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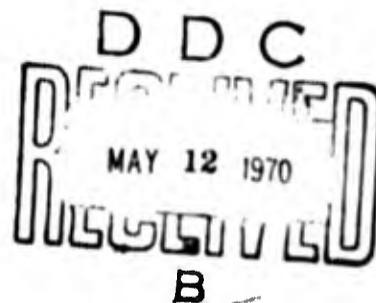


# **Aerospace Research Laboratories**

## **ELECTRIC ARCS IN TURBULENT FLOWS, IV**

GERHARD FRIND  
BEN LEE DAMSKY  
GENERAL ELECTRIC COMPANY  
PHILADELPHIA, PENNSYLVANIA

Contract No. F33615-67-C-1374  
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OFFICE OF AEROSPACE RESEARCH  
UNITED STATES AIR FORCE  
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## FOREWORD

This final report was prepared by the Laboratory Operation of the General Electric Company, Power Transmission Division, Philadelphia, Pennsylvania on Contract F 33 615 67 C 1374. The work was started in February, 1967 and the Report manuscript was completed in December, 1969.

The authors wish to acknowledge the able support of Mr. V. L. Harvey who assisted in design of the apparatus and in the collection of data.

The authors also thank T. J. Tuohy and Ping Tsao, who prepared and ran the computer program for calculating chemical equilibrium of a gas mixture.

We thank Dr. T. H. Lee for his continued support and encouragement over the five years involved in this program.

During this period there were helpful discussions and suggestions from Messrs. E. E. Soehngen, P. W. Schreiber, Dr. J. G. Skifstad, and Dr. H. O. Schrade of the Thermomechanics Branch of the Aerospace Research Laboratories.

Technical monitor for this report was Major K. H. Schumaker.

## ABSTRACT

Heat, mass, and momentum transfer were investigated experimentally in very long, cylindrical axial flow arcs in laminar and turbulent flow conditions. Special attention was given to assure that the end of the arc tube had a section of fully developed flow. In this section the number of significant terms in the energy balance equation is reduced and interpretation of experimental results is simplified. Strong evidence is presented for the existence of turbulence by the fact that all three transfer mechanisms are considerably accelerated in what we consider the "turbulent mode" and also by direct optical study of the arc by a light probe. The onset of turbulence was found with all our methods to occur around 1.5 gram/sec gas flow rate. A considerable number of measurements of heat and momentum transfer are presented in the following operating ranges: Gases - argon and nitrogen; Tube Diameters - 0.5 and 0.7 cm; Tube Length - 50 cm; Mass Flow Rates - 0.1 to 15 gram/sec; Ambient Pressures - 1 to 20 atmosphere; Currents 25 to 200 amperes (none higher).

## LIST OF SYMBOLS

$B$	Boltzmann's Constant
$D$	Tube Diameter
$DP$	Dissociation Energy
$E$	Electrical Gradient
$e$	Average Height of Wall Roughness
$F$	Potential Energy of Energy Level
$g$	Degeneracy of Energy Level
$I$	Spectral Intensity
$IP$	Ionization Energy
$L$	Length of Arc Tube
$m$	Mass of Electron
$N_0$	Density of Radiating Particles in Ground State
$P$	Pressure
$R$	Radius of Tube Containing Arc
$Re$	Reynolds Number
$r$	Radial Distance from Arc Axis
$S$	Radiation Rate Per Unit Volume
$T$	Absolute Temperature
$v$	Velocity
$W$	Axial Flow Rate
$x, y$	Cartesian Coordinate
$Z$	Partition Function
$\Delta T$	Amplitude of Temperature Fluctuation

LIST OF SYMBOLS (Cont'd)

$K$	Thermal Conductivity
$K^*$	Effective Thermal Conductivity Produced by Turbulent Mixing
$\lambda$	Wavelength
$\mu$	Kinematic Viscosity
$\varphi$	Excitation Energy
$\rho$	Density
$\sigma$	Electrical Conductivity
$\Omega$	Electrical Resistance

Subscripts

$D$	With Reference to Tube Diameter
$e$	Electron
$i$	Ion
$J$	Atomic Energy Level Index
$o$	Ground State

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## ELECTRIC ARCS IN TURBULENT FLOWS, IV

### 1. INTRODUCTION

#### 1.1 General Comments

This report is the last in a group of publications (1, 2, 3)\* devoted to studying the importance of turbulence for the transport of energy, momentum and mass in coaxial flow arcs. More specifically, work was concentrated on turbulent effects in an area which, in fluid dynamic studies, is usually called the fully developed section of coaxial flows. This, of course, is a very strong restriction on the broad objective of furthering understanding of arc turbulence in toto. Nevertheless, it will be seen from the results described in this paper that pursuit of even our limited objective offers formidable problems.

#### 1.2 Review of Work from Earlier Reports

To facilitate the task of following the zig-zag course of our investigation, we present in Table 1 a matrix containing what we consider the essential results obtained in all four reports, including this one. The following short discussion of this "matrix" will also familiarize the reader with our earlier work for further reference.

##### 1.21 Theoretical Aspects

As for the Theoretical Aspects of this investigation, we discussed in some detail the complexity of the physical processes(1, 2), demonstrated by way of the general energy balance equation. This discussion showed also that the concept of "fully developed conditions" cannot be used in the same way for electrical arcs as for low temperature fluids.

This is so because in hot plasma streams it is difficult to make the "flow" terms in the energy balance equation negligibly small. One must resort to experimental conditions of plasma velocities well under the speed of sound and of limited arc tube diameters (2). Under these strong limitations, fully developed conditions can be expected in the end section of long coaxial arc heaters and an energy balance with a significantly smaller number of terms can be retained (2, 3). This result was essential for our investigation because our first goal was to find, in axial flow arcs, experimental conditions in which the existence of radial turbulent energy

\* The numbers in parentheses are references; see page 54.

REPORT	THEORY	EQUIPMENT	VOLTAGE GRADIENT	PLASMA STRUCTURE	TEMPERATURE MEASUREMENTS	PRESSURE GRADIENTS
I	concept of coaxial fully developed flow as goal calculations show turbulence is reasonable	cascade discs with terminal pressure tank	low current measurements	measurements of flow velocities at low currents		
II	theory of where achieve fully developed flow calculation of energy balance	high voltage power supply	up to 80 amperes in 1 cm tube	photographic studies of 80 ampere arc at several flow rates		
III		high current power supply	200 amperes in argon and nitrogen	velocity measurements with image converter camera	temperature distribution determined by N <sub>2</sub> band line	11.2 atm pressure gradients
IV	higher pressure capability		wide variety of pressures and flow rates in argon extended measurements in nitrogen	amplitude and frequencies of light intensities fluctuations at different flow rates and currents	Sodium temperature point	Pressure gradients measured at many flow rates and pressures

TABLE I

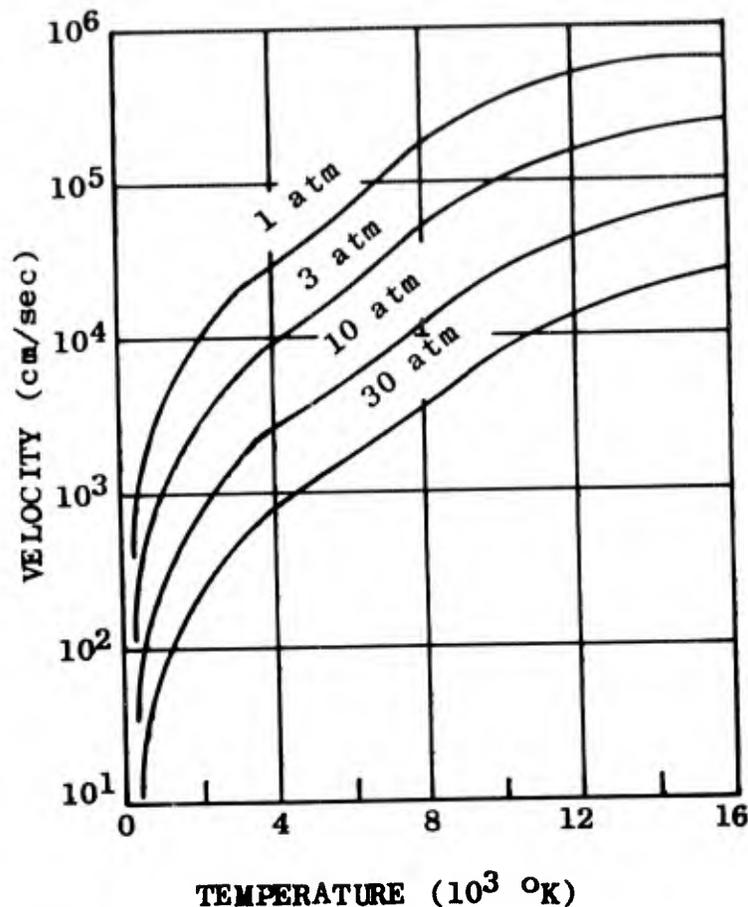


Figure 1-Estimate of Minimum Velocity for Turbulence.

Velocity for which the Reynolds number is equal to 2000 is plotted against temperature (assumed homogeneous).

Gas: Nitrogen  
 Pressure: as marked  
 Tube Diameter: 1 cm

transport could be proven by simple measurements. We felt that such a proof was too difficult a task if the usually strong flow terms could not be eliminated and the analysis correspondingly simplified.

The experiments in all four reports were therefore planned with the specific condition of fully developed flow in mind. This entails:

1. Use of velocities well subsonic (below 500 meter/sec)
2. Use of small diameter arc tubes (1.0 cm and less)
3. Concentration of measurements at the end section of a flow tube with as high L/D value (length/diameter) as possible (we used ratios from 70 to 150).

Another major aim of earlier theoretical work was the demonstration that turbulence could be reasonably expected under normal operating conditions of arc heaters. Figure 1 reproduces a calculation made in (1) assuming a constant temperature air plasma with pressures between 1 and 30 atm. It is seen that, at the higher pressures, Reynold numbers substantially higher than 2000 can be reached with subsonic velocities even at temperatures between 10,000 and 15,000°K, reasonable values for the core of high current arcs. If we now remember that the electric arc has layers of widely different temperatures, it appears likely that the outer, cooler layers of an arc with their higher density and lower viscosity may show turbulent effects even at quite moderate pressures and flow velocities.

This same calculation may be viewed in another light: using our plasma velocity measurements ((3) section 4.3) as a basis, the velocities of the 10 atmosphere curve in figure 1 were converted into flow rates to give figure 2. The assumptions were: a 10,000°K core temperature for the arc used in (3), and relative velocity and temperature distributions unaffected by changes in the central temperature. Neither assumption is accurate but the curve still serves as an estimate of the requirements for producing turbulence throughout an arc.

## 1.22 Equipment

Table I gives an idea of the gradual development of our experimental apparatus. The cascade tube of Maecker (4) was used from the beginning, but without the usual water cooling of the copper discs. We felt this omission was a necessary simplification if only for financial reasons. The number of discs in one tube exceeded 500 at times. This large number resulted from the length of tube required for fully developed flow and from the 0.05cm thickness of the individual discs. The discs were made so thin for two reasons; first, voltage gradients as high as 200 volt/cm were found in turbulent arcs, and second, it was expected that a strike-over to the copper walls would occur more easily on account of the violent plasma motions of turbulent conditions.

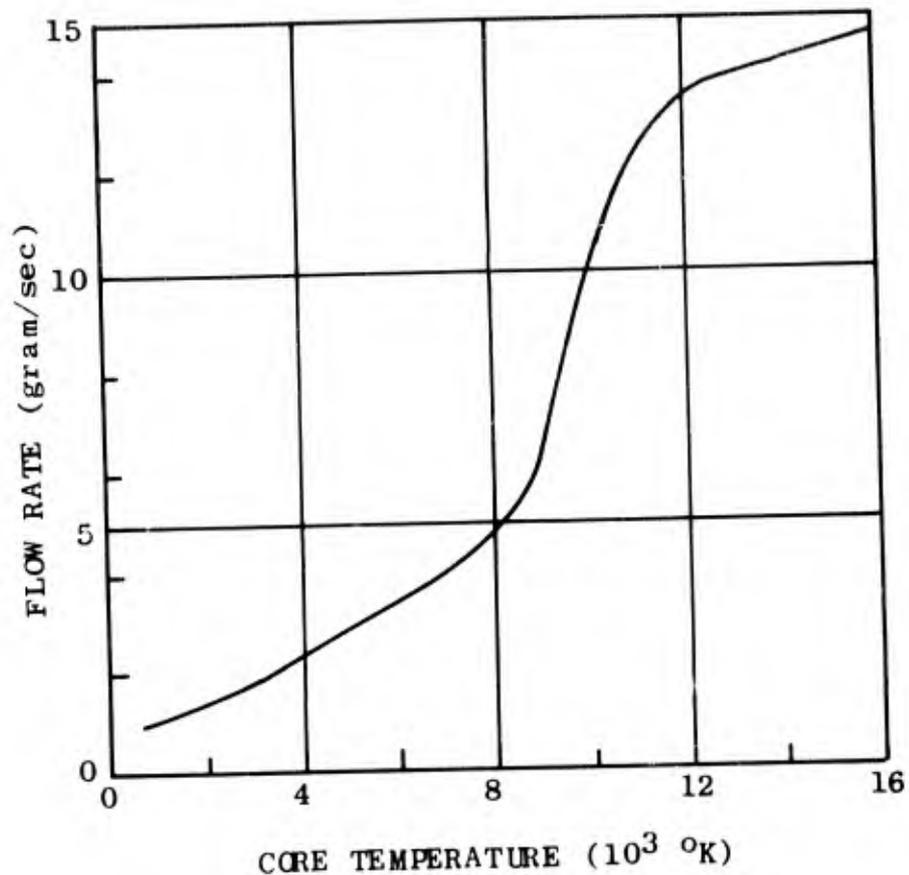


Figure 2—Estimate of Minimum Flow Rate for Turbulence in 0.7 cm Diameter Tube. The estimate of figure 1 was adapted using the velocity vs. flow rate measurements of (3).

Gas: Argon  
 Pressure: 10 atm  
 Tube Diameter: 0.7 cm

There are other advantages in using simple apparatus. To avoid overheating one has to make relatively short time arc runs ( $\sim 1/2$  sec), and so the electrical power source can be used at a multiple of its steady state capability (1, 2). Also under the thermal conditions of short arc runs, a segment of the copper cascade can be replaced with an uncooled quartz tube, a necessity for studying turbulent structure and flow velocity by high speed photography.

Turning to the disadvantages of our short time arc runs, one first notes the necessity for proof that steady state conditions have been reached with regard to time. It was found that the slowest process after starting of the arc was the establishment of steady state flow conditions along the arc heater. Flow was monitored by the pressure drop along the tube (3). We were careful in our experiments to evaluate measurements only with the plasma pressure close to steady state conditions. However, it must be reported that it was difficult to fulfill that condition equally well at all times. Development was slowest for tests combining high flow rate with high current.

### 1.23 Measurements in the Fully Developed Flow Region

#### 1.231 Heat Flux and Momentum Transfer

Perhaps the most convincing measurements in our group of reports were those of radial heat and momentum transfer in the fully developed section of the long, cylindrical arc tube. This was accomplished (2, 3) by measurements of the axial voltage and pressure gradients. The results of these measurements gave strong support to our conclusion that a laminar state of the arc exists at low flow rates (below about 1.5 gram/sec) while an increasingly turbulent character of the arc obtains at flow rates higher than that number. In our last report (3) the experimental range of voltage gradient measurements was expanded to currents up to 200 amps and to nitrogen in addition to argon; the strong effect of flow rate on pressure gradient was shown at 11.2 atm of argon with 125 amps.

#### 1.232 Plasma Structure

Insight into the plasma structure was gained in all our previous reports by high speed photography. Fig. 21-24 of (2) show convincingly the laminar character of the arcs at low flow rates and the increasingly turbulent character at high flow rates where the arc column is apparently "broken up" into plasma globules of irregular sizes and shapes. Of course, current is flowing steadily and relatively undisturbed even through a strongly temperature modulated plasma.

#### 1.233 Plasma Velocity

Although the measurement of the complete radial velocity distribution  $V(r)$  is a long range goal of our investigation, in our last report (3), an approximate measurement of the average plasma speed close to the axis of the tube was accomplished. We followed plasma globules as they prodeded along the tube with a high speed camera. Fig. 18 of (3) shows the plasma velocity as function of the mass rate of flow for argon.

### 1.234 Plasma Temperature

In our last report a temperature measurement was made on both a laminar and a turbulent 50 amp arc in 11.2 atm. of nitrogen. The choices of current and gas species were unfortunately not in the "main stream" of our investigation, which would have called for 125 amps and argon.

The reason for this inconsequence was our belief that a turbulent temperature measurement could best be made with the Milne method (5) which is insensitive against small fluctuations (6, 7). Furthermore nitrogen has well known molecular bands which radiate strongly even below 10,000°K which is the temperature range of interest. The atomic lines of Argon, on the other hand, have peak intensities at much higher temperatures and are difficult to use at 10,000°K.

The result of the measurement was a surprisingly close agreement between the laminar and turbulent temperature distributions  $T(r)$ , see Fig. 28 and Table III in (3).

### 1.3 Goals of This Investigation

One of the long range goals of our work on turbulent arcs is to find empirical relationships between energy, momentum and mass transfer and nondimensional properties of the temperature and velocity field. One more modest goal of this particular report is to quantify, as far as possible, such relationships as a function of the mass rate of flow, which in turn is approximately proportional to some "average" Reynolds number. This requires the collection of data at various combinations of test parameters so correlations of turbulent levels with operating conditions can be made. Other goals are to shed more light on turbulent structure and on the onset of turbulence, and further, to develop measuring methods suitable for a more precise determination of such important parameters as plasma temperature and velocity. With expanded facilities, the operating parameters can be extended to higher currents and pressures. Values chosen for testing should be as close as possible to typical parameters used in constricted arc heaters.

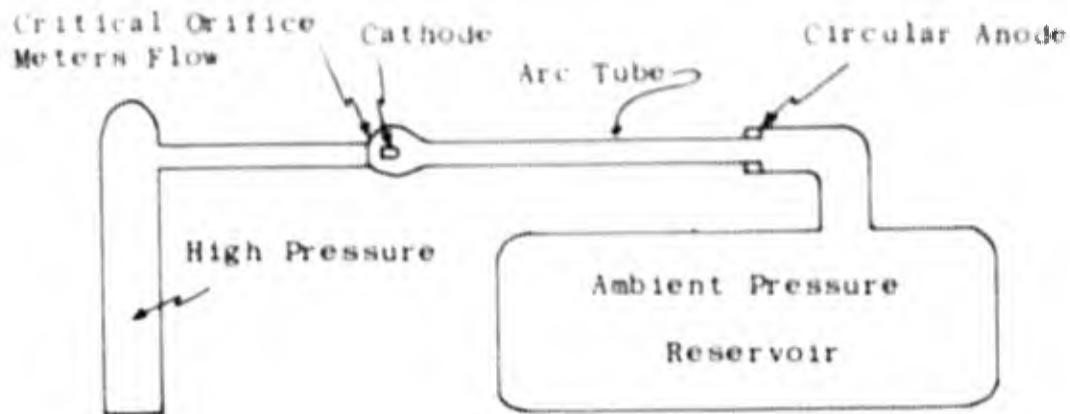


Figure 3-Schematic Diagram of Test Apparatus. The flow rate is set by the size of the critical orifice and the value of the upstream pressure which is three times the ambient pressure. Arc initiation is accomplished by exploding a 0.001cm diameter tungsten fuse wire stretched from anode to cathode. (See also Fig. 3 and 4 in (3)).

## 2. APPARATUS

The basic experimental arrangement has not changed since the last report, (3). A schematic diagram of the major equipment is shown in figure 3. We utilize the thermal storage capability of rectifiers, resistors, and cascade discs to withstand transient heating and so operated them at levels beyond their steady state capacity. The uncooled copper discs have been operated with heat fluxes of  $13\text{Kw}/\text{cm}^2$  for short periods. Transient operation ( $\leq 0.5$  sec.) makes it necessary to record data with Cathode ray oscillographs (CRO). Gas flow is metered by critical orifices whose upstream pressure is kept at three times the ambient reservoir pressure. Orifices were calibrated with an absolute accuracy of  $\pm 10\%$ ; relative accuracy is  $\pm 2\%$ . Because of the extremely high potential gradients produced by turbulent nitrogen, cascade discs were made of 0.5mm thick copper; for experiments in argon, 1mm discs were possible. Preparation of the discs included cleaning in an ultrasonic bath and etching a 0.02mm deep layer from the surface. This process insured cleanliness and removed the sharp edges which increase the chances for "strike over" at the walls. The discs were insulated by 0.25mm thick teflon washers which were recessed from the arc to prevent contamination. Before experiments, the discs were carefully aligned with a steel rod whose diameter matched the inner diameter of the discs.

## 3. EXPERIMENTS IN THE FULLY DEVELOPED FLOW REGION

### 3.1 Heat Transfer as Measured by Voltage Gradients

#### 3.11 Choice of Parameters

One of the long range objectives of our study is to find an experimental correlation of turbulent heat transfer and non dimensional quantities of the flow field such as some "average" Reynolds number.\* To accomplish this we need quantitative experimental data about the radial heat flux, the temperature distribution, and the velocity distribution. Although we do not have this data, but have measurements of potential gradient, flow rate, and average velocity in the arc core, we feel it is useful to correlate these data. Therefore we attempt in the following to correlate radial heat flux with axial flow rate which, in turn, is approximately proportional to the Reynolds number  $Re_D$  of the experiment.

For these measurements we chose to study the 125 ampere argon arc, which fills our tube rather well and for which a large amount of pressure gradient data was also taken. In the graphs which follow the directly measured quantities, arc voltage gradient, and flow rate are compared. Since the arc current is held constant at 125  $\pm$  4 amperes, the arc power is proportional to the voltage gradient and since, under the conditions of our experiments, the flow is fully developed in the end section of the flow tube, all of the arc power in this section is transmitted to the wall. The flow rate, on the other hand, would be strictly proportional to the Reynolds number only if increasing flow does not alter the temperature distribution and the relative velocity distribution. Our preliminary indications are that the temperature does not change radically with flow at constant current (3) so that flow rate may be a meaningful measure of relative Reynolds number.

There are five test variables in the heat transfer experiments: gas species; tube diameter; mass flow rate; ambient pressure; and current. In the tests described below four variables are held fixed and the fifth is varied. This systematic method helps to isolate and identify quantitatively the significance of each variable for heat loss. The ranges of these variables are:

\* An "average" Reynolds number for strongly non homogeneous flow conditions, such as found in our experiments, remains to be defined. See Eckert's paper (8) for an attempt in this direction.

gas species:	argon or nitrogen
tube diameter:	0.5 or 0.7 cm
mass flow rate:	0.1-15 gram/sec
ambient pressure:	1-20 atmosphere
current:	25-200 ampere (some higher)

The voltage gradient measurements which follow were taken in the terminal portion of our long tubes where fully developed flow conditions exist.

### 3.12 Methods

Potential gradients were measured by attaching resistive dividers ( $6 \times 10^6 \Omega$  input impedance) to eight cascade discs spaced along the arc tube. A typical current drawn from the plasma region would be  $3 \times 10^{-5}$  amperes. Data was recorded as CRO traces.

The potential gradients were almost invariably constant along the tube except in a few cm of the entrance region. As flow rate increased to several grams per second, the potential gradient line of the graph showed a significant voltage at the interception of the axis representing cathode position. This additional voltage resulted from a higher potential gradient in the first few cm of the cascade tube. The excess potential drop increased linearly with flow rate, finally exceeding 100 volts. Since the value of this potential drop corresponds roughly to the power required to heat the incoming gas flow to several thousand degrees Kelvin, we take this as the explanation of the effect and do not include this short segment of the arc constrictor length in the gradient measurements which follow.

A slightly different observation was made of the nitrogen arc below 100 amperes; see (3) for a discussion of this effect.

### 3.13 Results

#### 3.131 Variable Flow Rate

##### 3.1311 Argon

Figure 4 shows the effect of the flow rate on energy loss in a 15 atmosphere argon arc on a log-log plot. The slope is not constant but continues to increase at the highest flow rates, reaching the value of 0.88 at the end of the graph. This compares closely with the empirical dependence of heat transfer coefficient on Reynolds number found in (9) for low temperature gases, see figure 5. In such homogeneous, low temperature fluids, heat transfer rises rapidly with the onset of turbulence and eventually increases with the 0.80 power of Reynolds number.

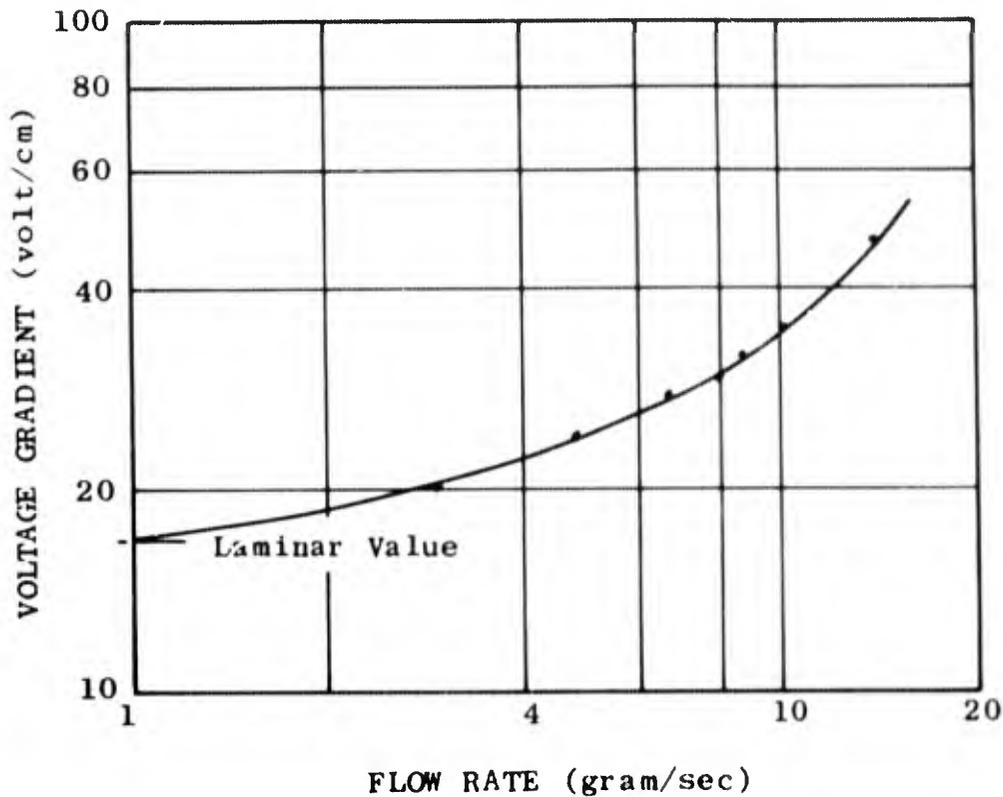


Figure 4-Voltage Gradient vs. Flow Rate at 15 atmospheres of Argon in a 0.7 cm Tube. The log-log plot shows an ever increasing slope.

Gas:	Argon
Pressure:	15 atm
Current:	124 $\pm$ 4 amp
Tube Diameter:	0.7 cm
Tube Length:	50 cm

Whereas our experimental data (figure 4) does not settle into a linear region on the log-log plot, the proper slope is approximated closely for the highest flow rates. Let us also bear in mind that we are plotting radial heat transfer against flow rate, which is not precisely proportional to Reynolds number but only approximately so. Tests at still higher flow rates might have indicated whether the slope of figure 4 would remain constant. However, the large axial pressure gradients involved at these flow rates would make the establishment of fully developed conditions questionable ((2), pages 7-9 and see also section 3.2 for measurements of pressure gradients).

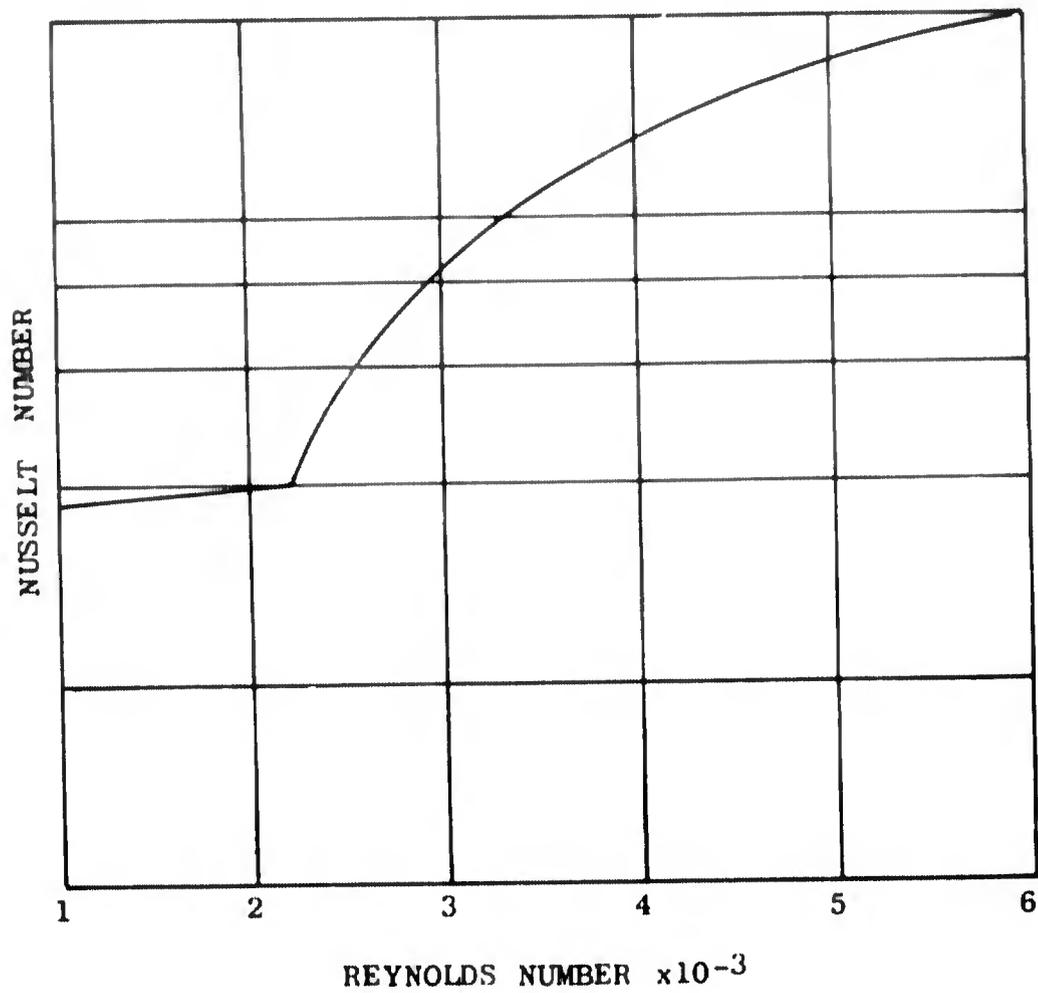


Figure 5-Heat Transfer in a Turbulent Tube (Homogeneous Fluid) The data are taken from (9).

The variation of the slope in figure 4 suggested an exponential dependence. This was checked in figure 6 which shows energy losses from a 6 atmosphere argon arc on a semilog plot. Three regions can be seen:

- I 0-1.2 gram/sec
- II 1.2-3.5 gram/sec
- III 3.5 gram/sec and up

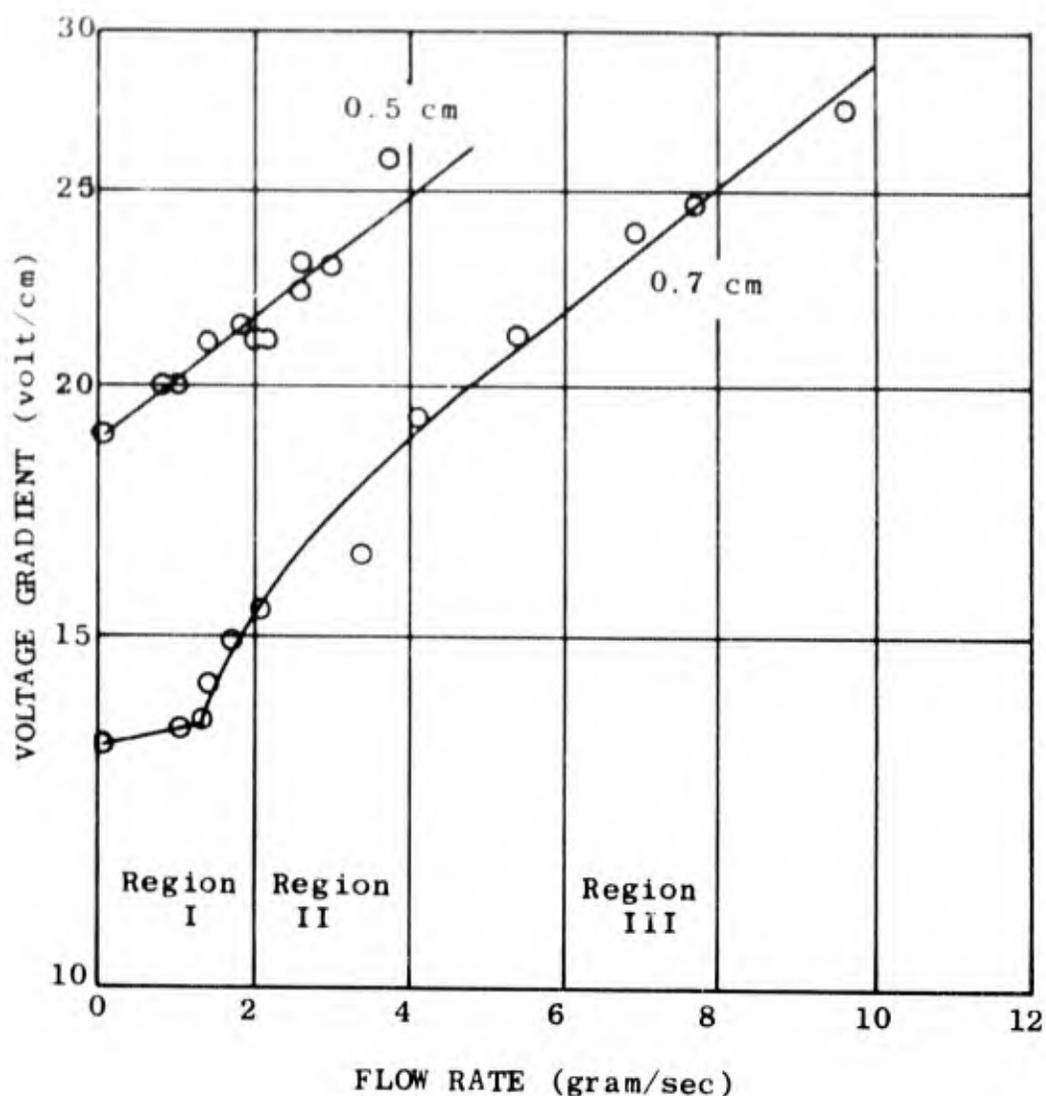


Figure 6-Voltage Gradient vs. Flow Rate at 6 atmospheres of Argon in 0.5 and 0.7 cm Tubes.

Gas: Argon  
 Pressure: 6 atm  
 Current: 125 amp  $\pm$ 4  
 Tube Diameter: 0.5 cm and 0.7 cm  
 Tube Length: 50 cm

In region I the energy loss is nearly independent of flow rate, as would be expected from laminar flow; in region II there is a sudden increase in energy loss, which could be explained by the onset of turbulence; and in region III the relationship becomes approximately linear on the semilog plot. Notice too that the lowest flow data point lies approximately on the extension of the straight line in region III.

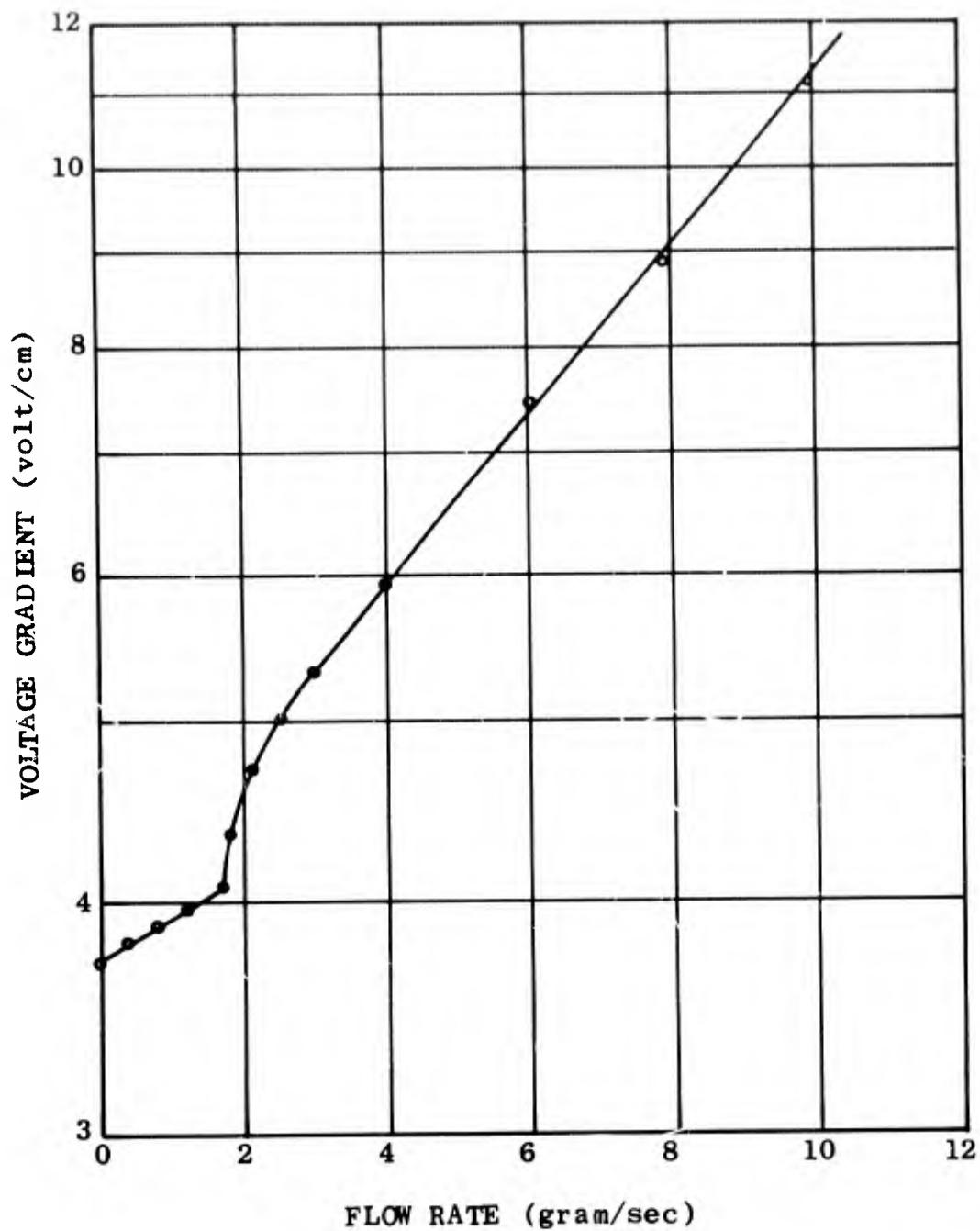


Figure 7-Voltage Gradient vs. Flow Rate at 1 Atmosphere of Argon in a 1 cm Tube, Data from Runstadler (10)

Gas:	Argon
Pressure:	1 atm
Current:	50 amp
Tube Diameter:	1 cm
Tube Length:	40 cm

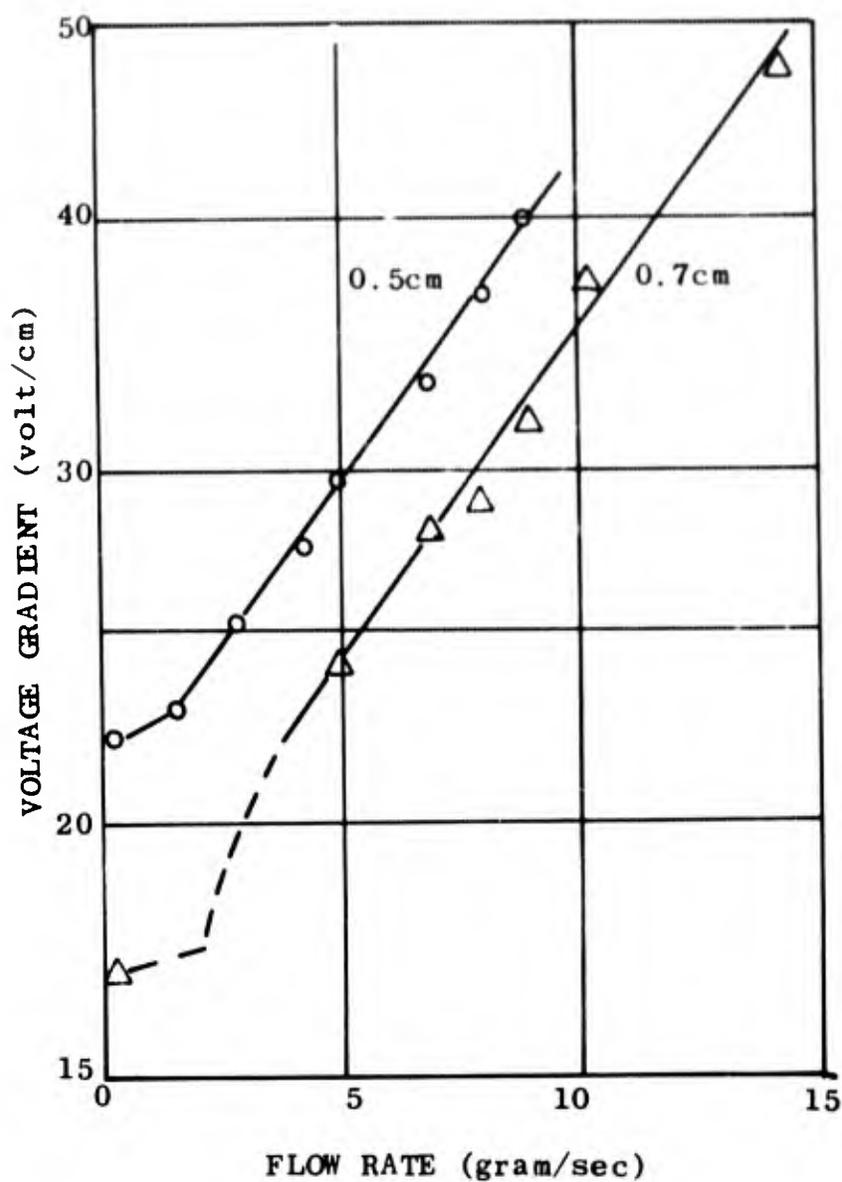


Figure 8-Voltage Gradient vs. Flow Rate at 15 Atmospheres of Argon in 0.5 and 0.7 cm Tubes

Gas: Argon  
 Pressure: 15 atm  
 Current: 125 ±4 amp  
 Tube Diameter: 0.5 and 0.7 cm  
 Tube Length: 50 cm

These same features can be seen in Runstadler's data on arc voltage measurements (10), figure 7. We replotted Runstadler's results on a semilog scale and there does seem to be an exponentially rising region at higher flow rates. Notice that his conditions are somewhat different (lower current and pressure, and a larger tube diameter), but his result is similar to ours.

It is interesting to see that a smaller arc tube diameter increases energy losses to the walls but maintains the same slope in region III. Figure 8 shows the results of tests in 0.5 and 0.7 cm I. D. tubes at 15 atmospheres of argon. The losses go up about 25% in the 0.5 cm tube but the two lines are parallel within experimental error. This same fact can be seen in figure 6 where the 6 atmosphere results are compared. The power ratio for 6 atmosphere tests are comparable. The power ratio at 6 atmospheres cannot be measured so accurately but is approximately 30%. Notice that the curves again have the same slope and this slope is identical with that found at 11.2 and 15 atmospheres, figure 9. The 20 atmosphere results are shown in figure 10.

If this study of heat transfer is compared with the study of pressure gradients, see section 3.2, it is immediately seen that near 1.5 gram/sec. both the pressure gradient and the energy loss begin to rise rapidly. This correlation again lends weight to our interpretation of this point as the onset of turbulence. Compare also Runstadler's value of 1 gram/sec. (10).

A summary of all the above results in a 0.7 cm I. D. constrictor is contained in figure 11. It has the additional advantage of showing the results on a linear plot as well as the semilog plot used in figure 6 through 10. The striking features of the curves are the strong effect which flow rate has on power loss (an effect which becomes more pronounced at higher flows) and the departure from the lower pressure data which occurs at an ambient pressure of 20 atmospheres. Tests at 1, 3 and 6 atmospheres do not extend to the highest flow rate because such steep pressure gradients would be produced\* that the flow terms in the energy balance equation could not be neglected and consequently the potential gradient would not be proportional to radial heat flux.

\* See section 3.2 for measurements of pressure gradients.

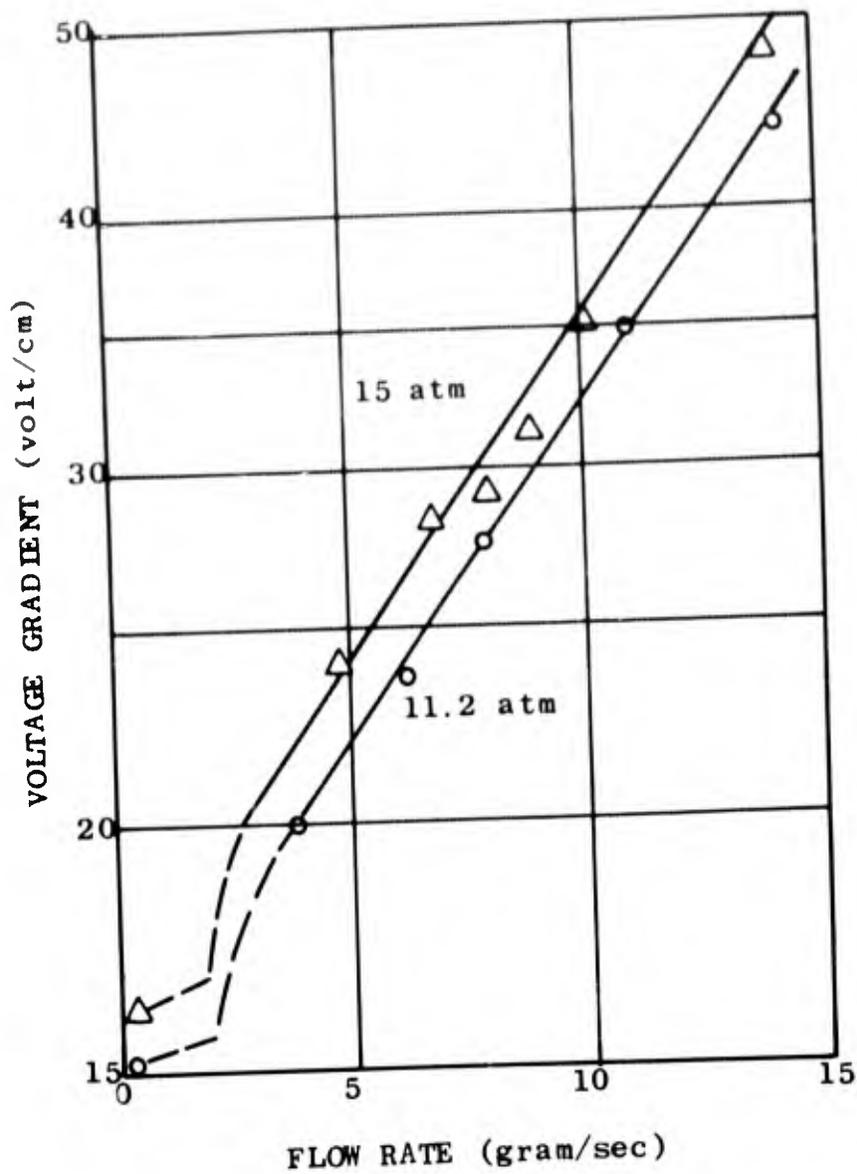


Figure 9-Voltage Gradient vs. Flow Rate at 11.2 and 15 Atmospheres of Argon in a 0.7 cm Tube

Gas:	Argon
Pressure:	11.2 and 15 atm
Current:	125 $\pm$ 4 amp
Tube Diameter:	0.7 cm
Tube Length:	50 cm

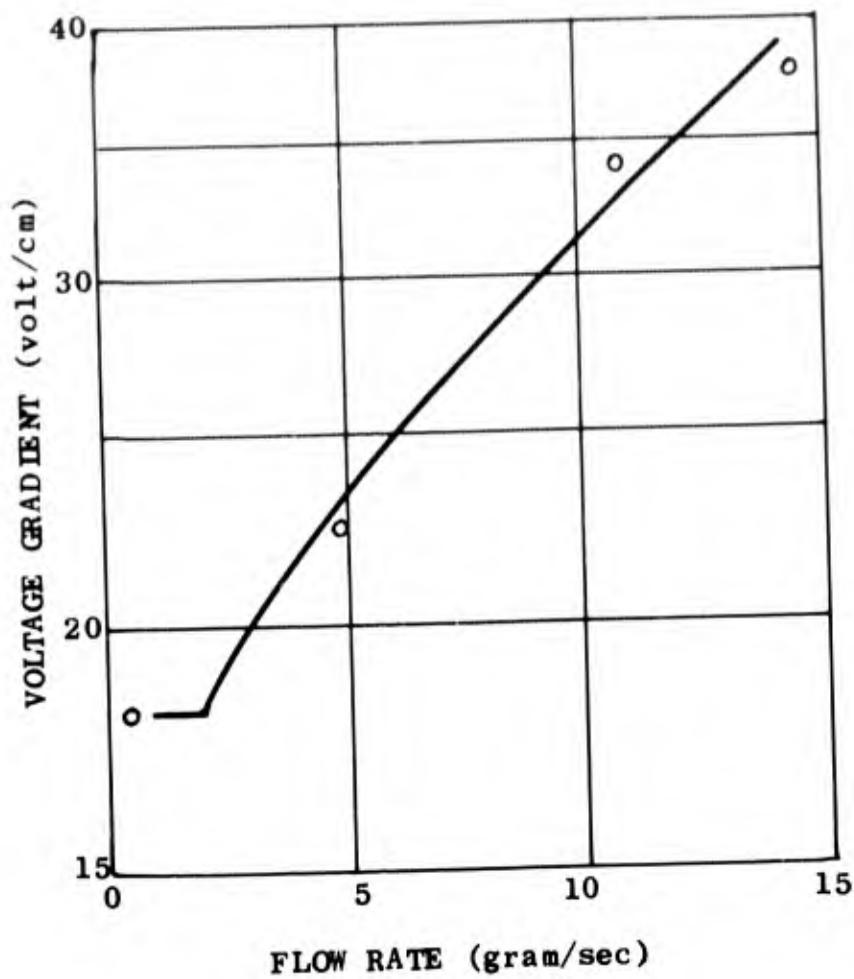


Figure 10-Voltage Gradient vs. Flow Rate at 20 Atmospheres of Argon in a 0.7 cm. Tube.

Gas:	Argon
Pressure:	20 atm
Current:	125 $\pm$ 4 amp
Tube Diameter:	0.7 cm
Tube Length:	50 cm

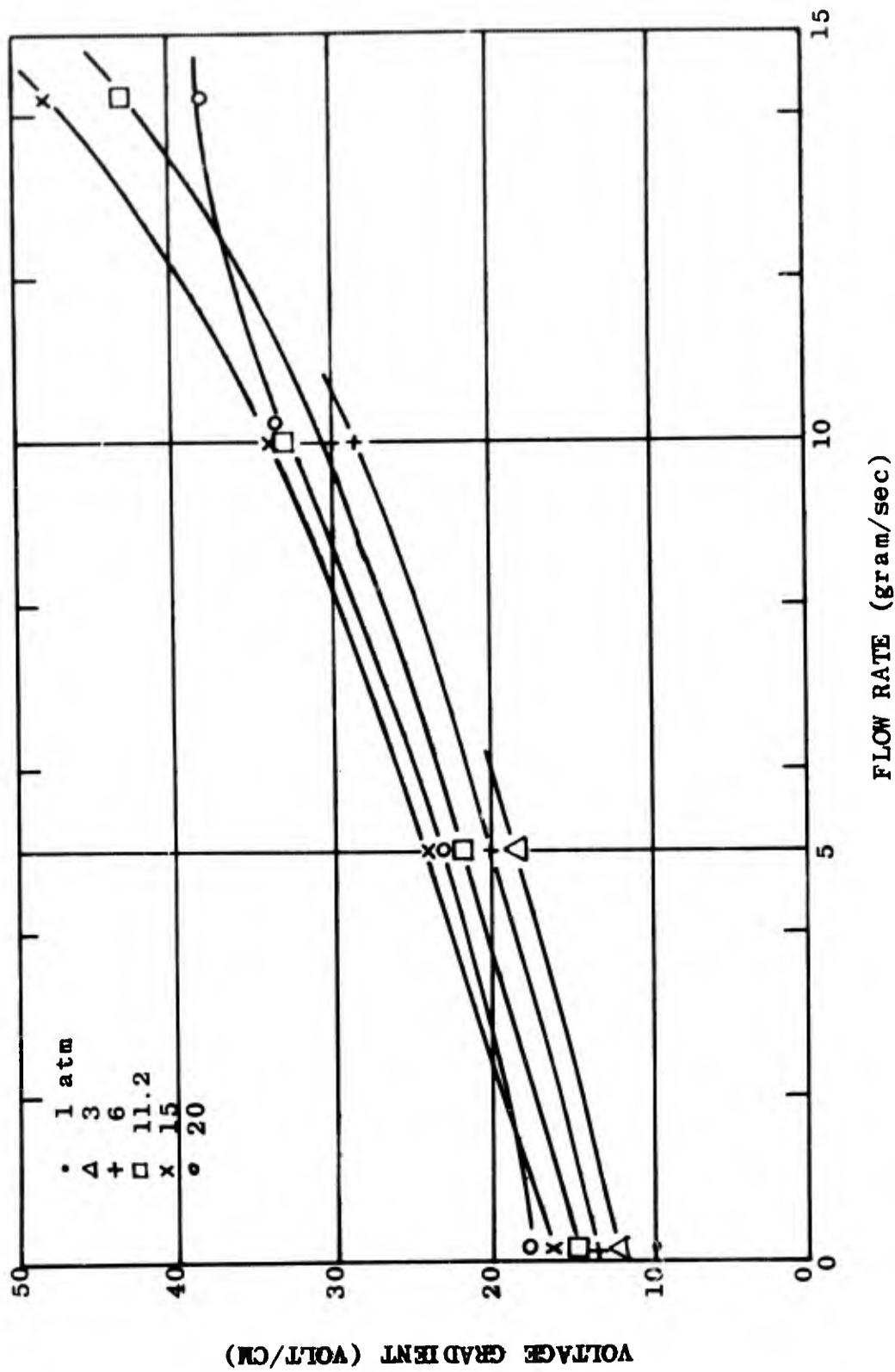


Figure 11-Voltage Gradient vs. Flow Rate in a 0.7 cm Tube, Argon  
Pressures from 1 to 20 Atmospheres

Gas: Argon  
 Pressure: 1, 3, 6, 11.2, 15, 20 atm  
 Current: 125 ± 4 amp  
 Tube Diameter: 0.7 cm  
 Tube Length: 50 cm

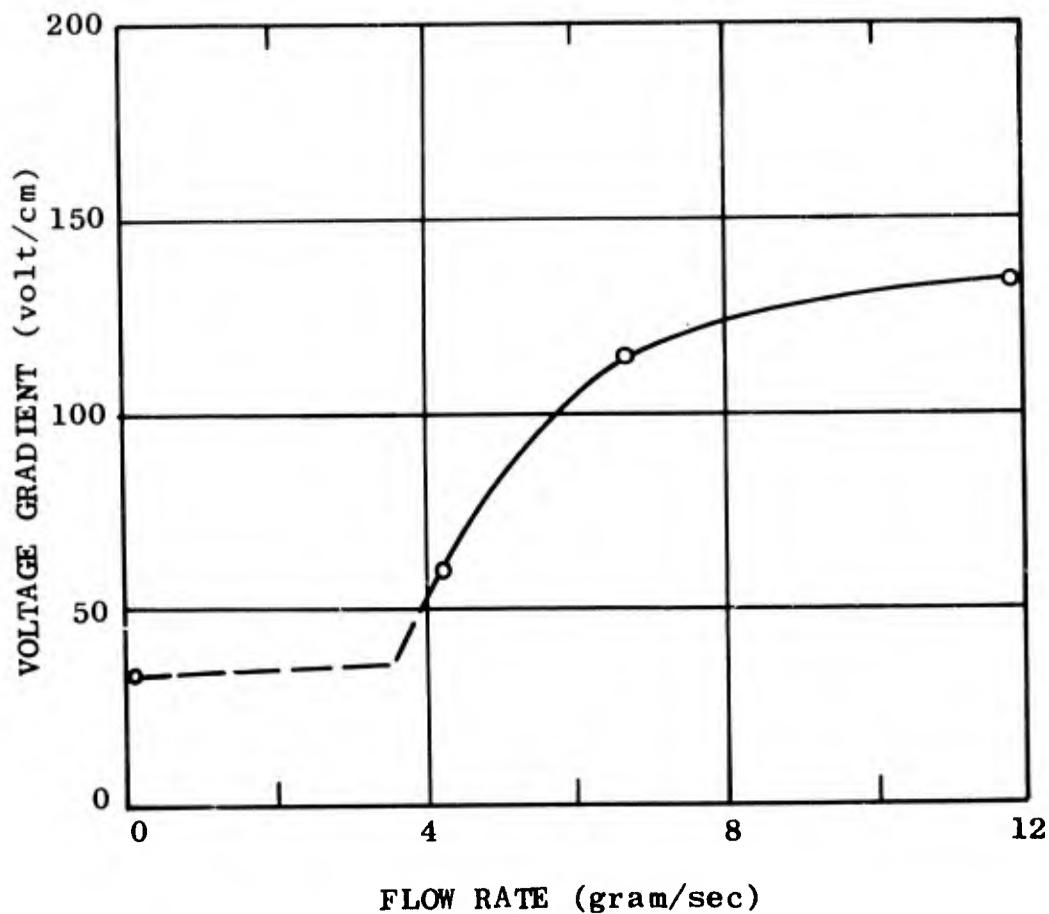


Figure 12-Voltage Gradient vs. Flow Rate at 11.2 Atmospheres of Nitrogen in a 0.7 cm Tube

Gas:	Nitrogen
Pressure:	11.2 atm
Current:	125 $\pm$ 4 amp
Tube Diameter:	0.7 cm
Tube Length:	50 cm

### 3.1312 Nitrogen

A short series of variable flow tests was made in nitrogen, again at a current level of 125 amperes. Figure 12 shows that there is a three-fold increase in voltage gradient for the range of flow rates studied, but the terminal slope of the graph is low. Obviously more measurements are necessary before firm conclusions can be drawn.

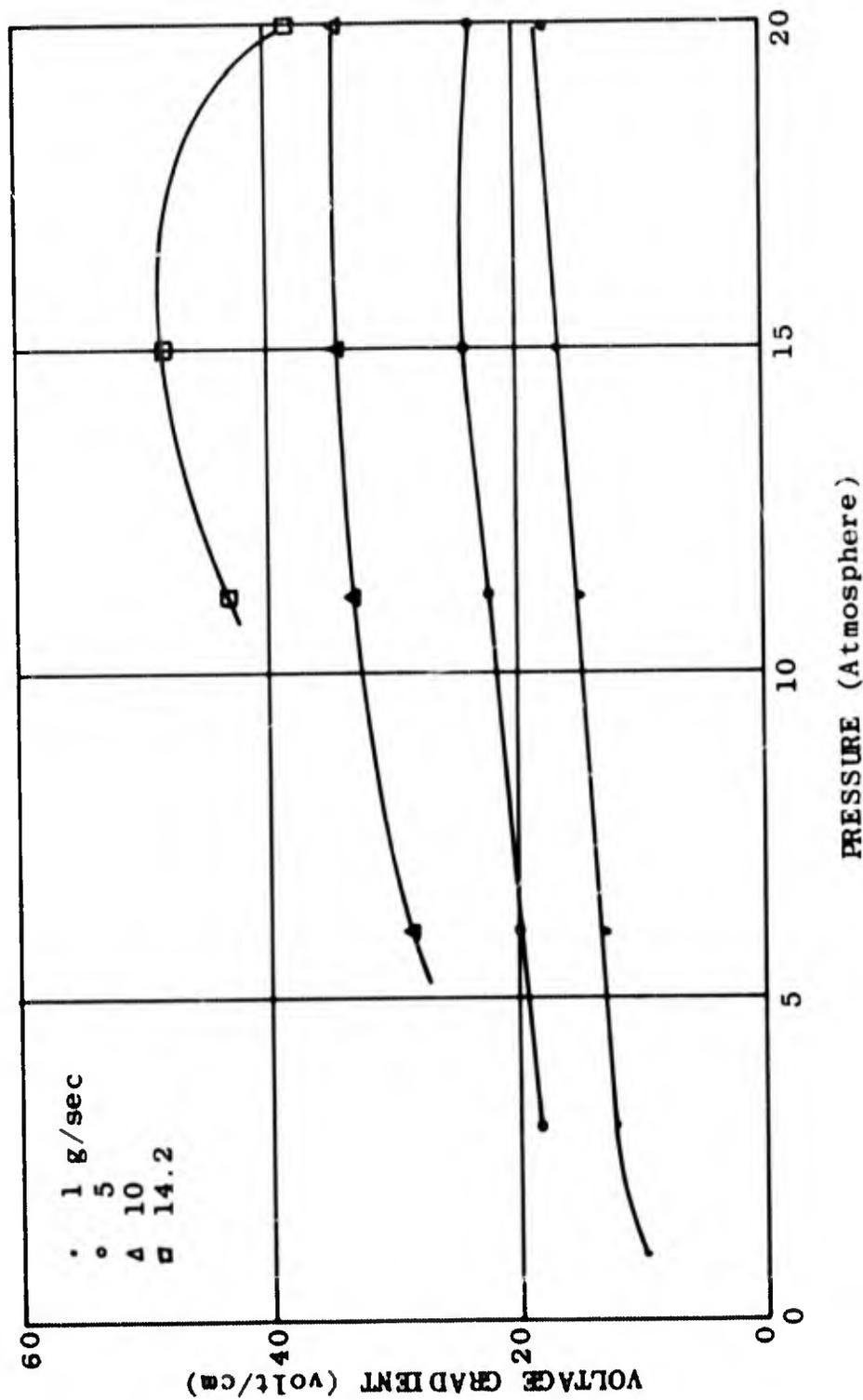


Figure 13-Voltage Gradient vs. Argon Pressure in a 0.7 cm Tube, Flow Rates from 0.1 to 14.2 gram/sec

Gas: Argon  
 Current: 125 ± 4 amp  
 Flow Rate: 0.1, 5, 10, 14.2 gram/sec  
 Tube Diameter: 0.7 cm  
 Tube Length: 50 cm

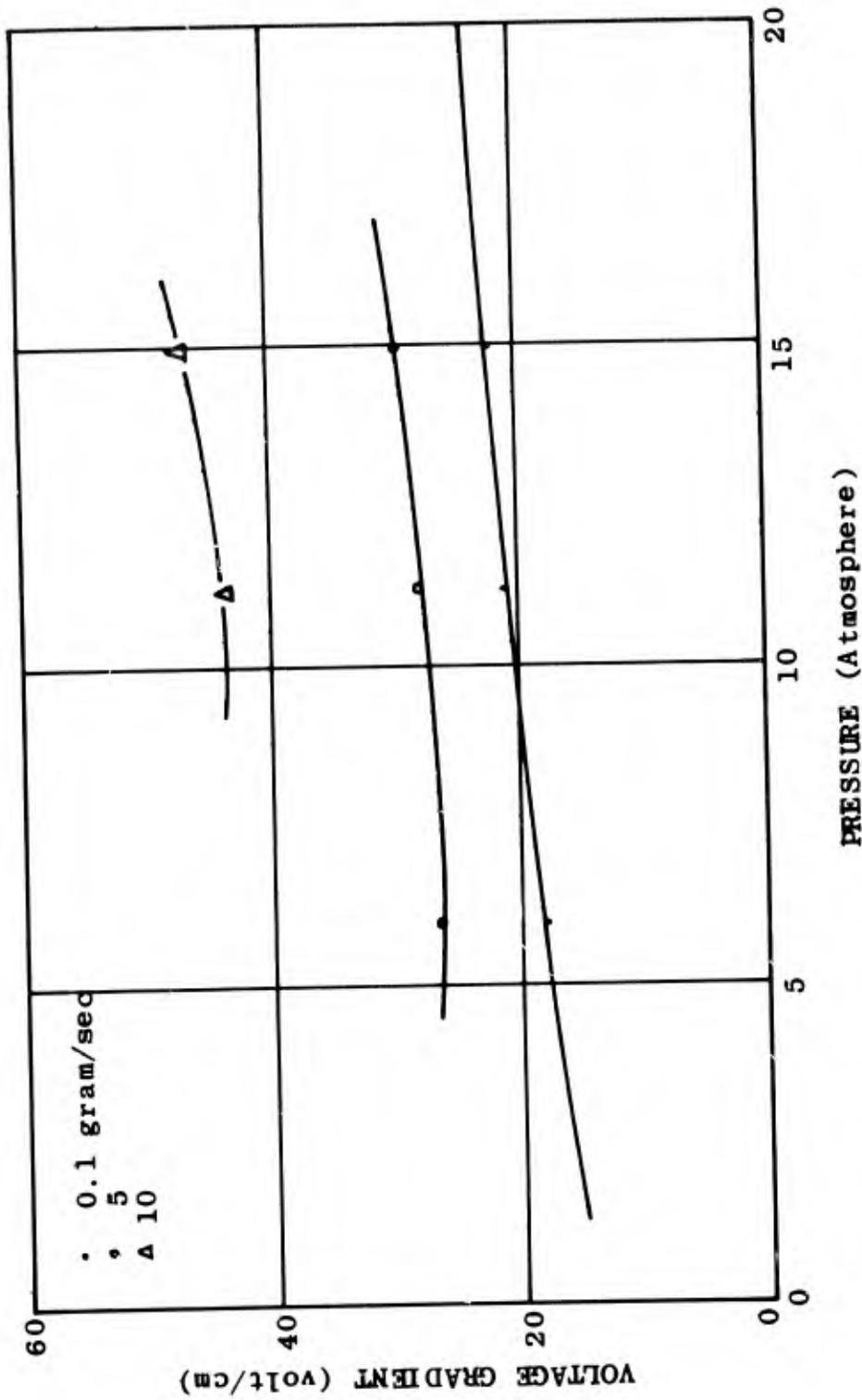


Figure 14-Voltage Gradient vs. Argon Pressure in a 0.5 cm Tube, Flow Rates from 0.1 to 10 gram/sec

Gas: Argon  
 Current: 125 ± 4 amp  
 Flow Rate: 0.1, 5, 10 gram/sec  
 Tube Diameter: 0.5 cm  
 Tube Length: 50 cm

### 3.132 Variable Pressure

#### 3.1321 Tube Diameter of 0.7 cm

The results of figure 11 are replotted in figure 13 to show the influence of pressure more clearly. Each line in the graph represents a constant flow rate result with variable ambient pressure. In addition, to the large flow effect, there is an interesting increase in the curve for low flow rate with pressure; since thermal conductivity is approximately independent of pressure, this rise must be due to increased radiation losses or to a higher core temperature required to achieve sufficient ionization. The behavior at 20 atmospheres is anomalous, especially at the highest flow rate.

#### 3.1322 Tube Diameter of 0.5 cm

Figure 14 is a repeat of the previous one with a smaller tube diameter. The laminar or low flow test series shows a higher power loss than the previous graph while the dependence of heat flux on flow rate is also greater. The same flow rate produces a greater increase in voltage gradient too. We feel this to be due to the smaller flow area which results in about twice the plasma velocity and thus in an increased Reynolds number for a given flow rate.

### 3.133 Variable Current

#### 3.1331 Argon

All previous results described in section 3.1 have been for 125 ampere arcs. It is important to determine that the effects seen at that current level exist at higher currents too. In the graphs which follow, ambient pressure is fixed and current varied; different data lines indicate different flow rates.

Figure 15 shows the results for 11.2 atmosphere argon in a 0.7 cm I. D. tube. Because a number of flow rates were studied a fairly complete picture emerges. The laminar curve shows a slow but steady increase in voltage gradient with current, as expected; the data taken with high flow rates show an initial decrease before beginning to rise again at higher currents. The initial decrease is quite probably the result of higher plasma temperatures which increase viscosity and diminish turbulence. There is no sign, however, that turbulent effects will disappear at any current reasonably extrapolated from this graph.

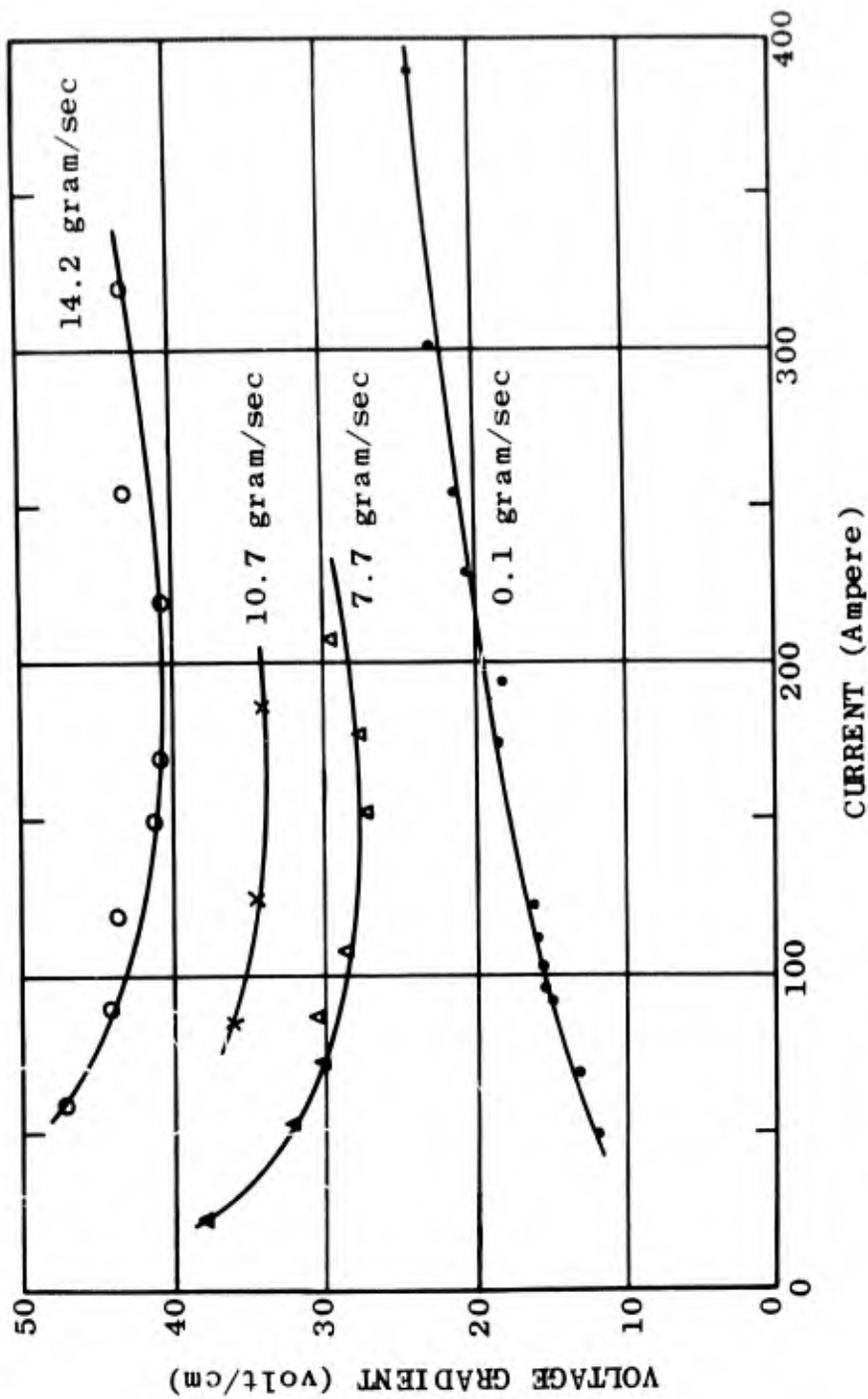


Figure 15-Characteristics of 11.2 Atmosphere Argon Arcs in a 0.7 cm Tube, Flow Rates from 0.1 to 14.2 gram/sec

Gas: Argon  
 Pressure: 11.2 atm  
 Flow Rate: 0.1, 7.7, 10.7, 14.2 gram/sec  
 Tube Diameter: 0.7 cm  
 Tube Length: 50 cm

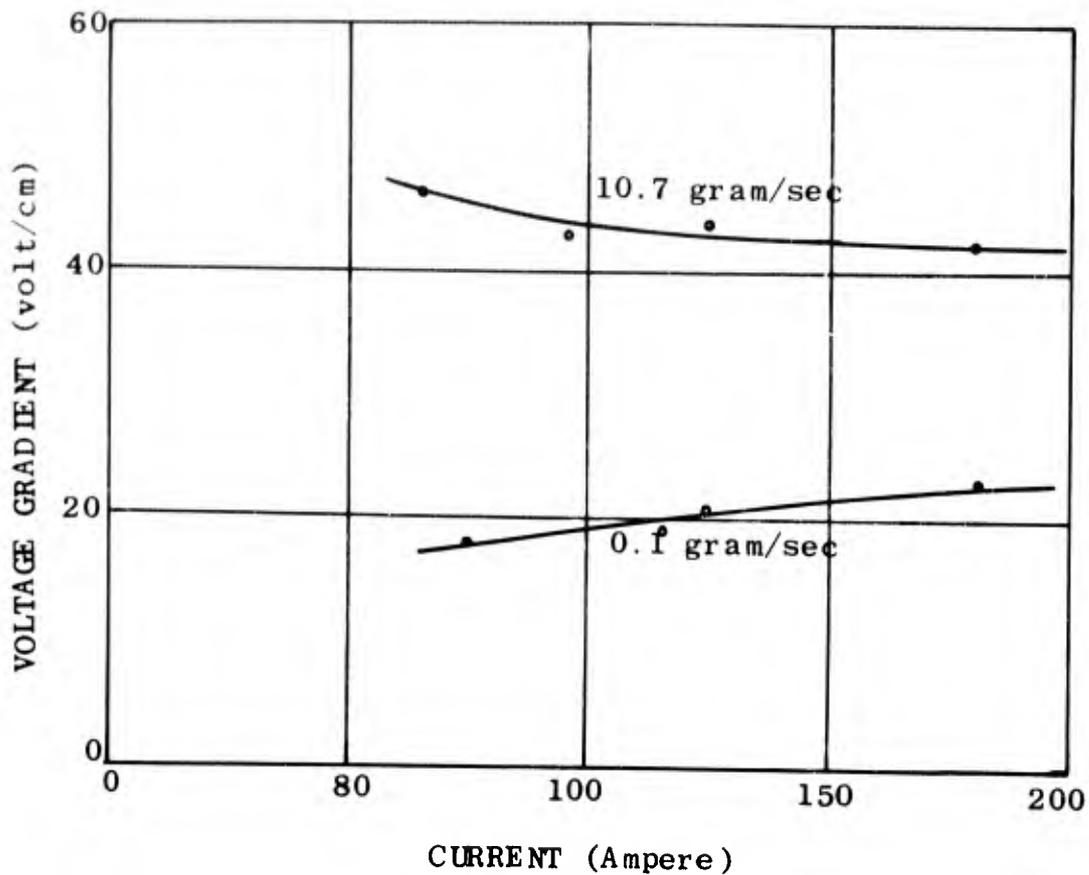


Figure 16-Characteristics of 11.2 Atmosphere Argon Arcs in a 0.5 cm Tube, Flow Rates of 0.1 and 10.7 gram/sec

Gas:	Argon
Pressure:	11.2 atm
Flow Rate:	0.1, 10.7 gram/sec
Tube Diameter:	0.7 cm
Tube Length:	50 cm

Figure 16, when compared with figure 15, shows again that the same flow rate in a smaller tube produces a higher power loss. This follows the pattern of "classical" turbulence in homogeneous fluids.

The work is extended to 15 and 20 atmospheres ambient respectively in figures 17 and 18. The independence of voltage gradient on current at 20 atmospheres is especially noteworthy. It shows no sign of declining at the highest current tested.

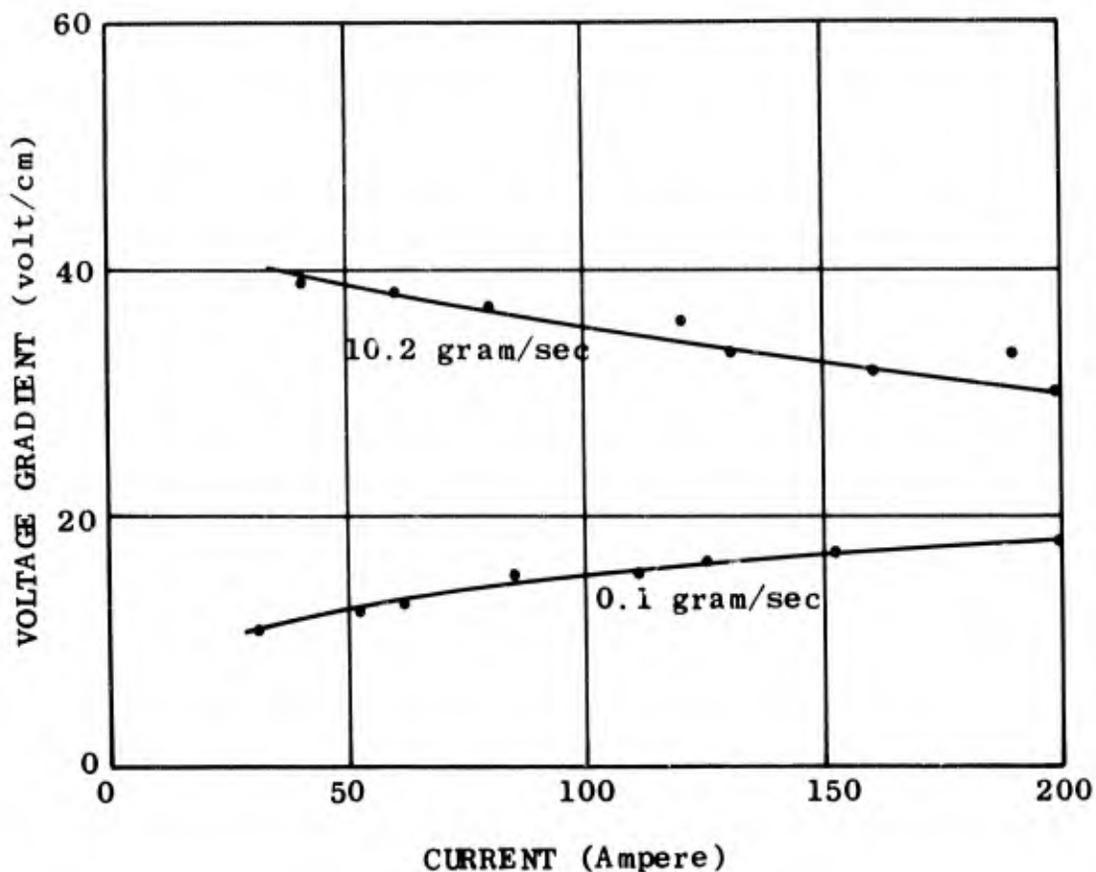


Figure 17-Characteristics of 15 Atmosphere Argon Arcs in a 0.7 cm tube, Flow Rates of 0.1 and 10.2 gram/sec

Gas: Argon  
 Pressure: 15 atm  
 Flow Rate: 0.1, 10.2 gram/sec  
 Tube Diameter: 0.7 cm  
 Tube Length: 50 cm

### 3.1332 Nitrogen

With the substitution of nitrogen for argon, figure 19 is comparable to figure 15. Four flow rates were tested over a wide range of currents. The laminar curve is only slightly dependent on current; the high flow rate curves decline markedly even above 50 amperes. The decrease in the separation of the turbulent and laminar curves at high currents is more pronounced than in argon. Notice however that the curve for the 12 gram/sec flow rate shows still more than a three fold increase over the 0.1 gram/sec value for heat transfer at the highest currents tested.

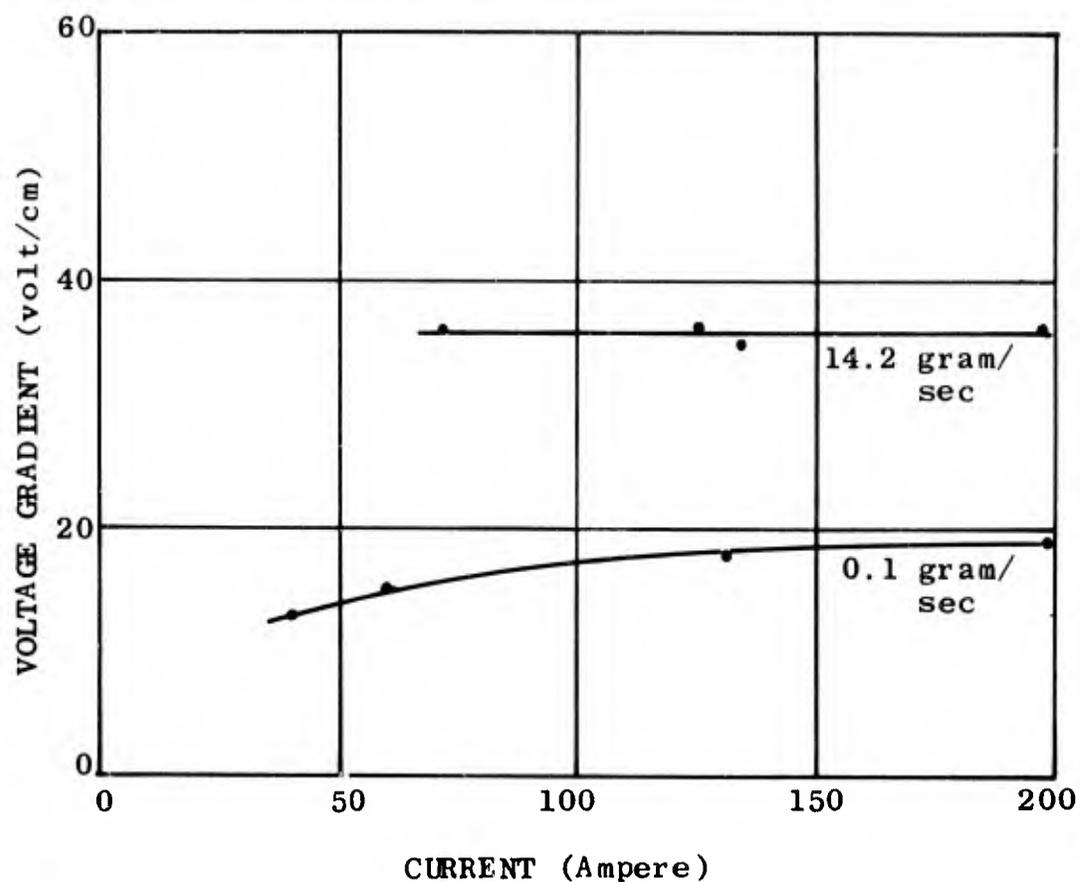


Figure 18-Characteristics of 20 Atmosphere Argon Arcs in a 0.7 cm Tube. Flow Rates of 0.1 and 14.2 gram/sec

Gas: Argon  
 Pressure: 15 atm  
 Flow Rate: 0.1, 14.2 gram/sec  
 Tube Diameter: 0.7 cm  
 Tube Length: 50 cm

### 3.134 Variable Tube Diameter

It is possible to gain an approximate idea of the influence of tube diameter on turbulent effects within the ranges of current and flow rate covered in this report. In the present investigation, tube diameters of 0.5 and 0.7 cm were used but the 1 cm tube has been studied for an earlier report (2). The

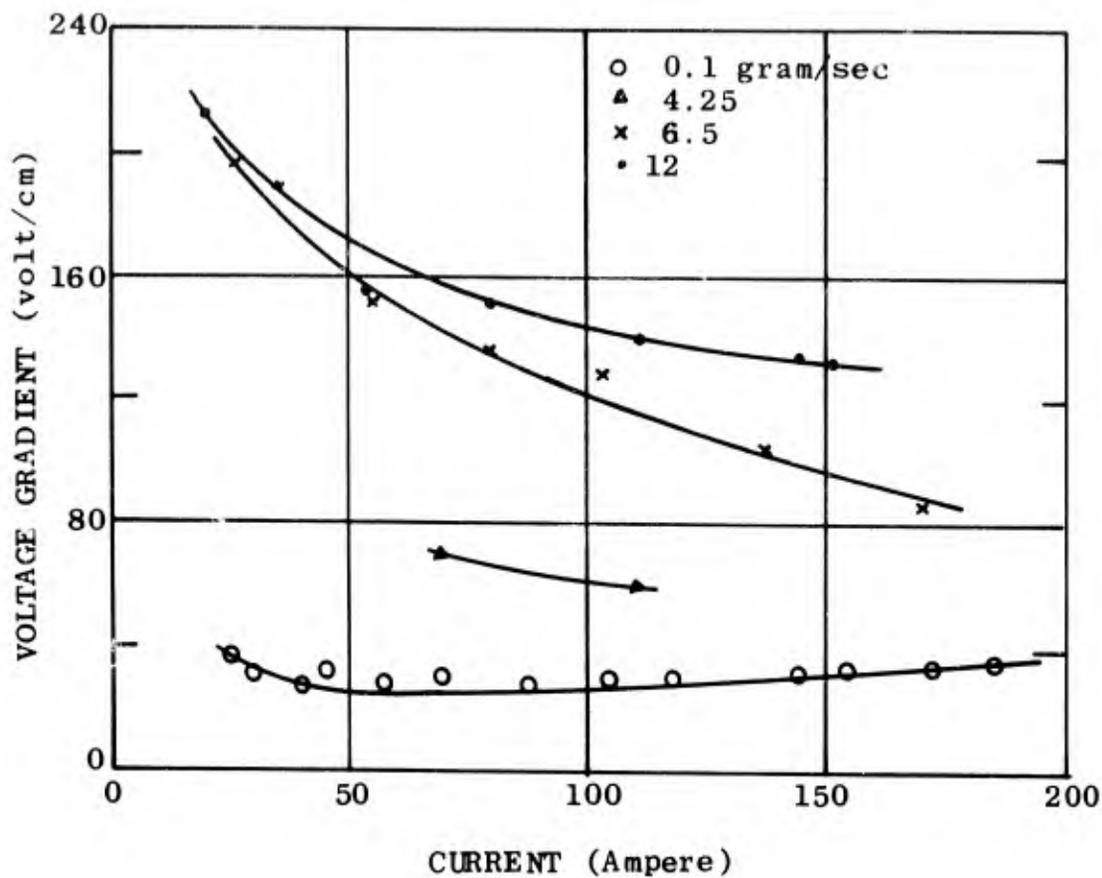


Figure 19-Characteristics of 11.2 Atmosphere Nitrogen Arcs in a 0.7 cm Tube, Flow Rates from 0.1 to 12 gram/sec

Gas:	Nitrogen
Pressure:	11.2 atm
Flow Rate:	0.1, 4.25, 6.5, 12 gram/sec
Tube Diameter:	0.7 cm
Tube Length:	50 cm

graph of figure 20 was put together by extrapolating the 1.0 cm diameter results (2) from 80 to 100 amperes. The results then do not represent direct measurements, but they should be accurate enough for a view of the major effect. Judging from figure 20, a tube 2 cm in diameter would show an appreciable increase in potential gradient only with a flow rate significantly above 15 gram/sec. while a tube 0.5 cm in diameter shows such an effect with a few gram/sec.

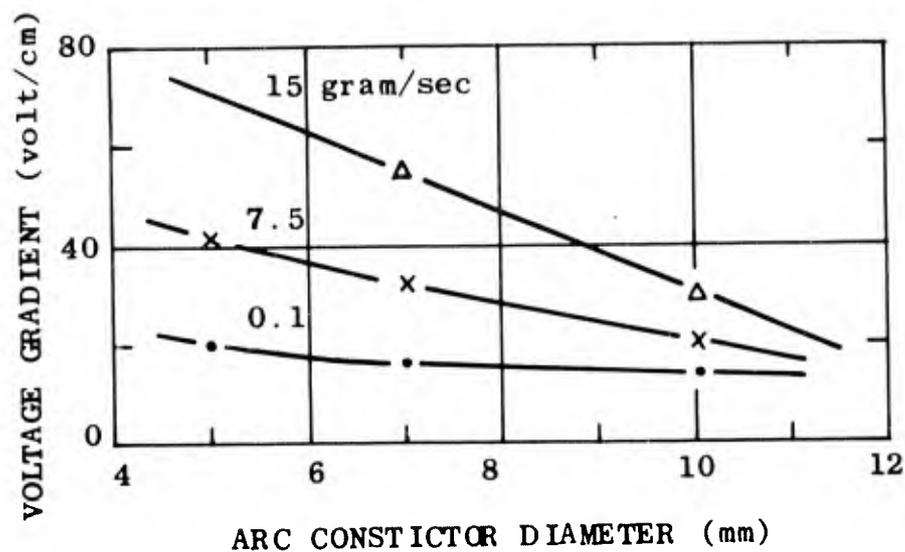


Figure 20-Voltage Gradient vs. Tube Diameter at 10 atmospheres of Argon, Flow Rates from 0.1 to 15 gram/sec

Gas: Argon  
 Flow Rate: 0.1, 7.5, 15 gram/sec  
 Pressure: 10 atm  
 Current: 100 amp  
 Tube Length: 50 and 75 cm

### 3.2 Momentum Transfer as Measured by Pressure Gradients

#### 3.21 Empirical Low Temperature Results

Earlier investigations of pressure drops along tubes with turbulent flows have resulted in experimentally known relationships for homogeneous fluids of moderate temperature. Although a comprehensive theory has not been developed we will summarize the important results below.

Taking the definition of friction factor, and rearranging the empirical relationships quoted in Rohsenow and Choi (11) we find that:

1)  $\frac{\Delta P}{\Delta L} \propto [Re]$  in laminar flow

and

2)  $\frac{\Delta P}{\Delta L} \propto \frac{\mu^2}{\rho} [Re]^{1.8}$  in turbulent flow

(D held constant)

but,

3)  $\frac{\Delta P}{\Delta L} \propto \frac{\mu^2}{\rho} [Re]^2$  in far turbulent flow

(D held constant)

where P = pressure  
 L = length of tube  
 $\mu$  = viscosity  
 $\rho$  = density  
 and D = tube diameter

The point of transition from 2) to 3) depends on the surface roughness ratio e/D for the tube wall which is quite high in our case. (This "rough" wall, e/D = 0.036, is the result of the 0.25 mm thick, recessed teflon spacers which isolate the individual copper discs of the cascade tube). The graph of (11) indicates that, in our tube, equation 3) should hold for  $Re_D \geq 30,000$ . It must be emphasized that the equations quoted above and the result for wall roughness are empirical—they satisfy data from experiments conducted on homogeneous fluids near room temperature.

### 3.22 Methods

In an earlier report (3) we presented the results of pressure gradient measurements on 125 ampere arcs at 11.2 atmospheres. The sensor response, the measuring technique, and the accuracy of this system were discussed there. Briefly, a quartz pressure-charge transducer is mounted at the end of a thin, electrically insulating tube which runs through a copper disc to the arc channel. The dead volume before the transducer is minimized so the response time of the system is held to a few milliseconds. Because of this limit we miss the high frequency fluctuations but are able to read the average pressure well. Measurements were made at 5 points along the tube

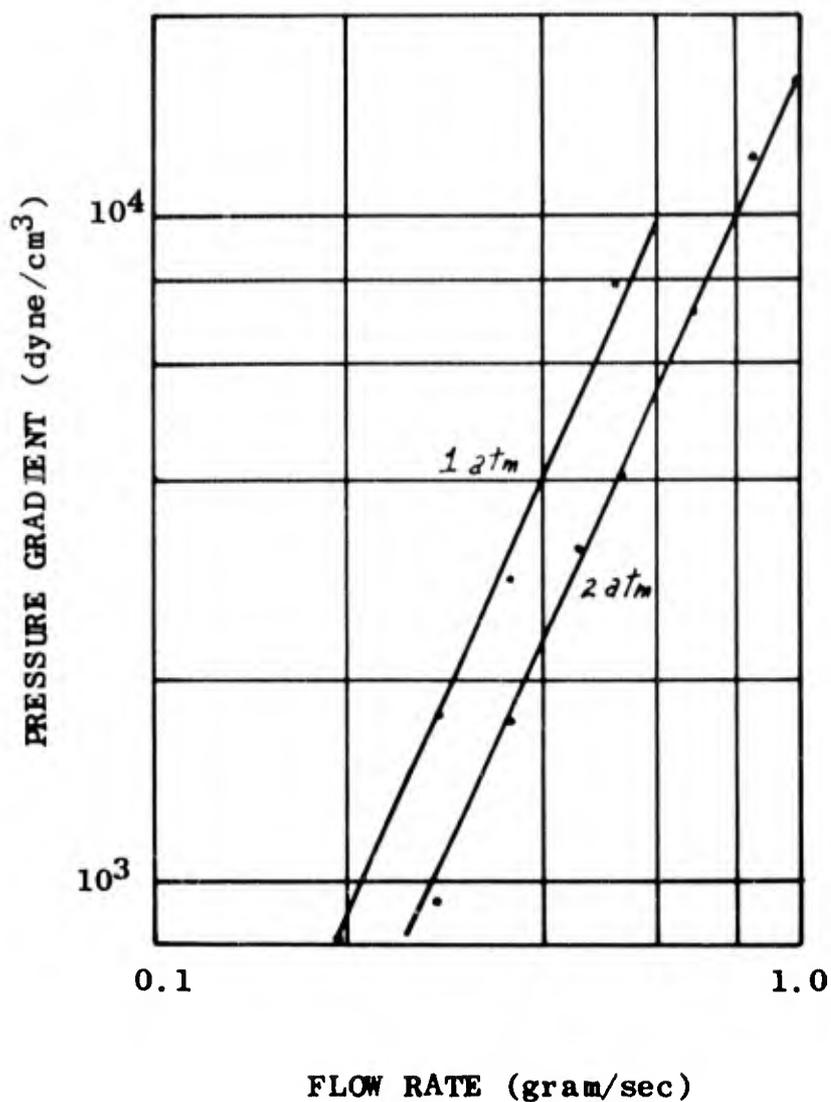


Figure 21-Pressure Gradient vs. Flow Rate, Comparing 1 and 2 Atmospheres of Argon without Arc

Gas:	Argon
Pressure:	1, 2 atm
Current:	None
Tube Diameter:	0.7 cm
Tube Length:	50 cm

simultaneously to establish the gradient. Preference was given to the fully developed flow region but no strong departures from linearity were seen in any segment of the tube. The current value of 125 amperes was chosen because the pressure gradient and plasma velocity were found to be largely independent of current in that area; it was also a convenient value for our apparatus.

From a log-log plot of our data we found a constant 1.8 power dependence of pressure gradient on flow rate, an indication that the flows were in the turbulent area, equation 2). While it would be desirable to extend these 11.2 atmosphere measurements to lower flow rates so the laminar region could be explored, the sensitivity of our instrumentation sets a limit. Consequently the ambient arc pressure was lowered to permit measurement of the pressure gradient at lower flow rates for which the flow is laminar. Equation 2) shows that the pressure gradient goes as  $1/\nu$  for a constant Reynolds number or, since:

$$4) \quad Re_D = \frac{4W}{\pi D \mu}$$

for constant flow rate, diameter, and viscosity ( $W$ ,  $D$ , and  $\mu$  respectively).

This variation of pressure gradient with ambient pressure can be seen in figure 21. The data was taken in our cascade tube with no arc burning. For a fixed flow rate, the two atmosphere case shows about half the pressure drop of the one atmosphere case.

### 3.23 Measurements in Hot Flow

A series of measurements with the arc burning is shown in figure 22, which is a log-log plot of pressure gradients against flow rates at ambient pressures of 1, 2, 3, and 11.2 atmospheres. The gas was argon in all cases; the arc current varied from 112 to 129 amperes. Since it was established in (3) that pressure gradients changed only slowly with currents above 100 amperes, the small current variation is negligible.

The major result from figure 22 is the expected transition from a laminar to turbulent dependence of pressure gradient on flow rate. Specifically the slope of the curve changes abruptly from 1.0 to approximately 1.8 which is the behavior predicted by equations 1) and 2). This transition occurs in the neighborhood of 1.5 gram/second, a point which agrees well with the one atmosphere experiments of Runstadler (10). His curve of friction factors, however, shows only marginal evidence of turbulence; by increasing pressure above one atmosphere we have shown both the laminar and the clearly turbulent branches of the curve.

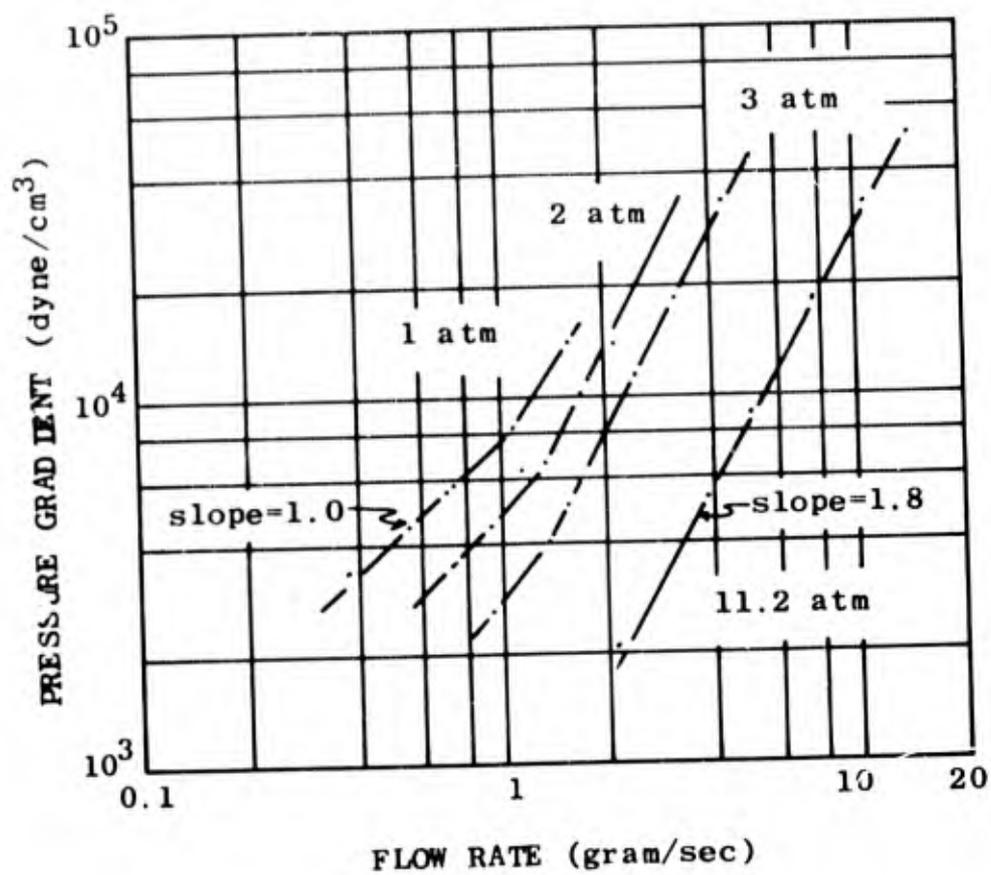


Figure 22-Pressure Gradient vs. Flow Rate, Argon  
Pressure from 1 to 11.2 Atmospheres

Gas:	Argon
Pressure:	1, 2, 4, 11.2 atm
Current:	125, +4, -13 amp
Tube Diameter:	0.7 cm
Tube Length:	50 cm

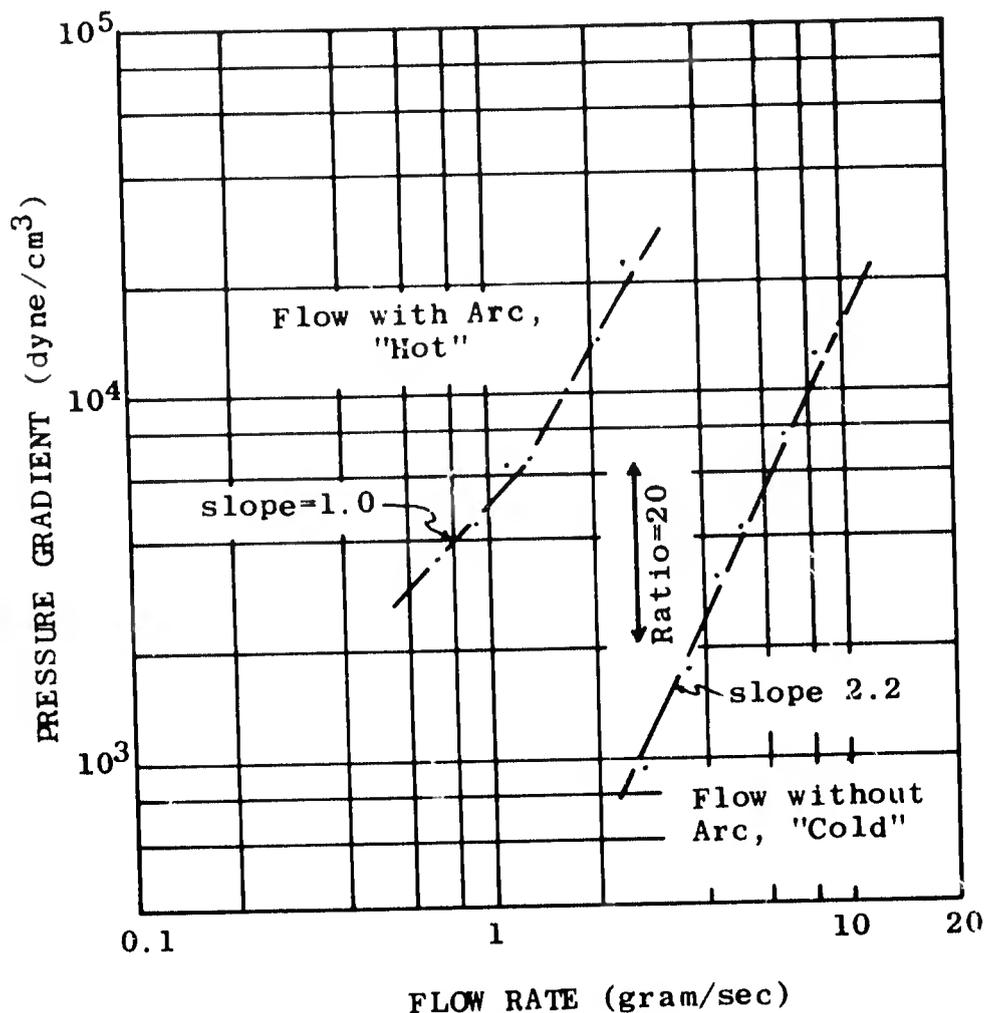


Figure 23-Pressure Gradient vs. Flow Rate, 2 Atmosphere Argon with and without an arc.

Gas:	Argon
Pressure:	2 atm
Current:	125 amp, none
Tube Diameter:	0.7 cm
Tube Length:	50 cm

The large effect which the presence of the arc has on the pressure gradient in the tube is evident from figure 23. At 3 gram/sec flow rate, the only area of overlap for the two curves, the ratio of the two pressure drops was found to be 20:1. But this measurement is at constant flow rate, not constant  $Re$ . It has been argued (10) that the proper choice of  $Re$  to describe the flow through an arc is the room temperature value. The graph above shows that, insofar as the pressure gradient is concerned, this selection is incorrect. The  $Re$  for this tube at 1.5 gram/sec is 13,000 if computed from room temperature fluid properties. The corresponding number for a 12,000°K model is 800. If one uses a more sophisticated choice of reference temperature, the average enthalpy value suggested by Eckert (9), the  $Re$  at 1.5 gram/sec is 1,400.

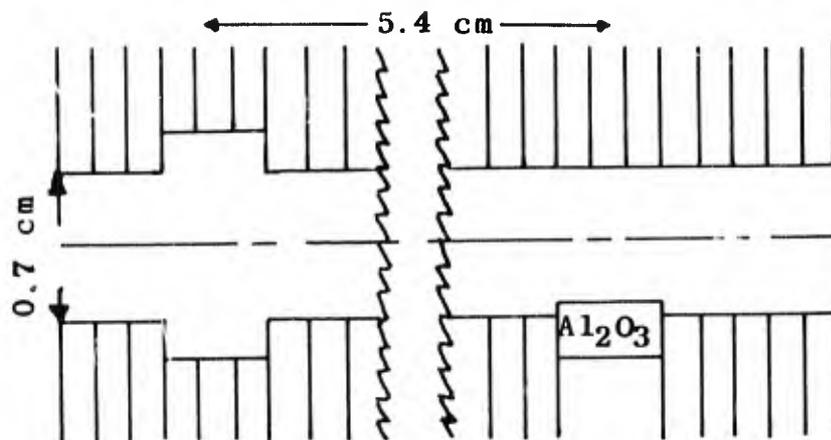


Figure 24-Diagram of Doping Method Used for Mass Transport Study. Gas flow is from right to left.

### 3.3 Mass Transfer as Measured by Radial Diffusion of Atoms

#### 3.31 Method

One of the distinguishing attributes of turbulent flow is the existence of macroscopic flow perpendicular to the direction of major flow. It was therefore of great interest for our investigation to make some direct tests on turbulent mass transfer in addition to the experiments on heat and momentum transfer. Our stigmatic Ebert spectrograph enables us to study such motion and we have utilized that ability. The diffusion of a dopant species in the arc can be compared when flow rates are high and low, thus giving an idea of the rate at which turbulent mixing proceeds compared with the rate of unassisted molecular diffusion.

In this study, dopant material was deliberately introduced, not symmetrically from the electrode, but from one side of the tube so that its diffusion toward an even concentration over the tube cross-section could be observed. This was done by placing a crystal of  $\text{Al}_2\text{O}_3$  at the edge of the arc and allowing it to ablate. Spectral observation was made through a window located 5.4 cm downstream from the point of introduction, (figure 24). Two "near resonance" Al I lines,  $\lambda = 3944 \text{ \AA}$  and  $3961.5 \text{ \AA}$ , are prominent enough for use even at low partial pressures.

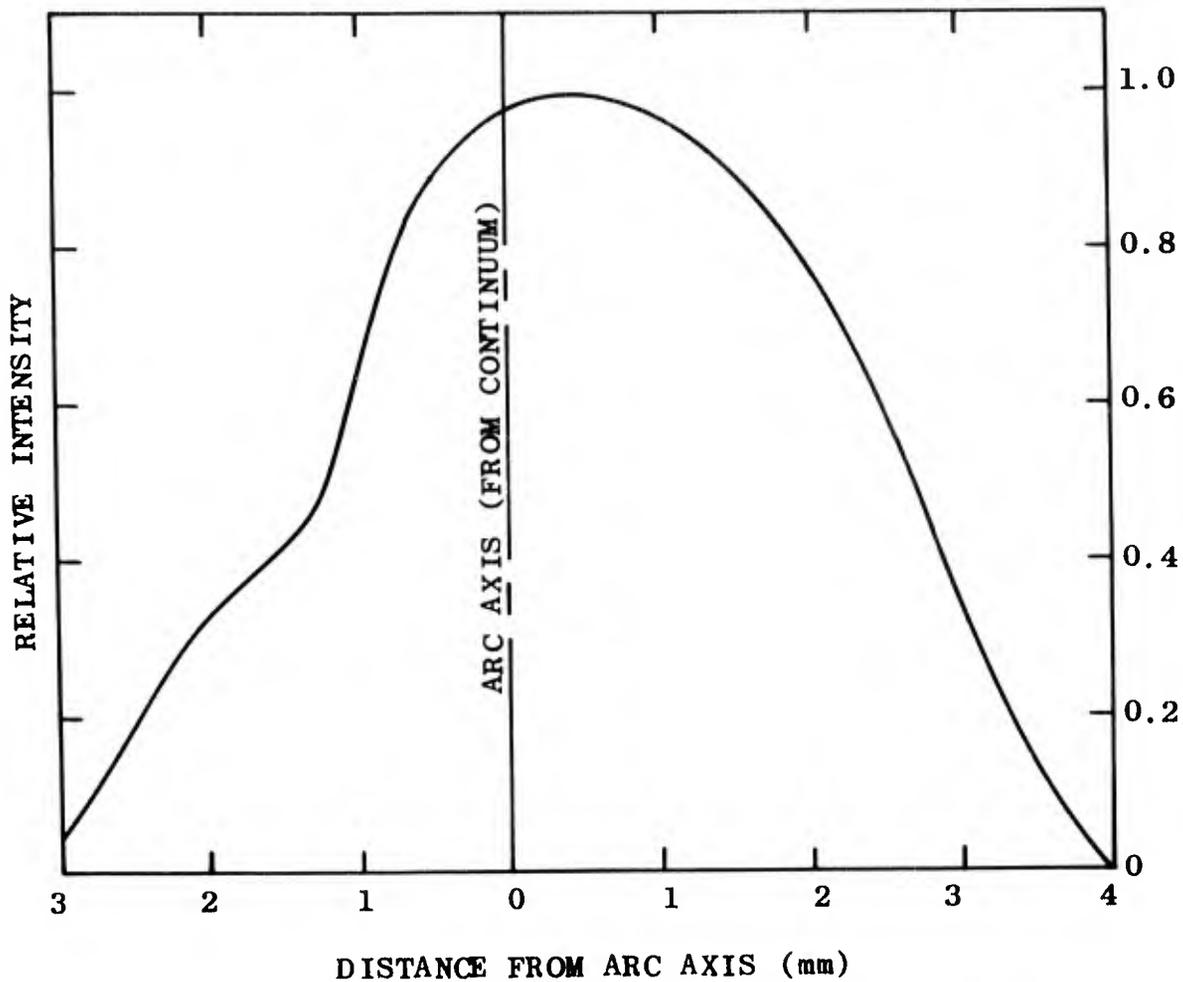


Figure 25-Line Intensity (Al I  $\lambda = 3961.5 \text{ \AA}$ )  
 Observed 5.4 cm Downstream from Introduction  
 Point; Flow Rate 0.3 gram/sec. The Aluminum  
 was introduced from the right side as seen  
 here.

Gas:	Argon
Pressure:	11.2 atm
Current:	50 ampere
Flow Rate:	0.3 gram/sec
Tube Diameter:	0.7 cm
Dopant Introduction Point:	41 cm from cathode
Window:	46.4 cm from cathode
Tube Length:	50 cm
Test #:	458

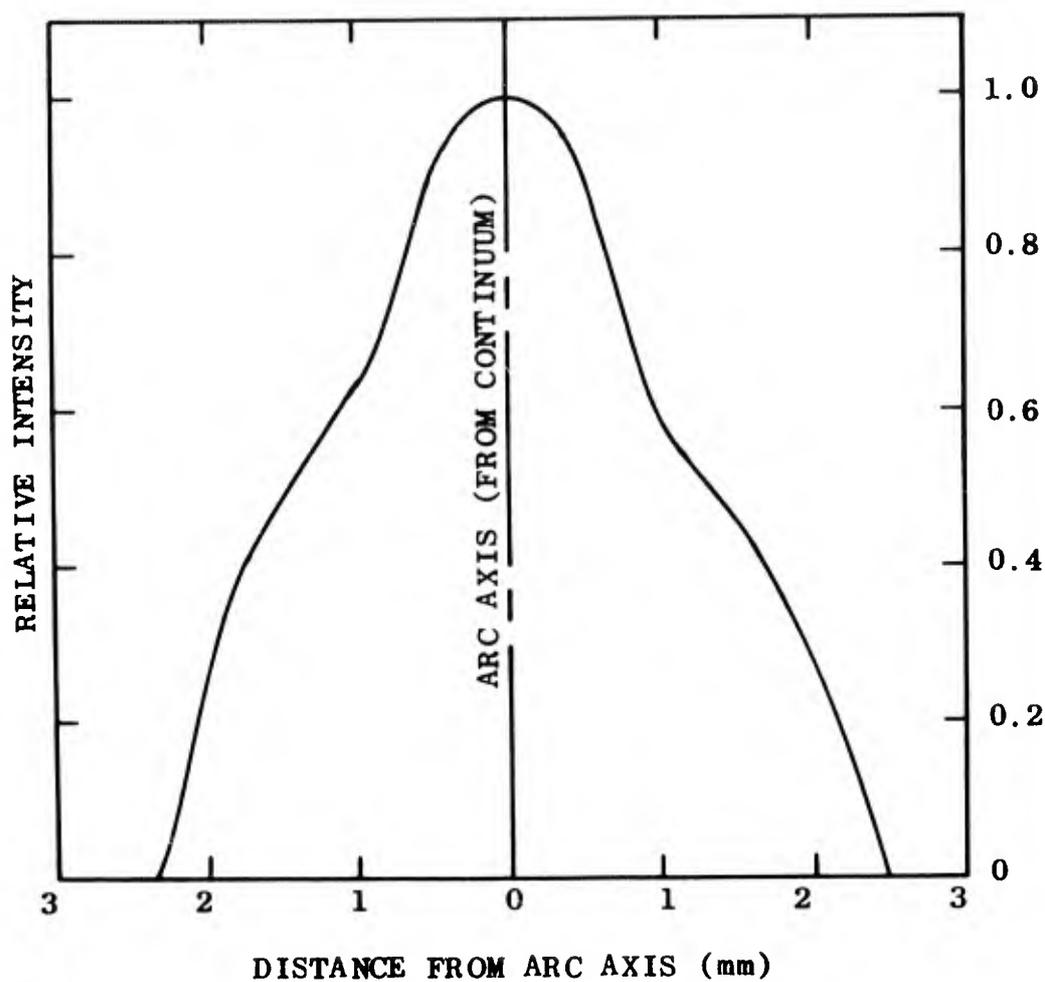


Figure 26-Line Intensity (Al I  $\lambda = 3961.5 \text{ \AA}$ )  
 Observed 5.4 cm Downstream from Intro-  
 duction Point, Flow Rate 5 gram/sec.  
 The aluminum was introduced from the  
 right side as seen here.

Gas:	Argon
Pressure:	11.2 atm
Current:	50 amp
Flow Rate:	5 gram/sec
Tube Diameter:	0.7 cm
Dopant Introduction Point:	41 cm from cathode
Window:	46.4 cm from cathode
Tube Length:	50 cm
Test #:	459

### 3.32 Results

Figures 25 and 26 show the intensities of the 3961.5 Å line in argon arcs of 50 amperes and 11.2 atm. with flow rates of 0.3 gram/sec. and 5 gram/sec. respectively.

These intensities are observed side-on and therefore represent values integrated through the varying depth of the arc, ( $I(x)$  values). The data are presented in the way because the graph of figure 25 is obviously strongly asymmetric and an Abel inversion cannot be performed without a much more sophisticated measurement as done by Olsen (12).

The arc axes marked on figures 25 and 26 were determined from the centers of the continuous radiation distribution which, significantly, was quite symmetric in both cases. In the high flow rate example, the maximum of the Al line intensity coincides with the maximum of the continuous radiation, as expected for a symmetrical distribution. But, for the laminar flow case, the line intensity distribution is strongly asymmetric, indicative of the asymmetric concentration of aluminum atoms.

In interpreting figs. 25 and 26 we feel that, in the laminar case, the time lapse between introduction of the aluminum and examination of the spectrum was obviously too short to achieve a symmetric distribution of the dopant species through the arc by molecular diffusion alone.

### 3.33 Discussion

In evaluating our observations of the last section, it is important to know whether the temperature distribution in these experiments was axially symmetric. If this were not so, small changes in temperatures might account for the observed line asymmetry rather than differences in the density of the aluminum atoms.

The arguments for accepting an unaltered temperature distribution and unaltered electrical and thermal conditions in the arc are:

1. The continuous radiation (which would also be quite sensitive to small changes in temperature) was symmetric.
2. Only a small percentage of the arc plasma is aluminum. Partial pressure of aluminum was estimated to be 50 dynes/cm<sup>2</sup> from an absolute measurement of the  $\lambda = 3944 \text{ \AA}$  line whose transition probability is known.
3. The arc voltage gradient remained unchanged when compared with tests containing no aluminum.

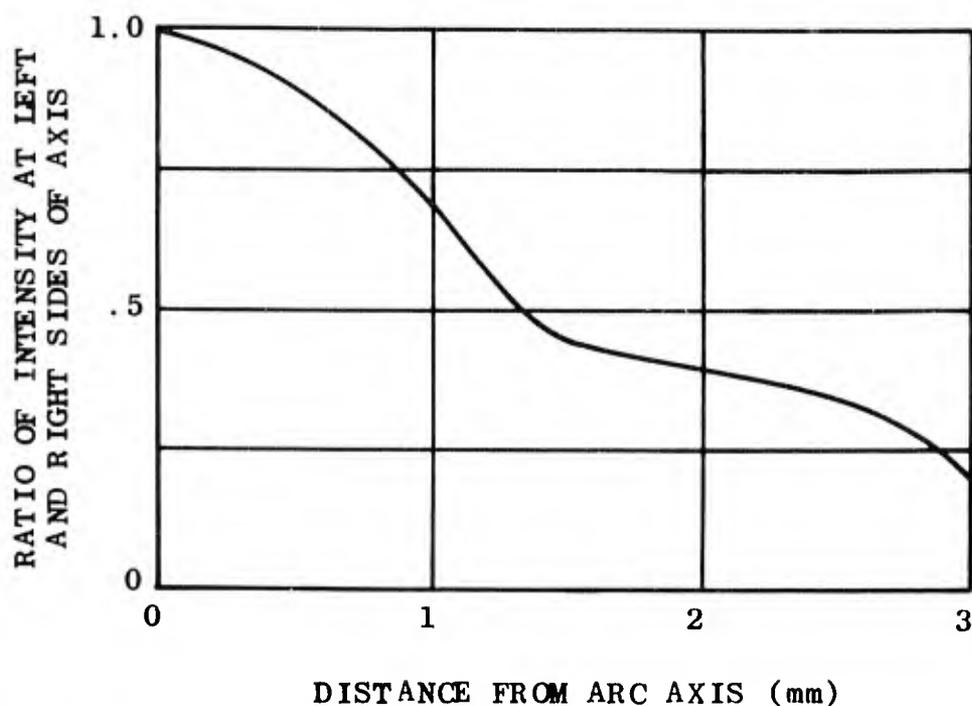


Figure 27-Ratio of Intensities at Two Sides of Arc Axis, Based on Figure 25. The aluminum distribution is clearly asymmetric. 0.3 gram/sec flow rate

To show the distribution of aluminum more clearly, we plot in figure 27 and 28 the intensity ratios at successive distances from the arc axis. The 5 gram/sec. test is essentially symmetric, Fig. 28 while the 0.3 gram/sec. test shows an ever increasing disparity as the edges of the plasma are approached, fig. 27.

It would, of course, be valuable if diffusion velocities could be derived from our experiments. However, an exact measure of diffusion velocities was not accomplished. Still, there are firm conclusions which can be drawn.

Figures 25 and 26 show clearly that diffusion is much more uniform at the higher flow rate, but a consideration of the axial flow velocities makes the comparison even more startling. If we take from our earlier plasma velocity measurements (3) average axial velocities of 9 meter/sec. and 150 meter/sec. for the two tests, the two respective diffusion times between introduction of dopant and measurement become 6 msec. and 360  $\mu$ sec. which differ by a factor of 15. Since the aluminum distribution is much more uniform in the high flow case, we can conclude that turbulence has increased the rate of diffusion of aluminum in argon at least a factor of 15.

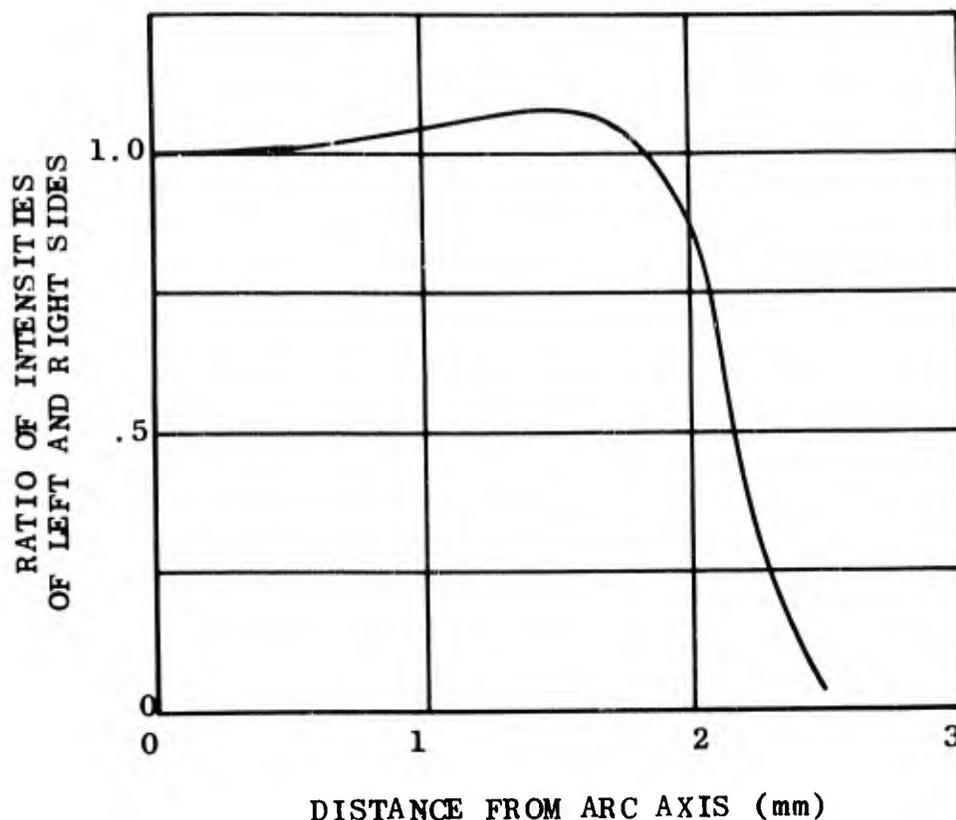


Figure 28-Ratio of Intensities at Two Sides of Arc Axis, Based on Figure 26. The aluminum distribution is symmetric except at the very edges. 5.0 gram/sec flow rate

### 3.4 Fluctuations of Light Intensity

#### 3.41 Method

Our earlier insight into the local three dimensional structure of turbulent arcs was mainly based on direct high speed photography of arcs (2, 3). Whereas the movies provided a striking confirmation of the non laminar character of high flow rate arcs, a quantitative measure of light fluctuations cannot be conveniently taken from such records.

Besides being a measure of plasma turbulence, a determination of the amplitude and frequency distribution of the light fluctuations is also important for a judgement of the accuracy of temperature measurements. As reported earlier (3), the critical problem for turbulent temperature measurements is that light intensity is linear with temperature only in restricted temperature ranges. If the limits of this linear range are exceeded, a mistake is introduced when the intensity average is accepted as a measure of the temperature average.

To measure light fluctuations a magnified image (x8) of the arc was projected on a photo-diode.\*

The light window was 45 cm downstream from the cathode in a 50 cm tube of 0.7 cm diameter. An f2.5 lens with 7.5 cm diameter was used for this purpose. Pertinent specifications of the photo-diode were:

Sensitivity:	0.5 amp/watt at 9,000Å
Quantum Efficiency:	70%
Rise Time:	5 x 10 <sup>-8</sup> or better
Linear Range:	7 decades
Spectral Sensitivity Range:	3dB points at 3,500Å and 11,300Å
Active Area:	5.1 mm <sup>2</sup>

The rise time includes a probe and the CRO: it was checked using a fast solid state laser with 9,000 Å output. Because of the magnification, the diode can resolve approximately 0.25 mm in the arc. The response of the diode, however, represents an integrated value running side on through the arc. The sensor was placed in the center of the arc image and so integrated light from a diameter which included the arc core.

In the tests of the following section, the CRO trace displaying the photo-diode response was delayed about 200 msec after arc ignition to allow time to establish steady state conditions. The trace then sweeps the screen rapidly providing a short sample (typically 500 μsec) of the fluctuations. Although repeat tests were made to establish reproducibility they proved to be redundant and are not shown here.



Gas:	Argon
Pressure:	11.2 atm
Current:	53 amp
Flow Rate:	3.75 gram/sec
Tube Diameter:	0.7 cm
Window:	45 cm from cathode
Tube Length:	50 cm
Test #:	468

Zero Intensity

1 msec

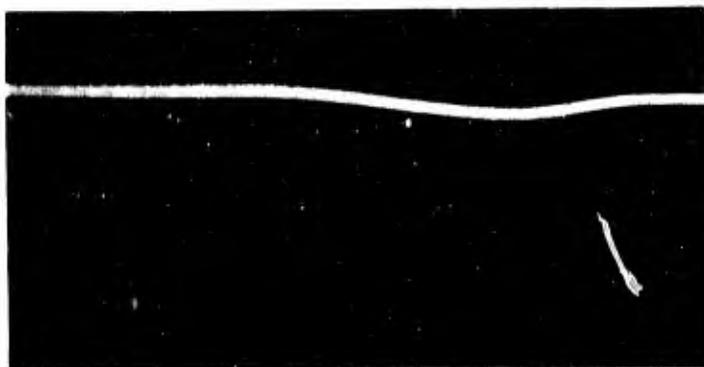
Figure 29-Light Fluctuations of a 50 Ampere Argon Arc with 3.75 gram/sec Flow Rate, Slow Sweep

\* The photo-diode was a SGD-100 silicon diffused model of Edgerton, Grier, and Germershausen, Boston, Mass.

### 3.42 Results

For a view of the slow fluctuations the oscilloscope was set to cover a 5 msec period with figure 29 as a result. The light comes from a 50 ampere argon arc in a 0.7 cm I. D. constrictor. The ambient pressure was 11.2 atm. and the flow rate 3.75 gram/sec.

From this figure one can see that, in the fully developed region of a high flow rate arc, there is a impressive level of fluctuations. The general nature of the trace is irregular with a mixture of frequencies. No signs of "snaking" of the arc column, which would create slow, regular swings in the average intensity are visible even at this low current.

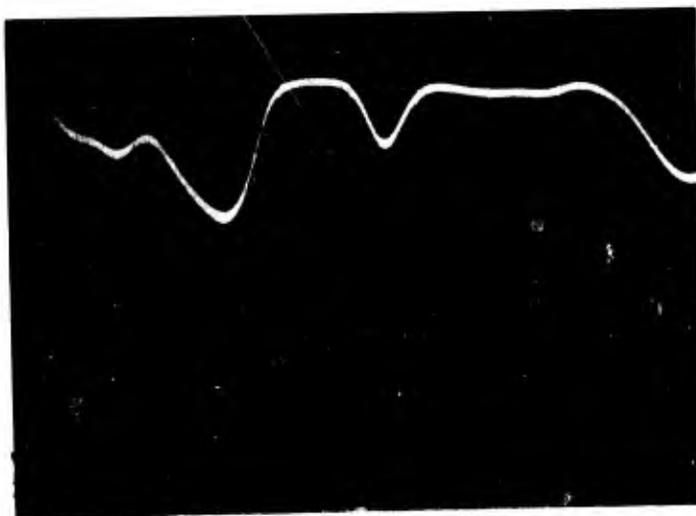


Gas: Argon  
Pressure: 11.2 atm  
Current: 54 amp  
Flow Rate: 0.3 gram/sec  
Tube Diameter: 0.7 cm  
Window: 45 cm from cathode  
Tube Length: 50 cm  
Test #: 476

Zero Intensity

100 $\mu$ sec

Figure 30-Light Fluctuations of a 50 Ampere Argon Arc with 0.3 gram/sec Flow Rate, Fast Sweep

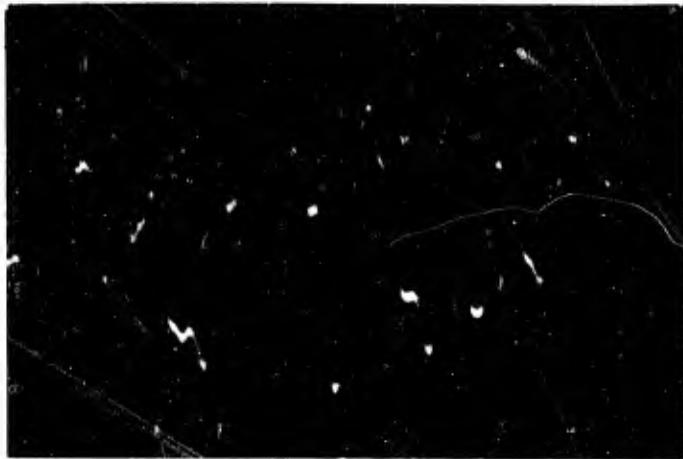


Gas: Argon  
Pressure: 11.2 atm  
Current: 53 amp  
Flow Rate: 1.1 gram/sec  
Tube Diameter: 0.7 cm  
Window: 45 cm from cathode  
Tube Length: 50 cm  
Test #: 475

Zero Intensity

100 $\mu$ sec

Figure 31-Light Fluctuations of a 50 Ampere Argon Arc with 1.1 gram/sec Flow Rate, Fast Sweep



Gas: Argon  
Pressure: 11.2 atm  
Current: 53 amp  
Flow Rate: 3.75 gram/sec  
Tube Diameter: 0.7 cm  
Window: 45 cm from cathode  
Tube Length: 50 cm  
Test #: 471

Zero Intensity

100 $\mu$ sec

Figure 32-Light Fluctuations of a 50 Ampere Argon Arc with 3.75 gram/sec Flow Rate, Fast Sweep



Gas: Argon  
Pressure: 11.2 atm  
Current: 53 amp  
Flow Rate: 10.8 gram/sec  
Tube Diameter: 0.7 cm  
Window: 45 cm from cathode  
Tube Length: 50 cm  
Test #: 474

Zero Intensity

100 $\mu$ sec

Figure 33-Light Fluctuations of a 50 Ampere Argon Arc with 10.8 gram/sec Flow Rate, Fast Sweep



Gas: Argon  
 Pressure: 11.2 atm  
 Current: 53 amp  
 Flow Rate: 15.5 gram/sec  
 Tube Diameter: 0.7 cm  
 Window: 45 cm from cathode  
 Tube Length: 50 cm  
 Test #: 473

Zero Intensity

100 $\mu$ sec

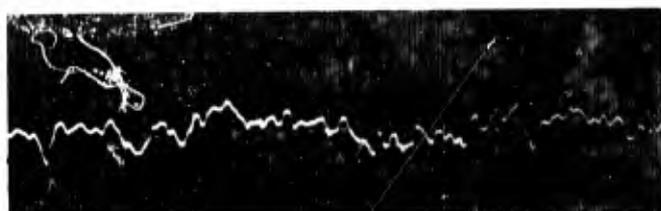
Figure 34-Light Fluctuations of a 50 Ampere Argon Arc with 15.5 gram/sec Flow Rate, Fast Sweep

For a correlation of figure 29 with our other measurements recall that the pressure gradient is on the turbulent branch of the curve (figure 22) and the heat transfer has nearly doubled over its laminar value (figure 15).

The series of figures 30 through 34 show this same arc with increasing flow rates; the time scale is fixed at 50 $\mu$ sec/division. At 0.3 gram/sec there are few fluctuations and they are slow and of low amplitude. At 1.1 gram/sec the fluctuations are faster and of appreciable amplitude. With 3.75 gram/sec the integrated light intensity almost disappears at times. At 10.8 and 15.5 gram/sec this trend continues: the light modulations increase both in amplitude and frequency.

The extreme intensity variations seen in figures 29 through 34 occur in a 50 ampere arc; a higher current arc which better fills the tube shows modulations with less amplitude at the same flow rate. In a 125 ampere arc with 7.7 gram/sec flow rate the RMS fluctuation amplitude is only about 20% of the DC average value. Compare figure 35 with figure 33 and notice that although the amplitude is suppressed, the frequency ranges are comparable.

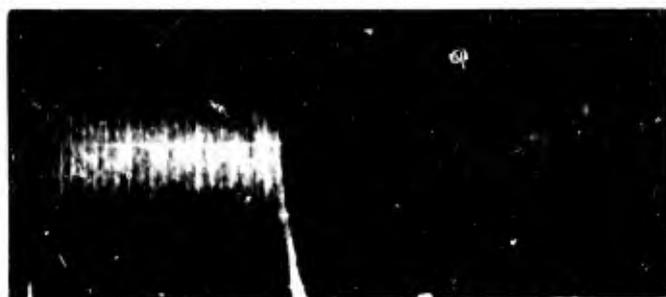
Another record of light intensity fluctuations over a longer duration is shown in figure 36. The slow sweep rate, 100 msec/division, shows the output of the photo-diode during the entire length of a test. Conditions are: 125 ampere argon arc with 3.75 gram/sec flow. There seems to be little change in intensity or fluctuation as the arc continues to burn.



100μsec

Gas: Argon  
 Pressure: 11.2 atm  
 Current: 125 amp  
 Flow Rate: 7.7 gram/sec  
 Tube Diameter: 0.7 cm  
 Tube Length: 50 cm  
 Test #: 457  
 Window: 45 cm from cathode  
 ← Zero Intensity

Figure 35-Light Fluctuations of a 125 Ampere Argon Arc with 7.7 gram/sec Flow Rate, Fast Sweep



200msec

Gas: Argon  
 Pressure: 11.2 atm  
 Current: 128 amp  
 Flow Rate: 3.75 gram/sec  
 Tube Diameter: 0.7 cm  
 Window: 45 cm from cathode  
 Tube Length: 50 cm  
 Test #: 453  
 ← Zero Intensity

Figure 36-Light Fluctuations of a 125 Ampere Argon Arc with 3.75 gram/sec Flow Rate, Slow Sweep

### 3.5 Temperature Measurements

#### 3.51 Methods

##### 3.511 Milne Lorentz Method

In an earlier report (3) measurements of the radial temperature distribution of both laminar and turbulent 11.2 atm 50 amp nitrogen arcs, burning in a 0.7 cm tube had been reported. These measurements were made by applying the Milne method to the molecular band  $\lambda = 3371 \text{ \AA}$  of nitrogen. However only the temperature of parts close to the arc axis could be measured and it was highly desirable to get more points at the lower temperatures of the outer mantle of the discharge.

Such a low temperature measurement has been performed earlier ( 13 ) by using the Milne peaking method on trace additives, specifically on sodium which has prominent resonance lines.

Although distortions of temperature measurements made with trace additives due to demixing by diffusion cannot be ignored, there were nevertheless several reasons for considering this measuring method. No argon or nitrogen line is suitable for any except the highest temperatures which require much higher currents. It can be argued that violent mixing of the turbulent arcs we wish to study should overwhelm the demixing caused by diffusion. In support of this, the image converter photographs of (3) indicate an eddy velocity of roughly 20 meter/sec which probably exceeds the diffusion velocity of sodium in nitrogen. Finally, if additives prove sufficiently accurate, a series of them could be used to give the radial position of lower temperature isotherms and thus provide a complete temperature distribution for arcs. For these reasons an application of the Milne peaking method on lines of trace elements was thought to be useful.

### 3.512 Calculation of Line Intensities for Milne-Lorentz Method

In this section the theoretical dependency of the Milne maximum of the sodium D lines on plasma temperature is determined. For this purpose the equilibrium atomic and molecular species concentrations for a nitrogen-copper-sodium system must be found at the temperatures of interest, and the temperature dependence of a sodium spectral line calculated on this basis.

The temperature variation of spectral line intensity is described by:

$$5) \quad I_j \propto \frac{N_0}{Z_0} \exp \left[ -\frac{F_j}{3T} \right]$$

In this equation both  $N_0$  and  $Z_0$  are also temperature dependent so equation 5) must be supported further.

Before we pursue our calculation, it may be remarked that it is the variation of these two factors,  $N_0$  and  $Z_0$ , which decreases the line intensity at high temperatures to produce the "Milne maximum".

Provided the atomic energy levels are known,  $Z_0(T)$  can be calculated from the definition of the partition function:

$$6) \quad Z_0(T) \equiv \sum_j g_j \exp \left[ -\frac{F_j}{3T} \right]$$

The temperature dependence of  $N_0$  is embodied in the Saha equation which describes the ionization of an atomic species:

$$7) \quad \frac{P_1 P_2}{P} = \frac{2(2\pi m)^{3/2}}{h^3} \frac{Z_1}{Z_0} \exp \left[ -\frac{IP}{3T} \right]$$

If only one atomic species is involved, then electrical

neutrality insures that

$$8) \quad P_e = P_i$$

and 7) can be solved in a straight forward manner.

The system we must consider, however, consists of the dominant gas, nitrogen or argon, the additive, sodium for instance, and a trace of copper.

As a result of this multicomponent gas composition equation 8) is no longer accurate. Electrons are contributed by the dominant gas and the copper vapor as well as the sodium so the situation is described by:

$$9) \quad P_e \text{ (total)} = P_e \text{ (additive)} + P_e \text{ (dominant gas)} + P_e \text{ (copper)}$$

The additive is chosen to have a considerably lower ionization temperature than the dominant gas so a "Milne" temperature point on the outside of the arc can be found. Near the radius of the Milne maximum there will be insignificant ionization of nitrogen or argon but copper has an ionization potential of 7.72 eV which is near enough to the 5.14 eV value of sodium to produce an appreciable ionization level. This additional source of electrons will be even more important if the arc contains more copper contamination than trace dopant and many water cooled copper cascades do produce a high partial pressure of copper. If the Saha equation 7) must be extended to higher temperatures, the influence of the dominant gas will be felt. Although the ionization potential of this gas may be 14 eV or higher, the density of the dominant gas is so much greater than the dopant that the former may be the major electron source. As a result of these additional gases, equation 7) must be applied to all electron sources and the set of equations solved simultaneously (see appendix). Our computer program includes an iterative solution to the parallel Saha equations, starting from an initial estimate to the total electron density. The situation is also complicated by the necessity to compute the dissociation of molecular nitrogen into atomic nitrogen because the two forms have different ionization potentials and thus contribute electrons in unequal degrees.

The appendix contains a simple block diagram of the computer program and some results for sodium additives in nitrogen at 5, 11.2 and 20 atmospheres. The mole fractions of each constituent are printed along with the spectral intensity of the sodium D lines. A low level of copper contamination was used because our short duration arc is relatively copper free. The results show that spectral intensity

does not continue to fall at temperatures above that of the Milne maximum but may level off or even increase.

If a non-resonance spectral line is chosen instead of the D lines, the temperature of the Milne maximum is shifted to a higher value by the change in the exponential factor of equation 5). Although the other factors in 5) remain as before,  $P_e$  grows rapidly with temperature and so will be much higher at the Milne temperature of a non resonance line. A higher electron pressure  $P_e$  can however broaden the Milne maximum or even eliminate it altogether. This can be seen by writing the Saha equation as:

$$10) \quad \left( \frac{P_i}{P_0} \right) P_e = S(T)$$

where  $S(T)$  is a function only of temperature. If we consider the two cases of single and multiple electron sources we see that  $P_e = P_i$  in the former but the latter it may be that  $P_e \gg P_i$  (sodium) and the ionization will be strongly suppressed. The ionization of the dopant atom is, as we recall, the process responsible for forming the Milne maximum so if ionization is suppressed, there will be no Milne maximum.

From this discussion it is clear that the Milne maximum will be more pronounced if

1. The Milne temperature is low and/or
2. The partial pressures of unwanted contaminants (e.g. copper) are low

Item 1. calls for a dopant with low ionization energy and a spectral line with low excitation energy. Item 2 calls for pure gas, clean cool walls, and well chosen electrode materials.

### 3.52 Apparatus

Sodium was introduced into the arc by doping a carbon cathode with NaCl and relying on the arc to ablate the material. So only a gross control over the doping level is achieved, but it proved fairly simple to produce suitable sodium D lines.

A magnified image of the arc was projected on the slit of a 3.4 meter Ebert Spectrograph producing a stigmatic spectrum on the photographic plate. Exposure time was 40 msec.

The response of the plate to different exposures was calibrated by using the anode crater of carbon arc as a constant intensity source and a variable aperture in the spectrograph to cover the range of exposure levels (14). The spectral plate was evaluated by reading local emulsion densities within the arc exposure with a microphotometer and comparing these readings with readings taken within the calibrated exposure of the carbon crater. The resulting values, which are integrated intensities viewed side-on, must undergo an Abel inversion before the radial light intensity  $I(r)$  is known. From the theoretical temperature dependence of spectral line intensity (see section 3.512) the intensities can then be converted to temperatures.

### 3.53 Results

Two 50 ampere 11.2 atm nitrogen arcs were seeded as described above and the radial position of the Na  $\lambda = 5896 \text{ \AA}$  line intensity maximum determined. Although the flow rate was 0.1 gram/sec in one case and 5.0 gram/sec in the other, the radial positions of the maxima were identical within our accuracy. From the known transition probability of the 5896  $\text{\AA}$  line, the partial pressure of sodium was estimated to be about 50 dyne/cm<sup>2</sup>. At this density the peak spectral intensity was calculated to be 3500°K, so this value is marked in figure 37.

For comparison, in this Fig. also the results of the temperature distribution measurements made with a nitrogen band line (3) on laminar and turbulent 50 ampere arcs are shown.

### 3.54 Discussion

As was shown in some detail (3), a temperature measurement in a turbulent arc using the Milne method is less accurate as large intensity fluctuations are present. Our measurements of the intensity fluctuations in a 0.7 cm I. D. 50 ampere turbulent nitrogen arc showed a very considerable level of fluctuations (section 3.4). This fact then limits the accuracy of these temperature measurements. Since the light fluctuations were found to be much smaller at currents of 125 amperes, a turbulent temperature measurement at that current level seems to be the next logical step.

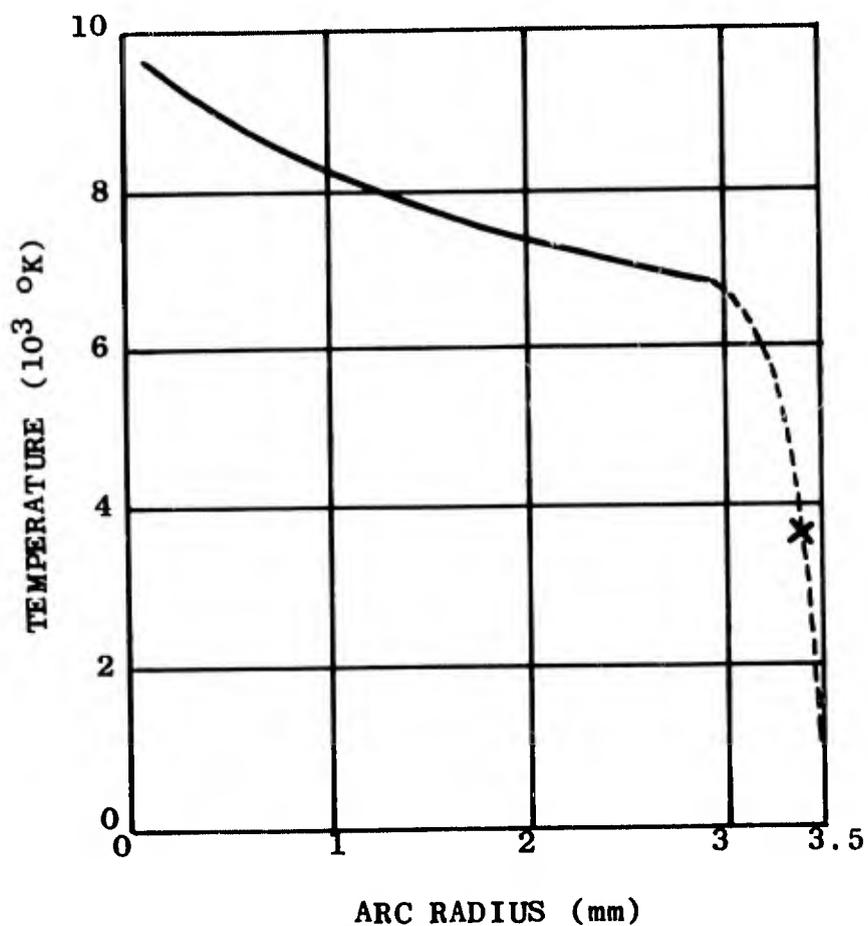


Figure 37-Radial Temperature Distribution of a 50 Ampere Nitrogen Arc. The line shows the result from a nitrogen molecular line; the "x" indicates the result of our sodium line measurement.

Gas:	Nitrogen
Current:	50 $\pm$ 2 amp
Pressure:	11.2 atm
Flow Rate:	0.1 and 5 gram/sec
Tube Diameter:	0.7 cm
Window Position:	44 cm from cathode
Tube Length:	50 cm
Shot #:	437 and 444

#### 4. CONCLUSIONS

The conclusions which follow result from experiments in the fully developed flow sections of long arc constrictors. From our experiments it is evident that turbulent effects exist in the developing region too, but it is difficult to separate the contributions of all the energy terms there, especially those of the flow terms.

##### 4.1 Existence of Turbulence

The strongest evidence we have presented in the previous sections to establish the existence of turbulence is the effect that increasing the flow rate has on heat transfer. It is interesting to note that our plasma temperature measurements tend to amplify this evidence. It was found (3) that the laminar and the turbulent temperature distributions of a 50 amp nitrogen arc were quite close. This implied that neither increased radiation nor greater ordinary thermal conduction can be responsible for the strongly increased radial heat flux in the turbulent case. Since our measurements are made in the fully developed flow section, the only remaining energy transport term to which the greatly increased losses can be assigned is turbulent thermal conduction.

Other effects we have seen which indicate turbulence include the behavior of axial pressure gradients at high flow rates, a significant enhancement of mass transport at only 5 gram/sec flow rate, and the three dimensional structure of the arc observed directly with high speed photography and indirectly with a localized light sensor. These individual effects become even more convincing when viewed together. Furthermore, the dependence of turbulent effects on experimental parameters (pressure, flow rate, tube diameter) follows a pattern similar to that of turbulence seen at low temperatures.

##### 4.2 Onset of Turbulence

The effects which establish the existence of turbulence agree closely as to the minimum conditions which show turbulent effects. For our 0.7 cm I. D. tube, we found the increase in heat transfer and the change in pressure gradient dependence both began between 1 and 2 gram/sec. The fluctuations in light intensity also begin in this area (see figures 31 and 32).

#### 4.3 Dependence of Turbulent Level on Parameters of Experiment

In addition to establishing turbulence, our aim is to provide useful data as to the level of the effect under differing operating conditions. In this way useful information for design specifications of plasma apparatus such as arc heaters is generated and values can be extrapolated even beyond the limits of our measurements.

The effects on heat flux of current, pressure, flow rate, and tube diameter were all studied. Pressure gradients were compared at different currents, pressures and flow rates. Light fluctuations were seen at different currents and flow rates. These measurements should form the foundation for design studies or more theoretical extensions of the present work.

#### 4.4 Measuring Techniques

In the process of making measurements several significant measuring techniques were developed. We saw that naturally occurring plasma structures could be utilized for measuring the plasma velocity near axis. The only requirements are high speed photographic equipment and a window in the copper cascade which does not disturb the arc yet exposes a sufficient length of the arc to view.

The sodium D lines were used as a measure of the radial position of the temperature corresponding to their maximum intensity. A close examination of the equations governing ionization and dissociation equilibria shows that the high electron density at the arc core may inhibit further ionization of dopant atoms and so alter the temperature dependence of their spectral line intensities. This complicating effect can be avoided by choosing dopant lines with expected maxima at temperatures so low that they lie well out from the arc axis.

Fluctuations of total visible light were so large for a 50 ampere 10 atmosphere argon arc at flow rates above 3 gram/sec in a 0.7 cm I. D. tube, that averaging errors must be significant. The same study showed that, if the current increased to 125 amperes, the arc was suitably stable for a measurement at flow rates below 10 gram/sec and possibly higher.

A method for observing mass transport was devised. The intense heat flux from the arc can be used to ablate low thermal conductivity materials placed at the side of the constrictor tube. Spectral observations made with stigmatic equipment focused at a downstream point can show the gradual diffusion of the ablated material through the arc.

## 5. SUGGESTIONS FOR FUTURE WORK

### 5.1 Measure Radiation Term

Our present measurement of the total radial heat flux in the fully developed section includes the radiation term with the conduction term. Direct measurement of the radiation term at several flow rates, currents and pressures will allow separation of their contributions.

### 5.2 Correlate Re with $T(r)$ and $v(r)$

Before an attempt can be made to correlate heat flux with some average Reynolds number, measuring methods for radial turbulent temperature and velocity distributions will have to be improved or new ones devised.

### 5.3 Investigate Turbulent Structure

An insight concerning the actual motion of plasma in an axial flow arc might be gained from photography at framing rates high enough to follow short term motions within a tube segment ( $10^6$  frames per sec). An alternate way to get a continuous view of plasma motions would be a streak camera.

### 5.4 Extend Parameters of Arc Operation

Because of the interest in use of high current devices, it would be desirable to extend the turbulence studies to higher currents, pressures, flow rates and tube diameters. Two obstacles are the necessity for long water cooled tubes and a power supply of about 10 megawatt.

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## APPENDIX

The program to compute the equilibrium concentrations of atomic and molecular species had to include the following processes:

1.  $N_2 \rightleftharpoons N + N$
2.  $N_2 \rightleftharpoons N_2^+ + e$
3.  $N \rightleftharpoons N^+ + e$
4.  $Cu \rightleftharpoons Cu^+ + e$
5.  $Na \rightleftharpoons Na^+ + e$

The extent to which the dissociation or ionization proceeds is described by a Saha equation (see equation 7) in text).

Process 1) is necessary to the others because it determines the concentration of N; the other processes are coupled through their common production of electrons. The sixth equation which allows the five Saha equations for the other five processes to be solved is:

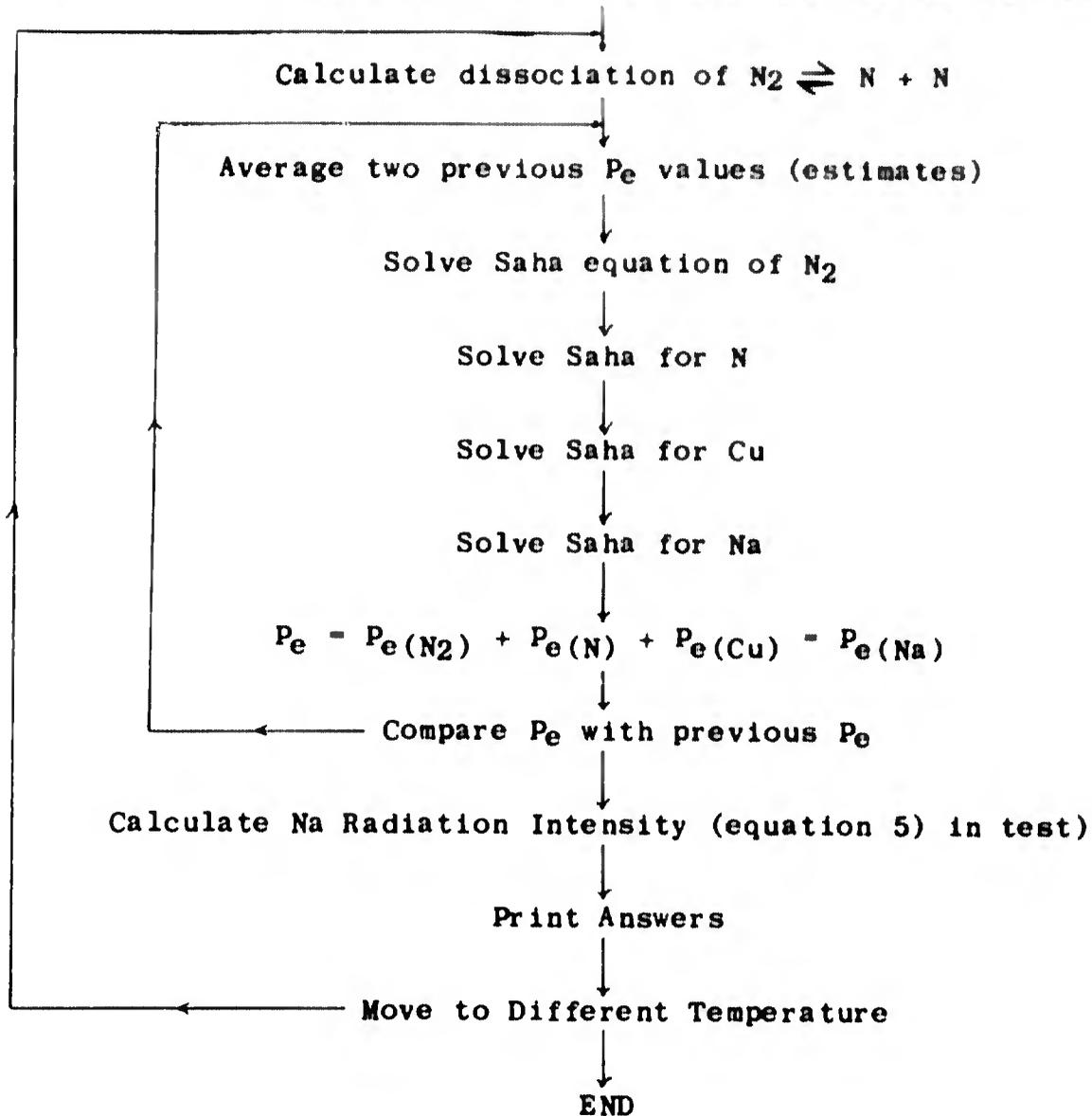
$$6) \quad P_e(\text{total}) = P_e(N_2) + P_e(N) + P_e(Cu) + P_e(Na)$$

Where the parenthetical subscript indicates the source of the electrons.

To solve these 6 simultaneous equations an iterative method was used. After solving the Saha equation for process 1), the ionization of each of the various species is computed assuming no other species contributes electrons. The degree of ionization thus calculated is too high since extra electrons suppress the ionization process. The total electron pressure is added up, see 6), giving an estimate of total electron pressure. These estimates are alternately too low and too high so, to speed the convergence, the Saha equations are solved again using the average of the two previous electron pressures. When additional repetitions of this cycle do not change the values, the computation is complete.

A Block diagram of the program is shown below:

Inputs: Total pressure, Cu partial pressure, Na partial pressure



Atomic and ionic partition functions were calculated summing only the first eight energy levels (or groups of closely spaced energy levels). Energies, degeneracies, and ionization potentials were taken from (15). The partition function of the nitrogen molecule was calculated using the methods of (16). Values of the molecular constants for the nitrogen molecular ion were taken from (17).

To test the completed program, it was run with negligible traces of copper and sodium in nitrogen and the computed species concentrations compared with the tabulations in (18). Agreement was better than 1%.

Some representative computations from the computer program are shown on the following pages. The printing format for each temperature "unit" is diagramed as follows:

Temp	Cu Atom	Na Atom	N <sub>2</sub> Molecule	N <sub>2</sub> Molec-Ion	electron
Temp	Cu Ion	Na Ion	N Atom	N Ion	Spec Intens
					5896 Å

These twelve items together describe the equilibrium state for the temperature and pressures listed. The temperature is listed at the left in °K. It is constant for two lines after which there is a skip before the temperature is increased. The temperature range covered is 2,000 to 10,000°K in steps of 200°K. The atomic and molecular species pressures are listed next, normalized against the total pressure. The total spectral intensity of a sodium D line is listed at the lower right in arbitrary units.

There are six compositions for which the equilibria are calculated. In order they are:

- I. 10 Atm N<sub>2</sub>; 1 x 10<sup>-10</sup> dyne/cm<sup>2</sup> Cu; 1 x 10<sup>-10</sup> d/cm<sup>2</sup> Na
- II. 11.2 Atm N<sub>2</sub>; 500 dyne/cm<sup>2</sup> Cu; 50 dyne/cm<sup>2</sup> Na
- III. 11.2 Atm N<sub>2</sub>; 500 dyne/cm<sup>2</sup> Cu; 150 dyne/cm<sup>2</sup> Na
- IV. 11.2 Atm N<sub>2</sub>; 500 dyne/cm<sup>2</sup> Cu; 500 dyne/cm<sup>2</sup> Na
- V. 5 Atm N<sub>2</sub>; 500 dyne/cm<sup>2</sup> Cu; 50 dyne/cm<sup>2</sup> Na
- VI. 20 Atm N<sub>2</sub>; 500 dyne/cm<sup>2</sup> Cu; 50 dyne/cm<sup>2</sup> Na

Each condition occupies 3 pages.

CHEMICAL COMPOSITION OF IMPURE NITROGEN

CONDITION: 1.0000E-10 DYNES CU; 1.0000E-10 DYNES NA  
1.0130E+07 TOTAL PRESSURE

TEMP TEMP (DEG K)	CU ATOM CU ION (FRACTION OF TOTAL PRESSURE)	NA ATOM NA ION (FRACTION OF TOTAL PRESSURE)	N2 MOLEC N ATOM (FRACTION OF TOTAL PRESSURE)	N2 MOL-ION N ION (FRACTION OF TOTAL PRESSURE)	ELECTRONS SPEC INTEN (RELATIVE)
2.00E+03	9.18E-18	5.12E-23	1.00E+00	2.59E-21	5.29E-18
2.00E+03	3.48E-19	4.94E-18	2.79E-10	2.97E-28	2.57E-21
2.20E+03	3.03E-18	6.57E-24	1.00E+00	4.90E-18	1.33E-17
2.20E+03	3.42E-18	4.94E-18	3.86E-09	4.48E-24	1.00E-21
2.40E+03	2.00E-18	1.15E-23	1.00E+00	2.72E-16	2.81E-16
2.40E+03	3.94E-18	4.94E-18	3.45E-08	1.41E-21	4.42E-21
2.60E+03	1.87E-18	2.71E-23	1.00E+00	5.57E-15	5.57E-15
2.60E+03	4.00E-18	4.94E-18	2.21E-07	1.25E-19	2.28E-20
2.80E+03	1.77E-18	5.70E-23	1.00E+00	7.36E-14	7.36E-14
2.80E+03	4.05E-18	4.94E-18	1.09E-06	5.79E-18	9.37E-20
3.00E+03	1.69E-18	1.08E-22	1.00E+00	6.94E-13	6.94E-13
3.00E+03	4.09E-18	4.94E-18	4.33E-06	1.62E-16	3.18E-19
3.20E+03	1.62E-18	1.88E-22	1.00E+00	4.97E-12	4.98E-12
3.20E+03	4.12E-18	4.94E-18	1.45E-05	3.01E-15	9.22E-19
3.40E+03	1.56E-18	3.07E-22	1.00E+00	2.85E-11	2.85E-11
3.40E+03	4.16E-18	4.94E-18	4.23E-05	3.97E-14	2.35E-18
3.60E+03	1.51E-18	4.72E-22	1.00E+00	1.35E-10	1.35E-10
3.60E+03	4.18E-18	4.94E-18	1.10E-04	3.95E-13	5.40E-18
3.80E+03	1.46E-18	6.94E-22	1.00E+00	5.43E-10	5.46E-10
3.80E+03	4.20E-18	4.94E-18	2.57E-04	3.09E-12	1.13E-17
4.00E+03	1.42E-18	9.81E-22	9.99E-01	1.91E-09	1.93E-09
4.00E+03	4.22E-18	4.94E-18	5.53E-04	1.98E-11	2.21E-17
4.20E+03	1.39E-18	1.34E-21	9.99E-01	5.99E-09	6.10E-09
4.20E+03	4.24E-18	4.94E-18	1.11E-03	1.06E-10	4.04E-17
4.40E+03	1.37E-18	1.78E-21	9.98E-01	1.69E-08	1.74E-08
4.40E+03	4.25E-18	4.93E-18	2.09E-03	4.88E-10	7.01E-17
4.60E+03	1.35E-18	2.32E-21	9.96E-01	4.37E-08	4.57E-08
4.60E+03	4.26E-18	4.93E-18	3.72E-03	1.96E-09	1.16E-16

CHEMICAL COMPOSITION OF IMPURE NITROGEN

CONDITION: 1.0000E-10 DYNES CU; 1.0000E-10 DYNES NA  
1.0130E+07 TOTAL PRESSURE

TEMP TEMP (DEG K)	CU ATOM CU ION (FRACTION OF TOTAL PRESSURE)	NA ATOM NA ION (FRACTION OF TOTAL PRESSURE)	N2 MOLEC N ATOM (FRACTION OF TOTAL PRESSURE)	N2 MOL-ION N ION (FRACTION OF TOTAL PRESSURE)	ELECTRONS SPEC INTEN (RELATIVE)
4.80E+03	1.33E-18	2.96E-21	9.94E-01	1.04E-07	1.11E-07
4.80E+03	4.27E-18	4.93E-18	6.33E-03	7.03E-09	1.85E-16
5.00E+03	1.32E-18	3.72E-21	9.90E-01	2.31E-07	2.54E-07
5.00E+03	4.27E-18	4.93E-18	1.03E-02	2.27E-08	2.84E-16
5.20E+03	1.32E-18	4.61E-21	9.84E-01	4.81E-07	5.48E-07
5.20E+03	4.28E-18	4.93E-18	1.62E-02	6.66E-08	4.25E-16
5.40E+03	1.32E-18	5.64E-21	9.75E-01	9.44E-07	1.12E-06
5.40E+03	4.28E-18	4.93E-18	2.45E-02	1.80E-07	6.20E-16
5.60E+03	1.32E-18	6.85E-21	9.64E-01	1.75E-06	2.20E-06
5.60E+03	4.28E-18	4.93E-18	3.60E-02	4.51E-07	8.84E-16
5.80E+03	1.33E-18	8.24E-21	9.49E-01	3.10E-06	4.15E-06
5.80E+03	4.27E-18	4.93E-18	5.15E-02	1.05E-06	1.24E-15
6.00E+03	1.34E-18	9.85E-21	9.28E-01	5.23E-06	7.54E-06
6.00E+03	4.27E-18	4.93E-18	7.17E-02	2.31E-06	1.70E-15
6.20E+03	1.35E-18	1.17E-20	9.03E-01	8.45E-06	1.33E-05
6.20E+03	4.26E-18	4.93E-18	9.74E-02	4.80E-06	2.30E-15
6.40E+03	1.37E-18	1.38E-20	8.71E-01	1.31E-05	2.26E-05
6.40E+03	4.25E-18	4.93E-18	1.29E-01	9.46E-06	3.07E-15
6.60E+03	1.38E-18	1.62E-20	8.32E-01	1.96E-05	3.74E-05
6.60E+03	4.25E-18	4.93E-18	1.68E-01	1.78E-05	4.05E-15
6.80E+03	1.39E-18	1.89E-20	7.87E-01	2.82E-05	6.02E-05
6.80E+03	4.24E-18	4.93E-18	2.13E-01	3.20E-05	5.26E-15
7.00E+03	1.40E-18	2.18E-20	7.34E-01	3.92E-05	9.45E-05
7.00E+03	4.24E-18	4.92E-18	2.65E-01	5.53E-05	6.74E-15
7.20E+03	1.41E-18	2.51E-20	6.76E-01	5.27E-05	1.45E-04
7.20E+03	4.23E-18	4.92E-18	3.24E-01	9.20E-05	8.55E-15
7.40E+03	1.41E-18	2.87E-20	6.13E-01	6.85E-05	2.16E-04
7.40E+03	4.23E-18	4.92E-18	3.87E-01	1.48E-04	1.07E-14

CHEMICAL COMPOSITION OF IMPURE NITROGEN

CONDITION: 1.0000E-10 DYNES CU; 1.0000E-10 DYNES NA  
1.0130E+07 TOTAL PRESSURE

TEMP TEMP (DEG K)	CU ATOM CU ION (FRACTION OF TOTAL PRESSURE)	NA ATOM NA ION (FRACTION OF TOTAL PRESSURE)	N2 MOLEC N ATOM (FRACTION OF TOTAL PRESSURE)	N2 MOL-ION N ION (FRACTION OF TOTAL PRESSURE)	ELECTRONS SPEC INTEN (RELATIVE)
7.60E+03	1.41E-18	3.26E-20	5.46E-01	8.62E-05	3.16E-04
7.60E+03	4.23E-18	4.92E-18	4.53E-01	2.30E-04	1.33E-14
7.80E+03	1.40E-18	3.67E-20	4.78E-01	1.05E-04	4.52E-04
7.80E+03	4.24E-18	4.92E-18	5.21E-01	3.47E-04	1.62E-14
8.00E+03	1.39E-18	4.11E-20	4.11E-01	1.25E-04	6.34E-04
8.00E+03	4.24E-18	4.92E-18	5.88E-01	5.09E-04	1.96E-14
8.20E+03	1.37E-18	4.57E-20	3.47E-01	1.44E-04	8.71E-04
8.20E+03	4.25E-18	4.91E-18	6.52E-01	7.27E-04	2.35E-14
8.40E+03	1.34E-18	5.05E-20	2.88E-01	1.61E-04	1.17E-03
8.40E+03	4.26E-18	4.91E-18	7.10E-01	1.01E-03	2.79E-14
8.60E+03	1.31E-18	5.54E-20	2.35E-01	1.76E-04	1.56E-03
8.60E+03	4.28E-18	4.91E-18	7.62E-01	1.38E-03	3.28E-14
8.80E+03	1.28E-18	6.06E-20	1.90E-01	1.88E-04	2.03E-03
8.80E+03	4.29E-18	4.91E-18	8.06E-01	1.84E-03	3.82E-14
9.00E+03	1.25E-18	6.59E-20	1.52E-01	1.98E-04	2.61E-03
9.00E+03	4.31E-18	4.90E-18	8.43E-01	2.42E-03	4.42E-14
9.20E+03	1.21E-18	7.14E-20	1.20E-01	2.05E-04	3.32E-03
9.20E+03	4.33E-18	4.90E-18	8.73E-01	3.11E-03	5.08E-14
9.40E+03	1.17E-18	7.70E-20	9.50E-02	2.10E-04	4.16E-03
9.40E+03	4.35E-18	4.90E-18	8.97E-01	3.95E-03	5.80E-14
9.60E+03	1.13E-18	8.29E-20	7.49E-02	2.13E-04	5.16E-03
9.60E+03	4.37E-18	4.89E-18	9.15E-01	4.95E-03	6.59E-14
9.80E+03	1.10E-18	8.91E-20	5.91E-02	2.14E-04	6.34E-03
9.80E+03	4.39E-18	4.89E-18	9.28E-01	6.13E-03	7.46E-14
1.00E+04	1.06E-18	9.54E-20	4.67E-02	2.13E-04	7.72E-03
1.00E+04	4.41E-18	4.89E-18	9.38E-01	7.50E-03	8.40E-14

CHEMICAL COMPOSITION OF IMPURE NITROGEN

CONDITION: 5.0000E+02 DYNES CU; 5.0000E+01 DYNES NA  
1.1350E+07 TOTAL PRESSURE

TEMP TEMP (DEG K)	CU ATOM CU ION (FRACTION OF TOTAL PRESSURE)	NA ATOM NA ION (FRACTION OF TOTAL PRESSURE)	N2 MOLEC N ATOM (FRACTION OF TOTAL PRESSURE)	N2 MOL-ION N ION (FRACTION OF TOTAL PRESSURE)	ELECTRONS SPEC INTEN (RELATIVE)
2.00E+03	4.41E-05	4.40E-06	1.00E+00	8.65E-30	1.42E-09
2.00E+03	5.58E-15	1.42E-09	2.63E-10	9.35E-37	2.47E-04
2.20E+03	4.41E-05	4.39E-06	1.00E+00	9.28E-27	6.25E-09
2.20E+03	9.43E-14	6.25E-09	3.64E-09	8.02E-33	7.50E-04
2.40E+03	4.41E-05	4.36E-06	1.00E+00	3.16E-24	2.17E-08
2.40E+03	1.00E-12	2.17E-08	3.26E-08	1.54E-29	1.88E-03
2.60E+03	4.41E-05	4.28E-06	1.00E+00	4.45E-22	6.23E-08
2.60E+03	7.54E-12	6.23E-08	2.09E-07	9.42E-27	4.03E-03
2.80E+03	4.41E-05	4.10E-06	1.00E+00	3.16E-20	1.53E-07
2.80E+03	4.32E-11	1.53E-07	1.03E-06	2.35E-24	7.56E-03
3.00E+03	4.41E-05	3.75E-06	1.00E+00	1.32E-18	3.26E-07
3.00E+03	2.02E-10	3.26E-07	4.09E-06	2.91E-22	1.24E-02
3.20E+03	4.41E-05	3.19E-06	1.00E+00	3.63E-17	6.10E-07
3.20E+03	8.16E-10	6.09E-07	1.37E-05	2.07E-20	1.75E-02
3.40E+03	4.40E-05	2.42E-06	1.00E+00	7.26E-16	9.96E-07
3.40E+03	2.99E-09	9.93E-07	4.00E-05	9.57E-19	2.08E-02
3.60E+03	4.40E-05	1.59E-06	1.00E+00	1.14E-14	1.42E-06
3.60E+03	1.04E-08	1.41E-06	1.03E-04	3.17E-17	2.03E-02
3.80E+03	4.40E-05	9.02E-07	1.00E+00	1.48E-13	1.79E-06
3.80E+03	3.45E-08	1.75E-06	2.43E-04	7.98E-16	1.65E-02
4.00E+03	4.38E-05	4.70E-07	9.99E-01	1.59E-12	2.08E-06
4.00E+03	1.08E-07	1.97E-06	5.23E-04	1.55E-14	1.19E-02
4.20E+03	4.34E-05	2.47E-07	9.99E-01	1.37E-11	2.38E-06
4.20E+03	3.02E-07	2.08E-06	1.05E-03	2.29E-13	8.35E-03
4.40E+03	4.26E-05	1.41E-07	9.98E-01	9.23E-11	2.85E-06
4.40E+03	7.23E-07	2.13E-06	1.97E-03	2.51E-12	6.23E-03
4.60E+03	4.11E-05	9.02E-08	9.96E-01	4.92E-10	3.62E-06
4.60E+03	1.47E-06	2.16E-06	3.52E-03	2.09E-11	5.05E-03

CHEMICAL COMPOSITION OF IMPURE NITROGEN

CONDITION: 5.0000E+02 DYNES CU; 5.0000E+01 DYNES NA  
1.1350E+07 TOTAL PRESSURE

TEMP TEMP (DEG K)	CU ATOM CU ION (FRACTION OF TOTAL PRESSURE)	NA ATOM NA ION (FRACTION OF TOTAL PRESSURE)	N2 MOLEC N ATOM (FRACTION OF TOTAL PRESSURE)	N2 MOL-ION N ION (FRACTION OF TOTAL PRESSURE)	ELECTRONS SPEC INTEN (RELATIVE)
4.80E+03	3.89E-05	6.26E-08	9.94E-01	2.17E-09	4.77E-06
4.80E+03	2.60E-06	2.17E-06	5.98E-03	1.38E-10	4.37E-03
5.00E+03	3.58E-05	4.59E-08	9.90E-01	8.29E-09	6.33E-06
5.00E+03	4.14E-06	2.18E-06	9.74E-03	7.68E-10	3.93E-03
5.20E+03	3.19E-05	3.47E-08	9.85E-01	2.83E-08	8.31E-06
5.20E+03	6.09E-06	2.19E-06	1.53E-02	3.71E-09	3.58E-03
5.40E+03	2.73E-05	2.66E-08	9.77E-01	8.90E-08	1.07E-05
5.40E+03	8.36E-06	2.19E-06	2.32E-02	1.60E-08	3.27E-03
5.60E+03	2.25E-05	2.06E-08	9.66E-01	2.60E-07	1.33E-05
5.60E+03	1.08E-05	2.19E-06	3.41E-02	6.31E-08	2.97E-03
5.80E+03	1.79E-05	1.61E-08	9.51E-01	7.09E-07	1.62E-05
5.80E+03	1.31E-05	2.19E-06	4.87E-02	2.27E-07	2.70E-03
6.00E+03	1.39E-05	1.29E-08	9.32E-01	1.79E-06	1.98E-05
6.00E+03	1.51E-05	2.20E-06	6.78E-02	7.45E-07	2.50E-03
6.20E+03	1.10E-05	1.10E-08	9.08E-01	4.03E-06	2.49E-05
6.20E+03	1.65E-05	2.20E-06	9.22E-02	2.16E-06	2.42E-03
6.40E+03	9.20E-06	1.01E-08	8.77E-01	8.04E-06	3.31E-05
6.40E+03	1.74E-05	2.20E-06	1.23E-01	5.46E-06	2.52E-03
6.60E+03	8.13E-06	1.01E-08	8.41E-01	1.42E-05	4.65E-05
6.60E+03	1.80E-05	2.20E-06	1.59E-01	1.21E-05	2.82E-03
6.80E+03	7.53E-06	1.06E-08	7.97E-01	2.27E-05	6.74E-05
6.80E+03	1.83E-05	2.20E-06	2.03E-01	2.42E-05	3.30E-03
7.00E+03	7.18E-06	1.15E-08	7.47E-01	3.39E-05	9.92E-05
7.00E+03	1.84E-05	2.20E-06	2.53E-01	4.47E-05	3.97E-03
7.20E+03	6.97E-06	1.27E-08	6.91E-01	4.76E-05	1.46E-04
7.20E+03	1.85E-05	2.20E-06	3.09E-01	7.77E-05	4.83E-03
7.40E+03	6.84E-06	1.41E-08	6.29E-01	6.38E-05	2.13E-04
7.40E+03	1.86E-05	2.20E-06	3.70E-01	1.28E-04	5.90E-03

CHEMICAL COMPOSITION OF IMPURE NITROGEN

CONDITION: 5.0000E+02 DYNES CU; 5.0000E+01 DYNES NA  
1.1350E+07 TOTAL PRESSURE

TEMP TEMP (DEG K)	CU ATOM CU ION (FRACTION OF TOTAL PRESSURE)	NA ATOM NA ION (FRACTION OF TOTAL PRESSURE)	N2 MOLEC N ATOM (FRACTION OF TOTAL PRESSURE)	N2 MOL-ION N ION (FRACTION OF TOTAL PRESSURE)	ELECTRONS SPEC INTEN (RELATIVE)
7.60E+03	6.74E-06	1.58E-08	5.64E-01	8.20E-05	3.06E-04
7.60E+03	1.87E-05	2.19E-06	4.35E-01	2.03E-04	7.19E-03
7.80E+03	6.64E-06	1.76E-08	4.97E-01	1.02E-04	4.34E-04
7.80E+03	1.87E-05	2.19E-06	5.02E-01	3.11E-04	8.71E-03
8.00E+03	6.54E-06	1.96E-08	4.30E-01	1.22E-04	6.04E-04
8.00E+03	1.88E-05	2.19E-06	5.69E-01	4.61E-04	1.05E-02
8.20E+03	6.43E-06	2.17E-08	3.66E-01	1.42E-04	8.27E-04
8.20E+03	1.88E-05	2.19E-06	6.33E-01	6.63E-04	1.25E-02
8.40E+03	6.31E-06	2.39E-08	3.06E-01	1.61E-04	1.11E-03
8.40E+03	1.89E-05	2.19E-06	6.92E-01	9.30E-04	1.48E-02
8.60E+03	6.17E-06	2.62E-08	2.52E-01	1.78E-04	1.47E-03
8.60E+03	1.89E-05	2.19E-06	7.45E-01	1.27E-03	1.74E-02
8.80E+03	6.01E-06	2.86E-08	2.05E-01	1.92E-04	1.92E-03
8.80E+03	1.90E-05	2.19E-06	7.91E-01	1.71E-03	2.02E-02
9.00E+03	5.85E-06	3.11E-08	1.65E-01	2.03E-04	2.47E-03
9.00E+03	1.91E-05	2.19E-06	8.30E-01	2.25E-03	2.34E-02
9.20E+03	5.68E-06	3.37E-08	1.31E-01	2.12E-04	3.14E-03
9.20E+03	1.92E-05	2.19E-06	8.62E-01	2.90E-03	2.69E-02
9.40E+03	5.50E-06	3.64E-08	1.04E-01	2.18E-04	3.93E-03
9.40E+03	1.93E-05	2.18E-06	8.88E-01	3.69E-03	3.07E-02
9.60E+03	5.33E-06	3.92E-08	8.26E-02	2.21E-04	4.88E-03
9.60E+03	1.94E-05	2.18E-06	9.08E-01	4.64E-03	3.49E-02
9.80E+03	5.15E-06	4.21E-08	6.53E-02	2.23E-04	5.99E-03
9.80E+03	1.94E-05	2.18E-06	9.23E-01	5.75E-03	3.95E-02
1.00E+04	4.98E-06	4.51E-08	5.18E-02	2.23E-04	7.30E-03
1.00E+04	1.95E-05	2.18E-06	9.34E-01	7.05E-03	4.45E-02

CHEMICAL COMPOSITION OF IMPURE NITROGEN

CONDITION: 5.0000E+02 DYNES CM 1.5000E+02 DYNES NA  
1.1350E+07 TOTAL PRESSURE

TEMP TEMP (DEG K)	CU ATOM CU ION (FRACTION OF TOTAL PRESSURE)	NA ATOM NA ION (FRACTION OF TOTAL PRESSURE)	N2 MOLEC N ATOM	N2 MOL-ION N ION	ELECTRONS SPEC INTEN (RELATIVE)
2.00E+03	4.41E-05	1.32E-05	1.00E+00	4.99E-30	2.45E-09
2.00E+03	3.22E-15	2.45E-09	2.63E-10	5.40E-37	7.43E-04
2.20E+03	4.41E-05	1.32E-05	1.00E+00	5.35E-27	1.08E-08
2.20E+03	5.44E-14	1.08E-08	3.64E-09	4.63E-33	2.25E-03
2.40E+03	4.41E-05	1.31E-05	1.00E+00	1.82E-24	3.77E-08
2.40E+03	5.79E-13	3.77E-08	3.26E-08	8.89E-30	5.66E-03
2.60E+03	4.41E-05	1.30E-05	1.00E+00	2.55E-22	1.09E-07
2.60E+03	4.33E-12	1.09E-07	2.09E-07	5.41E-27	1.22E-02
2.80E+03	4.41E-05	1.27E-05	1.00E+00	1.80E-20	2.69E-07
2.80E+03	2.46E-11	2.69E-07	1.03E-06	1.34E-24	2.34E-02
3.00E+03	4.41E-05	1.20E-05	1.00E+00	7.36E-19	5.84E-07
3.00E+03	1.13E-10	5.84E-07	4.09E-06	1.62E-22	3.97E-02
3.20E+03	4.41E-05	1.10E-05	1.00E+00	1.96E-17	1.13E-06
3.20E+03	4.40E-10	1.13E-06	1.37E-05	1.12E-20	6.01E-02
3.40E+03	4.40E-05	9.31E-06	1.00E+00	3.70E-16	1.95E-06
3.40E+03	1.53E-09	1.95E-06	4.00E-05	4.88E-19	8.01E-02
3.60E+03	4.40E-05	7.20E-06	1.00E+00	5.39E-15	3.01E-06
3.60E+03	4.88E-09	3.01E-06	1.03E-04	1.49E-17	9.22E-02
3.80E+03	4.40E-05	4.95E-06	1.00E+00	6.39E-14	4.15E-06
3.80E+03	1.49E-08	4.13E-06	2.43E-04	3.44E-16	9.06E-02
4.00E+03	4.40E-05	3.02E-06	99E-01	6.42E-13	5.14E-06
4.00E+03	4.37E-08	5.10E-06	23E-04	6.27E-15	7.63E-02
4.20E+03	4.38E-05	1.69E-06	9.99E-01	5.54E-12	5.88E-06
4.20E+03	1.23E-07	5.76E-06	1.05E-03	9.25E-14	5.72E-02
4.40E+03	4.34E-05	9.25E-07	9.98E-01	4.07E-11	6.47E-06
4.40E+03	3.25E-07	6.15E-06	1.97E-03	1.11E-12	4.07E-02
4.60E+03	4.25E-05	5.21E-07	9.96E-01	2.51E-10	7.12E-06
4.60E+03	7.71E-07	6.35E-06	3.52E-03	1.06E-11	2.92E-02

CHEMICAL COMPOSITION OF IMPURE NITROGEN

CONDITION: 5.0000E+02 DYNES CU; 1.5000E+02 DYNES NA  
1.1350E+07 TOTAL PRESSURE

TEMP TEMP (DEG K)	CU ATOM CU ION (FRACTION OF TOTAL PRESSURE)	NA ATOM NA ION (FRACTION OF TOTAL PRESSURE)	N2 MOLEC N ATOM (FRACTION OF TOTAL PRESSURE)	N2 MOL-ION N ION (FRACTION OF TOTAL PRESSURE)	ELECTRONS SPEC INTEN (RELATIVE)
4.80E+03	4.08E-05	3.14E-07	9.94E-01	1.29E-09	8.07E-01
4.80E+03	1.61E-06	6.45E-06	5.98E-03	8.18E-11	2.20E-02
5.00E+03	3.81E-05	2.05E-07	9.90E-01	5.55E-09	9.47E-01
5.00E+03	2.96E-06	6.51E-06	9.74E-03	5.13E-10	1.75E-02
5.20E+03	3.44E-05	1.42E-07	9.85E-01	2.07E-08	1.14E-01
5.20E+03	4.81E-06	6.54E-06	1.53E-02	2.71E-09	1.47E-02
5.40E+03	2.99E-05	1.03E-07	9.77E-01	6.91E-08	1.37E-01
5.40E+03	7.09E-06	6.56E-06	2.32E-02	1.24E-08	1.26E-02
5.60E+03	2.48E-05	7.63E-08	9.66E-01	2.10E-07	1.64E-01
5.60E+03	9.61E-06	6.57E-06	3.41E-02	5.10E-08	1.10E-02
5.80E+03	1.98E-05	5.78E-08	9.51E-01	5.91E-07	1.95E-01
5.80E+03	1.21E-05	6.58E-06	4.87E-02	1.98E-07	9.72E-03
6.00E+03	1.54E-05	4.51E-08	9.32E-01	1.53E-06	2.31E-01
6.00E+03	1.43E-05	6.59E-06	6.78E-02	6.39E-07	8.73E-03
6.20E+03	1.21E-05	3.71E-08	9.08E-01	3.58E-06	2.81E-01
6.20E+03	1.60E-05	6.59E-06	9.22E-02	1.92E-06	8.19E-03
6.40E+03	9.84E-06	3.31E-08	8.77E-01	7.38E-06	3.61E-01
6.40E+03	1.71E-05	6.59E-06	1.23E-01	5.01E-06	8.25E-03
6.60E+03	8.52E-06	3.20E-08	8.41E-01	1.34E-05	4.92E-01
6.60E+03	1.78E-05	6.59E-06	1.59E-01	1.14E-05	8.95E-03
6.80E+03	7.76E-06	3.29E-08	7.97E-01	2.19E-05	7.00E-01
6.80E+03	1.81E-05	6.59E-06	2.03E-01	2.33E-05	1.03E-02
7.00E+03	7.33E-06	3.52E-08	7.47E-01	3.31E-05	1.02E-01
7.00E+03	1.84E-05	6.59E-06	2.53E-01	4.37E-05	1.22E-02
7.20E+03	7.07E-06	3.86E-08	6.91E-01	4.68E-05	1.41E-01
7.20E+03	1.85E-05	6.59E-06	3.09E-01	7.64E-05	1.47E-02
7.40E+03	6.90E-06	4.28E-08	6.29E-01	6.31E-05	2.15E-01
7.40E+03	1.86E-05	6.59E-06	3.70E-01	1.27E-04	1.72E-02

CHEMICAL COMPOSITION OF IMPURE NITROGEN

CONDITION:  $5.0000E+02$  DYNES CU:  $1.5000E+02$  DYNES NA  
 $1.050E+07$  TOTAL PRESSURE

TEMP TEMP (°K)	O <sub>2</sub> ATOM O <sub>2</sub> ION (FRACTION OF TOTAL PRESSURE)	NA ATOM NA ION (FRACTION OF TOTAL PRESSURE)	N <sub>2</sub> MOLEC N ATOM (FRACTION OF TOTAL PRESSURE)	N <sub>2</sub> MOL-ION N ION (FRACTION OF TOTAL PRESSURE)	ELECTRONS SPEC INTEN (RELATIVE)
7.00E+02	2.78E-06	4.77E-08	5.64E-01	8.14E-05	3.09E-04
7.12E+02	2.86E-05	6.58E-06	4.35E-01	2.02E-04	2.17E-02
7.70E+02	6.67E-06	5.31E-08	4.97E-01	1.01E-04	4.36E-04
7.80E+02	1.87E-05	6.58E-06	5.02E-01	3.09E-04	2.63E-02
8.00E+02	6.57E-06	5.89E-08	4.30E-01	1.22E-04	6.06E-04
8.00E+02	1.87E-05	6.58E-06	5.69E-01	4.59E-04	3.16E-02
8.20E+02	6.45E-06	6.52E-08	3.66E-01	1.42E-04	8.29E-04
8.26E+02	1.88E-05	6.58E-06	6.33E-01	6.62E-04	3.76E-02
8.40E+02	6.32E-06	7.18E-08	3.06E-01	1.61E-04	1.11E-04
8.40E+02	1.89E-05	6.57E-06	6.92E-01	9.29E-04	4.45E-02
8.60E+02	6.18E-06	7.88E-08	2.52E-01	1.78E-04	1.48E-03
8.60E+02	1.89E-05	6.57E-06	7.45E-01	1.27E-03	5.22E-02
8.80E+02	6.02E-06	8.60E-08	2.05E-01	1.92E-04	1.90E-03
8.80E+02	1.90E-05	6.56E-06	7.91E-01	1.71E-03	6.00E-02
9.00E+02	5.85E-06	9.34E-08	1.65E-01	2.03E-04	8.47E-03
9.00E+02	1.91E-05	6.56E-06	8.30E-01	2.24E-03	7.02E-02
9.20E+02	5.68E-06	1.01E-07	1.31E-01	2.12E-04	3.14E-03
9.20E+02	1.92E-05	6.56E-06	8.62E-01	2.90E-03	8.07E-02
9.40E+02	5.51E-06	1.09E-07	1.04E-01	2.87E-04	3.94E-03
9.40E+02	1.93E-05	6.55E-06	8.88E-01	3.69E-03	9.22E-02
9.62E+02	5.33E-06	1.18E-07	8.26E-02	2.21E-04	4.88E-03
9.60E+02	1.94E-05	6.55E-06	9.08E-01	4.63E-03	1.05E-01
9.80E+02	5.15E-06	1.26E-07	6.53E-02	2.23E-04	6.00E-03
9.80E+02	1.94E-05	6.54E-06	9.23E-01	5.75E-03	1.18E-01
1.00E+04	4.98E-06	1.35E-07	5.18E-02	2.23E-04	7.30E-03
1.00E+04	1.95E-05	6.54E-06	9.34E-01	7.05E-03	1.33E-01

CHEMICAL COMPOSITION OF IMPURE NITROGEN

CONDITION: 5.0000E+02 DYNES CU; 5.0000E+02 DYNES NA  
1.1350E+07 TOTAL PRESSURE

TEMP TEMP (DEG K)	CU ATOM CU ION (FRACTION OF TOTAL PRESSURE)	NA ATOM NA ION (FRACTION OF TOTAL PRESSURE)	N2 MOLEC N ATOM (FRACTION OF TOTAL PRESSURE)	N2 MOL-ION N ION (FRACTION OF TOTAL PRESSURE)	ELECTRONS SPEC INTEN (RELATIVE)
2.00E+03	4.41E-05	4.40E-05	1.00E+00	2.73E-30	4.48E-09
2.00E+03	1.76E-15	4.48E-09	2.63E-10	2.96E-37	2.48E-03
2.20E+03	4.41E-05	4.40E-05	1.00E+00	2.93E-27	1.98E-08
2.20E+03	2.98E-14	1.98E-08	3.64E-09	2.53E-33	7.51E-03
2.40E+03	4.41E-05	4.39E-05	1.00E+00	9.94E-25	6.89E-08
2.40E+03	3.17E-13	6.89E-08	3.26E-08	4.86E-30	1.89E-02
2.60E+03	4.41E-05	4.37E-05	1.00E+00	1.39E-22	1.99E-07
2.60E+03	2.36E-12	1.99E-07	2.09E-07	2.95E-27	4.11E-02
2.80E+03	4.41E-05	4.31E-05	1.00E+00	9.76E-21	4.95E-07
2.80E+03	1.33E-11	4.95E-07	1.03E-06	7.26E-25	7.94E-02
3.00E+03	4.41E-05	4.19E-05	1.00E+00	3.95E-19	1.09E-06
3.00E+03	6.05E-11	1.09E-06	4.09E-06	8.71E-23	1.38E-01
3.20E+03	4.41E-05	3.98E-05	1.00E+00	1.03E-17	2.15E-06
3.20E+03	2.31E-10	2.15E-06	1.37E-05	5.86E-21	2.18E-01
3.40E+03	4.41E-05	3.63E-05	1.00E+00	1.88E-16	5.86E-06
3.40E+03	7.73E-10	3.86E-06	4.00E-05	2.47E-19	3.12E-01
3.60E+03	4.40E-05	3.15E-05	1.00E+00	2.58E-15	6.30E-06
3.60E+03	2.34E-09	6.29E-06	1.03E-04	7.14E-18	4.03E-01
3.80E+03	4.40E-05	2.53E-05	1.00E+00	2.83E-14	9.37E-06
3.80E+03	6.59E-09	9.37E-06	2.43E-04	1.52E-16	4.64E-01
4.00E+03	4.40E-05	1.86E-05	9.99E-01	2.59E-13	1.27E-05
4.00E+03	1.77E-08	1.27E-05	5.23E-04	2.53E-15	4.71E-01
4.20E+03	4.40E-05	1.25E-05	9.99E-01	2.06E-12	1.58E-05
4.20E+03	4.60E-08	1.58E-05	1.05E-03	3.44E-14	4.22E-01
4.40E+03	4.38E-05	7.73E-06	9.98E-01	1.44E-11	1.83E-05
4.40E+03	1.16E-07	1.82E-05	1.97E-03	3.92E-13	3.40E-01
4.60E+03	4.35E-05	4.56E-06	9.96E-01	8.91E-11	2.00E-05
4.60E+03	2.80E-07	1.97E-05	3.52E-03	3.78E-12	2.56E-01

CHEMICAL COMPOSITION OF IMPURE NITROGEN

CONDITION: 5.0000E+02 DYNES CU; 5.0000E+02 DYNES NA  
1.1350E+07 TOTAL PRESSURE

TEMP TEMP (DEG K)	CU ATOM CU ION (FRACTION OF TOTAL PRESSURE)	NA ATOM NA ION (FRACTION OF TOTAL PRESSURE)	N2 MOLEC N ATOM	N2 MOL-ION N ION	ELECTRONS SPEC INTEN (RELATIVE)
4.50E+03	4.28E-05	2.67E-06	9.94E-01	4.86E-10	2.13E-05
4.60E+03	6.32E-07	2.07E-05	5.98E-03	3.09E-11	1.86E-01
5.00E+03	4.14E-05	1.59E-06	9.90E-01	2.33E-09	2.26E-05
5.00E+03	1.34E-06	2.12E-05	9.74E-03	2.15E-10	1.36E-01
5.20E+03	3.89E-05	9.91E-07	9.85E-01	9.77E-09	2.41E-05
5.20E+03	2.56E-06	2.15E-05	1.53E-02	1.28E-09	1.03E-01
5.40E+03	3.53E-05	6.47E-07	9.77E-01	3.63E-08	2.61E-05
5.40E+03	4.39E-06	2.17E-05	2.32E-02	6.53E-09	7.96E-02
5.60E+03	3.05E-05	4.42E-07	9.66E-01	1.20E-07	2.87E-05
5.60E+03	6.76E-06	2.18E-05	3.41E-02	2.92E-08	6.39E-02
5.80E+03	2.52E-05	3.14E-07	9.51E-01	3.62E-07	3.18E-05
5.80E+03	9.43E-06	2.19E-05	4.87E-02	1.16E-07	5.28E-02
6.00E+03	1.99E-05	2.30E-07	9.32E-01	9.98E-07	3.54E-05
6.00E+03	1.21E-05	2.19E-05	6.78E-02	4.16E-07	4.46E-02
6.20E+03	1.54E-05	1.77E-07	9.08E-01	2.51E-06	4.01E-05
6.20E+03	1.43E-05	2.19E-05	9.22E-02	1.34E-06	3.89E-02
6.40E+03	1.21E-05	1.45E-07	8.77E-01	5.62E-06	4.74E-05
6.40E+03	1.60E-05	2.20E-05	1.23E-01	3.82E-06	3.61E-02
6.60E+03	9.91E-06	1.29E-07	8.41E-01	1.11E-05	5.96E-05
6.60E+03	1.71E-05	2.20E-05	1.59E-01	9.45E-06	3.61E-02
6.80E+03	8.61E-06	1.25E-07	7.97E-01	1.93E-05	7.95E-05
6.80E+03	1.77E-05	2.20E-05	2.03E-01	2.05E-05	3.89E-02
7.00E+03	7.86E-06	1.28E-07	7.47E-01	3.04E-05	1.11E-04
7.00E+03	1.81E-05	2.20E-05	2.53E-01	4.01E-05	4.42E-02
7.20E+03	7.41E-06	1.36E-07	6.91E-01	4.43E-05	1.57E-04
7.20E+03	1.83E-05	2.20E-05	3.09E-01	7.23E-05	5.19E-02
7.40E+03	7.12E-06	1.48E-07	6.29E-01	6.08E-05	2.23E-04
7.40E+03	1.85E-05	2.20E-05	3.70E-01	1.22E-04	6.19E-02

CHEMICAL COMPOSITION OF IMPURE NITROGEN

CONDITION: 5.0000E+02 DYNES CU; 5.0000E+02 DYNES NA  
1.1350E+07 TOTAL PRESSURE

TEMP TEMP (DEG K)	CU ATOM CU ION (FRACTION OF TOTAL PRESSURE)	NA ATOM NA ION (FRACTION OF TOTAL PRESSURE)	N2 MOLEC N ATOM (FRACTION OF TOTAL PRESSURE)	N2 MOL-ION N ION (FRACTION OF TOTAL PRESSURE)	ELECTRONS SPEC INTEN (RELATIVE)
7.60E+03	6.93E-06	1.63E-07	5.64E-01	7.94E-05	3.17E-04
7.60E+03	1.86E-05	2.19E-05	4.35E-01	1.97E-04	7.43E-02
7.80E+03	6.78E-06	1.80E-07	4.97E-01	9.95E-05	4.44E-04
7.80E+03	1.86E-05	2.19E-05	5.02E-01	3.04E-04	8.92E-02
8.00E+03	6.64E-06	1.99E-07	4.30E-01	1.20E-04	6.14E-04
8.00E+03	1.87E-05	2.19E-05	5.69E-01	4.53E-04	1.07E-01
8.20E+03	6.50E-06	2.19E-07	3.66E-01	1.41E-04	8.37E-04
8.20E+03	1.88E-05	2.19E-05	6.33E-01	6.55E-04	1.27E-01
8.40E+03	6.36E-06	2.41E-07	3.06E-01	1.60E-04	1.12E-03
8.40E+03	1.88E-05	2.19E-05	6.92E-01	9.22E-04	1.49E-01
8.60E+03	6.20E-06	2.64E-07	2.52E-01	1.77E-04	1.48E-03
8.60E+03	1.89E-05	2.19E-05	7.45E-01	1.27E-03	1.75E-01
8.80E+03	6.04E-06	2.88E-07	2.05E-01	1.91E-04	1.93E-03
8.80E+03	1.90E-05	2.19E-05	7.91E-01	1.70E-03	2.03E-01
9.00E+03	5.87E-06	3.12E-07	1.65E-01	2.02E-04	2.48E-03
9.00E+03	1.91E-05	2.19E-05	8.30E-01	2.24E-03	2.35E-01
9.20E+03	5.70E-06	3.38E-07	1.31E-01	2.11E-04	3.15E-03
9.20E+03	1.92E-05	2.19E-05	8.62E-01	2.89E-03	2.70E-01
9.40E+03	5.52E-06	3.65E-07	1.04E-01	2.17E-04	3.94E-03
9.40E+03	1.93E-05	2.18E-05	8.88E-01	3.68E-03	3.08E-01
9.60E+03	5.34E-06	3.93E-07	8.26E-02	2.21E-04	4.89E-03
9.60E+03	1.94E-05	2.18E-05	9.08E-01	4.63E-03	3.50E-01
9.80E+03	5.16E-06	4.21E-07	6.53E-02	2.23E-04	6.00E-03
9.80E+03	1.94E-05	2.18E-05	9.23E-01	5.74E-03	3.95E-01
1.00E+04	4.99E-06	4.51E-07	5.18E-02	2.23E-04	7.31E-03
1.00E+04	1.95E-05	2.18E-05	9.34E-01	7.04E-03	4.45E-01

CHEMICAL COMPOSITION OF IMPURE NITROGEN

CONDITION: 5.0000E+02 DYNES CU; 5.0000E+01 DYNES NA  
5.0650E+06 TOTAL PRESSURE

TEMP TEMP (DEG K)	CU ATOM CU ION (FRACTION OF TOTAL PRESSURE)	NA ATOM NA ION (FRACTION OF TOTAL PRESSURE)	N2 MOLEC N ATOM (FRACTION OF TOTAL PRESSURE)	N2 MOL-ION N ION (FRACTION OF TOTAL PRESSURE)	ELECTRONS SPEC INTEN (RELATIVE)
2.00E+03	9.87E-05	9.87E-06	1.00E+00	8.65E-30	3.17E-09
2.00E+03	1.25E-14	3.17E-09	3.94E-10	1.40E-36	2.47E-04
2.20E+03	9.87E-05	9.84E-06	1.00E+00	9.28E-27	1.40E-08
2.20E+03	2.11E-13	1.40E-08	5.45E-09	1.20E-32	7.50E-04
2.40E+03	9.87E-05	9.77E-06	1.00E+00	3.16E-24	4.86E-08
2.40E+03	2.25E-12	4.86E-08	4.88E-08	2.31E-29	1.88E-03
2.60E+03	9.87E-05	9.59E-06	1.00E+00	4.45E-22	1.40E-07
2.60E+03	1.69E-11	1.40E-07	3.12E-07	1.41E-26	4.03E-03
2.80E+03	9.87E-05	9.19E-06	1.00E+00	3.16E-20	3.42E-07
2.80E+03	9.68E-11	3.42E-07	1.54E-06	3.52E-24	7.56E-03
3.00E+03	9.87E-05	8.41E-06	1.00E+00	1.32E-18	7.31E-07
3.00E+03	4.53E-10	7.30E-07	6.12E-06	4.36E-22	1.24E-02
3.20E+03	9.87E-05	7.14E-06	1.00E+00	3.63E-17	1.37E-06
3.20E+03	1.83E-09	1.36E-06	2.05E-05	3.10E-20	1.75E-02
3.40E+03	9.87E-05	5.42E-06	1.00E+00	7.26E-16	2.23E-06
3.40E+03	6.70E-09	2.23E-06	5.98E-05	1.43E-18	2.08E-02
3.60E+03	9.87E-05	3.56E-06	1.00E+00	1.14E-14	3.18E-06
3.60E+03	2.32E-08	3.16E-06	1.55E-04	4.75E-17	2.03E-02
3.80E+03	9.86E-05	2.02E-06	1.00E+00	1.48E-13	4.00E-06
3.80E+03	7.73E-08	3.93E-06	3.63E-04	1.19E-15	1.65E-02
4.00E+03	9.82E-05	1.05E-06	9.99E-01	1.59E-12	4.65E-06
4.00E+03	2.42E-07	4.41E-06	7.82E-04	2.32E-14	1.19E-02
4.20E+03	9.74E-05	5.54E-07	9.98E-01	1.37E-11	5.34E-06
4.20E+03	6.77E-07	4.66E-06	1.57E-03	3.42E-13	8.35E-03
4.40E+03	9.55E-05	3.17E-07	9.97E-01	9.22E-11	6.40E-06
4.40E+03	1.62E-06	4.78E-06	2.95E-03	3.76E-12	6.23E-03
4.60E+03	9.21E-05	2.02E-07	9.95E-01	4.92E-10	8.12E-06
4.60E+03	3.28E-06	4.83E-06	5.26E-03	3.12E-11	5.05E-03

CHEMICAL COMPOSITION OF IMPURE NITROGEN

CONDITION: 5.0000E+02 DYNES CU; 5.0000E+01 DYNES NA  
5.0650E+06 TOTAL PRESSURE

TEMP TEMP (DEG K)	CU ATOM CU ION (FRACTION OF TOTAL PRESSURE)	NA ATOM NA ION (FRACTION OF TOTAL PRESSURE)	N2 MOLEC N ATOM (FRACTION OF TOTAL PRESSURE)	N2 MOL-ION N ION (FRACTION OF TOTAL PRESSURE)	ELECTRONS SPEC INTEN (RELATIVE)
4.80E+03	8.71E-05	1.40E-07	9.91E-01	2.17E-09	1.07E-05
4.80E+03	5.82E-06	4.87E-06	8.94E-03	2.07E-10	4.37E-03
5.00E+03	8.01E-05	1.03E-07	9.85E-01	8.26E-09	1.42E-05
5.00E+03	9.29E-06	4.88E-06	1.45E-02	1.15E-09	3.93E-03
5.20E+03	7.14E-05	7.76E-08	9.77E-01	2.82E-08	1.86E-05
5.20E+03	1.37E-05	4.90E-06	2.28E-02	5.53E-09	3.58E-03
5.40E+03	6.12E-05	5.94E-08	9.66E-01	8.83E-08	2.38E-05
5.40E+03	1.88E-05	4.91E-06	3.45E-02	2.39E-08	3.26E-03
5.60E+03	5.02E-05	4.57E-08	9.49E-01	2.58E-07	2.95E-05
5.60E+03	2.43E-05	4.91E-06	5.06E-02	9.44E-08	2.95E-03
5.80E+03	3.95E-05	3.53E-08	9.28E-01	7.07E-07	3.56E-05
5.80E+03	2.96E-05	4.92E-06	7.20E-02	3.44E-07	2.64E-03
6.00E+03	3.01E-05	2.75E-08	9.00E-01	1.81E-06	4.22E-05
6.00E+03	3.43E-05	4.92E-06	9.98E-02	1.15E-06	2.38E-03
6.20E+03	2.29E-05	2.23E-08	8.65E-01	4.25E-06	5.05E-05
6.20E+03	3.79E-05	4.92E-06	1.35E-01	3.49E-06	2.19E-03
6.40E+03	1.81E-05	1.94E-08	8.22E-01	8.83E-06	6.33E-05
6.40E+03	4.03E-05	4.93E-06	1.78E-01	9.27E-06	2.15E-03
6.60E+03	1.53E-05	1.82E-08	7.71E-01	1.61E-05	8.42E-05
6.60E+03	4.17E-05	4.93E-06	2.28E-01	2.15E-05	2.28E-03
6.80E+03	1.36E-05	1.84E-08	7.13E-01	2.61E-05	1.18E-04
6.80E+03	4.25E-05	4.93E-06	2.87E-01	4.41E-05	2.57E-03
7.00E+03	1.27E-05	1.95E-08	6.47E-01	3.87E-05	1.69E-04
7.00E+03	4.30E-05	4.93E-06	3.52E-01	8.21E-05	3.01E-03
7.20E+03	1.21E-05	2.12E-08	5.77E-01	5.33E-05	2.44E-04
7.20E+03	4.33E-05	4.93E-06	4.23E-01	1.42E-04	3.61E-03
7.40E+03	1.18E-05	2.33E-08	5.03E-01	6.93E-05	3.51E-04
7.40E+03	4.35E-05	4.92E-06	4.96E-01	2.33E-04	4.35E-03

CHEMICAL COMPOSITION OF IMPURE NITROGEN

CONDITION: 5.0000E+02 DYNES CU; 5.0000E+01 DYNES NA  
5.0650E+06 TOTAL PRESSURE

TEMP TEMP (DEG K)	CU ATOM CU ION (FRACTION OF TOTAL PRESSURE)	NA ATOM NA ION (FRACTION OF TOTAL PRESSURE)	N2 MOLEC N ATOM (FRACTION OF TOTAL PRESSURE)	N2 MOL-ION N ION (FRACTION OF TOTAL PRESSURE)	ELECTRONS SPEC INTEN (RELATIVE)
7.60E+03	1.15E-05	2.58E-08	4.30E-01	8.59E-05	5.00E-04
7.60E+03	4.36E-05	4.92E-06	5.69E-01	3.65E-04	5.24E-03
7.80E+03	1.12E-05	2.85E-08	3.59E-01	1.02E-04	7.00E-04
7.80E+03	4.38E-05	4.92E-06	6.39E-01	5.49E-04	6.29E-03
8.00E+03	1.09E-05	3.13E-08	2.94E-01	1.17E-04	9.66E-04
8.00E+03	4.39E-05	4.92E-06	7.04E-01	7.99E-04	7.49E-03
8.20E+03	1.06E-05	3.44E-08	2.37E-01	1.30E-04	1.31E-03
8.20E+03	4.40E-05	4.92E-06	7.61E-01	1.13E-03	8.85E-03
8.40E+03	1.03E-05	3.75E-08	1.87E-01	1.41E-04	1.75E-03
8.40E+03	4.42E-05	4.92E-06	8.09E-01	1.56E-03	1.04E-02
8.60E+03	1.00E-05	4.08E-08	1.46E-01	1.49E-04	2.29E-03
8.60E+03	4.43E-05	4.92E-06	8.49E-01	2.09E-03	1.21E-02
8.80E+03	9.69E-06	4.43E-08	1.13E-01	1.55E-04	2.96E-03
8.80E+03	4.45E-05	4.91E-06	8.81E-01	2.76E-03	1.40E-02
9.00E+03	9.36E-06	4.78E-08	8.76E-02	1.58E-04	3.79E-03
9.00E+03	4.47E-05	4.91E-06	9.05E-01	3.58E-03	1.60E-02
9.20E+03	9.03E-06	5.15E-08	6.75E-02	1.60E-04	4.78E-03
9.20E+03	4.48E-05	4.91E-06	9.23E-01	4.57E-03	1.83E-02
9.40E+03	8.70E-06	5.53E-08	5.20E-02	1.60E-04	5.96E-03
9.40E+03	4.50E-05	4.91E-06	9.36E-01	5.75E-03	2.08E-02
9.60E+03	8.38E-06	5.94E-08	4.02E-02	1.60E-04	7.37E-03
9.60E+03	4.52E-05	4.91E-06	9.45E-01	7.16E-03	2.36E-02
9.80E+03	8.07E-06	6.36E-08	3.12E-02	1.59E-04	9.03E-03
9.80E+03	4.53E-05	4.90E-06	9.51E-01	8.82E-03	2.66E-02
1.00E+04	7.77E-06	6.79E-08	2.44E-02	1.57E-04	1.10E-02
1.00E+04	4.55E-05	4.90E-06	9.54E-01	1.07E-02	2.99E-02

CHEMICAL COMPOSITION OF IMPURE NITROGEN

CONDITION: 5.0000E+02 DYNES CU; 5.0000E+01 DYNES NA  
2.0260E+07 TOTAL PRESSURE

TEMP TEMP (DEG K)	CU ATOM CU ION (FRACTION OF TOTAL PRESSURE)	NA ATOM NA ION (FRACTION OF TOTAL PRESSURE)	N2 MOLEC N ATOM (FRACTION OF TOTAL PRESSURE)	N2 MOL-ION N ION (FRACTION OF TOTAL PRESSURE)	ELECTRONS SPEC INTEN (RELATIVE)
2.00E+03	2.47E-05	2.47E-06	1.00E+00	8.65E-30	7.93E-10
2.00E+03	3.13E-15	7.93E-10	1.97E-10	7.00E-37	2.47E-04
2.20E+03	2.47E-05	2.46E-06	1.00E+00	9.28E-27	3.50E-09
2.20E+03	5.28E-14	3.50E-09	2.73E-09	6.00E-33	7.50E-04
2.40E+03	2.47E-05	2.44E-06	1.00E+00	3.16E-24	1.22E-08
2.40E+03	5.63E-13	1.22E-08	2.44E-08	1.15E-29	1.88E-03
2.60E+03	2.47E-05	2.40E-06	1.00E+00	4.45E-22	3.49E-08
2.60E+03	4.22E-12	3.49E-08	1.56E-07	7.05E-27	4.03E-03
2.80E+03	2.47E-05	2.30E-06	1.00E+00	3.16E-20	8.56E-08
2.80E+03	2.42E-11	8.56E-08	7.68E-07	1.76E-24	7.56E-03
3.00E+03	2.47E-05	2.10E-06	1.00E+00	1.32E-18	1.83E-07
3.00E+03	1.13E-10	1.83E-07	3.06E-06	2.18E-22	1.24E-02
3.20E+03	2.47E-05	1.79E-06	1.00E+00	3.63E-17	3.42E-07
3.20E+03	4.57E-10	3.41E-07	1.03E-05	1.55E-20	1.75E-02
3.40E+03	2.47E-05	1.36E-06	1.00E+00	7.26E-16	5.58E-07
3.40E+03	1.68E-09	5.56E-07	2.99E-05	7.17E-19	2.08E-02
3.60E+03	2.47E-05	8.89E-07	1.00E+00	1.14E-14	7.95E-07
3.60E+03	5.81E-09	7.89E-07	7.74E-05	2.37E-17	2.03E-02
3.80E+03	2.46E-05	5.05E-07	1.00E+00	1.48E-13	1.00E-06
3.80E+03	1.93E-08	9.81E-07	1.82E-04	5.97E-16	1.65E-02
4.00E+03	2.46E-05	2.64E-07	1.00E+00	1.59E-12	1.16E-06
4.00E+03	6.05E-08	1.10E-06	3.91E-04	1.16E-14	1.19E-02
4.20E+03	2.43E-05	1.39E-07	9.99E-01	1.37E-11	1.33E-06
4.20E+03	1.69E-07	1.16E-06	7.84E-04	1.71E-13	8.35E-03
4.40E+03	2.39E-05	7.93E-08	9.99E-01	9.23E-11	1.60E-06
4.40E+03	4.05E-07	1.19E-06	1.48E-03	1.88E-12	6.23E-03
4.60E+03	2.30E-05	5.05E-08	9.97E-01	4.93E-10	2.03E-06
4.60E+03	8.21E-07	1.21E-06	2.63E-03	1.56E-11	5.05E-03

CHEMICAL COMPOSITION OF IMPURE NITROGEN

CONDITION: 5.0000E+02 DYNES CU; 5.0000E+01 DYNES NA  
2.0260E+07 TOTAL PRESSURE

TEMP TEMP (DEG K)	CU ATOM CU ION (FRACTION OF TOTAL PRESSURE)	NA ATOM NA ION (FRACTION OF TOTAL PRESSURE)	N2 MOLEC N ATOM (FRACTION OF TOTAL PRESSURE)	N2 MOL-ION N ION (FRACTION OF TOTAL PRESSURE)	ELECTRONS SPEC INTEN (RELATIVE)
4.80E+03	2.18E-05	3.51E-08	9.96E-01	2.17E-09	2.67E-06
4.80E+03	1.45E-06	1.22E-06	4.48E-03	1.04E-10	4.38E-03
5.00E+03	2.00E-05	2.57E-08	9.93E-01	8.31E-09	3.55E-06
5.00E+03	2.32E-06	1.22E-06	7.30E-03	5.75E-10	3.93E-03
5.20E+03	1.79E-05	1.95E-08	9.89E-01	2.84E-08	4.66E-06
5.20E+03	3.41E-06	1.22E-06	1.15E-02	2.77E-09	3.59E-03
5.40E+03	1.53E-05	1.50E-08	9.83E-01	8.90E-08	5.99E-06
5.40E+03	4.67E-06	1.23E-06	1.74E-02	1.20E-08	3.29E-03
5.60E+03	1.27E-05	1.17E-08	9.74E-01	2.59E-07	7.53E-06
5.60E+03	6.00E-06	1.23E-06	2.56E-02	4.68E-08	3.01E-03
5.80E+03	1.02E-05	9.26E-09	9.63E-01	6.98E-07	9.35E-06
5.80E+03	7.25E-06	1.23E-06	3.67E-02	1.67E-07	2.78E-03
6.00E+03	8.11E-06	7.66E-09	9.49E-01	1.71E-06	1.18E-05
6.00E+03	8.28E-06	1.23E-06	5.12E-02	5.31E-07	2.65E-03
6.20E+03	6.67E-06	6.81E-09	9.30E-01	3.73E-06	1.54E-05
6.20E+03	9.01E-06	1.23E-06	6.99E-02	1.48E-06	2.68E-03
6.40E+03	5.77E-06	6.55E-09	9.07E-01	7.18E-06	2.15E-05
6.40E+03	9.45E-06	1.23E-06	9.32E-02	3.59E-06	2.92E-03
6.60E+03	5.25E-06	6.73E-09	8.78E-01	1.24E-05	3.11E-05
6.60E+03	9.71E-06	1.23E-06	1.22E-01	7.76E-06	3.37E-03
6.80E+03	4.96E-06	7.22E-09	8.44E-01	1.97E-05	4.61E-05
6.80E+03	9.86E-06	1.23E-06	1.56E-01	1.53E-05	4.02E-03
7.00E+03	4.79E-06	7.93E-09	8.04E-01	2.95E-05	6.87E-05
7.00E+03	9.95E-06	1.23E-06	1.96E-01	2.81E-05	4.90E-03
7.20E+03	4.69E-06	8.84E-09	7.58E-01	4.19E-05	1.02E-04
7.20E+03	1.00E-05	1.23E-06	2.42E-01	4.88E-05	6.02E-03
7.40E+03	4.62E-06	9.91E-09	7.06E-01	5.71E-05	1.49E-04
7.40E+03	1.00E-05	1.23E-06	2.94E-01	8.11E-05	7.39E-03

CHEMICAL COMPOSITION OF IMPURE NITROGEN

CONDITION: 5.0000E+02 DYNES CU; 5.0000E+01 DYNES NA  
2.0260E+07 TOTAL PRESSURE

TEMP TEMP (DEG K)	CU ATOM CU ION (FRACTION OF TOTAL PRESSURE)	NA ATOM NA ION (FRACTION OF TOTAL PRESSURE)	N2 MOLEC N ATOM	N2 MOL-ION N ION	ELECTRONS SPEC INTEN (RELATIVE)
7.60E+03	4.57E-06	1.11E-08	6.50E-01	7.51E-05	2.16E-04
7.60E+03	1.01E-05	1.23E-06	3.50E-01	1.30E-04	9.05E-03
7.80E+03	4.53E-06	1.25E-08	5.90E-01	9.55E-05	3.07E-04
7.80E+03	1.01E-05	1.23E-06	4.10E-01	2.01E-04	1.10E-02
8.00E+03	4.48E-06	1.39E-08	5.28E-01	1.18E-04	4.30E-04
8.00E+03	1.01E-05	1.23E-06	4.72E-01	3.01E-04	1.33E-02
8.20E+03	4.42E-06	1.55E-08	4.65E-01	1.42E-04	5.92E-04
8.20E+03	1.01E-05	1.23E-06	5.34E-01	4.39E-04	1.59E-02
8.40E+03	4.36E-06	1.72E-08	4.03E-01	1.66E-04	8.00E-04
8.40E+03	1.02E-05	1.23E-06	5.95E-01	6.23E-04	1.90E-02
8.60E+03	4.29E-06	1.89E-08	3.45E-01	1.89E-04	1.07E-03
8.60E+03	1.02E-05	1.22E-06	6.53E-01	8.65E-04	2.24E-02
8.80E+03	4.20E-06	2.08E-08	2.91E-01	2.10E-04	1.40E-03
8.80E+03	1.02E-05	1.22E-06	7.06E-01	1.18E-03	2.62E-02
9.00E+03	4.11E-06	2.27E-08	2.42E-01	2.29E-04	1.81E-03
9.00E+03	1.03E-05	1.22E-06	7.54E-01	1.56E-03	3.04E-02
9.20E+03	4.01E-06	2.47E-08	1.99E-01	2.45E-04	2.30E-03
9.20E+03	1.03E-05	1.22E-06	7.96E-01	2.05E-03	3.51E-02
9.40E+03	3.90E-06	2.68E-08	1.63E-01	2.58E-04	2.90E-03
9.40E+03	1.04E-05	1.22E-06	8.31E-01	2.63E-03	4.03E-02
9.60E+03	3.79E-06	2.89E-08	1.32E-01	2.68E-04	3.61E-03
9.60E+03	1.04E-05	1.22E-06	8.61E-01	3.33E-03	4.60E-02
9.80E+03	3.68E-06	3.11E-08	1.07E-01	2.75E-04	4.45E-03
9.80E+03	1.05E-05	1.22E-06	8.84E-01	4.16E-03	5.21E-02
1.00E+04	3.57E-06	3.34E-08	8.60E-02	2.80E-04	5.43E-03
1.00E+04	1.06E-05	1.22E-06	9.03E-01	5.14E-03	5.88E-02

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ABSTRACT Heat, mass, and momentum transfer were investigated experimentally in very long cylindrical axial flow arcs in laminar and turbulent flow conditions. Special attention was given to assure that the end of the arc tube had a section of fully developed flow. In this section the number of significant terms in the energy balance equation is reduced and interpretation of experimental results is simplified. Strong evidence is presented for the existence of turbulence by the fact that all three transfer mechanisms are considerably accelerated in what we condier the "turbulent mode" and also by direct optical study of the arc by a light probe. The onset of turbulence was found with all our methods to occur around 1.5 gram/sec gas flow rate. A considerable number of measurements of heat and momentum transfer are presented in the following operating ranges: Gases-argon and nitrogen; tube diameters-0.5 and 0.7 cm; tube length-50 cm; mass flow rates-0.1 to 15 gram/sec; ambient pressure - 1 to 20 atmosphere; currents 25 to 200 amperes (some higher).		

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Arc Plasma Heat Transfer Heat, Mass, Momentum Transfer Laminar Turbulent Arc Plasmas						