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ON THE DETERMINATION OF
CHARACTERISTIC VALUES FOR A CLASS
OF STURM-LIOUVILLE PROBLEMS

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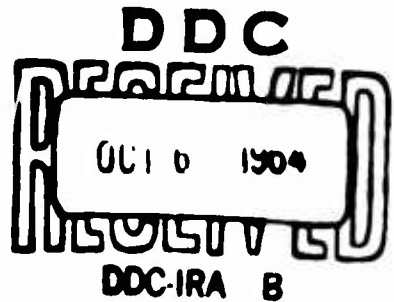
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SUMMARY

In this paper we present a new method for obtaining the characteristic values of the Sturm-Liouville problem $u'' + \lambda a(t)u = 0$, $u(0) = u(1) = 0$. The method yields upper and lower bounds and is particularly suitable for problems in which the first and second characteristic values are desired to a high degree of accuracy.

ON THE DETERMINATION OF CHARACTERISTIC VALUES
FOR A CLASS OF STURM-LIOUVILLE PROBLEMS

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1. INTRODUCTION

In this paper we are interested in the problem of determining the characteristic values of the Sturm-Liouville equation

$$(1) \quad u'' + \lambda a(t)u = 0, \quad u(0) = u(1) = 0.$$

It will be clear from what follows that the methods we discuss can be applied to questions of this type involving quite general boundary conditions, as long as the interval is finite.

There are, at present, a number of powerful techniques available for treating problems of this genre, based upon variational techniques, and upon matrix techniques applied to a finite difference version of the foregoing differential equation.

The variational approach depends upon the fact that if $a(t)$ satisfies a reasonable condition such as

$$(2) \quad 0 < a^2 \leq a(t) \leq b^2 < \infty, \quad 0 \leq t \leq 1,$$

then the characteristic values, $\lambda_1 < \lambda_2 < \dots$, are the respective relative minima of the functional

$$(3) \quad J(u) = \int_0^1 u'^2 dt / \int_0^1 a(t)u^2 dt$$

as u ranges over the space of functions for which the integrals exist and for which $u(0) = u(1) = 0$.

In particular,

$$(4) \quad \lambda_1 \leq \int_0^1 u^2 dt / \int_0^1 a(t)u^2 dt$$

for all functions $u(t)$ satisfying the prescribed boundary conditions. We thus have a means of obtaining upper bounds for λ_1 which turn out to be remarkably accurate even for simple choices of trial functions $u(t)$.

Another method is based upon using equations of the form

$$(5) \quad u_{n+2} - 2u_{n+1} + u_n + \lambda \Delta^2 a_n u_n = 0,$$

$u(0) = u(N) = 0$, and applying any of a number of methods used to derive the characteristic roots and vectors of a symmetric matrix. For a detailed discussion of these methods, and others, we refer to the book by Collatz, [2].

There is, however, a significant difference between a problem of this type, and the Sturm-Liouville problem described above. This is due to the fact that it is quite easy to find asymptotic solutions to (1) for large λ , and thus, approximate expressions for the higher characteristic values.

Let, for simplicity of notation $a(t) = q^2(t)$, then the Liouville transformation, cf. [1], p. 109, $s = \int_0^t q(t_1) dt_1$, converts

$$(6) \quad u'' + \lambda q^2(t)u = 0$$

into

$$(7) \quad \frac{d^2 u}{ds^2} + \frac{q'(t)}{q^2(t)} \frac{du}{ds} + \lambda u = 0.$$

The further transformation

$$(8) \quad v = u \sqrt{q(t)} = u e^{\frac{1}{2} \int_0^s \frac{q'(t)}{q^2(t)} ds}$$

converts (7) into

$$(9) \quad \frac{d^2 v}{ds^2} + \left[\lambda - \frac{1}{2} \frac{d}{ds} \left(\frac{a'(t)}{a^2(t)} \right) - \frac{1}{4} \left(\frac{a'(t)}{a^2(t)} \right)^2 \right] v = 0.$$

The new boundary conditions are

$$(10) \quad v(0) = 0, \quad v\left(\int_0^1 q(t) dt\right) = 0.$$

Writing (9) in the form

$$(11) \quad v''(s) + (\lambda + b(s))v(s) = 0,$$

we know that we can find asymptotic developments for $v(s)$ starting from the integral equation

$$(12) \quad v(s) = c_1 \cos \sqrt{\lambda} s + c_2 \sin \sqrt{\lambda} s - \int_0^s \frac{[\sin \sqrt{\lambda} (s-r)]}{\sqrt{\lambda}} b(r)v(r) dr$$

and iterating, cf. [1], pp. 55-62 for analogous treatment over the infinite interval. Approximate values of λ are now determined by means of the constraint $v\left(\int_0^1 q(t) dt\right) = 0$. Thus, the higher characteristic values have the principal term

$$(13) \quad \lambda_n \approx n^2 \pi^2 / \left(\int_0^1 q(t) dt\right)^2.$$

To obtain more precise results, we can use further terms of

the asymptotic series derived from (12), and we can combine this with numerical integration of (1).

It follows from these considerations that the greatest difficulty is experienced in obtaining accurate estimations of the first characteristic value. In many investigations this is all that is desired.

We wish to present a new method, suitable for hand or digital computer calculation, which furnishes monotone convergence, through sequences of upper and lower bounds, to the smallest characteristic value. Similar sequences can be used to obtain monotone convergence to products of the form $\prod_{i=1}^k \lambda_i$. The method has the advantage of permitting λ_1 to be determined to a high degree of accuracy.

To illustrate these techniques, we consider the equation

$$(14) \quad u'' + \lambda(1+t)u = 0, \quad u(0) = u(1) = 0,$$

which is connected with Airy's function, or Bessel functions of order $1/3$. The computations were performed with the assistance of Marvin Shapiro and Oliver Gross of The RAND Corporation.

2. THE EQUATION DETERMINING THE CHARACTERISTIC VALUES

Let us note in passing that the method we use is an application of an approach we have used, in various lecture courses on differential equations, to derive the fundamental results of Sturm-Liouville theory.

Consider the linear differential equation

$$(1) \quad u'' + \lambda a(t)u = 0, \quad u(0) = 0, \quad u'(0) = 1.$$

The solution of this initial value problem may be obtained over $0 \leq t \leq 1$ as a power series in λ in the form

$$(2) \quad u = t + \sum_{n=1}^{\infty} u_n(t)\lambda^n,$$

where the sequence of coefficient functions $\{u_n(t)\}$, $n = 1, 2, \dots$, may be determined by means of the recurrence relations

$$(3) \quad \begin{aligned} u_0(t) &= t, \\ u_n(t) &= -\int_0^t (t-s)u_{n-1}(s)a(s)ds, \quad n = 1, 2, \dots \end{aligned}$$

It is easy to see that u , as defined by (2), is an analytic function of λ for all finite λ for $0 \leq t \leq 1$. The roots of the equation

$$(4) \quad f(\lambda) = u(1) = 1 + \sum_{n=1}^{\infty} u_n(1)\lambda^n = 0$$

are the desired characteristic values.

3. DISCUSSION

Assuming that the sequence of coefficients is determined by means of either a hand or machine computation, a matter we will discuss again below, there is the problem of determining the first few roots of the equation in (2.4).

This is a problem which can be treated in several ways. It would seem that an efficient procedure would be to use the

sequences we shall describe presently to obtain reasonably accurate estimates for the characteristic values, and then use Newton's method, or a modification, to obtain very accurate values.

4. ANALYTIC PRELIMINARIES

Referring to the equation in (2.1), it is easy to see that

$$(1) \quad |u(t)| \leq e^{k\sqrt{|\lambda|}},$$

for $0 \leq t \leq 1$, where k is a constant. Consequently, the Weierstrass factorization of $f(\lambda)$ takes the form

$$(2) \quad f(\lambda) = \prod_{i=1}^{\infty} (1 - \lambda/\lambda_i).$$

As we know, $\lambda_n = O(n^2)$ as $n \rightarrow \infty$, in view of the assumptions we have made concerning $a(t)$ in (1.2).

Our aim is now, following the technique used by Newton to relate the sums of the powers of the roots and the elementary symmetric functions, which are the coefficients, to obtain relations for the sums

$$(3) \quad b_r = \sum_{i=1}^{\infty} \frac{1}{\lambda_i^r}, \quad r = 1, 2, \dots,$$

in terms of the coefficients $u_n(1)$.

It is clear that

$$\begin{aligned}
 (4) \quad \log f(\lambda) &= \sum_{r=1}^{\infty} \log \left(1 - \frac{\lambda}{\lambda_1^r}\right) = - \sum_{r=1}^{\infty} \frac{\lambda^r}{r} \left\{ \sum_{s=1}^{\infty} \frac{1}{\lambda_1^{rs}} \right\} \\
 &= - \sum_{r=1}^{\infty} \frac{\lambda^r b_r}{r},
 \end{aligned}$$

for $|\lambda| < \lambda_1$.

It is important then to obtain the coefficients of the expression of $\log f(\lambda)$. Although this can be done directly, it is easier to proceed as follows. Write

$$(5) \quad \log f(\lambda) = \sum_{k=1}^{\infty} c_k \lambda^k.$$

Then

$$(6) \quad \frac{f'(\lambda)}{f(\lambda)} = \sum_{k=1}^{\infty} k c_k \lambda^{k-1},$$

whence

$$(7) \quad \sum_{n=1}^{\infty} n u_n(1) \lambda^{n-1} = \left(\sum_{k=1}^{\infty} k c_k \lambda^{k-1} \right) \left(1 + \sum_{n=1}^{\infty} u_n(1) \lambda^n \right),$$

whence we obtain the well-known recurrence relations

$$(8) \quad n u_n = n c_n + \sum_{k=1}^{n-1} k c_k u_{n-k}.$$

These permit us to calculate the c_n in a very simple fashion once the sequence $\{u_n(1)\}$ has been determined, and thus the b_n .

5. INEQUALITIES

Let us now show that the sequence $\{b_k\}$ can be used to

obtain sequences which converge monotonically from above and below to the first characteristic value λ_1 .

Theorem 1. We have the inequalities

$$(1) \quad \frac{b_k}{b_{k+1}} > \lambda_1 > \frac{1}{b_k^{1/k}}, \quad k = 1, 2, \dots$$

The sequence $\{b_k/b_{k+1}\}$ is monotone decreasing; the sequence $\{1/b_k^{1/k}\}$ is monotone increasing, and

$$(2) \quad \lambda_1 = \lim_{k \rightarrow \infty} b_k/b_{k+1} = \lim_{k \rightarrow \infty} 1/b_k^{1/k}.$$

Proof. The monotone character of the ratio b_k/b_{k+1} follows directly from Schwarz's inequality, since

$$(3) \quad b_k^2 = \left(\sum_{i=1}^{\infty} \frac{1}{\lambda_1^k} \right)^2 = \left(\sum_{i=1}^{\infty} \frac{1}{\lambda_1^{\frac{k+1}{2}} \lambda_1^{\frac{k-1}{2}}} \right)^2$$

$$\leq \left(\sum_{i=1}^{\infty} \frac{1}{\lambda_1^{k+1}} \right) \left(\sum_{i=1}^{\infty} \frac{1}{\lambda_1^{k-1}} \right) \leq b_{k+1} b_{k-1}.$$

The monotone behavior of $b_k^{1/k}$ is a consequence of the well-known inequality

$$(4) \quad \left(\sum_{i=1}^{\infty} x_i \right) \geq \left(\sum_{i=1}^{\infty} x_i^2 \right)^{\frac{1}{2}} \geq \left(\sum_{i=1}^{\infty} x_i^3 \right)^{\frac{1}{3}} \geq \dots,$$

for any set of non-negative x_i .

The proof of the limiting relation is clear.

6. RATE OF CONVERGENCE

Since

$$(1) \quad \frac{b_k}{b_{k+1}} = \frac{\frac{1}{\lambda_1^k} \left[1 + \left(\frac{\lambda_1}{\lambda_2}\right)^k + \dots \right]}{\frac{1}{\lambda_1^{k+1}} \left[1 + \left(\frac{\lambda_1}{\lambda_2}\right)^{k+1} + \dots \right]} = \lambda_1 \left[1 + \left(\frac{\lambda_1}{\lambda_2}\right)^k - \left(\frac{\lambda_1}{\lambda_2}\right)^{k+1} + \dots \right],$$

we see that

$$(2) \quad \frac{b_k}{b_{k+1}} - \lambda_1 \approx \lambda_1 \left(\frac{\lambda_1}{\lambda_2}\right)^k$$

for large k .

Similarly,

$$(3) \quad b_k^{\frac{1}{k}} = \frac{1}{\lambda_1} \left(1 + \left(\frac{\lambda_1}{\lambda_2}\right)^k + \dots \right)^{\frac{1}{k}} \approx \frac{1}{\lambda_1} \left(1 + k \left(\frac{\lambda_1}{\lambda_2}\right)^k \right)$$

for large k .

It is to be expected that b_k/b_{k+1} will furnish a better approximation to λ_1 for large k .

7. DISCUSSION

For the case where $a(t) = 1$, $\lambda_1/\lambda_2 = 1/4$. Consequently, in general, the rate of convergence of these sequences will not be too rapid. There are two things we can do to obtain more accurate estimations of λ_1 . In the first place, we can use the root-squaring technique. Since

$$(1) \quad f(\lambda) = \prod_{i=1}^{\infty} \left(1 - \frac{\lambda}{\lambda_i} \right),$$

we see that

$$(2) \quad f_1(\lambda) = f(\sqrt{\lambda})f(-\sqrt{\lambda}) = \prod_{i=1}^{\infty} \left(1 - \frac{\lambda}{\lambda_i^2}\right).$$

Using the power series development for $f_1(\lambda)$ we obtain a sequence $\{b'_k\}$ with

$$(3) \quad \lim_{k \rightarrow \infty} \frac{b'_k}{b'_{k+1}} = \lambda_1^2,$$

and a rate of convergence depending upon $(\lambda_1/\lambda_2)^2$.

Alternatively, once we have an estimate for λ_1 with an accuracy of 1 in 10^{-8} , we can then turn to the power series for $f(\lambda)$ and use the Newton approximation technique,

$$(4) \quad \lambda_1^{(n+1)} = \lambda_1^{(n)} - \frac{f(\lambda_1^{(n)})}{f'(\lambda_1^{(n)})}.$$

This will yield a further approximation with accuracy of essentially 1 in 10^{-28} . Continued use of this technique is limited only by the number of $u_n(1)$ which are computed, and the accuracy of this computation. There is no difficulty involved in using this technique here, since we know from theoretical considerations that the roots of $f(\lambda)$ are simple.

8. INEQUALITIES FOR $\prod_{i=1}^{R+1} \lambda_i$

Similar upper bounds can be obtained for the products $\prod_{i=1}^{R+1} \lambda_i$, $R = 1, 2, \dots$

Consider the determinant

$$(1) \quad b_k^{(R)} = \begin{vmatrix} b_k & b_{k+1} & \cdots & b_{k+R} \\ \vdots & \vdots & \ddots & \vdots \\ b_{k+R} & b_{k+R+1} & \cdots & b_{k+2R} \end{vmatrix}, \quad R = 1, 2, \dots$$

It is not difficult to show that

$$(2) \quad \lim_{k \rightarrow \infty} \frac{b_k^{(R)}}{b_{k+1}^{(R)}} = \lim_{k \rightarrow \infty} (b_k^{(R)})^{-1/k} = \lambda_1 \lambda_2 \cdots \lambda_{R+1}.$$

To show that

$$(3) \quad \frac{b_k^{(R)}}{b_{k+1}^{(R)}} > \frac{b_{k+1}^{(R)}}{b_{k+2}^{(R)}}, \quad k = 1, 2, \dots,$$

for $R = 1, 2, \dots$, we use the well-known fact that the matrix

$$(4) \quad B_k^{(R)} = \begin{pmatrix} b_k & b_{k+1} & \cdots & b_{k+R} \\ \vdots & \vdots & \ddots & \vdots \\ b_{k+R} & b_{k+R+1} & \cdots & b_{k+2R} \end{pmatrix}$$

is positive definite for all k and R , and hence that $(B_k^{(R)})^{-1}$ is positive definite.

The sequence $(b_k^{(R)})^{-1/k}$ does not seem to have any simple monotonicity properties.

9. THE EQUATION $u'' + \lambda(1+t)u = 0$

Let us now illustrate some of the ideas discussed above by means of the equation

$$(1) \quad u'' + \lambda(1+t)u = 0, \quad u(0) = u(1) = 0.$$

The first problem we face is that of computing the sequence $\{u_n(t)\}$ by means of the recurrence relations of (2.3). Since $u(t)$ is an entire function of λ for $0 \leq t \leq 1$, the coefficients, $u_n(t)$, become quite small as n increases. If $a(t) \equiv 1$, the coefficient of λ^n is $(-1)^n/(2n+1)!$. Hence, if we are using a digital computer, even one with floating point arithmetic, it is necessary to renormalize. A very simple renormalization is one which sets

$$(2) \quad v_n(t) = (-1)^n(2n+1)! u_n(t).$$

Then

$$(3) \quad v_n(t) = \frac{1}{(2n+1)2n} \int_0^t (t-s)v_{n-1}(s)a(s)ds, \quad n = 1, 2, \dots,$$

$$v_0(t) = 1.$$

Since (3) is equivalent to the differential recurrence relation

$$(4) \quad v_n''(t) = \frac{a(t)v_{n-1}(t)}{2n(2n+1)}, \quad v_n(0) = v_n'(0) = 0,$$

we can use a Runge-Kutta integration procedure to obtain fairly accurate values of $v_n(1)$. A table is appended.

n	$v_n(1) = (-1)^n(2n+1)! u_n(1)$
0	1.000 000 000
1	1.499 999 92
2	2.238 094 66
3	3.333 330 15
4	4.960 358 93
5	7.378 146 87
6	10.971 261 4
7	16.310 824 0
8	24.244 529 3
9	36.028 967 6
10	53.522 379 4

The decision as to how many elements of sequence $\{u_n(1)\}$ to compute depends upon an a priori estimate of the magnitude of λ_1 , the time involved in the computation, the accuracy of the computation, and the accuracy with which λ_1 is desired.

Since $1 + t \geq 1$, we see that $\lambda_1 < \pi^2 < 10$. Hence the order of magnitude of the last term computed in the power series would be

$$(5) \quad \begin{aligned} u_n(1)\lambda_1^n &< \frac{53.5}{(21)!} 10^{10} < \frac{10^2 \cdot 10^{10}}{(20)!} \\ &< \frac{10^2 \cdot 10^{10}}{2^{20} \cdot 10^{20}} = \frac{10^{-8}}{2^{20}} \end{aligned}$$

(using Sterling's approximation). This is more than sufficient, considering the inaccuracy involved in numerical integration, for the determination of λ_1 , and is sufficient for the determination of $\lambda_2 \leq 4\pi^2$.

The next step is to compute the sequence of coefficients in $\log f(\lambda)$, namely $\{b_k\}$, using (4.8). The results are given below, together with the ratios b_k/b_{k+1} and the roots $b_k^{-1/k}$:

k	b_k	b_k/b_{k+1}	$b_k^{-1/k}$ (Slide Rule calculation)
1	25.000 0	9.921 26	4.00
2	251.984	6.958 90	6.30
3	3,621.03	6.632 47	6.51
4	54,595.5	6.567 79	6.54
5	831,261.0	6.553 06	6.55
6	12,685,100.0	6.549 54	—
7	193,679 x 10 ³	6.548 66	—
8	29,575 x 10 ⁵	—	—

For the purposes of using the Newtonian scheme mentioned above, (7.4), we see that b_4/b_5 and $b_4^{-1/4}$ yield sufficiently good initial approximations with an error of about 1 in 600. One or two applications of (7.4) would yield λ_1 to an accuracy sufficient for most purposes.

The convergence of the sequences for $\lambda_1 \lambda_2$ is much less rapid, as is to be expected. The results are

k	$b_k^{(1)} = b_k b_{k+2}^{-b_{k+1}^2}$	$b_k^{(1)}/b_{k+1}^{(1)}$	$(b_k^{(1)})^{-1/k}$
1	27,030.0	418.85	.37
2	645,330.0	219.85	12.45
3	29,353 x 10 ³	188.93	32.45
4	15,537 x 10 ⁵	179.34	89.5
5	86,634 x 10 ⁶	—	—

Using the value of λ_1 obtained above, we obtain a first approximation of $\lambda_2 \approx 27$. From the monotonicity of the ratios, we know that λ_2 is actually less than this. An application of Newton's approximation will yield a greatly improved result.

Note that λ_2 is sufficiently large so that the asymptotic techniques discussed in §1 can be used to provide an independent check of the accuracy of the first approximation to λ_2 .

10. ALTERNATE COMPUTATIONAL SCHEME FOR POLYNOMIAL COEFFICIENTS

In what has preceded, we have spoken in terms of numerical evaluation of the sequence $\{u_n(t)\}$. Although this procedure has the great advantage of straightforwardness and simplicity, via hand computation or digital computation, it suffers from the fact that errors of integration arise, and grow with each new member of the sequence.

Consequently, it is worth noting a special, but important, case in which we can avoid mechanical quadrature and carry out the entire operation by hand.

Suppose that $a(t)$ is a polynomial of the form

$$(1) \quad a(t) = a_0 + a_1 t + \dots + a_k t^k.$$

It will be clear then that the elements of the sequence $\{u_n(t)\}$ will also be polynomials. Furthermore, it is clear that $u_n(t)$ will have the form

$$(2) \quad u_n(t) = a_0 \frac{t^{2n+1}}{(2n+1)!} + \dots + a_{kn} t^{2n+k} + \dots$$

Using the recurrence relation of (2.3), we can then obtain linear recurrence relations for the sequence $\{a_{kn}\}$,
 $k = 1, 2, \dots, n = 1, 2, \dots$

There are a number of renormalization questions concerned with the effective calculation of the sequence, and asymptotic relations which can be used to speed the computation. A discussion of these would take us too far afield.

11. EXTENSION TO HIGHER ORDER EQUATIONS

Let us now consider the equation

$$(1) \quad u^{(4)} + \lambda a(t)u = 0$$

with the boundary conditions

$$(2) \quad u(0) = u'(0) = 0, \quad u(1) = u'(1) = 0.$$

Proceeding as above, we consider the solution, $u(t, \lambda)$, of the initial value problem

$$(3) \quad u(0) = 0, \quad u'(0) = 0, \quad u''(0) = c_1, \quad u'''(0) = c_2,$$

which we can write in the form

$$(4) \quad u = c_1 u_1(t, \lambda) + c_2 u_2(t, \lambda),$$

where u_1 and u_2 are determined by the initial conditions

$$\begin{aligned}
 (5) \quad & u_1(0) = 0 & u_2(0) &= 0 \\
 & u_1'(0) = 0 & u_2'(0) &= 0 \\
 & u_1''(0) = 1 & u_2''(0) &= 0 \\
 & u_1'''(0) = 0 & u_2'''(0) &= 1.
 \end{aligned}$$

As before, there is no difficulty in obtaining the power series developments in terms of λ for the functions u_1 and u_2 .

Applying the boundary conditions in (2), we obtain the simultaneous equations

$$\begin{aligned}
 (6) \quad & c_1 u_1(1, \lambda) + c_2 u_2(1, \lambda) = 0, \\
 & c_1 u_1'(1, \lambda) + c_2 u_2'(1, \lambda) = 0,
 \end{aligned}$$

whence the determining equation for λ is

$$(7) \quad f(\lambda) = \begin{vmatrix} u_1(1, \lambda) & u_2(1, \lambda) \\ u_1'(1, \lambda) & u_2'(1, \lambda) \end{vmatrix} = 0.$$

From here on, the argumentation is as before.

REFERENCES

1. Bellman, R., Stability Theory of Differential Equations,
McGraw-Hill, 1954.
2. Collatz, L., Numerische Behandlung von Differentialgleichungen,
Springer, Berlin, 1951.