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STOL TACTICAL AIRCRAFT INVESTIGATION, EXTERNALLY BLOWN FLAP. VOLUME II. DESIGN COMPENDIUM

Marshall H. Roe, et al

Rockwell International Corporation

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Air Force Flight Dynamics Laboratory

April 1973

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From an overall assessment of study results, it is concluded that the EBF concept provides a practical means of obtaining STOL performance for an MST with relatively low risk.

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STOL TACT:CAL AIRCRAFT INVESTIGATION-EXTERNALLY BLOWN FLAP

Volume II

Design Compendium

M.H.ROE D.J.RENSELAER R.A.QUAM ET AL

APRIL 1973

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FOREWORD

This report was prepared for the Prototype Division of the Air Force Flight Dynamics Laboratory by the Los Angeles Aircraft Division, Rockwell International. The work was performed as part of the STOL tactical aircraft investigation program under USAF contract F33615-71-C-1760, project 643A0020. Daniel E. Fraga, AFFDL/PTA, was the Air Force program manager, and Garland S. Oates, Jr., AFFDL/PTA, was the Air Force technical manager. Marshall H. Roe was the program manager for Rockwell.

This investigation was conducted during the period from 10 June 1971 through 9 December 1972. This final report is published in six volumes and was originally published as Rockwell report NA-72-868. This report was submitted for approval on 9 December 1972.

This technical report has been reviewed and is approved.

E. S. Cross J.

E. J. Cross, Jr. Lt Col, USAF Chief, Prototype Division

ABSTRACT

The basic objective of the work reported herein was to provide a broader technology base to support the development of a medium STOL Transport (MST) airplane. This work was limited to the application of the externally blown flap (EBF) powered lift concept.

The technology of EBF STOL aircraft has been investigated through analytical studies, wind tunnel testing, flight simulator testing, and design trade studies. The results obtained include development of methods for the estimation of the aerodynamic characteristics of an EBF configuration, STOL performance estimation methods, safety margins for takeoff and landing, wind tunnel investigation of the effects of varying EBF system geometry parameters, configuration definition to meet MST requirements, trade data on performance and configuration requirement variations, flight control system mechanization trade data, handling qualities characteristics, piloting procedures, and effects of applying an air cushion landing system to the MST.

From an overall assessment of study results, it is concluded that the EBF concept provides a practical means of obtaining STOL performance for an MST with relatively low risk. Some improvement in EBF performance could be achieved with further development - primarily wind tunnel testing. Further work should be done on optimization of flight controls, definition of flying qualities requirements, and development of piloting procedures. Considerable work must be done in the area of structural design criteria relative to the effects of engine exhaust impingement on the wing and flap structure.

This report is arranged in six volumes:

Volume I	- Configuration Definition
Volume II	- Design Compendium
Volume III	- Performance lethods and Takeoff und Landing Rules
Volume IV	- Analysis of Wind Tunnel Data
Volume V	- Flight Control Technology
Part I Part I Part I	 Control System Mechanization Trade Studies Simulation Studies/Flight Control System Validation Stability and Control Derivative Accuracy Requirements and Effects of Augmentation System Design
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Volume VI - Air Cushion Landing System Trade Study

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LIST OF SYMBOLS

196.54

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А	Aspect ratio
Aj	Nozzle exhaust area per engine, ft ²
b	Wing span, ft
с	Local chord, ft
c'	Wing chord on engine centerline, or midway between engines
	on four engine configurations, ft
ē	Mean aerodynamic chord, ft
Ĉ r	MAC of flapped area of wing, developed chord length. ft
C _D	Total drag coefficient, D/qS
$(\tilde{\Delta}C_{\rm D})_{\rm R}$	Ram drag coefficient, Dr/qS
C _D ;	Drag coefficient due to lift with jet reaction force removed
CL	Total lift coefficient, L/qS
CL	Maximum lift coefficient, LMAX/qS
Contact	Rolling moment coefficient, 2/qSb
Cm	Pitching moment coefficient, M/qSZ
C_	Yawing moment coefficient, M/qSb; section load coefficient
C ¹	Pressure coefficient, P/qS
С́я	Chordwise distance from intersection of jet reaction vector
	with reference plane to leading edge of chord at thrust
	centerline, ft
С _и	Blowing coefficient, T/qS
Cú	Sectional blowing coefficient, equal to $C_{\mu}/\lambda = C_{\mu}(\frac{S}{C_{\mu}})$
Dj	Diameter of exhaust nozzle, ft
Ej	Energy of jet exhaust per unit volume, lb/ft [*]
F	Aspect ratio correction factor
FA	Static axial jet reaction force, lb
PN	Static normal jet reaction force, 1b
F _R	Resultant static jet reaction force, 1b
٥F	Flap deflection angle, degrees. The angle formed by the wing
	reference line and a line bisecting the flap trailing
	edge closing angle (see Figure 12). Nearly all deflec-
	tion angles quoted in the data sources were revised to
	conform with the present definition.
g	Acceleration due to gravity, ft/sec4
h	Height of horizontal tail above wing reference plane, ft
Ĩ	Tail arm, distance from 0.25 Qy to 0.25 Q ₁ , ft
∆ℓ _R	Height of c.g. above centerline of engine inlet, ft
DeT.	Length of moment arm of engine thrust axis about moment
	reference center, ft
N	Normal torce or number of engines
° _N	Deflection angle of engine exhaust deflector, or nozzle,
-	døgtees

LIST OF SYMBOLS - Continued

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$\Delta_{\mathbf{p}}$	Exhaust pressure differential across nozzle exit, lb/ft ²
q	Freestream dynamic pressure, $1/2\rho V^2$, $1b/ft^2$
R	Radius of jet plume at flap trailing edge, ft
S	Wing reference area. ft^2
S'	Wing area affected by flap blowing, ft ²
t/c	Airfoil thickness ratio
Т	Nozzle exhaust static thrust, lb
δT	Angle between freestream and engine nozzle centerline, degrees
USB	Upper surface blowing
v	Velocity, ft/sec
Wa	Rate of engine airflow, 1b/sec
x	Chordwise dimension, ft
x _{c.p.}	Location of center of pressure, ft
XL.E.C'	Location of leading edge of c', ft
xL.E.C _F	Location of leading edge of c _F , ft
x _{REF} -	Location of moment reference center, ft
у	Spanwise dimension, ft
Z	Vertical dimension, ft
zF	Distance which flap trailing edge extends into jet plume
	perpendicular to engine centerline, ft
¢	Angle of attack, degrees or radians
r	Circulation
Y	Flight path angle, degrees
Δ	Incremental value
δ	Deflection angle, degrees
ε	Downwash angle, degrees
η	Jet turning efficiency factor
θ	Effective jet turning angle, radians
λ	Area ratio, S'/S
л	Wing sweep of 0.75 chord, degrees
ν	Lift slope correction for partial span blowing
π	Ratio of circumference of a circle to its diameter
ρ	Density, slugs/ it
ψ	Yaw angle, degrees
φ	$(C_{L\alpha})_{PO}/C_{L\alpha}$ cure contrare
	SUBSCRIPTS
()0	Baseline nover-off data
()ECT	Ectimated
	Everymental data
	Horizontal tail
()i	let exhaust
()) ())), F	Leading edge
()ΜΔΥ	Maximum value
τ <u></u> <u>μ</u> αγγγ	

SUBSCRIPTS (Continued)

-)MEAS. Measured value (
-)MIN. Minimum value

<u>)</u>0 Zero angle of attack ((

)PL Powered lift

)P.O. Power off

(

(

(

((

(

) R Due to jet reaction (

Ram drag) RD

About moment reference center)REF ((

)TEST Test data (

-)V At forward speed
- Due to jet deflection ()θ
 - Two dimensional value) ထ
 - Due to angle of attack)α
 -)r Circulation lift
 -)Ē Flaps

Section I

INTRODUCTION

The externally blown flap lift-propulsion system has recently evolved in design studies of advanced short takeoff and landing (STOL) aircraft. This lift-propulsion system is attractive because it generates a significant amount of wing lift at typical STOL speeds while it is not as complex as other lift-propulsion systems.

Because STOL aerodynamic data for this lift-propulsion system are relatively new it has been extremely difficult for the aerodynamicist to estimate effects of geometric perturbations that are normally encountered during preliminary design phases. It is therefore appropriate to collect and organize the basic aerodynamic information into a manual and to develop methods of estimation. This will give improved credence to STOL performance data for this lift-propulsion system and better comparison with the performance of other lift-propulsion systems. To provide such a manual is a purpose of the present design compendium.

The design compendium presents methods to determine the power effects on lift, drag, pitching moment, downwash, lateral and directional moments, and some dynamic derivatives. The methods are presented for possible inclusion in the DATCOM stability and control handbook. Increments due to power are given, so that these can be added to power off data from existing sections of DATCOM. Occasionally, when insufficient data were available to generate a method, some guidelines for the fairing of curves are given.

The general approach used here is based on jet flap theory with some empirically guided interpretations in application to the externally blown flap. Although the jet flap theory strictly applies only to inviscid, incompressible flow at small incidence angles, small jet turning angles, and a thin jet sheet, surprisingly good agreement is found with experimental data where these conditions are rather grossly exceeded.

The data and methods in this compendium are often given for the determination of characteristics at zero angle of attack, rather than for the angle at zero lift that is commonly used in conventional aerodynamics. The zero angle of attack is chosen because the zero lift angle results in extreme negative values for STOL airplanes in the power mode. These negative angles are generally unattainable in flight because of flow separation at the lower wing surface. On the other hand, STOL takeoff and landings are often made with angles of attack relatively close to zero.

Section II

MAJOR CORRELATING PARAMETERS

The forces developed by flap blowing are functions of the total jet momentum, the turning angle of the jet, and the geometry of the wing-flap system. The jet momentum is usually expressed as a coefficient, $C_{\mathcal{U}}$, and in conformance with jet flap theory the momentum exiting from the flap trailing edge is used in estimation methods. Where $C_{\mathcal{U}}$ is the jet nozzle momentum, the flap exit momentum is defined as $\gamma C_{\mathcal{U}}$, where γ is an efficiency factor accounting for energy losses incurred in the turning process. The turning angle, θ , is the effective angle at which the jet leaves the flap. Derivation of these parameters, and methods of estimation are discussed **balow**.

2.1 BLOWING COEFFICIENT

The lift and normal force at forward speed appear to be affected predominantly by the ratio of the energy of the freestream air and the energy exhausted by the blowing nozzle, as well as by geometric relations. The energy from the exhaust nozzle is, expressed in terms of unit exhaust volume:

$$E_{j} = \frac{1}{2} e_{j} V_{j}^{2}$$
2.1

The use of the freestream dynamic pressure $q = 1/2 \ Q \ V^2$ to nondimensionalize this energy yields the parameter

$$\frac{e_{j} \sqrt{j}^{2}}{e \sqrt{2}}$$

Herein, $\varrho_j V_j^2$ can be expressed in terms of the nozzle exhaust thrust. Denoting T_V as the nozzle exhaust thrust at forward speed, and introducing the symbol Aj as the nozzle exhaust area, a relation between $\varrho_j V_j^2$ and the nozzle thrust is obtained as follows:

$$T_{V} = (P_{j} \vee_{j} A_{j}) \vee_{j} + \Delta P A_{j}$$
 2.2

Herein Δ_p is the exhaust pressure differential across the nozzle exit, which is zero for unchoked exhaust typical for external blowing, so that

$$\frac{e_{j}v_{j}^{2}}{e_{v}v^{2}} = \frac{\tau_{v}/A_{j}}{2q}$$
2.3

or

$$\frac{P_{j}V_{j}^{2}}{P_{j}V^{2}} = \frac{T_{v}}{95} \left(\frac{5}{2A_{j}}\right)$$

Because for a given configuration the value of $\frac{S}{2Aj}$ is fixed, it is seen that the energy ratio is proportional to C_{AJ} :

$$\frac{P_{i}V_{i}^{2}}{P_{i}V_{i}^{2}} = \text{constant} \cdot C_{uv} \qquad 2.4$$

where, per definition

$$S_{\rm my} = \frac{T_{\rm v}}{9.5}$$
 2.5

The validity that the force characteristics correlate with the energy ratio, i.e., the blowing parameter, has been substantiated by several investigations wherein engine thrust and q were varied, and $C_{\mathcal{M}}$ held constant.

It should be noted that T_V is the exhaust thrust of the nozzle, and not the net engine thrust. The net engine thrust is equal to the ekhaust thrust minus the intake momentum drag. Only the exhaust thrust is used in this section because it is assumed that only the energy from the exhaust is determining the airfoil lifting characteristics from blowing regardless of the engine inlet flow characteristics.

The exhaust thrust, T_V , generally increases slightly with increase in speed. Often, the increase of T_V with speed is not quantitatively known, and therefore only the static thrust is taken as the reference thrust. Advantages of this procedure lie in the fact that the thrust used is a constant which makes it very suitable for use as a reference and no intake momentum drag exists in this condition. Thus, the use of the definition of the coefficient

$$c_{11} = \frac{1}{95}$$
 2.0

and a second a straight second and a second s

is favored, where T is the static thrust.

2.2 JET TURNING ANGLE

The jet turning angle, θ , is the effective direction at which the jet leaves the trailing edge of the flap system and defines the direction of the total reaction force vector, F_R (see Figure 1). The jet turning angle and reaction force components are determined for the static condition, since the effects of forward velocity on these values are **hot** known, and would be very difficult to define. By definition, then:

$$\theta = \tan^{-1} F_{\rm N}/F_{\rm A}$$
 2.7

2.8

 F_N * static normal force

 F_A = static axial force

$$F_R = (F_N^2 + F_A^2)^{1/2}$$
 = resultant reaction force.



Figure 1. Definition of Jet Turning Angle

2.3 EFFECTS OF FLAP IMPINGEMENT ON JET TURNING ANGLE

Direct application of the jet flap theory to the externally blown flap would assume that the jet turning angle, θ , is equal to the flap deflection angle, δ_F . This has proved to be a good assumption for the case where the flap system captures the entire jet efflux; however, where the flap intercepts only a part of the jet, the effective turning angle will be less than the flap angle. The following section presents an approach to estimating the jet turning angle for externally blown flaps, including those cases where there is less than full impingement of the jet exhaust on the flap.

Using a heuristic approach, it was assumed that the effective jet turning angle is related to the flap angle and to the portion of jet momentum that is intercepted by the flap. Following this reasoning correlations were made of experimental values of jet turning angle as a function of flap angle, θ/δ_F , and the extent of flap penetration into the jet. The results of these correlations are shown in Figures 3 through 5 for various types of flaps and wing sweep angles. In these figures the extent of flap penetration into the jet is given in terms of the ratio Z_F/R , where Z_F is the distance that the flap penetrates into the jet. The definition of this parameter, Z_F/R , is illustrated in Figure 2. The value of the radius, R, of the jet at the location of the flap trailing edge can be estimated by the relation:





Figure 2. Definition of Impingement Parameter

where D_j is the diameter of the exhaust nozzle, and X is the distance of the flap trailing edge behind the nozzle exhaust, Figure 2. The effective source length, 2.3 D_j , used here is based on a jet expansion envelope where the jet velocity is essentially zero at the edge of the jet wake profile. Other definitions such as one percent or five percent of the maximum jet velocity could be used, which would only change the scale of the Z_F/R parameter. Only an average spreading angle is used here, although according to Ribner, Reference 7, the spreading depends to some extent on the thrust coefficient.

The data from these correlation plots Figures 3 through 5, were used to develop a set of design curves, shown in Figure 6, for estimating jet turning angle values. These disign curves represent envelopes of the better performance data from the individual correlation plots. The data correlations have considerable scatter resulting from non-optimum flap geometry and variations in experimental technique. However, values of jet turning angle from these design curves have been used in substantiation calculations of the methods for prediction of aerodynamic characteristics of externally blown flaps in later sections of this report, with generally very good results (Section III, IV and V).

It is noted in the design curves, Figure 6, that the only flap geometry parameter that appears is the number of flap segments (double and triple slotted flaps). Other variables such as wing sweep, and perhaps aspect ratio, would be expected to influence the jet turning angle. However, the effects of these other variables are evidently of lower order and are lost in the data scatter.

Separate curves are shown for data with jet deflectors. These devices have been tested in several investigations and found to improve the lift augmentation of externally blown flaps (References 4, 8, 16, 20). These deflectors are simple flat plates that deflect the jet exhaust upward to increase the degree of impingement of the jet on the flap. In effect the deflectors are increasing the impingement ratio, $2_{\rm F}/T$, which has been shown above to influence the effective jet turning angle.

For purposes of predicting the effectiveness of jet deflectors, the procedure described above utilizing the impingement parameter, can not be used directly since the deflectors change the characteristic shape of the jet expansion envelope in some undefined manner so the impingement ratio of the deflected jet is not known. However, results of tests of deflectors from the referenced tests have provided effective turning angle data as shown in Figures 3 through 5, which are shown as a function of the impingement ratio of the undeflected jet, which can be determined.

The available data on the effectiveness of jet deflectors in terms of effective jet turning angle have been summarized in Figure 6. Here, a

recommended design curve for EBF systems utilizing jet deflectors is presented as a function of the impingement parameter of the undeflected jet. As indicated by this plot, there is no distinguishable difference between the effectiveness of double or triple slotted flaps when utilizing the deflector. The procedure for estimating the effective jet turning angle with deflectors is, then, to define the impingement ratio of the undisturbed jet, as described above, and then utilize the design curve labeled "with jet deflectors."

The effects of engine location relative to the flap, and the effects of nacelle incidence and nozzle deflection are all accounted for in this use of the impingement parameter.

Test data show nearly complete jet turning with flap immersion of Zp/R > 0.65 except for cases where flow separation over the flaps is suspected. It is concluded then, that in the design of the propulsion-lift system that the design goal should be to attain impingement of at least Zp/R = 0.6 for high lift performance.

Those data that fall significantly below the design curves are in some cases extreme flap angles $(75^{\circ} to 95^{\circ})$ which apparently did not have proper blowing to maintain flow attachment, and in other cases rather unconventional flap segment arrangements which evidently resulted in separation. In view of these results it would probably be prudent to not apply these methods to flap angles exceeding 65° unless experimental verification can be obtained.



Figure 3. Jet Turning Angle Ratio for Double Slotted Flaps - 0° Sweep







Figure 5. Jet Burning Angle Ratio for Triple Slotted Flaps - 25* Sweep



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2.4 JET TURNING EFFICIENCY

In applying the relationships of jet flap theory to the externally blown flap, the blowing momentum leaving the flap trailing edge should be used. This exiting momentum is defined as γC_{μ} , where the factor γ accounts for all of the losses in the jet turning process. The turning efficiency can be determined from static tests, where it is defined as:

$$\eta = Fiz/T$$
 2.10

where $F_R = (F_N^2 + F_A^2)^{1/2}$

 $F_{\rm N}$ = static normal force

 F_A = static axial force

T = nozzle thrust

It would be expected that the turning losses would depend on the degree of turning and the details of the flap geometry.

For use in prediction methods, data correlations have been made of the turning efficiency versus effective turning angle where effects of the number of flap slots (single, double, triple), and (f wing tweep are apparent. Considerable scatter in the test results exist; which no doubt is due to details of slot gaps, various amounts of separation on the flap segments, differences in flap segment contours, and experimental techniques (Figures 7 through 10). Curves have been established on Figure 11 for design use which are considered to be representative of better designed externally blown flap systems.



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Figure 7. Jet Turning Efficiency for Double Slotted Flaps - 0° Sweep



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Figure 9. Jet Turning Efficiency for Triple Slotted Flaps - 25° Sweep

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Section III

LIFT

The tail-off lift of an airplane with an externally blown flap is estimated by adding the effects of blowing to the power-off lift at zero angle of attack with flap deflected. The power-off lift id estimated by conventional methods, such as DATCOM, or is based on unpowered wind tunnel

(NOTE: Estimation of the power-off lift due to flap deflection at zero angle of attack can be included in the present methods, as will be discussed later; however, accuracy should be improved by use of conventional methods as stated above.)

$$C_{\rm L} = (C_{\rm L})_{\rm B} + (\Delta C_{\rm L})_{\rm PL}$$

$$3.1$$

where $(C_L)_B$ is the baseline lift at $\infty = 0$, flaps deflected, power-off; $(\Delta C_L)_{PL}$ powered lift increment due to jet deflection and angle of attack.

Estimation of the powered lift increment is based on the jet flap theory developed by Spence (Reference 10), which gives for the lift of a two-dimensional thin airfoil with a jet flap:

$$C_{L} = (C_{L_{\theta}}) \theta + (C_{L_{\infty}}) \sigma$$

-

where $(C_{L\theta})_{oo}$ is the two-dimensional lift gradient with jet deflection angle, θ , and $(C_{L,\infty}) \longrightarrow$ is the two-dimensional gradient of lift with angle of attack, ∞ .

3.2

In application of this theory to an airfoil of finite thickness and finite aspect ratio, the pressure lift terms are modified for thickness by the factor (1 + t/c), and for aspect ratio by a factor F, derived by Maskell and Spence (Reference 11). The derivation of the aspect ratio correction factor is based on the assumption of elliptical spanwise distribution of chord length and blowing coefficient. The resulting expression for the lift of a three-dimensional wing with a full span jet flap, as given in Reference (11), is:

$$C_{L} = F \left[(1 + \frac{t}{c}) \left\{ (C_{L\theta})_{\Theta \Theta} \theta + (C_{L\alpha c})_{\Theta \Theta} \Theta \right\} - \frac{t}{c} (\theta + \infty) C_{\mu} \right]$$
3.3

Further development by Williams, Butler, and Wood (Reference 12) to account for partial span blowing provided correction factors λ and μ for the jet deflection term and angle of attack term, respectively.

Adaptation of these methods for the jet flap to the case of the externally blown flap is based on the following considerations: The lift increment due to power is estimated by the jet flap methods and added to the power-off lift of the wing with flaps deflected; an calculation of the lift due to jet deflection the sine of the turning angle is used rather than the angle, since the theory was developed on the basis of small angles for which $\sin \theta \approx \theta$; the aspect ratio correction factor which is based on elliptical spanwise distribution of blowing, is not strictly applicable to the concentrated blowing technique of the EBF, but its use gives better results than the unpowered factor, A/(A + 2); the thickness correction factor is applied to the entire lift increment due to power rather than excluding the direct jet lift which results in a small discrepancy with the theory but improves the experimental correlations and simplifies the calculations.

With the above considerations, the expression for the lift increment due to power for an externally blown flap airplane is:

$$(\Delta C_L)_{pL} = F(1 + \frac{t}{c}) \left[\lambda (C_{L_{\theta}})_{\infty} \sin \theta + \mathcal{V} (C_{L_{\theta}})_{\infty} \infty \right] 3.4$$

It is more convenient to treat the total lift increment due to power in its components due to jet deflection and angle of attack:

$$(\Delta C_L)_{pL} = (\Delta C_L)_{\theta} + (\Delta C_L)_{\infty}$$
 3.5

$$(\Delta C_L)_{\theta} = F (1 + \frac{t}{c}) \lambda (C_{L_{\theta}}) \sigma \sin \theta$$
 3.6

$$(\Delta C_L)_{\infty} = F (1 + \frac{t}{c}) \mathcal{V} (C_L \infty)_{\infty} \infty \infty$$
 3.7

Further development of procedures for calculating these lift increments is given in Sections 3.2 and 3.3. A discussion of the geometrical parameters involved in these calculations is given in the next section.

3.1 GEOMETRIC PARAMETERS

The methods of estimation in the following sections make use of various geometrical parameters which will be defined here.

3.1.1 FLAP DEFLECTION ANGLE

The flap deflection angle, δ_F , as used herein for under the wing blowing is defined as the angle from the wing reference plane to the bisector of the trailing segment of the flap. For over the wing blowing the deflection angle is taken from the wing reference plane to the upper surface of the trailing edge of the flap. These angles are illustrated in Figure 12.



Figure 12. Definition of Flap Deflection Angle

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3.1.2 ASPECT RATIO FACTOR

An approximate relation for the aspect ratio correction factor, F, given in Reference (12), and replacing $C_{\mu\nu}$ with $\gamma_{\mu}C_{\mu\nu}$, is:

$$F \approx \frac{A + 2\eta \zeta_{u} f_{T}}{A + 2 + 0.604 (\eta \zeta_{u})^{t_{u}} + 0.876 \eta \zeta_{u}} \qquad 3.8$$

where C'_{u} is the sectional thrust coefficient = $C_{u} \left(\frac{S}{S'}\right)$.

A plot of this factor for various values of A and ηc_{\perp} is shown in Figure 13.

For the case of part span flaps, the aspect ratio is taken as that of the wing excluding half of the area outboard of the flaps.

3.1.3 EFFECTIVE WING AREA

The effective wing area, S', is defined as the area of the wing between the inboard station and outboard station of the flap (see Figure 14). If the wing has part span flaps the effective wing area includes onehalf of the area autboard of the flaps. This definition was derived from consideration of pressure data from Reference (14), which shows carry over on the wing outboard of the flap of about one half the loading on the flapped portion of the wing. Use of this definition in estimating characteristics of part span configurations has produced good correlations with test data (see Section 3.8). (NOTE: Limited data for upper surface blowing configurations indicate that only that portion of the wing surface that is in the exhaust flow should be included in the definition of S' for upper surface blowing. For example, S'/S for the configuration of Reference 22 is approximately 0.5.)

3.1.4 MEAN AERODYNAMIC CHORD OF EFFECTIVE WITH EA

The MAC of the effective wing area, $\overline{c_F}$, is determined by conventional methods applied to the effective wing area, S', as described above, but the chord length is taken as the developed chord length with the flaps extended as shown in Figure 15.

For full span flaps \overline{c}_F can be taken as the NAC of the basic wing, \overline{c} , corrected for flap chord extension.

3.1.5 WING THICKNESS RATIO

The wing thickness ratio, t/c, is taken as the thickness ratio of the MAC of the flapped portion of the wing.



Figure 13. Aspect Ratio Correction Factor

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Figure 14. Effective Wing Area



Figure 15. Definition of MAC With Flaps Extended

3.2 LIFT AT ZERO INCIDENCE

The lift increment due to power at zero angle of attack is defined by:

$$(\Delta C_L)_{\theta} = F (1 + \frac{C}{c}) \lambda (\partial C_L/\partial \theta) \cos \sin \theta$$
 3.6

where F, $r_{\rm L}$, t/c, and θ have been defined in previous sections, and

 λ = S'/S = area ratio defined in Section 3.1.3

 $(\partial C_{I}/\partial \theta)_{ee}$ = two-dimensional lift due to jet turning angle.

A plot of $(\partial C_L/\partial \theta) \sim as$ defined in Reference (10) is shown in Figure 16. The curve for $C_F/C = 0$, corresponding to a pure jet flap, should be used in these estimation procedures. An approximate interpolation formula for the $C_F/C = 0$ case, giver in Reference (11) is, substituting $\Lambda C'\mu$ for $C_{\mu\nu}$:

$$(\partial C_{L} / \partial \theta) \approx 4\pi \eta C' \mu \left\{ 1 + 0.151 (\eta C' \mu)^{1/2} + 0.139 \eta C' \mu \right\} \right]_{3.9}^{1/2}$$

where C' μ , again, is the sectional thrust coefficient = C μ (S/S').

NOTE: The power-off increment of lift due to flap deflection can be included in the methods of estimation by using the relation of $(\partial C_{\rm L}/\partial \theta) \infty$ as a function of flap chord as shown in Figure 16. In this case the ratio of flap chord to wing chord should be defined in the flap-extended condition. The lift increment calculated by this method should be added to the flaps-up, power-off, lift to obtain total lift at $\infty = 0$. It is believed, however, that other methods of estimation of the power-off lift due to flap deflection, such as DATCOM, or wind tunnel data, will give better accuracy.

The expression here for $(\Delta C_L)_{\theta}$ was developed from a small angle theory; however, it is used herein for jet turning angles, θ , exceeding one radian. It is therefore believed to be more appropriate to use sin θ rather than θ in the use of these expressions. This substitution has been made in all subsequent calculations.

It is more convenient in the pitching moment calculations to separate the components of the lift increment at zero incidence into the thrust reaction term, $\eta C \mu \sin \theta$, and the circulation lift correspondent (ΔC_L);

$$(\Delta C_{\rm L})_{\theta} = (\Delta C_{\rm L})_{\eta} + \eta C_{\mu} \sin \theta \qquad 3.10$$

then
$$(\Delta C_L)_{\mu} = (\Delta C_L)_{\theta} - \eta C_{\mu} \sin \theta$$

= $F(1 + t/c) \lambda (C_{L\theta})_{00} \sin \theta - \eta C_{\mu} \sin \theta$ 3.11

It should be noted that the factor F, and the two dimensional CL_{θ} are based on the sectional blowing coefficient, \mathcal{N} $C'\mu$. The thrust reaction term represents a discrete force and is defined by the blowing coefficient, $C\mu$, based on the wing reference area, and the static values of \mathcal{N} and θ .

Comparisons are shown in Figure 17 of estimated and experimental values of $(\Delta C_L)_{\theta}$ for a wide range of model configurations. The results show reasonably good agreement, being better than +10 percent in most cases. The lift estimation is quite sensitive to the parameters η and θ , so these must be determined as closely as possible. As discussed in Section II, the methods presented here for estimating η and θ tend to represent well designed lift systems that produce good performance; whereas some of the wind tunnel models used in these correlations perform below average.

It should be noted in these correlations that in those cases where the nozzle thrust was inclined relative to the wing to improve impingement on the flap, the measured lift data were corrected for the downward thrust reaction force by:

 $(\Delta C_L)_{\theta \in XP} = (\Delta C_L)_{\theta \in MEAS} + C \mu \sin \delta_T$

where d_T is the angle between the nozzle centerline and the reference plane.



Figure 16. $(\partial C_{I}/\partial \Theta)_{\infty}$ and $(\partial C_{L}/\partial \infty)_{\infty}$ for a Jet Flap



Figure 17. Correlation of Lift Increment due to Power at Zero Incidence

3.3 LIFT DUE TO ANGLE OF ATTACK

The powered lift increment due to angle of attack is estimated by:

$$(\Delta C_L)_{\alpha} = \left[F(1+\frac{t}{\epsilon})Y(\partial C_L/\partial x)_{\infty}\right] \alpha$$
 3.7

where F, η , θ and t/c have been defined in previous sections. The correction factor for partial span blowing, V, from Reference (12), is:

$$V = 5/5 + (1 - 5/5) 2\pi (\partial L_1 / \partial \alpha) = 3.12$$

This function is plotted in Figure 18.

 $(\partial C_L/\partial \alpha)_{\infty}$ = two dimensional lift due to angle of attack (Reference (10)). A plot of $(\partial C_L/\partial \alpha)_{\infty}$ versus $\gamma C'_{\perp}$ is shown in Figure 16. An approximate analytical expression from Reference (11), again substituting $\gamma C'_{\perp}$ for C_{\perp} , is:

$$(\partial (L/\partial x)_{\infty} = 2\pi [1 + 0.151 (\eta (J_{1}))^{2} + 0.219 \eta (J_{1})]_{3.13}$$

This procedure is of course limited to the linear range of lift with angle of attack. The extent of the linear range is strongly dependent on leading edge treatment, and varies also with flap deflection, thrust coefficient, aspect ratio, and sweep.

Correlations of estimated versus experimental values of the lift curve slope as defined by $(\Delta C_L)_{\alpha}$ are shown in Figure 19 for a wide range of test configurations. It appears that on the average the method of estimation slightly over-estimates: the value of the lift curve slope, although the agreement is within +10 percent for nearly all of the data.



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Figure 18. Lift Curve Slope Correction for Partial Span Blowing

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Figure 19. Correlation of Lift Curve Slope

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3.4 TOTAL LIFT ESTIMATION

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The recommended procedure for estimating the lift coefficient is to determine power-off lift at zero angle of attack by other conventional methods of estimation; then add the power effects due to jet turning, angle of attack and nozzle incidence:

$$C_{L} = (C_{L})_{B} + (\Delta C_{L})_{\theta} + (\Delta C_{L})_{\alpha} - C_{\mu} \sin \delta_{T} \qquad 3.14$$

where detailed description of the lift increments are given by:

 $(\Delta C_L)_{\theta}$ from Section 3.2

 $(\Delta C_L)_{old}$ from Section 3.3

3.5 MAXIMUM LIFT COEFFICIENT

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A relationship for the maximum lift coefficient of a two-dimensional wing with a jet augmented flap, based on the maximum lift of the unflapped wing, and assuming a leading edge stall for both wings, is given by D.N D.N. Foster in Reference (15), as:

$$C_{L_{MAX}} = \frac{C_{L_{MAX}}}{\delta, c_{\mu}} - \frac{\Delta C_{L_{\delta, c_{\mu}}}}{2} - \frac{\Delta C_{L_{c_{\mu}}}}{4} \qquad 3.15$$

where, in the nomenclature of Reference 15, C_{LMAX} is the maximum lift coefficient of the unflapped, unaugmented wing; C_{LMAX} $_{\delta}$, $_{C\mu}$ is the maximum lift coefficient of the wing with flap deflected and jet augmentation; $_{\Delta C_{L\delta}}$, $_{C\mu}$ is the increment of lift due to flap deflection and jet augmentation; and, $_{\Delta C_{LC\mu}}$ is the increment of lift due to jet augmentation only, with the flap already deflected

The increment in maximum lift coefficient due to flap deflection and jet augmentation is, then:

$$\begin{pmatrix} \Delta C_{L} \\ MAX \end{pmatrix} = C_{L} - C_{L} \\ MAX_{\delta}, c_{\mu} \\ = 1/2 \begin{bmatrix} \Delta C_{L} + 1/2 & \Delta C_{L} \\ \delta, c_{\mu} \\ \end{bmatrix}$$

by definition $\Delta C_{L_{\delta}, c_{\mu}} = \Delta C_{L_{\delta}} + \Delta C_{L_{c_{\mu}}}$ 3.16

so
$$\left(\Delta C_{L_{MAX}}\right)_{\delta, c_{\mu}}^{\cdot} = \frac{1/2}{\left[\Delta C_{L_{\delta}}^{\bullet} + \Delta C_{L_{c_{\mu}}}^{\bullet} + \frac{1/2\Delta C_{L_{c_{\mu}}}^{\bullet}}{c_{\mu}}\right]}$$

= $\frac{1/2}{\left[\Delta C_{L_{\delta}}^{\bullet} + \frac{3/2\Delta C_{L_{c_{\mu}}}^{\bullet}}{c_{\mu}}\right]}$ 3.17

Now according to Reference 24, the increment in maximum lift coefficient due to flap deflection only is equal to one-half the increment in lift at constant angle of attack due to flap deflection:

$$\binom{\Lambda C_{L_{MAX}}}{MAX}_{\delta} = \frac{1/2}{1/2} \frac{\Lambda C_{L_{\delta}}}{L_{\delta}}$$
 3.18

So, defining $(\Delta C_{LMAX})_{cu}$ as the increment in maximum lift coefficient due to jet augmentation only, with the flap already deflected:

$$\begin{pmatrix} \Delta C_{L} \\ MAX \end{pmatrix}_{C_{\mu}}^{\ast} = \begin{pmatrix} \Delta C_{L} \\ MAX \end{pmatrix}_{\delta_{i}}^{\ast} C_{\mu} \begin{pmatrix} \Delta C_{L} \\ MAX \end{pmatrix}_{\delta}^{\ast}$$

= $1/2 \begin{bmatrix} \Delta C_{L_{\delta}}^{\ast} + 3/2 \Delta C_{L_{C_{\mu}}} \end{bmatrix} - 1/2 \Delta C_{L_{\delta}}^{\ast} = 3/4 \Delta C_{L_{C_{\mu}}}^{\ast} 3.19$

In the derivation of $C_{L_{MAA}}$ in Reference 15, the term $\Delta C_{LC_{\mu}}$ is defined as the increment in lift due to jet augmentation at "constant angle of attack, and is then further interpreted as being at $\alpha = 0$. However, it is believed that the value of $(\Delta C_{L})_{C_{\mu}}$ to be used here should be defined at the angle of attack for stall with flap deflected and with blowing, as suggested by Moorhouse in Reference 23. The reason for this, on which the analysis of Reference 15 is based, is that the leading edge loading of a blown flapped wing is the same as that of an unblown flapped wing when the C_{L} of the unblown wing is equal to $C_{L} - 3/4$ ($\Delta C_{L})_{C_{\mu}}$ of the blown wing.

In applying this approach to the estimation of $C_{L_{MAX}}$ for the finite aspect ratio case, it is reasoned that the $(\Delta C_L)_{C_{\mu}}$ used in the calculation should be the two-dimensional value, since it is the section loading that determines $C_{L_{MAX}}$. However, it is usually the three-dimensional data that are available, so the three-dimensional value of $(\Delta C_L)_{C_{\mu}}$ should be modified for aspect ratio by the factor, 1/F. This procedure was first suggested by McRae in Reference 24.

The maximum lift increment due to flap blowing for the finite aspect ratio case is then:

$$\begin{pmatrix} \Delta C_{L} \\ MAX \end{pmatrix}_{PL} = \frac{3}{4F} (\Delta C_{L})_{PL}, \alpha_{MAX}$$
 3.20

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Now since the value of α_{MAX} is not known, the above expression has been defined in terms of known values, based on the simplified lift curves illustrated in Figure 22:

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$$(C_{L_{MAX}})_{PL} = \frac{3/(4F)[(\Delta C_{L})_{\theta} \diamond - (C_{L})_{PO}, \alpha = 0 \ (1-\phi)] + (C_{L_{MAX}})_{PO} - C_{\mu} \sin \delta T}{1 - \frac{3}{4F} \ (1-\phi)} \qquad 3.21$$

where:
$$\phi = \frac{(C_{L_{\alpha}})_{PO}}{C_{L_{\alpha}}}$$

and $-C_{\mu} \sin \delta_{T}$ is a correction to $C_{L_{MAX}}$ for those EBF systems which incorporate jet nozzle incidence to improve lift augmentation.

A correlation of experimental data versus estimates using equation 3.2.1 has been made representing many variations of aspect ratio, sweep angle, flap configuration, blowing coefficient, and leading edge devices. The results show a fairly consistent underestimation of $\Delta C_{L_{MAX}}$ of about 15 percent. Application of a multiplying factor of 1.15 to the first term of equation 3.2.1 was found to bring the data into reasonably good agreement, as shown in figure 20, although it caused the upper surface blown data to spread. The resulting expression for estimation of $C_{L_{MAX}}$ for EBF systems is, then:

$$\begin{pmatrix} C_{L_{MAX}} \end{pmatrix}_{PL} = \frac{\frac{3}{4F} \left\{ [1.15 \ (\Delta C_{L})_{\rho} \ \phi - \ (C_{L})_{PO}, \alpha = 0 \ (1-\phi)] + \ (C_{L_{MAX}})_{PO} \right\} - C_{\mu} \sin \delta_{T}}{1 - \frac{3}{4F} F \ (1-\phi)}$$

$$3.22$$

It should be remembered that this relation assumes similar type stall for both power on and power off, which is often not the case for powered lift wind tunnel tests.

It has been found from test data correlations that an adequate approximation for the estimation of ΔC_{LMAX} is:

$$\begin{pmatrix} \Delta C_{L} \\ MAX \end{pmatrix}_{PL} \approx 1/F (\Delta C_{L})_{\theta} - C_{\mu} \sin \delta_{T}$$
 3.23

This expression is much more amenable to quick 'hand calculated' estimations, and gives equally good correlation with test data as the relation given in equation 3.22.



Figure 20. Correlation of ΔC_{LMAX} Due to Power According to Equation 3.22



Figure 21. Correlation of $\Delta C_{L_{max}}$ Due to Power According to Equation 3.23 37

3.6 ANGLE OF ATTACK FOR MAXIMUM LIFT

The change in angle of attack at maximum lift due to jet augmentation can be estimated using the characteristics of idealized lift curves as illustrated in Figure 22, and the relations for $(CL_{MAX})_{PL}$ of the previous section.





The angle of attack at maximum lift power on is approximated by:

$$a_{MAX} = (a_{MAX})_{PO} + (\delta a)_{MAX}$$
 3.25

A correlation of estimated and experimental a_{MAX} data is shown in Figure 23. Reasonably good agreement is observed, with some notable exceptions. Again, the theoretical approach assumes a similar type of stall for both the power-on and power-off caseq, whereas examination of wind tunnel test results indicates obvious variances in the type of stall for



Figure 23. Correlation of Maximum Angle of Attack Data

different values of the blowing coefficient. When this condition exists agreement can not be expected between estimated and experimental data. It would be expected that the estimated values of α_{MAX} would be in better agreement with flight test results of a well designed full scale, EBF system, than with wind tunnel results of small scale models.

3.7 EFFECTS OF ASPECT RATIO

The theoretical aspect ratio correction factor, F, discussed in Section 3.1.2 was developed on the assumption that the jet blowing was distributed elliptically spanwise. In the case of the externally blown flap the blowing is concentrated at the engine nozzle and the lift distribution peaks in the vicinity of the thrust centerline, which violates considerably the assumptions under which the aspect ratio factor was derived. However, better results are obtained by using the aspect ratio factor, F, than are obtained by using the power-off factor, A/(A + 2).

Experimental data from Reference (13), shown in Figure 24, illustrates a negligible effect in varying aspect ratio by 40 percent. Estimated values of lift increment due to blowing for these same model configurations show a larger predicted effect of aspect ratio than that observed in the test data. The error in lift increment is not large, but these results indicate that the aspect ratio factor better accounts for the spanwise loading effects at aspect ratio 7 than at 10.

The comparison of experimental effects of aspect ratio with the theoretical prediction indicates that as the blowing coefficient is increased the effective aspect ratio is decreased. This trend would be expected due to higher loading of the inboard portion of the wing span with the higher blowing coefficients and with the typical EBF configuration.

It is believed that an aspect ratio correction factor could be developed based on loading distributions more typical of EBF systems, which would more accurately predict the aspect ratio effects on powered lift. However the factor, F, used herein is recommended for prediction methods until improved factors can be developed.



3.8 EFFECTS OF FLAP SPAN

Investigation of the effects of varying flap span on the lift augmentation due to flap blowing has shown that use of the effective wing area, S', as defined in Section 3.1.3 adequately accounts for flap span variations. As discussed in Section 3.1.3, available pressure data from Reference 14 indicate that the circulation lift due to blowing on externally blown flaps extends over the entire span of the flaps, regardless of the spanwise placing of the engines. Analysis of the pressure data in Reference (14) shows that the carryover loading on the wing area outboard of the 75 percent span flap is about one-half of that on the flapped portion of the wing. This observed characteristic would obviously not hold for extreme variations from the test configurations, but it is believed to be a valid assumption for reasonable STOL configurations.

The effects of flap span are thus included in the lift calculation through the dependency of the effective area, S', and the effective aspect ratio, on flap span. The correlation plots of lift at $\alpha = 0$, Figure 17, and $C_{L_{\alpha}}$, Figure 19, show equally good correlation of the part-span data as that of the full-span data.

Two sets of data were available (References 13 and 14) where all model geometry except the flap span was held constant. These models were tested with both full span and 75 percent span flaps. Results of these tests are shown in Figure 25. It is seen that the test data follow the trend of the estimated variation of lift with flap span ratio, although the absolute level of the zero sweep data is underestimated.



Figure 25. Correlation of Effects of Flap Span on Lift due to Power

3.9 EFFECTS OF WING SWEEP

In the present methods of estimation of lift augmentation due to flap blowing the effects of wing sweep appear only in the jet turning efficiency, n, as seen in Figures 7, 8, and 11. Figure 11 shows only a modest effect of sweep except for the zero sweep, double slotted flap case, which exhibits a markedly reduced efficiency. Referring to Figure 17 where estimated and experimental values of zero incidence lift are compared it is seen that the lift data for the unswept wing with double slotted flaps are generally underestimated.

It is possible that the turning efficiency data for the zero sweep wing which were taken from References 8 and 14, reflect some experimental technique which is not consistent with the other tests. Further tests of other models with low wing sweep values should be made to investigate the turning efficiency of double-slotted flapped wings at low sweep angles.

Tests were conducted in Reference 13 where only the wing sweep of the model was varied; flap deflection, nacelle geometry, and blowing coefficient were held constant. Results of these tests are shown in Figure 26 for sweep angles of 9° to 30°. The close agreement of estimated lift increments with the experimental data indicates that the prediction methods adequately account for sweep effects.



Figure 26. Correlation of Effects of Wing Sweep on Lift Due to Power

3.10 EFFECTS OF FLAP CONFIGURATION

The methods of prediction of lift increment due to power indicate a small benefit of triple slotted flaps over **d**ouble slotted flaps due to better turning effectiveness (see Figure 13). A comparison with experimental data from Reference (13), shown in Figure 27, shows that the predicted variation with flap angle is correct, but the differences due to double versus triple slotted flaps are no greater than the experimental scatter of the data. It is also observed that flap separation starts at flap angles above about 65 degrees for these tests at $C_{\rm LL} = 2$.

These results are in agreement with other data from Reference (6). There, single, double, and triple slotted flaps were investigated with a model having a straight wing. The various flap geometries are shown in Figure 28. The power-off lifts for these geometries are presented in Figure 29 for -= 0, showing a large improvement in lift in going from one gap to multi-gapped flap geometries. The large improvement above the single gap is probably due to a lessening of flow separation. Power effects from external blowing are presented in Figure 30 for the same flap geometries, using a blowing coefficient of $C_{II} = 1.0$. The single slotted flap shows the largest lift increment which probably results from ar elimination of the flow separation that might exist without the blowing. However, little difference is found between the double and triple slotted flap. Figure 31 shows the sum of the freestream lift and the increment due to power, i.e., the total lift. It is seen that the differences in total lift between double and triple slotted flaps are minor.

Also it should be noted that the comparison of the number of gaps is carried out at constant flap chord extensions. However, in case a simultaneous increase of flap chord is achieved together with a larger number of gaps, an increase of turning angle may be obtained because of a larger impingement parameter.

The prediction methods have no provision for flap gap variations. Test data from References (6) and (13) show that an optimum gap is 3 percent to 3.5 percent of wing chord (see Figure 32). It is assumed that any EBF design would utilize a near optimum gap in order to derive the maximum lift capability from the flap system, since there are no apparent reasons to use other gap limensions.



Figure 27. Effects of Flap Configuration on Lift Increment Due to Power





REF. G		
≪=0°		
	FLAP CHORD	Nº OF GAPS
v	SMALL	THREE
📀	SMALL	TWO
·- 0	SMALL	ONE
_ · ∆	LARGE	THREE
O	I ADGE	ANE

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Figure 29. Effect of Flap Configuration on Lift - Power Off









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Figure 32. Effect of Flap Gap on Optimum Lift

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3.11 EFFECTS OF LEADING EDGE DEVICES

The primary effect of the leading edge device is to increase C_{LMAX} by delaying leading edge stall to higher angles of attack. In application to an externally blown flap, test data indicate that the increment in C_{LMAX} due to the leading edge device is developed in the power-off condition, and no further increase in C_{LMAX} due to the leading edge device is realized for the power-on case.

In the present sethods, the effects of the leading edge device are estimated for the power-off case by DATCOM or other suitable methods. Power effects are then added as described in Sections 3.2 through 3.5.

Tests were made in Reference 13 of systematic variations of Krueger flap span from 100 percent span to 40 percent span. Figure 33 presents the lift increment due to flap blowing at zero incidence, and shows no effect of Krueger flap span. Figure 34 shows the effect of leading edge flap span on $C_{L_{MAX}}$ for the power-off and power-on cases. It is seen here that the estimated power effects on $C_{L_{MAX}}$, which assume no effect of the leading edge device, are in good agreement with the experimental data at all values of the leading edge flap span. The increment of $C_{L_{MAX}}$ due to power is shown in Figure 35 for two values of blowing coefficient.



Figure 33. Effect of Leading Edge Flap Span on Lift due to Power



REF. 13



Figure 34. Effect of Leading Edge Flap Span on $C_{L_{max}}$



Figure 35. Effect of Leading Edge Flap Span on ($\Delta C_{L_{max}}$) Due to Power



Figure 36. Alternate Nacelle Positions Tested in Reference 13

3.12 EFFECTS OF NACELLE LOCATION AND ENGINE ORIENTATION

The methods of estimation developed in this report account for the effects of nacelle location and engine orientation through the use of the "impingement ratio," Z_F/R , and the inclusion of direct force terms due to thrust axis incidence. Test results of a systematic variation of nacelle position (Figure 36) and nozzle deflection angle conducted in the wind tunnel tests of this program, substantiate these methods.

Test data are shown in Figure 37 of lift increment due to power versus engine nozzle angle for various nacelle locations. It is seen that the high nacelle location has better performance, and the aft nacelle location has poorer performance than the basic nacelle location. Also, there is no apparent benefit in spreading the engines spanwise any farther than the basic arrangement. These data show the very potent effects of deflecting the nozzle upward into the flaps. The nozzle angle effect is about the same for all nacelle locations.

These same data when shown as a function of the impingement ratio in Figure 38, tend to collapse about a single curve with the exception of the aft nacelle location data, which is about 10 percent below the other data. Static calibration data of the aft location were not taken, but it is believed that the aft location might produce a higher drag force which would result in a lower turning efficiency, thus explaining the reduced lifting effectiveness. It is concluded that the prediction methods adequately account for effects of nacelle location and thrust incidence, although they do not account for the penalty incurred by moving the nozzle closer to the flaps. Here again, it should be a design objective to locate the nozzle at or near the wing leading edge.






3.13 EFFECTS OF ASYMMETRIC THRUST

Experimental results of test simulating the loss of thrust on one engine of a four-engine airplane is shown on Figure 39. It is seen that the loss in lift due to one-fourth reduction in 7C is very nearly equal to one-fourth of the total lift. Since the powered lift increment is not linear with C_{μ} , the loss of lift due to an engine failure is not equivalent to the reduction in lift that would result from the reduction of C_{μ} by 25 percent uniformly on all engines.

The loss of lift due to an engine fialure with the thrust held constant on the remaining engines, can then be estimated by:

$$(\Delta C_L)_{\text{engine failure}} = \frac{1}{N} (\Delta C_L)_{\text{P.L. all engines}}$$
 3.26

where N is the total number of engines on the airplane.

The test data presented in Figure 39 show that the loss of an inboard engine results in slightly less lift loss than the loss of an outboard engine. This might be due to the outboard engine affecting a larger portion of the wing than the inboard engine. This effect is negligibly small.



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Section IV

DRAG

Defining drag as the total force in the wind direction, including engine thrust, the total drag coefficient is:

$$C_{\rm D} = C_{\rm Df} + C_{\rm Dj} - \gamma (C_{\rm LL} \cos (\theta + \alpha) + \Delta C_{\rm DR})$$
4.1

where CD_f is the power off minimum drag minus (CL at min. drag)²/TT A; CD_i is the induced drag due to lift (see Section 4.1); and ΔC_{DR} is ram drag.

The drag increment CD_f is estimated by conventional methods or obtained from unpowered wind tunnel model data, and is defined by:

$$C_{Df} = C_{D_{min}} - \frac{(C_{L} \otimes C_{D_{min}})^2}{TA}$$
4.2

The ram drag component is:

$$\Delta C_{DR} = \frac{V + V}{2} \qquad 4.3$$

where W_a = weight rate of flow of air through engine inlets, 1b/sec.,

obtained from engine data V = freestream velocity, fps

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4.1 DRAG DUE TO LIFT

The drag due to lift for a jet flap wing of finite aspect ratio has been derived in Reference 11 as:

$$C_{D_{i}} = C_{L}^{2} / (\pi A + 2 C_{\mu}),$$

where C_L is the total lift coefficient including the jet reaction term. Attempted correlations of this expression with wind tunnel test data of externally blown flap models have produced negative results. However, good correlations have been obtained by using the unpowered form of the induced drag expression, and removing the jet reaction force from the usual lift:

$${}^{C}D_{i} = \left[C_{L} - nC_{\mu} \sin(\theta + \alpha)\right]^{2} / \pi A \qquad 4.4$$

Experimental data representing various configurations and several different aspect ratios have been plotted in the form of

$$\left[C_{\rm L} - nC_{\mu}\sin(\theta + \alpha)\right]^2 \text{ versus } \left[C_{\rm D} + nC_{\mu}\cos(\theta + \alpha) - \Delta C_{\rm D_{\rm R}}\right]$$

in Figure 40. Ram drag and uncertainty in the values of n and θ (see Section 4.2) cause some dispersion of the data, but for one set of conditions held constant the slope is consistent with the assumed value of 1/A.

A summary of estimated induced drag compared to induced drag derived from test data by removing the drag at zero lift, ram drag, and thrust reaction, is shown in Figure 41. The agreement is reasonably good.



Figure 40a. Correlation of Induced Drag Data

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Figure 40b. Correlation of Induced Drag Data

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Figure 41. Summary of Induced Drag Data

4.2 THRUST REACTION FORCE

The thrust reaction term, $\eta C \mu c^{-(1)} + \infty$), is determined from static values of η and θ as discussed in Section II. In estimating the total drag force, this is the component most difficult to define accurately. Both η and θ are subject to variable details of the airplane and flap configuration, and methods are given only for estimating performance of a "well-designed" system. In attempting correlations with wind tunnel test data, good agreement depends on availability of good static calibration results for η and θ .

The deviation of experimental values of the jet turning efficiency from the design curves of Figure 11 are summarized for the various types of flaps and wing sweeps in Figure 42. Most of the experimental points fall within a $\Delta \eta$ of +.10. The sensitivity of estimated drag with the parameter η , $\Delta C_D / \Delta \eta$, varies from about 0.5 in the range of C_L 's of 4 to 5, to about 2.0 at high lift coefficients. So, the experimental scatter in η could result in deviations in the estimated drag coefficient of +.05

at nominal lift coefficients and +0.2 at high lift coefficients.



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Section V

PITCHINC MOMENT

A theoretical development of the pitching moment of a two-dimensional wing with a blown flap has been made by Spence in Reference (17). These methods assume uniform spanwise distributed blowing at the wing trailing edge or knee of the flap so they do not simulate the three dimensional externally blown flap sufficiently for direct application. However, the framework of the theoretical development can be used to guide an empirical approach to solution of the problem.

The method recommended here for estimation of the pitching moment characteristics was developed from consideration of pressure distribution data of externally blown flap tests from References (13) and (14). Typical chordwise and spanwise distributions from Reference (14) are reproduced here in Figures 43 and 44.

From these pressure and loading distribution data it has been concluded that the pitching moments due to an externally blown flap can be predicted by treating the three major components - the thrust reaction, circulation lift at zero incidence, and lift due to angle of attack individually. The pressure data indicate the following characteristics which have been applied in the present prediction methods.

1. Thrust Reaction Term

The lift component of the thrust reaction, $\gamma C_{\mathcal{U}} \sin \theta$, acts primarily on the flap with no noticeable carry-over forward of the flap. Spanwise, this force is centered about the thrust centerline, as illustrated in Figure 44.

2. Circulation Lift at Zero Incidence

In the chordwise direction this force is quite uniformly distributed, so the center of pressure would fall very close to mid-chord. This is in agreement with two-dimensional jet flap theory which predicts a .5c location of the center of pressure of the induced lift at zero incidence. The spanwise loading distribution of the circulation lift is also quite uniform even when the blowing is concentrated inboard, so the spanwise c.p. of this load can be assumed to be at the MAC of the wing. Data with partial span flaps show a carry-over on the panel outboard of the flaps that averages about half of the unit loading of the flapped portion of the wing.



Figure 43. Chordwise Pressure Distribution of an Externally Blown Flap



3. Lift Due to Angle of Attack

The additional load due to angle of attack has a conventional distribution, peaking near the leading edge, with a c.p. near the quarter chord. This loading is also uniformly distributed spanwise, so the MAC of the wing can be assumed to be the spanwise location of the c.p.

The tail-off pitching moment characteristics are estimated by adding the powered lift effects to the power off characteristics which are estimated by conventional methods, or are available from unpowered wind tunnel model tests. As discussed above, the total tail-off pitching moment is developed by summing the effects of the several moment inputs:

$$C_{m} = (C_{m})_{B} + (\Delta C_{m})_{R} + (\Delta C_{m})_{\Gamma} + (\Delta C_{m})_{\alpha} + (\Delta C_{m})_{RD} \qquad 5.1$$

where $(C_m)_B$ = basic power-off moment coefficient, variable with angle of attack, with flaps extended

- $(\Delta C_m)_p$ = moment coefficient due to thrust reaction
- $(\Delta C_m)_{\Gamma}$ = moment coefficient due to circulation lift at zero incidence
- $(\Delta C_m)_{\infty}$ = moment coefficient due to power effects on additional lift at angle of attack

 $(\Delta C_m)_{RD}$ = moment coefficient due to ram drag.

Each of these pitching moment increments will be discussed in the following sections.

5.1 THRUST REACTION MOMENT

By definition, the thrust reaction force, F_R , acts at an angle to the reference plane of $\theta = \tan^{-1} (F_R/F_A)$, where F_N is the normal force component of F_R , and F_A is the axial force component. At forward speed these force components are defined by the coefficients: $\gamma C_{\mathcal{I}\mathcal{I}}$, $\gamma C_{\mathcal{I}\mathcal{I}} \sin \theta$ and $\gamma C_{\mathcal{I}\mathcal{I}} \cos \theta$. Now, according to our analogy to the jet flap, this jet reaction force represents the momentum in the jet **sheet** leaving the trailing edge of the flap system at the angle θ . Therefore, the moment of this force can be found by extending the force vector from the trailing edge of the flap at angle θ , and determining its moment arm about the moment reference center. The spanwise location of the thrust reaction force, $\gamma \zeta \mu$, is taken to be at the engine centerline; or, in the case of a four-engine installation, midway between the two engines on one wing. Now to determine the length of the moment arm it is probably more convenient to work with the normal and axial thrust reaction components. By extending the vector, $\gamma \zeta \mu$. from the flap trailing edge at angle θ , the chordwise location of the intersection of this vector with the reference plane (the horizontal plane through the moment reference center) can be found (see figure 45). Then the pitching moment about the moment reference center due to the thrust reaction force is equal to the normal force component, $\gamma \zeta \mu$ sin θ , acting at the distance from the moment reference plane. The axial force component has no moment arm in this system. So, taking moments about the wing leading edge at the location of the engine centerline:

$$(\Delta C_m)_{R_{L,E_i}} = -\eta \zeta_u \sin \Theta \left(\frac{C_R}{C'}\right) \left(\frac{C'}{\overline{c}}\right)$$
 5.2

- where C_R = chordwise distance of intersection of thrust reaction vector and the reference plane from the wing leading edge
 - c' = length of wing chord at engine centerline
 - \overline{c} = MAC of wing.

These parameters are illustrated in Figure 45 for two possible locations of the moment reference center. This pitching moment can be transferred to the moment reference center by conventional methods:

$$(\Delta C_m)_{R_{REF.}} = -\eta \zeta_{u} \sin \Theta \left[\left(\frac{C_R}{c'} \right) \left(\frac{c'}{c} \right) \left(\frac{X_{REF} - X_{LE.C'}}{\overline{c}} \right) \right] 5.3$$





5.2 MOMENT DUE TO CIRCULATION LIFT AT ZERO INCIDENCE

As discussed above the circulation lift at zero incidence can be assumed to act at the mid chord of the MAC of the affected portion of the wing. An alternate approach is to define the center of pressure of the circulation lift due to power to be equal to the center of pressure of the power-off, zero incidence, flaps down lift. Rather extensive investigations of experimental results of various configurations have shown this correlation. This approach should give better accuracy, especially if unpowered wind tunnel model data are available. The pitching moment increment of the circulation lift at zero incidence about the moment reference center is then:

$$\left(\Delta C_{m}\right)_{R \in F} = \left(\Delta C_{L}\right)_{\Gamma} \frac{(C_{m})_{B}}{(C_{L})_{B}} 5.4$$

where $(C_m)_{B_O}$

- $= \text{ pitching moment coefficient about moment reference center,} \\ \text{power off, flaps down, } \ll = 0$
- $(C_L)_B$ = lift coefficient, power off, flaps down, $\ll = 0$
- $(\Delta C_L)_{\Gamma}$ = circulation lift at zero incidence as defined in Section 3.2.

5.3 PITCHING MOMENT DUE TO ANGLE OF ATTACK

The center of pressure of the additional lift due to angle of attack for a two-dimensional jet flap has been derived by Spence in Reference (10). His results can be approximated by the relation:

$$\frac{X_{c.p.}}{C} \approx .25 - .01 \eta Gm$$
5.5

so the center of pressure of the angle of attack term is very near the quarter chord. Experimental data confirm that the additional lift acts essentially at the quarter chord. The recommended procedure for estimating the incremental pitching moment due to angle of attack effects is, then:

$$(\Delta C_m)_{\alpha LE, } = - \left[(\Delta G_L)_{\alpha} - (\Delta C_L)_{\alpha PO} \right] \left[\frac{(\chi_{C,P})}{\bar{C}_F} \left(\frac{\bar{C}_F}{\bar{C}} \right) \right]$$
 5.6

where: $(\Delta C_L)_{\alpha}$ = lift increment due to angle of attack, defined in Section 3.3

$$(\Delta C_L)_{PO} = 1 \text{ ift increment due to angle of attack, power-off} \\ \approx 2 \pi F \propto (1+1/c) = 2\pi \propto (1+1/c) A \\ \frac{X_{CP}}{\overline{c_F}} = .25 - .01 \, \gamma \, C_{LL} \\ \overline{c_F} = MAC \text{ of affected wing area as defined in Section 3.1.}$$

This moment increment can be transferred to the moment reference center by conventional methods:

$$(\Delta C_m)_{\alpha \zeta} = -\left[(\Delta C_L)_{\alpha \zeta} - (\Delta C_L)_{\alpha \zeta} \right] \left[\frac{X_{CF}}{\overline{c}_F} \left(\frac{\overline{c}_F}{\overline{c}} \right) - \left(\frac{X_{REF} - X_{LE}\overline{c}}{\overline{c}} \right) \right] 5.7$$

5.4 RAM DRAG MOMENT

The pitching moment of the ram drag about the moment reference center is:

$$(\Delta C_m)_{RD} = -(\Delta C_D)_R \frac{\Delta L_R}{E} \qquad 5.8$$

where $(\Delta C_D)_{\mathbf{R}}$

 Δl_R

= ram drag, defined in Section 4.

= moment arm of engine inlet axis about moment reference center, positive when engine is below reference plane. In the stability axis system this dimension is variable with angle of attack.

5.5 TOTAL PITCHING MOMENT

The total pitching moment coefficient about the moment reference center is, then, the sum of the above increments:

$$C_{m_{REF}} = (C_{m})_{B} - \eta \zeta_{m} \sin \theta \left[\frac{C_{R}}{C'} \right] \frac{C'}{\overline{c}} - \left(\frac{X_{REF} - X_{L.E.C'}}{\overline{c}} \right) \right]$$

$$+ \left(\Delta C_{u} \right)_{R} \frac{(C_{m})_{B_{o}}}{(C_{L})_{B}} - \left[\left(\Delta C_{u} \right)_{A} - \left(\Delta C_{u} \right)_{R} \frac{\Delta L_{R}}{\overline{c}} \right]$$

$$- \left(\Delta C_{v} \right)_{R} \frac{\Delta L_{R}}{\overline{c}} \qquad 5.9$$

Correlations with experimental data using these methods have been made for several varied configurations with the results being shown in Figure 46. In general the agreement is very good, falling with + 10 percent. Those data showing the largest deviation are due to overestimation of the lift increment due to blowing. Again, this noticeable discrepancy in lift estimation is primarily due to the aspect ratio effects.



Figure 46. Correlation of Pitching Moment Coefficient Data

5.6 EFFECT OF LEADING EDGE DEVICES

The effect of wing leading edge blowing on the center of pressure location of the power effects at $\propto = 0$ was analyzed and found to be negligible, see Figure 47.

Also, the effect of Krueger flaps and leading edge droop was found to be negligible, see Figure 48.

PEF. 5

BASIC LEADING EDGE



Figure 47. Effect of Leading Edge Blowing on Center of Pressure Location of Power Effect



Figure 48. Effect of Leading Edge Configuration Variation on Center of Pressure Location of Power Effect

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Section VI

TAIL ENVIRONMENT EFFECTS

The tail environment characteristics are presented in this section as the downwash angle at zero angle of attack and the change of downwash angle with lift due to power. Next the downwash factor $(1 - de/d \ll)$ is discussed. In particular various configuration effects have been examined. A limited amount of data on the **dynamic** pressure environment at the airplane tail is presented herein.

6.1 DOWNWASH AT ZERO INCIDENCE

The method of Ross, Reference (18), for estimation of the downwash behind jet-flapped wings was evaluated for application to the externally blown flap. Although reasonable agreement of the theory with test results was found for small jet deflection angles, the theory significantly underestimated the downwash for larger jet deflection angles. This was contributed in Reference (12) to overestimation of the deflection of the jet wake due to jet angle, and due to roll-up of the vortex sheet, which is not accounted for in the theory. A further discrepancy in this theoretical treatment is the assumption of elliptical spanwise loading distribution which can be substantially violated with EBF configurations.

Since the linearized theory did not give satisfactory results in the region of interest, an empirical approach to providing means of estimating the downwash was taken. The limited amount of test data available for externally blown flaps were generalized and are presented in Figure 49 showing the relation of downwash angle at zero angle of attack with lift coefficient, tail height, and tail length. In this form the effects of flap deflection angle and blowing coefficient are reflected in the lift coefficient, so any independent effect of these variables on downwash angle is not defined. Other limitations of these data are that they are for the aspect ratio range of 7 to 8.

The aspect ratio effect on downwash angle, which might be expected to be a function of CL/A or $CLA/(TA + 2C_{cL})$ (from jet flap theory), could be incorporated in this type of data presentation if sufficient experimental data were available for development and substantiation of such a factor. Unfortunately, experimental data to investigate this relationship with aspect ratio were not available.

The faired curves of Figure 49 have been superimposed on test data from References (2), (4), and (13) in Figures 50 through 53 at the appropriate tail heights and tail lengths to illustrate the degree of experimental data variation.



Figure 49. Downwash Augle at Zero Incidence, Summary Curves



Due to Power



Figure S1. Effect of Tail Height and Engine Deflectors on Downwash Variation With Lift Due to Power at Zero Angle of Attack









6.2 DOWNWASH FACTOR

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Within the limitations of the parameters discussed above the downwash angle and the downwash factor $(1 - \partial \epsilon / \partial \alpha)$ can be estimated by the following procedure:

where $\boldsymbol{\epsilon}_{o}$ is the downwash angle at zero incidence, given by Figure 50.

To find $\partial e/\partial \kappa$ it was assumed that from a given baseline downwash angle corresponding to a certain lift coefficient, blowing coefficient and tail location, the downwash angle changes with angle of attack as C_L and tail height change with angle of attack. Expressing the change in downwash angle as:

$$\Delta \epsilon = \frac{\partial \epsilon}{\partial c_{L}} \Delta c_{L} + \frac{\partial \epsilon}{\partial z} \Delta z \qquad 6.2$$

and holding Gu constant and varying only angle of attack,

$$\Delta \epsilon = \frac{\partial \epsilon}{\partial \zeta} \Delta \zeta + \frac{\partial \epsilon}{\partial \zeta} \Delta z$$

Then for small changes in **e**(:

$$\frac{\partial \epsilon}{\partial \alpha} = \frac{\partial \epsilon}{\partial c_{L}} + \frac{\partial \epsilon}{\partial t} \frac{\partial t}{\partial \alpha} \qquad 6.3$$

where $\partial CL/\partial o L$ is the power-on lift curve slope as defined in Section 3.3, and $\partial L/\partial o L$ is equal to the tail length L/c (where Z is in units of chord length).

The slopes, $\partial \in / \partial C_{L}$ and $\partial \in / \partial \mathbb{R}$, have been derived from the plot of Figure 49 and shown versus C_{L} in Figure 54. It should be noted that these values are in degrees, whereas $\partial C_{L}/\partial \mathcal{A}$ and $\partial \mathcal{R}/\partial \mathcal{A}$ are per radian. Accounting for these differences in whits, the expression for $\partial \mathcal{E}/\partial \mathcal{A}$ is:

$$\frac{\partial \epsilon}{\partial x} = \frac{1}{57.3} \left[\frac{\partial \epsilon}{\partial c_{L}} C_{L} + \frac{\partial \epsilon}{\partial z} \frac{1}{2} \right] \qquad 6.4$$

A plot of the downwash factor calculated by this expression for a tail length of 3.5 and using typical values of $CL \propto$ is shown in Figure 55. Comparisons with experimental results from References (2), (4), (5), and (13) are shown in Figure 56. The average correlation is fairly good although some unaccounted configuration effects are apparent.

6.3 CONFIGURATION EFFECTS

The major configuration effect on downwash angle is the tail height as shown in the data of Figure 51. Tail length has a much lesser effect as seen in Figure 50. The data in Figure 52 indicate a small reduction in downwash angle when increasing the flap deflection from 50 degrees to 60 degrees. Also shown in this figure is the effect of rotating the engine exhaust nozzles from 15 degrees up to 15 degrees down; again, only a small reduction in downwash angle is observed. The effects of leading edge flap span and chord, and the effect of engine spacing is shown in Figure 53. The increased span leading edge flap with spread engines shows a small reduction in downwash angle. Unfortunately the individual effects of these two factors cannot be separated in the present data.

The most significant effects on the downwash factor $(1 - \partial \in /\partial \propto)$, other than tail height appear to be due to leading edge treatment. Figures 56(e) through (h) show rather large, unpredictable, effects of leading edge blowing, producing different effects with different locations of the outboard engines. The effects of leading edge flap span and chord on the downwash factor is illustrated in Figure 56(a). At the lower lift coefficients, reduction of leading edge flap span and chord appear to improve the downwash factor. At higher lift coefficients no noticeable effect is observed.

The effects of flap deflection angle, shown in Figure 56(c) indicate no observable effect on the downwash factor in increasing flap deflection angle from 50 to 60 degrees.

The erratic behavior of the data at the higher lift coefficients observed in Figures 56(h) and 56(i), is eviled by due to the spanwise location of the outboard engine. Both sets of data are from the same model with the engines in the "spread" configuration. Moving the outboard engine further inboard, or application of leading edge blowing, appears to improve the downwash factor to values compatible with the other configurations. Other data with spread engines from Reference (13), Figures 56(b), is quite comparable to that of Reference (5) with leading edge blowing, Figure 56(e).

The effects of varying tail length from L/c = 2.2 to 4.2 as shown in Figure 56(i), are negligible, which is in agreement with the methods of estimation.







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Figure 55. Sample Estimated Examinash Factor



Figure Só. Experimental Downwash Factor Data



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Figure S6. Experimental Dormwash Factor Data (Continued)






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The methods of estimation of the downwash angle and the downwash factor given here are believed to be adequate for typical EBF STOL configurations which utilize a leading edge device, large span flaps, and aspect ratio of 7 to 8. The most important configuration variable not accounted for here is the aspect ratio of the wing, or of the flapped portion of the wing. Additional data should be obtained of this variable and analyzed in the context of the approach used here for prediction methods.

6.4 TAIL DYNAMIC PRESSURE

The dynamic pressure ratio at the tail was obtained from Reference (2) and has been presented, verbatim, at the top of Figure 57. The dynamic pressure increase coefficient ($\Delta \frac{q}{T/S}$) may readily be determined from:

$$\frac{\Delta \varphi}{(T/5)} = \left(\frac{\varphi_{TAIL}}{\varphi} - 1\right) / C_{\mu}$$
6.5

The dynamic pressure increase coefficient in this case was a linear function of the gross thrust or blowing coefficient, C_{μ} .



Figure 57. Tail Dynamic Pressure Characteristics

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Section VII

ASYMMETRIC POWER EFFECTS ON LATERAL CHARACTERISTICS

7.1 ROLLING MOMENT DUE TO ENGINE FAILURE

The rolling moment due to engine failure is determined from the lift change due to the failure times a moment arm:

$$C_{LEF} = \Delta C_{LEF} \cdot \frac{1}{2} \left(\frac{\lambda}{b/2}\right)_{FF}$$
 7.1

Herein, the subscript EF denotes "engine failure", while Y denotes the lateral distance of the center of pressure of the forces due to engine failure.

The lateral distance can be taken as being equal to the lateral location of the centerline of the exhaust nozzle of the failed engine. Correlations of several test results, Figure 58, show some variation of this, both inboard and outboard; however, since the variation is reasonable over a large range of configurations it is recommended that the engine centerline be used as the c.p. of the failed engine forces. In the method presented hereafter the lateral c.p. location is assumed to be independent of α .

The lift change due to engine failure can be computed below the stall angle of attack from

$$\left(\Delta C_{L} \right)_{EF} = \left(\Delta C_{L} \right)_{NOP} - \left(\Delta C_{L} \right)_{OEO} = \frac{1}{N} \left(\Delta C_{L} \right)_{PL} 7.2$$

as defined in Section 3, where N is total number of engines and $(\Delta C_L)_{PL}$ is the lift increment due to powered lift with all engines operating.

The maximum change in lift at an angle of attack beyond stall can be substantially larger, as illustrated in Figure 59. The magnitude depends very much on the effect of power on the stall angle of attack, and probably also on the wing sweep angle. Data for 24-degree swept back wings and clustered engines are presented in Figure 60.

A comparison of test data with computations of the rolling moment due to engine failure, using the above described method, is given in Figures 61 and 62. It is seen that the method is useful for the prediction of the rolling moment versus angle of attack.

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Figure 58. Lateral Center of Pressure For Span Loading Due to Power With One Engine Not Operating

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Figure 59. Maximum Lift Loss Due to One Engine Out



Figure 60. Engine Failure Effect on Maximum Lift









7.2 YAWING MOMENT DUE TO ENGINE FAILURE

An attempt was made to use a similar technique to predict the yawing moment with asymmetric power by replacing the normal force, used for the rolling moment predictions, with the axial force or drag increase due to engine failure. The results show that the centroid of the outboard failed engine is near the engine centerline, but the centroid of the inboard failed engine is far outboard (see Figure 63). It is also observed that the center of pressure, or yaw centroid, moved outboard when the tail was added to the configuration; apparently the failure of an engine is inducing a strong crossflow which has a greater reaction on the fuselage in the case of an inboard engine failure, than for that of an outboard engine failure.

An investigation of the sidewash induced by the asymmetric vortex system resulting from a failed engine might yield a method for predicting the engine-out directional characteristics.





Section VIII

POWER EFFECTS ON AERODYNAMIC DERIVATIVES

8.1 LONGITUDINAL DERIVATIVES

8.1.1. CHANGE IN FORWARD VELOCITY DERIVATIVES

 $C_{X_{U}}$, $C_{Z_{U}}$, $C_{m_{1}}$ - These derivatives arise from three sources; Mach number effect, aeroelastic effects, and thrust or power effects. The first two of these sources are very small in the STOL flight region and may be neglected. The thrust or power effects, of concern in this section, however, can become quite important in some configurations and should be investigated in detail, as their effect is seen in both the period and damping of the phugoid oscillation.

In general, major contributions from power can be broken into three categories:

- 1. Direct effects
- 2. Inlet effects
- 3. Induced flow effects.

Both direct and inlet effects are readily handled once the engine characteristics are known through normal analysis. An example of a typical direct effect is, at low speeds when the net thrust varies greatly with velocity, noticeable $C_{m_{ll}}$ effects can occur if the thrust vector is highly offset from the center of gravity. The velocity-power effects from the inlet are usually negligible, as the resulting forces generated are basically a function of inflow angle, but power effects may arise here for other derivatives such as, $C_{m_{occ}}$, $C_{m_{ol}}$, C_{fp} , and C_{Mr} . In all cases caution should be used in dynamic analysis bookkeeping to avoid a double entry of power effects when thrust coefficient terms are included.

Induced flow effects are more difficult to assess. A successful prediction method, if test data on similar vehicles are available, is to plot a nondimensionalized contributor against the inverse of the blowing coefficient as illustrated in Figure 64. Since all the factors in the power-speed relationship, $1/C \mu p_E = qS/Tp_E$, are known and q is a function of velocity, a graphical differentiation when corrected by the appropriate constants will yield the desired change in forward velocity derivation due to induced flow power effects for that component. The total derivative is naturally the summation of all such contributing components.

8.1.2 CHANGE IN ANGLE OF ATTACK DERIVATIVES

 $C_{X_{\infty}}$, $C_{Z_{\infty}}$, $C_{m_{\infty}}$ - Power effects on these derivatives can be determined directly using prediction methods shown in Sections 3, 4, and 5. Of the contributing factors, $C_{L_{\infty}}$, $C_{D_{\infty}}$, and $C_{m_{\infty}}$, the first two normally exhibit a substantial positive increase, while the last $(C_{m_{\infty}})$ only a small, almost negligible positive amount.

8.1.3 RATE OF PITCH DERIVATIVES

 C_{Xq} - This derivative results basically from the increase in horizontal tail drag due to pitching about the center of gravity. Since it in itself is small and for first approximations set to zero, power effects are also assumed small and should be neglected.

 C_{Z_q} , C_{m_q} - The horizontal tail is again the main contributor to both of these derivatives. In STOL (low speed) flight it is basically the effect of the curved flight path which increases the horizontal's angle of attack producing a negative C_{Z_q} and C_{m_q} (pitch damping). Correlation of prediction methods and test data is difficult since results are always combined with linear acceleration (&) derivatives. It is reasonable to assume, however, that the theoretical horizontal tail contributions of;

$$\begin{bmatrix} C_{Z_{q}} \\ H \end{bmatrix} \stackrel{\simeq}{=} - \begin{bmatrix} \frac{\partial C_{L}}{\partial \left(\frac{4E}{2} \lor_{c}\right)} \\ \frac{\partial \left(\frac{4E}{2} \lor_{c}\right)}{\partial \left(\frac{2}{2} \lor_{c}\right)} \\ H \end{bmatrix} \stackrel{\simeq}{=} - \begin{bmatrix} C_{L_{q}} \\ H \end{bmatrix} \stackrel{=}{=} - \begin{bmatrix} C_{L_{q}} \\$$

still hold and estimations can thus be made knowing the variation of $CL_{\sigma C_{\mu}}$ or/and q_{μ}/q with power.

8.1.4 LINEAR ACCELERATION DERIVATIVES

 $C\chi_{\infty}^{*}$ - As in the rate of pitch case, resulting drag increases are usually neglected in this derivative since they are small in comparison with that of the total aircraft and power effects would thus be even a smaller contribution.

 CZ_{∞} , $C_{m_{\infty}}$ - Predictions similar to those in the rate of pitch derivatives can be made for the theoretical horizontal tail contribution. The potential for error in estimating these derivatives is increased due to the additional $\partial \mathcal{E}/\partial \infty$ variation with power factor

$$\begin{bmatrix} C_{\Xi_{\alpha}} \end{bmatrix}_{H} \stackrel{\cong}{=} - \begin{bmatrix} C_{L_{\varphi}} \end{bmatrix}_{H} \frac{\partial \mathcal{E}}{\partial \alpha} \qquad 8.3$$

8.1.5 LONGITUDINAL CONTROL DERIVATIVES

CX - This frequently neglected derivative is of little significance in longitudinal calculations. The power effects on it are of even less significance and should therefore also be neglected.

 C_{Z} ge, C_{m} ge - Power effects on these derivatives are readily determined knowing the variation of horizontal tail effectiveness with power through:

$$C_{z} \delta e = -C_{Loc_{H}} \frac{q_{H}}{q} \frac{S_{H}}{S} \frac{\partial \alpha C_{H}}{\partial \delta e} \qquad 8.5$$

$$C_{m} \delta e = \frac{le}{E} C_{z} \delta e \qquad 8.6$$

8.2 POWER EFFECT ON LATERAL-DIRECTIONAL DERIVATIVES

8.2.1 CHANGE IN SIDESLIP DERIVATIVES

Cyg - The major portion of the total derivative comes from both the vertical tail and fuselage, in an approximately even magnitude, for the vehicles investigated. As normally the case, an almost negligible contribution comes from the wing; flaps up or down. Power effects on this derivative are small. The increase (a larger negative value) is attributed to the increased flow generated by the engines and can be approximated by the velocity ratio at the vertical tail.

$$C_{Y_{\beta}} \approx (C_{Y_{\beta}})_{POWER} \frac{\frac{1}{2}}{\frac{1}{9}}_{eo}$$
 8.7

Large variations and even sign changes in this derivative can occur at low sideslip angles due to the vortex formed at the fuselage-wing-leading edge flap junction sweeping across the vertical tail. Caution should be exercised in interpreting test data with minimal points near the axis $(\beta \pm 4^{\circ})$ or other angles formed by extending similar discontinuities through the vertical tail.

 $C_{\ell,\beta}$ - With the main contribution of this derivative coming from the wing, variations due to power appear to be dependent on engine placement. Configurations with clustered engines or those causing a definite bulge in the wing lift distribution (with a trough across the fuselage) tend to produce negligible power effects, occasionally even decreasing the dihedral effect (positive $\Delta C_{\ell,\beta}$ pwr). However, vehicles with spread engines that tend to increase the lift distribution elliptically with power application exhibit the expected increase in dihedral effect. Standard handbook methods, such as DATCOM, which relate $(C_{\ell,\beta})_{WING}$ to lift coefficient, taking into account the proper sweep and aspect ratio factors, seem adequate for preliminary predictions.

 $C_{n\beta}$ - Contrary to non-STOL, unblown flap configurations where this derivative is primarily the balance between the large fuselage and vertical tail contributions, the wing is the predeminant factor with power application. The power effects on the vertical tail contribution may be neglected for first order approximations or increased by the q_t/q_{con} ratio if desired. Theoretical prediction methods for the power effects on the wing contribution to $C_{n\beta}$ need further investigation. Fair correlation has been obtained by modifying the expression from DATCOM for power-off $(C_{n,c})$ wing by the factor $(1+C_{\mu})^{1/2}$. The power-off term from DATCOM is:

$$\frac{(\Delta C n_{e})_{WING}}{C_{L}^{2}} = \frac{1}{57.3} \left[\frac{1}{4\pi A} - \frac{\pi n \Lambda}{\pi A (A - 4\cos\Lambda)} \left(\cos\Lambda - \frac{A}{2} - \frac{A^{2}}{8\cos\Lambda} + 6\frac{\overline{x}}{\overline{c}} \frac{\sin\Lambda}{A} \right) \right] 8.8$$

where \overline{x} is the distance from the moment reference (usually c.g.) to wing a.c., positive aft.

Another area which requires further investigation is the variation of yawing moment with asymmetric power. While only minor effects on $\mathcal{L}_{n,\mathcal{G}}$ have been noted, a large sidewash appears to exist with inboard engine failure due to change in the pressure field aft of the wing. This causes the inboard engine out case to be critical in yaw and may pose several design problems.

8.2.2 RATE OF ROLL DERIVATIVES

 Cy_p - This derivative is of very little significance in lateraldirectional dynamics and is frequently neglected is calculations. The small, garbled (sometimes negative, sometimes positive) power effect seen in available test data tends to corroborate data shown in Figure 7.1.2.1-3a of DATCOM. Assuming the major power effects are contributed by the wing, this figure indicates that for aspect ratios, sweeps, and taper ratios typical of STOL configurations the Cy_p/CL ratio is very close to zero and could very well be either positive or negative depending on the specific configuration.

 $C_{/P}$ - This derivative, quite important in lateral dynamics in roll damping, stems basically from the wing and is thus highly effected by power especially at high angles of attack as indicated by test data. While the theoretical analysis given by Thomas and Ross, in Reference (19), shows the damping-in-roll to increase linearly with C_j and be independent of angle of attack, a good analogy exists with the blown flap. It is suggested that the methods of this report be used for preliminary designs if similar configuration test data are unavailable.

 $C_{\rm Hp}$ - The major power effect contribution to this derivative arises from the wing. It is negative and directly proportional to the lift coefficient. The secondary contribution from the vertical tail may be positive or negative depending on geometry and angle of attack. However, since its isolated contribution is greatly altered by the complex sidewash produced by the rolling wing and it is only the small change in this sidewash which is affected by power, prediction of this secondary contribution does not seem warranted.

Power effects on this derivative are fairly important due to the influence on dutch roll damping. High power settings have been seen to double the negative value of this derivative which is not desired since it represents a reduction in dutch roll damping. With present day stability augmentation systems normally providing a C_{np} feedback, it becomes more important to predict the power effects for the determination of gains and closed loop analysis.

 C_{np} approximations can be obtained by first determining the power off derivative and then applying a $(C_L - n C_{\mu} \sin \theta)/C_{L,P,O}$ correction. The theoretical reversal shown in Reference (19) has not appeared in test data investigated.

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8.2.3 RATE OF YAW DERIVATIVES

 $C_{Y_{P}}$ - Power effects on this relatively unimportant derivative are small and probably arise from augmentation of the sidewash produced through rotation at the vertical tail. As it is common practice to neglect this derivative in calculations, it is assumed that these minor power effects are of even less importance and may thus also be neglected.

 $C_{\ell r}$ - Only very minor power effects have been noted in all test data surveyed, so it is concluded that power effects are negligible.

 C_{n_r} - In general terms, power effects improve this derivative (provide larger negative values), which is the main contributor to dutch roll damping and hence an important factor. However, various attempted prediction methods have yielded only mediocre results. Since the vertical tail is by far the largest contributor to the total derivative, wing-fuselage-leading edge flap vortex interference effects are probably playing an important role. It is suggested that if similar configuration test data are unavailable, the unpowered C_{n_r} predictions be increased by the dynamic pressure ratio, q_t/q_{∞} . The experimental data available show values of C_{n_r} at high power settings of 130 percent of the power-off value.

8.2.4 LATERAL ACCELERATION DERIVATIVES

 $CY_{\mathcal{B}}, C_{\mathcal{I}\mathcal{L}}, C_{\mathcal{B}\mathcal{B}}$ - These derivatives are extremely difficult to estimate or extract from test data regardless of power effects. If for some particular reason these derivatives cannot be neglected and a method is available for power-off estimation, then the power effects can be assumed proportional to the dynamic pressure ratio, $q_{\mathcal{L}}/q_{\infty}$, since these derivatives are a direct function of the sidewash time lag.

8.2.5 LATERAL-DIRECTIONAL CONTROL DERIVATIVES

It is necessary to evaluate each particular control system in its own realm. Since one engine-flap-aileron relationship will have a different lift distribution variation with power than another combination etc., a detailed discussion is not relevant. In general, nower effects on conventional surfaces, those not directly imminged upon by engine thrust, are small when compared to large forces produced by these surfaces through double hinges and other means. Neglecting the power effects thus seems reasonable and conservative in preliminary design work.





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Section IX

SAMPLE CALCULATIONS

A machine program, in Fortran IV, is given in the appendix for the calculation of lift, drag, and pitching moment, using the methods described in previous sections. The required input items and the output items are tabulated below.

INPUT	OUTPUT VS
	<u>C</u> <u><u><u><u></u></u><u>AND</u><u></u><u>AND</u><u></u><u>AND</u><u></u><u>A</u></u></u>
\mathcal{C}_{L}	CL CLMAX CD Cm (Δ CL) θ Δ CL/ Δ CC CDi (Δ CL) r (Δ CL) r (Δ Cm) RREF (Δ Cm) REF (Δ Cm) REF (Δ Cm) REF (Δ Cm) REF (Δ Cm) RD r_{C} r CMAX

The results of three test cases are presented in Figure 65 comparing estimated versus test data. One case is from the current NR-STAI test, Reference (13), of a model with A = 7, $\Lambda = 24^{\circ}$, with full span flaps. The other case, selected to illustrate a wide variation in configurations, is from Reference (14) of a model with A = 7, $\Lambda = 0^{\circ}$. and 75 percent flap span. The third illustrates triple slotted flaps.

A sample calculation using manual methods is given below, using the same test case from Reference (13) as used in the sample machine calculation case.

A listing of test cases used in development and substantiation of these methods, identifying the source of the data, is shown in the table following this section.

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SAMPLE CALCULATION

Geometric data: A = 7 25° = 24° Type of flap = double slotted Flap span = 100% wing span $S'/S = \lambda = .85$ t/c = .125 $\delta_{\rm T} = 18^{\circ}$ $\sigma_{\rm F} = 56^{\circ}$ $Z_{F}/R > 1.0$ **c** = 15.5 xREF - xLEc' = 6.6xREF = .25c $C_{R} = 14.6$ $x_{REF} - x^{LE} = 3.9$ CF = 19.6 AR = 4.7 Power-off wind tunnel data - or estimated power-off data: $(\alpha_{MAX})_{B} = 15^{\circ}$ $(C_L)_B = 2.25$ $(C_{LMAX})_B = 3.35$ **- 1**45 $CD_{f} = .145$ (ACD)_R = neglected for wind tunnel data comparisons $(C_m)_B = -.95 \text{ at } \propto = -5^\circ$ - .95 at $a = 0^{\circ}$ - .86 at $\propto = 10^{\circ}$ Calculations: $\theta/\delta_F = 1.0$ from Figure 6 $\theta = (\delta_F)(\theta/\delta_F) = 56^\circ$ $\eta = .76$ from Figure 11 2 .76 ncr 1.52 $\eta C'_{\mu} = \eta C_{\mu} (9/9')$.895 1.79 F, from Figure 13 at ηC_{μ} .738 .717 V, from Figure 18 at $\eta C'\mu$.962 .944 $(\partial C_{\rm L}/\partial \theta)_{\rm os}$, Figure 16 at $\eta C'_{\rm pc}$ 5.75 3.80 (OCL/Joc) . Figure 16 at nC'µ 10.05 8.11

3

2.28

2.68

.711

.937

7.32

11.55

Calculations - Continued

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	1	Сµ. 2	. 3
$(\Delta C_L)_{\theta} = F(1+t/c) \lambda (\partial C_L/\partial \theta)_{c_{\theta}} \sin \theta$	2.22	3.28	4.12
$C_{L} \propto F(1+t/c) \vee (\partial C_{L}/\partial c) \propto$	6.72	7.65	8.65
$(\Delta QL)_{\Gamma} = (\Delta CL)_{\theta} - \eta C\mu \sin \theta$	1.59	2.02	2.23
$C_{LO} = (C_L)_B + (\Delta C_L)_{\theta} - C_{\mu} \sin \delta_T$	4.16	4.91	5.44
$ \begin{pmatrix} (C_{LMAX}) = \frac{3}{4F} \begin{bmatrix} 1.15(\Delta C_{L}) \phi - (C_{L}) & (1-\beta) + (C_{LMAX}) \\ PL & 1 - \frac{3}{4F} & (1-\phi) \\ - C_{LL} \sin \delta_{T} \\ \Delta \alpha_{MAX} = \frac{(C_{LMAX})_{PL} - (C_{L})_{K=0}}{C_{LMAX}} $	5.87	7.16	8.32
$-\frac{(C_{LMAX})_{PO} - (C_L)_{PO, \alpha = 0}}{(C_{L\alpha})_{PO}}$	3.57	5.78	7.80
amax = (amax) po + damax	18.57	20.78	22.80
$(2C_L) \propto C_L \propto (\alpha C/57.3) \text{at } \alpha C = -5^{\circ}$	- - 59	67	70
$C_{L} = C_{L_{0}} + (\Delta C_{L})_{\alpha c} \qquad \text{at } \alpha c = -5^{\circ}$	1.17	1.34	1.51
$η_{C}^{C}\mu \sin(\theta + \alpha c) = -5^{\circ}$ $CL - η_{C}^{C}\mu \sin(\theta + \alpha c) = t - 5^{\circ}$ $CL - η_{C}^{C}\mu \sin(\theta + \alpha c) = t - 5^{\circ}$ 0° 0° 10° 10° 10° 10° 10°	4.16 5.33 .59 .63 .69 2.98 3.53 4.64	4.91 6.25 1.18 1.26 1.39 3.06 3.65 4.86	5.44 6.95 1.77 1.89 2.08 2.98 3.55 4.87
$\frac{1}{\pi A} = \frac{1}{\pi A} = \frac{1}$.405	.425	.405
$(AC_L) = 0, = \frac{A(1+t/c) 2\pi c}{(A+2)(57.5)}$.569 .980	.608 1.08	.569 1.08
$= -48 \text{ at } \text{ acc} = -5^{\circ}$ $= 0 \text{ at } \text{ acc} = 0$ $= .96 \text{ at } \text{ acc} = 10^{\circ}$ $= 0^{\circ} \text{ at } \text{ acc} = -5^{\circ}$ $= 0^{\circ} \text{ at } \text{ acc} = -5^{\circ}$ $= 0^{\circ} \text{ at } \text{ acc} = -5^{\circ}$ $= 0^{\circ} \text{ at } \text{ acc} = -5^{\circ}$ $= 0^{\circ} \text{ at } \text{ acc} = -5^{\circ} a$.479. .425 .309 .071 .289 .816	.906 .850 .618 335 095 .605	1.435 1.275 .926 ~.885 ~.561 .295

Calculations - Concluded	CIL			
$[C_R - (x_{REF} - x_{LEC})] / \overline{c} = .408$	1 1	2	3	
$(\Delta C_m)_{RREF} = -\eta C_{\mu} \sin \theta \left[C_R - (x_{REF} - x_{LEC}') \right] /2$	32	65	97	
$(\Delta C_m)_{\text{REF}} = -(\Delta C_L)_{\text{p}} [(C_m)_B/(C_L)_B]_{\alpha C} = 0$	67	85	94	
$x_{cp}/\bar{c}_{F} = .2501 \eta C \mu$.24	.23	.23	
$(x_{cp}/c_F)(c_F/c) - (x_{REF}-x_{LEcF})/c$.04	.05	.05	
$(\Delta C_m) \propto_{REF} = -\left[(\Delta C_L) \propto - (\Delta C_L) \propto P.O. \right]$				
$ \cdot \left[\left(\frac{x_{C,P}}{\overline{c}_F} \right) \left(\frac{\overline{c}_F}{\overline{c}} \right) - \frac{(x_{REF} - x_{LE}\overline{c}_F)}{\overline{c}} \right] \text{at } \alpha C = -5^{\circ} \\ 0^{\circ} \\ 10^{\circ} $.005 0 ~.008	.010 0 019	.011 0 027	
$(\Delta C_m)_{RD} = -(\Delta C_D)_R \Delta l_R/c$	0	0	0	
$C_{m} = (C_{m})_{B} + (\Delta C_{m})_{R} + (\Delta C_{m})_{P} + (\Delta C_{m})_{\infty}$ + $(\Delta C_{m})_{RD}$ $at \alpha = \begin{cases} -5^{\circ} \\ 0^{\circ} \\ 10^{\circ} \end{cases}$	-1.94 -1.94 -1.95	-2.44 -2.45 -2.47	-2.76 -2.77 -2.80	



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Figure 65 (a). Comparison of Estimated and Experimental Results 125

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Figure 65 (b). Comparison of Estimated and Experimental Results 126

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CASE	REF.	FIG.	$\Lambda_{\frac{c}{4}}$	A	\$ SPAN FLAPS	NO. FLAP SLOTS	NOMINAL FLAP ANGLE	NOZ. ANGLE OR DEFLEC.
la lb ld lf 2a 6u 9a 10a 11a 14a 21a 22a 23a 91a 61a 24a 25a 51a 71a 81a 51c 51d 101 102 103	-13 13 13 13 13 13 13 13 13 13 13 13 13 1	7 8 18 20 74/75 104 97/98 97/99 139/140 129 13 14 15 12a 8a 10 11 5b 8b 7 16 17 12 13 14	24 24 24 24 24 24 24 24 24 24 0 0 0 25 25 0 0 0 25 25 0 0 24 24 24 24 24 24 24 24 24 24 24 24	7 7 7 7 7 7 8 6.6 10 7 7 5.25 7.55 7.55 7.75 7.75 7.75 7.75	$ \begin{array}{r} 100\\ 100\\ 100\\ 100\\ 100\\ 100\\ 100\\ 100\\ 75\\ 100\\ 75\\ 100\\ 75\\ 100\\ 75\\ 100\\ 100\\ 100\\ 100\\ 100\\ 75 75 7 7 7 7 7 $	2 2 3 3 2 2 2 2 2 2 2 2 2 2 2 2 2 2 2 2	25/50 35/60 2.5/20/45 2.5/30/55 25/50 25/50 25/50 25/50 25/50 25/50 27.5/55 27.5/55 27.5/55 20/60 17.5/35 30/60 20/40 25/10/50 30/60 30/60 30/60 30/60	15° 15° 15° 15° 15° 15° 15° 15° 15° 15°

DEFINITION OF TEST CASES

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APPENDIX

FORTRAN IV FROGRAM FOR CALCULATING AERODYNAMIC CHARACTERISTICS OF EXTERNALLY BLOWN FLAP POWERED LIFT SYSTEMS

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นี้ได้มีการเหน่ามาตามีสำนักมากระบบไม่สนับการทำสามารถการเหลายาการเป็นการเป็น

INTRODUCTION

Aerodynamic characteristics of externally blown flaps are calculated by the program. The design compendium develops an analytical approach upon which the program is based.

The simple nature of the program means very little computer time is used. Fortran code was written to be easily understood and revised in support of analytical development.

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PROGRAM OPERATION

Data input consists of a set of six to twelve cards describing each case to be analyzed. After the cards are read in, the computations are performed followed by the printed and graphic output. After a case is finished the program attempts to process a new case terminating if there are no more input cards. Each case may contain as many as ten different angles of attack and ten different jet moment coefficients.

If graphic output is not desired, the main analytical routine will function by itself if the call subroutine "CRT CRT" is removed or the subroutine is replaced by a dummy routine.

ANALYTICAL METHODS

LIFT FORMULATIONS

Input data for lift

$$C_{\mu}$$
, λ OR S'/S, η , A, \pm/c , Θ , ∞ , $(C_{L})_{B}$, δT , $(C_{LMAX})_{B}$
Jet momentum coefficient based on effective wing area

$$C\mu' = \frac{C\mu}{\lambda}$$
(Eq. 1)

Two dimension gradient of lift with jet deflection angle

$$(C_{L} \Theta)_{\infty} = \left[4 \pi \eta C_{\mu} \left\{ 1 + \alpha_{151} (\eta C_{\mu})^{1/2} + 0.139 \eta C_{\mu} \right\} \right]^{1/2} (Eq. 2)$$

Two dimensional lift curve slope

$$(C_{L_{\infty}})_{\infty} = 2\pi \left[1 + 0.151 (n C_{\mu})^{1/2} + 0.219 n C_{\mu}^{\prime} \right]$$
 (Eq. 3)

Lift curve slope correction for partial span blowing

$$\mathcal{V} = \frac{5'}{5} + \left(1 - \frac{5'}{5}\right) \frac{2\pi}{(C \perp \infty)}$$
(Eq. 4)
Aspect ratio correction factor

n ractor

$$F = \frac{A + 2n C\mu}{A + 2 + 0.604 (n C\mu)^{1/2} + 0.876 n C\mu}$$
(Eq. 5)

Increment of lift coefficient due to circulation due to power

$$(\Delta C_{L})_{\Gamma} = \left[(1 + t/c) F \lambda (C_{L}_{\Theta})_{\infty} - \eta C_{\mu} \right] \sin \Theta^{(Eq. 6)}$$

Gradient of lift due to angle of attack

$$\frac{\Delta C_{L}}{\Delta \infty} = \left[F\left(1 + \frac{t}{c}\right) V\left(C_{L} \infty\right) \infty \right] \qquad (Eq. 7)$$

Increment of lift due to angle of attack, power on

$$(\Delta C_{L})_{\infty} = (\Delta C_{L} / \Delta \infty) \infty$$
 (Eq. 8)
Increment of lift due to jet turning

$$(\Delta C_{L})_{\Theta} = \left[(1 + t/c) F \lambda (C_{L\Theta})_{\Theta} - \eta C_{\mu} \right] \sin \Theta + \eta C_{\mu} \sin \Theta$$

$$= (\Delta C_{L})_{\mu} + \eta C_{\mu} \sin \Theta$$

$$(Eq. 9)$$

$$Total life coefficient$$

Total lift coefficient

$$C_{L} = (C_{L})_{B} + (\Delta C_{L})_{\Theta} + (\Delta C_{L})_{OC} - C_{\mu} \sin \delta T \quad (Eq. 10)$$

Maximum lift coefficient

$$(C_{L_{MAX}})_{PL} = \left\{ \frac{\int_{-\frac{1}{4F}}^{\frac{3}{4F}} \left[1.15 (\Delta C_{L})_{\theta} \phi^{-} (C_{L})_{PO, \alpha = 0} (1-\phi) \right] + (C_{L_{MAX}})_{B}}{1 - \frac{3}{4F} (1-\phi)} \right\} - C \mu \sin \delta_{T}$$

$$\phi = (C_{L\alpha})_{PO} / C_{L\alpha}$$
(Eq. 11)

Maximum angle of attack

$$\alpha_{MAX} = (\alpha_{MAX})_{B} + \frac{(C_{L_{MAX}})_{PL} - (C_{L})_{\alpha=0}}{C_{L_{\alpha}}} - \frac{(C_{L_{MAX}})_{B} - (C_{L})_{B,\alpha=0}}{(C_{L_{\alpha}})_{PO}}$$
(Eq. 12)

DRAG FORMULATIONS

Input data for drag

$$\pi$$
, C_{u} , θ , α , A , $(\Delta C_{D})_{R}$, $C_{D_{f}}$

 $C_{\rm L}$ from lift calculations

Coefficient of drag due to lift

$$C_{D_{i}} = \frac{\left[C_{L} - \eta C_{\mu} \sin \left(\Theta + \alpha\right)\right]^{2}}{\pi A} \qquad (Eq. 13)$$

Total drag coefficient

$$C_{D} = C_{D_{f}} + C_{D_{i}} - n C_{\mu} \cos(\Theta + \alpha C) + (\Delta C_{D})_{R} \quad (Eq. 14)$$

PITCHING MOMENT FORMULATIONS

Input data for pitching moment

$$n_{j}C_{\mu,j}\infty_{j}t/c_{j}X_{C,P,j}E_{F,C,j}(X_{REF}-X_{L,E,E_{F}})_{j}(C_{DR,j})$$

 $(C_{m})_{g,\Theta}, C_{R,j}(X_{REF}-X_{L,E,c})$
 $V_{j}(C_{L}\infty)_{\infty}, (\Delta C_{L})_{F,j}$ from lift calculations

Center of pressure of additional load due to angle of attack

$$X_{c,P} = (25 - 01 \eta C_P) \overline{c}_P \qquad (Eq. 15)$$

Increment of lift due to angle of attack due to power

$$(\Delta C_{L})_{\alpha, \beta 0} = \alpha c F(1+t/c) V [(C_{L}, \alpha)_{\infty} - 2\pi]$$
 (Eq. 16)

Pitching moment coefficient of power effects on lift due to angle of attack about the moment reference center

$$(\Delta C_{m})_{\alpha}_{RET} = -(\Delta C_{L})_{\alpha}_{PO} \left[\left(\frac{X_{C,P}}{\overline{c}_{F}} \right) \left(\frac{\overline{c}_{F}}{\overline{c}} \right) - \left(\frac{X_{REF} - X_{L,E}, \overline{c}_{F}}{\overline{c}} \right) \right]$$

Pitching moment coefficient of ram drag about the moment reference center

$$(\Delta C_m)_{RD} = -(\Delta C_D)_R \frac{\Delta L_R}{\overline{z}}$$
 (Eq. 18)

Pitching moment coefficient due to circulation lift due to power about the moment reference center

$$\left(\Delta C_{m} \right)_{\Gamma_{REF}} = \left(\Delta C_{L} \right)_{\Gamma} \left[\frac{(C_{m})_{B}}{(C_{L})_{B}} \right]_{\infty} = 0$$

$$(Eq. 19)$$

where $(G_m)_B$ is about moment reference.

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Pitching moment coefficient due to jet reaction about moment reference center

$$\left(\Delta Cm\right)_{RREF} = -\eta C_{\mu} \sin \Theta \left[\left(\frac{C_{R}}{C'}\right) \left(\frac{c'}{\overline{c}}\right) - \left(\frac{x_{REF} - x_{L.E.C'}}{\overline{c}}\right) \right] (Eq. 20)$$

Total pivching moment coefficient

$$C_{m} = (C_{m})_{B} + (\Delta C_{m})_{R} + (\Delta C_{m})_{T} + (\Delta C_{m})_{C} + (\Delta C_{m})_{RD} (Eq. 21)$$

NOMENCLATURE

SYMBOL	FORTRAN NAME	DESCRIPTION
A	A	Aspect Ratio
cc ₀	ALPHA	Angle of Attack, Degrees
∞C	ALPHAR	Angle of Attack, Radians
T	CBAR	Mean Aerodynamic Chord of Wing
Ĉŗ	CBARF	Mean Aerodynamic Chord of Wing with Flaps Extended
Ср	CD	Total Drag Coefficient, Power On
$C_{D_{f}}$	CDF	Minimum Drag Coefficient Extrapolated to Zero Lift
CDi	CDI	Coefficient of Drag Due to Lift
CL	CL	Total Lift Coefficient
$(\tilde{C}_{L,r})_{r}$	CLAI	Two Dimensional lift curve slope
$(C_{\rm I})_{\rm D}$	CLB	Total Lift Coefficient, Power Off
	CLMAX	Maximum Lift Coefficient
(CLARK) D	CLMAXB	Maximum Lift Coefficient, Power Off
$(CL_{\theta}) \infty$	CLTI	Two Dimensional Gradient of Lift with Jet Deflection Angle
Cm	CM	Total Pitching Moment Coefficient
(Üm) B	CMB	Total Pitching Moment Coefficient, Power Off
$(C_m)_B \propto = 0$	CMBAO	Total Pitching Moment Coefficient, Power Off Zero Angle of Attack
С и	CMU	Jet Momentum Coefficient, or Thrust Coefficient
C	CMUP	Jet Momentum Coefficient Based on Effective Wing Area
C _R	CR	Chordwise Distance from Intersection of Jet Reaction Vector with Reference Plane to Leading Edge of Chord at Thrust Centerline
$(\Delta C_D)_R$	DCDR	Ram Drag Coefficient
(ΔC_{I}) or	DCLA	Increment of Lift Due to Angle of Attack
DCL/DOC	DCLDA	Gradient of Lift Coefficient with Angle of Attack
(ACL)	DCLG	Increment of Lift Coefficient due to Circula- tion Due to Power
D CIMAX	DCLMAX	Increment in Maximum Lift Coefficient due to Power
(ΔCL) _θ	DCI.T	Increment of Lift Coefficient due to Jet Turning
$(\Delta C_m) \circ C_{REF}$	DOMA	Pitching Moment Coefficient of Power Effects on Lift Due to Angle of Attack About the Moment Reference Center
(acm)rref	DOMG	Pitching Moment Coefficient due to Circulation Lift Due to Power About the Moment Reference Center
(ACm) RREF	DCMR	Pitching Moment Coefficient due to Jet Reaction About Moment Reference Center
NOMENCLATURE

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SYMBOL	FORTRAN NAME	DESCRIPTION
$(\Delta C_m)_{RD}$	DCMRD	Pitching Moment Coefficient of Ram Drag About the Moment Reference Center
(∆CL)≪₽O	DDCLPO	Increment of Lift due to Angle of Attack due to Power
ΔLR	DLR	Moment Arm of Ram Drag about Moment Reference Center
Δ L _T	DLT	Moment Arm of Thrust about Moment Reference Center Positive if Nozzle Axis is below Moment Reference Center
TD	DT	Thrust Incidence Angle, Degrees
b _T	DTR	Thrust Incidence Angle, Radians
ΔXF	DXF	Chordwise Distance from Moment Reference Center to Flap Mid-Chord
∆x _G	DXG	Chordwise Distance from Moment Reference Center to Leading Edge of Mean Aerodynamic Chord of Flapped Wing Area
∆x _R	DXR	Chordwise Distance from Moment Reference Center to Leading Edge of Chord on Thrust Centerline
n	ETA	Jet Turning Efficiency
F	F	Aspect Ratio Correction Factor
入or S∕S	LAMDA	Area Ratio, Flapped Wing Area to Total Wing Area
	NA	Number of Alpha Values
	NC	Number of CMU Values
P	NU	Lift Curve Slope Correction for Partial Span Blowing
π	PI	Ratio of Circumference of a Circle to its Diameter
t/c	TC	Airfoil Thickness Ratio
θ	THETA	Jet Turning Angle, Degrees
0	THETAR	Jet Turning Angle, Radians
	TITLE	Title of Case
Х _{ср}	XCP	Center of Pressure of Additional Load due to Angle of Attack

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CASE 1A









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EQUIPMENT REQUIREMENTS

HARDWARE

The program has been run on an IBM System 370 model 165KJ computer with graphic output processed on an Information International Incorporated FR-80 simulating a Stromberg-Datagraphics S-C 4020. Runs have also been made on a Control Data Corporation Model 6600 without graphic output.

SOFTWARE

A fortran IV compiler is required to translate the source code into object form. Standard Fortran library routines are required for mathematical operations, card reading and printing output.

Graphics output requires several NR library routines that generate instructions for the S-C 4020. Some reprogramming may be necessary for graphic output at another computer facility.

FLOW CHART OF PROGRAM FUNCTIONS



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	J	ETA	JET TURNING E	FFICIENCY			0090
	ں	ETACMP	PRODUCT OF EI	TA AND CMUP			0610
· · · · · · · · · · · · · · · · · · ·	ں ر	ETACMU(I)	PRODUCT OF ET	TA AND CMU()			0620
	J	u.	FACTOR				0630
	J	FFF(10)	TEMPOKARY STU	CRAGE FOR F VAL	UES	-	
	J		INDEX ASSOCIA	ATED WITH CMU(0650
	ں ا	ر	INDEX ASSGCIA	ATED WITH ALPHA	()		0660
		X	SUBSCRIPT_ASS	SOCIATED WITH Z	ERO CMU		0670
	J		SUBSCRIPT ASS	SOCIATED WITH Z	ERO ALPHA		0680
	ں	LAMDA	AKEA RATIC, F	LAPPED WING AR	EA TO TOTAL WING A	REA	0690
	ပ ပ	NA	NUMBER OF ALF	PHA VALUES			0020
	ა	NC	NUMBER OF CHL	J VALUES			0110
	υ	NU	LIFT CURVE SI	OPE CORRECTION	FOR PARTIAL SPAN	BLOWING	0120
	C	- T- T-	RATIO UF CIRC	CUMFERENCE OF A	CIRCLE TO ITS DIA	METER	0130
	J	1 C	AIRFOIL THICH	(NESS RATIO			040
	J	THETA	JET TURNING A	ANGLE, DEGREES			0150
	ა	THETAR	JET TURNING A	ANGLE, RADIANS			0160
	ر	TITLE(18)	TITLE CF CASE	μ,			0170
	ن ر						0780
0001		COMMON	ALPHA (10)	. AMAX(10)	, CD(10,10) , CDI	(10, 10)	0800

FORTRAN	IV 6 LEVEL	źu	6.	NIA	DATE = 73110	00/26/45	
0002		COMMUN	(1(10-10)	1.1.1.0.X.4.1.J.			
0003		COMPON	CMUSTER			• CMB (10)	0810
4000		NUMMOU			VULA IVIU	• DCLDA(10)	0820
0005		COFACN	DCMK(12)	PUCEI (10)	FLCMA(10,10)	• DCMG(10)	0830
	J					• 17 [[E[[8]	0840
0000		DIMENSION	, FFF(1U)				0850
0007		REAL LAMD	NA + NU				0860
8000		DATA PI/	3.14155265 /				08 70
6000	60	READ(5,65) TITLÉ				0880
0010	65 65	FORMAT(18	A4)	• • • • •			0890
0011		KEAL (5,7	U)CLB.E1A.1H	ETA.TC.ET.A.LAM	DA "CRF. DYP. CP. C		0060
		L CEARF, CL	T, LLR, CLMAX5	, AMAX6		19V11V81	0160
2100	26	FORMAT (c	F12.0)				0000
0613	i	KEAL (5.7	5) NA, (ALPHA	(J), CMB(J), J=1	(NA)		0560
41.00	35	FURMAT (I	14,5512.0,30	/6F12.C))			
5100		KEAL (5°75	I NC . (CMU(I)	,DCER(I),I=1,NC			0000
9100		NKITE(6,2	99) IITLE				
2100	565	FURMAT(1H	1,1844)				0160
0. 18		WR17E (6,3	50)				0840
0010	350	FURMAT (32)	FUINPUT DATA	(PROGRAM VERSIC			0660
0.20		WRITE (6, 3	51)CLB				0001
0051		WRITE (6, 3	52}ETA	:		the bar and an	1010
0.022		WRITE(6,3	53)THETA				0201
0423		HRITE (0,3	54)1C				1020
00 24		MKI12 (6,3)	55107			-	
0025		HALTE (A 31	551A				1050
. 6326		WKITE (0,3)	57) LAMEA				10001
0027		WRITE(6,3!	58) C D F				0/01
0028		WRITE(6,39	59)DXR				1080
0629		WR1TE(6,30	EG)CR				0601
0030		WRITE(6,30	EIJCBAR				
0031		WR1TE(6,36	52)CXC				0111
. 80 32		#RITE(6,36	3)CEARF				1120
6000	-	HRITE(6,36	54) 5LT	•			1130
CO 34		WKITE(6,36	5)0LP				1140
00 35	~	KITE(6,36	2) CLMA XB				1150
0036	-	MRITE (5,36	S JAMAXB				1160
0037	-	#kITE(6,37	(C) (ALPHA(J)	() () () () () () () () () () () () () (0/11
. 6036	- - - - -	HLTE(6,37	TILCME (J).	(= *NA)			0811
0039	-	4RITE(6,37	2)(CMU (1),1	= I • NC)			1140
			•	• } > • • •			1200

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FORTRAN IV	6 ,LEVEL	20	MAIN	DATE = 73110	00/26/45	
0040		WRITE(6.373)(DCD	R(1).1=1.NC)			
00+1	351	FURMATICHOCLB	• F14.4)			100
0042	352	FORMATIOH ETA	• [] 4 • 4)			4 6
0043	555	FORMAT(8H THETA	•F14.4)	and a second deal of the second s		101
4400	354	FORMAT (BH TC	•F14.4)			
0045	325	FORMATISH DT	*F14.4)			24
0046	356	FORMATICH A	• F 14 • 4)			0 1 Vi Vi Vi Vi Vi Vi Vi Vi Vi Vi Vi Vi Vi
2470	357	FORMATISH LAMDA	·F14.4)			174
0048	358	- EURMAT (BH CDF	· F 14 • 4)			2001
9449	355	FORMAT (6H DXR	•F14 •41	raam oo ah		671
0020	360	FDRMAT(8H CR	· F 14 - 4)) ; ; ; ;
1400	361	FURMAT (BH CBAR	· · · · · · · · · · · · · · · · · · ·			
0052	362	FORMATION DXG	· · · · · · · · · · · · · · · · · · ·		•	2) (2) (2) (
0053	363	FCRMAT (8H CBARF	F14 4)			1) (1) (1) (
0054	364	EORMAT (Bh DLT				47 I
0055	365	FORMAT (BH DLK			ووجوانات والمحادثان المحاود والمحاولة والمحادة والمحادية والمحادية والمحادية	135
00 56	367	FCRMAT(6H CLMAX8	• F 14 • 4)			000
0057	308	FORMAT (8H AMAX8	• F 14 • 4)			
0058	370	FORMAT (SHOALPHA	• 7F14.4)			2 7 7
0059	371	FORMATIBH CMB	• 7F14 • 4)			500
0000	372	FORMATICHOCMU	•7F14.4)			5
0061	373	FORMAT (8H DCDR	+7F14.4)	and a second		
	ა	FIND ZERC	VALUE OF ALPHA			
0062		DO BCU J=1.NA				0 .
0063		IF (ALPHA (J) 800.	b 01.800			+ . + .
0064	801	CMBAD CMB(J)				1 1 1
00.65		ل = ا				4 i 4 i 1 i
0.766		GD TD 802				141
0067	600	CONTINUE				
UOAB						51
0069		CMBAD = CMB(1)				150
0070		WRITE(6.605)ALPH				151
1200	B 05	FCRMAT (45H # # #	A RADNING A A			152
	1	24H USING CHE H	DK ALPHA = .F14	+++ NU LENU ALKHA VALUE		153
0072	e 02	CONTINUE				
	، ت	FIND. ZERO.	VALUE OF CMU	:		
0073		DO 85G I = $1,NC$				
0074		IF(CMU(I))650,85	1,850			
57.5	851	K = I				1590
00.00	~	66 10 852				1600
						; } }

00/26/45 CATE = 73110

FORTRAN IV	G LÉVEL	ر ن	MAIN	UATE = 73110	00/26/45	
2 2 44	85C					1619
82.8						1620
	いられ					1630
	₽ \)	[[]] [] =] = NA			•	1640
00.61						1650
•	U					1660
2082	ŧ	CTK = DT + PI / 186	0.0			1670
0083		THETAR = THETA * PI	I / 180.6			1680
0084	:	ALPHAR = ALPHA(J) *	* PI / 180.0			1690
	د	LIFT SECTION				1700
0085		ETACMU(I) = ETA * C	(MU(I)			1710
30 80		CMUP = CMU(I) / LAM	IDA			1720
0087 0080		HIACMP # FIA # CMUP		151 + 17+/40++0 F +		1740
0008		ULII = 14+0 + FIF 1.330 + FTAC	6-8057 + 14+0 + 0+4 186114+0.5			1750
0089			1.0 + 0.151 * ETACM	4P**0.5 + 0.219 * ET	ACMP)	1760
0600		$hu = LAMEA + (1 \cdot C - C)$	- LAMDA] * 2.0 * PI	/ CLAI		1776
1600		F = (A + 2.0) + ETAC	MP / PI) / (A + 2.0	0 + 0.604 * ETACMP**	0.5 +	1780
1	-	1 0.876 * ETACMP)				06L1
0092		$FFF(\bar{i}) = F$				1800
0.93		0CLG(1) = (11,0 + T)	IC. * F * LAMDA * C	CLTI - ETACMU(I)*		1810
	~	L SIN(THETA	(K)			1820
7600		DCLDA(I) = F + (1, C)	> + TC) + NU + CLAI			1830
0095		vcla(1, J) = bclba(I)) * ALPHAR		i	
0096		DCLT(1) = DCLG(1) +	FETACMU(1) * SIN(TH	HETAR)		1850
2600	-	Cr(I+1) = Cre + nCr	_T(I) + UCLA(I,J)	CHU(I) * SIN(DTR)		1860
			and the second s			1991
0000	د		INIS * CHUCH * SINI	THETAR + A! PHAR]] **	A114 C	1896
00099		Cb(1, J) = CDF + CDI	(1.J) - ETACMU(1) *	COS(THETAR + ALPHA	R) +	1900
6 6 9 9	-	COR(I)	•			1910
	J	PITCHING MOMENT SE	CTION			1920
0100		XCP.= C6ARF .* 10.22	C.CL *. EIACMU(I)			1930
1010		DUCLPO = ALPHAR * (<pre>IF * (1.6 + TC) * NU</pre>	J * {CLAI - 2.0 * PI		1940
0102		OCMA(I,J) = -(DDCLP)	0) * ((XCP - DXG) /	CBAR)		1950
E010		DCMKD(1) = -DCDR(1)	H * DLR / CBAR		•	1960
0104		p(I) = p(I) + p(I)	CHEAU / CLB'			1970
0165		CCMR(I) = -ETACMU(I	() * SIN(THETAR) * (ICR / CBAR - DXR / C	BAR)	1980
01 (10		CH(1, J) = CMB(J) + .	DCMK(1) + DCMG(1) +	TUNNIA + TUNNIA	(T	1770
0107	200	CONTINUE				2000

2170 2180 2120 2140 2150 2150 2150 2200 2240 2250 2250 2280 2280 2300 2340 2380 2040 2070 2090 2110 2190 2230 2320 2330 2350 2360 2370 2390 2400 2020 2030 2010 2060 2080 2220 00/26/45 ÷ * 1.15 * DCLDA(K) / DCLDA(I) - CLB * (1.0-0CLDA(K) / DCLDA(I)) + CLMAX8) / (1.0 - 0.75/F * (1.0-0CLDA(K)/DCLDA(I)))-CMU(I) 1 ÷ AMAXIII = AMAX5 + (ICLMAXPUI) - CLUI+1) / DCLDAII) (CLMAX8 - CLB) / UCLDA(K)) * 180.0 / PI DATE = 73110((0.75 / F + (DCLT(I)))(I, J), I=1, NC) WRITE(6,301)(UCLA (I,J),I=1,NC) WRITE(6,308)(DCLG(1) ,I=1,NC) =1 *NC > (1, J), I=1, NC)), I=1,NC) (I + 7) + I=1 +NC) (I, J), I=1, NC) ,1=1,NC) • I = I • NC (I,J),I=1,NC • I=1,NC WRITE(6,500)(CMU(1), I = 1,NC) FOKMAT (8HOCMU ,7F14.4) • +7F14.4) FORMAT (THU ALPHA, F14.4) •7F14.4) • 7F14.41 * 7F14 .41 * 7F14.4 2F14.4 *7F14.4 +7F14.4 • 7F14.4 MAIN + 7F14.4 7514.4 *7F14.4 •7F14.4 WKITE(6,302)(ETACMU(1) * SIN(UTR) WRITE(6,311)(UCMRD (I WRITE (6,400) ALPHA(J) WEITE(6,315)(DCMK(1) WKITE(6.310)(DCMG(1) WRITE(6,312)(UCMA (WRITE(6.304)(UCLT(1) WRITE(6,300)(DCLDA (WRITE(6,299) TITLE ETACMU #KITE(5,305)(CUI #KITE(0,307)(CC WRITE(6,309)(CM HRITE(6,306)(CL FORMAT (8H DCLUA FORMAT.8H CDI DCMRD = JANC CCLG FCKMATINH ECLT FORMATISH CCMG FURMATION DCLA ECRMATION DCMR DCMA FORMAT (6H CC FORMAT (BH CM FCRMAT(BH CL u FURMAT (6H FORMAT (8H FORMAT (8H FCKMAT BH (I) dXVW (I) CONTINUE 00 700 S S ų, 300 305 301 2001 2001 100 308 315 400 sco 300 512 3 0 2 FORTRAN IV G LEVEL 6113 01 28 0137 0110 0115 0110 0118 0121 0127 5210 0130 GL31 0132 0133 0135 0135 0134 01-10 142 0143 0106010 0117 **U12**3 0123 **ul**24 01 25 6120 ÷ 110 0136 0113 0114 0122 6141 0111 0112 0

2410 2420 2420 2440 2440 2440 2440 2450 245		
00/26/45		· · · ·
DATE = 73110		i
0 MAIN WIINUE ITE(6,330)(CLMAXP(I),I=1,NC) ATE(6,331)(AMAX (I),I=1,NC) AMAT(6HOCLMAXP ,7F14.4) AMAT(6H AMAX ,7F14.4) AMAT(6H AMAX ,7F14.4) L CRTCRT(NA,NC) TL 60		
6 LEVEL 5 330 NEVEL 5 330 NEVEL 5 330 NEVEL 5 30 NEVEL 5 00 N		
FORTRAN IV 0144 0145 0145 01446 01448 01448 0150 0150		

FORTRAN IV G LEN	/EL 20	Ū	RTCRT	DATE = 73110	00/26/45	
0001	SUBROUT	INE CRICETINA.	1.57	•		1
0002	こともこし					2490
0003				• CD(10,10)	. CDI (10, 10)	2500
0004	CANON	Current of		•CM(10.10)	• CMB (10)	2510
CO.05				.DCLA(10.10)	• DCLDA(10)	2520
Children of the second s			*DUL11101	,DCMA(10,10)	• DCMG (10)	2530
2000			.DCMRD(10)	*ETACMU(10)	*TITLE(18)	2540
	LAILKNAL	- IABLIV, TABL	2V, TABL3V			
0000	CALL CHS	12V(2,2)				0002
0022	ISPACE =	. 12				2560
0010	CALL CAP	RAV(35)				2570
1100	CALL FRA					2580
0012	CALL RIT	E2V(100-1010-1	1.73.00.1.73.1	1111 C 11 C 1 C 1		2590
U	a	GT DCI TH VS ET		1 2 1 CE JNEAS 1)		2609
0013	CALL SET					2610
0014	CALL GRI					2620
0015			CALCE VALATION LA			2630
0010		CT10004071100				2640
100			PACE)			2650
0018						2660
0019						2670
0020						2680
(40.2.1		MIXIA Y I V				2690
ب د 1						2700
ر کوب		CT DCLDA VS ET	ACMU			2710
7775	CALL SET	MIV(600, 80,63	0, 40)			0110
	LALL GAL	D2V(-2. 0., 5.	· 01011.	1.1.1.1.1.21		1212
4 City 4	CALL NCU	CRT(770,595,1S	PACE)			2130
6226	CALL CLA	LBI (570.800.15	PACEN			2740
0.026	DC 210 I	= 1 • NC				2750
0027	X = ETACI	HU(1)				2760
0.02.6	X = DCLDI	4(I)				2770
0029	CALL POIN	41 V (X . Y . 1)				2780
0030 21	ID CONTINUE	•				2796
J	PLC	DI DCLMAX VS F	TAC MIS			2800
0431	CALL SET	11V(180.500.20	0.4501			2810
0032	CALL GRIT					2820
	CALL CLMC	RT(120-400-75)		4,4,4,4,4,4,4,2)		2830
0 434	CALL NCUC	RT1350 180 TC				2840
0.035	00 220 1					2850
0130						2860
0.37		P111				2870
						2880

3050 3130 3210 3280 3000 3010 3020 3020 3040 3060 3070 3080 3090 3100 3110 3120 3146 3150 3170 3180 3220 3230 0540 2910 2930 2940 2950 2960 2990 32.00 3250 2890 2920 2980 3160 3190 3260 2900 00/26/45 : ; ł ł -CALL GRI12V1-2.-10..30..0.12..10.11.1.2.-11.2.2.2.21 CALL GRILZV(-1, 0., 5., 0.,~5.,1.,1.,1.,1.,1.-1.,1.2) Call NCUCRT(770,180,15PACE) Call DLCMRT(570,400,15PACE) DATE = 73110CALL RITEZV(100,1010,1023,90,1,72,1,TITLE,NLAST) PLOT CL VS ALPHA į VCHARV(50,1,JX+60,JY-5,0,TABL2V) . CALL LABLY (ALPHA (J) . JX + 50 . JY . 4 . 1 . 2) CALL VCHARV(90,1,360,595,0.TABL2V) CALL PRINTV(6,0HSYMEDL, JX-20, JY) SETMIV(600, 80,200,450) CALL SETHIV(JEU, 500, 610, 40) CALL CLLCRT(150,800,15PACE) CRTCRT PLOT CLMAXP VS AMAX CALL POINTV(JX, JY, -K, ANY) PLOT DCM VS FTACHU CALL FOINTV(X,Y,-K) CALL PCINTS (X,Y,-I) CALL POINTV(X,Y,1) LO 265 1 = 1,NC PO 250 I =1.NC CALL FRAMEV(3) DO 260 J=1,NA DO 240 J=1 NA UC 230 I=1,NC = LLMAXP(3) JY = JY - 13 X = ALPHA(J)I LIXANA = V = CL(I ad JY = 140 = 780 COMTINUE CGATINUE CUNTINUE CONTINUE CONTINUE CALL CALL × 22 > 220 54 54 10 10 250 358 230 FURTRAN IV G LEVEL Ç U U. ł 8823 8823 8823 8056 0056 0900 0000 242 0400 **2000** 0066 0067 0072 1100 0047 0010 0000 0052 20.6.6. 0000 0200 5100 0400 543 543 9193 0.61 0062 0038 0000 1500 1709 42.00 1100

FURIRAN I	V G LEVEL 20 CRICRT DATE = 73110	0012414E	
0075			
607 6	CALL FULNIV(X,Y,L)	• •	3290
1700			3300
0078	CALL PRINTV(3.3HMAX.IX.IV)		3310
0079	265 CONFINUE		3320
	C PLOTICL VS CD		3330
0800	CALL SETMIV(600, 80.610, 40)		3340
Inna	CALL GRIDZV(~22?		3.150
0082	CALL CLLCRT(560.600.15PACF)		3360
0083	CALL CD0CRT(770.595.15PACE)		3370
0584			3380
0.485	00 270 I = 1.00		3390
0086	X = CO(1, J)	: : : : : : : : : : : : : : : :	3400
0087	v = CL(1, J)		3410
0088	CALL POINTV(X,Y,-I)		3420
0089	270 CONTINUE		3430
. 11600	280 CGNTINUE		3440
	C. PLOT CM VS ALPHA		3450
1600	CALL SETMIV(180,500,260,450)		3460
26.02	CALL GRID2V(-2,-10.,30.,-5.,2.,10.,1.,1.]11.2.2.		3470
0000	CALL VCHARV(90,1) 360, 180, 0, TABL2V)		3480
0005	CALL CMMCRT(150,400,ISPACE)		3490
0.00	DD 300 J=1,NA		0045
0.007	X = ALPHA(J)		0146
1 C C C C C C C C C C C C C C C C C C C	D0 296 1=1,NC		3520
0099		-	2020
01 00	290 CONTINUE		3550
0101	300 CONTINUE		3560
	C PLOT CL VS CM		3570
0102	CALL SETMIY(600. 80.200.250)		3580
0103	CALL GRID2V(-2,2,,-5,,0,,12,,1,,1,2,-1,-2, 2, 2)		3590
-0104	CALL CLLCRT (580, 400, ISPACE)	,	3600
0105	CALL CMMCRT(770,160,ISPACE)		3610
0107	JY = 140		3620
6108		• • • •	0000
0105	CALL PRINTV(6;6HSYMBOL,JX-20,JY)		3650
0110	DC 310 I = 1.0C		3660
0111	DG 320 $J = 1, NA$	•	3670
		ןיז <u>י</u>	3680

.

= CM(1,J) = CL(1,J) LL PQINTV(X.YI) NTINUE = JY - I3 LL PQINTV(JX.JYI.ANY) LL LABLV(ETACMU(I),JX+50,JY,4,1,2) NTINUE TURN	= CM(1,J) = CL(1,J) = CL(1,J) LL POINTV(X.YI) NTINUE = JY - I3 LL POINTV(JX.JYI.ANY) LL POINTV(JX.JYI.ANY) LL LABLV(ETACMU(I),JX+50,JY.4,1,2) NTINUE IURN 0 3 3	DATE = 73110	20 CRTCRT
LL POINTV(X,Y,-I) NTINUE = JY - 13 LL POINTV(JX,JY,-I,ANY) LL LABLV(ETACMU(I),JX+50,JY,4,1,2) NTINUE TURN	LL POINTV(X,Y,-1) NTINUE = JY - 13 LL POINTV(JX,JY,-1,ANY) LL POINTV(JX,JY,-1,ANY) LL LABLV(ETACMU(1),JX+50,JY,4,1,2) NTINUE TURN D		= CM(1,J) = C1(1,1)
NTINUE = JY - 13 LL POINTV(JX,JY,-I,ANY) LL LABLV(ETACMU(I),JX+50,JY,4,1,2) NTINUE TURN	NTINUE = JY - 13 LL POINTV(JX,JY,-I,ANY) LL LABLV(ETACMU(1),JX+50,JY,4,1,2) LL LABLV(ETACMU(1),JX+50,JY,4,1,2) NTINUE TURN D		ALL PRINTVIX.YI)
LL POINTV(JX,JY,-I,ANY) LL LABLV(ETACMU(I),JX+50,JY,4,1,2) NTINUE TURN	LL POINTV(JX,JY,-I,ANY) LL LABLV(ETACMU(I),JX+50,JY,4,1,2) NTINUE TURN D		DNTINUE (= JY - 13
LL LABLV(ETACMU(I),JX+50,JY,4,1,2) NTINUE TURN	LL LABLV(ETACMU(I),JX+50,JY,4,1,2) NTINUE TURN D		LL POINTV(JX, JY, -I, ANY)
		JY+4,1,2)	LL LABLV(ETACMU(I),JX+50,JY VIINHE
			URN
		•••••••••••••••••••••••••••••••••••••••	

FORTRAN	1	G LEVEI	- 20	NCUCRT	DATE = 73110	00/26/45	
0001			SUBROUTIN	E NCUCRT(IX.IY.ISPACE)			3790
0002			EXTERNAL	TABLIV, TABLZV			3800
		ა	CISPLA	Y NU LOWER CASE			3810
0003			CALL VCHAI	RV(90,1,1X,1Y,6,TABL2V)			3820
		J	DISPLAY C				3830
000+			I=IX+ISPAN				3840
0005			CALL VCHAI	RV (90,1,1,IY,19,TABL1V)			3850
0005			I=I+ISPAC	.11			3860
0007			CALL VCHAN	RV(90,1,1,IY-7,11,TABL2V)			3870
0008			RETURN		•		3680
6000			END				3890

3910 3920 3930 3940 3940 3940 3970 3970 3980 3900 4010 i 00/26/45 ; i = 73110 DATE ; I = IX + ISPACE CALL VCHARV(90,1,1,1Y,19,TABL1V) DISPLAY M LCWER CASE I=1+1SPACE CALL VCHARV(90,1,1,iY-10,36,TABL1V) CALL VCHARV (90,1,1X,1Y,3,TABL3V) : SUBRCUTINE DLCMRT (IX, IY, ISPACE) EXTERNAL TABLIV, TABL3V DISPLAY DELTA i DLCMRT : ¢ UISPLAY C : RETURN 20 END FORTRAN IV & LEVEL J J С О 0003 6008 0009 1000 0002 0000 0005 0000 0001

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FORTRAN	IV G LEVEL	. 20	DCLCRT	DATE = 73110	00/20/45	
0001		SUBROUTINE DCLCRT	(IX, IY, ISPACE)			
0002		EXTERNAL TABLIV, T	ABL2V, TABL3V			40
	C	DISPLAY OPEN PAREN	J.H.			Ó 4
0:03		CALL VCHARV (96.1.	IX, IY, 60, TABLIV)			40
	J	DISPLAY VELTA				40
0004		I = IX + ISPACE				4
5000		CALL VCHARV(50,1,1	, IY, 3, TABL3V)			40
	ა	DISPLAY C				40
0000		I = I + I SPACE				41
0007		CALL VCHARV(90,1,1	, IY, 19, TABLIV)			4 1
	J	DISPLAY L IN LOWER	CASE			41
0008		I=I+ISPACE			:	14
5000		CALL VCHARV(90,1,I	,IY-10,35,TABL1V)			4
	J	DISPLAY CLOSE PARE	Z11			41
0010		I=I+ISPACE				41
0011		CALL VCHARV(90,1,1	•IY,28,TABL1V)			4
	J	DISPLAY THETA IN LU	OWER CASE			41
0012	-	I=I+ISFACE	1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1	:		4
0013		CALL VCHARV(90,1,1	,IY-15,7,TABL2V)			42
GU14		RETURN				42
0615		END				42

FORTRAN IV G LEVEL	. 20 CODCRT	DATE = 73110	00/26/45	
0001 0002 C	SUBRCUTINE CDECRT(IX,IY,ISPACE) External Tabliv, Tabl2v, Tabl3v Draw CD			423 424
0003 C	CALL VCHARV(90,1,1X,1Y,19,TABLIV)			425
0.04 0005 0006	9ISPLAY C I = IX + ISPACE Call VCHARV(50,1,1,1Y-10,20,TABL1V) Return			000 000 000
0001	END			432
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	SUBROUTINE CLACRITIX.IV.ISDACEN	······································	
002	EXTERNAL TARITY TARITY TARITY		
C	DRAW CLART MULLY MULLY		
C	DICOLAV C		
P003	CALL VCHAPVION 1 1V 1V 10 TICL		
U	DICE TOTALY (709 19149119199 1456 14)		
004			
005	I = IX + ICDACE		
006	CALL VCHARV(90.1.1.1.36.TAB:111		
J	DISPLAY AI PHA		
007			
206	I = I + I CPACE		
009	CALL VCHARVION 1. 1. 1. 1. TABLOW		
010	RETURN		
011	END		

0001 SUBROUTINE CMMCRITIX, IY, ISFACE) 448 0002 EXTERNAL TABLLY, TABL2Y, TABL3Y 448 0003 C DISPLAY C 448 003 C OLAL VCHARY(9), I, IX, IY, 19, TABL1Y) 448 003 C ALL VCHARY(9), I, IX, IY, 19, TABL1Y) 453 003 C ALL VCHARY(9), I, IX, IY, 19, TABL1Y) 453 003 C ALL VCHARY(9), I, IX, IY, 19, TABL1Y) 453 0003 C ALL VCHARY(9), I, IX, IY, 19, TABL1Y) 453 0003 C ALL VCHARY(9), I, I, IY-10, 36, TABL1Y) 453 0004 ELUIAN ENU 453 0007 ENU ENU 453	FORTRAN	IV G LEVEL	20	CMMCRT	DATE = 73110	00/26/45	
003 C DISPLAY C 451 003 C DISPLAY C 451 005 C DISPLAY C 452 006 C I = IX + ISPACF 453 006 CALL VCHARV(90:11:1Y-10.36.TABLIV) 453 005 CALL VCHARV(90:11:1Y-10.36.TABLIV) 453 000 RELUN 456 0005 ELU 411.1Y-10.36.TABLIV) 455 0005 ELU 411.1Y-10.36.TABLIV) 455 0005 ELU 411.1Y-10.36.TABLIV) 455	0001 0002	i.	SUBROUTINE CMMCRT(IX External tabliv, tab Draw Cm	X,IY,ISFACE) Bl2V, Tabl3v			4 4 4 7 9 9 9 9 9
6104 C 1 = 1X + 1SACF 0005 Call VCHARV(90;1,1,1,Y-10,36,TABL1V) + 55 0007 END NETURN + 11,1Y-10,36,TABL1V) + 55 0007 END 0007 END 0007 + 54 455 0007 + 54 455 455 455 455 455 455 455	0003	ן ייני (ן	CALL VCHARV(90,1,1X,	.IY,19,TABL1V)			4 4
0005 CALL VCHARVI96111117-10.5611AELIV) 455	0004	ى	UISPLAY M I = IX + ISPACF			: : :	1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1
6007 ENU	0005 0006		CALL VCHARV(90,1,1,1 Return	IY-10,36,1ABLIV)			456
	0007		END		:		457
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4590 4600 4610 4620 4630 4650 4660 4670 4580 4640 ÷ ł : ! 00/26/45 i DATE = 73110i : ł CALL VCHARV(50,1,1,1Y-10,35,TABL1V) RETURN : ; SUFFCUTIME CLECRT(IX, IY, ISPACE) EXTERNAL TAELIV, TAELEV, TAELEV DRAW CL CISPLAY C CALL VCH.RV(54,1),IX,IY,19,TAELIV) i -----• • • • • CLLCRT : : DISPLAY L I = IX + ISPACE ł . ENL. ن د ı FORTRAN IV 6 LEVEL υu J ł 1 2000 1000 0003 0004 0005 0005 0005



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