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ELECTROCHEMISTRY OF FUEL CELL ELECTRODES

FINAL REPORT

Contract No. Nonr-3765(00)

ARPA ORDER NO. 302-62

Project Code 9800

Task No. NR 359-443

Prep od for:

OFFICE OF NAVAL RESEARCH Materials Sciences Division Washington 25, D.C.

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I. ABSTRACT

This report is essentially an extended abstract of the main results of the research program initiated in April 1962 under Contract Nonr-3765(00), ARPA Order No. 302-62, and carried out between that time and October, 1965. The program dealt with the electrochemistry of fuel cell reactions, and included studies of the hydrogen reaction, the oxidation of organic fuels, and oxygen reduction. The clarification of the role of the electrode in catalyzing fuel cell reactions was the central theme of this work.

In order to establish the relation of electrocatalytic activity to electrode composition it is necessary first to understand in detail the mechanism of the reaction of interest. The simplest reaction studied was the hydrogen reaction, and in this case, it was possible to show that electrocatalytic activity is dependent on the electronic structure rather than the surface composition of the electrode. The oxidation of organic fuels and oxygen reduction are more complicated reactions and the work with these reactions was directed primarily towards elucidating their kinetics and mechanisms.

A summary of the main experimental results, of the conclusions reached, and of the implications of these studies in the development of fuel cells is given in Section II of this report. Detailed abstracts of the work done under this contract, and published in various journals, are given in Section III.

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II. SUMMARY OF RESEARCH AND FINDINGS

A. Electrocatalysis

1. Electrocatalytic Activity and Electrode Composition

In attempting to understand, control, and eventually predict the dependence of electrocatalytic activity on electrode composition, it is convenient to distinguish between geometric, chemical, and electronic factors. The geometry of a catalyst surface is always important and in some cases crucial in determining the extent of interaction between reactants and substrate. The roles of the chemical composition of the surface and of the electronic structure of the bulk material are less directly perceived and more difficult to establish. It is particularly difficult to separate these factors and to establish unambiguously their relative importance. However, a distinction between these factors is extremely useful in providing a framework within which a rational search for superior electrode materials can be conducted.

By a "chemical factor" in electrocatalysis we mean: specific interactions between surface atoms and reactants which lead to distinct intermediate chemical compounds. In contrast, we speak of "electronic factors" if the bulk electronic structure of the catalyst or of the surface layer is important in the catalytic reaction. The chemical, or atomic, approach to electrocatalysis implies that reactivity is determined by the intrinsic chemical properties of the individual surface atoms, and is only slightly influenced by neighboring atoms in the crystal. This approach appears to be valid in certain cases: for example, in the chemisorption of oxygen on metalloid surfaces, and in the oxidation and dissolution of covalent intermetallic compounds. On the other hand, the electronic, or continuum, approach has been used extensively to interpret chemisorption and catalysis by transition metal alloys; and is particularly relevant to metals and semiconductors whose electronic structure (expressed, for example, by their electrical conductivity) appears to be intimately related to their catalytic properties.

A study of the relative contribution of atomic and of continuum factors in the electrocatalysis of the hydrogen reaction was undertaken using amalgam electrodes. The hydrogen reaction was chosen because it is relatively simple. The determination of the rates of individual steps in its mechanism is also straightforward. By using a dropping liquid amalgam electrode we were able to obtain a clean, reproducible surface whose chemical composition and electrical properties could be determined unambiguously.

The first alloy systems studied were homogeneous, liquid solutions of platinum and palladium in mercury. Both platinum and palladium are extremely active catalysts for the hydrogen reaction, whereas mercury is a very poor one. Using the exchange currents on the pure metals, we calculate that if platinum atoms were to exert their electrocatalytic effect independent of their surroundings, a nearly saturated amalgam (0. 03 mole % Pt) would have an exchange current of about 3 x 10^{-7} amp/cm^2 , as compared to only 1 x 10⁻¹² for pure mercury. Similarly, palladium amalgams (0.006 mole % Pd) are expected to have an exchange current of about 6 x 10^{-9} amp/cm². Surprisingly, our experiments showed that the hydrogen reaction on these amalgams proceeded at essentially the same rate as on pure mercury. In attempting to understand the complete absence of catalytic effects by platinum and palladium when these are dispersed in mercury, we examined the elect: ode capacity as a function of frequency. This measurement gives the extent of adsorption of hydrogen atoms formed as intermediates in the over-all reaction; the results showed that the discharge step on platinum atoms dissolved in mercury must be at least ten times slower than on platinum atoms in pure platinum.

These experiments showed that the electrocatalytic properties of platinum and palladium are greatly modified when they are present as dilute solutions in mercury, and therefore that the atomic approach to catalysis of this reaction is a poor approximation to reality. Thus they provided strong support for a continuum approach in which the electronic structure of platinum plays a dominant role.

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In order to understand in greater detail the continuum factors which determine electrocatalytic activity, we studied the hydrogen reaction using amalgams of indium and thallium, which have wide ranges of solubility in mercury, and of gallium. These studies were directed towards establishing the relation of the reaction rate to the surface charge on the electrode, the potential drop in the double layer, of the bulk electronic properties of the electrode (such as the work function). In addition, measurements of electrode capacity and interfacial tension permitted us to obtain the Gibbs surface excess of the amalgam components and to establish whether any relation exists between this quantity and catalytic activity.

The results in the case of the indium amalgams confirmed the continuum approach. Although the surface composition calculated from the Gibbs surface excess showed complex maxima and minima, the exchange current and enthalpy of activation were smooth monotonic functions of bulk composition. For comparisons, the electronic properties of the amalgams, such as electrical conductivity, temperature coefficient of resistivity, Hall Coefficient, and optical reflectivity, all are smooth monotonic functions of composition. These results confirm our previous conclusion that the actual surface composition of a homogeneous electrode has very little influence on the electrocatalytic properties of that electrode: these are determined by electronic properties which reflect the structure of the bulk of the material, and which are cooperative properties of the components, rather than independently additive ones.

One important conclusion which results from a rigorous analysis of the influence of electrode composition on the rate of the hydrogen evolution reaction concerns the influence of the work function of the electrode material on the exchange current for the reaction. From the knowledge that the rate-determining step is transfer of an electron from the electrode to a hydrogen ion at the surface, one might conclude that the work function is a primary factor in determining the rate, but such

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is not the case A cancellation arises if the zero-charge potential of the electrode is a linear function of the work function, with unit slope. This relation holds if the dipolar structure of the interface is independent of electrode composition, as has often been suggested in the literature; but the evidence to back it up has been sketchy, and of low accuracy. We have made accurate measurements of the zero-charge potentials of indium amalgams over the entire range of composition from 0 to 70% indium, and the correlation of these data with accurate measurements of the photo-electric work functions of these same amalgams is crucial in determining the influence which work function may have on electrocatal in the entire ranges in the dipolar structure of the amalgam and the aqueous solution at the metal-electrolyte interface.

2. The Role of the Double Layer in Electrode Reactions

All electrode reactions take place in a region within a few angstroms of the interface between metal and solution, and as we have seen, an understanding of the structure and properties of this region is crucial to a rational approach to electrocatalysis. We have carried out studies of the electrical double layer on indium and thallium amalgams to supplement our studies of the hydrogen reaction.

Two important questions have been answered by these studies. First: what effect does a shift in zero-charge potential of the electrode have on the rate of the hydrogen discharge reaction? We have shown that the effect of the diffuse double layer is relatively large, but that differences between amalgams are small primarily because the larger double layer of the more concentrated amalgams tends to cancel the effect of a smaller difference between the electrode potential and zerocharge potential.

Second: is the specific adsorption of ions on the electrode appreciably influenced by the electrode composition? We have shown, by studying the zero-charge potential and double layer capacity of indium amalgams in perchloric acid and sodium fluoride solutions, that although

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the fluoride ion is specifically adsorbed on mercury only at very large positive surface charges, it is strongly adsorbed on indium amalgams even at regative surface charges. The perchlorate ion was found to be much less strongly adsorbed on indium amalgams than was the fluoride ion, even though the reverse is true in mercury. These effects show that the composition of the electrode is crucial in determining specific adsorption, and that the properties of the double layer on electrodes cannot in general be predicted from results for mercury. Further studies with other electrode materials will make useful correlations possible, but the influence of electrode composition on the structure of the electrical double layer is a field which has barely been touched.

Our studies of electrocatalysis of the hydrogen reaction show therefore that the continuum electronic structure of the electrode is the principal factor in determining electrode activity: it controls the reaction rate by controlling the extent of adsorption of intermediates and by determining the composition of the double layer in any given solution. Therefore, research on fuel cell catalysts must emphasize the role of electronic structure and must concentrate on intermetallics and other compounds where a controlled variation of electronic structure can be achieved.

B. Kinetics of Oxidation of Organic Fuels

Meaningful correlations between electrode composition and electrode activity presuppose a detailed knowledge of the reaction mechanism. The kinetics of oxidation of organic fuels are complex and largely unknown. Therefore, we undercook to study a relatively simple reaction - the oxidation of HCOOH to CO_2 - which could be used as a model reaction in electrocatalytic studies with organic fuels.

A significant problem in studying these reactions is the difficulty of obtaining reproducible measurements. This is due to the difficulty of defining the surface structure and, most importantly, to the sluggishness of attaining a steady state for ad-species. To overcome this, very careful preparation and purification techniques were used and, in

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addition, a procedure of electrode pre-anodization was adopted to ensure a well-defined, clean surface. Both galvanostatic cleaning pulses, and potentiostatic pulses, which avoid any possible interference from O_2 evolved during the cleaning process, were used to pretreat the electrode. ÷.

The crucial and significant findings of the work with formic acid were: i) adsorption onto Pt from acid solutions of HCOOH is slow - much slower than the diffusion-limited rate; ii) the adsorption rate is potentialdependent; iii) the product of adscrption interfers with the oxidation process.

The self-poisoning characteristic of this reaction is of general significance to electrochemical oxidation of organic fuels since it is encountered with all organic fuels. Giner has described a species, which he calls "reduced CO_2 ", which is formed on the reductive adsorption of CO_2 on Pt. He found a qualitative similarity between this species and a species which is adsorbed on Pt from solutions of a number of simple oxygenated organic fuels. Recently, he has shown that the same species is found on working hydrocarbon electrodes. The studies carried out by us show that this conclusion may be too simple, but it does point up the intimate relation between these reactions and the reactions of hydrocarbons.

Further work was thus centered on efforts to identify this adsorbing "poison" for HCOOH oxidation. The experimental work was primarily concerned with trying to make an accurate determination of the characteristics of the film adsorbed on smooth Pt from HCOOH solutions and comparing them with those found for CO, which has been suggested as the poison. (This CO is supposed to originate from the decomposition of HCOOH).

The difficulty in characterizing Q_{HCOOH} (the charge to oxidize the species adsorbed from HCOOH solutions) is that the species can only be oxidized in a potential region where the electrode itself oxidizes. Thus, there are three faradaic processes which can contribute to the anodic charge passed during the oxidation of the adsorbate. These are: (1) oxidation of adsorbate, (2) oxidation of species in solution, and (5) surface oxidation of the electrode. In order to separate these and to

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determine (1), a new method was developed, involving the application of a series of pulses to the surface of the electrode. Then, an effective separation of these processes can be toade for HCOOH and CO, and accurate values of Q_{HCOOH} (285 μ coul/cm²) and Q_{CO} (365 μ coul/cm²) can thus be made. The difference between Q_{HCOOH} and Q_{CO} and a comparison of the oxidation kinetics showed that CO_{ads} and "HCOOH_{ads}", while similar, are not the same. The question of the identity of the adsorbing poison is, thus, still open.

These studies form a basis on which to attempt to establish the detailed kinetics of oxidation of organic fuels eventually leading to the efficient electrochemical oxidation of common fuels, similar to gasoline or diesel oil. It appears that platinum will not turn out to be the best catalyst for this electrochemical reaction and that alternative electrodes must be found. This suggests in turn that less drastic conditions than concentrated phosphoric acid and elevated temperatures must be developed, particularly if materials outside the noble metal group are to be explored.

C. Oxygen Reduction

In a number of fuel cell systems, such as H_2/O_2 , the oxygen electrode limits the attainable efficiency and power density. Oxygen reduction is a highly irreversible process with complex kinetics which are still largely unknown even for simple electrodes. The kinetics are further complicated by side reactions (e.g. H_2O_2 production) and by reactions with the substrate (surface oxide formation). In order to clear up at least some of the ambiguities in studies of this reaction, we undertook to examine the general role of surface oxide films in determining the kinetics of electrochemical reactions at electrodes. This work was extended later to include studies of the kinetics of formation and reduction of oxides and particularly of the factors which determine their steady state thickness and their electrochemical properties.

Aside from specific effects, e.g. changes in the energy of adsorption of reactants or of intermediates, superficial oxides may influence the kinetics of an electrochemical reaction by providing a barrier to electron

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transfer between the metal and an ion or molecule at the oxide/electrolyte interface. A simple way of examining the general electrochemical effects of surface oxides is to compare the kinetics of an elementary electron transfer reaction on oxide-covered and oxide-free electrodes. Specific effects, e.g. adsorption, should be at a minimum for the redox couple chosen. The ferric-ferrous couple was selected in this study because its kinetic behavior on oxide-free electrodes is straightforward. Overpotential measurements for this couple on passive nickel, iron, chromium, titanium, and iron-chromium alloys were carried out in solutions of fixed ionic strength but of varying pH. Tafel lines were generally obtained with exchange currents etween 10^{-7} and 10^{-5} amp/cm² and cathodic transfer coefficients of about 0.45. The anodic transfer coefficients were less, particularly with Ti and Fe electrodes. A limiting anodic current, which was unrelated to diffusion in solution, was observed under certain conditions on Fe and Ti. In general, the surface oxide film has rectifying properties, the easy direction of electron flow being from metal to solution. This rectification is additional to the usual faradaic rectification observed with most electrochemical reactions. The apparent transfer coefficients, calculated from the anodic and cathodic polarization curves, yield sums significantly less than unity. The results demonstrate the existence of a potential drop across the surface oxide film, the magnitude of the drop being dependent on the thickness and composition of the film.

These studies suggest that one condition for optimum performance of a material as an oxygen electrode is either the absence of an oxide film or the restriction to a minimum thickness (e.g. by suitable alloying or choice of electrolyte) of any surface oxide film which may be formed in the relevant potential region.

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III. ABSTRACTS OF WORK PERFORMED UNDER THIS CONTRACT

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A. Electrocatalysis

 Electrochemical Behavior of Nickel Compounds I. The Hydrogen Evolution Reaction on NiSi, NiAs, NiSb, NiS, NiTe₂, and Their Constituen. Elements

A. K. M. Shamsul Huq and Arthur J. Rosenberg

Submitted as Technical Memorandum No. 1 (April, 1963). Presented at the Boston Meeting of the Electrochemical Society, May, 1963. Published: J. Electrochem. Soc. 111, 270-278 (1964).

The objective of the studies was to clarify the influence of crystal environment on the electrochemical, and particularly the electrocatalytic, properties of a transition metal atom. We chose to examine several compounds of nickel with elements from groups IVB, VB, and VIB of the Periodic Table. The use of such compounds was favored for two reasons: the clear chemical differences between the component elements, and the highly ordered atomic coordination.

The cathodic polarization behavior of NiSi, NiAs, NiSb, NiTe₂, and NiS at room temperature was determined in 1M perchlorate solutions at pH 0. 04 and pH 10. 8 up to -1. 0v vs. the reversible hydrogen electrode. Measurements were also made on Ni, Si, Sb, Te, and on Ni electrodes exposed to As_2O_3 and H_2S . Sb, Te, Si, and Ni + As_2O_3 exhibit kinetic peculiarities which are probably attributable to hydride formation, but these features are not generally shown by the corresponding nickel compounds. With the exception of Si in acid solution, all of the electrodes show an extended Tafel line indicating the hydrogen evolution reaction with the following characteristics. In alkaline solution, the Tafei slope is greater than 100 mv for all metals, while the exchange current is between 10^{-5} and 10^{-8} amp/cm². In acid solution, Ni, Sb, NiSi, NiSb, and Ni + As_2O_3 have slopes close to 120 mv, while Te, NiTe₂, Ni + H_2S , and NiS show slopes between 50 and 80 mv. The exchange currents in $a^{-1}d$ solution vary from about 10^{-5} for Ni to about 10^{-11} amp/cm² for Te. The data are shown in Figs. 1-1 - 1-6 and Table 1-1.

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Fig. 1-i. Hydrogen evolution reaction on nickel in 1M HClO₄, pH 0.04 and in 1M NaClO₄ + 10⁻² NaOH, pH 10.8 at 25°C. Solution preelectrolyzed at 30 ma for 6 hr. Results of ten runs at pH 0.04 with mean b = 0.125v, $1_0 = 2 \times 10^{-6}$ amp/cm² and of aix runs at pH 10.8 with b = 0.105v, $1_0 = 1 \times 10^{-5}$ amp/cm².





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TABLE 1-1

Tafel Parameters for the Hydrogen Evolution Reaction

	pH	H 0.04	pH	<u>H 10. 8</u>	
		-log i _o ,		-log i _o ,	
	<u>b, mv</u>	amp/cm ²	<u>b, mv</u>	amp/cm ²	
Ni	125	5.3	105	5.0	
$Ni + As_2O_3$	120	8.0	no effec	no effect on Ni	
NiAs	56	7.3	155	6.1	
Si	no Tafe	l region	168	7.5	
NiSi	113	5.4	145	5.4	
Sb	170	5.7	260	5.3	
NiSb	108	7.6	245	5.0	
Ţe	48	11	-	-	
NiTe ₂	55	8	-	-	
$Ni + H_2S$	80	9.4	no effec	t on Ni	
NiS	62	7.2	100	4.9	

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Fig. 1-3. Cathodic polarization curves for NiAs and Ni + 10^{-3} M As₂O₃ in 1M HClO₄, pH 0.04, and in 1M NaClO₄ + 10^{-2} M NaOH, pH 10.8 at 25°C. At pH 0.04, I, Ni; II, NiAs, mean of 5 runs; IH, Ni with ad·lition of ~ 10^{-3} M As₂O₃ in solution; (A stirred, B, unstirred). At pH 10.8, IV, Ni, no effect of ~ 10^{-3} M As₂O₃ added to solution. V, NiAs (A, "up"; B, "down").



Fig. 1-4. Cathodic polarization curves for NiSb and Sb in 1M HClO₄, pH 0.04 and in 1M NaClO₄ + 10^{-2} M NaOH, pH 10.8 at 25°C. At pH = 0.04, i, Ni; H. NiSb; IH. Sb. At pH = 10.8, iV. Ni; V. NiSb; VI. Sb. (A, "up"; and B, "down").

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Fig. 1-5. Cathodic polarization curves in 1M HClO₄, pH 0.04 and in 1M NaClO₄ + $i0^{-2}$ M NaOH, pH i0.8 at 25°C. At pH = 0.04. 1, Ni; 11, NiS (odor of H₂S in the effluent gas); ⁻⁻⁻ Ni in presence of H₂S, added as Na₂S. At pH = 10.8, IV. Ni; V. NiS, no effect of addition of ~ 10^{-3} M Na₂S.



Fig. 1-6. Cathodic polarization curves in 1M HClO₄, pH 0.04, at 25°C. I. Ni; II. NiTe₂; III. Tc.

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In the cases of NiAs, NiTe₂, and NiSi in acid solution, there are potential regions where the current densities on the compounds are greater than on either of the component elements. This immediately rules out the simple atomic description since the surface concentration of either component, regardless of crystallographic orientation, does not exceed that of its elemental form. For the remaining cases under study, it is also evident that the over-all rate of the hydrogen evolution reaction is not a specific atomic property, and that interactions between the surface atoms or between adsorbed hydrogen atoms at neighboring sites on the crystal surface play a prominent role.

It was concluded, therefore, that the electrocatalysis of the hydrogen evolution reaction on nickel compounds cannot be explained in terms of the intrinsic chemical properties of the surface atoms. Further work along these lines, together with complementary studies of the electronic properties and gas phase chemisorption characteristics, is needed to elucidate the electronic factors governing electrocatalytic activity. 2. The Hydrogen Reaction on Nickel in Alkaline Solutions

A.C. Makrides

Submitted: J. Electrochem. Soc.

The hydrogen reaction on nickel was studied at a number of NaOH and Na₂SO₄ concentrations at 30° C. The exchange current, Tafel slope, stoichiometric number, and hydrogen oxidation current depend on the polarization history of the electrode. The initial electrode state, characterized by a relatively low exchange current and high Tafel slope and by a stoichiometric number of 2 is transformed, upon anodic-cathodic cycling, to a state with a higher exchange current, low Tafel slope, and a stoichiometric number of 1. These changes are attributed to the formation of a superficial oxide which is not easily reduced. This oxide is in some respects unique since it enhances the rate of cathodic hydrogen evolution although it inhibits hydrogen oxidation. The original electrode state can be regenerated by extreme cathodic polarization (> 100 ma/cm²).

3. Hydrogen Evolution at Dilute Platinum and Palladium Amalgam Electrodes

J. N. Butler and A. C. Makrides

Submitted as Technical Memorandum No. 7 (October, 1963). Presented at the Toronto Meeting of the Electrochemical Society May, 1964.

Published: Trans. Faraday Soc. 60, 938-946 (1965).

Studies of electrocatalysis by solid alloy electrodes are complicated by the problem of obtaining a clean, reproducible surface. Almost any operation designed to remove adsorbed material from the surface of a solid alloy changes the structure or composition of the surface to some extent. For this reason, studies of electrocatalysis by liquid alloys at a dropping electrode provide a more direct approach to the influence of alloy composition on electrode activity. We have made the present study to determine the extent to which platinum or palladium exerts an electrocatalytic effect for the evolution of hydrogen, when present as a dilute alloy in mercury, an inactive metal.

We have measured hydrogen overpotential and differential capacity on an amalgam containing 0.028 atom percent platinum and an amalgam containing 0.006 atom percent palladium. A conventional three-compartment pyrex cell was fitted with a capillary in place of the working electrode, and housed in an air thermostat. All measurements were made in 0.100 molar perchloric acid at 25°C. The electrolyte was saturated with hydrogen throughout the experiment. Extreme care was taken to eliminate any impurities in the solution.

Measurements of hydrogen overpotential were made using an electronic potentiostat to fix the potential of the drop with respect to a reversible hydrogen electrode in the same solution, and the current was displayed on an oscilloscope. From the known growth rate of the drop, the drop area was calculated. Galvanostatic measurements were made by fixing the current passing between the drop and a counter electrode,

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and measuring the potential of the drop with respect to a reference electrode with the oscilloscope. Both methods gave the same results for current densities higher than 10^{-4} amp/cm², and the results for mercury agreed with the best literature values.

Differential capacity measurements were made using a platinized platinum screen which surrounded the drop. The method was essentially that of Grahame [J. Am. Chem. Soc. <u>71</u>, 2975 (1949)]. Measurements were made at a fixed point in the drop lifetime, balancing a series resistance and capacitance decade in an impedance bridge. For mercury, the values agreed with those obtained by Grahame.

As can be seen from Figs.3-1 and 3-2, the hydrogen overvoltage and differential capacity values obtained with the platinum and palladium amalgams were identical with that of pure mercury within the experimental error. The differential capacity showed no frequency dependence between 500 cycles and 10 kilocycles.

Platinum and palladium are extremely active catalysts for the hydrogen evolution reaction, and mercury is a relatively poor one. 'The exchange current on platinum is 10^{-3} amp/cm² and on mercury 1.5 x 10^{-12} amp/cm². If platinum, atoms were to exert their electrocatalytic effect independently of their surroundings, an alloy containing 0.03 atom percent platinum in a catalytically inactive metal might be expected to have an exchange current about 0.03% of that observed on pure platinum, or 3×10^{-7} amp/cm². This is 2×10^{-5} times larger than the exchange current actually observed. For the palladium amalgam, since the exchange current is 4×10^{-3} larger than the value actually observed.

From the measurements of differential capacity, we can set an upper limit for the exchange current for the discharge reaction

 $H^{+} + e^{-} = H$ (ads).

This must be smaller than 1.6×10^{-5} amp/cm² for the platinum amalgam. The expected value for this quantity if platinum atoms behaved independently



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Fig. 3-1. Hydrogen overpotential or platinum imalgam (0.028 mole %), palladium amalgam (0.006 mole %), and pure mercury. O, platinum amalgam: \bullet , palladium amalgam: Δ , pure mercury. The Tafel line has a slope of 118 inv and an exchange current of 1.5 ± 10^{-12} A/cm².



Fig. 3-2. Differential capacity on platinum amalgam (0-028 mole %) and palladium unalgam (0-006 mole %). The line represents the data for pure mercury.

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of their surroundings is more than ten times as large. Furthermore, the over-all rate of the hydrogen evolution reaction at a platinum atom in the amalgam cannot be more than 100 times greater than the over-all rate of hydrogen evolution at a mercury atom, or a significantly lower over-potential would have been observed. When normalized to the fraction of the surface actually occupied by platinum atoms, this result shows that the over-all rate of hydrogen evolution at a platinum atom in the amalgam is approximately 10⁻⁷ of what it is at a platinum atom in pure platinum. An analogous series of arguments can be made regarding the palladium amalgam.

These experiments on the electrocatalysis of hydrogen evolution by dilute platinum and palladium amalgams have shown that the electrocatalytic properties of platinum and palladium are greatly modified when they are present as dilute solutions in mercury. This lends support to the idea that electrocatalysis is primarily a cooperative property of the continuum structure of ar alloy rather than an independently additive atomic property.

4. <u>Hydrogen Evolution at a Dropping Indium</u> Amalgam Electrode

James N. Butler and A. C. Makrides

Submitted as Technical Memorandum No. 10 (February, 1964). Presented in part at the Philadelphia Meeting of the American Chemical Society, April, 1964. Published: Trans. Faraday Soc. 60, 1664-1676 (1954).

We have continued our studies of electrocatalysis by liquid metal alloys, using indium amalgams. Indium-mercury alloys can be prepared which are liquid at room temperature and which contain up to 63.7 atcm percent indium. We have measured the hydrogen overpotential, the double-layer impedance, and the potential of zero charge on a series of these amalgams varying in composition from 0.0001% to 63.7%.

A conventional three-compartment pyrex cell was fitted with a capillary in place of the working electrode, and housed in an air thermostat. All measurements were made in 0.100 molar perchloric acid at 25° C. The electrolyte was saturated with hydrogen throughout the experiment. Extreme care was taken to eliminate any impurities in the solution.

Measurements of hydrogen overpotential were made using an electronic potentiostat to fix the potential of the drop with respect to a reversible hydrogen electrode in the same solution, and the current was displayed on an oscilloscope. From the known growth rate of the drop, the drop area was calculated. Galvanostatic measurements were made by fixing the current passing between the drop and a counter electrode, and measuring the potential of the drop with respect to a reference electrode with the oscilloscope. Both methods gave the same results for current densities higher than 10^{-4} amp/cm², and the results for mercury agreed with the best literature values. The results are shown in Fig. 4-1.

The potential of zero-charge was measured by the streaming electrode method. At these potentials, the capacity was frequency-

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Fig. 4-1. Tafel plot of experimental results for mercury, indium, and the amalgams.

- ▲, pure Hg, +, 8.4%: ●, 16.3% In: □, 30.4% In:
- ■, 52.2% ln: 0, 64.4% ln: ▲, pure in (solid)

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Fig. 4-2. Potential of zero charge shown as a function of the activity of indium in the amalgam.

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independent, indicating an ideally polarized electrode, and the values measured for mercury agreed with those in the literature. The values are shown as a function of the activity of indium in the amalgam in Fig. 4-2.

Using the values of the zero-charge potential and some preliminary measurements of the double-layer capacity, corrections were made for the effect of the diffuse double layer on the reaction, and the experimental results were found to fit the theoretical relation with a cathodic transfer coefficient $\alpha = 0.50$.

Because there is virtually no change in the impedance at large negative potentials as the composition is varied, and because the Tafel slope is independent of composition, the mechanism of the hydrogen evolution reaction on indium amalgams is probably the same as on mercury: the rate determining step being the discharge of a hydrogen ion.

Two extreme viewpoints may be taken in predicting the electrocatalytic properties of alloys. According to the "atomistic" viewpoint, the electrocatalytic properties of an alloy result from short-range interactions, and to the first approximation are a linearly additive property of the components of the alloy. At the other extreme, from the "continuum" viewpoint, the electrocatalytic properties of an alloy result from cooperative properties of the material, such as its electronic band structure. Both viewpoints have been useful in explaining electrocatalysis by alloys.

The rate of hydrogen evolution on indium amalgams at a given potential varies almost linearly with the composition; but this results from two factors: The potential across the solution-electrode interface, and the purely catalytic effects. By taking the current at the potential of zero charge, the effects of excess surface charge can be separated from the purely catalytic effects. The current at the potential of zero charge varies with composition over four orders of magnitude in a very nonlinear fashion. Thus, the purely catalytic properties of the amalgams have been shown to be a cooperative property, rather than a simple additive atomic one.

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5. Hydrogen Evolution at a Solid Indium Electrode

James N. Butler and Manfred Dienst

Submitted as Technical Memorandum No. 11 (April, 1964). Presented at the Washington Meeting of the Electrochemical Society, October, 1964; and at the Western States Combustion Meeting, Salt Lake City, October, 1964. Published: j. Electrochem. Soc. 112, 226-232 (1965).

Indium offers an unusual opportunity for a detailed study of hydrogen overvoltage, since its zero-charge point lies within the range of potentials where hydrogen evolution is the principal reaction. We have made accurate measurements of hydrogen overpotential on indium in 0.1 molar perchloric acid at temperatures from 21°C to 59.5°C. Typical results are shown in Fig. 5-1. These results have been analyzed in terms of the structure of the electrical double layer at the indiumelectrolyte interface, and the transition state theory of reaction rates.

The hydrogen evolution reaction on indium is very sensitive to traces of oxygen. In order to remove the oxide film initially present on the electrode and to obtain reproducible measurements of hydrogen overvoltage, it was necessary to polarize the electrode cathodically for 16 hours at 30°C. The oxide film was removed more rapidly at higher temperatures. Current-potential curves measured in oxygen-saturated solutions indicate that the hydrogen evolution reaction occurs at the same rate on clean or oxidized electrodes, but that oxygen reduction is more rapid on a clean electrode.

The slow step in the mechanism of hydrogen evolution is probably discharge of a hydrogen ion. The exchange current is approximately 10^{-11} amp/cm² at 30°C, and the cathodic transfer coefficient is 0.50. These values are similar to those observed on mercury and lead, for which the slow-discharge mechanism is well-established. Addition of iodide ion decreases the overvoltage; and the effect of acid concentration on overvoltage is also in accord with a slow-discharge mechanism.

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Fig. 5-*. Deviation of experimental current-potential curve from an arbitrary Tafel line with $\alpha = 0.500$ and $i_0 < 1.00 \times 10^{-11}$.



Fig. 5-2. Fit of experimental data to Eq. [2]. The temperature was 32^{9} C, and the capacity function (Eq. [1]) was used to calculate ϕ_{2} . The correlation coefficient for the least-squares straight line is 0, 9995. The best values for the parameters are $\alpha = 0, 472 \pm 0.003$ and $-\log i_{0}^{-1} = 16.82 \pm 0.05$.

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The capacity of the electrode was measured and found to be given approximately by the equation

K = 8.9 + 2.75
$$\phi_0$$
 + 1.25 ϕ_0^2 (1)

where K is the integral capacity $(\mu f/cm^2)$ and ϕ_0 is the potential with respect to the zero charge point of indium. ϕ_2 , the potential of the outer Helmholtz plane, was calculated using this capacity function and the experimental data was fitted to the Frumkin equation

$$\eta = \frac{1-\alpha}{\alpha}\phi_2 - \frac{RT}{\alpha F} \ln (i/i_0')$$
 2)

where i_0' is the corrected exchange current and is given by:

$$ln i_0' = (1 - \alpha) ln[H^+] + \frac{\alpha F}{RT} g_H + const.$$

In these equations, η is the hydrogen overpotential, α is the cathodic transfer coefficient, $[H^+]$ is the concentration of hydrogen ion in the electrolyte, and g_{H} is the free energy of adsorption of a hydrogen atom on the electrode. The values of the parameters α and i_{0} are given in Table 5-1. The fit of a typical set of data is shown in Fig. 5-2.

The temperature dependence of the current at constant potential gives the enthalpy of activation for the discharge reaction. At the reversible hydrogen potential, this is 14.1 ± 0.5 kcal/mole, and at the zero charge point it is 10 ± 0.2 kcal/mole. Using a theoretical curve calculated by Parsons and Bockris [Trans. Faraday Soc. <u>47</u>, 914 (1951)] using the transition state theory, the heat of adsorption of hydrogen atoms on indium is estimated from these values to be 50 ± 10 kcal/g atom.

TABLE 5-1

Parameters of the Current-Potential Curve for Hydrogen Evolution on Indium (Errors are 95% Confidence Limits)

	corrected ^(a)		uncorrected ^(b)		
Temp °C	α	- log i'	α	- log i	
21.0	. 389 ± . 011	10. 19 ±14	. 445 ± . 009	$10.74 \pm .12$	
25.0	. 449 ± . 003	10.86±.04	. 498 ± . 002	11. 15 + . 03	
30.0	. 489 ± . 004	10.83±.06	. 530 🔔 . 004	11. 19 ± . 07	
41.0	. 495 ± . 005	10.77 ±.07	. 538 ± . 005	11. 14 + . 07	
50.5	. 50 2 ± . 006	10.44 + .08	. 546 ± . 006	10.81 ± .09	
59.5	. 510 ± . 006	10. 22 ± . 07	. 552 + . 006	10. 57 ± . 08	

a) using the capacity given by equation (1)

b) assuming ϕ_2 - 0 in equation (2)

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5. <u>...ctivity Coefficients of Liquid Indium-Mercury</u> <u>Amalgams at 25^oC</u>

James N. Butler

Submitted as Technical Memorandum No. 9 (January, 1964). Published: J. Phys. Chem. 68, 1828-1833 (1964),

Although indium amalgams have been used in a number of electrochemical investigations, no accurate measurements of the activity coefficients in these liquid solutions are readily available. During the course of investigations at this laboratory on electrocatalysis by indium amal_bams, we had occasion to measure these activity coefficients. The results are summarized in this paper.

Potentiometric measurements were made using the following cell:

 $\ln(Hg)/0.01 \text{ m HC1G}_4, 0.01 \text{ m In}_{2}(SO_4)_{3}/Pt, H_{2}(g)$

at 25.0°C. Amalgams were prepared from triple-distilled mercury and 99.999% pure indium by weighing out the required quantities of mercury and indium and combining them under an atmosphere of nitrogen. For amalgams containing from 1 to 30% indium, the potential could be reproduced to within 0.01 mv.

The measurements of cell potential and the activity coefficients calculated from these measurements are given in Table 6-1. These were found to fit the equation

$$\log \gamma_{\rm In} = B (1 - X_{\rm Hg}^2)$$

which describes regular solutions, for which the entropy of mixing is ideal. In equation (1), γ_{In} is the activity coefficient of indium in the amalgam, referred to infinite dilution, B is a temperature-dependent parameter, and X_{Hg} is the mole fraction of mercury in the amalgam.

Using literature data for the heat of fusion and heat capacity of indium, together with the activity of pure indium, a value B = 1.50 at $25^{\circ}C$

Mole % indium	Cell emf e, mv.	Activity based on Henry's law, mole fraction	Activity, coefficient
100. 0 ^a	266.7	21.6	
70.0 ^b	266.8	21.8	31.2
64.41	264.30	16.37	25.4
63.60	263.03	14.07	22.1
52.22	258.95	8.73	16.7
50.00	258.31	8.10	16.2
39.98	251.80	3.78	9.46
30.40	245.50	1.812	5.96
30.00 ^c	245.35	1.780	5.94
20. 00 ^C	237.06	0.678	3.38
16.26	234.70	0.514	3.16
10.00 ^C	226.46	0.196	1.96
8.41	223.31	0.1358	1.612
5.00 [°]	217.00	0.0723	1.447
3.88	214.40	0.0480	1.238
1.00°	201.44	0.01057	1.059
0.920	200.68	0.00968	1.051
). 09 40	179.0	0. J00946	1.006
0		0	1.000

Emf.	Measurements	and Calculated	Activity	^v Coefficients.	25.0°

TABLE 6-1

^aPure solid indium. ^bSaturated amalgam with exclusions solid indium present: ^CMost precise measurements, ± 0.01 mv.
was calculated. The values of B obtained directly from the activity coefficient data were compared with those calculated from the emf data of other workers^(1, 2). The solutions appear to be regular throughout the composition range, as can be seen from Fig. 6-1; and the temperature dependence of the coefficient B is consistent with a zero value for the entropy of mixing (Fig. 6-2).

In both the dependence of the activity coefficient on composition, and in the ideal entropy of mixing, mercury-indium mixtures behave approximately like regular solutions. In one respect, however, they are not completely regular. A large negative volume change $(1 \text{ cm}^3/\text{mc}^2)$ is observed on mixing mercury and indium, and its magnitude suggests that a negative excess entropy of approximately 5 joules/deg. should be observed. It is possible that the compensating positive entropy arises from repulsive interactions between partially shielded In⁺ ions in the metal.

The large activity coefficients (indicating large negative deviations from Raoult's law), as well as the large negative heat of mixing, indicate that indium-mercury interactions are much stronger than mercury-mercury or indium-indium interactions. Whether or not an actual compound species such as InHg exists in the liquid phase cannot be unan.biguously decided from thermodynamic measurements alone, but such a compound certainly exists in the solid phase. Phase diagram studies show that in the composition range where liquid amalgams can be obtained at room temperature, two compounds, lnHg and InHg₆, are probably formed; but only InHg, melting at -18.5^o, is agreed upon by all investigators. The solid solution range of InHg is only about 2 mole %, which indicates the possibility of specific bond formation between In and Hg in the liquid.

Kozin² observed discontinuities in the slope of his emf -concentration curves at compositions corresponding to $InHg_3$ and InHg, which he interpreted as evidence for the formation of stoichiometric compounds in the liquid phase. Such discontinuities are possible, and when observed have been interpreted as indicating strong ordering in the liquid phase. However, only Kozin's data of part 1 show the breaks prominently. His

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Fig. 6-1. Fit of experimental activity coeffi fent data to the empirical equation log $\Upsilon_{1n} = B(1 - X_{HZ}^{-2})$. The best value of B obtained from our data (circles) is 1.54, from Sunden's data (crosses) is 1.46, from Kozin's data (part 1, triangles) is 1.64 and (part 11, squares) 1.53. The value of B calculated from the heat of fusion of indium and the activity coefficient of pure solid indium is 1.50.



Fig. 6-2. Temperature dependence of the coefficient B. The curve was calculated from the known heat of inixing of indium and mercury, with the assumption that the excess entropy of the solution was zero.

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data of part II show them to a much smaller extent and at compositions considerably different from $InHg_3$ and InHg. Neither our data nor Sunden's data show any such breaks. Since the variations which are interpreted as breaks in the slope of the emf --concentration curve are probably within experimental errors (see Fig. 6-1), they cannot be considered adequate evidence for the formation of compounds in the liquid phase.

The consistent description of the thermodynamic properties of these solutions by the theory of regular solutions argues against the formation of any specific compound other than InHg. According to statistical mechanical theory, zero excess entropy is obtained with a single potential of interaction with nearest neighbors of a given kind, and compound formation as such need not be postulated.

References:

- 1. N. Sunden, Z. Elektrochem. 57, 100 (1953).
- 2. L. F. Kozin, Tr. Inst. Khim, Nauk, Akad. Nauk Kaz. SSR <u>9</u>, 71, 81 (1962).

James N. Butler, Mary L. Mcehan, and A. C. Makrides
Submitted as Technical Memorandum No. 15 (September, 1964)
Presented at the Chicago Meeting of the American Chemical
Society, September, 1964.
Published: J. Electroanal. Chem. 9, 237-250 (1965).

Indium amalgams provide an excellent system for studying the effect of electrode composition on the structure and properties of the electrical double layer. At room temperature, liquid amalgams containing up to 70 mole % indium can be prepared, and accurate capacity measurements can be made at the clean, reproducible surface obtainable at a dropping electrode.

Differential capacity was measured in perchloric acid solutions at 25°C, by the method of Grahame []. A. C. S. <u>71</u>, 2975 (1949)]. A platinized platinum screen surrounding the drop was used as a potentialfixing electrode. Measurements were made at a fixed point in the drop lifetime, balancing a series resistance and capacitance decade in an impedance bridge. The area of the drop was calculated from the measured mass flow rate and density of the amalgam, with correction for back-pressure due to interfacial tension.

The potential of zero charge was measured by observing the potential at which no current flowed between the drop and screen. At this potential, the polarization resistance and capacity was frequency independent, indicating that the electrode was ideally polarized. Under these conditions the current results entirely from charging the double layer as the drop changes area, and hence the zero-current point corresponds to the zero-charge point.

At very small concentrations of indium (Fig. 7-1), the capacity of the double-layer is nearly identical to pure mercury; but at concentrations above 0.001%, the diffusion-controlled rate of dissolution of indium exceeds the rate of dissolution of mercury, and a large pseudo-capacity

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Fig. 7-1. Differential capacity of indium amalgams containing up to 16.3 mole % In, at i kc. (0, 15.3%; \triangle , 3.9%; \diamondsuit , 1.0%; \Box , 0.1%; **a**, 0.01%; X, 0.005%; \div , 0.002% In: •, 0.001% In (and pure Hg). The potentials given on the figure are 0.065 v more negative than those in the tables. because of the different reference electrode.



Fig. 7-2. Differential capacity of concentrated In amalgams at 10 kc. Δ , 64%; +, 40%; O, 20%.

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and polarization resistance is observed at potentials positive with respect to the zero-charge point. At negative potentials, the capacity is identical to that of mercury.

At higher concentrations of indium (Fig. 7-2), the zero-charge point shifts to more negative values. The zero-charge point of a 70% indium amalgam is approximately 0. 4 volts more negative than the zerocharge point for mercury. If capacity curves are compared at the same rational potential (potential with respect to the zero-charge point), the shape of the curves at negative values of potential remains relatively constant, but the capacity increases by about 30% as the composition changes from pure mercury to 70% indium.

The capacity curves were integrated to obtain the surface charge (Fig. 7-3). At negative values of surface charge, the behaviour of capacity as a function of surface charge is similar to the behaviour of capacity as a function of rational potential. At positive values of surface charge, however, the capacity goes through a maximum as the concentration of indium increases, as a result of changes in anion adsorption in the inner part of the double-layer. The capacity at $+10 \,\mu \, \text{coul/cm}^2$ for 65% indium amalgam is approximately twice that for mercury at the same surface charge, but the curve shows the same "hump" (attributed to changes in the dielectric properties of water) near zero surface charge in both cases.

The values obtained for the potential drop across the diffuse double layer were used in a revised interpretation of our previous measurements (part 4 of this report) of hydrogen overvoltage on indium amalgams and the influence of the electrode composition on the rate of the discharge reaction was shown to be similar to that obtained from the uncorrected current-potential curves. The Gouy-Chapman theory was adequate for calculation of the potential drop across the diffue double-layer, as indicated by the excellent fit of the experimental data to the theoretical current-potential curve.

In Fig. 7-4 we have compared the values of corrected exchange current (log i_0) obtained from a least-squares fit of the Frumkin equation

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Fig. 7-3. Comparison of capacity as a function of surface charge for mercury (X) and 63.6% In amaigam (+). Note that the "hump" occurs at approximately the same surface charge.



Fig. 7-4. Exchange current for the hydrogen evolution reaction (corrected for double-layer effects) as a function of amalgam composition. The crosses were calculated with the transfer coefficient, α , fixed at 0.500. The circles were calculated from a least-squares fit of Frumkin's equation to the data. The vertical bars indicate 95% confidence limits.

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 $\| f \|_{L^2} \leq 1$

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with values calculated assuming a fixed value of the transfer coefficient $\alpha = 0.500$, as was done in part 5 of this report. The values of $-\log i_0'$ calculated assuming $\alpha = 0.500$, show an almost linear dependence on indium concentration, whereas the values obtained by the least-squares fit show a rather non-linear dependence, and perhaps even go through a maximum at about 20% In.

This preliminary study leaves many questions still unanswered. Although we have assumed that there is no specific adsorption of ions in the Helmholtz layer, this is not rigorously correct, and perchlorate ion is almost certainly adsorbed to some extent at positive surface charges. Whether or not this specific adsorption is entirely responsible for the large change in capacity with indium concentration at constant surface charge (Fig. 7-3), we cannot unambiguously decide at the present time. The capacity difference might also result from differences in the In-Hg dipole layer in the metallic phase, or from differences in dipole orientation of water in the aqueous phase.

James N. Butler

Submitted as Technical Memorandum No. 19 (April, 1965). Published: J. Phys. Chem. 69, 3817-3825 (1965).

One of the factors which may influence the rate of an electrocatalytic reaction at an alloy electrode is the chemical composition of the surface. Because of the different surface energies of the components of an alloy, the surface composition will, in general, be different from the bulk composition and will change in a way which minimizes the total surface energy. It is difficult to determine the surface composition of a solid electrode, but, for a liquid electrode, the measurement of interfacial tension as a function of bulk composition can give important information about the surface composition.

The interfacial tension is related to the surface composition by the Gibbs augorption isotherm

$$\Gamma_{i} = -\frac{\partial \gamma}{\partial \mu_{i}}$$
(1)

where γ is the interfacial tension, μ_i is the chemical potential of a component in one phase, and Γ_i is the surface excess of that component with respect to a standard component (usually Hg in amalgams and H₂O in aqueous electrolytes). The surface excess is not a true surface concentration but is the total difference in composition of the surface region from that of the bulk. The region of integration extends from the physical surface to where the composition equals that of the bulk. Although thermodynamic measurements cannot give the true surface composition, they can, with additional assumptions, give a quantity which approximates the surface composition.

Interfacial tension was measured by the sessile drop method, using a traveling microscope which could be read to ± 0.00002 in. A complete set of measurements of one drop consisted of approximately

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20 readings of vertical coordinates, together with the 40 corresponding left and right horizontal coordinates. From these data, the interfacial tension was calculated by fitting the experimental points to a theoretical curve for the shape of a sessile drop, using a computer program (see part 9 of this report). The results obtained are given in Table 8-1.

Although interfacial tension can be measured directly as a function of potential and composition, more precise results can be obtained by integrating measurements of double-layer capacity. If the composition of both the amalgam and the aqueous phase are held constant, the differential capacity is the second derivative of the interfacial tension with respect to potential

$$C = 0.1 \frac{\partial^2 \gamma}{\partial \phi^2}$$
(2)

where C is the differential capacity in μ f./cm.², γ is the interfacial tension in erg/cm.², and ϕ is the potential with respect to a reference electrode.

We have made measurements of zero-charge potential and doublelayer capacity for indium amalgams in 0.1 M HC10₄ solutions at 25° (see part 7 of this report). With the additional experimental values of interfacial tension at the zero-charge potential, measured in the present work, we can obtain by integration the complete field of data required to calculate surface excess as a function of potential and electrode composition.

Figure 8-1 shows the surface excess curves calculated from the slopes of interfacial tension-activity curves, which were obtained by integrating the capacity measurements. Note that the surface is deficient in indium, the surface deficiency passes through a maximum at approximately 20 mole % indium, and an increase in bulk composition beyond this point causes the surface deficiency to decrease in magnitude. As with thallium amalgams (see part 12 of this report) the surface excess of indium amalgams was found to be potential dependent.

One of the principal objects of this investigation was to relate the surface composition of indium amalgams to the observed electrocatalytic

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			Parameters	s of sess shape	ile drop	
Mole % In	Potential, ^a	Density, g/cm ³	β	<u>b, cm</u>	Std. dev. ^b	, ergs/cm ²
0.0	-0.165	12.534	9.88	0.586	1.21	426.0
5 : 97	-0.310	12.168	5.15	0.433	0.82	434.5
10.00	-0.333	11.92	6.68	0.500	0.86	437.9
20.00	-0.388	11.305	12.22	0.707	1.46	454.0
39. 83	-0.480	10.03	3.75	0.432	0.85	491.0
63.02	-0.558	8.52	2.86	0.418	0.55	510.8

Interfacial Tension of Indium Amalgams in 0. 1M HClO₄ at 25°

^aZero-charge potential, with respect to a reversible H_2 electrode in 0.1M HClO₄. ^bStandard deviation of experimental points from theoretical curve, in microns.

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Fig. 8-1. Surface excess calculateo from the slopes of interfacial tension-activity curves. The scatter of the points gives some indication of the random errors involved.

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activity. The rate of the hydrogen evolution reaction, as reflected by the exchange current i_0' (see parts 5 and 7 of this report), is a smooth monotonic function of composition. In the region from 10 to 30% indium, no salient features appear which might be correlated with the maximum surface deficiency seen in Fig. 8-1.

Two alternative explanations may be made of this lack of correlation: the first, that the surface composition is essentially that of the bulk, and the second, that the electrocatalytic properties are more closely related to the bulk electronic properties of the alloy than to the surface composition.

If the composition profile in a direction perpendicular to the surface is so slowly varying that the observed surface deficiency is spread over more than 20 atomic layers, the surface composition would be essentially that of the bulk. Such a diffuse surface structure seems unlikely. No evidence for specific long-range ordering in liquid Hg-In amalgams has been found by X-ray diffraction or thermodynamic methods (see part 6 of this report). The mean tree path of electrons in mercury-indium alloys is only about 2.5 interatomic distances, so that appreciable space charge is unlikely. The diffusion coefficients of metals in amalgams are comparable to the diffusion coefficients of solutes in nonassociated solutions. Thus a more reasonable estimate would be that the surface deficiency arises in the first four or five atomic layers.

In contrast to the complex way in which the surface composition (Fig. 8-1) varies, the electronic properties of mercury-indium alloys show relatively simple behavior. The resistivity, temperature coefficient of resistivity, and Hall coefficient all decrease monotonically as the concentration of indium increases. The optical reflectivity increases monotonically with increasing indium concentration. None of these properties agrees precisely with the values calculated from a free-electron theory, even though pure mercury and pure indium both are well described by such a theory. The only property which shows a comples behavior is the thermcelectric power, which goes through an

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unexplained maximum at approximately 3 atom % indium. From our results (parts 4, 7, and 10 of this report) we have found that the exchange current and probably the enthalpy of activation for the nydrogen evolution reaction decrease monotonically as the concentration of indium increases.

We have previously mat a dichotomy between the "atomic approach" and the "continuum approach" to electrocatalysis. The atomic approach assumes that the electrocatalytic properties of an alloy result from independently additive properties of the surface atoms; the continuum approach assumes that the electrocatalytic activity of an alloy is a cooperative property and may bear little relation to the properties calculated on the basis of the atomic approach. In our studies of hydrogen evolution at platinum and palladium amalgam electrodes (part 3 of this report), we showed that the continuum approach was probably closer to the truth since the atomic approach gave an answer which was many orders of magnitude different from the experimental results.

The results of the present study on indium amalgams also indicate that the continuum approach is the fruitful one. We have seen that the electronic properties of these alloys are virtually all simple monotonic functions of the bulk composition, as are the electrocatalytic properties. The surface composition, on the other hand, shows very complex behavior and does not correlate well with electrocatalytic activity. It would thus seem likely that electrocatalytic activity is nore closely related to the cooperative electronic properties of an alloy rather than to its surface composition.

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9. <u>A Curve-Fitting Method for Calculating Interfacial</u> <u>Tension by the Sessile Drop Method</u>

James N. Butler and Burton H. Blcom

Submitted in part as Technical Memorandum No. 19 (April, 1965). Published: Surface Science 4, [1-17 (1966).

The equilibrium shape of a cylindrically symmetric sessile drop (Fig. 9-1) is governed by the equation

$$\frac{b}{R} = \left(\frac{b}{x}\right)\sin\phi = 2 + \beta\left(\frac{z}{b}\right)$$
(1)

where

$$\beta = \frac{gb^2}{\gamma} (d_1 - d_2)$$
⁽²⁾

In these equations, R is the radius of curvature in the vertical plane at the point (x, z), b is the radius of curvature at the origin, ϕ is the angle made by the radius of curvature with the z axis, g is the acceleration of gravity, d_1 and d_2 are the densities of the denser and lighter fluid, respectively, on each side of the interface, and γ is the interfacial tension.

This form of the equation is especially useful because the linear dimensions R, x, and z appear as ratios to the maximum radius of curvature b. Thus, a given value for the constant B gives the same drop shape, regardless of drop size, and changing b influences only the size of the drop without altering its shape. Of course, changing the size of the drop while holding the surface tension and density constant causes both β and b to change simultaneously.

Both R and ϕ can be expressed in terms of derivatives of z with respect to x, and Eq. 1 can be transformed into the nonlinear differential equation

$$b \frac{d^2 z}{dx^2} + \frac{b}{x} \left\{ 1 + \left(\frac{dz}{dx}\right)^2 \right\} \frac{dz}{dx} = \left(2 + \frac{\beta z}{b}\right) \left\{ 1 + \left(\frac{dz}{dx}\right)^2 \right\}^{3/2}$$
(3)

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Fig. 9-1. The shape of a typical sessile drop. The shaded area indicates the surface on which the drop rests. The curve was calculated from the equations in the text assuming $\beta = 5,397$ and b = 0.5766 cm. The crosses indicate the experimental measurements obtained for a 63,0% indum amalgam.

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This equation cannot be solved in terms of any conventional functions. Some solutions of this equation have been tabulated by Bashforth and Adams ["An attempt to Test the Theories of Capillary Action", Cambridge University Press, 1883].

For our purposes, however, we want to fit a number of experimental measurements of x, z pairs to the theoretical drop contour and to obtain the best values of b and β consistent with these experimental values. Neither the tables of Bashforth and Adams nor the empirical approximations derived by various other workers are sufficiently accurate to assure that the best possible fit of the parameters is obtained.

We have, therefore, devised a computer program which generates the function $z(x, b, \beta)$ defined by Eq. 3 by an iterative solution of the differential equation. The parameters b and β are then adjusted so that the best fit is obtained for the experimental data. The interfacial tension is then calculated using Eq. 2. An example of the fit obtained to experimental measurements is shown in Fig. 9-1.

10. <u>Temperature Dependence of Hydrogen Overvoltage</u> on Indium Amalgams

James N. Butler and Mary L. Meehan

Submitted as Technical Memorandum No. 18 (March, 1965). Published: Trans. Faraday Soc. 61, 2521-2530 (1965).

This paper presents a study of the temperature dependence of hydrogen overvoltage for amalgam electrodes containing 10, 20, 40 and 63% indium. From these results, the enthalpy of activation, which is closely related to the heat of adsorption of a hydrogen atom on the electrode surface, can be obtained, and its relation to the composition of the electrode determined. The method described previously (part 4 of this report with refinements such as charging current corrections) was used, and current-potential curves were measured for the four amalgams over the temperature range from $15^{\circ}C$ to $65^{\circ}C$. A typical set of data is shown in Fig. 10-1.

The experimental values were fitted to the Frumkin equation

$$\eta = \frac{1 - \alpha}{\alpha} \phi_2 + \frac{RT}{\alpha F} \ln\left(\frac{i}{i_o'}\right) \tag{1}$$

by the least-squares method described in part 5 of this report. The potential ϕ_2 of the outer Helmholtz plane was calculated from capacity measurements made at 25°C, assuming no specific adsorption of perchlorate ion, as described in part 7 of this report. Eq. 1 was cast into the form of a straight line, whose slope was equal to the cathodic transfer coefficient α , and whose intercept gave the corrected exchange current i_0 . In Eq. (1), R is the gas constant, T is the absolute temperature, and F is the Faraday constant.

The variations in α for the various current-potential curves cause large variations in the least-squares values of $-\log i_0'$, which obscure the influence of temperature and composition. To show the effects of temperature and composition more clearly, we have calculated

TABLE 10-1

Enthalpy of Activation

(Errors are 95% Confidence Limits)

amalgam <u>% ln</u>	ΔH_0^*	ΔH_z^* (a)	$\Delta H^{\ddagger} (\eta = 1.000 \text{ V})$
0.0 ^(b)	$22.3 \pm .4$	20.3 <u>+</u> .4	11.0 <u>+</u> .4
10.0	20.8 \pm 1.6	16.5 ± 1.6	11.2 ± 1.2
20.0	21.9 <u>+</u> .8	$17.0 \pm .8$	11.1 ± .4
40.0	21.9 <u>+</u> .6	16.0 <u>+</u> .6	11.2 ± .2
63.0	19.5 ± 1.5	12.5 ± 1.5	9.0 <u>+</u> .7
63.0 ^(c)	21.1 ± 1.2	14.2 ± 1.2	9.9 <u>+</u> 1.0

(a) extrapolated using the Frumkin equation with $\alpha = 0.500$. ΔH_z^{\ddagger} includes a contribution from the shift of η_z with temperature.

(b) calc. from data of Post and Hiskey.

(c) points at 56 and 66° discarded.

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average values of $-\log i_0'$, using a fixed value of α (0.500). These values are shown in Fig. 10-2.

The differences between values of $-\log i_0^+$ with α fixed show the variation of overvoltage with temperature and composition as clearly as would the change in η at constant i, but, in addition, contain corrections for the influence of the diffuse double layer, and provide a self-consistent extrapolation to the reversible potential.

The temperature dependence of the corrected exchange current gives a quantity which is independent of parameters related to the potential:

$$\Delta H_0^{\ddagger} = -R \frac{d \ln i_0}{d(1/T)} = \alpha h_H + \alpha F \frac{d(Q/T)}{d(1/T)}, \qquad (2)$$

where h_{H} is the enthalpy of adsorption of a hydrogen atom on the electrode surface and Q is a parameter characteristic of the electrode material. Thus, changes in the heat of adsorption of hydrogen atoms on the surface would be expected to show up as changes in the enthalpy of activation ΔH_{o}^{\ddagger} . Fig. 10-2 shows the plot of log i_o' (with $\alpha = 0.500$) as a function of 1/T. The least-squares straight lines gave the values of ΔH_{o}^{\ddagger} which are listed in Table 10-1, together with values of ΔH_{z}^{\ddagger} , the enthalpy of activation at the zero-charge potential, and the values for mercury calculated from the data of Post and Hiskey [J. Am. Chem. Soc. <u>72</u>, 4203 (1950)].

These results show that it is probable that the enthalpy of activation ΔH_0^{\ddagger} is nearly independent of the amalgam composition; although a decrease of approximately 1 kcal/mole on going from mercury to the 63% amalgam is consistent with the data. Fig. 10-3 compares the quite significant trend of exchange current at constant temperature, with the almost insignificant trend of enthalpy of activation, as the composition of the amalgam is varied.



Fig. 10-1. Temperature dependence of hydrogen overvoltage on 40.0% indium amalgam.



Fig. 10-2. Temperature dependence of exchange current. The llnes were fitted by the method of least-squares, and were used to calculate the values of $\Delta = H_0^{\dagger}$ listed in Table 10-1.



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Fig. 10-3. Comparison of the variation in exchange current with composition and the variation in enthalpy of activation with composition.

According to the Frumkin theory of slow discharge reactions, the corrected exchange current is given by

$$\ln i'_{O} = (1 - \alpha) \ln [H^{+}] - \frac{\alpha}{RT} g_{H} + Q, \qquad (3)$$

where H^+ is the concentration of hydrogen ion in the bulk of the solution, g_H is the free energy of desorption of a hydrogen atom from the electrode surface, and Q is a parameter which may depend on the composition of the electrode and the temperature.

If the change in $-\log i_0'$ were due entirely to a change in g_H , the observed shift of 0.55 in $-\log i_0'$ would correspond to a change in g_H of approximately 1.4 kcal/mole. If the entropy of adsorption of hydrogen atoms on the surface of the electrode were independent of composition, this variation could be ascribed entirely to a change in the enthalpy of adsorption h_H . Then eq. (2) predicts the ΔH_0^{\ddagger} changes by approximately 0.7 kcal/mole provided that the temperature coefficient of Q is independent of composition. This estimate is consistent with the observed change in ΔH_0^{\ddagger} .

Because of the large experimental error in ΔH_0^{\ddagger} , however, we cannot unambiguously state how much of the variation in $-\log i_0^{\dagger}$ can be ascribed to a variation in h_H . The enthalpy of adsorption of hydrogen atoms appears to be the primary factor, but contributions from the entropy of adsorption of hydrogen atoms, from the variation in the dipole structure of the double-layer at the interface, and from changes in the surface composition of the amalgam compared to its bulk composition, may all influence the rate of hydrogen evolution to a lesser extent.

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11. The Potential of Zero Charge on Indium Amalgams in Perchlorate Solutions

James N. Butler

Presented at the Buffalo Meeting of the Electrochemical Society, October, 1965.

Submitted to J. Phys. Chem.

Indium amalgams provide one of the few systems other than mercury for which the zero-charge potentials can be measured by the streaming electrode method. We have made measurements on amalgams containing from 0 to 64% indium in perchloric acid solutions from 0.01 to 1.0 molar. These data are summarized in Fig. 11-1, and confirm the measurements made earlier (see parts 4 and 7 of this report).

To test whether the values obtained by the streaming electrode inethod corresponded to true zero-charge potentials, direct measurements of surface charge on the amalgams were made by observing the transient current as a spherical drop electrode was formed. By a suitable extrapolation to zero time, the surface cnarge can be measured directly [J. N. Butler and M. L. Meehan, J. Phys. Chem. <u>69</u>, 4051 (1965)]. Zero-charge potentials obtained by this method are also shown in Fig. 11-1, and agree with the streaming electrode measurements.

For an ideally polarized electrode, the shift of zero-charge potential with electrolyte composition is given by the equation

$$\left(\frac{\partial E_z}{\ln a_+}\right)_q = \frac{RT}{F} \left[1 + 2\left(\frac{\partial q_-}{\partial q}\right)_a\right]$$

where E_z is the potential with respect to an external, fixed reference electrode, a_{\pm} is the mean activity of the 1-1 electrolyte, q is the surface charge on the electrode, q_{\pm} is the charge in solution contributed by anions, R is the gas constant, T is the absolute temperature, and F the faraday constant. If there is no specific adsorption of ions at the electrode, the charge in the solution is entirely in the diffuse double-

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Fig. 11-1. Zero charge potential of indium amalgains in perchloratic acid solution.



Fig. 11-2. Specific adsorption of perchlorate ion on indium amalgama.

layer, and the Gouy-Chapman theory may be used to calculate the dependence of q_ on q. At the potential of zero charge, this theory gives simply

$$\left(\frac{\partial q_{-}}{\partial q}\right)_{q=0} = -\frac{1}{2}$$

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which says that E_z is independent of the electrolyte composition. Such a result is obtained for sodium fluoride solutions.

Our results for perchlorate solutions are shown in Fig. 11-2 where we have plotted the coefficient $-(\partial q_{-}/\partial q)$. The fact that it is larger than 0.5 indicates that the perchlorate ion is specifically adsorbed on the electrode, even at the potential of zero-charge. Note, however, that the indium amalgams show less specific adsorption of perchlorate than does mercury.

From measurements of the temperature dependence of the zerocharge potential in a cell without liquid junction, it is possible to obtain an entropy associated with the electrical double layer at the interface. For mercury in 1.0 M KC1, where there is considerable specific adsorption, this term is 0.72 mv/deg; for mercury in 0.1 M HC10₄, where there is less specific adsorption, this term is 0.68 mv/deg; and for indium amalgams (20% and 40%) in 0.1 M HC10₄, this term is 0.55 to 0.58 mv/deg. These results agree, in indicating a lower degree of specific adsorption on the amalgam than on mercury, with the variation of zero charge potential with electrolyte composition.

These results show that the large differences in double layer structure between mercury and the amalgams (see parts 7 and 8 of this report) must be attributed primarily to the orientation of solvent molecules in the aqueous phase and electronic dipoles in the matallic phase; and not to the specific adsorption of ions from the solution.

12. The Electrical Double Layer on Thallium Amalgam Electrodes

James N. Butler

Submitted as Technical Memorandum No. 16 (October, 1964). Published: J. Electroanal. Chem. 9, 149-162 (1965).

In connection with our study of hydrogen overvoltage on thallium amalgams, we have π ade some measurements of double-layer capacity of thallium amalgam electrodes in 0.1 M HClO₄ at 25°C. We have also made a critical survey of the literature concerning the electrical double layer on thallium amalgams. In this paper, our experimental results are presented and compared with the literature values.

The methods were described previously (see part 7 of this report). The streaming-electrode method gives different values of the zero-charge potential depending on whether thallium ion is present in the solution or not. The values of Frumkin and Gorodetzkaya [Z. Physik. Chem. 136, 451 (1928)] obtained in solutions containing the equilibrium amount of $T1^+$ ion, were taken to be the most accurate, since the integrated capacity curves agreed with experimental electrocapillary curves, as can be seen in Fig. 12-1.

Our experimental values of capacity are shown in Fig. 12-2. By integrating these, and combining them with the data of Frumkin and Gorodetzkaya, values of interfacial tension as a function of potential and composition were obtained. From this field of data, the surface excess of thallium in the amalgam was calculated as described in part 8 of this report. The results are shown in Fig. 12-3. Note that the surface excess is negative: the surface of the amalgam is deficient in thallium.

Although the surface excess of thallium is known, the absolute value of the surface concentration cannot be obtained without knowledge of the shape of the concentration profile in a direction perpendicular to the surface. If all the thallium were removed from the first atomic layer, but the concentration were equal to the bulk concentration out to that point, the observed surface deficiency could be accounted for almost exactly. Of course, such a concentration profile is absurd; but if one assumes that the

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Fig. 12-2. Differential capacity of thallium amalgam electrodes: •, experimental data; O. zero-charge point. (a), 40, 5% T1; (b), 31.2% T1; (c), 10, 1% T1; (d), pure Hg.





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surface deficiency arises primarily in the first four or five atomic layers, one finds that the surface concentration is one-half to three-fourths of the bulk concentration, depending on the profile chosen.

Thus, it is rather unlikely that the observed similarity in hydrogen overvoltage between mercury and thallium amalgams (see part 13 of this report) is due entirely to a depletion of thallium at the surface. Although the surface concentration of thallium is certainly lower than the bulk concentration, the ratio of surface to bulk concentration is probably greater than 0.5, and is not very different for the different amalgams.

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13. Hydrogen Overvoltage at a Dropping Thallium Amalgam Electrode

James N. Butler, Evelyn A. Barron-Apps, and Mary L. Meehan Submitted as Technical Memorandum No. 17 (March, 1965) Published: Trans. Faraday Soc. 61, 2509-2520 (1965).

To extend our studies of hydrogen overvoltage on indium amalgams to other liquid metal alloy systems, we have made a detailed study of hydrogen overvoltage on 10% and 40% thallium amalgams, including the temperature dependence. Using data on the properties of the electrical double layer (part 12 of this report) we have taken into account the influence of the diffuse double layer on the rate of hydrogen evolution.

The experimental method was that used in the studies of indium amalgams $\sqrt{11}$ s 4 and 10 of this report). As can be seen from Fig. 13-1, only slight differences were observed between mercury and the thallium amalgams. The experimental measurements were fitted to the Frumkin equation, and the transfer coefficient was found to vary between 0.5 and 0.58, showing little trend with temperature or composition.

To observe the variations in the rate of hydrogen discharge free from the influence of the diffuse double layer, the transfer coefficient was fixed at $\alpha = 0.50$, and the corrected exchange current calculated. These results are shown in Fig. 13-2. From this temperature dependence, the enthalpy of activation was calculated, and the values obtained are summarized in Table 13-1.

Mercury and 10% thallium amalgam show enthalpies of activation which are the same within experimental error, but the values are clearly higher for the 40% thallium amalgam. The fact that discarding the two extreme points (in the 10% amalgam data) can change the values of enthalpy of activation by nearly 1.5 kcal/mole casts some doubt on the validity of the statistical limits of error. Thus, although the statistical analysis indicates a substantial probability that the enthalpy of activation is higher on 40% thallium amalgam than on mercury, the increase is not large, and may be within the experimental errors.



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Fig. 13-1. Variation in hydrogen overvoltage with composition at 25°C.



Fig. 13-2. Variation of exchange current with temperature and composition. The lines were fitted by the method of least-squares, and were used to calculate the values of ΔH_0^{\ddagger} listed in Table 13-1.

TABLE 13-1

Enthalpy of Activation

	ΔH_0^{\ddagger}	ΔH_z^{\ddagger} (a)	ΔH^{\ddagger} $(\eta = 1.00 \text{ V})$
mercury ^b	22.3 <u>+</u> 0.4	20.4 ± 0.4	11.0 ± 0.4
10.0% Tl amalgam	21.2 <u>+</u> 1.1	16.0 <u>+</u> 1.1	9.1 <u>+</u> 1.0
10.0% T1 ^C	22.7 <u>+</u> 1.2	17.5 ± 1.2	10.6 ± 0.3
40.6% Tl	23.1 ± 0.3	16.4 \pm 0.3	11.7 ± 0.5

^a extrapolated using Frumkin's equation with $\alpha = 0.500$. No correction has been made for the shift of η_z with temperature.

^b calc. from data of Post and Hiskey [J. Am. Chem. Soc. <u>72</u>, 4203 (1950)]. ^c points at 15° and 66° discarded.

14. Hydrogen Evolution and the Electrical Double Layer on Gallium and Gallium Alloy Electrodes

James N. Butler and Mary L. Meehan

Presented at the Cleveland Meeting of the Electrochemical Society, May, 1966. Submitted to Trans. Faraday Soc.

Although studies of hydrogen evolution and the electrical double layer⁽¹⁾ have been made on pure gallium, no studies of this type have been extended to alloys of gallium which are liquid near room temperature. We have made measurements of hydrogen overwoltage in 0.1 M $HC10_4$ at dropping liquid metal electrodes consisting of gallium, and alloys of gallium with mercury and indium. Although gallium-mercury alloys are liquid at 30° C, in the range from 2 atom % Hg to 97.5 atom % Hg the alloy separates into two liquid phases. Alloys of gallium with indium containing up to 20% In are liquid at 30° C.

Our overvoltage measurements at 1 ma/cm^2 on the various alloys are summarized in Table 14-1, together with previous measurements on mercury, indium, and indium amalgams for comparison. The introduction of either indium or mercury raises the overvoltage on gallium. The increase in overvoltage when 20% In is added to Ga is more than 70% of the increase between Ga and pure indium. The increase in overvoltage observed on an alloy containing 1. 64% Hg is nearly 90% of the increase betwee. Ga and pure mercury. In contrast, introduction of 0. 5% of gallium in mercury has an almost indetectible effect.

Measurements of surface charge⁽²⁾ in the electrical double layer on gallium-mercury and gallium-indium alloys showed small but significant changes from the properties of the electrical double layer on gallium. Using these values, corrections were made for the effect of the diffuse double layer on the hydrogen discharge reaction and corrected exchange currents and transfer coefficients were calculated. The variations of exchange current were qualitatively similar to the variations of overvoltage in Table 14-1, and the transfer coefficients varied between 0.5 and 0.6 with no particular trend.

TABLE 14-1

Hydrogen C	in 0.1 M HClO ₄	lpy of Activation
Electrode	Overvoltage at 1 ma/cm ² 35°C	∆H ₀ ‡ kcal
Ga	.710	17.8 <u>+</u> 1
20% In-Ga	. 903	22.9 <u>+</u> 1
In (s)	. 9 5	14.1 (?)
0.122% Hg-Ga	.747	
0. 36%	. 796	
1.64%	. 973	
99. 65%	1.020	
Hg	1.027	22.3 \pm 0.4
63% in-Hg	. 974	19.5 <u>+</u> 1.5

The temperature dependence of the corrected exchange current on gallium and the 20% indium-gallium alloy were measured, and the values obtained for the enthalpy of activation are summarized in Table 14-1. Note that although the enthalpy of activation on the indium-gallium alloy is nearly the same as for mercury, the overvoltage at a constant current is 100 mv lower - thus indicating large differences in the structure of the transition state which are probably attributable to the electronic structure of the metallic phase.

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 A. N. Frumkin, N. S. Polyanovskaya, and N. B. Grigor'ev, Dokl. Akad. Nauk SSSR <u>157</u>, 1455 (1964); Electrochim. Acta, 10, 793 (1965).
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15. On the Relation Between Heat of Adsorption of Hydrogen and Hydrogen Overpotential

A. C. Makrides

Submitted to J. Electrochem. Soc.

Suggested relations^(1, 2) between the heat of adsorption of hydrogen, ΔH_{ads} , as calculated from an equation due to Pauling^(3, 4) and the hydrogen overpotential were examined using recent, accurate data for hydrogen evolution on Pb, In, Hg, Ga, Zn, T1, and Cd. It was found that no relation in fact exists, the overpotential being unrelated to the calculat dsorption energy. The absence of such a correlation may be due to (a) a fundamental failure of the theoretical model on which it is ball or (b) a failure of the Pauling equation, or (c) negles of the change in the adsorption energy of water for different metales. A qualitative argument is advanced to show that item (c) cannot be neglected and that it may in fact account for the absence of the expected dependence of overpotential on ΔH_{ads} .

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b. Kinetics of Oxidation of Organic Fuels

16. Adsorption and Oxidation of Formic Acid on Smooth Platinum Electrodes in HC10₄ Solutions

S. B. Brummer and A. C. Makrides

Submitted as Technical Memorandum No. 8 (November, 1963). Published: J. Phys. Chem. 68, 1448-1459 (1964).

The adsorption and oxidation of formic acid, a model organic fuel, on smooth Pt electrodes at 40° C was studied in HC10₄ solutions with anodic and cathodic chronopotentiometry and with potentiostatic, steady-state techniques. The oxidation rate at each potential was determined as a function of electrode preparation, pH (-0.3 to + 1.55), and formic acid concentration (10⁻³ to 1 M).

Anodic pre-treatment activates an electrode for subsequent oxidation of HCOOH. However, the oxidation rate at a given potential declines with time since pre-anodization (Fig. 16-1). This is because of the slow adsorption of a substance which interferes with the subsequent oxidation of HCOOH. The kinetics of adsorption of this blocking substance were determined with cathodic transients (H-atom co-deposition) and are much slower than would be expected from diffusion of formic acid molecules up to the electrode (Fig. 16-2). This suggests that the poisoning species is formed as the result of a reaction of the formic acid. The oxidation of HCOOH proceeds on the surface not occupied by this poisoning species, and is first-order in bare surface (Fig. 16-3).

The extent of coverage with this blocking species declines with increase in potential (see part 17 of this report) and is essentially zero at 0.7 volts. Anodic charging curves, when the adsorbate was oxidized, also showed that this adsorbate is unlikely to be HCOOH.

Currents for HCOOH oxidation were determined at each potential of interest at various times after pre-anodization. Current-potential curves follow the Tafel relation (Fig. 16-4) but the tafel slope decreases with time after pre-anodization. The limiting value is close to 2.3RT/F. Above 0. 45 volts, the current becomes independent of potential (Fig. 16-4).



Fig. 16-1. Decline of rate of oxidation, at various potentials, as function of time since cleaning for 1 M HCOOH in 2 M $HCIO_4$.



Fig. 16-2. Adsorption of HCOOH from 1 M HClO_4 , at 0.45 v., as a function of time.

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Fig. 16-3. Oxidation of HCOCH in 1 M HCIO_4 at various potentials as a function of the coverage of the electrode.



Fig. 16-4. Log current vs. potential for oxidation of 1 M HCOOH at pH 1.9.

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This current plateau is not limited by diffusion in the solution. Similarly, the pH dependence of the oxidation rate at fixed potential verses the standard hydrogen electrode reaches a steady value at times longer than two minutes after the pre-anodization. This can be expressed as $(\partial \log i/\partial pH)_{\psi} = 1.0$. The order of the reaction with respect to HCOOH is unity.

The following mechanism was suggested to account for these observations:

$$HCOOH^{soln} \Rightarrow HCOOH^{double \ layer}$$
 (1)

HCOOH<sup>double layer
$$\rightarrow$$
 H_{ads} + COOH_{ads} (2)</sup>

$$COOH_{ads} \stackrel{fast}{\to} CO_2 + H^+ + e$$
 (3)

$$H_{ads} \xrightarrow{fast} H^+ + e$$
 (4)

If it is assumed that the active surface area, θ_A , is divided between free sites, θ_F , those covered with hydrogen atoms, θ_H , and those covered with some adsorbed blocking substance (not COOH), θ_P , then

$$\theta_{\rm F} = \theta_{\rm A} - \theta_{\rm H} - \theta_{\rm P} \tag{5}$$

From (2)

$$i = \vec{k}(HCOOH)^{\text{double layer}} \theta_{F}$$
 (6)

where \vec{k} is a rate constant, and reaction (2) is assumed to take place only on free, active sites. θ_A is taken to be independent of potential, and θ_P

Data of Fig. 16-3 indicate that some sites on the surface are more active for HCOOH oxidation.

was shown above to be independent of potential in the Tafel region, although not at higher potentials. If θ_H is determined by the equilibrium

$$H^+ + e \rightleftharpoons H \tag{7}$$

and if, in HCOOH solutions, it is given by

$$e_{\rm H} = K({\rm H}^{+}) \exp - \left(\frac{\psi F}{{\rm RT}}\right)$$
 (8)

where K is the equilibrium constant of reaction and ψ is the potential of the electrode with respect to a hydrogen electrode in the same solution, then

$$i = \vec{k} (HCOOH)^{d. l.} \left[\theta_A - \theta_P - K(H^+) \exp \left(\frac{\psi F}{RT} \right) \right]$$
(9)

It is clear that eq. (9) offers an approximate explanation of all the presently known facts. The Tafel slope, dE/d log i, is approximately 60 mv., and ($\partial \log i/\partial pH$) is about 1.0 in the Tafel region. In addition, when $\theta_H \rightarrow 0$, as it is for $E_{\psi} \gtrsim 0.35 \text{ v}$, a limiting current is observed. The order of the reaction with respect to HCOOH is unity via eq. (1) and this equation also accounts for the apparent decline of the surface activity as the potential increases. The slowly adsorbing substance is formally included as θ_p , and the fact that θ_p is not very sensitive to concentration indicates why an order of unity with respect to HCOOH is observed even in the Tafel region. In the limiting current region, the pH dependence is small and may reflect some second-order effect on the rate of reaction (2).

17. The Use of Large Anodic Gaivanostatic Transients to Evaluate the Maximum Adsorption on Platinum from Formic Acid Solutions

S. B. Brummer

Submitted as Technical Memorandum No. 13 (August. 1964).

Presented in part at the Washington Meeting of the Electro-Chemical Society, October, 1964. Published: J. Phys. Chem. 69, 562-571 (1965).

In order to investigate the nature of the adsorbed species which poisons the low potential oxidation of HCOOH, further studies of the adsorption process on smooth Pt at 40° C were made with large anodic and cathodic galvanostatic transients.

During the application of a large anodic current, i_a , three processes contribute to the charge passed, Q_a . These are (I) oxidation of adsorbate, (II) oxidation of a species in solution, and (III) surface oxidation of the electrode. By applying a cathodic current, i_c , at various times during the anodic transient, it is possible to measure III directly and, from the charge passed in deposition of H atoms prior to H₂ evolution, one can follow changes in the concentration of adsorbed HCOOH. By varying i_a , one can assess the effect of II. It is found that the cortributions of processes I, II, and III can be effectively separated. This method is illustrated in Fig. 17-1.

Typical results from the application of this novel technique are shown in Fig. 17-2. Here, we plot the fraction of the surface available for hydrogen deposition, θ_H , as a function of the anodic stripping charge. The dotted region at high θ_H indicates where electrode oxidation is first found, during the anodic transient. In general, θ_H increases linearly with Q_a , indicating that the adsorbate has a constant oxidation state over its whole range of coverage: two electrons per site are released when the adsorbate is oxidized to CO₂. The value of Q_a (with minor corrections) extrapolated to $\theta_H = 1$ (a 'clean' electrode) is taken as Q_{HCOOH}^{max} , the

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Fig. 17-2. Variation of $\theta_{\mbox{H}}$ with $\mbox{Q}_{\mbox{a}}$ in 0.1 M HCOOH. The measurements were taken from 0.3 v., 2 min. after cleaning pulse.

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maximum adsorption from HCOOH solutions. This value (285 microcoulombs/real cm²) is too large to correspond to HCOOH itself, particularly as the adsorbed species has a maximum coverage of 73% of the surface (Fig. 17-3). Electrode coverage with this species which poisons HCOOH oxidation declines as the potential is raised above approximately 0.35 volts. The decline is more rapid at lower HCOOH concentrations. Coverage is essentially zero by approximately 0.7 volts.



Fig. 17-3. Variation of $\theta_{\rm H}$ with electrode potential in HCOOH solutions. The measurements were taken 2 min. after cleaning pulse, usually with 50-ma. /cm.² pulses.

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18. <u>Galvanostatic Transient Studies of Carbon Monoxide</u> <u>Adsorption on Platinum Electrodes</u>

S. B. Brummer and J. l. Ford

Submitted as Technical Memorandum No. 14a (August, 1964). Presented in part at the Washington Meeting of the Electrochemical Society, October, 1964.

Published: J. Phys. Chem. <u>69</u>, 1355-1362 (1965).

In order to test suggestions that the adsorbed species which poisons the low potential oxidation of HCOOH is CO formed from decomposition of the HCOOH, the adsorption of CO on smooth Pt electrodes from solutions of 1 N HClO₄ saturated with 1 atm of CO at 40⁶ was studied using anodic and cathodic galvanostatic transients. It is found that when a large anodic current, i_a , is applied to the electrode, the potential increases rapidly to ~ 1.2 v. (vo). H⁺/H₂ in the same solution), falls to ~ 1.0 v, slowly starts to increase again, and finally becomes almost steady. The charge passed in the various regions of this transient is insensitive to the current density in the range 2.3 to 147 ma/cm² and, also, to the starting potential, in the range 0.1 to 0.8 v.

As in the case of HCOOH, three processes contribute to the charge passed, Q_a , during the application of such a current. These are (I) oxidation of adsorbate, (II) oxidation of a species in solution, and (III) surface oxidation of the electrode. By applying a cathodic current, i_c , at various times during the anodic transient, it is possible to measure (III) directly. From the charge passed in deposition of H atoms prior to H_2 evolution, Q_H , one can follow changes in the concentration of adsorbed CO. The effect of (II) can be assessed by varying i_a .

It is found that $\theta_H(Q_H/Q_H^{max})$ increases almost linearly with Q_a and this relation is independent of i_a (Fig. 18-2). Oxidation of the electrode is impeded, even when the adsorbed CO has been removed, but does occur during the progress of (I) (Fig. 18-2). In addition, it is a function of i_a . Thus, the coincidence of the $\theta_H - Q_a$ curves in Fig. 18-3 for the different current densities, does not mean that we are solely observing

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Fig. 18-1. Anodic galvanostatic transients taken from 0.3 v., 2 min. after cleaning pulse, with 1 atm. of CO, at ~ 50 ma. /cm.². At 25° , point B is not very well defined, but at 40° quite a sharp break is observed. The overshoot also is more pronounced at 40° .



Fig. 18-2. Variation of $\theta_{\rm H}$ with $\Theta_{\rm a}$ for various current densities. Transients were taken from 0.3 v., 2 min. after cleaning pulse, with 1 atm. of C.O. Anodic current densities: Θ , 2.35 ma./cm.²; Θ , 7.35 ma./cm.²; Θ , 73.5 ma./cm.²; +, 147 ma./cm.²: cathodic current density, 147 ma./cm.². $\theta_{\rm H}$ values are slightly corrected according to dotted line in Figure 3a.



Fig. 18-3. Variation of Q_0 with Q_a for various current densities. The transients were taken after 2 min. at 0.3 v. with 1 atm. of CO. Anodic current densities: 0, 2.35 ma./cm.²; 0, 7.35 ma./cm.²; •, 73.5 ma./cm.²; +, 147 ma./cm.²; cathodic current density, 147 ma./cm.².

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(I); there must also be a contribution from (II). The sum of (II) and (III) is evidently independent of i_a , and (II) is not limited by diffusion. Q_a is corrected for (II) and (III) from the measurements.

Then, it is found that 390 μ coulombs/real cm² are required to oxidize all of the adsorbed CO. Of this, 60 μ coulombs/cm² is weakly adsorbed and involves three electrons per Pt surface atom for oxidative removal. The remainder, from $\theta_{\rm H} = 0.1$ to $\theta_{\rm H} = 1.0$, requires 1.77 electrons per site. This would correspond to ~23% of the adsorbed CO in a doubly bonded or "bridged" structure. From cathodic charging curves it is found that the electrode is less than 2% bare at 0.3 v. The rate of adsorption of CO at low potentials is limited by diffusion in solution.

19. Comparison of Adsorbed Formic Acid and Carbon Monoxide on Platinum Electrodes

S. B. Brummer

Submitted as Technical Memorandum No. 14b (August, 1964). Presented in part at the Washington Meeting of the Electrochemical Society, October, 1964. Published: J. Phys. Chem. <u>69</u>, 1363-1365 (1965).

A comparison is made between the properties of the adsorbed films formed on smooth platinum in HCOOH and CO-saturated solutions of $HC10_4$. Both the coulometric behavior and the oxidation kinetics of the adsorbed films (compared at constant surface coverage, Fig. (19-1) suggest that while they are similar they are not the same.



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Fig. 19-1. The variation of potential with surface bareness during a 14.7 ma. /cm. 2 anodic galvanostatic transient. HCOOH measurement was taken in a 0.2 M solution in 1 N HClO₄. The CO measurement was made with 1 N HClO₄ saturated with 1 atm. of CO.

- C. Oxygen Reduction
 - 20. Electrochemical Behavior of Nickel Compounds II. Anodic Dissolution and Oxygen Reduction in Perchlorate Solutions

A. K. M. S. Huq, A. J. Rosenberg, and A. C. Makrides

Submitted as Technical Memorandum No. 3 (April, 1963). Presented at the Boston Meeting of the Electrochemical Society, May, 1963,

Published: J. Electrochem. Soc. 111, 278-286 (1964).

The anodic dissolution of Ni, Si, and Sb and of the compounds iSi, NiSb, NiAs, and NiS was studied in molar perchlorate solutions at pH 0.04 and 10.8. The results are shown in Figs. 20-1 - 20-7 and Table 20-1.

In all cases except with Si in acid solution, active dissolution at low anodic potentials is followed by a transition to a passive state where me anodic current density is nearly independent of potential. In acid solution, the critical potentials of NiSi, NiAs, and NiS are less positive than that of Ni and the critical currents are less by two to three orders of magnitude. The critical potential of NiSb is about the same as for Sb, but the critical current is substantially less. The anodic current in the passive region for the compounds is greater than for passive Ni. NiSi, NiAs, and NiS show a transpassive region in which Sb, As, and S are probably oxidized to ions of a high valence state. In alkaline solutions, NiSi, NiSb, and NiS show an active-passive transition at the same potential as pure Ni. However, the rate of oxidation of these compounds in the passive state is again considerably greater than for passive Ni. It is suggested that in acid solution the passive film on nickel compounds is either a mixed oxide or an oxide having the basic structure of the second element; in alkaline solution, nickel oxide is formed first but it probal'y has a different structure than the oxide formed on passive Ni. The verence between acid and alkaline solutions arises from the difference in the stability of nickel oxide in the two electrolytes. A stable, passive oxide on Ni is formed in acid solution only when the dissolution rate is sufficiently large to cause nickel hydroxide to precipitate at the

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Fig. 20-2. Polarization curves for Sb. Fig. 2A, pH 0.04 (IM $HCIO_4$): curve I, cathodic hydrogen evolution; curve II, anodic current for dissolution and passivation, toward increasing potentials; curve III, anodic currents, toward decreasing potentials; curve IV, cathodic currents, toward decreasing potentials, following curve III; O_2 reduction not observed. Fig. 2B, pH 10.8 (IM $NaCIO_4$): curves I and IV, cathodic processes including hydrogen evolution; curve II, anodic dissolution and passivation, toward increasing potentials, stirring-dependent current in the active/passive transition region; curve III, reverse of curve II; no O_2 reduction observed.

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Fig. 20-3. Polarization curves for Si. Fig. 3A, pH 0.04 (1M $IICIO_4$): Fig. 3B, at pH 10.8 (1M $NaCiO_4$): passivation in alkaline solution: no O_2 reduction observed.



Fig. 20-4. Polarization curves for NiAs. Fig. 4A, pil, 0.04 (1M HClO_4): curve I, cathodic hydrogen evolution: curve II, anotic current for dissolution and passivation of NiAs; curve $1 \cdot O_{2^2}$ cathodic O_2 reduction, Tafei region masked by anodic dissolution of NiAs. Fig. 4B, pH 10.8 (1M NaClO₄): curve 1, cathodic hydrogen evolution; curve II, anodic dissolution of NiAs, toward increasing potentials; curve III, same as II, toward decreasing potentials; curve $1 \cdot O_2$, O_2 reduction.

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Fig. 20-6. Polarization curves for NiS. Fig. 6A, pH 0. 04 (1M HCHO₄): curve I, cathodic hydrogen evolution: curve IV, oxide reduction and hydrogen evolution: curve II, (toward increasing potentials) anodic dissolution and passivation current: curve III, (toward decreasing potentials): curve I-O₂, reduction currents for O₂-reduction. Fig. 6B, pH 10.8 (1M NaClO₄): curves I, II, III, and IV have same meanings as in Fig. 6A.

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Fig. 20-7. Polarization curves for NiSi. Fig. 7A, pH 0.04 (1M HClO₄): curves I and IV, cathodic hydrogen evolution: curves II and III, anodic currents for dissolution and passivation of NiSi: curve II, (traced toward increasing potentials) shows active/passive transition; curve I-O₂, O₂ reduction, masked by anodic dissolution of electrode. Fig. 7B, pH 10.8 (1M NaClO₄): curves I, II, III, and IV have same meanings as in Fig. 7A; curve II shows two transitions: curve I-O₂ shows pronounced O₂ reduction current.

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TABLE 20-1

Summary	of	Results
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Molar	Perchlorate	e Solutions	of pH 0.04

	Passivation potential <u> </u>	Critical C.D. i _t , amp/cm ²	C.D. in passive region i _p , amp/cm ²	Passive potential range $\Delta E_{p}, v_{$	$\Delta E_t / \Delta pH$
Ni	+0.3	3×10^{-1}	$\sim 10^{-6}$	0.4 to 1.4	-0.074
Sb	+0.4	4.5×10^{-3}	3.5×10^{-3}	0.4 to 0.9	-0.041
Si			$\sim 10^{-3}$		
NiAs	+0.2	2.3 x 10 ⁻⁵	$\sim 10^{-5}$		
NiSb	+0.45	7×10^{-3}	7×10^{-4}	0.8 to 1.8	-0.082
NiS	-0.10	1.5×10^{-4}	1.5×10^{-4}	-0.1 to 0.7	-0.026
NiSi	÷0. 10	4.5×10^{-3}	$\sim 10^{-3}$	0.4 to 1.8	-0.050

<u>1M NaClO₄ at pH 1...8</u>

	Passivation potential E _t , <u>v</u>	Critical C. D. $-i_t$, $\frac{amp/cm^2}{2}$	C. D. in passive regin _i _p , amp/cm ²	Passive poten- tial range $\Delta E_{p}, v_{}$
Ni	-0.5	10 ⁻⁶	7 x 10 ⁻⁷	
Sb	-0.04	5×10^{-4} (unstirred) 5 x 10^{-3} (stirred)	d) $\sim 10^{-4}$	
Si	-0.16	1.7×10^{-6}	1.6 x 10 ⁻⁶	-0.16 - 2.0
NIAS		<10 ⁻⁶	<10 ⁻⁶	
NiSb	-0.44	3×10^{-4}	3×10^{-4}	-0.44 - 0.37
	(0. 56)	$(2 \times 10^{-3})^{-3}$	(2 x 10 ⁻³)	(0.56 - 1.0)
NiS	-0.39	6×10^{-4}	3×10^{-6}	-0.39 - 0.56
NiSi	-0.44	1.7×10^{-4}	1.2×10^{-4}	

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electrode surface, a condition not fulfilled during the dissolution of any of the compounds. In general, the oxide of the companion element or a mixed oxide is more stable than nickel oxide in acid solution at low potentials. This situation is reversed in alkaline solution, and in this case the initial passive film is nickel oxide.

The order of electrocatalytic activity for oxygen reduction is Ni > NiAs > NiSi in acid solution, and NiS > Ni > NiSi ~ NiSb > NiAs in alkaline solution. The order in the activity of the materials remains essentially independent of pH except that NiSi is more active than NiAs in alkali and that the reverse is true in acid. A significant observation is that O_2 reduction takes place in the potential region in which the materials are generally passive and are covered with surface oxides. The observed differences between these compounds indicate that the presence of the negative element (Si, S, Sb, As) modifies the nature of the surface oxide to produce changes in the electrocatalytic activity.

21. Kinetics of Redox Reactions on Passive Electrodes

A.C. Makrides

Submitted as Technical Memorandum No. 2 (April. 1963). Presented in part at the Washington Meeting of the Electrochemical Society, October, 1964. Published: J. Electrochem. Soc. 111, 392-400 (1964).

The formation of a surface oxide is an important factor in the kinetic. of oxygen evolution and oxygen reduction on practically all metallic electrodes. In general, a superficial oxide is expected to change the specific interaction of oxygen (or of reaction intermediates) with the electrode to alter the potential distribution between metal and electrolyte, and to modify the dissolution kinetics of the metal. The change in the kinetics of dissolution has been studied in detail, since it is the characteristic phenomenon in passivity. However, relatively little work has been done on the effects of oxide films on electrochemical reactions other than corrosion reactions. The general features of electrode processes on superficially oxidized metals are examined in this paper using as a model a simple, one-electron redox reaction.

Aside from specific effects, e.g., changes in the energy of adsorption of reactants or of intermediates, superficial oxides may influence the kinetics of an electrochemical reaction by providing a barrier to electron transfer between the metal and an ion or molecule at the oxide/electrolyte interface. With thick films, the oxide is the electrode, the underlying metal serving merely as a contact. The chemical character of the bulk oxide and its electrical properties, e.g., its semiconducting properties, enter then into the description of its electrode characteristics. At the other extreme, very thin oxide films may be considered as simply providing a barrier through which electrons must tunnel in order to participate in the reaction. The upper limit for tunneling is about 30 Å; conduction through the oxide is necessary for electron transfer with thicker films. However, if the film is only 100-200Å, its electrical characteristics are influenced strongly by the underlying

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metal, and this interaction may be important in the description of the electrode behavior of superficially oxidized electrodes.

A simple way of examining the general electrochemical effects of surface oxides is to compare the kinetics of an elementary electron transfer reaction on oxide-covered and oxide-free electrodes. Specific effects, e. g., adsorption, should be at a minimum for the redox couple chosen. The Fe⁺⁺⁺/Fe⁺⁺ couple was selected in this study because its kinetic behavior on oxide-free electrodes is straightforward. Furthermore, since its reversible potential falls in about the middle of the potential region over which a number of common metals and alloys are passive, both the oxidation of ferrous ion and the reduction of ferric ion can be followed without gross interference from corrosion reactions.

Overpotential measurements for the Fe⁺⁺⁺/Fe⁺⁺ couple on passive Ni, Fe, and Ti, were carried out in solutions of fixed ionic strength but of varying pH. The results are given in Figs. 21-1 - 21-4and Table 21-I.

Tafel lines were generally obtained with exchange currents between 10^{-7} and 10^{-5} amp/cm² and cathodic transfer coefficients of about 0. 45. The anodic transfer coefficients were less, particularly with Ti and Fe electrodes. A limiting anodic current, which was unrelated to diffusion of Fe⁺⁺ ion in solution, was observed under certain conditions on Fe and Ti. The contribution of ionic current to the total current through the film is negligible in most cases. In general, the passive film has rectifying properties, the easy direction of electron flow being from metal to solution. This rectification is additional to the usual Faradaic rectification observed with most electrochemical reactions. The apparent transfer coefficients, calculated from the anodic and cat' odic polarization curves, yield sums significantly less than unity.

The polarization behavior of the Fe^{+++}/Fe^{++} couple on passive Ni, and Ti can be summarized by an expression of the form (potentials in units of RT/F)

$$i = i_{O} \{ \exp[-\alpha_{P}(E - Erev)] - \exp[\beta_{P}(E - Erev - V)] \}$$

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Fig. 21-1. Polarization curves for the Fe^{+++}/Fe^{++} couple (0.05M) on Ni electrodes in M MgSO₄ at pH = 2.3 and at 30^oC. Current densities corresponding to the full circles are calculated from $i_{ox} = i_{appl} + i_{red}$ where i_{red} is given by the extrapolated cathodic curve.



Fig. 21-2. Polarization curves for the $Fe^{+1} + /Fe^{++}$ couple (0.05M) on Ti electrodes in M MgSO₄ at pH = 2.3 and at 20^oC. Current densities corresponding to the full circles are calculated from $i_{0x} = i_{appl} + i_{red}$ where i_{red} is given by the extrapolated cathodic curve.



Fig. 21-3. Polarization curves for the Fe^{+++}/Fz^{++} couple on Fe electrodes at 40° C. The solution was 0.30M in Fe⁺⁺⁺, 0.03M in Fe⁺⁺, and M/20 in H₂SO₄ and M/10 in Na₂SO₄. The anodic curve was obtained after the electrode was kept for 30 min at 0.20v vs. Erev for Fe⁺⁺⁺/Fe⁺⁺.



Fig. 21-4. Anodic polarization curves for the Fe^{+++}/Fe^{++} couple on iron in M MgSO₄ at pH = 1, 45. The Fe^{+++} concentration was 0.3M. The current was kept constant at each value shown until the potential changed by less than 1 mv in 5 r. n.

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				TABLE 21-	1			
				Nickel ^a				
			i_ extravolated	$i_{o} \operatorname{from}\left(\frac{\partial i}{\partial E}\right)$				
pH	Т, °С	Ecorr	from cathodic 2 curve, amp/cm ²	(i _{o)}	$(i_{\rm c})_{\rm a}$	QC	aa	
2.3	10	-0.0003	2.6 x 10 ⁻⁵	2.5 x 10 ⁻⁵	2.5 x 10 ⁻⁵	0.415	0. 484	
	20	-0.0006	4.55 x 10 ⁻⁵	4.4 x 10 ⁻⁵	4.4 x 10 ⁻⁵	0.40	0.495	
	20	-0.0012	7.6 x 10 ⁻⁵	7.45 x 10 ⁻⁵	7. 15 x 10 ⁻⁵	0.41	0. 49	
0.35	30	-0.0006	8.6 x 10 ⁻⁵	8.ύ x 10 ⁻⁵	7.5 x 10 ⁻⁵	0. 46	(0. 60) (0. 375)	
0.00	0	-0.0003	5.0 x 10 ⁻⁵	-		0.47	1 8	
	10	-0.0008	8.0 x 10 ⁻⁵	1	1	0.47	1	
	20	-0.0005	11. x 10 ⁻⁵	10. 2 x 10^{-5}	8.0 x 10 ⁻⁵	0.47	!	
	30	-0.001	14. x 10 ⁻⁵	12.8 x 10 ⁻⁵	7.0 x 10 ⁻⁵	0.47	(0.33)	
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^a All solutions $C_{Fe}^{+++} = C_{Fe}^{++-+} = 0.050M$ and M MgSO₄.

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Here V is the potential drop within the oxide and is, in general, a function of the current density. The results suggest that V is proportional to E, i.e.

$$V = (E - Erev)$$

In this case

$$i = i_0 \{ \exp[-\alpha_r (E - Erev)] - \exp[(\beta_r) (1 - \gamma) (E - Erev)] \}$$

Consequently

$$\alpha_{c} + \alpha_{a} = (\alpha_{r}) + \beta_{r} (1 - \gamma) = 1 - \gamma \beta_{r}$$

with an apparent stoichiometric number of

$$v = \frac{+}{i} i_0 / \left(\frac{\partial i}{\partial E}\right) (E - Erev) = 0 = 1 - \gamma \beta_r$$

These equations are followed by passive Ni at pH = 2.3 and by Ti at pH = 1.5 and for $T \le 20^{\circ}$ C. The coefficient _Y is 0.15 for Ni and 0.60 for Ti. Fe, Ni at low pH, and Ti at higher temperatures show a more complicated behavior on anodic polarization.

These results demonstrate the existence of a potential drop across the passive, oxide film, the magnitude of this potential being dependent on the thickness and composition of the film.

22. Kinetics of the Ferric-Ferrous Reaction on

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A.C. Makrides

Submitted as Technical Memorandum No. 4 (June, 1963).

Presented in part at the Washington Meeting of the Electrochemical Society, October, 1964. Published: J. Electrochem. Soc. 111, 400-407 (1964).

The kinetics of the Fe⁺⁺⁺/Fe⁺⁺ couple on superficially oxidized electrodes present a number of interesting features. In general, the evchange currents are low (~ 10^{-6} amp/cm²), the transfer coefficients have unusual values and sums less than unity, and the anodic and cathodic curves are generally asymmetrical. These features appear to be basic to an electrode covered with a thin oxide (i. e. , a passive electrode) and not to be related to the specific nature of the redox couple. They are probably connected with electronic processes within the surface oxide and are expected, therefore, to occur with any redox reaction taking place at such an electrode. In particular, they undoubtedly play a major role in oxygen evolution and reduction on electrodes which form surface oxides in the relevant potential region.

A study of the kinetics of the Fe⁺⁺⁺/Fe⁺⁺ reaction on passive Fe, Ni, and Ti electrodes was presented in a previous paper (1). This redox couple was chosen as a model reaction because its kinetics are straight-forward on oxide-free electrodes and because both oxidation of Fe⁺⁺ and reduction of Fe⁺⁺⁺ can be followed over a relatively wide potential range.

A detailed study of the redox kinetics of the Fe^{+++}/Fe^{++} reaction on Fe-Cr alloys is presented here. These alloy electrodes, which show the same general features as other superficially oxidized electrodes, are convenient experimentally, since they have small dissolution rates in acid solutions. In addition to the usual experimental quantities required for a more or less complete characterization of a redox reaction (Tafel parameters, reaction order, pH dependence, and temperature coefficient), the



Fig. 22-1. Polarization curves of the re^{i++}/Fe^{i++} couple (0.050M) on passive Fe-Cr in M MgSO₄ at pH = 2.2. Solid circles are calculated from $i_a = i_c \pm i_{appl}$ where i_c is the extrapolated cuthodic current and l_{appl} the externally applied current (30^oC).



Fig. 22-2. Current at a fixed potential as a function of the concentration of ferrous or ferric ions $__{--}$ oxidation of Fe⁺⁺ at 0. 600v vs. SCE; _____ reduction of Fe⁺⁺⁺ at 0. 200v vs. SCE (30^oC).

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Fig. 22-3. Temperature coefficient of exchange cu_1 rent (extrapolated from the cathodic curve) for $\text{Fe}^{+++}/\text{Fe}^{++}$ (0.50M) in M MgSO₄.



Fig. 22-4. Polarization curves of the Fe^{+++}/Fe^{++} couple (0.10M) on passive Fe-Cr in M MgSO₄ at different pH (30^oC).

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Fig. 22-5. Anodic polarization curve for $Fe^{\tau/t}$ exidation on passive Pe-Cr alloy after cathodic pulsing. Measurements were made 24 hr after pulsing.



Fig. 22-6. Effect of anodic (dark circle) and cathodic (open circle) pulsing on polarization curve of the Fe⁺⁺⁺/Fe⁺⁺ couple (0, 10M) on passive Fe-Cr in M MgSO₄ at pH = 1.35. Curve corresponding to (dash half da \measuredangle circle dash) was obtained with electrodes which had not been pulsed (30°C). Arrows show points at which potential begins to drift in indicated direction. Measurements in these potential ranges were made rapidly and do not correspond to a steady state.

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TABLE 22-1

	Fe ⁺⁺⁺ /Fe ⁺⁺ or 	n Passive Fe-0 30.0°C	Cr	
	T = 30	0.0°C		
Fe ⁺⁺⁺ , M/''.er	Fe ⁺⁺ , M/liter	pH ^a	αC	<u>aa</u>
5.0×10^{-3}	5.0 x 10 ⁻³	0.30	0. 40	0.37
5.0×10^{-2}	5.0×10^{-2}	0.30	0.38	0.37
1.0×10^{-1}	1.0×10^{-1}	1.35	0. 43	0.36
$5 0 \times 10^{-2}$	5.0×10^{-2}	1.5	0.40	0.35
1.0×10^{-1}	1.0×10^{-1}	1.7	0.41	0.35
1.0×10^{-1}	1.0×10^{-1}	2.1	0.38	0.30
1.4×10^{-2}	5.0×10^{-3}	0.60 ^b	0.42	0.38
4.1 x 10^{-2}	5.0×10^{-3}	0. 60 ^b	0.40	0.38
5.0×10^{-2}	5.0×10^{-2}	0.30	0. 40	$\{ \begin{array}{c} 0.58^{\rm C} \\ 0.37 \end{array} \}$
5.0×10^{-2}	5.0×10^{-2}	1.35	0. 42	0.58 ^c

Avg $\alpha c = 0.40 \pm 0.01$ $\alpha a = 0.36 \pm 0.02$

 a All solutions M in $\ensuremath{\text{MgSO}}_4$ except where noted.

^b No MgSO₄ added.

^c Two Tafel slopes on anodic polarization; data not included in average αc or αa .

dependence of the kinetics on the electrochemical conditioning of the surface film (e.g., "deforming") was also determined. The results are presented in Figs. 22-1 - 22-6 and Table 22-I.

Tafel behavior is generally observed with exchange currents of the order of 10^{-6} amp/cm^2 (C_{Fe}+++ = C_{Fe}++ = 0.050M), a cathodic transfer coefficient of 0.40, and an anodic transfer coefficient of 0.36. The electrochemical reaction order is unity both for oxidation of Fe⁺⁺ and reduction of Fe⁺⁺⁺; the heat of activation at the reversible potential is 10.8 kcal/mole.

The results show that an appreciable fraction of the total potential drop on anodic polarization occurs within the passive film. This potential drop is probably associated with changes in the average oxidation state of cations between the inner and outer layers of the surface oxide. The electrochemical characteristics of the film depend on the pH of the solution, probably because of the migration of protons into the passive film. A change in the kinetics of Fe⁺⁺ oxidation which occurs in the transpassive region of the Fe-Cr alloy is due to the production of Fe⁺³ or Cr⁺⁶ ions in the inner oxide layer.

23. Surface Oxidation of Gold Electrodes

S. B. Brummer and A. C. Makrides

Submitted as Technical Memorandum No. 6 (October, 1963). Presented at the New York Meeting of the Electrochemical Society, October, 1963. Published: J. Electrochem. Soc. 111, 1122-1128 (1964).

Most metals form stable surface oxides in the potential range where oxygen can be reduced. Therefore, a study of the electrochemical properties of surface oxides is a necessary step in the detailed study of the mechanism of oxygen reduction.

The surface oxides formed on electropolished gold by potentiostatic anodization in the range 1.2 - 1.85 v vs. H^+/H_2 were studied by galvanostatic reduction at current densities between 10 and 1000 $\mu a/cm^2$. Molar perchlorate solutions of pH 0.06 to 2.8 were employed.

Typical cathodic chronopotentiograms, taken after 5 minutes of constant potential anodization, are shown in Fig. 23-1. A notable feature of these curves is that the oxide is reduced at virtually a constant potential, although its "thickness" is only 1-2 monolayers.

The amount of oxide on the electrode increases with potential above 1.3 volts (Fig. 23-2), but is independent of pH at a fixed potential against the reversible hydrogen electrode (RHE) as can be seen from Fig. 23-3.

Current-potential curves constructed from the chronopotentiograms determined at various current densities follow Tafel curves with a slope of 39 to 42 mv/decade (Fig. 23-4). The Tafel slope is independent of the potential of formation of the oxide E_a , and of pH, but the exchange current of the reduction depends both on E_a and on pH. The order of the reduction reaction with respect to hydrogen ions is 1.4.

The following mechanism was proposed to account for the main experimental findings:


Fig. 23-1. Typical cathodic chi onopotentiogram in 1N HClO₄ after 5 min of ano.iization at controlled potential. Arrew in upper figure indicates ~ 0.8 monolayer of O^{2-} remaining on surface. Arrow in lower figure indicates the potential corresponding to the reduction of the oxide.



Fig. 23-2. Variation of cathodic charge with the potential of anodization in $1N \text{ HClO}_4$. Open and closed circles are from different experiments in the present study. Crosses are from the results of Laitinen and Chao (J. Electrochem Soc. <u>108</u>, 726(1961)). The arrow indicates a monolayer.

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Fig. 23-4. Current-potential curves for reduction of films formed at various potentials.

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AuOOH + H⁺ + e
$$\stackrel{\text{fast}}{\rightleftharpoons}$$
 AuO + H₂O (10)

$$AuO + H^+ + e \xrightarrow{slow} AuOH$$
 (11)

AuOH - H⁺ +
$$e \frac{fast}{r}$$
 Au + H₂O (12)

where AuO is written as the Au^{II} oxidation state and AuOH as the Au^{I} oxidation state, but either of these may be hydrated. The rate equation is

$$i = 3k_{II} (AuO) (H^+) \exp -\frac{\alpha \psi F}{RT}$$
(13)

where α , the transfer coefficient, is ~0.5, and ψ is the potential. We can solve for (AuO) by assuming that reaction (I) is fast and is in equilibrium at all potentials. Then

$$k_{I}^{A}(AuOOH) (H^{+}) \exp - \frac{\beta \psi F}{RT} = k_{I}^{A}(AuO) (H_{2}O) \exp \frac{(1 - \beta) \psi F}{RT}$$
(14)

and

(AuO) = K (AuOOH) (H+) exp -
$$\frac{\psi F}{RT}$$
 (15)

where K is a constant equal to k_I/k_I (H₂O). Then the final rate equation is

i =
$$3k_{II}K$$
 (AuOOH) (H⁺)² exp - $\frac{(1+\alpha)\psi F}{RT}$ (16)

If the surface oxide AuOOH is thought of as a well-defined and essentially homogeneous solid, its activity will remain constant although its concentration declines during reduction. This would account for the initial flatness of the galvanostatic reduction curve. If α , in Eq. 16, is 0.5, the Tafel slope is 39 mv, in good agreement with experiment.

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The pH dependence, $(\Im \log i / \Im pH)_E$, as given by Eq. 16 is -2. This result is only in moderate agreement with the experimental observation of $(\Im \log i / \Im pH)_E = -1.39$.

Further experiments with this system have showed that certain cations in solution, viz. Mg^{+2} , considerably alter the kinetics of reduction of the oxide although the amount of oxide at any given potential is essentially the same. These effects probably result from specific adsorption of cations which modifies the double layer structure at the oxide-solution interface and may be crucial in determining the electrocatalytic properties of superficial oxides present during the cathodic reduction of O₂.

24. Ageing Effects on Thin Anodic Oxide Films on Gold in N HC10₄

S. B. Brummer

Submitted as Technical Memorandum No. 20 (June, 1965). Presented at the Toronto Meeting of the Electrochemical Society, May, 1964.

Published: J. Electrochem. Soc. <u>112</u>, 633-636 (1965).

The properties of anodic oxide films on Au in 1N HC10₄ were studied as a function of time (2 sec to 5 min) and of potential of formation (1. 45 to 1. 85 v vs. Pt, H^+/H_2 in the same solution).

Over the range observed (10-20% of the total oxide) the oxide grows slowly with time, apparently according to Elovich kinetics (Fig. 24-1) although the constants are potential-dependent.

The oxides formed at longer times at a given potential are harder to reduce (Fig. 24-2). This "ageing" Effect is greatest at the lower potentials of formation despite the greater thickness of the oxide at the higher potentials and despite the larger change in the amount of oxide at the higher potentials.

The triangular sweep method show is not be used in the quantitative investigation of these thin oxides, partly because the oxide is formed over a range of potentials, but mainly because of this ageing effect. Thus, oxides formed during a triangular sweep are easier to reduce than expected from the potential of formation, because of lack of ageing. They are also less homogeneous (reduced over a wider range of potentials) as shown in Fig. 24-3.



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Fig. 24-1. Growth of oxide at various potentials (mv vs. R. H. E.) as a function of time.



Fig. 24-2. Current-potential curves for the reduction of oxides, as a function of time and potential of formation.

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Fig. 24-3. Current-potential curves for the reduction of oxides, made with linear potential sweeps of 0.14 v/sec. Curve A, oxide reduced after 5 min at 1500 mv; curve B, oxide formed and reduced during triangular potential sweep, reversed at 1500 mv; curve C, triangular sweep reversed at ~ 1600 mv, such that $Q_A = Q_{C'}$.

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¹³ ABSTRACT This report is an extended abstra dealt with electrochemistry of fuel ce hydrogen reaction, the oxidation of or clarification of the role of the electro the central theme of this work. In order to establish the relation composition it is necessary first to un reaction of interest. The simplest re and in this case, it was possible to sh dependent on the electronic structure electrode. The oxidation of organic fi plicated reactions and the work with the towards elucidating their kinetics and A summary of the main experime and of the implications of these studies given in Section II of this report. Det this contract, and published in various	ict of the main ll reactions, and ganic fuels, and de in catalyzing of electrocatal inderstand in de eaction studied low that electroc rather than the uels and oxyget hese reactions mechanisms. ental results, o in the develo tailed abstracts s journals, are	results nd included ind oxyg g fuel of ytic active tail the was the catalyne was di was di of the co pment s of the e given	s on a program which uded studies on the gen reduction. The cell reactions was tivity to electrode e mechanism of the e hydrogen reaction, tic activity is ce composition of the ction are more com- rected primarily onclusions reached, of fuel cells is work done under in Section III.						

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