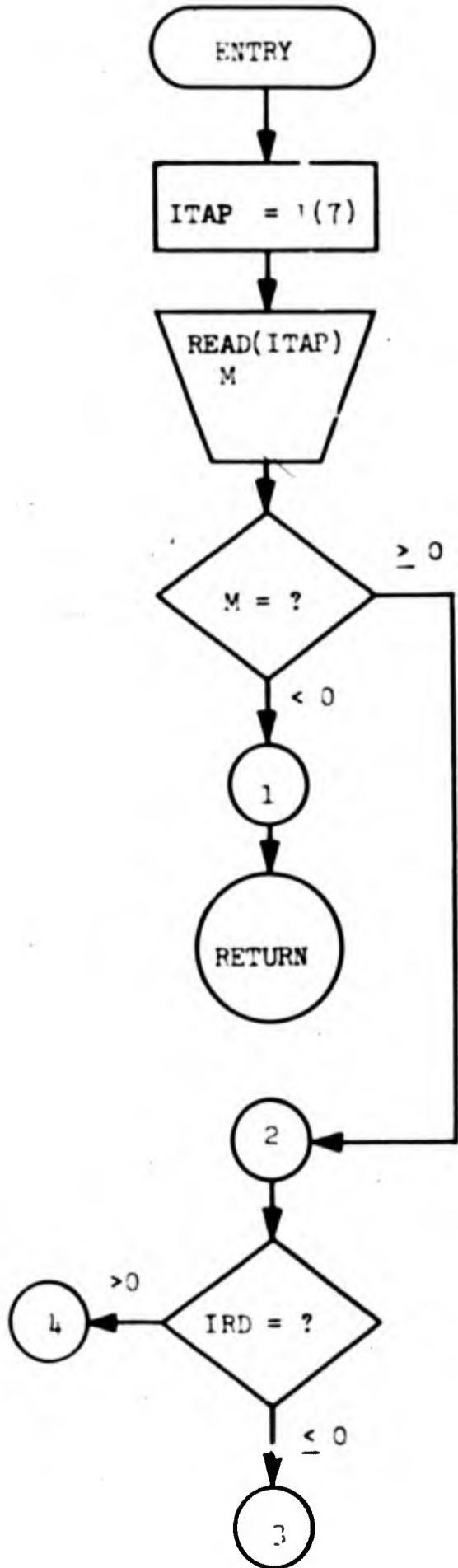
=0 - read data in COMMON/INPUT/.

>0 - read data required for a specific segment.

ITAP - variable used for temporary storage of the input tape unit number.

M - variable indicating the end of a case, when a negative value is read.

C. FLOW CHART



ENTER SUBROUTINE INTAPE.

SAVE THE INPUT TAPE NUMBER.

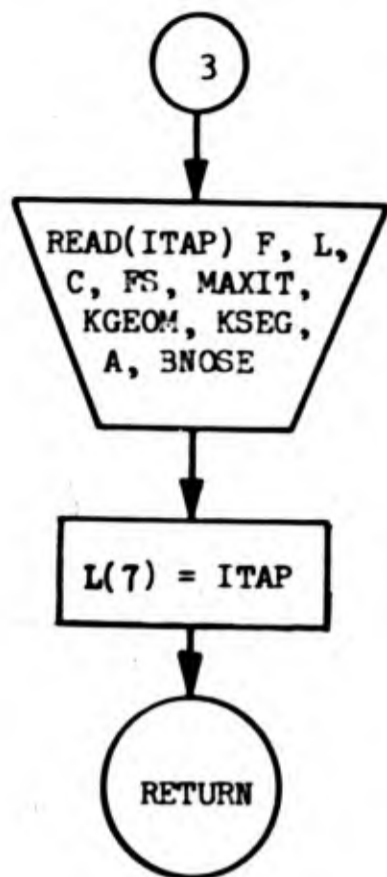
READ END OF CASE INDICATOR

END OF CASE.

RETURN.

THIS FILE CONTAINS DATA
FOR ANOTHER SEGMENT.

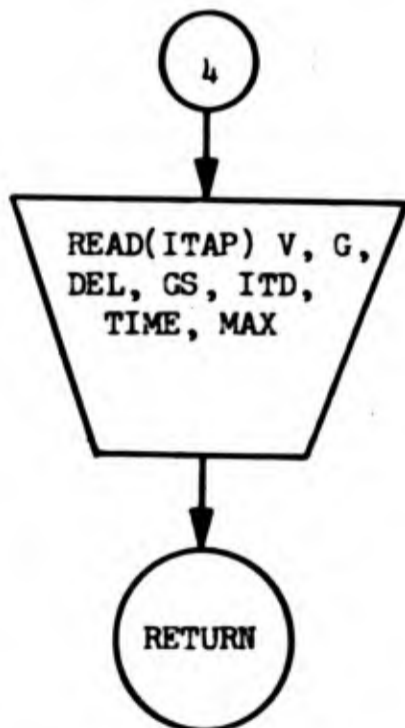
WHAT TYPE OF DATA IS TO
BE READ ?



READ INPUT DATA AND CONTROL
VARIABLES.

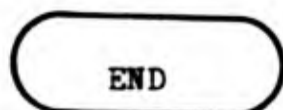
RECOVER INPUT TAPE NUMBER.

RETURN.



READ FLOW FIELD AND
SEGMENT GEOMETRY.

RETURN.



END OF SUBROUTINE INTAPE.

0. LISTING

```

C          SUBROUTINE INTAPE (IRD,M)
C
C          THIS ROUTINE READS THE FLOW FIELD DATA FROM TAPE
C
COMMON /1/ V(15,40,4),O(40,4),P(3,40)
COMMON /2/ G(1,15)
COMMON /3/ MAX(6)
COMMON /4/ IDUM(2),KGFCM,A(6),KSEG
COMMON /5/ DEL(5)
COMMON /6/ FS(3)
COMMON /9/ ISEG,GS(3),RNOSE
COMMON /11/ MAXIT,ISTOP,ITD,TTIME
COMMON /INPUT/ F(13),L(15)
COMMON /CONST/ NUM(2),C(3)
ITAP=L(7)
PFAD (ITAP) M
IF (M) 1,2,2
RETURN
1
2
3
4
IF (IRD) 3,3,4
PFAD (ITAP) F,L,C,FS,MAXIT,KGFCM,KSEG,A,RNOSE
L(7)=ITAP
RETURN
PFAD (ITAP) V,G,OFL,GS,ITD,TIME,MAX
RETURN
END
INTAP  1
INTAP  2
INTAP  3
INTAP  4
INTAP  5
INTAP  6
INTAP  7
INTAP  8
INTAP  9
INTAP 10
INTAP 11
INTAP 12
INTAP 13
INTAP 14
INTAP 15
INTAP 16
INTAP 17
INTAP 18
INTAP 19
INTAP 20
INTAP 21
INTAP 22
INTAP 23
INTAP 24
INTAP 25

```

12. SUBROUTINE INTEG

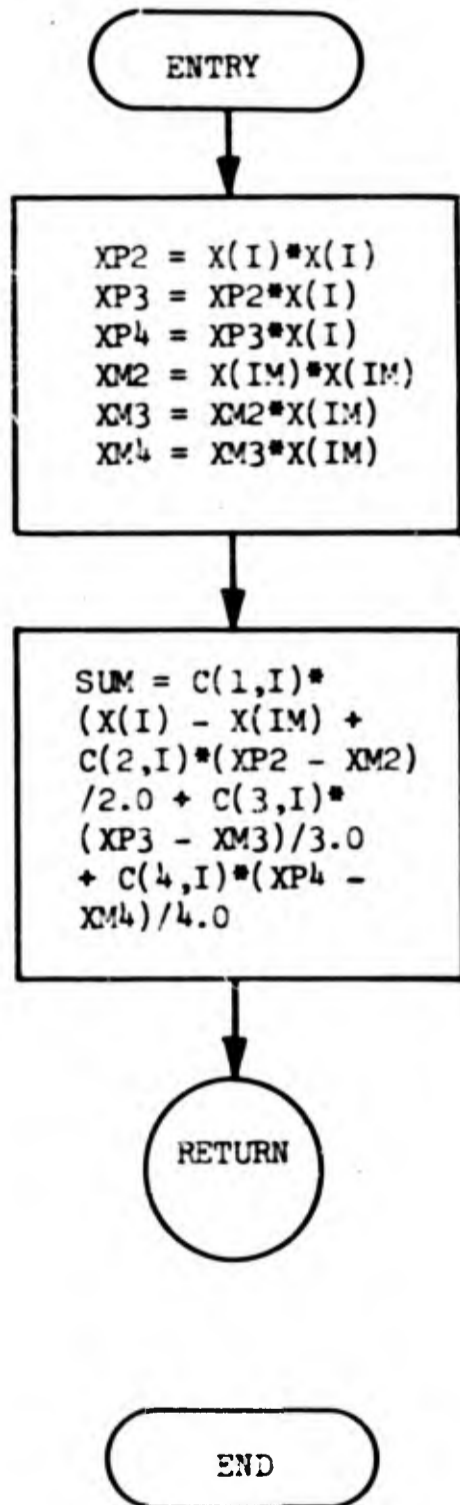
A. PURPOSE

This routine uses the coefficients for a cubic fit generated by subroutine COEFF to integrate the function between 2 adjacent values of the independent variable.

B. VARIABLE LIST

- C - array of cubic coefficients generated in subroutine COEFF.
- I - index of the point in the X array corresponding to the upper limit of integration.
- IM - index of the point in the X array corresponding to the lower limit of integration.
- SUM - value of the integral.
- X - array containing the values of the independent variable.
- XM2 - $[X(IM)]^2$
- XM3 - $[X(IM)]^3$
- XM4 - $[X(IM)]^4$
- XP2 - $[X(I)]^2$
- XP3 - $[X(I)]^3$
- XP4 - $[X(I)]^4$

C. FLOW CHART



ENTER SUBROUTINE INTEG.

COMPUTE X^2 , X^3 AND X^4
AT THE POINTS IM AND I.

INTEGRATE THE FUNCTION
DEFINED BY THE CUBIC
COEFFICIENTS FROM $X(IM)$
TO $X(I)$.

RETURN.

END OF SUBROUTINE INTEG.

P. LISTING

```

C
C
C
C
SUBROUTINE INTEG (IM,I,X,SUM)
C
C   THIS ROUTINE INTEGRATES A FUNCTION USING THE SPLINE COEFFICIENTS
C   FROM COEFF.
C
COMMON /SPLINE/ C(4,15)
DIMENSION X(15)
XP2=X(I)*X(I)
XP3=XP2*X(I)
XP4=XP3*X(I)
XM2=X(IM)*X(IM)
XM3=XM2*X(IM)
XM4=XM3*X(IM)
SUM=C(1,I)*X(I)-X(IM))+C(2,I)*(XP2-XM2)/3.0+C(3,I)*(XP3-XM3)/3.0+
+C(4,I)*(XP4-XM4)/4.0
RETURN
END
INTEG 1
INTEG 2
INTEG 3
INTEG 4
INTEG 5
INTEG 6
INTEG 7
INTEG 8
INTEG 9
INTEG 10
INTEG 11
INTEG 12
INTEG 13
INTEG 14
INTEG 15
INTEG 16
INTEG 17

```

13. SUBROUTINE INTERP

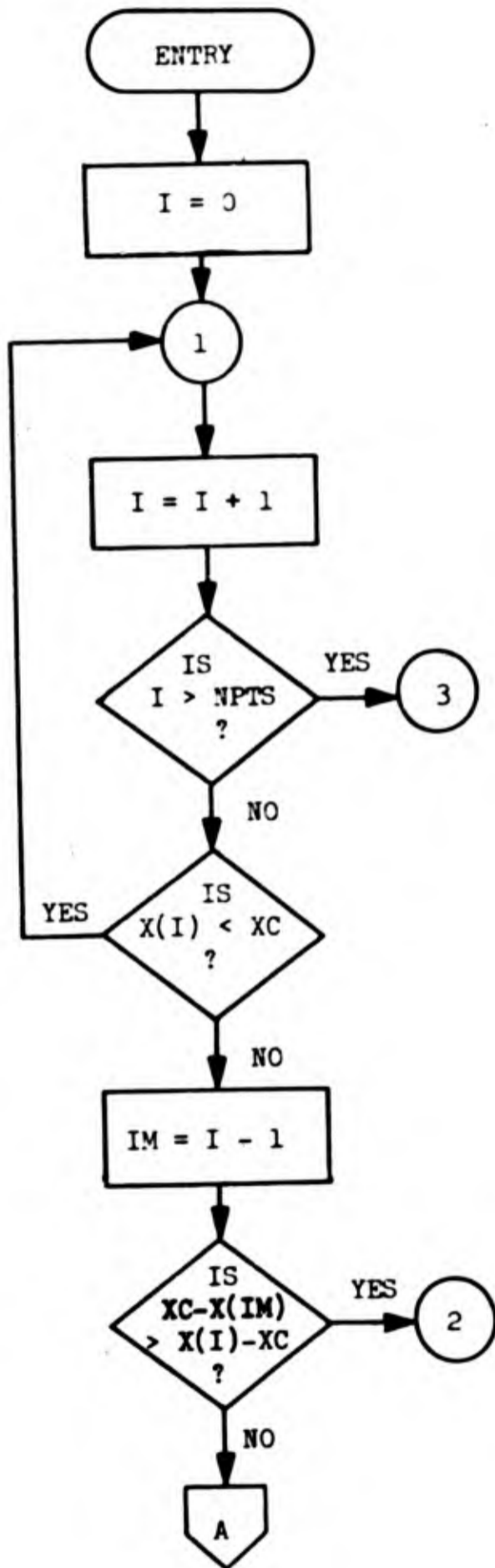
A. PURPOSE

This subroutine interpolates between data points for the values of a function and its derivative, using the cubic fit coefficients generated by subroutine COEFF. An input table of independent variable points is searched to locate the point closest to the desired point of interpolation. If the table is exceeded by a significant amount, an error message is printed and the current case is terminated.

B. VARIABLE LIST

- C - array of cubic fit coefficients.
- DEL - difference between the values of the independent variable at the last 2 points.
- DYDX - $\frac{dy}{dx}$
- I - index specifying the location in the input array, X.
- IM - I - 1.
- ISTOP - error indicator.
- NPTS - number of points in the input array, X.
- X - array containing the input table.
- CX - value of the independent variable at which y and $\frac{dy}{dx}$ are to be determined.
- YC - value of y at XC.

C. FLOW CHART



ENTER SUBROUTINE INTERP.

INITIALIZE I.

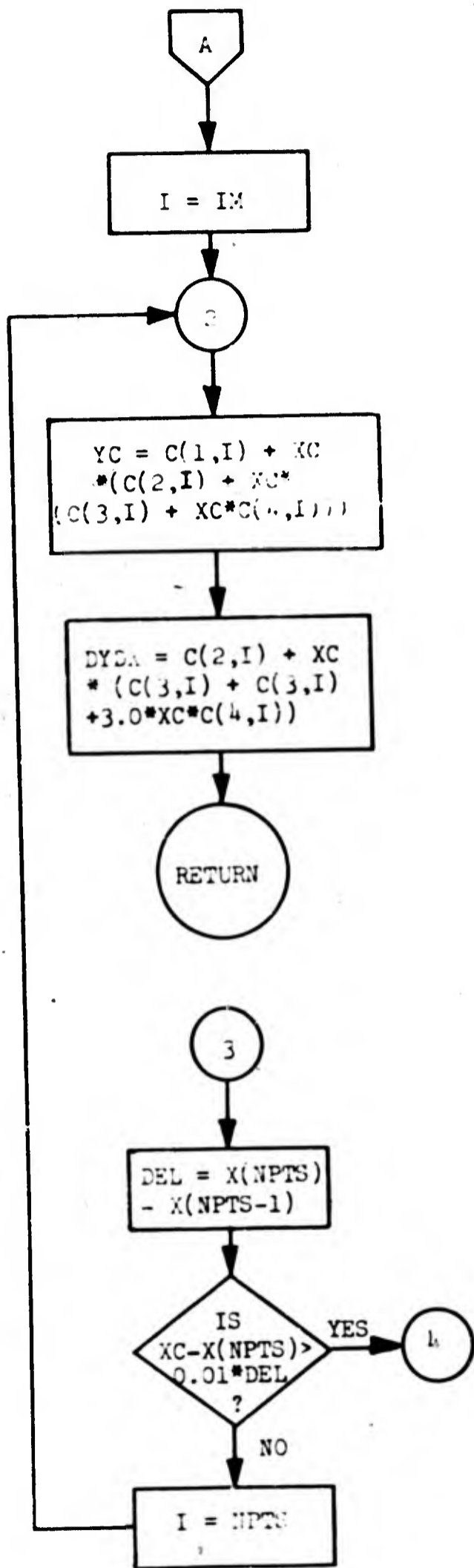
INCREMENT I.

DOES I EXCEED THE
TABLE SIZE ?

IS $X(I) < XC$?

STORE $I - 1$.

IS XC CLOSER TO $X(I)$
THAN TO $X(IM)$?



COMPUTE YC.

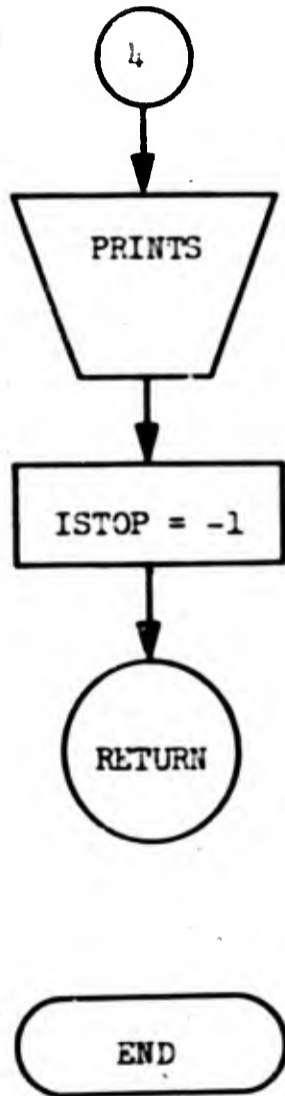
COMPUTE DYDX.

RETURN.

EXCEEDED TABLE SIZE.

CHECK TO SEE IF XC EXCEEDS THE TABLE BY A SIGNIFICANT AMOUNT.

SET I TO USE THE LAST POINT IN THE TABLE.



PRINTS ERROR MESSAGE.

SET ERROR INDICATOR.

RETURN.

END OF SUBROUTINE INTERP.

0. LISTING

```

SUBROUTINE INTERP (X,XC,YC,NPTS)
C
C * THIS ROUTINE INTERPOLATES FOR THE VALUE OF A FUNCTION USING
C THE SPLINE COEFFICIENTS FROM COEFF.
C
COMMON /11/ TDUMMY,ISTOP
COMMON /SPLINE/ C(4,15)
DIMENSION X(15)
I=0
I=I+1
IF (I.GT.NPTS) GO TO 3
IF (X(I).LT.XC) GO TO 1
IM=I-1
IF (XC-X(IM).GT.X(I)-XC) GO TO 2
I=IM
YC=C(1,I)+XC*(C(2,I)+XC*(C(3,I)+XC*(C(4,I))))
DYDX=C(2,I)+XC*(C(3,I)+C(3,I)+3.0*XC*(C(4,I)))
RETURN
DFL=X(NPTS)-X(NPTS-1)
IF (XC-X(NPTS).GT.0.01*DFL) GO TO 4
I=NPTS
GO TO 2
PRINT 5
ISTOP=-1
RETURN
FORMAT (1H1,17X,26HEPPOE IN SUBROUTINE INTERP/17X,19HDATA RANGE EX
*CEEDEN/17X,14HJOB TERMINATED)
END

```

```

INTER 1
INTER 2
INTER 3
INTER 4
INTER 5
INTER 5
INTER 7
INTER 8
INTER 9
INTER 10
INTER 11
INTER 12
INTER 13
INTER 14
INTER 15
INTER 16
INTER 17
INTER 18
INTER 19
INTER 20
INTER 21
INTER 22
INTER 23
INTER 24
INTER 25
INTER 26
INTER 27
INTER 28
INTER 29

```

14. SUBROUTINE LOAD

A. PURPOSE

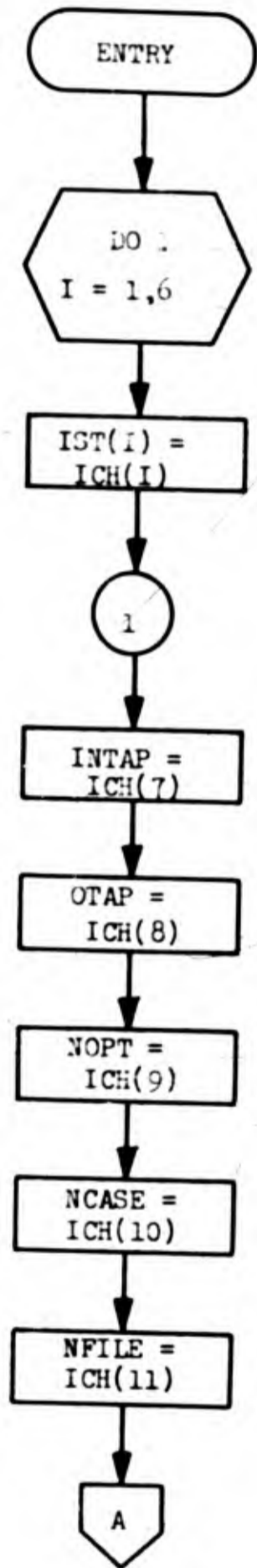
This subroutine loads data from a magnetic tape for use as initial conditions. The principal purpose of this option is to allow additional iterations to be performed on a previous case which did not achieve a steady state. It is also useful when unsteady flows are being considered, allowing the solution to be advanced in time in stages to obtain printed outputs at the desired intervals. Also, if the supplied initial conditions are inadequate for a particular problem, this option will allow different initial conditions to be supplied by the user without requiring alterations of the computer code. Case data variables that may be changed from the values used on the previous case are saved before the read operation and are recovered before returning to the main program. Two error checks are accomplished when this option is used. First, the number of iterations to be performed must have been specified (the steady-state option is not permitted - see section VII). Also, the input and output magnetic tape unit numbers must be different. This requirement has been included to prevent an accidental write operation which could erase data acquired on a previous run, possibly a multibase run which required extensive computer time. When either error is sensed, an error message is printed and the case is terminated. The case and segment numbers where computation should start are supplied in the input data (see section VII).

B. VARIABLE LIST

Variables previously described in the description of the main program have been omitted from this list.

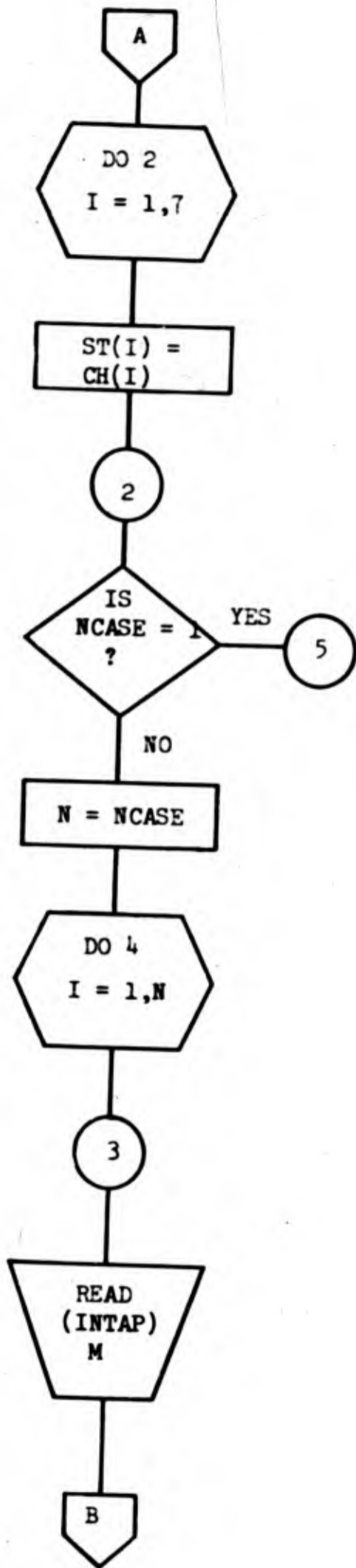
- CH - array containing values of floating point variables in COMMON block /INPUT/ that are to be used rather than the values read from magnetic tape.
- I - DO loop index.
- ICH - same application as CH for integer variables.
- IST - array used for temporary storage of the ICH array.
- M - indicator written on the tape that specified when the last segment has been completed.
- N - variable specifying the number of cases and the number of files to be skipped.
- ST - array used for temporary storage of the CH array.

C. FLOW CHART



ENTER SUBROUTINE LOAD.

SAVE ALL INTEGER CASE DATA
VARIABLES WHICH MAY BE
CHANGED UNDER THIS OPTION



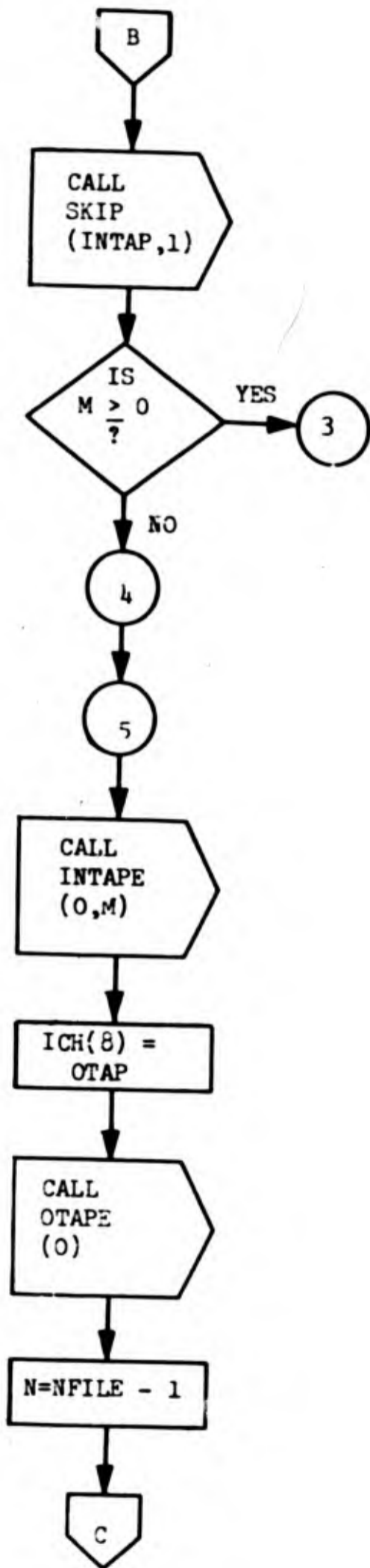
SAVE ALL FLOATING POINT VARIABLES WHICH MAY BE CHANGED UNDER THIS OPTION.

IS THE CASE TO BE LOADED THE FIRST CASE STORED ON THE TAPE ?

COMPUTE THE NUMBER OF CASES TO BE SKIPPED.

SKIP OVER N CASES.

READ VARIABLE USED TO SIGNAL THE END OF A CASE.



SKIP 1 FILE ON THE TAPE.

ARE THERE MORE FILES FOR
THE CURRENT CASE ?

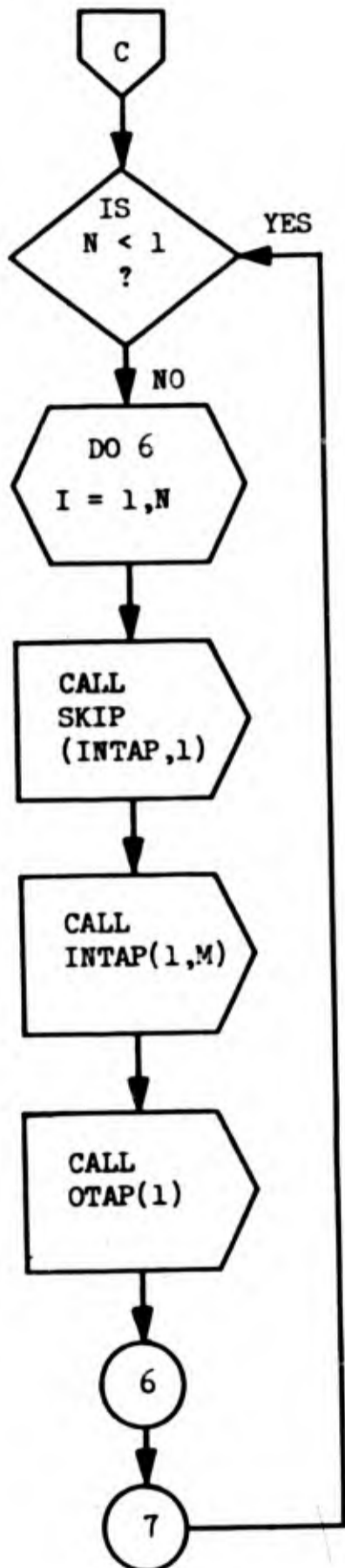
CORRECT CASE HAS BEEN
LOCATED.

READ CASE DATA PREVIOUSLY
USED FOR THE CASE.

RECOVER THE OUTPUT MAGNETIC
TAPE UNIT NUMBER.

WRITE THE CASE DATA ON
MAGNETIC TAPE.

COMPUTE THE NUMBER OF
SEGMENTS TO BE SKIPPED.



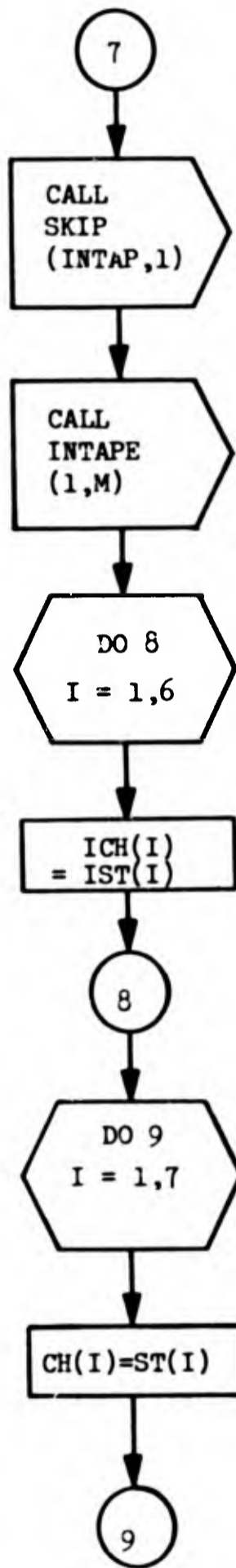
SHOULD COMPUTATIONS START
AT THE FIRST SEGMENT ?

SKIP AN END OF FILE MARK.

READ FLOW FIELD DATA.

STORE DATA ON THE OUTPUT
TAPE.

FIRST SEGMENT TO BE TREATED
HAS BEEN LOCATED.

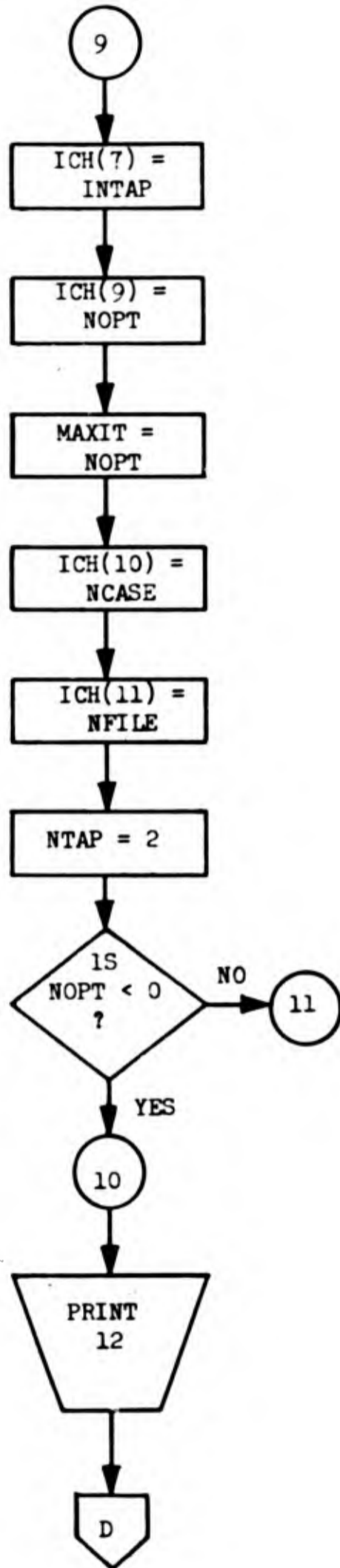


SKIP AN END OF FILE MARK.

READ FLOW FIELD DATA.

RECOVER INTEGER CASE
DATA VARIABLES.

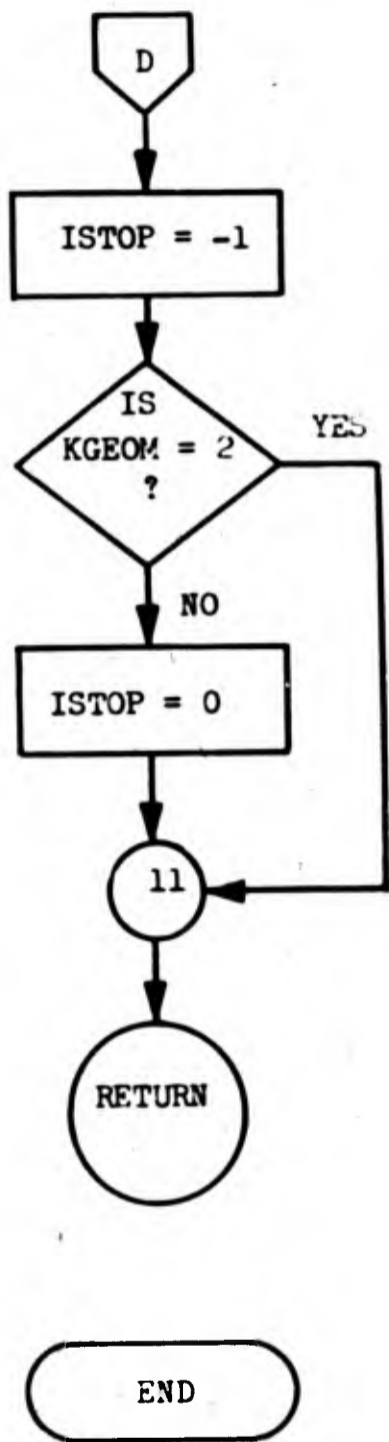
RECOVER FLOATING POINT
CASE DATA VARIABLES.



RECOVER REMAINING CASE
DATA VARIABLES AND SET
REQUIRED CONSTANTS.

ERROR IN CASE DATA ?

PRINT ERROR MESSAGE.



SET ERROR INDICATOR.

ARBITRARY BODY OPTION ?

ALL DATA FOR THIS CASE WAS
READ. CHANGE ERROR
INDICATOR.

RETURN.

END OF SUBROUTINE LOAD.

RESULTS - LOAD

		LOAD	1
		LOAD	2
		LOAD	3
		LOAD	4
		LOAD	5
		LOAD	6
		LOAD	7
		LOAD	8
		LOAD	9
		LOAD	10
		LOAD	11
		LOAD	12
		LOAD	13
		LOAD	14
		LOAD	15
		LOAD	16
		LOAD	17
		LOAD	18
1	DO 1 T=1.6	LOAD	19
	IST(T)=TCH(T)	LOAD	20
	INTAP=ICH(7)	LOAD	21
	INTAP=ICH(8)	LOAD	21
	INTAP=ICH(9)	LOAD	22
	INTAP=ICH(10)	LOAD	23
	INTAP=ICH(11)	LOAD	24
	INTAP=ICH(11)	LOAD	25
2	DO 2 T=1.7	LOAD	26
	ST(T)=CH(T)	LOAD	27
	IF (NOCASE.EQ.1) GO TO 5	LOAD	28
	N=CASE-1	LOAD	29
	DO 4 T=1.8	LOAD	30
3	READ (INTAP) *	LOAD	31
	CALL SKIP (INTAP,1)	LOAD	32
	IF (N.CE.1) GO TO 3	LOAD	33
4	CONTINUE	LOAD	34
5	CALL INTAPE (1,M)	LOAD	35
	ICH(1)=STAP	LOAD	36
	CALL STAPE (1)	LOAD	37
	N=FILE-1	LOAD	38
	IF (N.LT.1) GO TO 7	LOAD	39
	DO 6 T=1.8	LOAD	40
	CALL SKIP (INTAP,1)	LOAD	41
	CALL INTAPE (1,M)	LOAD	42
6	CALL STAPE (1)	LOAD	43
7	CALL SKIP (INTAP,1)	LOAD	44
	CALL INTAPE (1,M)	LOAD	45
	DO 8 T=1.6	LOAD	46
8	ICH(1)=IST(T)	LOAD	47
	DO 9 T=1.7	LOAD	48
9	CH(T)=ST(T)	LOAD	49
	TCH(7)=INTAP	LOAD	50
	TCH(8)=INTAP	LOAD	51
	TCH(9)=INTAP	LOAD	52
	MAXIT=NORT	LOAD	53

	TCH(10)=NCASE	LCAD	53
	TCH(11)=NFILE	LCAD	54
	NTAP=2	LCAD	55
	IF (NCPT) 17,11,11	LCAD	56
10	PRINT 12	LCAD	57
	ISTOP=-1	LCAD	58
11	DUM(5)=CH(1)/(CH(2)*CH(5))	LOAD	59
	RETURN	LCAD	61
C		LCAD	61
12	FORMAT (/30X,19HERROR IN INPUT DATA/30X,25HNTAP=2 REQUIRES NCPT.G/	LCAD	62
	*F.)/30X,15HCASE TERMINATED)	LCAD	63
	END	LCAD	64

15. SUBROUTINE NMESH

A. PURPOSE

This subroutine establishes the nodal spacing across the shock layer. The nodal spacing used is selected according to the following order of priority:

- (1) The value, if any, specified in the case data.
- (2) The value, if any, specified in the body geometry data.
- (3) Empirical relations depending on the flow conditions, the initial guess for shock shape and the body bluntness parameter.

The number of nodes selected is checked against the maximum number allowed by the dimension statements. If too many have been selected, the maximum number allowed is used. If this occurs, it will be indicated by a printed message in the output for the segment. The flow field variables are interpolated from the old nodal geometry to the new nodal geometry using first and second control differences in a truncated Taylor series.

B. VARIABLE LIST

Variables previously defined in the description of the main program have been omitted from this list.

ANS - value of n_s to be used to establish the nodal geometry under option (1).

D - array used for temporary storage of the interpolated flow field variables.

DE - variable used to store the value of $\Delta\eta$ used in the old nodal spacing.

DEL - $\frac{\eta - \eta_0}{2\Delta\eta}$, used for interpolation.

DEL2 - $\frac{1}{2} \left[\frac{\eta - \eta_0}{\Delta\eta} \right]^2$, used for interpolation.

I1 - index of the first streamwise station to be considered.

I2 - index of the last streamwise station to be considered.

K - index for nodes across the shock layer in the old nodal geometry.

KM - K-1

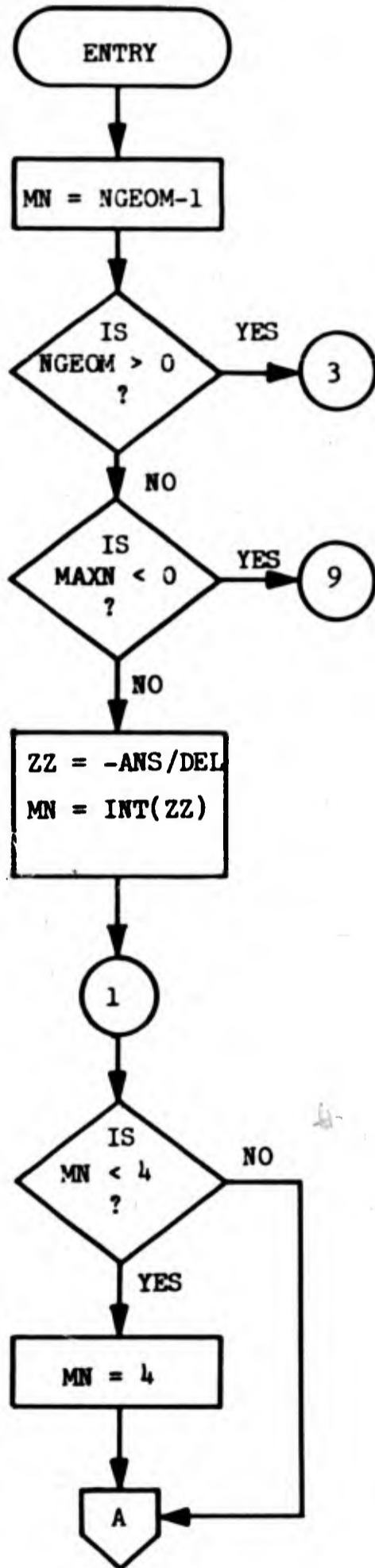
KP - K+1

M - index specifying the proper flow parameters in the D array.

NST - number of nodes across the shock layer used in the old nodal geometry.

- VE - variable containing a term required for the first central difference.
- VE2 - variable containing a term required for the second central difference.
- ZZ - dummy variable used for intermediate calculations.

C. FLOW CHART



ENTER SUBROUTINE NMESH

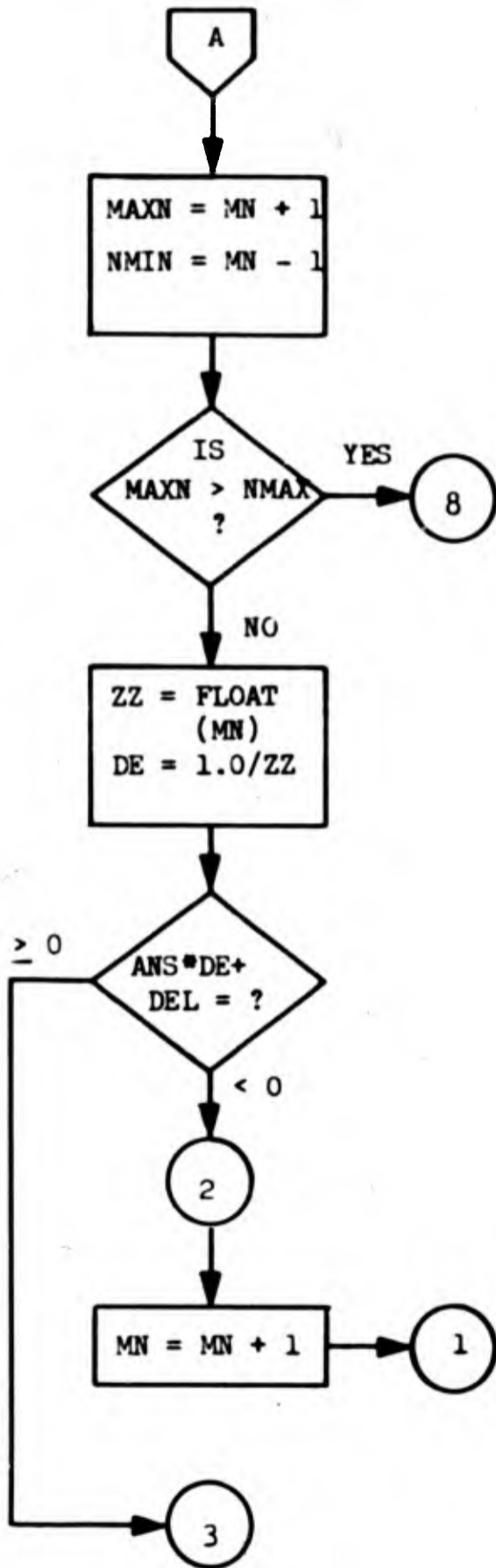
WAS MAXN SPECIFIED IN THE CASE DATA ?

WAS MAXN SPECIFIED WHEN THE BODY GEOMETRY DATA WAS READ ?

ESTIMATE MN.

IS MN < 4 ?

REQUIRE MN > 4.



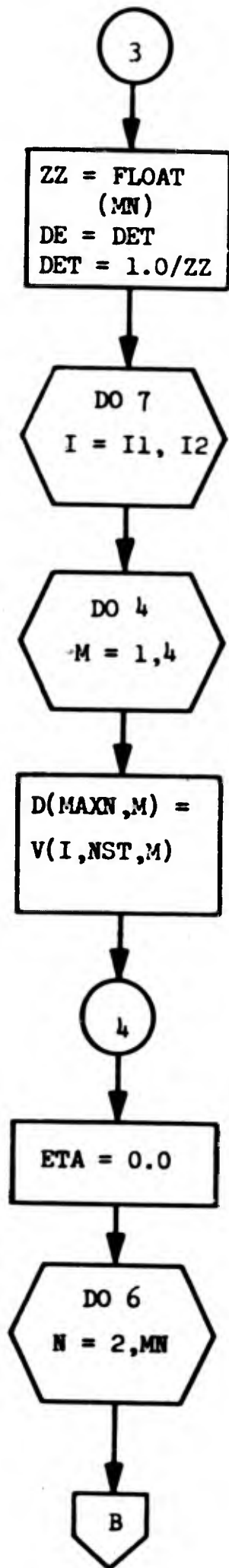
COMPUTE MAXN AND NMIN.

HAVE DIMENSIONS BEEN EXCEEDED ?

COMPUTE A NEW $\Delta\eta$.

IS THIS $\Delta\eta$ SMALL ENOUGH TO PROVIDE THE REQUIRED $\Delta\eta$?

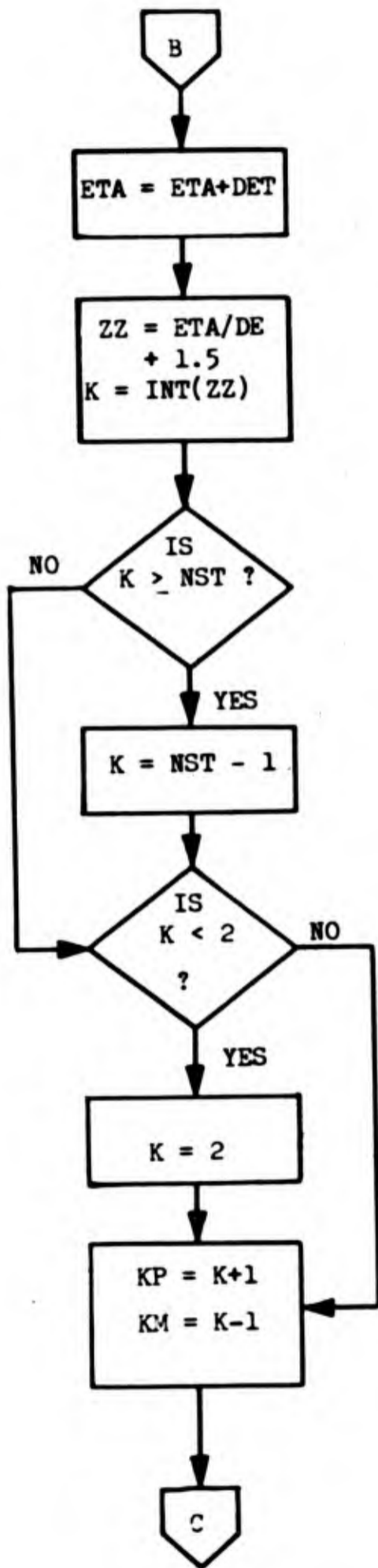
INCREMENT MN.



SAVE THE RESIDENT VALUE OF $\Delta\eta$ AND COMPUTE THE NEW VALUE.

STORE SHOCK CONDITIONS AT THE NEW $N = \text{MAXN}$ LOCATION.

INITIALIZE η .



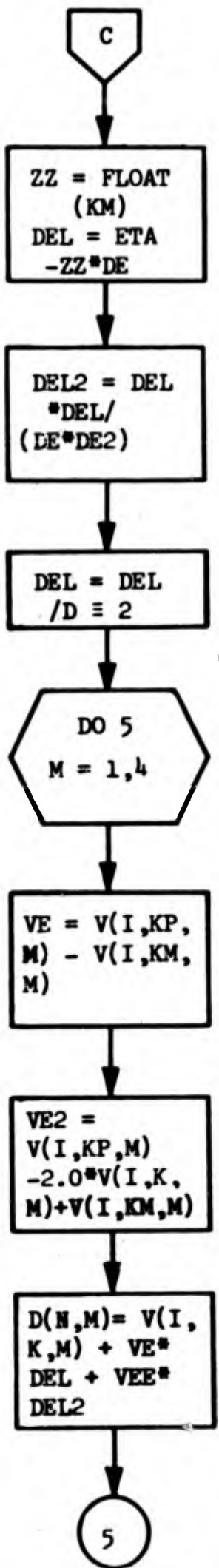
INCREMENT η .

ESTIMATE THE VALUE OF THE INDEX OF THE NODE IN THE RESIDENT NODAL GEOMETRY CLOSEST TO THE REQUIRED VALUE OF η .

IS $K \geq NST$?

IS $K < 2$?

COMPUTE KP AND KM .



COMPUTE $\eta_N - \eta_K$

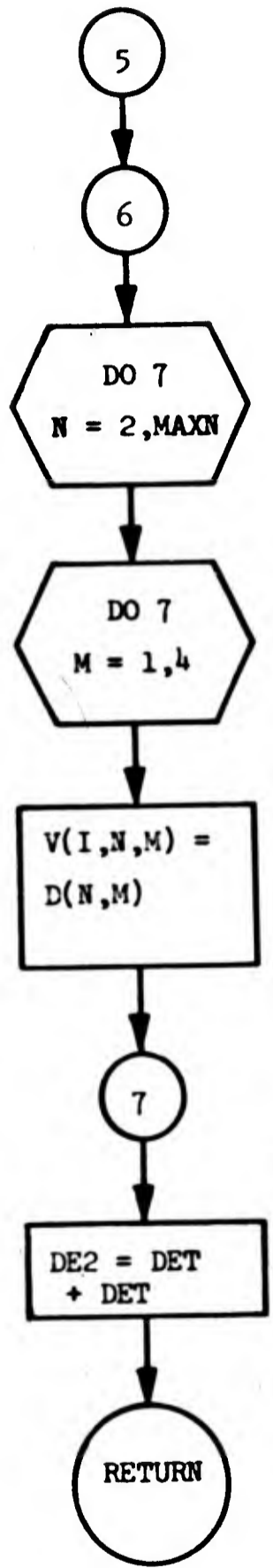
COMPUTE $\frac{(\eta_N - \eta_K)^2}{2 (\Delta\eta)^2}$

COMPUTE $\frac{\eta_N - \eta_K}{2\Delta\eta}$

COMPUTE $f(\eta+\Delta\eta) - f(\eta-\Delta\eta)$

COMPUTE $f(\eta+\Delta\eta) - 2f(\eta) + f(\eta-\Delta\eta)$

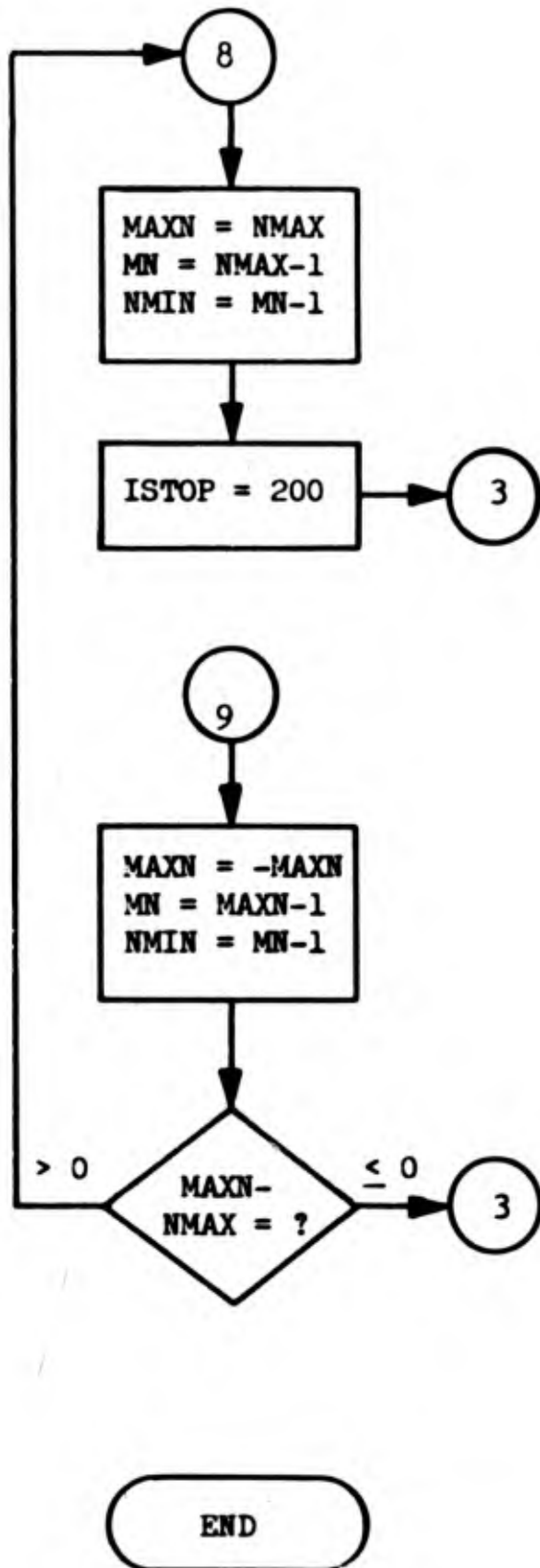
$f = f_0 + f_\eta \Delta\eta + \frac{1}{2} f_{\eta\eta} (\Delta\eta)^2$



TRANSFER ρ , U, V AND S FROM
TEMPORARY TO PERMANENT STORAGE.

COMPUTE $2\Delta\eta$.

RETURN.



DIMENSIONS EXCEEDED.

USE MAXIMUM NUMBER OF
NODES AVAILABLE.

SET ERROR INDICATOR TO
PROVIDE AN ERROR MESSAGE
DURING PRINTOUT.

MAXN WAS SPECIFIED WHEN
BODY GEOMETRY WAS READ.

USE THE VALUE READ.

EXCEEDED DIMENSIONS ?

END OF SUBROUTINE NMESH.

7. LISTING

	SUBROUTINE NPESH (ANS,I1,I2,DEL,NST)	NPESH 1
C		NPESH 2
C	THIS ROUTINE COMPUTES THE PROPER NUMBER OF NCDFS IN THE NORMAL	NPESH 3
C	DIRECTION TO BE EMPLOYED BASED ON DATA FROM RESET.	NPESH 4
C		NPESH 5
	COMMON /1/ V(15,4),D(4,4),P(3,4)	NPESH 6
	COMMON /3/ MAXI,MI,IMIN,MAXN,MN,NMIN	NPESH 7
	COMMON /5/ DXI,DEI,DX2,DE2,DT	NPESH 8
	COMMON /11/ MAXIT,ISTOP,ITD,TIME	NPESH 9
	COMMON /INPUT/ A(13),IV(13),NGEOM	NPESH 10
	COMMON /MESH/ IMAX,NMAX	NPESH 11
	MN=NGEOM-1	NPESH 12
	IF (NGEOM.GT.0) GO TO 3	NPESH 13
	IF (MAXN.LT.0) GO TO 9	NPESH 14
	ZZ=-ANS/DEL	NPESH 15
	MN=INT(ZZ)	NPESH 16
1	IF (MN.LT.4) MN=4	NPESH 17
	MAXN=MN+1	NPESH 18
	NMIN=MN-1	NPESH 19
	IF (MAXN.GT.NMAX) GO TO 8	NPESH 20
	ZZ=FLOAT(MN)	NPESH 21
	DE=1.0/ZZ	NPESH 22
	IF (ANS*DE+DEL) 2,3,3	NPESH 23
2	MN=MN+1	NPESH 24
	GO TO 1	NPESH 25
3	ZZ=FLOAT(MN)	NPESH 26
	DE=DEI	NPESH 27
	DEI=1.0/ZZ	NPESH 28
	DO 7 I=I1,I2	NPESH 29
	DO 4 M=1,4	NPESH 30
4	D(MAXN,M)=V(I,NST,M)	NPESH 31
	ETA=T.M	NPESH 32
	DO 6 N=2,MN	NPESH 33
	ETA=ETA+DEI	NPESH 34
	ZZ=ETA/DE+1.5	NPESH 35
	K=INT(ZZ)	NPESH 36
	IF (K.GE.NST) K=NST-1	NPESH 37
	IF (K.LT.2) K=2	NPESH 38
	KP=K+1	NPESH 39
	KM=K-1	NPESH 40
	ZZ=FLOAT(KM)	NPESH 41
	DEL=ETA-7Z*DE	NPESH 42
	DFL2=DEL*DEL/(DE*DE2)	NPESH 43
	DFL=DFL2/DF2	NPESH 44
	DO 5 M=1,4	NPESH 45
	VE=V(I,KP,M)-V(I,KM,M)	NPESH 46
	VE2=V(I,KP,M)-2.0*V(I,K,M)+V(I,KM,M)	NPESH 47
5	D(N,M)=V(I,K,M)+VE*DEL+VE2*DFL2	NPESH 48
6	CONTINUE	NPESH 49
	DO 7 N=2,NMAX	NPESH 50
	DO 7 M=1,4	NPESH 51
7	V(I,N,M)=D(N,M)	NPESH 52

1 MAY-1954
2 MAY-1954
3 MAY-1954
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24 MAY-1954
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26 MAY-1954
27 MAY-1954
28 MAY-1954
29 MAY-1954
30 MAY-1954
31 MAY-1954
END

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NYFCH 62
NYFCH 63
NYFCH 64

16. SUBROUTINE OTAPE

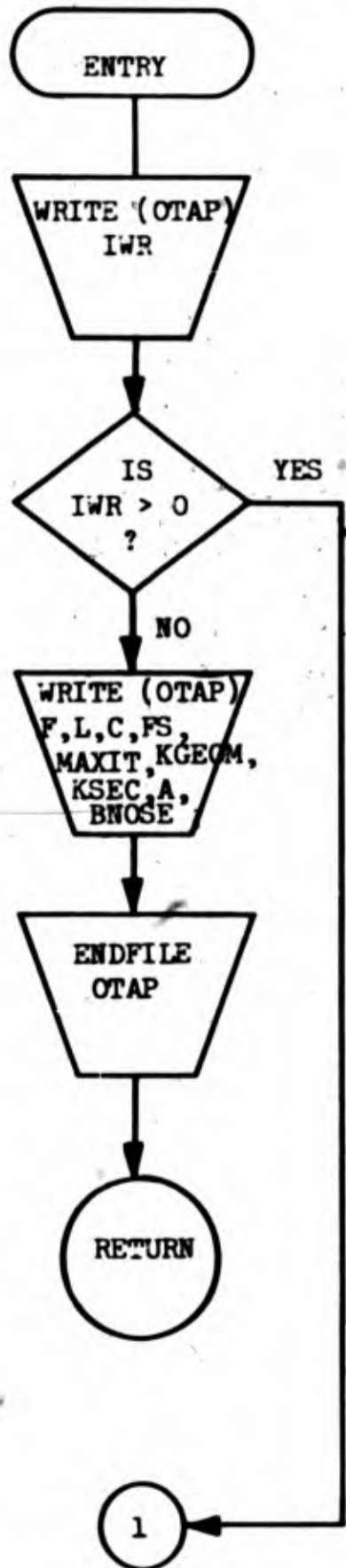
A. PURPOSE

This subroutine performs all output operations to the magnetic tape unit specified for storing the flow field data. By isolating the READ and WRITE statements in subroutines, access to this logic from several locations in the computer code is possible, and it is easily assured the binary read and write operations are consistent. A control variable, IWR, specifies whether case data or flow field data are to be written on tape. This argument should not be negative since it is written on the tape. A negative value would be interpreted as the end of a case if the data were later recovered from tape.

B. VARIABLE LIST

Most variables in this routine are written on tape only. Their definitions can be found in the description of the main program by matching locations in the various COMMON blocks. The control variable, IWR, has been described above. OTAP is the magnetic tape unit number on which the data is to be stored.

C. FLOW CHART



ENTER SUBROUTINE OTAPE.

WRITE INDICATOR ON TAPE.

IS THIS ENTRY TO WRITE
FLOW FIELD DATA FOR A
SEGMENT ?

WRITE CASE DATA AND
CONSTANTS.

ENDFILE OTAP.

RETURN.



WRITE FLOW FIELD DATA
FOR THIS SEGMENT.

ENDFILE OTAP.

RETURN.

END OF SUBROUTINE OTAPE.

D. LISTING

```
C
C
C
      SURROUTINE OTAPE (IWR)
      THIS ROUTINE WRITES THE FLOW DATA ON TAPE.
      COMMON /1/ V(15,40,4), N(40,4), P(3,40)
      COMMON /2/ G(133)
      COMMON /3/ MAX(6)
      COMMON /4/ ICDM(2), KGEOM, A(6), KSEG
      COMMON /5/ NFL(3)
      COMMON /6/ FS(3)
      COMMON /9/ ISEG, GS(3), RNOSE
      COMMON /11/ MAXIT, ISTOP, ITD, TIME
      COMMON /INPUT/ F(13), L(15)
      COMMON /CONST/ DUM(2), C(3)
      INTEGER OTAP
      EQUIVALENCE (L(8), OTAP)
      WRITE (OTAP) IWR
      IF (IWR.GT.1) GO TO 1
      WRITE (OTAP) F, L, C, FS, MAXIT, KGEOM, KSEG, A, RNOSE
      END FILE OTAP
      RETURN
      WRITE (OTAP) V, G, DEL, GS, ITD, TIME, MAX
      END FILE OTAP
      RETURN
      END
```

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CTAPE 20
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CTAPE 23
CTAPE 24
CTAPE 25
```

17. SUBROUTINE RESET

A. PURPOSE

This subroutine schedules the calculation of initial conditions and body geometry parameters for all segments except the first. This data may be generated by calls to other routines or read from magnetic tape, depending on the option specified in the input data (see section VII). A maximum value of $\Delta\eta$ is computed for use in subroutine NMESH. It is based on the signs of the pressure gradient and the normal velocity component of the first node above the body on the previous segment. If the number of nodes across the shock layer has been specified in the input data, this quantity will be ignored.

B. VARIABLE LIST

Variables previously defined in the description of the main program have been omitted from this list.

DUM - array containing the value of ξ at the starting line of a segment and the segment length, measured along the body.

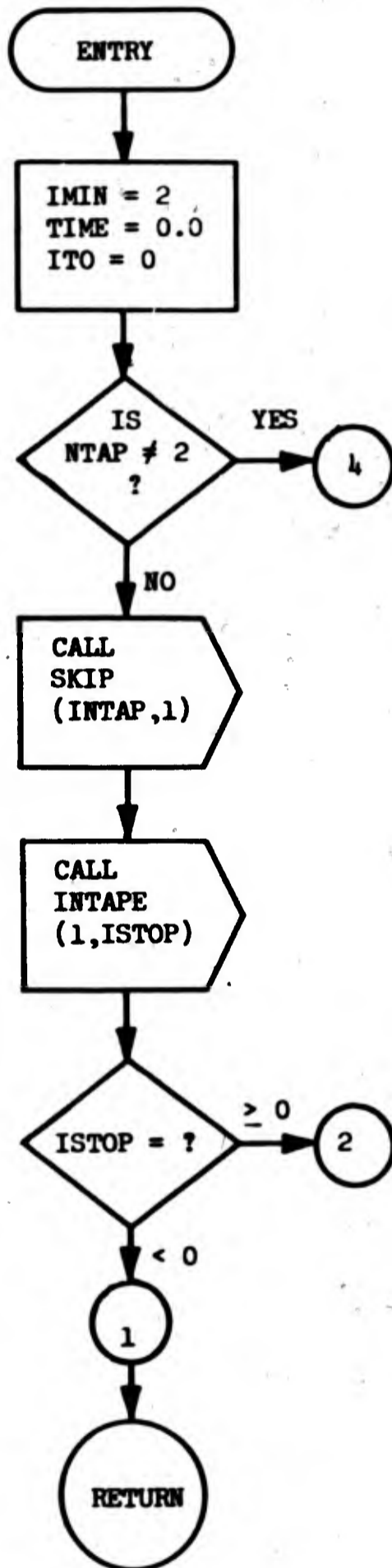
IST - index of the last streamwise node on the previous segment.

M - dummy index.

NST - number of nodes across the shock layer used on the previous segment.

ZZ - dummy variable used for intermediate computations.

C. FLOW CHART



ENTER SUBROUTINE RESET.

INITIALIZE VARIABLES FOR
A NEW SEGMENT.

NEW CASE ?

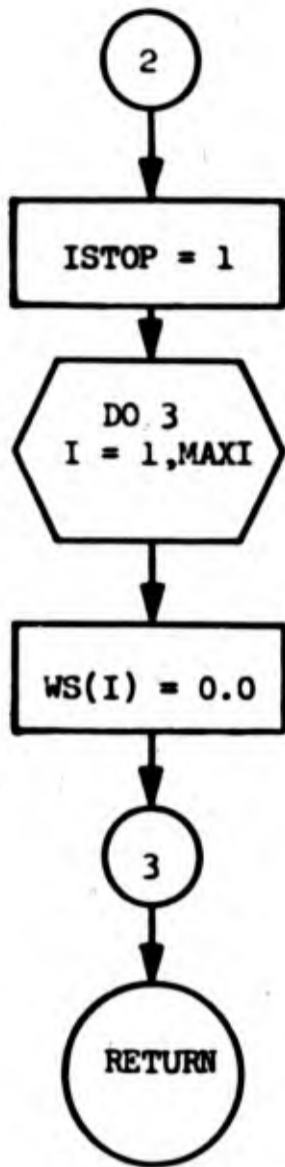
SKIP AN END OF FILE MARK.

READ FLOW FIELD DATA
FROM TAPE.

END OF CASE ?

YES.

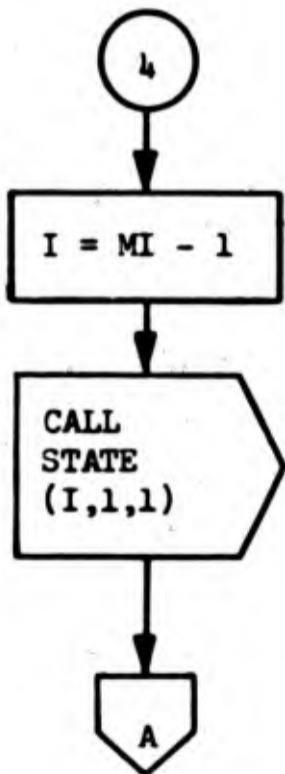
RETURN.



SET ERROR INDICATOR.

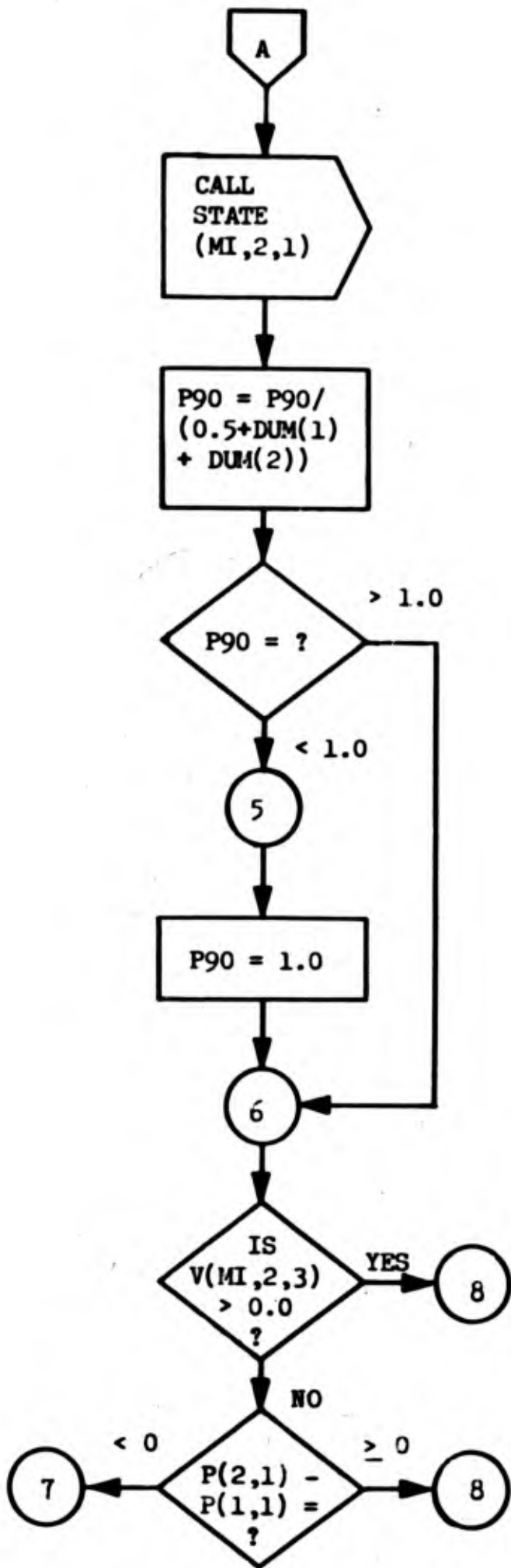
SET SHOCK VELOCITY TO ZERO.

RETURN.



NEW CASE.

COMPUTE BODY PRESSURE
AT MI - 1.



COMPUTE BODY PRESSURE AT MI.

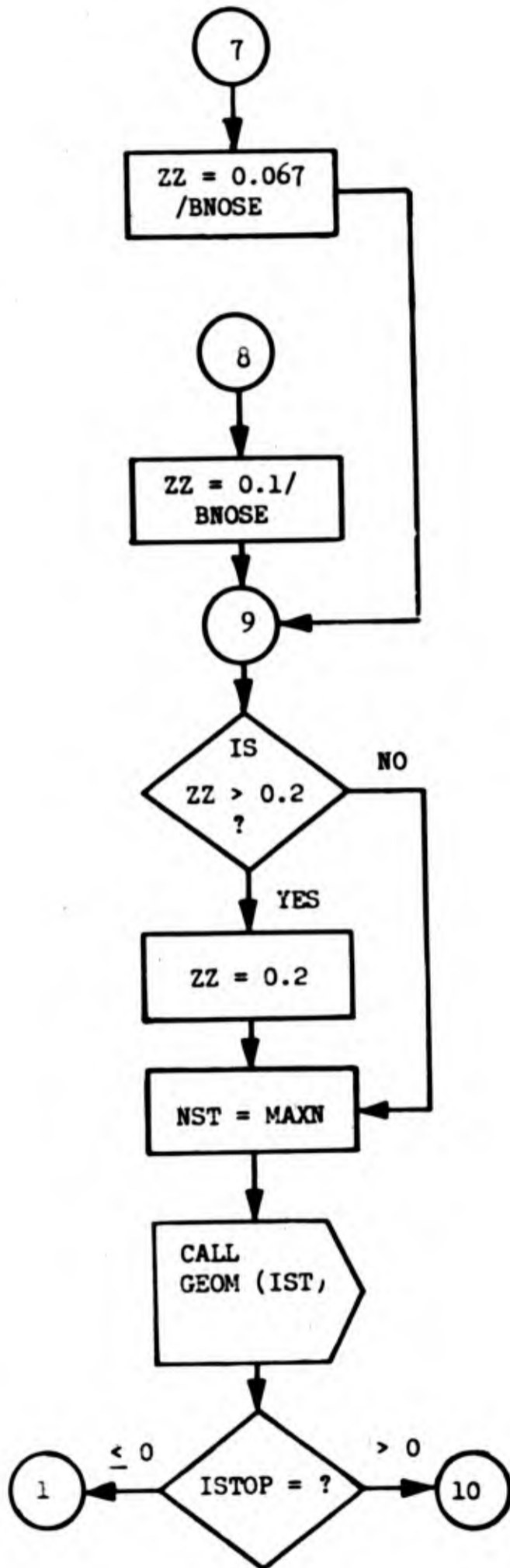
COMPUTE P90.

IS P90 < 1.0 ?

REQUIRE P90 ≥ 1.0

IS V POSITIVE AT THE FIRST NODE ABOVE THE BODY AT THE END OF THE LAST SEGMENT ?

WHAT IS THE SIGN OF THE PRESSURE GRADIENT ?



FINE NODAL SPACING
REQUIRED.

COARSE NODAL SPACING
ALLOWABLE.

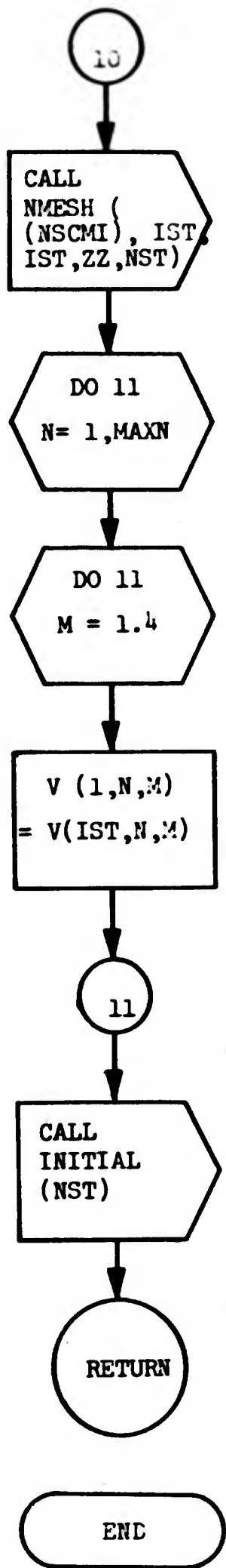
WILL $\Delta\eta$ BE > 0.2 ?

REQUIRE $\Delta\eta \leq 0.2$

SAVE THE VALUE OF
MAXN

COMPUTE BODY GEOMETRY
FOR THE NEXT SEGMENT.

ERROR SENSED IN
SUBROUTINE GEOM ?



NO ERRORS.

ESTABLISH NODAL SPACING
ACROSS THE SHOCK LAYER.

TRANSFER THE VALUES OF THE FLOW
PARAMETERS AT THE END OF THE
PREVIOUS SEGMENT TO THE I = 1
LOCATION IN THE ARRAY, V.

COMPUTE INITIAL CONDITIONS FOR
OTHER NODES.

RETURN.

END OF SUBROUTINE RESET.

1. LISTING

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SUBROUTINE RESET
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18. SUBROUTINE RESTART

A. PURPOSE

This subroutine recovers data stored on magnetic tape to restart an unfinished case. Since it is difficult to estimate the computation time required for a particular problem, this capability is necessary. Should the job be terminated prior to completion, computations can be resumed with very little loss of computer time. The magnetic tape is searched for the correct case, the case data input is read and files are skipped until the flow field data for the last completed segment is read. Computations resume for the next segment.

B. VARIABLE LIST

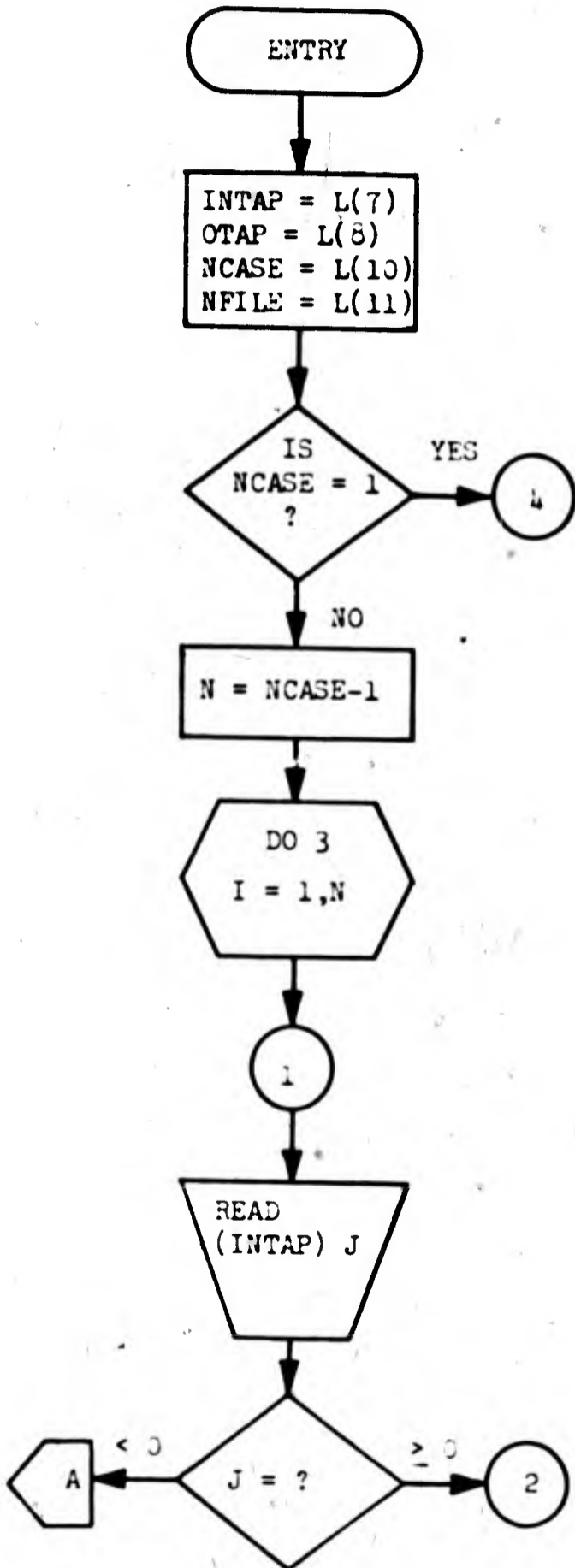
Variables previously defined in the description of the main program have been omitted from this list.

J - variable read from magnetic tape. A negative value indicates the end of a case.

L - array containing all integer variables in COMMON block /INPUT/.

M - dummy variable in a subroutine call statement.

C. FLOW CHART



ENTER SUBROUTINE RESTART.

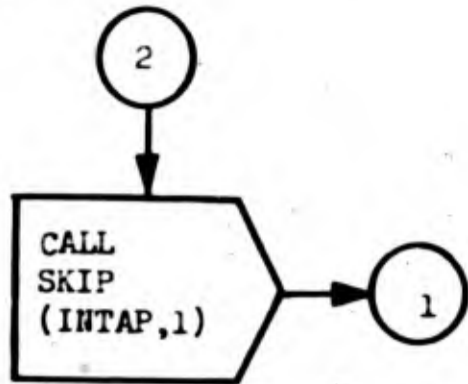
SAVE CASE DATA VARIABLES
WHICH ARE TO BE CHANGED
FROM THEIR VALUES ON TAPE.

IS THE UNFINISHED CASE
THE FIRST ONE STORED
ON THE TAPE ?

DO LOOP TO SKIP ALL CASES
PRIOR TO THE UNFINISHED
CASE.

READ END OF CASE INDICATOR.

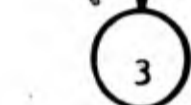
WAS THIS THE LAST FILE
OF THE CASE ?



SKIP END OF FILE MARK
AND RETURN TO 1 TO READ
ADDITIONAL FILES.



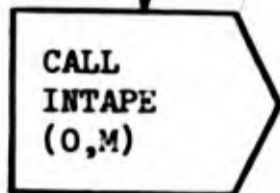
LAST FILE IN THIS CASE.



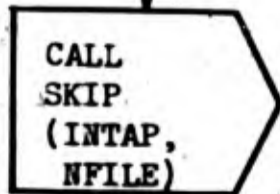
SKIP END OF FILE MARK.



CORRECT CASE HAS BEEN
LOCATED.

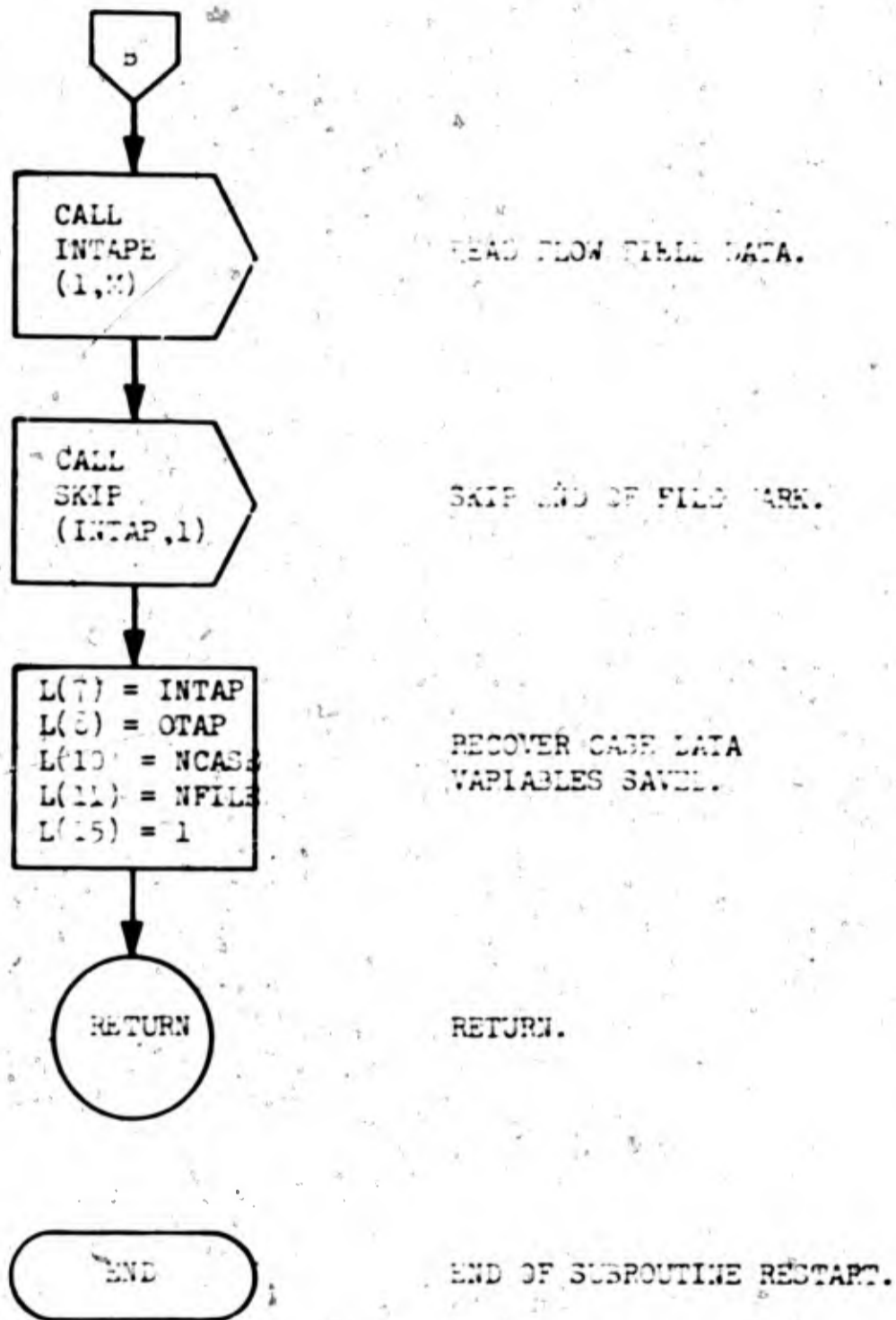


READ CASE DATA AND
CONSTANTS FROM TAPE.



SKIP FILES TO THE
ONE CONTAINING THE
LAST SEGMENT COMPLETED.





D. LISTING

```

1 SURROUTINE RESTART
2 THIS ROUTINE READS THE DATA FROM TAPE TO RESTART AN
3 UNFINISHED CASE.
4
5 COMMON /1/ V(15,4),D(40,4),P(3,40)
6 COMMON /2/ G(105)
7 COMMON /3/ MAX(6)
8 COMMON /4/ ICUM(2),KGFOM,A(6),KSEG
9 COMMON /5/ DEL(5)
10 COMMON /6/ FS(3)
11 COMMON /9/ ISEG,DUM(4)
12 COMMON /11/ MAXIT,ISTOP,ITD,TIME
13 COMMON /INPUT/ F(13),L(15)
14 COMMON /CONST/ DUMM(5)
15 INTEGER NTAP
16 INTAP=L(7)
17 GTAP=L(8)
18 NCASE=L(10)
19 NFILE=L(11)
20 IF (NCASE.EQ.1) GO TO 4
21 N=NCASE-1
22 DO 3 I=1,N
23 READ (INTAP) J
24 IF (J) 2,2
25 CALL SKIP (INTAP,1)
26 GO TO 1
27 CALL SKIP (INTAP,1)
28 CALL INTAPE (0,M)
29 CALL SKIP (INTAP,NFILE)
30 CALL INTAPE (1,M)
31 CALL SKIP (INTAP,1)
32 L(7)=TNTAP
33 L(8)=NTAP
34 L(10)=NCASE
35 L(11)=NFILE
36 L(15)=1
37 RETURN
38 END

```

RESTART 1
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 RESTART 39

19. SUBROUTINE SHOCK

A. PURPOSE

This subroutine performs all computations for nodes at the shock using the method presented in section IV.

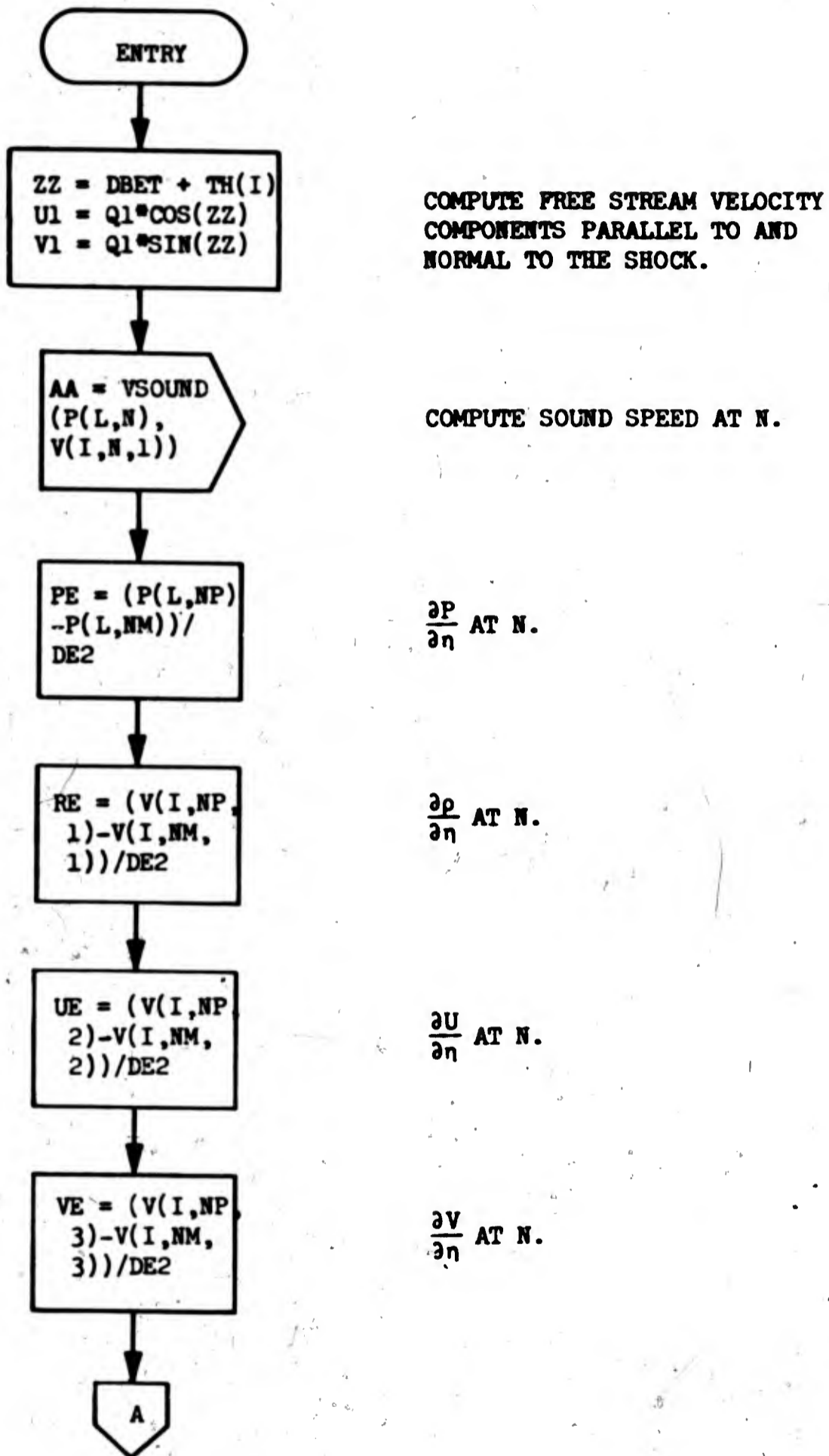
B. VARIABLE LIST

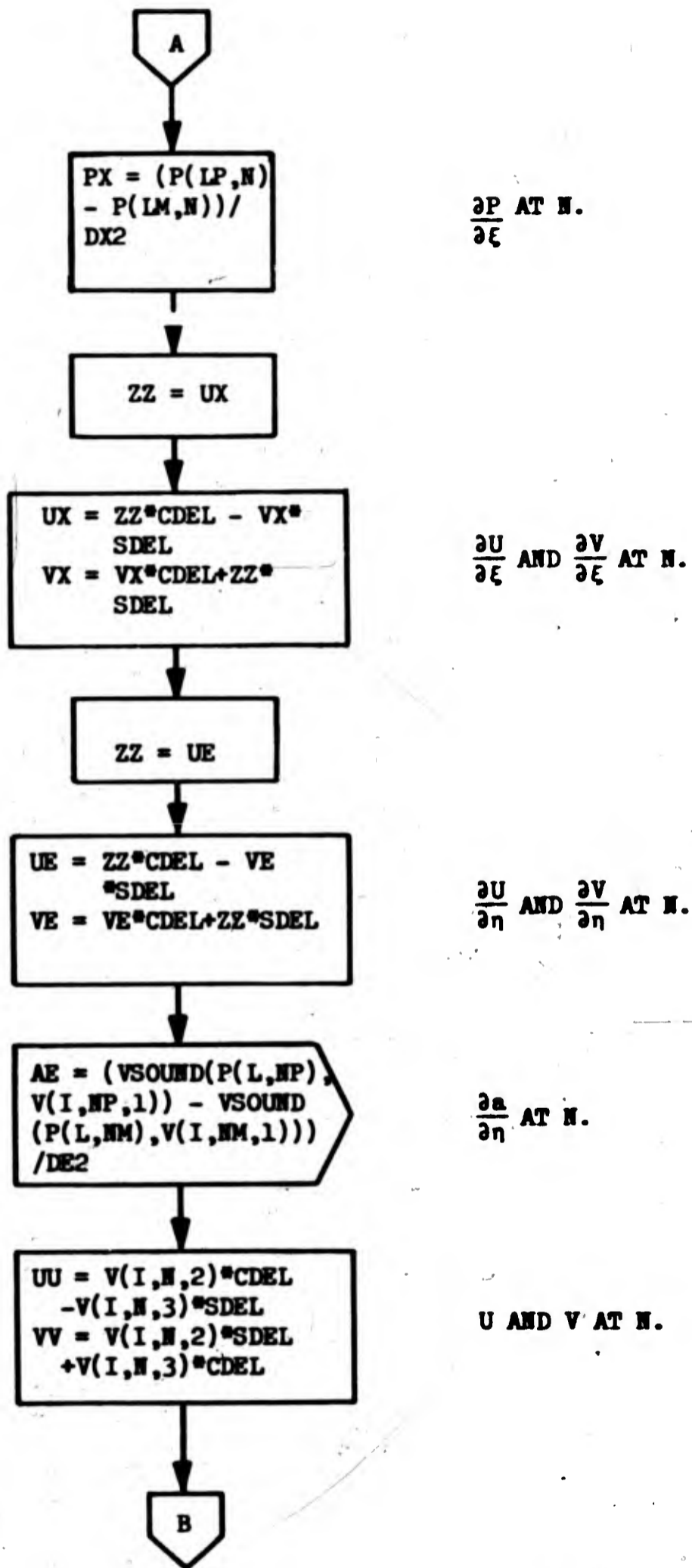
Variables previously defined in the description of the main program have been omitted from this list.

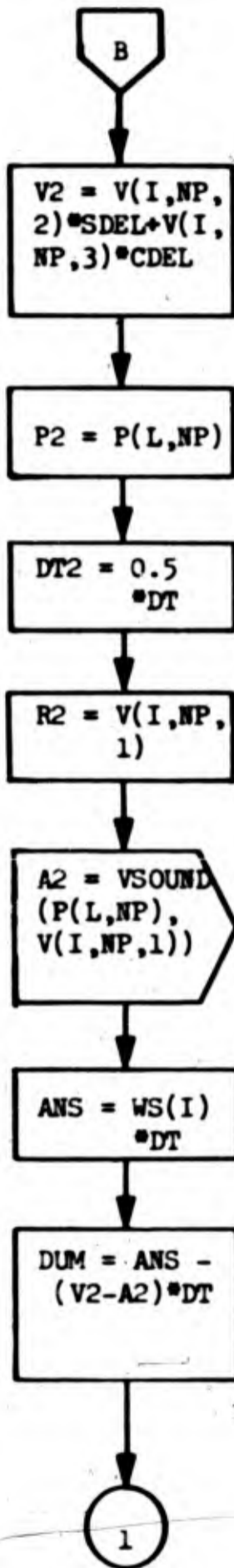
- A - weighting function, equation (45).
- AA - sound speed at the node defined by N.
- AE - $\frac{\partial a}{\partial \xi}$ at the node defined by N.
- ANS - $w_s \Delta r$
- AS - average of the sound speeds at points 2 and 4 (Fig. 4).
- ATSTAG - Assigned a value of 2.0 at the symmetry axis, and a value of 1.0 at all other points.
- AX - $\frac{\partial a}{\partial \xi}$ at the node defined by N.
- A1 - value of A at the shock.
- A2 - sound speed at point 2 (Fig. 4).
- A4 - sound speed at point 4 (Fig. 4).
- B - weighting function, equation (46).
- B1 - value of B at the shock.
- DE - difference in η between the node defined by N and other points at which interpolation is required.
- DS - distance parallel to the shock, used to compute A and B.
- DSO - intermediate value of DS, used to ensure that DS is non-zero.
- DTZ - $\Delta r/2$
- DUM - dummy variable used for intermediate computations.
- DX - difference in ξ between the node defined by N and points at which interpolation is required.
- N - NP-1
- NP - number of nodes across the shock layer. Identifies the node at the shock.
- NM - NP-2
- NPLUS - distance between points 2 and 4 (Fig. 4).
- PE - $\frac{\partial p}{\partial \eta}$ at the node defined by N.
- PX - $\frac{\partial p}{\partial \xi}$ at the node defined by N.

- P2 - pressure at point 2 (Fig. 4).
- P3 - intermediate value of P2.
- P4 - pressure at point 4 (Fig. 4).
- RE - $\frac{\partial \rho}{\partial \eta}$ at the node defined by N.
- RHO - average of the values of ρ at points 2 and 4 (Fig. 4).
- RX - $\frac{\partial \rho}{\partial \xi}$ at the node defined by N.
- RO - value of ρ used to compute A.
- UE - $\frac{\partial U}{\partial \eta}$ at the node defined by N.
- UU - U at the node defined by N.
- UX - $\frac{\partial U}{\partial \xi}$ at the node defined by N.
- UO - value of U used to compute A.
- U1 - free-stream value of U.
- U4 - value of U at point 4 (figure 4).
- VE - $\frac{\partial V}{\partial \eta}$ at the node defined by N.
- VV - value of V at the node defined by N.
- VX - $\frac{\partial V}{\partial \xi}$ at the node defined by N.
- VO - value of V used to compute B.
- V1 - free-stream value of V.
- V2 - value of V at point 2 (Fig. 4).
- V4 - value of V at point 4 (Fig. 4).
- WSH - local shock velocity.
- ZZ - dummy variable used for intermediate computations.

C. FLOW CHART







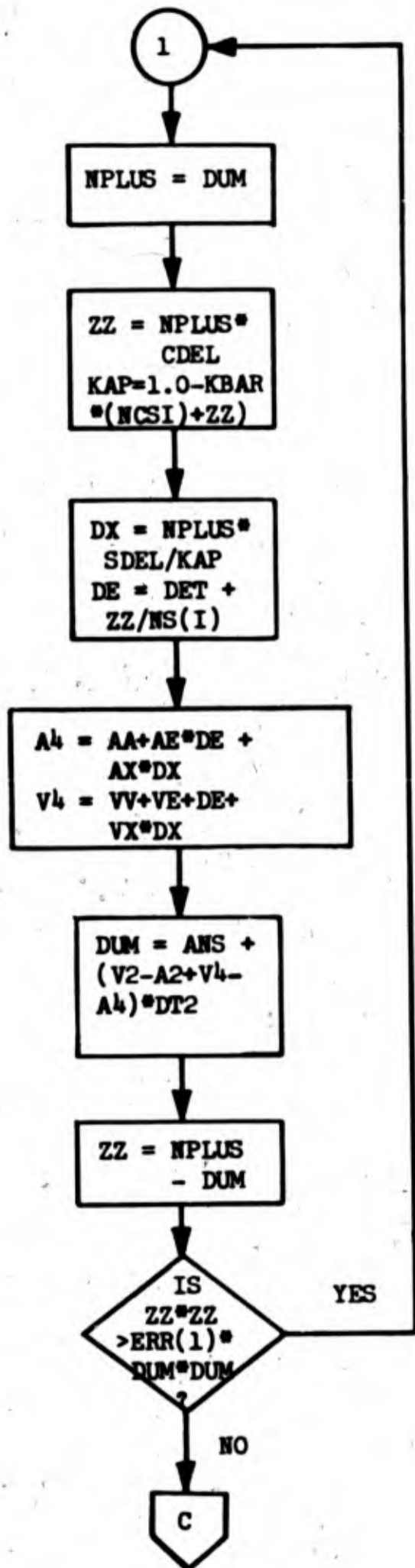
COMPUTE V AT THE SHOCK.

STORE P AT THE SHOCK.

STORE ρ AT THE SHOCK.

COMPUTE a AT THE SHOCK.

COMPUTE X_N .



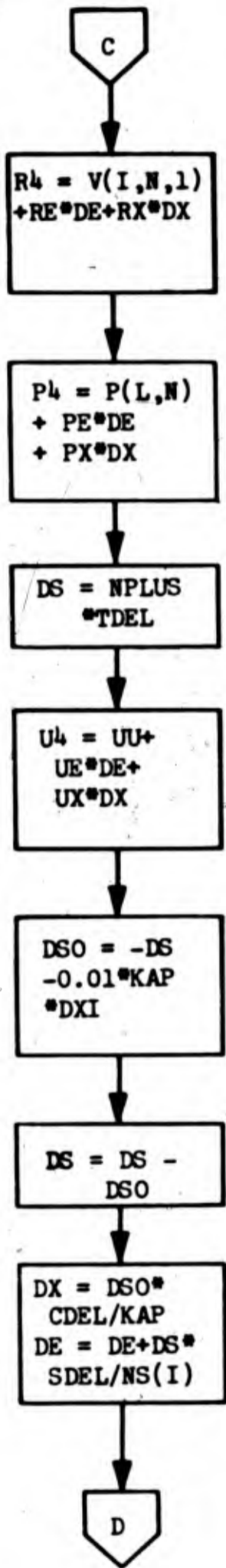
STORE x_N .

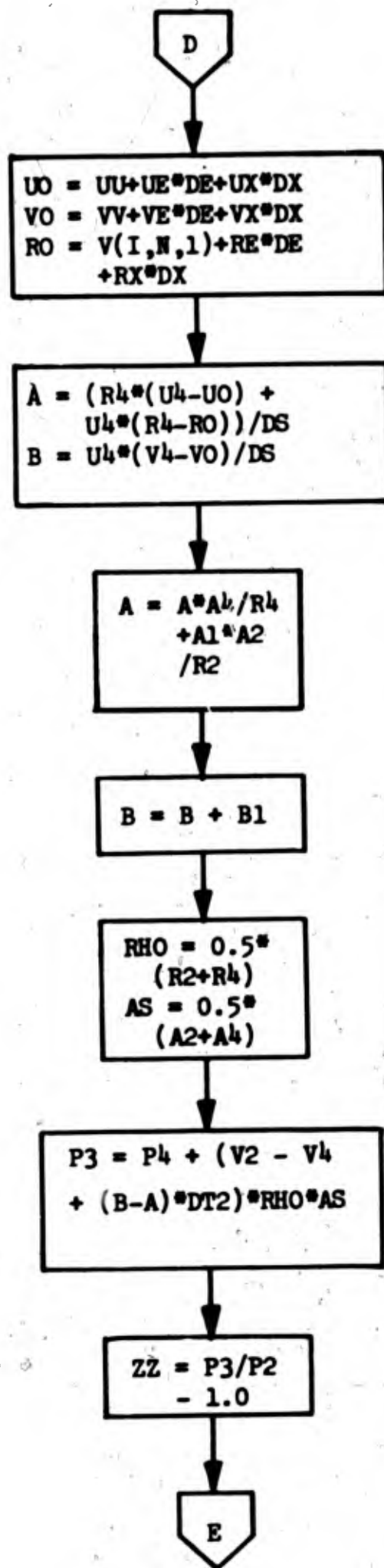
COMPUTE $\Delta \xi$ AND $\Delta \eta$
BETWEEN x_N AND THE
POINT (I, N)

COMPUTE a_4 AND v_4 .

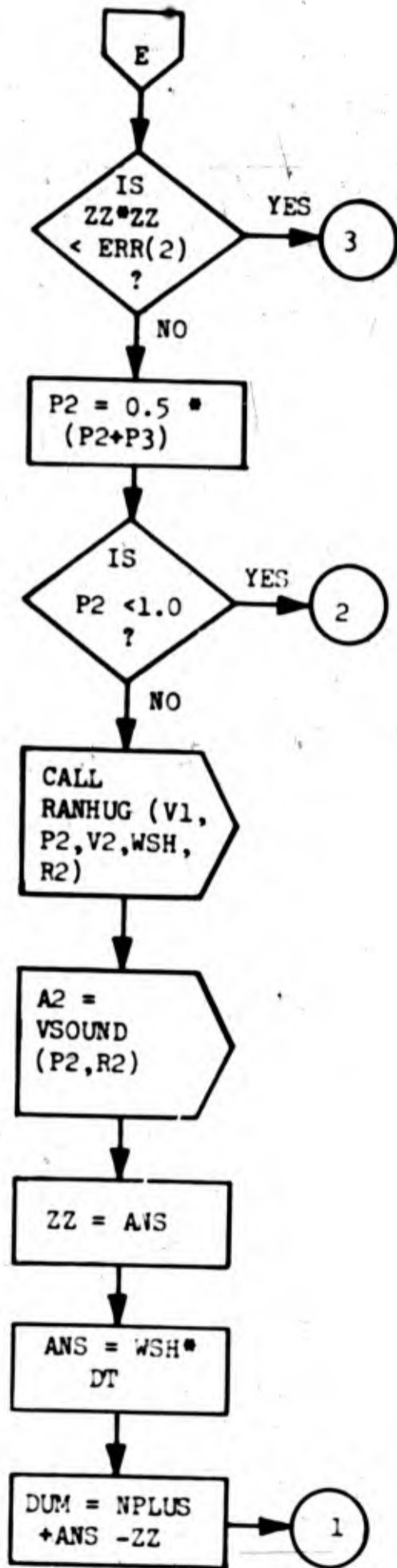
COMPUTE NEW x_N .

IS THE DIFFERENCE BETWEEN
SUCCESSIVE VALUES OF x_N
TOO LARGE ?





EQUATION (49)



DO TWO SUCCESSIVE VALUES OF SHOCK PRESSURE AGREE WITHIN THE ALLOWED ERROR ?

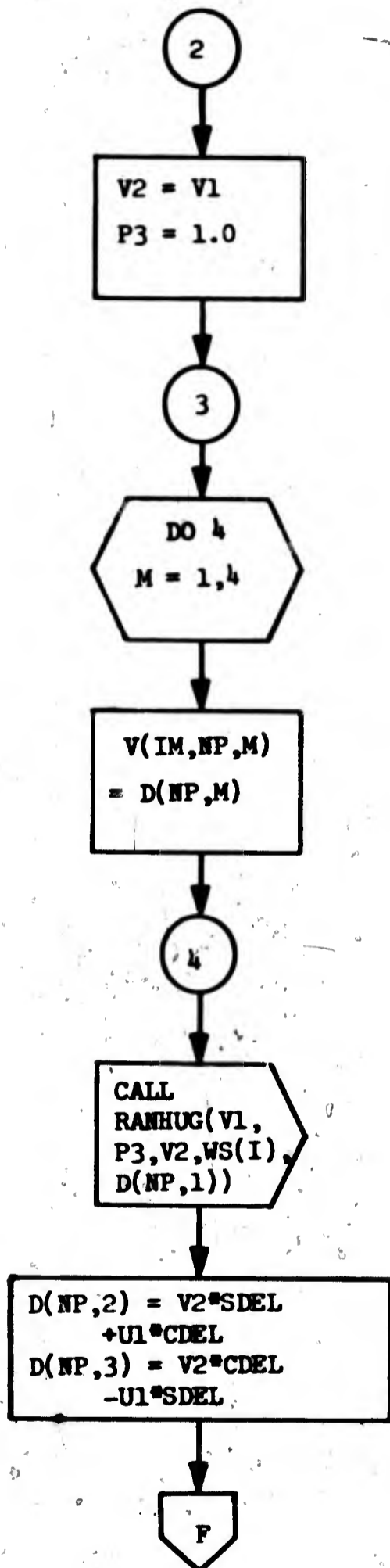
AVERAGE THE TWO VALUES.

IS THE SHOCK PRESSURE < 1.0 ?

SOLVE RANKINE-HUGONIOT RELATIONS.

COMPUTE NEW SOUND SPEED AT THE SHOCK.

COMPUTE NEW VALUE OF NPLUS.



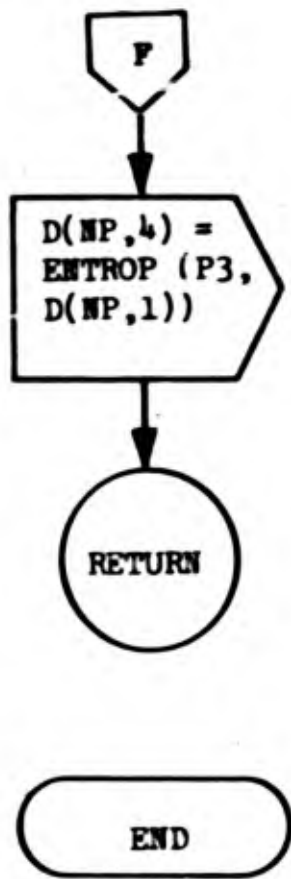
PRESSURE AT THE SHOCK WAS
< 1.0. SHOCK IS A MACH WAVE.

SET NORMAL VELOCITY AND PRESSURE
AT THE SHOCK EQUAL TO FREE
STREAM VALUES.

TRANSFER ρ , U, V AND S FROM
TEMPORARY TO PERMANENT STORAGE
AT IM.

SOLVE RANKINE-HUGONIOT
RELATIONS.

COMPUTE U AND V.



COMPUTE S AT THE SHOCK.

RETURN.

D. LISTING

0
0
0
0

```

SUBROUTINE SHOCK (ATSTAG)
    THIS ROUTINE TREATS THE SHOCK NODES USING THE METHOD SUGGESTED
    BY MARIETTI AND ROBERTI, ATAA J. NO. 4, 2176-2141 (1966).

COMMON /1/ V(15,4),Z(4),P(3,4)
COMMON /2/ PE(15),ZB(15),TH(15),CTH(15),K(15),NS(15),MS(15)
COMMON /3/ MAXI,MI,IMIN,MP,N,NM
COMMON /5/ DX1,DET,DX2,DF2,DT
COMMON /6/ Q1,VS1,HT1,A1,Q1
COMMON /7/ I,IC,IM,L,LP,LM
COMMON /8/ PX,LX,VX,AX,KBAR,C2J
COMMON /12/ DEFT,COEL,SOEL,TOEL
COMMON /INPUT/ GAM1,AMACH1,THFTC,ZRN,TCOIT(2),P1,RHC1,PN,AMUD,T1,F
*CO(2),IPR(6),INTAP,OTAP,NCPT,NCASE,NFTLE,NSEG,IGFCM,NGEOM,NTAP
COMMON /CONST/ PI,PI2,GAMA,GAMB
REAL NS,K,KAP,NPLUS,KBAR
Z7=COET+TH(I)
U1=Q1*COS(ZZ)
V1=Q1*SIN(ZZ)
AA=VSCUND(P(L,A),V(I,N,1))
PC=(P(L,MP)-P(L,NM))/DF2
PE=(V(I,MP,1)-V(I,NM,1))/DE2
UF=(V(I,MP,2)-V(I,NM,2))/DE2
VE=(V(I,MP,3)-V(I,NM,3))/DE2
PX=(P(LP,N)-P(LM,N))/DX2
ZZ=UX
UX=Z7*COEL-VX*SOEL
VX=VX*SOEL+Z7*SOEL
Z7=UF
UF=Z7*COEL-VE*SOEL
VF=VE*SOEL+Z7*SOEL
AF=(VSCUND(P(L,MP),V(I,MP,1))-VSCUND(P(L,NM),V(I,NM,1)))/DE2
UU=V(I,N,2)*COEL-V(I,N,3)*SOEL
VV=V(I,N,2)*SOEL+V(I,N,3)*COEL
V2=V(I,MP,2)*SOEL+V(I,MP,3)*COEL
P2=P(L,MP)
DT2=.5*DT
P2=V(I,MP,1)
A2=VSCUND(P2,Z2)
ANS=MS(I)*DT
DUM=ANS-(V2-A2)*DT
NPLUS=DUM
Z7=NPLUS*COEL
KAP=1.-KBAR*(NS(I)+Z7)
PX=NPLUS*SOEL/KAP
DE=DET+Z7/NS(I)
A4=AA+AF*DE+AX*DX
V4=VV+VF*DE+VX*DX
DUM=ANS-(V2-A2+V4-A4)*DT2
ZZ=NPLUS-DUM
IF (Z7*.GT.ERR(1)*DUM*DUM) GO TO 1
    
```

SHOCK 1
SHOCK 2
SHOCK 3
SHOCK 4
SHOCK 5
SHOCK 6
SHOCK 7
SHOCK 8
SHOCK 9
SHOCK 10
SHOCK 11
SHOCK 12
SHOCK 13
SHOCK 14
SHOCK 15
SHOCK 16
SHOCK 17
SHOCK 18
SHOCK 19
SHOCK 20
SHOCK 21
SHOCK 22
SHOCK 23
SHOCK 24
SHOCK 25
SHOCK 26
SHOCK 27
SHOCK 28
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SHOCK 31
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SHOCK 34
SHOCK 35
SHOCK 36
SHOCK 37
SHOCK 38
SHOCK 39
SHOCK 40
SHOCK 41
SHOCK 42
SHOCK 43
SHOCK 44
SHOCK 45
SHOCK 46
SHOCK 47
SHOCK 48
SHOCK 49
SHOCK 50
SHOCK 51
SHOCK 52

1

```

P3=U(I,M,1)+P3*DE+DX*DY
P4=P(L,M)+P4*DE+DX*DY
P5=U(L,M)*DE
U4=U1+U2*DE+UX*DY
P31=-P3-1.)*I*NS*DY
P5=P5-NS)
DY=P5*DE/KA=
P3=P3+P5*DE/NS(I)
U4=U1+U2*DE+UX*DY
V1=V1+V2*DE+VX*DY
P3=V(I,M,1)+P3*DE+DX*DY
A=(P4*(U4-U2)+L4*(P4-P1))/DS
P=U4*(V4-V1)/DS
A=(A*A4/P4+11*P2/P2)*ATSTAG
P=P+3)
PHO=1.5*(P4+P2)
A4=1.5*(A4+A2)
P2=P4+(V2-V4+(P-1)*DT2)*RHO*AS
P2=P2/P2-1.)
IF (P2.LT.EPP(P)) GO TO 3
P2=1.5*(P2+P2)
IF (P2.LT.1.7) GO TO 2
CALL PAMHUG (V1,P2,V2,WSH,P2)
P2=V2*U1*(P2,P2)
P2=ANS
ANS=4*P2
PUM=UPLUR+ANS-22
GO TO 1
2
V1=V1
P2=1.7
3
P2 4 M=1,4
4
V(I,M,M)=P(AN,M)
CALL PAMHUG (V1,P2,V2,WS(I),P(NP,1))
P(NP,2)=V2*SCFL+U1*DEFL
P(NP,3)=V2*DEFL-U1*SCFL
P(NP,4)=ENTRCE(P2,P(NP,1))
PSTUN
=AN

```

```

SHOCK 57
SHOCK 58
SHOCK 59
SHOCK 60
SHOCK 61
SHOCK 62
SHOCK 63
SHOCK 64
SHOCK 65
SHOCK 66
SHOCK 67
SHOCK 68
SHOCK 69
SHOCK 70
SHOCK 71
SHOCK 72
SHOCK 73
SHOCK 74
SHOCK 75
SHOCK 76
SHOCK 77
SHOCK 78
SHOCK 79
SHOCK 80
SHOCK 81
SHOCK 82
SHOCK 83
SHOCK 84
SHOCK 85
SHOCK 86
SHOCK 87
SHOCK 88
SHOCK 89
SHOCK 90
SHOCK 91

```


20. SUBROUTINE SKIP

A. PURPOSE

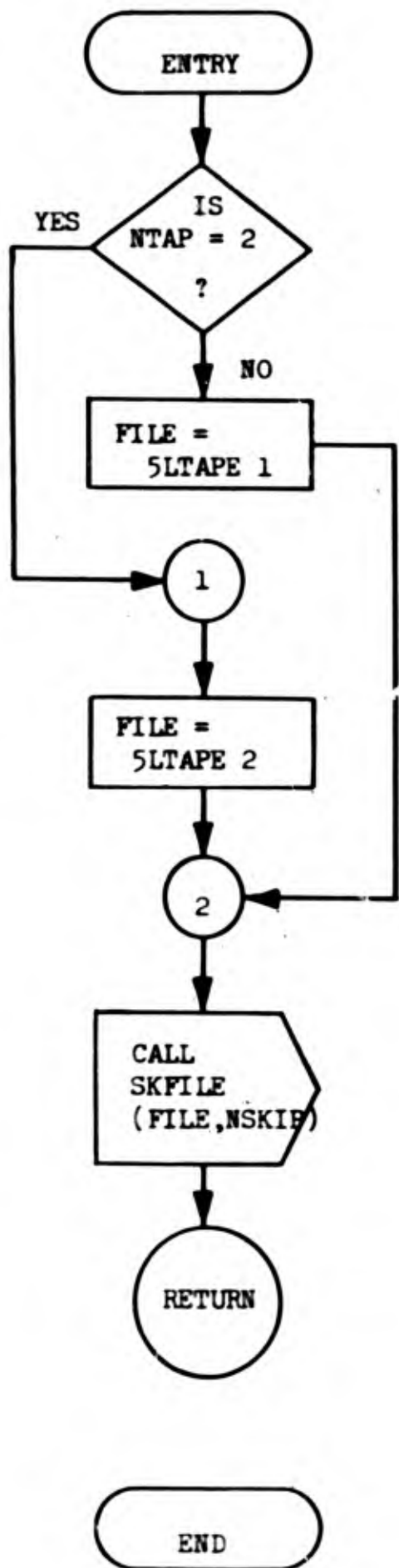
This subroutine has been included to isolate a call to a CDC machine routine which skips files on a magnetic tape. This machine routine requires a Hollerith argument which might not be accepted by the equivalent routine available to the user. Provisions for skipping files only on tapes 1 and 2, have been provided. This could easily be generalized to apply to other tape numbers.

B. VARIABLE LIST

NTAP - magnetic tape unit number.

NSKIP - number of files to be skipped.

C. FLOW CHART



ENTER SUBROUTINE SKIP.

IS NTAP = 2 ?

SET ARGUMENT TO SKIP FILES ON TAPE 1.

SET ARGUMENT TO SKIP FILES ON TAPE 2.

SKIP NSKIP FILES.

RETURN.

END OF SUBROUTINE SKIP.

D. LISTING

```
3      SUBROUTINE SKIP (NTAP,NSKIP)
C      THIS IS A DUMMY ROUTINE. IT IS USES A CDC MACHINE ROUTINE
C      TO SKIP FILES ON A MAGNETIC TAPE UNIT. NTAP=TAPF NUMBER, NSKIP=
C      NUMBER OF FILES TO BE SKIPPED.
C
C      IF (NTAP.EQ.2) GO TO 1
C      FILE=5LTAPE1
C      GO TO 2
C      FILE=5LTAPE2
C      CALL SKFILE (FILF,NSKIP)
C      RETURN
C      END
```

```
SKIP 1
SKIP 2
SKIP 3
SKIP 4
SKIP 5
SKIP 6
SKIP 7
SKIP 8
SKIP 9
SKIP 10
SKIP 11
SKIP 12
SKIP 13
```

21. EQUATION OF STATE ROUTINES

A. PURPOSE

These routines provide the required equation of state computations. Perfect gas equations of state are used. Real gas properties could be incorporated into the various routines. The routines available are:

SUBROUTINE STATE - computes the sound speed from values of density and entropy.

FUNCTION VSOUND - computes the sound speed from values of density and pressure.

FUNCTION TEMP - computes the temperature from values of density and pressure. Used only for printed output.

FUNCTION HSP - computes the enthalpy from values of density and pressure. Used only for printed output.

FUNCTION ENTROP - computes the entropy from values of density and pressure.

SUBROUTINE RANHUG - solves the Rankine-Hugoniot relations for the shock.

B. VARIABLE LIST

Variables previously defined in the description of the main program have been omitted from this list.

ENTROP - entropy.

HSP - enthalpy.

MDOT - mass flux, \dot{m}

P - pressure.

P2 - pressure behind a shock.

RGAS - gas constant.

RHO - density.

RHO2 - Density behind a shock.

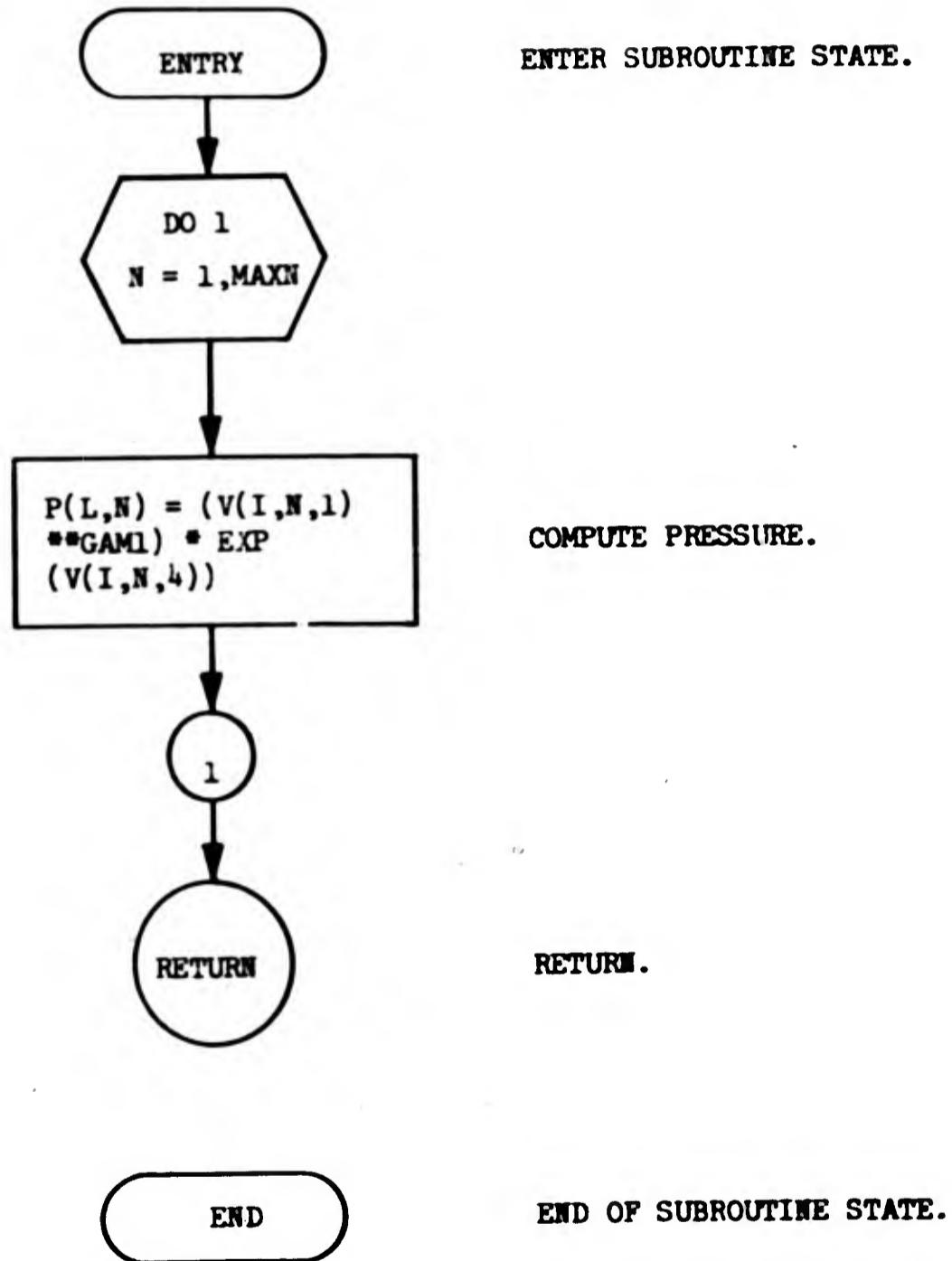
TEMP - temperature.

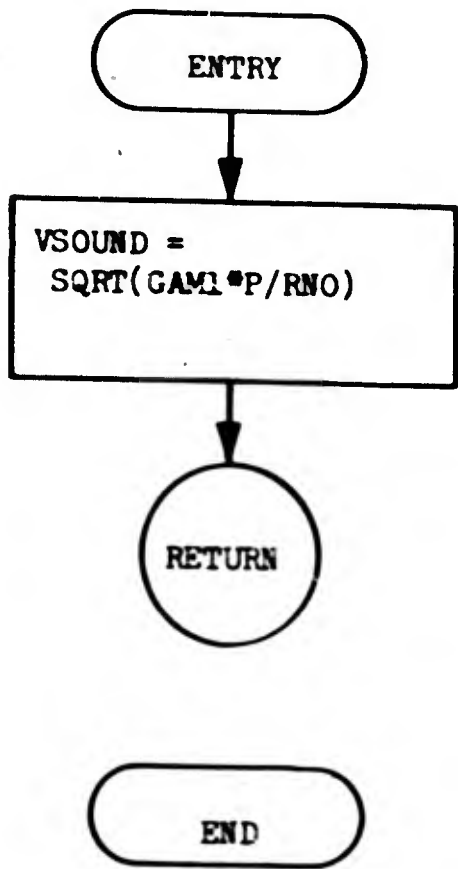
VSOUND - sound speed.

V1 - free-stream velocity component normal to the shock.

V2 - velocity component normal to the shock and behind the shock.

C. FLOW CHART



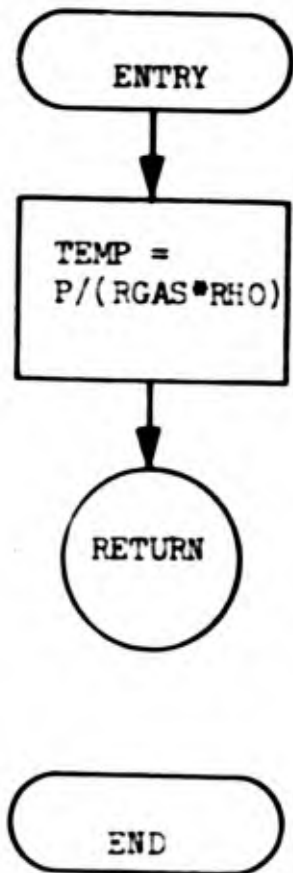


ENTER FUNCTION VSOUND.

COMPUTE SOUND SPEED.

RETURN.

END OF FUNCTION VSOUND.

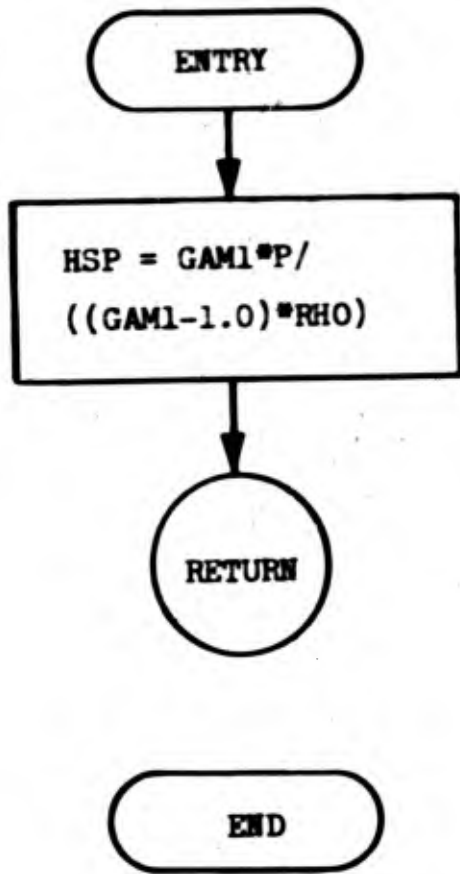


ENTER FUNCTION TEMP.

COMPUTE TEMPERATURE.

RETURN.

END OF FUNCTION TEMP.

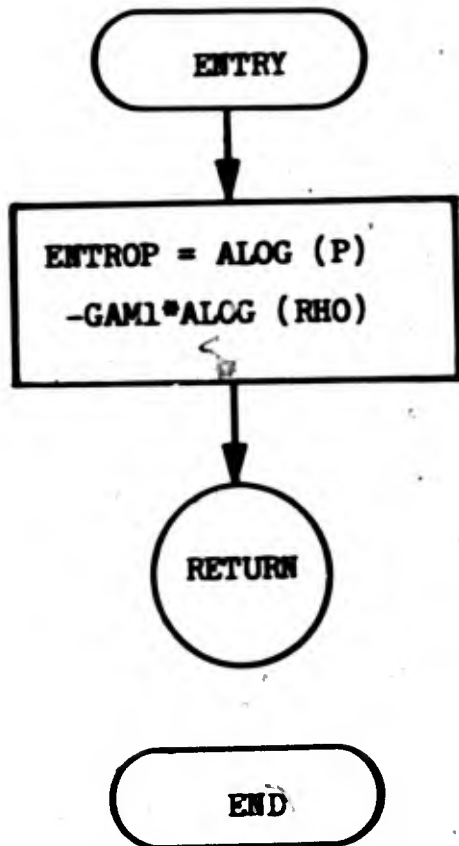


ENTER FUNCTION HSP.

COMPUTE SPECIFIC ENTHALPY.

RETURN.

END OF FUNCTION HSP.



ENTER FUNCTION ENTROP.

COMPUTE ENTROPY.

RETURN.

END OF FUNCTION ENTROP.

ENTRY

ENTER SUBROUTINE RANHUG.

$MDOT = \sqrt{GAMA * P2 + GAMB}$

COMPUTE m .

$V2 = V1 + (1.0 - P2) / MDOT$

COMPUTE V AT THE SHOCK.

$W = V1 - MDOT$

COMPUTE W_s .

$RHO2 = MDOT / (V2 - W)$

COMPUTE ρ AT THE SHOCK.

RETURN

RETURN.

END

END OF SUBROUTINE RANHUG.

7. LISTING

```

SUBROUTINE STATE (I,L,MAXN)
C
C THIS ROUTINE COMPUTES PRESSURE, GIVEN DENSITY AND ENTROPY,
C USING A PERFECT GAS EQUATION OF STATE.
C
COMMON /V/ V(15,4),C(4,4),P(3,4)
COMMON /INPUT/ GAM1
DO 1 N=1,MAXN
1 P(L,N)=(V(I,N,1)**GAM1)*EXP(V(I,N,4))
RETURN
END
STATE 1
STATE 2
STATE 3
STATE 4
STATE 5
STATE 6
STATE 7
STATE 8
STATE 9
STATE 10
STATE 11
STATE 12

```

```

FUNCTION VSCALD (P,RHO)
C
C THIS ROUTINE COMPUTES THE SOUND SPEED USING THE PERFECT GAS
C RELATION.
C
COMMON /INPUT/ GAM1
VSCALD=SQRT(GAM1*P/RHO)
RETURN
END
VSCUN 1
VSCUN 2
VSCUN 3
VSCUN 4
VSCUN 5
VSCUN 6
VSCUN 7
VSCUN 8
VSCUN 9

```

```

FUNCTION TEMP (P,RHO)
C
C THIS ROUTINE COMPUTES THE TEMPERATURE USING A PERFECT GAS
C RELATION (USED ONLY FOR PRINTOUT).
C
COMMON /CONST/ PI,PI2,GAMA,GAMP,RCAS
TEMP=P/(RCAS*PI)
RETURN
END
TEMP 1
TEMP 2
TEMP 3
TEMP 4
TEMP 5
TEMP 6
TEMP 7
TEMP 8
TEMP 9

```

```

FUNCTION HSD (P,RHO)
C
C THIS ROUTINE COMPUTES THE STATIC ENTHALPY USING A PERFECT GAS
C RELATION (USED ONLY FOR PRINTOUT).
C
COMMON /INPUT/ GAM1
HSD=GAM1*P/((GAM1-1.0)*RHO)
RETURN
END
HSD 1
HSD 2
HSD 3
HSD 4
HSD 5
HSD 6
HSD 7
HSD 8
HSD 9

```

FUNCTION ENTROP (P,RHO)

THIS ROUTINE COMPUTES THE ENTROPY, GIVEN THE PRESSURE AND DENSITY, USING A PERFECT GAS EQUATION (USED IN SHOCK).

COMMON /INPUT/ GAM1
ENTROP=ALOG(P)-GAM1*ALOG(RHO)
RETURN
END

ENTRO 1
ENTRO 2
ENTRO 3
ENTRO 4
ENTRO 5
ENTRO 6
ENTRO 7
ENTRO 8
ENTRO 9

SUBROUTINE PANKUG (V1,P2,V2,W,RHO2)

THIS ROUTINE SOLVES THE PANKINE-HUGONIOT RELATIONS FOR A NORMAL SHOCK USING PERFECT GAS RELATIONS.

COMMON /CONST/ PI,PI2,GAMA,GAMB
REAL MNDT
MNDT=SQRT(GAMA*P2+GAMB)
V2=V1+(1.0-P2)/MNDT
W=V1-MNDT
RHO2=MNDT/(V2-W)
RETURN
END

PANKU 1
PANKU 2
PANKU 3
PANKU 4
PANKU 5
PANKU 6
PANKU 7
PANKU 8
PANKU 9
PANKU 10
PANKU 11
PANKU 12
PANKU 13

SECTION VII

COMPUTER CODE USER'S INSTRUCTIONS

1. INPUT DATA

The computer code provided in section VI performs most of the decision making and scheduling required to apply the numerical method developed in section IV. Since a certain degree of uncertainty is always involved when applying a time-dependent method to the blunt body problem, a very general input capability has been incorporated. Should the internal logic or empirical constants prove inadequate for a particular problem, the user has several options available to over-ride the internal logic with direct input parameters. It should be noted that the internal logic has been adequate for all problems considered in this study. To avoid unnecessarily complicated data input while retaining the general input capability, all problem constants, empirical constants and control variables are loaded by a sorting routine, SUBROUTINE DATA. These data will be referred to as case data variables. An arbitrary number of these variables may be input in any order. The principal advantage of this data loading procedure is that initial resident values supplied for most of the case data variables by a DATA statement in the main program can be used. Consequently, the case data input required is usually quite brief, typically 2 to 4 variables, while up to 23 variables can be input if necessary. Case data input is also simplified for multicase runs, since only variables whose resident values are to be changed need be included in the case data input. Some caution is required when relying on resident values from a previous case. The computer code accomplishes numerous error checks while loading the input data. When an error is sensed, the case under consideration is terminated with a printed error message and a search for the first data card of the next case is initiated. This will usually salvage a multicase run, but some case data variables may not have been loaded. Thus, reliance on a resident value loaded in a previous case could result in solving the wrong problem. Since the case data input is usually quite brief, some repetitions of data input for subsequent cases is advisable on long multi-case runs. Also, when relying on resident values from a previous case, the user must be aware of the source of these values. Usually, the resident values originate from the DATA statement mentioned previously and card input. However, under certain options described later in this section,

many case data variables are loaded from magnetic tape. When sphere-cone vehicles are to be considered, the case data completely specify the problem. Other body shapes are considered by inputting body geometry data. Basically, this data consists of values of r , z , Θ and K at points on the body surface. The body geometry data input is described in more detail later in this section.

2. CASE DATA INPUT

The computer code first loads the case data in SUBROUTINE DATA. The case data packet organization is as follows:

First card: FORMAT (8A10)

Field 1 (columns 1-10) - the word TITLE, left justified in the field.

This signals the start of the case data.

Fields 2-8 (columns 11-80) - an alphanumeric title which will appear on the printed output.

Cards 2, 3, . . . : FORMAT (A10, E10.4, 6I10)

Field 1 (columns 1-10) - the variable name, left justified in the field. The order in which variables are input is arbitrary.

Field 2 (columns 11-20) - the value of the variable if it is floating point form.

Field 3 (columns 21-30) - the value of the variable if it is integer form.

Fields 4-8 (columns 31-80), 10 per field) - additional integer values for the dimensional variable, IPR(6).

Last card: FORMAT (A10, E10.4, 6I10)

Field 1 (columns 1-10) - the word LAST, left justified in the field.

This signals the end of the case data.

The complete list of case data variable names and their definitions will now be given. An (I) and (F) will be used to signify integer and floating point variables, respectively. Initial resident values will be specified when they have been provided.

M - (F).

>0 - free-stream Mach number.

<0 - free-stream velocity. The absolute value will be used with the values of P_1 and RH_01 (below) to compute the free stream Mach number. Consequently, the units of these 3 variables must be compatible.

- Z/RN - (F). Specifies the body geometry option to be used. Initialized to 0.
 >0 - Specifies the axial length, in nose radii, for a sphere-cone vehicle. No body geometry data will be read.
 ≤0 - Arbitrary body geometry data will be read.
- THETA - (F) cone half-angle (degrees). Used only when Z/RN>0.
- GAMMA - (F) ratio of specific heats. Initialized to 1.40.
- NTAP - (I) specifies the type of initial conditions to be used. Initialized to 0.
 >0 - initial shock shape assigned close to the body and at a constant distance from the body. Used to model a transient flow essentially starting impulsively from rest. Should not be used when NSEG >1.
 =0 - initial shock shape selected empirically. Attempts to approximate the steady state shock shape as the initial profile.
 =1 - restart option (discussed below).
 =2 - load option (discussed below).
- MAXIT - (I). specifies the iteration procedure to be used. Initialized to -1.
 <0 - steady-state solution using equation (63) and the requirement that the non-dimensional shock velocities be $<10^{-3}$ as criteria for a steady state.
 =0 - no iterations are performed, but a complete printed output is provided. Useful for checking initial conditions or for obtaining dimensional output for a completed case.
 >0 - specifies the number of iterations to be performed for all segments. This option is required when NTAP=2.
- NSEG - (I). specifies the number of segments to be considered. Must be input when the body geometry is read in. A value greater than 0 will over-ride the axial length specification provided for the sphere-cone vehicle option. Normally should be ≤0 for this option. Initialized to 0.
- IPR - (I).-dimensional variable providing up to 6 additional printed outputs. Since a printed output is provided after the last iteration, specifying that iteration number here will give 2 identical printed outputs and should be avoided. A negative value, or a value less than a preceding value will suppress further printed output except for the final. IPR(1)=0 will result in printing the initial conditions. All values initialized to -5.

MAXIT 0 - iteration numbers, in ascending order, at which printed outputs should be provided.

MAXIT=0 - no effect

MAXIT<0 - numbers between 0 and 10, in ascending order, specifying times when additional printed outputs should be supplied, where a value of 10 corresponds to the predicted time for steady state generated by equation (63).

- P1 - (F).- free-stream pressure. Initialized to 1.0 (thus will give non-dimensional output).
M>0 - used only to supply dimensional printed output.
M<0 - same as above plus used to compute the free-stream Mach number (see discussion of M, above).
- RH01 - (F).- free-stream density. Same application as P1 and units must be compatible with P1. Initialized to 1.0 (thus, will give nondimensional output).
- T1 - (F).- free-stream temperature. Used only to provide dimensional printed output. Units must be compatible with P1 and RH01. Initialized to 1.0 (thus, will give nondimensional output).
- RN - (F).- all length parameters are multiplied by this scale factor to provide dimensional printed output. Units are arbitrary. Normally, RN would be equal to the radius of curvature at the stagnation point. Initialized to 1.0 (thus, will give nondimensional output).
- MU - (F).- constant, μ_0 , in equation (62) specifying the magnitude of the stabilizing terms. Initialized to 0.04.
- MAXI - (I). specifies the number of nodes to be used per segment in the streamwise direction. A value >0 will over-ride other possible specifications for this quantity (see body geometry data for an exception). Normally, an empirical specification (sphere-cone vehicles) or a specification supplied in the body geometry input would be used. Initialized to 0.
- MAXN - (I). specifies the number of nodes to be used across the shock layer. A value >0 will over-ride all other specifications of this quantity and the empirical specification generated by the computer code. Normally, this specification is not used. Initialized to 0.
- INTAP - (I). specifies the magnetic tape unit number to be used for input of initial conditions. Used only for NTAP>0. See discussion of

- tape handling, below. Initialized to 2.
- OTAP - (I). specifies the magnetic tape unit number to be used for storing the completed solutions for all segments. See discussion of tape handling, below. Initialized to 1.
- NCASE - (I). specifies the case number to be used in supplying initial conditions, where the cases are considered to be numbered in ascending order starting with 1 as they appear on the tape. Used only when NTAP>0 (see discussion of Restart and Lead options). Initialized to 1.
- NFILE - (I). must be equal to 1 unless NTAP>0. Initialized to 1.
- NTAP=1 : Specifies the number of the last segment completed for an incomplete case. Restart initiates on the next segment.
- NTAP=2 : Specifies the segment number at which computation should start. All segments following this one will also be treated.
- TCRIT1 - (F). constant, C_1 , in Equation (63). Initialized to 2.0.
- TCRIT2 - (F). constant, C_2 , in Equation (63). Initialized to 10.0.
- ERRORN - (F). convergence criterion for the iteration on n in the shock point computations. Initialized to 10^{-5} .
- ERRORP - (F). convergence criterion for the iteration on P in the shock point computations. Initialized to 10^{-8} .

3. BODY GEOMETRY DATA INPUT

The body geometry data are loaded by SUBROUTINE GEOM when the case data variable $Z/RN \leq 0$. These data cards are arranged in packets, one for each segment considered, and are supplied immediately after the case data cards. An exact specification of the body geometry at each point for the present coordinate system is supplied by values for r, z, θ and K. It is recognized that not all of these parameters will be known for some problems of interest. However, it has been left to the user to select approximate numerical techniques for generating these parameters when necessary. This computer code supplies an exact solution for known body geometrics. Since explicit expressions for arc length are seldom available, even for analytical body profiles, numerical procedures have been incorporated to generate the required parameters at equal spaced intervals along the body, based on their values at a representative series of points. For example, it is easy to define all of the required parameters at any point on an ellipse,

but considerably more difficult to compute their values at equally spaced intervals along the body. Since the geometry specifications for the representative series of points are exact, extremely accurate numerical integrations and interpolations are possible to obtain the equal spaced data. When computations start at a segment other than the first, $NTAP > 0$, all body geometry data packets can be included. The computer code skips the data for all segments up to the one at which computations start. Specifications for the number of nodes to be used along each coordinate direction can be specified for each segment considered. An initial shock shape can be read in. This option has been included since the user could encounter a problem for which the empirical relation used to generate initial shock shape is inadequate. The required parameters can be input at equal spaced intervals along the body, for direct use in the numerical solution, or at a representative series of points on the body. Each body geometry packet is arranged as follows:

Card 1 : FORMAT (5I10, 2E10.4, A10)

Field 1 (columns 1-10) - M. The segment number.

Field 2 (columns 11-20) - NPTS. The number of data cards to follow for the current segment. If no additional specifications for the number of nodes to be used along the body are provided, the value of NPTS will be used as this specification. Also, see field 7 of this card.

Field 3 (columns 21-30) - MAXI. The total number of streamwise nodes to be used. Specifications of this variable in the case data will over-ride the present input. Also, see field 7 of this card.

Field 4 (columns 31-40) - IP. The number of nodes to be used across the shock layer. Specification of MAXN in the case data will over-ride the present input. If no specification for this parameter is supplied, the computer code will determine one empirical based on the variable input into field 6 of this card.

Field 5 (columns 41-50) - KSH.

KSH = 0. Empirical relations used to determine initial shock shape.

KSH \neq 0. Initial shock shape supplied on the following data cards.

Field 6 (columns 51-60) - BNOSE. A nose bluntness parameter. Used to establish the nodal spacing across the shock layer when it is determined empirically. The usual definition applies, e.g., for an elliptical nose BNOSE is the ratio of the semi-major axes squared. A value ≤ 0 is ignored. For each case, BNOSE is initialized to 1.0. For downstream segments where nose bluntness effects have damped out, BNOSE should be set to 1.0.

Field 7 (columns 61-70) - DX2.

$DX2 \leq$ - The body geometry data to follow will be specified for a representative series of body points, not necessarily equally spaced along the body. Numerical procedures will be used to obtain equally spaced data.

$DX2 > 0$ - The body geometry data to follow will be specified for points equally spaced along the body with a nodal spacing equal to the value read in this field. The number of stream-wise nodes to be used will be established by the value of NPTS. Other specifications for this parameter will be ignored. Thus, the body geometry data will be used directly in the numerical solution.

Field 8 (columns 71-80) - ZZ. Any entry into this alphanumeric field will signify a discontinuity in body curvature at the upstream boundary of the segment. In this case, the body geometry data on card 2 will be identical to the data on the last card for the previous segment, except for the new value of body curvature. If this field is blank, no data for the upstream boundary should be input since all parameters will be saved from the previous segment.

Cards 2, 3 . . . , (NPTS+1): FORMAT (8E10.4)

Field 1 (columns 1-10) - RB(I). The body radius, measured normal to the symmetry axis, for the Ith point.

Field 2 (columns 11-20) - ZB(I). The value of Z, measured from the stagnation point along the symmetry axis, for the Ith point.

Field 3 (columns 21-30) - TH(I). The angle, Θ , between the symmetry axis and a tangent to the body for the Ith point.

TH(I) > 0 - units are degrees.

TH(I) < 0 - units are radians.

Field 4 (columns 31-40) - $K(I)$. The body curvature at the I th point.

Field 5 (columns 41-50) - $NST(I)$. The initial value of n_s at the I th point (must be <0 when used). Ignored if $KSH = 0$.

The above data packets are repeated for each segment, 1 to NSEG. The maximum values of NPTS and MAXI (Fields 2 and 3 of Card 1) are limited by the dimensioned variable sizes (currently 15 streamwise nodes are available). MAXI must not exceed this value since it includes the additional downstream node used for linear extrapolation and, for segments other than the first, a node for the constant upstream starting line. In establishing NPTS, more care is required. When $DX2 \leq 0$ (card 1, Field 7), the total number of body geometry data points (including the upstream boundary point for segments other than the first) must not exceed 15. When $DX2 > 0$, the value of MAXI will be established from NPTS. Therefore, the total number of body geometry data points (including the upstream boundary point for segments other than the first) must not exceed 14. That is, all body geometry points input will be actual computation points. Thus, an additional point will be required for the linear extrapolation at the downstream boundary. Should the user exceed the dimension sizes, the case will be terminated with a printed error message.

4. MAGNETIC TAPES

The computer code employs two magnetic tape units. The tape unit number specified by the case data variable INTAP is used as the source of initial conditions when $NTAP > 0$. The tape unit number specified by OTAP is used to store the flow field data for each segment as soon as the solution at that segment is completed. The computer code rewinds tape unit number INTAP immediately after loading the case data. Tape unit number OTAP is never rewound by the computer code. Thus, for multi-case runs, the 2 tape unit numbers normally should not be equal. However, when $NTAP \leq 0$, the computer code does not treat this situation as an error since it is often useful to have $INTAP = OTAP$. When $NTAP > 0$, the computer code terminates the case if the following restrictions are violated:

$NTAP = 1$: Requires $INTAP = OTAP$

$NTAP = 2$: Requires $INTAP \neq OTAP$

5. RESTART OPTION

When the case data variable $NTAP = 1$, the computer code will restart an unfinished case using the data stored on magnetic tape. The case data variables

NTAP, INTAP, OTAP, NCASE and NFILE will retain the values present after the case data was loaded. All other case data variables will be read from magnetic tape for the unfinished case. As previously noted, INTAP = OTAP is required. This means the data up to and including the last segment completed (NFILE) will be skipped on the tape and the remainder of the solution will be loaded on the same magnetic tape. The option to restart an unfinished case is a necessity for a time-dependent method. When a steady state solution is being generated, it is difficult to estimate the computation time required, since it is quite dependent on the body geometry and reentry conditions. Recognizing that the computation time allotted by the user or available to him may not be sufficient, the restart option has been incorporated to prevent excessive waste of computer time.

6. LOAD OPTIONS

When the case data variables NTAP = 2, the computer code loads all initial conditions from magnetic tape unit INTAP. All data for the case up to but excluding segment number NFILE is transferred to tape unit number OTAP with no computation. As previously noted, INTAP \neq OTAP is required. The case data variables NTAP, MAXIT, IPR, INTAP, OTAP, NCASE, NFILE, P1, RH01, T1, RN, MU, ERRORN and ERRORP retain the values present from the case data input operation. All other case data variables are read from magnetic tape. The above list of variables includes all of the case data variables the user can change without altering the basic problem stored on magnetic tape. The variable MAXIT \geq 0 is required, i.e., the number of iterations to be performed on each segment must be specified. The principal purpose of this option is to allow the user to perform additional iterations on a case when a satisfactory approximation to the steady-state flow was not realized on a previous attempt. Rather than starting the case over with new values of TCRIT1 and TCRIT2, the user can use the previous solution as an initial guess for additional iterations. It should be noted that none of the cases attempted in this study required use of this option. Other applications for this option can be identified. By setting MAXIT = 0 and specifying P1, RH01, T1 and RN, dimensional output can be obtained from the nondimensional data stored on magnetic tape for a completed case while no iterations will be performed. Also, this option could be used to input initial conditions, should the empirical methods provided for this purpose prove inadequate for a particular problem.

7. SEGMENT LENGTHS

The choice of the length of the segment to be used can significantly affect the computation time required for the problem. The length of the first segment

will, generally, be specified by the location of the sonic line, since the downstream segment boundary must lie in supersonic flow. The lengths of other segments are arbitrary. It should be apparent that the use of shorter segments will reduce the computation time required for the problem. If a segment is treated with a single computational node (MAXI = 3), the length of the segment will have little effect on the computation time. It is easily shown that this is the optimum choice for the number of nodes per segment. Even for this optimum choice, reasonably short segments are recommended to provide greater numerical accuracy. Also, should the case require more computation time than was allotted by the user, shorter segments will result in more efficient restarts. That is, shorter segments require less computation time per segment. Consequently, less computation time will be lost when a restart is necessary.

8. PRINTED OUTPUT

The computer code provides a comprehensive printed output. After the case data is loaded, the case title and all case data variables are printed. After computations for a segment are completed, and at intermediate iterations when requested by the user, complete flow field outputs are provided by SUBROUTINE DMPOUT. This printed output may be nondimensional or dimensional (see case data variables P1, RHO1, T1 and RN). The following nomenclature applies to the printed output:

TIME	- nondimensional time.
S	- streamwise distance.
N	- distance normal to the body, directed away from the body.
RHO	- density.
U	- streamwise velocity component.
V	- velocity component normal to the body and directed outward from the body.
(S-S1)/CV	- nondimensional entropy.
P	- pressure.
M	- Mach number.
T	- temperature.
DH/H1	- error in total enthalpy.
RB	- body radius.
THETA	- θ
K	- body curvature.
ZB	- value of Z at a body point.

- RS - shock radius.
- ZS - value of z at the shock point.
- WS - shock velocity, directed outward along the normal coordinate.
- CP - pressure coefficient at a body point.
- P/PO - ratio of the body pressure to the stagnation pressure behind the normal shock.

When the nodal spacing normal to the body is computed by the computer code, the computed number of nodes can exceed the dimension sizes of the arrays. When this occurs, the maximum number allowed will be used and a printed message will appear in the standard output. Currently, up to 40 nodes across the shock layer can be used. This number has been found to be adequate for typical problems even when the empirical relations predict that more are required. When this number of nodes is exceeded by the empirical determination, the shock layer is quite thick and the gradients of the flow parameters normal to the body are, generally, quite small.

9. SAMPLE PROBLEMS

Input data and computer outputs for 2 sample problems are provided on the following pages. The first problem considers a Mach 6 flow about an elliptically blunted cylinder with a nose bluntness parameter of 2.25. This case utilizes the arbitrary body shape option. The second case considers a Mach 4.0 flow about a sphere-cone vehicle with a cone half-angle of 40° . The internal logic for sphere-cone geometry is used for this case. These 2 cases were run as a single multicase job.

TITLE SAMPLE PROBLEM 1 - ELLIPSE-CYLINDER - B/A=1.5

M 6.0
 NSEG 3
 LAST 2.25

	1	15	10	
1.	0.	3.0000E	014.4444E-01	
7.4534E-01	21.7441E-03	3.7549E+01	4.4592E-01	
1.4904E-01	7.9762E-03	4.5154E+01	4.5039E-01	
2.4361E-01	1.4037E-02	4.2645E+01	4.5402E-01	
3.7415E-01	7.2299E-02	4.0147E+01	4.6707E-01	
4.7263E-01	5.0950E-02	7.7514E+01	4.8393E-01	
5.6723E-01	7.4255E-02	7.4766E+01	5.0320E-01	
6.6176E-01	1.0258E-01	7.1954E+01	5.2767E-01	
7.5637E-01	1.3641E-01	6.8732E+01	5.5849E-01	
8.5094E-01	1.7644E-01	6.5337E+01	5.9718E-01	
9.4534E-01	2.2361E-01	6.1579E+01	6.4601E-01	
1.0399E+00	2.7933E-01	5.7327E+01	7.0524E-01	
1.1345E+00	3.4579E-01	5.2373E+01	7.8872E-01	
1.2290E+00	4.2667E-01	4.6397E+01	8.9507E-01	
1.3235E+00	5.2941E-01	3.8650E+01	1.0397E+00	

	2	14	6	
1.3476E+00	5.6078E-01	3.6254E+01	1.0949E+00	
1.3696E+00	5.9216E-01	3.3923E+01	1.1299E+00	
1.3956E+00	6.2353E-01	3.1364E+01	1.1745E+00	
1.4274E+00	6.5490E-01	2.8479E+01	1.2181E+00	
1.4643E+00	6.8627E-01	2.6363E+01	1.2604E+00	
1.4990E+00	7.1765E-01	2.3921E+01	1.3009E+00	
1.5220E+00	7.4902E-01	2.1252E+01	1.3388E+00	
1.4634E+00	7.8039E-01	1.8657E+01	1.3739E+00	
1.4732E+00	8.1176E-01	1.6040E+01	1.4056E+00	
1.4914E+00	8.4314E-01	1.3401E+01	1.4334E+00	
1.4991E+00	8.7451E-01	1.0743E+01	1.4568E+00	
1.4933E+00	9.0588E-01	8.0711E+00	1.4754E+00	
1.4971E+00	9.3725E-01	5.3873E+00	1.4890E+00	
1.5000E+00	9.6862E-01	2.7035E+00	1.5000E+00	

	3	15	6	
1.5000E+00	1.0000E+00	0.0000E+00	0.	
1.5000E+00	1.0296E+00	0.0000E+00	0.	
1.5000E+00	1.0571E+00	0.0000E+00	0.	
1.5000E+00	1.0857E+00	0.0000E+00	0.	
1.5000E+00	1.1143E+00	0.0000E+00	0.	
1.5000E+00	1.1429E+00	0.0000E+00	0.	
1.5000E+00	1.1714E+00	0.0000E+00	0.	
1.5000E+00	1.2000E+00	0.0000E+00	0.	
1.5000E+00	1.2286E+00	0.0000E+00	0.	
1.5000E+00	1.2571E+00	0.0000E+00	0.	
1.5000E+00	1.2857E+00	0.0000E+00	0.	
1.5000E+00	1.3143E+00	0.0000E+00	0.	
1.5000E+00	1.3429E+00	0.0000E+00	0.	
1.5000E+00	1.3714E+00	0.0000E+00	0.	
1.5000E+00	1.4000E+00	0.0000E+00	0.	

TITLE SAMPLE PROBLEM 2 - SPHERE-CONE VEHICLE

M 4.0
 Z/RN 1.4
 THETA 40.0
 NSEG 0
 LAST

CHANGE K

AFWL TIME-DEPENDENT FLOW FIELD PROGRAM FOR AXISYMMETRIC BLUNT BODIES
 AFWL-TR-70-16, R.M. AUNGIER

CASE TITLE - SAMPLE PROBLEM 1 - ELLIPSE-CYLINDER - B/A=1.5

```

***** INPUT DATA *****
INTAP      2      OTAP      1      MAXIT      -1
NCASE      1      NFILE      1      NSEG      3
MAXI       0      MAXN       0      NTAP      0
IPR(1)    -5      IPR(2)    -5      IPR(3)    -5
IPR(4)    -5      IPR(5)    -5      IPR(6)    -5
GAMMA     1.400E+00  M      0.000E+00  THETA     0.
Z/RN      0.      TCRIT1    2.000E+00  TCRIT2    1.000E+01
P1        1.000E+00  RH01     1.000E+00  RN        1.000E+00
MU        4.000E-02  T1       1.000E+00  ERRORN    1.000E-05
ERRORP    1.000E-08
  
```

Table with columns: S, R, Z, T, N, RMO, U, V, W, P, I, DM/MI. Includes values for S, R, Z, T, N, RMO, U, V, W, P, I, DM/MI.

Table with columns: S, R, Z, T, N, RMO, U, V, W, P, I, DM/MI. Includes values for S, R, Z, T, N, RMO, U, V, W, P, I, DM/MI.

Table with columns: S, R, Z, T, N, RMO, U, V, W, P, I, DM/MI. Includes values for S, R, Z, T, N, RMO, U, V, W, P, I, DM/MI.

Table with columns: S, R, Z, T, N, RMO, U, V, W, P, I, DM/MI. Includes values for S, R, Z, T, N, RMO, U, V, W, P, I, DM/MI.

SZ RJ# ZJ# IMETA= K#	.73 .32 .05 09.92 1.23	0. 1.194/E-02 2.377/E-02 3.700/E-02 4.754/E-02 5.443/E-02 7.132/E-02 9.322/E-02 9.203/E-02 1.181/E-01 1.307/E-01 1.420/E-01	RHO 2.0731E+00 3.1057E+00 5.1399E+00 7.1672E+00 9.1918E+00 1.1233E+01 3.2271E+01 5.2372E+01 7.2494E+01 9.2409E+01 1.1017E+02 1.3075E+02 1.4204E+02	J 1.0225E+00 1.5245E+00 1.901E+00 1.7304E+00 1.703E+00 1.4231E+00 1.4744E+00 1.330E+00 1.4479E+00 2.441E+00 2.1917E+00 2.1807E+00 2.2207E+00	V -0. -1.0292E-01 -2.0733E-01 -3.090E-01 -4.0995E-01 -5.0877E-01 -6.0727E-01 -7.0522E-01 -8.0414E-01 -9.0305E-01 -1.0195E-00 -1.1104E-00 -1.2119E-00	(S-S1)/CV 1.4073E+00 1.4073E+00 1.4073E+00 1.3958E+00 1.3890E+00 1.3890E+00 1.3767E+00 1.3700E+00 1.3631E+00 1.3562E+00 1.3491E+00 1.3417E+00 1.3342E+00 1.3262E+00	P 3.4679E+01 3.4609E+01 3.4948E+01 4.0032E+01 4.0013E+01 3.4904E+01 3.9800E+01 3.4696E+01 3.9470E+01 3.9174E+01 3.8825E+01 3.8342E+01 3.7915E+01	M 4.4123E-01 5.1325E-01 5.1333E-01 5.3573E-01 5.5029E-01 5.7930E-01 6.0449E-01 6.3105E-01 6.0445E-01 6.9058E-01 7.2337E-01 7.5726E-01 7.9212E-01	T 7.121E+00 7.7471E+00 7.7732E+00 7.7732E+00 7.7414E+00 7.7081E+00 7.0031E+00 7.0270E+00 7.0270E+00 7.0270E+00 7.0470E+00 7.0470E+00 7.0470E+00	DM/M1 .0011J .0011J .0011J .0011J .0011J .0011J .0011J .0011J .0011J .0011J .0011J .0011J .0011J
SZ RJ# ZJ# IMETA= K#	.41 .40 .03 01.42 1.34	0. 1.2345E-02 2.4770E-02 3.7155E-02 4.9540E-02 6.1925E-02 7.4310E-02 8.6694E-02 9.9079E-02 1.1144E-01 1.2345E-01 1.3623E-01 1.4802E-01	RHO 4.0454E+00 4.7749E+00 4.8101E+00 4.7044E+00 4.9144E+00 4.4949E+00 4.9992E+00 5.3357E+00 5.0688E+00 5.0947E+00 5.1244E+00 5.1452E+00 5.1641E+00	U 2.1237E+00 2.1041E+00 2.2030E+00 2.2436E+00 2.2804E+00 2.3231E+00 2.3603E+00 2.4032E+00 2.4463E+00 2.4894E+00 2.5325E+00 2.5756E+00 2.6187E+00	V -0. -9.0947E-02 -1.9215E-01 -2.1237E-01 -3.1409E-01 -4.1581E-01 -5.1753E-01 -6.1925E-01 -7.2097E-01 -8.2269E-01 -9.2441E-01 -1.0261E-00 -1.0815E+00	(S-S1)/CV 1.4073E+00 1.3975E+00 1.3877E+00 1.3779E+00 1.3681E+00 1.3583E+00 1.3485E+00 1.3387E+00 1.3289E+00 1.3191E+00 1.3093E+00 1.2995E+00 1.2897E+00	P 3.2500E+01 3.5122E+01 3.0120E+01 3.0359E+01 3.0598E+01 3.6673E+01 3.6756E+01 3.0782E+01 3.0782E+01 3.0782E+01 3.0782E+01 3.0782E+01 3.0782E+01	M 0.2302E-01 0.0005E-01 0.4211E-01 0.9912E-01 7.1810E-01 7.3822E-01 7.6070E-01 7.8494E-01 8.1100E-01 8.4117E-01 8.7134E-01 9.0151E-01 9.3168E-01	T 7.0700E+00 7.0200E+00 7.0200E+00 7.0200E+00 7.0200E+00 7.0200E+00 7.0200E+00 7.0200E+00 7.0200E+00 7.0200E+00 7.0200E+00 7.0200E+00 7.0200E+00	DM/M1 .0011J .0011J .0011J .0011J .0011J .0011J .0011J .0011J .0011J .0011J .0011J .0011J .0011J
SZ RJ# ZJ# IMETA= K#	.57 .55 .16 4.424 1.93	0. 1.4041E-02 2.4011E-02 3.4022E-02 4.4033E-02 5.4044E-02 6.4055E-02 7.4066E-02 8.4077E-02 9.4088E-02 1.0409E-01 1.0409E-01 1.0409E-01	RHO 3.5242E+00 3.7230E+00 3.4574E+00 3.9742E+00 4.0909E+00 4.2076E+00 4.3243E+00 4.4410E+00 4.5577E+00 4.6744E+00 4.7911E+00 4.9078E+00 5.0245E+00	U 3.1305E+00 3.1479E+00 3.1653E+00 3.2028E+00 3.2403E+00 3.2778E+00 3.3153E+00 3.3528E+00 3.3903E+00 3.4278E+00 3.4653E+00 3.5028E+00 3.5403E+00	V -0. -4.7223E-02 -1.0072E-01 -1.0072E-01 -1.0072E-01 -1.0072E-01 -1.0072E-01 -1.0072E-01 -1.0072E-01 -1.0072E-01 -1.0072E-01 -1.0072E-01 -1.0072E-01	(S-S1)/CV 1.4073E+00 1.3904E+00 1.3734E+00 1.3564E+00 1.3394E+00 1.3224E+00 1.3054E+00 1.2884E+00 1.2714E+00 1.2544E+00 1.2374E+00 1.2204E+00 1.2034E+00	P 2.4476E+01 2.7338E+01 2.0114E+01 2.0793E+01 2.0793E+01 2.7476E+01 2.4944E+01 2.4944E+01 2.4944E+01 2.4944E+01 2.4944E+01 2.4944E+01 2.4944E+01	M 1.0133E+00 1.0202E+00 1.0271E+00 1.0340E+00 1.0409E+00 1.0478E+00 1.0547E+00 1.0616E+00 1.0685E+00 1.0754E+00 1.0823E+00 1.0892E+00 1.0961E+00	T 0.5130E+00 0.7473E+00 6.7703E+00 6.7703E+00 6.7703E+00 6.7703E+00 6.7703E+00 6.7703E+00 6.7703E+00 6.7703E+00 6.7703E+00 6.7703E+00 6.7703E+00	DM/M1 0M/M1 0.0021J .0010J .0010J 0.0124 0.0134

S	67	N	PHO	U	V	(S-S1)/CV	P	M	T	DM/MI
0.	.54	0.	2.0492E+00	3.7022E+00	-0.	1.4073E+00	1.5978E+01	1.5429E+00	0.3314E+00	.00041
1.6341E-01	.24	1.0712E-02	2.4461E+00	3.8443E+00	5.4770E-02	1.3085E+00	1.7336E+01	1.3106E+00	0.0312E+00	.00030
2.4212E-01	54.00	3.5714E-02	3.0428E+00	3.3024E+00	7.0755E-02	1.3685E+00	1.1659E+01	1.2974E+00	0.1322E+00	.00111
3.2042E-01	2.34	5.0376E-02	3.2224E+00	3.7837E+00	4.1818E-02	1.3469E+00	1.4748E+01	1.2907E+00	0.1412E+00	.00101
4.0853E-01		0.7104E-02	3.4024E+00	3.7474E+00	4.1369E-02	1.3232E+00	2.0452E+01	1.2933E+00	0.1217E+00	.00271
4.923E-01		4.3960E-02	3.5807E+00	3.3112E+00	7.3925E-02	1.2972E+00	2.1823E+01	1.3050E+00	0.1442E+00	.00301
5.7194E-01		1.0075E-01	3.7632E+00	3.6579E+00	7.2692E-02	1.2688E+00	2.2742E+01	1.3245E+00	0.1413E+00	.00444
6.5305E-01		1.1754E-01	3.9506E+00	3.3023E+00	0.5925E-02	1.2308E+00	2.3619E+01	1.3200E+00	0.1777E+00	.00512
		1.3434E-01	4.1416E+00	3.9694E+00	5.6013E-02	1.2070E+00	2.4447E+01	1.3609E+00	0.3023E+00	.00501
		1.5113E-01	4.3422E+00	4.0407E+00	4.7223E-02	1.1742E+00	2.5273E+01	1.4154E+00	0.3217E+00	.00577
		1.6792E-01	4.5392E+00	4.1238E+00	3.0216E-02	1.1404E+00	2.6094E+01	1.4502E+00	0.3249E+00	.00627
		1.8471E-01	4.7347E+00	4.2129E+00	2.3066E-02	1.1050E+00	2.6879E+01	1.4941E+00	0.3541E+00	.00634
		2.0152E-01	4.9500E+00	4.3296E+00	-1.4401E-02	1.0690E+00	2.7334E+01	1.5372E+00	0.3822E+00	.00700

SMUCK AND BUDY PARAMETERS

S	68	THETA	K	MS	ZS	CP	P/PU
0.	0.	9.0000E+00	1.0000E+00	-1.1408E-14	-1.3576E-01	1.4241E+00	1.0032E+00
1.6341E-01	3.3431E-02	6.5239E+01	1.0126E+00	9.2825E-02	-1.3299E-01	1.6010E+00	9.9373E-01
2.4212E-01	1.3434E-01	3.0477E+01	1.0518E+00	1.8544E-01	-1.2231E-01	1.7300E+00	9.0334E-01
3.2042E-01	3.0402E-02	7.5404E+01	1.1214E+00	2.7781E-01	-1.0492E-01	1.8614E+00	9.1580E-01
4.0853E-01	5.4745E-01	6.9919E+01	1.2231E+00	3.6951E-01	-7.9223E-02	1.9534E+00	8.4750E-01
4.923E-01	8.6740E-02	0.3823E+01	1.3804E+00	4.6124E-01	-4.0626E-02	2.0369E+00	7.5930E-01
5.7194E-01	1.2716E-01	5.0033E+01	1.5107E+00	5.5290E-01	-3.1874E-03	2.1310E+00	6.3010E-01
6.5305E-01	1.7637E-01	4.8595E+01	1.6237E+00	6.4316E-01	-3.5425E-03	2.2315E+00	5.2202E-01
	2.3520E-01	5.8660E+01	2.3333E+00	7.4555E-01	-4.5911E-03	2.3443E+00	3.4131E-01

S	RH	ZH	META	K	THETA	RS	ZS	MS	CP	P/PU	DM/M1
7.1092E-01	0.2073E-01	2.8237E-01	3.0429E+01	3.0472E+00	3.0429E+01	6.1792E-01	1.6655E-01	-1.6403E-04	4.0460E-01	2.3415E-01	.0J33
7.6920E-01	0.4509E-01	3.3396E-01	2.1078E+01	3.0187E+00	2.1078E+01	9.0180E-01	2.3514E-01	2.5619E-04	2.3424E-01	1.4960E-01	.0J32
8.2546E-01	0.6139E-01	3.3412E-01	1.0953E+01	3.2874E+00	1.0953E+01	9.8227E-01	3.2675E-01	7.7719E-04	1.1761E-01	0.3594E-02	.0J23
8.8276E-01	0.6737E-01	4.44574E-01	-1.3344E-00	3.3750E+00	-1.3344E-00	1.1090E+00	4.4574E-01	-4.5285E-04	4.2869E-02	4.4323E-02	.0J22
											.0J21
											.0J20
											.0J19
											.0J18
											.0J17
											.0J16
											.0J15
											.0J14
											.0J13
											.0J12
											.0J11
											.0J10
											.0J09
											.0J08
											.0J07
											.0J06
											.0J05
											.0J04
											.0J03
											.0J02
											.0J01

SHOCK AND BODY PARAMETERS

S	RH	ZH	META	K	THETA	RS	ZS	MS	CP	P/PU	DM/M1
7.1092E-01	0.2073E-01	2.8237E-01	3.0429E+01	3.0472E+00	3.0429E+01	6.1792E-01	1.6655E-01	-1.6403E-04	4.0460E-01	2.3415E-01	.0J33
7.6920E-01	0.4509E-01	3.3396E-01	2.1078E+01	3.0187E+00	2.1078E+01	9.0180E-01	2.3514E-01	2.5619E-04	2.3424E-01	1.4960E-01	.0J32
8.2546E-01	0.6139E-01	3.3412E-01	1.0953E+01	3.2874E+00	1.0953E+01	9.8227E-01	3.2675E-01	7.7719E-04	1.1761E-01	0.3594E-02	.0J23
8.8276E-01	0.6737E-01	4.44574E-01	-1.3344E-00	3.3750E+00	-1.3344E-00	1.1090E+00	4.4574E-01	-4.5285E-04	4.2869E-02	4.4323E-02	.0J22
											.0J21
											.0J20
											.0J19
											.0J18
											.0J17
											.0J16
											.0J15
											.0J14
											.0J13
											.0J12
											.0J11
											.0J10
											.0J09
											.0J08
											.0J07
											.0J06
											.0J05
											.0J04
											.0J03
											.0J02
											.0J01

S= .91
 RJ= .07
 ZJ= .41
 IMETA= 0.
 KE= J.

S	RHO	U	V	(S-S1)/CV	P	M	I	U/M
0.	2.1109E-01	5.4223E+00	-0.	1.4073E+00	1.0842E+00	2.0174E+00	3.17E+00	.0101
2.271E-12	7.1000E-01	5.7007E+00	2.6579E-01	1.3961E+00	2.0009E+00	2.0372E+00	3.0224E+00	.0120
4.3432E-12	1.0004E-01	5.0272E+00	6.4274E-01	1.3833E+00	3.2234E+00	2.4230E+00	3.7541E+00	.0112
6.4144E-12	1.0157E+00	5.3539E+00	9.7320E-01	1.3634E+00	4.7135E+00	2.3101E+00	3.4310E+00	.0111
9.0404E-12	1.1591E+00	5.2133E+00	1.1923E+00	1.3530E+00	4.7430E+00	2.2533E+00	3.4324E+00	.0123
1.1354E-11	1.2947E+00	5.1307E+00	1.7824E+00	1.3295E+00	5.4439E+00	2.1706E+00	4.2337E+00	.0134
1.3670E-11	1.4350E+00	4.9494E+00	1.5023E+00	1.3067E+00	6.1247E+00	2.1377E+00	4.2072E+00	.0140
1.5911E-11	1.5894E+00	4.7232E+00	1.0191E+00	1.2816E+00	6.7774E+00	2.1040E+00	4.1311E+00	.0147
1.8173E-11	1.7040E+00	4.4606E+00	1.7233E+00	1.2545E+00	7.3977E+00	2.0732E+00	4.0493E+00	.0154
2.0444E-11	1.7474E+00	4.1213E+00	1.1177E+00	1.2253E+00	7.9917E+00	2.0407E+00	4.0340E+00	.0161
2.2711E-11	1.7303E+00	3.7717E+00	1.4938E+00	1.1944E+00	8.5902E+00	2.0093E+00	4.0377E+00	.0168
2.4974E-11	1.2608E+00	4.7747E+00	1.4710E+00	1.1604E+00	9.1733E+00	2.1010E+00	4.3141E+00	.0175
2.7254E-11	2.2401E+00	4.7604E+00	2.0401E+00	1.1239E+00	9.7554E+00	2.1200E+00	4.2737E+00	.0182
2.9531E-11	2.4456E+00	4.7712E+00	2.1184E+00	1.0843E+00	1.0345E+01	2.1473E+00	4.2232E+00	.0189
3.1812E-11	2.6252E+00	4.7827E+00	2.1902E+00	1.0418E+00	1.0945E+01	2.1773E+00	4.1842E+00	.0196
3.4074E-11	2.4204E+00	4.4011E+00	2.2606E+00	9.9811E-01	1.1502E+01	2.2131E+00	4.1331E+00	.0203
3.6345E-11	3.0335E+00	4.4243E+00	2.3302E+00	9.6840E-01	1.2207E+01	2.2517E+00	4.1243E+00	.0210
3.8617E-11	3.2658E+00	4.3523E+00	2.3992E+00	9.4890E-01	1.2842E+01	2.3035E+00	3.9444E+00	.0217
4.0894E-11	3.7516E+00	4.0852E+00	2.4659E+00	9.4804E-01	1.3579E+01	2.3756E+00	3.3614E+00	.0224
4.3166E-11	3.4010E+00	4.4201E+00	2.5388E+00	7.9588E-01	1.4372E+01	2.4404E+00	3.7413E+00	.0231
4.5432E-11	4.1003E+00	4.4031E+00	2.6063E+00	7.4104E-01	1.5124E+01	2.4907E+00	3.0847E+00	.0238
4.7703E-11	4.3470E+00	5.0264E+00	2.6574E+00	6.8349E-01	1.5699E+01	2.5411E+00	3.5733E+00	.0245

S= .91
 RJ= .67
 ZJ= .53
 IMETA= 0.
 KE= J.

S	RHO	U	V	(S-S1)/CV	P	M	I	U/M
0.	5.2032E-01	5.9323E+00	-0.	1.4073E+00	1.0366E+00	2.0271E+00	3.1474E+00	.0141
2.4317E-12	6.5924E-01	5.7704E+00	2.3890E-01	1.3955E+00	2.0099E+00	2.0547E+00	3.3402E+00	.0132
4.4014E-12	7.2224E-01	5.5541E+00	4.1772E-01	1.3833E+00	2.0772E+00	2.0842E+00	3.5402E+00	.0144
7.2927E-12	8.9062E-01	5.5173E+00	7.0312E-01	1.3634E+00	3.3404E+00	2.4272E+00	4.7513E+00	.0130
9.7209E-12	1.0222E+00	5.3549E+00	9.4603E-01	1.3509E+00	3.9817E+00	2.3412E+00	3.1952E+00	.0142
1.2154E-11	1.1541E+00	5.2687E+00	1.1453E+00	1.3304E+00	4.0252E+00	2.2703E+00	4.0070E+00	.0141
1.4702E-11	1.2850E+00	5.1049E+00	1.3097E+00	1.3042E+00	4.2233E+00	2.2034E+00	4.2413E+00	.0147
1.7022E-11	1.4160E+00	5.1463E+00	1.4482E+00	1.2831E+00	4.5713E+00	2.1973E+00	4.1402E+00	.0153
1.9454E-11	1.5444E+00	5.0230E+00	1.5680E+00	1.2556E+00	4.4735E+00	2.1773E+00	4.1402E+00	.0167
2.1885E-11	1.6736E+00	4.9736E+00	1.6745E+00	1.2257E+00	4.7049E+00	2.1607E+00	4.1902E+00	.0160
2.4317E-11	1.8234E+00	4.9374E+00	1.7711E+00	1.1936E+00	7.0649E+00	2.1044E+00	4.1944E+00	.0160
2.6749E-11	1.9699E+00	4.8139E+00	1.8606E+00	1.1588E+00	6.2316E+00	2.1724E+00	4.1740E+00	.0167
2.9181E-11	2.1254E+00	4.7016E+00	1.9441E+00	1.1204E+00	5.1107E+00	2.1002E+00	4.1474E+00	.0174
3.1612E-11	2.2940E+00	4.4946E+00	2.1263E+00	1.0795E+00	3.4114E+00	2.2233E+00	4.1727E+00	.0179
3.4044E-11	2.4771E+00	4.3002E+00	2.1052E+00	1.0349E+00	1.0023E+01	2.2403E+00	4.0401E+00	.0186
3.6476E-11	2.6772E+00	4.3204E+00	2.1824E+00	9.8729E-01	1.0675E+01	2.2000E+00	3.9747E+00	.0194
3.8907E-11	2.4964E+00	4.1412E+00	2.2543E+00	9.6729E-01	1.1315E+01	2.3201E+00	3.3002E+00	.0201
4.1334E-11	3.1365E+00	4.3041E+00	2.3334E+00	9.3720E-01	1.2038E+01	2.3034E+00	3.4244E+00	.0208
4.3771E-11	3.1365E+00	4.3041E+00	2.3334E+00	9.3720E-01	1.2038E+01	2.3034E+00	3.4244E+00	.0215
4.6203E-11	3.3986E+00	4.3462E+00	2.4074E+00	8.3137E-01	1.2731E+01	2.4223E+00	3.7401E+00	.0222
4.8634E-11	3.6319E+00	5.0343E+00	2.4852E+00	7.7534E-01	1.3514E+01	2.4725E+00	3.0614E+00	.0229
5.1066E-11	4.1122E+00	5.0774E+00	2.5611E+00	7.1631E-01	1.4312E+01	2.5447E+00	3.0071E+00	.0236
5.3498E-11	4.3275E+00	5.1370E+00	2.6192E+00	6.5216E-01	1.4927E+01	2.6242E+00	3.4442E+00	.0243

S= 1.02
RB= .67
ZB= .53
TMLTA= 0.
K= 0.

S= 1.00
RB= .67
ZB= .52
TMLTA= 0.
K= 0.

U.	1.0340E-01	2.9031E+00	1.4073E+00	1.4073E+00	15-S11/CV	1.4073E+00	1.209E+00	2.7091E+00	3.1971E+00	DM/MI
1.	2.5747E-02	5.9371E+00	1.3945E+00	1.3945E+00	1.4073E+00	1.3945E+00	2.0755E+00	2.0800E+00	3.1527E+00	.0353
5.	5.1719E-02	5.7303E+00	1.3825E+00	1.3825E+00	1.3825E+00	1.3825E+00	2.3335E+00	2.6266E+00	3.4119E+00	.0063
7.	7.7634E-02	5.0242E+00	4.4767E-01	1.3679E+00	1.3679E+00	1.3679E+00	2.4900E+00	2.7125E+00	3.5972E+00	.0062
1.	1.0340E-01	5.5032E+00	7.4758E-01	1.3508E+00	1.3508E+00	1.3508E+00	3.4350E+00	2.4312E+00	3.7311E+00	.0020
1.	1.2932E-01	5.4014E+00	9.5532E-01	1.3309E+00	1.3309E+00	1.3309E+00	4.0111E+00	2.3120E+00	3.9372E+00	.0014
1.	1.5522E-01	5.3000E+00	1.1305E+00	1.3084E+00	1.3084E+00	1.3084E+00	4.7878E+00	2.2702E+00	3.4973E+00	.0073
1.	1.8104E-01	5.2241E+00	1.4291E+00	1.2872E+00	1.2872E+00	1.2872E+00	5.1632E+00	2.2534E+00	4.0321E+00	.0053
2.	2.0610E-01	5.1573E+00	1.8206E+00	1.2555E+00	1.2555E+00	1.2555E+00	5.7326E+00	2.2304E+00	4.3321E+00	.0047
2.	2.3118E-01	5.1041E+00	1.9448E+00	1.2253E+00	1.2253E+00	1.2253E+00	6.4713E+00	2.2333E+00	4.5053E+00	.0104
2.	2.5470E-01	5.0533E+00	1.8547E+00	1.1924E+00	1.1924E+00	1.1924E+00	7.4215E+00	2.2384E+00	4.6202E+00	.0101
2.	2.7471E-01	5.0330E+00	1.7544E+00	1.1765E+00	1.1765E+00	1.1765E+00	8.7215E+00	2.2756E+00	3.3450E+00	.0127
3.	3.0441E-01	5.0201E+00	1.4479E+00	1.1172E+00	1.1172E+00	1.1172E+00	1.0744E+00	2.2298E+00	3.3232E+00	.0107
3.	3.3631E-01	5.0144E+00	1.3369E+00	1.0744E+00	1.0744E+00	1.0744E+00	9.2263E+00	2.3070E+00	3.3064E+00	.0106
3.	3.7181E-01	5.1171E+00	2.0229E+00	1.0281E+00	1.0281E+00	1.0281E+00	9.6629E+00	2.3428E+00	3.7976E+00	.0063
3.	4.0805E-01	5.0292E+00	2.1065E+00	9.7860E-01	9.7860E-01	9.7860E-01	1.1529E+00	2.3859E+00	3.7200E+00	.0067
4.	4.4392E-01	5.0472E+00	2.1846E+00	9.2045E-01	9.2045E-01	9.2045E-01	1.1230E+00	2.4344E+00	3.6562E+00	.0067
4.	4.8392E-01	5.0712E+00	2.2699E+00	8.7201E-01	8.7201E-01	8.7201E-01	1.1230E+00	2.4809E+00	3.5506E+00	.0042
4.	5.2889E-01	5.1015E+00	2.3537E+00	8.1526E-01	8.1526E-01	8.1526E-01	1.1966E+00	2.4809E+00	3.4531E+00	.0033
4.	5.7153E-01	5.1372E+00	2.4341E+00	7.5561E-01	7.5561E-01	7.5561E-01	1.2754E+00	2.5478E+00	3.4531E+00	.0033
5.	5.1734E-01	5.1766E+00	2.5168E+00	6.9175E-01	6.9175E-01	6.9175E-01	1.3573E+00	2.6187E+00	3.4531E+00	.0033
5.	5.4326E-01	5.2376E+00	2.5747E+00	6.2257E-01	6.2257E-01	6.2257E-01	1.4217E+00	2.7038E+00	3.3304E+00	.0000

N	0.	5.8627E+00	1.5-S11/CV	1.4073E+00	1.6506E+00	2.7497E+00	DM/MI
2.	2.7347E-02	5.6221E+00	1.4073E+00	1.4073E+00	1.4073E+00	2.7097E+00	.0351
5.	5.4693E-02	5.7918E+00	1.3934E+00	1.3934E+00	1.3934E+00	2.7097E+00	.0062
7.	7.4495E-02	5.7028E+00	1.3813E+00	1.3813E+00	1.3813E+00	2.6824E+00	.00917
1.	1.0399E-01	5.6003E+00	1.3666E+00	1.3666E+00	1.3666E+00	2.5827E+00	.00636
1.	1.3073E-01	5.5008E+00	1.3496E+00	1.3496E+00	1.3496E+00	2.5006E+00	.00603
1.	1.6408E-01	5.4311E+00	1.3294E+00	1.3294E+00	1.3294E+00	2.4411E+00	.00744
1.	1.9143E-01	5.3401E+00	1.3074E+00	1.3074E+00	1.3074E+00	2.3872E+00	.00814
2.	2.1875E-01	5.2731E+00	1.2821E+00	1.2821E+00	1.2821E+00	2.3477E+00	.00661
2.	2.4612E-01	5.2100E+00	1.2541E+00	1.2541E+00	1.2541E+00	2.3209E+00	.00942
2.	2.7347E-01	5.1763E+00	1.2234E+00	1.2234E+00	1.2234E+00	2.3048E+00	.00992
3.	3.0041E-01	5.1402E+00	1.1988E+00	1.1988E+00	1.1988E+00	2.2908E+00	.01033
3.	3.2816E-01	5.1272E+00	1.1730E+00	1.1730E+00	1.1730E+00	2.3022E+00	.01072
3.	3.5551E-01	5.1167E+00	1.1455E+00	1.1455E+00	1.1455E+00	2.3150E+00	.01068
3.	3.8285E-01	5.1195E+00	1.1204E+00	1.1204E+00	1.1204E+00	2.3366E+00	.01037
4.	4.1020E-01	5.1264E+00	1.0920E-01	1.0920E-01	1.0920E-01	2.3603E+00	.00963
4.	4.3755E-01	5.1444E+00	9.6920E-01	9.6920E-01	9.6920E-01	2.4033E+00	.00963
4.	4.6494E-01	5.1670E+00	9.1509E-01	9.1509E-01	9.1509E-01	2.4409E+00	.00632
4.	4.9224E-01	5.1900E+00	8.5834E-01	8.5834E-01	8.5834E-01	2.4909E+00	.00630
5.	5.1954E-01	5.2371E+00	7.9481E-01	7.9481E-01	7.9481E-01	2.5335E+00	.00370
5.	5.4693E-01	5.2740E+00	7.3578E-01	7.3578E-01	7.3578E-01	2.5814E+00	.00370
5.	5.7429E-01	5.3260E+00	6.8202E-01	6.8202E-01	6.8202E-01	2.6344E+00	.00318
5.	6.0168E-01	5.3800E+00	6.3574E-01	6.3574E-01	6.3574E-01	2.6936E+00	.0000
5.	6.2907E-01	5.4350E+00	5.9574E-01	5.9574E-01	5.9574E-01	2.7700E+00	.0000

SHOCK AND BODY PARAMETERS

S	9.2721E-01	R3	4.8886E-01	Z3	0.	K	0.	MS	4.7019E-07	CP	2.7153E-02	P/PJ	3.5970E-02
	9.7167E-01		5.3333E-01		0.		0.		3.0975E-06		2.5203E-02		3.4760E-02
	1.0161E+00		5.7777E-01		0.		0.		-1.9347E-06		2.1740E-02		3.0184E-02
	1.0005E+00		6.2222E-01		0.		0.		6.2222E-01		3.1747E-02		3.9520E-02

CASE TITLE - SAMPLE PROBLEM 2 - SPHERE-CONE VEHICLE

***** INPUT DATA *****

INTAP	2	UTAP	1	MAXIT	-1
NCASE	1	NFILL	1	NSEG	0
MAXI	0	MAXN	0	NTAP	0
IPR(1)	-5	IPR(2)	-5	IPR(3)	-5
IPR(4)	-5	IPR(5)	-5	IPR(6)	-5
MAIMA	1.4000E+00	M	4.0000E+00	THTA	+.0000E+01
Z/RN	1.4000E+00	TCRIT1	2.0000E+00	TCRIT2	1.0000E+01
Q1	1.0000E+00	RF01	1.0000E+00	RN	1.0000E+00
MU	4.0000E-02	T1	1.0000E+00	ERRRN	1.0000E-05
ERRORP	1.0000E-06				

STATEMENT NUMBER 1		ITERATION NUMBER 304		TIME= 4.46/4E+00	
S=	0.	RHO	0.	(S-S1)/CV	7.9001E-01
RB=	0.	U	0.	V	-0.
ZJ=	0.	PHO	5.0344E+00	V	-1.7524E-01
IMETA=	10.00	N	2.9632E-02	V	-3.5174E-01
K=	1.00	U	4.9149E+00	V	-5.1831E-01
		N	4.4849E-02	V	-8.3525E-01
		U	1.1857E-01	V	-8.5522E-01
		N	1.4416E-01	V	-1.1353E+00
		U	1.7774E-01	V	-0.
S=	.12	RHO	4.9278E+00	(S-S1)/CV	7.9001E-01
RB=	.01	U	4.9140E+00	V	-0.
ZJ=	.01	N	3.5573E-02	V	-1.6724E-01
IMETA=	11.67	U	6.1146E-02	V	-3.4072E-01
K=	1.00	N	9.1718E-02	V	-5.3060E-01
		U	4.8405E+00	V	-8.7617E-01
		N	1.2229E-01	V	-8.4511E-01
		U	1.5290E-01	V	-7.8322E-01
		N	1.4344E-01	V	-7.4130E-01
S=	.23	RHO	4.6752E+00	(S-S1)/CV	7.9001E-01
RB=	.04	U	4.6752E+00	V	-0.
ZJ=	.04	N	3.1775E-02	V	-1.5571E-01
IMETA=	73.33	U	6.3549E-02	V	-3.2530E-01
K=	1.00	N	9.5324E-02	V	-4.8190E-01
		U	1.2710E-01	V	-6.4791E+00
		N	1.5847E-01	V	-8.0408E-01
		U	1.9065E-01	V	-9.5907E-01
S=	.44	RHO	4.2897E+00	(S-S1)/CV	7.9001E-01
RB=	.04	U	4.3544E+00	V	-0.
ZJ=	.04	N	6.4071E-02	V	-1.4577E-01
IMETA=	65.00	U	1.0211E-01	V	-2.9514E-01
K=	1.00	N	1.3614E-01	V	-4.3373E-01
		U	1.7018E-01	V	-5.7358E-01
		N	2.0421E-01	V	-7.1727E-01
S=	.56	RHO	3.7821E+00	(S-S1)/CV	7.9001E-01
RB=	.16	U	3.8997E+00	V	-0.
ZJ=	.16	N	7.4754E-02	V	-1.2425E-01
IMETA=	56.67	U	1.1214E-01	V	-2.5319E-01
K=	1.00	N	1.4952E-01	V	-3.6910E-01
		U	1.8649E-01	V	-4.4458E-01
		N	2.2427E-01	V	-5.9888E-01
S=	.73	RHO	3.2651E+00	(S-S1)/CV	7.9001E-01
RB=	.25	U	3.4461E+00	V	-0.
ZJ=	.25	N	4.1857E-02	V	-9.1762E-02
IMETA=	40.33	U	1.3714E-02	V	-1.8818E-01
K=	1.00	N	1.2557E-01	V	-2.6566E-01
		U	1.6743E-01	V	-3.4330E-01
		N	2.0928E-01	V	-4.1901E-01
		U	2.5114E-01	V	-5.0472E-01

S=	07	N	RMO	U	V	(S-S1)/CV	P	M	T	DM/M1
RB=	.77	0.	2.639E+00	2.5535E+00	-0.	7.9001E-01	8.5763E+00	1.1997E+00	3.2431E+00	.00362
Z0=	.30	4.9523E-02	2.8810E+00	2.5701E+00	-3.2730E-02	7.5044E-01	9.3211E+00	1.2079E+00	3.2344E+00	.00214
THETA=	0.00	9.9045E-02	3.12A4E+00	2.6042E+00	-8.0467E-02	7.0917E-01	1.0039E+01	1.2246E+00	3.2071E+00	.00550
K=	1.00	1.4857E-01	3.3605E+00	2.5540E+00	-1.1734E-01	6.6648E-01	1.0627E+01	1.2649E+00	3.1623E+00	.00611
		1.9809E-01	3.5936E+00	2.7346E+00	-1.5973E-01	6.2144E-01	1.1159E+01	1.3139E+00	3.1023E+00	.00543
		2.4761E-01	3.8335E+00	2.8253E+00	-1.9545E-01	5.7393E-01	1.1649E+01	1.3731E+00	3.0347E+00	.00367
		2.9714E-01	4.0510E+00	2.7451E+00	-2.5110E-01	5.2285E-01	1.1934E+01	1.4540E+00	2.9219E+00	.00400

SHOCK AND BODY PARAMETERS

S	07	K	THETA	RS	ZS	MS	CP	P/PU
0.	1.4493E-01	1.0000E+00	9.0000E+01	0.	-1.7774E-01	-0.8414E-07	1.4015E+00	1.0022E+00
1.4544E-01	1.4493E-01	1.0000E+00	1.1667E+01	1.7152E-01	-1.7094E-01	-7.4010E-07	1.7455E+00	9.7544E-01
2.9089E-01	2.8680E-01	1.0000E+00	7.3333E+01	3.4148E-01	-1.4063E-01	-8.0570E-07	1.6152E+00	9.0613E-01
4.3633E-01	4.2262E-01	1.0000E+00	8.5000E+01	5.0892E-01	-9.1387E-02	-9.2682E-07	1.4217E+00	8.3327E-01
5.8178E-01	5.4951E-01	1.0000E+00	5.6667E+01	6.7275E-01	-2.2866E-02	-1.1032E-06	1.1772E+00	6.7343E-01
7.2722E-01	6.6480E-01	1.0000E+00	4.8333E+01	8.3175E-01	8.5360E-02	-1.3157E-06	9.4184E-01	5.4818E-01
8.7266E-01	7.6604E-01	1.0000E+00	4.0000E+01	9.9366E-01	1.6622E-01	-1.7338E-06	6.7640E-01	4.0708E-01

TIME= 1.7201E+00

ITERATION NUMBER 318

SEGMENT NUMBER 2

TIME= 1.7201E+00

S=	1.02	N	PHO	U	V	(S-S1)/CV	P	M	T	DM/MI
RU=	.80	0.	2.797E+00	2.4530E+00	-0.	7.9001E-01	9.3032E+00	1.1372E+00	3.1222E+00	.00391
ZB=	.47	0.1624E-02	2.9632E+00	2.0417E+00	-1.5102E-01	7.1020E-01	9.3675E+00	1.2579E+00	3.1000E+00	.00932
IMETA=	40.00	1.6325E-01	3.1486E+00	2.8063E+00	-2.5720E-01	6.4170E-01	9.0575E+00	1.3708E+00	3.1231E+00	.01471
K=	0.	2.4487E-01	3.5295E+00	2.7672E+00	-3.2219E-01	5.6150E-01	1.0248E+01	1.4803E+00	2.4330E+00	.00506
		3.2644E-01	3.9261E+00	3.1415E+00	-3.7204E-01	4.7191E-01	1.0877E+01	1.6003E+00	2.7703E+00	.00000
S=	1.10	N	PHO	U	V	(S-S1)/CV	P	M	T	DM/MI
RU=	.95	0.	2.9067E+00	2.3921E+00	-0.	7.9001E-01	9.6149E+00	1.1002E+00	3.1704E+00	.00142
ZB=	.58	0.6726E-02	3.0379E+00	2.6443E+00	-1.3967E-01	7.0359E-01	9.5755E+00	1.2605E+00	3.1523E+00	.01102
IMETA=	40.00	1.7345E-01	3.1983E+00	2.8801E+00	-3.0994E-01	6.1609E-01	9.4287E+00	1.4288E+00	2.7843E+00	.01149
K=	0.	2.6018E-01	3.4671E+00	3.1006E+00	-4.1018E-01	5.2419E-01	9.6297E+00	1.5861E+00	2.7774E+00	.00594
		3.4691E-01	3.7998E+00	3.3011E+00	-5.0098E-01	4.2436E-01	9.4078E+00	1.7476E+00	2.6074E+00	.00300
S=	1.31	N	PHO	U	V	(S-S1)/CV	P	M	T	DM/MI
RU=	1.05	0.	2.9381E+00	2.3742E+00	-0.	7.9001E-01	9.9630E+00	1.0887E+00	3.3910E+00	.00123
ZB=	.69	9.0449E-02	3.1279E+00	2.6431E+00	-1.1889E-01	6.9219E-01	9.8619E+00	1.2543E+00	3.1523E+00	.01121
IMETA=	40.00	1.8090E-01	3.2858E+00	2.9300E+00	-3.1449E-01	5.9017E-01	9.5411E+00	1.4645E+00	2.9038E+00	.01207
K=	0.	2.7135E-01	3.4949E+00	3.1931E+00	-4.5254E-01	4.8878E-01	9.3990E+00	1.6620E+00	2.6833E+00	.00592
		3.5140E-01	3.7177E+00	3.3886E+00	-5.8558E-01	3.9534E-01	9.3329E+00	1.6342E+00	2.5106E+00	.00300
S=	1.43	N	PHO	U	V	(S-S1)/CV	P	M	T	DM/MI
RU=	1.14	0.	2.9611E+00	2.3504E+00	-0.	7.9001E-01	1.0072E+01	1.0798E+00	3.4010E+00	.00124
ZB=	.80	9.3527E-02	3.2023E+00	2.6531E+00	-1.3952E-01	6.8091E-01	1.0078E+01	1.2636E+00	3.1474E+00	.01143
IMETA=	40.00	1.8705E-01	3.3911E+00	2.9773E+00	-2.9883E-01	5.6661E-01	9.7398E+00	1.4922E+00	2.3722E+00	.01160
K=	0.	2.8054E-01	3.5821E+00	3.2531E+00	-4.6137E-01	4.6177E-01	9.3959E+00	1.7098E+00	2.6377E+00	.00477
		3.7411E-01	3.6842E+00	3.4230E+00	-8.1953E-01	3.8410E-01	9.1138E+00	1.6879E+00	2.4737E+00	.00300

SHOCK AND BODY PARAMETERS

S	4d	Zd	N	IMETA	RS	ZS	MS	CP	P/PO
1.0101E+00	4.2937E-01	4.6103E-01	0.	4.0000E+01	1.1096E+00	2.5870E-01	1.4347E-00	7.4136E-01	4.4154E-01
1.1036E+00	4.5332E-01	5.8005E-01	0.	4.0000E+01	1.2188E+00	3.5700E-01	-1.3495E-00	7.4701E-01	4.0285E-01
1.3090E+00	1.0405E+00	6.9146E-01	0.	4.0000E+01	1.3237E+00	4.5890E-01	2.2344E-00	8.0027E-01	4.7290E-01
1.4544E+00	1.1400E+00	8.0246E-01	0.	4.0000E+01	1.4266E+00	5.6241E-01	-8.4036E-00	8.1004E-01	4.7807E-01

SEGMENT NUMBER 3		ITERATION NUMBER 328		TIME= 1.7741E+00																
S=	1.60	N	0.	(S-S1)/CV	7.9001L-01	P	1.0251E+01	M	1.0624E+00	I	3.4157E+00	Dt/M1	0.0120							
R3=	1.23	RMO	2.4986E+00	U	2.3309E+00	V	-0.						0.0120							
ZB=	.91	N	9.6446E-02	U	2.0636E+00	V	-1.0316E-01						0.0120							
TMETA=	40.00	N	1.9219E-01	U	3.0141E+00	V	-2.8092E-01						0.0120							
K=	0.	N	2.8914E-01	U	3.2902E+00	V	-4.5856E-01						0.0120							
		N	3.6578E-01	U	3.4333E+00	V	-6.3027E-01						0.0120							
S=	1.70	N	0.	RMO	3.0376E+00	U	2.3054E+00	V	-0.	(S-S1)/CV	7.9001L-01	P	1.04439E+01	M	1.0624E+00	I	3.4157E+00	Dt/M1	0.0120	
R3=	1.33	N	9.9439E-02	RMO	3.3233E+00	U	2.5730E+00	V	-9.7818E-02											0.0120
ZB=	1.03	N	1.9838E-01	N	3.6094E+00	U	3.0491E+00	V	-2.6519E-01											0.0120
TMETA=	40.00	N	2.9832E-01	N	3.7302E+00	U	3.3246E+00	V	-4.4739E-01											0.0120
K=	0.	N	3.9776E-01	N	3.6819E+00	U	3.4229E+00	V	-6.2197E-01											0.0120
S=	1.84	N	0.	RMO	3.0763E+00	U	2.2702E+00	V	-0.	(S-S1)/CV	7.9001L-01	P	1.0625E+01	M	1.0300E+00	I	3.4573E+00	Dt/M1	0.0120	
R3=	1.42	N	1.0259E-01	RMO	3.6122E+00	U	2.0957E+00	V	-9.3202E-02											0.0120
ZB=	1.14	N	2.0519E-01	N	3.7103E+00	U	3.0009E+00	V	-2.5263E-01											0.0120
TMETA=	40.00	N	3.0778E-01	N	3.6217E+00	U	3.3446E+00	V	-4.231E-01											0.0120
K=	0.	N	4.1018E-01	N	3.7007E+00	U	3.4049E+00	V	-6.0264E-01											0.0120
S=	2.04	N	0.	RMO	3.1069E+00	U	2.2571E+00	V	-0.	(S-S1)/CV	7.9001L-01	P	1.0773E+01	M	1.0244E+00	I	3.4670E+00	Dt/M1	0.0120	
R3=	1.51	N	1.0615E-01	RMO	3.4357E+00	U	2.7133E+00	V	-8.7908E-02											0.0120
ZB=	1.25	N	2.1231E-01	N	3.7923E+00	U	3.1172E+00	V	-2.4520E-01											0.0120
TMETA=	40.00	N	3.1840E-01	N	3.6821E+00	U	3.3540E+00	V	-4.1675E-01											0.0120
K=	0.	N	4.2402E-01	N	3.7213E+00	U	3.3840E+00	V	-5.1140E-01											0.0120

SHOCK AND BODY PARAMETERS

S	Rd	K	TMETA	MS	Z5	MS	CP	P/PJ
1.5494E+00	1.2335E+00	9.1430E-01	0.	1.5290E+00	0.0072E-01	-1.5032E-00	1.20J1E-J1	4.1027E-J1
1.7453E+00	1.3270E+00	1.0257E+00	0.	1.6317E+00	7.7004E-01	-1.3942E-00	1.4271E-J1	4.1047E-J1
1.8300E+00	1.4200E+00	1.1371E+00	0.	1.7348E+00	8.7334E-01	-3.2755E-00	1.2747E-J1	4.1403E-01
2.0362E+00	1.5140E+00	1.2745E+00	0.	1.8392E+00	9.7361E-01	5.0525E-00	3.7201E-J1	4.1135E-J1

S= 2.23
 RD= 1.64
 ZB= 1.40
 THETA= 40.00
 K= 0.

SEGMENT NUMBER	ITERATION NUMBER	174	TIME= 3.4110E-01	DM/MI				
N	RHO	U	V	(S-S1)/CV	P	M	F	CP
0.	3.1147E+00	2.2524E+00	-0.	7.9001E-01	1.0011E+01	1.0210E+00	3.4711E+00	.00100
6.9240E-02	3.4197E+00	2.6540E+00	-6.9417E-02	6.5152E-01	1.0720E+01	1.2672E+00	3.1372E+00	.01310
1.7850E-01	3.7624E+00	3.0230E+00	-1.0323E-01	5.1740E-01	1.0724E+01	1.5101E+00	2.3504E+00	.00930
2.6774E-01	3.9433E+00	3.2627E+00	-2.9964E-01	4.2837E-01	1.0477E+01	1.6900E+00	2.6570E+00	.00220
3.5099E-01	3.9261E+00	3.3640E+00	-4.3240E-01	3.9310E-01	1.0053E+01	1.7914E+00	2.5600E+00	.00090
4.4624E-01	3.7522E+00	3.3530E+00	-5.4975E-01	4.0744E-01	9.5710E+00	1.7900E+00	2.5507E+00	.00001

SHOCK AND BODY PARAMETERS

S	RD	ZB	THETA	RS	ZS	MS	CP	P/PO
2.2339E+00	1.6410E+00	1.4000E+00	0.	1.9829E+00	1.1132E+00	3.6455E-00	0.7600E-01	5.1316E-01

SECTION VIII

CONCLUSIONS

1. THE DIFFERENCING SCHEME

The differencing scheme developed in this study incorporates a complete decoupling of the magnitude of the stabilizing terms from the finite difference mesh sizes used. Consequently, different mesh sizes for each coordinate direction and variations in mesh size can be employed. The mesh sizes can be selected based on finite difference accuracy requirements without regard to the stabilizing term influence. Since the stabilizing terms can be specified arbitrarily small, the ability to achieve a valid mathematical model of the inviscid flow problem is ensured. The stability arguments used to establish the differencing scheme should be applicable to any coordinate system. Thus, the differencing scheme is not limited to the blunt body problem. It should be useful for any problem for which the time-dependent technique is applicable.

2. SURFACE BOUNDARY CONDITIONS

The technique developed for treating points on the body surface is exact in the limit of standard finite difference approximations. This technique is applicable to any point for which a streamline and a coordinate line are coincident, in particular, the body points with a body oriented coordinate system. This technique has significantly improved the accuracy of the blunt body solution. Since the general class of body shapes required to realize the objectives of this study can be achieved with body oriented coordinates, this exact method has been employed. However, definite restrictions on the magnitude of a negative body curvature result from this technique. The restriction to smooth body contours is fairly academic, since any finite difference solution for a body with corners will necessarily assume a finite curvature at the discontinuities (e.g., see Ref. 11).

3. COMPUTER EFFICIENCY

The segmented computation procedure developed greatly enhances the computer efficiency obtainable with a time-dependent technique. The computer storage required is minimal regardless of the body length. Computation times are reduced by factors in excess of 1000 when typical high performance vehicles are considered. It has been shown that even the relatively simple axisymmetric flow problem requires the type of computation time reduction realized by this procedure if a unified

time-dependent solution is to be economically feasible. An equivalent three-dimensional solution without this type of improved computer efficiency would certainly be impossible with the computers currently available. This procedure should enable the aerodynamicist to employ the time-dependent technique as a practical solution procedure for complex problems such as nonequilibrium and three-dimensional flows. Certainly, this procedure or another equally effective approach is required to transform the time-dependent method from an interesting theoretical solution to a useful computational method.

4. THE BLUNT BODY SOLUTION

A practical computational method has been presented to treat axisymmetric blunt body flows. The computer code presented can consider a very general class of body geometries, sufficient to satisfy the objectives of this study. Since additional stabilizing terms are introduced into the governing equations, reference to this technique as an exact solution may be properly questioned. However, the differencing scheme assures that the effect of these terms can be made negligible relative to the numerical errors. Consequently, this question is rather academic. For practical purposes, the present technique is a direct, exact and unified solution procedure. This is believed to be the first technique reported which contains all of these features and still may be employed as a practical computational method. Extensive comparison with other computational methods has been accomplished. The present method is clearly more accurate than other time-dependent techniques reported. In fact, the accuracy is comparable to other methods currently regarded as standard solution procedures. A calculation for a typical high performance reentry vehicle has been presented. The computation time and computer storage required was quite reasonable. This is believed to be the first reported solution for a problem of this magnitude generated with a unified time-dependent method.

APPENDIX I
STABILITY CONSIDERATIONS

1. RESULTS OF REFERENCE 1

In reference 1 it was shown that useful results could be obtained from a stability analysis of a simple linearized one-dimensional momentum equation of the form

$$u_t + Ku_x = \mu u_{xx} \quad (71)$$

Employing a forward difference approximation for the partial derivative with respect to time and appropriate first and second central difference approximations for the spatial partial derivatives, the von Neumann (Ref. 5) stability analysis was used. In this analysis, u is replaced by

$$u \rightarrow u_0 + (\delta u) \text{EXP}(\alpha t - i\beta \Delta x)$$

where the second term introduces an error into u .

The differencing scheme will be stable if the errors decay exponentially in time, i.e.,

$$|e^{\alpha t}| \leq 1 \quad (72)$$

The details of the analysis are given in reference 1. Only the results of that analysis will be given here

$$\frac{K^2(\Delta t)^2}{(\Delta x)^2} \sin^2 \theta + \left[1 - \frac{2\mu \Delta t}{(\Delta x)^2} (1 - \cos \theta) \right]^2 \leq 1 \quad (73)$$

where

$$\theta = \beta \Delta x \quad (74)$$

It is easily seen that if μ vanishes, the central difference scheme must be unstable for a finite, positive Δt , i.e.,

$$\frac{K^2(\Delta t)^2}{(\Delta x)^2} \sin^2 \theta \leq 0 \quad (75)$$

Using a special relation for μ

$$\mu = \nu \frac{(\Delta x)^2}{2\Delta t} \quad (76)$$

It was found (Ref. 1) that

$$\Delta t \leq \frac{\sqrt{\nu \Delta x}}{K} \quad (77)$$

should result in a stable differencing scheme where ν is a constant and

$$0 < \nu \leq 1 \quad (78)$$

K was replaced by $(u+a)$ to agree with the Courant-Friedricks-Lewy (CFL) stability criterion, which must always be satisfied. Experience with this expression showed that it was extremely accurate when applied to the multidimensional flow problem. The impressive success of this analysis in reference 1 prompted the attempt to generalize equation (77) discussed in section IV. The computational results presented in section V clearly demonstrate the validity of using this analysis.

2. A STABILITY ANALYSIS FOR BACKWARD DIFFERENCES.

A similar stability analysis will now be presented where u_x is approximated by a backward difference

$$u_x \rightarrow [u(x,t) - u(x-\Delta x,t)] / \Delta x \quad (79)$$

Introducing this difference approximation into equation (71) it is easily shown that

$$\left| 1 - \frac{2\mu\Delta t}{(\Delta x)^2} (1 - \cos \theta) - \frac{K\Delta t}{\Delta x} (1 - \cos \theta + i \sin \theta) \right| \leq 1 \quad (80)$$

is the appropriate stability criterion.

By taking the absolute value of this complex expression, it is easily reduced to

$$\Delta t \leq \frac{\Delta x}{\frac{2\mu}{\Delta x} + K} \quad (81)$$

In this case, if μ vanishes, a stable differencing scheme is still possible when K is positive. This accounts for the success of linearly extending the solution in the supersonic region to terminate the solution domain (Refs. 1, 11-15). It is easily shown that this procedure results in a backward difference in the streamwise coordinate with no stabilizing terms relative to that direction, i.e., the second derivatives vanish. The question of stability is more complicated when K vanishes. However, to achieve agreement with the CFL stability criterion, equation (81) must be written

$$\Delta t \leq \frac{\Delta x}{\frac{2\mu}{\Delta x} + u + a} \quad (82)$$

If μ and u vanish, we expect

$$\Delta t \leq \frac{\Delta x}{a} \quad (83)$$

will result in a stable differencing scheme. This will permit a backward difference approximation, together with the requirement that the normal velocity component vanish on the body, to be used to impose surface boundary conditions. Similar arguments using a forward difference approximation for u_x result in

$$\Delta t \leq \frac{\Delta x}{\frac{2\mu}{\Delta x} - K} \quad (84)$$

which allows μ to vanish when K is negative. As a result, the artificial dissipative terms could vanish in the streamwise direction at the stagnation streamline for axisymmetric flow.

APPENDIX II

THE UNSTEADY SHOCK PROBLEM

1. THE METHOD OF GODUNOV

Time-dependent blunt body flows with discontinuous shock waves were considered by Godunov (Ref. 9) using the one-dimensional Riemann problem to describe the wave interactions near the shock surface. Reference 16 contains an excellent description of this approach. The relevant wave phenomena are illustrated in figure 21. W_s is the shock velocity and V_s is the flow velocity behind a normal shock wave. The left running disturbance is the bow shock. The right running disturbance defined by $\frac{dx}{dt} = V+a$ may be either a shock or a finite expansion wave. The wave defined by $\frac{dx}{dt} = V_s$ is an entropy line. The shock jump conditions are expressed by

$$\dot{m} = \left(\frac{\gamma+1}{2} P_s + \frac{\gamma-1}{2} \right)^{1/2} \quad (85)$$

$$\dot{m} V_s + P_s = \dot{m} V_\infty + P_\infty \quad (86)$$

where \dot{m} is the mass flow across the wave, the subscripts s and ∞ refer to conditions behind the shock and in the free stream, respectively, and nondimensional variables with a perfect gas equation of state have been used. The Godunov method relies

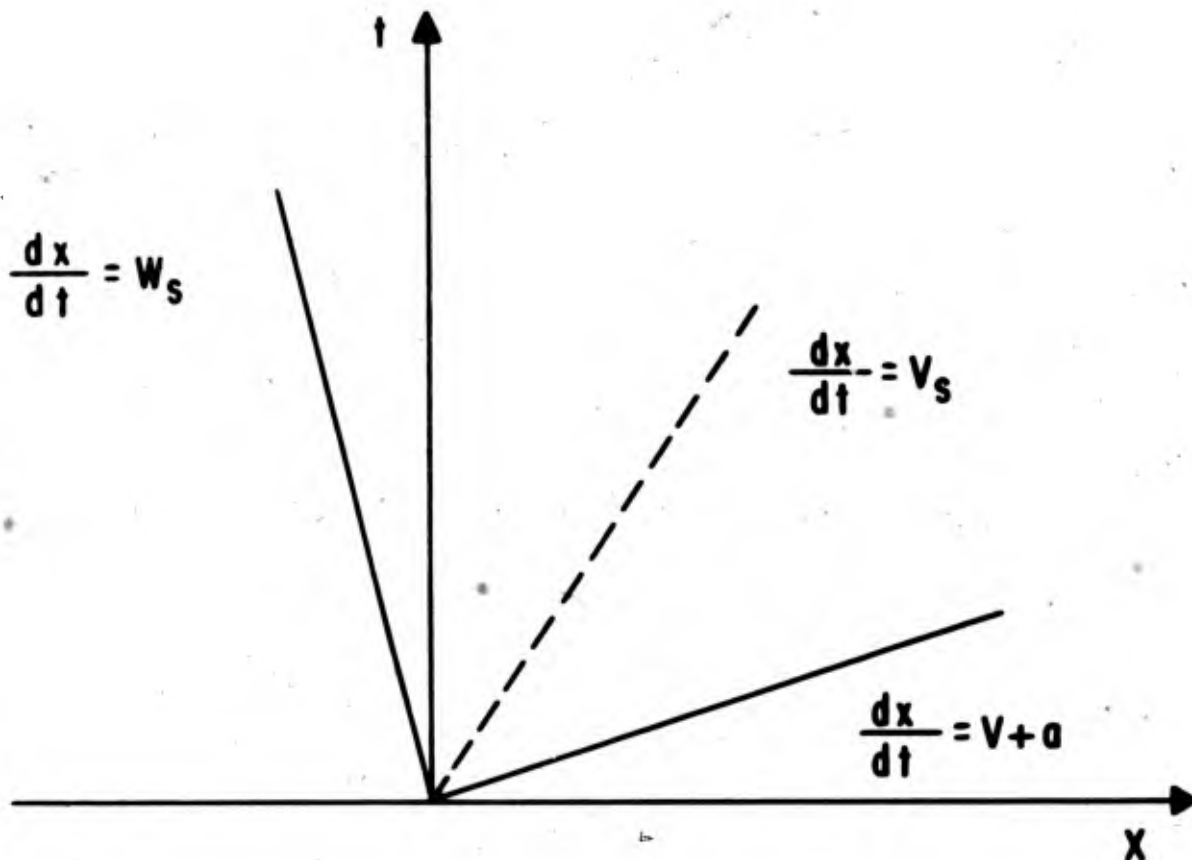


Figure 21. The One-Dimensional Riemann Problem

on two independent solutions for P_s based on independent jump relations for right and left running waves. An iteration procedure is used to obtain acceptable agreement between the two predictions for P_s . The validity of the method appears questionable for the general blunt body problem. Specifically, when the stagnation streamline for a spherical body is considered, this method does not model the problem correctly. This region is characterized by increasing pressure and decreasing velocity as the body is approached from the shock. In the limit of a steady state flow, the flow inside the shock is isentropic. Then, \dot{m} across the right running disturbance is given by

$$\dot{m} = -\delta(\gamma P \rho)^{1/2} \frac{1 - \zeta}{1 - \zeta} \delta \quad (87)$$

where

$$\delta = \frac{\gamma - 1}{2\gamma} \quad (88)$$

ζ is the ratio of the pressures behind and in front of the wave and P and ρ are the pressure and density in front of the wave. Since ζ is less than unity, \dot{m} must be negative. Then, equation (86) requires that the velocity in front of the right running disturbance be greater than V_s . This is not in agreement with the physical situation where V decreases with increasing X . To avoid this difficulty, some attempt to consider multidimensional effects is required.

2. THE METHOD OF MORETTI AND ABBETT

A method that includes multidimensional effects in an approximate manner was suggested by Moretti and Abbett (Ref. 13). The governing equations in the shock fixed coordinate system shown in figure 3 are

$$\frac{\partial \rho}{\partial t} + \rho \frac{\partial V}{\partial x} + V \frac{\partial \rho}{\partial x} = - \frac{\partial \rho u}{\partial y} \quad (89)$$

$$\frac{\partial V}{\partial t} + V \frac{\partial V}{\partial x} + \frac{1}{\rho} \frac{\partial P}{\partial x} = - u \frac{\partial V}{\partial y} \quad (90)$$

and S is presumed constant inside the shock. The right hand terms in equations (89) and (90) are considered to be constant

$$A = \frac{\partial \rho u}{\partial y} \quad (91)$$

$$B = U \frac{\partial V}{\partial y} \quad (92)$$

For a rectilinear shock, A and B would vanish. It is reasonable to expect constant values for A and B to adequately model the curved shock problem. Multiplying both sides of equation (89) by $\frac{a}{\rho}$

$$\frac{a}{\rho} \frac{\partial \rho}{\partial t} + v \frac{a}{\rho} \frac{\partial \rho}{\partial x} + a \frac{\partial v}{\partial x} = -A \frac{a}{\rho} \quad (93)$$

and rearranging equation (90)

$$\frac{\partial v}{\partial t} + v \frac{\partial v}{\partial x} + \frac{a^2}{\rho} \frac{\partial \rho}{\partial x} = -B \quad (94)$$

where

$$\frac{\partial P}{\partial x} = \left(\frac{\partial P}{\partial \rho} \right)_S \frac{\partial \rho}{\partial x} = a^2 \frac{\partial \rho}{\partial x} \quad (95)$$

for isentropic flow. Adding and subtracting equations (94) and (95)

$$\begin{aligned} \frac{\partial v}{\partial t} + (v \pm a) \frac{\partial v}{\partial x} \pm \frac{a}{\rho} \left[\frac{\partial \rho}{\partial t} \right. \\ \left. + (v \pm a) \frac{\partial \rho}{\partial x} \right] = - (B \pm A \frac{a}{\rho}) \end{aligned} \quad (96)$$

The \pm notation should be used consistently through the remaining equations.

Defining

$$\frac{d}{dt} = \frac{\partial}{\partial t} + (v \pm a) \frac{\partial}{\partial x} \quad (97)$$

Equation (96) can be expressed as

$$\frac{dv}{dt} \pm \frac{a}{\rho} \frac{d\rho}{dt} = - \left[B \pm A \frac{a}{\rho} \right] \quad (98)$$

This indicates that the directional derivative defined by equation (97) is the significant operator. This operator corresponds to a derivative along the characteristic directions

$$\frac{dx}{dt} = v \pm a \quad (99)$$

Noting that

$$\frac{d\rho}{dt} = \left(\frac{\partial \rho}{\partial P} \right)_S \frac{dP}{dt} = \frac{1}{a^2} \frac{dP}{dt} \quad (100)$$

for isentropic flow

$$\frac{dP}{dt} = \rho a \left[\frac{dv}{dt} \pm (B - A \frac{a}{\rho}) \right] \quad (101)$$

is obtained. Equation (101) represents the compatibility relation associated with the characteristic directions specified by equation (99). The pressure predictions from the Rankine-Hugoniot shock jump relations and from equation (101) can be used in an iterative scheme to obtain the conditions behind a normal shock. It is easily seen that a clear parallel exists between this approach and the method of Godunov, the difference being that multi-dimensional effects have been considered in an approximate manner.

3. APPLICATION ON THE SYMMETRY AXIS

When the method of reference 13 is applied on the symmetry axis, equations (91) and (92) must be modified. The unsteady shock problem in axisymmetric flow becomes a three-dimensional problem at this point. It is easily shown that if

$$A = \frac{\partial \rho u}{\partial x} + \frac{\partial \rho W}{\partial z} \quad (102)$$

$$B = U \frac{\partial V}{\partial y} + W \frac{\partial V}{\partial z} \quad (103)$$

where z is perpendicular to x and y and W is the velocity component along z , all equations developed above are valid at the symmetry axis. For axisymmetric flow, the second term in equation (102) is equal to the first term and the second term in equation (103) vanishes. Consequently, the only modification required is to multiply equation (91) by 2 when computations are performed on the symmetry axis.

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