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Volume 1
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Mechanical Technology Incorporated

in fulfillment of

Contract No. NOBS-86914
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Volume 1

Spring and Damping Coefficients For The Tilting-Pad Journal Bearing

by

Jorga Lund

Mechanical Technology Incorporated
Preface

This report is the first of three volumes which are the product of an analytical and experimental investigation into the effect of a hydraulically supported, tilting-pad, journal bearing on the attenuation of noise originating from rotor unbalance.

The study is a part of a larger program to reduce structure-borne noise originating in high-speed rotating equipment onboard Navy vessels. The introduction of additional flexibility at the journal bearings was proposed to reduce noise transmission, and in the case under study the flexibility is provided by hydraulic pistons and accumulators. In order to evaluate and optimize the additional flexibility, this is a comprehensive investigation of the dynamics of a rotor-bearing system and how the system behavior is affected by the dynamic properties of both the hydrodynamic oil film and the flexible supports.

The second volume of this study will be concerned with the force transmission of (1) two-mass rotor bearing system and (2) uniform rotor bearing system. The third volume will include (1) general rotor analysis and (2) experimental correlation.

Mechanical Technology Incorporated was primarily responsible for the analytical portion of the investigation while Westinghouse Electric Corp. designed and conducted the experimental test.
# TABLE OF CONTENTS

<table>
<thead>
<tr>
<th>Section</th>
<th>Page</th>
</tr>
</thead>
<tbody>
<tr>
<td>ABSTRACT</td>
<td>1</td>
</tr>
<tr>
<td>INTRODUCTION</td>
<td>2</td>
</tr>
<tr>
<td>RESULTS</td>
<td>3</td>
</tr>
<tr>
<td>DISCUSSION OF RESULTS</td>
<td>7</td>
</tr>
<tr>
<td>NUMERICAL EXAMPLE</td>
<td>8</td>
</tr>
<tr>
<td>DISCUSSION OF ANALYSIS</td>
<td>9</td>
</tr>
<tr>
<td>ANALYSIS</td>
<td>10</td>
</tr>
<tr>
<td>CONCLUSIONS</td>
<td>19</td>
</tr>
<tr>
<td>RECOMMENDATIONS</td>
<td>20</td>
</tr>
<tr>
<td>REFERENCES</td>
<td>21</td>
</tr>
<tr>
<td>APPENDIX A: Computer Program PN0078:</td>
<td>22</td>
</tr>
<tr>
<td>Spring and Damping Coefficients for the Tilt-</td>
<td></td>
</tr>
<tr>
<td>ing Pad Journal Bearing.</td>
<td></td>
</tr>
<tr>
<td>APPENDIX B: Computer Program PN0131:</td>
<td>37</td>
</tr>
<tr>
<td>Interpolation of Partial Arc Bearing Force</td>
<td></td>
</tr>
<tr>
<td>Derivatives for Use in Tilting Pad Bearing</td>
<td></td>
</tr>
<tr>
<td>Calculations.</td>
<td></td>
</tr>
<tr>
<td>APPENDIX C: Computer Program PN0091:</td>
<td>50</td>
</tr>
<tr>
<td>Numerical Calculation of the Incompressible Full</td>
<td></td>
</tr>
<tr>
<td>or Partial Journal Bearing.</td>
<td></td>
</tr>
<tr>
<td>TABLE 1</td>
<td>74</td>
</tr>
<tr>
<td>FIGURES</td>
<td>76</td>
</tr>
<tr>
<td>NOMENCLATURE</td>
<td>92</td>
</tr>
</tbody>
</table>
ABSTRACT

A method for calculating spring and damping coefficients for the tilting pad journal bearing is presented. The analysis includes the effect of pad inertia. Numerical results are given in form of design curves for the centrally pivoted 4 shoe, 5 shoe, 6 shoe and 12 shoe journal bearing. A comparison with test results is included.
INTRODUCTION

This report is part of a study of the force transmitted from a rotating shaft to its bearing pedestals where the force is generated by the presence of an unbalance in the rotor. The shaft is supported in fluid film bearings and since the fluid film possesses both flexibility and damping the resulting rotor motion, and therefore the transmitted force, are directly influenced by the properties of the bearing film. Thus it is necessary to determine the bearing stiffness and damping, expressed in terms of spring and damping coefficients, in order to calculate the transmitted force. The bearing spring and damping coefficients depend on the operating conditions and the bearing configuration and have been determined for several bearing types in a preceding investigation, see Ref. 2. In connection with the present continuation of the investigation it has become necessary to examine yet another bearing type, namely the tilting pad bearing.

References 1 and 2 show that although proper bearing design procedures can effect a reduction in the transmitted force, bearing instability imposes an upper limit on the attainable force attenuation. Hence, to further improve the attenuation other methods must be employed. Two possibilities seem open: a) the use of an inherently stable bearing type such as the tilting pad bearing, or b) by vibration isolation of the bearing housing. In order to achieve a substantial force attenuation the latter possibility must be chosen. However, this suggests a very soft bearing support which in turn implies a potential risk for bearing instability. Hence, from both points of view the tilting pad journal bearing offers advantages. For these reasons an analysis of this bearing has been undertaken and it is the purpose of the present report to give the results of the investigation in form of design curves for the spring and damping coefficients. A second report will deal with the force attenuation due to the flexible bearing support.
RESULTS

Numerical calculations have been performed for four bearing geometries: the 4-shoe, the 5-shoe, the 6-shoe, and the 12-shoe tilting pad journal bearing. All four configurations have centrally pivoted pads. The results are given in form of graphs, Figs. 1 to 19, where the spring and damping coefficients are given as functions of the bearing Sommerfeld number. The following factors are investigated:

a) bearing length-to-diameter ratio
b) the direction of the static load
c) the inertia of the shoes
d) the preload of the bearing (for vertical rotors)

The results are given in dimensionless form:

- dimensionless spring coefficient: $\frac{CK_{xx}}{W}$, $\frac{CK_{yy}}{W}$
- dimensionless damping coefficient: $\frac{C_{xx}C}{W}$, $\frac{C_{yy}C}{W}$

Sommerfeld Number: $S = \frac{\mu ND L}{W} \left( \frac{R}{C} \right)^2$

where:

- $C$ - Pad radius of curvature minus journal radius, inch
- $W$ - Bearing reaction, lbs.
- $D$ - Journal diameter, inch
- $R$ - Journal radius, inch
- $L$ - Bearing length, inch
- $\mu$ - Lubricant viscosity, lbs.sec/in.$^2$
- $N$ - Rotational speed, RPS
- $\omega$ - $2\pi N$, angular speed, rad/sec.
- $K_{xx}, K_{yy}$ - Spring coefficients, lbs/in.
- $C_{xx}, C_{yy}$ - Damping coefficients, lbs.sec/in.
Figures 1 to 12 are intended as design curves for horizontal rotors. They cover the three most commonly used bearing geometries for three values of the L/D-ratio. The first 9 graphs have the static load direction between the bottom pads whereas for the next 3 graphs the load vector passes through the pivot of the bottom pad. Since the bearing geometries are symmetric about the vertical load line and shoe inertia is neglected the cross-coupling spring and damp coefficients vanish. Furthermore, for the 4-shoe bearing there is also symmetry around a horizontal line such that coefficients in the horizontal and the vertical direction are identical, i.e. \( K_{xx} = K_{yy} \) and \( C_{xx} = C_{yy} \). This does not hold for the 5-shoe and the 6-shoe bearing.

On each graph is shown a curve labeled "Critical Mass". It gives the value of the pad inertia necessary to incur resonance of the pad motion, resonance being defined by a 90 degree phase angle between journal and pad motion (or in other words, at resonance the phase angle between a zero-inertia pad and a pad with a "critical mass" is 90 degrees). Hence, for a particular bearing design the value of the dimensionless "pad mass" can be calculated:

\[
\text{dimensionless mass} = \frac{\text{CWM}}{(\mu DL CR^2)^2}
\]

where:

\[
\begin{align*}
M &= \frac{L}{R_p^2} \quad \text{equivalent pad mass,} \quad \text{lbs. sec}^2 \text{ in.} \\
I &= \text{Transverse mass moment of inertia of pad,} \quad \text{lbs. in. sec}^2 \\
R_p &= \text{Pad radius, inch.} \\
W &= \text{Bearing reaction,} \quad \text{lbs.} \\
D &= \text{Journal diameter, inch} \\
R &= \text{Journal radius, inch} \\
L &= \text{Bearing length, inch} \\
C &= \text{Pad clearance, inch} \\
\mu &= \text{Lubricant viscosity,} \quad \text{lbs. sec/in.}^2
\end{align*}
\]

Entering the appropriate graph with this value the corresponding "critical" Sommerfeld number, and, therefore, the resonant speed, can be determined.
Figures 13 to 16 illustrate the effect of pad inertia on the spring and damping coefficients for a particular bearing geometry, namely the 4-shoe bearing, L/D = .75, with the load passing between the two bottom pads. The calculations are performed for six values of the dimensionless pad mass. Since the pads are given inertia the cross-coupling terms no longer vanish as for Figs. 1 to 12.

Figures 17 to 19 apply to a vertical rotor with zero bearing eccentricity ratio. In order to have stiffness the bearing must be preloaded. The preload is defined by:

\[ \text{Preload} = (1 - C'/C) \]

where:

\[ C = \text{difference between radius of curvature of pad and journal radius, inch.} \]
\[ C' = \text{difference between radius of pivot point circle and journal radius, inch.} \]

Hence, \( C'/C = 1 \) corresponds to no preload whereas \( C'/C = 0 \) means that the journal touches the pad.

Since the bearing reaction \( W \) is zero for a vertical rotor the coefficients must be made dimensionless in a form different from the horizontal rotor results. The following form is chosen:

\[ \text{dimensionless spring coefficient} = \frac{K_{xx}}{\mu NL (\frac{R}{C})^3} \]
\[ \text{dimensionless damping coefficient} = \frac{\omega C_{xx}}{\mu NL (\frac{R}{C})^3} \]
\[ \text{dimensionless critical mass} = \frac{M_{\text{crit}} N}{\mu L (\frac{R}{C})^3} \]

where the symbols are defined above. Note, that for a vertical rotor \( K_{xx} = K_{yy} \) and \( C_{xx} = C_{yy} \).
All the given results are based on partial arc fluid film force derivative computed by finite difference calculations on a computer using program PN0091, Appendix C. Fluid film rupture is included. Each partial arc (in total 10) is calculated for 8 eccentricity ratios: $e = 0.01, 0.1, 0.2, 0.35, 0.5, 0.65, 0.8, 0.95$. A summary of the derivatives is given in Table 1.
DISCUSSION OF RESULTS

A comparison between calculated and measured results is shown in Figs. 20 and 21. The test results are taken from Ref. 4. Even if Ref. 4 gives experimental data for values of the Sommerfeld number ranging from 0 to 16, Ref. 5 indicates that the actual measurements are limited to a range from .15 to 1.2 and that outside this range extrapolation has been used. Based on these considerations the agreement seems reasonably good within the normal range of bearing operation. It should also be noted that the test measurements are obtained with vibration amplitudes of up to 50 percent of the minimum film thickness whereas the theoretical calculations assume very small amplitudes.

The most unusual aspect of the theoretical results is the sudden loss of bearing stiffness when approaching pad resonance. However, the effect may be less drastic than shown. Since at resonance the "critical mass" is zero the effect of pad inertia cannot be neglected. Then Figs. 13 to 16 show that the cross-coupling terms become important. Hence, the overall coefficient may not be zero (see Eq. (10), Ref. 1). Pad resonance occurs in general at high Sommerfeld numbers where the eccentricity ratio is small and thus, the out-of-phase between pad and journal motion may not be too serious. In addition the fluid film damping is large.

Finally, it should be noted that the fluid film force calculations include film rupture. When the bearing is flooded with oil the film can sustain pressures below the ambient and rupture does not take place. Instead the film may cavitate. Therefore, if the present results are applied to the calculation of a flooded bearing there may be minor errors at low Sommerfeld numbers.
NUMERICAL EXAMPLE

Let a bearing have a diameter $D = 6$ inch and length $L = 3$ inch, i.e., $L/D = .5$. The radial pad clearance is $C = .0045$ inch and the pad center coincides with the bearing center so that there is no preload. The speed is 3600 RPM, i.e., $N = 60$ RPS, and the corresponding oil viscosity is 14 centipoise, i.e., $\mu = 2.0 \cdot 10^{-6}$ lbs.sec/in$^2$. With a bearing reaction $W = 4,000$ lbs. the Sommerfeld number becomes $S = .244$. When the load is between the pads and pad inertia is neglected the dynamic coefficients for the 4-shoe bearing are obtained from Fig. 7.

\[
\frac{C}{W} = 1.15 \quad K_{xx} = 1.02 \cdot 10^6 \text{ lbs/inch}
\]

\[
\frac{C}{W} = 5.4 \quad C_{xx} = 12,700 \text{ lbs.sec/inch}
\]

The dimensionless critical mass is $5.25 \cdot 10^{-3}$ which gives a critical pad mass moment of inertia $I = .69$ lbs. in.sec$^2$. This corresponds approximately to a shoe weight of 180 lbs. which is much larger than the actual weight. Hence, there is no danger of resonance. For a 6-shoe bearing the results are obtained from Fig. 2 as:

\[
K_{xx} = 4.3 \cdot 10^6 \text{ lbs/inch} \quad K_{yy} = 1.4 \cdot 10^6 \text{ lbs/inch}
\]

\[
C_{xx} = 10,400 \text{ lbs.sec/inch} \quad C_{yy} = 3,400 \text{ lbs.sec/inch}
\]

critical pad moment of inertia = $18.5$ lbs.in.sec$^2$

corresponding to a shoe weight of 12,500 lbs.
DISCUSSION OF ANALYSIS

The theoretical equation governing the fluid film behavior is the well-known Reynolds equation. Due to its complexity it is normally solved numerically on a computer and the resulting force is a non-linear function of the eccentricity ratio and the journal center velocity. For an exact solution Reynolds equation should be solved simultaneously with the equation of motion for the rotor, leading to quite involved calculations. For practical purposes it is convenient instead to assume that the rotor amplitude is sufficiently small to allow replacing the fluid film forces by their gradients around the steady state operating eccentricity. Thus, the forces become proportional to the vibratory amplitude and velocity, the coefficients for proportionality being denoted spring and damping coefficients. Introducing an x-y-coordinate system the fluid film forces can be written as:

\[ F_x = -K_{xx} \dot{x} - C_{xx} \dot{x} - K_{xy} \dot{y} - C_{xy} \dot{y} \]
\[ F_y = -K_{yx} \dot{x} - C_{yx} \dot{x} - K_{yy} \dot{y} - C_{yy} \dot{y} \]

such that there are 4 spring coefficients and 4 damping coefficients.

For a fixed pad the spring and damping coefficients can be computed from the gradients of the fluid film force as shown in the Analysis Section or in Refs. 1, 2, 3 and 6. The gradients are obtained by numerical differentiation of computer results.

The analysis of the tilting pad bearing assumes the spring and damping coefficients for the fixed pad to be known. Thus, for an arbitrary rotor motion it becomes possible to establish the equation of motion for the pad itself, including the inertia of the pad. Coupling this equation with the above equations for the fluid film force yields the spring and damping coefficients for the tilting pad. A summation over all pads results in the combined spring and damping coefficients for the complete tilting pad journal bearing.
ANALYSIS

Steady State Equilibrium

Referring to Fig. 22, let the center of the tilting pad bearing be \( O_B \). The bearing is made up of a number of tilting pads, the arbitrary pad having the pivot point \( P \) located in an angle \( \psi \) from the vertical load line. The pad is free to tilt around the pivot point which for convenience is assumed to be located on the surface of the pad. The steady state position of the journal center is \( O_J \) which is the origin of two fixed coordinate systems: the \( x-y \)-system (\( x \)-axis vertical downward, \( y \)-axis horizontal) and the \( \xi-\eta \)-system (the \( \xi \)-axis passing through the pivot point). The location of \( O_J \) with respect to the bearing center is given by the eccentricity \( O_B O_J = e = C e_o \) and the attitude angle \( \varphi \). The steady state position of the pad center is designated \( O_n \) such that the journal center eccentricity with respect to the pad is determined by \( O_n O_J = e = C e_o \). The corresponding attitude angle is \( \varphi \). The journal radius is \( R \), the radius of the pad is \( O_n P = R + C \) and the radius of the circle passing through all pivot points with center in \( O_B \) is \( O_B P = R + C' \).

The point \( O_n \) is the pad center with no tilting. Hence, \( O_n O_n \) is a circular arc, or for small motions, a line perpendicular to \( O_n P \). Projecting \( O_J \) on \( O_n P \) yields:

\[
(1) \quad e \cos \varphi = \frac{C'}{C} - e_o \cos (\psi - \varphi_o)
\]

This equation contains three unknowns, namely \( e, \varphi \) and \( \varphi_o \) since \( e_o \) is the independent variable. The second equation derives from the requirement that the force on the pad passes through the pivot point which establishes a relationship between \( e \) and \( \varphi \) (i.e. the journal center locus with respect to the pad). The third relationship is the requirement that the total horizontal force component (in the \( y \)-direction), summed over all pads, is zero:

\[
(2) \quad \sum_{\text{all pads}} F \sin \psi = 0
\]
Eq. (2) is usually used to determine \( \varphi_0 \) by trial-and-error as follows: for a particular case \( \xi_0, \xi' \) and \( \psi \) are known in Eq.(1). For several assumed values of \( \varphi_0 \) calculate \( \xi \cos \varphi \) for each pad, determine the pad forces \( F \) from available pad data and plot \( \Sigma F \sin \psi \) as a function of \( \varphi_0 \). The zero-point determines the desired value of \( \varphi_0 \).

This procedure is at best tedious. It becomes very complicated if the pivot point is not in the center of the pad in which case the pad force \( F \) can be a multivalued function of \( \xi \cos \varphi \). However, when the pivot points are located symmetrically with respect to the vertical load line through the bearing center \( O_0 \), then \( \varphi_0 = 0 \). A further simplification arises when the pivot point is in the center of the pad since the pad force \( F \) then is uniquely determined by \( \xi \cos \varphi \). These two conditions apply to all the numerical calculations in the present report but is not a necessary assumption for the analysis to be valid. The analysis requires only that \( \xi \) and \( \varphi \) are known for each pad. How these values have been arrived at is immaterial in so far as the analysis is concerned.

**Fixed Pad Coefficients**

In order to determine the spring and damping coefficients for the complete tilting pad bearing it is necessary to know the forces and their derivatives for each pad as if the pad was fixed. Under steady state conditions the journal center has the eccentricity ratio \( \xi \) and the attitude angle \( \varphi \) with respect to the pad center. The fluid film pad force has the components \( F_\xi \) and \( F_\gamma \) and under steady state conditions the resultant force passes through the pivot point, i.e. \( F_\xi = -F \) and \( F_\gamma = 0 \), see Fig. 22. Thus \( F \) denotes the load on the pad. Usually the force is resolved along the radial and the tangential directions with the components \( F_\xi \) and \( F_\gamma \), respectively, see Fig. 22. Hence:

\[
\begin{bmatrix}
  F_\xi \\
  F_\gamma 
\end{bmatrix}
= \begin{bmatrix}
  -F \\
  0 
\end{bmatrix}
= -\begin{bmatrix}
  \cos \varphi & \sin \varphi \\
  \sin \varphi & -\cos \varphi 
\end{bmatrix}
\begin{bmatrix}
  F_\xi \\
  F_\gamma 
\end{bmatrix}
\]
For an infinitesimal small motion around the steady state position the dynamic forces become from Eq. (3):

\[
\begin{align*}
\begin{bmatrix} dF_x \\ dF_y \end{bmatrix} &= - \begin{bmatrix} \cos \varphi & \sin \varphi \\ \sin \varphi & -\cos \varphi \end{bmatrix} \begin{bmatrix} dF_x + F_t \, d\varphi \\ dF_y - F_r \, d\varphi \end{bmatrix}
\end{align*}
\]

The infinitesimal dynamic motion of the journal center is described by the coordinates \((\xi, \eta)\):

\[
\xi = d(e \cos \varphi) \quad \eta = d(e \sin \varphi)
\]

or

\[
\begin{align*}
\begin{bmatrix} de \\ ed\varphi \end{bmatrix} &= \begin{bmatrix} \cos \varphi & \sin \varphi \\ -\sin \varphi & \cos \varphi \end{bmatrix} \begin{bmatrix} \xi \\ \eta \end{bmatrix}
\end{align*}
\]

The velocities transform similarly:

\[
\begin{align*}
\begin{bmatrix} de \\ ed\varphi \end{bmatrix} &= \begin{bmatrix} \cos \varphi & \sin \varphi \\ -\sin \varphi & \cos \varphi \end{bmatrix} \begin{bmatrix} \dot{\xi} \\ \dot{\eta} \end{bmatrix}
\end{align*}
\]

The dynamic force components \(dF_x\) and \(dF_y\) may be expressed in terms of the dynamic amplitudes. From Reynolds equation it can be shown that the fluid film force \(F\) can be written (Refs. 1,2):

\[
F = \lambda \omega (1 - 2 \, \dot{\varphi}) \cdot f(\varepsilon, \psi, (\dot{\xi})/(1 - 2 \, \dot{\varphi}))
\]

where:

\[
\lambda = \frac{\mu R L}{\pi} \left(\frac{R}{C}\right)^2
\]

\[
f = \frac{1}{S_p}
\]

\[
S_p = \frac{M N D L}{F} \left(\frac{R}{C}\right)^2 \quad \text{(Pad Sommersfeld Number)}
\]

Therefore:

\[
dF = \lambda \omega \left\{ (1 - 2 \, \dot{\varphi}) \left[ \frac{\partial F}{\partial \varepsilon} \, d\varepsilon + \frac{\partial F}{\partial \psi} \, d\psi + \frac{\partial F}{\partial (\dot{\xi}/(1 - 2 \, \dot{\varphi}))} \, d\left(\dot{\xi}/(1 - 2 \, \dot{\varphi})\right) \right] - \frac{2 f}{\omega e} \, ed\dot{\varphi} \right\}
\]
Now:
\[
d\left(\frac{\dot{\xi}}{1-2\frac{\dot{\xi}}{\dot{\phi}}}\right) = \frac{1}{1-2\frac{\dot{\xi}}{\dot{\phi}}} d\left(\dot{\xi}\right) + \frac{2}{\omega^2} d\left(\frac{\dot{\phi}}{\omega}\right) = d\left(\dot{\phi}\right)
\]
because at the equilibrium position \(\dot{\xi} = \dot{\phi} = 0\). Hence, Eq.(11) reduces to:

(12) \[dF = \frac{1}{\epsilon} \lambda \omega \left\{ \frac{\partial \xi}{\partial \phi} \frac{d\phi}{d\xi} \frac{de}{d\phi} + \frac{1}{\omega} \frac{d\xi}{d\phi} \frac{d\phi}{d\xi} \right\} = \frac{d\phi}{d\xi} \frac{de}{d\phi} - \frac{1}{\omega} \frac{d\phi}{d\xi} \frac{d\phi}{d\xi} \frac{d\phi}{d\xi} \frac{de}{d\phi}
\]

Eq.(12) applies to both \(F_r\) and \(F_t\). Let the corresponding dimensionless forces be denoted \(f_r\) and \(f_t\) as defined by Eq. (7). Thus, by substitution of Eq. (12) into Eq. (4):

(13) \[\begin{bmatrix} df_r \\ df_t \end{bmatrix} = -\frac{1}{\epsilon} \lambda \omega \begin{bmatrix} \cos \phi & \sin \phi \\ -\sin \phi & -\cos \phi \end{bmatrix} \begin{bmatrix} \frac{d\xi}{d\phi} & \frac{d\phi}{d\xi} \\ \frac{d\phi}{d\phi} & \frac{d\phi}{d\phi} \end{bmatrix} \begin{bmatrix} de \\ d\phi \end{bmatrix} + \frac{1}{\omega} \begin{bmatrix} \frac{d\phi}{d\phi} & \frac{d\phi}{d\phi} \\ \frac{d\phi}{d\phi} & \frac{d\phi}{d\phi} \end{bmatrix} \begin{bmatrix} de \\ d\phi \end{bmatrix}
\]

Define the fixed pads spring and damping coefficients by:

(14) \[dF_r = -K_{f_r} \ddot{\xi} - C_{f_r} \dot{\xi} - K_{f_t} \dot{\xi} - C_{f_t} \dot{\xi}
\]

To determine the 8 coefficients, substitute Eq. (5) and (6) into Eq.(13) and collect the terms in accordance with Eq.(14) to get:

(15) \[K_{f_r} = \frac{1}{\epsilon} \lambda \omega \left[ \frac{d\xi}{d\phi} \cos \phi - \frac{d\phi}{d\phi} \sin \phi \right] - \left(\frac{d\xi}{d\phi} - \frac{d\phi}{d\phi} \cos \phi \sin \phi - \frac{f_{t}}{\epsilon} \sin \phi \right]
\]

(16) \[\omega C_{f_r} = \frac{1}{\epsilon} \lambda \omega \left[ \frac{d\xi}{d\phi} \cos \phi + \frac{d\phi}{d\phi} \cos \phi \sin \phi \right] - \frac{2f_t}{\epsilon} \sin \phi
\]

(17) \[K_{f_t} = \frac{1}{\epsilon} \lambda \omega \left[ \frac{d\phi}{d\phi} \cos \phi + \frac{d\phi}{d\phi} \sin \phi \right] + \left(\frac{d\phi}{d\phi} + \frac{d\phi}{d\phi} \cos \phi \sin \phi + \frac{f_{t}}{\epsilon} \cos \phi \right]
\]

(18) \[\omega C_{f_t} = \frac{1}{\epsilon} \lambda \omega \left[ \frac{d\phi}{d\phi} \sin \phi + \frac{d\phi}{d\phi} \cos \phi \sin \phi \right] + \frac{2f_t}{\epsilon} \cos \phi
\]

(19) \[K_{f_r} = \frac{1}{\epsilon} \lambda \omega \left[ -\frac{d\phi}{d\phi} \cos \phi - \frac{d\phi}{d\phi} \sin \phi \right] + \left(\frac{d\phi}{d\phi} + \frac{d\phi}{d\phi} \cos \phi \sin \phi + \frac{f_{t}}{\epsilon} \sin \phi \right]
\]

(20) \[\omega C_{f_r} = \frac{1}{\epsilon} \lambda \omega \left[ -\frac{d\phi}{d\phi} \cos \phi + \frac{d\phi}{d\phi} \cos \phi \sin \phi - \frac{2f_t}{\epsilon} \sin \phi \right]
\]
where:

(23) \( f_s = \frac{1}{2} C_p \)

(24) \( f_\gamma = 0 \)

(25) \( \frac{1}{2} \lambda \omega = \frac{1}{2} F S_p \)

and all forces and derivatives are calculated for the given steady state position, defined by \( \epsilon \).

**Tilting Pad Coefficients**

Referring to Fig. 22, \( \eta_n \) is the steady state position of the pad center. Under dynamic load the pad center oscillates around \( \eta_n \) with the amplitude \( \eta_P \) such that \( \eta_P / R_P \) represents the dynamic tilting angle of the pad (\( R_p \) is the distance from the actual pad pivot to the pad center). The moment on the pad from the fluid film pressure is \( -R_P dF_\gamma \). Therefore, if the mass moment of inertia of the pad is \( I \) the equation of motion becomes:

\[
I \frac{\dddot{\eta}_P}{R_P} = -R_P dF_\gamma
\]

or

(26) \[
\frac{I}{R_P^2} \dddot{\eta}_P = M \dddot{\eta}_P = -dF_\gamma
\]

where:

(27) \[
M = \frac{I}{R_P^2}
\]
Thus, under dynamic conditions $\gamma$ should be replaced by $(\gamma - \gamma_p)$ in Eq. (14). In order to eliminate $\gamma$, substitute the expression for $dF_\gamma$ into Eq. (26):

$$M \ddot{\eta}_p = K_{\eta_p} \ddot{\xi} + C_{\eta_p} \dot{\xi} + K_{\gamma \eta_p} (\gamma - \gamma_p) + C_{\gamma \eta_p} (\gamma - \gamma_p)$$

Let the dynamic motion of the journal center around the steady state position be harmonic:

$$(\xi, \eta) e^{i\omega t}$$

Hence, solve Eq. (28):

$$\gamma - \gamma_p = -\frac{(K_{\eta_p} + i\omega C_{\eta_p}) \xi + M\omega^2 \eta}{K_{\eta_p} - M\omega^2 + i\omega C_{\eta_p}} = -\left[(K_{\eta_p} + i\omega C_{\eta_p}) \xi + M\omega^2 \eta\right] (p - iq)$$

where:

$$p = \frac{K_{\eta_p} - M\omega^2}{(K_{\eta_p} - M\omega^2)^2 + (\omega C_{\eta_p})^2}$$

$$q = \frac{\omega C_{\eta_p}}{(K_{\eta_p} - M\omega^2)^2 + (\omega C_{\eta_p})^2}$$

Thus, replacing $\gamma$ by $(\gamma - \gamma_p)$ Eq. (14) becomes:

$$dF_\xi = -(K_{\xi} + i\omega C_{\xi}) \dot{\xi} - (K_{\xi} + i\omega C_{\xi}) \eta$$

$$dF_\gamma = -(K_{\gamma} + i\omega C_{\gamma}) \dot{\xi} - (K_{\gamma} + i\omega C_{\gamma}) \eta$$

where $K'_{\xi \xi}, C'_{\xi \xi}$ etc., are the spring and damping coefficients for the tilting pad and given by:

$$K'_{\xi \xi} = K_{\xi \xi} - (pK_{\xi \xi} + q\omega C_{\xi \xi})K_{\xi \xi} - (qK_{\xi \xi} - p\omega C_{\xi \xi})C_{\xi \xi}$$

$$\omega C'_{\xi \xi} = \omega C_{\xi \xi} - (pK_{\xi \xi} + q\omega C_{\xi \xi})C_{\xi \xi} + (qK_{\xi \xi} - p\omega C_{\xi \xi})K_{\xi \xi}$$

$$K'_{\xi \gamma} = -M\omega^2 (pK_{\xi \gamma} + q\omega C_{\xi \gamma})$$
If the pad has no inertia, i.e. $M=0$, only $k'$ and $\omega C'$ remain as would be expected.

**Bearing Spring and Damping Coefficients**

Having determined the spring and damping coefficients for the individual tilting pads it remains to combine them into the overall bearing coefficients. The coordinate system for the bearing is the x-y-system, see Fig. 22, with the coordinate transformation:

$$\begin{align*}
\begin{bmatrix}
\xi \\
\eta
\end{bmatrix} &= -
\begin{bmatrix}
\cos \psi & \sin \psi \\
-\sin \psi & \cos \psi
\end{bmatrix}
\begin{bmatrix}
x \\
y
\end{bmatrix} \\
\begin{bmatrix}
dF_x \\
dF_y
\end{bmatrix} &= -
\begin{bmatrix}
\cos \psi & -\sin \psi \\
\sin \psi & \cos \psi
\end{bmatrix}
\begin{bmatrix}
dF_\xi \\
dF_\eta
\end{bmatrix}
\end{align*}$$

(42) \hspace{1cm} (43)

where $dF_x$ and $dF_y$ are the dynamic forces measured in the x-y-system. The bearing spring and damping coefficients are defined by:

$$
\begin{align*}
dF_x &= -K_{xx} x - C_{xx} \dot{x} - K_{xy} y - C_{xy} \dot{y} \\
dF_y &= -K_{yx} x - C_{yx} \dot{x} - K_{yy} y - C_{yy} \dot{y}
\end{align*}$$

(44)

Thus, substituting Eq. (42) and (43) into Eq. (33) and grouping terms in accordance with Eq. (44) yields:

$$K_{xx} = K_{ff}' \cos^2 \psi + k_{ff}' \sin^2 \psi - (k_{ff} + k_{ff}') \cos \psi \sin \psi$$

(45)
A summation over all the pads making up the bearing gives the bearing spring and damping coefficients. If the pads have no inertia the equations simplify to:

\begin{align}
    k_{yx} &= k_{ff} \cos^2 \psi - k_{ff} \sin^2 \psi + (k_{ff}' - k_{ff}) \cos \psi \sin \psi \\
    \omega C_{yx} &= \omega C_{ff}' \cos^2 \psi - \omega C_{ff} \sin^2 \psi + (\omega C_{ff}' - \omega C_{ff}) \cos \psi \sin \psi \\
    k_{yy} &= k_{ff} \cos^2 \psi + k_{ff}' \sin^2 \psi + (k_{ff}' + k_{ff}) \cos \psi \sin \psi \\
    \omega C_{yy} &= \omega C_{ff}' \cos^2 \psi + \omega C_{ff} \sin^2 \psi + (\omega C_{ff}' + \omega C_{ff}) \cos \psi \sin \psi
\end{align}

Thus, for symmetry around the x-axis and no pad inertia the cross-coupling terms disappear.

**Pad Motion**

The pad motion is given by Eq. (30) which can also be written:

\begin{align}
    \eta_p &= \frac{(K_{ff} + i\omega C_{ff}) f + (K_{ff} + i\omega C_{ff}) \gamma}{K_{ff} - M\omega^2 + i\omega C_{ff}} \\
    \eta_i &= \frac{(K_{ff} + i\omega C_{ff}) f}{K_{ff} + i\omega C_{ff}}
\end{align}

Let \( \eta_i = \eta_p \) for \( M = 0 \), i.e. for no pad inertia:

\begin{align}
    \eta_i &= \eta_p \frac{K_{ff} + i\omega C_{ff}}{K_{ff} + i\omega C_{ff}} f
\end{align}
Then:

\[ \frac{\eta_p}{\eta_o} = \frac{k_{\eta\eta} + i\omega C_{\eta\eta}}{k_{\eta\eta} - M\omega^2 + i\omega C_{\eta\eta}} = 1 + \frac{M\omega^2}{k_{\eta\eta} - M\omega^2 + i\omega C_{\eta\eta}} = 1 + \rho M\omega^2 - i\omega M\omega^2 \]

or:

\[ \frac{|\eta_p|}{\eta_o} = \sqrt{(1 + \rho M\omega^2)^2 + (\omega M\omega^2)^2} \]

\[ \arg(\eta_p) - \arg(\eta_o) = \tan^{-1}\left(\frac{\omega M\omega^2}{1 + \rho M\omega^2}\right) \]

\[ = \tan^{-1}\left(\frac{\omega C_{\eta\eta} M\omega^2}{k_{\eta\eta}(k_{\eta\eta} - M\omega^2) + (\omega C_{\eta\eta})^2}\right) \]

Eq. (58) gives the amplitude ratio, i.e. the magnification factor, and Eq. (59) gives the phase angle lag with respect to an inertialess pad. Thus the two equations indicate how well the pad follows the shaft motion.

The phase angle becomes 90° when:

\[ M_{\text{crit}} \omega^2 = \frac{k_{\eta\eta}^2 + (\omega C_{\eta\eta})^2}{k_{\eta\eta}} \]

Hence, if the pad mass satisfies Eq. (60) there will be a resonance of the pad motion.

It is convenient to use Eq. (60) to establish a value for \( M \), designated the critical mass. In dimensionless form Eq. (60) may be written:

\[ \frac{\omega M_{\text{crit}}}{\sqrt{(\mu Bu)^2}} = \frac{1}{4\pi^2 S} \left(\frac{Ck_{\eta\eta}}{W}\right)^2 + \left(\frac{C\omega C_{\eta\eta}}{W}\right)^2 \]

where \( S \) is the bearing Sommerfeld number.
CONCLUSIONS

1. An analytical method has been established for calculating the spring and damping coefficients of the tilting pad journal bearing. The coefficients can be used directly in computing the critical speed, the response and the transmitted force of a rotor.

2. Numerical results for several bearing configurations have been obtained and are presented in form of design curves. A comparison with test results shows fair agreement.

3. The results indicate a sudden reduction in stiffness when approaching resonance of the pad motion. The implications of this behavior have not been assessed.
RECOMMENDATIONS

1. Since the force attenuation obtainable with a tilting pad journal bearing is not necessarily limited by oil whip, a study should be undertaken on how to optimize the attenuation by a proper choice of bearing dimensions, notably the clearance.

2. The present analysis considers the motion of the shoes around an axial axis, i.e. the "rolling" of the shoes, a more complete investigation should incorporate both "pitch" and "yaw" of the shoe.

3. Although a tilting pad journal bearing is considered "inherently stable" this holds true only in the idealized case where the shoes have no inertia and the pivots are frictionless. Hence, the stability limit of a tilting pad bearing should be studied, especially the stability of the shoe motion.
REFERENCES


APPENDIX A

COMPUTER PROGRAM PNO078: SPRING AND DAMPING COEFFICIENTS FOR THE TILTING PAD JOURNAL BEARING
APPENDIX A

Computer Program PNO078: Spring and Damping Coefficients for the Tilting Pad Journal Bearing

The program calculates the spring and damping coefficients for the tilting pad journal bearing based on the analysis given in this report. The program consists of the partial derivatives of forces, the position of the pads and the steady state bearing eccentricity ratio. The forces and their derivatives can be obtained from PNO091, Appendix C, and data reduced by PNO131, Appendix B. The output from the latter program serves directly as input to PNO078.

Input Data

The program is written for the IBM 1620 computer, 40K memory storage, with input and output on punched cards.

Card 1 and 2

Descriptive text, Col. 2 to 52

Card 3

7 : (I5)

Word 1 (NP) gives the number of pads. NP <= 12
Word 2 (NEB) gives the number of bearing eccentricity ratios for which calculations are performed. There is no upper limit on NEB.
Word 3 (NM) gives the number of calculations performed for each bearing eccentricity ratio with different values of the pad inertia. The maximum value of NM is 12. If the pad inertia is neglected, set NM = 1.
Word 4 (NF) controls the meaning of the input items $f_r$ and $f_t$ in the pad data ($f_r$ = radial force component, $f_t$ = tangential force component). If NF = 0, then the input value for $f_r$ is the pad Sommerfeld number and the input value for $f_t$ is the horizontal pad force $f_t$ (usually $f_t = 0$). If NF  0, then $f_r$ input and $f_t$ input = $f_t$. 
Word 5 (MC) controls the way in which the pad inertia data is given as input. If MC = -1, each pad will be given its own inertia. If MC = 0, each pad has zero inertia. If MC = +1, all the pads have the same inertia.

Word 6 (MS) controls the form in which the pad inertia is made dimensionless in the input. Denote the input value \( \bar{M} \). Then:

- If \( MS = -1 \), \( \bar{M} = \frac{1}{S} \frac{MN^2}{W/C} \)
- If \( MS = 0 \), \( \bar{M} = \frac{MN^2}{W/C} \)
- If \( MS = +1 \), \( \bar{M} = \frac{1}{S} \frac{MN^2}{W/C} = \frac{MWC}{[\mu DL(\xi)]^2} \)

where:
- \( C \) Radial pad clearance, inch
- \( D \) Journal diameter, inch
- \( L \) Bearing length, inch
- \( I \) Mass moment of inertia of pad around longitudinal axis, lbs-in-sec²
- \( M = I/R^2 \), equivalent pad mass, lbs-sec²/in.
- \( N \) Rotational speed, RPS
- \( R \) Journal radius, inch
- \( S = \frac{\mu NDL}{W} \left( \frac{R}{C} \right)^2 \), overall bearing Sommerfeld Number
- \( W \) Total bearing reaction, lbs.
- \( \mu \) Lubricant viscosity, lbs-sec/in²

In general it is convenient to use the last form of \( \bar{M} \) (i.e. for \( MS = +1 \)) since it is independent of speed.

Word 7 (NPR). If NPR = 1, the output includes the calculated values of the pad spring and damping coefficients in addition to the overall bearing coefficients. If NPR = 0, the pad coefficients are not given, only the bearing coefficients.

List of Pad Position Angles

This input list defines the position of the pads by means of the angle \( \Psi \) measured from the vertical to the pivot point of the pad (See Fig.22). \( \Psi \) is measured in degrees. There must be NP-values (word 1, card 3), 4 values per card.
List of Equivalent Pad Masses if MC = +1
4(E15.7)

This input list is used when it is desired to study the effect of pad inertia in the calculations and all pads have identical inertia. This form of the input list can only be supplied when MC=+1 (word 5, card 3). The pad inertia is given in form of the dimensionless equivalent pad mass $\overline{M}$. The definition of $\overline{M}$ depends on the value of MS (word 6, card 3). The input list must contain as many values of $\overline{M}$ as given by NM (word 3, card 3), 4 values per card. Thus it is possible to obtain the dynamic bearing coefficients as functions of the pad inertia.

List of Pad Forces and Derivatives

The spring and damping coefficients are calculated from the pad forces and their derivatives. These quantities can be computed by means of PNO091, Appendix C, and data-reduced by program PNO131, Appendix B. The output cards from PNO131 can be used directly to make up the present input list. The input list consists of a number of sets of complete bearing data, identified by a bearing eccentricity ratio. There are as many sets as given by NEB (word 2, card 3). Each set comprises the forces and their derivatives for each pad. To illustrate:

1st Pad
- Card 1 (3(E15.7) ) $\varepsilon_0$, $\varphi$, C/C
- Card 2a (4(E15.7) ) $\varepsilon$, $\varphi$, $f_r$, $f_x$, $\delta f_x/\delta \varphi$, $\delta f_x/\delta \varepsilon$
- Card 3a (4(E15.7) ) $\delta f_x/\delta \delta f_x$, $\delta f_x/\delta \delta \varphi$
- Card 4a (4(E15.7) )

$\varepsilon_0$

2nd Pad
- Card 2b (4(E15.7) ) $\varepsilon$, $\varphi$, $f_r$, $f_x$

1st Card 4b (4(E15.7) )

Last Pad
- Card 4x (4(E15.7) )

Card 1 (3(E15.7) )

2nd Card 2 (4(E15.7) )

$\varepsilon_0$ 1st Pad
where

- $\epsilon_o$ = bearing eccentricity ratio, see Fig. 22
- $\phi_o$ = bearing attitude angle, degrees, see Fig. 22
- $C'/C$ = ratio of pivot circle clearance to pad clearance
- $\epsilon$ = pad eccentricity ratio, see Fig. 22
- $\phi$ = pad attitude angle
- $f_r$ = dimensionless radial force component
- $f_t$ = dimensionless tangential force component
- $\omega$ = angular speed of journal, rad/sec

"Card 1" above is actually not used by the program in the calculations. It serves only the purpose of identification. Furthermore, it should be noted that $\epsilon$ cannot be zero since the calculations include $f_r/\epsilon$ and $f_t/\epsilon$.

**List of Equivalent Pad Masses if MC = -1**

4(E15.7)

This input is used when the pads have different inertia. This form of the input list can only be supplied when MC=-1 (word 5, card 3). The pad inertia is given in the form of the dimensionless equivalent pad mass $\bar{M}$ as defined through the value of MS (word 6, Card 3). An input list must be given for each bearing eccentricity ratio, i.e. following "Card 4x" for each $\epsilon_o$ above. The list is made up of NM - sets (word 3, card 3) and each set contains NP values of $\bar{M}$ (word 1, card 3). To illustrate:

1st $\epsilon_o$

```
| Card 4x | \( \delta f_r/\partial (\epsilon_o) \) | \( \delta f_t/\partial (\epsilon_o) \) |

| \( \bar{M}_{pad 1} \) | \( \bar{M}_{pad 2} \) | \( \bar{M}_{pad 3} \) | \( \bar{M}_{pad 4} \) |
| \( \bar{M}_{pad 5} \) | --- | --- | --- |
| \( \bar{M}_{pad 6} \) | --- | --- | --- | --- |
| \( \bar{M}_{pad NP} \) | --- | --- | --- | --- |
```

List of Equivalent Pad Masses for MC = -1. NM sets

2nd $\epsilon_o$

```
| Card 1 | $\epsilon_o$ | $\phi_o$ | $C'/C$ |
```

List of Pad Forces and Derivatives (continued)
Output Data

The output first repeats the input data for identification and checking purposes. Then follow the results for each bearing eccentricity ratio. The text identifying the numerical values is explained below except where the meaning is obvious:

SOMMERF.NO. - When part of pad results: $S_n = \mu_{NDL} \left( \frac{R}{C} \right)^2 (F=\text{total force on pad})$

When part of bearing results: $S = \mu_{NDL} \left( \frac{R}{C} \right)^2 (W=\text{total force on bearing})$

$F - T = f \gamma$

$F \times I = -1/S_p$, where $S_p$ is the pad Sommerfeld Number.

$K_{11}, K_{12}, K_{21}, K_{22} = K_{11}'', K_{12}'', K_{21}'', K_{22}''$ i.e. the 4 spring coefficients for the fixed pad.

$W_{C11}, W_{C12}, W_{C21}, W_{C22} = \omega_{C11}', \omega_{C12}', \omega_{C21}', \omega_{C22}'$ i.e. the 4 damping coefficients for the fixed pad.

The 8 coefficients are dimensionless in the form:

$K_{ff} = \frac{F}{F} K_{ff}$

$\omega_{Cff} = \frac{F}{F} \omega_{Cff}$

where $F$ is the total force on the pad. Thus $K_{11}$ is calculated as:

$K_{11} = S_p \left[ \frac{\partial^2}{\partial \xi^2} \cos^2 \phi - \frac{\partial^2}{\partial \epsilon^2} \sin^2 \phi - \left( \frac{\partial^2}{\partial \xi^2} - \frac{\partial^2}{\partial \epsilon^2} \right) \cos \phi \sin \phi - \frac{\epsilon \xi}{\xi} \right]$

and similarly for the other coefficients, see Eq. (15).

$\text{CRIT.MASS} = \frac{\omega_{CMR}}{(\mu_{NDL} \left( \frac{R}{C} \right)^2)}$

where part of pad results.

$\text{CRIT.MASS} = \frac{\omega_{CMR}}{(\mu_{NDL} \left( \frac{R}{C} \right)^2)}$

where part of bearing results.

$K_{D11}, K_{D12}, K_{D21}, K_{D22} = K_{11}', K_{12}', K_{21}', K_{22}'$ i.e. the spring coefficients for the tilting pad.

$\omega_{CD11}, \omega_{CD12}, \omega_{CD21}, \omega_{CD22} = \omega_{C11}', \omega_{C12}', \omega_{C21}', \omega_{C22}'$ i.e. the damping coefficients for the tilting pad.

The coefficients are dimensionless in the form:

$K_{ff}' = \frac{F}{F} K_{ff}'$

etc.
KXX, KXY, KYX, KYY = K_{xx}, K_{xy}, K_{yx}, K_{yy}, i.e. the spring coefficients for the complete bearing.

WCXX, WCXY, WCYX, WCYY = \omega C_{xx}, \omega C_{xy}, \omega C_{yx}, \omega C_{yy}, i.e. the damping coefficients for the complete bearing.

The coefficients are dimensionless in the form: \( K_{xx} = \frac{C_{xx}}{W} \), \( \omega C_{xy} = \frac{\omega C_{xy}}{W} \), etc. where \( W \) is the total force on the bearing. Thus for instance \( K_{xx} \) is calculated internally in the program directly on the basis of the input values for the pad forces and their derivatives and the final result is multiplied by the bearing Sommerfeld Number \( S \).

DIMENSIONLESS PAD MASS - is the pad mass as defined through MS (word 6, card 3, input).

AMPLITUDE = \left| \frac{\eta_p}{\eta_0} \right|, i.e. the ratio between the actual pad angular amplitude and the angular amplitude of a massless pad, see Eq. (58). Hence, the ratio may be thought of as a magnification factor.

PHASE ANG = \arg(\eta_o) - \arg(\eta_p), i.e. the phase angle in degrees between the angular motion of a massless pad and the angular motion of the actual pad.

CALC. MASS = \begin{align*}
\frac{1}{S} \frac{M_0^2}{W/C} & \quad \text{Four different forms of the dimensionless pad inertia. For details, see explanation in "Input Data", Card 3, Word 6.} \\
\frac{1}{S} \frac{MN^2}{W/C} & \\
\frac{MN^2}{W/C} & \\
\frac{1}{S^2} \frac{MN^2}{W/C} = \frac{MWC}{(\mu DLR_C)^2} &
\end{align*}

CRIT. SP.RATIO = \frac{\omega_{\text{crit}}}{\omega} / M = \frac{\omega_{\text{crit}}}{\omega}, i.e. the ratio between the resonant pad frequency \( \omega_{\text{crit}} \) and the operating angular journal speed for the given Sommerfeld Number.
INPUT FORM FOR
PN0078: Spring and Damping Coefficients for the Tilting Pad Journal Bearing.

IBM 1620 Computer, FORTRAN I

Card 1: Text, Col. 2-52

Card 2: Text, Col. 2-52

Card 3: 7 (I5)

1. NP. Number of pads, NP \leq 12
2. NEB. Number of bearing eccentricity ratios
3. NM. Number of calculations including pad inertia per bearing eccentricity ratio. NM \leq 12.
4. NF. If NF = 0, input value of \( f_r \) is pad Sommerfeld number and input value of \( f_t \) is horizontal pad force \( f_t \). If NF = 1, \( f_r = f_t \) and \( f_t = f_t' \).
5. MC. If MC = -1: each pad is given its own inertia
   If MC = 0: all pads have zero inertia
   If MC = +1: all pads have same inertia
6. MS If MS = -1: dimensionless pad inertia = \( \frac{1}{S} \frac{MN^2}{W/C} \)
   If MS = 0: dimensionless pad inertia = \( \frac{MN^2}{W/C} \)
   If MS = +1: dimensionless pad inertia = \( \frac{1}{S^2} \frac{MN^2}{W/C} = \frac{MWC}{(\mu DL)^2} \)
6. NPR If NPR = 0: output does not include results for individual pads
   If NPR = 1: output includes results for individual pads

List of Pivot Point Angles: 4(E15.7)

Give the position angle, degrees, of the pivot point for each pad, in total NP - values, 4 values per card.


List of Dimensionless Pad Inertia, MC = +1: \( 4 \cdot (E15.7) \)

Use this list only when \( MC = +1 \). Give the dimensionless pad inertia (same for all pads) according to value of MS, in total NM - values, 4 values per card:

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List of Bearing and Pad Data: 4 \((E15.7)\)

Give NEB-sets of cards. If \( MC \neq -1 \), each set contains \((1+3 \cdot NP)\) cards.
If \( MC = -1 \), each set contains \((1 + 3NP + MN \cdot \{NP\}\) cards:

\[
\begin{array}{cccc}
\varepsilon & \Phi & \frac{C}{C} & \\
\frac{d\varepsilon}{d\Phi} & \frac{d\Phi}{d\varepsilon} & \frac{d\xi}{d\Phi} & \frac{d\xi}{d\varepsilon} \\
\end{array}
\]

1st Pad

\[
\begin{array}{cccc}
\varepsilon & \Phi & f_r & f_t \\
\frac{d\varepsilon}{d\Phi} & \frac{d\Phi}{d\varepsilon} & \frac{d\xi}{d\Phi} & \frac{d\xi}{d\varepsilon} \\
\end{array}
\]

2nd Pad

\[
\begin{array}{cccc}
\varepsilon & \Phi & f_r & f_t \\
\frac{d\varepsilon}{d\Phi} & \frac{d\Phi}{d\varepsilon} & \frac{d\xi}{d\Phi} & \frac{d\xi}{d\varepsilon} \\
\end{array}
\]

3rd Pad

\[
\begin{array}{cccc}
\varepsilon & \Phi & f_r & f_t \\
\frac{d\varepsilon}{d\Phi} & \frac{d\Phi}{d\varepsilon} & \frac{d\xi}{d\Phi} & \frac{d\xi}{d\varepsilon} \\
\end{array}
\]

4th Pad

\[
\begin{array}{cccc}
\varepsilon & \Phi & M_{pad} & \varepsilon \\
\frac{d\varepsilon}{d\Phi} & \frac{d\Phi}{d\varepsilon} & \frac{dM_{pad}}{d\Phi} & \frac{dM_{pad}}{d\varepsilon} \\
\end{array}
\]

Pad Inertia

\[
\begin{array}{cccc}
\varepsilon & \Phi & M_{pad} & NP \\
\frac{d\varepsilon}{d\Phi} & \frac{d\Phi}{d\varepsilon} & \frac{dM_{pad}}{d\Phi} & \frac{dM_{pad}}{d\varepsilon} \\
\end{array}
\]

Pad Inertia

If \( MC = -1 \)
2nd Bearing Eccentricity Ratio

\[ \frac{C}{C} \]

1st pad

2nd pad

3rd pad

4th pad

Pad inertia if \( MC = 1 \)

Pad

\[ \begin{align*}
\tilde{M}_{\text{pad1}} & \quad \tilde{M}_{\text{pad2}} \\
\tilde{M}_{\text{padNP}} & \quad \text{Pad}
\end{align*} \]
DIMENSION PVA(20), CSQ(20), SN(20), CSQ(20), SNQ(20), CSN(20), P(20)
DIMENSION B(8), GXX(20), GXY(20), GYY(20), HXX(20), HXY(20)
DIMENSION HYX(20), HYY(20), CMS(20)

50 READ 200
   READ 201
   READ 202, NP, NEB, NM, NF, MC, MS, NPR
   PUNCH 203
   PUNCH 200
   PUNCH 201
   PUNCH 204
   PUNCH 205, NP, NEB, NM, NF, MC, MS, NPR
   PUNCH 206
   DO 51 I = 1, NP
   READ 207, PVA(I), PVA(I+1), PVA(I+2), PVA(I+3)
   51 PUNCH 207, PVA(I), PVA(I+1), PVA(I+2), PVA(I+3)
   IF (MC) 54, 54, 52
   DO 53 I = 1, NM
   READ 207, P(I), P(I+1), P(I+2), P(I+3)
   54 DO 55 I = 1, 8
   55 B(I) = 0.0
   DO 100 I = 1, NP
      C2 = 0.017453293*PVA(I)
      C1 = COS(C2)
      C2 = SIN(C2)
      CS(I) = C1
      SN(I) = C2
      CSQ(I) = C1*C1
      SNQ(I) = C2*C2
   100 CSN(I) = C2*C1
   DO 157 I = 1, NEB
   READ 208, ECB, ATB, CRT
   PUNCH 210
   PUNCH 209, ECB, ATB, CRT
   BSN = 0.0
   FH = 0.0
   DO 120 J = 1, NP
   READ 207, ECP, ATP, FR, FT, DFRE, DFTE, DFTA, DFRS, DFTS
   C1 = 0.017453293*ATP
   CSP = COS(C1)
   SNP = SIN(C1)
   CSPQ = CSP*CSP
   SNPQ = SNP*SNP
   CS2 = SNP*CSP
   IF (NF) 101, 104, 101
   101 C1 = FR
   C2 = FT
   FXI = -FR*CSP - FT*SNP
   FT = -FR*SNP + FT*CSP
   FR = SQRT (FR*FR + C2*C2)
   IF (FR) 103, 102, 103
   102 FR = 1.0E90
   GO TO 105
   103 FR = 1.0/FR
   GO TO 105
   104 FXI = -1.0/FR
   105 BSN = BSN + FXI*CS(J) - FT*SN(J)
   FH = FH - FXI*SN(J) - FT*CS(J)
   PUNCH 211, J
   PUNCH 212
   PUNCH 213, ECP, ATP, FR, FT, DFRE
   PUNCH 214
PUNCH 213, DFTE, DFRA, DFTA, DFRS, DFTS

106 PUNCH 215
PUNCH 213, C1, C2, FXI
107 C1 = (DFRE + DFTA) * CS2
C2 = (DFTE - DFRA) * CS2
C3 = DFRS * CS2
C4 = DFTS * CS2
IF (ECP) 109, 108, 109

108 C5 = 0.0
C6 = 0.0
GO TO 110
109 C5 = FXI / ECP
C6 = FT / ECP
110 AXX = DFRE * CSPQ - DFTA * SNPQ + C2 - C6 * SNP
AXY = DFRA * CXPQ + DFTE * SNPQ + C1 + C6 * CSP
AYX = -DFTE * CSPQ - DFRA * SNPQ + C1 + C5 * SNP
AYY = -DFTA * CSPQ + DFRE * SNPQ - C2 - C5 * CSP
C5 = 2.0 * C5
C6 = 2.0 * C6
BXX = DFRS * CSPQ + C4 - C5 * SNP
BXY = DFTS * SNPQ + C3 + C5 * CSP
BYX = -DFTS * CSPQ + C3 - C6 * SNP
BYY = DFRS * SNPQ - C4 + C6 * CSP
IF (AYY) 56, 55, 56

55 C10 = 0.0
GO TO 57
56 C10 = (AYY + BYY) / (AYY * BYY) / 39.478418
57 CMS(J) = C10.
IF (NPR) 111, 112, 111
111 C1 = FR * AXX
C2 = FR * AXY
C3 = FR * AYY
C4 = FR * BXX
C5 = FR * BXY
C7 = FR * BYX
C8 = FR * BYY
C9 = C10 / FR
PUNCH 216
PUNCH 213, C1, C2, C3, C4
PUNCH 217
PUNCH 213, C5, C6, C7, C8, C9

112 IF (MC) 119, 113, 119
113 C1 = AYY * AYY + BYY * BYY
IF (C1) 115, 114, 115
114 PUNCH 218, J
GO TO 120
115 C2 = AXY * AXY - BXY * BYX
C3 = AXY * BXY + AYY * BXY
C4 = AXX - (AYY * C2 + BYY * C3) / C1
C5 = BXX - (AYY * C3 - BYY * C2) / C1
C1 = CSQ(J)
C2 = SQ(J)
C3 = SNQ(J)
B(1) = B(1) + C4 * C1
B(2) = B(2) + C4 * C3
B(3) = B(2)
B(4) = B(4) + C4 * C2
B(5) = B(5) + C5 * C1
B(6) = B(6) + C5 * C3
B(7)=B(6)
B(8)=B(8)+C5*C2
IF (NPR) 117,118,117
117 PUNCH 219,C4,C5
118 GO TO 120
119 GXX(J)=AXX
GYX(J)=AXY
GYX(J)=AYX
GYY(J)=AYY
HXX(J)=BXX
HXY(J)=BXY
HYX(J)=BYX
HYY(J)=BYY
120 CONTINUE
BSN=1.0/BSN
IF (MC) 123,121,123
121 DO 60 J=1,8
60 B(J)=BSN*B(J)
PUNCH 220
PUNCH 213,B(1),B(2),B(3),B(4),B(5)
PUNCH 221
PUNCH 213,B(6),B(7),B(8),BSN,FH
PUNCH 228
DO 158 J=1,NP
C9=CMS(J)/BSN
158 PUNCH 229,J,C9
DO 122 J=1,8
122 B(J)=0.0
GO TO 157
123 DO 156 K=1,NM
IF (MC) 124,126,126
124 DO 125 J=1,NP,4
125 READ 207,P(J),P(J+1),P(J+2),P(J+3)
GO TO 127
126 PM=P(K)
PUNCH 223,PM
127 DO 151 J=1,NP
IF (MC) 128,129,129
128 PM=P(J)
129 PZ=39.478418*PM
IF (MS) 130,131,132
130 PC=PZ
P1=PM
P2=BSN*PM
P3=PM/BSN
GO TO 133
131 PC=PZ/BSN
P1=PM/BSN
P2=PM
P3=P1/BSN
GO TO 133
132 PC=PZ*BSN
P1=PM*BSN
P2=P1*BSN
P3=PM
133 AXX=GXX(J)
AXY=GXY(J)
AYX=GYX(J)
AYY=GYY(J)
BXX=HXX(J)
BXY=HXY(J)
BYX = HYX(J)
BYY = HYY(J)
C3 = AYY * PC
C4 = C3 * C3 + BYY * BYY
IF (C4) 137, 134, 137
134 IF (AYY) 136, 135, 136
135 PUNCH 218, J
GO TO 151
136 PUNCH 222, J
GO TO 151
137 C1 = C3 / C4
C2 = BYY / C4
C3 = 1.0 + C1 * PC
C4 = C2 * PC
ETP = SQRT (C3 * C3 + C4 * C4)
IF (C3) 146, 138, 146
138 IF (BYY) 142, 139, 142
139 IF (AYX + BYX) 141, 140, 141
140 PUNCH 218, J
GO TO 151
141 ETP = 0.0
PHA = 180.0
GO TO 150
142 IF (C4) 145, 143, 144
143 PHA = 0.0
GO TO 150
144 PHA = 90.0
GO TO 150
145 PHA = -90.0
GO TO 150
146 C4 = C4 / C3
PHA = 57.29578 * ATAN (C4)
IF (C3) 147, 150, 150
147 IF (C4) 148, 150, 149
148 PHA = PHA + 180.0
GO TO 150
149 PHA = PHA - 180.0
150 C5 = -PC * C3
C6 = PC * C4
C7 = C1 * AXY + C2 * BXY
C8 = C1 * BXY - C2 * AXY
GR = AXX - C7 * AXY + C8 * BYX
GE = BXX - C7 * BYX - C8 * AXY
DR = -PC * C7
DE = -PC * C8
C3 = -PC * (C1 * AXY + C2 * BYX)
C4 = -PC * (C1 * BYX - C2 * AXY)
C8 = CMS(J) / BSN
IF (P3) 161, 160, 161
160 C9 = 1.0E+09
GO TO 162
161 C9 = C8 / P3
162 PUNCH 211, J
PUNCH 224
PUNCH 213, GR, DR, C3, C5, GE
PUNCH 225
PUNCH 213, DE, C4, C6, ETP, PHA
PUNCH 226
PUNCH 213, PC, P1, P2, P3
PUNCH 227
PUNCH 213, C8, C9
C8 = CSN(J)
C1 = (DR+C3)*C8
C2 = (DE+C4)*C8
C7 = (GR-C5)*C8
C8 = (GE-C6)*C8

AXX = CSN(J)
BXX = SNQ(J)

B(1) = B(1) + GR*AXX + C5*BXX - C1
B(2) = B(2) + DR*AXX - C3*BXX + C7
B(3) = B(3) + C3*AXX - DR*BXX + C7
B(4) = B(4) + C5*AXX + GR*BXX + C1
B(5) = B(5) + GE*AXX + C6*BXX - C2
B(6) = B(6) + DE*AXX - C4*BXX + C8
B(7) = B(7) + C4*AXX - DE*BXX + C8
B(8) = B(8) + C6*AXX + GE*BXX + C2

IF (BSN) 152, 154, 152
DO 153 J=1, 8
B(J) = BSN*B(J)
END
APPENDIX B

COMPUTER PROGRAM PN0131: INTERPOLATION OF PARTIAL ARC BEARING FORCE DERIVATIVES FOR USE IN TILTING PAD BEARING CALCULATIONS
APPENDIX B


In order to calculate the spring and damping coefficients for the tilting pad journal bearing it is necessary to know the force derivatives of the pads making up the complete bearing. Since these derivatives are calculated numerically in a computer by finite difference methods they are usually only available for a limited number of eccentricity ratios. When the pad is incorporated in the tilting pad bearing its eccentricity ratio is determined by the bearing eccentricity ratio (see Eq. (B6)) and it, therefore, becomes necessary to know the force derivatives at other eccentricity ratios than the ones for which calculations have been performed. It is rather costly and impractical to obtain this data by performing additional calculations with the finite difference computer program and instead interpolation of the existing data can be employed. It has been found that to do this manually (by curve plotting) is time consuming. Hence, a computer program has been written for the purpose of performing the interpolation. The computer program accepts as input a list of the force derivatives as computed from the finite difference program (i.e. PNO091, Appendix C). In addition the number of pads and their position must be given. Then the computer program calculates the force derivatives for each pad for a specified number of bearing eccentricity ratios. The output can be used directly as input for the tilting pad bearing program (i.e. PNO078, Appendix A).

Analysis

The interpolation is based on quadratic curve fitting. Let the function to be interpolated be denoted $q$, which is given for a number of discreet eccentricity ratio values by the input. Then:

$$(B1) \quad q = A_i (\varepsilon - \varepsilon_i)^2 + B_i (\varepsilon - \varepsilon_i) + q_i \quad \varepsilon_{i-1} \leq \varepsilon \leq \varepsilon_{i+1}$$
where:

\[
A_i = \frac{\Delta_2 (q_{i+1} - q_i) - \Delta_1 (q_i - q_{i-1})}{\Delta_1 + \Delta_2}
\]

\[
B_i = \frac{\Delta_2 (q_{i+1} - q_i) + \Delta_1 (q_i - q_{i-1})}{\Delta_1 + \Delta_2}
\]

\[
\Delta_1 = \varepsilon_i - \varepsilon_{i-1}
\]

Thus, for each \(\varepsilon\)-interval, except the first and the last, \(q\) is determined by two equations. The average value is used:

\[
q = \frac{1}{2} \left[ A_i (\varepsilon - \varepsilon_i)^2 + B_i (\varepsilon - \varepsilon_i) + q_i + A_{i+1} (\varepsilon - \varepsilon_{i+1})^2 + B_{i+1} (\varepsilon - \varepsilon_{i+1}) + q_{i+1} \right]
\]

There are nine functions to be interpolated. Since they all (except the attitude angle) become infinite when \(\varepsilon = 1\) it is found necessary to normalize the functions before interpolation takes place. The normalization is chosen as follows:

\[
\varepsilon_L \leq \varepsilon \leq \varepsilon_{L+1} \quad \text{No normalization}
\]

\[
\varepsilon_L \leq \varepsilon \leq \varepsilon_n \quad \text{Normalization used:}
\]

\[
\left\{ \begin{array}{l}
E_t, \frac{\partial E_t}{\partial \varepsilon}, \frac{\partial^2 E_t}{\partial \varepsilon^2} \text{ is divided by } \left(1 - \varepsilon^2 \right)^2 \\
E_p, \frac{\partial E_p}{\partial \varepsilon}, \frac{\partial^2 E_p}{\partial \varepsilon^2} \text{ is divided by } \left(1 - \varepsilon^2 \right)^2 \\
\frac{\partial \phi}{\partial \varepsilon}, \frac{\partial^2 \phi}{\partial \varepsilon^2} \text{ is divided by } \frac{\varepsilon^2}{1 - \varepsilon^2}
\end{array} \right.
\]

Here \(\varepsilon_L\) and \(\varepsilon_n\) are the first and last eccentricity ratio, respectively, in the input list. \(\varepsilon_L\) is specified by the input as described later.

Let the tilting pad bearing have \(m\) pads whose pivot points are located the angle \(\psi\) from the vertical, measured in the direction of rotation. Then the eccentricity ratio of any pad is determined by:

\[
\varepsilon \cos \phi = 1 - \frac{E'}{E} - \varepsilon_e \cos (\psi - \phi_0)
\]

where \(\varepsilon_e\) is the bearing eccentricity ratio, \(\phi_0\) is the bearing attitude angle, \((1 - \varepsilon')\) is the preload, \(\varepsilon\) is the pad eccentricity ratio and \(\phi\) is the pad attitude angle. The input specifies \(\varepsilon_e, \phi_0, \varepsilon'\) and \(\psi\) and \(\phi\) is given as a function of \(\varepsilon\) in the earlier mentioned input list.
Thus, $\varepsilon \cos \varphi$ can be calculated from the input list and by interpolation Eq. (B6) can be solved to find $\varepsilon$ for each pad. With $\varepsilon$ determined the 8 force parameters can be calculated by interpolation using Eq. (B4) and (B5).

**Input Data**

The program is written for the IBM 1620 computer, 40K memory storage, with input and output punched on cards.

**Card 1**
Descriptive text, Column 2 to 52

**Card 2**
5 X (IXI4)

**Word 1** (NEP) gives the number of pad eccentricity ratios for which the force derivatives are provided in the pad data list. NEP $\leq$ 15.

**Word 2** (LM) gives the value of $L$ as used in Eq. (B5). If the pad eccentricity ratios in the pad data list are $\varepsilon_1, \varepsilon_2, \ldots, \varepsilon_n$ then the pad data are divided by a scale factor internally in the program for $\varepsilon \geq \varepsilon_0$. Since the scale factor is zero for $\varepsilon = 0$ it has been found necessary to set the scale factor equal to 1 for $\varepsilon < \varepsilon_0$. In the determination of $L$ the most critical parameter is $\varepsilon$, which is usually negative for $\varepsilon \leq 0.5$. Hence, it is suggested to give $L$ such a value that $\varepsilon_0 + \varepsilon \leq 0.6$.

**Word 3** (NBE) gives the number of tilting pad bearing eccentricity ratios. There is no limit to NBE.

**Word 4** (NA) gives the number of pads. NA $\leq 20$.

**Word 5** (INP) If INP=0, more input data follows the present set of input. If INP $\neq 0$, the present set of input is the last set.
Pad Data List

5 · (1XE13.6)

The pad data is usually obtained from a finite difference solution of the Reynolds equation, see Appendix C. The pad forces, force derivatives and attitude angle are calculated for a number of pad eccentricity ratios $\varepsilon_1, \varepsilon_2, \ldots, \varepsilon_n$ and supplied as input to the present computer program. For each eccentricity ratio two cards are given:

First card: $\varepsilon \hspace{0.1cm} \varphi \hspace{0.1cm} f_r \hspace{0.1cm} \frac{df_r}{d\varepsilon} \hspace{0.1cm} \frac{df_r}{d\varphi}$

Second card: $\varepsilon \hspace{0.1cm} \frac{df_t}{d\varepsilon} \hspace{0.1cm} \frac{df_t}{d\varphi} \hspace{0.1cm} \varepsilon \varphi \hspace{0.1cm} \varepsilon \varphi$

where:

$\varepsilon$ = pad eccentricity ratio

$\varphi$ = pad attitude angle

$f_r = \frac{F_r}{L(2R)}$, dimensionless radial force

$f_t = \frac{F_t}{L(2R)}$, dimensionless tangential force

$\dot{\varepsilon}$ = radial journal center velocity, sec$^{-1}$

Here, $\varphi, f_r$ and $f_t$ are calculated for the given $\varepsilon$ at the steady state position such that the total force passes through the pivot point.

Hence, the pad data list consists of 2 · (NEP) cards (see word 1, card 2). The list is given in sequence such that $\varepsilon_1 < \varepsilon_2 < \ldots < \varepsilon_{\text{NEP}}$

List of Pivot Point Angles

5 · (1XE13.6)

This input list gives the values of the angle $\psi$ measured from the vertical to the pivot point of each pad in the direction of rotation, see Fig. 22. The angle is measured in degrees. There are (NA)-values in total (word 4, card 2), 5 values per card.
List of Bearing Eccentricity Ratios
3 - (E15.7)

The position of the journal center with respect to the bearing center is defined through the eccentricity ratio $E_0$ and the attitude angle $\varphi_0$ see Fig. 22. $E_0$ is calculated with respect to the pad clearance (pad radius minus journal radius) and $\varphi_0$ is in degrees, measured from the vertical to the line of centers in the direction of rotation. In the general case (off-center pivot position, unsymmetric arrangement of the pads) a considerable effort is involved in determining $\varphi_0$ but if the pads are centrally pivoted and located symmetrically around the vertical then $\varphi_0 = 0$.

There must be an input card for each bearing eccentricity ratio, in total (NBE) cards. Each card contains three words:

Word 1 gives the bearing eccentricity ratio $E_0$
Word 2 gives the bearing attitude angle $\varphi_0$ in degrees
Word 3 gives the ratio $C'/C$ where $C'$ is the radius of the pivot point circle minus the journal radius and $C$ is the pad radius minus the journal radius.

Output Data

The output gives first the values of the input data for identification and checking purposes. Then follow 6 lines of text describing the format of the output data and thereafter the numerical values of the output data arranged as follows (FORMAT 4(E15.7)):
Note that two values of the pad attitude angle $\varphi$ are given. The first value (i.e. on card la, lb etc) is obtained by interpolation of the pad data list whereas the second value (i.e. on card 2a, 2b etc) is calculated from $\varphi = \tan^{-1}(4_k/P_r)$. It is the latter value that should be actually used and by comparison with the first $\varphi$-value some indications are obtained of the accuracy of the interpolation.

The output cards can be used directly as input to the tilting pad bearing computer program (PN0078, Appendix A) by removing the 3 blank cards in front of each bearing eccentricity ratio calculation (card 1,2 and 3 above) and by removing the first card of each pad data output set (card la, lb etc. above).

Returning to Eq. (B6) it can be seen that geometrical considerations limit the range of the bearing eccentricity ratio $\varepsilon$. The program takes care of this automatically as follows. For a given $\varepsilon, c/C$ and $\psi$ the right hand side of Eq. (B6) may be calculated. Denote the result $(E \cos \varphi)_{calc}$. Three cases may arise:

a) If $(E \cos \varphi)_{calc} \leq (E \cos \varphi)_{calc} \leq |\varepsilon|$, Eq. (B6) is solved by interpolation. Here $\varepsilon$, is the first eccentricity ratio value in the pad data list.

b) If $(E \cos \varphi)_{calc} < (E \cos \varphi)_{calc}$, the computer writes: "NO LOAD, PVT. ANG=X.XX, $E \star \cos(A)=X.XX$" indicating that the particular pad, identified by its pivot angle $\psi$ has a $(E \cos \varphi)_{calc}$-value less than the lowest value in the pad data list. Strictly speaking there may be solutions in this case but assuming that $\varepsilon$, in general is almost zero it is safest not to allow extrapolation.

c) If $(E \cos \varphi)_{calc} \geq |\varepsilon|$, the computer writes: "E=1, PVT. ANG=X.XX, $E \star \cos(A)=X.XX$" indicating that for the particular pad the journal is either touching or beyond the pad surface. The calculations proceed with the next pad.
INPUT FORM FOR

PN0131: Interpolation of Partial Arc Bearing Force Derivatives for Use in Tilting Pad Bearing Calculations

IBM 1620 Computer, FORTRAN I

Card 1 (Text Col.2-52)

Card 2: 5(1X14)

1. NEP: Number of pad eccentricity ratios in pad data list. NEP ≤ 15
2. LM: \( \varepsilon \leq \varepsilon_{LM+1} = \text{pad data not normalized} \)
\( \varepsilon \geq \varepsilon_{LM} = \text{pad data is normalized} \) see Eq. (B5)
3. NBE: Number of bearing eccentricity ratios
4. NA: Number of pads, NA ≤ 20
5. INP: If INP=0: more input follows. If INP≠0, last set of input.

Pad Data List: 5(1XE13.6)

Give 2 cards per pad eccentricity ratio, in total 2 · NEP cards:
Pivot Point Angles: 5(1XE13.6)
Give the position angle(degrees) of the pivot point for each pad,
in total NA-values, 5 values per card:

— — — — —  

Bearing Eccentricity Ratio: 3(1E15.7)
Give one card per bearing eccentricity ratio, in total NBE cards:

— — — — —  \( \varepsilon_c, \varphi_c, c'/c \)  
— — — — —  
— — — — —  
— — — — —  
— — — — —  
— — — — —  
— — — — —  
— — — — —  
— — — — —  

DIMENSION P(11,15),PN(11,15),A(11,15),B(11,15),PA(20),Q(11)

49 READ 101
READ 102,NEP,LM,NBE,NA,INP
PUNCH 100
PUNCH 101
PUNCH 104
PUNCH 103,NEP,LM,NBE,NA,INP
PUNCH 107
DO 47 I=1,20
47 PA(I)=0.0
DO 48 I=1,11
Q(I)=0.0
DO 48 J=1,15
A(I,J)=0.0
B(I,J)=0.0
P(I,J)=0.0
48 PN(I,J)=0.0
KM=LM+1
NM=NEP-1
DO 55 J=1,NEP
PUNCH 108,J
DO 50 I=1,6,5
READ 109,P(I,J),P(I+1,J),P(I+2,J),P(I+3,J),P(I+4,J)
50 PUNCH 109,P(I,J),P(I+1,J),P(I+2,J),P(I+3,J),P(I+4,J)
C1=P(I,J)
C2=P(2,J)
C2=0.017453293*C2
C2=COSF(C2)
P(I1,J)=C1*C2
IF(J-LM) 55,51,51
51 C2=1.0-C1*C1
C3=C1/C2
C4=SQRT(F(C2))
C4=C3/C4
C2=C3*C3
DO 52 I=3,5
PN(I,J)=P(I,J)/C2
52 PN(I+3,J)=P(I+3,J)/C4
PN(9,J)=P(9,J)/C3
PN(10,J)=P(10,J)/C3
PN(2,J)=P(2,J)
55 CONTINUE
PUNCH 110
DO 56 I=1,NA,5
READ 109,PA(I),PA(I+1),PA(I+2),PA(I+3),PA(I+4)
56 PUNCH 109,PA(I),PA(I+1),PA(I+2),PA(I+3),PA(I+4)
PUNCH 113
PUNCH 114
PUNCH 115
PUNCH 116
PUNCH 117
PUNCH 118
DO 65 I=2,10
Y2=P(I,1)
Y3=P(I,2)
X2=P(I,1)
X3=P(I,2)
DO 65 J=2,NM
Y1=Y2
Y2=Y3
X1=X2
X2=X3
X3=P(I,J+1)
IF(J=K) 60,61,62
60 Y3=P(I,J+1)
GO TO 64
61 Y1=PN(I+1)
62 Y3=PN(I,J+1)
64 DT1=X2-X1
DT2=X3-X2
C1=Y2-Y1
C2=Y3-Y2
C3=X3-X1
C4=DT1/DT2
A(I,J)=(C2*C4-C1)/(DT1*C3)
B(I,J)=(C2*C4+C1/C4)/C3
65 CONTINUE
Y2=P(I+1)
Y3=P(I+2)
X2=P(I+1)
X3=P(I+2)
DO 66 J=2,NM
Y1=Y2
Y2=Y3
Y3=P(I,J+1)
X1=X2
X2=X3
X3=P(I,J+1)
DT1=X2-X1
DT2=X3-X2
C1=Y2-Y1
C2=Y3-Y2
C3=X3-X1
C4=DT1/DT2
A(I,J)=(C2*C4-C1)/(DT1*C3)
B(I,J)=(C2*C4+C1/C4)/C3
66 DO 239 K=1,NBE
READ 105,EB,ATB,CPC
PUNCH 112
PUNCH 105,EB,ATB,CPC
ATR=017453293*ATB
DO 238 L=1,NA
PAN=PA(L)
PAR=017453293*PAN
ECS=COSF(PAR-ATR)
ECS=10-CPC-EB*ECS
IF(ECS-P(I,J)) 202,205,201
201 PUNCH 119,L,PAN,ECS
GO TO 238
202 IF(ECS-P(I,J)) 203,203,204
203 PUNCH 120,L,PAN,ECS
GO TO 238
204 IF(ECS-P(I,J)) 206,200,200
200 KC=NEP
GO TO 209
205 KC=2
GO TO 209
206 DO 208 J=2,NEP
C1=P(I,J)
IF(ECS-C1) 207,207,208
207 KC=J
GO TO 209  
208 CONTINUE  
209 AR=A(I, KC)  
BR=B(I, KC)  
CR=P(I, KC)  
XR=P(I, KC)  
KL=KC-1  
AL=A(I, KL)  
BL=B(I, KL)  
CL=P(I, KL)  
XL=P(I, KL)  
C1=ECS-XR  
C1=CR+C1*(BR+C1*AR)  
C2=ECS-XL  
C2=CL+C2*(BL+C2*AL)  
IF(KC=2) 210 210 210  
210 C2=C1  
GO TO 213  
211 IF(NEP-KC) 212 212 212  
212 C1=C2  
213 EP=(C1+C2)*2.0  
C2=1.0-EP*EP  
C3=EP/C2  
C4=SQRTF(C2)  
C4=C3/C4  
C2=C3*C3  
DO 234 I=2,10  
AL=A(I, KL)  
BL=B(I, KL)  
XL=P(I, KL)  
AR=A(I, KC)  
BR=B(I, KC)  
XR=P(I, KC)  
IF(I=2) 220 220 220  
220 C1=1.0  
GO TO 226  
221 IF(I=5) 222 222 222  
222 C1=C2  
GO TO 226  
223 IF(I=8) 224 224 224  
224 C1=C4  
GO TO 226  
225 C1=C3  
226 QR=EP-XR  
QR=QR*(BR+QR*AR)  
QL=EP-XL  
QL=QL*(BL+QL*AL)  
IF(KC=LM) 227 227 227  
227 QR=QR+P(I, KC)  
228 QL=QL+P(I, KL)  
IF(KC=2) 229 229 229  
229 QL=QR  
GO TO 233  
230 QR=C1*(QR+PN(I, KC))  
IF(KC=KM) 228 228 228  
231 QL=C1*(QL+PN(I, KL))  
IF(KC=NEP) 233 233 233  
232 QR=QL  
233 Q(I)=(QR+QL)/2.0  
234 CONTINUE  
C2=Q(3)
IF (C2) 236,235,236
235 C1=Q(2)
GO TO 237
236 C1=ATANF(Q(6)/C2)
C1=57.295780*C1
237 PUNCH 117, PAN.ECS,Q(2)
PUNCH 106, EP,C1,C2,Q(6)
PUNCH 106, Q(4)*Q(7)*Q(9)*Q(10)
PUNCH 106, Q(5)*Q(8)
238 CONTINUE
239 CONTINUE
IF (INP) 240,49,240
240 STOP
100 FORMAT(49H1PNO131-INTERPOLATION OF PART.ARC DATA J*LUND 8-9,3H-63)
101 FORMAT(49H0)
102 FORMAT((XI4*1XI4*1XI4*1XI4*1XI4)
103 FORMAT((2XI4*6XI4*6XI4*6XI4*6XI4)
104 FORMAT((9H0NO.PAD E4X5HLIMIT3X8HNO.BRG.E3X7HNO.ANG.5X5HINPUT)
105 FORMAT(E15.7,E15.7,E15.7)
106 FORMAT(E15.7,E15.7,E15.7,E15.7)
107 FORMAT((49H0E ATT FR DFRDE DFRDEW FT DFTDE DFTDEW DFRED A DFT*3HEDA)
108 FORMAT(3HOE(91291H))
109 FORMAT((1XE13.6,1XE13.6,1XE13.6,1XE13.6,1XE13.6)
110 FORMAT((13H0PIVOT ANGLES)
111 FORMAT((1HO,E14.7,E15.7,E15.7)
112 FORMAT((//1HO)
113 FORMAT((14H0OUTPUT FORMAT)
114 FORMAT((8H EPS.BRG5X7HATT.BRG5X4HCP/C)
115 FORMAT((38H PIVOT ANG E*COS(A) ATT.PAD(INTP))
116 FORMAT((8H EPS.PAD5X7HATT.PAD5X2HFR8X2HFT)
117 FORMAT((7H DFR/DE6X6HDFT/DE6X7HDFR/EDA3X7HDFT/EDA)
118 FORMAT((23H DFR/D1(E/W) DFT/D(E/W))
119 FORMAT((4H0PAD I2,17H NO LOAD,PVT.ANG=,E13.6,10H,E*COS(A)={},E13.6)
120 FORMAT((4H0PAD I2,13H E=1,PVT.ANG=,E13.6,10H,E*COS(A)={},E13.6)
END
APPENDIX C

COMPUTER PROGRAM PN0091: NUMERICAL CALCULATION OF THE INCOMPRESSIBLE FULL OR PARTIAL JOURNAL BEARING
NUMERICAL CALCULATION OF
THE INCOMPRESSIBLE FULL OR
PARTIAL JOURNAL BEARING

by

Jorgen W. Lund

Prepared for

Mechanical Technology Incorporated
# TABLE OF CONTENTS

<table>
<thead>
<tr>
<th>Section</th>
<th>Page</th>
</tr>
</thead>
<tbody>
<tr>
<td>ABSTRACT</td>
<td>44</td>
</tr>
<tr>
<td>INTRODUCTION</td>
<td>45</td>
</tr>
<tr>
<td>ANALYSIS</td>
<td>47</td>
</tr>
<tr>
<td></td>
<td></td>
</tr>
<tr>
<td>Nomenclature</td>
<td>47</td>
</tr>
<tr>
<td>Analysis</td>
<td>48</td>
</tr>
<tr>
<td>COMPUTER INPUT INSTRUCTIONS</td>
<td>53</td>
</tr>
<tr>
<td>COMPUTER OUTPUT</td>
<td>58</td>
</tr>
<tr>
<td>COMPUTER OPERATION</td>
<td>59</td>
</tr>
<tr>
<td>COMPUTER TIME</td>
<td>59</td>
</tr>
<tr>
<td>704 INPUT FORMS</td>
<td>61</td>
</tr>
<tr>
<td>1620 INPUT FORMS</td>
<td>63</td>
</tr>
</tbody>
</table>
This report describes the computer program PN0091: "Hydrodynamic Incompressible Journal Bearing." The program exists in 2 versions, one for the IBM 704 computer and one for the IBM 1620 computer. The report contains a description of the analysis and the instructions for using the programs. The program solves the dynamical incompressible Reynolds equation in dimensionless form by finite difference equations and gives results for the pressure distribution, the load and the attitude angle. The rupture of the oil film is included. The bearing may be full 360° or a partial arc.
INTRODUCTION

The dynamical Reynolds equation is a partial differential equation relating the pressure in the fluid film to the journal rotational speed and the squeeze film velocity. The film is assumed to be isothermal so that the viscosity is constant. The equation is used in dimensionless form and is given as Eq. (4) page 4.

The bearing is taken to be operating in air. Therefore, no oil can enter the bearing from the ends and the film pressure cannot be lower than ambient. For this reason the film contracts in the diverging part of the bearing with a uniform pressure of $P_a$. This is denoted film rupture and is taken into account.

Reynolds equation is transformed into finite difference form and solved by iteration. The convergence of the iteration process is established in two ways, one in which the relative difference between two iterations is checked against a preassigned error and one which compares the result of each iteration against an extrapolated value, i.e., a form of absolute convergence.

The program provides for the calculation of force derivatives with respect to eccentricity ratio and the two squeeze film velocity components. Such derivatives are needed in computing spring and damping coefficients for the bearing.

If desired the program will iterate on the attitude angle in order to determine the attitude angle corresponding to no horizontal force component.
ANALYSIS

NOMENCLATURE

C  Diametrical clearance, inch
D  Journal diameter, inch
F  Force, lbs.

\( f_r, f_t \)  Radial and tangential force components, dimensionless, see Figure 1.
\( f_v, f_h \)  Vertical and horizontal force components, dimensionless, see Figure 1

h  Dimensionless film thickness

i, j  Finite difference coordinates, axial and circumferential, see Figure 2.

k  Iteration number
L  Bearing length, inch

m  Number of circumferential subdivisions, see Figure 2
n  Number of axial subdivisions, see Figure 2

N  Journal speed, RPS
P  Dimensionless pressure (above ambient)
R  Journal radius, inch
S  Sommerfeld number

t  Time, seconds

\( U_1, U_2 \)  Journal and bearing surface speed, inch/sec.

x, z  Circumferential and axial coordinates, dimensionless

\( y_k \)  Sum of all pressures after the k'th iteration, see Eq. (10)

\( y_{k\infty} \)  Extrapolated pressure sum after the k'th iteration, see Eq. (11)

\( \alpha \)  Attitude angle, degrees

\( \dot{\alpha} \)  Tangential speed of journal center, rad/sec.

\( \delta_A \)  Absolute convergence limit, see Eq. (12)

\( \delta_R \)  Relative convergence limit, see Eq. (9)

\( \Delta_A \)  Error in absolute convergence, see Eq. (12)

\( \Delta_R \)  Error in relative convergence, see Eq. (9)

\( \varepsilon \)  Eccentricity ratio

\( \dot{e} \)  Radial speed of journal center, sec\(^{-1}\)

\( \theta \)  Circumferential angular coordinate, see Figure 1.

\( \mu \)  Viscosity, lbs-sec/in\(^2\)

\( \omega \)  Angular speed of journal, rad/sec.
For an incompressible lubricant Reynolds equation is:

\[ \frac{\partial}{\partial x} \left[ \frac{h^3}{\mu} \frac{\partial P}{\partial x} \right] + \frac{\partial}{\partial z} \left[ \frac{1}{\mu} \frac{\partial P}{\partial z} \right] = \frac{1}{2} \frac{\partial}{\partial x} \left[ \bar{h}(U_1 + U_2) \right] + \frac{\partial \bar{h}}{\partial t} \]

To make dimensionsless set:

\[ x = \frac{x}{D}, \quad z = \frac{z}{L}, \quad h = \frac{h}{C}, \quad \bar{P} = \frac{P}{\mu N (E)^3} \]

Furthermore:

\[ U_1 = \pi D N \quad U_2 = 0 \]

Also set:

\[ h = \frac{1}{2} + \frac{\xi}{2} \cos(\theta - \alpha) \]

where \( \theta = \frac{x}{R} = 2x \)

Under isothermal conditions \( \mu \) is constant and the dimensionless Reynolds equation becomes:

\[ \frac{\partial}{\partial x} \left[ h^3 \frac{\partial P}{\partial x} \right] + (\bar{P})^2 \frac{\partial}{\partial z} \left[ h^3 \frac{\partial P}{\partial z} \right] = 12 \pi \left[ \frac{\xi}{\alpha} \cos(\theta - \alpha) + \left( \frac{\xi}{\omega} - \frac{\xi}{2} \right) \sin(\theta - \alpha) \right] \]

The corresponding finite difference equation is:

\[ P_{ij} = \frac{h_{i+1,j}^{3} P_{i+1,j} + h_{i-1,j}^{3} P_{i-1,j} + (P)^2 (\Delta x)^2 \left[ h_{i+1,j}^{3} + h_{i-1,j}^{3} P_{i+1,j} + h_{i+1,j}^{3} P_{i-1,j} \right] - 12 \pi \Delta x \left[ \frac{\xi}{\alpha} \cos(\theta - \alpha) + \left( \frac{\xi}{\omega} - \frac{\xi}{2} \right) \sin(\theta - \alpha) \right]}{h_{i+1,j}^{3} + h_{i-1,j}^{3} + (P)^2 (\Delta x)^2 \left[ h_{i+1,j}^{3} + h_{i-1,j}^{3} \right]} \]

where the finite difference mesh is as follows:

![Figure 2.](image-url)
The pressure at the bearing end is ambient (i.e. 0). At the leading and the trailing edge it is given by the input (partial arc only). At the bearing centerline there is symmetry. Thus the boundary conditions to Eq. (5) become:

\[
P_{ij} = 0 \\
P_{m+2,j} = P_{ij}
\]

(6)

\[
P_{ij} = \begin{cases} 
  P_{m+1} & \text{full 360° bearing} \\
  P_{\text{out}} \text{ (input)} , i= 2,3,\ldots,(n+1) & \text{partial arc} 
\end{cases}
\]

If \((P_{ij})_{\text{calculated}} < 0\), then \(P_{ij} = 0\)

The pressure is initially set equal to zero and Eq. (5) is then solved by iterations.

The load carrying capacity is calculated from:

(7) \[ f_r = -2 \int_0^{\theta_m} P \cos(\theta-\alpha) \, d\alpha d\theta = -2 \Delta x \Delta z \sum_{ij} P_{ij} \cos(\theta_j - \alpha) \]

(8) \[ f_\theta = 2 \int_0^{\theta_m} P \sin(\theta-\alpha) \, d\alpha d\theta = 2 \Delta x \Delta z \sum_{ij} P_{ij} \sin(\theta_j - \alpha) \]

The indicated summations are computed by Simpson rule of integration, i.e. \( \frac{1}{3} \Delta \left[ P_1 + 4P_2 + 2P_3 + 4P_4 + \cdots \right] \) and if either \( m \) or \( n \) are odd the first interval is integrated by \( \frac{1}{6} \Delta \left[ 5P_1 + 8P_2 - P_3 \right] \). In addition and corrections are made at the boundary to the ruptured film.

**Iteration Convergence and Extrapolation**

The convergence of the pressure iteration is tested by two methods, which shall be denoted relative convergence and absolute convergence. The relative convergence is tested as follows: after the \( k \)'th pressure iteration compute:
where $\delta_R$ is given by the input. When $\Delta_R$ becomes zero or negative then relative convergence has been achieved.

To compute the absolute convergence the following criteria is chosen. After each iteration the sum of the pressures are computed. Let the result for the $k$'th iteration be denoted:

$$ y_k = \sum_i \sum_j P_{ij}^{(n)} $$

Assume that after infinitely many iterations this sum will be $y_{k=\infty}$ and set:

$$ y_k = y_{k=\infty} - A e^{-B_k} \quad (A \text{ and } B \text{ constants}) $$

from which:

$$ y_{k=\infty} = y_k + \frac{(y_k-y_{k-1})^2}{2y_{k-1}-y_{k-2}-y_k} $$

$y_{k=\infty}$ is calculated after each iteration (exceptions, see later) and the absolute convergence is computed from:

$$ \frac{|y_{k=\infty} - y_k|}{y_k} - \delta_A = \Delta_A $$

where $\delta_A$ is given by the input. When $\Delta_A$ becomes zero or negative absolute convergence has been achieved.

In order to obtain complete convergence, both relative and absolute convergence must be satisfied. Equations (9), (11) and (12) are printed as output after each iteration. Eq. (11) serves an additional purpose, namely to extrapolate the pressure distribution. When $y_k$ becomes a smooth curve and starts to level off then a new extrapolated pressure distribution is calculated from:

$$ P_{ij} = \frac{y_{k=\infty}}{y_k} \cdot P_{ij}^{(n)} $$
and the pressure iterations proceed from this new distribution either until convergence is achieved or until $y_k$ again becomes sufficiently smooth that Eq. (13) may be used once more.

It remains to define the criteria for "smoothness". As applied to Eq. (11), $y_{k=0}$ is not calculated for the first 2 iterations or for the 2 iterations following a pressure extrapolation. In these cases the computer output shows:

$$y_{k=0} = 10^{-38}$$

$$\Delta_A = 1.0$$

Furthermore, $y_{k=0}$ is not calculated if the $y_k$-curve is not monotonically increasing with decreasing gradient, i.e.:

\[
\begin{align*}
&\text{a/ } (2y_{k-1} - y_{k-2} - y_k) > 0 \\
&\text{b/ } (y_k - y_{k+1}) \geq 10^{-6} y_k \\
&\text{c/ } k \leq 2, (k-k_{\text{extrapolation}}) \leq 2 \\
&\text{d/ } (2y_k - y_{k-2} - y_k) < 0 \\
&\text{e/ } (2y_{k-1} - y_{k-2} - y_k) = 0, (y_k - y_{k+1}) \neq 0 \\
&\text{f/ } (2y_{k-1} - y_{k-2} - y_k) > 0, (y_k - y_{k+1}) \leq 0 \\
&\text{g/ } (2y_{k-1} - y_{k-2} - y_k) = 0, (y_k - y_{k+1}) = 0 \\
&\text{h/ } (2y_{k-1} - y_{k-2} - y_k) > 0, 0 < (y_k - y_{k+1}) \leq 10^{-6} y_k \\
\end{align*}
\]

Graphically:

Similarly, Eq. (13) is not applied until $y_{k=0}$ is "smooth". Basically this is determined by the input item $\delta_{EX}$ which specifies how small the relative difference between two consecutive $y_{k=0}$ must be for extrapolation to be performed. In addition the $y_{k=0}$-curve must be increasing. Hence, the criteria for extrapolation becomes:
When these equations are satisfied \( \text{Eq.}(13) \) is performed, otherwise not.

These Equations (15) to (18) are all used in the 704-version of the program. In the 1620-version Eq. (18,a) is not included and \( \delta_{\text{EX}} \) in Eq. (18,b) is built into the program, given the value \( \delta_{\text{EX}} = .005 \). Furthermore, in Eq. (16) \( y_{k=0} = 10^{38} \) is replaced by \( y_{k=0} = 10^{40} \).
Input

The input form is shown on page 17. Below follow the instructions for preparing the input.

Card 1 and 2

These are heading cards to be used for identification and must always be given. Punch text in column 2 to 72 (2 to 52 in 1620 version).

Card 3

L/D Ratio This is the ratio between bearing length and journal diameter.

Inlet Angle $\Theta_{\text{in}}$. This is the angle in degrees measured from the vertical reference line to the leading edge of the bearing pad in the direction of journal rotation. For a full $360^\circ$ bearing $\Theta_{\text{in}}$ may be any value as long as $\Theta_{\text{out}} - \Theta_{\text{in}} = 360^\circ$.

Outlet Angle, $\Theta_{\text{out}}$. This is the angle in degrees measured from the vertical reference line to the trailing edge of the bearing pad in the direction of journal rotation. In a full $360^\circ$ bearing $\Theta_{\text{out}}$ may be any value as long as $\Theta_{\text{out}} - \Theta_{\text{in}} = 360^\circ$.

Inlet Pressure $P_{\text{in}}$. This is the dimensionless pressure along the leading edge of the bearing pad. If the actual pressure is $P$ psig, then

$$P_{\text{in}} = \frac{P}{\mu N L}$$

Thus, for ambient pressure $P_{\text{in}} = 0$. In the finite difference representation the pressure along the leading edge is set equal to $P_{\text{in}}$ except at the end of the bearing ($i=1$) where the pressure is ambient.

Card 4

Outlet Pressure $P_{\text{out}}$. This is the dimensionless pressure along the trailing edge of the bearing pad. For comments, see above.

Relative Convergence Limit $\delta_p$. This limit determines when relative convergence has been achieved i.e. when the relative difference between two consecutive pressure iterations is less than $\delta_p$. For details, see Eq. (9) page 6. A typical value is $\delta_p = .001$ or .002.
Absolute Convergence Limit $\delta_A$  

This limit determines when absolute convergence has been achieved, i.e., when the difference between a given pressure iteration and the corresponding extrapolated final value becomes less than $\delta_A$. For details, see Eq. (12), page 6. A typical value may be $\delta_A = 0.002$ to $0.005$. It should be noted that absolute convergence is much more strict than relative convergence i.e. $\delta_A$ should be 2 or 3 times greater than $\delta_R$ for consistency. Hence, $\delta_A = 0.002$ is rather small, especially if the mesh has many points (150 to 200 and more). Considering that $\delta_A = 0.002$ means an absolute accuracy of $0.2\%$ which in most cases is much finer than needed, and also that the method of solution is approximate anyway, $\delta_A$ should not be made any smaller than actually needed. The number of pressure iterations increases rapidly with decreasing $\delta_A$ and for most practical cases $\delta_A = 0.005$ should give sufficient accuracy.

Relaxation Factor for Pressure Calculation $f_p$  

In order to increase (or decrease) the rate of convergence of the pressure iterations a relaxation factor may be used. To illustrate assume that in the k'th iteration the pressure $P_i^{(k)}$ has been computed from Eq. (5) for any meshpoint. Then the actually stored value is not $P_i^{(k)}$ but

$$P_i = P_i^{(k-1)} + f_p (P_i^{(k)} - P_i^{(k-1)})$$

Hence, $f_p > 1$ will accelerate the convergence and $f_p < 1$ will decelerate the convergence. Setting $f_p = 1$ results in no relaxation. It should be noted that the program includes an additional feature for speeding up the calculation, namely the pressure extrapolation given by Eq. (13). If extrapolation is desired it is questionable if relaxation should also be used.

Card 5

1. **Number of Circumferential Subdivisions, m**  
   As shown in Figure 2, the circumferential length of the bearing is subdivided into m increments $\Delta X = \frac{180 (\Theta_{out} - \Theta_{in})}{m}$. Thus the number of meshpoints in the circumferential direction is $(m+1)$.

2. **Number of Axial Subdivisions, n**  
   As shown on Figure 2, the bearing half-length is subdivided into n increments $\Delta Z = \frac{1}{2n}$. Thus the number of meshpoints in the axial direction is $(n+1)$.

3. **Full or Partial Bearing**  
   If this item is 0 the program assumes that the bearing is full $360^\circ$ such that the pressure along the leading edge is equal
to the pressure along the trailing edge with a continuous gradient. Hence, item 4, Card 3, and item 1, Card 4, are ignored.

If this item is 1 the program assumes that the bearing is partial arc.

4. **Number of Pressure Iterations** In order to ensure that the computer does not use an excessive amount of time to meet the convergence limits this item sets a limit for the maximum number of pressure iterations. If this number is exceeded the computer gives the results obtained up to this point.

5. **Number of E-w-Input** Frequently the load and other quantities are wanted as a function of eccentricity ratio $E$ keeping other parameters such as $L/D$ ratio and arc length the same. The program provides for giving as many values of $E$ as desired and the present item specifies how many cases there are, i.e. how many of Card 6 (or pairs of Card 6-Card 7).

6. **Input Control** If this item is zero the program will return after completed calculation to read in an additional input set. If this item is 1 the program assumes this to be the last input set and will go to normal stop after completed calculation. Note, that this of course does not refer to the above mentioned $E$-$w$-input which is regarded as a part of an input set.

An input set starts with Card 1.

7. **Extrapolation Limit, $\delta_{gx}$** (704-version only). After each pressure iteration, say the $k$'th, all pressures are summed up. Denote the sum $y_k$. Based on 3 consecutive values of $y_k$ an exponential extrapolation is used to calculate the value of the sum after infinitely many iterations, denoted $y_{\infty}$. When the relative difference between three consecutive values of $y_{\infty}$ is less than $\delta_{gx}$ an extrapolation of the pressure distribution is performed by multiplying each pressure by $y_{\infty}/y_k$. For details, see Eq. (13) and (18), page 64. A few remarks are needed. It is obvious that the intention of the extrapolation is to speed up the convergence of the pressure iteration but it is difficult to ascertain that this is always the case. When the absolute convergence limit $\delta_A$ is not too tight the extrapolation may reduce the number of iterations by more than 50 percent but when $\delta_A$ is small it is more difficult to evaluate the effect. After an extrapolation the pressure distribution is no longer in "equilibrium" and this reflects clearly in the pressure sums $y_k$ and
very strongly in the extrapolated sums $y_k$. In general $y_k$ will be decreasing right after an extrapolation such that $y_k$ cannot be computed and therefore is given the arbitrary value $10^{38}$ ($10^{30}$ in 1620-version) in the output. After a certain number of iterations $y_k$ will slowly start to increase again and, when smooth enough, a new extrapolation will be performed, repeating the cycle. When $\delta_A$ is small and the number of meshpoints is large these cycles may be rather long which tends to reduce the benefit of the extrapolation. If it is not desired to extrapolate, set $\delta_{EX} = 0.0$. If extrapolation is desired, an example of its value is $\delta_{EX} = .005$.

Card 6

There should be as many Card 6 as given by item 5, Card 5. Each Card 6 with MATT=+1 (item 5 below) should be followed by a Card 7.

Eccentricity Ratio, $E$ This is the ratio between the journal eccentricity and the radial clearance, see Figure 1.

Attitude Angle, $\alpha$ This is the angle in degrees measured from the vertical reference line to the line of centers (i.e. the line connecting bearing and journal centers) in the direction of rotation, see Figure 1.

Radial Squeeze Film Velocity, $\xi/\omega$ This is the dimensionless velocity of the journal center along the line of centers, see Figure 1. However, if item 5, Card 6, is MATT=-1 then this item means $\Delta(\frac{\xi}{\omega})$ for use in calculating force derivatives with respect to $\xi/\omega$.

Tangential Squeeze Film Velocity, $E\xi/\omega$ This is the dimensionless velocity of the journal center perpendicular to the line of centers, see Figure 1.

MATT If MATT=-1 a Card 7 must be given. The program will perform a number of complete pressure calculations, changing the attitude angle in order to make the horizontal force component zero. Thereafter an additional 4 complete calculations will be performed to calculate $\partial f/\partial c$ and $\partial f/\partial E(\frac{\xi}{\omega})$

If MATT=0 no Card 7 is required. Only one complete calculation will be performed based on the input value for the attitude angle.

If MATT=+1 a Card 7 is required after Card 6. The program will iterate on the attitude angle in order to eliminate the horizontal force component. However, force derivatives will not be computed.
Card 7 can only be given when MATT=±1 (see item 5 above)

Eccentricity Ratio Increment \( \Delta \varepsilon \) When MATT=+1 this item is not used.
When MATT=-1 \( \Delta \varepsilon \) is used in computing the derivatives with respect to \( \varepsilon \).
Upon completion of the attitude angle iterations the program proceeds to calculate the change in bearing load as a function of \( \varepsilon \), keeping all other parameters constant. Denoting the original eccentricity ratio \( \varepsilon_0 \) (item 1, Card 6) the program performs calculations for \( \varepsilon = \varepsilon_0 + \Delta \varepsilon \) and for \( \varepsilon = \varepsilon_0 - \Delta \varepsilon \). Hence, the results may be used to calculate \( \frac{\partial \delta}{\partial \varepsilon} = \frac{\Delta \varepsilon}{2 \Delta \varepsilon} \)

Convergence Limit for Attitude Angle Iteration, \( \delta_\alpha \). \( \delta_\alpha \) is in degrees and determines when the attitude angle iteration has converged. After the \( k \)'th iteration the program computes the value of the attitude angle as

\[
\alpha^{(k+1)} = \tan^{-1}(\delta_\varepsilon)
\]

If \( |\alpha^{(k+1)} - \alpha^{(k)}| \leq \delta_\alpha \) the calculation has converged (or if item 4 below is exceeded) and the program proceeds to calculate derivatives. Otherwise a new trial value of \( \alpha \) is determined as described below and a new calculation is performed. A typical value is \( \delta_\alpha = 0.01^\circ \)

Relaxation Factor for Attitude Angle Iteration \( \delta_\varepsilon \). As described above, a new pressure calculation is performed when \( |\alpha^{(k+1)} - \alpha^{(k)}| > \delta_\alpha \). This new calculation is based on an \( \alpha \)-value determined as follows: the first calculation uses the original value for \( \alpha \) (item 2, Card 6). The two next calculations use

\[
\alpha^{(k+1)} = \alpha^{(k)} + \delta_\varepsilon (\alpha^{(k+1)} - \alpha^{(k)})
\]

Any additional calculations uses an \( \alpha \)-value determined by parabolic interpolation. \( \delta_\alpha \) is actually a function of eccentricity ratio, L/D-ratio, bearing geometry etc. It usually has a maximum around \( \varepsilon = 0.5 \) where \( \delta_\alpha = 2.5 \) to 3. In general, use \( \delta_\alpha = 1.0 \), to 1.5.

Maximum Number of Attitude Angle Iterations This item limits the number of attitude angle iterations so that an excessive number of iterations is not performed even if convergence has not been obtained. Normally 5 to 6 iterations are needed if \( \delta_\alpha = 0.01^\circ \).
Output

The output is in general self-explanatory. First come the two heading cards followed by 3 or 4 lines listing the input. Thereafter the convergence of the pressure iteration is printed as a list with 6 columns. The first column gives the number of the iteration, the second column is $\Delta R$ from Eq. (9), the third column is $\Delta A$ from Eq. (12), the fourth column is $y_k$ from Eq. (10), the fifth column is the extrapolated sum $y_{new}$ from Eq. (11) or from Eq. (15) to (17) and the sixth column gives a count of the number of iterations after last pressure extrapolation. In the 1620-version column 5 and 6 are interchanged. Note, that for convergence $\Delta R$ and $\Delta A$ must be zero or negative.

Next follows the pressure distribution listed as:

```
leading edge
rotation
trailing edge
```

The output format provides for a maximum of 7 columns. Hence, if $(n+1)$ exceeds 7 the pressure field is broken up as shown above. The last line gives the dimensionless radial force component $f_r$ (Eq. (7)), the tangential force component $f_t$ (Eq. (8)), the vertical force component $f_v$, the horizontal force component $f_h$, the total force $f_{total}$, the Sommerfeld Numbers and the calculated attitude angle $\alpha_{calc}$ where:

$$f_v = -2 \int_0^L P \cos \theta \, dx \, dz = f_r \cos \alpha + f_t \sin \alpha$$

$$f_h = 2 \int_0^L P \sin \theta \, dx \, dz = -f_r \sin \alpha + f_t \cos \alpha$$

$$f_{total} = \sqrt{f_r^2 + f_t^2}$$

$$S = \frac{1}{f_{total}} = \frac{2AN}{F/IL} \left( \frac{D}{C} \right)^2$$

$$\alpha_{calc} = \tan^{-1} \left( \frac{f_t}{f_r} \right)$$
Hence, to convert the dimensionless force $f$ to actual force $F$;

$$F = A D L N \left( \frac{P}{\xi} \right) f \quad \text{lbs}$$

Similarly, the actual pressure $\bar{P}$ in psig is obtained from the dimensionless pressure $P$ by

$$\bar{P} = A N \left( \frac{P}{\xi} \right) P \quad \text{psig}$$

**Computer Operation**

The input is on cards. In the 704-version the output is given on Tape 3, which is the only tape used by the program. Upon completion of calculation Tape 3 is given an "End of File", but not rewound, and the computer stops with 77777 in the address field. The stop is at 33617. The program is compiled with part of LSTG's TOP-system, i.e. senseswitch 1 and 2 down. The 1620-version is written in FORTRAN I, both input and output is on cards.

**Computer Time**

The two most important factors in determining the computer time is the number of meshpoints $m \times n$ and the absolute convergence limit $\delta_A$. However, it is difficult to evaluate the effect of the latter, it is approximately correct to state that the computer time is proportional to $(m \times n)^2$. This is due to two contributions: a) for each iteration the time is of course directly proportional to the number of meshpoints. This time may be estimated from:

1000 mesh point calculations (i.e. of $P_y$) per 5 seconds, 

b) the number of iterations needed for convergence increases approximately proportional to the number of meshpoints. Example: for $\delta_A = .002$ and $\delta_{EX} = .005$ the average number of iterations is:

<table>
<thead>
<tr>
<th>No. iterations</th>
<th>time</th>
</tr>
</thead>
<tbody>
<tr>
<td>$m \times n = 12 \times 6$</td>
<td>46.6</td>
</tr>
<tr>
<td>$m \times n = 16 \times 6$</td>
<td>65</td>
</tr>
<tr>
<td>$m \times n = 36 \times 10$</td>
<td>Approx. 200</td>
</tr>
</tbody>
</table>
The first two numbers are averaged over several values of $\epsilon$. It should be noted that $\epsilon$ also affects the number of iterations. As an example, take the same conditions as above. Then:

<table>
<thead>
<tr>
<th>$\epsilon$</th>
<th>$m \times n = 12 \times 6$</th>
<th>$m \times n = 16 \times 6$</th>
</tr>
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<tr>
<td>$= .1$</td>
<td>44</td>
<td>70</td>
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<tr>
<td>$= .2$</td>
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<td>75</td>
</tr>
<tr>
<td>$= .4$</td>
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<td>74</td>
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<tr>
<td>$= .6$</td>
<td>47</td>
<td>55</td>
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<tr>
<td>$= .8$</td>
<td>31</td>
<td>55</td>
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</table>

From the above both computer time and number of iterations may be estimated.
INPUT FOR PNO091 (IBM 704)
HYDRODYNAMIC INCOMPRESSIBLE JOURNAL BEARING

Card 1. Text, Col. 2-72: ____________________________________________________________

Card 2. Text, Col. 2-72: ____________________________________________________________

Card 3
(1P4E15.7)

- - - - - - - E - - - L/D -Ratio
- - - - - - - E - - - \( \theta_{in} \), Inlet Angle, degrees
- - - - - - - E - - - \( \theta_{out} \), Outlet Angle, degrees
- - - - - - - E - - - \( P_{in} \), Inlet Pressure, dimensionless

Card 4
(1P4E15.7)

- - - - - - - E - - - \( P_{out} \), Outlet Pressure, dimensionless
- - - - - - - E - - - \( \delta_{R} \), Relative Convergence Limit
- - - - - - - E - - - \( \delta_{A} \), Absolute Convergence Limit
- - - - - - - E - - - \( f_{p} \), Relaxation Factor for Pressure Calculation.

Card 5
(6I5,1PE15.7)

- - 1. m, Number of Circumferential Subdivision (m \( \leq 200 \))
- - 2. n, Number of Axial Subdivision on Half length (n \( \leq 100 \))
- - 3. 0: Full 360° Bearing  1: Partial Arc Bearing
- - 4. Maximum Number of Pressure Iterations
- - 5. Number of \( \varepsilon \)-\( \delta \)-Input
- - 6. 0: More Input sets follow  1: Last Input set
- - - - - - - E - - - 7. \( \delta_{EV} \), Extrapolation Limit
Card 6
(1P4E15.7, 15)

- . . . . . . . E . . . . Eccentricity Ratio
- . . . . . . . E . . . . α, Attitude Angle, degrees
- . . . . . . . E . . . . $\frac{\varepsilon}{\delta}$, Radial Squeeze Film Velocity, dimensionless
- . . . . . . . E . . . . $\frac{\varepsilon}{\psi}$, Tangential Squeeze Film Velocity, dimensionless

MATT, -1: Attitude Angle Iteration + Force Derivatives
0: Single Calculation
+1: Attitude Angle Iteration, no Derivatives

Card 7
(1P3E15.7, 15)

Only if Card 6, item 5, MATT=+1

- . . . . . . . E . . . . $\Delta \varepsilon$, Eccentricity Ratio Increment for $\frac{\delta f}{\delta \varepsilon}$
- . . . . . . . E . . . . $\delta_{\alpha}$, Convergence Limit for Attitude Angle Iteration, degrees.
- . . . . . . . E . . . . $\phi_{\alpha}$, Relaxation Factor for Attitude Angle Iteration

Maximum Number of Attitude Angle Iterations.

Note: An Input Set consists of Card 1 to Card 5 plus as many pairs of Card 6-Card 7 as given by item 5, Card 5.
**INPUT FOR PNOO91 (IBM 1620)**

**HYDRODYNAMIC INCOMPRESSIBLE JOURNAL BEARING**

**Card 1.** Text, Col. 2-52: ____________________________________________

**Card 2.** Text, Col. 2-52: ____________________________________________

**Card 3**

(1XE14.7, 1XE14.7, 1XE14.7, 1XE14.7)

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<td>θ₀ᵣ, Outlet Angle, degrees</td>
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<tr>
<td></td>
<td></td>
<td>Pᵢ₀, Inlet Pressure, dimensionless</td>
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</tbody>
</table>

**Card 4**

(1XE14.7, 1XE14.7, 1XE14.7, 1XE14.7)

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<td>Pᵢᵣ, Outlet Pressure, dimensionless</td>
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<td></td>
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<td>δᵣᵢ, Relative Convergence Limit</td>
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<tr>
<td></td>
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<td>δᵣₒ, Absolute Convergence Limit</td>
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<td>fₓᵢᵢ, Relaxation Factor for Pressure Calculation</td>
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**Card 5**

(1XI4, 1XI4, 1XI4, 1XI4, 1XI4, 1XI4, 1XI4)

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<td>n, Number of Axial Subdivision on Half length (n ≤ 12)</td>
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<td>Maximum Number of Pressure Iterations</td>
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<td>5.</td>
<td>Number of θ-Input</td>
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Card 6
(1XE14.7, 1XE14.7, 1XE14.7, 1XE14.7, 1X14)

- - - - - - - E - - -
ε, Eccentricity Ratio
- - - - - - - E - - -
α, Attitude Angle, degrees
- - - - - - - E - - -
vr, Radial Squeeze Film Velocity, dimensionless
- - - - - - - E - - -
vθ, Tangential Squeeze Film Velocity
- MATT, -1: Attitude Angle Iteration + Force Derivatives.
0: Single Calculation.
+1: Attitude Angle Iteration, no Derivatives.

Card 7
(1XE14.7, 1XE14.7, 1XE14.7, 1X14)
Only if Card 6, item 5, MATT=1

- - - - - - - E - - -
Δε, Eccentricity Ratio Increment for \( \frac{d\epsilon}{d\alpha} \)
- - - - - - - E - - -
Δα, Convergence Limit for Attitude Angle Iteration, degrees
- - - - - - - E - - -
ζ, Relaxation Factor for Attitude Angle Iteration
- Maximum Number of Attitude Angle Iterations

Note: An Input Set consists of Card 1 to Card 5 plus as many pairs of Card 6 – Card 7 as given by item 5, Card 5.
### Table I

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</tbody>
</table>

| 30° ARC | 476.0 | .01 | 87.38 | .55×10⁻⁵ | .006979 | .0131 | .0021 | .0177 | -.203 | .013 |
| 43.8 | .1 | 71.23 | .00703 | .0799 | .0633 | .0207 | .228 | .16 | -.194 | .0813 |
| 20.4 | .2 | 56.84 | .0268 | .186 | .2 | .0410 | .275 | .268 | -.171 | .155 |
| .25 .573 | 8.63 | .35 | 42.78 | .0851 | .502 | .593 | .0787 | .445 | .487 | -.146 | .272 |
| 3.76 | .5 | 33.29 | .222 | 1.37 | 1.56 | .146 | .185 | .886 | -.0874 | .472 |
| 1.44 | .65 | 25.88 | .625 | 4.87 | 4.92 | .303 | 2.2 | 1.94 | .19 | .934 |
| .383 | .8 | 19.68 | 2.46 | 29.2 | 24.8 | .879 | 9.46 | 6.0 | 2.31 | 2.37 |
| .0262 | 95 | 12.15 | 37.3 | 1410 | 611.0 | 8.04 | 247.0 | 76.0 | 27.3 | 12.3 |

| 209.0 | .01 | 86.76 | 2.69×10⁻⁴ | .0274 | .0219 | .00478 | .479 | .0532 | -.467 | .0402 |
| 20.3 | .1 | 67.19 | .0191 | .217 | .185 | .0456 | .505 | .381 | -.413 | .256 |
| 9.16 | .2 | 52.14 | .067 | .476 | .565 | .0862 | .597 | .669 | -.317 | .443 |
| 3.91 | .35 | 38.44 | .2 | 1.2 | 1.57 | .159 | .919 | 1.14 | -.139 | .741 |
| 3.0 | .4 | 35.18 | .272 | 1.61 | 2.12 | .192 | 1.09 | 1.37 | -.0468 | .89 |
| .5 1.15 | 1.73 | .5 | 29.68 | .501 | 3.17 | 3.93 | .286 | 1.73 | 2.0 | .185 | 1.22 |
| .951 | .6 | 23.13 | .952 | 6.76 | 7.81 | .447 | 3.0 | 3.17 | .707 | 1.88 |
| .678 | .65 | 23.11 | 1.36 | 10.6 | 11.6 | .579 | 4.26 | 4.18 | 1.26 | 2.38 |
| .188 | .8 | 17.63 | 5.06 | 59.1 | 53.9 | 1.61 | 17.2 | 12.5 | 7.48 | 6.02 |
| .0158 | .95 | 11.79 | 62.0 | 2111 | 1040 | 12.9 | 365.0 | 129.0 | 68.8 | 27.7 |

| 157.0 | .01 | 86.16 | 4.25×10⁻⁴ | .0431 | .0271 | .00636 | .637 | .0865 | -.623 | .0667 |
| 15.3 | .1 | 65.0 | .0276 | .315 | .287 | .0591 | .663 | .551 | -.518 | .388 |
| 6.97 | .2 | 49.92 | .0924 | .662 | .831 | .11 | .77 | .923 | -.359 | .643 |
| .75 1.72 | 3.0 | .35 | 36.65 | .267 | 1.61 | 2.21 | .199 | 1.16 | 1.53 | -.079 | 1.03 |
| 1.34 | .5 | 28.25 | .656 | 4.17 | 5.4 | .353 | 2.15 | 2.62 | .459 | 1.7 |
| .53 | .65 | 22.03 | 1.75 | 13.6 | 15.6 | .708 | 5.2 | 5.38 | 2.17 | 3.24 |
| .15 | .8 | 16.90 | 6.36 | 72.5 | 70.1 | 1.93 | 20.3 | 15.9 | 11.0 | 8.08 |
| .0136 | .95 | 11.7 | 71.9 | 2360 | 1230 | 14.9 | 407.0 | 150.0 | 91.9 | 34.6 |

<p>| 60° ARC | 371.0 | .01 | 87.85 | 1.01×10⁻³ | .0105 | .0268 | .00269 | .269 | .0167 | -.286 | .0124 |
| 35.8 | .1 | 72.11 | .00859 | .0981 | .0876 | .0266 | .289 | .173 | -.246 | .0923 |
| 16.0 | .2 | 58.45 | .0327 | .228 | .25 | .0533 | .332 | .324 | -.224 | .178 |
| .25 .477 | 6.79 | .35 | 44.54 | .105 | .612 | .740 | .103 | .564 | .603 | -.191 | .313 |
| 3.01 | .5 | 35.0 | .272 | 1.66 | 1.94 | .19 | 1.07 | 1.08 | -.11 | .533 |
| 1.19 | .65 | 27.76 | .746 | 5.71 | 5.98 | .393 | 2.7 | 2.28 | .224 | .96 |
| .333 | .8 | 21.67 | 2.79 | 31.6 | 28.8 | 1.11 | 11.5 | 6.94 | 2.18 | 2.17 |
| .0349 | .95 | 13.11 | 39.1 | 1440 | 620.0 | 9.11 | 266.0 | 81.7 | 20.0 | 9.39 |</p>
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<th>(B)</th>
<th>(s)</th>
<th>(\varepsilon)</th>
<th>(\alpha)</th>
<th>(f_r)</th>
<th>(\frac{\Delta f}{\Delta s})</th>
<th>(\frac{\Delta f}{\Delta (\varepsilon^2)})</th>
<th>(f_t)</th>
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Fig. 1 Six 50° Tilting Pads, Centrally pivoted, L/D = 25, 
L/b = .573, C'/C = 1. Load between pads. No pad 
terms.
$S = \frac{\mu_{DL}(\frac{R}{C})^2}{W} \left[ \frac{C}{W} \right]$
Fig. 3 Six 50° Tilting Pads, Centrally pivoted, L/D = 0.75, L/B = 1.719, C'/C = 1. Load between pads. No pad inertia.
Fig. 4  Five 60° Tilting Pads, Centrally pivoted, $L/D = .25$, $L/B = .477$, $C'/C = 1$. Load between pads. No pad inertia.
Fig. 5 Five 60° Tilting Pads, Centrally pivoted, L/D = .5, L/B = .955, c'/c = 1. Load between pads. No pad inertia.
Fig. 6 Five 60° Tilting Pads, Centrally pivoted, L/D = .75, L/B = 1.432, C'/C = 1. Load between pads. No pad inertia.
Fig. 7  Four 80° Tilting Pads, Centrally pivoted, L/D = .5, L/B = .716, C'/C = 1. Load between pads. No pad inertia.
Fig. 8 Four 80° Tilting Pads, Centrally pivoted, L/D = .75, L/B = 1.074, C'/C = 1. Load between pads. No pad inertia.
Fig. 9 Four 80° Tilting Pads, Centrally pivoted, L/D = 1.0,
L/B = 1.432, C'/C = 1. Load between pads. No pad inertia.
Fig. 10 Six 50° Tilting Pads, Centrally pivoted, L/B = .5,
L/B = 1.146, C'/C = 1. Load on pad. No pad inertia.
Fig. 11 Five 60° Tilting Pads, Centrally pivoted, L/D = .5, L/B = .955, C'/C = 1. Load on pad. No pad inertia.
Fig. 12: Four 80° Tilting Pad, Generally Allocated T/D = .75.
L/B = 1.074, C/C = 1, Load on pad. No pad inertia.
Fig. 13 Four 80° Tilting Pads, Centrally pivoted, L/D = .75, L/B = 1.074, C'/C = 1. Load between pads. Effect of pad inertia.
Fig. 14 Four 80° Tilting Pads, Centrally pivoted, L/D = .75,
L/B = 1.074, C'/C = 1. Load between pads. Effect of
load (horizontal)
Fig. 15 Four 80° Tilting Pads, Centrally pivoted, L/D = .75,
L/W = 1.074, C'/C = 1. Load between pads. Effect
of pad inertia.
Fig. 16 Four 80° Tilting Pads, Centrally pivoted, L/D = .75, L/B = 1.074, C'/C = 1. Load between pads. Effect of pad inertia.
Fig. 17: Twelve 25° Elliptic Pads, Centrally pivoted, L/D = 25. Vertical rotor, Effect of pre-load (LC/OC).
Fig. 18 Six 50° Tilting Pads, Centrally pivoted, L/B = .75, L/B = 1.146. Vertical Rotor. Effect of pre-load (1-C'/C).
Fig. 19 Four 80° Tilting Pads, Centrally pivoted, L/D = .75,
Fig. 20 Six 50° Tilting pads, Centrally pivoted, L/D = .75,
L/B = 1.719, C'/C = 1. Load between pads. Comparison
with Test Data.
Fig. 21 Four 80° Tilting Pads, Centrally pivoted, L/D = .75, L/B = 1.074, C'/C = 1. Load between pads. Comparison with Test Data.
Fig. 22 Coordinate Systems for Analysis.
NOMENCLATURE

C Pad clearance (radius of curvature of pad minus journal radius) inch.

C' Pivot circle clearance (radius of pivot circle minus journal radius)

C_{xx}, C_{yy}, C_{yx}, C_{yy} Bearing damping coefficients, lbs sec/in.

C_{ff}, C_{ff}', C_{ff}, C_{ff}' Fixed pad damping coefficients, lbs sec/in.

C_{ff}', C_{ff}, C_{ff}', C_{ff}' Tilting pad damping coefficients, lbs sec/in.

D Journal diameter, inch

e Journal center eccentricity with respect to pad center, inch

\( e_c \) Journal center eccentricity with respect to bearing center, inch

F Load on pad, lbs.

F_r, F_t Radial and tangential components of pad load, lbs.

F_y, F_\gamma Components in \( \xi \) and \( \gamma \)-direction of pad load, lbs (see Fig. 22)

f = F/\omega, dimensionless pad force

f_r, f_t Radial and tangential components of dimensionless pad force

f_y, f_\gamma Components in \( \xi \) and \( \gamma \)-direction of dimensionless pad force

I Transverse mass moment of inertia of shoe around pivot, lbs in sec^2

K_{xx}, K_{xy}, K_{yx}, K_{yy} Bearing spring coefficients, lbs/in.

K_{ff}, K_{ff}', K_{ff}, K_{ff}' Fixed pad spring coefficients, lbs/in.

K_{ff}', K_{ff}, K_{ff}', K_{ff}' Tilting pad spring coefficients, lbs/in.

L Bearing length, inch

M = I/R^2, equivalent pad mass, lbs sec/in

M_{crit} Value of equivalent pad mass to cause pad motion resonance, lbs sec/in

N Rotational speed of journal, RPS

p, q Coefficients defined by Eq. (31) and (32)

R Journal radius, inch
\( R_p \)  
Radius from pad center to actual pivot point of pad, inch

\( S \)  
\( = (\mu NDL/W) \cdot (R/C)^2 \), bearing Sommerfeld number

\( S_p \)  
\( = (\mu NDL/F) \cdot (R/C)^2 \), pad Sommerfeld number

\( W \)  
Bearing load, lbs.

\( x, y \)  
Coordinates of journal center with respect to the bearing, see Fig. 22, inch

\( \varepsilon \)  
\( = e/C \), eccentricity ratio with respect to the pad center

\( \varepsilon_o \)  
\( = e_o/C \), eccentricity ratio with respect to the bearing center

\( \eta_p \)  
Amplitude for pad center motion, see Fig. 22, inch

\( \eta \)  
Amplitude of the center of a massless pad, inch.

\( \lambda \)  
\( = (\mu RL/\pi) \cdot (R/C)^2 \) bearing coefficient

\( \mu \)  
Lubricant viscosity, lbs.sec./in\(^2\)

\( \xi, \eta \)  
Coordinates of journal center with respect to the pad, see Fig. 22, inch

\( \varphi \)  
Attitude angle with respect to the pad load line, radians

\( \varphi_o \)  
Attitude angle with respect to the bearing load line, radians

\( \psi \)  
Angle from vertical (negative x-axis) to pad pivot point, see Fig. 22, inch

\( \omega \)  
Angular speed of shaft, radian/sec.