Transport of Nordic Seas overflow water into and within the Irminger Sea: An eddy-resolving simulation and observations

X. Xu,¹ W. J. Schmitz Jr.,² H. E. Hurlburt,³ P. J. Hogan,³ and E. P. Chassignet⁴

Received 16 April 2010; revised 16 September 2010; accepted 18 October 2010; published 18 December 2010.

Results from a climatologically forced, eddy-resolving (1/12°) Atlantic simulation using the Hybrid Coordinate Ocean Model help clarify some presently unresolved connections between volume transports of Nordic Seas overflow water at key locations in the northernmost North Atlantic Ocean. The model results demonstrate that, in addition to the known westward flow through the Charlie Gibbs Fracture Zone (CGFZ), some Iceland Scotland overflow water (ISOW) flows westward through gaps in the Reykjanes Ridge north of the CGFZ into the Irminger Sea, and some flows southward along the eastern flank of the Mid-Atlantic Ridge into the West European Basin. These results provide insights into the well-known inconsistency between observed westward transport of ISOW through the CGFZ (2.4 Sv) and the transports upstream at Southeast of Iceland section (3.2 Sv) and downstream in the western Irminger Sea (4.5 Sv). Although the portion of the simulated ISOW that flows through CGFZ is about 500 m deeper than observed, the model results also show two ISOW pathways of this flow into the Irminger Sea, one northward along the western flank of the Reykjanes Ridge and the other westward before turning north–eastward on the western side of the Irminger Basin. Comparisons with the long-term moored instrument database in the Irminger Sea show that the model-based mean circulation is in reasonable agreement with observed volume transports of overflow water and that it gives approximately correct temperature and salinity characteristics.


1. Introduction

It has been known for some time that flow over the Greenland-Iceland-Scotland (GIS) Ridge provides regional sources of Nordic Seas overflow water (NSOW) into the northern North Atlantic Ocean [see, e.g., Worthington and Wright, 1970, Plate 29]. The “western source” is through the Denmark Strait [see Ross, 1984; Girton et al., 2001, for example]. After flowing over the sill, and entraining ambient water, this bottom-trapped Denmark Strait Overflow Water (DSOW) has been observed to continue down the slope in the western Irminger Sea [e.g., Dickson and Brown, 1994; Dickson et al., 2008; Bacon and Saunders, 2010]. The “eastern sources” are more complex [Østerhus et al., 2008, Figure 18.1], involving bottom-trapped flows over the Iceland-Faroe Ridge (IFR) and out through the Faroe Bank Channel (FBC). A minor contribution also spills over the Wyville-Thomson Ridge. After flowing over the IFR and the FBC sills, and entraining ambient water, the resulting Iceland-Scotland overflow water (ISOW) has been observed to flow southwestward along the northwestern slope of the Iceland Basin [Saunders, 1996], and westward through the Charlie-Gibbs fracture zone (CGFZ) [Saunders, 1994]. It has been suggested that this flow through the CGFZ subsequently turns northward into the Irminger Sea, joining DSOW there [e.g., Dickson and Brown, 1994; Saunders, 2001].

DSOW and ISOW are two important components of North Atlantic Deep Water (NADW), which in addition includes Labrador Sea Water (LSW) and modified Antarctic Bottom Water (AABW). Knowledge of the circulation pathways and volume transports of these overflow water masses is therefore fundamental for a general description of the associated meridional overturning circulation in the Atlantic Ocean (often referred to as the AMOC). For related observational studies considering the source of inflow water from the North Atlantic Ocean, water transformations in the Nordic Seas, and the flow exchange over the GIS Ridge, the reader is referred to reviews by McCartney and Mauritzen [2001], Mauritzen [1996a, 1996b], Hansen and Østerhus [2000],
### Abstract
Transport of Nordic Seas overflow water into and within the Irminger Sea: An eddy-resolving simulation and observations

### Security Classification
- a. REPORT: unclassified
- b. ABSTRACT: unclassified
- c. THIS PAGE: unclassified
- 17. LIMITATION OF ABSTRACT: Same as Report (SAR)
- 18. NUMBER OF PAGES: 20
- 19a. NAME OF RESPONSIBLE PERSON: unclassified
and Hansen et al. [2008]. Our focus here is on the overflow water transports downstream of the GIS Ridge, but before entering the subpolar gyre.

[4] Overflow water in this area has been relatively well documented compared to further downstream. There are questions that require further clarification, however. The direct measurements of ISOW transport by Saunders [1994, 1996] lead to unresolved, although not unanticipated, questions concerning a possible branch of overflow water flowing through gaps in the Reykjanes Ridge north of the CGFZ. Questions concerning the transport of ISOW from the Iceland Basin into and within the Irminger Sea have been put forth in reviews by Saunders [2001] and Hansen and Østerhus [2000]. Specifically, (1) What is the circulation of ISOW in the Iceland Basin? (2) Does all ISOW flow westward through the CGFZ into western basins? (3) What is the flow pathway for ISOW from the CGFZ into the Irminger Sea and where (in the Irminger Sea) does it merge with DSOW? Furthermore, the recent overflow transport estimates of Bacon and Saunders [2010], along with Dickson et al. [2008], add a new perspective to the interpretation of the basic overflow transports at key sections in the western Irminger Sea. The key interest here is to use numerical model results as an aid to observations in addressing these questions.

[5] While the spreading of overflow plumes close to the Denmark Strait or the FBC has been adequately simulated in some regional models [e.g., Shi et al., 2001; Käse et al., 2003; Riemenschneider and Legg, 2007], their representation in global or basin-scale models has always been a challenge. For a climate model with a horizontal resolution of ~1°, an alternative (probably more effective) approach to explicitly simulating the overflow is to parameterize it as a marginal sea boundary condition [Legg et al., 2009]. For an eddy-resolving model (~0.1° or finer), the main difficulty is to parameterize the entrainment process and that depends strongly on the choice of vertical coordinate [Griffies et al., 2000]. A level model, by virtue of its stepped seabed, gives rise to numerically induced, unrealistic downslope spillage and consequent vertical convection, which mixes away the cold overflow water at a much too shallow depth. In contrast, an isopycnic model, in the absence of entrainment parameterization, contains insufficient diapycnal mixing and consequently the overflow properties are retained over unrealistically long distances [Willebrand et al., 2001]. Transport estimates based on the long-term moored instrument arrays of Dickson and Brown [1994], Saunders [1994, 1996], Dickson et al. [2008], and Bacon and Saunders [2010] have formed a valuable benchmark for evaluating model performance in simulating overflow transport, which has not been well documented in the past. Some, but not all, of these transport data were used in a recent numerical study by Chang et al. [2009]. Their prime interest is to investigate the impact of horizontal resolution (from 1° to 1/12°) and associated topography representation on simulating NSOW pathways. A detailed evaluation of the model-based overflow transports and water properties using the updated data is also a key interest of this investigation.

[6] The presentation is organized as follows. After this introduction, a brief summary of the key NSOW transport measurements and a description of model configurations are presented in section 2. The main results related to mean overflow transport, as determined by a basin-scale, eddy-resolving simulation are presented in section 3, and time evolution/variation of the model solution is discussed in section 4. In section 5 the questions above are discussed within the combined framework of observations and model results.

2. Observations and the Ocean Model

2.1. A Brief Summary of Key NSOW Transport Measurements

[7] An essential component of the observational basis for overflow water in our area of interest is the volume transport estimated from five substantially instrumented, and “comparatively long” moored current meter arrays. The locations of these mooring arrays (numbered 1 through 5) are shown as sections of orange circles in Figure 1. These transport measurements are summarized in Table 1 and the reader is referred to Appendix A for a detailed account. Southeast of Cape Farewell there are also many studies considering overflow transports primarily based on hydrographic data [e.g., Bacon, 1997, 1998; Kieke and Rhein, 2006; Lherminier et al., 2007; Sarafanov et al., 2009; Holliday et al., 2009]. Some of these studies suggest that the (baroclinic) transports of overflow water contain a significant decadal variation, which seems to be in conflict with the results based on long term moored instruments [see Dickson et al., 2008, their section 19.5]. Since our interest is mainly in the mean state, we will focus the discussion on directly measured transports based on mooring arrays.

[8] The Reykjanes Ridge, which separates the Iceland Basin and the Irminger Sea, is a key topographic feature in the area of prime interest. A zoomed depiction of the bottom topography surrounding the Reykjanes Ridge is shown in Figure 2. It also depicts section 6 (denoted as a white line), across which model-determined transports through gaps in the Reykjanes Ridge are calculated (see section 3.1).

2.2. Model Configurations

[9] The HYbrid Coordinate Ocean model [HYCOM, Bleck, 2002; Chassignet et al., 2003] is the ocean general circulation model (OGCM) used in this investigation. The vertical coordinate in HYCOM is isopycinal in the stratified open ocean and makes a dynamically smooth and time-dependent transition to terrain-following in shallow coastal regions and to fixed pressure levels in the surface mixed layer and/or unstratified seas. In doing so, the model ideally combines the advantages of the different coordinate types in optimally simulating coastal and open ocean circulation features.

[10] The primary model results for this investigation come from a basin-scale Atlantic simulation with horizontal resolution of 1/12° at the Equator (and about 5 km resolution in our area of interest). The computational domain extends from 28°S to near 80°N at the Fram Strait meridionally and from 98°W to 36°E zonally. The horizontal grid is a Mercator projection south of 47°N, and north of this latitude it uses a bipolar Arctic cap where the two poles are shifted over land to avoid a singularity at the North Pole (in a global simulation). Vertically, the simulation contains 32 hybrid layers with density referenced to 2000 m (σt): 28.10, 28.90, 29.70, 30.50, 30.95, 31.50, 32.05, 32.60, 33.15, 33.70, 34.25,
34.75, 35.15, 35.50, 35.80, 36.04, 36.20, 36.38, 36.52, 36.62, 36.70, 36.77, 36.83, 36.89, 37.02, 37.06, 37.10, 37.17, 37.30, 37.42, and 37.48 kg m$^{-3}$. Layers 25–29 are connected with the overflow water. We have run sensitivity experiments with more layers in the overflow water, with similar results.

The model topography is based on the 2′ Naval Research Laboratory (NRL) digital bathymetry database, which combines the global topography based on satellite

Figure 1: Bathymetry of the northern North Atlantic Ocean based on version 12.1 of Smith and Sandwell [1997]. The contour intervals are 500 m, with 2000 and 4000 m emphasized by bold black lines. Orange circles identify the locations of five long-term current meter mooring arrays that have been used to estimate the volume transports of Nordic Seas overflow water. 1, TTO [Dickson and Brown, 1994]; 2, Angmagssalik [Dickson et al., 2008]; 3, Southeast of Cape Farewell [Bacon and Saunders, 2010]; 4, Southeast of Iceland [Saunders, 1996]; 5, Charlie-Gibbs fracture zone [Saunders, 1994]. Blue dots show the locations of the hydrographic TTO section by Livingston et al. [1985]. The area outlined by the dashed black rectangle is shown in Figure 2.

Table 1. Transports (in Sv) of Nordic Seas Overflow Water Based on Moored Current Meter Arrays at Five Locations Downstream of the GIS Ridge

<table>
<thead>
<tr>
<th>Mooring Array Number</th>
<th>Name</th>
<th>Transports</th>
<th>Deployment Period(s)</th>
<th>M/I</th>
<th>Key Reference</th>
</tr>
</thead>
<tbody>
<tr>
<td>1</td>
<td>TTO</td>
<td>5.1$^b$</td>
<td>12 Jul 1990 to 30 Jul 1991</td>
<td>5/12</td>
<td>Dickson and Brown [1994]</td>
</tr>
<tr>
<td>2</td>
<td>ANG</td>
<td>7.3$^c$, 4.0$^f$</td>
<td>1986–1990, 1997–2005</td>
<td>8/24</td>
<td>Dickson et al. [2008]</td>
</tr>
<tr>
<td>3</td>
<td>SECF</td>
<td>9.0$^o$, 4.5$^g$</td>
<td>6 Sep 2005 to 30 Aug 2006</td>
<td>9/32</td>
<td>Bacon and Saunders [2010]</td>
</tr>
<tr>
<td>5</td>
<td>CGFZ</td>
<td>2.4$^b$</td>
<td>1 Jun 1988 to 30 Sep 1989</td>
<td>8/16</td>
<td>Saunders [1994]</td>
</tr>
</tbody>
</table>

$^a$TTO, Transient Tracer in the Ocean; ANG, Angmagssalik; SECF, Southeast of Cape Farewell; SEI, Southeast of Iceland; CGFZ, Charlie-Gibbs fracture zone. M/I are numbers of mooring locations and total instruments. The ANG Array includes multiple deployments with variable M/I numbers. Listed numbers are based on the most extensive deployment covering 3 July 1988 to 29 June 1989.

$^b$The transports are evaluated for water below $\sigma_\theta$ of 27.80 kg m$^{-3}$.

$^c$The transports are evaluated for water below $\sigma_\theta$ of 27.85 kg m$^{-3}$. 
altimetry of Smith and Sandwell [1997] with several high-resolution regional databases (see http://www7320.nrlssc.navy.mil/DBDB2_WWW for documentation). A 9-point smoother was applied once to reduce inaccurate model response at poorly resolved scales. The simulation was integrated for 20 years after initialization from rest and January temperature (T) and salinity (S) from the Generalized Digital Environmental Model (GDEM) [Carnes, 2009] ocean climatology. Within a buffer zone, extending about 3° from the northern and southern boundaries, the model T and S are restored to the same (monthly) climatology with an e-folding time scale of 5–60 days. The horizontal diffusion and viscosity parameters are listed in Table 2. For vertical/diapycnal mixing the model uses the K-profile parameterization of Large et al. [1994] in the surface mixed layer and in the ocean interior.

[12] The surface forcing is from a monthly climatology based on the European Center for Medium-Range Weather Forecasts reanalysis (ERA40) [Uppala et al., 2005]. To better simulate the surface mixed layer, submonthly wind anomalies from the Fleet Numerical Meteorology and Oceanography Center 3 hourly 0.5° Navy Operational Global Atmospheric Prediction System (NOGAPS) for the year 2003 are added to interpolated monthly means. Surface heat flux is given by

$$Q_{net} = Q_{shortwave} + Q_{longwave} + Q_{latent} + Q_{sensible},$$  \hspace{1cm} (1)$$

in which the surface short- and long-wave radiations are from ERA40, and the latent and sensible heat fluxes are calculated using the modeled sea surface temperature (SST) and the bulk formulae of Kara et al. [2005]. The latter provide negative feedback that increases/decreases the net

![Figure 2](image.png)

Figure 2. Enlargement of the area within the dashed black rectangle drawn on Figure 1. The white solid line indicates a model Reykjanes Ridge section 6, along which model-based volume transport was calculated. Red triangles show the location of three current meter moorings discussed by Shor et al. [1980].

| Table 2. Horizontal Diffusion Parameters Used in the Basin-Scale Atlantic Simulation |
|---------------------------------|------------------|
| Diffusion Parameter            | Value            |
| Laplacian deformation-dependent viscosity coefficient | 0.05             |
| Laplacian viscosity for momentum | 20 m² s⁻¹        |
| Biharmonic diffusion velocity for momentum | 1.0 cm s⁻¹      |
| Biharmonic diffusion velocity for layer thickness | 1.0 cm s⁻¹       |
| Laplacian diffusion velocity for temperature, salinity | 0.5 cm s⁻¹     |
heat flux if model SST is too cold/warm. Surface freshwater flux is treated as a virtual salinity flux,

$$F_{\text{saln}} = \frac{E}{C_0} - \frac{P}{C_0} \times \frac{R}{C_0} \times (S_{\text{SSS}} + \frac{V}{C_2} S_{\text{clim}} / C_0 S_{\text{SSS}}),$$

in which $E$, $P$, and $R$ represent evaporation, precipitation, and river runoffs and the sea surface salinity ($SSS$) is restored to monthly climatology with a restoring strength ($V$) of 15 m/30 days. The salinity difference (between model and climatology) in $SSS$ restoring is clipped to be 0.5 psu. The idea of reduced salinity restoring is to diminish its damping effect on ocean fronts. The reader is referred to section 3 and Appendix B of Griffies et al. [2009] for an insightful discussion of surface thermohaline fluxes.

3. Model-Based Mean Transports of NSOW

We now consider model results from the eddy-resolving Atlantic simulation described in section 2.2. Specifically, how does our numerical experiment help to clarify presently unresolved connections between overflow transports, and how do model results compare to available observations? The simulation was integrated for 20 years after initialization (from climatology) and the time average of the final five years is used as the “mean” state here. We first discuss the flow over the Reykjanes Ridge north of the CGFZ in section 3.1, then follow the general spreading pathway of ISOW from the Iceland Basin, through the CGFZ, and into the Irminger Sea (3.2–3.4). The circulation in the area south of Cape Farewell is presented in section 3.5.

### 3.1. ISOW Over the Reykjanes Ridge

The observed transport through the CGFZ (2.4 Sv) based on direct measurements by Saunders [1994] has presented a long-standing unresolved conundrum, because it is too small to account for the ISOW found upstream and downstream [e.g., Hansen and Østerhus 2000]. The topic of a possible westward flow of ISOW over the Reykjanes Ridge north of the CGFZ has been mentioned by a few authors (see section A2). Here we use results based on an eddy-resolving numerical experiment, with about 5 km resolution around 57°N, and model-data comparisons to demonstrate the likelihood that some ISOW flows into the Irminger Sea through gaps in the Reykjanes Ridge.

The model-based mean circulation of ISOW surrounding the Reykjanes Ridge and the CGFZ area is illustrated in Figure 3. As has long been realized, the CGFZ is the principle conduit carrying ISOW from the Iceland Basin to the western side of the Mid-Atlantic Ridge (MAR). It is not the sole conduit, however. As shown in the zoomed Figure 3 (right), some model-based ISOW flows into the Irminger Sea over the shallow Reykjanes Ridge near 59°N, some through the relatively deep Bight Fracture Zone near 57°N, and some through a couple of unnamed gaps/valleys between 55°N and 56°N. Along section 6 over the Reykjanes Ridge (denoted as the white line in Figure 3), the model-based mean normal velocity and the accumulated volume transports are displayed in Figure 4. The mean velocity of ISOW over the Ridge is about 1–2 cm s⁻¹ westward. The model-
determined mean transport for water below 27.80 kg m

−3 is 1.2 Sv, half of the value observed through the CGFZ.

To further demonstrate that a westward flow of
ISOW over the Reykjanes Ridge is possible, we have
plotted the salinity distribution based on historical bottle
data collected in the World Ocean Database (Figure 5). The
high salinity water sitting over the Reykjanes Ridge south of
about 58°N, with density greater than 27.80 kg m
−3, is a clear signature of ISOW. The highest salinity of ISOW is
found in the Bight Fracture Zone near 57°N. For compari-
on, the model-based salinity distribution is also shown in
Figure 5. There are differences. For example, the observa-
tions show a subsurface low salinity core associated with
LSW in the CGFZ, while it is found over the Reykjanes
Ridge in model. Nevertheless, the overall structure and the
saline ISOW over the Reykjanes Ridge roughly agrees with
observations.

3.2. ISOW in the Iceland Basin

Along a section colocated with the mooring array
deployed southeast of Iceland (SEI hereafter), the model-
based vertical distribution of mean salinity, mean normal
velocity, and the accumulated volume transport are illus-
trated in Figure 6a. Here the simulated ISOW is seen as a
bottom trapped layer of saline water on the Icelandic slope,
occupying a depth range from 1300 m to the maximum
depth of this section (about 2300 m). The maximum ISOW
salinity is about 35.04 psu. Relatively fresh LSW with a
salinity minimum of about 34.94 psu is centered at a depth
of 1600 m in the middle of the Iceland Basin. The warm and
saline MNAW (7°C–8°C and 35.1–35.3 psu) occupies the
upper 600–1000 m above the halocline (also the thermo-
cline). While the maximum salinity of modeled MNAW is
about 0.1 psu higher than observed, this three-layer salinity
structure agrees well with observations [see Saunders, 1996,
Figure 2b]. Strong currents are simulated within the ISOW
on the mid Icelandic slope (20–25 cm s
−1 near the 1500 m
contour) and decrease toward the deeper part of the section
(2–5 cm s
−1 at depths greater than 2000 m). This simulated
horizontal pattern and flow magnitude roughly agrees with
observations [see Saunders, 1996, Table 2 and Figure 4].
The model-based accumulated volume transport has a
maximum of 3.3 Sv for \( \sigma_0 \geq 27.80 \) kg m
−3. A transport of
3.2 Sv was observed by Saunders [1996].

Spreading of the modeled deep ISOW in the Iceland
Basin, as defined by volume transport per unit width (in m²
s
−1) below \( \sigma_0 \) of 27.85 kg m
−3, is displayed in Figure 7.
Southwest of Iceland, the modeled ISOW separates into
three main branches. The westernmost (shallowest) branch
flows parallel to the Reykjanes Ridge and contributes to the
flow across the Reykjanes Ridge north of the CGFZ (also
see Figure 3). The middle branch follows along a depth
contour of about 2500 m and contributes mainly to the flow
through the CGFZ. The deepest branch flows into the dee-
pest part of the Iceland Basin and does not contribute to the
westward flow. These branches of the ISOW plume closely
follow contours of the bottom topography in the Iceland
Basin. Superimposed on Figure 7 are the crests of three
depositional drifts reproduced from Figure 5.4 of Saunders
[2001]. These sedimentation patterns, called the Björn,
Björnsson, and Gardar Drifts, are thought to be formed by the
interaction of ISOW with sediment in the Iceland Basin (for
more detail see Bianchi and McCave [2000]).

3.3. Flow Through the Charlie-Gibbs Fracture Zone

Figure 6b depicts the model-based vertical distribu-
tion of the mean salinity and normal velocity along the

Figure 5. Vertical distribution of salinity (in psu) along the
Reykjanes Ridge and across the Charlie–Gibbs fracture zone
based on (top) observations and (bottom) the model. The ob-
servations are from the historical ocean station data col-
lected in the World Ocean Database 2005 (WOD05, http://
www.nodc.noaa.gov). The insert map illustrates the location
of selected stations.
CGFZ array near longitude 35°W. The northern channel of the CGFZ in the model has a sill depth of 3428 m, compared to 3675 m observed [Shor et al., 1980]. The simulated ISOW has a maximum salinity of about 35.0 psu close to the bottom. This is slightly more saline than observed [Saunders, 1994, Figure 2] and the depth of the salinity maximum is also about 500 m deeper. The modeled isopycnal 27.80 kg m$^{-3}$ gradually deepens from about 1500 m at 53°N to 2000 m at 51.5°N in agreement with observations. The model-determined westward flow (Figure 6b) is stronger in the northern channel (about 10 cm s$^{-1}$) than in the southern channel (2–5 cm s$^{-1}$). This flow distribution and magnitude compare quite well with observations [Saunders, 1994, Table 3 and Figure 6].

Within the area of mooring coverage, the accumulated volume transports of the simulated westward flow are 2.4 and 2.1 Sv, for water below $\sigma_\theta$ of 27.80 and 27.85 kg m$^{-3}$, respectively (Figure 6b, bottom).

Figure 8 illustrates the mean circulation of model-based deep ISOW ($\sigma_\theta \geq 27.85$ kg m$^{-3}$) in the vicinity of the CGFZ. The net southward transport of the modeled deep ISOW is 2.5 Sv across the southern Iceland Basin near 54°N (not shown). About 1.9 Sv of this deep ISOW flows through various gaps west of 32°W and turns westward in the northern channel of the CGFZ. The rest, which flows southward east of 32°W, turns eastward in the eastern part of northern channel. This eastward flow is found above the sill of the northern channel and therefore is unlikely a result of topographic blocking. The model-based flow in the southern channel of the CGFZ is also westward and contains a high-salinity signature similar to that in the northern channel (see Figure 6b). Figure 8, however, suggests that this westward flow is not fed from the Iceland Basin, but rather is a localized recirculation feature. Also, the small westward transport (0.3 Sv) between 27.80 and 27.85 kg m$^{-3}$ is fed from an eastward flow in the southern part of the CGFZ. Therefore, the model-determined “net” transport of ISOW from the Iceland Basin through CGFZ is 1.9 Sv westward with $\sigma_\theta$ larger than 27.85 kg m$^{-3}$. This is 0.5 Sv less than the 2.4 Sv based on data [Saunders, 1994].

[21] The model results also show eastward flows that extends below $\sigma_\theta$ of 27.80 kg m$^{-3}$ south of about 52.5°N and all the way to the bottom near 51.5°N (Figure 6b). Below 27.85 kg m$^{-3}$, the eastward flow eventually feeds into the northern channel of the CGFZ through several gaps around 32°W and merges with the eastward flow contributed from the Iceland Basin. East of about 28°W the combined eastward flow (about 1 Sv) turns southward and flows along the eastern flank the MAR. If this southward flow is realistic, it would account for the ISOW tracer signals observed in the West European Basin (WEB) by Fleischmann and Rhein [2000] and Fleischmann et al. [2001]; see also
Based on the observed tracer signals, Fleischmann et al. [2001] estimated a southward transport of 2.4–3.5 Sv (corresponding to $\sigma_\theta/C_{18} \geq 27.80$ kg m$^{-3}$) from the Iceland Basin into the WEB. The model results yield a net southward transport of 2.6 Sv into the WEB. About 2 Sv are from west of the CGFZ and most of it has $\sigma_\theta/C_{18} < 27.85$ kg m$^{-3}$. It seems possible that via recirculations and mixing, the eastward flow through the CGFZ is tagged with the tracer signals of ISOW. We also note that in addition some eastward flow with $\sigma_\theta$ between 27.80 and 27.85 kg m$^{-3}$ flows northward into the Iceland Basin and adds to the ISOW transport over the Reykjanes Ridge.

### 3.4. Circulation Within the Irminger Sea

For simplicity, we use the location of the Southeast of Cape Farewell (SECF hereafter) mooring array as a southern boundary of the Irminger Sea. The model-determined mean circulation in the Irminger Sea is shown for two layers in Figure 9. For $\sigma_\theta \geq 27.85$ kg m$^{-3}$ (Figure 9a), the circulation is dominated by the southward flowing DSOW along the western boundary. The flow of deep ISOW into and in the Irminger Sea roughly follows the deepest trough of the basin. For $27.80 \leq \sigma_\theta < 27.85$ kg m$^{-3}$ (Figure 9b), this upper overflow layer circulation is generally cyclonic around the basin and contains inputs from the Denmark Strait, from east of the Reykjanes Ridge, as well as from recirculation south of Cape Farewell. To shed some light on the location(s) where ISOW joins DSOW, we draw a straight white line in Figure 9 from the Denmark Strait southwestward along the Irminger Sea trough roughly in parallel with the Reykjanes Ridge. The accumulated volume transports along this line are shown in Figure 10. The model results suggest that the crossing (from east to west) is approximately evenly distributed for $\sigma_\theta \geq 27.85$ kg m$^{-3}$. In contrast, significantly more crossing takes place north of the ANG line for $27.80 \leq \sigma_\theta < 27.85$ kg m$^{-3}$. One reason is that the less dense ISOW through the gaps in the Reykjanes Ridge flows around the northern end of the Irminger Sea.

The distributions of mean salinity and velocity across the basin are shown in Figure 11 along model sections that are colocated with the ANG and SECF arrays. The model-
based salinity contains a structure similar to that observed in the Irminger Sea (see Figure A1 in section A3, for an example). The model flows are generally strong and south-westward along the western boundary, weak and northeastward through the eastern half of the sections. Within the western boundary region, the current speeds decrease from the (shallow) continental slope toward the deeper part of the sections. The easternmost mooring of both arrays shows northward flow, see Figure 12 of Dickson and Brown [1994] for the ANG array [see also Saunders, 2001, Figure 5.6.5] and Figure 11 of Bacon and Saunders [2010] for the SECF array. The reversal in flow direction marks the eastern edge of the western boundary current, and the simulated flow reversal along these two arrays agrees well with observations (see Figure 9). Vertically, the moorings of these two arrays contain multiple instruments (up to six in the ANG array), and therefore to some extent resolve the vertical structure of the velocity within the overflow layer. A direct comparison of the velocity profiles is presented in Figure 12 for both arrays. The model-based profiles reproduce to some degree the observed bottom intensification of the mean speed at the deep mooring locations, but not at shallow stations.

Along the western boundary at the ANG section, the model-determined volume transports are 7.8 and 4.6 Sv for $\sigma_\theta \geq 27.80$ kg m$^{-3}$, compared to 7.3 and 4 Sv based on observations [Dickson et al., 2008, Table 19.2]. At the SECF section, these two values are 9.4 and 5.6 Sv from the model results, in comparison to 9.0 and 4.5 Sv based on data [Bacon and Saunders, 2010]. At this point, we summarize the model-determined circulation in terms of a transport budget. The transport of DSOW into the Irminger Sea through the Denmark Strait Sea is 3.2 Sv for $\sigma_\theta \geq 27.80$ kg m$^{-3}$ and 2.9 Sv for $\sigma_\theta \geq 27.85$ kg m$^{-3}$. Downstream about halfway through the basin at the ANG section, the “net” transports across the entire Irminger Sea increase to 4.0 and 3.4 Sv for water below 27.80 and 27.85 kg m$^{-3}$, respectively. This increase is due to the entrainment of ambient water. Focusing on the SECF section, the model-determined western boundary transport of 9.4 Sv for $\sigma_\theta \geq 27.80$ kg m$^{-3}$ includes about 4.0 Sv DSOW and 3.1 Sv ISOW that flows through gaps in the Reykjanes Ridge and through the northern channel of the CGFZ. The remainder is due to recirculation (about 2.3 Sv, discussed below). For $\sigma_\theta \geq 27.85$ kg m$^{-3}$, the model-determined boundary transport of 5.6 Sv is the sum of 3.4 Sv DSOW and 2.2 Sv ISOW that flows over the Reykjanes Ridge and through the northern channel of the CGFZ.
3.5. Flow West of the CGFZ and South of Cape Farewell

West of the CGFZ, the ISOW (that has passed through the CGFZ) is generally assumed to flow northward into the eastern side of the Irminger Sea, eventually turning westward and merging with/overriding the DSOW flowing southwestward. But details of the spreading pathway(s) remain unanswered. The model-based mean circulation of deep ISOW in the area between the CGFZ and SECF mooring arrays is displayed in Figure 9. In the longitude range 36°–39°W, part of the westward flowing deep ISOW splits off to form a northward branch along the western flank of the Reykjanes Ridge. This flow merges with westward currents through gaps in the Reykjanes Ridge and penetrates into the northern Irminger Sea. The rest of the deep ISOW, however, follows a deeper and more western route, and after passing a complex topographic feature around 45°W, 56°N (called Gloria Drift), eventually reaches the area south of Cape Farewell. There the southwestward flowing Deep Western Boundary Current (DWBC) and the northeastward flowing ISOW are nearly adjacent. Several closed cyclonic and anticyclonic recirculations are formed, and mixing between two water masses is to be expected. A small portion of the deep ISOW may join the southwest flow of DWBC into the Labrador Sea, while the rest continues to flow eastward to northeastward into the Irminger Sea.

The model results also depict a large scale cyclonic recirculation west of the northward deep ISOW flow in the Labrador Sea and east of the southeastward DWBC along the Labrador coast. Below \( \sigma_\theta \geq 27.85 \text{ kg m}^{-3} \) (Figure 9a), the recirculation occupies the deepest area (>3500 m) from 51 to 58°N mostly within the Labrador Sea. In the overflow layer between \( \sigma_\theta \geq 27.80 \text{ and } 27.85 \text{ kg m}^{-3} \) (Figure 9b), the recirculation includes a well-defined eastward to northeastward “jet” south of Cape Farewell penetrating into the Irminger Sea. Transport of this jet is about 2.2 Sv, consistent with the transport balance discussed in section 3.4. This recirculation component, if realistic, raises a question regard-
ing the source of the volume transports that have been estimated along the western Irminger Sea. Dickson et al. [2008, p. 462] mentioned that the transports of DSOW (based on hydrography) appear to increase downstream from sections south of the ANG array to southeast of Cape Farewell and suggested that this increase might be due to a recirculation loop within the Irminger Sea. We should also note that in this area recirculation is a commonly observed above the overflow water in the regime of LSW [e.g., Lavender et al., 2000, 2005; Faure and Speer, 2005]. The model results suggest recirculations may exist deep in the overflow water as well.

4. Temporal Variation of the Model-Based NSOW

We now shift our focus from long-time averages to temporal variation of the model results. A 20 year integration of the simulation is relatively short in terms of reaching a model equilibrium, especially for water properties. This warrants some documentation of how the model-based NSOW evolves over the integration period. Besides, the observed transports typically are based on moored instrument records of about 1 year. It is therefore useful to examine if the model results contain nondeterministic, long-term variations that are not due to external forcing.

4.1. Overflow Sources at the Denmark Strait and Through the FBC

Dickson et al. [2008, p. 462] mentioned that the transports of DSOW (based on hydrography) appear to increase downstream from sections south of the ANG array to southeast of Cape Farewell and suggested that this increase might be due to a recirculation loop within the Irminger Sea. We should also note that in this area recirculation is a commonly observed above the overflow water in the regime of LSW [e.g., Lavender et al., 2000, 2005; Faure and Speer, 2005]. The model results suggest recirculations may exist deep in the overflow water as well.

Figure 11. (top) Model mean salinity (psu), (middle) normal velocity (cm s⁻¹), and (bottom) accumulated volume transports (Sv) along two mooring arrays in the Irminger Sea: (a) Angmagssalik and (b) Southeast of Cape Farewell. Orange triangles mark the approximate mooring locations.

After about 5 years of integration, the simulated DSOW at the Denmark Strait reached an equilibrium state, with a transport value of about 3 Sv, with T of 0–1.2°C and S of slightly less than 34.90 psu is a typical estimate.

The denominator in above equations denotes the volume transport, in which v, L, z_u, and z_l represent the normal velocity, the horizontal span, the upper, and lower interfaces of the overflow water. Transport and water properties of DSOW at or near the Denmark Strait sill have been topics of many studies [see Ross, 1984; Dickson and Brown, 1994; Rudels et al., 1999; Girton et al., 2001; Macrander et al., 2005]. A transport value of about 3 Sv, with T of 0–1.2°C and S of slightly less than 34.90 psu is a typical estimate. After about 5 years of integration, the simulated DSOW at the Denmark Strait reached an equilibrium state, with a
Figure 12. Comparison of simulated mean along slope velocities (solid line) with data from two mooring arrays in the western Irminger Sea: (a) Angmagssalik, data from Figure 19.6 of Dickson et al. [2008] and (b) Southeast of Cape Farewell, data from Table 2 of Bacon and Saunders [2010].

Figure 13. Model-based mean vectors of volume transport (per unit width, m² s⁻¹) in the area south of Cape Farewell and west of the Charlie-Gibbs fracture zone for (a) σθ ≥ 27.85 kg m⁻³ and (b) 27.80 ≤ σθ < 27.85 kg m⁻³.
transport of $3.2 \pm 0.4$ Sv, $T$ of $0.6 \pm 0.1°C$, and $S$ of $34.88 \pm 0.01$ psu. The standard deviation values are based on the monthly means of the final 5 years of integration. Overall, the model-based DSOW at the Denmark Strait agrees well with observations. The modeled transport contains a seasonal cycle that is not a characteristic of observations.

The overflow water out through the FBC has been well-quantified based on long-term, moored instrument data [see Hansen and Østerhus, 2007]. It has a volume transport of $1.9 \pm 0.3$ Sv, with $T$ and $S$ of $0.25°C$ and $34.93$ psu, respectively. The observed transport contains a seasonal variation which accounts for about 10 percent of its total value. The model-determined transport though the FBC, stable throughout the entire simulation, agrees well with observation. It also includes a seasonal cycle of similar magnitude and phase (a minimum in February and a maximum in August). The model-determined $T$ and $S$ never reached a clear equilibrium state, however, and the mean values are about $1.5°C$ and $0.1$ psu higher than observed. The resulting $\sigma_0$ is slightly lighter than observed (by $0.03$ kg m$^{-3}$), due to the opposite effects of $T$ and $S$ on density.

4.2. Variation of the Model-Based NSOW at Mooring Arrays

The time evolution of model-based ISOW at the SEI mooring array (not shown) is similar to that through the FBC. The model-determined transport ($3.3 \pm 0.43$ Sv) also contains a seasonal variation, which is not seen in the mooring results by Saunders [1996]. The transport increase, when compared to $2.2$ Sv through the FBC, is mainly due to entrained ambient water and IFR overflow water. Consequently, the average temperature is increased to $3.1 \pm 0.1°C$. The average salinity is only slightly higher, however, because the modeled overflow source water (through the FBC) has high salinity similar to that of the entrained water.

Further downstream over the Reykjanes Ridge and through the CGFZ, the time evolution of the model-determined transport and $T/S$ are illustrated in Figure 15. The transport over the Reykjanes Ridge is highly variable during the first 5 years. It decreased about $1$ Sv over the next 10 years, before leveling off in the final 5 years of the integration (Figure 15a). The decrease is largely due to the decreasing eastward flow from southern part of the CGFZ into the Iceland Basin. Through the CGFZ (Figure 15b), the model-based transport is stabilized after about 5 years of integration. The monthly transport values are more variable than upstream and a seasonal cycle no longer exists. Rather, the variation appears to be associated with the horizontal location and vertical extension of the eastward flow through the southern part of the CGFZ, somewhat similar to variations discussed by Saunders [1994] and Schott et al. [1999]. Both $T$ and $S$ exhibit a small trend that continues through the end of the 20 year simulation.

At the ANG and SECF arrays in the western Irminger Sea, the time evolution of model-based transports and $T/S$ are shown in Figure 16 for the two overflow layers. The transport evolution at the ANG array (Figure 16a) is similar to that of DSOW at the Denmark Strait, but without a clear seasonal variation. The downstream increase in transport is due to entrainment of ambient water as well as to contributions from the east (ISOW) and south (recirculation). As a result of the added transport, $T$ and $S$ rise significantly.
Figure 15. Time evolution of the model-based volume transport, average T and S of ISOW (a) over the Reykjanes Ridge and (b) through the northern channel of CGFZ. Lines, dots, and squares are results based on monthly means, annual means, and 5 year means, respectively.

Figure 16. Time evolution of the model-based overflow characteristics at the (a) ANG and (b) SECF arrays in the western Irminger Sea. Gray and black colors denote overflow layers above and below 27.85 kg m$^{-3}$. Lines, dots, and squares are results based on monthly means, annual means, and 5 year means, respectively.
tributions of DSOW and ISOW from the north and east, the flows southward into the Web along the eastern flank of the Reykjanes Ridge north of the CGFZ, and some ISOW flows westward into the Irminger Sea through gaps in addition to the westward flow through the CGFZ, some numerical simulation. The model results suggest that (1) in question, we have used a basin‐too small, when compared to upstream at the SEI array ports. Specifically, the transport of ISOW through CGFZ is considerably smaller than the transport of DSOW, after it flows over the GIS Ridge but we have primarily considered the initial development phase of the overflow water, after it flows over the GIS Ridge but before it spreads into the subpolar gyre and beyond.

Further downstream at the SECF array (Figure 16b), the transport increases further, but T and S remains almost the same. This implies that the change of T and S for overflow in the western Irminger Sea is largely due to entrainment mixing, and that occurs primarily upstream of the ANG array.

5. Summary and Discussion

[33] Nordic Seas Overflow Water, specifically DSOW and ISOW, are key components of the North Atlantic Deep Water. Knowledge of the circulation and volume transports of the overflow water masses is therefore of fundamental importance in understanding the associated meridional overturning circulation of the Atlantic Ocean. In this study we have primarily considered the initial development phase of the overflow water, after it flows over the GIS Ridge but before it spreads into the subpolar gyre and beyond.

[34] Key elements of the observational database for overflow water transport in this region are acquired from five substantially instrumented, long-term moored current meter arrays. In considering the regional circulation as sampled by these arrays, there exists an unresolved question regarding the consistency/balance of these observed transports. Specifically, the transport of ISOW through CGFZ is too small, when compared to upstream at the SEI array and to downstream at the SECF array. To help address this question, we have used a basin-scale, eddy-resolving numerical simulation. The model results suggest that (1) in addition to the westward flow through the CGFZ, some ISOW flows westward into the Irminger Sea through gaps in the Reykjaness Ridge north of the CGFZ, and some ISOW flows southward into the Web along the eastern flank of the MAR south of the CGFZ and (2) in addition to the contributions of DSOW and ISOW from the north and east, the overflow water transport at the SECF array includes some recirculation from the south.

[35] The general model-based circulation is summarized in Figure 17 and features that are pertinent to the specific questions, listed in section 1, are discussed below. Downstream of the SEI array in the Iceland Basin, flow of the model-based ISOW contains three main branches, or components, which are closely associated with bottom bathymetry (Figure 7). These branches eventually feed the different ISOW pathways: westward through gaps in the Reykjanes Ridge, westward through the northern channel of the CGFZ, and southward along the eastern flank of the MAR. In the southern part of the CGFZ, the model results suggest an eastward flow below 27.80 kg m⁻³. Part of this eastward flow turns northward and flows into the Iceland Basin, while the rest turns southward and flows along the eastern flank of the MAR.

[36] West of the CGFZ, the model results indicate two flow pathways for the deep ISOW from the CGFZ into the Irminger Sea. Some ISOW turns northward immediately west of the CGFZ and flows along the western flank of the Reykjanes Ridge. The remaining flow follows a more complicated route to reach the area south of Cape Farewell (before flowing into the Irminger Sea). There the north-eastward flowing deep ISOW and the southward flowing DWBC are adjacent and lateral interaction/mixing between them would be expected. A small amount of ISOW may join the DWBC without flowing into the Irminger Sea. Within the Irminger Sea, the deep ISOW with σₒ ≥ 27.85 kg m⁻³ and larger merges with DSOW fairly uniformly (as a function of latitude). In contrast, more water in the layer of σₒ between 27.80 and 27.85 kg m⁻³ merges with DSOW north of the ANG array than south (Figure 10).

[37] The model results in addition depict recirculations in the overflow water (Figure 13). Below 27.85 kg m⁻³, a large scale recirculation is present in the deepest part (≥3500 m) of the Labrador Sea. Some small scale recirculations are also present in the area south and southeast of Cape Farewell and in the southern part of the CGFZ. Between 27.80 and 27.85 kg m⁻³, a prominent component of the recirculation is depicted as an eastward to northeastward jet south of Cape Farewell that penetrates into the Irminger Sea. This model-based recirculation jet accounts for about 2.2 Sv out of the 9.4 Sv total transport for composite overflow water calculated through the SECF array.

[38] The results obtained from moored instrument arrays have provided a valuable benchmark for evaluating the performance of eddy-resolving, large-scale ocean models in simulating the downstream circulation of the overflow water source. The evaluation in this study includes detailed comparisons with observations along the mooring arrays, including the model-based mean volume transport, the vertical/lateral distribution of the water properties (especially salinity) associated with different water masses, as well as the velocity profiles of the overflow plume in the western Irminger Sea. These comparisons suggest that the model-based circulation is generally consistent with the available data. We have also investigated the temporal evolution of the model-based NSOW transports as well as T/S properties. While the T/S has some long-term trend, the transports are quite stable, especially in the final 5 years of the integration. Figures 14–16 suggest that a 1 year mean transport represents the long-
term mean well, probably in part due to the use of climatological forcing.

[39] There are, however, several aspects of the model results that need further improvement. The model-based ISOW through the CGFZ is slightly denser than observed and consequently flows at about 500 m deeper. This suggests that the diapycnal mixing in the model is too weak when the overflow water, after passing over the sill, descends along the northern and western slope of the Iceland Basin. This notion is similar to the experience with isopycnic model simulation of overflow water that is summarized in Willebrand et al. [2001], and an entrainment parameterization may be the key to improving the results. Another issue is that the modeled ISOW source in the FBC is warmer, more saline, and slightly less dense than observed. The exact reason is not clear. Diapycnal mixing and the T/S properties of MNAW and overflow source water in the Nordic Seas might be responsible. The model also does (does not) reproduce the observed bottom intensification of the overflow velocity profiles on the deep (shallow) part of continental slope in the western Irminger Sea. Increasing the vertical resolution does not improve this result. The transition of overflow water from pressure coordinates at the sill toward isopycnic layers along the slope might be important.

[40] Further numerical studies, with tools like passive tracers and/or numerical drifters to separate flows of different origins, would be helpful in better quantifying the circulation features described. We should also note that the simulation we have discussed does not simulate the northward flow of Antarctic Bottom Water (AABW). The model-based overflow water consequently misses the contribution from modified AABW, or Lower Deep Water [LDW, McCartney 1992], in the northern North Atlantic. As noted in section A3, LDW mixes with the ISOW that flows westward through the CGFZ and its (high silicate) signature is also present in the northern Irminger Sea. New experiments with improved representation of AABW/LDW are needed.

Appendix A: Specific Comments on the Observational Base for Overflow Transports

[41] Appendix A is a detailed discussion of the database relevant to overflow transports in the western Irminger Sea and in the Iceland Basin, as listed in Table 1 in section 2.1. It also includes a comment concerning a possible influence on the ISOW characteristics in the Irminger Sea related to a contribution by Lower Deep Water (LDW) [McCartney, 1992]. The composite overflow water has been typically defined as having potential density ($\sigma_\theta$) of 27.80 kg m$^{-3}$ and larger. In the recent determinations and discussions of DSOW transports in the western Irminger Sea, Dickson et al. [2008] and Bacon and Saunders [2010] identify the boundary between DSOW and ISOW as having $\sigma_\theta$ of 27.85 kg m$^{-3}$. As pointed out by these authors, the hydrographic data used in determining $\sigma_\theta$ boundaries are typically acquired on setting, retrieval, and/or turn-around research cruises. As a result, the thickness (of overflow water) is less well determined in the mean than the velocity measured by using extensive long-term moored array data acquisition at dense temporal sampling intervals. We should also note that these typically yearlong moored instrument arrays at different sites cover different periods over a total span of about two decades, and while a long-term trend in overflow transport has not clearly been detected, interannual variability has.

A1. Overflow Transports in the Western Irminger Sea

[42] Array 1 (in Figure 1) is colocated with the hydrographic “Transient Tracers in the Ocean” (TT0) section described by Livingston et al. [1985]. The salinity distribution along this TTO section has been used to identify different water masses in the Irminger Sea and a modified version is presented later in this discussion. The estimated composite overflow transport of 5.1 Sv is based on a moored instrument array with a record length of 368 days [Dickson and Brown, 1994]. They also found a similar value (5.2 Sv) 160 km upstream at Dohlm Bank, based on another moored instrument array of shorter duration. The transport of DSOW alone is not available at the TTO section, but Table 19.1 of Dickson et al. [2008] implies that it is close to that found at the ANG section.

[43] Array 2 is called the ANG array after its location offshore of Angmagassalik, Greenland. On their Figure 19.6, Dickson et al. [2008] have plotted the cross-sectional distribution of the mooring locations along with mean isopycnals of $\sigma_\theta = 27.80$ and 27.85 kg m$^{-3}$. Most of their current meters are located between the 27.85 kg m$^{-3}$ contour and the bottom, sampling the bottom-intensified flow associated with DSOW. Estimates of mean speeds in the regime of ISOW, identified with 27.80 $\leq \sigma_\theta < 27.85$ kg m$^{-3}$, are more sparsely located. The mean (downslope) current speeds for the period 1986 to 2005 are also listed in their Figure 19.6. Note that there is a large gap during 1990–1997, but these results nevertheless are very long and constitute a tremendous achievement in their acquisition. The estimates of transport (see their Table 19.2) through this ANG Array are 4 Sv for DSOW, and 7.3 Sv for the composite overflow. In their Table 19.1, Dickson et al. [2008] also estimated the interannual variability of DSOW transport. These estimates were based on a composite of three hydrographic sections shown in their Figure 19.10 (the center one being ANG). The transports varied from 2.5 to 5 Sv from 1997 through 2005, with an average of 3.6 Sv.

[44] Array 3 runs southeast of Cape Farewell, Greenland (here called SECF). In their Figure 2, Bacon and Saunders [2010] plotted a cross-sectional view of the mooring locations in relationship to the mean depth of the density surfaces ($\sigma_\theta = 27.80$ and 27.85 kg m$^{-3}$). Similar to the ANG array, most of the current meters were deployed at depths below $\sigma_\theta$ of 27.85 kg m$^{-3}$ for a similar reason (a focus on bottom-trapped currents). The roughly yearlong transports were found to be 4.5 Sv for DSOW, and 9 Sv for the composite overflow. The contribution of ISOW is therefore also about 4.5 Sv. For continuity, this 4.5 Sv of ISOW should arrive at the SECF section from the Iceland Basin, east of the Reykjanes Ridge. Bacon and Saunders [2010] also discussed the temporal variability of their records, and suggest that some interannual variability is present, but not dominant. Interannual variability in DSOW transport is also shown at the ANG array [Dickson et al., 2008, Figure 19.7]
and at the Denmark Strait sill [Macrander et al., 2005, Figure 2].

[45] A transport of 9 Sv for $\sigma_0 \geq 27.80$ kg m$^{-3}$ at the SECF array is rather low in comparison with the previous estimate of 13.3 Sv across a section due south of Cape Farewell by Clarke [1984], as adopted by Dickson and Brown [1994]. The direct current measurements by Clarke [1984] include six current meters on three moorings for 60 days. So there could be a sampling issue in comparing this sparse data set with more substantially instrumented long-term direct measurements. The transport estimates by Clarke [1984] also involved CTD data taken along a section running south of Cape Farewell to Flemish Cap. The issue with respect to estimating the location of $\sigma_0$ surfaces could also apply to Clarke’s data set. The contour for $\sigma_0 = 27.80$ kg m$^{-3}$ is about 800 m in the immediate vicinity of Cape Farewell in Figure 4c of Clarke [1984], compared to about 1400 m in Figure 2 of Bacon and Saunders [2010].

On his p.128, Clarke [1984] noted that his (transport) values are “useful estimates of the probable order of magnitudes of the transports of various currents and water types” as indicated by his data. Bacon and Saunders [2010] in their section 4.3 discussed the geostrophic reference level choices by Clarke [1984] in detail and concluded that the difference between 9 and 13 Sv (Bacon and Saunders actually suggest 13–16 Sv as a defensible range of values) could be due to temporal variability on interannual time scales.

[46] Our model results suggest the existence of considerations other than interannual variability for the differences in the overflow transports being considered. One possibility could be a change in the flow regime between south and southeast of Cape Farewell. Relative to the SECF section in the Irminger Sea, the overflows south of Cape Farewell involve a more pronounced presence of horizontal recirculations (on a variety of horizontal scales) from the Labrador Sea that could play a role in the transport difference. Within the Labrador Sea and close to the bottom, there may exist a large scale horizontal recirculation of overflow water [see Clarke, 1984, Figures 2 and 4]. This recirculation could also connect with the different overflow transports observed by Schott et al. [2004] east of the Grand Banks of Newfoundland and by Fischer et al. [2004] at the exit of the Labrador Sea. This speculation has not been suggested elsewhere as far as we know, but seems worthy of further examination.

A2. Overflow Water in the Iceland Basin

[47] Array 4 (in Figure 1) is located southeast of Iceland (here called SEI). The estimate of the mean transport for ISOW is 3.2 Sv at $\sigma_0 \geq 27.80$ kg m$^{-3}$. Here Saunders [1996], his Figure 3 observed a well-defined bottom-trapped mean flow up on the western slope between the 1300 and 2200 m depth contours, while in contrast only a small mean flow is seen in the deep Iceland Basin. For comments on ISOW flowing into and within the deep Iceland Basin [see Hansen and Østerhus, 2000, Figures 43 and 49]. Saunders [1996] stated that the flow at the SEI array included only a very small contribution from over the IFR. However, Saunders [2001] changed the composition of the total flux of 3.2 Sv of ISOW to be 1.6 Sv of FBC overflow water, mixed with 0.8 Sv of entrained MNAW (7–8.5°C and 35.1–35.3 psu) and 0.8 Sv of Modified East Icelandic Water (MEIW, 1–3°C and 34.7–34.9 psu), and remarked that MEIW could only have come from over the IFZ. A more detailed description of MEIW as an overflow water mass crossing the IFZ may be found in the work of Hansen and Østerhus [2000, section 5.2].

[48] Array 5 was located along a north–south section across the CGFZ near 35°W. In his Figure 5, Saunders [1994] illustrates the distribution of mooring locations relative to the isopycnal depth of 27.80 kg m$^{-3}$ (27.85 kg m$^{-3}$ is not shown), and to the salinity contour of 34.94 psu. The salinity distribution, also shown in his Figure 2, is the actual control used to identify ISOW. The estimated mean transport of 2.4 Sv is at $S \geq 34.94$ psu but could be applied to $\sigma_0 \geq 27.80$ kg m$^{-3}$ as well. The transport is dominated by temporal variability [Saunders, 1994, Figure 9]. The author speculated that the comparatively large amplitude overflow reversals observed might be connected to latitudinal migrations of the North Atlantic Current over the CGFZ, a speculation that was later found to be the case by subsequent observations [Schott et al., 1999]. In an earlier study of deep flow within the CGFZ, Shor et al. [1980] found a cold freshwater type beneath ISOW with a flow tendency to the east, at least partially similar to the observations by Saunders [1994] and Schott et al. [1999].

[49] Hansen and Østerhus [2000] emphasized that a westward transport of 2.4 Sv in CGFZ seems too small for the total flow of ISOW over the GIS Ridge. It is also smaller than the 3.2 Sv observed at the upstream SEI section [Saunders, 1996]. The problem is augmented if flow through the CGFZ also contains a contribution from LDW. To reconcile this discrepancy, Hansen and Østerhus [2000] suggested a wide variety of possible ISOW paths in their Figure 54. One of these paths goes through a gap in the Reykjanes Ridge near 55°N to the north of the CGFZ. It turns out that the estimate for ISOW transport in the CGFZ by Shor et al. [1980] also yields 2.4 Sv to the west, seeming rather strong support for the surprisingly low transport found by Saunders [1994]. The size of the transport estimate by Shor et al. [1980] was not discussed by Saunders [1994] or Saunders [2001], probably due to the relatively limited long-term direct measurements involved.

[50] Interestingly, Shor et al. [1980] in their Figure 1 also suggest that some ISOW could leak westward through fracture valleys near 56.5°N into the Irminger Sea. Their premise was influenced by Vogt and Johnson [1973]. Based on a seismic reflection profile parallel to the axis of the Reykjanes Ridge, Vogt and Johnson [1973] suggested that the Reykjanes Ridge was broken by “many” transverse valleys between 52 and 57°N. They further pointed out that the asymmetry of the depositional sediments therein might imply that westward flow (of ISOW) crosses the Reykjanes Ridge. Over time, a few other authors in addition to Shor et al. [1980] and Hansen and Østerhus [2000] have questioned a 2.4 Sv transport of ISOW through the CGFZ and have noted it presents a conundrum because it is too small relative to observed transports of overflow water in the western Irminger Sea. Some of these other authors have suggested (but not demonstrated from either direct current measurements, nor tentatively illustrated by model results) a flow of ISOW through the gaps in the Reykjanes Ridge into the Irminger Sea.
A3. A Comment on an LDW Contribution

[51] McCartney [1992] suggested that modified Antarctic Bottom Water (AABW), or Lower Deep Water (LDW), penetrates into the northernmost part of the eastern North Atlantic Ocean, and mixes with the overflow transport measured along the western Irminger Sea. The presence of LDW in the Iceland Basin, identified with high silicate content, and the interaction with ISOW there were also discussed by van Aken and de Boer [1995] and van Aken and Becker [1996]. Because of LDW influence, the silicate content of the ISOW through the CGFZ (13–17 \( \mu \text{mol l}^{-1} \)) is significantly higher than that at the SEI array (9–10 \( \mu \text{mol l}^{-1} \)) [see Saunders, 1996, Figure 2d].

[52] In order to help identify a possible influence of LDW on the water masses in the Irminger Sea, the distributions of salinity and silicate along the TTO section are presented in Figure A1. The location of the section is shown as blue dots in Figure 1. The salinity plot was originally published by Livingston et al. [1985, Figure 2a] and was modified in Dickson and Brown [1994, Figure 3]. Here the salinity contours are colorized with some added text indicating different water masses. The vertical stacking of DSOW, ISOW, LSW, as well as the saline MNAW, in the upper level Irminger Current (IRC) and East Greenland Current (EGC) is relatively well known. Here we comment that the broad area of ISOW (S \( \geq 34.94 \) psu) contains different levels of silicate (or tritium) [see Livingston et al., 1985, Figure 2b]. The deep portion (of ISOW) is associated with a high silicate core of 12–13 \( \mu \text{mol l}^{-1} \). The shallow portion that exists over the western slope of the Reykjanes Ridge between 1100 and 2000 m has a silicate content of 10–11 \( \mu \text{mol l}^{-1} \). One possibility, consistent with our numerical results showing ISOW flowing over the Reykjanes Ridge, is that the shallow portion of the ISOW is perhaps composed of diluted IFR overflow water and entrained MNAW and that this water flows over the Reykjanes Ridge into the Irminger Sea without much LDW influence; while the deep portion of the ISOW is FBC water mixed with LDW that flows through the CGFZ before turning northward into the Irminger Sea.

Acknowledgments. This work is a contribution to the project “U.S.-GODAE: Global Ocean Prediction using the HYbrid Coordinate Ocean Model (HYCOM)” funded under the National Ocean Partnership Program and to the 6.1 project “Global Remote Littoral Forcing via Deep Water Pathways” funded by the Office of Naval Research under program element 601153N. Simulations were performed on the supercomputer at the Naval Oceanographic Office, Stennis Space Center, MS, provided by the U.S. Defense Department High Performance Computing Modernization Program. The authors thank A. Wallcraft and E. J. Metzger from NRL/SSC for help on the configurations of numerical experiments. Figures 4a and A1b are plotted using Ocean Data View (http://odv.awi.de). This is Naval Research Laboratory publication number NRL/JA/7304‐10‐251, and it has been approved for public release.

References


E. P. Chassignet, Center for Ocean-Atmospheric Prediction Studies, Florida State University, 2035 E. Paul Dirac Dr., 200 RM Johnson Bldg., Tallahassee, FL 32306-2840, USA.

P. J. Hogan and H. E. Hurlburt, Oceanography Division, Naval Research Laboratory, Building 1009, Stennis Space Center, MS 39529, USA.

W. J. Schmitz Jr., Harte Research Institute, Texas A&M University, 6300 Ocean Dr., Corpus Christi, TX 78412-5869, USA.

X. Xu, Department of Marine Science, University of Southern Mississippi, 1020 Balch Blvd., Stennis Space Center, MS 39529, USA. (xiaobiao.xu@usm.edu)