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THE INVERSE PROBLEM OF THE ANNULAR AIRFOIL
AND DUCTED PROPELLER

by

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NOTATION

a	Duct chord
b(z)	Annular-airfoil inner-surface ordinate
c(z)	Mean line ordinate of the duct section measured from the nose-tail line
C_p	$[p(x_d, z) - p_\infty] / \frac{1}{2} \rho V^2$, pressure coefficient
E(k)	Complete elliptic integral of the second kind
h	$(a/2 R_d)$ chord-diameter ratio of the duct
K(k)	Complete elliptic integral of the first kind
k	Modulus of the elliptic integrals
$p(x_d, z)$	Local pressure on the annular airfoil
p_∞	Ambient pressure at infinity
q	Ring-source strength
R_d	Duct radius
(r, ϕ , z)	Cylindrical coordinates
s(z)	Half-thickness ordinate of the duct section
u(z)	Annular-airfoil outer-surface ordinate
V	Free-stream velocity
w_a	Axial component of induced velocity
w_o	Component of free-stream velocity in direction of duct axis
w_r	Radial component of induced velocity
w_{x_d}	Taylor wake fraction at the duct
x	Radial coordinate nondimensionalized by the propeller radius
α	Angle of attack of a duct section
γ	Ring-vortex strength
ρ	Mass density

Subscripts

d	Duct	q	Ring source
p	Propeller	γ	Ring vortex

Note: Many functions are defined in the text.

ABSTRACT

A computer program is presented that calculates the annular airfoil shape from a given pressure distribution. A brief review of the theory of the inverse problem of the annular airfoil is also presented. The distortion of the duct shape by the presence of an axisymmetric body or a propeller may be taken into consideration. Calculations show that for a given pressure distribution, the propeller loading and location affect the duct shape.

ADMINISTRATIVE INFORMATION

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1. INTRODUCTION

Annular airfoils have been used for many years¹ as shroud rings around marine propellers. These ducts have been used for the acceleration of the velocity at the propeller (Kort nozzles); this has the effect of increasing the efficiency when the unit is heavily loaded. Impetus in the last few years has been given for their use to increase the thrust of propellers during takeoff for short and vertical takeoff aircraft. Also, in naval architecture, ducts that decelerate the flow at the propeller (pump-jets) have found application for delaying cavitation on the propeller.

A number of experimental and theoretical investigations have been conducted on the annular airfoil and the ducted propeller. A general review of these studies has been made by Sacks and Burnell² so only the pertinent and more recent results will be reviewed here. In most of the theoretical investigations, a distribution of ring vortices and ring sources is used which lie on a cylinder of diameter representative of the duct and of length equal to the duct length.³⁻⁷ The use of this mathematical model implies that the boundary conditions are linearized and are satisfied on the representative cylinder and not on the duct surface. A

¹References are listed on page 46.

nonlinear theory for the annular airfoil has been presented by Chaplin⁸ where the annular airfoil is represented by a system of ring vortices lying on the surface of the airfoil. He includes both the static and the free-flight cases.

In the references just given, the direct problem of the annular airfoil is considered, i.e., given an annular airfoil shape, determine the pressure distribution and forces on the duct. Reference 9, however, presents the theory for the inverse problem, i.e., given a pressure distribution, determine the annular airfoil shape.

Both the direct and the inverse problems require the solution of a singular integral equation, the first for the ring vortex distribution, and the second for the ring source distribution. Another approach given in References 10, 11,* 12, and 13 is to assume the ring vortex strength and, if the effect of thickness is considered, to assume the thickness distribution. The disadvantages of this procedure are that the circulation must be specified, which is not a physical property of the annular airfoil, and it is not possible to tell a priori whether the pressure distribution or the shape will be satisfactory.

The usefulness of the inverse problem is to delineate annular airfoil shapes which will operate satisfactorily for a given flow condition. This is to say that in the presence of a propeller producing a given thrust, a duct shape can be determined which will not separate in air or water, nor cavitate in water. Both separation and cavitation are, of course, real fluid effects, so some criteria for the pressure distribution must be specified which will indicate satisfactory operation. For instance, for cavitation a minimum pressure coefficient would be assumed, and for separation a maximum rate of change of the pressure coefficient would be assumed.

This report presents a computer program and representative type calculations based on the inverse problem presented in Reference 9. The program allows the inclusion of an axisymmetric perturbation velocity so

* There is a difference in definition of the direct and inverse problem between References 9 and 11.

that the duct shape can be determined in the presence of a propeller or any body which is axisymmetric to the duct. Only the average effect of the propeller can be considered, however, as including the effect of a finite number of blades would imply that the annular airfoil can change shape as the propeller rotates.

The following discussion is divided into three main sections. The development of the theory is reviewed briefly, then the computer program is described and presented, and finally, computed duct shapes are presented for comparison. Calculations are also included which show how the shape of the duct changes with variation in propeller loading.

2. LINEARIZED THEORY OF THE ANNULAR AIRFOIL

Although the linearized theory of the annular airfoil has been adequately developed in the reference cited, a brief sketch of the development will be repeated here for the sake of completeness.

As was stated in the Introduction, the method of singularities is used for the representation of the flow field about the annular airfoil. The mathematical model used is a distribution of ring vortices and ring sources lying on a cylinder having a length equal to the length of the annular airfoil and a diameter representative of the diameter of the annular airfoil. The use of this model for the inverse problem necessitates a number of assumptions:

1. The real fluid is inviscid and incompressible and no separation occurs on the annular airfoil.
2. Body forces such as gravity are negligible.
3. The freestream flow is axisymmetric and steady, but an axisymmetric disturbing velocity may exist. An implication here is that the static condition is not considered.
4. The annular airfoil is axisymmetric and has finite length.
5. The distribution of ring vortices and ring sources along a cylinder does, indeed, represent the annular airfoil. This implies that the boundary conditions are linearized.

2.1 BOUNDARY CONDITIONS

The coordinate system used is a cylindrical system (r, ϕ, z) whose polar-axes are located at the trailing edge and z -axis coincides with the centerline of the annular airfoil. For convenience, the axial coordinate is nondimensionalized by the chord a , and the radial coordinate by the representative radius R_d , of the annular airfoil. Figure 1 gives the coordinate system and Figure 2 is a delineation of the system.

The pressure distribution on the annular airfoil is assumed known and the cross section of the airfoil is calculated in terms of a camber distribution $c(z)$, a thickness distribution $s(z)$, and an angle of attack α , as follows:

$$\begin{aligned}u'(z) &= c'(z) + s'(z) + \tan \alpha \\b'(z) &= c'(z) - s'(z) + \tan \alpha\end{aligned}\tag{2.1.1}$$

where $u'(z)$ is the slope of the outer surface and $b'(z)$ is the slope of the inner surface of the annular airfoil.

The boundary conditions to be satisfied are that the normal velocity must be zero on the surface of the annular airfoil and the Kutta condition must be satisfied at the trailing edge. In linearized theory¹⁴ this means that the radial velocity at the representative cylinder must be equal to the slope of the section, or

$$\frac{w_r}{V}(x_d \pm 0, z) = -(1 - w_{x_d}) [c'(z) \pm s'(z) + \tan \alpha]\tag{2.1.2}$$

where for convenience for the naval architect the wake $(1 - w_{x_d}) = \frac{w_0}{V}$ is introduced. At the trailing edge the radial velocity is zero, or

$$\frac{w_r}{V}(x_d \pm 0, 0) = 0$$

In the symbol \pm , the $+$ sign denotes the outer surface of the annular airfoil and the $-$ sign denotes the inner surface.

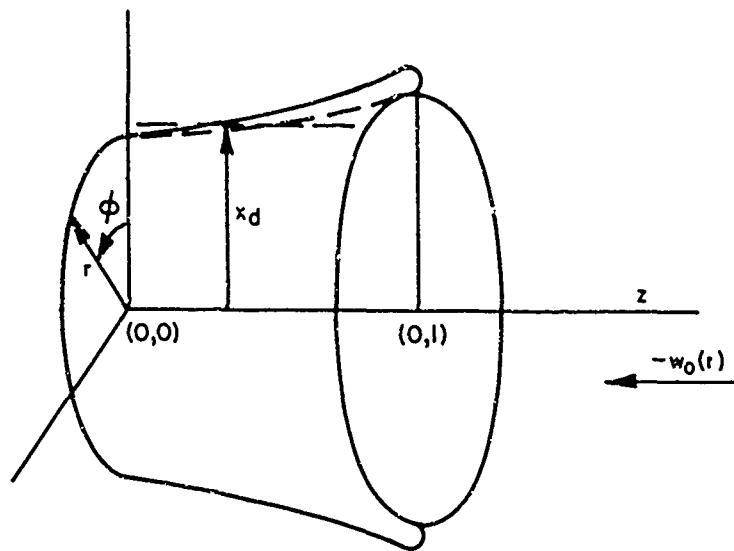


Figure 1 - The Annular Airfoil Coordinate System

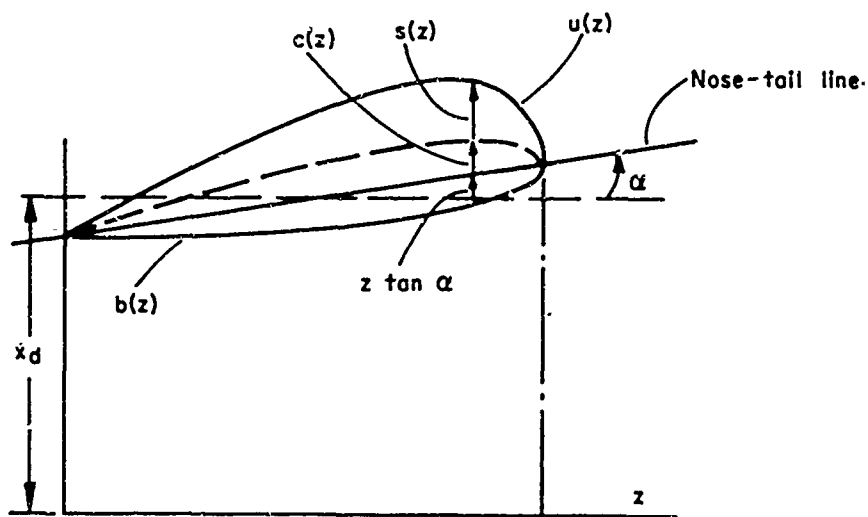


Figure 2 - Delineation of the Annular Airfoil Section

2.2 DERIVATION OF THE EQUATION FOR THE THICKNESS DISTRIBUTION

To calculate the thickness distribution, camber distribution, and the angle of attack of the annular airfoil from the pressure distribution, it is necessary to obtain the strength of the ring vortices and ring sources. The ring vortex and ring source strengths must be of sufficient magnitude that they induce radial velocities that satisfy the given boundary conditions. By substituting the equations for the radial velocity induced by the ring vortices and sources into Equation [2.1.1], an expression for the slope of the thickness distribution is obtained in terms of known quantities.

The nondimensional elementary circulation of the ring vortices is represented by $\gamma(z)$ and the strength of the ring sources by $q(z)$. Since the velocities are linear, they are additive; thus the radial velocities along the annular airfoil are found to be

$$\frac{w_r}{V}(x_d \pm 0, z) = \left[\frac{w_r}{V}(x_d, z) \right]_{\gamma} + \left[\frac{w_r}{V}(x_d \pm 0, z) \right]_q + \left[\frac{w_r}{V}(x_d, z) \right]_p \quad (0 \leq z \leq 1) \quad [2.2.1]$$

$[w_r(x_d, z)]_{\gamma}$ = the radial velocity induced on the annular airfoil by the ring-vortex system,

$[w_r(x_d \pm 0, z)]_q$ = the radial velocity induced on the annular airfoil by the ring-source system, and

$[w_r(x_d, z)]_p$ = the radial velocity induced on the annular airfoil by a propeller or any other singularities in the flow.

The expressions for these velocity components have been derived in Reference 14 and are as follows:

$$\left[\frac{w_r}{V}(x_d, z) \right]_{\gamma} = -\frac{1}{4\pi} \int_0^1 \frac{\gamma(z') k}{(z - z')} \left\{ 4h^2 (z - z')^2 [K(k) - E(k)] - 2E(k) \right\} dz' \quad [2.2.2]$$

$$\left[\frac{w_r}{V} (x_d \pm 0, z) \right]_q = \frac{h}{2\pi} \int_0^1 q(z') k [K(k) - E(k)] dz' \pm \frac{1}{2} q(z) \quad [2.2.3]$$

where $h = \frac{a}{2R_d}$, $k^2 = \frac{1}{h^2(z - z')^2 + 1}$, and $K(k)$ and $E(k)$ are complete elliptic integrals of the first and second kind, respectively. The symbol \oint denotes a Cauchy principal value integral.

By substituting these expressions in Equations [2.1.2] and [2.2.1], the following equation is obtained.

$$\begin{aligned} & - \frac{1}{4\pi} \oint_0^1 \frac{\gamma(z')}{(z-z')} k \{4h^2(z-z')^2 [K(k) - E(k)] - 2E(k)\} dz' \\ & + \frac{h}{2\pi} \int_0^1 q(z') k [K(k) - E(k)] dz' \pm \frac{1}{2} q(z) = \\ & - \left(1 - w_{x_d}\right) [c'(z) \pm s'(z) + \tan \alpha] - \left[\frac{w_r}{V} (x_d, z) \right]_p \end{aligned} \quad [2.2.4]$$

Since the integrals occurring in this equation have only one sign and since the radial velocity induced by the propeller on the annular airfoil does not change sign from one side of the foil to the other, it may be concluded that the \pm signs go with the \mp signs and hence

$$q(z) = - 2 \left(1 - w_{x_d}\right) s'(z) \quad [2.2.5]$$

This equation gives the relationship between the ring source strength and the slope of the thickness distribution. The next step is to find the total velocity tangent to the airfoil in terms of the ring vortex strength, ring source strength, and the strength of any other singularity present in the flow such as a propeller. Since the annular airfoil is theoretically replaced by a cylinder, the velocity in question is just the axial velocity. The pressure distribution is related to the induced velocities by means of the linearized Bernoulli equation, i.e.,

$$C_p = \frac{p(x_d, z) - p_\infty}{\frac{1}{2} \rho V^2} \approx 2 \frac{w_a}{V} (x_d, z) \quad [2.2.6]$$

where $p(x_d, z)$ = the local pressure on the annular airfoil,
 p_∞ = ambient pressure at infinity,
 $w_a(x_d, z)$ = the axial component of induced velocity on the annular airfoil, and
 ρ = the mass density.

Within linearized theory all velocities are additive; thus the total axial velocity is given by

$$\left[\frac{w_a}{V} \right]_{\text{total}} = \left[\frac{w_a(z)}{V} \right]_\gamma + \left[\frac{w_a(z)}{V} \right]_q + \left[\frac{w_a(z)}{V} \right]_p - \frac{w_o}{V} \quad [2.2.7]$$

where

- $\left[\frac{w_a(z)}{V} \right]_\gamma$ = the axial velocity induced on the annular airfoil by the ring-vortex system,
- $\left[\frac{w_a(z)}{V} \right]_q$ = the axial velocity induced on the annular airfoil by the ring-source system,
- $\left[\frac{w_a(z)}{V} \right]_p$ = the axial velocity induced on the annular airfoil by a propeller or any other singularity in the flow, and
- $\left[\frac{w_o}{V} \right]$ = speed of advance of the duct.

From Equation [2.2.2] the perturbation velocity distribution w_a is obtained as

$$\frac{w_a(z)}{V} = \left[\frac{w_a(z)}{V} \right]_{\text{total}} + \left[\frac{w_o}{V} \right] = \left[\frac{w_a(z)}{V} \right]_\gamma + \left[\frac{w_a(z)}{V} \right]_q + \left[\frac{w_a(z)}{V} \right]_p \quad [2.2.8]$$

It is, of course, not possible to determine the section shape in the presence of a propeller with a finite number of blades since the pressure distribution is essentially time-dependent; however, it is possible to consider the average effect of the propeller which corresponds to an infinitely bladed propeller. Under this restriction the duct circulation distribution, defined as $\gamma(z)$, is a function of z only and justifies the use of the following relationship from Reference 14 for the axial velocity induced on the duct by the ring vortex distribution and ring

source distribution, respectively:

$$\left[\frac{w_a(x_d \pm 0, z)}{V} \right]_y = \frac{h}{2\pi} \int_0^1 \gamma(z') k [K(k) - E(k)] dz' \mp \frac{1}{2} \gamma(z) \quad [2.2.9]$$

$$\left[\frac{w_a(x_d, z)}{V} \right]_q = \frac{2}{\pi} \int_0^1 \frac{q(z')}{(z-z')} k E(k) dz' \quad [2.2.10]$$

or in terms of Equation [2.2.5]

$$\left[\frac{w_a(x_d, z)}{V} \right]_q = - \frac{(1-w_{x_d})}{\pi} \int_0^1 \frac{s'(z')}{(z-z')} k E(k) dz'$$

If Equations [2.2.9] and [2.2.10] are substituted into Equation [2.2.8],

$$\begin{aligned} \left[\frac{w_a(z)}{V} \right]_{\pm} &= \frac{1}{2} C_p = \frac{h}{2\pi} \int_0^1 \gamma(z') k [K(k) - E(k)] dz' \mp \frac{1}{2} \gamma(z) \\ &- \frac{(1-w_{x_d})}{\pi} \int_0^1 \frac{s'(z') k E(k)}{(z-z')} dz' + \left[\frac{w_a(z)}{V} \right]_p \end{aligned} \quad [2.2.11]$$

A mean velocity is defined as

$$\left[\frac{w_a(z)}{V} \right]_{\text{mean}} = \frac{1}{2} \left\{ \left[\frac{w_a(z)}{V} \right]_{+} + \left[\frac{w_a(z)}{V} \right]_{-} \right\} = \frac{1}{4} [C_{p+} + C_{p-}] \quad [2.2.12]$$

and substitution of Equation [2.2.11] into this equation gives

$$\begin{aligned} \left[\frac{w_a(z)}{V} \right]_{\text{mean}} &= \frac{h}{2\pi} \int_0^1 \gamma(z') k [K(k) - E(k)] dz' \\ &- \frac{(1-w_{x_d})}{\pi} \int_0^1 \frac{s'(z') k E(k)}{(z-z')} dz' + \left[\frac{w_a(z)}{V} \right]_p \end{aligned} \quad [2.2.13]$$

A velocity difference is also defined as

$$\left[\frac{w_a(z)}{V} \right]_{\text{diff}} = \left\{ \left[\frac{w_a(z)}{V} \right]_{+} - \left[\frac{w_a(z)}{V} \right]_{-} \right\} = \frac{1}{2} [C_{p+} - C_{p-}] \quad [2.2.14]$$

and substitution into Equation [2.2.11] gives

$$\left[\frac{w_a(z)}{V} \right]_{\text{diff}} = -\gamma(z) \quad [2.2.15]$$

Since $\left[\frac{w_a(z)}{V} \right]_{\text{mean}}$ and $\left[\frac{w_a(z)}{V} \right]_{\text{diff}}$ are known from the pressure distribution, substitution of Equation [2.2.15] into [2.2.13] gives a singular integral equation for the slope of the thickness distribution in terms of known quantities:

$$\int_0^1 \frac{s'(z) k E(k)}{(z - z')} dz' = \frac{-\pi}{(1-w_{x_d})} \left\{ \left[\frac{w_a(z)}{V} \right]_{\text{mean}} - \left[\frac{w_a(z)}{V} \right]_p \right\} - \frac{h}{2(1-w_{x_d})} \int_0^1 \left[\frac{w_a(z)}{V} \right]_{\text{diff}} k [K(k) - E(k)] dz' \quad [2.2.16]$$

2.3 REDUCTION OF THE THICKNESS DISTRIBUTION EQUATION FOR SOLUTION ON THE IBM 7090

To facilitate solving the expression for the slope of the thickness distribution, Equation [2.2.16] is reduced to a Fredholm equation of the second kind by using a method given by Muskhelishvili.¹⁵ The resulting equation contains singularities and is handled by a method discussed in Reference 16.

For convenience, Equation [2.2.16] is rewritten as

$$\int_0^1 g(z - z') \frac{s'(z')}{(z - z')} dz' = H(z) \quad [2.3.1]$$

where $g(z - z') = k E(k)$

$$\text{and } H(z) = -\frac{h}{2(1-w_{x_d})} \int_0^1 \left[\frac{w_a(z')}{V} \right]_{\text{diff}} k [K(k) - E(k)] dz' - \frac{\pi}{(1-w_{x_d})} \left\{ \left[\frac{w_a(z)}{V} \right]_{\text{mean}} - \left[\frac{w_a(z)}{V} \right]_p \right\}$$

Now the term $g(0) \frac{s'(z')}{(z - z')}$ is added to and subtracted from the integrand of Equation [2.3.1] giving

$$g(0) \int_0^1 \frac{s'(z')}{(z - z')} dz' + \int_0^1 [g(z - z') - g(0)] \frac{s'(z')}{(z - z')} dz' = H(z)$$

or since $g(0) = 1$

$$\int_0^1 \frac{s'(z')}{(z - z')} dz' = H(z) - \int_0^1 [g(z - z') - 1] \frac{s'(z')}{(z - z')} dz' \equiv \bar{H}(z) \quad [2.3.2]$$

This equation is now in the form of a Cauchy integral equation because the right side is free from singularities. As found in Reference 9, this equation has a unique inverse given by

$$s'(z) = \frac{1}{\pi \sqrt{z(1-z)}} \left\{ \frac{1}{\pi} \int_0^1 \frac{\sqrt{z'(1-z')}}{(z' - z)} \bar{H}(z') dz' + 2 \int_0^1 s'(z') dz' \right\}$$

Since the airfoil section must be closed, $\int_0^1 s'(z') dz' = 0$ and by using equation [2.3.2] a Fredholm equation of the second kind is obtained

$$s'(z) = \frac{1}{\sqrt{z(1-z)}} \left\{ -f(z) + \frac{1}{\pi^2} \int_0^1 G(z, z') s'(z') dz' \right\} \quad [2.3.3]$$

where

$$f(z) = -\frac{1}{\pi^2} \int_0^1 \frac{\sqrt{z'(1-z')}}{(z' - z)} H(z') dz'$$

and

$$G(z, z') = \frac{1}{\pi} \int_0^1 \frac{\sqrt{z''(1-z'')}}{(z'' - z)(z'' - z')} [1 - g(z' - z'')] dz''$$

Equation [2.3.3] cannot be handled in its present form since there exists a square root singularity at $z = 0$ and $z = 1$. To remove the singularities and, as will be seen later, to facilitate evaluation of $f(z)$ and $G(z, z')$ the following change of variable is used:

$$z = \frac{1}{2} (1 + \cos \theta)$$

Since

$$s'(z) = -\frac{2 s'[(1/2)(1 + \cos \theta)]}{\sin \theta}$$

the following equation is obtained from Equation [2.3.3]:

$$s^{**}(\theta) = f^*(\theta) + \int_0^1 G^*(\theta, \theta') s^{**}(\theta') d\theta' \quad [2.3.4]$$

where

$$f^*(\theta) = -\frac{1}{2\pi^2} \int_0^\pi \frac{H^*(\theta') \sin^2 \theta' d\theta'}{(\cos \theta' - \cos \theta)}$$

$$G^*(\theta, \theta') = \frac{1}{\pi^2} \int_0^\pi \frac{[1 - k E(k)] \sin^2 \theta'' d\theta''}{(\cos \theta'' - \cos \theta) (\cos \theta'' - \cos \theta')}$$

and

$$k^2 = \frac{1}{\frac{1}{4} h^2 (\cos \theta'' - \cos \theta')^2 + 1}$$

The symbol * is used to simplify notation; for example,

$$f(z) = f\left[\frac{1}{2}(1 + \cos \theta)\right] \equiv f^*(\theta)$$

The integrals $f^*(\theta)$ and $G^*(\theta, \theta')$ are Cauchy principal value integrals and to evaluate, part of the integrand of each is expanded in a Fourier cosine series, i.e.,

$$H^*(\theta') \sin^2 \theta' = a_0 + \sum_{n=1}^{\infty} a_n \cos n\theta'$$

where

$$a_0 = \frac{1}{\pi} \int_0^\pi H^*(\theta') \sin^2 \theta' d\theta'$$

$$a_n = \frac{2}{\pi} \int_0^\pi H^*(\theta') \sin^2 \theta' \cos n\theta' d\theta'$$

then

$$f^*(\theta) = \frac{1}{2\pi} \sum_{n=1}^{\infty} a_n \frac{\sin n\theta}{\sin \theta} \quad [2.3.5]$$

and

$$\frac{[1 - k E(k)] \sin^2 \theta''}{\cos \theta'' - \cos \theta'} = b_0(\theta') + \sum_{n=1}^{\infty} b_n(\theta') \cos n\theta''$$

where

$$b_0(\theta') = \frac{1}{\pi} \int_0^\pi \frac{[1 - k E(k)] \sin^2 \theta''}{\cos \theta'' - \cos \theta'} d\theta''$$

$$b_n(\theta') = \frac{2}{\pi} \int_0^\pi \frac{[1 - k E(k)] \sin^2 \theta''}{\cos \theta'' - \cos \theta'} \cos n\theta'' d\theta''$$

then

$$G^*(\theta, \theta') = \frac{1}{\pi} \sum_{n=1}^{\infty} b_n(\theta') \frac{\sin n\theta}{\sin \theta} \quad [2.3.6]$$

The kernel $G^*(\theta, \theta')$ is now of the degenerate type and Equation [2.3.4] can be solved by the method applicable to this type of kernel. Details of the method are given in Reference 17. If this procedure is followed, the following equation for the slope of the thickness distribution is obtained.

$$s^{*\prime}(\theta) = f^*(\theta) + \frac{1}{\pi} A_1 + \frac{1}{\pi} A_2 \frac{\sin 2\theta}{\sin \theta} + \dots + \frac{1}{\pi} A_n \frac{\sin n\theta}{\sin \theta} \quad [2.3.7]$$

where

$f^*(\theta)$ is given by Equation [2.3.5].

And A_n is given by the following set of simultaneous equations:

$$\begin{aligned} A_1 (1 - c_{11}) - A_2 c_{12} - \dots - A_n c_{1n} &= d_1 \\ - A_1 c_{21} + A_2 (1 - c_{22}) - \dots - A_n c_{2n} &= d_2 \\ \cdot &\cdot \\ \cdot &\cdot \\ - A_1 c_{n1} - A_2 c_{n2} - \dots + A_n (1 - c_{nn}) &= d_n \end{aligned} \quad [2.3.8]$$

where

$$c_{ij} = \frac{1}{\pi} \int_0^\pi b_i(\theta') \frac{\sin j\theta'}{\sin \theta'} d\theta' \quad (i, j = 1, 2, 3 \dots n)$$

$$d_i = \int_0^\pi b_i(\theta') f^*(\theta') d\theta' \quad (i = 1, 2, 3 \dots n)$$

Expression [2.3.7] is the expression solved on the computer with a finite sum replacing the infinite sum in Equation [2.3.5]. The c_{ij} 's and d_i 's may be evaluated as given, so the A_i 's follow immediately from [2.3.8]. However, to evaluate the a_n 's in Equation [2.3.5], the singularity must be removed in $H^*(\theta')$. This is accomplished by using the change of variables $t^3 = \theta'' - \theta'$ where θ'' is the variable of integration for $H^*(\theta')$. The thickness distribution now follows immediately by integrating Equation [2.3.7] from 0 to θ .

2.4 DERIVATION OF THE EQUATION FOR THE CAMBER DISTRIBUTION AND ANGLE OF ATTACK

To calculate the camber distribution and the angle of attack, it is only necessary to integrate Equation [2.2.4] after substituting for the known quantities. Equation [2.2.4] is rewritten and substitution for the source strength q and vortex strength γ is made from Equations [2.2.5] and [2.2.15], respectively.

$$\begin{aligned}
 c'(z) + \tan \alpha = & - \frac{h^2}{\pi(1 - w_{x_d})} \int_0^1 \left[\frac{w_a(z')}{V} \right] \text{diff} (z - z') k [K(k) - E(k)] dz' \\
 & + \frac{1}{2\pi(1 - w_{x_d})} \int_0^1 \left[\frac{w_a(z')}{V} \right] \text{diff} \frac{k E(k)}{(z - z')} dz' \quad [2.4.1] \\
 & + \frac{h}{\pi} \int_0^1 s'(z') k [K(k) - E(k)] dz' - \frac{1}{(1 - w_{x_d})} \left[\frac{w_r}{V} (x_d, z) \right]_p
 \end{aligned}$$

If Equation [2.4.1] is integrated with respect to z from 0 to 1, the tangent of the angle of attack is obtained since $\int_0^1 c'(z) dz = 0$. Equation [2.4.1] is the desired expression for the slope of the camber distribution.

2.5 SIMPLIFICATION OF THE CAMBER DISTRIBUTION EQUATION FOR SOLUTION ON THE IBM 7090

To facilitate solving the expression for the slope of the camber distribution, Equation [2.4.1], a change of variable is made. The resulting equation contains various singularities which are handled by a method suggested in Reference 15..

With $z = \frac{1}{2} (1 + \cos \theta)$ Equation [2.4.1] has the form

$$\begin{aligned}
 \frac{-2c^*(\theta)}{\sin \theta} + \tan \alpha = & - \frac{h^2}{4\pi(1 - w_{x_d})} \int_0^1 \left[\frac{w_a^*(\theta')}{V} \right] \text{diff} (\cos \theta - \cos \theta') \\
 & \cdot k [K(k) - E(k)] \sin \theta' d\theta' \\
 & + \frac{1}{2\pi(1 - w_{x_d})} \int_0^1 \left[\frac{w_a^*(\theta')}{V} \right] \text{diff} \frac{k E(k) \sin \theta' d\theta'}{\cos \theta - \cos \theta'}
 \end{aligned}$$

$$-\frac{h}{\pi} \int_0^\pi s^{*'}(\theta') k[K(k) - E(k)] d\theta' \quad [2.5.1]$$

$$\frac{1}{(1 - w_{x_d})} \left[\frac{w_r^*(x_d, \theta)}{V} \right]_p$$

where the notation * and the form of k are given in Equation [2.3.4].

Because of the presence of the term $\cos \theta - \cos \theta'$ the first integral in Equation [2.5.1] has no singularities and may, therefore, be evaluated immediately by use of the computer.

The second integral in Equation [2.5.1] is a Cauchy principal value integral and is first simplified by a technique given in Reference 15; the term

$$\left[\frac{w_a^*(\theta)}{V} \right]_{\text{diff}} \frac{\sin \theta}{\cos \theta - \cos \theta'}$$

is added to and subtracted from the integrand:

$$\frac{1}{2\pi(1 - w_{x_d})} \left\{ \int_0^\pi \frac{\left[\frac{w_a^*(\theta)}{V} \right]_{\text{diff}} kE(k) \sin \theta' - \left[\frac{w_a^*(\theta)}{V} \right]_{\text{diff}} \sin \theta}{\cos \theta - \cos \theta'} d\theta' \right. \\ \left. + \left[\frac{w_a^*(\theta)}{V} \right]_{\text{diff}} \sin \theta \int_0^\pi \frac{d\theta'}{\cos \theta - \cos \theta'} \right\} \quad [2.5.2]$$

The second integral in this expression is zero so only the first integral needs to be considered.

The behavior of the integrand is investigated as $\theta' \rightarrow \theta$ by L'Hospital's rule and the following limit is obtained:

$$\frac{d}{d\theta} \left\{ \frac{\left[\frac{w_a^*(\theta)}{V} \right]_{\text{diff}} \sin \theta}{\sin \theta} \right\} \quad [2.5.3]$$

Now with this knowledge of the behavior of the integrand of expression

[2.5.2], the second integral in Equation [2.5.1] is evaluated by integrating [2.5.2] on the computer and using [2.5.3] evaluated numerically for the value of the integrand at $\theta' = \theta$.

The third integral in Equation [2.5.1] is simplified as follows. First this equation is integrated by parts with

$$u = k [K(k) - E(k)]$$

$$dv = s^*(\theta') d\theta'$$

Then $v = s^*(\theta')$ and after differentiation and considerable simplification (see page 107 of Reference 18),

$$du = -k^3 \frac{h^2}{4} (\cos \theta - \cos \theta') \sin \theta' \left\{ K(k) - E(k) + E(k) \frac{k^2}{1-k^2} \right\} d\theta'$$

Now by using the fact that $s^*(\theta') \big|_{\theta'=0} = s^*(\theta') \big|_{\theta'=\pi} = 0$

the following equation is obtained

$$\int_0^\pi s^*(\theta') k [K(k) - E(k)] d\theta' = \frac{h^2}{4} \int_0^\pi s^*(\theta') k^3 (\cos \theta - \cos \theta')$$

$$\sin \theta' [K(k) - E(k)] d\theta'$$

$$+ \int_0^\pi \frac{s^*(\theta') k^3 \sin \theta' E(k)}{\cos \theta - \cos \theta'} d\theta' \quad [2.5.4]$$

The first integral on the right-hand side of Equation [2.5.4] can be evaluated on the computer immediately; the second, however, must be simplified as in the second integral of Equation [2.5.1]. Thus the term

$$\frac{s^*(\theta) \sin \theta}{\cos \theta - \cos \theta'}$$

is added to and subtracted from the integrand obtaining

$$\int_0^\pi \frac{s^*(\theta') k^3 \sin \theta' E(k) - s^*(\theta) \sin \theta}{\cos \theta - \cos \theta'} d\theta'$$

since, as before, $\int_0^\pi \frac{d\theta'}{\cos \theta - \cos \theta'} d\theta' = 0$

As in the integral [2.5.2] the behavior of the integrand as $\theta' \rightarrow \theta$ is investigated and $s^{*\prime}(\theta) + s^*(\theta) \cot \theta$ is obtained as the value of the limit. Thus the complete expression [2.5.4] may now be evaluated numerically.

All three integrals of Equation [2.5.1] are now in a form suitable for numerical evaluation, so $\tan \alpha$ may be found by integrating [2.5.1] from 0 to π and recalling that $\int_0^\pi c^{*\prime}(\theta) d\theta = 0$. The slope of the camber distribution then follows immediately, and integration from 0 to θ gives the camber distribution.

3. COMPUTER PROGRAM

The calculations based on the theory presented in Section 2 have been programmed for the IBM 7090 high-speed computer. Input consists of a pressure distribution on the inner and outer surfaces of the annular airfoil, the chord-diameter ratio of the annular airfoil, and a wake fraction. It is also possible to make calculations with a propeller or any axisymmetric body located in the annular airfoil; however, axial and radial components of the induced velocity must be included.

The output consists of the thickness distribution, the camber distribution, and the angle of attack of the annular airfoil. It takes approximately 8 minutes on the IBM 7090 high-speed computer to obtain these results.

The input-output format and the FORTRAN listing of the computer program are discussed in the following sections.

3.1 INPUT FORMAT

The first input card contains an integer which represents the number of cases to be run. This number is entered in Columns 1-4 under an I4 format. The only limit to the number of cases to be run is a consideration of computer time.

The second card is a problem identification card. A one (1) must appear in Column 1 and any alphanumeric characters may appear in Columns 2 through 72. Even if no identification is desired, the card must be included with a one (1) in Column 1.

The third card contains 10 parameters in the format 2F10.6,8I4.

They are as follows:

1. The chord-diameter ratio.
2. The wake fraction.
3. The maximum number of Fourier coefficients to be allowed in the evaluation of the kernel function $G^*(\theta, \theta')$, Equation [2.3.4]. The user may specify any number not exceeding 40. The program calculates Fourier coefficients until either a convergence criteria of 10^{-6} has been met or the number of coefficients equals the number specified in this parameter. The program normally uses between 20 and 25 terms so the user may specify 40 since only as many terms are calculated as are needed.
4. The number of ordinates used in the integral evaluations. The integrations are performed using Simpson's rule so this number must be odd. The maximum allowed by the dimension of the program is 101. Normally, a value of 51 is more than adequate and if time is a factor, 25 would probably suffice.
5. The number of input pressure data points. This number may not exceed 37. If possible, to facilitate interpolation, the user should specify closer spaced data in regions where the pressure curves have steep slopes.
6. The number of points to be used for the harmonic analysis of the Cauchy principle value integral $f^*(\theta)$, Equation [2.3.4]. The maximum allowed by the dimension of the program is 200 and this is also the suggested value.
7. The number of harmonics used in the harmonic analysis described in parameter 6. This number must be strictly less than one-half the number specified in 6. Normally, a value of 45 is more than adequate.
8. The maximum number of ordinates used in integrating $s^*(\theta)$ and $c^*(\theta)$ to obtain the thickness and camber distributions. The number must be greater than or equal to 51 and must be odd. A value of 61 has proven to be satisfactory. There is no upper bound on this parameter.
9. The parameter signifying the presence of a propeller or any axisymmetric body in the annular airfoil. The parameter must have the values:
 - 1 if a propeller or body is present,
 - 2 if a propeller or body is not present.

10. The parameter controlling the punching of one-half the thickness and camber distributions on IBM cards by the machine. This parameter must have the values:

- 1 if no data is to be punched,
- 2 if data is to be punched.

The format of the punched data is discussed in the following paragraphs. A discussion of the effects of varying parameters 3 and 7 may be found in Reference 16.

The fourth card is the first card containing pressure data. The value of parameter 5 on the third card determines the number of these cards. Each card contains the abscissa of the pressure distribution and four pressure data terms in a format of 5F14.8. The fourth card contains the abscissa $x = 0.0$ and the appropriate pressure data at the leading edge of the annular airfoil. The fifth card contains the next abscissa and its corresponding pressure data. The abscissa value may be arbitrarily spaced in the interval $[0,1]$ as long as the first one is 0.0 and they are strictly increasing. The second term on each card is the pressure on the outer surface of the annular airfoil. The third is the pressure on the inner surface of the annular airfoil.

If no propeller is present, i.e., a two (2) has been given as parameter 9 on the third card, the fourth and fifth terms on each card are omitted; however, if a propeller is present, i.e., a one (1) has been given as parameter 9 on the third card, then the fourth and fifth terms must be specified. The fourth term is the axial induced velocity and the fifth is the radial induced velocity. An example showing the input data for a ducted propeller is shown in Appendix A.

All cards after the first card must be repeated for each case even though some of the data may be the same.

3.2 OUTPUT FORMAT

The first page of output contains the input data. The second page contains the angle of attack of the annular airfoil in degrees and a table consisting of one-half the thickness and of the camber distribution from the leading edge of the foil, 0.0, to the trailing edge, 1.0, in increments of 0.05. The output obtained from the input data of the ducted propeller given in Section 3.1 is also shown in Appendix A.

If parameter 10 on the third card is a two (2), data will be punched on IBM cards by the program. The first three cards contain a total of 21 numbers, ranging from 0.0 to 1.0 in increments of 0.05, which are the abscissa values of one-half the thickness distribution. The next three cards contain the values of one-half the thickness distribution at the abscissa values given the first three cards. The next three cards, which are the abscissa values of the camber distribution, are the same as the first three cards. The last three cards contain the values of the camber distribution at the abscissa values given on the preceding three cards. All this output is punched in a format of 9F8.5.

3.3 FORTRAN LISTING

The FORTRAN listing of the computer program is given in Appendix B. In addition to the subroutines furnished automatically by the Bell Monitor System on the IBM 7090, the binary coding of the following subroutines available from SHARE must be added to the FORTRAN listing: BE-ELIP, AMGMHA, AMMATI, LACBRT, AQALLI, and VG-AS + C. The program takes about eight minutes to run under the Bell Monitor System.

4. RESULTS OF CALCULATIONS

4.1 ANNULAR AIRFOIL

A number of calculations were made to compare calculated duct shapes with actual shapes. Two annular airfoils from Reference 7 were considered, Duct II and the BTZ duct, and the pressure distributions for these ducts are given in Tables 1 and 2, respectively. There are two pressure distributions given for each duct; one is the experimental distribution and the other, the linearized theoretical distribution as calculated by the method of Reference 7. Tables 3 and 4 give the tabulated data for the thickness, camber, and angle of attack as calculated and the actual shape for each of the ducts. Figures 3 and 4 show a comparison of the section shapes for Duct II and the BTZ duct, respectively. The ordinates have been expanded to accentuate the differences.

A comparison of the results show that for either the linearized or experimental pressure distribution the calculated thickness is a few

TABLE 1

Pressure Distribution for Duct II

1 - z	Experimental		Linearized Theory	
	C _p +	C _p -	C _p +	C _p -
Leading Edge	1.000	1.000	1.000	-1.000
0.0019	0.400	0.400	0.563	-0.676
0.0076	0.151	0.151	0.117	-0.247
0.0170	0.030	0.030	-0.077	-0.062
0.0302	-0.081	0.058	-0.195	0.051
0.0469	-0.160	0.088	-0.272	0.121
0.0670	-0.199	0.101	-0.319	0.162
0.0904	-0.217	0.110	-0.343	0.179
0.1170	-0.234	0.112	-0.350	0.178
0.1464	-0.250	0.109	-0.345	0.166
0.1786	-0.275	0.101	-0.334	0.148
0.2132	-0.305	0.100	-0.323	0.129
0.2500	-0.321	0.107	-0.313	0.114
0.2887	-0.331	0.119	-0.309	0.107
0.3290	-0.338	0.131	-0.311	0.108
0.3706	-0.339	0.137	-0.320	0.118
0.4132	-0.345	0.135	-0.334	0.134
0.4564	-0.351	0.145	-0.351	0.155
0.5000	-0.355	0.146	-0.367	0.177
0.5436	-0.350	0.152	-0.381	0.198
0.5868	-0.360	0.159	-0.388	0.215
0.6294	-0.369	0.165	-0.387	0.227
0.6710	-0.359	0.161	-0.376	0.233
0.7113	-0.345	0.175	-0.354	0.233
0.7500	-0.330	0.171	-0.321	0.230
0.7868	-0.308	0.179	-0.279	0.224
0.8214	-0.270	0.191	-0.229	0.218
0.8536	-0.235	0.192	-0.174	0.213
0.8830	-0.169	0.192	-0.116	0.209
0.9096	-0.100	0.192	-0.058	0.209
0.9330	-0.032	0.190	-0.002	0.211
0.9532	0.032	0.190	0.051	0.215
0.9698	0.085	0.192	0.098	0.220
0.9830	0.140	0.195	0.139	0.224
0.9924	0.160	0.197	0.170	0.224
0.9981	0.184	0.199	0.192	0.218
1.0000	0.200	0.200	1.000	1.000

TABLE 2

Pressure Distribution for the BTZ Duct

l - z	Experimental		Linearized Theory	
	C _{p+}	C _{p-}	C _{p+}	C _{p-}
Leading Edge	1.000	1.000	1.000	-1.000
0.0019	0.700	0.700	-0.115	-0.178
0.0076	0.120	0.250	-0.120	-0.166
0.0170	-0.037	0.070	-0.118	-0.167
0.0302	-0.085	-0.025	-0.112	-0.170
0.0469	-0.100	-0.100	-0.101	-0.173
0.0670	-0.110	-0.145	-0.103	-0.178
0.0904	-0.100	-0.170	-0.098	-0.184
0.1170	-0.100	-0.182	-0.095	-0.190
0.1464	-0.100	-0.200	-0.093	-0.198
0.1786	-0.100	-0.215	-0.093	-0.207
0.2132	-0.102	-0.226	-0.094	-0.217
0.2500	-0.110	-0.241	-0.097	-0.228
0.2887	-0.110	-0.255	-0.101	-0.239
0.3290	-0.110	-0.245	-0.105	-0.249
0.3706	-0.113	-0.280	-0.110	-0.258
0.4132	-0.118	-0.280	-0.114	-0.256
0.4564	-0.119	-0.280	-0.117	-0.269
0.5000	-0.112	-0.270	-0.119	-0.270
0.5436	-0.110	-0.260	-0.119	-0.266
0.5868	-0.115	-0.250	-0.116	-0.258
0.6294	-0.106	-0.221	-0.110	-0.244
0.6710	-0.090	-0.192	-0.100	-0.225
0.7113	-0.060	-0.145	-0.087	-0.200
0.7500	-0.037	-0.110	-0.070	-0.169
0.7868	-0.100	-0.061	-0.049	-0.135
0.8214	0.011	-0.030	-0.025	-0.097
0.8536	0.039	-0.005	0.002	-0.057
0.8830	0.058	0.020	0.029	-0.017
0.9096	0.087	0.050	0.058	0.023
0.9330	0.062	0.030	0.086	0.061
0.9532	0.085	0.065	0.113	0.096
0.9698	0.120	0.100	0.137	0.126
0.9830	0.142	0.142	0.159	0.152
0.9924	0.164	0.164	0.176	0.173
0.9981	0.180	0.180	0.193	0.191
1.0000	1.000	1.000	1.000	1.000

TABLE 3

Section Shape for Duct II

l - z	Design		Calculated from Experimental C_p		Calculated from Theoretical C_p	
	s(z)	c(z)	s(z)	c(z)	s(z)	c(z)
0	0.00000	0.00000	0.0000	0.0000	0.0000	0.0000
0.05	0.02095	0.01085	0.0185	0.0075	0.0210	0.0125
0.10	0.02920	0.01793	0.0271	0.0140	0.0296	0.0204
0.15	0.03528	0.02348	0.0338	0.0193	0.0359	0.0254
0.20	0.04002	0.02797	0.0393	0.0240	0.0408	0.0289
0.25	0.04364	0.03162	0.0435	0.0281	0.0446	0.0315
0.30	0.04637	0.03454	0.0466	0.0314	0.0474	0.0339
0.35	0.04832	0.03681	0.0489	0.0341	0.0494	0.0362
0.40	0.04952	0.03846	0.0504	0.0360	0.0507	0.0383
0.45	0.05000	0.03952	0.0513	0.0373	0.0512	0.0400
0.50	0.04962	0.04000	0.0514	0.0379	0.0510	0.0413
0.55	0.04846	0.03988	0.0507	0.0379	0.0500	0.0418
0.60	0.04653	0.03914	0.0494	0.0374	0.0482	0.0413
0.65	0.04383	0.03774	0.0474	0.0359	0.0456	0.0396
0.70	0.04035	0.03557	0.0442	0.0336	0.0421	0.0367
0.75	0.03612	0.03248	0.0402	0.0303	0.0377	0.0325
0.80	0.03110	0.02811	0.0351	0.0260	0.0324	0.0272
0.85	0.02532	0.02170	0.0289	0.0207	0.0262	0.0209
0.90	0.01877	0.01434	0.0215	0.0141	0.0192	0.0140
0.95	0.01143	0.00685	0.0130	0.0069	0.0115	0.0069
1.00	0.00000	0.00000	0.0000	0.0000	0.0000	0.0000
$\alpha = 0$		$\alpha = 0.0048$ degrees			$\alpha = -0.0720$ degrees	

TABLE 4

Section Shape for the BTZ Duct

1 - z	Design		Calculated from Experimental C_p		Calculated from Theoretical C_p	
	s(z)	c(z)	s(z)	c(z)	s(z)	c(z)
0.00	0.00000	0.00000	0.0000	0.0000	0.0000	0.0000
0.05	0.01257	0.00000	0.0114	0.0001	0.0129	0.0001
0.10	0.01752	0.00000	0.0172	-0.0003	0.0178	0.0001
0.15	0.02119	0.00000	0.0213	-0.0006	0.0214	0.0002
0.20	0.02401	0.00000	0.0245	-0.0007	0.0242	0.0002
0.25	0.02618	0.00000	0.0270	-0.0007	0.0265	0.0002
0.30	0.02782	0.00000	0.0288	-0.0007	0.0284	0.0003
0.35	0.02899	0.00000	0.0301	-0.0008	0.0298	0.0003
0.40	0.02971	0.00000	0.0309	-0.0008	0.0307	0.0004
0.45	0.03000	0.00000	0.0310	-0.0007	0.0311	0.0004
0.50	0.02985	0.00000	0.0304	-0.0005	0.0310	0.0003
0.55	0.02925	0.00000	0.0292	-0.0003	0.0302	0.0003
0.60	0.02815	0.00000	0.0273	0.0000	0.0287	0.0003
0.65	0.02611	0.00000	0.0246	0.0003	0.0266	0.0003
0.70	0.02316	0.00000	0.0211	0.0006	0.0237	0.0002
0.75	0.01953	0.00000	0.0174	0.0009	0.0202	0.0002
0.80	0.01543	0.00000	0.0134	0.0011	0.0161	0.0002
0.85	0.01107	0.00000	0.0090	0.0005	0.0116	0.0001
0.90	0.00665	0.00000	0.0052	0.0002	0.0070	0.0001
0.95	0.00262	0.00000	0.0023	0.0000	0.0027	0.0000
1.00	0.00000	0.00000	0.0000	0.0001	0.0000	0.0000
$\alpha = 0$		$\alpha = 0.0973$ degrees			$\alpha = -0.0097$ degrees	

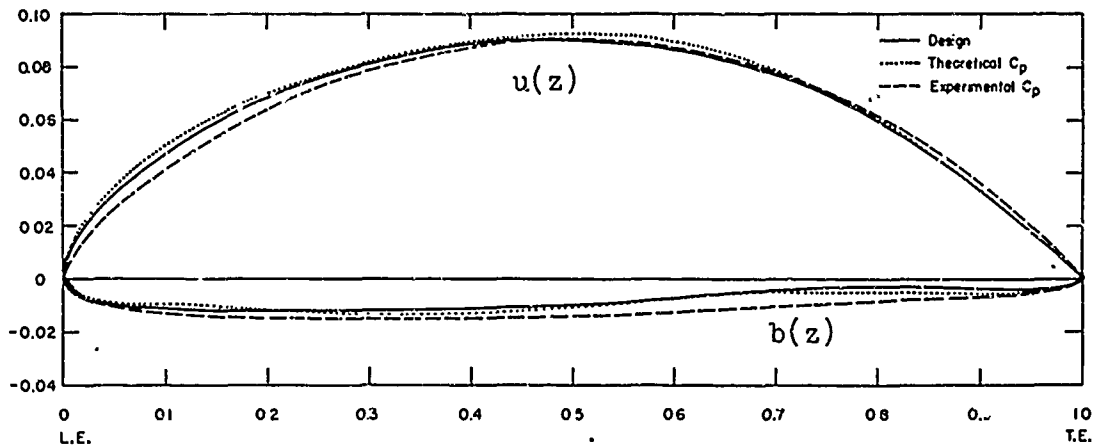


Figure 3 - Comparison of the Computed Shapes from Theoretical and Experimental Pressure Distributions with Duct II

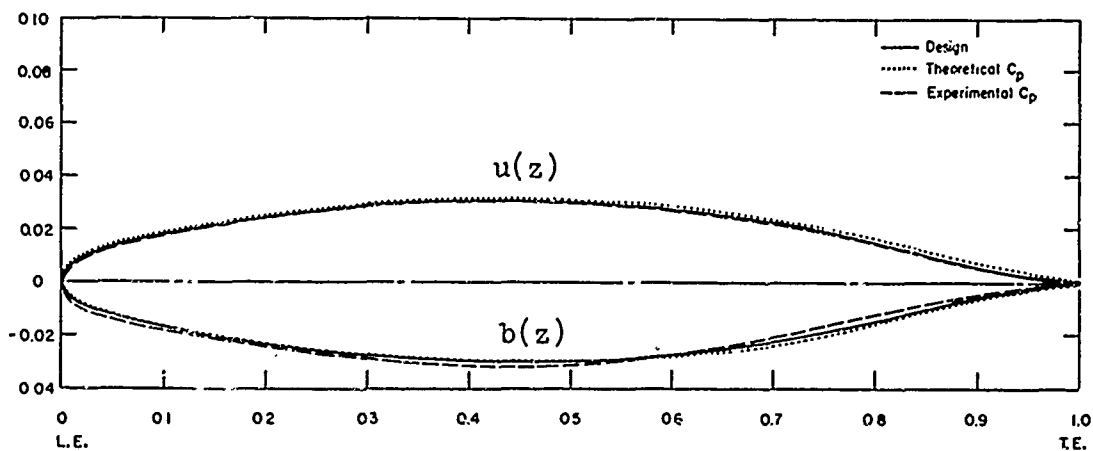


Figure 4 - Comparison of the Computed Shapes from Theoretical and Experimental Pressure Distributions with the BTZ Duct

percent larger than the design thickness. It is not possible to draw any general conclusions about the accuracy of predicting the thickness except it is reasonable and probably within the accuracy of either the linearized or experimental pressure distributions. This is all that can be expected.

The camber appears to differ by a somewhat greater magnitude than the thickness for Duct II. In fact, for the linearized pressure distribution the calculated camber is somewhat greater than the actual camber and for the experimental pressure distribution, the camber is somewhat less. The section angle of attack is less than the design angle of zero degrees for the linearized pressure distribution and slightly greater for the experimental pressure distribution. Both the camber and angle of attack show the effect of accuracy of determining the pressure distribution near the leading edge. For the linearized pressure distribution, the pressure of the duct is calculated to be infinite at the leading edge, whereas for the experimental pressure distribution the closest point measured to the leading edge was 0.025 of the chord. In either case the true pressure distribution near the leading edge is not known.

For the BTZ duct the camber and angle of attack calculated from the linearized pressure distribution are certainly within the accuracy that can be expected. The angle of attack is within 0.01 of a degree and the camber is only 1 percent of the thickness and should be considered negligible. Surprisingly, the calculated camber and angle of attack do not appear to be as accurate from the experimental pressure distribution as from the linearized. There are only two experimental points (one inside and the other outside) within 10 percent of the leading edge and both the calculated results of Reference 7 and the present calculations indicate that the point on the inside of the duct is incorrect.

The foregoing results of the ducts by themselves are presented to give an indication of the accuracy of the procedure. In reality, it makes little sense to use a linearized theory to calculate a shape from a linearized pressure distribution. The usefulness of the program presented in this report is to design a duct for a given pressure distribution when the duct is in the presence of a propeller or some other axisymmetric body. This allows a duct shape to be chosen which will not cavitate or separate.

TABLE 5

Pressure Distribution for a NACA 66-010 Thickness Distribution
and Propeller Induced Velocities

$1 - z$	c_{p+}	c_{p-}	$(\frac{w}{V})p/c_T$	$(\frac{w}{V})p/c_T$
Leading Edge	1.000	1.000	-0.0382	-0.0157
0.0050	0.104	0.104	-0.0384	-0.0158
0.0075	0.028	0.028	-0.0386	-0.0160
0.0125	-0.023	-0.023	-0.0388	-0.0168
0.0250	-0.078	-0.078	-0.0389	-0.0176
0.0500	-0.125	-0.125	-0.0410	-0.0180
0.0750	-0.154	-0.154	-0.0426	-0.0199
0.1000	-0.174	-0.174	-0.0449	-0.0200
0.1500	-0.198	-0.198	-0.0480	-0.0230
0.2000	-0.215	-0.215	-0.0525	-0.0280
0.2500	-0.226	-0.226	-0.0564	-0.0697
0.3000	-0.236	-0.236	-0.0555	-0.1470
0.3500	-0.243	-0.243	-0.0400	-0.1660
0.4000	-0.249	-0.249	-0.0330	-0.1760
0.4500	-0.255	-0.255	-0.0210	-0.1860
0.5000	-0.261	-0.261	-0.0000	-0.1950
0.5500	-0.265	-0.265	0.0210	-0.1860
0.6000	-0.270	-0.270	0.0330	-0.1760
0.6500	-0.250	-0.250	0.0400	-0.1660
0.7000	-0.190	-0.190	0.0555	-0.1470
0.7500	-0.121	-0.121	0.0564	-0.0697
0.8000	-0.052	-0.052	0.0525	-0.0280
0.8500	0.021	0.021	0.0480	-0.0230
0.9000	0.096	0.096	0.0449	-0.0200
0.9500	0.179	0.179	0.0410	-0.0180
1.0000	0.271	0.271	0.0382	-0.0157

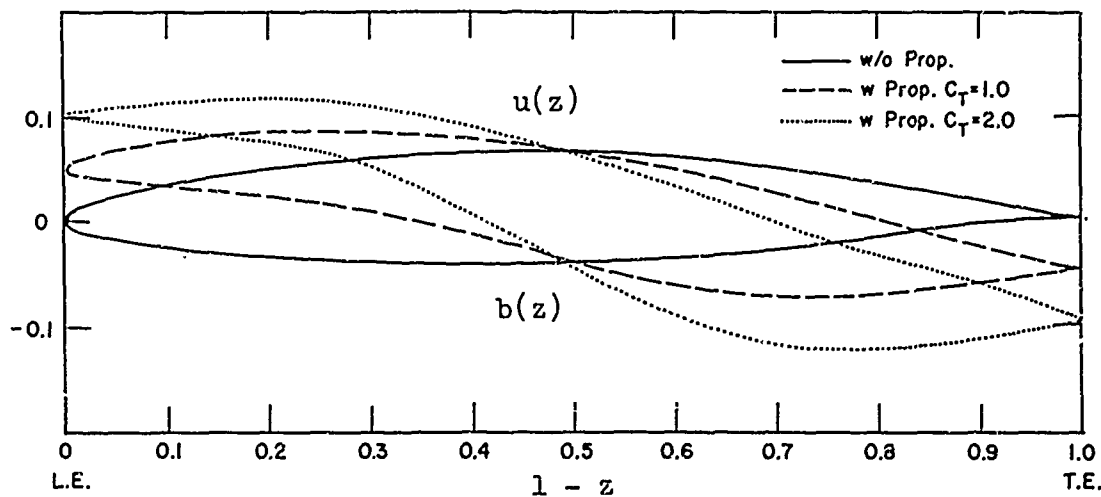


Figure 5 - Shape of Duct with Pressure Distribution
Corresponding to a NACA 66-010 Thickness
Form with and without a Propeller

4.2 DUCTED PROPELLER

The shape of the duct of a ducted propeller can be calculated by the program by inputting the steady axial and radial velocities induced by the propeller on the duct. These velocities for use in this manner are discussed in References 9 and 14 and have been tabulated in Reference 19.

Calculations for duct shape were made for a pressure distribution corresponding to the distribution on a NACA 66-010 basic thickness form,²⁰ Table 5. Two values of the propeller thrust loading coefficient C_T were assumed, i.e., $C_T = 1.0$ and $C_T = 2.0$.

$$C_T = \frac{T}{\frac{\rho}{2} \pi R^2 V^2}$$

where

- R = propeller radius,
- T = propeller thrust,
- V = velocity, and
- ρ = mass density.

Propeller-induced velocities which were input into the program are shown in Table 5. For these calculations the tip clearance was assumed to be one percent of the propeller radius and the propeller location was assumed to be at the midchord of the duct. Also, for convenience, the thrust loading coefficient was based on the average velocity at the propeller and not the free-stream velocity as would be the usual case.

Results of these calculations are tabulated in Table 6 and plotted in Figure 5. Also plotted in this figure is the shape of the duct without the propeller. Not only is the camber and angle of attack of each duct changed considerably but the thickness is also changed. The biggest effect is on the angle of attack which goes from essentially zero to 5.35 degrees for a $C_T = 1.0$ and from zero to 10.48 degrees for a $C_T = 2.0$. Also the cambers of all the ducts are S-shaped.

With the propeller in the duct, the duct sections are thinner near the leading edge and thicker toward the trailing edge than for the duct alone. This effect increases with propeller loading. Also, the maximum duct

TABLE 6

Shape of Duct with and without a Propeller

1 - z	Without Propeller		With Propeller $C_T = 1.0$		With Propeller $C_T = 2.0$	
	s(z)	c(z)	s(z)	c(z)	s(z)	c(z)
0.00	0.0000	0.0000	0.0000	0.0000	0.0000	0.0000
0.05	0.0192	0.0020	0.0133	0.0054	0.0073	0.0088
0.10	0.0284	0.0039	0.0203	0.0106	0.0122	0.0174
0.15	0.0352	0.0056	0.0257	0.0157	0.0163	0.0259
0.20	0.0404	0.0070	0.0302	0.0207	0.0201	0.0343
0.25	0.0445	0.0083	0.0343	0.0243	0.0240	0.0402
0.30	0.0477	0.0093	0.0382	0.0244	0.0287	0.0394
0.35	0.0500	0.0101	0.0422	0.0222	0.0344	0.0343
0.40	0.0514	0.0106	0.0457	0.0191	0.0401	0.0276
0.45	0.0521	0.0109	0.0490	0.0154	0.0460	0.0199
0.50	0.0518	0.0109	0.0519	0.0109	0.0520	0.0108
0.55	0.0507	0.0106	0.0539	0.0061	0.0571	0.0016
0.60	0.0485	0.0101	0.0542	0.0015	0.0600	-0.0070
0.65	0.0445	0.0092	0.0525	-0.0029	0.0604	-0.0151
0.70	0.0389	0.0081	0.0485	-0.0072	0.0582	-0.0225
0.75	0.0322	0.0068	0.0426	-0.0089	0.0529	-0.0247
0.80	0.0250	0.0054	0.0353	-0.0080	0.0455	-0.0214
0.85	0.0177	0.0039	0.0272	-0.0062	0.0368	-0.0164
0.90	0.0106	0.0025	0.0188	-0.0044	0.0270	-0.0112
0.95	0.0044	0.0012	0.0104	-0.0022	0.0164	-0.0056
1.00	0.0000	0.0000	0.0000	0.0000	0.0000	0.0000
$\alpha = 0.13$ degrees			$\alpha = 5.35$ degrees		$\alpha = 10.48$ degrees	

thickness increases with propeller loading. Both effects are a consequence of assuming the propeller location to be at midchord.

If the propeller induces a velocity in the same direction as the free-stream velocity everywhere on the duct, the thickness is decreased everywhere. This would be the case when the propeller is located at or aft of the duct trailing edge. Conversely, if the propeller induces a velocity in the opposite direction to the free-stream velocity everywhere on the duct, the thickness is increased everywhere. This would be the case when the propeller is located at or forward of the duct leading edge.

The calculations for a duct shape with a propeller were based on the assumption that the inside and outside pressures were the same. Additional calculations were made for a thrust loading coefficient of two, $C_T = 2.0$, and the pressure distribution given in Table 7. The pressure distributions correspond to a NACA 66-010 thickness distribution with a NACA $a = 0.8$ mean line of 4-percent camber. In one case the camber is negative (toward the inside of the duct) and in the other, positive which has the effect of shifting the negative pressure from the inside to the outside of the duct.

Results of these calculations are shown in Table 8. The pressure distribution with a negative pressure on the inside of the duct results in a decrease in duct thickness whereas that with a positive pressure on the inside of the duct results in an increase in duct thickness. This is a consequence of the change in sign of $\left[\frac{w}{v} a(z)\right]_{\text{diff}}$ in Equation [2.2.16]. In fact, the effect of this term is so large that the calculations show that to achieve the pressure distribution chosen for the calculation with a negative pressure inside the duct, the thickness would have to be negative near the leading edge. Since ducts of this type present construction problems, it can only be concluded that such a pressure distribution cannot be achieved for the propeller loading and location assumed.

As in Table 6, both cambers given in Table 8 are somewhat S-shaped. The difference being that for the negative pressure inside the duct, the camber is generally negative and for the positive pressure inside the duct, the camber is generally positive. It should be noted that the angle of attack varies less than a degree between the two ducts.

TABLE 7

Pressure Distribution of a NACA 65-010 Thickness Form
with a NACA $\alpha = 0.8$ Mean Line of 4 Percent
Positive and Negative Camber

1 - z	Positive Camber		Negative Camber	
	c_{p+}	c_{p-}	c_{p+}	c_{p-}
Leading Edge	1.000	1.000	1.000	1.000
0.0050	-0.234	0.387	0.387	-0.234
0.0075	-0.323	0.324	0.324	-0.323
0.0125	-0.378	0.283	0.283	-0.378
0.0250	-0.445	0.236	0.236	-0.445
0.0500	-0.501	0.195	0.195	-0.501
0.0750	-0.533	0.172	0.172	-0.533
0.1000	-0.558	0.154	0.154	-0.558
0.1500	-0.585	0.133	0.133	-0.585
0.2000	-0.603	0.120	0.120	-0.603
0.2500	-0.615	0.111	0.111	-0.615
0.3000	-0.628	0.101	0.101	-0.628
0.3500	-0.636	0.096	0.096	-0.636
0.4000	-0.644	0.090	0.090	-0.644
0.4500	-0.649	0.086	0.086	-0.649
0.5000	-0.656	0.080	0.080	-0.656
0.5500	-0.662	0.076	0.076	-0.662
0.6000	-0.667	0.073	0.073	-0.667
0.6500	-0.644	0.090	0.090	-0.644
0.7000	-0.575	0.141	0.141	-0.575
0.7500	-0.496	0.199	0.199	-0.496
0.8000	-0.416	0.257	0.257	-0.416
0.8500	-0.237	0.250	0.250	-0.237
0.9000	-0.067	0.245	0.245	-0.067
0.9500	-0.103	0.252	0.252	-0.103
1.0000	1.000	1.000	1.000	1.000

TABLE 8

Shape of Duct with a Propeller for Positive
and Negative Inside Pressures

1 - z	From Pressure Distribution with Positive Camber		From Pressure Distribution with Negative Camber	
	s(z)	c(z)	s(z)	c(z)
0.00	0.0000	0.0000	0.0000	0.0000
0.05	0.0292	0.0212	-0.0086	-0.0030
0.10	0.0428	0.0381	-0.0103	-0.0022
0.15	0.0532	0.0531	-0.0110	0.0001
0.20	0.0619	0.0668	-0.0110	0.0037
0.25	0.0697	0.0770	-0.0100	0.0056
0.30	0.0774	0.0795	-0.0075	0.0016
0.35	0.0853	0.0771	-0.0035	-0.0059
0.40	0.0924	0.0723	0.0010	-0.0143
0.45	0.0991	0.0656	0.0063	-0.0231
0.50	0.1054	0.0569	0.0122	-0.0325
0.55	0.1100	0.0474	0.0177	-0.0414
0.60	0.1118	0.0376	0.0215	-0.0489
0.65	0.1104	0.0275	0.0233	-0.0551
0.70	0.1056	0.0170	0.0230	-0.0597
0.75	0.0972	0.0107	0.0203	-0.0580
0.80	0.0857	0.0085	0.0161	-0.0496
0.85	0.0717	0.0062	0.0110	-0.0376
0.90	0.0561	0.0037	0.0064	-0.0251
0.95	0.0378	0.0027	0.0027	-0.0132
1.00	0.0000	0.0000	0.0000	0.0000

$\alpha = 11.01$ degrees $\alpha = 9.95$ degrees

CONCLUSIONS

This report presents a computer program for the inverse problem of the annular airfoil. As a result of calculations made with this program the following conclusions can be made:

1. With the restrictions of linearized theory, duct shapes can be determined for desired pressure distributions even in the presence of a propeller.
2. For a given pressure distribution, the presence of the propeller at the duct midchord tends to increase the section angle of attack, make the camber S-shaped, and move the point of maximum thickness toward the trailing edge.
3. For a positive pressure on the inside of the duct, the duct thickness is increased. For a negative pressure on the inside of the duct and a positive pressure outside the duct thickness is decreased.
4. The propeller location and loading have important effects on the duct shape.
5. The computer program is quite versatile and will facilitate the design of ducted systems.

ACKNOWLEDGMENTS

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APPENDIX A - Continued

OUTPUT

NACA 66010 CT=1

CALCULATION OF THE THICKNESS DISTRIBUTION AND THE CAMBER DISTRIBUTION USING THE FOLLOWING DATA.

CHORD DIAMETER RATIO 0.80000000 WAKE FRACTION 1.00000000

ABSCISSA OF THE PRESSURE DISTRIBUTION

INNER PRESSURE

INDUCED VELOCITY TERMS
AXIAL RADIAL

-0.	1.0000000	1.0000000	-0.0382000	-0.0157200
0.00500	0.1040000	0.1040000	-0.0384000	-0.0158000
0.00750	0.0280000	0.0280000	-0.0386000	-0.0160000
0.01250	-0.0230000	-0.0230000	-0.0388000	-0.0163000
0.02500	-0.0780000	-0.0780000	-0.0389000	-0.0176000
0.05000	-0.1250000	-0.1250000	-0.0410000	-0.0180000
0.07500	-0.1540000	-0.1540000	-0.0426000	-0.0199000
0.10000	-0.1740000	-0.1740000	-0.0449000	-0.0200000
0.15000	-0.1980000	-0.1980000	-0.0480000	-0.0230000
0.20000	-0.2150000	-0.2150000	-0.0525000	-0.0280000
0.25000	-0.2260000	-0.2260000	-0.0564000	-0.0297000
0.30000	-0.2360000	-0.2360000	-0.0555000	-0.0147000
0.35000	-0.2430000	-0.2430000	-0.0400000	-0.0166000
0.40000	-0.2490000	-0.2490000	-0.0330000	-0.0176000
0.45000	-0.2550000	-0.2550000	-0.0210000	-0.0186000
0.50000	-0.2610000	-0.2610000	-0.	-0.0195000
0.55000	-0.2650000	-0.2650000	0.0210000	-0.0186000
0.60000	-0.2700000	-0.2700000	0.0330000	-0.0176000
0.65000	-0.2500000	-0.2500000	0.0400000	-0.0166000
0.70000	-0.1900000	-0.1900000	0.0555000	-0.0147000
0.75000	-0.1210000	-0.1210000	0.0564000	-0.00697000
0.80000	-0.0320000	-0.0320000	0.0525000	-0.0280000
0.85000	0.0210000	0.0210000	0.0480000	-0.0230000
0.90000	0.0960000	0.0960000	0.0449000	-0.0200000
0.95000	0.1790000	0.1790000	0.0410000	-0.0180000
1.00000	0.2710000	0.2710000	0.0382000	-0.0157000

APPENDIX A - Continued
 OUTPUT

THE FOLLOWING TABLE IS THE DISTRIBUTION OF THE DESIRED FOIL
 AT A 5.34884 DEGREE ANGLE OF ATTACK.

ABSCISSA	THICKNESS	CAMBER
LEADING EDGE	0.	0.
0.05	0.013274	0.005393
0.10	0.020276	0.010640
0.15	0.025718	0.015720
0.20	0.030249	0.020667
0.25	0.034266	0.024266
0.30	0.038213	0.024375
0.35	0.042190	0.022269
0.40	0.045733	0.019137
0.45	0.049012	0.015378
0.50	0.051928	0.010850
0.55	0.053906	0.006127
0.60	0.054231	0.001526
0.65	0.052461	-0.002947
0.70	0.048536	-0.007205
0.75	0.042579	-0.008922
0.80	0.035285	-0.008012
0.85	0.027224	-0.006227
0.90	0.018823	-0.004351
0.95	0.010426	-0.002209
TRAILING EDGE	0.000053	-0.000003

APPENDIX B - FORTRAN LISTING OF COMPUTER PROGRAM

```

FOR
DIMENSION POUT(37),PIN(37),PAX(37),PRAD(37),THETA(37),THETAB(101)
1,THICK(101),OMEGA(37),CPRIME(101),C(21),S(21),Z(21),X(37)
2,DUM(21)
READ 199,NCASE
199 FORMAT (I4)
DO 198 ICASE = 1,NCASE
READ 99
99 FORMAT (72H1
1
READ 10, H,WXD,M,N,NZ,NP,NH,IJK,J1,IPUN
10 FORMAT (2F10.6,8I4)
READ 11, (X(I),POUT(I),PIN(I),PAX(I),PRAD(I), I=1,NZ)
11 FORMAT (5F14.8)
WAKE = 1.0 - WXD
PRINT 99
PRINT 13, H,WAKE
13 FORMAT (98H0 CALCULATION OF THE THICKNESS DISTRIBUTION AND THE CAM
1BER DISTRIBUTION USING THE FOLLOWING DATA. //22H CHORD DIAMETER R
2ATIO 5X,15H WAKE FRACTION /7X,F12.8,9X,F12.8//17H ABSCISSA OF TH
3E 10X,16H OUTER PRESSURE 5X,16H INNER PRESSURE 10X,24H INDUCED VEL
4OCITY TERMS /23H PRESSURE DISTRIBUTION 52X,7H AXIAL 6X,8H RADIAL
5 //)
PRINT 14, (X (I),POUT(I),PIN(I),PAX(I),PRAD(I),I=1,NZ)
14 FORMAT (1H 8X,F8.5,14X,F12.7,9X,F12.7,12X,F12.7,F14.7)
DO 765 I=1,N
765 PRAD(I) = -PRAD(I)
PI = 3.14159265
DN = N-1
DELTA = PI/DN
NJN = N-3
DO 115 I=1,NZ
X(I) = 1.0 - X(I)
XARG = 2.0*X(I) - 1.0
THETA(I) = ACOSF(ABSF(XARG))
IF (XARG) 114,115,115
114 THETA(I) = PI - THETA(I)
115 CONTINUE
THETAB(1) = 0.0
DO 116 I=2,N
116 THETAB(I) = THETAB(I-1) + DELTA
18 IF (ABSF(POUT(I) -PIN(I)) - .0001) 210,210,211
210 IZ = 1
GO TO 212
211 IZ = 2
212 DO 25 I=IZ,NZ
OMEGA(I) = PI*(.25*(POUT(I) + PIN(I)) - PAX (I))
25 PAX (I) = .5*(POUT(I) - PIN(I))*SINF(THETA(I))
C PAX IS NOW THE VELOCITY DIFFERENCE TERM TIMES SINE( THETA ).
GO TO (27 ,225),IZ
225 OMEGA(1) = -PI*PAX(1)
DO 224 I=IZ,NZ
THETA(I-1) = THETA(I)
224 PAX(I-1) = PAX(I)
NZ1 = NZ-1
CALL DISCOT (0.0,0.0,THETA,PAX,PAX,-020,NZ1,0,PAXT)
DO 231 I=1,NZ
POUT(I) = THETA(I)
231 PIN(I) = PAX(I)
DO 230 I=2,NZ

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      THETA(I) = POUT (I-1)
230 PAX(I) = PIN(I-1)
      THETA(1) = 0.0
      PAX(1) = PAXT
27 Z(1) = 0.0
      DO 30 I=2,21
30 Z(I) = Z(I-1) + .05
      CALL THKDIS (DELTA, DN, H, M, N, NJN, NZ, PI, THICK, YV18 230
1PAX, OMEGA, NP, NH, THETA, THETAB, Z, S, IJK)
      CALL CAMBER (DELTA, DN, H, N, NJN, NZ, PI, THICK, WAKE, PAX, PRAD, J1, THETA
1B, THETA, ATTACK, CPRIME, S, Z)
      CALL SOLVE (IJK, PI, THETAB, CPRIME, N, Z, C)
      PRINT 98, ATTACK
98 FORMAT (67H1 THE FOLLOWING TABLE IS THE DISTRIBUTION OF THE DESIRE
1D FOIL AT A F9.5,25H DEGREE ANGLE OF ATTACK. /15X,9H ABSCISSA 1
20X,10H THICKNESS 8X,8H CAMBER //)
      PRINT 150, S(21), C(21)
150 FORMAT (11X,14H LEADING EDGE 9X,F10.6,F16.6)
      DO 777 I=2,20
      JIP = 22-I
777 PRINT 97, Z(I), S(JIP), C(JIP)
97 FORMAT (10X,F12.2,12X,F10.6,F16.6)
      PRINT 151, S(1), C(1)
151 FORMAT(10X,15H TRAILING EDGE 9X,F10.6,F16.6)
      IF (IPUN - 1) 198,198,599
599 PUNCH 600, (Z(I), I=1,21)
      DO 610 I=1,21
      JIP = 22- I
610 DUM(I) = C(JIP)
      PUNCH 600, (DUM(I), I=1,21)
      DO 611 I=1,21
      JIP = 22- I
611 DUM(I) = S(JIP)
      PUNCH 600, (Z(I), I=1,21)
      PUNCH 600, (DUM(I), I=1,21)
600 FORMAT (9F8.5)
198 CONTINUE
      CALL ENDJOB YV18 310
      END YV18 320
      FOR YV181010
      SUBROUTINE THKDIS (DELTA, DN, H, M, N, NJN, NZ, PI, THICK, PAX, OMEGA, NP, NH
1, THETA, THETAB, Z, S, IJK)
      DIMENSION ALPHA(40,40), B(40,101), GAMMA(40,40), FM(40), FTHETB(101),
1OMEGA(37), THETAB(101), THICK(101), PAX(37), THETA(37)
      EQUIVALENCE (GAMMA, ALPHA)
      ERROR = 0.00001
      CALL KERNEL (B, DELTA, H, M, N, PI, ERROR, KEEP) YV070250
C LOOP FOR CALCULATING A SERIES OF F(THETA BAR) VALUES BY INCREMENTING YV070260
C THETA BAR FROM 0.0 TO PI. YV070270
      DO 52 I=1, N
      IF (I-1) 52, 52, 51 YV070290
51 IF(I-N) 54, 52, 52 YV070315
54 SINB = 1.0/(2.0*PI*SINF(THETAB(I))) YV070330
52 CALL USEGM (THETAB(I), SINB, I, PI, NP, N, DN, H, THETA, PAX, NZ, OMEGA, NH,
1FTHETB(I))
      CALL CALCAL (ALPHA, B, DELTA, KEEP, N, THETAB, PI) YV070430
      CALL CALCFM (B, DELTA, FM, FTHETB, KEEP, N) YV070440
C LOOPS FOR PREPARATION OF COEFFICIENTS OF LINEAR SYSTEM FOR MATINV YV070520
61 DO 70 I=1, KEEP YV070530
      DO 70 J=1, KEEP YV070530

```

```

IF (I-J) 71,72,71
72 GAMMA(I,J) = 1.0 - ALPHA(I,J)/PI YV070540
GO TO 70
71 GAMMA(I,J) = -ALPHA(I,J)/PI YV070560
70 CONTINUE
CALL MATINV (GAMMA,KEEP,FM,1,X,1D) YV070570
IF(2-ID) 75,76,75
76 PRINT 77 YV070590
77 FORMAT (117H0COEFFICIENT MATRIX IS SINGULAR. THE INTEGRAL EQUATIOYV070600
1N IS EITHER INSOLVABLE OR HAS AN INFINITE NUMBER OF SOLUTIONS. ) YV070610
CALL ENDJOB
C LOOPS FOR CALCULATION OF THE SLOPE OF THE THICKNESS DISTRIBUTION. YV070700
75 DO 90 I=1,N
IF (I-1) 92,92,85
85 IF(I-N) 86,93,93
92 THICK (I) = FTHETC(I) + FM(1)/PI
DO 150 J = 2,KEEP
DJ = J
150 THICK (I) = THICK (I) + FM(J)*DJ/PI
GO TO 90
93 THICK (I) = FTHETB(I) + FM(1)/PI
DO 149 J = 2,KEEP
DJ = J
149 THICK (I) = THICK (I) - (-1.0)** J*FM(J)*DJ/PI
GO TO 90
86 SINT= SINF(THETAB(I)) YV070720
SUM = FM(1) YV070730
DO 91 J=2,KEEP YV070740
DJ = J YV070750
91 SUM = SUM + FM(J)*SINF(DJ*THETAB(I))/SINT YV070760
THICK (I) = FTHETB(I) + SUM/PI
90 CONTINUE
CALL SOLVE (IJK,PI,THETAB,THICK ,N,Z,S)
RETURN YV181080
END YV181090
FOR YV14 10
SUBROUTINE KERNEL (B,DELTA,H,M,N,PI,ERROR,KEEP)
DIMENSION B(40,101),CONSK(101),EE(101),FB(101),THETA1(101),
1THETA5(201),FBCOS(201) YV140030
C LOOP FOR VARYING THETA 1 IN CALCULATION OF FOURIER COEFFICIENTS FOR YV140040
C APPROXIMATION OF THE KERNEL. YV140050
KEEP = 1 YV140060
DO 200 I=1,N YV140070
IN = 1 YV140080
C LOOP FOR CALCULATION OF ACTUAL FOURIER COEFFICIENTS -- B(THETA 1). YV140090
DO 202 L=1,M YV140100
DL = L YV140110
II = 1 YV140120
IF (I-1) 700,700,710 YV140130
700 THETA1(1) = 0.0 YV140140
C THETA1 VARIES FROM 0.C TO PI, YV140150
GO TO 755 YV140160
710 THETA1(I) = THETA1(I-1) + DELTA YV140170
755 COS1 = COSF(THETA1(I)) YV140180
IF (10-L) 740,741,741 YV140190
741 GO TO (745,760),II YV140200
745 IN = 2 YV140210
NN = 51 YV140220
NNN = 51 YV140230
ANGLE = PI/50.0 YV140240

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FIRST = 0.0	YV140250
JJJ = 1	YV140260
GO TO 750	YV140280
740 IF (20-L) 742,743,743	YV140290
743 GO TO (746,746,760),IN	YV140300
746 IN = 3	YV140310
DO 201 J=1,50	YV140320
LESS = 103 - 2*J	YV14 325
MINUS = 52-J	YV14 326
THETA5(LESS) = THETA5(MINUS)	YV140330
201 FB(LESS) = FB(MINUS)	YV140340
NN = 100	YV140350
NNN = 101	YV140360
ANGLE = PI/100.0	YV140370
FIRST = ANGLE	YV140380
JJJ = 2	YV140400
GO TO 750	YV140410
742 GO TO (747,747,747,760),IN	YV140420
747 IN = 4	YV140430
DO 205 J=1,100	YV140440
LESS = 203 - 2*J	YV14 445
MINUS = 102 - J	YV14 446
THETA5(LESS) = THETA5(MINUS)	YV140450
205 FB(LESS) = FB(MINUS)	YV140460
NN = 200	YV140470
NNN = 201	YV140480
ANGLE = PI/200.0	YV140490
FIRST = ANGLE	YV140500
JJJ = 2	YV140520
750 CALL INTGRL (ANGLE ,CONSK,COS1,EE,FB,FIRST,I,II,L, NN,NNN, YV140530	
1THETA,1(I),THETA5,JJJ,M,H)	
760 DO 203 J=1,NNN	YV140560
203 FBCOS(J) = FB(J)*COSF(DL*THETA5(J))	YV140570
C INTEGRAND VALUES FOR INTEGRATION WITH RESPECT TO THETA5	YV140580
SUMB = FBCOS(1) + 4.0*FBCOS(NNN-1) + FBCOS(NNN)	YV140590
NJN = NNN-3	YV140600
DO 204 J=2,NJN,2	YV140610
204 SUMB = SUMB + 4.0*FBCOS(J) + 2.0*FBCOS(J+1)	YV140620
B(L,I) = 2.0*ANGLE*SUMB/(3.0*PI)	YV140670
C B IS THE VALUE OF THE FOURIER COEFFICIENT FOR EACH VALUE OF THETA 1.	YV140680
IF (ABS(B(L,I)) - ERROR) 790,790,202	YV140690
C CHECK FOR THE CONVERGENCE OF THE FOURIER COEFFICIENTS	YV140700
202 CONTINUE	YV140710
KEEP = M	
GO TO 200	YV140770
790 L = L+1	
DO 771 LL=L,M	YV140840
771 B(LL,I) = 0.0	YV140850
KEEP1 = L - 1	
IF (KEEP - KEEP1) 795,200,200	
795 KEEP = KEEP1	
200 CONTINUE	YV140860
RETURN	YV140870
END	YV140880
FOR	
SUBROUTINE USEGM (THETAB,SINB,I,PI,NP,N,ON,H,THETA,WDIFFR,NZ,OMEGA	
1,NH,SUM)	
DIMENSION THETA4(101), WDIFFR(37),OMEGA(37),ANSWER(101),	YV21 030
1T(101),F(101),R(507),A(100),THETA(37)	YV21 040
IF (I-1) 501,501,500	YV21 050

501 LIFT = 1	YV21 060
NP1 = NP/2	
NP2 = NP1 + 1	
P = NP1	YV21 070
DELANG = P1/P	YV21 080
THETA4(1) = 0.0	YV21 090
ANSWER(1) = 0.0	YV21 100
C LOOP FOR CALCULATING SERIES OF FUNCTIONAL VALUES FOR HARMONIC ANALYSIS	YV21 110
DO 510 J = 2, NP2	YV21 120
THETA4(J) = THETA4(J-1) + DELANG	YV21 130
TLOW = -CUBERTF(P! - THETA4(J))	YV21 140
TUP = CUBERTF(THETA4(J))	YV21 150
DELTAT = (TUP - TLOW)/DN	YV21 160
SINZ = SIN(THETA4(J))**2	YV21 170
C LOOP FOR EVALUATION OF INTEGRATION WITH RESPECT TO T.	YV21 180
DO 515 K=1,N	YV21 190
IF (K-1) 516,516,517	YV21 200
516 T(1) = TLOW	YV21 210
GO TO 555	YV21 220
517 T(K) = T(K-1) + DELTAT	YV21 230
555 IF (ABS(T(K)) - 0.0001) 556,556,557	YV21 240
556 F(K) = 0.0	YV21 250
GO TO 515	YV21 260
557 X = THETA4(J) - T(K)**3	YV21 270
COSZ = COS(THETA4(J))	YV21 280
COSX = COS(X)	YV21 290
C CALCULATION OF ELLIPTIC CONSTANT K	YV21 300
CONK = SQRTF(1.0/(.25*(H*(COST - COSX))**2 + 1.0))	YV21 310
C SUBROUTINE FOR CALCULATION OF ELLIPTIC INTEGRALS.	YV21 320
CALL WORK (CONK, THETA4(J), T(K), H, EK, EE, DIF)	YV21 330
C SUBROUTINE FOR INTERPOLATION OF VELOCITY DIFFERENCE VALUES.	YV21 350
CALL DISCOT (X, X, THETA, WDIFFR, WDIFFR, -120, NZ, 0, WDIFFR)	
C CALCULATION OF INTEGRAND VALUES FOR INTEGRATION WITH RESPECT TO T.	YV21 370
F(K) = H*WDIFR*CONK*DIF*3.0*T(K)**2 /4.0	85
515 CONTINUE	YV21 390
SUMT = F(1) + 4.0*F(N-1) + F(N)	YV21 400
NTN = N-3	YV21 410
DO 560 L=2,NTN,2	YV21 420
560 SUMT = SUMT + 4.0*F(L) + 2.0*F(L+1)	YV21 430
C SUBROUTINE FOR INTERPOLATION OF MEAN VELOCITY MINUS AVERAGE INDUCED	YV21 450
C VELOCITY VALUES	YV21 460
CALL DISCOT (THETA4(J), THETA4(J), THETA, OMEGA, OMEGA, -120, NZ, 0, ZOMEGA)	
510 ANSWER(J) = (DELTAT*SUMT/3.0 + ZOMEGA)*SINZ	YV21 480
DO 520 J=2, NP1	
NMINUS = NP - J + 2	
520 ANSWER(NMINUS) = ANSWER(J)	
C SUBROUTINE FOR CALCULATING FOURIER COEFFICIENTS.	YV21 530
CALL GMHAS (NP, NH, ANSWER, R(507))	YV21 540
MAX = 507	YV21 550
DO 580 L=1, NH	YV21 560
IL = MAX-5*L	YV21 570
580 A(L) = R(IL)	YV21 580
500 SUM = A(1)/(2.0*PI)	YV21 630
LEFT=LIFT	
GO TO (570, 571), LEFT	YV21 640
570 DO 572 L=2, NH	YV21 650
DL = L	YV21 660
572 SUM = SUM + DL*A(L)/(2.0*PI)	YV21 670
LIFT = 2	YV21 680

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GO TO 599
571 IF (I-N) 575,574,574
575 DO 576 L=2,NH
DL = L
576 SUM = SUM + A(L)*SINF(DL*THETAB)*SINB
GO TO 599
574 DO 577 L=2,NH
DL = L
577 SUM = SUM - (-1.0)**L*DL*A(L)/(2.0*PI)
599 RETURN
END
FOR
SUBROUTINE CALCAL(ALPHA,B,DELTA, KEEP,N,THETA1,PI
DIMENSION ALPHA(40,40),B(40,101),THETA1(101)
NN=(N-3)/2
DO 352 LL=1,KEEP
DO 350 L=1,KEEP
DL=L
SUMAL = B(LL,1)*DL + 4.0*B(LL,N-1)*SINF(DL*THETA1(N-1))/SINF(THETA1(VV070120
11(N-1))- B(LL,N)*DL*COSF(DL*PI)
DO 351 J=1,NN
351 SUMAL = SUMAL + 4.0*B(LL,2*J)*SINF(DL*THETA1(2*J))/SINF(THETA1(2*J(VV070160
1)) + 2.0*B(LL,2*J+1)*SINF(DL*THETA1(2*J+1))/SINF(THETA1(2*J+1))
350 ALPHA(LL,L) = DELTA*SUMAL/3.0
352 CONTINUE
RETURN
END
FOR
SUBROUTINE CALCFM(B,DELTA,FM,FTHET1, KEEP,N
DIMENSION B(40,101),FM(40),FTHET1(101)
NN=(N-3)/2
DO 301 L=1,KEEP
SUMFM=B(L,1)*FTHET1(1)+4.0*B(L,N-1)*FTHET1(N-1)+B(L,N)*FTHET1(N)
DO 300 I=1,NN
300 SUMFM=SUMFM+4.0*B(L,2*I)*FTHET1(2*I)+2.0*B(L,2*I+1)*FTHET1(2*I+1)
301 FM(L) = DELTA*SUMFM/3.0
RETURN
END
FOR
SUBROUTINE CHEAT (THETA4,T,H,DIF)
CON = H*(T**6/2.0*COSF(THETA4) - T**3*SINF(THETA4))
CK = CON/SQRTF(4.0+CON**2)
IF (CK) 930,931,930
930 ALAM = LOGF(ABSF(4.0/CK))
DIF = ALAM - 1.0 - ALAM*CK**2/4.0
RETURN
931 DIF = 0.0
RETURN
END
FOR
SUBROUTINE WORK (CONK,THETA4,T,H,EK,EE,DIF)
IF (CONK - 0.995) 940,940,941
940 CALL ELLIP (0.0,CONK,ZZ,EK,EE)
GO TO 981
941 IF (CONK - 0.999999) ,942,943
942 CALL VELLIP (CONK,EE,EK)
981 DIF = EK - EE
RETURN
943 CALL CHEAT (THETA4,T,H,DIF)
EK = DIF + 1.0

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EE = -1.0	YV07 110
945 RETURN	YV07 120
END	YV07 130
FOR	YV07 010
SUBROUTINE EELLIP (CONSK,EE)	YV07 020
IF (CONSK - 0.995) 940,940,941	YV07 030
940 CALL ELLIP (0.0,CONSK,Z,ZZ,ZZZ,EE)	YV07 040
RETURN	YV07 050
941 YK = 1.0 - CONSK**2	YV07 060
IF (YK) 942,942,943	YV07 070
943 ALAM = LOGF(4.0/SQRTF(YK))	YV07 080
EE = 1.0+(2.0*ALAM-1.0)*YK/4.0 + 3.0*(ALAM-13.0/12.0)*YK**2/16.0	YV07 090
RETURN	YV07 100
942 EE = 1.0	YV07 110
RETURN	YV07 120
END	YV07 130
LOAD BATCH	
FOR	
SUBROUTINE VELLIP (XK,EE,EK)	YV01 10
C EVALUATION OF COMPLETE ELLIPTIC INTEGRALS WHEN K SQUARED IS GREATER	YV01 20
C THAN .99 BUT NOT EQUAL TO 1.0 BY ASYMPTOTIC SERIES.	YV01 30
YK = 1.0 - XK**2	YV01 40
ALAM = LOGF(4.0/SQRTF(YK))	YV01 50
EE = 1.0 + (2.0*ALAM-1.0)*YK/4.0 + 3.0*(ALAM-13.0/12.0)*YK**2/16.0	YV01 60
EK = ALAM + (ALAM-1.0)*YK/4.0 + 9.0*(ALAM-7.0/6.0)*YK**2/64.0	YV01 70
RETURN	YV01 80
END	YV01 90
FOR	YV33
SUBROUTINE SOLVE(JK,PI,THETAB,F,N,Z,STORE)	YV33
DIMENSION THETAB(101),F(101),STORE(21),Z(21)	YV33
IJK=JK	YV33
DO 101 I=1,20	YV33
DIJK=IJK-1	YV33
ZARG=2.0*Z(I)-1.0	YV33
ANG=ACOSF(ABS(ZARG))	YV33
IF(ZARG)134,135,135	YV33
134 ANG=PI-ANG	YV33
135 DEL=ANG/DIJK	YV33
CALL DISCOT(DEL,DEL,THETAB,F,F,-120,N,0,RESULT)	YV33
STORE(I)=F(I)+4.0*RESULT	YV33
IJK=IJK-1	YV33
DO 100 J=3,IJK,2	YV33
DJ=J-1	YV33
T=DJ*DEL	YV33
CALL DISCOT(T,T,THETAB,F,F,-120,N,0,RESULT)	YV33
STORE(I)=STORE(I)+2.0*RESULT	YV33
T=(DJ+1.0)*DEL	YV33
CALL DISCOT(T,T,THETAB,F,F,-120,N,0,RESULT)	YV33
100 STORE(I)=STORE(I)+4.0*RESULT	YV33
CALL DISCOT(ANG,ANG,THETAB,F,F,-120,N,0,RESULT)	
STORE(I)=(STORE(I)+RESULT)*DEL/3.0	YV33
101 IJK=IJK-1	YV33
STORE(21)=0.0	YV33
RETURN	YV33
END	YV33
FOR	YV18D 10
SUBROUTINE CWORK (CONK,TERM,EK,EE,DIF)	YV18D 20
IF (CONK - 0.995) 940,940,941	YV18D 30
940 CALL ELLIP (0.0,CONK,ZZ,ZZZ,CX,EE)	YV18D 40
GO TO 981	YV18D 50

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941 IF (CONK - 0.999999) 942,942,943                                YV18D 60
942 CALL VELLIP (CONK,EE,EK)                                       YV18D 70
981 DIF = EK - EE                                                 YV18D 80
    RETURN                                                         YV18D 90
943 IF (TERM) 944,945,944                                         YV18D100
945 DIF = 1.0                                                       YV18D110
    GO TO 946                                                       YV18D120
944 TERM2 = TERM**2                                               YV18D130
    ALAM = LOGF(ABSF(4.0*SQRTF(TERM2 + 1.0)/TERM))                YV18D140
    DIF = ALAM - 1.0 - ALAM*TERM2/(4.0*(TERM2 + 1.0))           YV18D150
946 EK = DIF + 1.0                                                YV18D160
    EE = -1.0                                                       YV18D170
    RETURN                                                         YV18D180
    END                                                             YV18D190
    FOR                                                             YV18F 10
    FOR                                                             YV320010
    SUBROUTINE CAMBER (DELTA,DN,H,N,NJN,NZ,PI,THICK,WAKE,WDIFFR,PRAD YV320020
    I,J1,THETAB,THETA,ATTACK,CPRIME,S,Z)
    DIMENSION THICK(101),WDIFFR(37),PRAD(37),THETAB(101),THETA(37), YV320040
    1CPRIME(101),CINT1(101),CINT2(101),CINT3(101),          VDIFR(101), YV320050
    2S(21),Z(21)
    N1 = N-1                                                       YV320080
    RADIAL = 0.0                                                   YV320090
    DO 50 I=2,N1                                                  YV320100
    CALL INTRL1 (CINT1(I),DELTA,H,                                NJN,NZ,PI,WAKE YV320110
    1,WDIFFR,THETA,N,THETAB,VDIFR,I)                             YV320120
    CALL MUSK (DELTA,THETA,WDIFFR,NZ,N,H,NJN,CINT3(I),THETAB,
    1 I,WAKE,PI,VDIFR)                                           YV320140
    CALL INTRL2 (N1,I,THETAB,H,Z,S,DELTA,THICK,NJN,N,CINT2(I),PI)
    GO TO (48,50),J1                                             YV320170
48 CALL DISCOT (THETAB(I),THETAB(I),THETA,PRAD,PRAD,-120,NZ,0,RADIAL) YV320180
50 CPRIME(7-I) = (CINT1(I) + CINT2(I) + CINT3(I) - RADIAL/WAKE ) * YV320190
    1SINF(THETAB(I))/2.0                                         YV320200
    DO 49 I=2,N1                                                  YV320210
49 THETAB(I-1) = THETAB(I)                                       YV320220
    N2 = N-2                                                       YV320230
    CALL DISCOT (0.0,C.0,THETAB,CPRIME,CPRIME,-020,N2,0,FIRST) YV320240
    DO 51 I=1,N2                                                  YV320250
    NMINUS = N-I                                                  YV320260
    NLESS = N1-I                                                  YV320270
    THETAB(NMINUS) = THETAB(NLESS)                               YV320280
51 CPRIME(NMINUS) = CPRIME(NLESS)                                 YV320290
    THETAB(1) = 0.0                                               YV320300
    CPRIME(1) = FIRST                                             YV320310
    CALL DISCOT (PI,PI,THETAB,CPRIME,CPRIME,-020,N1,0,CPRIME(N)) YV320320
    TANA = CPRIME(1) + 4.0*CPRIME(N1) + CPRIME(N)                YV320330
    DO 52 I=2,NJN,2                                              YV320340
52 TANA = TANA + 4.0*CPRIME(I) + 2.0*CPRIME(I+1)                YV320350
    TANA = DELTA*TANA/3.0
    ATTACK = 57.295780*ATANF(TANA)                                YV320370
    DO 53 I=1,N                                                  YV320380
53 CPRIME(I) = -CPRIME(I) + SINF(THETAB(I))*TANA/2.0           YV320390
    RETURN                                                         YV320400
    END                                                             YV320410
    FOR                                                             YV33
    SUBROUTINE INTRL2(N1,I,THETAB,H,Z,S,DELTA,THICK,NJN,N,SUMI2,PI)
    DIMENSION THETAB(101),F2(101),THICK(101),STHETA(101),S(Z1),Z(21) YV33
    IF(I-2)290,290,291                                           YV33
290 DO 292 J=2,N1                                                YV33
    ZTHETA=.5*(1.0+COSF(THETAB(J)))                               YV33

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292 CALL DISCOT(ZTHETA,ZTHETA,Z,S,S,-130,21,0,STHETA(J))          YV33
291 DO 200 J=2,N1
    IF(I-J)210,211,210
211 F2(J)=0.0                                                    YV33
    GO TO 200                                                    YV33
210 PART=COSF(THETAB(I))-COSF(THETAB(J))                        YV33
    PARTH=.5*H*PART                                             YV33
    CONK2=SQRTF(1.0/(PARTH**2+1.0))                             YV33
    CALL CWORK(CONK2,PARTH,EK,EE,DIF)                            YV33
    F2(J)=STHETA(J)*CONK2**3*PART*DIF*SINF(THETAB(J))          YV33
200 CONTINUE                                                    YV33
    SUMI2=4.0*F2(N-1)                                           YV33
    DO 201 J=2,NJN,2                                           YV33
201 SUMI2=SUMI2+4.0*F2(J)+2.0*F2(J+1)                            YV33
    SUMI2=.25*H**2*DELTA*SUMI2/3.0                              YV33
    DO 202 J=1,N
    IF(J-1)250,250,251
250 F2(J) = -STHETA(I)*SINF(THETAB(I))/(COSF(THETAB(I)) - COSF(THETAB(
    1J)))                                                         YV33
    GO TO 202                                                    YV33
251 IF(N-J)250,250,252
252 IF(I-J)253,254,253
254 F2(J) = THICK(I) + STHETA(I)*COSF(THETAB(I))/SINF(THETAB(I))
    GO TO 202                                                    YV33
253 PART=COSF(THETAB(I))-COSF(THETAB(J))                        YV33
    PARTH=.5*H*PART                                             YV33
    CONK2=SQRTF(1.0/(PARTH**2+1.0))                             YV33
    CALL ELLIP(CONK2,EE)                                         YV33
    F2(J)=(STHETA(J)*SINF(THETAB(J))*CONK2**3*EE
    1 - STHETA(I)*SINF(THETAB(I)))/PART
202 CONTINUE                                                    YV33
    STORE=F2(1)+4.0*F2(N-1)+F2(N)                                YV33
    DO 203 J=2,NJN,2                                           YV33
203 STORE=STORE+4.0*F2(J)+2.0*F2(J+1)                            YV33
    STORE= DELTA*STORE/3.0                                       YV33
    SUMI2=-H*(SUMI2+STORE)/PI                                    YV33
    RETURN
    END                                                            YV33
    FOR                                                            YV30 10
    SUBROUTINE INTRL1 (CINT1,DELTA,H,                            NJN,NZ,PI,WAKE,WDIFFRYV30 20
    1,THETA, N,THETAB, VDIFR,I)                                YV30 30
    DIMENSION WDIFFR(37),THETA(37),THETAB(101),F1(101),VDIFR(101)
    IF (I-2) 202,202,203
202 DO 212 J=1,N
212 CALL DISCOT (THETAB(J),THETAB(J),THETA,WDIFFR,WDIFFR,-020,NZ,0
    1,VDIFR(J))
203 DO 200 J=1,N
    IF (I-J) 232,231,232
231 F1(J) = 0.0
    GO TO 200
232 PART = COSF(THETAB(I)) - COSF(THETAB(J))
    PART1 = .5*PART*H
    CONK1 = SQRTF(1.0/(PARTH**2 + 1.0))
    CALL CWORK (CONK1,PARTH,EK,EE,DIF)
    F1(J) = VDIFR(J)*PART*CONK1*DIF
200 CONTINUE
    CINT1 = F1(1) + 4.0*F1(N-1) + F1(N)
    DO 250 I=2,NJN,2
250 CINT1 = CINT1 + 4.0*F1(I) + 2.0*F1(I+1)
    CINT1 =-H**2*DELTA*CINT1/(12.0*PI*WAKE)

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RETURN
END
FOR
SUBROUTINE INTGRL (ANGLE ,CONSK,COS1,EE,FB,FIRST,I,II,L,
INNN,THETA1,THETA5,JJJ,M,H)
DIMENSION CONSK(101),THETA5(101),EE(101),FB(101)
C LOOP FOR CALCULATION OF INTEGRAND VALUES FOR INTEGRATION WITH
C RESPECT TO THETA 5
DO 101 J=JJJ,NN,JJJ
IF (J-JJJ) 800,800,810
800 THETA5(J) = FIRST
C THETA5 VARIES FROM 0.0 TO PI
GO TO 811
810 THETA5(J) = THETA5(J-1) + ANGLE
811 DEN = COSF(THETA5(J)) - COS1
IF (ABS(DEN) - 0.000001) 850,850,856
850 FB(J) = 0.0
GO TO 101
856 CONSK(J) = SQRTF(1.0/(0.25*(H*DEN)**2 + 1.0))
C CALCULATION OF ELLIPTIC CONSTANT K.
CALL EELLIP (CONSK(J),EE(J))
FB(J) = (1.0 - 1.0/COSK(J) * EE(J))*(SINF(THETA5(J)))**2/DEN
C PARTIAL INTEGRAND VALUE FOR INTEGRATION WITH RESPECT TO THETA5.
101 CONTINUE
RETURN
END
FOR
SUBROUTINE MUSK (DELTA, THETA,WDIFFR,NZ,N,H,NJN,CINT3,THETAB,
I,WAKE,PI,VDIFR)
DIMENSION THETA(37),WDIFFR(37),F3(101),VDIFR(101),THETAB(101)
DO 10 J=1,N
IF (I-J) 11,12,11
12 ABOVE = THETAB(I) + .01
CALL DISCOT (ABOVE,ABOVE,THETA,WDIFFR,WDIFFR,-020,NZ,0,FAB)
BELOW = THETAB(I) - .01
CALL DISCOT (BELOW,BELOW,THETA,WDIFFR,WDIFFR,-020,NZ,0,FBEL)
F3(J) = (FAB - FBEL)/(.02*SINF(THETAB(I)))
GO TO 10
11 DENOM = COSF(THETAB(I))- COSF(THETAB(J))
CCONK = SQRTF(1.0/(0.25*(H*DENOM)**2 + 1.0))
CALL EELLIP (CCONK,EE)
F3(J) = (VDIFR(J)*CCONK*EE - VDIFR(I))/DENOM
10 CONTINUE
CINT3 = F3(1) + 4.0*F3(N-1) + F3(N)
DO 40 J= 2,NJN,2
40 CINT3 = CINT3 + 4.0*F3(J) + 2.0*F3(J+1)
CINT3 = DELTA*CINT3/(6.0*PI*WAKE)
RETURN
END

```

```

YV30 320
YV150010
NN, YV150020
YV150040
YV150050
YV150060
YV150070
YV150080
YV150090
YV150100
YV150110
YV150120
YV150130
YV150140
YV150170
YV150190
YV150200
YV150210
YV150220
YV150230
YV150240
YV150250
YV07 850
YV150310
YV31 10
YV31 20
YV31 40
YV31 50
YV31 90
YV31 110
YV31 170
YV31 180
YV31 190
YV31 210
YV31 220
YV31 230
YV31 240
YV31 250
YV31 260
YV31 320

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