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INSTALLATION AND TEST OF A RECTILINEAR CASCADE

CHARLES C. ROSE DAROLD L. GUTTORMSON I IV I P IV ite School Montrie, Chitoria INSTALLATION AND TEST

OF A

RECTILINEAR CASCADE

by

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Lieutenant Commander, United States Navy

and

Darold L. Guttormson

Captain, United States Marine Corps

Submitted in partial fulfillment of the requirements for the degree of

> MASTER OF SCIENCE IN AERONAUTICAL ENGINEERING

United States Naval Postgraduate School Monterey, California

Library U. S. Naval Postgraduate School Montecey, California

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from the

United States Naval Postgraduate School

ABSTRACT

The purpose of this thesis is to investigate the capabilities of the rectilinear cascade and its associated equipment as installed in the new Propulsion Laboratory of the U. S. Naval Postgraduate School, Monterey, California. The various parameters are developed from the momentum theorem and results are compared with NACA data. Computer programming is included to rapidly evaluate future test data obtained from various blade configurations. The purpose of these tests is to gain more precise knowledge of the physical nature of the flow through cascades and to establish design data for the accurate prediction of the performance of exial-flow turbomachines, turbines as well as pumps and compressors. These efforts, it is hoped, will result eventually in more efficient machines for many applications.

ACKNOWLEDGMENT

The writers wish to express their appreciation to Professor M. H. Vavra for his assistance in this investigation.

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TABLE OF SYMBOLS

drag coefficient, dimensionless $C_{\rm D}$ lift coefficient, dimensionless CL D drag force, lb./ft. d.R. force exerted by fluid on test blade, lb./ft. J conversion unit, ft.lb./B.T.U. \mathbf{L} lift force, 1b./ft. Mach number, dimensionless M integer number of blade spacings, dimensionless \mathbf{N} absolute total pressure, 1b./ft.² P gas constant, ft./OR. R temperature, °R. T v velocity of fluid. ft./sec. speed of sound, ft./sec. a С chord of test blade, ft. specific heat at constant pressure, B.T.U./1b.ºR. cp specific heat at constant volume, B.T.U./1b.ºR. CV mass flow rate, lb.sec./ft. dms gravitational constant. ft./sec.² g test blade height, ft. h i unit vector in tangential direction. dimensionless j unit vector in axial direction, dimensionless n unit vector normal to control surface, dimensionless blade opening or throat, ft. 0 absolute static pressure, 1b./ft.2 p dynamic pressure, 1b./ft.2 q blade spacing, ft. S

t unit vector orthogonal to n in direction of shear stress, dimensionless distance in tangential direction, ft. \mathbf{x} distance along blade height, ft. У gas flow angle, deg. α angle between i and \overline{V} , deg. B 8 ratio of specific heats, dimensionless -5 loss coefficient, dimensionless (subscripted) angle between j and \overline{V} , deg. Θ (non-subscripted) angle between i and dR. deg. Θ 3 x/Ns distance in tangential direction/blade spacing, dimensionless density, 1b.-sec. 2/ft.4 P solidity of cascade, c/s, dimensionless ď shear stress, 1b./ft.2 3 Ø stagger angle, angle between chord line and axial direction, deg. Ψ velocity coefficient, dimensionless Subscripts: axial direction 2 plenum р S static conditions total t tangential direction u 0 station 0, far ahead of cascade 1 station 1, measuring plane ahead of cascade 2 station 2, measuring plane after cascade 3 station 3. far aft of cascade 00 station ∞ , vectorial mean of conditions at stations 0 & 3

1. Introduction

There are many steps in the design of axial-flow turbomachinery and these steps should be performed in a systematic, orderly manner in order to arrive at a method of design which is based on a proper understanding of the flow phenomena that occur in turbomachines. As a first analysis one is concerned with the two-dimensional flow of an ideal fluid; that is a fluid which is homogeneous, continucus, imcompressible, and nonviscous. There are presently several analytical methods of determining the solution of such conditions. When finite blade spacing is considered, these analytical solutions become extremely laborious except for trivial blade shapes, such as flat plates. Therefore with the assumption of an ideal fluid it becomes advantageous to use a more direct approach such as the electrical analogy method described by F. R. Fahland and L. L. Hawkins, Ref. 1.

The introduction of the compressibility and viscosity of the working fluid requires experimental methods of flow analysis since there do not exist analytical methods to evaluate these conditions. The first step in this direction is the use of a rectilinear cascade as described by J. R. Eshman, Ref. 2.

Eshman's investigation was conducted on a small working model of the present large cascade rig which is the subject of this thesis. As might be expected the results obtained from the model cascade were rather poor and limited in value.

The large rectilinear cascade rig is located in the Cascade Laboratory of the new Propulsion Laboratories of the Department of Aeronautics, U. S. Naval Postgraduate School, Monterey, California. The cascade rig, shown in Fig. 1, was designed by Professor M. H. Varva of the Department of Aeronautics, at the U. S. Naval Postgraduate School. An overall view of the Cascade Laboratory is shown in Fig. 2.

Construction of the new Propulsion Laboratories started in the summer of 1962. Fruin-Colnon of San Mateo, California was the general contractor for this construction. The Cascade Laboratory and the cascade rig were accepted in December, 1963. The cascade rig was built by Kenney Engineering Corporation of Pasadena, California. The intrumentation was by General Electric of Ontario, California and Coleman Electronic Systems of Santa Ana, California. These companies were subcontractors to Fruin-Colnon.

The purpose of this cascade test rig is to determine the performance of various blade shapes under conditions of different solidities, inlet angles and incidence angles. This rig makes it possible to undertake systematic tests of cascades for various applications with the aim to evaluate and improve different blade shapes which will lead to the design of improved turbines, compressors, and pumps.

The purpose of this thesis was to investigate whether the cascade rig would perform its designed function. This was accomplished by determining flow deflection and the lift and drag coefficients of a known blade shape and comparing





these values with those published by NACA in Ref. 3. At the same time, parameters more useful to the turbomachinery engineer, such as loss coefficient, will be presented.

A further purpose of this thesis was to indicate improvements and modifications in order to bring the test rig into operating order.

Coupled with the above work, was the testing and utilization of an automatic data logging system. The purpose of this system is to eliminate, or to avoid the laborious task of collecting and transcribing the data into a form compatible with a computer solution. This system utilizes transducers to transform pressure inputs into electrical signals, which are then collected along with other data and punched on paper tape available as input to a computer program. Without the use of this automatic data logging system along with computer solutions the rig could not be utilized to its capabilities because of the long periods required to evaluate the data.

2. Description of Cascade Rig

A general overall view of the cascade is shown in Fig. 3. The rig is essentially an open cycle wind tunnel. Air is brought from the cutside down through a 5 by 12 ft. stack into the blower room. The air then enters a doubleentry blower.

The blower is a Turbovane Fan, number 470997, size 1001, design 2, manufactured by B. F. Sturtevant Co. of



Boston, Massachusetts. This blower is driven by a 60 cycle General Electric Tri-Clad Induction Motor, number 5 M 6346 AA1, rated at 700 horsepower at 4160 volts, 86 amperes, and 1780 RPM. The motor has three speeds which are controlled by a three position switch located on the main power control cabinet. This cabinet is located on the ground floor of the Laboratory as shown in Fig. 3. The air exhausts through the diffusor into the noise attenuator section. The diffusor is 11 feet long with a 2 ft. 2 in. by 3 ft. 5 in. inlet and 5 ft. by 5 ft. outlet. The noise attenuator section is 14.5 ft. high, 7 ft. long and 12 ft. wide. Fig. 4 is a schematic drawing of the noise attenuators. The air then enters the plenum chamber from which the air is supplied to the bell mouth. The plenum chamber is 14.5 ft. high, 12 ft. wide and 21.5 ft. long. The bell mouth is a nine foot by five foot rectangle at the entrance and the sides converge along a hyperbolic curve to a 60 inch by 10 inch rectangle at the exit.

At this point the air is guided into the test section by 29 inlet guide vanes. These IGV's are constructed of aluminum with a blade height of 10 inches and a chord of approximately $5\frac{1}{4}$ inches. Presently there are five different blade shapes available. These blades have an angle of 0, ± 30 and ± 60 degrees between the tangents from the camber line at the leading and trailing edges, Fig. 5. When mounted these blades may be rotated individually to produce the desired flow exit angle. For this thesis the positive 60 degree



Top View

Front View







Angle Notation \emptyset = Stagger Angle

blades were used with the trailing edges set at angles of between 40 and 45 degrees from the vertical depending on the run. Fig. 6 shows a typical setting.

The test section is bounded on the long side by two parallel side walls. One wall acts as the structural support and all blades and attachments are affixed to this wall. The other wall is removable in order to change the position of the end walls or to change blades or blade position. A three ton over-head electric crane is used for the removal of this wall. The removable wall is located by means of dowels to simplify the assembly.

The end walls which bound the short side of the working section guide the flow to the cascade which is to be tested. These walls may be set independently of each other and therefore may not necessarily be parallel. These walls may be set at any angle in order to provide the proper flow direction. Each wall has an upper and a lower portion which are hinged together. This allows the upper portion to be independently set at an arbitrary angle. For testing turbine blades, as in this work, the upper portions are "folded" back in the opposite direction of the lower portions depending on the test blade shape and setting. For this thesis the lower walls were set at 45 degrees and the upper walls were set at -49 degrees. These conditions are also shown in Fig. 6. The lower walls telescope so that the hinge points may be adjusted to various vertical positions.

For testing compressor blades the upper portion of the



end walls would be nearly parallel to the lower portion. The air flow would be deflected between 15 to 25 degrees in a compressor blade test instead of the 30 to 110 degrees as in a turbine blade test. The test blade shape used for these tests was a reaction turbine blade with a camber angle of 95 degrees and a maximum thickness of 15 per cent. This type of blade was established by NACA and used for tests in the cascade investigation conducted by Dunavant and Erwin at the Langley Aeronautical Laboratory, Langley Field, Virginia, Ref. 3. The blades tested at the present cascade rig were cast of epoxy and nickel plated. They are 10 inches high and have a chord of 8.5 inches.

The blades are mounted into a sliding blade holder. This sliding blade holder is required since the blades must be located between the adjustable end walls. For this investigation the test blades were set with a stagger angle as defined in Fig. 5, of =25 degrees and with a spacing of six inches. The spacing is dictated by the blade holder used. Holders with four and ten inch spacings are also available.

The air, after passing through the test section, exhausts into the laboratory. The air can then either exhaust through open windows or be recirculated. In order to recirculate the air, the louvers in the intake stack would be closed and the doors between the laboratory and stack opened.

Inlet guide vanes mounted on the blower are controlled

pneumatically from the laboratory. These vanes are shown in Fig. 7 along with the three speed motor used to drive the blower. By running the blower at one of its three speeds and properly adjusting the inlet guide vanes many volume flow rates may be obtained up to approximately 100,000 cubic feet per minute.

Table I shows the range of pressures, velocities and Reynolds numbers available for the various inlet guide vane positions and motor speeds. Reynolds number is based on a chord of 8.5 inches.

TABLE I

Performance Range

| Pressure | | | Motor | Blower | Velocity | Revnolds . | |
|----------|--------|---------|-------|--------|----------|-------------|--|
| Total | Static | Dynamic | Speed | Vanes | (ft/sec) | Number x 10 | |
| 5.6 | 1.1 | 4.5 | slow | closed | 140 | ,63 | |
| 10.3 | 2.0 | 8.3 | slow | open | 191 | .855 | |
| 11.5 | 2.0 | 9.5 | med. | closed | 204 | .916 | |
| 26.6 | 4.0 | 22.6 | med. | open | 314 | 1.41 | |
| 16,6 | 2.7 | 13.9 | fast | closed | 247 | 1.11 | |
| 46.7 | 8.7 | 38.0 | fast | open | 407 | 1.83 | |

of The Cascade Rig

Pressures - inches of water gage

This cascade rig is unique in many respects. This is due to its advantages and versatility as compared to conventional cascade rigs.



Fig. 8 shows the cross section of a conventional cascade rig. One of the difficulties encountered in this type of cascade rig is the control of the boundary layer ahead of the cascade. Fig. 8 shows that the air passing along the lower wall of the channel must travel a greater distance before reaching the blades than the air passing along the upper wall. Since the thickness of the boundary layer increases with the distance traveled along a wall, the thickness of the boundary layer is greater at the lower blades than at the upper blades. Hence boundary layer removal is required. In the present cascade rig the distance along each wall to the test section is equal. This eliminates the dissimilarity in boundary layer removal.



Fig. 8 Cross Section Of A Conventional Cascade Rig.

This cascade, however, has provisions for removing the boundary layer along the side walls. By conducting tests with and without boundary layer removal, studies can be made concerning the effects of boundary layer on cascades. Hence information can be obtained on the effects of boundary layer in turbomachines.

In all conventional cascade rigs, the air flow is idealized in that conditions of static and total pressure, velocity and flow angle are constant at the entrance of the test section. Thus the effects of inlet guide vanes or preceding rows of blades, as encountered in turbomachinery, are not considered. Due to the inlet guide vane section incorporated in the present cascade, the flow entering the test section approximates the flow in turbomachines.

Any axial flow compressor or turbine blade shape can be tested in this cascade rig. This is accomplished by positioning the inlet guide vanes and the lower and upper portions of the end walls to achieve the proper flow angles. The solidity may be varied by changing blade holders. The present blade holders provide blade spacings of 4, 6, 8, and 10 inches. The blade incidence angle can easily be changed by rotating the blades in the blade holder. All of these factors give the cascade rig a versatility not attained by any conventional cascade rig.

The blower used with the rectilinear cascade rig can also be used for testing circular cascades. Referring to Fig. 9, the doors between the intake stack and the blower room would be closed and the doors between the plenum and the discharge stack would be opened. Doors at the entrance to the bell mouth of the rectilinear cascade rig would also



- A Circular Cascades
- B Noise Attenuating Baffles
- C Rectilinear Cascade
- M Motor



Fig. 9 Schematic Of Circular Cascade Configuration

be closed. Air would then be drawn through the circular cascade by the blower and exhausted out the discharge stack. In this case the blower could provide negative pressures up to approximately 46 inches of water.

Circular cascades with 5.5 ft. outer diameters can be tested at the locations shown on Fig. 9. A unique feature of this arrangement is that the ground level portion of the Cascade Laboratory acts as the plenum chamber for the circular cascades. Hence the flow entering the circular cascades would be more uniform than flow obtained by blowing the air through the cascades.

3. Description of Measuring and Recording Equipment

Static taps in the two side walls of the cascade give an indication of the uniformity of the air flow through the test section. These taps are approximately three inches apart and are located 13 inches ahead of the blades and three inches aft of the blades, as shown in Fig. 10. Plastic tubes connect the static ports with water-filled multitube manometer banks. These banks, as shown in Fig. 11, are 96 inches high and are graduated in tenths of an inch. Green colored water was used in the manometers in order that photographs can be taken during a run for later interpretation.

Two pitot-static probes traverse the test section; one ahead and the other after the blades. The lower probe traverses the test section 13 inches ahead of the center of the test blades. The upper probe can traverse the test





section either 10 or 16 inches behind the center of the test blades. These two upper positions for the probe provide a means by which mixing losses can be determined. This would be accomplished by comparing the results of two succeeding tests differing only in the position of the upper probe. Mixing losses are discussed in section 5 of this thesis.

These United Sensor, YC-120, probes have a three-hole, prism shaped measuring section. The probes measure total and static pressure and yaw angle. They are insensitive to pitch angle up to 20° and are usable up to Mach 0.7. These probes were supplied by United Sensor and Control Corporation of East Hartford, Connecticut. Each probe is mounted in a traverse, Fig. 12, which is located on the stationary side wall of the cascade rig as shown on Fig. 13. The traverse is electrically controlled from an operating and measuring console shown in Fig. 14. The probe extends into the test section through a metal runner and a 1/2 inch slot. This metal runner moves with the traverse and seals the slot in order to prevent air leakage.

Each probe is mounted in a carriage, which moves the probe in and out of the test section (y-direction). The probe can be rotated 90 degrees, in either direction, automatically or manually. This is accomplished manually by switches on the console. The number of degrees rotated are indicated on the console. When on automatic the probe is rotated by a servo.

Also mounted on each traverse are four transducers,






which convert the pressures obtained from the probes into electrical signals. One transducer is used for the total pressure and two for the static pressure. Positive static pressures are transmitted from one of the transducers and negative pressures from the other.

The fourth transducer is used to supply power to a servo which rotates the pitot tube into the flow. This differential transducer measures the difference between the two static pressures obtained from the probe and converts this difference into an electrical signal. This signal drives the servo, which rotates the probe in a direction depending on the sign of the signal. When the pressures become equal no power is supplied to the servo by the differential transducer. With this condition the probe is directed exactly into the air flow and measures the actual total and static guage pressures.

This unique system was designed by Professor M. H. Vavra as an improvement to the tedious and time consuming conventional methods. Conventional methods require the probe to be rotated manually until the difference in static pressures are zero as measured on a differential micromancmeter. Due to the inertia of the fluid columns in the micromanometer, this is a very insensitive system and therefore time consuming. This method would require five to ten minutes to locate the flow direction whereas this would be accomplished within five to ten soconds using the electrical system. Fig. 13 shows the electrical leads that connect the console with the upper and lower traverses.

All controls for positioning and monitoring the probes are contained on the right side of the console shown in Fig. 14. A view of the right side only is found in Fig. 15. Only one of the probes can be controlled and monitored at any one time. Two sets of rotating counters indicate the probe's position in the x and y directions to one hundredth of an inch. The angle that the probe is rotated from the vertical position is indicated by a voltmeter graduated every five degrees clockwise and counter-clockwise. The monitoring system is a separate and distinct system from that used for recording.

The left side of the console as shown in Fig. 16, contains all controls for recording the data. All data received from the probes and the carriages are recorded automatically on punched and printed tapes. Fig. 17 shows the Friden paper tape punch. This machine is installed in a drawer that is closed during normal operations in order to protect the punched tape.

The console uses a digitizer to convert the signals received from the measuring equipment into a certain number of counts. There are 1000 counts (000-999) available for indicating measured values. Each of the six measured quantities $(x, y, \theta, P, p$ positive and p negative) for both upper and lower traverses have gain controls so that the counts may indicate in any dimension. Normally for the pressures, the counts will indicate 0.1 inches of water pressure. An example would be 131 counts which would be







13.1 inches of water. The flow angle, Θ , is scaled so that a reading of 500 counts is zero degrees, 50 is 90 degrees clockwise (-) and 950 is 90 degrees counter-clockwise (+), as viewed from the console. Each count is therefore 0.2 of a degree. An example would be a count of 350. 350 - 500= -150 and 0.2 x \sim 150= -30.0 degrees (CW). The y position is indicated to the hundredth of an inch from 0 to 9.99 inches.

The x position is also indicated to the hundredth of an inch (000-500). In order to have this accuracy it was necessary to recycle the measuring equipment every five inches. Thus after every five inches the counts go to zero and five inches are added to x in the computer program for the CDC 1604. This computer is made by the Control Data Corporation of Minneapolis, Minnesota.

Data can also be manually punched and printed on tape. Presently this is used for such information as plenum temperature, barometric pressure, run number, number of blade spacings, and a code to indicate end of input data.

Connected into the various pressure lines leading to the transducers are tubes which are connected to two manometer boards. One for each of the two probes (upper and lower). On these boards P, $p_{,\Delta}p$ and q are monitored. The main use of these manometers would be in calibrating the conscle readings.

Calibration of the console can be carried out between tests at arbitrary intervals. The procedures used to calibrate the pressure components are given as an example.

All of the pressure components can be calibrated at the same time by connecting their pressure leads to a manifold. With atmospheric pressure on the manifold, the component being calibrated is set to zero using the zero rheostat for that component. A known pressure is then put on the system and the gain control adjusted to give the counts required. The pressure must then be released and the zero rechecked. This procedure may have to be repeated several times. For more accurate readings the known pressure should be equal or greater than those expected during the tests. Checks on the calibration of each component can be made at any time by depressing the calibration switch for that component. The digitizer will then indicate the full scale deflection. This value will not change unless the component is no longer properly calibrated.

Total temperature is obtained from a temperature probe that extends into the plenum chamber. This information is manually punched on the paper tape.

4. Static Pressure Distribution

The momentum theorem, as discussed in Section 5, is used to determine the force exerted by the fluid on the test blades. This requires surveys of the flow conditions ahead of and behind the test blades. Theoretically the volume of fluid that is surveyed behind the blades should be the same as that surveyed ahead of the blades. Actually it is only necessary to survey a volume of air that has had the same

changes in momentum as that surveyed ahead. For that reason it is desired to have a wide regular flow profile entering the test section. A regular flow being defined as a steadyuniform flow except for wakes that are equal and evenly spaced. An accurate survey can then be conducted if the width of the regular flow profile entering the test section exceeds the length of the survey by one spacing. For these tests a static pressure gradient of less than two percent of the test section dynamic pressure was considered permissible.

The cascade was set up as shown in Fig. 6. The lower end walls were set at $\pm 45^{\circ}$, the upper end walls at $\pm 49^{\circ}$, the test blade stagger angle was set at $\pm 25^{\circ}$, and the trailing edges of the inlet guide vanes were set at $\pm 45^{\circ}$. The angle of the inlet guide vanes was determined from the relation:

 $\alpha_2^* = f(\cos^{-1}o/s)$ of Ref. 4, where o is the blade opening and s is the spacing or pitch. Fig. 5, of this reference indicates the relationship between gas outlet angles and $\cos^{-1}o/s$ is linear for straight-backed blades operating at Mach numbers less than 0.5.

Ref. 4 limits this relationship to a Reynolds number of approximately $2 \ge 10^5$ but there seems to be no reason why this relation should not be valid for higher values. There-fore this relation was used to initially position the inlet guide vanes.

Operation of the cascade under these conditions produced a pressure gradient in the x direction ahead of the test blades, Fig. 18. This pressure variation amounted to



20.6% of the test section dynamic pressure and therefore is not permissible.

To correct this situation the inlet guide vanes were set to 40°. In addition, gaskets were installed on the free end of each blade to prevent the passage of air between the blade tips and the wall. Frior to the installation of the gaskets a clearance of 3/32 inch existed between the removable side wall and the test blades. The gaskets eliminated this clearance and form a tight bond between the blade tips and the wall. As a result of the above alterations it was found that although a pressure gradient was still present, it had lessened to 12.9% of the test section dynamic pressure but had reversed its direction, Fig. 19.

The fact that the gradient had reversed indicated that the required inlet guide vane angle was between 40° and 45°. The guide vanes were reset to 42° and a fine mesh wire screen (20 holes/in., 0.016 in. wire diam.) was installed at the exit of the bell mouth immediately ahead of the guide vanes. Operation of the cascade in this configuration produced the pressure distribution as shown of Fig. 20. This figure indicates an undesirable pressure variation near the center of the cascade. It was also found from pressure readings that the total pressure had been reduced from 50 inches of water to 29 inches of water. This is a 42% reduction in total pressure caused by installing the screen just in front of the inlet guide vanes.

For the next run the fine mesh screen was removed. The-





guide vanes, test blades, and wall angles were maintained at the same position as in the preceding test. A comparison of Figs. 20 and 21 indicates the variation in operating with and without a fine mesh screen in front of the inlet guide vanes. These figures show that the general shape of the curve remains the same. The variation of the static pressure is 10.1% of the test section dynamic pressure with the screen and 15.2% without the screen.

Due to the excessive loss of total pressure caused by the location of fine mesh screen in this position, it was necessary to remove the screen. However, to reduce the static pressure variation to within 2% of the test section dynamic pressure the flow pattern had to be corrected.

Reference to Fig. 3 shows that the diffusor from the fan exits to a plenum chamber after passing through sound reducing baffles. To prevent the major portion of the air flow from proceeding out of the diffusor and through only the center part of the sound reducing system, a partial screen was installed on the upstream end of the sound reducers. This screen covers only the portion of the sound reducing baffles (described in section 2) that is directly in the flow from the diffusor. The purpose of the screen is to offer a greater resistance to air passing through the center portion of the sound attenuators and thereby forcing the flow to the edges. In this way the flow into the plenum would be more uniform.

Prior to the next run the inlet guide vanes and test blades were removed to eliminate any effects that the blading could have on the flow. The upper and lower end walls were positioned vertically. With this configuration the static pressure as indicated on the multitube manon-



eter boards were fluctuating by two inches of water. These fluctuations occur: ed at such a rapid rate that it was impossible to take readings. These pressure fluctuations could have been caused onl; by flow irregularities in the plenum chamber. Hence it was necessary to reduce the pressure fluctuations by installing additional screens.

At the entrance to the bell mouth from the plenum chamber a set of two screens were installed. The screens were both fine mesh (20 holes/ in., 0.016 in. diam. wire) placed $3\frac{1}{2}$ inches apart. This arrangement produced a symmetric static pressure distribution with a maximum static pressure variation of 0.35 inches of water, Fig. 22. This variation was 1.55% of the test section dynamic pressure and was within the desired limits stated previously.

Locating the screens at the entrance to the bell mouth produced an 8% total pressure drop from 50 to 46 inches of water. This is a considerable improvement over the 42% total pressure drop caused by placing one screen at the bell mouth exit. The reason for this improvement is that the air flow has a lower velocity at the bell mouth entrance, therefore the screens offer less resistance.

5. Evaluation of Test Data

Measurements of the air flow through the test section are taken just ahead of the test blades at station 1 and just behind the test blades at station 2. The data obtained at these stations must be used to find the conditions that would exist far ahead of the test blades at station 0 and far behind the blades at station 3. This must be done in order to account for the mixing loss.

Mixing loss is a nonreversible process in which fluids of



different velocities are mixed, resulting in a loss of available energy. Therefore, there is a loss in total pressure due to the equalizing of the flow velocities after the air passes through the test blades. A velocity variation exists behind the test blades due to the blade wakes. The equalizing of the flow occurs because of viscous effects and turbulent mixing. These mixing losses are accounted for by analyzing the flow between stations 0 and 3 since uniform conditions are assumed to exist at these locations.

Appendix A shows how the total pressure, static pressure, the flow angle, and the flow velocity at stations 0 and 3 are determined from the measured data.

In order to find these values the momentum theorem is applied between stations 1 and 0 and between stations 2 and 3. The momentum theorem makes it possible to evaluate the force that the flow exerts on the test blades. This force is determined by equation 5(21) of Ref. 5 as:

$$\vec{R} = \int_{S_1} dm_{S_1} \vec{V}_1 - \int_{S_2} dm_{S_2} \vec{V}_2 + \int_{S_1} -\vec{m} p_1 dS_1 + \int_{S_2} -\vec{m}_2 p_2 dS_2 + \int_{S_1} -\vec{\tau}_1 \tau_1 dS_1 + \int_{S_2} -\vec{\tau}_2 \tau_2 dS_2$$
(1)

In the derivation of equations, Appendix A, the last two integrals of equation (1) are ignored. This does not introduce excessive errors as the velocity variations ahead of and after the test section are regular, Figs. 23 and 24. With a regular velocity variation the effects of shear stresses (γ) are small. It is for this reason the last two integrals of equation (1) may be omitted.

The elemental force vector dR of equation (1) may be represented by





$$\overline{dR} = i dR_u + j dR_a \tag{2}$$

Equation (2) shows the force in the tangential direction represented by the subscript u and in the axial direction by the subscript a. On the cascade the vertical direction (j) is axial and the horizontal direction (i) is tangential. The sign convention is as follows:



From this:

$$\Theta = \tan^{-1} \left(\frac{dR_u}{dR_a} \right) \tag{3}$$

The measured values at stations 1 and 2 are total pressures, static pressures, and flow angles. To be found are the values of total pressures, static pressures, flow angles and velocities at stations 0 and 3. This is done in Appendix A by applying the momentum theorem (1) and the equation of continuity.

The equation of continuity is:

$$m_{s} = \int \rho_{1} dx \, V_{1} \cos \theta_{1} = \rho_{3} \Delta x \, V_{0} \cos \theta_{0}$$

$$(4)$$

$$m_{s} = \int \rho_{2} dx \, V_{2} \cos \theta_{2} = \rho_{3} \Delta x \, V_{3} \cos \theta_{3}$$

Since a Mach number of 0.4 is obtained in the cascade rig, changes in density across the test section were taken into consideration.

The pressures and flow angles vary as a function of position in the x direction. At station 1 this is due to the presence of the wakes from the inlet guide vanes, Fig. 23. At station 2 this variation is caused by the test blade wakes, Fig. 24. Therefore it is necessary to integrate these values over the area tested. The trapezoidal area method of integration is used in the computer program to evaluate these integrals. The advantage of this method is that unequal intervals may be taken as required. In the blade wake a great number of readings must be taken as the velocity distribution changes rapidly with small changes in x. Between the blade wakes the flow velocity does not change as rapidly and a larger interval can be taken. In the blade wake it is necessary to take readings at intervals as small as 0.01 inch, whereas outside of the wakes it is necessary to take readings only every 0.10 inch.

The trapezoidal area method is as follows:



The variable of integration may be made non-dimensional by :

$$\int = \frac{x}{N_s} ; \qquad N_s = \Delta x \qquad (6)$$

Where: s = blade spacing; N = integer number of blade spacings.

The value for the density of air at stations 0 and 3 is found by:

$$P_{o} = \int_{o}^{o} P_{1} df ; \qquad P_{3} = \int_{o}^{o} P_{2} df \qquad (7)$$

where g and g are found by the following equation:

$$\rho = \frac{1}{a_p^2} \int_0^{\infty} \mathcal{P}(\vartheta - \frac{\Delta \mathcal{P}}{\mathcal{P}}) d\xi \quad ; \quad a_p^2 = \vartheta g \mathcal{R} \mathcal{T}_p \tag{8}$$

Equation (8) of Appendix A is as follows:

$$P_{o} = \frac{1}{\Delta x} \int_{\mathcal{X}_{g}}^{\mathcal{X}_{a}} P_{i} dx + \frac{1}{\Delta x} \int_{\mathcal{X}_{g}}^{\mathcal{X}_{a}} P_{i} V_{i}^{2} \cos^{2} \theta_{i} dx - \frac{1}{P_{o}} \left[\frac{1}{\Delta x} \int_{\mathcal{X}_{g}}^{\mathcal{X}_{a}} P_{i} V_{i} \cos \theta_{i} dx \right]^{2} (9)$$

This may be re-written by substitution in the last two integrals of: $\rho V^2 = 2 \Delta P$ (10)

where:
$$\Delta P = P - p$$

and:

$$PV = \sqrt{2\rho \Delta P} = \sqrt{\frac{2\Delta P}{a_{p}^{2}}(PV - \Delta P)}$$
(11)

giving:

$$P_{o} = \int_{0}^{\infty} P_{o} d\xi + 2 \int_{0}^{\infty} \Delta P \cos^{2}\theta_{o} d\xi - \frac{2}{\rho_{o}} a_{p}^{2} \left[\int_{0}^{\infty} \sqrt{\Delta P(\gamma P - \Delta P)} \cos \theta d\xi \right]^{2} (12)$$

Equation (10) of Appendix A is:

$$\tan \Theta_{0} = \frac{P_{0} \Delta \chi}{2} \frac{\int_{\chi_{R}}^{\pi_{R}} P_{1} V_{1}^{2} \sin (2\theta_{1}) d\chi}{\left[\int_{\chi_{R}}^{\pi_{R}} P_{1} V_{1} \cos \theta_{1} d\chi\right]^{2}}$$
(13)

Using (10) and (11) this may be re-written in the form:

$$\Theta_{o} = Tan^{-1} \frac{P_{o} \alpha_{p}^{2}}{2} \frac{\int_{o} \Delta P_{i} \sin(2\theta_{i}) df}{\left[\int_{o}^{i} \sqrt{\Delta P_{i} (P_{i} \mathcal{Y} - \Delta P_{i}) \cos \theta_{i} df}\right]^{2}}$$
(14)

Equation (11) of appendix A is:

$$V_{o}^{2} = \frac{1}{4} \frac{\left[\int_{\mathcal{X}_{g}}^{\mathcal{T}_{g}} P_{i} V_{i}^{2} \sin(2\theta_{i}) dy\right]^{2}}{\left[\int_{\mathcal{X}_{g}}^{\mathcal{T}_{g}} P_{i} V_{i} \cos\theta_{i} dy\right]^{2}} + \frac{1}{P_{o}^{2} \Delta \pi^{2}} \left[\int_{\mathcal{X}_{g}}^{\mathcal{T}_{g}} P_{i} V_{i} \cos\theta_{i} dy\right]^{2}$$
(15)

Again using (10) and (11) this gives:

$$V_{o}^{2} = \frac{a_{p}^{2}}{2} \frac{\left[\int_{0}^{\infty} \Delta P_{i} \sin\left(2\theta_{i}\right) d\xi\right]^{2}}{\left[\int_{0}^{\omega} \left[\left[\int_{0}^{\omega} \Delta P_{i} (\gamma P_{i} - \Delta P_{i}) \cos\theta_{i} d\xi\right]^{2} + \frac{2}{a_{p}^{2} \int_{0}^{\omega} \left[\int_{0}^{\omega} \left[\Delta P_{i} (\gamma P_{i} - \Delta P_{i}) \cos\theta_{i} d\xi\right]^{2}\right]^{2}} + \frac{2}{a_{p}^{2} \int_{0}^{\omega} \left[\int_{0}^{\omega} \left[\Delta P_{i} (\gamma P_{i} - \Delta P_{i}) \cos\theta_{i} d\xi\right]^{2}\right]^{2}}$$
(16)

Total pressure at station 0 is found from:

$$P_{o} = \beta_{o} + \frac{\beta_{o}}{2} V_{o}^{2}$$
(17)

The above procedure is repeated to find the conditions at station 3 by substituting the subscript 3 for 0 and 2 for 1.

These equations determine the total and static pressures, the angle and the velocity of flow at stations far upstream and downstream from the test section where uniform conditions exist.

The forces on the blades per unit blade height were found as follows: From equation (1):

$$d\vec{R} = m_{so}(\vec{i}V_{0}\sin\theta_{0} + \vec{j}V_{0}\cos\theta_{0}) - m_{S3}(\vec{i}V_{3}\sin\theta_{3} + \vec{j}V_{3}\cos\theta_{3}) (18) + \vec{j}(\rho_{0} - \rho_{3})S$$

Resolving into components:

$$dR_{*} = m_{so} V_{o} \sin \theta_{o} - m_{s_{3}} V_{s} \sin \theta_{3} = \int_{0}^{o} S V_{o}^{2} \cos \theta_{o} \sin \theta_{o} - \int_{3}^{o} S V_{3} \cos \theta_{s} \sin \theta_{3}$$
(19)

$$dR_{a} = m_{s0} V_{0} \cos \theta_{0} - m_{s3} V_{3} \cos \theta_{3} + (p_{0} - p_{3}) S$$

$$dR_{a} = P_{0} S V_{0}^{2} \cos^{2} \theta_{0} - P_{3} S V_{3}^{2} \cos^{2} \theta_{3} + (p_{0} - P_{3}) S$$
(20)

Therefore:

$$dR_{u} = \frac{1}{2} S(P_{0}V_{0}^{2} \sin 2\theta_{0} - P_{3}V_{3}^{2} \sin 2\theta_{3})$$
(21)

$$\int R_{a} = S(p_{0}V_{0}^{2}c_{re}^{2}\theta_{0} - p_{3}V_{3}^{2}c_{re}^{2}\theta_{3} + p_{0} - p_{3})$$
(22)

The lift and drag forces on the blade are determined from these values. This is accomplished by solving for $\overline{V_{\infty}}$.

For individual airfoils V_{∞} is the velocity at an infinite distance from the airfoil. In a rectilinear cascade V_{∞} is the mean vectorial velocity shown on Fig. 25 and represented by:

$$\overline{V}_{\infty} = \frac{1}{2} \left(\overline{V}_0 + \overline{V}_3 \right); \quad V_a = V \sin \theta \quad ; \quad V_u = V \cos \theta \quad (23)$$

$$V_{\infty} = \frac{1}{2} \sqrt{(Y_{ao} + V_{a3})^2 + (V_{uo} + V_{u3})^2}$$
(24)

$$\Theta_{\infty} = Tan^{-1} \frac{(V_{uo} + V_{u3})}{(V_{a0} + V_{a3})}$$
(25)



Lift is the component of the force acting on the blades that is perpendicular to $\overline{V_{\infty}}$. Drag is the force component that is parallel to $\overline{V_{\infty}}$. From Fig. 25, β_{∞} is defined as the angle from the horizontal to $\overline{V_{\infty}}$. The angle θ between \overline{dR} and its horizontal component is known from equation (3). The sum of the lift and drag vectors must equal the force \overline{dR} on the blade.

$$\begin{aligned}
\vec{R} &= \vec{L} + \vec{D} \\
\vec{L} &= \vec{dR} \sin (\beta_{\infty} - \Theta) \\
\vec{D} &= \vec{dR} \cos (\beta_{\infty} - \Theta)
\end{aligned}$$
(26)

Therefore all of the vectors on Fig. 25 are known.

In order to non-dimensionalize lift and drag the coefficients of lift and drag are used. The coefficient of lift is defined as the lift force divided by the product of the test blade chord and dynamic pressure. The dynamic pressure is based on the velocity V_{∞} .

$$C_{L} = \frac{L}{\frac{i}{z} \rho V_{\infty}^{2} c}$$
(27)

Likewise:

$$C_{\rm D} = \frac{D}{\frac{1}{2} \rho V_{\rm so}^2 c}$$
(28)

where:

$$\mathcal{P} = \frac{1}{2} \left(\mathcal{P}_{0} + \mathcal{P}_{3} \right)$$

NACA in Ref. 3 bases dynamic pressure on the velocity V_0 and the density ρ_0 . However, since V_{∞} is the factor that fixes the lift and drag vectors it is more logical to define the coefficients by this value.

The total loss in the cascade is the difference in total pressures far upstream and downstream of the cascade. This total loss is given by the loss coefficient 5. The coefficient 5 includes the mixing and frictional losses in stationary cascades for both accelerated and decelerated flows. Therefore the total loss of the cascade, as developed in Appendix A, is defined as:

$$5 = \frac{\left(P_{o} + \frac{1}{2}P_{o}V_{o}^{2}\right) - \left(P_{3} + \frac{1}{2}P_{o}V_{3}^{2}\right)}{\left(P_{o} + \frac{1}{2}P_{o}V_{o}^{2}\right) - P_{3}}$$
(29)

6. Discussion of Results

The results calculated from the test run are listed in Table II. These results are compared with NACA TN 3802 results, Ref. 3, in Table III.

Table II is the output of the computer program for the test run. The computer program is included in Appendix B. Table II indicates that the flow accelerates slightly from 312.49 ft./sec. ahead of the test blades to 322.32 ft./sec. after the test blades. The inlet flow angle was 43.63 deg. and the outlet flow angle was -46.18 deg. The value for V_{∞} was 224.86 ft./sec. at an angle of -2.16 deg. These vectors are shown on Fig. 25.

The total force per test blade is 120.58 lb./ft. This total force is composed of a lift force of 120.07 lb./ft. and a drag force of 11.10 lb./ft.

The coefficient of lift is 2.8223 and the coefficient of drag is 0.2609. These coefficients are based on V and the average of the fluid densities at stations 0 and 3.

The flow deflection is the difference between the inlet flow angle and the outlet flow angle. For the test run this value was 43.63 deg. minus -46.18 deg. or 89.81 deg.

Fig. 15(5) of Ref. 5 is a plot of the variation of velocity

TABLE II

CASCADE LABORATORY U.S. NAVAL POSTGRADUATE SCHOOL

TEST RUN NO. 100

BAROMETRIC PRESS. 29.98 IN. HG.

AMBIENT TEMP. 66.5 DEG. F.

| VO | THETA O | V3 | THETA 3 | V INF | THETA INF | BETA |
|----------|---------|----------|---------|----------|-----------|-------|
| (FT/SEC) | (DEG) | (FT/SEC) | (DEG) | (FT/SEC) | (DEG) | (DEG) |
| 312.49 | 43.63 | 322,32 | -46.18 | 224.86 | -2.16 | 92.17 |

| AXIAL TANGENTIAL TUTAL | |
|---------------------------------------|---------|
| FORCE FORCE FORCE THETA LIFT | DRAG |
| (LB/FT) (LB/FT) (LB/FT) (DEG) (LB/FT) | (LB/FT) |
| | |
| 15.62 119.56 120.58 7.44 120.07 | 11.10 |

NON-DIMENSIONAL COEFFICIENTS LIFT DRAG LOSS

2.8223 .2609 .1334

LEGEND

V0 = V₀, velocity at station 0 THETA 0 = Θ_0 , angle between \overline{j} and \overline{V}_0 V3 = V₃, velocity at station 3 V INF = V₀, vectorial mean of \overline{V}_0 and \overline{V}_3 THETA INF = Θ_0 , angle between \overline{j} and \overline{V}_{∞} BETA = A_0 , angle between \overline{i} and \overline{V}_{∞} AXIAL FORCE = $d\overline{R}_a$, \overline{j} component of force per blade TANGENTIAL FORCE = $d\overline{R}_u$, \overline{i} component of force per blade TOTAL FORCE = $d\overline{R}_a$, force per blade THETA = Θ , angle between \overline{i} and $d\overline{R}$ THETA 3 = Θ_3 , angle between \overline{j} and \overline{V}_3

TABLE III

COMPARISON OF RESULTS

| | TN 3802 | TN 3802 | USNPGS |
|-----------------|---------------------|---------------------|---------------------|
| o | 1.8 | 1.5 | 1.415 |
| e _l | 45° | 450 | 43.63° |
| θ _c | 950 | 950 | 950 |
| θ | 94.40 | 93.00 | 89.81° |
| C_{Lo} | 1.239 | 1.362 | 1.453 |
| C _{Do} | .0236 | .0227 | .1343 |
| Re | 0.5x10 ⁶ | 0.5x10 ⁶ | 1.4x10 ⁶ |

LEG END

 θ = Turning angle of air

 $\Theta_c = Angle between tangents to camber line at leading and trailing edges$

coefficient with respect to flow deflection. Velocity coefficient Ψ is related to the loss coefficient \mathcal{F} by: $\mathfrak{K} = 1 - \Psi^2$

Using this relation Ref. 5 gives a loss coefficient of 0.1338 for wide blades. Table II indicates a loss coefficient of 0.1334.

In the comparison on Table III it must be taken into consideration that slightly different test conditions existed between the present cascade and the cascade of Ref. 3. The solidity of the present cascade was 1.415 whereas 1.5 is the closest value with which a comparison could be made with TN 3802. The 1.8 solidity values of TN 3802 are included to show the effect of solidity on the various parameters.

From Table III the inlet flow angle of the present cascade was 43.63 deg. This is 1.37 deg. less than the inlet angle of TN 3802. A lower inlet angle indicates a lower angle of attack on the blades. The turning angle of the flow was also found to be 3.19 deg. less than the 1.5 solidity case and 4.59 deg. less than the 1.8 solidity case. For a given inlet angle, cascades of high solidity have greater turning angles than low solidity cascades as the flow is guided more exactly. Fig. 51 of Ref. 3 indicates the deviation of the leaving flow direction for this blade profile at 1.5 and 1.8 solidity. A greater angle between leaving flow direction and trailing-edge meanline direction exists for the lower solidity. For a leaving flow angle of 46.18 deg. this reference gives a deviation of 5.5 deg. for solidity of 1.5. This indicates that the actual turning angle of the TN 3802 blade should be 89.5 deg. The turning angle for a cascade of 1.415 solidity should be less than 89.5 deg. However, a turning angle of 89.81 deg. was obtained. This discrepancy may be due to the interpolation error in reading Fig. 51 of Ref. 3.

The lift and drag coefficients are based on the upstream dynamic

pressure in order to compare them with values from TN 3802.

The lift coefficients of TN 3802 are 1.239 for a solidity of 1.8 and 1.362 for a solidity of 1.5. A lift coefficient of 1.453 for a solidity of 1.415 was obtained for the present cascade.

At given inlet flow angles and solidities, the flow deflection is a measure of blade loading. Table III indicates that with decreasing solidity there is a decrease in turning angle and an increase in the lift coefficient. This increase in the lift coefficient appears to be a discrepancy as a smaller turning angle should produce less lift. However, separation must occur at the higher turning angles and thus reduce the lift coefficient.

The coefficients of drag from Table III are 0.0236 for 1.8 solidity and 0.0227 for 1.5 solidity of the NACA cascade. The coefficient of drag for the present cascade is 0.1343.

A cascade with a large number of blades will have large frictional effects on the surfaces of these blades. In other words, skin friction drag increases as solidity increases. This is shown on Table III where the drag coefficient increases from 0.0227 to 0.0236 as solidity increases from 1.5 to 1.8. However, the drag coefficient for the present cascade is 0.1343 for a solidity of 1.415.

To analyze these results it is necessary to examine the two cascade rigs involved. In the NACA cascade the boundary layer was removed for reasons stated in Section 2. In the present cascade the boundary layer was not removed.

The removal of the boundary layer makes the flow more nearly ideal and reduces the secondary losses. Secondary losses are caused by the interaction of the blade ends with the boundary layer on the walls. The

secondary losses predominate over all other losses at large angles of deflection.

Fig. 24 is a plot of $1/2 \oint_2 V_2^2$. This figure shows the large velocity deficit that is present after the test blades. This velocity deficit is included in the loss coefficient 5 of Table II. The loss coefficient is related to the drag coefficient by equation 13(11) of Ref. 5. This equation is for compressor blades but is also applicable to turbine blades. Equation 13(11) is rewritten as:

$$C_{D} = \frac{5}{0} \frac{\cos^{3} \Theta_{0}}{\cos^{3} \Theta_{3}}$$

Using this equation with the results of Table II gives a coefficient of drag of 0.1926.

The test Reynolds number for NACA TN 3802 was 0.5×10^6 . The test Reynolds number for the present cascade was 1.4×10^6 . This is a small variation in Reynolds number and could not be a factor in the comparison of results.

Fig. 23 is a graph showing the variation of dynamic pressure (P_1-p_1) with respect to x. These measurements were taken on the lower traverse (station 1). The curve passes through a maximum and minimum dynamic pressure at two inch intervals. The minimum q values represent the wakes of the inlet guide vanes which are spaced two inches apart.

Fig. 24 is a graph of dynamic pressure (P_2-P_2) variation with respect to distance in the x direction for the upper traverse. On this figure the minimum and maximum values of q occur at six inch intervals. The minimum dynamic pressure is the center of a wake from the test blades which are spaced six inches apart. The spread of the points on Fig. 24 is an indication of the non-uniformity of the flow after passing

through the test blades. This non-uniformity is more pronounced on Fig. 24 than on Fig. 23 as the measurements are taken closer to the blades. Fig. 23 then shows how the flow tends to become more uniform as it progresses downstream.

Fig. 26 is a plot of static pressure variation with respect to the x distance. This figure shows the variation for both the upper and lower traverses. The readings of the static pressure by the measuring equipment are accurate to ± 0.1 inch of water pressure. The maximum variation of static pressure on the lower level was 0.2 inch of water from 2.2 to 2.4 inches. In other words the possible error is 50% of the maximum variation. For this reason the curve does not present a good picture of the static pressure distribution. However it does show that there is no static pressure gradient throughout the test section.

The upper level curve of Fig. 26 indicates two areas of high static pressures at 4.5 inches and 10.5 inches. These peaks indicate the wakes of the test blades or areas of reduced dynamic pressure. The peaks occur at the test blade spacing of six inches. The variations on the pressure side of the wake are due to the wakes of the inlet guide vanes.

Fig. 27 is a graph of flow angle variation with respect to x distance. The lower graph on this figure shows the variation of flow angle as measured by the lower traverse. The accuracy of the measuring equipment for flow angle is ± 0.2 deg. The maximum variation of the flow angle was 1.2 deg. from 44.4 to 43.2. This is a 16.7% error and reduces the value of the graph to a point where no definite statements may be made.

The upper level indicates maximum flow angles at a six inch

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interval. This again is due to the wakes following the test blades. As the value for x increases the flow angles are read from the pressure side of the blade to the suction side. The slope of the flow angle curve is symmetric with respect to the peak. This indicates that the flow experiences the same angle change with respect to distance on each side of the wake. The slight variations in the flow angle on the pressure side of each peak at 2.0 and 8.0 inches are due to the inlet guide vane wakes.

The variations in the outlet flow angles and the static pressures behind the test blades show the effects of the inlet guide vane wakes. This also is the cause of the scatter of points in Fig. 24, the plot of downstream dynamic pressure. This shows the advantage of this unique test rig in approximating the flow obtained in turbomachinery as compared to conventional cascade rigs.

In section 4, Static Pressure Distribution, it is indicated that the inlet guide vanes were set at 42 deg. in order to obtain the acceptable static pressure distribution.

With the vanes set at 42 deg. the inlet flow angle was found to be 43.63 deg. The relation: $\alpha_2^* = \cos^{-1} \%$

as mentioned in section 4 indicated that the guide vanes should be set at 45 deg. This relation can be used to find a first approximation for inlet guide vane settings. However, the position of the inlet guide vanes must be adjusted to give a static pressure variation of less than 2.0% of the test section dynamic pressure and the required inlet flow angle. Therefore, the above equation must be used only as an approximation. The inlet flow angle and static pressure distribution are

factors that must be considered in positioning the inlet guide vanes.

7. General Discussion

During the earlier tests adverse pressure gradients were present in the test section. It was determined that these gradients were caused by flow irregularities that entered the bell mouth from the plenum. The main contributor to the irregularities was the sound attenuating baffles shown in Fig. 4. The size and shape of these baffles are completely unsatisfactory. The blunt upstream end of each baffle is $8\frac{1}{2}$ inches in width, extends from the floor to the ceiling, and is $3\frac{1}{2}$ inches horizontally from the next baffle. The downstream end of each baffle is slightly less blunt but this is not considered sufficient to prevent separation. Also the area reduction caused by the baffles (3.4:1) accelerates the air at a position where deceleration is desired.

These baffles create large turbulent wakes which prevent the velocity from being equal over the cross section of the plenum. They were installed without regard to the aerodynamic conditions as an afterthought which proved to have bad consequences. These baffles are not needed as the noise is produced by the Mach 0.4 air passing through the cascade rig and not by the blower.

During the installation of the screen on the sound reducing baffles it was discovered that the diffusor was warped. This was caused when the outer frames were welded onto the diffusor walls. These frames are located about every three feet along the diffusor. The walls have deflections of about one inch between each pair of frames. This warping causes additional irregularities in the air flow.

The original design included a bell mouth at the exit of the diffusor.
In addition, screen attachments were to be installed on the walls at the bell mouth exit. The installation of several screens at this position would have ensured an equal velocity distribution over the cross section of the plenum. The installation of the sound baffles eliminated the space available for this arrangement.

Installing two wire screens at the entrance of the cascade rig bell mouth gave the best results in smoothing out the flow irregularities. This configuration produced a static pressure variation of only 1.55% of test section dynamic pressure and a total pressure loss of only 8%.

There was some difficulty in setting the inlet guide vanes to the proper angles. This was due to the fact that the vanes are not identical. Some fixture is required to determine which blades have the proper shape and alignment. A similar fixture is required for each test blade shape.

The automatic data logging system performed as desired.

The Friden paper tape punch used in the automatic data logging system uses one inch, 8 level paper tape. This tape has eight bits per frame, i.e. eight holes per line. Standard systems use 7/8 inch, 7 level tape which has seven bits per frame. This was the system used by the Control Data Corporation 1604 computer located at the U.S. Naval Postgraduate School. The computer was modified by the Computer Center to be able to read 8 level tapes.

The computer program used for the data reduction was written in Fortran 60 language. A difficult part of the program was the coding for the reading of the punched paper tape. This part of the program was written in Symbolic Fortran by Mrs. Carol Hayworth, a programer for the

Computer Center. The computer program is listed and explained in Appendix B.

The paper tape supplied with the Friden punch was pink, oiled tape. This proved too transparent for the photo-electric reader to properly decode. To solve this problem the punched paper tape was colored black with a quick drying marker ink. Future data will be punched on black paper tape.

8. Conclusions

The results obtained in this work indicated that the cascade rig will perform its function properly but modifications are necessary. It can be expected, therefore, that the rectilinear cascade test rig will become a valuable research toolwhich can be used to establish more precise design data for existing and future cascades. Such data will make it possible to design axial-flow machines more accurately and save considerable development time.

The data logging system proved to be a labor and time conserving device. For example, an average test run would cover at least three blade spacings. The x distance involved would vary depending on the blade spacing. Measurements of distances, flow angles, and total and static pressures must be taken on an average every 0.05 inch. This would require roughly 700 readings of each value for the six inch blade spacing on each plane in the y direction. Without the punched paper tape this data must be transcribed to punched cards for the computer input. This would involve about 30 man hours of labor to do what can be accomplished in two man hours with the data logging system.

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9. Recommendations

It is recommended that additional modifications be made to the cascade rig in order to improve the flow in the test section. Static pressure variation of less than one percent should be obtained with a zero pressure gradient.

The following are the steps required to obtain this improved flow: Eliminate the effects of the warped diffusor. This could be accomplished by lining the diffusor with plywood walls. These wooden walls could be attached to the metal walls with countersunk bolts.

Then the diffusor should be modified with an addition of a hyperbolic divergence bell mouth. This was in the original design but was eliminated in the latter stages of planning when it was decided to install the sound baffles.

At the exit of the bell mouth a large number of screens could be attached to the side walls of the chamber. These screens would offer greater resistance to the higher velocities and thus smooth out disturbances in the flow. With these modifications the air would leave the modified diffusor section at a lower velocity and more uniformly.

These modifications require the removal of the sound baffles in the plenum chamber. Diffusors of the type described have been used with success in wind tunnels. Since the sound baffles have a very unfavorable aerodynamic shape, their removal would produce a considerable improvement of the flow conditions ahead of the cascade. These baffles are not required as the noise is created by the Mach 0.4 air passing through the cascade rig and not by the blower.

To measure the total pressure in the plenum chamber it is recommended that a copper tube of 1/8 in. inner diameter be installed in

£3

the conduit which is used for the total temperature probe.

Fixtures should be constructed to check the profiles and alignments of the test blades and the inlet guide vanes. In each case the profile should be checked at several positions along the blade height.

For low and medium speed tests, the data logging system should be recalibrated to measure the static pressures to 0.01 inches of water. This will provide more accurate results and better plots of the static pressure distribution.

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Appendix A

x = distance along cascade axis

 $\Theta =$ flow angle

DERIVATION OF EQUATIONS

1. SYMBOLS

- p = static pressure
- P = total pressure
- $q = \frac{1}{2} \rho V^2$
- V = velocity

Indices:

- (0) far ahead of cascade
- (1) measuring plane ahead of cascade
- (2) " after cascade
- (3) far aft of cascade
- 2. CONDITIONS FAR AHEAD OF CASCADE



Measured Values: P_i ; p_j ; θ_i as functions of x To be determined: P_o , p_o , θ_o , V_o From Eq. 5(21) Ref. 5 for $\vec{R} = 0$, ignoring shear stress integrals γ

$$\int_{(0)} dm_{s} \vec{V}_{o} - \int_{(1)} dm_{s} \vec{V}_{i} + \vec{j} \int_{(0)} p_{o} dS - \vec{j} \int_{(1)} p_{i} dS = 0 \qquad (1)$$

for A h = unity perpendicular to plane of Fig. 1

$$m_{s}\overline{V_{o}} - \int_{x_{d}}^{x_{d}} p_{,} dy V_{,} \cos \theta_{,} = \overline{j} \int_{x_{d}}^{x_{d}} p_{,} dy - \overline{j} p_{o} \Delta x \qquad (2)$$

From continuity

$$m_{s} = \int_{\mathcal{A}} \rho_{s} d\gamma V_{s} \cos \Theta_{s} = \rho_{o} \Delta \gamma V_{o} \cos \Theta_{o} \quad (3)$$

From Fig. Al

$$\overline{V_o} = \overline{i} \ V_o \ \sin \Theta_o + \overline{j} \ V_o \ \cos \Theta_o \tag{4}$$

$$\overrightarrow{V}_{i} = \overrightarrow{i} \quad V_{i} \quad \sin \theta_{i} \quad + \quad \overrightarrow{j} \quad V_{i} \quad \cos \theta_{i} \tag{5}$$

From Eqs. 2, 4, and 5

$$m_{s}(\vec{x} \vee_{o} \sin \theta_{o} + \vec{j} \vee_{o} \cos \theta_{o}) - \int_{\mathcal{X}_{a}}^{\mathcal{X}_{a}} P_{i} d\varphi \vee_{i} \cos \theta_{i}(\vec{x} \vee_{i} \sin \theta_{i})$$
$$+ \vec{j} \vee_{i} \cos \theta_{i}(\vec{y}) = \vec{j} \int_{\mathcal{X}_{a}}^{\mathcal{X}_{a}} P_{i} d\varphi - \vec{j} \nabla_{o} \Delta \chi \qquad (6)$$

From equality of i components of Eq. 6

$$m_{s} V_{o} \sin \theta_{o} = \int_{\mathcal{H}_{e}}^{\mathcal{H}_{e}} \rho_{i} d\gamma V_{i}^{2} \sin \theta_{i} \cos \theta_{i}$$

$$m_{s} V_{o} \sin \theta_{o} = \frac{1}{2} \int_{\mathcal{H}_{e}}^{\mathcal{H}_{e}} \rho_{i} d\gamma V_{i}^{2} \sin (2\theta_{i})$$

With Eq. 3

$$V_{0} - sin \Theta_{0} = \frac{1}{2} \frac{\frac{\pi_{e}}{\pi_{e}}}{\int_{\mathcal{X}_{e}} \rho_{i} V_{i} \cos \Theta_{i} dy}$$
(7a)

From equality of j components of Eq. 6

or, with Eq. 3

$$V_{o} \cos \theta_{o} = \frac{\frac{\gamma_{e}}{\int P_{i} V_{i}^{2} \cos^{2} \theta_{i} dy + \int P_{i} dy - P_{o} \Delta \chi}{\frac{\gamma_{e}}{\int P_{i} V_{i} \cos \theta_{i} dy}}$$
(7b)

 $\frac{\frac{1}{1}\int_{P_{1}}^{P_{2}}P_{1}V_{1}\cos\theta,d\gamma}{\Delta\gamma} = \frac{\int_{R_{2}}^{P_{2}}P_{1}V_{1}^{2}\cos^{2}\theta,d\gamma + \int_{T_{4}}^{P_{2}}P_{1}d\gamma - P_{0}\Delta\gamma}{\int_{R_{2}}^{T_{2}}P_{1}V_{1}\cos\theta,d\gamma} = \frac{\chi_{2}}{\chi_{4}}$

or

 $\frac{1}{\rho} \left[\int_{-\infty}^{\infty} P_{i} V_{i} \cos \theta_{i} dy \right]_{\Delta X}^{2} = \int_{-\infty}^{\infty} P_{i} V_{i} \cos^{2} \theta_{i} dy + \int_{-\infty}^{\infty} P_{i} dy - P_{o} \Delta x$

Thus

$$P_{0} = -\frac{1}{P_{0}} \left[\frac{1}{\Delta \gamma} \int_{\gamma_{e}}^{\gamma_{e}} P_{i} V_{i} \cos \Theta_{i} d\gamma \right]^{2} + \frac{1}{\Delta \chi} \int_{\gamma_{e}}^{\gamma_{e}} P_{i} V_{i}^{2} \cos^{2} \Theta_{i} d\gamma + \frac{1}{\Delta \gamma} \int_{\gamma_{e}}^{\gamma_{e}} P_{i} d\gamma$$

or



From Eq. 3

$$V_{o} \operatorname{Cre} \Theta_{o} = \frac{i}{\rho_{o} \Delta \gamma} \int P_{i} V_{i} \operatorname{Cre} \Theta_{i} d\gamma \qquad (9)$$

From Eqs. 7a and 9

$$\tan \Theta_{0} = \frac{P_{0} \Delta \gamma}{2} \frac{\int_{\chi_{R}}^{\chi_{R}} P_{1} V_{1}^{2} \sin(2\Theta_{1}) d\gamma}{\left[\int_{\chi_{R}}^{\chi_{R}} P_{1} V_{1} \cos \Theta_{1} d\gamma\right]^{2}}$$
(10)

Since

$$V_o^2 = V_o^2 \sin^2 \theta_o + V_o^2 \cos^2 \theta_o$$

From Eqs. 7a and 9

$$V_{o}^{2} = \frac{1}{4} \frac{\left[\int_{Y_{d}}^{Y_{d}} P_{1}V_{1}^{2} \sin\left(2\Theta_{1}\right) dy \right]^{2}}{\left[\int_{Y_{d}}^{Y_{d}} P_{1}V_{1} \cos\Theta_{1} dy \right]^{2}} + \frac{1}{\frac{1}{P_{o}^{2}} \Delta \chi^{2}} \left[\int_{Y_{d}}^{Y_{d}} P_{1}V_{1} \cos\Theta_{1} dy \right]^{2}}$$
(11)

3. CONDITIONS FAR DOWNSTREAM OF CASCADE



From Eq. 5(21) similar to Eq. 1

$$\int_{(2)} dm_{s} V_{2} - \int_{(3)} dm_{s} V_{3} + \vec{j} \int_{(2)} p_{2} ds - \vec{j} \int_{(3)} p_{3} ds = 0$$

$$\int_{(2)} dm_{s} V_{2} - m_{s} V_{3} + \vec{j} \int_{(2)} p_{2} d\chi - \vec{j} p_{3} \Delta \chi = 0 \quad (12)$$

From continuity

$$m_{s} = \int \int_{2} V_{2} \cos \Theta_{2} d\chi = \int_{3} \Delta \chi V_{3} \cos \Theta_{3}$$
(13)

From Fig. A2

Į

$$\overline{V_2} = i V_2 \sin \theta_2 + j V_2 \cos \theta_2$$
(14)

$$\overrightarrow{V_3} = \overrightarrow{i} \quad V_3 \sin \Theta_3 + \overrightarrow{j} \quad V_3 \cos \Theta_3$$
(15)

From Eqs. 12, 14, 15

$$\begin{array}{l}
\overset{\mathcal{H}_{2}}{\underset{\mathcal{H}_{3}}{\int}} P_{2} d_{\chi} V_{2} \cos \theta_{2} \left(\overline{i} V_{2} \sin \theta_{2} + \overline{j} V_{2} \cos \theta_{2} \right) \\
\overset{\mathcal{H}_{3}}{\underset{\mathcal{H}_{3}}{\int}} - m_{s} \left(\overline{i} V_{3} \sin \theta_{3} + \overline{j} V_{3} \cos \theta_{3} \right) \\
&= + \overline{j} P_{3} \Delta \chi - \overline{j} \int_{\mathcal{H}_{3}}^{\mathcal{H}_{2}} P_{2} d\chi \qquad (16)
\end{array}$$

For the i components of Eq. 16

$$\frac{1}{2} \int_{\mathcal{P}_2}^{\mathcal{P}_2} V_2^2 \sin(2\Theta_2) dy = m_s V_3 \sin \Theta_3$$

$$\frac{1}{\mathcal{P}_2} \int_{\mathcal{P}_2}^{\mathcal{P}_2} V_2^2 \sin(2\Theta_2) dy = m_s V_3 \sin \Theta_3$$

With Eq. 13

$$V_{3} \sin \theta_{3} = \frac{1}{2} \frac{\int_{X_{2}}^{X_{2}} P_{2} V_{2}^{2} \sin (2\theta_{2}) dx}{\int_{X_{2}}^{X_{2}} P_{2} V_{2} \cos \theta_{2} dx}$$
(17)

From the
$$\vec{j}$$
 components of Eq. 16 and with Eq. 13

$$\int_{A_{e}}^{A_{e}} p_{2} V_{2}^{2} \cos^{2} \Theta_{2} d\chi - m_{s} V_{3} \cos \Theta_{3} = p_{3} \Delta \chi - \int_{A_{e}}^{A_{e}} p_{2} d\chi$$

$$V_{3} \cos \Theta_{3} = \frac{\int_{4}^{4} P_{2} V_{2}^{*} \cos^{2}\Theta_{2} d\chi + \int_{\chi_{4}}^{\chi_{4}} P_{2} d\chi - P_{3} \Delta \chi}{\int_{4}^{2} P_{2} V_{2} \cos \Theta_{2} d\chi}$$
(18)

From Eq. 13

$$V_3 \cos \theta_3 = \frac{1}{\rho_3 \Delta \gamma} \int_{\gamma_8}^{\gamma_8} \rho_2 V_2 \cos \theta_2 d\gamma \qquad (19)$$

Similar to Eq. 8

$$\mathcal{P}_{3} = \frac{1}{\Delta \chi} \int_{\mathcal{H}_{4}}^{\mathcal{H}_{e}} \mathcal{P}_{2} d\chi + \frac{1}{\Delta \chi} \int_{\mathcal{H}_{4}}^{\mathcal{H}_{e}} \mathcal{P}_{2} V_{2} \cos^{2} \Theta_{2} d\chi - \frac{1}{\mathcal{P}_{3}} \left[\frac{1}{\Delta \chi} \int_{\mathcal{H}_{2}}^{\mathcal{H}_{e}} \mathcal{P}_{2} V_{2} \cos \Theta_{2} d\chi \right]^{2}$$
(20)

Similar to Eq. 10, from Eqs. 17 and 19

$$\tan \theta_3 = \frac{\beta_3 \Delta \chi}{2} \frac{\int_{\chi_{\pm}}^{\chi_{\pm}} \rho_2 V_2 \sin(2\theta_2) d\chi}{\left[\int_{\chi_{\pm}}^{\chi_{\pm}} \rho_2 V_2 \cos \theta_2 d\chi\right]^2}$$
(21)

From Eqs. 17 and 19

$$V_{3}^{2} = \frac{1}{4} \frac{\left[\int_{R_{2}}^{\pi_{e}} P_{2} V_{2}^{2} \sin\left(2\Theta_{2}\right) dx \right]^{2}}{\left[\int_{R_{2}}^{\pi_{e}} P_{2} V_{2} \cos\Theta_{2} dx \right]^{2}} + \frac{1}{P_{3}^{2} \Delta \gamma^{2}} \left[\int_{R_{p}}^{\pi_{e}} P_{2} V_{2} \cos\Theta_{2} dx \right]^{2}$$
(22)

4. THERMODYNAMIC RELATIONS AT MODERATE MACH NUMBERS.



Fig. A3 P= measured total pressure (psia) p= measured static pressure (psia) V= actual velocity (ft./sec) $C_p = \frac{R}{J} \frac{\delta}{\delta'-1}$ T_p^{*} constant total temperature = plenum temperature (^oR) T= static temperature (^oR)

$$\frac{\sqrt{2}}{2 \sqrt{3} \sqrt{c_{\rho}}} = T_{\rho} - T = T_{\rho} \left[1 - \left(\frac{\sigma}{p}\right)^{\frac{\gamma}{s}} \right]$$
(23)

Let

 $\Delta \mathcal{P} = \mathcal{P} - \mathcal{P} \tag{24}$

With

$$p = P - \Delta P$$

$$V^{2} = 2gR\frac{\delta}{\delta - 1}T_{p}\left[1 - \left(1 - \frac{\Delta P}{P}\right)^{\frac{\delta' - 1}{P}}\right]$$

Let

$$a_p^2 = g R \delta T_p =$$
velocity of sound squared in plenum

Then

$$V^{2} = \frac{2}{\delta' - 1} a_{p}^{2} \left[1 - \left(1 - \frac{41P}{P} \right)^{\frac{\delta' - 1}{\delta'}} \right]$$
(25)

With

$$(1-\gamma)^{\frac{y'-1}{y'}} = 1 - \frac{y'-1}{y'} \gamma + \left(\frac{y'-1}{y'}\right) \left(\frac{y'-1}{y'}-1\right) \frac{\gamma}{2}^{2}$$
$$= 1 - \frac{y'-1}{y'} \gamma - \left(\frac{y'-1}{y'}\right) \left(\frac{1}{2y'}\right) \gamma^{2}$$
$$= 1 - \frac{y'-1}{y'} \gamma \left(1 + \frac{1}{2y'}\gamma\right) \qquad (25^{\circ})$$

Eq. 25 becomes

$$V^{2} = \frac{2}{\gamma-1} \alpha_{p}^{2} \left[1 - 1 + \frac{\delta-1}{\gamma} \frac{\Delta P}{P} \left(1 + \frac{1}{2\gamma} \frac{\Delta P}{P} \right) \right]$$

With

$$K_{v} = 1 + \frac{1}{2y} \frac{\Delta P}{P} = 1 + \frac{1}{2y} \frac{P - P}{P}$$
 (26)

there is

$$V^{2} = \frac{2}{y-1} a_{p}^{2} \left[\frac{y-i}{y} \frac{\Delta P}{P} K_{v} \right] = \frac{2}{y} a_{p}^{2} \frac{\Delta P}{P} K_{v}$$
$$= \frac{2}{y} a_{p}^{2} \left(\frac{P-P}{P} \right) K_{v}$$
(27)

The density $/_{\tau}^{\circ}$ at P, T_p is

$$P_t = \frac{P}{R T_r q}$$

From pv = const or $\rho \propto \rho''$, the density ρ at the static

temperature T and the static pressure p is

$$\rho = \rho_{\overline{\tau}} \left(\frac{\sigma}{\overline{P}}\right)^{\prime s}$$

With Eq. 24

$$P = P_{\tau} \left(\frac{P - \Delta P}{P} \right)^{\prime} = P_{\tau} \left(1 - \frac{\Delta P}{P} \right)^{\prime}$$

From

$$(1-\gamma)'' = 1 - \frac{1}{3'}\gamma + (\frac{1}{3'})(\frac{1}{3'}-1)\frac{\gamma}{2}$$
$$= 1 - \frac{1}{3'}\gamma - \frac{1}{3'}(\frac{3'-1}{23'})\gamma^{2}$$
$$= 1 - \frac{1}{3'}\gamma(1 + \frac{3'-1}{23'}\gamma)$$

Then

$$P = \frac{P}{RT_{\theta}g} \left[1 - \frac{1}{\delta'} \frac{\Delta P}{P} \left(1 + \frac{\delta'-1}{2\delta'} \frac{\Delta P}{P} \right) \right]$$

$$= \frac{P}{RT_{P} g \gamma} \left[\gamma - \frac{\Delta P}{P} \left(1 + \frac{\gamma - 1}{2\gamma} \frac{\Delta P}{P} \right) \right]$$

Let

$$K_{p} = 1 + \frac{\delta' - 1}{2\delta'} \frac{\Delta P}{P} = 1 + \frac{\delta' - 1}{2\delta'} \frac{P - \delta'}{P}$$
(28)

Then

$$\rho = \frac{P}{a_{\rho}^{2}} \left[\gamma - \frac{P - P}{P} K_{p} \right]$$
(29)

For measured values of P and p the velocity V, corresponding to P and p is therefore obtained from Eqs. 27 and 26. The value of ρ that has to be introduced in the equations of paragraphs 2 and 3 is given by Eqs. 28 and 29. To be determined are also the densities ρ_0 and ρ_3 . They can be determined with small error with the relations

$$P_{0} = \frac{\int_{x_{e}}^{x_{e}} \rho_{i} \, d\mu}{\Delta \pi}$$
(30)
$$P_{3} = \frac{\int_{x_{e}}^{x_{e}} \rho_{z} \, d\mu}{\Delta \pi}$$
(31)

A more accurate way or check of these values could be carried out by using Eqs. 30 or 31 for the determination of, say, V_0 , p_0 , θ_0 and then checking Eq. 3 to see whether continuity is satisfied. This can be followed up by a method of successive approximations. A similar procedure could be used for β_3 .

Assume now that p_0 , V_0 , ρ_0 and p_3 , V_3 , ρ_3 have been determined by the equations of paragraphs 2 and 3. To be established is now the loss

coefficient 5 of the cascade.



Fig. A4

By definition, for turbine cascades,

$$\mathcal{J} = \frac{T_3 - T_3^*}{T_p - T_3^*} = \frac{T_p - \frac{\sqrt{3}}{2}gJc_p - T_3^*}{T_p - T_3^*}$$

or

$$5 = 1 - \frac{V_3^2}{2 g J c_p (T_p - T_3^*)}$$
(32)

There is

$$\mathcal{T}_{\mathcal{P}} - \mathcal{T}_{3}^{*} = \mathcal{T}_{\mathcal{P}} \left[I - \left(\frac{\mathcal{C}_{3}}{\mathcal{P}_{0}} \right)^{\frac{\mathcal{C}_{1}}{\mathcal{C}}} \right]$$

Further

$$T_{P} - T_{O} = T_{P} \left[I - \left(\frac{P_{O}}{P_{O}} \right)^{\frac{N-1}{N}} \right] = \frac{V_{O}^{2}}{2 g J c_{P}}$$

hence

$$\frac{P_o}{P_o} = \left[1 - \frac{V_o^2}{2gJc_p T_p}\right]^{\frac{N}{N-1}} = \left[1 - \frac{V_o^2}{2gR\frac{N}{N-1} T_p}\right]^{\frac{N}{N-1}}$$

$$= \left[1 - \frac{N_o^2}{2gR\frac{N}{N-1} T_p}\right]^{\frac{N}{N-1}}$$

Thus

$$P_{o} = P_{o} \left[1 - \frac{y'-1}{z} \frac{V_{o}^{2}}{a_{F}^{2}} \right]^{-\frac{y'}{y'-1}}$$
(33)

With

(

$$\begin{aligned} & -\frac{\delta'}{\delta'-1} = 1 + z \frac{\delta'}{\delta'-1} + \left(-\frac{\delta'}{\delta'-1}\right) \left(-\frac{\delta'}{\delta'-1} - 1\right) \frac{z^2}{2} \\ & = 1 + z \frac{\delta'}{\delta'-1} + \left(\frac{\delta'}{\delta'-1}\right) \left(\frac{2\delta'-1}{\delta'-1}\right) \frac{z^2}{2} \\ & = 1 + z \frac{\delta'}{\delta'-1} \left(1 + \frac{2\delta'-1}{2(\delta'-1)}z\right) \end{aligned}$$

there is

$$P_{o} = \vartheta_{o} \left[1 + \frac{\vartheta_{-1}}{z} \frac{V_{o}^{2}}{a_{p}^{2}} \frac{\vartheta_{-1}}{\vartheta_{-1}} \left(1 + \frac{2\vartheta_{-1}}{2(\vartheta_{-1})} \frac{\vartheta_{-1}}{z} \frac{V_{o}^{2}}{a_{p}^{2}} \right) \right]$$
$$= \vartheta_{o} \left[1 + \frac{\vartheta_{-1}}{z} \frac{V_{o}^{2}}{a_{p}^{2}} \left(1 + \frac{2\vartheta_{-1}}{4} \frac{V_{o}^{2}}{a_{p}^{2}} \right) \right]$$

With

$$K_{p} = 1 + \frac{2\delta' - 1}{4} \frac{V_{o}^{2}}{a_{p}^{2}}$$
(34)

$$P_{o} = P_{o} \left[1 + \frac{\delta'}{2} \frac{V_{o}^{2}}{a_{p}^{2}} K_{p} \right]$$
(35)

Eq. 32 rewritten

$$I - \frac{\zeta}{\zeta} = \frac{V_{3}^{2}}{2g^{2}R\frac{\vartheta}{\vartheta-1}\frac{\tau_{p}}{T_{p}}(1-T_{3}^{*}/T_{p})} = \frac{\frac{y'_{-1}}{z} - \frac{V_{3}^{2}}{\alpha_{p}^{2}}}{1 - T_{3}^{*}/T_{p}}$$
(36)

$$\frac{T_{3}^{*}}{T_{p}} = \left(\frac{\ell_{3}}{\ell_{o}}\right)^{\frac{\chi'_{-1}}{\vartheta}} = \left(\frac{\ell_{3}}{\ell_{o}}\frac{\ell_{o}}{\ell_{o}}\right)^{\frac{\chi'_{-1}}{\vartheta}} = \left(\frac{\ell_{3}}{\ell_{o}}\right)^{\frac{\chi'_{-1}}{\vartheta'}} = \left(\frac{\tau_{3}}{\ell_{o}}\right)^{\frac{\chi'_{-1}}{\vartheta'}} \left[\left(\frac{\tau_{o}}{T_{p}}\right)^{\frac{\chi'_{-1}}{\vartheta'}}\right]^{\frac{\chi'_{-1}}{\vartheta'}}$$

$$= \left(\frac{\ell_{3}}{\ell_{o}}\right)^{\frac{\chi'_{-1}}{\vartheta'}} \frac{T_{o}}{T_{p}} = \left(\frac{\ell_{3}}{\ell_{o}}\right)^{\frac{\chi'_{-1}}{\vartheta'}} \frac{T_{p} - \frac{V_{o}^{2}}{2g^{3}} \zeta_{p}}{T_{p}}$$

$$= \left(\frac{\ell_{3}}{\ell_{o}^{2}}\right)^{\frac{\chi'_{-1}}{\eta'}} \left[1 - \frac{\chi'_{-1}}{2} - \frac{V_{o}^{2}}{\alpha_{p}^{2}}\right]$$
(37)

Let

$$P_{0} - P_{3} = \Delta P_{3}$$

$$P_{3} = P_{0} - \Delta P_{3}$$

$$\frac{P_{3}}{P_{0}} = 1 - \frac{\Delta P_{3}}{P_{0}}$$
(38)

From Eq. 25

$$\left(\frac{\mathcal{P}_3}{\mathcal{P}_0}\right)^{\frac{\lambda'-1}{\ell'}} = \left(1 - \frac{\Delta \mathcal{P}_3}{\mathcal{P}_0}\right)^{\frac{\lambda'-1}{\ell'}} = 1 - \frac{\lambda'-1}{\lambda'} \frac{\Delta \mathcal{P}_3}{\mathcal{P}_0} \left(1 + \frac{1}{2\lambda'} \frac{\Delta \mathcal{P}_3}{\mathcal{P}_0}\right)$$

With

$$K_{s} = 1 + \frac{1}{28} \frac{\Delta P_{s}}{P_{o}} = 1 + \frac{1}{28} \left(\frac{P_{o} - P_{3}}{P_{o}} \right)$$
 (39)

$$\left(\frac{\rho_3}{\rho_0}\right)^{\frac{\beta-1}{\beta}} = 1 - \frac{\beta'-1}{\beta'} \frac{\Delta P_s}{\rho_0} K_s = 1 - \frac{\beta'-1}{\beta'} \left(\frac{\rho_0 - \rho_3}{\rho_0}\right) K_s$$
(40)

Then, from Eqs. 36, 37, and 40

$$I - S = \frac{\frac{\delta - I}{2} - \frac{V_3^2}{\alpha_p^2}}{I - \left[I - \frac{\delta - I}{\delta} - \frac{\delta P_3}{\rho_0} K_3\right] \left[I - \frac{\delta - I}{2} - \frac{V_0^2}{\alpha_p^2}\right]}$$

$$= \frac{\frac{\sqrt[3]{1-1}}{2} \frac{\sqrt[3]{2}}{a_{p}^{2}}}{1-1+\frac{\sqrt[3]{1-1}}{\sqrt[3]{p}} \frac{\Delta Ps}{P_{o}} K_{s} + \frac{\sqrt[3]{1-1}}{2} \frac{\sqrt[3]{1-1}}{a_{p}^{2}} - \frac{(\sqrt[3]{1-1})^{2}}{2\sqrt[3]{1-1}} \frac{\Delta Ps}{P_{o}} K_{s} \frac{\sqrt[3]{1-1}}{a_{p}^{2}}}{\frac{\sqrt[3]{1-1}}{2\sqrt[3]{1-1}} \frac{\sqrt[3]{1-1}}{\sqrt[3]{1-1}} \frac{\Delta Ps}{P_{o}} K_{s} \frac{\sqrt[3]{1-1}}{a_{p}^{2}}}{\frac{\sqrt[3]{1-1}}{2\sqrt[3]{1-1}} \frac{\sqrt[3]{1-1}}{\sqrt[3]{1-1}} \frac{\Delta Ps}{P_{o}} K_{s} \frac{\sqrt[3]{1-1}}{a_{p}^{2}}}{\frac{\sqrt[3]{1-1}}{2\sqrt[3]{1-1}} \frac{\sqrt[3]{1-1}}{\sqrt[3]{1-1}} \frac{\sqrt[3]{1-1}}{\sqrt[3]{1-1}} \frac{\sqrt[3]{1-1}}{\sqrt[3]{1-1}} \frac{\Delta Ps}{P_{o}} K_{s} \frac{\sqrt[3]{1-1}}{\sqrt[3]{1-1}} \frac{\sqrt[3]{1-1}}{\sqrt[3]{1-1}} \frac{\sqrt[3]{1-1}}{\sqrt[3]{1-1}} \frac{\Delta Ps}{P_{o}} K_{s} \frac{\sqrt[3]{1-1}}{\sqrt[3]{1-1}} \frac{\sqrt[3]{1-1}}{\sqrt[3]{1-1}} \frac{\sqrt[3]{1-1}}{\sqrt[3]{1-1}} \frac{\sqrt[3]{1-1}}{\sqrt[3]{1-1}} \frac{\Delta Ps}{P_{o}} \frac{\sqrt[3]{1-1}}{\sqrt[3]{1-1}} \frac$$

$$\begin{split} & \mathcal{J} = I - \frac{\frac{\delta^{l-1}}{2} - \frac{V_3^2 / \alpha_P^2}{\alpha_P^2}}{\frac{\delta^{l-1}}{\delta'} - \frac{\Delta P_3}{\rho_0} - \frac{\delta^{l-1}}{2} - \frac{\delta^{l-1}}{2\delta'} - \frac{\delta P_3}{\rho_0} - \frac{\delta P_3}{\delta'} - \frac{\delta P_3}{\delta'}$$

or

$$\mathcal{S} = \frac{\frac{\cancel{\delta}^{-1}}{\cancel{\delta}^{1}} \frac{\triangle \cancel{\delta}^{2}}{\cancel{\delta}^{2}} K_{s} \left(1 - \frac{\cancel{\delta}^{-1}}{2} \frac{\cancel{V_{o}}^{2}}{\cancel{\alpha_{p}}^{2}}\right) + \frac{\cancel{\delta}^{-1}}{2} \frac{1}{\cancel{\alpha_{p}}^{2}} \left(\cancel{V_{o}}^{2} - \cancel{V_{s}}^{2}\right)}{\frac{\cancel{\delta}^{-1}}{\cancel{\delta}^{1}} \frac{\triangle \cancel{\delta}^{2}}{\cancel{\delta}^{2}} K_{s} \left(1 - \frac{\cancel{\delta}^{-1}}{2} \frac{\cancel{V_{o}}^{2}}{\cancel{\alpha_{p}}^{2}}\right) + \frac{\cancel{\delta}^{-1}}{2} \frac{\cancel{V_{o}}^{2}}{\cancel{\alpha_{p}}^{2}}$$

or

$$S = \frac{\frac{\beta_{o} - \beta_{3}}{\beta_{o}} K_{s} \left(1 - \frac{\gamma_{-1}}{2} \frac{V_{o}^{2}}{a_{p}^{2}}\right) + \frac{\gamma_{0}}{2} \frac{1}{a_{p}^{2}} \left(V_{o}^{2} - V_{3}^{2}\right)}{\frac{\beta_{o} - \beta_{3}}{\beta_{o}^{2}} K_{s} \left(1 - \frac{\gamma_{-1}}{2} \frac{V_{o}^{2}}{a_{p}^{2}}\right) + \frac{\gamma_{0}}{2} \frac{V_{o}^{2}}{a_{p}^{2}}}$$
(41)

where K_s is given by Eq. 39. It can be noted that in this expression there occur only p_0 , p_3 , V_0 , and V_3 .

5. RELATIONS FOR INCOMPRESSIBLE FLOWS From Eqs. 26 and 27 for $K_v = 1$

$$V^{2} = 2 g R T_{p} \left(\frac{P - p}{P} \right) = \frac{2}{P_{t}} \left(P - p \right)$$

$$= \frac{2 g R T_{p} Y}{Y P} \left(P - p \right) = \frac{2}{Y} a_{p}^{2} \frac{\left(P - p \right)}{P}$$

$$(42)$$

where

In this equation P is the total pressure at a station and T_p the plenum temperature.

From Eq. 28 and 29 for F_D 1, with Eq. 43

$$\rho = \frac{P}{a_{p}^{3}} \left[\gamma - \frac{P - p}{P} \right]$$
(44)

From Eqs. 34 and 35, with $K_p = 1$

$$P_{o} = P_{o} + P_{o} \frac{y}{z} \frac{V_{o}}{gy RT_{p}} = P_{o} + \frac{V_{o}^{2}}{z} \frac{P_{o}}{gRT_{p}}$$

It is seen that the total pressure P_{O} is obtained from p_{O} and V_{O} by

$$P_o = \rho_{so} \frac{V_o}{2} + \rho_o \tag{45}$$

where

$$P_{so} = \frac{P_o}{g R T_p}$$
(46)

Hence it is more accurate to take the static density \int_{SO}^{∞} for the determination of the total pressure.

From Eqs. 39 and 41 for $K_s = 1$

$$\mathcal{J} = \frac{\frac{P_{o} - P_{3}}{P_{o}} \left(1 - \frac{\mathcal{Y}_{-1}}{2} \frac{V_{o}^{2}}{g \mathcal{Y} R T_{p}}\right) + \frac{\mathcal{Y}}{2} \frac{1}{g \mathcal{Y} R T_{p}} \left(Y_{o}^{2} - V_{3}^{2}\right)}{\frac{P_{o} - P_{3}}{P_{o}} \left(1 - \frac{\mathcal{Y}_{-1}}{2} \frac{V_{o}^{2}}{g \mathcal{Y} R T_{p}}\right) + \frac{\mathcal{Y}}{2} \frac{V_{o}^{2}}{g \mathcal{Y} R T_{p}}}{\frac{V_{o}}{g \mathcal{Y} R T_{p}}}$$

$$=\frac{(P_{o}-P_{3})(I-\frac{\vartheta'-I}{2}\frac{V_{o}^{2}}{\vartheta'RT_{P}})+\frac{\vartheta_{o}}{2\vartheta'RT_{P}}(V_{o}^{2}-V_{3}^{2})}{(P_{o}-P_{3})(I-\frac{\vartheta'-I}{2}\frac{V_{o}^{2}}{\vartheta'RT_{P}})+\frac{P_{o}}{2\vartheta'RT_{P}}V_{o}^{2}}$$

On page 80 it is seen that

$$\frac{T_{o}}{T_{p}} = 1 - \frac{y_{-1}}{2} \frac{V_{o}^{2}}{a_{p}^{2}} = 1 - \frac{y_{-1}}{2} \frac{V_{o}^{2}}{gy_{R}T_{p}}$$

Hence

$$\mathcal{J} = \frac{(\mathcal{P}_{o} - \mathcal{P}_{3})\frac{T_{o}}{T_{P}} + \frac{\mathcal{P}_{o}}{2\mathcal{P}RT_{P}}\left(V_{o}^{2} - V_{3}^{2}\right)}{(\mathcal{P}_{o} - \mathcal{P}_{3})\frac{T_{o}}{T_{P}} + \frac{\mathcal{P}_{o}}{2\mathcal{P}RT_{P}}\left(V_{o}^{2}\right)}$$

With

$$\int_{0}^{0} = \frac{P_{0}}{9 R T_{0}}$$
(47)

$$\mathcal{J} = \frac{\left(\mathcal{P}_{o} + \frac{1}{2}\mathcal{P}_{o} \vee_{o}^{2}\right) - \left(\mathcal{P}_{3} + \frac{1}{2}\mathcal{P}_{o} \vee_{3}^{2}\right)}{\left(\mathcal{P}_{o} + \frac{1}{2}\mathcal{P}_{o} \vee_{o}^{2}\right) - \mathcal{P}_{3}}$$
(48)

Hence to be consistent the static density of Eq. 47 should be introduced here as well as in the relations of paragraphs 2 and 3.

APPENDIX B

COMPUTER PROGRAM

Introduction

Program sheets B-1 through B-6 list the computer program which was used in conjunction with the CDC 1604 computer to reduce the test data. This program utilizes Fortran 60 language for the main program and Symbolic Fortran for the subroutine.

The main program consists of two parts. The subroutine sends integers read from the punched paper tape to the first part of the main program. In this part of the program the integers are identified and stored in the proper arrays of the computer memory. The second part of the program then evaluates the input data and computes the lift, drag, and loss coefficients.

Input Procedures

Manual inputs to the punch paper tape include the run number, barometric pressure, plenum temperature, number of blade spacings surveyed, and length of blade spacing. In addition, a code is manually punched to indicate the end of the input data. A five digit number is available for manual entry. The left most digit is used as a code to identify the manual input data. The remaining four digits are available for the value of the data. The coding of the manual input data is as follows:

- 1 Run Number (3 digit number)
- 2 Barometric Pressure (in. Hg. to 0.01 in.)
- 3 Plenum Temperature (deg. F. to 0.1 deg.)
- 4 Number of Spacings Surveyed (1 digit number)

5 Length of Spacing (1 or 2 digit number)

9 End of Input Data

Manual information is punched by depressing the manual Data record button after rotating the dials of the manual control to the proper code and value. In the case where less than all four digits are required the value must be inserted to the extreme right. For example, a temperature of 66.5 degrees is inserted as 30665.

The data collected automatically by the data logging system is recorded whenever the scan button is depressed. A six digit number is used for this automatic data. The left three integers identify the information and the right three integers indicate the value of the data.

Listed below is the coding used by the data logging system to identify the automatic input data:

| 300 | x, lower |
|-----|-------------------------------------|
| 400 | y, lower |
| 500 | Θ , lower |
| 610 | P, total pressure, lower |
| 620 | p, positive static pressure, lower |
| 060 | -p, negative static pressure, lower |

A one in the third column indicates values for the upper traverse. The paper tape is one inch wide and uses an eight level frame to indicate an integer. The levels are numbered from the right. The first four levels use binary to indicate a decimal integer. The values of the levels are 1, 2, 4, and 8. Funches in the first and third levels would be the integer five (1+4). The fifth and six levels are always punched. A punch in the seventh level indicates the first digit of a five digit manual input. The eighth level is used to indicate the first digit

of a six digit automatic input.

Discussion of Program

In discussing the program, the term "sequence number" will refer to the identification number in the right hand six columns of the program. The term "statement number" will refer to the assigned number in the left five columns.

Subroutine Paper starts at sequence number 002600 and continues through the end cards. This routine first determines whether the first frame it reads is that of a manual or automatic input. It sets INDIC to one if manual input or to zero if automatic input. If the seventh level is punched (manual input) the subroutine reads that frame and the next four frames. After setting INDIC to one, a return is made to statement number 1 of the main program. At the next statement INDIC is tested to determine if the data is manual input. If manual it goes to statement number 7 where the manual operations start. The code (first integer) is checked at the "Assigned Go To" statement where the value formed by the four other integers is assigned to the proper constant.

If the input is automatic the first three integers must be checked to determine which array the value is to be assigned. This value is formed with the last three integers.

After the punched tape has been completely read the main program continues at statement number 19. At this point the program must add five inches to the values of x where the counters of the data logging system recycled to zero. This is done within the two outer "DO" loops between statements 19 and 950.

Between statements 950 and 101 the values of x are adjusted so that

the first x value is set equal to zero. This is done as an aid in the plotting of the input data.

While the punch tape is being read, the values read are printed as shown in Fig. B-1. This provides a format which is easy to check with the printed paper tape of the data logging system. At sequence number 001140 the input data is again printed but this time in the proper units as shown in Tables B-I and B-II. These tables contain the actual input data used to obtain the results listed in Table II.

One restriction on the program is that the value for M must always be equal to, or greater than, the value for L. This means that more data points must be taken on the upper traverse than on the lower. This is the normal situation due to the wakes at the upper station and therefore is not a problem.

After converting the pressures into pounds per square foot and angles into radians, the program performs the required integrations starting after statement 100.

The integrals of the lower traverse are defined as follows:

SUML = $\int_{0}^{1} P_{i} d\xi$ SUML = $\int_{0}^{1} \Delta P_{i} d\xi$ SUML = $\int_{0}^{1} P_{i} d\xi$ SUML = $\int_{0}^{1} \Delta P_{i} \cos^{2} \Theta_{i} d\xi$ SUML = $\int_{0}^{1} \sqrt{\Delta P_{i} P_{i} \gamma - \Delta P_{i}^{2}} \cos \Theta_{i} d\xi$

The integrals for the upper traverse are obtained by substituting U for L on the left side of the above equations and 2 for 1 on the right side.

The remainder of the program computes the lift, drag, and loss coefficients as indicated in Section 4 of this thesis. Starting at sequence number 002290 a table of results is printed as shown in Table B-III. This table is the same as Table II.

Definitions

The following symbols are followed by L to indicate lower traverse or U to indicate upper traverse:

| FORTRAN SYMBOLS | DEFINITION |
|-----------------|--|
| DELP | total pressure minus static pressure (P-p) |
| PS | static pressure (p) |
| PSN | negative values of static pressure (-p) |
| PSP | positive values of static pressure (+p) |
| PT | total pressure (P) |
| THETA | angle of flow (Θ) |
| x | distance in tangential direction |
| Y | distance along blade height |

The following symbols are followed by A to indicate station O and B to indicate station 3:

FORTRAN SYMBOLS

DEFINITION

| PS | static pressure (p) |
|-------|--------------------------------------|
| PT | total pressure (P) |
| RHO | fluid density (P) |
| THETA | angle of flow (⊖) |
| V | velocity of fluid |
| VCT | product of velocity and cosine theta |
| VST | product of velocity and sine theta |

The following symbols are used without a suffix:

| FORTRAN SYMBOLS | DEFINITION |
|-----------------|--|
| A | conversion unit, 5.2022 , $(1b/ft^2)/(in H_20)$ |
| AP2 | speed of sound squared, σ_{gRT_p} ,(ft ² /sec ²) |
| В | conversion unit, 2gJc _p ,(ft ² / ^o R sec ²) |
| BETAIN | angle from i to \overline{V}_{∞} , (deg) |
| BN | integer number of blades, (dimensionless) |
| BP | barometric pressure, (in Hg) |
| C | test blade chord, (ft) |
| CD | coefficient of drag based on V_∞ and RHO, (dimensionless) |
| CL | coefficient of lift based on V_{∞} and RHO, (dimensionless) |
| CD1 | coefficient of drag based on V_0 and RHOA, (dimensionless) |
| CLI | coefficient of lift based on V_0 and RHOA, (dimensionless) |

| FORTRAN SYMBOLS | DEFINITION |
|-----------------|---|
| DR | total force per test blade, (1b/ft) |
| DRA | axial component of total force per test blade, (1b/ft) |
| DRAG | drag per test blade, (1b/ft) |
| DRU | tangential component of total force per test blade, (lb/ft) |
| G | gravitational constant, (ft/sec ²) |
| GAMMA | ratio of specific heats, (dimensionless) |
| K | run number, (dimensionless) |
| L | number of lower data points, (dimensionless) |
| М | number of upper data points, (dimensionless) |
| R | product of gas constant and gravitational constant, (ft 2 / OR sec 2) |
| RHO | average of fluid density at stations 0 and 3, (1b \sec^2/ft^4) |
| S | blade spacing, (in) |
| TA | ambient temperature, (^O R) |
| \mathbf{TF} | ambient temperature, (°F) |
| TLIFT | lift per test blade, (lb/ft) |
| THETAIN | angle from \overline{j} to $\overline{V_{\infty}}$, (deg) |
| VA2 | velocity at station 0 squared, (ft^2/sec^2) |
| VB2 | velocity at station 3 squared, (ft^2/sec^2) |
| VIN | velocity at station ∞, (ft/sec) |
| XLO | initial lower traverse value of x, (in) |
| XUO | initial upper traverse value of x, (in) |
| ZETA | loss coefficient, (dimensionless) |

, DELPL 1), NC 3) *106+IN(4) * 10+IN(5) 4,15,16,17,18,19),11 DZ ,26),I 2 NI +0 3 M S 34,34 54,54 VC BOX 2)*] N, INDIC) DN N -950.133, -950.)53. * PSNL (L Z Z O INCI +0 2 ω кВР 2 Ô 66 1 II. Ĥ PRINI PRINI zo DRMA O 0 II H ŧ ۵. I=ZO -JOB+ROSE 000 Au H ž 11 0 0 pt ц. 2 \mathbf{c} -**02-86** <u>_</u> S ろういろ hÔ

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8-2

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Stretau(I), PIU(I), PSPU(I), PSNU(I), I=1, M)
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J([+1])/2. * **2+ cc. SF(THETAU(I))**2+ cc. XU(I)// (BN * S) * GAM4 - DELPU(I)**2) * 3 AMA - DELPU(I)**2) * 1* 3 AMA - DELPU(I)**2) * 0PP] * COSF(THETAU(I+1) * XL(I)) XL(I)) (I)**2) LPL(I+1)**2) * COSF(THE TAL(I+1))) XU(I)/(BN*S) 1) - XU(I)) / (BN*S) * COSF 1 I))+ DELPL (I+1) BN * S) xL(I+1)- XL(I *2+ DELPL(I+1) S) + 1) - XL((I + I) NX) RHCB) (L+1) NX + - DEL' . ب ×* **4** Oc) (2 * (x I+1))/2. THET DELPU(I+1))/2. * (DELPU(I+1))/2 RHOHA Ш О AMMA - DI + SAMMA -Z ω -N 1))/2 × 40 * 14 ELPL(\$ PTU(I+1)) + DELPU(ŧ 20 SUML12)/AP2 SUML12)/AP2 SUML4 SUML4 42/RH0A 142/RH0A 142/RH0A 142/RH0B + * SU42) + T A Z + PSUC + -0 ×U(I+ ΗÏ (DELPU(I) * × PTL(I) + (DELPL(I) S E I PSL DELP XL ÷ SUML1 SUMU1 5 ā+ d O - 4 + 2 +2 + A 5000 14P2 SUMUT F SUMU2 SUMU3 4 GAMM/ SURI 0 STNF(2 CONTINUE SUMUTINUE SUMU12 SUMU12 SUMU2 SU 0 DI S UMI SUMI -0 [+] ral () Sort vu 300 I = 1. L NC 00 (XU 11 0 " S .0=0 11 X HETAU(n n 11 II $\overline{\boldsymbol{\omega}}$ SUMU4 UMU12 11 UMU2 UMU3 THET UML 0PP 1 ž SUMI M 90 300

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0021 00 COEFFIC IENTS/29X4HL IFT3X4HDRAG8X4HLGS MODE 2 2 TAPE IN BUFF. WORD 5X 4HTEMP 5X 2HBN 5X 1HS /) F7.0, F7.1) LéX6HTHETAL5X3HPTL7X4HPSPL6X44PSNL HXUEX2HYU&X6HTHETAU5X3HPTU7X4HPSPU&X44PSNU F10.5, 5X 5HCD1 = F10.5) TIFIED MANUAL INPJT CODE USED IN PUNCH TAP P UNCH CHAR SAV = 1, 1N = 7) 777775778,MASK2=7777777777776778,MASK3=7 5,6,7,8 PUNCHES ~ NOT 8 PUNCH GO TEST 7 PUNCH READER CHARAC TERS T READER CHAR. MODE CHAR • ALL BUT 8 8 PUNCH. F NOT 7 TAPE CC AN READ IN 5 THE PT P AP ER NZ A SLEAR / SELECT EXIT D (1A41) SET LA READ I CLEAR IF NO CLEAR IND IC) SAL **CRMAT(29X28HNDN-CIMENSICNAL** 5X 2HEP F8.1.F A31 IN1, N). 118). ZETA . 18). 28) MASK3) 146). HAR), SCL(MASK2) A3) HAR+4), STA(N). 1.2, F , SCL (MASK1) ENI1(4). 4 2 X K K -1 EXF7 ICEI CD 4 LJ(1A5) PAPER N1(6) ~ J P C F C JP(ົວແ R+5) · SC z XRSAV AR AR HAR HAR) ď CHAR GH AH + + ORMAT ORMAT ORMAT **DRMA** UBRCI ORMA CRMA W LOA LOA AZ 4 ENA UN UN 22 A L **V**O Sol A O **V**O 110 1 A 4 405 A21 1A3 1 A 6 2ERR 101 1A2 **441** σ I A 1 1A31

98

B-6

Initial Print Out

of Input Data



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| Table | B-I |
|-------|-----|
|-------|-----|

Input Data

Lower Traverse

| RUN NO. | BP | TEMP | a | BN | S | | ÷. ÷ | | |
|---|--|--------------------------------------|---|---------------------------------------|----|--|------|--|------|
| 100 | 29.98 | 66.5 | | 2. 6 | .0 | | | | |
| POINT | XL | YL | | THETAL | | PTL | | PSPL | PSNL |
| 12345678900123456789001234567890012345678900123456789000000000000000000000000000000000000 | 000076646162840082698843026175848280827080510672636 111122223333333344445555566667777788888999990000111111111111111111111 | 000000000000000000000000000000000000 | | атттттттттттттттттттттттттттттттттттт | | 20000000000000000000000000000000000000 | | 22222222222222222222222222222222222222 | |

| | | | | | | ~ |
|--|---|--|--|---|------|--|
| POINT | XU | YU | THETAU | PTU | PSPU | PSNU |
| 1234567890123456789012345678901234567890123456789012345678901234555555555555555555555555555555555555 | 0258624656456444232480642767454120099877099007088077 1111122222222222345678901234456780012344667990 111122222222222234567890123445678001344667990 667990 | 55555555555555555555555555555555555555 | $\begin{array}{c} 0 \\ 0 \\ 0 \\ 0 \\ 0 \\ 0 \\ 0 \\ 0 \\ 0 \\ 0 $ | 24 + 30 24 + 30 24 + 40 24 + 40 24 + 40 24 + 40 24 + 40 20 20 20 20 20 20 20 20 20 20 20 20 20 | | -21-2222222-2222-2222-2222-2222-2222-2 |

Table B-II

Input Data

Upper Traverse

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4C9245626258461839688138C2055291221066668222246486588

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Table B-II (Continued)

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Table B-III

CASCADE LABORATORY U.S. NAVAL POSTGRADUATE SCHOOL

TEST RUN NO. 100 BARCMETRIC PRES. = 29.98 IN. HG. AMBIENT TEMP. = 66.5 DEG. F.

| (FT/SEC) | THETA O (DEG) | (FT/SEC) | THETA 3 (DEG) | V INF (FT/SEC) | THETA INF (DEG) | BETA (DEG) |
|--------------------------|------------------|-----------------------------|----------------------------|-------------------|--------------------|---------------|
| 312.49 | 43.63 | 322.32 | -46.18 | 224.86 | -2.16 | 92.17 |
| | • | | | | | 1.0 |
| AXIAL FORCE (LB/FT | TANGEN FOR | TIAL TO CE FO FT) (LB | TAL RCE THE /FT) (DE | TA LI G) (LB | FT DRA | AG. /FT) |

| 15.62 119.56 | 120.58 | 7.44 | 120.07 | 11.10 |
|--------------|--------|------|--------|-------|

NCN-CIMENSIONAL COEFFICIENTS LIFT DRAG LOSS

~

| 2 | 0002 | 2400 | 1221 |
|----|-----------|--------|-------|
| ۷. | • 6 2 2 3 | • 2009 | +1334 |

