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MEMORANDUM REPORT NO. 1983

JUNE 1969

AN EXPERIMENTAL INVESTIGATION AT SUPERSONIC MACH NUMBERS OF BASE DRAG OF VARIOUS BOATTAIL SHAPES WITH SIMULATED BASE ROCKET EXHAUST

Joseph Huerta

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JHuerta/lca Aberdeen Proving Ground, Md. June 1969

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ABSTRACT

An experimental investigation was conducted to determine the effect of boattail geometry and simulated base rocket nozzle flow on overall base drag. All tests were performed at Mach numbers 2.50, 3.00 and 3.50. Reasonable prediction of base drag for the reported boattail shapes can be made by using approximate and empirical equations. A conical boattail produced less drag than the other boattail shapes investigated. The maximum nozzle stagnation to free-stream pressure ratio reported is 315.

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	LIST OF SYMBOLS
a [*] j	sonic velocity at jet exit conditions
^m i	mass flow based on $A_{\underline{B}}$ and free stream conditions
^m j	mass flow through sustainer nozzle
р _b	average base pressure
p _c	sustainer nozzle stagnation pressure
Pj	jet exit static pressure
Pi	local static pressure
P _∞	free-stream static pressure
p 2	static pressure downstream of a corner
đ	free-stream dynamic pressure
r	radius of body at point x
S .	rectangular coordinate in streamline direction
x	coordinate measured from vertex of body
x ₂	position of a corner
A _B	maximum body area
А _р	base area
A _j	jet exit area
AL	increment of area associated with the local static pressure
$c_{\rm DB}$	base drag coefficient
C _{DBT}	boattail drag coefficient
c,	characteristic coordinate along Mach line
Ъ	base diameter
D _j	jet exit diameter
М	Mach number

LIST OF SYMBOLS (Continued)

м_ј jet exit Mach number M*j dimensionless jet velocity, V_j/a_j^* M_ free-stream Mach number Ро tunnel stagnation pressure free-stream Reynolds number Re V_j jet exit velocity v. local velocity ratio of specific heats Y δ. flow deflection angle flow deflection angle before a corner δ1 flow deflection angle after a corner δ₂

Mach angle (arcsin 1/M)

μ

I. INTRODUCTION

Methods for reducing missile drag, especially base drag, have been of considerable practical importance. As a consequence, the Army Missile Command (MICOM), Redstone Arsenal, Alabama, has been conducting a continual program to investigate techniques for minimizing missile base drag. As a part of this continuing investigation, the present tests were conducted in the supersonic wind tunnel of the Ballistic Research Laboratories (BRL) to determine the effect of six boattail configurations on drag. Base drag was measured in the presence of a sustainer rocket motor exhausting cold flow through the base. A comparison of experimental data and a second-order shock-expansion method for boattail-pressure distribution was made. An empirical technique was employed to estimate "power on" base drag and these results are compared with measured data.

The basic model was a strut-mounted body of revolution. The section forward of the boattail consisted of a tangent-ogive nose followed by a cylindrical center body. Pressure data were obtained along the boattail surface and model base at nominal test Mach numbers of 2.50, 3.00 and 3.50 for each boattail shape with the model at zero angle of attack. Sustainer nozzle pressure was also varied for each boattail. The ratio of nozzle exit diameter to body diameter was 0.200 and the nozzle exit Mach number was 2.7 for fully developed flow or "power on" condition.

II. APPARATUS

A. Wind Tunnel

Supersonic Wind Tunnel No. 1 of the Ballistic Research Laboratories was used to acquire the data. This is a two-dimensional continuously operated, variable density, closed circuit tunnel having a test section 15 inches high and 13 inches wide. Two flexible steel plates form the upper and lower walls of the nozzle section. A complete description of this facility and its flow characteristics are given in References 1 and 2.

B. Test Model

The model was a body of revolution attached to the upper test section wall by a swept back support strut at zero angle of attack. It consisted of a tangent-ogive nose attached to a cylindrical center body and ending with a boattail. Figure 1 shows a photograph of the model installation in the test section and the schematic of Figure 2 presents the pertinent dimensions of the model and support strut. The overall length of the model was 15.0 inches and the diameter was 2.500 inches. The support strut was a modified double wedge section, 0.435 inch thick with a 3.125 inches chord. The exit diameter of the sustainer nozzle was 0.500 inch or 0.200 calibers and the exit Mach number was 2.7. Model pressure taps were placed opposite the model support strut and located on the cylindrical surface, boattail surface and at the base. The photograph of Figure 3 presents a view of these pressure taps.

Six boattail configurations were investigated in this program. All of these boattails were one caliber in length. Figures 4 and 5 define the geometric shapes of these boattails which consisted of conical, concave, convex and reverse curvature configurations.

C. Instrumentation

The model pressures were transmitted through metal tubing to a pressure scanner system. The pressure scanner (Figure 6) is a solid metal block which is channeled such that, with the aid of pneumatically operated valves, seals and stepping switches, pressures can be measured in sets of seven pressures per cycle. This scanner unit is necessary since the tunnel automatic data acquisition system is capable of handling only eight inputs per sampling time. Pressures from the scanner were measured by a group of pressure transducers of suitable pressure range. The electrical output signals from these transducers were converted by an automatic readout system to digital readings which were recorded automatically by typewriter. A complete description of the scanner is given in Reference 3.

The sustainer nozzle total pressure and tunnel total pressure were monitored continuously throughout the period of data acquisition. The sustainer nozzle stagnation pressure was measured in its plenum chamber.

A schlieren system with camera provides continuous visual indication, as well as photographs, of the flow conditions in the test section.

III. TEST PROCEDURE AND CONDITIONS

The flow in the wind tunnel was established at a low stagnation pressure and then the pressure was raised to the desired level. The sustainer nozzle pressure was adjusted to the proper value. With all flow conditions at steady state, model pressure readings were obtained with the pressure scanner system. Before each set of model pressures were recorded, it was ascertained that these pressures were stabilized in the scanner system by observing their respective dial readings on the control panel. Pressure readings were recorded when all dial motion ceased. After making a complete pressure sampling cycle, the jet pressure was adjusted to the next value, and the model pressure reading cycle was repeated. The available line pressure for the sustainer nozzle was 300 psia. The variation in tunnel stagnation and sustainer nozzle pressure was less than 0.1 percent.

The average tunnel operating conditions are listed in Table I for the three Mach numbers. The specific humidity was maintained below 0.0002 pound of water per pound of air. Schlieren photographs were taken of the flow field for selected model configurations.

Table I. Tunnel Operation Conditions

Mach No.	P _o <u>rsia</u>	q psi	Re x 10 ⁻⁶ per inch
2.50	28.88	7.40	0.50
3.00	40.45	6.94	0.50
3.50	52.69	5.92	0.50

A narrow band of fine grit was placed on the nose surface near the tip to insure a turbulent boundary layer at all times.

IV. DATA REDUCTION

The raw pressure data were converted into the required ratio and coefficient form by a computer. The data acquisition system automatically recorded the raw data from the pressure transducers in digital typewritten and coded punched tape form. The punched tape was used to obtain punched cards which were fed into the computer for reduction and tabulation of the raw data into the desired form.

The local pressures were weighted by area to determine the boattail and base drag coefficients. The boattail and base drag coefficients were calculated from the following general relationship

 $C_{D} = \frac{\frac{1}{2} \left[\left(p_{l_{i}} - p_{m} \right) A_{i} \right]}{q A_{B}} \qquad \begin{array}{l} n = number of \\ local pressures \\ and associated \\ incremental areas \end{array}$

where p, is the local pressure. The area, A, used to calculate boattail drag is the local area, A,, projected on the vertical plane. The overall error in determining the model static pressures was approximately ± .003 psi.

V. DISCUSSION OF THEORETICAL AND EMPIRICAL CALCULATIONS

The equation for the streamwise pressure gradient is given as

 $\frac{\partial p}{\partial s} - \frac{2\gamma p}{\sin 2\mu} \frac{\partial \delta}{\partial s} = \frac{1}{\cos \mu} \frac{\partial p}{\partial C_{1}}$

In obtaining the Prandtl-Meyer relation, the pressure is considered constant along the first-family Mach lines. Therefore, the right hand part of the above equation becomes zero and the solution of the integrated equation is the Prandtl-Meyer relation. The second-order shock expansion method imposes a different condition on the pressure gradient equation to obtain an approximate solution (Reference 5).

First consider a body of revolution with a pointed nose and supersonic flow on its surface. Drawing a series of tangent lines to this body contour results in an approximate shape consisting of a series of frustums with a cone vertex. Now on this new geometric shape, it can be assumed that the flow direction does not change within a single element. Now it can be stated that $\partial\delta/\partial s = 0$ and the pressure gradient equation reduces to

$$\frac{\partial p}{\partial s} = \frac{1}{\cos \mu} \frac{\partial p}{\partial C_1}$$

The solution to the above differential equation is the second-order shock expansion method which is

$$p = p_{f} - (p_{f} - p_{2}) e^{-T_{f}}$$

where

pf = pressure on a cone tangent to the body at the same point as a frustum element

$$\eta = \left(\frac{\partial p}{\partial s}\right)_{2} \frac{\mathbf{x} - \mathbf{x}_{2}}{(\mathbf{p}_{f} - \mathbf{p}_{2})\cos\delta_{2}}$$
$$\left(\frac{\partial p}{\partial s}\right)_{2} = \frac{B_{2}}{r} \left(\frac{\Omega_{1}}{\Omega_{2}}\sin\delta_{1} - \sin\delta_{2}\right) + \frac{B_{2}}{B_{1}}\frac{\Omega_{1}}{\Omega_{2}} \left(\frac{\partial p}{\partial s}\right)_{1}$$

$$B = \frac{\gamma p M^2}{2 (M^2 - 1)} = \frac{q}{M^2 - 1}$$

and

$$\Omega = \frac{1}{M} \left[\frac{1 + (\frac{\gamma - 1}{2}) M^2}{\frac{\gamma + 1}{2}} \right]^{\frac{\gamma + 1}{2(\gamma - 1)}}$$

The above equation for Ω is the one dimensional area ratio. Now the second-order equation gives the pressure variation along the surface of the body.

Only the boattail surface pressures are of interest here, so the foregoing second-order equation is further simplified by imposing the following conditions:

- 1. $(\partial p/\partial s)_1 = 0$
- 2. $p_1 = p_{\infty}$ or $p_1 = p_{\infty}$
- 3. $\mathbf{x} = 0$ at the start of the boattail which eliminates the dimension \mathbf{x}_{2}

The following sketch shows that $\delta_1 = 0^\circ$, and δ_2 is a negative angle.



The boattail pressure distributions were calculated from the following set of simplified equations.

$$p = p_{\infty} - (p_{\infty} - p_2) e^{-\Pi}$$

$$\Pi = \left(\frac{\partial p}{\partial s}\right)_2 \frac{p_{\infty} - p_2 \cos \delta_2}{(p_{\infty} - p_2) \cos \delta_2}$$

$$\left(\frac{\partial p}{\partial s}\right)_2 = -\frac{B_2}{r} \sin \delta_2$$

$$B_2 = \frac{\gamma p_2 M_2^2}{2 (M_2^2 - 1)}$$

The presence of nozzle flow at the base adds to the complexity of the flow mechanism at the base which is difficult to analyze. Consequently, empirical results have been sought by correlation of available experimental data (Reference 6). A reasonable base pressure prediction

for a base with nozzle flow can be made by use of the following empirical equation:

$$\frac{p_{b}}{p_{m}} = \frac{1}{M_{j}^{*}} \left[\frac{3.5}{1 + 2.5 (A_{b}/A_{B})} \right] \left[0.19 + 1.28 \left(\frac{R_{mf}}{1 + R_{mf}} \right) \right]$$

where

P -	(m v) _j	_	Y _j P _j A _j M _j ²
nnf =	(m V) _{co}	-	$\gamma_{\infty} p_{\infty} A_{B} M_{\infty}^{2}$

The base drag coefficient is then calculated by use of the equation:

 $C_{D_{B}} = \frac{2}{\gamma_{\infty} M_{\infty}^{2}} \left(1 - \frac{p_{b}}{p_{\infty}}\right) \left(\frac{A_{b} - A_{j}}{A_{B}}\right)$

VI. PRESENTATION AND DISCUSSION OF DATA

Figures 7 through 9 present data of the boattail portion of all test configurations. The curves in Figures 10, 11 and 12 indicate results of base measurements for these configurations. A comparison of total drag (boattail and base drag) for configurations 6, 7 and 8 is presented in Figure 13. The test data presented are for test Mach numbers of 2.50, 3.00 and 3.50.

The boattail surface pressure distributions were determined for the six boattail configurations and typical distributions are presented in Figure 7 as a function of distance in calibers from the boattail corner. For boattail configurations 6, 7 and 8, there are included, for comparison, pressure distribution curves obtained by the method of characteristics at Mach numbers 2.50 and 3.00 and by the second-order shock expansion method of Reference 5 for all three Mach numbers. Configuration 11 data are presented for Mach numbers 2.50 and 3.00 with the method of characteristic curves included for comparison. Either method appears to predict the pressure distribution of these boattails well for at least the last two-thirds of the boattail surface. The wide differences between experimental points and the curves at the beginning of the boattail must be due to the boundary layer present at the corner.

The data of Figure 8 provides evidence that the flow from the sustainer nozzle has negligible effect on the boattail drag. The boattail drag coefficient of each boattail configuration appears unchanged for all base flow conditions.

The variation in boattail drag with Mach number is shown in Figure 9. It appears that the conical boattail configuration 6 offers the least drag when compared to other boattail shapes having the same base diameter.

The data of Figure 10 indicate the influence of the base nozzle flow on base pressure for the case of fully developed flow in the nozzle. Results for configurations 6, 7 and 8 are presented. The results of configuration 6 apply to configurations 9, 10 and 11 as well. The method of Reference 6 for predicting base pressure was utilized in calculating the curves shown with these data for comparison. Reasonable predictions of base pressure can be made with this empirical method. However, it appears that the ratio of nozzle diameter to base diameter has a strong influence on accuracy of the prediction. The results in Figure 11 are the base pressure data of Figure 10 reduced to base drag coefficient form and the comments made on Figure 10 data apply to these also. The base drag as a function of the sustainer jet flow parameter is displayed in Figure 12. These curves show that the base drag for configurations 6, 9, 10 and 11 is the same as was pointed out above.

A comparison of total drag (boattail and base drag) for the three conical boattails and a square base configuration from Reference 7 is shown in Figure 13. Only slight differences in total drag are indicated for the conical boattails. However, the square base configuration has about 60 percent more drag than the conical boattail configurations.

The schlieren photographs of Figure 14 show the flow field at the base for configuration 6 at Mach numbers 2.50 and 3.50 for base jet flow off and on.

VII. CONCLUSIONS

From this investigation, it can be concluded that less total drag results from a conical boattail configuration than any of the other contours tested.

Boattail pressure distribution can be determined over at least the last two-thirds of a conical boattail surface using the approximate shock-expansion method of Reference 5 which is much easier to use than the method of characteristics.

A reasonable prediction of base drag with rocket flow at the base can be made using the empirical method of Reference 6.



L SURFACE	DISTANCE FROM BASE	3366	2.964	2.564	2434	2.138	1.842	1.546	1250	0954	Q658	Q362	0066	
MODE	TAP	-	2	n	4	2	9	2	ø	6	⁰	11	12	



SUNIO

5

•

TANGEN



Figure 2. Model Dimensions and Boattail Tap Locations



MODEL DIMENSIONS



2.500



FFFF

3.125













BOATTAIL NO.6



BOATTAIL NO.7





BOATTAIL NO. 8





BOATTAIL NO.10



BOAT TAIL NO.11



Figure 5. Boattail Dimensions and Base Tap Locations





























6/Pa



(e) Boattail No. 11 Figure 7. Concluded

٥	CONFIGURATION	NO. 6
×	CONFIGURATION	NO. 7
+	CONFIGURATION	NO. 8
•	CONFIGURATION	NO. 9
Δ	CONFIGURATION	NO. 10
۵	CONFIGURATION	NC. II





Figure 8. Effect of Sustainer Nozzle Operating Pressure on Boattail Drag













(b) Configuration No. 7

Figure 10. Continued







(a) Configuration No. 6







(b) Configuration No. 7

Figure 11. Continued



(c) Configuration No. 8

Figure 11. Concluded





Figure 12. Effect of Jet Mass Flow on Base Drag



(b) Configuration Nos. 6, 9, 10 and 11

Figure 12. Concluded









 $P_c / P_{co} = 0$

K / P_∞ = 168

M= 2.50





 $p_{c} / p_{\infty} = 315$

M = 3.50

Figure 14. Schlieren Photographs of Base Flow Field for Boattail No. 6

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14. LINK D LINK A REY WORDS LINK C ROLE WT ROLE WT ROLE W T Supersonic Flow Base Drag Base Rocket Boattail Drag Base Pressure Boattail Pressure Unclassified Security Classification