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Adsorption of Pyridine at the Au(311) - Solution Interface

by

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ADSORPTION OF PYRIDINE AT THE AU(311)-SOLUTION INTERFACE

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Abstract

The energetics of pyridine adsorption onto a Au(311) single crystal electrode surface have been investigated using chronocoulometry. Adsorption parameters such as the relative Gibbs surface excesses, the Gibbs energies of adsorption, and electrosorption valencies have been determined. The data have been analyzed in terms of both the electrode potential and surface charge density. Over the potential region investigated (-0.8 V to +0.6 V), pyridine has been found to adsorb in the vertical orientation with the nitrogen end of the pyridine molecule facing the Au(311) surface. The limiting surface concentration for this orientation was found to be $6.6 \pm 0.3 \times 10^{-10}$ mol cm$^{-2}$. The Gibbs energies of adsorption were calculated using the Henry's Law isotherm. At the potential of maximum adsorption, $\Delta G^0_{\text{max}}$ was found to be equal to -44 kJ mol$^{-1}$. The absolute value for $\Delta G^0_{\text{max}}$ is fairly large indicating that pyridine is weakly chemisorbed onto the Au(311) electrode surface.
Introduction

This is the fifth paper in a series which is devoted to the study of the influence of crystallographic orientation of gold electrodes on the adsorption of pyridine from aqueous electrolyte solutions. In previous communications the adsorption of pyridine onto polycrystalline gold (1,2), Au(100) (3), Au(110) (4), and Au(111) (5) have been described. In this contribution the adsorption of pyridine onto a Au(311) single crystal electrode surface is reported. Chronocoulometry has allowed quantitative characterization of pyridine adsorption onto Au(311). Thus, adsorption isotherms, Gibbs energies of adsorption, and electrosorption valencies will be presented. The orientation of the pyridine molecules on the Au(311) electrode surface has been determined and will be discussed. Also, the adsorption of pyridine on the (311) and (110) faces of gold will be compared.

Experimental

The experimental procedures used in this work have been described in preceding communications (1-4). The working electrode was a Au(311) single crystal rod (99.99%, Johnson Matthey) with a cross sectional area of 0.0659 cm². Prior to each experiment the working electrode was flamed and then quenched with Milli-Q water.

All solutions were prepared from Milli-Q water (Waters). The supporting electrolyte was 0.1 M KClO₄. The KClO₄ (ACS Certified, Fisher) was purified by first calcinating, then recrystallizing twice from Milli-Q water and drying. Pyridine (ACS
Certified, Fisher) was used without further purification. Pyridine solutions ranging in concentration from $5 \times 10^{-6}$ M to $2 \times 10^{-3}$ M were investigated. Solutions were prepared from a 0.1 M pyridine stock solution. All solutions were deaerated with argon and during the experiment argon was passed over the top of the solution. Potentials were measured with respect to the external saturated calomel reference electrode (SCE). The counter electrode was a gold coil. The experiments were carried out at 25 $\pm$ 1 °C.

The instrumentation and the sequence of the different steps involved in the data acquisition and processing have been described in previous papers. The experimental strategy involved characterization of the surface by recording cyclic voltammetry and differential capacity curves and quantitative determination of the electrode charge density, $\sigma_M$, from chronocoulometric experiments (1,3,6,7). These charge density data allowed the determination of the relative Gibbs excesses, the Gibbs energies of adsorption, and the electrosorption valencies.

Cyclic Voltammetry

Figure 1 shows cyclic voltammograms which have been recorded for Au(311) in the absence (a) and presence (b) of pyridine. The cyclic voltammograms presented here span the potential region -0.8 V to +1.1 V and were recorded for a sweep rate of 20 mV s$^{-1}$. It was always observed that cyclic voltammograms which superimpose from one cycle to the next could be obtained after 3 to 5 cycles when recorded immediately after contact between the crystal and the electrolyte solution was made. No distortion of
the curves due to the creeping effect or the presence of oxygen could be detected. This suggests that the solutions being investigated were always of high purity and free of oxygen.

Figure 1a shows that the double layer region extends from -0.8 V to +0.6 V. Therefore, pyridine adsorption studies on Au(311) will be restricted to this potential region. For potentials more negative than -0.8 V hydrogen evolution on Au(311) begins while for potentials more positive than +0.6 V oxidation of the gold surface takes place. Two peaks at +0.8 V and +1.02 V characterize the oxide forming region. On the reverse, negative scan, one can see that the reduction process is characterized by two peaks at +0.31 V and +0.73 V. Overall, the cyclic voltammogram presented here for Au(311) in the presence of 0.1 M KClO₄ is in good agreement with the cyclic voltammogram reported by Hamelin (8) for Au(311) in the presence of 0.01 M NaF. However, the peak at +0.73 V was not observed by Hamelin. This peak is connected with a change of pH in the vicinity of the electrode. It appears if the electrolyte solution is not buffered and/or stirred. We did not stir our solution while Hamelin did. This explains the difference between the present results and those of Hamelin.

Shown in Fig. 1b is a cyclic voltammogram which was recorded for Au(311) in the presence of 3 x 10⁻⁴ M pyridine + 0.1 M KClO₄. In the double layer region pyridine adsorption/desorption peaks can be seen. These peaks first become visible once the pyridine concentration reaches 7 x 10⁻⁵ M. The oxide formation region, as Fig. 1b
shows, is now characterized by a single broad peak at +0.97 V. A shallow shoulder can be seen on the negative side of this peak. Changes in the appearance of the oxide formation region can be seen for pyridine concentrations as low as $7 \times 10^{-6}$ M.

**Differential Capacity**

Presented in Fig. 2 are several differential capacity curves which have been recorded for Au(311) in both the absence and presence of pyridine. In each case the sweep rate was 5 mV s$^{-1}$ and only the scan which was recorded in the direction of positive potentials is shown. On the reverse, negative scan, hysteresis was observed over the entire potential region investigated irrespective of whether pyridine was absent or present in the solution. The differential capacity curve recorded for Au(311) in the presence of 0.1 M KClO$_4$ displays a diffuse layer minimum at +0.020 ± 0.010 V which corresponds to the potential of zero charge (pzc); this value is in very good agreement with the value of +0.01 ± 0.01 V reported recently by Lecoeur et al. (9).

Figure 2 also shows differential capacity curves which were recorded for Au(311) in the presence of pyridine. The differential capacity curve recorded for the $10^{-4}$ M pyridine solution displays a broad peak at negative electrode potentials. When the pyridine concentration is increased to $10^{-3}$ M (or $2 \times 10^{-3}$ M) this peak becomes sharp and a shoulder appears on its negative side. Both the position and the height of the peak are sensitive to the bulk pyridine concentration suggesting that this peak is characteristic of the adsorption/desorption of pyridine. At positive electrode potentials,
and for the two highest pyridine concentrations presented in this figure, a broad shallow hump is visible. It will be shown later that within this potential region the surface concentration of pyridine has reached a maximum saturation value, $\Gamma_{\text{max}}$.

For the most negative potentials shown in Fig. 2, and for bulk pyridine concentrations $\leq 2 \times 10^{-3}$ M, the differential capacity curves merge with the capacity curve recorded for the pure supporting electrolyte solution. This behaviour indicates that, at the negative potential limit (-0.8 V), the pyridine molecules are desorbed from the Au(311) electrode surface. In the present work pyridine adsorption studies were restricted to pyridine concentrations in the range $5 \times 10^{-6}$ M to $2 \times 10^{-3}$ M. The lower limit is determined by the slow mass transport of the pyridine molecules to the electrode surface. The upper limit is determined by the requirement that the pyridine molecules are desorbed from the electrode surface before the onset of the hydrogen evolution reaction.

**Charge-Potential Plots**

In order to acquire quantitative data for the adsorption process potential step experiments were carried out. These involved the following steps: (1) the electrode was initially held at a potential (initial potential) at which pyridine adsorption takes place for a period of time long enough to establish a state of equilibrium between the interface and the bulk of the solution (approximately 1.5 min.); (2) the potential was stepped to the value $E = -0.8$ V (final potential) at which a total desorption takes place (see Fig. 2) and the current transient corresponding to the charging of the interface was recorded. The
charging current was digitally integrated and the absolute charge density, $\sigma_M$, at the initial potential was calculated using the procedure described in refs. 1, 3, and 7.

Shown in Fig. 3 are the absolute charge density-potential curves which have been obtained for Au(311) in both the absence and presence of pyridine. Pyridine concentrations ranging from $5 \times 10^{-6}$ M to $2 \times 10^{-3}$ M are shown in this figure. The reproducibility of the charge data was always checked and found to be within 1%. In the region which corresponds to the most negative electrode potentials studied, the $\sigma_M$ - E curves for the various pyridine solutions are observed to merge with the curve obtained for the supporting electrolyte. This indicates that, for $E < -0.7$ V, the pyridine molecules are completely desorbed from the Au(311) electrode surface. This result is consistent with the differential capacitance measurements presented earlier in Fig. 2. For more positive values of $E$ the absolute charge densities become strongly dependent upon the bulk pyridine concentration. Within this region the surface concentration of pyridine is an increasing function of the electrode potential. Following this region, the absolute charge densities, which are now positive, are essentially concentration independent indicating that the surface concentration of pyridine reached a maximum (saturation) value $\Gamma_{\text{max}}$. This region is characterized by a quasi-plateau. Initially, in this section of the curve, $\sigma_M$ depends almost linearly on the electrode potential. Extrapolation of this segment of the curve to $\sigma_M = 0$ allowed us to determine the shift of the potential of zero charge, $E_N$. The value for $E_N$ was found to be -0.73 ± 0.025 V. It can be shown that this value for
$E_N$ indicates that the pyridine molecules are orientated on the Au(311) electrode surface in the vertical position with the nitrogen atom facing the metal surface (3-5). The section of the $\sigma_M$ versus E curve, for which $\sigma_M$ is independent of the bulk pyridine concentration, intersects the curve which was obtained for the supporting electrolyte at a unique value of the electrode potential and charge density. These values, which correspond to the potential and charge of maximum adsorption are: $+0.39 \pm 0.01$ V and $+27.6 \pm 0.4$ $\mu$C cm$^{-2}$, respectively.

**Film and Surface Pressure Curves**

Following the procedure outlined in ref. 1, the charge density versus potential curves were integrated and the pressure of the film of adsorbed pyridine ($\pi$) calculated. Shown in Fig. 4 are the film pressure versus electrode potential curves which have been obtained for the various pyridine concentrations studied. All of the curves are bell shaped and display a well defined maximum.

Independently, the surface pressures, $\phi$, were calculated as a function of the charge density at the metal surface using the Parsons function $\xi$ (i.e., $\xi = \gamma + \sigma_M E$ and $\phi = \xi_{\phi = 0}$, $\xi_{\phi = 1}$) (10). In this way the thermodynamic analysis of the data was carried out also using charge as the independent electrical variable. Shown in Fig. 5 are plots of $\phi$ versus $\sigma_M$ for the various pyridine concentrations indicated. The surface pressure curves, like the film pressure curves presented above in Fig. 4, are bell shaped. Thus, the results in Figs. 4 and 5 are qualitatively similar. As expected, and for a given pyridine
concentration, the maximum value observed on the surface pressure curve is very similar to the maximum value observed on the corresponding film pressure curve.

The film and surface pressures reach fairly large values (> 100 mN m⁻¹) for moderate to high pyridine concentrations. This indicates that the zero coverage Gibbs energy of adsorption and/or the energy of attractive lateral interactions are also large.

**Adsorption Isotherms**

By differentiating the film pressure versus natural logarithm of the bulk pyridine concentration curves, the relative Gibbs excesses were calculated. Shown in Fig. 6a are plots of the relative Gibbs excess as a function of the applied potential for the various pyridine concentrations studied. Shown in Fig. 6b are the adsorption isotherms (\( \Gamma \) versus \( \ln(c) \)) which have been obtained at constant potential. Figures 7a and 7b show the corresponding curves determined using charge density as the independent electrical variable. The \( \Gamma \) versus \( E \) or \( \sigma_M \) plots have a sigmoidal shape with one inflection point indicating that the adsorbed pyridine molecules assume only one orientation at the Au(311) electrode surface. The curves display a long well defined plateau which corresponds to a limiting surface concentration equal to \( 6.6 \pm 0.3 \times 10^{-10} \) mol cm⁻². This value is in good agreement with the packing density expected for a surface covered by a monolayer of pyridine molecules orientated in the vertical position.

The adsorption isotherms determined at constant potential (Fig. 6b) are steep. In contrast, the isotherms plotted at constant \( \sigma_M \) (Fig. 7b) have very small slopes.
Consequently, the shape of the adsorption isotherm is essentially determined by the choice of the electrical variable. Any attempt to discuss these curves in terms of molecular models leading to an equation of an adsorption isotherm would give contradictory conclusions. Therefore, no isotherm has been used to describe the present results.

**Gibbs Energy of Adsorption**

In the potential range -0.5 V to 0 V the Gibbs energies of adsorption, $\Delta G^0$, were calculated from the initial slopes of the $\pi$ versus mole fraction of pyridine plots using the expression for Henry’s Law as explained in ref. 3. A value of $\Gamma_{\text{max}}$ equal to $6.6 \times 10^{-10} \text{ mol cm}^{-2}$ was used in the calculations. For potentials more positive than 0 V, the surface concentration of pyridine is equal to $\Gamma_{\text{max}}$ even for the lowest bulk pyridine concentrations investigated. Under these conditions the Henry’s Law coefficients could not be determined from the slopes of the $\pi$ versus mole fraction of pyridine plots. To evaluate the Gibbs energies of adsorption at more positive potentials the $\pi$ values were plotted against $RT\ln(c)$ at constant $E$ for potentials in the region 0 V to +0.5 V. The plots were linear and parallel to each other. These plots were then shifted by a distance $RT\Delta\ln(c)$ along the $RT\ln(c)$ axis to give the best overlap with the curve for $E = 0$ V. The magnitude of $RT\Delta\ln(c)$ was then taken as equal to the difference between the Gibbs energies at $E = 0$ V and at a given potential $E$. 
Shown in Fig. 8 is a plot of $\Delta G^0$ versus E. The standard state is unit mole fraction of pyridine in the bulk of the solution and unit coverage at the surface (unsymmetrical choice of standard state (12)). The data presented in Fig. 8 span the potential region -0.7 V to +0.5 V. As Fig. 8 shows, the $\Delta G^0$ versus E relationship is parabolic in nature. The maximum value observed for $\Delta G^0$ is -44 kJ mol$^{-1}$. This fairly large value for $\Delta G^0_{\text{max}}$ indicates that the pyridine molecules are weakly chemisorbed onto the Au(311) electrode surface.

The electrosorption valency, $\gamma'$, is equal to the first derivative of the $\Delta G^0$ versus E plot. Independently, the electrosorption valency can also be determined from the dependence of the charge density on the surface concentration taken at constant electrode potential by using the following relationship (13):

$$
\gamma' = \frac{1}{F} \left( \frac{\partial \mu^0}{\partial E} \right)_T - \frac{1}{F} \left( \frac{\partial \Delta G^0}{\partial E} \right)_T - \frac{1}{F} \left( \frac{\partial \sigma_m}{\partial \Gamma} \right)_E
$$

Therefore, the slopes of the $\sigma_m$ versus $\Gamma$ plots can be compared to the first derivative of the $\Delta G^0$ versus E plot to check the consistency of our results.

Presented in Fig. 9 are plots of $\sigma_m$ versus $\Gamma$ which have been obtained at constant electrode potential. The data presented in this figure span the potential region -0.55 V to -0.05 V. As this figure shows, quasilinear relationships are observed although the spread of the experimental points is fairly large. The electrosorption valencies determined from the slopes of these plots are shown in Fig. 10. Independently, the $\Delta G^0$
versus E data of Fig. 8 were fitted by a polynomial of the 2nd order and numerically differentiated. The electrosorption valencies determined by this method are also plotted in Fig. 10. The two sets of data agree with each other fairly well. This shows that our data display an internal consistency and that no major errors were made in the data processing.

Shown in Fig. 11 is a plot of $\Delta G^0$ versus $\sigma_M$. For charge densities in the region -20 to 0 $\mu$C cm$^{-2}$ the Gibbs energies of adsorption were calculated from the initial slopes of the surface pressure versus mole fraction of pyridine plots obtained at constant charge density as outlined in ref. 14. The Gibbs energies for $\sigma_M > 0$ were obtained by plotting the $\phi$ values versus RT ln(c) and subsequently shifting the curves corresponding to a given value of $\sigma_M$ along the RT ln(c) axis to give the best overlap with the curve determined for $\sigma_{\text{max}}$ (+26 $\mu$C cm$^{-2}$). The Gibbs energy at $\sigma_{\text{max}}$ was taken to be equal to the value of $\Delta G^0$ at $E_{\text{max}}$ as shown in Fig. 8. The data presented in Fig. 11 show that the Gibbs energies of adsorption depend on $\sigma_M$ in a pseudo parabolic way.

The consistency of our results can once again be checked by taking into consideration the following relationship obtained by cross differentiating the electrocapillary equation:
According to the above expression, the first derivative of the Gibbs energy of adsorption with respect to the charge density (taken at constant \( \Gamma \)), should be equal to the slope of the electrode potential versus the surface excess plot (taken at constant \( \sigma_M \)). Figure 12 shows a plot of \( \Delta_2^M \phi \) versus \( \Gamma \) taken at constant \( \sigma_M \). \( \Delta_2^M \phi \) is the potential drop across the inner layer region of the double layer which is given by the following equation:

\[
\Delta_2^M \phi - E - E_{pzc} - \phi_2
\]  

where \( E_{pzc} \) is the potential of zero charge and \( \phi_2 \) is the potential drop across the diffuse layer which is obtained from the classical Gouy-Chapman theory. At constant \( \sigma_M \) the derivative of \( \Delta_2^M \phi \) with respect to \( \Gamma \) is equal to the derivative of \( E \) with respect to \( \Gamma \). The plots presented in Fig. 12 are nonlinear. However, the limiting slopes of these curves taken at \( \Gamma = 0 \) progressively decrease with negative charge at the surface. The initial slopes of these curves are compared to the result of the numerical differentiation of the \( \Delta G^0 \) versus \( \sigma_M \) curve in Fig. 13. Although the scatter of the experimental points is large, satisfactory agreement between the two sets of data is observed indicating that no major errors were made in the data processing.

**Summary and Conclusions**

Data characterizing the adsorption of pyridine onto a Au(311) single crystal electrode surface have been presented. The data have been analyzed in terms of both
potential and charge. The important adsorption parameters are summarized in Table 1. The results show that pyridine adsorbs onto the Au(311) electrode surface in the vertical orientation over the entire potential (charge) region investigated. From the shift of the potential of zero charge one can conclude that the nitrogen end of the pyridine molecule is facing the metal surface. The fairly large values observed for the Gibbs energies of adsorption indicate that the pyridine molecules are weakly chemisorbed onto the Au(311) electrode surface. The consistency of our results has also been checked using expressions for the partial derivatives determined by cross differentiation of the electrocapillary equation. It can be concluded that our data are free from major errors which may have arisen as a result of data processing.

To conclude this paper we would like to stress briefly the similarities between pyridine adsorption at the Au(311) and Au(110) surfaces. The structure of the two surfaces consists of monatomic steps separated by two atom wide terraces. Using the TLK model, they can be denoted as (110) = 2(111) x (111) and (311) = 2(111) x (100) = 2(100) x (111). Consequently, the two planes can be considered as having the same type of steps but different terraces, or alternatively, different steps but the same terraces. The structural differences between the two planes are therefore small.

In an earlier contribution we reported on the adsorption of pyridine onto an Au(110) surface (4). The adsorption parameters obtained from that study are presented in Table 1 where they are compared to the data obtained here for the Au(311) plane.
Pyridine molecules assume only one - vertical orientation with the nitrogen end facing the metal on both the negatively and positively charged surfaces. The magnitude of all the adsorption parameters listed in Table 1 are comparable for the two planes. Clearly, pyridine adsorption onto the Au(311) and Au(110) surfaces is very similar and virtually unaffected by the small structural differences between these single crystal faces. Incidentally, similarities between pyridine adsorption at the (110) and (311) surfaces of silver have been observed recently in this laboratory (15). A more detailed discussion of the effect of surface crystallography on pyridine adsorption at gold electrodes will be presented in our next contribution in which results from adsorption studies on the (100), (110), (111), and the (311) planes of gold will be compared.

Acknowledgements

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References


Figures

Fig. 1. Cyclic voltammograms recorded for (a) 0.1 M KClO₄ and (b) 0.1 M KClO₄ + 3 x 10⁻⁴ M pyridine. Sweep rate = 20 mV s⁻¹.

Fig. 2. Differential capacity curves recorded for Au(311) in the presence of the supporting electrolyte (0.1 M KClO₄) as well as the various pyridine concentrations indicated. Sweep rate = 5 mV s⁻¹ and the ac modulation frequency is 25 Hz.

Fig. 3. Charge density-potential curves which have been obtained for the following pyridine concentrations: (●) 2 x 10⁻³ M; (○) 10⁻³ M; (×) 3 x 10⁻⁴ M; (□) 10⁻⁴ M; (x) 7 x 10⁻⁵ M; (☆) 3 x 10⁻⁵ M; (△) 2 x 10⁻⁵ M; (◇) 10⁻⁵ M; (□) 5 x 10⁻⁶ M; (●) 0.

Fig. 4. Film pressure versus electrode potential curves for the following pyridine concentrations: (1) 2 x 10⁻³ M; (2) 10⁻³ M; (3) 3 x 10⁻⁴ M; (4) 10⁻⁴ M; (5) 7 x 10⁻⁵ M; (6) 3 x 10⁻⁵ M; (7) 2 x 10⁻⁵ M; (8) 10⁻⁵ M; (9) 5 x 10⁻⁶ M.

Fig. 5. Surface pressure versus absolute charge density curves obtained for the following pyridine concentrations: (O) 2 x 10⁻³ M; (●) 10⁻³ M; (△) 3 x 10⁻⁴ M; (☆) 10⁻⁴ M; (□) 7 x 10⁻⁵ M; (■) 3 x 10⁻⁵ M; (△) 2 x 10⁻⁵ M; (◇) 10⁻⁵ M; (○) 5 x 10⁻⁶ M.

Fig. 6. (a) Surface concentration-potential curves obtained for the following pyridine concentrations: (1) 2 x 10⁻³ M; (2) 10⁻³ M; (3) 3 x 10⁻⁴ M;
(4) $10^{-4}$ M; (5) $7 \times 10^{-5}$ M; (6) $3 \times 10^{-5}$ M; (7) $2 \times 10^{-5}$ M; (8) $10^{-5}$ M and
(b) the corresponding adsorption isotherms which have been obtained at
the following electrode potentials: (1) -0.55 V; (2) -0.50 V; (3) -0.45 V;
(4) -0.40 V; (5) -0.35 V; (6) -0.30 V; (7) -0.25 V; (8) -0.20 V; (9) -0.15 V;
(10) -0.10 V; (11) -0.05 V.

Fig. 7.
(a) Surface concentration-charge density curves for the following pyridine
concentrations: (1) $10^{-3}$ M; (2) $3 \times 10^{-4}$ M; (3) $10^{-4}$ M; (4) $7 \times 10^{-5}$ M; (5)
$3 \times 10^{-5}$ M; (6) $2 \times 10^{-5}$ M; (7) $10^{-5}$ M and (b) the corresponding
adsorption isotherms which have been obtained for the various charge
densities in $\mu$C cm$^{-2}$ indicated.

Fig. 8.
Plot of $\Delta G^0$ versus E as determined from the Henry’s Law isotherm.

Fig. 9.
Plots of $\sigma_M$ versus $\Gamma$ obtained at the following electrode potentials:

(●) -0.55 V; (O) -0.45 V; (■) -0.35 V; (□) -0.25 V; (★) -0.20 V; (Δ) -0.15 V;
(Õ) -0.10 V; (◊) -0.05 V.

Fig. 10.
Dependence of the electrosorption valencies on the electrode potential. (□)
determined from the slope of the $\sigma_M$ versus $\Gamma$ plots and (O) by numerical
differentiation of the $\Delta G^0$ versus E curve corresponding to Henry’s Law.

Fig. 11.
Plot of $\Delta G^0$ versus $\sigma_M$ as determined from the Henry’s law isotherm.

Fig. 12.
Plots of $\Delta_2^M \phi$ versus $\Gamma$ for the following charge densities indicated: (O) -20
$\mu$C cm$^{-2}$; (●) -16 $\mu$C cm$^{-2}$; (□) -12 $\mu$C cm$^{-2}$; (■) -8 $\mu$C cm$^{-2}$;
Fig. 13. Comparison of the dependence of $(\Delta_2 \alpha_M^\sigma)$ on the charge density determined from the initial slopes of the plots presented in Fig. 12. (O) and by numerical differentiation of the $\Delta G^0$ versus $\alpha_M$ curve presented in Fig. 11. (■).
Table 1. Adsorption parameters for pyridine adsorbed onto the (311) and (110) faces of gold.

<table>
<thead>
<tr>
<th>Parameters</th>
<th>Au(311)</th>
<th>Au(110)</th>
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<tbody>
<tr>
<td>$E_{\text{max}} / \text{V}$</td>
<td>+0.39</td>
<td>+0.23</td>
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<tr>
<td>$E_{\text{pzc}} / \text{V}$</td>
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<td>$10^{10} \Gamma_{\text{max}} / \text{mol cm}^{-2}$</td>
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<td>6.1</td>
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<td>$\Delta G^0_{\text{max}} / \text{kJ mol}^{-1}$</td>
<td>-44</td>
<td>-42</td>
</tr>
</tbody>
</table>
Fig. 3

$\sigma_M / \mu C \ cm^{-2}$

$E / V \ (SCE)$
Fig. 5

$\phi \text{ mN m}^{-1}$

$\sigma_M \text{ } \mu \text{C cm}^{-2}$
Fig. 9

$\sigma_m / \mu C \text{ cm}^{-2}$ vs. $10^{10} \Gamma / \text{ mol cm}^{-2}$
\[ 10^{-8} \left( \frac{d\Delta_2^M}{d\Gamma} \right) \sigma_M / V \text{ mol}^{-1} \text{ cm}^2 \]
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