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PART I—APPLICATION OF THE SUBSONIC KERNEL FUNCTION TO NONPLANAR LIFTING SURFACES

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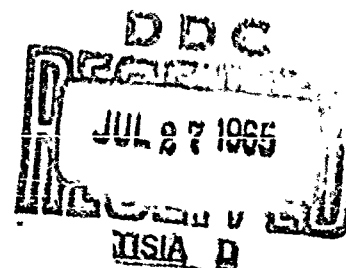
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The Space and Information Systems Division
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FOREWORD

This report covers the research conducted by the Space and Information Systems Division of North American Aviation, Inc., Downey, California, for the Aerospace Dynamics Branch, Vehicle Dynamics Division, Air Force Flight Dynamics Laboratory, Wright-Patterson Air Force Base, Ohio, under Contract No. AF 33(657)-10399.

The work was performed to advance the state of the art of flutter prevention for flight vehicles as part of the Research and Technology Division, Air Force Systems Command's exploratory development program. This research was conducted under Project No. 1370 "Dynamic Problems in Flight Vehicles," and Task No. 137003, "Prediction and Prevention of Dynamic Aerothermoelastic Instabilities." Mr. James Olsen of the Vehicle Dynamics Division, Air Force Flight Dynamics Laboratory was the Project Engineer.

Mr. L. V. Andrew was the Program Manager for North American Aviation. Mr. H. T. Vivian, under the guidance of Dr. H. Ashley, laid out the form of the solution and wrote the computer program. Dr. E. R. Rodemich made many significant contributions to the numerical analysis.

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This technical documentary report has been reviewed and is approved.

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ABSTRACT

In this part, equations for pressure distributions and generalized aerodynamic forces are derived for a thin nonplanar lifting surface in simple harmonic motion at subsonic speeds. A digital computer program, written in Fortran IV, is also presented herein. The computer program will generate up to a ten by ten matrix of generalized aerodynamic forces when given data for the geometry of a planar lifting surface with a folded planar tip, the flight Mach number, the reduced frequency of motion, and some control constants. Control surface deflections are not accounted for in this study.

The kernel function method given by Watkins, Runyan, and Woolston (Reference 1), which relates the pressure distribution to the downwash on a planar lifting surface, has been extended and applied to a nonplanar lifting surface. Hsu's technique (Reference 4) of employing Gaussian quadrature formulas is used when integrating the product of the kernel function and the lift function over the planform area.

Recommendations are made to extend the method to account for blunted leading edges and the accompanying airfoil thickness and to account for control surface deflections.

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SYMBOLS

a	Local speed of sound (constant in linear theory)
b_0	Root semichord
$b(s)$	Local semichord
C_p	Pressure coefficient
\bar{C}_p	Time independent factor of oscillatory part of C_p
i, j, k	Unit vectors parallel to the coordinate axes
i	$\sqrt{-1}$
k	Reduced frequency, $\omega b_0 / U_\infty$
k_1	kr_1
K	Kernel function
K_0, K_1, K_2	Modified Bessel functions of the second kind
M	Local Mach number (same as free stream Mach number, M_∞ , in linear theory)
n	Coordinate measured normal to wing surface
$n(x, s, t)$	Contribution to n of elastic deformation of the wing
$n_t(x, s)$	Contribution to n of wing thickness
$\bar{n}(x, s)$	Time independent factor of $n(x, s, t)$
p	Pressure
\vec{q}	Fluid velocity
R	Gas constant
R	$\sqrt{(x-\xi)^2 + \beta^2 r_1^2}$

List of Symbols continued on next page.

r_1	$\sqrt{(y-\eta)^2 + (z-\zeta)^2}$
s	Curvilinear coordinate on the wing (page 10)
T	Absolute temperature
t	Time
u, l	Subscripts indicating upper and lower wing surfaces
u, v, w	Components of perturbation velocity
U, W	x- and z-components of \vec{V}
U_∞	$ \vec{V} $ (speed at infinity)
\vec{V}	Uniform fluid velocity at infinity
\bar{W}	Time independent factor of normal velocity
x, y, z	Cartesian coordinates
x_0	$x - \xi$
y_0	$y - \eta$
β	$\sqrt{1-M^2}$
γ	Ratio of specific heats
$\gamma(s)$	Local angle between wing surface and xy-plane
$\Delta \bar{p}$	Time independent factor of pressure difference between wing surfaces, $\bar{p}_l - \bar{p}_u$
ξ, η, ζ	Cartesian coordinates
$\bar{\xi}$	Coordinate on the wing (page 22)
ϕ	Velocity potential
φ	Perturbation velocity potential
$\bar{\varphi}$	Time independent factor of φ
$\bar{\psi}$	Acceleration potential
ρ	Fluid density
ω	Angular frequency (radians per unit time)

I. INTRODUCTION

The first published numerical method for solving the subsonic pressure distribution problem for planar lifting surfaces undergoing simple harmonic motion was developed at NASA's Langley Research Center by Watkins, Runyan, and Woolston (Reference 1). Watkins, et al., presented two methods of handling the numerical integration of the kernel function in the region where high-order singularities exist. Both methods involved a dense concentration of integration points in the neighborhood of the singularity. Using these methods, it is possible to obtain downwash integrals, in terms of the pressure-loading coefficients, at any arbitrary set of points on the surface (e. g., at all the kinematic downwash points known from previously determined vibration mode data). However, in order to reduce the running time on the computer, the downwash integrals were obtained at a selected set of collocation points, such as those at intersections of quarter, half, and three-quarter chord stations and like half-span stations. When downwashes were matched exactly (and thus, boundary conditions) at these collocation points, responsibility was placed upon the user to evaluate the kinematic downwashes there. A least-square error surface fitted to the mode data was commonly used to evaluate them. Furthermore, if the user desired that the boundary conditions be satisfied at a greater number of points, it was necessary that he use a correspondingly greater number of loading functions.

Procedures were then described by Rodden and Revell (Reference 2) and the correct form of the equations were presented by Fromme (Reference 3) for calculating pressure-loading coefficients which match a greater number of kinematic downwashes than coefficients, in the sense that the sum of squares of amplitudes of differences of complex numbers are minimized. Since it was still the responsibility of the user to evaluate the kinematic downwashes at the collocation point, least-square error procedures were used twice: once implicitly and once explicitly.

Hsu (Reference 4) significantly advanced the logical development of the kernel function approach when he established an optimum set of collocation and integration points. He started with the previously established chordwise pressure functions based on steady-state, two-dimensional, incompressible aerodynamics, and with spanwise loading functions, based on steady-state lifting-line theory. He concluded that there is sufficient reason to believe that these functions display the proper characteristics near the edges of lifting surfaces oscillating in a compressible fluid.

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Returning to the two-dimensional case, Hsu established that if the chordwise distribution of modal deflections (and thus downwashes) is accurately represented by a polynomial of degree $2N-1$ and is approximated by a polynomial of degree $N-1$, then the integral for the sectional load is evaluated with zero error by a N -point Gaussian quadrature if the difference between the accurate and the approximate representation of the downwashes is made equal to zero at each of the N points (i. e., the chordwise collocation stations). Conversely, still for the two-dimensional case, Hsu established that if the product of the pressure function and the kernel, divided by the Jacobi-Gauss weight factor (which produces the square root singularity at the leading edge), is accurately represented by a polynomial of degree $2N-1$, then the integral for the downwash at any one of the collocation stations is evaluated with zero error by a N -point Gaussian quadrature. These N points are then made the chordwise integration stations.

For the spanwise direction, using lifting-line theory, Hsu similarly established M -spanwise collocation stations and $M + 1$ interdigitated spanwise integration stations plus the conditions under which the Gaussian quadrature can be used with zero error.

It is important to note that the kernel of the integral equation for the downwashes in unsteady, three-dimensional, compressible flow cannot be accurately represented by a polynomial of finite degree. It is equally important to note, however, that, because of the edge characteristics of the pressure and loading functions, the Gaussian quadratures employed at Hsu's optimum point set evaluate the integrals with the least squared error for a given number of integration points. We have yet to match the boundary conditions using Hsu's method.

The downwash matching problem in Hsu's approach is basically the same as in Watkin's approach; we merely have a more logical choice of points at which to match them. In the examples Hsu used to demonstrate his approach, he chose to use the same number of pressure-loading functions as collocation points. However, the approach is not dependent upon that choice. If a smaller number of pressure-loading functions are used, then the procedures described by Rodden, Revell, and Fromme may be used to compute pressure-loading coefficients which yield a minimum sum of squares of amplitudes of differences in downwashes.

A need has arisen for application of the kernel function method to non-planar lifting surfaces on future aerospace vehicles. Application is also required to more conventional non-planar surfaces such as T-tail, V-tail, and wing-vertical tail combinations.

Professor H. Ashley outlined the application to the folded tip configuration. A computer program based on Ashley's work was developed for steady-state flow by L. Johnson, et al., of the Los Angeles Division of North American Aviation, Inc.

The work reported herein is based on Professor Ashley's outline. However, the expression for the kernel has been greatly simplified by Dr. E. R. Rodemich of North American Aviation, Inc., Space and Information Systems Division.

II. FUNDAMENTAL EQUATIONS OF FLUID MOTION

Consider a body immersed in a compressible, nonviscous, perfect fluid and assume the fluid flow to be isentropic and irrotational. Under these conditions, a velocity potential ϕ , exists:

$$\vec{q} = \nabla \phi \quad (1)$$

where \vec{q} is the velocity vector of a fluid element and ∇ is the gradient operator (See Reference 4.) Also under these conditions, the isentropic (constant entropy) pressure-density relationship is valid. Thus,

$$a^2 = \left(\frac{\partial p}{\partial \rho} \right)_s = \gamma R T \quad (2)$$

Other equations which govern the flow are the continuity equation for conservation of mass

$$\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \vec{q}) = 0 \quad (3)$$

and Euler's equations for conservation of momentum

$$\frac{D\vec{q}}{Dt} = -\frac{1}{\rho} \nabla p \quad (4)$$

where, $\frac{D}{Dt}$, the substantial derivative, is

$$\frac{D}{Dt} = \frac{\partial}{\partial t} + \vec{q} \cdot \nabla \quad (5)$$

These equations may be combined, as described in Reference 5, to yield the nonlinear, unsteady flow equation

$$\nabla^2 \phi - \frac{1}{a^2} \left[\frac{\partial^2 \phi}{\partial t^2} + \frac{\partial q^2}{\partial t} + (\vec{q} \cdot \nabla) \frac{q^2}{2} \right] = 0 \quad (6)$$

Consider, then, that the fluid motion consists of a perturbation superimposed on a uniform stream velocity $\vec{V} = U\vec{i} + W\vec{k}$ parallel to the xz -plane of a rectangular Cartesian coordinate system. Then the velocity potential may be expressed as the sum of a uniform part and a perturbation part

$$\phi = Ux + Wz + \varphi \quad (7)$$

and, similarly, the velocity vector becomes

$$\vec{q} = \vec{V} + \nabla\varphi = \nabla\phi \quad (8)$$

The pressure coefficient at any point in an isentropic flow field is

$$C_p = \frac{p - p_\infty}{\frac{1}{2} \rho_\infty U_\infty^2} \quad (9)$$

where $p - p_\infty$ is the difference between local pressure and free-stream pressure, $U_\infty = |\vec{V}|$, and $1/2 \rho_\infty U_\infty^2$ is the free-stream dynamic pressure. From Kelvin's equation (Reference 5) for isentropic flow

$$C_p = \frac{2}{\gamma M_\infty^2} \left\{ \left[1 + \frac{\gamma - 1}{2} M_\infty^2 \left(1 - \frac{\vec{q} \cdot \vec{q} + 2 \frac{\partial \varphi}{\partial t}}{U_\infty^2} \right) \right]^{\frac{\gamma}{\gamma - 1}} - 1 \right\} \quad (10)$$

A complete statement of the fundamental problem requires specification of the boundary conditions. The boundary conditions at infinity depend upon the free-stream velocity. When it is less than the speed of sound in the fluid, the disturbances to the flow die out and are not felt at infinity. When it is greater than the sonic speed, then in the region where disturbances are felt, even at infinity, the component of flow due to the disturbance is directed away from the source of disturbance and otherwise the free-stream flow is undisturbed. The boundary conditions at the surface of the body require that the flow be tangent to the surface everywhere on the body. This condition is satisfied by the equation

$$\frac{D}{Dt} \psi(x, y, z, t) = 0 \quad (11)$$

where

$$B(x, y, z, t) = 0 \quad (12)$$

is the equation for the position of the surface at any time t , and the substantial derivative D/Dt is defined by Equation 5.

III. LINEARIZED EQUATIONS OF MOTION

Linearization of the equations of motion is not dependent upon an explicit form of the body equation, Equation 12, so long as the normal derivatives of the equation are everywhere nearly perpendicular to the free-stream direction. Thin lifting surfaces at small angle of attack satisfy this condition and are treated herein and in Parts 2 and 4 of this report. The special considerations required for thick bodies and high angles of attack are treated in Parts 3 and 5. The following development is, therefore, restricted to thin airfoils.

We first obtain the specialized form of Equation 7 when the uniform stream velocity lies along the x-axis; i. e., $W = 0$ and, therefore, $\vec{V} = U\vec{i}$, $U_\infty = U$. The velocity potential is

$$\phi = Ux + \varphi \quad (13)$$

and the velocity vector of a fluid element becomes

$$\vec{q} = (U + u)\vec{i} + v\vec{j} + w\vec{k} \quad (14)$$

where

$$u = \frac{\partial \varphi}{\partial x}, \quad v = \frac{\partial \varphi}{\partial y}, \quad \text{and} \quad w = \frac{\partial \varphi}{\partial z}$$

The perturbation velocities u , v , and w are assumed to be much smaller than the free-stream velocity; i. e., $u, v, w \ll U$.

The linearization procedure when applied to Equation 10 yields the fully linearized pressure coefficient

$$C_p = -\frac{2}{U_\infty^2} \left(\frac{\partial}{\partial t} + U \frac{\partial}{\partial x} \right) \varphi \quad (15)$$

and when applied to Equation 6 yields the fully linearized unsteady flow equation

$$\frac{\partial^2 \varphi}{\partial x^2} + \frac{\partial^2 \varphi}{\partial y^2} + \frac{\partial^2 \varphi}{\partial z^2} - \frac{1}{a_\infty^2} \left[U^2 \frac{\partial^2 \varphi}{\partial x^2} + 2U \frac{\partial^2 \varphi}{\partial x \partial t} + \frac{\partial^2 \varphi}{\partial t^2} \right] = 0 \quad (16)$$

Next, we write the body equation for a thin, nonplanar lifting surface (Figure 1), in terms of a curvilinear coordinate

$$s = s(y)$$

s represents the integral of distance along the line of the mean position of the airfoil from the centerline to y ,

$$s(y) = \int_0^y ds \quad (17)$$

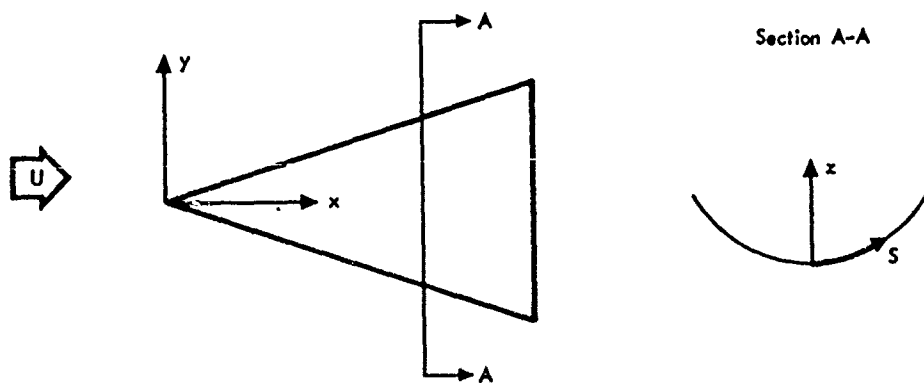


Figure 1. Generalized Curvilinear Planform

The position of the surface in terms of s and n (the normal to S), is separated into two parts; one part for the upper (or inner) surface, and the other for the lower (or outer) surface

$$B_u(x, n, s, t) = n - n_\tau(x, s) - n(x, s, t) \quad (18a)$$

$$B_l(x, n, s, t) = n + n_\tau(x, s) - n(x, s, t) \quad (18b)$$

where $n_\tau(x, s)$ represents the thickness of the airfoil and $n(x, s, t)$ represents the elastic deflection of the airfoil. In accordance with Equation 11, we use the operator

$$\nabla = i \frac{\partial}{\partial x} + j' \frac{\partial}{\partial s} + k' \frac{\partial}{\partial n} \quad (19)$$

on Equation 18a to get

$$\begin{aligned} \nabla B_u(x, n, s, t) = & -i \left[\frac{\partial}{\partial x} n_\tau(x, s) + \frac{\partial}{\partial x} n(x, s, t) \right] \\ & - j' \left[\frac{\partial}{\partial s} n_\tau(x, s) + \frac{\partial}{\partial s} n(x, s, t) \right] + k' \end{aligned}$$

Substitution into

$$\frac{D}{Dt} B_u(x, s, n, t) = \frac{\partial B_u}{\partial t} + (U i + \nabla \varphi) \cdot \nabla B_u = 0 \quad (20)$$

of the equation for the upper surface, after higher order terms have been discarded, gives

$$\frac{\partial \varphi}{\partial n} = \left(\frac{\partial}{\partial t} + U \frac{\partial}{\partial x} \right) n(x, s, t) + U \frac{\partial}{\partial x} n_\tau(x, s) \quad (21)$$

The same procedure, for the lower surface, gives

$$\frac{\partial \varphi}{\partial n} = \left(\frac{\partial}{\partial t} + U \frac{\partial}{\partial x} \right) n(x, s, t) - U \frac{\partial}{\partial x} n_T(x, s) \quad (22)$$

Finally, we restrict the analysis to that class of problems in which the effects of thickness on the time dependent forces can be neglected. By letting

$$n_T(x, s) = 0 \quad (23)$$

we get a single expression for the boundary condition

$$\frac{\partial \varphi}{\partial n} = \left(\frac{\partial}{\partial t} + U \frac{\partial}{\partial x} \right) n(x, s, t) \quad (24)$$

It is evident from Equation 24 that, when the motion of the surface is simple harmonic motion,

$$n(x, s, t) = \bar{n}(x, s) e^{i\omega t} \quad (25)$$

then,

$$\varphi(x, s, n, t) = \bar{\varphi}(x, s, n) e^{i\omega t} \quad (26)$$

Substitution of Equations 25 and 26 into Equations 15, 16, and 24 gives

$$\bar{C}_p = - \frac{2}{U_\infty^2} \frac{D\bar{\varphi}}{Dt} \quad (27)$$

$$\nabla^2 \bar{\varphi} = \frac{1}{a_\infty^2} \frac{D^2 \bar{\varphi}}{Dt^2} \quad (28)$$

$$\frac{\partial \bar{\varphi}}{\partial n} = \frac{D\bar{n}}{Dt} \quad (29)$$

where

$$\frac{D}{Dt} \equiv U \frac{\partial}{\partial x} + i\omega \quad (30)$$

and

$$\nabla^2 = \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial s^2} + \frac{\partial^2}{\partial n^2} \quad (31)$$

To this degree of approximation, $M = M_\infty$ and $a = a_\infty$.

IV. THE ACCELERATION POTENTIAL

PLANAR WINGS

Equation 27 shows that the calculation of pressure requires taking derivatives of the velocity potential, and Equation 29 states the boundary condition which must be satisfied. Watkins, Runyan and Woolston (Reference 1) solved this problem for planar surfaces in terms of a series of pressure functions, or acceleration potential functions ($\bar{\psi}$), where,

$$\bar{\psi}(x, y, z) = \frac{D}{Dt} \bar{\phi}(x, y, z) \quad (32)$$

Integration of Equation 32 gives the general expression for the velocity potential in terms of the acceleration potential:

$$\bar{\phi}(x, y, z) = \frac{1}{U} e^{-i\omega \frac{x}{U}} \int_{-\infty}^x e^{i\omega \frac{\lambda}{U}} \bar{\psi}(\lambda, y, z) d\lambda \quad (33)$$

if the velocity at infinity is $\bar{V} = U\mathbf{i}$. The velocity potential ϕ due to a pulsating doublet satisfies Equation 28 (in the planar case $s = y$ and $n = z$); and, because the order of operators is interchangeable, the acceleration potential ψ also satisfies Equation 28.

A complete discussion of the application of boundary conditions in the planar case is given in Section 6-4 of Reference 5. The application in the nonplanar case is discussed in less detail in the following text.

NONPLANAR WINGS

In the nonplanar case, the acceleration potential at a point (x, s, N) due to a pulsating doublet located at the point (ξ, σ, n) , or (ξ, η, ζ) in the direction of n is

$$\bar{\psi}(x, s, N) = -A \frac{\partial}{\partial n} \left[\frac{e^{i\omega \left[\frac{M}{a_{\infty}^2} (x-\xi) - \frac{R}{a_{\infty}^2} \right]}}{R} \right] \quad (34)$$

where

$$\frac{\partial}{\partial n} = \cos \gamma (\eta) \frac{\partial}{\partial \zeta} - \sin \gamma (\eta) \frac{\partial}{\partial \eta} \quad (35)$$

$$R = \sqrt{(x - \xi)^2 + \beta^2 [(y - \eta)^2 + (z - \zeta)^2]} \quad (36)$$

and $\gamma(\eta)$ is angle between the wing and the xy-plane at the point (ξ, η, ζ) . The perturbation velocity potential may be built up from a distribution of doublets of acceleration potential over the wing. If A is the infinitesimal doublet strength at (ξ, σ, η) , the contribution to $\bar{\phi}$ from the doublet at this point, from Equations 33 and 34, is

$$\Delta \bar{\phi}(x, s, N) = \frac{-A}{U} \frac{\partial}{\partial n} e^{-i\omega \frac{x - \xi}{U}} \int_{-\infty}^{x - \xi} e^{\frac{i\omega \left[\frac{\lambda}{U} + \frac{M\lambda}{a_\infty^2 \beta^2} - \frac{R'}{a_\infty^2 \beta^2} \right]} } \frac{d\lambda}{R'} \quad (37)$$

where

$$R' = \sqrt{\lambda^2 + \beta^2 [(y - \eta)^2 + (z - \zeta)^2]}$$

and the velocity component normal to the surface at (x, s, N) is

$$\Delta \bar{w}(x, s, 0) = \lim_{N \rightarrow 0} \frac{\partial}{\partial N} \Delta \bar{\phi}(x, s, N) \quad (38)$$

where

$$\frac{\partial}{\partial N} = \cos \gamma (y) \frac{\partial}{\partial z} - \sin \gamma (y) \frac{\partial}{\partial y} \quad (39)$$

Note that when the operator $\partial/\partial n$ is applied in Equation 37, the partials $\partial/\partial \zeta$ and $\partial/\partial \eta$ may be replaced by $-\partial/\partial z$ and $-\partial/\partial y$, respectively. Substitution of Equation 37 into Equation 38 gives

$$\Delta \bar{w}(x, s, 0) = \frac{A}{U} e^{-i\omega \frac{x - \xi}{U}} \lim_{N \rightarrow 0} P \int_{-\infty}^{x - \xi} e^{\frac{i\omega(\lambda - MR')/U\beta^2}{R'}} d\lambda \quad (40)$$

where the operator P is

$$P = \left[\cos \gamma (y) \frac{\partial}{\partial z} - \sin \gamma (y) \frac{\partial}{\partial y} \right] \left[\cos \gamma (\eta) \frac{\partial}{\partial z} - \sin \gamma (\eta) \frac{\partial}{\partial y} \right] \quad (41)$$

and finally

$$A = \frac{\Delta \bar{p} d\xi d\sigma}{4\pi \rho} \quad (42)$$

and $\Delta \bar{p}$ is the complex amplitude of the difference in pressure on the upper and lower sides of the surface at (ξ, σ) ,

$$\Delta \bar{p}(\xi, \sigma) = \bar{p}_u(\xi, \sigma) - \bar{p}_l(\xi, \sigma) \quad (43)$$

and $d\xi d\sigma$ is the incremental area of the doublet sheet.

The normal wash at $(x, s, 0)$ given by Equation 40 is that due to a point pressure doublet at $(\xi, \sigma, n = 0)$. The total normal wash is the integral over the surface of all the pressure doublets,

$$\frac{\bar{w}(x, s, 0)}{U} = \frac{-1}{4\pi \rho U^2} \iint \Delta \bar{p}(\xi, \sigma) K(x - \xi, s, \sigma, \omega, M) d\xi d\sigma \quad (44)$$

where \iint denotes the Mangler formula for evaluating infinite integrals (Reference 7), and the kernel of the integral equation is (omitting the arguments ω, M for brevity)

$$K(x_o, s, \sigma) = \lim_{\substack{n \rightarrow 0 \\ N \rightarrow 0}} \left\{ e^{-i \frac{\omega x_o}{U}} P \int_{-\infty}^{x_o} \frac{e^{i \omega (\lambda - MR')/U\beta^2}}{R'} d\lambda \right\} \quad (45)$$

where

$$R' = \sqrt{\lambda^2 + \beta^2 r_1^2}$$

$$r_1 = \sqrt{y_o^2 + z_o^2}$$

and

$$x_o = x - \xi, y_o = y - \eta, \text{ and } z_o = z - \zeta$$

Now, by putting $k_1 = \omega r_1 / U$ and $v = \lambda / \beta r_1$; then, by putting $u = -\frac{v - M \sqrt{1+v^2}}{\beta}$ the integral in Equation 45 may be written

$$I_o \equiv \int_{-\infty}^{x_o} \frac{e^{i\omega(\lambda - MR')/U\beta^2}}{R'} d\lambda = \int_{u_1}^{\infty} \frac{e^{-ik_1 u}}{\sqrt{1+u^2}} du \quad (46)$$

where

$$u_1 = -\frac{x_o - MR}{\beta^2 r_1}$$

By breaking up the interval of integration into three subintervals and in the first two integrals letting $u = w/i$

$$I_o = -i \int_0^1 \frac{e^{-k_1 w}}{\sqrt{1-w^2}} dw + \int_1^{\infty} \frac{e^{-k_1 w}}{\sqrt{w^2-1}} dw - \int_0^{u_1} \frac{e^{-ik_1 u}}{\sqrt{1+u^2}} du$$

or,

$$I_o = K_o(k_1) - i \int_0^1 \frac{e^{-k_1 w}}{\sqrt{1-w^2}} dw - \int_0^{u_1} \frac{e^{-ik_1 u}}{\sqrt{1+u^2}} du \quad (47)$$

where $K_o(k_1)$ is the modified Bessel function of the second kind and zeroth order of the argument k_1 .

To obtain the analytic form of the kernel, we write the operator P (Equation 41) in a more convenient form, taking advantage of the fact that I_o is a function of only r_1 when x_o is held constant

$$PI_o = \cos [\gamma(y) - \gamma(\eta)] \left(\frac{1}{r_1} \frac{\partial I_o}{\partial r_1} \right) + [z_o \cos \gamma(y) - y_o \sin \gamma(y)] [z_o \cos \gamma(\eta) - y_o \sin \gamma(\eta)] \left(\frac{1}{r_1} \frac{\partial}{\partial r_1} \right) \left(\frac{1}{r_1} \frac{\partial I_o}{\partial r_1} \right)$$

The resulting kernel is

$$\begin{aligned}\bar{K}(x_o, s, \sigma) &= s_{TIP}^2 K(x_o, s, \sigma) \\ &= e^{-i \frac{\omega x_o}{U}} \left\{ \frac{T_1 K_1(x_o, s, \sigma) + T_2 K_2(x_o, s, \sigma)}{r_1^2} \right\} \quad (48)\end{aligned}$$

where

$$\begin{aligned}T_1 &= \cos [v(\underline{s}) - v(\underline{\sigma})] \\ T_2 &= \left[\frac{z_o}{r_1} \cos v(\underline{s}) - \frac{y_o}{r_1} \sin v(\underline{s}) \right] \left[\frac{z_o}{r_1} \cos v(\underline{\sigma}) - \frac{y_o}{r_1} \sin v(\underline{\sigma}) \right] \quad (48a)\end{aligned}$$

$$\begin{aligned}K_1(x_o, s, \sigma) &= -k_1 K_1(k_1) - \frac{x_o}{R} e^{-k_1 u_1} + ik_1 \int_0^1 \frac{w e^{-k_1 w}}{\sqrt{1-w^2}} dw \\ &\quad + ik_1 \int_0^{u_1} \frac{u e^{-ik_1 u}}{\sqrt{1+u^2}} du \\ K_2(x_o, s, \sigma) &= k_1^2 K_2(k_1) + \left(\frac{2x_o}{R} + \frac{\beta^2 r_1^2 x_o}{R^3} + i \frac{k_1 x_o (Mr_1 + Ru_1)}{R^2} \right) e^{-ik_1 u_1} \\ &\quad - ik_1 \int_0^1 \frac{w e^{-k_1 w}}{\sqrt{1-w^2}} dw - ik_1^2 \int_0^1 \frac{w^2 e^{-k_1 w}}{\sqrt{1-w^2}} dw \\ &\quad - ik_1 \int_0^{u_1} \frac{u e^{-ik_1 u}}{\sqrt{1+u^2}} du + k_1^2 \int_0^{u_1} \frac{u^2 e^{-ik_1 u}}{\sqrt{1+u^2}} du \quad (48b)\end{aligned}$$

$$r_1 = s_{TIP} \sqrt{y_o^2 + z_o^2}$$

$$R = \sqrt{x_o^2 + \beta^2 r_1^2}$$

$$u_1 = - \frac{x_o - MR}{\beta^2 r_1}$$

$$k_1 = \frac{\omega r_1}{U} \quad (48c)$$

and $K_1(k_1)$ and $K_2(k_1)$ are modified Bessel functions of the second kind and first and second orders. The sub-bar indicates division by s_{TIP} , e.g., $\bar{r}_1 = r_1/s_{TIP}$.

V. THE BOUNDARY CONDITIONS

The remainder of the problem is to match the boundary conditions; i. e., to find a pressure-loading function $\Delta \bar{p}(\xi, \sigma)$ which, when inserted into Equation 45 and integrated over the surface, yields the kinematic down-washes at selected points on the surface $w(x, s, 0)$.

In subsonic flow, the behavior of the pressure distribution is known in the area of the wing edges from a few of the exact solutions in lifting surface theory. In the neighborhood of the leading edge, the pressure should behave as

$$\lim_{\delta \rightarrow 0} \sqrt{\frac{1}{\delta}}$$

In the neighborhood of the trailing edge and all edges parallel to the free-stream direction, the pressure should behave as

$$\lim_{\delta \rightarrow 0} \sqrt{\delta}$$

where δ is the distance to the wing edge. Both Hsu and Watkins employ a linear superposition of functions that satisfy these conditions. Hsu's function differs from Watkins only in that for any given number of terms in the series, Hsu's terms are linear combinations of Watkins terms. We use a normalized form of the function given by Watkins:

$$\Delta \bar{p}(\xi, \sigma) \equiv \frac{\rho U^2/2}{b(\sigma)} \sqrt{1 - \sigma^2} \sum_{n=0}^N \sum_{m=0}^M a_{nm} \sigma^m f_n(\xi) \quad (49)$$

where

$$f_0(\xi) \equiv \sqrt{\frac{1 - \xi}{1 + \xi}}$$

$$f_n(\xi) \equiv \sqrt{1 - \xi^2} U_n(\xi); 1 \leq n$$

$$U_1(\tilde{\xi}) = 1.0$$

$$U_2(\tilde{\xi}) = -2\tilde{\xi}$$

$$U_n(\tilde{\xi}) = -(2\tilde{\xi} U_{n-1} + U_{n-2}); 3 \leq n$$

and

$$\tilde{\xi} \equiv \left(\xi - \frac{\xi_{LE} + \xi_{TE}}{2} \right) / b(\sigma)$$

$$b(\sigma) = \frac{\xi_{TE} - \xi_{LE}}{2}$$

The a_{nm} 's are unknown pressure coefficients to be determined by matching the kinematic downwashes at the selected points (x_j, s_r) on the surface. Substitution of Equation 49 into Equation 44 leads to the matrix equation given by Rodden and Revell (Reference 2) Equation 39, for the point set x_j, s_r

$$\left[\frac{\bar{w}_i}{U} \right] = [D_{nm}^i] [a_{nm}] \quad (50)$$

where, in this case,

$$D_{nm}^i = \frac{1}{8\pi} \int_{-1}^1 \sqrt{1 - \sigma^2} \sigma^m \int_{-1}^1 f_n(\tilde{\xi}) \bar{K}(x_j - \xi, s_r, \sigma) d\tilde{\xi} d\sigma \quad (51)$$

We now reexamine the fundamentals of the problem before proceeding to evaluate the integrals in Equation 51 and thence to solve Equation 50.

One of the basic reasons for development of the kernel function method is that pressure distributions over a continuous lifting surface are smooth continuous functions that can be represented with reasonable accuracy by a series of analytic functions. We point out that $f_n(\tilde{\xi})$ can be written

$$f_0(\tilde{\xi}) = \frac{1 - \tilde{\xi}}{\sqrt{1 - \tilde{\xi}^2}}$$

$$f_n(\tilde{\xi}) = \frac{1 - \tilde{\xi}^2}{\sqrt{1 - \tilde{\xi}^2}} U_n(\tilde{\xi}); 1 \geq n$$

and, therefore, the inner integral in Equation 51 may be written

$$I_1 = \int_{-1}^1 \frac{P_n(\tilde{\xi})}{\sqrt{1 - \tilde{\xi}^2}} d\tilde{\xi} \quad (52)$$

where

$$P_0(\tilde{\xi}) = (1 - \tilde{\xi}) \bar{K}(x_j - \tilde{\xi}, \underline{s}_r, \underline{\sigma})$$

$$P_n(\tilde{\xi}) = (1 - \tilde{\xi}^2) U_n(\tilde{\xi}) \bar{K}(x_j - \tilde{\xi}, \underline{s}_r, \underline{\sigma}); 1 \leq n$$

Now, we assume for the moment that the kernel $\bar{K}(X_j - \xi, s_r, \sigma)$ can be represented with reasonable accuracy by a polynomial in ξ . Then, $P_n(\tilde{\xi})$ is also a polynomial in $\tilde{\xi}$, and the Chebyshev-Gauss quadrature formula may be used to obtain the exact value of the integral expression (52); i. e.,

$$\int_{-1}^1 \frac{P_n(\tilde{\xi})}{\sqrt{1 - \tilde{\xi}^2}} d\tilde{\xi} = \sum_{k=1}^K \frac{\pi}{K} P_n(\tilde{\xi}_k) + E \quad (52a)$$

where

$$E = \frac{2\pi}{2^{2K} (2K)!} P_n^{(2K)}(\lambda)$$

and,

$$|\lambda| < 1.0$$

The error term E is zero if P_n is a polynomial of degree $\leq 2K - 1$.

A more accurate formula which utilizes the fact that $P_n(1.0) = 0$ is used by Hsu (and by us)

$$f_0(\tilde{\xi}) = \sqrt{\frac{1 - \tilde{\xi}}{1 + \tilde{\xi}}}$$

$$f_n(\tilde{\xi}) = \sqrt{\frac{1 - \tilde{\xi}}{1 + \tilde{\xi}}} (1 + \tilde{\xi}) U_n(\tilde{\xi})$$

Then, the inner integral in Equation 51 may be written

$$I_1 = \int_{-1}^1 \sqrt{\frac{1-\tilde{\xi}}{1+\tilde{\xi}}} F_n(\tilde{\xi}) d\tilde{\xi} \quad (53)$$

where

$$F_0(\tilde{\xi}) = \bar{K}(x_j - \tilde{\xi}, \underline{s}_r, \underline{\sigma})$$

$$F_n(\tilde{\xi}) = (1 + \tilde{\xi}) U_n(\tilde{\xi}) \bar{K}(x_j - \tilde{\xi}, \underline{s}_r, \underline{\sigma}), \quad 1 \leq n$$

This is evaluated by the L-point Jacobi-Gauss quadrature with the weight function $\sqrt{(1-\tilde{\xi})/(1+\tilde{\xi})}$ (see Reference 8, Chapter 8). The resulting formula

$$I \cong \sum_{k=1}^L W_k F_n(\tilde{\xi}_k) \quad (54)$$

is exact if $F_n(\tilde{\xi})$ is a polynomial of degree $\leq 2L-1$, which corresponds to degree $2L$ for $P_n(\tilde{\xi})$. Putting $\tilde{\xi} = -\cos \theta$, the polynomials

$$\phi_m(\tilde{\xi}) = \frac{\cos\left(m + \frac{1}{2}\right)\theta}{\cos\frac{1}{2}\theta}$$

are orthogonal with respect to the weight function $\sqrt{(1-\tilde{\xi})/(1+\tilde{\xi})}$

$$\int_{-1}^1 \phi_m(\tilde{\xi}) \phi_n(\tilde{\xi}) \sqrt{\frac{1-\tilde{\xi}}{1+\tilde{\xi}}} d\tilde{\xi} = \begin{cases} 0, & m \neq n \\ \pi, & m = n \end{cases}$$

This is easily verified by expressing the integral in terms of θ . Referring to formulas in Reference 8, Chapter 8, Section 8.4, it can be shown, using these polynomials, that Equation 54 takes the form

$$\int_{-1}^1 \sqrt{\frac{1-\tilde{\xi}}{1+\tilde{\xi}}} F_n(\tilde{\xi}) d\tilde{\xi} = \frac{2\pi}{2L+1} \sum_{k=1}^L H_k F_n(\tilde{\xi}_k) \quad (55)$$

where

$$H_k = (1 - \tilde{\xi}_k)$$

and

$$\tilde{\xi}_k = -\cos\left(\frac{2k-1}{2L+1} \pi\right)$$

In two-dimensional, steady, incompressible flow, there is an optimum set of chordwise collocation stations (\tilde{x}_j) for the determination of sectional lift, depending upon the order of the polynomial required to adequately represent the downwash distribution

$$\tilde{x}_j = -\cos\left(\frac{2j}{2N+1} \pi\right), j = 1, 2, \dots, N. \quad (56)$$

Since the behavior of the integrand for the chordwise loading is apt to exhibit similar characteristics near the surface edges, it is inferred that this set should also yield the best approximation in three-dimensional, unsteady, compressible flow. Note that the number of collocation points is not required to be the same as the number of integration points. As will be seen later, it is only necessary that the total number of downwash collocation points be equal to or greater than the number of pressure coefficients a_{nm} . When N chordwise integration stations are used, the quadrature used to evaluate the inner integral of Equation 51 is exact for integrands represented by a polynomial of degree $\leq 2N - 1$.

Hsu shows that an optimum set of interdigitated spanwise collocation stations and integration stations exists for evaluation of the outer integral in Equation 51. By reasoning similar to that used to establish the chordwise collocation stations, it was established that the optimum spanwise collocation stations are

$$\frac{s}{-r} = -\cos \frac{r}{M+1} \pi, r = 1, 2, \dots, M. \quad (57)$$

It was observed that the quadrature for the integral of difference between the actual and polynomial approximation of the spanwise loading is zero when the actual loading is precisely represented by a polynomial of degree $\leq 2M - 1$, and the polynomial approximation is of degree $= N - 1$.

Then, by substitution of Equation 55 into Equation 51,

$$D_{nm}^i = \frac{1}{8\pi} \int_{-1}^1 \frac{\sqrt{1 - \sigma^2}}{(\underline{s}_r - \sigma)^2} G_{nm}(\tilde{x}_j, \underline{s}_r, \sigma) d\sigma \quad (58)$$

where

$$G_{om} = \frac{2\pi}{2L+1} \sum_{k=1}^L \sigma^m (1 - \tilde{\xi}_k) (\underline{s}_r - \sigma)^2 \overline{K}(\underline{x}_j - \underline{\xi}_k, \underline{s}_r, \sigma) \quad (59)$$

$$G_{nm} = \frac{2\pi}{2L+1} \sum_{k=1}^L \sigma^m (1 - \tilde{\xi}_k^2) U_n(\tilde{\xi}_k) (\underline{s}_r - \sigma)^2 \overline{K}(\underline{x}_j - \underline{\xi}_k, \underline{s}_r, \sigma); 1 \leq n$$

Hsu established the form of the Gaussian quadrature and the spanwise integration stations. The difficulties of the singularity of the kernel at $\sigma = \underline{s}_r$ and the difficulty of differentiation with respect to σ (he uses the steady-state lifting line formula to derive the form of the quadrature) are avoided by removal of the singularity at $\sigma = \underline{s}_r$ and then by an integration by parts.

We first integrate by parts to get

$$D_{nm}^i = \frac{1}{8\pi} \int_{-1}^1 \frac{\frac{\partial}{\partial \sigma} \left[\sqrt{1 - \sigma^2} G_{nm}(\tilde{x}_j, \underline{s}_r - \sigma, \omega, M) \right]}{(\underline{s}_r - \sigma)} d\sigma$$

which corresponds to Equation 58 in Reference 9. The Gaussian quadrature formula is developed and shows that when the number of integration stations is one greater than the number of collocation stations, and if they are interdigitated in the prescribed way

$$D_{nm}^i = \frac{1}{8\pi} \left\{ \sum_{p=1}^{M+1} \frac{\pi}{M+1} \frac{(1 - \sigma_p^2) G_{nm}(\tilde{x}_j, \underline{s}_r, \underline{\sigma}_p)}{(\underline{s}_r - \underline{\sigma}_p)^2} - \pi (M+1) G_{nm}(\tilde{x}_j, \underline{s}_r, \underline{s}_r) \right\} \quad (60)$$

where

$$\underline{s}_r = -\cos \frac{r\pi}{M+1}$$

and

$$\underline{\sigma}_p = -\cos \frac{2p-1}{2(M+1)} \pi$$

$$r = 1, 2, \dots, M$$

Evaluation of the second term in the brackets requires the observation that the multiplier of K_2 in Equation 48 goes to zero whenever the collocation point is in the plane of the doublet sheet located at the integration point. Therefore, the finite part of the integral of the K_2 term is zero, and the entire contribution comes from the K_1 term. In this case, $\gamma(\underline{s}) = \gamma(\underline{\sigma})$ and $r_1^2 = (\underline{s}_r - \underline{\sigma})^2 = 0$.

It can be shown that

$$K_1(x - \xi, \underline{s}_r, \underline{s}_r) = \begin{cases} -2, & x > \xi \\ 0, & x < \xi. \end{cases}$$

Thus, the chordwise integral which defines $G_{nm}(x_j, \underline{s}_r, \underline{s}_r)$ is

$$G_{nm}(x_j, \underline{s}_r, \underline{s}_r) = -2 \int_{-1}^{\tilde{x}_j} \underline{\sigma}^m f_n(\tilde{\xi}) e^{-i \frac{\omega}{U} (x_j - \xi)} d\tilde{\xi}$$

If the range of integration is extended to $\tilde{\xi} = 1$ by making the integrand zero for $\tilde{\xi} > \tilde{x}_j$, the integral cannot be well approximated by a polynomial because it has a jump discontinuity at $\tilde{\xi} = \tilde{x}_j$. To overcome this difficulty, we write

$$G_{nm}(x_j, s_r, s_r) = -2 \int_{-1}^{\tilde{x}_j} \sigma^m f_n(\tilde{\xi}) \left[e^{-i \frac{\omega}{U} (x_j - \tilde{\xi})} - 1 \right] d\tilde{\xi} \quad (61a)$$

$$-2 \int_{-1}^{\tilde{x}_j} \sigma^m f_n(\tilde{\xi}) d\tilde{\xi}.$$

The second term here depends on the integrals

$$h_n(\tilde{x}_j) = \int_{-1}^{\tilde{x}_j} f_n(\tilde{\xi}) d\tilde{\xi}$$

which may be evaluated exactly. We have

$$h_0(\tilde{x}_j) = \int_{-1}^{\tilde{x}_j} \sqrt{\frac{1-\tilde{\xi}}{1+\tilde{\xi}}} d\tilde{\xi} = \frac{\pi}{2} + \sin^{-1} \tilde{x}_j + \sqrt{1-\tilde{x}_j^2}$$

$$h_1(\tilde{x}_j) = \int_{-1}^{\tilde{x}_j} \sqrt{1-\tilde{\xi}^2} d\tilde{\xi} = \frac{1}{2} \left(\frac{\pi}{2} + \sin^{-1} \tilde{x}_j + \tilde{x}_j \sqrt{1-\tilde{x}_j^2} \right)$$

This list may be extended as far as it is needed.

The first integral in Equation 61a is considered as an integral over $-1 < \tilde{\xi} < 1$ with zero integrand when $\tilde{\xi} > \tilde{x}_j$, and Equation 55 is applied. The result is

$$G_{om}(x_j, s_r, s_r) = -2 \cdot \frac{2\pi}{2L+1} \sum_{\substack{k \geq 1, \\ \tilde{\xi}_k < \tilde{x}_j}} \sigma^m (1 - \tilde{\xi}_k) \left[e^{-i \frac{\omega}{U} (x_j - \tilde{\xi}_k)} - 1 \right] \\ - 2 \sigma^m g_o(\tilde{x}_j) \quad (61b)$$

$$G_{nm}(x_j, s_r, s_r) = -2 \frac{2\pi}{2L+1} \sum_{\substack{k \geq 1, \\ \tilde{\xi}_k < \tilde{x}_j}} \sigma^m (1 - \tilde{\xi}_k^2) U_n(\tilde{\xi}_k) \left[e^{-i \frac{\omega}{U} (x_j - \xi_k)} - 1 \right]$$

$$- 2\sigma^m g_n(\tilde{x}_j), \quad n \geq 1.$$

Equations 59 and 61 are used to evaluate the D_{nm}^i 's given by Equation 60. Equation 50 is then used to determine the pressure coefficients a_{nm} to match the downwashes (i.e. the \bar{w}_i/U 's) at the collocation points (x_j, s_r) . Once the pressure coefficients are determined, the generalized forces are computed. A polynomial expression for the i^{th} modal deflections normal to the surface,

$$n^{(i)}(\xi, \sigma) = \sum_{\nu=0}^N \sum_{\mu=0}^M b_{\nu\mu}^{(i)} \xi^{\nu} \sigma^{\mu}$$

and Equation 49 are substituted into the equation

$$Q_{ij} = \int_{-s_{\text{TIP}}}^{s_{\text{TIP}}} \int_{x_{\text{LE}}}^{x_{\text{TE}}} n^{(i)}(\xi, \sigma) \Delta \bar{p}^{(j)}(\tilde{\xi}, \sigma) d\tilde{\xi} d\sigma$$

to get

$$Q_{ij} = s_{\text{TIP}}^2 (1/2 \rho U^2) \sum_{n=0}^N \sum_{m=0}^M \sum_{\nu=0}^N \sum_{\mu=0}^M a_{nm}^{(j)} b_{\nu\mu}^{(i)} \Delta Q_{nm\nu\mu} \quad (62)$$

where

$$\Delta Q_{nm\nu\mu} = \int_{-1}^1 \int_{-1}^1 \sqrt{1 - \sigma^2} \sigma^{m+u} \xi^{\nu} \sqrt{\frac{1 - \tilde{\xi}}{1 + \tilde{\xi}}} F_n(\tilde{\xi}) d\tilde{\xi} d\sigma$$

The quadrature formula given by Equation 55 with L set equal to N may be used to evaluate the inner integral. For evaluation of the outer integral, a quadrature formula with the weight function $\sqrt{1 - \sigma^2}$, analogous to Equation 60 for a nonsingular integral, may be used. This is not practical for a wing

with folded tip, for which a different representation should be used on each plane part of the surface. Here, the spanwise integral over one part of the wing may have a square root factor at one end of the interval of integration, or at neither end. Such integrals may best be evaluated by a suitable application of ordinary Gaussian quadrature with a weight function of 1.0. The points and weights for this quadrature method are not given by simple formulas such as Equation 55. They are listed in many places, (e. g. Reference 5).

VI. APPLICATION TO A WING WITH A FOLDED TIP

The planform may be any continuous surface. However, the computer program was developed to treat either a planar or nonplanar planform of the type shown in Figure 2. (Only one-half the planform is shown.) To facilitate modifications of the program, comment cards are placed throughout the program to indicate where changes may be made to handle other nonplanar surfaces like that of the Paraglider.

In the application of the kernel function method to the planform shown in Figure 2, the computer program calculates from the equations of the leading and trailing edges the collocation and integration points for which the integrands of Equation 51 must be evaluated. For demonstration purposes, we calculate the collocation and integration points for values of $L = 4$, $N = 6$, and $M = 10$ and show the results in Figure 3.

$$x_{LE} = s \tan \lambda_{LE}, \quad x_{LETIP} = s_{FL} \tan \lambda_{LE} + (s - s_{FL}) \tan \lambda_{LETIP}$$

$$x_{TE} = 2b_o + s \tan \lambda_{TE}, \quad x_{TETIP} = 2b_o + s_{FL} \tan \lambda_{TE} + (s - s_{FL}) \tan \lambda_{TETIP}$$

$$b(s) = 1/2 (x_{TE} - x_{LE})$$

$$x_m = 1/2 (x_{LE} + x_{TE})$$

$$x = x_m + b(s) \tilde{x}$$

$$\xi = \xi_m + b(\sigma) \tilde{\xi}$$

$$\tilde{x}_j = -\cos \frac{2j}{2N+1} \pi, \quad j = 1, 2, \dots, N.$$

$$\tilde{\xi}_k = -\cos \frac{2k-1}{2L+1} \pi, \quad k = 1, 2, \dots, L.$$

$$s = \underline{s} s_{TIP}$$

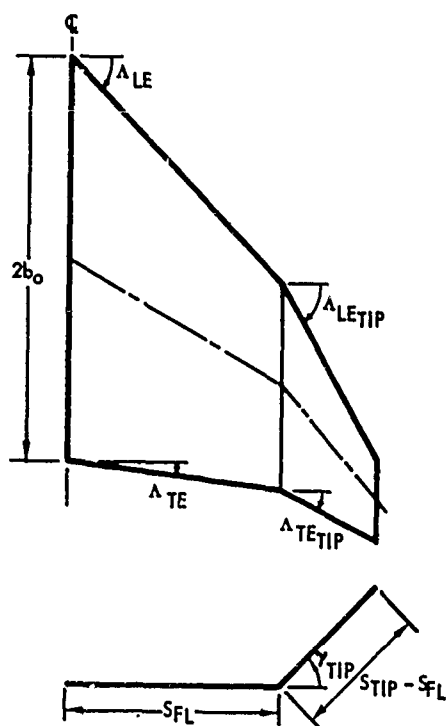


Figure 2. Planar Wing With Planar Symmetrically Folded Tips

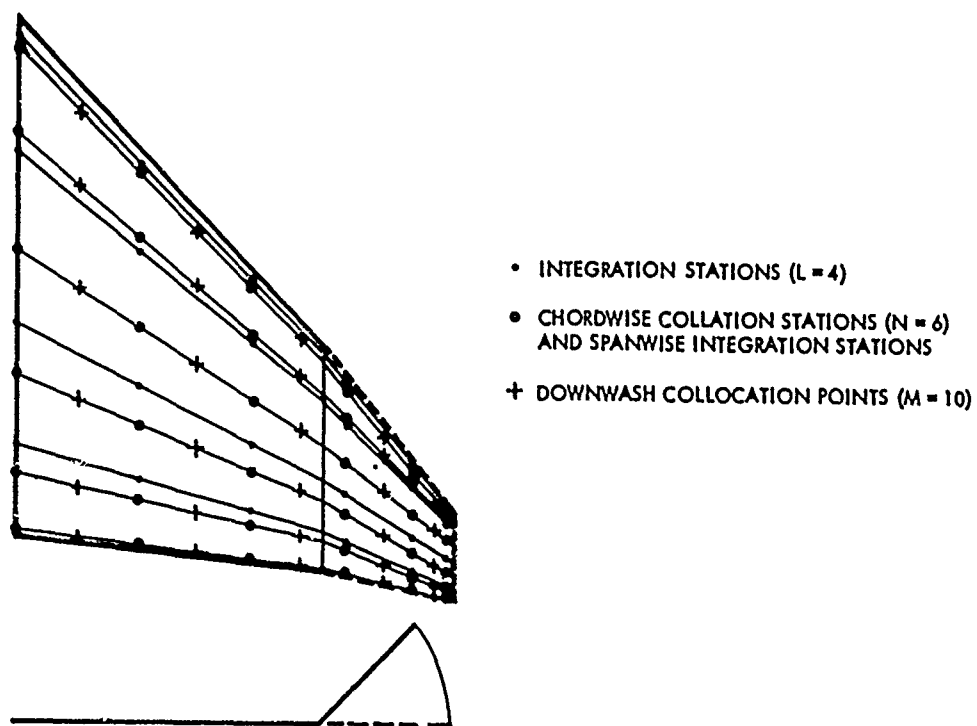


Figure 3. An Optimum Set of Collocation and Integration Points

$$\underline{s}_m = -\cos \frac{m\pi}{M+1}, \quad m = 1, 2, \dots, M.$$

$$\underline{\sigma}_p = -\cos \frac{2p-1}{2(M+1)} \pi, \quad p = 1, 2, \dots, (M+1).$$

We have also constructed a table of equations (Table 1), for the functions y_0 , z_0 , r_1 , T_1 , and T_2 for use in the expression for the kernel, Equations 48 through 48c. In Table 1, the subscript F_L indicates the fold line.

A difficulty is encountered in the spanwise integration because the kernel function has a finite discontinuity at the fold line. For example, note in Table 1 the change in T_1 and T_2 for receiving or collocation points on the wing as the sending or integration points shift from the port tip to the wing and from the wing to the starboard tip.

If we consider the kernel as a function of $\tilde{\xi}$ and $\underline{\sigma}$ for fixed values of x and \underline{s} ,

$$q(\tilde{\xi}, \underline{\sigma}) = \bar{K}(x - \xi, \underline{s}, \underline{\sigma}),$$

this function may be broken up into a simple discontinuous part g^{**} and a part g^* which is continuous across the fold lines:

$$g(\tilde{\xi}, \underline{\sigma}) = g^*(\tilde{\xi}, \underline{\sigma}) + g^{**}(\tilde{\xi}, \underline{\sigma}) \quad (63)$$

To do this, define

$$g^{**}(\tilde{\xi}, \underline{\sigma}) = \begin{cases} g(\tilde{\xi}, \underline{\sigma}_{F_L}^+) - g(\tilde{\xi}, \underline{\sigma}_{F_L}^-), & \underline{\sigma} > \underline{\sigma}_{F_L} \\ 0, & -\underline{\sigma}_{F_L} < \underline{\sigma} < \underline{\sigma}_{F_L} \\ g(\tilde{\xi}, -\underline{\sigma}_{F_L}^-) - g(\tilde{\xi}, -\underline{\sigma}_{F_L}^+), & \underline{\sigma} < -\underline{\sigma}_{F_L} \end{cases} \quad (64)$$

and then define g^* by Equation 63.

More explicitly, for $\underline{\sigma} > \underline{\sigma}_{F_L}$, define

$$\xi_{F_L} = b(\underline{\sigma}_{F_L}) \tilde{\xi} + \frac{1}{2} \left[\xi_{LE}(\underline{\sigma}_{F_L}) + \xi_{TE}(\underline{\sigma}_{F_L}) \right]$$

Table 1. Spanwise Parameter in the Kernel Function

Receiving Points (x, s)									
Port Tip				Wing			Starboard Tip		
$S < -Y_{FL}$ and $\gamma(S) = -\gamma_{TIP}$				$-Y_{FL} < S < Y_{FL}$ and $\gamma(S) = 0$			$Y_{FL} < S$ and $\gamma(S) = \gamma_{TIP}$		
Sending Points (ξ, σ)									
Functions	PT	W	ST	PT	W	ST	PT	W	ST
	$\gamma(\sigma) = -\gamma_{TIP}$	$\gamma(\sigma) = 0$	$\gamma(\sigma) = \gamma_{TIP}$	$\gamma(\sigma) = \gamma_{TIP}$	$\gamma(\sigma) = 0$	$\gamma(\sigma) = \gamma_{TIP}$	$\gamma(\sigma) = -\gamma_{TIP}$	$\gamma(\sigma) = 0$	$\gamma(\sigma) = \gamma_{TIP}$
Y_0	$a_0 \xi$	$-\sigma^2 + a_0^2 \xi$	$(a_0^2 - \sigma_0^2) \xi - 2\gamma_{FL}$	$a_0^2 - \sigma^2 \xi$	a_0	$a_0 - \sigma_0 \xi$	$(a_0 - \sigma^2) \xi + 2\gamma_{FL}$	$-\sigma_0 + a_0 \xi$	$a_0 \xi$
z_0	$-a_0 \beta$	$-a_0^2 \beta$	$-(a_0 + \sigma) \beta$	$\sigma^2 \beta$	0	$-\sigma_0 \beta$	$(a_0 + \sigma^2) \beta$	$a_0 \beta$	$a_0 \beta$
r_1	$ a_0 $	$\sqrt{a_0^2 + \sigma^2}$	$\sqrt{[(a_0^2 - \sigma_0^2) \xi - 2\gamma_{FL}]^2 + [(a_0^2 + \sigma_0^2) \beta]^2}$	$\sqrt{a_0^2 + \sigma^2}$	$ a_0 $	$\sqrt{a_0^2 + \sigma^2}$	$\sqrt{[(a_0 - \sigma^2) \xi + 2\gamma_{FL}]^2 + [(a_0 + \sigma^2) \beta]^2}$	$\sqrt{a_0^2 + \sigma^2}$	$ a_0 $
T_1	1.0	ξ	$\xi^2 - \beta^2$	ξ	1.0	ξ	$\xi^2 - \beta^2$	ξ	1.0
T_2	0	$a_0^2 \beta^2$	$4(a_0 \xi + \gamma_{FL})(a_0^2 \xi - \gamma_{FL}) \beta^2$	$a_0^2 \beta^2$	0	$a_0^2 \beta^2$	$4(a_0^2 \xi - \gamma_{FL})(a_0 \xi + \gamma_{FL}) \beta^2$	$a_0^2 \beta^2$	0
$a_0^2 = a_0 + \gamma_{FL}$, $\sigma^2 = \sigma + \gamma_{FL}$, $a_0 = a - \gamma_{FL}$, $\sigma_0 = \sigma - \gamma_{FL}$, $a_0 = a \sin \gamma_{TIP}$, $\beta = \cos \gamma_{TIP}$									

Then

$$g^{**}(\tilde{\xi}, \underline{\sigma}) = \bar{K}^{ST} (x - \xi_{FL}, \underline{\sigma}, \underline{\sigma}_{FL}) - \bar{K}^W (x - \xi_{FL}, \underline{\sigma}, \underline{\sigma}_{FL})$$

in which

$$v(\underline{\sigma}) = v_{TIP} \text{ in } \bar{K}^{ST} \quad (ST \sim \text{starboard tip})$$

$$v(\underline{\sigma}) = 0 \text{ in } \bar{K}^W \quad (W \sim \text{wing})$$

A similar formula applies when $\underline{\sigma} < -\underline{\sigma}_{FL}$. g^{**} may be written in a form which indicates that it is independent of σ in each tip region:

$$g^{**}(\tilde{\xi}, \underline{\sigma}) = \begin{cases} g_{ST}^{**}(\tilde{\xi}), & \underline{\sigma} > \underline{\sigma}_{FL} \\ 0, & -\underline{\sigma}_{FL} < \underline{\sigma} < \underline{\sigma}_{FL} \\ g_{PT}^{**}(\tilde{\xi}), & \underline{\sigma} < -\underline{\sigma}_{FL} \end{cases} \quad (65)$$

With the use of Equations 63 and 65, Equation 51 may be rewritten as

$$\begin{aligned} D_{mn}^i &= \frac{1}{8\pi} \int_{-1}^1 \sqrt{1 - \underline{\sigma}^2} \underline{\sigma}^m \int_{-1}^1 f_n(\tilde{\xi}) g^*(\tilde{\xi}, \underline{\sigma}) d\tilde{\xi} d\underline{\sigma} \\ &+ \frac{1}{8\pi} \int_{-1}^1 f_n(\tilde{\xi}) g_{ST}^{**}(\tilde{\xi}) \int_{\underline{\sigma}_{FL}}^1 \sqrt{1 - \underline{\sigma}^2} \underline{\sigma}^m d\underline{\sigma} d\tilde{\xi} \\ &+ \frac{1}{8\pi} \int_{-1}^1 f_n(\tilde{\xi}) g_{PT}^{**}(\tilde{\xi}) \int_{-1}^{-\underline{\sigma}_{FL}} \sqrt{1 - \underline{\sigma}^2} \underline{\sigma}^m d\underline{\sigma} d\tilde{\xi} \end{aligned}$$

The first of these three double integrals is calculated according to Equation 60. In the others, the inner integral may be evaluated exactly. The constants

$$u_m = \int_{\frac{\sigma}{F_L}}^1 \sqrt{1 - \sigma^2} \sigma^m d\sigma$$

are given by formulas

$$u_0 = \frac{1}{2} \left[\frac{\pi}{2} - \sin^{-1} \frac{\sigma}{F_L} + \frac{\sigma}{F_L} \sqrt{1 - \frac{\sigma^2}{F_L^2}} \right]$$

$$u_1 = \frac{1}{3} (1 - \frac{\sigma^2}{F_L^2})^{3/2}$$

etc.

In terms of these constants

$$D_{mn}^i = \frac{1}{8\pi} \int_{-1}^1 \sqrt{1 - \sigma^2} \sigma^m \int_{-1}^1 f_n(\tilde{\xi}) g^*(\tilde{\xi}, \sigma) d\tilde{\xi} d\sigma$$

$$+ \frac{1}{8\pi} u_m \int_{-1}^1 f_n(\tilde{\xi}) \left[g_{ST}^{**}(\tilde{\xi}) + (-1)^m g_{PT}^{**}(\tilde{\xi}) \right] d\tilde{\xi}$$

The last integral in this formula is evaluated by Equation 55.

In the evaluation of O_{ij} , given by Equation 62 for the plane case, it is assumed that modes numbers i and j are either both symmetric in σ , or both antisymmetric. Then the contribution to the integral for $\sigma < 0$ is the same as the contribution for $\sigma > 0$. Let the deflection in the i^{th} mode be given by

$$n^{(i)}(\xi, \sigma) = \begin{cases} \sum_{\nu=0}^N \sum_{\mu=0}^M c_{\nu\mu}^{(i)} \xi^{\nu} \sigma^{\mu}, & 0 < \sigma < \sigma_{FL} \\ \sum_{\nu=0}^N \sum_{\mu=0}^M d_{\nu\mu}^{(i)} \xi^{\nu} \sigma^{\mu}, & \sigma > \sigma_{FL} \end{cases}$$

Then

$$O_{ij} = 2 \cdot y_{TIP}^2 \left(\frac{1}{2} \rho U^2 \right) \sum_{n=0}^N \sum_{m=0}^M \sum_{\nu=0}^N \sum_{\mu=0}^M a_{nm}^{(j)} \left\{ c_{\nu\mu}^{(i)} I_{nm\nu\mu}^{(1)} + d_{\nu\mu}^{(i)} I_{nm\nu\mu}^{(2)} \right\}$$

in which

$$I_{nm\nu\mu}^{(1)} = \int_0^{\sigma_{FL}} \sqrt{1 - \sigma^2} \sigma^{m+\mu} \int_{-1}^1 \xi^\nu \sqrt{\frac{1 - \tilde{\xi}}{1 + \tilde{\xi}}} F_n(\tilde{\xi}) d\tilde{\xi} d\sigma$$

$$I_{nm\nu\mu}^{(2)} = \int_{\sigma_{FL}}^1 \sqrt{1 - \sigma^2} \sigma^{m+\mu} \int_{-1}^1 \xi^\nu \sqrt{\frac{1 - \tilde{\xi}}{1 + \tilde{\xi}}} F_n(\tilde{\xi}) d\tilde{\xi} d\sigma$$

Note that the inner integrals are the integrals of polynomials in $\tilde{\xi}$ multiplied by $\sqrt{(1 - \tilde{\xi})/(1 + \tilde{\xi})}$, by virtue of the relation

$$\xi = b(\sigma)\tilde{\xi} + \frac{1}{2} \left[\xi_{LE}(\sigma) + \xi_{TE}(\sigma) \right]$$

Hence, the inner integral may be evaluated exactly by either Equation 52a or Equation 55 if enough points are used in the formulas. For the limits $\nu \leq 5$, $n \leq 4$ used in the computer program, six points are sufficient for Equation 52a, five points for Equation 55. Equation 52a was used with $K = 6$. (This choice was made arbitrarily; it would be just as good to use Equation 55.)

In the integrations over σ , six-point Gaussian integration with weight function 1.0 was used. The basic formula is

$$\int_0^1 f(v) dv = \sum_{\ell=1}^6 h_\ell f(v_\ell) \quad (66)$$

exact for $f(v)$ a polynomial of degree at most 11. The constants occurring in this formula are given in the subroutine FORCE. They may be derived from a table given by Scarborough (Reference 6, p. 148).

In $I_{nm\nu\mu}^{(1)}$, Equation 66 is applied by putting

$$\underline{\sigma} = \underline{\sigma}_{F_L} v$$

The resulting expression is

$$I_{nm\nu\mu}^{(1)} = \frac{\pi}{6} \underline{\sigma}_{F_L} \sum_{\ell=1}^6 h_{\ell} \sqrt{1 - \underline{\sigma}_{\ell}^2} \underline{\sigma}_{\ell}^{m+\mu} \sum_{k=1}^6 \xi_k(\underline{\sigma}_{\ell})^{\nu} (1 - \tilde{\xi}_k^2)^{F_n}(\tilde{\xi}_k)$$

in which

$$\underline{\sigma}_{\ell} = \underline{\sigma}_{F_L} v_{\ell}$$

$$\xi_k(\underline{\sigma}_{\ell}) = b(\underline{\sigma}_{\ell}) \tilde{\xi}_k + \frac{1}{2} \left[\xi_{LE}(\underline{\sigma}_{\ell}) + \xi_{TE}(\underline{\sigma}_{\ell}) \right]$$

In $I_{nm\nu\mu}^{(2)}$, the transformation

$$\underline{\sigma} = 1 - (1 - \underline{\sigma}_{F_L}) v^2$$

was used. This makes the v -integrand behave like a polynomial at the ends of the interval. The resulting formula is

$$I_{nm\nu\mu}^{(2)} = \frac{\pi}{3} (1 - \underline{\sigma}_{F_L}) \sum_{\ell=1}^6 h_{\ell} v_{\ell} \sqrt{1 - \underline{\sigma}_{\ell}^2} \underline{\sigma}_{\ell}^{m+\mu} \sum_{k=1}^6 \xi_k(\underline{\sigma}_{\ell})^{\nu} (1 - \tilde{\xi}_k^2)^{F_n}(\tilde{\xi}_k)$$

in which

$$\underline{\sigma}_{\ell} = 1 - (1 - \underline{\sigma}_{F_L}) v_{\ell}^2$$

$$\xi_k(\underline{\sigma}_{\ell}) = b(\underline{\sigma}_{\ell}) \tilde{\xi}_k + \frac{1}{2} \left[\xi_{LE}(\underline{\sigma}_{\ell}) + \xi_{TE}(\underline{\sigma}_{\ell}) \right]$$

DESCRIPTION OF THE COMPUTER PROGRAM

A functional diagram of the computer program is given in Figure 4. With the exception of two subroutines named MSIMEC and MSIMER, all of the programs are written in Fortran IV. These two subroutines, written in machine language, are used for complex and real matrix inversion, respectively. There are certain limitations related to the various other subprograms, which are listed below.

Subprogram	Limitations
Subroutine Data	
NCC	The number of chordwise collocation stations must be ≤ 10 .
NCS	The number of spanwise collocation stations must be ≤ 9 .
NDATA	The number of sets of data must be ≤ 10 .
N	The number of chordwise pressure modes ≤ 5 .
M	The number of spanwise pressure modes ≤ 5 .
Subroutine Zen	
MODES	This is the number of modes used in the calculation of generalized forces, and must be ≤ 10 .
NPTS	For a planar wing, this is the number of points at which the deflection is given in the horizontal surface and must be ≤ 66 . For a nonplanar wing, there must be ≤ 66 points for the deflections in the horizontal surface and ≤ 66 points for the deflections in the vertical surface.

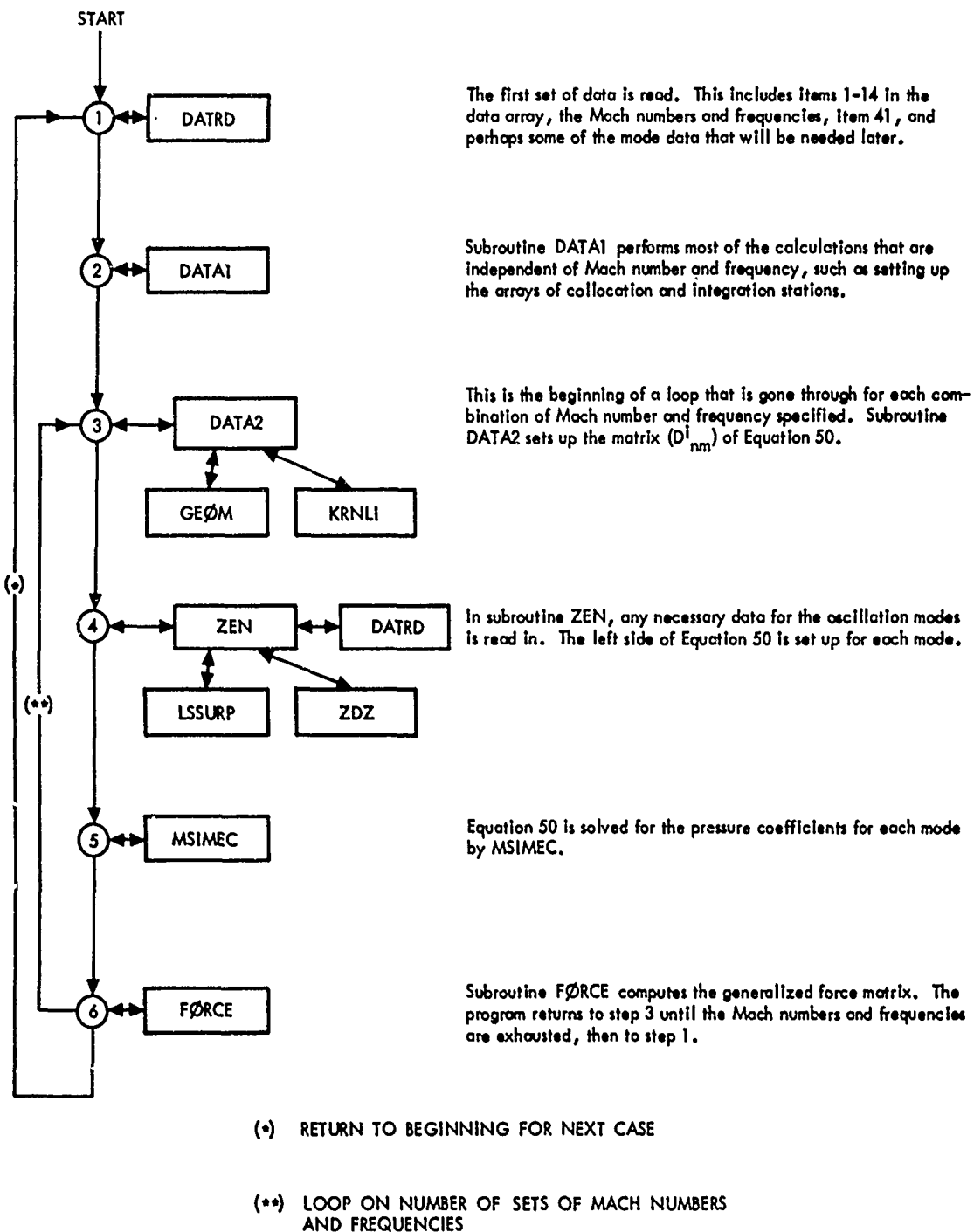


Figure 4. Functional Flow Diagram—Main Program

USE OF THE COMPUTER PROGRAM

The following rules apply to optimum use of the computer program:

1. Before attempting to use the computer program compute the coefficients of polynomials of minimum order in x and y that will adequately represent the modal deflection distributions. Least-square fitted surfaces are very useful for this purpose, since weighting factors may be used to obtain a better fit in special regions.
2. Set the number of chordwise collocation points M equal to one plus the highest of the orders in x ; and, unless there is special reason to reduce the number of chordwise integration points, set L equal to M .
3. Set the number of spanwise collocation points R equal to one plus the highest of the orders in y . This establishes the number of pressure coefficients a_{nm} at $M \times R$. Their values are computed by matching exactly the downwashes at the $M \times R$ spanwise collocation points.

The input data is read by the subroutine DATRD. Use of this subroutine requires that, on each data card, the first 72 columns are six fields of width 12, as indicated on the sample data sheets (Figure 5). The first field contains an integer giving the location in the data array in which the number in the second field is to be stored. The numbers in the remaining fields are stored in consecutive locations. If a field is blank, the corresponding location in the data array is unchanged. DATRD reads any number of cards. A minus sign in column 1 indicates the last card to be read; if this minus sign is not present, DATRD continues with the next card. The storage locations of the data on a card are not affected by the sign in column 1. All floating point numbers must be written with decimal points. All integers must be at the right of their fields.

The data array is set up as follows:

- | | |
|----------|--|
| 1. N | The number of chordwise pressure modes |
| 2. M | The number of spanwise pressure modes |
| 3. NCC | The number of chordwise collocation points |
| 4. NCS | The number of spanwise collocation points |
| 5. NDATA | The number of sets of values of Mach number and frequency to be used. |
| 6. NSYM | Indicator for symmetric (NSYM = +1) or antisymmetric (NSYM = -1) modes of oscillation. |

- | | | |
|--------|---|---|
| 7. | SFOLD | Distance spanwise from wing center line to fold line |
| 8. | S'TIP | Semispan |
| 9. | BO | One-half of the root chord |
| 10. | ALFA1 | The fold angle (in degrees) |
| 11. | λ_{LE} | The sweep angle of the leading edge (in degrees) |
| 12. | λ_{TE} | |
| 13. | λ_{LETIP} | |
| 14. | λ_{TETIP} | |
| 21-30. | Values of Mach number. | NDATA of these must be entered. |
| 31-40. | Values of reduced frequency,
$\omega \cdot BO/U$. | NDATA of these must be entered. |
| 41. | NMOD | The number of modes |
| 42. | JD | Indicator for the type of input data for a mode. If JD=1, the deflections are given at a set of points on the wing (and tip). If JD=2, the coefficients of polynomials for the deflection of the wing and tip are given. If JD=0, the current mode and subsequent modes are not given by data. They are the same as the corresponding modes which were used for the previous frequency and Mach number. |
| 43. | NPTSW | Number of points on the wing at which deflections are given (used if JD=1). |
| 44. | NPTST | Number of points on tip at which deflections are given. |
| 51-71. | Deflection coefficients on the wing. | The coefficients are stored as follows: |

$$a_{00} + a_{10}x + a_{20}x^2 + a_{30}x^3 + a_{40}x^4 + a_{50}x^5 + y(a_{01} + a_{11}x + a_{21}x^2$$

$$+ a_{31}x^3 + a_{41}x^4) + y^2(a_{02} + a_{12}x + a_{22}x^2 + a_{32}x^3)$$

$$+ y^3(a_{03} + a_{13}x + a_{23}x^2) + y^4(a_{04} + a_{14}x) + a_{05}y^5$$

- 76-96. Deflection coefficients on the tip. Same storage rule as above, except all locations are increased by 25.
98. Indicator that no more modes are to be read after the present one.
(For current and subsequent frequencies and Mach numbers.)
- 101-299. Deflection data at points on the wing, in the order $x_1, s_1, n_1, x_2, s_2, n_2$, etc.
- 301-499. Deflection data at points on the tip.

Items 1-41 must be read in the first set of data. There may be an additional set of data for each mode, for each Mach number and frequency case. After the indicator DAT(98) has been given a non-zero value, no more deflection data will be read. After JD has been given the value zero, no data will be read for the higher numbered modes.

The data in Figure 5 is for a 60° triangular wing folded at 75 percent semispan, at an angle of 30° . The root chord is 5.0 feet, making $BO = 2.5$. Three modes: plunge, pitch, and a third nonrigid mode are considered. The Mach number is 0.7, and six frequencies: 10, 20, 30, 40, 50, and 60 cps are used. The speed of sound is taken to be 1000 ft./sec. This gives reduced frequencies of 0.157, 0.314, 0.471, 0.628, 0.785, and 0.942.

Three spanwise and three chordwise pressure modes, six chordwise and eight spanwise collocation stations are specified.

In the set of cards numbered 15 - 30, which give deflection data, some of the cards have been omitted. Otherwise, this is a complete set of data for a computer run.

FORTRAN FIXED IO DIGIT DECIMAL DATA

DECK NO. _____ PROGRAMMER _____ DATE _____ PAGE 1 of 5 JOB NO. _____

NUMBER	IDENTIFICATION	DESCRIPTION	DO NOT KEY PUNCH
1	1		
13		N	
25		H	
37		NCC	
49		NCS	
61		NDATA	
1	2		
13		NSYM	
25		SEOLD	
37		STIP	
49		BO	
61		Fold Angle	
1	3		
13			
25			
37			
49			
61			
1	4		
13			
25			
37			
49			
61			

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Figure 5. Sample Data Sheets (Sheet 1 of 5)

FORTRAN FIXED IO DIGIT DECIMAL DATA

DECK NO. _____		PROGRAMMER _____		DATE _____		PAGE 2 of 5		JOB NO. _____	
NUMBER		IDENTIFICATION		DESCRIPTION		DO NOT KEY PUNCH			
1	2 6	[Shaded Box]		Last Mach Number					
13	0 7								
25									
37									
49									
61		73 80							
1	3 1	[Shaded Box]		Reduced Frequencies					
13	0 1 5 7								
25	0 3 1 4								
37	0 4 7 1								
49	0 6 2 8								
61	0 7 8 5	73 80							
1	3 6	[Shaded Box]							
13	0 9 4 2								
25									
37									
49									
61		73 80							
1	4 1	[Shaded Box]		Minus sign indicates last card in first set of data					
13	3								
25									
37									
49									
61		73 80							
FORM 114-C-17 REV. 7-55-VELLUM		8							

Figure 5. Sample Data Sheets (Sheet 2 of 5)

FORTRAN FIXED IO DIGIT DECIMAL DATA

DECK NO. _____ PROGRAMMER _____		DATE _____	PAGE 3 of 5	JOB NO. _____
NUMBER	IDENTIFICATION	DESCRIPTION	DO NOT KEY PUNCH	
1				
13				
25				
37				
49				
61				
1				
13				
25				
37				
49				
61				
1				
13				
25				
37				
49				
61				
1				
13				
25				
37				
49				
61				
1				
13				
25				
37				
49				
61				
1				
13				
25				
37				
49				
61				
1				
13				
25				
37				
49				
61				

Figure 5. Sample Data Sheets (Sheet 3 of 5)

FORTRAN FIXED IO DIGIT DECIMAL DATA

DECK NO. _____ PROGRAMMER _____ DATE _____ PAGE 4 of 5 JOB NO. _____

NUMBER	IDENTIFICATION	DESCRIPTION DO NOT KEY PUNCH
1	73. 80	
13		
25		
37		
49		
61		
1	73. 80	
13		
25		
37		
49		
61		
1	73. 80	JD
13		
25		NPTSW
37		
49		NPTST
61		
1	73. 80	
13		
25		
37		
49		
61		
1	73. 80	Deflection data on the wing begin.
13		X1
25		S1
37		N1
49		X2
61		S2
1	73. 80	
13		
25		
37		
49		
61		
1	73. 80	
13		
25		
37		
49		
61		

Figure 5. Sample Data Sheets (Sheet 4 of 5)

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FORTRAN FIXED 10 DIGIT DECIMAL DATA

DECK NO. _____ PROGRAMMER _____ DATE _____ PAGE 5 of 5 JOB NO. _____

NUMBER	IDENTIFICATION	DESCRIPTION DO NOT KEY PUNCH
1	1.6 6	End of deflection data on the wing
13	1 2.0 6	N22
25		
37		
49		
61	73. 80	
	2.6	
1	3.0 1	Deflection data on the tip begin
13	4. 2 5	X1
25	2. 3	S1
37	1. 3 8	N1
49	4. 5	X2
61	2. 3	S2
	2.7	
1	3.1 6	End of deflection data on the tip
13	4. 8	X6
25	2. 7	S6
37	1. 4 4	N6
49		
61	73. 80	
	3.0	
1	9.8	Indicator that modes are not to be read in hereafter
13	1	
25		
37		
49		
61	73. 80	
	3.1	

FORM 114-C-17 REV. 7-53-YELLUM

Figure 5. Sample Data Sheets (Sheet 5 of 5)

VII. RESULTS

The computer program was applied to the rigid modes of two wings.

ASPECT RATIO 2.0 RECTANGULAR WING FOLDED AT 80 PERCENT SEMISPAN

$C_{L\alpha}$ is plotted as a function of Mach number in Figure 6. The dashed curve is a plot of the approximation

$$C_{L\alpha} = \frac{2\pi A.R.}{2 + \sqrt{4 + (A.R.)^2 (1-M^2)}}$$

where A.R. is the aspect ratio. This is formula (6-31) of Reference 5, for the case of a rectangular wing.

TRIANGULAR WING WITH FOLDED TIPS

The configuration used was a triangular wing with a sweep angle of 65 degrees, folded at 60 percent semispan.

Figures 7 and 8 show $C_{L\alpha}$ and $C_{m\alpha}$ as functions of fold angle (at $M = 0.8$), and as functions of Mach number. Figure 9 is a plot of unsteady generalized forces for rigid oscillations of the wing in the pitching mode.

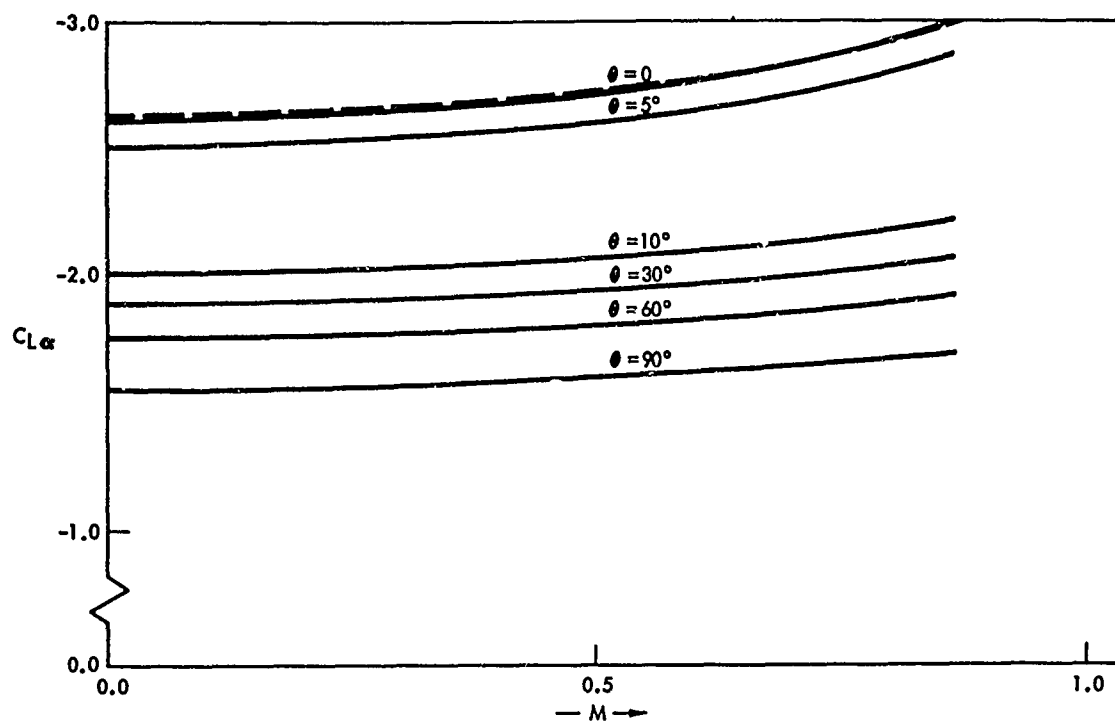


Figure 6. Lift and Moment Coefficients Vs Mach Number for Aspect Ratio 2.0 Rectangular Wing at Various Fold Angles

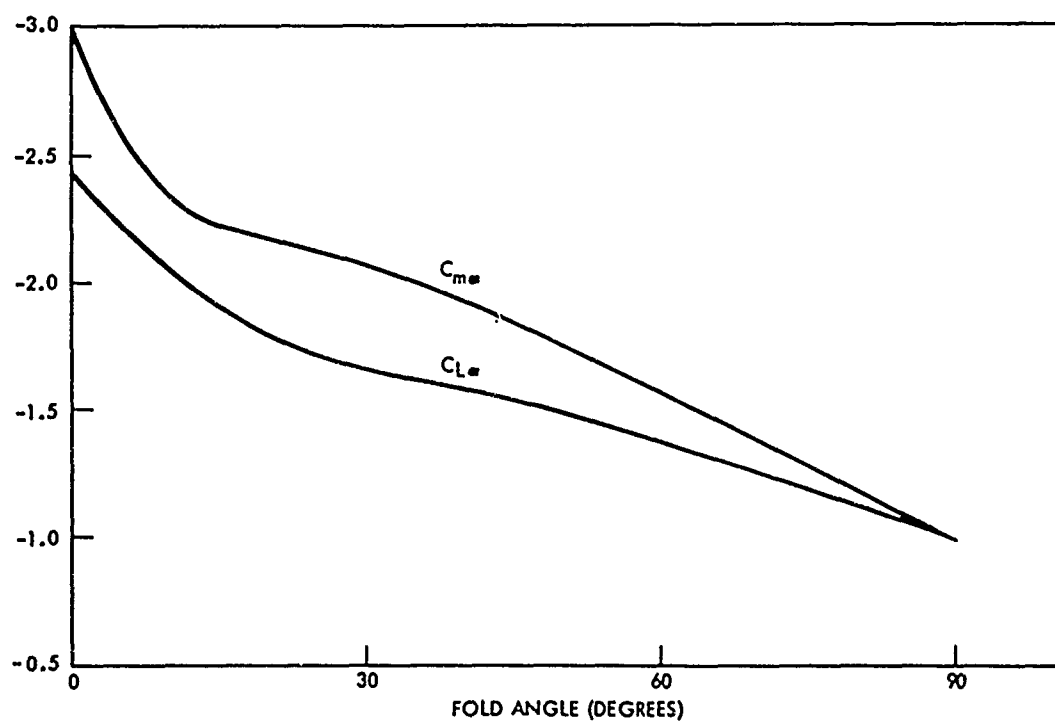


Figure 7. Lift and Moment Coefficients Vs Fold Angle for 65° Triangular Wing at $M_\infty = 0.8$

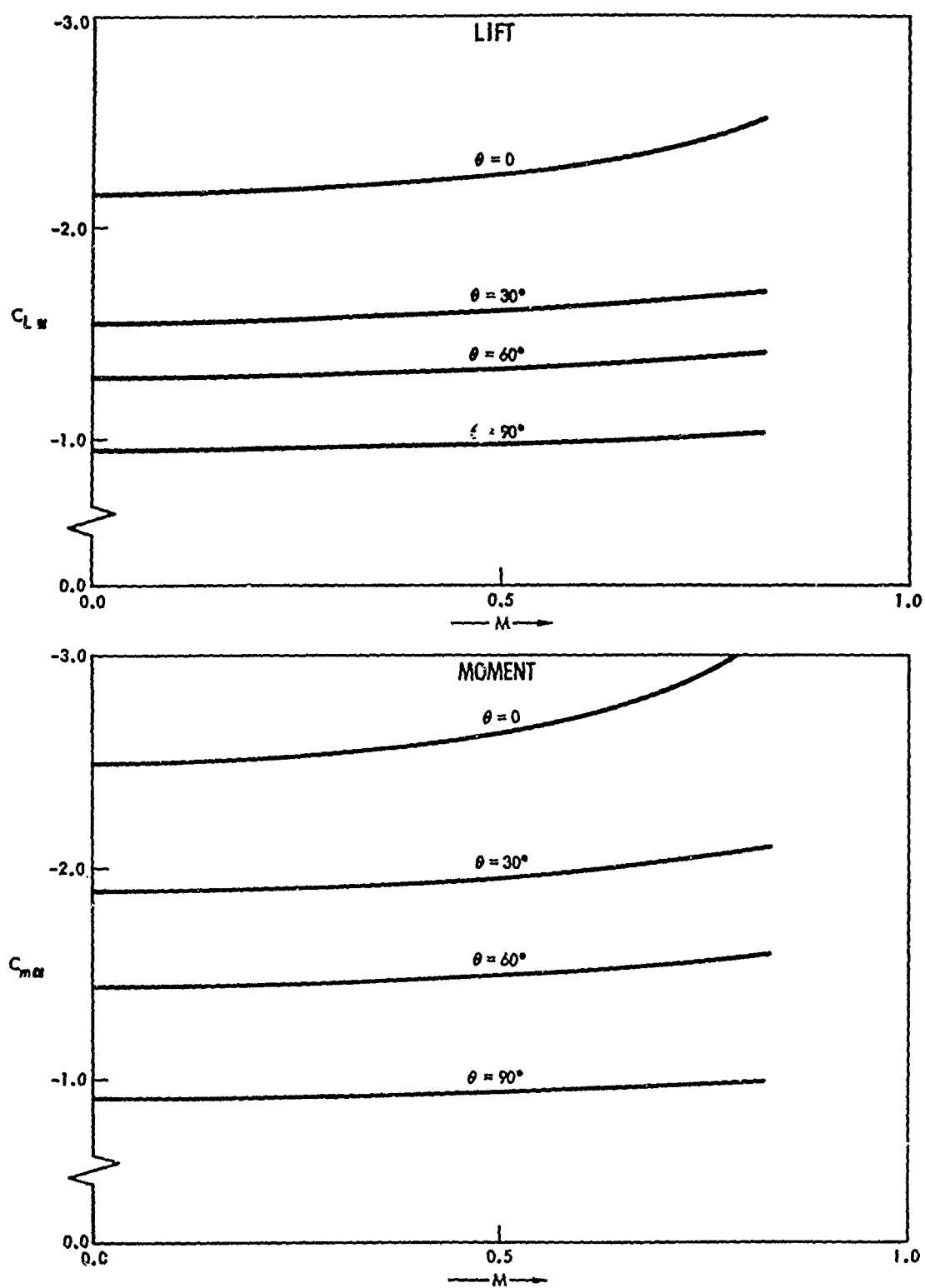
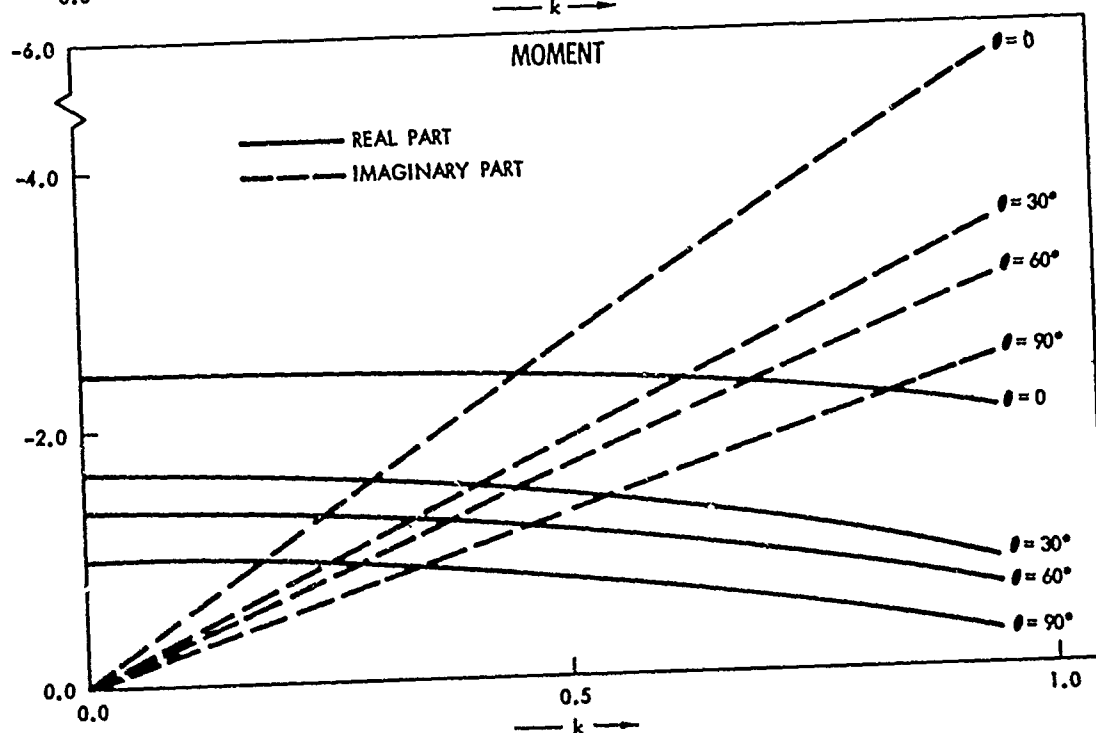
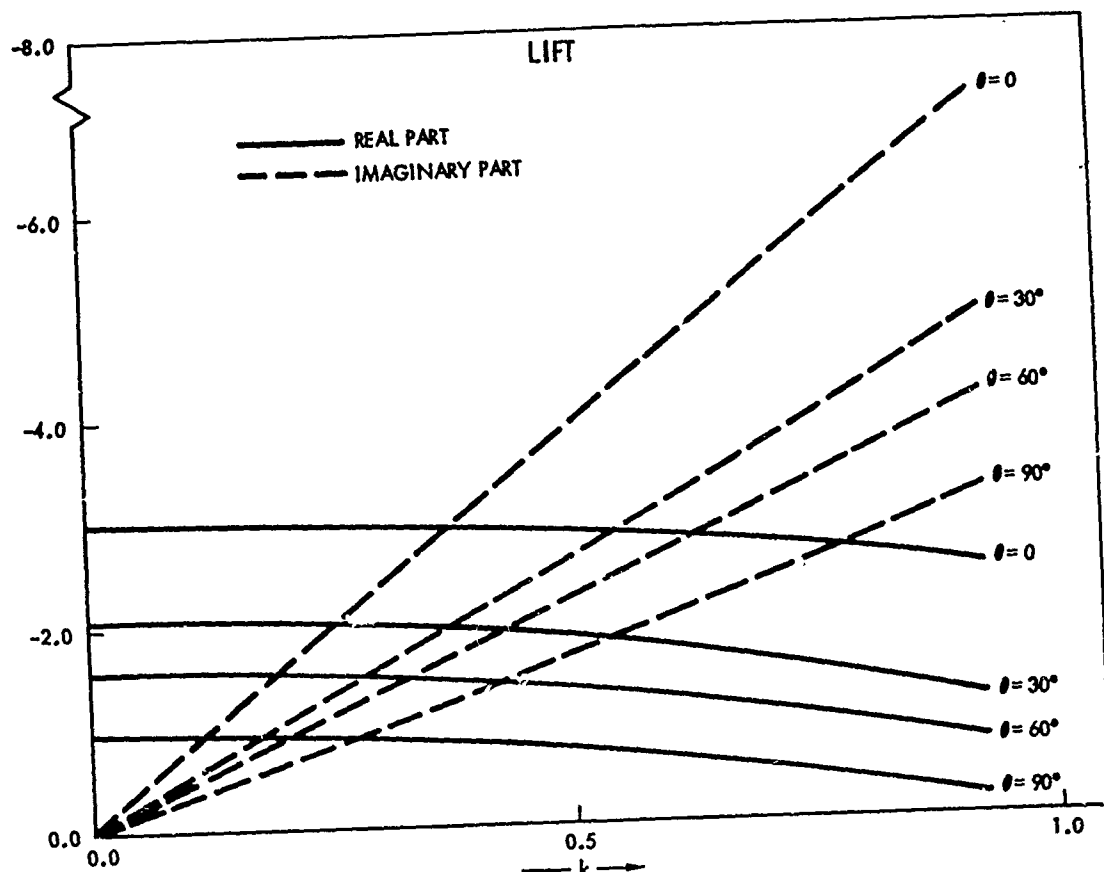


Figure 8. Lift and Moment Coefficients Vs Mach Number for 65° Triangular Wing at Various Fold Angles



Coefficients of Lift and Moment Due to Pitch Vs Reduced Frequency for a
65° Triangular Wing at $M_\infty = 0.8$ and Various Fold Angles

VIII. CONCLUSIONS AND RECOMMENDATIONS

The results obtained by the nonplanar kernel function method show the expected trend with increasing fold angle, which agrees with the observed experimental trend. For zero fold angle, the method reduces to that already used by Hsu (Reference 4).

Possible extensions of the method include the treatment of more general configurations, or of other specific configurations, such as the T-tail. Also, any generalizations proposed for the planar case should be considered here, such as the problem of a nonplanar wing with a control surface.

The formula that was used for the kernel function (page 19) should be useful in any future developments using the three-dimensional kernel function. The previously available formula was much longer. A special case of that formula is given in Reference 12. The use of the simplified kernel function, together with Hsu's method of integration, results in greatly reduced computer running times. This makes it practical to use the kernel function method as a tool in the preliminary analysis of new wing configurations.

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APPENDIX. COMPUTER PROGRAM LISTINGS

MP

```

COMMON W(50,25),WI(50,25),A(2,25,25),B(2,25,10),XSV(10,5)
COMMON X(10),Y(10),PSI(10),ETA(10),ZXY(21,10),ZZZ(21,10)
COMMON /C0M1/ N,M,NR,NC,NDATA,NM0D,NSYM,NDEX,NCC,NCS1,NCST
COMMON /C0M1/FM,FR,BO,SF,ST
COMMON /CDAT/DAT(500)
DO 1 I=1,100
1 DAT(I)=0.0
2 CALL DATRD(DAT)
CALL DATA
DO 200 ND=1,NDATA
FR=DAT(ND+30)
FM=DAT(ND+20)
CALL DATA2
CALL ZEN
DO 35 I=1,NC
DO 35 J=1,NC
A(1,I,J)=0.0
A(2,I,J)=0.0
DO 35 K=1,NR
A(1,I,J)=A(1,I,J)+W(K,I)*W(K,J)+WI(K,I)*WI(K,J)
A(2,I,J)=A(2,I,J)+W(K,I)*W(K,J)-WI(K,I)*W(K,J)
35 CONTINUE
K=MSIMEC(25,NC,NM0D,A,B)
IF (K.NE.1) GO TO 99
WRITE (6,650)
650 FORMAT(1H18X,78HPRESSURE = STIP/B(S) * SUM OF A(N,M)*F(N,PSI)*SQRT00000270
1(1.0-(S/STIP)**2)*(S/STIP)**)
IF (NSYM.LT.0) GO TO 10
WRITE (6,651)
651 FORMAT(1H+86X,5H(M-1))
GO TO 15
10 WRITE (6,652)
652 FORMAT(1H+86X,1HM)
15 LINES=1
LINLIM=35-NC/2
DO 40 I1=1,NM0D
WRITE (6,700) I1

```

00000020
00000030
00000040
00000050
00000060
00000070
00000080
00000090
00000100
00000110
00000120
00000130
00000140
00000150
00000160
00000170
00000180
00000190
00000200
00000210
00000220
00000230
00000240
00000250
00000260
00000270
00000280
00000290
00000300
00000310
00000320
00000330
00000340
00000350
00000360
00000370
00000380

MP

```

700 FORMAT(1H-20X,41HPRESSURE COEFFICIENTS A(N,M) FOR MODE NO.12/1H0
1 2(9X,5HN M6X,9HREAL PART7X,10HIMAG PART )/1H )
J1=1
JCARR=1
DO 19 NPR=1,N
DO 19 MPR=1,M
IF (JCARR.EQ.0) GO TO 17
JCARR=0
WRITE (6,705) NPR,MPR,B(1,J1,I1),B(2,J1,I1)
GO TO 19
17 JCARR=1
WRITE (6,706) NPR,MPR,B(1,J1,I1),B(2,J1,I1)
705 FORMAT(11I,14,1P2E16.5)
706 FORMAT(1H+52X,2I4,1P2E16.5)
19 J1=J1+1
LINES=LINES+6*NC/2
IF (LINES.LE.LINLIM) GO TO 40
WRITE (6,708)
708 FORMAT(1H1)
LINES=0
40 CONTINUE
IF (LINES.EQ.0) GO TO 24
LINES=LINES+7+2*NM0D**2
IF (LINES.LE.40) GO TO 23
WRITE (6,708)
GO TO 24
23 WRITE (6,709)
709 FORMAT(1H0)
24 WRITE (6,45) FR,FM,DAT(10)
45 FORMAT(1H010X,26HGENERALIZED FORCES -- K =F7.4,6H, M =F7.4,
1 15H, FOLD ANGLE =F8.4,5H DEG.)
WRITE (6,46)
46 FORMAT(1H05X,5HMGDES/4X,11H0SC. DEFL.8X,9HREAL PART10X,9HIMAG PAR00000710
1T10X,10HABS. VALUE6X,11HPHASE ANGLE)
CALL FORCE
200 CONTINUE
GO TO 2

```

MP

99 WRITE (6,98), ND
98 FORMAT(57H1 PRESSURE COEFFICIENTS CANNOT BE FOUND FOR DATA CASE
1.12)
GO TO 200
END

00000760
0000000770
00000780
00000790
00000800

DATA1

SUBROUTINE DATA1

```

COMMON W(50,25),WI(50,25),A(2,25,25),B(2,25,10),XSV(10,5)
COMMON X(10),Y(10),PSI(10),ETA(10),ZXY(21,10),ZZZ(21,10)
COMMON /C0M1/ N,M,NR,NC,NDATA,NM0D,NSYM,NDEX,NCC,NCS1,NCST
COMMON /C0M1/ FM,FR,BO,SF0LD,STIP
COMMON /C0M2/ Q(5),DELP(5,10),NIC,NIS,ZIS,ZIP,ZCP,SF0ST
COMMON /C0M2/ BF,BMF,C0S1,SIN1,CT32P,CT32M,CT54P,CT54M
COMMON /CDAT/DAT(500)
COMMON /C0NST/ PI,PI0VR2,RAD,ZI(10)

```

N=DAT(1)

M=DAT(2)

NCC=DAT(3)

NCS=DAT(4)

NDATA=DAT(5)

NSYM=DAT(6)

```

NCS = N0. 0F COLLOCATION PTS. SPANWISE
NCC = N0. 0F COLLOCATION PTS. CHORDWISE
NIS = N0. 0F INTEGRATION PTS. SPANWISE
NIC = N0. 0F INTEGRATION PTS. CHORDWISE
NDEX = (1-NSYM)/2

```

NIC = NCC

NIS = NCS + 1

ZIS = NIS

BO=DAT(9)

SF0LD=DAT(7)/BO

STIP=DAT(8)/BO

DO 15 I=1,10

X(I)=D0TS

Y(I)=D0TS

PSI(I)=D0TS

ETA(I)=D0TS

SB02=STIP**2

WRITE (6,20) NCS,NCC,M,N,NDATA,NDATA

```

20  FORMAT(1H1,54X,10HINPUT DATA/21X,16HCOLLOCATION PTS., 5X,14HPRESSU00001150
1RE MODES, 5X, 9HMACH NOS., 5X, 6HFREQS./21X,12,5H SPAN,13,6H CHORD00001160
2,4X,12,5H SPAN,13,6H CHORD,9X,12,11X,12)
13  SF0ST = SF0LD/STIP

```

```

00000820
00000830
00000840
00000850
00000860
00000870
00000880
00000890
00000900
00000910
00000920
00000930
00000940
00000950
00000960
00000970
00000980
00000990
00001000
00001010
00001020
00001030
00001040
00001050
00001060
00001070
00001080
00001090
00001100
00001110
00001120
00001130
00001140
00001150
00001160
00001170
00001180

```

DATA1

ALFA1=DAT(10)/RAD
 ALFA2=(90.0-DAT(11))/RAD
 ALFA3=(90.0-DAT(12))/RAD
 ALFA4=(90.0-DAT(13))/RAD
 ALFA5=(90.0-DAT(14))/RAD
 EVALUATE TIP INTEGRALS
 F= SQRT(1.0-SF0ST**2)
 THE1=ATAN(SF0ST/F)
 Q(1)=1.0-(THE1+F*SF0ST)/PI0VR2
 F=F**3/PI
 Q(2)=F/3.0
 DO 1607 IL=3,5
 F=F*SF0ST
 Q(IL)=(F+ZI(IL-2)*Q(IL-2))/ZI(IL+1)

C

1607 CONTINUE
 C GENERATE COLLOCATION AND INTEGRATION STATIONS

ZCP=2.0*PI/FL0AT(2*NCC+1)

ZIR=S802*ZCP/8.0

F=0.5*ZCP

DO 27 J=1,NCC

PSI(J)=-COS(F)

W=PSI(J)

DO 26 I=1,N

26 DELP(I,J)=FXI(W,I)*ZIR

I=NCC+1-J

X(I)=-W

27 F=F+ZCP

ZS=PI/ZIS

F=-ZS+PI0VR2

DF=0.5*ZS

DO 29 J=1,NCS

Y(J)=SIN(F)

W=F+DF

ETA(J)=SIN(W)

29 F=F-ZS

ETA(NIS)=-ETA(1)

WRITE (6,35)

00001190
 00001200
 00001210
 00001220
 00001230
 00001235
 00001240
 00001250
 00001260
 00001270
 00001280
 00001290
 00001300
 00001310
 00001320
 00001330
 00001340
 00001350
 00001360
 00001370
 00001380
 00001390
 00001400
 00001410
 00001420
 00001430
 00001440
 00001450
 00001460
 00001470
 00001480
 00001490
 00001500
 00001510
 00001520
 00001530
 00001540

DATA1

```

35  FORMAT(1H0,27X,47H,CALCULATED COLLOCATION AND INTEGRATION STATIONS/00001550
    1H ,24X,1HX,15X,3HPSI,15X,1HY,15X,3HETA)
    NLIN=MAX0(NIS,NCC)
    WRITE (6,36) (X(I),PSI(I),Y(I),ETA(I),I=1,NLIN)
36  FORMAT(13X,4FI7.4)
    WRITE (6,560)
    WRITE (6,565) (DAT(I+20),DAT(I+30),I=1,NDATA)
    C
    EVALUATE QUANTITIES FOR LATER USE
    C0S1=C0S(ALFA1)
    SIN1=SIN(ALFA1)
    NCS1=(NCS+1)/2
    IF (NSYM.LT.0) NCS1=NCS/2
    CTN2=0.5*C0S(ALFA2)/SIN(ALFA2)
    CTN3=0.5*C0S(ALFA3)/SIN(ALFA3)
    CTN4=0.5*C0S(ALFA4)/SIN(ALFA4)
    CTN5=0.5*C0S(ALFA5)/SIN(ALFA5)
    CT32P=CTN3+CTN2
    CT32M=CTN3-CTN2
    CT54P=CTN5+CTN4
    CT54M=CTN5-CTN4
    BF =1.0+SF0LD*CT32M
    BMF=1.0+SF0LD*CT32P
    NR=NCC*NCS1
    NC=M*N
    DO 82 J=1,10
    IF (Y(J).LE.SF0ST) GO TO 84
82  CONTINUE
84  NCST=J-1
    ZIP=16.0/SB02
    RETURN
560  FORMAT(1H0,35X, 9HMACH NOS. ,15X,6HFREQS.)
565  FORMAT(1H ,36X,F8.4,14X,F8.4)
    END

```

DAT2

```

SUBROUTINE DATA2
COMMON WASH(50,25),WASHI(50,25),A(2,25,25),B(2,25,10),XSV(10,5)
COMMON X(10),Y(10),PSI(10),ETA(10),ZXY(21,10),ZZ(21,10)
COMMON /COM1/ N,M,NR,NC,NDATA,NMOD,NSYM,NDEX,NCC,NCS1,NCST
COMMON /COM1/ FM,FR,BO,SFOLD,STIP
COMMON /COM2/ Q(5),DELP(5,10),NIC,NIS,ZIS,ZIP,ZCP,SFQST
COMMON /COM2/ BF,BMF,      C0S1,SIN1,CT32P,CT32M,CT54P,CT54M
COMMON /CDAT/DAT(500)
COMMON /CONST/ PI,PIQVR2,RAD,ZI(10)
DIMENSION RA(10),RB(10),RC(5),FS(5),FT(5),FU(5),FV(5),FW(5),FX(5)
DIMENSION XU(2),XX(2),XY(2),XZ(2)
EVALUATION OF MATRIX OF INTEGRALS
DO 38 J=1,1250
  WASH(J,1) = 0.0
  WASHI(J,1) = 0.0
  BETA2=1.0-FM**2
  IU = 0
  DO 1000 IQ=1,NCS1
    CALL GEOM(Y(IQ),XI1,XI2,C02,SI2,YY,ZY)
    DO 1000 IP = 1,NCC
      XA=X(IP)*XI2+XI1
      XSV(IP,IQ) = XA
      IU = IU+1
      XARG=ZI(IP)*ZCP
      F=SQRT(1.0-X(IP)**2)
      RC(1)=2.0*(F+XARG)
      G=-2.0*X(IP)*F
      RC(2)=XARG-0.5*G
      DO 1604 IC=3,N
        H=-F-2.0*X(IP)*G
        RC(IC)=F/ZI(IC-2)-H/ZI(IC)
        F=G
      G=H
1604  DO 1640 IN = 1,NIS
58      YE=ETA(IN)
      CALL GEOM(YE,XI3,XI4,C04,SI4,YETA,ZETA)
      SY=1.0

```

DAT2

IF (YE.GE.0.0) GO TO 54
SY=SIGN(1.0,NSYM)
YE=SIGN(YE,NSYM)

54 CONTINUE

YO=YY-YETA
ZOO=ZY-ZETA
DO 1580 IJ=1,N
FS(IJ) = 0.0
FT(IJ) = 0.0
FU(IJ) = 0.0

1580 FV(IJ)=0.0

DO 1600 IL = 1,NIC
XB=PSI(IL)*XI4+XI3

XO = XA - XB

51 CALL KRNL1(FR,XO,YO,ZOO,FM,BETA2,C02,SI2,C04,SI4,XX)
IF (SI4.EQ.0.0) GO TO 93

C ETA IS OUTBOARD OF FOLDLINE,CALCULATE KPRIME.

96 IF(ETA(IN))I01,I01,I02

101 SA=YY+SFOLD

GO TO 103

102 SA=YY-SFOLD

103 X1=XA-PSI(IL)*BF-BMF

CALL KRNL1(FR,X1,SA,ZY,FM,BETA2,C02,SI2,C04,SI4,XY)

CALL KRNL1(FR,X1,SA,ZY,FM,BETA2,C02,SI2,1.0,0.0,XZ)

C KPRIME = XX - XY + XZ

XU(1) = XY(1) - XZ(1)

XU(2) = XY(2) - XZ(2)

XX(1) = XX(1) - XU(1)

XX(2) = XX(2) - XU(2)

GO TO 94

93 XU(1) = 0.0

XU(2) = 0.0

94 DO 1590 IJ=1,N

FS(IJ)=FS(IJ)+XX(1)*DELP(IJ,IL)

FT(IJ)=FT(IJ)+XX(2)*DELP(IJ,IL)

FU(IJ)=FU(IJ)+XU(1)*DELP(IJ,IL)

1590 FV(IJ)=FV(IJ)+XU(2)*DELP(IJ,IL)

00002250
00002260
00002270
00002280
00002290
00002300
00002310
00002320
00002330
00002340
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00002380
00002390
00002400
00002410
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00002430
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00002450
00002460
00002470
00002480
00002490
00002500
00002510
00002520
00002530
00002540
00002550
00002560
00002570
00002580
00002590
00002600
00002610

DAT2

1600 CONTINUE
C CHORDWISE INTEGRATION IS COMPLETE.

IA = 0
W=(1.0-YE**2)/ZIS
IF (NDEX.NE.0) W=W*YE
DO 1625 IL=1,M
ILN=IL+NDEX
DO 1620 IM = 1,N
IA = 1 + IA
WASH(IU,IA)=WASH(IU,IA)+FS(IM)*W
WASHI(IU,IA)=WASHI(IU,IA)+FT(IM)*W
IF (IN.NE.1.AND .IN.NE.NIS) GO TO 1620
WASH(IU,IA)=WASH(IU,IA)+FU(IM)*Q(ILN)*SY
WASHI(IU,IA)=WASHI(IU,IA)+FV(IM)*Q(ILN)*SY

1620 CONTINUE

1625 W=YE*W

1640 CONTINUE

DO 1601 LG=1,IP
ARG=FR*(XA-PSI(LG)*XI2-XI1)
RA(LG)=ZIP*(COS(ARG)-1.0)
RB(LG)=-ZIP*SIN(ARG)

1601 CONTINUE

C RA = REAL, RB = IMAG

DO 1646 IJ=1,N

FW(IJ) = 0.0

FX(IJ) = 0.0

DO 1646 IK=1,IP

FW(IJ) = FW(IJ) + DELP(IJ,IK)*RA(IK)

FX(IJ) = FX(IJ) + DELP(IJ,IK)*RB(IK)

1646 CONTINUE

IA = 0

YP=ZIS/8.0

IF (NDEX.NE.0) YP=YP*Y(IQ)

DO 1649 IL = 1,M

DO 1648 IM = 1,N

IA = 1+IA

WASH(IU,IA) = WASH(IU,IA) - (FW(IM)+RC(IM))*YP

00002620
00002630
00002640
00002650
00002660
00002670
00002680
00002690
00002700
00002710
00002720
00002730
00002740
00002750
00002760
00002770
00002780
00002790
00002800
00002810
00002820
00002830
00002840
00002850
00002860
00002870
00002880
00002890
00002900
00002910
00002920
00002930
00002940
00002950
00002960
00002970
00002980

DATA2

```

1648 WASHI(IU,IA) = WASHI(IU,IA)-FX(IM)*YP
1649 YP=YP*Y(IQ)
1000 CONTINUE
1010 WRITE (6,530) FR,FM
      DO 1015 K = 1,IU
      WRITE (6,555)
      WRITE (6,535)(WASH(K,L),WASHI(K,L),L=1,IA)
1015 CONTINUE
      RETURN
530 FORMAT(1H121X,30HMATRIX FOR REDUCED FREQUENCY =F7.4,
      1 16H, MACH NUMBER =F7.4)
535 FORMAT(3(5X,1P2E15.4,1H))
555 FORMAT(1H )
      END

```

00002990
00003000
00003010
00003020
00003030
00003040
00003050
00003060
00003080
00003090
00003100
00003110
00003120
00003130

GEOM

```

SUBROUTINE GEOM(Y,X1,X2,C0,SI,YY,ZY)
COMMON /C0M1/ N,M,NR,NC,NDATA,NM0D,NSYM,NDEX,NCC,NCS1,NCST
COMMON /C0M1/FM,FR,BO,SF0LD,STIP
COMMON /C0M2/ Q(5),DELP(5,10),NIC,NIS,Z'IS,ZIP,ZCP,SF0ST
COMMON /C0M2/BF,BMF, C0S1,SIN1,CT32P,CT32M,CT54P,CT54M
EVALUATION OF GEOMETRICAL QUANTITIES AT A GIVEN CHORD
Y1=STIP*ABS(Y)
F=Y1-SF0LD
IF (F) 2,2,4
2 C0=1.0
SI=0.0
YY=Y1
ZY=0.0
X1=1.0+Y1*CT32P
X2=1.0+Y1*CT32M
GO TO 6
4 C0=C0S1
SI=SIN1
YY=SF0LD+F*C0S1
ZY=F*SIN1
X1=BMF+F*CT54P
X2=BF+F*CT54M
6 IF (Y.GE.0.0) RETURN
SI=-SI
YY=-YY
RETURN
END

```

00003150
00003160
00003170
00003180
00003190
00003195
00003200
00003210
00003220
00003230
00003240
00003250
00003260
00003270
00003280
00003290
00003300
00003310
00003320
00003330
00003340
00003350
00003360
00003370
00003380
00003390
00003400

KRNL

```

SUBROUTINE KRNL(C,X,Y,ZO,CM,B2,C01,S11,C02,S12,AKERN)
DIMENSION Z(6),H(6),AKERN(2)
DATA H/0.08566225,0.18038079,2*0.23395697,0.18038079,0.08566225/
DATA Z/0.96623476,0.83060469,0.61930959,0.38069041,0.16939531,
10.03376524/
R2 = Y*Y+Z0*Z0
R = SQRT(R2)
CK1 = CK*R
G1 = 0.0
G2 = 0.0
G3 = 0.0
G4 = 0.0
G5 = 0.0
G6 = 0.0
S2 = X*X + B2*R2
S = SQRT(S2)
U1 = (CM*S-X)/(B2*R)
UK = CK1*U1
DO 20 I = 1,6
UZ = U1*Z(I)
UZ2 = UZ**2
F = UK*Z(I)
C0 = COS(F)
S1 = SIN(F)
F = H(I)/SQRT(1.0+UZ2)*UZ*U1
G3 = G3 + F*C0
G4 = G4 - F*S1
F = UZ*F
G5 = G5 + F*C0
G6 = G6 - F*S1
V = 1.0 - Z(I)**2
F = H(I)*2.0*V*EXP(-CK1*V)/SQRT(1.0+V)
G1 = G1 + F
G2 = G2 + V*F
G7 = G1 + G3
XS = X/S
C0 = COS(UK)*XS

```

5

20

KRNL

```

SI = SIN(UK)*XS
IF(CK.NE.0.0) GO TO 22
F14 = 1.0
F15 = 2.0
GO TO 23
22 F14 = CK1*BESEL(CK1,1,3)
    F15 = CK1*CK1*BESEL(CK1,2,3)
23 F11 = -(CK1*G4+F14+C0)
    F12 = CK1*G7+SI
    F = 2.0+B2*R2/S2
    FP = UK+CK*CM*R2/S
    F21 = CK1*G4+CK1*CK1*G5+F15+C0*F +SI*FP
    F22 = CK1*(-G7+CK1*(G6-G2))+C0*FP-SI*F
    XK = CK*X
    C0 = C0S(XK)
    SI = SIN(XK)
    F=(C01*C02+SI1*SI2)/R2
    FP=(Z0*C01-Y*SI1)*(Z0*C02-Y*SI2)/R2**2
    G1 = F*F11+FP*F21
    G2 = F*F12+FP*F22
    AKERN(1) = -C0*G1-SI*G2
    AKERN(2) = SI*G1-C0*G2
    RETURN
    END
00003790
00003800
00003810
00003820
00003830
00003840
00003850
00003860
00003870
00003880
00003890
00003900
00003910
00003920
00003930
00003940
00003950
00003960
00003970
00003980
00003990
00004000
00004010
00004020

```

ZEN

```

SUBROUTINE ZEN
COMMON W(50,25),WI(50,25),A(2,25,25),B(2,25,10),XSV(10,5)
COMMON X(10),Y(10),PSI(10),ETA(10),ZXY(21,10),ZZZ(21,10)
COMMON /COM1/ N,M,NR,NC,NDATA,NMOD,NSYM,NDEX,NCC,NCS1,NCST
COMMON /COM1/ FM,FR,BO,SF,ST
COMMON /CDAT/DAT(500)
DIMENSION XH(66),SH(66),ZH(66),XCH(50),YCH(50),Z(50),DZ(50)
DIMENSION WD(2,100)
EQUIVALENCE (XH,A),(SH,A(67)),(ZH,A(133)),(XCH,A(200)),(Z,A(250))
EQUIVALENCE (YCH,A(300)),(DZ,A(350)),(WD,A(400))
LINES=50
NMOD=DAT(41)
IDEF=0
DO 50 I=1,NMOD
IF (DAT(98).NE.0.0) GO TO 20
IF (IDEF.NE.0) GO TO 20
CALL DATRD(DAT)
JD=DAT(42)
IF (JD-1) 2,4,16
2 IDEF=1
GO TO 20
4 READ DEFLECTION VALUES AT POINTS
L=101
DO 8 J=1,NPTS
XH(J)=DAT(L)
SH(J)=DAT(L+1)
ZH(J)=DAT(L+2)
8 L=L+3
NDH=5
9 CALL LSSURP(XH,SH,ZH,NPTS,NDH,G,ZXY(1,1),IND)
IF (IND.EQ.0) GO TO 10
NDH=NDH-1
GO TO 9
10 NPTS=DAT(44)
L=301
DO 11 J=1,NPTS

```

ZEN

XH(J)=DAT(L)
SH(J)=DAT(L+1)
ZH(J)=DAT(L+2)

11 L=L+3
NDV=5

12 CALL LSSURP(XH,SH,ZH,NPTS,NDV,G,ZZZ(1,I),IND)
IF (IND.EQ.0) GO TO 20
NDV=NDV-1
GO TO 12

C REA: POLYNOMIAL COEFFICIENTS

16 DO 17 J=1,21
ZXY(J,I)=DAT(J+50)
17 ZZZ(J,I)=DAT(J+75)

C FIND UPWASHES
20 K=0

DO 24 I2=1,NCST
DO 24 J=1,NCC
K=K+1

XCH(K)=80*XSX(J,I2)
24 YCH(K)=DAT(8)*Y(I2)
LIND=8+K

IF (LINES+LIND.LE.40) GO TO 25
WRITE (6,41)

41 FORMAT(1H1)

LINES=0

25 LINES=LINES+LIND

WRITE (6,43)

WRITE(6,42) I,(ZZZ(J,I),J=1,21)
42 FORMAT(1H+5X,26HDEFLECTION COEFFICIENTS ON7X,8HMODE NO.12/1H /
1 (1P7E14.4))

43 FORMAT(1H032X,4HTIP,)

CALL ZDZ(XCH,YCH,K,ZZZ(1,I),Z,DZ)
DO 26 I2=1,K

WD(1,I2)=DZ(I2)*80
26 WD(2,I2)=FR*Z(I2)

I3=NCST+1
L=0

00004410
00004420
00004430
00004440
00004450
00004460
00004470
00004480
00004490
00004495
00004500
00004510
00004520
00004525
00004530
00004540
00004550
00004560
00004570
00004580
00004590
00004600
00004610
00004620
00004630
00004640
00004650
00004660
00004670
00004680
00004690
00004700
00004710
00004720
00004730
00004740
00004750

ZEN

```

D0 27 I2=I3,NCS1
D0 27 J=1,NCC
L=L+1
XCH(L)=80*XS(V(J,I2))
27 YCH(L)=DAT(8)*Y(I2)
LIND=8+L
IF (LINES+LIND.LE.40) GO TO 28
WRITE (6,41)
LINES=0
28 LINES=LINES+LIND
WRITE (6,44)
44 FORMAT(1H032X,5HWINING,)
WRITE (6,42) I,(ZXY(J,I),J=1,21)
CALL ZDZ(XCH,YCH,L,ZXY(1,I),Z,DZ)
J2=K+1
D0 29 I2=1,L
WD(1,J2)=DZ(I2)*80
WD(2,J2)=FR*Z(I2)
29 J2=J2+1
D0 33 J=1,NC
B(1,J,I)=0.0
B(2,J,I)=0.0
D0 33 K=1,NR
B(1,J,I)=B(1,J,I)+W(K,J)*WD(1,K)+WI(K,J)*WD(2,K)
33 B(2,J,I)=B(2,J,I)+W(K,J)*WD(2,K)-WI(K,J)*WD(1,K)
50 CONTINUE
RETURN
END

```

00004760
00004770
00004780
00004790
00004800
00004810
00004820
00004830
00004840
00004850
00004860
00004870
00004880
00004890
00004900
00004910
00004920
00004930
00004940
00004950
00004960
00004970
00004980
00004990
00005000
00005010
00005020
00005030

F0RC

SUBROUTINE F0RC

```

COMMON W(50,25),WI(50,25),A(2,25,25),AP(2,25,10),XSV(10,5)
COMMON X(10),Y(10),PSI(10),ETA(10),ZXY(21,10),ZZZ(21,10)
COMMON /CGM1/ N,M,NR,NC,NDATA,NMOD,NSYM,NDEX,NCC,NCS1,NCST
COMMON /CGM1/FM,FR,80,SF,ST
COMMON /CGM2/ Q(5),DELP(5,10),NIC,NIS,ZIS,ZIP,ZCP,SF0ST
COMMON /CGM2/BF,BMF, CGS1,SIN1,CT32P,CT32M,CT54P,CT54M
COMMON /CGNST/ PI,PI0VR2,RAD,ZI(10)
DIMENSION AW(5,6,10),AT(5,6,10),AS(10)
EQUIVALENCE (A,AW),(A(301),AT)
DIMENSION FRC(2)
DIMENSION U(6),V(6),H(6)
DATA H/0.085662246,0.18038079,2*0.23395697,0.18038079,0.085662246/
DATA V/0.96623476,0.83060459,0.61930959,0.38069041,0.16939531,
10.03376524/
C WEIGHTS AND POINTS FOR 6 POINT GAUSSIAN QUADRATURE ON (0,1)
DATA U/-0.965926,-0.707107,-0.258819,0.258819,0.707107,0.965926/
DATA PI6/0.52359877/
C POINTS FOR 6 POINT GAUSS QUADRATURE ON (-1,1) WITH WEIGHT FUNCTION
1/SQRT(1-U**2). WEIGHTS ARE PI/6
INDEX=NDEX-1
DO 2 I=1,300
AT(I,1,1)=0.0
2 AW(I,1,1)=0.0
4 CS=ST-SF
BT=BF+DS*CT54M
DO 12 L=1,2
5 DO 12 I=1,6
GO TO (6,7),L
6 SG=SF*V(I)
XM=1.0+SG*CT32P
B =1.0+SG*CT32M
RS=SF
GO TO 8
7 SG=ST-DS*V(I)**2
RS=2.0*DS*V(I)
SG1=SG-SF

```

F0RC

```

XM=BMF+SG1*CT54P
B = BF+SG1*CT54M
8 RS = RS*PI6*H(I)*SQRT(ST**2-SG**2)
D0 9 K=1,10
AS(K)=RS
9 RS=SG*RS
D0 12 J=1,6
Z=XM+B*U(J)
Z IS X
D0 12 L1=1,N
FAC=FXI (U(J),L1)
G0 T0 (23,24),L
23 D0 11 L2=1,6
D0 10 K=1,10
10 AW(L1,L2,K)=AW(L1,L2,K)+FAC*AS(K)
11 FAC=Z*FAC
G0 T0 12
24 D0 21 L2=1,6
D0 20 K=1,10
20 AT(L1,L2,K)=AT(L1,L2,K)+FAC*AS(K)
21 FAC=Z*FAC
12 CONTINUE
AREA=SF*(1.0+BF)+DS*(BF+BT)
D0 30 N1=1,NM0D
U0 30 N2=1,NM0D
FRC(1)=0.0
FRC(2)=0.0
I2 = 1
G0=1.0
D0 17 J=1,6
G=GO
G0=GO*BO
JL=7-J
D0 17 I = 1,JL
L3 = 1
F1=ZXY(I2,N2)*G
F2=ZZZ(I2,N2)*G

```

C

00005420
00005430
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00005470
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00005490
00005500
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00005640
00005650
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00005670
00005680
00005690
00005700
00005710
00005720
00005730
00005740
00005750
00005760
00005770
00005780

F0RC

```
G=G*BO
IF (F1.EQ.0.0 .AND. F2.EQ.0.0) GO TO 17
Z=1.0
DO 19 K=1,M
  L2=K+J+IDEX
DO 18 L1=1,N
  F=F1*AW(L1,I,L2)+F2*AT(L1,I,L2)
  F=F*Z
DO 16 M1=1,2
  16 FRC(M1)=FRC(M1)+F*AP(M1,L3,N1)
  18 L3=L3+1
  19 Z=Z/ST
  17 I2 = 1 + I2
  F1=FRC(1)/AREA
  F2=FRC(2)/AREA
  F3=SQRT(F1**2+F2**2)
  F4=DOTS
  IF (F3.NE.0.0) F4=ATAN(F2,F1)*RAD
  30 WRITE (6,47) N1,N2,F1,F2,F3,F4
  47 FORMAT(1H02I6,1P3E19.5,0P1F16.4)
  RETURN
END
```

00005790
00005800
00005810
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00005840
00005850
00005860
00005870
00005880
00005890
00005900
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00005920
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00005940
00005950
00005960
00005970
00005980
00005990
00006000

LSSURP

```

SUBROUTINE LSSURP(X,Y,Z,N,ND,NG,A,IND)
DIMENSION X(1),Y(1),Z(1),A(21),C(21,21),B(21)
DIMENSION G(21),F(21)
DIMENSION JLIM(6)

```

```

C      A POLYNOMIAL WITH THE COEFFICIENTS (A) IS FITTED BY LEAST SQUARES
C      TO THE VALUES (Z) AT THE POINTS (X,Y)
C      DEGREE ND IS GIVEN, MUST BE AT MOST 5
C      NUMBER OF POINTS N IS UNRESTRICTED
C      IND=0 FOR SUCCESSFUL FIT

```

```

      IND=0
      ILIM=ND+1
      IF(ILIM.GT.6) ILIM=6
      IJLIM=ILIM+1

```

```

      DO 2 I=1,ILIM
      JLIM(I)=IJLIM-I

```

```

      2 NT=0

```

```

      DO 4 I=1,ILIM
      NT=NT+JLIM(I)

```

```

      DO 5 I=1,NT

```

```

      B(I)=0.0

```

```

      DO 5 J=1,NT

```

```

      C(J,I)=0.0

```

```

      DO 8 K=1,N

```

```

      YP=1.0

```

```

      L=1

```

```

      DO 7 I=1,ILIM

```

```

      XYP=YP

```

```

      JL=JLIM(I)

```

```

      DO 6 J=1,JL

```

```

      G(L)=XYP

```

```

      XYP=X(K)*XYP

```

```

      6 L=L+1

```

```

      7 YP=Y(K)*YP

```

```

      DO 8 I=1,NT

```

```

      B(I)=B(I)+G(I)*Z(K)

```

```

      DO 8 J=1,NT

```

```

      C(J,I)=C(J,I)+G(J)*G(I)

```

```

00006020
00006030
00006040
00006050
00006060
00006070
00006080
00006090
00006100
00006110
00006120
00006130
00006140
00006150
00006160
00006170
00006180
00006190
00006200
00006210
00006220
00006230
00006240
00006250
00006260
00006270
00006280
00006290
00006300
00006310
00006320
00006330
00006340
00006350
00006360
00006370
00006380

```

LSSURP

```

00006390
00006400
00006410
00006420
00006430
00006440
00006450
00006460
00006470
00006480
00006490
00006500
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00006520
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00006620
00006630
00006640
00006650
00006660
00006670
00006680
00006690
00006700
00006710
00006720
00006730
00006740
00006750

      DO 10 I=1,NT
      F(I)=0.0
      DO 9 J=1,NT
      9 F(I)=AMAX1(F(I),ABS(C(J,I)))
      DO 10 J=1,NT
      10 C(J,I)=C(J,I)/F(I)
      K=MSIMER(21,NT,1,C,B)
      IF (K.EQ.1) GO TO 15
      11 DO 12 I=1,ILIM
      IP=ILIM+1-I
      IF (JLIM(IP)+IP.EQ.IJLIM) GO TO 13
      12 CONTINUE
      IJLIM=IJLIM-1
      IF (IJLIM.GT.1) GO TO 11
      IND =1
      RETURN
      13 JLIM(IP)=JLIM(IP)-1
      IF (JLIM(IP).EQ.0) IJLIM=IP-1
      GO TO 3
      15 NG=NT
      K=1
      L=1
      DO 18 I=1,ILIM
      JL=JLIM(I)
      JM=JL+1
      JN=7-I
      DO '6 J=1,JL
      A(L)=B(K)/F(K)
      K=K+1
      16 L=L+1
      IF (JN.LE.JL) GO TO 18
      DO 17 J=JM,JN
      A(L)=0.0
      17 L=L+1
      18 CONTINUE
      IF (L.GT.21) RETURN
      DO 19 K=L,21

```

00006760
00006770
00006780

19 A(K)=0.0
RETURN
END

ZDZ

```

SUBROUTINE ZDZ(X,Y,N,A,Z,DZ)
DIMENSION C(7),X(300),Y(300),Z(300),DZ(300)
DIMENSION A(21)
DATA C/0.0,1.0,2.0,3.0,4.0,5.0,6.0/
FINDS VALUES OF Z,DZ/DX AT THE COLLOCATION POINTS
DO 50 IN = 1,N
  Z(IN) = 0.0
  DZ(IN) = 0.0
  YP = 1.0
  NG = 1
  DO 20 I=1,6
    JL=7-I
    XYP = YP
    DO 10 J = 1,JL
      Z(IN) = Z(IN) + A(NG)*XYP
      XYP = X(IN)*XYP
      IF(J.EQ.1) GO TO 5
      DZ(IN) = DZ(IN) + C(J)*A(NG)*XYQ
      XYQ = X(IN)*XYQ
    GO TO 10
  XYQ = YP
  NG = NG + 1
  YP = Y(IN)*YP
  CONTINUE
  WRITE (6,101)
  FORMAT(1H0,17X,1HX,20X,1HY,20X,1HZ,20X,2HDZ)
  DO 100 I = 1,N
    WRITE (6,105) X(I),Y(I),Z(I),DZ(I)
    FORMAT(1H ,10X,1PE15.4,5X,1PE15.4,6X,1PE15.4,7X,1PE15.4)
  RETURN
END

```

00006800
00006810
00006820
00006830
00006835
00006840
00006850
00006860
00006870
00006880
00006890
00006900
00006910
00006920
00006930
00006940
00006950
00006960
00006970
00006980
00006990
00007000
00007010
00007020
00007030
00007040
00007050
00007060
00007070
00007080
00007090

CNSI'S

BLOCK DATA
COMMON /CONST/ PI,PIOVR2,RAD,ZI(10)
USEFUL CONSTANTS
DATA PI/3.14159265/,PIOVR2/1.57079633/,RAD/57.2957795/
DATA ZI/1.,2.,3.,4.,5.,6.,7.,8.,9.,10./
END

00007110
00007120
00007125
00007130
00007140
00007150

C

```

C      FUNCTION FXI(U1,K)
      EVALUATES THE KTH CHORDWISE PRESSURE MODE AT U1
      U=U1
      IF (1-K) 2,4,8
      2  FXI=1.0-U*U
      IF (K-5) 3,5,8
      3  IF (K-3) 8,5,7
      4  FXI=1.0-U
      GO TO 8
      5  FXI=-2.0*U*FXI
      IF (3-K) 6,8,8
      6  FXI=FXI*(4.0*U*U-2.0)
      GO TO 8
      7  FXI=FXI*(3.0-4.0*FXI)
      8  RETURN
      END

```

```

00007170
00007175
00007180
00007190
00007200
00007210
00007220
00007230
00007240
00007250
00007260
00007270
00007280
00007290
00007300
00007310

```


K=MSIMER(N,L,LB,A,B)

SIMR

STD	A12+2
STD	A20
STD	A22
STD	A25+1
STD	A25+2
STD	A33
STD	A36+1
STD	A36+2
SXD	A52,1
TXI	*+1,1,1
SCD	A32,1
TXI	*+1,1,-2
SCD	A22+1,1
CLA*	5,4
PAX	0,7
SXA	1A6-1,7
SXA	A14,7
SXA	1A21-1,7
SXA	A26,7
SXA	A37,7
CLA*	4,4
PAX	0,1
TXL	E1,1,0
TXH	E1,1,**
SXA	A2,1
SXA	A5-1,1
SXD	A9,1
SXD	A12,1
TXI	*+1,1,-1
SXD	A7,1
SXD	A36,1
SXD	A18,1
SXD	A38-1,1
SXD	A21-1,1
SXD	A23,1
SXD	A25,1
SXD	A31,1

A52

00050380
00050390
00050400
00050410
00050420
00050430
00050440
00050450
00050460
00050470
00050480
00050490
00050500
00050510
00050520
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00050570
00050580
00050590
00050600
00050610
00050620
00050630
00050640
00050650
00050660
00050670
00050680
00050690
00050700
00050710
00050720
00050730
00050740

K=MSIMER(N,L,LB,A,B)

SIMR

1A6

FMP B,6
STG B,6
LDQ T
TNX 2A6,4,1
TXI 1A6,6,-N
TNX A7,1,1
TXI **1,3,1
TXI A2,5,1
TXH B7,2,L -1

A7

*
*
*

SEARCH FOR MAXIMUM PIVOT IN COLUMN

PXA 0,2
PAX 0,1
PXA 0,3
PAX 0,6
PXA 0,0
LDQ A,6
LRS 0
TLQ **3

A8

XCA A10,1
SXA **1,1,1
TXI **2,1,L
TXH A8,6,-1
TXI TOL
LDQ **2
TLQ E3
TRA **1
AXT **1,2
SXD A17,1,**

A9

A10

ROW INTERCHANGE

PXA 0,3
PAX 0,6
PAX 0,7

00051120
00051130
00051140
00051150
00051160
00051170
00051180
00051190
00051200
00051210
00051220
00051230
00051240
00051250
00051260
00051270
00051280
00051290
00051300
00051310
00051320
00051330
00051340
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00051370
00051380
00051390
00051400
00051410
00051420
00051430
00051440
00051450
00051460
00051470
00051480

K=MSIMER(N,L,LB,A,B)

SIMR

	SCD	*+1,1	00051490
	TXI	*+1,7,**	00051500
	PXA	0,2	00051510
	PAX	0,4	00051520
A11	CLA	A,6	00051530
	LDQ	A,7	00051540
	STQ	A,7	00051550
	STQ	A,6	00051560
	TXI	*+1,4,1	00051570
A12	TXH	A13,4,L	00051580
	TXI	*+1,6, -N	00051590
	TXI	A11,7, -N	00051600
A13	PXA	0,5	00051610
	PAX	0,6	00051620
	PAX	0,7	00051630
A14	AXT	LB,4	00051640
	SCD	*+1,1	00051650
	TXI	*+1,6,**	00051660
A15	CLA	B,7	00051670
	LDQ	B,6	00051680
	STQ	B,6	00051690
	STQ	B,7	00051700
	TNX	A17,4,1	00051710
A16	TXI	*+1,6, -N	00051720
	TXI	A15,7, -N	00051730
*			00051740
*			00051750
*			00051760
A17	CLA	=1.0	00051770
	FDP	A,3	00051780
	STQ	AM	00051790
A18	TXH	A21,2,L -1	00051800
	PXA	0,2	00051810
	PAX	0,4	00051820
	PXA	0,3	00051830
	PAX	0,6	00051840
A20	TXI	*+1,6, -N	00051850

DIVISION OF ROW BY PIVOT

K=MSIMER(N,L,LB,A,B)

SIMR

	LDQ	A,6	00051860
	FMP	AM	00051870
	STG	A,6	00051880
	TXI	*+1,4,1	00051890
	TXL	A20,4,L -1	00051900
A21	PXA	0,5	00051910
	PAX	0,6	00051920
	AXT	LB,4	00051930
1A21	LDQ	B,6	00051940
	FMP	AM	00051950
	STG	B,6	00051960
	TNX	*+2,4,1	00051970
2A21	TXI	1A21,6, -N	00051980
*			00051990
*			00052000
*			00052010
			00052020
			00052030
			00052040
			00052050
			00052060
			00052070
			00052080
			00052090
			00052100
			00052110
			00052120
A22	TXI	*+1,7, -N	00052130
	TXI	*+1,3, -N+1	00052140
	TXI	A28,7	00052150
A23	SXA	A26,2,L -1	00052160
	TXH	A27,3	00052170
	SXA	A,6	00052180
A24	LDQ	A,7	00052190
	FMP		00052200
	CHS		00052210
	FAD		00052220
	STG		00052230

ROW REDUCTION

K=MSIMER(N,L,LB,A,B)

SIMR

A25	TXI	**1,4,1	00052230
	TXH	A27,4,L -1	00052240
	TXI	**1,3, -N	00052250
	TXI	A24,7, -N	00052260
A27	AXT	**3	00052270
A26	AXT	LB,4	00052280
	AXT	**7	00052290
	TXI	**1,7,1	00052300
	SXA	**2,7	00052310
1A26	LDQ	A,6	00052320
	FMP	B,5	00052330
	CHS		00052340
	FAD	B,7	00052350
	STG	B,7	00052360
	TNX	3A26,4,1	00052370
2A26	TXI	**1,5, -N	00052380
	TXI	1A26,7, -N	00052390
3A26	AXT	**5	00052400
	TNX	A29,1,1	00052410
A28	AXT	**7	00052420
	PXA	0,2	00052430
	PAX	0,4	00052440
	TXI	**1,3,1	00052450
	TXI	A23,6,1	00052460
A29	AXT	**3	00052470
A31	TXH	A43,2,L -1	00052480
	PXA	0,2	00052490
	PAX	0,1	00052500
	PAX	0,4	00052510
	PXA	0,3	00052520
	PAX	0,6	00052530
	PAX	0,7	00052540
A32	TXI	**1,3, -N-1	00052550
	TXI	**1,6,-1	00052560
A33	TXI	**1,7, -N	00052570
	SXA	A37+1,5	00052580
	SXA	A41,5	00052590

K=MSIMER(N,L,LB,A,B)

SIMR

A34	SXA	A40,3	00052600
A35	SXA	A39,7	00052610
	SXA	A38,3	00052620
	LDQ	A,6	00052630
	FMP	A,7	00052640
	CH		00052650
	FAL	A,3	00052660
	STG	A,3	00052670
	TXI	*+1,4,1	00052680
A36	TXH	A37,4,L -1	00052690
	TXI	*+1,3, -N	00052700
	TXI	A35,7, -N	00052710
A37	AXT	LB,4	00052720
	AXT	**7	00052730
	TXI	*+1,7,-1	00052740
	SXA	*-2,7	00052750
1A37	LDQ	A,6	00052760
	FMP	B,5	00052770
	CHS		00052780
	FAD	B,7	00052790
	STG	B,7	00052800
	TXI	A41,4,1	00052810
2A37	TXI	*+1,5, -N	00052820
	TXI	1A37,7, -N	00052830
A41	AYT	**5	00052840
	TXI	*+1,1,1	00052850
	TXH	A40,1,L -1	00052860
A38	AXT	**3	00052870
A39	AXT	**7	00052880
	PXA	0,2	00052890
	PAX	0,4	00052900
	TXI	*+1,3,-1	00052910
	TXI	A34,6,-1	00052920
A40	AXT	**3	00052930
	TXI	*+1,5,-1	00052940
	TXI	A7,2,1	00052950
A45	CLA	=1	00052960

K=MSIMER(N,L,LB,A,B)

SIMR

00052970
00052980
00052990
00053000
00053010
00053020
00053030
00053040
00053050
00053060
00053070
00053080
00053090
00053100
00053110
00053120
00053130
00053140
00053150
00053160
00053170
00053180
00053190
00053200
00053210
00053220
00053230

TRA MSIMER+1

BRANCH FOR LAST ROW

CLA A,3
SSP
LDQ TOL
TLQ AL7

ERROR BRANCHES

CLA =2
TRA MSIMER+1
CLA =3
TRA MSIMER+1

STORAGE

GCT 151400000000
EQU E.1
EQU E.2
EQU 0
EQU 0
EQU 0
EQU 0
EQU 0
EQU 0
END

*
*
*
B7

*
*
*
E3
E1
*
*
*
TOL
AM
T
A
B
L
N
LB

K=MSIMEC(N,L,LB,A,B)

SIMC

STD 2A37+1
 STD A12+2
 STD A20
 STD A22
 STD A25+1
 STD A25+2
 STD A33
 STD A36+1
 STD A36+2
 STD A52+1
 TXI *+1+1,2
 SCD A32+1
 TXI *+1+1,-4
 SCD A22+1,1
 CLA* 5,4
 PAX 0,7
 SXA 1A6-1,7
 SXA A14,7
 SXA 1A21-1,7
 SXA A26,7
 SXA A37,7
 CLA* 4,4
 ALS 1
 PAX 0,1
 TXL E1,1,1
 TXH E1,1,*
 SXA A2,1
 SXA A5-1,1
 SXD A9,1
 SXD A12,1
 TXI *+1,1,-2
 SXD A7,1
 SXD A36,1
 SXD A18,1
 SXD A38-1,1
 SXD A21-1,1
 SXD A23,1

A52

00053620
 00053630
 00053640
 00053650
 00053660
 00053670
 00053680
 00053690
 00053700
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 00053930
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 00053950
 00053960
 00053970
 00053980

K=MSIMEC(N,L,LB,A,B)

SIMC

00053990
00054000
00054010
00054020
00054030
00054040
00054050
00054060
00054070
00054080
00054090
00054100
00054110
00054120
00054130
00054140
00054150
00054160
00054170
00054180
00054190
00054200
00054210
00054220
00054230
00054240
00054250
00054260
00054270
00054280
00054290
00054300
00054310
00054320
00054330
00054340
00054350

A25,1
A31,1
A51,1
A51+1,1
*+1,1,2
6,4
0,3
7,4
0,5
*+1,3,2L -2
*+1,5,2L -2
NORMALIZATION OF ROWS
AXT 2L,2
PXA 0,3
PAX 0,6
PAX 0,7
PXA 0,0
LDQ A,6
LRS 0
TLQ *+2
XCA A+1,6
LDQ 0
LRS *+2
XCA A4,2,2
TNX A3,6,2N
TXI E1
TZE T
STG =1.0
CLA T
FDP T
STQ 2L,2
AXT A,7
FMP A,7
STG A,7

A51

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*

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A2

A3

A4

A5

K=MSIMEC(N,L,LB,A,B)

SIMC

	LDQ	T
	FMP	A+1,7
	STG	A+1,7
	LDQ	T
	TNX	A6,2,2
	TXI	A5,7,2N
A6	PXA	O,5
	PAX	O,6
	AXT	LB,4
1A6	FMP	B,6
	STG	B,6
	LDQ	T
	FMP	B+1,6
	STG	B+1,6
	LDQ	T
	TNX	2A6,4,1
2A6	TXI	1A6,6,2N
	TNX	A7,1,2
	TXI	**1,3,2
	TXI	A2,5,2
A7	TXH	B7,2,2L -2
*		
*		
*		

SEARCH FOR MAXIMUM PIVOT IN COLUMN

	PXA	O,2
	PAX	O,1
	PXA	O,3
	PAX	O,6
	PXA	O,0
	LDQ	A,6
	LRS	O
	TLQ	**+3
	XCA	
	SXA	A10,1
	LDQ	A+1,6
	LRS	O
	TLQ	**+3

00054360
00054370
00054380
00054390
00054400
00054410
00054420
00054430
00054440
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00054460
00054470
00054480
00054490
00054500
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00054640
00054650
00054660
00054670
00054680
00054690
00054700
00054710
00054720

K=MSIMEC(N,L,LB,A,B)

SIMC

A9	XCA	A10,1	00054730
	SXA	**1,1,2	00054740
	TXI	**2,1,2L	00054750
	TXH	A8,6,-2	00054760
	TXI	T0L	00054770
	LDQ	**2	00054780
	TLQ	E3	00054790
A10	TRA	**1	00054800
	AXT	**1,2	00054810
	SXD	A17,1,**	00054820
	TNX		00054830
*			00054840
*			00054850
*			00054860
	PXA	0,3	00054870
	PAX	0,6	00054880
	PAX	0,7	00054890
	SCD	**1,1	00054900
	TXI	**1,7,**	00054910
	PXA	0,2	00054920
A11	PAX	0,4	00054930
	CLA	A,6	00054940
	LDQ	A,7	00054950
	STQ	A,7	00054960
	STQ	A,6	00054970
	CLA	A+1,6	00054980
	LDQ	A+1,7	00054990
	STQ	A+1,7	00055000
	STQ	A+1,6	00055010
A12	TXI	**1,4,2	00055020
	TXH	A13,4,2L	00055030
	TXI	**1,6,2N	00055040
	TXI	A11,7,2N	00055050
A13	PXA	0,5	00055060
	PAX	0,6	00055070
	PAX	0,7	00055080
A14	AXT	LB,4	00055090

ROW INTERCHANGE

K=MSIMEC(N,L,LB,A,B)

SIMC					
	SCD		++1,1		00055100
	TXI		++1,6,**		00055110
A15	CLA		B,7		00055120
	LDQ		B,6		00055130
	STG		B,6		00055140
	STQ		B,7		00055150
	CLA		B+1,7		00055160
	LDQ		B+1,6		00055170
	STG		B+1,6		00055180
	STQ		B+1,7		00055190
A16	TNX		A17,4,1		00055200
	TXI		++1,6,2N		00055210
	TXI		A15,7,2N		00055220
*					00055230
*					00055240
*					00055250
A17	NZT		A,3		00055260
	TRA		B1		00055270
	LDQ		A,3		00055280
	FMP		A,3		00055290
	STG		T		00055300
	LDQ		A+1,3		00055310
	FMP		A+1,3		00055320
	FAD		T		00055330
	STG		T		00055340
	CLA		A,3		00055350
	FDP		T		00055360
	STQ		AM		00055370
	CLA		A+1,3		00055380
	FDP		T		00055390
	STQ		AN		00055400
A18	TXH		A21,2,2L -2		00055410
	PXA		O,2		00055420
	PAX		O,4		00055430
	PXA		O,3		00055440
	PAX		O,6		00055450
A20	TXI		++1,6,2N		00055460

DIVISION OF ROW BY PIVOT

K=MSIMEC(N,L,LB,A,B)

SIMC

LDQ	A,6	00055470
FMP	AN	00055480
STG	T	00055490
LDQ	A+1,6	00055500
FMP	AM	00055510
FSB	T	00055520
LDQ	A+1,6	00055530
STG	A+1,6	00055540
FMP	AN	00055550
STG	T	00055560
LDQ	A,6	00055570
FMP	AM	00055580
FAD	T	00055590
STG	A,6	00055600
TXI	*+1,4,2	00055610
TXL	A20,4,2L -2	00055620
PXA	0,5	00055630
PAX	0,6	00055640
AXT	LB,4	00055650
LDQ	B,6	00055660
FMP	AN	00055670
STG	T	00055680
LDQ	B+1,6	00055690
FMP	AM	00055700
FSB	T	00055710
LDQ	B+1,6	00055720
STG	B+1,6	00055730
FMP	AN	00055740
STG	T	00055750
LDQ	B,6	00055760
FMP	AM	00055770
FAD	T	00055780
STG	B,6	00055790
TXN	*+2,4,1	00055800
TXI	1A21,6,2N	00055810
2A21		00055820
*		00055830
*		

ROW REDUCTION

K=MSIMEC(N,L,LB,A,B)

SIMC

	PXA	0,2	00055840
	PAX	0,1	00055850
	PAX	0,4	00055860
	TNX	A31,1,2	00055870
	PXA	0,3	00055880
	PAX	0,6	00055890
	PAX	0,7	00055900
	STA	A29	00055910
	SXA	A26+1,5	00055920
	SXA	3A26,5	00055930
	TXI	*+1,6,2	00055940
A22	TXI	*+1,7,2N	00055950
	TXI	*+1,3,2N -2	00055960
	SXA	A28,7	00055970
A23	TXH	A26,2,2L -2	00055980
	SXA	A27,3	00055990
A24	LDQ	A,6	00056000
	FMP	A,7	00056010
	STG	T	00056020
	LDQ	A+1,6	00056030
	FMP	A+1,7	00056040
	FSB	T	00056050
	FAD	A,3	00056060
	STG	A,3	00056070
	LDQ	A,6	00056080
	FMP	A+1,7	00056090
	STG	T	00056100
	LDQ	A+1,6	00056110
	FMP	A,7	00056120
	FAD	T	00056130
	CHS		00056140
	FAD	A+1,3	00056150
	STG	A+1,3	00056160
	TXI	*+1,4,2	00056170
A25	TXH	A27,4,2L -2	00056180
	TXI	*+1,3,2N	00056190
			00056200

K=MSIMEC(N,L,LB,A,B)

SIMC

A27	TXI	A24,7,2N	00056210
A26	AXT	**3	00056220
	AXT	LB,4	00056230
	AXT	**7	00056240
	TXI	**1,7,2	00056250
	SXA	*-2,7	00056260
1A26	LDQ	A,6	00056270
	FMP	B,5	00056280
	STG	T	00056290
	LDQ	A+1,6	00056300
	FMP	B+1,5	00056310
	FSB	T	00056320
	FAD	B,7	00056330
	STG	B,7	00056340
	LDQ	A,6	00056350
	FMP	B+1,5	00056360
	STG	T	00056370
	LDQ	A+1,6	00056380
	FMP	B,5	00056390
	FAD	T	00056400
	CHS	B+1,7	00056410
	FAJ	B+1,7	00056420
	STG	3A26,4,1	00056430
2A26	TXI	**1,5,2N	00056440
	TXI	1A26,7,2N	00056450
3A26	AXT	**5	00056460
	TNX	A29,1,2	00056470
A28	AXT	**7	00056480
	PXA	O,2	00056490
	PAX	O,4	00056500
	TXI	**1,3,2	00056510
	TXI	A23,6,2	00056520
A29	AXT	**3	00056530
A31	TXH	A43,2,2L -2	00056540
	PXA	O,2	00056550
	PAX	O,1	00056560
			00056570

K=MSIMEC(N,L,LB,A,B)

SIMC

	PAX	0,4
	PXA	0,3
	PAX	0,6
	PAX	0,7
A32	TXI	*+1,3,2N +2
	TXI	*+1,6,-2
A33	TXI	*+1,7,2N
	SXA	A37+1,5
	SXA	A41,5
	SXA	A40,3
	SXA	A39,7
A34	SXA	A38,3
A35	LDQ	A,6
	FMP	A,7
	STG	T
	LDQ	A+1,6
	FMP	A+1,7
	FSB	T
	FAD	A,3
	STG	A,3
	LDQ	A,6
	FMP	A+1,7
	STG	T
	LDQ	A+1,6
	FMP	A,7
	FAD	T
	CHS	A+1,3
	FAD	A+1,3
	STG	*+1,4,2
A36	TXI	A37,4,2L -2
	TXH	*+1,3,2N
	TXI	A35,7,2N
	TXI	LB,4
A37	AXT	**7
	AXT	*+1,7,-2
	TXI	*-2,7
	SXA	

00056580
00056590
00056600
00056610
00056620
00056630
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00056670
00056680
00056690
00056700
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00056860
00056870
00056880
00056890
00056900
00056910
00056920
00056930
00056940

K=MSIMEC(N,L,LB,A,B)

SIMC

1A37	LDQ	A,6	00056950
	FMP	B,5	00056960
	STG	T	00056970
	LDQ	A+1,6	00056980
	FMP	B+1,5	00056990
	FSB	T	00057000
	FAD	B,7	00057010
	STG	B,7	00057020
	LDQ	A,6	00057030
	FMP	B+1,5	00057040
	STG	T	00057050
	LDQ	A+1,6	00057060
	FMP	B,5	00057070
	FAD	T	00057080
	CHS		00057090
	FAD	B+1,7	00057100
	STG	B+1,7	00057110
	TNX	A41,4,1	00057120
2A37	TXI	**+1,5,2N	00057130
	TXI	1A37,7,2N	00057140
A41	AXT	**+5	00057150
	TXI	**+1,1,2	00057160
	TXH	A40,1,2L -2	00057170
A38	AXT	**+3	00057180
A39	AXT	**+7	00057190
	PXA	0,2	00057200
	PAX	0,4	00057210
	TXI	**+1,3,-2	00057220
	TXI	A34,6,-2	00057230
A40	AXT	**+3	00057240
	TXI	**+1,5,-2	00057250
	TXI	A7,2,2	00057260
A43	CLA	=1	00057270
	TRA	MSIMEC+1	00057280
B1	STZ	AM	00057290
	CLA	=1.0	00057300
	FDP	A+1,3	00057310

K=MSIMEC(N,L,LB,A,B)

SIMC

STQ AN
TRA A18

BRANCH FOR LAST ROW

CLA A,3
SSP
LOQ TOL
TLQ A17+2
CLA A+1,3
SSP
TLQ A17

ERROR BRANCHES

CLA =2
TRA MSIMEC+1
CLA =3
TRA MSIMEC+1

STORAGE

GCT 151400000000

EQU 0
EQU 0
EQU 0
EQU 0
EQU 0
END

00057320
00057330
00057340
00057350
00057360
00057370
00057380
00057390
00057400
00057410
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00057440
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00057460
00057470
00057480
00057490
00057500
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00057530
00057540
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00057560
00057570
00057580
00057590
00057600
00057610
00057620
00057630

N.A.A. SUBROUTINE LIBRARY

BESEL	700		
TSX	B,4		BESE2690
XCA			BESE2700
FMP	FRE+18		BESE2710
FAD	FRE+16		BESE2720
STG	FRE+16		BESE2730
CAS	FRE+17		BESE2740
TRA	K2		BESE2750
TRA	K3		BESE2760
STG	FRE+17	K2	BESE2770
CLA	FRE+2		BESE2780
FAD	FL01		BESE2790
STG	FRE+2		BESE2800
TRA	K1		BESE2810
STG	FRE+18	K3	BESE2820
CLM			BESE2830
SLW	FRE+2		BESE2840
SLW	FRE+16		BESE2850
CLA	FL01		BESE2860
STG	FRE+15		BESE2870
TSX	C,4	K4	BESE2880
XCA			BESE2890
FMP	FRE+15		BESE2900
FAD	FRE+16		BESE2910
STG	FRE+16		BESE2920
CLS	FRE+15		BESE2930
STG	FRE+15		BESE2940
CLA	FRE+2		BESE2950
FAD	FL01		BESE2960
STG	FRE+2		BESE2970
FSB	FRE+1		BESE2980
THI	K4		BESE2990
CLA	FRE+14	K6	BESE3000
LBT			BESE3010
TRA	K7		BESE3020
CLA	FL01		BESE3030
TRA	K8		BESE3040
CLS	FL01	K7	BESE3050

73538 BESEL 700

8E5E32060
8E5E3207C
8E5E32080
8E5E32090
8E5E32100
8E5E32110
8E5E32120
8E5E32130
8E5E32140
8E5E32150
8E5E32160
8E5E32170
8E5E3218C
8E5E32190
8E5E32200
8E5E32210
8E5E32220
8E5E32230
8E5E32240
8E5E32250
8E5E32260
8E5E32270
8E5E32280
8E5E32290
8E5E32300
8E5E32310
8E5E32320
8E5E32330
8E5E32340
8E5E32350
8E5E32360
8E5E32370
8E5E32380
8E5E32390
8E5E32400
8E5E32410
8E5E32420

N.A.A. SUBROUTINE LIBRARY

BESEL	700	STG	FRE+7	BESE3430
		FSB	FLG1	BESE3440
		FDP	FRE+11	BESE3450
		STQ	FRE+8	BESE3460
		CLA	FRE+8	BESE3470
	KA4	FAD	FRE+16	BESE3480
		STG	FRE+16	BESE3490
		CAS	FRE+17	BESE3500
		TRA	KA5	BESE3510
		TRA	KA6	BESE3520
	KA5	STG	FRE+17	BESE3530
		CLA	FRE+3	BESE3540
		FAD	FLG1	BESE3550
		STG	FRE+3	BESE3560
		TSX	D,4	BESE3570
		TRA	KA4	BESE3580
	KA6	FAD	FLG1	BESE3590
		XCA		BESE3600
		FMP	FRE+21	BESE3610
		STG	FRE+19	BESE3620
		CLS	FRE	BESE3630
		FAD	FLG88	BESE3640
		TMI	KA7	BESE3650
		CLA	FRE+19	BESE3660
		TRA	ESC	BESE3670
	KA7	LDQ	FRE+19	BESE3680
	KA8	FMP	EM88	BESE3690
		TRA	ESC	BESE3700
	*		FACTORIAL	BESE3710
	*		M IN ACC	BESE3720
	*		LEAVE M/ IN ACC	BESE3730
	FACT	TNZ	FACT+3	BESE3740
		CLA	FLG1	BESE3750
		TRA	I,4	BESE3760
		STG	FRE+4	BESE3770
		STG	FRE+5	BESE3780
		CLA	FRE+4	BESE3790

N.A.A. SUBROUTINE LIBRARY

BESEL	700	FSB	FL01	BESE3800
		TZE	FACT+12	BESE3810
		STG	FRE+4	BESE3820
		LDQ	FRE+5	BESE3830
		FMP	FRE+4	BESE3840
		TRA	FACT+4	BESE3850
		CLA	FRE+5	BESE3860
		TRA	1,4	BESE3870
			SUBR.	BESE3880
			X/2,N+2S/S/N+S/	BESE3890
			X/2,N AND S IN STORAGE	BESE3900
			RESULT LEFT IN ACC	BESE3910
			FRE+10,4	BESE3920
		SXD	FRE+1	BESE3930
		CLA	FRE+2	BESE3940
		FAD	FRE+7	BESE3950
		STG	FRE+2	BESE3960
		FAD	FRE+8	BESE3970
		STG	FRE+2	BESE3980
		CLA	FACT,4	BESE3990
		TSX	FRE+9	BESE4000
		STG	FRE+7	BESE4010
		CLA	FRE+7	BESE4020
		TSX	FACT,4	BESE4030
		XCA		BESE4040
		FMP	FRE+9	BESE4050
		STG	FRE+9	BESE4060
		CLA	FRE+8	BESE4070
		TNZ	A1	BESE4080
		CLA	FL01	BESE4090
		TRA	A2	BESE4100
		LDQ	FRE+6	BESE4110
		FMP	FRE+8	BESE4120
		STG	FRE+24	BESE4130
		CALL	EXP(FRE+24)	BESE4140
		FDP	FRE+9	BESE4150
		STQ	FRE+9	BESE4160
		CLA	FRE+9	

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 A1
 A2

N.A.A. SUBROUTINE LIBRARY

BESEL	700	
LXD	FRE+10,4	BESE4170
TRA	1,4	BESE4180
	PARENTHEICAL SUBR	BESE4190
	X/2,N,AND S	BESE4200
	LEAVE RESULT IN ACC.	BESE4210
SXD	FRE+10,4	BESE4220
CLA	FRE+2	BESE4230
TSX	B1,4	BESE4240
CHS		BESE4250
STQ	FRE+8	BESE4260
CLA	FRE+1	BESE4270
FAD	FRE+2	BESE4280
TSX	B1,4	BESE4290
CHS		BESE4300
FAD	FRE+8	BESE4310
FAD	2GAM	BESE4320
STQ	FRE+8	BESE4330
LDQ	FRE+6	BESE4340
FMP	FLQ2	BESE4350
FAD	FRE+8	BESE4360
LXD	FRE+10,4	BESE4370
TRA	1,4	BESE4380
	SUMMATION OF 1/Y FROM R EQUALS 1 TO LIMIT	BESE4390
STQ	FRE+3	BESE4400
TZE	1,4	BESE4410
CLM		BESE4420
SLW	FRE+9	BESE4430
CLA	FLQ1	BESE4440
FDP	FRE+3	BESE4450
STQ	FRE+11	BESE4460
CLA	FRE+11	BESE4470
FAD	FRE+9	BESE4480
STQ	FRE+9	BESE4490
CLA	FRE+3	BESE4500
FSB	FLQ1	BESE4510
TZE	B1+15	BESE4520
STQ	FRE+3	BESE4530

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* B

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* B1

N.A.A. SUBROUTINE LIBRARY

BESEL	700		
TRA	B1+4		BESE4540
CLA	FRE+9		BESE4550
TRA	1,4		BESE4560
	SUBROUTINE FOR LAST TERM OF SERIES OF YN AND KN		BESE4570
	X/2,N AND S IN STORAGE		BESE4580
	LEAVE RESULT IN ACCUMULATOR		BESE4590
	FRE+1		BESE4600
CLA	FL01		BESE4610
FSB	C3		BESE4620
TPL			BESE4630
ZAC			BESE4640
TRA	1,4		BESE4650
SXD	FRE+10,4		BESE4660
CLA	FRE+2		BESE4670
TSX	FACT,4		BESE4680
STQ	FRE+7		BESE4690
CLA	FRE+1		BESE4700
FSB	FRE+2		BESE4710
FSB	FL01		BESE4720
TSX	FACT,4		BESE4730
FDP	FRE+7		BESE4740
STQ	FRE+7		BESE4750
CLA	FRE+2		BESE4760
FAD	FRE+2		BESE4770
FSB	FRE+1		BESE4780
TNZ	C1		BESE4790
LDQ	FL01		BESE4800
TRA	C2		BESE4810
STQ	FRE+8		BESE4820
LDQ	FRE+6		BESE4830
FMP	FRE+8		BESE4840
STQ	FRE+24		BESE4850
CALL	EXP(FRE+24)		BESE4860
XCA			BESE4870
FMP	FRE+7		BESE4880
LXD	FRE+10,4		BESE4890
TRA	1,4		BESE4900
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FRE+3
FRE+18
FRE+11
FRE+9
FLQ2
FRE+9
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FLQ1
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5. AUTHOR(S) (Last name, first name, initial) Vivian, H. T. Andrew, L. V.		
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13. ABSTRACT Equations for pressure distributions and generalized aerodynamic forces are derived for a thin nonplanar lifting surface in simple harmonic motion at subsonic speeds. A digital computer program, written in Fortran IV, is also presented. The computer program will generate up to a ten by ten matrix of generalized aerodynamic forces when given data for the geometry of a planar lifting surface with a folded planar tip, the flight Mach number, the reduced frequency of motion, and some control constants. Control surface deflections are not accounted for in this study. The kernel function method given by Watkins, Runyan, and Woolston (Reference 1), which relates the pressure distribution to the downwash on a planar lifting surface, has been extended and applied to a nonplanar lifting surface. Hsu's technique (Reference 4) of employing Gaussian quadrature formulas is used when integrating the product of the kernel function and the lift function over the planform area. Recommendations are made to extend the method to account for blunted leading edges and the accompanying airfoil thickness and to account for control surface deflections.		