**Title and Subtitle:**
Soot and NOx Emissions and Combustion Characteristics of Low Heat Rejection Direct Injection Diesel Engines

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**Performing Organization:**
Greensboro, NC 27411

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**Abstract:**
Performance and emissions data were gathered on a normally aspirated single-cylinder DI engine with various combinations of ceramic coatings installed. Thin ceramic thermal barrier coatings were applied to the piston crown and bowl, the head and valves, and the cylinder liner. The coated piston and head were run singly and in combination with the cylinder liner to investigate the effects of these different coated surfaces on emissions and performance. Coating the piston crown alone results in generally lower cylinder pressure, lower brake specific fuel consumption and lower NOx emission compared to the baseline engine. Soot emission is typically increased below 2000 RPM and decreased above 2000 RPM. Coating the head alone reduces cylinder pressure, but generally increases specific fuel consumption and NOx and soot emission. The KIVA-II code was used to model the Hydra engine with the thermal coatings. The computer modeling has led to an understanding of the effect of coating the piston on NO production. The hotter piston crown warms the intake air, shortening ignition delay and decreasing the ratio of premixed to diffusion combustion, ultimately resulting in lower peak cylinder temperature and reduced NOx. The KIVA-II results agree reasonably well with the experimental data for cylinder pressure and NO and soot emission.
FINAL REPORT

SOOT AND NOx EMISSIONS AND COMBUSTION CHARACTERISTICS
OF LOW HEAT REJECTION DIRECT INJECTION DIESEL ENGINES

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By

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SUMMARY

An experimental and analytical study on the effect of thin ceramic coatings on soot and NOx emissions and performance of a direct injection diesel engine was conducted jointly between North Carolina A&T State University and North Carolina State University.

Performance and emissions data were gathered on a normally aspirated Ricardo Hydra single cylinder DI engine with various combinations of ceramic coatings installed. Thin ceramic thermal barrier coatings were applied to the piston crown and bowl, the head and valves, and the cylinder liner. The coated piston and head were run singly and in combination with the cylinder liner to investigate the effects of these different coated surfaces on emissions and performance for two different pure hydrocarbon fuels, hexadecane and dodecane. Coating the piston crown alone results in generally lower cylinder pressure, lower brake specific fuel consumption and lower NOx emission compared to the baseline engine. Soot emission is typically increased below 2000 RPM and decreased above 2000 RPM. Coating the head alone reduces cylinder pressure, but generally increases specific fuel consumption and NOx and soot emission.

The analytical portion of the study involved modifications to the KIVA-II code and its use to model the Hydra engine with the thermal coatings. Modifications to the code include incorporation of an eddy-break-up combustion model, to replace the standard Arrhenius single reaction model. A time dependent combustion surface temperature was also incorporated to simulate the effects of thermal barrier coatings on cylinder temperature, pressure and NOx production. A soot model was added to the code, following the work of Magnussen and Hjertager. The EBU model gives better results for the diffusional portion of the combustion process, but fails to adequately model the premixed combustion, typically resulting in a lower predicted peak cylinder pressure (and temperature) than predicted by the single reaction model and shown by experiments. Consequently, the EBU model also under-predicts NO emission to a greater degree than the SR model. The KIVA-II modeling has led to an understanding of the effect of coating the piston on NO production. The hotter piston crown warms the intake air, shortening ignition delay and decreasing the ratio of premixed to diffusion combustion, ultimately resulting in lower peak cylinder temperature and reduced NO. The KIVA-II results agree reasonably well with the experimental data for cylinder pressure and NO and soot emission.
LIST OF PUBLICATIONS

Technical Papers Published


Papers in Preparation

"Effect of Selective Insulation on the Combustion, Performance and NO Emission of DI Diesel Engines"

"Modeling of Soot and NO Emissions in a DI Diesel Engine: KIVA-II Versus Experiments"

Theses


LIST OF PERSONNEL AND DEGREES EARNED

David E. Klett, Project Director, Professor of Mechanical Engineering, North Carolina A&T State University

E. M. Affy, Co-Investigator, Professor of Mechanical Engineering, North Carolina State University

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M. Vakil, Graduate Research Assistant, MS, North Carolina A&T State University, Expected 5/94

Nigel Brooks, Undergraduate Research Assistant, BS, North Carolina A&T State University, Expected 12/94

Martin Ashley, Undergraduate Research Assistant, BS, North Carolina A&T State University, 5/92

Raphael Jones, Undergraduate Research Assistant, North Carolina A&T State University, 5/93

Postoria, Aguirre, Undergraduate Research Assistant, North Carolina A&T State University, 5/92
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This work was supported by the U. S. Army Research Office under grant number DAAL03-89-G-0103. D. E. Klett was the project director and principal investigator at North Carolina A&T State University. E. M. Afffy was the principal investigator at North Carolina State University, and M. Saad was a research associate on the project. The authors would like to express their appreciation to David Mann of the Army Research Office for his assistance, advice and encouragement during the duration of this project.
1. INTRODUCTION

1.1 PROBLEM STATEMENT

The objective of this work was to study the effects of thin ceramic coatings applied to various internal engine surfaces on the performance and emissions of a direct injection (DI) diesel engine. The study included an experimental investigation coupled with computer modeling using the KIVA II code to gain insight into the combustion effects caused by higher surface temperatures resulting from the use of ceramic coatings. The project was a joint effort between North Carolina A&T State University, where the experiments were carried out, and North Carolina State University, where the bulk of the computer modeling was performed.

The Army, among a number of other agencies and corporations, has maintained an interest in the use of ceramics in diesel engines for a number of years, and many studies have been conducted on various aspects of this topic [1-8]. This study focused on the use of KIVA II to model the combustion process in a particular DI diesel engine with thin ceramic coatings applied to the piston crown, the cylinder head inner surface and the cylinder liner. Experiments were performed on a single-cylinder Ricardo Hydra DI diesel research engine to provide a data base for comparison with the modeled results. Of particular interest were the effects of the coatings on brake specific fuel consumption, NOx emission and soot emission. Two pure hydrocarbon fuels, Dodecane \((\text{C}_{12}\text{H}_{26})\) and Hexadecane \((\text{C}_{16}\text{H}_{34})\), were used for the experiments and in the combustion modeling. These fuels were chosen for purposes of reproducibility of experimental results and because the properties are available to facilitate the modeling.

1.2 BACKGROUND

The use of ceramic coatings on combustion chamber surfaces influences the performance and exhaust emissions of DI diesel engines [1-4]. This can be attributed to changes in the combustion process brought about by the modification of the thermal boundary conditions due to the lower thermal conductivity of ceramics relative to metals. The added thermal resistance of thin low conductivity coatings increases the surface temperature, and higher surface temperatures affect the combustion process in a variety of ways, depending on which combustion chamber surfaces are coated.
To better understand the effect of insulating the combustion chamber walls on the performance and exhaust emissions of the engine, a brief review of the combustion process in DI diesel engines is first presented.

The combustion process in DI diesel engines involves two modes of combustion, usually referred as the premixed and the diffusion combustion modes [9]. Premixed combustion occurs early in the process when fuel which has evaporated and mixed with air during the ignition delay period auto ignites. This mode is accompanied by a high rate of heat release which produces a rapid rate of pressure rise. When the premixed fuel/air mixture is depleted, diffusion combustion, characterized by a lower rate of heat release, takes over and controls the remainder of the combustion process.

Modification of the engine design, or variation of its operating conditions, will affect the combustion process in such a way as to vary the ratio of the time duration of the two modes, $t_p/t_d$, where $t_p$ represents the duration of the premixed mode and $t_d$ represents the duration of the diffusion mode. This ratio plays an important role in determining the effect of engine parameters and other factors on the performance and emissions of the engine. Factors which increase the ignition delay period will increase the premixed combustion duration while decreasing the diffusion combustion duration, causing an increase in the ratio $t_p/t_d$. This may lead to higher peak cylinder pressure and temperature which may improve thermal efficiency and reduce CO and unburned hydrocarbon (UHC) emissions at the expense of increasing NOx emissions of the engine. Large increases in $t_p/t_d$ will cause a high rate of pressure rise and lead to objectionable diesel knock. Factors such as engine speed, advanced injection timing and the use of low cetane fuels contribute to longer ignition delay and the increase in $t_p/t_d$. Reducing the ignition delay period causes the premixed combustion duration to decrease while increasing the diffusion combustion duration, i.e. reduces $t_p/t_d$, prolonging the combustion process. Large decreases in $t_p/t_d$ may cause the diffusion combustion mode to dominate which may lead to loss of power, decrease in thermal efficiency and possible deterioration of engine exhaust emissions. Preheating the intake air and insulating combustion chamber surfaces are among the factors which can decrease ignition delay and $t_p/t_d$.

Several experimental studies on low heat rejection (LHR) DI naturally aspirated diesel engines [1,3,4,5,6] showed that engine performance and exhaust emissions suffered from lengthening of the combustion process which is indicative of an extreme reduction in the ratio $t_p/t_d$. 

2
1.3 PERTINENT LITERATURE REVIEW

Few investigators have studied the effect of insulating various combustion chamber surfaces in DI diesel engines and the results have been contradictory in some cases. Miyairi, et. al [3] studied the effect of selective insulation of the cylinder head, piston crown, and cylinder liner, using thick monolithic ceramic inserts, on the performance and emission characteristics of a single-cylinder, normally-aspirated DI diesel engine. They showed that fuel economy and NO emissions of the engine were improved by insulating the cylinder head and liner, but were made worse by insulating the piston crown. Part of the degradation in BSFC with the insulated piston was attributed to the increased reciprocating mass due to the heavy monolithic ceramic piston crown.

Assanis, et. al [2] conducted a series of tests on a supercharged DI diesel engine with and without piston surface insulation to determine the effect of ceramic coating the piston crown on engine performance and emissions. In their study, they emphasized the significance of the heat release profile, and indicated that insulating the piston with a thin coating of PSZ resulted in better engine efficiency and reduced emissions over the baseline engine.

Dickey [1] studied the effect of applying thin ceramic coatings to all combustion chamber surfaces in a supercharged single-cylinder Caterpillar 1Y-540 DI diesel engine. The results showed decreased thermal efficiency, but also decreased specific NOx and UHC for the ceramic coated engine relative to the baseline engine, especially at higher loads.

Daby, et. al [10] conducted tests on a single-cylinder, uncooled, pressurized-intake, DI diesel engine equipped with thick monolithic ceramic inserts on the head and cylinder liner (above top ring travel). The piston was a low-heat-pass articulated design with a ferrous metal top and an aluminum skirt. These tests were conducted to investigate the behaviour of alternative fuels (low cetane, higher density) in conjunction with an uncooled engine. Basline tests were conducted with Phillips diesel control fuel D-2 in the uncooled engine and in a baseline cooled metal engine. The baseline engine, however, had a different stroke and compression ratio than the ceramic engine, and direct comparisons of performance and emissions for the two engines would not be meaningful.
2. DESCRIPTION OF EXPERIMENTS

2.1 EXPERIMENTAL APARATUS

The experimental data used for comparison with the KIVA results were obtained from a Ricardo Hydra single-cylinder DI diesel research engine. Table 1 provides the basic specifications of the engine and Figure 2.1 provides a cross-sectional view of the combustion chamber showing the toroidal bowl piston. Measured data included load, speed, crank angle, needle lift, cylinder and fuel line pressure, fuel mass flow rate, air volume flow rate, intake and exhaust gas temperatures, coolant and lube oil temperatures, exhaust smoke opacity and NO and NOx emissions.

Engine loading was via a digitally-controlled DC motoring dynamometer interconnected to the utility grid through a KTK power control and signal conditioning system. Speed, load and injection timing could be remotely controlled from an instrumentation panel which provides readouts of these parameters plus coolant, lubricant, intake and exhaust manifold temperatures. Additional digital thermocouples were used to monitor room air temperature, NOx sampling line temperature and air filter intake temperature. A sling psychrometer was used to determine combustion air humidity. The engine coolant temperature and oil temperature could be individually controlled with a pair of thermostats that regulate the flow of laboratory water through heat exchangers that provide cooling for the engine fluids.

Exhaust smoke opacity was measured with a USPHS Diesel Smokemeter permanently mounted in the exhaust pipe 75 mm downstream from the exhaust port. The smoke meter was calibrated against a set of neutral density gelatin filters. Oxides of nitrogen were measured with a Thermo Environmental Instruments Model 10AR chemiluminescent NO-NOx analyzer in conjunction with a Model 800 heated sample conditioning unit and a heated sampling line. The NOx system was calibrated periodically against 2000 ppm NO in nitrogen calibration gas.

Intake air flow was measured with a Meriam Model 50MC2-2F Laminar Flow Element equipped with a digital manometer to read the pressure differential, and an inclined water manometer for calibration purposes. Fuel mass flow rate was monitored with an AVL Model 730 gravimetric fuel balance.
The cylinder pressure, fuel pressure, needle lift and crank angle were recorded with a Data Precision DATA 6000 digital wave form analyzer at a data rate of 25 KHz. The data was stored on diskettes for subsequent computer analysis. Table 2 gives the test matrix of load, speed and injection timings used for the experiments.

Ignition delay and injection duration were determined from the digital wave form data. The 25 KHz data rate translates to 40 microseconds per data point, which then represents the uncertainty for both the ignition delay and the injection duration.

Figure 2.1
Cross Section of Ricardo Hydra DI Diesel Combustion Chamber
<table>
<thead>
<tr>
<th>TABLE 1 - SPECIFICATIONS OF RICARDO HYDRA DI DIESEL ENGINE</th>
</tr>
</thead>
<tbody>
<tr>
<td><strong>Engine</strong></td>
</tr>
<tr>
<td>Number of Cylinder</td>
</tr>
<tr>
<td>Bore</td>
</tr>
<tr>
<td>Stroke</td>
</tr>
<tr>
<td>Swept Volume</td>
</tr>
<tr>
<td>Maximum Speed</td>
</tr>
<tr>
<td>Maximum Power</td>
</tr>
<tr>
<td>Maximum Cylinder Pressure</td>
</tr>
<tr>
<td>Compression Ratio</td>
</tr>
<tr>
<td>Connecting Rod Length</td>
</tr>
<tr>
<td>Squish Height</td>
</tr>
<tr>
<td>Swirl</td>
</tr>
<tr>
<td>Valve Timing:</td>
</tr>
<tr>
<td>Intake Opens</td>
</tr>
<tr>
<td>Intake Closes</td>
</tr>
<tr>
<td>Exhaust Opens</td>
</tr>
<tr>
<td>Exhaust Closes</td>
</tr>
<tr>
<td><strong>Injection System</strong></td>
</tr>
<tr>
<td>Injector Pump</td>
</tr>
<tr>
<td>Nozzle</td>
</tr>
<tr>
<td>Nozzle Opening Pressure</td>
</tr>
<tr>
<td>Injector</td>
</tr>
<tr>
<td>Lift Pump</td>
</tr>
<tr>
<td><strong>Dynamometer</strong></td>
</tr>
<tr>
<td>Manufacturer</td>
</tr>
<tr>
<td>Type</td>
</tr>
<tr>
<td>Rating</td>
</tr>
<tr>
<td>Max. Speed</td>
</tr>
<tr>
<td>Control</td>
</tr>
</tbody>
</table>
2.2 CERAMIC COATINGS

The piston, cylinder head, valves and cylinder liner were sent to Adiabatics, Inc. in Columbus, IN for coating. The piston crown and bowl received a 0.25 mm slurry-sprayed coating of partially stabilized zirconia (PSZ) consisting of 85 percent partially calcium stabilized cubic zirconia, 10 percent tungsten cobalt chrome powder, and 5 percent chrome oxide. The head and valves were coated with a 0.5 mm thermal barrier coating incorporating 5 percent hollow alumina spheres in a slurry of 65 percent silica, 15 percent PSZ, 7 percent tungsten chrome powder, and 8 percent chrome oxide.

The cylinder liner was bored 0.2 mm over the entire length plus an additional 0.75 mm in the region above the top ring reversal (TRR). The region above the TRR was plasma sprayed with yttria stabilized zirconia, and then the entire length of the liner was given a 0.2 mm wear coat of slurry sprayed PSZ. The coatings were sealed and densified with a chrome oxide based drain cast slurry.

2.3 TEST MATRIX

A full set of data runs were conducted with each of the two fuels, hexadecane and dodecane, and for each of four different engine builds utilizing different combinations of coated components including: (1) baseline (no coatings), (2) coated piston alone, (3) coated head alone, and (4) coated piston, head and liner together. A full set of runs consisted of the test matrix shown in Table 2 involving four engine speeds, four loads and four injection timings, for a total of 64 runs per engine build, per fuel (512 total runs). After initially investigating the effect of engine coolant temperature on NOx production and observing the expected increase in NOx with increasing temperature, as shown in Figure 2, it was decided to use a constant coolant temperature of 353 K and constant oil temperature of 323 K. A sample data set for hexadecane for one engine speed is shown in Table 3.

<table>
<thead>
<tr>
<th>INJECTION TIMING (DEG BTDC)</th>
<th>16</th>
<th>18</th>
<th>20</th>
<th>22</th>
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<td>3.91</td>
<td>4.47</td>
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<tr>
<td>Speed (RPM)</td>
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<td>2000</td>
<td>2500</td>
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</tr>
<tr>
<td>--------------------------------</td>
<td>------------</td>
<td></td>
<td></td>
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<tr>
<td></td>
<td>1</td>
<td>2</td>
<td>3</td>
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<tr>
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<td>8</td>
<td>8</td>
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<td>20</td>
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<tr>
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<td>WATER TEMP. K</td>
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<td>353</td>
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<td>515</td>
<td>519</td>
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<tr>
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<td>47</td>
<td>46</td>
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<tr>
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<td>.234</td>
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<td>.335</td>
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<tr>
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<td>1750</td>
<td>1900</td>
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<td>1125</td>
<td>1250</td>
<td>1450</td>
<td>1600</td>
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<tr>
<td>FUEL INJ. g/min</td>
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<td>8</td>
<td>8.3</td>
<td>8.6</td>
</tr>
<tr>
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<td>0.72</td>
<td>0.8</td>
<td>0.84</td>
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<tr>
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<td>FUELLING RATE *10⁻³ (g/inj)</td>
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<td>1106</td>
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<td>353</td>
<td>353</td>
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<td>40</td>
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<td>.428</td>
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<td>11.9</td>
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3. EXPERIMENTAL RESULTS

3.1 PRESENTATION OF DATA

Plots of data for all cases investigated are provided in the appendices. This includes data for fuel consumption (Appendix A), NOx (Appendix B), soot (Appendix C), ignition delay (Appendix D) and exhaust gas temperature (Appendix E). All of the data was included in the appendices in order to have it organized and available for scrutiny by interested parties. A page typically contains four plots, representing four different engine speeds, of the above variables versus load or timing for a common value of the remaining independent variable. Figures for hexadecane fuel precede those for dodecane fuel in each appendix. The following provides a discussion of general trends observed. For convenience, certain example data plots from the appendices are reproduced in the following sections to illustrate these trends. Thus one can rely simply on the examples used here for illustration, or refer to the appendices for additional in-depth evaluation of all cases.

3.2 EFFECT OF COATINGS ON CYLINDER PRESSURE

The effect of coating various combustion chamber surfaces on the cylinder pressure is illustrated in Figure 3.1. The four plots in this figure are for different injection timings. All of the data is for hexadecane fuel, with an engine speed of 1500 RPM and a load of 447 KPa BMEP (the maximum load used for these tests). In all cases, the baseline engine produces the highest cylinder pressure followed in descending order by the coated head, all surfaces coated, and coated piston cases. Some decrease in maximum cylinder pressure can be attributed to decreased volumetric efficiency that occurs with the coated surfaces (all tests were performed under normally aspirated conditions), but changes in the combustion process caused by the surface coatings accounts for most of the decrease as discussed in detail in Chapter 4. This is evident since operating the engine with all of the surfaces coated results in the lowest volumetric efficiency, while coating only the piston crown consistently produces the lowest pressure. Coating only the piston effects the combustion process to the greatest degree of all the insulation schemes and produces the lowest fuel consumption and NOx emission as described in the following sections.
Figure 3.1

Cylinder Pressure With Hexadecane Fuel at 1500 RPM and High Load
3.3 EFFECT OF COATINGS ON BSFC

Figure 3.2 compares the brake specific fuel consumption (BSFC) for the four different engine builds as a function of load (BMEP) for an injection timing of 20 DBTDC for hexadecane fuel. Figure 3.3 makes the same comparison for dodecane fuel. Each graph in these figures represents a different engine speed. The injection timing value of 20 DBTDC was chosen arbitrarily for purposes of these comparisons. The same trends are present for the other injection timings tested. The full set of BSFC data is included in Appendix A. The effect of varying the injection timing is to increase BSFC with advancing timing as indicated in Figure 3.4. The general trend exhibited in Figures 3.2 and 3.3 is that specific fuel consumption for the coated-piston case is better than baseline for a number of operating conditions, particularly under low load and at high rpm. The coated-head case produces lower fuel consumption for a few combinations of operating conditions and higher than baseline for others. When all of the surfaces are coated (piston, head and liner), fuel consumption is always worse than baseline, probably due to the reduced volumetric efficiency.

What appear to be anomalous readings in several instances, e.g. the low fuel consumption for the coated piston case at 1000 RPM and a load of 335 KPa and the high consumption for the coated head case at 1000 RPM and 391 KPa, one might at first glance attribute to experimental scatter. However, a review of the entire set of BSFC data in Appendix A show these apparent anomalies to be remarkably consistent at the same engine speed throughout the injection timing range and for both fuels. The strong consistency of trends in the fuel consumption data adds confidence to the validity of the measurements. Uncertainty in the fuel consumption data is estimated to be ±10 gr/kwh indicated by the uncertainty bracket shown in Figure 3.2(a). Apparently, certain combinations of operating conditions (speed, timing and load) can work together to produce either unusually low or unusually high fuel consumption. The lowest overall fuel consumption for hexadecane (270 gr/kwh) and dodecane (260 gr/kwh) is produced by the baseline engine at 2000 rpm, 16 DBTDC and a load of 391 KPa.
Figure 3.2

BSFC Versus Load for Hexadecane Fuel at 20 Degrees BTDC
Figure 3.3

BSFC Versus Load for Dodecane Fuel at 20 Degrees BTDC
Figure 3.4

BSFC Versus Timing for Dodecane Fuel at High Load
3.4 EFFECT OF COATINGS ON NOx

Figures 3.5 and 3.6 compare the NOx emission in PPM for each of the four engine builds as a function of load for hexadecane and dodecane, respectively. Again, there are separate plots for each engine speed, and the example data shown is for a timing of 20 DBTDC. Complete data plots for all runs are provided in Appendix B. Uncertainty in the NOx data is estimated at ±50 PPM, approximately represented by the size of the data symbols in the figures.

Ceramic coating the piston crown has a very favorable effect on the NOx emission of this engine. The decrease in NOx for hexadecane at 20 DBTDC timing ranges from zero percent at high speed and low load to 25 percent at low speed and high load. The decrease under high load conditions ranges from 20 to 34 percent depending on injection timing, regardless of engine speed. Possible reasons for the decreased NOx production are related to the effect that the hotter piston surface has on the combustion process as discussed in the next chapter in conjunction with the KIVA modeling results. Lowered NOx emission coupled with improved fuel economy would appear to make coating of the piston crown a desirable engine modification.

Coating of the cylinder head by itself, on the other hand, does not significantly reduce NOx emission and under some operating conditions causes an increase. This is especially true with dodecane fuel, where coating the head causes increased NOx under most conditions. The unimproved or degraded NOx emission, coupled with similarly unimproved or degraded fuel consumption, provide little incentive to ceramic coat the cylinder head.

Operating the engine with coatings on all combustion chamber surfaces, i.e. piston, head and cylinder liner, produces NOx emission values that fall between those for the piston alone and for the head alone. One might expect this result based on the reasoning that the beneficial effects of coating the piston are offset somewhat by the often adverse effects of coating the head. In addition, the insulating effect of coating all three surfaces results in slightly higher cylinder gas temperatures throughout the combustion cycle which will tend to increase NOx as shown by the KIVA modeling results discussed in Chapter 4. Figure 3.7 shows the effect of varying injection timing on NOx emission at high load and four engine speeds for dodecane fuel. The trend is for increasing NOx with advancing injection timing as would be expected due to the increased ignition delay and premixed combustion fraction resulting in higher peak cylinder pressure and temperature.
Figure 3.5

NOx Versus Load for Hexadecane Fuel at 20 Degrees BTDC
Figure 3.6

NOx Versus Load for Dodecane Fuel at 20 Degrees BTDC

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Figure 3.7

NOx Versus Injection Timing for Dodecane Fuel at High Load
3.5 EFFECT OF COATINGS ON IGNITION DELAY PERIOD

Figures 3.8 and 3.9 show examples of ignition delay versus injection timing at light load for hexadecane and dodecane fuel respectively. Data for all cases is provided in Appendix C. Ignition delay was measured with the digital waveform analyzer used to collect the data by placing the cursor at the point on the needle lift trace where the needle opens and then advancing to the point on the pressure trace where rapid pressure rise begins. At the maximum 25 KHz data rate of the instrument, the time between data points is 0.04 ms, and this is assumed to represent the uncertainty in the ignition delay measurements. The uncertainty of ±0.04 ms is indicated by the brackets in the figures.

The ignition delay values for both fuels lie between 0.4 and 1.1 ms depending mainly on speed and injection timing, and to a lesser extent on load. These values are typical for high cetane fuels.

At low speeds the coated piston appears to decrease the ignition delay relative to the baseline engine as predicted by the KIVA results discussed in the next chapter. The shortened ignition delay is the reason proposed in Chapter 4, based on computer modeling, for reduced cylinder gas temperature and pressure and, consequently, reduced NOx relative to the baseline. In many cases the values of ignition delay for the coated-piston case and the baseline case are within the limits of uncertainty and are therefore inconclusive. In a few cases, the measured ignition delay for the coated-piston case is clearly longer than that for the baseline case, e.g. for hexadecane at 2500 RPM and 335 KPa BMEP (Figure C2(d), Appendix C). For this same case, the NOx emission of the coated-piston engine is slightly less than the baseline engine, in contradiction to our theory for decreased NOx. The conclusion that can be drawn from this is that ignition delay is not the only controlling parameter in the NOx reaction and additional combustion phenomena, of which we do not have a complete understanding at this time, are affected by the ceramic coatings. The lack of in-cylinder gas temperature measurements hinders our ability to completely understand the effect of coatings on the total combustion process.

For the cases of coated head alone and all surfaces coated, no unequivocal statements can be made concerning the effect on ignition delay. For hexadecane fuel, coating the head produces the longest ignition delay values under low load conditions, while coating all the surfaces produces the longest values under high load conditions. For dodecane fuel, coating all surfaces generally produces the longest ignition delays regardless of load.
Figure 3.8

Ignition Delay Versus Timing for Hexadecane Fuel at 223 KPa BMEP
Figure 3.9

Ignition Delay Versus Timing for Dodecane Fuel at 223 KPa BMEP
3.6 EFFECT OF COATINGS ON SOOT EMISSION

Soot concentration values in gr/m$^3$ were calculated from the photocell output of the USPHS smoke opacity meter by calibrating the meter output with a set of neutral density filters and using the concentration versus opacity data from Reference 11. The measurement uncertainty for the soot concentration values is unknown, but the data is useful for observing relative effects of the different coating schemes on soot production.

Figure 3.10 is representative of most of the soot data. It provides soot concentration versus load for hexadecane fuel at 22 DBTDC injection timing for four engine speeds. Most of the trends displayed in this figure hold true for the rest of the soot data which are presented in Appendix D. Coating all of the combustion chamber surfaces has a deleterious effect on soot emission, particularly at lower speeds and higher loads. The effect on soot emission of coating the piston alone depends on engine speed. Coating the piston results in higher soot production than baseline for some engine speeds and lower at others. Likewise, the effect of coating the head alone depends on engine speed. At 1500 RPM the coated head produces relatively low soot emissions, while at 2500 RPM it produces relatively high amounts of soot.

Possible explanations of these results based on computer simulations are not available at this time. The soot modeling described in the following chapter has so far been applied only to the baseline engine. While the results for the baseline case are encouraging, the soot model has not yet been implemented in conjunction with the coated surface temperature model in order to predict the effects of various surface coatings on soot emission. This work is continuing and will be published in technical paper form when completed.
Figure 3.10

Soot Concentration Versus Load for Hexadecane at 22 DBTDC
3.7 EFFECT OF COATINGS ON EXHAUST GAS TEMPERATURE

Typical exhaust gas temperature data are shown in Figure 3.11 which shows exhaust temperature versus load for hexadecane at 20 DBTDC injection timing. The same trends are exhibited by the rest of the exhaust temperature data provided in Appendix E. From these figures several interesting phenomena can be noted. Coating the head by itself has little effect on the exhaust temperature relative to baseline except at low speed where it actually reduces exhaust temperature. This trend is predicted by the KIVA results (e.g. see Figure 4.27). A study by Morel, et al. [12] indicated that for a non-insulated diesel engine, 49 percent of the heat loss from the hot cylinder gas occurs through the piston, 32 percent occurs through the head and the remaining 19 percent occurs through the liner. Thus, insulating the head should have less impact on the total heat transfer and the exhaust gas temperature. Peak cylinder gas temperatures are less than baseline due to the effect of the higher surface temperature on the combustion process, as discussed in Chapter 4.

The exhaust temperature for the coated-piston case is typically higher than baseline due to the reduced heat transfer through the piston to the oil. Since approximately half of the heat transfer takes place through the piston, adding a low conductivity coating to the piston crown should significantly reduce heat transfer and increase exhaust temperature.

Coating all of the surfaces produces the highest exhaust temperatures as would be expected due to the larger reduction in heat loss during the expansion stroke when all surfaces are insulated.
Figure 3.11

Exhaust Gas Temperature Versus Load for Hexadecane at 20 DBTDC
4. COMPUTER MODELING

4.1 BACKGROUND

The KIVA-II code, developed at Los Alamos National Laboratory [13], solves either the 2-D or the 3-D unsteady equations of motion of a chemically reactive mixture of ideal gases including the dynamics of a liquid fuel spray and the coupling between the spray and the gas. The code was developed with application to internal combustion engines specifically in mind, and contains a number of features to facilitate this application, such as gas flow, liquid fuel injection, spray dynamics, evaporation, heat transfer, combustion, species transport, and mixing.

KIVA-II was chosen for this work because of the documentation that is available [13,14,15] and the fact that it is being widely used by other researchers in industry, academia and national laboratories to model combustion and emissions in both spark ignition and diesel engines, e.g. Amsden et al. [16], O’Rourk and Amsden [17], Gentry et al. [18], Naber and Reitz [19], Kuo and Reitz [20], Reitz and Diwaker [21], Pinchon [22], Zellat [23], Gibson et al. [24], Varnavas and Assanis [25]. In addition, Ramos [26], and Markatos [27] documented KIVA in their recent books.

The original KIVA-II code employs the Arrhenius single reaction combustion model (SR model). This model is suitable for premixed combustion systems such as spark ignition engines. However, in diesel engines the combustion process is mainly diffusional after a short premixed period. The mixing-controlled eddy-break-up model (EBU model) developed by Magnussen et al. [28-31] provides a better physical representation of the diffusional combustion process.

In this study, KIVA-II was modified to include the EBU combustion model, and results from both models are compared with experiments. New model constants were introduced for the original Arrhenius combustion model. These constants were used for different combinations of load, injection timing and speed and were found to give good agreement between the predicted combustion and experiments for both hexadecane and dodecane fuel.

As part of this work, a soot model was also added to KIVA-II. The model for the rate of soot formation developed by Tesner et al. [32,33] was implemented to calculate the amount of soot formation, and the model of soot combustion developed by Magnussen
and Hjertager [28] concerning the behavior of soot in turbulent flames, was implemented to compute the oxidation rate of soot.

4.2 GENERAL STRUCTURE OF KIVA-II CODE

The gas-phase solution procedure in KIVA-II is based on a finite volume method called the ALE (Arbitrary Lagrangian-Eulerian) method which facilitates calculations with curved and changing boundaries. Spatial differences are formed on a finite-difference mesh that subdivides the computational region into a number of small cells that are hexahedrons. The code features a Stochastic Particle Technique, an efficient and accurate method for solving the spray dynamics, based on the Monte Carlo and discrete particle methods. The transient solution is marched out in a sequence of finite time increments called cycles or timesteps. On each cycle the values of the dependent variables are calculated.

The KIVA-II computer program consists of a set of subroutines controlled by a short main program. It was written specifically for use on the CRI Cray family of computers, operating under the Cray Time Sharing System (CTSS) and using the Cray FORTRAN (CFT and CFT77) compilers.

All computational runs were made on the CRAY Y-MP at the North Carolina Supercomputing Center (NCSC). KIVA-II contains statements peculiar to the CFT and CFT77 compilers that permit vectorization of many of the loops in the subroutines. The vectorization of the code improved the run time by a factor of five. A typical run for the Ricardo Hydra engine simulation, including the soot model, beginning at 138 degrees BTDC and ending at 90 degrees ATDC, required about 1200 CPU seconds on the Cray Y-MP.

4.3 COMPUTATIONAL PARAMETERS

For simplicity, a 2-D axisymmetric simulation of the combustion in the Ricardo Hydra direct injection diesel engine was used. Figure 4.1 shows the KIVA-II computational mesh of the Ricardo Hydra combustion chamber at 138° BTDC (point of intake valve closure), and again at TDC. The mesh has 23 cells in the radial direction, 1 cell in the azimuthal direction and 24 cells in the axial direction (nx=23, ny=1, nz=24).
Figure 4.1
Computational Mesh Used to Model Ricardo Hydra DI Engine
The calculations were started when the intake valve closed and ended at 90 degrees crank angle ATDC. The initial species was air which was assumed to be 21 percent oxygen and 79 percent nitrogen. The initial densities of the oxygen and nitrogen were assigned so that the initial pressure in the combustion chamber as the intake valve closed was 96 Kpa. The initial air temperature was assumed to be 310 K.

To calculate the dynamics of the spray, the model requires the radii and velocities of the droplets at the injector [15]. The magnitude of the injection velocity, \( V_{\text{inj}} \), is determined from

\[
V_{\text{inj}} = \frac{m_{\text{inj}}}{\rho_f \pi r_{\text{inj}}^3 t_{\text{inj}}^{2/3}}
\]

where:

- \( m_{\text{inj}} \) is the injected fuel mass, g/cycle.
- \( \rho_f \) is the fuel density, g/cm³
- \( r_{\text{inj}} \) is the injector nozzle radius, cm
- \( t_{\text{inj}} \) is the injection duration, s.

The sauter mean radius was assumed to be 9.5 micron. The model requires the injection angle, injected fuel mass, fuel density, and injection duration. These values were taken from the experimental data.
4.4 COMBUSTION MODELS

4.4.1 Single Reaction Model

The single reaction Arrhenius model is the standard combustion model built into the KIVA-II code. Oxidation of the fuel is assumed to be modeled by a one-step global kinetics scheme. The complete set of chemical reactions is given below.

Two different fuels were used in both the experiments and the modeling, viz hexadecane \((C_{16}H_{34})\) and dodecane \((C_{12}H_{36})\). The oxidation reaction for the two fuels are given by:

\[
2 \, C_{16}H_{34} + 49 \, O_2 \rightarrow 32 \, CO_2 + 34 \, H_2O
\]

and

\[
2 \, C_{12}H_{26} + 37 \, O_2 \rightarrow 24 \, CO_2 + 26 \, H_2O
\]

The reaction rate for these oxidation reactions is given by:

\[
\omega = k_{f_i} \, [\text{Fuel}]^a \, [\text{Oxidizer}]^b
\]

The coefficient, \(k_{f_i}\), is assumed to be of a generalized Arrhenius form,

\[
k_{f_i} = A_{f_i} \, T^{\zeta_{f_i}} \exp \left(-\frac{E_{f_i}}{T}\right)
\]

where:

- \(A_{f_i}\) is the Arrhenius coefficient
- \(\zeta_{f_i}\) is the forward temperature exponent
- \(E_{f_i}\) is the activation temperature, K.

The most common choice of parameters in the global rate expression is \(a = 1\) and \(b = 1\), assuming that the overall reaction is approximately first order with respect to both...
fuel and oxidizer. The other parameters, \( A_1 \) and \( E_1 \), can be adjusted so that the predicted Pressure-Crank Angle diagram matches the experimental value. It was found in this study that the best agreement with experiments was obtained with values of the fuel exponent \( a = 0.25 \) and the oxidizer exponent \( b = 1.5 \) for both hexadecane and dodecane. The best values of the Arrhenius coefficient and activation temperature that gave satisfactory results for both fuels were found to be \( A_1 = 9.0 \times 10^9 \) and \( E_1 = 1.5 \times 10^4 \) K respectively.

4.4.2 Eddy-Break-Up Model

In diffusion flames, fuel and oxygen occur in separate eddies. Because the chemical reactions in most cases are very fast, it can be assumed that the rate of combustion will be determined by the rate of intermixing of fuel and oxygen eddies on a molecular scale: in other words, by the rate of dissipation of the eddies. Because fuel and oxygen appear as fluctuating intermittent quantities, there will be a relationship between the fluctuations and the mean concentration of the species. Consequently, the rate of dissipation can be expressed by the mean concentration of the reacting species. The EBU model gives a physically based representation of combustion, because it takes into account the effect of turbulence on the mean chemical reaction rates. In this study, the model of Magnussen and Hjertager was used [28]. This model relates the rate of combustion to the rate of dissipation of turbulent eddies and expresses the rate of reaction by the mean concentration of a reacting species, the turbulent kinetic energy, and the rate of dissipation of this energy. This model differs from the EBU model of Spalding [34,35] in relating the dissipation of eddies to the mean concentration of intermittent quantities instead of the concentration fluctuations. This is advantageous in view of the uncertainty in determining the concentration fluctuations of the reacting species.

Accordingly, in the case where the adequate presence of fuel is the controlling factor, the rate of combustion of fuel can be expressed by:

\[
R_f = A \cdot C_f \cdot \left( \frac{\varepsilon}{k} \right) \tag{4.6}
\]

where:
- \( A \) is a constant
- \( C_f \) is the local time-mean fuel concentration
- \( k \) is the turbulent kinetic energy \( \text{cm}^2/\text{s}^2 \)
- \( \varepsilon \) is the rate of dissipation of turbulent kinetic energy \( \text{cm}^2/\text{s}^3 \)

32
In regions of the flame where the time mean concentration of fuel is high and the oxygen concentration is low, oxygen will be the reacting species that controls the rate of combustion. Accordingly, the rate of combustion in this case can be expressed by:

\[ R_f = A \cdot \frac{C_o}{r_f} \cdot \left( \frac{\varepsilon}{k} \right) \]  

(4.7)

where:

- \( C_o \) is the local time-mean oxygen concentration
- \( r_f \) is the stoichiometric oxygen requirement to burn 1 kg fuel.

In premixed turbulent flames, fuel and oxygen will occur in the same eddies. These eddies will be separated by eddies containing hot combustion products. The rate of combustion will, in this case, be determined by the same mechanism described above. However, an extra equation which takes care of the dissipation of the hot eddies must be added in cases where the concentration of hot combustion product is low. The combustion rate can then be expressed by:

\[ \kappa_f = A \cdot B \cdot \frac{C_p}{1 + r_f} \cdot \left( \frac{\varepsilon}{k} \right) \]  

(4.8)

where:

- \( B \) is a constant
- \( C_p \) is the local time-mean concentration of reaction products.

Thus, the EBU model is applicable to diffusion as well as to premixed flames. The equation that gives the lowest reaction rate is the one that determines the local rate of combustion. Therefore, combining equations (4.6-4.8), the rate of combustion can be expressed by:

\[ R_f = A \cdot \min \left( C_o, \frac{C_o}{r_f}, \frac{B \cdot C_p}{1 + r_f} \right) \cdot \left( \frac{\varepsilon}{k} \right) \]  

(4.9)

For the model constants, Magnussen suggested the values of \( A = 4.0 \) and \( B = 0.5 \), but Gosman [36] used \( A = 20 \) and \( B = 2.5 \), while Pinchon [22] used \( A = 16 \) and \( B = 2 \) and Varnavas [25] used \( A = 0.5 \) and \( B = 0.5 \). It was found in this study that the original values used by Magnussen (\( A = 4, B = 0.5 \)) gave the best overall agreement with the experimental data.
The EBU model has no provision for initiating combustion. To circumvent this problem, the SR chemical kinetic model was employed up to the point of auto-ignition after which the calculation procedure was switched to the EBU model. Switching from the chemical kinetics model to the EBU model depends on whether the ratio of the local turbulent time scale, $\tau_m$, to the local kinetic time scale, $\tau_r$, called the Damköhler number [23] is less than or greater than unity.

When $\tau_m / \tau_r < 1$, the process is controlled by chemistry since the chemical reaction time scale is much larger than the turbulent time scale.

When $\tau_m / \tau_r > 1$, the process is mixing controlled since the time required for chemical reactions is negligible compared to the time required for mixing.

The two time scales are defined by the relations:

$$\tau_m = \frac{k}{\varepsilon}$$

and

$$\tau_r = A \cdot \rho \exp(-E/R_o T)$$

where:

- $\tau_m$ is the turbulence mixing time scale
- $\tau_r$ is the chemical reaction time scale
- $A$ is the pre-exponential coefficient in the kinetics formula
- $\rho$ is the density
- $E$ is the activation energy
- $R_o$ is the universal gas constant
- $k$ is the turbulent kinetic energy
- $\varepsilon$ is the dissipation rate.

It was decided to switch from the SR to the EBU model of combustion after the ignition delay period when the combustion becomes controlled by mixing. The cylinder gas temperature was used to trigger the switch. A cell gas temperature of 1500 K was chosen to signal the switch to the EBU model. This cell temperature is generally exceeded within one time-step of the onset of combustion. By contrast, the maximum motoring temperature (no combustion) is about 950 K.
Although there is much evidence to support the dependence of the combustion rate on turbulent mixing, it is important to note that the mixing controlled EBU model is inadequate to predict combustion near walls, self-ignition, pollutant formation, and lean and rich flammability limits. These phenomena are controlled by chemical kinetics.

4.4.3 Results of Combustion Modeling

Comparisons of the P-θ diagrams obtained for the baseline engine with both the Eddy-Break-Up model and the single reaction model with experimental results have been made for several different loads and injection timings, and some representative results are presented here. As seen in Figures 4.2 and 4.3 for hexadecane fuel under two different loads for the same timing and engine speed, the cylinder pressure rises faster (rate of heat release is greater) during premixed combustion for the SR Model compared to the EBU Model. Beginning at the point of ignition, the rate of heat release shown by the EBU Model is representative of diffusion combustion. There appears to be no premixed pressure spike. The EBU Model predicts slightly lower peak pressure than the SR Model in all cases examined. A similar trend is noted for dodecane fuel as well, as shown in Figures 4.4 and 4.5.

It is also evident from the P-θ diagrams that the predicted ignition delay is always longer than the experimental value by approximately 3-5 degrees. Recall that the EBU Model has no mechanism for autoignition, and that when using the EBU Model we rely on the Arrhenius SR Model to initiate combustion. We thus see identical ignition delay periods for both models. Ignition delay in the SR model is primarily a function of the activation temperature constant, \( E_n \). The choice of activation temperature that gave the best overall results for the SR Model over a wide range of operating conditions, i.e. \( E_n = 1.5 \times 10^4 \), resulted in slightly longer ignition delay periods.

Figures 4.6 and 4.7 show the predicted cylinder pressure traces for three different injection timings for the EBU Model and SR Model, respectively. The maximum cylinder pressure increases with advancing injection timing as it should, due to a larger proportion of premixed combustion.

Figure 4.8 shows the predicted cylinder gas temperature versus crank angle for both combustion models. The EBU Model predicts slower gas temperature rise and a lower peak gas temperature than the SR Model, in keeping with the lower peak cylinder pressures.
Figure 4.2
Cylinder Pressure Versus Crank Angle For Hexadecane Fuel
1500 RPM, 18 Deg BTDC, 223 KPa BMEP

Figure 4.3
Cylinder Pressure Versus Crank Angle For Hexadecane Fuel
1500 RPM, 18 Deg BTDC, 391 KPa BMEP

Figure 4.4
Cylinder Pressure Versus Crank Angle For Dodecanne Fuel
1500 RPM, 16 Deg BTDC, 223 KPa BMEP

Figure 4.5
Cylinder Pressure Versus Crank Angle For Dodecanne Fuel
1500 RPM, 18 Deg BTDC, 223 KPa BMEP
Figure 4.6
Cylinder Pressure Versus Crank Angle For Hexadecane Fuel
EBU Model for Different Injection Timings

Figure 4.7
Cylinder Pressure Versus Crank Angle for Hexadecane Fuel
SR Model for Different Injection Timings

Figure 4.8
Cylinder Temperature Versus Crank Angle For Hexadecane
Single Reaction Model Versus Eddy-Break-Up Model
4.5 NITRIC OXIDE MODELING

4.5.1 NO Model

Nitric oxide (NO) is formed in diesel engines by the high temperature oxidation of atmospheric nitrogen. In oxygen-rich (lean) regions, the elementary process for the formation reaction of NO is expressed as:

\[
2N_2 + O_2 \xrightleftharpoons[k_f]{k_r} 2N + 2NO \tag{4.10}
\]

\[
2O_2 + N_2 \xrightleftharpoons[k_f]{k_r} 2O + 2NO \tag{4.11}
\]

In excess-fuel (rich) regions, the oxygen concentration required in equation (4.10) is low, and the reaction of N and OH (which is formed during the decomposition of the fuel) becomes important. This reaction is given by:

\[
N_2 + 2OH \xrightleftharpoons[k_f]{k_r} 2NO + 2H \tag{4.12}
\]

The elementary reaction processes given by equations (4.10) and (4.11) are known as the Zeldovich mechanism [37]. When the elementary reaction process represented by equation (4.12) is further added, the total reaction is called the extended Zeldovich mechanism [38].

Assuming a steady state approximation for the nitrogen atom concentration, the rate expression fits the extended zeldovich mechanism:

\[
\frac{d}{dt} [NO] = 2 k_f [O] [N_2] \left( \frac{1 - [NO]^2 / K [O_2] [N_2]}{1 + k_{b_i} [NO] / ( k_{f_i} [O_2] + k_{f_i} [OH])} \right) \tag{4.13}
\]

where:

\[
[ ] \quad \text{is the species concentration}
\]

\[ k_f \quad \text{is the forward reaction coefficient for the reaction } r \]

\[ k_{b_i} \quad \text{is the backward reaction coefficient for the reaction } r \]

\[ K \quad \text{is the dissociation equilibrium constant} = ( k_{f_i} / k_{b_i} ) ( k_{f_i} / k_{b_i} ) \]

The reaction rate constants used in KIVA-II are [39]:

38
\[ k_{f_1} = 1.5587 \times 10^{14} \exp(-67627/T) \]

\[ k_{b_1} = 7.5 \times 10^{12} \]

\[ k_{f_2} = 2.6484 \times 10^{14} \exp(-59418/T) \]

\[ k_{b_2} = 1.6 \times 10^{14} \exp(19678/T) \]

\[ k_{f_3} = 2.123 \times 10^{14} \exp(-57020/T) \]

\[ k_{b_3} = 0.0 \]

The required \textit{OH} concentrations were computed from the following set of equilibrium reactions:

\[ H_2 \rightleftharpoons 2H \]

\[ O_2 \rightleftharpoons 2O \]

\[ N_2 \rightleftharpoons 2N \]

\[ H_2 + O_2 \rightleftharpoons 2OH \]

\[ O_3 + 2H_2O \rightleftharpoons 4OH \]

\[ 2CO + O_2 \rightleftharpoons 2CO_2 \]

The following relation was used to convert the \textit{NO} values in gm/cycle, as calculated by KIVA, to concentration in particles per million (ppm) for comparison with experiments.

where:

\[ \dot{m}_{\text{total}} \] is the total mass flow rate (air + fuel injected)

\[ rps \] is the revolution per second.
4.5.2 Results of NO Modeling

Figures 4.9 and 4.10 compare the NO emission versus crank angle for both models with the experimental values for hexadecane fuel and two different injection timings. The EBU model predicts lower NO emissions than the SR model due to the lower peak cylinder temperature. The values predicted by the SR model are in better agreement with the experimental values, although they are still somewhat below the measured values. Figures 4.11 and 4.12 show similar NO emission results for dodecane fuel. For dodecane, the SR model gives very good agreement with the measured NO values.

Figures 4.13 - 4.15 show the effect of injection timing on NO emission for hexadecane for three different loads, comparing experimental results with both models. Both models underpredict NO, but produce the correct trend of increasing NO emission with advanced timing. Figure 4.16 shows similar data for dodecane at low load.
Figure 4.9
NO Emission Versus Crank Angle for Hexadecane Fuel
1500 RPM, 16 deg BTDC, 223 KPa BMEP

Figure 4.10
NO Emission Versus Crank Angle for Hexadecane Fuel
1500 RPM, 18 deg BTDC, 391 KPa BMEP

Figure 4.11
NO Emission Versus Crank Angle for Dodecane Fuel
1500 RPM, 16 deg BTDC, 223 KPa BMEP

Figure 4.12
NO Emission Versus Crank Angle for Dodecane Fuel
1500 RPM, 18 deg BTDC, 223 KPa BMEP
Figure 4.13
NO Emission Versus Injection Timing For Hexadecane Fuel
1500 RPM, 223 KPa BMEP

Figure 4.14
NO Emission Versus Injection Timing For Hexadecane Fuel
1500 RPM, 335 KPa BMEP

Figure 4.15
NO Emission Versus Injection Timing For Hexadecane Fuel
1500 RPM, 391 KPa BMEP

Figure 4.16
NO Emission Versus Injection Timing For Dodecane Fuel
1500 RPM, 223 KPa BMEP
4.6 SOOT MODELING

4.6.1 Review of Existing Soot Models

Most models appearing in the literature for the formation and oxidation of soot are empirical and much of the data used in the development of the models are based on laboratory burner flames rather than diesel engines. The literature on soot has been reviewed by several researchers, including Haynes and Wagner [40], and most recently by Morel and Keibar [41]. Although many useful observations have been made, no generally applicable soot emission model has yet been devised.

Khan, et al. [42,43] presented a soot formation correlation based on high pressure diesel engine data. The parameters used in this model are local values of temperature, unburned hydrocarbon concentration, and equivalence ratio. A corresponding soot combustion scheme was not suggested. The rate of soot formation was described as:

\[
\frac{ds}{dt} = \phi_f^3 P_f \exp(-2000/T) \tag{4.15}
\]

where:

- \( s \) is the soot mass
- \( \phi_f \) is the equivalence ratio in the formation zone
- \( P_f \) is the partial pressure of fuel in the formation zone
- \( T \) is the unburned temperature.

Nagle and Strickland [44] proposed the following relation for the rate of soot oxidation:

\[
\frac{ds}{dt} = -\frac{6}{\rho_s d_s} \frac{s}{P_o} f(T, P_o) \tag{4.16}
\]

where:

- \( P_o \) is the partial pressure of oxygen
- \( \rho_s \) is the density of the soot particle
- \( d_s \) is the diameter of the soot particle
- \( f(T, P_o) \) is a complex Arrhenius type expression.

Lee [45] proposed the following soot oxidation model:

\[
\frac{ds}{dt} = d_s^2 P_o T^{-1/2} \exp(-20,000/T) \tag{4.17}
\]
Fenimore and Jone's [46] suggested that the rate of soot oxidation is given by:

$$\frac{ds}{dt} = -P_s^{1/2} P_{H_2O}^{1/2} T^{-1/2} \exp(-19000 / T) \quad (4.18)$$

Hiroyasu and Kadota [47] proposed the following equations for soot formation and oxidation:

$$\frac{dm_f}{dt} = A_f m_f^{1/3} \exp(-E_f / RT)$$

$$\frac{dm_o}{dt} = A_c m_o^{1/3} \exp(-E_o / RT)$$

$$\frac{dm_s}{dt} = \frac{dm_f - dm_o}{dt} \quad (4.19)$$

where:

- $m_f$ is the mass of formed soot
- $m_o$ is the mass of oxidized soot
- $m_f$ is the mass of fuel vapor
- $m_s$ is the mass of soot.

$E_f$ and $E_o$ are assumed to be 1.25E+04 Kcal/Kmol and 1.4E+04 Kcal/Kmol respectively. $A_f$ and $A_c$ are constants which are determined so as to match the computed soot with the experimental results.

Morel [41] recommended empirical correlations for soot formation and oxidation directly linked to the combustion model. The soot model incorporates only those variables calculated by the combustion model, i.e., temperature, equivalence ratio, partial pressure of certain species, etc. According to Morel the amount of soot formation is described by

$$\frac{ds}{dt} = A_t \dot{m}_d \exp(-A_t / T_f) / (1 + 4.76 Y_{o_1})^3 \quad (4.19)$$

and the subsequent soot oxidation in the burned zone is described by

$$\frac{ds}{dt} = -B_s s / (P_s d_s) \exp(-B_s / T_f) \quad (4.20)$$

where:

- $s$ is the mass of the soot
- $Y_{o_1}$ is the mole fraction of available oxygen in the actively burning zone
- $\dot{m}_d$ is the rate of fuel burned in the diffusion burning model
- $A_t$ is a constant = 0.38
\( A_1 \) is a constant = 5000
\( T_f \) is the temperature in the formation zone
\( \rho_s \) is the soot density = 900 Kg/m\(^3\) (400 - 1500 Kg/m\(^3\))
\( d_s \) is the diameter of the elementary soot particle entering the burned zone = 0.012 - 0.032 \( \mu \)m
\( P_s \) is the partial pressure of \( O_2 \) in the burned zone
\( B_1 \) is a constant = 0.015
\( B_2 \) is a constant = 5000.

Tesner [32,33] suggested a kinetic scheme resulting in soot formation based on a chain-type process involving radical nuclei from which soot particles will later grow.

Magnussen and Hjertager [28] used Tesner's soot formation model to analyze a turbulent combustion flow, and added a model of soot oxidation which considers the behavior of soot in turbulent flames. Although there are a number of adjustable parameters in the Tesner model, Magnussen and Hjertager only had to modify Tesner's coefficients slightly to model their measurements of soot formation in turbulent acetylene/air flames. These models are described in more detail below.

4.6.2 Soot Model Adapted to KIVA

It was decided to use the soot formation model developed by Tesner, et al. [32,33] and the soot oxidation model developed by Magnussen and Hjertager [28] due to the fact that they take into consideration the behavior of soot in turbulent flames and, therefore, are suitable for use in conjunction with the eddy-break-up combustion model.

It has been assumed, following Tesner, that soot is formed from a gaseous fuel in two stages, where the first stage represents formation of radical nuclei, and the second stage represents soot particle formation from these nuclei. The rate of formation of radical nuclei is expressed by:

\[
R_{nf} = n_o + (f-g) \ n - g_o \cdot n \cdot N \quad \text{(part/m}^3\text{s)} \quad (4.21)
\]

where:
\( f \) is a linear branching coefficient
\( g \) is a linear termination coefficient
\( g_o \) is the coefficient of linear termination of soot particles
\( n \) is the concentration of radical nuclei (part/m\(^3\))
\( N \) is the concentration of soot particles (part/m\(^3\))
\( n_o \) is the rate of spontaneous formation of radical nuclei.

\( n_o \) is expressed by:

\[
n_o = a_o \cdot c_f \cdot e^{(-E/RT)} \quad \text{(part/m}^3\text{s)} \quad (4.22)
\]
where:

- \( a \) is a constant
- \( c_f \) is the mass concentration of fuel (Hx/m³)
- \( E \) is the activation temperature (K)
- \( T \) is the absolute temperature (K)
- \( R \) is the universal gas constant.

The rate of soot particle formation is expressed by:

\[
R_{sf} = m_p \left( a - b N \right) n \quad \text{(kg/m}^3\text{/s)}
\]

where:

- \( a \) & \( b \) are constants
- \( m_p \) is the mass of the soot particle (Kg/part). It is calculated assuming soot particles to be spherical and having an average diameter \( d_s = 200 \) Å and density \( \rho_s = 2000 \) Kg/m³.

In this model, the soot formation is limited to regions of unburnt fuel.

Soot oxidation occurs when soot particles come into contact with oxygen. The soot combustion model developed by Magnussen and Hjertager takes into account the behavior of soot in turbulent flames.

For regions where the local mean soot concentration is low compared to the oxygen concentration, the rate of soot combustion can be expressed as:

\[
R_{sc} = A \cdot C_s \left( \frac{\varepsilon}{k} \right) \quad \text{(kg/m}^3\text{/s)}
\]

where:

- \( A \) is a constant
- \( C_s \) is the local mean soot concentration (kg/m³)
- \( k \) is the turbulent kinetic energy cm²/s²
- \( \varepsilon \) is the rate of dissipation of turbulent kinetic energy cm²/s³.

In regions where the oxygen concentration is low, the oxygen will limit the rate of soot combustion. The soot must also compete for oxygen with the unburned fuel. In this case, the rate of soot combustion is given by:

\[
R_{sc} = A \left( \frac{C_s}{r_s} \right) \left( \frac{\varepsilon}{k} \right) \left( \frac{C_s r_s}{C_s r_s + C_f r_f} \right) \quad \text{(kg/m}^3\text{/s)}
\]

where:

- \( r_s \) is the stoichiometric oxygen requirement to burn 1 kg soot
- \( r_f \) is the stoichiometric oxygen requirement to burn 1 kg fuel
- \( C_{s_i} \) is the local time mean oxygen concentration (kg/m³)
- \( C_f \) is the local time mean fuel concentration (kg/m³).
The equation that gives the lower reaction rate determines the local rate of soot combustion. Thus, combining equations (4.24) and (4.25) leads to the expression:

$$R_{sc} = A \frac{\varepsilon}{k} \min \left[ C_i, \left( \frac{C_{h}}{r_{i}} + \frac{C_{d}}{r_{j}} \right) \right] \text{ (kg/m}^3\text{s)} \quad (4.26)$$

The local number of radical nuclei can be assumed to be reduced by combustion according to:

$$R_n = - \frac{n}{C} \text{ (part/m}^3\text{s)} \quad (4.27)$$

The concentration rate of soot particles is given by the difference of the soot formation rate and the oxidation rate, i.e.

$$\frac{dN}{dt} = R_{sf} - R_{so} \text{ (part/m}^3\text{s)} \quad (4.28)$$

Similarly, the concentration rate of radical nuclei is given by:

$$\frac{dn}{dt} = R_{af} - R_{ao} \text{ (part/m}^3\text{s)} \quad (4.29)$$

The simultaneous solution of the above two differential equations is obtained using a fourth order Runge-Kutta method.

4.6.3 Soot Model Parameters

Referring to studies on soot formation in atmospheric flames and shock tubes showed that the conditions under which soot formation occurs, as well as the quantity of soot formed, vary little for most hydrocarbons. Therefore, the constants appearing in the soot formation equation recommended by Tesner et al. for acetylene-oxygen flames have been used in this investigation along with the constants recommended by Magnussen [28] for the soot oxidation equations. The values of these constants are given below:

\[
\begin{align*}
a &= 1.6e5 \\
f-g &= 100 \\
g_o &= 1.0e-15 \\
b &= 8e-14 \\
a_o &= 7.254e30
\end{align*}
\]
The local mean soot concentration, $C_s$, (kg/m$^3$), is defined as:

$$C_s = \rho_s \text{ (kg/m}^3\text{)} \cdot \text{vol}_s \text{ (m}^3\text{/part)} \cdot N \text{ (part/m}^3\text{)}$$

where:

- $\rho_s$ is the density of the soot particle (kg/m$^3$)
- $\text{vol}_s$ is the volume of a soot particle (m$^3$/part)
- $N$ is the concentration of soot particles (part/m$^3$).

The mass of a soot particle, $m_p$ (kg/part), is defined as:

$$m_p = \rho_s \text{ (kg/m}^3\text{)} \cdot \text{vol}_s \text{ (m}^3\text{/part)}$$

$$= \frac{2000 \text{ (kg/m}^3\text{)} \cdot (4/3)\pi (d_s/2)^3 \text{ (m}^3\text{/part)}}{= 8.37758e-21 \text{ (kg/part)}}$$

Therefore, the local mean soot concentration, $C_s$, (kg/m$^3$), is given by:

$$C_s = 8.37758e-21 \text{ (kg/part)} \cdot N \text{ (part/m}^3\text{)}$$

The stoichiometric oxygen requirement to burn 1 kg of soot, $r_s$, is calculated from the following chemical reaction [48]:

$$\text{Soot} + O_2 \rightarrow CO_2$$

Therefore, $r_s$ is given by:

$$r_s = \frac{32}{12} = 2.66667$$

The soot model calculates the soot concentration in (kg/m$^3$). The calculation stops at 90 degree crank angle ATDC where the calculated average cylinder gas temperature is 990 K, and the average cylinder pressure is 4.98 atm. For comparison with the experimental data, this value is converted to exhaust gas conditions using the following relation, assuming ideal gas conditions.

$$C_s \text{ (at STP) = } C_s \text{ (at 90° ATDC) } \frac{T_{\text{cyl}}}{T_{\text{exh}}} \frac{P_{\text{exh}}}{P_{\text{cyl}}} \text{ (kg/m}^3\text{)}$$

Appendix F lists the soot model FORTRAN source code, and Appendix G shows the modifications made to the KIVA-II subroutine NEWCYC, and subroutine GLOBAL associated with the soot model.
4.6.4 Results of Soot Modeling

Figures 4.17 and 4.18 show the computed soot formation, soot oxidation, and the soot concentration, which is the difference between the formation and oxidation, for hexadecane and dodecane at 1500 RPM, 18 DBTDC and a load of 3.91 bar BMEP. Figures 4.19 and 4.20 show the computed soot concentration versus crank angle for hexadecane and dodecane at three different BMEP of 2.24, 3.35, and 3.91 bar at 1500 RPM and an injection timing of 18 DBTDC. The model produces the correct trend of increasing soot concentration with increasing load.

The rate of soot concentration versus crank angle can be divided into three regions. First, soot begins to form in the fuel rich mixture at the beginning of combustion and during the injection process where the fuel-air ratio at the inner zone of the spray is high. Then the compression stroke continues until the piston reaches TDC, and the expansion stroke begins. During this period, the air motion (swirl) is low and the rate of soot formation is higher than the rate of soot oxidation. This continues to about 20 degrees ATDC and the soot concentration rises sharply. At 20 degrees ATDC, the air motion continues increasing at a higher rate as the piston continues moving to about 40 degrees ATDC. This results in increasing the rate of soot oxidation. It is observed that the maximum soot concentration occurs at about 40 degrees crank angle ATDC. After 40 degrees crank angle ATDC, the air motion continues to increase at a higher rate and the small amount of remaining fuel continues to burn. On the other hand, the soot formed earlier comes in contact with oxygen, resulting in a high oxidation rate, and, as seen from the figures, the rate of soot concentration continues to decrease to about 60 degrees ATDC. Finally, the mixture reaches equilibrium and the amount of soot concentration remains constant until the end of the expansion stroke.

Figures 4.21 and 4.22 compare the experimental soot emission for the baseline engine with the computed results. The Figures show the effect of load on the soot emission at STP condition (T=289 K, P=1 atm) for hexadecane and dodecane respectively. The soot model produces the correct trend of increasing soot concentration with increasing load. As shown in the figures, the model gives fair agreement with the measured soot values.
**Figure 4.17**
Soot Formation and Oxidation Versus Crank Angle for Hexadecane

**Figure 4.18**
Soot Formation and Oxidation Versus Crank Angle for Dodecane
Figure 4.19
Soot Concentration Versus Crank Angle for Hexadecane

Figure 4.20
Soot Concentration Versus Crank Angle for Dodecane
Figure 4.21
Soot Emission Versus Load for Hexadecane

Figure 4.22
Soot Emission Versus Load for Dodecane
4.7 THERMAL MODELING OF CERAMIC COATINGS

Two of the input parameters required for KIVA-II are the surface temperature of the combustion chamber walls and the intake air temperature. For the baseline case, the temperature of all surfaces was assumed constant and uniform at 425 K, and the intake air temperature was assumed to be 315 K. The effect of applying a thin ceramic coating to combustion chamber surface is to increase both the average surface temperature and the variation in surface temperature relative to the baseline metallic surface. To model this effect, the coated surfaces were assumed to have a cyclic temperature profile, as depicted in Figure 4.23, consisting of a 100 K variation centered on an average value of 475 K. This temperature profile is based on results of experimental and analytical studies related to ceramic coated surface temperatures as cited in references 3, 6, and 7. The uncoated surfaces of the combustion chamber were kept at the same value as for the baseline model, i.e. 425 K, following a study by Smavik [8] who observed negligible temperature rise in the uninsulated components. The intake air temperature at the start of compression was also increased to account for both the mixing with higher exhaust gas temperature residuals caused by insulated surfaces and heat transfer from the higher temperature surfaces to the intake air. For the coated-piston case and the all-surface-coated case, 325 K was assumed for the intake air, while for the coated-head case, 320 K was assumed. These values were based on the experimental volumetric efficiency data.

4.8 THERMAL MODELING RESULTS

KIVA results are presented in Figures 4.24-4.27 for four different sets of operating conditions. All results shown are for hexadecane fuel at 1500 RPM. Figure 4.24 is for a load of 8 NM (223 KPa BMEP) and a timing of 22 DBTDC. Figure 4.25 is for a load of 12 NM (335 KPa BMEP) and a timing of 20 DBTDC. Figure 4.26 is for a load of 12 NM (335 KPa BMEP) and a timing of 22 DBTDC. Figure 4.27 is for a load of 14 NM (391 KPa BMEP) and a timing of 22 DBTDC. Each figure includes plots of cylinder pressure, cylinder gas temperature, mass of burnt fuel and NO concentration for the specified load and injection timing. Each plot compares the three ceramic coated cases with the baseline case. The effect of the various insulation schemes on pressure, temperature, rate of fuel combustion and NO emission is discussed in the following sections.
4.8.1 Cylinder Gas Pressure

The $P\theta$ diagrams in Figures 4.24(a) - 4.27(a) are simulations from KIVA for the baseline case and the three insulated engines. In each case the peak cylinder pressure is highest for the baseline engine, followed by the coated-head case, the all-surface-coated case and the coated-piston case. The coated-piston case always produces the lowest peak cylinder pressure. These KIVA results agree well with the experimental $P\theta$ diagrams shown in Figure 3.1. In general, the peak cylinder pressure is a function of the ratio $t_p/t_d$ (premixed combustion duration/diffusion combustion duration). The higher the ratio, the higher the peak cylinder pressure, and vice versa. This leads to the conclusion that selective insulation of the combustion chamber walls can modify the combustion characteristics of the engine, causing the peak cylinder pressure to vary accordingly. It also implies that the ignition delay for the coated-piston engine is the shortest and the ignition delay for the baseline engine is the longest of the four engine builds investigated. This is substantiated by the experimental data in some cases but not in others, as discussed in Section 3.5.

4.8.2 Cylinder Gas Temperature

KIVA-II calculates the gas temperature in each computational cell at each time step. Relations were added in this study to further calculate the mass average temperature of the entire cylinder contents based on an energy balance at each time step. Figures 4.24(b) - 4.27(b) compare the calculated average gas temperature versus crank angle for the baseline and three selective insulation schemes. The following general observations can be made from these figures: 1) During compression, the cylinder gas temperature is higher for the coated cases than for the baseline, 2) the gas temperature for the various ceramic coated cases begins to rise earlier but at a slower rate than the baseline case, 3) the maximum average cylinder gas temperature is lowest for the coated-piston case and highest for the baseline case, with the coated-head and totally insulated cases lying in between.

These observations may be attributed to the effect of the selective coating schemes on the ignition delay period. The initial higher gas temperature is due to the higher surface temperature of the ceramic surfaces. The earlier rising temperature is indicative of a shorter ignition delay which results in a smaller premixed combustion fraction and a longer diffusion combustion period resulting in a lower average peak cylinder temperature and a higher gas temperature during the expansion stroke. The results discussed above for peak cylinder pressure are also in support of this reasoning.
4.8.3 Mass of Burnt Fuel

Figures 4.24(c) - 4.27(c) depict the mass of fuel burnt during the combustion process. In general, it can be noted that the rapid rate of combustion associated with the premixed combustion phase occurs earlier and consumes less fuel for all of the coated cases relative to the baseline case, indicative of shortened ignition delay and decreased t_i/t_d. This is in keeping with the lower cylinder pressure and temperature predictions. Typically it is the coated-piston case that displays the lowest rate of combustion.

4.8.4 Nitric Oxide Emission

The rate of NO formation during combustion for the baseline and three selective insulation schemes is shown in Figures 4.24(d) - 4.27(d). It is noted that the coated-piston case typically produces the lowest NO emission consistent with the lowest rate of combustion and cylinder temperature. The baseline engine typically produces the highest NO emission except for the case shown in Figure 4.27(d) for 14 NM (391 KPa BMEP) load and 22 DBTDC. In this case the coated head produced a higher NO value than did the baseline. The experimental values of NO for these cases are indicated by the data points shown along the right side Y-axis. Reasonable agreement between experimental and predicted values is shown in most cases and the correct relative values of the coated cases versus baseline are obtained except for the low load condition shown in Figure 4.24(d) where agreement is not obtained for the insulated head case. Figures 4.28 and 4.29 further illustrate reasonable agreement between experimental values of NO and those predicted with KIVA, especially at higher values of load. At low loads KIVA underpredicted NO by about 25 percent.

Figure 4.23 Assumed Coated Surface Temperature Versus Crank Angle Profile
Figure 4.24

KIVA Results for 1500 RPM, 22 DBTDC and 223 kPa BMEP Load (8 NM Torque)
Figure 4.25
KIVA Results for 1500 RPM, 20 DBTDC and 335 KPa BMEP Load (12 NM Torque)
Figure 4.26(a)
Cylinder Gas Pressure Versus Crank Angle

Figure 4.26(b)
Cylinder Gas Temperature Versus Crank Angle

Figure 4.26(c)
Mass of Fuel Burnt Versus Crank Angle

Figure 4.26(d)
Nitric Oxide Concentration Versus Crank Angle

Figure 4.26

KIVA Results for 1500 RPM, 22 DBTDC and 335 KPa BMEP Load (12 NM Torque)
Figure 4.27(a)  
Cylinder Gas Pressure Versus Crank Angle

Figure 4.27(b)  
Cylinder Gas Temperature Versus Crank Angle

Figure 4.27(c)  
Mass of Fuel Burnt Versus Crank Angle

Figure 4.27(d)  
Nitric Oxide Concentration Versus Crank Angle

Figure 4.27  
KIVA Results for 1500 RPM, 22 DBTDC and 391 KPa BMEP Load (14 NM Torque)
Figure 4.28
Experimental Results For NO Emission Versus Load
For Hexadecane at 1500 RPM and 22 DBTDC Injection Timing

Figure 4.29
KIVA Results For NO Emission Versus Load
For Hexadecane at 1500 RPM and 22 DBTDC Injection Timing
REFERENCES


APPENDIX A

BSFC Data Plots
Figure A1

BSFC Versus Load for Hexadecane Fuel at 16 Degrees BTDC
Figure A2

BSFC Versus Load for Hexadecane Fuel at 18 Degrees BTDC
Figure A3

BSFC Versus Load for Hexadecane Fuel at 20 Degrees BTDC
Figure A4

BSFC Versus Load for Hexadecane Fuel at 22 Degrees BTDC

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Figure A5

BSFC Versus Load for Dodecane Fuel at 16 Degrees BTDC
Figure A6

BSFC Versus Load for Dodecane Fuel at 18 Degrees BTDC
Figure A7

BSFC Versus Load for Dodecane Fuel at 20 Degrees BTDC
Figure A8

BSFC Versus Load for Dodecane Fuel at 22 Degrees BTDC
APPENDIX B

NOx Data Plots
Figure B1

NOx Versus Load for Hexadecane Fuel at 16 Degrees BTDC
Figure B2
NOx Versus Load for Hexadecane Fuel at 18 Degrees BTDC
Figure B3

NOx Versus Load for Hexadecane Fuel at 20 Degrees BTDC
Figure B4

NOx Versus Load for Hexadecane Fuel at 22 Degrees BTDC
Figure B5

NOx Versus Load for Dodecane Fuel at 16 Degrees BTDC
Figure B6

NOx Versus Load for Dodecane Fuel at 18 Degrees BTDC

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Figure B7
NOx Versus Load for Dodecane Fuel at 20 Degrees BTDC

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Figure B8

NOx Versus Load for Dodecane Fuel at 22 Degrees BTDC
APPENDIX C

Ignition Delay Data Plots
Figure C1

Ignition Delay Versus Timing for Hexadecane Fuel at 223 KPa BMEP
Figure C2

Ignition Delay Versus Timing for Hexadecane Fuel at 335 KPa BMEP

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Figure C3

Ignition Delay Versus Timing for Hexadecane Fuel at 391 KPa BMEP

1000 RPM & 391 KPa BMEP

1500 RPM & 391 KPa BMEP

2000 RPM & 391 KPa BMEP

2500 RPM & 391 KPa BMEP

Fig C3(a)

Fig C3(b)

Fig C3(c)

Fig C3(d)
Figure C4

Ignition Delay Versus Timing for Hexadecane Fuel at 447 KPa BMEP
Figure C5

Ignition Delay Versus Timing for Dodecane Fuel at 223 KPa BMEP

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Figure C6

Ignition Delay Versus Timing for Dodecane Fuel at 335 KPa BMEP
Figure C7

Ignition Delay Versus Timing for Dodecane Fuel at 391 KPa BMEP
Figure C8

Ignition Delay Versus Timing for Dodecane Fuel at 447 KPa BMEP
APPENDIX D

Soot Concentration Data Plots
Figure D1

Soot Concentration Versus Load for Hexadecane at 16 DBTDC
Figure D2

Soot Concentration Versus Load for Hexadecane at 18 DBTDC
Figure D3

Soot Concentration Versus Load for Hexadecane at 20 DBTDC
Figure D4
Soot Concentration Versus Load for Hexadecane at 22 DBTDC
Figure D5

Soot Concentration Versus Load for Dodecane at 16 DBTDC
Figure D6

Soot Concentration Versus Load for Dodecane at 18 DBTDC
Figure D7

Soot Concentration Versus Load for Dodecane at 20 DBTDC
Figure D8

Soot Concentration Versus Load for Dodecane at 22 DBTDC
Figure E1
Exhaust Gas Temperature Versus Load for Hexadecane at 16 DBTDC
Figure E2

Exhaust Gas Temperature Versus Load for Hexadecane at 18 DBTDC
Figure E3

Exhaust Gas Temperature Versus Load for Hexadecane at 20 DBTDC
Figure E4

Exhaust Gas Temperature Versus Load for Hexadecane at 22 DBTDC
Figure E5

Exhaust Gas Temperature Versus Load for Dodecane at 16 DBTDC

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Figure E6

Exhaust Gas Temperature Versus Load for Dodecane at 18 DBTDC
Figure E7

Exhaust Gas Temperature Versus Load for Dodecane at 20 DBTDC
Figure E8

Exhaust Gas Temperature Versus Load for Dodecane at 22 DBTDC

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