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 $^{\omega}4\,0\,2\,0\,6\,1$ ւ հեյзեն s/142/62/005/006/006/01 (5) E192/E382 Rapoport, G.N. and Zhurakhovskiy, V.A. AUTHORS: TITLE: (G'Theory of phasochronous devices of type "O" with helical electron beams PERIODICAL: (5)Izvestiya vysshikh uchebnykh zavedeniy, Radiotekhnika, v. 5, no. 6, 1962, 707 - 713 TEXT: The helical electron beam is controlled by a constant magnetic field $B_z = B_o$ and moves in a high-frequency field described by: $\vec{E}_{\pm s} = \vec{E}_{\pm s} \quad Ae \\ \vec{\pm}s = \vec{E}_{\pm s} \quad Ae \\ \vec{\pm}s = \vec{H}_{\pm s} \quad Ae \\ \vec{H}_{\pm s} = \vec{H}_{\pm s} \quad Ae$ 5 (1). The waves of the field propagate in a uniform cylindrical waveguide in the direction +z . [" is the "hot" propagation constant and γ is a propagation constant in the absence of an electron beam. The axis of the helix z is parallel to the axis of the waveguide Card 1/4 s/142/62/005/006/006/011 E192/E382 Theory of phasochronous devices The motion of the electrons in the presence of the field is des- $\left\{ \right\}$ cribed by: $\mathbf{\dot{v}}_{\mathbf{x}} = -\eta \mathbf{E}_{\mathbf{x}} + \eta \mathbf{v}_{\mathbf{z}} \mathbf{B}_{\mathbf{y}} - \eta \mathbf{v}_{\mathbf{y}} (\mathbf{B}_{\mathbf{z}} + \mathbf{B}_{\mathbf{o}})$ (4a) $\mathbf{\ddot{v}}_{\mathbf{v}} = - \eta \mathbf{E}_{\mathbf{v}} + \eta \mathbf{v}_{\mathbf{x}} (\mathbf{B}_{\mathbf{z}} + \mathbf{B}_{\mathbf{o}}) - \eta \mathbf{v}_{\mathbf{z}} \mathbf{B}_{\mathbf{x}}$ $\mathbf{\dot{v}}_{\mathbf{z}} = - \eta \mathbf{E}_{\mathbf{z}} + \eta \mathbf{v}_{\mathbf{v}} \mathbf{B}_{\mathbf{x}} - \eta \mathbf{v}_{\mathbf{x}} \mathbf{B}_{\mathbf{v}}$ (46) . The solutions of Eqs. (4a) are in the form: $x = \widetilde{X}(z, \alpha) + x_1(z, \alpha)$ $y = \widetilde{Y}(z, \alpha) + y_1(z, \alpha)$ where: $\widetilde{X}(z, \alpha) = X(t(t_{\alpha}, z), t_{\alpha})$ magnitudes for the alternating components Y produced X and The by the afferent transit times of the electrons are much higher Card 2/4

Theory of phasochronous devices

than the components x_1 and y_1 ; the latter can therefore be neglected. An equation for the starting current of the system is derived and it is shown that this is similar to the small signal scattering equation of a travelling-wave tube or a backward-wave tube. If the spread of the electron velocities $\Delta v_{\pi} = 0$ is taken

into account, the scattering equation becomes:

$$(x + ib)(x^2 + 4QC) - i = 0$$

where

$$\left(\sqrt{\frac{2}{2}}\right)^2 = 4$$
 QC

The following notation is adopted in Eq. (21):

$$\begin{aligned} \varepsilon/C &\equiv b \\ \delta/C &\equiv \mu \\ \mu - \mathbf{i}b &\equiv \mathbf{x} \end{aligned}$$
 (17)

E192/E382

where ε is the detuning parameter. Eq. (22) is similar to that Card $\frac{3}{4}$

S/142/62/005/006/006/011 E192/E382

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(22)

(21)

Theory of phasochronous devices E192/H

of a travelling-wave tube with space charge. The effect of the variations in the magnetic field is also taken into account and it is found that 1% variation can lead to the doubling of the starting b current.

ASSOCIATION:	Institut radiotekhnicheskikh problem AN USSR (Institute of Radio-engineering Problems, AS UkrSSR)	
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