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INVISCID BLUNT BODY SHOCK LAYERS

Two-Dimensional Symmetric and Axisymmetric Flows

by

Gino Moretti

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POLYTECHNIC INSTITUTE OF BROOKLYN

DEPARTMENT of AEROSPACE ENGINEERING and APPLIED MECHANICS June 1968

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PIBAL REPORT NO. 68-15

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INVISCID BLUNT BODY SHOCK LAYERS

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(Two-Dimensional Symmetric and Axisymmetric Flows)

by

Gino Moretti

This research was conducted under the sponsorship of the Office of Naval Research under Contract No. Nonr 839(34), Project No. NR 061-135.

Polytechnic Institute of Brooklyn

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INVISCID BLUNT BODY SHOCK LAYERS

(Two-Dimensional Symmetric and Axisymmetric Flows)

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Gino Moretti

Polytechnic Institute of Brooklyn

ABSTRACT

The results of a time-dependent computation of blunt body shock layers for two-dimensional symmetric and axisymmetric flows are presented in a systematic form for a range of values of the free stream Mach number and bodies of different shapes and variable bluntness.

A brief discussion of the relevant features of the computational technique is given. The results are presented in a graphical form; the graphs have been produced by the computer.

This research was conducted under the sponsorship of the Office of Naval Research under Contract No. Nonr 839(34), Project No. NR 061-135.

Professor, Department of Aerospace Engineering and Applied Mechanics.

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1. Introduction

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The main features of the flow field about the blunt nose of determined flying at supersonic speed are, qualitatively, wellwhown. A complete, detailed description of the flow, however, is hard to obtain. Formulae for the determination of some relevant parameters are available, which generally rely on simplifying assumptions. More detailed results can be achieved through the use of humerical techniques on high-speed computers.*

One of these techniques, which was considered for the first time in 1965 (Refs. 1 and 2) appeared as well-suited for practical purposes. As in most numerical procedures, the flow field is computed at the nodal points of a mesh. The mesh covers the shock layer only. Its fineness is controlled by two parameters in plane and axisymmetric problems, and by three parameters in threedimensional problems. Obviously, the finer the mesh, the longer the computational time required to obtain the solution.^{**} One important advantage of the technique resides in the fact that, even when an extremely coarse mesh is used, the values at the nodal points may be sufficiently accurate to provide a useful

^{*} Hayes and Probstein's book on Hypersonic Flow (Academic Press, 1967) contains an exhaustive description and a critical analysis of practically all methods published in the open literature until ~965.

^{**} In problems involving two space parameters, halving the mesh size in both directions lengthens the computational time by a factor of 8. A factor of 16 is related to the halving of the mesh size in all three directions in a three-dimensional problem.

preliminary description of the flow field. If the mesh is coarse, the computational time is of the order of a few seconds.^{*} By increasing the number of nodal points, not only more detailed information is obtained but the accuracy is increased. A computation which requires about two minutes of machine operation yields results whose accuracy is generally greater than that of the most sophisticated experiments. In this connection, it may be noted that the use of a very coarse mesh is conceptually similar to the application of the method of integral relations (Ref. 3) with one or two strips. However, the present technique does not require a reformulation of the equations if the fineness of the mesh is increased, but merely a change in the two, or three, integers which define the number of nodal points.

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Therefore, it was considered appropriate to perform a systematic series of computations with a twofold purpose:

1) to test the technique in the widest possible range of body shapes and flow properties, and

2) to provide a parametric compilation of cases which are of practical interest, in the hope that it could be used as a quick reference for neighboring cases.

Part of the results are published in the present report. Since the object of the report is a compilation and discussion of

All time estimates are based on actual computations performed on the CDC 6600 computer.

results, only a few words will be spent to describe the procedure used to obtain them. In what follows, the main features of the technique will be recalled from the previous communications referred to above and the adoption of new frames of reference will be justified.

2. Outline of the Computational Technique

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For the sake of simplicity, we will focus our attention on the two-dimensional case. Let AB (Figs. 1 and 2) be a section of the body, CD a section of the shock wave, EF the sonic line. In Refs. 1



and 2, the equations of motion are written in a Cartesian frame (x,y). A natural boundary of the region to be computed is then

the closed line ABDCA, which has a segment, BD, parallel to the x-axis. Such a choice is possible only if (i) the body slope is positive up to a point inside the supersonic region, and (ii) the sonic line does not intersect the upper boundary. Both limitations are particularly severe if the free stream Mach number is low.

The present computations have been performed after reformulating the equations of motion in a polar (r, θ) frame, as in Fig. 2. In the axisymmetric case, a spherical frame (r, θ, ϕ) is used, with the polar axis along the body centerline. In any meridional plane the section of the flow field appears again as in Fig. 2.



As in Ref. 2, the region to be computed is first mapped onto a rectangle (Fig. 3) by linearly stretching the r-coordinate between shock and body. New coordinates, ζ , Y and T are introduced, defined by

$$\zeta = \frac{r - r_{body}}{r_{shock} - r_{body}}$$
$$Y = \pi - \theta$$
$$T = t$$

where t is the time.

To compute interior points, the equations of motion are reformulated in terms of the new independent variables. Then the derivatives of the equations of motion with respect to ζ , Y and T are formally evaluated. From such a system of equations the first and second order derivatives of density (ρ), velocity components (u and v), and entropy (S) with respect to T can be expressed as functions of first and second order derivatives of the same parameters with respect to ζ and Y. The increment of any parameter $f(f=\zeta,u,v,S)$ in a time step, t, is computed as

The time step at each nodal point is taken equal to

where $q = (u^2 + v^2)^{\frac{1}{2}}$, a is the local speed of sound, and Δs is the length of the shorter local side of a mesh quadrangle in the physical plane. At the end of each step, the local increments are linearly interpolated to a common time step, chosen as the smallest

of the local time steps.

Shock and body points are computed differently. In both cases the pressure is determined by a modified method of characteristics, as outlined in Ref. 2. At shock points a complete system of equations is obtained by considering the Rankine-Hugoniot conditions for a moving shock. The system is then solved by iterating on a preliminary guess until the relative error in the velocity component normal to the shock is less than a prescribed tolerance, ε .

At body points, two additional conditions are obtained by writing that the entropy is constant for a moving particle and by using the momentum equation for the tangential velocity. Again, the complete system is solved by iterating on a preliminary guess until the error in the distance between the body and the initial point on the characteristic is less than ϵ .

The computation is started by assuming a parabolic shape of the shock and prescribing a linear distribution of Mach numbers on the body. The shock is initially assumed at rest. The values of pressure (p), ρ , u, v, and S are computed behind the shock and at the body, and linearly distributed at the interior points of the mesh.

3. Presentation of the Numerical Results.

The computational program has options to output partial results at any stage as well as at the end of the run. These include the mesh coordinates, the pressure, density, velocity components,

entropy and Mach number at each mesh point. To get a direct feeling for the properties of the flow field, plots are necessary. Such plots, together with some typical numerical information, are obtained by processing the final output of the programs on a Stromberg-Carlson 4020 cathode tube display machine. Two pages of plots are printed, as shown in this report. It will be recalled that only two-dimensional symmetric and axisymmetric flows of a perfect gas are considered at this time.

In the first page, the physical nature of the flow is described by the free stream Mach number and the value of the ratio of specific heats, v. The shape of the body and the extent of the computed region are shown in the figures of the second page and in the lower left figure in the first page. The number of mesh intervals is indicated in the first page; the first number denotes the number of intervals between shock and body and the second number denotes the number of intervals along the body.

As an example, in Fig. 4 a mesh is shown which corresponds to the legend "4 BY 6 MESH". The mesh points on the body are marked by short lines pointing toward the origin of coordinates. The origin is the intersection of the upper boundary line and the centerline of the body.



FIG. 4

The standoff distance is expressed in arbitrary units. The abscissa of the stagnation point is expressed in the same units. The latter can be measured on the drawings, starting from the origin, and the unit length can thus be determined. As a rule, the unit length has a simple geometrical meaning in relationship with the geometry of the body. If the cross-section of the nose is a circle, the unit length is its radius; if the cross-section is an ellipse, the unit length is its major semi-axis; if the cross-section is a rectangle with a rounded shoulder, the unit length is the height of the rectangle. Finally, if the crosssection is a parabola, it is defined by

$$y^2 = 2(x+x_0) + 4$$

₽

Also, in the first page the number of time steps and the value of ϵ (called "TOLERANCE") used in the computation of shock and body points are shown. The values of the pressure printed in the first page are referred to the free stream pressure. The theoretical and computed values of the pressure at the stagnation point are printed, together with the relative error,

$$E = \frac{p_{st(comp)}^{-p} t(theor)}{p_{st(theor)}}$$

The temperatures printed in the first page are referred to the free stream temperature. Again, the theoretical and computed values of the temperature at the stagnation point are printed, together with the relative error,

$$E = \frac{T_{st(comp)} - T_{st(theor)}}{T_{st(theor)}}$$

The computed values of pressure, density and temperature at the body point where M=1, divided by the computed values of pressure, density and temperature at the stagnation point respectively are also printed, with their relative errors with respect to the theoretical ratios.

In the right-hand side of the first page, pressure, Mach number, and temperature along the body surface are plotted. The pressure is scaled to the computed pressure at the stagnation point and the temperature is scaled to the computed temperature at the stagnation point. The abscissae of these plots are angles, α (in degrees) measured between the centerline and a line joining the body point to the origin

(Fig. 5).

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In the bottom left of the first page and in the second page the body and shock geometries are repeated five times. In the figure of the first page, some streamlines are drawn. In a steady motion the



FIG. 5

streamlines can be defined either as lines of constant entropy or as lines of constant total pressure. When the first is used, S is defined as

$$S = ln(p/p_{o}) - \gamma ln(\rho/\rho_{o})$$

and the lines of constant entropy are spaced by 1/20 of the value of S at the stagnation point. Sometimes such a spacing does not provide enough information. In this case, lines of constant total pressure are used, spaced by 1/20 of the value of p at the stagnation point. As a proof of the equivalence of the two definitions, Fig. 6 shows, on the left, lines of constant entropy and, on the right, lines of constant total pressure for the same case, a twodimensional flow about a circle at a free stream Mach number of 4. When superposed, the two figures match perfectly.



STREAML INES

CONSTANT TOTAL PRESSURE LINES

The other figures are self-explanatory. To interpret them quantitatively, it must be kept in mind that the values of the Mach number on the M=constant lines are spaced .1 apart. The sonic line is drawn heavier than the other lines. The isobars correspond to constant values of p/p_{st} , .05 apart. The isopycnics correspond to constant values of $1/2_{st}$, .05 apart. The isotherms correspond to constant values of $1/2_{st}$, .05 apart. The isotherms correspond to constant values of T/T_{st} , .025 apart.

4. Discussion

In the present report, only results for the two-dimensional and the axisymmetric problems and for a perfect gas at constant Y are presented. In addition, one single program has been used for the computation of all the two-dimensional cases and one single program for all the axisymmetric cases, * regardless of the geometry of the body. In other words, no special effort was made to achieve a greater accuracy in cases of a challenging geometry (very blunt ellipses, flat-faced bodies) and we wished to explore the range of acceptability of such basic programs.

A. <u>Mesh-size effects</u>

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To study the effect of mesh-size on accuracy, a two-dimensional flow about a circle at M_{∞} =4 was computed five times, using the following meshes:

These programs are labelled 2E and 2F, respectively.

Run No.	199	200	201	202	203
Mesh	2 x 4	3 x 5	5 x 8	7 x 12	10 x 16
Total number of steps	160	240	400	560	800
Total time at final step	6.354	7.590	7.861	7.395	7.948

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(the time is scaled to $r_0 \sqrt{\rho_{\infty}} / p_{\infty}$, where r_0 is the radius of the circle and ρ_{∞} and p_{∞} are the free stream values of density and pressure, respectively). In Figs. 7 and 8, the time history of some representative parameters is shown for the coarsest and the finest mesh of the set above. In the top and bottom parts of the figures respectively, the standoff distance and the pressure at the stagnation point are plotted (the latter is scaled to the free stream pressure). The middle part is the plot of the logarithm (base 10) of the difference between the maximum and minimum values of the velocity of the shock points, divided by the free stream velocity. Since the time-dependent computation aims at reaching a steady state, the first and third functions should asymptotically tend to a constant value, and the second function should tend to $-\infty$. One can see from the graphs that the stagnation point pressure and the standoff distance reach a steady state in a relatively short time, whereas the shock wave is never perfectly at rest. This last feature is a result of the finiteness of the mesh, of the tolerance accepted in the iterations, and of the limited capacity of the computer. However, for all practical purposes, a



Fig.

1.3



PROBRAM 26. RUN NO 283 . H. 4.88

Fig. 8

shock wave whose velocity is five orders of magnitude smaller than the free stream velocity is a steady configuration.

No major differences can be noted between Fig. 7 and Fig. 8, despite the strong difference in mesh-size. The final results of all the cases mentioned above show, obviously, a better and smoother definition of values as the mesh-size is made finer and finer. The general trend of the curves, however, is the same in all cases. In addition, from a quantitative point-of-view, the differences between different cases at each point are surprisingly small.

One can conclude that, in a case where the body shape is fairly smooth, preliminary estimates can be made with a fast, inexpensive run using a very coarse mesh. To make the point clearer, the dependence of some relevant parameters on mesh-size is shown in Fig. 9. The first plot shows the logarithm (base 10) of the relative error in pressure and temperature, as computed at the stagnation point. These errors seem to level off at a value of about 10^{-4} , probably because of the worsening of round-off effects with decreasing mesh-size. The second plot shows similar errors in p, c, and T as computed at the critical point on the body. Here it must be noted that the critical point itself and all the attached values are computed by linear interpolation between adjacent points on the body. Therefore, the accuracy is affected in different ways for different mesh-sizes (the relative location of the critical point in a mesh interval is not the same in two different cases), and this explains a greater scattering in the plot. A line has

been drawn to represent an estimate of the errors and it shows the same trend as the two lines in the first plot. The third plot shows the standoff distance whose third significant figure seems rather hard to stabilize. Finally, the fourth plot shows the location of the critical points on the body and on the shock (measured in the polar frame of reference in degrees).

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FIG. 9

To get an idea of the price one has to pay for a greater accuracy, here is a comparison of computational times for the five cases above, taking the time for case 199 as the unit:

Run No.		199	200	201	202	203
Computational	time	1	2.5	10	28	73.5
(the time for	Run 203	is about	3 minute	es on	the CDC	6600).

In order to test the mesh-size effect on a body of a more complicated geometry, runs No. 169 and 270 were made. Their time histories are shown in Figs. 10 and 11, respectively. At the end of run 169 (400 steps), the standoff distance is not yet perfectly stabilized. A better situation is achieved in run 270, only because the computation runs for a longer physical time. In the final plots, the pressure and temperature distributions on the body are not substantially different between the two runs. The Mach number distribution on the body is quite different in the region of high curvature. This results from a seemingly intrinsic difficulty in computing velocities on the body where the curvature is high. More comments on this matter will be found under D. The general trends of constant Mach number lines, isobars, isopycnics, isotherms, and streamlines is the same in both cases and let us once again draw the conclusion that, for a preliminary, inexpensive evaluation of the shock layer, a coarse mesh can be used.

B. <u>Tolerance effect</u>

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Some runs were made with different values of the tolerance in the shock and body point iterations, namely, $\varepsilon = 10^{-4}$, 10^{-5} , and 10^{-6} . No appreciable differences were observed, and $\varepsilon = 10^{-5}$ was adopted uniformly throughout the present set of runs.

C. Number of time steps necessary to achieve convergence

In all runs, plots as in Fig. 10 were made. As a general rule, the stagnation pressure is the first parameter to become stabilized.



PROGRAM 22+ RUN NO 188 + M= 6.88

Fig. 10



Fig. 11



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Fig. 12



Fig. 13

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However, if one is interested in obtaining an accurate description of the entire flow field, it would be a mistake to use the stagnation pressure as a criterion for stopping the computation. The standoff distance takes much longer to stabilize. Moreover, the logarithmic plot of the shock wave velocity range shows that an overall stabilization of the shock wave is achieved only after a long time has elapsed. There are evidently a number of wavelets travelling over the computational region, which sometimes are hard to damp. From a physical point-of-view, an inviscid flow does not provide any mechanism for the damping of such wavelets, except on the shock wave itself. We find a similar behavior in our numerical computations since their artificial viscosity is purposely kept very low.

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No attempt is made at this time to analyze the propagation and damping of wavelets in the numerical computations in relation to the propagation and damping of sound waves in the present problem. Here we simply show the patterns obtained in a specific case, a two-dimensional flow field about a 5:1 ellipse at $M_{\odot} = 20$. Fig. 12 is obtained by stopping the computation at step 1000. The constant Mach number lines and the isobars seem to be pretty smooth, but the isopycnics and the isotherms are full of ripples. Fig. 13 is obtained at step 2000. Now all wavelets are practically damped. For the sake of completeness, Fig. 14 shows the time history of the run up to step 2000. Between step 1000 and step 2000 the shock velocity range drops by an order of magnitude.

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In a practical case, one must decide when to stop a computation by compromising between his need for accuracy and the available computational time. If a computation has to be stopped prematurely, smoother curves can be obtained by a crude averaging of values. Fig. 15 shows how the plots of Fig. 12 look like after averaging. Of course, these curves do not exactly fall on top of the curves of Fig. 13; however, they are qualitatively correct and probably good for all practical purposes.

Another important conclusion has been reached through the present numerical experimentation. For an arbitrarily prescribed set of initial conditions, the first phase of the transient is always very active. All the major changes take place in the first 5% of the total time necessary to achieve smooth results. Therefore, the discussion above is practically independent of the choice of initial conditions.

D. Mach number and bluntness effects

To study the effect of bluntness on accuracy, we have run several cases for bodies with an elliptical nose, the ratio of the two axes of the ellipse being taken as a measure of the bluntness. An ellipse with a bluntness of 1 is a circle; an ellipse with a bluntness of 6 is almost a flat-faced body with a rounded shoulder.

However, it turned out that a study of the bluntness effect could not be achieved independently of a study of the free stream Mach number effect. Fig. 16 shows the Mach number and bluntness



PROBRAM 22. RUN NO 317 . M-28.88

Fig. 14



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Fig. 15

parameters for the cases presented in this report (for both the two-dimensional and the axisymmetric problem).



No computations were made for a bluntness parameter greater than 6. * Below 6, some of the cases had to be discarded because the results were evidently poor (the relative errors mentioned in section 3 were too high; the plots were far from being smooth). Some other cases did not even run to completion. In general, it seems that, for a given free stream Mach number, the results worsen with increasing bluntness; the program fails for values of

See, however, a preliminary survey of flat-faced bodies under E.

the bluntness parameter above a limiting value. Such a limiting value is a function of the free stream Mach number; the lower the Mach number, the lower the limiting bluntness.

These limitations were, in a way, expected. Neither program has been tailored for handling high curvature effects on the body. The truncation errors due to the linear interpolations performed in computing the values on the wall grow excessively if the curvature of the particle paths and the consequent rate of change of physical parameters are too high. Such errors propagate within the shock layer and are steadily generated at each computational step. If we accept the hypothesis, mentioned under C, that the damping of the error waves is mostly due to a dissipative mechanism at the shock, we should conclude that a weak shock is less capable of producing damping than a strong shock. At a low Mach number, not only the error waves in the shock layer are harder to eliminate (and may even become unstable) but the shock itself is more sensitive to such perturbations and tends to wrinkle. It is interesting to note that, at very low Mach numbers, the shock wave is actually extremely sensitive to all perturbations, and it is hard to keep it stable in an experiment. However, it is not the intention here to suggest any quantitative correlation between the natural phenomenon and the present numerical effects. This is, at least, premature. It should be noted, also, that at a low Mach number the disturbance field to be computed becomes very large in comparison

to the body size. Consequently, the overall accuracy tends to deteriorate. We would rather say that it is surprising how far one can force the bluntness, and how close to 1 the Mach number can be taken, with still acceptable results, and that this can be done without increasing the number of mesh points over 200 and without providing any special treatment for high curvature walls.

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How to improve the situation at low Mach numbers and high bluntness is, at this time, hard to say. Any modification to the program which increases the artificial viscosity is not advisable. With a mesh as coarse as the ones used here (and let us recall that the object of the present program is to minimize the computation time), the artificial viscosity is bound to deface the whole flow pattern (when the flow tends to become steady, the time derivatives become smaller and smaller and eventually their effect is nullified by an equal and opposite contribution of the artificial terms).

One can note that the Mach number on the wall tends to oscillate more than the pressure and the temperature. This indicates that the weakest computation is that of the velocity on the body. In the present programs, the velocity is computed by using one of the momentum equations. The energy equation has been tried instead, with no success. The latter defines the square of the modulus of the velocity by the difference between the Lagrangian and the Eulerian derivatives of pressure. Truncation errors can

occasionally make such a square become negative, so that the computation halts. Some attempts to damp the oscillations in velocity at the wall only resulted in a general catastrophic worsening of the computation, except perhaps, in the case of a sharp corner where the wall region before the corner is practically disconnected from the wall region behind the corner. The matter is being studied further.

In the range of validity of these programs, some interesting parametric results can be obtained, which confirm and extend the ones available in Hayes and Probstein's book. Fig. 17 shows different shapes of shock waves and sonic lines for two-dimensional elliptical bodies at different Mach numbers; Fig. 18 does the same in the axisymmetric case. Each part of these figures deals with a given Mach number to show the bluntness effect. Figs. 19 and 20 show the Mach number effect on the shock wave and the sonic line for a parabolic body in the two-dimensional and the axisymmetric case, respectively. Figs. 21 and 22 show the Mach number effect on the shock wave and the sonic line for a circular nose in the two-dimensional and the axisymmetric case, respectively.





FIG.18

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FIG. 19

FIG. 20





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5. List of Cases Reported

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In this report, the results for the following cases are shown: Two-dimensional problem (program 2E):

A. Study of the mesh size effect

A1 A2 A3 A4 A5	Run No.	199 200 201 202 203	circular com. 1, 4
A6 A7	Run No.	169 270	flat-faced body well a rounded shoulder, M -6

B. Mach number effect on a circular nose

в1	Run No. 320	$M_{\infty} = 20$
B2	221	10
в3	(see A5)	* ~†
B4	222	3
В5	423	2
в6	324	1.7
B7	325	1.5

C. Elliptical nose at $M_{\infty} = 20$

C1	(see Bl)	bluntness	parameter	=	1
C 2	Run No. 215				2
С3	319				3
C4	316				4
C5	317				5
C6	318				6

D. Elliptical nose at $M_{\infty} = 10$

Dl	(see B2)	bluntness	parameter	=	1
D2	Run No. 230				2
D3	331				3
D4	332				4
D 5	233				5
D6	234				6

E. Elliptical nose at $M_{\infty} = 4$

El	(see A5)	bluntness I	parameter	=	1
E2	Run No. 210				2
E3	211				3
E4	212				4
E5	313				5

F. Elliptical nose at $M_{\infty} = 2$

Fl	(see B5)	bluntness parameter	=	1
F2	Run No. 340			2
F3	341			3

G. Parabolic nose

Gl	Run No.	450	M=20
G2		451	10
G3		352	4

Axisymmetric problem (program 2F):

B. Mach number effect on a spherical nose

вl	Run NO.	120	M=20
В2		121	10
в3		103	4
В4		122	3
в5		223	2
в6		124	1.7
B7		125	1.5

C. <u>Ellipsoid</u>, at $M_{\infty} = 20$

C1	(see Bl)	bluntness parameter =	1
C2	Run No. 115		2
С3	119		3
C4	116		4
C5	117		5
C6	118		6
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D. <u>E</u>	llipsoid,	at $M_{\infty} = 1$.0				
	ъl	(see B2)		blunt	ness	parameter =	1
	DI	USCE DE!	130				2
	D2	Run NO.	131				3
	D3	-					4
	D4	-	132				5
	D 5		133				6
	D6		134				
E.]	Ellipsoid,	at $M_{\infty} = 0$	4				_
	-1	(coo B3)	blun	tness	parameter =	1
	EI	(See D)	,				2
	E2	Run NO.	111				3
	E3		111				4
	E4		112				5
	E5		113				
			2				
F.	Ellipsoid	, at $M = \infty$	2				
	Fl	(see B5)	blun	tness	s parameter =	
	F2	Run No.	140				2
	F3		141				3
	• •	_					
G.	Paraboloi	<u>d</u>					
	Gl	Run No.	150	м _∞	20		
	G2		151		10		
	G3		152		4		
				a r oi	nded	shoulder. M	=10.5
н.	Flat-face	d cylinde	er with	<u>a 100</u>	indea		20
	(·=ratio	of the sl	noulder	radi	us to	the cylinder	r radius)
	Hl	Run No.	80	• =	.50		
	H2		81		.25		
	112 113		82		.10		
	113		83		.05		
	<u>n4</u>		00				
The	H-series	of runs	has been	n mad	e to	test the com	putational
rechniqu	e against	a recent	ly issu	ed se	t of	experimental	data
(Ref . 5	and 6).	The comp	uted pr	essur	e dis	stributions o	n the body,

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when expressed as a function of the arc-length, fall exactly on

top of the experimental curves. The shock waves also fit exactly the experimental shapes. Some of the computed patterns are rather irregular, particularly when σ is very small, but this should be expected, as we said in section 4. Note, for example, that the body shape assumed by the machine in case H4 is only a rough approximation to a flat-faced cylinder with a rounded shoulder, $\sigma = .05$.

6. Acknowledgements

The initial work on the two-dimensional and the axisymmetric problem for a perfect gas was performed at the General Applied Science Laboratories under the sponsorship of the Advanced Research Projects Agency in 1965 (Refs. 1 and 2). The extension to threedimensional flows of a perfect gas was also performed at the General Applied Science Laboratories under the sponsorship of the Sandia Corporation, in 1966 (Ref. 4). More recently, the Sandia Corporation sponsored additional work, as a result of which the problem was reformulated in new frames of reference; the pertinent analysis, performed last year at the General Applied Science Laboratories, is not available in the open literature. The computations shown in the present report make use of the latter frames of reference, with some additional features added in the last few

months. The latter research, as well as the parametric study partially contained in the present report, has been performed at the Polytechnic Institute of Brooklyn under the sponsorship of the Office of Naval Research under Contract Nonr 839(34).

The plots have been obtained by transferring the pertinent information (as computed by the CDC 6600 machine at the Courant Institute of the New York University) to the Stromberg-Carlson 4020 plotter of the Polytechnic Institute of Brooklyn and by using an additional plotting program. I am glad to acknowledge the dedicated and efficient assistance of Mr. Martin Tillinger of the Polytechnic Institute of Brooklyn in the delicate manipulation necessary to obtain these plots.

7. <u>References</u>

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- 4. Moretti, G. and Bleich, G., Three-dimensional flow around blunt bodies. AIAA J., <u>5</u>, 1557, 1967.
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6. Inouye, M., Marvin, J.G., and Sinclair, A.R., Comparison of experimental and theoretical shock shapes and pressure distributions on flat-faced cylinders at Mach 10.5. NASA TN D-4397, 1968.

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 CRITICAL PRESSURE RATIO
 0.5755
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10 BY 16 MESH. DOD STEPS TOLERANCE: 0.00010

FREE STREAM HACH NUMBERs 4.00. BANNAS 1.40

THEOR. STADUATION PRESSURE: 21.000 CON". STADUATION PRESSURE: 21.072 RELATIVE ERROR: 0.00017

THEOR. STADNATION TEMPERATURE: 4.200 CONF. STADNATION TEMPERATURE: 4.200 RELATIVE ERIDR: 0.00003

 CRITICAL PRESSURE RATIO: 0.5203
 CREL. ERROR: 0.0007 3

 CRITICAL DENSITY
 RATIO: 0.6340
 [REL. ERROR: 0.0001 3

 CRITICAL TEMPERAT.RATIO: 0.0303
 (REL. ERROR: 0.0006 3

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TWO DIMENSIONAL STMMETRIC BLUNT BODY RUN NO 169



STANDOFF DISTANCE=0.8394 ABSCISSA OF STAGNATION POINT=-1.8394



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THE DIMENSIONAL SYMMETRIC BLUNT BODY - RUN NO 270

10 BY 18 MESH. 800 STEPS TOLERANCE . 0.00010

FREE STREAM MACH NUMBER: 5.00. BAMMAR 1.40

THEOR. STAGNATION PRESSURE: 48.815 COMP. STAGNATION PRESSURE: 48.790 Relative error: 0.00053

THEOR. STAGNATION TEMPERATURE: 0.200 Comp. Stagnation temperature: 0.199 Relative error: 0.00015

CRITICAL PRESSURE RATIO= 0.6008 (REL. ERROR= 0.1487) CRITICAL DENSITY RATIO= 0.6075 (REL. ERROR= 0.1003) CRITICAL TEMPERAT.RATIO= 0.8700 (REL. ERROR= 0.0440)

STANDOFF DISTANCE=0.8777 ABSCISSA OF STAGMATICN POINT= - 1.00000



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TWO DIMENSIONAL SYMMETRIC BLUNT BODY RUN NO 221



FREE STREAM MACH NUMBER=10.00+ BAMMA= 1.40

THEOR. STAGNATION PRESSURE=129.217 COMP. STAGNATION PRESSURE=129.237 RELATIVE ERROR= 0.00016

THEOR. STAGNATION TEMPERATURE= 21.000 COMP. STAGNATION TEMPERATURE= 21.001 RELATIVE ERROR= 0.00005

CRITICAL PRESSURE RATIO= 0.5301 (REL. ERROR= 0.0035) CRITICAL DENSITY RATIO= 0.6355 (REL. ERROR= 0.0024) CRITICAL TEMPERAT.RATIO= 0.6342 (REL. ERROR= 0.0011)

STANDOFF DISTANCE=0.4005 ABSCISSA OF STAGNATION POINT=-1.50000



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TWO DIMENSIONAL SYMPETHIC BUUNT BODY - RUN NO 325

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10 BY 20 MESH, 1000 STEPS TOLERANCE= 0. 000010

FREE STREAM MACH NUMBER=20.00+ GAMMA= 1.40

THEOR. STAGNATION PRESSURE=515.484 Comp. Stagnation pressure=515.524 Relative error= 0.00009

THEOR. STAGNATION TEMPERATURE: 81.000 Comp. Stagnation temperature: 81.002 Relative error: 0.00002

CRITICAL PRESSURE RATIO=0.5242 (REL. ERROR= 0.0077) CRITICAL DENSITY RATIO=0.6301 (REL. ERROR= 0.0030) CRITICAL TEMPERAT.RATIO=0.0319 (REL. ERROR= 0.0017)

STANDOFF DISTANCE=0.5090 AD30133A OF STACHATION POINT= -1.30000



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TWO DIMENSIONAL SYMMETRIC BLUNT BODY RUN NO 319

 10 BY 20 MESH, 2000 STEPS
 TOLERANCE= 0.000010

 FREE STREAH HACH NUMBER-20.00. OANHA= 1.40

 THEOR. STAONATION PRESSURE=515.494

 CONP. STAONATION PRESSURE=515.370

 RELATIVE EIROR= 0.00022

 THEGR. STAONATION TEMPERATURE= 01.000

 COMP. STAONATION TEMPERATURE= 01.000

 CRITICAL PRESSURE RATIO=0.5773
 (REL. ERAGR= 0.0340.)

 CRITICAL DENSITY RATIO=0.6009
 (REL. ERAGR= 0.0203.)

 STANDOPP DISTANCE=0.00005
 STAONATION POINT= -1.13333









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TWO DIMENSIONAL SYMMETRIC DEUNT DODY - RUN NO 310

10 BY 20 MESH+ 2000 STEP3 TOLERANCE=0.000010

FREE STREAM MACH NUNDER=20.00+ DAMMA= 1.40

THEOR. STACHATION PRESSURE=515.434 COMP. STACHATION PRESSURE=515.302 Relative Error= 0.00035

THEOR. STAGNATION TEMPERATURE: 01.000 COMP. STAGNATION TEMPERATURE: 00.932 Relative error: 0.00010

CRITICAL PRESSURE RATIO=0.5004 (REL. ERROR= 0.1100) CRITICAL DENSITY RATIO=0.6707 (REL. ERROR= 0.0706) CRITICAL TEMPERAT.RATIO=0.6005 (REL. ERROR= 0.0422)

STANDOFF DISTANCE=0.7764 ADDCI55A OF STAGNATION POINT= 0.00007



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TWO DIMENSIONAL SYNMETRIC BLUNT DODY - RUN NO 230

10 BY 20 ME3H. 1000 STEPS TOLERANCE=0.000010

FREE STREAM MACH NUMBER:10.00+ BAMMA: 1.40

THEOR. STACNATION PRESSURE=129.217 Comp. Stacnation Pressure=129.217 Relative Ergor= 0.00000

THEOR. STACNATION TEMPERATURE: 21.000 CONF. STABNATION TEMPERATURE: 21.000 Relative epror: 0.00000

CRITICAL PRESSURE RATIO:0.5237 (REL. ERADR: 0.0007.) CRITICAL DENSITY RATIO:0.6207 (REL. ERADR: 0.0007.) CRITICAL TENDERAT.RATIO:0.0017 (REL. ERROR: 0.0020.)

STANDOFF DISTANCE=0.5012 AU3CISSA OF STACNATICH POINT= 1.30000



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TWO DIMENSIONAL SYMMETRIC PLUNT DODY RUN NO 332

10 BY 20 MESH, 1000 STEPS TOLERANCE= 0.000010

FREE STREAM MACH NUMBER=10.00+ GAMMA= 1.40

THEOR, STAGNATION PRESSURE=123.217 Comp. Stagnation Pressure=129.152 Relative error= 0.00050

THEOR. STABNATION TEMPERATURE= 21.000 COMP. STAGNATION TEMPERATURE= 20.907 Relative error= 0.00014

CRITICAL PRESSURE RATIO=0.5539 (REL. ERROR= 0.0405) CRITICAL DENSITY RATIO=0.6554 (REL. ERROR= 0.0339) CRITICAL TEMPERAT.RATIO=0.0452 (REL. ERROR= 0.0142)

STANDOFF DISTANCE+0.7613 AdSCISSA OF STADNATION POINT= -0.37000



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THO DIMENSIONAL SYMMETRIC DUNT BODY - RUN NO 223

10 BY 20 MESH+ 1000 STEPS TOLERANCE+0.000010 FREE STREAM MACH NUMBER+10.00+ GAMMA+ 1.40 THEOR, STACHATION PRESSURE+120.217

COMP. STADIATION PRESSURE=129.073 RELATIVE ERROR= 0.00112

THEOR, STACHATION TEMPERATURE: 21.000 Comp. Stachation temperature: 20.003 Relative error: 0.00032

 CRITICAL PRESSURE RATIO=0.5554
 (HEL. ENROR= 0.0532)

 CRITICAL DENSITY
 RATIO=0.6514
 (HEL. ENROR= 0.0275)

 CRITICAL TEMPERAT.RATIO=0.6341
 (HEL. ERROR= 0.0249)

STANDOFF DISTANCE+0.7759 ADSCISSA OF STACNATION POINT+ +1.00000



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THE DIMENSIONAL SYMMETRIC BLUNT BODY - RUY NO 234

10 BY 20 MESH. 1000 STEPS TOLERANCE= 0.000010

FREE STREAM MACH NUMBER+10.00+ BANMA+ 1.40

THEOT, STADMATION PRESSURE=129.217 Comp. Stachation pressure=129.045 Relative error= 0.00133

THEOR, STAGNATION TEMPERATURE: 21.000 COMP. STAGNATION TEMPERATURE: 20.992 Relative error: 0.00030

CRITICAL PRESSURE RATIO: 0.5877 (REL. ERROR: 0.1123) CRITICAL DENGITY RATIO: 0.6772 (REL. ERROR: 0.0502) CRITICAL TEMPERAT.RATIO: 0.6378 (REL. ERROR: 0.0114)

STANDOFF DISTANCE+0.7007 ABSCISSA OF STADNATION FOINT+**0.80037



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10 DY 20 MESH. 1000 STEPS TOLERANCE= 0. 000010

FREE STREAM MACH NUMBER: 4.00. DAMMA: 1.40

THEOR, STACHATION PRESSURE: 21.009 COMP. STACHATION PRESSURE: 21.055 Relative error: 0.00034

THEOR, STADNATION TEMPERATURE= 4.200 COMP. STADNATION TEMPERATURE= 4.109 Relative error= 0.00018

CRITICAL PRESSURE RATIO= 3.5200 (REL. ERROR= 0.0141) CRITICAL DENSITY RATIO= 0.6274 (REL. ERROR= 0.0103) CRITICAL TEMPERAT.RATIO= 3.6301 (REL. ERROR= 0.0029)

STANDOFF DISTANCE=0.7643 AUGCIESA C/ STADNATION POINT= "1.30000



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STANDOFF DISTANCE=0.0741 AUGCISSA OF STAGNATION POINT= 1.13333



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10 BY 20 HESH+ 2000 STEPS TOLERANCE= 0.000010
FREE STREAM MACH NUMBER= 2.00+ BANMA= 1.40
THEOR. STAGNATION PRESSURE= 5.640
COHP. STAGNATION PRESSURE: 5.607
RELATIVE ERROR= 0.00599
THEOR. STAGNATION TEMPERATURE: 1.800
COMP. STAGNATION TEMPERATURE= 1.797
RELATIVE ERROR= 0.00171
CRITICAL PRESSURE RATIO= 0.5332 (REL. ERROR= 0.0094
CRITICAL DENSITY RATIO= 0.6375 (REL. ERROR= 0.0055
CRITICAL TEMPERAT, RATIO: 0.0365 [REL. ERROR: 0.0030
STANDOFF DISTANCE=1.6597
ADSCISSA OF STAGNATION POINT= - 0.70000



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TWO DIMENSIONAL SYMMETRIC BLUNT BODY - RUN NO 450

• BY 13 MESH, BDD STEPS TOLERANCE=0.000010 FREE STREAM MACH NUMBER=20.00. GAMMA= 1.40 THEOR. STAGNATION PRESSURE=515.404 COMP. STAGNATION PRESSURE=515.757 RELATIVE ERROR= 0.00053 THEOR. STAGNATION TEMPERATURE= 01.000 COMP. STAGNATION TEMPERATURE= 01.000 STANDOFP DISTANCE=0.4644 ADSCISSA OF STAGNATION POINT= -2.00000



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TVO DIMENSIONAL SYMMETRIC BLUNT BODY RUN NO 451

S BY 13 MESH. BOD STEPS TOLERANCE= 0.000010

FREE STREAM MACH NUMBER=10.00. BAMMA= 1.40

THEOR. STAGNATION PRESSURE=129,217 CONP. STAGNATION PRESSURE=129.287 RELATIVE ERROR= 0.00055

THEOR. STAGNATION TEMPERATURE= 21.000 COMP. STAGNATION TEMPERATURE= 21.003 RELATIVE ERROR= 0.00015

CRITICAL PRESSURE RATIO= 0.5241 (REL. ERFOR= 0.0070) CRITICAL DENSITY RATIO= 0.6303 (REL. ERFOR= 0.0057) CRITICAL TEMPERAT.RATIO= 0.8315 (REL. ERFOR= 0.0022)

STANDOFP DISTANCE=0.4914 ABSCISSA OF STADNATION FOINT= -2.00000



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TWO DIMENSIONAL SYMMETRIC BLUNT BODY RUN NO 352

BY 13 MESH. 1400 STEPS TOLERANCE . 0.00010

PREE STREAM MACH NUMBER: 4.00. DAMMA: 1.40

THEOR. STADNATION PRESSURE: 21.068 COMP. STADNATION PRESSURE: 21.079 RELATIVE ERROR: 0.00052

THEOR. STABNATION TEMPERATURE: 4.200 COMP. STABNATION TEMPERATURE: 4.201 RELATIVE ERROR: 0.00015

 CRITICAL PRESSURE RATIO= 0.5218
 (REL. ERROR= 0.0123)

 CRITICAL DENSITY
 RATIO= 0.6283
 (REL. ERROR= 0.0089)

 CRITICAL TEMPERAT.RATIO= 0.8304
 (REL. ERROR= 0.0035)

STANDOFF DISTANCE=0.7130 ABSCISSA OF STAGNATION POINT= -2.00000







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THEOR. STACNATION PRESSURE=129.217 COMP. STACNATION PRESSURE=129.240 RELATIVE EFROR= 0.00010

THEOR. STAGNATION TEMPERATURE= 21.000 Comp. Stagnation temperature= 21.001 Relative Epror= 0.00005

CRITICAL PRESSURE RATIO: 0.5305 (REL. ERROR: 0.0042) CRITICAL DENSITY RATIO: 0.6340 (REL. ERROR: 0.0016) CRITICAL TEMPERAT.RATIO: 0.8355 (REL. ERROR: 0.0026)

STANDOFF DISTANCE=0.1359 ABSCISSA OF STAGNATION POINT="1.50000



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18 BY 17 MESH. 1008 STEPS TOLERANCE= 0. 800010

PREE STREAM MACH NUMBER=10.00. DAMMA= 1.40

THEOR. STAGNATION PRESSURE=129.217 COMP. STAGNATION PRESSURE=129.202 RELATIVE ERROR= 0.00011

THEOR. STAGNATION TEMPERATURE: 21.000 COMP. STAGNATION TEMPERATURE: 20.939 RELATIVE ERROR: 0.00003

CRITICAL PRESSURE RATIO= 0.5595 (REL. ERROR= 0.0590) CRITICAL DENSITY RATIO= 0.6590 (REL. ERROR= 0.0379) CRITICAL TEMPERAT.RATIO= 0.6503 (REL. ERROR= 0.0204)

STANDOFF DISTANCE-0.3487 ABSCISSA OF STADNATION POINT= -1.00000

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 THEOR, STADNATION PRESSURE+129.217

 COMP, STADNATION PRESSURE+129.253

 RELATIVE ERROR+ 0.00020

 THEOR, STADNATION TEMPERATURE+ 21,000

 0.2

 COMP, STAONATION TEMPERATURE+ 21,002

 RELATIVE ERROR+ 0.00000

 0.0

 CRITICAL PRESSURE RATIO+0.5194 (REL. ERROR+ 0.0100)

 CRITICAL DENSITY RATIO+0.6732 (REL. ERROR+ 0.0122)

 CRITICAL TEMPERAT.RATIO+0.6735 (REL. ERROR+ 0.0148)

STANDOFF DISTANCE+0.1504 ABBCISSA OF STADNATION FOINT+-2.00000



STREAML INES





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TEMPERATURE DISTRIBUTION ON THE BODY





STREAML INES







60.0

60.0

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TEMPERATURE DISTRIBUTION ON THE GODY

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AXIS METRIC BLUNT BOOY

RUN NO 83

10 DY 20 MESH+ 1000 STEPS - TOLERANCE+0.000010

FREE STREAM MACH HUNDER: 10,50 - DAMMAR 1.40

THEOR. STACHATION PROSOURE+142.414 COMP. STACHATICH PRESSURE+142.034 RELATIVE ERROR+ 0.00207

THEO4, STADUATION TEMPERATURE: 23.030 COMP. STADUATION TEMPERATURE: 23.032 Relative Error: 0.00077

CHITICAL PRESSURE RATIO=0.7433 (REL. ERROR= 0.4079) CRITICAL DENSITY RATIO=0.0034 (REL. ERROR= 0.2705) CRITICAL TEMPERAT.RATIO=0.0235 (REL. ERROR= 0.1032)

STANDOFF DISTANCE+0.4616 Adscissa of Stadnation Point+ - 1.00000



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